

VARIATIONAL DERIVATION OF MODAL-NODAL  
FINITE DIFFERENCE EQUATIONS IN  
SPATIAL REACTOR PHYSICS

by

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ABSTRACT

A class of consistent coarse mesh modal-nodal approximation methods is presented for the solution of the spatial neutron flux in multigroup diffusion theory. The methods are consistent in that they are systematically derived as an extension of the finite element method by utilizing general modal-nodal variational techniques. Detailed subassembly solutions, found by imposing zero current boundary conditions over the surface of each subassembly, are modified by piecewise continuous Hermite polynomials of the finite element method and used directly in trial function forms. Methods using both linear and cubic Hermite basis functions are presented and discussed.

The proposed methods differ substantially from the finite element methods in which homogeneous nuclear constants, homogenized by flux weighting with detailed subassembly solutions, are used. However, both schemes become equivalent when the subassemblies themselves are homogeneous.

One-dimensional, two-group numerical calculations using representative PWR nuclear material constants and 18-cm subassemblies were performed using entire subassemblies as coarse mesh regions. The results indicate that the proposed methods can yield comparable if not superior criticality measurements, comparable regional power levels, and extremely accurate subassembly fine flux structure with little increase of computational effort in comparison with existing coarse mesh methods.

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. . . Until one is committed there is hesitancy, the chance to draw back, always ineffectiveness. Concerning all acts of initiative (and creation), there is one elementary truth, the ignorance of which kills countless ideas and splendid plans: that the moment one definitely commits oneself, then Providence moves too. All sorts of things occur to help one that would never otherwise have occurred. A whole stream of events issues from the decision, raising in one's favor all manner of unforeseen incidents and meetings and material assistance, which no man could have dreamt would have come his way. I have learned a deep respect for one of Goethe's couplets:

Whatever you can do, or dream you can, begin it.  
Boldness has genius, power, and magic in it.

— W.H. Murray



Great importance attaches to the material comforts of life, and equanimity, unconcern, security are all sacrificed to them. The American lives even more for his goals, for the future, than the European. Life for him is always becoming, never being.

– Albert Einstein 1921

"When someone is seeking," said Siddhartha, "it happens quite easily that he only sees the thing that he is seeking; that he is unable to absorb anything, because he is only thinking of the thing he is seeking, because he has a goal, because he is obsessed with his goal. Seeking means: to have a goal; but finding means: to be free, to be receptive, to have no goal. You, O worthy one, are perhaps indeed a seeker, for in striving toward your goal, you do not see many things that are under your nose."

– Hermann Hesse

## TABLE OF CONTENTS

Abstract	2
Acknowledgments	3
List of Figures	7
List of Tables	10
 Chapter 1. Introduction	 11
1.1 Preface	11
1.2 The Time-Independent, Multigroup Diffusion Theory Equations	13
1.3 Solution Methods	15
1.3.1 Nodal Methods	16
1.3.2 Modal Methods	19
1.3.3 Modal-Nodal Methods	25
 Chapter 2. Variational Derivation of Finite Difference Approximations in Time-Independent Multigroup Diffusion Theory	 27
2.1 Calculus of Variations Applied to Diffusion Theory	28
2.2 Discontinuous Trial Functions	30
2.3 The Finite Element Approximation Methods	34
2.3.1 The Conventional Finite Difference Equations	35
2.3.2 Multichannel Polynomial Synthesis	37
2.3.3 The Linear Basis Function Approximation	46
2.3.4 The Cubic Hermite Basis Function Approximation	47

Chapter 3. Development of a Consistent Coarse Mesh Approximation Method	52
3.1 Formulation	52
3.2 The Proposed Linear Basis Function Approximations	54
3.3 The Proposed Cubic Hermite Basis Function Approximation	60
Chapter 4. Numerical Solution Techniques	65
4.1 Solution Methods and Matrix Properties	65
4.2 Calculational and Programming Techniques	77
Chapter 5. Numerical Results	84
5.1 Nuclear Constants and Subassembly Geometry	84
5.2 Subassembly Detailed Solutions and Homogenized Nuclear Constants	86
5.3 Case Studies and Results	95
5.3.1 Case 1	98
5.3.2 Case 2	113
5.3.3 Case 3	118
5.3.4 Case 4	124
Chapter 6. Conclusions and Recommendations	130
6.1 Characteristics of the Proposed Approximation Methods	130
6.2 Applicability and Limitations	132
6.3 Recommendations for Future Work	133
References	134
Biographical Note	138

Appendix A. Table of Symbols	140
Appendix B. Difference Equation Coefficients Resulting from Use of the Finite Element Approximation Methods	143
B.1 Coefficients of the Conventional Finite Difference Equations	143
B.2 Coefficients of the Linear Finite Element Method Equations	144
B.3 Coefficients of the Cubic Hermite Finite Element Method Equations	145
Appendix C. Difference Equation Coefficients Resulting from Use of the Proposed Approximation Methods	147
C.1 Coefficient-Integrands of the Proposed Approximation Method Equations Using Linear Basis Functions	148
C.2 Coefficient-Integrands of the Proposed Approximation Method Equations Using Cubic Hermite Basis Functions	150
Appendix D. Description of the Computer Programs	153
D.1 Description of Program REF2G	153
D.2 Description of Program LINEAR	158
D.3 Description of Program CUBIC	163
D.4 Description of Program ANALYZE	164
Appendix E. Sample Input and Output Data Blocks for Programs REF2G, LINEAR, CUBIC, and ANALYZE (Included in only the first six copies)	168
E.1 REF2G	
E.2 LINEAR	
E.3 CUBIC	
E.4 ANALYZE	
Appendix F. Source Listings of the Programs (Included in only the first six copies)	183
F.1 REF2G	184
F.2 LINEAR	242
F.3 CUBIC	287
F.4 ANALYZE	350

## LIST OF FIGURES

<u>No.</u>		
1.1	Illustration of One-Dimensional Multichannel Synthesis	23
1.2	Illustration of One-Dimensional Overlapping Multichannel Synthesis	23
2.1	Conventional Nodal Finite Difference Approximation Trial Function Forms	36
2.2	Matrix Form of the Conventional Finite Difference Equations	38
2.3	Basis Functions of Equations 2.23 for $N = 0, 1, 2,$ and $3$	43
2.4	Cubic B Spline $\Omega_k^B(z)$	45
2.5	Cubic Hermite Basis Functions $\Omega_k^{H_1}(z)$ and $\Omega_k^{H_2}(z)$	45
2.6	Matrix Form of the Linear Finite Element Method Approximation	48
2.7	Matrix Form of the Cubic Hermite Finite Element Method Approximation	51
4.1	Solution of $\mathbf{A}\underline{F} = \frac{1}{\lambda}\mathbf{B}\underline{F}$ Using the Fission Source Power Iteration Method Without Fission Source Renormalization	66
4.2	Subassembly Notations and Detailed Solutions	79
5.1	Subassembly Configuration Geometries	87
5.2	Mesh Geometry in Half a Subassembly	88
5.3	Subassembly Detailed Flux Solutions for the One-Group Case	89
5.4	Subassembly Detailed Flux Solutions for the Two-Group Case	90
5.5	Subassembly Detailed Adjoint Flux Solutions for the Two-Group Case	91
5.6	Geometry of the Four Case Studies Composed of Types of Subassemblies	96

<u>No.</u>		
5.7	Case 1: One-Group Results Using Linear Basis Function Approximations and 18-cm Coarse Mesh Regions	99
5.8	Case 1: One-Group Results Using Cubic Hermite Basis Function Approximations and 18-cm Coarse Mesh Regions	100
5.9	Case 1: Two-Group Fast Results Using Linear Basis Function Approximations and 18-cm Coarse Mesh Regions	101
5.10	Case 1: Two-Group Thermal Results Using Linear Basis Function Approximations and 18-cm Coarse Mesh Regions	102
5.11	Case 1: Two-Group Fast Results Using Cubic Hermite Basis Function Approximations and 18-cm Coarse Mesh Regions	103
5.12	Case 1: Two-Group Thermal Results Using Cubic Hermite Basis Function Approximations and 18-cm Coarse Mesh Regions	104
5.13	Case 1: Two-Group Fast Results Using Linear Basis Function Approximations and 9-cm Coarse Mesh Regions	105
5.14	Case 1: Two-Group Thermal Results Using Linear Basis Function Approximations and 9-cm Coarse Mesh Regions	106
5.15	Case 1: Two-Group Fast Results Using Cubic Hermite Basis Function Approximations and 9-cm Coarse Mesh Regions	107
5.16	Case 1: Two-Group Thermal Results Using Cubic Hermite Basis Function Approximations and 9-cm Coarse Mesh Regions	108
5.17	Case 1: Two-Group Fast Results Using Cubic Hermite Finite Element Approximations and 18-cm Coarse Mesh Regions	111
5.18	Case 1: Two-Group Thermal Results Using Cubic Hermite Finite Element Approximations and 18-cm Coarse Mesh Regions	112
5.19	Case 2: Two-Group Fast Results Using Linear Basis Function Approximations and 18-cm Coarse Mesh Regions	114

<u>No.</u>		
5.20	Case 2: Two-Group Thermal Results Using Linear Basis Function Approximations and 18-cm Coarse Mesh Regions	115
5.21	Case 2: Two-Group Fast Results Using Cubic Hermite Basis Function Approximations and 18-cm Coarse Mesh Regions	116
5.22	Case 2: Two-Group Thermal Results Using Cubic Hermite Basis Function Approximations and 18-cm Coarse Mesh Regions	117
5.23	Case 3: Two-Group Fast Results Using Linear Basis Function Approximations and 18-cm Coarse Mesh Regions	119
5.24	Case 3: Two-Group Thermal Results Using Linear Basis Function Approximations and 18-cm Coarse Mesh Regions	120
5.25	Case 3: Two-Group Fast Results Using Cubic Hermite Basis Function Approximations and 18-cm Coarse Mesh Regions	121
5.26	Case 3: Two-Group Thermal Results Using Cubic Hermite Basis Function Approximations and 18-cm Coarse Mesh Regions	122
5.27	Case 4: Two-Group Fast Results Using Linear Basis Function Approximations and 18-cm Coarse Mesh Regions	125
5.28	Case 4: Two-Group Thermal Results Using Linear Basis Function Approximations and 18-cm Coarse Mesh Regions	126
5.29	Case 4: Two-Group Fast Results Using Cubic Hermite Basis Function Approximations and 18-cm Coarse Mesh Regions	127
5.30	Case 4: Two-Group Thermal Results Using Cubic Hermite Basis Function Approximations and 18-cm Coarse Mesh Regions	128
F.1	Structure of Program REF2G	185
F.2	Structure of Program LINEAR	243
F.3	Structure of Program CUBIC	288
F.4	Structure of Program ANALYZE	351

## LIST OF TABLES

<u>No.</u>		
3.1	Matrix Order N of the Proposed Linear Basis Function Approximations as a Function of the Imposed Boundary Conditions	59
5.1	Representative Two-Group, 18-cm, PWR Subassembly Regional Nuclear Constants	85
5.2	Representative One-Group, 18-cm, PWR Subassembly Regional Nuclear Constants	85
5.3	Homogenized Subassembly One-Group Nuclear Constants	92
5.4	Homogenized Subassembly Two-Group Nuclear Constants	92
5.5	Test Results Using Three Consecutive Type A Subassemblies	94
5.6	Results of Case 1	109
5.7	Two-Group Results of Case 2	113
5.8	Results of Case 3	123
5.9	Results of Case 4	129
D.1	Sample Storage Requirements and Execution Times of the Programs for Two-Group Results	154



## Chapter 1

### INTRODUCTION

#### 1.1 Preface

The large variety of approximation methods and techniques used in computational reactor analysis and simulation has caused the area of numerical reactor physics to become one of the most exciting areas in applied nuclear reactor physics today. The application of numerical analysis is most important in two phases of reactor design; feasibility studies and safety analysis. The primary consideration of the reactor physicist has been and must continue to be the safety of the reactor during and after any foreseeable nuclear accident. A realistic safety analysis can be obtained only if all the physical processes occurring within the reactor can be adequately described and related. Since all of these processes can be shown to be dependent upon the neutron density distribution throughout the reactor core, a detailed solution of the spatial neutron flux is vital.<sup>1</sup>

The dynamic characteristics of a reactor strongly depend upon the spatial approximation and solution of the neutron flux. Approximation methods utilizing gross averaging of the flux near localized strong absorption and production regions, such as cruciform control rods or small water channels, can lead to inaccurate results. Large errors may result from the use of such methods in spatial kinetics problems such as depletion and xenon oscillation calculations. Much attention has therefore been focused upon approximation methods

which can obtain detailed spatial neutron flux distributions within large reactor cores.

The Boltzmann neutron transport equation<sup>2</sup> is considered to be a sufficiently detailed description of the physical processes occurring within a nuclear reactor, and naturally is most difficult to solve. The P-1 and diffusion theory approximations<sup>3</sup> greatly simplify the transport equation into more tractable equations which have been found to approximate adequately the flux distributions for most large-core reactors such as PWR, BWR, and LMFBR core geometries. The advent of high speed digital computers has enabled widespread use of diffusion theory because of its simple mathematical form and straightforward numerical solution techniques inherent with its use.

The treatment of the spatial approximation in diffusion theory is the primary concern of this report. There is in existence an increasingly abundant variety of such approximation methods currently in use. Fine mesh methods,<sup>4</sup> for example, can yield very accurate results through the use of extremely large numbers of unknowns. However, such methods may well exceed the storage capacity of present day computers, as well as being exceedingly costly. Coarse mesh methods and particularly synthesis techniques,<sup>5</sup> on the other hand, have recently become attractive as the number of unknowns can be drastically reduced, although the accuracy of many of these methods is in doubt.

The purpose of this report is twofold: first, to present the general development of variational approximation methods used to derive

difference approximations to the neutron diffusion equation; and second, to extend this development in order to develop systematically a class of consistent coarse mesh approximation methods which can approximate accurately the detailed spatial neutron flux and can also be easily incorporated into present day computer codes. As this report will deal only with the spatial approximation, the inclusion of time dependence will be set aside for future study.

## 1.2 The Time-Independent, Multigroup Diffusion Theory Equations

The energy discretized multigroup P-1 approximation to the Boltzmann neutron transport equation excluding time dependence can be written in standard group notation for each energy group  $g$  as follows:<sup>3</sup>

$$\underline{j}_g(r) + D_g(r) \underline{\nabla} \phi_g(r) = 0 \quad (1.1a)$$

$$\underline{\nabla} \cdot \underline{j}_g(r) + \Sigma_g(r) \phi_g(r) - \sum_{\substack{g'=1 \\ g' \neq g}}^G \Sigma_{gg'}(r) \phi_{g'}(r) = \frac{1}{\lambda} \chi_g \sum_{g'=1}^G \nu \Sigma_{fg'}(r) \phi_{g'}(r) \quad (1.1b)$$

where the group index  $g$  runs from the highest energy group, 1, to the lowest energy group,  $G$ . The symbols and notation used throughout this report are summarized in Appendix A. Equations 1.1 are the standard P-1 equations which relate the vector neutron current  $\underline{j}_g(r)$  for each energy group  $g$  with the scalar neutron flux  $\phi_g(r)$ . The current may be eliminated via Fick's law, Eq. 1.1a, in order to obtain the multigroup diffusion equation:

$$-\underline{\nabla} \cdot D_g(r) \underline{\nabla} \phi_g(r) + \Sigma_g(r) \phi_g(r) - \sum_{\substack{g'=1 \\ g' \neq G}}^G \Sigma_{gg'}(r) \phi_{g'}(r) = \frac{1}{\lambda} \chi_g \sum_{g'=1}^G \nu \Sigma_{fg'}(r) \phi_{g'}(r) \quad (1.2)$$

Equation 1.2 can be written in operator matrix notation as

$$-\underline{\nabla} \cdot \mathbb{D}(\mathbf{r}) \underline{\nabla} \Phi(\mathbf{r}) + [\mathbb{M}(\mathbf{r}) - \mathbb{T}(\mathbf{r})] \Phi(\mathbf{r}) = \frac{1}{\lambda} \mathbb{B}(\mathbf{r}) \Phi(\mathbf{r}) \quad (1.3)$$

where  $\mathbb{D}$ ,  $\mathbb{M}$ ,  $\mathbb{T}$ , and  $\mathbb{B}$  are  $G \times G$  group matrices defined by

$$\mathbb{D}(\mathbf{r}) = \text{Diag}[D_1(\mathbf{r}) \dots D_g(\mathbf{r}) \dots D_G(\mathbf{r})] \quad (1.4a)$$

$$\mathbb{M}(\mathbf{r}) = \text{Diag}[\Sigma_1(\mathbf{r}) \dots \Sigma_g(\mathbf{r}) \dots \Sigma_G(\mathbf{r})] \quad (1.4b)$$

$$\mathbb{T}(\mathbf{r}) = \begin{bmatrix} 0 & -\Sigma_{12}(\mathbf{r}) & \dots & -\Sigma_{1G}(\mathbf{r}) \\ -\Sigma_{21}(\mathbf{r}) & 0 & \dots & -\Sigma_{2G}(\mathbf{r}) \\ \vdots & \vdots & \ddots & \vdots \\ \vdots & \vdots & \ddots & \vdots \\ -\Sigma_{G1}(\mathbf{r}) & -\Sigma_{G2}(\mathbf{r}) & \dots & 0 \end{bmatrix} \quad (1.4c)$$

$$\mathbb{B}(\mathbf{r}) = \begin{pmatrix} \chi_1 \\ \vdots \\ \vdots \\ \chi_G \end{pmatrix} \begin{bmatrix} \nu \Sigma_{f1}(\mathbf{r}) & \dots & \nu \Sigma_{fG}(\mathbf{r}) \end{bmatrix} \quad (1.4d)$$

and  $\Phi(\mathbf{r})$  is the group flux vector

$$\Phi(\mathbf{r}) = \text{Col}[\phi_1(\mathbf{r}) \dots \phi_G(\mathbf{r})] \quad (1.4e)$$

In problems where no upscattering is present,  $\Sigma_{gg'}(\mathbf{r}) = 0$  for  $g < g'$ , and  $\mathbb{T}$  becomes  $G \times G$  lower triangular.

It is also convenient to define the group current vector  $\underline{\mathbf{J}}(\mathbf{r})$

$$\underline{\mathbf{J}}(\mathbf{r}) = \text{Col}[\underline{j}_1(\mathbf{r}) \dots \underline{j}_G(\mathbf{r})] \quad (1.4f)$$

and the  $G \times G$  group absorption, scattering and production matrix  $\underline{\mathbf{A}}(\mathbf{r})$

$$\underline{\mathbf{A}}(\mathbf{r}) = \mathbb{M}(\mathbf{r}) - \mathbb{T}(\mathbf{r}) - \frac{1}{\lambda} \mathbb{B}(\mathbf{r}) \quad (1.4g)$$

Equations 1.1 and 1.2 may then be written simply as

$$\underline{J}(\mathbf{r}) + \underline{D}(\mathbf{r}) \underline{\nabla} \Phi(\mathbf{r}) = 0 \quad (1.5a)$$

$$\underline{\nabla} \cdot \underline{J}(\mathbf{r}) + \underline{\Lambda}(\mathbf{r}) \Phi(\mathbf{r}) = 0 \quad (1.5b)$$

and

$$-\underline{\nabla} \cdot \underline{D}(\mathbf{r}) \underline{\nabla} \Phi(\mathbf{r}) + \underline{\Lambda}(\mathbf{r}) \Phi(\mathbf{r}) = 0 \quad (1.6)$$

respectively. These forms of the group diffusion equations will be used throughout this report. The boundary conditions on  $\Phi(\mathbf{r})$  are of the homogeneous Neumann or Dirichlet type,<sup>6</sup> while the normal component of the current  $\underline{J}(\mathbf{r})$  is required to be continuous across all internal interfaces.

### 1.3 Solution Methods

All of the solution methods which can be employed in order to obtain approximate solutions to the time-dependent, multigroup diffusion equations may be conveniently classified as belonging in the area of either nodal analysis or modal analysis, or a combination of the two: modal-nodal analysis. The principal concept in each of these analyses is that the neutron flux, a continuous function of many variables, may be approximated as a set of unknown coefficients and/or functions of possibly fewer variables. The ultimate goals of such approximation methods are to produce easily solvable coupled equations which relate the unknowns to each approximation and yield results of acceptable accuracy at a low cost. Various commonly used methods and their drawbacks are discussed below.

### 1.3.1 Nodal Methods

Nodal methods involve the local approximation of an average flux at points called nodes, where each node represents a distinct region within the reactor in which the average flux is defined. An ordered set of nodes connected by a grid of mesh lines is then used to approximate the spatial flux behavior. The accuracy of such methods is generally governed by the internodal coupling or neutron current approximation inherent in each method.

