# Design and Control of Photoflash Capacitor Charging Circuits 

by
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Submitted to the Department of Electrical Engineering and Computer Science
in partial fulfillment of the requirements for the degree of Masters of Engineering in Electrical Engineering and Computer Science at the

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#### Abstract

This thesis develops an optimal strategy for charging photoflash capacitors. Photoflash capacitors need to be charged to voltages as high as 350 V in low-voltage batterypowered portable devices. With the decreasing size of digital cameras, existing solutions are too large. This thesis will study the operation and losses of a flyback capacitor charger. Specifically, the thesis will focus on minimizing the solution size, given an input current, in addition to keeping efficiency acceptable.


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## Chapter 1

## Introduction

### 1.1 Background

This thesis develops an optimal strategy for charging photoflash capacitors. Two ICs developed by Albert Wu at Linear Technology, LT3420 [2]and LT3468 [1], inspired the ideas presented in this thesis. These ICs implement two different charging strategies, both focusing on shrinking the solution size while improving the efficiency over previous charging methods. Only one other significant research paper, by Sokal, has been written on charging capacitors. In [5], Sokal comes to a conclusion on the fastest and most efficient method to charge a capacitor given a maximum peak switch current. The techniques developed in this thesis are most applicable to charging photoflash capacitors in digital cameras.

Before the wide spread use of electronics, cameras used individual flash bulbs or flash bars to produce a 40 ms pulse of intense white light from a chemical reaction. About 40 years ago, professional photographers started to use electronic flashes with a much shorter 1 ms pulse of white light, generated using a Xenon bulb. Electronic flashes were not used extensively until the last 10 years when all but the cheapest cameras utilize them as standard equipment. With improvements in technology in the last ten years, cameras have decreased considerably in size. The smallest digital camera is the size of a $\frac{1}{4}$ inch thick credit card. Cell phones now feature built-in digital cameras, which may someday incorporate an integrated Xenon flash. With
space as a premium in camera cases, the size of the flash capacitor charger is critical to its utility.

Currently, popular methods to charge a photoflash capacitor include the selfoscillating forward converter and the micro-controlled flyback converter. The selfoscillating forward converter is comparatively the most cost effective, since it requires only a few discrete transistors. However, it is inefficient at low voltages and the custom multi-winding transformers are too bulky for the increasingly feature-packed digital and film cameras. This type of charger is only found in disposable and inexpensive film cameras. Most other cameras produce a flash with a microprocessor-controlled flyback converter. It controls the gate of a power switch with a pre-programmed set of switch on and off times with only the ability to detect the primary current and the output voltage. Since the controller cannot sense secondary current, the controller has to store an algorithm to calculate an appropriate switch off-time. If the secondary current decays to zero before the end of the off-time, the circuit will remain idle which increases the peak-current requirements, thus it enlarges the transformer and the switch.

Clearly, a more effective method of charging capacitors exists. This thesis will explore the variable frequency control methods explained in [5], but apply them to the rapidly expanding market of digital cameras.

### 1.2 Overview

To generate a flash, a Xenon flash bulb requires a special capacitor charged to a high voltage. This thesis will study methods to charge this photoflash capacitor from a low input voltage with a power limit. The required capacitance is determined by the size of the flashbulb. [7] Without variations in efficiency, the charge time is set by the final output voltage, the output capacitance, and the input power. As a result of this dependency, this makes solution size and efficiency the most important features of a capacitor charger. This thesis will focus on analyzing flyback capacitor charger performance normalized to a given input power specification.

A flyback converter consists of a transformer, switch, diode, output capacitor and control circuitry. Figure 2-1 shows the configuration of a flyback converter. The transformer serves as the energy storage device. With the switch on, the transformer's magnetizing inductance magnetizes from the input power source. The input voltage determines the rate of magnetization. When the core of the transformer becomes close to saturation, the switch turns off, forcing the transfer of current from the transformer's primary winding to the secondary winding. The current from the secondary winding flows through the diode to the output capacitor. Depending on the charging strategy, the switch might turn on to terminate the current to the output capacitor before the secondary current falls to zero. The secondary winding has more turns than the the primary winding to limit the voltage the switch has to withstand.


Figure 1-1: Generic flyback converter.

In a flyback converter with ideal components, the amount of magnetizing inductance would not affect the charge time of the charger, but would only determine the operating frequency. With stray capacitances on the switch node and core losses, efficiency decreases with increased operating frequency. Therefore, a magnetizing inductance should be determined to keep the operating frequency below a maximum operating frequency. This maximum operating frequency would be dependent on the type of core material used and the amount of stray capacitance on the switch node. Although losses from the core and the stray capacitance decrease with larger magnetizing inductance, more turns are needed for both the primary and secondary
windings, thereby increasing the winding resistance. An accurate model of these losses is needed to determine the optimum amount of magnetizing inductance.

All the loss terms for a flyback charger may be easily derived analytically as a function of the output voltage. These equations could be added together analytically, but would result in a large, un-intuitive equation. Instead, MATLAB is used to plot, sum and integrate these equations numerically. MATLAB is also capable of converting a power loss in terms of $V_{\text {out }}$ to a total charge cycle efficiency. This thesis will rely on MATLAB to plot total efficiency versus parameters such as magnetizing inductance. The calculations done in MATLAB will focus the experimentation and be correlated with actual data afterwards.

For the experiments, a flyback controller was built with adjustable primary and secondary current limits. The primary and secondary currents are measured with sense resistors and op amps. With control over both current limits, the controller is capable of keeping the maximum input current constant with all the charging strategies. The controller is also capable of turning the switch on by monitoring the switch node voltage instead of the secondary current. The flyback capacitor charger is flexible enough to use a wide range of transformers. These transformers have different magnetizing inductances, turns ratios, winding window allocations, and core gap lengths. A TDK EPC10 core is used for all of the experiments. [6]

### 1.3 Organization

In Chapter 2, the thesis explains the operation of a flyback capacitor charger and the benefits of variable-frequency operation. The components are also discussed briefly. In Chapter 3, the flyback chargers losses are modeled analytically, along with the charge time. Chapter 4 outlines the techniques used in MATLAB to compute losses with the analytical models. From there, Chapter 5 describes the construction and testing of a flyback capacitor charger. Chapter 6 analyzes the transformer in detail. Chapter 7 compares the experimental results with the simulations and also suggests optimal values for components. Chapter 8 is an overview of the components used to
create a flash in a digital camera. Finally, conclusions and suggestions for further work are discussed in Chapter 9.

## Chapter 2

## Operation of a Flyback Converter

### 2.1 Theory

A flyback converter, as shown in Figure 2-1, consists of a transformer, a power transistor, a diode, and an output capacitor. The following description of a flyback converter is valid for one that regulates or charges. The switch turns on to allow the current in the magnetizing inductance of the transformer to reach a peak value, $I_{\text {lim }}$, as shown in Figure 2-2. The slope of the current in the charging pulse is constant over the charging cycle. When the switch turns off, the magnetizing inductance delivers current to the output capacitor through the secondary winding; this time period is known as the flyback period. The peak secondary current is N times smaller than the primary current, as shown in Figure 2-3. As the output voltage increases, the secondary current decreases faster.

$$
\begin{equation*}
P_{s w}=\frac{1}{T} \int_{0}^{T} R_{s w} I_{s w}^{2} d t=d I_{l i m}^{2} R_{s w}\left[\alpha+\frac{1}{3}(1-\alpha)^{2}\right] \tag{2.1}
\end{equation*}
$$

Most regulating flyback converters operate in a constant-frequency control mode. With a constant-frequency, the steady-state duty cycle is determined solely by the input voltage, output voltage, and the turns ratio. However, with a light load, the converter enters discontinuous mode and the duty cycle relationship is no longer valid. Discontinuous mode occurs when the magnetizing current falls to zero before


Figure 2-1: Flyback converter.


Figure 2-2: Primary current waveform.


Figure 2-3: Secondary current waveform.
the switch turns on again. In discontinuous mode, the duty cycle controls the average current to the output capacitor. In lieu of duty cycle control, many regulators control duty cycle implicitly by controlling the peak current in the primary winding which allows converters to operate in either continuous or discontinuous mode. Converters use a sense resistor between the emitter of the switch and ground to sense the peak current. The peak current limit is adjusted by sensing if the output voltage is above or below the set output voltage.

Constant-frequency control works efficiently with a constricted output voltage range, but in the charging of a capacitor, the output voltage ranges from 0 volts to the final output voltage which could be as high as 500 volts. The voltage across the secondary winding varies drastically, resulting in off-times that vary by a 500:1 ratio. At low voltages, the duty cycle will become very small and will approach the minimum on-time of the controller. Once the minimum on-time is reached the part can no longer return the magnetizing current to the level at the start of the switch cycle. The magnetizing current will increase with every switching cycle. To limit current, the regulator will need to be capable of skipping cycles to let the secondary current fall below the current limit which will in turn reduce the switching frequency. At high output voltages, the secondary current falls fast compared to switch on-time. As a result, the secondary current falls to zero before the end of the switching period, leaving the circuit in an idle state, which leads to higher peak currents for a given input power. At both low and high output voltages, undesirable operation occurs when implementing constant-frequency control for charging capacitors.

To operate more efficiently in capacitor charging, the flyback converter should operate with a variable frequency. Without a set switching frequency, the circuit determines when to end the flyback period. As in the constant-frequency case, the switch turns off once the primary winding current reaches a current limit. One method to determine when to terminate the flyback period involves sensing the secondary winding current. The switch is turned back on once the current falls to a fraction of the current limit. This technique is shown in Figure 2-4 where $\alpha$ is the ratio of the secondary current to the primary current. In [5], Sokal and Redl discuss flyback
charging circuits. They conclude that an $\alpha$ close to unity, producing flat current pulses to the output capacitor, minimizes peak and RMS currents, thus reducing losses associated with parasitic resistances and current-carrying requirements of the switch, transformer and the diode. In contrast with their findings, the Linear Technology converter LT3468 switches when the secondary winding current falls to zero [1]. This charging method may use a smaller inductor and reduces the losses due to parasitic capacitances of the transformer on the collector of the switch.


Figure 2-4: Magnetizing inductor current.

While charging, the flyback capacitor charger needs to be able to sense when the output voltage reaches the desired value. A resistive voltage divider connected to the output is commonly used in regulators. With a finite resistance voltage sense amplifier connected to the output of the voltage divider, the resisters cannot be made arbitrarily large, therefore a substantial current can flow through the resistors when the output is near its final value. This loss is unacceptable in battery operated devices. Not only does it lower the efficiency of the flyback capacitor charger, the capacitor loses its charge from the end of the charging period till the user presses the flash button. Linear Technology has patented a method to avoid this problem by sensing the voltage on the primary winding during the flyback period [3]. When the switch is off, the diode is conducting and the output voltage is across the secondary winding. The switch node sees the input voltage plus the output voltage divided by the turns ratio. By subtracting the input voltage with a circuit, the output voltage is available to the control circuitry without power dissipation from the output voltage.

At high voltages, the flyback period, or off-time, becomes very short. $t_{o f f}=$ $\frac{L_{\text {sec }} I_{\text {lim }}(1-\alpha)}{V_{\text {out }} N}$. For a comparator to sense this voltage during the flyback period, there is a minimum off-time based on the speed of the comparator. For the output voltage
sense to work correctly, the inductance of the secondary winding has to satisfy the following relationship: $L_{\text {sec }}>\frac{t_{\text {off }} V_{\text {final }} N}{I_{\text {lim }}(1-\alpha)}$. Without considering efficiency, this inequality limits the minimum size of the transformer.

### 2.2 Transformer

The transformer is often the most complicated component in a flyback converter, and often accounts for the majority of losses. In a flyback transformer, the magnetizing inductance acts as the main energy storage device. The transformer acts as a coupled inductor, since current never flows through both windings simultaneously, thus never obeying the current relationship of an ideal transformer. The turns ratio of the transformer serves two main purposes: to protect the power switch from the high output voltage, and to decrease the rate of decay of the magnetizing current. The turns ratio should be kept to a minimum to reduce the amount of winding area used by the secondary winding.

As the main energy storage device, the magnetizing inductance value affects the operating frequency of the flyback converter. By increasing the magnetizing inductance, the switching frequency decreases linearly. The lower frequency reduces frequency-dependent losses. By increasing magnetizing inductance, more turns are needed around the core in both the primary and secondary windings. However, the windings still need to fit in the same winding window. This leads to the need for longer wires while decreasing the winding wire's width, consequently increasing the DC winding resistance and the associated losses.

