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Citation: Zhao, Benzhong, Christopher W. MacMinn, and Ruben Juanes. "Residual Trapping, Solubility Trapping and Capillary Pinning Complement Each Other to Limit CO₂ Migration in Deep Saline Aquifers." Energy Procedia 63 (2014): 3833–3839.

As Published: <http://dx.doi.org/10.1016/j.egypro.2014.11.412>

Publisher: Elsevier

Persistent URL: <http://hdl.handle.net/1721.1/101642>

Version: Final published version: final published article, as it appeared in a journal, conference proceedings, or other formally published context

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GHGT-14

Residual trapping, solubility trapping and capillary pinning complement each other to limit CO₂ migration in deep saline aquifers

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Abstract

We derive a theoretical model for the post-injection migration of a CO₂ gravity current in a confined, sloping aquifer under the influence of residual trapping, solubility trapping, and capillary pinning. The resulting model consists of two coupled partial differential equations that describe the local thickness of the buoyant CO₂ current and the thickness of the mound of brine saturated with dissolved CO₂ as a function of time. We apply this model to a representative geological formation and provide estimates of the lifetime of the buoyant CO₂ current and its maximum migration distance. Our analysis shows that residual trapping, solubility trapping, and capillary pinning complement each other in limiting the ultimate migration distance of CO₂ gravity currents. The relative contribution of residual trapping, solubility trapping, and capillary pinning varies as a function of the injection volume. Our model can be used as a screening tool to evaluate the potential of deep saline aquifers for large-scale CO₂ sequestration.

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Peer-review under responsibility of the Organizing Committee of GHGT-12

Keywords: Geologic storage; capillary trapping; capillary pinning; convective dissolution; post-injection; sharp-interface; migration distance

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1. Introduction

Gravity currents in porous media have attracted much interest in the context of geological carbon dioxide (CO₂) storage, where supercritical CO₂ is captured from the flue gas of power plants and injected underground into deep saline aquifers [1-4]. Saline aquifers are geologic layers of permeable rock located 1 to 3 km below the surface, saturated with salty groundwater, and, in general, lined on top by a layer of much less permeable “caprock” [5]. They have been considered as ideal storage sites for CO₂ (see, e.g., [3,4]) due to their geologic prevalence around the globe, as well as their large storage capacity [6,7]. At reservoir conditions, CO₂ is less dense than the ambient groundwater. As a result, it will migrate upward due to buoyancy and spread along the caprock. Capillarity can be important in the spreading and migration of the buoyant CO₂ after injection, because the typical pore size is very small (~10 - 100 μm), but the impact of capillarity on these flows is not well understood.

In previous studies [8,9], we studied the impact of capillarity on the buoyant spreading of a finite-size current of non-wetting fluid into a dense, wetting fluid in a vertically confined aquifer. We demonstrated via simple table-top experiments in packings of glass beads that capillary pressure hysteresis pins a portion of the fluid-fluid interface. We showed that capillary pressure hysteresis in this system is caused by the fundamental difference in the pore-scale invasion patterns between drainage and imbibition, and is present even in the absence of contact angle hysteresis. The length of the pinned interface scales with the relative importance of capillary and gravity forces. The horizontal extent of the pinned portion of the interface grows with time and this is responsible for ultimately limiting the spreading of the buoyant current to a finite distance. This is in stark contrast to miscible gravity currents, which in the absence of diffusion, spreads forever. We developed a theoretical model that captures the evolution of immiscible gravity currents and predicts the maximum migration distance.

Other physical mechanisms such as residual trapping, where tiny blobs of CO₂ are immobilized by capillary forces, and solubility trapping, where CO₂ dissolves into the ambient groundwater, also play important roles in the migration of CO₂ gravity currents. The goal of this paper is to derive a theoretical model for the post-injection migration of a CO₂ gravity current in a confined, sloping aquifer under the combined influence of residual trapping, solubility trapping, and capillary pinning.

