Detection of Fracture Orientation Using Azimuthal Variation of P-wave AVO Responses

by

Maria Auxiliadora Pérez

Submitted to the Department of Earth, Atmospheric, and Planetary Sciences
in partial fulfillment of the requirements for the degree of Master of Science

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Signature of Author

Department of Earth, Atmospheric, and Planetary Sciences
January 17, 1997

Certified by

Professor M. Nafi Toksöz
Director, Earth Resources Laboratory
Thesis Advisor

Accepted by

Professor Thomas H. Jordan
Department Head
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ABSTRACT

Azimuthally-dependent P-wave AVO (amplitude variation with offset) responses can be related to open fracture orientation and have been suggested as a geophysical tool to identify fracture orientation in fractured oil and gas reservoirs. A field experiment recently conducted over a fractured reservoir in the Barinas Basin (Venezuela) provides data for an excellent test of this approach. Three lines of data were collected in three different azimuths, and three component receivers were used. The distribution of fractures in this reservoir was previously obtained using measurements of shear wave splitting from P-S converted waves from the same dataset (Ata and Michelena, 1995). In this work, we use P-wave data to see if the data can yield the same information using azimuthal variation of P-wave AVO responses. Results obtained from the azimuthal P-wave AVO analysis corroborate the results previously obtained using P-S converted waves. Additionally, we modelled P-wave responses over the different azimuths (perpendicular and parallel to the fracture strike) using logs from nearby wells to design the earth model. The results obtained from the modeling show the same trends as the field data. This analysis with field data is an example of the high potential of P-waves to detect fracture effects on seismic wave propagation.

Thesis supervisor: M. Nafi Toksöz
Title: Professor of Geophysics

Thesis co-supervisor: Richard L. Gibson, Jr.
Title: Research Scientist
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Finally, I would like to dedicate this thesis to my husband Tony and my children, Daniel and Fabiana, for their dedication, patience and love, during the course of my graduate years.
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Chapter 1

Introduction

1.1 Background and Motivation

The detection of fractured zones and determining their orientation is an important part of reservoir development and enhanced oil recovery (EOR) projects. S-waves have been shown to be very effective in detecting azimuthal anisotropy and, more precisely, of fracture induced anisotropy. They are considered more reliable than P-waves for two main reasons. First, an anisotropic medium allows propagation of two quasi-shear waves with different polarizations. Measurements of travel times of these waves in a single propagation direction (vertical, for example) allows a definitive identification of anisotropy. In contrast, a single quasi-compressional wave exists. Its velocity can be affected by heterogeneity as well as anisotropy, and its travel time must be examined in many directions. Second, the polarization of the two quasi-shear waves in a fractured reservoir is related to fracture orientation, one perpendicular to cracks, and the other polarized parallel to the cracks. Therefore, crack orientation, as well as anisotropy, can be determined with a shear wave experiment. However, the acquisition and processing of S-waves is very expensive compared to conventional P-wave data. Therefore, the characterization of fractured reservoirs using P-waves instead of S-waves is an important exploration problem that has attracted much attention from ex-
ploration geophysicists and reservoir engineers interested in fracture detection and analysis. Some theoretical works show that P-waves are sensitive to fractures or cracks (Crampin, 1980). In particular, P-wave reflection AVO has been suggested as an indicator of azimuthal anisotropy (Mallick and Frazer, 1991; Rüger and Tsvankin, 1995; Strahilevitz and Gardner, 1995). There are also some field studies of azimuthally-dependent P-wave AVO responses related to fractured reservoirs (Lefeuvre, 1994; Lynn et al., 1995).

Amplitude Variation with Offset (AVO) analysis is based on the variation of reflection and transmission coefficients with incident angle and the corresponding increasing offset (Castagna, 1993a). Different equations have been presented in the literature for the reflection coefficients (Knott, 1899; Zoeppritz, 1919; Aki and Richards, 1980; Waters, 1981). These equations are very complex. A simplified approximation of the reflection coefficient for isotropic media was presented by Shuey (1985), for restricted angles of incidence, the equation is

$$R_{pp}(\theta) = R_p + B \sin^2(\theta)$$  \hspace{1cm} (1.1)

where $R_{pp}(\theta)$ is the reflection coefficient at angle $\theta$, $R_p$ is the reflection coefficient at normal incidence (also called the AVO intercept), and $B$ is called the AVO gradient which is mainly influenced by variation in Poisson’s ratio ($\sigma$):

$$B = A_0 * R_p + \frac{\Delta\sigma}{(1 - \sigma)^2}$$  \hspace{1cm} (1.2)

$$A_0 = B_0 - 2(1 + B_0) * \frac{(1 - 2 * \sigma) / (1 - \sigma)}{(1 - \sigma)}$$  \hspace{1cm} (1.3)

$$B_0 = \frac{((\Delta V_p)/V_{pa}) / (((\Delta V_p)/V_{pa}) + (\Delta \rho / \rho_a))}{(1)}$$  \hspace{1cm} (1.4)

where

$$\Delta V_p = V_{p2} - V_{p1}$$

$$V_{pa} = (V_{p2} + V_{p1}) / 2$$

$$\Delta \rho = \rho_2 - \rho_1$$
\[ \rho_a = \frac{\rho_2 + \rho_1}{2} \]

\[ \Delta \sigma = \sigma_2 - \sigma_1 \]

\[ \sigma_a = \frac{\sigma_2 + \sigma_1}{2}. \]

Therefore, \( R_p \) dominates the reflection coefficient at small angles, whereas \( \Delta \sigma \) and consequently \( V_p/V_s \) contrast dominates at larger angles. However, the term \( B \) in Shuey’s equation 1.1 is still complicated. Thomsen (1990) suggests that Wright’s (1986) reflection coefficient equation has a simpler expression for the term \( B \) (AVO gradient)

\[ B = \frac{1}{2} \left( \frac{V_p}{V_p} \right) - \left( \frac{2 \cdot V_{sa}/V_p}{\Delta \mu/\mu_a} \right) \]

where

\[ V_{sa} = \frac{V_{s2} + V_{s1}}{2} \]

\[ \Delta \mu = \mu_2 - \mu_1 \]

\[ \mu_a = \frac{\mu_2 + \mu_1}{2}. \]

Therefore, the equation for the reflection coefficient in this case is given by

\[ R_{pp}(\theta) = R_p + \frac{1}{2} \left( \frac{\Delta V_p}{V_p} \right) - \left( \frac{2 \cdot V_{sa}/V_p}{\Delta \mu/\mu_a} \right) \]

\[ \cdot \sin^2(\theta). \]

It can be observed in equation 1.6 that any change in \( V_p \) clearly affects the resulting reflection coefficient. A detailed description of AVO theory is presented in Appendix A.

