

1 A Twenty-Year Dynamical Oceanic Climatology: 1994-2013.

2 Part 2: Velocities, Property Transports, Meteorological
3 Variables, Mixing Coefficients

4 *Draft Version 1.3**

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10 June 13, 2017

11 **Abstract**

12 The World Ocean Circulation Experiment (WOCE) was created to produce the first
13 truly climatologically useful picture of the ocean circulation and its variability. This goal is
14 addressed here from the state estimate of the Estimating the Circulation and Climate of the
15 Ocean (ECCO) consortium, which uses almost all of the data obtained during WOCE and
16 its aftermath along with the much improved general circulation modeling capabilities. A
17 dynamically and data-consistent, state estimate is available depicting the ocean and its ice-
18 cover over a 24-year time-span, globally, from the sea surface to the sea floor. The resulting
19 time-dependent 20-year long climatology includes temperature, salinity, surface elevation,
20 bottom pressure, sea-ice, and three components of velocity. Accompanying the state estimate
21 are modified estimates of meteorological forcing-fields, ocean interior mixing coefficients, and
22 initial conditions. Much spatial structure persists through the two-decade averaging. Results
23 here are primarily pictorial in nature, intended to give the wider community a sense of what
24 is now available and useful and where more detailed analysis would be fruitful. An extended
25 reference list is included.

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1 Introduction: The State Estimate (Mostly Repeated from Introduction to Part 1)

Purpose

One of the central goals of the World Ocean Circulation Experiment (WOCE) was to produce the first truly global time-varying estimate of the circulation over approximately a decade, an estimate that would be useful in defining the major climatologically important ocean elements. The Estimating the Circulation and Climate of the Ocean (ECCO) project was formed near the start of the WOCE field program so as to address this goal using both the conventional and newly-deploying WOCE observation system, along with the rapidly advancing general circulation modelling capability (Stammer et al., 2002). In this paper, and in subsequent Parts, this WOCE goal is addressed by defining a time-dependent climatology over the 20-year (bidecadal) interval 1994-2013. Little or no dynamical or kinematical interpretation is provided—that is left to other authors and times.

Various oceanic climatologies are in use by the oceanographic and climate dynamics communities. They serve as tests of models, as initial conditions, and as a basic descriptor of the ocean. Definitions of climatologies vary widely both in terms of how they were formed and the durations they represent. Here we describe a 20-year average modern climatology from a dynamically consistent model that also has a consistent fit to the majority of global data between 1992 and 2015 (Wunsch and Heimbach, 2013; Fukumori et al., 2017). The climatology is based upon the ECCO version 4 state estimate (Forget et al., 2015). It derives from a least-squares fit of the evolving MITgcm (Marshall et al., 1997; Adcroft et al., 2004; Forget et al., 2015) to the numerous and diverse global observations. A summary would be that all of the Argo, altimetry, the CTD hydrography appearing in the WOCE Climatology and successors (Gouretski and Koltermann, 2004; Talley et al., 2016), all extant, bias error-corrected XBTs, the considerable elephant seal profile data (Roquet et al., 2013), GRACE mission mean and time-dependent geoids, satellite-measured sea surface temperature and salinity, and the ECMWF¹ ERA-interim reanalysis of the meteorological variables (Dee et al., 2011, 2014), have been included, with the fits inferred to be adequate relative to the estimated uncertainties of the data. (Atmospheric reanalyses should not be considered “data”, however.)

Previous climatologies, e.g. Levitus et al. (1982) and its later incarnations as the NOAA World Ocean Atlas, or Gouretski and Koltermann (2004) have usually been based only upon temperature and salinity averages and over much longer time intervals than employed here. Other climatologies (e.g., AchutaRao et al., 2007) have focussed on the upper 700 or 1000m

¹European Centre for Medium Range Weather Forecasts

59 and relied heavily on XBT measurements. Ishii et al. (2005) is a climatology of the sea surface
60 temperature. As such, all these suffer from the very great inhomogeneities of data distribution
61 prior to the WOCE period and a series of untestable statistical hypotheses (see e.g., Kennedy
62 et al., 2011; Wunsch, 2016; Boyer et al., 2016). This present climatology differs from earlier
63 ones most obviously in its production of the three-dimensional, time-varying, three components
64 of velocity and of a self-consistent surface meteorology, as determined at the model time-step,
65 $\Delta t \approx 1$ h. Use of *any* fluid climatology confronts one basic problem: that the resulting time or
66 space-time average fields do not satisfy any simply derivable equations of motion—requiring a
67 variety of turbulence closure schemes—and the relationships among the different variables can
68 be complicated and poorly known. Here, time/space means of fluid quantities are based upon
69 the uniform average of fields exactly satisfying the model equations at each model time-step (at
70 present, 1 hour) and grid-point. Some authors have used ocean general circulation models fit
71 to data in methods analogous to those in meteorology and commonly known as “reanalyses.”
72 These, unfortunately, are usually not property conserving (heat, salt, momentum, etc.) and
73 thus unsuitable for global-scale climate calculations (see e.g., Wunsch and Heimbach, 2013; and
74 Fig. 1 of Stammer et al., 2016).

75 Some sketches of global-scale analyses of earlier multi-decadal ECCO estimates have been
76 published starting with Stammer et al. (2002). Among them, an earlier 16-year global time-
77 average was described by Wunsch (2011), with a focus on the accuracy of Sverdrup balance, and
78 Wunsch and Heimbach (2014) discussed the heat content changes. Liang et al. (2015, 2017)
79 describe the vertical redistribution of heat and Forget and Ponte (2015) the regional sea level
80 changes. Forget (2010) presented an 18-month estimate from an earlier ECCO state estimate.
81 In general, the present solution differs only subtly from those previously used, with the chief
82 differences being ascribed to the inclusion of more data over a longer duration, inclusion of
83 geothermal heating (see Piccuch et al. 2015), improvements in the handling of sea ice, and
84 where appropriate, separate uncertainties for time-average and time-anomaly measurements.
85 Solutions are generally robust, as much of the volume of ocean in the model state vector is in
86 near-geostrophic balance with the density field at all times longer than a few days.

87 By choosing the period following 1992, a much more nearly uniform global data coverage
88 is obtained than was possible earlier. Chief among the remaining data inhomogeneities are the
89 intensification of the Argo float profile data availability after about 2005.

90 Any temporally averaged state will be considerably smoother than states which are sampled
91 more or less as “snapshots.” Thus classical ship-borne hydrographic sections (e.g., Fuglister,
92 1960 or the various WOCE Atlases) show many small-scale features which vanish on averaging.
93 Suppressed features include internal waves, tides, and geostrophically balanced eddy motions.

94 Meandering currents, such as the off-shore Gulf Stream, are broader and smoother than in any
95 near-synoptic estimate. In addition, fluid regions that are only marginally or poorly resolved
96 numerically (particularly boundary currents), will be smoother than even a true 20-year average
97 would be. Nonetheless, even a 20-year average leaves remarkably many structures much smaller
98 than the basin-scale in the estimated circulation.

99 No model with a nominal horizontal grid-spacing of 1° of longitude can resolve small-scale
100 circulation features, which include the important boundary currents. Nonetheless, the near-
101 geostrophy of the bulk of the ocean supports the conjecture that to the extent that a successful
102 fit to the interior temperature, salinity, and altimetric fields and surface boundary conditions, has
103 been obtained, the boundary currents will be forced by the interior flows to carry the appropriate
104 amount of mass (volume), temperature, etc. so as to satisfy the basic overall conservation laws.
105 This conjecture, upon which we rely, but which is tested elsewhere, can be regarded as a re-
106 statement of that used by Stommel and Arons (1960) in their discussion of deep boundary
107 currents—whose existence and structure was fixed by the mass and property requirements of
108 the interior flow—even though they were not dynamically resolved.

109 As with any estimation problem, a crucial element in the determination of the best values
110 lies with the use of realistic error estimates for *all* of the data that are being fit. For a full
111 discussion of the error estimate used here, reference must be made to the literature. Temperature
112 measurements are described by Forget and Wunsch (2007) and Abraham et al. (2013). Altimetry
113 accuracies are discussed by Fu and Haines (2013) and Forget and Ponte (2015). For the gravity
114 data from the GRACE mission, see Quinn and Ponte (2008). Satellite surface salinities are
115 addressed by Vinogradova et al. (2014). Meteorological variable accuracies are described e.g.,
116 by Chaudhuri et al. (2014, 2016).

117 This paper is *not* an in-depth analysis of *any* features of the global ocean circulation. It
118 is instead mainly visually descriptive—a suggestive pictorial subsample—intended primarily to
119 serve as an invitation to the wider community to exploit it by demonstrating various products.
120 With the widespread recognition that a steady-state ocean never exists, attention turns instead
121 to the temporal changes over the estimation period. Here for descriptive purposes, a few pictures
122 of changes year-by-year for 20 years, by 20-year averages for each month, and by season, are
123 displayed. All results can readily be calculated month-by-month at the expense of using a larger
124 volume of numbers.

125 Results here are intended mainly to be indicative of possibilities and an invitation to use,
126 rather than being the most precise or accurate possible. Thus for example, the heat capacity,
127 c_p and the mean density, $\bar{\rho}$ are treated as constant in calculations of heat uptake even though
128 both are (weak) functions of position.

129 *The State Estimate*

130 The ECCO state estimate is obtained from the *free-running* MITgcm after the adjustment of
131 the control parameters required to fit the data. In the least-squares methodology with Lagrange
132 multipliers (see Wunsch and Heimbach, 2013), the entire interval 1992-2015 has been fit to
133 the data. Parameters adjusted include the three-dimensional, top-to-bottom, initial conditions,
134 internal mixing coefficients, and the surface meteorology. At any given time in the estima-
135 tion interval, the solution represents data both preceding and *following* that date so that the
136 equations are always satisfied while coming as close to the data as possible within uncertainty
137 estimates. The 20-year period 1994-2013 has been chosen for averaging as sufficiently distant
138 from the poorly constrained earlier years before the high accuracy altimetry begins in late 1992
139 and the time of the then non-existent data following 2016. The period corresponds to that of
140 complete coverage by satellite altimetry, the WOCE CTD survey, and the interval after about
141 2005 when the Argo array became fully-deployed. All data, plus the ECMWF estimate, have
142 been assigned uncertainties that include both instrumental and natural noise. After adjustment
143 of the parameters, the state estimates are the solution to a forward model satisfying all basic
144 conservation requirements. Structurally, it is no different from any other unconstrained model
145 estimate except that its residual data misfits are fully known.

146 No state estimate is definitive or “correct”; they are “best-estimates” for the present time:
147 data are continuously added, both from more recent years and previously omitted earlier val-
148 ues; estimated data errors are sometimes revised; models are improved; and in all situations,
149 minimizing iterations are ongoing. Values shown here are obtained from ECCO version 4 as of
150 mid-November 2016.

151 Undoubtedly the state estimate has residual systematic errors at some level, particularly
152 in data-poor regions and times. To some extent, these will be removed when considering only
153 temporal changes in the state over the 20-years and these latter are given some emphasis.
154 Uncertainty estimates remain an amorphous problem: much of the variability in the model
155 represents deterministically evolving elements. Stochastic elements are introduced by weather,
156 some longer-period meteorological variability, and by elements of the initial-conditions best
157 regarded as random. Because the true probability distributions are not known, discussion of
158 estimate uncertainties is postponed to an intended Part 4.

159 A full description of the many features of the 20-year global ocean circulation requires a
160 book-length publication, if not a library. The strategy here is to sketch the gross hydrographic
161 and circulation features and to do a limited comparison to a few of the special regions (boundary
162 currents, mixed-layer, etc.) to provide some of the flavor of the differences between a moderate-
163 duration, nearly homogeneous, average and both the more common limited-time analyses usually

164 available (classical synoptic hydrographic sections), as well as the far more data-inhomogeneous
165 published climatologies.

166 With time-mean fields being spatially and temporally smoother than in nominally synoptic
167 measurements, second order quantities such as the time averages e.g., $\langle \mathbf{v} \rangle \langle T \rangle \neq \langle \mathbf{v}T \rangle$, where
168 $\langle \cdot \rangle$ denotes a space-time average, and the difference between them may be very large. Much
169 of physical oceanography has been based upon the unstated assumption that quasi-synoptic
170 measurements represented the mean motion. Thus e.g., the calculation of Sverdrup balance, or
171 of “abyssal recipes”, are implicitly steady-state results, despite the common use of individual
172 hydrographic stations or sections. Here true 20-year average estimates are now possible. This
173 description and discussion thus largely focusses on the properties of single variables, T, u , etc.,
174 their 20-year means and estimates of the deviation from those means. As Part 2, this paper
175 describes the three dimensional Eulerian velocity field and the estimated (that is, adjusted)
176 meteorological forcing. The hydrographic fields and related properties are discussed in Part 1.
177 Most emphasis is placed on the global fields. A number of higher resolution, regional versions,
178 of the state estimate exist (e.g., Gebbie et al., 2006; Mazloff et al., 2010), and a high northern
179 latitude version is forthcoming (An Nguyen, in preparation, 2017), but these estimates are not
180 further discussed here.

181 All of the ECCO system output described here is available in Matlab[®] form at: <http://mit.ecco->
182 [group.org/opendap/diana/h8_i48/contents.html/2](http://mit.ecco-group.org/opendap/diana/h8_i48/contents.html/2) as 20-year means, 20-separate annual means,
183 20-year average individual months, and 20-year average seasonal means (DJF, MAM, JJA, SON)
184 on a grid in 50 vertical levels, of thickness plotted in Fig. 1. Many studies are best done in
185 isopycnal-like coordinate systems; but the present description is confined to calculations in geo-
186 metrical (latitude-longitude-depth) coordinates, with the interpolations to isopycnals postponed
187 (but see Speer and Forget, 2013 for a mode water discussion).

188 2 Eulerian Horizontal Velocities

189 *Misfits*

190 As described in Part 1 (ECCO Consortium, 2017), a misfit can be computed between the
191 state estimate and any particular data type. Here, Fig. 3 displays the misfit to some of the
192 TOGA-TAO equatorial current meter array data (Hayes et al., 1991) annual means to the state
193 estimate. Note that in this case, the data were *not* used as constraints on the state estimate,
194 and are thus a completely independent test. At shallower depths (not shown), the consistency
195 between the two estimates is even better.

²Or contact Carl Wunsch directly (cwunsch@mit.edu) for data or advice.

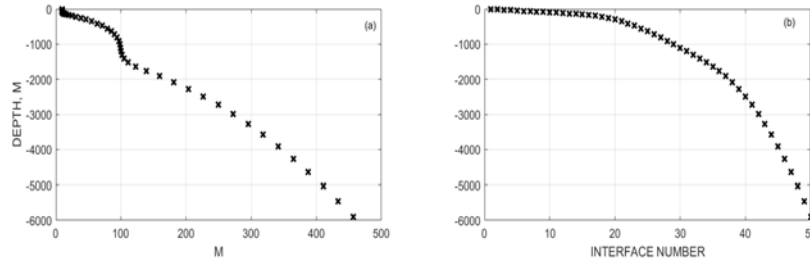


Figure 1: (a) Level thicknesses; (b) level depths in the ECCO version 4 of the MITgcm.

