# HELICOPTER ROTOR BLADE LOADING CALCULATIONS USING AN AXISYMMETRIC VORTEX SHEET AND THE FREE WAKE METHOD 

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# HELICOPTER ROTOR BLADE LOADING CALCULATIONS USING AN AXISYMMETRIC VORTEX SHEET AND THE FREE WAKE METHOD 

by<br>Kevin Huh<br>Submitted to the Department of Aeronautics and Astronautics on January 26, 1988, in partial fulfillment of the requirements of the Degree of Master of Science in Aeronautics and Astronautics


#### Abstract

The wake geometry and the loading characteristics for a hovering helicopter are studied using an axisymmetric vortex sheet and the free wake analysis. The structure of the rolled-up vortex wake is calculated using a filament model as suggested by Kantelis \& Widnall. The model represents the spiral portion of the sheet as an equivalent single vortex with higher order terms. The entire sheet is rediscretized at each time step to minimize errors due to uneven sheet stretching. The rollup produces distinct regions of high vorticity concentrations, which are then replaced with discrete vortices to simplify further calculations.

The lifting line theory is used to predict the rotor bound circulation distribution, and the extended Miller's free wake analysis is applied to the intermediate and the far wakes. Both the load distributions and the tip vortex locations show good agreements with experimental data. It appears using three vortices to represent amalgamated trailing vorticies in the intermediate wake is adequate.


Thesis Supervisor: Sheila E. Widnall
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# NOMENCLATURE 

## Roman Letters

| $a$ | Pullin's scaling constant |
| :--- | :--- |
| $b_{n}$ | Pullin's inner spiral constant |
| $c$ | Core diameter |
| $c$ | Chord lenth |
| $C_{l}$ | Section lift coefficient |
| $d_{n}$ | Pullin's inner spiral constant |
| $E$ | Elliptic integral of the second kind |
| $F$ | Elliptic integral constant |
| $f_{i}$ | Interpolated function value for rediscretization |
| $g_{i}$ | Rediscretization parabola |
| $K$ | Elliptic integral of the first kind |
| $K_{1}$ | Constant defined in L-L theory |
| $K_{2}$ | Constant defined in L-L theory |
| $k_{0}$ | Runge-Kutta constant |
| $k_{1}$ | Runge-Kutta constant |
| $k_{2}$ | Runge-Kutta constant |
| $k_{3}$ | Runge-Kutta constant |
| $l_{0}$ | Runge-Kutta constant |
| $l_{1}$ | Runge-Kutta constant |


| $l_{2}$ | Runge-Kutta constant |
| :---: | :---: |
| $l_{3}$ | Runge-Kutta constant |
| $L$ | Lift |
| $m_{0}$ | Lift-curve slope |
| $n$ | Number of intermediate wakes vorticies |
| $N$ | Number of near wake vorticies |
| $P$ | Impulse |
| $r$ | Radius, radial direction |
| $r_{1}$ | Least distance from ( $r, z$ ) to the ring |
| $r_{2}$ | Greatest distance from (r,z) to the ring |
| $r_{i}$ | Radius of vortex $\boldsymbol{i}$ |
| $r_{i}^{\text {old }}$ | Previous radius of vortex $i$ |
| $r_{i}^{\text {new }}$ | Current radius of vortex i |
| $\tilde{\boldsymbol{r}}$ | Radius of vortex inside spirali |
| $\vec{r}$ | Position vector |
| $\dot{r}$ | Total velocity in the radial direction |
| $R$ | Radius of rotor |
| $s$ | Arc-length parameter |
| $S_{i}$ | Shearing velocity at ring i |
| $T$ | Kinetic energy |
| $\vec{s}$ | Position vector |
| $\vec{t}_{i}$ | Tangent vector |
| $\vec{u}$ | Velocity vector |
| $u_{r}$ | Velocity in the radial direction |
| $u_{\boldsymbol{\phi}}$ | Velocity in the azimuthal direction |
| $u_{x}$ | Velocity in the axial direction |
| $u_{z \text { sind }}$ | Self-induced component of velocity |
| $\boldsymbol{u}_{\boldsymbol{x}}$ | Velocity in the $\tilde{x}$ direction |


| $u_{\tilde{y}}$ | Velocity in the $\tilde{y}$ direction |
| :--- | :--- |
| $u_{z \text { sind }}$ | Self-induced component of velocity |
| $U_{m a x}$ | Max velocity |
| $\vec{V}$ | Velocity vector |
| $w$ | Downwash |
| $w_{w a k e}$ | Downwash due to the wake |
| $w_{t v}$ | Downwash due to trailing vortex sheet |
| $x_{m a x}$ | Maximum translation |
| $\vec{x}$ | Position vector |
| $\vec{x}^{\prime}$ | Position vector used for Biot-Savart law |
| $\tilde{x}$ | Local horizontal axis for the spiral |
| $\tilde{y}$ | Local verticle axis for the spiral |
| $\boldsymbol{z}$ | Axial direction |
| $z_{i}$ | Axial position of ring $\boldsymbol{i}$ |
| $z_{i}^{o l d}$ | Previous axial position of vortex $\boldsymbol{i}$ |
| $z_{i}^{n e w}$ | Current axial position of $\boldsymbol{i}$ |
| $\dot{\boldsymbol{z}}$ | Total velocity in the axial direction |

## Greek Letters

| $\alpha$ | Angle of attack |
| :--- | :--- |
| $\alpha_{\text {geomt }}$ | Geometric angle of attack |
| $\alpha_{i n d}$ | Induced angle of attack |
| $\alpha_{\text {rot }}$ | Rotation angle for spiral modelling |
| $\boldsymbol{\gamma}_{\boldsymbol{i}}$ | Vorticity of vortex $\boldsymbol{i}$ |
| $\Gamma$ | Circulation |

$\Gamma_{i} \quad$ Circulation of vortex $i$
$\psi \quad$ Variable of integration for elliptic integrals
$\Psi \quad$ Streamfunction
$\omega \quad$ Underrexation parameter, stream funtion constant
$\Omega \quad$ Blade rotation rate
$\Theta_{\max } \quad$ Maximum spiral angle
$\boldsymbol{\theta}_{\boldsymbol{s}}$ Rotation angle defined by Kantelis
$\tilde{\theta} \quad$ Rotation angle for coordinate transformation
$\eta \quad$ Vortex ring calculation constant, radial direction
$\phi \quad$ Blade passing angle
$\tau \quad$ Pullin's inner spiral constant
$\tau_{a} \quad \tau$ at A, constant position on the sheet
$\rho \quad$ Fluid density
$\sigma \quad$ Solidity
A Initial portion to be lumped to tip

## Chapter 1

## Introduction

### 1.1 Brief History

An accurate prediction methodology for the aerodynamic flow over a lifting rotor has been a continuous problem to engineers for over forty years. The main difficulty lies in the complexity of the flow field about a rotating airfoil. In particular, due to the close interaction of the blade and the tip vortex of the previous blade, any methodology must be able to accurately predict the rollup schedule of the trailing vortex sheet. (Figure 1.1) A precursor to such a condition is the necessity of having an accurate wake model, since a trailing vortex sheet of a rotor convects within its own wake.

The early methods of prediction involved vortex methods originally developed for propellers and adapted to lifting rotors. These classical theories assumed that the wake may be treated as a rigid, non-contracting helical vortex cylinder. While this assumption is consistent for high speed propellers for which the induced velocities are small with respect to the axial translational speed, it is clearly wrong for a hovering helicopter for which the only velocities are the induced velocities. The momentum theory predicts that the induced velocity on the rotor blade is half of the velocity in the fully developed wake, and therefore the wake must contract.[14] Hence the fixed non-contracting wake theory is inapplicable to an analysis of a hovering rotor, and a more accurate description
of the vortex wake is required. There exists currently two distinct methods of modelling the helicopter vortex wake.

One such method is called the prescribed wake method [13]. These methods involve empirical formulas involving blade geometry and the thrust coefficients to obtain vortex wake geometries. They have proven to be quite efficient and accurate for conventional rotor configurations. However, the method's dependence on empiricism means that given an unconventional blade geometry, the results would become unreliable. The scheme's greatest fallacy is its inability to model the physics of vorticity transport.

A second scheme which is applicable given a general geometry is called the Free Wake Method. [20,14,15,21,19] These schemes are all based on the fact that vorticies move with the fluid. The wake is modelled as vortex sheets or filaments and allowed to convect in space according to the Biot-Savart relationship in a Lagrangian fashion until geometric and loading convergence is achieved. Although the free wake method is completely general, they are also computationally quite expensive owing to the large number of calculations required for vortex transportation. In addition, the question of numerical stability of these methods in three-dimensional flow has not been adequately addressed. The apparent convergences of current three-dimensional free-wake calculations are likely due to the coarse discretizations used because of limitations of current computers. There is a need for a test case which does not suffer from such three-dimensional instabilities, and thus can be used to study their convergence.

A simpler, yet physically viable approach to the rotor wake problem is the simplified free wake theory first proposed by Miller [14]. This scheme reduces
the number of vortex velocity calculations required by amalgamating the vortex sheet into two or three axisymmetric vorticies after $1 / 4$ blade rotation. The amalgamation process is based on the Donaldson et al. [5] method first proposed for fixed wing aircrafts. Simply stated, Donaldson's rollup criteria for a complex wake states that separate vortices will be formed between the local minima of the curve $|d \Gamma / d r|$ vs $r$, and thus the tip vortex will be of maximum strength containing all of the negative vorticity in the region.[5] In addition, Miller simplified the vorticity representation in the far field by replacing the sequence of far-wake finite vorticies with distributed axisymmetric vortex sheets or cylinders. The calculations obtained seem to show good agreements with empirical data and required only a fraction of computations required of the full free wake technique.

Roberts \& Murman [19] extended Miller's work by coupling the simplified free wake technique with an Euler solver in the near field of the rotor. This work offered the advantages of using an Euler scheme as opposed to potential methods, namely permitting distributed vorticity. This allows one to gain insights into the structure of the vortex core. In addition, Roberts \& Murman were able to avoid excessive number of computations and vorticity dissipation by implementing a novel perturbation scheme. The calculations obtained showed good agreement with empirical data. However, Roberts used the same rollup model proposed earlier by Donaldson et al. [5], and questions remain as to its validity for hovering rotors.

### 1.2 Motivation for Study

The analysis included in the current report expands on the simplified free wake method to include the near wake rollup model for an axisymmetric vortex
sheet proposed by Kantelis \& Widnall. [11] Previous works of rotor in hover have essentially ommitted a detailed analysis of the structure of rolled-up vorticity in the wake. Hence, the objectives of the study are:

- To formulate an accurate prediction method for the rotor blade bound circulation distribution and the tip vortex geometries.
- To offer physical insights of the transportation of vorticity in the wake, in particular, the amalgamation of inboard vorticity.
- To assess the accuracy of the rollup criteria proposed by Donaldson et al. [5] as used in other free wake techniques.


Figure 1.1: Blade-Vortex interaction

## Chapter 2

## Theoretical Development

### 2.1 Physical Model

The physical model in study here is one of an isolated rotor blade in hover position as illustrated in Figure 2.1; its associated wake is shown in Figure 2.2. The fluid is assumed to be steady, incompressible, and inviscid. There are two important consequences of these physical assumptions:

- Inviscidness and incompressibility of the fluid implies vorticity associated with a particular fluid element must remain constant.
- A rotor in hover implies no mean flow in the axial or radial directions, and hence all the perturbation velocities in the medium is generated by the vorticity in the fluid.


### 2.2 Wake Modelling

### 2.2.1 Miller's Simplified Free Wake Technique

A model which embodies the above two physical conditions is the simplified free wake technique, first used by Miller [14] and referred to earlier in the introduction. A detailed descriptions are available in other refrences [14,15,19,21],
and only a brief summary will be presented here.

The bound circulation is computed using lifting-line theory. The trailing vorticity from the blade is then transported into the wake using discrete vortex filaments according to Helmholtz's vortex theorems for a steady, inviscid flow; vortex lines must follow streamlines in the flow.[19] The vortex trajectories, therefore, must obey the following equations:

$$
\begin{equation*}
\frac{u_{r i}}{d r_{i}}=\frac{u_{\phi i}}{r_{i} d \phi}=\frac{u_{z i}}{d z_{i}} \tag{2.1}
\end{equation*}
$$

where $u_{r}, u_{\phi}$, and $u_{z}$ are the radial, azimuthal, and axial velocity components in the frame of reference of the rotor (Figure( 2.3)), and $i$ is the filament index. The essence of the free wake analysis is the solution of the Equations 2.1 for the $N$ discrete vorticies, where $N$ represents the total number of vortex filaments in the wake. Several approximations are made to simplify the solution procedure. First, the contribution of the rotor blade bound circulation to the induced velocity in the wake is neglected; the contributions of each blade cancel when averaged over the azimuth. [19] Next, the inclinations of the spiral wake are ignored. Miller [14] showed that the change in inclination angle only refines the downwash on the blade to the second-order. By ignoring the inclination effect of the wake, one may ignore the perturbations in the azimuthal direction, thus $u_{\phi}$ becomes,

$$
\begin{equation*}
u_{\phi i}=r_{i} \Omega \tag{2.2}
\end{equation*}
$$

where $\Omega$ is the angular speed of the rotor. Substituting into Equations yields,

$$
\begin{equation*}
d r_{i}=\frac{u_{r i}}{\Omega} d \phi \tag{2.3}
\end{equation*}
$$

$$
\begin{equation*}
d z_{i}=\frac{u_{z i}}{\Omega} d \phi \tag{2.4}
\end{equation*}
$$

Neglecting the azimuthal velocity perturbations is equivalent to treating the flow in a given azimuthal plane to be axisymmetric, and therefore, the helical vortex filaments may be replaced by vortex rings. A further simplification is made by computing the velocities and solving for the position of the vorticies only in the azimuthal plane containing the rotor blade, $\phi=0$. Each vortex ring is constructed by taking $\pi$ segments of the helical filaments from each blade and replacing them with vortex rings at their mean locations, as shown in Figure( 2.5).

The total number of vorticies is reduced by dividing the wake into three distinct sections: Near wake attached to the blade, an intermediate wake, and a far wake as shown in Figure( 2.4).

The near wake represents vorticity transported between blade passing time of $\phi=0$ to $\phi=\frac{\pi}{2}$ and is approximated by a series of straight semi-infinite vortex filaments extending from the blade nodes.

It is assumed that the near wake rolls-up almost immediately according to Betz's theory of conservation of momentum into vortex rings. [4] The rollup criterion is based on the Donaldson et al. [5] approximation of dividing the vortex sheet along the local minima of the curve $|d \Gamma / d r|$ vs $r$. Hence, a tip vortex will form from the amalgated trailing vortex sheet starting at the tip to the maximum circulation position. A second vortex will form from this position to where $\frac{d \Gamma}{d r}$ again changes abruptly. The remaining circulation to the root is rolled-up into a third vortex.

The intermediate wake consists of four sets of vortex rings. Each ring is formed by taking $\pi$ segments of the helical filaments as explained above. The first represents vorticity transported from blade passing time $\phi=\frac{\pi}{2}$ to $\phi=\frac{3 \pi}{2}$, the second from time $\phi=\frac{3 \pi}{2}$ to $\phi=\frac{5 \pi}{2}$, the third from time $\phi=\frac{5 \pi}{2}$ to $\phi=\frac{7 \pi}{2}$, and the fourth from time $\phi=\frac{7 \pi}{2}$ to $\phi=\frac{9 \pi}{2}$. (Figure 2.5) The radii and vertical spacings are determined from the velocities induced by the wake, averaged over the spacing between vorticies.

To transport the vorticity from blade passing time $\phi=\frac{9 \pi}{2}$ out to $\phi=$ $\infty$, as per required for steady state hover conditions, two semi-infinite vortex cylinders are placed one vortex spacing below the last set of intermediate wake vorticies. The outboard cylinder is placed directly underneath the tip vortex, and its strength, $\frac{d \Gamma}{d z}$, is determined by dividing tip vortex circulation $\Gamma$ by the spacing between the last two intermediate wake tip vorticies. To determine the location and strength of the inner vortex cylinder, first the centroids of the inner portion from the last two intermediate vortex sheets are computed, then the determination is completed as per tip vortex cylinder.

Since the problem is nonlinear, an iterative approach is used for its solution. At each iteration step the velocities induced at the intermediate wake vortex rings are computed from all the vorticity in the wake, including the rolled-up near wake. A new wake geometry followed by a new bound circulation of the blade is then computed. The new loading now determines the strengths of the vorticies in the wake, and the process is repeated until a suitable convergence criterion is satisfied for the bound circulation distribution.

### 2.2.2 Extended Simplified Free Wake Theory

The analysis given here expands on the ideas given by Miller by adding more detail to the near wake. As explained in the introduction, to correctly predict the blade loading distribution, it is essential that the analysis must accurately forecast the position and strength of the tip vortex at the first pass. Previous free wake methods relied upon Donaldson's [5] rollup model for the tip vortex determination, and it is felt that this rollup model may not be accurate for rotor wakes. Hence, a model which more closely resembles the fluid physics at the near wake seems necessary. The near wake in the current method is refined by direct calculation of the vortex wake rollup during the interval to the first blade passage ,and by more rigorous adhesion to vortex physics as explained by Betz[4]; namely conservation of circulation, conservation of impulse, and conservation of kinetic energy.

The near wake begins attached to the rotor blade and consists of $n$ number of vortex rings, where $\boldsymbol{n}$ is typically $50-100$. As in Miller's method, the attached near wake represents vorticity transport from time $\phi=0$ to $\phi=\pi / 2$. At the end of the time period, $\phi=\pi / 2$, the near wake vorticies are allowed to convect freely in space until time $\phi=3 \pi / 2$ at which time a new wake is formed. It is assumed that a new vortex ring is formed at each $\pi$ rotation (for a two-bladed rotor), and hence the near wake represents vorticity transport from blade passing time $\phi=0$ to $\phi=3 \pi / 2$.(Figure 2.6) At the end of the transport, the vortex sheet is split into computationally determined number of intermediate wake vorticies according to a shear velocity criterion to be discussed.

As in Miller's method there are four levels of intermediate wakes. Thus the first level represents time $\phi=3 \pi / 2$ to $\phi=5 \pi / 2$, the second from $\phi=5 \pi / 2$ to
$\phi=7 \pi / 2$ and so on. The locations of the intermediate wake vortex rings are determined from velocity calculations of all the vorticies in the wake, and no averaging is performed.

The far wake position and strength are computed as in Miller's technique.

The problem still requires an iteration process. The positions of each vortex rings in the near and intermediate wakes are integrated in a Runge-Kutta fashion from blade passing time of $\phi=\pi / 2$ to $\phi=3 \pi / 2$; at which point a new bound circulation distribution, and a new vortex sheet would be formed on the rotor blade. The process is repeated until a convergence criterion is satisfied in the wake geometry and the bound circulation distribution. The iteration process will be discussed in more detail in the following chapter concerning numerical procedures.

### 2.2.3 Velocity Field

The vortex filament trajectories are computed by integrating the Equations 2.3 and 2.4. The perturbation velocities at the vortex locations are induced by all the discrete vorticies in the wake, and since vortex rings are curved this incldes the control vortex itself. The Biot-Savart law expresses the velocity field $\vec{V}(\vec{x})$ induced by a vortex line $\vec{\Gamma}^{\prime}$ as, [3]

$$
\begin{equation*}
\vec{V}(\vec{x})=-\frac{1}{4 \pi} \int_{\sigma} \frac{\vec{s} \times \vec{\Gamma}^{\prime}}{s^{3}} d A\left(\vec{x}^{\prime}\right) \tag{2.5}
\end{equation*}
$$

where $\vec{s}=\vec{x}-\vec{x}^{\prime}, \vec{\Gamma}^{\prime}$ is the strength of the sheet at the point $\vec{x}$, and $\sigma$ extends over the entire sheet.

For the special case of axisymmetric vortex rings, Equation 2.5 is greately simplified. Given a vortex ring of radius $r$ and circulation $\Gamma$, the radial and axial components of the induced velocity at a radius $\eta$ and axial distance $z$ from the plane of the ring (Figure 2.7) are given by Lamb [12] to be:

$$
\begin{align*}
& u_{r}=\frac{\Gamma}{4 \pi} \int_{0}^{2 \pi} \frac{z r \cos \psi d \psi}{\left(\eta^{2}+r^{2}+z^{2}-2 r \eta \cos \psi\right)^{\frac{3}{2}}}  \tag{2.6}\\
& u_{z}=\frac{\Gamma}{4 \pi} \int_{0}^{2 \pi} \frac{r(r-\eta \cos \psi) d \psi}{\left(\eta^{2}+r^{2}+z^{2}-2 r \eta \cos \psi\right)^{\frac{3}{2}}} \tag{2.7}
\end{align*}
$$

The Equations 2.6 and 2.7 may further be simplified by expressing them in terms of complete elliptic integrals of the first and second kind,

$$
\begin{gather*}
K=\int_{0}^{\frac{\pi}{2}} \frac{d \psi}{\left(1-k^{2} \sin ^{2} \psi\right)^{\frac{1}{2}}} d \psi  \tag{2.8}\\
E=\int_{0}^{\frac{\pi}{2}}\left(1-k^{2} \sin ^{2} \psi\right)^{\frac{1}{2}} d \psi \tag{2.9}
\end{gather*}
$$

where

$$
k^{2}=\frac{4 r \eta}{(r+\eta)^{2}+z^{2}}
$$

is the argument of the elliptic integrals $K$ and $E$. In terms of the elliptic integrals, the velocity components $u_{r}$ and $u_{z}$ may be written in the following form:

$$
\begin{equation*}
u_{r}=\frac{\Gamma}{4 \pi} \frac{z}{2 \eta} \sqrt{\frac{k^{2}}{r \eta}}\left\{E \frac{2-k^{2}}{1-k^{2}}-2 K\right\} \tag{2.10}
\end{equation*}
$$

$$
\begin{equation*}
u_{z}=\frac{\Gamma}{4 \pi} \sqrt{\frac{k^{2}}{r \eta}}\left\{K-E\left[\frac{1-\frac{k^{2}}{2}\left(1+\frac{r}{\eta}\right)}{\left(1-k^{2}\right)}\right]\right\} . \tag{2.11}
\end{equation*}
$$

In addition, the elliptic integrals $K$ and $E$ may be expressed in a series form.[19] These series are rapidly convergent, and only the first four terms are kept for computational purposes.

$$
\begin{align*}
K= & F+\frac{1}{2}^{2}(F-1)\left(1-k^{2}\right)+\left(\frac{1 \cdot 3}{2 \cdot 4}\right)^{2}\left(F-1-\frac{2}{3 \cdot 4}\right)\left(1-k^{2}\right)^{2}+ \\
& \left(\frac{1 \cdot 3 \cdot 5}{2 \cdot 4 \cdot 6 \cdot}\right)^{2}\left(F-1-\frac{2}{3 \cdot 4}-\frac{2}{5 \cdot 6}\right)\left(1-k^{2}\right)^{3}+\cdots  \tag{2.12}\\
E= & 1+\frac{1}{2}\left(F-\frac{1}{2}\right)\left(1-k^{2}\right)+\left(\frac{1}{2}\right)^{2} \frac{3}{4}\left(F-\frac{2}{2}-\frac{1}{3 \cdot 4}\right)\left(1-k^{2}\right)^{2}+  \tag{2.13}\\
& \left(\frac{1.3}{2 \cdot 4}\right)^{2} \frac{5}{6}\left(F-\frac{2}{2}-\frac{2}{3 \cdot 4}-\frac{1}{5 \cdot 6}\right)\left(1-k^{2}\right)^{3}+\cdots
\end{align*}
$$

where

$$
F=\ln \frac{4}{\sqrt{1-k^{2}}}
$$

The velocity induced by the far wake votex cylinders are given by:

$$
\begin{align*}
& u_{r c y l}=\frac{1}{4 \pi} \frac{d \Gamma}{d z} \int_{z}^{\infty} \int_{0}^{2 \pi} \frac{z r \cos \psi d \psi d z}{\left(\eta^{2}+r^{2}+z^{2}-2 z \eta \cos \psi\right)^{\frac{3}{2}}}  \tag{2.14}\\
& u_{z c y l}=\frac{1}{4 \pi} \frac{d \Gamma}{d z} \int_{z}^{\infty} \int_{0}^{2 \pi} \frac{r(r-\eta \cos \psi) d \psi d z}{\left(\eta^{2}+r^{2}+z^{2}-2 z \eta \cos \psi\right)^{\frac{3}{2}}} \tag{2.15}
\end{align*}
$$

Here, $\frac{d \Gamma}{d z}$ is the vortex cylinder strength, $z$ is the axial distance from the end of the cylinder to the point at which the velocity is required, $r$ is the cylinder radius, and $\eta$ is the radius at the point of evaluation. By integrating the above equations over $z$ we obtain,

$$
\begin{gather*}
u_{r c y l}=\frac{1}{4 \pi} \frac{d \Gamma}{d z} \int_{0}^{2 \pi} \frac{r \cos \psi d \psi}{\left(\eta^{2}+r^{2}+z^{2}-2 r \eta \cos \psi\right)^{\frac{1}{2}}}  \tag{2.16}\\
u_{z c y l}=\frac{1}{4 \pi} \frac{d \Gamma}{d z} \int_{0}^{2 \pi} \frac{r(r-\eta \cos \psi)}{\eta^{2}+r^{2}-2 r \eta \cos \psi}\left[1-\frac{z}{\left(\eta^{2}+r^{2}+z^{2}-2 r \eta \cos \psi\right)^{\frac{1}{2}}}\right] d \psi \tag{2.17}
\end{gather*}
$$

The Equation 2.16 can be rewritten in terms of the complete integrals of the first and the second kind,

$$
\begin{equation*}
u_{r c y l}=\frac{1}{2 \pi k} \frac{d \Gamma}{d z} \sqrt{\frac{r}{\eta}}\left[K\left(2-k^{2}\right)-2 E\right] \tag{2.18}
\end{equation*}
$$

where $k, K$, and $E$ are as defined previously. The equation for the axial component of velocity due to a vortex cylinder cannot be expressed solely in terms of the complete integrals of the first and second kind, but contains elliptic integrals of the third kind. Therefore, it is more convenient to evaluate the velocity by numerical integration of Equation 2.17 using the trapezoidal rule.[19] The approach used here is to use sixty integration steps in the $\psi$ direction. As Roberts et al. points out, care must be taken when $\eta=r$, since the denominators of both fractions in the integrand vanish at the endpoints of the integration, $\psi=0$ and $\psi=2 \pi$. However, the integrand has a well defined limit of zero as $\psi \rightarrow 0,2 \pi$. The integrand is ignored at the endpoints when $r=\psi$, and the expression is numerically evaluated for each of the remaining steps.

