

Waves and Instabilities in the Presence of Motion

I. Galilean Transformation

Waves and Instabilities in the Presence of Material Motion

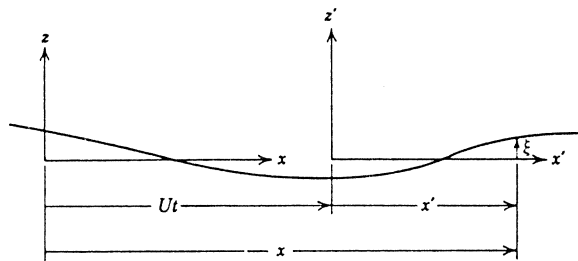


Fig. 10.2.1 The fixed coordinate system x, y, z or the moving coordinate system x', y', z' can be used to define a given position along the equilibrium axis of the string. Because the moving coordinate system has the same velocity U in the x -direction as the string, the string appears to move only in the z' -direction when viewed from the primed frame.

$$\begin{aligned} X &= x' + Ut' \\ z &= z' \\ t &= t' \end{aligned} \Rightarrow \frac{\partial \tilde{z}(x,t)}{\partial t'} = \frac{\partial \tilde{z}}{\partial t} \frac{\partial t}{\partial t'} + \frac{\partial \tilde{z}}{\partial x} \frac{\partial x}{\partial t'} \quad \underbrace{\frac{\partial t}{\partial t'}}_1 \quad \underbrace{\frac{\partial x}{\partial t'}}_U$$

$$\frac{\partial \tilde{z}}{\partial t'} = \frac{\partial \tilde{z}}{\partial t} + U \frac{\partial \tilde{z}}{\partial x}$$

$$\begin{aligned} \frac{\partial^2 \tilde{z}}{\partial t'^2} &= \frac{\partial}{\partial t} \left[\frac{\partial \tilde{z}}{\partial t} + U \frac{\partial \tilde{z}}{\partial x} \right] + U \frac{\partial}{\partial x} \left[\frac{\partial \tilde{z}}{\partial t} + U \frac{\partial \tilde{z}}{\partial x} \right] \\ &= \left(\frac{\partial}{\partial t} + U \frac{\partial}{\partial x} \right)^2 \tilde{z} \end{aligned}$$

$$\begin{aligned} \frac{\partial^2 \tilde{z}(x,t)}{\partial x'^2} &= \frac{\partial \tilde{z}}{\partial x} \frac{\partial x}{\partial x'} + \frac{\partial \tilde{z}}{\partial t} \frac{\partial t}{\partial x'} \\ &= \frac{\partial \tilde{z}}{\partial x} \Rightarrow \frac{\partial^2 \tilde{z}}{\partial x'^2} = \frac{\partial^2 \tilde{z}}{\partial x^2} \end{aligned}$$

$$m \frac{\partial^2 \xi}{\partial t^2} = T \frac{\partial^2 \xi}{\partial x^2} + F_{ext}$$

$$m \left(\frac{\partial}{\partial t} + U \frac{\partial}{\partial x} \right)^2 \xi = T \frac{\partial^2 \xi}{\partial x^2} + F_{ext}$$

Non-dispersive

II. Fast and Slow Waves ($F_{ext} = 0$)

$$\xi(x, t) = \text{Re} \xi e^{j(\omega t - kx)}$$

$$\left(\frac{\partial}{\partial t} + U \frac{\partial}{\partial x} \right)^2 \xi = v_s^2 \frac{\partial^2 \xi}{\partial x^2} ; \quad v_s^2 = T/m$$

$$+(\omega - kU)^2 = +k^2 v_s^2$$

$$\omega = k(U \pm v_s)$$

$$k = \frac{\omega}{U \pm v_s}$$

$$= \frac{\omega(U \mp v_s)}{U^2 - v_s^2}$$

$$= n \pm \gamma$$

$$n = \frac{\omega U}{U^2 - v_s^2}, \quad \gamma = \frac{\omega v_s}{U^2 - v_s^2}$$