{interfaces_la

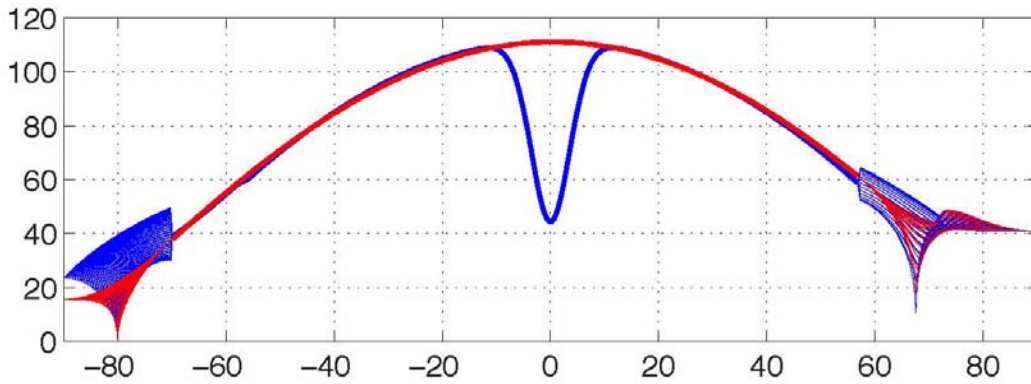


Figure 2: Latitude (blue curve) and longitude spacing in kilometers as a function of latitude (from Forget et al., 2015). Higher latitude spacing exists near the equator. At high latitudes the more complex grid leads to a distribution of spacings (see Figs. 1, 2 of Forget et al., 2015). Most of the high latitude southern region is land.

{fig03-eccov4_

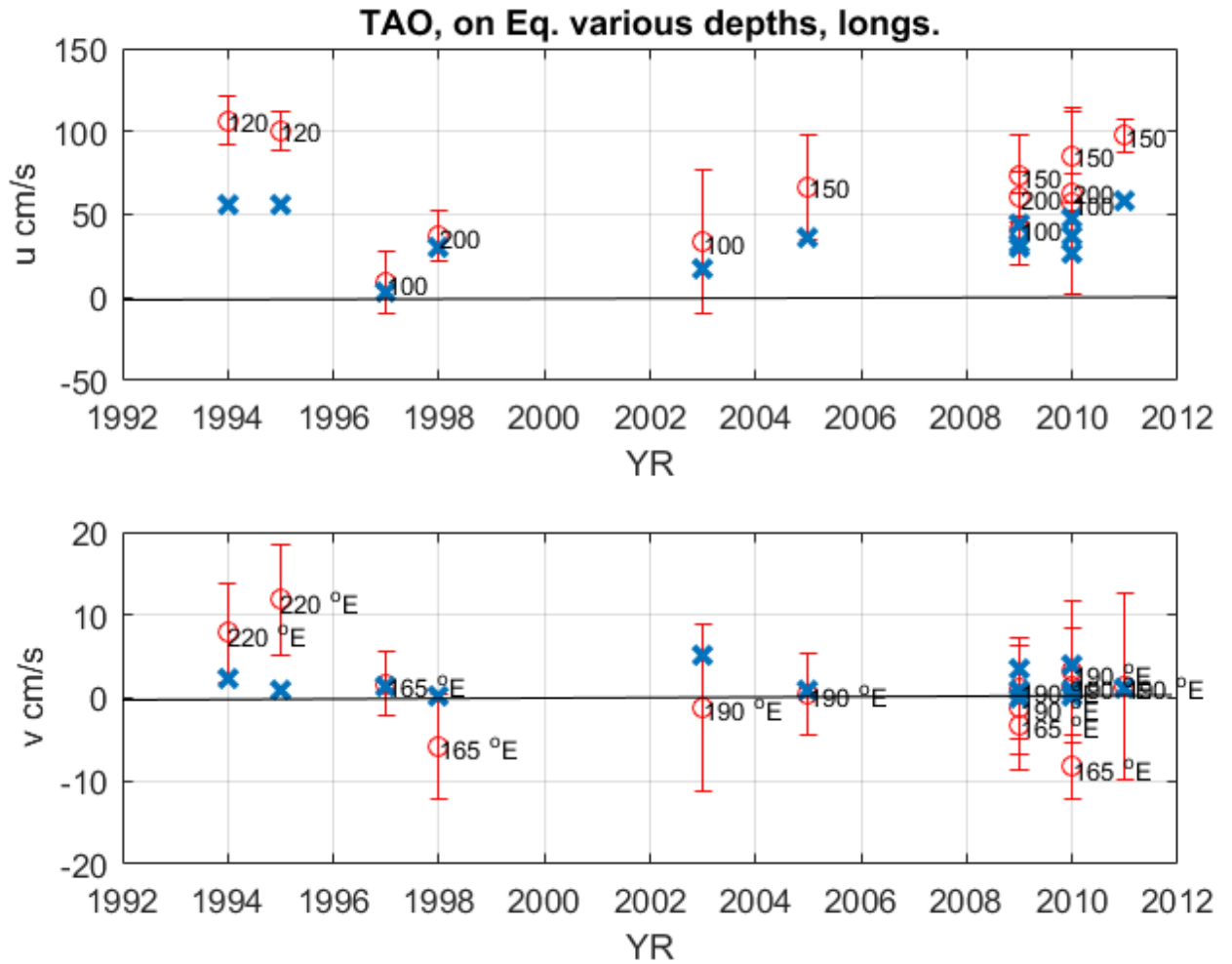


Figure 3: Upper panel shows the u component from the TAO array on the equator at various depths (red symbols) with standard errors. 'x' denotes the corresponding ECCO state estimate annual mean. Values are within one standard error. Labels are the water depth. Lower panel shows the same result for the v component. Now the labels indicate the longitude of the measurement.

{tao_anmeans_

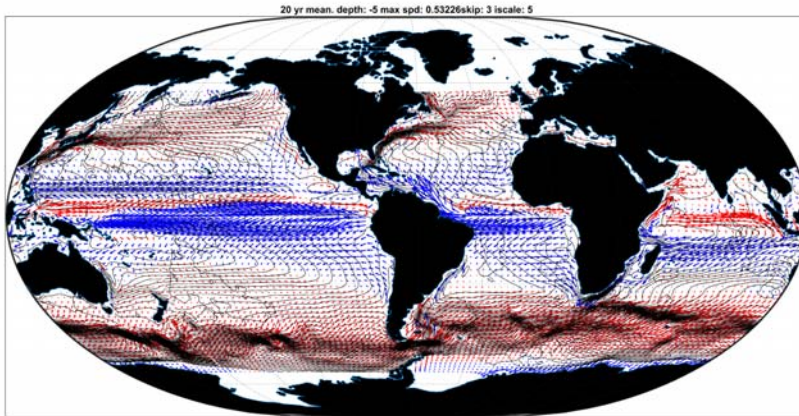


Figure 4: The 20-year average Eulerian flow at 5m depth superimposed upon the time-mean surface elevation, η . Red arrows have an eastward component, blue a westward one. Largest value here (longest arrow) correspond to 40cm/s. In the centers of gyres, particularly, the ageostrophic component of flow visually crosses the surfaces of constant elevation.

{quiver_map_5m

196 *Time Means*

197 Figs. 4-8 depict the 20-year Eulerian mean flow fields as arrow plots at four depths. A
 198 number of distinct, expected features can be seen. These include the strongly divergent (to
 199 north and south) flows on the equator, the western boundary currents and their extensions as
 200 well as the Antarctic Circumpolar Current. All of these flows are broader and smoother than is
 201 familiar from attempts at instantaneous depictions. The corresponding pressure field contours
 202 are also shown as a visual guide.

203 The time average zonal flow on the equator is displayed in Fig. 9 with a conspicuous equator-
 204 ial undercurrent; and the average meridional flow across the equator is in Fig. 10. Time average
 205 zonal flow in the Drake Passage is shown in Fig. 11 with a net transport of 146Sv, close to most
 206 published values (Meredith et al., 2011), but in contrast to the much larger transport claimed
 207 by Donohue et al. (2016), the difference probably owing to the strong assumptions made there.
 208 The estimated value here is necessarily consistent with the near-geostrophic interior flows both
 209 to the west and east of the passage. Mild annual variations in the transport are depicted below.
 210 Fig. 12 shows the remarkably complex meridional mean flow at 60°S, a latitude passing through
 211 the Drake Passage.

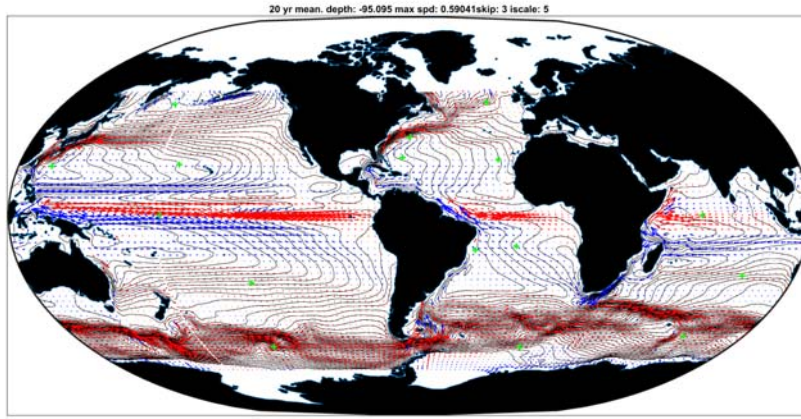


Figure 5: Twenty-year average of the mean horizontal flow at 95m superimposed on the time-mean sea surface elevation. Largest value is 59 cm/s. Vectors more closely follow the elevation lines than does the velocity at 5m in Fig. 4. Note the strong eastward flow on the equator as compared to the near-surface values.

{quiver_map_10

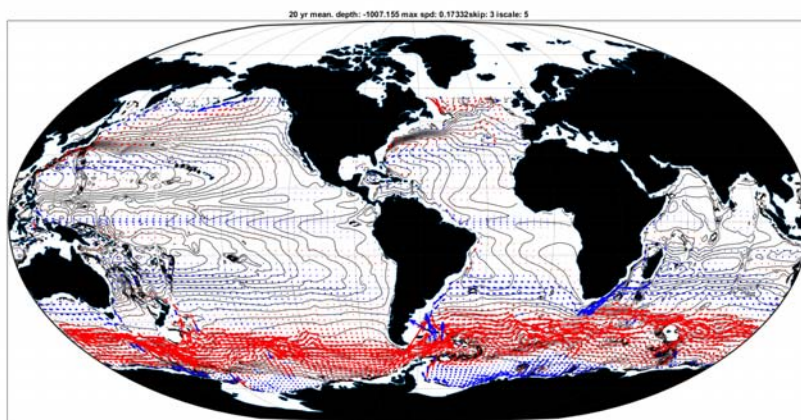


Figure 6: Twenty-year mean flow at 1000m (compare Ollitrault and Colin de Verdiere, 2014). Largest value shown is 17 cm/s, but arrow lengths are saturated in the Southern Ocean. Weak banding is visible in the tropics generally. The corresponding hydrostatic pressure field at this depth is shown.

{quiver_map_20

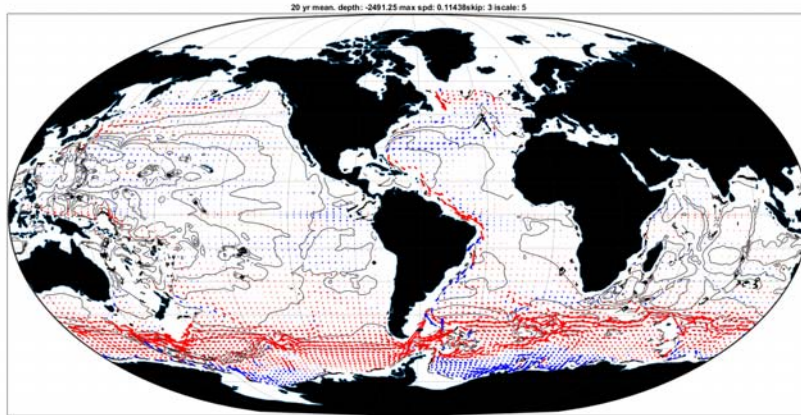


Figure 7: Same as Fig. 4 except at 2500m . Largest arrow corresponds to 13 cm/s. The Atlantic deep western boundary current and the Southern Ocean eastward flow are the most conspicuous features.

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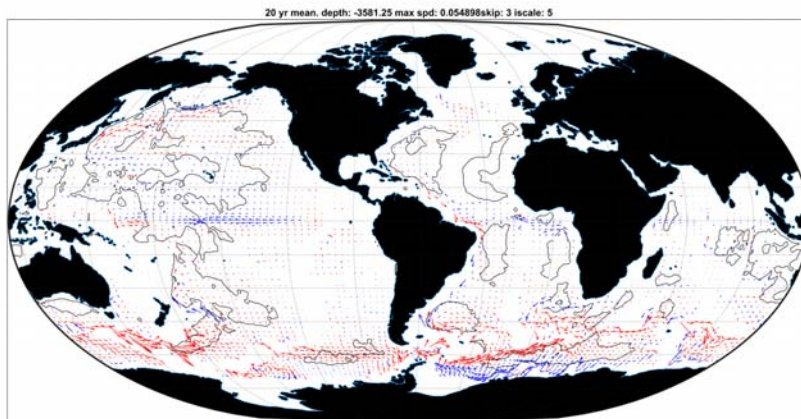


Figure 8: Twenty-year average horizontal flow at 3600m with the 5000m contour and not the pressure field. Largest arrow is 5.5 cm/s.

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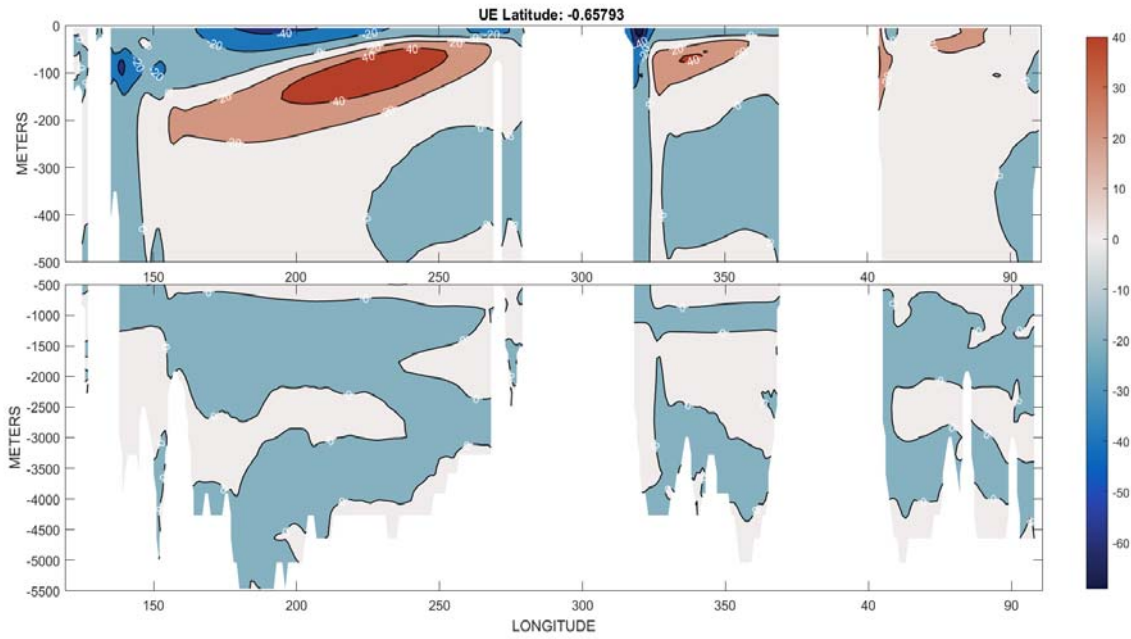


Figure 9: Twenty-year average Eulerian zonal flow, u , along the equator in all three oceans (cm/s). The eastward flowing equatorial undercurrent is visible in the Pacific and Atlantic Oceans, as is a zonal westward flow below.

