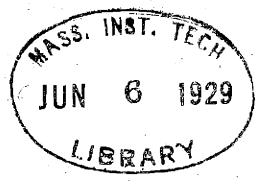


*Chem. eng'g dept.  
Thesis  
1929*



The Effect of Heat Transfer  
on Friction Factors in Fanning's Equation.

by

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B. S. in Chemical Engineering, University of Maine 1927

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Submitted in Partial Fulfillment  
of the Requirements for the  
Degree of Master of Science in  
Chemical Engineering Practice  
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June 1, 1928.

Professor A. L. Merrill,  
Secretary of the Faculty,  
Massachusetts Institute of Technology,  
Cambridge, Mass.

Dear Sir-

We submit herewith a thesis entitled  
"The Effect of Heat Transfer on the Friction  
Factors in Fanning's Equation" in partial ful-  
fillment of the requirements for the degree of  
Master of Science in Chemical Engineering Prac-  
tice.

Respectfully submitted,

\_\_\_\_\_  
\_\_\_\_\_  
\_\_\_\_\_

164917

### ACKNOWLEDGMENT

The authors wish to express their appreciation to Professor W. H. McAdams for his ever available suggestions and criticisms.

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The Effect of  
Heat Transfer on the Friction Factor  
in Fanning's Equation.

While it is a quite definitely established fact, that the Fanning equation<sup>(1)</sup> fits all cases of isothermal flow of fluids through pipes, some data in a thesis by White<sup>(2)</sup> seems to indicate that the actual pressure drop as measured during the heating of a fluid in an interchanger is considerably less than one would predict by the use of this equation. But as White was not primarily interested in the pressure drop, he did not attempt to correlate his observations or to explain this deviation. And his results cover too small a range to warrant using them as a basis of calculation. Many industries now make use of fluid to fluid heat exchangers and similar types of apparatus, so this thesis was undertaken with a view to providing a more reliable basis for the calculation of pressure drops in the design of such equipment.

The apparatus required for this investigation is quite similar to that required in the actual determination of heat transfer coefficients. So it was

thought advisable to take sufficient data to make possible the calculation of both friction factors and heat transfer coefficients. To make sure that the experimental method and apparatus gave results in accord with the established curve of friction factors, several runs were made both with water and with oil covering a range of the modulus  $\frac{DuS}{\mu}$  from 0.02 to 5.0. Thus this range includes both viscous and turbulent motion and the critical range that lies between these two. Also as another check on the method used, the actual heat transfer coefficients have been figured and compared with the work of Morris and Whitman<sup>(3)</sup>. The greater part of the work, however, was the study of the effect of heating and cooling on the actual value of the friction factor for fluids flowing in horizontal pipes.

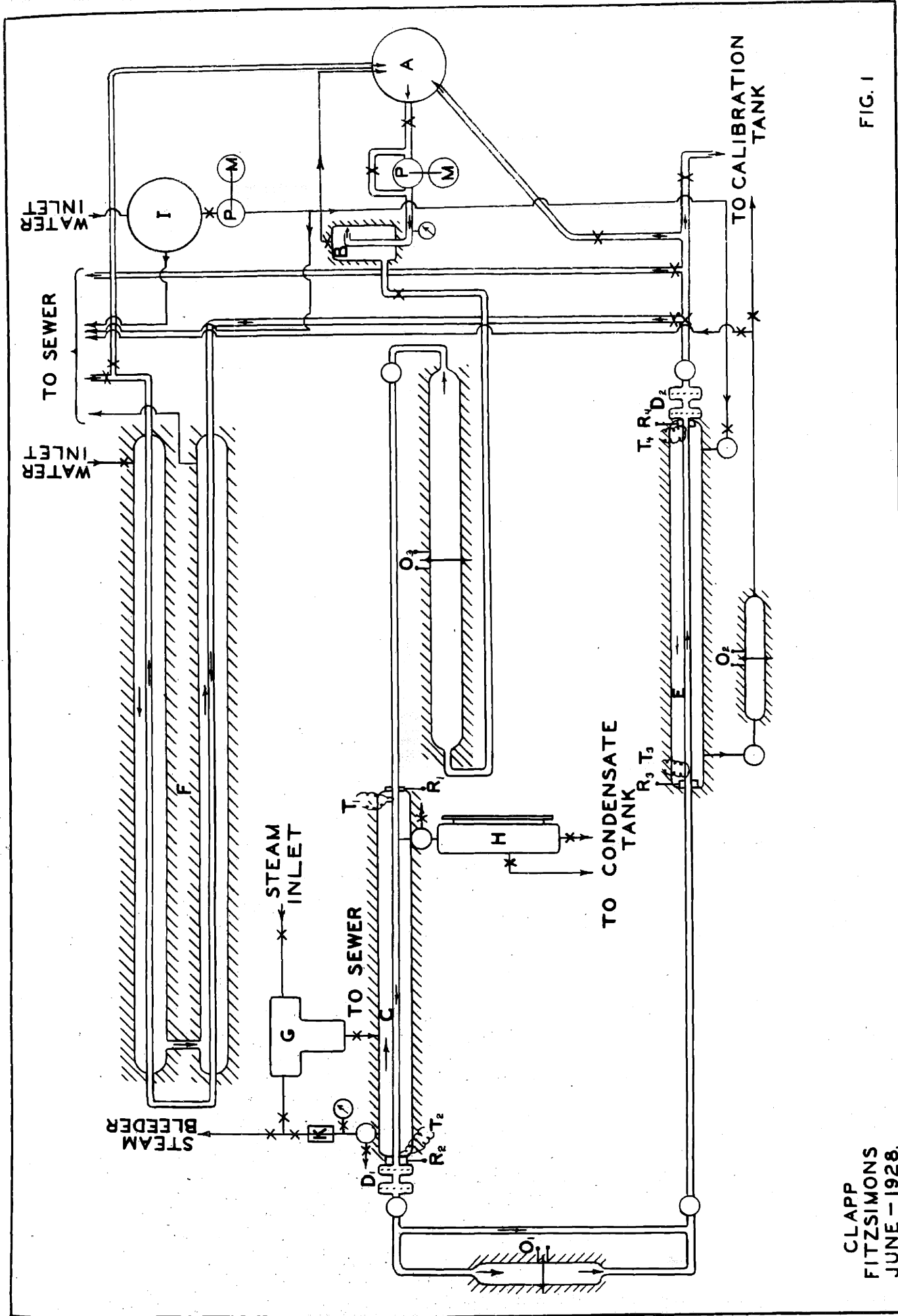
### APPARATUS AND METHOD

The apparatus used for the determination of friction factors and heat transfer coefficients is shown diagrammatically in Figure 1, while the important dimensions are given in Figure 2. It consists of two main units, (C) the heater and, (E) the cooler. The tube through which the fluid under consideration flowed was alike in both units. It was a standard 3/8-inch I. P. S. copper tube with internal and external diameters of 0.494 and 0.675 inches respectively. To allow any eddy currents, caused by bends, contractions or enlargements in the pipe, to become negligible; the fluid passed through a straight section of the tube over one hundred diameters long before entering either the heating or cooling zone. This "calming" section was a continuation of the standard copper tube. A turbulence chamber<sup>(D<sub>1</sub>-2)</sup> was placed at the exit of both units to thoroughly mix the stream coming from the exchanger and thus insure that the temperature read would be a true average of the stream. These turbulence chambers consisted of two slotted copper discs placed a short distance apart with the slots in one disc at right angles to those

DESCRIPTION FIGURE I

- A - Oil storage tank
- B - Air separator
- C - Heating unit
- D - Turbulence chambers
- E - Cooling unit
- F - Auxiliary cooler
- G - Steam separator
- H - Steam Condensate trap
- I - Cooling water storage tank
- K - Gas-fired superheater
- M - Motors
- O - Orifice leads
- P - Pumps
- R - Piezometer rings
- T - Thermocouples





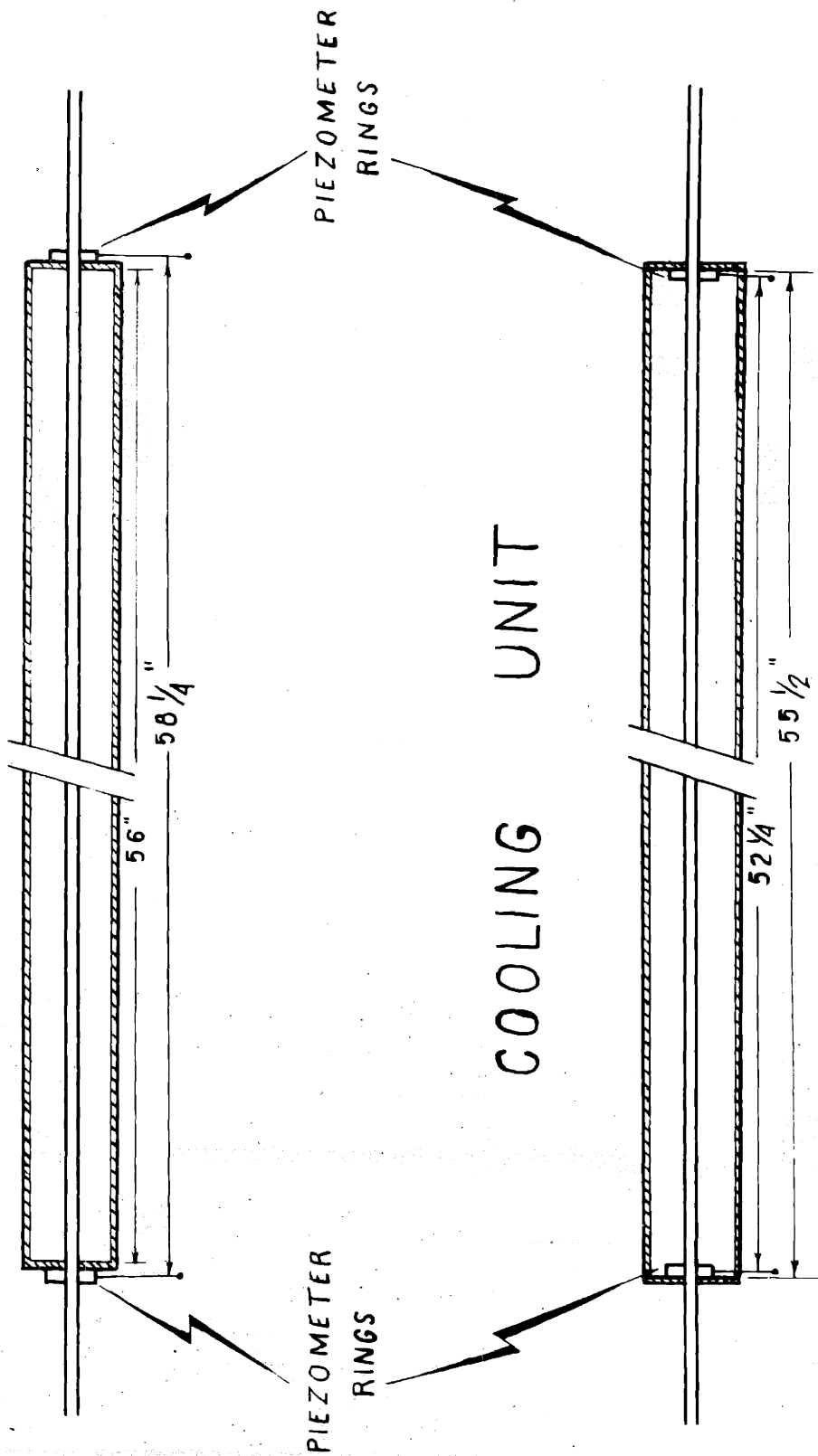
CLAPP  
 FITZSIMONS  
 JUNE - 1928.