### 2.3 Power Switch

In the test circuit, a 2A MOSFET is used to control the primary current. The MOSFET is subjected to DC drain-source voltage equal to the output voltage divided by the turns ratio. The leakage inductance also creates a high voltage on the drain of the MOSFET. When the switch turns off, the leakage inductance continues to source
current into the drain of the MOSFET. The energy in the inductance charges the capacitance of the switch causing a voltage spike. The voltage spike becomes larger with more leakage inductance, but remains constant throughout the charging cycle. This voltage spike could reach as high as $I_{\text {lim }} \sqrt{\frac{L_{\text {leak }}}{C_{p}}}$. The capacitance, $C_{p}$, comes from the switch's capacitance, and the primary winding's capacitance. The switch needs to be capable of withstanding this voltage spike.

### 2.4 Diode

The diode blocks current from flowing from the output capacitor back into the transformer. The diode serves as the second switch in the topology. The secondary current turns the switch on after the MOSFET turns off. When the switch is turned back on, the diode blocks current from flowing into the transformer. To block this current, the diode withstands a reverse voltage of $V_{o u t}+N V_{i n}$. The most important property of the diode in this application is its DC reverse breakdown voltage. The parasitic capacitance adds to the problem of reverse breakdown voltage. The parasitic capacitance on the secondary winding is charged to the output voltage. At this point, the capacitance is in parallel with the secondary winding's leakage inductance. With $V_{i n}$ across the primary, the parasitic capacitance sees $-N V_{\text {in }}$ on the other side of the leakage inductance, as shown in Figure 2-5. This produces a damped second-order response on the secondary winding with an amplitude of $V_{\text {out }}+N V_{\text {in }}$ with a steady state voltage of $-N V_{i n}$. With the damping, the voltage does not swing down completely to the negative amplitude, but does increase the requirement of the dynamic blocking voltage of the diode substantially.

### 2.5 Boundary Mode Operation

Boundary mode operation constitutes a major difference from continuous conduction mode, and the following section will detail these differences. Continuous conduction mode (CCM) indicates that the inductor current or magnetizing current of the trans-
former is always positive. In contrast, discontinuous conduction mode (DCM) is when the current in the inductance falls to zero. Furthermore, with both the switch and diode off, the switch voltage rings. The energy from the parasitic capacitance of the switch, transformer, and the diode transfers to the inductance, and forms a parallel resonance tank. At low output current levels, most fixed-frequency converters enter DCM. In a variable frequency power converter, as the one described above, it does not make sense for the circuit to remain idle in DCM, since it is capable of turning the switch on at anytime, unless a reduction in input current is wanted. If the switch has a fixed current limit, this idle time would lower the output power capabilities of the switch.

With a variable-frequency converter, there is the option of allowing the parasitic capacitance to ring to zero before turning the switch on opposed to turning the switch on immediately after the current reaches zero. This mode of operation is called boundary mode or edge of DCM. Boundary mode brings higher efficiency by recycling the energy from the parasitic capacitance instead of dissipating the energy in the switch resistance, and is also known as zero-voltage switching. With high Q capacitors and inductors, all the energy from the capacitance is recovered. In actuality, a fraction of the energy is dissipated in parasitic resistances. Since this capacitance loss is the dominant loss at higher output voltages, boundary mode could possibly result in significant improvements in efficiency over a converter in CCM.


Figure 2-5: Second order network when switch turns on.

In addition, the diode is turned off when the current through it is zero, known as zero-current switching. Zero-current switching does not improve the efficiency at all since the reverse recovery loss is not a significant factor in the efficiency. Boundary mode decreases the power output of the converter in a slightly different way than a converter in DCM. The ring of the capacitance does not take much time compared to the operating frequency of the converter. However, the current in the magnetizing inductance becomes negative when storing the energy from the parasitic capacitance. When the switch turns on, the current in the magnetizing inductance takes a fraction of the on-time to reverse the negative current in the magnetizing inductance.

### 2.6 Linear Technology Flyback Capacitor Chargers

The LT3468 operates in boundary mode operation. In contrast, the LT3420 is a continuous mode controller. The LT3420 was the first part to be released as a capacitor charger for photoflash applications. The LT3420 miniaturized the components traditionally needed in a photoflash capacitor charger, but also suffered from some unexpected problems. The part operates by sensing both the primary and secondary currents and switches when those currents reach their limits. The LT3420 enjoyed fast charge times with a low peak switch current. Although the LT3420 benefited from its continuous operation, the LT3420 had large losses due to the parasitic capacitance of the transformer, and also required a large magnetizing inductance to keep the operating frequency low. The LT3468 was designed to solve the problems that plagued the LT3420. The LT3468 improves upon the previous design with three major improvements. Instead of sensing the secondary current, the part switches on when the switch pin rings down to the input voltage. The current change every cycle is much larger than the LT3420, thus resulting in either a reduced switching frequency or the freedom to lower the magnetizing inductance. The LT3468 takes advantage of the power savings of boundary mode operation. More information is available about
these parts in [2], and [1].

## Chapter 3

## Modeling a Flyback

### 3.1 Introduction

To better understand the tradeoffs with components in a flyback converter, the losses need to be accurately modelled. There are four forms of power loss in a flyback converter: switch loss, transformer loss, parasitic capacitor loss, and diode loss. While most of the losses can be modelled as an energy loss per cycle or a power loss, the manufacturer core loss data is given as a power loss, so to maintain consistency, power loss is used throughout. Unlike most power converters, a flyback capacitor charger is never in steady state. The power in and out of the circuit varies with output voltage, as well as the power loss terms calculated in the following sections. The most efficient method to understand the losses below is to graph them over $V_{\text {out }}$ with MATLAB. While this method produces graphs that are easily correlated with data collected in lab, the graph is misleading since the flyback charger spends more time at higher voltages. To more accurately model the capacitor charger, an equation is derived to give the amount of time spent per $\Delta V$, or $\frac{d t}{d v}$. By multiplying this quantity by power loss, the energy lost per $\Delta V$, or $\frac{d E}{d v}$ is calculated. By integrating this equation over $V \max$, the total energy lost per charge cycle is used to compare a capacitor charger while different parameters such as the turns ratio, or the magnetizing inductance are varied. Also in this chapter, the charge time will also be modeled.

### 3.2 Losses associated with the switch

### 3.2.1 Switch Resistance Losses

In the test circuit describe in the thesis, the switch is a MOSFET. In contrast, the parts made by Linear Technology use an integrated bipolar junction transistor. These two transistors can be modelled as an ideal switch with series resistance. Using a resistance, instead of modeling it with a $V_{c e}$ saturation voltage, more accurately reflects the switch plus simplifies calculations since its in series with the primary winding resistance.

$$
\begin{equation*}
P_{s w}=\frac{1}{T} \int_{0}^{T} R_{s w} I_{s w}^{2} d t=d I_{l i m}^{2} R_{s w}\left[\alpha+\frac{1}{3}(1-\alpha)^{2}\right] \tag{3.1}
\end{equation*}
$$

The loss from the switch resistance is calculated as the time average of the equation $P=I^{2} R$, or the $I_{r m s}^{2} R$. With this equation and the current waveform in Figure 3-2, the power loss in the switch is calculated. As $\alpha$ approaches one, the circuit loses three times the amount of power in the switch with only twice the amount of power in, or equivalently a decrease in charge time by half without considering the loss in efficiency.


Figure 3-1: Magnetizing inductor current.


Figure 3-2: Primary current waveform.

### 3.2.2 Losses due to Rise and Fall Time of Switch

With non-zero rise and fall times, the switch dissipates energy as current and voltage exist at the same time. Figure 3-3 shows a simple model of the switch turning off. As the switch turns off, the switch voltage rises linearly to $\frac{V_{\text {out }}}{N}$ before the current falls linearly to zero from its initial value of $I_{\text {lim }}$. The switch turn on is the opposite process with the current rising linearly before the voltage falls linearly, as shown in Figure 3-4. The rise and fall time energy loss is the area of the multiplication of the current waveform and the voltage waveform. By multiplying the energy loss by frequency, the power loss is given by

$$
\begin{align*}
P_{f} & =\left(\frac{V o u t}{N}\right)\left(I_{\text {lim }}\right) t_{f} \cdot f  \tag{3.2}\\
P_{r} & =\left(\frac{V o u t}{N}\right)\left(\alpha I_{l i m}\right) t_{r} \cdot f \tag{3.3}
\end{align*}
$$



Figure 3-3: Switch turn off waveform.


Figure 3-4: Switch turn on waveform.

### 3.3 Losses from Transformer

The transformer contributes a majority of the losses in the flyback converter. The thin copper wire used for the windings has significant resistance. The loss from the winding is known as the DC winding resistance loss. At higher frequencies, the windings may suffer additional losses from proximity and skin effect. These two losses will not be modelled because they are highly dependent on the winding method, which cannot be closely controlled in my thesis, and also they do not contribute a significant loss compared to other loss terms. Losses in the core encompasses another fraction of the energy loss in the transformer. The copper losses and the core losses translate into heat lost inside the transformer, resulting in a considerable increase in the transformer's temperature and causing it to be the only component to become noticeably hot.

### 3.3.1 Loss from Leakage Inductance

The core is responsible for transferring flux between the windings on the transformer. Even though the permeability of the core is much higher than air, some flux still leaks into the air, thus not coupling into the secondary. This leads to additional inductance in series with the windings and the magnetizing inductance. Leakage inductance is the name given to this parasitic inductance. A method of measuring the leakage inductance is presented in Chapter 6. The primary leakage inductance causes a voltage spike when the switch turns off. The leakage inductance forms a second-order circuit with the capacitance on the switch node. This transient might exceed the maximum allowable voltage the switch can withstand. In most flyback converters, a snubber network clamps the voltage on the switch node. A snubber dissipates an energy greater than the amount stored in the leakage inductance per switch cycle. Because space is limited in a flyback capacitor charger, the switch is designed to handle the voltage transient caused by the leakage inductance. With no snubber, the energy in the leakage inductance rings briefly, but most of the energy is eventually transferred to the output. On the secondary side, the leakage inductance
is not a problem because it discharges through the diode to the output capacitor. The power loss from the leakage inductance is given by

$$
\begin{equation*}
P_{l e a k}=\frac{1}{2} L_{l e a k} I_{l i m}^{2} \chi f \tag{3.4}
\end{equation*}
$$

$\chi$ is a factor much less than one. Leakage inductance was not seen experimentally to make a difference in efficiency, but caused substantial ringing in the secondary winding current.

### 3.3.2 Loss from DC Winding Resistance



Figure 3-5: Secondary current waveform.

DC resistance is the simplest loss to understand in a transformer. The finite conductivity of copper results in a parasitic resistance in each of the windings. The resistance is given by $R=\frac{l_{t} n}{A \rho}$, where $l_{t}$ is the average length per winding, $n$ is the number of windings, $A$ is the cross-sectional area of the wire, and $\rho$ is the conductivity of copper. The power loss is given by $P=I^{2} R$, where $I$ is shown in Figure 3-2 for the primary winding and Figure $3-5$ for the secondary winding. The power loss equations reduce to the following:

$$
\begin{gather*}
P_{d c p}=d I_{l i m}^{2} R_{p}\left[\alpha+\frac{1}{3}(1-\alpha)^{2}\right]  \tag{3.5}\\
P_{d c s}=(1-d) I_{l i m}^{2} \frac{R s}{N^{2}}\left[\alpha+\frac{1}{3}(1-\alpha)^{2}\right] \tag{3.6}
\end{gather*}
$$

### 3.3.3 Core Loss

Core loss consists of two remagnetization losses: hysteresis loop loss and eddy current loss. In most textbooks, these losses are considered separate, but in reality they cannot be separated. In [9], the authors explain the origin of a combined remagnetization loss. Manufacturer's publish the core loss with a sinusoidal waveform. In a flyback converter, the excitation waveform is a square wave. The paper introduces a simple way to modify the Steinmetz equation to use non-sinusoidal waveforms.