{ue_20yr_secti

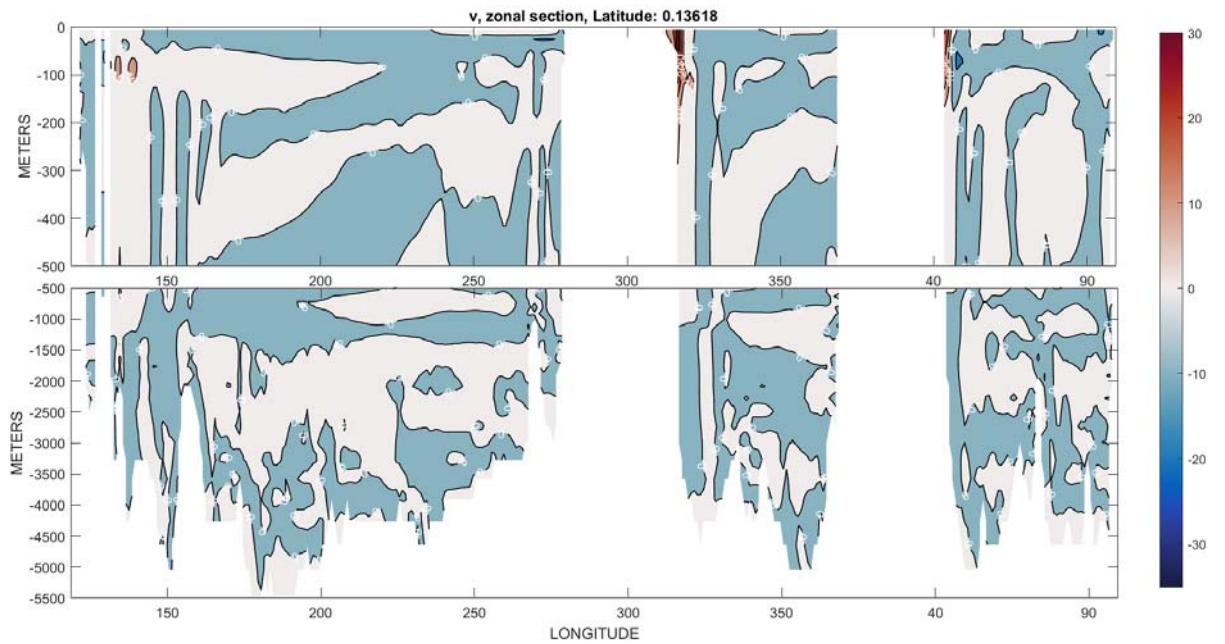


Figure 10: Twenty-year average mean Eulerian meridional velocity, v , at the equator (cm/s).

{vn_equatorial

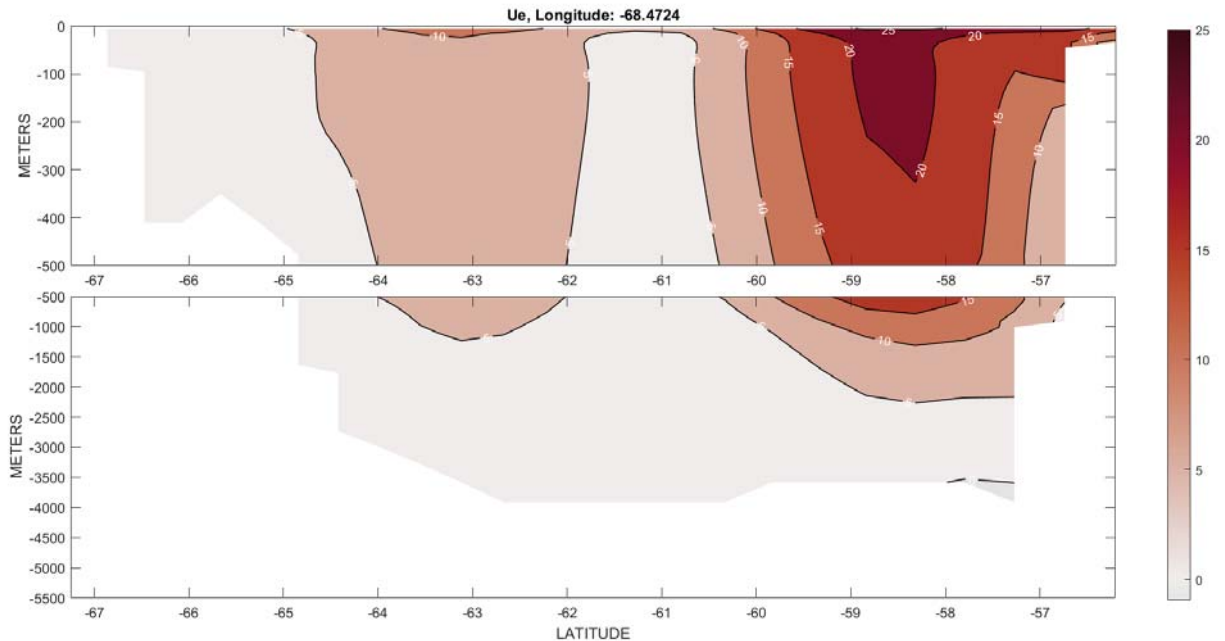


Figure 11: Twenty-year average *zonal* flow, u in the Drake Passage at 70°W . The 20 year average transport is 146 Sv.

{zonalflow_20y

212 3 Time-Dependent Flows

213 The oceanic flow field varies on all time scales from seconds to the age of the ocean. In Figs.
 214 13-15 are shown the anomalies of Eulerian velocity about the 20-year mean at 5 m.

215 A few representative anomalies of the annual average meridional component, v , are shown
 216 in Figs. 16-18 across the equator. Such results become part of the story of tropical variability
 217 including the ENSO cycle.

218 Oceanic kinetic energy is one of its basic physical properties. Fig. 19 displays the logarithm
 219 of the 5m depth value of the kinetic energy in one year (2004). As expected, some variation in
 220 total kinetic energy (top-to-bottom) for each of the 20 years as well as that for the abyssal layer
 221 (3600m to the bottom) can be seen in Fig. 20. The slow overall increase over 20 years and the
 222 decay in the abyss are not easily testable.

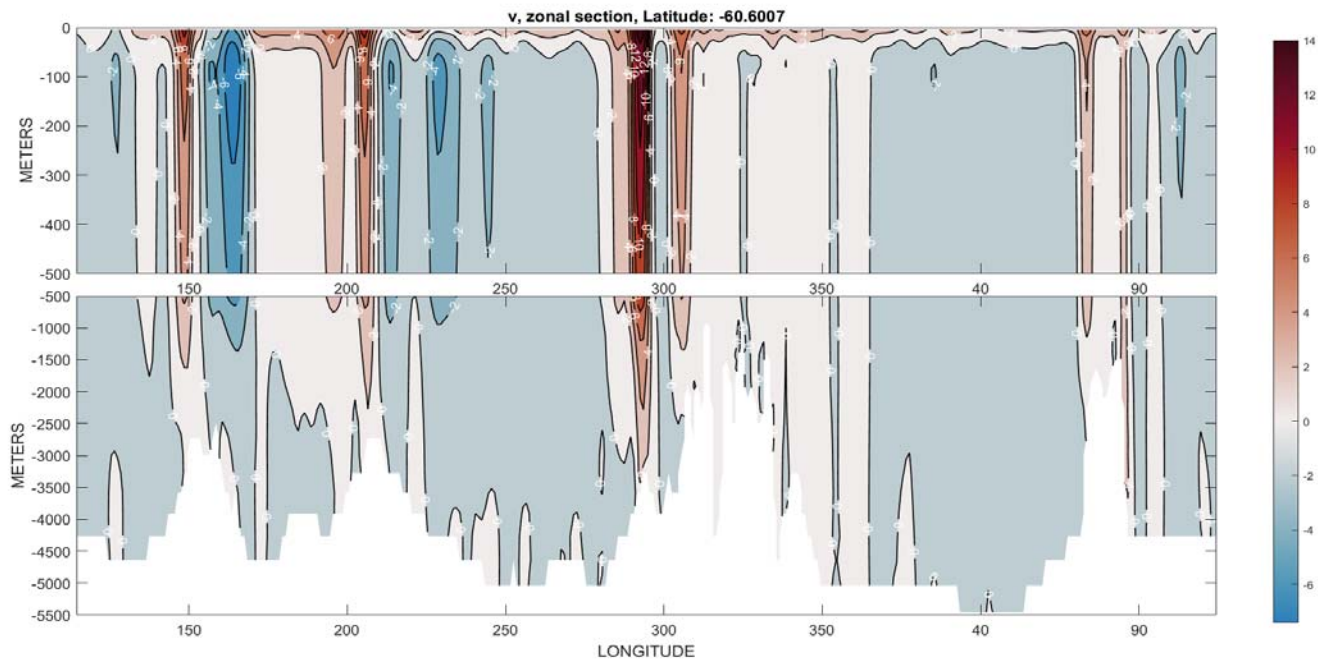
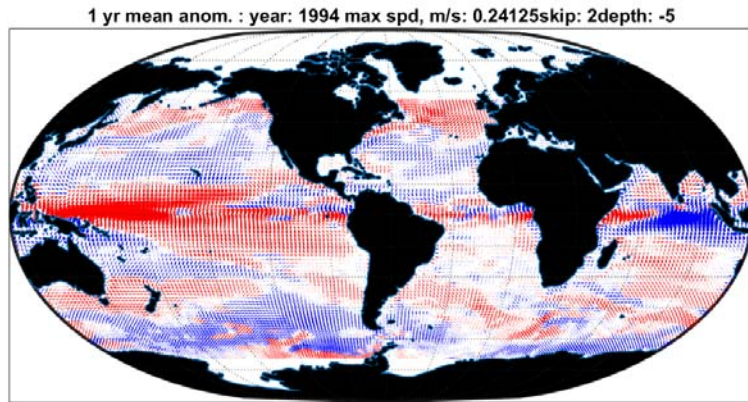


Figure 12: Twenty-year mean meridional velocity, v , in a section through the Drake Passage. A conspicuously variable structure survives 20-years of averaging.

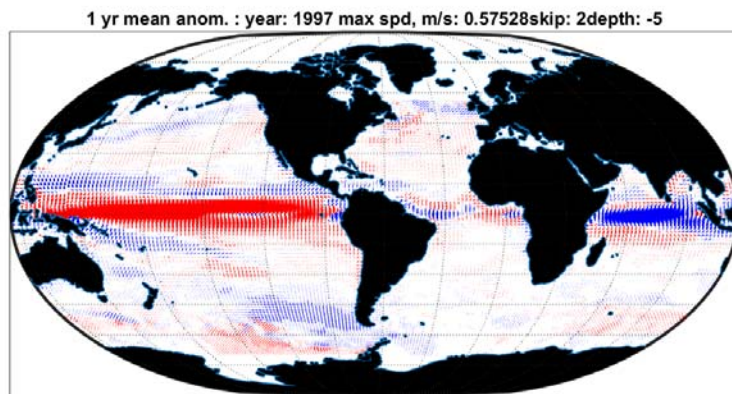
{vn_drakepassa



anom yr 1994.tif

Figure 13: Anomaly of the 5m horizontal flow in 1994, again with red arrows having an eastward component. Largest arrow is 24 cm/s.

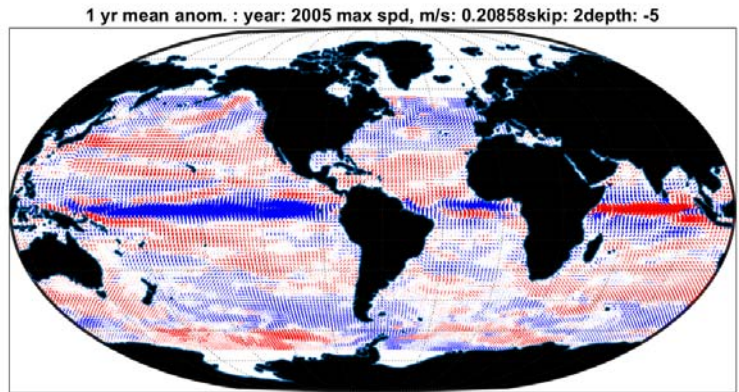
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Figure 14: Same as Fig. 13 except for 1997 with the largest arrow at 58 cm/s.

{quiver_anom_y



anom yr 2005.tif

Figure 15: Same as Fig. 13 except for 2005 with the largest value be 21 cm/s.

{quiver anom y

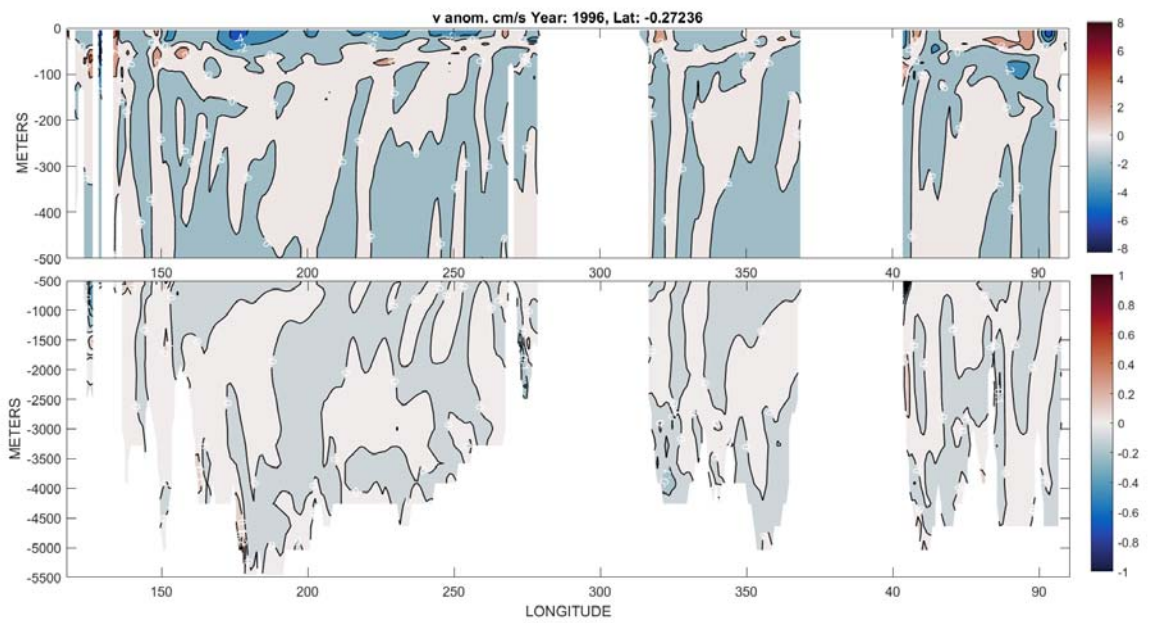


Figure 16: Anomaly of meridional flow across the equator in 1996 (cm/s).