FIG. I

CLAPP  
FITZSIMONS  
JUNE  
1916

HEATING UNIT

COOLING UNIT



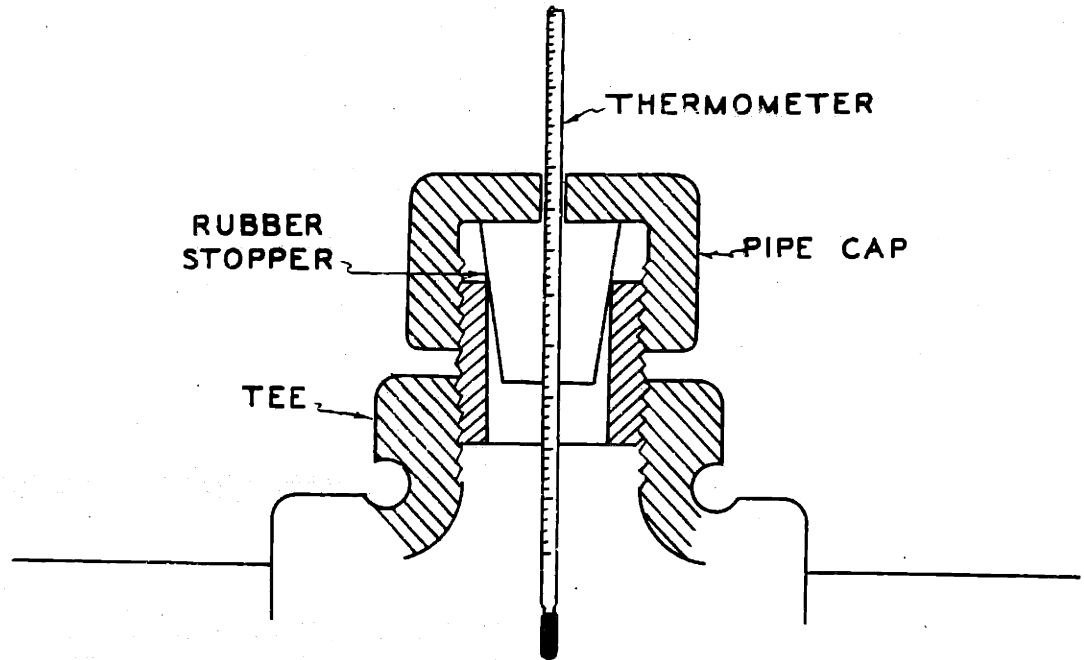
in the second. It was assumed that the changes in velocity and direction of the cooling water leaving the jacket would mix it sufficiently to give the same result. The exchangers and "calming" sections were all thoroughly lagged beyond all thermometer wells.

To allow room for the thermometer bulbs the 3/8-inch pipe was enlarged into a 1-inch tee to form the well. Each thermometer was fitted with a rubber stopper and held firmly in place by means of a pipe cap as shown in Figure 3A. All thermometers were compared with a set of German Standard thermometers P. T. R. 89859-60 (1922) and found to be sufficiently accurate for this investigation. Thermometers used in measuring the oil temperatures at the entrance and exit to bath units had a scale of 0-100°C in 17.5 inches, graduated in tenths making possible a reading to within 0.02 degrees. Thermometers used for measuring steam temperatures had a scale of 100-150°C in 13.5 inches, graduated in tenths. Those used for measuring cooling water temperatures had a scale of 0-30°C in 9.5 inches, graduated in tenths, making possible a reading to within 0.02 degrees.

Measurements of the static pressure were made in all cases with a specially designed piezometer ring,

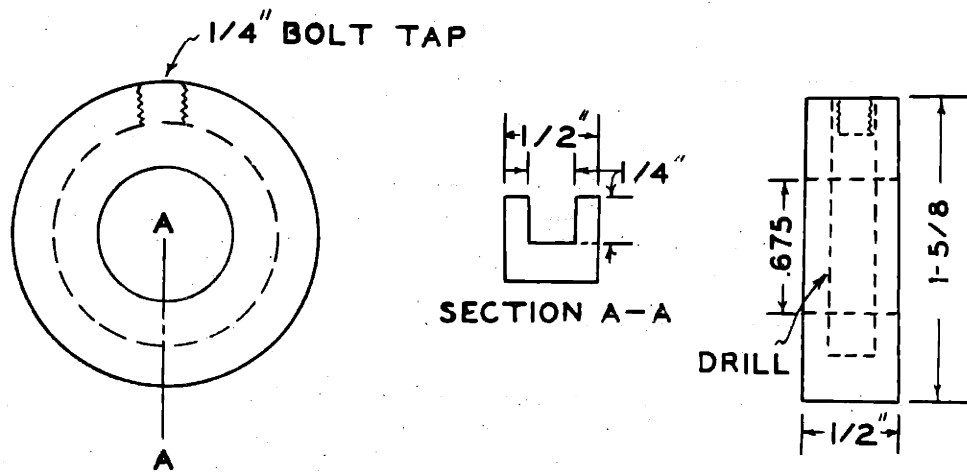
FIG. 3A

### THERMOMETER WELL



### PIEZOMETER RING

FIG. 3B



CLAPP  
FITZSIMONS  
JUNE - 1928.

Figure 3-B. In the heating unit, due to the high temperatures maintained in the jacket it was necessary to place the rings on the outside as close as possible to the entrance and exit of the jacket. In the cooler it was possible to place the rings on the inside. The length under consideration, therefore, varied slightly with the unit, that in the heater and cooler being 58.25 inches and 52.25 inches respectively. While the length exposed to the heating medium was 56 inches and that to the cooling medium was 55.5 inches.

The outside wall temperature of the pipe was measured at the entrance and exit of each unit by means of a No. 22 gauge copper-ideal thermocouple soldered into slots in the tube wall. A Leeds, Northrop student potentiometer was used, the scale of which was divided into tenths of a millivolt making possible a reading within five hundredths of one degree Centigrade. As a cold junction, ice in distilled water was used, and all thermocouples were calibrated in place.

As this is a continuous system, calibrated orifices were used in measuring the velocity of the fluid. For low velocities, ( $O_1$ ) orifice No. 1 was used and for high velocities ( $O_3$ ) orifice No. 3 was used. By means of a by-pass around No. 1 and screw clamps on the leads

running from the orifice to the instrument board the same manometer was used on both orifices. Orifice No. 1 consisted of chambers of standard 2-inch iron pipe and a 0.023-inch galvanized iron plate with the hole made by a 0.25-inch reamer. The static pressure openings were located at 1.6 and 0.8 inches from the plate on the upstream and downstream sides respectively. Orifice No. 3 consisted of chambers of standard 3-inch iron pipe and a 0.023 -inch galvanized iron plate with the hole made by a 0.375-inch reamer. The static pressure openings were located at  $2 \frac{6}{16}$  and  $1 \frac{3}{16}$  inches from the plate on the upstream and downstream sides respectively. ( $O_2$ ), Orifice No. 2, used for measuring the flow of the cooling water was identical with No. 1. All the static pressure outlets were of  $\frac{1}{4}$ -inch O. D. heavy walled brass tube and were situated at the bottom of the chambers to prevent any air collecting in them. By a proper arrangement of manometers it was possible to read the pressure drop across either the heater or the cooler, directly in inches of the fluid flowing for low drops and in inches of mercury for the higher values.

To keep a constant velocity of the cooling water through the jacket it was found necessary to fill a large storage tank from the water main and then pump the water

through the jacket with a Gould pump run by a constant speed motor, The amount of water through the jacket being regulated by means of a by-pass.

When making runs with water it was necessary to install an air separator (B) to remove as much air as possible from the water before running it through the heater. This was very effective with water but because of the greater viscosity of oil it was necessary to submerge the circulating pump in a bath of oil to prevent air getting into the system and then the oil being circulated was heated to and held at a high temperature until practically all of the dissolved air had escaped from the storage tank. The circulating pump was a Kinney rotary pump run by a constant speed motor. The amount of fluid through the apparatus was regulated by means of a by-pass.

For low temperature runs an auxiliary cooler (F) was used to keep the oil at a constant temperature but for the high temperature runs the fluid flowing was shunted directly to the storage tank after leaving the cooling unit without passing through the auxiliary cooler.

In order to be sure of a supply of dry saturated steam in the heating unit all steam lines were

thoroughly logged and enough steam was bled from the line to insure the operation of the steam separator. Also a small gas-fired super-heater was installed between the separator and the heating unit and operated so as to furnish dry saturated steam in the runs in which heat balances were desired. By means of an expansion valve the steam pressure could be varied from atmospheric to 50-lb. gauge. Heat balances were made, in all runs using water, as a check on the method and were found to be within a very few percent in all cases, v. Table I, showing the heat lost to the surroundings to be negligible. It was necessary to collect the condensate formed in the heating unit in cold water to prevent flashing. In the case of the runs made with oil, because of the poor coefficients less condensate was formed. While this tends to make the percentage lost slightly larger it may still be assumed negligible. And as the length of each run was quite short the amount of condensate formed could not be measured with any accuracy and no attempt was made to figure heat balances on these runs. Likewise, because of the large amount of cooling water used its rise in temperature was too small to justify the calculation of heat balances on this unit.