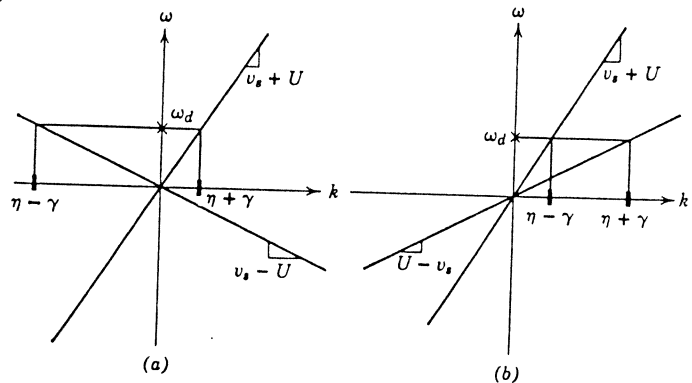


Fig. 10.2.8 Dispersion relations for waves on a string moving with the equilibrium velocity U in the positive x -direction. Although the geometry for the subsonic and supersonic cases is the same as that for the characteristic lines in Fig. 10.2.2, the axes here are ω - k , whereas in Fig. 10.2.2 they are x - t : (a) $U < v_s$; (b) $U > v_s$.

Supersonic Boundary Conditions ($U > v_s$): $\xi(0, t) = 0$

$$\frac{\partial \xi}{\partial x}(0, t) = \Delta_0 \cos \omega_d t$$

$$\begin{aligned} \xi(x, t) &= \text{Re} A e^{j(\omega_d t - nx)} \sin \gamma x \\ &= \frac{\Delta_0}{\gamma} \cos(\omega_d t - nx) \sin \gamma x \end{aligned}$$

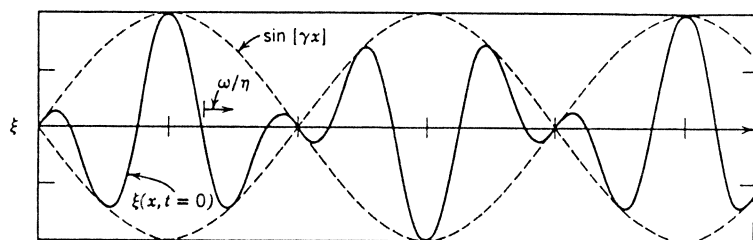


Fig. 10.2.9 Fast and slow waves interfere to form beats in space when $U > v_s$. The deflection is as the string appears if illuminated once each period of excitation $2\pi/\omega_d$.

III. Evanescence and Oscillation with Convection

$$(\omega - kU)^2 = k^2 v_s^2 + \omega_c^2$$

$$\omega^2 - 2kU\omega + k^2 U^2 = k^2 v_s^2 + \omega_c^2$$

$$k^2 (U^2 - v_s^2) - 2kU\omega + \omega^2 - \omega_c^2 = 0$$

$$k = \frac{\omega U \pm \sqrt{4\omega^2 v_s^2 - 4(U^2 - v_s^2)(\omega^2 - \omega_c^2)}}{2(U^2 - v_s^2)}$$

$$= n \pm \gamma$$

$$n = \frac{\omega U}{U^2 - v_s^2}, \quad \gamma = \frac{\sqrt{\omega^2 v_s^2 + (U^2 - v_s^2)\omega_c^2}}{(U^2 - v_s^2)}$$