{vanom_1996_la

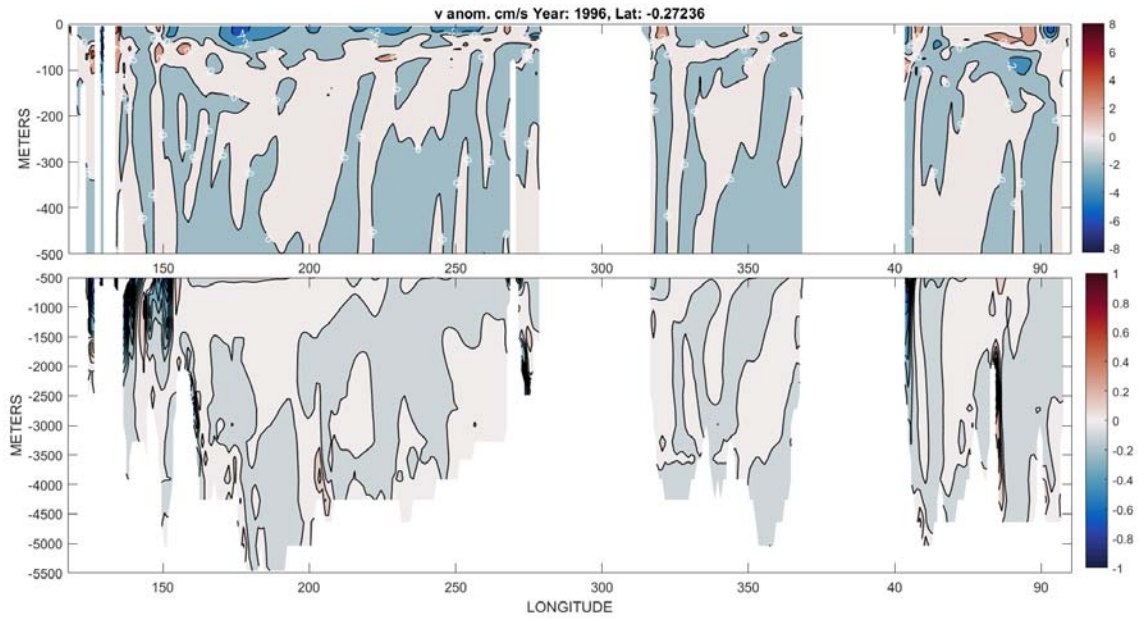


Figure 17: Anomaly of meridional flow across the equator in 1998 (cm/s)—an El Niño year.

{vanom_1998_la

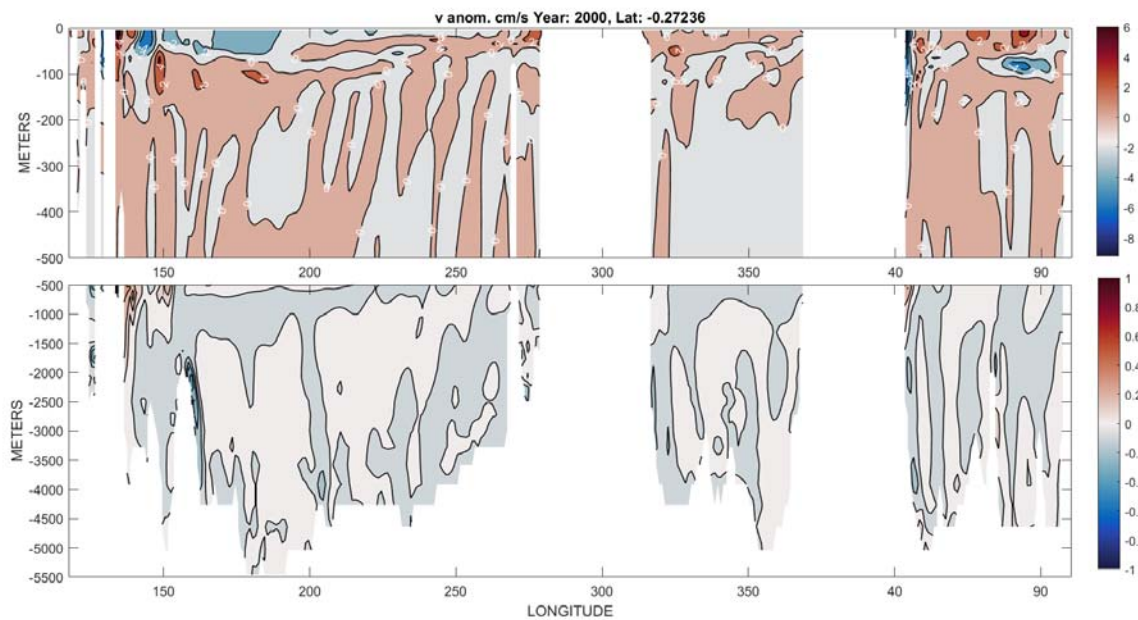


Figure 18: Anomaly of meridional velocity, v , (cm/s) at the equator in 2000.

{vanom_2000_la

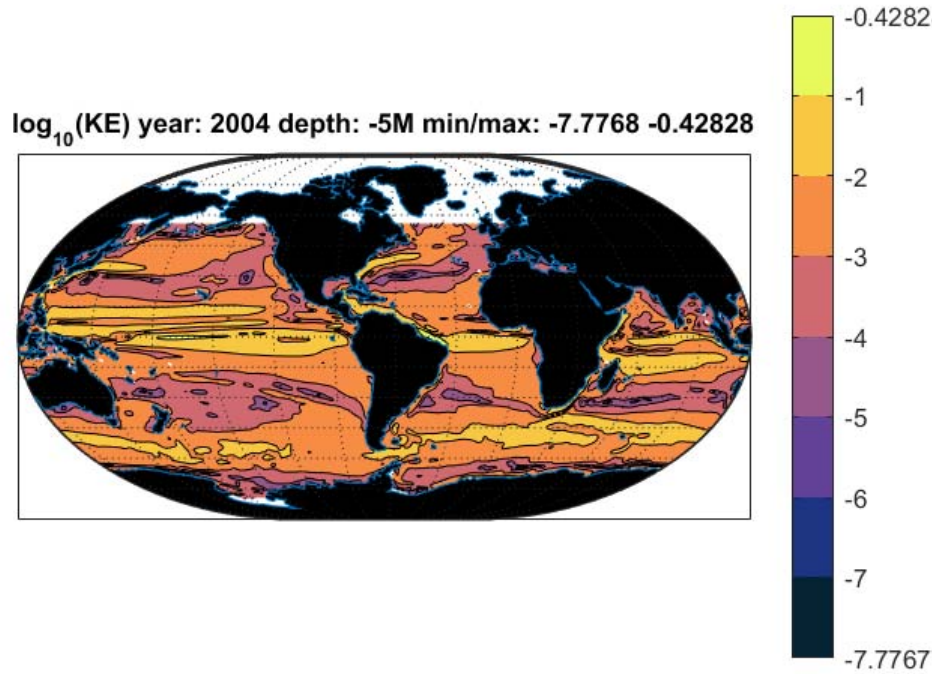


Figure 19: Logarithm of the Eulerian horizontal kinetic energy/unit mass at 5m averaged over 2004. Other years are visually similar, differing in details.

{ke_5m_2004.ti

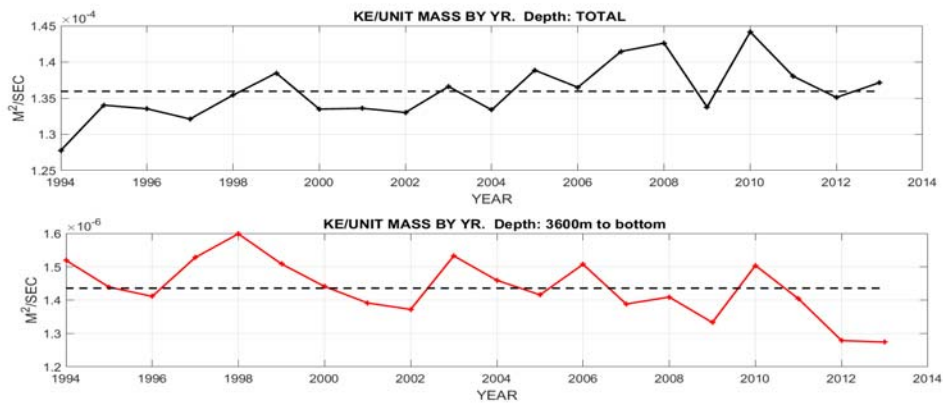


Figure 20: (Upper panel) Total (top-to-bottom) but excluding the northern high latitudes, kinetic energy/kg by year. El Niño year 1998-99 is prominent early in the record. A weak upward trend might be real. (Lower panel) Kinetic energy/unit mass by year in the layer 3600m to the bottom. Note the scale change from the upper panel.

{ke_total&3600

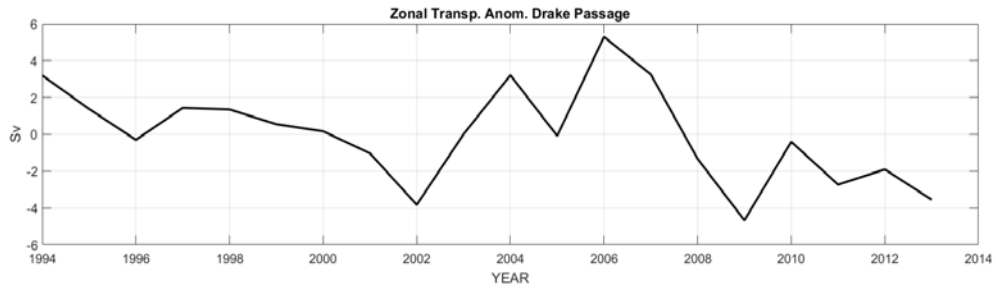


Figure 21: Anomaly (Sv) of transport integrated across the Drake Passage for each year.

{yearly_trans_

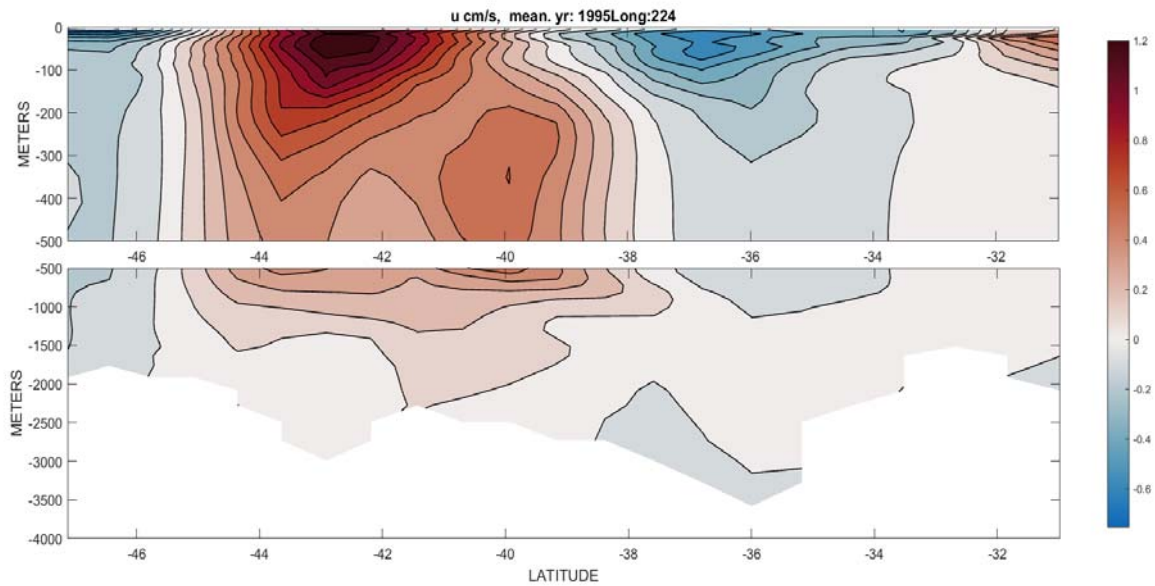


Figure 22: Anomaly of the zonal flow in the Drake Passage in 1995 (cm/s).

{u_drakepassag

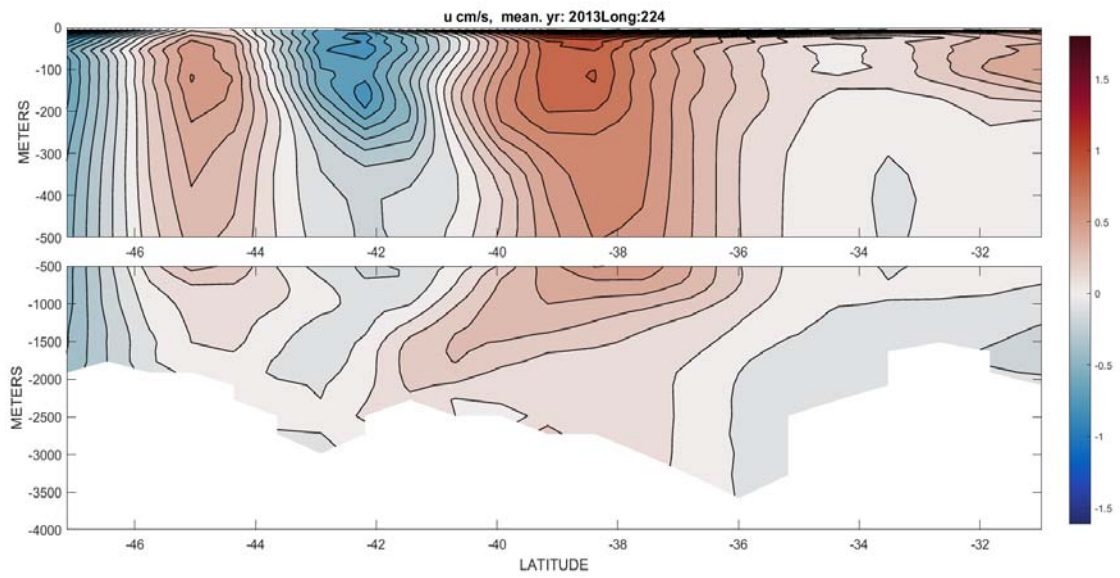


Figure 23: Anomaly of the zonal flow (cm/s) through Drake Passage in 2013.

{u_drakepassag

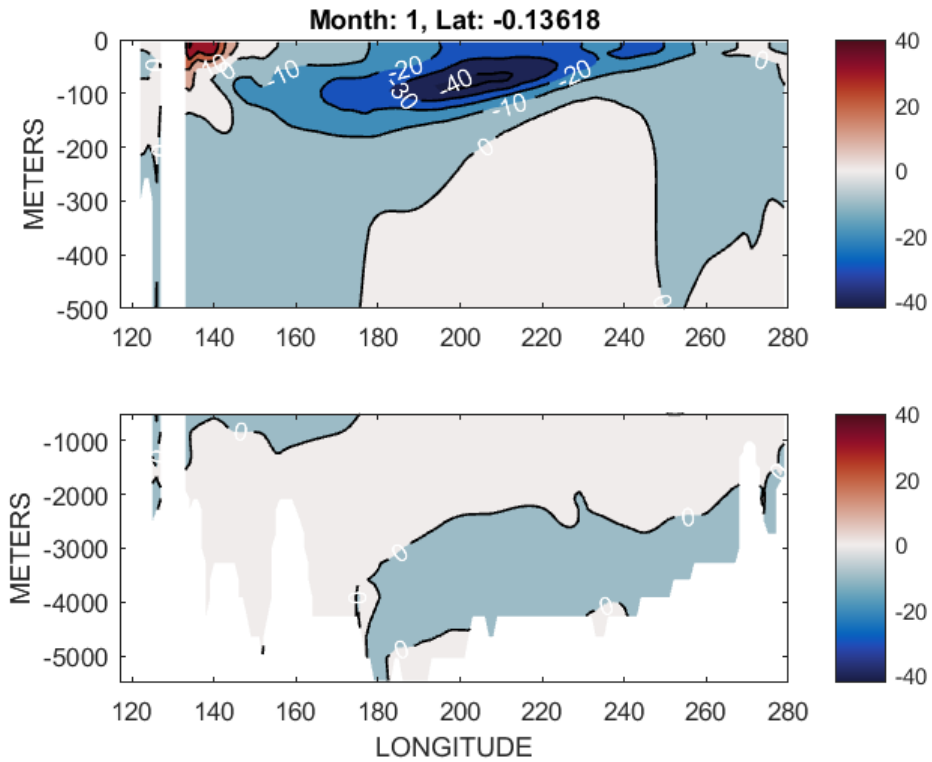


Figure 24: Twenty-year mean zonal flow anomaly (cm/s) on the equator in January in the Pacific Ocean.