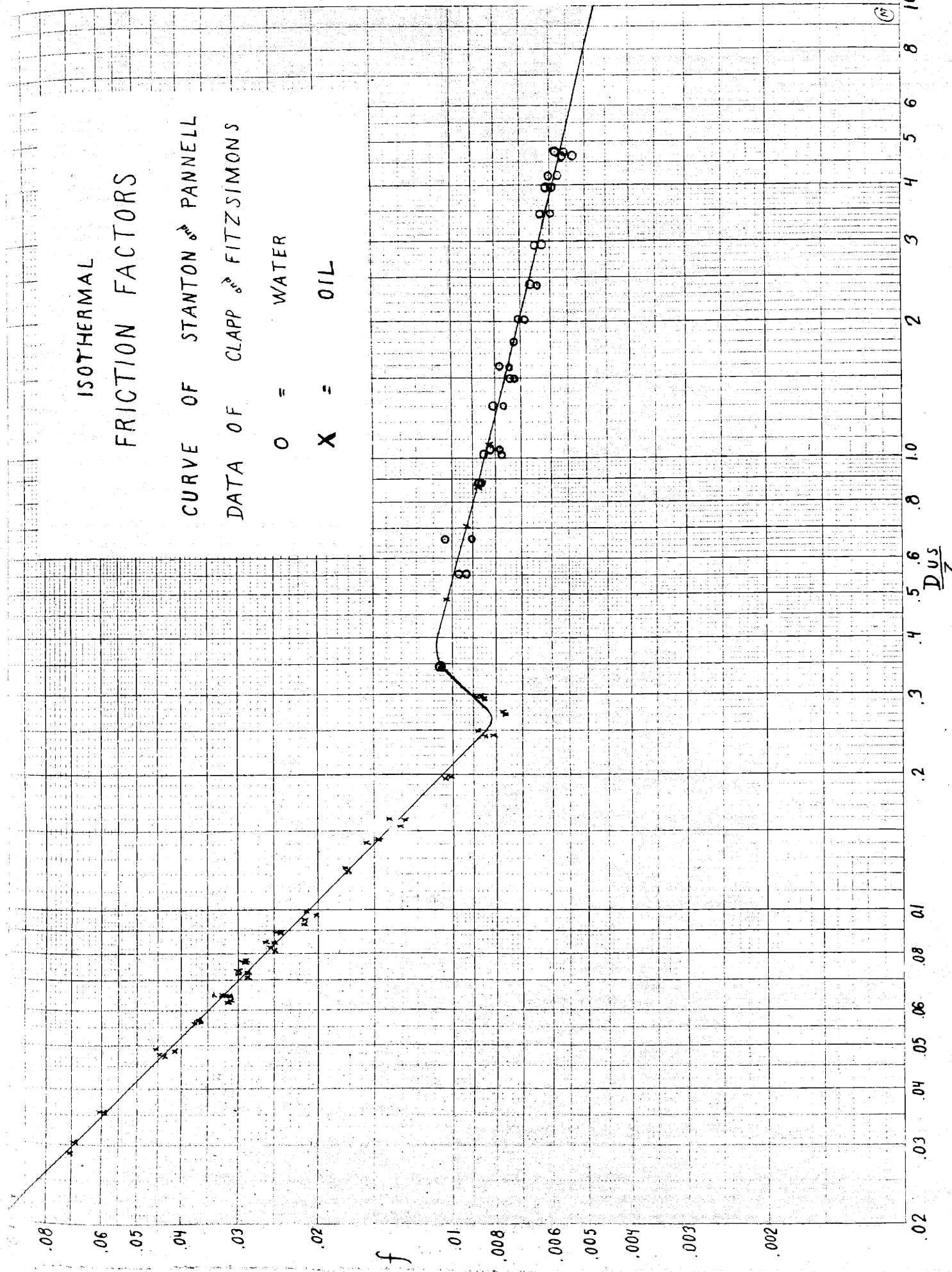


The standard tubes in both main units were cleaned each day while using water as a fluid. Also before each run it was necessary to remove any trace of air that collected in the manometer arms. When using water direct from the mains and discharging it without recirculation quite a bit of air collected in these arms, but when using oil with the circulating pump submerged this was eliminated.

DISCUSSION OF RESULTS

Very little information was available on this problem. The data taken by White<sup>(2)</sup> have been recalculated by Spurdle<sup>(4)</sup> and seem to indicate that the actual pressure drop is less than the value predicted by the Fanning equation, when the fluid is being heated. And there were no data available on the actual value of the pressure drop during cooling. Saph and Schoder have done a large amount of work in establishing a relation between the friction factor ( $f$ ) and the modulus  $\frac{(DuS)}{Z}$  for isothermal conditions. This work is summarized in a plot by Spurdle. Stanton and Pannell<sup>(6)</sup> have also done considerable work on isothermal conditions and their results agree with the curve established by Saph and Schoder. Therefore this curve may be considered a correct representation of the conditions existing during isothermal flow. In this investigation a number of isothermal runs were made with both oil and water and it was found that the friction factors calculated from these agreed very nicely with the established curve over its entire length. Figure 4. Thus showing that the experimental method and apparatus would produce results in accord with previous work. These runs were made before

ISOTHERMAL  
FRICTION FACTORS  
CURVE OF STANTON & PANNELL  
DATA OF CLAPP & FITZSIMONS  
O = WATER  
X = OIL



the pump was immersed in a bath of the liquid under consideration, to eliminate the last traces of air from the system. Hence it is thought that results could be obtained after the pump was immersed that would agree even more closely with Stanton and Pannell's curve. Velocite "B" was the oil used in all oil runs.

The results obtained seem to show that the pressure drop is dependent to quite an extent on the viscosity of the film sliding along the pipe wall. That is, for the heating of fluids this film is much hotter than the main stream and therefore has a greater fluidity, causing the main stream to move through the pipes as a slug slipping over the less viscous film, with an actual pressure drop less than that predicted by the Fanning equation in which the viscosity of the main body of the fluid is used in determining the friction factor ( $f$ ). While for the cooling of fluids the opposite takes place, the film being much colder than the main stream has a greater viscosity and retards the flow, causing a greater pressure drop than predicted.

Two methods of allowing for this film effect were proposed. First, since the average viscosity of the main body does not give correct results, by choosing some other viscosity such as the film viscosity, or per-

haps some mean between the two it would still be possible to use the Fanning equation with the friction factors for isothermal flow. But this method presents one big disadvantage. In order to determine the temperature of the film, trial and error work is necessary to find the correct pipe temperature to use.

The second method, one which is more desirable in engineering practice because of its simplicity, is as follows. Since the viscosity is a function of the temperature and this film effect is dependent on the difference between the viscosity of the main body and that of the film, it is obvious that it should be possible to allow for the film effect by taking into consideration the temperature difference between the average temperature of the main body of the fluid being heated or cooled and the inside pipe wall temperature. As the velocity of a fluid is increased, however, the type of flow changes from viscous to turbulent. In the turbulent range, while there is still a thin film slipping along the pipe wall, it is very much thinner and hence the effect referred to above is evident to a less extent. As the turbulence is made more violent at still higher velocities it wears this film down extremely thin and the film effect becomes negligible.

Any attempt to express this varying effect of the temperature difference in mathematical form would give rise to quite complicated equations. Therefore, it was deemed advisable to allow for it by constructing a series of curves showing the friction factor which, when substituted in the Fanning equation will give the actual pressure drop for the various temperature differences. These curves are plotted from the data obtained in runs H<sub>1</sub>-66, v. Table II. In the viscous range, they appear to be nearly parallel to the isothermal curve while in the turbulent range they approach it as predicted. Since the properties of the fluid are taken at the average temperature of the main body of the fluid, this set of curves present a simple solution to the problem, one quite easy to use in calculations for design work.

HEAT TRANSFER COEFFICIENTS

As the main object of this thesis was the investigation of the pressure drop, the conditions under which some runs were made did not make the data collected suitable for calculating heat transfer coefficients. Therefore, for heating, only such runs that had a rise in the temperature of the main body of the fluid of over  $2.5^{\circ}\text{C}$ , in flowing through the tube, were used. And for cooling only those having a fall in temperature of over  $1^{\circ}\text{C}$  were used.

The coefficients were calculated and plotted as suggested by Morris and Whitman, Figure 6. While all of the points for both the heating and cooling of oil in turbulent range seem to be a bit lower than their curve, they are on the whole in quite good agreement. For heating water, however, the points seem to be much lower than their suggested curve. As there was no apparent reason for this, there seemed to be a question as to where the curve should be drawn. Some data obtained by Baldwin and Sherwood<sup>(7)</sup> in 1924 on the heating of water were recalculated and plotted on the same coordinates. These points also fall some distance below

the curve of Morris and Whitman and seem to be in good agreement with the results obtained in this investigation. Hence it is suggested that the slope of the curve be reduced in this range.

The data obtained on the cooling of water are all that are available at this time. They seem to indicate that the coefficients for cooling water are approximately the same as for heating it.

The calculations of these coefficients are given in Table V.



### CONCLUSION

Friction factors for both heating and cooling of fluids in the viscous and turbulent ranges have been determined. The results seem to be correlated best by a series of curves above and below the established isothermal curve and symmetrical with it.

Heat transfer coefficients for oil and water have been determined. The results on oil seem to agree with the curve of Morris and Whitman. In their article they propose one curve for heating and that value, 75 per cent of these be used for cooling. The results obtained for water, however, are much lower. Results of Baldwin and Sherwood on water are also lower, indicating that the curve proposed by Morris and Whitman should have a flatter slope in this range.

Figure 5 shows  $(f)$  plotted against  $\frac{(DuS)}{z}$  for isothermal, heating and cooling conditions.

Figure 6 shows  $\frac{(hD)}{k} / \left(\frac{cz}{k}\right)^{0.37}$  plotted against  $\frac{(DuS)}{z}$  for heating and cooling. The curve shown is Morris and Whitman's for heating, and cooling is 75 per cent of the heating value.

Figure 7 shows White's data plotted on the curves adopted in this investigation. All of his values seem to run low, but as this is also the case with his isothermal runs some of the deviation is obviously due to errors in his experimental method. Because of this it seems advisable to use the curves as drawn and any errors resulting from this will simply increase the factor of safety.