The first step in using the Steinmetz equation is to calculate the ac peak flux density In the manufacturer's data, power loss density is plotted against peak ac flux density with sinusoidal excitation at different frequencies. To find peak ac flux density, the change in current per cycle needs to be found with the following:

$$
\begin{equation*}
\Delta I=\frac{1}{2}(1-\alpha) I_{l i m} . \tag{3.7}
\end{equation*}
$$

After the change in current is found, the peak ac flux density is found by the following equation.

$$
\begin{equation*}
\Delta B=\frac{\Delta I A_{l} n}{A_{e}}=\frac{\Delta I L}{n A_{e}} \tag{3.8}
\end{equation*}
$$

Where $A_{l}$ is nF per turns squared of the $\operatorname{core}\left(A_{l} \approx \frac{l_{g}}{\mu_{0} A_{c}}\right), n$ is the number of turns for the primary winding, and $A_{e}$ is the effective cross-sectional area of the core.

The core power loss is approximated by the Steinmetz equation. By using the published data, $K_{f e 0}, \alpha$, and $\xi$ are determined by fitting the following equation to the manufacturer's plot of core loss data.

$$
\begin{equation*}
P_{f e}=K_{f e 0}(\Delta B)^{\beta} f_{e q}^{\xi} V_{e} \tag{3.9}
\end{equation*}
$$

The frequency used in the above equation is not the switching frequency of the flyback charger, but a modified frequency from [9] or [10]. In a capacitor charger, the modified frequency takes the following form.

$$
\begin{equation*}
f_{e q}=\frac{2 f}{\pi^{2} d(1-d)} \tag{3.10}
\end{equation*}
$$

### 3.3.4 Transformer's Parasitic Capacitance Loss

While not directly a loss in the transformer, the transformer has a significant amount of capacitance between the windings and between the opposing ends of the primary and secondary windings. In continuous mode, this capacitance energy is dissipated across the switch when it turns on during every switching cycle. In boundary mode, the energy is transferred to the magnetizing inductance of the transformer, but during this transfer a portion of the energy is lost. The only way to determine the amount of energy in this capacitance is by observing a flyback capacitor charger in operation. In discontinuous mode, the capacitance forms a second-order network with the magnetizing inductance and rings. By measuring the frequency and the magnetizing inductance, the total capacitance on the switch pin can be calculated. This total capacitance not only accounts for all the parasitic capacitance in the transformer, but also the diode's capacitance and the switch's capacitance. The formula to calculate the total parasitic capacitance is shown below along with a scope photo of the fall-time, Figure 3-6.

$$
\begin{equation*}
C_{p a r a}=\frac{\left(4 t_{t f}\right)^{2}}{4 \pi^{2} L_{p r i}} \tag{3.11}
\end{equation*}
$$

In the equation, $t_{t f}$ is the fall-time of the flyback waveform. It is also measured in the scope photo, Figure 3-6.

### 3.4 Diode Losses

While the diode is in forward conduction, the power loss is approximately the forward voltage drop times the current. In the case of a flyback capacitor charger, the current through the diode cannot be approximated as constant. The power equation needs to be integrated over a switching cycle and divided by the time period of a switching cycle. This results in the following equation.

$$
\begin{equation*}
P_{\text {diode }}=V_{f} I_{\text {lim }}(1+\alpha) \frac{V_{\text {in }}+V_{\text {out }}}{2 V_{\text {out }}\left(V_{\text {out }}+N V_{\text {in }}\right)} \tag{3.12}
\end{equation*}
$$

The forward diode drop does not contribute a significant loss above 25 V .
Another loss occurs in the diode when it turns off. The diode stores a small amount of charge when conducting forward current. The diode conducts current in the opposite direction to remove this charge. The amount of time it takes is called the reverse recovery time. Modern diodes that only conduct small amounts of current typically have very fast reverse recovery times. The reverse recovery current is proportional to the forward current of the diode at turn off. In the diode used in the test circuit, a Vishay GSD2004S, the reverse recovery time $\left(t_{r r}\right)$ is 50 nS and the reverse recovery current is 3 mA with a 30 mA forward current prior to the turn off. By using a very conservative estimation using the following equation to calculate the


Figure 3-6: Scope shot: Ch4 is primary switch pin.


Figure 3-7: Diode reverse recovery current.
power loss, $V_{\text {out }} I_{F} t_{r r} f$, the reverse recovery loss is not significant compared to the other losses and will not be modeled.

### 3.5 Charge time

There are many different approaches to calculate charge time. To start with the simplest method, the input current over the charge cycle can be approximated as constant. This is a fairly accurate representation in the test circuit above 100V. With this one assumption, the charge time can be found with the following equation.

$$
\begin{equation*}
t_{\text {charge }}=\frac{C_{\text {load }} V_{\text {out }}^{2}}{I_{\text {in }} V_{\text {in }} \mu} \tag{3.13}
\end{equation*}
$$

In the equation above, $\mu$ is the total efficiency of the circuit. This model of the charge time is relatively simple and is not that useful, except to understand on a first order how parameters influence charge time.

A more complete model is derived by integrating $\frac{\Delta t}{\Delta v}$ over the charging voltage range. Instructions on how to calculate $\frac{\Delta t}{\Delta v}$ are in Chapter 4, Modeling in MATLAB. This integration results in the following equation.

$$
\begin{equation*}
t_{\text {charge }}=\int_{0}^{V_{\text {out }}} \frac{\Delta t}{\Delta v} d v=\frac{C V_{\text {out }}}{I_{\text {lim }} \mu}\left(\frac{V_{\text {out }}}{V_{\text {in }}}+2 N\right) \frac{1-\alpha}{1-\alpha^{2}} \tag{3.14}
\end{equation*}
$$

This equation shows the effects of changing $\alpha$ and the other parameters. With an $\alpha$ close to 1 , the charge time decreases by half over an $\alpha$ of 0 .

The last two methods have assumed a constant efficiency over the charge cycle. The efficiency varies by up to $10 \%$ over the charge cycle, so the previous methods would be inaccurate. While this can be done numerically with an efficiency plot, there are no benefits because charge time cannot be modeled to this accuracy because of circuit delays. There are two major delays not accounted for in the models above. The first major delay is the amount of time it takes for the primary winding current to decrease, and transfer to the secondary winding. Another delay is the amount of time it takes for the switch to turn back on. These delays will be explained in more
detail in Chapter 5.

## Chapter 4

## Modeling in MATLAB

### 4.1 Introduction

This thesis uses MATLAB to numerically calculate the losses for a flyback capacitor charger. The vector operations are used extensively, along with the analytical expressions in Chapter 3, to calculate the losses. These vectors are capable of calculating these loss equations over the range of $V_{\text {out }}$.

### 4.2 Calculating Individual Losses

The first step in developing a model to evaluate the performance of a flyback capacitor charger is to plot each of the individual loss term versus output voltage. These individual losses are shown in Figure 4-1 and Figure 4-2. Each of these individual loss terms are checked for obvious errors. A high power loss in any of these terms generates heat, which is easy to check for in lab. The two major loss terms correspond with the two components which become warm during operation, therefore assuring reasonable values for each of the individual power losses.

Each of the losses needs the correct behavior over output voltage range. There are four different types behavior over $V_{\text {out }}$ out of the nine loss terms. The primary winding resistance (Pdcp), the switch resistance (Psw), and the leakage inductance (Pleak) increase with the duty cycle of flyback capacitor charger. The duty cycle, or
the proportion of time the switch is on, increases quickly at lower voltages and stays relatively constant over 100 V . The diode loss (Pdiode), and the secondary winding loss (Pdcs) are proportional to current through the secondary winding. The average current through the secondary side of the circuit is proportional to the complement of the duty cycle, and determines the loss in these two secondary side components. The loss due to the parasitic capacitance of the transformer increases quadratically with $V_{\text {out }}$, because the energy stored in this capacitance is proportional to $V_{\text {out }}^{2}$. The rise and fall time losses from the switch are proportional to the operating frequency. The core loss is proportional to frequency to 1.72 power with the TDK core.


Figure 4-1: Breakdown of losses from a typical flyback charger (Part 1 of 2).

Subtracting the sum of all these losses from the input power calculates the output power. Efficiency simply equals $\frac{P_{\text {out }}}{P_{i n}}$; a plot of efficiency is generated, as shown in Figure 4-3. This plot shows the decrease in efficiency at higher output voltages caused mainly by the losses due to parasitic capacitances on the switch. At higher output voltages, the operating frequency increases. Consequently, the frequencydependent losses increase at higher output voltages. The parasitic capacitance loss
increases quadratically with output voltage and is the main cause of efficiency decrease at higher output voltages. The plot shown in Figure 4-3 is relatively flat because of adequate magnetizing inductance, keeping operating frequency low. Operating frequency should be kept low enough to keep the parasitic capacitance from being the dominant loss term over the DC losses in the switch and the primary winding.

To plot efficiency versus a parameter such as magnetizing inductance, we need to convert the efficiency plot into total efficiency. The efficiency curve is deceiving since the charger spends more time at higher voltages. By starting with power loss in terms of $V_{\text {out }}$, we can multiply this with $\frac{\Delta t}{\Delta V}$. The first step in calculating $\frac{d t}{d V}$ is to find the output voltage increase per switching cycle as a function of the output voltage. The amount of energy added to the output capacitor each cycle is the energy held in the magnetizing inductance. This leads to the following equation.

$$
\begin{equation*}
\frac{1}{2} L_{p} I_{l i m}^{2}\left(1-\alpha^{2}\right)=\frac{1}{2} C(V+\Delta V)^{2}-\frac{1}{2} C V^{2} \tag{4.1}
\end{equation*}
$$

By solving for $\Delta V$ and ignoring second-order terms, we arrive at the following


Figure 4-2: Breakdown of losses from a typical flyback charger (Part 2 of 2).
equation.

$$
\begin{equation*}
\Delta V=\frac{L_{p} I_{\text {lim }}^{2}}{2 C_{\text {out }} V_{\text {out }}} . \tag{4.2}
\end{equation*}
$$

$\Delta t$ is simply the reciprocal of the the cycle frequency, or $t_{o n}+t_{o f f}$. By dividing these two terms, we arrive at

$$
\begin{equation*}
\frac{\Delta t}{\Delta V}=\frac{2 C V_{\text {out }}}{I_{\text {lim }}} \frac{1-\alpha}{1-\alpha^{2}}\left[\frac{1}{V_{\text {in }}-V_{\text {sat }}}+\frac{N}{V_{\text {out }}+V_{d}}\right] \tag{4.3}
\end{equation*}
$$

After multiplying the power loss curve with (4.3), we integrate over this new curve, giving us the energy lost during a charge. An integral is impossible to do with sampled data, so the integral is approximated by summing the multiplication of the value of the efficiency by the distance between efficiency data points for all the efficiency data points. The total efficiency is given by energy out divided by the energy in. The energy out is equal to the energy stored in the capacitor, $\frac{1}{2} C V^{2}$ and the energy in is given by the energy out plus the energy lost in charging. By creating a MATLAB function with this as an output, we may plot efficiency as parameters are changed.


Figure 4-3: Efficiency curve for flyback converter with alpha=0 and $\mathrm{L}=24 \mathrm{uH}$.

Table 4.1 lists all the inputs to this function and a short description, while the code is listed in Appendix A.

| Variable Name | Description |
| :--- | :--- |
| Al | Henries per turns squared |
| Wa | Winding Window Area |
| MLT | Mean Length per Turn |
| Ve | Effective Volume of Core |
| Ae | Effective Cross-sectional Area of Core |
| Bex | $\beta$ Value in Core Power Loss Equation |
| fex | $\xi$ Value in Core Power Loss Equation |
| n | Number of Turns for Primary Winding |
| N | Turns ratio |
| Iin | Average Input Current |
| alpha | Sets Secondary Current Limit |
| Cload | Load Capacitance |
| Vin | Input Voltage |
| Vmax | Final Output Voltage |
| leakpercent | Leakage Inductance is this Fraction of Magnetizing Inductance |
| primarywinding | Fraction of Winding Window Dedicated to Primary Winding |

Table 4.1: Table of inputs to total efficiency function for flyback charger.

As an example, Figure 4-4 shows a sweep of magnetizing inductance for a typical flyback capacitor charger. Each inductance uses the same core and winding window area. As the the inductance increases, the number of turns on both the primary and secondary windings increases, so therefore the cross-section area of the wire needs to be smaller to fit within the allocated winding window. The function accounts for this new cross-sectional area by calculating the resistance per length of the wire and multiplying by the required length of the winding based on the mean length per turn information given by the core manufacturer.