#### A. Conventional Finite Difference Equations

The common finite difference equations used in diffusion theory can be derived using Taylor series expansion, variational techniques, or box integration methods about each spatial node.<sup>7</sup> The second-order diffusion term at each node is replaced by three-point difference equations relating consecutive nodes in each spatial direction. The resulting band-structured matrix equations exhibit many advantageous mathematical properties and can be solved with the use of simple solution algorithms.

The attractiveness of these difference equations is further enhanced by the fact that, for properly posed problems (including proper boundary conditions), the approximation can be shown to converge to the solution of the differential equation as the mesh size approaches zero. Also, the accuracy of the approximation can be shown to be in general of order  $\theta(h)$ ,<sup>8</sup> thus error estimates for the approximation are available. It is for these reasons that these equations are frequently invoked as "exact" solutions to diffusion equation problems. The main disadvantage, however, is that, as the

number of nodes increases, the amount of labor and cost involved in order to obtain an accurate solution increases geometrically. A point of diminishing returns is then quickly reached where further accuracy is prohibitively expensive. Another disadvantage is that any known physical insight or *a priori* detailed flux behavior cannot be used with this approximation.

A formal derivation of the conventional difference equations is given in section 2.3 of Chapter 2.

### B. Gross Coupling Models

In gross coupling or coarse mesh nodal techniques an attempt is made to decrease drastically the number of nodes needed for solution without significantly decreasing solution accuracy. Many such methods have been proposed by postulating various forms of neutronic coupling or communication interaction between nodes.

#### 1. Phenomenological Model<sup>9,10,11</sup>

From a physical viewpoint, the reactor can be divided into several distinct regions, each represented by a node located somewhere in that region. Equations of balance relating state variables of interest (average neutron flux, regional power, etc.) can then be written for each region and between region nodes. Internodal coupling is governed by a set of coefficients, say  $p_{ij}$ , which may account for the number of neutrons born in region  $i$  which appear in region  $j$ . A set of algebraic equations can then be written which describe the coupled core dynamics of the nodal interactions.

The principal drawback of such methods lies in the definition of the interaction parameters  $p_{ij}$ . Although the describing equations of the phenomenological model can be directly formulated from diffusion theory,<sup>12</sup> the method of calculating the coefficients  $p_{ij}$  remains unclear. However, the physical simplicity of this model has made it very appealing in coupled kinetics methods development. Much of the work in this field is based on deriving approximations which reduce to this simple conceptual model.

## 2. Effective D/L Coupling<sup>13</sup>

These methods are very similar to finite difference approximations in that the structural forms of the resulting difference equations are identical. In order to compensate for the use of large internodal mesh spacing, the reactor constants, and the diffusion coefficients in particular, may be altered so that they correspond in an average sense to those obtained from a fine mesh calculation.<sup>14</sup> In this way it is hoped that the gross internodal coupling will be sufficiently improved to compensate for the large mesh spacings.

It has been shown that such methods can indeed improve internodal coupling for large mesh regions; however, the results are generally not satisfactory since the coupling constants are dependent in an unpredictable way on changes in the properties of the nodes.

## 3. Fission Source Coupling<sup>15</sup>

The assumption that the reactor flux can be separated into partial region fluxes due to nodal fission sources permits a consistent derivation of nodal coupled kinetics equations from multigroup



diffusion theory. Fission modes can be found from detailed flux solutions which are then used to account for internodal coupling. This method gives reasonably accurate results for fast and thermal reactor transients, although the number of nodes necessary to achieve an accurate solution must increase as the form of the spatial flux becomes more detailed.

#### 4. Multichannel Coupling<sup>16</sup>

By partitioning the reactor into regions called channels and allowing only adjacent channel-to-channel interactions, coupling coefficients  $p_{ij}$  can be found which represent the net leakage of neutrons from channel  $i$  into channel  $j$  in terms of the corresponding averaged channel fluxes. The coupling coefficients can be calculated using diffusion theory or variational techniques which yield the diffusion equations as stationary conditions. This model is appealing in that it can be shown to reduce to the conventional difference equations when a regular grid of small channel regions is used.

The above examples of gross coupling models are generally unsatisfactory because they require the use of average fluxes defined within large regions of the reactor. More acceptable results are obtained by utilizing known or *a priori* detailed spatial flux shapes in the regions in the approximation method.

#### 1.3.2 Modal Methods<sup>17</sup>

Modal methods imply an extensive rather than local approximation to the spatial neutron flux. In general the flux is represented by a combination of known functions defined over the regions of interest

with unknown functions as mixing coefficients. Depending upon the approximation employed, relationships among these coefficients can be derived which are hopefully simpler to solve than the original equation.

#### A. Helmholtz Modes<sup>18</sup>

The diffusion equation for a completely homogeneous reactor formally has an infinite solution set of eigenvalues and corresponding orthogonal eigenfunctions, called Helmholtz modes, which satisfy the homogeneous boundary conditions. For the general case of a heterogeneous reactor, the spatially dependent flux can be approximated as a linear combination of these modes. The major difficulty with this approach is that a large number of modes is required in order to approximate the solution flux, and thus the appeal for this simplistic modal approach is quickly lost.

#### B. Lambda<sup>19</sup> and Omega Modes<sup>20</sup>

Although included in the class of modal approximations, these methods require the use of known spatial solutions for time-dependent analysis. Lambda modes belong to the set of detailed flux solutions of the time-independent diffusion equations which correspond to different lambda eigenvalues.

A set of detailed flux solutions can also be found from the time-dependent diffusion equations by allowing the time-dependent flux to be separable and given in the form  $e^{\omega t}$ . The solutions of the resulting equations, called  $\omega$  modes, correspond to different omega eigenvalues.

Both of these methods have successfully been used in the transient analysis of coupled nodal kinetics.

### C. Synthesis Methods<sup>5, 21</sup>

The use of synthesis techniques for the derivation of modal approximations is the most exciting and fastest growing area of reactor analysis methods development. This can be attributed to the fact that all diffusion theory approximation schemes, both modal and nodal, and those including time dependence, can be ultimately derived from one single variational principle. Each approximation scheme is therefore dependent solely upon the form of the trial functions used to represent the flux, current, and weighting functions (or adjoint functions) in the synthesis procedure. The outstanding advantage of the synthesis method is that knowledge of *a priori* detailed flux shapes or other physical insights can be incorporated directly into the approximation method.

#### 1. Multichannel Synthesis<sup>22</sup>

This method may be viewed as a modal extension of the multichannel gross coupling method. Assuming the flux to be separable in its variables  $(x, y, z)$ , the number of unknowns can be reduced by specifying detailed flux shapes in any dimension. A common example assumes that in each channel,  $k$ , of the reactor the flux trial function,  $U_k(x, y, z)$ , can be expressed as the product of a known transverse flux,  $\psi_k(x, y)$ , with an unknown spatially dependent axial flux,  $\rho_k(z)$ , as:

$$U_k(x, y, z) = \rho_k(z) \psi_k(x, y) \quad (1.7)$$

The specification of the flux in two dimensions reduces the problem to an approximation involving only one dimension. If, however, the flux is approximated by a full spatial solution times an unknown constant,

$$U_k(x, y, z) = F_k \psi_k(x, y, z) \quad (1.8)$$

the method reduces to an approximation similar to the multichannel nodal method. Figure 1.1 illustrates the resulting flux shape characteristics of such a method for the one-dimensional case.

The major disadvantage of these multichannel synthesis methods lies in the fact that in general the flux is discontinuous at channel interfaces.<sup>23,24</sup> Therefore the adjacent channel coupling currents, which then must be continuous across these interfaces,<sup>†</sup> are defined in terms of averaged channel fluxes. Although these methods can produce detailed flux distributions in each channel, their accuracy appears to be not much better than nodal multichannel methods because of the averaged gross neutronic coupling requirements inherent in these methods.<sup>25</sup>

## 2. Overlapping Multichannel Synthesis<sup>26,27</sup>

The interchannel neutronic coupling can be improved by requiring that the flux trial functions be continuous across channel interfaces. This can be accomplished by modulating the known expansion functions,  $\psi_k$ , by piecewise continuous normalized polynomial functions,  $p_k$ , which are nonzero only within coupled channels

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<sup>†</sup>Variational techniques used with diffusion theory in general do not allow the flux and current to be simultaneously discontinuous. Further clarification is given in section 2.2 of Chapter 2.

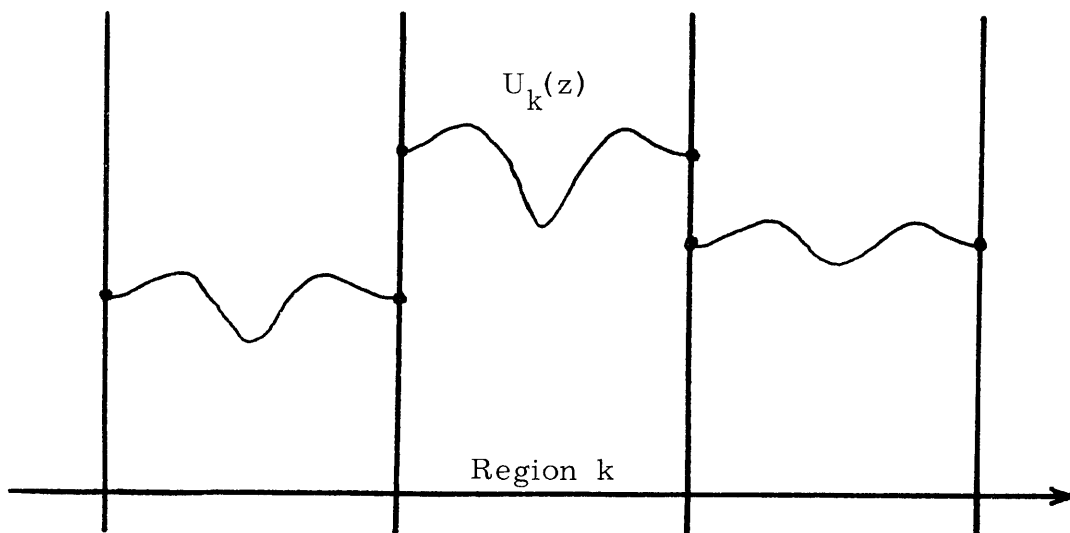


Figure 1.1. Illustration of One-Dimensional Multichannel Synthesis

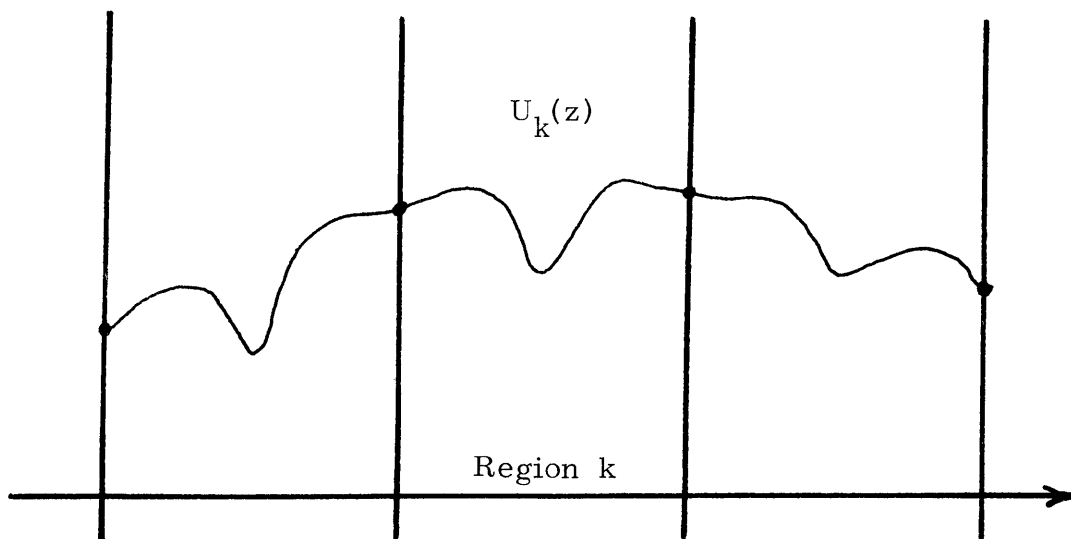


Figure 1.2. Illustration of One-Dimensional Overlapping Multichannel Synthesis

of interest, providing the expansion functions are continuous over all channels for which the corresponding polynomial functions are nonzero. Such polynomials are required to be normalized to unity at the coupling interface and zero along the external boundary of the channels in order to preserve flux trial function continuity.

In one dimension represented by the continuous variable  $z$  and  $K$  mesh regions bounded by the nodes  $z_k$  where  $k = 1$  to  $K+1$ , for example, the simple linear functions

$$p_k(z) = \begin{cases} \frac{z - z_{k-1}}{z_k - z_{k-1}} & z_{k-1} \leq z \leq z_k \\ \frac{z_{k+1} - z}{z_{k+1} - z_k} & z_k \leq z \leq z_{k+1} \\ 0 & \text{otherwise} \end{cases} \quad (1.9)$$

satisfy these conditions. The flux can then be approximated as

$$U(z) = \sum_{k=1}^K F_k p_k(z) \psi_k(z) \quad (1.10)$$

where the set of  $F_k$ 's are the unknowns of the method. The resulting flux shape characteristics of this approximation are illustrated in Figure 1.2.

Approximations based on this synthesis method are dependent upon the class of overlapping polynomial functions used as well as the form of the current trial functions employed. The form of the current is extremely important in that it specifies the coupling interaction between regions and in this sense governs the usefulness and accuracy

of the approximation. Work performed with this method to date has used current trial functions of a form similar to those of the flux trial function. Although the results of these investigations have been encouraging, such methods do not reduce to more simple known approximation methods. In addition, the band-structured matrix equations which arise from the use of such methods do not exhibit mathematical properties desired of such approximation schemes and may be difficult and costly to solve.

### 1.3.3 Modal-Nodal Methods

Approximation methods have also been developed in which the flux has a known extensive definition, or shape, and the unknowns are local flux values averaged in accordance with their corresponding extensive definition. Such modal-nodal methods retain all of the advantages of modal methods while generally reducing the number of unknowns and producing matrix equations which have desirable mathematical properties for numerical approximation and solution.

The finite element method is the best example of a modal-nodal approximation. Greater accuracy than that of conventional difference techniques can be obtained by allowing the flux in each region of interest to be represented as a polynomial which is continuous at region interfaces. The forms of the flux approximations and the resulting difference equations which arise from the use of the finite element method are described in detail in section 2.3 of Chapter 2.

The purpose of this report is to present an original and consistent class of modal-nodal coarse mesh approximation methods which retain given or known detailed flux structure within the regions of interest, while providing detailed neutronic coupling between adjacent regions. These methods are consistent in that they are derived from a general variational principle and are a systematic extension of the finite element method as applied to diffusion theory reactor analysis.

For purposes of simplicity, the methods will be developed for the case of one-dimensional, time-dependent, multigroup diffusion theory, although it is expected that these methods can be extended to the general spatially dependent kinetics problem with relative ease.

The remainder of this report is organized as follows. Chapter 2 summarizes the use of variational principles and synthesis techniques in time-independent diffusion theory. The difference equations of the finite element methods applied in one dimension are derived using modal-nodal trial function forms in order to illustrate the use of these techniques. The forms of the proposed approximation methods are given in Chapter 3. The resulting finite difference equations are presented and boundary conditions discussed for approximation methods involving both linear and cubic Hermite basis functions. The numerical properties of the resulting matrix equations, as well as their numerical solution scheme, and useful programming techniques are discussed in Chapter 4. Chapter 5 presents results of the proposed methods for four representative one-dimensional PWR configurations, and compares the results with those of coarse mesh finite element methods. Finally, Chapter 6 presents conclusions and recommendations as well as comments concerning the possibility of extending the proposed methods to multidimensional geometries.



## Chapter 2

VARIATIONAL DERIVATION OF FINITE DIFFERENCE  
APPROXIMATIONS IN TIME-INDEPENDENT  
MULTIGROUP DIFFUSION THEORY

The application of variational calculus to the describing equations of physical systems is perhaps the most general and powerful method of obtaining approximate solutions in mathematical physics. Variational methods seek to combine known "trial functions" into approximate solutions through the use of a variational functional which characterizes the equations of the system.

Essentially, variational methods consist of first finding a characteristic functional whose first-order variation when set to zero yields the describing equations of the system as its Euler equations. A class of trial functions, given in terms of known functions and unknown coefficients (or functions), is then chosen to approximate the solutions of the describing equations. These trial functions are then substituted into the variational functional, and its first variation is set to zero. Allowing arbitrary variations in all of the trial function unknowns results in a set of relationships among the unknowns. These relationships when solved then yield the "best" obtainable approximate solution within the space of trial functions given.

Variational methods can be thought of as a class of weighted residual methods since "weighting functions" appear in the functional and in the equations that result from setting the first variation of the functional to zero. The weighting functions are determined by the

form of the functional itself; or equivalently, by the set of Euler equations selected to describe the system. In non-self adjoint problems, the adjoint equations are generally included in the set of Euler equations. The inclusion of corresponding "adjoint trial functions" in the functional results in adjoint weighting in the variation equations and allows greater approximation flexibility of the variational method.

## 2.1 Calculus of Variations Applied to Diffusion Theory

The time-independent multigroup diffusion equations as given by Eq. 1.3 can be written as

$$\mathbb{H} \Phi = \frac{1}{\lambda} \mathbb{B} \Phi \quad (2.1a)$$

where

$$\mathbb{H} = -\underline{\nabla} \cdot \mathbb{D} \underline{\nabla} + \mathbb{M} - \mathbf{T} \quad (2.1b)$$

Since the multigroup diffusion equations are not self-adjoint, it is convenient to introduce the adjoint diffusion equations

$$\mathbb{H}^* \Phi^* = \frac{1}{\lambda} \mathbb{B}^* \Phi^* \quad (2.2a)$$

where  $\mathbb{H}^*$  and  $\mathbb{B}^*$  are the adjoint operators corresponding to  $\mathbb{H}$  and  $\mathbb{B}$ , respectively, and are defined as:<sup>3</sup>

$$\mathbb{H}^* = \mathbb{H}^T = -\underline{\nabla} \cdot \mathbb{D} \underline{\nabla} + \mathbb{M} - \mathbf{T}^T \quad (2.2b)$$

$$\mathbb{B}^* = \mathbb{B}^T \quad (2.2c)$$

since  $\mathbb{D}$  and  $\mathbb{M}$  are diagonal.  $\Phi^*$  is the group adjoint flux vector, or importance vector, which must obey the same boundary conditions as  $\Phi$ .

The exact solutions  $\Phi(r)$  and  $\Phi^*(r)$  of the diffusion equations and the adjoint diffusion equations can be approximated by flux and adjoint flux trial functions denoted as  $U(r)$  and  $U^*(r)$  using a variational functional of the form

$$\mathcal{F}_1 [U, U^*] = \frac{1}{\lambda} = \frac{\int_{\mathbf{R}} U^{*\text{T}} \mathbf{H}U \, dr}{\int_{\mathbf{R}} U^{*\text{T}} \mathbf{B}U \, dr} \quad (2.3)$$

where it is assumed that the group-theory flux trial function vectors  $U^*$  and  $U$  as well as the group current vectors  $\mathbf{D}\nabla U^*$  and  $\mathbf{D}\nabla U$  are everywhere continuous, and that  $U^*$  and  $U$  vanish outside the reactor region  $\mathbf{R}$ . Allowing arbitrary trial function variations, denoted by  $\delta U^*$  and  $\delta U$ , making  $\mathcal{F}_1$  stationary first with respect to  $U^*$  and then with respect to  $U$  results<sup>29</sup> in the following equations:

$$\int_{\mathbf{R}} \delta U^{*\text{T}} \left[ \mathbf{H}U - \frac{1}{\lambda} \mathbf{B}U \right] \, dr = 0 \quad (2.4a)$$

$$\int_{\mathbf{R}} \left[ U^{*\text{T}} \mathbf{H} - \frac{1}{\lambda} U^{*\text{T}} \mathbf{B} \right] \delta U \, dr = 0 \quad (2.4b)$$

The above equations, containing the desired Euler equations, are the equations upon which the approximation method is based.

A significant characteristic of this approximation form is the property of exact solution reproduction. Although general choices of the trial functions  $U$  and  $U^*$  result in approximate eigenvalues which may differ substantially from the exact solution eigenvalue, the exact solutions, when chosen within the given class of trial functions, are yielded as the result of the approximation along with the exact solution eigenvalue.

The nature of the above approximation depends solely upon the forms of the flux trial functions given. Each trial function can be defined in terms of unknown coefficients (or functions) and known functions. Independent variation of the unknown coefficients of the adjoint trial function in Eq. 2.4a will yield the "best" flux solution obtainable for that class of flux and adjoint flux trial functions given. The corresponding "best" adjoint flux solution can be found in an analogous manner using Eq. 2.4b. These techniques are illustrated in the next section.