Since a curved vortex has a velocity influence upon itself, the vortex rings must have a self-induced velocity component. From Lamb [12], a vortex ring of strength $\Gamma$, radius $r$, and a core radius of $c$ has a self-induced velocity of,

$$
\begin{equation*}
u_{z \operatorname{sind}}=\frac{\Gamma}{4 \pi r}\left(\ln \frac{8 r}{c}-\frac{1}{4}\right) \tag{2.19}
\end{equation*}
$$

The core radius $c$ is computed from the rolled-up near wake by conserving the kinetic energy of the vortex sheet; the details of the process will be discussed in the next section.

Using the above formulas for the induced velocities of each vortex ring and cylinder in the wake, the total induced velocity at each vortex location is determined. These velocities are integrated to solve for the new vortex locations at time $\phi=\phi+\Delta \phi$.

### 2.2.4 Near Wake Model

Kantelis \& Widnall [11] discretized an isolated axisymmetric vortex sheet and tracked its rollup schedule. They showed that a simple straightforward application of the vortex filament model yields a chaotic result. As in previous two-dimensional cases, for example Fink \& Soh [6,7], the vortex rings begin to rollup as expected, but quickly accumulate large errors as the rings begin to cross over each other. A typical example of this is shown in Figure (2.8) modelled by 20 rings. Kantelis et al. also showed that a simple increase in the number of vortex rings does not in itself alleviate the problem.

Kantelis et al. suggests that the cause of the errors are two-fold:

- The inablity of a finite number of discrete vorticies to model the singularity at the sheet edge.
- As the sheet stretches unevenly, the initially equally spaced vortex elements become unevenly spaced which can introduce large errors in the velocity calculations. This is to be expected since Fink \& Soh [6,7] ) re-
ported similiar errors for two-dimensinal vortex sheets. They suggested that the errors are due to the neglect of a logarithmic term in the solution of the Cauchy principle value integral representing the velocity. When the vortex is centered in the segment, the logarithmic term vanishes. As the vortex approaches the edges of the segment, however, the contribution of the logarithmic term becomes large and can not be neglected.

Hence, to model the time history of the near wake rollup correctly, a different method which reduces the above two errors must be introduced. A method which reduces the severity of the errors at the sheet edge in a discrete manner has been suggested by, among others, Hoeijmakers \& Vaastra [8] and by Kantelis \& Widnall [11]. The method replaces the inner portion of the spiral with a single tip vortex ring such that the total circulation and centroid of the system is preserved; amalgamation criterion used by Kantelis is to limit the spiral portion of the sheet to a fixed number of loops. (Figure( 2.9)) In other words, if a vortex ring spiraled past the threshold loop number, it would be amalgamated into the tip vortex. The new amalgamated vortex strength is then computed by summing the two vorticies' circulation, and the location of the tip vortex is computed by conserving total impulse, i.e. the total impulse before amalgamation must equal the total impulse after amalgamation. For an axisymmetric vortex ring system the total impulse is given by Lamb [12] to be,

$$
\begin{equation*}
\frac{P}{\rho}=\pi \sum_{i=1}^{N} \Gamma_{i} r_{i}^{2} \tag{2.20}
\end{equation*}
$$

where $P$ is the impulse, $\rho$ is the density , $N$ is the number of vorticies in the near wake, $\Gamma_{i}$ is the circulation, and $r_{i}$ is the radius of the vortex ring.

Similiarly, the core size of the newly amalgamated vortex is computed by conserving the kinetic energy of the system. Lamb [12] gives the the total kinetic energy of the flow to be,

$$
\begin{equation*}
T=-\pi \rho \int_{0}^{\infty} \int_{-\infty}^{\infty} \Psi \gamma d z d r \tag{2.21}
\end{equation*}
$$

where $\gamma$ is the vorticity and $\Psi$ is the streamfunction. The streamfunction $\Psi$ is again given by Lamb [12] to be,

$$
\begin{equation*}
\Psi(r, z)=-\frac{\Gamma}{2 \pi}\left(r_{1}+r_{2}\right)[K(\omega)-E(\omega)] \tag{2.22}
\end{equation*}
$$

where:
$r_{1}=$ least distance from $(r, z)$ to the ring
$r_{2}=$ greatest distance from $(r, z)$ to the ring
$\omega=\left(r_{1}-r_{2}\right) /\left(r_{1}+r_{2}\right)$
$K=$ Complete elliptic integral of the first kind
$E=$ Complete elliptic integral of the second kind
$\boldsymbol{\Psi}, \Gamma$ as defined previously

Since the flow is irrotational everywhere except in the interior of the vortex rings, the energy can be expressed as a summation over all the rings

$$
\begin{equation*}
T=-\pi \rho \sum_{i=1}^{N} \iint \Psi_{i} \gamma_{i} d z d r \tag{2.23}
\end{equation*}
$$

where the integration is performed over the filament core sections. For each core section $i$, the streamfunction is written as the sum of the contributions of all of
the vortex rings including that of ring $i$ itself.

$$
\begin{equation*}
\Psi_{i}=\sum_{j=1}^{N} \Psi_{i, j} \tag{2.24}
\end{equation*}
$$

where $\Psi_{i, j}$ is the streamfunction at ring $i$ due to ring $j$. Then the kinetic energy becomes,

$$
\begin{equation*}
T=-\pi \rho \sum_{i=1}^{N}\left[\iint \Psi_{i, i} \gamma_{i} d z d r+\iint \sum_{j=1, j \neq i} N \Psi_{i, j} \gamma_{i} d z d r\right] \tag{2.25}
\end{equation*}
$$

where $\Psi_{i, i}$ is the streamfunction at ring $i$ due to itself. The first term in the brackets is given by Lamb as,

$$
\begin{equation*}
\iint \Psi_{i, i} \gamma_{i} d z d r=-\frac{k_{i}^{2} R_{i}}{2 \pi}\left[\ln \frac{8 R_{i}}{c_{i}}-\frac{7}{4}\right] \tag{2.26}
\end{equation*}
$$

where $R_{i}$ is the radius and $c_{i}$ is the core size of the vortex ring. The second term is integrated numerically assuming that $\Psi_{i, j}$, given by Equation 2.22, is constant over the core section of the ring.(Figure( 2.10))

In addition to the tip amalgamation, extra higher-order velocity terms are required for both within and without the spiral since the spiral is asymmetric. Kantelis computes the asymmetric terms of the spiral using Pullin's [18] numerical integration of Kaden's [10] similarity solution for a parabolically loaded two-dimensional semi-infinite sheet. The assumption here is that a typical helicopter load corresponds closely to a parabolic load at the sheet edge, and based on previous calculations this assumption appears to be a reasonable one. For a detailed explanation of the higher-order terms consult Kantelis \& Widnall [11] ; only the final result will be given here. The velocity at the tip vortex ring due to the higher-order asymmetric terms are,

$$
\begin{equation*}
u_{\tilde{x}}=-0.000332 \tau_{A}^{-2.383} \tag{2.27}
\end{equation*}
$$

$$
\begin{equation*}
u_{\tilde{y}}=-0.004885 \tau_{A}^{-1.543} \tag{2.28}
\end{equation*}
$$

where ( $\tilde{x}, \tilde{y}$ ) is the rotating coordinate system fixed to the spiral center. $\tau_{A}$ corresponds to $\tau=a^{4} t / \Gamma^{3}$ at position $A$; where $A$ is the reference starting position of the spiral, and $a$ is an arbitrary scaling factor used by Pullin[18]. (Figure(2.11)) To obtain the velocity components with respect to ( $r, z$ ) rather than ( $\tilde{x}, \tilde{y})$ a coordinate transformation is performed such that,

$$
\begin{align*}
& u_{r}=\left(u_{\tilde{x}} \cos \alpha_{r o t}-u_{\tilde{y}} \sin \alpha_{r o t}\right)  \tag{2.29}\\
& u_{z}=\left(u_{\tilde{x}} \sin \alpha_{r o t}+u_{\tilde{y}} \cos \alpha_{r o t}\right) \tag{2.30}
\end{align*}
$$

The angle $\alpha_{\text {rot }}$ is dependent on time and is the angle through which the ( $\left.\tilde{x}, \tilde{y}\right)$ coordinate system must be rotated in order to match the point $A$ with the actual sheet edge. (Figure( 2.12) ) The sheet edge position is determined during the rollup calculations.

The velocity at the remaining portion of the near wake due to the asymmetry in the spiral rollup is,

$$
\begin{gather*}
u_{\bar{\theta}}=-\frac{\partial \Psi}{\partial \tilde{r}}+\frac{\Gamma}{2 \pi \tilde{r}}  \tag{2.31}\\
u_{\tilde{r}}=\frac{1}{r} \frac{\partial \Psi}{\partial \tilde{\theta}} \tag{2.32}
\end{gather*}
$$

where $(\tilde{r}, \tilde{\theta})$ is the corrdinate system fixed to the center of the spiral as shown in Figure( 2.12). The asymmetric stream function, $\Psi$, is computed numerically,

$$
\begin{equation*}
\Psi=\sum_{n=1}^{\infty}\left[b_{n} \tau^{m_{n}} \frac{\sin n \tilde{\theta}}{r^{n}}+d_{n} \tau^{l_{n}} \frac{\cos n \tilde{\theta}}{r^{n}}\right] \tag{2.33}
\end{equation*}
$$

Computationally, only the first seven terms in the series are used and their values are,

| n | $b_{n}$ | $m_{n}$ | $d_{n}$ | $l_{n}$ |
| :---: | :---: | :---: | :---: | :---: |
| 1 | 0.005301 | -1.386 | 0.0010070 | -2.373 |
| 2 | 0.001310 | -1.363 | 0.0001550 | -2.282 |
| 3 | 0.000575 | -1.345 | 0.0000657 | -2.409 |
| 4 | 0.000319 | -1.330 | 0.0000312 | -2.361 |
| 5 | 0.000202 | -1.316 | 0.0000157 | -2.173 |
| 6 | 0.000139 | -1.301 | 0.0000102 | -2.221 |
| 7 | 0.000100 | -1.283 | 0.0000115 | -2.807 |

The velocity components in the usual ( $r, z$ ) system are again computed by coordinate transformation

$$
\begin{align*}
& u_{r}=\left(u_{\tilde{r}} \cos \theta_{s}-u_{\tilde{\theta}} \sin \theta_{s}\right)  \tag{2.34}\\
& u_{z}=\left(u_{\tilde{r}} \sin \theta_{s}+u_{\vec{\theta}} \cos \theta_{s}\right) \tag{2.35}
\end{align*}
$$

where $\theta_{s}=\tilde{\theta}-\pi / 2+\alpha_{\text {rot }}$ and is taken positive in the usual sense with $\theta_{s}=0$ on the positive $r$ axis.

The Pullin's numerical solution was also used as an initial condition such that the numerical simulation may be started at some time, $\phi=\pi / 2+\phi_{0}$, where $\phi_{0}>0$. The outer portion of the sheet will already have rolled up into a logarithmic spiral which can be approximated using the Pullin model described earlier in the section. This approach eliminates the problem of attempting to model the initial motion of the tip singularity in a discrete manner with a series of vortex elements. [11]

The logarithmic spiral of the sheet edge at time $\phi=\pi / 2+\phi_{0}$ is replaced by a finite tip vortex at position $(r, z)=\left(a_{0}, b_{0}\right)$ with its associated circulation and higher order terms. The position vector ( $a_{o}, b_{o}$ ) is,

$$
\begin{aligned}
& a_{o}=0.308\left(a \phi_{0}\right)^{\frac{2}{3}} \\
& b_{0}=0.489\left(a \phi_{o}\right)^{\frac{2}{3}}
\end{aligned}
$$

where $a$ is the Pullin scaling factor. The strength of the initial tip vortex is determined as,

$$
\Gamma_{t i p o}=-\Gamma_{r=1-a_{0}}(\phi=\pi / 2)
$$

The core size and the higher order terms are computed as described earlier. For simplicity the entire sheet, except the tip, is assumed to be flat.(Figure( 2.13))

The large distortions in the vortex ring motion caused by the uneven sheet stretching is reduced by rediscretizing the near wake vortex sheet at every time step to ensure that the vortex elements remain equally spaced. Fink \& Soh $[6,7]$ have used this rediscretizing technique for a two-dimensional sheet with good results. The mechanics of such a procedure is as follows: (Ref.[11])

- After each time step a continuous model of the near wake is reconstructed by passing a smooth curve through all of the vortex elements.
- A functional representation of the circulation distribution along the sheet is obtained.
- The total arc length for the sheet is computed using the cubic spline technique and then divided by the number of vortex elements to yield a new
sheet segment length. The vortex elements are now placed at the centers of these segments.
- The vortex rings' strengths are determined from the cubic spline technique.

The cubic spline technique used here tends to bound the interpolated curve and reduces the appearance of wiggles seen in other types of splines.

The geometry and circulation of each portion of the sheet between any two vortex rings are described parametrically by an arc length parameter. For each ring interval $i$, the near wake vortex sheet is modelled as,

$$
\begin{align*}
& r_{i}=r_{i}(s)  \tag{2.36}\\
& z_{i}=z_{i}(s)  \tag{2.37}\\
& \Gamma_{i}=\Gamma_{i}(s) \tag{2.38}
\end{align*}
$$

where $\boldsymbol{r}_{\boldsymbol{i}}$ and $z_{\boldsymbol{i}}$ represent radial and axial components of geometry respectively, $\Gamma_{i}$ is the circulation, and $s$ is the arc length along the sheet.

The spline is constructed for a sheet segment between vortex rings $i$ and $i+1$ as shown in Figure( 2.14). The first step is to construct a parabola, $g_{i}(s)$, through points $i-2, i-1$, and $i$ and another parabola, $g_{i+1}(s)$, through the points $i-1, i$, and $i+1$. The interpolated function, $f_{i}(s)$, is then computed by linearly interpolating the two functions $g_{i}(s)$ and $g_{i+1}(s)$

$$
\begin{equation*}
f_{i}(s)=g_{i}(s) \frac{s_{i+1}-s}{s_{i+1}-s_{i}}+g_{i+1}(s) \frac{s-s_{i}}{s_{i+1}-s_{i}} \tag{2.39}
\end{equation*}
$$

where $s_{i}$ and $s_{i+1}$ are the arc lengths at the points $i$ and $i+1$ respectively. This generates a cubic function for $f_{i}$ which has a slope at the end points that match the two parabolas $g_{i}$ and $g_{i+1}$. The cubic function $f_{i}$ may represent any of the three dependent variables $\boldsymbol{r}_{i}, z_{i}$, or $\Gamma_{i}$. Repeated over all of the near wake vortex sheet, this procedure yields a continuous function over the entire sheet that has continuous first derivatives and is relatively free of the extraneous curvature seen in some other types of splines.[11]

Using the above technique, the near wake vortex sheet is allowed to convect freely from blade passing time $\phi=\pi / 2$ to $\phi=3 \pi / 2$. At the end of the time period, the effect of a typical near wake vortex on the rotor blade is small, and therefore, the near wake is amalgamated into $n$ number of intermediate wake vorticies, where $n \ll N$.

### 2.2.5 Intermediate Wake Formation

From observations it is clear that the rotor vortex wake filaments convect nonuniformly in space. There are natural tendencies for the vorticies to split the vortex sheet and form distinct amalgamated vorticies. Therefore, at the end of the near wake time period the near wake is divided into $n$ number of amalgamated intermediate vorticies by analyzing the shearing velocities along the sheet. The shearing velocity $S_{i}$ along the sheet is defined to be,

$$
\begin{equation*}
S_{i}=\vec{u}_{i} \cdot \vec{t}_{i} \tag{2.40}
\end{equation*}
$$

where $\vec{u}_{i}$ is the induced velocity and $\vec{t}_{i}$ is the unit tangent vector at the vortex location $\boldsymbol{i}$. The tangent vector $\vec{t}_{\boldsymbol{i}}$ is computed in a central-differencing fashion such that

$$
\vec{t}_{i}=\frac{\vec{r}_{i+1}-\vec{r}_{i-1}}{\left|\vec{r}_{i+1}-\vec{r}_{i-1}\right|}
$$

and $\vec{u}_{i}$ is computed as explained previously.

With this definition of the shearing velocity, it becomes apparent that the vortex sheet will divide itself along the positions where $S_{i}$ changes signs in such a manner that the vortex sheet at that point is being stretched apart. Therefore, to determine if a point is a dividing point or not, one needs to consider:

- if $S_{i}<0$ and $S_{j}>0, j>i$, then the vorticies from $i$ to $j$ are travelling away from one another, and this is a point of sheet division.
- if $S_{i}>0$ and $S_{j}<0, j>i$, then the vorticies from $i$ to $j$ are travelling towards one another, and this is a point of sheet concentration.

During sheet rollup, these conditions do not remain true for all time at a particular point. However, after the rollup is substantially complete, the concentrations of vorticity are widely separated, and the condition for sheet division may be applied with some confidence. We have not made a study of how the sheet division condition varies with time as the wake rolls up. The aforementioned criteria was originally suggested by Yates[22], but Yates' application was to the initial wake of a wing rather than to a later rolled-up configuration. The intermediate wake vorticies are formed by amalgamating vorticies between the points where the vorticies are convecting away from each other. The procedure is identical to that of the tip vortex amalgamation; the conservation of circulation determines the strength, the conservation of impulse determines the location, and the conservation of kinetic energy determines the core size.

The vortex rings which constitute the near and intermediate wakes and the vortex cylinders which constitute the far wake are allowed to convect freely from time $\phi=\pi / 2$ to $\phi=3 \pi / 2$. Then a new bound circulation distribution along the rotor blade is computed. The next section will discuss the lifting-line theory employed here to calculate the blade loading.

### 2.3 Rotor Blade Load Modelling

Consider the Kutta-Zhukovsky theorem for a blade element

$$
\begin{equation*}
\vec{L}(r)=\rho \vec{U}(r) \times \vec{\Gamma}(r) \tag{2.41}
\end{equation*}
$$

where $\vec{L}$ is the lift vector, $\vec{U}$ is the free-stream velocity, and $\vec{\Gamma}$ is the bound vortex.

From Thin-Airfoil theory we obtain,

$$
\begin{equation*}
L(r)=\frac{1}{2} \rho U(r)^{2} c(r) C_{l}(r) \tag{2.42}
\end{equation*}
$$

where $c(r)$ here is the section chord length and $C_{l}(r)$ is the section lift coefficient. The lifting-line theory assumes that at any station along the blade the local flow is two-dimensional, and hence the section lift coefficient may be expressed as,

$$
C_{l}(r)=m_{o}\left(\alpha_{g e o m t}(r)-\alpha_{i n d}(r)\right)
$$

Here $m_{o}$ is the lift-curve slope, $\alpha_{g e o m t}(r)$ is the geometric angle of attack, and $\alpha_{\text {ind }}(r)$ is the induced angle of attack. In terms of these definitions, the bound circulation distribution may be expressed as,

$$
\begin{equation*}
\Gamma(r)=\frac{1}{2} U(r) c(r) m_{o}\left(\alpha_{g e o m t}(r)-\alpha_{i n d}(r)\right) \tag{2.43}
\end{equation*}
$$

where for a rotor blade,

$$
U(r)=\Omega r
$$

and $\Omega$ is the angular speed of the rotor blade.

For a small angle of attack $\alpha_{\text {geomt }}(r)$, the induced angle of attack $\alpha_{\text {ind }}(r)$ may be expressed as,

$$
\begin{equation*}
\alpha_{i n d}=\frac{w(r)}{U(r)} \tag{2.44}
\end{equation*}
$$

where $w(r)$ is the local downwash velocity with a positive sense in the $-z$ direction. Thus if the downwash velocity can be found, the lifting-line theory gives us the bound circulation distribution.

The downwash distribution, $w(r)$, is composed of two distinct components: the downwash contribution from the wake produced by previous blade passages, and the downwash contribution from the attached trailing vortex sheet. That is,

$$
\begin{equation*}
w(r)=w_{w a k e}(r)+w_{t v}(r) \tag{2.45}
\end{equation*}
$$

where $w_{\text {wake }}(r)$ is the previous wake contribution and $w_{t v}(r)$ is the attached trailing vortex sheet contribution. $w_{\text {wake }}(r)$ is the induced axial component of velocity on the blade computed from the near wake, the intermediate wake, and the far wake; and $w_{t v}(r)$ is the induced axial component of velocity computed from the attached vortex sheet. The attached vortex sheet represents the vortex
transportation from time $\phi=0$ to $\phi=\pi / 2$, and thus consists of $1 / 2$ vortex rings. The half vortex rings are replaced with complete vortex rings of half circulation. [14] In addition, since the strength of the attached trailing vortex sheet is a function of the bound circulation, and hence is not known apriori, the final bound circulation distribution is computed in an iterative fashion. The details of this process is discussed in the next chapter concerning numerical procedures.

The validity of the current lifting-line theory was tested against an elliptically loaded airfoil. Cosine spacing technique with 250 angular divisions were used to increase the accuracy.[17] The downwash obtained agreed well with the theory; from $r=0$ to $r=0.985$, the deviations were less than $1 \%$.


Figure 2.1: Isolated rotor and its trailing vortex sheet


Figure 2.2: Vortex wake of a hovering rotor


Figure 2.3: Vortex velocity components


Figure 2.4: Wake divisions in Miller's free wake technique


Figure 2.5: Formulation of vortex ring model


Figure 2.6: A typical wake geometry in motion


Figure 2.7: Coordinates for vortex ring velocity calculations


Figure 2.8: Straightforward Application of Discrete Model


Figure 2.9: Inner spiral model


1. Strength of tip filament found by conserving circulation
2. Location of tip filament found by conserving impulse
3. Tip filament core radius found by conserving kinetic energy

Figure 2.10: Tip vortex amalgamation procedure


Figure 2.11:Numerical integration of inner spiral model


Figure 2.12: Inner spiral model rotation

Sheet at $\phi=\phi_{0}$

is modelled as:
tip vortex ring with finite core

and higher order terms
$\Lambda$

Figure 2.13: Initial representation of the near wake


$$
f_{i}(s)=g_{i}(s) \frac{s_{i+1}-s}{s_{i+1}-s_{i}}+g_{i+1}(s) \frac{s-s_{i}}{s_{i+1}-s_{i}}
$$

$$
\begin{aligned}
g_{i}(s) & =\text { parabola through points } i-1, i, i+1 \\
& =a_{0}+a_{1}\left(s-s_{i}\right)+a_{2}\left(s-s_{i}\right)^{2}
\end{aligned}
$$

$$
\begin{aligned}
g_{i+1}(s) & =\text { parabola through points } i, i+1, i+2 \\
& =b_{0}+b_{1}\left(s-s_{i+1}\right)+b_{2}\left(s-s_{i+1}\right)^{2}
\end{aligned}
$$

Figure 2.14: Cubic spline for sheet rediscretization

## Chapter 3

## Numerical Methods

### 3.1 Time Integration

The new positions of the vortex rings in the near and intermediate wakes are calculated by numerical integrations of the velocity vectors shown in the previous chapter. For each vortex ring $i$ a total induced velocity from all the vorticies, including itself, is computed:

$$
\begin{align*}
& \dot{r}_{i}=\sum_{j=1}^{N} u_{r j}  \tag{3.1}\\
& \dot{z}_{i}=\sum_{j=1}^{N} u_{z j} \tag{3.2}
\end{align*}
$$

The new $r_{i}$ and $z_{i}$ position components are computed by integrating the above two equations. Two different types of the integration procedures were considered. The first is just the simple Euler scheme, i.e. a first order approximation to a Taylor Series.

$$
\begin{align*}
& r_{i}^{n+1}=r_{i}^{n}+\dot{r}_{i}^{n} \Delta \phi  \tag{3.3}\\
& z_{i}^{n+1}=z_{i}^{n}+\dot{z}_{i}^{n} \Delta \phi \tag{3.4}
\end{align*}
$$

where $\Delta \phi$ is the time step. In addition, the four stage Runge-Kutta time step-
ping sheme was also considered, [11]

$$
\begin{align*}
& r_{i}^{n+1}=r_{i}^{n}+\frac{k_{0}+2 k_{1}+2 k_{2}+k_{3}}{6}  \tag{3.5}\\
& z_{i}^{n+1}=z_{i}^{n}+\frac{l_{0}+2 l_{1}+2 l_{2}+l_{3}}{6} \tag{3.6}
\end{align*}
$$

where

$$
\begin{gathered}
k_{0}=\Delta \phi \dot{r}_{i}\left(r_{i}^{n}, z_{i}^{n}\right) \\
k_{1}=\Delta \phi \dot{r}_{i}\left(r_{i}^{n}+k_{0} / 2, z_{i}^{n}+l_{0} / 2\right) \\
k_{2}=\Delta \phi \dot{r}_{i}\left(r_{i}^{n}+k_{1} / 2, z_{i}^{n}+l_{1} / 2\right) \\
k_{3}=\Delta \phi \dot{r}_{i}\left(r_{i}^{n}+k_{2}, z_{i}^{n}+l_{2}\right)
\end{gathered}
$$

and

$$
\begin{gathered}
l_{0}=\Delta \phi \dot{z}_{i}\left(r_{i}^{n}, z_{i}^{n}\right) \\
l_{1}=\Delta \phi \dot{z}_{i}\left(r_{i}^{n}+k_{0} / 2, z_{i}^{n}+l_{0} / 2\right) \\
l_{2}=\Delta \phi \dot{z}_{i}\left(r_{i}^{n}+k_{1} / 2, z_{i}^{n}+l_{1} / 2\right) \\
l_{3}=\Delta \phi \dot{z}_{i}\left(r_{i}^{n}+k_{2}, z_{i}^{n}+l_{2}\right)
\end{gathered}
$$

For an accurate representation of the near wake it is essential that the numerical integration methodolgy conserve the three invariants: circulation, impulse,
and kinetic energy. The Euler scheme is attractive because of its simplicity, however, it is also quite inaccurate for large time steps. Figure ( 3.1) and Figure (3.2) show the change in fluid impulse with time as predicted by the two methods. The Euler scheme requires a time step of $O(0.001)$ to acheive the same accuracy of conservation of impulse as the Runge-Kutta scheme does in a time step of $O(0.01)$. Since the Runge-Kutta scheme requires four times as many calculations per time step as the Euler scheme, but is more accurate by $O(10)$, it was decided that the Runge-Kutta scheme was more efficient; all time integrations are done in a Runge-Kutta fashion.