Figure 4-4: Efficiency versus magnetizing inductance.

## Chapter 5

## Design, Construction and Testing of a Flyback Capacitor Charger

### 5.1 Introduction

Designing a flyback capacitor charger controller is straightforward. Without the complicated feedback loops of regulating DC/DC converters, the flyback capacitor charger is driven by three main events. The design of a flyback controller can be done by hand initially and requires no equations. The controller to be built will operate with a variable frequency, the advantages of this are discussed in a previous section, and the basic operation is as follows. The primary current needs to be monitored and once an adjustable current limit is met, the switch turns off. Then, as the secondary current feeds the output capacitor, the current decreases to the secondary current limit, and the switch turns back on. The output voltage needs to monitored to check if an adjustable final output voltage has been reached. With the two adjustable current limits, any $\alpha$ is possible while keeping the average input current constant.

### 5.2 Design

Most of the controller circuitry is digital, except the portion that determines the current in the primary and secondary windings. To determine the current on the


Figure 5-1: Flyback capacitor charger test circuit.
primary, a low-value sense resistor is placed between ground and the source of the MOSFET. A non-inverting operation amplifier configuration is used to measure the current across the sense resistor. This amplified version of the sense resistor voltage is compared with the adjustable primary current limit reference voltage with a comparator. When the current reaches the current limit, the comparator outputs high. Similarly, the voltage on the secondary winding is measured with a sense resistor between the the secondary winding and ground. The current on the secondary winding is in the opposite direction, requiring the use of an inverting operation amplifier configuration. A comparator compares the output of the op amp with the secondary current limit voltage, so that the output goes high when the secondary current is less than the current limit.

After the primary and secondary currents are in digital form and are ready to be interfaced to the digital portion of the circuit. The digital portion of the circuit consists of one-shots, S-R latches, AND gates, and OR gates. The whole circuit, in Figure $5-1$, is relatively simple in its operation with one exception. Once the controller is started with a rising edge on the net labelled "ENABLE", the switch turns on and the primary current in the transformer ramps up. The primary current will eventually trigger the primary current limit comparator and reset a latch. The output of the latch will then force the switch off. The comparator is connected to the latch through an AND gate, which has the other input connected to an inverted one-shot that triggers when the switch turns on. A current spike occurs after the switch turns on caused by the stray capacitance on the switch node. The one-shot disables the primary current comparator to turn the switch off. When the switch turns off, the energy stored in the core releases into the output capacitor. The secondary current declines to the secondary current limit and the comparator goes high, and this positive edge on the comparator signal triggers a one-shot. The one-shot turns on the latch that determines the state of the switch. At the beginning of the charge, a latch is set to tell the circuit to charge. When the final output voltage is reached, this latch turns off. This latch's output is connected to an AND gate with the latch that determines the state of the switch. The circuit uses the reflected output voltage
on the primary winding during the flyback time period to determine if the capacitor is charged.


Figure 5-2: Simulated primary and secondary currents of test circuit.

We use LT1800s for the operational amplifier and we use LT1720s for the comparators. The op amp was chosen since it has an acceptable slew rate. The one-shot is made by using an inverter and an AND gate, shown in Figure 5-1. The one-shot connected to the enable pin does not require a specific time length. However, the one used after the switch turns on, needs to have a duration long enough to blank the stray capacitor current, around 100nS. The S-R latches are J-K Flip-flops with preset and reset. The clock and the J-K inputs are tied to ground and only the preset and reset inputs are used. A LTC1693, a CMOS gate drive, is used to drive the MOSFET.

After designing the flyback controller on paper, it was tested in Spice. One major error was found in the paper design. A one-shot after the secondary current comparator was necessary. Although the magnetizing current will always remain above the secondary current limit, the secondary current drops to zero during the switch on period, therefore the secondary comparator output is high. When the primary current limit is reached, both inputs of the S-R latch are high, which is an undefined state. A rising edge event from the secondary output comparator is unique to the secondary current crossing the current limit from a higher current. A one-shot is the ideal circuit to capture this rising edge and turn the switch on. After finding this error, the circuit simulated in Spice as expected. The one-shot used in blanking the initial primary current was combined with this new one shot, since they fired at the same time.

### 5.3 Construction

While Spice simulations are useful for debugging purposes, actual testing in lab is necessary to make real performance measurements. Since most of the components are only available in surface mount packages, layout software was utilized to expedite routing of the copper board in-house. Constructing the board consists of determining the component packages, figuring out special requirements for traces, and paying attention to large switching current paths. The backside of the copper board is the ground plane. Many of the digital interconnects, required external wiring. A bypass capacitor was added near each of the voltage pins of the digital and analog parts used in the design.

When the board layout was complete, a routing machine was used to make the board. This process proceeded smoothly. To put the final touches on the board, the excess copper was removed with a soldering iron and tweezers. First, the digital logic for the one-shots were placed on the board. Because these were designed from scratch extensive tests were done to verify their performance. A major problem was detected with the first design, as shown in Figure $5-1$ without the included diode. The one-shot needs a time in the low state to reset. The short off-time of the switch does not allow the one-shot to reset, so the design was modified with a diode to quickly charge the capacitor to its high state. After completion of the one-shots, the rest of the digital logic was connected. The next step was to place the analog components. This portion was straight forward and there was no easy way to test their individual functionality. After all components were properly assembled, the circuit was probed. The output of the op amps were probed to show the primary and secondary currents. Nothing worked on the first attempt. A couple of wiring errors were then found by reexamining the circuit. After fixing these errors, the circuit charged the capacitor.

### 5.4 Debugging

Further testing with a load to operate the flyback with a steady state output voltage. The circuit would operate initially, but then the output voltage would collapse. After a careful inspection of voltages at the time of the collapse, the collapse was linked to noise in the primary and secondary current sensing circuitry. The adjustable voltage levels for the current limits picked up noise from external sources and would cause the comparator to change states. A premature trigger of the primary current limit and a high $\alpha$ allows the circuit to enter an invalid state where the secondary current never exceeds the secondary current limit, thus not triggering the one-shot to turn the switch back on. A quick solution to the problem was to add more capacitance to the voltage limit inputs of the comparators and minimize the length of the wires feeding into these inputs.

### 5.5 Boundary Mode Operation

After studying the possible benefits of boundary mode operation, a circuit was added to allow the controller to operate in boundary mode. Instead of turning the switch on when the secondary current falls below the limit, the switch monitors when the switch pin falls below $V_{i n}$. The ringing settles at $V_{i n}$. At low voltages, the amplitude of the ring is small, and the switch pin voltage falls slightly below $V_{i n}$. To add a noise margin, the comparator trips at a voltage slightly above $V_{i n}$ to guarantee the switch turns on, but below the lowest possible flyback period voltage. The circuit is shown in Figure 5-3. The resistive dividers move the comparator trip point slightly above $V_{i n}$. They also lower the inputs to the comparator to keep it within its common-mode range. The diode also protects the comparator by limiting the voltage seen at the input of the comparator to a diode drop above $V_{i n}$. The one-shot is already present in the existing circuitry. The secondary current comparator usually connects to the input of the one-shot. This input can be switched back and forth to change the circuit from boundary mode operation to continuous operation.

### 5.6 Final Product

The output of the flyback capacitor charger is shown in Figure 5-4. The input current waveform is filtered with a large bypass capacitor to show the average input current. The average input current stays constant over the charge cycle in this example using a 16 uH magnetizing inductance, an output capacitance of 150 uF and an input voltage of 3.5 V . The output voltage increases as the square root of the time elapsed, since the energy input is constant and energy storage in a capacitor is proportional to voltage squared.

The final constructed circuit uses two separate power supplies. One power supply is for the digital logic, the gate driver, comparators, and the operational amplifiers and the other is for the energy to be transferred to the capacitor. The general architecture of the circuit makes it capable of accepting any input voltage, but low voltages suffer from high losses in efficiency. The maximum input voltage is set at 10 V by the ceramic input capacitor, but could easily accommodate higher voltages. The MOSFET is rated at 2 A with a breakdown voltage of 60 V . This MOSFET allows the primary


Figure 5-3: Circuit diagram of the boundary mode controller.
current limit to be as high as 2 A and the output voltage to reach 600 V with 10 turn transformer. With a low leakage inductance transformer, the $\alpha$ may be set as high as 0.9. The maximum input power is 16 W with an input voltage of 10 V , current limit of 2 A , and an $\alpha$ of .9. But running at this power for an extended period of time would need adequate heat sinking. The charge time for this circuit follows equation 5.1. The plot in Figure $5-4$ shows a charge time of 6.7 s . The equation predicts a charge time of 7.2 s . The charge time is slightly higher due to dielectric absorbtion in the capacitor, lowering the value of the capacitance with a quick charge. Experimental charge times predict other charge times with different photoflash capacitance values nicely by scaling.

$$
\begin{equation*}
t_{\text {charge }}=\frac{C_{\text {load }} V_{\text {out }}^{2}}{I_{\text {in }} V_{\text {in }} \mu} \tag{5.1}
\end{equation*}
$$



Figure 5-4: Scope shot: Ch1 is output voltage, and Ch3 is input current(.25A/div).

## Chapter 6

## Transformer Optimization

### 6.1 Introduction

This chapter focuses on the transformer and its impact on the performance of a flyback capacitor charger. A brief outline of the requirements of the transformer is presented in the first section. Within these requirements, the design still remains flexible. In the model developed to simulate the capacitor charger, the winding resistances and core loss are modeled. There is no easy way to model some of the parasitic effects of the transformer such as leakage inductance, winding capacitance, and proximity loss. The tradeoffs of these parameters are discussed without the use of simulations.

### 6.2 Transformer Basics

As discussed before in previous sections, the transformer serves a dual role in a flyback capacitor charger. The transformer protects the switch from the high output voltage, and it stores energy in the core. While the transformer is the simplest component to manufacture in a flyback, the transformer has the largest impact on efficiency. In addition, a transformer is also the largest component in a flyback capacitor charger. Therefore, the transformer is the most important component to optimize and analyze in depth.

The transformer in the flyback capacitor charger has two windings. The primary
winding is connected to the switch on the input side. The secondary winding is connected to the output capacitor through a diode. To protect the switch from high collector-to-emitter voltages, the secondary winding usually needs ten times the amount of turns as the primary winding in typical photoflash applications. In contrast with a forward converter, the flyback converter intentionally stores energy in the transformer's magnetizing inductance. In most transformers, the magnetizing inductance is made as high as possible using an un-gapped core. A flyback converter uses a gapped-core transformer with energy stored as a magnetic field in the air gap. The amount of magnetizing inductance in the core is most important to the operation of the flyback converter because it determines the operating frequency of the flyback capacitor charger along with the primary and secondary current limits. The magnetizing inductance can be measured with an impedance analyzer on the primary and secondary windings, but this measurement includes the leakage inductance. The approximate turns ratio is found by taking the square root of the ratio of these two inductances. When the flyback waveform is used to determine when the output has reached its final value, the switch off portion needs to be long enough for a speedlimited comparator to trigger once the output voltage is reached. A minimum offtime will be specified by the controller, and this corresponds to a minimum secondary magnetizing inductance.

In an inductor, the winding is wrapped around the core as closely as possible to keep the flux in the core. But the core's permeability is only several orders of magnitude larger than air so some flux is leaked into the surrounding air. When this happens in a transformer, the flux leaked into the air creates an inductance in series with the transformer, as shown in Figure 6-1. Leakage inductance is the worst when the primary and secondary windings are poorly coupled. Poor coupling occurs when flux from one winding has significant room to go between itself and the other winding. Coupling becomes worse with a winding area with a small width, since it leads to the use of many layers. These layers create a lot of space between the primary and secondary windings. A winding window with a large width is best to lower the amount of layers, thus decreasing leakage inductance. To improve leakage inductance, the
primary and secondary may be interleaved. It is typically possible to decrease leakage inductance by half with interleaving. In the flyback capacitor charger, multiple wires may be used for the primary winding and each of these windings could be interleaved with the secondary winding. This technique is difficult to do by hand for prototypes and is best left to transformer manufacturers.

While interleaving will reduce leakage inductance, it will increase the capacitance between the windings. This capacitance will increase the total lumped capacitance from the switch node and the secondary winding node, which can be analyzed as a reflected capacitance on the primary switch node. In continuous operation mode, the capacitance on the switch node is charged to the transformer's step-downed output voltage when the switch turns on. This capacitance discharges through the closed switch. At lower output voltages, the amount of energy lost is low, but it increases with the square of the output voltage, and becomes the dominate loss term at higher voltages.