{equator_jan_s

223 3.1 Annual Cycle

224 The annual cycle dominates the atmospheric climate system, with a similar strong response
 225 in the very upper levels of the ocean. Simple Rossby wave theory (e.g., Gill and Niiler, 1973;
 226 Wunsch, 2015) shows that the vertical penetration of the baroclinic response to annual forcing
 227 at the surface is very restricted, but a bit deeper on the equator. An example of the mean annual
 228 cycle, shown as the 20-year average of the monthly anomaly of u , along the equatorial section
 229 in the Pacific Ocean is displayed in Figs. 24-27 for a few months.. Although the response in the
 230 upper 100 m is far larger than at depth, a detectable annual cycle in u exists to the sea floor.
 231 Note that interpretation of the upper ocean structures requires use of the mean flow in Fig. 9,
 232 as a positive anomaly will weaken the westward-going near-surface South Equatorial Current,
 233 and amplify the eastward moving Undercurrent.

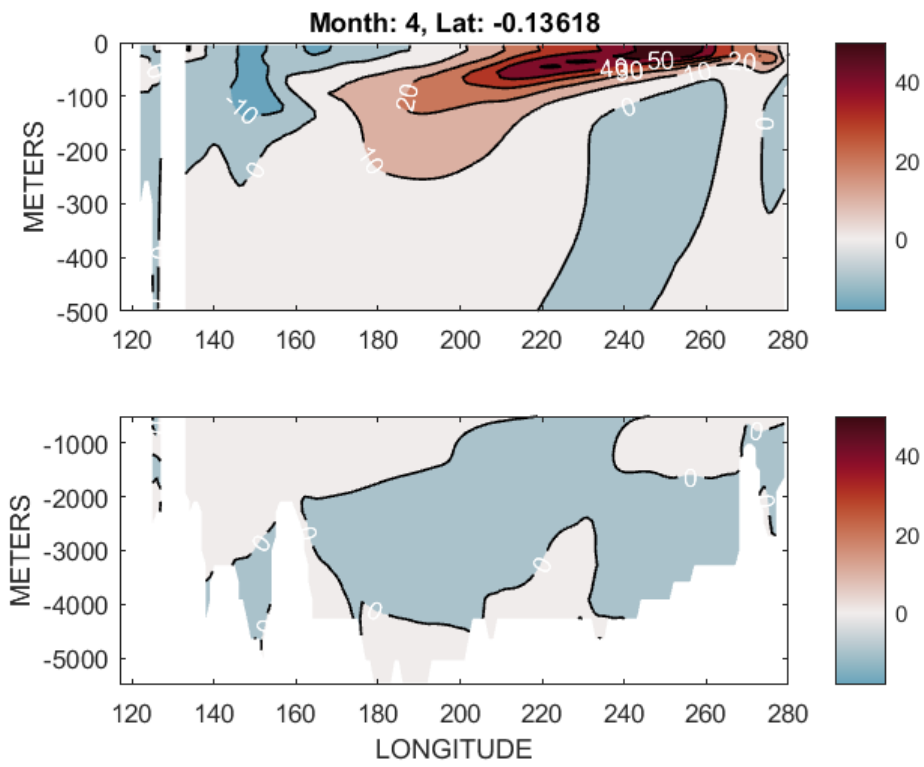


Figure 25: Zonal flow anomaly (cm/s) on the equator, mean April.

{equator_apr_s

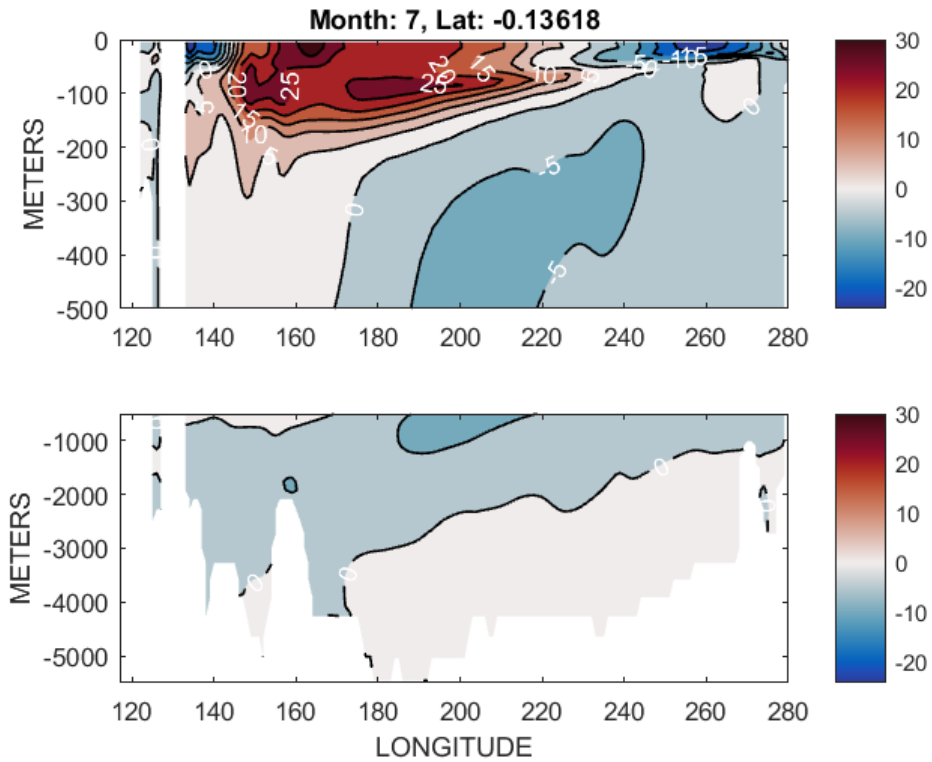


Figure 26: Zonal flow anomaly (cm/s) on the equator, mean July.

{equator_jul_s

234 3.2 Meridional Transports

235 One example of a 20-year time mean flow is shown in Fig. 28 at 30°S in the Pacific Ocean.
 236 These are readily computed monthly, seasonally etc. for any location.

237 When integrated through the entire longitude range of 360°, time-average oceanic mass
 238 conservation requires that the top-to-bottom meridional transports must vanish up to the di-
 239 vergence contained in net average evaporation plus runoff minus precipitation. The resulting
 240 global mean, accumulating integral is shown in Fig. 29. Residual imbalance, an estimate of
 241 the average evaporation minus precipitation appears in Fig. 30, but whose properties will be
 242 discussed elsewhere. An earlier result is by Stammer et al. (2004).

243 3.3 Property Transports

244 The state estimate provides a comprehensive set of output fields on the native grid which permit
 245 accurate property transport calculations, consistent with Griffies et al. (2016). As noted already,
 246 transport properties involving time mean products such as $\langle vT \rangle$ are expected to be different
 247 from values computed from the time means of each, $\langle v \rangle \langle T \rangle$. Thus Fig. 31 displays the depth,

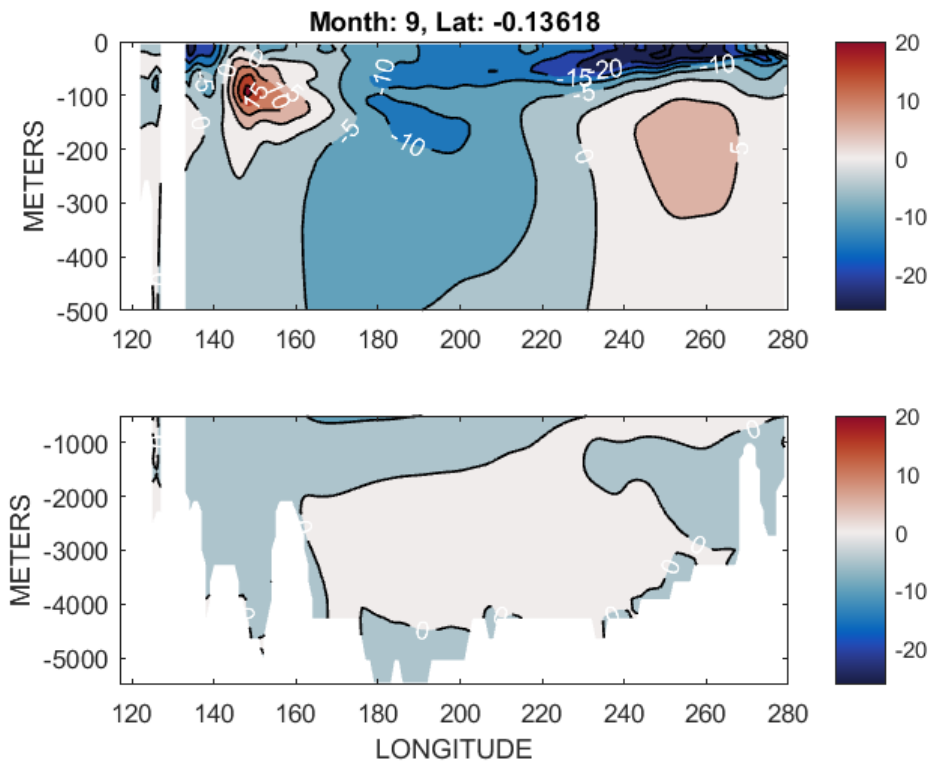


Figure 27: Zonal flow on the equator, mean September.

{equator_sep_s

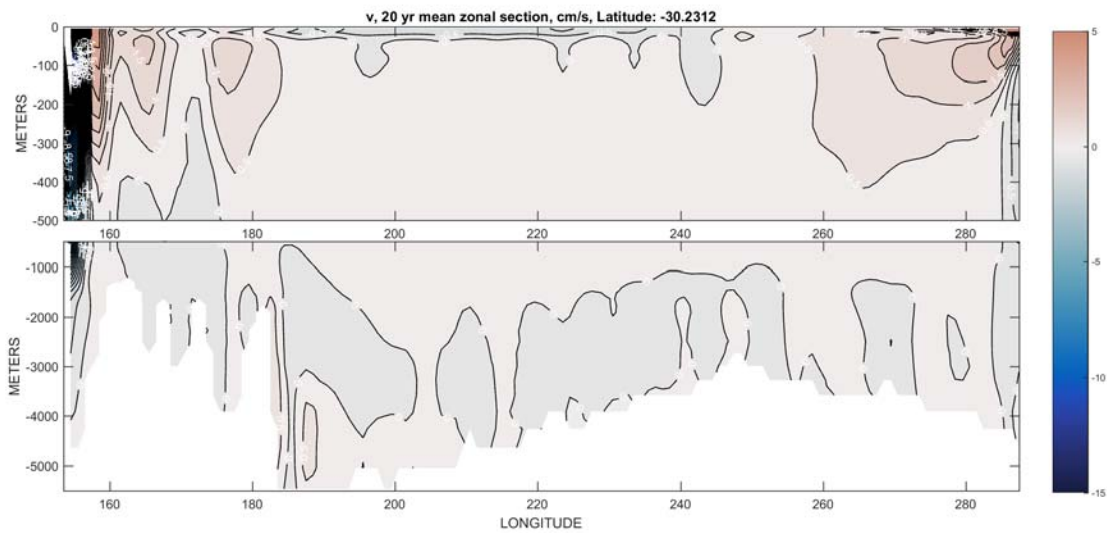


Figure 28: Twenty-year average meridional flow at 30°S in the Pacific Ocean. Intense flow in the East Australia Current and a flow reversing with depth along the coast of South America are visible. As in many such sections, weak deep flow reversals occur throughout.

{vn_20yrmean_3

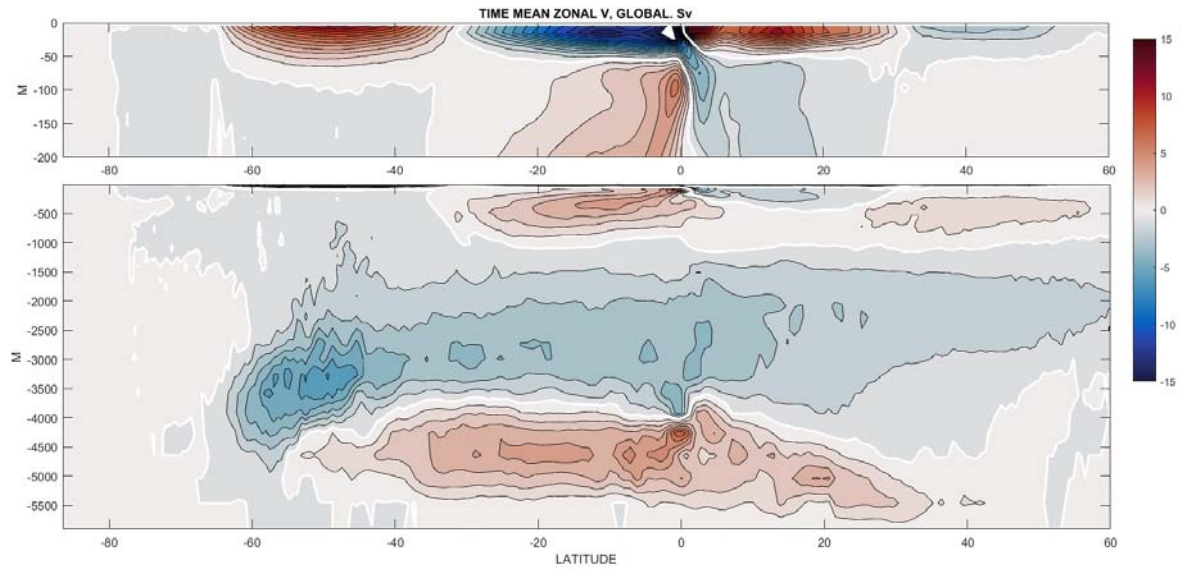


Figure 29: Zonal integral of vertically accumulating meridional transport in Sverdrups. (Not a stream function.) The values at the bottom necessarily almost vanish. See Fig. 30.

{zonal_integra

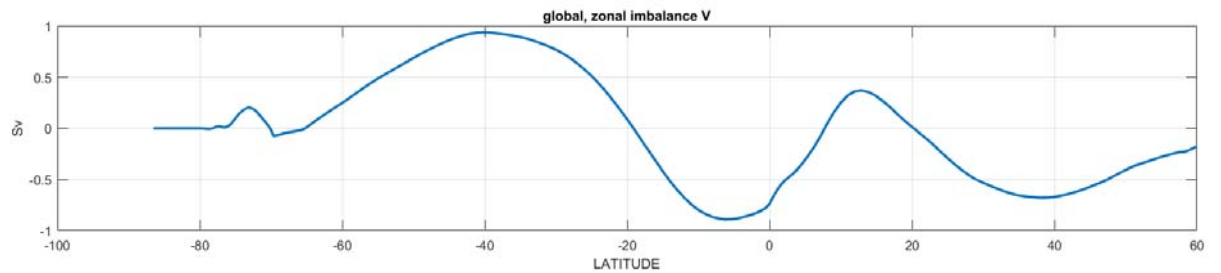


Figure 30: Integral, top-to-bottom, of the meridional transport as a 20-year mean. Bottom value of Fig. 29. Divergence is an estimate of the average evaporation minus precipitation.