### 3.2 Iteration Procedure

The convergence criteria must be met for both rotor blade loading and its related wake geometry. Furthermore, since the flow is steady, given a particular blade bound circulation distribution, there must exist a unique force-free wake geometry associated with it. To avoid numerical difficulties the two quantities are converged separately. That is, the wake rollup is converged before the blade loading is updated to respond to the changed blade downwash.

The vortex ring trajectories are calculated from time $\phi=\pi / 2$ to $\phi=3 \pi / 2$, at which point a new wake is created. The convergence criteria is now applied to the wake geometry; if the wake geometry has converged, a new rotor blade loading and its associated attached trailing vortex sheet is computed, if not, then no bound circulation computations are done, and the old trailing vortex sheet is again used. The vortex rings are now allowed to convect until time $\phi=5 \pi / 2$ at which point the convergence criteria for the wake geometry is again applied. The procedure continues similiarly until a convergence in both the rotor blade
loading and its wake is achieved.

Thus, there are a total of three distinct iterations performed,

- Lifting-line theory and matching the total downwash distribution with the blade loading
- Bound circulation
- Wake geometry


### 3.2.1 L-L theory and $\Gamma \longleftrightarrow w$ matching

From section(2.3) the bound circulation $\Gamma(r)$ is given as,

$$
\begin{equation*}
\Gamma(r)=K_{1}(r)\left(K_{2}(r)-\frac{w(r)}{U(r)}\right) \tag{3.7}
\end{equation*}
$$

where

$$
K_{1}(r)=\frac{1}{2} U(r) c(r) m_{o}
$$

and

$$
K_{2}(r)=\alpha_{g e o m t}(r)
$$

But $w(r)=w_{\text {wake }}(r)+w_{t v}(r)$ and is a function of $\Gamma(r)$. Thus $\Gamma(r)$ cannot be solved directly and an underrelaxation method is employed to iterate between $\Gamma(r)$ and $w(r)$. At each iteration (except the first) $w_{t v}(r)$ is computed from the attached trailing vortex sheet, a new $\Gamma(r)$ is computed, followed by the trailing
vortex sheet strength; the process is then repeated until a convergence criterion is satisfied, i.e.

$$
\begin{gathered}
\gamma^{k}(r) \longrightarrow w^{k+1}(r) \\
w^{k+1}(r) \longrightarrow \Gamma^{k+1}(r) \\
\Gamma^{k+1^{\prime}}(r)=\Gamma^{k^{\prime}}(r)+\omega\left(\Gamma^{k+1}(r)-\Gamma^{k^{\prime}}(r)\right) \\
\Gamma^{k+1^{\prime}}(r) \longrightarrow \gamma^{k+1}(r)
\end{gathered}
$$

where $\gamma(r)$ is the strength of the trailing vortex sheet, $k$ is the iteration level and $\omega=0.4$ is the relaxation parameter. The iteration begins with $\gamma^{0}(r)=w_{t v}^{0}(r)=$ 0 and ends when

$$
\begin{equation*}
\frac{\Gamma^{k+1^{\prime}}(r)-\Gamma^{k^{\prime}}(r)}{\Gamma^{k+1^{\prime}}(r)} \leq 0.001 \tag{3.8}
\end{equation*}
$$

### 3.2.2 Bound Circulation Iteration

Miller [14] showed that to eliminate large oscillations in the bound circulation computations, it is necessary to employ an underrelaxation scheme. That is,

$$
\begin{equation*}
\Gamma^{n e w}(r)=\Gamma^{o l d}(r)+\omega\left(\Gamma^{\prime}(r)-\Gamma^{o l d}(r)\right) \tag{3.9}
\end{equation*}
$$

where $\Gamma^{\prime}(r)$ is the value computed from the lifting-line theory just shown in the previous section, and $\omega=0.1$ is the relaxation parameter. The iteration ends when

$$
\begin{equation*}
\frac{\Gamma^{\text {new }}(r)-\Gamma^{\text {old }}(r)}{\Gamma^{\text {new }}(r)} \leq 0.001 \tag{3.10}
\end{equation*}
$$

### 3.2.3 Wake Geometry Iteration

It became apparent during computations that without some sort of dampening added to the vortex trajectories, the wake geometries would become chaotic with a large number of vortex pairings. To counteract this tendency an underrelaxation technique was added to the geometry calculations for the intermediate vortex rings,

$$
\begin{align*}
& r_{i}^{\text {new }}=r_{i}^{\text {old }}+\omega\left(r_{i}^{\prime}-r_{i}^{\text {old }}\right)  \tag{3.11}\\
& z_{i}^{\text {new }}=z_{i}^{\text {old }}+\omega\left(z_{i}^{\prime}-z_{i}^{\text {old }}\right) \tag{3.12}
\end{align*}
$$

where $r_{i}^{\prime}$ and $z_{i}^{\prime}$ are the integrated values and $\omega=0.1$ is the relaxation parameter. Again the iteration ends when

$$
\begin{equation*}
\frac{r_{i}^{n e w}-r_{i}^{o l d}}{r_{i}^{\text {new }}} \leq 0.001 \tag{3.13}
\end{equation*}
$$

and

$$
\begin{equation*}
\frac{z_{i}^{\text {new }}-z_{i}^{\text {old }}}{z_{i}^{\text {new }}} \leq 0.001 \tag{3.14}
\end{equation*}
$$



Figure 3.1: Fluid Impulse for the Euler Scheme


Figure 3.2: Fluid Impulse for the Runge-Kutta Scheme

## Chapter 4

## Application

### 4.1 Rotor Blade Geometry

The validity of the method is demonstrated on a straight two-bladed rotor as described below:[9]

| Blade radius, $R$ | 1.045 m |
| :--- | :--- |
| Angular speed, $\Omega$ | $73.3 \mathrm{rad} / \mathrm{sec}$ |
| Chord, $c / R$ | 0.0729 |
| Solidity, $\sigma$ | 0.0464 |
| Twist | $-11^{\circ}$ |

where solidity is

$$
\sigma=\frac{2 c}{\pi R}
$$

and $c$ is the dimensional chord length.

The distances and velocities were nondimensionalized as follows,

$$
\vec{x}=\frac{\vec{x}}{R}
$$

$$
\vec{u}=\frac{\vec{u}}{\Omega R}
$$

and thus nondimensional time is the blade passing angle $\phi$.

### 4.2 Initial Wake Placement and Input Values

Since the flow is steady, the simulation at $\phi=\pi / 2+\phi_{0}$ must start within a prescribed wake. The approach here is to use Miller's [14] final blade bound circulation distribution and wake positions as the initial conditions. See Figure( 4.1) and Figure( 4.2).

The input parameters for the simulation were as follows:

$$
\begin{array}{ll}
\Delta x_{\max }=0.01 & \text { Maximum translation in one time step } \\
\Theta_{\max }=2 \pi & \text { Near wake amalgamation angle } \\
N=60 & \text { Number of near wake vorticies } \\
\Lambda_{o}=0.01 & \text { Outer portion to be lumped at } \phi=\pi / 2+\phi_{o}
\end{array}
$$

At each time step maximum velocity, $U_{m a x}$, is computed; the incremental time $\Delta \phi$ then follows as,

$$
\Delta \phi=\frac{U_{\max }}{\Delta x_{\max }}
$$

The choice of $\Delta x_{\max }=0.01$ yields an average value of $\Delta \phi_{\text {avg }} \approx 0.05$, which is about the maximum time step allowed while conserving impulse/kinetic energy.
$\boldsymbol{\theta}_{\max }$ of $2 \pi$ is equivalent to one spiral turn amalgamation; more spiral turns unnecessarily increase the computational time. The choice of $N=60$ and $\Lambda_{o}=$ 0.01 was reached at after some numerical experimentation. It is felt that the given choices would allow a sufficient resolution without excessive sacrifices in computational speed.

One percent of the outer portion of the near wake rolled into the tip vortex filament results in a tip filament strength of about $30 \%$ of the total near wake strength (of one sign). The large value is a result of the high concentration of vorticity near the edge of the sheet. As determined using the methods described in section(2.2.4) the initial tip vortex location is found to be $r=0.987$ and $z=0.011$ with a corresponding time $\phi_{0}=0.00963$.

### 4.3 Results

The computed bound circulation distribution and its related wake are compared with Ballard's [2] published results in Figure( 4.3) and Figure( 4.4). The computed circulation distribution compares relatively well with the experimental values, although the peak value seems a bit excessive and outboard of the empirical result. The cause of the discrepancies near the rotor tip cannot be ascertained exactly, but three possible causes are:

- The three-dimensionality of the flow around the tip caused a breakdown in the lifting-line theory.
- The non-linear interactions of the tip vortex formation on the blade prior to the trailing edge.
- Relatively large viscous effects due to large velocity gradients present near the tip affected the experimental results.

One possible solution for the first cause would be to use a lifting-surface technique.

The vortex positions in the wake also compare reasonably well with the experimental values. The flagged values shown in Figure( 4.4) are the experimentally observed tip vortex locations. It should be noted that the inboard vorticies have not yet been observed. The computed tip vortex locations are outboard of the experimental values, and this is to be expected since the peak in the loading is also outboard of the observed values.

The current method required 191 iterations to converge. To demonstrate the process of the intermediate wake formation, the vorticies were convected for one additional iteration from time $\phi=191 \pi$ to time $\phi=192 \pi$. The initial conditions for this iteration are shown in Table( 4.1) and Figure( 4.5) and the final positions are shown in Figure( 4.4). At this point, the shear velocity, $S_{i}$, were computed along the near wake to determine the dividing points for the intermediate wake. From Figure( 4.6) and Figure( 4.7) one can see that the sheet divides itself at three different positions. Vorticies from arc length $\approx 0$ to arc length $\approx 0.6$ form a root vortex, vorticies from arc length $\approx 0.6$ to arc length $\approx 1.1$ form a center vortex, and remaining vorticies out to arc length $\approx$ 1.4 form the tip vortex. This formation of the first level of intermediate wake is shown in Figure( 4.8), where the amalgamation process was previously described in Chapter 2 to conserve the three invariants: circulation, impulse, and kinetic energy.

By applying Donaldson's rollup criteria to the loading curve one should conclude that three vorticies are necessary to transport the trailing vorticity. The direct rollup calculations also show a formation of three vorticies. The tip vortex strength shows an amalgamation of primarily negative vorticity. Donaldson's tip vortex strength is computed by summing only the negative vorticity in Table( 4.1) which leads to $\Gamma_{\text {tip }}=-0.023428$. The current analysis indicates a tip vortex strength of $\Gamma_{\text {tip }}=-0.022220$. There is a difference of approximately $5 \%$ between the values, indicating that the tip vortex is somewhat weaker than Donaldson's model would indicate due to the rollup of counter-signed vorticity.

Figure( 4.9) shows the distribution of vorticity vs. arc length. From arc length $\approx 0.9$ to arc length $\approx 1.4$ there is little or no vorticity, and one can conclude that the strong tip vortex has essentially separated from the remainder of the sheet. In addition, the figure shows two distinct areas of vorticity concentrations inboard of the tip. The tip and the two inboard vorticies combined would, infact, represent approximately $85 \%$ of the total vorticity (of both signs) in the sheet. Therefore, replacing the vortex sheet by a small number of concentrated vorticies should be an accurate representation.

### 4.4 Computer Usuage

All calculations were done on the MIT-CFD Alliant FX/8 vector computer. In a full vector/concurrency mode each iteration required approximately 10 minutes of CPU time for completion. The present method required 191 iterations for convergence, i.e. about 32 hours of CPU time. Such a large requirement in computer time is a deterrent for the use of the present method, especially since there are other methods $[14,15,13,21]$ which require far fewer hours of CPU
usuage with comparable results in bound circulation distribution. However, the methodolgy used here was developed to assess the effect of free wake rollup on rotor load predictions and also to validate some of the simpler models of wake geometry as Miller's $[14,15]$ method rather than as an engineering tool. It should be noted that all attempts were made to make the computer code as efficient as possible, including considerable manual vectorization.


Figure 4.1: Initial load distribution from Miller


Figure 4.2: Initial wake geometry


Figure 4.3: Final load distribution


Figure 4.4: Final wake geometry


Figure 4.5: Initial wake geometry for final iteration


Figure 4.6: Shear strength distribution for near wake


Figure 4.7: Velocity vectors for near wake


Figure 4.8: New intermediate wake formation


Figure 4.9: Vorticity distribution in near wake

| 1 | RADIUS | HEIGHT | GAMMA |
| :---: | :---: | :---: | :---: |
| 0 | $0.100000 E+00$ | -0.190906E-02 | 0.000000 E+00 |
| 1 | $0.107555 \mathrm{E}+00$ | -0.190906E-02 | $0.132221 \mathrm{E}-02$ |
| 2 | $0.122665 E+00$ | -0.190906E-02 | 0.110944E-02 |
| 3 | $0.137775 \mathrm{E}+00$ | -0.190906E-02 | $0.969350 \mathrm{E}-03$ |
| 4 | $0.152886 \mathrm{E}+00$ | -0.190906E-02 | $0.839926 \mathrm{E}-03$ |
| 5 | 0.167996E+00 | -0.190906E-02 | $0.742824 E-03$ |
| 6 | $0.183106 E+00$ | -0.190906E-02 | $0.660581 \mathrm{E}-03$ |
| 7 | 0.198216E+00 | -0.190906E-02 | $0.603929 \mathrm{E}-03$ |
| 8 | 0.213326E+00 | -0.190906E-02 | $0.568352 \mathrm{E}-03$ |
| 9 | $0.228436 E+00$ | -0.190906E-02 | $0.539236 \mathrm{E}-03$ |
| 10 | 0.243547E+00 | -0.190906E-02 | $0.513880 \mathrm{E}-03$ |
| 11 | $0.258857 E+00$ | -0.190906E-02 | $0.488983 \mathrm{E}-03$ |
| 12 | 0.273767E+00 | -0.190906E-02 | $0.458082 \mathrm{E}-03$ |
| 13 | 0.288877E+00 | -0.190906E-02 | $0.432115 \mathrm{E}-03$ |
| 14 | 0.303987E+00 | -0.190906E-02 | $0.414475 \mathrm{E}-03$ |
| 15 | 0.319097E+00 | -0.190906E-02 | $0.397019 \mathrm{E}-03$ |
| 16 | 0.334208E+00 | -0.190906E-02 | 0.375556E-03 |
| 17 | $0.349318 E+00$ | -0.190906E-02 | $0.356263 \mathrm{E}-03$ |
| 18 | $0.364428 E+00$ | -0.190906E-02 | $0.336615 \mathrm{E}-03$ |
| 19 | $0.379538 E+00$ | -0.190906E-02 | $0.322777 \mathrm{E}-03$ |
| 20 | $0.394648 \mathrm{E}+00$ | -0.190906E-02 | $0.310860 \mathrm{E}-03$ |
| 21 | $0.409758 \mathrm{E}+00$ | -0.190906E-02 | $0.298928 \mathrm{E}-03$ |
| 22 | $0.424869 E+00$ | -0.190906E-02 | 0.285679E-03 |
| 23 | $0.439979 E+00$ | -0.190906E-02 | $0.284327 \mathrm{E}-03$ |
| 24 | $0.455089 E+00$ | -0.190906E-02 | 0.252877E-03 |
| 25 | $0.470199 E+00$ | -0.190906E-02 | $0.242041 \mathrm{E}-03$ |
| 26 | $0.485309 \mathrm{E}+00$ | -0.190906E-02 | $0.226987 \mathrm{E}-03$ |
| 27 | $0.500419 \mathrm{E}+00$ | -0.190906E-02 | $0.223110 \mathrm{E}-03$ |
| 28 | $0.515530 E+00$ | -0.190906E-02 | 0.221499E-03 |
| 29 | $0.530640 \mathrm{E}+00$ | -0.190906E-02 | $0.218379 \mathrm{E}-03$ |
| 30 | $0.545750 E+00$ | -0.190906E-02 | $0.190482 \mathrm{E}-03$ |
| 31 | $0.560860 \mathrm{E}+00$ | -0.190906E-02 | $0.157134 \mathrm{E}-03$ |
| 32 | $0.575970 \mathrm{E}+00$ | -0.190906E-02 | $0.129644 \mathrm{E}-03$ |
| 33 | $0.591081 \mathrm{E}+00$ | -0.190906E-02 | $0.108640 \mathrm{E}-03$ |
| 34 | 0.608191E+00 | -0.190906E-02 | 0.858419E-04 |
| 35 | $0.621301 \mathrm{E}+00$ | -0.190906E-02 | 0.671865E-04 |
| 36 | $0.636411 \mathrm{E}+00$ | -0.190906E-02 | $0.587981 \mathrm{E}-04$ |
| 37 | 0.651521E+00 | -0.190906E-02 | $0.484763 \mathrm{E}-04$ |
| 38 | $0.666631 \mathrm{E}+00$ | -0.190906E-02 | $0.425968 \mathrm{E}-04$ |
| 39 | $0.681742 \mathrm{E}+00$ | -0.190906E-02 | $0.393866 \mathrm{E}-04$ |
| 40 | $0.696852 \mathrm{E}+00$ | -0.190906E-02 | 0.324324E-04 |
| 41 | $0.711962 \mathrm{E}+00$ | -0.190906E-02 | 0.243736E-04 |
| 42 | $0.727072 \mathrm{E}+00$ | -0.190906E-02 | 0.162050E-04 |
| 43 | $0.742182 \mathrm{E}+00$ | -0.190906E-02 | 0.149012E-04 |
| 44 | $0.757292 \mathrm{E}+00$ | -0.190906E-02 | 0.286410E-04 |
| 45 | $0.772403 \mathrm{E}+00$ | -0.190906E-02 | $0.441615 E-04$ |
| 46 | $0.787513 \mathrm{E}+00$ | -0.190906E-02 | $0.921236 E-04$ |
| 47 | $0.802623 E+00$ | -0.190906E-02 | $0.156075 \mathrm{E}-03$ |
| 48 | $0.817733 \mathrm{E}+00$ | -0.190906E-02. | 0.293590E-03 |
| 49 | $0.832843 \mathrm{E}+00$ | -0.190906E-02 | $0.488207 E-03$ |
| 50 | $0.847953 \mathrm{E}+00$ | -0.190906E-02 | $0.854593 E-03$ |
| 51 | $0.863084 \mathrm{E}+00$ | -0.190906E-02 | $0.138038 \mathrm{E}-02$ |
| 52 | $0.878174 \mathrm{E}+00$ | -0.190906E-02 | $0.171804 \mathrm{E}-02$ |
| 53 | $0.893284 \mathrm{E}+00$ | -0.190906E-02 | 0.172476E-02 |
| 54 | $0.908394 \mathrm{E}+00$ | -0.190906E-02 | 0.131142E-02 |
| 55 | $0.923504 E+00$ | -0.190906E-02 | $0.325283 E-03$ |
| 56 | $0.938614 \mathrm{E}+00$ | -0.190906E-02 | -0.343490E-03 |
| 57 | $0.953725 E+00$ | -0.190906E-02 | -0.911472E-03 |
| 58 | $0.968835 \mathrm{E}+00$ | -0.190906E-02 | -0.234807E-02 |
| 59 | $0.983945 \mathrm{E}+00$ | -0.190906E-02 | -0.118641E-01 |
| 60 | $0.995580 \mathrm{E}+00$ | $0.510344 \mathrm{E}-02$ | -0.796055E-02 |

Table 4.1: Initial near wake values for final iteration

## Chapter 5

## Conclusions

A free wake analysis using a lifting-line technique is combined with the Kantelis \& Widnall's model of a vortex sheet rollup to analyze the flow-field about a hovering rotor. The computational results are compared with experimental data for circulation distribution and wake geometry.

Miller's free wake technique is extended by using a high resolution near wake. A naive application of the vortex sheet convection was shown to lead to large numerical errors due to 1) uneven vortex sheet stretching and 2) inaccurate modelling of the singularity at the tip. The errors due to vortex sheet stretching are reduced by rediscretizing the sheet at every time step to ensure that the vortex elements remain equally spaced. The singularity at the sheet edge is resolved by using a special model at the tip. The model replaces the outermost portion of the sheet with a single vortex of equivalent circulation, impulse and kinetic energy. The model also includes some higher order terms to account for the asymmetry of the spiral.

The intermediate wake is formed from an analysis of the shearing forces on the near wake. The shearing force distribution was calculated to find the division points along the sheet; the sheet was then amalgamated to form the intermediate wake vorticies. The amalgamtion process was equivalent to that of the near wake tip vortex, i.e. conserve the three invariants, circulation, impulse
and kinetic energy.

The validity of the model was tested by comparing it with known experimental values. The circulation distribution showed an overall good agreement with some discrepancies arising near the tip. The errors may be due to the breakdown of the model because of 1) three-dimensionality of the flow, 2) non-linearities near the tip, or 3) viscous effects.

The comparison of direct rollup calculation with those given by Donaldson [5] yielded mixed results. The strength of the tip vortex is somewhat less than that predicted by Donaldson's criterion due to the incorporation of counter-signed vorticity during the rollup process. The analysis of the shearing velocity distribution indicated the same number of amalgamated vorticies as that predicted by Donaldson.