Fig. 5

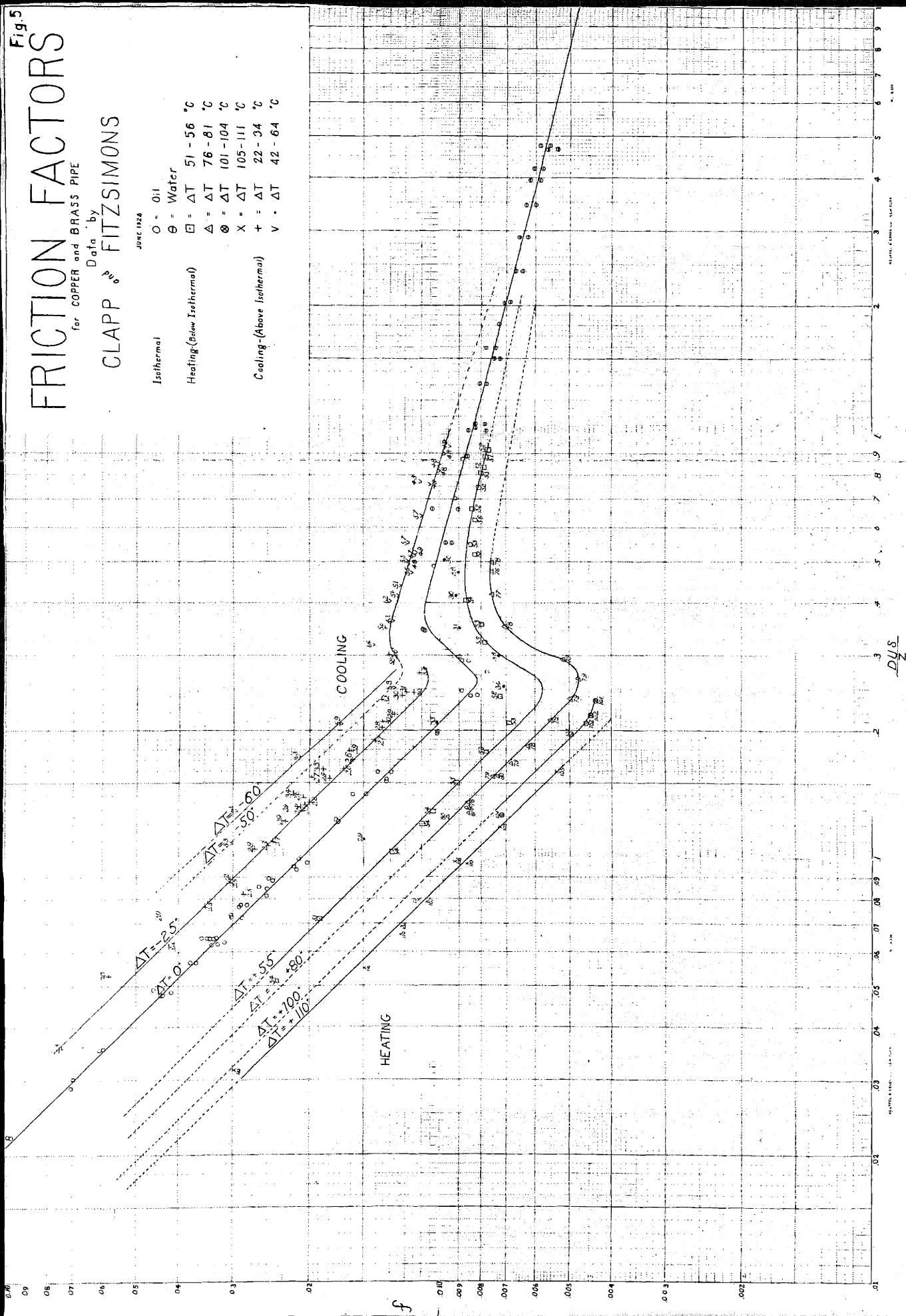
# FRICITION FACTORS

for COPPER and BRASS PIPE

Data by  
**CLAPP and FITZSIMONS**

JUNE 1926

- = 0.1
  - ⊖ = Water
  - = ΔT 51 - 56 °C
  - △ = ΔT 76 - 81 °C
  - ⊙ = ΔT 101 - 104 °C
  - × = ΔT 105 - 111 °C
  - + = ΔT 22 - 34 °C
  - ∇ = ΔT 42 - 64 °C
- Isenthal
- Heating (Below Isenthal)
- Cooling (Above Isenthal)



10-1111

REYNOLDS NUMBER

$DVs$

10-1111

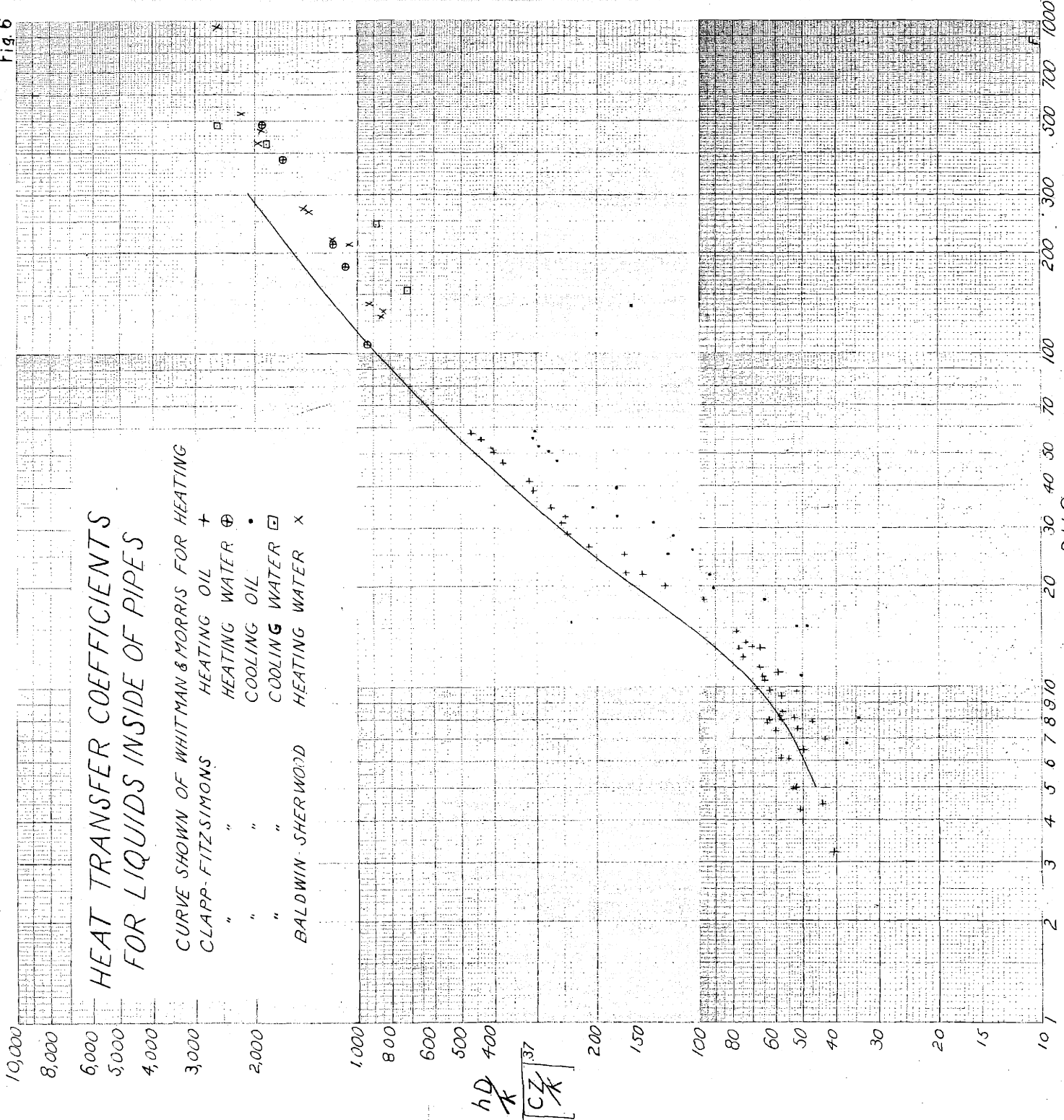
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Fig. 6

# HEAT TRANSFER COEFFICIENTS FOR LIQUIDS INSIDE OF PIPES

CURVE SHOWN OF WHITMAN & MORRIS FOR HEATING  
 CLAPP-FITZSIMONS HEATING OIL +  
 " " HEATING WATER ⊕  
 " " COOLING OIL •  
 " " COOLING WATER ⊠  
 BALDWIN-SHERWOOD HEATING WATER x



$$\frac{hD}{k} = \frac{CZ}{k}^{0.37}$$

$$DVU = \frac{DUP}{Z}$$

Fig. 7

# FRICITION FACTORS

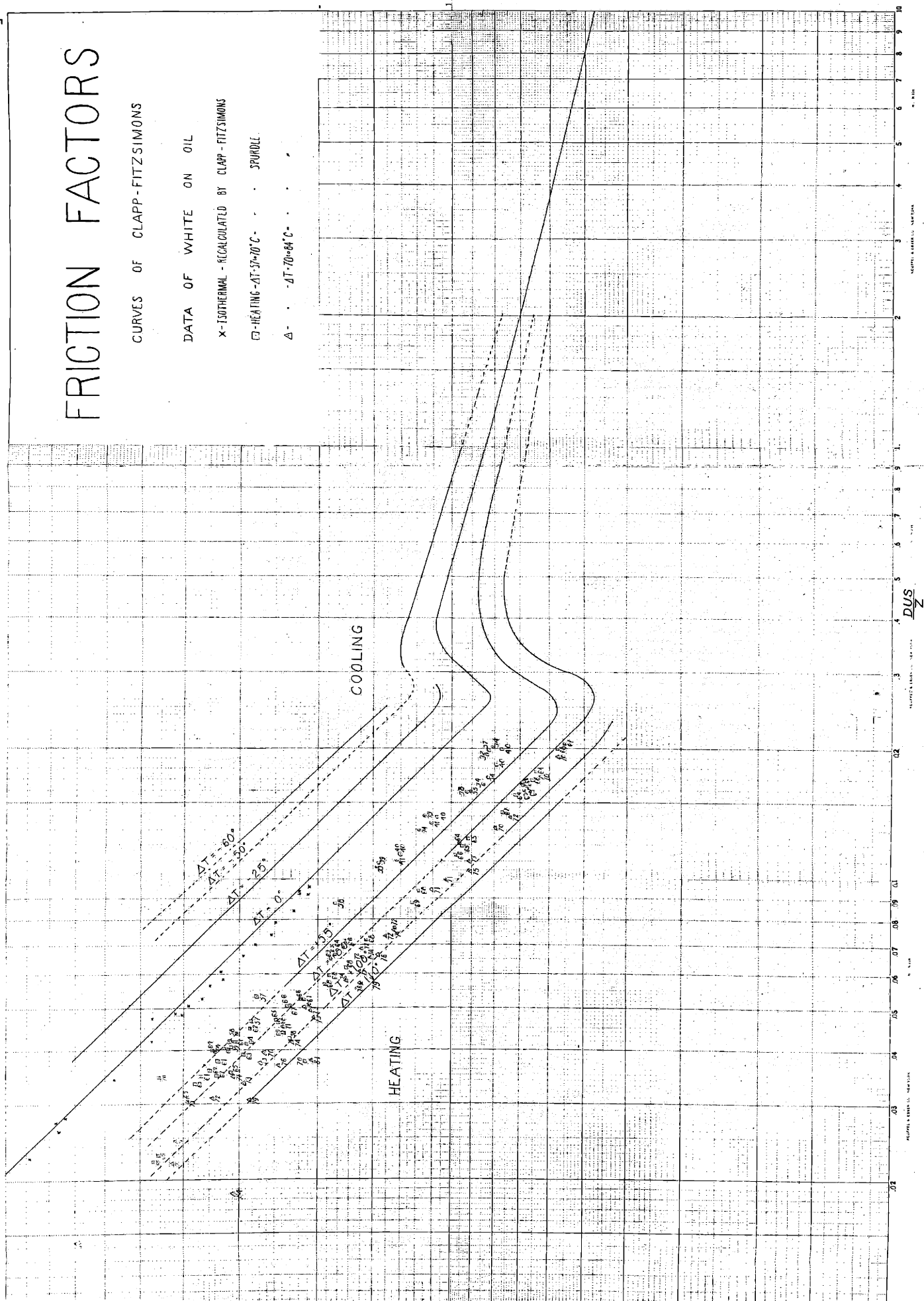
CURVES OF CLAPP-FITZSIMONS

DATA OF WHITE ON OIL

X-ISOTHERMAL - RECALCULATED BY CLAPP-FITZSIMONS

□ - HEATING -  $\Delta T = 10^{\circ}C$  . . . SPURDILL

Δ - . . .  $\Delta T = 70^{\circ}C$  . . .



