Search for Muon Neutrino and Antineutrino Disappearance in MiniBooNE


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The MiniBooNE Collaboration reports a search for \( \nu_\mu \) and \( \bar{\nu}_\mu \) disappearance in the \( \Delta m^2 \) region of 0.5–40 eV\(^2\). These measurements are important for constraining models with extra types of neutrinos, extra dimensions, and CPT violation. Fits to the shape of the \( \nu_\mu \) and \( \bar{\nu}_\mu \) energy spectra reveal no evidence for disappearance at the 90% confidence level (C.L.) in either mode. The test of \( \bar{\nu}_\mu \) disappearance probes a region below \( \Delta m^2 = 40 \text{ eV}^2 \) never explored before.

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Neutrino oscillations have been observed and have been confirmed in a variety of experiments, such as the Sudbury Neutrino Observatory (SNO) and the Maine-BooNE experiment. These observations have been consistent with three generations of neutrinos and a unitary mixing matrix. Complicating this picture, the Liquid Scintillator Neutrino Detector (LSND) experiment observed an excess of \( \bar{\nu}_e \) in a \( \bar{\nu}_e \) beam [1], indicating a possible third \( \Delta m^2 \) around 1 eV\(^2\) thus requiring more than three neutrino generations or other exotic physics. Recently, the MiniBooNE experiment [2] excluded two-neutrino appearance-only oscillations (98% C.L.) as an explanation of the LSND excess if oscillations of neutrinos and antineutrinos are the same.

Exotic physics models [3–6], including sterile neutrinos, extra dimensions, and CPT violation, have been proposed to explain the LSND observation. Some of these models can also accommodate the MiniBooNE \( \nu_e \) appearance oscillation results. These models are testable with measurements of \( \nu_\mu \) and \( \bar{\nu}_\mu \) disappearance which constrain any nonstandard oscillations of \( \nu_\mu \rightarrow \nu_x \). As described in this Letter, the MiniBooNE Collaboration has performed
based simulation, which was tuned using MiniBooNE
174 kA. The horn uses a positive current to focus
lium target surrounded by a magnetic horn pulsed at
1 cm diameter, 71 cm long (1.7 interaction length) beryl-
Neutrino Beam, produced by 8 GeV protons incident on a
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modified to include updated
MiniBooNE detector [9] is 541 m. The MiniBooNE
only disappearance fit is performed. The \( \nu_\mu \) flux to the
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Neutrino Beam, produced by 8 GeV protons incident on a
1 cm diameter, 71 cm long (1.7 interaction length) beryl-
ium target surrounded by a magnetic horn pulsed at 174 kA. The horn uses a positive current to focus \( \pi^+ \) and \( K^+ \) mesons for the neutrino mode sample, and a negative current to focus \( \pi^- \) and \( K^- \) for the antineutrino mode sample. The mesons that pass through a 60 cm diameter collimator 259 cm downstream of the target
decay in a 50 m long tunnel to produce the \( \nu_\mu \) beam. The Booster Neutrino Beam flux [10] is determined using a
GEANT4 [11] based beam simulation which has been further
modified to include updated \( p\)-Be particle production data
[12,13].

The distance from the proton interaction target to the
MiniBooNE detector [9] is 541 m. The MiniBooNE
detector is a 12 m diameter spherical tank filled with 800
tons of mineral oil (\( C_2 H_8 \)). The detector is separated
into an inner region filled with 1280 inward-facing 8 in.
photomultiplier tubes (PMTs), and an optically iso-
alted outer region used to reject cosmic-ray induced
events. Charged particles produced in neutrino inter-
actions emit primarily Cherenkov light, though a small
amount of scintillation light is also produced. Light and
particle production and propagation in the
MiniBooNE detector is modeled using a GEANT3 [14]
based simulation, which was tuned using MiniBooNE and external data.

Neutrino interactions are simulated with the v3 NUANCE
event generator [15]. Prior to selection, approximately
42% of all events in MiniBooNE are charged current
quasielastic (CCQE) scattering, and 22% are charged
current single charged pion production (CC1 \( \pi^+/\pi^- \)) in both the neutrino and antineutrino mode.

