

# THE VISCOSITY OF HUMAN BLOOD AT HIGH HEMATOCRITS

by

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S.B. (1981) Massachusetts Institute of Technology

SUBMITTED TO THE DEPARTMENT OF  
CHEMICAL ENGINEERING IN PARTIAL  
FULFILLMENT OF THE REQUIREMENTS FOR THE  
DEGREE OF

MASTER OF SCIENCE

IN CHEMICAL ENGINEERING

at the

MASSACHUSETTS INSTITUTE OF TECHNOLOGY

February 1986

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## **The Viscosity of Human Blood at High Hematocrits**

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**James Douglas Oliver III**

Submitted to the Department of Chemical Engineering on  
February 19, 1986 in partial fulfillment of the requirements  
for the degree of Master of Science in Chemical Engineering.

### **Abstract**

The viscosity of human blood was measured in both plasma and saline suspensions with a Haake Rotovisco RV100 viscometer. Hematocrits of samples ranged from below physiologic to greater than 98, and shear rates ranged from 0.03 to 300  $\text{sec}^{-1}$ . Even at high concentrations, blood retained its fluidity, with viscosity at high shear rates on the order of 50-100 cP. At high hematocrits, where cell-cell crowding is the dominant rheological mechanism, the type of suspending fluid had a negligible effect on the absolute viscosity, but led to an artificially higher relative viscosity for the saline suspensions. At lower hematocrits and lower shear rates, where aggregation and cell-protein interactions are important, the absolute viscosity of plasma suspensions was higher. The emulsive properties of the red cell, however, created a lower relative relative viscosity for plasma suspensions than for cells in saline. The data was fit to the three-parameter Quemada equation, and empirical curves were developed to represent blood viscosity as a function of shear rate and hematocrit over the range of experimentation. Examination of the rheological parameters provided some indication as to the thresholds of influence for the important mechanisms of blood viscosity. Most noteworthy, it appears that at hematocrits of 65 to 80 and higher, the erythrocytes begin to behave as a continuous phase and blood viscosity becomes independent of the suspending fluid. The mathematical expressions are of interest in the design of clinical cross-flow membrane plasmapheresis devices, in which high concentrations are established near the filter.

Thesis Supervisor:     **Clark K. Colton**  
                          Title:     **Bayer Professor of Chemical Engineering**

## Acknowledgments

This is the hardest part to write, where I must attempt to thank within the space of a few paragraphs the people who have made this possible. I gratefully acknowledge the support of the GEM, Prince, and Kleberg Foundations and the Amoco Oil Corporation during the work of this thesis.

Many thanks to Professor Clark K. Colton, for his guidance and patience and his insistence on quality. Thanks also to Gail Gamboa and Barbara Driscoll for their help and for making life so pleasant around the office.

Thanks to Andrew Zydney for his always excellent advice and the best of luck to him in his career. Also to W. Mark Saltzman, fellow vampire, for sharing jokes and saving me a lot of time along the way. Mark Green of the Ceramics Laboratory was a big help in mastering the vagaries of viscometry, even if he might have gotten queasy at the sight of blood. David Yarmush's assistance with albumin labeling was greatly appreciated.

The Medical Engineering/Medical Physics program of the Division of Health Sciences and Technology consistently admits an odd assortment of people who know how to make long hours pass quickly. If they ever graduate, there will be some serious damage inflicted on the world. I also thank countless friends at M.I.T. and Harvard Medical School, too many to name here.

To my many teachers, and especially to Mrs. Bright, who kept me off the streets, I give my eternal gratitude.

Christine and Ed have put up with me for a long time and will have to continue to do so. The chances are that they'll probably never read this, so I'll admit that they have been better friends than I deserve.

My little sisters, Jeanne and June, probably don't realize how old and wise they've gotten over the past few years, or how much I depend on them, but if they ever find out, I'm in big trouble.

And finally, to my parents, for whom the words just do not exist.



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## Chapter 1

# INTRODUCTION

The study of the viscosity of blood traces its roots back to the dawn of rheology, to the pioneering days of Poiseuille [1843], and with advances in scientific instrumentation reached its peak in the 1960s with the numerous works of Merrill, Cokelet, Dintenfass, Copley, Chien, and many others. In more recent times, attention has turned to the microscopic phenomena underlying the observed physical behavior, such as the properties of the red cell membrane and the mechanisms of the cell-cell and cell-fluid interactions [Skalak *et al.*, 1981a and 1981b; Secomb *et al.*, 1983a and 1983b]. Despite the complex rheological nature of blood suspensions, they in many ways are an ideal system for the study of non-Newtonian fluid mechanics: The particles are well-defined and almost identical to one another, they deform in a well-characterized manner, and the extent of their aggregation and deformation can be controlled through the system properties. The fluid-like nature of the erythrocytes and their aggregability give blood suspensions properties similar to emulsions and polymers.

Since most blood viscosity work has been carried out with an eye toward understanding physiological and clinical-pathological conditions, the vast majority of data exists at cell concentrations around 40-50 volume per cent. A limited amount of data for concentrated ( $>80\%$ ) suspensions exists (see Chapter 2) which was gathered primarily as a means of assessing properties of the red cell, but which also had applications to the pathophysiological effects of polycythemia, where hematocrits of up to 80 to 85 per cent may be recorded at high altitudes [Harris and Kellermeyer, 1970].

Recent research on extracorporeal blood treatment has provided a new incentive for a greater understanding of blood viscosity in concentrated suspensions. An analysis of the ultrafiltrate flux and blood flow characteristics of cross-flow membrane plasmapheresis devices has led to conclusions that the local cell concentration near the filter is at least 95% and may be up to or greater than 98% [Zydney and Colton, 1982; Zydney, 1985]. In order to generate accurate expressions for the mass and momentum profiles, it is necessary to expand the present data base at high hematocrits and to develop a means for representing blood viscosity as a function of both hematocrit and shear rate over the broadest possible range of applicability.

The scope of this thesis, then, will be to present systematic work on blood and red cell suspension viscosity as a function of hematocrit with particular emphasis on hematocrits above 90 and shear rates in the region of interest for plasmapheresis (50-500  $\text{sec}^{-1}$ ). The data will then be fitted to mathematical models from the literature and combined with empirical formulations of the relevant parameters in order to derive an expression that will provide general usage. By examination of the parameter variation with hematocrit, an assessment of the microscopic behavior of red cell rheology can be made.



## Chapter 2

### BACKGROUND

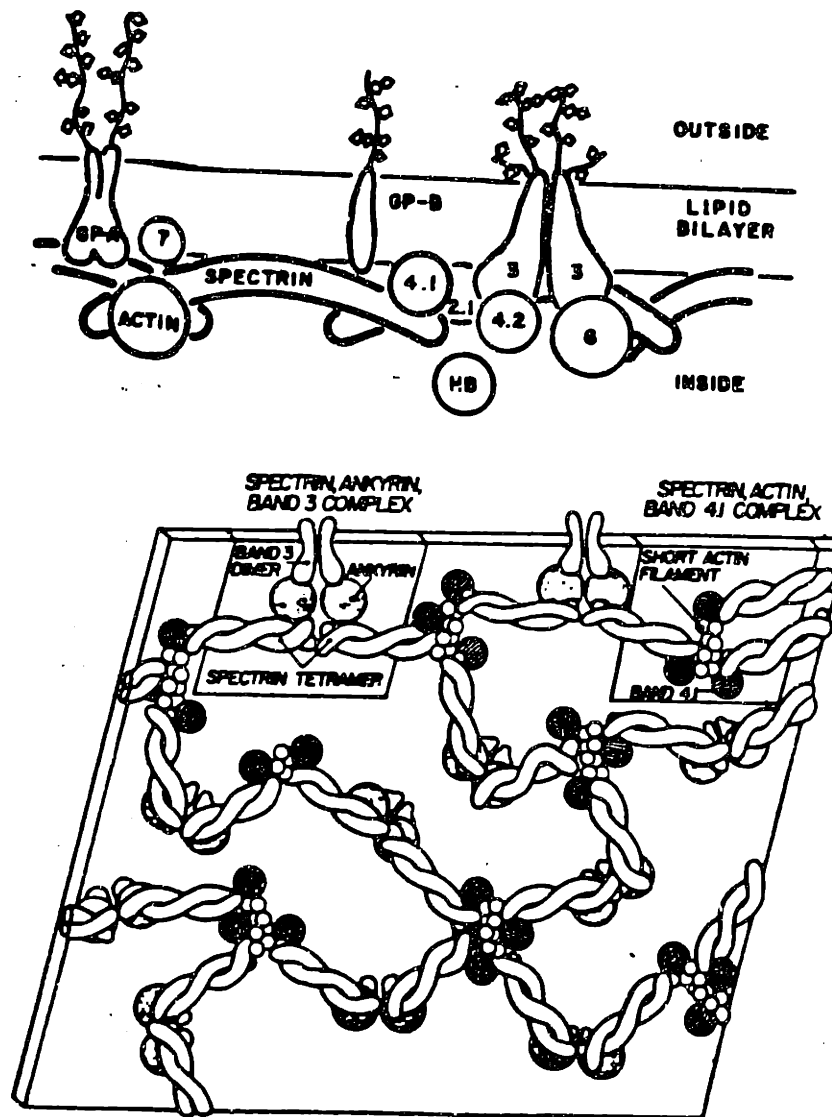
#### 2.1 Composition of Blood and Structure of the Red Blood Cell

Blood is composed of the corpuscles--erythrocytes (red blood cell, RBC), leukocytes (white blood cells, WBC), and platelets--suspended in a fluid medium, the plasma. Erythrocytes in normal human blood make up greater than 99 per cent of the corpuscular volume, outnumbering leukocytes by 500-1000 to 1 and platelets by approximately 20 to 1. Therefore, the cellular contributions to the physical characteristics of blood come almost exclusively from the red cells.

By weight, the plasma is 90 per cent water, 7 per cent proteins, 1 per cent inorganic salts and 2 per cent low molecular weight organic solutes. The plasma proteins are principally albumin (4.5 gm.%, mol. wt. = 69,000), globulins (2.5 gm.%, mol. wt. = 35,000 to 1,000,000), and fibrinogen (0.3 gm. %, mol. wt. = 330,000).

The red blood cell under normal conditions is typically a biconcave disk 6-9  $\mu\text{m}$  in diameter, 1  $\mu\text{m}$  thick in the center and increasing to 2-2.5  $\mu\text{m}$  at the periphery. The average cell volume is 90  $\mu\text{m}^3$  and the surface area is about 140  $\mu\text{m}^2$ .

The red cell membrane is believed to be a two-layer structure consisting of an inner, rigid shell of proteins covered by a more fluid phospholipid bilayer (see Figure 2-1). Choline phospholipids (phosphatidyl choline and sphingomyelin) are found primarily in the outer surface of the bilayer, while the inside half is composed mostly of amino phospholipids (phosphatidyl ethanolamine and phosphatidyl serine).



**Figure 2-1:** Schematic of the red cell membrane.

*Top*, cross-section illustrating lipid bilayer, integral membrane proteins, and cytoskeleton proteins. *Bottom*, structure of the membrane skeleton.  
From Lux, [1981].

Interspersed throughout the bilayer are molecules of glycolipids and unesterified cholesterol along with integral membrane proteins which penetrate or span the bilayer. Peripheral membrane proteins (spectrin, ankyrin, red cell actin, and band 4.1) bind to or are absorbed to the integral proteins and to each other, providing an internal microfibrillar "cytoskeleton" for the membrane. The exact nature of these connections is the subject of much current research but it appears that a complex of spectrin-actin-band 4.1 is responsible for the intraskeletal connections while a spectrin-ankyrin-band 3 complex provides the bonds to the lipid bilayer [Lux, 1981].

It has been proposed that the red cell membrane can be modeled as a solid-liquid composite [Evans and Hochmuth, 1977]. The solid cytoskeleton provides the membrane's resistance to shear deformation at constant surface area. It is also responsible for the elastic reversibility and the yield behavior of the membrane [Evans and Hochmuth, 1978]. The liquid-like lipid bilayer resists changes in membrane surface area. The ratio of the elastic modulus ( $\equiv \frac{\Delta force}{\Delta dimension}$ ) for surface area changes to that for shear deformation at constant surface area is on the order of  $10^4$ - $10^5$  [Waugh and Evans, 1976; Evans *et al.*, 1976].

The cytoplasm of the erythrocyte is predominantly a 34 per cent hemoglobin-in-water solution. Adult hemoglobin (hemoglobin A) consists of four globin polypeptide chains each bound to a separate heme group. It has a molecular weight of 64,458. The main function of the red blood cell is to transport oxygen bound to the hemoglobin. The physical characteristics of blood and of the red blood cell are summarized in Appendix A.

Blood is stored in liquid state at 4 °C under mandate from the Food and Drug Administration. The time limits for classification of indated blood are also under federal regulations. For citrate phosphate dextrose, the time is 21 days; for citrate phosphate dextrose-adenine, it is 35 days; and for adenine-saline, it is 42 days.

While stored, blood undergoes a variety of changes which include loss of platelets, granulocytes, plasma proteins, enzymes, and seepage of intracellular electrolytes [Sherwood, 1981]. Additionally, as cells begin to age, the content of membrane lipids decreases and the mean corpuscular hemoglobin concentration increases. These factors increase the viscosity of erythrocytes with time [Erslev and Atwater, 1963; Usami *et al.*, 1971].

## **2.2 Rheological Behavior of Blood**

### **2.2.1 Rheological Properties of Plasma**

The earliest studies of plasma viscosity indicated that it was a non-Newtonian fluid with shear-thinning properties [Copley *et al.*, 1960; Wells and Merrill, 1961; Cerny *et al.*, 1962], but subsequent data [Charm and Kurland, 1962; Copley and Scott-Blair, 1962; Merrill *et al.*, 1963a] showed that indeed plasma did conform to Newtonian behavior, with a viscosity of 1.6 cP at 20 °C. The apparent shear rate dependence of the original measurements was attributed to interactions at the fluid-air interface of cone-and-plate viscometers and of Couette viscometers without guard rings [Copley, 1971; Copley and King, 1972]. This will be addressed later in the section on viscometry (Section 3.3).

Serum, the fluid remaining after blood is allowed to clot *in vitro*, is essentially a defibrinogenated plasma solution. It shows none of the non-Newtonian behavior attributed to plasma's surface effects. Based on this observation, it has been suggested that the plasma proteins and principally fibrinogen may form an interlocking network which is irreversibly degraded at shear rates greater than 5  $\text{sec}^{-1}$  [Charm and Kurland, 1974]. Plasma does exhibit some additional non-Newtonian characteristics such as drag reduction in turbulent flow [Chmiel, 1974].

Using a sphere-in-sphere viscometer to minimize surface effects, Chmiel and Walitza [1980] have confirmed Newtonian behavior and have described the temperature dependence of plasma viscosity by

$$\mu(T) = \mu(T_0) \cdot \exp\left(\frac{420(T_0 - T)}{(T + 162)(T_0 - 162)}\right) \quad (2.1)$$

where  $T$  is the temperature of interest in degrees Kelvin,  $T_0$  is the reference temperature at which the value  $\mu(T_0)$  is known. This temperature dependence is much the same as that of water.

The effect of protein concentration on plasma viscosity has been looked at in some investigations, but has not been well characterized [Eastham and Morgan, 1965; Wells, 1965; Mayer, 1966; Mayer *et al.*, 1966; Rand *et al.*, 1970]. Serum typically has a viscosity 20% less than that of plasma. Bayliss [1952] proposes an equation for correlating plasma viscosity  $\mu_p$  to total protein concentration

$$\frac{\mu_p}{\mu_w} = \frac{1}{1 - b \cdot c_{\text{protein}}} \quad (2.2)$$

where  $\mu_w$  is water viscosity,  $b$  is a constant, and  $c_{\text{protein}}$  is the protein concentration in gms/100 ml. Eastham and Morgan [1965] found that increasing albumin concentration to 11 per cent increased viscosity from 1.6 to 2 centipoises at 25 °C while fibrinogen at 4 percent increased the viscosity to greater than 3 centipoises. MacKenzie and co-workers [1970] found a proportional increase in viscosity with an increase in globulin concentration. In summary, there appears to be an increase in viscosity with an increase in plasma proteins, although the viscosity is least sensitive to changes in albumin concentration. There is evidence that in some clinical pathological states, such as severe leukemia or macroglobulinemia, shear-thinning behavior of plasma may be present [Dintenfass, 1965b and 1966].

### 2.2.2 Rheological Behavior of Red Blood Cell Suspensions

Since the plasma is essentially Newtonian, and since leukocytes and platelets have a negligible effect on flow properties, the singularities of blood rheology stem from the behavior of the erythrocytes. At low shear rates the cell reversible form stacked primary aggregates, called rouleaux, which can in turn form larger secondary aggregates (see Figure 2-2). The rouleaux disintegrate with increasing shear. The mechanism behind the aggregation phenomenon is not yet understood, although electrostatic attraction, plasma surface tension, and molecular bridging between membranes have all been proposed [Fahraeus and Lindqvist, 1931; Casteneda *et al.*, 1965; Merrill *et al.*, 1966]. Experiments suggest that fibrinogen concentration is the most important determinant of cellular aggregation [Merrill, *et al.*, 1963b and 1965b].

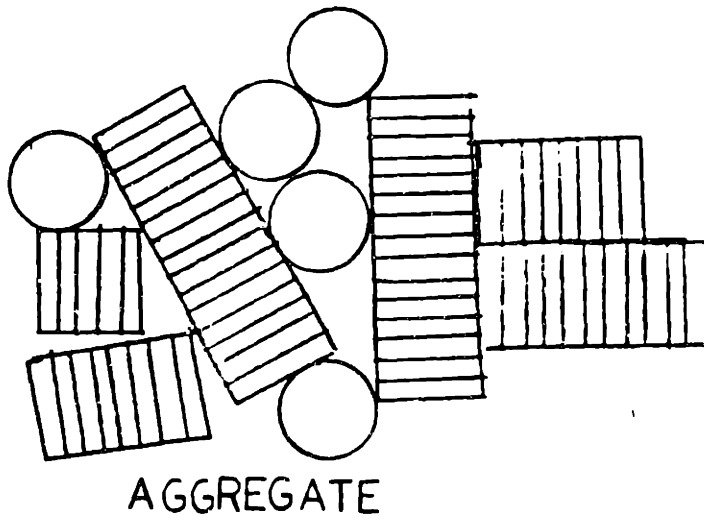
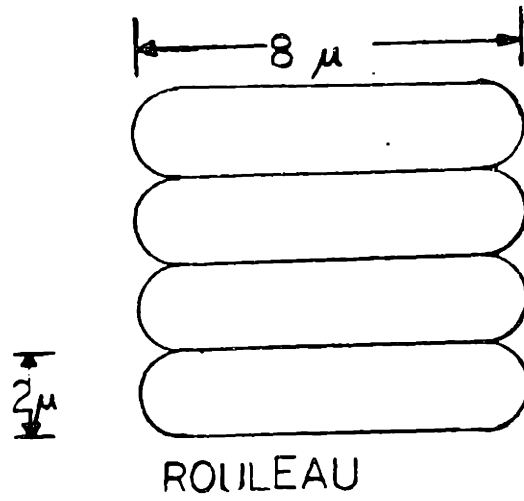
Skalak and co-workers [1977] have developed a general formulation for energy conservation in rouleau aggregation and disaggregation:

$$\frac{dW}{dt} = \frac{dE}{dt} + \frac{dK}{dt} + \frac{dD}{dt} - \gamma \frac{dA_c}{dt} \quad (2.3)$$

where  $W$  is work done by external forces (surface and body);  $E$  is elastic stored energy (strain energy);  $K$  is kinetic energy;  $D$  is dissipated energy (including viscous energy of solids, fluids, and the process itself); and  $-\gamma \frac{dA_c}{dt}$  is the rate of energy required to separate the contact surfaces, with  $\gamma$  being the surface energy and  $A_c$  the common contact area. For static equilibrium it can be shown that

$$\gamma = \frac{\partial E}{\partial A_c}. \quad (2.4)$$

Using a known strain energy function, the surface energy can be estimated to be in the range of  $10^{-4}$  dyne/cm [Skalak *et al.*, 1981b]. There is presently no method of proceeding from equation (2.3) and deriving an expression for rouleau formation and disaggregation in dynamic states [Skalak *et al.*, 1981a].



**Figure 2-2:** Red blood cell rouleau and aggregate.

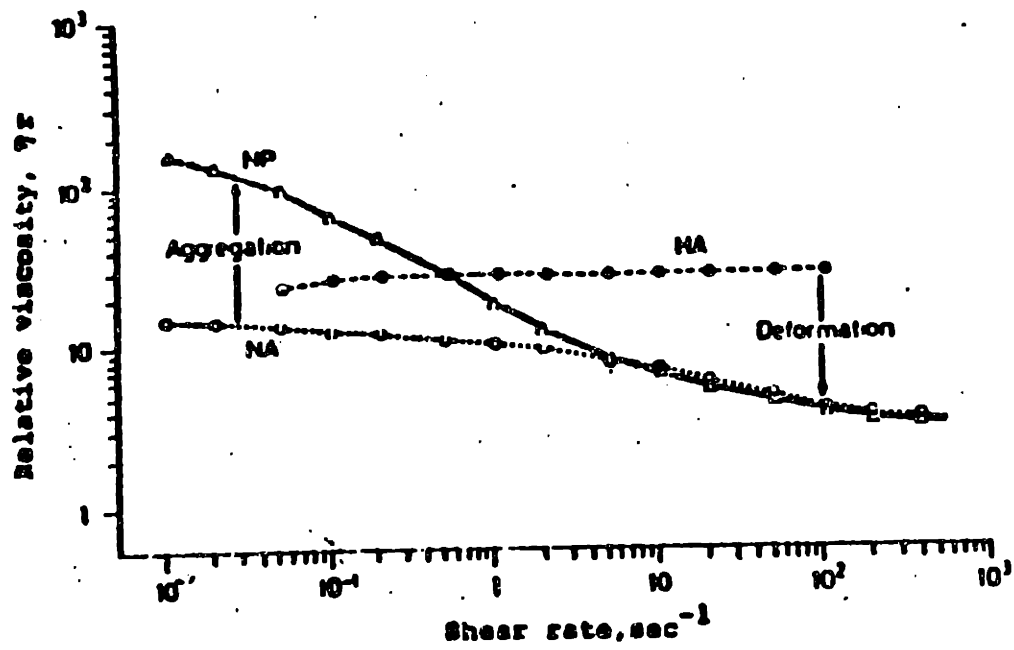
At very high shear rates, on the order of  $50-100 \text{ sec}^{-1}$ , the cells undergo deformation and orient themselves with the flow in order to achieve the smallest possible effective volume. Since the cell membrane is highly resistant to changes in surface area but relatively receptive to shear deformation, they assume the shape of prolate ellipsoids. The membrane rotates cyclically around the cellular contents in a "tank-treading" motion, which establishes an internal velocity profile in the cytoplasm [Schmid-Schonbein *et al.*, 1971]. The circulation of fluid inside the membrane has the effect of aligning the cell along the external streamlines and it also reduces the effective volume of the cell. The increase in deformation with shear for the red cell follows the theoretical development for the behavior of a fluid drop developed by Rumscheidt and Mason [1961b].

The effects of these structural changes on the viscosity are shown schematically in Figure 2-3, from Chien [1970], where three sets of data are represented. One curve is for normal blood in plasma, one is for normal blood cells in an 11% albumin solution and the third is for gluteraldehyde-hardened blood cells in 11% albumin.

The albumin solution lacks the proteins necessary for aggregate formation, principally fibrinogen. At low shear rates, therefore, the difference in viscosity between the plasma-suspended RBC and the albumin-suspended normal RBC is due to the effects of aggregation. Similarly, at high shear rates, the difference between the hardened RBC viscosity and the normal RBC viscosity represents deformational effects. The hardened RBC in albumin solution undergo no microstructural changes and thus viscosity is relatively constant with shear rate, with a degree of shear-thickening behavior common to suspensions of rigid particles or of polymers.

For suspensions of rigid particles, the viscosity of the suspension is directly proportional to the viscosity of the continuous phase. In other words, the relative





**Figure 2-3:** Effects of aggregation, deformation, and orientation of the red cell on viscosity.

$H=45$ ,  $\eta_f=1.2$  cP. NP=normal cells in plasma. NA= normal cells in 11% albumin. HA=hardened cells in 11% albumin. From Chien, 1970.

viscosity  $\eta_r$  is independent of the viscosity of pure suspending fluid  $\eta_f$ , where  $\eta_r \equiv \eta/\eta_f$ . For blood suspensions, however, this is not the case: Experimental work has shown that  $\eta_r$  decreases with increasing  $\eta_f$  [Dintenfass, 1965a]. Blood in this respect is analagous to an emulsion, for which Taylor [1932] and Oldroyd [1953] have shown that, for dilute emulsions, the viscosity is a function of the ratio of the droplet viscosity to that of the non-miscible suspending fluid. Taylor's expression for the bulk viscosity of a dilute emulsion is

$$\eta_r = 1 + 2.5cT \quad (2.5)$$

where  $c$  is the volume fraction of the disperse phase and  $T$  is the Taylor coefficient equal to  $\frac{\eta_i + 0.4\eta_f}{\eta_i + \eta_f}$ , where  $\eta_i$ , the internal viscosity, is the viscosity of the pure disperse phase of the emulsion. This coefficient results from the internal velocity profile established in the fluid drop from the shear and normal forces of the continuous phase. The higher the viscosity of the suspending fluid relative to that of the droplet, the greater is the magnitude of the effect. For  $\eta_i \approx 50\eta_f$ , the droplets behave as rigid particles. For blood cells, the internal viscosity is not a true liquid viscosity but also includes effects from the red cell membrane on the cytoplasm, and therefore  $\eta_i$  has a shear rate dependence.

Secomb and co-workers [1983a] have developed a model for the tank-treading motion of the RBC membrane. This model is based on the following observations:

1. Cell membranes labeled with Heinz bodies or latex spheres show tank-treading motion in high-viscosity dextran solutions.
2. All areas of the membrane move with uniform frequency.
3. Tank-treading motion is seen in whole blood, but not in dilute RBC/plasma suspensions.

From observations (1) and (3), they conclude that the tank-treading motion results from physical interactions with other cells or with very viscous media. This

also supports the hypothesis of a shear-thinning internal viscosity.

The effects of anticoagulation have been characterized by Rosenblum [1968], who classifies anticoagulants into those which shrink the red cell, such as citrate and oxalate, and those that do not change cell size, such as acid-citrate dextrose (ACD). The former increase blood viscosity while the latter show no effect. Cokelet [1963] was also unable to find any effect of anticoagulation with ACD other than those resulting from dilution of the plasma proteins.

The effects of plasma osmolarity were studied by Meiselman *et al.* [1967] who found that cells in hypotonic plasma solutions had an increased hematocrit, a decreased viscosity, and an increased yield stress. Previously, Rand and Burton [1964] had found that cells washed with saline had much less rigid membranes in hypertonic saline solutions. This apparent contradiction in results has been related to the loss of membrane proteins and lipoproteins from cell washing [Meiselman *et al.*, 1967].

## **2.3 Previous Experimental Work**

### **2.3.1 Experimental Work on Blood Viscosity at Normal Hematocrits**

The study of blood viscosity literally dates to the beginnings of viscometry itself. The French physician Jean Poiseuille, credited with the pioneering work on fluid flow through glass capillaries [1843], initially had addressed himself to the relationship of blood flow to pressure, temperature, and vessel diameter. It was only after encountering inconsistencies with blood, which he attributed to variations in blood composition, that he turned to water and other fluids.

Hagenbach [1860] first established the theoretical derivation of Poiseuille's law for laminar flow through a capillary, and first included the concept of viscosity:

$$\mu = \frac{\pi \Delta p}{8QL} R^4 \quad (2.6)$$

where  $\Delta p$  is the pressure difference along the capillary of length  $L$  and radius  $R$ , and  $Q$  is the volumetric flow rate of the fluid.

Attempts to correlate blood flow with this equation persisted through the remainder of the 19<sup>th</sup> century [Ewald, 1877; Lewy, 1897] but systematic work did not begin until the early 20<sup>th</sup> century with the work of duPre'Denning and Watson [1906], who studied blood anticoagulated with potassium oxalate in capillary viscometers and the effects of hematocrit, temperature, anticoagulant concentration, and the size of the capillary bore. Interestingly, they observed an increase in blood viscosity at diameters of 300  $\mu\text{m}$  or less.

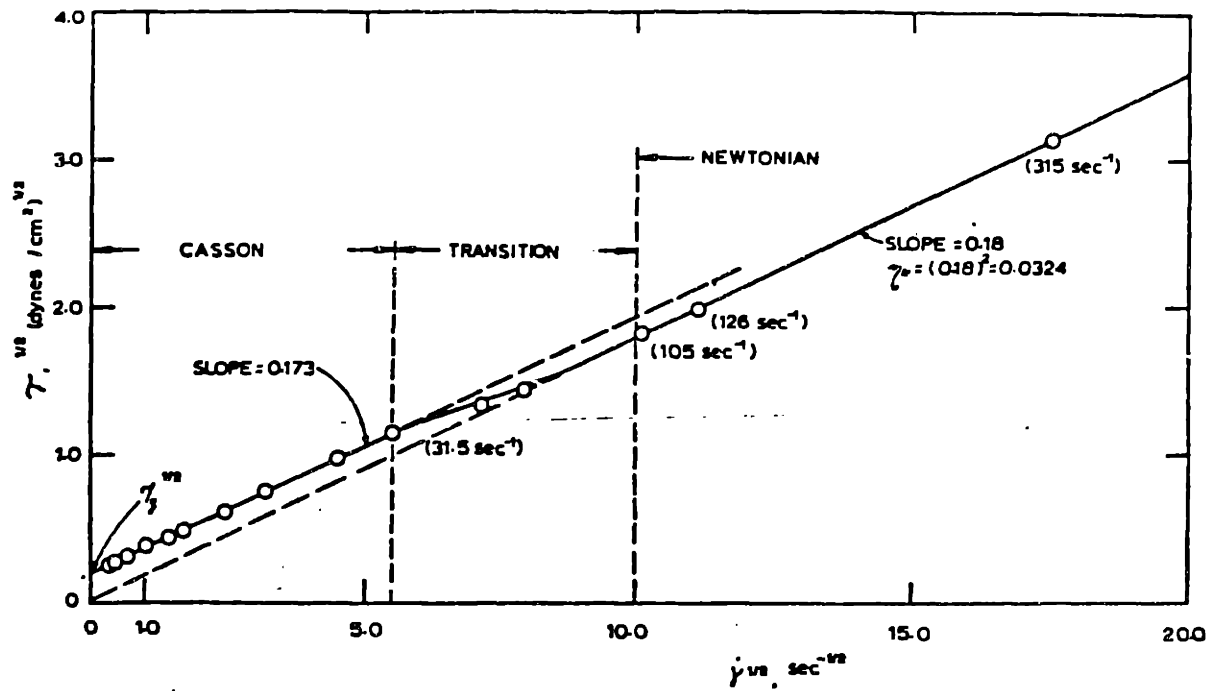
Other investigations in the early 1900s followed a similar tack and include studies of blood of different animals [Burton-Opitz, 1906 and 1911; Welsh, 1911; Hess, 1911 and 1920; Trevan, 1918; Langstroth, 1919; Berczeller and Wastl, 1924]. Nygaard and co-workers [1935] studied the blood viscosities from  $H=14$  to 58 and proposed a combination of a linear and a hyperbolic equation to model the hematocrit dependence.

Of significant importance is the work of Fahraeus and Lindqvist [1931], who showed that in capillaries of diameters less than 300  $\mu\text{m}$ , Poiseuille flow is not observed and the apparent viscosity of blood decreases as the tube diameter decreases, in contrast to the results of duPre'Denning and Watson. This phenomenon, known as the Fahraeus-Lindqvist effect, results from the tendency of red cells to migrate away from the vessel walls, establishing a cell-free layer of fluid around the flowing cells. This layer has a lubricating effect which grows more significant as the tube diameter decreases, resulting in detectable alterations in the bulk viscosity.

In the 1940s and 1950s activity in the field of blood rheology further increased [Copley *et al.*, 1942; Moll, 1943; Riehl, 1943; Bingham and Roepke, 1944a and 1944b], but it was not until the early 1960s that investigators were able to study blood at low rates of shear (less than  $100 \text{ sec}^{-1}$ ), where its non-Newtonian behavior is evident. This became possible through the employment of cone-and-plate and concentric-cylinder (Couette) viscometers with advanced instrumentation, which were more suitable for this purpose than capillary viscometers. Dintenfass [1962a] claims the first report on the shear-thinning behavior of blood; that is, that the viscosity of blood reversibly decreases with increasing shear rate. (Dintenfass uses the term *thixotropic* throughout his work to describe this behavior, but in the classical sense thixotropy defines behavior under non-steady-state conditions: A thixotropic material is one which experiences a transient decrease in viscosity upon the application of a constant shear stress [Goodeve and Whitfield, 1938; Bird *et al.*, 1960]. Dintenfass [1968a] argues that the difference between shear-thinning and thixotropy is only a question of time scale and instrumentation limits; nevertheless, the classical definitions will be retained in this work.)

Cokelet, Merrill, and co-workers conducted extensive studies into the low-shear rate behavior of blood [Cokelet, 1963; Cokelet *et al.*, 1963, Merrill *et al.*, 1963a and 1965a]. Typical data for normal blood is shown in Figure 2-4.

Figure 2-4 is a plot of the square root of shear stress ( $\tau$ ) versus the square root of shear strain rate ( $\dot{\gamma}$ ). Three regimes are noted: A region at high shear rates (greater than  $100 \text{ sec}^{-1}$ ) governed by Newtonian behavior, a non-Newtonian region below  $20 \text{ sec}^{-1}$  and a transition region in-between the two. The low shear rate range (down to  $0.1 \text{ sec}^{-1}$ ) is correlated to a semiempirical equation first derived by Casson [1958] to describe printing ink rheology and later proposed by Reiner and Scott-Blair [1959] to be applicable to blood:



**Figure 2-4:** Casson plot for normal human blood.

H=40, T=37 °C.  $\tau_y \approx 0.04$  dynes/cm<sup>2</sup>. From Merrill, 1969.

$$\sqrt{\tau} = \sqrt{\tau_y} + \sqrt{\mu_N \dot{\gamma}} \quad (2.7)$$

where  $\tau_y$  is a yield shear stress and  $\mu_N$  is the asymptotic Newtonian viscosity at high shear rates.

Values for  $\tau_y$  have been obtained by a variety of experimental methods, all of which have the potential for introducing significant instrument artifact. These methods include extrapolation of Casson plots to zero shear rate [Cokelet *et al.*, 1963; Merrill *et al.*, 1966], torque decay [Merrill *et al.*, 1965a], pressure decay [Merrill *et al.*, 1965a], force on an oscillating plate [Benis and Lacoste, 1968], and sedimentation [Charm and Kurland, 1967]. The values for  $\tau_y$  measured by these methods range from 0.003 to 0.20 dyne/cm<sup>2</sup>.

Red cells in Ringer or isotonic saline solution or suspended in plasma suspensions lacking fibrinogen do not exhibit a yield shear stress [Cokelet, 1963; Merrill *et al.*, 1963b and 1965c]. Addition of a slight amount of fibrinogen gives conformity to the Casson equation, and the magnitude of the yield stress varies with fibrinogen concentration [Merrill *et al.* 1963b, 1965c, 1966].

The existence of a yield shear stress for blood, and thus the validity of the Casson equation, is disputed by some researchers [Chien *et al.*, 1970; Chmiel, 1974] who believe it is an artifact extrapolated from the limits of the instrumentation. Experiments conducted at shear rates from 0.1 to 5 x 10<sup>-3</sup> sec<sup>-1</sup> using a wide gap Couette system showed significant deviation from the Casson equation which indicated the absence of a yield shear stress [Chmiel, 1974]. The questions about yield shear stress concern a no-flow situation of principally academic interest. They are of no concern at the shear rates of interest for plasmapheresis devices and require measurements out of the range of the instrumentation used here. The yield shear stress will not be considered at any length in this work.

### 2.3.2 Experimental Work on Blood Viscosity at High Hematocrits

As early as 1924, there were reports that the viscosity of packed cells was 8 to 90 times that of plasma [Berczeller and Wastl, 1924]. More recently, the viscosity of packed red cells became of interest when it was realized that blood does not behave as a suspension of rigid particles [Dintenfass, 1962b]. A 65% suspension of clay or microscopic glass spheres would have a viscosity 10 to 100 times that of blood at the same cell concentration.

There is a dearth of data on blood viscosity at high hematocrits. What little data that does exist shows a considerable variance in reported values, as plotted in Figure 2-5. Chien *et al.* [1966] have the most complete base of data to date on the variation of viscosity over a wide range of hematocrits and shear rates. Three types of blood preparations--whole blood, defibrinated blood, and Ringer-suspended cells--were measured over hematocrits ranging from 0 to 95. Their data is shown in Figures 2-6 to 2-8 fitted to fifth-order polynomial curves of the form,

$$\ln (\eta) = \sum_{i=0}^5 a_i H^i \quad (2.8)$$

whose constants  $a_i$  are given in Table 2-I. At low hematocrits ( $H < 5$ ), Newtonian behavior is observed for all suspensions. Newtonian behavior persists up to  $H=15$  when fibrinogen is absent and up to  $H=30$  when all plasma proteins are absent. The viscosity differences between suspensions is the greatest at low shear rates, supporting the contention that protein-mediated aggregation dominates rheology at low shear rates. Viscosity decreases when plasma proteins are removed, but at high hematocrits, the differences between the suspensions become negligible as the volume of non-cellular medium vanishes.

The work of Chien *et al.* is probably the most referred-to source in the literature and their equations the most used for examining blood viscosity as it



Viscosity vs. Shear rate  
Literature data at high hematocrits

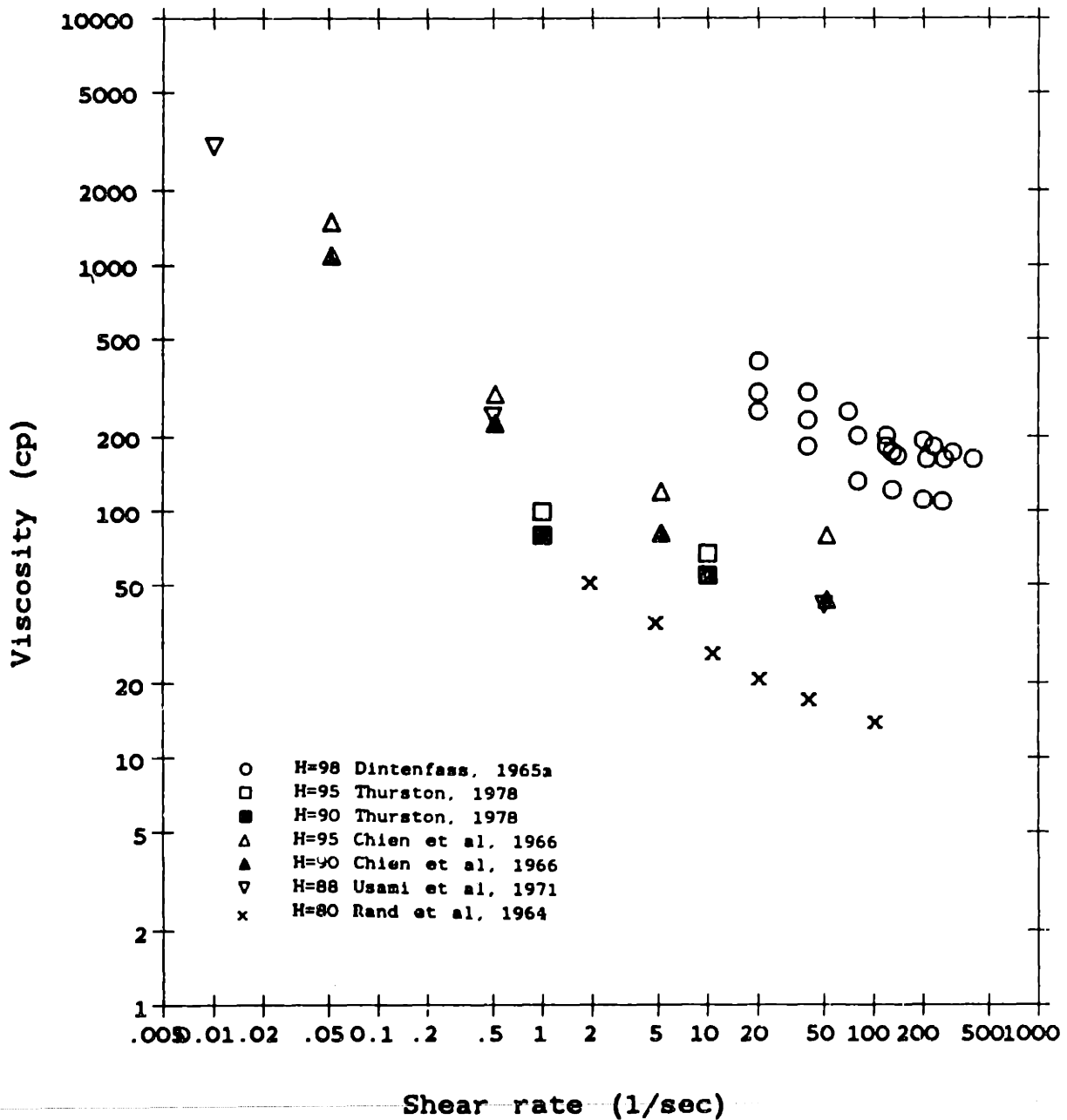


Figure 2-5: Literature data for blood viscosity at high hematocrits.

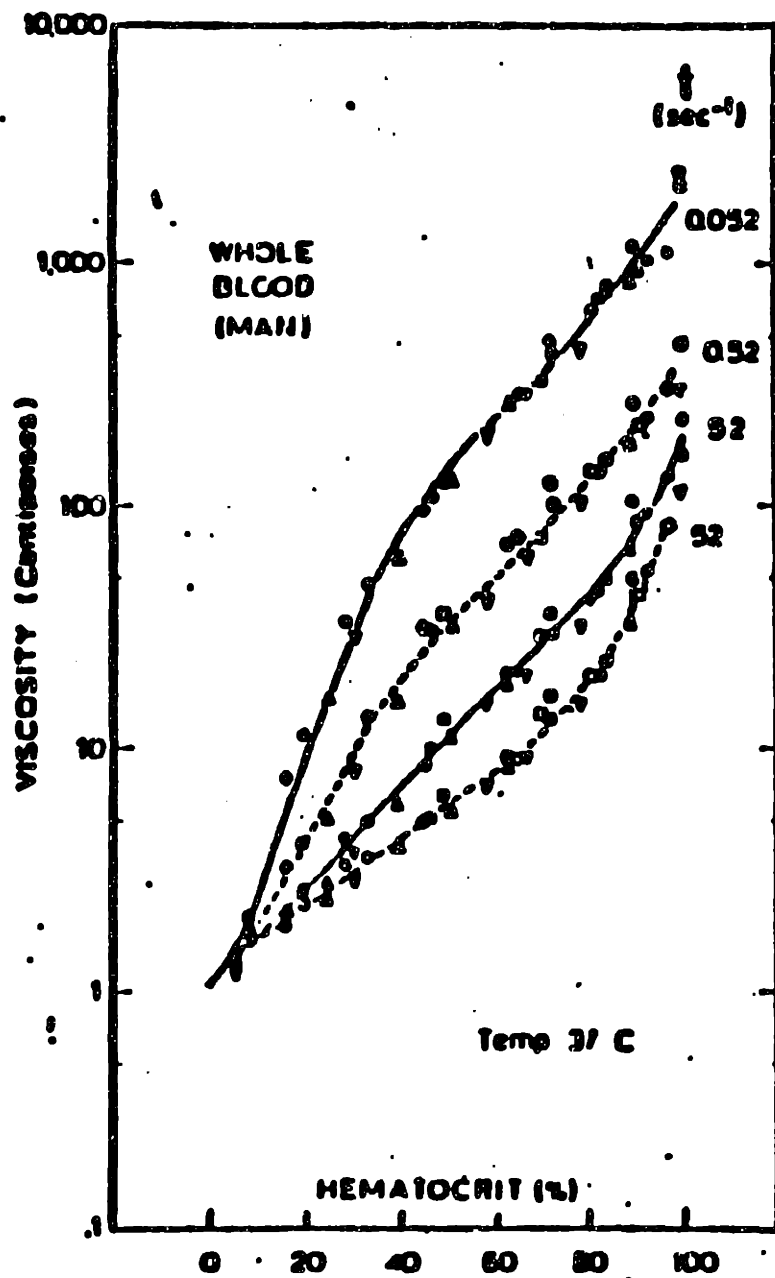
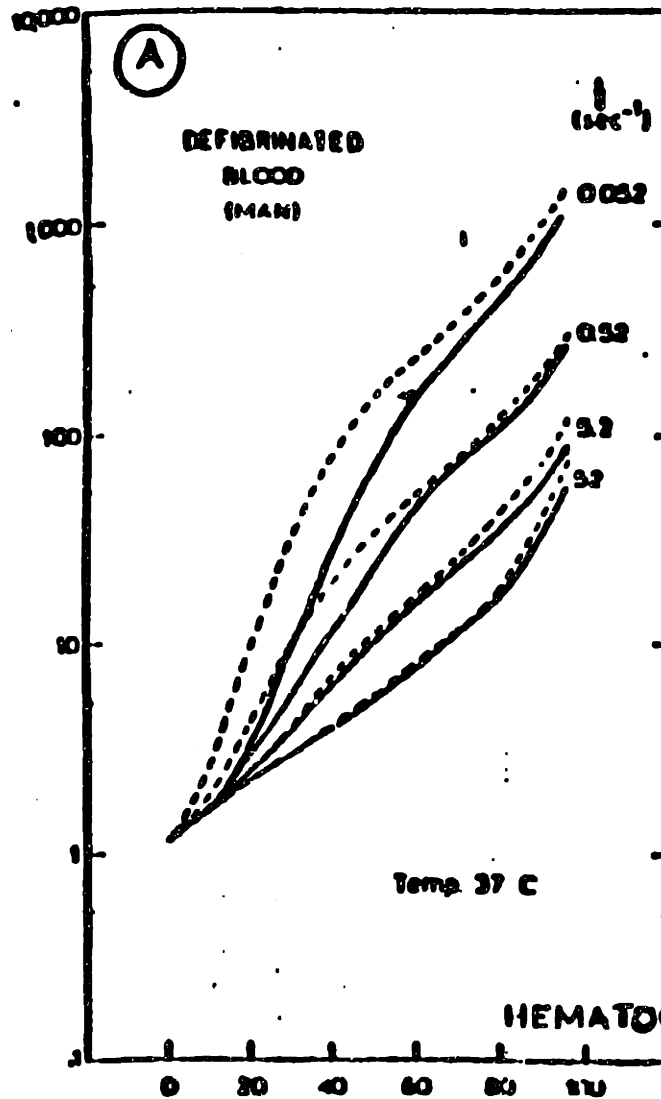


Figure 2-6: Data of Chien *et al.* [1966] for whole blood.



**Figure 2-7:** Data of Chien *et al.* [1966] for defibrinated blood.

Dashed lines represent whole blood curves.

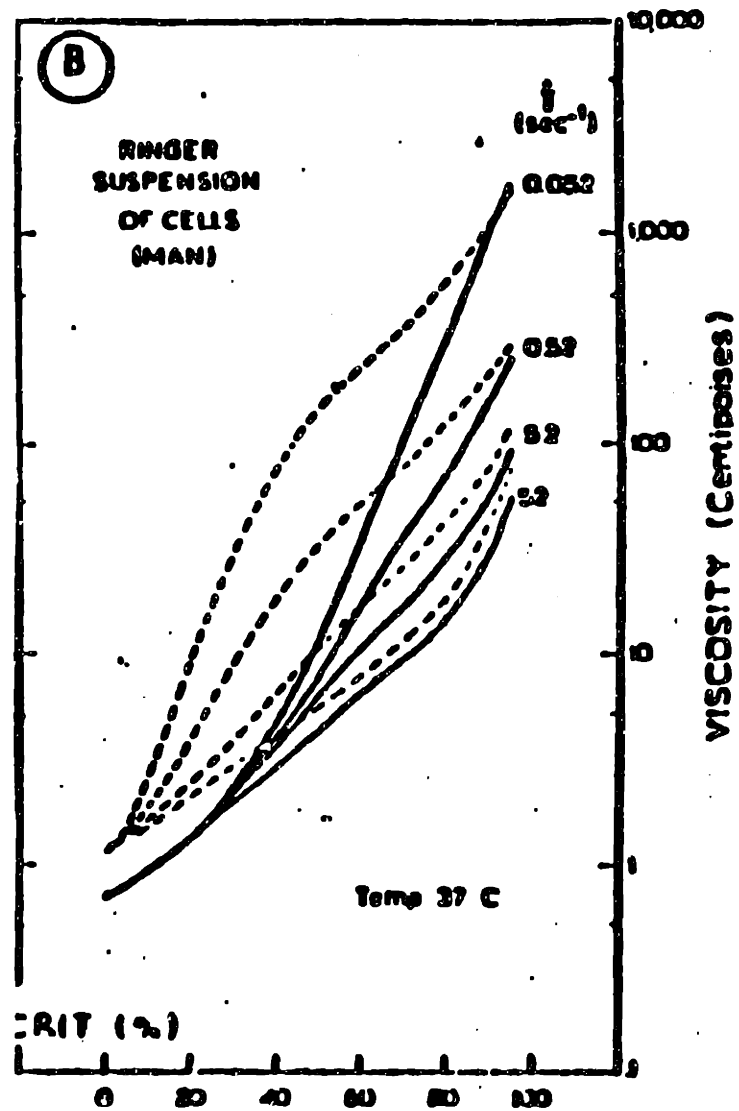


Figure 2-8: Data of Chien *et al.* [1966] for Ringer's suspension.

Dashed lines represent whole blood curves.

Constants for polynomial fit of data of Chien et al, 1966.						
Solution type/ Shear rate (1/s)	a0	a1	a2	a3	a4	a5
vb 52	0.2028	2.928e-02	-1.570e-04	1.3850e-05	-2.8130e-07	1.788e-09
vb 5.2	0.1817	2.423e-02	9.850e-04	-1.3720e-05	2.9900e-08	3.530e-10
vb .52	0.1499	2.599e-02	3.598e-03	-9.4670e-05	9.5240e-07	-3.335e-09
vb .052	0.1342	4.230e-02	5.820e-03	-1.6425e-04	1.6958e-06	-6.035e-09
db 52	0.1336	5.467e-02	-1.369e-03	4.4390e-05	-5.2420e-07	2.381e-09
db 5.2	0.0813	4.245e-03	-4.300e-04	2.4560e-05	-4.0510e-07	2.097e-09
db .52	0.1260	4.130e-02	-9.300e-05	3.2280e-05	-6.0980e-07	3.147e-09
db .052	0.0681	-3.440e-03	4.383e-03	-7.3530e-05	4.4460e-07	-7.000e-10
rs 52	-0.3651	4.131e-02	-9.700e-04	4.0600e-05	-6.1720e-07	3.186e-09
rs 5.2	-0.3644	2.160e-02	3.140e-04	1.7630e-05	-4.1340e-07	2.397e-09
rs .52	-0.3658	4.237e-02	-1.508e-03	6.7960e-05	-8.9610e-07	3.908e-09
rs .052	-0.3760	3.763e-02	-6.450e-04	3.8020e-05	-4.0730e-07	1.369e-09

"vb"-whole blood "db"-defibrinated whole blood  
"rs"-Ringer's suspension

**Table 2-I:** Constants for polynomial fit of Chien data.

From Chien *et al.*, 1966.

varies with cell concentration and shear rate. An inconsistency with the Chien polynomials is that the  $a_0$  values for any one type of suspension are not equal, and thus non-Newtonian behavior is predicted for pure plasma and saline. The purely empirical constants from these polynomials give little insight into the physical phenomena being observed. The fact that their ranges extend only to hematocrits of 95 and shear rates of  $52 \text{ sec}^{-1}$  points out the need for more extensive data at the high-shear, high-hematocrit spectrum as well as a renewed effort to develop a set of equations satisfactory for predicting blood viscosity over a broad range of conditions.

An obvious approach to expanding the applicability of the Chien equations is to derive another set of empirical equations relating the coefficients  $a_i$  in equation (2.8) to shear rate. Constraints may be imposed such as requiring that as  $H \rightarrow 0$ ,  $\eta \rightarrow \eta_f$  for all shear rates (i. e., setting  $a_0$  to a constant for each set of curves) and that Newtonian behavior is approached at high shear rates. Such an undertaking was made as a part of this work and the results are reported on in Chapter 4.

## **2.4 Theoretical Models and Mathematical Correlations**

Because of the complex nature of the cell-cell and cell-fluid interactions in blood and red blood cell suspensions, to date it has not been possible to explain their rheology from the standpoint of a unified theory. Such a theory would have to incorporate the various aspects of shear rate and concentration dependence described in the previous sections of this chapter. To recap:

1. As shear rates approach zero, non-Newtonian behavior results from the domination of cell-protein interactions and subsequent increased aggregation.
2. At high shear rates, cells deform and orient along the streamlines, there

is tank-treading motion of the membrane and establishment of an internal velocity profile. Newtonian behavior is observed.

3. The fluid nature of the erythrocyte makes for substantial deviations in behavior of RBC suspensions versus that of rigid particle suspensions. For blood, the relative magnitude of the viscosity is smaller at a given concentration. Even at greater than 95 volume percent, blood suspensions are quite fluid. Analogies can be drawn to emulsions, with the cellular phase represented by a non-Newtonian internal viscosity.
4. Suspensions are Newtonian in the limit of dilute concentrations.
5. At high concentrations, cell crowding effects predominate and viscosity becomes independent of the suspending fluid.

Most proposed models for blood are extensions of models for suspensions of rigid particles or for emulsions. It is the intent here to briefly mention several approaches.

The Casson equation, equation (2.7), was originally developed to explain the behavior of mutually attractive particles which form structures that diminish in size as shear rate increases. In addition, the structure which exists at zero shear must be subjugated to a yield stress before any flow develops. The Casson equation does not explain blood rheology at low hematocrits, where no yield stress is observed.

Many empirical equations have been modifications of an equation first derived by Einstein [1906] for dilute suspensions of rigid, non-interacting particles,

$$\eta_r = \frac{1}{1-\alpha \cdot c} \quad (2.9)$$

where  $\alpha$  is a shape factor equal to 2.5 for spheres. At hematocrits less than 5, where Newtonian behavior is observed, equation (2.9) describes whole blood adequately. Jeffrey [1923] extended Einstein's theory to particles of ellipsoid shape, from which Taylor [1932] derived his expression for emulsion viscosity, equation (2.5). Dintenfuss [1968b], starting with Taylor's equation, proceeded by a method where a

concentrated emulsion is equated to a hypothetical continuous phase of equal viscosity, and then sequential addition of disperse phase to the emulsion is assumed to follow Taylor's equation. The result is

$$\eta_r = \frac{1}{(1-c \cdot T)^{2.5}} \quad (2.10)$$

Values for  $\eta_i$  from this equation are between 1 and 6 cP [Dintenfass, 1968b]. This is smaller than the reported 6-8 cP viscosity of a 34 per cent hemoglobin-in-water solution, which approximates the makeup of the erythrocyte cytoplasm [Quemada, 1977].

The methodology used to derive this equation is valid only at shear rates high enough to prevent aggregation, so that the cells exist as independent units. Based on attempts at using it with experimental data, it will hold for more dilute suspensions ( $H < 85$ ) only with very viscous suspending fluids. For example, in order to measure  $\eta_i$  at  $H=38$ , Dintenfass [1968b] uses a suspending fluid viscosity  $\eta_f > 9$  cP. The limitations of equation (2.10) would follow from the concept of a shear-thinning internal viscosity and from the previously mentioned observations of Secomb *et al.* [1983a]. Only at high shear rates would the emulsion properties of blood become apparent from deformation, orientation, and tank-treading. The tank-treading is dependent on viscous stresses from either other cells or from highly viscous media [Secomb *et al.*, 1983a].

Quemada [1977, 1978a, and 1978b] has proposed an interesting equation for concentrated suspensions based on the principle of minimum energy dissipation. His derivation is strictly valid for rigid, attracting spheres but is seen to adequately describe blood rheology. The expression for the relative viscosity as a function of both shear rate and concentration is

$$\eta_r = \left[ 1 - \frac{c}{2} \left( \frac{k_0 + k_\infty \sqrt{\gamma_r}}{1 + \sqrt{\gamma_r}} \right) \right]^{-2} \quad (2.11)$$

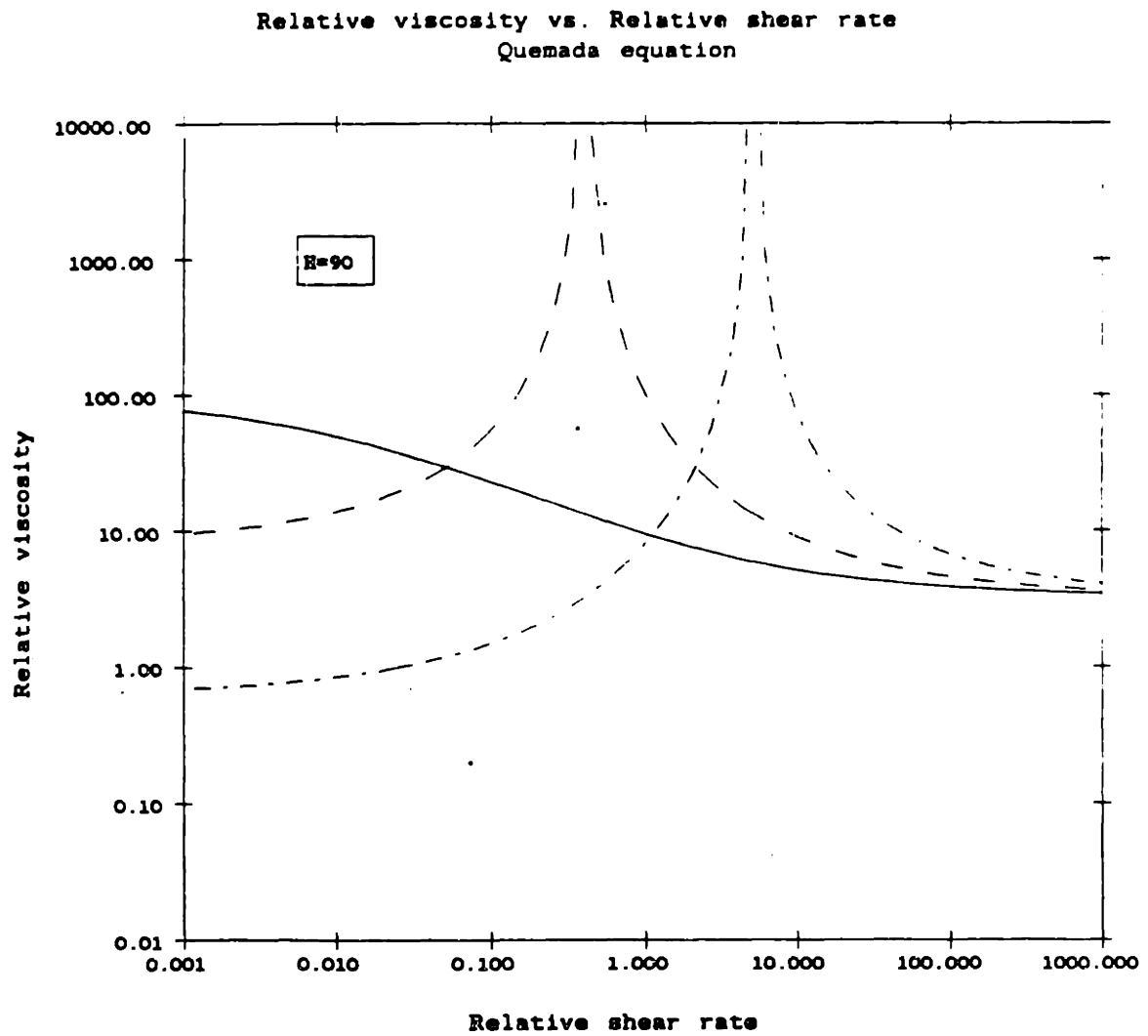


or

$$\eta_r = \left[ 1 - \frac{c}{2} \left( k_\infty + \frac{k_0 - k_\infty}{1 + \sqrt{\gamma_r}} \right) \right]^{-2}$$

where  $k_0$  and  $k_\infty$  are dimensionless "intrinsic viscosities" defined as  $\left( \frac{d\eta_r}{dc} \right)_{c \rightarrow 0}$  at zero and high shear rate, respectively. The term in parentheses in equation (2.11) is the intrinsic viscosity  $k(c, \gamma)$  at the shear rate of interest. For rigid particles, the intrinsic viscosity is inversely proportional to the packing concentration  $c_p$ , which is the maximum concentration to which the effective particles existing at the shear rate may be packed. In other words, the intrinsic viscosity is directly proportional to the volume of an effective particle. For deformable fluid particles such as erythrocytes, the intrinsic viscosity empirically includes any deviations from rigid spheres, and the relationship of intrinsic viscosity to particle volume is more qualitative. The dimensionless relative shear rate  $\gamma_r$  is equal to  $\gamma/\gamma_c$  where  $\gamma_c$  is a theoretical "critical shear rate" for aggregate formation. Aggregation is assumed to be governed by Brownian collisions while disaggregation is controlled by shear-induced collisions. In the limit of a dilute concentration of rigid spheres,  $\gamma_c$  is roughly equivalent to the rotational diffusion coefficient  $\frac{D}{a^2}$ , where  $D$  is the translational diffusivity and  $a$  is the radius of the particle. For blood cell aggregates,  $\gamma_c$  may be on the order of a Maxwell relaxation time  $\frac{E_i}{\eta_i}$ , where  $E_i$  is the elastic modulus of the aggregate. A complete derivation of the Quemada equation is given in Appendix C.

Since the Quemada expression will be employed extensively in the analysis of this work, it will be worthwhile at this point to examine this equation in some detail, first from a purely mathematical standpoint and then later in the context of the physical meaning of its parameters. Sets of curves for the Quemada equation are

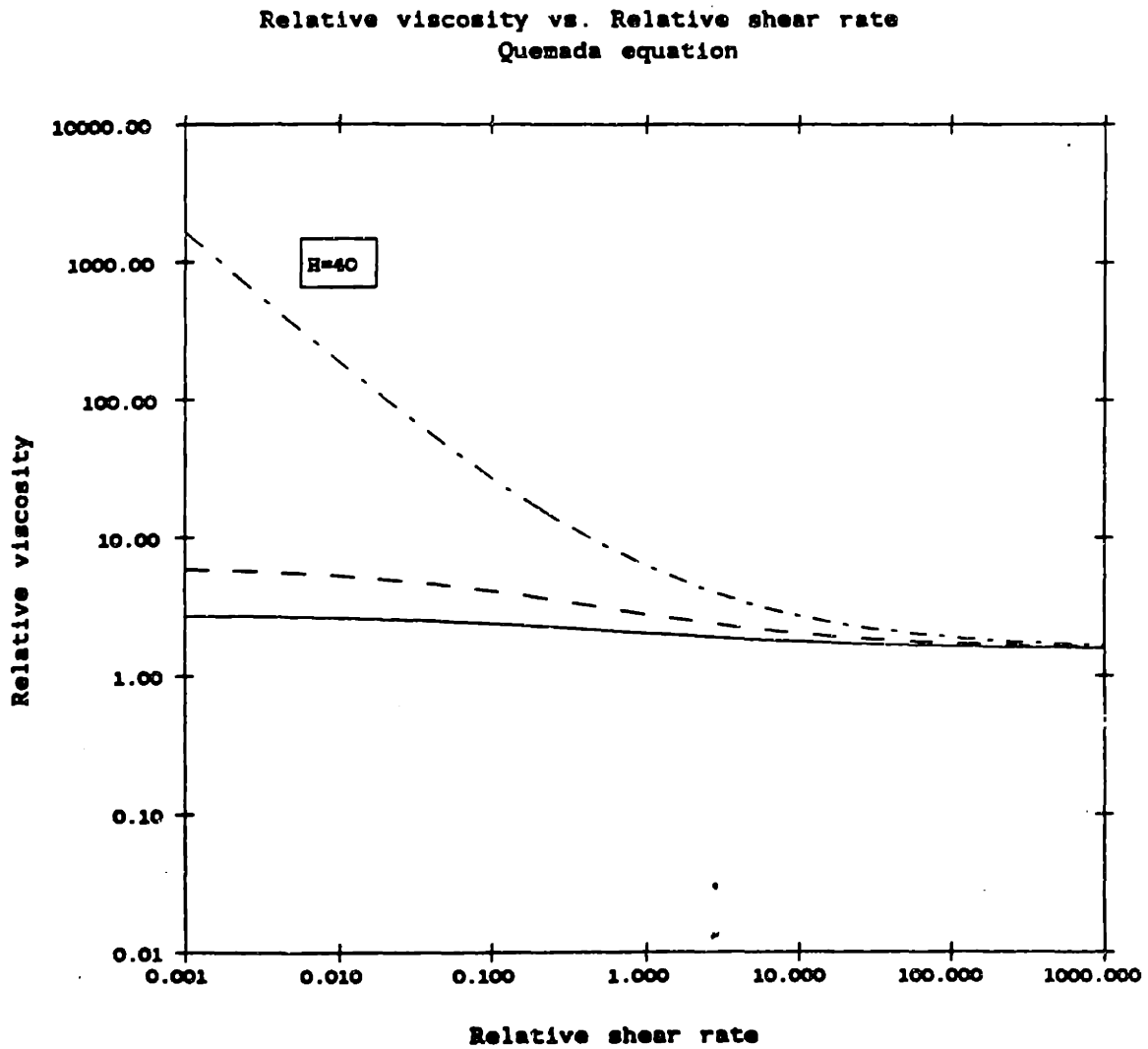


**Figure 2-9:** Plots of the Quemada equation at  $H=90$ .

Solid lines:  $k_{\infty}=1, k_0=2$

Dashed lines:  $k_{\infty}=1, k_0=3$

Dot/dashed lines:  $k_{\infty}=1, k_0=5$

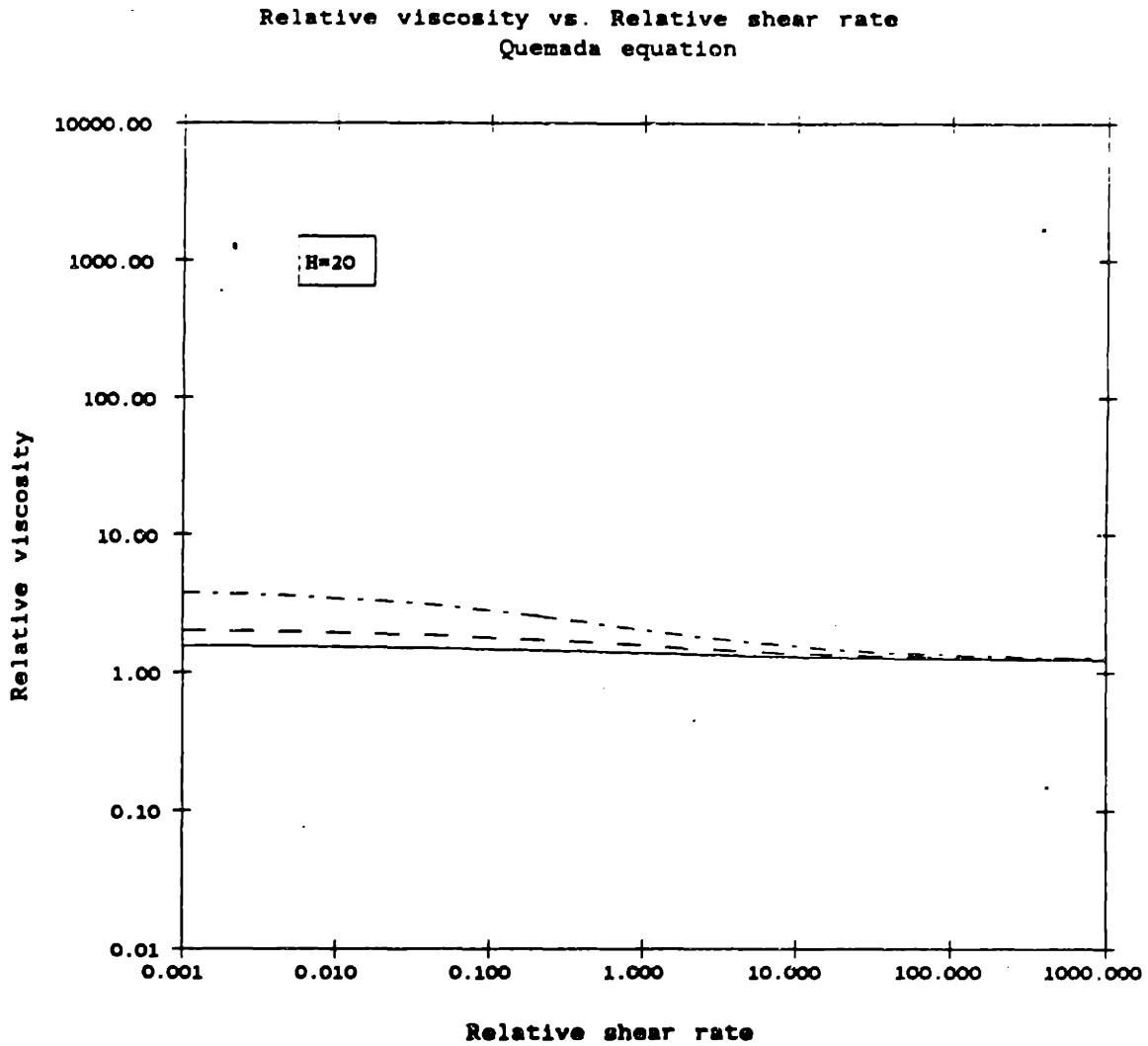


**Figure 2-10:** Plots of the Quemada equation at  $H=40$ .

Solid lines:  $k_{\infty}=1, k_0=2$

Dashed lines:  $k_{\infty}=1, k_0=3$

Dot/dashed lines:  $k_{\infty}=1, k_0=5$



**Figure 2-11:** Plots of the Quemada equation at  $H=20$ .

Solid lines:  $k_{\infty}=1, k_0=2$

Dashed lines:  $k_{\infty}=1, k_0=3$

Dot/dashed lines:  $k_{\infty}=1, k_0=5$

shown in Figures 2-9 to 2-11 for hematocrits of 90, 40, and 20. From these curves it is evident that  $k_\infty$  and  $k_0$  provide the asymptotic values of the viscosity, while the critical shear rate scales the x-axis over the range of shear rates examined. A low value for  $\gamma_c$  shifts the relative shear rate to the right, toward the  $k_\infty$  asymptote. Alternatively, the critical shear rate can be thought of as an analogue to the Michaelis-Menten rate constant: when  $\gamma = \gamma_c$ ,  $k = \frac{k_0 + k_\infty}{2}$ .

As shown in the graph for  $H=90$ , there exists the potential for a singular solution at high hematocrit and high  $k_0$ , where the viscosity is infinite. This occurs when the term in brackets in equation (2.11) is zero, and

$$\sqrt{\gamma_r} = \frac{\frac{c}{2}k_0 - 1}{1 - \frac{c}{2}k_\infty}. \quad (2.12)$$

This has a real-valued solution only when  $\sqrt{\gamma_r} > 0$ , and thus the criterion for the existence of a singular point is

$$\frac{1 - \frac{c}{2}k_0}{1 - \frac{c}{2}k_\infty} < 0. \quad (2.13)$$

Since  $k_0 > k_\infty$  for shear-thinning suspensions, this reduces to

$$\frac{c}{2}k_0 > 1. \quad (2.14)$$

This case has no physical meaning, as viscosity would never follow this behavior, but may result when  $k_0$  values are extrapolated from data at higher shear rates.

Quemada fit his equation to experimental data from the literature [Quemada, 1978b]. The values derived for the rheological parameters from these fits are shown in Table 2-II and the corresponding equations are plotted in Figure 2-12. He has also [Quemada, 1981] fit his equation to the data of Chien *et al.*, 1966, but unfortunately presented his results graphically, in terms of  $c_p$ , without reporting the

Source	H	$k_{\infty}$	$k_0$	$\gamma_c$ (sec <sup>-1</sup> )
1	41.9	1.8	4.68	0.69
2a	40	1.84	4.65	2.23
2b	45	2.07	4.33	1.8
3a	45	1.78	4.20	5.0
3b	45	1.78	3.29	25.0
3c	45	3.62	-	-
4	88	1.83	2.26	2.16

**Table 2-II:** Fits of literature data to Quemada equation.

Source codes: 1-Merrill *et al.*, 1965c, whole blood

2-Schmid-Schonbein *et al.*, 1971b, whole blood

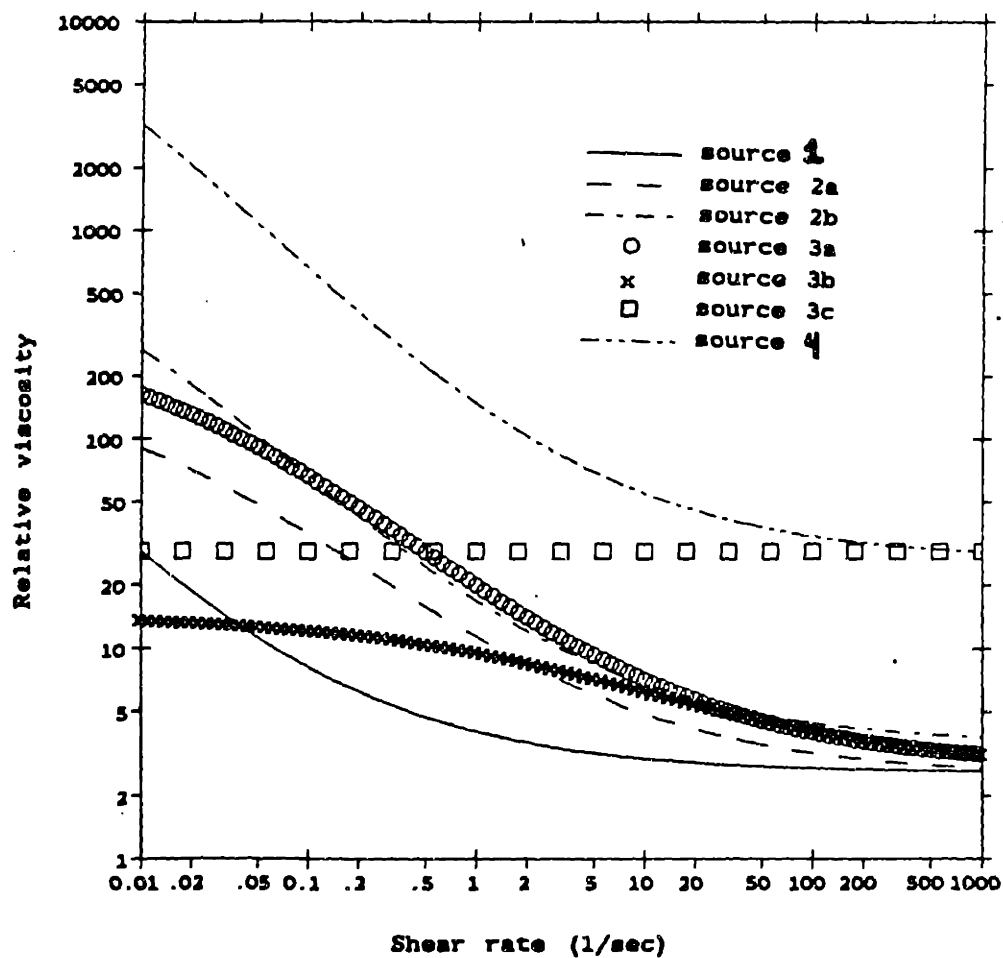
3-Chien, 1970: (a) normal cells in plasma, (b) normal cells in albumin,  
(c) hardened cells in albumin

4-Usami *et al.*, 1971, packed cells in plasma

$\eta_f = 1.2$  cP for all samples.

From Quemada, 1978b.

Relative viscosity vs. Shear rate  
Plots of literature data as fitted by Quemada



**Figure 2-12:** Plots of Quemada equation for literature sources given in Table 2-II.

values obtained. As part of the present work, fits were made to the Chien data and are discussed in Chapter 5.

Values for  $k_{\infty}$  in Table 2-II are very close to 1.8 save for the hardened cells. Ideally,  $k_{\infty}$  can be thought of as being affected by anything which would affect either the internal viscosity of the cell or the effective hydrodynamic volume of the cell at high shear, since  $k_{\infty}$  is related to the packing concentration of effective particles (see Appendix C). It should therefore hold that as long as the nature of the cells, the hematocrit, and the suspending fluid viscosity are the same, then  $k_{\infty}$  should be constant. The value obtained at  $H=88$  is similar to the others in Table 2-II, but in analyzing data from Chien *et al.* [1966], Quemada [1981] noted that at low hematocrits,  $k_{\infty}$  decreased as hematocrit increased, reached a minimum at  $H=75$ , and then increased. His proposed explanation for this is to imagine that below hematocrits of around 40, the effective particle volume at high shear is the hydrodynamic volume of an individual cell, which is reduced by crowding effects as hematocrit increases. At about  $H=40$ , these effective volumes begin to overlap, and at  $H=75$  an inversion phenomenon typical to emulsions happens as the continuous solution phase becomes sequestered between cells and converts into a dispersed phase, while the erythrocytes become the continuous phase.

Comparison of the values of  $k_0$  give an indication of the effects of aggregation. The absence of plasma proteins leads to less aggregation at zero shear rate and higher packing concentrations as described in Appendix C. In fitting the data of Chien *et al.*, Quemada [1981] noted an almost linear decrease of  $k_0$  with hematocrit, which he ascribed to successive effects of increased aggregation and then deformation, both which result in increasing the "compactness" of the red cell aggregates.

The values for  $\gamma_c$  in Table 2-II show a large jump from 5 to 25  $\text{sec}^{-1}$  in going



from plasma to albumin suspensions. If, as in polymer chains,  $\gamma_c$  is assumed to be proportional to the elastic modulus  $E_i$  as shown in Appendix C, and  $E_i \sim a^{-3}$  where  $a$  is the effective radius of the rouleau, then effects which decrease aggregation would increase  $\gamma_c$ .

In the limit as  $c \rightarrow 0$ , it is reasonable to assume that  $\gamma_c \rightarrow \infty$  since there is less aggregation and the time constant for rotational diffusion will be small. The term in parentheses in equation (2.11) becomes  $1 - \frac{c}{2}k_0$ . In the fits of Quemada for the data of Chien *et al.*, however, it appears that both  $k_0$  and  $k_\infty$  are discontinuous at the boundary of Newtonian behavior and that the Newtonian intrinsic viscosity is some intermediate value between the two.

A few other methods will be quickly touched upon. Linear viscoelastic models for blood attempt to model both the dissipative (viscous) and energy-storing (elastic) properties of blood. Thurston [1972] was the first to delve into this arena of biorheology. There have been many proposed constitutive equations for blood, but none have been found to be fully satisfactory in terms of representing all of the material function of blood (shear viscosity, normal stress differences, and oscillatory material functions). If interest is limited to just the shear viscosity, however, several models predict blood behavior reasonably well, although all previous work has been at hematocrits in the physiologic range [Walitza *et al.*, 1979; Riha, 1980]. These models are limited by their complexity and their inability to be related to physical parameters.

Princen [1983] has proposed a method of calculating yield shear stress for foams and concentrated emulsions based on a geometric analysis of the shear stress and membrane tension of hexagonally closed-packed drops. His analysis does not extend to void fractions above 10 per cent.

Secomb *et al.* [1983b] have developed very complicated expressions for the relative viscosity of packed red cells based on a model for tank-treading motion of the cells. Their model shows good agreement at high shear rates but is very poor at both lower hematocrits and low shear rates, where the tank-treading effect is less important rheologically.

## **Chapter 3**

# **MATERIALS AND METHODS**

### **3.1 Blood Sampling and Preparation**

The blood samples used for the viscosity measurements were either drawn from the author or were units of whole blood or of packed cells obtained from the American Red Cross. Those drawn from the author were used in preliminary measurements and consisted of approximately 7 ml of blood drawn into Vacutainers (Becton-Dickinson no. 6460) with sodium citrate anticoagulant. Each unit was drawn into either citrate phosphate dextrose (CPD), citrate phosphate dextrose-adenine<sub>1</sub> (CPD-A<sub>1</sub>), or adenine-saline<sub>1</sub> (AS<sub>1</sub>) anticoagulant-preservative and was used before its expiration date. Typically these units were obtained approximately 14 days after they had been collected from donors. A unit of whole blood consists of 450 ml of blood drawn into 63 ml of anticoagulant-preservative and a unit of packed red cells consists approximately of 250 ml of blood concentrated to 60-75% cell volume. Units were stored refrigerated at 4 °C.

Saline suspensions of blood were made from the units of packed red cells. Ten 15-ml Pyrex centrifuge tubes were filled with 8 to 10 ml each of packed cells and centrifuged for four hours at 900*g* in a refrigerated centrifuge (International Equipment Company, model PR6). The supernatant and buffy coat were removed and discarded and the remaining cells (about 5-6 ml) were washed in an equal volume of Eagle's solution and centrifuged again for thirty minutes to an hour. The cells were washed and centrifuged twice more. The final centrifugation lasted for 12 hours and, after removal of the supernatant, Eagle's solution was added to adjust to

the desired hematocrit. The recipe for Eagle's solution is given in Appendix B.

Suspensions of red blood cells in plasma were prepared by spinning down 10 ml aliquots for 12 hours. The supernatant was retained and the cells were not washed in Eagle's solution. Instead, the supernatant was centrifuged for 30 minutes to settle any solid materials, which were discarded, and the plasma was used to adjust the hematocrits of the red cell samples.

### **3.2 Hematocrit Determination**

Hematocrit determination was done by the microhematocrit method. Two 25  $\mu$ l Microcap (Drummond Scientific Company) glass capillary tubes were filled one-half to two-third full with the sample in question. One end of each tube was sealed with Seal-Ease (Clay-Adams, Parisippany, NJ), and the tubes were spun in a Readacrit microcentrifuge (Clay-Adams no. 0591, Parisippany, NJ) at 6300g for 5 minutes. The length of the red cell column and the total length of material in the tube was read by graduations built into the centrifuge. For microhematocrits below 98, if agreement was not within 0.5 hematocrit units, a third measurement was done. For microhematocrits above 98, the agreement criteria between tubes was within 0.3 hematocrit units and the values were recorded to the nearest tenth.

It is well known that a column of packed red cells contains a small amount of plasma trapped between the cells, which in the case of microhematocrit determination means that the reading is higher than the actual volume percentage of cells. Methods using radioactive-labeled albumin have been employed to quantify the amount of plasma trapped [Chien *et al*, 1965; Thomas and Janes, 1968], and an attempt was made to adapt one of these methods to determine with greater accuracy the hematocrit of the high concentration samples. The nature of these samples, their high viscosity and adhesiveness to the sides of glass tubes, made it

impossible to measure out accurate volumes as required by the Chien method and the method of Thomas and Janes proved unsuitable for samples of extremely low plasma to cell ratio. Values of the plasma trapping correction factor (equal to 1 minus the porosity) for the microhematocrit method are in the range of .99 [Chien *et al*, 1965] at cell volumes of 44.5%. Since the porosity of a centrifuged column decreases with increasing length of cell column, the correction factor should be even higher for highly concentrated suspensions.

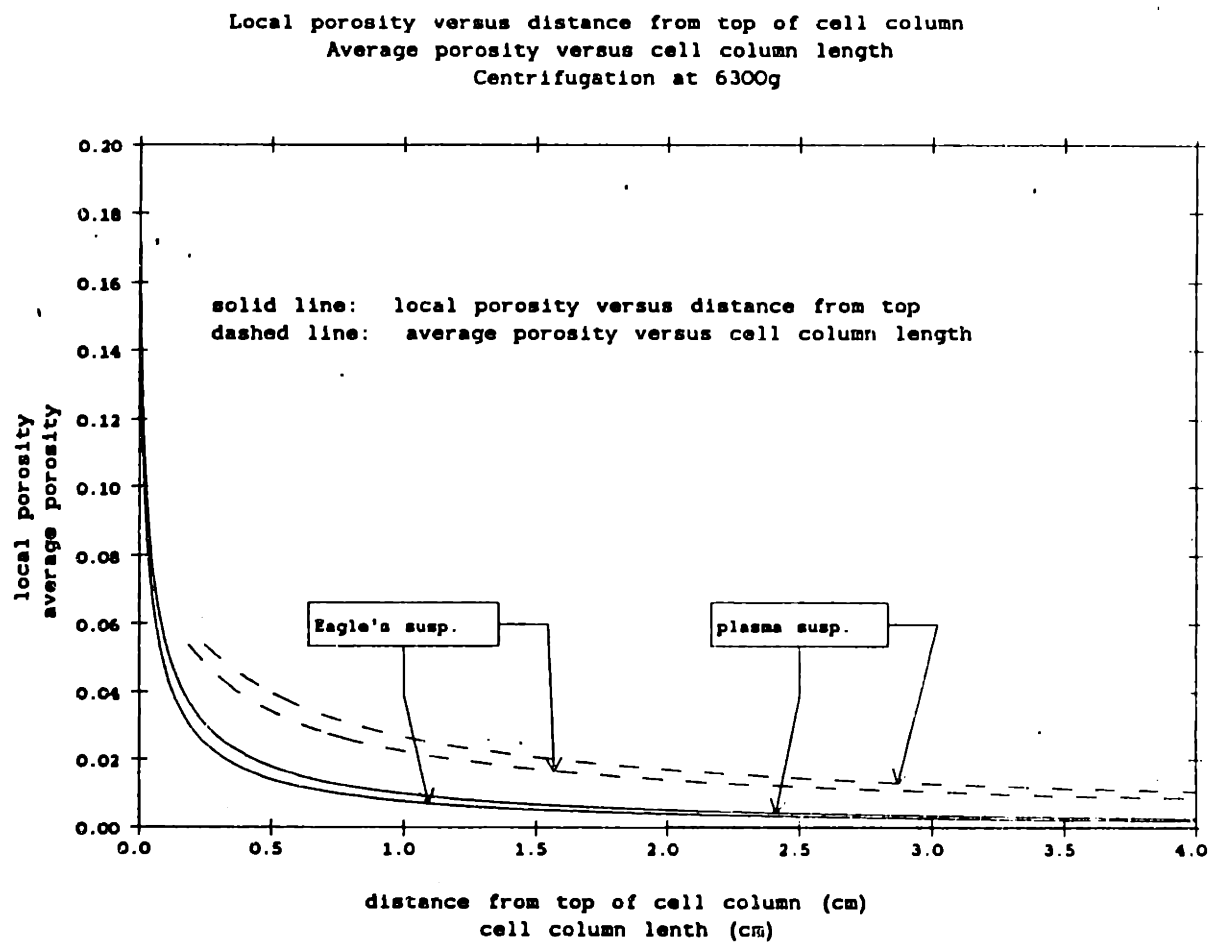
The relationship between the steady-state porosity of a centrifuged column of packed cells and the compressive pressure of centrifugation has been examined by Saltzman [1984] and Zydney *et al*. [1985]. Using a differential force balance on a layer of cells in the column, the compressive pressure  $\pi$  at any point in the column is given by

$$\pi = \frac{1}{(1-\epsilon)} \int_0^y (\rho_c - \rho_f) \omega^2 R d[(1-\epsilon)y] \quad (3.1)$$

where  $\epsilon$  is the porosity,  $\rho_c$  and  $\rho_f$  are the cell and suspending fluid densities, respectively,  $\omega$  is the rotational velocity of the centrifuge and  $R$  is the arm length of the centrifuge. An empirical expression relating compressive pressure to porosity derived by Zydney *et al*. is

$$K\pi = \left(\frac{\epsilon_0}{\epsilon}\right) + \left(\frac{\epsilon}{\epsilon_0}\right) - 2 \quad (3.2)$$

where  $K$  is a parameter fitted from experimental data and  $\epsilon_0$  is the porosity at zero compressive pressure, equal to 0.2. From experimental data at 592g and 925g  $K$  is determined to be  $4.0 \times 10^{-5} \text{ cm}^2/\text{dyne}$  or  $5.3 \times 10^{-2} (\text{mm Hg})^{-1}$ . Assuming that  $K$  is a property only of the material being centrifuged, these equations may be used to provide a value for fluid trapping in the microcentrifuge. Solution of equations (3.1) and (3.2) at 6300g for plasma and Eagle's suspensions are shown in Figure 3-1 for both the local porosity and the average porosity, defined as



**Figure 3-1:** Plots of  $\epsilon$  as a function of position in the cell column and of  $\bar{\epsilon}$  as a function of cell column length.

$$\bar{\epsilon} = \frac{1}{L} \int_0^L \epsilon \, dy \quad (3.3)$$

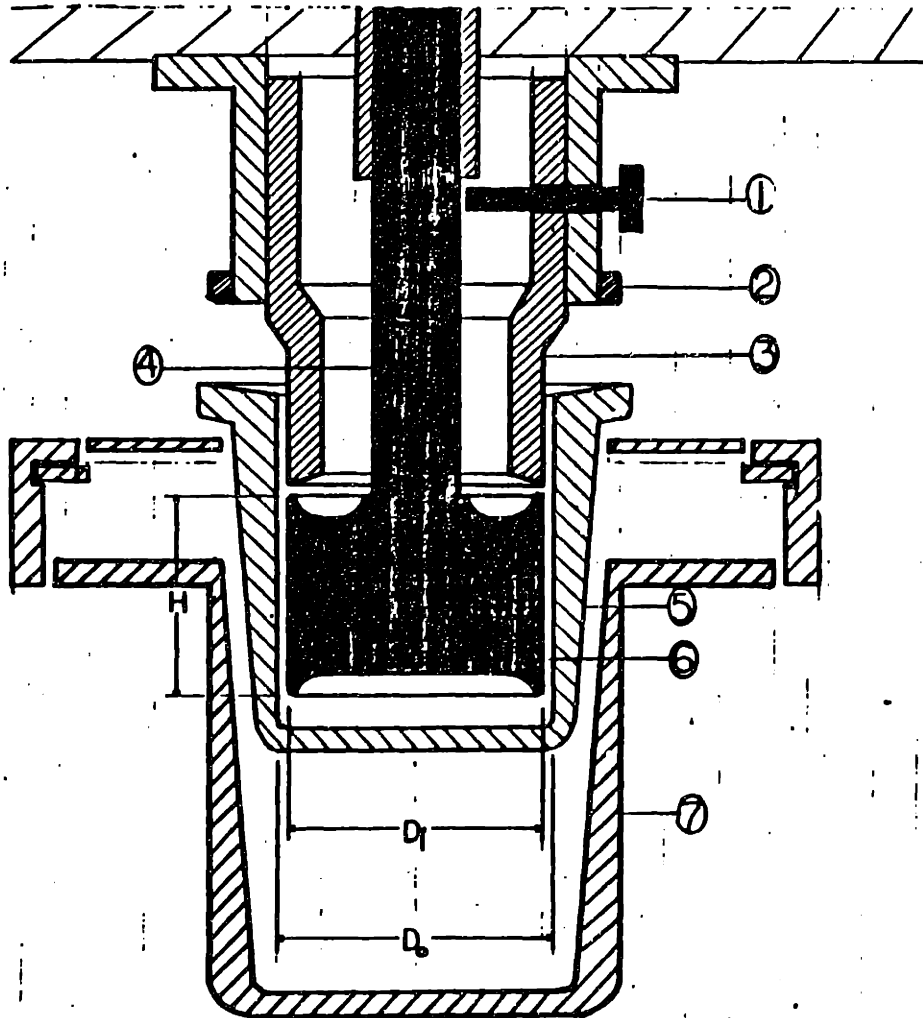
Values for the average porosity at the typical microcentrifuge cell column length of 3 cm are 0.01 for Eagle's solution and 0.012 for plasma, which is in agreement with the value reported by Chien *et al.* [1965]. The microhematocrit values were multiplied by  $(1-\bar{\epsilon})$  to obtain the hematocrits reported in the results. Because of these adjustments, the hematocrit  $H$  and cell volume fraction  $c$  are used interchangeably throughout this work.