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RECOMMENDATIONS

For future work on this problem it would be well to equip the storage tank, for the oil, with a steam coil to maintain constant temperature at the high levels.

The circulating pump should be kept immersed in the fluid under consideration.

It would be desirable to have more turbulent flow of the cooling water in the jacket. This could be obtained in two ways. First, by enlarging the orifice to permit a greater velocity. This would require a large amount of cooling water and make it practically impossible to obtain heat balances on this unit. Second, by reducing the cross sectional area of the annular space in the jacket. This would lead to difficulties in construction.

An attempt should be made to use a wider range of temperature differences in the cooling unit.

Also it is advised that a study be made of the effect of a high temperature difference in the turbulent range for bath units. To do this, it would be necessary to use an oil with a lower viscosity than the one studied here.

APPENDIX A



NOMENCLATURE

- $A_c$  = Area of cooling surface  
 $A_H$  = Area of heating surface  
 $c$  = Specific heat  
 $D$  = Diameter of the pipe in inches  
 $f$  = Friction factor in the Fanning's Equation  
 $g$  = Gravitational constant = 32.2 ft/sec<sup>2</sup>  
 $h$  = Heat transfer coefficient for fluids expressed in B. t. u. per hour per sq. ft. per degree F temperature difference  
 $\Delta H$  = Liquid head in inches  
 $J$  = Fluidity =  $\frac{1}{Z}$   
 $k$  = Thermal conductivity; for oil a constant = 0.078  
 $L$  = Length of pipe in inches  
 $Q/\theta$  = C. h. u. or B. t. u. per hour  
 $R_o$  = Reading of Orifice manometer in inches of mercury in contact with oil  
 $S$  = Specific gravity  
 $S_o$  = " " at orifice  
 $S_T$  = " " at average temperature of fluid in pipe  
 $T_o$  = Temperature of fluid at orifice °C  
 $T_T$  = Average temperature of fluid in tube °C.  
 $T_1$  = Temperature of steam entering jacket  
 $T_2$  = " " " leaving "

- $T_3$  = Temperature of cooling water entering jacket  
 $T_4$  = " " " " leaving "  
 $t_1$  = " " fluid entering heater  
 $t_2$  = " " " leaving "  
 $t_3$  = " " " entering cooler  
 $t_4$  = " " " leaving cooler  
 $t'$  = Average temperature of fluid in tube °F  
 $\Delta T$  = The effective temperature difference through the fluid film  
 $u$  = Velocity in feet per second  
 $Z$  = Viscosity of fluid in centipoises relation to water  
 $\rho$  = density of fluid = 62.3S

CALIBRATIONS

The diameter of the standard tubes was measured by filling each tube with water, the ends being carefully sealed and weighing the amount of water. In bath cases they were as specified 0.494 inches I. D.

All orifices were calibrated by weighing the amount of water pumped through at a constant velocity in a given period of time. This was done at several different velocities for each orifice. When this data was plotted it was found to be within two per cent of the curve calculated from the theoretical equation for sharp-edged orifices. Figure 8, Table IV.

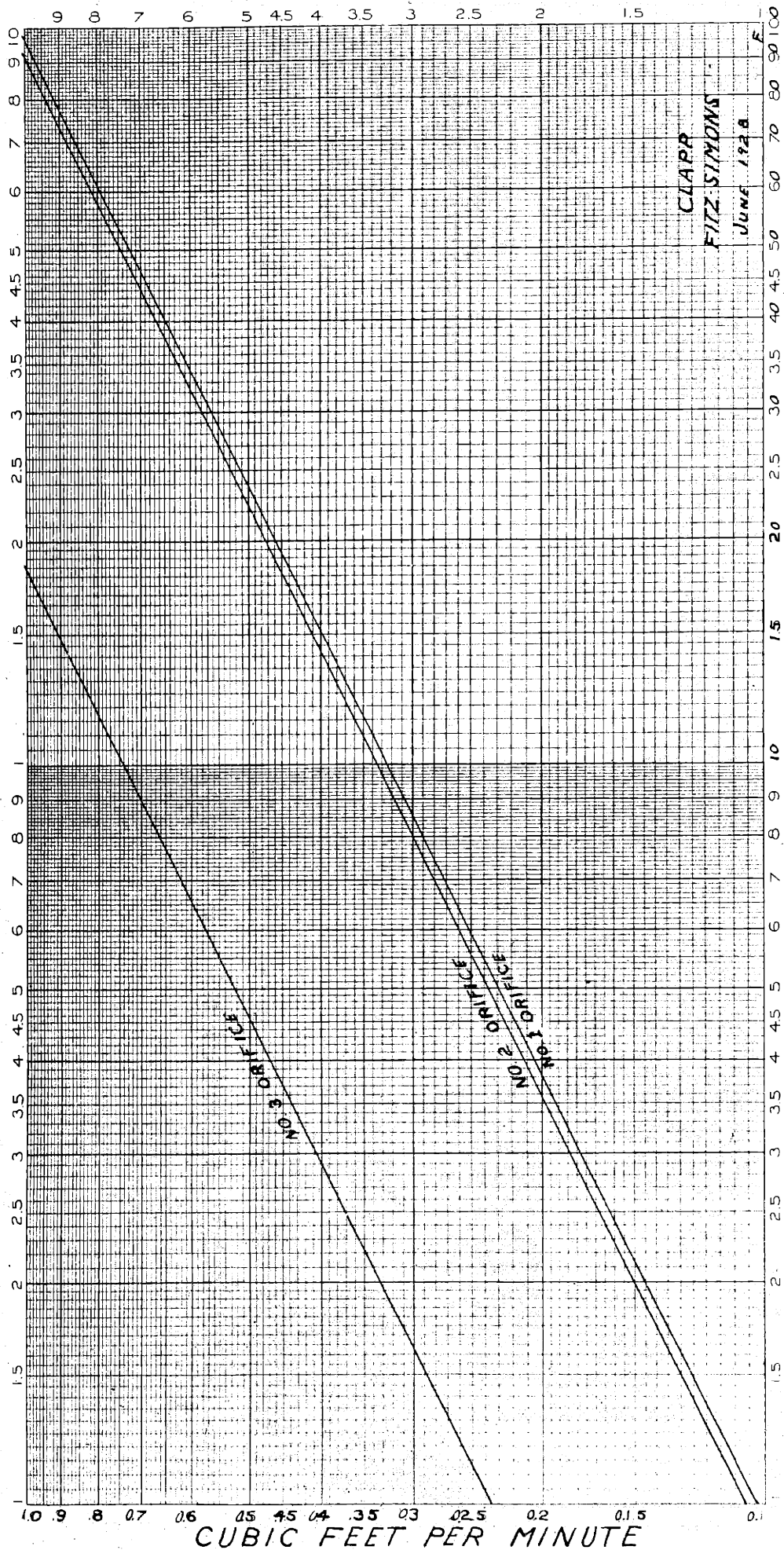
Thermocouples in the cooling unit were calibrated by running water at constant temperature through the system until equilibrium was reached in the unit. This was done at several different temperature levels. Thermocouples in the heating unit were calibrated by admitting steam, at different pressures, into the jacket and taking the reading of the pressure by a calibrated gauge after equilibrium was established between the pipe and jacket, Figure 9.

It was necessary to keep the pressure at the discharge end of the pump below 40 lb. gauge, as this

was the maximum pressure the circulating system would withstand. All pressure gauges used were calibrated with an Ashton Tester in the Mechanical Engineering Department.

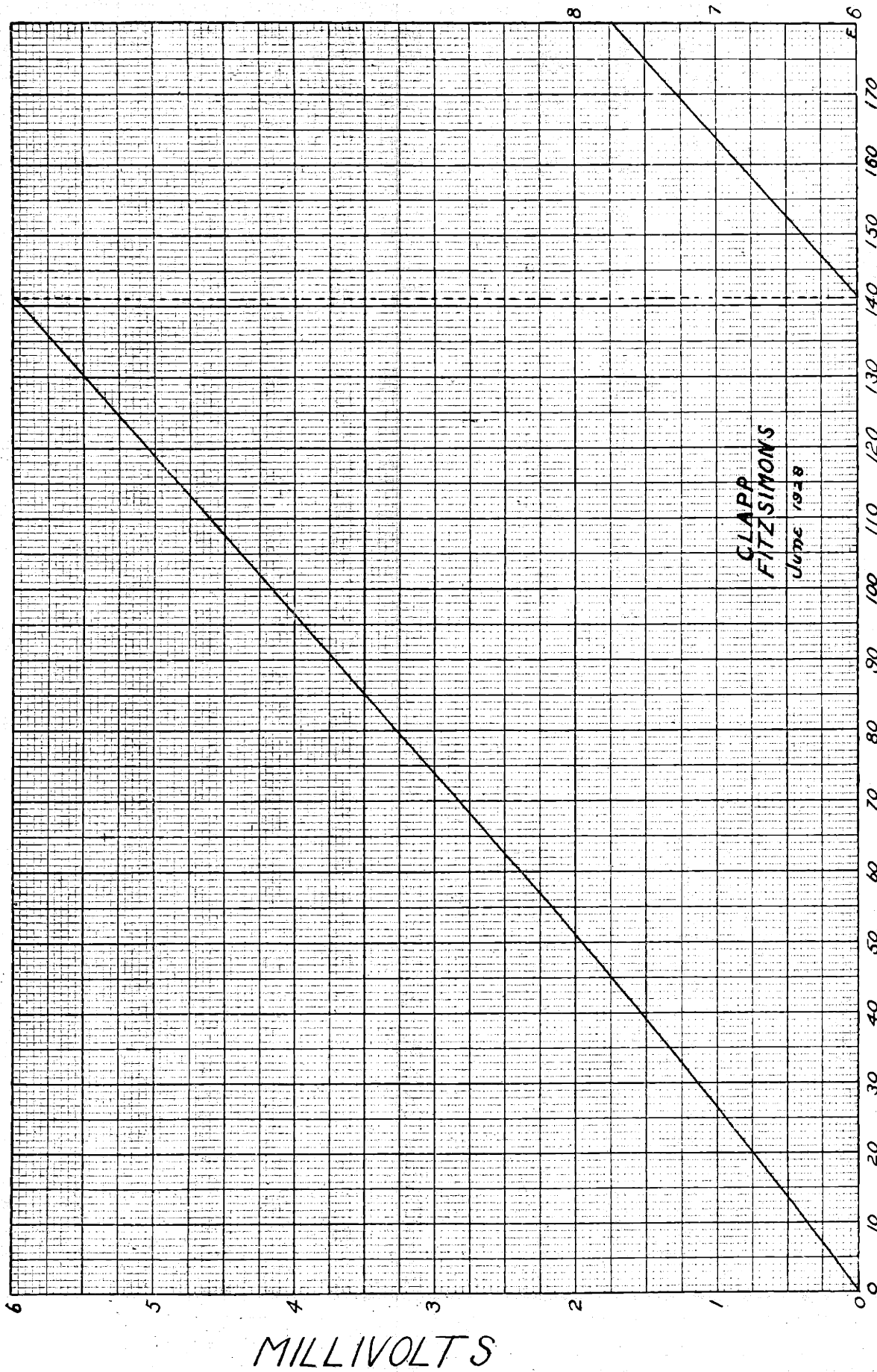
Due to a change in design it was necessary to replace the standard copper tube in the heater with a similar one. This was done between runs C<sub>8</sub> and D<sub>1</sub>.