Gassmann’s (1951) equations predict a large drop in P-wave velocity and a small increase in S-wave velocity when even a small amount of gas is introduced into the pore space of a compressible brine-saturated sand. This drop causes a drop in \( V_p/V_s \) resulting in AVO anomalies. Oil and water are often assumed to have similar acoustic properties, and to be indistinguishable using seismic methods. However, gas or light hydrocarbons can go into solution in crude oils, and can dramatically alter the velocities (Castagna, 1993b). Therefore,
depending on the gas-oil-ratio (GOR), the ratio $V_p/V_s$ can be strongly affected, likewise the AVO response.

We analyze the fracture reservoir data with the following conceptual model (Figures 1-1 and 1-2). If the experiment is conducted parallel to fracture orientation, the fractures should have minimal influence on the reflection properties, regardless of the angle of incidence. This is because the P-wave particle motion will always be parallel to the thin cracks. However, if the line is oriented more perpendicular to the fractures, at large angles of incidence, the reflection coefficients will be affected strongly (Lynn et al., 1995). At large angles of incidence, the P-wave velocity is expected to be affected by the acoustic properties of the fluid filling the fractures when the wave propagation is perpendicular to the fractures, while it is less affected when the wave propagation is parallel to the fractures.

If no azimuthal anisotropy exists, the AVO response will be the same in all directions, while in the presence of anisotropy, the AVO response will vary depending on the source-receiver azimuth. In our work, we will apply AVO analysis to a known fractured reservoir, which has azimuthal anisotropy associated with the aligned, vertical fractures. In this case, Wright’s (1986) equation will not be directly applicable for the study of P-wave reflections because it assumes isotropic formations. However, for P-wave AVO, the reflection coefficient at normal incidence (AVO intercept) is independent of azimuth. We therefore assume that the reflection coefficient at small angles of incidence will still have the same form as equation 1.6, but that the precise, explicit form of the AVO gradient term $B$ will depend on the experimental azimuth. Some initial theoretical results for P-wave reflection coefficients for symmetry planes and AVO attributes in azimuthally anisotropic media (Rüger, 1996) confirm this hypothesis. For incident angle $\theta$ and azimuthal angle $\phi$ the reflection coefficient is given by

$$R_{pp}(\theta, \phi) = R_p + (1/2) \times ((\Delta V_p/V_{pa}) - ((2 \times V_{sa}/V_{pa})^2 \times (\Delta \mu/\mu_a))$$

$$+ (\Delta \delta + 2 \times (2 \times V_{sa}/V_{pa}) \times \Delta \omega) \times \cos^2(\phi) \times \sin^2(\theta).$$

For the particular case perpendicular to the symmetry axis $\phi = 90$ (parallel to the fracture orientation), equation 1.7 reduces to the approximation given by equation 1.6 for the
isotropic case, and for any other case, any change in $V_p$ will also affect the reflection coefficient.

Some examples of calculated reflection coefficients for models of azimuthally anisotropic reservoirs illustrate this behavior. In Figures 1-3 and 1-4 we show the reflection coefficients for two models with a fractured layer over an isotropic layer with different velocity contrasts. Model 1 corresponds well to the model where the azimuthal P-wave AVO response is found in our field data further on in our study. We use Hudson’s (1981) model to obtain the elastic constants for the fractured media. Then, we calculate the P-wave reflection coefficients solving the boundary conditions at the boundary of the two layers. The difference in P-wave AVO can be observed in the figures for lines parallel and perpendicular to the fractures. Thus, we can demonstrate that P-wave AVO is possibly a good indicator of azimuthal anisotropy. We expect to obtain similar results from our P-wave AVO analysis over the field data.

In a previous work, Ata and Michelen (1995) estimated the fracture orientation in a fracture reservoir using splitting measurements from the P-S converted data. The purpose of our study is to attempt to identify azimuthal P-wave AVO over the same fractured oil reservoir and compare it to previous independent results obtained using P-S converted waves to see if we can obtain similar results to identify fracture orientation.

1.2 Outline

In this chapter, we studied the azimuthal variation of AVO of field data in order to evaluate its applicability to the production geophysics environment.

In Chapter 2, we present a brief description of the area were the data were acquired, some details about the data acquisition, P and P-S converted wave data processing, AVO analysis results, and some discussion.
In Chapter 3, we present some results of the ray trace modeling using nearby well logs to build an earth model. We also compare the results obtained from the modeling to the field data.

Finally, Chapter 4 summarizes the results obtained in this work. Additionally, we describe some future related work.

Various appendices complement this thesis.
Figure 1-1: Conceptual model. If the source-receiver azimuth is perpendicular to fracture orientation, at large angles of incidence, the reflection coefficients will be affected strongly by the hydrocarbons if there is sufficient gas in the oil.
Figure 1-2: Effects of fractures in wave propagation. At large angles of incidence, the P-wave velocity is expected to be affected by the acoustic properties of the fluid filling the fractures when the wave propagation is perpendicular to the fractures, while it is less affected when the wave propagation is parallel to the fractures.
Layer 1: Fractured layer
- $V_p = 4.99$ Km/s
- $V_s = 2.78$ Km/s
- $\rho = 2.7$ g/cm$^3$

Layer 2: Isotropic layer
- $V_p = 4.60$ Km/s
- $V_s = 2.65$ Km/s
- $\rho = 2.7$ g/cm$^3$

Figure 1-3: Reflection coefficients vs. offset for a fractured medium over an isotropic media (Model 1).
Layer 1: Fractured layer
- $V_p = 4.99$ Km/s
- $V_s = 2.78$ Km/s
- $\rho = 2.7$ g/cm$^3$

Layer 2: Isotropic layer
- $V_p = 5.30$ Km/s
- $V_s = 2.65$ Km/s
- $\rho = 2.7$ g/cm$^3$

Figure 1-4: Reflection coefficients vs. offset for a fractured medium over an isotropic media (Model 2).
Chapter 2

Field Data: Maporal Field

2.1 Characteristic of the Field

The Maporal field is located in the north-central part of the Barinas-Apure Basin. Structurally, the Maporal field is a dome slightly extended in the NE direction. The geological setting is composed mainly of nearly flat-lying sediments. Two fault systems are present in the area. One runs southeast-northwest and the other northeast-southwest. A brief description of the main formations present in the zone is shown in Figure 2-1. The target zone is the member ‘O’ of the ‘Escandalosa’ formation. It is a fractured limestone at a depth of approximately 3000 m (2.32 s). The fractures in the reservoir are filled with crude oil of approximately 28 API number. There is also evidence that some production comes from the ‘P’ and ‘R’ members of the stratigraphic section. These are composed mainly of sandstones; fractures also exist in these two members.