### 3.3 Viscometry Measurements

#### 3.3.1 Techniques of Viscometry

Before discussing the particulars of the equipment and methods used for measuring blood viscosity, it will be helpful to first review some general considerations of blood viscometry. The rotating coaxial cylinder (Couette) viscometer is in widespread use for these applications. It consists of an inner cylinder placed in an outer cup, with a gap between the two where the sample of fluid to be assessed is placed (see Figure 3-2). Either the inner or the outer cylinder may be rotated, and the viscosity of the fluid being sheared causes a momentum flux which develops a torque which is measured on the cylinder not rotated. If the gap width between the cylinders is narrow (less than 5%) compared to the radius of the cup, the fluid will undergo simple, steady, laminar shear and the velocity profile within the gap can be approximated by a straight line. The shear rate in  $\text{sec}^{-1}$  in the annulus is given by

$$\gamma = 2\omega \frac{R_o^2}{R_o^2 - R_i^2} \quad (3.4)$$



**Figure 3-2:** The Haake ZA30 sensor system, an example of a coaxial cylinder viscometer.

1. Knurled screw to hold inner cylinder. 2. Clamping ring to hold guard ring.
3. Guard ring. 4. Inner cylinder. 5. Outer cylinder (beaker).
6. Gap where sample is located. 7. Beaker housing



where  $\omega$  is the angular velocity of the rotating cylinder in radians/sec, and  $R_o$  and  $R_i$  are the radii of the outer and inner cylinders, respectively. The torque across the gap is constant, so that if end effects are neglected, the shear stress in dyne/cm<sup>2</sup> exerted on the inner cylinder is given by

$$\tau = \frac{M}{2\pi h R_i^2} \quad (3.5)$$

where  $M$  is the torque exerted in g-cm and  $h$  is the length of the cylinder. Thus knowledge of  $\omega$ ,  $M$ , and the geometric parameters of the system leads to measurement of  $\gamma, \tau$ , and hence,  $\eta = \tau/\dot{\gamma}$ .

The interpretation of rheological data for blood is complicated by several effects, artifacts, and sources of error, some of which are common to all fluids while others are peculiar to blood. Included in the former are end and edge effects resulting from flows above and below the inner cylinder. These can be minimized by modification of the inner cylinder, either by recessing the ends or converting the bottom end into a conical configuration [Meiselman and Cokelet, 1973].

The air-fluid interface at the top of the viscometer can be a source of significant error. Even a simple Newtonian sample such as water exhibits surface effects from its high surface tension if a free meniscus develops in the gap. When plasma viscosities were initially measured [Copley *et al.*, 1960; Wells and Merrill, 1961; Cerny *et al.*, 1962] in viscometers without guard rings, shear-thinning behavior was reported due to the formation of a wax-like "skin" of lipids, lipoproteins, or denatured proteins. Employment of a guard ring such as that shown in Figure 3-2 eliminates the surface effects by preventing fluid which lies on the air-fluid interface from influencing the torque reading.

The presence of the walls of the cylinder necessarily introduces a inhomogeneity into the distribution of cells in the gap, since the center of any one

cell cannot lie closer to the wall than its radius. This wall exclusion effect was studied mathematically and experimentally by Vand [1948], who used glass spheres in coaxial and capillary viscometers. His model assumed a "rectangular" concentration profile of a layer of pure suspending fluid next to the walls and a bulk region of constant concentration which was later used by Quemada [1977] to derive his expression for blood viscosity, equation (2.11) (see Appendix C). In a monodisperse system, Vand found a particle-free layer with a width of 1.1 times the particle radius which resulted in a decrease in the measured wall shear stress below that expected for a homogeneous suspension. Cokelet *et al.* [1963] found that torque values of 2 to 5% higher could be obtained by the use of grooves 66 $\mu$ m deep in the walls of the cylinders. The necessity of the grooved surfaces is disputed by Chien *et al.* [1966], who note that their presence does not always prevent formation of the plasma layer [Merrill *et al.*, 1963a]. Also, it is desirable to obtain data at conditions which represent flow systems of interest (*e. g.*, blood vessels), and these do not have regularly grooved walls. The cylinders of the Haake viscometer used in this study had smooth surfaces.

The applicability of equations (3.4) and (3.5) depends on the maintenance of laminar, unidimensional (tangential) flow in the gap. If the inner cylinder is the rotor, centrifugal forces can introduce secondary flows and Taylor vortices and as the rotation rate increases, eventually lead to turbulent flow. For an annulus width of 1 mm and a outer radius of 2.5 cm, the shear rate at which secondary flow occurs for plasma at 37 °C is about 150 sec<sup>-1</sup> and for blood is about 300-450 sec<sup>-1</sup>. When the outer cylinder is the rotor, as in the case of the present study, inertial forces tend to stabilize the flow and the shear rate at which secondary flow begins is around 10<sup>5</sup> sec<sup>-1</sup> [Meiselman and Cokelet, 1973].

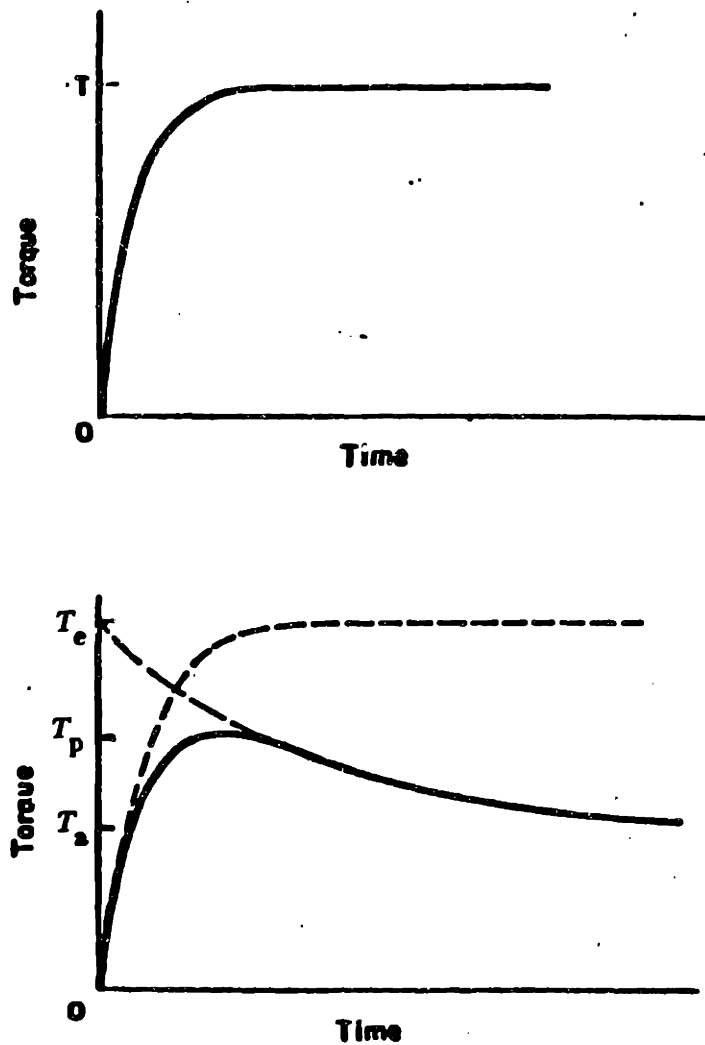
Other considerations are more specific to blood. The difference in density

between the erythrocytes (approximately  $1.09 \text{ g/cm}^3$ ) and plasma (approximately  $1.02 \text{ g/cm}^3$ ) results in a tendency of red cells to settle in plasma solutions. A homogeneous sample can be maintained in the coaxial cylinder viscometer by having the region above the inner cylinder covered with a pool of blood [Meiselman and Cokelet, 1973].

The torque-time curves for blood show two regimes of behavior: At shear rates above  $1 \text{ sec}^{-1}$  the torque reading rises quickly (on the order of a second) to a constant value which is maintained until the shear is stopped. At shear rates above  $10 \text{ sec}^{-1}$  the time for this rise is negligible. At rates below  $1 \text{ sec}^{-1}$  a more gradual rise to a maximum is observed followed by a decay to an equilibrium value. The time for this rise can be as long as a minute or more at shear rates of  $0.01 \text{ sec}^{-1}$ . Direct visualization of the annulus shows that at low shear rates, the red cells migrate away from the outer and possibly from the inner walls, forming a substantial cell-free layer near the walls [Cokelet *et al.*, 1963; Meiselman, 1965]. It has been hypothesized that the slow rise is due to transient effects from initiation of the shear and that the decay represents migration of the cells away from the cylinder walls, which can be represented mathematically as an exponential decay [Cokelet and Smith, 1973] (see Figure 3-3). The argument thus is that the correct torque is one which is extrapolated to zero time from the decaying exponential curve. As will be discussed in the next section, there was no torque decay at the lowest observable shear rate of the present study. A theoretical model for blood flow at low shear rates in a coaxial cylinder viscometer is given by Bloor [1982].

### **3.3.2 Viscometer**

Rheological measurements were performed with a Haake Rotovisco RV100 with the CV100 measuring system (Haake, Inc., Saddle Brook, NJ). The instrument



**Figure 3-8:** Torque-time curves for blood.

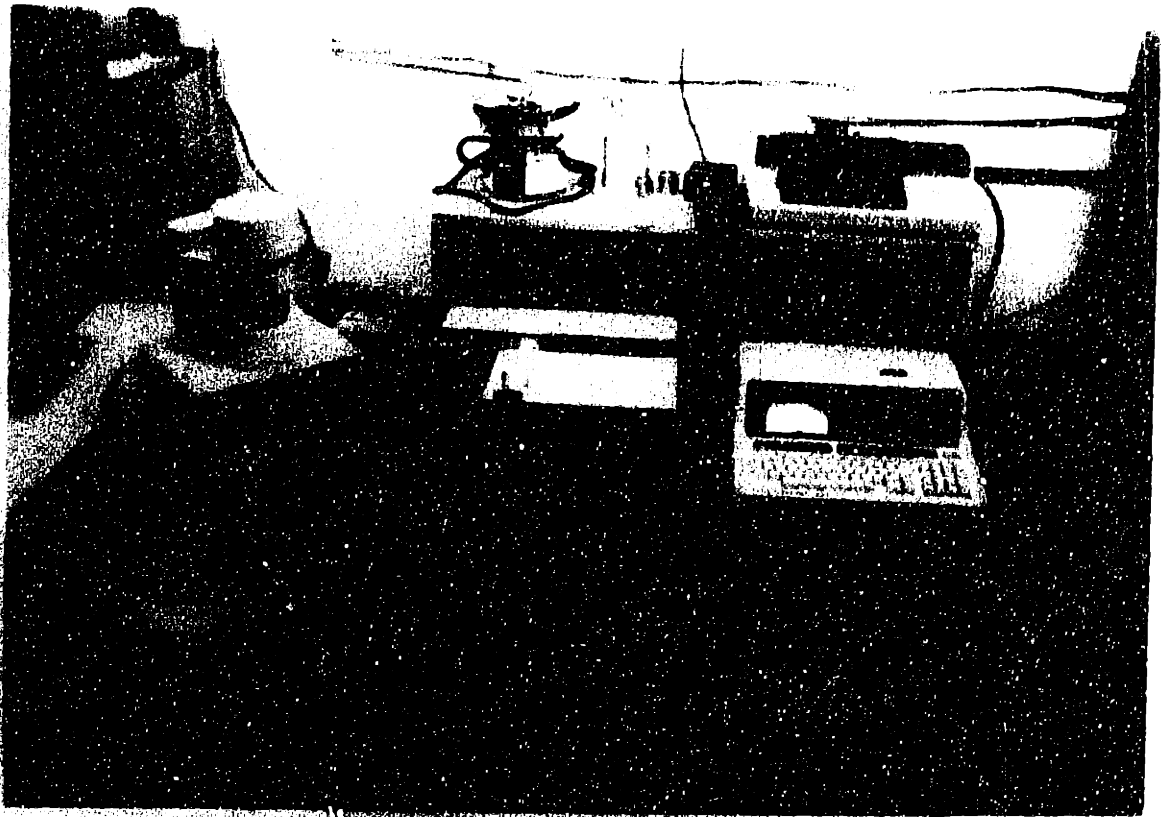
*Top*, shear rates greater than  $1 \text{ sec}^{-1}$ . *Bottom*, shear rates less than  $1 \text{ sec}^{-1}$ .

$T_a$  is the asymptotic equilibrium torque value measured,  $T_p$  is the peak torque, and  $T_e$  is the torque extrapolated to time zero. From Cokelet and Smith [1973].

is the property of the Ceramics Laboratory of the Department of Materials Science and Engineering at M.I.T. and is pictured in Figure 3-4. The ZA30 sensor system was used, which is a stainless steel coaxial cylinder system (see Figure 3-2, page 52). The outer cup is driven by a motor, which provides greater stabilization against secondary flows as described previously. The viscosity of the fluid being sheared causes a momentum flux which develops a torque on the inner cylinder, deflecting measuring springs in the drive unit and producing an electrical signal proportional to the torque. The relevant dimensions for the ZA30 are listed in Table 3-1. The inner cylinder slides into the measuring shaft of the unit and is held on the shaft by a knurled screw. It is mechanically positioned and centered by an air bearing in the measuring system which is supplied with compressed air from an air supply unit (Haake model LV100). The air pressure in the supply unit reservoir is maintained between 5.5 and 3.5 bar. The feed pressure to the bearing is 2.5 bar.

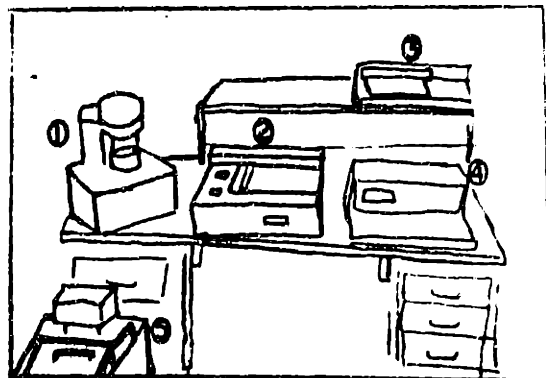
The top and bottom surfaces of the inner cylinder are not flat but are recessed from the outer edge to minimize end effects as shown in Figure 3-2. A guard ring, secured on the measuring shaft by a clamping ring, is employed in all measurements to minimize artifacts from the fluid-air interface. A refrigerated bath and circulator (Haake model F3-K) supplies a water jacket in the measuring system with a supply of distilled water, which is maintained at a constant temperature by a balance of controlled heating against continuous cooling. The water jacket is used to set sample temperature and to remove dissipated heat from the sample created by shearing.

The viscometer is linked to an HP-85 Desk Top Computer (Hewlett-Packard Corporation, Corvallis, OR), which in turn is connected to an HP 7470A Graphics Plotter. The viscometer interface is a Haake DVM 7615 four-channel digital voltmeter with three control lines connected to an HP 82940A I/O interface.



**Figure 3-4: Haake Rotovisco RV100 viscometer with accesories as described in text.**

- 1. CV100 measuring system**
- 2. RV100 viscometer**
- 3. Graphics Plotter**
- 4. HP-85 computer**
- 5. LV100 air supply unit**



Inner cylinder diameter ( $D_i$ )	27.83 mm
Inner cylinder length (h)	26.9 mm
Outer cylinder diameter ( $D_o$ )	30 mm
Gap width	1.085 mm
Sample volume	5 ml

**Table 3-I:** Geometric parameters for the ZA30 sensor sytem.

A Haake rotational software cartridge (part #1077922) is used to program the computer, which can then be used to set the viscometer parameters and control the measuring operation. The viscometer may be run in either a  $\tau$  vs. time mode at constant shear rate, or in a  $\tau$  vs. shear rate mode with a maximum range of 0 to 300  $\text{sec}^{-1}$ . The software reads 100 data points evenly distributed over time or over the selected range of shear rates and can print digital values or plot data on the Graphics Plotter.

### 3.3.3 Viscometry Methods

Enough sample to cover the inner cylinder, about 5 to 7 ml in volume, is carefully mixed and poured into the cup, which is covered with a piece of Parafilm and placed into the measuring unit for about two minutes to allow temperature equilibration at 37 °C. While the sample is warming, the guard ring and inner cylinder are mounted following the procedure recommended by the instrument instructions: The clamping ring is loosened and the guard ring slides into the housing, resting at the upper stop, and is clamped into position. The inner cylinder slides into the measuring shaft and is secured by the knurled screw. The guard ring is now loosened and allowed to rest on the upper face of the cylinder. The clamping ring is again tightened, which lifts the guard ring to a level of 0.1 to 0.2 mm above the inner cylinder. For highly viscous samples it is necessary to elevate the guard ring slightly more.

The shaft is lowered into the cup and then briefly raised up to overcome surface tension effects in the annulus and to observe whether the sample covers the gap between the guard ring and the inner cylinder. The sample is just visible at the border of the outer cylinder, providing a small pool of blood to maintain homogeneity in the annulus [Meiselman and Cokelet, 1973]. The exposed shaft is



encased with a Plexiglas housing and then the entire measuring system is covered with a large cardboard box to eliminate artifacts from outside air currents.

A fresh piece of plotting paper is placed in the recorder and calibrated. The scale parameters for the recorder are set according to the particular sensor system used. For the ZA 30 these values are 0.3 Pa/scale division for the  $\tau$  axis and 3  $\text{sec}^{-1}$ /scale division for the  $\dot{\gamma}$  axis. The mode of operation is set to automatic reversal, that is, to step up continuously from 0 to 300  $\text{sec}^{-1}$  over a period of one minute, to hold at 300  $\text{sec}^{-1}$  for six seconds, and then to step down back to 0  $\text{sec}^{-1}$ . (There is an associated hysteresis that is reported on in Chapter 4.) The full-scale factor for the  $\tau$  axis of the recorder was set to the minimum value possible. The run is initiated.

Upon completion, a hard copy of the data points is obtained from the HP-85. Repeat measurements are made from 0 to 30  $\text{sec}^{-1}$  and from 0 to 3  $\text{sec}^{-1}$  with appropriate modifications in the full-scale setting for the  $\tau$  axis. The sample is then removed and discarded and the cup, inner cylinder, and guard ring are washed out with hot water.

In order to insure that the shear induced by the instrument would not cause hemolysis sufficient to affect the measurements, a series of samples had their viscosities measured first at a constant shear rate of 0.3  $\text{sec}^{-1}$ , then were subjugated to constant shear rates of 300  $\text{sec}^{-1}$  for a period of a minute, and then viscosities were remeasured at 0.3  $\text{sec}^{-1}$ . No significant change was noted in the values at the low shear rate. Secomb *et al.* [1983a], in the direct microscopic observation of concentrated red cells at shear rates up to 1585  $\text{sec}^{-1}$ , did not report on any significant hemolysis.

An assessment of the torque-time curves for concentrated suspensions of blood showed no transient effects such as those seen in Figure 3-3 at shear rates of 0.3

sec<sup>-1</sup>, which is the lowest constant shear rate at which the viscometer could be run. There was agreement between those viscosity values obtained in the constant shear rate mode and the corresponding value obtained with continuous variation of shear rates.

### 3.4 Data Analysis

The data was fitted to several functions, both linear and non-linear. A general optimization routine was written in FORTRAN and run on the VAX-11 computer of the Whitaker College Computer Facility. The program was adapted from one by Zydney [1985] and is given in Appendix D. The general principals behind the data analysis will be presented here, as developed by Zydney [1985].

Determination of the optimal set of parameters  $b_1, b_2, \dots, b_P$  for a function  $Y$  which describes a set of data of  $N$  points can be equated to minimizing  $E$ , the value of the sum of squared residuals

$$E = \sum_{i=1}^N [y_i - y_{i,exp}]^2 \quad (3.6)$$

where  $y_i$  is the calculated value from the function  $Y$  at the same set of conditions as the experimental value  $y_{i,exp}$ .

Minimizing  $E$  is equivalent to finding values of  $b_i$  such that the gradient of  $E$  is zero. In mathematical terms, this is represented as a set of  $P$  coupled equations

$$F_j(b) = \frac{\partial E}{\partial b_j} = 2 \sum_{i=1}^N [y_i - y_{i,exp}] \frac{\partial y_i}{\partial b_j} = 0 \quad (3.7)$$

The method of steepest descent [Himmelblau, 1970] is used to find the set of parameters which minimize  $E$ . The set of equations represented by equation (3.7) can be solved by a Newton-Raphson iteration,

$$\mathbf{b}^{j+1} = \mathbf{b}^j - (\mathbf{A}^j)^{-1} \mathbf{F}^j \quad (3.8)$$

where  $\mathbf{A}$  is the jacobian maxtrix with components

$$A_{ij} = \frac{\partial F_i}{\partial b_j} \quad (3.9)$$

The derivatives of  $y_i$  are evaluated numerically as

$$\frac{\partial y_i}{\partial b_j} = \frac{y_i(b_j + \Delta b) - y_i(b_j)}{\Delta b} \quad (3.10)$$

The derivatives of  $F$  are found in a similar fashion.

After the data has been fit to the equation, various regression parameters are calculated from the covariance matrix:

$$\text{cov}(b) = \frac{E}{N-P} (\mathbf{X}^t \mathbf{X})^{-1} \quad (3.11)$$

where  $\mathbf{X}$  is the matrix whose elements are

$$X_{ij} = \frac{\partial y_i}{\partial b_j} \quad (3.12)$$

and  $\mathbf{X}^t$  is the transpose of  $\mathbf{X}$ . The variance of the fitted parameter  $b_i$  is given by the diagonal elements of the covariance matrix.

$$\text{var}(b_i) = \text{cov}(b_i, b_i) \quad (3.13)$$

The correlation coefficient between any two parameters  $b_i$  and  $b_j$  is

$$\rho_{ij} = \frac{\text{cov}(b_i, b_j)}{[\text{var}(b_i) \cdot \text{var}(b_j)]^{0.5}} \quad (3.14)$$

The 95% confidence limits for a function of the fitted parameters  $Y=Y(b)$  is evaluated first by finding the variance of  $Y$  from the variances and covariances of the  $R$  parameters [Freund and Minton, 1979]:

$$\text{var}(Y) = \sum_{i=1}^R \left( \frac{\partial Y}{\partial b_i} \right)^2 \cdot \text{var}(b_i) + 2 \sum_{i=2}^R \sum_{j < i}^{R-1} \left( \frac{\partial Y}{\partial b_i} \right) \left( \frac{\partial Y}{\partial b_j} \right) \text{cov}(b_i, b_j) \quad (3.15)$$

where the second term on the right represents sums over all possible combinations of  $i$  and  $j$  such that  $i > j$ . The standard deviation  $\sigma_Y$  at any point on the function  $Y$  is given by the square root of the variance. The 95% confidence interval at any point on  $Y$  is given by

$$Y - t_{0.025} \sigma_Y < Y < Y + t_{0.025} \sigma_Y \quad (3.16)$$

where  $t_{0.025}$  is the area under one tail of the Student t-distribution.

To avoid weighing the fits toward minimizing residuals at high values of  $\eta$ , the residuals were weighted by the reciprocal of the calculated value. The actual form of equation (3.6) used is

$$E = \sum_{i=1}^N \left[ \frac{\eta_i - \eta_{i,exp}}{\eta_i} \right]^2 \quad (3.17)$$

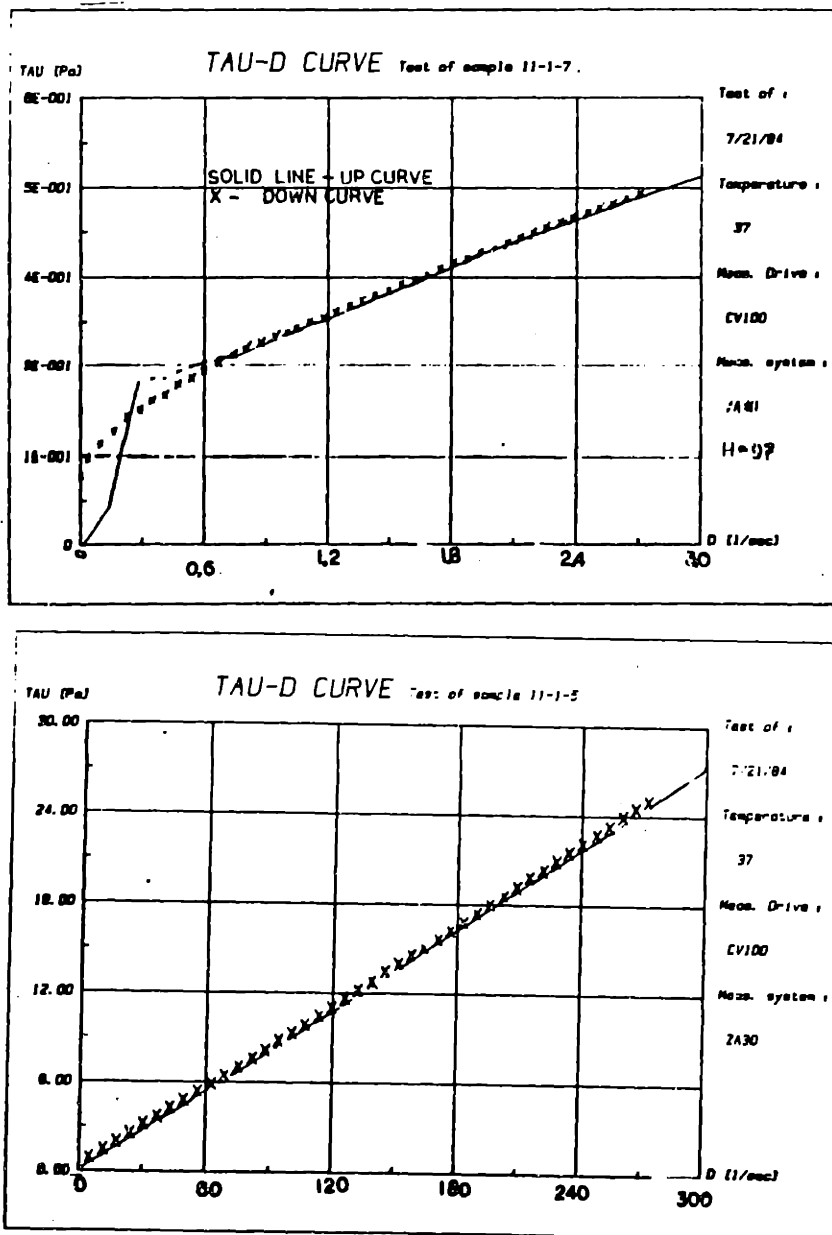
## Chapter 4

# RESULTS

### 4.1 General comments

Examples of the raw data from the Haake viscometer as printed by the Graphics Plotter are shown in Figure 4-1. Each sample was run in three ranges of operation: high shear from 0 to 300  $\text{sec}^{-1}$ , medium from 0 to 30  $\text{sec}^{-1}$ , and low from 0 to 3  $\text{sec}^{-1}$ . The plot of shear stress vs. shear rate in the lowest range shows two discernible artifacts: (a) a period in the step-up curve representing a start-up lag during which time the cylinder is rotating but the transducer as yet records no reading, and (b) a hysteresis between the step-up and step-down curves, with the step-down curve representing values on the order of 1 to 3 per cent higher than the step-up curve. This hysteresis, which existed to a smaller extent at higher shear rates, is caused by the time lag between the shear rate imposed and the shear rate measured. Its magnitude diminished as the rate of change of shear rate was decreased, but it could never be completely eliminated. The HP-85 computer printed out a hard copy of 100 data points for each run, and a representative sampling of points were selected from these lists for fitting. Approximately every fifth point from both the up and the down curves was used, averaging twenty data points obtained per run. From the up curve, no points were obtained from the region of the start-up artifact, and in compensation there was a higher rate of selection of points from the down curve in this region.

Measurement of plasma and saline viscosities at 37 °C showed a small but significant shear-thickening behavior which could have resulted from the



**Figure 4-1: Plots of  $\tau$  vs.  $\dot{\gamma}$  from H-P Graphics Plotter interfaced with Haake viscometer, H=97.**

Solid line=up curve. X=down curve.

development of secondary flows in the viscometer [Merrill, 1969]. The viscosity measured for Eagle's solution ranged from 0.85 to 0.9 cP. These values are somewhat higher than literature values of 0.69 cP for water at 37 °C, and the discrepancy is probably due to viscometer limitations in measurement of low-viscosity samples. In order to maintain consistency with other published data, 0.69 cP is used as the pure fluid viscosity for all saline solutions reported here. Plasma viscosities were more variable, ranging from 1.03 to 1.21 cP, which are comparable to typical literature values of 1.2 cP. For plasma suspensions the individual plasma values were used for calculating relative viscosity.

A series of measurements using freshly drawn blood obtained from a single source (the author) was used in order to check the reproducibility of the experimental method. Viscosity differences between different samples at the same hematocrit were on the order of 5% or less at hematocrits greater than 95. At hematocrits above 97 the sensitivity to concentration was enough to require that hematocrit be measured to the nearest tenth of a unit. At lower hematocrits it was sufficient to measure hematocrit to within a unit.

Tables of the data used in this work are presented in Appendix F. For definitional purposes, throughout the remainder of this work a *unit* of blood refers to all of the blood contained in the original bag obtained from the Red Cross. Each unit provided one or more *batches* of blood. A batch of blood is defined as an amount of blood removed from a unit and processed at one particular time, which is then divided into several *samples* of varying hematocrit. Sample identification is by the batch number followed by the sample number, *e. g.*, 19-6. The unit identifications for each sample are given in Appendix F.

Variations in viscosity from batch-to-batch and from unit-to-unit were shear-dependent, with the greatest differences at very low shear rates. Different batches

from the same unit generally had comparable values at all shear rates, within 5% of each other. Samples from different units were within 5% at high shear rates, but could be separated by up to a factor of two or more at shear rates below  $0.1 \text{ sec}^{-1}$ . At shear rates above  $0.1 \text{ sec}^{-1}$ , viscosity differences were usually small. Some of the variability can be traced to the age of the blood samples, and this is expanded on later in this chapter.

## 4.2 Viscosity of RBC suspended in saline solution

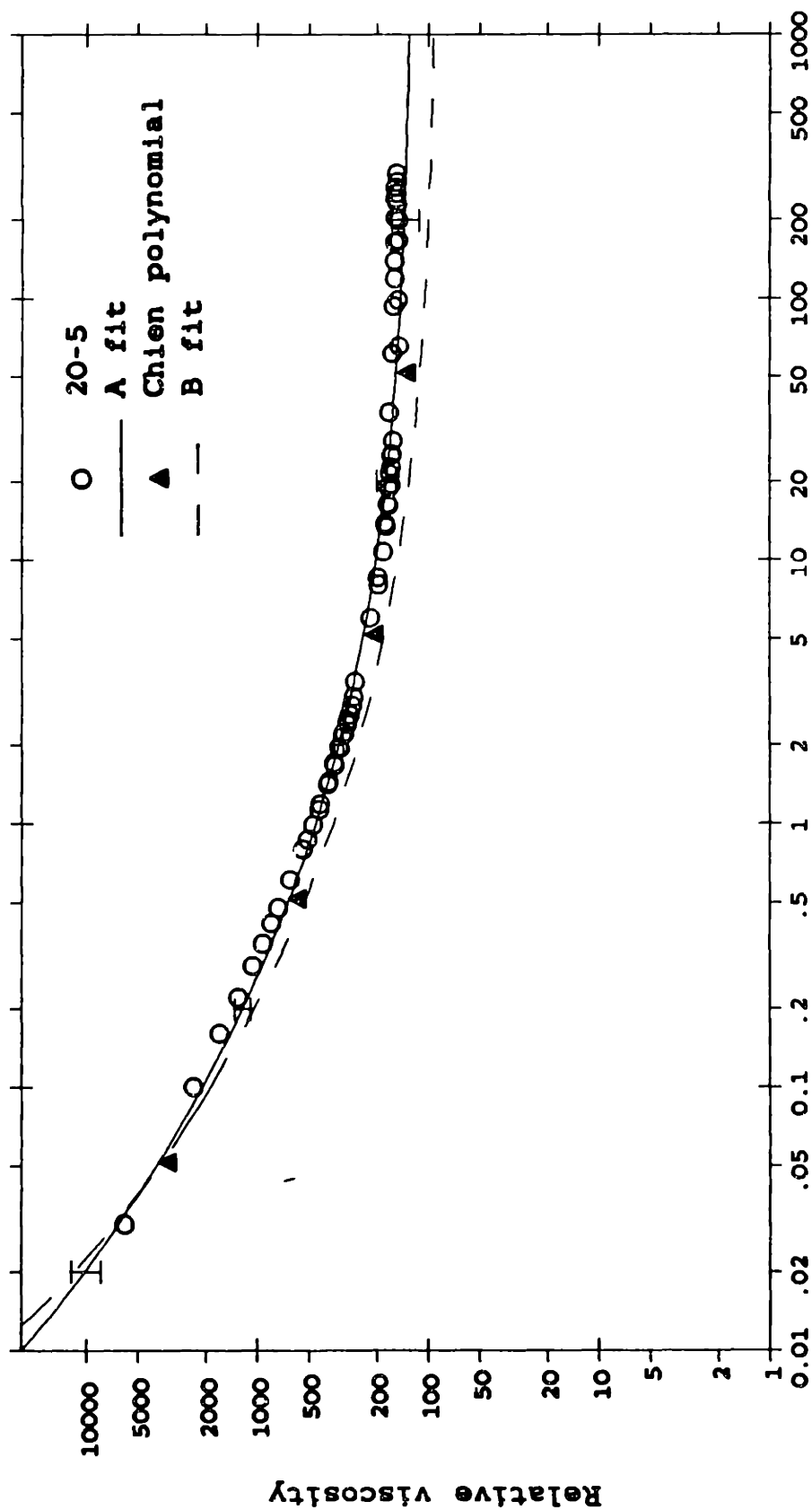
The data for RBC suspensions in Eagle's solution are presented in Figures 4-2 to 4-25 (pages 69 to 92). Each plot is one of relative viscosity ( $\eta_r$ ) versus shear rate ( $\dot{\gamma}$ ) at a different hematocrit. Hematocrits range from 98.4 to 7. Relative viscosities were calculated by dividing the measured viscosity by 0.88 cP, which was used as the pure Eagle's solution viscosity for all samples. On any one figure, data from different samples are represented by different symbols. Different samples from the same unit have the same designation with regard to the symbol being open or filled. Thus, in Figure 4-10, the two batches used are from the same unit, while in Figure 4-11, the two batches are from different units. There is no correlation of symbols between different figures.

The filled triangles in each figure are data values derived from the Chien *et al.* [1966] equation, (2.8), for Ringer's solution suspensions. Since the Chien *et al.* data are for viscosity and not relative viscosity, an estimate for  $\eta_r$  was made by taking the average values of the polynomials at zero hematocrit. In this manner, for Ringer's saline,  $\eta_r = 0.692 \text{ cP}$ . The solid curves with error bars and the dashed curves represent fits of the data to the Quemada equation, equation (2.11), and will be discussed in detail in Chapter 5.

On visual inspection, the points on each plot may appear to be segregated into



Relative viscosity vs. Shear rate  
RBC in Eagle's solution, H=98.4



Shear rate (1/sec)

Figure 4-2: Eagle's solution data, H=98.4.

Relative viscosity vs. Shear rate  
RBC in Eagle's solution, H=97.2

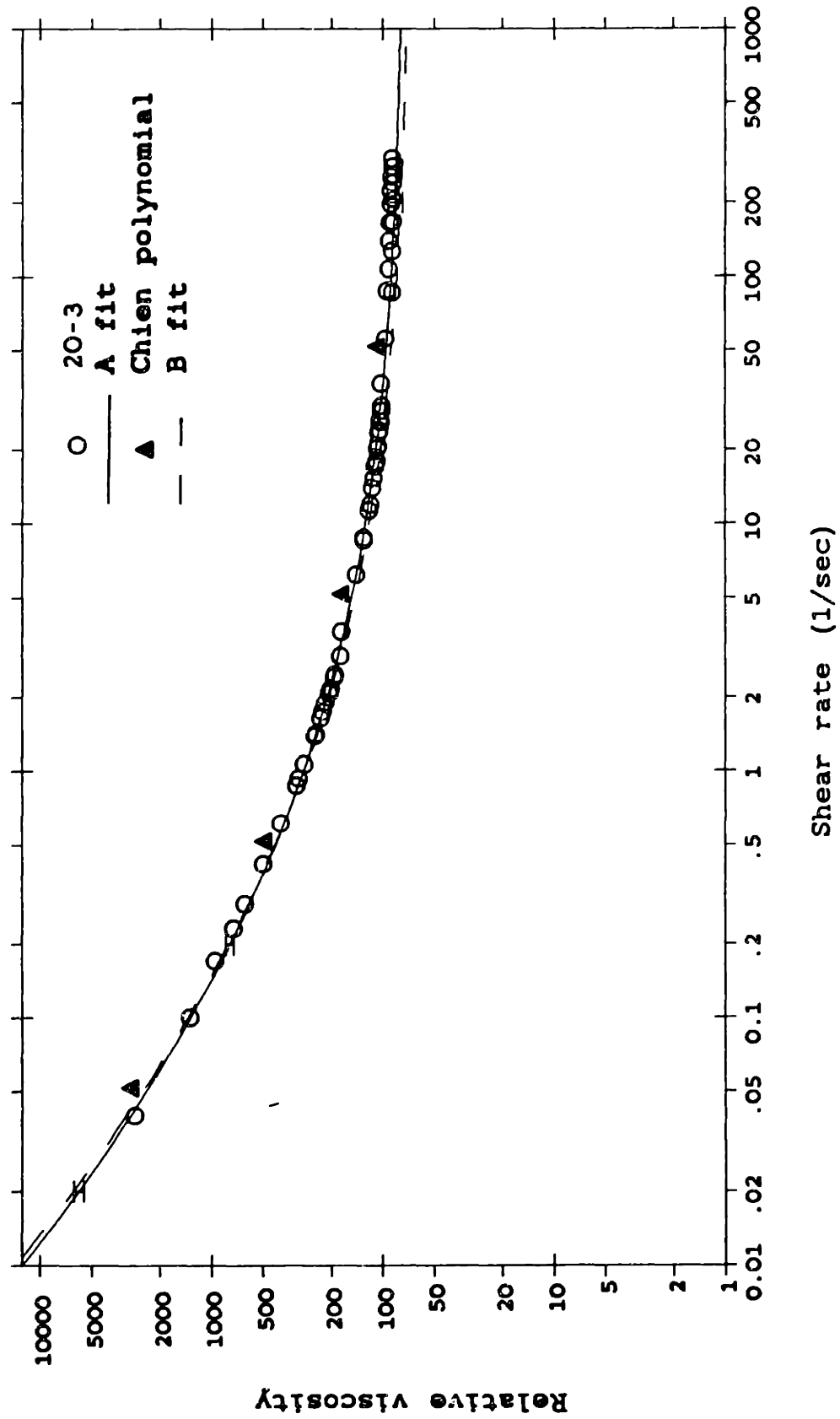
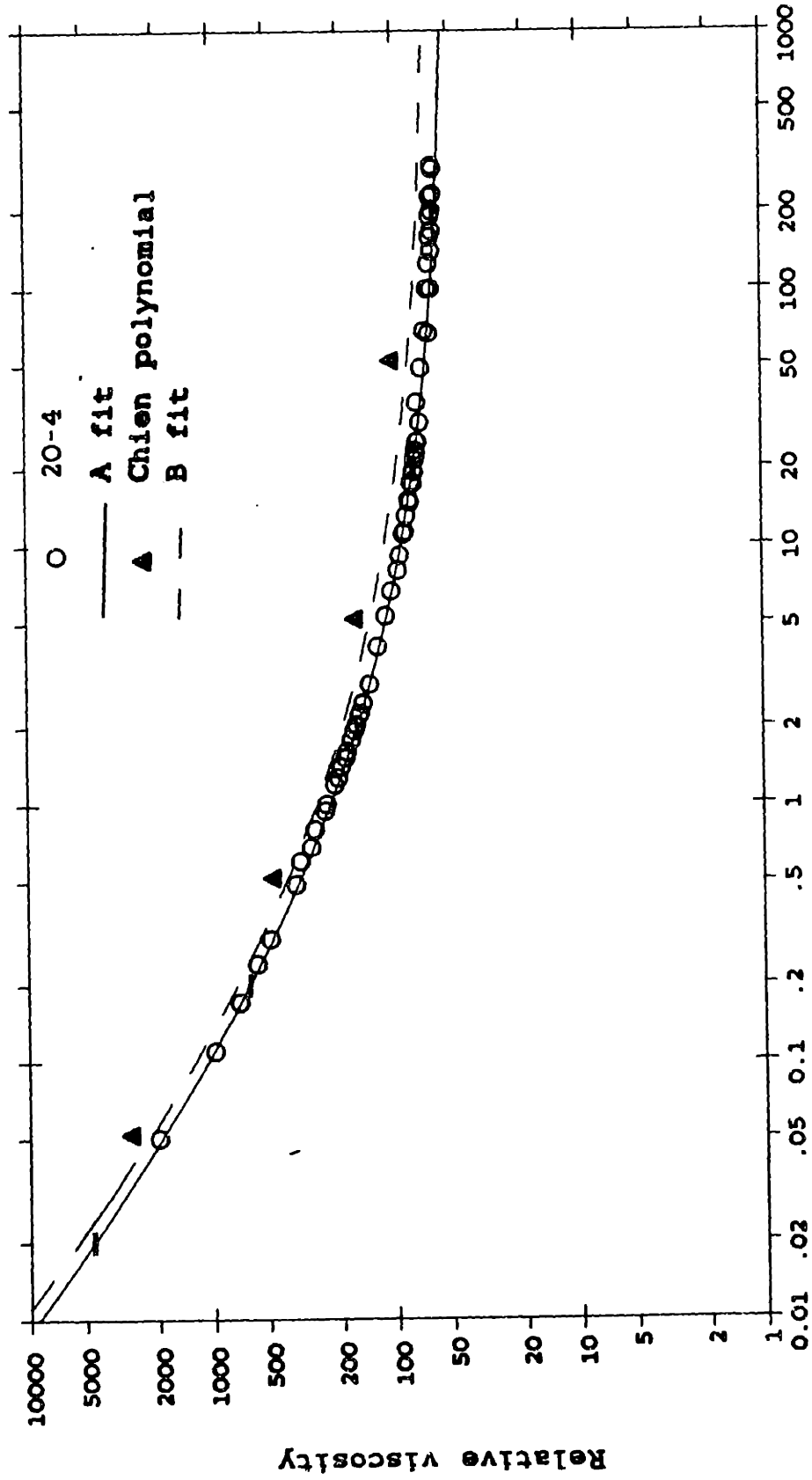


Figure 4-3: Eagle's solution data, H=97.2.

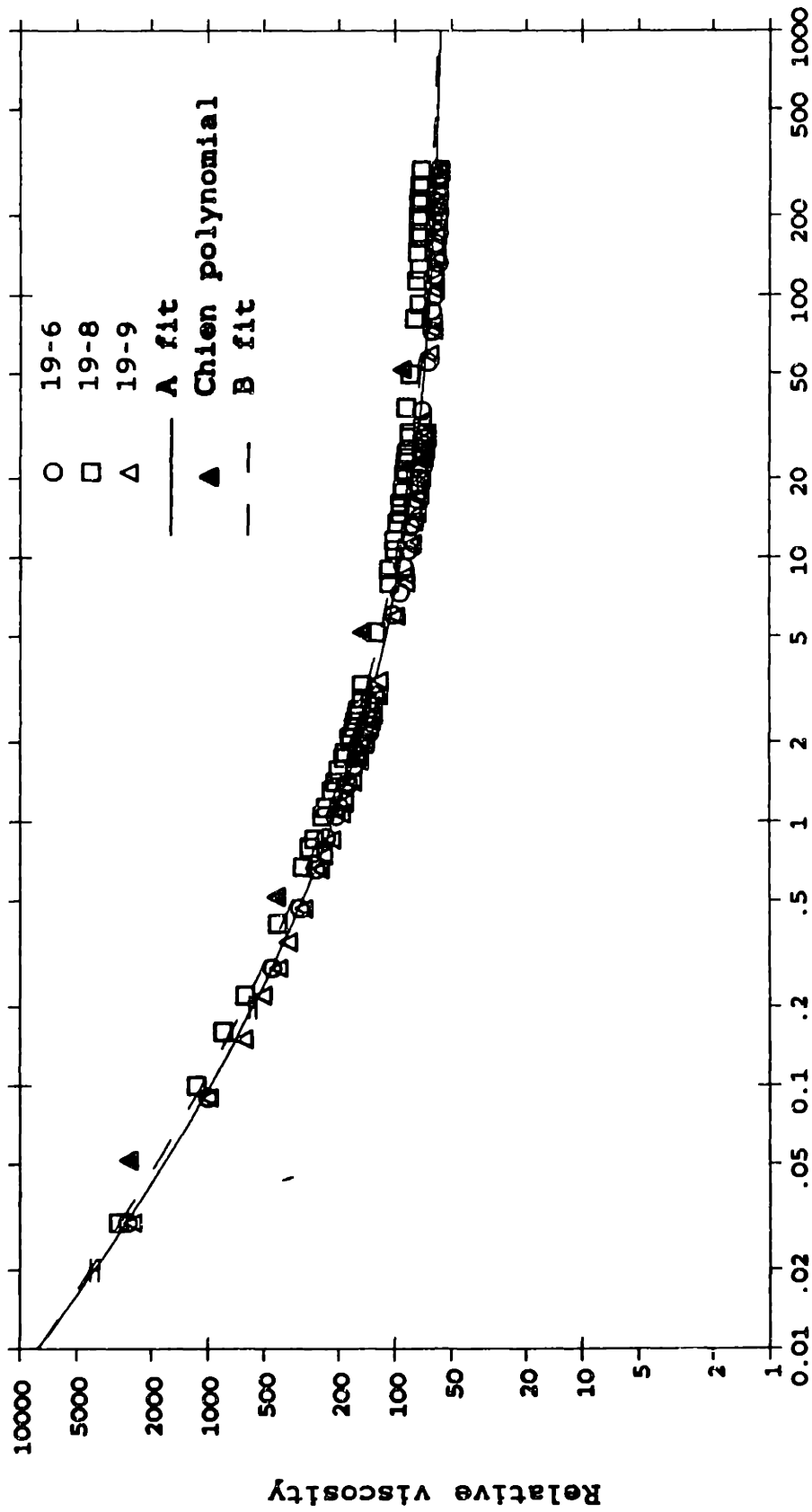
Relative viscosity vs. Shear rate  
RBC in Eagle's solution, H=96.8



Shear rate (1/sec)

Figure 4-4: Eagle's solution data, H=96.8.

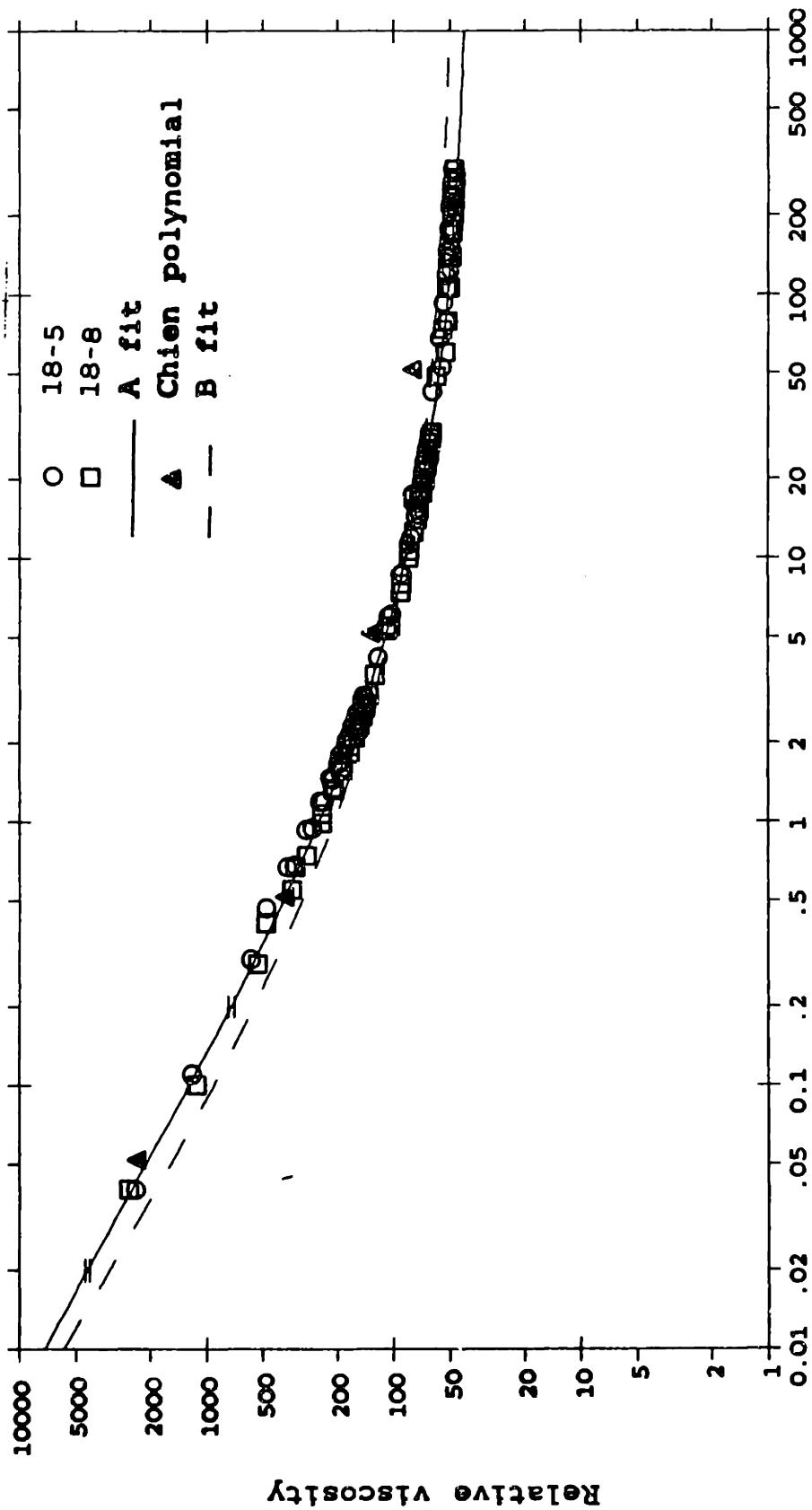
Relative viscosity vs. Shear rate  
RBC in Eagle's solution, H=96



Shear rate (1/sec)

Figure 4-5: Eagle's solution data, H=96.

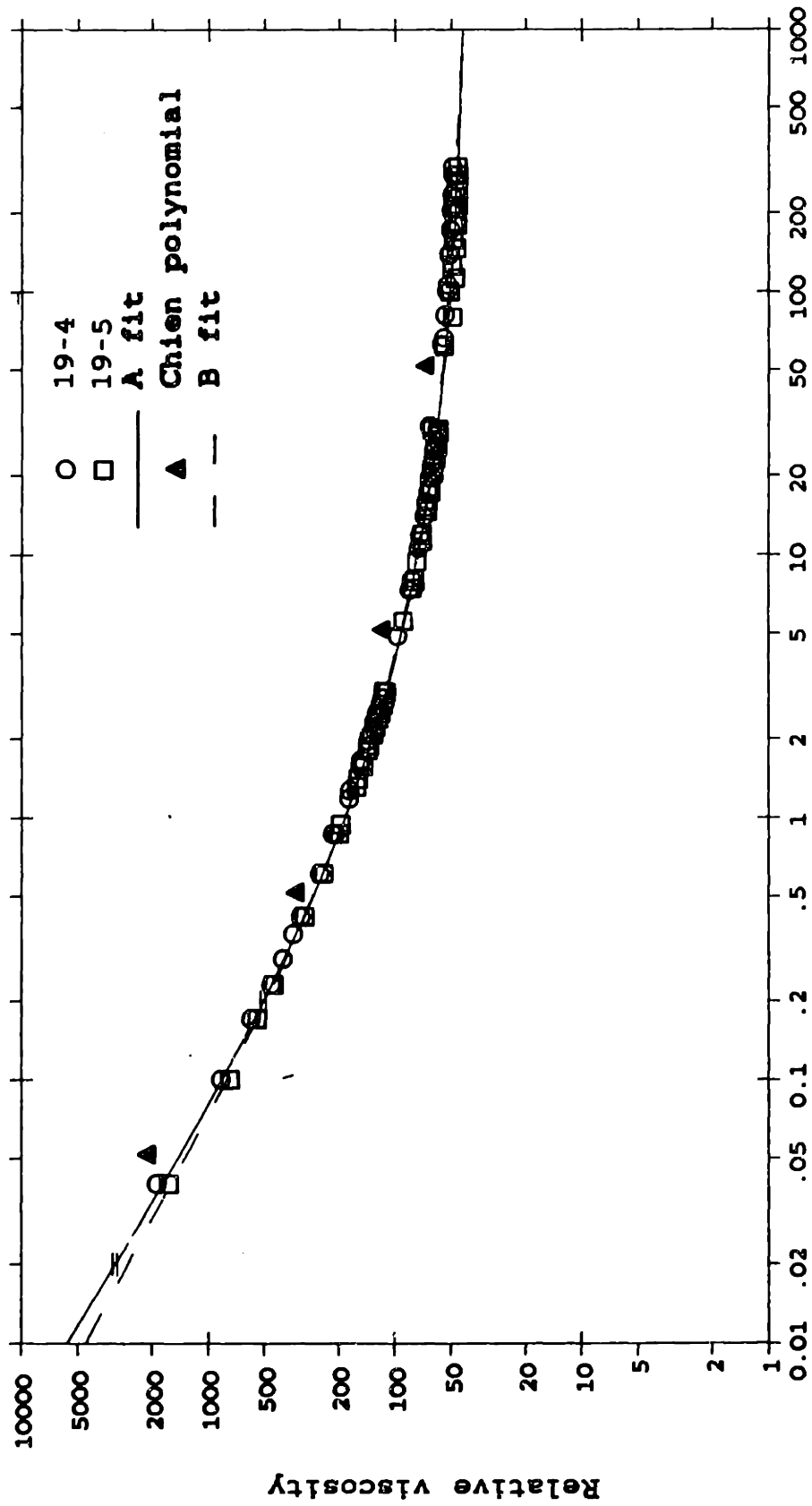
Relative viscosity vs. Shear rate  
RBC in Eagle's solution, H=95



Shear rate (1/sec)

Figure 4-6: Eagle's solution data, H=95.

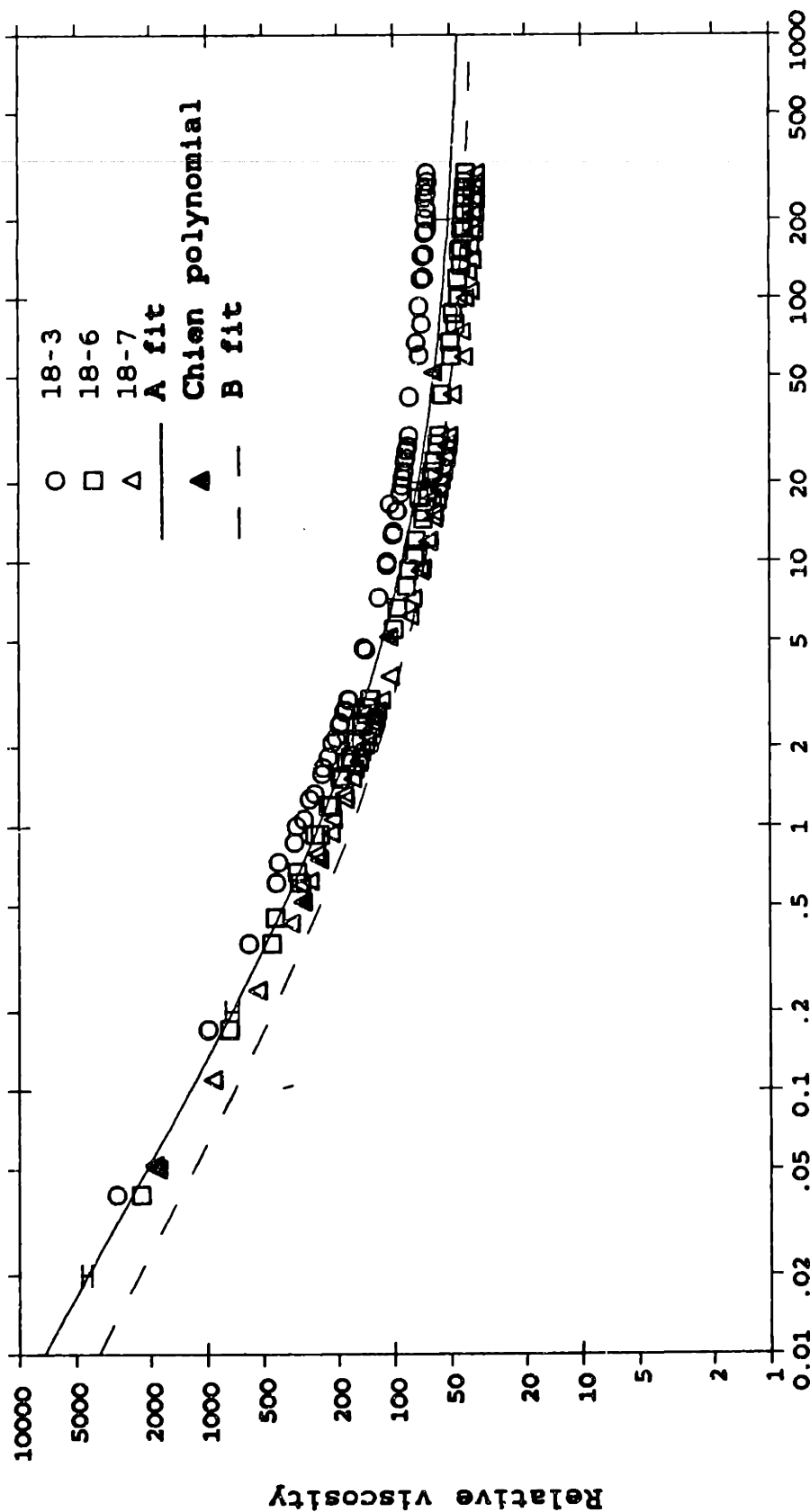
Relative viscosity vs. Shear rate  
RBC in Eagle's solution, H=94



Shear rate (1/sec)

Figure 4-7: Eagle's solution data, H=94.

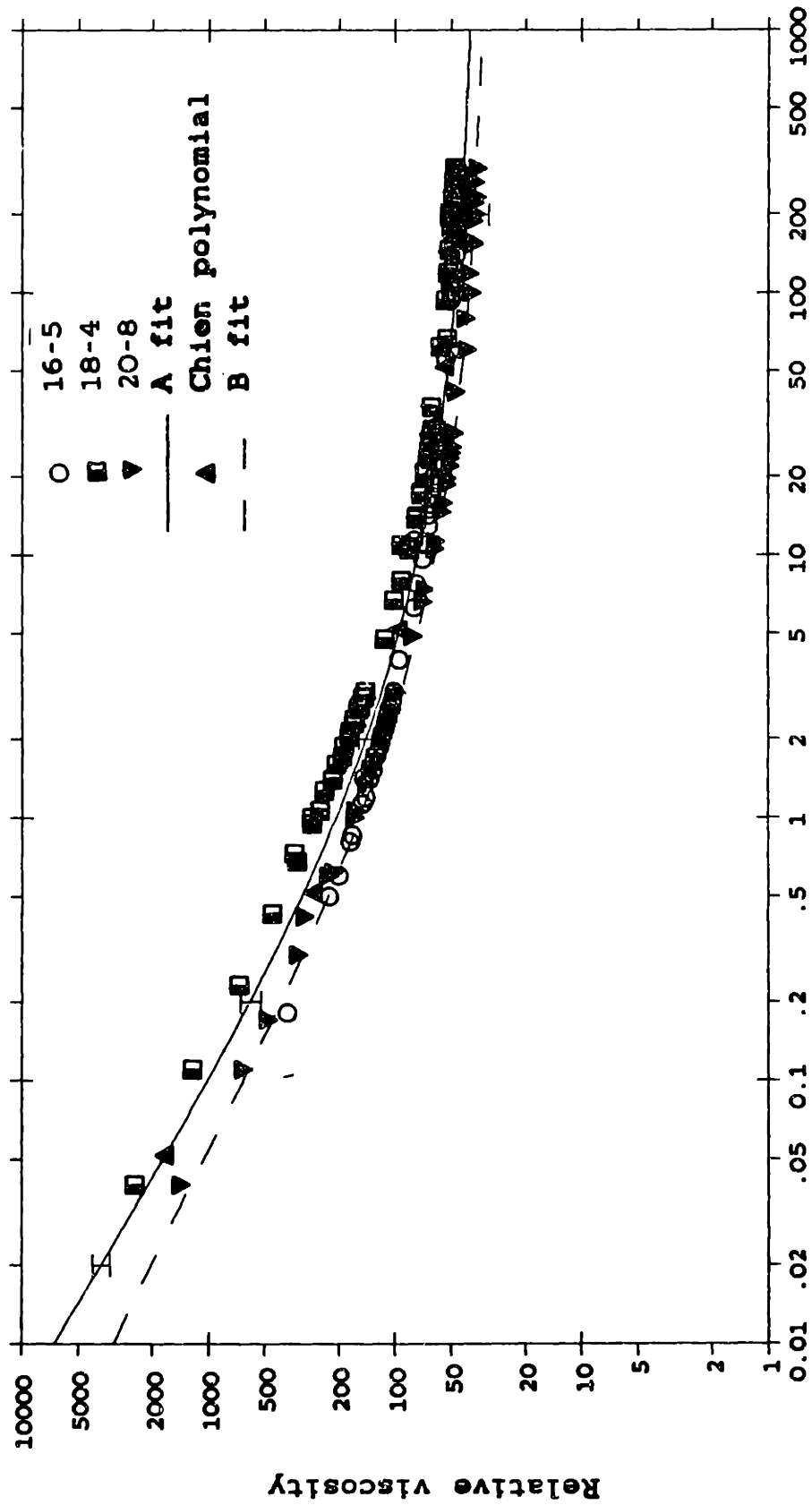
Relative viscosity vs. Shear rate  
RBC in Eagle's solution, H=93



Shear rate (1/sec)

Figure 4-8: Eagle's solution data, H=93.

Relative viscosity vs. Shear rate  
RBC in Eagle's solution, H=92

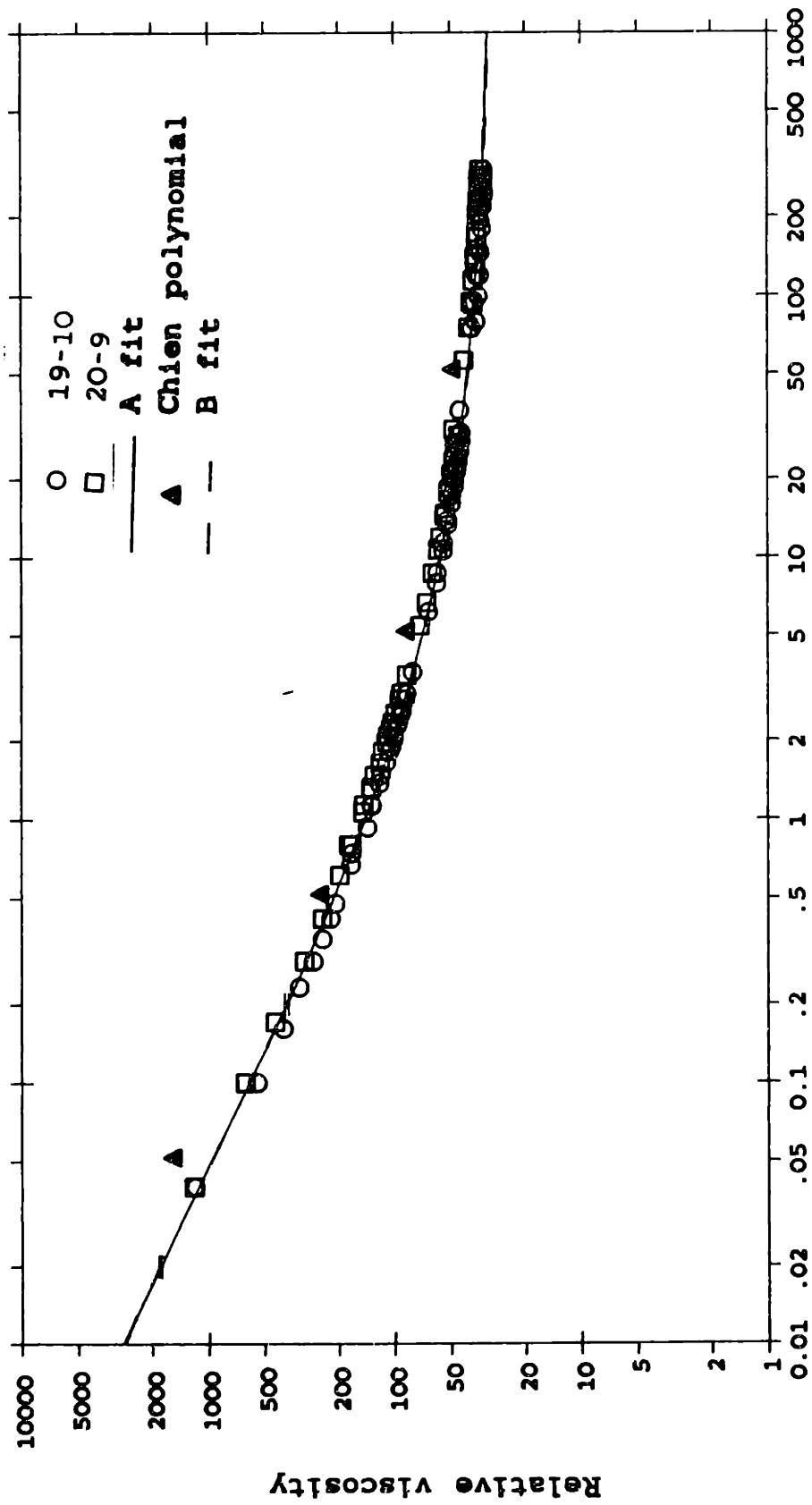


Shear rate (1/sec)

Figure 4-9: Eagle's solution data, H=92.



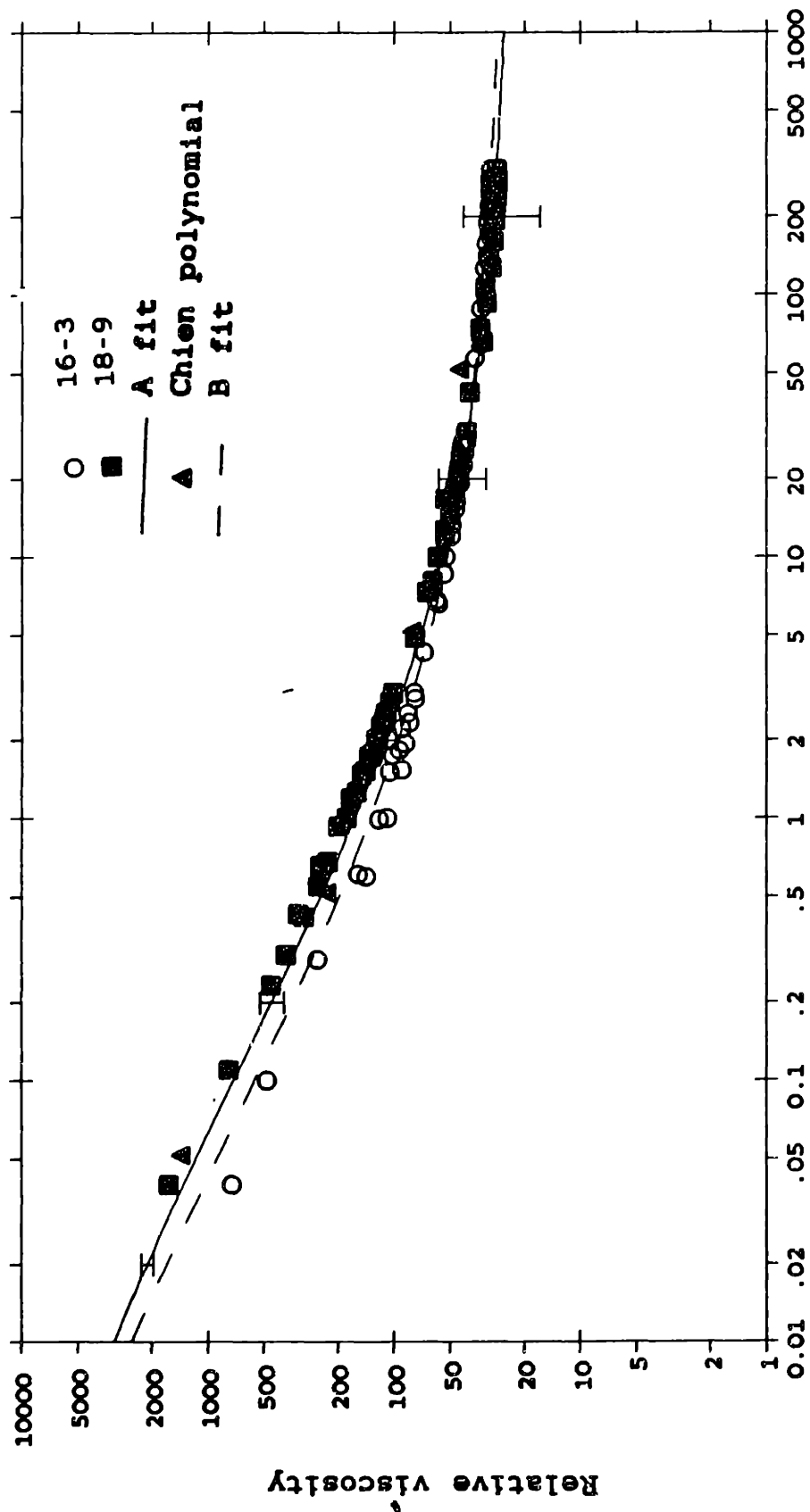
Relative viscosity vs. Shear rate  
RBC in Eagle's solution, H=91



Shear rate (1/sec)

Figure 4-10: Eagle's solution data, H=91.

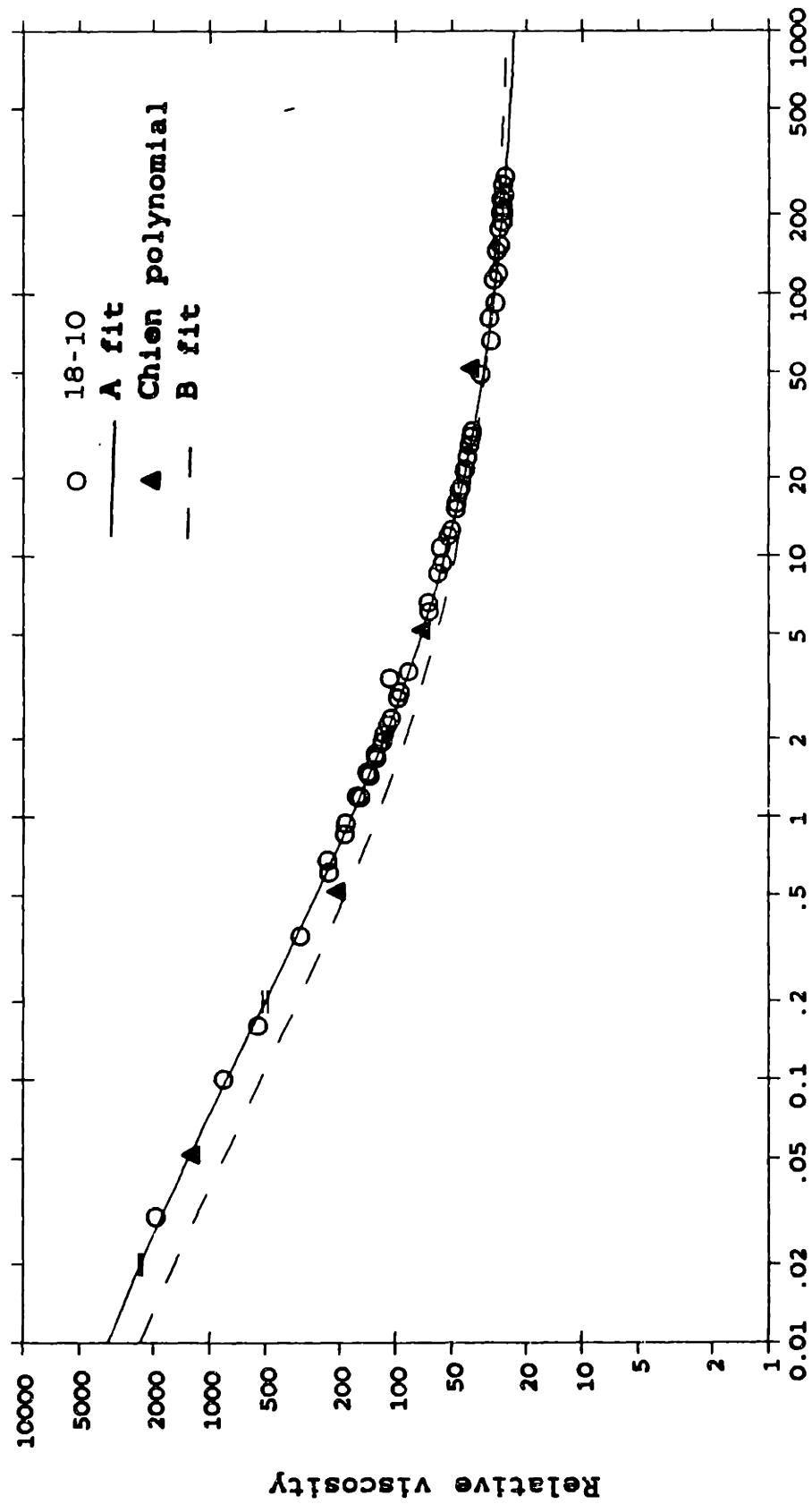
Relative viscosity vs. Shear rate  
RBC in Eagle's solution,  $H=90$



Shear rate (1/sec)

Figure 4-11: Eagle's solution data,  $H=90$ .

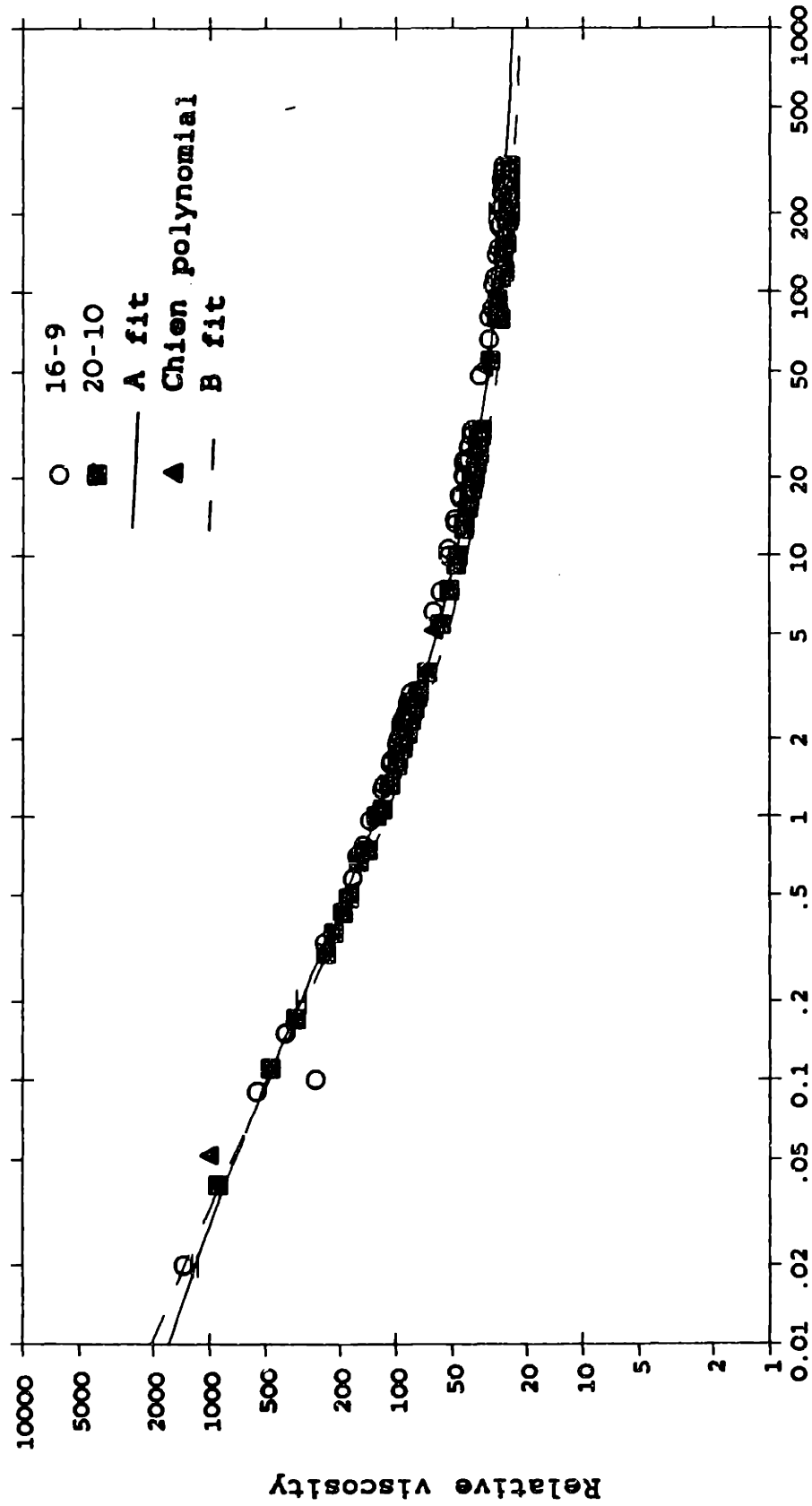
Relative viscosity vs. Shear rate  
RBC in Eagle's solution, H=89



Shear rate (1/sec)

Figure 4-12: Eagle's solution data, H=89.

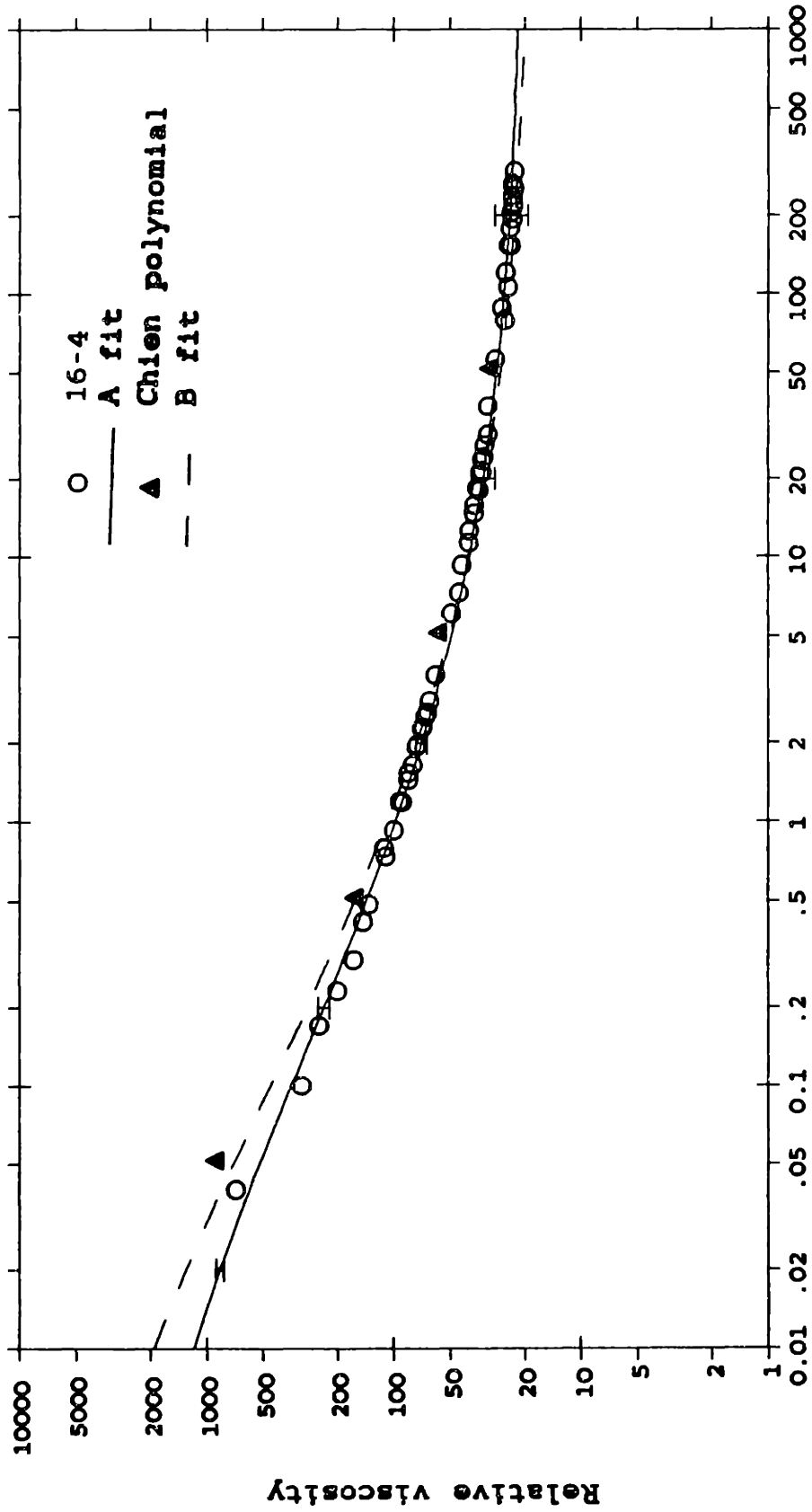
Relative viscosity vs. Shear rate  
RBC in Eagle's solution, H=87



Shear rate (1/sec)

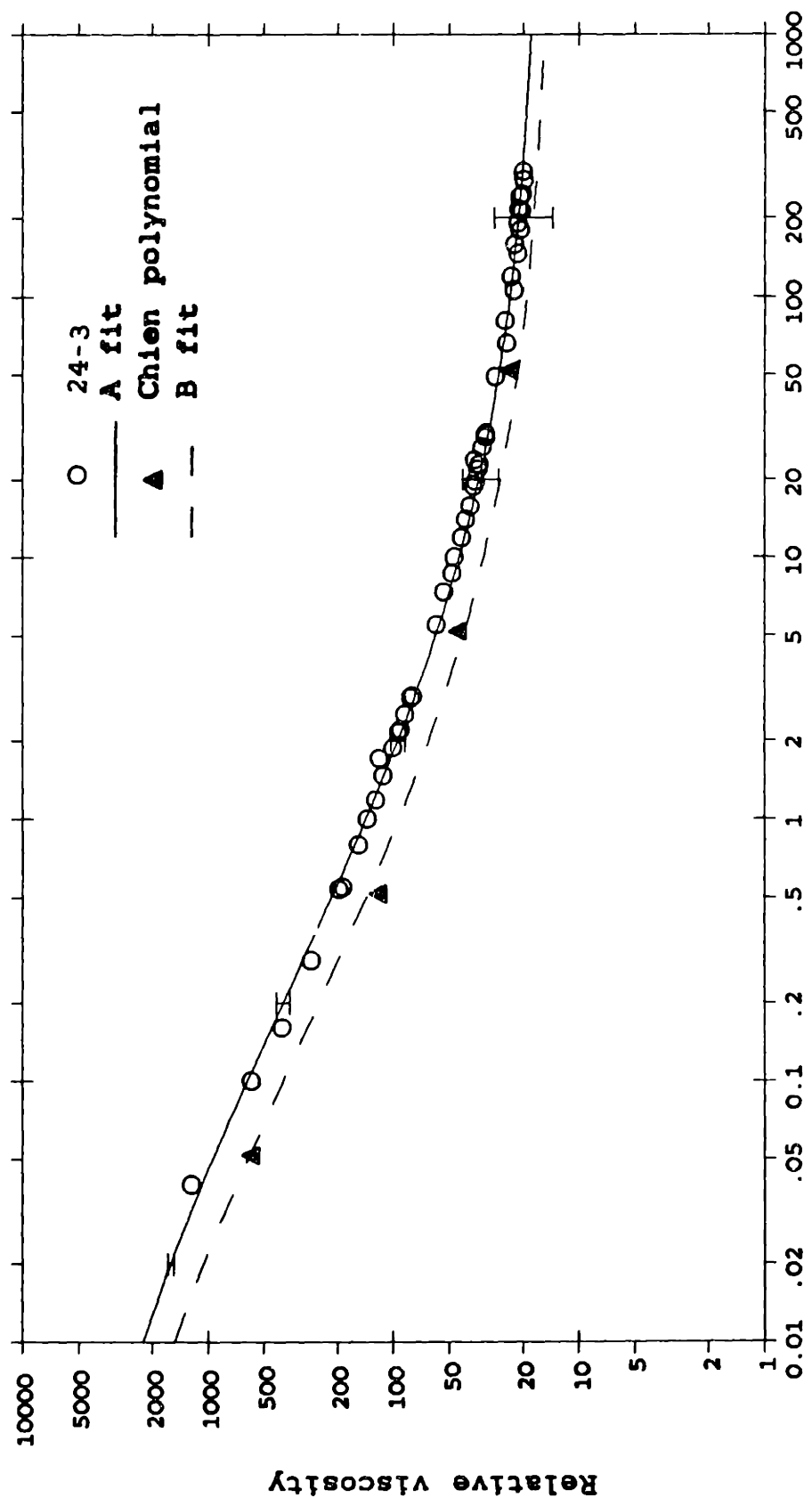
Figure 4-13: Eagle's solution data, H=87.