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## APPENDIX

Program listings

```
    program hover
        include '/user/huh/rollup/common.f'
            common /th/ theta(0:900),thmax
            COMMON /E1/ TAUO,BDUM,CDUM,DDUM,TAU2LOOP
            COMMON /E2/ DUM1,DUM2,DUM3,DUM4,DUM5,AE2
            DIMENSION SGAM(0:900),DDGM(0:900)
            REAL*4 LAMBDA,K
            CHARACTER PLTITL*80,TITLE*80
            character*60 gtitle,ltitle
            integer indexvort(10)
            logical lgo,lrite,lgo2,lrite2
            real tmp(2)
            data lgo,lrite,lgo2,lrite2/.true.,.false.,.true.,.false./
            data iaxi,ipanel,iredisc,itrs,intkod,iload/1,0,1,0,1,0/
C
C NRING = number of rings in the near wake (counting starts at 0)
C NSMALL = number of vortex rings which form the sheet, NSMALL=NRING-1
C
C Ring Indicies
C
C Index Meaning
C
C 0 Degenerate ring on the central axis of the ring system.
C 1 - NSMALL These are the "real" vortex rings, i.e., they have non-
```

```
C (inclusive) zero circulation. These all lie on the sheet.
C
C NSMALL This ring forms the outer boundary of the sheet.
C
C
C
C
C
C
C
C
READ(5,*)IWR
read(5,*) inew
READ(5,*)DXMAX,nitrs,nload
READ (5,*) THMAX
READ (5,*)FRAC
READ (5,*)NU2,nu22
if(inew .eq. 0)then
    READ(5,*)NRING
    READ(5,*)LINEUP
    READ(5,*)MINSECT
    READ(5,*)MAXSEGS
    NSMALL=NRING-1
endif
TIMEAB=9.0E20
DISTIN=9.0E5
DISTMX=DISTIN**2
C
c open files for output
C
```

```
open(unit=15,file='loading.dat',form='formatted',status='new')
```

open(unit=15,file='loading.dat',form='formatted',status='new')
open(unit=9,file='time.dat',status='new',form='formatted')
open(unit=9,file='time.dat',status='new',form='formatted')
open(unit=12,file='geomt.dat',form='formatted',status='new')

```
open(unit=12,file='geomt.dat',form='formatted',status='new')
```

```
            open(unit=26,file='velvec.dat',form='formatted',status='new')
            open(unit=19,file='newgeomt.dat',form='formatted',status='new')
            open(unit=7,form='unformatted',status='new',file='sheettemp.dat')
            call seta
            TIME(1)=((2.0*PI**2)/(9.0*AE2))*frac**(1.5)
            call experiment
C
c if inew = 0 then start new else restart
C
    if(inew .eq. 0)then
c
c set circulation from Miller's method
c
            call initwake
            call initcirc(frac)
            call ritecirc
                    tint = pi
                    istep = 1
                    TIMENOW=TIME(1)
            else
                    open(unit=16,file='restart.dat',form='unformatted',status='old')
                    call readinput(tint,itpre)
                    istep = 2
                    timenow = time(2)
                    close(16)
        endif
            CALL movem(istep)
        call riteinput
C
C--Write out the input data
C
C
C----------------------------------------------------------------------------
C-------- MAIN TIME LOOP ------------------------------------------------
C------------------------------------------------------------------------
```

```
C
    234 continue
        if(.not. lgo)goto 99
        if ((timenow-tint) .ge. 0.)goto99
        istep = istep + 1
C
C--Compute velocities at step ISTEP which
C corresponds to time (ISTEP-1)*"DT" and also update positions
C
            CALL MOVEM(ISTEP)
            WVTIP=WVEL(NRING)
            UVTIP=UVEL(NRING)
            CTIP=CORE (NRING)
            NIN=NSMALL
            NOUT=NSMALL
            GMM=GAMMA (NRING)
            CALL REDISC(RADIUS,HEIGHT,GAMMA,NIN,NOUT,GMM,SUML,DELTAL)
            WVEL (NRING)=WVTIP
            UVEL (NRING)=UVTIP
            CORE (NRING)=CTIP
            WVTIP=WVEL (NRING)
            UVTIP=UVEL(NRING)
            CTIP=CORE (NRING)
            CALL LUMPS
            WVEL (NRING)=WVTIP
            UVEL (NRING)=UVTIP
            CORE(NRING)=CTIP
C
C--Now update the core size, since the sheet has stretched
C
            DO 691 II=1,NSMALL
    691 CORE(II)=WPAN (II)/2.0
C
    if(istep .ge. 50000)then
                type*,'istep .gt. 50000',istep,timenow
```

```
                goto 57
        endif
            goto 234
    9 9 ~ C O N T I N U E ~
C
C------------------------------------------------------------------------
C--------- END OF MAIN TIME LOOP
C----------------------------------------------------------------------
C
    itrs = itrs + 1
            write(9,*) , '
            write(9,17)'pi itr# = ',itrs,' time = ',timenow
            write(9,*) , '
17 format(a,i6,a,f10.5)
        tint = tint + pi
c
c write out results
C
            if(mod((itrs),nu2) .eq. 0 .or. itrs .eq. nitrs)then
            lrite = .true.
            endif
C
c redivide the sheet, calculate new loading and wake geometry
C
            call seta
            call divide(indexvort,nvort)
            call newwake(indexvort,nvort,lrite,lgo)
c
c check if wake converged
c
    if(lgo .and. itrs .lt. nitrs)then
            call oldsheet(istep)
            call riterst(tint,itrs+itpre)
            goto 234
            endif
```

```
            if(itrs .ge. nitrs)goto57
            iload = iload + 1
            write(9,*) ' ,
            write(9,18)'load itr # = ',iload
            write(9,*) ' '
18
            format(a,i6)
            if(mod(iload,nu22) .eq. 0)lrite2=.true.
C
c if wake converged, recalculate new loading on blade
c
    call circulation(lgo2,lrite2)
            call makesheet(frac,istep)
            call riteinput
c
c check if loading has converged
c
            if(lgo2 . and. iload .lt. nload)then
                lgo = .true.
                call riterst(tint,itrs+itpre)
                    goto 234
            endif
5 7
                    continue
            type*,'Total # of pi iterations = ',itrs+itpre
            type*,'Total # of load iterations = ',iload
            call riterst(tint,itrs+itpre)
                    END
```

C
c this is the include file of common blocks c
COMMON /B1/ RADIUS (0:900), HEIGHT (0:900), WVEL(0:900), UVEL(0:900),
\$
\$ GAMMA (0:900), CORE (0:900), TIME (0:50000), TENERG (0:50000), TIMPUL(0:50000)

```
COMMON /B2/ NRING.PI,DXMAX,INTKOD,NSMALL,VOLTIP,IAXI
COMMON /B9/ IWR
COMMON /PANEL/ IPANEL,DRPAN(0:900),DZPAN(0:900),WPAN (0:900)
COMMON /C1/ LINEUP,MAXSEGS,MINSECT
COMMON /D1/ DISTMX
COMMON /SEGS/ AXLSEG (0:900)
COMMON /F1/ ALPLAST,GMTO
COMMON /ZZ/ ZXO,ZYO
COMMON /G1/ TIPVAX
COMMON /NOW/ TIMENOW,BIGRREF
    common /wake/zinter(10,5),rinter(10,5),ginter(10,5),cinter (10,5)
& ,zfar(10),rfar(10),gfar(10),uveli(10,5),wveli(10,5)
& ,ninter(5),nfar,level,volinter(10,5)
    common /blade/pitch(900),chord(900),span(900),aslope,sigma,nspan
COMMON /RK/ RADO(0:900),HEIO(0:900),RK(0:900,4),HK(0:900,4)
    common /rkwake/ zint0(10,5),rint0(10,5),rki(10,5,4),hki(10,5,4)
    common /loading/circold(900),circ(900)
    common /oldwake/roldi (10,5),zoldi (10,5),goldi(10,5)
    common /exper/expx(28), expy(28)
```

    SUBROUTINE ANGLES(RS,ZS,NTOTSUB,THSUB)
    C
C--Get angles relative to tip for each subinterval
C
include '/user/huh/rollup/common.f'
real $r(0: 900), z(0: 900)$
equivalence ( $r$, radius)
equivalence ( $z$, height)
DIMENSION RS (0:8100), ZS (0:8100), THSUB (0:8100)
REAL*4 MAG
C
$R T=R$ (NRING) !tip values of radius

```
ZT=Z(NRING) ! and height
C
    A=-1.0
B=0.0
C=RS(0)-RT
D=ZS(0)-ZT
DOT =A*C+B*D
CROSS=A*D-B*C
MAG=SQRT ((A*A+B*B)*(C*C+D*D))
ARG=DOT/MAG
IF (ABS(ARG) .GT. 1.0) ARG=SIGN(1.0,ARG)
THSUB(0)=SIGN(1.0,CROSS)*ACOS(ARG)
C
            DO 100 I=1,NTOTSUB
            A=RS(I-1)-RT
            B=ZS}(I-1)-Z
            C=RS(I)}-R
            D=ZS(I)-ZT
            DOT =A*C+B*D
            CROSS }=A*D-B*
            MAG=SQRT ((A*A+B*B)*(C*C+D*D))
            ARG=DOT/MAG
            IF (ABS(ARG) .GT. 1.0) ARG=SIGN(1.0,ARG)
    100 THSUB (I)=THSUB(I-1)+SIGN(1.0,CROSS)*ACOS(ARG)
C
            RETURN
            END
            SUBROUTINE CENSPIR(VXP,VYP)
C--Calculates the velocity induced at the spiral center due to the
C inner spiral region
    include '/user/huh/rollup/common.f'
```

```
COMMON /E1/ TAUO,B,C,D,TAU2LOOP
COMMON /E2/ AO,A1,A2,A3,RHOSTAR,A
C
TAUOLOOP=0.24016
BGROLOOP=0.32094
BIGRPRM=BGROLOOP*(A*TIME(1))**(2.0/3.0)
CC=A**(4.0/3.0)*TIMENOW**(1.0/3.0)/BIGRPRM
ALFR=alplast-pi/2.
TAUA=TAUOLOOP+ALFR*(0.60702015)/(2.0*PI)
C
ALF=alplast-pi/2.
IF (ALF .LT. 4.0*PI) THEN
        ZX=ZXO
        ZY=ZYO
ELSE IF (ALF .GT. 6.0*PI) THEN
        ZX=0.0
        ZY=0.0
ELSE
        ZX=ZXO-ZXO*(ALF-4.0*PI)/(2.0*PI)
        ZY=ZYO-ZYO*(ALF-4.0*PI)/(2.0*PI)
ALF=alplast-pi/2
ENDIF
C
VXTILD=-0.00029091*ZX*CC*TAUA**(-2.30525)
VYTILD=-0.00424515*ZY*CC*TAUA**(-1.45439)
VXP=VXTILD*COS(ALFR)-VYTILD*SIN(ALFR)
VYP=VXTILD*SIN(ALFR)+VYTILD*COS(ALFR)
C WRITE(6,*)'NRING,UVEL(NRING),WVEL(NRING)'
C WRITE(6,*) NRING,UVEL(NRING),WVEL(NRING)
C WRITE(6,*)'TAUOLOOP,BGROLOOP,A,TIME(1),TIMENOW,BIGRPRM'
C WRITE(6,*) TAUOLOOP,BGROLOOP,A,TIME(1),TIMENOW,BIGRPRM
C WRITE(6,*)'CC,ALFR,TAUA,ZX,ZY,VXTILD,VYTILD,VXP,VYP'
C WRITE(6,*) CC,ALFR,TAUA,ZX,ZY,VXTILD,VYTILD,VXP,VYP
C
RETURN
```


## END

```
    SUBROUTINE circulation(1go,1rite)
C
C--This routine calculates new loading on the blade based on the
C lifting-line theory.
C
    include '/user/huh/rollup/common.f'
    REAL*4 LAMBDA,K,k1
    logical lgo,lrite
    real wvnear(900),wvint(10,5),wvfar(10),wvelt(103),circt(103)
    character*60 gtitle,title
    common /temp1/a0,a1,a2,a3,a4,b0,b1,b2,b3,b4
    common /temp2/c0,c1,c2,c3,c4,d0,d1,d2,d3,d4
    data omega,epsilon/0.1,0.001/
CVD$R ASSOC
C
C--Loop over all rings, except center ring, get induced velocity at each
C
    DO 10 I=2,nspan
        W=span(i)
        Z=0.
        DO 20 J=1,nring
            WP=RADIUS (J)
            ZP=HEIGHT (J)
            if(wp .le. 0.18 .or. zp .gt. O.)goto20
            R1=SQRT((Z-ZP)**2+(W-WP)**2)
            R2=SQRT((Z-ZP)**2+(W+WP)**2)
            LAMBDA=(R2-R1)/(R2+R1)
            FK=1.0-LAMBDA**2
            K=AO+FK*(A1+FK*(A2+FK*(A3+FK*A4)))
    $ +(BO+FK*(B1+FK*(B2+FK*(B3+FK*B4))))*ALOG(1.0/FK)
```

```
                    E=CO+FK*(C1+FK*(C2+FK*(C3+FK*C4)))
        F1=((W-WP)/R1+(W+WP)/R2)*(K-E)
        F2=2.0*E*LAMBDA/(1.0-LAMBDA**2)
        F3=(W*(R1**2-R2**2)+WP*(R1**2+R2**2))/(R1*R2*(R1+R2))
        WV=GAMMA (J)*(F1+F2*F3)/(2.0*PI*W)
        WVnear (j)=WV
        CONTINUE
```

```
                        +(D0+FK*(D1+FK*(D2+FK*(D3+FK*D4))))*ALOG(1.0/FK)
c
c calculate influence of intermediate wake
c
        do 50 j1 = 2,level
            do 50 i1 = 1,ninter(j1)
            wp=rinter(i1,j1)
            zp=zinter(i1,j1)
            r1=sqrt((z-zp)**2+(w-wp)**2)
            r2=sqrt((z-zp)**2+(w+wp)**2)
            lambda=(r2-r1)/(r2+r1)
            fk=1.0-1ambda**2
            k = a0+fk*(a1+fk*(a2+fk*(a3+fk*a4)))
    $ +(b0+fk*(b1+fk*(b2+fk*(b3+fk*b4))))*alog(1.0/fk)
            e = c0+fk*(c1+fk*(c2+fk*(c3+fk*c4)))
    $ +(d0+fk*(d1+fk*(d2+fk*(d3+fk*d4))))*alog(1.0/fk)
        f1=((w-wp)/r1+(w+wp)/r2)*(k-e)
        f2=2.0*e*lambda/(1.0-lambda**2)
        f3=(w*(r1**2-r2**2)+wp*(r1**2+r2**2))/(r1*r2*(r1+r2))
        wv=ginter(i1,j1)*(f1+f2*f3)/(2.0*pi*w)
        wvint(i1,j1)=wv
        continue
c
c calculate influence of far wake cylinder
C
        dphi = pi/20. !40 integration steps
        nphi = nint(2.*pi/dphi) - 1
        do 671 ifar = 1,nfar
```

```
    zp = abs(z - zfar(ifar))
    wsum = 0.
    do }601\mathrm{ iphi = 1,nphi
        phi = float(iphi)*dphi
        x3 = rfar(ifar)*rfar(ifar) +w*w+zp*zp-
            2.*rfar(ifar)*w*cos(phi)
    dwfar = dphi*rfar(ifar)*(rfar(ifar)-w*cos(phi))/
    &
    wsum = wsum + dwfar
601 continue
    x3 = rfar(ifar)*rfar(ifar)+w*w+zp*zp-2.*rfar(ifar)*w
    if((x3-zp*zp) .ne. 0.)then
        dwfar = dphi*rfar(ifar)*(rfar(ifar)-w)/(x3-zp*zp)*(1.-zp/sqrt(x3))
        wsum = wsum + dwfar
    endif
    wfar = wsum/4./pi*gfar(ifar)
    wvfar(ifar)= wfar
671 continue
    wvelt(i) = 0.
    do i1 = 1,nring
        wvelt(i) = wvelt(i) + wvnear(i1)
    enddo
    do j1 = 2,level
        wvelt(i) = wvelt(i) + wvint(1,j1)+wvint(2,j1)+wvint(3,j1)
    enddo
    wvelt(i) = wvelt(i) + wvfar(1) + wvfar(2)
C
C--get velocity of higher order terms in tip region
C
CALL DOUBLET( \(\mathbf{W}, \mathrm{Z}, \mathrm{VXPRM}, \mathrm{VYPRM}, \mathrm{VRPRM}, \mathrm{VTHPRM}\) )
WVELt (i) \(=\) WVELt ( \(i\) ) +VYPRM
10 CONTINUE
call itercirc (wvelt, circt)
do \(i=1\), nspan
if(span(i).le. 0.11 .or. span(i) .ge. 0.985 )then
```

```
                circ(i) = circold(i)
            else
                circ(i) = circold(i) + omega*(circt(i)-circold(i))
            endif
enddo
lgo = .false.
do 200 i = 50,nspan
        if(circ(i) .eq. O.)goto200
    dchange = abs((circ(i)-circold(i))/circ(i))
        if(dchange .gt. epsilon)then
            lgo = .true.
            goto 201
        endif
        continue
        Irite = .true.
        type*,'Load convergence at time = ',timenow
c
        201
c
c reset old values
c
    do i = 1,nspan
            circold(i) = circ(i)
        enddo
        if(lrite)then
    C
    c write out loading values
c
    gtitle='LOADING'
    title='RADIUS~GAMMA TIME = '
    write(15,*)gtitle
    write(15,*)2
    write(15,*)52
    write(15,2)title,timenow
        format(a20,f9.5)
```

```
    write(15,*)1
    write(15,*)0.,1.,0.,3.e-02
    write(15,*)2
    write(15,*)nspan
CVD$ NOCONCUR
CVD$ NOVECTOR
    do 90 i = 1,nspan
        write(15,*) span(i),circ(i)
    90 continue
    write(15,*)4
    write(15,*)28
CVD$ NOCONCUR
CVD$ NOVECTOR
    do i}=1,2
        write(15,*) expx(i), expy(i)
    enddo
    lrite = .false.
    endif
C
C--done
C
    RETURN
    END
    SUBROUTINE divide(indexvort,nvort)
C
C--This routine calculates the division positions of the near wake
C
    include '/user/huh/rollup/common.f'
    common /oldgeom/radiusold(0:900), heightold(0:900)
        COMMON /TH/ THETA(0:900),THMAX
    real tanr(900), tanz(900)
```

```
character*60 gtitle,title
integer indexvort(10)
CVD$R ASSOC
c
    do i=1,nring
        radiusold(i) = radius(i)
        heightold(i) = height(i)
        enddo
c
        do 5 i = 1,10
        indexvort(i) = 0
        continue
c
c calculate tangent vector to the sheet, central differencing
C
2
        do 2i=1,nsmall
        tanr(i)=radius(i+1) - radius(i-1)
        tanz(i)=height(i+1) - height(i-1)
        str = sqrt(tanr(i)**2 + tanz(i)**2)
        tanr(i)=tanr(i)/str
        tanz(i)=\operatorname{tanz(i)/str}
        continue
C
c calculate the inner product of tangent vector and the velocity vector
c where they change the sign is the dividing point
c
        nvort = 1
c gtitle = 'SHEAR STRENGTH VS. ARC LENGTH'
c write(2,*)gtitle
c write(2,*) 1
c write(2,*) 60
c write(2,*) 'x-UNIT TANGENT*VELOCITY~'
c WRITE(2,*) 1
c WRITE(2,*) 0,1.6,-.3,.3
c WRITE(2,*)2
```

c

C
c WRITE (2,*) Nring
ch do $i=2$, nring

```
    WRITE(2,*) NSMALL
    predot = tanr(1)*uvel(1) + tanz(1)*wvel(1)
    write(2,*)axlseg(1),predot
    do 90 i = 2,nsmall
        curdot = tanr(i)*uvel(i) + tanz(i)*wvel(i)
        write(2,*)axlseg(i),curdot
        if((curdot * predot) .lt. 0. .and. curdot .It. O.)then
            indexvort(nvort) = i
            nvort = nvort + 1
        endif
        predot = curdot
        continue
        indexvort(nvort) = nring + 1
    dr = radius(nring)-radius(nsmall)
        dz = height(nring)-height(nsmall)
    axlseg(nring) = sqrt(dr*dr+dz*dz)+axlseg(nsmall)
    write(2,*)gtitle
        write(2,*) 1
    write(2,*) 60
    write(2,*) 'x~VORTICITY~'
    WRITE(2,*) 1
        write(3,*)axlseg(i),gamma(i)
    enddo
RETURN
eND
```

SUBROUTINE DOUBLET (W, Z,VXPRM,VYPRM,VRPRM, VTHPRM)

C--Calculates the higher order velocity components in the spiral region C

C $W$ radius of control point
C $Z \quad$ height of control point
C VXPRM $V$ sub $x$ prime in the notes, radial velocity component
C VYPRM $V$ sub $y$ prime in the notes, axial velocity component
C GMTO strength of tip at start of run
C
include '/user/huh/rollup/common.f'
COMMON /E1/ TAUO, BDUM, CDUM, DDUM, TAU2LOOP
COMMON /E2/ AO,A1,A2,A3,RHOSTAR,A
CVD\$R ASSOC
DIMENSION $B(7,2), D(7,2) \quad!B(c o e f .$, exponent)
DATA B /. $0053006800, .0013100600, .000574976$, . 000319442 ,
\$ .0002020870, .0001387070, . 000100337 ,
\$ -1.38557 , $-1.36294,-1.34587$, -1.33015 ,
\$ -1.31621 , -1.30109 , -1.28252 /
DATA D /. $0010074100, .0001546680, .0000656961, .0000312179$,
\$ .0000157322, .0000101938, . 0000115224 ,
\$ -2.37327 , -2.28212 , -2.40883 , -2.36051
\$ -2.17250 , -2.22138 , -2.80713 /
BIGRREF=0.0966063
NCOEF=7
C
BIGRPRM=BIGRREF*(A*TIME(1))**(2.0/3.0) !constant in time RPRIME=SQRT( (W-RADIUS (NRING)) $* * 2+(Z-H E I G H T(N R I N G)) * * 2)$
R=RPRIME/BIGRPRM
SINTST=(Z-HEIGHT(NRING))/RPRIME
COSTST $=(W$-RADIUS(NRING))/RPRIME
THETAST = ATAN2 (SINTST, COSTST)
VRPRM=0.0
VTHPRM=0.0
C WRITE ( $6, *)^{\prime}$ ALPLAST, ALPROT, BIGRREF,BIGRPRM, A,TIME(1)'
C WRITE (6,*) ALPLAST,ALPROT,BIGRREF,BIGRPRM,A,TIME(1)

```
C WRITE(6,*)' CC,W,RADIUS(NRING),Z,HEIGHT(NRIMG),NRING'
C WRITE(6,*) CC,W,RADIUS(NRING),Z,HEIGHT(NRIMG),NRING
C WRITE(6,*)' RPRIME,R,SINTST,COSTST,THETAST,GAMMA(NRING),GMTO'
C WRITE(6,*) RPRIME,R,SINTST,COSTST,THETAST,GAMMA(NRING),GMTO
c***
    if(gmt0 .eq. 0.)gmt0=gamma(nring)
        DO 10 N=1,NCOEF
            if(gmt0 .eq. 0.)then
            rrr = 1.
        else
            RRR=GAMMA (NRING)/GMTO
                                    ladjust for varying tip strength
        endif
    TAUA=TAU2LOOP+(alplast-pi/2.)*(0.60702015)/(2.0*PI)
    BN=B(N,1)*TAUA**B(N,2)
    DN=D(N,1)*TAUA**D(N,2)
    DDVR=(BN*COS(N*(THETAST+PI/2.0-(alplast-pi/2.)))
    $ -DN*SIN(N*(THETAST+PI/2.0-(alplast-pi/2.))))
    $ *(N/R**(N+1))*RRR*(A**(4.0/3.0)*TIMENOW**(1.0/3.0)/BIGRPRM)
        DDVTH=(BN*SIN(N*(THETAST+PI/2.0-(alplast-pi/2.)))
    $ +DN*COS(N*(THETAST+PI/2.0-(alplast-pi/2.))))
    $ *(N/R**(N+1))*RRR*(A**(4.0/3.0)*TIMENO;***(1.0/3.0)/BIGRPRM)
C WRITE(6,*)' I,BETA(I),RRR,RR,DDVR,DDVTH'
C WRITE(6,*) I,BETA(I),RRR,RR,DDVR,DDVTH
    VRPRM =VRPRM+DDVR
10 VTHPRM=VTHPRM+DDVTH
    VXPRM=VRPRM*COSTST-VTHPRM*SINTST
    VYPRM=VRPRM*SINTST+VTHPRM*COSTST
C
ALF=alplast-pi/2.
IF (ALF .LT. 4.0*PI) THEN
        ZX=ZXO
        ZY=ZYO
ELSE IF (ALF .GT. 6.0*PI) THEN
        ZX=0.0
        ZY=0.0
```

```
    ELSE
        ZX=ZXO-ZXO*(ALF-4.0*PI)/(2.0*PI)
        ZY=ZYO-ZYO*(ALF-4.0*PI)/(2.0*PI)
    ALF=alplast-pi/2.
    ENDIF
    VXPRM=VXPRM*ZX
    VYPRM=VYPRM*ZY
C IF (IWR .LE. 9) THEN
C WRITE(6,*)' VXPRM,VYPRM',VXPRM,VYPRM
C WRITE(IWR,*)' TIMENOW,VXPRM,VYPRM',TIMENOW,VXPRM,VYPRM
C ENDIF
C
RETURN
END
SUBROUTINE ELLIP (RK, SUMK,SUME)
    common /temp1/a0,a1,a2,a3,a4,b0,b1,b2,b3,b4
    Common /temp2/c0,c1,c2,c3,c4,d0,d1,d2,d3,d4
C
C--Computes complete elliptic integrals, function of RK.
C Uses a four term series approximation.
C SUMK = K(RK) - FIRST KIND
C SUME = E(RK) - SECOND KIND
C
C
            FK=1.0-RK**2
            SUMK = AO+FK*(A1+FK*(A2+FK*(A3+FK*A4)))
            $ +(BO+FK*(B1+FK*(B2+FK*(B3+FK*B4))))*ALOG (1.0/FK)
            SUME = CO+FK*(C1+FK*(C2+FK*(C3+FK*C4)))
            $ +(DO+FK*(D1+FK*(D2+FK*(D3+FK*D4))))*ALOG(1.0/FK)
C
    RETURN
```