A number of wells exist in the field, which provide good background information on reservoir and fracture properties in the area, and which can be used for correlation studies. The well data includes sonic, dipole(shear), caliper, resistivity, gamma ray and other logs.

Knowing the orientation of the maximum horizontal stress in the field is important to deter-
mine the fractures that are more likely to be open or to be closed. Some maps of maximum horizontal stress and fracture strike in the field have been estimated using well logs (Figures 2-2 and 2-3). The maximum horizontal stress runs southeast-northwest and different fracture systems have been identified. Based on the theory that open fractures tend to align along the direction of maximum regional stress, and consequently that perpendicular fractures may be closed, the fracture system of interest should be oriented southeast-northwest.

As we can see, the Maporal field is well suited for this study. The existence of fracturing is confirmed from in situ borehole observations, and we also have a preliminary estimate of orientation.

2.2 Data Acquisition

Three 10 Km three-component seismic lines were recorded over the area of interest along three different azimuths. The lines have a single intersection point. As was mentioned above, two systems of normal faults are present in the area. The azimuths of two lines (lines 1 and 3) were almost parallel to the fault systems, and the third line (line 2) almost bisected them and formed an angle of approximately 40 degrees with line 1. Figure 2-4 is an illustration of survey geometry with respect to the fault systems. A charge of one kilogram explosives at 10 m depth was used for the source. The source interval was set at 51 m and the geophone group interval at 17 m with a linear array between stations. The near offset trace was set at 17 m and the far offset extended to 3600 m, which satisfies the suggested minimum offset for an azimuthal AVO study (at least the same as the depth of the target zone). With this geometry we can expect a maximum incident angle of approximately 32 degrees for the target zone. A summary of the geometry of the survey and recording parameters is given in Table 2.1.
<table>
<thead>
<tr>
<th>GEOMETRY</th>
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<tbody>
<tr>
<td>Number of traces</td>
<td>216</td>
</tr>
<tr>
<td>Total number of 3C receivers</td>
<td>648</td>
</tr>
<tr>
<td>Near offset trace</td>
<td>17 m (55ft)</td>
</tr>
<tr>
<td>Far offset trace</td>
<td>3672 m (12,000ft)</td>
</tr>
<tr>
<td>Number of elements per string</td>
<td>6</td>
</tr>
<tr>
<td>Separation of geophone elements</td>
<td>3.4 m (11ft)</td>
</tr>
<tr>
<td>Total geophone spread</td>
<td>10 Km (33,000ft)</td>
</tr>
<tr>
<td>Total number of shotpoints per line</td>
<td>124</td>
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<tr>
<td>Shotpoint space interval</td>
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<tr>
<td>Charge depth</td>
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<td>Charge size</td>
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<tr>
<td>Source offset</td>
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<tr>
<td>RECORdING</td>
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<tr>
<td>High cut</td>
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<tr>
<td>Notch filter</td>
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</tr>
<tr>
<td>Record length</td>
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</tr>
</tbody>
</table>

Table 2.1: Survey parameters

2.3 Previous Studies

A study of S-wave birefringence, or splitting, was previously conducted using this dataset by Ata and Michelena (1995). Due to the simple structure of the area, they used the asymptotic approximation (Tessmer and Behle, 1988; Tessmer et al., 1990) to calculate the
Common Conversion Point (CCP). Then, they applied rotation analysis to align observed data to the principal axes of symmetry and therefore to estimate the fracture orientation. They also made travel time studies ($S_1$ and $S_2$ modes) to estimate the fracture density. As a result of this study a map showing the fracture orientation at the reservoir level was generated and is shown in Figure 2-5. Ata and Michelena (1995) concluded that lines 1 and 2 are nearly-perpendicular to fracture orientation, and line 3 is nearly-parallel to fracture orientation. We compare these results to our analysis.

2.4 AVO Study

The first goal of our work is to perform AVO analysis of the P-wave data. Specifically, we focus on the intersection point of the three seismic lines. At this point, all lines should be looking at the same point in the subsurface and we will compare the differences in the azimuthal response.

2.4.1 P-wave Data

2.4.1.1 Data Processing

The basic objective of the data processing for AVO analysis is to preserve relative amplitude for all offsets at all times for any CDP gather. Additionally, relative amplitudes among all depth points need to be preserved. In the case of stacked data, the process of CDP stacking cancels many types of noise. However, in the prestack domain the processor cannot count on this tool. Many processing schemes for AVO analysis have been presented in the literature and some of them are presented in Appendix B.

The same data processing sequence was applied to all lines using a basic but robust scheme in order to conserve the reflectivity variation with offset. A subset of the data was extracted from the original one to consider the area of the reservoir near the intersection point. The
number of shots selected guaranteed maximum fold over the cross point (36 traces). Special emphasis in the processing was made around the cross point of the lines, at approximately CDP 224 of the reduced dataset. The processing sequence applied to all lines was:

1. Spherical divergence correction:
   The data was received with spherical divergence applied. Therefore, no correction was necessary at this point.

2. Coherent noise suppression:
   F-k filtering was applied to the data in order to suppress coherent noise such as ground roll and guided waves.

3. Velocity analysis and statics iterations:
   Several iterations of semblance velocity analysis and statics corrections were applied to obtain a better quality in the NMO-corrected CDP gathers. The semblance velocity analysis at the interception point for the three lines is similar and is presented in Figure 2-6. The maximum power autostatics method was applied to all lines in order to maximize the CDP stack power.

4. NMO correction:
   Standard NMO correction was applied to all CDPs. It is known that the NMO stretch decreases frequency at far offsets, which affects the amplitudes. However, this distortion is significant at shallow depths and at a large offset. The target zone in our study is relatively deep and the stretch problem does not significantly affect the results. NMO-corrected gathers around the intersection point for the three lines are presented in Figures 2-7, 2-8 and 2-9.

5. AVO analysis:
   AVO analysis was made over the three lines. A detailed description of the procedures and results of the AVO analysis are presented in the next section.

6. Stacking:
Finally, we stack the three lines. The results obtained for the three lines are presented in Figures 2-10, 2-11 and 2-12. In these figures we can corroborate the simple structure of the area described above.