The search for oscillations is conducted with a sample
of CCQE events because of the high statistics and purity. The
reconstructed neutrino energy \( E_{\nu_\mu}^{QE} \) is calculated assum-
ing the target nucleon is at rest:

\[
E_{\nu_\mu}^{QE} = \frac{2(M_n - E_B)E_\mu - (E_B^2 - 2M_nE_B + M + M_n^2)}{2[(M_n - E_B) - E_\mu + p_\mu \cos\theta_\mu]},
\]

(1)

where \( \Delta M = M_n^2 - M_p^2; M \) indicates the muon, proton, or
neutron mass with appropriate subscripts; \( E_B \) is the nu-
eleon binding energy; \( E_\mu(p_\mu) \) is the reconstructed muon
energy (momentum); and \( \theta_\mu \) is the reconstructed muon
scattering angle with respect to the neutrino beam direc-
tion. A small correction is applied in both data and simu-
lation to account for the biasing effects of Fermi smearing.
At 300 MeV, the muon energy resolution is 7% and the
angular resolution is 5 degrees. The average \( E_{\nu_\mu}^{QE} \) resolution
is 11% for CCQE events [16].

A CCQE event sample is selected by identifying a single
muon in the detector and its associated decay electron,
using the same criteria as in the previous measurement of
CCQE model parameters on carbon [16]. Timing informa-
tion from the PMTs allows the light produced by the initial
nu neutrino interaction (first “subevent”) to be separated from
the light produced by the decay electron (second subevent).
The timing and charge response of the PMTs is then used
to reconstruct the position, kinetic energy, and direction
vector of the primary particle within each subevent.

Exactly two subevents are required in the analysis (the
muon and its decay electron). Requiring both subevents
to have fewer than six PMT hits anywhere in the veto
region rejects 99.99% of all cosmic-ray interactions [9].
The first subevent must be in coincidence with a beam
pulse, and have greater than 200 inner tank PMT hits,
to eliminate electrons from stopped cosmic-ray muon decays.

The mean emission point of the Cherenkov light along the
track for the first subevent must be less than 500 cm from
the center of the tank. The second subevent must have
fewer than 200 inner tank PMT hits, to
eliminate electrons from stopped cosmic-ray muon decays.

The selection criteria yield 190 454 data events with
\( 0 < E_{\nu_\mu}^{QE} < 1.9 \) GeV for \( 5.58 \times 10^{20} \) protons on target in the
neutrino mode sample, and 27 053 data events for
\( 3.39 \times 10^{20} \) protons on target in the antineutrino mode sample.

According to the selection, the neutrino mode sample is
74% pure CCQE, and the antineutrino mode sample is 70%
pure CCQE. The primary background (\( \sim 75% \)) for both the
\( \nu_\mu \) and \( \bar{\nu}_\mu \) samples is CC1 \( \pi^+ \) events where the outgoing
pion is unobserved (due, e.g., to absorption in the nucleus).

Though the neutrino mode sample has \(<1\% \bar{\nu}_\mu \) content, the
beam in the antineutrino mode contains a substantial contribution of $\nu_\mu$ due to the higher $\pi^+$ production at the target and the higher $\nu_\mu$ cross section. The antineutrino mode is predicted to have 25% $\nu_\mu$ content.

The CCQE cross section depends on the axial vector form factor, which is commonly assumed to have a dipole form as a function of four-momentum transfer ($Q^2$) with one adjustable parameter, $M_A$, the axial mass. Global fits to the world’s neutrino scattering data on deuterium yield $M_A = 1.015$ GeV [17]. However, recent results from K2K ($M_A = 1.14 \pm 0.11$ GeV, on carbon [18]; $M_A = 1.20 \pm 0.12$ GeV, on oxygen [19]) and MiniBooNE ($M_A = 1.23 \pm 0.12$ GeV, on carbon [16]) suggest a higher effective value of $M_A$ for nuclear targets. In addition, the level of Pauli blocking was adjusted in the MiniBooNE analysis, using a parameter $\kappa = 1.019$, to better reproduce the experimental data at low $Q^2$ [16]. The effect of $M_A$ and $\kappa$ on the $Q^2$ shape is pronounced, but oscillations would provide relatively little $Q^2$ distortion; this means that a spurious value of $M_A$ or $\kappa$ cannot be caused by underlying oscillations. The MiniBooNE CCQE $\nu_\mu$ analysis which produced a value of $M_A = 1.23$ assumed no oscillations, and therefore those values of $M_A$ and $\kappa$ should not be used in a disappearance analysis of the same data set. Consequently, the lower values of $M_A$ and $\kappa$ are used with conservative uncertainties which span the difference between the deuterium and nuclear target results ($M_A = 1.015 \pm 0.20$ GeV, $\kappa = 1.000 \pm 0.019$). The disappearance limits obtained were insensitive to the values of $M_A$ and $\kappa$ used. With $M_A = 1.015$ GeV and $\kappa = 1.000$, the ratio of detected events to predicted events in MiniBooNE is $1.31 \pm 0.26$ for neutrinos and $1.18 \pm 0.18$ for antineutrinos. The ratio for neutrinos reported in Ref. [16] is lower because higher values of $M_A$ and $\kappa$ were used there.