Fig. 8



ORIFICE CURVES

Fig. 9



CLAPP  
FITZSIMONS  
June 1928

TEMPERATURE - DEG. CENTIGRADE  
THERMOCOUPLE CALIBRATION CURVE

DESCRIPTION OF RUNS

- Series A - Calibration of Orifice No. 1 and No. 3  
" B - " " " " No. 2  
" C - Heating and Cooling of water  
" D - Isothermal on Water  
" E - " " Oil (Velocite "B")  
" H - Heating and Cooling of Oil (Velocite "B")

The standard tubes under observation were thoroughly cleaned at regular intervals which using water, between runs C<sub>4</sub> and C<sub>5</sub>; C<sub>6</sub> and C<sub>7</sub>; D<sub>11</sub> - D<sub>12</sub>. Also they were cleaned again after all the water had been removed from the apparatus, before refilling it with Velocite "B".

COMPUTATIONS

It was assumed that the mercury in the manometers and the oil or water in contact with it was at an average temperature of 25° C., at which temperature the specific gravity of mercury is 13.53 and that of oil 0.89.

$$\text{Velocity in tube} = \frac{(\text{cu. ft. per sec. at orifice}) (\text{S.g. at orifice})}{(\text{cross section area of pipe}) (\text{S.gr. at ave. temp. of fluid})}$$

$$u = \frac{\text{ft.}^3/\text{min.}}{(60) \frac{(\pi) (0.494)^2}{(4) (144)}} \times \frac{S_o}{S_T} = \frac{1}{0.0795} \times \text{ft.}^3/\text{min.} \times \frac{S_o}{S_T}$$

$$f_{\text{(for new heater)}} = \frac{gD \Delta H}{24 L u^2} = \frac{(32.2)(0.494)}{(24)(58.25)} \times \frac{\Delta H}{u^2} = 0.01137 \frac{\Delta H}{u^2}$$

$$f_{\text{(for cooler)}} = \frac{(32.2)(0.494)}{(24)(52.25)} \times \frac{\Delta H}{u^2} = 0.01267 \frac{\Delta H}{u^2}$$

$$f_{\text{(for old heater)}} = \frac{(32.2)(0.494)}{(24)(57.75)} \times \frac{\Delta H}{u^2} = 0.01267 \frac{\Delta H}{u^2}$$

$$A_H = \frac{(56) (\pi)}{144} \times 0.494 = 0.603 \text{ sq. ft.}$$

$$A_C = \frac{(55.5) (\pi)}{144} \times 0.494 = 0.598 \text{ sq. ft.}$$



To convert inches of mercury to inches of oil.

$$\text{Inches of Hg.} \times \frac{13.53 - 0.89}{S_T} = \text{inches of Hg.} \times \frac{12.64}{S_T}$$

To convert inches of cold oil in manometer to inches of oil at the average temperature of the fluid in the tube

$$\Delta H \text{ (hot oil)} = \Delta H \text{ (cold oil)} \frac{0.89}{S_T}$$

The orifice curves are plotted cubic feet per minute vs. inches of mercury in contact with water.

To convert reading of mercury in contact with oil to inches of mercury in contact with water

$$R_o \times \frac{13.53 - 0.89}{S_T} \times \frac{1}{13.53 - 1} \times \frac{S_T}{S_o} = R_o \times \frac{1.008}{S_o}$$

$$c = \left( \frac{t' + 670}{2030} \right) (2.1 + \text{sp.gr. at } 60^\circ \text{ F.})$$

$$c = \frac{0.593}{10^3} (t' + 670)$$

## FLUID PROPERTIES

### SPECIFIC HEAT - Fig. 10

For oil the specific heat was calculated by means of the Whitman and Fortsch<sup>(8)</sup> equation.

### SPECIFIC GRAVITY - Fig. 11

The specific gravity of the oil was determined by means of a specific gravity bottle.

### THERMAL CONDUCTIVITY - Fig. 12

The Morris and Whitman procedure was followed in determining the modulus of heat transfer coefficients and their value of  $(k)$  for oil was taken,  $k = 0.078$ . For water, however, the Jakob<sup>(9)</sup> equation was used.

### VISCOSITY - Fig. 13.

The viscosity of Velocite "B" was determined by means of a Saybolt viscosimeter over the entire range from 0-100°C. The thermometer used to determine these temperatures was that used to measure  $t_1$ .

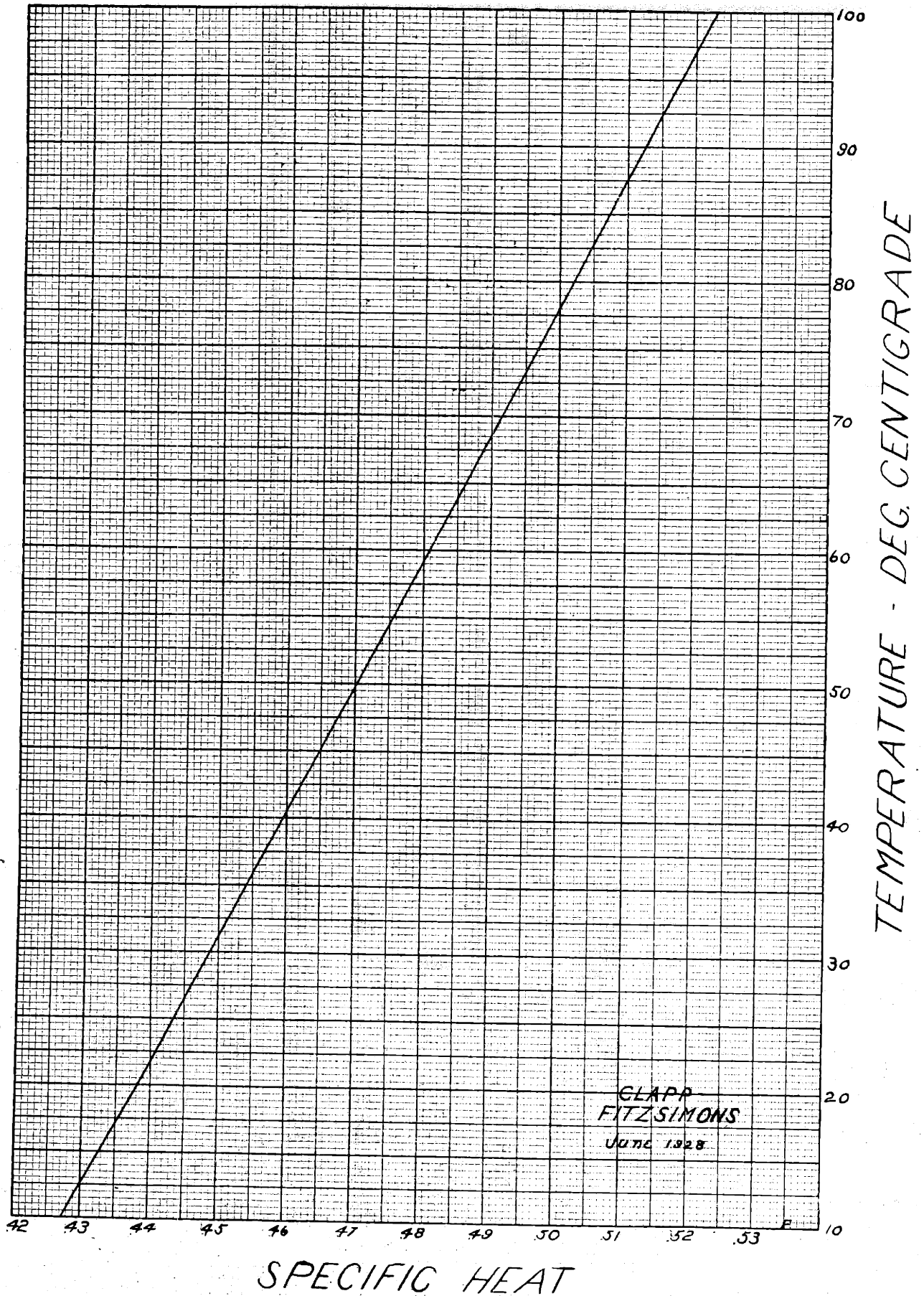
### EFFECTIVE TEMPERATURE DIFFERENCE

In the runs made with water where the change in temperature difference was relatively great a logarithmic mean value was used. The change in the

temperature differences when oil was used was relatively small so that the arithmetic mean was used. Since a standard copper tube was used the inside pipe wall temperature may be assumed to be the same as the outside wall temperature. In the cooling jacket the arithmetic mean of the two end thermocouples was used when the ratio of the two was less than 2, while for other cases the logarithmic mean was used.

The specific gravity, viscosity and fluidity of water was taken from E. C. Bingham's<sup>(10)</sup> book on "Fluidity and Plasticity."