Another functional incorporating the flux and adjoint flux diffusion equations can be defined as

$$\mathcal{F}_2[U^*, U] = \int_{\mathbf{R}} U^{*\text{T}} [\mathbb{H}U - \frac{1}{\lambda} \mathbb{B}U] dr \quad (2.5)$$

Although the forms of the above functionals differ, it can be shown that both produce the same variation equations, Eqs. 2.4, when made stationary. The form of  $\mathcal{F}_2$  and its first variation are much less complex than the form and first variation of  $\mathcal{F}_1$ . For these reasons, functionals of the form of  $\mathcal{F}_2$  will be used in this report.

## 2.2 Discontinuous Trial Functions

The addition of discontinuous flux trial functions into the class of allowable trial functions for use in diffusion theory variational methods greatly enhances and generalizes the versatility of such methods.<sup>25</sup> However, special provisions must be made in the approximation method itself in order that such trial functions can be properly used.<sup>23, 24, 30</sup> In order to account for the discontinuities in the flux (and in general also the current) trial functions, it is necessary to

include special terms specifying continuity conditions directly within the approximation method. This can be accomplished through the use of a variational functional whose Euler equations include the P-1 equations and continuity conditions for both flux and current. A general functional of this type which allows discontinuous flux, current, and adjoint trial functions can be derived from previous work<sup>30, 31</sup> and is given as follows:

$$\begin{aligned} \mathcal{F}[U^*, U, \underline{V}_+, \underline{V}_-, \alpha, \beta] = & \int_R \{ U^{*\text{T}} [\underline{\nabla} \cdot \underline{V}_+ + \underline{A}U] + \underline{V}_+^{*\text{T}} \cdot [\underline{\nabla}U + \underline{D}^{-1}\underline{V}_-] \} dr \\ & + \int_{\Gamma} \hat{n} \cdot \{ [U_+^{*\text{T}} \alpha + U_-^{*\text{T}} (I-\alpha)] (\underline{V}_+ - \underline{V}_-) \\ & + [\underline{V}_+^{*\text{T}} \beta + \underline{V}_-^{*\text{T}} (I-\beta)] (U_+ - U_-) \} ds \end{aligned} \quad (2.6)$$

where  $U^*$ ,  $U$ ,  $\underline{V}_+$ , and  $\underline{V}_-$  are the group flux and group current approximations to  $\Phi^*$ ,  $\Phi$ ,  $\underline{J}_+$ , and  $\underline{J}_-$ , respectively, and where the first integral extends over the volume  $R$  of the reactor and the second extends over all interior surfaces  $\Gamma$  upon which discontinuities are defined.  $\hat{n}$  is the unit vector perpendicular to interior surfaces, and quantities evaluated on sides of surfaces toward which  $\hat{n}$  is pointing are denoted with the subscript (+). Quantities evaluated on sides of surfaces from which  $\hat{n}$  is pointing are denoted with the subscript (-).  $\alpha$  and  $\beta$  are in general  $G \times G$  undefined variable matrices, and  $I$  is in general a  $G \times G$  unit matrix, which allow a general treatment of the discontinuities.

The restrictions generally imposed upon trial functions for use in functionals of this type are the following:

1. The trial functions must be piecewise continuous.
2. The trial functions  $U$  and  $\underline{V}^*$  as well as  $U^*$  and  $\underline{V}$  are not allowed to be discontinuous at the same point.
3. The components of  $U^T \underline{V}^*$  and  $U^{*T} \underline{V}$  normal to the exterior surface of the reactor must vanish.

Due to restriction 2, the general quantities  $\alpha$  and  $\beta$  always cancel and are never used within these approximation methods.

The first variation of  $\mathcal{F}$  can be found in a straightforward manner, and can be simplified to the following form which indicates the desired P-1 and adjoint P-1 equations and the trial function continuity conditions as Euler equations:

$$\begin{aligned}
\delta \mathcal{F} = & \int_{\mathbf{R}} \left\{ \delta U^{*T} [\underline{\nabla} \cdot \underline{V} + \mathbf{\Lambda} U] + \delta \underline{V}^{*T} \cdot [\underline{\nabla} U + \mathbf{D}^{-1} \underline{V}] \right. \\
& + [-\underline{\nabla} U^{*T} + \underline{V}^{*T} \mathbf{D}^{-1}] \cdot \delta \underline{V} + [-\underline{\nabla} \cdot \underline{V}^{*T} + U^{*T} \mathbf{\Lambda}] \delta U \Big\} dr \\
& + \int_{\Gamma} \hat{n} \cdot \left\{ \delta U^{*T} (\underline{V}_+ - \underline{V}_-) + \delta \underline{V}^{*T} (U_+ - U_-) \right. \\
& \quad \left. + (U_-^* - U_+^*)^T \delta \underline{V} + (\underline{V}_-^* - \underline{V}_+^*) \delta U \right\} ds
\end{aligned} \tag{2.7}$$

In most applications, only approximations to the flux and current solutions are desired. In such instances variations in only the adjoint trial functions need be taken. Setting the first variation of  $\mathcal{F}$  equal to zero under these conditions and imposing the above trial function restrictions results in the following variation equation for flux and current approximations:

$$\begin{aligned}
& \int_{\mathbf{R}} \{ \delta U^{*\text{T}} [\underline{\nabla} \cdot \underline{V} + \mathbf{\Lambda}U] + \delta \underline{V}^{*\text{T}} \cdot [\underline{\nabla}U + \mathbb{D}^{-1}\underline{V}] \} d\mathbf{r} \\
& + \int_{\Gamma} \hat{\mathbf{n}} \cdot \{ \delta U^{*\text{T}} (\underline{V}_+ - \underline{V}_-) + \delta \underline{V}^{*\text{T}} (U_+ - U_-) \} ds = 0
\end{aligned} \tag{2.8}$$

The above approximation can also be expressed independently of adjoint trial functions. If the adjoint trial functions are defined as

$$U^* = U \tag{2.9a}$$

$$\underline{V}^* = -\underline{V} \tag{2.9b}$$

then Eq. 2.8 reduces to the Rayleigh-Ritz Galerkin method, a weighted residual method based upon flux weighting.

Regardless of the choice of weighting, the variation equations can be further simplified for those approximation methods which require the currents to obey explicitly Fick's laws:

$$\underline{V} = -\mathbb{D}\underline{\nabla}U \tag{2.10a}$$

$$\underline{V}^* = +\mathbb{D}\underline{\nabla}U^* \tag{2.10b}$$

Under these conditions the variation equations for discontinuous flux and discontinuous current trial functions reduce to

$$\begin{aligned}
& \int_{\mathbf{R}} \{ \delta U^{*\text{T}} \mathbf{\Lambda}U - \delta \underline{V}^{*\text{T}} \cdot \mathbb{D}^{-1}\underline{V} \} d\mathbf{r} \\
& + \int_{\Gamma} \hat{\mathbf{n}} \cdot \{ (\delta U_-^* - \delta U_+^*) \underline{V} + \delta \underline{V}^{*\text{T}} (U_+ - U_-) \} ds = 0
\end{aligned} \tag{2.11}$$

If in addition the flux is required to be everywhere continuous, the variation equations reduce to the appealing forms

$$\int_{\mathbf{R}} \{ \delta U^{*\text{T}} \mathbf{\Lambda}U - \delta \underline{V}^{*\text{T}} \cdot \mathbb{D}^{-1}\underline{V} \} d\mathbf{r} = 0 \tag{2.12a}$$

or equivalently

$$\int_{\mathbf{R}} \{ \delta U^{*T} \mathbf{A} U + (\nabla \delta U^*) \cdot \mathbb{D}(\nabla U) \} dr = 0 \quad (2.12b)$$

Variation equations 2.11 and 2.12 are the approximation equations which are used with the finite element methods and the proposed approximation methods.

### 2.3 The Finite Element Approximation Methods <sup>32,33</sup>

This section introduces the notation and techniques used in conjunction with the modal-nodal variational analysis of the finite element method approximations in one-dimensional multigroup diffusion theory. These fundamentals are presented in these simple approximations before applying them to the more general proposed approximation method in the next chapter.

The one-dimensional problem is defined by the continuous variable  $z$  and divided into  $K$  adjoining regions which are in general inhomogeneous. Each region  $k$  is bounded by nodes  $z_k$  and  $z_{k+1}$  and has width  $h_k = z_{k+1} - z_k$ . It is convenient to define the dimensionless variable  $x$  within each region  $k$  as

$$x = \frac{z - z_k}{h_k} \quad (2.13a)$$

so that region  $k$  can be described in terms of  $z$  as

$$z_k \leq z \leq z_k + h_k = z_{k+1} \quad (2.13b)$$

or equivalently in terms of  $x$  as

$$0 \leq x \leq 1 \quad (2.13c)$$

for each of the regions  $k, k = 1$  to  $K$ . This notation will be used throughout this report.



### 2.3.1 The Conventional Finite Difference Equations

The conventional nodal flux-averaged, three-point, finite difference equations of one-dimensional diffusion theory can be derived from Eq. 2.8 using discontinuous flux and current multigroup column vector trial functions of the following form:<sup>7,8 †</sup>

$$\begin{aligned}
 & \left. \begin{aligned} U(z) = F_k \\ U^*(z) = F_k^* \end{aligned} \right\} \left\{ \begin{aligned} & \left\{ z_{k-1} + \frac{1}{2}h_{k-1} \right\} < z < \left\{ z_k + \frac{1}{2}h_k \right\}; \quad k=1 \text{ to } K+1. \\ & z_1 \quad \text{if } k = 1 \\ & 0 \quad \text{otherwise} \end{aligned} \right. \\
 & \left. \begin{aligned} V(z) = G_k \\ V^*(z) = G_k^* \end{aligned} \right\} \left\{ \begin{aligned} & z_k < z < z_k + h_k; \quad k = 1 \text{ to } K. \\ & 0 \quad \text{otherwise} \end{aligned} \right.
 \end{aligned} \tag{2.14}$$

The forms of these trial functions are illustrated in Figure 2.1.

Inserting these trial functions into variation equation 2.8 results in the equation

$$\begin{aligned}
 & \delta F_1^{*T} \left\{ \int_0^{\frac{1}{2}} \Lambda_1 F_1 h_1 dx + G_1 - G_0 \right\} \\
 & + \sum_{k=2}^K \delta F_k^{*T} \left\{ \int_{\frac{1}{2}}^1 \Lambda_{k-1} F_k h_{k-1} dx + \int_0^{\frac{1}{2}} \Lambda_k F_k h_k dx + G_k - G_{k-1} \right\} \\
 & + \delta F_{K+1}^{*T} \left\{ \int_{\frac{1}{2}}^1 \Lambda_K F_{K+1} h_k dx + G_{K+1} - G_K \right\} \\
 & + \sum_{k=1}^K \delta G_k^{*T} \left\{ \int_0^1 \mathbb{D}_k^{-1} G_k h_k dx + F_{k+1} - F_k \right\} = 0
 \end{aligned} \tag{2.15}$$

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<sup>†</sup>Shifting the domain of definition of the trial functions results in other approximation schemes with equivalently averaged nuclear constants.<sup>34</sup>

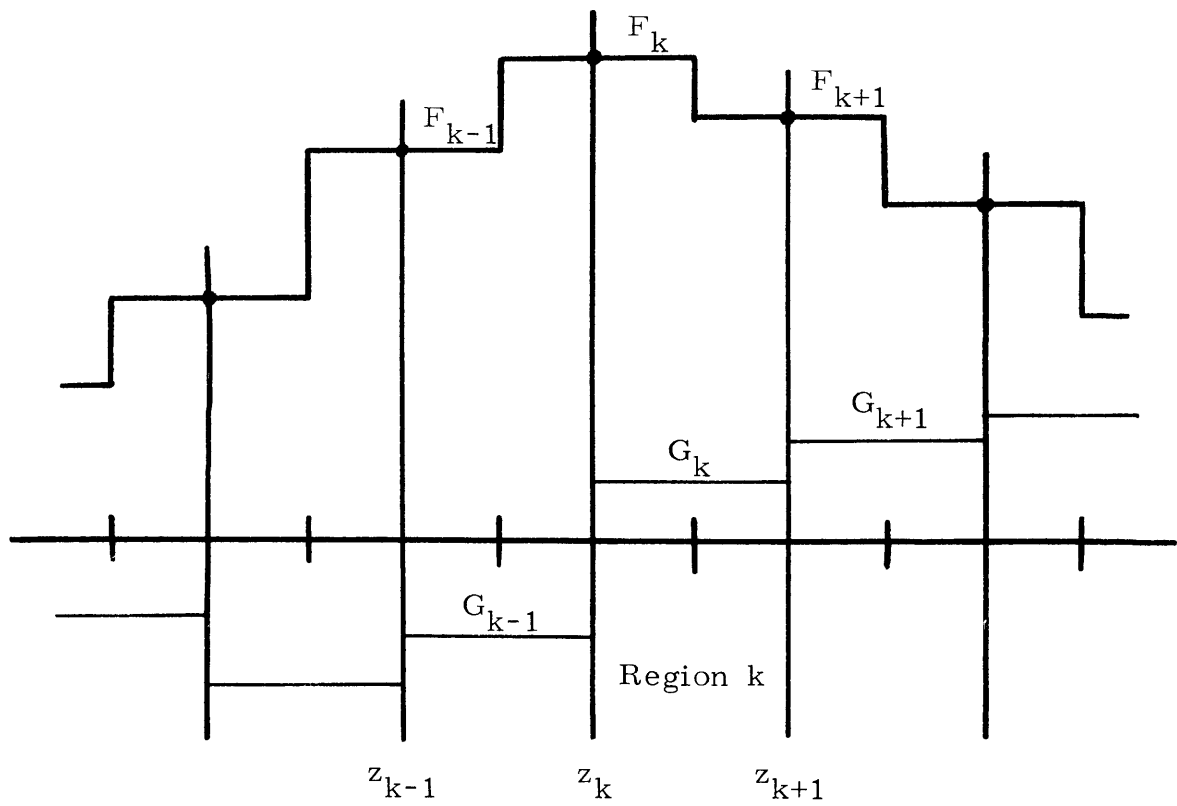


Figure 2.1. Conventional Nodal Finite Difference Approximation Trial Function Forms

Independent variation of all  $F_k^*$  and  $G_k^*$  then results in a system of  $2K+1$  equations and  $2K+3$  unknowns (including  $G_0$  and  $G_{K+1}$ ). The choice of boundary conditions supplies the missing equations. Zero flux boundary conditions can be imposed by setting  $F_1 = F_{K+1} = 0$ , which also requires  $\delta F_1^* = \delta F_{K+1}^* = 0$  thereby eliminating  $G_0$  and  $G_{K+1}$ , and results in a system of  $2K-1$  equations and  $2K-1$  unknowns. Symmetry boundary conditions can be imposed on the left by  $G_0 = -G_1$  and on the right by  $G_{K+1} = -G_K$ , resulting in a system of  $2K+1$  equations and  $2K+1$  unknowns.

Elimination of all  $G_k$ ,  $k=1$  to  $K$ , results in the standard three-point difference equations

$$b_1 F_1 + c_1 F_2 = 0 \quad (2.16a)$$

$$a_k F_{k-1} + b_k F_k + c_k F_{k+1} = 0; \quad k=2 \text{ to } K \quad (2.16b)$$

$$a_{K+1} F_K + b_{K+1} F_{K+1} = 0 \quad (2.16c)$$

where Eqs. 2.16a and 2.16c are used for the cases of symmetry boundary conditions. The  $G \times G$  matrix coefficients  $\{a_k, b_k, c_k\}$  are of the form  $A - \frac{1}{\lambda} B$  and are defined assuming homogeneous regional nuclear constants in section 1 of Appendix B. The matrix form of Eqs. 2.16 for the use of zero flux boundary conditions on the left and symmetry on the right is illustrated in Figure 2.2.

### 2.3.2 Multichannel Polynomial Synthesis

The one-dimensional neutron flux  $\Phi_k(z)$  defined as nonzero only within each region  $k$  for each region ( $k=1$  to  $K$ ) can be approximated within each region as a polynomial of order  $N$  by the power series

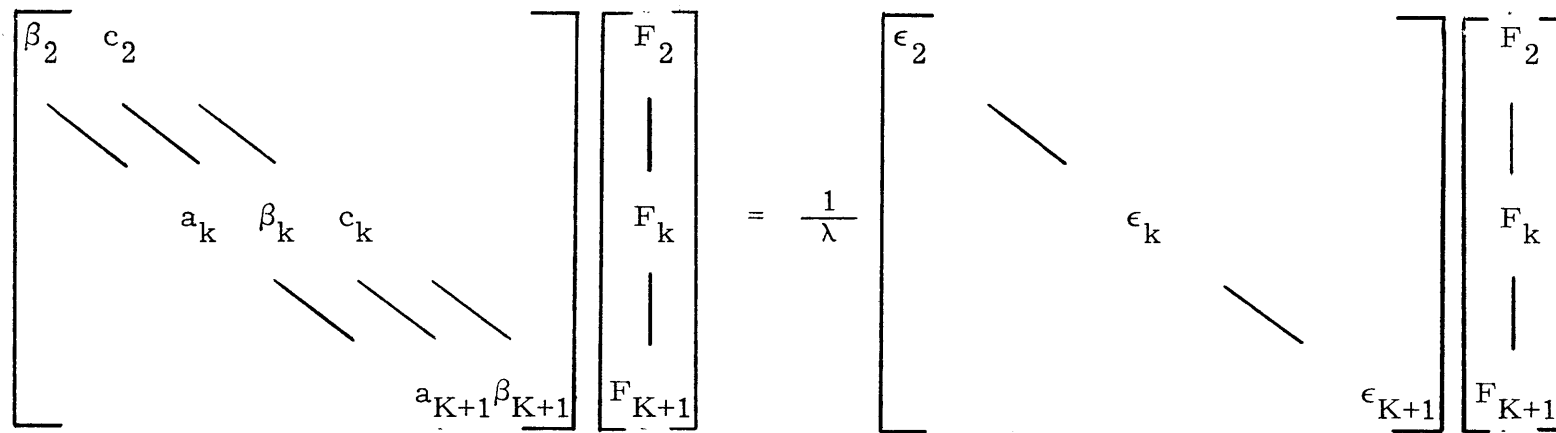


Figure 2.2. Matrix Form of the Conventional Finite Difference Equations. Boundary conditions chosen are zero flux on the left and symmetry on the right.

$$a_k F_{k-1} + b_k F_k + c_k F_{k+1} = 0; \quad k = 2 \text{ to } K.$$

$$a_{K+1} F_K + b_{K+1} F_{K+1} = 0.$$

where:  $b_k = \beta_k - \frac{1}{\lambda} \epsilon_k; \quad k = 2 \text{ to } K+1.$

$$U_k^{(N)}(z) = \sum_{i=0}^N a_{k,i} x^i \quad (2.17)$$

where the distinction between  $z$  and  $x$  is understood since  $0 \leq x \leq 1$  within each region  $k$ . Such approximations are not useful in diffusion theory because: (1) the resulting matrix equations relating the  $a_{k,i}$ 's contain full matrices similar to Hilbert matrices which may be very difficult to solve; and (2) such matrices are almost always highly singular and may produce numerical instabilities in the solution method. These difficulties can be eliminated by employing polynomials in the trial functions in the following form:

$$U_k^{(N)}(z) = \sum_{i=0}^N p_i^{(N)}(x) F_{k+\frac{i}{N}} \quad (2.18)$$

where the  $p_i^{(N)}(x)$  are polynomials in  $x$  of degree  $N$ . This form is convenient because for a particular selection of the  $p_i^{(N)}(x)$  the unknowns  $F_{k+\frac{i}{N}}$  can be defined as the approximate flux solution evaluated at points  $z_i + \frac{i}{N}$  within region  $k$ . For high order approximations,  $i > 0$ , the flux can be made continuous by imposing the following restrictions on  $p_i^{(N)}(x)$ :

$$p_i^{(N)}\left(\frac{\ell}{N}\right) = \begin{cases} 1 & \ell = i \\ 0 & \ell \neq i \end{cases} \text{ for } \ell = 0 \text{ to } N \quad (2.19)$$

The specific polynomial flux approximations of this form through degree  $N=3$  are given below:

$$U_k^{(0)}(x) = F_k \quad (2.20a)$$

$$U_k^{(1)}(x) = (1-x)F_k + xF_{k+1} \quad (2.20b)$$

$$U_k^{(2)}(x) = (1-3x+2x^2)F_k + (4x-4x^2)F_{k+\frac{1}{2}} + (-x+2x^2)F_{k+1} \quad (2.20c)$$

$$U_k^{(3)}(x) = \left(1 - \frac{11}{2}x + 9x^2 - \frac{9}{2}x^3\right)F_k + \left(9x - \frac{45}{2}x^2 + \frac{27}{2}x^3\right)F_{k+\frac{1}{3}} \\ + \left(-\frac{9}{2}x + 18x^2 - \frac{27}{2}x^3\right)F_{k+\frac{2}{3}} + \left(x - \frac{9}{2}x^2 + \frac{9}{2}x^3\right)F_{k+1} \quad (2.20d)$$

An immediate drawback of these approximations lies in the definitions of the corresponding current trial functions. Given a flux polynomial approximation of degree  $N$ , polynomial approximations for the current can be of order zero through  $N$ , and may even be of higher order than the flux approximation. Each set of chosen trial function pairs ultimately results in a characteristic complex band-structured matrix problem which may or may not have desirable numerical solution properties and is usually very difficult to solve.