The capacitance between the windings is distributed throughout both of the windings and cannot be well modeled with a lumped capacitance linking the two windings. If the secondary winding is put on top of the primary winding, the capacitance is greater on the section of the winding directly on top of the primary winding. This pin of the secondary winding should be connected to ground to minimize the effect of the interwinding capacitance. In experiments, the efficiency decreases by at least $5 \%$ if the preferential transformer connection is not used.

Losses in the core and the windings are discussed in Chapter 3. The primary winding DC resistance loss is the greatest out of these losses. The duty cycle of the charger is relatively constant above 100 V where is spends most of its time. Therefore, the amount of power lost in the primary and secondary windings is constant over the charge cycle. While decreasing the primary winding resistance helps efficiency, its returns are marginal because the switch's on-resistance is in series with the winding and is usually much higher in a well-designed transformer. In addition, a larger gauge primary winding increases the leakage inductance.

### 6.3 Hand-winding Transformers

Transformers from manufacturers are wound by machine. A machine is able to place the windings in a similar way each time. This results in nearly identical transformers with the same leakage inductance and capacitance. It is impossible to achieve this consistency in placing the windings by winding the transformer by hand. In the lab, in which this thesis was carried out, there is a transformer winding machine, but it does not have the capabilities of directing the windings. This type of machine will still be considered winding a transformer by hand since its only function is to spin the bobbin and keep a count of the number of turns.

There are not many documents or textbooks written on the art of winding transformers. The only way to learn is by experimentation. At first, all the transformers made by hand had a significant fraction of their magnetizing inductance as leakage inductance. The high leakage inductance was linked to the secondary winding, which is significantly harder to wind with ten times the amount of turns. The wires for the secondary winding should be done slowly and adjacent turns should not be on top of each other but just to the side. The wire should be swept from side to side in the bobbin at a good pace to not bunch adjacent wires.


Figure 6-1: Transformer model.

### 6.4 Measuring Inductance Values for Transformer $\pi$ Model

A good way to understand the effects of a transformer is to model it with the $\pi$ model, shown in Figure 6-1. The turns ratio is calculated by exciting the primary side with a sinusoid and measuring the magnitude of the resulting sinusoid, or by counting the number of primary and secondary windings.

$$
\begin{equation*}
N=\frac{V_{p}}{V_{s}}=\frac{N_{p}}{N_{s}} \tag{6.1}
\end{equation*}
$$

Measuring the magnetizing inductance and the leakage inductances is more complicated. The coupling coefficient is the first step in calculating these values. The coupling coefficient is found with the following equation.

$$
\begin{equation*}
k=\sqrt{1-\frac{L_{\text {short }}}{L_{\text {open }}}} \tag{6.2}
\end{equation*}
$$

With an impedance analyzer, measure the inductance of the primary winding with the secondary side open. This is equal to $L_{\text {open }}$. Now with the secondary side shorted, measure the inductance of the primary winding with the secondary side shorted. This is the value $L_{\text {short }}$ in the above equation. The primary leakage inductance is equal to $(1-k) \cdot L_{\text {open }}$, and the secondary leakage inductance is equal to $(1-k) \cdot L_{\text {open }} \frac{1}{N^{2}}$. The magnetizing inductance is equal to $k \cdot L_{\text {open }}$. The primary and secondary winding resistances are measured with an ohm meter.

### 6.5 Effects of Leakage Inductance

As discussed in previous sections, leakage inductance causes a spike as shown in Figure 6-2. When the switch turns off, the majority of the energy stored in the core transfers to the secondary winding, but not all the flux from the primary winding is linked to the secondary winding. This inductance is known as the leakage inductance.

The leakage inductance current flows into the switch node and charges the switch's capacitance. This results in a voltage spike approximated by $V_{l e a k}=I_{l i m} \sqrt{\frac{L_{\text {leak }}}{C_{s w}}}$. The current reverses and the energy stored in the switch capacitance flows through the transformer to the input, because there is no voltage across the primary side of the ideal transformer at this point. Simultaneously, the magnetizing current is linearly charging the parasitic capacitance of the transformer to the output voltage on the secondary side. As soon as the primary terminal voltage of the ideal transformer increases from the charging of the parasitic capacitance, the switch ring current becomes transferred to the output. As seen in Figure 6-2, once the switch voltage reaches its final value, the ringing dies off completely. This explains how the leakage inductance energy is recovered. Although high leakage inductance does not result in substantially higher losses, the voltage spike caused by the leakage inductance requires a switch with a higher collector-emitter voltage rating. Since the voltage spike is related inversely to the the capacitance on the switch node, an external capacitance on the switch pin lowers the peak of the leakage inductance voltage spike. This capacitance does not add substantially to the equivalent capacitance on the switch pin when determining the capacitor loss, since the transformer contributes almost all the capacitance on the switch pin.

Along with the voltage spike on the switch node, one of the main problems with leakage inductance is the ringing it causes on the secondary winding. This ringing, along with the current in the secondary winding, is shown in Figure 6-3. In this example, the output voltage is at 250 V and a transformer with $16 \mu \mathrm{H}$ of magnetizing inductance is used. At the peak of the negative ring, the diode has 430 V across it, and this voltage across the diode can be as high as 600 V when the output voltage is at its final value. In our test circuit, we use two diodes with a DC blocking voltage of 200 V , but are capable of handling voltage spikes up to 650 V . Not accounting for this ringing could cause the diode to go into reverse breakdown and blow up while in operation. We have tested this phenomenon in lab and this occurs at 375 V for a single diode. With two diodes, the circuit will not breakdown under any conditions.

Another effect of the leakage inductance, occurring with the above phenomenon,


Figure 6-2: Scope shot: Ch3 is primary current (AC coupled, . $1 \mathrm{~A} /$ div), and Ch4 is switch voltage.


Figure 6-3: Scope shot: Ch2 is secondary winding Pin, and Ch3 is secondary winding current.
is the primary current ring when the switch turns on. When the switch turns on, the secondary current falls to zero, and the diode turns off. When the diode turns off, the capacitance from the diode holds the secondary winding pin at the output voltage. Normally at this point, the magnetizing inductance would have 3 V across it from the input voltage and the current in the inductance would increase linearly. But this operation does not account for the effects of the leakage inductance and the ringing on the secondary winding pin. In addition to forming a second-order network, the leakage inductances from the primary and secondary form an impedance divider. The output of this impedance divider is the node of the magnetizing inductance, which cannot be probed experimentally. The ringing on the secondary produces a current through the primary winding as shown in Figure 6-4. The current on the secondary during this time period is shown in Figure 6-5. When the ringing dies, the voltage drop across the primary winding's leakage inductance becomes zero, and the magnetizing current begins to increase.


Figure 6-4: Scope shot: Ch3 is real primary current (ac coupled, 1A/div), and Ch4 is primary switch pin.

### 6.6 Effects of the Transformer's Capacitance

As discussed above, the transformer's parasitic capacitance is one of the major losses in continuous operation flyback capacitor chargers, and becomes a major loss with boundary mode with small magnetizing inductances. This energy is discharged through the switch in continuous mode. In boundary mode, this energy is transferred to the magnetizing inductance with some of the energy lost in the transformer's resistance.

The transformer's capacitance is charged up when the switch turns off. The charging of this capacitance causes a delay between the transfer of current between the primary winding to the secondary winding. Two nodes in the circuit increase with the turn off of the switch. The secondary winding pin increases from $-V_{i n} N$ to $V_{\text {out }}$. The switch pin increases as a step-downed version of the secondary winding pin, therefore increasing to $\frac{V_{\text {out }}}{N}$. Since these two voltage increases are not instantaneous, these two points in the circuit have parasitic capacitance. The secondary capacitance may be reflected to the switch pin and lumped with the capacitance on the switch pin. The slope of the voltage on the switch when the switch turns off is equal to


Figure 6-5: Scope shot: Ch3 is real secondary current (inverted, $100 \mathrm{~mA} /$ div), and Ch4 is primary switch pin.
$\frac{I_{\text {lim }}}{C_{\text {para }}}$. This slope is shown in Figure 6-6 with the primary and secondary currents. The slope of the switch voltage cannot be clearly seen with the ring caused by the leakage inductance and the capacitance of the switch, but can be easily measured by dividing the voltage rise by the amount of time it takes. The value of the capacitance found corresponds nicely with the the amount of capacitance found by measuring the frequency of the fall of the switch voltage when the switch turns on by the formula as described in the Chapter about the Flyback Operation.


Figure 6-6: Scope shot: Ch1 is amplified primary current(1A/div), Ch2 is amplified secondary current( $100 \mathrm{~mA} /$ div), and Ch4 is primary switch pin.

### 6.7 Energy Storage Requirements

A transformer is required to store $\frac{1}{2} L_{p r i} I_{\text {lim }}^{2}$ in the magnetizing inductance. To make a fair comparison when comparing energy storage requirements as a function of $\alpha$, input current and operating frequency should remain constant. Since the switch offtime has very little influence on the operating frequency past 100 V and simplifies calculations greatly, the on-time will be held constant instead of operating frequency. The on-time is equal to $\frac{L_{p r i}}{V_{i n}} I_{\text {lim }}(1-\alpha)$. $I_{\text {lim }}$ decreases as $\alpha$ increases if input current is held constant. Input current is equal to $\frac{(1+\alpha) I_{l i m}}{4-2 d}$. By eliminating the $I_{l i m}$ term
and solving for the needed magnetizing inductance to keep $t_{o n}$ constant, the following equation is found.

$$
\begin{equation*}
L_{p r i}=\frac{t_{o n} V_{i n}}{(4-2 d) I_{i n}} \frac{1+\alpha}{1-\alpha} \tag{6.3}
\end{equation*}
$$



Figure 6-7: Magnetizing inductance increase with $\alpha$.
In this equation, $\alpha$ increases the inductance by a factor graphed in Figure 6-7. At an $\alpha$ of .2, the required inductance is already $50 \%$ greater. As $\alpha$ increases, the current limit decreases, but not enough to cancel out the increase in the inductance as demonstrated in the following energy equation.

$$
\begin{equation*}
E=\frac{t_{\text {on }} V_{\text {in }} I_{\text {in }}(4-2 d)}{2\left(1-\alpha^{2}\right)} \tag{6.4}
\end{equation*}
$$

This equation shows that an $\alpha$ of zero will result in the lowest energy storage
requirement for the transformer. This is an important result for determining what type of charge scheme to use. With just this result, an $\alpha$ of zero charge scheme appears optimal in reducing the size of the transformer.

## Chapter 7

## Experimental Results

### 7.1 Introduction

The experimental results are a combination of measurements from lab and simulations performed with MATLAB. Winding transformers has limitations mentioned in the previous chapter, so simulations are relied upon to expand the experimental findings. To do so, the simulations are compared with the experimental results to show consistent simulation results from different transformers and charging schemes. Simulations are then used to show the affect of changing the magnetizing inductance, charge scheme, turns ratio, and the size of the core.

### 7.2 Correlation of Simulated and Measured Data

The following efficiency curves show the experimental data with the simulated data. The MATLAB simulator code is listed in Appendix A. The simulator uses the transformer's parameters measured by an impedance analyzer. The primary input current is adjusted to keep the average maximum input current below 500 mA . For alpha values other than zero, the secondary current limit is raised as the primary current is lowered until the ratio of these currents is alpha and the maximum average input current is 500 mA . An $\alpha$ of zero signifies boundary mode. An $\alpha$ of .1 is approximately continuous mode with the current falling to zero at voltages higher than 100 V because
of small delays in turning the switch on.
The boundary mode simulation only loses $30 \%$ of the energy stored in the parasitic capacitance on the switch pin, whereas the continuous simulations lose all the energy. The $30 \%$ factor is based on the amount of the capacitor's energy transferred to the magnetizing inductance. The primary current limit for boundary mode is also adjusted to a higher value in the simulation because of the time it takes to reverse the negative magnetizing current caused by the capacitance.