In addition, we made a spectral analysis over the data; the center frequency was found to be around 25 Hz.

As it is known, deconvolution is a process that improves the temporal resolution by compressing the seismic wavelet. It is often used to isolate seismic events, which is important for AVO analysis. We point out that we did not use deconvolution in order to conserve amplitude information. However, the results obtained from our data processing correspond well to the results obtained from our ray tracing modeling, which is shown in Chapter 3.

Finally, it is important to note the consistency of the processing among all lines, which is supported by the good tie between the three lines (see Figure 2-13). The top of the fractured reservoir is located at 2320 ms and the base at 2370 ms. A higher frequency can be observed at the reservoir level for the line parallel to fracture orientation (line 3). However, these results are consistent with the ones obtained by Lynn et al. (1996), where the frequency content is related to source-receiver azimuth relative to fracture strike.

2.4.1.2 AVO Analysis

AVO analysis was made for each line and the AVO gradient and intercept were obtained. For AVO analysis, the PROMAX module AVO Attribute Stack was used. AVO Attribute Stack is a tool used to analyze the relative amplitudes within a pre-stack, NMO-corrected CDP gather. A common problem with AVO analysis is residual NMO on the CDP gathers resulting from imperfect velocity specification. In order to compensate for this problem, this tool allows one to select the amplitudes within polarity gates (see Figure 2-14), reducing the requirement for exact NMO application. Polarity gates are determined for partial stack of the data in the CDP gather. Contiguous stack amplitudes with a common polarity define a
polarity gate. The data within each polarity gate is condensed, selecting either the maximum or the average amplitude on each trace within the gate. We used the maximum option in our analysis. A robust least squares line is fitted to the amplitudes of each trace, and the AVO intercept and AVO gradient are computed from the fit to the data (see Figure 2-15).

An AVO anomaly with a large positive gradient was found near the base of the Escandalosa formation in lines 1 and 2, perpendicular to fracture orientation (see Figures 2-16 and 2-17). This AVO anomaly is not present in line 3, parallel to fracture orientation (see Figure 2-18). The AVO gradient versus AVO intercept graph for CDP gathers of each line around the cross point shows the high positive gradients for the lines perpendicular to the fracture, while they are smaller for the line parallel to the fracture (see Figure 2-19).

The AVO intercept for the bottom of the reservoir is very small for all lines (see Figures 2-20, 2-21 and 2-22). Theoretically, it should be the same for all lines no matter the source-receiver azimuth relative to fracture strike. However, some small differences were observed that could be caused by some noise in the data.

A detailed example of the behavior of the amplitudes for the cross point CDP gather for each line is shown in Figures 2-23a, 2-23b and 2-23c. These figures also show that an increase of the amplitude with the offset is larger for the lines perpendicular to the orientation of the fractures (lines 1 and 2) than for the line parallel to the orientation of the fractures (line 3).

2.4.1.3 Discussion

To start the discussion of the obtained AVO response, we present two of the five rules empirically established by Koefoed (1955) about the effects of Poisson's ratio on the reflection coefficients of plane waves:

1. “When the underlying medium has the greater longitudinal P-wave velocity and the other relevant properties of the two strata are equal to each other, an increase of
Poisson’s ratio for the underlying medium causes an increase of the reflection coefficient at the larger angles of incidence.” Shuey (1985) points out that the qualifications concerning the medium properties are not necessary for this and the next rule.

2. “When, in the above case, Poisson’s ratio for the incident medium is increased, the reflection coefficient at the larger angles of incidence is thereby decreased.”

The azimuthal difference in AVO gradients can be explained by a low Vp/Vs ratio for the lines perpendicular to the fracture due the influence of the fluid filling the fractures, resulting in a large positive contrast in Poisson’s ratio with the lower layer, and a very positive AVO gradient anomaly. For the line parallel to the fracture, the Vp/Vs ratio is less affected by the fluid filling the fractures and is greater than the previous case. This corresponds to rule number 2 and therefore the AVO anomaly is not present.

Until this point, we have been analyzing the results with the isotropic idea in mind. As we mentioned before, it has been shown that the reflection coefficient response depends on the source-receiver azimuth related to fracture orientation. However, there are only a few studies (e.g., Rüger, 1996) that relate AVO attributes to crack parameters. What seems to be clear is that the P-wave AVO gradient is affected by fracture-induced azimuthal anisotropy. In order to analyze our results using azimuthal differences in the reflection coefficients due to source-receiver azimuths, we calculate some reflection coefficients curves in two azimuths (one parallel and the other perpendicular to fracture orientation), and compare them to our results. We use information about the velocities (Vp and Vs) and densities at the bottom of the reservoir from nearby well logs to build our two-layer model. Layer 1 is a fractured layer and layer 2 is an isotropic one (see Figure 2-24). We use Hudson’s (1981) theory to estimate the elastic constants, using a fracture density of 0.1, and a bulk modulus for fluid filling the fractures that approximates an oil of 28 API (Batzle and Wang, 1992), as the one we have in the reservoir. Thomsen’s (1986) anisotropic coefficients for the fractured layer are shown in Table 2.2. Different results were obtained changing only Vs, and consequently the relation of the Vp/Vs ratio between the two layers (see Figures 2-25, 2-26 and 2-27). Model 1 corresponds very well to the results obtained from our data where the AVO gradient is
bigger for the lines perpendicular to fractures (lines 1 and 2) than for the line parallel to the fractures (line 3).

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<td>γ</td>
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<tr>
<td>δ</td>
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Table 2.2: Thomsen's anisotropic coefficients for the fractured layer in models 1, 2 and 3

No azimuthal AVO anomaly is apparently observed at the top of the fractured reservoir. We calculate the reflection coefficient for this interface (see Figure 2-28), using data from well logs (see Figure 2-29). Thomsen's anisotropic coefficient for the fractured layer in this model, are presented in Table 2.3. We can observe a difference between the two lines (perpendicular and parallel to fractures), however it is smaller than for the bottom of the reservoir, and could also be reduced for some tuning effects. We do not get the anomaly we expect from the top of the reservoir from our field data, but we show in Chapter 3 that this behavior can be reproduced by synthetic seismograms.