Fig. 10



CALCULATED FROM WHITMAN-FORTSCH EQUATION [IND. ENG. CHEM. 18, 795]

Fig. 11

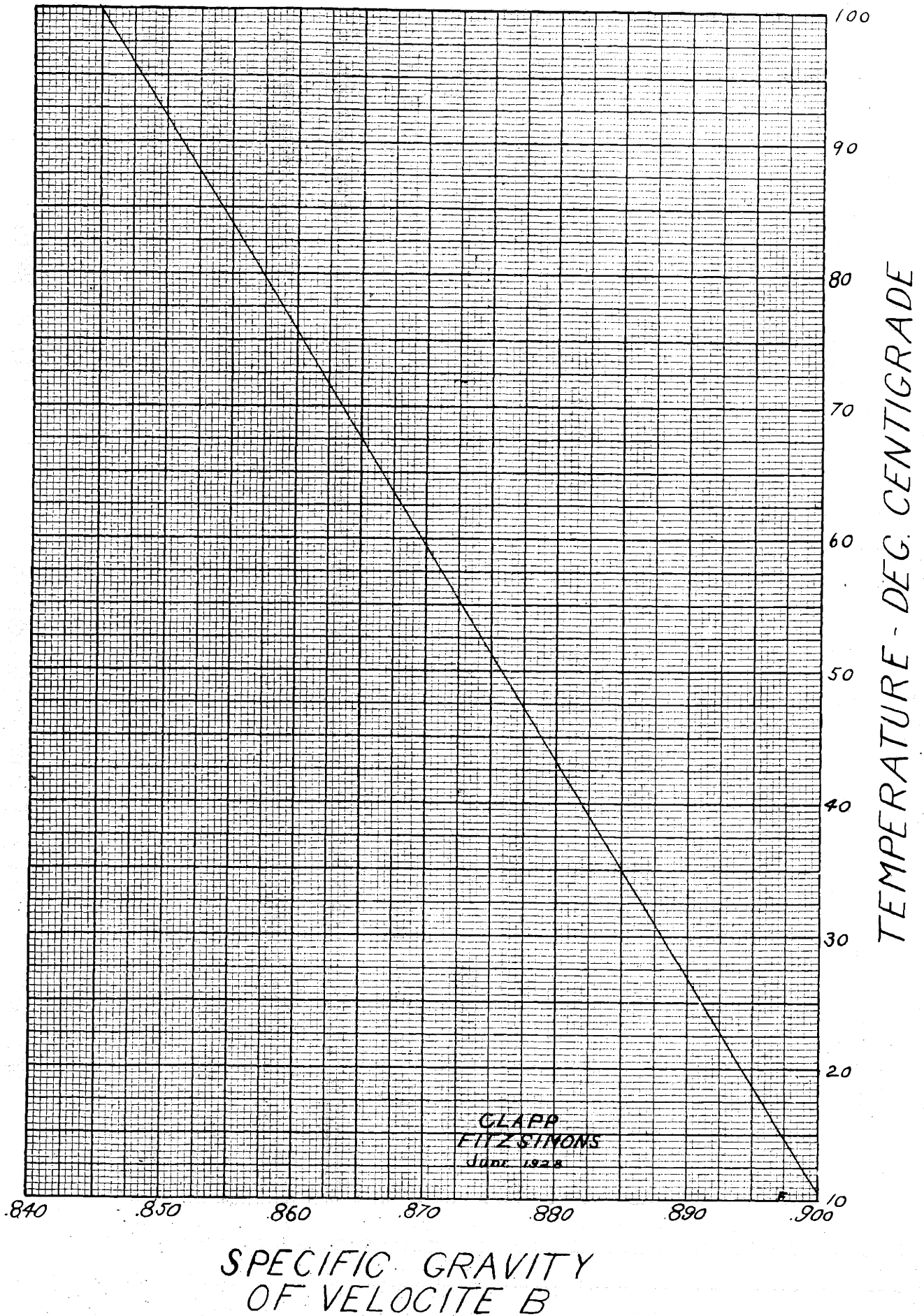
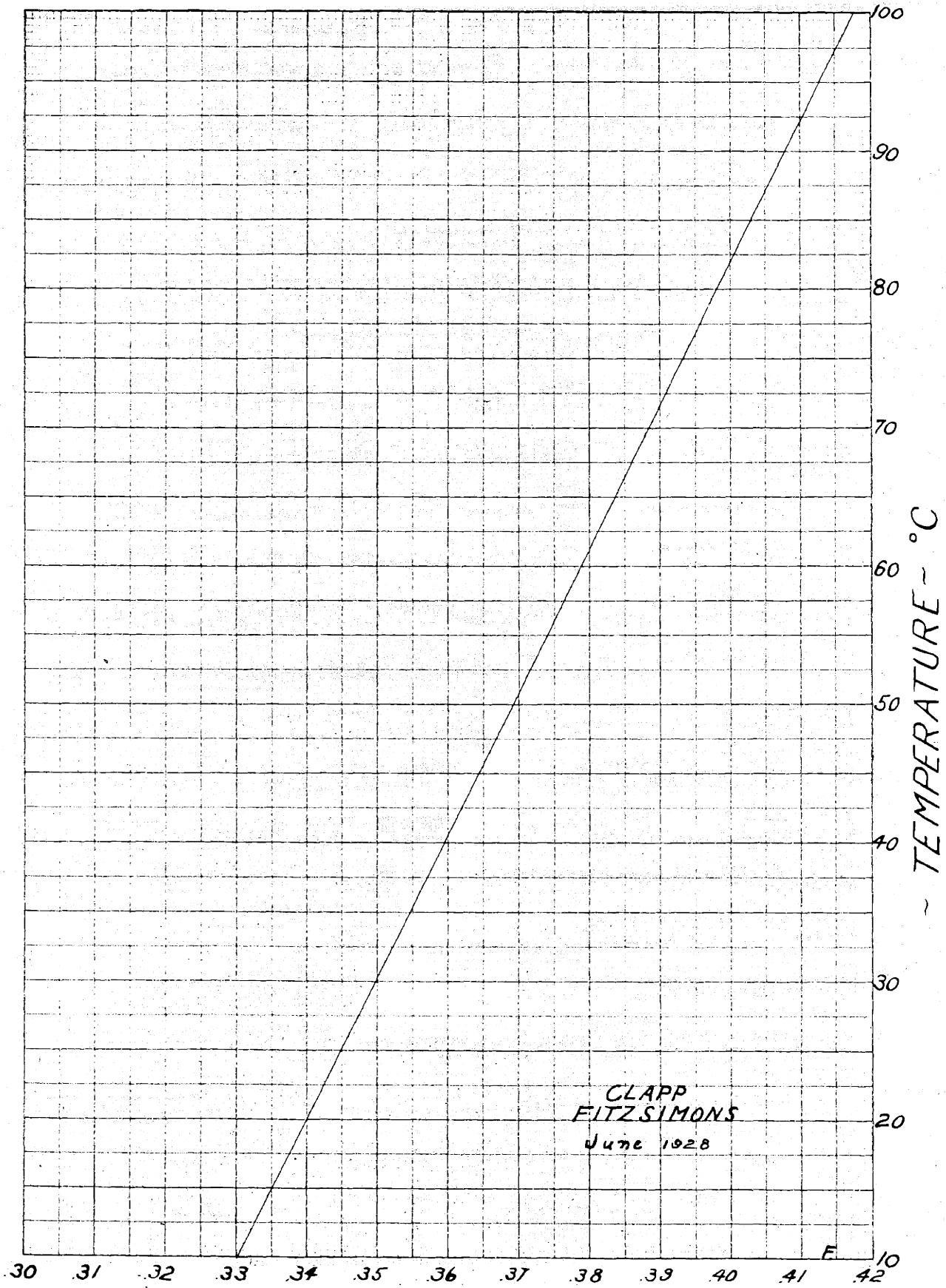


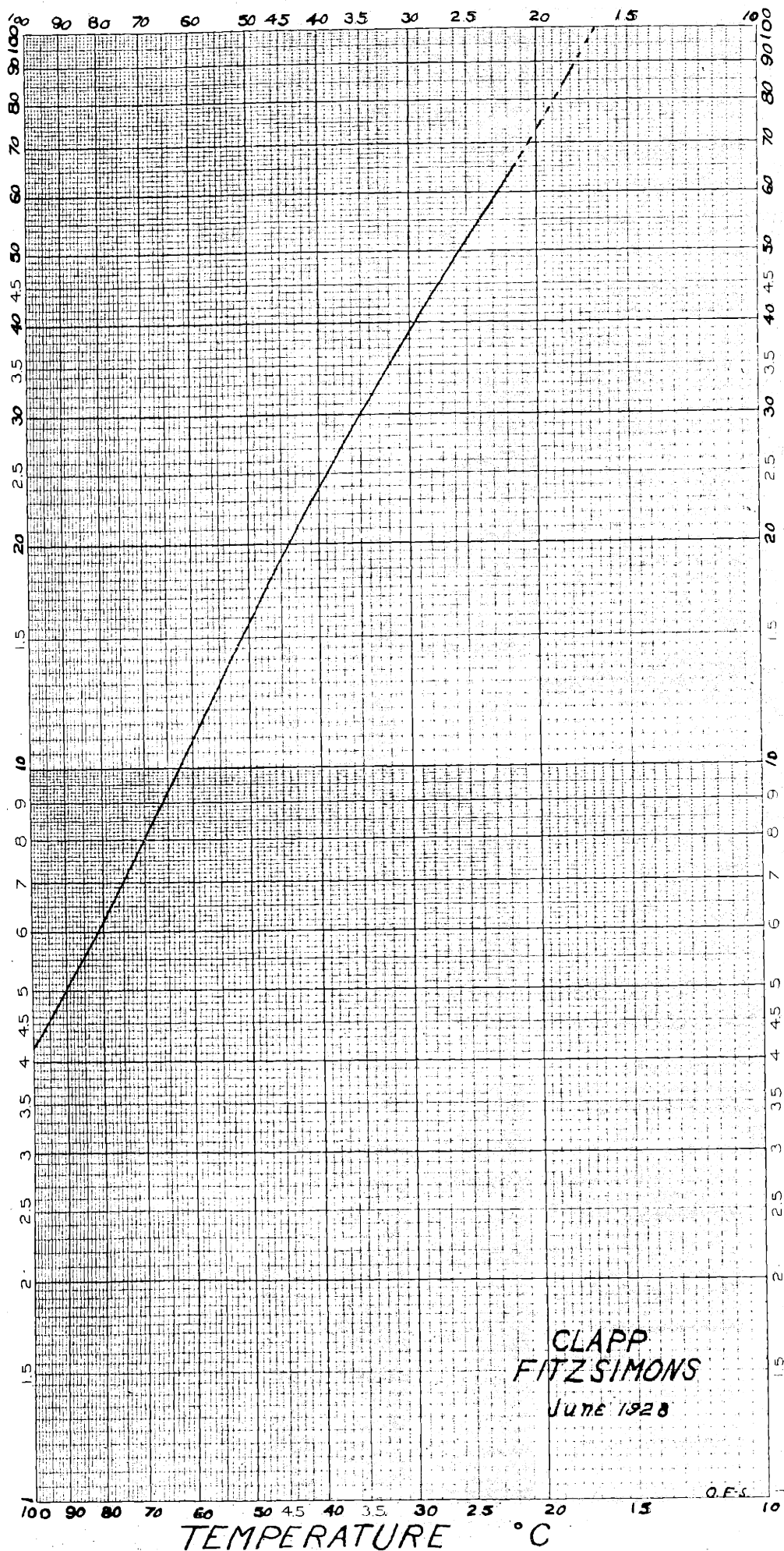
Fig. 12



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Fig. 13



Z = VISCOSITY OF VELOCITE B IN CENTIPOISES

TEMPERATURE °C

APPENDIX B



HEAT BALANCESHeater

Run No.	lb./min fluid	Temp. Rise °C	B. t. u. gained by fluid	lb. of steam	Latent Heat	B. t. u. lost by steam	B. t. u. unaccounted for	%
02	13.0	52.35	26,150	16.0	955.3	15,300	-10,850	71.0
3	23.6	36.30	25,350	23.5	"	22,450	+ 900	4.0
4	26.2	23.00	13,070	13.94	"	13,310	+ 240	1.8
5	48.0	27.95	28,900	29.94	"	28,600	+ 300	1.05
6	30.6	25.14	22,180	24.45	"	23,400	- 1220	5.22
7	24.1	12.38	15,000	14.12	949.0	13,420	+ 1580	11.78

Cooler

Run No.	lb./min fluid	Temp. Fall °C	B. t. u. lost by fluid	lb./min cooling water	Temp. Rise °C	B. t. u. gained by cooling water	B. t. u. unaccounted for	%
2	13.0	5.2	2690	19.2	3.55	2620	+ 20	0.77
3	23.6	2.35	1512	16.4	3.48	1557	+ 35	2.97
4	26.2	4.8	2730	26.8	4.58	2660	- 70	2.56
5	38.0	0.9	931	13.3	3.05	875	- 56	6.02
6	30.6	3.4	3000	13.45	7.86	3050	+ 50	1.67
7	24.1	9.99	- - -	32.3	18.95	30,700	- - -	- - -



Table III

WATER DATA

Table with columns: Run No., Description, Liquid - Deg Cent, Heater, Cooling Water, Steam, Friction Drop, Orifice Reading, etc. Includes sub-sections for RECORDED DATA, HEATER, and COOLER.