Figure 7-1 and Figure 7-2 show the simulation data and the test data on the same efficiency curve. The points on the two curves match fairly well and the slope of the curves also correspond, although the curves are not exactly the same. The most likely source of errors in the boundary mode is the modeling of the parasitic capacitance loss. The $30 \%$ factor used to determine the amount of energy lost in the capacitance is not constant over the output voltage range. More energy is lost during the higher output voltages due to the primary resistance because the magnetizing current becomes much more negative. This resistance loss scales with $I^{2}$ and becomes greater with the more negative peak current. Figure 7-2 shows a much stronger correlation between the two sets of data.

By comparing the efficiencies from Figure 7-1 and Figure 7-2, the advantage of boundary mode over continuous conduction mode is clearly shown. At 300 V , the difference between the two modes is $6 \%$. The boundary mode scheme does use a peak current of 1.38 A compared to a peak current of 1.24 A for the continuous conduction mode scheme. With everything else kept the same, an increase of $10 \%$ in peak current provides a $6 \%$ increase in efficiency. In addition to the loss savings of recovering energy from the parasitic capacitance, the boundary mode scheme saves energy from a lower operating frequency, 165 kHz compared to 200 kHz . The advantage of boundary mode is most profound at lower magnetizing inductance values.

Figure 7-3, Figure 7-4 and 7-5 show a similar degree of correlation using a 16 uH transformer as the previous set of efficiency curves. The operating frequency of the circuit is much less than the previous 10 uH transformer, therefore the losses due to the parasitic capacitance is much less, but the peak efficiency drops slightly because of
more primary winding resistance. The efficiency difference between boundary mode and continuous mode is much less in this case, around $2 \%$.

For the $24 u H$ transformer, the efficiency curves are done for four different charging schemes. All the charging schemes are well predicted by the simulation. As can be seen from the curves, efficiency decreases with higher alpha due to the higher operating frequency discharging and charging the parasitic capacitance. The boundary mode scheme with a $24 u H$ transformer performs much better than the rest of the continuous conduction mode charging schemes.

| Magnetizing Inductance | Alpha | Total Charge Eff |
| :--- | :--- | :--- |
| 10 uH | 0 | $78 \%$ |
| 10 uH | .1 | $75.6 \%$ |
| 10 uH | .2 | $72.4 \%$ |
| 16 uH | 0 | $77.2 \%$ |
| 16 uH | .1 | $75.4 \%$ |
| 16 uH | .2 | $75.5 \%$ |
| 24 uH | 0 | $77.1 \%$ |
| 24 uH | .2 | $77.0 \%$ |
| 24 uH | .4 | $75.0 \%$ |
| 24 uH | .6 | $69.7 \%$ |

Table 7.1: Total efficiency using capacitor energy method.

### 7.3 Magnetizing Inductance

Magnetizing inductance plays an important role in a capacitor charger. The size of transformer is also strongly a function of its magnetizing inductance. In addition to storing the energy for the flyback capacitor charger, it determines the operating frequency. A low magnetizing inductance causes a high operating frequency, leading to large losses in continuous conduction mode operation. A large magnetizing inductance requires a large amount of turns, which increases the primary and secondary winding resistances with a fixed winding area. Figure 7-10 shows the effect of changing magnetizing inductance on efficiency. The plot displays a continuous conduction mode capacitor charger and a discontinuous capacitor charger. In both


Figure 7-1: Efficiency versus output voltage for $\mathrm{L}=10 \mathrm{uH}$ and $\alpha=0$.


Figure 7-2: Efficiency versus output voltage for $\mathrm{L}=10 \mathrm{uH}$ and $\alpha=0.1$.


Figure 7-3: Efficiency versus output voltage for $\mathrm{L}=16 \mathrm{uH}$ and $\alpha=0$.


Figure 7-4: Efficiency versus output voltage for $\mathrm{L}=16 \mathrm{uH}$ and $\alpha=0.1$.


Figure 7-5: Efficiency versus output voltage for $\mathrm{L}=16 \mathrm{uH}$ and $\alpha=0.2$.


Figure 7-6: Efficiency versus output voltage for $\mathrm{L}=24 \mathrm{uH}$ and $\alpha=0$.


Figure 7-7: Efficiency versus output voltage for $\mathrm{L}=24 \mathrm{uH}$ and $\alpha=0.2$.


Figure 7-8: Efficiency versus output voltage for $\mathrm{L}=24 \mathrm{uH}$ and $\alpha=0.4$.
configurations, an optimal amount of magnetizing inductance exists to maximize the efficiency. In this case, the simulations uses a model of a TDK EPC10 core. The optimal amount of magnetizing inductance for a core depends on many transformer characteristics, such as winding area and $A_{l}$. Using MATLAB is a cumbersome way to calculate the optimal amount of magnetizing inductance; most designers would prefer an analytical expression. As mentioned above, the capacitance loss is greatest at low magnetizing inductance, and the loss from the primary winding resistance is greatest at higher magnetizing inductances. The peak efficiency occurs approximately when these two losses equal each other. In continuous conduction mode, the optimal magnetizing inductance is approximated by the following formula.

$$
\begin{equation*}
L_{p r i}=\frac{V_{i n} C_{s w} V_{\text {mid }}^{2}}{1.85 * I_{l i m}^{3}\left(R_{s w}+R_{p r i}\right)\left(\alpha+\frac{1}{3}(1-\alpha)^{2}\right)} \tag{7.1}
\end{equation*}
$$

This formula has its weaknesses, since it includes the primary resistance of the transformer and the parasitic capacitance of the transformer. This formula needs to be used in an iterative fashion to solve it exactly. By sizing the primary winding


Figure 7-9: Efficiency versus output voltage for $\mathrm{L}=24 \mathrm{uH}$ and $\alpha=0.6$.
correctly, the primary resistance can be specified before winding the transformer, but the parasitic capacitance has to be approximated and cannot be calculated a priori.

Boundary mode operation cannot be approximated in this fashion, since the parasitic capacitance loss is less substantial. Other loss terms, such as core loss and the rise and fall time losses also increase with higher operating frequencies. An analytical expression would be of little use with the added complication of all these loss terms. Figure 7-10 shows that the needed magnetizing inductance for discontinuous mode is less than CCM by about $20 \%$. As a rough estimate, the optimal value for a CCM operation capacitor charge can be found with Equation 7.1 and reduced slightly. The fall at lower magnetizing inductance is not nearly as steep for boundary mode, therefore the designer may go to a lower magnetizing inductance without losing much efficiency.

The simulation optimizes the transformer's windings for each magnetizing inductance. In actuality, there are not the perfect size wires for each of the magnetizing inductances, therefore the windings would be slightly smaller, resulting in more winding resistance. In addition, the primary winding has to lay flat within the winding window to keep the packing fairly tight and leakage inductance low. Tight packing depends on the number of turns for the primary and the width of the wire. Since a tightly packed primary winding is important to the operation of the charger, the primary windings amount of turns and the size should be optimized for packing around the optimal value. The optimal value should be used as a guide, not as a tight design constraint.

### 7.4 Alpha Comparisons

Since proving the simulation can be relied upon, simulations will be used to show efficiencies for different charging strategies with optimal transformer characteristics. The simulations will calculate the resistance of the primary and secondary windings with the winding window area, mean length per turn, packing factor, primary to secondary window allocation, and the number of turns.

In the previous section, the simulation was shown to be reliable in predicting the efficiency behavior of a flyback capacitor charger. With the aid of simulation, the optimal charging strategy will be studied. From the efficiency curves above, when alpha is increased while keeping input power constant, the efficiency decreases at higher output voltages. While this supports that low alpha has higher efficiency, the same transformer is used for all the charging schemes, not one optimized for that particular charging scheme. When alpha is increased, the primary current limit decreases along with the magnetic flux density, allowing the transformer to have a smaller gap length or more windings. Without the aid of a core manufacturers equipment, changing the gap length is difficult. Even changing the amount of windings is not trivial, since the winding wire size needs to shrink. In simulation, these tasks can be done easily, allowing many points of data.

To study the optimal charging strategy, alpha will be swept from 0 to .9 using MATLAB. As mentioned briefly above, the transformer may be optimized by changing the gap length while leaving the amount of turns untouched, or by changing


Figure 7-10: Inductance versus efficiency.
the amount of turns while leaving the gap length constant. The first method keeps the resistance of the primary and secondary windings while increasing the magnetizing inductance linearly with decreasing peak primary current. The second method increases the resistance of the primary and secondary windings linearly with turns ratio, but the magnetizing inductance increases quadratically with decreasing peak primary current. Figure 1 shows efficiency versus alpha with a variable gap length. Figure 2 shows efficiency versus alpha with a variable amount of turns. Coincidentally, the efficiency drops similarly in the two methods. As alpha increases, the operating frequency becomes too great with either of these methods and the loss due to the capacitance dominates.

These two graphs clearly show that an alpha of .4 should not be exceeded with this particular core. The efficiency at high values of alpha improves if a larger core capable of higher magnetizing inductance is used, but the extra space for the larger core is a higher design cost than the savings in peak current for the switch.


Figure 7-11: Efficiency versus alpha with variable gap length.


Figure 7-12: Efficiency versus alpha with fixed gap length.


Figure 7-13: Maximum frequency versus alpha with variable gap length.

### 7.5 Turns Ratio

Four transformers were wound to explore the affects of a higher turns ratio on efficiency. The parasitic capacitance from the transformer has a large impact on the efficiency, even using boundary mode operation. The turns ratio has a strong affect on the amount of capacitance reflected to the switch pin from the secondary windings parasitic capacitance. Any capacitance on the secondary side is multiplied by the square of the turns ratio when reflected to the primary side. But the turns ratio also decreases the output voltage reflected upon the primary side when the switch is off. In regards to the capacitance loss on the switch, the effects of varying the turns ratio cancel each other and do not change this loss term. A simulation sweep of turns ratio is not possible because of these changes in the capacitance on the switch pin. Table 7.2 lists the turns ratio equivalent capacitance on the switch and a factor proportional to the amount of energy stored in this capacitance at 300 V for the four transformers. The energy factors do not show any correlation with the amount of turns and their variations likely result from winding differences.


Figure 7-14: Maximum frequency versus alpha with variable turns.

| Turns Ratio | Equivalent Switch Capacitance | Cap. Energy Factor |
| :--- | :--- | :--- |
| 10 | 1.9 uF | 1710 |
| 13 | 3.55 uF | 1863 |
| 16 | 4.3 uF | 1511 |
| 20 | 6.5 uF | 1462 |

Table 7.2: Equivalent switch capacitance effects with turns ratio.


Figure 7-15: Efficiency curves with different turns ratios.

Figure 7-15 shows the efficiency curves for a capacitor charger with different transformer turns ratios. The transformers are wound with 10 uH of magnetizing inductance, while also using the same wire sizes for all the turns ratios. Thus, the primary winding resistance is constant, but the secondary winding resistance increases linearly with the turns ratio. The efficiency curves are relatively flat, since they are all operating in boundary mode. The higher turns ratio transformers suffer from higher losses at lower voltages. The secondary resistance becomes much more important at lower output voltages, because the secondary current decreasing more slowly causes the diode conduction period to be a larger portion of the switch cycle. This effect is shown clearly in the individual power loss figures, Figure 4-1 and Figure 4-2. The efficiency curves vary by only one-percent, which is unsubstantial in a practical application and could fluctuate with the measurement techniques used. Therefore, turns ratio does not dramatically increase the efficiency of the capacitor charge. Along with the increased area used in the winding window, the higher turns ratio transformers increase the charge time as shown in Equation 3.14, or increase the primary peak current to compensate for additional charge time.

### 7.6 Scaled Transformer Core

From the previous data, the boundary mode scheme outperforms the continuous conduction mode converters with any value of alpha. To focus directly on this type of converter, an TDK EPC10 is scaled using MATLAB to find the optimal power to volume ratio. Although a manufacturer does not create a transformer by keeping the aspect ratio the same for all the dimensions, this gives a good estimate of the size the transformer can be made without efficiency suffering. The volume factor in the figures is cubed when multiplied with volume, squared when multiplied with area, and just multiplied with length. The volume factor effects the cross-section area, volume of the core, mean length per turn, winding window area, and the maximum number of turns. The inductance per turn squared is kept constant, so gap length is multiplied by the volume factor squared. Since the EPC10 has plenty of energy


Figure 7-16: Efficiency versus volume factor.


Figure 7-17: Efficiency per unit volume versus volume factor.


Figure 7-18: Magnetizing inductance versus volume factor.