Relative viscosity vs. Shear rate  
RBC in Eagle's solution, H=86



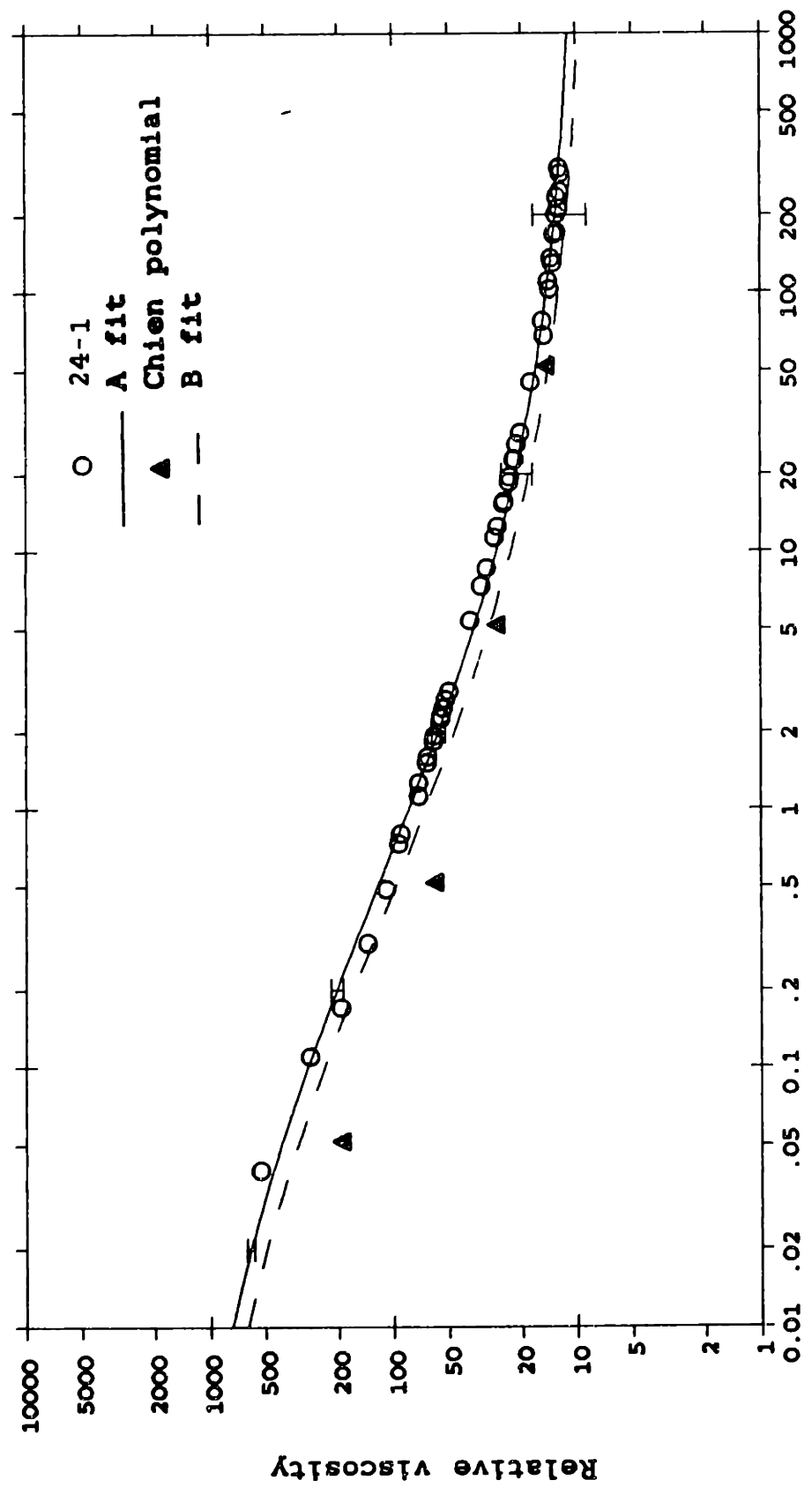
Shear rate (1/sec)  
Figure 4-14: Eagle's solution data, H=86.

Relative viscosity vs. Shear rate  
RBC in Eagle's solution, H=82



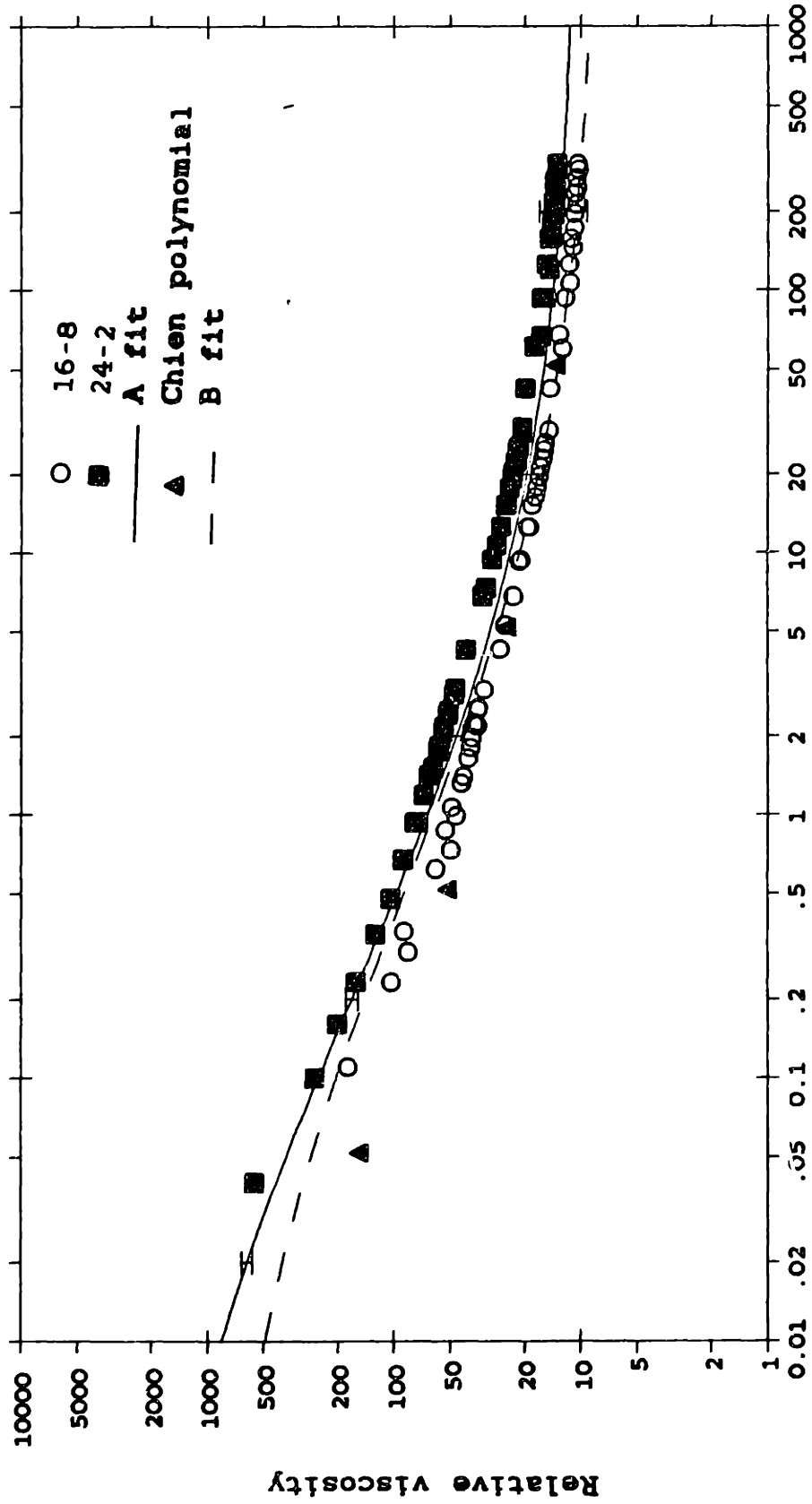
Shear rate (1/sec)  
Figure 4-15: Eagle's solution data, H=82.

Relative viscosity vs. Shear rate  
RBC in Eagle's solution, H=72



Shear rate (1/sec)  
Figure 4-16: Eagle's solution data, H=72.

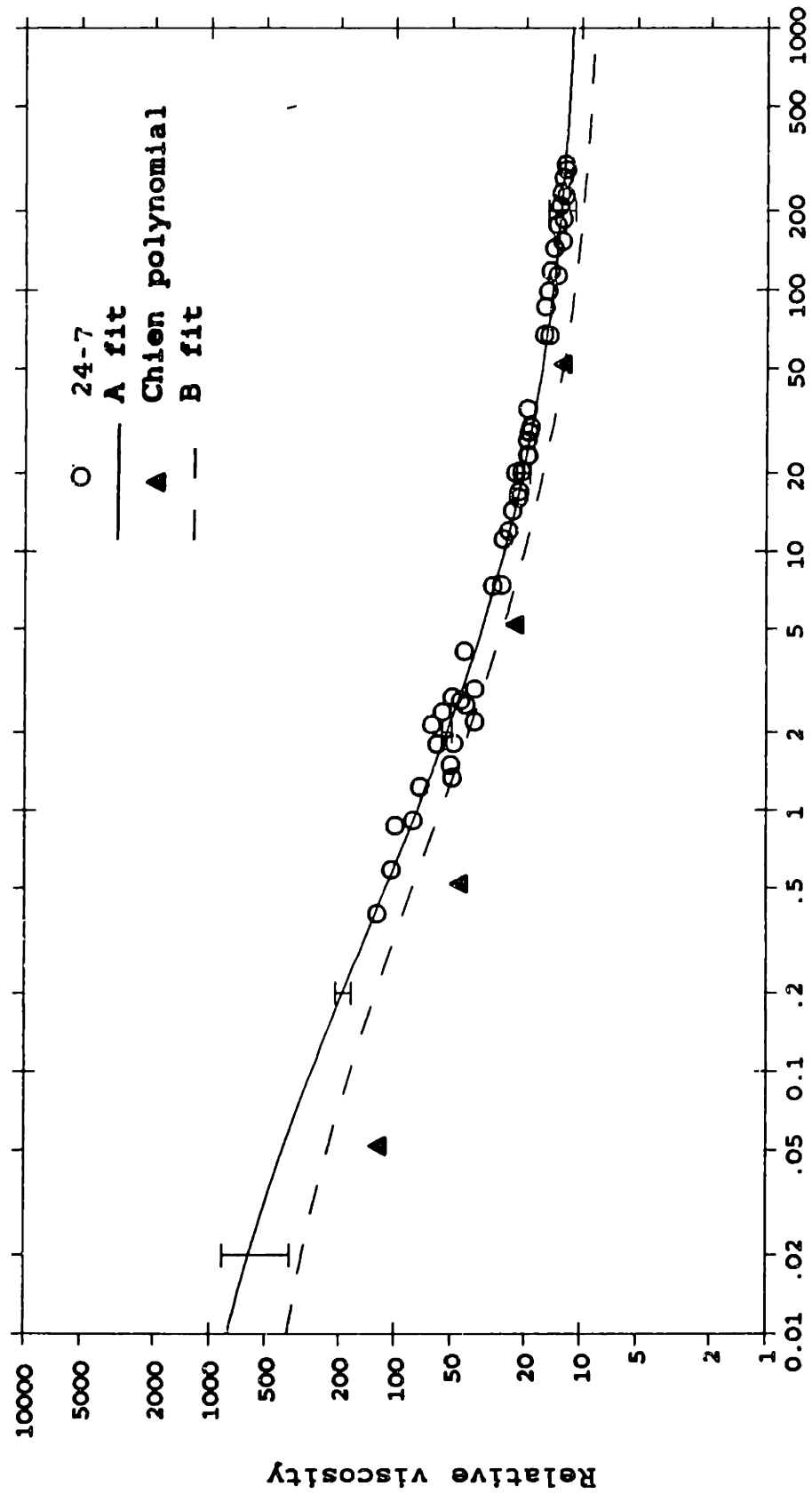
Relative viscosity vs. Shear rate  
RBC in Eagle's solution, H=70



Shear rate (1/sec)  
Figure 4-17: Eagle's solution data,  $\eta_0=70$ .

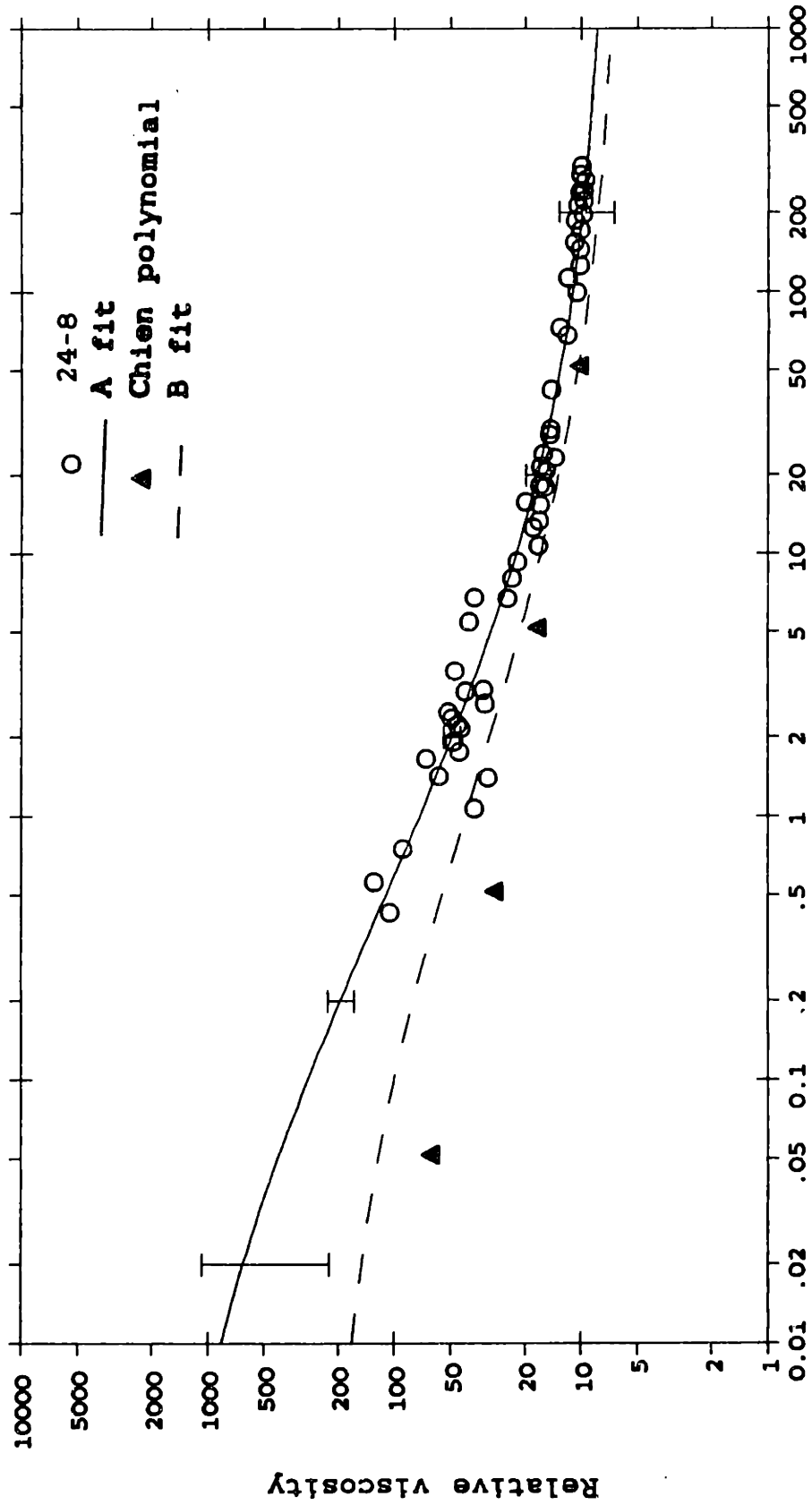


Relative viscosity vs. Shear rate  
RBC in Eagle's solution, H=68



Shear rate (1/sec)  
Figure 4-18: Eagle's solution data, H=68.

Relative viscosity vs. Shear rate  
RBC in Eagle's solution, H=62



Shear rate (1/sec)  
**Figure 4-19:** Eagle's solution data, H=62.

Relative viscosity vs. Shear rate  
RBC in Eagle's solution, H=59

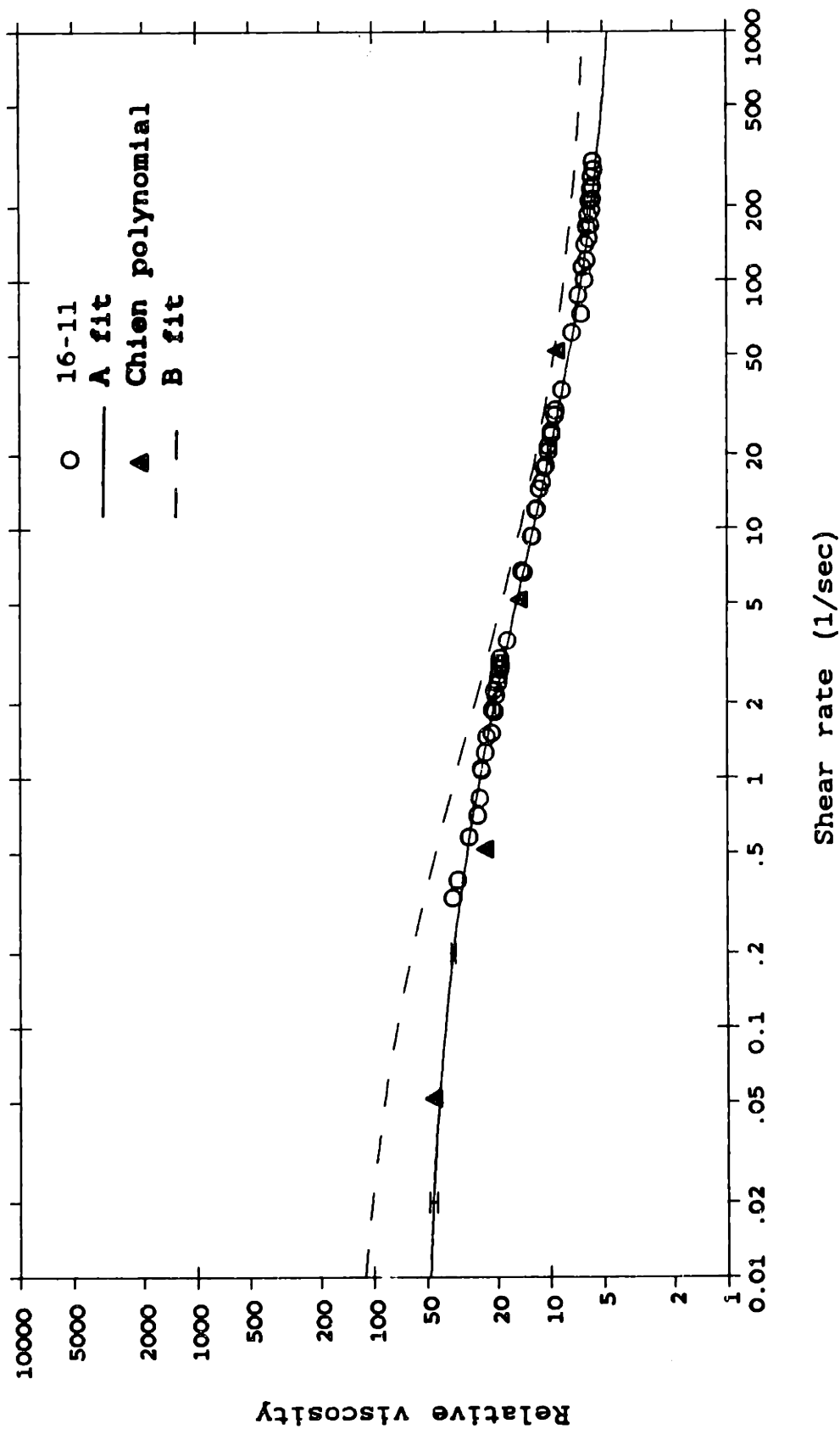
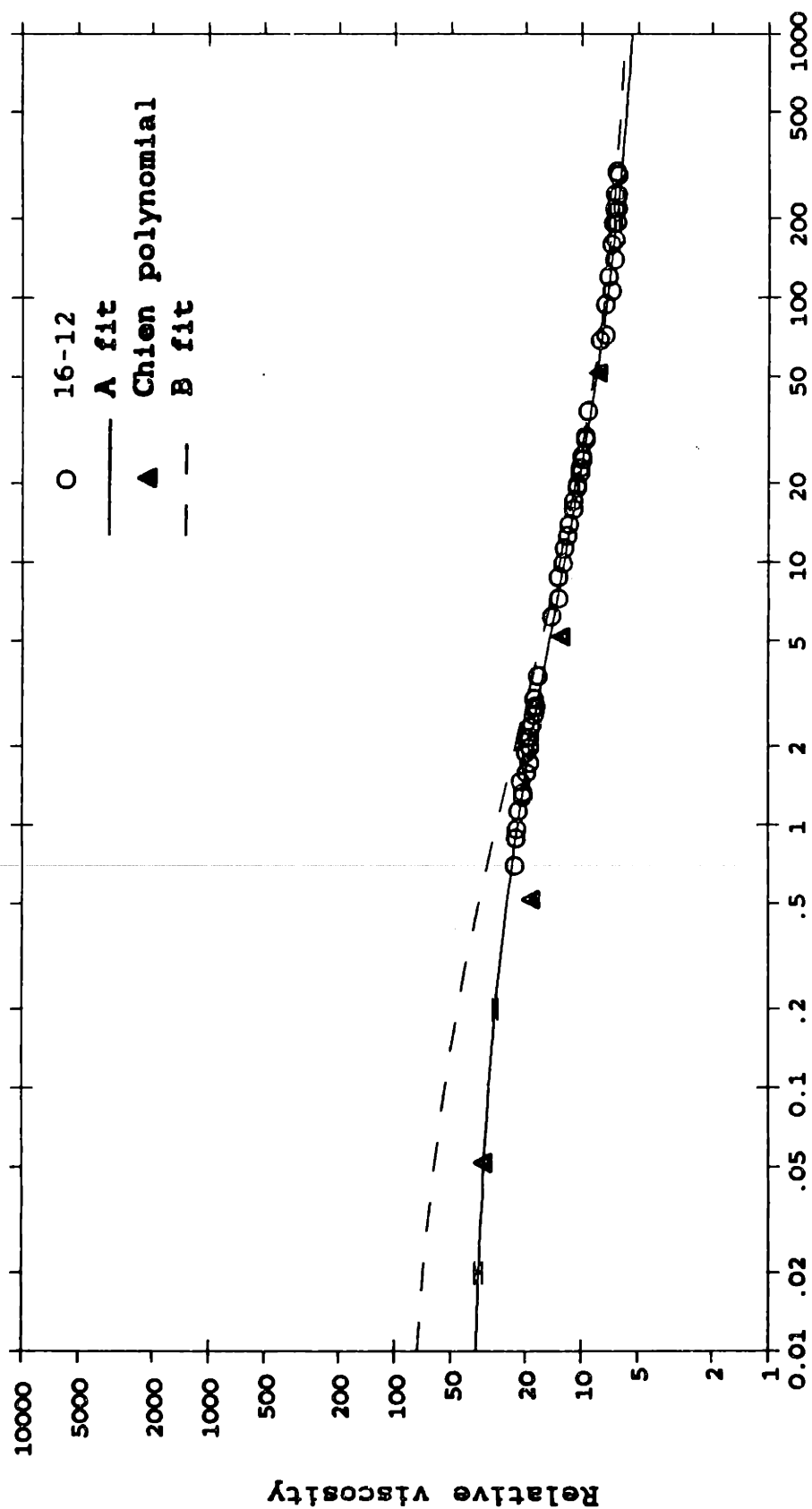


Figure 4-20: Eagle's solution data, H=59.

Relative viscosity vs. Shear rate  
RBC in Eagle's solution, H=56



Shear rate (1/sec)  
Figure 4-21: Eagle's solution data, H=56.

Relative viscosity vs. Shear rate  
RBC in Eagle's solution, H=48

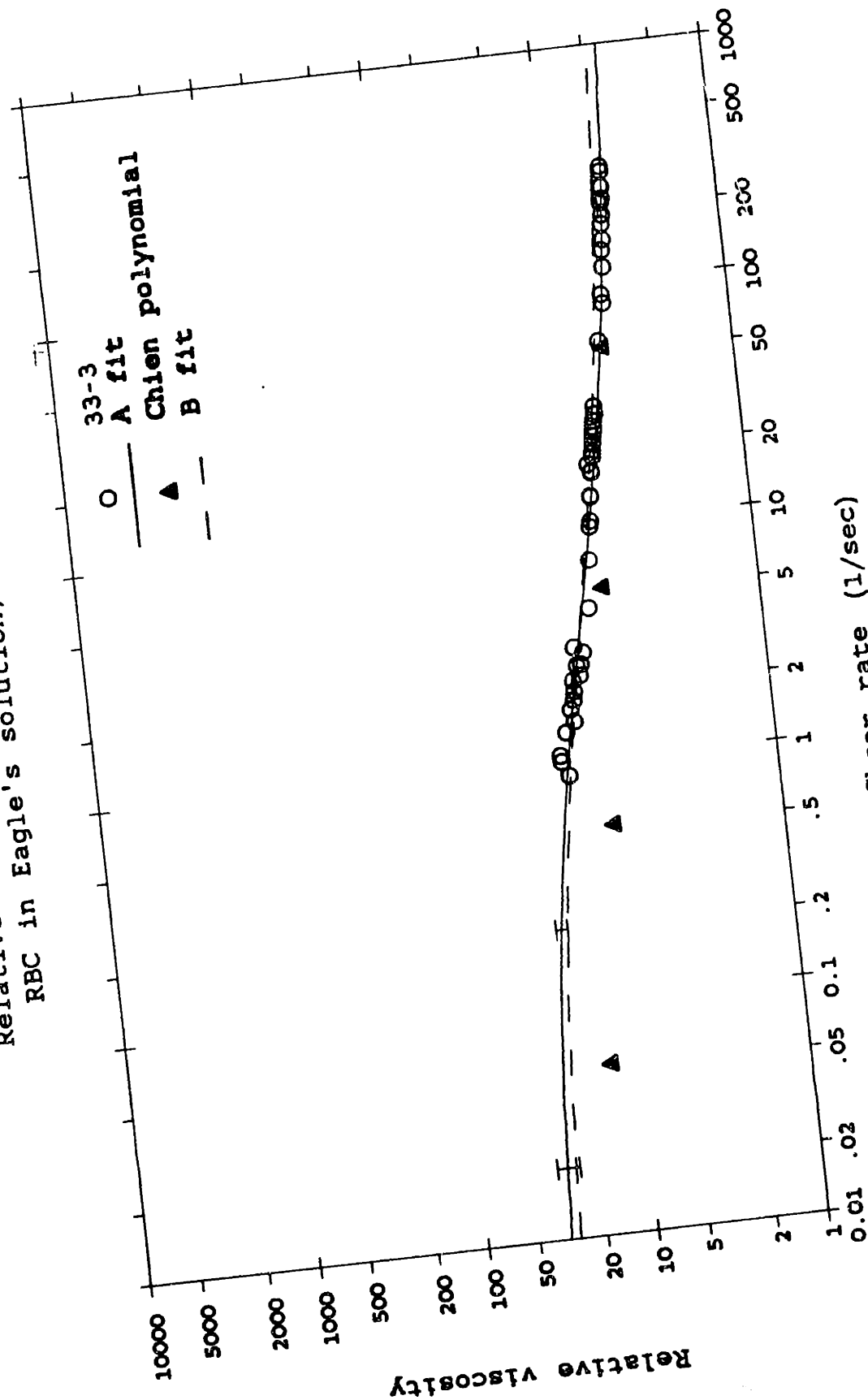
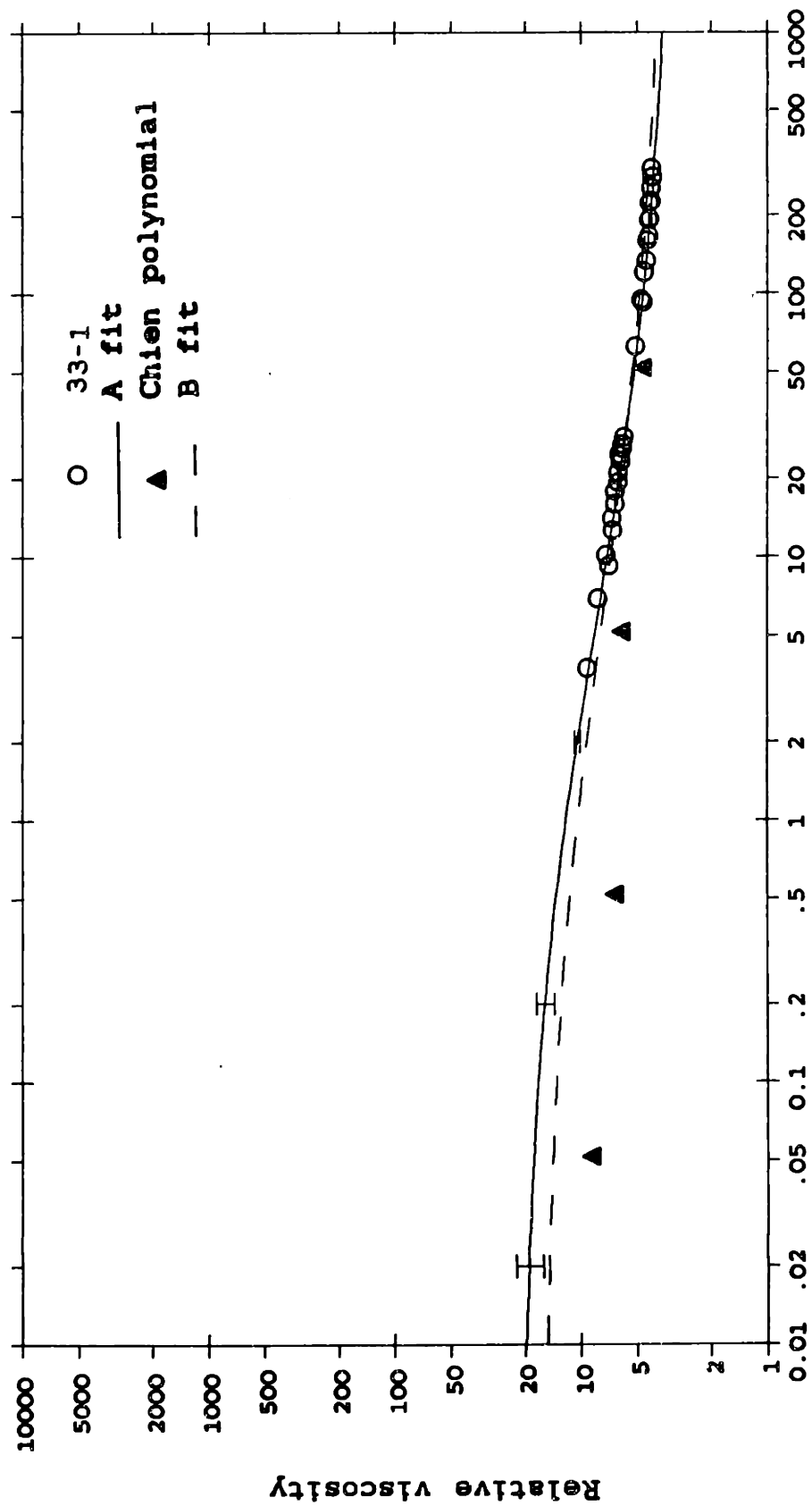


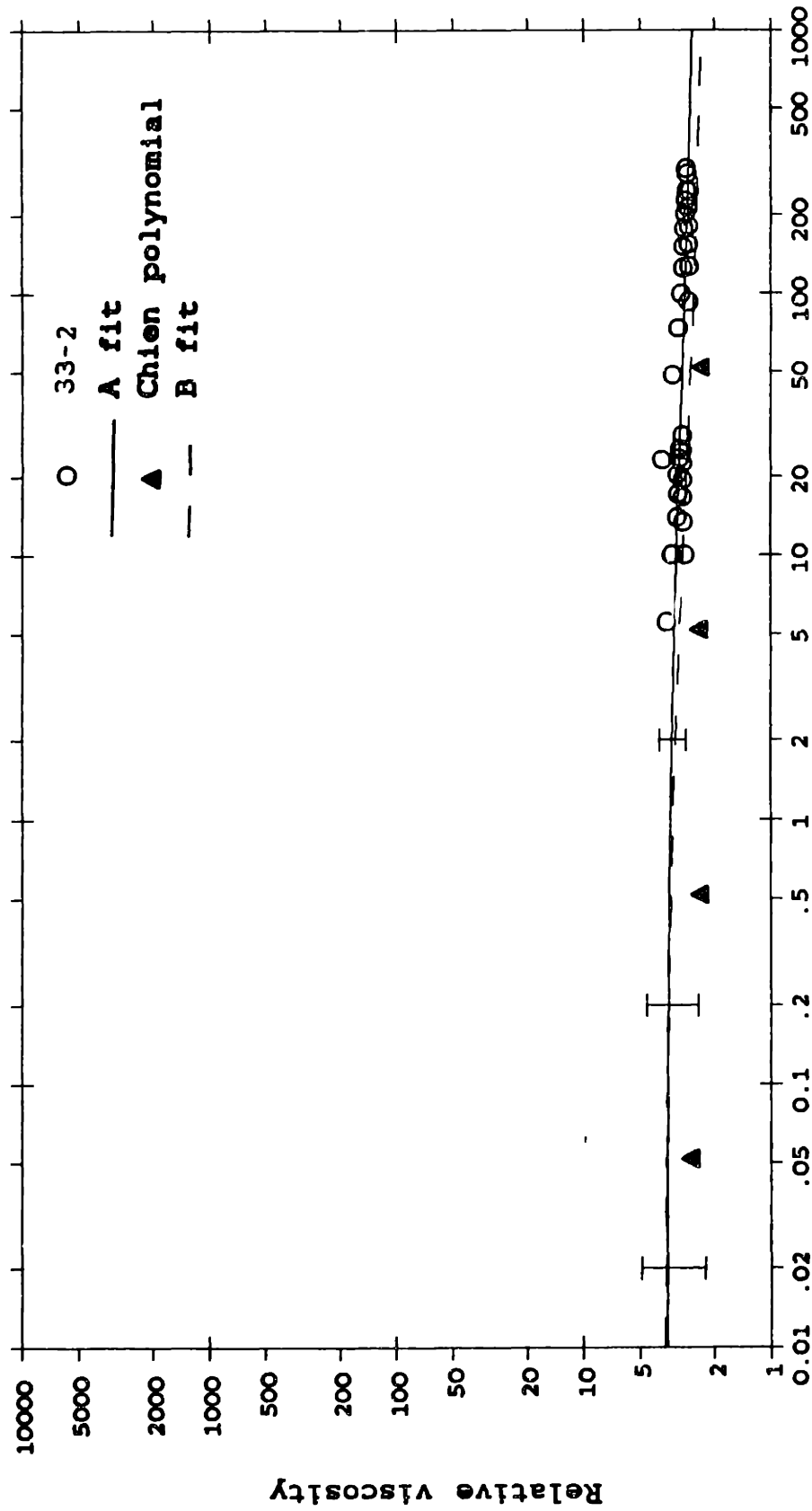
Figure 4-22: Eagle's solution data, H=48.

Relative viscosity vs. Shear rate  
RBC in Eagle's solution, H=42



Shear rate (1/sec)  
Figure 4-23: Eagle's solution data, H=42.

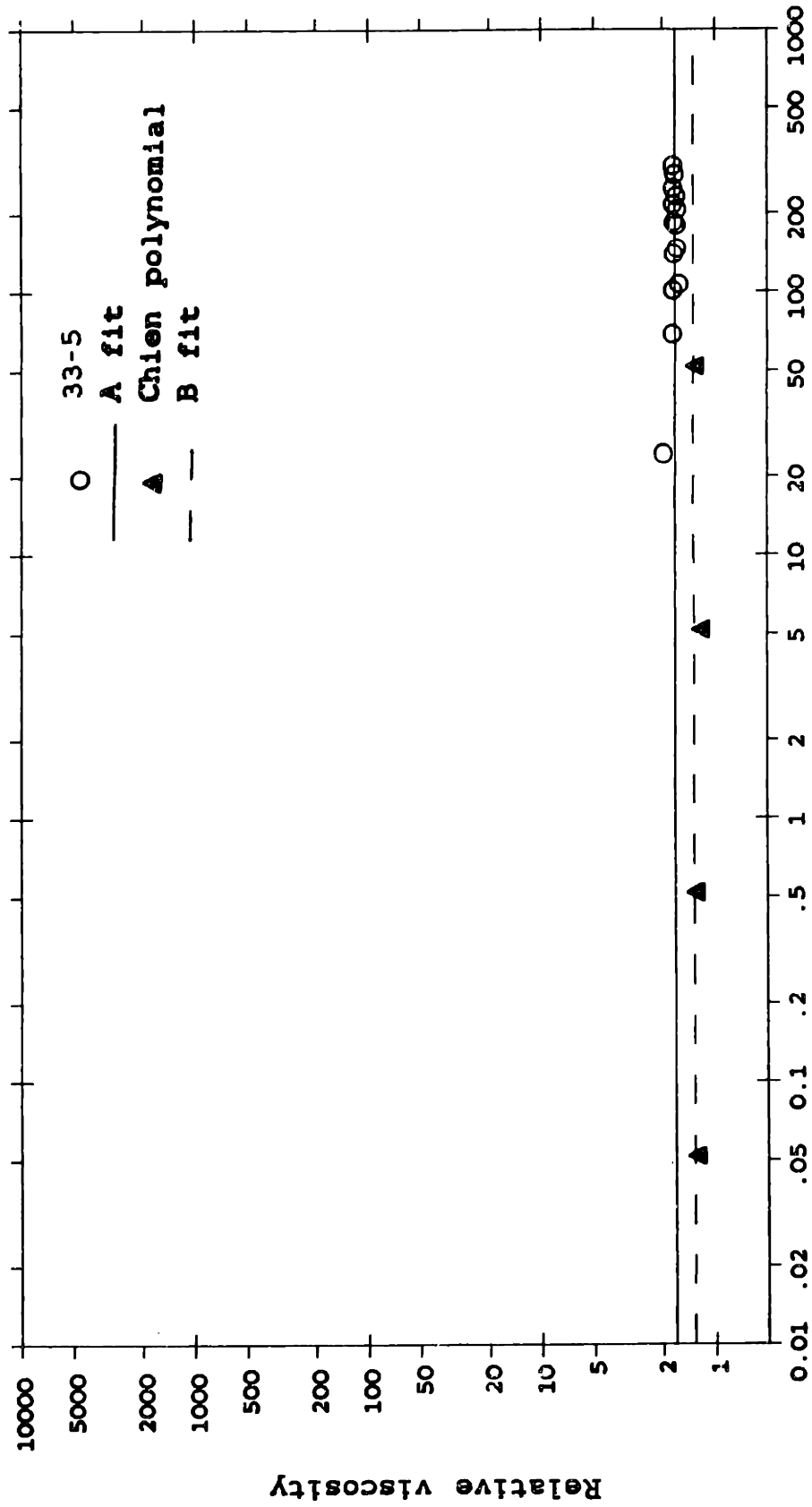
Relative viscosity vs. Shear rate  
RBC in Eagle's solution, H=25



Shear rate (1/sec)

Figure 4-24: Eagle's solution data, H=25.

Relative viscosity vs. Shear rate  
RBC in Eagle's solution, H=7



Shear rate (1/sec)  
Figure 4-25: Eagle's solution data, H=7.



Relative viscosity vs. Shear rate  
RBC in plasma, H=98.1

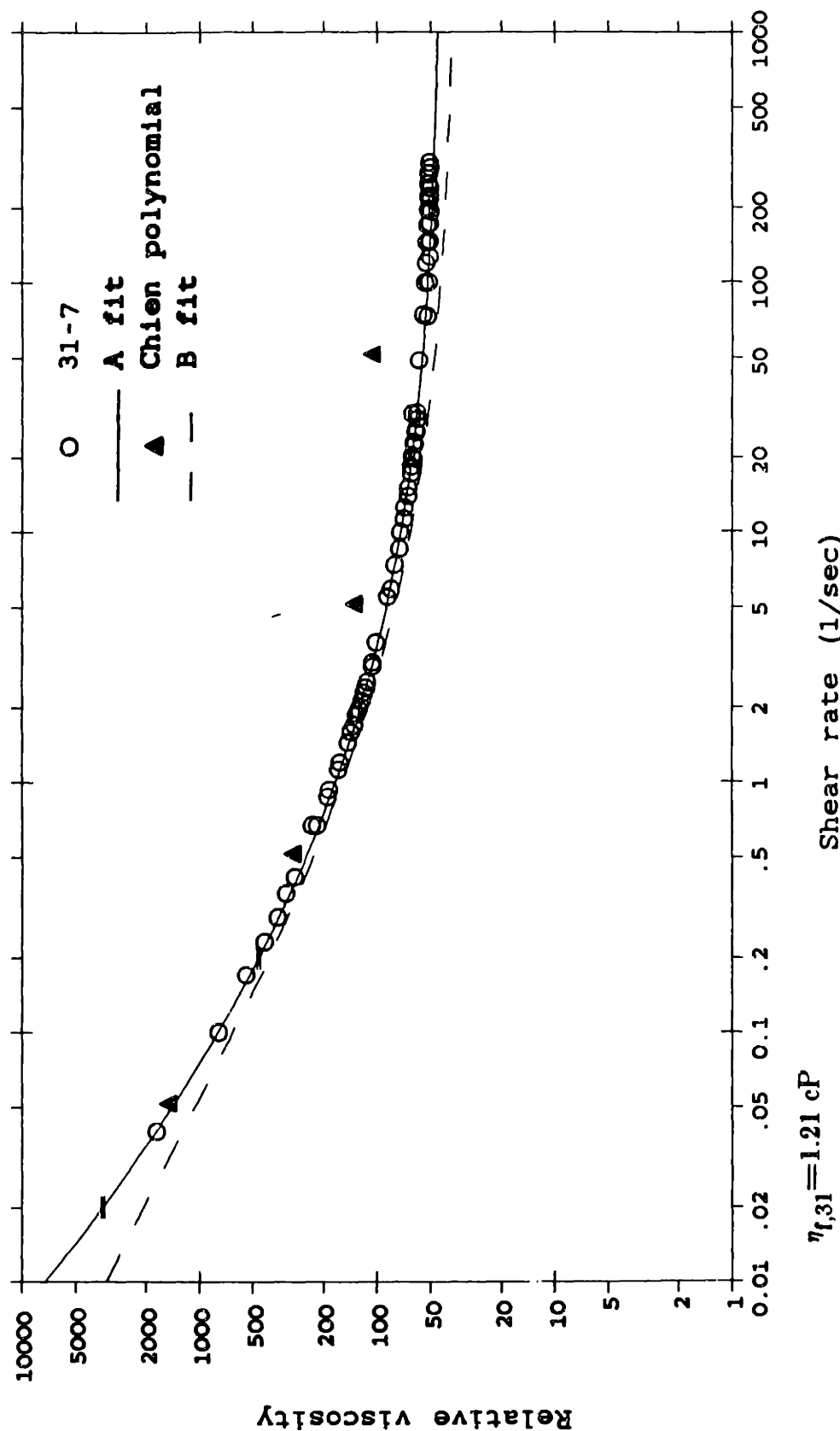


Figure 4-26: Plasma data, H=98.1.

Relative viscosity vs. Shear rate  
RBC in plasma, H=97.6

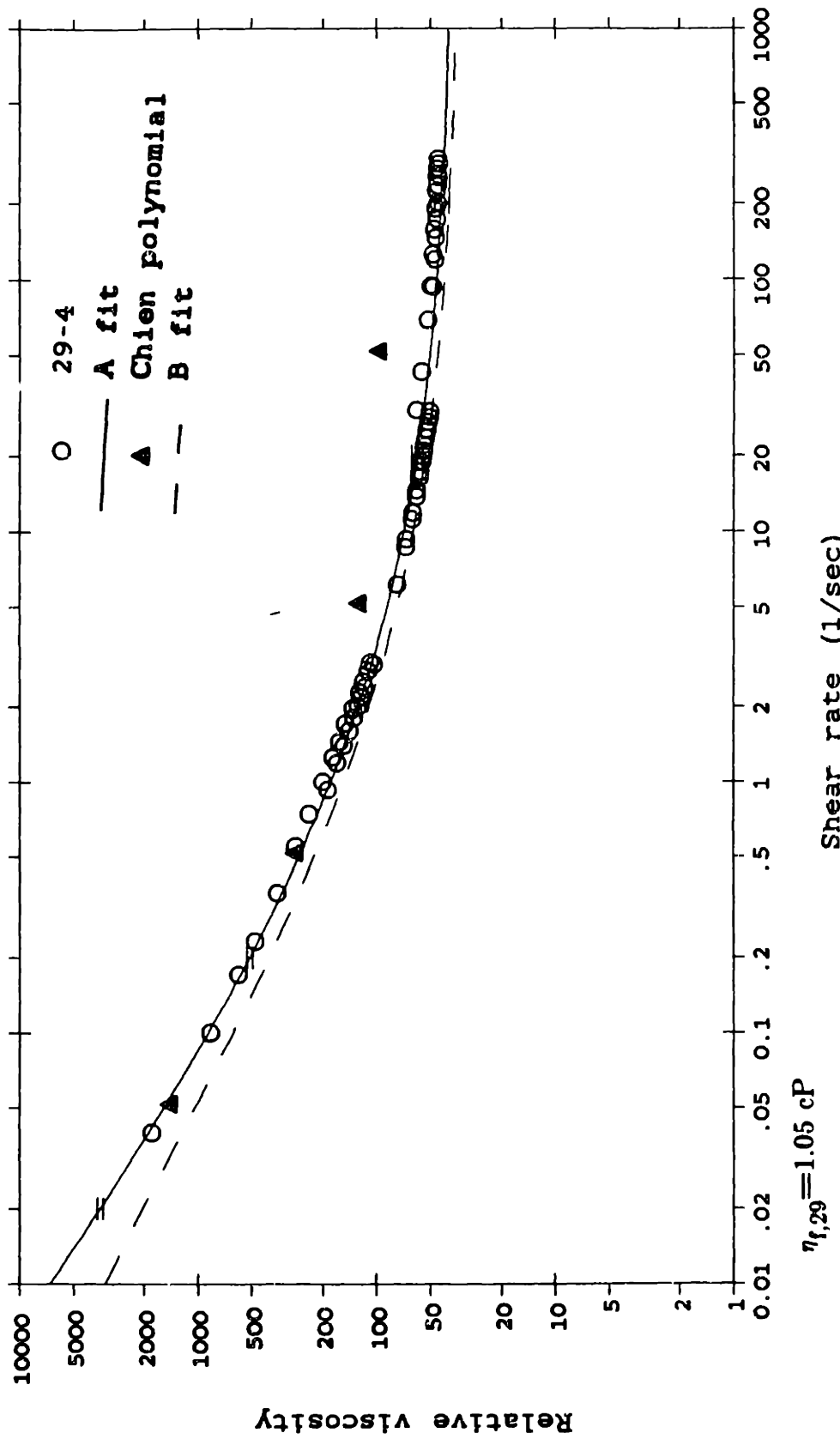
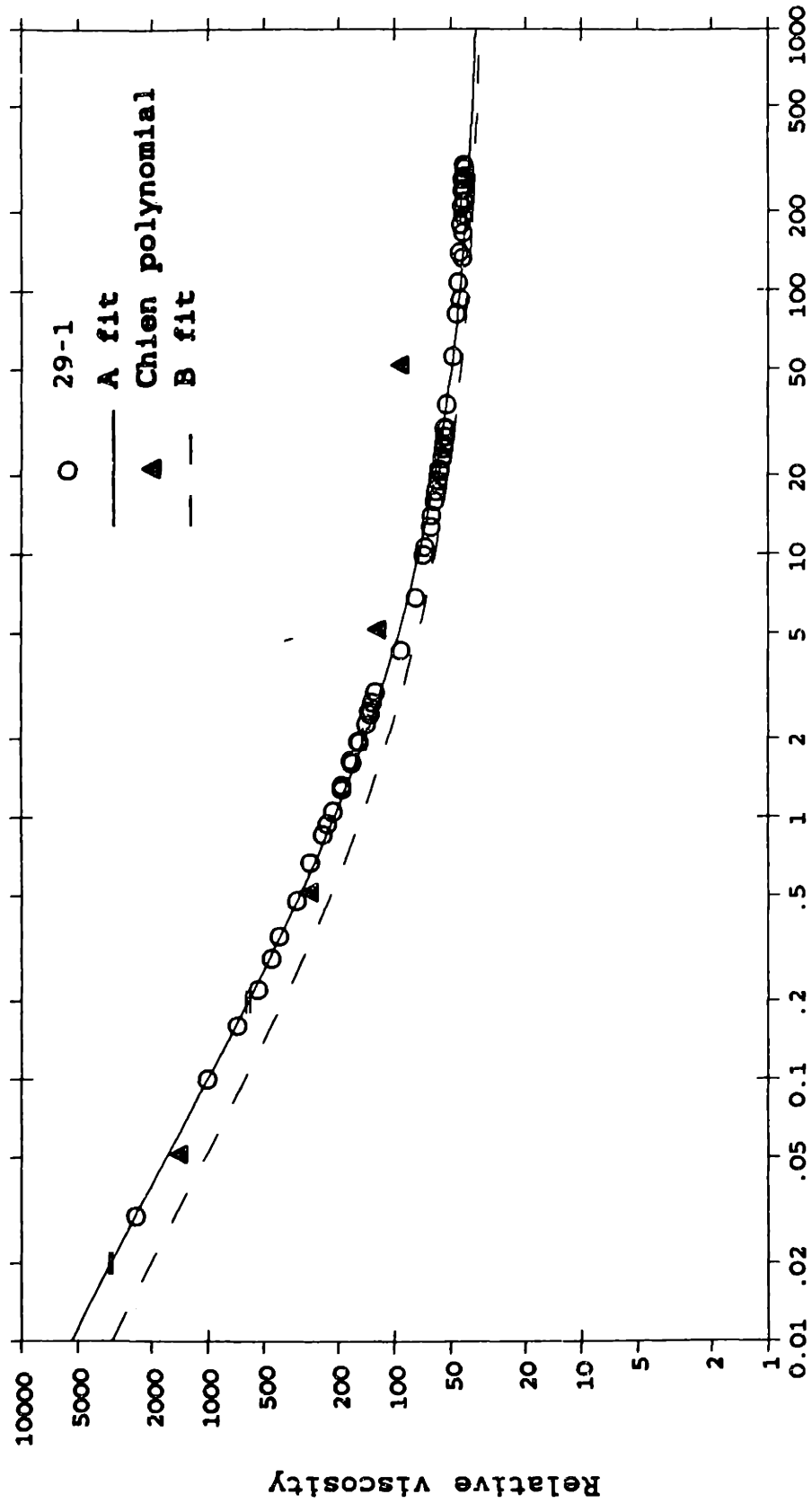


Figure 4-27: Plasma data, H=97.6.

Relative viscosity vs. Shear rate  
RBC in plasma, H=97.3



$\eta_{r,29} = 1.05 \text{ cP}$

Shear rate (1/sec)

Figure 4-28: Plasma data, H=97.3.

Relative viscosity vs. Shear rate  
RBC in plasma, H=97

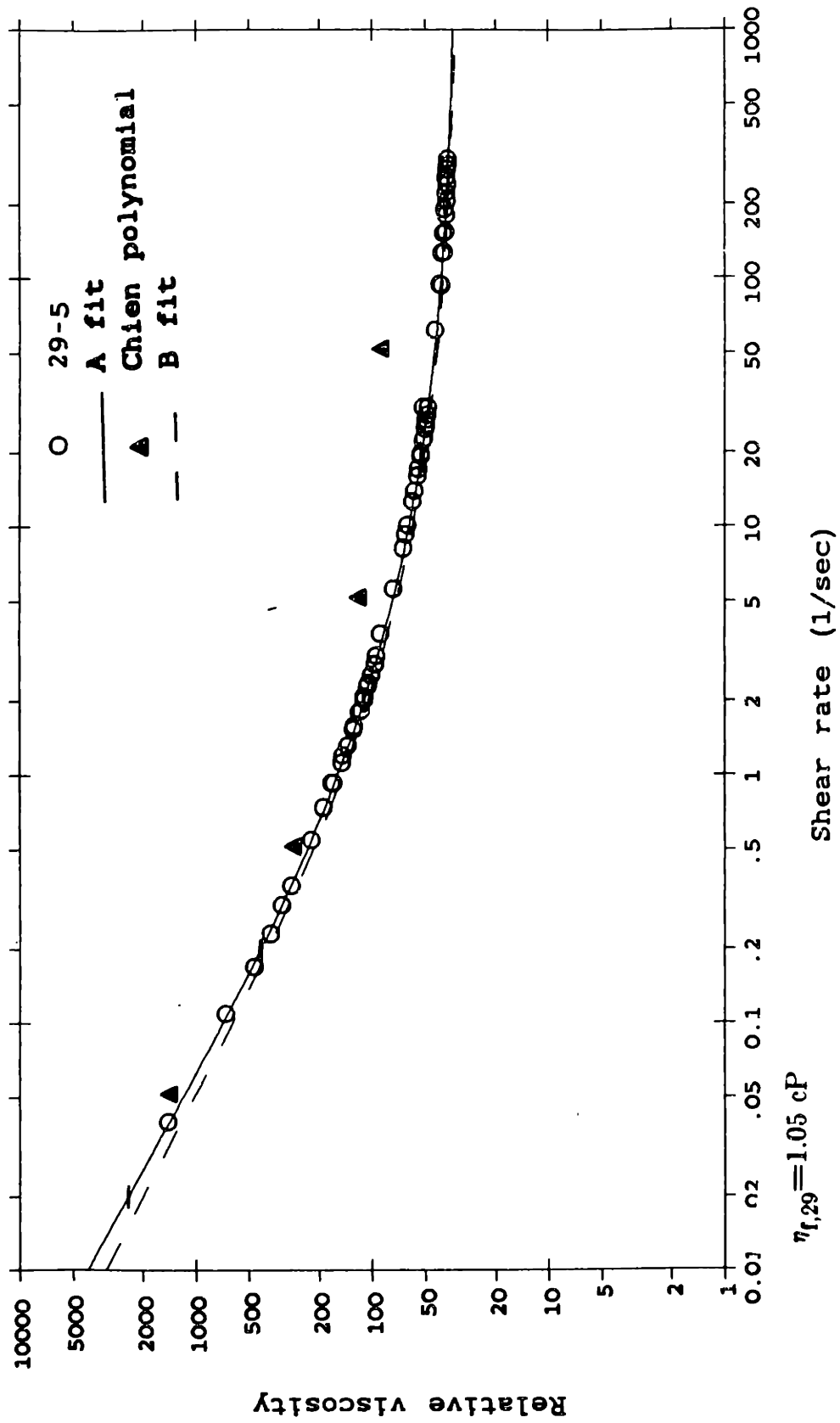
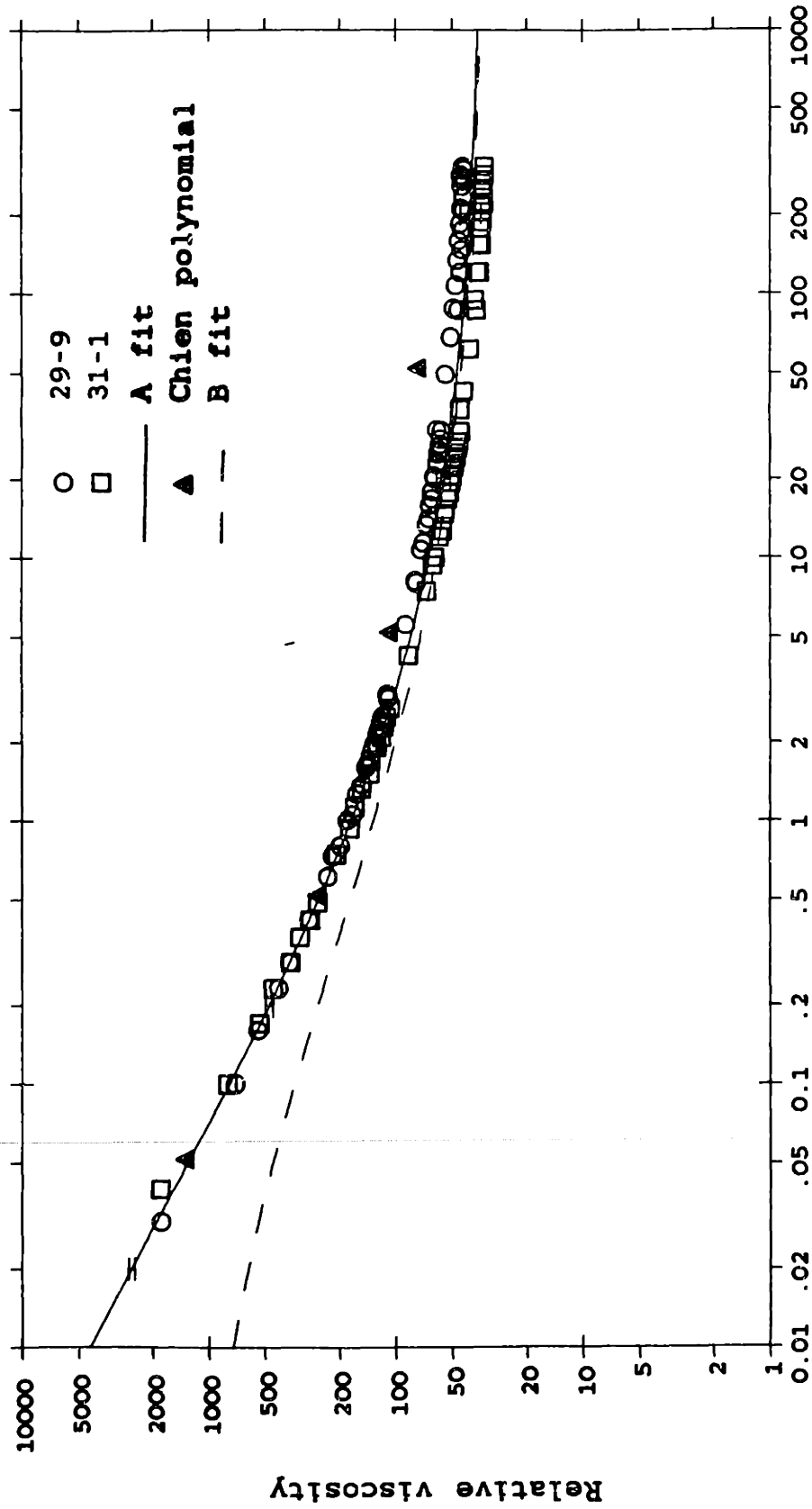


Figure 4-29: Plasma data, H=97.

Relative viscosity vs. Shear rate  
RBC in plasma, H=96



$\eta_{t,29}=1.05 \text{ cP}$   $\eta_{t,31}=1.21 \text{ cP}$

Shear rate (1/sec)

Figure 4-30: Plasma data, H=96.

Relative viscosity vs. Shear rate  
RBC in plasma, H=95

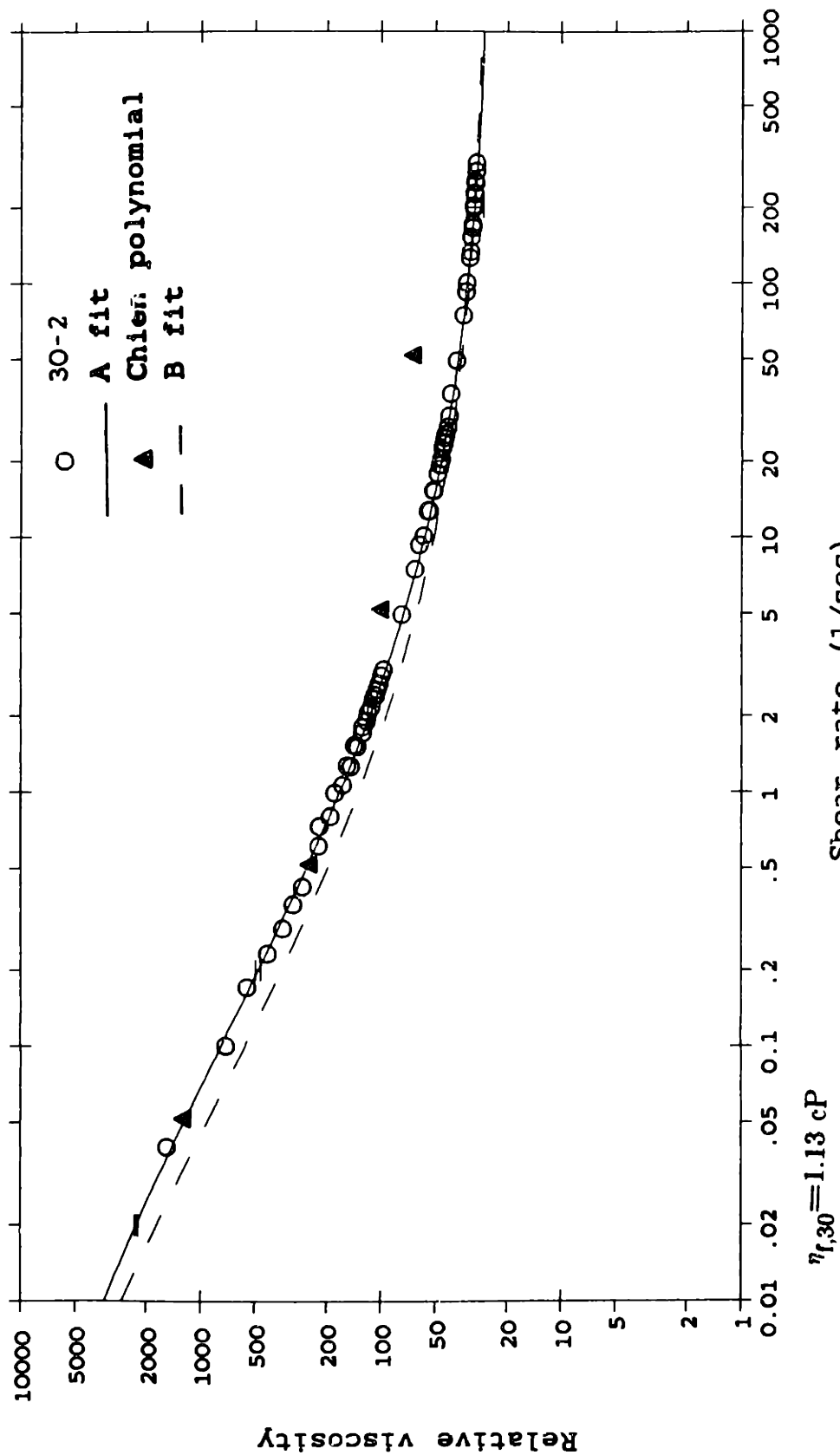


Figure 4-31: Plasma data, H=95.

Relative viscosity vs. Shear rate  
RBC in plasma, H=94

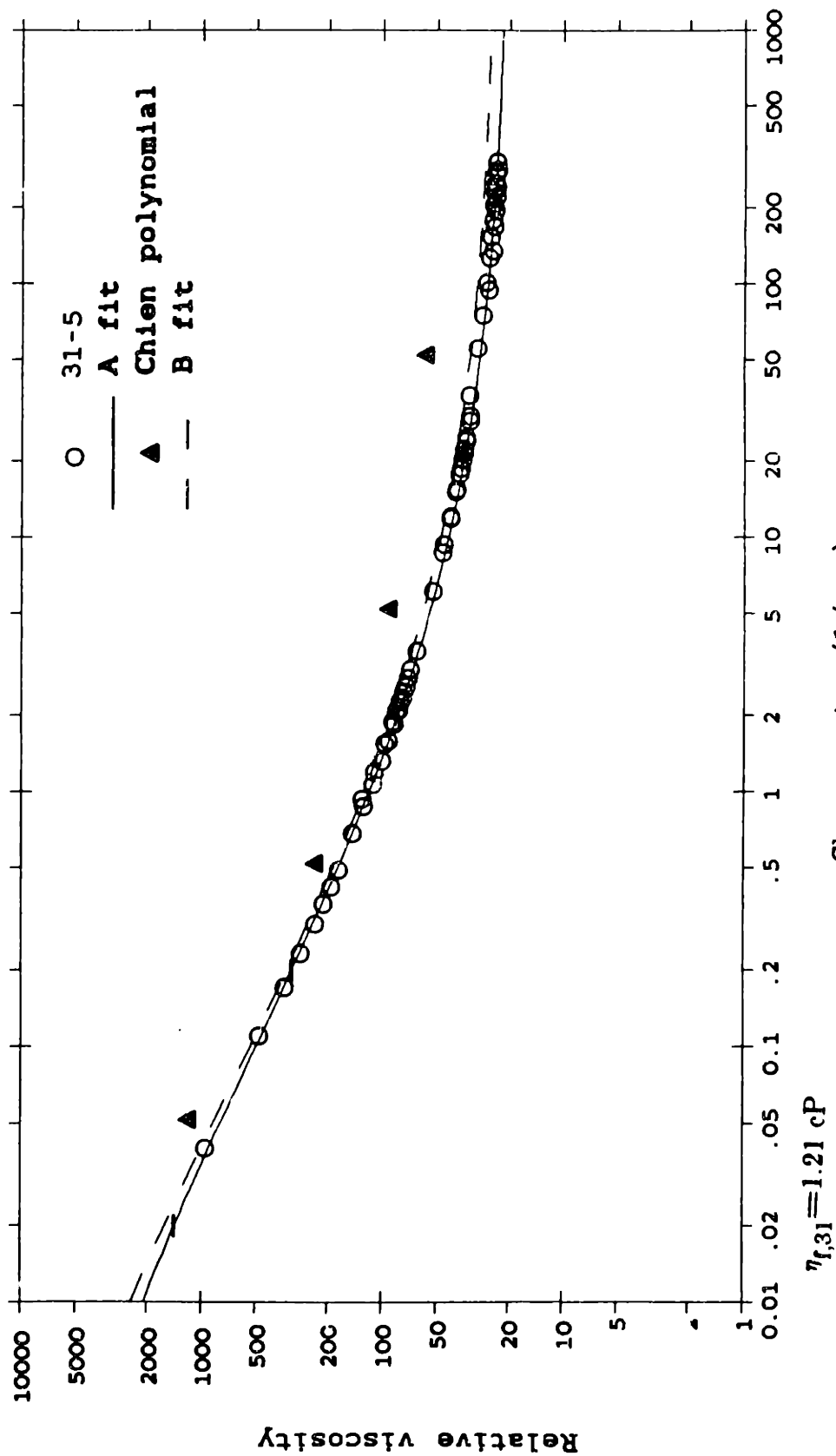
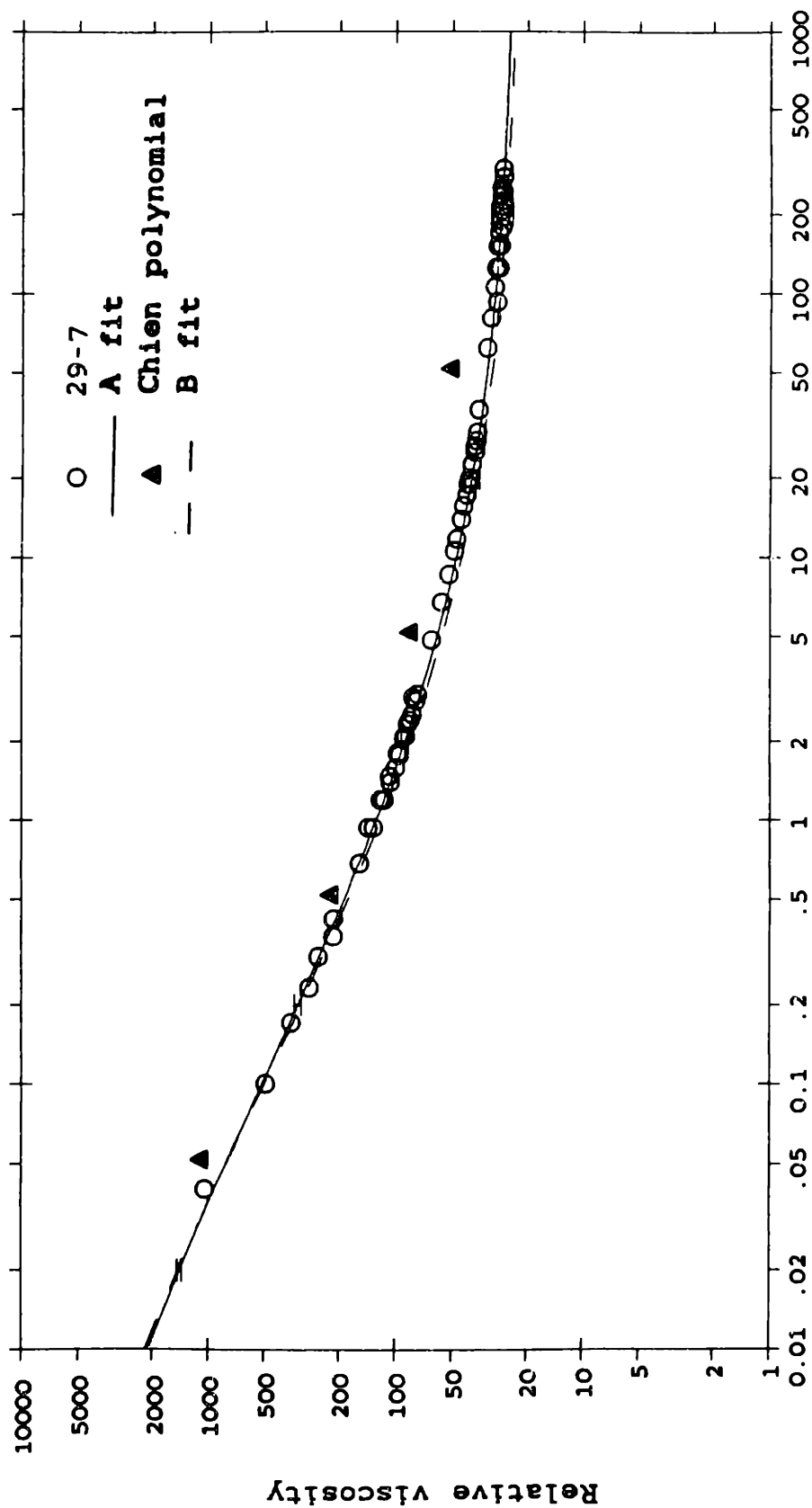


Figure 4-32: Plasma data, H=94.

Relative viscosity vs. Shear rate  
RBC in plasma, H=93



$\eta_{l,29} = 1.05 \text{ cP}$

Shear rate (1/sec)

Figure 4-33: Plasma data, H=93.



Relative viscosity vs. Shear rate  
RBC in plasma, H=92

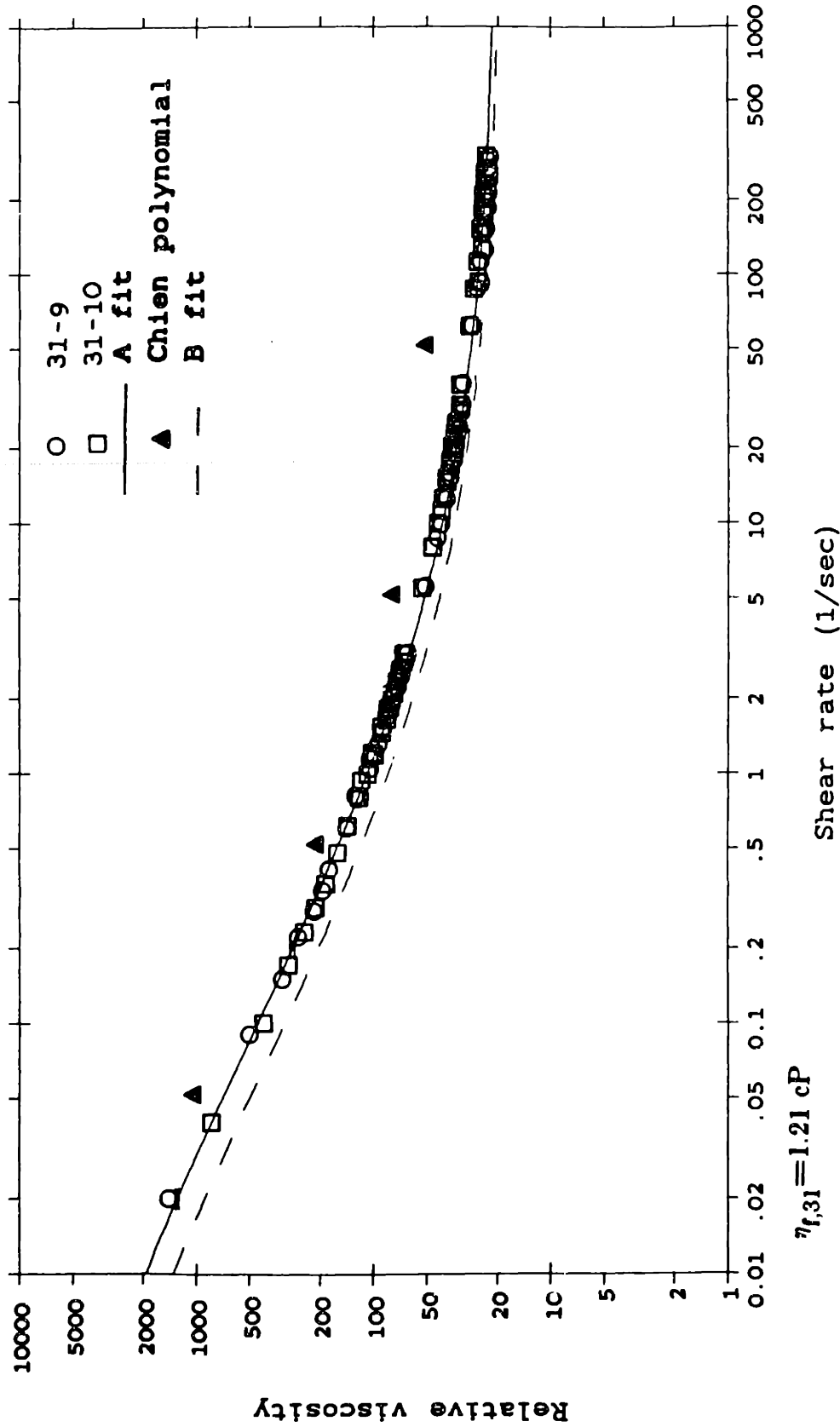


Figure 4-34: Plasma data, H=92.

Relative viscosity vs. Shear rate  
RBC in plasma, H=91

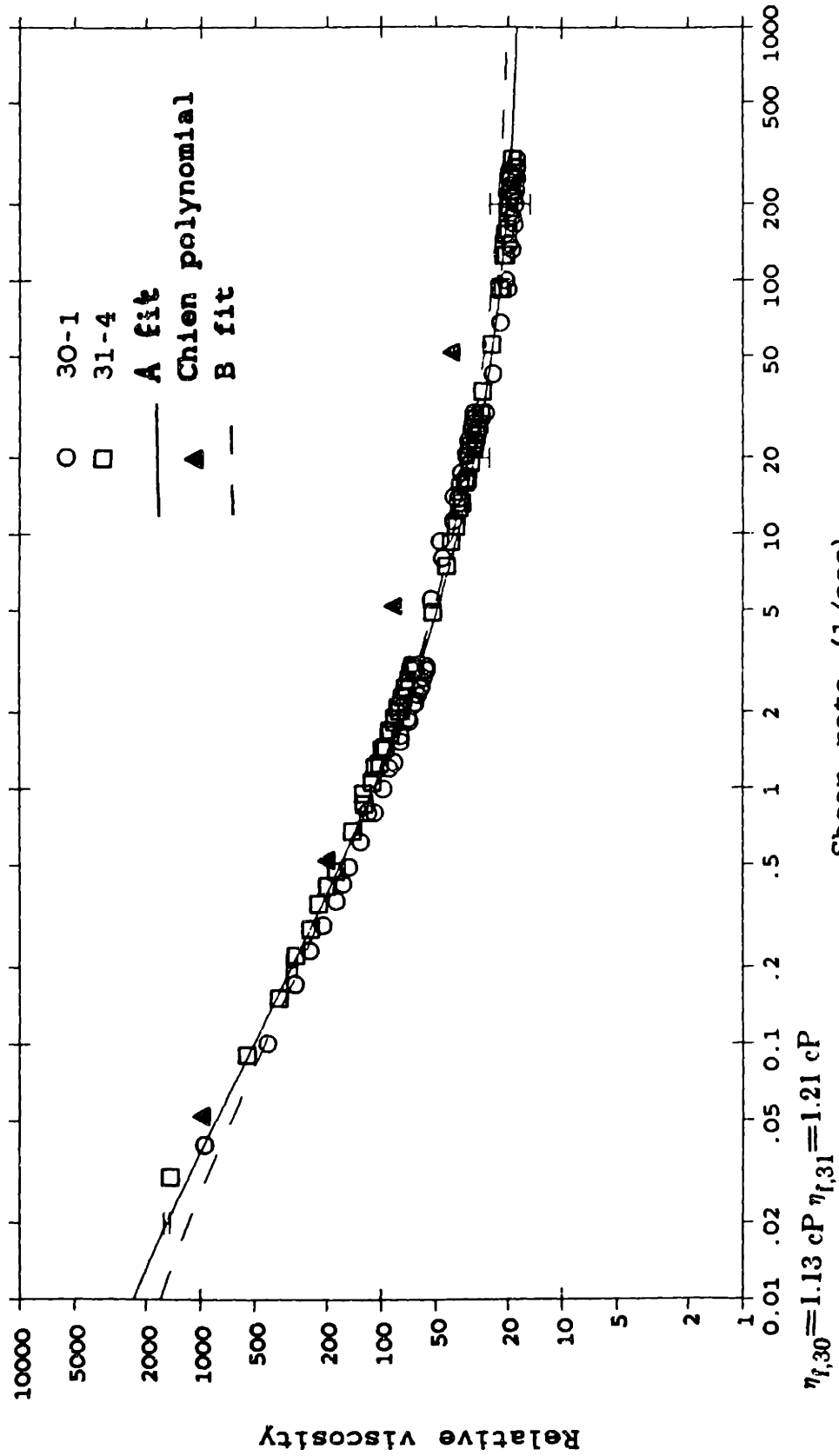
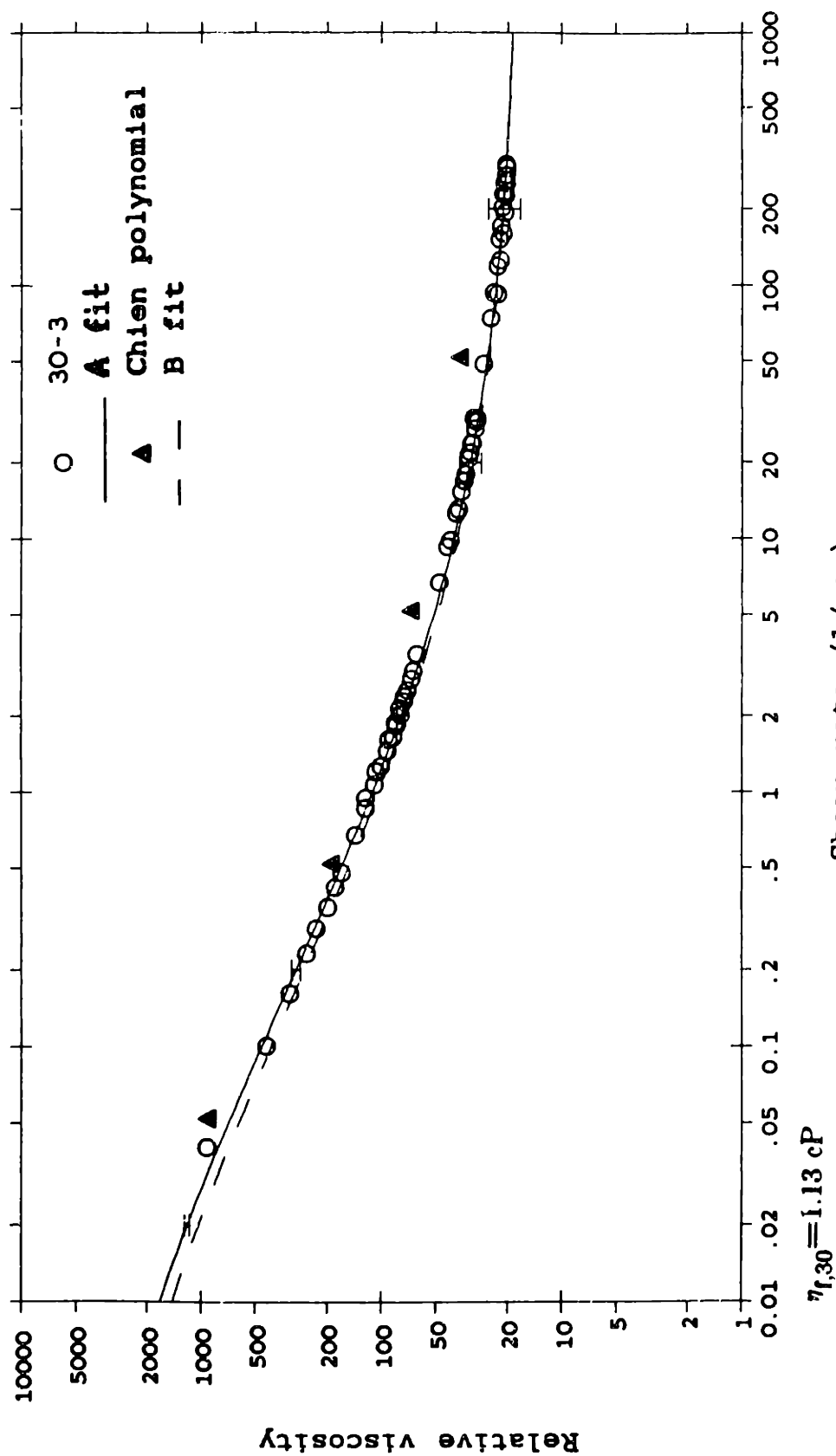


Figure 4-35: Plasma data, H=91.

Relative viscosity vs. Shear rate  
RBC in plasma, H=90



Shear rate (1/sec)

Figure 4-36: Plasma data, H=90.

Relative viscosity vs. Shear rate  
RBC in plasma, H=89

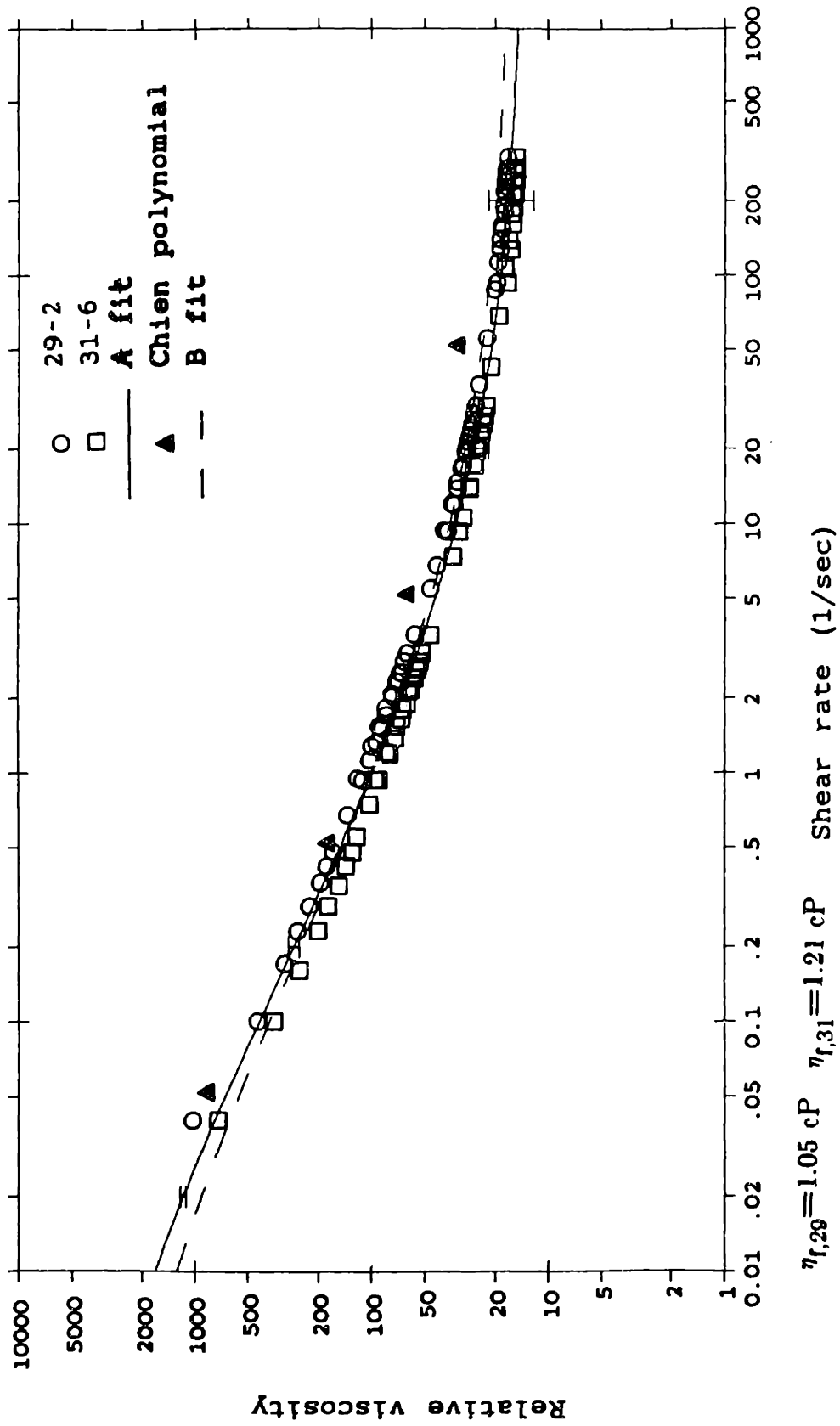
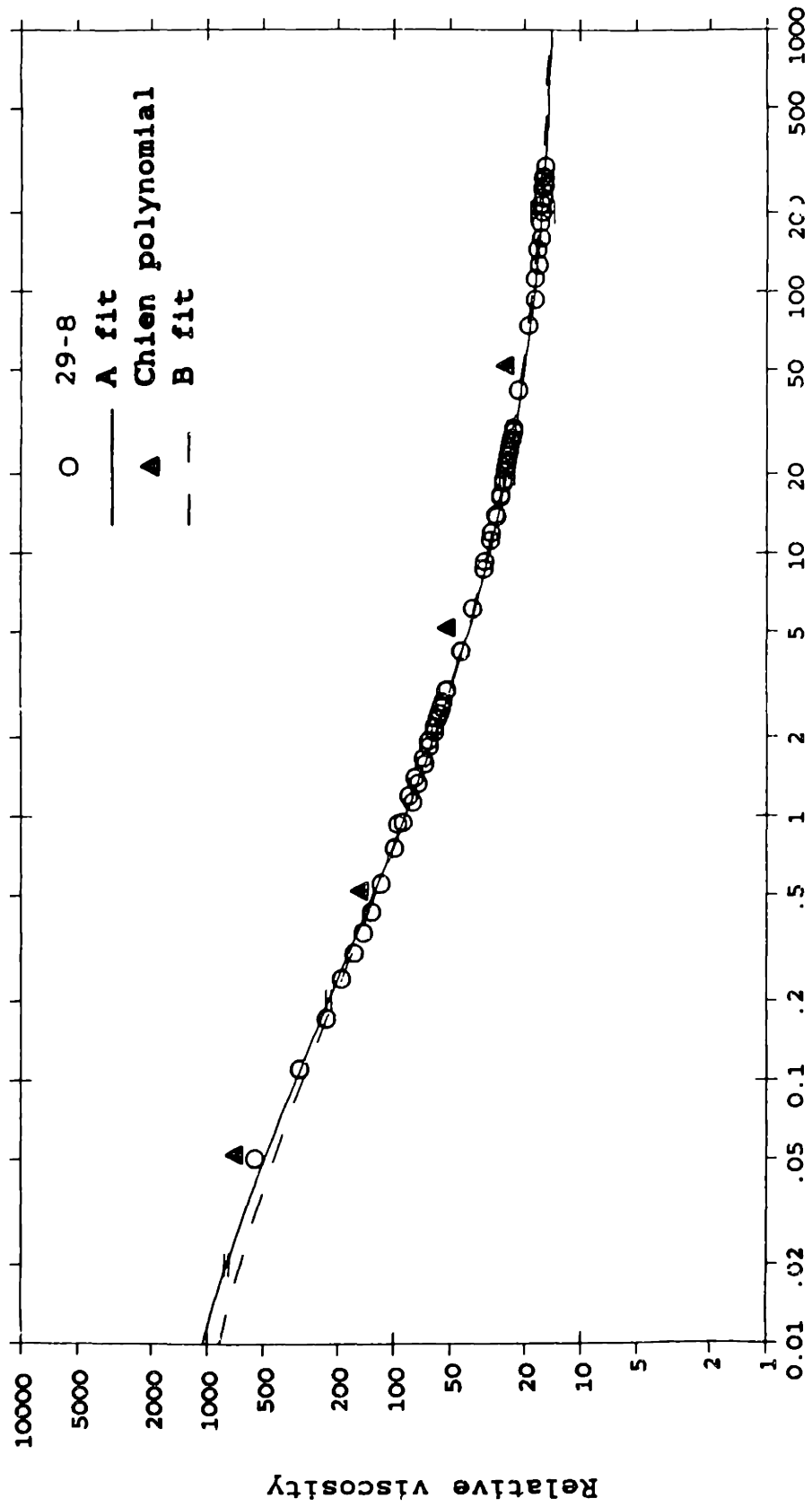


Figure 4-37: Plasma data, H=89.