For the disappearance search, systematic uncertainties are included for the underlying neutrino flux prediction, neutrino interaction cross section, and detector response. The method used to estimate the uncertainties due to the underlying neutrino flux prediction and detector model is identical to the method used in previous MiniBooNE results [2,20]. The uncertainties on the cross section include uncertainties on the CCQE cross section and CC1 $\pi^+$ background. The latter is estimated using the MiniBooNE CC1 $\pi^+$ data sample. Systematic uncertainties produce correlated errors between $E_{\nu}^{\text{QE}}$ bins that are included by developing a covariance matrix in the same manner as in previous MiniBooNE oscillation analyses [2,20]. This covariance matrix includes separate normalization and shape-only error contributions. For the shape-only disappearance search, the prediction is normalized to data, and just the shape-only covariance matrix is used.

The disappearance search uses the Pearson’s $\chi^2$ test to determine allowed regions in the $\Delta m^2 - \sin^2 2\theta$ plane. The $\chi^2$ is calculated from a comparison of the data, $d_i$, in the $E_{\nu}^{\text{QE}}$ bin $i$, to a prediction $p_i(\Delta m^2, \sin^2 2\theta)$ for 16 bins. The prediction assumes a two-flavor $\nu_\mu \rightarrow \nu_e$ disappearance characterized by one large mass splitting ($\Delta m^2 \equiv \Delta m^2_{\nu_{\mu h}}$) between the light neutrino mass states $k$, which participate in standard three-neutrino oscillations, and $h$, the heavier neutrino state, and one oscillation amplitude $\sin^2 2\theta = 4|U_{\mu h}|^2(1 - |U_{\mu h}|^2)$, where $|U_{\mu h}|^2$ is the muon flavor content of the heavy state $h$:

$$\chi^2 = \sum_{i,j} (d_i - N_{pj}) M_{ij}^{-1}(d_j - N_{pj}).$$

where $M_{ij}$ is the shape-only error matrix, and $N$ is a factor which normalizes the prediction to the total number of observed events in data. All neutrino events in the prediction, including the CC1 $\pi^+$ background events, are allowed to oscillate in the fit based on the incident neutrino energy and distance traveled. The 90% C.L. limit corresponds to $\chi^2 > 23.5$ for 16 degrees of freedom (DF). The sensitivity is a fit to an unoscillated prediction including all statistical and systematic uncertainties.

The top plot of Fig. 1 (Fig. 2) shows the $E_{\nu}^{\text{QE}}$ spectrum after selection cuts for the neutrino (antineutrino) data and
the prediction, assuming no oscillations (null hypothesis) with diagonal elements of the error matrix. The dominant systematics arise from the neutrino flux (production of $\pi^+$ from $p$-Be interactions) and CCQE cross section uncertainties; uncertainties at low energy are larger because of the substantial CC1 $\pi^+$ background and uncertainties in the CCQE cross section in this region. As shown in Fig. 1, the individual bin errors are large, but adjacent bins are nearly fully correlated. The $\chi^2$ between the data and the null hypothesis is 17.78 (16 DF, 34% probability) for the neutrino mode sample, which is consistent with no oscillations at the 90% C.L. The top plot of Fig. 3 shows the 90% C.L. sensitivity and limit curves for the neutrino mode sample. The minimum $\chi^2 = 12.72$ (13 DF, 47% probability) at $\Delta m^2 = 17.5$ eV$^2$, $\sin^2 2\theta = 0.16$. The probability distribution and number of degrees of freedom for the $\chi^2$ statistic are determined from an analysis of a set of simulated data samples, as suggested in Ref. [21].