Such problems can be eliminated by noting that the use of variational analysis attempts to force the current approximation to obey Fick's law. The obvious solution is direct use of Fick's law in the trial function forms

$$V_k(z) = -\mathbb{D}_k(z) \frac{d}{dz} U_k(z) \quad (2.21)$$

which results in simple band-structured matrix equations relating only flux unknowns. The use of current polynomial approximations of order  $N-1$  as given in Eqs. 2.20 with flux approximations of order  $N$ , however, does not improve the situation.

The accuracy of these difference equations can be found first by eliminating all non-integer subscripted unknowns, then expanding the resulting three-point difference equations in a Taylor series about

node  $k$ , and comparing results to the exact three-point difference solution known for the one-dimensional case.<sup>7, 35</sup> By comparison of terms containing equal powers of  $h_k$ , it can be shown that the  $N=1$  and  $N=2$  polynomial approximations are accurate to order  $\theta(h^2)$  while the  $N=3$  approximation is accurate to order  $\theta(h^3)$ .

The approximation of a function by a polynomial of order  $N$  leads immediately to the concept of basis functions. The  $N+1$  polynomial functions which multiply the  $N+1$  unknowns in Eq. 2.18 form a basis for the approximation and can be called basis functions. The simplicity of basis functions becomes apparent in an error analysis of the approximation as follows. An approximate solution  $U^{(N)}(z)$  of order  $N$  to the exact one-dimensional solution  $\Phi(z)$  can be expressed as

$$U^{(N)}(z) = \sum_{k=1}^K \Phi(z_k) \Omega_k^{(N)}(z) \quad (2.22a)$$

where  $\Omega_k^{(N)}(z)$  is a basis function of order  $N$  centered about node  $z_k$ . By Taylor series expansion about any node, it can be shown<sup>36</sup> that if  $\Omega_k(z)$  satisfies:

$$\sum_{k=1}^K z_k^\alpha \Omega_k^{(N)}(z) = \left(\frac{z}{h_k}\right)^\alpha \quad \text{for } |\alpha| \leq N \quad (2.22b)$$

then  $U^{(N)}(z)$  is an approximation to  $\Phi(z)$  accurate to order  $\theta(h_k^{N+1})$ .

Basis functions found using Eqs. 2.22 are unique for each  $N$  and generally extend over surrounding regions. The forms of the basis functions for  $N \leq 3$  are summarized below and illustrated in Figure 2.3. Since the following basis functions are symmetric, only the right half,  $z \geq z_k$ , is expressly given.

$$N = 0: \quad \Omega_k^{(0)}(z) = \begin{cases} 1 & z_k \leq z < z_k + \frac{1}{2}h_k \\ 0 & \text{otherwise} \end{cases} \quad (2.23a)$$

$$N = 1: \quad \Omega_k^{(1)}(z) = \begin{cases} (1-x) & z_k \leq z \leq z_{k+1} \\ 0 & \text{otherwise} \end{cases} \quad (2.23b)$$

$$N = 2: \quad \Omega_k^{(2)}(z) = \begin{cases} \frac{3}{4} - x^2 & z_k \leq z \leq z_k + \frac{1}{2}h_k \\ \frac{1}{2} \left( \frac{3}{2} - x \right)^2 & z_k + \frac{1}{2}h_k \leq z \leq z_{k+1} \\ \frac{1}{2} \left( \frac{1}{2} - x \right)^2 & z_{k+1} \leq z \leq z_{k+1} + \frac{1}{2}h_{k+1} \\ 0 & \text{otherwise} \end{cases} \quad (2.23c)$$

$$N = 3: \quad \Omega_k^{(3)}(z) = \begin{cases} \frac{1}{36} (30 - 54x^2 + 28x^3) & z_k \leq z \leq z_{k+1} \\ \frac{1}{36} (4 - 24x + 30x^2 - 11x^3) & z_{k+1} \leq z \leq z_{k+2} \\ \frac{1}{36} (-1 + 3x - 3x^2 + x^3) & z_{k+2} \leq z \leq z_{k+3} \\ 0 & \text{otherwise} \end{cases} \quad (2.23d)$$

where  $0 \leq x \leq 1$  within each region  $k$  in the above cases.

Use of these basis functions results in approximate solutions which are continuous for  $N \geq 1$  and whose derivatives  $dU^{(N)}(z)/dz$  through  $d^{N-1}U^{(N)}(z)/dz^{N-1}$  are also continuous. In high order approximations in diffusion theory, it is advantageous to retain flux and current continuity and employ basis functions defined over two adjacent regions in order to produce three-point difference equations. This can be accomplished in the  $N=3$  approximations with the cubic Hermite basis functions.



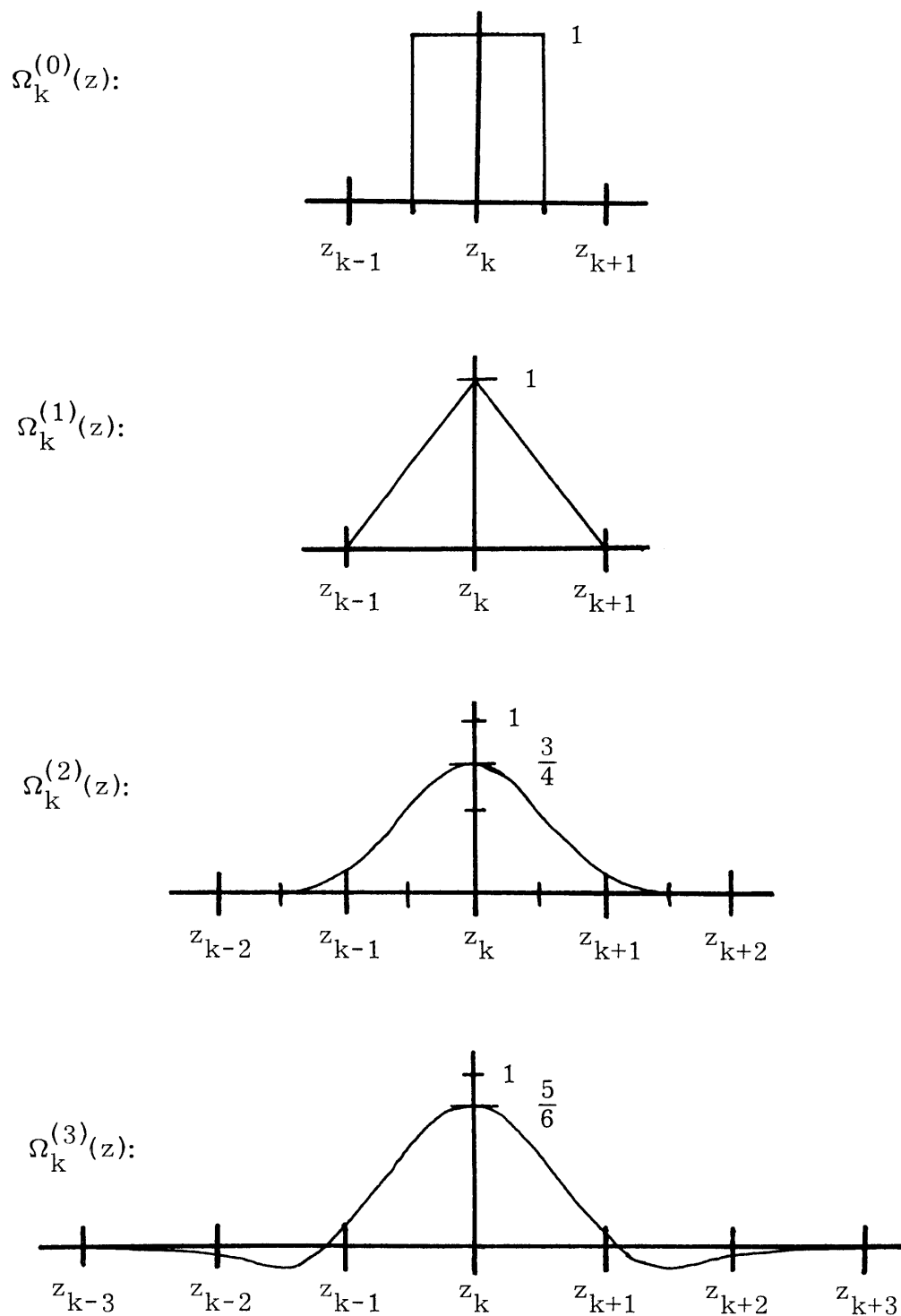


Figure 2.3. Basis Functions of Eqs. 2.23  
for  $N = 0, 1, 2,$  and  $3$

The above cubic basis function  $\Omega_k^{(3)}(z)$  can be constructed from a combination of either cubic B splines,  $\Omega_k^B(z)$ , or cubic Hermite polynomials,  $\Omega_k^{H_1}(z)$  and  $\Omega_k^{H_2}(z)$ , as follows: <sup>37</sup>

$$\Omega_k^{(3)}(z) = -\frac{1}{6} \Omega_{k-1}^B(z) + \frac{4}{3} \Omega_k^B(z) - \frac{1}{6} \Omega_{k+1}^B(z) \quad (2.24)$$

where:

$$\Omega_k^B(z) = \frac{2}{3} \Omega_k^{H_1}(z) + \frac{1}{6} \Omega_{k+1}^{H_1}(z) - \frac{1}{2} \Omega_{k+1}^{H_2}(z) \quad (2.25)$$

The forms of these cubic B and Hermite polynomials are given below and illustrated in Figures 2.5 and 2.6. Again, only the right half of the functions are expressly given as  $\Omega_k^B$  and  $\Omega_k^{H_1}$  are symmetric, while  $\Omega_k^{H_2}$  is antisymmetric.

$$\Omega_k^B(z) = \begin{cases} \frac{2}{3} - x^2 + \frac{1}{2}x^3 & z_k \leq z \leq z_{k+1} \\ \frac{1}{6}(1-x)^3 & z_{k+1} \leq z \leq z_{k+2} \\ 0 & \text{otherwise} \end{cases} \quad (2.26)$$

$$\Omega_k^{H_1}(z) = \begin{cases} 1 - 3x^2 + 2x^3 & z_k \leq z \leq z_{k+1} \\ 0 & \text{otherwise} \end{cases} \quad (2.27a)$$

$$\Omega_k^{H_2}(z) = \begin{cases} x - 2x^2 + x^3 & z_k \leq z \leq z_{k+1} \\ 0 & \text{otherwise} \end{cases} \quad (2.27b)$$

where again  $0 \leq x \leq 1$  in each region  $k$ .

The fact that the cubic Hermite polynomials form a basis for the cubic basis functions and extend over only two adjacent regions makes them very attractive for use in diffusion theory approximation methods.

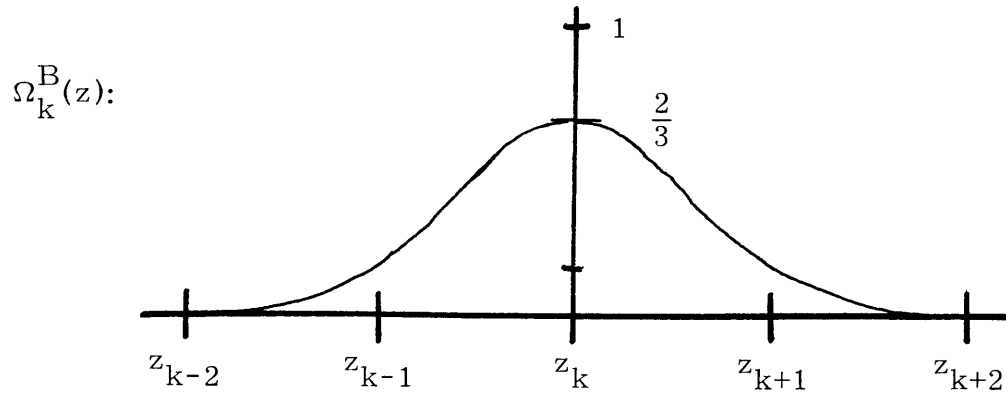


Figure 2.4. Cubic B Spline  $\Omega_k^B(z)$  of Eq. 2.26

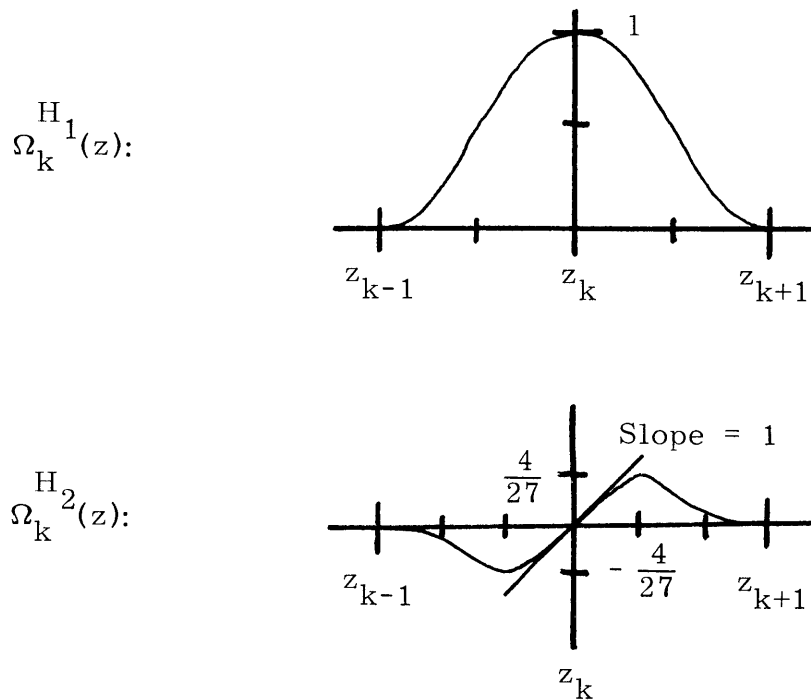


Figure 2.5. Cubic Hermite Basis Functions

$\Omega_k^{H_1}(z)$  and  $\Omega_k^{H_2}(z)$  of Eqs. 2.27

### 2.3.3 The Linear Basis Function Approximation

The group flux trial functions defined as nonzero within each region  $k$  can be expressed in modal-nodal form in terms of linear basis functions as

$$U_k(z) = (1-x)F_k + xF_{k+1} \quad (2.28a)$$

;  $k = 1$  to  $K$ .

$$U_k^*(z) = (1-x)F_k^* + xF_{k+1}^* \quad (2.28b)$$

where  $F_k$  is the approximate group flux column vector at node  $z_k$  and  $0 \leq x \leq 1$  with each region  $k$ . Although the flux trial functions are continuous, the current trial functions defined within each region by Eqs. 2.10 are not, and are given by

$$V_k(z) = \frac{1}{h_k} \mathbb{D}_k(x) [F_k - F_{k+1}] \quad (2.28c)$$

;  $k = 1$  to  $K$ .

$$V_k^*(z) = \frac{1}{h_k} \mathbb{D}_k(x) [F_{k+1}^* - F_k^*] \quad (2.28d)$$

Insertion of these trial function forms into variation equation 2.12a results in the equation

$$\sum_{k=1}^K h_k \int_0^1 \left\{ [(1-x)\delta F_k^* + x\delta F_{k+1}^*]^T \mathbf{\Lambda}_k(x) [(1-x)F_k + xF_{k+1}] + [\delta F_k^* - \delta F_{k+1}^*]^T \frac{1}{h_k} \mathbb{D}_k(x) [F_k - F_{k+1}] \right\} dx = 0 \quad (2.29)$$

Allowing arbitrary variations in all  $F_k^*$  results in a system of  $K+1$  equations and  $K+1$  unknowns which can be written as:

$$b_1 F_1 + c_1 F_2 = 0 \quad (2.30a)$$

$$a_k F_{k-1} + b_k F_k + c_k F_{k+1} = 0 \quad ; \quad k = 2, K. \quad (2.30b)$$

$$a_{K+1} F_K + b_{K+1} F_{K+1} = 0 \quad (2.30c)$$

where the  $G \times G$  matrix coefficients  $\{a_k, b_k, c_k\}$  are of the form  $A - \frac{1}{\lambda} B$  and are defined assuming homogeneous regional nuclear constants in section 2 of Appendix B. Zero flux boundary conditions can be imposed by use of only Eq. 2.30b with  $F_1 = F_{K+1} = 0$ , while symmetry boundary conditions require the use of the other equations as well. The matrix form of these equations for the boundary conditions of zero flux on the left and symmetry on the right is given in Figure 2.6.

#### 2.3.4 The Cubic Hermite Basis Function Approximation<sup>38, 39</sup>

The cubic Hermite polynomials can be incorporated into modal-nodal flux trial functions which allow continuous flux and continuous current by defining the flux trial functions within each region  $k$  as

$$U_k(z) = (1-3x^2+2x^3)F_k + (3x^2-2x^3)F_{k+1} \\ + (-x+2x^2-x^3) \frac{\theta}{h_k} \mathbb{D}_k^{-1}(x)G_k + (x^2-x^3) \frac{\theta}{h_k} \mathbb{D}_k^{-1}(x)G_{k+1} \quad (2.31a)$$

$$U_k^*(z) = (1-3x^2+2x^3)F_k^* + (3x^2-2x^3)F_{k+1}^* \\ + (-x+2x^2-x^3) \frac{\theta}{h_k} \mathbb{D}_k^{-1}(x)G_k^* + (x^2-x^3) \frac{\theta}{h_k} \mathbb{D}_k^{-1}(x)G_{k+1}^* \quad (2.31b)$$

where  $k = 1$  to  $K$ .

$F_k$  is again the approximate group flux solution vector at node  $z_k$ , and  $G_k$  is proportional to the approximate group current solution vector at node  $z_k$ . Application of Fick's law defines the current trial functions for each  $k$  as



$$\begin{aligned}
V_k(z) &= \frac{1}{h_k} \mathbb{D}_k(x)(6x-6x^2) [F_{k+1} - F_k] \\
&\quad + (1-4x+3x^2)\theta G_k + (-2x+3x^2)\theta G_{k+1}
\end{aligned} \tag{2.31c}$$

$$\begin{aligned}
V_k^*(z) &= \frac{1}{h_k} \mathbb{D}_k(x)(6x-6x^2) [F_k^* - F_{k+1}^*] \\
&\quad + (-1+4x-3x^2)\theta G_k^* + (2x-3x^2)\theta G_{k+1}^*
\end{aligned} \tag{2.31d}$$

Continuity of flux and current are automatically guaranteed since

$$\begin{aligned}
U_k(0) &= U_{k-1}(h_{k-1}) = F_k \\
U_k^*(0) &= U_{k-1}^*(h_{k-1}) = F_k^*
\end{aligned} \tag{2.32}$$

$$V_k(0) = V_{k-1}(h_{k-1}) = \theta G_k$$

$$V_k^*(0) = V_{k-1}^*(h_{k-1}) = -\theta G_k^* \dagger$$

The normalization constant  $\theta$  is introduced in order to produce stiffness matrices having small condition numbers and can be chosen such that  $\frac{\theta}{D_k(0)} \approx 1$ .

Insertion of these trial function forms into variation equation 2.12a results in a lengthy equation which can be written as follows:

---

<sup>†</sup>Such a choice of  $-G_k^*$  allows the matrix of coefficients to be positive definite. Cf., Chapter 4.

$$\begin{aligned}
& \delta F_1^{*T} \{ b1_1 F_1 + b2_1 G_1 + c1_1 F_2 + c2_1 G_2 \} \\
& \delta G_1^{*T} \{ b3_1 F_1 + b4_1 G_1 + c3_1 F_2 + c4_1 G_2 \} \\
& + \sum_{k=2}^K \delta F_k^{*T} \{ a1_k F_{k-1} + a2_k G_{k-1} + b1_k F_k + b2_k G_k + c1_k F_{k+1} + c2_k G_{k+1} \} \\
& + \sum_{k=2}^K \delta G_k^{*T} \{ a3_k F_{k-1} + a4_k G_{k-1} + b3_k F_k + b4_k G_k + c3_k F_{k+1} + c4_k G_{k+1} \} \\
& \delta F_{K+1}^{*T} \{ a1_{K+1} F_K + a2_{K+1} G_K + b1_{K+1} F_{K+1} + b2_{K+1} G_{K+1} \} \\
& \delta G_{K+1}^{*T} \{ a3_{K+1} F_K + a4_{K+1} G_K + b3_{K+1} F_{K+1} + b4_{K+1} G_{K+1} \} = 0.
\end{aligned} \tag{2.33}$$

where the  $G \times G$  matrix coefficients  $\{a1, \dots, c4\}$  are of the form  $A - \frac{1}{\lambda} B$  and are defined assuming homogeneous regional nuclear constants in section 3 of Appendix B.