Relative viscosity vs. Shear rate  
RBC in plasma, H=86

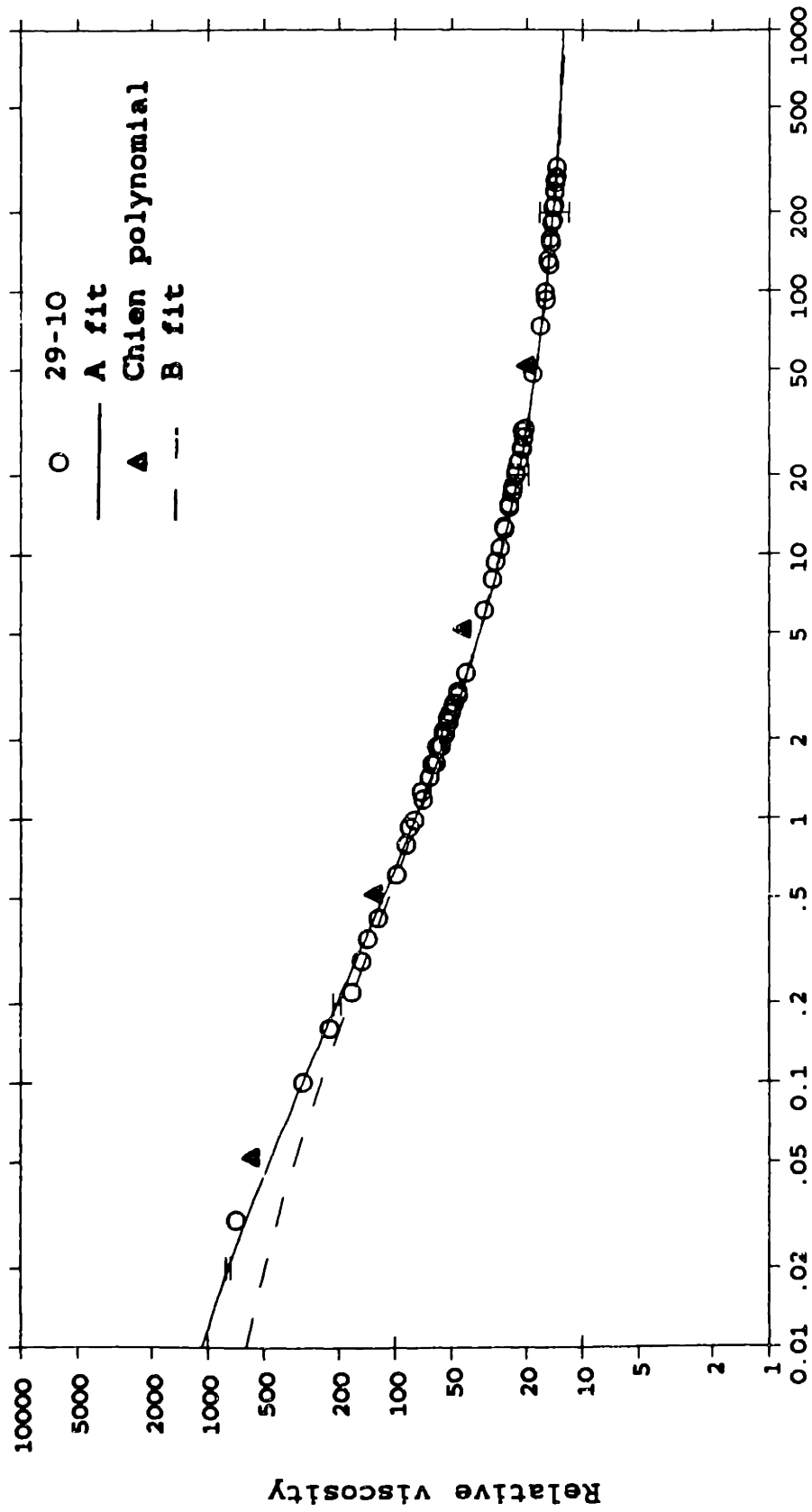


$\eta_{r,29}=1.05$  cP

Shear rate (1/sec)

Figure 4-38: Plasma data, H=86.

Relative viscosity vs. Shear rate  
RBC in plasma, H=83



$\eta_{l,29} = 1.05 \text{ cP}$

Shear rate (1/sec)

Figure 4-39: Plasma data, H=83.

Relative viscosity vs. Shear rate  
RBC in plasma, H=82

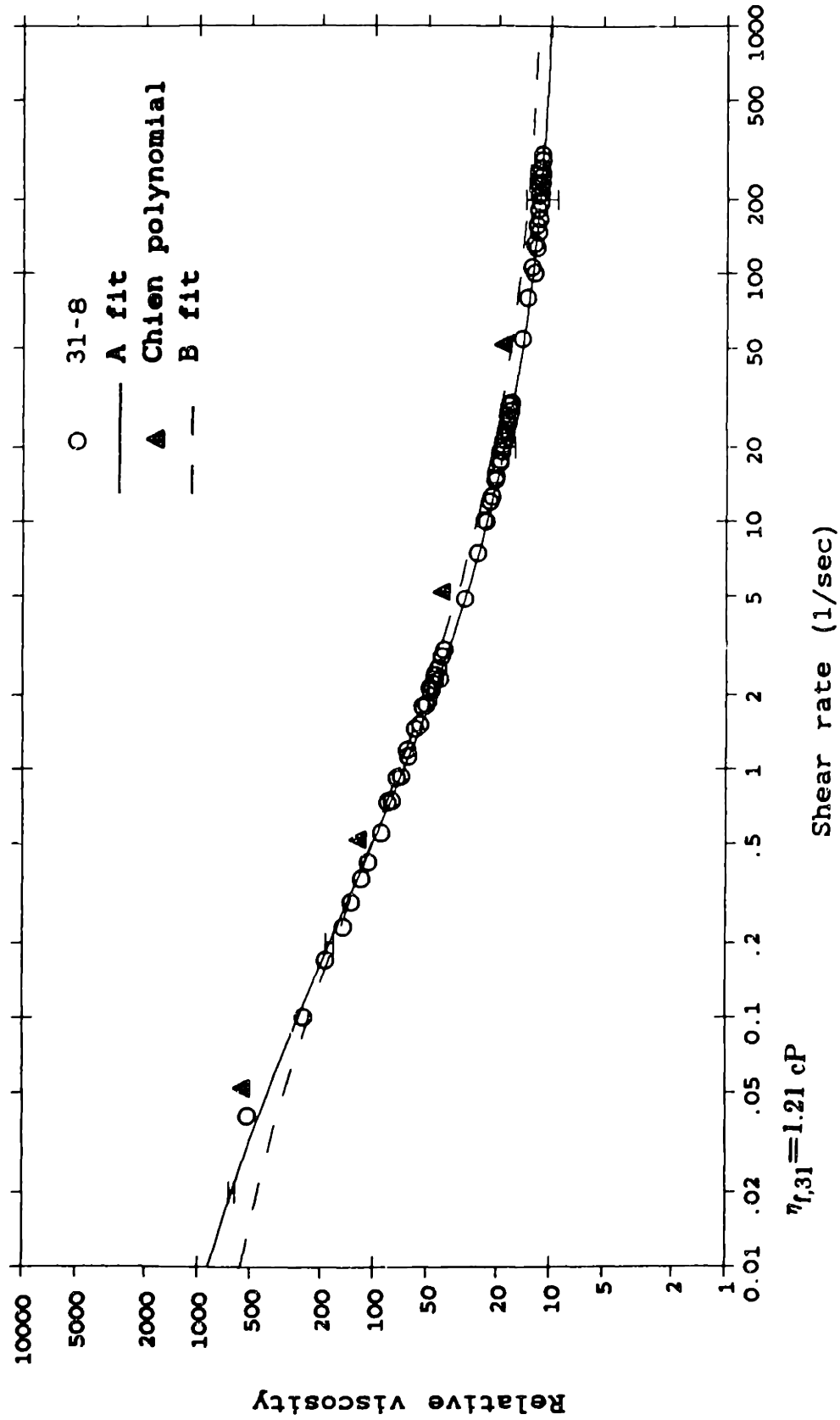


Figure 4-40: Plasma data, H=82.

Relative viscosity v.s Shear rate  
RBC in plasma, H=73

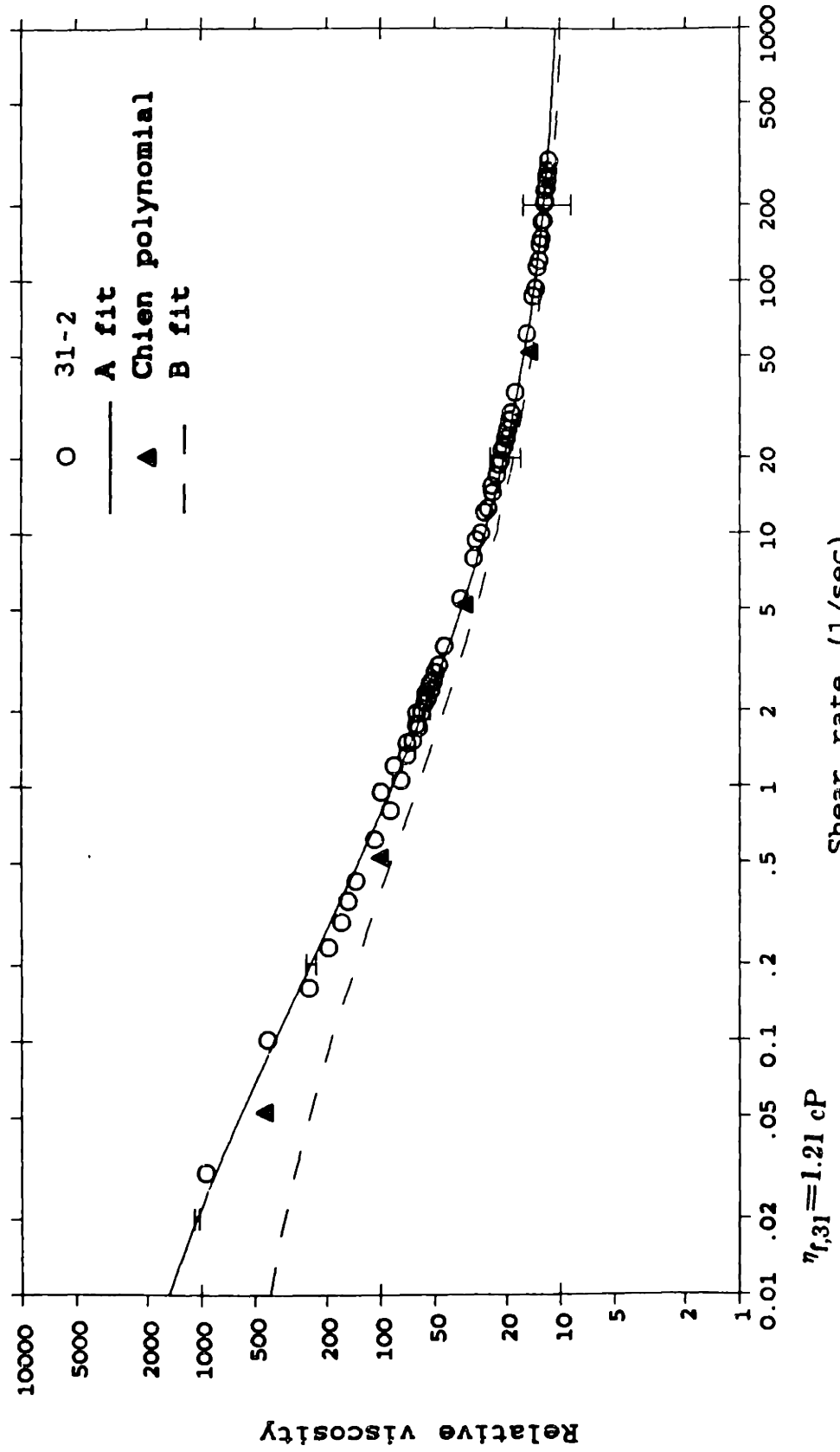
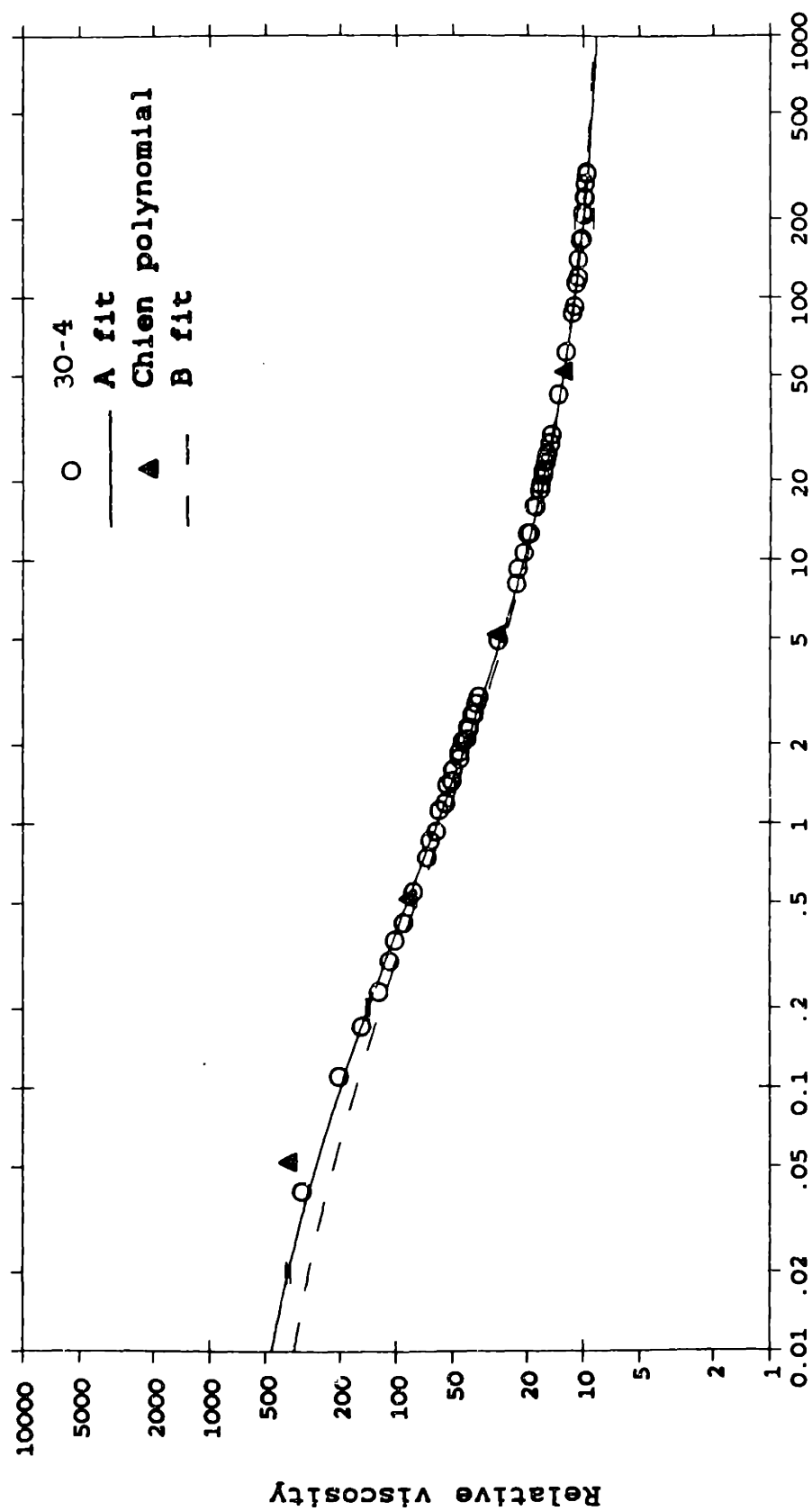


Figure 4-41: Plasma data, H=78.



Relative viscosity vs. Shear rate  
RBC in plasma, H=75



$\eta_{t,30} = 1.13 \text{ cP}$

Shear rate (1/sec)  
Figure 4-42: Plasma data, H=75.

Relative viscosity vs. Shear rate  
RBC in plasma, H=66

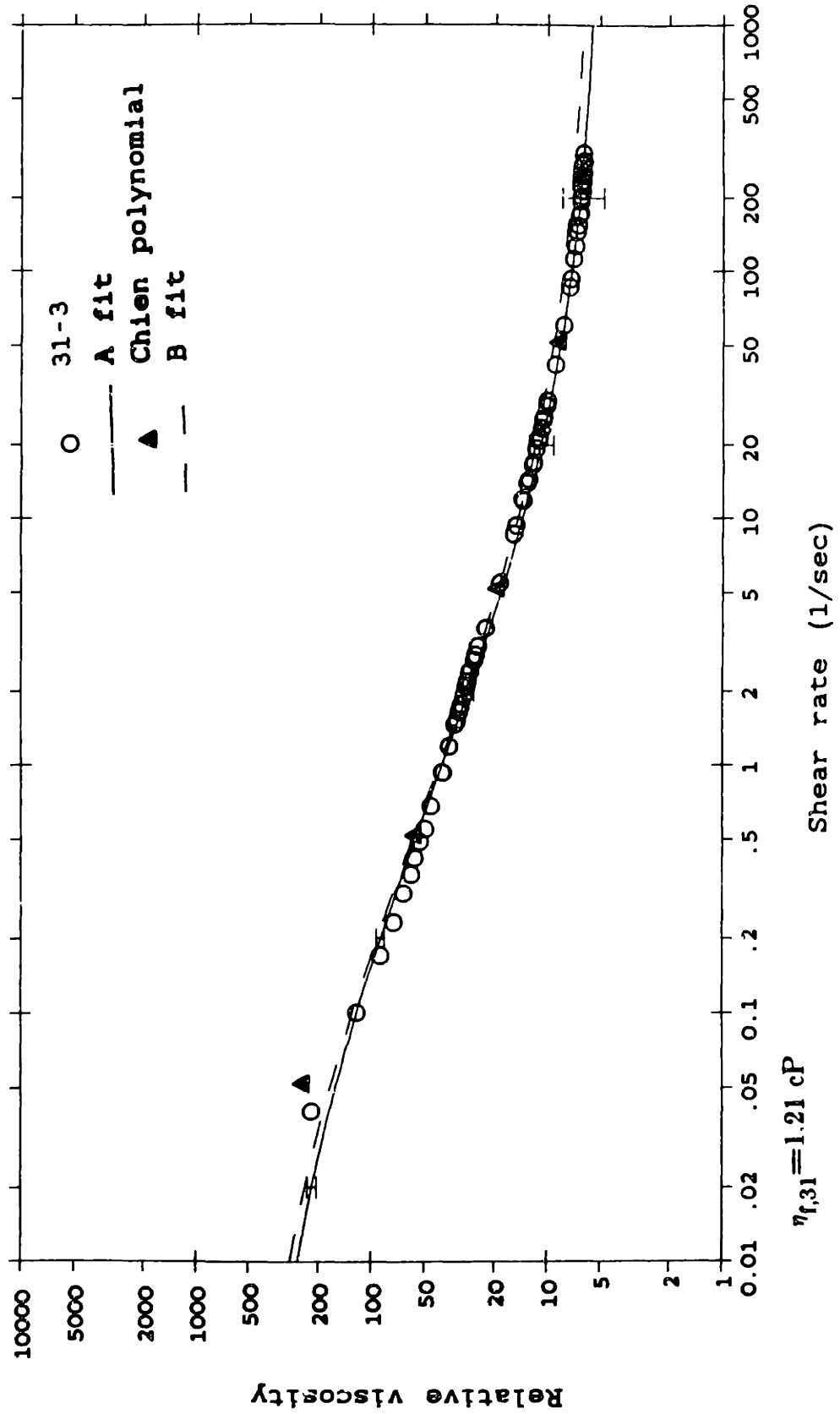
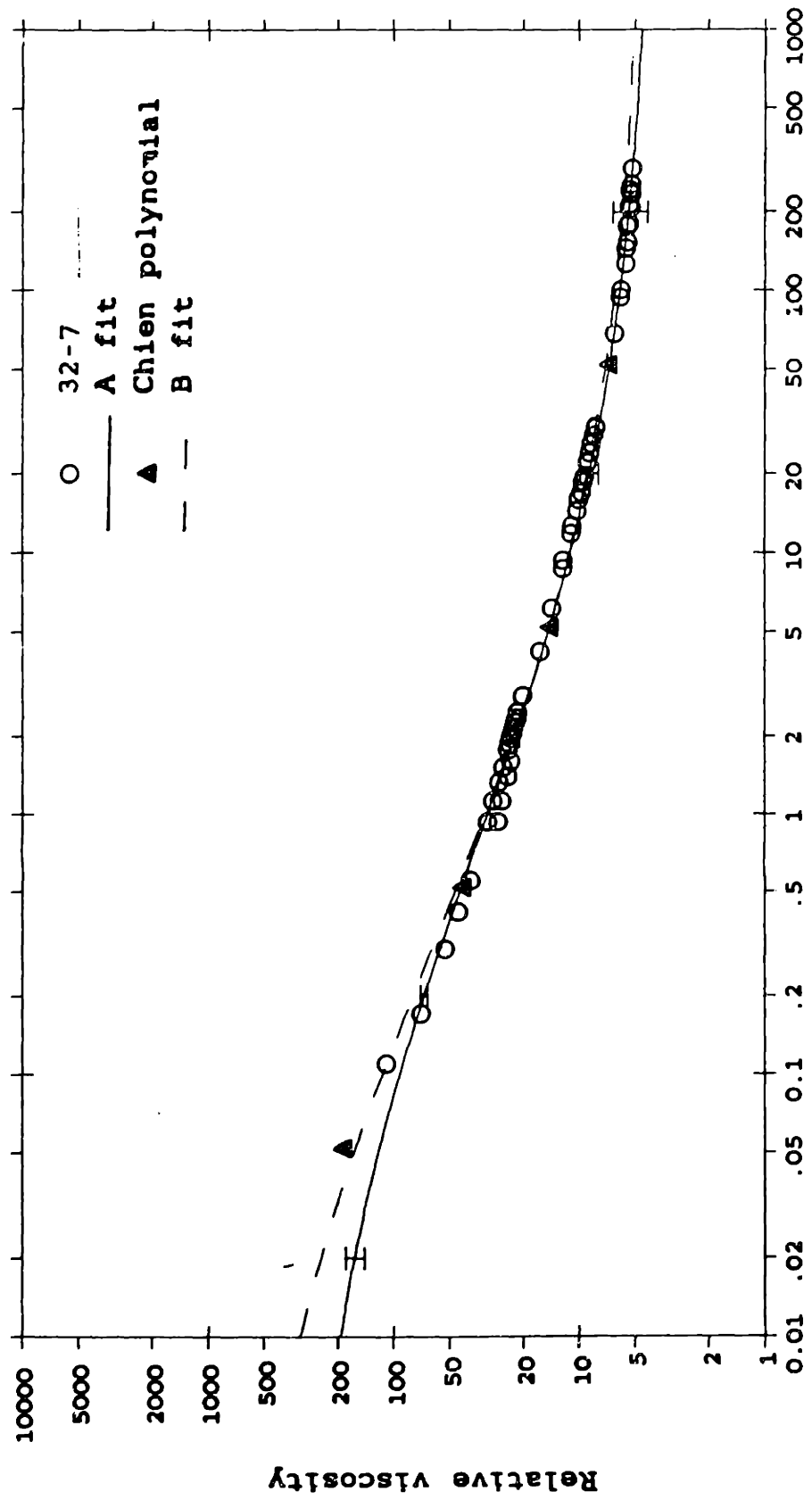


Figure 4-43: Plasma data, H=66.

Relative viscosity vs. Shear rate  
RBC in plasma, H=59



$\eta_{t,32} = 1.06 \text{ cP}$

Shear rate (1/sec)  
Figure 4-44: Plasma data, H=59.

Relative viscosity vs. Shear rate  
RBC in plasma, H=50

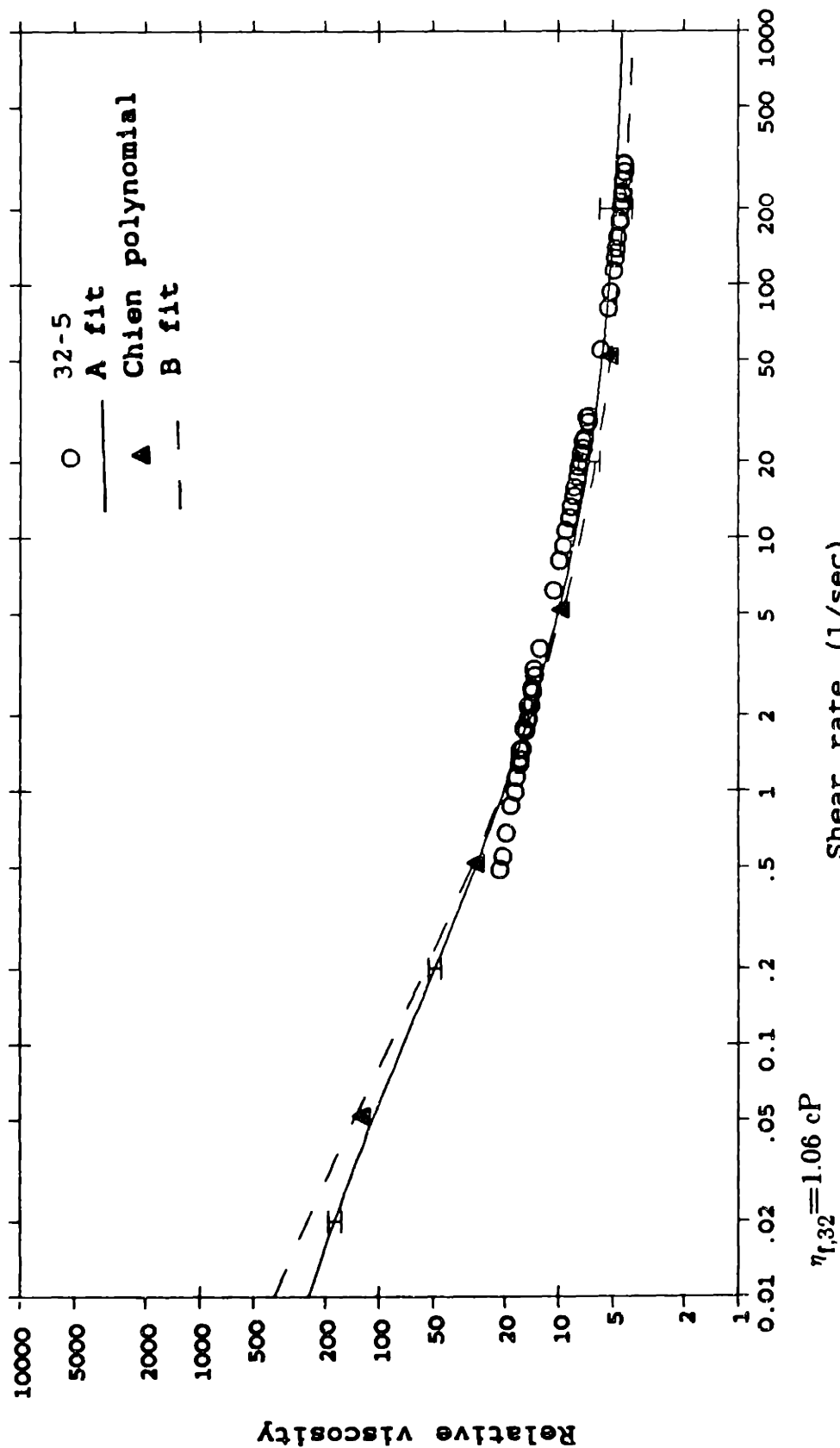
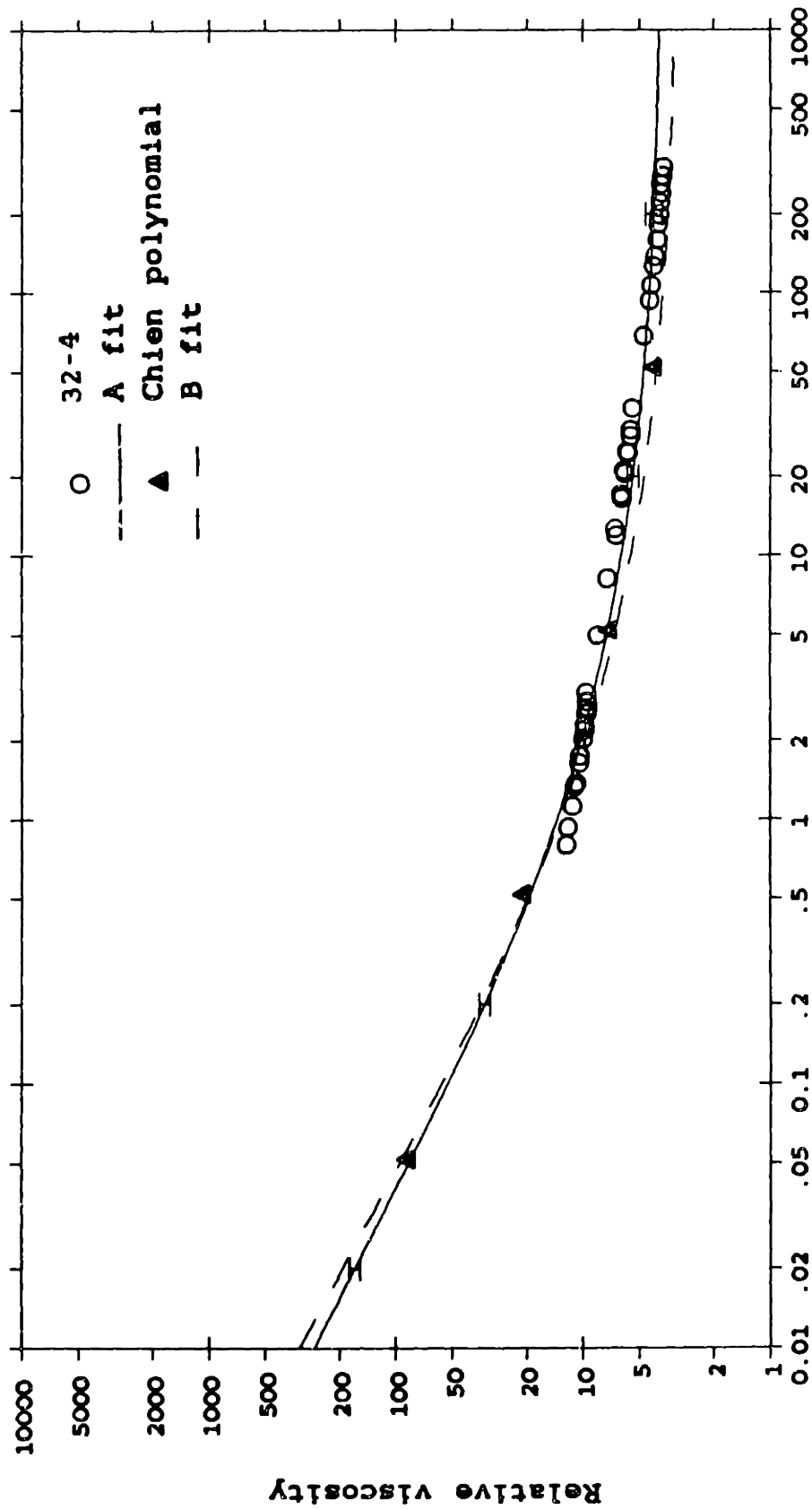


Figure 4-45: Plasma data, H=50.

Relative viscosity vs. Shear rate  
RBC in plasma, H=44



$\eta_{1,32} = 1.06 \text{ cP}$

Shear rate (1/sec)

Figure 4-46: Plasma data, H=44.

Relative viscosity vs. Shear rate  
RBC in plasma, H=34

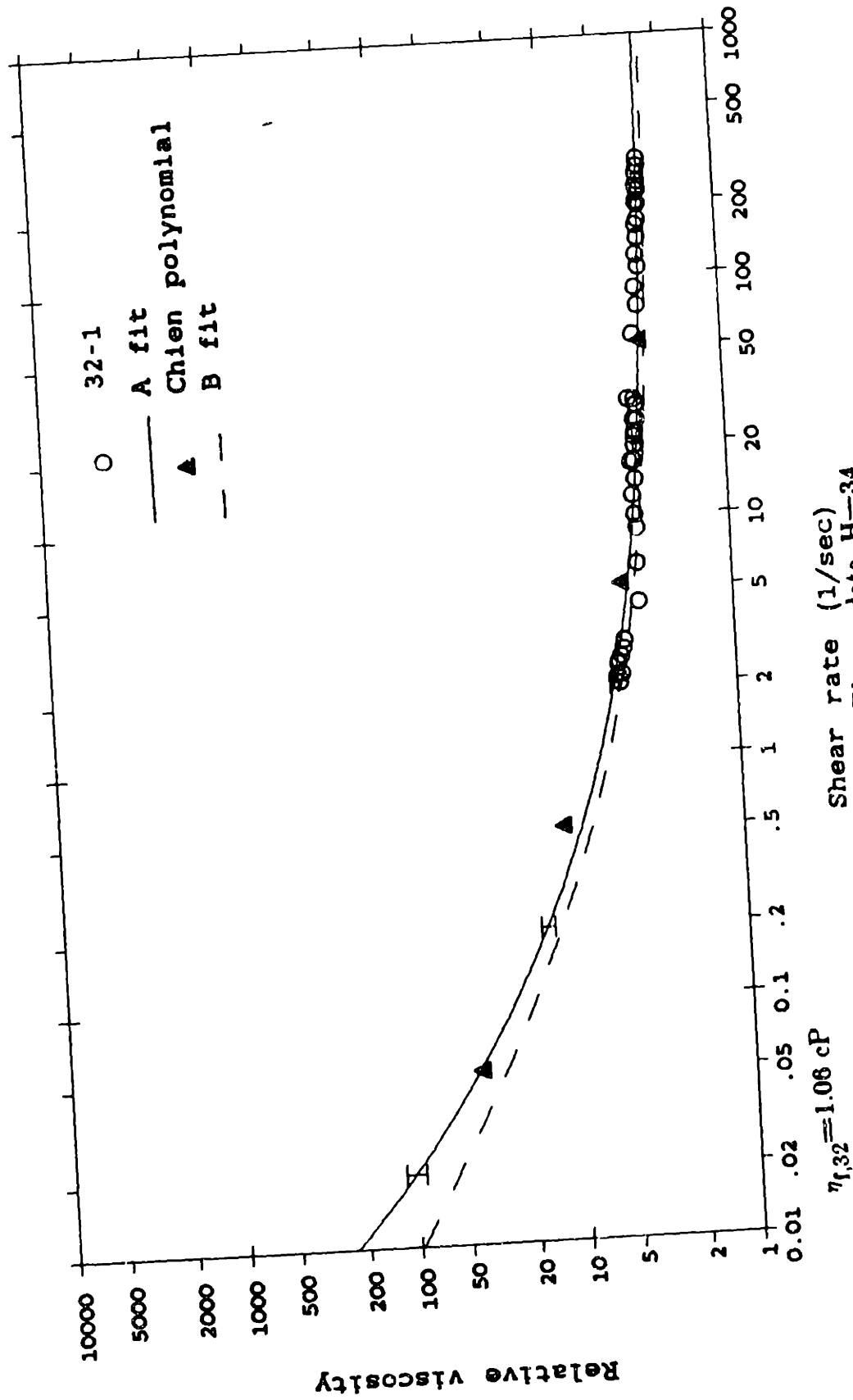


Figure 4-47: Plasma data, H=34.

Relative viscosity vs. Shear rate  
RBC in plasma, H=24

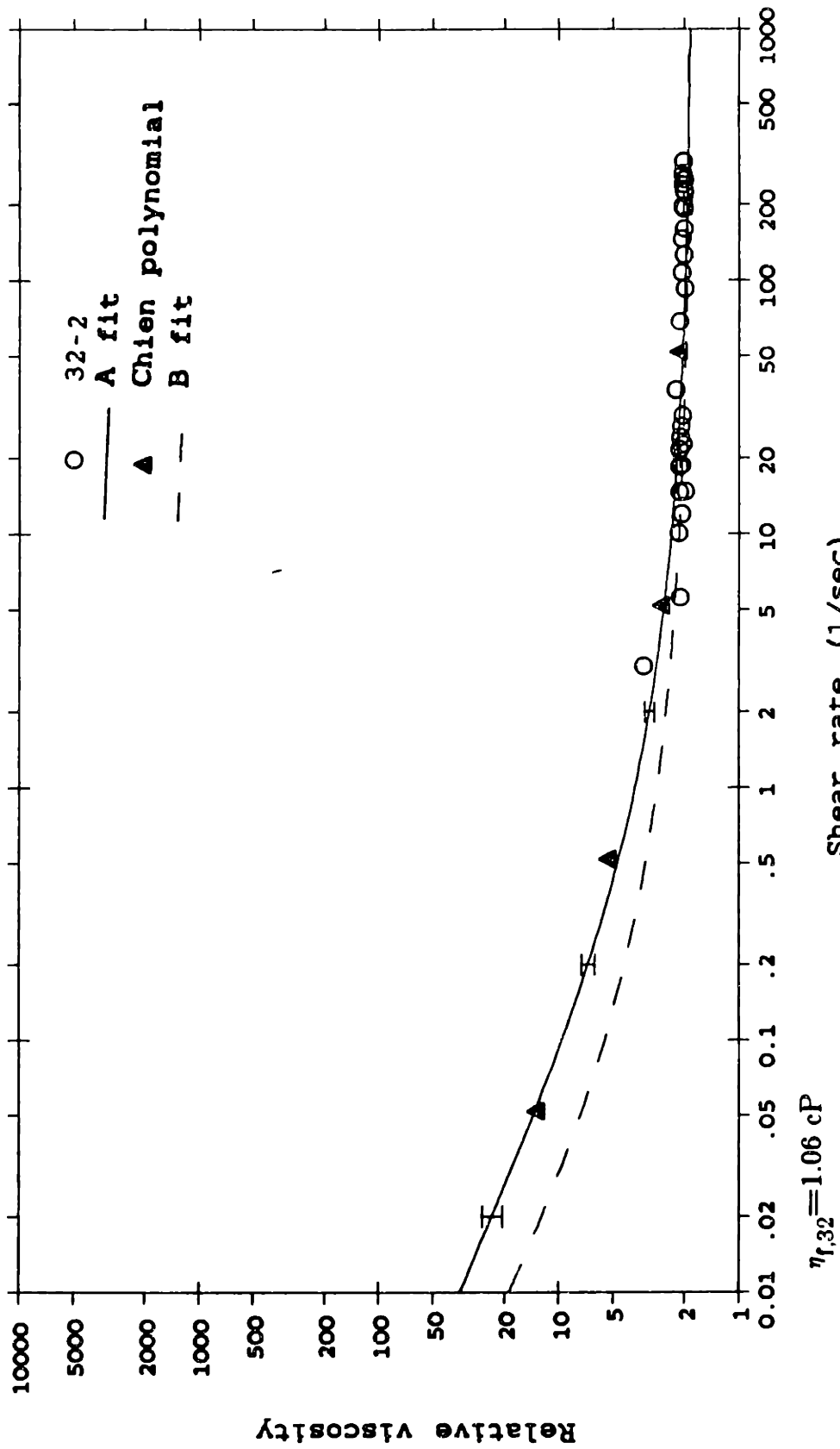
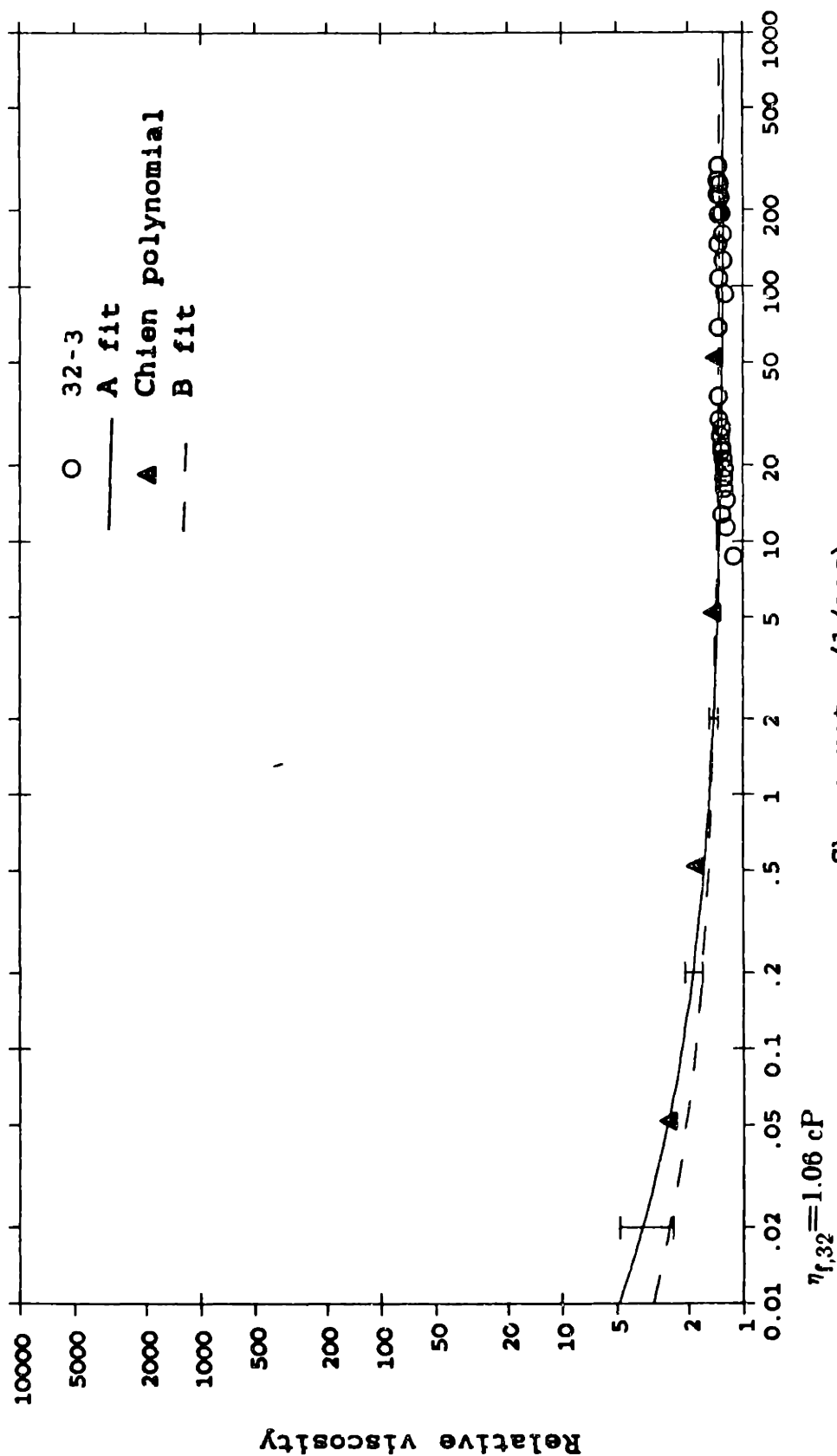


Figure 4-48: Plasma data, H=24.

Relative viscosity vs. Shear rate  
RBC in plasma, H=11



Shear rate (1/sec)

Figure 4-49: Plasma data, H=11.



three groups along the x-axis. This is an artifact resulting solely from the choice of the ranges of operation used and the selection of data points.

The data follows the trends for blood viscosity reported in the literature and reviewed in Chapter 2. Shear-thinning behavior is observed in all cases. At low shear rates, the viscosity seems to approach infinity at hematocrits above 90 and seems to level to an asymptote at hematocrits below 60, but there is not sufficient data to make a judgement as to the existence of a yield stress. Asymptotic values for viscosity are approached at the highest obtainable shear rate ( $300 \text{ sec}^{-1}$ ) for all plots, but only in the most concentrated suspensions ( $H > 94$ ) do they actually appear to be reached. The sensitivity of the viscometer limits the minimum shear rate at which data can be obtained: Above  $H=70$ , this minimum  $\dot{\gamma}$  is on the order of  $0.1 \text{ sec}^{-1}$  or less; above  $H=50$ , it is between  $0.1$  and  $1 \text{ sec}^{-1}$ ; below  $40$ , the minimum is around  $1 \text{ sec}^{-1}$ .

A general pattern of increasing viscosity with increasing hematocrit is observed, with  $\frac{d\eta}{dH}$  greatest at the highest concentrations. For example, at  $\dot{\gamma} = 100 \text{ sec}^{-1}$ , an increase in hematocrit from 89 to 97 (a 6.7% increase) increases viscosity by over 70%; while increasing hematocrit from 42 to 48 (a 14% increase) increases viscosity by only about 15%. Even at hematocrits above 98, however, blood is very fluid.

#### **4.3 Viscosity of RBC suspended in plasma**

The data for RBC suspended in plasma solutions are shown in Figures 4-26 to 4-49 (pages 93 to 116). The range of hematocrits is from 98.1 to 11. Relative viscosities were calculated from the individual plasma viscosities for each sample, which are tabulated in Appendix F. The symbol conventions are the same as for the saline suspensions. For the data of Chien *et al.*,  $\eta_f$  was calculated as 1.18 cP in a

manner analagous to that used for Ringer's saline. The descriptive comments applied to Eagle's solution suspensions in the previous section qualitatively apply to plasma suspensions. Because the plasma samples are for the most part more viscous than their corresponing saline samples, data could be obtained at somewhat lower shear rates but again, below hematocrits of about 40 it was not possible to get data over a broad spectrum of shear rates.

#### 4.4 Comparison of saline and plasma data

Data from the two types of RBC suspensions are plotted together as  $\eta_r$  vs. H in Figure 4-50 and as  $\eta$  vs. H in Figure 4-51. Values are shown at shear rates of 200, 20, 2, and 0.2 sec<sup>-1</sup>. In order to express values for all hematocrits at precisely each of these shear rates, the points shown are not actual data but are values from the Quemada equation fits (Fit A on the individual data plots) for each hematocrit. These fits are close enough to the actual data values that there would be no discernible difference on Figures 4-50 and 4-51. For the absolute viscosity, plasma suspension  $\eta_r$  was multiplied by a value of 1.1 cP for average plasma viscosity. Points are shown only where data was actually taken, thus there is no representation in this figure of the low hematocrit, low shear rate regions.

There is a greater amount of scatter from a monotonic relationship with hematocrit at low shear rates, where the data is less reliable, and a greater scatter in the saline suspensions versus the plasma suspensions. The relative values for Eagle's suspensions are consistently higher than those for plasma, although on an absolute basis, the plasma suspension viscosities are higher. For example, at H=96, the ratio of the relative viscosity in saline to that in plasma is between 1.29 and 1.49 at the shear rates given. For H=94 the range is 1.87 to 1.56; for H=90, 1.38 to 1.40; H=86, 1.06 to 1.40. There is no clear trend of the ratios increasing or

Comparison of plasma and Eagle's suspensions

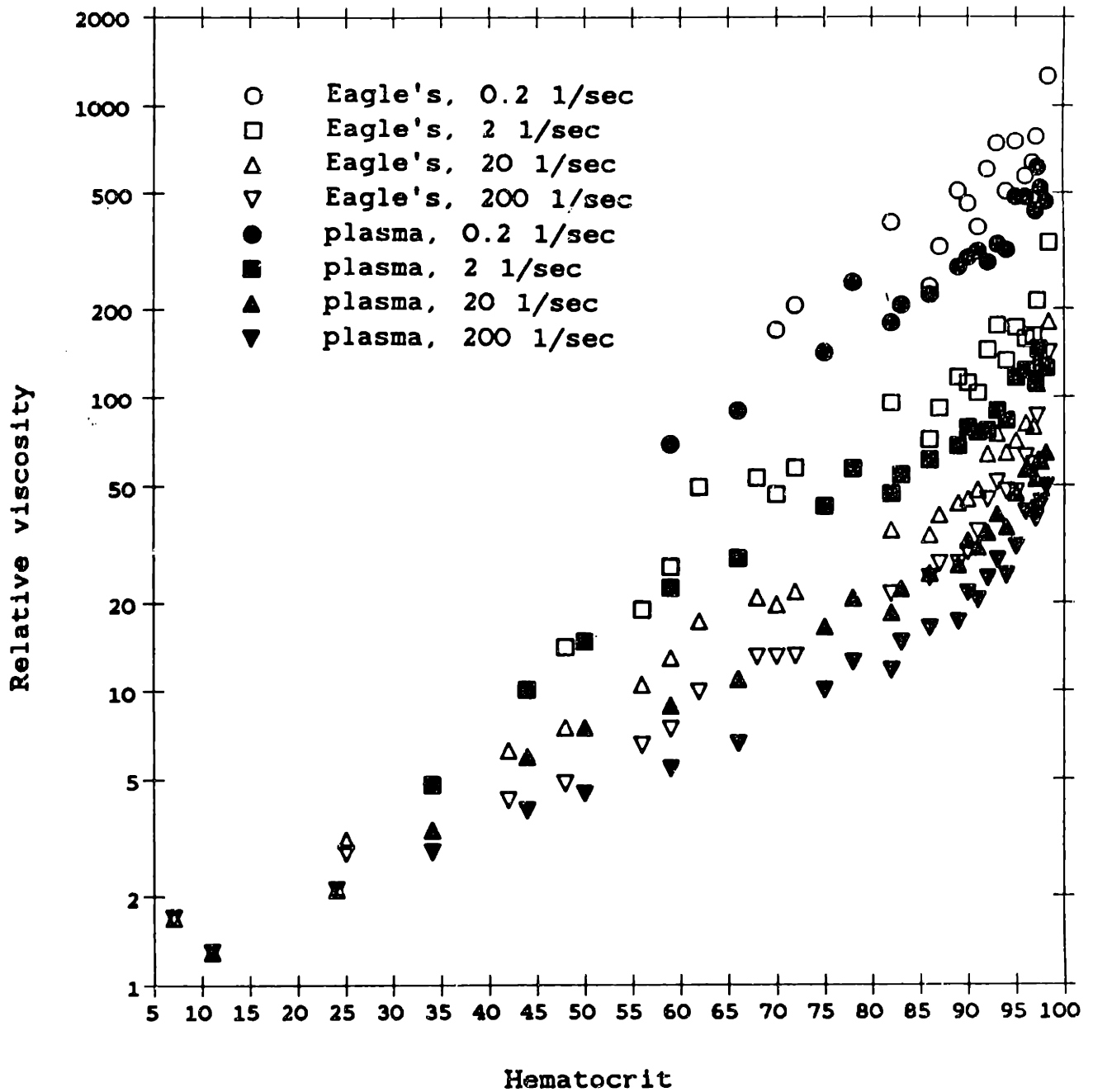


Figure 4-50: Relative viscosity vs. Hematocrit for Eagle's and plasma suspensions

Comparison of plasma and Eagle's suspensions

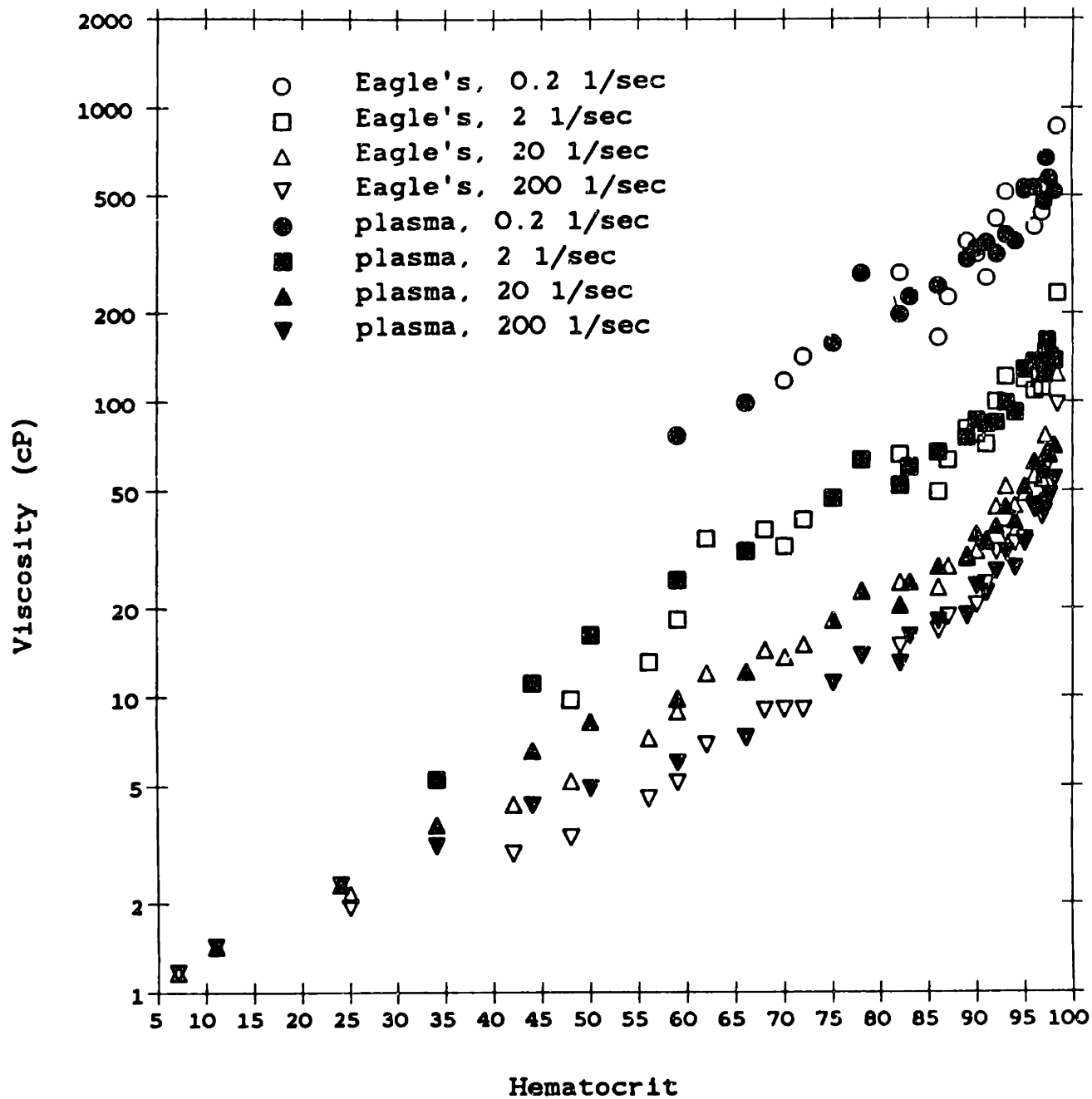


Figure 4-51: Absolute viscosity vs. Hematocrit for Eagle's and plasma suspensions

decreasing with shear rate. On a scale of absolute viscosities, the ratios just given would be decreased by a factor of about 0.6.

Figure 4-51 demonstrates how, on an absolute scale, the differences between saline and plasma suspensions decrease in the limit of high cell concentration. Above  $H=80$ , the data cannot be clearly distinguished between suspensions, particularly at high shear rates, and it appears at this point that an inversion phenomenon has taken place, with the cells behaving physically as the continuous phase. By an analysis of the behavior of the parameters for equation (2.11), Quemada [1981] has also suggested the importance of inversion, as will be discussed further in Chapter 5.

From Figures 4-50 and 4-51, it can be seen that, while the absolute viscosities become equal, the differences in relative viscosity are significant at high hematocrits. Viscosity becomes independent of  $\eta_f$  as the amount of suspending fluid becomes small, and the implication is that relative viscosity becomes a meaningless parameter in the limit of high hematocrits.

The decrease in relative viscosity with increasing viscosity of the continuous phase is a well-known aspect of the emulsion behavior of blood as described in Chapter 2. Dintenfass [1968b] suggests a modified form of the Taylor equation for emulsions, equation (2.5), to calculate the internal viscosity of the red cell:

$$\eta_r = \frac{1}{(1-cT)^{2.5}} \quad (4.1)$$

where  $T(\eta_i, \eta_f)$  is the Taylor parameter described for equation (2.5) on page 22 and  $c$  is the cell volume fraction. The shear rate dependence of equation (4.1) comes from the internal viscosity of the red cell, which represents not only the cytoplasmic viscosity but also includes the effects of the red cell membrane. As discussed previously in Section 2.4, this equation has limited applicability in that it is only valid when the shear rate is high enough to prevent rouleau formation. Fits of the

data at  $\dot{\gamma}=52 \text{ sec}^{-1}$  yield values for  $\eta_i$  around 2 cP, which is comparable to values found in the literature [Dintenfass, 1971]. The Dintenfass equation is used in the next section to provide an evaluation of how viscosity is related to changes in the red cell properties.

#### **4.5 Effect of erythrocyte age and anticoagulant**

It has been known for some time that the viscosity of blood increases with age [Erslev and Atwater, 1963; Usami *et al.*, 1971]. Erythrocytes have been shown to have increased hemoglobin concentration as they get older, and this along with possible loss of membrane deformability may increase the intrinsic viscosity of the cell.

The blood units obtained from the Red Cross were stored in a refrigerator at a temperature of about 40 °F. There was a noticeable increase in the viscosity of the samples with longer storage times. This effect was largest at low shear rates. For example, the plot of  $H=92$  in Eagle's solution, Figure 4-9, page 76, shows data for three samples from three units: Two (16-5 and 20-8) had their viscosities measured 10 days before the expiration date (11 and 15 days after collection, respectively) and one (18-4) was measured 2 days before expiration (19 days after collection). At  $300 \text{ sec}^{-1}$  the oldest sample has the highest viscosity, but its value is almost identical to that of one of the fresher samples and is 1.3 times that of the other. As shear rate decreases, the difference in viscosity increases substantially. At  $\dot{\gamma} = 0.3 \text{ sec}^{-1}$  the oldest sample has a viscosity 1.75 times that of the other two. This is not attributable to storage in the laboratory refrigerator because both samples 16-5 and 18-4 were measured two days after being obtained from the Red Cross.

Relevant samples which illustrate the effect of storage time are described in

sample	suspending fluid	anti- coagulant	H	d.s.c.	d.b.e.	v.s.f. (cP)
16-5	Eagle's	CPD	92	11	10	0.69
20-8	Eagle's	CPD-A1	92	14	10	0.69
18-4	Eagle's	CPD	92	25	2	0.69
19-10	Eagle's	CPD-A1	91	22	13	0.69
20-9	Eagle's	CPD-A1	91	25	10	0.69
16-3	Eagle's	CPD	90	11	10	0.69
18-9	Eagle's	CPD	90	19	2	0.69
16-9	Eagle's	CPD	87	11	10	0.69
20-10	Eagle's	CPD-A1	87	25	10	0.69
16-8	Eagle's	CPD	70	11	10	0.69
24-2	Eagle's	CPD-A1	70	35	0	0.69
29-9	plasma	AS-1	96	9	33	1.05
31-1	plasma	AS-1	96	11	31	1.21
30-1	plasma	AS-1	91	10	32	1.13
31-4	plasma	AS-1	91	11	31	1.21
29-2	plasma	AS-1	89	9	33	1.05
31-6	plasma	AS-1	89	11	31	1.21

d.s.c.--days since collection  
d.b.e.--days before expiration  
v.s.f.--viscosity of suspending fluid

**Table 4-I:** Description of samples which illustrate effect of erythrocyte age.

sample	visc. (cP) @ 0.3 1/sec	rel. visc.	visc. (cP) @ 3 1/sec	rel. visc.	visc. (cP) @ 30 1/sec	rel. visc.
16-5	208	301.4	69.3	100.4	38.7	56.1
20-8	223	323.2	66.0	95.7	32.5	47.1
18-4	392	568.1	99.3	143.9	43.2	62.6
19-10	184	266.7	59.0	85.5	29.6	42.9
20-9	206	298.6	63.0	91.3	31.0	44.9
16-3	175	253.6	53.0	76.8	27.5	39.9
18-9	260	376.8	89.0	129.0	27.7	40.1
16-9	170	246.4	56.0	81.2	27.0	39.1
20-10	161	233.3	51.0	73.9	23.6	34.2
16-8	58	84.1	23.0	33.3	10.2	14.8
24-2	95	137.7	32.0	46.4	14.0	20.3
29-9	382	363.8	114.0	108.6	59.2	56.4
31-1	438	362.0	126.0	104.1	54.0	44.6
30-1	233	206.2	62.0	54.9	34.0	30.1
31-4	293	242.1	80.0	66.1	33.6	27.8
29-2	232	221.0	65.0	61.9	26.6	25.3
31-6	210	173.6	61.5	50.8	26.7	22.1

sample	visc. (cP) @ 300 1/sec	rel. visc.	int. visc. (cP) @ 30 1/sec	int. visc. (cP) @ 300 1/sec
16-5	31.2	45.2	2.49	2.08
20-8	24.7	35.8	2.15	1.70
18-4	32.5	47.1	2.74	2.15
19-10	22.5	32.6	2.16	1.69
20-9	23.5	34.1	2.25	1.76
16-3	20.4	29.6	2.20	1.67
18-9	19.0	27.5	2.21	1.56
16-9	18.0	26.1	2.89	1.86
20-10	16.7	24.2	2.49	1.72
16-8	7.0	10.1	6.47	2.33
24-2	9.1	13.2	can't calculate	4.45
29-9	45.0	42.9	2.74	2.27
31-1	40.0	33.1	2.69	2.16
30-1	20.0	17.7	2.58	1.59
31-4	22.6	18.7	2.58	1.79
29-2	17.6	16.8	2.36	1.58
31-6	18.0	14.9	2.38	1.61

**Table 4-II: Viscosities of samples from Table 4-I.**



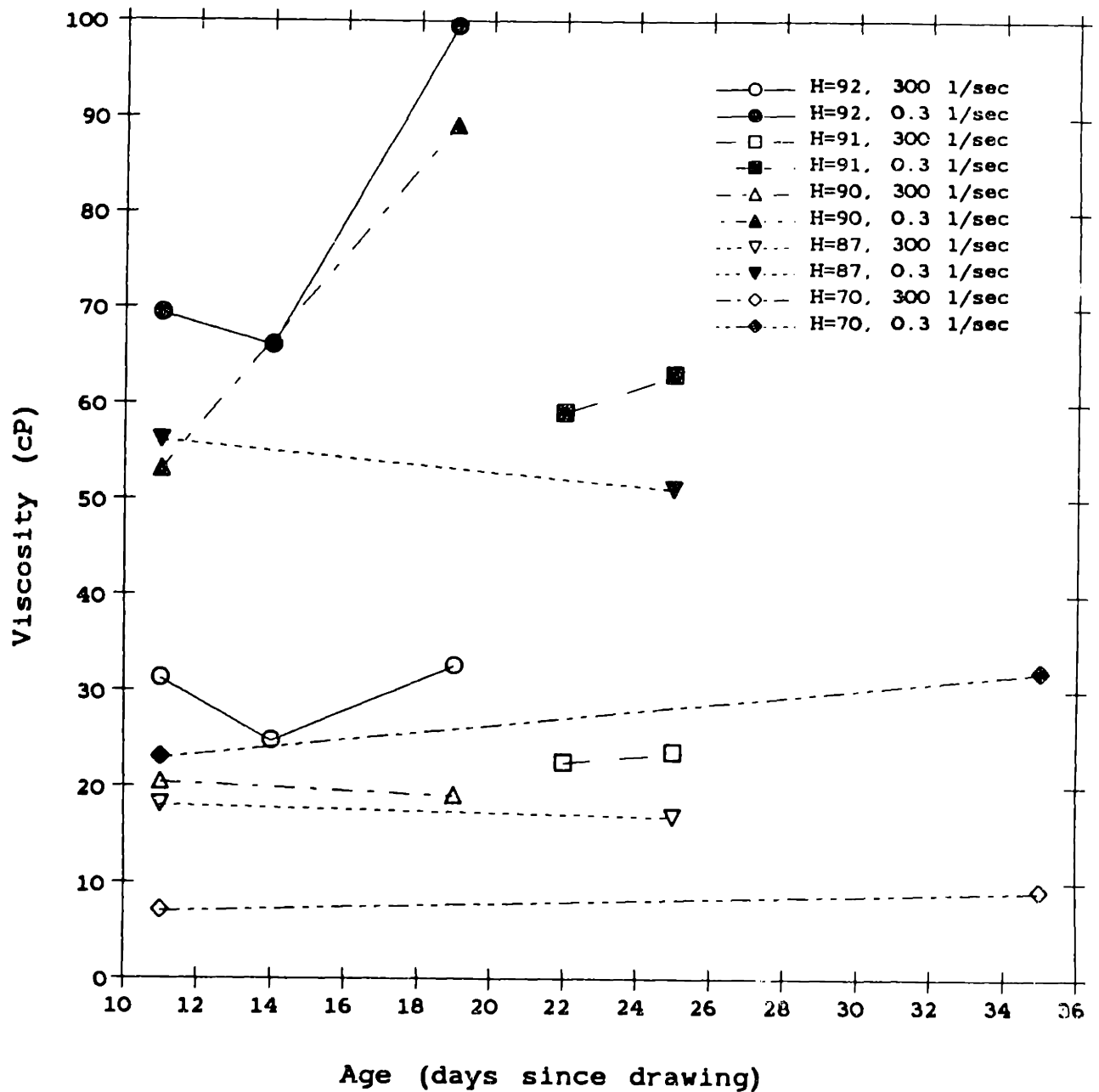
Table 4-I and their viscosities at selected shear rates are given in Table 4-II. A complete catalog of all sample ages and anticoagulants is given in Appendix F.

Similar age effects are observed in the Eagle's suspensions with  $H=91$ , 90 and 70. The samples for  $H=91$  were obtained from the same unit, one measured 22 days after collection, the other 25 days after collection. For  $H=90$ , the difference is 11 days versus 19 days. At  $H=70$ , a sample 11 days after collection is compared to one measured on its date of expiration, 35 days after collection. In all of these cases the older blood had higher viscosity. These results are summarized graphically in Figure 4-52. The increases in viscosity at  $0.3 \text{ sec}^{-1}$  suggest an increase in the aggregation of the red cells, since this is the major determinant of suspension viscosity at low shear rate. The data for  $H=92$  and  $H=87$ , which show decrease in viscosity with age, are probably related to the type of anticoagulant-preservative solution, as CPD- $A_1$ -preserved suspensions consistently have lower viscosities than CPD-preserved suspensions.

There are analogous increases in viscosity with age for plasma suspensions (see Figure 4-53), and these are compounded by the fact that the plasma itself grows more viscous with time. The plasma data at  $H=96$ , 91, and 89 all come from the same Red Cross unit and were measured over the course of three days, during which time the plasma viscosity increased from 1.05 to 1.13 to 1.21 cp. These samples were preserved in  $AS_1$ , which has a storage time of 42 days.

The Dintenfass equation, (4.1), was used to calculate the internal viscosity for the samples listed above for which age comparisons were made. These values are shown in Table 4-II. For saline suspensions they provide no additional information since  $\eta_i$  is a constant for all saline samples, but for plasma suspensions they provide an indication of the relative changes in the cell fluidity versus those in the plasma. The values for  $\eta_i$  in Table 4-II indicate either slight increases or no change in the

Viscosity vs. Age of blood  
Eagle's solution suspensions



**Figure 4-52:** Effect of sample age on viscosity of Eagle's solution suspensions.

Viscosity vs. Age of blood  
Plasma suspensions

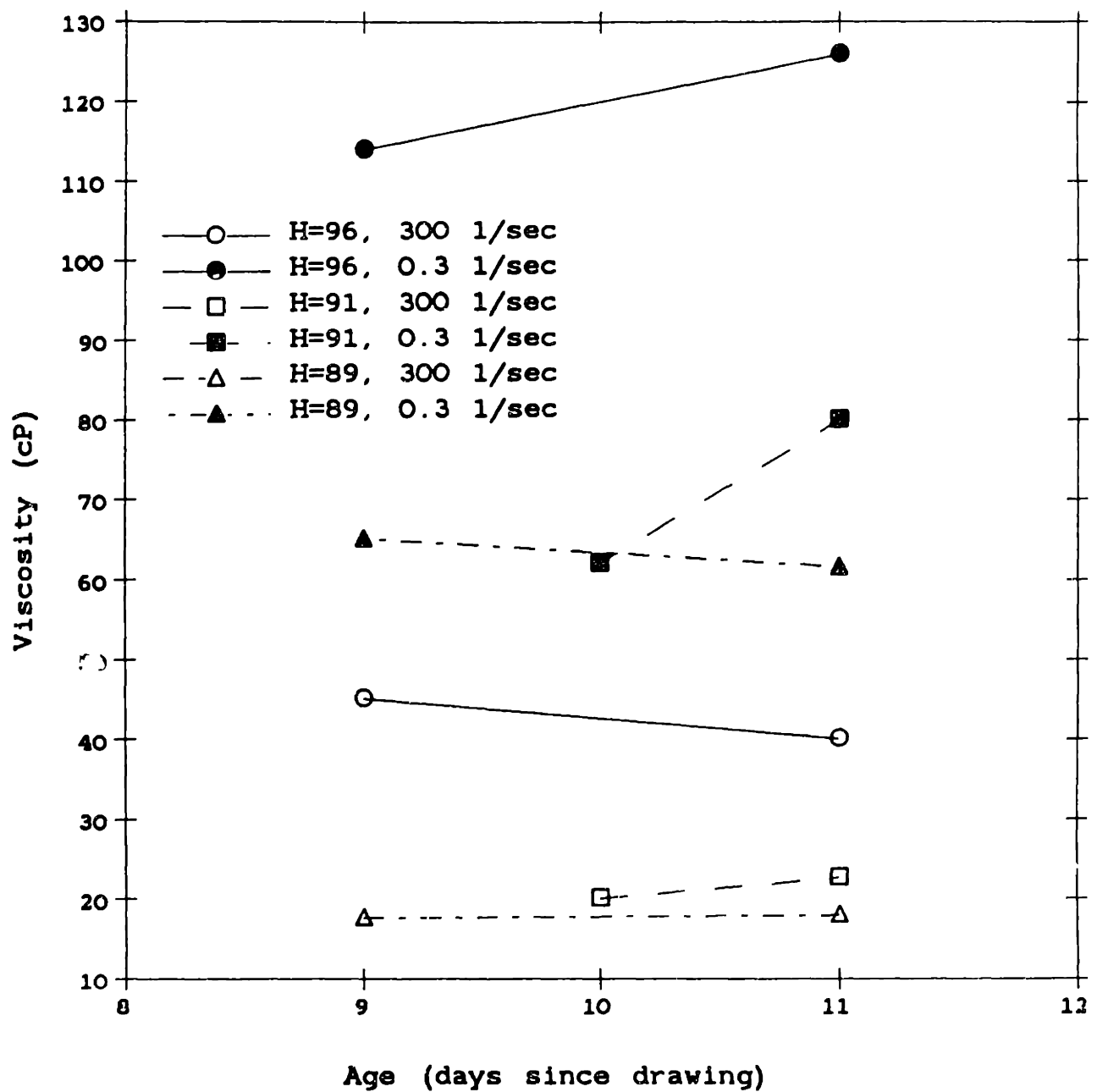


Figure 4-53: Effect of sample age on viscosity of plasma suspensions.

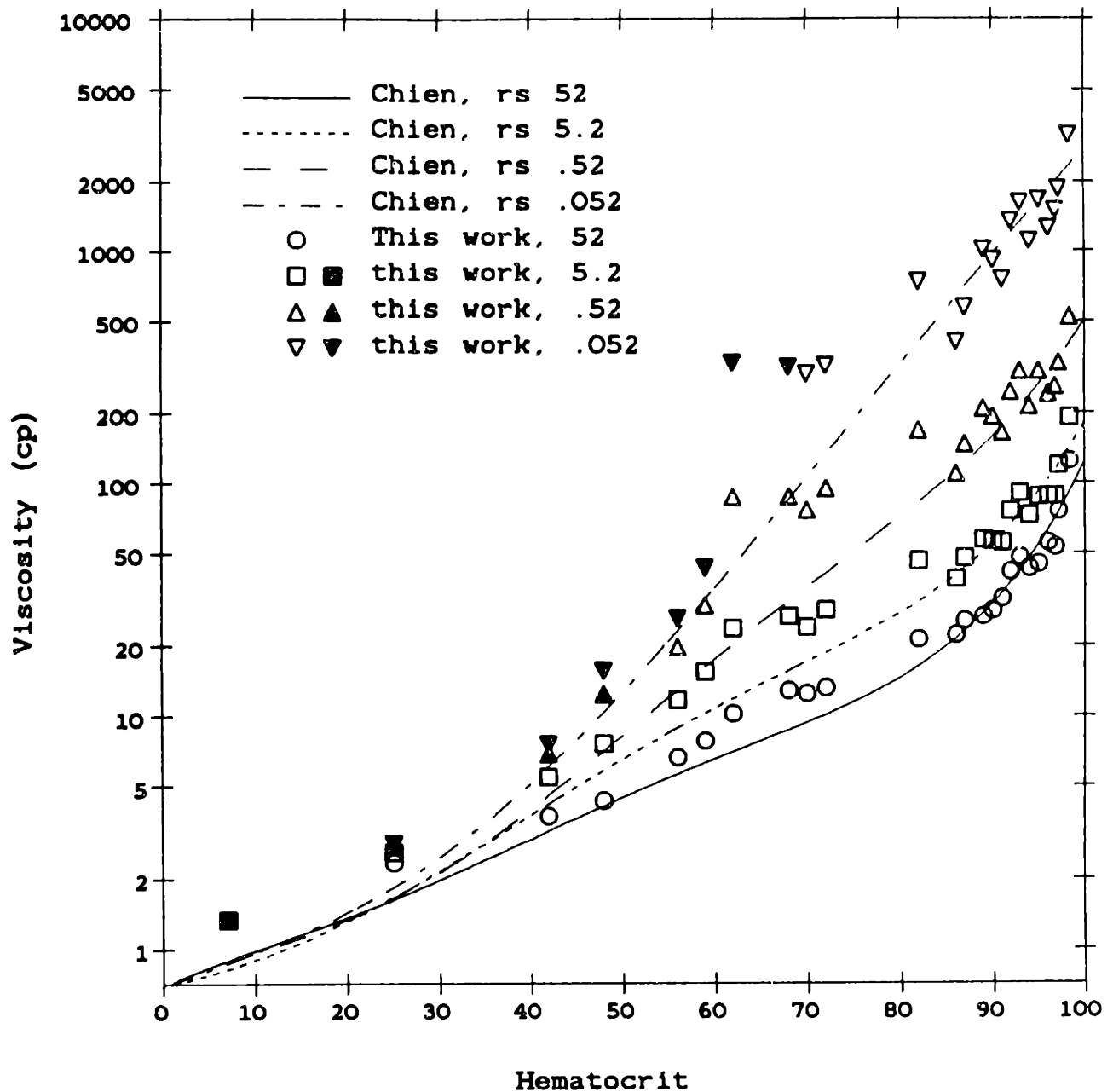
internal viscosity with time, but the Dintenfass equation is not applicable at low shear rates, where the age effects would be most noticeable.

To summarize, with all other things being equal, there is a rise in the viscosity of saline suspensions with erythrocyte age that is a reflection solely of the internal viscosity of the cell and which are more striking at low shear rates. Longer-lasting anticoagulant-preservative solutions appear to retard the age effects, to the extent that they are relatively insignificant. For plasma suspensions, the internal viscosity may rise slightly with storage in AS<sub>1</sub>, and the plasma viscosity increases steadily with age.

#### **4.6 Comparison to data of Chien *et al.*, 1966.**

The data from Figure 4-51 is replotted versus curves from the equations of Chien *et al.* [1966] in Figures 4-54 and 4-55. Since Chien *et al.* reports only absolute viscosities, the values for plasma suspensions in this study have been adjusted by multiplying them by a value of 1.1 cP for average plasma viscosity, which is an almost undetectable difference in these plots. The general trends of the two sets of data are similar. It must be remembered that the data points in these figures which represent the work of this paper are actually derived values from the curve fits, and therefore in the case of the lower hematocrits (less than 59 for plasma, less than 70 for saline) at low shear rates ( $0.052 \text{ sec}^{-1}$  and  $0.52 \text{ sec}^{-1}$ ) are only extrapolations of data. Filled symbols are used on Figures 4-54 and 4-55 to represent these extrapolations. The data for saline are comparable to each other at high hematocrits; below 85, the Chien data are consistently lower by factors of 2 or more. On Figures 4-20 to 4-23, the Chien values display a difference curvature with regards to shear rate dependence than the data of this work. The Chien points can be described as exhibiting a concavity, while the data shows convexity. More recent

Comparison of Chien data with this work  
RBC in saline solution

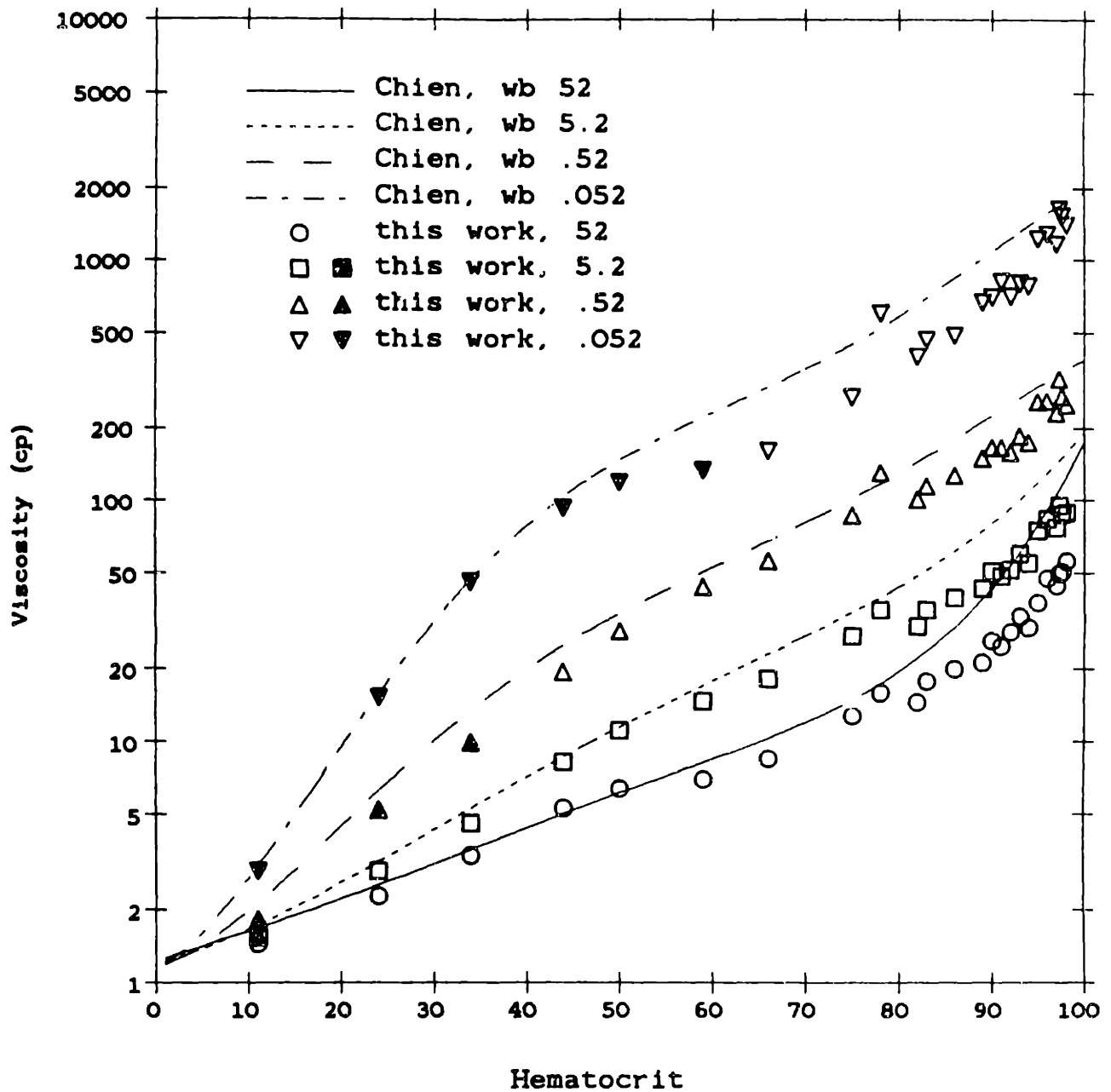


**Figure 4-54:** Comparison of data of this work to equations of Chien *et al.*, 1966 for saline suspensions.

"rs"—Ringer's Solution

Numbers in legend represent shear rate in  $\text{sec}^{-1}$ .

Comparison of Chien data with this work  
RBC in plasma



**Figure 4-55:** Comparison of data of this work to equations of Chien *et al.*, 1966 for plasma suspensions.

"wb"--whole blood

Numbers in legend represent shear rate in  $\text{sec}^{-1}$ .

data by Chien [1970], shown in Figure 2-3, also has a convex curvature. As will be discussed in Chapter 5, the convexity is the more physically reasonable behavior from the Quemada model. Perhaps by nature of their form, the Chien *et al.* [1966] equations, which model the hematocrit dependence by fifth-order polynomials, introduce a concavity artifact when plotted versus shear rate.

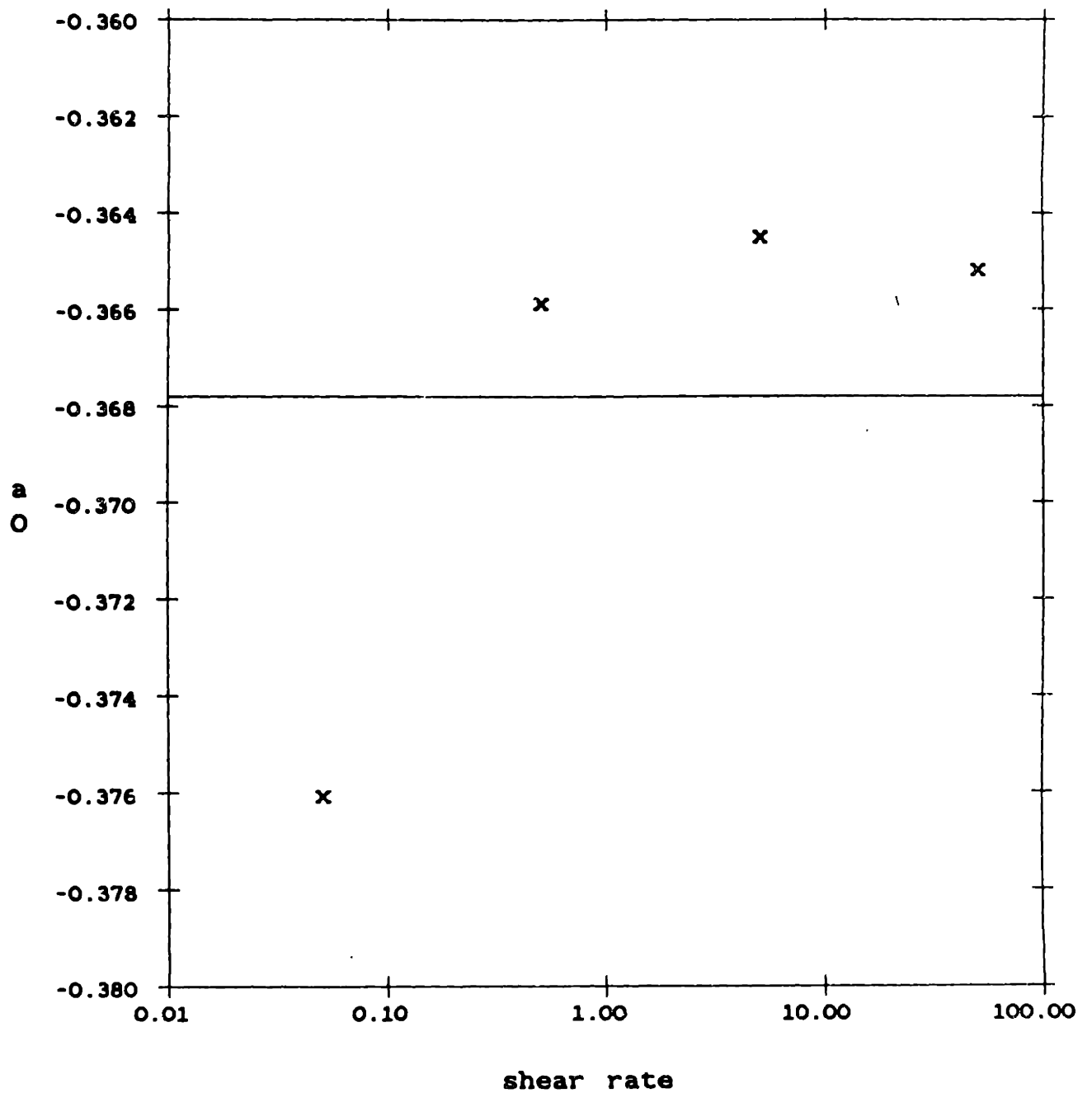
For plasma suspensions, the data of this work is lower than the Chien *et al.* values at the combination of high hematocrits and high shear rates. The magnitude of this difference ranges from 10-30% to factors of 2 or greater. This could be a reflection of the difference between the freshly drawn blood used by Chien *et al.* and the stored units used in this work, although the stored units would be expected to have a higher viscosity. Alternatively, dilution of plasma proteins by the anticoagulant-preservative could lower the viscosity, but this would have its largest impact at low shear rates. At hematocrits less than 80, the two sets of data are quite compatible with one another.

#### 4.7 Global refitting of Chien *et al.* polynomials

It was desirable to examine whether the constants given in Table 2-I for equation (2.8) could be fitted to functions of shear rate, and thus the utility of the Chien polynomials expanded to a wider range. The values for saline suspensions oscillate greatly with shear rate (see Figure 4-56(a-f)) and can be best approximated by polynomials of  $\ln(\dot{\gamma})$  of the form:

$$a_i = \sum_{j=0}^3 \alpha_{ij} [\ln(\dot{\gamma})]^j \quad (4.2)$$

The constants  $\alpha_{ij}$  for these equation are given in Table 4-III. When such equations are plotted versus shear rate at a single hematocrit, they are seen to provide a good fit of data at high hematocrits (see Figure 4-57). As hematocrit decreases, the curve



**Figure 4-56:** Plots of constants for Chien polynomial for saline versus shear rate.

(a) Plot for  $a_0$  in Ringer's solution.



Chien's rs a1

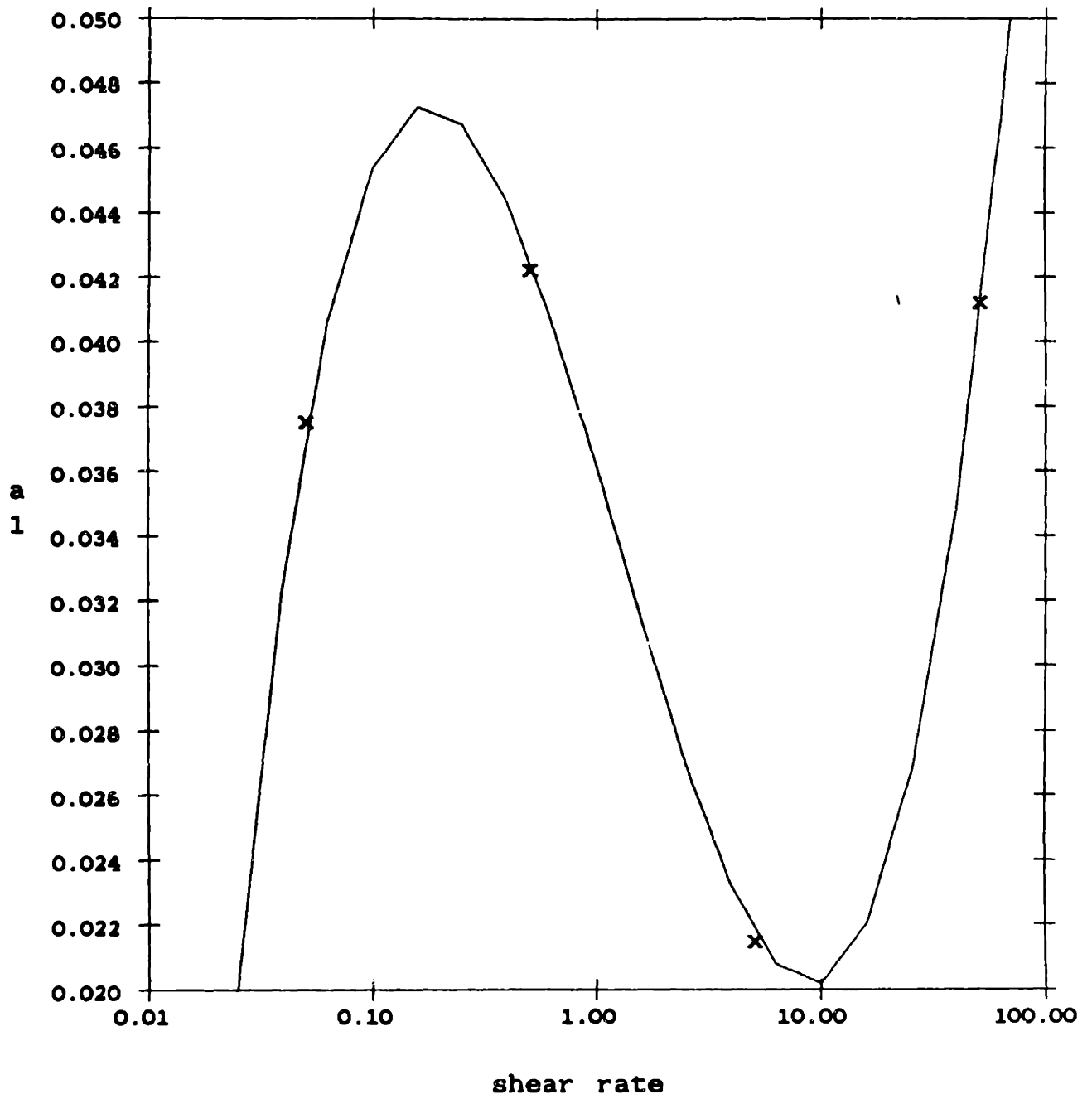


Figure 4-56(b) Plot for  $a_1$  in Ringer's solution.