## END

```
    subroutine experiment
C
c a subroutine to set the experimental loading positions
C
    include '/user/huh/rollup/common.f'
    data expx/.99,.9733334,.9666666,.96,.95,.9533334,.926666,.91,.9,
& . 8993333,.8813334,.8733333,.8766666,.85,.8499667,.8266,
& .8,.8003334,.78,.773334,.7731334,.75,.7,.65,.623333,.6,
& .5733333,.55/
data expy/1.093333,1.36,1.6,1.653333,1.84,1.893333,1.986666,2.16,
& 2.1573335,2.213333,2.066666,2.0933334,2.0266669,1.8933,
& 1.8533333,1.76,1.72,1.762669,1.7186668,1.68,1.733333,
& 1.6266668,1.5733335,1.506668,1.493335,1.32,1.36,1.2933/
    do i=1,28
            expy(i) = expy(i) * . 01
        enddo
        return
        end
```

    SUBROUTINE GETANG
    C
C--This routine computes the angles of the segments relative to
C the horizontal axis
C
C
include '/user/huh/rollup/common.f'
COMMON /TH/ THETA(0:900),THMAX

DIMENSION DT (0:900)
C
C--Get increment for first angle
C
DT(0) = ATAN ( (HEIGHT (1)-HEIGHT (0)) /(RADIUS (1)-RADIUS (0)))
$\operatorname{THETA}(0)=D T(0)$
C
C--Get THETA for the rest
C
DO $10 \mathrm{I}=1$, NSMALL-1
$X M=R A D I U S(I-1)$
$X I=$ RADIUS (I)
$X P=R A D I U S(I+1)$
YM=HEIGHT (I-1)
YI=HEIGHT (I)
$\mathrm{YP}=\mathrm{HEIGHT}(\mathrm{I}+1)$
C
C--Get sign of cross product of Vi and $\mathrm{Vi}+1$
C
SCROSS $=S I G N(1.0,(X I-X M) *(Y P-Y I)-(X P-X I) *(Y I-Y M))$
VMAGS $=\operatorname{SQRT}(((X I-X M) * * 2+(Y I-Y M) * * 2) *((X P-X I) * * 2+(Y P-Y I) * * 2))$
$\mathrm{DOT}=(\mathrm{XI}-\mathrm{XM}) *(\mathrm{XP}-\mathrm{XI})+(\mathrm{YI}-\mathrm{YM}) *(\mathrm{YP}-\mathrm{YI})$
C
C--Watch out for numerical inaccuracy in argument of arccos
C
$D T(I)=0.0$
C WRITE(6,*)' I, DOT/VMAGS:',I,DOT/VMAGS
IF (ABS (DOT/VMAGS) .LE. 1.0) DT(I)=SCROSS*ACOS(DOT/VMAGS)
$\operatorname{THETA}(I)=\operatorname{THETA}(I-1)+D T(I)$
10 CONTINUE
C DO $22 I=0$, NSMALL-1
C 22 WRITE $(6, *)^{\prime}$ I,DT(I), THETA(I):',I,DT(I), THETA(I)
C
RETURN
END

```
    SUBROUTINE GETDT(DT)
            include '/user/huh/rollup/common.f'
C
C--Find max velocity and DT, DT=(ABS(DXMAX) if DXNAX .LT. 0.0)
C
            IF (DXMAX .LT. O.O) THEN
                    DT=ABS (DXMAX)
                ELSE
                VMAX=UVEL(0)**2+WVEL(0)**2
                DO 401 I=1,NRING
                    VRING=UVEL(I)**2+(WVEL(I))**2
                    IF (VRING .GT. VMAX) VMAX=VRING
    401 CONTINUE
                VMAX=SQRT (VMAX)
                DT=DXMAX/VMAX
            ENDIF
            RETURN
            END
            SUBROUTINE GETK(N)
                    include '/user/huh/rollup/common.f'
C
C--Computes the 4 K parameters used in the Runge-Kutta routine.
C The K's are just the velocity components.
C
    DO 10 I=0,NRING
        RK(I,N)=UVEL (I)
        HK(I,N)=WVEL(I)
```

10 CONTINUE

```
do 20 j = 1,level
do 20 i =1,ninter(j)
                rki(i,j,n) = uveli(i,j)
                hki(i,j,n) = wveli(i,j)
```

20
continue
C
RETURN
END
FUNCTION GETRAD(ILOC,R1,R2)
C
C--This function determines the location of the vortex ring within
C a segment, if $I L O C=1$ then the ring is centered, if $I L O C=2$ then the
C ring is positioned such that the total impulse is maintained
C
IF (ILOC .EQ. 1) THEN
GETRAD $=($ R1 1 R2 $) / 2.0$

## ELSE

$\mathrm{F} 1=((\operatorname{SQRT}(1.0-\mathrm{R} 1 * * 2)) * * 3-(\mathrm{SQRT}(1.0-\mathrm{R2} 2 * 2)) * * 3) * 8.0 / 3.0$
F2=4.0*R1**2*SQRT(1.0-R1**2)
F3=4.0*R2**2*SQRT(1.0-R2**2)
F4=4.0*(SQRT(1.0-R1**2)-SQRT(1.0-R2**2))
GETRAD $=\operatorname{SQRT}(($ F1+F2-F3)/F4)
ENDIF
C
RETURN
END

SUBROUTINE IMPULS(TIMP)
C
C--Computes total impulse of the vortex filament system for
C either axi or 2-D. (Note: only computing the "x" or "radial"

C component of the 2-D impulse.
C
include "/user/huh/rollup/common.f'
CVD\$R ASSOC
C
TIMP $=0.0$
DO $10 \mathrm{I}=1$,NRING
10
TIMP $=$ TIMP + PI $*$ GAMMA ( $I$ ) $*$ RADIUS (I) $* * 2$
C
RETURN
END

SUBROUTINE INDVEL
C
C AXISYMMETRIC FILAMENTS (RINGS)
C--This routine calculates the induced velocity at each ring position
C due to all the other rings and due to itself.
c--For all wake vorticies. (near,intermediate and far)
C
include '/user/huh/rollup/common.f'
REAL*4 LAMBDA, K, k1
real uvnear (900), wvnear (900), uvint $(10,5)$,wvint $(10,5)$
real uvfar(10),wvfar(10)
common /temp1/a0, a1, a2, a3, a4, b0, b1, b2, b3, b4
common /temp2/c0,c1,c2,c3,c4,d0,d1,d2,d3,d4
DATA AO, A1, A2, A3, A4, BO, B1, B2, B3, B4/
$\$ 1.38629436112, .09666344259, .03590092383$,
\$ . $03742563713, .01451196212, .5$
$\$ .12498593597, .06880248576, .03328355346$,
\$ .00441787012/
DATA CO,C1, C2,C3,C4,DO,D1,D2,D3,D4/
\$ $1.0 \quad . .44325141463, .06260601220$,

```
$ .04767383546, .01736506451, .0
$ .24998368310, .09200180037, .04069697526,
$ .00526449639/
```

CVD\$R ASSOC
C
C-Loop over all rings, except center ring, get induced velocity at each C

DO $10 \mathrm{I}=1$,NRING
W=RADIUS (I)
$\mathrm{Z}=\mathrm{HE}$ IGHT (I)
do $\mathrm{il}=1, \mathrm{nring}$
unnear (i1) $=0$.
wvnear(i1) $=0$.
enddo
C
C--Loop over all rings again (except control ring) and determine the C contribution of each to the induced velocity at ring $I$.
C
DO $20 \mathrm{~J}=1, \mathrm{i}-1$
WP=RADIUS (J)
ZP=HEIGHT (J)
$R 1=\operatorname{SQRT}((Z-Z P) * * 2+(W-W P) * * 2)$
$\operatorname{R2}=\operatorname{SQRT}((\mathrm{Z}-\mathrm{ZP}) * * 2+(W+W P) * * 2)$
LAMBDA $=($ R2 $2-R 1) /($ R2 $2+$ R1 $)$
FK=1.0-LAMBDA**2
$K=A 0+F K *(A 1+F K *(A 2+F K *(A 3+F K * A 4)))$
$\$ \quad+(B 0+F K *(B 1+F K *(B 2+F K *(B 3+F K * B 4))) * A L O G(1.0 / F K)$
$\mathrm{E}=\mathrm{CO}+\mathrm{FK} *(\mathrm{C} 1+\mathrm{FK} *(\mathrm{C} 2+\mathrm{FK} *(\mathrm{C} 3+\mathrm{FK} * \mathrm{C} 4)))$
\$ $\quad+(\mathrm{DO}+\mathrm{FK} *(\mathrm{D} 1+\mathrm{FK} *(\mathrm{D} 2+\mathrm{FK} *(\mathrm{D} 3+\mathrm{FK} * \mathrm{D} 4)))) * A L O G(1.0 / F K)$
$\mathrm{F} 1=((\mathrm{W}-\mathrm{WP}) / \mathrm{R} 1+(\mathrm{W}+\mathrm{WP}) / \mathrm{R} 2) *(\mathrm{~K}-\mathrm{E})$
$\mathrm{F} 2=2.0 * \mathrm{E} * \mathrm{LAMBDA} /(1.0-\mathrm{LAMBDA} * * 2)$
$\mathrm{F} 3=(\mathrm{W} *(\mathrm{R} 1 * * 2-\mathrm{R} 2 * * 2)+\mathrm{WP} *(\mathrm{R} 1 * * 2+\mathrm{R} 2 * * 2)) /(\mathrm{R} 1 * \mathrm{R} 2 *(\mathrm{R} 1+\mathrm{R} 2))$
$\mathrm{WV}=\mathrm{GAMMA}(\mathrm{J}) *(F 1+\mathrm{F} 2 * \mathrm{~F} 3) /(2.0 * \mathrm{PI} * W)$
WVnear ( j$)=$ WV
$\mathrm{F} 4=(\mathrm{Z}-\mathrm{ZP}) *(\mathrm{R} 1+\mathrm{R} 2) /(\mathrm{R} 1 * \mathrm{R} 2)$

```
                    F5=K-E*(1.0+LAMBDA**2)/(1.0-LAMBDA**2)
                    UV=GAMMA (J) *F4*F5/(2.0*PI*W)
UVnear(j)=-UV
```

$\mathrm{F} 1=((\mathrm{W}-\mathrm{WP}) / \mathrm{R} 1+(\mathrm{W}+\mathrm{WP}) / \mathrm{R} 2) *(\mathrm{~K}-\mathrm{E})$
$\mathrm{F} 2=2.0 * \mathrm{E} * \mathrm{LAMBDA} /(1.0-\mathrm{LAMBDA} * * 2)$
$\mathrm{F} 3=(\mathrm{W} *(\mathrm{R} 1 * * 2-\mathrm{R} 2 * * 2)+\mathrm{WP} *(\mathrm{R} 1 * * 2+\mathrm{R} 2 * * 2)) /(\mathrm{R} 1 * \mathrm{R} 2 *(\mathrm{R} 1+\mathrm{R} 2))$
$\mathrm{WV}=\mathrm{GAMMA}(\mathrm{J}) *(\mathrm{~F} 1+\mathrm{F} 2 * \mathrm{~F} 3) /(2.0 * \mathrm{PI} * \mathrm{~W})$
WVnear ( j ) $=$ WV
$\mathrm{F} 4=(\mathrm{Z}-\mathrm{ZP}) *(\mathrm{R} 1+\mathrm{R} 2) /(\mathrm{R} 1 * \mathrm{R} 2)$
$\mathrm{F} 5=\mathrm{K}-\mathrm{E} *(1.0+\mathrm{LAMBDA} * * 2) /(1.0-\mathrm{LAMBDA} * * 2)$
$\mathrm{UV}=\operatorname{GAMMA}(\mathrm{J}) * \mathrm{~F} 4 * \mathrm{~F} 5 /(2.0 * \mathrm{PI} * W)$
$\operatorname{UVnear}(\mathrm{j})=-\mathrm{UV}$
CONTINUE
CONTINUE
DO $40 \mathrm{~J}=\mathrm{i}+1$, nring
$W P=$ RADIUS ( J )
ZP=HEIGHT (J)
$\mathrm{R} 1=\mathrm{SQRT}((\mathrm{Z}-\mathrm{ZP}) * * 2+(\mathrm{W}-\mathrm{WP}) * * 2)$
R2 $2=\operatorname{SQRT}((Z-Z P) * * 2+(W+W P) * * 2)$
LAMBDA $=($ R2-R1 $) /($ R2 + R1 $)$
$\mathrm{FK}=1.0-\mathrm{LAMBDA} * * 2$
$K=A 0+F K *(A 1+F K *(A 2+F K *(A 3+F K * A 4)))$ $+(\mathrm{BO}+\mathrm{FK} *(\mathrm{~B} 1+\mathrm{FK} *(\mathrm{~B} 2+\mathrm{FK} *(\mathrm{~B} 3+\mathrm{FK} * \mathrm{~B} 4)))) * \mathrm{ALOG}(1.0 / \mathrm{FK})$
$\mathrm{E}=\mathrm{C} 0+\mathrm{FK} *(\mathrm{C} 1+\mathrm{FK} *(\mathrm{C} 2+\mathrm{FK} *(\mathrm{C} 3+\mathrm{FK} * \mathrm{C} 4)))$
$+(\mathrm{DO}+\mathrm{FK} *(\mathrm{D} 1+\mathrm{FK} *(\mathrm{D} 2+\mathrm{FK} *(\mathrm{D} 3+\mathrm{FK} * \mathrm{D} 4)))) * \operatorname{ALOG}(1.0 / \mathrm{FK})$

```
C
c
calculate influence of intermediate wake c
```

```
do 50 j1 = 1,level
```

do 50 j1 = 1,level
do 50 i1 = 1,ninter(j1)
do 50 i1 = 1,ninter(j1)
wp=rinter(i1,j1)
wp=rinter(i1,j1)
zp=zinter(i1,j1)
zp=zinter(i1,j1)
r1=sqrt((z-zp)**2+(w-wp)**2)
r1=sqrt((z-zp)**2+(w-wp)**2)
r2=sqrt((z-zp)**2+(w+wp)**2)
r2=sqrt((z-zp)**2+(w+wp)**2)
lambda=(r2-r1)/(r2+r1)

```
    lambda=(r2-r1)/(r2+r1)
```

```
        fk=1.0-1ambda**2
        k=a0+fk*(a1+fk*(a2+fk*(a3+fk*a4)))
    $ +(b0+fk*(b1+fk*(b2+fk*(b3+fk*b4))))*alog(1.0/fk)
        e = c0+fk*(c1+fk*(c2+fk*(c3+fk*c4)))
            +(d0+fk*(d1+fk*(d2+fk*(d3+fk*d4))))*alog(1.0/fk)
        11=((w-wp)/r1+(w+wp)/r2)*(k-e)
        f2=2.0*e*lambda/(1.0-lambda**2)
        I3=(w*(r1**2-r2**2)+wp*(r1**2+r2**2))/(r1*r2*(r1+r2))
        wv=ginter(i1,j1)*(f1+f2*f3)/(2.0*pi*w)
        wvint(i1,j1)=wv
        f4=(z-zp)*(r1+r2)/(r1*r2)
        f5=k-e*(1.0+lambda**2)/(1.0-1ambda**2)
        uv=ginter(i1,j1)*f4*f5/(2.0*pi*w)
        uvint(i1,j1)=-uv
        continue
c
c calculate influence of far wake cylinder
C
    dphi = pi/20. !40 integration steps
    nphi = nint(2.*pi/dphi) - 1
    do 67 ifar = 1,nfar
    zp = abs(z - zfar(ifar))
    wsum = 0.
    do 60 iphi = 1,nphi
        phi = float(iphi)*dphi
        x3 = rfar(ifar)*rfar(ifar)+w*w+zp*zp-
    &
        dwfar = dphi*rfar(ifar)*(rfar(ifar)-w*cos(phi))/
                                    (x3-zp*zp)*(1.-zp/sqrt(x3))
        wsum = wsum + dwfar
        continue
    x3 = rfar(ifar)*rfar(ifar)+w*w+zp*zp-2.*rfar(ifar)*w
    if((x3-zp*zp).ne. 0.)then
        dwfar = dphi*rfar(ifar)*(rfar(ifar)-w)/(x3-zp*zp)*(1.-zp/sqrt(x3))
        wsum = wsum + dwfar
```

```
            endif
            wfar = wsum/4./pi*gfar(ifar)
            wvfar(ifar) = wfar
    k1 = 4.*w*rfar(ifar)/((w+rfar(ifar))**2+zp*zp)
        if (kl .eq. 1.)then
            k=10.
            e=1.
            goto 62
        endif
        f1 = alog(4./sqrt(1.-k1))
        e=1.+.5*(f1-.5)*(1.-k1)+3./16.*(f1-1.-1./12.)*(1.-k1)**2
    &
                +0.1171875*(f1-1.-1./6.-1./30.)*(1. -k1)**3
            k = f1 + . 25*(f1-1.)*(1.-k1) + 9./64*(f1-1.-1./6.)*(1.-k1)**2
            &
                +0.09765625*(&1-1.-1./6.-1./15.)*(1.-k1)**3
            continue
    ufar = (sqrt(rfar(ifar)/w/k1))*(k*(2.-k1)-2.*e)*.5/pi
    ufar = abs(ufar)*gfar(ifar)
    uvfar(ifar) = ufar
67
            continue
    uvel(i) = 0.
    wvel(i) = 0.
    do i1 = 1,nring
        uvel(i) = uvel(i) + uvnear(ii)
        wvel(i) = wvel(i) + wvnear(i1)
    enddo
    do j1 = 1,level
        uvel(i) = uvel(i) + uvint(1,j1)+uvint(2,j1)+uvint(3,j1)
        wvel(i) = wvel(i) + wvint(1,j1)+wvint(2,j1)+wvint(3,j1)
    enddo
    uvel(i) = uvel(i) + uvfar(1) + uvfar(2) +uvfar(3)
    wvel(i) = wvel(i) + wvfar(1) + wvfar(2) +wvfar(3)
c
C--Calculate final velocity including self-induced component
C
        WV=GAMMA (I)*(ALOG (8.0*W/CORE(I))-.25)/(4.0*PI*W)
```

$$
\text { WVEL }(I)=\text { WVEL }(I)+W V
$$

C
C--get velocity of higher order terms in tip regicn, if not tip ring now C

IF (I .NE. NRING) THEN

CALL DOUBLET (W,Z,VXPRM,VYPRM,VRPRM,VTHPRM)
WVEL ( $I$ ) =WVEL ( $I$ ) +VYPRM
$\operatorname{UVEL}(I)=$ UVEL $(I)+V X P R M$
ENDIF
C
10 CONTINUE
C
C--velocity at the intermediate and far wakes
C

```
        DO 101 I=1,level
```

        do 101 ii = 1 , ninter \((i)\)
            \(W=r i n t e r(i i, i)\)
            \(\mathrm{Z}=\) zinter (ii,i)
            do \(\mathrm{j} 1=1\), level
            do \(i 1=1, \operatorname{ninter}(j 1)\)
                    uvint \((i 1, j 1)=0\).
                    wvint(i1,j1) \(=0\).
                enddo
            enddo
    C
C-Loop over all rings again (except control ring) and determine the
$C$ contribution of each to the induced velocity at intermediate wake ii,i C

```
    DO 201 J=1,nring
```

        WP=RADIUS (J)
        ZP=HEIGHT (J)
        \(\mathrm{R} 1=\mathrm{SQRT}((\mathrm{Z}-\mathrm{ZP}) * * 2+(W-W P) * * 2)\)
        \(R 2=S Q R T((Z-Z P) * * 2+(W+W P) * * 2)\)
        LAMBDA \(=(\) R2-R1 \() /(\) R2 + R1 \()\)
        FK=1.0-LAMBDA**2
    ```
            K=AO+FK*(A1+FK*(A2+FK*(A3+FK*A4)))
        F4=(Z-ZP)*(R1+R2)/(R1*R2)
        F5=K-E*(1.0+LAMBDA**2)/(1.0-LAMBDA**2)
        UV=GAMMA (J)*F4*F5/(2.0*PI*W)
        UVnear(j)=-UV
        CONTINUE
```

$\$$ $\mathrm{E}=\mathrm{CO}+\mathrm{FK} *(\mathrm{C} 1+\mathrm{FK} *(\mathrm{C} 2+\mathrm{FK} *(\mathrm{C} 3+\mathrm{FK} * \mathrm{C} 4)))$
$\$$ $\mathrm{F} 1=((\mathrm{W}-\mathrm{WP}) / \mathrm{R} 1+(\mathrm{W}+\mathrm{WP}) / \mathrm{R} 2) *(\mathrm{~K}-\mathrm{E})$ F2 $=2.0 * E * L A M B D A /(1.0-L A M B D A * * 2)$ $\mathrm{F} 3=(W *(R 1 * * 2-R 2 * * 2)+W P *(R 1 * * 2+R 2 * * 2)) /(R 1 * R 2 *(R 1+R 2))$ WV $=\operatorname{GAMMA}(J) *(F 1+F 2 * F 3) /(2.0 * P I * W)$ WVnear $(j)=W V$
c C

```
calculate influence of intermediate wake
    do 501 j1 = 1,level
    do 501 i1 = 1,ninter(j1)
        if(i1 .eq. ii .and. j1 .eq. i)goto501
            wp=rinter(i1,j1)
            zp=zinter(i1,j1)
            r1=sqrt((z-zp)**2+(w-wp)**2)
            r2=sqrt((z-zp)**2+(w+wp)**2)
            lambda=(r2-r1)/(r2+r1)
            fk=1.0-1ambda**2
            k=a0+fk*(a1+fk*(a2+fk*(a3+fk*a4)))
        +(b0+fk*(b1+fk*(b2+fk*(b3+fk*b4))))*alog(1.0/fk)
                e = c0+fk*(c1+fk*(c2+fk*(c3+fk*c4)))
                +(d0+fk*(d1+fk*(d2+fk*(d3+fk*d4))))*alog(1.0/fk)
                f1=((w-wp)/r1+(w+wp)/r2)*(k-e)
                f2=2.0*e*lambda/(1.0-1ambda**2)
                f3=(w*(r1**2-r2**2)+wp*(r1**2+r2**2))/(r1*r2*(r1+r2))
                wv=ginter(i1,j1)*(f1+f2*f3)/(2.0*pi*w)
                wvint(i1,j1)=wv
```

```
f4=(z-zp)*(r1+r2)/(r1*r2)
f5=k-e*(1.0+lambda**2)/(1.0-1ambda**2)
uv=ginter(i1,j1)*f4*f5/(2.0*pi*w)
uvint(i1,j1)=-uv
```

    501
        continue
    c
c calculate influence of far wake cylinder
C

```
    dphi = pi/20. !40 integration steps
    nphi = nint(2.*pi/dphi) - 1
    do 671 ifar = 1,nfar
    zp = abs(z - zfar(ifar))
    wsum = 0.
    do 601 iphi = 1,nphi
        phi = float(iphi)*dphi
        x3 = rfar(ifar)*rfar(ifar) +w*w+zp*zp-
        * 2.*rfar(ifar)*w*cos(phi)
        dwfar = dphi*rfar(ifar)*(rfar(ifar)-w*cos(phi))/
                                    (x3-zp*zp)*(1.-zp/sqrt(x3))
            wsum = wsum + dwfar
        continue
        x3 = rfar(ifar)*rfar(ifar) +w*w+zp*zp-2.*rfar(ifar)*w
        if((x3-zp*zp) .ne. 0.)then
            dwfar = dphi*rfar(ifar)*(rfar(ifar)-w)/(x3-zp*zp)*(1.-zp/sqrt(x3))
            wsum = wsum + dwfar
        endif
        wfar = wsum/4./pi*gfar(ifar)
        wvfar(ifar) = wfar
        k1 = 4.*w*rfar(ifar)/((w+rfar(ifar))**2+zp*zp)
        if (kl .eq. 1.)then
            k = 10.
            e = 1.
            goto 621
        endif
        f1 = alog(4./sqrt(1.-k1))
```

        \&
    601

```
            e = 1. + . 5*(11-. 6)*(1.-k1) +3./16.*(11-1.-1./12.)*(1.-k1)**2
    &
                        +0.1171875*(f1-1.-1./6.-1./30.)*(1.-k1)**3
    k=f1 + . 25*(f1-1.)*(1.-k1) + 9./64*(f1-1.-1./6.)*(1.-k1)**2
    &
621
    continue
    ufar = (sqrt(rfar(ifar)/w/k1))*(k*(2.-k1)-2.*e)*.5/pi
    ufar = abs(ufar)*gfar(ifar)
    uvfar(ifar) = ufar
671
            continue
    uveli(ii,i) = 0.
    wveli(ii,i) = 0.
    do i1 = 1,nring
        uveli(ii,i) = uveli(ii,i) + uvnear(i1)
        wveli(ii,i) = wveli(ii,i) + wvnear(i1)
    enddo
    do j1 = 1,level
        uveli(ii,i)=uveli(ii,i)+uvint(1,j1)+uvint(2,j1)+uvint(3,j1)
        wveli(ii,i)=wveli(ii,i)+wvint(1,j1)+wvint(2,j1)+wvint(3,j1)
    enddo
    uveli(ii,i) = uveli(iii,i) + uvfar(1) + uvfar(2) +uvfar(3)
    wveli(ii,i) = wveli(ii,i) + wvfar(1) + wvfar(2) +wvfar(3)
C
C--Calculate final velocity including self-induced component
C
    WV=Ginter(ii,I)*(ALOG(8.0*W/Cinter(ii,I))-.25)/(4.0*PI*W)
    WVELi(ii,I)=WVELi(ii,I)+WV
C
C--get velocity of higher order terms in tip region, if not tip ring now
C
CALL DOUBLET (W, Z, VXPRM, VYPRM, VRPRM, VTHPRM)
                                    WVELi(ii,I)=WVELi(ii,I)+VYPRM
            UVELi(ii,I)=UVELi(ii,I)+VXPRM
C
    101 CONTINUE
C
```