<table>
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<tr>
<td>δ</td>
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</table>

Table 2.3: Thomsen's anisotropic coefficients for the fractured layer in model 4
Figure 2-1: Lithology of the reservoir.
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B- AVO gradient, slope

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Figure 2-24: Reflection coefficients for models 1, 2 and 3 are calculated at * point.
Layer 1: Fractured layer

\[ V_p = 4.99 \text{ Km/s} \]
\[ V_s = 2.78 \text{ Km/s} \]
\[ \rho = 2.5 \text{ g/cm}^3 \]

Layer 2: Isotropic layer

\[ V_p = 4.75 \text{ Km/s} \]
\[ V_s = 3.345 \text{ Km/s} \]
\[ \rho = 2.2 \text{ g/cm}^3 \]

Figure 2-25: Reflection coefficients vs. angle at the bottom of the reservoir (Model 1).
Layer 1: Fractured layer
\[ V_p = 4.99 \text{ Km/s} \]
\[ V_s = 2.78 \text{ Km/s} \]
\[ \rho = 2.5 \text{ g/cm}^3 \]
Layer 2: Isotropic layer
\[ V_p = 4.75 \text{ Km/s} \]
\[ V_s = 2.845 \text{ Km/s} \]
\[ \rho = 2.2 \text{ g/cm}^3 \]

Figure 2-26: Reflection coefficients vs. angle at the bottom of the reservoir (Model 2).
Layer 1: Fractured layer
  $V_p = 4.99 \text{ Km/s}$
  $V_s = 2.78 \text{ Km/s}$
  $\rho = 2.5 \text{ g/cm}^3$

Layer 2: Isotropic layer
  $V_p = 4.75 \text{ Km/s}$
  $V_s = 2.345 \text{ Km/s}$
  $\rho = 2.2 \text{ g/cm}^3$

Figure 2-27: Reflection coefficients vs. angle at the bottom of the reservoir (Model 3).
Figure 2.28: Reflection coefficient for model 4 is calculated at * point.
Layer 1: Isotropic layer
   $V_p = 3.345 \text{ Km/s}$
   $V_s = 1.580 \text{ Km/s}$
   $\rho = 2.4 \text{ g/cm}^3$

Layer 2: Fractured layer
   $V_p = 4.99 \text{ Km/s}$
   $V_s = 2.78 \text{ Km/s}$
   $\rho = 2.5 \text{ g/cm}^3$

Figure 2-29: Reflection coefficients vs. angle at the top of the reservoir (Model 4)
Chapter 3

Modeling

AVO analysis shows an anomaly near the bottom of the Escandalosa formation (reservoir). In order to interpret the results and confirm our conclusions from AVO analysis, we built a realistic earth model at the cross point of the three lines using well log data. Synthetic seismograms were computed using ray tracing to simulate the field data. The results show that the P-wave AVO can reproduce the type of anomaly observed in the field data.

3.1 Well Logs

The data used in this study correspond to a producing area. Consequently, many wells are present in the area, most of them with different types of logs. Three different wells were studied to build an earth model that is realistic and that reproduce seismic properties of the relevant formations. The wells were well MA-1, well MA-17 and well TOR-2X (see Figure 3-1). As was mentioned previously, the geology of the area is very simple, and all main formations are present across it. This allows a good correlation between wells in the interpretation process.

Well MA-17 has a dipole sonic log, which is very convenient because it yields both S and
P-wave velocities (see Figure 3-2). Unfortunately, well MA-17 does not extend to the bottom of the reservoir where the AVO anomaly was found in the field data. Well MA-1 does have a resistivity log that extends to the basement. P-wave velocities were estimated from the resistivity log using Faust’s (1953) approximation. Faust found an empirical formula for velocity in terms of depth of burial \( z \) and formation resistivity \( R \)

\[
V_p = 900(zR)^{1/6}
\]  

(3.1)

\( V_p \) being in \( m/s \), \( z \) in \( m \), and \( R \) in \( \Omega.m \). Then, densities were estimated using Gardner’s rule. Gardner et al. (1974) graphed velocity against density and found that the major sedimentary lithologies defined a relatively narrow swath across the graph. They determined an empirical equation relating velocity and density, often called Gardner’s rule:

\[
\rho = aV^{1/4}
\]  

(3.2)

where density \( \rho \) is in \( g/cm^3 \), \( a = 0.31 \) when velocity is in \( m/s \) and \( a = 0.23 \) when \( V \) is in \( ft/s \). Estimated density and P-wave velocity logs are shown in Figure 3-3. Finally, S-wave velocities were obtained from P-wave velocities using different ratios \( V_p/V_s \), between 1.5 and 2.0. Well TOR-2X does not belong to the Maporal field, but it is located in a field nearby in the Barinas basin. The geology is very similar to all main formations at the Maporal field also present in this area. Well TOR-2X has a sonic log (see Figure 3-4) that extends to the basement, and was used to calibrate the velocities obtained from the resistivity log in well MA-1. Velocities for each formation were compared in order to validate the approximated ones obtained in well MA-1.

### 3.2 The Model

A flat layer structure was assumed for the earth model of the Maporal field. This is a realistic model of the real geological structure. Additionally, our study was concentrated at the cross point of the three lines, and no other structural details were needed. An eight layer model was built (see Figure 3-5) using the well logs presented in the previous section. The elastic properties of formations Gobernador, Burguita and Quevedo were averaged and
modeled as a single layer. The ‘O’ member of the Escandalosa formation (reservoir), has the primary fracture concentration, and was modeled as an anisotropic layer due to vertical fractures, as is evidenced from the televiewer logs of nearby wells. There is also evidence from the televiewer logs of the presence of fractures in the members ‘P’ and ‘R’ of the same Escandalosa formation. Therefore, these two members were also modeled as vertical fractured layers, even though the fracture density modeled was smaller than for the ‘O’ member. All other layers were modeled as isotropic layers.

We computed the elastic constants for the anisotropic layers using Hudson’s (1981) theory for the properties of the fractured layer. Different experiments were made changing the fluid content in the fractures to analyze the effect of fluid in the P-wave AVO seismic response. In our study we label as light oil those with a high API number, which have higher gas content, and heavy oils, which have low API. The API number is about 100 for light condensates and nearly five for very heavy oils. The fractures in the reservoir are filled with crude oil of approximately 28 API, which we call medium oil. In our experiments we used the following as fluids: gas, light oil, medium oil, heavy oil and water. To calculate the elastic constants for the different fluids filling the fractures in the media, we use an approximation of the bulk modulus for different API numbers (Batzle and Wang, 1992). A detailed table including computational parameters and model material properties, for the fracture layers (‘O’ member and ‘P’ and ‘R’ sandstones), used for all cases (gas, light oil, medium oil, heavy oil and water) are given in Appendix C.