If $U > v_s$, γ real (propagation without decay)
 $U < v_s$, γ can be imaginary

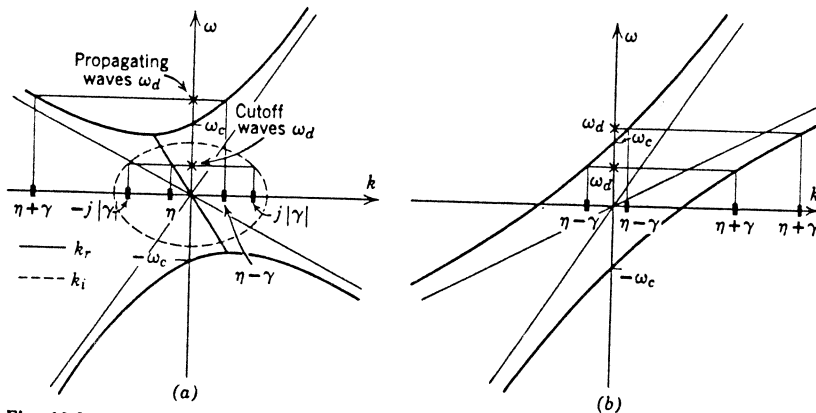


Fig. 10.2.10 Dispersion relations for system of Section 10.1.2, with the addition of a longitudinal velocity U . When U exceeds v_s , both waves propagate without decay, regardless of the frequency, and the effect of the convection is to eliminate the evanescence. (a) $U < v_s$; (b) $U > v_s$.

IV Convective Instability or Wave Amplification

$$(\omega - kU)^2 = k^2 v_s^2 - \omega_c^2$$

$$\omega = kU \pm \sqrt{k^2 v_s^2 - \omega_c^2}$$

$$k = \frac{\omega U \pm \sqrt{\omega^2 v_s^2 - (U^2 - v_s^2)(\omega^2 + \omega_c^2)}}{U^2 - v_s^2}$$

$$k = n \pm j\gamma$$

$$n = \frac{\omega U}{U^2 - v_s^2}, \quad \gamma = \frac{\sqrt{(U^2 - v_s^2)\omega_c^2 - \omega^2 v_s^2}}{U^2 - v_s^2}$$

k complex (γ real) if $\omega^2 < \frac{(U^2 - v_s^2)}{v_s^2} \omega_c^2$

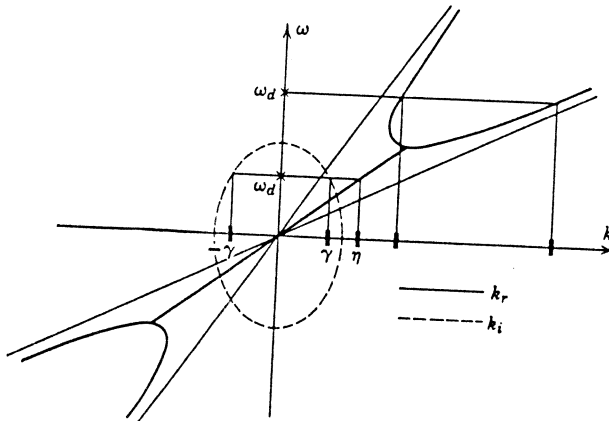


Fig. 10.2.13 Dispersion plot for supersonic string with a destabilizing magnetic force that shows complex values of k for real values of ω .

Dynamics of Electromechanical Continua

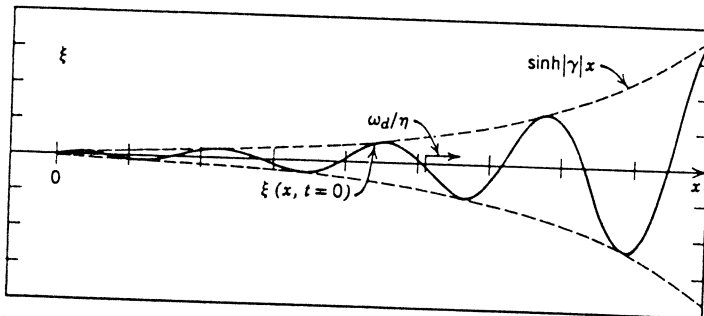


Fig. 10.2.14 Instantaneous view of string deflections when they are convectively unstable.