```
C--Get velocity contribution at tip due to inner spiral
C
    CALL CENSPIR(VXP,VYP)
    UVEL(NRING)=UVEL (NRING)+VXP
    WVEL(NRING) =WVEL (NRING) +VYP
C
C--Get induced velocity at center by extrapolating the velocities of
C points 1 and 2. Fit a parabola (for the velocity) through points
C 0, 1, and 2 by matching the values at points 1 and 2 and setting
C the slope=0 at r=0. Then the velocity at r=0 (i.e., the center
C point) is just the constant part of the parabola.
C
    UVEL(0)=0.0
            WVEL (0) = (WVEL (1) *RADIUS (2) **2-WVEL (2) *RADIUS (1) **2)
        $ /(RADIUS(2)**2-RADIUS(1)**2)
C
C--done
C
RETURN
END
```

```
        subroutine initcirc(frac)
C
c a subroutine to determine the initial circulation distribution
c from Miller
c also sets blade geometry.
C
    include '/user/huh/rollup/common.f'
    real circt(103),r(103)
    real lambda,k
    data aslope,sigma/6.15,.0464/
CVD$R ASSOC
c
    nspan = 103
    open(unit=21,file='initload.dat',status='old',form='formatted')
    do i = 1,103
        read(21,*)r(i),circt(i)
        span(i) = r(i)
            enddo
            close(21)
            do 10 i = 1,nspan
        pitch(i) = 0.123 + 0.19*(1.-span(i))
        chord(i) = . 07288495
1 0
        continue
c
            nsmall = nring
            call initsheet(0.)
            dr = .9/float(nsmall)
            do 40 i = 1,nsmall
            r1 = float(i-1)*dr+.1
            r2 = float(i)*dr+. 1
            circ1 = spl1(103,r,circt,r1)
            circ2 = spl1(103,r,circt,r2)
            gamma(i) = circ2 - circ1
4 0
        continue
```

```
call toteng(tenerg1)
c
c calculate initial vortex sheet strength
c
        nsmall = nsmall -1
            call initsheet(frac)
            dr = (.9-frac)/float(nsmall)
            do 50 i = 1,nsmall
                    r1 = float(i-1)*dr+.1
                    r2 = float(i)*dr+.1
                    circ1 = spl1(103,r,circt,r1)
                    circ2 = spl1(103,r,circt,r2)
                    gamma(i) = circ2 - circ1
50 continue
            HEIGHT (NRING)=0.825*frac
            redg = 1.-frac
            radius(nring) = 1.-(0.52*frac)
            gamma(nring) = -spl1(103,r,circt,redg)
            call rke(tenerg2,nring,nring)
            sum3 = 0.
            DO 80 I=1,NSMALL
            R1=SQRT((HEIGHT (I)-height (nring))**2+
            * (RADIUS(I)-Radius(nring))**2)
            R2=SQRT((HEIGHT(I)-Height(nring))**2+
            & (RADIUS(I)+Radius(nring))**2)
            LAMBDA=(R2-R1)/(R2+R1)
            CALL ELLIP(LAMBDA,K,E)
            sum3=sum3+GAMMA (I)*gamma(nring)*(R1+R2)*(K-E)/2.0
                    8 0 ~ C O N T I N U E ~
            senergy = tenerg1 - tenerg2-2.*sum3
            czzz = exp(2.*senergy/(gamma(nring)**2*radius(nring))+1.75)
            core(nring) = 8.*radius(nring)/czzz
C
C--Get volume of tip core
C
```

```
        VOLTIP=2.0*PI*PI*RADIUS(NRING)*CORE(NRING)**2
C
c--Center ring
C
        RADIUS(0)=0.1
        HEIGHT (0) =0.0
        DO 8315 I=0,NRING
    8315 HEIGHT (I) = HEIGHT (I)-1.0*TIME (1)
        GMTO=GAMMA (NRING)
        do 19 i = 1,nspan
            circold(i) = spl1(103,r,circt,span(i))
19 continue
        do i = 1,nspan
            circ(i) = circold(i)
        enddo
        write(7)(radius(i), i=0,nring)
        write(7)(height(i),i=0,nring)
        write(7)(core(i), i=0,nring)
        write(7)(gamma(i),i=0,nring)
        write(7)(axlseg(i),i=0,nring)
        return
        end
            subroutine initsheet(frac)
c
c
    sets the strength and the location of the trailing vortex sheet
    C
    include '/user/huh/rollup/common.f'
    COMMON /E2/ DUM1,DUM2,DUM3,DUM4,DUM5,AE2
    data iloc/1/
    XO=FRAC
    AO=0.520*XO
```

$B O=0.825 * X 0$
C
C--Set core size
C
CCORE $=(1.0 /$ NSMALL $) / 4.0$
C
C--Compute radius, height, and ring circulation for constant spacing C
$D R=(.9-F R A C) / f$ loat (NSMALL)
expand $=1$.
DO $10 \mathrm{I}=1$, NSMALL
HEIGHT ( $I$ ) $=0.0$
CORE ( $I$ ) =CCORE
$\mathrm{R} 1=(\mathrm{I}-1) * \mathrm{DR}+.1$
$\mathrm{R} 2=\mathrm{I} * \mathrm{DR}+.1$
RADIUS (I) =EXPAND*GETRAD (ILOC , R1,R2)
10 CONTINUE
C
C-Initialize the arc length array, AXLSEG(I) is the total arc length
$C$ to the end of segment $I$. (Note: segment $I$ is the segment which
C contains ring or panel I)
C
$\operatorname{AXLSEG}(0)=0.0$
DO $11 \mathrm{I}=1$, NSMALL-1
$\operatorname{AXLSEG}(I)=\operatorname{AXLSEG}(I-1)$
\$ +SQRT( (RADIUS (I+1)-RADIUS (I))**2
\$ +(HEIGHT(I+1)-HEIGHT (I))**2)
11 CONTINUE
AXLSEG (NSMALL) $=$ AXLSEG (NSMALL -1$)$
\$ +SQRT( (RADIUS (NSMALL) -RADIUS (NSMALL-1)) **2
\$ +(HEIGHT(NSMALL)-HEIGHT (NSMALL-1)) $* * 2$ )
return
end

```
        subroutine initwake
C
c sets the initial intermediate and far wake fron Miller's output
C
        include '/user/huh/rollup/common.f'
CVD$R ASSOC
level = 4
nfar = 2
do 134 i = 1,level
                ninter(i) = 3
    134
        continue
C
c root vortex
        zinter(1,1) = -0.147389
        zinter(1,2) = -0.329837
        zinter(1,3) = -0.524966
        zinter(1,4) = -0.720314
        rinter(1,1) = 0.414289
        rinter(1,2) = 0.41414
        rinter(1,3) = 0.405735
        rinter(1,4) = 0.392395
        ginter(1,1:4)=0.00970935
c
c center vortex
c
        zinter(2,1) = -0.245905
        zinter(2,2) = -0.482224
        zinter(2,3) = -0.7527273
        zinter(2,4) = -0.9638
        rinter(2,1) = 0.718209
        rinter(2,2) = 0.6980455
        rinter(2,3) = 0.688636
        rinter(2,4) = 0.672146
```

```
ginter(2,1:4)=0.0059792
```

C
.
tip vortex

```
    zinter(3,1) = -0.0612717
    zinter(3,2) = - 0.211619
    zinter(3,3) = -0.3981818
    zinter(3,4) = -0.572313
    rinter(3,1) = 0.872496
    rinter(3,2) = 0.797164
    rinter (3,3) = 0.7789773
    rinter(3,4) = 0.769287
    ginter(3,1:4) = - 0.020662807
```

    \(\$\)
                zinter (ninter(level-1), level-1))
    \(\operatorname{rfar}(n f a r)=\operatorname{rinter}(\) ninter (level), level)
    zfar(nfar) \(=\) zinter(ninter(level), level)-zdis
    \(\operatorname{gfar}(n f a r)=\) ginter(ninter(level), level)/zdis
    hsum1 \(=0\).
    zsum1 \(=0\).
    do \(60 \mathrm{i}=1\), ninter(level-1)-1
            hsum1 = hsum1 +
    \& ginter(i,level-1)*rinter(i,level-1)*rinter(i,level-1)
            zsum1 = zsum1 +
    \& \(\quad \operatorname{ginter}(i, l e v e l-1) * z i n t e r(i, l e v e l-1) *\) rinter (i,level-1)**2
            continue
    \(z 1=z s u m 1 / h s u m 1\)
    hsum2 \(=0\).
    zsum2 \(=0\).
    gsum2 \(=0\).
    do \(70 \mathrm{i}=1\), ninter(level)-1
        hsum2 = hsum2 +
    ```
    &
    &
70
c
c set core sizes
c
    do i=1,3
        cinter(i,1) = rinter(i,1)*.025
        volinter(i,1) = 2.*pi*pi*rinter(i,1)*cinter(i,1)**2
    enddo
    do j = 2,level
        do i = 1,ninter(j)
            volinter(i,j) = volinter(i,1)
                cinter(i,j) = sqrt(volinter(i,j)/(2.*rinter(i,j)))/pi
        enddo
    enddo
    do j = 1,level
        do i = 1, ninter(j)
            roldi(i,j) = rinter(i,j)
            zoldi(i,j) = zinter(i,j)
            goldi(i,j) = ginter(i,j)
        enddo
    enddo
    return
    end
```

subroutine itercirc (wwake, circt)
c
c a subroutine to iterate for bound circulation
C
parameter (maxitrs $=5000$ )
include '/user/huh/rollup/common.f'
common/temp1/a0, a1, a2, a3, a4, b0, b1, b2, b3, b4
common /temp2/c0,c1,c2,c3,c4,d0,d1,d2,d3,d4
real wwake (250), circt(250)
real $g(250), r(250)$, wtotal(250), cold(250), cnew(250)
logical go
data epsilon,omega/0.001,0.4/
data go/.true./
CVD\$R ASSOC
do $i=1$, nspan
if(span(i).lt. 0.11 .or. span(i) .gt. 0.985 )then
cold(i) $=$ circold $(i)$
else
alphaeff $=$ pitch $(i)+$ wwake (i) $/$ span(i)
cold(i) $=0.5 *$ chord $(i) * a s l o p e * s p a n(i) * a l p h a e f f$
endif
enddo
do $i=1, n s p a n-1$
$r(i)=0.5 *(\operatorname{span}(i)+\operatorname{span}(i+1))$
$g(i)=\operatorname{cold}(i+1)-\operatorname{cold}(i)$
enddo
itrs $=1$
do while(go .and. itrs .lt. maxitrs)
go $=. f a l s e$.
DO $I=2$,nspan
$W=\operatorname{span}(i)$
$\mathrm{Z}=0$.
wvelt $=0$.
DO $\mathrm{J}=2$, nspan-1
$W P=r(J)$

```
        ZP=0.
        R1=SQRT((Z-ZP)**2+(W-WP)**2)
        R2=SQRT((Z-ZP)**2+(W+WP)**2)
            LAMBDA=(R2-R1)/(R2+R1)
        FK=1.0-LAMBDA**2
        K = AO+FK*(A1+FK*(A2+FK*(A3+FK*A4)))
        +(B0+FK*(B1+FK*(B2+FK*(B3+FK*B4))))*ALOG(1.0/FK)
        E = C0+FK*(C1+FK*(C2+FK*(C3+FK*C4)))
        +(DO+FK*(D1+FK*(D2+FK*(D3+FK*D4))))*ALOG(1.0/FK)
        F1=((W-WP)/R1+(W+WP)/R2)*(K-E)
        F2=2.0*E*LAMBDA/(1.0-LAMBDA**2)
        F3=(W*(R1**2-R2**2)+WP*(R1**2+R2**2))/(R1*R2*(R1+R2))
            WV=G(J)*(F1+F2*F3)/(2.0*PI*W)
        wvelt = wvelt + wv
    enddo
    wvelt = 0.5*wvelt
    wtotal(i) = wwake(i) + wvelt
enddo
do i = 2,nspan
    if(span(i) .le. 0.11 .or. span(i) .ge. 0.985)then
        cnew(i) = cold(i)
    else
        alphaeff = pitch(i) + wtotal(i)/span(i)
        ctemp = 0.5*chord(i)*aslope*span(i)*alphaeff
        cnew(i) = cold(i) + omega*(ctemp-cold(i))
    endif
enddo
ispan = 1
do while((.not. go) .and. ispan .lt. nspan)
    if(cnew(ispan) .gt. 1.e-05)then
            test = (cnew(ispan)-cold(ispan))/cnew(ispan)
        if(abs(test) .ge. epsilon)go = .true.
    endif
    ispan = ispan + 1
enddo
```

```
            do i = 1,nspan-1
                g(i) = cnew(i+1) - cnew(i)
                cold(i) = cnew(i)
            enddo
            cold(nspan) = cnew(nspan)
            itrs = itrs + 1
enddo
do i = 1,nspan
                circt(i) = cnew(i)
    enddo
    return
    end
    SUBROUTINE LUMPS
C
C--This subroutine looks at the rolled up part of the sheet and
C determines if part of it should be truncated and lumped into
C the tip vortex ring.
C
C
    include '/user/huh/rollup/common.f'
    COMMON /TH/ THETA(0:900),THMAX
    REAL*4 LAMBDA,K
        real en2i(900), en2(900), en3(900)
CVD$R ASSOC
CVD$R CNCALL
C
C-- Get segment angles relative to horizintal axis
C
    CALL GETANG
C
C--See if any rings exceed the cut-off angle THMAX
C
```

```
            M=0 ! M will contain the # of rings to be lumped
            I=NSMALL-1 ! Highest segment number
    20 IF (THETA(I) .LT. THMAX) GO TO 10
    M=M+1
    I=I-1
    GO TO 20
C
C--Now M contains the # of rings to be lumped into the core
C
    10 IF (M .EQ. 0) GO TO 999 !Return
C
            IF (IWR .EQ. 11) THEN
                IF (IAXI .EQ. 1) CALL TOTENG(TTKE)
                CALL IMPULS(TTIM)
            WRITE(11,4387)TTKE, TTIM
    4387 FORMAT(' START LUMPS: TTKE,TTIM',2F12.6)
                ENDIF
C
C--Have M rings to lump, first find new position of tip by conserving
C impulse.
C
C GTIP = new tip circulation
C RTIP = new tip radius
C HTIP = new tip height
C
            GTIP=0.0
            DO 30 I=NSMALLL-M+1,NRING
        30 GTIP=GTIP+GAMMA (I)
C
                    RT2=0.0
                    DO 40 I=NSMALL-M+1,NRING
    40 RT2=RT2+GAMMA(I)*RADIUS (I)**2
            RTIP=SQRT(RT2/GTIP)
            HTIP=0.0
            DO 50 I=NSMALL-M+1,NRING
```

```
    50 HTIP=HTIP+GAMMA (I)*HEIGHT(I)*RADIUS(I)**2
        HTIP=HTIP/RT2
        CTIP=0.0
C
C Only get new core size for the axi case
C--Get new tip core size by conserving kinetic energy.
C First compute total kinetic energy of system.
C
    CALL TOTENG(TKE)
C------------------------------------------------------------------------
C---DEBUG----------------------------------------------------------------
c WRITE(6,*)' M,GTIP,RTIP,HTIP,TKE',M,GTIP,RTIP,HTIP,TKE !
C-----------------------------------------------------------------------------
C
C--Now get the various sums needed to compute TTBAR
C
            SUM1=0.0
            DO 60 I=1,NSMALL-M
            TIBAR=GAMMA(I)**2*RADIUS(I)*(ALOG(8.0*RADIUS(I)/CORE (I))
            $ -1.75)/2.0
        60 SUM1=SUM1+TIBAR
C
            DO 70 I=1,NSMALL-M
                do j = 1,nsmall-m
                    en2(j) = 0.
            enddo
            if(i .eq. 1)goto73
            DO 72 J=1,i-1
                        R1=SQRT((HEIGHT (I) -HEIGHT(J))**2+(RADIUS(I)-RADIUS (J))**2)
                    R2=SQRT((HEIGHT (I) -HEIGHT(J))**2+(RADIUS(I)+RADIUS(J))**2)
                    LAMBDA=(R2-R1)/(R2+R1)
                    CALL ELLIP(LAMBDA,K,E)
                    en2(j)=GAMMA(I)*GAMMA (J)*(R1+R2)*(K-E)/2.0
        72 CONTINUE
    73 continue
```

```
                    if(i .eq. (nsmall-m))goto75
            DO 74 J=i+1,NSMALL-M
                R1=SQRT((HEIGHT (I)-HEIGHT(J))**2+(RADIUS(I)-RADIUS(J))**2)
                R2=SQRT((HEIGHT(I)-HEIGHT(J))**2+(RADIUS(I)+RADIUS (J))**2)
                LAMBDA= (R2-R1)/(R2+R1)
                    CALL ELLIP(LAMBDA,K,E)
                    en2(j)=GAMMA (I)*GAMMA (J)*(R1+R2)*(K-E)/2.0
    7 4 \text { CONTINUE}
75 continue
        en2i(i) = 0.
        do j = 1,nsmall-m
            en2i(i) = en2i(i) + en2(j)
        enddo
70 continue
        SUM2=0.0
            do i = 1,nsmall-m
            sum2 = sum2 + en2i(i)
        enddo
C
    DO 80 I=1,NSMALL-M
            R1=SQRT((HEIGHT(I)-HTIP)**2+(RADIUS(I)-RTIP)**2)
            R2=SQRT((HEIGHT (I) -HTIP)**2+(RADIUS(I) +RTIP)**2)
            LAMBDA= (R2-R1)/(R2+R1)
            CALL ELLIP(LAMBDA,K,E)
            en3(i)=GAMMA(I)*GTIP*(R1+R2)*(K-E)/2.0
    8 0 ~ C O N T I N U E ~
        SUM3=0.0
            do i = 1,nsmall-m
            sum3 = sum3 + en3(i)
            enddo
C
            SUM4=SUM3
C
C--Get energy contribution of tip.
C
```

```
    TTBAR=TKE-SUM1-SUM2-SUM3-SUM4
C
C--Now get core size corrosponding to this energy.
C
    CZZZ=EXP(2.0*TTBAR/(GTIP**2*RTIP)+1.75)
    CTIP=8.0*RTIP/CZZZ
C
C--Put tip variables into main arrays
C
    RADIUS(NRING)=RTIP
    HEIGHT'(NRING)=HTIP
    GAMMM(NRING)=GTIP
    CORE (NRING)=CTIP
C
C--Compute new volume of tip ring
C
    VOLTIP=2.0*PI*PI*RADIUS(NRING)*CORE(NRING)**2
C
C--Now rediscretize remaining portion of sheet.
C
    IF (IWR .EQ. 11) THEN !this only good for M=1
        RADIUS(NRING-1)=RADIUS(NRING)
            HEIGHT (NRING-1)=HEIGHT (NRING)
            GAMMA (NRING-1)=GAMMA (NRING)
            CORE(NRING-1)=CORE(NRING)
            NRING=NRING-1
            CALL TOTENG(TTKE)
            CALL IMPULS(TTIM)
            NRING=NRING+1
            WRITE(11,4388)TTKE,TTIM
4388 FORMAT(' LUMPS BEFORE REDISC, TTKE,TTIM',2F12.6)
            ENDIF
            NIN=NSMALL-M
            NOUT=NSMALL
c
GMM=GMAX-GAMMA (NRING)
```

GMM=GAMMA (NRING)
CALL REDISC(RADIUS, HEIGHT , GAMMA , NIN, NOUT , GMM, SUML, DELTAL)
IF (IWR .EQ. 11) THEN
IF (IAXI .EQ. 1) CALL TOTENG(TTKE)
CALL IMPULS(TTIM)
WRITE (11, 4389)TTKE, TTIM
4389 FORMAT(' LUMPS AFTER REDISC , TTKE,TTIM', 2F12.6)
ENDIF
C

C--done
C
999 RETURN
END
subroutine makesheet(frac,istep)
c
c a subroutine to create a new near wake
C
include '/user/huh/rollup/common.f'
real lambda,k
CVD\$R ASSOC
cCVD\$R CNCALL
nsmall = nring
call initsheet(0.)
$d r=.9 / f l o a t$ (nsmall)
do $40 \mathrm{i}=1$, nsmall
$r 1=$ float(i-1)*dr+. 1
$r 2=$ float(i)*dr+. 1
$\operatorname{circ} 1=\operatorname{spl1}($ nspan, span, circ, r1)
circ2 = spl1 (nspan, span, circ, r2)
$\operatorname{gamma}(i)=\operatorname{circ} 2-c i r c 1$
40
continue

```
            call toteng(tenerg1)
            nsmall = nsmall -1
            call initsheet(frac)
            dr = (.9-frac)/float(nsmall)
            do 50 i = 1,nsmall
                    r1 = float(i-1)*dr+.1
                    r2 = float(i)*dr+.1
                    circ1 = spl1(nspan,span,circ,r1)
                    circ2 = spl1(nspan, span,circ,r2)
                    gamma(i) = circ2 - circ1
5 0
                    continue
        HEIGHT(NRING)=0.825*frac
        redg = 1.-frac
        radius(nring) = 1.-0.52*frac
        gamma(nring) = -spll(nspan,span,circ,redg)
        call rke(tenerg2,nring,nring)
        sum3 = 0.
        DO 80 I=1,NSMALL
        R1=SQRT((HEIGHT(I)-height(nring))**2+
            * (RADIUS(I)-Radius(nring))**2)
    R2=SQRT((HEIGHT(I)-Height(nring))**2+
    * (RADIUS(I)+Radius(nring))**2)
    LAMBDA=(R2-R1)/(R2+R1)
    CALL ELLIP(LAMBDA,K,E)
    sum3=sum3+GAMMA(I)*gamma(nring)*(R1+R2)*(K-E)/2.0
    80 continue
    senergy = tenerg1 - tenerg2-2.*sum3
    czzz = exp(2.*senergy/(gamma(nring)**2*radius(nring))+1.75)
    core(nring) = 8.*radius(nring)/czzz
C
C--Get volume of tip core
C
    VOLTIP=2.0*PI*PI*RADIUS(NRING)*CORE(NRING)**2
C
C--Center ring
```

C

```
    RADIUS(0)=0.1
    HEIGHT(0)=0.0
    DO }8315\textrm{I}=0,NRIN
8315 HEIGHT(I)=HEIGHT(I)-1.0*TIME (1)
    GMTO=GAMMA(NRING)
time(istep) = time(istep) + time(1)
close(7)
open(unit=7,file='sheettemp.dat',status='new',form='unformatted')
write(7)(radius(i),i=0,nring)
write(7)(height(i), i=0,nring)
write(7)(core(i),i=0,nring)
write(7)(gamma(i),i=0,nring)
write(7)(axlseg(i),i=0,nring)
return
end
```

subroutine movem(istep)
include '/user/huh/rollup/common.f'
CVD\$R ASSOC
C
C--Computes the new wake positions using the
C Runge-Kutta time stepping
C
C--First save original data
C