The choice of either zero flux,  $F_k = 0$  as well as  $\delta F_k^* = 0$ , or zero current,  $G_k = 0$  as well as  $\delta G_k^* = 0$ , boundary conditions for  $k = 1$  or  $K + 1$  along with arbitrary variations of the remaining  $F_k^*$  and  $G_k^*$  results in a system of  $2K$  equations and  $2K$  unknowns. Figure 2.7 illustrates the matrix form of such a system for the case of zero flux on the left and zero current on the right boundary conditions.

The basis functions and approximation techniques presented in this section are applied to the proposed approximation methods in the next chapter. Also, various techniques for treating zero flux and symmetry boundary conditions are discussed. The matrix properties of the equations resulting from the above finite element approximations and their solution methods are discussed in Chapter 4.





## Chapter 3

DEVELOPMENT OF A CONSISTENT COARSE MESH  
APPROXIMATION METHOD3.1 Formulation

The finite element methods have been shown<sup>32, 33</sup> to approximate accurately flux solutions and criticality measurements of multigroup diffusion theory when applied to problems allowing homogeneous nuclear material within the mesh regions. Use of such homogeneous material, while simplifying the calculation of the matrix elements (since numerical integrations are not required), may result in limiting the region mesh sizes allowed unless some type of homogenization procedure is used. If the mesh spacing is chosen such that some or all mesh regions are heterogeneous, then direct application of the variational techniques given in Chapter 2 results in weight averaging the nuclear constants with products of the basis functions and their derivatives, as given by the approximation. Although such a procedure is a direct application of the finite element technique, the accuracy of such methods depends upon the placement of the mesh regions and may vary significantly as their placement is altered.

A more useful homogenization procedure which is commonly used in reactor diffusion theory analysis allows the nuclear material within each mesh region to be homogenized by flux weighting with an assumed flux shape determined *a priori* within that region in order (hopefully) to preserve reaction rates.

In large reactors the core can be thought of as composed of a lattice of heterogeneous fuel subassemblies containing fuel, clad, coolant channels, and/or absorption control rods. Each subassembly can be divided into several distinct homogeneous regions whose few-group microcell macroscopic nuclear constants are found by multi-group energy-dependent calculations.<sup>40</sup> Detailed subassembly solutions,  $\psi_k(r)$ , are then found for each subassembly  $k$  by assuming that the current on the boundary of the subassemblies is zero. Flux weighting the nuclear material in each subassembly with the corresponding detailed subassembly solution for each subassembly region then results in regional homogeneous nuclear constants  $\langle \Sigma_k \rangle$  which may better approximate the physics of the region.

$$\langle \Sigma_k \rangle = \frac{\int_k \psi_k(r) \Sigma_k(r) dr}{\int_k \psi_k(r) dr} \quad (3.1)$$

Proper use of detailed flux weighted constants can lead to accurate criticality measurements, but the detailed *a priori* fine flux structure within each region is lost since it appears only in cross-section homogenization and not in the approximation. Attempts to retain the fine flux structure have only recently been proposed in several multichannel synthesis approximations.<sup>27, 41, 42, 43</sup>

Unfortunately, each of these approximations are approximations in themselves and do not reduce to desirable approximations if the detailed flux solutions are themselves constant, as would be the case in large homogeneous regions.

Just as the discontinuous multichannel synthesis approximation method can be shown to reduce to low order difference equations (of the type which could result using the finite element method with constant or flat basis functions) when constant trial functions are used, approximation methods are presented below which retain the given detailed flux structure and also reduce to the higher order finite element approximations. The use of linear or cubic Hermite basis functions in the approximation provides flux continuity and results in better approximation accuracy.

The approximations are presented and discussed for the case of one-dimensional, multigroup diffusion theory. Extension to higher dimensions remains a problem that will require some further study. The approximations which are the linear basis functions are considered in the next section, while the approximations using the cubic Hermite basis functions are considered in section 3.3.

### 3.2 The Proposed Linear Basis Function Approximations

The proposed approximation method utilizing linear basis functions and defined as nonzero within each mesh region  $k$ ,  $k=1$  to  $K$ , is given by the following modal-nodal trial function forms:

$$U_k(z) = \psi_k(x) [ \psi_k^{-1}(0)(1-x)F_k + \psi_k^{-1}(1)x F_{k+1} ] \quad (3.2a)$$

$$U_k^*(z) = \psi_k^*(x) [ \psi_k^{*-1}(0)(1-x)F_k^* + \psi_k^{*-1}(1)x F_{k+1}^* ] \quad (3.2b)$$

$$V_k(z) = \eta_k(x) [ \psi_k^{-1}(0)(1-x)F_k + \psi_k^{-1}(1)x F_{k+1} ] \\ + \frac{1}{h_k} \mathbb{D}_k(x) \psi_k(x) [ \psi_k^{-1}(0)F_k - \psi_k^{-1}(1)F_{k+1} ] \quad (3.2c)$$

$$\begin{aligned}
V_k^*(z) = & \eta_k^*(x) [ \psi_k^{*-1}(0)(1-x)F_k^* + \psi_k^{*-1}(1)x F_{k+1}^* ] \\
& + \frac{1}{h_k} \mathbb{D}_k(x) \psi_k^*(x) [ \psi_k^{*-1}(1) F_{k+1}^* - \psi_k^{*-1}(0) F_k^* ]
\end{aligned} \tag{3.2d}$$

where:

$$x = (z - z_k) / h_k \tag{3.2e}$$

and  $0 \leq x \leq 1$ , as  $z_k \leq z \leq z_{k+1}$ , for each region  $k = 1$  to  $K$ .

$F_k$  is the unknown approximate group flux column vector at node  $z_k$ , and  $\psi_k$ ,  $\psi_k^*$ ,  $\eta_k$ , and  $\eta_k^*$  are  $G \times G$  diagonal matrices composed of the detailed group flux  $\psi_{g,k}(z)$  and group current  $\eta_{g,k}(z)$  solutions, and their adjoints, defined as nonzero only within region  $k$ . Because of the variable transformation between  $z$  and  $x$ ,  $\psi_k(0)$  represents  $\psi_k(z_k)$ , and  $\psi_k(1)$  represents  $\psi_k(z_{k+1})$ ; neither of which, for the moment, is allowed to be zero for any region. The detailed current solutions are given from the detailed flux solutions by Fick's law as

$$\eta_k(z) = -\mathbb{D}_k(z) \frac{d\psi_k(z)}{dz} \tag{3.3a}$$

$$\eta_k^*(z) = +\mathbb{D}_k(z) \frac{d\psi_k^*(z)}{dz} \tag{3.3b}$$

As a result, the current trial functions are related to the flux trial functions by analogous expressions.

Continuity of the flux is imposed by the form of the trial functions since

$$U_k(0) = U_{k-1}(h_{k-1}) = F_k \tag{3.4a}$$

$$U_k^*(0) = U_{k-1}^*(h_{k-1}) = F_k^* \tag{3.4b}$$

The current trial functions, however, are discontinuous. It is evident

by comparison to Eqs. 2.28, that this approximation reduces to the linear basis function finite element method if the detailed flux solutions for each group are taken to be constant.

Insertion of these trial function forms in Eq. 2.12a results in the following variation equation:

$$\begin{aligned}
& \sum_{k=1}^K h_k \int_0^1 \left\{ \psi_k^{*T}(x) \psi_k^{*T^{-1}}(0)(1-x) \delta F_k^{*T} \Lambda_k(x) U_k(x) \right. \\
& + \psi_k^{*T}(x) \psi_k^{*T^{-1}}(1)x \delta F_{k+1}^{*T} \Lambda_k(x) U_k(x) \\
& + \left[ \eta_k^{*T}(x) \psi_k^{*T^{-1}}(0)(1-x) - \frac{1}{h_k} \mathbb{D}_k(x) \psi_k^{*T}(x) \psi_k^{*T^{-1}}(0) \right] \delta F_k^{*T} \mathbb{D}_k^{-1}(x) V_k(x) \\
& \left. + \left[ \eta_k^{*T}(x) \psi_k^{*T^{-1}}(1)x + \frac{1}{h_k} \mathbb{D}_k(x) \psi_k^{*T}(x) \psi_k^{*T^{-1}}(1) \right] \delta F_{k+1}^{*T} \mathbb{D}_k^{-1}(x) V_k(x) \right\} dx = 0
\end{aligned} \tag{3.5}$$

This equation can be written in the form

$$\begin{aligned}
& \delta F_1^{*T} [b_1 F_1 + c_1 F_2] \\
& + \sum_{k=2}^K \delta F_k^{*T} [a_k F_{k-1} + b_k F_k + c_k F_{k+1}] \\
& + \delta F_{K+1}^{*T} [a_{K+1} F_K + b_{K+1} F_{K+1}] = 0
\end{aligned} \tag{3.6}$$

where the  $G \times G$  matrix coefficients  $\{a_k, b_k, c_k\}$  are integral quantities of the form  $A - \frac{1}{\lambda} B$  and are defined in detail in section 1 of Appendix C.

External zero flux boundary conditions are easily imposed by setting  $F_1 = F_{K+1} = 0$ . This requires that  $F_1^*$  and  $F_{K+1}^*$  must then also be zero, which in turn requires the  $\delta F_1^*$  and  $\delta F_{K+1}^*$  coefficients in

Eqs. 3.5 and 3.6 to vanish. Allowing independent variations in the remaining  $F_k^*$ ,  $k = 2$  to  $K$ , results in a matrix problem of the form illustrated in Figure 2.6 which would contain  $K-1$  equations and  $K-1$  unknowns.

Zero current boundary equations are found using symmetry considerations. If, for example, a zero current or symmetry boundary condition is imposed on the right at  $z_{K+1}$ , then a "boundary condition equation" can be derived by assuming a pseudo-region  $k = K + 1$  of width  $h_K$  having mirror image properties of region  $K$  about  $z_{K+1}$  with corresponding symmetric flux and antisymmetric current properties of the detailed flux and current solutions. These properties in pseudo-region  $K+1$  can be related to properties of region  $K$  as a function of  $x$  in each region as

$$\mathbb{D}_{K+1}(x) = \mathbb{D}_K(1-x) \quad (3.7a)$$

$$\Lambda_{K+1}(x) = \Lambda_K(1-x) \quad (3.7b)$$

and

$$U_{K+1}(x) = U_K(1-x) \quad (3.8a)$$

$$U_{K+1}^*(x) = U_K^*(1-x) \quad (3.8b)$$

$$V_{K+1}(x) = -V_K(1-x) \quad (3.8c)$$

$$V_{K+1}^*(x) = -V_K^*(1-x) \quad (3.8d)$$

The addition of pseudo-region  $K+1$  to the summation in Eq. 3.5 results in the calculation of coefficients  $a_{K+1}$  and  $b_{K+1}$  in Eq. 3.6. Detailed definitions of the  $G \times G$  zero current coefficient matrices  $b_1$ ,  $c_1$ ,  $a_{K+1}$ , and  $b_{K+1}$ , all of which vanish for the case of zero flux boundary conditions, are also given in Appendix C.1. If symmetry is imposed

on both sides of the problem, independent variations in  $F_k^*$  for  $k = 1$  to  $K + 1$  result in a matrix problem of  $K + 1$  equations and  $K + 1$  unknowns of similar form as illustrated in Figure 2.6.

Other boundary conditions may be imposed on the approximation, including albedo and reflector boundary conditions, which specify the flux to current ratio at the boundaries. Such conditions will always lead to a variation equation of the form of Eq. 3.6, where in general the matrix coefficients  $a_2$ ,  $b_2$ ,  $b_K$ , and  $c_K$  as well as the boundary coefficients  $b_1$ ,  $c_1$ ,  $a_{K+1}$ , and  $b_{K+1}$  will have modified definitions.

A serious drawback of the approximation given by Eqs. 3.2 is that it does not allow the use of detailed flux solutions containing explicit zero flux boundary conditions. For this reason the exact solution,  $\psi_k(z) = \Phi_k(z)$  for all  $k$ , is excluded from the class of admissible trial function forms. However, such detailed solutions can be allowed by modifying the trial function forms in the boundary regions. If a detailed solution  $\psi_1(z)$  is given in the first region with the zero flux condition  $\psi_1(z_1) = 0$ , for example, the trial functions of Eqs. 3.2 could be modified for region  $k=1$  as

$$U_1(z) = \psi_1(x) \psi_1^{-1}(1) F_2 \quad (3.9a)$$

$$U_1^*(z) = \psi_1^*(x) \psi_1^{*-1}(1) F_2 \quad (3.9b)$$

$$V_1(z) = \eta_1(x) \psi_1^{-1}(1) F_2 \quad (3.9c)$$

$$V_1^*(z) = \eta_1^*(x) \psi_1^{*-1}(1) F_2^* \quad (3.9d)$$

In this way, the imposed zero flux boundary condition is explicitly given by  $\psi_1(z)$  rather than in the form of the trial function. Similar



trial functions can be given for an explicit zero flux boundary condition in the last region,  $k=K$ .

The use of these special trial functions in the boundary regions alters the definitions of the matrix coefficients  $b_2$  and  $b_K$  as given in Eq. 3.6. Detailed definitions of these coefficients when these special trial functions are used are also included in Appendix C.

Regardless of the types of boundary conditions imposed, Eq. 3.6 results in an  $N \times N$  matrix problem of the form

$$\underline{\mathbf{A}}\underline{\mathbf{F}} = \frac{1}{\lambda} \underline{\mathbf{B}}\underline{\mathbf{F}} \quad (3.10)$$

where  $\underline{\mathbf{A}}$  and  $\underline{\mathbf{B}}$  are independent of  $\lambda$ . The order  $N$  of the matrix equations is dependent upon the chosen boundary conditions, and is given for various choices in Table 3.1.

Table 3.1. Matrix Order  $N$  of the Proposed Linear Basis Function Approximations as a Function of the Imposed Boundary Conditions.

- 1 – Explicit or Implicit Zero Flux
- 2 – Symmetry

Boundary Condition Type		Matrix Order
on Left	on Right	$N$
1	1	$G \times (K-1)$
1	2	$G \times K$
2	1	$G \times K$
2	2	$G \times (K+1)$

### 3.3 The Proposed Cubic Hermite Basis Function Approximation

The proposed modal-nodal approximation method utilizing the cubic Hermite polynomials

$$\begin{aligned}
 p_1(x) &= 1 - 3x^2 + 2x^3 \\
 p_2(x) &= 3x^2 - 2x^3 \\
 p_3(x) &= -x + 2x^2 - x^3 \\
 p_4(x) &= x^2 - x^3
 \end{aligned} \tag{3.11}$$

and their negative derivatives

$$\begin{aligned}
 q_1(x) &= 6x - 6x^2 \\
 q_2(x) &= -6x + 6x^2 \\
 q_3(x) &= 1 - 4x + 3x^2 \\
 q_4(x) &= -2x + 3x^2
 \end{aligned} \tag{3.12}$$

and defined as nonzero only within each mesh region  $k$ , is given by the below regional trial function forms.  $F_k$  and  $G_k$  are again the unknown group column approximate flux and current solutions at  $z_k$  respectively, and the remaining symbols have been previously defined. As in the cubic Hermite finite element method described in section 2.3.4 of Chapter 2,  $\theta$  is an optional normalization parameter.

$$\begin{aligned}
 U_k(z) &= \psi_k(x) [ \psi_k^{-1}(0) p_1(x) F_k + \psi_k^{-1}(1) p_2(x) F_{k+1} \\
 &\quad + h_k \theta \mathbb{D}_k^{-1}(0) \psi_k^{-1}(0) p_3(x) G_k + h_k \theta \mathbb{D}_k^{-1}(1) \psi_k^{-1}(1) p_4(x) G_{k+1} ]
 \end{aligned} \tag{3.13a}$$

$$\begin{aligned}
U_k^*(z) = & \psi_k^*(x) [ \psi_k^{*-1}(0) p_1(x) F_k^* + \psi_k^{*-1}(1) p_2(x) F_{k+1}^* \\
& + h_k \theta \mathbb{D}_k^{-1}(0) \psi_k^{*-1}(0) p_3(x) G_k^* + h_k \theta \mathbb{D}_k^{-1}(1) \psi_k^{*-1}(1) p_4(x) G_{k+1}^* ] \quad (3.13b)
\end{aligned}$$

$$\begin{aligned}
V_k(z) = & \eta_k(x) [ \psi_k^{-1}(0) p_1(x) F_k + \psi_k^{-1}(1) p_2(x) F_{k+1} \\
& + h_k \theta \mathbb{D}_k^{-1}(0) \psi_k^{-1}(0) p_3(x) G_k + h_k \theta \mathbb{D}_k^{-1}(1) \psi_k^{-1}(1) p_4(x) G_{k+1} ] \\
& + \mathbb{D}_k(x) \psi_k(x) [ \frac{1}{h_k} \psi_k^{-1}(0) q_1(x) F_k + \frac{1}{h_k} \psi_k^{-1}(1) q_2(x) F_{k+1} \\
& + \theta \mathbb{D}_k^{-1}(0) \psi_k^{-1}(0) q_3(x) G_k + \theta \mathbb{D}_k^{-1}(1) \psi_k^{-1}(1) q_4(x) G_{k+1} ] \quad (3.13c)
\end{aligned}$$

$$\begin{aligned}
V_k^*(z) = & \eta_k^*(x) [ \psi_k^{*-1}(0) p_1(x) F_k^* + \psi_k^{*-1}(1) p_2(x) F_{k+1}^* \\
& + h_k \theta \mathbb{D}_k^{-1}(0) \psi_k^{*-1}(0) p_3(x) G_k^* + h_k \theta \mathbb{D}_k^{-1}(1) \psi_k^{*-1}(1) p_4(x) G_{k+1}^* ] \\
& - \mathbb{D}_k(x) \psi_k^*(x) [ \frac{1}{h_k} \psi_k^{*-1}(0) q_1(x) F_k^* + \frac{1}{h_k} \psi_k^{*-1}(1) q_2(x) F_{k+1}^* \\
& + \theta \mathbb{D}_k^{-1}(0) \psi_k^{*-1}(0) q_3(x) G_k^* + \theta \mathbb{D}_k^{-1}(1) \psi_k^{*-1}(1) q_4(x) G_{k+1}^* ] \quad (3.13d)
\end{aligned}$$

Again, for the moment,  $\psi_k(0)$  and  $\psi_k(1)$  and their adjoints are not allowed to be zero in any region.

The forms of these trial functions impose both flux and current continuity since:

$$U_k(z_k) = U_{k-1}(z_{k-1} + h_{k-1}) = F_k \quad (3.14a)$$

$$U_k^*(z_k) = U_{k-1}^*(z_{k-1} + h_{k-1}) = F_k^* \quad (3.14b)$$

$$V_k(z_k) = V_{k-1}(z_{k-1} + h_{k-1}) = G_k \quad (3.14c)$$

$$V_k^*(z_k) = V_{k-1}^*(z_{k-1} + h_{k-1}) = -G_k^* \quad (3.14d)$$

where it is assumed that at region mesh points the detailed current solutions are zero:

$$\eta_k(z_k) = \eta_k(z_k+h_k) = 0 \quad (3.15a)$$

$$\eta_k^*(z_k) = \eta_k^*(z_k+h_k) = 0 \quad (3.15b)$$

for all regions  $k=1$  to  $K$ .

Also, by comparison to Eqs. 2.31, it is evident that this approximation reduces to the cubic Hermite finite element method when the  $\psi_k$ 's are constant, and the  $\eta_k$ 's are correspondingly zero.

The insertion of the above trial function forms into variation equation 2.12a results in a lengthy equation which can be simplified to the following form:

$$\begin{aligned} & \delta F_1^{*T} [b1_1 F_1 + b2_1 G_1 + c1_1 F_2 + c2_1 G_2] \\ & + \delta G_1^{*T} [b3_1 F_1 + b4_1 G_1 + c3_1 F_2 + c4_1 G_2] \\ & + \sum_{k=2}^K \left\{ \delta F_k^{*T} [a1_k F_{k-1} + a2_k G_{k-1} + b1_k F_k + b2_k G_k + c1_k F_{k+1} + c2_k G_{k+1}] \right. \\ & \left. + \delta G_k^{*T} [a3_k F_{k-1} + a4_k G_{k-1} + b3_k F_k + b4_k G_k + c3_k F_{k+1} + c4_k G_{k+1}] \right\} \\ & + \delta F_{K+1}^{*T} [a1_{K+1} F_K + a2_{K+1} G_K + b1_{K+1} F_{K+1} + b2_{K+1} G_{K+1}] \\ & + \delta G_{K+1}^{*T} [a3_{K+1} F_K + a4_{K+1} G_K + b3_{K+1} F_{K+1} + b4_{K+1} G_{K+1}] = 0 \end{aligned} \quad (3.16)$$

where the detailed definitions of the twelve integral  $G \times G$  matrix coefficients  $\{a1_k, \dots, c4_k\}$  of the form  $A - \frac{1}{\lambda} B$  for all  $k$  are given in section 2 of Appendix C.