### 3.3 Synthetic Data

#### 3.3.1 Synthetic Seismograms

We computed synthetic seismograms for the eight layer model, using 3D paraxial ray tracing (Gibson et al., 1991). Two sets of seismograms were calculated for each case (gas, light oil, medium oil, heavy oil and water), one perpendicular to fracture orientation and
the other parallel to it. Seismograms were calculated in shot gathers, using 25 receivers that guarantee the same minimum (17 m) and maximum offset (3672 m) as the field data. Because we have a horizontal layer model, this geometry is equivalent to a CDP gather at the cross point. The wavelet was extracted from the field data in order to use a similar one in the modeling process. An explosive type of source was used and the wavelet correspond to the function \( \text{symmetrical} - \text{exponential} \times \text{cosine} \). A center frequency of 25 Hz was used to calculate the seismograms, as was found in the field data. Geometrical spreading was neglected in the calculations, so no additional processing was required to eliminate this effect. We processed each gather and obtained the NMO corrected gathers. The synthetic seismograms for one model (parallel to fractures filled with medium oil) were compared to the field data by copying the trace obtained by stacking the synthetic CDP gather to create a small section. The good correlation of the seismic events between the modeled data and the field data is shown in Figure 3-6. Some mismatches are observed between the first seismic event and the top of the reservoir because the formations Gobernador, Burguita and Quevedo were averaged as a single layer.

### 3.3.2 AVO Analysis

AVO analysis was made for each gather using the same AVO Attribute Stack PROMAX tool, that was previously used with the field data. The AVO gradient was obtained for each gather. Figures 3-7 and 3-8 show the AVO gradient for the wave propagation perpendicular and parallel to the fractures filled with gas, Figures 3-9 and 3-10 show the case of fractures filled with light oil, Figures 3-11 and 3-12 show the case of fractures filled with medium oil and Figures 3-13 and 3-14 show the case of fractures filled heavy oil. Finally, Figures 3-15 and 3-16 show the results of the AVO gradient obtained for the case of water filling the fractures. These figures show that the AVO gradient at the bottom of the reservoir is greater for the lines perpendicular to the fractures compared to the lines parallel to them. It can also be observed that this difference is greater when the gas content in the fluid filling the fractures is also larger.
Detailed graphs of the P-wave amplitudes vs. offset for the bottom of the reservoir were obtained from the synthetic seismograms for all of the cases (Figure 3-17). The percentage differences in the AVO gradient, between the lines perpendicular and parallel to fractures for all cases, are presented in Figure 3-18. The difference in AVO gradient is greater as the gas content in the fluid increases.

3.4 Discussion

We compare the P-wave AVO analysis over the synthetic seismograms around the bottom of the reservoir to the results obtained from the field data (see Figures 3-17 and 3-19). Even the absolute values are not the same, they exhibit similar trend near the bottom of the reservoir, where the AVO gradient is bigger for the lines perpendicular to fracture orientation than for the line that is parallel to it. This behavior is obtained using an oil to fill the fractures with a similar API to the one produced in the field.

Although we know that the main concentration of fractures is at the top of the reservoir ('O' member), the AVO anomaly is observed around the seismic event that has been interpreted as the bottom of the reservoir. However, because the 'O' member and the 'S' shale are very thin layers, and also the 'P' and 'R' sandstones conform an small one, some tuning effects can change the waveforms and the amplitude content (see Figure 3-20). The reflections associated with the 'O'member start in the seismic time section at the "depth" (time) of the 'O' member, but extend below it, being superposed with other events from 'P' and 'R' members (top and base). Likewise, the reflections from the top and the bottom of the 'O' member is combined together in the seismic section. Therefore, one of the major goals of the ray tracing modeling was to check if the response of the 'combination' of reflections from the different interfaces reproduce the observed pattern of AVO variations of the field data. The results obtained from the synthetic data confirm the same results obtained from the field data. We also did some ray calculations using a similar model, but with fractures only in the 'O' member, and they yield to fairly similar results. Some deconvolution could
help to improve the vertical resolution and identify individual reflections, but we decided not to do so to try to better preserve accurate amplitudes.

We showed, with field and synthetic data, that the P-wave AVO response depends on the orientation of the shot line with respect to the fractures. These results demonstrate that P-wave AVO may be useful to detect azimuthal anisotropy and potential fracture orientation.

No AVO gradient anomaly is obtained at the top of the reservoir from the synthetic data. This confirms the results obtained from the field data.
Figure 3-1: Well logs used to build an earth model at the cross point of the three lines.
Figure 3-2: P-wave and S-wave velocities obtained from a dipole sonic log in well MA-17.
Figure 3-3: Density and sonic logs obtained from a resistivity log in well MA-1, using Gardner’s and Faust’s approximations.
Figure 3-4: Sonic log in well TOR-2X.
Figure 3-5: The eight-layer model used to generate the synthetic seismograms.
Figure 3-6: Comparison of stacked synthetic seismogram to field data. The synthetic seismogram (middle) corresponds to the propagation parallel to fractures filled with medium oil. Some mismatches are observed between the first event and the top of the reservoir because the formations Gobernador, Burguita and Quevedo were averaged as a single layer.
Figure 3-7: Synthetic AVO gradient for a line perpendicular to fracture orientation. Fractures filled with gas.
Figure 3-8: Synthetic AVO gradient for a line parallel to fracture orientation. Fractures filled with gas.
Figure 3-9: Synthetic AVO gradient for a line perpendicular to fracture orientation. Fractures filled with light oil.
Figure 3-10: Synthetic AVO gradient for a line parallel to fracture orientation. Fractures filled with light oil.
Figure 3-11: Synthetic AVO gradient for a line perpendicular to fracture orientation. Fractures filled with medium oil.
Figure 3-12: Synthetic AVO gradient for a line parallel to fracture orientation. Fractures filled with medium oil.
Figure 3-13: Synthetic AVO gradient for a line perpendicular to fracture orientation. Fractures filled with heavy oil.
Figure 3-14: Synthetic AVO gradient for a line parallel to fracture orientation. Fractures filled with heavy oil.
Figure 3-15: Synthetic AVO gradient for a line perpendicular to fracture orientation. Fractures filled with water.
Figure 3-16: Synthetic AVO gradient for a line parallel to fracture orientation. Fractures filled with water.
Figure 3-17: P-waves amplitudes vs. offset for the bottom of the reservoir. Fractures filled with (a) gas, (b) light oil, (c) medium oil, (d) heavy oil, and (e) water.
Figure 3-18: Percentage differences in AVO gradients. Perpendicular - Parallel to fracture orientation
Figure 3-19: Comparison of P-waves amplitudes vs. offset at the bottom of the reservoir - Field Data
Figure 3-20: Tuning effect. The reflection from the top and the base of the layers (a) interfere, and result in a composed signal (b). Waveshapes and amplitudes differ from the original signal.
Chapter 4

Discussion and Conclusions

From the azimuthal dependent AVO analysis at the crosspoint of three 2D seismic lines, over a fractured reservoir, the following conclusions can be made:

1. In the presence of fracture induced azimuthal anisotropy, the P-wave AVO response may depend upon azimuth.

2. The results obtained using azimuthal-dependent P-wave AVO analysis are consistent with the results obtained in a previous study with the same field data using P-S converted waves.