Figure 7-19: Primary peak current versus volume factor.
storage at its original size, the number of windings has to be limited to keep the inductance value at 16 uH , shown in the previous section as the optimal inductance value for this particular core in boundary mode. The MATLAB model averages the negative current in the primary inductance caused by the boundary mode scheme with the primary current to keep the input current at 500 mA . With a small magnetizing inductance, the negative current becomes comparable to the original peak current, therefore peak current has to increase substantially.

Figure 7-18 shows how the inductance scales with the volume factor. The maximum turns causes the inductance to be constant from a volume factor of .55 to 1 . As a result the efficiency, shown in Figure 7-16, stays relatively constant for this range of volume factor. At about 10 uH , the efficiency begins to fall because the increase in operating frequency causes the capacitance loss to become the dominate loss. Along with the increase in frequency, the main reason for the sharp fall in efficiency is the sharp increase in primary peak current, shown in Figure 7-19. The Efficiency volume factor is calculated with the following equation, where e is the efficiency, $P_{i n}$ is the input power, and core volume is the effective volume of the core.

$$
\begin{equation*}
\frac{e^{3} P_{\text {in }}}{\text { core volume }} \tag{7.2}
\end{equation*}
$$

The power to unit volume factor is plotted against the volume factor in Figure 7-17. The plot shows a peak power density at a volume factor of 0.3 . A volume factor of 0.3 corresponds to a magnetizing inductance of 1.7 uH and a total charge efficiency of $65 \%$. This value of magnetizing inductance operates the converter at a frequency above 600 kHz . While this frequency would not work with the LT3468, since it needs a longer off-time to compare the flyback voltage waveform, it works in the modified test circuit. The modified test circuit uses zener diodes to detect when the final output voltage is reached. Since scaled versions of the EPC10 do not exist, a EPC10 core was used to produce 2 uH with only a small section of the winding window. The efficiency of the converter was $62 \%$, very close to the predicted value.

Two scope photos of the 2 uH transformer are shown in Figure 7-20 and Figure


Figure 7-20: Scope shot: Ch1 is secondary current waveform(100mA/div), and Ch4 is switch waveform.


Figure 7-21: Scope Shot: Ch1 is primary current(1A/div), and Ch3 switch waveform.

7-21. In the first oscilloscope screen shot, the ringing secondary winding current is shown with the switch pin voltage. Unlike continuous converters that compare the secondary current waveform to a reference to determine when to turn the switch on, the boundary mode converter does not need a slow and clean waveform to determine when to switch. This allows more flexibility with the amount of magnetizing inductance and the amount of leakage inductance.

### 7.7 Experimental Conclusion

The experimental results conclusively support the use of the boundary mode control scheme. From the results, a boundary mode controller is shown to operate more efficiently with the same transformer than a continuous mode controller. Even with a continuous controller scheme, the highest efficiency is achieved at an $\alpha$ of zero. In addition, the boundary mode controller is easier to implement since it only needs to monitor the primary winding current. The continuous mode controllers do allow lower peak current levels for the transformer and the switch. However, as shown in Chapter 6, the lower of peak current in the transformer does not offset the amount magnetizing inductance needs to increase to keep operating frequency constant.

The turns ratio of the transformer should only be high enough to protect the switch from the output voltage. Most of the capacitance from the transformer is from the secondary side. Therefore, lowering the voltage on the primary side by increasing the turns ratio does not lower the amount of energy stored because the capacitance is higher when reflected to the primary side. This all leads to using the lowest turns ratio possible while still protecting the switch. The lower turns ratio require less secondary turns, which frees winding window area for thicker primary or secondary wires.

Choosing the optimal amount of magnetizing inductance depends on many factors. With a boundary mode control scheme, the process is much simplified because the value of the magnetizing inductance has less of an impact on the performance of the capacitor charger. The volume factor experiment for boundary mode control shows
that there is an optimal size for the transformer. This value results in a marginally acceptable efficiency, but results in a good trade-off for the space conscience camera manufacturer. These optimal values are currently too low for the comparator in Linear Technology's LT3468. For this controller, the magnetizing inductance should be chosen so the off-time is the smallest possible time the comparator can work consistently with.

## Chapter 8

## Flash Unit

### 8.1 Introduction

### 8.2 Self-Oscillating Capacitor Charger



Figure 8-1: Self-oscillating capacitor charger circuit diagram.

The first widely used method to charge a photoflash capacitor was a self-oscillating capacitor charger. This capacitor charger is perfect for disposable cameras where the product is highly cost sensitive. The main components are a high turns ratio transformer, a transistor, a diode, an output capacitor, and an input capacitor. The circuit diagram can be seen in Figure 8-1. [11] The circuit starts with current following through the $2.2 k \Omega$ resistor through the F winding and into the $0.01 \mu F$ capacitor. As soon as this capacitor reaches the threshold of the transistor, the transistor turns on and current flows through the P winding. The P winding drives the
current in the S winding, but the S winding determines the voltage drop across the P winding. Since the turns ratio between the P and S winding is approximately 125, the voltage across the P winding is the current output voltage divided by 125 . The turns ratio of the transformer should be kept close to the beta of the transistor since the base current is supplied by the S winding and the P winding is the collector current of the transistor. The transformer will eventually saturate, causing the S winding current to go to zero and turning the switch off. Unlike the flyback capacitor charger, the core does not need to store much energy, since the current is directly transferred between the primary and secondary windings; never being stored in the magnetizing inductance. With the switch off, the magnetizing current of the transformer core will reset itself with the positive voltage across the F winding. This positive voltage will again create a current across the F winding into the $0.01 \mu F$ capacitor, and the switch will turn on to start the cycle again. As the output voltage becomes higher, the voltage across the P winding will become almost as large as the battery input. The transistor collector-to-emitter voltage will become small and eventually the transistor will not be capable of conducting current. At this point, the charging of the capacitor stops and this is the maximum output voltage of the charger. Therefore, the final output voltage on the output capacitor will be a function of the input voltage.

While this circuit requires a minimum amount of expensive components, it also has many disadvantages compared to a flyback capacitor charger. As mentioned above, the final output voltage will decrease linearly with the input voltage. A typical battery will lose over $30 \%$ of its value over its lifetime. Even though most Xenon flash bulb used in cameras only require a minimum voltage of 225 V and can handle a maximum voltage of up to 350 V , they operate more efficiently at their nominal voltage, 330 V . This operating range is only slightly larger than the variations in output voltage of a self-oscillating capacitor charger using alkaline batteries. The other major drawback of this solution is its size. The transformer required has three windings with the 1:125 primary-to-secondary winding ratio. Peak currents through the primary winding will go as high as 5A. A battery cannot supply a current this high so a large input capacitor is needed to average the current consumption.

In addition to the size and output variations, the self-oscillating capacitor charger will suffer large power losses in the transistor at low output voltages of the charge cycle. The collector-to-emitter voltage is equal to the input voltage minus the output voltage divided by the turns ratio. The power loss is the collector-to-emitter voltage times the input current. The charger spends more time at higher voltages during the charge cycle, but there is a major advantage of using a flyback capacitor charger which uses the transistor as a switch. As a switch in the flyback capacitor charger, the collector-to-emitter voltage is determined by placing a on-resistance in series with an ideal switch.

### 8.3 Xenon Bulb

Instead of light emitting from a filament as in a incandescent, the Xenon gas ionizes and emits light. Michael Faraday discovered light emissions from a tube with lowpressure gas in 1838. Geissler continued Faraday's work and made the first practical low-pressure gas tubes with Krypton, which were called Geissler tubes. In the 1920s, the Seguin brothers were the first to adapt the Geissler tube to a working stroboscope. The Seguin brothers drastically improve the firing of the bulb by separating the main energy storage capacitor from the ignition circuit. This allowed more precise control over the large amount of energy stored in the capacitor. They used their stroboscopes to aid their work on airplane engines.

MIT professor Harold Edgerton pioneered the use of Xenon bulbs for high speed photography in the 1930's and is known as the inventor of the modern-day Xenon flash. For his research as a doctoral student, he became interested in using Xenon strobe lights to allow the human eye to see the movement of fast mechanical machinery. While most in industry saw strobe technology as a novelty, including General Electric, Edgerton saw a great need for it and developed a commercial strobe light. Soon, he became interested in the use of electronic flashes for use in high speed photography. As a pioneer in the field, Edgerton created stunning images of bullets impacting objects such as apples. These photographs were a combination of cutting
edge technology and art and were exhibited in the Museum of Modern Art in New York and the Boston Museum of Fine Art.

To trigger a Xenon bulb, a high voltage waveform ionizes the Xenon gas in the bulb. This high voltage waveform just needs to be next to the bulb and not electrically connected to either the cathode or anode. Once the gas ionizes it becomes conductive. A large capacitor connected in parallel to the bulb provides the energy for the flash. To create a short duration flash, 320 V needs to be stored in a capacitor with hundreds of microfarads of capacity. The length of the is proportional with the capacitance.


Figure 8-2: Xenon triggering circuit.

The standard circuit for triggering a Xenon flash bulb is shown in Figure 8-2. The capacitor, C4, becomes charged to the output voltage. When the switch is pressed, the capacitor, C 4 , is in parallel with the transformer, T2. The transformers magnetizing inductance forms a second-order network, and produces a decaying sine waveform. The transformer multiplies this sine waveform by the turns ratio. The generated waveform needs to be greater than 4 kV and the frequency greater than 500 KHz . [7] An example of the waveform is shown in Figure 8-3. The tube starts flashing when the first trigger pulse occurs. The current in the bulb is shown with the trigger waveform in Figure 8-4. A Xenon flashbulb manufacturer will provide typical capacitance and voltage values for the output capacitor. More luminance can be achieved with more capacitance or a higher voltage, however the life expectancy of the bulb shortens.


Figure 8-3: Xenon triggering waveform.


Figure 8-4: Xenon bulb current.

### 8.4 IGBT

An Insulated Gate Bipolar Transistor(IGBT) is the perfect device to control the current through the Xenon flashbulb once triggered. Current IGBTs for strobe applications can pulse 150A with only a 4V gate drive and no gate/base current. These devices are perfect for implementing a red-eye reduction feature on a digital camera.

Red-eye occurs when intense light is reflected off the retina of the eye. The retina reflects the blood vessels nourishing the eye. This occurs when a picture is taken from a camera with a flash and lens very close. More specifically, if the lens to the eye to the flash form an angle less than 5 degrees red eye will certainly occur if the subject is looking directly into the lens. With a point-and-shoot camera, the flash is very close to the lens. One way to reduce red-eye is to reduce the size of the iris. One method to decrease the size of the iris is to simply turn on all the lights in the room. Another way to reduce the size of the iris is to have a series of pre-flashes before the picture is taken. An IGBT can switch the current off in the bulb to create smaller flashes before the main photoflash.


Figure 8-5: IGBT circuit.

Figure 8-5 shows a typical implementation of a IGBT strobe circuit. [12] The IGBT goes in series with the Xenon bulb. The current generation of IGBTs require a 4 V gate drive, but this will decrease to 2.5 V in the next generation. The IGBT should be turned on before the Xenon bulb is triggered, and can be shut off anytime after the triggering. Figure 8-6 shows the triggering voltage, gate voltage on IGBT, and the current through the IGBT. The current through a IGBT designed for strobing can
be as high as a 150 A . These IGBTs can also block up to 400 V between the collector and emitter without any substantial gate drive current.


Figure 8-6: Illustrative IGBT waveforms.

## Chapter 9

## Conclusion

### 9.1 Summary

This thesis investigated the design of a flyback capacitor charger as used in a camera's photoflash circuitry. Since a limited about of power is available in a portable device, this thesis focused on the efficiency and size of the flyback capacitor charger solution. The experimental results and simulation data show a boundary mode controller is able to work with the smallest possible transformer with acceptable efficiency. This is an exciting result for camera manufacturer's with the decreasing size of digital cameras. A boundary mode controller saves power by recycling the energy in the parasitic capacitance of the transformer and operating at a lower frequency for a given magnetizing inductance value.

### 9.2 Further Work

As a battery is drained, its internal resistance increases, lowering the amount of current it can supply. Some camera designers would like the flyback input current to decrease with the battery energy level. The most obvious way to accomplish lowering the input current is to lower the primary peak current. However this increases the frequency of the flyback charger, and under utilizes the transformer. A way around this problem is to delay the turn on of the switch when it reaches discontinuous
mode. The delay decreases the frequency of the converter, but the amount of energy transferred to the output capacitor is the same per cycle. This feature of a flyback converter would increase the battery life of a digital camera.