Boundary conditions for either zero flux or symmetry are easily imposed by setting either  $F_k$  or  $G_k$ , respectively, to zero with  $k=1$  for the conditions on the left at  $z_1$  or  $k=K+1$  for conditions on the right at  $z_{K+1}$ . The corresponding variations for  $k=1$  and  $k=K+1$  then vanish. Allowing arbitrary variations in the remaining  $F_k^*$  and  $G_k^*$  in Eq. 3.16 results in a system of  $2K$   $G \times G$  matrix equations relating  $2K$   $G$  column vector unknowns, as illustrated in Figure 2.7.

Explicit zero flux boundary conditions imposed by  $\psi_1(z_1) = 0$  or  $\psi_K(z_{K+1}) = 0$  can be incorporated into the approximation by modifying the trial function definitions in the boundary regions. The modified flux trial functions in the first region, for example, are

$$\begin{aligned} U_1(x) &= \psi_1(x) [ \psi_1^{-1}(1) F_{2+h_1} \theta \mathbb{D}_1^{-1}(0) \psi_1^{-1}(0) p_3(x) G_{1+h_1} \theta \mathbb{D}_1^{-1}(1) \psi_1^{-1}(1) p_4(x) G_2 ] \\ U_1^*(x) &= \psi_1^*(x) [ \psi_1^{*-1}(1) F_{2+h_1}^* \theta \mathbb{D}_1^{-1}(0) \psi_1^{*-1}(0) p_3(x) G_{1+h_1}^* \theta \mathbb{D}_1^{-1}(1) \psi_1^{*-1}(1) p_4(x) G_2^* ] \end{aligned} \quad (3.17)$$

where the current trial functions are again given by Fick's laws. Use of modified trial function forms of this type in the boundary regions results in  $2K$  equations with different definitions of  $c_{3_1}$ ,  $a_{2_2}$ ,  $b_{1_2}$ ,  $b_{2_2}$ , and  $b_{3_2}$  as well as  $b_{1_K}$ ,  $b_{2_K}$ ,  $c_{2_K}$ , and  $a_{3_{K+1}}$ , which are also included in Appendix C.

Other boundary condition restrictions may be imposed on this approximation, but the matrix form of the resulting difference equations will remain unchanged. Only the coefficients defined for  $k=1, 2, K$ , and  $K+1$  will in general be altered.

The matrix equations resulting from this approximation can always be written as

$$\mathbf{A} \underline{\mathbf{F}} = \frac{1}{\lambda} \mathbf{B} \underline{\mathbf{F}} \quad (3.18)$$

where  $\mathbf{A}$  and  $\mathbf{B}$  are  $(G \times 2K)$  by  $(G \times 2K)$  matrices, independent of  $\lambda$ , and  $\underline{\mathbf{F}}$  is the  $G \times K$  column vector of unknowns containing both  $F_k$  and  $G_k$  column vectors for each  $k$ . The matrix properties and solution methods of the matrix equations derived in these proposed approximation methods are discussed in the next chapter.

## Chapter 4

## NUMERICAL SOLUTION TECHNIQUES

The matrix properties of the difference equations resulting from both the proposed approximations and the finite element methods in one dimension, as well as the solution schemes used to solve these equations, are summarized in the following section. Various calculational and programming techniques used in conjunction with these approximation methods and their solution schemes are presented and discussed in section 4.2.

4.1 Solution Methods and Matrix Properties

The matrix equations which result from the approximations given in this report are of the form

$$\underline{A}\underline{F} = \frac{1}{\lambda} \underline{B}\underline{F} \quad (4.1)$$

and are solved using the fission source power iteration method without fission source renormalization.<sup>7,44</sup> The method of solution is illustrated schematically in Figure 4.1. Other definitions of the iteration eigenvalue  $\lambda^{(i)}$  can be found elsewhere.<sup>45</sup>

Figure 4.1 illustrates that an outer iteration solution scheme<sup>46</sup> is used, and that the geometry and nuclear properties of the reactor are not altered. Since the fission source is not normalized by the iteration eigenvalue during the iterations,  $\lambda^{(i)}$  converges to the effective multiplication factor,  $k_{\text{eff}}$ , of the problem. Had fission source renormalization been included, by  $\underline{S}^{(i)} = \underline{B}\underline{F}^{(i)}/\lambda^{(i-1)}$  for example, then  $\lambda^{(i)}$

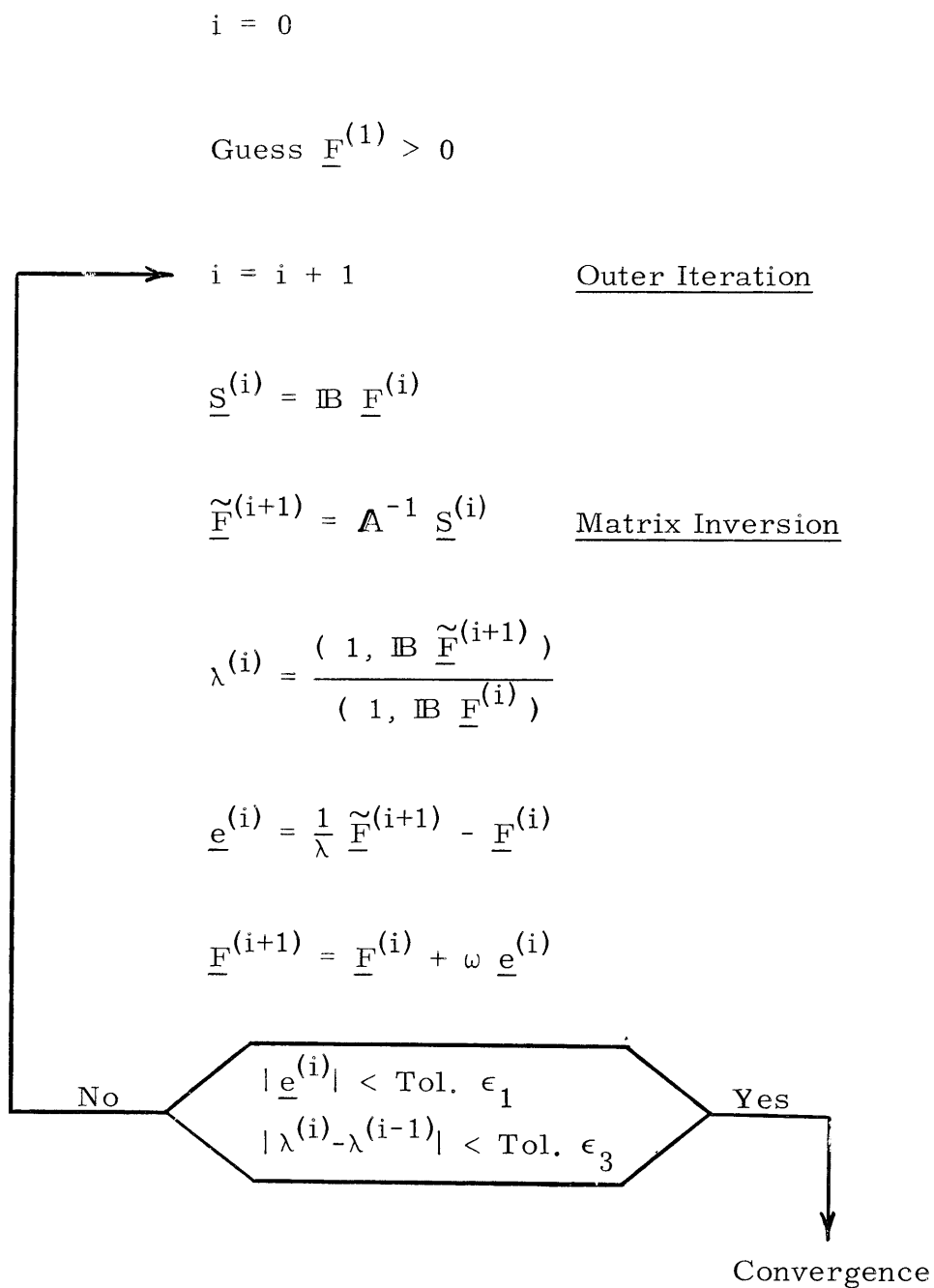


Figure 4.1. Solution of  $\underline{A} \underline{F} = \frac{1}{\lambda} \mathbb{B} \underline{F}$  Using the Fission Source Power Iteration Method Without Fission Source Renormalization.



would converge to unity. The  $k_{\text{eff}}$  of the problem would then be simply the product of all of the iteration eigenvalues.<sup>46, 47, 48</sup>

The matrix inversions required within the iteration scheme were performed directly. Although overrelaxation methods are usually employed only in iterative matrix inversion schemes (or inner iterations),<sup>49</sup> an overrelaxation parameter  $w$ ,  $1 \leq w \leq 2$ , is available in the outer iteration in order to hasten the convergence of the solution vector.

The power method is very appealing to neutron diffusion flux calculations because it converges to the largest or fundamental eigenvalue  $|\lambda_0| > |\lambda_i|$ ,  $i \neq 0$ , and the corresponding eigenvector  $\underline{F}_0$  of the given matrix problem. The convergence rate is governed by the dominance ratio, defined as  $\max_{i \neq 0} |\lambda_i/\lambda_0|$ , in such a way that smaller ratios result in faster convergence. Although the power method will always converge when  $\lambda_0$  is positive and unique, specific matrix properties of  $\mathbf{A}$  and  $\mathbf{B}$  are sufficient but not always necessary to insure convergence to a positive  $k_{\text{eff}}$  and everywhere positive neutron flux approximation.<sup>50</sup>

In many problems the order of  $\mathbf{A}$  may be quite large, and solution methods which require the direct inversion of  $\mathbf{A}$  may not be practical. For the purposes of this report, as in most multigroup calculational schemes, neutron up-scattering will not be permitted. The inversion of  $\mathbf{A}$  is then performed by successive group-iteration techniques.

The equations given in Eq. 4.1 have been defined as ordered first by spatial indexing followed by group indexing within each spatial index. It is convenient to reorder these equations so that they are ordered first by group indexing followed by spatial indexing within each group.

After reordering, Eq. 4.1 can be written as

$$(\mathbb{L} + \mathbb{M}) \underline{\mathbf{F}} = \mathbf{T} \underline{\mathbf{F}} + \frac{1}{\lambda} \mathbb{B} \underline{\mathbf{F}} \quad (4.2a)$$

where

$$\mathbf{A} = \mathbb{L} + \mathbb{M} - \mathbf{T} \quad (4.2b)$$

and:  $\mathbb{L}$ , the stiffness matrix, results from leakage;  $\mathbb{M}$ , the mass matrix, results from absorption;  $\mathbf{T}$  is the group-to-group scattering transfer matrix; and  $\mathbb{B}$  is the fission source production matrix.

Assuming  $K$  spatial unknowns in each of the  $G$  groups,  $\mathbb{L}$  and  $\mathbb{M}$  are  $G \times K$  block diagonal matrices composed of  $G$   $K \times K$  matrices  $\mathbb{L}_g$  and  $\mathbb{M}_g$  of the form

$$\mathbb{L} = \text{Diag}[\mathbb{L}_1, \dots, \mathbb{L}_G] \quad (4.3a)$$

$$\mathbb{M} = \text{Diag}[\mathbb{M}_1, \dots, \mathbb{M}_G] \quad (4.3b)$$

and  $\mathbf{T}$  and  $\mathbb{B}$  are in general full block matrices composed of  $G^2$   $K \times K$  matrices  $\mathbf{T}_{gg'}$  and  $\mathbb{B}_{gg'}$ , respectively. Since only downscattering is permitted,  $\mathbf{T}$  becomes lower block triangular;  $\mathbf{T}_{gg'} = 0$  whenever  $g' \geq g$ . The matrix inversion,  $\underline{\tilde{\mathbf{F}}}^{(i+1)} = \mathbf{A}^{-1} \underline{\mathbf{S}}^{(i)}$ , can then be solved for the GK unknowns

$$\underline{\tilde{\mathbf{F}}}^{(i+1)} = \text{Col} \left[ \underline{\tilde{\mathbf{F}}}_1^{(i+1)} \quad \dots \quad \underline{\tilde{\mathbf{F}}}_G^{(i+1)} \right] \quad (4.4)$$

by solving successively the following system of group equations:

$$\left[ \begin{array}{l} \text{Do for } g = 1 \text{ to } G: \\ \left| \begin{array}{l} \underline{\mathbf{S}}_g^{(i)} = \sum_{g'=1}^G (\mathbf{T}_{gg'} + \mathbb{B}_{gg'}) \underline{\mathbf{F}}_{g'}^{(k)} \\ \text{where: } k = \begin{cases} i+1; & g' < g \\ i; & g' \geq g \end{cases} \\ \underline{\tilde{\mathbf{F}}}_g^{(i+1)} = (\mathbb{L}_g + \mathbb{M}_g)^{-1} \underline{\mathbf{S}}_g^{(i)} \end{array} \right. \end{array} \right. \quad (4.5)$$

where the updating of the group fission source by the iteration index  $k = i + 1$  for  $g' < g$  generally enables a faster rate of convergence of the outer iteration than  $k = i$ .

The desirable convergence properties of a positive eigenvalue and everywhere positive flux solution when using the above group iteration method depend upon the properties of the  $K \times K$  spatial matrices for each group  $g$ :  $\mathbb{L}_g$ ;  $\mathbb{M}_g$ ; and  $\mathbb{T}_{gg'}$  and  $\mathbb{B}_{gg'}$ , for  $g' = 1$  to  $G$ . Using the Perron-Frohenius theorem,<sup>50</sup> it can be shown that if  $\mathbb{T}_{gg'}$  and  $\mathbb{B}_{gg'}$  are all nonnegative for each group  $g$  and  $\mathbb{L}_g$  and  $\mathbb{M}_g$  are both Stieltjes or S-type matrices<sup>†</sup> for each group  $g$ , then the power method will converge to a positive eigenvalue,  $\lambda_0 > 0$ , and a corresponding positive eigenvector,  $\underline{F}_0 > 0$ . These matrix properties naturally depend upon the form of the spatial approximations employed and generally differ for different approximation schemes.

The conventional finite difference approximation has become popular because the spatial matrices which arise from its use exhibit these desirable properties regardless of the size of the mesh regions chosen. The spatial matrices resulting from the linear finite element method, however, are known to exhibit these properties only if the mesh size is restricted by

$$h_k \leq \max_{g=1 \text{ to } G} \{ \sqrt{6} \ell_{g,k} \} \quad (4.6)$$

where  $\ell_{g,k}$  is the diffusion length,  $\ell_{g,k}^2 = D_{g,k} / \Sigma_{g,k}$ , for group  $g$  in mesh region  $k$ . The spatial matrices resulting from the cubic Hermite

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<sup>†</sup>A Stieltjes matrix is a real, irreducible, positive definite matrix with nonpositive off-diagonal elements. <sup>7, 50</sup>

finite element method do not exhibit these "desirable" characteristics. Since the eigenvector  $\underline{F}$  contains current as well as flux unknowns, convergence to an all-positive solution vector is not desirable.

The properties of the spatial matrices for each group  $g$  resulting from the proposed approximation methods can be found by generalizing the proposed trial function forms in each group as

$$U_g(z) = \sum_{k=1}^K \left[ \underline{P}_{g,k}^{-T}(x) \underline{F}_{g,k} + \underline{P}_{g,k}^{+T} \underline{F}_{g,k+1} \right] \quad (4.7a)$$

$$U_g^*(z) = \sum_{k=1}^K \left[ \underline{P}_{g,k}^{-*T}(x) \underline{F}_{g,k}^* + \underline{P}_{g,k}^{+*T} \underline{F}_{g,k+1}^* \right] \quad (4.7b)$$

where the  $\underline{F}_{g,k}$  are in general column vectors of length  $N$  given by:

$$\underline{F}_{g,k} = F_{g,k} \quad (N = 1) \quad (4.8a)$$

for the linear basis function approximations, and

$$\underline{F}_{g,k} = \text{Col} [ F_{g,k}, G_{g,k} ] \quad (N = 2) \quad (4.8b)$$

for the cubic Hermite basis function approximations. Similar definitions hold for the  $\underline{F}_{g,k}^*$ . The  $\underline{P}_{g,k}^{\pm}(x)$  are column vectors of length  $N$  whose elements are functions of  $z$  (or  $x$ ) defined as nonzero only within region  $k$  which provide the basis for the approximations. The definitions of the  $\underline{P}_{g,k}^{\pm}(x)$  for the proposed approximations are given as follows:

$N = 1$ ; Linear Basis Functions:

$$\underline{P}_{g,k}^{-}(x) = (1-x) \psi_{g,k}^{-1}(0) \psi_{g,k}(x) \quad (4.9a)$$

$$\underline{P}_{g,k}^{+}(x) = x \psi_{g,k}^{-1}(1) \psi_{g,k}(x) \quad (4.9b)$$

$N = 2$ ; Cubic Hermite Basis Functions:

$$\underline{P}_{g,k}^{-}(x) = \text{Col}[p_1(x)\psi_{g,k}^{-1}(0)\psi_{g,k}(x), h_k \theta p_3(x)D_{g,k}^{-1}(0)\psi_{g,k}^{-1}(0)\psi_{g,k}(x)] \quad (4.10a)$$

$$\underline{P}_{g,k}^{+}(x) = \text{Col}[p_2(x)\psi_{g,k}^{-1}(1)\psi_{g,k}(x), h_k \theta p_4(x)D_{g,k}^{-1}(1)\psi_{g,k}^{-1}(1)\psi_{g,k}(x)] \quad (4.10b)$$

where  $\psi_{g,k}(x)$  and  $D_{g,k}(x)$  are the detailed flux solutions and diffusion coefficients of group  $g$ , the polynomials  $p_1(x)$  through  $p_4(x)$  are defined in Eqs. 3.11, and  $0 \leq x \leq 1$  within each region  $k$ . Similar definitions hold for the  $\underline{P}_{g,k}^{\pm*}(x)$ .

Equations 4.7 can be written in matrix form as

$$\underline{U}_g(z) = \underline{IP}_g(x)\underline{F}_g \quad (4.11a)$$

$$\underline{U}_g^*(z) = \underline{IP}_g^*(x)\underline{F}_g^* \quad (4.11b)$$

where:

$$\underline{F}_g = \text{Col}(F_{g,1}, \dots, F_{g,K+1}) \quad (4.12a)$$

and  $\underline{IP}_g(x)$  is the  $K$  by  $N(K+1)$  matrix defined by

$$\underline{IP}_g(x) = \begin{bmatrix} \underline{P}_{g,1}^{-T}(x) & \underline{P}_{g,1}^{+T}(x) & & 0 \\ & & & \\ & & & \\ & & & \\ & 0 & & \underline{P}_{g,K}^{-T}(x) & \underline{P}_{g,K}^{+T}(x) \end{bmatrix} \quad (4.12b)$$

$\underline{IP}_g^*(x)$  is defined similarly. Insertion of these trial function forms into variation equation 2.12b for each group  $g$  results in

$$\delta \underline{F}_g^{*T} \int_K \left( \dot{\underline{IP}}_g^{*T} \underline{D}_g \dot{\underline{IP}}_g + \underline{IP}_g^{*T} \sum_{g'=1}^G \underline{\Lambda}_{gg'} \underline{IP}_{g'} \underline{F}_{g'} \right) dz = 0 \quad (4.13)$$

where  $\underline{D}_g$  and  $\underline{\Lambda}_{gg'}$  are  $K \times K$  diagonal matrices of the form

$$\underline{D}_g = \underline{D}_g(x) = \text{Diag} [D_{g,1}(x), \dots, D_{g,K}(x)] \quad (4.14a)$$

and

$$\underline{\Lambda}_{gg'} = \underline{\Lambda}_{gg'}(x) = \text{Diag} [\Lambda_{gg',1}(x), \dots, \Lambda_{gg',K}(x)] \quad (4.14b)$$

The quantity  $\dot{\underline{IP}}_g$  represents the derivative of  $\underline{IP}_g(z)$  with respect to  $z$ , and the integration over  $K$  denotes integration over the entire range of  $z$ ;  $z_1 \leq z \leq z_{K+1}$ .