{global_imbala

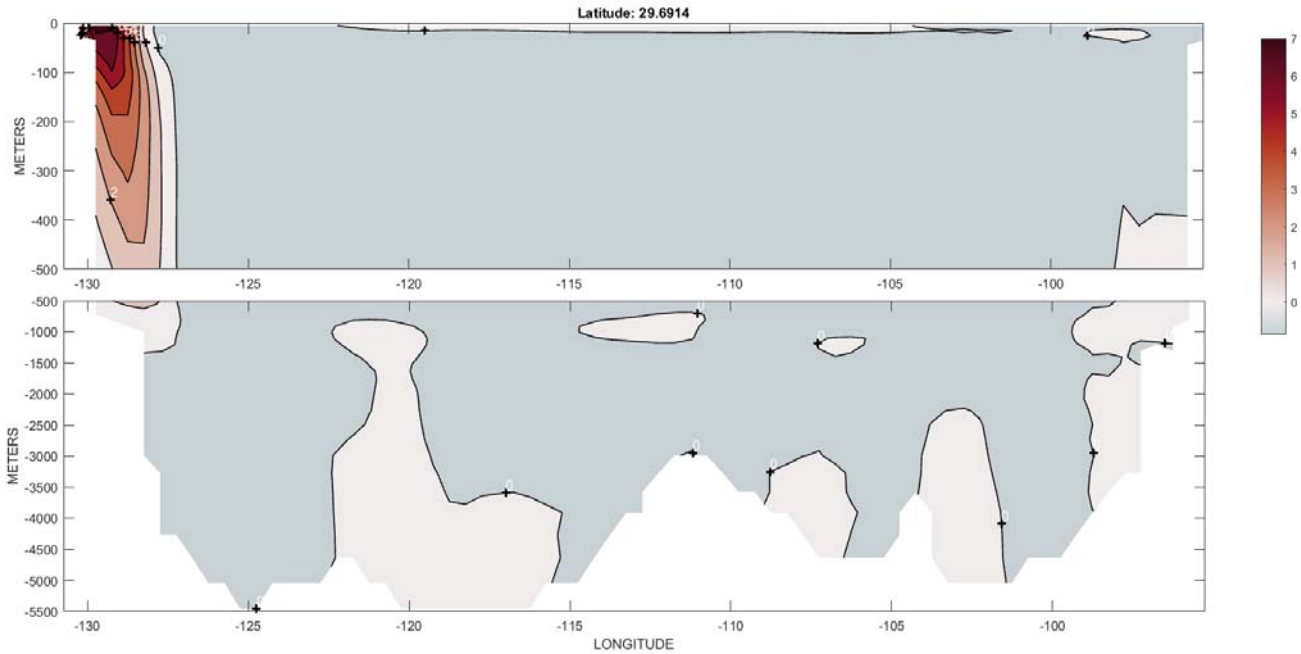


Figure 31: Product of the twenty-year means $\langle \bar{v} \rangle \langle \bar{T} \rangle$ at 30°N in the North Atlantic ($\text{m/s } ^\circ\text{C}$) with a reference temperature of 0°C . Corresponding heat transport is 0.6PW in contrast to values computed from quasi-synoptic sections of about 1.3PW (e.g., Bryden and Imawaki, 2003). Southward transport in the weak flowing interior is non-negligible.

{vn_theta_sect

248 longitude contributions of $\langle v \rangle \langle T \rangle$ 30°N in the North Atlantic, producing an equivalent heat
 249 transport of 0.6PW , smaller than estimates based e.g., on monthly or single section data (e.g.,
 250 Bryden and Imawaki, 2001; Picuch and Ponte, 2012, Table 2). As with many of the multi-
 251 decadal results, these values are best interpreted as quantitatively descriptive, and as serving as
 252 tests of unconstrained results from different models.

253 The corresponding values in the Pacific Ocean at 30°N are negligible (not shown) with a
 254 northward temperature transport mainly in the Kuroshio nearly cancelled by the interior return
 255 flow.

256 4 Vertical Velocities

257 *Eulerian Means*

258 Vertical velocities in the ocean are almost never measured directly, but must be computed
 259 diagnostically from the horizontal flow divergences. The result for the 20-year average at 105m
 260 can be seen in Fig. 32 and is a useful surrogate for the Ekman pumping. (See Roquet et al.,

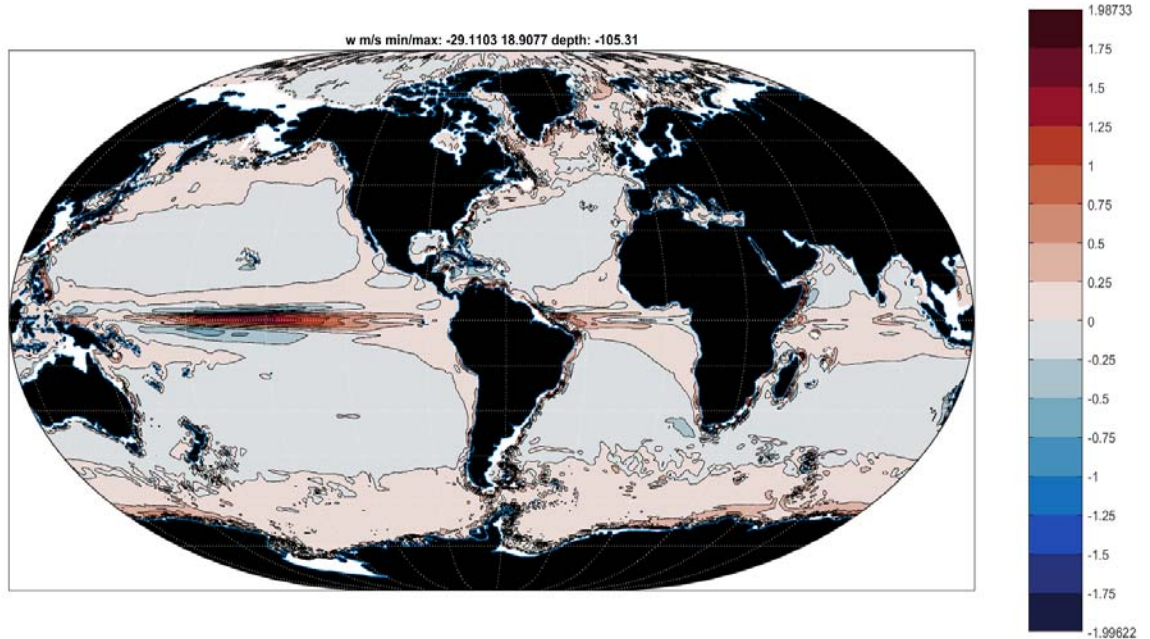


Figure 32: Twenty-year average Eulerian vertical velocity, w , (m/s) at 105m depth. Intense upwelling is apparent on the equator in all oceans, at high latitudes, and in traditional coastal upwelling regions.

{map_w_105m_20

261 2011 for an explicit discussion of the latter.) Main features are the subtropical and subpolar
 262 gyres as well as the powerful upwelling on the equator and the upwelling zones on the eastern
 263 margins. Fig. 33 shows the same result, but at 720m. At greater depths, e.g. 2000m (Fig.
 264 34), the influence of bottom topography has begun to dominate and the complexity of w defies
 265 simple description. Liang et al. (2017) provide a fuller discussion.

266 The mean annual cycle of w at 105m is shown in Figs. 35-38 and can be regarded as a
 267 quantitative estimate of the cycle in Ekman pumping.

268 5 Meteorological Variables

269 Meteorological forcing at the sea surface is part of the state estimate control vector—that is, the
 270 a priori windstress, surface air temperatures, specific humidity, shortwave downwelling radiation,
 271 and precipitation are modified along with other elements of the control vector so that the model
 272 is as consistent as possible with the oceanographic data. Comparatively small adjustments are
 273 made to the values obtained from the Dee et al. (2014) ERA-Interim atmospheric “reanalysis.”

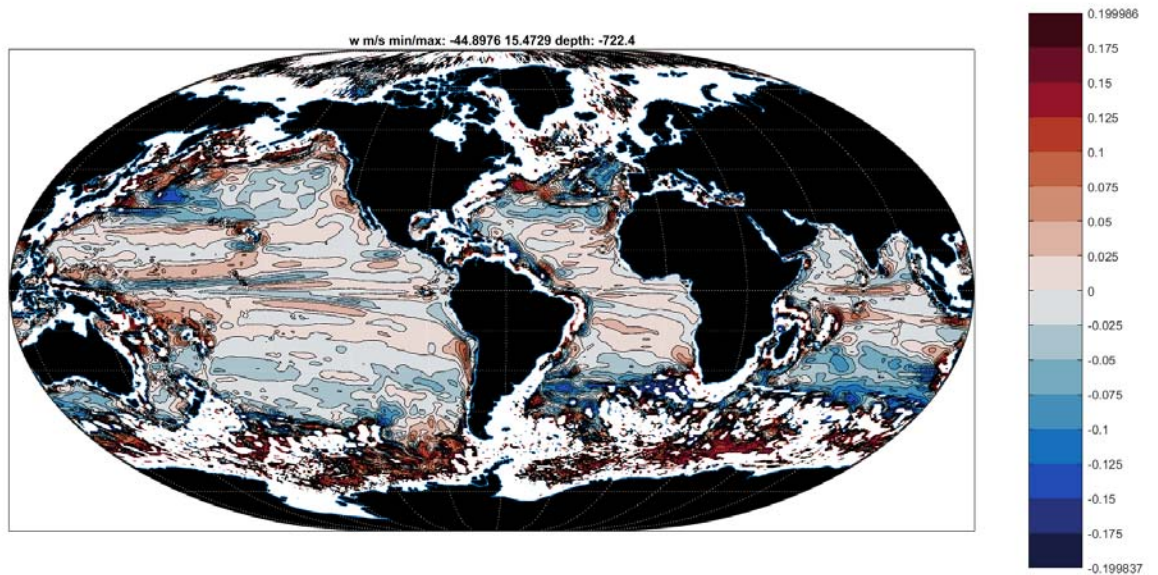


Figure 33: Twenty-year average vertical velocity, w , (10^5m/s) at 720m. The most conspicuous mid-latitude feature is the zonal banding, with a small residual of the large-scale surface gyres still visible. The Southern Ocean stands out as a region of extremely intense values of w of both signs (extreme values have been truncated there).

{map_w_720m_20

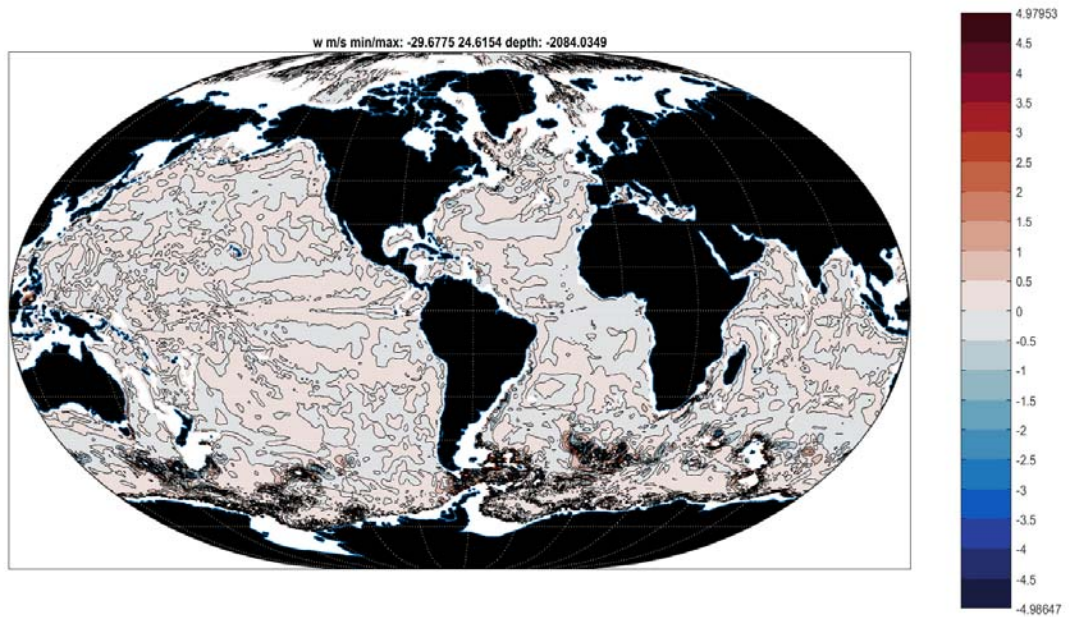


Figure 34: Twenty-year mean Eulerian w at 2100m (10^5m/s). At this depth, the complex structures induced by topography come to dominate the patterns. Some extreme values near topographic features have been omitted. See Liang et al. (2017).

{map_w_2084m_2

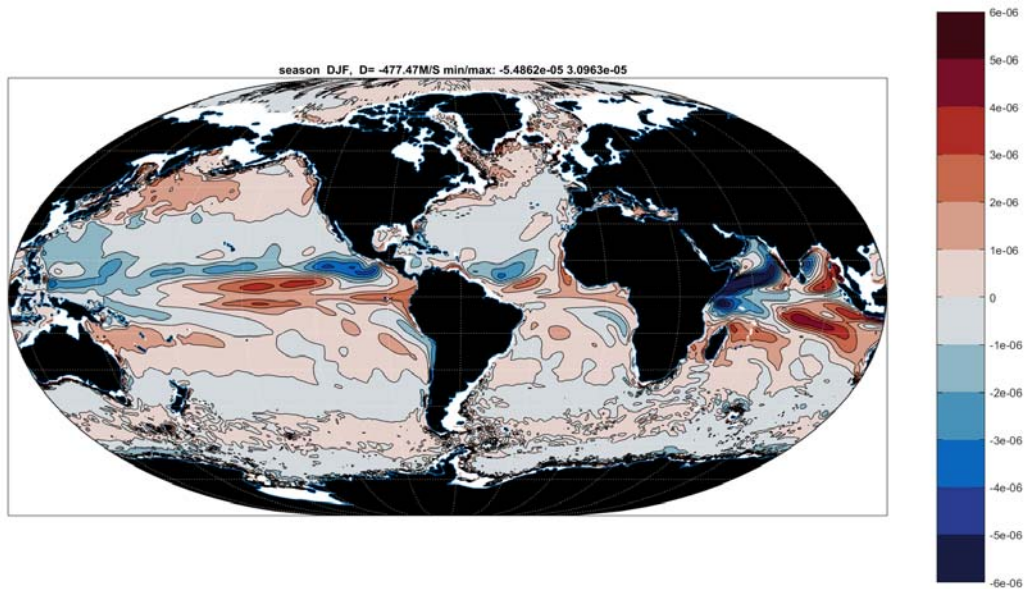


Figure 35: Twenty-year seasonal anomaly of w at 105m DJF.