3. The modeling results, using information from nearby wells, show the same trend as the field data.

4. Azimuthal P-wave AVO analysis can aid in detecting azimuthal anisotropy and fracture orientation.

One limitation in this work was that 2D data allowed the analysis of only one surface location (intersection point of the three seismic lines). There is a 3D survey over the same area, and in the future we will do azimuth-dependent AVO analysis using this dataset to confirm the results obtained in the present work.
Appendix A

AVO Theory

(From Castagna, 1993a)

The variation of reflection and transmission coefficients with the incident angle (and corresponding increasing offset) is referred to as offset-dependent-reflectivity and is the fundamental basis for amplitude-versus-offset (AVO) analysis.

In exploration geophysics, we rarely deal with simple isolated interfaces. Nevertheless, it is easier to understand offset-dependent reflectivity as the partitioning of energy at just an interface. We must be careful, however, to remember that most reflections we see are a superposition of events from a series of layers, and will have a more complex AVO behavior than will be shown here.

Consider two semi-infinite isotropic homogeneous elastic media in contact at a plane interface. Further, consider an incident compressional plane wave impinging upon this interface. The plane-wave assumption is valid at source-to-receiver distances, which are much longer than the wavelength of the incident wave and are generally acceptable for precritical reflection data at exploration depths and frequencies. As shown in Figure A-1, reflection at an interface involves energy partition from an incident P-wave to (1) a reflected P-wave, (2) a transmitted P-wave, (3) a reflected S-wave, and (4) a transmitted S-wave. The angles for
incident, reflected, and transmitted rays synchronous at the boundary are related according
to Snell’s law,

\[ p = \frac{\sin(\theta_1)}{V_{P1}} = \frac{\sin(\theta_2)}{V_{P2}} = \frac{\sin(\sigma_1)}{V_{S1}} = \frac{\sin(\sigma_2)}{V_{S2}} \]

(A.1)

where,

- \( V_{P1} \) = P-wave velocity in medium 1,
- \( V_{P2} \) = P-wave velocity in medium 2,
- \( V_{S1} \) = S-wave velocity in medium 1,
- \( V_{S2} \) = S-wave velocity in medium 2,
- \( \theta_1 \) = incident P-wave angle,
- \( \theta_2 \) = transmitted P-wave angle,
- \( \sigma_1 \) = reflected S-wave angle,
- \( \sigma_2 \) = transmitted S-wave angle,

and \( p \) is the ray parameter.

The P-wave reflection coefficient as a function of incident angle \( R_{pp}(\theta_1) \) is defined as the
ratio of the amplitude of the reflected P-wave to that of the incident P-wave. Similarly,
the P-wave transmission coefficient \( T_{pp}(\theta_1) \) is the ratio of the amplitude of the transmitted
P-wave to that of the incident P-wave.

At normal incidence, there are no converted S-waves and the P-wave reflection coefficient
\( R_p \) is given by

\[ R_p = \frac{I_{P1} - I_{P2}}{(I_{P1} + I_{P2})} \]

(A.2)

where,
$I_{P2} = \text{impedance of medium 2} = \rho_2 V_{P2}$,

$\rho_2 = \text{density of medium 2},$

$I_{P1} = \text{impedance of medium 1} = \rho_1 V_{P1},$

$\rho_1 = \text{density of medium 1}.$

However, the approximation for the reflection coefficient at non-normal incidence $R_{pp}(\theta)$ is much more complicated, and many equations have been presented in the literature (e.g., Aki and Richards, 1980; Shuey, 1985; Wright, 1986; Hilterman, 1989).

All too often, the tremendous potential of AVO analysis as a prospecting tool has not been realized partially because of the wide variety of complications, problems, and pitfalls involved in isolating and interpreting the offset-dependent-reflectivity (ODR). Table A.1 shows the myriad of factors affecting seismic amplitudes. For most AVO analysis techniques, all factors other than reflection coefficient variation with offset must be considered undesired noise (which should be attacked by processing or properly accounted for) or must be appropriately comprehended as a signal by the processing/analysis techniques.

---

A. Desired information (signal)

1. reflection coefficients versus angle of incidence

B. Potential information (considered noise for some methods; signal for others)

1. composite reflections from multiple interfaces

2. tuning caused by NMO convergence

3. mode conversions

C. Factors without offset dependence (noise)

1. random noise

2. instrumentation
3. source/receiver coupling

D. Factor with offset dependence (noise)

1. source/receiver directivity including ghosting and array response
2. emergence angles
3. coherent noise
4. spherical spreading
5. processing distortions, NMO errors and stretch
6. inelastic attenuation and anisotropy
7. transmission coefficients and scattering above target
8. structural complexity

---

Table A.1: Factor affecting seismic amplitudes

Successful AVO analysis requires a degree of integration beyond that which has been absolutely necessary for conventional seismic interpretation: well logging, rock physics, reservoir engineering, seismic data acquisition and processing, seismic modeling and inversion, conventional seismic interpretation, and a healthy dose of geological concepts and constraints must be incorporated in the analysis.