Chien's  $rs a_2 \times 10^4$

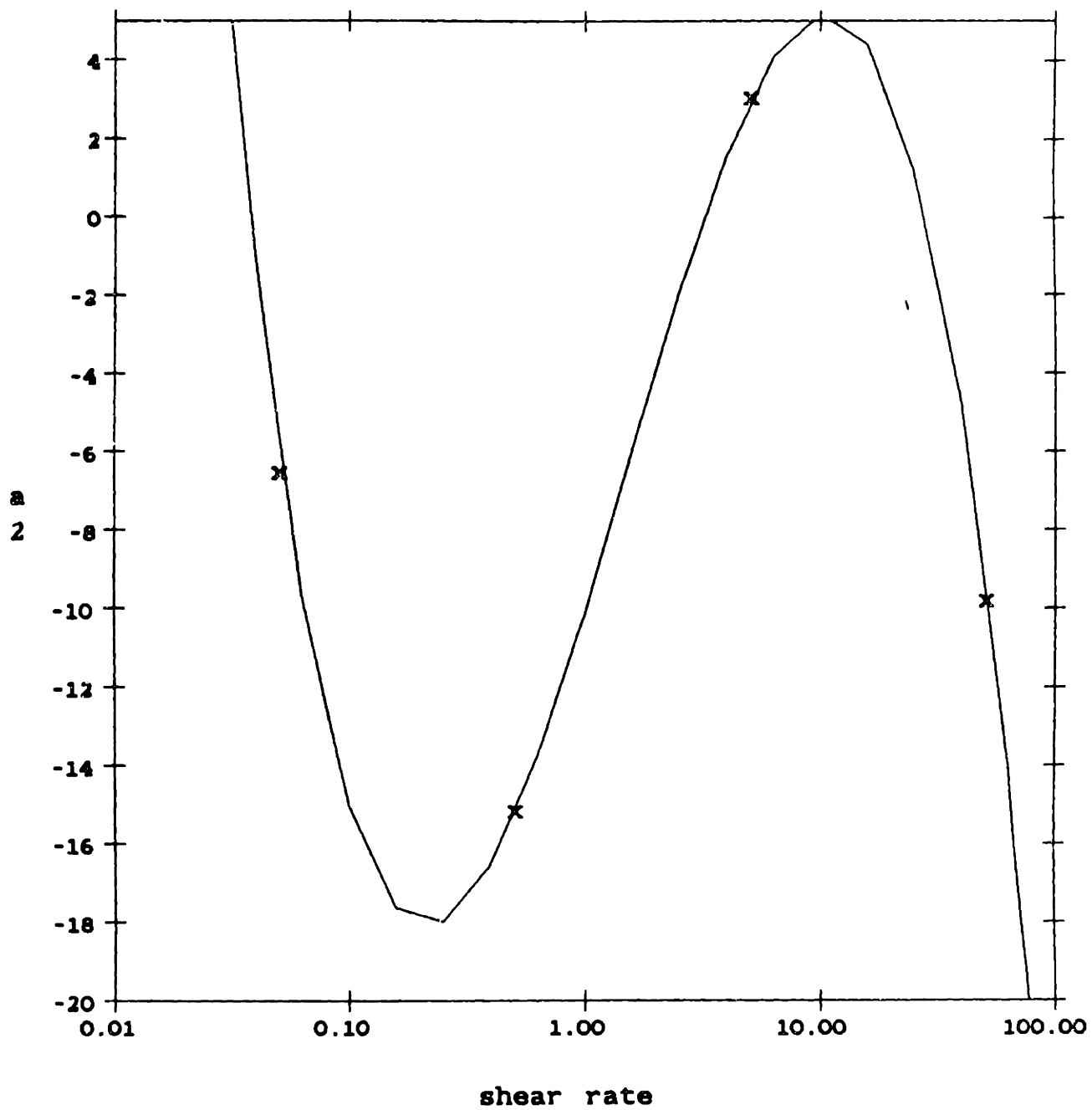
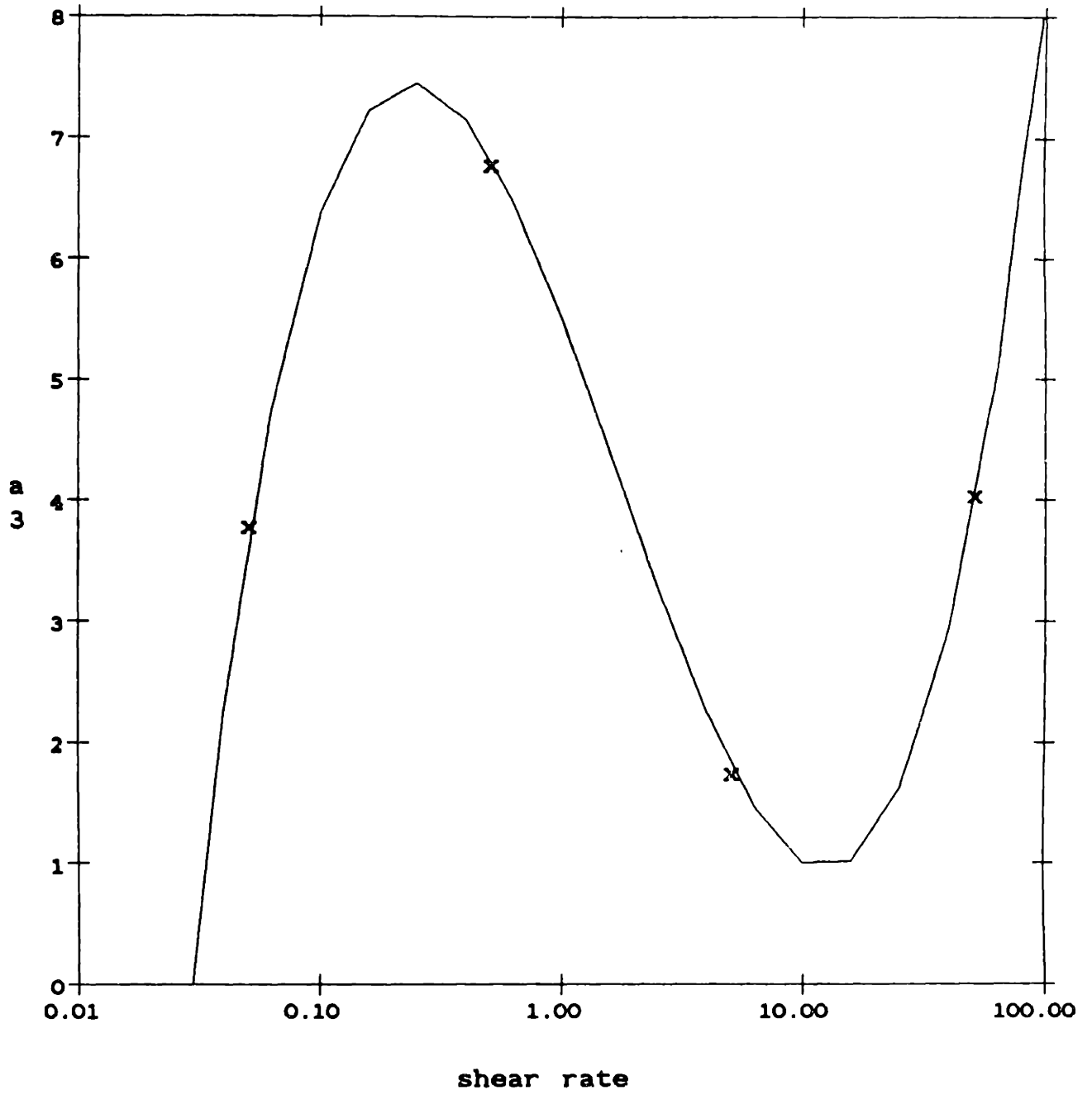


Figure 4-56(c): Plot for  $a_2$  in Ringer's solution.

Chien's  $rs \ a_3 \times 10^5$



**Figure 4-56(d):** Plot for  $a_3$  in Ringer's solution.

Chien's  $rs a_4 \times 10^7$

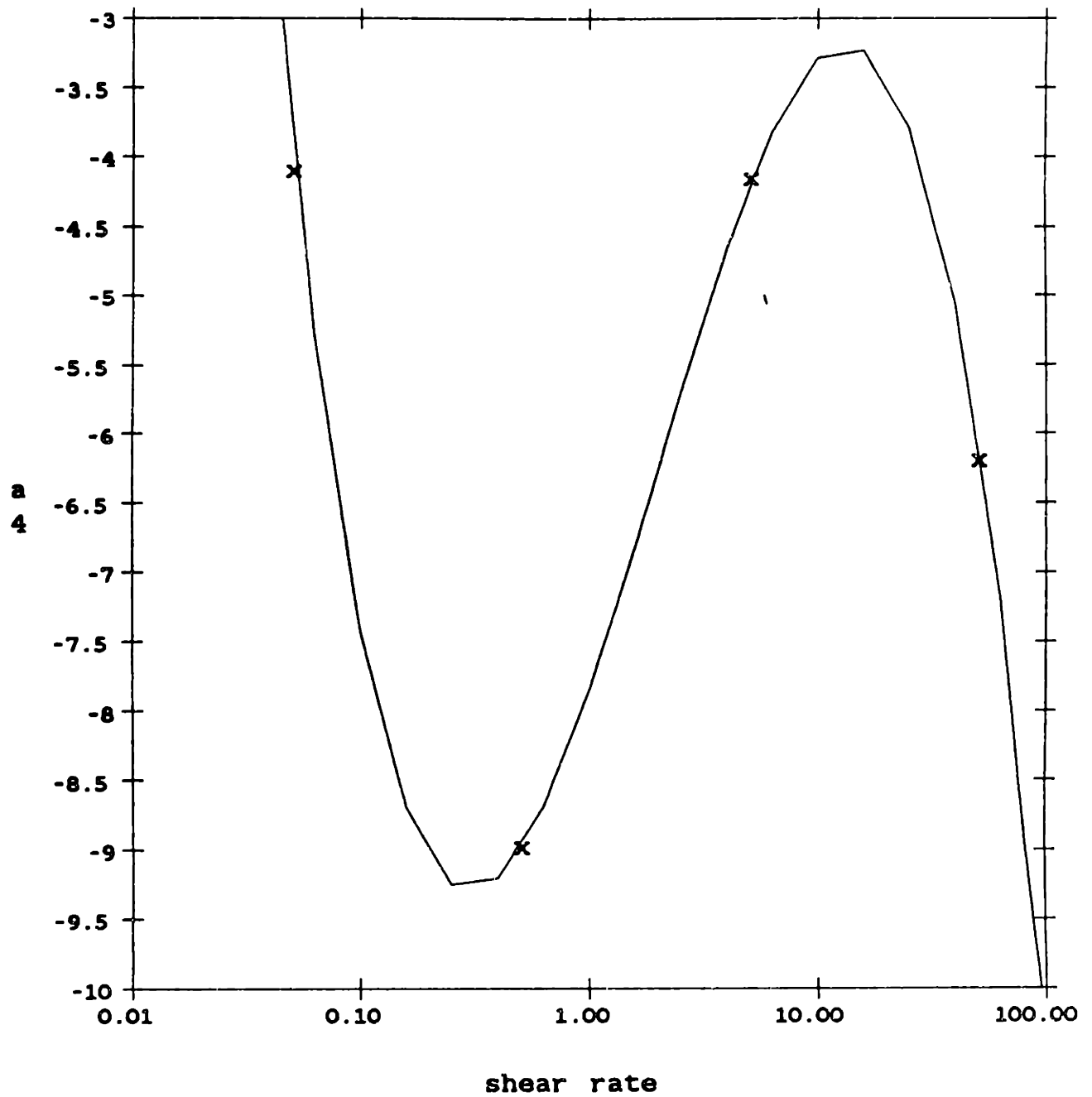


Figure 4-56(e): Plot for  $a_4$  in Ringer's solution.

Chien's  $rs$   $a_5 \times 10^9$

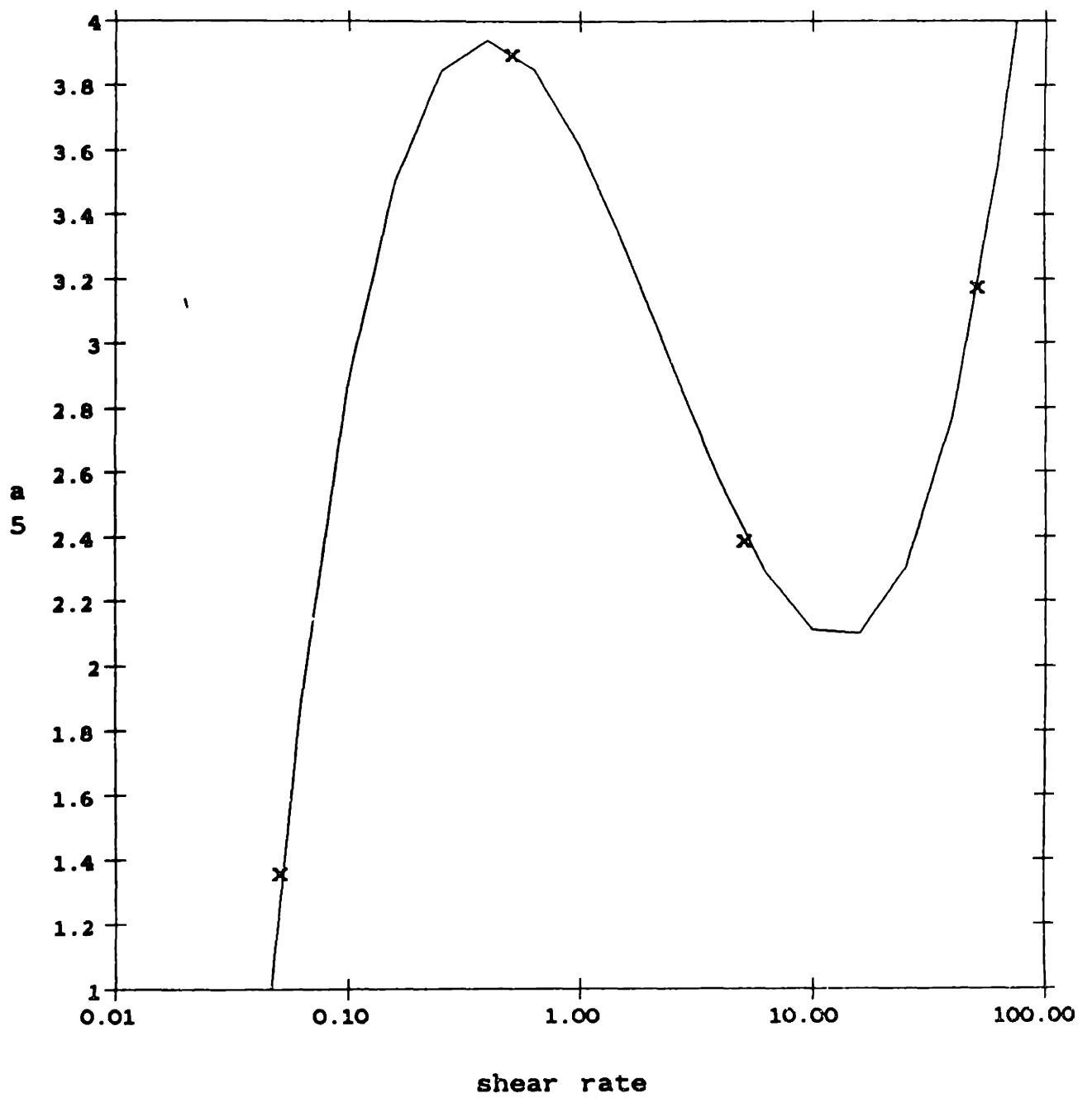


Figure 4-56(f): Plot for  $a_5$  in Ringer's solution.

	j=0	j=1	j=2	j=3
i=0	-3.678e-01			
i=1	3.619e-02	-1.025e-02	-6.384e-04	9.009e-04
i=2	-1.012e-03	8.572e-04	9.811e-05	-7.906e-05
i=3	5.514e-05	-2.275e-05	-3.457e-06	2.097e-06
i=4	-7.856e-07	2.095e-07	4.721e-08	-2.264e-08
i=5	3.614e-09	-6.247e-10	2.119e-10	8.669e-11

**Table 4-III:** Constants for refit saline Chien equation (4.2).

Relative viscosity vs. Shear rate  
RBC in Eagle's solution, H=96

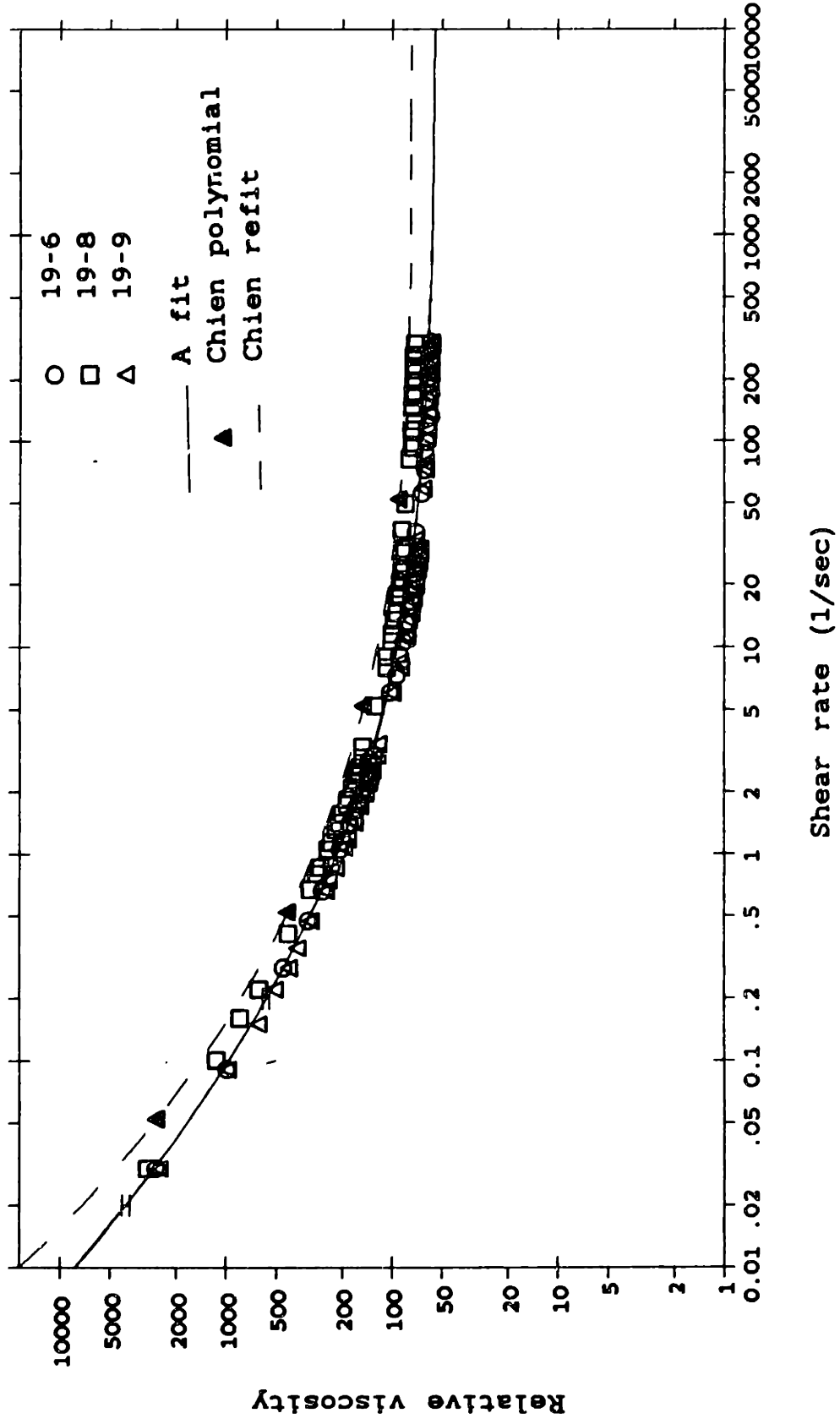
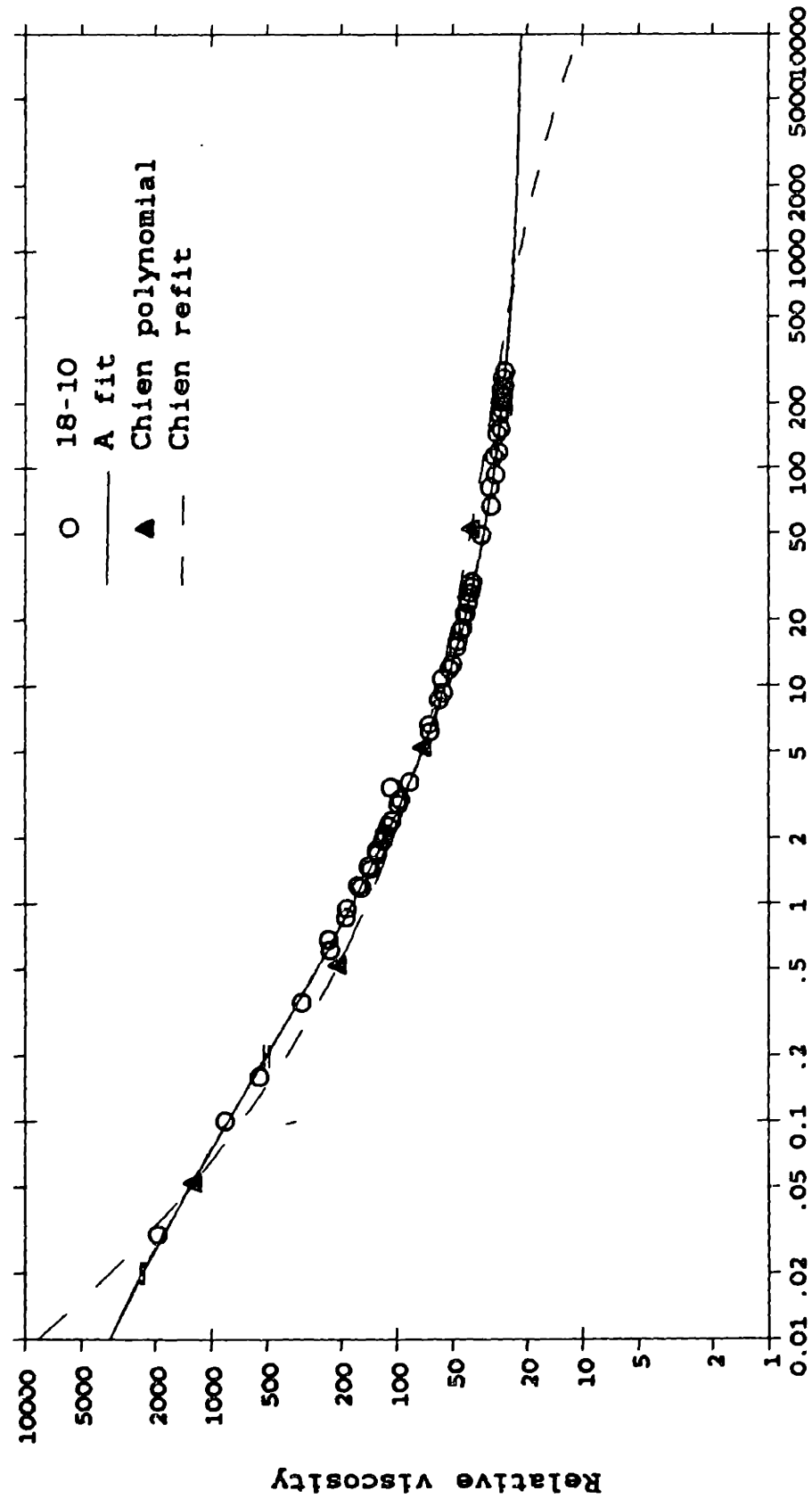


Figure 4-57: Plot of refit saline Chien equation versus shear rate at H=96.

Relative viscosity vs. Shear rate  
RBC in Eagle's solution, H=89



Shear rate (1/sec)

Figure 4-58: Plot of refit saline Chien equation versus shear rate at H=89.



Relative viscosity vs. Shear rate  
RBC in Eagle's solution, H=48

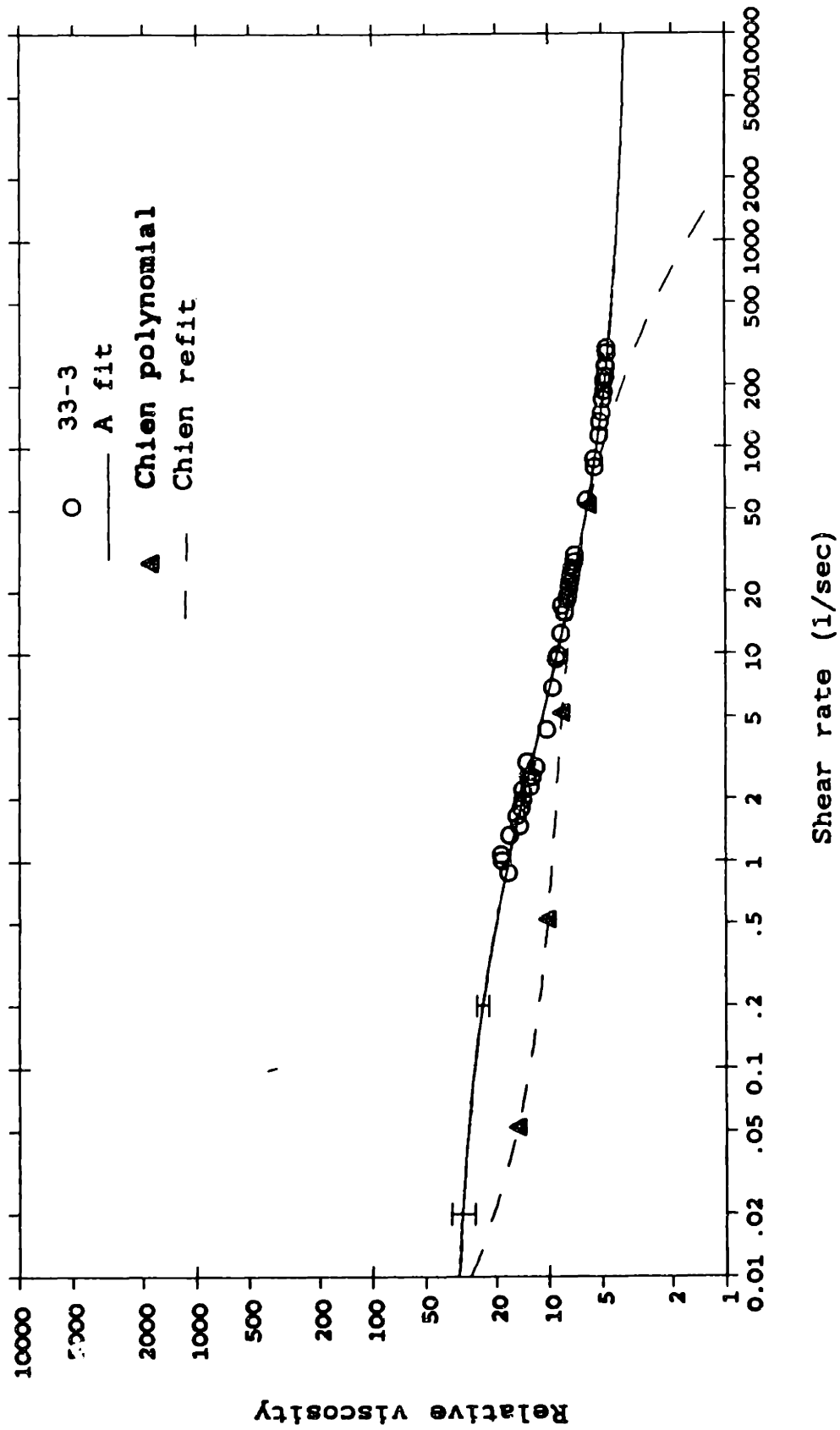


Figure 4-59: Plot of refit saline Chien equation versus shear rate at H=48.

becomes less adequate, veering downward at high shear rates (see Figures 4-58 and 4-59). The constants for plasma suspensions actually suggest asymptotic values at high shear rates (see Figure 4-60(a-f)). The following equations are plotted in Figure 4-60(a-f) as providing a reasonable representation of the functional dependence of the constants:

$$a_0 = 0.1672 \quad (4.3)$$

$$a_1 = -\frac{8.513 \times 10^{-4}}{\gamma} + 0.02595 \quad (4.4)$$

$$a_2 = \frac{-6.243 \times 10^{-3}\gamma}{(0.8106 + \gamma)} + 6.153 \times 10^{-3} \quad (4.5)$$

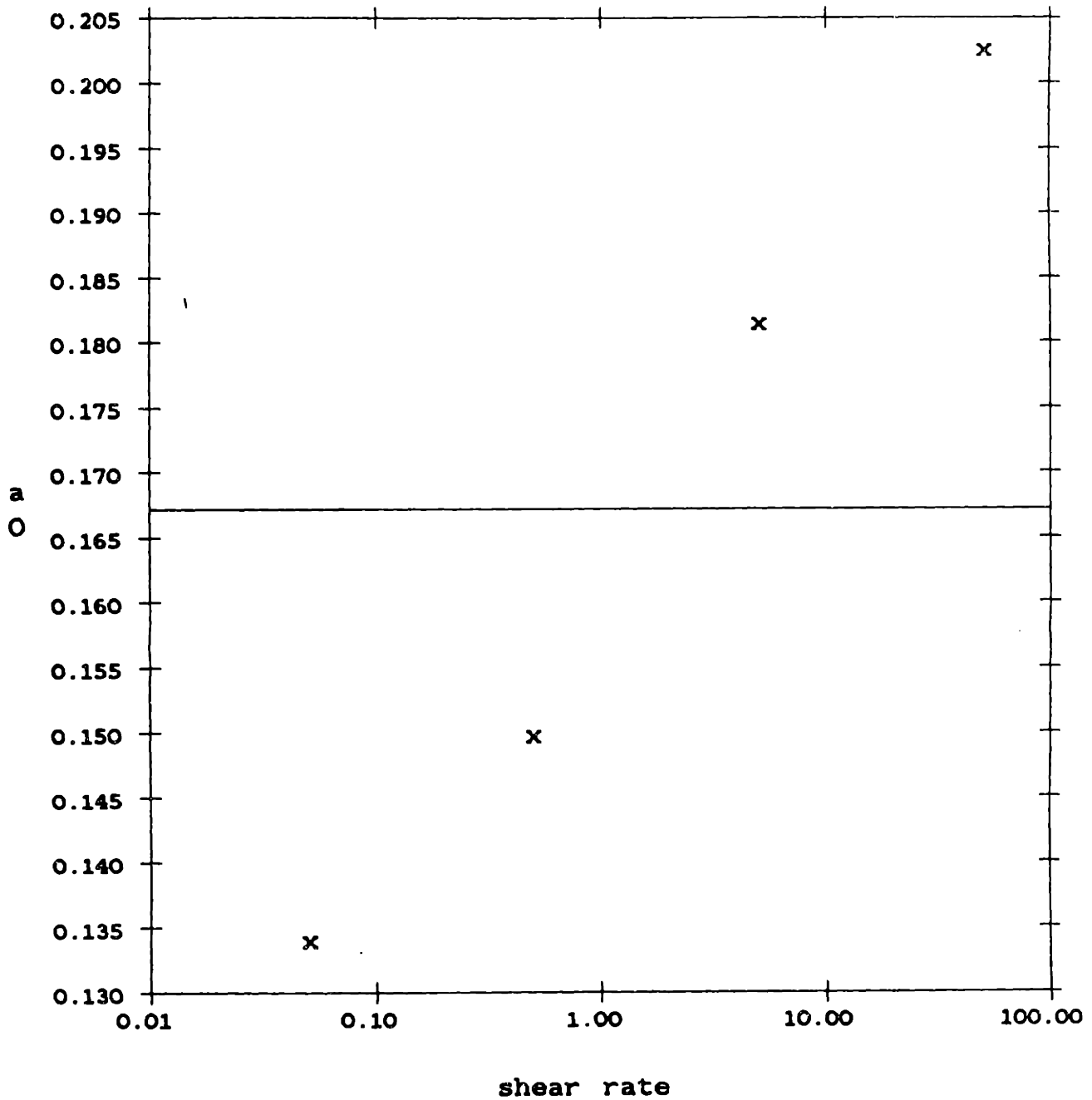
$$a_3 = 1.732 \times 10^{-4} \exp\left(\frac{-0.4886}{\gamma}\right) - 1.639 \times 10^{-4} \quad (4.6)$$

$$a_4 = \frac{1.238 \times 10^{-5}}{\ln(1000 \cdot \gamma)} - 1.326 \times 10^{-6} \quad (4.7)$$

$$a_5 = 7.539 \times 10^{-9} \exp\left(\frac{-0.5584}{\gamma}\right) - 6.009 \times 10^{-9} \quad (4.8)$$

These equations do not prove satisfactory, however, as it is not possible even to retrieve the original plots from them except at low hematocrits (see Figure 4-61). Once again the need is emphasized for a better representation of the data, in a more general form.

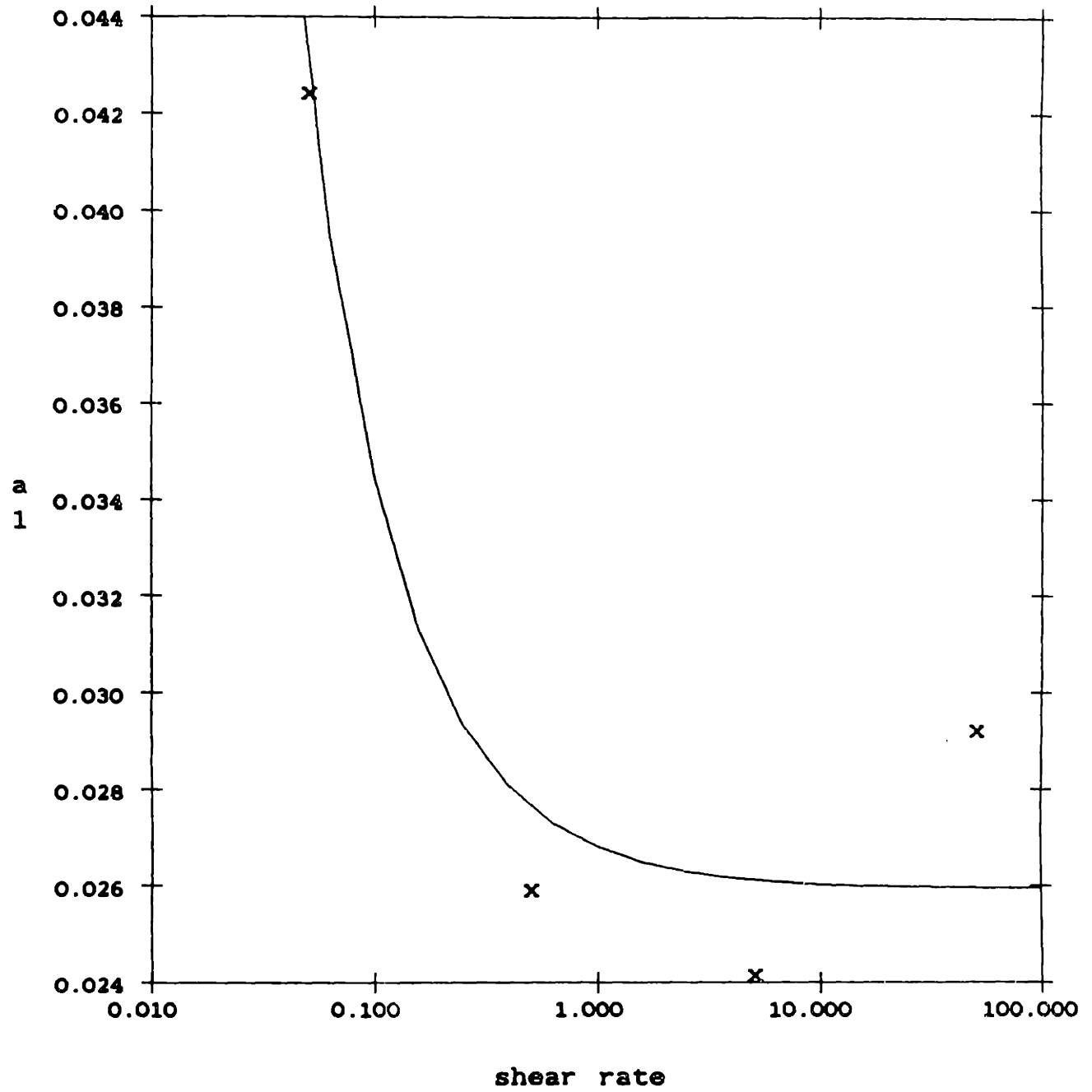
Chien's  $wb$   $a_0$



**Figure 4-60:** Plots of constants for Chien polynomial for plasma versus shear rate.

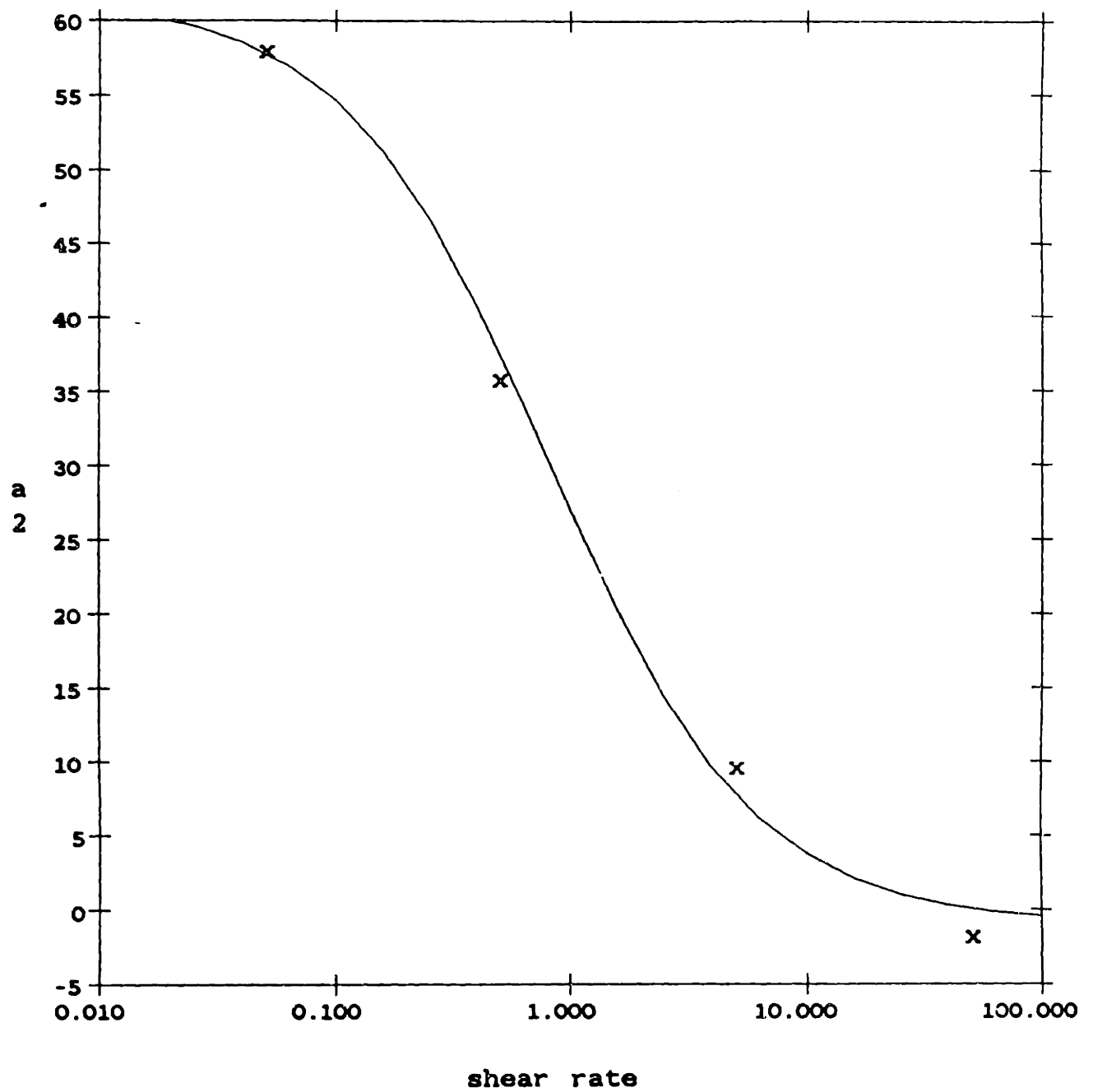
(a) Plot for  $a_0$  in plasma.

Chien's  $w_b$   $a_1$



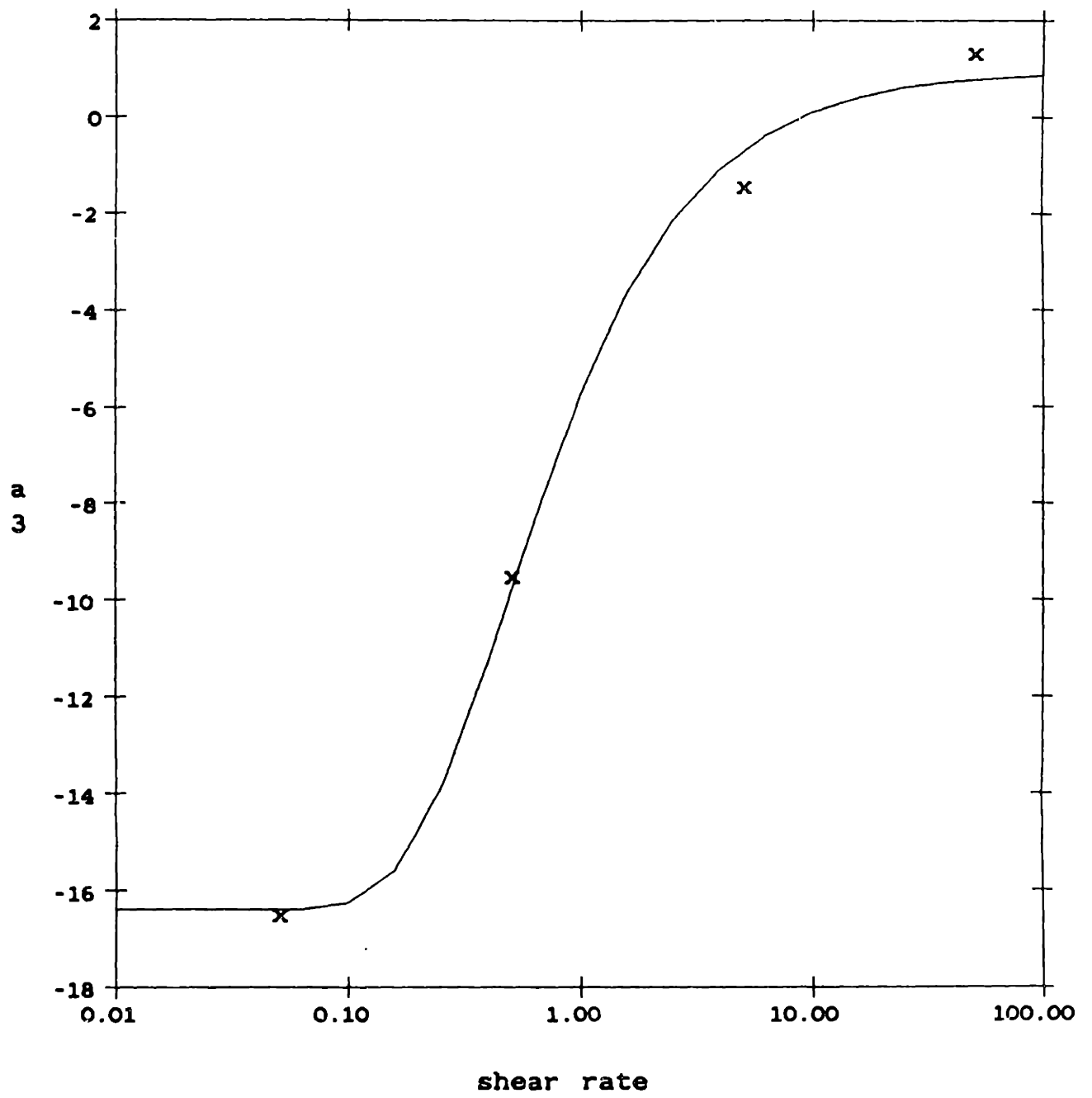
**Figure 4-60(b):** Plot for  $a_1$  in plasma.

Chien's  $wb a_2 \times 10^4$



**Figure 4-80(c):** Plot for  $a_2$  in plasma.

Chien's  $wb a_3 \times 10^5$



**Figure 4-60(d):** Plot for  $a_3$  in plasma.

Chien's  $wb a_4 \times 10^7$

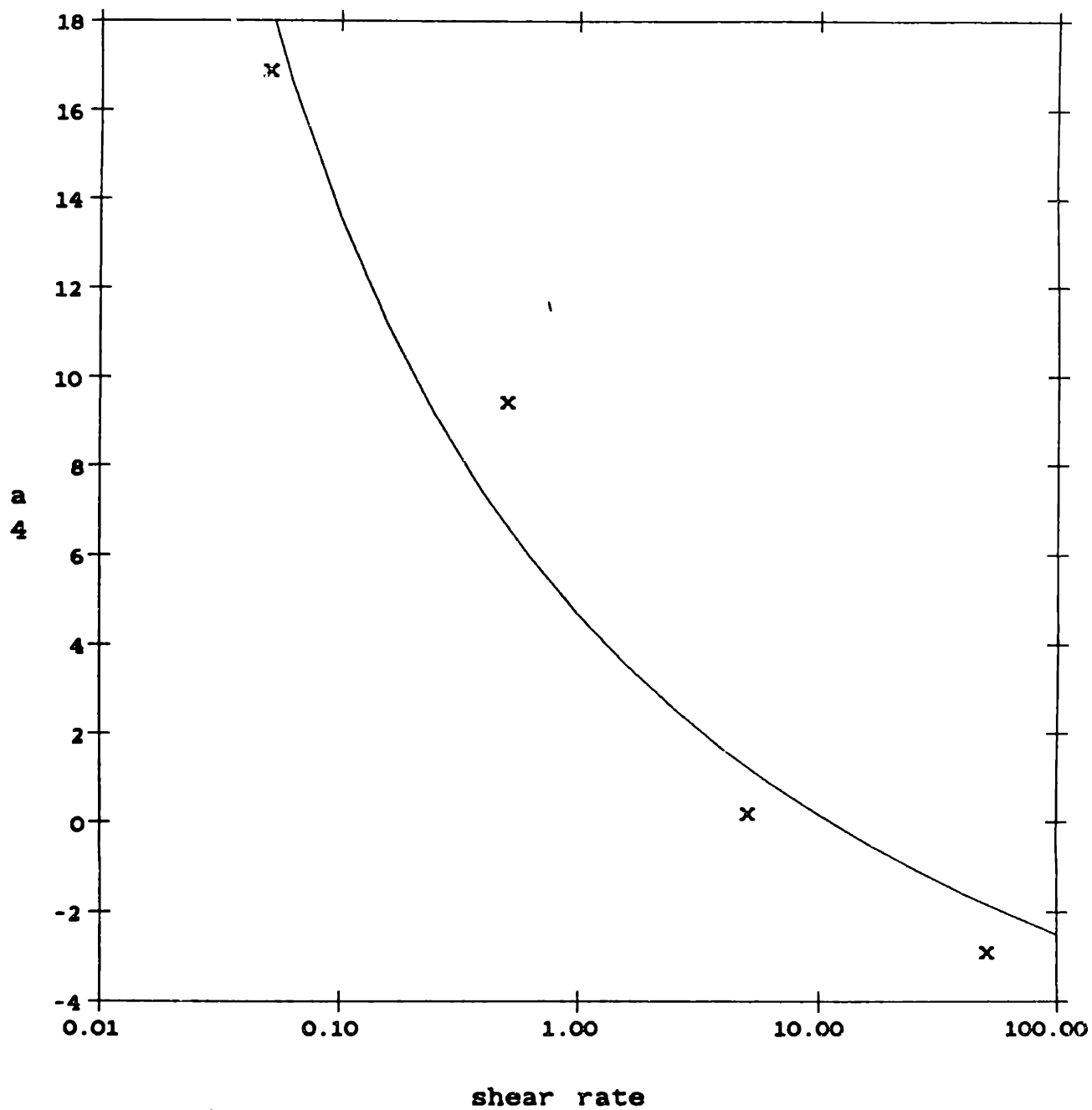
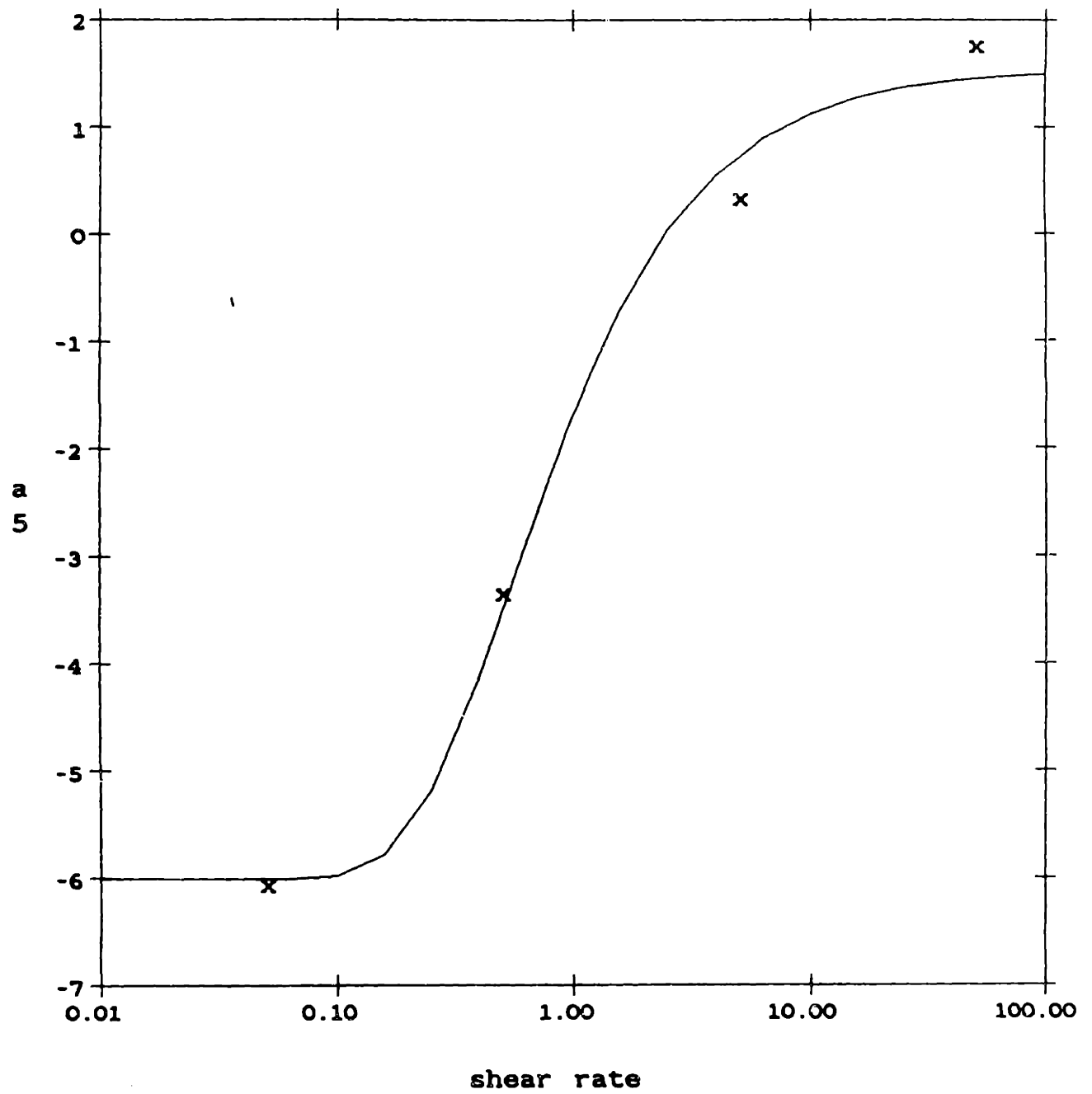


Figure 4-60(e): Plot for  $a_4$  in plasma.

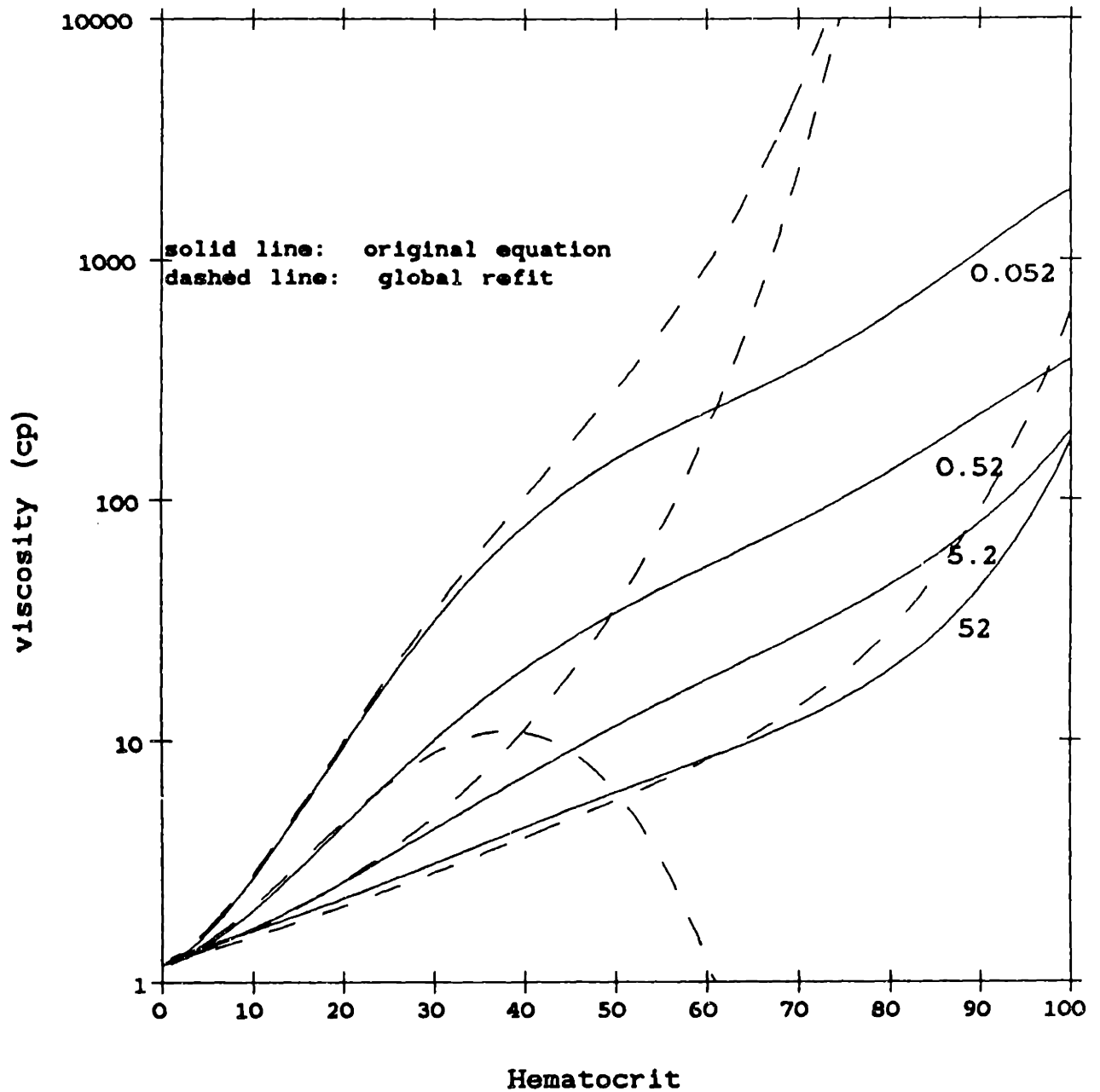
Chien's  $wb a_5 \times 10^9$



**Figure 4-60(f):** Plot for  $a_5$  in plasma.



Global refitting of Chien polynomials  
for whole blood



**Figure 4-61:** Plot of refit equations (4.3) to (4.8) for plasma suspensions.  
Numbers represent shear rates in  $\text{sec}^{-1}$ .

## Chapter 5

### DISCUSSION

#### 5.1 Curve fits to data

The data from the previous chapter was fitted to the three-parameter Quemada expression, equation (2.11), using the method and computer program given in Section 3.4 and Appendix D. Equation (2.11) was rearranged to the following form for ease in calculation:

$$\eta_r = \left[ 1 - \frac{c}{2} \left( k_\infty + \frac{k_0 - k_\infty}{1 + \sqrt{\gamma/\gamma_c}} \right) \right]^{-2} \quad (5.1)$$

This form of the equation minimizes the frequency of appearance of the shear rate term, whose fit, as will be seen, has the highest degree of uncertainty. It will be recalled that the term in parentheses is a general intrinsic viscosity  $k(c, \gamma)$ . The squared residuals were weighted by  $y_i^{-2}$ , where  $y_i$  is the predicted value of  $\eta_r$  from the equation, in order to assign equal relative weight to each point regardless of the magnitude of  $\eta_r$ . (When the term *relative weighting* is used henceforth, it will refer to a fit by this method. If there is additional weighting used, it will be so specified in each case.) In all cases the expression was able to provide a good representation of the data, and in almost every case the sum of squared weighted residuals using this equation was smaller than for any of several empirical equations attempted, the best of which were a fourth-order polynomial fit of  $\ln(\eta)$  to  $\ln(\gamma)$  and a fit of  $\ln(\eta)$  to the form  $b_1 + \frac{b_2}{\gamma b_3}$  where  $b_1$ ,  $b_2$  and  $b_3$  are constants.

The Quemada equation, when fitted to the low hematocrit plasma data from this work, often did not provide reasonable extrapolations to low shear rates where

data was unavailable. In order to provide a better behavior in the low hematocrit--low shear rate regime, the data points from the Chien *et al.*, [1966] equations were added to the plasma data where there was good agreement between the two sets. Therefore, for plasma suspension at H=50, 44, 34, 24, and 11 the fits are to this work plus the four points of Chien *et al.* at each hematocrit. In the case for Eagle's solution at H=7, the data was fit to a constant from which a value of  $k_{\infty}$  was calculated.

Statistical parameters for the goodness-of-fit of the equations are given in Table 5-I. The standard weighted error of regression is defined as  $\sqrt{\frac{E}{N-P}}$ , where E is the sum of squared weighted residuals as defined by equation (3.17), N is the number of data points and P is the number of parameters. It will be noted that the absolute magnitude of the correlation coefficients involving the critical shear rate are all very close to 1, suggesting a relative insensitivity of the fit to the value of  $\gamma_c$ . As will be recalled from the analytical discussion of the Quemada equation in Chapter 2,  $\gamma_c$  does not change the overall shape of the Quemada curve but instead acts as a scaling factor for the x-axis. Thus, for data that is relatively constant or that expresses asymptotic behavior at only one end, it will be more difficult to pinpoint the value of  $\gamma_c$ , as seen by the  $\rho_{23}$  values for plasma at low and high hematocrits ( $\rho_{13}$  values are practically constant for all cases because asymptotic approach to  $k_{\infty}$  is always seen). This uncertainty in  $\gamma_c$  will be expressed in relatively higher standard deviations for the value of its fit.

The high magnitudes of the correlation coefficients might also suggest an overspecification of the problem for any one curve. It is quite possible that a data at one hematocrit may be effectively represented by a two parameter equation, but three parameters are needed to model behavior over a large range of hematocrits.

The rheological parameters  $k_{\infty}$ ,  $k_0$ , and  $\gamma_c$  generated by the fits are given

Hematocrit	number of data points	std. wghtd. regress. err.	corr. coeff. 12	corr. coeff. 13	corr. coeff. 23
98.4	60	7.651e-02	0.7437	-0.9413	-0.9225
97.2	55	3.342e-02	0.6890	-0.9251	-0.9075
96.8	60	2.879e-02	0.7147	-0.9379	-0.9069
96.0	170	1.013e-01	0.6957	-0.9301	-0.9065
95.0	129	4.210e-02	0.7297	-0.9573	-0.8901
94.0	112	3.153e-02	0.6876	-0.9382	-0.8893
93.0	181	2.096e-01	0.7366	-0.9544	-0.8986
92.0	161	1.757e-01	0.7166	-0.9506	-0.8905
91.0	122	4.011e-02	0.6809	-0.9410	-0.8776
90.0	118	1.212e-01	0.7102	-0.9610	-0.8678
89.0	54	3.829e-02	0.7056	-0.9721	-0.8443
87.0	120	7.184e-02	0.6203	-0.9541	-0.8067
86.0	57	5.163e-02	0.6543	-0.9385	-0.8587
82.0	47	5.303e-02	0.6715	-0.9692	-0.8235
72.0	46	5.949e-02	0.6517	-0.9658	-0.8091
70.0	107	1.789e-01	0.6385	-0.9513	-0.8261
68.0	48	9.703e-02	0.6850	-0.9142	-0.9127
62.0	53	1.638e-01	0.7176	-0.9397	-0.9046
59.0	57	3.020e-02	0.6533	-0.9307	-0.8642
56.0	54	2.435e-02	0.7450	-0.9321	-0.9233
48.0	46	4.270e-02	0.7650	-0.9269	-0.9406
42.0	30	1.689e-02	0.7872	-0.9131	-0.9644
25.0	32	6.469e-02	0.9061	-0.9740	-0.9745
7.0	14	4.355e-03			
Hematocrit	number of data points	std. wghtd. regress. err.	corr. coeff. 12	corr. coeff. 13	corr. coeff. 23
98.1	65	2.584e-02	0.6980	-0.9283	-0.9085
97.6	63	6.425e-02	0.6985	-0.9423	-0.8919
97.3	59	4.792e-02	0.7018	-0.9659	-0.8550
97.0	60	1.471e-02	0.7015	-0.9463	-0.8859
96.0	122	9.660e-02	0.6694	-0.9489	-0.8604
95.0	63	2.514e-02	0.7077	-0.9673	-0.8541
94.0	62	2.112e-02	0.6850	-0.9572	-0.8545
93.0	62	2.943e-02	0.6732	-0.9492	-0.8590
92.0	122	3.445e-02	0.6610	-0.9590	-0.8322
91.0	124	8.631e-02	0.7064	-0.9699	-0.8479
90.0	62	3.614e-02	0.6788	-0.9642	-0.8361
89.0	124	1.004e-01	0.6931	-0.9694	-0.8378
86.0	65	3.351e-02	0.6708	-0.9573	-0.8418
83.0	65	3.515e-02	0.6661	-0.9638	-0.8251
82.0	68	3.697e-02	0.6518	-0.9631	-0.8161
78.0	64	5.593e-02	0.6997	-0.9781	-0.8225
75.0	59	2.728e-02	0.6263	-0.9662	-0.7844
66.0	65	5.593e-02	0.5921	-0.9614	-0.7654
59.0	54	5.786e-02	0.6677	-0.9543	-0.8420
50.0	59	1.437e-02	0.7276	-0.9382	-0.9028
44.0	45	1.193e-02	0.7828	-0.9162	-0.9608
34.0	44	4.518e-02	0.7449	-0.8327	-0.9880
24.0	31	6.291e-02	0.6682	-0.7513	-0.9886
11.0	31	3.329e-03	0.6812	-0.7019	-0.9982

**Table 5-I:** Standard weighted error of regression and correlation coefficients for Quemada fits to data.

*Top*, values for saline suspensions. *Bottom*, values for plasma suspensions.

For correlation coefficients: 1= $k_{\infty}$ , 2= $k_0$ , 3= $\gamma_c$

Hematocrit	k-inf. A fit	st. dev. k-inf.	k-O A fit	st. dev. k-O	gam.-crit. A fit	st. dev. gam. crit.
98.4	1.849	1.081e-02	2.039	2.085e-03	0.7572	1.873e-01
97.2	1.819	7.028e-03	2.063	1.582e-03	0.8711	1.123e-01
96.8	1.774	3.090e-03	2.069	6.329e-04	1.2130	5.526e-02
96.0	1.800	7.707e-03	2.087	1.346e-03	0.9387	1.064e-01
95.0	1.764	8.821e-03	2.100	8.093e-04	2.6770	2.519e-01
94.0	1.794	8.269e-03	2.126	1.321e-03	1.3300	1.337e-01
93.0	1.817	1.881e-02	2.146	2.033e-03	2.3230	4.980e-01
92.0	1.815	3.395e-02	2.171	4.001e-03	1.9850	7.161e-01
91.0	1.790	8.140e-03	2.184	1.148e-03	1.7760	1.442e-01
90.0	1.761	6.705e-02	2.206	5.263e-03	3.3170	1.715e+00
89.0	1.752	1.739e-02	2.231	7.749e-04	4.4220	4.876e-01
87.0	1.806	2.802e-02	2.268	1.940e-03	2.8530	5.700e-01
86.0	1.807	3.368e-02	2.289	4.791e-03	2.0640	5.662e-01
82.0	1.829	6.995e-02	2.412	3.550e-03	5.1660	1.898e+00
72.0	1.884	9.638e-02	2.709	5.382e-03	6.1900	2.276e+00
70.0	1.969	8.117e-02	2.802	8.218e-03	3.3140	1.146e+00
68.0	2.005	5.583e-02	2.874	2.499e-02	4.9410	1.767e+00
62.0	2.014	1.505e-01	3.155	4.397e-02	7.5150	4.714e+00
59.0	1.853	3.547e-02	2.983	7.337e-03	23.4500	3.234e+00
56.0	1.896	2.604e-02	3.004	9.297e-03	22.0800	2.836e+00
48.0	1.947	7.059e-02	3.476	3.734e-02	13.7800	4.065e+00
42.0	2.099	2.987e-02	3.730	5.228e-02	14.5500	2.937e+00
25.0	2.906	4.546e-01	3.787	4.281e-01	52.8800	2.248e+02
7.0	6.605	1.489e-01				

Hematocrit	k-inf. A fit	st. dev. k-inf.	k-O A fit	st. dev. k-O	gam.-crit. A fit	st. dev. gam. crit.
98.1	1.728	3.730e-03	2.045	8.178e-04	0.91790	4.823e-02
97.6	1.711	9.668e-03	2.050	1.387e-03	1.48300	1.663e-01
97.3	1.700	8.846e-03	2.048	5.806e-04	2.56200	2.110e-01
97.0	1.698	3.011e-03	2.056	4.183e-04	1.59300	5.226e-02
96.0	1.724	1.395e-02	2.076	1.456e-03	1.80500	2.566e-01
95.0	1.678	1.495e-02	2.091	1.084e-03	3.26600	3.905e-01
94.0	1.651	7.985e-03	2.108	7.744e-04	2.55200	1.575e-01
93.0	1.699	1.909e-02	2.130	2.149e-03	2.33000	3.797e-01
92.0	1.684	1.435e-02	2.154	1.012e-03	2.21600	2.241e-01
91.0	1.649	4.420e-02	2.181	2.759e-03	3.11400	8.285e-01
90.0	1.682	3.542e-02	2.195	2.781e-03	3.29200	7.587e-01
89.0	1.630	5.729e-02	2.221	3.682e-03	3.88700	1.208e+00
86.0	1.674	2.977e-02	2.284	2.885e-03	3.75000	6.440e-01
83.0	1.696	4.127e-02	2.370	2.621e-03	3.57000	7.088e-01
82.0	1.625	5.459e-02	2.391	3.734e-03	4.38000	1.013e+00
78.0	1.722	8.138e-02	2.532	3.026e-03	5.30400	1.556e+00
75.0	1.688	3.757e-02	2.577	2.343e-03	6.30400	8.458e-01
66.0	1.653	1.266e-01	2.888	8.337e-03	6.99900	2.291e+00
59.0	1.695	1.245e-01	3.202	1.589e-02	6.80300	2.108e+00
50.0	2.018	1.275e-01	3.865	2.414e-02	2.07700	6.568e-01
44.0	2.181	7.282e-02	4.484	2.987e-02	0.95380	1.924e-01
34.0	2.219	7.624e-02	6.092	1.265e-01	0.30890	8.253e-02
24.0	2.150	9.283e-02	7.913	2.719e-01	0.23780	7.731e-02
11.0	2.023	1.434e-01	17.410	1.087e+01	0.01124	2.345e-02

**Table 5-II:** Quemada parameters for data of present study.

*Top*, values for saline suspensions. *Bottom*, values for plasma suspensions.

along with their standard deviations in Table 5-II for Eagle's solution and plasma suspensions. The curves generated by these parameters are plotted as solid lines designated as "Fit A" in the results graphs of Chapter 4. The hematocrit dependence of the parameters, to be discussed in the sections to follow, are used to generate the dashed lines designated as "Fit B" in Chapter 4. For comparison, fits were also made of the Quemada equation to the polynomials of Chien *et al.* for both saline and plasma at intervals of five hematocrit units. The values obtained from these fits are shown in Table 5-III.

As discussed in Chapter 2, the potential for singular points of infinite viscosity exists as the cell volume fraction approaches unity, dependent on the magnitude of  $\frac{c}{2}k_0$ . Some of the fits obtained had singular shear rates, and these are documented in Table 5-IV, although in all of these cases they occur well out of the limits of the data.

The standard deviations in Tables 5-II and 5-III are notable in that those for the critical shear rate are large relative to the value itself. At low hematocrits, the standard deviation may be several times the parameter value for  $\gamma_c$  and it is also larger for  $k_0$ . These observations underscore the fact that data at very low shear rates is needed to accurately fit these two constants.

## 5.2 Dependence of $k_\infty$ on hematocrit and plasma proteins

The A fit values for  $k_\infty$  range from 1.752 to 6.605 for saline and 1.625 to 2.219 for plasma. On the whole, these values are compatible with the values around 1.8 fit by Quemada to other data (see Table 2-II, page 42), but lower because the data from this work extended to higher shear rates. The saline values for the A fit are plotted versus hematocrit along with the values obtained from the Ringer solution Chien polynomials in Figure 5-1. Since the data of Chien *et al.* does not

Hematocrit	k-infinity	st. dev. k-infinity	k-0	st. dev. k-0	gamma. crit.	st. dev. gamma. crit.
100	1.829	3.85e-03	2.005	2.74e-03	0.763400	1.37e-01
95	1.834	3.38e-03	2.112	1.84e-03	1.049000	9.23e-02
90	1.836	1.33e-02	2.227	6.17e-03	1.321000	3.05e-01
85	1.847	2.50e-02	2.348	1.07e-02	1.534000	5.01e-01
80	1.872	3.77e-02	2.474	1.62e-02	1.672000	6.83e-01
75	1.915	5.15e-02	2.604	2.34e-02	1.742000	8.68e-01
70	1.975	6.72e-02	2.736	3.34e-02	1.770000	1.08e+00
65	2.098	8.61e-02	2.865	4.75e-02	1.779000	1.35e+00
60	2.131	1.09e-01	2.985	6.74e-02	1.785000	1.72e+00
55	2.219	1.38e-01	3.091	9.48e-02	1.786000	2.23e+00
50	2.313	1.71e-01	3.175	1.31e-01	1.745000	2.87e+00
45	2.413	2.02e-01	3.234	1.80e-01	1.545000	3.40e+00
40	2.525	2.11e-01	3.280	2.57e-01	0.972800	2.90e+00
35	2.653	1.69e-01	3.389	5.33e-01	0.219800	9.26e-01
30	2.750	1.10e-01	5.054	1.82e+01	0.002042	4.21e-02
25	2.876	7.05e-02				
20	2.867	6.61e-02				
15	2.884	9.72e-02				
10	2.963	1.71e-01				
5	3.162	2.90e-01				

Hematocrit	k-infinity	st. dev. k-infinity	k-0	st. dev. k-0	gamma. crit.	st. dev. gamma. crit.
100	1.822	8.00e-03	2.054	4.00e-02	0.079000	6.70e-02
95	1.818	1.00e-03	2.110	1.00e-03	0.656100	2.44e-02
90	1.794	1.60e-02	2.210	7.00e-03	1.725000	4.36e-01
85	1.764	3.70e-02	2.328	1.00e-02	2.897000	1.05e+00
80	1.741	5.80e-02	2.463	1.40e-02	3.745000	1.59e+00
75	1.735	7.70e-02	2.618	1.80e-02	4.086000	1.85e+00
70	1.747	9.00e-02	2.798	2.00e-02	3.993000	1.76e+00
65	1.775	9.30e-02	3.009	2.30e-02	3.647000	1.46e+00
60	1.813	8.70e-02	3.257	2.30e-02	3.203000	1.05e+00
55	1.856	7.10e-02	3.551	2.10e-02	2.746000	6.53e-01
50	1.898	4.60e-02	3.902	1.50e-02	2.307000	3.12e-01
45	1.938	1.10e-02	4.324	4.00e-03	1.891000	5.70e-02
40	1.974	3.00e-02	4.835	1.40e-02	1.494000	1.06e-01
35	2.010	7.50e-02	5.458	4.50e-02	1.122000	1.85e-01
30	2.052	1.21e-01	6.222	9.78e-02	0.790000	1.99e-01
25	2.111	1.63e-01	7.156	1.94e-01	0.508900	1.73e-01
20	2.205	1.99e-01	8.286	3.89e-01	0.289600	1.30e-01
15	2.365	2.28e-01	9.676	8.81e-01	0.013150	8.47e-02
10	2.645	2.51e-01	12.304	3.64e+00	0.028990	3.92e-02
5	2.982	4.89e-01	31.373	3.01e+02	0.000528	1.39e-02

Table 5-III: Quemada parameters for data of Chien *et al.*, [1966].

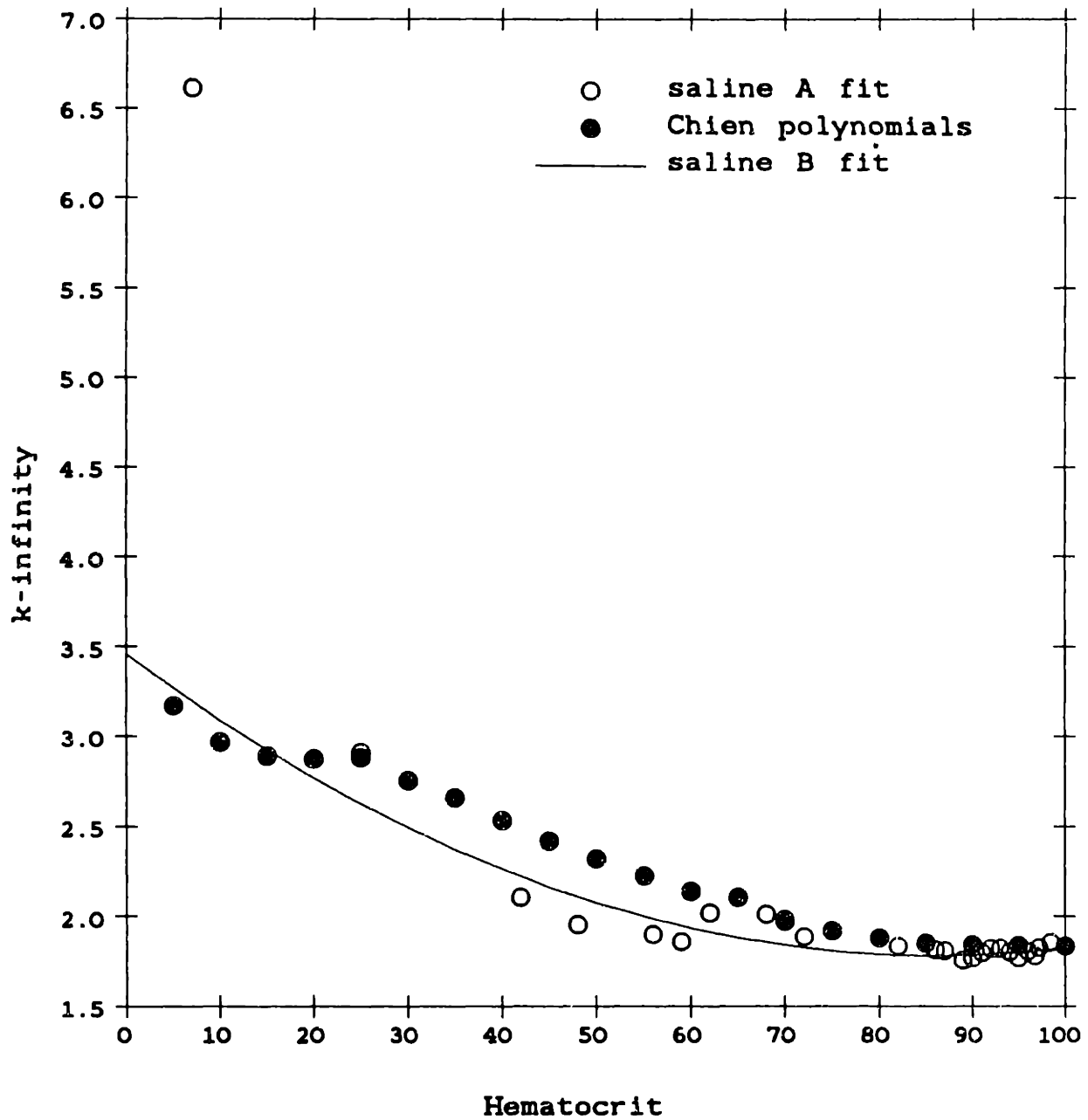
Top, values for saline suspensions. Bottom, values for plasma suspensions.

Fit	Suspension	H	$\gamma(\text{sec}^{-1})$
A	saline	98.4	$1.25 \times 10^{-3}$
A	saline	97.2	$5.70 \times 10^{-4}$
A	saline	96.8	$9.75 \times 10^{-5}$
A	saline	96	$1.67 \times 10^{-4}$
A	plasma	98.1	$3.73 \times 10^{-4}$
A	plasma	97.6	$8.71 \times 10^{-6}$
B	saline	98.4	$2.68 \times 10^{-3}$
B	saline	97.2	$5.17 \times 10^{-4}$
B	saline	96.8	$2.33 \times 10^{-4}$
B	saline	96	$7.88 \times 10^{-6}$

**Table 5-IV:** Singular shear rates of curve fits.



k-infinity vs. hematocrit  
saline suspensions



**Figure 5-1:** Plot of  $k_{\infty}$  versus hematocrit for saline suspensions.

extend to high shear rates, it is to be expected that the derived  $k_{\infty}$  values from that work will be of higher magnitude than those from the A fit. Similar behavior is shown by the  $k_{\infty}$  values for plasma suspensions in Figure 5-2.

The B fits in Figures 5-1 and 5-2 represent fits of the open symbols to function of cell concentration. For both saline and plasma suspensions the fits are to second-order polynomials:

$$k_{\infty} = \sum_{i=1}^3 k_{\infty i} \cdot c^{(i-1)} \quad (5.2)$$

where  $k_{\infty i}$  are constants and  $c$  is the volume concentration of cells. The values for the best fits of equation (5.2) and their standard deviations are, for saline,

$$\begin{aligned} k_{\infty 1} &= 3.458 \pm 0.251 \\ k_{\infty 2} &= -3.894 \pm 0.730 \\ k_{\infty 3} &= 2.260 \pm 0.500 \end{aligned}$$

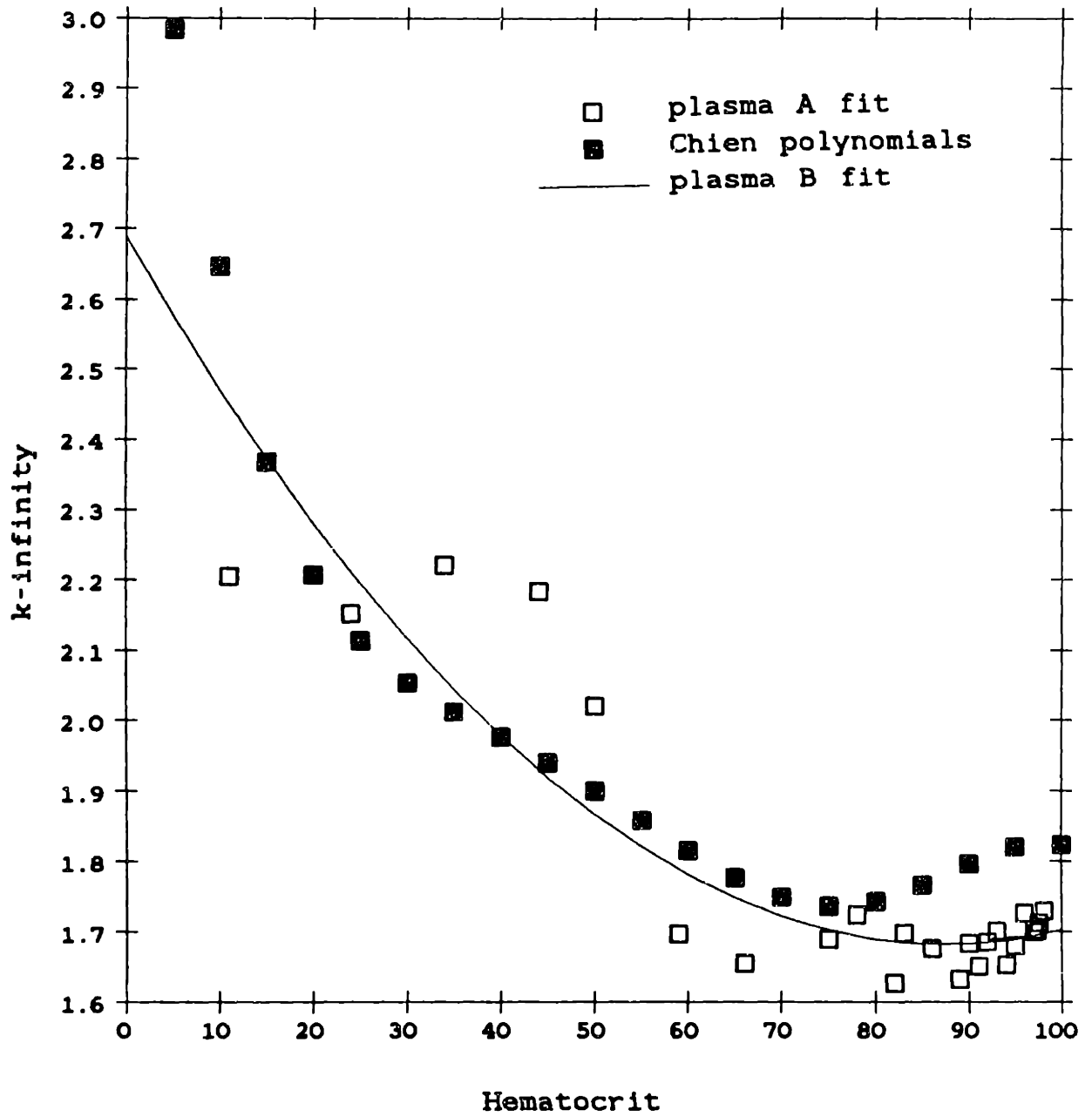
and for plasma,

$$\begin{aligned} k_{\infty 1} &= 2.688 \pm 0.117 \\ k_{\infty 2} &= -2.300 \pm 0.433 \\ k_{\infty 3} &= 1.314 \pm 0.353. \end{aligned}$$

Both fits were relatively weighted, and the values at hematocrits of 95 and above were triple weighted to provide better correlation at high hematocrit. For the same reason, the saline  $k_{\infty}$  value at  $H=7$  was not included in the fit. The importance of accuracy at high hematocrit is seen by the sensitivity of the asymptotic relative viscosity at high shear rates, where  $k=k_{\infty}$ . The sensitivity is much greater at higher cell volume fractions because the magnitude of  $k_{\infty}$  does not change a great deal with  $H$ . For example, at  $c=0.95$ ,  $k_{\infty} \approx 1.7$ , and  $\frac{\partial \eta_r}{\partial k} = \frac{c}{(1-k \cdot c/2)^3} = 133$ . At  $c=0.45$ ,  $k_{\infty} \approx 2.0$ , and  $\frac{\partial \eta_r}{\partial k} = 2.7$ .

The saline and plasma fits are plotted together on Figure 5-3. Plasma values are consistently lower than those for saline. This is not due to cell-protein

k-infinity vs. hematocrit  
plasma suspensions



**Figure 5-2:** Plot of  $k_{\infty}$  versus hematocrit for plasma suspensions.

k-infinity vs. hematocrit  
saline vs. plasma suspensions

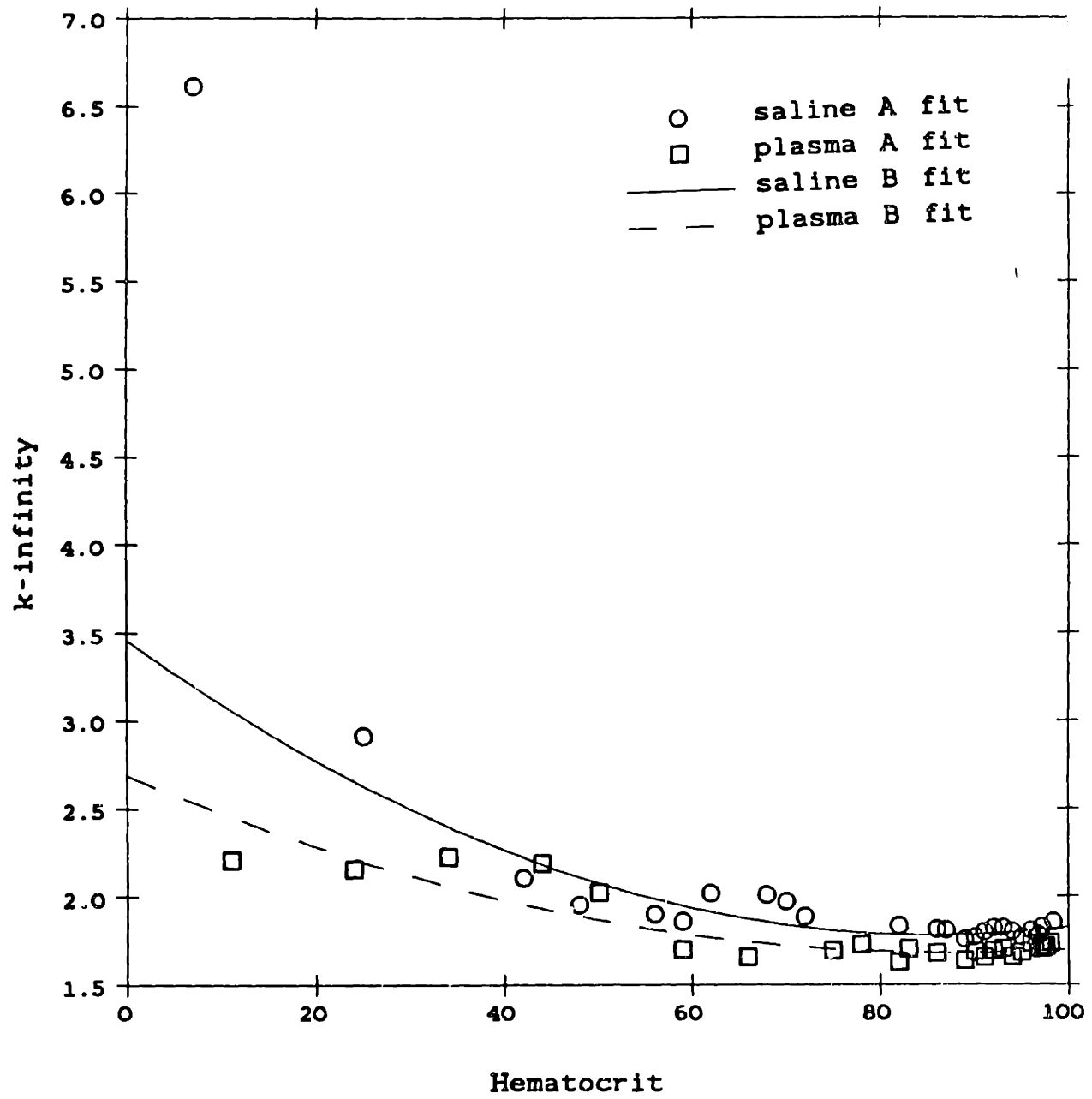


Figure 5-3: Comparison of  $k_{\infty}$  values for saline and plasma suspensions.

interactions, because at high shear rates aggregation effects are not important, but results directly from the differences in  $\eta_f$  and the emulsive properties of the erythrocytes. For rigid particles,  $k_\infty$  is independent of  $\eta_f$ ; but for emulsions, the intrinsic viscosity will clearly have a functional dependence on the internal viscosity, and if  $\eta_i$  is lowered in more viscous media, then  $k_\infty$  will also be lowered. Quemada [1978b] obtained no difference in  $k_\infty$  from data [Chien, 1970] for erythrocytes in plasma and Ringer's albumin solution (see Figure 2-3 and Table 2-II), but these solutions all had the same  $\eta_f$ , which is not the case here.

As noted in Chapter 4, as  $c \rightarrow 1$ , the absolute viscosities of plasma and saline solutions become equal, and, barring the apparently negligible differences in cell properties from saline washing, the definition of the relative viscosity leads to contrived distinctions. These distinctions are manifested in the difference in  $k_\infty$  values at high hematocrits, as they are in the values for  $k_0$  at high hematocrits. Thus  $k_{\infty, \text{saline}}$  and  $k_{\infty, \text{plasma}}$  can be related as a function of  $\eta_f$ . Since  $\eta_{\text{saline}} = \eta_{\text{plasma}}$  at high H:

$$\frac{\eta_{f, \text{plasma}}}{\eta_{f, \text{saline}}} = \frac{(1 - \frac{c}{2} k_{\infty, \text{plasma}})^2}{(1 - \frac{c}{2} k_{\infty, \text{saline}})^2} \quad (5.3)$$

This reduces to

$$k_{\infty, \text{plasma}} = k_{\infty, \text{saline}} \cdot \left( \frac{\eta_{f, \text{plasma}}}{\eta_{f, \text{saline}}} \right)^{0.5} - \frac{2}{c} \left( \left( \frac{\eta_{f, \text{plasma}}}{\eta_{f, \text{saline}}} \right)^{0.5} - 1 \right) \quad (5.4)$$

The values at high hematocrit comply with this relationship.

Quemada [1981], in fitting the plasma data of Chien *et al.* [1966], noted that  $k_\infty$  reached a minimum at about  $H=75$  and proposed that this represented the point of inversion of the emulsion, where the erythrocytes begin to behave physically as a continuous phase. Using the first-order approximation that the intrinsic viscosity scales as effective particle volume (see Appendix C), a physical

picture may be constructed in which the effective particle volume at low concentration and high shear rate is the hydrodynamic volume of a single red cell. The effective volume is reduced by crowding as hematocrit increases, and eventually a point is reached where the effective particle is no longer a single cell, but a volume of sterically interacting cells. At this point inversion takes place. With further crowding, the volume of this effective particle no longer changes, but the number of cells it contains increases, leading to increases in intrinsic viscosity. There is not clear-cut evidence of a minimum  $k_{\infty}$  in the results (Fit A) of this work, but there is a trend at hematocrits above 90 of  $k_{\infty}$  increasing with H. The polynomial curve fits to the data (Fit B) reach minimums at about  $H=86-87$ .