{mapw_105m_sea

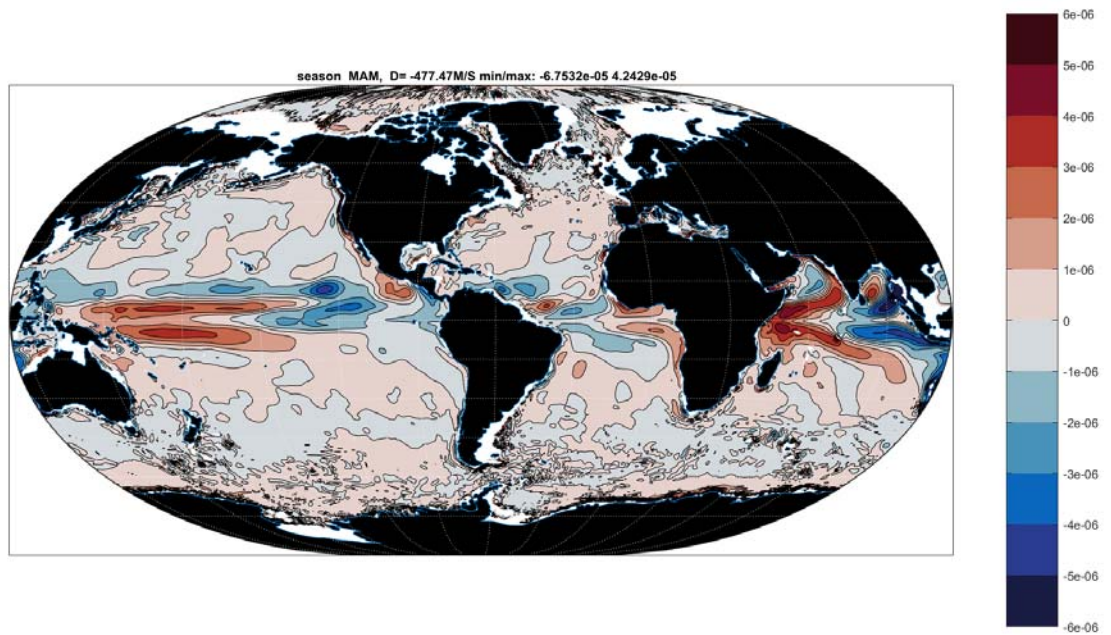


Figure 36: Anomaly of w , 105m March, April, May. (m/s, not multiplied by 10^5)

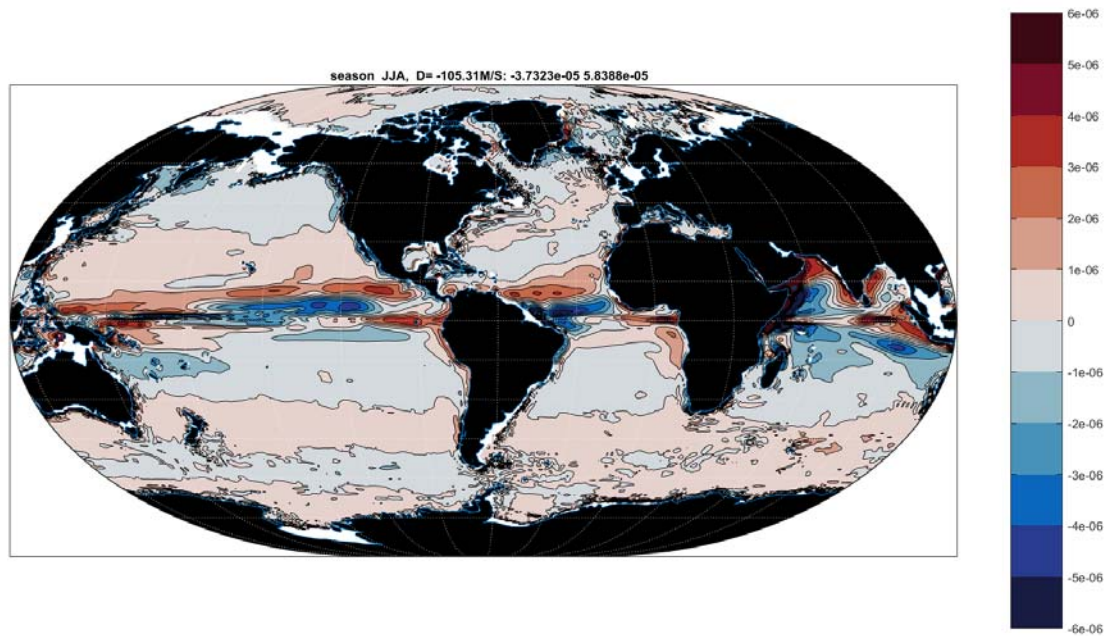


Figure 37: Anomaly of w (m/s) at 105m, June, July, August.

{mapw_105m_sea

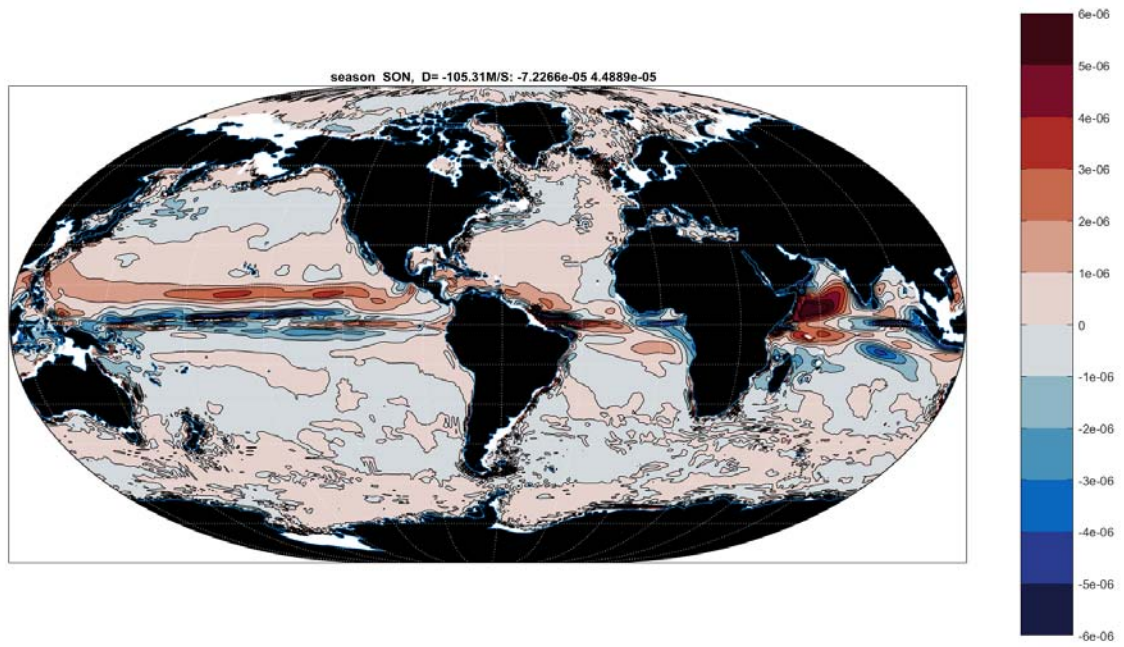


Figure 38: SON anomaly of $w,105$ m (m/s).

{mapw_105m_sea

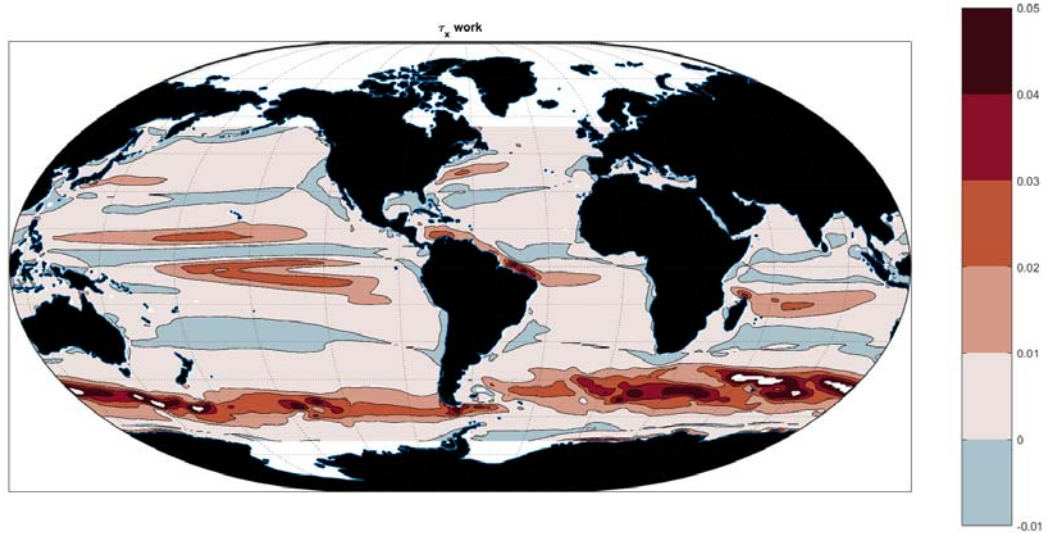


Figure 39: Twenty-year average misfit (here the inferred *correction*) to the time-mean τ_x (N/m^2). The state estimate is obtained by correcting the time-dependent Dee et al. (2014) estimates by a time-varying version of this correction when the model is run forward.

{misfit_tau_x_m

274 That reanalysis is not provided with explicit uncertainty estimates, but these have been discussed
 275 by Chaudhuri et al. (2014, 2016).

276 The adjustment (the “misfit” to the reanalysis) to the separate zonal and meridional esti-
 277 mates (τ_x, τ_y) are displayed in Figs. 39, ?? for the 20-year average. A generalization is that
 278 fitting to oceanic data strengthens both components of τ at high latitudes, and tends to weaken
 279 them in the subtropics and tropics. The global realism of these adjustments remains to be
 280 tested. Similar charts can be made for monthly, annual, or seasonal, etc. misfits.

281 The 20-year average wind-stress as adjusted by the state estimate calculation is shown in
 282 Fig. 41. On the large-scale the conventional easterly and westerly wind bands are all prominent.
 283 Its curl is shown in Fig. 42 and can be compared to Fig. 32, keeping in mind that the Ekman
 284 pumping, $w_E = \nabla \times (\tau / \bar{\rho} f)$.

285 The rate of wind working on the surface flow (not just the geostrophic component) is readily
 286 computed from the products $W_x^{(1)} = \langle \tau_x \rangle \langle u(z=5) \rangle, W_y^{(1)} = \langle \tau_y \rangle \langle v(z=5) \rangle$ in Figs. 43, 44 al-
 287 though as discussed earlier, these are only a part of the respective second order products $\langle \tau_x u \rangle,$
 288 $\langle \tau_y v \rangle$, and can only be interpreted as the work done by the mean wind on the mean surface flow.
 289 Omitting high ice-covered latitudes, thus the spatial average value is $W_x^{(1)} = 0.0043 \text{ W/m}^2$ and
 290 $W_y^{(1)} = -0.00025 \text{ W/m}^2$ which integrate to a total rate of working of about 1.6 TW. Monthly or

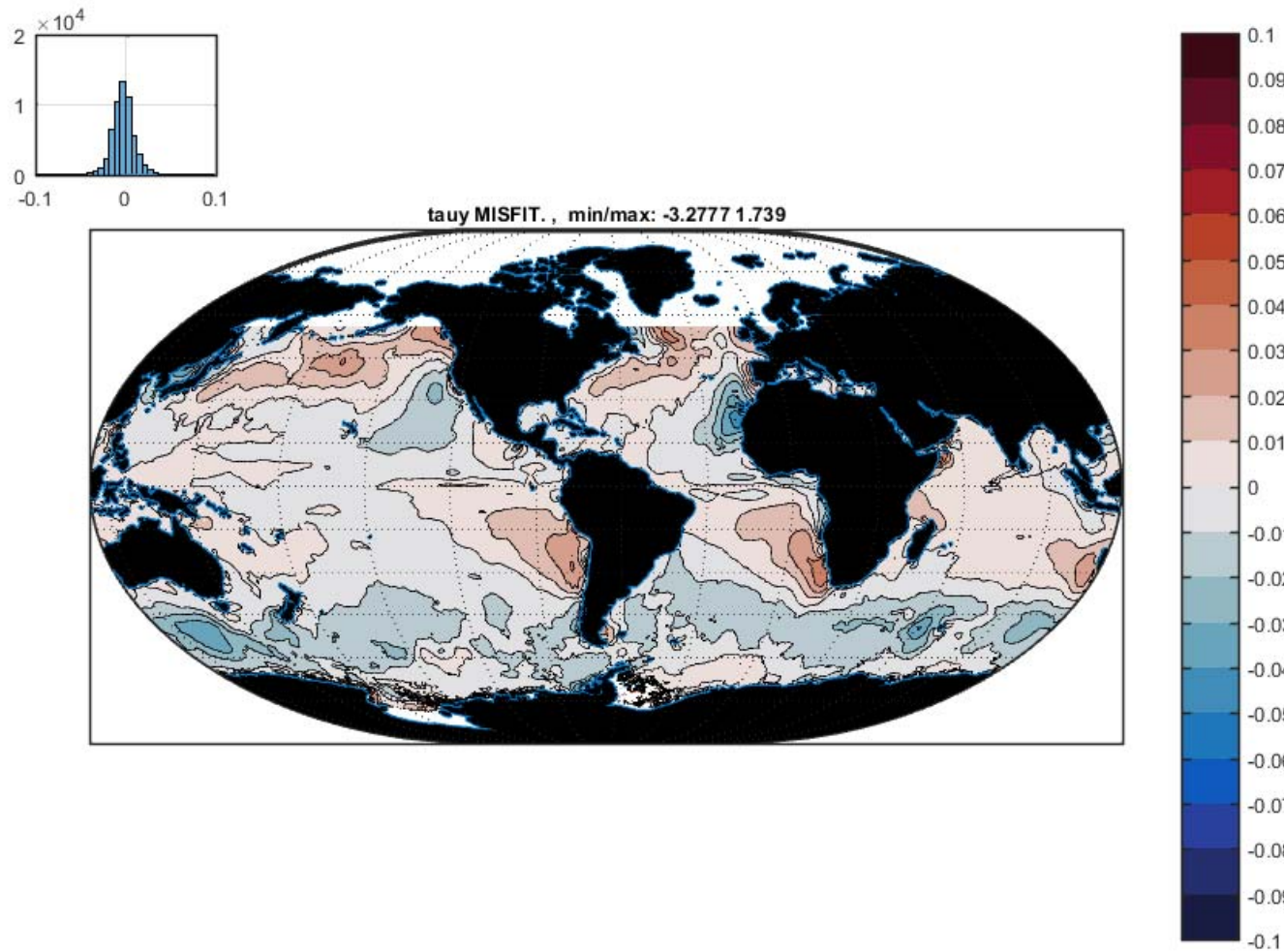


Figure 40: Twenty-year average “misfit” or correction to the time-mean τ_y (N/m^2).