ALPO=ALPLAST
DO 20 I=0,NRING
RADO(I) $=$ RADIUS (I)
HEIO (I) = HEIGHT (I)
20 continue
do $25 \mathrm{j}=1$ 1.level

```
                                    do 25 i = 1,ninter(j)
                                    rintO(i,j)=rinter(i,j)
                                    zint0(i,j)=zinter(i,j)
25
        continue
C
C--Now check for time step method
C
C
    CALL PULANG(NRING-1,BEFORE) !angle at last panel center at start
    TIMENOW=TIME(ISTEP)
    CALL indvel
    CALL GETDT(DT)
    CALL GETK(1)
    CALL UPDATE(DT/2.0,1)
    TIPVAX=WVEL(NRING)
C
    CALL PULANG(NRING-1,ANOW) !get new angle and compute the change
    ALPLAST=ALPO+(ANOW-BEFORE) !and add it to the angle of the edge
    TIMENOW=TIME (ISTEP) +DT/2.0
    CALL indvel
    CALL GETK(2)
    CALL UPDATE(DT/2.0,2)
C
    CALL PULANG(NRING-1,ANOW)
    ALPLAST=ALPO+(ANOW-BEFORE)
    TIMENOW=TIME (ISTEP) +DT/2.0
    CALL indvel
    CALL GETK(3)
    CALL UPDATE(DT,3)
C
    CALL PULANG(NRING-1,ANOW)
    ALPLAST=ALPO+(ANOW-BEFORE)
    TIMENOW=TIME(ISTEP) +DT
    CALL indvel
    CALL GETK(4)
```

```
C
C--Update positions for Runge-Kutta
C
                    DO 10 I=0,NRING
                    RADIUS(I)=RADO(I)+(DT/6.0)*(RK(I,1)+2.0*RK(I,2)+
        $
                        2.0*RK(I,3)+RK(I,4))
            HEIGHT(I)=HEIO(I)+(DT/6.0)*(HK(I,1) +2.0*EK(I,2)+
        $
                        2.0*HK(I,3)+HK(I,4))
        10 CONTINUE
            DO 101 j = 1,level
            do 101 i = 1,ninter(j)
                    Rinter(I,j)=RintO(I,j)+(DT/6.0)*(RKi(I,j,1)+2.0*RKi(I,j,2)+
            $
                        2.0*RKi(I,j,3)+RKi(I,j,4))
                            zinter(I,j)=zintO(I,j)+(DT/6.0)*(HKi(I,j,1)+2.0*HKi(I,j,2)+
            $
                        2.0*HKi(I,j,3)+HKi(I,j,4))
101 CONTINUE
c
c far wake
c
    rfar(2) = rinter(ninter(4),4)
    zdis = abs(zinter(ninter(4),4)-zinter(ninter(3),3))
        zfar(2) = zinter(ninter(4),4)-zdis
        gfar(2) = ginter(ninter(4),4)/zdis
        hsum1 = 0.
        zsum1 = 0.
        do 60 i = 1,ninter(3)-1
            hsum1 = hsum1 + ginter(i,3)*rinter(i,3)*rinter(i,3)
            zsum1 = zsum1 + ginter(i,3)*zinter(i,3)*rinter(i,3)**2
60 continue
        z1 = zsum1/hsum1
        hsum2 = 0.
        zsum2 = 0.
        gsum2 = 0.
        do 70 i = 1,ninter(4)-1
            hsum2 = hsum2 + ginter(i,4)*rinter(i,4)*rinter(i,4)
```

```
            zsum2 = zsum2 + ginter(i,4)*zinter(i,4)*rinter(i,4)**2
                gsum2 = gsum2 + ginter(i,4)
    70 continue
        z2 = zsum2/hsum2
        rfar(1) = sqrt(hsum2/gsum2)
        zdif = abs(z2-z1)
        zfar(1) = z2 - zdif
        gfar(1) = gsum2/zdif
C
c--compute new core radii for intermediate vorticies
C
            do 27 j = 1, level
                do 27 i = 1,ninter(j)
                        cinter(i,j) =SQRT(VOLinter(i,j)/
            &
                                    (2.0*Rinter(i,j)))/PI
        continue
    27
C
C--Compute the new core radius for the tip ring and update time
C
    CORE(NRING)=SQRT(VOLTIP/(2.0*RADIUS(NRING)))/PI
    TIME(ISTEP+1)=TIME (ISTEP) +DT
C
    RETURN
    END
        subroutine newwake(indexvort, nvort,lrite,lgo)
C
c subroutine to form new 1st level intermediate wake from rolled-up
c near wake
C
    include '/user/huh/rollup/common.f'
    integer indexvort(10)
```

```
    integer ivstart(10)
    character*60 gtitle,title
    real rtemp(10),ztemp(10),gtemp(10)
    real rfirst(3),zfirst(3),gfirst(3),lambda,k
    logical lrite,lgo,lstart
    data omegaw,epsilon/1.0,0.001/
    data rmin/0.2/
CVD$R ASSOC
    title = 'r/R"z/R"time = '
    istart = 1
    do 10 i = 1,nvort
        gtemp(i) = 0.
        hsum = 0.
        zsum = 0.
        lstart = .false.
        do 20 j = istart,indexvort(i)-1
            if(radius(j) .ge. rmin)then
                if(.not. lstart)then
                    ivstart(i) = j
                    lstart = .true.
            endif
            gtemp(i) = gtemp(i) + gamma(j)
            hsum = hsum + gamma(j)*radius(j)*radius(j)
            zsum = zsum + gamma(j)*height(j)*radius(j)*radius(j)
                endif
20 continue
        if( abs(gtemp(i)) .le. 1.e-05
    & .or. abs(hsum) .le.1.e-05)then
                ztemp(i) = 1.
            goto10
        endif
        rtemp(i) = sqrt(hsum/gtemp(i))
        ztemp(i) = zsum/hsum
        istart = indexvort(i)
10 continue
```

c
c discard vorticies if they are above the plane of the rotor
c
$i=1$ continue
if(ztemp(i) .ge. 0. .or. rtemp(i).le. rmin)then do $17 \mathrm{j}=\mathrm{i}$, nvort-1
$\operatorname{rtemp}(\mathrm{j})=\operatorname{rtemp}(\mathrm{j}+1)$
$z \operatorname{temp}(j)=z \operatorname{temp}(j+1)$
gtemp $(\mathrm{j})=$ gtemp $(\mathrm{j}+1)$
ivstart( j ) $=$ ivstart ( $\mathrm{j}+1$ )
17
continue nvort $=$ nvort -1
else $i=i+1$
endif
if(i .ne. nvort)then goto 15
endif
c
c repack arrays
c

$$
\text { do } 40 \mathrm{j}=4,2,-1
$$

do $40 i=1$, ninter $(j-1)$
rinter $(i, j)=\operatorname{roldi}(i, j)+\operatorname{omegaw}(\operatorname{rinter}(i, j-1)-r o l d i(i, j))$
zinter(i,j) $=$ zoldi(i,j) + omegaw*(zinter(i,j-1)-zoldi(i,j))
ninter ( j ) $=$ ninter $(\mathrm{j}-1)$
40 continue
c
c determine 1st level of intermediate wake from rolled-up near wake
c
call toteng(tke)
groot $=$ gtemp (1)
hroot $=$ gtemp(1)*rtemp(1)**2

```
    zroot = gtemp(1)*ztemp(1)*rtemp(1)**2
    rfirst(1) = sqrt(hroot/groot)
    zfirst(1) = zroot/hroot
    gfirst(1) = groot
    call rke(eroot,ivstart(1),ivstart(2)-1)
    sum3 = 0.
        DO I=1,ivstart(1)-1
            R1=SQRT((HEIGHT(I)-zfirst(1))**2+
        (RADIUS(I)-rfirst(1))**2)
            R2=SQRT((HEIGHT(I)-zfirst(1))**2+
        (RADIUS(I)+rfirst(1))**2)
            LAMBDA=(R2-R1)/(R2+R1)
        CALL ELLIP(LAMBDA,K,E)
        sum3=sum3+GAMMA(I)*gfirst(1)*(R1+R2)*(K-E)/2.0
    enddo
    DO I=ivstart(2),nring
        R1=SQRT((HEIGHT(I)-zfirst(1))**2+
            (RADIUS(I)-rfirst(1))**2)
        R2=SQRT((HEIGHT(I)-zfirst(1))**2+
            (RADIUS(I)+rfirst(1))**2)
        LAMBDA= (R2-R1)/(R2+R1)
        CALL ELLIP(LAMBDA,K,E)
        sum3=sum3+GAMMA(I)*gfirst(1)*(R1+R2)*(K-E)/2.0
enddo
senergy = tke - eroot - 2.*sum3
czzz = exp(2.*senergy/(gfirst(1)**2*rfirst(1))+1.75)
croot = 8.*rfirst(1)/czzz
volinter(1,1) = 2.*pi*pi*rfirst(1)*croot**2
gcent = 0.
hcent = 0.
zcent = 0.
do i = 2,nvort-1
    gcent = gcent + gtemp(i)
    hcent = hcent + gtemp(i)*rtemp(i)**2
    zcent = zcent + gtemp(i)*ztemp(i)*rtemp(i)**2
```

```
    enddo
    rfirst(2) = sqrt(hcent/gcent)
    zfirst(2) = zcent/hcent
    gfirst(2) = gcent
    call rke(ecent,ivstart(2),nsmall)
    sum3 = 0.
        DO I=1,ivstart(2)-1
                            R1=SQRT((HEIGHT(I)-zfirst(2))**2+
        (RADIUS(I)-rfirst(2))**2)
            R2=SQRT((HEIGHT(I)-zfirst(2))**2+
        (RADIUS(I)+rfirst(2))**2)
            LAMBDA=(R2-R1)/(R2+R1)
        CALL ELLIP(LAMBDA,K,E)
        sum3=sum3+GAMMA(I)*gfirst(2)*(R1+R2)*(K-E)/2.0
    enddo
    R1=SQRT((HEIGHT(nring)-zfirst(2))**2+
        (RADIUS(nring)-rfirst(2))**2)
    R2=SQRT((HEIGHT(nring)-zfirst(2))**2+
        (RADIUS(nring)+rfirst(2))**2)
    LAMBDA=(R2-R1)/(R2+R1)
    CALL ELLIIP(LAMBDA,K,E)
    sum3=sum3+GAMMA(nring)*gfirst(2)*(R1+R2)*(K-E)/2.0
    senergy = tke - ecent - 2.*sum3
    czzz = exp(2.*senergy/(gfirst(2)**2*rfirst(2))+1.75)
    ccent = 8.*rfirst(2)/czzz
    volinter(2,1) = 2.*pi*pi*rfirst(2)*ccent**2
    rfirst(3) = rtemp(nvort)
    zfirst(3) = ztemp(nvort)
    gfirst(3) = gtemp(nvort)
    volinter(3,1) = voltip
    do i = 1,3
    rinter(i,1) = roldi(i,1) + omegaw*(rfirst(i) - roldi(i,1))
    zinter(i,1) = zoldi(i,1) + omegaw*(zfirst(i) - zoldi(i,1))
    ginter(i,1) = gfirst(i)
    enddo
```

```
    ninter(1) = 3
c
c
set vortex strengths equal to each other
c
    do j=2,level
            do i = 1,ninter(j)
                ginter(i,j) = ginter(i,1)
                volinter(i,j) = volinter(i,1)
            enddo
        enddo
c
c set new core sizes
c
    do j = 1,level
            do i = 1,ninter(j)
                    cinter(i,j) = sqrt(volinter(i,j)/(2.*rinter(i,j)))/pi
            enddo
        enddo
c
c check for convergence of tip vortex
C
    lgo = .false.
    do j = 1,level
        disnew = rinter (3,j)**2 + zinter(3,j)**2
        disold = roldi(3,j)**2 + zoldi(3,j)**2
        test = abs((disnew-disold)/disold)
        if(test .gt. epsilon )then
            lgo = .true.
            goto 33
            endif
        enddo
        Irite = .true.
        type*,'wake convergence at time =',timenow
        title = 'r/R`z/R~CONVERGED WAKE time = '
33
        continue
```

C

C
C

C

$$
\text { do } j=1, \text { level }
$$

$$
\text { do } \mathbf{i}=1, \text { ninter }(j)
$$

$$
\operatorname{roldi}(i, j)=\operatorname{rinter}(i, j)
$$

$$
\operatorname{zoldi}(i, j)=\text { zinter }(i, j)
$$

$$
\operatorname{goldi}(i, j)=\operatorname{ginter}(i, j)
$$

enddo

> enddo
if(lrite)then
call prntout
gtitle $=$ 'New geometry'
write(19,*) gtitle
write(19,*) level
write (19,*) 52
write $(19,12)$ title,timenow
format (a30,f9.6)
write (19,*) 1
write(19,*) 0.,1.5,-1.18,. 02
write(17,*)'New geometry time $=$ ',timenow
write (17,*)'level = ', level
do $30 \mathrm{j}=1$, level
write (19,*) 6
write (17,*)' ,
write(17,*)'for level\#', j
write (17,*)'\# of vorticies $=$ ', ninter (j)
write (19,*) ninter (j)
do $30 \mathrm{i}=1$, ninter $(\mathrm{j})$
write(19,*) rinter(i, $),$ zinter (i, $j$ )
write(17,*)i,rinter(i,j),zinter(i,j),ginter(i,j)
continue
write (17,*)' '
write(17,*)' far wake'
write (17,*)'nfar =', nfar
do $99 \mathrm{i}=1$, nfar
write(17,*)i,rfar(i),zfar(i),gfar(i)
continue
lrite = .false.
endif
return
end
subroutine oldsheet(istep)
c
c a subroutine to read in old near wake values if the geometry has c not converged
C
include '/user/huh/rollup/common.f'
rewind(7)
read(7) (radius ( $i$ ) , $i=0$, nring)
read (7) (height ( $i$ ), $i=0$, nring)
read (7) (core (i), $i=0, n r i n g)$
read(7) (gamma(i), $i=0, n r i n g)$
$\operatorname{read}(7)(\operatorname{axlseg}(i), i=0, n r i n g)$
gmt0=gamma(nring)
time(istep) $=$ time (istep) +time (1)
voltip $=2 . * p i * p i * r a d i u s(n r i n g) * c o r e(n r i n g) * * 2$
return
end
subroutine prntout
c
c a subroutine to printout the final results
c

```
include '/user/huh/rollup/common.f'
    common /oldgeom/radiusold(0:900), heightold(0:900)
    character*60 gtitle,Ititle,title
    real expwx(4), expwy(4)
    data expwx/.77,.783334,.81,.86/
```

        data expwy/-.5966666,-.4133333,-.243333,-7.6666594e-02/
    C
c write out wake geometry
c
12
gtitle $=$ 'WAKE POSITIONS'
ITITLE $=$ 'RADIUS ${ }^{\circ}$ HEIGHT~TIME $=$ '
write(12,*) gtitle
write(12,*) level+5
write(12,*) 62
write (12,12) ITITLE,timenow
format(a21,f9.5)
WRITE(12.*) 1
write(12,*) 0.,1.5,-1.18,. 02
write(12,*) 2
write (12,*) 2
write(12,*) 0.1,0.0
write(12,*) 1.0,0.0
write (12,*) 6
write(12,*)nring
CVD $\$$ NOVECTOR
CVD $\$$ NOCONCUR
do $151 \mathrm{i}=1$, nring
write(12,*)radius(i), height(i)
151
continue
CVD $\$$ NOVECTOR
CVD\$ NOCONCUR
do $163 j=1$, level
write(12,*)6
write(12,*)ninter(j)
CVD\$ NOVECTOR
CVD $\$$ NOCONCUR
do $183 i=1$, ninter $(j)$
write(12,*) rinter(i,j),zinter(i,j)
continue
183
continue

```
        data expwy/-.5966666,-.4133333,-.243333,-7.6666594e-02/
```

c
c
c
gtitle $=$ 'WAKE POSITIONS'
1TITLE = 'RADIUS~HEIGHT ${ }^{\text {TIME }}=$ '
write(12,*) gtitle
write(12,*) level+5
write(12,*) 52
write(12,12) 1TITLE, timenow
12
format(a21,f9.5)
WRITE(12,*) 1
write(12,*) 0.,1.5,-1.18,. 02
write(12,*) 2
write(12,*) 2
write(12,*) 0.1,0.0
write(12,*) 1.0,0.0
write(12,*) 6
write(12,*)nring
CVD\$ NOVECTOR
CVD\$ NOCONCUR
do 151 i $=1$, nring write(12,*)radius(i), height(i)
151 continue
CVD\$ NOVECTOR
CVD\$ NOCONCUR

```
do 163 j = 1,level
        write(12,*)6
        write(12,*)ninter(j)
```

CVD\$ NOVECTOR
CVD $\$$ NOCONCUR
do $183 \mathrm{i}=1$, ninter $(\mathrm{j})$
write(12,*) rinter(i,j),zinter(i,j)
183
163
continue continue

> write $(12, *) 4$
> write $(12, *) 4$

CVD\$ NOVECTOR
CVD\$ NOCONCUR

```
do i = 1,4
    write(12,*)expwx(i), expwy(i)
enddo
write(12,*)2
write(12,*)2
write(12,*)rfar(1),zfar(1)
write(12,*)rfar(1),-1.18
write(12,*)2
write(12,*)2
write(12,*)rfar(2),zfar(2)
write(12,*)rfar(2),-1.18
```

c
c write out velocity vectors
c

13
gtitle='Velocity vectors'
title='r/R~z/R~time $=$ '
write(26,*) gtitle
write(26,*) 52
write $(26,13)$ title,timenow
format(a15,f9.5)
write (26,*) 1
write(26,*) 0.,1.,-.6,. 15
write (26,*) nsmall
write (26,*) 2
CVD\$ NOVECTOR
CVD\$ NOCONCUR

```
            do 27 i = 1,nsmall
        write(26,*) radius(i),height(i),uvel(i),wvel(i)
        continue
```

    27
    CVD\$ NOVECTOR
CVD $\$$ NOCONCUR

```
    do 271 i = 1,nsmall
        write(26,*) radius(i),height(i),0.,0.
        continue
    return
    end
```

```
SUBROUTINE PULANG(NIN,THNIN)
    include '/user/huh/rollup/common.f'
    real r(0:900), z(0:900)
    equivalence(r,radius)
    equivalence(z,height)
COMMON /TEMP/ BETA(0:900)
REAL*4 MAG
```

C
$\mathrm{RT}=\mathrm{R}$ (NRING)
$\mathrm{ZT}=\mathrm{Z}$ (NRING)
C
$A=-1.0$
$\mathrm{B}=0.0$
$\mathrm{C}=\mathrm{R}(0)-\mathrm{RT}$
$D=Z(0)-Z T$
$D O T=A * C+B * D$
CROSS $=A * D-B * C$
$\operatorname{MAG}=\operatorname{SQRT}((A * A+B * B) *(C * C+D * D))$
ARG=DOT/MAG
IF (ABS (ARG) . GT. 1.0) $A R G=\operatorname{SIGN}(1.0, A R G)$
$\operatorname{BETA}(0)=\operatorname{SIGN}(1.0, \operatorname{CROSS}) * \operatorname{ACOS}(\operatorname{ARG})$
btemp = beta(0)
C
DO $100 \mathrm{I}=1$, NIN
$A=R(I-1)-R T$
$B=Z(I-1)-Z T$

```
    C=R(I)-RT
    D=Z(I)-ZT
    DOT=A*C+B*D
    CROSS=A*D-B*C
    MAG=SQRT((A*A+B*B)*(C*C+D*D))
    ARG=DOT/MAG
    IF (ABS(ARG).GT. 1.0) ARG=SIGN(1.0,ARG)
    BETA(I)=btemp+SIGN(1.0,CROSS)*ACOS(ARG)
        btemp = beta(i)
    100
        continue
C
    THNIN=BETA(NIN)
C
RETURN
END
```

subroutine readinput(tint,itrs)
c
c a subroutine to read in initial values for restart C
include '/user/huh/rollup/common.f'
CVD\$R ASSOC
c
read(16) nring, lineup,minsect,maxsegs
read(16) time(2),itrs
read(16) tint
nsmall $=$ nring-1
read(16) level,nfar
read(16) (ninter(i), $i=1,1 e v e l)$
do $\mathbf{i}=1,3$
do $j=1,4$
read(16)rinter(i,j)
read(16)zinter ( $i, j$ )
read(16)ginter ( $i, j$ )
read(16)cinter(i,j)
enddo
enddo
do $i=1,3$
volinter(i,1) $=2 . * \operatorname{pi*pi} *$ rinter(i,1)*cinter(i,1)**2 do $j=2,4$
volinter(i,j) $=\operatorname{volinter}(i, 1)$ enddo
enddo
read(16)(rfar(i), $i=1, n f a r)$
read(16)(zfar(i),i=1,nfar)
read(16)(gfar(i), $i=1, n f a r)$
do $i=1,3$

$$
\text { do } j=1,4
$$

$\operatorname{roldi}(i, j)=\operatorname{rinter}(i, j)$
zoldi(i,j) $=$ zinter (i,j)

```
            goldi(i,j) = ginter(i,j)
            enddo
enddo
read(16)aslope,sigma
read(16)nspan
read(16)(span(i),i=1,nspan)
read(16)(pitch(i),i=1,nspan)
read(16)(chord(i), i=1,nspan)
read(16)(circold(i), i=1,nspan)
read(16)(radius(i), i=0,nring)
read(16)(height(i),i=0,nring)
read(16)(gamma(i), i=0,nring)
read(16)(core(i), i=0,nring)
read(16)(axlseg(i),i=0,nsmall)
voltip = 2.*pi*pi*radius(nring)*core(nring)**2
gmt0 = gamma(nring)
write(7)(radius(i), i=0,nring)
write(7)(height(i), i=0,nring)
write(7)(core(i),i=0,nring)
write(7)(gamma(i),i=0,nring)
write(7)(axlseg(i),i=0,nring)
return
end
```

SUBROUTINE REDISC(R,Z,DGAM,NIN,NOUT, GMAX, SUML, DELTAL)