2. Mathematical Model

The study of gravity currents under the general framework of sharp interface models has received renewed attention in the context of geologic CO₂ sequestration in the past decade. Sharp interface models are able to simulate flow over large distances due to their relatively low computational cost compared to full numerical simulations, and they can be formulated to include the trapping mechanisms encountered in CO₂ storage. In particular, residual trapping is accounted for by incorporating a discontinuous coefficient on the accumulation term of the sharp interface model [10-12]. Solubility trapping can be modeled by incorporating convective dissolution as a constant flux of CO₂ per unit length of CO₂-brine interface [13-17]. The rate of convective dissolution slows down as the brine underneath the CO₂ current becomes saturated with dissolved CO₂, a process known as convective shutdown [18-20]. *Hidalgo et al.* (2013) [17] demonstrated that this decrease in dissolution rate can be effectively captured by modeling the CO₂ current and the CO₂-rich brine layer as a coupled system, and assuming that convective dissolution shuts down once the layer of brine saturated with CO₂ fills the region beneath the buoyant CO₂. *Zhao et al.* (2014) [9] developed a model that captures capillary pinning of immiscible gravity currents and predicts the maximum migration distance.

Here, we present a theoretical model that incorporates residual trapping, solubility trapping and capillary pinning. The complete derivation of the model and the details of the underlying assumptions will be given in a future publication. The model consists of two coupled partial differential equations that describe the evolution of the buoyant CO₂ current and the mound of brine saturated with dissolved CO₂ (hereafter referred to as the dense mound)

$$\begin{aligned} & \frac{\partial h_g}{\partial t} + \frac{\kappa \lambda_g \Delta \rho_{gw} g}{(1 - S_{wc} - S_{gr} \tilde{R}) \phi} \frac{\partial}{\partial x} \left[(1 - f_g) h_g (\sin \theta - \cos \theta) \frac{\partial h_g}{\partial x} - \frac{1}{\Delta \rho_{gw} g \cos \theta} \frac{\partial P_c}{\partial x} \right] \\ & + \frac{\kappa \lambda_d \Delta \rho_{wd} g}{(1 - S_{wc} - S_{gr} \tilde{R}) \phi} \frac{\partial}{\partial x} \left[f_d h_d (\sin \theta + \cos \theta) \frac{\partial h_d}{\partial x} \right] = -\frac{q_d}{\phi} \varepsilon, \end{aligned} \quad (1)$$

and

$$\begin{aligned} & \frac{\partial h_d}{\partial t} - \frac{\kappa \lambda_g \Delta \rho_{gw} g}{\phi} \frac{\partial}{\partial x} \left[f_d h_g (\sin \theta - \cos \theta) \frac{\partial h_g}{\partial x} - \frac{1}{\Delta \rho_{gw} g \cos \theta} \frac{\partial P_c}{\partial x} \right] \\ & - \frac{\kappa \lambda_d \Delta \rho_{wd} g}{\phi} \frac{\partial}{\partial x} \left[(1 - f_d) h_d (\sin \theta + \cos \theta) \frac{\partial h_d}{\partial x} \right] = \frac{q_d (1 - S_{wc} - S_{gr} \tilde{R})}{\phi \chi_v} \varepsilon, \end{aligned} \quad (2)$$

where h_g and h_d are the local thickness of the CO₂ current and the thickness of the dense mound. κ is the intrinsic permeability of the aquifer, ϕ is the porosity of the aquifer, and θ is the slope of the aquifer. $\lambda_g = k_{rg} / \mu_g$ and $\lambda_d = k_{rd} / \mu_d$ are the mobilities of the CO₂ current and dense mound respectively. μ_g, μ_d and k_{rg}, k_{rd} are the dynamic viscosities and the relative permeabilities of the respective fluid phases. $\Delta \rho_{gw}$ is the density difference between the buoyant CO₂ and the ambient brine, $\Delta \rho_{wd}$ is the density difference between the ambient brine and the dense mound. S_{wc} is the connate water saturation, which represents the fraction of pore space occupied by immobile brine. Similarly, S_{gr} is the residual gas saturation, which represents the fraction of pore space occupied by trapped blobs of CO₂. The discontinuous coefficient \tilde{R} captures the fact that residual trapping only occurs in areas where the bulk CO₂ current has been displaced by brine through imbibition. Hence, $\tilde{R} = 1$ if $\partial h_g / \partial t < 0$, and $\tilde{R} = 0$ if $\partial h_g / \partial t \geq 0$. The local capillary pressure across the CO₂/brine interface P_c is given by the Young-Laplace equation and $P_c \sim \gamma / d$, where γ is the interfacial tension between CO₂ and brine, d is the mean grain size of the porous media. We take P_c to be constant and equal to a characteristic imbibition capillary pressure where the CO₂ current is displaced by brine. Similarly, we take P_c to be constant and equal to the drainage capillary pressure where the CO₂ current displaces the ambient brine [8,9]. Along the pinned interface, P_c transitions from the imbibition capillary pressure to the drainage capillary pressure. This transition in capillary pressure is offset by the change in hydrostatic pressure along the pinned interface, with a corresponding difference in interface height of $\Delta h_c = \Delta P_c / \Delta \rho_{gw} g \cos \theta$. Hence, the length of the pinned interface $\Delta h_c \sim \text{Bo}^{-1} = \gamma / \Delta \rho_{gw} g H d$, where H is the thickness of the aquifer. The fractional flow functions f_g and f_d are given by