The Quemada equation demonstrates a better approach to an asymptotic value at high shear rate than any of the empirical fits attempted, but the curves continue to descend after  $300 \text{ sec}^{-1}$ . Comparison of data values at  $300 \text{ sec}^{-1}$  versus the limiting high shear rate asymptote for the Quemada equation,  $(1 - \frac{c}{2} k_{\infty})^{-2}$ , are shown in Table 5-V for both the A and B fits. The average difference between data and equation values is on the order of 15-20%, although in some cases can be greater than 40%. These results are shown graphically in Figures 5-4 and 5-5. From these figures it can be seen that the difference between the Chien *et al.* values at  $52 \text{ sec}^{-1}$  and the data at  $300 \text{ sec}^{-1}$  is large enough to bring into question any use of the Chien values as the limiting viscosity. For saline, ratio of the Chien value to the data at  $300 \text{ sec}^{-1}$  is about 1.1 to 1.2 over most of the hematocrits. For plasma, the ratio is greater than 2 for the most concentrated suspensions and about 1.1 at physiologic hematocrit.

Hematocrit	Rel. viscosity at 300 1/sec	A fit asymptote	% change from data	B fit asymptote	% change from data
98.4	151.3	122.7	-18.9	87.0	-42.5
97.2	86.7	74.4	-14.2	68.1	-21.5
96.8	59.3	50.0	-15.6	63.2	6.6
96.0	57.3	54.1	-5.6	55.1	-3.9
95.0	47.6	38.1	-20.0	47.1	-1.2
94.0	47.1	40.7	-13.7	40.8	-13.4
93.0	47.2	41.6	-11.9	35.8	-24.1
92.0	43.6	36.7	-15.9	31.8	-27.1
91.0	33.4	29.0	-13.0	28.5	-14.7
90.0	27.5	23.2	-15.6	25.7	-6.5
89.0	25.8	20.6	-20.2	23.4	-9.4
87.0	25.3	21.8	-14.0	19.7	-22.2
86.0	22.3	20.1	-9.8	18.2	-18.3
82.0	19.8	16.0	-19.3	13.9	-29.9
72.0	12.1	9.7	-20.2	8.5	-29.6
70.0	11.2	10.3	-7.6	7.9	-29.6
68.0	12.1	9.9	-18.4	7.3	-39.4
62.0	9.8	7.1	-27.7	6.0	-38.4
59.0	7.1	4.9	-31.5	5.5	-22.2
56.0	6.4	4.5	-29.0	5.1	-20.7
48.0	4.6	3.5	-23.4	4.1	-10.8
42.0	4.1	3.2	-22.0	3.5	-14.3
25.0	2.8	2.5	-11.9	2.2	-20.9
7.0	1.7	1.7	-0.5	1.3	-25.4
Hematocrit	Rel. viscosity at 300 1/sec	A fit asymptote	% change from data	B fit asymptote	% change from data
98.1	49.8	43.0	-13.6	35.4	-28.8
97.6	44.6	36.7	-17.7	33.5	-25.0
97.3	41.8	33.4	-20.0	32.3	-22.6
97.0	36.8	32.1	-12.7	31.3	-14.9
96.0	39.6	33.6	-15.1	28.2	-28.8
95.0	29.5	24.3	-17.7	25.6	-13.3
94.0	23.2	19.9	-14.1	23.3	0.5
93.0	26.5	22.7	-14.4	21.4	-19.3
92.0	21.3	19.7	-7.6	19.7	-7.5
91.0	18.2	16.0	-11.9	18.2	0.2
90.0	20.0	16.9	-15.4	16.9	-15.3
89.0	15.9	13.3	-16.6	15.8	-0.6
86.0	15.3	12.7	-16.7	13.1	-14.7
83.0	13.6	11.4	-16.2	11.0	-18.9
82.0	11.1	9.0	-19.1	10.5	-5.6
78.0	11.5	9.3	-19.4	8.7	-24.6
75.0	9.5	7.4	-21.8	7.6	-19.5
66.0	6.2	4.8	-21.9	5.5	-10.7
59.0	5.1	4.0	-21.6	4.5	-12.1
50.0	4.2	4.1	-3.0	3.5	-16.3
44.0	3.7	3.7	-0.1	3.0	-18.3
34.0	2.7	2.6	-4.5	2.4	-12.4
24.0	2.0	1.8	-9.2	1.9	-7.3
11.0	1.3	1.3	-2.6	1.3	2.8

**Table 5-V: High shear rate asymptotes.**

*Top*, saline suspensions. *Bottom*, plasma suspensions.

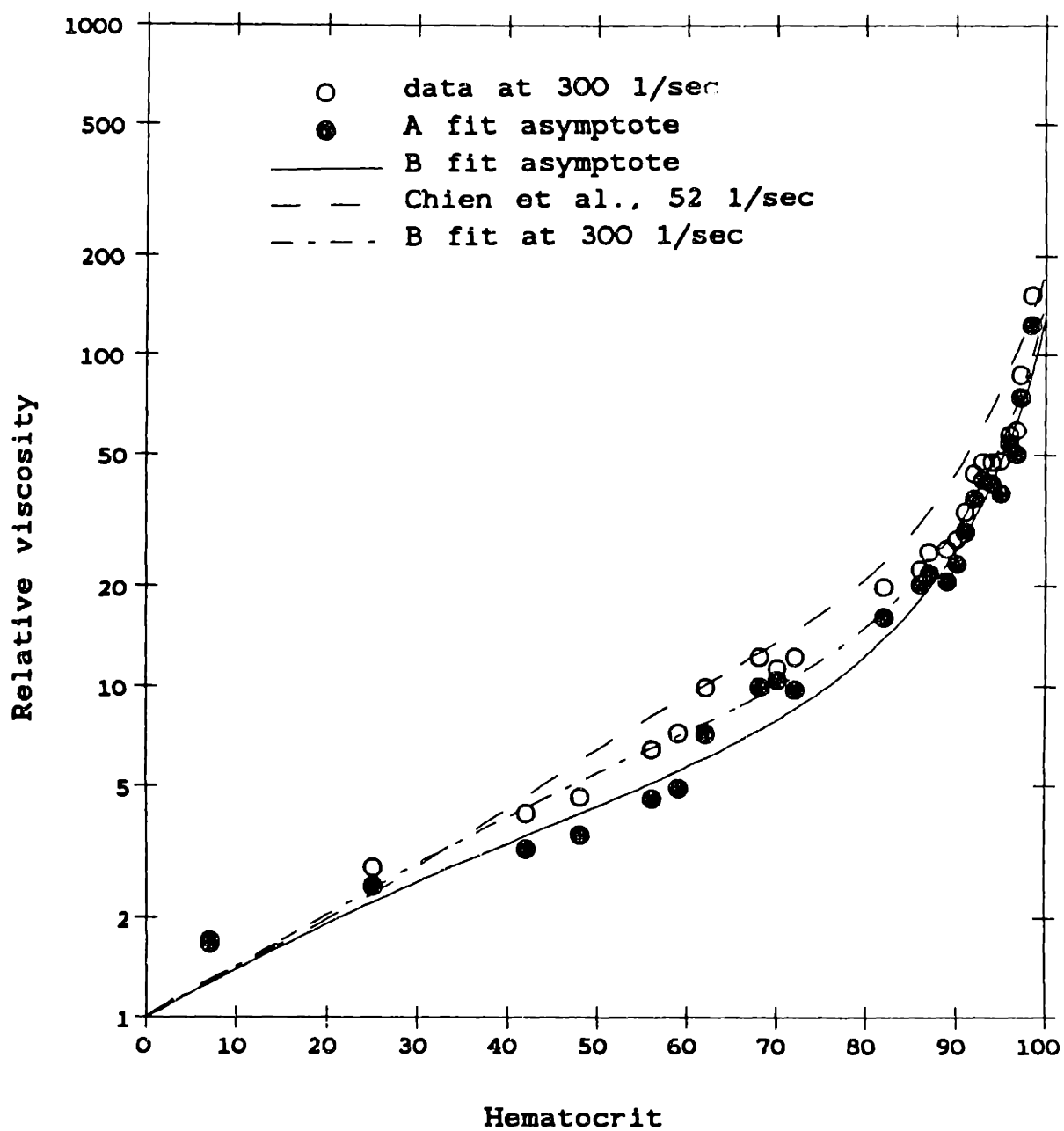
Average change from data for saline =  $-16.7\% \pm 7.3$  (A fit)

=  $-20.0\% \pm 12.2$  (B fit)

Average change from data for plasma =  $-13.9\% \pm 6.3$  (A fit)

=  $-13.9\% \pm 9.1$  (B fit)

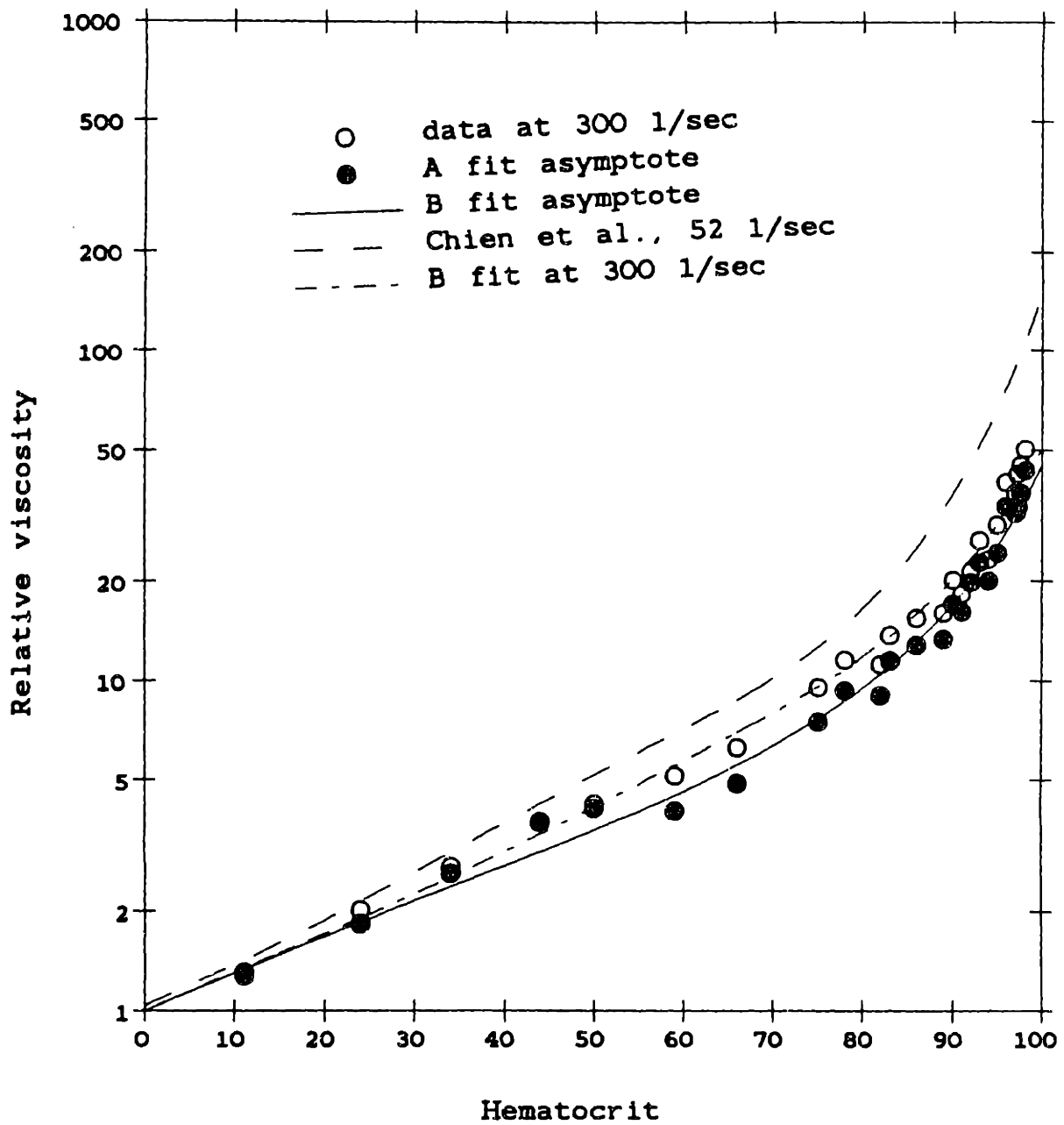
# High shear rate asymptotes for saline fits



**Figure 5-4:** High shear rate asymptotes for saline fits.



High shear rate asymptotes  
for plasma fits



**Figure 5-5:** High shear rate asymptotes for plasma fits.

### 5.3 Dependence of $k_0$ on hematocrit and plasma proteins

The values for  $k_0$  of 4.48 in plasma at H=44, 2.22 in plasma at H=89, and 3.48 in saline at H=48 were comparable to values previously fit by Quemada (see Table 2-II. The A fit values for  $k_0$  are plotted versus hematocrit in Figure 5-6 for saline and in Figure 5-7 for plasma. Again, included in both plots are values from fits of the Quemada equation to the Chien polynomials. There is very close agreement between the A fit  $k_0$  and the Chien  $k_0$ , particularly at high hematocrits, and this follows from the comparable viscosity values obtained at low shear rate and high hematocrits.

The B fits in Figure 5-6 represents a relatively weighted fit of the open symbol to a third-order polynomial, to which an exponential term has been added for plasma:

$$k_0 = \sum_{i=1}^4 k_{0i} \cdot c^{(i-1)} + k_{05} \cdot \exp(k_{06} \cdot c) \quad (5.5)$$

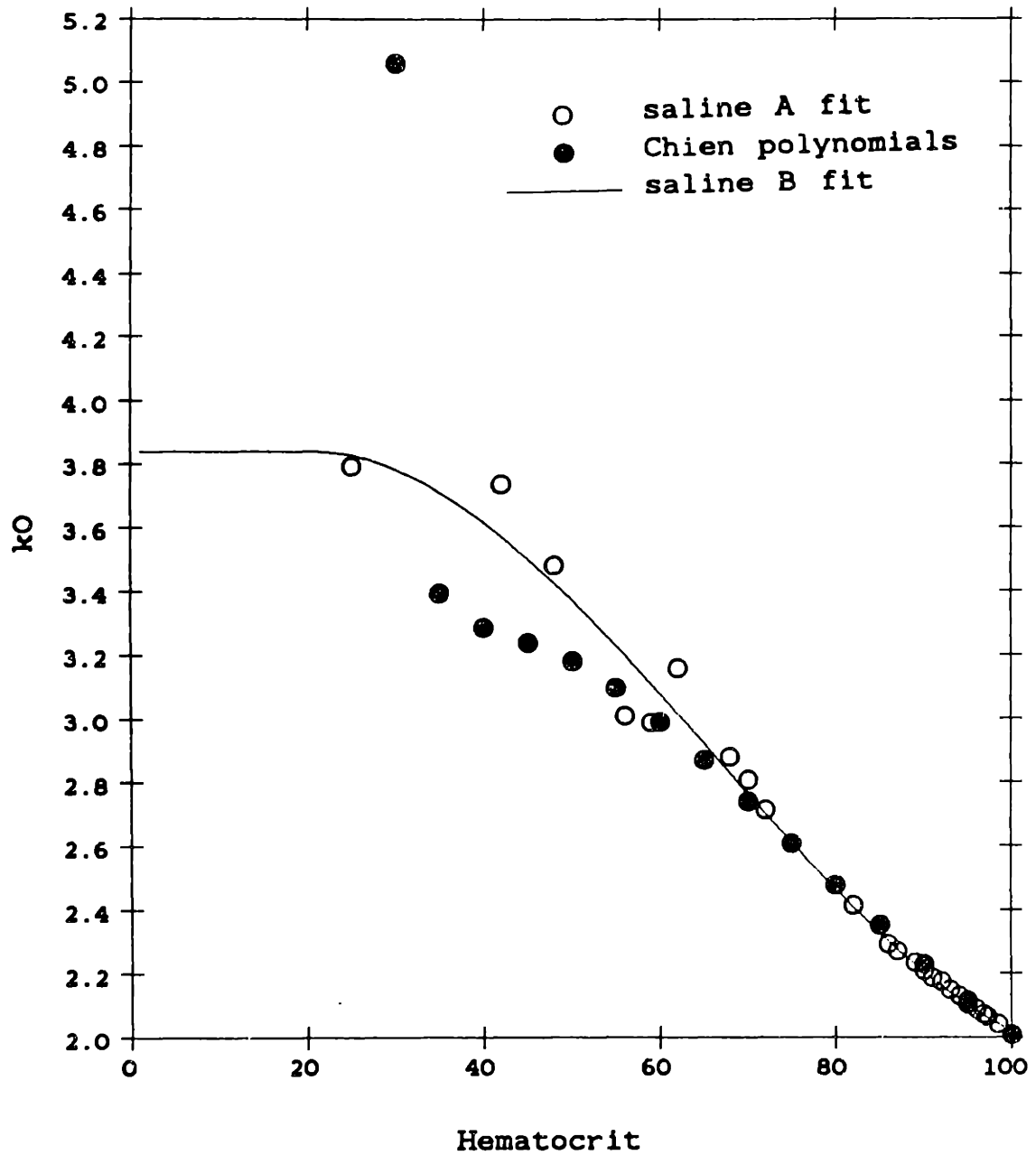
The best-fit values and standard deviations for saline are

$$\begin{aligned} k_{01} &= 3.503 \pm 0.413 \\ k_{02} &= 3.467 \pm 2.140 \\ k_{03} &= -9.976 \pm 3.470 \\ k_{04} &= 5.026 \pm 2.841. \\ k_{05} &= 0 \\ k_{06} &= 0. \end{aligned}$$

At  $c < 20$ , the data is fit to a constant  $k_0 = 3.838$ . The polynomial reaches a maximum at about  $c = 0.2$ , and there is no physical reason to expect  $k_0$  to do this. It seems physically reasonable to expect  $k_0$  to approach a constant at low hematocrits in saline, as will be explained shortly. Attempts were made to fit to equations which level off to a constant as  $c$  approaches zero, such as equations of the form  $k_{01} \tanh(k_{02}c + k_{03})$ , but these did not prove satisfactory.

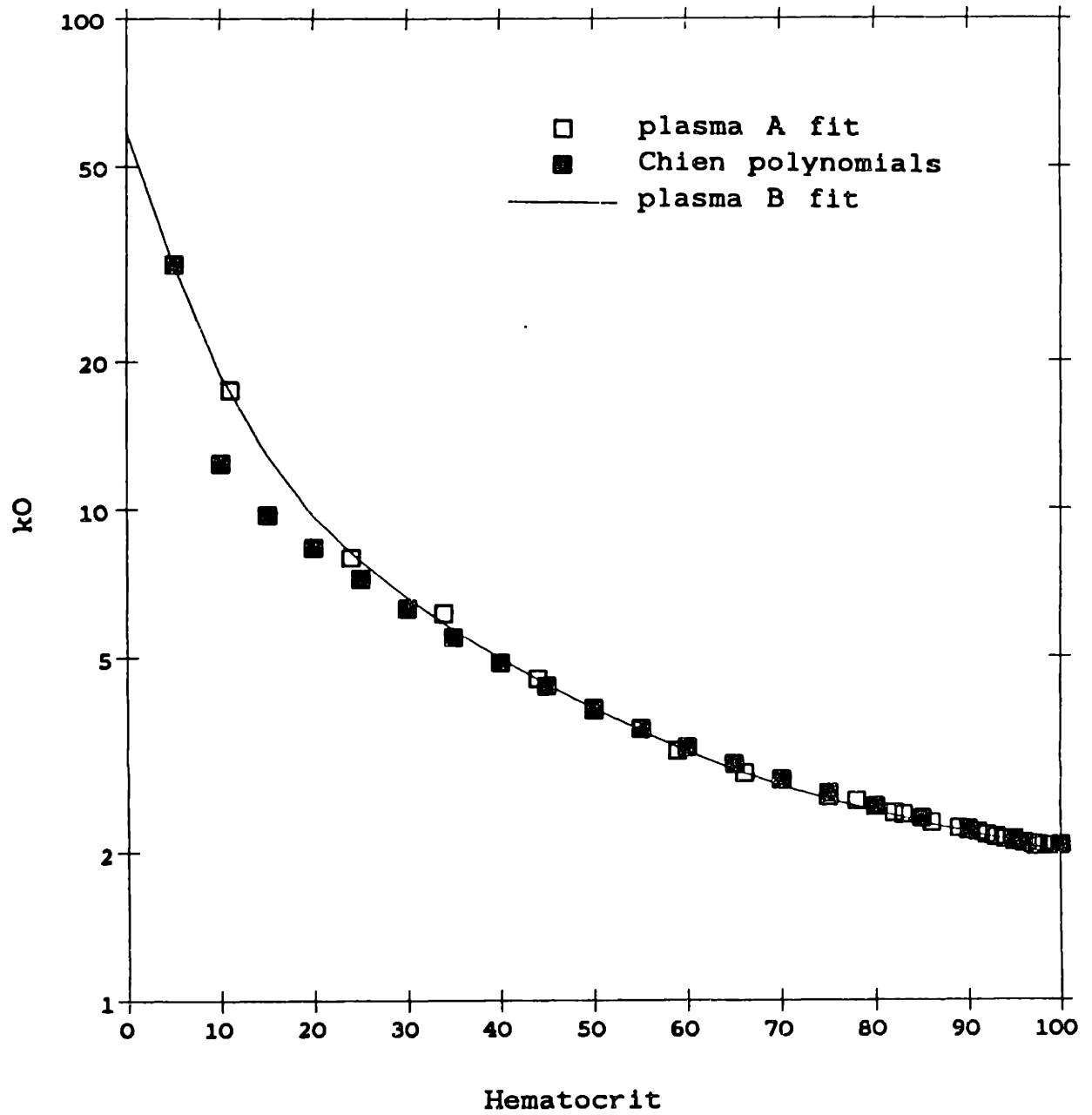
For the plasma fit, the exponential term is required to model the dependence

$k_0$  vs. hematocrit  
saline suspensions



**Figure 5-6:** Plot of  $k_0$  versus hematocrit for saline suspensions.

$k_0$  vs. hematocrit  
plasma suspensions



**Figure 5-7:** Plot of  $k_0$  versus hematocrit for plasma suspensions.

at low hematocrits. The best-fit values for plasma are

$$k_{01} = 13.28 \pm 0.60$$

$$k_{02} = -31.93 \pm 2.67$$

$$k_{03} = 32.13 \pm 3.74$$

$$k_{04} = -11.49 \pm 0.02$$

$$k_{05} = 45.46 \pm 1.06$$

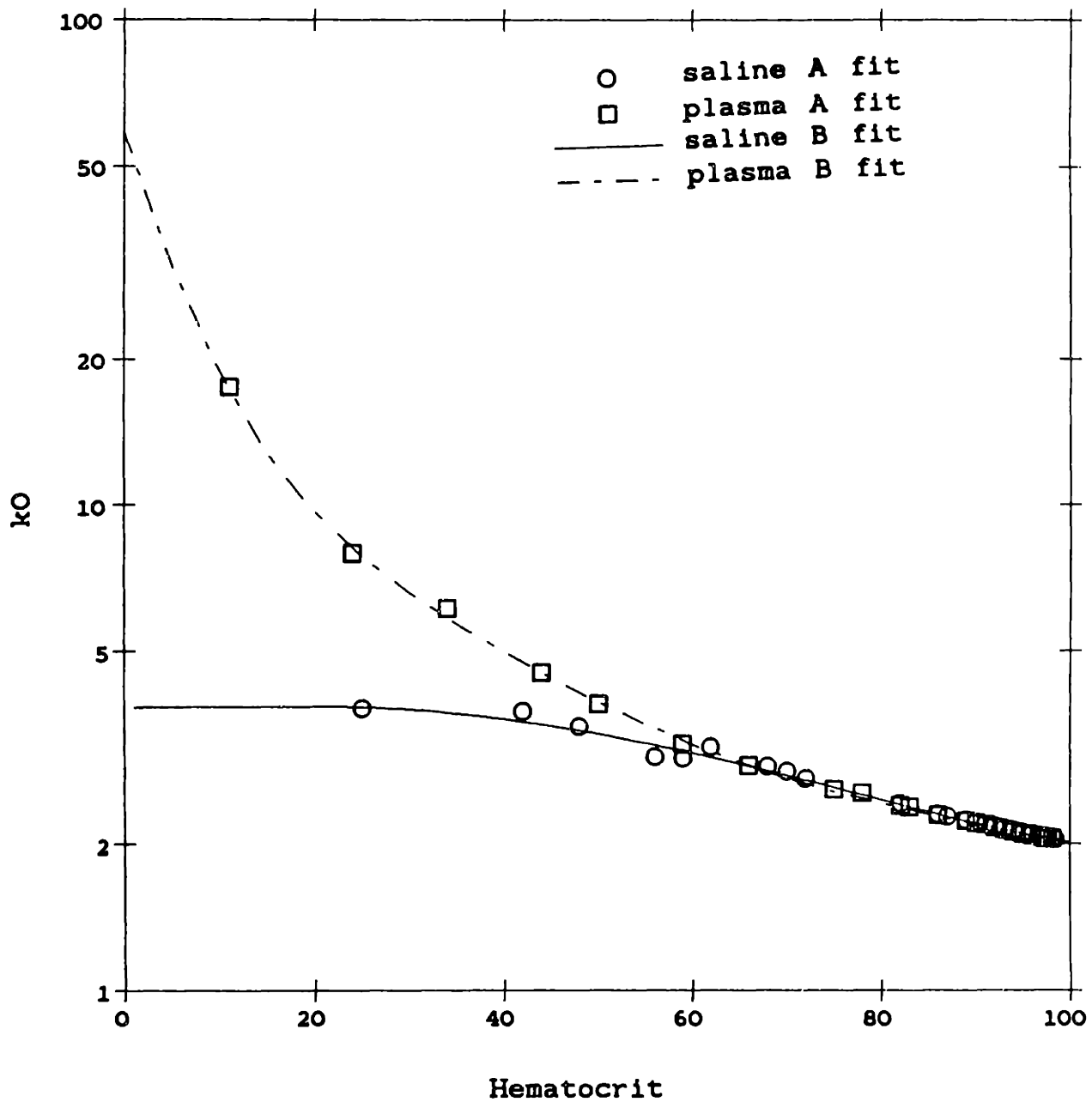
$$k_{06} = -16.78 \pm 0.65.$$

The plasma B fit is a relatively weighted fit to the open squares in Figure 5-7 plus the single filled square at  $H=5$ . The Chien polynomial value was included because these values closely followed the data at low hematocrit and because this is a critical region for fitting the exponential term. Values at hematocrits of 90 and above were additionally weighted by a factor of 10 to insure agreement at high  $H$ . This additional weighting is needed to avoid the appearance of singular points, where  $\frac{c}{2}k_0 > 1$ , in the B fits.

The sensitivity of the asymptotic zero-shear-rate viscosity to the value of  $k_0$  is an expression for  $\frac{\partial \eta_r}{\partial k}$  equivalent to that shown for  $k_\infty$  in the last section, with  $k$  now equal to  $k_0$ . For saline suspensions, this sensitivity is highest at high hematocrits. For plasma suspensions, the sensitivity is highest at low hematocrits because of the large values for  $k_0$  at low hematocrits. For example, for both suspensions, at  $c=0.95$ ,  $k_0 \approx 2$  and  $\frac{\partial \eta_r}{\partial k} = 780$ . For saline at  $c=0.45$ ,  $k_0 \approx 3.2$  and  $\frac{\partial \eta_r}{\partial k} = 20.5$ . For plasma at  $c=0.45$ ,  $k_0 \approx 4.4$  and  $\frac{\partial \eta_r}{\partial k} = 450,000$ . Insurance must also be taken to avoid singular points in the refit Quemada equation where  $\frac{c}{2}k_0 - 1 > 0$ .

Saline and plasma values are plotted together on Figure 5-8. The values are almost identical at high hematocrits, down to about  $H=65$ . With further decrease in hematocrit, the saline values appear to level off, while the plasma values show a large increase. At zero shear rate, the emulsive and shear deformation effects are

kO vs. hematocrit  
saline vs. plasma suspensions



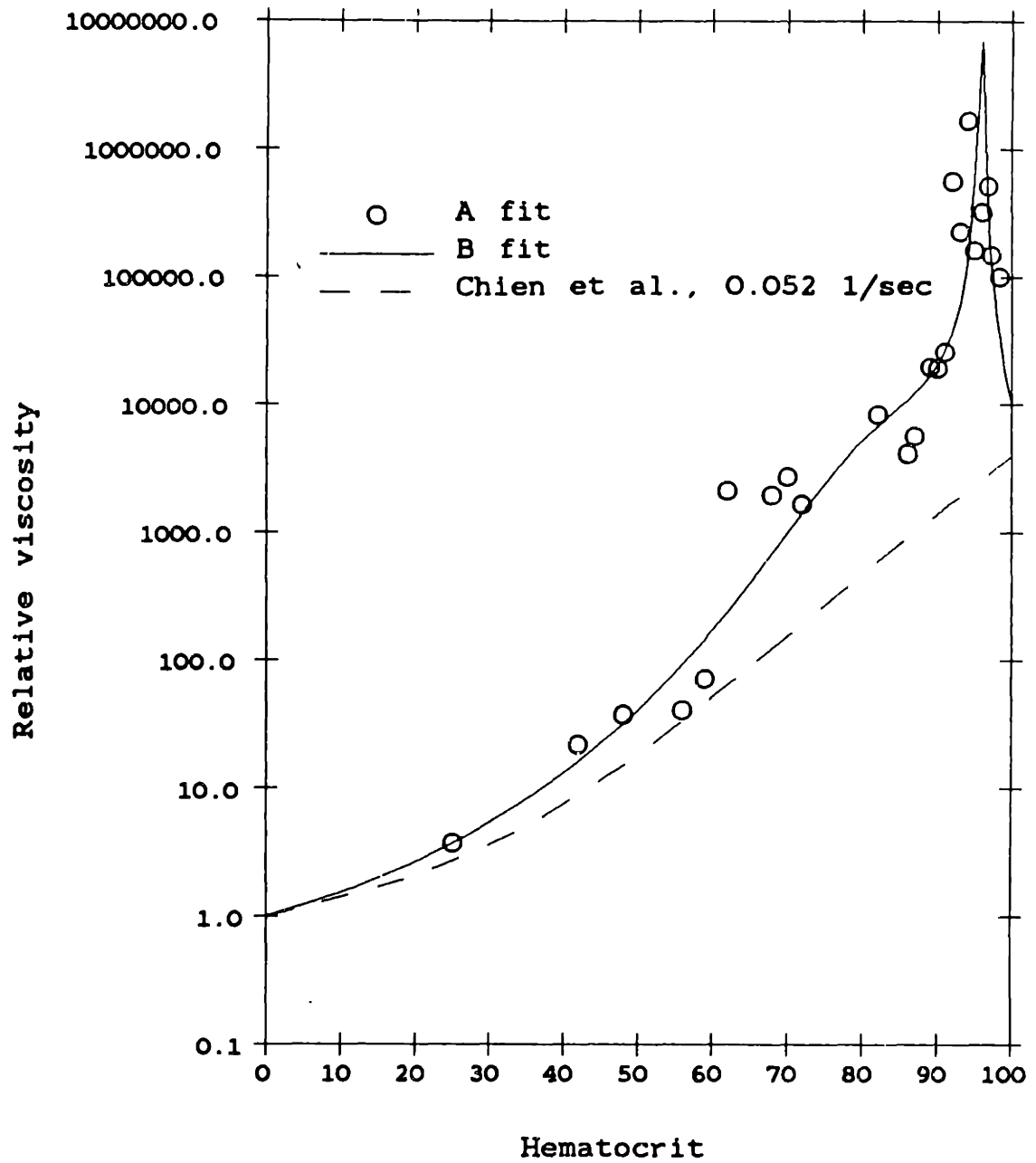
**Figure 5-8:** Comparison of  $k_0$  values for saline and plasma suspensions.

minimized, and thus  $k_0$  should principally be an indicator of cell-cell interactions, either protein-mediated (aggregation) or steric (crowding). At low hematocrits, crowding is not seen, and the difference between saline and plasma  $k_0$  is caused by rouleaux formation. Where plasma and saline values are equal, at high hematocrits, should be where crowding effects predominate over aggregation. An analysis by Chien *et al.* [1966] suggested that at a hematocrit of 60, crowding interactions begin to influence viscosity relative to aggregation and that at  $H=80$ , they begin to be the dominant influence. This is also probably related to the inversion phenomenon. The merging of the B fit lines in Figure 5-8 around  $H=65$  would seem to lend some credence to this hypothesis.

At low shear rate, and at hematocrits low enough so that crowding is not a factor, normal cells in saline solution should approximate the behavior of hardened cells in saline. The maximum packing concentration  $c_p$  for hardened erythrocytes in saline has been determined experimentally to be about 0.6 [Chien *et al.*, 1971]. Using the relationship derived in Appendix C for rigid particles,  $k_0=2/c_p=3.33$ , which compares reasonably well with the asymptotic value of 3.838 from the B fit.

The A and B fit values for viscosity at zero shear rate are plotted in Figures 5-9 and 5-10 for saline and plasma suspensions, respectively. Unlike for high shear rates, a smooth limiting envelope can not be generated from the  $k_0$  values. The  $\gamma=0$  curve for saline is not well-behaved at high hematocrits, and the  $\gamma=0$  curve for plasma is obviously meaningless. The message behind these figures is that, except for low hematocrit saline suspensions, the Quemada equation is extremely sensitive to  $k_0$  in the limit of low shear rate. If one wishes to use the Quemada equation at low shear rates,  $k_0$  values must be derived from fits to data that clearly reach an asymptote at low  $\gamma$ . Values for  $k_0$  which have been extrapolated from data at higher shear rates, such as those in this work, will do a satisfactory job of

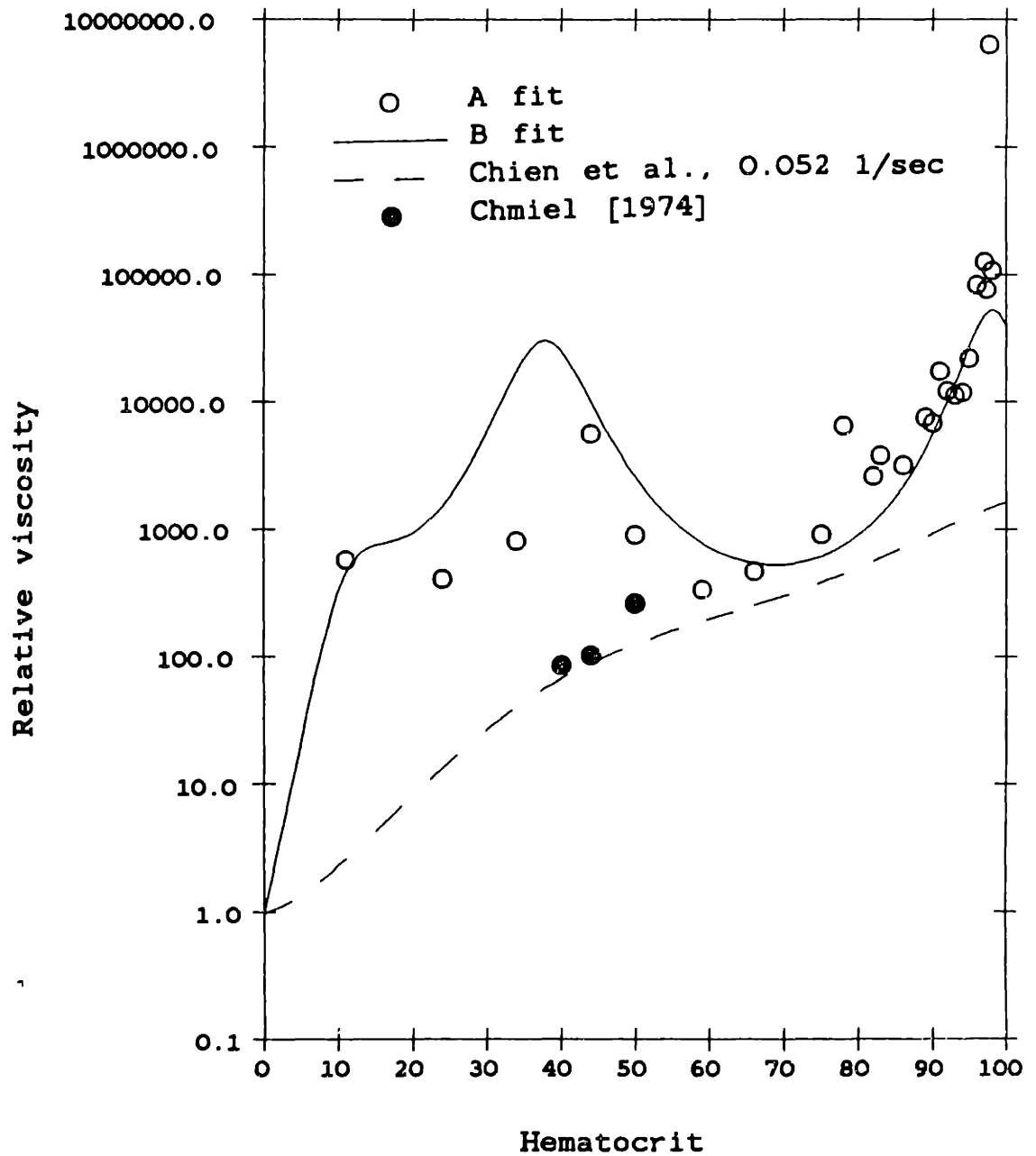
Viscosity at zero shear rate  
for saline fits



**Figure 5-9:** Zero shear rate asymptotes for saline fits.



Viscosity at zero shear rate  
for plasma fits



**Figure 5-10:** Zero shear rate asymptotes for plasma fits.

modeling overall behavior, but they will not be accurate enough for predicting viscosity at low shear rates.

At very high hematocrits, the difference between saline and plasma  $k_0$  can probably also be related by equation (5.4), replacing  $k_\infty$  with  $k_0$ . An attempt was made to use this equation to derive new plasma  $k_0$  values from the Fit A saline values, in hopes that the new values would provide a more realistic  $\gamma = 0$  curve. The highest hematocrit at which a singular point does not result from the saline A fit is 95. Using the saline value 2.1 at  $H=95$ , equation (5.4) gives a plasma  $k_0$  of 2.099, and  $\eta_r=10^5$ . Similarly, at  $H=94$ , the saline value of 2.126 results in a plasma value of 2.126 and  $\eta_r=10^6$ . This indicates that at high hematocrits, the saline A fit values are inaccurate also, at least on the scale of determining low shear rate viscosity.

As an example of just how accurately  $k_0$  must be determined, Chmiel [1974] collected data for whole blood at shear rates of 0.1 to 0.005  $\text{sec}^{-1}$  in a wide-gap Couette system, in which an asymptotic approach to a constant viscosity at low shear rate was demonstrated. Assuming a  $\eta_r$ , which is not explicitly given in the paper, of 1.2 cP, relative viscosities at  $H=50$ , 44, and 40 correspond to  $k_0$  values of 3.75, 4.09 and 4.45, respectively. The parameter values from this study of 3.87 at  $H=50$  and 4.48 at  $H=44$  are both comparable, within 3 and 10 per cent of the Chmiel values, respectively, and yet the calculated viscosities at  $\gamma=0$  are an order of magnitude or more higher. For comparison, Chmiel's data is shown on Figure 5-10 as filled circles.

From an analytical standpoint, the difficulty stems from the fact that for plasma,  $\frac{c}{2}k_0$  is always approximately equal to 0.99. Thus a smooth plot of  $(1-\frac{c}{2}k_0)^{-2}$  depends on accurately representing small deviations from zero.

#### 5.4 Dependence of $\gamma_c$ on hematocrit and plasma proteins

The values for the A fits of  $\gamma_c$  for saline are plotted versus hematocrit in Figure 5-11. Also plotted are values derived from fits of the Quemada equation to the Chien *et al.* polynomials. There is a good deal of scatter in the A fit values, but the trend seems to be for  $\gamma$  becoming large as hematocrit decreases. The values from the Chien polynomials decay toward zero as hematocrit decreases. The result of these opposing trends is seen in the different curvatures of the two sets of data along the  $\gamma$  axis (see Figures 4-20 to 4-23). The data from this work and the A fits are convex, while the Chien points indicate a concave curvature.

From a strictly analytical viewpoint, it can be seen how  $\gamma_c$  affects this curvature. It will be recalled from the discussion in Section 2.4 that the Quemada equation results in a sigmoidal curve along the x-axis ( $\gamma$  axis). As  $\gamma_c$  increases, the curve is shifted to the right, and the curvature at a fixed shear rate becomes more convex. As  $\gamma_c$  decreases, the reverse happens.

The question exists whether the high  $\gamma_c$  values are artifacts resulting from secondary flows in the viscometer at higher shear rates and low hematocrits. Jerrard [1950] has developed charts for assessing the onset of instability in rotating-cylinder viscometers in which the outer cylinder rotates, such as the Haake viscometer used in this work. Based on the geometry of the Haake viscometer, the critical shear rate in  $\text{sec}^{-1}$  for secondary flow is  $3.5 \times 10^5 \nu$ , where  $\nu$  is the kinematic viscosity of the fluid in stokes. For water, plasma, and red cells in saline at  $H=42$ , the calculated values are 2300, 4100, and 10,000  $\text{sec}^{-1}$ , respectively. It would seem that secondary flow is not a consideration, and yet shear-thickening behavior was noted in the measurement of both saline and plasma viscosities, which is normally an indication of some instability.

gamma-critical vs. hematocrit  
saline suspensions

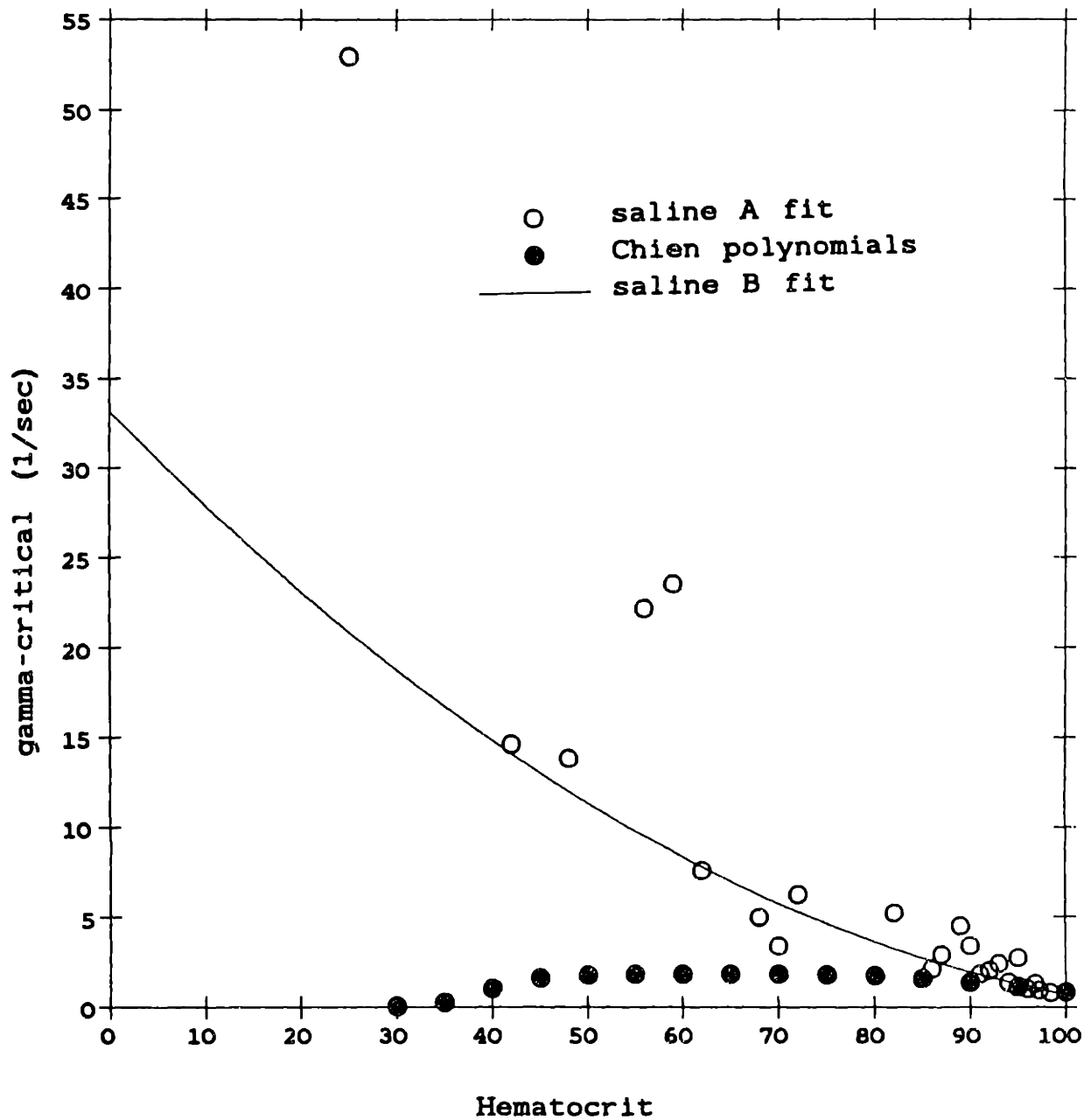


Figure 5-11: Plot of  $\gamma$  versus hematocrit for saline suspensions.

It is possible that the concavity of the Chien *et al.* values is forced from the nature of the fifth-order polynomials used to calculate them, and that the actual data of Chien [et al.] would correspond to the data of this work. Later data by Chien [1970] shows convex curvature similar to the A fits (see Figures 2-3 and 2-11). Quemada [1978b] fit a  $\gamma_c$  value of  $25 \text{ sec}^{-1}$  to this later data of Chien, which agrees with the values shown in Figure 5-11.

Plasma values for  $\gamma_c$  are plotted in Figure 5-12. Both the Chien polynomials and the A fits indicate  $\gamma_c \approx 0$  at  $H=0$ , then an increase to a maximum at about  $H=70-75$ , and subsequent decrease with increasing  $H$ . Correspondingly, the A fits are concave at both high and low hematocrits, but around  $H=70$  (see Figures 4-42 to 4-44) a more sigmoidal curve is seen.

The B fit in Figure 5-11 represents a relatively weighted fit of the saline A fit values to the form

$$\gamma_c = \sum_{i=1}^3 \gamma_{ci} \cdot c^{(i-1)} \quad (5.6)$$

where the best-fit values are

$$\gamma_{c1} = 33.16 \pm 11.95$$

$$\gamma_{c2} = -54.79 \pm 28.57$$

$$\gamma_{c3} = 22.33 \pm 16.93.$$

A second-order polynomial was fit to the Chien values in Figure 5-11. This fit, when substituted for the saline B fit curve from this study, resulted in a shifting of the curve to the left and predicted relative viscosities up to an order of magnitude lower than the data at low shear rates. For plasma, the B fit is of the form

$$\gamma_c = \gamma_{c1} \cdot \exp(-(\gamma_{c2}c + \gamma_{c3})^2) \quad (5.7)$$

where by a relatively weighted fit, the values are

gamma-critical vs. hematocrit  
plasma suspensions

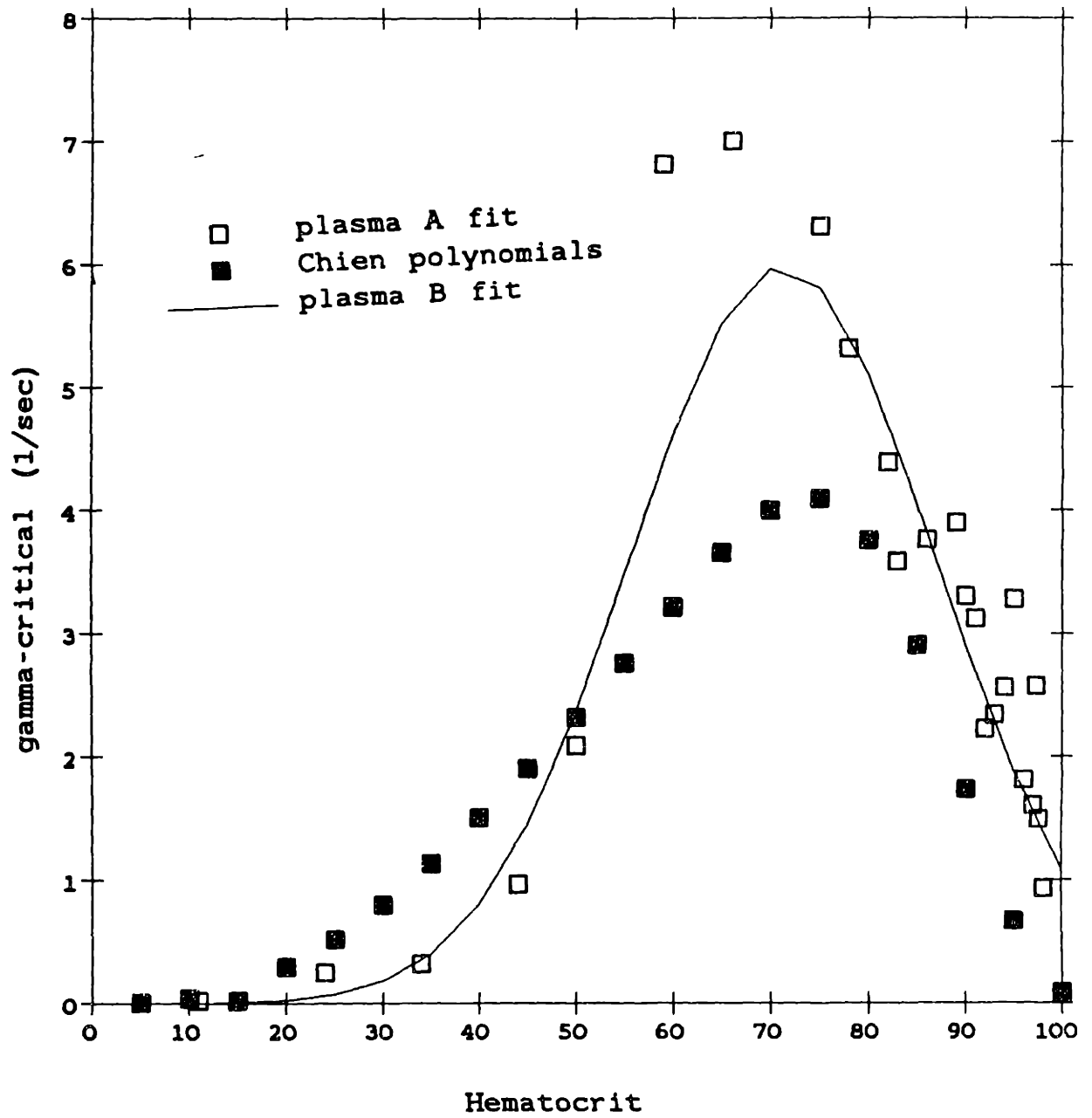


Figure 5-12: Plot of  $\gamma_c$  versus hematocrit for plasma suspensions.

$$\gamma_{c1} = 5.979 \pm 0.580$$

$$\gamma_{c2} = -4.536 \pm 0.162$$

$$\gamma_{c3} = 3.232 \pm 0.098.$$

The saline and plasma values are plotted together on Figure 5-13. The values from the two suspensions converge with increasing hematocrit at about  $H=70$  and are comparable at high concentrations, with the saline values somewhat higher.

From a physical interpretation,  $\gamma_c$  is analagous to a rotational diffusion constant, with a dependence  $\gamma_c \sim a^{-3}$ , where  $a$  is the particle radius (see Appendix C). Thus effects which increase effective particle size will decrease  $\gamma_c$ . The two such mechanisms for increasing effective particle size are aggregation and crowding. In saline, the cells do not aggregate, and so there is a simple trend of  $\gamma_c$  decreasing with hematocrit. In plasma suspensions, one could imagine that initially as concentration increases, crowding acts to diminish rouleau and aggregate size, and  $\gamma_c$  increases. When  $\gamma_c$  reaches a maximum that roughly corresponds to inversion, crowding acts to pack cells together and  $\gamma_c$  decreases.

### 5.5 Limiting behavior of parameters

A major concern in globally refitting the Quemada expression over all hematocrits is the limiting behavior of the equation, and specifically the intrinsic viscosity  $k = k_\infty + \frac{k_0 - k_\infty}{1 + \sqrt{\gamma_r}}$  as the cell volume concentration goes to zero. Since the equation was derived for concentrated dispersions, it may be appropriate to say that it is not valid in the limit  $c \rightarrow 0$ . Nevertheless, the equation can be fit successfully to low concentration data, and the parameters do appear to approach asymptotic values at low hematocrits.

Newtonian behavior should be recovered where  $k$  is not a function of shear

gamma-critical vs. hematocrit  
saline vs. plasma suspensions

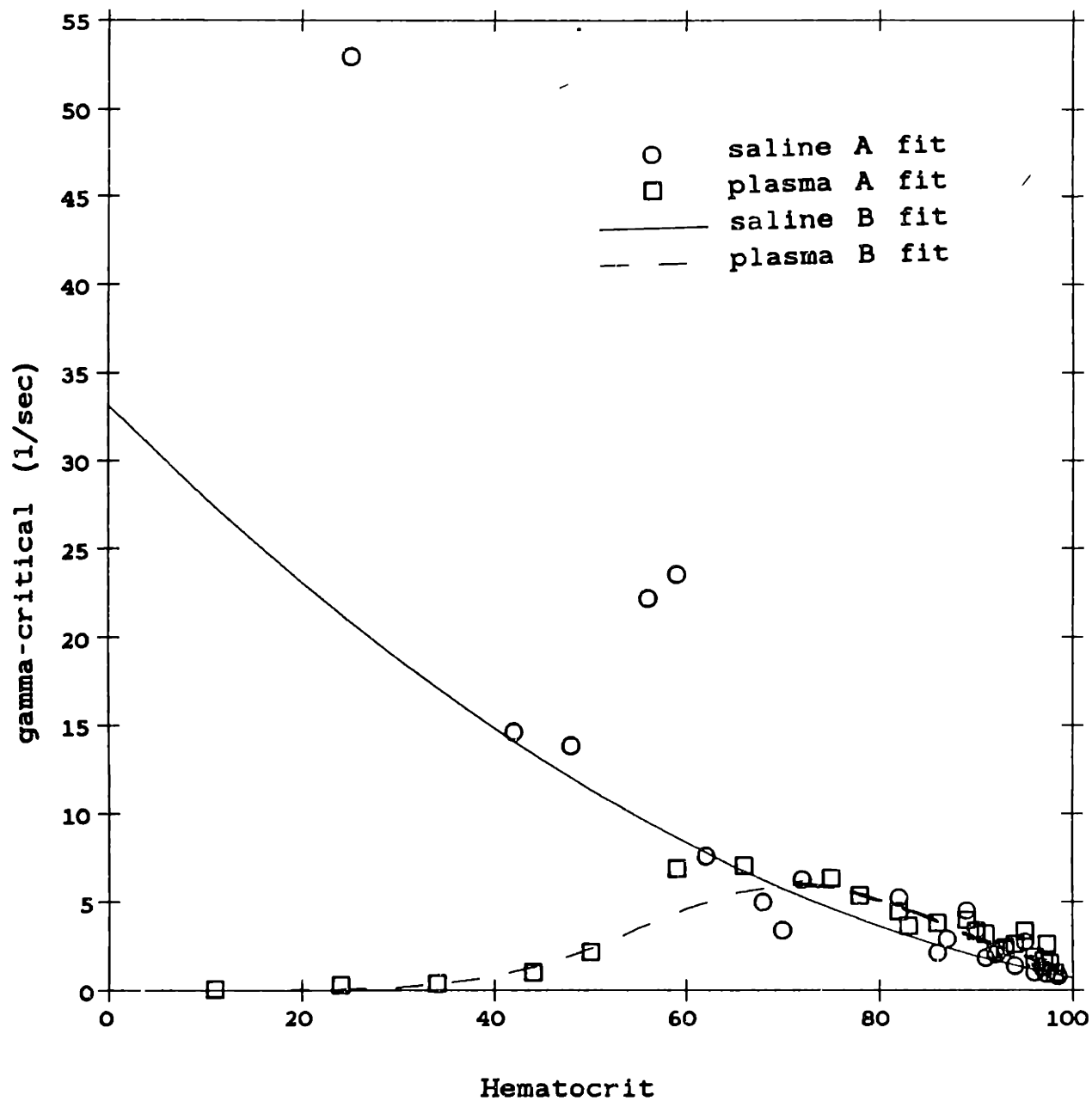


Figure 5-13: Comparison of  $\gamma_c$  for saline and plasma suspensions.



rate. There are three possible mathematical ways by which  $k$  can become independent of shear rate, and these will be considered in turn, first from a purely analytical standpoint, and then by how the analysis relates to data and to physical interpretation:

- $k_{\infty} = k_0$ . Quemada [1981] proposed that at the Newtonian thresholds  $H=30$  for saline suspensions and  $H=5$  for plasma suspensions,  $k_{\infty}$  and  $k_0$  become equal at some intermediate value. The data does not support this, as the difference between the two parameters is greatest at low hematocrits for both plasma and saline suspensions. Additionally, from a conceptual standpoint, it would seem that even in the absence of aggregation and cell-cell crowding, orientation and deformation would lead to a finite distinction between intrinsic viscosities at low and high shear rates.
- $\gamma_c \rightarrow \infty, k \rightarrow k_0$ . This behavior is demonstrated by the Eagle's solution data of this work and appears to correlate with the analysis of  $\gamma_c$  as inversely proportional to particle radius, although if this were strictly true, some asymptotic finite value would be expected. Also, the limit of  $k \rightarrow k_0$  is troublesome, because it indicates that as long as  $\gamma_c$  is finite, the viscosity will decrease at large enough  $\gamma$ . A more realistic stating of this limit that would satisfy these concerns is to combine this criterion with the first case, so that  $\gamma_c$  approaches a large, but finite, value and  $k_{\infty} \rightarrow k_0$ .
- $\gamma_c \rightarrow 0, k \rightarrow k_{\infty}$ . This behavior is demonstrated by both sets of Chien data and by the plasma data of this work after incorporation of the Chien data, and it is also coupled with a large increase in  $k_0$ , which causes the curves to increase sharply at low  $H$ . It appears to be

incompatible with the idea of  $\gamma_c$  as a rotational diffusion constant dependent solely on effective particle radius, although it seems logical that as concentration decreases, non-Newtonian behavior will be expressed at increasingly lower shear rates. A possible explanation is that there is an initial increase in the frequency of particle collisions as cell concentration increases, with a resulting enhancement of aggregate formation, followed by a decrease when steric effects limit mobility. It is interesting to note that the maximum in  $\gamma_c$  in all cases occurs around  $H=70-75$ , which again corresponds with the analysis of Chien *et al.* [1966] as the advent of prominent effects from cell crowding.

## 5.6 Comparison of global fit equations to data

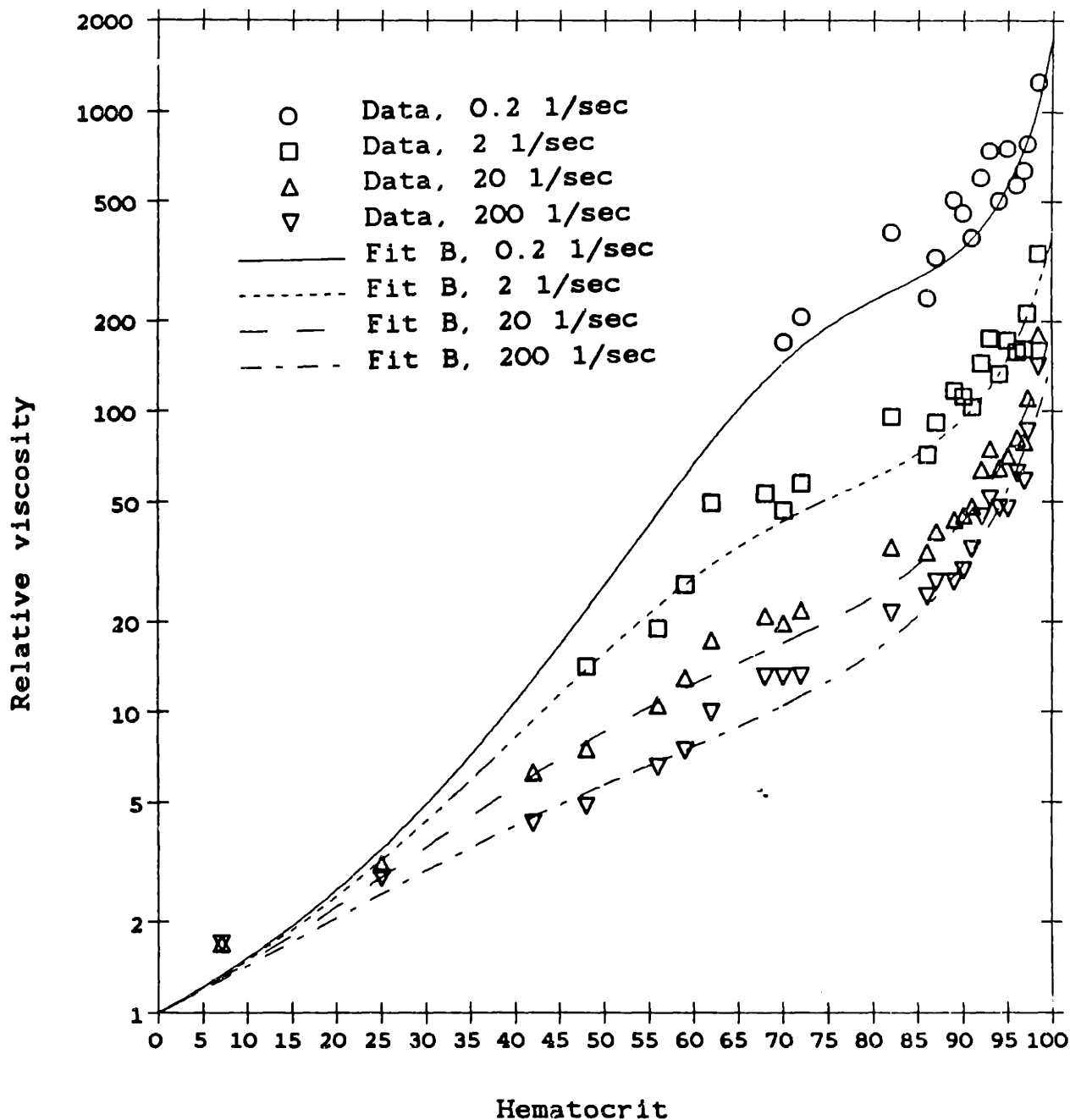
The B fit values for the rheological parameters are shown in Table 5-VI. These values are obtained from the equations discussed in the previous three sections, and are used to plot the dashed curves in Figures 4-2 to 4-49. As shown in Figures 5-14 to 5-17, the B fits provide a reasonable representation of the data, whether plotted as a function of hematocrit or of shear rate. The fits are best at shear rates greater than  $0.5 \text{ sec}^{-1}$ . The plasma curves provide better fits to data at shear rates down to  $0.05 \text{ sec}^{-1}$ , despite the crossing over of the plasma curves of Figure 5-17 and the relatively smooth behavior of the saline curves in Figure 5-15. Even in the absence of data, the saline curves probably adequately model viscosity at low hematocrits and low shear rates, since the behavior is Newtonian. At hematocrits greater than 98, both the saline and plasma fits tend to underestimate the data. When plotted against the Chien polynomials (see Figures 5-18 and 5-19) the same comments can be made as for the comparison of the polynomials to the data of this work (see Figures 4-54 and 4-55).

Hematocrit	k-inf. B fit	k-O B fit	gam.-crit. B fit
98.4	1.815	2.044	0.858
97.2	1.808	2.063	1.001
96.8	1.806	2.070	1.047
96.0	1.803	2.084	1.141
95.0	1.798	2.102	1.262
94.0	1.795	2.122	1.388
93.0	1.791	2.142	1.519
92.0	1.788	2.163	1.653
91.0	1.786	2.184	1.793
90.0	1.784	2.207	1.936
89.0	1.782	2.230	2.084
87.0	1.781	2.278	2.394
86.0	1.781	2.303	2.556
82.0	1.785	2.409	3.247
72.0	1.826	2.704	5.287
70.0	1.840	2.766	5.749
68.0	1.855	2.828	6.228
62.0	1.912	3.016	7.774
59.0	1.947	3.108	8.607
56.0	1.986	3.199	9.480
48.0	2.110	3.425	12.006
42.0	2.221	3.572	14.087
25.0	2.626	3.825	20.858
7.0	3.838	3.699	29.434

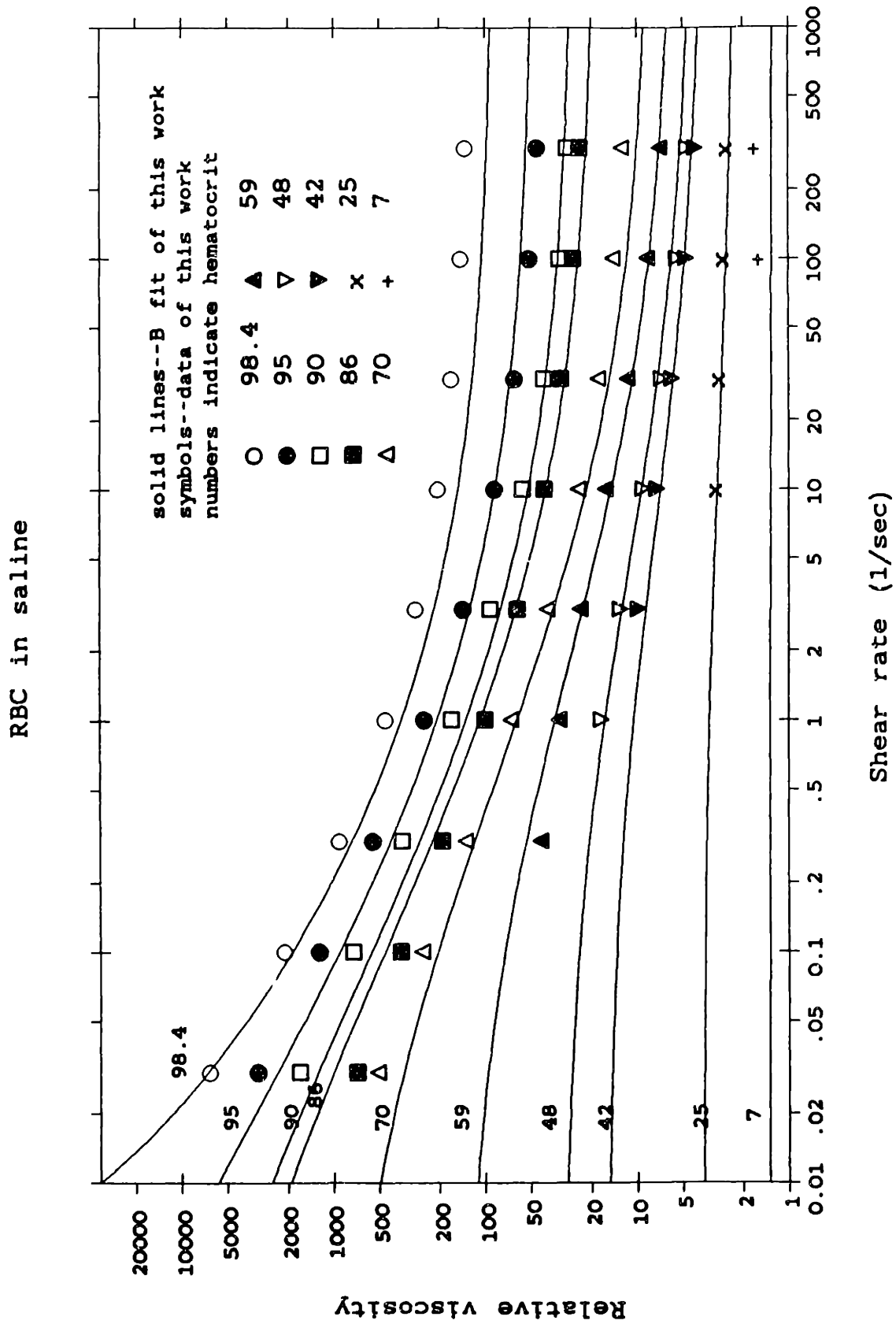
Hematocrit	k-inf. B fit	k-O B fit	gam.-crit. B fit
98.1	1.696	2.030	1.357
97.6	1.695	2.040	1.433
97.3	1.694	2.046	1.480
97.0	1.693	2.052	1.528
96.0	1.691	2.073	1.696
95.0	1.689	2.093	1.874
94.0	1.687	2.112	2.062
93.0	1.685	2.132	2.259
92.0	1.684	2.152	2.466
91.0	1.683	2.172	2.680
90.0	1.682	2.192	2.901
89.0	1.682	2.212	3.127
86.0	1.682	2.275	3.822
83.0	1.684	2.343	4.501
82.0	1.686	2.366	4.714
78.0	1.693	2.470	5.444
75.0	1.702	2.558	5.809
66.0	1.742	2.899	5.649
59.0	1.788	3.268	4.390
50.0	1.867	3.922	2.361
44.0	1.930	4.501	1.297
34.0	2.058	5.838	0.344
24.0	2.212	8.119	0.060
11.0	2.451	17.319	0.003

**Table 5-VI:** B fit values of Quemada parameters for Eagle's solution and plasma.  
*Top*, values for Eagle's solution. *Bottom*, values for plasma.

Comparison of saline B fit to data



**Figure 5-14:** Comparison of saline global fit to data as a function of hematocrit.



**Figure 5-15:** Comparison of saline global fit to data as a function of shear rate.

Comparison of plasma B fit to data

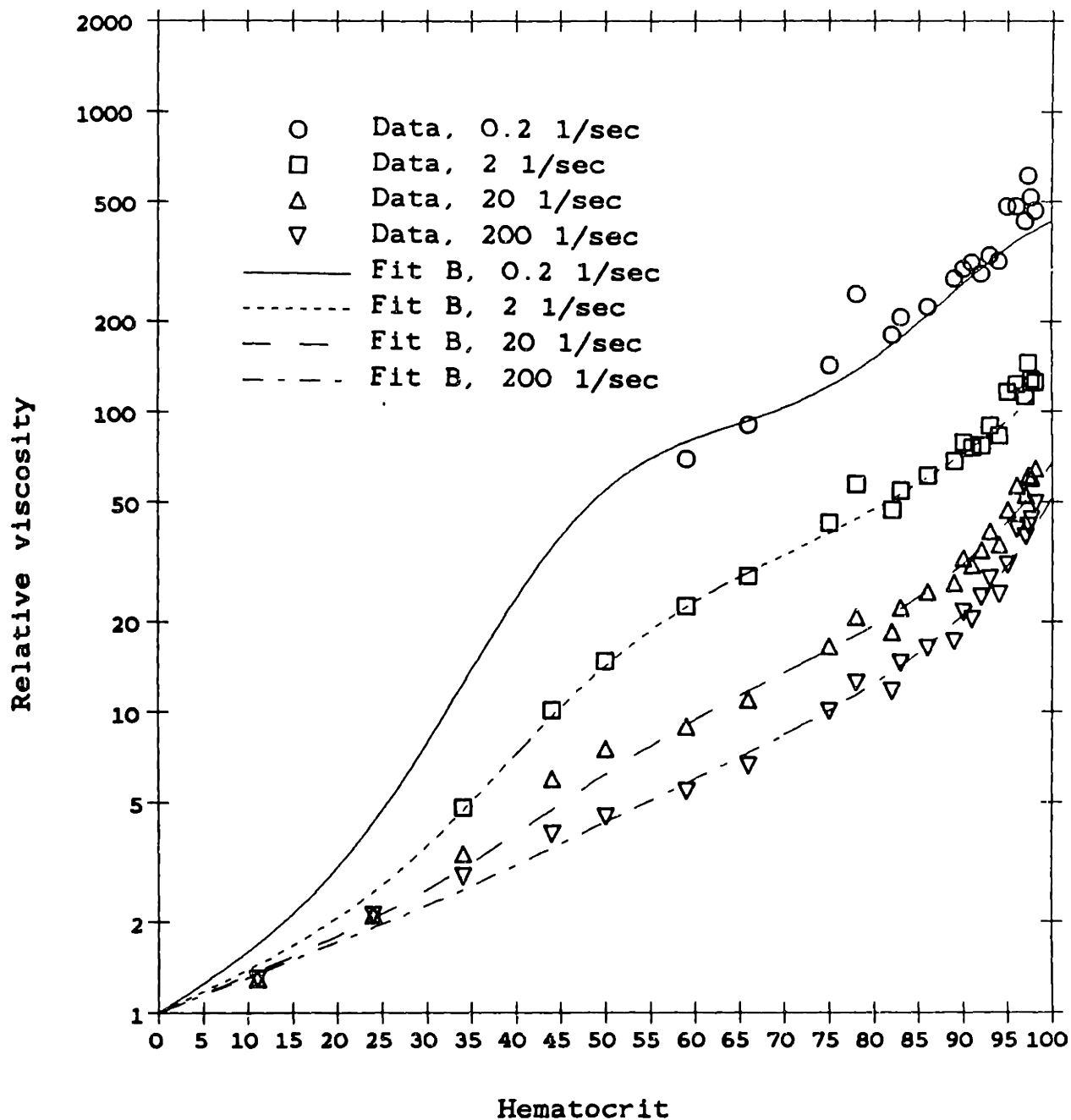


Figure 5-16: Comparison of plasma global fit to data as a function of hematocrit.

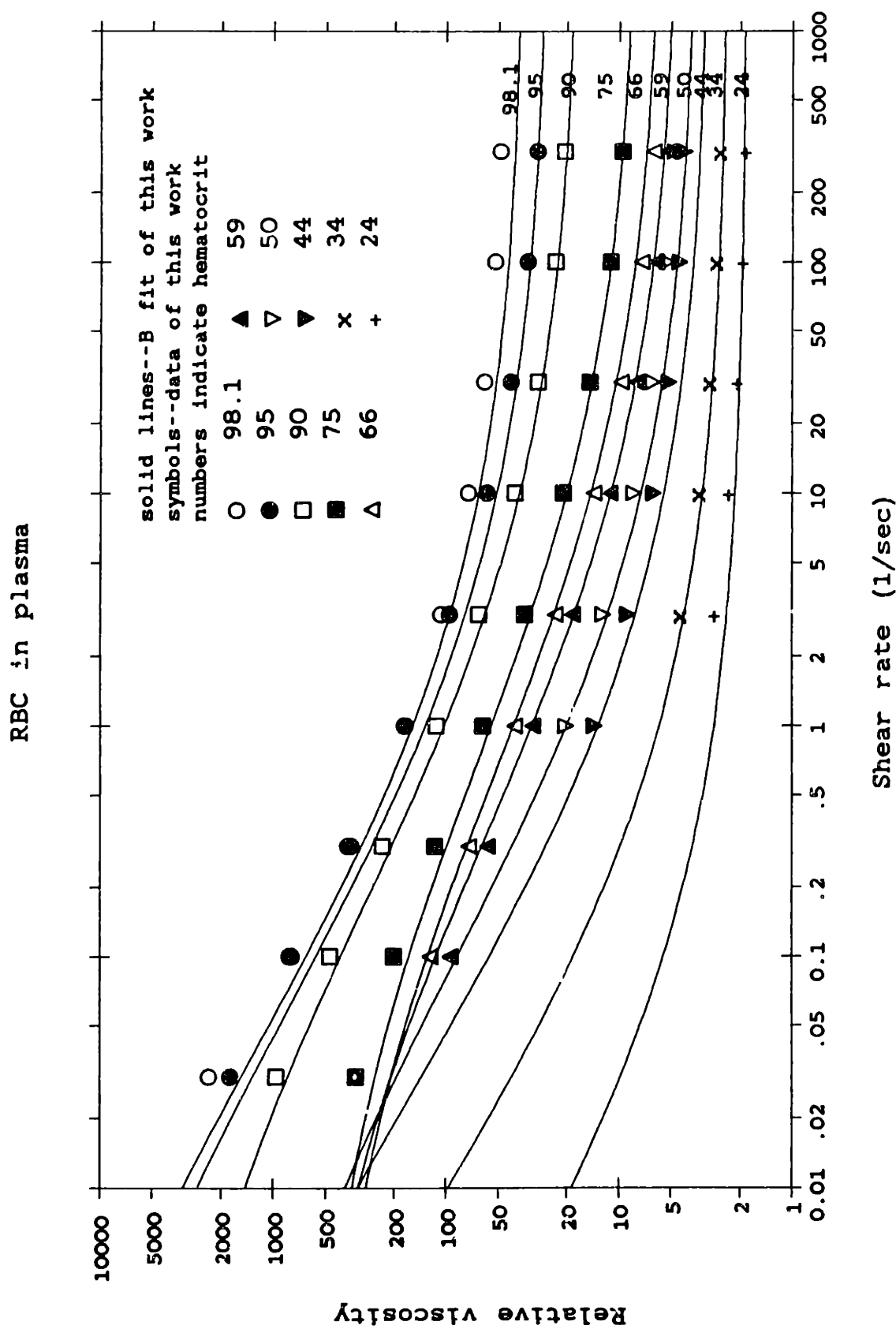


Figure 5-17: Comparison of plasma global fit to data as a function of shear rate.

RBC in saline

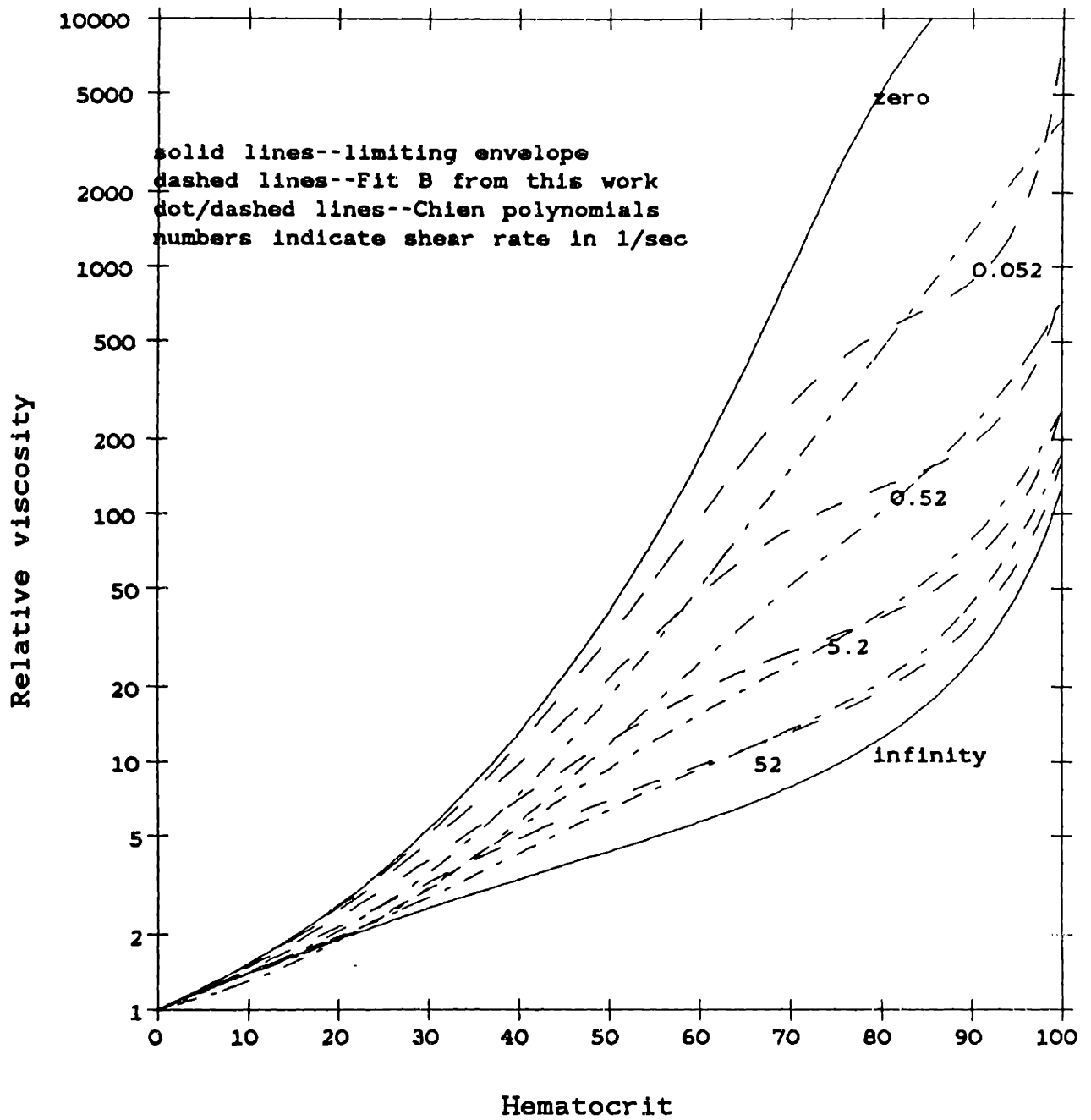
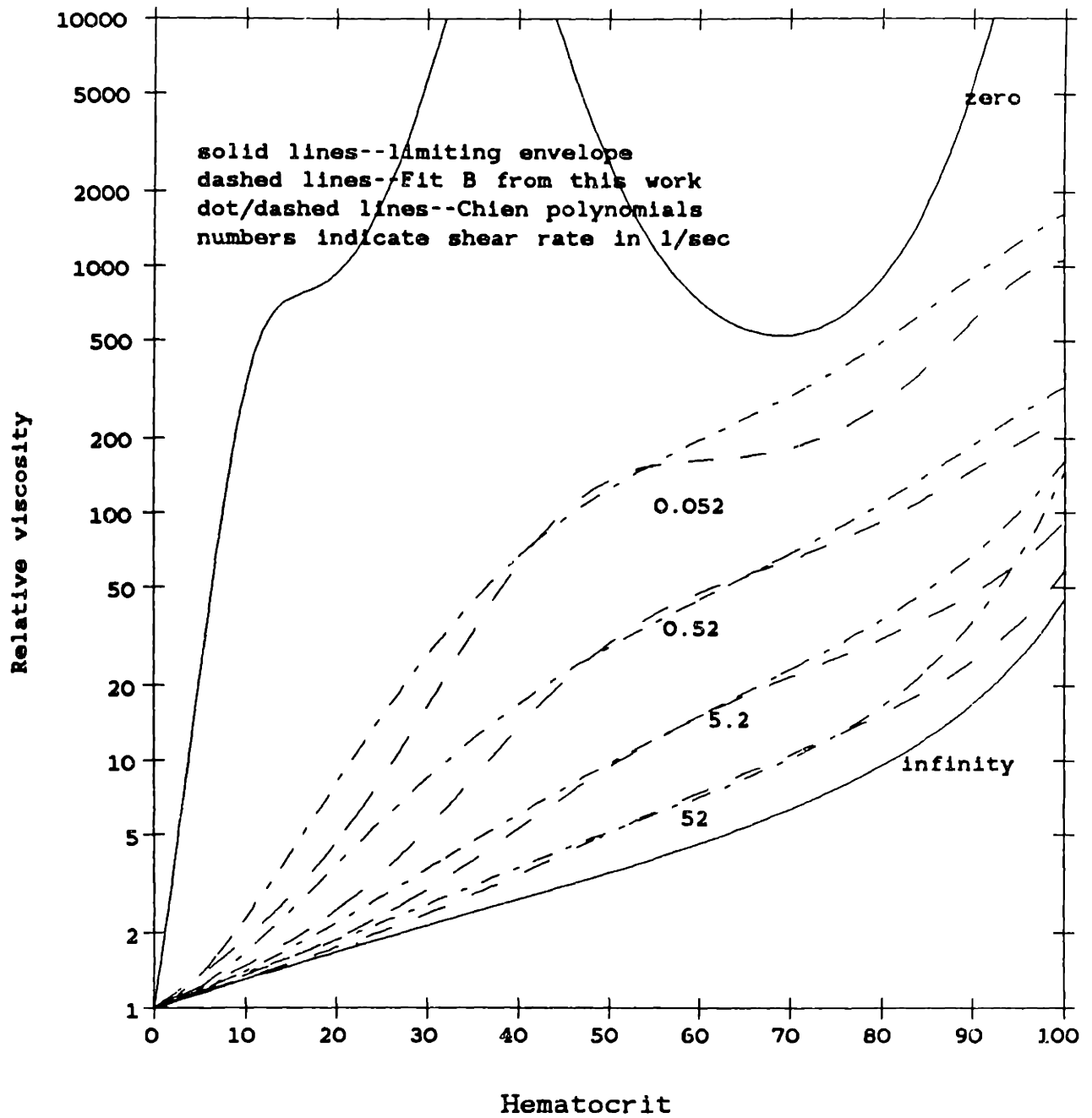


Figure 5-18: Comparison of saline global fit to Chien polynomials.



RBC in plasma



**Figure 5-10:** Comparison of plasma global fit to Chien polynomials.

# Absolute viscosity Comparison of B fits for saline and plasma

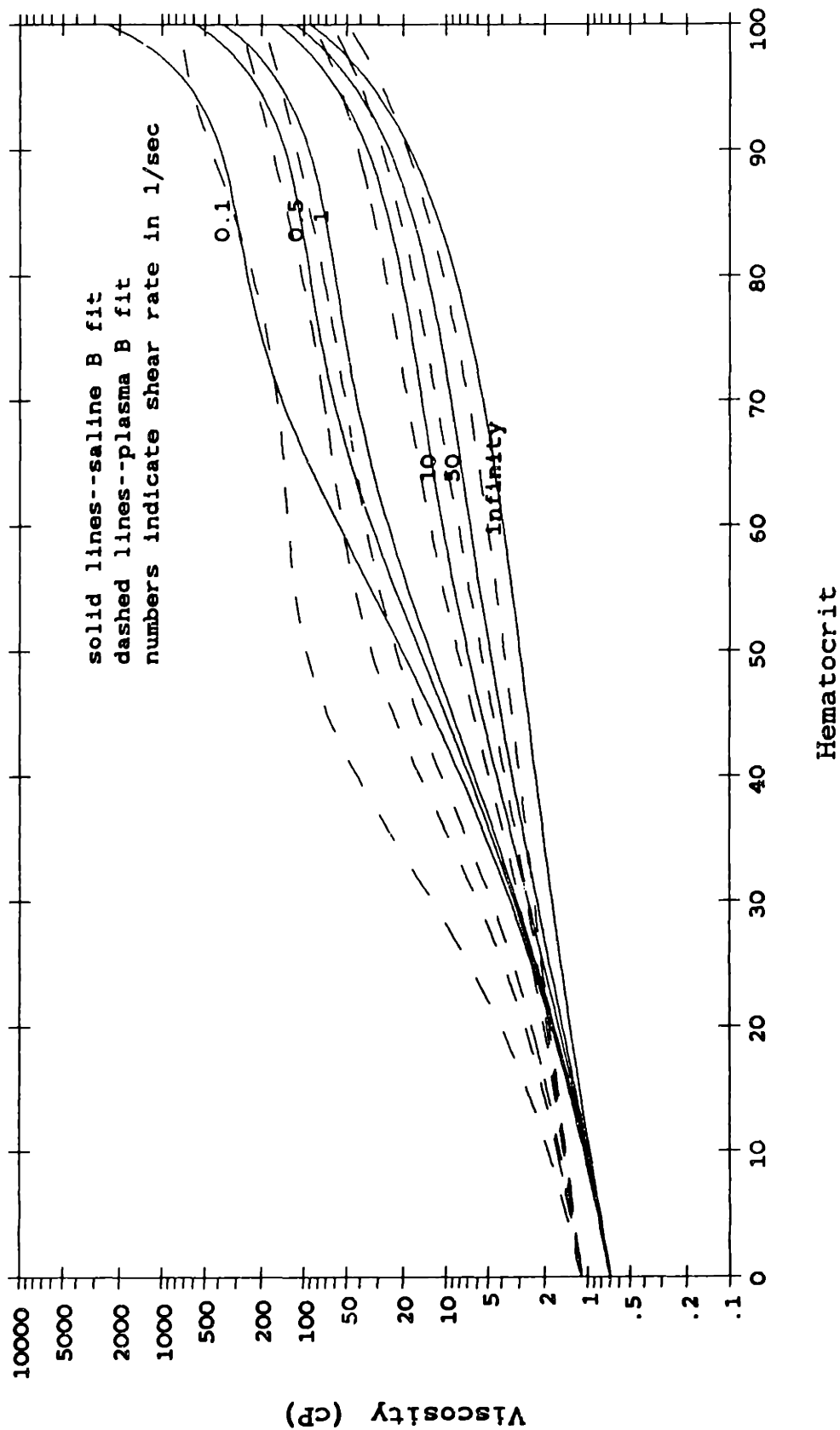


Figure 5-20: Comparison of B fit curves for plasma and saline.

The B fit curves for plasma and saline are compared at various shear rates in Figure 5-20. In the limit of high hematocrit, the B fit curves begin to approach each other in a manner similar to the data (see Figure 5-21). Plasma suspensions maintain slightly higher viscosities until about  $H=92-95$  when the curves cross. This crossing is primarily due to the saline data at  $H=98.4$ , where the viscosity was much higher than for any other sample.

The final global fits (B fits) are summarized in Table 5-VII. The equations are plotted alone as functions of hematocrit and shear rate in Figures 5-22 and 5-23 for saline and in Figures 5-24 and 5-25 for plasma.

## **5.7 Summary, conclusions, and recommendations**

Data collected at hematocrits above 90 shows the still-fluid nature of packed cells, emphasizing that blood is more like an emulsion than a suspension of rigid particles. At high hematocrits, the difference in viscosity between cells in plasma and cells in saline solution is small, but it increases at lower hematocrits and lower shear rates, indicating the effects of the plasma proteins on aggregation.