$D_{g,k}$  and  $\Lambda_{gg',k}$ ,

$$\Lambda_{gg',k} = \left( \Sigma_{tg} - \Sigma_{gg'} - \frac{1}{\lambda} \chi_g \nu \Sigma_{fg'} \right)_{\text{in region } k} \quad (4.14c)$$

are the group material constants in mesh region  $k$ , and are usually dependent on  $x$ .  $\underline{\Lambda}_{gg'}$  can thus be conveniently expanded as

$$\underline{\Lambda}_{gg'} = \underline{\Lambda}_g^A - \underline{\Lambda}_{gg'}^S - \frac{1}{\lambda} \underline{\Lambda}_{gg'}^F \quad (4.14d)$$

Allowing arbitrary variations in each element of  $\underline{F}_g^*$  for each group  $g$  in Eq. 4.13 results in the matrix equations

$$(\underline{L}_g + \underline{M}_g) \underline{F}_g = \sum_{g'=1}^G \left( \underline{T}_{gg'} + \frac{1}{\lambda} \underline{B}_{gg'} \right) \underline{F}_{g'}; \quad g = 1 \text{ to } G \quad (4.15)$$

as described in Eqs. 4.1 through 4.5, where:

$$\underline{L}_g = \int_K \dot{\underline{IP}}_g^{*T} \underline{D}_g \dot{\underline{IP}}_g dz \quad (4.16a)$$

$$\underline{M}_g = \int_K \underline{IP}_g^{*T} \underline{\Lambda}_g^A \underline{IP}_g dz \quad (4.16b)$$

$$\mathbb{T}_{gg'} = \int_K \mathbb{P}_g^{*\text{T}} \mathbf{\Lambda}_{gg'}^S \mathbb{P}_{g'} dz \quad (4.16c)$$

$$\mathbb{B}_{gg'} = \int_K \mathbb{P}_g^{*\text{T}} \mathbf{\Lambda}_{gg'}^F \mathbb{P}_{g'} dz \quad (4.16d)$$

These matrices are  $N(K+1)$  by  $N(K+1)$  block tridiagonal of similar form whose  $N \times N$  submatrices are integrals of  $N \times N$  dyads. The  $k^{\text{th}}$  row of the product  $\mathbf{\Lambda}_{gg'}^F \mathbb{P}_{g'}$  is, for example:

$$\begin{aligned} [\mathbf{\Lambda}_{gg'}^F \mathbb{P}_{g'}]_k &= \left\{ \int_0^1 h_{k-1} \Lambda_{gg',k-1}(x) \underline{P}_{g,k-1}^{+*}(x) \underline{P}_{g',k-1}^{-\text{T}}(x) dx \right\} F_{g',k-1} \\ &+ \left\{ \int_0^1 h_{k-1} \Lambda_{gg',k-1}(x) \underline{P}_{g,k-1}^{+*}(x) \underline{P}_{g',k-1}^{+\text{T}}(x) dx \right. \\ &+ \left. \int_0^1 h_k \Lambda_{gg',k}(x) \underline{P}_{g,k}^{-*}(x) \underline{P}_{g',k}^{-\text{T}}(x) dx \right\} F_{g',k} \\ &+ \left\{ \int_0^1 h_k \Lambda_{gg',k}(x) \underline{P}_{g,k}^{-*}(x) \underline{P}_{g',k}^{+\text{T}}(x) dx \right\} F_{g',k+1} \quad (4.17) \end{aligned}$$

These matrix relationships allow presentation of the following matrix properties.

Theorem 1:  $\mathbb{L}_g$  and  $\mathbb{M}_g$  are guaranteed to be positive definite whenever the detailed weighting functions  $\psi_{g,k}^*(z)$  have a similar shape to that of the detailed flux solutions  $\psi_{g,k}(z)$ , as given by:

$$\psi_{g,k}^*(z) = C_{g,k} \psi_{g,k}(z) \quad (4.18)$$

where  $C_{g,k}$  is a positive constant for each energy group  $g$  and each region  $k$ .

Proof: Under these conditions,  $\underline{P}_{g,k}^{*\pm} = \mathbb{C}_{g,k} \underline{P}_{g,k}^{\pm}$ ;

hence,

$$\mathbb{I}\mathbb{P}_g^* = \mathbb{C}_g \mathbb{I}\mathbb{P}_g \quad (4.19a)$$

where

$$\mathbb{C}_g = \text{Diag}(C_{g,1}, \dots, C_{g,K}) \quad (4.19b)$$

First consider  $\mathbb{I}\mathbb{M}_g$ . Given any arbitrary constant nonzero vector  $\underline{q}$ ,

$$\underline{q}^T \mathbb{I}\mathbb{M}_g \underline{q} = \int_K (\mathbb{I}\mathbb{P}_g \underline{q})^T \mathbb{C}_g^T \underline{\Lambda}_g^A \mathbb{I}\mathbb{P}_g \underline{q} \, dz \quad (4.20)$$

and since  $\mathbb{C}_g$  and  $\underline{\Lambda}_g^A$  are both positive block diagonal matrices with diagonal submatrices, their product can be factored into

$$\mathbb{C}_g^T \underline{\Lambda}_g^A = \left[ (\underline{\Lambda}_g^A \mathbb{C}_g)^{\frac{1}{2}} \right]^T (\mathbb{C}_g \underline{\Lambda}_g^A)^{\frac{1}{2}} \quad (4.21)$$

Therefore,

$$\underline{q}^T \mathbb{I}\mathbb{M}_g \underline{q} = \int_K \left[ (\mathbb{C}_g \underline{\Lambda}_g^A)^{\frac{1}{2}} \mathbb{I}\mathbb{P}_g \underline{q} \right]^T \left[ (\mathbb{C}_g \underline{\Lambda}_g^A)^{\frac{1}{2}} \mathbb{I}\mathbb{P}_g \underline{q} \right] \, dz \quad (4.22a)$$

$$= \int_K \mathbb{R}^T \mathbb{R} \, dz \quad (4.22b)$$

which is always greater than zero for arbitrary nonzero  $\underline{q}$ . Hence, by definition,  $\mathbb{I}\mathbb{M}_g$  is positive definite. A similar proof holds for  $\mathbb{I}\mathbb{L}_g$  using Eq. 4.16a.

The following corollaries immediately result.

Corollary 1: If Rayleigh-Ritz Galerkin weighting,  $\underline{U}^* = \underline{U}$ , is used in the approximation, then  $\mathbb{I}\mathbb{L}_g$  and  $\mathbb{I}\mathbb{M}_g$  are positive definite.



Corollary 2:  $\mathbb{L}_g$  and  $\mathbb{M}_g$  resulting from the finite element methods using linear and cubic Hermite basis functions are both positive definite.

Corollary 3: If  $\mathbb{L}_g$  and  $\mathbb{M}_g$  are positive definite, then so is the matrix  $(\mathbb{L}_g + \mathbb{M}_g)$ . These three matrices are then also symmetric.

It is also interesting to note the properties of  $\mathbb{L}_g$  and  $\mathbb{M}_g$  for cases of symmetry; that is, when the material properties and detailed flux solutions are symmetric about the center of each coarse mesh region  $k$ . Such symmetry occurs in regular repeating reactor geometries, and is denoted by:

$$D_{g,k}(x) = D_{g,k}(1-x) \quad (4.23a)$$

and

$$\Lambda_{g,k}(x) = \Lambda_{g,k}(1-x) \quad (4.23b)$$

Hence,

$$\psi_{g,k}(x) = \psi_{g,k}(1-x) \quad (4.23c)$$

and

$$\eta_{g,k}(x) = -\eta_{g,k}(1-x) \quad (4.23d)$$

and similarly for the weighting fluxes and currents. Under such conditions, the  $\underline{P}_{g,k}^+(x)$  and  $\underline{P}_{g,k}^-(x)$  support functions can be found by inspection of Eqs. 4.9 and 4.10 to obey the following symmetries:

For  $N = 1$ :

$$\underline{P}_{g,k}^{\mp}(x) = \underline{P}_{g,k}^{\pm}(1-x) \quad (4.24a)$$

$$\dot{\underline{P}}_{g,k}^{\mp}(x) = -\dot{\underline{P}}_{g,k}^{\pm}(1-x) \quad (4.24b)$$

For  $N = 2$ :

$$\underline{P}_{g,k}^{\mp}(x) = \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix} \underline{P}_{g,k}^{\pm}(1-x) \quad (4.24c)$$

$$\dot{\underline{P}}_{g,k}^{\mp}(x) = - \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix} \dot{\underline{P}}_{g,k}^{\pm}(1-x) \quad (4.24d)$$

where the following symmetries of the polynomials defined in Eqs. 3.11 have been used:

$$\begin{aligned} p_1(x) &= p_2(1-x) \\ p_3(x) &= -p_4(1-x) \\ q_1(x) &= -q_2(1-x) \\ q_3(x) &= q_4(1-x) \end{aligned} \quad (4.25)$$

Similar identities with identical signs hold for the weighting quantities  $\underline{P}_{g,k}^{\pm*}(x)$ .

Theorem 2: For cases of symmetry, as given above, the matrices  $\mathbb{L}_g$ ,  $\mathbb{M}_g$ ,  $\mathbb{T}_{gg'}$ , and  $\mathbb{B}_{gg'}$  are all symmetric regardless of the relation of  $\psi_{g,k}^*(x)$  to  $\psi_{g,k}(x)$ .

Proof: Referring to Eq. 4.17,  $\mathbb{L}_g$ , for example, is symmetric only if

$$\int_0^1 h_k D_{g,k}(x) \dot{\underline{P}}_{g,k}^{+*}(x) \dot{\underline{P}}_{g,k}^{-T}(x) dx = \int_0^1 h_k D_{g,k}(x) \dot{\underline{P}}_{g,k}^{-*}(x) \dot{\underline{P}}_{g,k}^{+T}(x) dx \quad (4.26)$$

This can be shown for any  $N$  by changing variables in one of the integrals from  $x$  to  $1-x'$  and using the symmetry properties of Eqs. 4.23 and 4.24. Similar proofs hold for the other matrices.

It is unfortunate that the above symmetry conditions do not allow direct proof that  $\mathbb{L}_g$  and  $\mathbb{M}_g$  have positive diagonal elements and are also diagonally dominant (for at least one row) for arbitrary positive and symmetric detailed flux solutions. Under such conditions,  $\mathbb{L}_g$  and  $\mathbb{M}_g$  would then be positive definite, since they are block tridiagonal with nonzero diagonal elements and hence irreducible. Instead, these conditions can be used to obtain a set of algebraic equations which, for completely arbitrary detailed flux solutions, must be satisfied in order that  $\mathbb{L}_g$  and  $\mathbb{M}_g$  be positive definite.

The requirement that  $\mathbb{L}_g$  and  $\mathbb{M}_g$  be positive definite is useful only in the inversion of  $(\mathbb{L}_g + \mathbb{M}_g)$ . Although the inversion can always be performed using Gaussian elimination techniques, the property of positive definiteness allows the use of Cholesky's method, discussed in the next section, which is faster and requires less computer storage.

#### 4.2 Calculational and Programming Techniques

Calculation of the one-dimensional subassembly detailed fluxes, currents, and adjoint solutions, as well as detailed "exact" or reference solutions, were performed using the program REF2G described in section 1 of Appendix D. Assuming subassembly  $k$  to be divided into  $N$  homogeneous intervals at nodes  $t_i$  and of width  $hs_i$ , the program uses fine mesh linear finite element approximations to calculate the detailed flux solutions for each group. Omitting group subscripts, the detailed flux solution for each group in subassembly  $k$  is represented by a set of  $N+1$  points

$$\psi_k(x) = \{ \psi_{k,i} : i = 1, N+1 \} \quad (4.27)$$

where  $\psi_k(x)$  is linear between points. Detailed group current solutions  $\eta_k(x)$  are represented by a set of  $N$  points

$$\eta_k(x) = \{ \eta_{k,i} : i = 1, N \} \quad (4.28)$$

which are found from the converged flux solutions by Fick's law

$$\eta_{k,i} = \frac{1}{hs_i} D_{k,i} [ \psi_{k,i} - \psi_{k,i+1} ]; \quad i = 1, N \quad (4.29)$$

where  $D_{k,i}$  is the diffusion constant homogeneous in interval  $i$ .

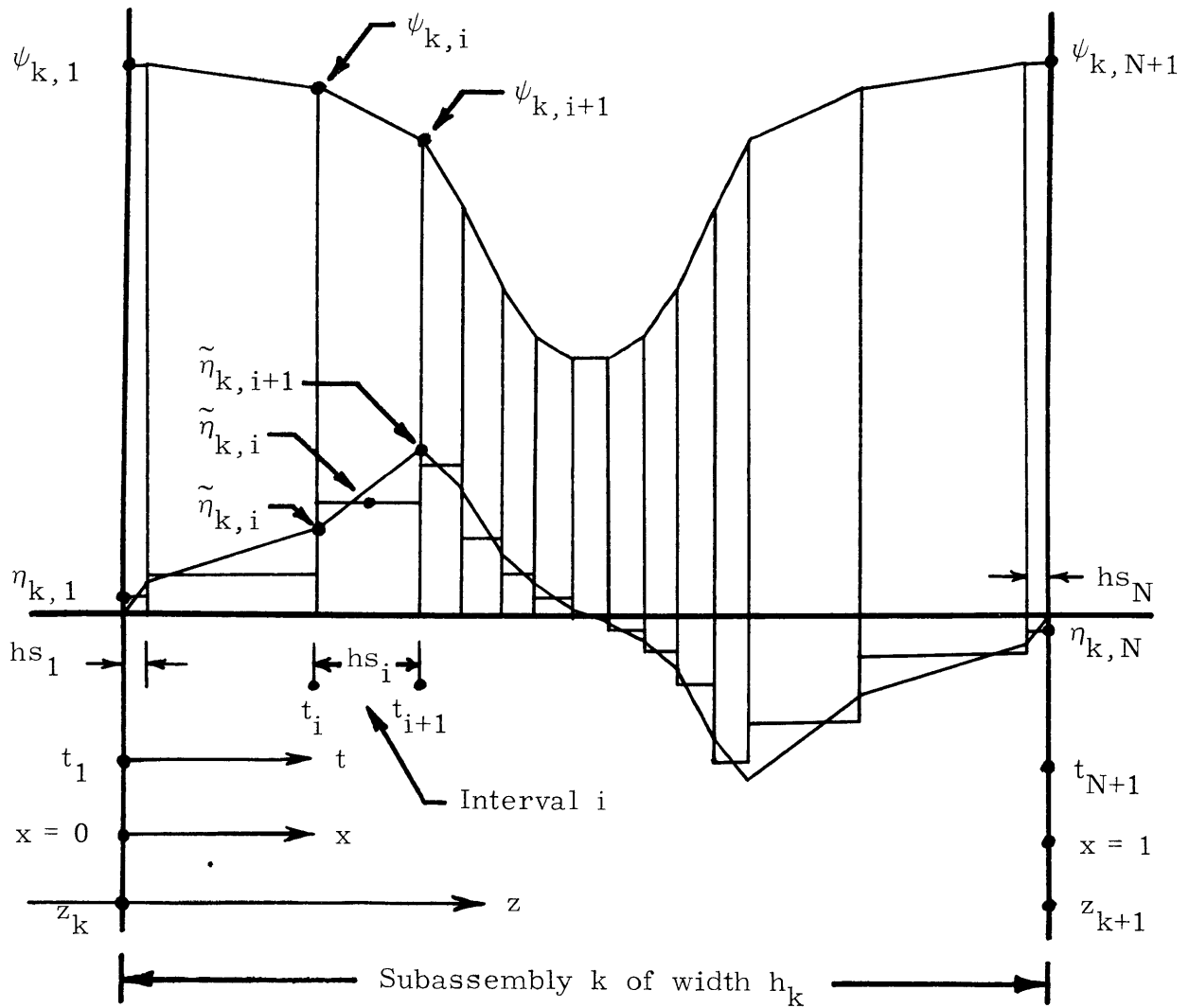
$\eta_k(x)$  is of constant value,  $\eta_{k,i}$ , within each interval. The forms of these solutions are illustrated in Figure 4.2.

In order to approximate the symmetry boundary conditions imposed on the detailed subassembly flux solutions, small intervals  $hs_1$  and  $hs_N$  are defined at the edges of each subassembly. The detailed current solutions can then be made to have zero boundary values by setting  $\eta_{k,1}$  and  $\eta_{k,N}$  to zero. However, since the currents in each interval are defined as inversely proportional to the mesh size, the calculated boundary currents using this scheme may not be small enough to be negligible.

Explicit zero current boundary conditions can be imposed on the detailed current solutions by transforming the above discontinuous current  $\eta_k(x)$  into a continuous current solution  $\tilde{\eta}_k(x)$  represented by a set of  $N+1$  points

$$\tilde{\eta}_k(x) = \{ \tilde{\eta}_{k,i} : i = 1 \text{ to } N+1 \} \quad (4.30)$$

where  $\tilde{\eta}_k(x)$  is linear between points, as also illustrated in Figure 4.2. By seeking to minimize the mean square error between  $\eta_k(x)$  and  $\tilde{\eta}_k(x)$  within each interval  $i$ , variational techniques yield the following set of  $N-1$  equations for each group:



$$x = \frac{1}{h_k} (z - z_k)$$

$0 \leq x \leq 1$ ; for each subassembly  $k$ ;  $k = 1$  to  $K$ .

$$y = \frac{1}{hs_i} (t - t_i)$$

$0 \leq y \leq 1$ ; for each interval  $i$  in subassembly  $k$ ;  $i = 1$  to  $N$ .

Figure 4.2. Subassembly Notations and Detailed Solutions

$$\begin{aligned}
& \left(\frac{1}{6} h s_{i-1}\right) \tilde{\eta}_{k,i-1} + \frac{1}{3} \left(h s_{i-1} + h s_i\right) \tilde{\eta}_{k,i} + \left(\frac{1}{6} h s_i\right) \tilde{\eta}_{k,i+1} \\
& = \left(\frac{1}{2} h s_{i-1}\right) \eta_{k,i-1} + \left(\frac{1}{2} h s_i\right) \eta_{k,i} ; \quad i = 2 \text{ to } N
\end{aligned} \tag{4.31}$$

These equations, given the  $\eta_{k,i}$  from Eq. 4.29, are easily solved for the  $\tilde{\eta}_{k,i}$ ,  $i = 2$  to  $N$ , where  $\tilde{\eta}_{k,1}$  and  $\tilde{\eta}_{k,N+1}$  are set to zero. Both forms of the detailed currents,  $\eta_k(x)$  and  $\tilde{\eta}_k(x)$ , are allowed for use in the proposed approximation methods.

The proposed methods using linear and cubic Hermite basis functions have been programmed into computer codes LINEAR and CUBIC which are described respectively in sections 2 and 3 of Appendix D.

The matrix elements required for use in the approximation methods are integrals of products of subassembly detailed solutions and polynomial functions. These integrals are calculated, for each index  $k$ , from the basic integral unit

$$\text{BIU}_k = \int_0^1 f_k(x) g_k(x) C_k(x) x^n h_k dx \tag{4.32}$$

where the functions  $f_k(x)$  and  $g_k(x)$  represent flux and/or current solutions for same or different groups. These functions may be either constant within each interval

$$f_k(x) = \{f_{k,i} : i = 1 \text{ to } N\} \tag{4.33}$$

or of linear form within each interval

$$f_k(x) = \{(1-y)f_{k,i} + yf_{k,i+1} : i = 1 \text{ to } N\} \tag{4.34}$$

where  $y = \frac{1}{h s_i} (t - t_i)$ , as defined in Figure 4.2.  $C_k(x)$  represents a

group nuclear constant which is homogeneous in each interval

$$C_k(x) = \{ C_{k,i} : i = 1 \text{ to } N \} \quad (4.35)$$

and  $n$  is a positive integer exponent in the range  $0 \leq n \leq 6$ . Since the following remarks concern only subassembly  $k$ , the index  $k$  is dropped for simplicity.

The basic integral unit can be broken into integrals over each interval by transforming variables from  $x$  to  $y$ . The result

$$\text{BIU}_k = \frac{1}{h_k^n} \sum_{i=1}^N h s_i \int_0^1 f_i(y) g_i(y) C_i(t_i + h s_i y)^n dy \quad (4.36)$$

can be integrated analytically by expanding  $(t_i + h s_i y)^n$  into a binomial series. The results of these integrations for any  $n$  depend only on the given forms of  $f(x)$  and  $g(x)$ , and are summarized in Table 4.1.

The coarse mesh flux-weighting homogenization calculations were performed using the above basic integral unit with  $n = 0$ . In these calculations a linear form of  $f(x)$ , representing the detailed subassembly flux solutions from REF2G, and a constant value of  $g(x) = 1$  were used.

Once the elements of the matrices of the approximation methods have been formed, considerable computer storage can be saved by collapsing the sparse band-structured matrices into full matrix form using row index transformations. In this way, a  $N \times N$  tridiagonal matrix  $\mathbb{L}$  resulting from the use of linear basis functions can be stored as the  $N \times 3$  matrix  $\mathbb{L}'$  by

$$(\mathbb{L}')_{ik} = (\mathbb{L})_{ij} \quad (4.37)$$

where  $k = j + 2 - i$ , and  $k$  values outside  $1 \leq k \leq 3$  are omitted.

Table. 4.1. Calculation of the Basic Integral Unit in Subassembly k.