$$f_g = \frac{h_g \lambda_g}{h_g \lambda_g + h_d \lambda_d + \lambda_w (H - h_g - h_d)}, \quad f_d = \frac{h_d \lambda_d}{h_g \lambda_g + h_d \lambda_d + \lambda_w (H - h_g - h_d)}. \quad (3)$$

q_d is the volumetric rate of convective dissolution per unit area of fluid-fluid interface, and is given by

$$q_d = \alpha \chi_v \Delta \rho_{wd} g \kappa \lambda_d, \quad (4)$$

where α is a constant roughly equal to 0.01 [21] and χ_v is the solubility of CO₂ in brine, expressed as the volume of free-phase CO₂ that can be dissolved per unit volume of brine saturated with CO₂ [7,14]. The discontinuous coefficient ε captures the fact that convective dissolution shuts down locally where the CO₂ saturated dense mound

fills the aquifer beneath the buoyant CO₂ current [15,17]. Hence, $\varepsilon = 1$ if $h_g + h_d < H$, and $\varepsilon = 0$ if $h_g + h_d = H$.

We write equation (1) and equation (2) in its dimensionless form by scaling h_g , h_d and x by the aquifer height H , and we scale t by the characteristic time

$$T_c = \frac{H\phi}{\kappa\lambda_d\Delta\rho_{wd}g}. \quad (5)$$

The dimensionless equations are then given by

$$\frac{\partial \tilde{h}_g}{\partial \tilde{t}} + \frac{1}{(1-S_{wc}-S_{gr}\tilde{R})} \frac{\partial}{\partial \tilde{x}} \left[\delta(1-\tilde{f}_g)\tilde{h}_g(\sin\theta - \cos\theta \frac{\partial \tilde{h}_g}{\partial \tilde{x}} - \frac{\partial \tilde{h}_c}{\partial \tilde{x}}) + \tilde{f}_g\tilde{h}_d(\sin\theta + \cos\theta \frac{\partial \tilde{h}_d}{\partial \tilde{x}}) \right] = -\tilde{D}, \quad (6)$$

$$\frac{\partial \tilde{h}_d}{\partial \tilde{t}} - \frac{\partial}{\partial \tilde{x}} \left[\delta\tilde{f}_d\tilde{h}_g(\sin\theta - \cos\theta \frac{\partial \tilde{h}_g}{\partial \tilde{x}} - \frac{\partial \tilde{h}_c}{\partial \tilde{x}}) + (1-\tilde{f}_d)\tilde{h}_d(\sin\theta + \cos\theta \frac{\partial \tilde{h}_d}{\partial \tilde{x}}) \right] = \frac{\tilde{D}}{\chi_v}(1-S_{wc}-S_{gr}\tilde{R}), \quad (7)$$

where we have defined two additional dimensionless parameters

$$\delta = \frac{\lambda_g\Delta\rho_{gw}}{\lambda_d\Delta\rho_{wd}}, \quad \tilde{D} = \frac{q_d}{\kappa\lambda_d\Delta\rho_{wd}g}\varepsilon. \quad (8)$$

3. Model Application

We now use the model to investigate the interplay between residual trapping, solubility trapping, and capillary pinning in the context of large-scale carbon sequestration. *MacMinn & Juanes* (2013) [16] studied the effects of residual trapping and solubility trapping and demonstrated that the combination of the two mechanisms always traps the CO₂ much more effectively than either mechanism alone. Here, we study the impact of capillary pinning as an additional trapping mechanism by modeling the post-injection migration of a CO₂ current in the Paluxy Sandstone located in the East Texas Basin. The aquifer and the CO₂/brine properties associated with the Paluxy Sandstone are presented in *Szulcowski et al.* (2012) [20] and summarized in Table 1 below. Specifically, we compare the lifetime of the spreading CO₂ current and its maximum migration distance with and without the influence of capillary pinning.