The complex interactions of blood cells and plasma, summarized in Section 2.4, constitute the challenge of developing theoretical expressions for blood rheology. To date, theory can only be taken so far, and then at some point one must incorporate empirical fits into the theory.

The three-parameter Quemada equation has been fit individually to data at each hematocrit and has proven to be an effective model on this basis. The parameters were empirically fitted to functions of hematocrit, and in this way a set of equations has been derived which can be used with the Quemada expression to represent blood viscosity over a wide range of hematocrits and shear rates,

Absolute viscosity at high hematocrits  
Comparison of B fits for saline and plasma

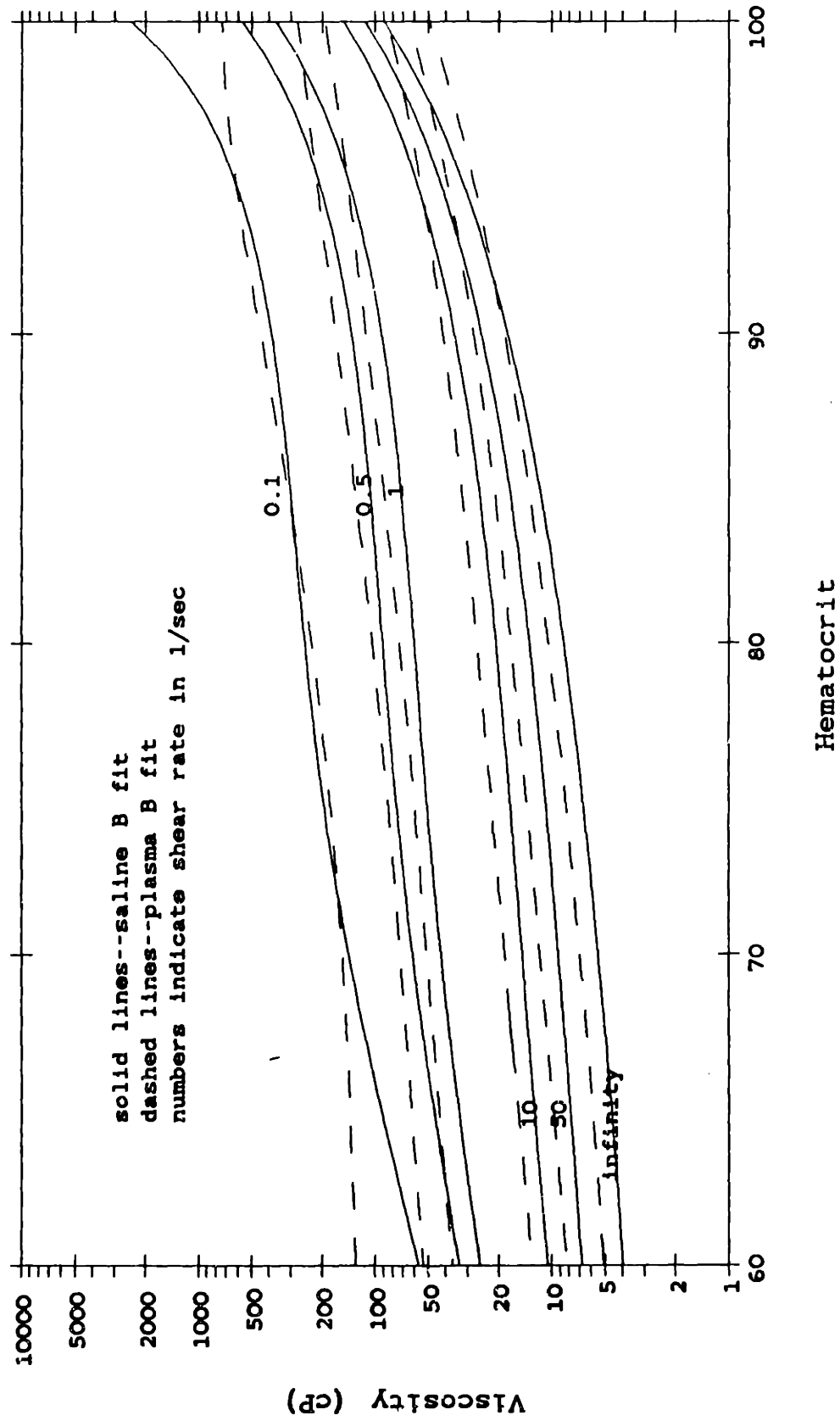


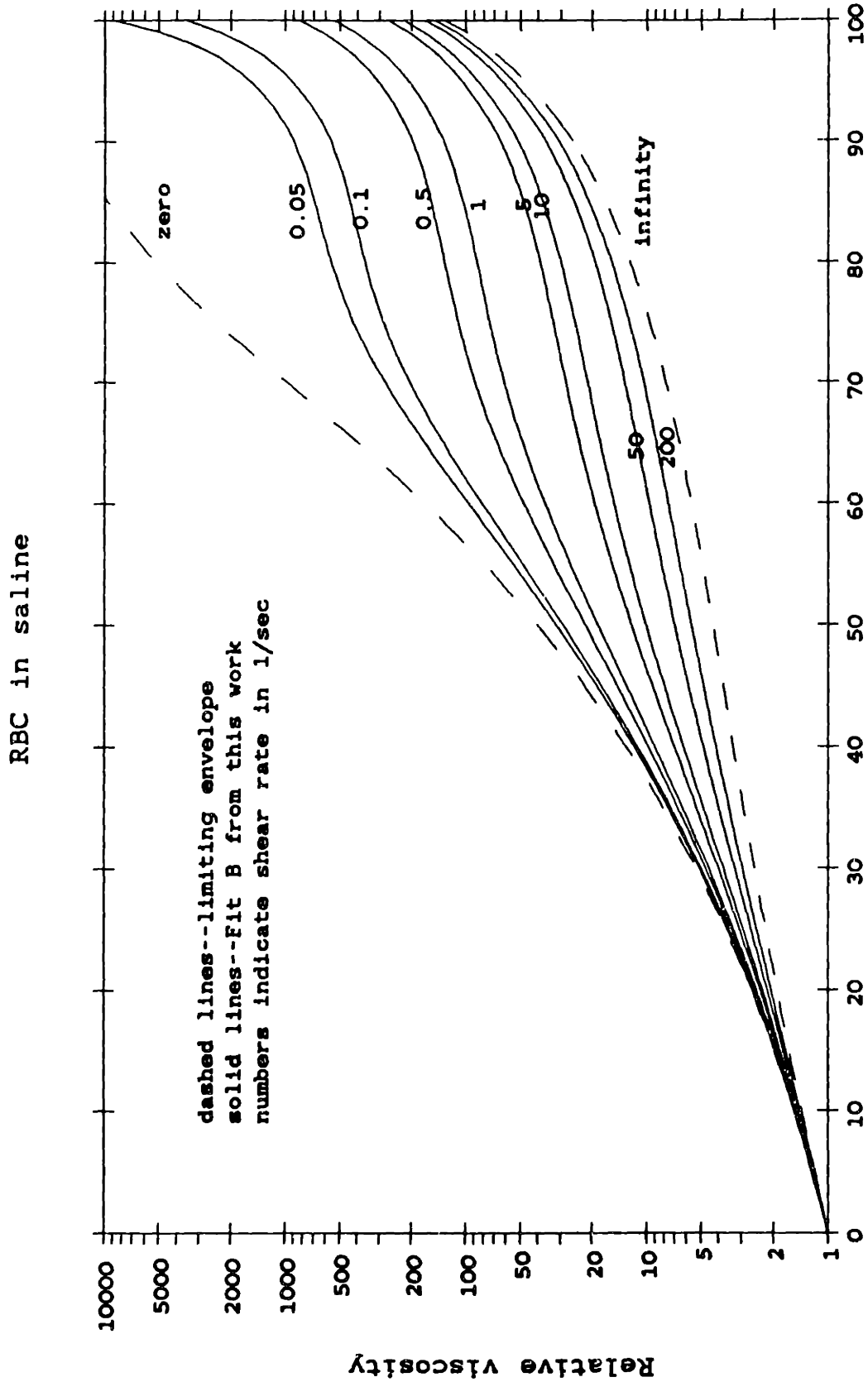
Figure 5-21: High hematocrit viscosities from the B fit for saline and plasma.

QUEMADA EQUATION

$$\eta_r = \left[ 1 - \frac{c}{2} \left( k_\infty + \frac{k_0 - k_\infty}{1 + \sqrt{\gamma/\gamma_c}} \right) \right]^{-2}$$

Parameter Equation	Constants	
	saline	plasma
$k_\infty = \sum_{i=1}^3 k_{\infty i} \cdot c^{(i-1)}$	$k_{\infty 1} = 3.458$ $k_{\infty 2} = -3.894$ $k_{\infty 3} = 2.260$	$k_{\infty 1} = 2.688$ $k_{\infty 2} = -2.300$ $k_{\infty 3} = 1.314$
$k_0 = \sum_{i=1}^4 k_{0i} \cdot c^{(i-1)} + k_{05} \exp(k_{06} \cdot c)$  (Note: for saline and $c < 20$ , $k_0 = 3.838$ )	$k_{01} = 3.503$ $k_{02} = 3.467$ $k_{03} = -9.976$ $k_{04} = 5.026$ $k_{05} = 0$ $k_{06} = 0$	$k_{01} = 13.28$ $k_{02} = -31.93$ $k_{03} = 32.13$ $k_{04} = -11.49$ $k_{05} = 45.46$ $k_{06} = -16.78$
saline $\gamma_c = \sum_{i=1}^3 \gamma_{ci} \cdot c^{(i-1)}$  plasma $\gamma_c = \gamma_{c1} \exp(-(\gamma_{c2} c + \gamma_{c3}))^2$	$\gamma_{c1} = 33.16$ $\gamma_{c2} = -54.79$ $\gamma_{c3} = 22.33$  $\gamma_{c1} = 5.979$ $\gamma_{c2} = -4.536$ $\gamma_{c3} = 3.232$	

**Table 5-VII:** Summary of global fits for saline and plasma suspensions.



**Figure 5-22:** Plot of saline global equations versus hematocrit.

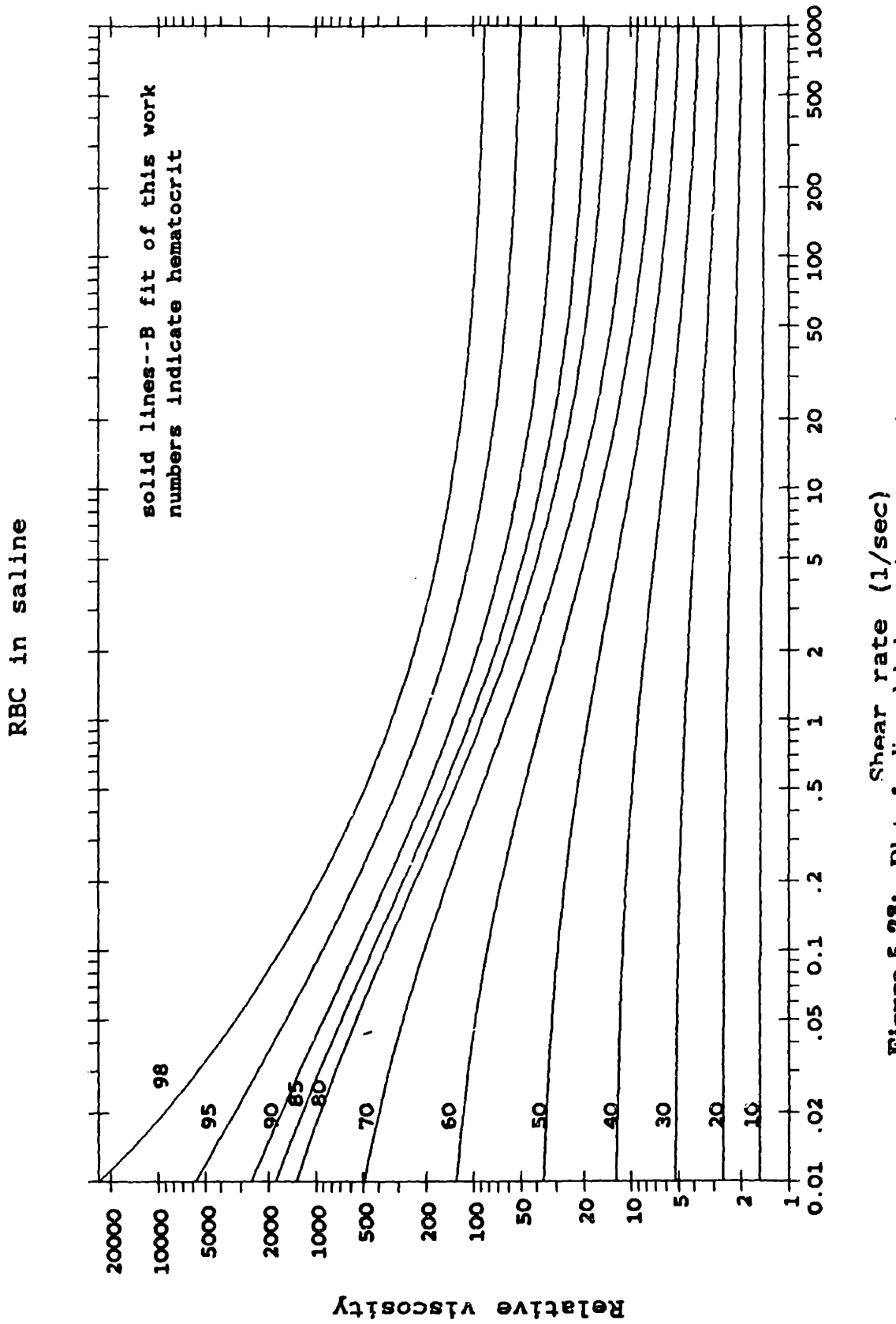
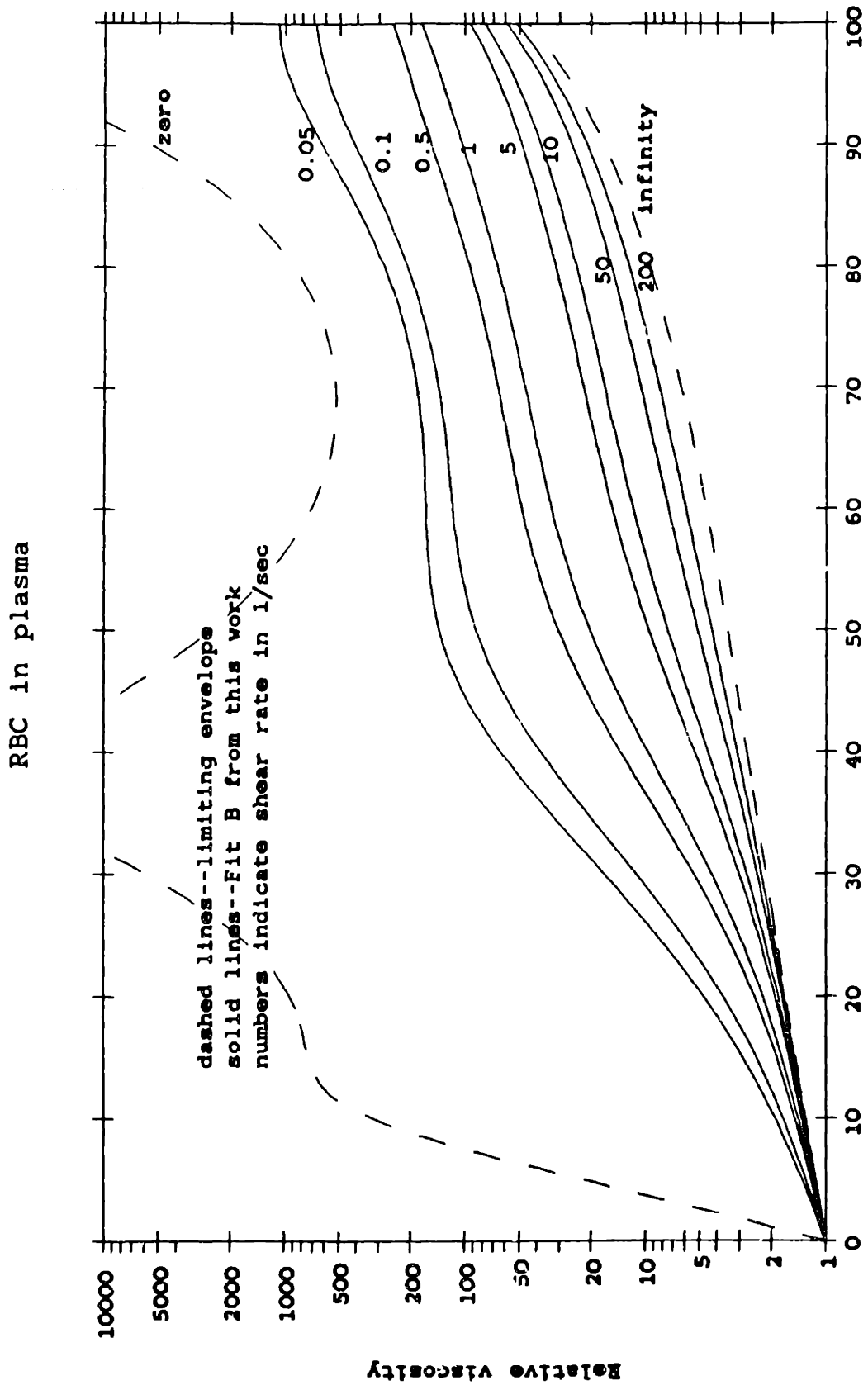


Figure 5-23: Plot of saline global equations versus shear rate.



Hematocrit

Figure 5-24: Plot of plasma global equations versus hematocrit.



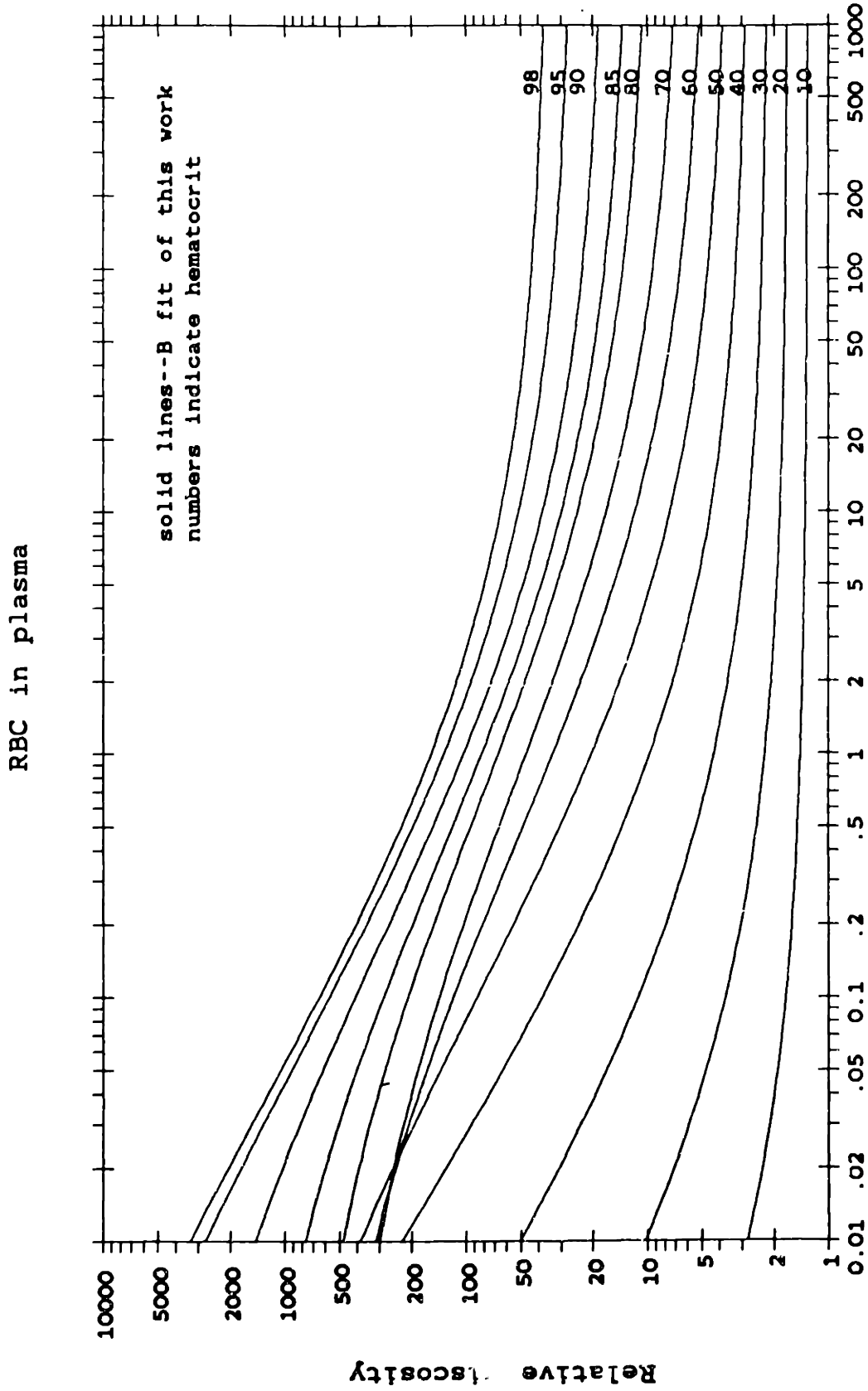


Figure 5-25: Plot of plasma global equations versus shear rate.

particularly those of interest for cross-flow membrane plasmapheresis phenomena. The data of this work did not extend to low enough shear rates for plasma suspensions to reasonably model viscosity dependence at  $\dot{\gamma} < 0.1 \text{ sec}^{-1}$  globally.

The difficulties in extending the Quemada equation to all hematocrit lies in the extreme sensitivity of the equation to the value of  $k_0$  at low shear rates. Ideally, determination of the parameters requires data at only three shear rates:  $\dot{\gamma} = 0, \infty$ , and  $\dot{\gamma}_c$ , which is defined as the point where  $k = \frac{k_\infty + k_0}{2}$ . It would be interesting to see if a fit to only three shear rates could actually model behavior everywhere. Data from a device such as a wide-gap Couette viscometer should be obtained to extend the low shear rate utility of the equations.

In extending the Quemada equation to all hematocrits and shear rates, the present development uses ten constants for saline solution and twelve constants for plasma suspensions. The complexity of this fits seems reasonable when one considers that Chien *et al.* required six constants to model viscosity behavior over all hematocrits at a single shear rate. Interpretation of the meaning of the parameter dependence suffers from the fact that they are conceptual properties than physically measureable quantities.

The Quemada expression does an overall satisfactory job of describing blood viscosity and is probably the most promising way to develop a set of global equations for every hematocrit and every shear rate. A major drawback is that it does not flatten to its limiting high shear rate value as quickly as one would like. One adjustment that could be made is in the power dependence of the square root term, which was obtained from an empirical representation  $\theta = \dot{\gamma}_r^p$  where  $p = 0.5$  (see Appendix C). Perhaps a larger value could be assigned to  $p$ , or  $p$  could be allowed to vary with hematocrit.

The hematocrit dependence of the two limiting viscosities shows an added variation due to the suspending fluid viscosity. A possible modification to the equation is to include the emulsive properties of erythrocytes into the idea of the intrinsic viscosity and thus make the rheological parameters independent of  $\eta_f$ . This can be done by incorporating some sort of Taylor coefficient, with a simplifying assumption of constant internal viscosity.

Most formulations for representing viscosity are for  $\eta_r$ , and thus there is an implicit dependence of  $\eta$  on  $\eta_f$  at all hematocrits. Ideally,  $\eta$  should become independent of  $\eta_f$  at high hematocrits, and a truly general expression for viscosity should provide for this. With regards to this, the Quemada model seems to indicate that cell crowding effects and inversion becomes important at hematocrits of 65 to 80 and above, as this seems to be the region where all three rheological parameters become similar for saline and plasma suspensions.

## Appendix A

### Physical Properties of Human Blood and the Red Blood Cell

#### Properties of whole blood

##### Whole blood

specific gravity(25/4 °C)	1.056
---------------------------	-------

##### Plasma

specific gravity(25/4 °C)	1.0239
---------------------------	--------

viscosity (37 °C)	1.2 cP
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##### Plasma proteins

###### albumin

m.w. 69,000

4.5 g/100 ml plasma

###### globulins

m.w. 35,000-1,000,000

2.5 g/100ml plasma

###### fibrinogen

m.w. 330,000

0.3 g/100ml plasma

##### Erythrocytes

###### concentration

###### male

$5.2 \times 10^6 / \text{mm}^3$  whole blood

42 volume %

###### female

$4.7 \times 10^6 / \text{mm}^3$  whole blood

38 volume %

##### Leukocytes

###### concentration

$7.0 \times 10^3 / \text{mm}^3$  whole blood

0.1 to 0.5 volume %

##### Platelets

###### concentration

$3.0 \times 10^5 / \text{mm}^3$  whole blood

0.1 to 0.6 volume %

## **Properties of Red Blood Cell**

### **Red blood cell**

specific gravity	1.098
mean corpuscular volume	85-100 $\mu\text{m}^3$
surface area	140 $\mu\text{m}^2$
diameter	8.4 $\mu\text{m}$
center thickness	1 $\mu\text{m}$
peripheral thickness	2.4 $\mu\text{m}$
mean corpuscular hemoglobin concentration	31-35 g/100 ml packed red cells

*Sources:* Beck [1981] and Guyton [1981].

## **Appendix B**

### **Modified Eagle's Solution Recipe**

In one liter of distilled water, dissolve

NaCl	6.2 gm
KCl	0.36 gm
NaH <sub>3</sub> PO <sub>4</sub> ·H <sub>2</sub> O (sodium phosphate monobasic)	0.13 gm
NaHCO <sub>3</sub>	2.0 gm
CaCl <sub>2</sub>	0.18 gm
MgCl <sub>2</sub> ·6H <sub>2</sub> O	0.15 gm
dextrose	0.9 gm

Adapted from Eagle [1959].

## Appendix C

### Derivation of Quemada Equation

Quemada [1977, 1978a, 1978b] examined the steady-state viscometric flow of concentrated disperse systems. He noted that in such flows the concentration of the dispersed phase is not uniform but that an approximation to two-phase flow is developed: an axial region of high concentration and a region near the wall of low concentration, which in a sense provides a region of lubricant flow. What follows here describes his application of the principle of minimum energy dissipation, which states that with fixed system properties (concentration, temperature, flow conditions, fluid characteristics), the steady-state solution is that which minimizes the rate of energy dissipation (or, more generally, the rate of entropy production). The development assumes many features which are strictly valid only for suspensions of uniform, non-deformable rigid spheres but, as seen in the text (Section 2.4), with appropriate modifications the final result correlates well with experimental data for blood suspensions.

Consider simple viscometric flow in direction  $x$  between two plates separated by a distance  $2Y$  along axis  $y$ , where  $v_y = v_z = 0$ ,  $v_x = v_x(y)$  and no-slip conditions at the walls. The concentration profile in the channel is  $c = c(y)$ . The viscous dissipation rate term in the conservation of momentum equation is [Bird *et al.*, 1960]

$$\left. \frac{DU}{Dt} \right|_{\text{viscous}} = \underline{\tau} : \nabla \mathbf{v} = \frac{1}{2} \eta \sum_i \sum_j \left[ \left( \frac{\partial v_i}{\partial x_j} + \frac{\partial v_j}{\partial x_i} \right) - \frac{2}{3} (\nabla \cdot \mathbf{v}) \delta_{ij} \right]^2 \quad (C.1)$$

where  $U$  is the internal energy,  $\frac{D}{Dt}$  is the substantial time derivative,  $\underline{\tau}$  is the stress tensor, and  $\delta_{ij}$  is the Kronecker delta function. For viscometric flow with an incompressible fluid,  $(\nabla \cdot \mathbf{v}) = 0$ , and this reduces to

$$\tau \cdot \nabla \mathbf{v} = \frac{1}{2} \eta \left( \frac{\partial v_x}{\partial y} \right)^2 \quad (C.2)$$

Generalizing this to include dissipation from mass fluxes and integrating over the surface element  $s_x$  normal to the x-direction gives

$$\left. \frac{DU}{Dt} \right|_{\text{viscous}} = \int_{s_x} \left[ \frac{1}{2} \eta \left( \frac{\partial v_x}{\partial y} \right)^2 + \frac{1}{2} D \left( \frac{\partial c}{\partial y} \right)^2 \right] ds_x \quad (C.3)$$

where  $D$  is the diffusivity of the dispersed particles.

Equation (C.3) is to be minimized subject to the following constraints:

1. Flow rate constraints

$$\text{Fluid flow constraints } Q_f = \int_{s_x} (1-c) \cdot v \cdot ds_x$$

$$\text{Particle flow constraints } Q_p = \int_{s_x} c \cdot v \cdot ds_x$$

2. Mean particle concentration

$$\bar{c} = \frac{\int_{s_x} c \cdot ds_x}{s_x}$$

3. Upper and lower limits of concentration,  $c_w \leq c \leq c_b$ . Using a method

of optimum control theory, this constraint can be written as

$$\psi^2 = (c_b - c)(c - c_w)$$

so that  $\psi$  has an imaginary solution when  $c$  lies outside of the limits.

Equation (C.3) is minimized by the use of Lagrange multipliers [Wylie, 1975].

The auxiliary function to be solved is

$$F = \int_{s_x} \left\{ \frac{1}{2} \eta v_x'^2 + \frac{1}{2} D c'^2 + \lambda_1 (1-c) v_x + \lambda_2 c v_x + \lambda_3 c + A(x) [\psi^2 - (c - c_b)(c - c_w)] \right\} ds_x \quad (C.4)$$

where  $\lambda_1$ ,  $\lambda_2$ ,  $\lambda_3$ , and  $A(x)$  are the multipliers, and the "prime" symbol denotes



partial differentiation with respect to  $x$ . Solution of  $F$  for extrema satisfies the set of equations:

$$\frac{\partial}{\partial x} \left( \frac{\partial F}{\partial f'} \right) - \frac{\partial F}{\partial f} = 0 \quad (C.5)$$

where  $f = v_x, c$ , and  $\psi$ .

At this point Newtonian behavior in the viscometer is assumed and the equations obtained are

$$(\eta v_x')' = \lambda_1(1-c) + \lambda_2 c \quad (C.6)$$

$$(Dc')' = \frac{1}{2} v_x^2 \frac{\partial \eta}{\partial c} + (\lambda_2 - \lambda_1) v_x + \lambda_3 + A(2c - c_b - c_w) \quad (C.7)$$

$$0 = A\psi \quad (C.8)$$

Quemada [1977] notes that other works using the principle of minimum energy dissipation have yielded similar sets of equations. Equation (C.6) can be considered as a generalization of the momentum transport equation, since it reduces to the Navier-Stokes equations for the case of a homogeneous ( $c=0$ ) incompressible fluid, with  $\lambda_1$  equal to the pressure gradient. Equation (C.7) is a generalized mass transport equation, and equation (C.8) provides the boundaries of the admissible domain.

Further solution is based on the observation that the transition between the concentration in the bulk and the concentration near the wall is relatively sharp. Also, for particles with radius  $\geq 1 \mu m$ , Brownian diffusion is negligible compared to viscous drag.

Equations (C.6) and (C.7) are now decoupled by setting the left hand side of equation (C.7) to zero and using an empirical "rectangular" concentration profile similar to the one used by Vand [1948] to describe the viscometric flow of a suspended solution of rigid spheres:

$$c = \begin{cases} c_b & \text{for } 0 \leq y \leq \beta Y \\ c_w & \text{for } \beta Y \leq y \leq Y \end{cases}$$

where  $\beta < 1$ .

Satisfaction of extrema at the break point  $y_b = \beta Y$  is accounted for by Erdmann-Weierstrass corner conditions [Tolle, 1975]:

$$\begin{aligned} \frac{\partial F}{\partial f'_i} \Big|_{y_{b+}} &= \frac{\partial F}{\partial f'_i} \Big|_{y_{b-}} \\ F - \sum_i \frac{\partial F}{\partial f'_i} f'_i \Big|_{y_{b+}} &= F - \sum_i \frac{\partial F}{\partial f'_i} f'_i \Big|_{y_{b-}} \end{aligned}$$

Additionally, the velocity must be continuous at  $y_b$ . Defining  $\Delta f = f|_{y_{b+}} - f|_{y_{b-}}$ , these conditions become

$$\Delta v_x = 0 \quad (C.9)$$

$$\Delta(\eta v'_x) = 0 \quad (C.10)$$

$$\Delta(Dc') = 0 \quad (C.11)$$

$$\Delta \frac{\partial F}{\partial \psi'} = 0 \quad (C.12)$$

$$\Delta \left( \frac{1}{2} \eta v_x'^2 - (\lambda_1(1-c)v_x + \lambda_2 c v_x + \lambda_3 c) \right) = 0 \quad (C.13)$$

These equations can now be used to solve for the velocity profile. Assuming two-phase flow with  $c_b, v_b$  in the core and  $c_w, v_w$  near the wall and applying equations (C.6), (C.9), and (C.10), the result is the same as that derived for any two-phase Newtonian flow:

$$v_w(y) = \frac{Y^2 \Delta p}{4L\eta_w} \left( 1 - \frac{y^2}{Y^2} \right) \quad (C.14)$$

$$v_b(y) = \frac{Y^2 \Delta p}{4L\eta_w} (1 - \beta^2) + \frac{Y^2 \Delta p}{4L\eta_b} \left( \beta^2 - \frac{y^2}{Y^2} \right) \quad (C.15)$$

Solving equations (C.7) and (C.13) for the concentration relationship gives

$$\left( \frac{d\phi}{dc} \right)_b + \left( \frac{d\phi}{dc} \right)_w = \frac{2}{c_b - c_w} (\phi_b - \phi_w) \quad (C.16)$$

where  $\phi$  is a relative viscosity based on the pure fluid viscosity,  $\phi \equiv \eta_f/\eta \equiv 1/\eta_r$ . In the limiting case where  $c_w \rightarrow 0$ ,

$$\frac{d\phi}{dc} - k(c, \gamma) = \frac{2}{c}(\phi - 1)$$

where  $k \equiv -\left(\frac{d\phi}{dc}\right)_{c \rightarrow 0} = \left(\frac{d\eta_r}{dc}\right)_{c \rightarrow 0}$  is the intrinsic viscosity.

The intrinsic viscosity is a commonly used empirical parameter in polymer rheology. Deviations from the Einstein value  $k=2.5$  from equation (2.9) are used for determination of dispersion characteristics such as polydispersity, solvation, particle ellipticity, and molecular weight (see, for example, Hiemenz [1977], Chapter 2).

Equation (C.17) is integrated using the condition that in the limit as  $c$  approaches its packing concentration  $c_p$ ,  $\eta \rightarrow \infty$ ,  $\frac{d\eta}{dc} \rightarrow \infty$  and the result is

$$\eta_r = \frac{1}{\phi} = \left(1 - \frac{1}{2}kc\right)^{-2} \quad (C.18)$$

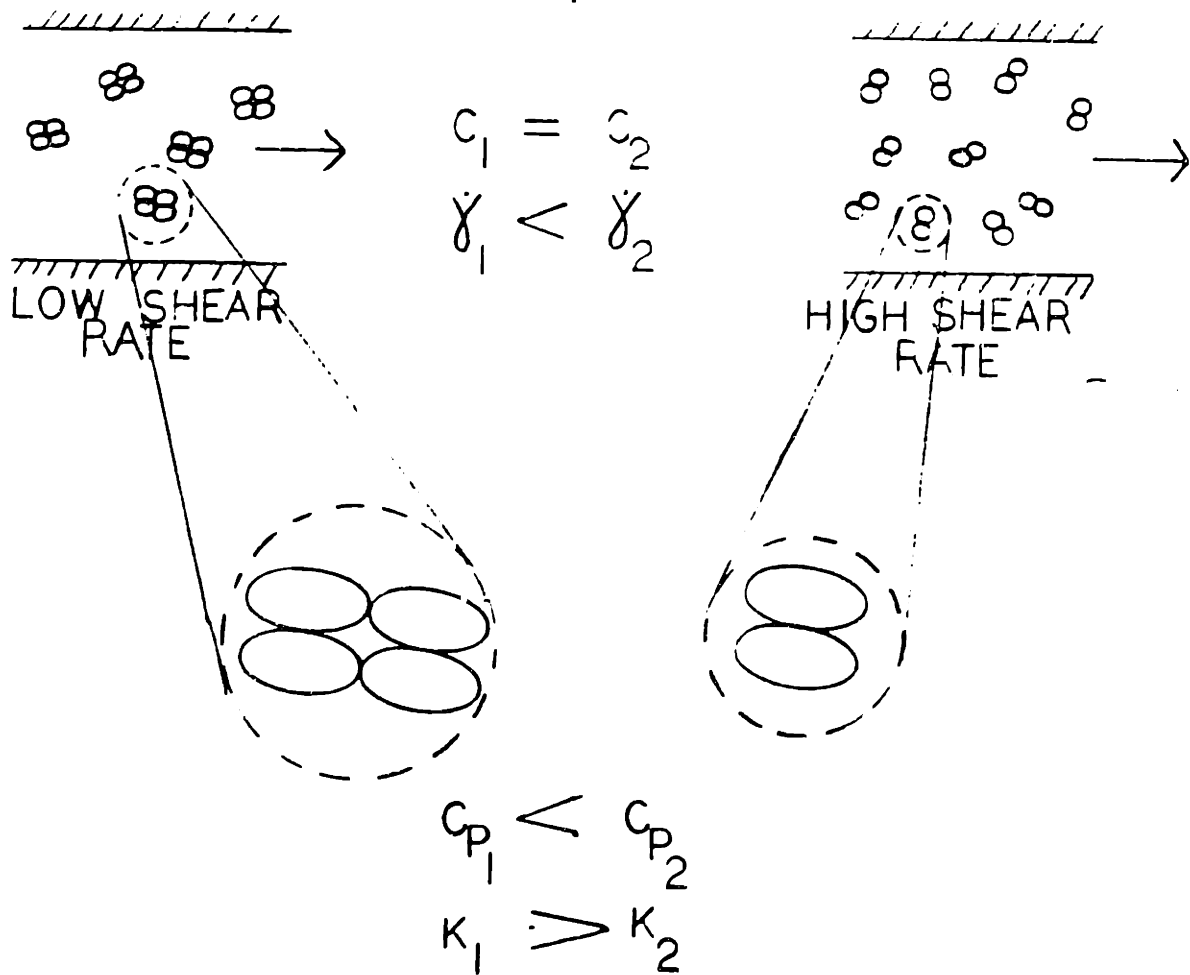
This expression is rigorously valid only for Newtonian viscosities, and the boundary condition as  $c \rightarrow c_p$  is true only for rigid particles. Quemada argues, however, that in highly concentrated systems the two-phase flow model is still valid and that with appropriate modifications this equation is applicable to blood suspensions. The form of equation (C.18) is very similar to that of other equations modified from the Einstein equation, (2.9).

The non-Newtonian aspects of blood are related to changes in the particle morphology (from aggregation, plasma trapping, deformation, crowding, etc.) at changing concentration and shear rate. Equation (C.18) can be empirically modified to incorporate these effects in one of two ways:

First,  $k$  is defined as before, but  $c$  is now the volume concentration of *effective* particles, where an effective particle is defined as a statistically-averaged "typical" aggregation of individual units, including any trapped fluid. If cells aggregated without deforming or trapping fluid, then  $c$  would be constant no matter

the degree of aggregation. With fluid trapping, aggregation would lead to effective particles with a larger volume than the individual cells, and  $c$  would increase by a factor of  $(1-\epsilon)^{-1}$ , where  $\epsilon$  is the porosity of the effective particle.

The alternative and more convenient method is to define  $c$  as before, equal to the volume percent of solids, but to now empirically incorporate the effective volume changes into  $k$ . To visualize this, two samples of suspensions of rigid, attracting particles can be imagined, each at identical concentration  $c$ , one at shear rate  $\dot{\gamma}_1$  and the other at  $\dot{\gamma}_2 > \dot{\gamma}_1$ .



**Figure C-1:** Hypothetical comparison of two samples with identical solid concentrations at different shear rates.

As shown in Figure C-1, at  $\dot{\gamma}_1$ , the average effective particle consists of four

units, and at  $\gamma_2$ , the average effective particle is two units. Because of the extra fluid trapping, the particles at  $\gamma_1$  can not be packed as tightly as those at  $\gamma_2$ . If the packing concentration  $c_p$  is defined as the maximum volume concentration of solids achievable by packing the effective particles as closely as possible, then  $c_{p1} < c_{p2}$ . Also, as  $c \rightarrow c_p$ ,  $\eta \rightarrow \infty$  and from equation (C.18),

$$c_p = \frac{2}{k(c, \gamma)} \quad (C.19)$$

Since in the above example  $c_1 = c_2$ , the effect of increasing  $\gamma$  at constant  $c$  is to lower  $k$ , consistent with shear-thinning behavior. If the individual units could deform elastically and orient themselves along fluid streamlines, these effects would result in a further reduction of  $k$  with shear rate.

Thus, for relatively rigid particles, the intrinsic viscosity is proportional to the volume occupied by a single effective particle. For fluid particles such as red blood cells, equation (C.18) and especially equation (C.19) are true only in a qualitative sense, as the boundary condition of infinite viscosity at high concentration is obviously not justifiable. Nevertheless, Quemada [1981] has used this model to estimate the size and structure of red cell aggregates as a function of hematocrit and shear rate.

Continuing with the development for rigid spheres, it is at this point necessary to address the concentration and shear rate dependence of the intrinsic viscosity  $k$ . Simplifying to the case of constant concentration, Quemada proposes an expression for  $k$  based on a two limit relaxation model. Defining  $k_0 = k(\gamma = 0)$  and  $k_\infty = k(\gamma \text{ large, but not high enough to disrupt individual units})$ , the model assumes aggregate formation dependent on collision frequency and aggregate disruption by shear. At a constant concentration,

$$\text{for aggregation, } \left(\frac{dk}{dt}\right)_A = \frac{1}{\tau_A}(k_0 - k) \quad (C.20)$$

$$\text{for disaggregation, } \left(\frac{dk}{dt}\right)_D = \frac{1}{\tau_D}(k - k_\infty) \quad (C.21)$$

$$\frac{dk}{dt} = \frac{1}{\tau_A}(k_0 - k) - \frac{1}{\tau_D}(k - k_\infty). \quad (C.22)$$

where  $\tau_A$  and  $\tau_D$  represent the time constants for aggregation and disaggregation.

At steady-state,

$$k = \frac{k_0 + \theta k_\infty}{1 + \theta} \quad (C.23)$$

where  $\theta = \tau_A / \tau_D$ .

For dilute suspensions of rigid spheres, the time constant for aggregation,  $\tau_A$ , is assumed to be the inverse of the Brownian collision frequency as derived by Smoluchowski [1917]:

$$\tau_A = \frac{\pi \eta a^3}{2c\kappa T} \quad (C.24)$$

where  $a$  is the sphere radius,  $\kappa$  is Boltzmann's constant, and  $T$  is absolute temperature. This expression is comparable to both the time constant for translation of the particle over a distance equal to its radius and to the time constant for rotation of a sphere through  $\pi$  radians.

The time constant for disaggregation due to shear,  $\tau_D$ , is assumed to be the inverse of the shear-induced collision frequency developed by Goldsmith and Mason [1967]:

$$\tau_D = \frac{\pi \gamma}{8c} \quad (C.25)$$

Combining equations (C.24) and (C.25),

$$\theta = \gamma_r \quad (C.26)$$

where the relative shear rate  $\gamma_r = \gamma / \gamma_c$  and the critical shear rate  $\gamma_c$  is  $\frac{\kappa T}{4\eta a^3}$ . For

dilute suspensions of rigid spheres  $\gamma_c$  is roughly equivalent to the rotational diffusion coefficient. Equivalent ratios to  $\theta$  appear frequently in the literature of polymer rheology. Batchelor [1976] described this ratio as a Peclet number  $\frac{\gamma a^2}{D}$  which represents the time constant for diffusional rotation divided by the time constant for shear-induced motion.

For non-rigid particles such as blood, the relaxation model must also incorporate time constants for changes in particle shape. Assuming that a basic functional relationship  $\theta(\gamma_r)$  exists as it does for rigid particles, it then becomes necessary to formulate an expression for  $\gamma_c$  for deformable bodies. Using the Maxwell relaxation time, Quemada suggests

$$\gamma_c = \frac{E_i}{\eta_i} \quad (C.27)$$

where  $E_i$  is the elastic modulus and  $\eta_i$  is the internal viscosity of the particle. Since  $E_i \sim a^3$ , the variance with particle size is the same as for spheres.

Using equation (C.26) as a starting point, Quemada derived a semi-empirical relationship for concentrated suspensions:

$$\theta = \gamma_r^p \quad \text{for } 0 \leq p \leq 1 \quad (C.28)$$

where for blood  $p \approx 0.5$ . Substituting equations (C.23) and (C.28) into equation (C.18) and rearranging leads to the final form of the equation:

$$\eta = \left(1 - \frac{c}{2} \cdot \frac{k_0 + k_\infty \sqrt{\gamma_r}}{1 + \sqrt{\gamma_r}}\right)^{-2} \quad (C.29)$$

where  $k_0$  and  $k_\infty$  are functions of concentration and  $\gamma_r$  is a function of concentration and directly proportional to the shear rate. In this work, the concentration dependences of the three rheological parameters  $k_\infty$ ,  $k_0$ , and  $\gamma_c$  are found from empirical fits to data.

To summarize, the Quemada equation is derived from a model of a concentrated suspension of rigid, attracting spheres in which two-phase viscometric flow is observed. Deviation of particle behavior from this model, such as deformability and fluidity, are incorporated empirically into the parameters.

The Quemada expression does not predict a yield shear stress *per se*, but in the limit of high concentration,  $c \rightarrow c_p$ , equation (C.29) can be shown to yield the Casson relation, equation (2.7) [Quemada, 1981]. In the limit of low shear,  $\dot{\gamma} \ll \dot{\gamma}_c$ , no yield shear stress is obtained and Newtonian behavior is recovered from the equation. Quemada argues that, while it is possible that blood may indeed have a yield shear stress, it is of almost negligible magnitude due to the smallness of the cohesive forces between cells. In the limit of high shear,  $\dot{\gamma} \gg \dot{\gamma}_c$ , a pseudo-yield stress can be extrapolated to zero shear rate.



## **Appendix D**

### **Computer Program for Data Fitting**

This is the FORTRAN computer program described in Section 3.4 for regression of data to non-linear functions. It is adapted from Zydney [1985].

program datafit

```
double precision b(5), bl(5), bo(5), f(5), fl(5), a(5,5),
1 ainv(5,5), yl(200), dy(200,5), y(200), ee, el, ss,
1 xtx(5,5), c(5,5), gamma(200), eta(200), delz(5),
1 zval, zvalnew
character*5 type, hct
```

c ASSIGN DIFFERENTIAL MULTIPLIERS AND CONVERGENCE REQUIREMENTS

```
d1=0.999
d2=.99
c1=.001
c2=.001
```

c READ IN DATA FROM DATAFILE

```
print *, 'enter sample type and hematocrit'
read *, type, hct
open (unit=4,
1 file=type(:len(type))//'dat'//hct(:len(hct))//'.dat',
1 status='old')
do 10 i=1,200
    read (4,501) gamma(i), eta(i)
501    format(2(f10.3))
    if (gamma(i) .eq. -100. .and. eta(i) .eq. -100.) goto 20
10    continue
20    n=i-1
```

c INPUT INITIAL GUESSES

```
print *, 'cell volume fraction and number of unknowns'
read *, h, ir
do 30 i=1,ir
    print *, 'initial guess for unknown #',i,'?'
    read *, b(i)
30    continue
```

c LOOP TO EVALUATE RESIDUALS AND JACOBIAN

```
35    do 130 ii=1, ir+1
        do 80 jj=1, ir+1
c        CALCULATE SUM OF SQUARED RESIDUALS (SS)
c        AND SUM OF WEIGHTED SQUARED RESIDUALS (EE)
            ee=0.
            ss=0.
            do 40 i=1,n
                call predictval(gamma(i),y(i),b,h)
                ee=ee+((y(i)-eta(i))/y(i))**2.
                if (ii .eq. 1 .and. jj .eq. 1)
                    ss=ss+(y(i)-eta(i))**2.
            endif
40        continue
```

c PRINT VALUE OF SUM OF SQUARES BEFORE DERIVATIVES ARE EVALUATED

```

        if (ii .eq. 1 .and. jj .eq. 1) then
            print *, 'sum of weighted sqd residuals=', ee
        endif
c      IF THIS IS THE FIRST PASS, THEN SET LAST VALUES (BL) EQUAL TO
c      CURRENT GUESSES (B) AND DO NOT EVALUATE DERIVATIVE
        if (jj .eq. 1) then
            el=ee
            do 50 i=1,ir
                bl(i)=b(i)
50          continue
            goto 70
c      OTHERWISE, DETERMINE THE DERIVATIVE VALUE (F) FOR THIS B
        else
            f(jj-1)=(ee-el)/(b(jj-1)-bl(jj-1))
c      RESET B VALUES
            do 60 i=1,ir
                b(i)=bl(i)
60          continue
        endif
c      MAKE SMALL CHANGE IN THE NEXT B AND START LOOP TO EVALUATE
c      NEXT F
70          b(jj)=d1*b(jj)
80          continue
c      BEGIN DETERMINATION OF JACOBIAN (A)
c      ON FIRST PASS, MAKE OLD VALUES EQUAL TO CURRENT VALUES
        if (ii .eq. 1) then
            do 90 i=1,ir
                bo(i)=bl(i)
                fl(i)=f(i)
90          continue
            goto 110
c      OTHERWISE, EVALUATE JACOBIAN FOR THIS F AND B
        else
            do 100 i= 1, ir
                a(i,ii-1)=(f(i)-fl(i))/(b(ii-1)-bo(ii-1))
100         continue
        endif
c      RESET B VALUES
110         do 120 i=1,ir
            b(i)=bo(i)
120         continue
c      MAKE DIFFERENTIAL CHANGE IN NEXT B
            b(ii)=d2*b(ii)
130         continue
c      END OF LOOP TO CALCULATE RESIDUALS AND JACOBIAN

c      INVERT JACOBIAN AND RESET B VALUES
        call invert(ir,a,ainv)
        do 140 i=1,ir
            b(i)=bo(i)
140         continue

```

```
c MAKE NEXT GUESSES FOR B
  do 160 j=1,ir
    do 150 i=1,ir
      b(j)=b(j)-f1(i)*ainv(j,i)
150      continue
160    continue

c CHECK FOR CONVERGENCE
  do 170 i=1,ir
    if (abs(1-b(i)/bo(i)) .gt. c1) goto 180
    if (abs(f(i)) .gt. c2) goto 180
170  continue

c ON CONVERGENCE, GO TO FINAL OUTPUT STAGE
  goto 200

c IF NOT CONVERGED, OUTPUT NEXT GUESS
180  do 190 i=1,ir
      print *, 'next guess for b(',i,')=',b(i)
190  continue
      goto 35

c FINAL OUTPUT
200  do 260 jj=1, ir+1
      do 210 i=1,n
        call predictval(gamma(i),y(i),b,h)
210      continue
c EVALUATE STANDARD DEVIATIONS FOR B VALUES
c   EVALUATE COMPONENTS OF X MATRIX (DY)
      if (jj .gt. 1) then
        do 220 i=1,n
          dy(i,jj-1)=(y(i)-y1(i))/(b(jj-1)-bl(jj-1))
220        continue
      else
        do 230 i=1,n
          y1(i)=y(i)
230        continue
        do 240 i=1,ir
          bl(i)=b(i)
240        continue
      endif
      do 250 i=1,ir
        b(i)=bl(i)
250      continue
        b(jj)=d1*b(jj)
260  continue
c   EVALUATE COVARIANCE MATRIX (C)
  do 280 i=1,ir
    do 280 j=1,ir
      s=0.
      do 270 k=1,n
        s=s+dy(k,i)*dy(k,j)
```

```

270             continue
                xtx(i,j)=s
280             continue
290             continue
                call invert(ir,xtx,c)
                do 300 i=1,ir
                    print *, 'standard deviation of b(',i,')=',sqrt(c(i,i)*ss/(n-ir))
300             continue

c CALCULATION OF CORRELATION COEFFICIENTS
                do 307 i=1,ir
                    do 305 j=1,ir
                        corcoef=c(i,j)/sqrt(c(i,i)*c(j,j))
                        print *, 'coeff. ',i,'-',j,'=',corcoef
305             continue
307             continue

c CALCULATION OF 95% CONFIDENCE LIMITS EVALUATED AT GAMMA=.02,.2,2,20,200
c             GET STUDENT T VALUE AND ESTABLISH VALUES AT WHICH TO CALCULATE
                print *, 'input student t value'
                read*, tval
                do 360 ij=-2,2
                    gam=2.*10.**ij
c             CALCULATE DERIVATIVES
                    call predictval(gam,zval,b,h)
                    do 310 i=1, ir
                        bo(i)=b(i)
310             continue
                    do 330 ii=1,ir
                        b(ii)=d2*b(ii)
                        call predictval(gam,zvalnew,b,h)
                        delz(ii)=(zval-zvalnew)/(bo(ii)-b(ii))
                        do 320 jj=1,ir
                            b(jj)=bo(jj)
320             continue
330             continue

c             CALCULATE VARIANCE IN EQUATION AT THIS VALUE
                    v=0.
                    do 350 i=1,ir
                        do 340 j=1,ir
                            v=v+delz(i)*delz(j)*ss/(n-ir)*c(i,j)
340             continue
350             continue
c             OUTPUT CONFIDENCE LIMITS
                print *, '95% conf. int. at shear rate',gam,'=',tval*sqrt(v)
360             continue
                end

```

c SUBROUTINE TO CALCULATE PREDICTED VALUES BASED ON THEORY

```

subroutine predictval(xx,yy,b,h)
double precision b(5), xx, yy
yy=1./(1.-(b(1)+(b(2)-b(1))/(1+sqrt(xx/b(3))))*h/2.)*2.
return
end

```

c SUBROUTINE TO INVERT MATRIX

```

subroutine invert(ir,rorigmatx,rinvmatx)
double precision rorigmatx(5,5), rinvmatx(5,5)
do 3 i=1,ir
    do 2 j=1,ir
        rinvmatx(i,j)=0.
        if (i .eq. j) rinvmatx(i,j)=1.
2      continue
3    continue
do 60 j=1,ir
    i=j-1
10   i=i+1
    if (i .gt. ir) then
        print *, 'singular matrix'
        stop 100
    endif
    if (rorigmatx(i,j) .eq. 0) goto 10
    do 20 k=1,ir
        q=rorigmatx(j,k)
        qq=rorigmatx(i,k)
        rorigmatx(j,k)=qq
        rorigmatx(i,k)=q
        q=rinvmatx(j,k)
        qq=rinvmatx(i,k)
        rinvmatx(j,k)=qq
        rinvmatx(i,k)=q
20   continue
    t=1/rorigmatx(j,j)
    do 30 k=1,ir
        rorigmatx(j,k)=t*rorigmatx(j,k)
        rinvmatx(j,k)=t*rinvmatx(j,k)
30   continue
    do 50 l=1,ir
        if (l .eq. j) goto 50
        t=-rorigmatx(l,j)
        do 40 k=1, ir
            rorigmatx(l,k)=rorigmatx(l,k)+t*rorigmatx(j,k)
            rinvmatx(l,k)=rinvmatx(l,k)+t*rinvmatx(j,k)
40       continue
50       continue
60   continue
return
end

```

## Appendix E

### Notation

The dimensions given are in terms of mass ( $M$ ), length ( $L$ ), time ( $t$ ), and temperature ( $T$ ). When no dimensions are listed, the units are dimensionless, or they may vary.

$a$	radius of particle or rouleau, $L$ .
$A_c$	contact area between cells in rouleau, $L^2$ .
$A_{ij}$	component of jacobian matrix $= \frac{\partial F_i}{\partial b_j}$ .
$a_i$	constants for Chien polynomials.
$b_i$	parameter of function $Y$ to be fit to data.
$c$	cell volume fraction $= H/100$ .
$c_b$	volume fraction of particles in bulk region.
$c_p$	packing concentration $= 2/k$ .
$c_{\text{protein}}$	protein concentration, $ML^{-3}$ .
$c_w$	volume fraction of particles in wall layer.
$D$	dissipated energy of rouleau formation, $ML^2t^{-2}$ .
$\mathcal{D}$	translational diffusivity, $L^2t^{-1}$ .
$E$	elastic stored energy in rouleau, $ML^2t^{-2}$ .
	sum of squared residuals.
$E_i$	elastic modulus, $ML^{-1}t^{-2}$ .
$F$	auxiliary equation from method of Lagrange multipliers.
$F_i$	component of the matrix which represents the gradient of $E = \frac{\partial E}{\partial b_i}$ .
$f_i$	independent variables of $F$ .
$H$	hematocrit $= 100 \cdot c$ .
$h$	length of capillary, $L$ .
$K$	kinetic energy of rouleau, $ML^2t^{-2}$ .
	fitted parameter for centrifugation equation (3.2), $M^{-1}Lt^2$ .
$k$	intrinsic viscosity.

$k_0$	intrinsic viscosity at zero shear rate.
$k_{0i}$	parameter to fit $k_0$ as a function of $c$ .
$k_\infty$	intrinsic viscosity at high shear rate.
$k_{\infty i}$	parameter to fit $k_\infty$ as a function of $c$ .
$L$	length of capillary or channel, $L$ .
$M$	torque exerted on inner cylinder of viscometer, $ML$ .
$N$	number of data points to fit to equation $Y$ .
$P$	number of parameters for equation $Y$ .
$p$	pressure, $ML^{-1}t^{-2}$ .
$Q$	volumetric flow rate, $L^3t^{-1}$ .
$Q_f$	volumetric flow rate of fluid, $L^3t^{-1}$ .
$Q_p$	volumetric flow rate of particles, $L^3t^{-1}$ .
$R$	radius of capillary, $L$ .
$R_i$	radius of inner cylinder of viscometer, $L$ .
$R_o$	radius of outer cylinder of viscometer, $L$ .
$S$	total entropy of fluid volume, $ML^2t^{-2}T^{-1}$ .
$s_x$	integrated surface area normal to $x$ -direction, $L^2$ .
$T$	absolute temperature, $T$ .
$t$	time, $t$ .
$T_0$	reference temperature, $T$ .
$U$	total internal energy of fluid volume, $ML^2t^{-2}$ .
$V$	total volume of fluid element, $L^3$ .
$v$	velocity, $Lt^{-1}$ .
$W$	work done by external surface and body forces on rouleau, $ML^2t^{-2}$ .
$x$	rectangular coordinate, $L$ .
$x_i$	rectangular coordinate, $L$ .
$X_{ij}$	component of matrix with elements $\frac{\partial Y_i}{\partial b_j}$ .
$Y$	function to be fit to data.
$y$	rectangular coordinate, $L$ .
$y_i$	calculated values from function $Y$ corresponding to experimental



	value $y_{i,\text{exp}}$ .
$y_{i,\text{exp}}$	experimental value corresponding to calculated value $y_i$ .
<i>Greek symbols</i>	
$\alpha$	shape factor for Einstein equation (2.9).
$\gamma$	surface energy of cell membrane, $\text{Mt}^{-2}$ .
$\dot{\gamma}$	shear rate, $\text{t}^{-1}$ .
$\dot{\gamma}_c$	critical shear rate for Quemada equation (2.11), $\text{t}^{-1}$ .
$\dot{\gamma}_{ci}$	parameter to fit $\dot{\gamma}_c$ as a function of $c$ .
$\dot{\gamma}_r$	relative shear rate $= \dot{\gamma} / \dot{\gamma}_c$ .
$\delta_{ij}$	Kronecker delta function.
$\epsilon$	local porosity $= 1 - c$ .
$\bar{\epsilon}$	average porosity in a cell column.
$\epsilon_0$	local porosity at zero compressive pressure.
$\eta$	non-Newtonian viscosity, $\text{ML}^{-1}\text{t}^{-1}$ .
$\eta_b$	viscosity of suspension in bulk region, $\text{ML}^{-1}\text{t}^{-1}$ .
$\eta_f$	viscosity of suspending fluid phase, $\text{ML}^{-1}\text{t}^{-1}$ .
$\eta_r$	relative viscosity $= \eta / \eta_f$ .
$\eta_w$	viscosity of suspension in wall layer, $\text{ML}^{-1}\text{t}^{-1}$ .
$\theta$	ratio $\tau_A / \tau_D$ .
$\kappa$	Boltzmann's constant, $\text{ML}^2\text{t}^{-2}\text{T}^{-1}$ .
$\Lambda$	Lagrange multiplier.
$\lambda_i$	Lagrange multiplier.
$\mu$	Newtonian viscosity, $\text{ML}^{-1}\text{t}^{-1}$ .
$\mu_N$	Casson Newtonian viscosity, $\text{ML}^{-1}\text{t}^{-1}$ .
$\mu_p$	plasma viscosity, $\text{ML}^{-1}\text{t}^{-1}$ .
$\mu_w$	water viscosity, $\text{ML}^{-1}\text{t}^{-1}$ .
$\nu$	kinematic viscosity $= \eta / \rho$ , $\text{L}^2\text{t}^{-1}$ .
$\pi$	3.14159 . . . .
	compressive pressure, $\text{ML}^{-1}\text{t}^{-2}$ .
$\rho_c$	density of erythrocyte, $\text{ML}^{-3}$ .
$\rho_f$	density of suspending fluid, $\text{ML}^{-3}$ .

$\sigma_y$	standard deviation of function Y.
$T$	Taylor coefficient for equation (2.5).
$\tau$	shear stress, $ML^{-1}t^{-2}$ .
$\underline{\tau}$	stress tensor.
$\tau_A$	time constant for rouleau aggregation, t.
$\tau_D$	time constant for rouleau disaggregation, t.
$\tau_y$	yield shear stress, $ML^{-1}t^{-2}$ .
$\phi$	relative viscosity= $1/\eta_r$ .
$\psi$	control theory constraint, $\psi^2=(c_b-c)(c-c_w)$ .
$\omega$	angular velocity, radians/t.

## **Appendix F**

### **Data From Viscosity Measurements**

Listed in this appendix are the tables of data used in Chapter 4. Data is listed by hematocrit, and the beginning of different samples within a hematocrit is noted in the left-hand column. The beginning of data for a new hematocrit is delineated by a new set of column headings. A description of the units of blood and the various batches and samples obtained from the units is given in Table F-I. The Eagle's solution data begins on page 225, the plasma data begins on page 237.

Unit expiration date	anti- coagulant	Batch	suspending fluid	suspending viscosity (cP)	days since drawing	days before expiration
27-AUG-84	CPD	16	Eagle's	0.69	11	10
11-SEP-84	CPD	18	Eagle's	0.69	25	2
27-SEP-84	CPD-A1	19 20	Eagle's Eagle's	0.69 0.69	22 25	13 10
21-DEC-84	CPD-A1	24	Eagle's	0.69	35	0
06-JUN-85	AS-1	29 30 31	plasma plasma plasma	1.05 1.13 1.21	9 10 11	33 32 31
06-AUG-85	AS-1	32 33	plasma Eagle's	1.06 0.69	10 12	32 30

**Table F-1:** Catalog of units and batches used in this work.

**Table F-II:** Data tables for Eagle's solution suspensions.

0 sample and hematocrit	1 shear rate (1/sec)	2 viscosity (cP)	3 relative viscosity	32	20 49	72 918	135 099
1 20-3	0 03	4066 071	5892 837	33	23 39	71 800	104 058
2 H=98 4	0 10	1609 939	2333 243	34	23 78	70 623	102 642
3	0 16	1126 194	1646 658	35	23 95	70 045	101 514
4	0 22	886 369	1287 491	36	24 44	69 202	100 293
5	0 29	731 250	1059 783	37	29 96	69 168	100 243
6	0 35	654 872	920 104	38	35 72	69 771	101 117
7	0 42	564 648	816 359	39	55 80	66 129	95 839
8	0 48	515 058	746 432	40	64 22	60 021	86 987
9	0 61	437 054	633 368	41	87 54	63 849	92 529
10	0 80	369 657	555 735	42	106 68	62 711	90 886
11	0 87	346 768	502 552	43	123 85	59 666	86 472
12	0 99	322 404	457 252	44	138 42	52 007	89 865
13	1 13	294 673	427 065	45	163 95	61 537	99 184
14	1 18	293 125	424 819	46	185 34	59 514	86 252
15	1 40	263 096	381 299	47	195 78	61 140	88 609
16	1 43	262 084	379 832	48	205 20	59 415	81 109
17	1 66	240 520	348 590	49	221 10	60 868	83 214
18	1 69	242 420	351 333	50	238 20	59 395	86 000
19	1 93	224 486	325 342	51	252 96	60 614	87 846
20	1 94	224 007	327 546	52	257 97	59 367	86 039
21	2 19	210 956	305 703	53	277 80	59 406	86 076
22	2 20	212 410	307 841	54	299 79	59 812	86 684
23	2 39	202 811	293 929	55			
24	2 45	202 481	293 451				
25	2 59	196 779	285 187				
26	2 79	190 366	275 893				
27	3 00	186 040	269 623				
28	3 43	182 386	264 328				
29	5 98	148 946	215 864				
30	8 01	133 720	193 797				
31	8 53	134 528	195 055				
32	10 66	125 098	181 301				
33	13 32	119 265	173 717				
34	13 64	121 914	176 687				
35	15 98	112 576	168 951				
36	16 17	119 452	171 670				
37	18 73	112 117	168 286				
38	19 29	112 637	169 591				
39	21 29	114 402	165 900				
40	22 60	111 953	162 164				
41	25 11	112 598	163 128				
42	25 27	110 668	160 388				
43	28 59	109 798	159 128				
44	36 39	115 829	167 868				
45	61 74	110 593	160 230				
46	66 03	101 227	146 756				
47	93 57	108 240	156 870				
48	98 88	101 760	147 478				
49	119 10	107 134	155 276				
50	138 15	103 794	154 716				
51	163 65	106 324	154 093				
52	165 09	102 471	148 509				
53	198 18	102 664	143 788				
54	201 75	105 678	153 591				
55	231 09	102 554	149 180				
56	239 88	103 653	151 120				
57	250 92	103 109	150 433				
58	265 44	105 515	155 920				
59	277 38	103 374	149 817				
60	299 76	104 353	151 236				
0 sample and hematocrit	1 shear rate (1/sec)	2 viscosity (cP)	3 relative viscosity	32	20 49	72 918	135 099
1 20-3	0 04	1920 301	2783 045	32	4 06	84 072	109 943
2 H=97 2	0 10	908 883	1317 222	33	5 32	75 540	109 910
3	0 17	651 079	943 592	34	6 61	70 318	101 910
4	0 23	505 837	732 097	35	7 98	65 103	94 352
5	0 29	434 286	629 400	36	9 14	63 404	91 890
6	0 42	339 943	492 671	37	11 06	60 369	87 401
7	0 61	268 692	389 409	38	11 28	59 224	85 802
8	0 87	218 062	316 032	39	12 97	57 885	83 891
9	0 93	211 825	306 993	40	14 58	55 469	80 680
10	1 06	194 783	282 294	41	14 89	56 119	81 332
11	1 38	169 747	246 010	42	17 24	53 676	77 791
12	1 40	168 679	244 462	43	17 44	54 331	78 741
13	1 63	156 552	224 887	44	19 22	52 449	76 013
14	1 73	152 379	220 859	45	19 34	53 101	76 958
15	1 89	146 224	211 919	46	21 27	52 067	75 459
16	2 06	140 298	203 300	47	22 54	51 007	73 923
17	2 14	137 544	199 339	48	23 17	51 275	74 312
18	2 40	131 414	190 455	49	25 09	50 588	73 316
19	2 45	129 911	188 277	50	25 20	49 931	72 364
20	2 92	120 592	174 771	51	29 87	48 604	70 441
21	3 65	118 914	172 339	52	35 67	50 464	73 426
22	6 20	97 244	146 533	53	48 45	47 889	69 404
23	9 60	87 749	127 172	54	48 45	47 889	69 404
24	8 75	87 697	127 097	55	66 18	43 038	62 359
25	11 29	82 094	118 977	56	67 59	45 699	64 230
26	11 91	80 730	117 000	57	99 12	42 058	60 954
27	13 64	78 578	113 891	58	99 39	43 882	63 597
28	15 21	76 612	111 032	59	124 89	43 137	62 517
29	17 04	75 436	109 328	60	138 81	41 470	60 101
30	17 86	74 227	107 575	61	156 78	42 537	61 648
31	20 21	73 352	106 307	62	165 24	41 253	59 787
0 sample and hematocrit	1 shear rate (1/sec)	2 viscosity (cP)	3 relative viscosity	32	20 49	72 918	135 099
1 19-6	0 03	1007 609	2618 723	32	168 64	42 118	61 041
2 H=96	0 09	673 443	976 004	33	198 33	41 050	59 493
3	0 28	307 006	444 936	34	226 77	41 815	60 601
4	0 47	219 480	318 087	35	231 30	40 934	59 325
5	0 66	178 826	259 168	36	291 33	40 778	59 099
6	0 87	157 898	224 838	37	296 76	41 103	59 570
7	1 05	140 092	203 032				
8	1 14	134 747	195 286				
9	1 37	122 744	177 890				
10	1 40	121 600	176 232				
11	1 62	112 673	163 294				
12	1 73	109 396	158 345				

[illegible]

38	19 88	47 781	14 212	10	0 87	147 649	212 494
39	21 41	47 162	14 211	11	0 87	141 476	214 490
40	22 51	46 349	14 210	12	1 18	120 492	174 626
41	23 60	45 280	14 209	13	1 27	116 947	172 357
42	24 51	45 422	14 208	14	1 27	103 980	150 645
43	25 55	45 565	14 207	15	1 66	103 231	149 610
44	27 51	44 635	14 206	16	1 93	94 690	137 232
45	29 79	43 553	14 205	17	2 06	93 052	134 859
46	29 99	43 652	14 204	18	2 32	67 937	127 445
47	42 81	42 156	14 203	19	2 46	85 320	132 652
48	37 98	37 627	14 202	20	2 52	84 688	122 707
49	37 37	38 570	14 201	21	2 83	80 015	117 654
50	35 58	35 658	14 200	22	3 00	79 065	114 597
51	33 69	33 319	14 199	23	4 91	66 209	104 614
52	105 31	34 519	14 198	24	7 36	56 913	92 493
53	119 19	35 658	14 197	25	8 09	55 640	90 638
54	125 53	33 944	14 196	26	10 63	50 831	73 668
55	144 69	34 916	14 195	27	11 90	49 699	72 317
56	152 07	33 537	14 194	28	13 95	47 516	68 654
57	176 49	34 277	14 193	29	15 72	46 431	67 271
58	178 53	33 146	14 192	30	17 26	45 150	65 415
59	211 47	32 934	14 191	31	19 54	44 372	64 307
60	214 68	32 741	14 190	32	19 90	43 669	63 286
61	237 87	32 753	14 189	33	22 09	43 222	62 641
62	259 14	32 341	14 188	34	23 19	42 200	61 159
63	259 69	32 647	14 187	35	24 43	42 187	61 141
64	259 73	32 800	14 186	36	25 84	41 403	61 004
65	259 79	32 803	14 185	37	27 16	41 552	60 220
66	0 04	1773 984	14 184	38	28 49	40 584	59 817
67	0 10	771 171	14 183	39	29 98	40 489	58 680
68	0 29	361 933	14 182	40	30 60	44 216	64 081
69	0 41	327 658	14 181	41	62 31	38 372	55 613
70	0 55	229 284	14 180	42	64 23	37 130	53 816
71	0 57	226 639	14 179	43	81 48	36 752	53 307
72	0 74	198 450	14 178	44	100 59	35 759	51 868
73	0 99	164 545	14 177	45	105 59	35 409	51 217
74	1 07	164 356	14 176	46	138 81	34 709	50 320
75	1 31	141 541	14 175	47	139 82	34 463	49 946
76	1 33	143 599	14 174	48	170 73	34 144	49 455
77	1 56	126 913	14 173	49	172 68	34 100	49 420
78	1 60	130 002	14 172	50	205 05	33 632	49 052
79	1 69	127 936	14 171	51	205 05	33 753	48 917
80	1 82	117 169	14 170	52	231 36	33 556	48 557
81	1 92	116 687	14 169	53	234 30	33 470	48 507
82	2 07	109 593	14 168	54	277 53	32 226	48 158
83	2 26	105 254	14 167	55	299 79	33 594	48 687
84	2 33	103 351	14 166	56	0 04	1106 818	1604 054
85	2 52	100 464	14 165	57	0 10	525 809	754 776
86	2 52	99 799	14 164	58	0 17	374 607	542 039
87	2 71	95 888	14 163	59	0 23	303 121	439 306
88	2 79	95 598	14 162	60	0 42	266 752	299 655
89	3 00	92 653	14 161	61	0 61	164 328	212 127
90	3 60	85 917	14 160	62	0 67	135 724	154 701
91	5 33	73 746	14 159	63	0 94	133 429	153 315
92	5 51	70 811	14 158	64	1 31	108 943	125 955
93	7 42	62 859	14 157	65	1 40	107 574	121 604
94	7 76	62 208	14 156	66	1 57	100 575	115 761
95	9 97	56 758	14 155	67	1 80	95 447	109 329
96	13 60	56 027	14 154	68	1 63	92 978	114 751
97	12 52	52 708	14 153	69	2 06	69 419	129 593
98	13 91	51 240	14 152	70	2 20	86 373	123 106
99	15 08	50 030	14 151	71	2 39	83 750	121 318
100	15 30	51 462	14 150	72	2 52	81 564	119 209
101	17 63	47 484	14 149	73	2 59	81 039	117 448
102	17 28	47 676	14 148	74	2 71	79 221	114 813
103	19 55	46 777	14 147	75	2 86	77 347	112 397
104	19 66	46 412	14 146	76	3 00	76 329	110 627
105	22 11	45 436	14 145	77	3 03	78 168	112 587
106	22 49	44 998	14 144	78	3 27	61 712	69 418
107	24 65	44 342	14 143	79	3 48	55 598	60 577
108	25 14	43 866	14 142	80	3 55	53 678	77 794
109	27 19	43 436	14 141	81	3 99	52 204	75 658
110	28 45	42 649	14 140	82	11 26	48 521	70 320
111	29 97	42 399	14 139	83	11 42	48 616	70 458
112	48 69	40 234	14 138	84	14 56	45 498	65 939
113	60 30	35 970	14 137	85	15 11	45 880	64 493
114	74 10	37 247	14 136	86	17 21	43 822	63 510
115	79 98	34 771	14 135	87	17 66	44 336	64 252
116	105 93	35 316	14 134	88	20 20	43 065	62 411
117	106 56	33 742	14 133	89	20 51	42 257	61 242
118	137 88	34 258	14 132	90	22 76	42 103	61 013
119	139 50	33 060	14 131	91	23 14	41 353	60 762
120	169 77	33 539	14 130	92	23 31	41 226	60 762
121	172 59	32 513	14 129	93	23 78	40 429	58 593
122	199 05	32 231	14 128	94	29 10	39 578	57 359
123	201 51	33 665	14 127	95	29 94	43 056	57 443
124	225 39	32 044	14 126	96	29 98	39 578	57 504
125	233 28	32 697	14 125	97	61 74	37 026	52 661
126	231 85	31 924	14 124	98	79 62	33 421	49 436
127	298 81	32 474	14 123	99	99 63	34 584	50 122
128	478 28	31 819	14 122	100	112 47	32 232	46 699
129	499 73	31 944	14 121	101	125 40	33 684	48 817
			14 120	102	145 59	31 733	45 990
			14 119	103	157 52	33 047	47 894
			14 118	104	179 68	31 420	45 551
			14 117	105	182 76	32 649	47 293
			14 116	106	204 26	32 247	46 943
			14 115	107	211 74	31 289	45 317
			14 114	108	231 00	31 058	46 514
			14 113	109	244 68	31 192	45 296
			14 112	110	265 41	31 947	46 274
			14 111	111	277 44	31 185	45 175
			14 110	112	299 79	31 432	45 554

1 sample and hematocrit	1 shear rate (1/sec)	2 viscosity (cP)	3 relative viscosity
1 19-4	0 04	1284 252	1861 235
2 H=94	0 10	584 751	847 465
3	0 17	403 085	534 181
4	0 23	312 462	432 757
5	0 29	271 553	393 990
6	0 36	239 163	345 164
7	0 42	217 190	314 768
8	0 42	217 190	314 768
9	0 61	172 901	240 561



0 sample and hematocrit	1 shear rate (1/sec)	2 viscosity (cP)	3 relative viscosity				
1 18-3	0.04	2032.837	2975.155	100	42.81	37.431	54.572
2 H=93	0.17	671.174	972.716	101	60.09	33.125	48.007
3	0.36	406.339	568.897	102	68.25	34.044	49.339
4	0.61	289.932	420.191	103	79.86	31.649	45.868
5	0.73	281.421	407.357	104	87.39	32.612	47.264
6	0.87	231.572	333.612	105	99.69	30.639	44.399
7	1.00	225.930	327.435	106	119.16	30.979	44.897
8	1.06	204.925	296.993	107	132.81	29.614	42.919
9	1.26	192.593	279.120	108	151.05	30.079	43.591
10	1.32	180.570	261.646	109	152.55	29.213	42.308
11	1.57	164.564	238.499	110	182.85	29.377	42.575
12	1.66	162.638	225.707	111	185.61	28.641	41.509
13	1.82	152.491	221.001	112	208.47	28.983	42.004
14	2.06	145.473	210.800	113	212.07	28.377	41.126
15	2.14	139.519	202.201	114	240.00	28.613	41.668
16	2.40	132.357	192.112	115	245.10	28.060	40.667
17	2.45	131.956	191.241	116	265.59	28.386	41.139
18	2.71	124.320	180.174	117	278.16	27.836	40.342
19	2.72	123.271	178.654	118	299.79	27.920	40.464
20	2.72	123.146	181.371	119 18-7	0.05	1243.750	1802.536
21	2.98	119.119	172.636	120 H=93	0.11	631.016	914.516
22	3.00	119.468	173.113	121	0.24	366.163	330.674
23	4.63	96.766	140.241	122	0.43	243.017	252.199
24	4.68	98.673	143.004	123	0.60	213.673	209.674
25	7.32	82.297	119.271	124	0.62	192.668	279.229
26	9.74	73.988	107.229	125	0.75	171.303	248.265
27	9.95	73.778	106.925	126	0.79	176.511	255.813
28	12.60	68.866	99.806	127	0.94	140.579	215.322
29	12.91	68.162	98.786	128	1.06	144.759	209.796
30	15.46	64.882	94.032	129	1.26	124.434	180.339
31	15.52	65.401	94.784	130	1.32	125.714	182.194
32	16.44	71.390	103.406	131	1.51	113.175	164.032
33	18.02	62.649	90.796	132	1.65	110.627	160.369
34	19.25	61.643	89.236	133	1.75	105.394	152.745
35	20.59	60.945	88.326	134	1.77	103.872	150.539
36	21.91	59.725	86.556	135	1.92	101.595	147.259
37	23.78	59.309	85.926	136	2.02	96.749	140.216
38	24.58	59.510	84.787	137	2.18	93.911	136.103
39	26.35	58.071	84.161	138	2.34	89.254	129.354
40	27.23	57.207	83.609	139	2.45	88.657	128.468
41	29.98	56.637	82.083	140	2.53	86.141	124.842
42	42.03	55.889	80.999	141	2.58	86.356	125.154
43	60.21	49.452	71.670	142	2.71	84.228	122.070
44	67.47	51.645	74.844	143	2.72	83.414	120.690
45	79.80	48.195	69.846	144	2.98	79.996	115.676
46	92.73	49.563	71.850	145	3.66	71.575	103.732
47	118.32	48.162	69.800	146	6.21	56.042	81.220
48	119.28	46.730	67.725	147	7.29	54.035	78.212
				148	9.25	48.979	70.984
49	143.91	47.409	68.703	149	9.39	47.684	69.107
50	145.83	46.184	66.954	150	11.91	44.647	64.706
51	175.74	46.654	67.616	151	11.94	44.006	63.777
52	178.95	45.624	66.172	152	14.50	41.465	62.094
53	201.24	46.266	67.052	153	15.21	41.185	59.626
54	218.61	45.238	65.562	154	17.05	39.565	57.241
55	239.37	45.858	66.461	155	17.34	42.388	61.432
56	251.61	45.006	65.226	156	18.51	38.793	56.234
57	264.84	45.635	66.138	157	20.23	37.763	54.729
58	278.07	44.877	65.039	158	21.83	37.033	53.671
59	299.76	45.086	65.342	159	22.14	36.933	53.526
60 18-6	0.04	1532.837	2221.532	160	24.69	35.962	52.119
61 H=93	0.17	514.795	746.080	161	25.13	35.624	51.629
62	0.36	307.736	446.023	162	27.23	35.237	51.060
63	0.61	292.053	423.265	163	27.74	34.812	50.452
64	0.61	217.294	314.919	164	29.98	34.391	49.842
65	0.67	223.792	324.336	165	42.78	32.819	47.554
66	0.93	180.032	260.916	166	59.61	28.636	41.501
67	0.93	169.102	245.075	167	74.67	29.450	42.661
68	1.19	146.613	212.483	168	100.02	28.044	40.643
69	1.20	152.643	221.222	169	105.93	26.423	38.294
70	1.51	129.041	187.016	170	125.52	27.187	39.401
71	1.53	131.739	190.924	171	138.90	25.724	37.281
72	1.73	122.506	177.545	172	157.35	26.462	38.200
73	1.82	116.428	168.736	173	176.37	26.144	37.090
74	2.08	108.225	156.848	174	178.56	25.193	36.512
75	2.13	109.401	158.552	175	208.29	25.731	37.291
76	2.39	100.914	146.252	176	211.56	24.750	36.159
77	2.46	101.466	147.052	177	238.08	24.773	35.903
78	2.65	96.151	139.349	178	240.03	25.391	35.799
79	2.79	94.399	136.810	179	265.59	25.189	36.506
80	3.00	90.909	131.752	180	277.65	24.657	35.735
81	3.00	91.393	132.454	181	299.76	24.740	35.855
82	5.54	68.211	98.857				
83	6.65	65.005	94.210	0 sample and	1 shear rate	2 viscosity	3 relative
84	8.08	58.076	84.168	hematocrit	(1/sec)	(cP)	viscosity
85	9.29	56.294	81.586	1 18-5	0.18	257.264	372.846
86	10.64	52.525	76.123	2 H=92	0.50	154.573	224.019
87	11.94	51.219	74.375	3	0.60	136.866	196.357
88	14.45	47.509	68.854	4	0.81	119.049	172.525
89	15.25	47.128	68.301	5	0.86	115.457	167.329
90	17.40	49.483	71.714	6	1.13	102.704	149.846
91	17.65	44.725	64.862	7	1.19	98.208	142.330
92	18.55	44.307	64.213	8	1.39	93.143	134.990
93	20.83	42.648	61.809	9	1.41	100.851	146.161
94	21.86	42.128	61.055	10	1.52	89.143	129.193
95	23.84	41.140	59.623	11	1.71	85.707	124.213
96	24.01	41.241	59.770	12	1.85	82.936	120.167
97	27.80	39.563	57.308	13	1.96	81.160	117.623
98	29.97	38.955	56.457	14	2.18	77.212	111.901
99	29.97	39.026	56.559				

15	2	28	75	329	113	912	109	231	63	32	463	47	648		
16	2	44	74	420	117	855	110	258	66	32	368	48	939		
17	2	60	73	534	118	571	111	255	62	33	104	49	877		
18	2	77	70	579	119	289	112	291	18	32	315	45	803		
19	3	60	69	543	120	787	113	299	79	32	578	47	214		
20	3	97	65	129	121	290	114	0	04	450	600	1375	612		
21	4	30	53	930	122	159	115	0	11	439	385	256	7901		
22	7	78	52	641	123	291	116	0	17	324	561	470	378		
23	9	59	49	608	124	446	117	0	30	223	317	313	646		
24	10	94	46	177	125	822	118	0	42	205	143	247	304		
25	11	31	33	846	126	938	119	0	60	154	287	213	604		
26	12	88	45	398	127	794	120	0	62	144	252	209	104		
27	14	14	45	395	128	790	121	1	00	111	159	161	752		
28	15	53	43	540	129	101	122	1	06	110	175	159	674		
29	16	69	43	735	130	284	123	1	32	46	359	134	216		
30	18	84	41	965	131	814	124	1	53	91	498	132	606		
31	19	88	42	176	132	125	125	1	68	87	701	127	103		
32	22	14	40	746	133	125	126	1	70	85	555	123	991		
33	22	43	41	267	134	807	127	1	92	81	298	117	823		
34	24	79	39	571	135	784	128	2	02	78	771	114	161		
35	24	62	40	212	136	278	129	2	23	75	962	110	090		
36	26	75	39	420	137	130	130	2	33	74	361	107	770		
37	26	13	39	836	138	443	131	2	52	71	897	104	199		
38	31	41	38	706	139	696	132	2	59	70	915	102	775		
39	32	97	38	379	140	346	133	2	98	66	744	96	730		
40	36	36	40	667	141	372	134	4	67	64	418	78	867		
41	35	89	36	608	142	555	135	6	64	48	300	73	000		
42	39	94	34	134	143	470	136	7	42	46	748	67	751		
43	43	96	34	153	144	497	137	10	58	41	623	60	323		
44	48	42	32	770	145	493	138	11	24	41	249	59	781		
45	123	85	33	159	146	557	139	14	54	39	177	55	329		
46	139	05	32	039	147	433	140	15	69	37	630	54	536		
47	154	07	32	492	148	490	142	18	51	35	960	52	145		
48	172	17	31	713	149	961	143	19	52	35	764	51	832		
49	199	57	32	173	150	628	144	21	82	34	578	50	113		
50	205	11	31	490	151	638	145	23	34	34	407	49	865		
51	221	57	31	901	152	713	146	24	46	33	759	48	926		
52	231	27	31	268	153	345	147	25	69	33	632	48	742		
53	246	84	31	745	154	607	148	29	09	32	537	47	155		
54	257	67	31	191	155	204	149	41	55	31	413	46	271		
55	270	76	31	027	156	610	150	60	27	27	659	40	578		
56	0	04	1695	350	2457	812	151	79	74	26	349	41	084		
57	0	11	825	676	1126	664	152	99	03	26	824	38	305		
58	0	23	466	337	676	575	153	117	00	26	824	38	877		
59	0	43	309	232	448	162	154	15	00	25	422	35	843		
60	0	68	229	018	330	461	155	15	00	25	075	37	790		
61	0	73	236	659	342	955	156	16	00	25	630	37	145		
62	0	95	189	790	275	043	157	199	11	24	674	36	194		
63	1	30	190	181	275	625	158	219	99	25	365	36	761		
64	1	06	171	725	240	891	159	232	05	24	809	35	955		
65	1	26	162	892	236	075	160	264	45	25	105	36	384		
66	1	38	147	460	210	710	161	298	32	24	568	35	406		
67	1	59	140	515	200	574	162	0	sample and	linear rate	viscosity	viscosity			
68	1	70	120	663	169	367	163	hematocrit	(1/sec)	(cP)	viscosity				
69	1	86	128	096	155	643	164								
70	2	02	119	459	175	124	165	1	19-10	0	04	805	691	1157	668
71	2	12	119	372	175	003	166	2	M=91	0	10	372	537	559	909
72	2	27	113	071	165	871	167	3		0	16	269	301	390	291
73	2	36	112	165	162	552	168	4		0	23	222	193	322	019
74	2	59	106	818	124	800	169	5		0	29	184	877	267	938
75	2	63	104	284	151	136	170	6		0	35	165	765	240	209
76	2	85	101	517	147	126	171	7		0	42	149	426	215	559
77	2	87	99	843	144	700	172	8		0	48	141	319	204	810
78	3	00	99	029	143	520	173	9		0	67	116	429	168	738
79	4	77	77	687	112	590	174	10		0	74	115	128	166	852
80	6	71	69	150	100	217	175	11		0	93	95	511	138	422
81	7	96	62	585	91	293	176	12		1	12	89	934	130	194
82	10	51	57	165	82	848	177	13		1	13	90	963	131	251
83	10	67	57	677	83	590	178	14		1	37	81	816	118	574
84	10	92	63	324	91	774	179	15		1	47	50	315	116	399
85	13	68	52	610	76	246	180	16		1	63	75	129	108	683
86	13	97	52	652	76	307	181	17		1	80	72	805	105	514
87	16	87	49	618	71	910	182	18		1	88	70	644	102	383
88	17	28	49	392	71	583	183	19		2	01	69	799	99	709
89	20	07	47	466	68	791	184	20		2	04	68	433	99	178
90	20	59	47	174	68	368	185	21		2	26	66	087	95	778
91	23	22	45	743	66	294	186	22		2	33	64	846	93	980
92	24	53	45	336	65	704	187	23		2	46	63	497	92	025
93	25	87	44	527	64	532	188	24		2	58	62	005	89	862
94	27	06	44	414	64	369	189	25		2	59	59	698	86	475
95	28	52	43	358	63	128	190	26		3	00	58	699	85	071
96	29	98	43	270	62	710	191	27		3	62	54	262	78	641
97	36	42	43	410	62	913	192	28		6	16	44	659	64	780
98	61	92	38	978	56	490	193	29		7	94	40	460	58	694
99	86	57	35	940	52	087	194	30		8	70	40	055	59	051
100	93	60	35	603	50	048	195	31		10	58	37	269	54	042
101	106	11	34	097	49	416	196	32		11	26	37	160	54	000
102	119	13	33	470	51	806	197	33		13	25	35	176	50	980
103	144	69	34	750	50	362	198	34		13	77	33	295	51	149
104	159	60	33	075	47	935	199	35		15	86	33	618	48	751
105	176	55	34	112	49	935	200	36		16	99	33	510	46	681
106	195	57	33	624	49	029	201	37		18	51	32	457	47	054
107	198	66	32	664	47	329	202	38		19	52	32	455	47	094
108	227	37	33	441	48	465	203	39		21	15	31	516	45	678
								40		22	08	31	623	45	898
								41		23	13	30	528	44	838
								42		24	63	30	395	44	970
								43		25	11	30	143	44	120
								44		27	15	30	365	44	007
								45		27	76	29	631	43	233
								46		29	50	29	630	42	926
								47		36	27	30	356	43	994
								48		74	55	26	237	38	027