The boundary mode technique also has potential applications with regulating circuits. Many of the benefits stated in this thesis would apply to a controller with a fixed output. The operating frequency would increase substantially at low output currents, but this problem can be overcome.

## Appendix A

## MATLAB Code

```
%Closs.m
%MATLAB Code to Calculate Efficiency Curve, Individual Power Losses
%and Charge Time
Vout = linspace(1, 320,1000);
dV = Vout(2)-Vout(1);
Ilim=1.3;
alpha = .0;
N=10.1;
Vin=3.3;
Vdiode = 1.2;
Vsat = .30;
Vmax = 320;
Cload=100e-6;
Cpara = 2.1e-9;
Rsw = .35;
Lleakfactor = . 3e-6;
%Core Winding
Al = 40e-9;
n=16;
L = n^2*Al;
%Core Geometry
Ae = 9.39e-6;
Ve = 167e-9;
Ec = . 5*24e-6*1.4^2;
```

```
%Winding Resistances
Rp = . 22;
Rs = 32.6;
tf = 10e-9;
tr = 15e-9;
figure(1)
cntr=0;
%Frequency Calculations
ton = L*Ilim*(1-alpha)/(Vin-Vsat);
toff = L*Ilim*N*(1-alpha)./(Vout+Vdiode)+t_blanking;
duty = ton./(ton+toff);
freq = 1./(ton+toff);
%Modified frequency to use sine wave core loss data
freq_m = (2/pi^2).*freq./(duty.*(1-duty));
%Core Loss Data Calculations
di = .5*(1-alpha)*Ilim;
dB = di*Al*n/Ae;
P}=5.17e-2.*Ve.*dB^2.45.*freq_m.^1.72
%DC Resistance in transformer
Pdcs = (1-duty).*Ilim^2.*(alpha+((1-alpha)^2)/3)*Rs/N^2;
Pdcp = (duty).*Ilim^2.*(alpha+((1-alpha)^2)/3)*(Rp);
Ip = (Pdcp./Rp).^.5;
Is = (Pdcs./Rs).^.5;
%Losses in switch
Psw = (duty).*Ilim^2.*(alpha+((1-alpha)^2)/3)*(Rsw);
Pf= (Vout./N)*Ilim*tf.*freq./2;
Pr = (Vout./N)*Ilim*alpha*tr.*freq./2;
%Loss from Leakage
Pleakfactor = . 5*Ilim^2*Lleak*freq;
%Parasitic Capacitor Loss
Pcloss = .5*Cpara.*(Vout./N).^2.*freq;
%Diode losses
Pdiode = Vdiode*Ilim*(1+alpha).*((Vin-Vsat).*(Vout+Vdiode))
    ./(2.*Vout.*(Vout+Vdiode+N*(Vin-Vsat)));
```

```
%Power from battery
Pin = .5*Ilim.*duty.*(alpha+1)*Vin;
%Iin =Pin./Vin;
```

```
Plosslf = Pdcs+Pdcp;
Plosshf = Pcloss+P+Pf+Pr;
Ploss = Pdcs+Pdcp+P+Pcloss+Pdiode+Pf+Pr+Pleak+Psw;
Eff = (Pin-(Ploss))./Pin;
%Efficiency Curve Plot
plot(Vout,Eff);
axis([100 320 . 6 . 85]);
%Individual Power Loss Plot
figure(2)
plot(Vout, Pdcp,Vout,Pdcs,Vout,Pleak,Vout, Pdiode,Vout,
    Pcloss,Vout,Pf,Vout, Pr,Vout, P, Vout,Psw);
legend('Pdcp','Pdcs','Pleak','Pdiode','Pcloss',
    'Pf','Pr','Pcore','Psw')
%Charge Time Calculation
nu=.75;
Cload=100e-6;
tb=300e-9;
Lpri=10e-6;
tcharge = (Cload*Vmax./Ilim./nu)*(Vmax./Vin+tb*Vmax/Lpri/Ilim+2*N)
```

```
%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%
%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%
%losscalc.m
%A function to calculates the efficiency for
%a capacitor charger
function [eff,data] = losscalc(Al, Wa, MLT, Ve, Ae, Bex, fex, n,
N, Ilim, alpha, Cload, Vin, Vmax, leakpercent, primarywinding)
Vout = linspace(1, 320,1000);
dV = Vout(2)-Vout(1);
Vdiode = 1.2;
Vsat = .3;
t_blanking = 0;
```

```
alpha1 = primarywinding;
alpha2 = 1-alpha1;
tf = 10e-9;
tr = 15e-9;
Rsw = . 30;
%Set to . }3\mathrm{ for Boundary Mode and 1 for Continuous
boundaryeff = .3;
Cpara = 2e-9;
C = boundaryeff*Cpara;
```

$\mathrm{L}=\mathrm{n}^{\wedge} 2 * \mathrm{Al}$;
data(1)=L;
\%Not used since leakage energy is mostly recovered
Lleak = L*leakpercent;
$\mathrm{Ku}=.3$;
$p=1.724 e-8$;
\%Negative Current from Boundary Mode
$I_{c}=(1$-boundaryeff $) . *($ Vout.$/ N) *(\text { Cpara/L) })^{\wedge} .5 ;$
\%Winding Resistances calculated from winding window data
$\mathrm{Rp}=\mathrm{p} * \mathrm{n}^{\wedge} 2 * \mathrm{MLT} / \mathrm{Wa} / \mathrm{Ku} / \mathrm{al}$ pha1
Rs $=p *(n * N)^{\wedge} 2 * M L T / W a / K u / a l p h a 2$
\%Wire thickness
aw1=alpha1*Ku*Wa/n;
aw2=alpha2*Ku*Wa/ (n*N);
data(2)=aw1;
data (3) =aw 2 ;
\%data (4)=Ilim;
\%data(5)=Ilim*alpha;
\%Frequency Calculations
ton $=\mathrm{L} . *(\mathrm{Ilim}+\mathrm{Ic}) . *(1-\mathrm{alpha}) . /(\mathrm{Vin}-\mathrm{Vsat})$;
toff $=$ L*Ilim*N*(1-alpha)./(Vout+Vdiode)+t_blanking;
duty $=$ ton./(ton+toff);
freq $=1 . /($ ton+toff $)$;
data (4) $=\max (f r e q)$;
\%Modified frequency to use sine wave core loss data
freq_m $=\left(2 /\right.$ pi^2 $\left.^{2}\right) . * f r e q . /(d u t y . *(1-d u t y)) ;$

```
%Core Loss Data Calculations
di = .5*(1-alpha)*Ilim;
dB = di*Al*n/Ae;
P=5.17e-2.*Ve.*dB^2.45.*freq_m.^1.72;
%DC Resistance in transformer
Pdcs = (1-duty).*Ilim^2.*(alpha+((1-alpha)^2)/3)*Rs/N^2;
Pdcp = (duty).*Ilim^2.*(alpha+((1-alpha)^2)/3)*(Rp);
Ip = (Pdcp./Rp).^.5;
Is = (Pdcs./Rs).^.5;
%Losses in switch
Psw = (duty).*Ilim^2.*(alpha+((1-alpha)^2)/3)*(Rsw);
Pf= (Vout./N)*Ilim*tf.*freq./2;
Pr = (Vout./N)*Ilim*alpha*tr.*freq./2;
%Loss from Leakage
Pleak = .5*Ilim^2*Lleak*freq;
%Parasitic Capacitor Loss
Pcloss = .5*C.*(Vout./N).^2.*freq;
%Diode losses
Pdiode = Vdiode*Ilim*(1+alpha).*((Vin-Vsat).*(Vout+Vdiode))
./(2.*Vout.*(Vout+Vdiode+N*(Vin-Vsat)));
%Power from battery
Pin = .5*Ilim.*duty.*(alpha+1)*Vin;
Plosslf = Pdcs+Pdcp;
Plosshf = Pcloss+P+Pf+Pr;
Ploss = Pdcs+Pdcp+P+Pcloss+Pdiode+Pf+Pr+Pleak+Psw;
Eff = (Pin-(Ploss))./Pin;
dv_dt=2.*Cload.*Vout./Ilim.*(1/Vin+N./Vout).*((1-alpha)/(1-alpha^2));
temp = dv_dt.*(Ploss);
intarea = dV*sum(temp);
EnergyC = . 5*Cload*Vmax^2;
eff = (EnergyC)./(EnergyC+intarea);
data(6)=eff;
\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\% \%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\%\% \%ploteff
```

\%Generates Efficiency Sweeps for different Parameters such as alpha $\%$ and magnetizing inductance.
clear L;
clear eff;
clear Pin
cntr $=0$;
\%for alpha = 0:.01:.9,
for $n=0: 1: 30$,
cntr=cntr+1;
$\mathrm{u} 0=1.26 \mathrm{e}-6$;
Al $=40 \mathrm{e}-9$;
$\mathrm{Wa}=3.2 \mathrm{e}-6$;
MLT $=17.5 \mathrm{e}-3$;
$\mathrm{Ve}=167 \mathrm{e}-9$;
$\mathrm{Ae}=9.4 \mathrm{e}-6$;
Bex $=2.45$;
fex = 1.72;
$\% \mathrm{n}=16$;
$\mathrm{N}=10$;
Iin = .6;
Ilim $=1.3$
alpha $=0$;
Cload $=100 \mathrm{e}-6$;
Vin = 3.3;
Vmax $=320$;
leakpercent = .03;
primarywinding $=.7 ; \%$ Winding window allocation for primary wind.
$\mathrm{lg}=\mathrm{u} 0 * \mathrm{Ae} / \mathrm{Al}$;
Bsat $=400 \mathrm{e}-3$;
\%Bsat = Al*n*1.2/Ae;
\%Ilim $=2 *$ Iin/(alpha+1)/.85;
\%Al = Bsat*Ae/n/Ilim
\%n= Bsat*Ae/Al/Ilim
[eff(cntr) $j k($ cntr,: $)]=$ losscalc(Al, Wa, MLT, Ve, Ae, Bex, fex, n, N, Ilim, alpha, Cload, Vin, Vmax, leakpercent, primarywinding);
alphalst(cntr) = alpha;
ncntr(cntr) $=$ n^2*Al;
end
figure(1)
plot(ncntr,eff,'b');

```
%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%
%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%
%ploteff2
%Generates efficiency data for the scaled core.
clear L;
clear eff;
clear Pin
cntr = 0;
for vf = .1:.01:1,
    cntr=cntr+1;
    u0 = 1.26e-6;
    Al = 40e-9;
    Wa = 3.2e-6;
    MLT = 17.5e-3;
    Ve=167e-9;
    Ae=9.4e-6;
    Bex = 2.45;
    fex = 1.72;
    nmax = 20;
    N = 10;
    Iin = .5;
    alpha = 0;
    Cload = 100e-6;
    Vin = 3.3;
    Vmax = 320;
    leakpercent = 0;
    primarywinding = .7;%Winding window allocation for primary wind.
    lg = u0*Ae/Al;
    Bsat = 400e-3;
    Pin = Iin*Vin;
    Vtot = 280.9e-3*vf^3;
    MLTf = vf.*MLT;
    Aef = vf^2*Ae;
    lgf = vf^2*lg;
    Alf = u0*Aef/lgf;
    Waf = vf^2*Wa;
    Vef = vf^3*Ve;
    Ilim = 1.4;
    nsat = Bsat*Aef/Ilim/Alf;
    n = min(nmax,nsat);
    L(cntr) = n^2*Al;
```

```
    Ilimvf(cntr) = 2*Iin/.9+(23)*(2e-9/L(cntr))^.5;
    [eff(cntr) jk(cntr,:)] = losscalc2(Alf, Waf, MLTf, Vef, Aef,
        Bex, fex, n, N, Ilimvf(cntr), alpha, Cload, Vin, Vmax,
        leakpercent, primarywinding,Pin);
    toffvf(cntr) = L(cntr)*Ilim*N*(1-alpha)./(300);
    vflist(cntr) = vf;
    factorvf(cntr) = eff(cntr)^3.*Pin./Vtot;
end
figure(1)
plot(vflist,factorvf);
figure(2)
plot(vflist,eff);
figure(3)
plot(vflist,L);
figure(4)
plot(vflist,Ilimvf);
```


## Appendix B

Board Layout


Figure B-1: Board Layout

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