We assume that CO₂ is injected uniformly along a linear, 200 km array of wells, and we model the evolution of the planar cross-sectional profile of the resulting CO₂ current. The shape of the CO₂ current at the end of injection can be described by the analytical equation below [11,12,22],

$$\tilde{h}_g(\tilde{x}) = \frac{\sqrt{\frac{lM_{gw}}{H\tilde{x}} - 1}}{M_{gw} - 1}, \quad (9)$$

where $M_{gw} = \mu_w / \mu_g$ is the viscosity ratio between CO₂ and the ambient brine, l is the characteristic length scale of the injected CO₂. The product of l and H equals to the initial cross-sectional area of the CO₂ current. We use the solution to Equation 9 as the initial condition for our model.

Table 1. Model input parameters for the Paluxy Sandstone.

Paluxy Sandstone	Bo^{-1}	κ	ϕ	θ	S_{gr}	H	$\Delta\rho_{gw}$	$\Delta\rho_{wd}$	μ_g	μ_d	χ_v
	[-]	m ²	[-]	[°]	[-]	[m]	[kg/m ³]	[kg/m ³]	[mPa-s]	[mPa-s]	[-]
	1.89E-02	3E-13	0.20	1	0.3	15	400	6	0.06	0.8	0.05

Figure 1 shows the post-injection migration of 7.2 Mt of injected CO₂ in the Paluxy Sandstone. The CO₂ current is subject to residual trapping and solubility trapping here, but the effects of capillary pinning are not included. In this scenario, the CO₂ current reaches a maximum migration distance of 8.2 km at approximately 200,000 years, after which it starts to retreat due to continued dissolution. Figure 2 shows the post-injection migration of 7.2 Mt of CO₂ in the Paluxy Sandstone, under the same conditions as the simulation shown in Figure 1, except that capillary pinning is now included in the simulation. In this case, capillary pressure hysteresis pins an increasingly larger portion of the CO₂/brine interface as the CO₂ current migrates, until the entire current becomes pinned. Comparison with Figure 1 shows that the additional contribution of capillary pinning helps stop the migration of the CO₂ current much earlier (64,000 years), and at a shorter distance (5.9 km).

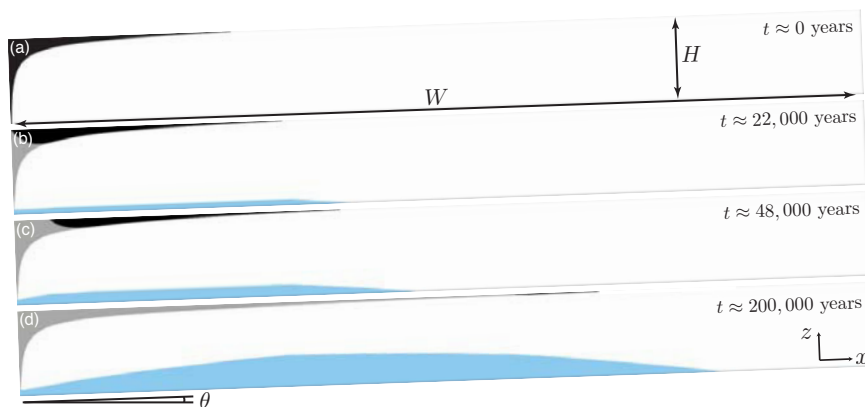


Fig. 1. Simulation of post-injection migration of 7.2 Mt tonnes of injected CO₂ in the Paluxy Sandstone, without including the effects of capillary pinning. After injection stops, the buoyant CO₂ current (black) migrates up-dip, leaving behind a trail of residually trapped CO₂ blobs in its wake (grey). CO₂ is also dissolved into the ambient brine due to convective mixing, forming a dense mound of brine saturated with dissolved CO₂ (blue). CO₂ and brine are assumed to be completely miscible here and capillarity is ignored. Note that the vertical scale of the aquifer is greatly exaggerated in the figure. The aspect ratio of the model domain is $W/H = 800$. The CO₂ current reaches its maximum migration distance after approximately 200,000 years (d), after which it starts to retreat due to continued dissolution.