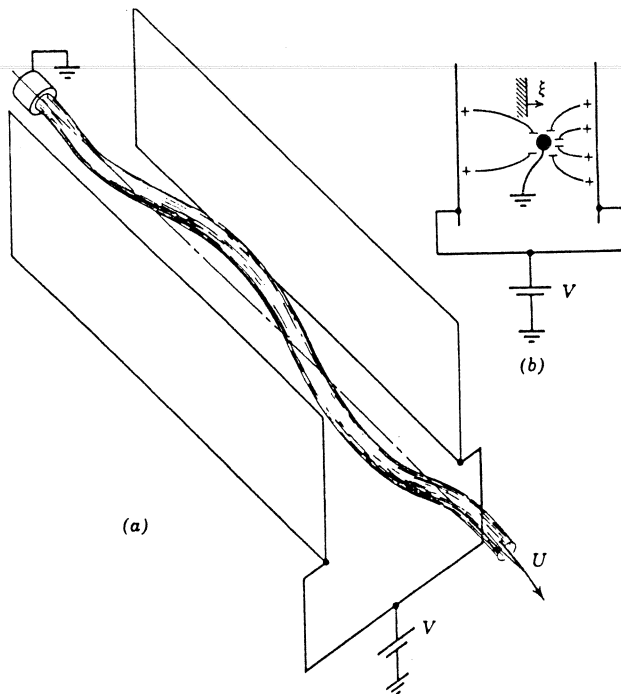


Fig. 10.2.15 (a) A jet of water has a streaming velocity U midway between plane-parallel electrodes. A potential difference V is applied between the plates and the jet. (b) As the jet is deflected toward one of the plates, an unbalance in the electric force of attraction tends to carry it even farther in the direction of deflection. The instability resulting from this force is washed downstream to form the convective instability sketched in Fig. 10.2.16.

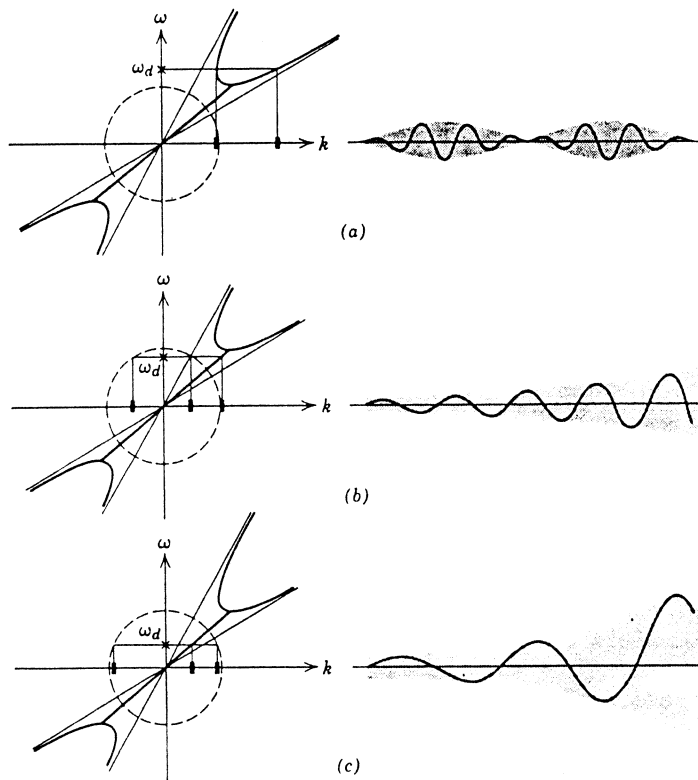


Fig. 10.2.16 Sinusoidal excitation of jet shown in Fig. 10.2.15. Driving frequency decreases from (a) to (c). This experiment together with animated ω - k plots can be seen in the film "Complex Waves II," produced for the National Committee on Electrical Engineering Films by the Education Development Center, Newton, Mass.