AVO inversion, as many other geophysical techniques, is not unique. However, we should not let nonuniqueness in AVO inversion blind us to the tremendous utility of AVO analysis in general. After all, let us not forget that generating structure maps from reflection times is an inherently nonunique process. Ultimately, geophysics is a science of anomalies, and anomalies are defined by our expectations. There is no rigorous solution and our expectations are inexact; nevertheless, AVO does have the ability to reduce risk and to ‘illuminate’ new prospects which were previously overlooked. It is important to point out that all this analysis over the AVO technique corresponds to isotropic media. However, AVO for anisotropic media has been recently studied by several different authors, and should be considered for any AVO analysis.
Figure A-1: Reflection and transmission at an interface between two infinite elastic half-spaces for an incident P-wave.
Appendix B

Processing Schemes Suggested for AVO Analysis

Some different schemes have been presented in the literature for data processing in order to preserve amplitude information for AVO analysis (from Castagna, 1993a):

| Ostrander (1984): | 1. spherical divergence correction |
|                 | 2. exponential gain correction     |
|                 | 3. minimum-phase spiking deconvolution |
|                 | 4. velocity analysis                |
|                 | 5. NMO correction                   |
|                 | 6. trace equalization                |
|                 | 7. horizontal trace summing         |

| Chiburis (1984): | 1. mild f-k multiple suppression |
|                 | 2. spherical divergence and NMO correction |
Long and Richgels (1985):

3. whole-trace equalization
4. flattening on a consistent reference event
5. horizontal trace summing
6. peak amplitudes picked interactively
7. smoothed least-squares curve fitting
8. despiking of outliers
9. curve refitting
10. results clipped and smoothed

1. spherical divergence, attenuation, and emergence angle corrections determined by regression analysis
2. variation in source strength, receiver sensitivity, overall offset amplitudes, consistently modelled and resolved
3. deviations rejected as noisy
4. inverse amplitude correction
5. spectral surface array attenuation correction
6. accurate NMO applied using very fine velocity sampling and geologically consistent velocity contouring
7. surface consistent deconvolution for land data
8. CDP residual statics carefully monitored for incorrect shifts due to AVO effects
9. minimal muting
10. common offset spatial filtering
11. linear or non-linear fit for AVO coefficients

Yu (1985):

1. apply exponential gain, suppress coherent noise, and remove gain
2. offset compensation
3. deconvolution
4. NMO correction
5. surface consistent correction
6. partial trace sum
7. bandpass
8. section dependent equalization

Todd (1986):

1. spatial averaging
2. gapped deconvolution
3. geometrical divergence correction
4. global velocity analysis and NMO correction
5. partial stacking
6. correction for average amplitude variation in time and offset
Appendix C

Modeling Parameters for the Fractured Layers

Modeling parameters for the fractured layers used for all cases (gas, light oil, medium oil, heavy oil and water) are presented below. The fractured layers are:

1. The member 'O' of the Escandalosa formation (reservoir).

2. An average of the 'P' and 'R' sandstones of the Escandalosa formation (reservoir).

We use these parameters, to calculate the elastic constants for the fractured layers, using Hudson's (1981) theory.

Case 1: Gas filling the fractures

<table>
<thead>
<tr>
<th></th>
<th>$V_P$</th>
<th>$V_S$</th>
<th>$\rho$</th>
<th>Bulk modulus of fluid filling the fractures</th>
<th>Shear modulus of fluid filling the fractures</th>
<th>Fracture density</th>
<th>Aspect ratio</th>
</tr>
</thead>
<tbody>
<tr>
<td>1</td>
<td>3513 m/s</td>
<td>1887 m/s</td>
<td>2.5 g/cm$^3$</td>
<td>0.0</td>
<td>0.0</td>
<td>0.1</td>
<td>0.001</td>
</tr>
<tr>
<td>2</td>
<td>4994 m/s</td>
<td>2784 m/s</td>
<td>2.7 g/cm$^3$</td>
<td>0.0</td>
<td>0.0</td>
<td>0.05</td>
<td>0.001</td>
</tr>
</tbody>
</table>
Case 2: Light Oil filling the fractures

<table>
<thead>
<tr>
<th></th>
<th>$V_p$</th>
<th>$V_s$</th>
<th>$\rho$</th>
<th>Bulk modulus of fluid filling the fractures</th>
<th>Shear modulus of fluid filling the fractures</th>
<th>Fracture density</th>
<th>Aspect ratio</th>
</tr>
</thead>
<tbody>
<tr>
<td>1</td>
<td>3513 m/s</td>
<td>1887 m/s</td>
<td>2.5 g/cm$^3$</td>
<td>0.2</td>
<td>0.0</td>
<td>0.1</td>
<td>0.001</td>
</tr>
<tr>
<td>2</td>
<td>4994 m/s</td>
<td>2784 m/s</td>
<td>2.7 g/cm$^3$</td>
<td>0.2</td>
<td>0.0</td>
<td>0.05</td>
<td>0.001</td>
</tr>
</tbody>
</table>

Case 3: Medium Oil filling the fractures

<table>
<thead>
<tr>
<th></th>
<th>$V_p$</th>
<th>$V_s$</th>
<th>$\rho$</th>
<th>Bulk modulus of fluid filling the fractures</th>
<th>Shear modulus of fluid filling the fractures</th>
<th>Fracture density</th>
<th>Aspect ratio</th>
</tr>
</thead>
<tbody>
<tr>
<td>1</td>
<td>3513 m/s</td>
<td>1887 m/s</td>
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<td>0.0</td>
<td>0.1</td>
<td>0.001</td>
</tr>
<tr>
<td>2</td>
<td>4994 m/s</td>
<td>2784 m/s</td>
<td>2.7 g/cm$^3$</td>
<td>1.0</td>
<td>0.0</td>
<td>0.05</td>
<td>0.001</td>
</tr>
</tbody>
</table>

Case 4: Heavy Oil filling the fractures

<table>
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<th></th>
<th>$V_p$</th>
<th>$V_s$</th>
<th>$\rho$</th>
<th>Bulk modulus of fluid filling the fractures</th>
<th>Shear modulus of fluid filling the fractures</th>
<th>Fracture density</th>
<th>Aspect ratio</th>
</tr>
</thead>
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<td>0.001</td>
</tr>
<tr>
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<td>2784 m/s</td>
<td>2.7 g/cm$^3$</td>
<td>1.8</td>
<td>0.0</td>
<td>0.05</td>
<td>0.001</td>
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</tbody>
</table>

Case 5: Water filling the fractures

<table>
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<tr>
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<th>$V_p$</th>
<th>$V_s$</th>
<th>$\rho$</th>
<th>Bulk modulus of fluid filling the fractures</th>
<th>Shear modulus of fluid filling the fractures</th>
<th>Fracture density</th>
<th>Aspect ratio</th>
</tr>
</thead>
<tbody>
<tr>
<td>1</td>
<td>3513 m/s</td>
<td>1887 m/s</td>
<td>2.5 g/cm$^3$</td>
<td>2.0</td>
<td>0.0</td>
<td>0.1</td>
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</tr>
<tr>
<td>2</td>
<td>4994 m/s</td>
<td>2784 m/s</td>
<td>2.7 g/cm$^3$</td>
<td>2.0</td>
<td>0.0</td>
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<td>0.001</td>
</tr>
</tbody>
</table>