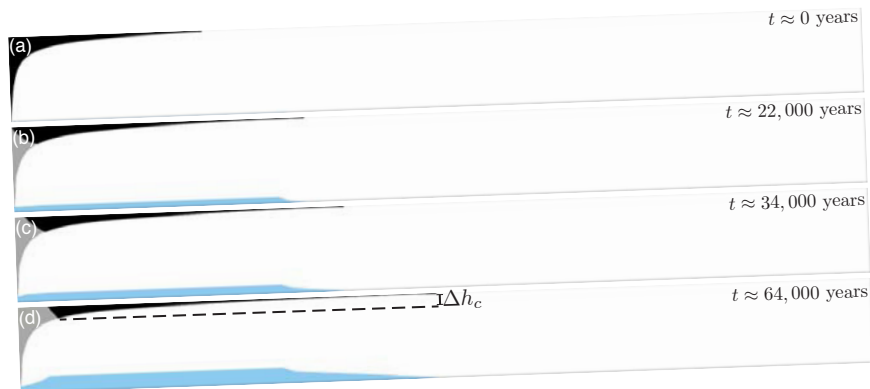


Fig. 2. Simulation of post-injection migration of 7.2 Mt tonnes of injected CO₂ in the Paluxy sandstone, including the effects of capillary pinning. Capillary pressure hysteresis pins an increasingly larger portion of the CO₂/brine interface as the CO₂ current migrates, until the entire current becomes pinned (d). Complete pinning of the CO₂/brine interface corresponds to a CO₂ current thickness of Δh_c , at which point the current lacks the necessary hydrostatic pressure to overcome capillary pressure hysteresis and stops spreading. Comparison with Figure 1 shows that the additional contribution of capillary pinning help stop the migration of the CO₂ current earlier ($t \approx 64,000$ years), and at a shorter distance.

When the CO₂ current reaches its maximum migration distance, the previously mobile CO₂ is either residually trapped, dissolved into the ambient brine, or immobilized by capillary pinning. We investigate the varying contribution of each of the trapping mechanisms at the end of the mobile CO₂ current’s lifetime, as a function of the amount of CO₂ injected (Figure 3). We find that the majority of the CO₂ is trapped by capillary pinning. However, capillary pinning does become less important as a larger amount of CO₂ is injected, whereas solubility trapping becomes more important. This can be attributed to the fact that convective dissolution occurs only at the CO₂-brine interface. A larger amount of injected CO₂ creates a longer CO₂-brine interface, which enhances solubility trapping.

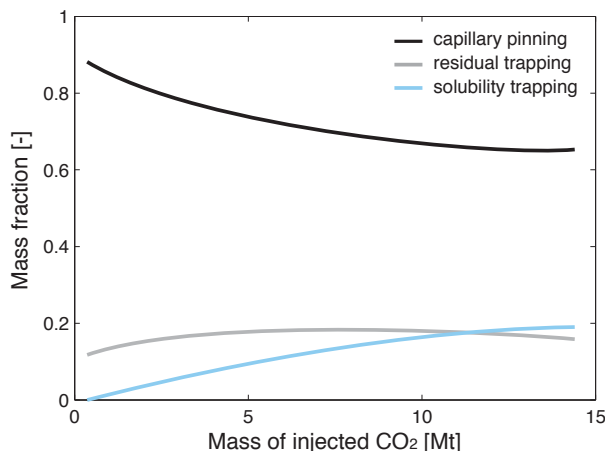


Fig. 3. Mass fraction of CO₂ immobilized by capillary pinning, residual trapping, and solubility trapping, as a function of the amount of CO₂ injected.

We have presented a model for the post-injection migration of CO₂ gravity currents in a confined, sloping aquifer, subject to residual trapping, solubility trapping, and capillary pinning. We have shown that residual trapping, solubility trapping, and capillary pinning complement each other to limit the maximum migration distance, as well as the duration of migration of a buoyant CO₂ current. The relative contribution of residual trapping, solubility trapping, and capillary pinning varies as a function of the total amount of CO₂ injected.

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