The bottom plot in Fig. 1 shows the ratio of data to the null hypothesis and three oscillation scenarios. The shape distortion for $\Delta m^2 = 0.5$ eV$^2$ is very different from $\Delta m^2 = 3.0$ eV$^2$. The $\chi^2$ therefore changes rapidly as a function of $\Delta m^2$, resulting in rapid changes in the 90% C.L. sensitivity and limit curves (Fig. 3) for small differences in $\Delta m^2$. Similar features are also seen in previous disappearance analyses [7,8].

The $\bar{\nu}_\mu$ disappearance analysis proceeds in the same manner as the $\nu_\mu$ analysis, except that only the $\bar{\nu}_\mu$ events are allowed to oscillate in the fit and the $\nu_\mu$ events are kept fixed. This determines the limit on a model where the $\bar{\nu}_\mu$ can oscillate but the $\nu_\mu$ cannot. A model where both $\nu_\mu$ and $\bar{\nu}_\mu$ oscillate with equal oscillation probability versus energy would produce a limit very similar to the neutrino mode limit.

During antineutrino data taking, two absorber plates inadvertently fell vertically into the decay volume at 25 m and were later removed, creating three distinct data taking periods with zero, one, or two absorbers in the beam line. The event rate was predicted to be 13% (20%) lower for one (two) plate(s) in the beam. One (two) absorber plate(s) were in the beam for 16.8% (18.1%) of the antineutrino data taking. Beam line monitoring systems indicated when each plate dropped. Because the changes to the beam line are understood, a separate simulation was run with the appropriate number of absorber plates in the beam line. Figure 2 shows the $E_y^{\text{QE}}$ distribution for the antineutrino mode sample. The $\chi^2$ of the null hypothesis is 13.7, 8.2, 15.2, 10.29 (16 DF) for the zero, one, and two absorber plate and total data, respectively. The antineutrino mode data are also consistent with no oscillations at the 90% C.L., so the bottom plot of Fig. 3 shows the 90% C.L. sensitivity and limit curves for the antineutrino disappear-

![FIG. 2. Same convention as Fig. 1 for the antineutrino mode sample. The background neutrino events are also shown (dotted line in top plot). Minimum $\chi^2 = 5.43$ (11 DF) is at $\Delta m^2 = 31.3$ eV$^2$, $\sin^2 2\theta = 0.96$.](image)

![FIG. 3. The top plot shows the sensitivity (dashed line) and limit (solid line) for 90% C.L. for neutrino disappearance in MiniBooNE. Previous limits by CCFR (dark grey) and CDHS (light grey) are also shown. The bottom plot uses the same convention for antineutrino disappearance.](image)
ance fit to all antineutrino data; the limit curves for the individual absorber data periods were found to be consistent within errors to the total.

In addition to the two-neutrino oscillation fits described above, two 3 + 2 sterile neutrino models are tested. A 3 + 2 model assumes two heavy neutrinos mix with the lighter neutrinos, with two mass splittings and two oscillation amplitudes. Global fits to existing appearance and disappearance data yield regions of allowed 3 + 2 model parameters. The best fit points in these allowed regions are tested with MiniBooNE data by forming a $\chi^2$ with a prediction assuming the full 3 + 2 oscillation formalism with the best fit 3 + 2 model parameters. The best fit 3 + 2 model in Ref. [4] is consistent with both MiniBooNE $\nu_\mu$ and $\bar{\nu}_\mu$ data. However, the $\nu_\mu$ data are incompatible with the best fit 3 + 2 model in Ref. [5] at the 90% C.L. with $\chi^2 = 24.7$ (16 DF).

In summary, MiniBooNE observes no evidence for $\nu_\mu$ or $\bar{\nu}_\mu$ disappearance at the 90% C.L. in the $\Delta m^2$ region of 0.5–40 eV$^2$. The test of $\bar{\nu}_\mu$ disappearance probes a region unexplored by previous experiments.

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