{misfit_tauy_2

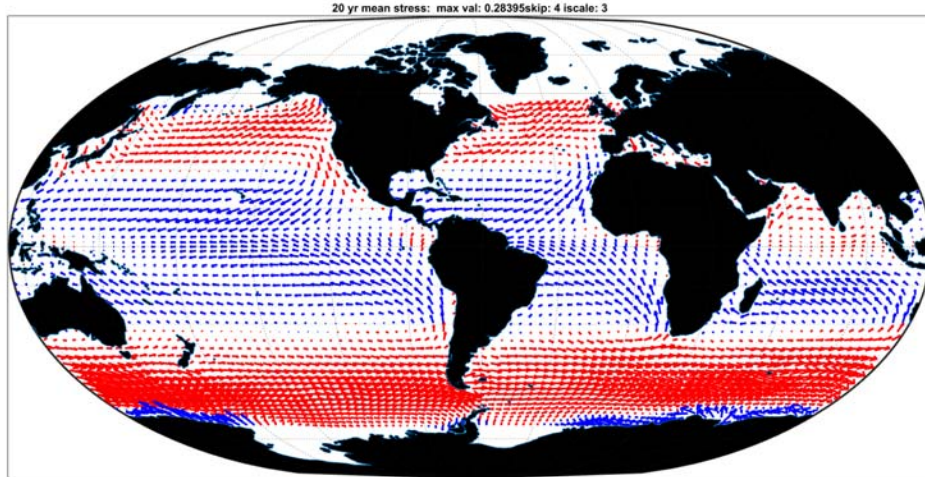


Figure 41: The 20-year average wind stress vectors (N/m^2) after adjustment by the state estimate calculation.

{quiver_tau_ar

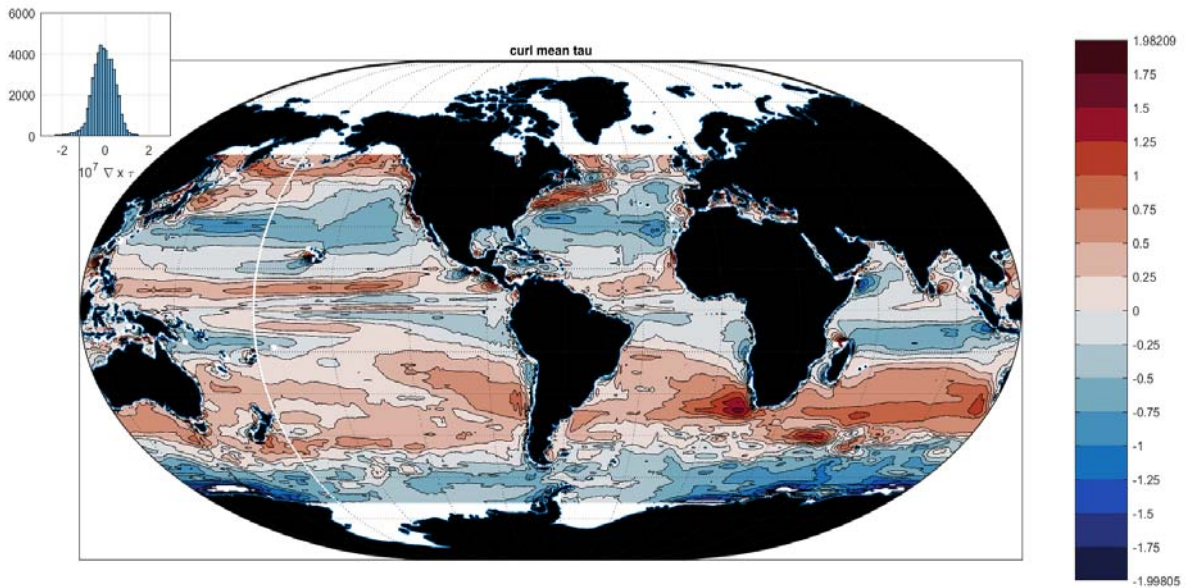


Figure 42: Vertical component of the curl of the 20-year average wind stress in Fig. 41.

{curl_20yearme

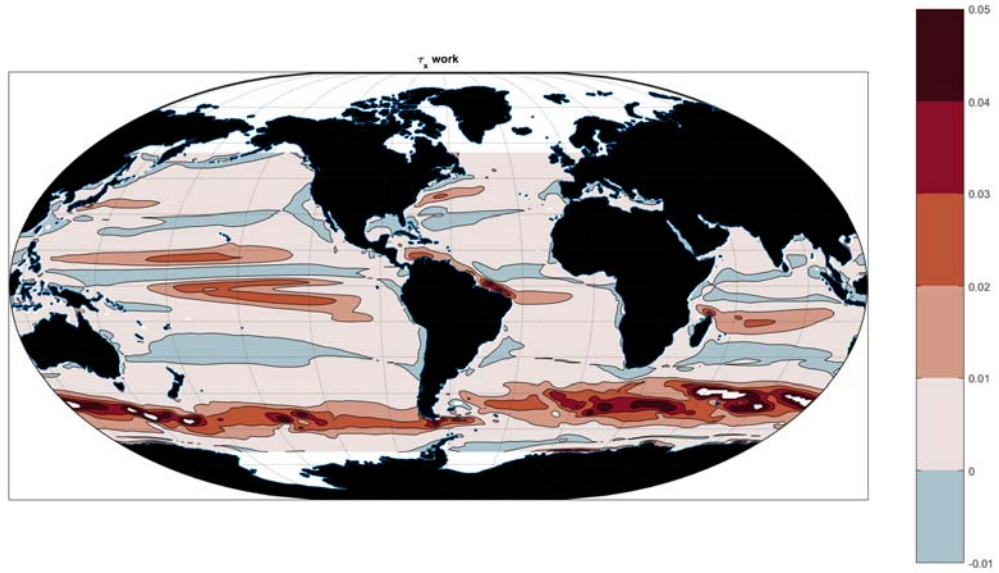


Figure 43: Wind work by the 20-year zonal average wind on the 20-year average surface velocity. (W/m^2) {taux_work_map}

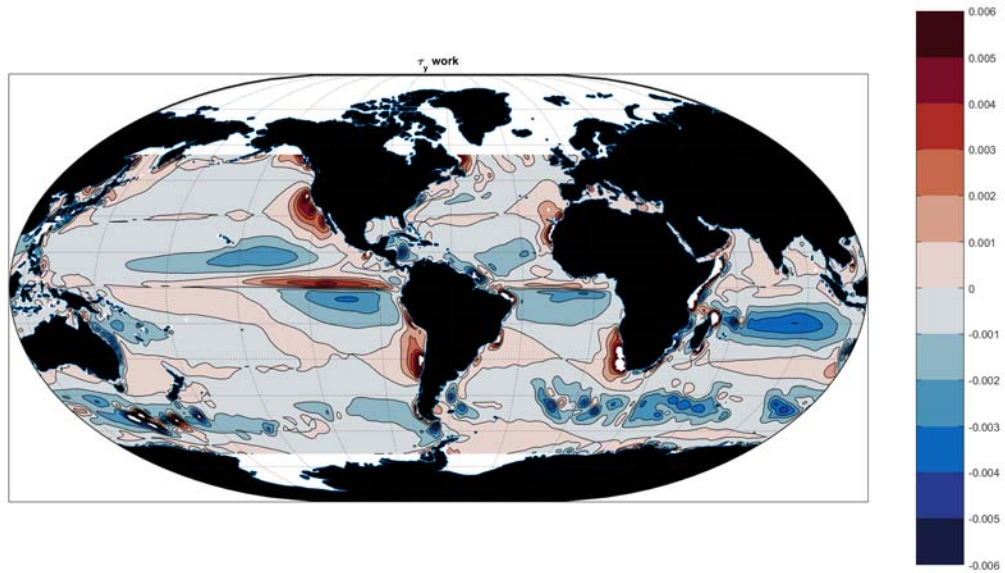


Figure 44: Rate of work on the time-mean sea surface velocity (W/m^2) of the meridional component of the wind stress. Note the change in scale from Fig. 43. Coastal upwelling regions tend to dominate. {tauy_work_map}

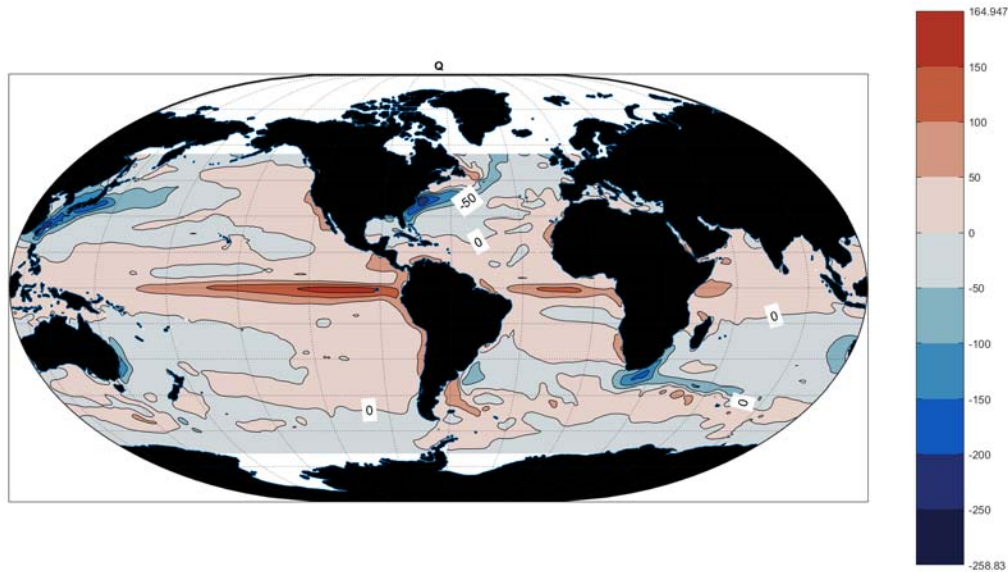


Figure 45: Twenty-year average estimated net heat exchange with the atmosphere (W/m^2) with positive values indicating a flux into the ocean.

{q_20yearmean.

291 seasonal or annual values of the rate of working can readily be computed from the climatology,
 292 but pursuit of this subject is left for elsewhere (see Zhai et al., 2012).

293 *Heat Exchange*

294 The 20 year average heat exchange, Q , with the atmosphere is depicted in Fig. 45 and its
 295 20-year average seasonal anomalies in Fig. 46. Qualitatively, these are all conventional, with
 296 heat gain in the tropics and major heat loss over the western boundary currents. Liang and Yu
 297 (2016) have compared these and related fields to reanalyses and OAFflux/CERES, showing a
 298 greater consistency with observations than do other estimates.

299 **6 Eddy Contributions**

300 As described by Forget et al. (2015), the model contains a variety of parameterizations intended
 301 to mimic the influence of eddies, waves and a variety of physical processes not properly resolved
 302 by the present model grid. Most of these formulas include empirical parameters varying horizon-
 303 tally, with depth, and in some cases, time. A full depiction of all of them would be overwhelming
 304 in the present context. As one example of what is now possible, Fig. 47 depicts the so-called bo-
 305 lus velocity at 722m derived from the Gent and McWilliams (1990) parameterization (cf. Ferrari

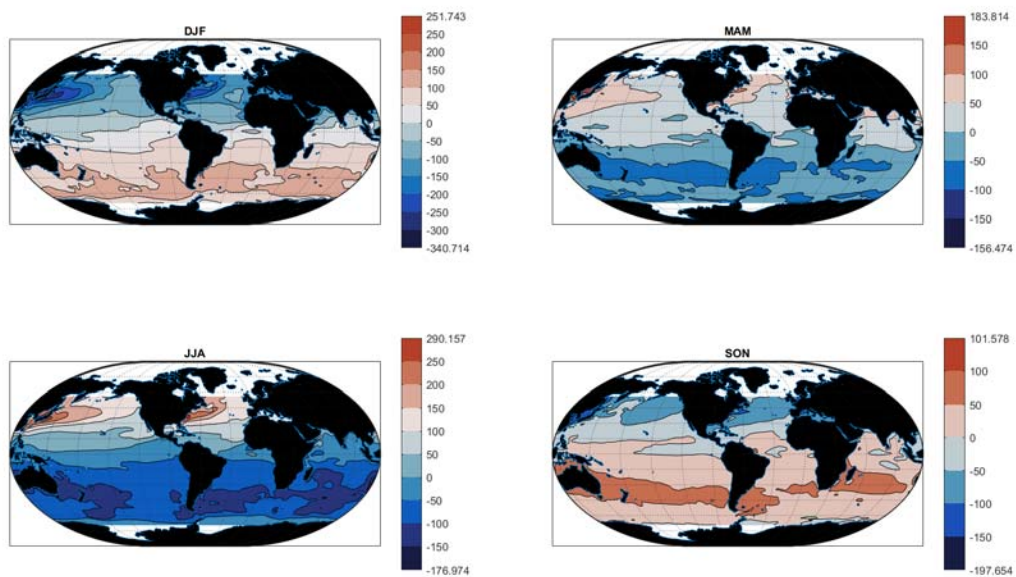


Figure 46: Anomaly of Q (W/m²) by season. Note changes in color scales.

{q_anom_4seasons}

306 and Plumb, 2003; Ferreira et al., 2005; Young 2012). As expected, a complex pattern results,
 307 one dependent upon the stability properties of the parameterized eddy field. On average, as
 308 compared to the Eulerian mean velocities, the relative kinetic energy in the bolus velocities is
 309 very small (about 0.5%) of the total. These results too, vary with year, month etc., but are not
 310 further displayed here.

311 *Acknowledgments.* Supported by NASA through the ECCO Consortium through contracts
 312 with MIT, AER and JPL.

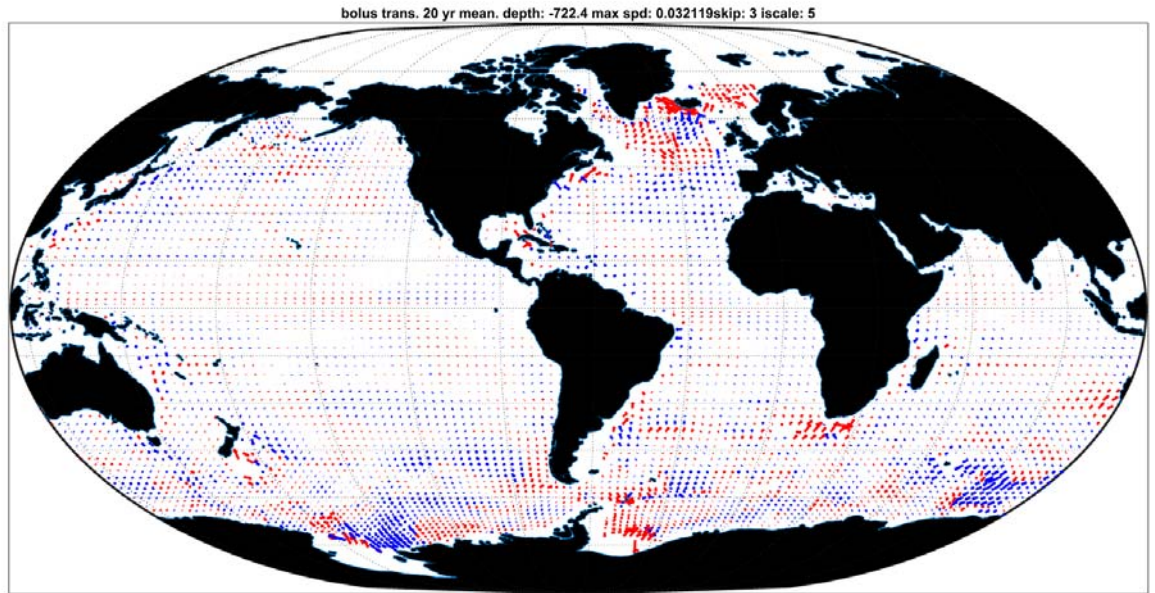


Figure 47: The time mean bolus velocities (u_{bolus}, v_{bolus}) at 722m (m/s).

{quiver_bolus_

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