49	79 38	24 754	35 871	24	11 93	33 553	1 628
50	93 57	25 351	36 755	25	13 21	33 352	1 335
51	99 15	24 115	34 942	26	13 23	31 722	1 315
52	118 77	23 693	34 332	27	16 39	31 493	1 542
53	119 10	24 547	35 575	28	19 20	30 033	1 526
54	145 20	23 316	33 791	29	19 58	30 244	1 832
55	150 93	23 942	34 699	30	22 49	29 048	1 128
56	178 32	22 905	33 196	31	22 76	29 367	1 561
57	189 12	23 441	33 972	32	23 13	28 385	1 128
58	217 95	22 698	32 896	33	23 30	28 438	1 504
59	233 53	22 987	33 314	34	26 46	27 252	1 344
60	237 78	22 947	32 749	35	28 45	27 675	1 109
61	250 98	22 532	32 635	36	29 98	27 568	1 454
62	265 56	22 831	33 088	37	27 39	24 435	1 133
63	284 13	22 411	32 460	38	64 80	23 333	1 815
64	299 76	22 538	32 664	39	98 32	23 234	1 672
65	0 04	814 179	1179 970	40	103 66	21 894	1 730
66	0 10	430 792	624 336	41	125 46	22 143	1 691
67	0 17	299 457	433 996	42	140 16	21 297	1 865
68	0 29	206 721	299 596	43	156 48	21 607	1 314
69	0 42	165 338	239 620	44	171 63	20 958	1 374
70	0 61	133 759	193 854	45	187 38	21 230	1 768
71	0 80	120 187	174 184	46	202 98	20 692	1 988
72	0 60	116 075	168 225	47	218 40	20 962	1 380
73	1 06	100 823	146 120	48	234 33	20 561	1 799
74	1 13	99 470	144 159	49	249 36	20 777	1 112
75	1 31	90 782	131 568	50	271 95	20 463	1 657
76	1 46	87 129	126 274	51	274 11	20 652	1 930
77	1 63	81 561	118 204	52	299 76	20 476	1 675
78	1 80	78 693	114 048	53	0 04	1106 294	1 602 325
79	1 95	75 423	109 309	54	0 11	523 944	1 319
80	2 06	73 936	107 154	55	0 23	312 580	1 504
81	2 20	71 224	103 223	56	0 50	263 101	1 319
82	2 32	69 853	101 236	57	0 42	207 431	1 106 825
83	2 52	67 396	97 675	58	0 43	222 737	1 303 097
84	2 52	67 412	97 699	59	0 55	174 333	1 312 091
85	2 85	63 821	92 494	60	0 66	169 683	1 319 311
86	3 00	62 800	91 014	61	0 68	154 574	1 314 032
87	3 53	58 895	85 355	62	0 93	135 081	1 319 779
88	5 44	49 625	71 420	63	1 00	122 219	1 319 129
89	6 68	45 919	66 406	64	1 19	115 969	1 319 969
90	8 62	42 339	61 361	65	1 25	107 180	1 319 333
91	10 64	39 515	57 268	66	1 45	101 731	1 319 455
92	11 79	38 636	55 994	67	1 51	96 837	1 319 343
93	14 35	36 662	53 133	68	1 71	89 942	1 319 255
94	14 60	36 258	52 504	69	1 72	92 352	1 319 843
95	17 54	34 963	50 671	70	1 76	88 515	1 319 293
96	17 92	34 389	49 839	71	1 99	84 978	1 319 157
97	20 56	33 249	48 259	72	2 02	82 724	1 319 890
98	20 72	33 548	48 620	73	2 25	79 899	1 319 796
99	23 21	32 398	46 954	74	2 54	75 186	1 319 525
100	23 26	32 740	47 449	75	2 34	76 813	1 319 123
101	25 19	31 807	46 097	76	2 51	75 046	1 319 755
102	25 79	32 020	46 406	77	2 53	73 885	1 319 051
103	29 15	30 865	44 732	78	2 78	71 069	1 319 369
104	30 63	32 566	47 197	79	2 98	68 636	1 319 472
105	36 10	29 556	41 286	80	3 00	68 450	1 319 214
106	37 09	27 167	39 372	81	4 89	52 221	1 319 683
107	42 40	25 584	37 078	82	7 31	45 267	1 319 604
108	94 26	26 273	39 077	83	8 05	42 013	1 319 888
109	113 40	25 608	37 113	84	9 95	39 771	1 319 539
110	132 00	24 648	35 722	85	11 87	36 291	1 319 596
111	138 90	25 043	36 294	86	12 59	36 283	1 319 554
112	164 94	24 218	35 099	87	14 43	34 022	1 319 307
113	170 79	24 548	35 577	88	15 25	33 829	1 319 028
114	198 06	23 909	34 651	89	16 65	36 126	1 319 357
115	202 59	24 226	35 110	90	16 98	32 353	1 319 888
116	224 52	23 697	34 343	91	17 69	32 076	1 319 487
117	228 03	24 023	34 816	92	19 32	31 020	1 319 557
118	244 29	23 600	34 207	93	20 53	30 712	1 319 510
119	253 53	23 861	34 581	94	21 44	30 221	1 319 799
120	272 40	23 789	34 477	95	23 17	29 636	1 319 951
121	284 01	23 413	33 932	96	23 99	29 348	1 319 555
122	299 76	23 536	34 110	97	25 81	28 640	1 319 577
				98	26 55	28 397	1 319 445
				99	28 46	27 929	1 319 477
				100	29 98	27 719	1 319 172
				101	42 15	26 548	1 319 475
				102	66 48	22 676	1 319 864
				103	73 98	23 297	1 319 764
				104	92 79	21 403	1 319 019
				105	105 84	21 797	1 319 590
				106	125 91	20 407	1 319 575
				107	137 79	20 945	1 319 355
				108	159 00	19 830	1 319 739
				109	169 62	20 348	1 319 490
				110	192 06	19 437	1 319 171
				111	201 45	19 911	1 319 697
				112	218 58	19 108	1 319 853
				113	233 10	19 588	1 319 288
				114	239 35	19 394	1 319 672
				115	265 08	19 536	1 319 023
				116	271 26	18 912	1 319 409
				117	291 15	18 810	1 319 261
				118	299 70	18 959	1 319 477

0 sample and hematocrit	1 shear rate (1/sec)	2 viscosity (cP)	3 relative viscosity
1 16-3	0 04	508 547	737 025
2 H=90	0 10	325 994	472 455
3	0 29	175 934	254 977
4	0 60	96 020	137 159
5	0 61	106 365	154 442
6	0 99	81 871	118 654
7	1 00	74 198	107 533
8	1 50	72 067	104 474
9	1 53	61 822	89 599
10	1 74	69 606	100 875
11	1 82	63 782	92 438
12	1 93	58 937	85 416
13	2 20	62 028	89 896
14	2 33	55 874	80 977
15	2 52	57 048	82 678
16	2 86	52 672	76 326
17	3 00	52 980	76 786
18	4 30	46 427	68 010
19	6 64	39 372	57 061
20	6 84	40 029	58 028
21	7 77	42 471	61 552
22	8 61	36 388	53 026
23	10 01	35 606	51 893

0 sample and hematocrit	1 shear rate (1/sec)	2 viscosity (cP)	3 relative viscosity				
1 18-10	0.03	1315 741	1905 871	28	25 83	27 876	28 430
2 H=87	0.10	571 473	828 222	29	28 49	26 897	28 971
3	0.18	273 221	540 900	30	24 97	26 786	28 791
4	0.35	221 129	370 477	41	48 30	24 242	29 131
5	0.61	155 919	255 970	42	56 48	21 674	31 416
6	0.88	123 226	227 313	43	50 10	21 674	31 416
7	0.35	127 729	185 114	44	45 19	20 884	30 257
8	0.94	126 849	147 859	45	135 60	20 574	29 817
9	1.18	105 785	133 312	46	112 68	20 056	29 067
10	1.20	109 965	125 376	47	127 28	19 753	28 629
11	1.43	93 946	126 154	48	145 65	19 427	28 155
12	1.47	96 486	139 808	49	175 62	19 120	27 710
13	1.67	85 667	125 604	50	178 83	18 938	27 466
14	1.69	86 496	125 357	51	207 54	18 794	27 238
15	1.74	87 777	127 213	52	211 83	18 609	26 970
16	1.94	80 247	116 300	53	235 23	18 391	26 624
17	2.07	79 062	114 612	54	239 19	18 563	26 408
18	2.26	73 096	106 835	55	233 09	18 266	26 472
19	2.39	72 545	105 158	56	264 75	18 439	26 725
20	2.86	66 737	96 720	57	284 67	18 116	26 255
21	3.00	65 162	94 435	58	299 72	18 184	26 354
22	3.39	58 012	84 075	59	0.04	610 204	884 354
23	4.14	45 029	65 274	60	0.11	321 008	465 209
24	4.62	43 432	65 843	61	0.17	233 919	359 013
25	4.61	40 418	60 577	62	0.30	161 647	234 271
26	4.23	38 096	55 213	63	0.36	147 103	213 193
27	4.77	39 136	56 719	64	0.43	131 783	190 990
28	11 42	35 509	51 432	65	0.49	122 454	177 470
29	12 50	34 257	49 648	66	0.67	107 824	156 267
30	15 05	32 265	46 761	67	0.74	96 491	139 847
31	15 89	31 986	45 357	68	1.00	86 565	125 487
32	17 67	30 903	44 787	69	1.06	80 718	117 055
33	17 24	30 367	44 010	70	1.32	72 961	105 741
34	21 17	29 164	42 267	71	1.33	74 095	105 650
35	21 41	29 000	42 059	72	1.57	67 559	97 480
36	22 81	28 142	40 786	73	1.66	67 100	97 159
37	22 96	28 174	40 832	74	1.83	63 103	91 509
38	28 10	27 406	39 717	75	1.92	62 412	90 452
39	28 45	26 902	38 841	76	2.08	59 138	85 707
40	29 97	25 564	38 400	77	2.19	59 103	85 657
41	49 02	23 607	34 503	78	2.34	56 539	81 941
42	66 12	21 121	30 610	79	2.58	54 590	79 123
43	80 76	21 545	31 225	80	2.59	54 213	78 570
44	92 43	20 326	29 023	81	2.85	52 349	75 868
45	112 65	20 320	29 449	82	3.00	51 291	74 335
46	118 00	19 280	27 942	83	3.56	45 430	67 290
47	144 48	19 591	28 393	84	3.47	45 221	66 842
48				85	3.77	35 079	50 839
49	151 60	18 596	27 386	86	4.25	32 481	47 074
50	176 29	18 078	27 000	87	4.92	31 897	46 228
51	184 63	18 059	26 959	88	12 56	29 463	42 700
52	201 68	18 752	27 117	89	15 11	29 279	42 433
53	211 28	18 110	26 246	90	15 22	27 830	40 333
54	227 22	15 484	25 789	91	16 31	27 517	39 890
55	237 69	17 429	25 984	92	17 86	26 664	38 643
56	259 05	18 211	26 393	93	18 86	26 487	38 387
57	277 29	17 749	25 723	94	19 86	25 907	37 546
				95	21 42	25 651	37 175
				96	22 49	25 101	36 378
				97	23 96	24 944	36 151
				98	24 48	24 581	35 825
				99	25 50	24 400	35 267
				100	28 45	23 765	34 442
1 16-9	0.02	935 000	1352 072	101	29 58	24 736	34 849
2 H=87	0.09	381 033	551 251	102	29 97	23 634	34 252
3	0.10	184 511	267 407	103	54 96	21 337	30 923
4	0.15	267 689	397 955	104	79 41	18 870	27 546
5	0.33	164 163	237 917	105	80 40	19 826	28 748
6	0.58	116 624	169 020	106	93 12	19 340	28 629
7	0.71	109 996	159 414	107	112 17	18 747	27 199
8	0.77	102 078	147 939	108	119 10	17 909	25 955
9	0.96	93 886	126 067	109	137 70	18 209	26 390
10	1.01	87 819	127 274	110	152 22	17 442	25 278
11	1.27	80 819	117 129	111	156 96	17 907	25 952
12	1.32	79 515	115 239	112	185 19	17 107	24 793
13	1.59	73 187	106 068	113	188 82	17 545	25 428
14	1.63	72 584	105 194	114	218 28	16 874	24 458
15	1.90	67 245	97 457	115	220 59	17 271	25 020
16	2.00	66 013	95 671	116	244 74	16 743	24 265
17	2.22	63 634	92 223	117	246 12	17 106	24 791
18	2.31	64 013	92 772	118	265 26	17 002	24 641
19	2.31	61 850	89 638	119	277 68	16 646	24 125
20	2.47	61 029	88 449	120	299 79	16 731	24 248
21	2.62	58 863	85 338				
22	2.72	58 931	85 407				
23	2.97	56 329	82 361				
24	3.13	43 061	62 407				
25	7 29	39 284	56 913				
26	9 95	35 747	51 807				
27	10 57	35 823	51 917				
28	13 25	32 716	47 414				
29	13 77	33 059	47 912				
30	16 56	30 853	44 714				
31	16 97	31 146	45 139				
32	19 88	29 443	42 671				
33	20 15	29 763	43 135				
34	22 70	28 791	41 726				
35	22 83	29 763	43 135				
36	23 18	28 241	40 929				
37	25 83	27 596	39 994				

0 sample and hematocrit	1 shear rate (1/sec)	2 viscosity (cP)	3 relative viscosity
1 24-1	0 04	362 590	525 490
2 H=72	0 11	194 284	221 574
3	0 17	131 194	190 136
4	0 30	94 248	126 591
5	0 49	74 338	107 736
6	0 73	53 562	72 119
7	0 80	52 230	70 188
8	1 12	44 720	72 058
9	1 26	42 758	72 171
10	1 31	44 924	65 107
11	1 59	44 327	64 242
12	1 83	41 136	59 617
13	1 92	40 630	59 834
14	2 23	37 542	54 409
15	2 27	37 500	54 348
16	2 45	36 369	52 709
17	2 66	35 525	51 486
18	2 85	33 902	49 133
19	3 41	25 912	27 554
20	7 32	22 434	32 313
21	8 61	21 192	30 713
22	11 28	19 059	27 622
23	12 44	18 331	26 567
24	15 26	16 965	24 597
25	15 62	16 908	24 504
26	18 57	15 828	22 954
27	19 45	15 676	22 748
28	22 56	14 809	21 462
29	22 63	14 518	21 649
30	25 82	14 367	20 891
31	28 53	13 729	19 805
32	45 18	12 102	17 939
33	68 28	10 237	14 815
34	77 34	10 444	15 136
35	101 70	9 462	13 713
36	109 47	9 695	14 051
37	128 75	9 110	13 233
38	135 03	9 331	13 523
39	167 40	9 005	12 051
40	167 80	8 701	12 610
41	199 65	8 824	12 728
42	210 30	8 492	12 337
43	231 57	8 660	12 551
44	243 96	8 344	12 093
45	284 49	8 233	11 932
46	297 99	8 369	12 129

0 sample and hematocrit	1 shear rate (1/sec)	2 viscosity (cP)	3 relative viscosity
1 24-3	0 04	839 680	1216 932
2 H=82	0 10	396 165	571 152
3	0 16	272 445	374 546
4	0 29	188 843	273 686
5	0 54	133 856	193 994
6	0 55	128 070	185 609
7	0 80	104 906	152 038
8	1 00	94 473	136 417
9	1 18	85 417	123 793
10	1 46	77 817	112 778
11	1 70	80 847	117 170
12	1 86	68 732	99 612
13	2 14	63 528	92 070
14	2 19	62 806	91 024
15	2 52	58 919	85 390
16	2 92	54 746	79 342
17	2 94	54 031	76 306
18	3 50	40 153	58 193
19	7 36	36 443	52 814
20	8 68	32 861	47 625
21	10 01	31 905	46 239
22	11 86	29 226	42 357
23	13 97	27 933	40 483
24	15 68	26 448	38 330
25	18 60	25 228	36 562
26	19 51	24 699	35 795
27	21 91	23 868	34 591
28	22 70	23 584	34 180
29	23 67	24 842	35 003
30	26 51	22 621	32 784
31	29 18	21 740	31 507
32	29 98	21 768	31 548
33	49 17	19 280	27 942
34	61 21	16 652	24 133
35	81 00	17 019	24 665
36	105 75	15 305	22 181
37	119 16	15 768	22 823
38	145 38	14 579	21 129
39	157 41	15 028	21 780
40	173 38	14 169	20 535
41	189 30	14 612	21 177
42	211 53	13 906	20 154
43	214 71	14 371	20 628
44	240 12	14 174	20 542
45	244 59	13 707	19 865
46	277 53	13 555	19 645
47	299 82	13 618	19 736

0 sample and hematocrit	1 shear rate (1/sec)	2 viscosity (cP)	3 relative viscosity
1 16-8	0 11	120 563	174 729
2 H=70	0 23	70 732	102 510
3	0 30	57 719	85 491
4	0 36	60 183	87 222
5	0 62	40 984	59 397
6	0 73	33 814	49 006
7	0 87	36 047	52 242
8	0 99	31 862	46 177
9	1 06	33 568	48 678
10	1 31	26 827	40 228
11	1 39	29 025	42 065
12	1 63	27 331	39 610
13	1 79	26 641	38 610
14	1 95	26 210	37 984
15	2 18	24 495	35 500
16	2 21	25 493	36 946
17	2 51	24 168	35 046
18	2 55	24 311	35 233
19	2 98	22 444	32 557
20	4 28	18 569	26 912
21	5 30	17 360	25 159
22	6 82	15 704	22 759
23	9 26	14 512	21 032
24	9 36	14 256	20 661
25	12 55	12 866	18 646
26	12 57	13 142	19 046
27	15 22	12 366	17 922
28	16 39	11 990	17 377
29	17 87	11 776	17 067
30	19 56	11 373	16 483
31	21 17	11 172	16 191
32	22 75	10 888	15 780
33	24 47	10 663	15 454
34	25 94	10 591	15 349
35	29 09	10 159	14 723
36	42 09	9 979	14 452
37	60 06	8 566	12 414
38	67 53	8 863	12 845
39	93 03	8 288	12 012
40	106 26	7 806	11 313
41	124 92	7 877	11 416
42	145 86	7 507	10 860
43	156 69	7 630	11 070
44	172 41	7 360	10 667
45	194 91	7 411	10 741
46	212 01	7 202	10 439
47	233 01	7 281	10 552
48	245 04	7 131	10 335

49	264 90	7 203	14 439	49	143 58	9 609	12 419
50	254 87	7 046	13 241	40	152 79	8 659	12 549
51	299 97	7 389	13 272	41	175 44	9 236	13 336
52	299 73	7 091	13 272	42	165 91	8 594	12 455
53 24-2	0 04	382 258	53 897	43	207 33	9 905	13 905
54 H=70	0 10	182 440	29 406	44	225 57	8 332	12 075
55	0 18	138 322	100 771	45	232 74	8 733	12 657
56	0 23	109 737	159 039	46	264 57	8 592	12 452
57	0 35	85 508	123 625	47	285 09	8 232	11 930
58	0 48	71 038	102 954	48	299 82	8 367	12 126
59	0 67	60 628	57 567	0 sample and 1 shear rate 2 viscosity 3 relative hematocrit (1/sec) (cP) viscosity			
60	0 67	61 676	39 235	-----			
61	0 93	50 402	73 046	1 24-8	0 43	72 067	104 445
62	0 93	52 701	76 378	2 H=62	0 56	67 709	127 114
63	1 18	47 425	58 732	3	0 75	61 207	88 891
64	1 40	44 392	64 236	4	1 07	29 202	36 670
65	1 50	42 089	60 997	5	1 39	21 564	31 252
66	1 73	39 604	57 694	6	1 41	39 418	57 128
67	1 82	39 007	55 532	7	1 65	45 849	66 448
68	2 06	37 135	53 815	8	1 75	30 474	44 165
69	2 14	36 923	53 512	9	1 90	32 871	47 639
70	2 39	35 219	51 042	10	2 15	30 133	43 671
71	2 46	34 910	50 594	11	2 22	31 235	45 268
72	2 85	32 586	47 226	12	2 35	33 865	48 790
73	3 00	32 196	46 661	13	2 48	35 028	50 765
74	4 25	28 011	40 996	14	2 67	28 266	38 570
75	6 80	22 928	33 243	15	2 96	28 339	41 071
76	7 32	21 917	31 764	16	3 00	22 640	32 812
77	9 34	20 405	29 575	17	3 55	32 433	47 033
78	10 63	19 142	27 742	18	5 47	27 197	39 416
79	12 53	18 120	26 362	19	6 70	16 803	24 352
80	15 09	17 133	24 830	20	6 75	25 421	36 837
81	15 26	16 614	24 515	21	9 02	16 027	23 223
82	17 63	16 579	23 726	22	9 29	14 953	21 671
83	19 22	15 708	22 896	23	10 67	11 561	16 755
84	20 17	15 717	22 778	24	12 47	12 299	17 757
85	22 09	15 249	22 100	25	13 32	11 473	15 628
86	22 53	15 023	21 772	26	15 29	11 207	16 250
87	24 64	14 849	21 514	27	15 66	13 541	18 527
88	25 17	14 594	21 151	28	17 35	10 622	17 394
89	28 49	14 035	20 341	29	19 22	11 271	16 335
90	29 98	13 970	20 246	30	20 58	10 496	15 212
91	42 18	13 485	19 543	31	21 41	11 205	15 239
92	61 32	12 049	17 535	32	23 73	9 434	13 672
93	66 87	11 050	16 072	33	23 96	10 883	15 772
94	93 03	10 950	15 928	34	28 51	10 032	14 539
95	93 12	10 450	15 188	35	28 98	9 925	14 384
96	119 55	10 040	14 551	36	42 30	8 823	14 236
97	124 95	10 339	15 057	37	67 83	8 107	11 749
98	156 96	10 036	14 501	38	72 96	8 906	12 907
99	159 15	9 646	13 980	39	99 57	7 150	10 362
100	182 52	9 773	14 171	40	112 71	6 046	10 661
101	192 27	9 410	13 638	41	125 04	6 936	10 052
102	212 13	9 321	13 599	42	144 33	6 922	10 032
103	214 44	9 553	13 298	43	152 43	7 367	10 677
104	258 53	9 189	13 317	44	169 86	6 826	9 893
105	259 02	9 385	13 601	45	185 58	7 346	10 646
106	284 88	9 052	13 119	46	195 18	6 649	9 626
107	289 79	9 136	12 241	47	212 10	7 140	10 348
0 sample and 1 shear rate 2 viscosity 3 relative hematocrit (1/sec) (cP) viscosity				-----			
1 24-7	0 40	84 771	122 257	48	220 48	6 611	9 581
2 H=68	0 59	71 596	123 762	49	229 47	6 899	9 999
3	0 87	67 903	99 410	50	239 85	6 628	9 606
4	0 91	94 226	78 598	51	265 56	6 456	9 357
5	1 23	49 976	72 429	52	275 04	6 856	9 936
6	1 33	33 574	48 658	53	299 76	6 773	9 816
7	1 49	34 342	49 771	0 sample and 1 shear rate 2 viscosity 3 relative hematocrit (1/sec) (cP) viscosity			
8	1 79	40 298	53 403	-----			
9	1 81	32 863	47 628	1 16-11	0 33	30 822	44 684
10	2 13	43 278	52 722	2 H=59	0 39	29 094	42 165
11	2 19	25 356	26 748	3	0 58	25 052	36 207
12	2 38	37 765	54 732	4	0 71	22 420	32 493
13	2 53	28 580	41 142	5	0 83	21 678	31 417
14	2 64	50 296	43 937	6	1 07	21 133	30 628
15	2 73	33 447	43 474	7	1 09	21 231	30 770
16	2 93	25 349	36 738	8	1 26	20 162	29 220
17	3 12	28 871	41 842	9	1 46	19 725	28 587
18	7 30	20 386	29 545	10	1 51	18 496	26 805
19	7 35	18 171	25 335	11	1 64	17 950	26 014
20	11 13	17 811	25 813	12	1 88	18 088	26 649
21	12 00	16 724	24 281	13	2 15	17 517	25 387
22	14 33	15 987	23 175	14	2 25	17 804	25 803
23	15 98	14 685	21 572	15	2 40	17 139	24 839
24	16 88	14 692	21 253	16	2 56	16 891	24 465
25	16 99	15 356	22 255	17	2 72	16 560	24 000
26	20 06	14 219	22 867	18	2 91	16 478	23 891
27	23 25	13 153	22 062	19	2 87	16 586	24 038
28	23 43	13 243	22 193	20	3 00	16 665	24 152
29	26 45	13 286	22 255	21	3 56	15 143	21 946
30	28 59	13 020	22 870	22	6 65	12 155	17 616
31	29 97	12 727	22 445	23	6 75	12 586	18 241
32	35 10	13 222	19 162	24	9 29	10 970	15 899
33	36 92	10 081	14 610	25	9 30	11 025	15 971
34	37 05	10 792	15 641	26	11 85	10 415	15 094
35	46 04	10 711	15 523	27	11 94	10 352	15 002
36	58 85	10 349	14 999	28	14 40	9 458	14 431
37	112 07	9 249	13 404	29	15 26	9 628	13 951
38	117 99	10 003	14 497	30	17 59	9 294	13 475
				31	17 90	9 179	13 303

32	20 53	8 854	12 832	15	4 30	6 9850	10 123
33	21 40	8 855	12 833	16	6 66	6 4600	9 362
34	23 64	8 507	12 329	17	9 39	6 1250	8 977
35	24 59	8 513	12 328	18	9 91	6 0100	8 710
36	26 46	8 129	11 701	19	12 56	5 7890	8 390
37	26 97	8 055	11 674	20	12 58	5 7954	8 399
38	26 09	7 398	10 722	21	15 76	5 5540	8 049
39	31 62	6 475	9 384	22	15 87	5 5230	8 064
40	33 26	5 733	8 338	23	17 10	5 7620	8 352
41	36 97	5 985	8 674	24	18 32	5 3550	7 761
42	36 72	5 475	7 935	25	19 16	5 2860	7 661
43	112 50	5 667	8 213	26	20 89	5 1900	7 522
44	119 43	5 325	7 717	27	22 46	5 1140	7 412
45	128 06	5 443	7 868	28	23 42	5 0920	7 380
46	145 77	5 186	7 516	29	24 44	5 0320	7 294
47	163 65	5 316	7 704	30	25 97	4 9820	7 220
48	165 57	5 110	7 406	31	27 74	4 8770	7 068
49	182 91	5 240	7 594	32	29 97	4 8530	7 033
50	192 03	5 015	7 268	33	35 35	4 1540	6 020
51	208 44	5 145	7 451	34	39 77	3 7160	5 366
52	211 77	4 979	7 216	35	67 21	3 7560	5 443
53	233 67	5 058	7 330	36	112 77	3 5060	5 081
54	238 29	4 935	7 152	37	131 79	3 4790	5 042
55	259 17	5 012	7 264	38	145 80	3 3790	4 897
56	277 92	4 684	7 076	39	170 13	3 3570	4 865
57	299 70	4 905	7 109	40	185 46	3 2760	4 748
0 sample and 1 shear rate 2 viscosity 3 relative				41	208 35	3 2760	4 748
hematocrit (1/sec) (cP) viscosity				42	218 55	3 2190	4 665
1 1a-12	0 70	15 530	22 512	43	244 86	3 1850	4 616
2 H=36	0 89	15 267	22 126	44	246 37	3 1250	4 674
3	0 96	15 047	21 807	45	284 49	3 1550	4 572
4	1 13	14 754	21 383	46	299 73	3 1720	4 597
5	1 28	14 067	20 287	0 sample and 1 shear rate 2 viscosity 3 relative			
6	1 32	14 045	20 358	hematocrit (1/sec) (cP) viscosity			
7	1 46	14 426	20 907	1 33-1	3 76	6 532	9 133
8	1 57	13 395	19 413	2 H=42	6 94	5 521	8 001
9	1 71	12 951	18 772	3	9 24	4 808	6 968
10	1 68	13 462	19 510	4	10 11	4 604	7 235
11	1 97	12 834	18 600	5	12 54	4 617	8 000
12	2 13	13 373	19 410	6	13 93	4 472	8 591
13	2 15	12 942	18 757	7	15 85	4 454	8 454
14	2 37	12 709	18 419	8	17 75	4 226	8 212
15	2 40	12 566	18 212	9	19 15	4 269	8 216
16	2 62	12 208	17 693	10	20 94	4 158	8 056
17	2 72	12 089	17 520	11	23 10	4 168	8 011
18	2 81	12 016	17 414	12	24 11	4 241	8 106
19	3 00	12 181	17 654	13	24 51	4 058	8 831
20	3 66	11 641	16 871	14	25 75	4 061	8 806
21	6 21	9 807	14 215	15	26 65	4 075	8 806
22	7 26	9 041	13 151	16	28 39	3 475	8 036
23	7 75	8 095	12 151	17	32 64	3 161	6 591
24	9 70	8 560	12 450	18	42 34	3 233	4 500
25	11 30	8 294	12 155	19	44 50	3 105	4 368
26	12 54	8 091	11 703	20	119 94	3 014	4 354
27	12 85	7 948	11 519	21	132 00	2 945	4 269
28	15 85	7 490	10 855	22	158 31	2 949	4 204
29	17 64	7 492	10 858	23	164 94	2 901	4 215
30	19 16	7 179	10 604	24	190 05	2 861	4 146
31	19 59	7 185	10 412	25	191 55	2 809	4 108
32	21 80	6 930	10 043	26	221 05	2 810	4 072
33	22 77	6 899	9 960	27	224 49	2 822	4 090
34	24 45	6 696	9 704	28	253 65		
35	25 32	6 778	9 873	29	277 14		
36	29 06	6 453	9 352	30	299 76		
37	29 97	6 454	9 354	0 sample and 1 shear rate 2 viscosity 3 relative			
38	37 26	6 204	9 156	hematocrit (1/sec) (cP) viscosity			
39	69 12	5 417	7 951	1 33-2	3 53	2 451	3 552
40	72 48	5 076	7 348	2 H=25	6 97	1 955	2 853
41	94 47	5 091	7 378	3	9 98	2 285	3 312
42	105 42	4 744	6 875	4	13 28	1 975	2 891
43	119 88	4 895	7 034	5	13 82	2 152	3 119
44	138 42	4 569	6 622	6	16 60	2 001	2 900
45	158 10	4 712	6 829	7	17 01	2 134	3 093
46	164 85	4 495	6 514	8	19 24	2 002	2 901
47	189 99	4 628	6 707	9	20 18	2 142	3 104
48	191 40	4 426	6 429	10	22 53	2 012	2 916
49	215 46	4 571	6 625	11	23 04	2 576	3 733
50	217 77	4 400	6 377	12	23 36	2 050	3 029
51	244 20	4 273	6 338	13	25 18	2 003	2 903
52	247 23	4 515	6 543	14	25 28	2 064	2 991
53	240 43	4 261	6 320	15	29 48	2 001	2 900
54	299 73	4 432	6 394	16	48 39	2 264	3 281
0 sample and 1 shear rate 2 viscosity 3 relative				17	73 89	2 152	3 090
hematocrit (1/sec) (cP) viscosity				18	83 00	1 861	2 697
1 33-3	0 87	11 5050	16 674	19	99 27	2 051	2 972
2 H=48	1 00	12 5410	18 175	20	124 71	1 465	2 891
3	1 06	12 5630	18 252	21	125 07	1 843	2 671
4	1 32	11 3440	16 441	22	150 27	1 962	2 872
5	1 46	9 9990	14 246	23	152 55	1 865	2 703
6	1 64	10 3370	14 601	24	175 89	1 972	2 858
7	1 79	9 7970	14 199	25	178 98	1 867	2 706
8	1 95	9 5760	13 878	26	201 36	1 962	2 843
9	2 18	9 5400	13 826	27	211 95	1 889	2 738
10	2 27	8 6340	12 513	28	226 83	1 933	2 686
11	2 51	8 8680	12 892	29	244 92	1 892	2 742
12	2 53	8 6170	12 199	30	246 03	1 892	2 742
13	2 84	8 0500	11 679	31	284 70	1 914	2 774
14	3 00	9 0480	13 113	32	299 76		

0 sample and hematocrit	1 shear rate (1/sec)	2 viscosity (cP)	3 relative viscosity
1 33-5	24 21	1 341	1 943
2 H=7	28 76	1 195	1 732
3	100 45	1 177	1 706
4	105 87	1 086	1 574
5	138 75	1 171	1 697
6	145 47	1 121	1 625
7	178 56	1 137	1 646
8	163 27	1 173	1 700
9	204 93	1 118	1 620
10	215 19	1 181	1 712
11	231 21	1 128	1 635
12	246 90	1 185	1 717
13	277 35	1 151	1 668
14	299 70	1 178	1 707



**Table F-III:** Data tables for plasma suspensions.

0 sample, hct.	1 shear rate	2 viscosity	3 relative	31	32	33	34	35	36	37	38	39	40	41	42	43	44	45	46	47	48	49	50	51	52	53	54	55	56	57	58	59	60	61	62	63	64	65				
plasma visc.	(1/sec)	(cP)	viscosity																																							
1 31-7	0.04	2084.400	1724.298	34	35	36	37	38	39	40	41	42	43	44	45	46	47	48	49	50	51	52	53	54	55	56	57	58	59	60	61	62	63	64	65							
2 H=98.1	0.10	924.471	745.679	35	36	37	38	39	40	41	42	43	44	45	46	47	48	49	50	51	52	53	54	55	56	57	58	59	60	61	62	63	64	65								
3 1.21 cP	0.17	636.364	542.450	36	37	38	39	40	41	42	43	44	45	46	47	48	49	50	51	52	53	54	55	56	57	58	59	60	61	62	63	64	65									
4	0.23	514.660	425.339	37	38	39	40	41	42	43	44	45	46	47	48	49	50	51	52	53	54	55	56	57	58	59	60	61	62	63	64	65										
5	0.29	433.674	358.410	38	39	40	41	42	43	44	45	46	47	48	49	50	51	52	53	54	55	56	57	58	59	60	61	62	63	64	65											
6	0.36	386.667	319.560	39	40	41	42	43	44	45	46	47	48	49	50	51	52	53	54	55	56	57	58	59	60	61	62	63	64	65												
7	0.42	347.103	286.864	40	41	42	43	44	45	46	47	48	49	50	51	52	53	54	55	56	57	58	59	60	61	62	63	64	65													
8	0.47	277.183	229.079	41	42	43	44	45	46	47	48	49	50	51	52	53	54	55	56	57	58	59	60	61	62	63	64	65														
9	0.47	260.676	215.433	42	43	44	45	46	47	48	49	50	51	52	53	54	55	56	57	58	59	60	61	62	63	64	65															
10	0.67	226.073	186.839	43	44	45	46	47	48	49	50	51	52	53	54	55	56	57	58	59	60	61	62	63	64	65																
11	0.93	222.923	184.234	44	45	46	47	48	49	50	51	52	53	54	55	56	57	58	59	60	61	62	63	64	65																	
12	1.12	197.697	163.386	45	46	47	48	49	50	51	52	53	54	55	56	57	58	59	60	61	62	63	64	65																		
13	1.20	192.832	159.263	46	47	48	49	50	51	52	53	54	55	56	57	58	59	60	61	62	63	64	65																			
14	1.44	174.098	143.883	47	48	49	50	51	52	53	54	55	56	57	58	59	60	61	62	63	64	65																				
15	1.59	168.233	139.034	48	49	50	51	52	53	54	55	56	57	58	59	60	61	62	63	64	65																					
16	1.69	160.638	132.739	49	50	51	52	53	54	55	56	57	58	59	60	61	62	63	64	65																						
17	1.86	153.886	126.831	50	51	52	53	54	55	56	57	58	59	60	61	62	63	64	65																							
18	1.95	151.486	123.193	51	52	53	54	55	56	57	58	59	60	61	62	63	64	65																								
19	2.12	143.930	120.603	52	53	54	55	56	57	58	59	60	61	62	63	64	65																									
20	2.27	142.229	117.545	53	54	55	56	57	58	59	60	61	62	63	64	65																										
21	2.39	138.593	114.540	54	55	56	57	58	59	60	61	62	63	64	65																											
22	2.52	133.784	112.218	55	56	57	58	59	60	61	62	63	64	65																												
23	2.92	127.621	105.472	56	57	58	59	60	61	62	63	64	65																													
24	3.00	127.060	105.008	57	58	59	60	61	62	63	64	65																														
25	3.62	122.020	100.843	58	59	60	61	62	63	64	65																															
26	5.34	103.200	86.942	59	60	61	62	63	64	65																																
27	5.97	99.799	82.479	60	61	62	63	64	65																																	
28	7.44	93.683	79.079	61	62	63	64	65																																		
29	8.60	89.784	74.202	62	63	64	65																																			
30	9.98	88.311	72.984	63	64	65																																				
31	11.24	84.283	69.637	64	65																																					
32	12.53	83.603	69.093	65																																						
33	13.89	80.235	64.326	0 sample, hct.	1 shear rate	2 viscosity	3 relative	31	32	33	34	35	36	37	38	39	40	41	42	43	44	45	46	47	48	49	50	51	52	53	54	55	56	57	58	59	60	61	62	63	64	65
34	15.11	80.298	64.362	1 29-1	0.03	2301.786	2282.653																																			
35	17.20	76.971	63.612	2 H=97.3	0.10	1038.509	989.036																																			
36	18.30	77.341	63.918	3 1.05 cP	0.16	712.404	678.482																																			
37	19.84	73.064	62.034		0.22	555.674	529.213																																			
38	20.21	73.861	62.693		0.29	473.069	450.542																																			
39	22.48	73.483	60.731		0.35	428.400	408.000																																			
40	22.76	74.268	61.379	7	0.48	347.712	331.154																																			
41	25.12	72.116	59.600	8	0.67	292.233	278.336																																			
42	25.31	73.035	60.360	9	0.86	249.948	238.046																																			
43	28.43	70.827	58.533	10	0.94	236.474	223.213																																			
44	29.64	75.911	62.736	11	1.05	222.273	211.689																																			
45	29.97	70.768	58.486	12	1.27	199.386	189.891																																			
46	48.72	69.458	57.403	13	1.31	197.704	188.290																																			
47	73.29	62.053	51.285	14	1.60	173.117	166.778																																			
48	74.16	63.838	54.428	15	1.63	177.716	169.233																																			
49	99.66	64.028	52.916		1.93	160.006	152.387																																			
50	99.72	61.312	50.671	16	1.94	161.925	154.214																																			
51	118.63	63.287	52.303	17	2.26	146.730	139.762																																			

0 sample, hct. plasma visc	1 shear rate (1/sec)	2 viscosity (cP)	3 relative viscosity	36	37	38	39	40	41	42	43	44	45	46	47	48	49	50	51	52	53	54	55	56	57	58	59	60																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																												
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29-5	0.04	1484.507	1413.814	22.49	22.73	24.64	25.14	26.55	27.80	29.97	30.06	49.11	68.13	86.07	87.27	106.32	119.16	131.82	145.47	157.29	172.05	182.82	204.96	208.29	231.36	232.75	264.46	271.65	277.31	290.85	299.70	0.04	2177.444	1799.540	951.163	706.085	0.10	951.163	641.230	543.937	438.940	393.792	344.690	312.801	249.716	211.097	199.091	183.771	164.858	163.929	163.562	152.417	143.600	139.019	134.496	128.497	102.914	82.050	76.731	9.95	74.118	70.385	67.788	66.007	64.971	63.529	62.265	60.901	59.908	58.626	58.341	57.357	57.188	56.158	55.321	53.897	54.310	53.930	52.131	52.289	44.682	45.191	42.937	43.698	42.573	41.935	41.952	41.200	41.467	40.782	41.207	40.430	40.921	40.755	40.160	40.236	39.253	6.255	58.169	56.023	54.551	53.695	52.503	51.459	50.331	49.511	48.451	48.216	47.402	47.263	46.412	45.885	44.194	44.684	44.570	43.083	39.908	36.927	37.348	35.485	34.114	33.184	34.637	34.671	34.074	34.704	34.055	33.413	33.819	33.602	33.190	32.253	1.00	1705.263	1509.082	1307.580	1144.673	1013.135	920.181	844.790	780.024	729.455	684.703	644.494	604.275	564.224	524.275	484.275	444.275	404.275	364.275	324.275	284.275	244.275	204.275	164.275	124.275	84.275	44.275	0.000	201.207	178.059	159.082	144.673	124.135	104.181	84.790	64.024	44.455	24.703	4.494	189.624	164.424	144.275	124.275	104.275	84.275	64.275	44.275	24.275	4.275	0.000	201.207	178.059	159.082	144.673	124.135	104.181	84.790	64.024	44.455	24.703	4.494	189.624	164.424	144.275	124.275	104.275	84.275	64.275	44.275	24.275	4.275	0.000	201.207	178.059	159.082	144.673	124.135	104.181	84.790	64.024	44.455	24.703	4.494	189.624	164.424	144.275	124.275	104.275	84.275	64.275	44.275	24.275	4.275	0.000	201.207	178.059	159.082	144.673	124.135	104.181	84.790	64.024	44.455	24.703	4.494	189.624	164.424	144.275	124.275	104.275	84.275	64.275	44.275	24.275	4.275	0.000	201.207	178.059	159.082	144.673	124.135	104.181	84.790	64.024	44.455	24.703	4.494	189.624	164.424	144.275	124.275	104.275	84.275	64.275	44.275	24.275	4.275	0.000	201.207	178.059	159.082	144.673	124.135	104.181	84.790	64.024	44.455	24.703	4.494	189.624	164.424	144.275	124.275	104.275	84.275	64.275	44.275	24.275	4.275	0.000	201.207	178.059	159.082	144.673	124.135	104.181	84.790	64.024	44.455	24.703	4.494	189.624	164.424	144.275	124.275	104.275	84.275	64.275	44.275	24.275	4.275	0.000	201.207	178.059	159.082	144.673	124.135	104.181	84.790	64.024	44.455	24.703	4.494	189.624	164.424	144.275	124.275	104.275	84.275	64.275	44.275	24.275	4.275	0.000	201.207	178.059	159.082	144.673	124.135	104.181	84.790	64.024	44.455	24.703	4.494	189.624	164.424	144.275	124.275	104.275	84.275	64.275	44.275	24.275	4.275	0.000	201.207	178.059	159.082	144.673	124.135	104.181	84.790	64.024	44.455	24.703	4.494	189.624	164.424	144.275	124.275	104.275	84.275	64.275	44.275	24.275	4.275	0.000	201.207	178.059	159.082	144.673	124.135	104.181	84.790	64.024	44.455	24.703	4.494	189.624	164.424	144.275	124.275	104.275	84.275	64.275	44.275	24.275	4.275	0.000	201.207	178.059	159.082	144.673	124.135	104.181	84.790	64.024	44.455	24.703	4.494	189.624	164.424	144.275	124.275	104.275	84.275	64.275	44.275	24.275	4.275	0.000	201.207	178.059	159.082	144.673	124.135	104.181	84.790	64.024	44.455	24.703	4.494	189.624	164.424	144.275	124.275	104.275	84.275	64.275	44.275	24.275	4.275	0.000	201.207	178.059	159.082	144.673	124.135	104.181	84.790	64.024	44.455	24.703	4.494	189.624	164.424	144.275	124.275	104.275	84.275	64.275	44.275	24.275	4.275	0.000	201.207	178.059	159.082	144.673	124.135	104.181	84.790	64.024	44.455	24.703	4.494	189.624	164.424	144.275	124.275	104.275	84.275	64.275	44.275	24.275	4.275	0.000	201.207	178.059	159.082	144.673	124.135	104.181	84.790	64.024	44.455	24.703	4.494	189.624	164.424	144.275	124.275	104.275	84.275	64.275	44.275	24.275	4.275	0.000	201.207	178.059	159.082	144.673	124.135	104.181	84.790	64.024	44.455	24.703	4.494	189.624	164.424	144.275	124.275	104.275	84.275	64.275	44.275	24.275	4.275	0.000	201.207	178.059	159.082	144.673	124.135	104.181	84.790	64.024	44.455	24.703	4.494	189.624	164.424	144.275	124.275	104.275	84.275	64.275	44.275	24.275	4.275	0.000	201.207	178.059	159.082	144.673	124.135	104.181	84.790	64.024	44.455	24.703	4.494	189.624	164.424	144.275	124.275	104.275	84.275	64.275	44.275	24.275	4.275	0.000	201.207	178.059	159.082	144.673	124.135	104.181	84.790	64.024	44.455	24.703	4.494	189.624	164.424	144.275	124.275	104.275	84.275	64.275	44.275	24.275	4.275	0.000	201.207	178.059	159.082	144.673	124.135	104.181	84.790	64.024	44.455	24.703	4.494	189.624	164.424	144.275	124.275	104.275	84.275	64.275	44.275	24.275	4.275	0.000	201.207	178.059	159.082	144.673	124.135	104.181	84.790	64.024	44.455	24.703	4.494	189.624	164.424	144.275	124.275	104.275	84.275	64.275	44.275	24.275	4.275	0.000	201.207	178.059	159.082	144.673	124.135	104.181	84.790	64.024	44.455	24.703	4.494	189.624	164.424	144.275	124.275	104.275	84.275	64.275	44.275	24.275	4.275	0.000	201.207	178.059	159.082	144.673	124.135	104.181	84.790	64.024	44.455	24.703	4.494	189.624	164.424	144.275	124.275	104.275	84.275	64.275	44.275	24.275	4.275	0.000	201.207	178.059	159.082	144.673	124.135	104.181	84.790	64.024	44.455	24.703	4.494	189.624	164.424	144.275	124.275	104.275	84.275	64.275	44.275	24.275	4.275	0.000	201.207	178.059	159.082	144.673	124.135	104.181	84.790	64.024	44.455	24.703	4.494	189.624	164.424	144.275	124.275	104.275	84.275	64.275	44.275	24.275	4.275	0.000	201.207	178.059	159.082	144.673	124.135	104.181	84.790	64.024	44.455	24.703	4.494	189.624	164.424	144.275	124.275	104.275	84.275	64.275	44.275	24.275	4.275	0.000	201.207	178.059	159.082	144.673	124.135	104.181	84.790	64.024	44.455	24.703	4.494	189.624	164.424	144.275	124.275	104.275	84.275	64.275	44.275	24.275	4.275	0.000	201.207	178.059	159.082	144.673	124.135	104.181	84.790	64.024	44.455	24.703	4.494	189.624	164.424	144.275	124.275	104.275	84.275	64.275	44.275	24.275	4.275	0.000	201.207	178.059	159.082	144.673	124.135	104.181	84.790	64.024	44.455	24.703	4.494	189.624	164.424	144.275	124.275	104.275	84.275	64.275	44.275	24.275	4.275	0.000	201.207	178.059	159.082	144.673	124.135	104.181	84.790	64.024	44.455	24.703	4.494	189.624	164.424	144.275	124.275	104.275	84.275	64.275	44.275	24.275	4.275	0.000	201.207	178.059	159.082	144.673	124.135	104.181	84.790	64.024	44.455	24.703	4.494	189.624	164.424	144.275	124.275	104.275	84.275	64.275	44.275	24.275	4.275	0.000	201.207	178.059	159.082	144.673	124.135	104.181	84.790	64.024	44.455	24.703	4.494	189.624	164.424	144.275	124.275	104.275	84.275	64.275	44.275	24.275	4.275	0.000	201.207	178.059	159.082	144.673	124.135	104.181	84.790	64.024	44.455	24.703	4.494	189.624	164.424	144.275	124.275	104.275	84.275	64.275	44.275	24.275	4.275	0.000	201.207	178.059	159.082	144.673	124.135	104.181	84.790	64.024	44.455	24.703	4.494	189.624	164.424	144.275	124.275	104.275	84.275	64.275	44.275	24.275	4.275	0.000	201.207	178.059	159.082	144.673	124.135	104.181	84.790	64.024	44.455	24.703	4.494	189.624	164.424	144.275	124.275	104.275	84.275	64.275	44.275	24.275	4.275	0.000	201.207	178.059	159.082	144.673	124.135	104.181	84.790	64.024	44.455	24.703	4.494	189.624	164.424	144.275	124.275	104.275	84.275	64.275	44.275	24.275	4.275	0.000	201.207	178.059	159.082	144.673	124.135	104.181	84.790	64.024	44.455	24.703	4.494	189.624	164.424	144.275	124.275	104.275	84.275	64.275	44.275	24.275	4.275	0.000	201.207	178.059	159.082	144.673	124.135	104.181	84.790	64.024	44.455	24.703	4.494	189.624	164.424	144.275	124.275	104.275	84.275	64.275	44.275	24.275	4.275	0.000	201.207	178.059	159.082	144.673	124.135	104.181	84.790	64.024	44.455	24.703	4.494	189.624	164.424	144.275	124.275	104.275	84.275	64.275	44.275	24.275	4.275	0.000	201.207	178.059	159.082	144.673	124.135	104.181	84.790	64.024	44.455	24.703	4.494	189.624	164.424	144.275	124.275	104.275	84.275	64.275	44.275	24.275	4.275	0.000	201.207	178.059</

12	1.060	182.468	161.476				
13	1.250	169.353	146.370				
14	1.260	173.124	153.207	49		125.16	31.057
15	1.500	151.986	134.501	50		133.11	29.707
16	1.520	156.679	138.654	51		150.69	30.338
17	1.700	142.276	125.908	52		166.05	29.075
18	1.790	142.388	126.007	53		175.98	29.714
19	1.890	134.182	118.745	54		192.60	28.687
20	2.050	132.048	116.857	55		201.40	29.299
21	2.140	126.959	112.353	56		219.09	28.381
22	2.320	123.728	109.494	57		226.95	28.972
23	2.390	120.316	106.474	58		238.83	28.234
24	2.400	121.402	107.435	59		246.18	28.781
25	2.520	118.321	104.709	60		271.44	28.555
26	2.650	114.522	101.347	61		278.25	28.023
27	2.850	111.078	98.299	62		299.64	28.132
28	3.000	108.388	95.919				
29	4.940	86.278	76.352				
30	7.490	72.527	64.183				
31	9.280	68.429	60.734				
32	10.040	65.202	57.701				
33	12.580	61.445	54.376				
34	12.590	60.262	53.329				
35	15.130	56.811	50.275				
36	15.230	57.573	50.950				
37	17.680	54.326	48.076				
38	19.200	53.148	47.034				
39	20.210	52.205	46.199				
40	22.490	50.834	44.986				
41	23.390	50.224	44.446				
42	24.480	49.540	43.841				
43	25.310	49.259	43.592				
44	27.130	48.020	42.496				
45	29.970	47.092	41.674				
46	36.480	46.234	40.915				
47	49.170	42.941	38.001				
48	74.610	39.445	34.907				
49	92.190	37.761	33.617				
50	100.140	37.478	33.166				
51	125.520	36.281	32.107				
52	131.820	35.858	31.733				
53	151.050	35.323	31.426				
54	164.730	34.901	30.886				
55	170.160	35.025	30.996				
56	197.730	34.247	30.307				
57	201.960	34.456	30.492				
58	224.220	33.902	30.002				
59	227.370	34.136	30.209				
60	250.620	33.591	29.727				
61	252.840	33.873	29.976				
62	276.960	33.354	29.517				
63	299.700	33.377	29.537				
0 sample, hct. 1 shear rate 2 viscosity 3 relative							
plasma visc. (1/sec) (cP) viscosity							
1 29-7	0.04	1077.778	1026.455				
2 H=92	0.10	513.043	488.612				
3 1.05 cP	0.17	375.090	357.229				
	0.23	298.960	284.724				
	0.30	268.293	255.517				
	0.36	220.903	210.384				
	0.42	222.656	212.053				
	0.68	160.752	153.097				
	0.93	144.380	137.505				
	0.93	135.167	128.730				
	1.19	124.680	118.743				
	1.19	119.025	113.357				
	1.38	110.207	104.959				
	1.46	111.550	106.238				
	1.57	103.151	98.239				
	1.76	98.423	93.746				
	1.79	101.039	96.228				
	2.06	93.491	89.039				
	2.08	91.644	87.299				
	2.32	88.849	84.618				
	2.40	86.422	82.307				
	2.52	84.759	80.723				
	2.53	84.719	80.685				
	2.85	80.286	76.463				
	2.93	83.265	79.300				
	3.00	79.181	75.410				
	4.85	67.079	63.885				
	6.76	59.458	56.627				
	8.63	54.141	51.563				
	10.63	50.635	48.224				
	11.84	49.151	46.810				
	13.93	46.457	44.245				
	15.67	45.165	43.014				
	17.24	43.761	41.677				
35		18.85	42.879				
36		19.88	42.040				
37		22.52	40.742				
38		22.67	40.934				
39		25.16	39.533				
40		26.48	39.379				
41		27.80	38.435				
42		29.97	38.154				
43		36.42	37.463				
44		61.95	33.850				
45		80.94	32.357				
46		92.67	30.185				
47		106.50	31.031				
48		125.67	29.119				
49		125.67	30.408				
50		151.29	29.827				
51		152.22	28.652				
52		170.37	29.382				
53		178.62	28.317				
54		189.45	29.083				
55		204.99	28.090				
56		214.95	28.765				
57		224.91	27.955				
58		233.94	28.369				
59		246.74	27.838				
60		253.08	28.407				
61		277.68	27.725				
62		299.70	27.832				
0 sample, hct. 1 shear rate 2 viscosity 3 relative							
plasma visc. (1/sec) (cP) viscosity							
1 31-9	0.02	1694.284	1401.972				
2 H=92	0.09	600.342	496.150				
3 1.21 cP	0.15	391.849	323.842				
	0.22	318.104	262.898				
	0.28	257.725	212.996				
	0.34	233.772	193.200				
	0.41	212.841	175.902				
	0.60	168.373	139.151				
	0.79	146.268	120.883				
	0.81	150.595	124.454				
	1.04	125.381	103.621				
	1.14	124.447	102.849				
	1.30	112.248	92.767				
	1.47	108.170	89.397				
	1.49	105.344	87.061				
	1.68	99.375	82.128				
35							
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17	1 73	100.346	82 931				
18	1 87	94.320	77 950	116	207 90	28 693	23 713
19	2 06	91.471	75 996	117	225 36	27 891	23 050
20	2 19	87.819	72 578	118	233 43	28 385	23 439
21	2 33	86.752	71 779	119	245 19	27 732	22 919
22	2 44	84.201	69 589	120	252 66	28 180	23 289
23	2 53	83.197	68 758	121	264 85	27 539	22 760
24	2 79	79.632	65 812	122	299 70	27.646	22 848
25	3.00	77.345	64 087				
26	3.02	77.237	63 952	0 sample, hct.	1 shear rate	2 viscosity	3 relative
27	5 57	60.857	50 295	plasma visc	(1/sec)	(cP)	viscosity
28	8 74	52.095	43 054	1 30-1	0.04	1059.701	937 788
29	9 95	49.653	41 036	2 H=91	0.10	475.793	421 056
30	12.57	46.408	38 354	3 1.13 cP	0.17	334.234	295 782
31	12.57	46.779	38 660		0.23	277 143	245 259
32	15.23	44.049	36 404		0.29	233.978	207 060
33	15.76	43.993	36 358		0.36	199.079	176 176
34	17.87	42.237	34 907		0.42	182 717	161.696
35	18.95	41.917	34 642		0.49	168 850	149 425
36	20.50	40.824	33 739		0.61	146 090	129 283
37	22.14	40.487	33 460		0.80	132.657	117 396
38	23.81	39 253	32 440		0.80	122.741	108 620
39	24.68	39 325	32 500		0.99	109 469	96 875
40	27.77	38 029	31 429		1.19	101 910	90 184
41	29.97	37 608	31 081		1.25	94 957	84 033
42	36.42	37.395	31 070		1.31	87 682	77 395
43	61.98	33 340	27 554		1.59	87 427	77 369
44	87.45	31.314	25 879		1.82	77 988	69 016
45	92.10	29 498	24 379		1.86	79 800	70 619
46	112 77	29 976	24 774		2.12	74 147	65 617
47	125.25	28 340	23 421		2.14	72 606	64 253
48	144.60	28 946	23 922		2.32	70 199	62 123
49	151.71	27.819	22 991		2.46	67 771	59 974
50	182.79	28 118	23 238		2.52	67 048	59 335
51	184 71	27 267	22 535		2.71	65 300	57 788
52	208.29	27 752	22 936		2.91	62 858	55 627
53	211.17	27 004	22 317		2.97	76 438	67 644
54	237 63	26 784	22 136		3.00	62 847	55 617
55	240 00	27 340	22 595		3.52	58 474	51 747
56	250 92	26.686	22 055		8.05	50 838	44 989
57	264 15	25 630	22 008		9 34	52 473	46 436
58	265 47	27 104	22 400		11.24	44 290	39 195
59	297 24	26 463	21 870		13 79	41 581	36 797
60 31-10	0 04	987 302	815 952		13.97	44 064	38 995
61 H=92	0 10	498 827	412 254		15.69	40 422	35 772
62 1.21 cP	0 17	359 328	297 131		17 29	40 212	35 586
63	0 23	293 864	242 863		20.15	37 459	33 166
64	0 29	252 875	208 988		20 58	37 393	33 091
65	0 36	223 038	184 329		23 22	36 408	32 219
66	0 48	189 745	156 814		23 33	35 952	31 816
67	0 61	166 912	137 944		25 74	35 310	31 248
68	0 80	143 350	118 471		25 85	35 292	31 232
69	0 93	139 614	115 383		27 14	34 610	30 628
70	0 99	128 070	105 843		28 52	34 314	30 366
71	1 18	117 275	96 921		29 97	34 184	30 251
72	1 20	121 518	100 428		30 12	29 482	26 090
73	1 44	106 461	87 984		42 78	26 999	23 893
74	1 53	107 350	88 719		68 25	24 484	21 667
75	1 63	100 147	82 766		92 49	22 251	19 691
76	1 79	99 381	82 133		100 05	22 909	20 273
77	1 82	96 410	79 678		132 09	21 111	18 682
78	2 06	93 325	77 128		138 24	21 853	19 339
79	2 08	90 939	75 156		165 12	20 394	18 225
80	2 32	88 127	72 822		176 46	21 192	18 754
81	2 33	87 252	72 109		198 15	20 257	17 927
82	2 52	84 805	70 087		214 62	20 758	18 370
83	2 58	83 459	68 474		224 64	20 085	17 774
84	2 78	80 899	66 859		250 92	19 984	17 685
85	2 78	80 595	66 607		259 08	20 432	18 081
86	3 00	78 797	65 121		277 38	19 895	17 606
87	3 52	63 308	52 321		299 70	19 995	17 695
88	7 77	55 102	45 539	31-4	0 03	1750.000	1446 281
89	8 06	55 452	45 826	H=91	0 09	658 938	544 593
90	9 98	51 684	42 714	1 21 cP	0 15	438 566	362 451
91	11.28	49 375	40 806		0 22	354 658	293 106
92	12.53	48 480	40 046		0 28	293 008	242 155
93	14 58	45 947	37 973		0 35	264 935	218 955
94	15 10	46 146	38 137		0 41	236 503	195 457
95	17 24	43 866	36 253		0 47	212 873	175 928
96	17 65	44 322	36 630		0 67	173 354	143 268
97	19 22	42 545	35 161		0 86	149 823	123 822
98	20 17	42 824	35 392		0 94	151 149	124 917
99	22 52	41 239	34 082		1 05	134 154	110 871
100	22 73	41 480	34 281		1 20	129 661	107 158
101	24 51	40 416	33 402		1 24	122 298	101 073
102	25 28	40 494	33 666		1 40	118 079	97 586
103	28 47	38 862	32 117		1 43	113 556	93 831
104	29 97	38 735	32 012		1 62	106 498	88 015
105	35 91	38 772	32 043		1 67	108 155	89 384
106	61 41	34 455	28 475		1 81	100 878	83 370
107	86 79	32 354	26 735		1 86	100 997	83 469
108	93 12	30 570	25 264		2 00	95 351	78 966
109	112 29	30 946	25 372		2 06	96 105	79 426
110	126 18	29 280	24 196		2 26	91 429	75 561
111	150 60	29 797	24 626		2 26	90 519	74 809
112	152 70	28 745	23 756		2 45	86 837	71 766
113	172 53	28 430	23 496		2 46	87 630	72 421
114	182 43	29 066	24 025		2 92	83 848	69 312
115	205 47	28 042	23 175		3 00	80 100	65 198
			89			79 281	65 521

90	3.02	74.763	63.440	14	1.31	97.458	92.817
91	4.92	61.829	51.098	15	1.50	90.212	85.916
92	7.47	51.888	42.883	16	1.53	93.410	88.962
93	9.28	48.997	40.493	17	1.69	85.167	81.111
94	10.66	45.650	37.727	18	1.80	85.903	81.812
95	12.59	43.943	36.318	19	2.01	78.201	74.477
96	13.20	42.354	35.003	20	2.06	79.444	75.661
97	15.89	40.427	33.411	21	2.33	74.297	70.759
98	16.39	39.610	32.736	22	2.33	72.725	69.262
99	18.93	37.958	31.370	23	2.52	70.639	67.275
100	19.20	37.422	31.340	24	2.52	69.574	66.261
101	21.48	36.741	30.364	25	2.79	67.589	64.370
102	22.50	36.285	29.988	26	3.00	65.259	62.151
103	23.38	35.859	29.636	27	3.57	58.814	56.013
104	24.49	35.361	29.224	28	5.48	48.020	45.733
105	25.93	34.865	28.831	29	6.76	44.021	41.925
106	27.12	34.248	28.304	30	9.30	38.451	36.620
107	29.77	33.476	27.664	31	9.33	39.994	38.090
108	36.48	32.689	27.016	32	11.85	35.310	33.629
0 sample, Act.	1 shear rate	2 viscosity	3 relative	33	11.97	36.292	34.564
plasma visc	(1/sec)	(cP)	viscosity	34	13.76	32.622	32.021
1 30-3	0.04	1037.398	918.051	35	14.63	33.525	31.929
2 H=98	0.10	484.084	428.393	36	16.61	32.054	30.528
3 1.13 cP	0.16	360.221	318.780	37	16.96	31.385	29.890
4	0.23	291.029	257.548	38	19.26	30.369	28.923
5	0.29	256.481	226.974	39	20.13	29.759	28.342
6	0.35	222.589	196.981	40	21.25	29.449	28.047
7	0.35	222.589	196.981	41	22.69	28.719	27.351
8	0.42	201.935	178.704	42	23.89	28.403	27.050
9	0.48	185.305	163.987	43	25.24	27.871	26.544
10	0.67	155.372	137.497	44	27.86	27.150	25.857
11	0.86	136.439	120.742	45	29.98	26.688	25.417
12	0.94	136.337	120.652	46	36.15	25.768	24.541
13	1.06	122.150	108.087	47	55.29	23.087	21.988
14	1.20	119.770	105.991	48	87.03	21.010	20.010
15	1.25	112.013	99.127	49	92.70	20.437	19.464
16	1.44	103.735	91.801	50	112.53	20.008	19.055
17	1.60	101.353	89.693	51	125.76	19.358	18.456
18	1.63	96.409	85.318	52	138.00	19.391	18.468
19	1.82	92.121	81.523	53	152.22	18.841	17.944
20	1.86	93.763	82.976	54	157.02	19.049	18.142
21	2.01	87.325	77.279	55	178.59	18.411	17.534
22	2.13	88.655	78.456	56	188.91	18.500	17.695
23	2.27	83.503	73.896	57	205.02	18.123	17.260
24	2.39	82.498	73.007	58	214.41	18.315	17.443
25	2.52	79.586	70.430	59	224.94	17.945	17.090
26	2.79	76.467	67.670	60	239.76	18.106	17.244
27	3.00	74.324	65.773	61	251.22	17.763	16.917
28	3.51	70.445	62.341	62	258.90	17.966	17.110
29	6.70	53.135	47.022	63	271.02	17.644	16.804
30	9.32	47.989	42.468	64	299.70	17.613	16.774
31	9.87	45.913	40.631	65 31-2	0.04	877.600	725.289
32	12.62	42.906	37.970	66 H=89	0.10	425.000	351.240
33	13.05	41.665	36.872	67 1.21 cP	0.16	304.762	251.869
34	15.28	40.143	35.525	68	0.23	240.682	198.911
35	16.89	38.625	34.181	69	0.29	210.072	173.613
36	17.93	36.050	33.673	70	0.35	182.149	150.536
37	20.73	34.460	32.265	71	0.42	167.001	138.017
38	21.88	33.947	31.812	72	0.48	152.861	126.331
39	23.87	34.984	30.959	73	0.55	145.360	120.132
40	23.90	35.053	31.020	74	0.74	122.353	101.118
41	27.17	32.701	29.824	75	0.93	113.962	94.167
42	29.16	33.114	29.304	76	0.93	108.573	89.730
43	29.94	34.369	30.415	77	1.18	94.367	77.989
44	29.97	32.943	29.153	78	1.20	97.923	80.928
45	48.96	30.300	26.814	79	1.37	87.413	72.242
46	74.43	27.670	24.487	80	1.53	84.154	71.202
47	92.31	25.528	22.591	81	1.63	80.520	66.545
48	93.42	26.557	23.502	82	1.79	79.569	65.760
49	118.92	25.505	22.571	83	1.88	75.287	62.221
50	125.49	24.384	21.579	84	2.12	73.020	60.347
51	150.81	24.707	21.865	85	2.14	70.892	58.588
52	158.67	23.700	20.973	86	2.39	68.674	56.757
53	169.98	24.294	21.499	87	2.39	67.757	55.998
54	191.67	23.227	20.555	88	2.59	65.928	54.486
55	201.63	23.799	21.061	89	2.65	65.005	53.723
56	224.64	22.923	20.286	90	2.78	63.158	52.197
57	227.16	23.494	20.791	91	3.00	61.688	50.982
58	251.07	22.751	20.134	92	3.56	55.173	45.598
59	252.66	23.243	20.569	93	7.39	41.162	34.018
60	270.96	22.669	20.061	94	9.29	38.110	31.496
61	290.76	22.570	19.973	95	10.56	35.947	29.708
62	299.73	22.700	20.088	96	13.75	33.227	27.460
0 sample, Act.	1 shear rate	2 viscosity	3 relative	97	13.92	33.096	27.352
plasma visc	(1/sec)	(cP)	viscosity	98	14.95	31.274	25.846
1 29-2	0.04	1063.780	1013.124	99	17.23	31.071	25.679
2 H=89	0.10	456.509	434.770	100	19.49	30.125	24.897
3 1.05 cP	0.17	322.364	307.013	101	20.54	29.501	24.381
4	0.23	272.405	259.433	102	22.67	28.864	23.855
5	0.29	232.856	221.768	103	23.18	28.435	23.500
6	0.36	202.277	192.645	104	25.16	27.764	22.945
7	0.42	185.877	177.026	105	25.21	28.079	23.206
8	0.48	170.329	162.218	106	27.11	27.590	22.802
9	0.67	142.584	135.794	107	27.14	27.274	22.540
10	0.93	117.200	111.619	108	29.78	26.718	22.081
11	0.94	124.225	118.310	109	42.42	25.389	20.983
12	1.12	106.056	101.006	110	48.01	22.805	18.867
13	1.27	103.806	98.863	111	95.03	20.477	16.923
				112	106.14	20.916	17.286

103	126.33	19.405	16.103			
114	137.91	20.067	16.584			
115	159.33	18.989	15.693	17	1.86	61.413
116	176.13	19.392	16.026	18	1.88	58.639
117	185.82	18.694	15.451	19	2.07	55.876
118	207.99	19.003	15.705	20	2.13	57.347
119	212.22	18.490	15.281	21	2.32	53.764
120	238.62	18.293	15.118	22	2.39	53.900
121	239.76	18.694	15.450	23	2.51	51.915
122	264.93	18.158	15.007	24	2.70	50.216
123	271.59	18.463	15.259	25	2.72	49.972
124	299.70	18.143	14.994	26	2.92	47.995
				27	3.00	47.998
0 sample, hct.	1 shear rate	2 viscosity	3 relative	28	3.56	43.067
plasma visc.	(1/sec)	(cP)	viscosity	29	6.11	34.740
				30	8.01	31.225
1 29-8	0.05	573.718	546.398	31	9.33	30.164
2 H=82	0.11	332.065	315.252	32	10.56	28.457
3 1.05 cP	0.17	236.379	225.123	33	12.48	26.870
4	0.24	197.089	187.704	34	12.64	27.051
5	0.30	168.091	160.087	35	15.03	25.389
6	0.36	149.630	142.505	36	15.29	25.443
7	0.43	135.133	128.698	37	16.95	24.504
8	0.55	121.038	115.074	38	17.93	24.206
9	0.75	102.570	97.686	39	19.91	23.358
10	0.93	98.448	93.760	40	20.14	23.319
11	0.94	92.603	88.193	41	22.55	22.593
12	1.13	82.190	78.276	42	22.68	22.517
13	1.19	84.145	82.043	43	25.19	21.810
14	1.32	76.990	73.324	44	25.24	21.938
15	1.39	79.771	75.972	45	27.85	21.299
16	1.57	71.034	67.631	46	29.43	21.509
17	1.65	72.193	68.735	47	29.97	20.955
18	1.83	66.710	63.533	48	48.45	19.009
19	1.92	67.402	64.192	49	73.92	17.290
20	2.08	62.660	59.674	50	92.76	16.349
21	2.18	62.606	59.625	51	99.45	16.486
22	2.34	60.326	57.453	52	123.88	15.586
23	2.45	58.848	56.046	53	131.31	15.810
24	2.59	57.368	54.636	54	152.34	15.213
25	2.71	56.215	53.538	55	156.75	15.445
26	2.98	53.849	51.285	56	182.19	15.165
27	3.00	53.554	51.004	57	185.37	14.873
28	4.22	45.043	42.898	58	207.63	14.947
29	6.13	38.894	37.042	59	211.95	14.692
30	8.68	33.846	32.234	60	238.26	14.518
31	9.30	33.727	32.121	61	239.37	14.739
32	11.23	31.266	29.777	62	258.09	14.402
33	11.93	30.995	29.519	63	264.96	14.612
34	13.77	29.078	27.693	64	271.47	14.361
				65	297.73	14.272
35	13.92	29.341	27.944			
36	16.32	27.568	26.255	0 sample, hct.	1 shear rate	2 viscosity
37	16.57	27.688	26.368	plasma visc.	(1/sec)	(cP)
38	18.56	26.717	25.449			3 relative
39	18.88	26.587	25.321	1 31-8	0.04	viscosity
40	20.55	26.033	24.783	2 H=82	0.10	625.185
41	21.42	25.352	24.430	3 1.21 cP	0.17	300.289
42	22.53	25.415	24.505		0.23	224.417
43	23.33	25.123	23.927		0.29	177.403
44	24.51	24.858	23.455		0.36	160.529
45	25.24	24.634	23.461		0.42	140.586
46	27.14	24.287	23.100		0.55	128.196
47	27.17	24.066	22.920		0.73	108.515
48	29.15	23.688	22.560		0.74	100.330
49	29.97	23.682	22.554		0.92	94.161
50	42.03	22.091	21.029		0.93	88.774
51	73.80	19.654	18.718		1.12	83.881
52	93.18	18.259	17.370		1.19	76.578
53	112.02	18.278	17.408		1.45	77.834
54	126.24	17.395	16.567		1.51	70.481
55	143.88	17.640	16.800		1.78	65.644
56	159.18	16.943	16.136		1.82	63.312
57	182.07	17.087	16.273		2.08	60.293
58	198.84	16.551	15.763		2.11	56.658
59	213.84	16.765	15.967		2.27	57.777
60	225.21	16.358	15.579		2.31	54.309
61	245.70	16.528	15.750		2.38	50.584
62	251.73	16.208	15.436		2.53	54.258
63	270.81	16.395	15.614		2.84	51.728
64	271.52	16.136	15.368		3.00	49.335
65	298.05	16.019	15.256		4.85	48.038
0 sample, hct.	1 shear rate	2 viscosity	3 relative		7.40	36.623
plasma visc.	(1/sec)	(cP)	viscosity		9.95	30.998
					9.97	27.864
1 29-10	0.03	731.532	696.697			28.203
2 H=80	0.10	322.981	307.601		11.94	26.488
3 1.05 cP	0.16	229.268	218.350		12.50	25.732
4	0.22	175.368	167.017		14.60	24.641
5	0.29	156.250	148.810		15.05	24.329
6	0.35	144.103	137.241		17.24	23.309
7	0.42	127.220	121.162		17.59	23.252
8	0.61	101.535	96.700		19.21	22.914
9	0.80	89.755	85.481		20.14	22.301
10	0.93	87.058	82.912		21.20	21.879
11	0.99	80.967	77.111		22.68	21.537
12	1.18	73.393	70.089		23.18	21.319
13	1.26	74.727	71.169		24.60	21.079
14	1.43	67.232	64.030		25.16	20.826
15	1.60	65.796	62.463		26.51	20.662
16	1.62	62.696	59.710		27.81	20.308
					29.22	20.334
					29.80	19.941
					29.97	19.976

49	34.63	17.490	14.455	10	0.86	73.341	65.081
50	80.04	16.267	13.444	11	0.93	68.391	60.523
51	99.84	14.934	12.342	12	1.12	65.348	57.830
52	105.37	15.416	12.740	13	1.19	60.878	53.874
53	126.30	14.477	11.764	14	1.39	59.139	52.335
54	131.04	14.892	12.307	15	1.44	56.073	49.624
55	146.07	14.202	11.737	16	1.39	55.600	49.204
56	156.36	14.334	12.012	17	1.76	50.962	45.099
57	165.99	13.962	11.539	18	1.85	51.240	45.345
58	181.86	14.236	11.765	19	2.05	49.070	43.425
59	192.30	13.752	11.365	20	2.08	47.054	41.641
60	207.39	14.017	11.584	21	2.27	45.504	40.269
61	212.22	13.649	11.280	22	2.31	45.945	40.659
62	226.33	13.882	11.473	23	2.58	43.451	38.452
63	231.93	13.356	11.203	24	2.59	42.928	37.989
64	245.55	13.769	11.379	25	2.04	41.448	36.680
65	251.79	13.481	11.141	26	3.00	40.311	35.850
66	264.60	13.656	11.286	27	4.93	31.832	28.170
67	284.88	13.396	11.071	28	8.10	25.444	22.519
68	299.67	13.415	11.087	29	9.24	25.000	22.124
0 sample, hct.	plasma visc	1 shear rate (1/sec)	2 viscosity (cP)	3 relative viscosity	10.65	23.083	20.427
1 31-2	0.03	1123.276	928.327	30	12.55	22.151	19.403
2 H=78	0.10	507.879	419.735	31	12.57	21.556	19.076
3 1.21 cP	0.16	300.738	248.544	32	15.76	19.977	17.679
4	0.23	233.201	192.728	33	15.87	20.352	18.011
5	0.29	197.098	162.891	34	18.32	19.060	16.867
6	0.35	162.073	150.474	35	19.19	18.837	16.670
7	0.42	162.923	134.647	36	20.87	18.218	16.122
8	0.61	128.416	106.129	37	21.82	18.114	16.032
9	0.80	105.115	86.872	38	23.40	17.705	15.668
10	0.94	118.269	97.743	39	24.47	17.484	15.473
11	1.05	91.028	75.230	40	25.31	17.307	15.316
12	1.20	100.100	82.727	41	27.75	16.785	14.854
13	1.31	84.878	70.167	42	29.97	16.437	14.546
14	1.47	84.677	69.981	43	42.48	15.085	13.350
15	1.50	78.305	64.715	44	41.65	13.752	12.170
16	1.69	72.869	60.222	45	87.03	12.796	11.324
17	1.73	75.225	62.169	46	92.40	12.422	10.993
18	1.93	74.825	61.839	47	112.47	12.195	10.792
19	1.94	70.489	58.255	48	118.89	11.870	10.504
20	2.13	67.287	55.609	49	138.00	11.817	10.458
21	2.20	64.943	53.672	50	163.56	11.337	10.210
22	2.33	66.241	54.765	51	165.15	11.353	10.047
23	2.39	62.065	51.293	52	204.75	11.059	9.787
24	2.53	62.870	51.959	53	208.08	11.205	9.916
25	2.58	60.272	49.812	54	237.78	10.898	9.644
26	2.79	58.424	48.450	55	239.76	11.039	9.769
27	2.99	56.079	46.366	56	271.62	10.923	9.666
28	3.55	52.646	43.509	57	297.42	10.690	9.460
29	5.45	42.713	35.300	58	299.70	10.761	9.523
30	8.00	35.896	29.666	59	0 sample, hct.	plasma visc	1 shear rate (1/sec)
31	9.36	34.875	28.822	60	1 31-3	0.04	262.590
32	9.91	32.889	27.181	61	2 H=66	0.10	144.828
33	12.01	31.364	25.921	62	3 1.21 cP	0.17	106.798
34	12.45	29.942	24.745	63		0.23	89.406
35	14.36	28.371	23.447	64		0.30	78.362
36	15.31	28.425	23.492	65		0.36	70.701
37	16.91	26.780	22.132	66		0.42	67.472
38	18.63	26.270	21.711	67		0.46	63.194
39	19.46	25.478	21.056	68		0.55	59.148
40	21.28	24.971	20.437	69		0.69	54.674
41	22.03	24.406	20.170	70		0.93	46.658
42	23.72	23.877	19.733	71		0.93	46.972
43	23.92	23.788	19.660	72		1.19	43.119
44	25.85	23.249	19.214	73		1.19	42.799
45	27.88	22.701	18.761	74		1.46	39.773
46	29.97	22.265	18.401	75		1.51	38.765
47	36.00	21.175	17.500	76		1.65	37.781
48	61.50	18.220	15.058	77		1.76	36.543
49	86.85	16.870	13.942	78		1.92	35.453
50	93.09	16.376	13.536	79		2.08	34.223
51	112.41	15.997	13.221	80		2.18	33.828
52	119.55	15.591	12.885	81		2.38	32.649
53	138.06	15.443	12.763	82		2.40	32.436
54	146.07	15.089	12.470	83		2.65	31.098
55	169.86	14.963	12.366	84		2.78	30.351
56	172.59	14.799	12.231	85		3.00	29.723
57	201.66	14.628	12.087	86		3.58	26.711
58	205.56	14.409	11.908	87		5.48	23.167
59	227.04	14.425	11.921	88		8.67	18.345
60	231.96	14.228	11.759	89		9.35	17.927
61	251.85	14.080	11.636	90		11.84	16.317
62	258.87	14.223	11.755	91		11.99	16.413
63	271.65	13.981	11.555	92		13.98	15.446
64	298.05	13.872	11.464	93		14.40	15.176
0 sample, hct.	plasma visc	1 shear rate (1/sec)	2 viscosity (cP)	3 relative viscosity	16.61	14.496	11.980
1 30-4	0.04	256.081	215.116	36	16.94	14.296	11.815
2 H=75	0.11	226.816	200.722	37	19.26	13.752	11.365
3 1.13 cP	0.17	171.228	151.529	38	20.71	13.261	11.042
4	0.23	138.540	122.602	39	21.25	13.262	10.960
5	0.30	121.169	107.229	40	23.23	12.829	10.602
6	0.36	113.882	100.781	41	25.30	12.860	10.628
7	0.42	101.909	90.185	42	25.21	12.570	10.388
8	0.55	90.331	79.929	43	25.85	12.504	10.334
9	0.74	76.212	67.444	44	28.52	12.021	9.935
				45	29.97	11.876	9.815
				46	41.88	10.774	8.904
				47	60.90	9.680	8.000
				48	86.34	8.940	7.388



49	93 00	8 810	7 281	24	10 63	9 306	8 963
50	111 78	8 481	7 009	25	11 89	9 112	8 596
51	126 12	8 259	6 826	26	13 17	8 943	8 437
52	143 89	8 119	6 710	27	14 33	8 617	8 125
53	152 50	7 998	6 610	28	15 71	8 505	8 024
54	169 23	7 945	6 566	29	17 19	8 192	7 728
55	172 50	7 850	6 488	30	18 90	8 063	7 607
56	192 27	7 745	6 401	31	19 82	7 884	7 438
57	201 09	7 773	6 424	32	21 44	7 789	7 348
58	212 07	7 650	6 322	33	22 49	7 614	7 183
59	226 59	7 674	6 342	34	23 98	7 578	7 149
60	231 87	7 575	6 260	35	24 47	7 464	7 042
61	245 79	7 611	6 290	36	28 43	7 152	6 747
62	251 61	7 499	6 198	37	29 70	7 253	6 842
63	271 02	7 549	6 239	38	29 97	7 149	6 744
64	278 04	7 442	6 150	39	55 14	6 069	5 725
65	299 70	7 448	6 155	40	80 52	5 518	5 206
0 sample. hct. 1 shear rate 2 viscosity 3 relative				41	92 97	5 353	5 050
plasma visc. (1/sec) (cP) viscosity				42	112 32	5 131	4 841
1 32-7	0 11	114 731	108 237	43	126 06	5 032	4 747
2 H=59	0 17	74 956	70 713	44	137 94	4 939	4 659
3 1.06 cP	0 30	55 288	52 253	45	152 52	4 864	4 569
4	0 42	46 851	44 199	46	176 22	4 755	4 486
5	0 55	40 437	38 167	47	179 01	4 738	4 470
6	0 93	28 715	27 090	48	201 66	4 675	4 410
7	0 93	32 712	30 860	49	212 07	4 628	4 366
8	1 12	27 501	25 944	50	227 22	4 617	4 356
9	1 12	30 547	28 818	51	231 90	4 576	4 317
10	1 32	28 435	26 825	52	259 11	4 559	4 301
11	1 39	25 335	23 901	53	278 10	4 487	4 233
12	1 51	26 965	25 439	54	297 84	4 459	4 207
13	1 59	24 546	23 157	55	299 70	4 488	4 234
14	1 76	25 336	23 902	0 sample. hct. 1 shear rate 2 viscosity 3 relative			
15	1 85	24 960	23 547	plasma visc. (1/sec) (cP) viscosity			
16	1 95	24 412	23 030	1 32-4	0 80	12 805	12 080
17	2 05	23 917	22 563	2 H=44	0 93	12 589	11 876
18	2 14	23 706	22 364	3 1.06 cP	1 12	11 896	11 223
19	2 25	23 249	21 933	4	1 33	11 599	10 942
20	2 33	22 860	21 566	5	1 37	11 342	10 700
21	2 45	22 568	21 291	6	1 63	10 945	10 325
22	2 84	21 117	19 922	7	1 73	10 812	10 200
23	4 19	17 022	16 058	8	2 01	10 412	9 823
24	6 10	14 732	13 899	9	2 19	10 273	9 692
25	8 66	12 815	12 090	10	2 26	10 208	9 630
26	9 30	12 771	12 048	11	2 52	10 049	9 480
27	11 84	11 518	10 866	12	2 52	9 999	9 433
28	12 62	11 467	10 818	13	2 64	9 993	9 427
29	14 38	10 761	10 152	14	2 78	9 976	9 411
30	15 92	10 548	9 951	15	3 00	10 028	9 460
31	16 94	10 188	9 611	16	4 97	8 731	8 237
32	18 57	9 971	9 407	17	8 15	7 739	7 301
33	19 47	9 740	9 125	18	11 89	6 928	6 536
34	21 88	9 413	8 880	19	12 60	7 029	6 631
35	22 02	9 409	8 876	20	16 50	6 464	6 098
36	23 92	9 171	8 552	21	17 06	6 349	6 178
37	25 82	8 966	8 458	22	20 47	6 244	5 891
38	27 82	8 724	8 230	23	20 88	6 298	5 942
39	29 81	8 511	8 029	24	24 45	6 004	5 664
40	30 12	8 506	8 025	25	24 70	6 062	5 719
41	48 16	6 721	6 341	26	28 41	5 809	5 480
42	53 27	6 230	5 877	27	29 97	5 801	5 473
43	99 99	4 184	3 834	28	36 15	5 693	5 371
44	126 18	3 877	3 544	29	67 98	4 949	4 649
45	144 63	3 829	3 499	30	92 76	4 607	4 346
46	152 58	3 710	3 387	31	106 14	4 532	4 275
47	176 34	3 686	3 364	32	125 94	4 374	4 124
48	179 01	3 592	3 273	33	137 85	4 331	4 086
49	205 50	3 509	3 197	34	158 94	4 213	3 975
50	208 23	3 584	3 268	35	182 52	4 158	3 923
51	231 90	3 450	3 142	36	196 72	4 077	3 846
52	239 94	3 519	3 207	37	220 62	4 058	3 828
53	251 67	3 419	3 112	38	238 32	3 980	3 755
54	291 30	3 368	3 064	39	258 84	3 989	3 763
0 sample. hct. 1 shear rate 2 viscosity 3 relative				40	277 45	3 914	3 692
plasma visc. (1/sec) (cP) viscosity				41	299 73	3 912	3 691
1 32-5	0 49	22 222	20 964	0 sample. hct. 1 shear rate 2 viscosity 3 relative			
2 H=50	0 55	21 281	20 076	plasma visc. (1/sec) (cP) viscosity			
3 1.06 cP	0 68	20 432	19 275	1 32-1	1 99	3 074	4 787
4	0 87	19 290	18 198	2 H=34	2 15	3 125	4 844
5	0 99	18 250	17 217	3 1.06 cP	2 10	4 853	4 578
6	1 13	17 920	16 906	4	2 40	3 079	4 792
7	1 26	17 159	16 188	5	2 45	3 074	4 787
8	1 32	17 058	16 092	6	2 59	4 879	4 603
9	1 44	17 016	16 053	7	2 78	4 700	4 434
10	1 46	16 581	15 642	8	3 00	4 624	4 362
11	1 72	15 854	14 958	9	4 32	3 720	3 509
12	1 76	16 293	15 373	10	6 24	3 709	3 499
13	1 92	15 438	14 564	11	8 78	3 638	3 432
14	2 14	15 379	14 508	12	9 94	3 702	3 492
15	2 18	14 874	14 022	13	11 96	3 786	3 572
16	2 45	14 594	13 768	14	13 90	3 629	3 424
17	2 53	14 708	13 875	15	16 42	3 791	3 576
18	2 84	14 109	13 310	16	16 55	3 646	3 440
19	3 00	14 303	13 493	17	19 17	3 549	3 367
20	3 62	13 292	12 340	18	20 88	3 557	3 356
21	4 17	11 156	9 718	19	21 81	3 524	3 325
22	8 07	10 301	9 198	20	24 46	3 430	3 236
23	9 25	9 750	9 198	21	24 70	3 554	3 333

22	28.43	3.426	3.232
23	29.97	3.487	3.290
24	30.03	3.801	3.586
25	55.41	3.433	3.239
26	73.11	3.131	2.934
27	87.18	3.212	3.030
28	106.05	3.041	2.869
29	119.04	3.097	2.922
30	139.20	2.988	2.819
31	157.20	3.024	2.853
32	165.48	2.955	2.788
33	192.06	2.924	2.758
34	195.39	2.976	2.808
35	218.43	2.903	2.739
36	227.28	2.944	2.777
37	238.26	2.892	2.728
38	259.11	2.917	2.752
39	277.80	2.867	2.705
40	299.76	2.876	2.713

0 sample, hct. 1 shear rate 2 viscosity 3 relative  
plasma visc. (1/sec) (cP) viscosity

1 32-3	3.00	3.913	3.314
2 H=24	5.56	2.218	2.092
3 1.06 cP	10.01	2.245	2.118
4	11.91	2.162	2.040
5	14.46	2.234	2.108
6	14.55	2.068	1.951
7	18.58	2.235	2.108
8	18.51	2.159	2.037
9	21.47	2.227	2.101
10	22.47	2.154	2.013
11	24.02	2.219	2.093
12	26.50	2.195	2.071
13	29.07	2.143	2.022
14	36.78	2.334	2.202
15	68.64	2.228	2.102
16	92.61	2.082	1.964
17	106.83	2.175	2.052
18	125.67	2.103	1.984
19	144.96	2.169	2.046
20	158.91	2.107	1.988
21	191.94	2.111	1.992
22	195.84	2.154	2.032
23	224.91	2.113	1.993
24	240.27	2.156	2.034
25	251.37	2.112	1.992
26	265.83	2.157	2.035
27	297.75	2.119	1.999

0 sample, hct. 1 shear rate 2 viscosity 3 relative  
plasma visc. (1/sec) (cP) viscosity

1 32-3	8.72	1.170	1.104
2 H=11	11.27	1.289	1.216
3 1.06 cP	12.61	1.359	1.262
4	14.45	1.287	1.214
5	15.91	1.328	1.253
6	17.63	1.329	1.254
7	19.20	1.308	1.234
8	20.82	1.340	1.264
9	22.50	1.367	1.290
10	23.36	1.362	1.285
11	25.90	1.387	1.308
12	27.76	1.366	1.289
13	29.98	1.413	1.333
14	36.78	1.431	1.350
15	68.67	1.429	1.348
16	92.70	1.311	1.237
17	106.83	1.416	1.336
18	125.61	1.337	1.251
19	145.02	1.422	1.342
20	158.67	1.356	1.272
21	189.57	1.424	1.340
22	191.82	1.371	1.293
23	224.73	1.387	1.308
24	227.67	1.437	1.356
25	251.25	1.401	1.322
26	259.56	1.444	1.362
27	297.42	1.422	1.342

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