OIL DATA

Table with columns: Run No., Description, Liquid - Deg Cent, Heater, Friction Drop, Orifice Reading, etc. Includes sub-sections for RECORDED DATA, HEATER, and COOLER.

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## Calibration of Orifices

Series	Run No.	Description	Average Manometer Reading	Time in min	lb. of water	lb. per min.	Temp of water	ft <sup>3</sup> /min at Orifice
A	1	Orifice No. 2	5.55	12.95	200	15.44	5°C	0.248
	2	"	2.28	20.27	200	9.87	"	0.158
	3	"	7.47	11.18	200	17.89	"	0.287
	4	"	20.54	6.76	200	29.57	"	0.475
	5	"	11.05	9.23	200	21.66	"	0.348
	6	"	5.43	13.10	200	15.27	"	0.245
	7	"	15.03	7.89	200	25.34	"	0.407
	8	"	15.27	7.81	200	25.60	"	0.411
B	1	Orifice No. 1-A	1.21	14.14	100	7.07	"	
	2	"	2.68	9.41	100	10.63	"	
	3	"	4.91	14.02	200	14.27	"	
	4	"	8.34	10.86	200	18.42	"	
	5	"	12.37	8.91	200	22.46	"	
	6	"	16.75	7.68	200	26.06	"	
	7	Orifice No. 1	7.36	5.85	100	17.10	"	0.274
	8	"	13.55	4.30	100	23.26	"	0.373
	9	"	9.65	5.09	100	19.64	"	0.315
	10	"	2.08	5.45	50	9.18	"	0.147
	11	"	4.35	7.58	100	13.20	"	0.212
	101	Orifice No. 3	2.28	4.50	100	22.20	"	0.356
	102	"	5.50	4.39	150	34.20	"	0.549
	103	"	9.22	5.64	250	44.40	"	0.713
	104	"	14.16	5.50	300	54.60	"	0.876
	105	"					"	

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Table V

# HEAT TRANSFER COEFFICIENTS-CALCULATIONS

Series	Run	Fluid Flowing	Temp change in fluid °C	k	c	lb./hr.	Chu/hr.	ΔT °C	h	hD/K	cZ/h	(cZ/K) <sup>1/2</sup>	hD/K / (cZ/h) <sup>1/2</sup>	Dup/±
C	2	Heating water	52.3	0.351	1.00	780	40800	77.8	870	1222	2.07	1.309	934	107.3
	3	"	38.3	0.349	"	1416	51400	82.3	1040	1472	2.29	1.359	1085	181
	4	"	23.0	0.379	"	1572	36100	51.0	1172	1530	1.208	1.072	1428	346
	5	"	28.0	0.344	"	2880	60500	87.5	1530	1635	2.63	1.430	1143	328
	6	"	25.1	0.375	"	1836	46200	50.1	1390	1828	1.304	1.103	1636	382
	7	"	12.4	0.412	"	1445	17950	21.2	1400	1680	0.716	0.885	1900	485
	8	"	21.3	0.403	"	895	14850	27.7	888	1090	0.814	0.927	1175	211
H	6	Heating Velocity "B"	7.03	0.078	0.458	1041	3360	109	51.1	324	147.5	6.36	51.0	4.26
	7	"	6.6	"	0.460	1117	3390	107	524	391	138.3	6.20	53.4	4.92
	8	"	5.85	"	0.461	1215	3530	106	555	352	130.8	6.07	58.0	6.11
	9	"	5.2	"	0.463	1540	3710	104	592	375	121.2	5.90	63.6	7.83
	10	"	6.0	"	0.455	1350	3690	111	55.1	349	164.0	6.60	52.8	4.97
	11	"	5.3	"	0.457	1525	3700	110	55.9	354	153.0	6.43	55.0	6.03
	12	"	4.9	"	0.459	1735	3870	108	59.2	375	144.0	6.29	59.7	7.36
	13	"	4.4	"	0.463	1880	4040	105	64.0	405	123.8	5.95	68.1	9.86
	14	"	4.0	"	0.465	2250	4190	103	67.7	428	113.6	5.76	74.3	12.23
	15	"	3.9	"	0.466	2340	4260	102	69.3	438	111.6	5.72	76.6	12.97
	16	"	3.6	"	0.466	2410	4050	102	65.9	417	110.2	5.70	73.1	13.55
	17	"	3.6	"	0.467	2510	4210	101	69.2	436	106.0	5.62	78.0	16.66
	18	"	3.2	"	0.474	2500	3790	79.1	79.4	503	86.0	5.20	96.8	18.27
	19	"	2.6	"	0.473	2340	2880	79.0	60.5	383	89.2	5.27	72.7	16.53
	20	"	2.7	"	0.472	2170	2770	79.0	58.2	368	91.2	5.32	69.3	14.90
	21	"	3.15	"	0.473	1860	2770	78.9	58.3	369	86.7	5.26	70.2	13.15
	22	"	3.45	"	0.472	1267	2060	78.7	43.4	275	90.0	5.28	52.1	9.67
	23	"	3.95	"	0.472	1205	2250	77.7	48.0	344	90.4	5.30	57.3	8.42
	24	"	6.25	"	0.514	1805	5800	55.5	173	1097	31.6	3.59	306	26.8
	25	"	6.25	"	0.518	2510	8140	52.1	259	1640	30.1	3.53	465	57.2
	26	"	6.0	"	0.518	2390	7440	51.2	241	1527	29.8	3.52	435	55.2
	27	"	6.1	"	0.518	2250	7110	51.7	228	1445	29.8	3.53	409	51.8
	28	"	6.2	"	0.517	2210	7080	52.8	223	1413	30.2	3.53	400	50.1
	29	"	6.2	"	0.517	2050	6570	52.4	208	1318	30.3	3.53	373	46.6
	30	"	5.9	"	0.519	1785	5450	51.5	175.5	1112	29.8	3.51	316	41.3
	31	"	6.1	"	0.517	1525	4810	52.8	151.0	958	30.5	3.54	271	34.2
	32	"	5.9	"	0.517	1421	4340	52.2	136.0	875	30.2	3.53	248	32.2
	33	"	5.2	"	0.517	1113	2990	53.6	92.4	586	30.3	3.53	166	25.0
	34	"	5.0	"	0.515	926	2390	55.4	71.5	453	31.6	3.59	126	20.0
	35	"	5.6	"	0.517	1025	2970	53.0	92.9	588	31.4	3.58	165	22.0
	36	"	3.2	"	0.514	709	1165	56.3	34.3	217	32.3	3.82	80.0	15.0
	37	"	4.9	"	0.513	518	1288	57.0	37.4	237	32.3	3.82	65.5	10.8
	38	"	6.4	"	0.512	384	1260	58.0	35.9	227	32.5	3.83	62.5	7.96
C	2	Cooling Water	5.2	0.351	1.00	780	7310	21.0	582	819	1.448	1.147	713	154
	3	"	2.35	0.345	"	1416	5990	19.2	757	1084	1.726	1.227	884	243
	4	"	5.8	0.360	"	1572	13600	22.2	1023	1402	1.140	1.050	1335	387
	5	"	0.9	0.339	"	2880	6690	67.2	1662	2420	2.060	1.307	1852	427
	7	"	10.0	0.377	"	1445	26500	24.4	1613	2370	0.790	0.816	2590	481
H	1	Cooling Velocity "B"	4.25	0.078	0.460	208	406	23	29.5	187	140.0	6.23	30.0	0.304
	24	"	3.85	"	0.516	1805	3390	57	89.3	629	31.1	3.57	176	39.7
	25	"	3.8	"	0.519	2510	4950	49	169.0	1070	29.8	3.50	306	58.2
	26	"	3.9	"	0.519	2390	4850	48	166.5	1070	29.4	3.48	308	55.9
	27	"	4.0	"	0.519	2250	4670	48	162.0	1028	29.4	3.48	296	57.6
	28	"	3.9	"	0.518	2210	4430	48	154.0	976	29.8	3.51	276	50.7
	29	"	4.0	"	0.518	2050	4250	49	145.0	918	29.8	3.51	261	47.3
	30	"	4.4	"	0.519	1785	4070	43	158.0	1000	29.4	3.48	287	41.7
	31	"	5.0	"	0.517	1525	3940	57	115.5	728	29.7	3.51	207	34.5
	32	"	4.5	"	0.518	1421	3320	57	97.3	616	30.3	3.53	175	32.3
	33	"	4.6	"	0.517	1113	2690	67	68.5	434	30.4	3.53	123	25.1
	34	"	4.2	"	0.514	926	2000	64	52.2	330	32.3	3.82	91.2	19.7
	37	"	3.5	"	0.513	518	920	63	24.4	154.5	32.3	3.88	50.2	10.76
	38	"	4.0	"	0.513	384	780	63	19.6	124.0	32.5	3.62	34.7	8.05
	40	"	3.6	"	0.503	467	730	53	29.0	145.6	39.8	3.91	37.2	6.77
	41	"	1.6	"	0.500	506	675	59	19.1	121.0	43.2	4.03	30.0	12.98
	42	"	2.1	"	0.500	1027	1078	58	31.0	196.5	43.3	4.03	48.8	15.10
	43	"	2.9	"	0.503	1730	2520	55	76.5	485	39.4	3.90	125	29.30
	44	"	2.8	"	0.502	1560	2150	56	65.3	412	40.3	3.93	105	25.8
	45	"	2.9	"	0.501	1362	1965	56	58.6	371	41.4	3.96	93.7	21.7
	46	"	2.4	"	0.501	1160	1405	58	40.5	257	42.0	3.99	64.4	18.5
	47	"	3.0	"	0.504	2360	2720	55	82.7	254	38.8	3.87	13.6	31.0

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