C--This does curve fitting using blended parabolas on $R_{2} Z$, DGAM C

C
include '/user/huh/rollup/common.f'
COMMON/STAR/ICIRC(10), FCIRC(10) , SSTAR (10) , GSTAR(10) , NSTAR
DIMENSION R(0:900), Z $(0: 900)$, DGAM $(0: 900), S(0: 900), Q R(0: 900)$,
\$
QZ $(0: 900), A 0(0: 900), A 1(0: 900), A 2(0: 900), A 3(0: 900)$,

```
                BO(0:900),B1(0:900), B2(0:900), B3(0:900),C0(0:900),
                C1(0:900),C2(0:900),C3(0:900), GMID (0:900) ,SMID (0:900),
                Q(0:900),RS(0:8100),ZS(0:8100),AL(0:8100).
                    $ THSUB(0:8100),ARCG(0:900),GPL(0:900)
                        DIMENSION ROLD(0:900),ZOLD(0:900),DGAMOLD(0:900)
                        LOGICAL LINEAR ILINEAR = .TRUE. for linear circ. fit
                    DATA LINEAR/.FALSE./ ILINEAR = .FALSE. for cubic circ. fit
                    DATA NSUB/9/
C
CVD$R ASSOC
CVD$R CNCALL
    ILOOPER=0
C
    DO 1 I=0,NRING
    ROLD(I)=R(I)
        ZOLD(I)=Z(I)
        1 DGAMOLD(I)=DGAM(I)
C
C--Get straight line distances between points, S(I) is the total
C distance along the straight lines to node I
C
        S(0)=0.0
        DO 2O I=1,NIN
    20 S(I)=S(I-1)+SQRT((Z(I)-Z(I-1))**2+(R(I)-R(I-1))**2)
C
C--Get value of s for the trailing extended segment, this value
C of S is out at the presumed end of the sheet.
C
    S(NIN+1)=S(NIN)+0.5*(S(NIN)-S(NIN-1))
C
C--Get slopes at each interior node, QR(I)=D(R)/DS at node I
C
    DO 30 I=1,NIN-1
    S1=(S(I+1)-S(I))/(S(I-1)-S(I))
    S2=1.0/S1
```

```
        S3=1.0/(S(I+1)-S(I-1))
        QR(I)=((R(I-1)-R(I))*S1-(R(I+1)-R(I))*S2)*S3
    30 QZ(I)=((Z(I-1)-Z(I))*S1-(Z(I+1)-Z(I))*S2)*S3
C
C--Coef's for interior segments
C
            DO 40 I=1,NIN-2
            S1=1.0/(S(I+1)-S(I))
            AO(I) =R(I)
            A1(I)=QR(I)
            A2(I)=S1*(-QR(I+1)-2*QR(I)+3*S1*(R(I+1)-R(I)))
            A3(I)=S1*S1*(QR(I+1)+QR(I)-2*S1*(R(I+1)-R(I)))
            BO(I)=Z(I)
            B1(I)=QZ(I)
            B2(I)=S1*(-QZ(I+1)-2*QZ(I)+3*S1*(Z(I+1)-Z(I)))
        40 B3(I)=S1*S1*(QZ(I+1)+QZ(I)-2*S1*(Z(I+1)-Z(I)))
C
C--Coef's for first segment
C
    S1=S(1)-S(0)
    S2=S(2)-S(0)
    S3=1.0/(S(2)-S(1))
        AO(0)=R(0)
        A1(0)=((R(1)-R(0))*S2/S1-(R(2)-R(0))*S1/S2)*S3
        A2(0)=((R(2)-R(0))/S2-(R(1)-R(0))/S1)*S3
        A3(0)=0.0
        BO(0)=Z(0)
        B1(0)=((Z(1)-Z(0))*S2/S1-(Z(2)-Z(0))*S1/S2)*S3
        B2(0)=((Z(2)-Z(0))/S2-(Z(1)-Z(0))/S1)*S3
        B3(0)=0.0
C
C--Coef's for last segment
C
\(\mathrm{N}=\mathrm{NIN}\)
S1 \(=\mathbf{S}(\mathrm{N})-\mathrm{S}(\mathrm{N}-1)\)
```

```
S2=S(N-2)-S(N-1)
S3=1.0/(S(N-2)-S(N))
AO(N-1)=R(N-1)
A1(N-1)=((R(N)-R(N-1))*S2/S1-(R(N-2)-R(N-1))*S1/S2)*S3
A2(N-1)=((R(N-2)-R(N-1))/S2-(R(N)-R(N-1))/S1)*S3
A3(N-1) =0.0
BO(N-1)=Z(N-1)
B1(N-1)=((Z(N)-Z(N-1))*S2/S1-(Z(N-2)-Z(N-1))*S1/S2)*S3
B2(N-1)=((Z(N-2)-Z(N-1))/S2-(Z(N)-Z(N-1))/S1)*S3
B3(N-1)=0.0
C
C--Get the angle of the ring with index NSMALL wrt the tip core
C
CALL PULANG(NIN,THNIN)
C
C--Get the "A" coefficient for the Kadin spiral
C
    RNU=XNU(THNIN)
    IF (ABS(SIN(THNIN)) .GE. O.5) THEN
            A=(Z(NRING)-Z(NIN))/(THNIN**(-RNU)*SIM(THIIN))
                                ELSE
            A=(R(NRING)-R(NIN))/(THNIN**(-RNU)*COS(THFIN))
                    ENDIF
C
C--Get array of R and Z at the subinterval locations, and compute
C the total length of the curve
C
C--First do the blended parabola segments
C
    RS(0)=AO(0)
    ZS(0)=B0(0)
    AL(0)=0.0
    DO 50 I=0,NIN-1
    DO 50 KSUB=1,NSUB
    J=I*NSUB+KSUB
```

```
        T=(S(I+1)-S(I))*FLOAT(KSUB)/FLOAT(NSUB)
        RS(J)=AO(I) +T*(A1 (I) +T*(A2(I) +T*A3(I)))
        ZS(J)=BO(I)+T*(B1(I)+T*(B2(I)+T*B3(I)))
    60 AL(J)=AL(J-1)+SQRT((RS(J)-RS(J-1))**2+(ZS (J)-ZS(J-1))**2)
C
C--The next few sections find the value of theta at the end of the
C trailing extended segment, i.e., gets the theta corrosponding
C to arclength S(NIN+1)
C
C--First, before doing the numerical integration, need to get a
C value to use for delta theta
C
        DS=S(NIN+1)-S(NIN)
        RNU=XNU(THNIN)
        DTG=DS/(SQRT(RNU**2+THNIN**2)*A*THNIN**(-RNU-1.0))
        DTG=DTG/2.0
C
C--Now have delta theta (i.e., DTG) which should yield several
C steps in the numerical integration from S(NIN) to S(NIN+1).
C
C--Now do the numerical integration and search for theta(NIN+1)
C
    ITERS=0
    SEND=S(NIN+1)-S(NIN)
    STH=0.0
    TH=THNIN
    64 ITERS=ITERS+1
            IF (ITERS .GT. 10000) STOP 'ITERS .GT. 10000 LOOP 63'
            RNU=XNU(TH)
            R1=R(NRING) - A*COS (TH)*TH**(-RNU)
            Z1=Z(NRING) - A*SIN (TH)*TH**(-RNU)
            RNU=XNU(TH+DTG)
            R2=R(NRING)-A*COS(TH+DTG)*(TH+DTG)**(-RNU)
            Z2=Z(NRING)-A*SIN(TH+DTG)*(TH+DTG)**(-RNU)
            DS=SQRT((R2-R1)**2+(Z2-Z1)**2)
```

```
            IF (STH+DS .GE. SEND) GO TO 63
            STH=STH+DS
            TH=TH+DTG
            GO TO }6
    63 FRAC=(SEND-STH)/DS
            THNIN1=TH+FRAC*DTG
C
C--Now can get the sub-interval locations over the extended spiral
C segment.
C
            DTH=(THNIN1 - THNIN)/NSUB
            DO 65 KSUB=1,NSUB
            THETA=THNIN +KSUB *DTH
            J=NIN*NSUB+KSUB
            RNU=XNU(THETA)
            RS(J)=R(NRING) - A*THETA**(-RNU)*COS (THETA)
            ZS(J)=Z(NRING) -A*THETA**(-RNU)*SIN(THETA)
            AL(J)=AL(J-1)+SQRT((RS(J)-RS(J-1))**2+(ZS (J)-ZS(J-1))**2)
            6 5 ~ C O N T I N U E ~
C
C--Total length of curve is SUML
C
    SUML=AL ((NIN+1)*NSUB)
C
    6 8 9 \text { CONTINUE}
C
C--Check for lineup of spiral sheet segments
C
            IF (LINEUP .EQ. 1) THEN
                GO TO 2010
                                lline up the segments
            ELSE
                GO TO 2020 !for equal sized segments
            ENDIF
C
C--This section for equal sized segments
```

```
C
C--Get new R and Z points at distance DELTAL apart, these are the
C centroids
C
    2020 DELTAL=SUML/NOUT
            XLSEG=DELTAL/2.0
            AXLSEG(0)=0.0
            I=1
            DO 60 J=1,NSUB*(NIN+1)
        80 IF (XLSEG .LT. AL(J-1) .OR. XLSEG .GT. AL(J)) GO TO 60
            FRAC=(XLSEG-AL(J-1))/(AL(J)-AL(J-1))
            R(I)=abs(FRAC*(RS(J)-RS(J-1))+RS(J-1))
            Z(I)=FRAC*(ZS(J)-ZS(J-1))+ZS(J-1)
            AXLSEG(I)=XLSEG+DELTAL/2.0 !running total of arc length to
                                    ! segment edge.
            WPAN(I)=DELTAL/2.0 !panel half-width
            I=I+1
            XLSEG=XLSEG+DELTAL
            IF (I .GT. NOUT) GO TO 70
            GO TO 80
            6 0 \text { CONTINUE}
            70 CONTINUE
            GO TO 2030
C
C--This part is for lined up segments in the spiral region
C
    2010 NOUT1=NOUT
    NTOTSUB=NSUB*(NIN+1)
    CALL ANGLES(RS,ZS,NTOTSUB,THSUB)
    ALPLAST=THSUB(NTOTSUB) !save angle at edge of sheet for DOUBLET
    IF (THSUB(NTOTSUB) .LE. 9.0*PI/2.0) THEN !ALFSPIR is the angle
        ALFSPIR=THSUB(NTOTSUB)-3.0*PI/2.0 lat which the "spiral"
    ELSE
        ALFSPIR=3.0*PI
    ENDIF
        part of the sheet begins
    !(i.e., the equal angular
    !segments)
```

IF (THSUB(NTOTSUB) .LT. 7.0*PI/2.0) GO TO 2020 !check for at
!least 1.5 loop

STOTAL=SUML
$I=1$
3513 IF (THSUB(I) . GT. ALFSPIR) GO TO 3512
$I=I+1$
GO TO 3513
3512 SFLAT=AL(I)
SSPIRAL=STOTAL-SFLAT
R3PIQ2=SQRT ((RS (I) -ROLD (NRING)) **2+(ZS(I)-ZOLD(NRING))**2)
RAVERAGE $=0.0$
DO $3511 \mathrm{~J}=\mathrm{I}, \mathrm{NTOTSUB}$
3511 RAVERAGE=RAVERAGE
\$ +SQRT((RS(J)-ROLD(NRING))**2+(ZS(J)-ZOLD (NRING))**2)
RAVERAGE=RAVERAGE/(NTOTSUB-I+1)
7126 DENOM=1.0+(SSPIRAL*R3PIQ2)/(SFLAT*RAVERAGE)
NSEG=NOUT1 !\# of sheet segments to be output
NSEGFLAT=INT(NSEG/DENOM+0.5) lestimate the \# of segments on
NSEGSPIR=NSEG-NSEGFLAT ! the flat and spiral parts
RNLOOP $=($ THSUB (NTOTSUB) -ALFSPIR)/(2.0*PI) !\# of loops in spiral
NTHSECT=INT ( (FLOAT (NSEGSPIR)/RNLOOP) +0.5) !\# of theta sections
IF (NOUT1 .LT. MAXSEGS) THEN
IF (NTHSECT .LT. MINSECT) THEN
ILOOPER=ILOOPER +1
IF (ILOOPER .GT. MAXSEGS) STOP ' ILOOPER .GT. MAXSEGS'
NOUT1 $=$ NOUT1 +1
GO TO 7126
ENDIF
ENDIF
c
if(nthsect .eq. O)return

DTHETA $=2.0 * P I / \mathrm{NTHSECT}$
$I=1$
ISEG=NSEG+1-I
THEDGE=THSUB (NTOTSUB) -DTHETA !edge of current segment (small theta)
IF (THEDGE .LT. 7.0*PI/2.0) GO TO 2020 !check for at least 1.5 loop

```
    THMID=THEDGE+DTHETA/2.0 !theta for centroid of current segment
    AXLSEG(ISEG)=SUML
    DO 3060 J=NTOTSUB,1,-1
    3080 IF (THMID .GT. THSUB(J) .OR. THMID .LT. THSUB(J-1)) GO TO 3050
        FRAC=(THMID-THSUB(J-1))/(THSUB(J)-THSUB(J-1))
    R(ISEG)=abs(FRAC*(RS(J)-RS(J-1))+RS(J-1))
    Z(ISEG)=FRAC*(ZS(J)-ZS(J-1))+ZS(J-1)
    THMID=THMID-DTHETA
3050 IF (THEDGE .GT. THSUB(J) .OR. THEDGE .LT. THSUB(J-1)) GO TO 3060
    FRAC=(THEDGE-THSUB(J-1))/(THSUB(J)-THSUB(J-1))
    AXLSEG(ISEG-1)=FRAC*(AL(J)-AL(J-1))+AL(J-1)
    THEDGE=THEDGE-DTHETA
    IF (THEDGE .LT. ALFSPIR) GO TO 3070 !done with spiral part
    I=I+1
    ISEG=NSEG+1-I
    GO TO 3080
3060 CONTINUE
3070 CONTINUE II contains number of segments in the spiral region
    NSEGSPIR=I !exact value
    NSEGFLAT=NSEG-NSEGSPIR
C
C now find segments on flat part of sheet
C
C--get segment stretching parameters
C
    SUMI=0.0
    DO 3761 I=1,NSEGFLAT-1
    3761 SUMI=SUMI+I
    DSLAST=AXLSEG(ISEG)-AXLSEG(ISEG-1)
    HHH=(AXLSEG(ISEG-1)-NSEGFLAT*DSLAST)
    $
        /(SUMI-NSEGFLAT*(HSEGFLAT-1))
    DSFIX=DSLAST-HHH*(NSEGFLAT-1)
C
    AXLSEG(0)=0.0
    I=1
```

```
    DSI=DSFIX +HHH* (I-1)
    XLSEG=DSI/2.0 larc length to filament
    EDGE=DSI
                                larc length to outboard edge of segment
DO 3061 J=1,NTOTSUB
    3081 IF (XLSEG .LT. AL(J-1) .OR. XLSEG .GT. AL(J)) GO TO 3061
    FRAC=(XLSEG-AL(J-1))/(AL(J)-AL(J-1))
    R(I)=abs(FRAC*(RS(J)-RS(J-1))+RS(J-1))
    Z(I)=FRAC*(ZS(J)-ZS(J-1))+ZS(J-1)
    AXLSEG(I)=EDGE
    I=I+1
    DSI=DSFIX + HHH* (I-1)
    XLSEG=EDGE+DSI/2.0
    EDGE=EDGE+DSI
    IF (I .GT. NSEGFLAT) GO TO 3071
    GO TO 3081
    3 0 6 1 ~ C O N T I N U E ~
    3071 CONTINUE
C--set WPAN for all segments
    DO 3110 ISEG=1,NSEG
    WPAN(ISEG)=(AXLSEG(ISEG)-AXLSEG(ISEG-1))/2.0
    3110 CONTINUE
    R(NOUT1+1)=ROLD(NRING)
    Z(NOUT1+1)=ZOLD(NRING)
    DGAM(NOUT1+1)=DGAMOLD (NRING)
    NOUT=NOUT1
    NRING=NOUT1+1
    NSMALL=NRING-1
    GO TO 2030
C
    2 0 3 0 ~ C O N T I N U E ~
C
C--If panel case, estimate inclinations (only nec. for R-K)
C
IF (IPANEL .EQ. 1) THEN
        DO 920 I=1,NSMALL-1
```

```
                    SL=-SQRT((R(I)-R(I-1))**2+(Z(I)-Z(I-1))**2)
                    SR= SQRT((R(I)-R(I+1))**2+(Z(I)-Z(I+1))**2)
                A1P=((SL*SR)/(SL-SR))
\$
            DRPAN(I)=A1P
            DZPAN(I)=B1P
    920 CONTINUE
            A2P=((R(I+1)-R(I))/SR-(R(I-1)-R(I))/SL)/(SR-SL)
            B2P=((Z(I+1)-Z(I))/SR-(Z(I-1)-Z(I))/SL)/(SR-SL)
            DRPAN(NSMALL)=A1P+A2P*SR
            DZPAN(NSMALL)=B1P+B2P*SR
                ENDIF
C
C--Get values of S at the midpoints of the original ring locations
C
            SMID (0)=0.0
            SMID(NIN)=SUML
            DO 90 I=1,NIN-1
            90 SMID(I)=(AL(NSUB*(I+1))+AL(NSUB*I))/2.0
C
C--Get values of circulation at the midpoints of the original
C ring locations
C
                    GMID(0)=GMAX
                    DO 100 I=1,NIN
    100 GMID(I)=GMID(I-1)+DGAM(I)
C
C--Curve fit for circulation, test for linear or cubic fit
C
    IF (LINEAR) THEN
C
C--The following block is for a linear fit of circulation
C
```

```
                    DO 411 I=0,NIN-1
                CO(I)=GMID(I)
                C1(I)=(GMID(I+1)-GMID(I))/(SMID (I+1)-SMID (I))
                C2(I) =0.0
    411 C3(I)=0.0
C
            ELSE
C
C--The following block is for a blended parabola fit of circulation
C
C--Get slopes at each interior point
C
            DO 110 I=1,NIN-1
                    S1=(SMID(I+1)-SMID(I))/(SMID(I-1)-SMID(I))
    110 Q(I)=((GMID(I-1)-GMID(I))*S1-(GMID (I+1)-GMID(I))/S1)
            $ /(SMID(I+1)-SMID(I-1))
C
C--Coef's for interior points
C
            DO 120 I=1,NIN-2
            S1=1.0/(SMID(I+1)-SMID(I))
            G1=(GMID(I+1)-GMID(I))*S1
            CO(I)=GMID(I)
            C1(I)=Q(I)
            C2(I) =S1*(-Q(I+1)-2*Q(I)+3*G1)
            120 C3(I)=S1*S1*(Q(I+1)+Q(I)-2*G1)
C
C--Coef's for first segment
C
    S1=SMID (1)-SMID (0)
    S2=SMID (2)-SMID (0)
    S3=1.0/(SMID (2)-SMID (1))
    CO(0)=GMID (0)
    C1(0)=S3*((GMID (1)-GMID (0))*S2/S1-(GMID (2)-GMID (0))*S1/S2)
    C2(0)=S3*((GMID(2)-GMID (0))/S2-(GMID (1)-GMID (0))/S1)
```

```
        C3(0)=0.0
    C
C--Coef's for last segment
C
S1=SMID(NIN-2)-SMID(NIN-1)
S2=SMID(NIN)-SMID(NIN-1)
S3=1.0/(SMID(NIN-2)-SMID(NIN))
G1=GMID(NIN-2)-GMID(NIN-1)
G2=GMID(NIN)-GMID(NIN-1)
CO(NIN-1)=GMID(NIN-1)
C1(NIN-1)=S3*(G2*S1/S2-G1*S2/S1)
C2(NIN-1)=S3*(G1/S1-G2/S2)
C3(NIN-1) =0.0
C
    ENDIF
C
    651 CONTINUE
C
C--Interpolate for new delta circulation values.
C XLSEG will now be an array containing the arc length to the
C edge of segment I, call it AXLSEG(I)
C
    I=1
    GPREV=GMAX
    DO 130 J=1,NIN
    150 IF(AXLSEG(I) .LT. SMID(J-1) .OR. AXLSEG(I) .GT. SMID(J))GO TO 130
    T=AXLSEG(I)-SMID(J-1)
    GNEW=CO(J-1)+T*(C1(J-1)+T*(C2(J-1)+T*C3(J-1)))
    DGAM(I)=GNEW-GPREV
    GPREV=GNEW
    I=I+1
    IF (I .GT. NOUT-1) GO TO 140
    GO TO 150
    130 continue
    140 CONTINUE
```

```
    T=SMID(NIN)-SMID(NIN-1)
    GNEW=CO(NIN-1)+T*(C1(NIN-1)+T*(C2(NIN-1)+T*C3(NIN-1)))
    DGAM (NOUT)=GNEW-GPREV
    C
    C--Done
C
            RETURN
            END
            subroutine ritecirc
C
c a subroutine to output final load distribution
C
    include '/user/huh/rollup/common.f'
    real circt(103),r(103)
    character*60 gtitle,title
    open(unit=21,file='initload.dat',status='old',form='formatted')
    do i = 1,103
        read(21,*)r(i),circt(i)
    enddo
    close(21)
    open(unit=8,file='initsheet.dat',form='formatted',status='new')
    gtitle='INIT-SHEET'
    title = 'RADIUS VORT'
    write(8,*) gtitle
    write(8,*) 1
    write(8,*) 52
    write(8,*) title
    write(8,*) 1
    write(8,*) 0.,1.,-5.e-03,1.e-03
    write(8,*) 2
    write(8,*) nring
```

```
    do 65 i = 1,nring
    write(8.*) radius(i),gamma(i)
6 5
    continue
    close(8)
    gtitle='MILLER LOADING'
    title='RADIUS GAMMA'
    write(15,*)gtitle
    write(15,*)2
    write(15,*)52
    write(15,*)title
    write(15,*)1
    write(15,*)0.,1.,0.,3.e-02
        write(15,*)2
    write(15,*)103
        do 133 i = 1,103
        write(15,*)r(i),circt(i)
        continue
    write(15,*)4
    write(15,*)28
    do i = 1,28
        write(15,*)expx(i), expy(i)
    enddo
    return
    end
    subroutine riteinput
c
c a subroutine to write out initial near wake values
C
CVD$R ASSOC
    include '/user/huh/rollup/common.f'
    WRITE(IWR,*)'
```

```
        WRITE(IWR,*)'
        WRITE(IWR,*)' '
        write(iwr,*)'TIME = ',timenow
    WRITE(IWR,*)'
```



```
    WRITE(IWR,*)' '
    WRITE(IWR, 3243)
3243
3244
    FORMAT(1X,I3,3x,e14.6,3x,e14.6,3x,e14.6)
    sum = 0.
    do i = 1,nring
    sum = sum+gamma(i)
    enddo
    write(iwr,*) 'total circulation = ',sum
    return
    end
    subroutine riterst(tint,itrs)
c
c a subroutine to write out values for future restart runs
C
```

```
include '/user/huh/rollup/common.f'
```

include '/user/huh/rollup/common.f'
open(unit=16,file='restart.dat',form='unformatted',status='new')
open(unit=16,file='restart.dat',form='unformatted',status='new')
write(16) nring,lineup,minsect,maxsegs
write(16) nring,lineup,minsect,maxsegs
write(16) timenow,itrs
write(16) timenow,itrs
write(16) tint
write(16) tint
write(16) level,nfar
write(16) level,nfar
write(16) ninter(1:level)
write(16) ninter(1:level)
do i = 1,3
do i = 1,3
do j = 1,4

```
    do j = 1,4
```

```
    write(16) rinter(i,j)
    write(16) zinter(i,j)
    write(16) ginter(i,j)
    write(16) cinter(i,j)
enddo
enddo
write(16)rfar(1:nfar)
write(16)zfar(1:nfar)
write(16)gfar(1:nfar)
write(16)aslope,sigma
write(16)nspan
write(16) span(1:nspan)
write(16)pitch(1:nspan)
write(16)chord(1:nspan)
write(16)circ(1:nspan)
write(16)radius(0:nring)
write(16)height(0:nring)
write(16)gamma(0:nring)
write(16)core(0:nring)
write(16)axlseg(0:nsmall)
close(16)
return
end
```

SUBROUTINE rke(ENERG,istart,iend)
C
C-Computes total energy of the vortex ring system
C
include '/user/huh/rollup/common.f'
CVD\$R ASSOC
REAL*4 LAMBDA,K
C

```
        SUM1=0.0
        SUM2=0.0
        if(istart .eq. 1)goto40
        DO 200 I=1,istart-1
        W=RADIUS(I)
    Z=HEIGHT(I)
    DO 20 J=1,istart-1
        if(i .eq. j)goto20
            WP=RADIUS (J)
            ZP=HEIGHT (J)
            R1=SQRT((Z-ZP)**2+(W-WP)**2)
            R2=SQRT((Z-ZP)**2+(W+WP)**2)
            LAMBDA= (R2-R1)/(R2+R1)
            CALL ELLIP(LAMBDA,K,E)
            SUM1=SUM1+GAMMA (I)*GAMMA (J)* (R1+R2)* (Y-E)/(2.0*PI)
    CONTINUE
    DO 200 J=iend+1,nring
        WP=RADIUS (J)
        ZP=HEIGHT (J)
        R1=SQRT((Z-ZP)**2+(W-WP)**2)
        R2=SQRT((Z-ZP)**2+(W+WP)**2)
        LAMBDA= (R2-R1)/(R2+R1)
        CALL ELLIP(LAMBDA,K,E)
        SUM1=SUM1+GAMMA (I)*GAMMA (J)*(R1+R2)*(I.-E)/(2.0*PI)
        CONTINUE
            continue
        DO 300 I=iend+1,nring
        W=RADIUS(I)
        Z=HEIGHT(I)
    if(istart .eq. 1)goto60
        DO 70 J=1,istart-1
        WP=RADIUS(J)
        ZP=HEIGHT(J)
        R1=SQRT((Z-ZP)**2+(W-WP)**2)
        R2=SQRT((Z-ZP)**2+(W+WP)**2)
```

```
                                    LAMBDA= (R2-R1)/(R2+R1)
                                    CALL ELLIP(LAMBDA,K,E)
                                    SUM1 =SUM1 +GAMMA (I)*GAMMA (J)*(R1+R2)*(K-E)/(2.0*PI)
    CONTINUE
    continue
    DO 80 J=iend+1,nring
        if(i .eq. j)goto 80
            WP=RADIUS(J)
            ZP=HEIGHT (J)
            R1=SQRT((Z-ZP)**2+(W-WP)**2)
            R2=SQRT((Z-ZP)**2+(W+WP)**2)
            LAMBDA= (R2-R1)/(R2+R1)
            CALL ELLIP(LAMBDA,K,E)
            SUM1 =SUM1 +GAMMA (I)*GAMMA (J)*(R1+R2)*(K-E)/(2.0*PI)
        CONTINUE
            continue
SUM1=SUM1*PI
C
    700
    800
DO 900 I=iend+1,nring
            W=RADIUS (I)
    900
    SUM2=SUM2+GAMMA(I)**2*W*(ALOG(8.0*W/CORE (I))-1.75)/(2.0*PI)
SUM2=SUM2*PI
C
ENERG=SUM2+SUM1
C
RETURN
END
```

```
    SUBROUTINE SETA
    include '/user/huh/rollup/common.f'
    COMMON /E2/ AO,A1,A2,A3,RHOSTAR,A
    COMMON /E1/ TAUO,BDUM,CDUM,DDUM,TAU2LOOP
C
    ALPLAST=PI/2.0
    TAUO=0.045
    BIGRREF=0.0968063
    TAU2LOOP=1.4542
    A=2.0*SQRT (2.0)/PI
C
    RETURN
    END
    SUBROUTINE TOTENG(ENERG)
C
C--Computes total energy of the vortex ring system
C
    include '/user/huh/rollup/common.f'
CVD$R ASSOC
    REAL*4 LAMBDA,K
C
    SUM1=0.0
    SUM2=0.0
    DO 20 I=1,NRING
        W=RADIUS(I)
        Z=HEIGHT (I)
        DO 20 J=1,NRING
            IF (I .NE. J) THEN
                WP=RADIUS(J)
                ZP=HEIGHT (J)
```

```
                    R1=SQRT((Z-ZP)**2+(W-WP)**2)
                    R2=SQRT((Z-ZP)**2+(W+WP)**2)
                    LAMBDA=(R2-R1)/(R2+R1)
                CALL ELLIIP(LAMBDA,K,E)
                    SUM1=SUM1+GAMMA (I)*GAMMA (J)* (R1+R2)*(K-E)/(2.0*PI)
                    ENDIF
    2O CONTINUE
        SUM1=SUM1*PI
C
        DO 30 I=1,NRING
            W=RADIUS(I)
        SUM2=SUM2+GAMMA (I)**2*W*(ALOG(8.0*W/CORE(I))-1.75)/(2.0*PI)
            continue
        SUM2=SUM2*PI
C
        ENERG=SUM2+SUM1
C
        RETURN
        END
        SUBROUTINE UPDATE(H,N)
            include '/user/huh/rollup/common.f'
C
C--Computes new positions of the rings based on the time step H}\mathrm{ and
C the velocity contained in RK and HK.
C
            DO 10 I=O.NRING
            RADIUS (I)=RADO(I) +H*RK(I,N)
            HEIGHT (I) = HEIO (I) +H*HK (I N N
        10 CONTINUE
            do 15 j = 1,level
            do 15 i = 1,ninter(j)
```

```
                    rinter(i,j) = rint0(i,j) + h*rki(i,j,n)
                    zinter(i,j) = zintO(i,j) + h*hki(i,j,n)
    1 5
        continue
C
C--Compute the new core radius for the tip ring
C
    CORE(NRING)=SQRT(VOLTIP/(2.0*RADIUS(NRING)))/PI
C
C--Compute the new core radius for intermediate wake
C
    do 17 j = 1,level
                do 17 i = 1,ninter(j)
                Cinter(i,j)=SQRT(VOLinter(i,j)/(2.0*Rinter(i,j)))/PI
    17
        continue
C
    RETURN
    END
    FUNCTION XNU(THETA)
    DATA PI/3.141592654/
    XNU=2.0/3.0
    IF (THETA .LT. 2.0*PI) THEN
        XNU=(1.0/(2.0*PI-1.0))-(1.0/(THETA-1.0))+2.0/3.0
                                    ENDIF
    RETURN
    END
```