1. Constant  $f(x)$  and constant  $g(x)$  in each interval  $i$ :

$$\text{BIU}_k = \frac{1}{h_k^n} \sum_{i=1}^N \left\{ \text{hs}_i f_i g_i C_i \sum_{\ell=0}^n b_{\ell}^n t_i^{n-\ell} \text{hs}_i^{\ell} \frac{1}{(\ell+1)} \right\}$$

2. Linear  $f(x)$  and constant  $g(x)$  in each interval  $i$ :

$$\text{BIU}_k = \frac{1}{h_k^n} \sum_{i=1}^N \left\{ \text{hs}_i g_i C_i \sum_{\ell=0}^n b_{\ell}^n t_i^{n-\ell} \text{hs}_i^{\ell} \left[ \frac{f_i}{(\ell+1)(\ell+2)} + \frac{f_{i+1}}{(\ell+2)} \right] \right\}$$

3. Linear  $f(x)$  and linear  $g(x)$  in each interval  $i$ :

$$\text{BIU}_k = \frac{1}{h_k^n} \sum_{i=1}^N \left\{ \text{hs}_i C_i \sum_{\ell=0}^n b_{\ell}^n t_i^{n-\ell} \text{hs}_i^{\ell} \left[ \frac{2f_i g_i}{(\ell+1)(\ell+2)(\ell+3)} + \frac{f_i g_{i+1} + f_{i+1} g_i}{(\ell+2)(\ell+3)} + \frac{f_{i+1} g_{i+1}}{(\ell+3)} \right] \right\}$$

where

$$b_{\ell}^n = \frac{n!}{\ell! (n-\ell)!}$$

is the binomial series coefficient.



Similarly, a  $N \times N$  matrix  $\mathbb{L}$  of half-band width equal to three, which results from the use of cubic Hermite basis functions, can be stored as the  $N \times 6$  matrix  $\mathbb{L}'$ , as given above, where in this case:

$$k = j - i + 3 + \begin{cases} 1 & \text{for } i \text{ odd} \\ 0 & \text{for } i \text{ even} \end{cases} \quad (4.38)$$

and  $k$  values outside  $1 \leq k \leq 6$  are omitted.

In those cases where the mass and stiffness matrices are both positive definite, Cholesky's method of matrix factorization,

$$\mathbb{L} = \mathbb{G} \mathbb{G}^T \quad (4.39)$$

where  $\mathbb{L}$  is positive definite and  $\mathbb{G}$  is lower triangular, can be used to solve the matrix inversion for each group in the power method. The matrix elements  $g_{ij} = (\mathbb{G})_{ij}$  are calculated from the elements  $\ell_{ij} = (\mathbb{L})_{ij}$  by the following algorithm:<sup>51</sup>

$$\left[ \begin{array}{l} \text{For each } j = 1 \text{ to } N: \\ \quad g_{jj} = \left[ \ell_{jj} - \sum_{k=1}^{j-1} g_{jk}^2 \right]^{\frac{1}{2}} \\ \quad \text{For each } i = j+1 \text{ to } N: \\ \quad \quad g_{ij} = \left[ \ell_{ij} - \sum_{k=1}^{j-1} g_{ik} g_{jk} \right] / g_{jj} \end{array} \right. \quad (4.40)$$

Similar algorithms of a more complex form are used in the computer codes in conjunction with the matrix collapsing schemes given above.

## Chapter 5

## NUMERICAL RESULTS

5.1 Nuclear Constants and Subassembly Geometry

The effectiveness and accuracy of the proposed approximation methods were examined using one-dimensional, one- and two-group reactor configurations composed of representative PWR fuel subassemblies. Four separate subassemblies with identical geometry but different number constants are considered. Each subassembly is represented as an 18-cm, homogeneous fuel region of low, medium, or high enrichment, surrounding a 1-cm centrally located absorption rod or water channel. Two-group regional nuclear constants used to represent such PWR subassembly geometries are given in Table 5.1,<sup>52</sup> where all fission neutrons are assumed to be born in the fast group. These constants were collapsed into representative one-group constants using the standard infinite medium group reduction procedure for two groups:

$$\langle \Sigma \rangle_{1G} = \frac{1}{(1+\alpha)} (\Sigma_1 + \alpha \Sigma_2) \quad (5.1)$$

where  $\Sigma_1$  and  $\Sigma_2$  are macroscopic cross sections for the fast and thermal groups, respectively, and  $\alpha$  is the infinite medium thermal to fast flux ratio.<sup>53</sup> The resulting one-group regional constants for the fuel and rod regions are given in Table 5.2, where the flux ratios of the three fuel regions have been averaged in order to collapse the absorption rod constants.

Table 5.1. Representative Two-Group, 18-cm, PWR  
Subassembly Regional Nuclear Constants.  
 $\chi_1 = 1.0$ ;  $\chi_2 = 0.0$ .

Region Material	$\Sigma_T$	$\nu\Sigma_f$	D	$\Sigma_{21}$
Fuel A: Low w/o	.0259	.00485	1.396	.0179
	.0532	.0636	.388	
Fuel B: Medium w/o	.0260	.00553	1.397	.0172
	.0710	.102	.389	
Fuel C: High w/o	.0261	.00659	1.399	.0168
	.0832	.129	.387	
Absorption Rod	.0452	0.0	1.0	0.0
	.959	0.0	1.0	
Water	.0383	0.0	1.63	.0380
	.0108	0.0	.275	

Fast group constants appear first for each region material, followed by thermal group constants. Fission neutrons are assumed to be born in the fast group only.

Table 5.2. Representative One-Group, 18-cm, PWR  
Subassembly Regional Nuclear Constants.

Region Material	$\Sigma_T$	$\nu\Sigma_f$	D
Fuel A: Low w/o	.0329	.0199	1.14
Fuel B: Medium w/o	.0348	.0244	1.20
Fuel C: High w/o	.0357	.0272	1.23
Absorption Rod	.235	0.0	1.0
Water	.0136	0.0	.414

Four subassembly configurations, labeled A through D, were used in the one- and two-group test configurations, and are illustrated in Figure 5.1. Subassemblies labeled A, B, and C contain homogeneous fuel of low, medium, and high enrichment, respectively, surrounding the 1-cm absorption rod while subassembly D contains low enriched homogeneous fuel surrounding a 1-cm water channel.

## 5.2 Subassembly Detailed Solutions and Homogenized Nuclear Constants

The detailed flux and current solutions for each subassembly were found using the computer code REF2G with symmetry boundary conditions and a 68-mesh region per subassembly geometry as indicated in Figure 5.2. The resulting one-group detailed flux solutions for each subassembly are shown in Figure 5.3. The resulting two-group detailed flux and adjoint flux solutions for each subassembly are shown in Figures 5.4 and 5.5, respectively.

Homogeneous subassembly group constants for use in the finite element approximations were found by flux weighting the group cross sections in each subassembly by the corresponding subassembly detailed group flux solutions. The resulting homogenized one-group constants for each subassembly are given in Table 5.3, and the resulting homogenized two-group constants are given in Table 5.4. The results of homogenizing the diffusion coefficient as the transport cross section,  $1/\langle 1/D \rangle$ , as well as by direct homogenization,  $\langle D \rangle$ , are included in the tables. The results of both schemes were found to differ at most by only 2%. The directly homogenized diffusion coefficients,  $\langle D \rangle$ , were used in the finite element approximations.

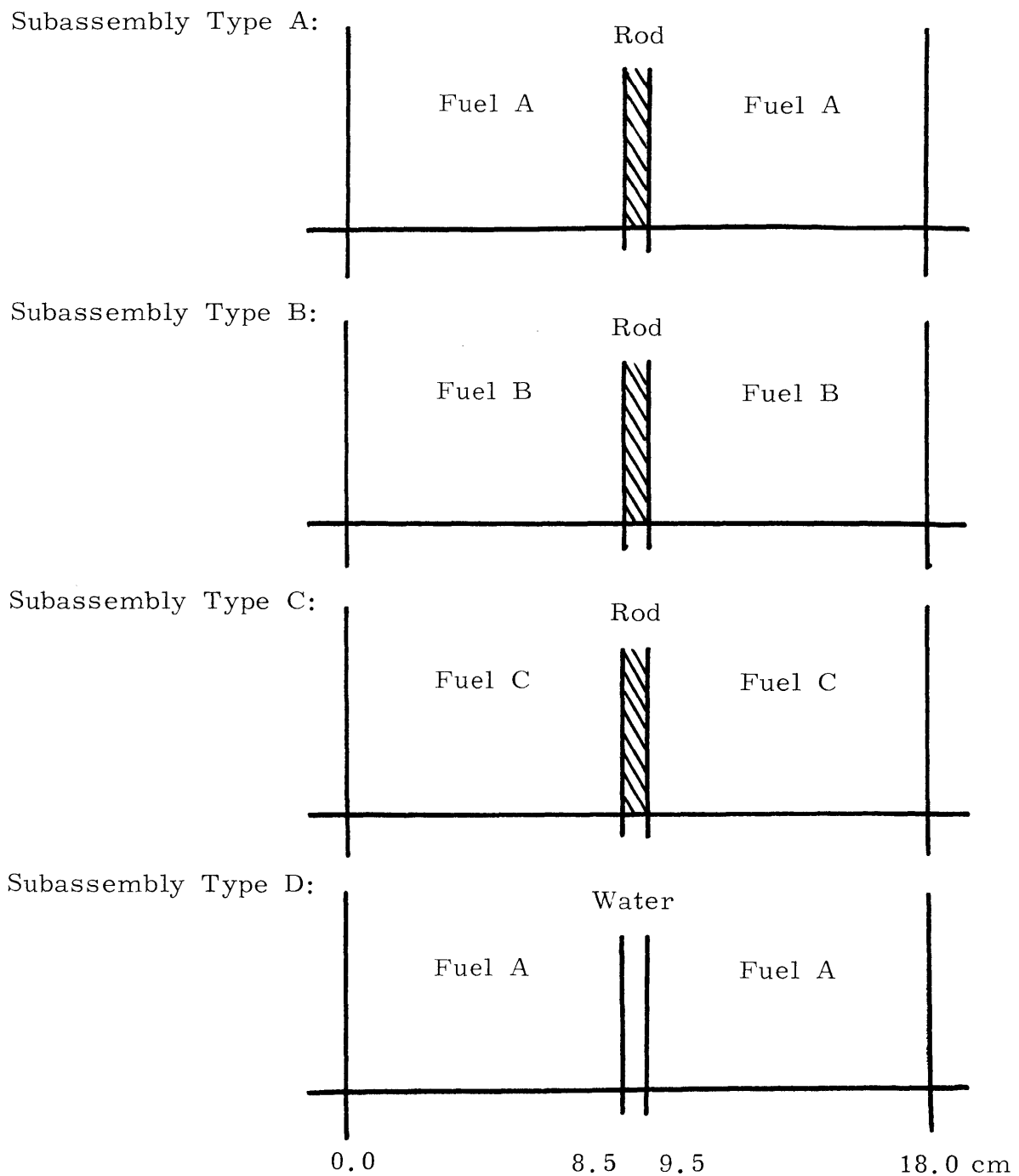
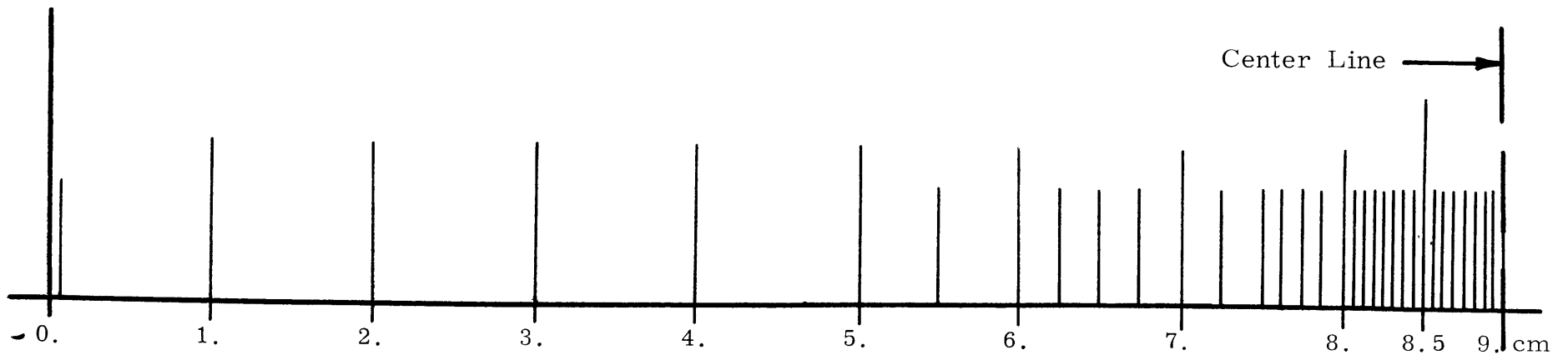


Figure 5.1. Subassembly Configuration Geometries



Symmetric Partitioning:

$$1 \left( \frac{1}{16} \text{ cm} \right) + 1 \left( \frac{15}{16} \right) + 4 (1) + 2 \left( \frac{1}{2} \right) + 6 \left( \frac{1}{4} \right) + 4 \left( \frac{1}{8} \right) + 8 \left( \frac{1}{16} \right) + 8 \left( \frac{1}{16} \right)$$

Figure 5.2. Mesh Geometry in Half a Subassembly.  
Detailed flux and current solution calculations use this 68 intervals/subassembly geometry in each subassembly type.

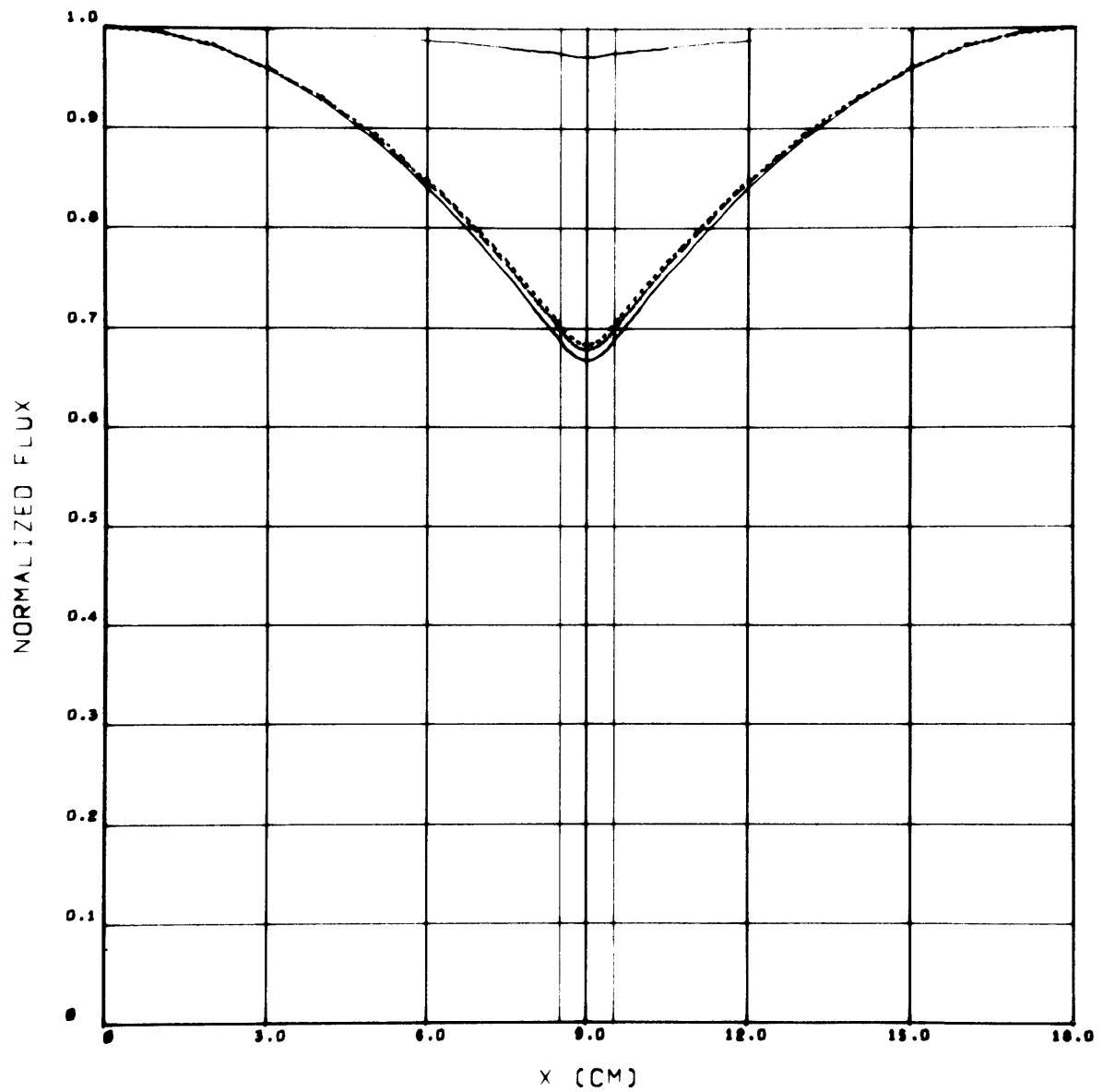


Figure 5.3. Subassembly Detailed Flux Solutions  
for the One-Group Case

Subassembly Type A ————— (lower curve)  
 Subassembly Type B - - - - -  
 Subassembly Type C .....  
 Subassembly Type D ————— (upper curve)

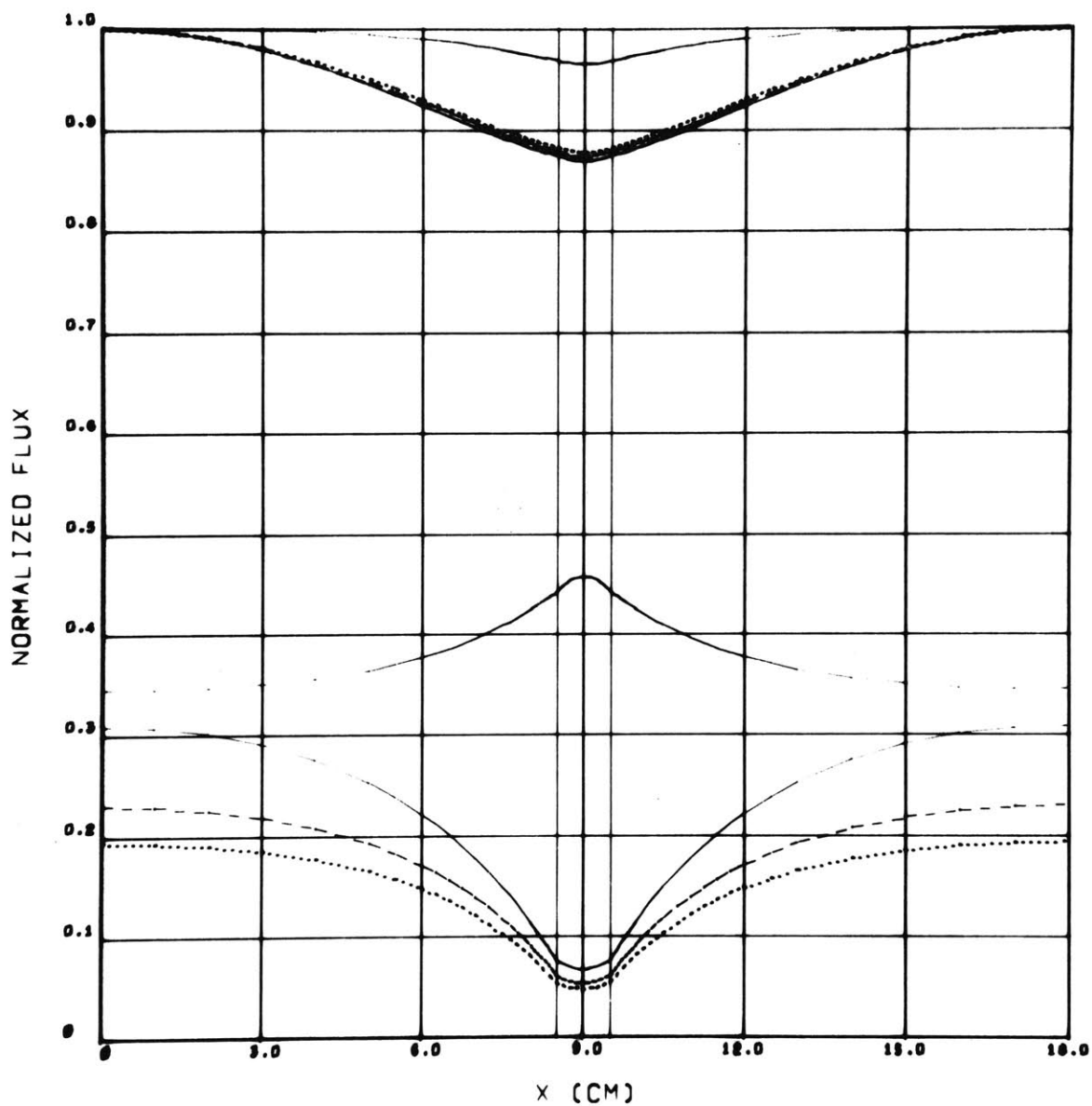


Figure 5.4. Subassembly Detailed Flux Solutions for the Two-Group Case

The fluxes are normalized by fast flux values so that the thermal fluxes appear in the lower portion of the figure.

Subassembly Type A ————— (lower curves)  
 Subassembly Type B - - - - -  
 Subassembly Type C .....  
 Subassembly Type D ————— (upper curves)



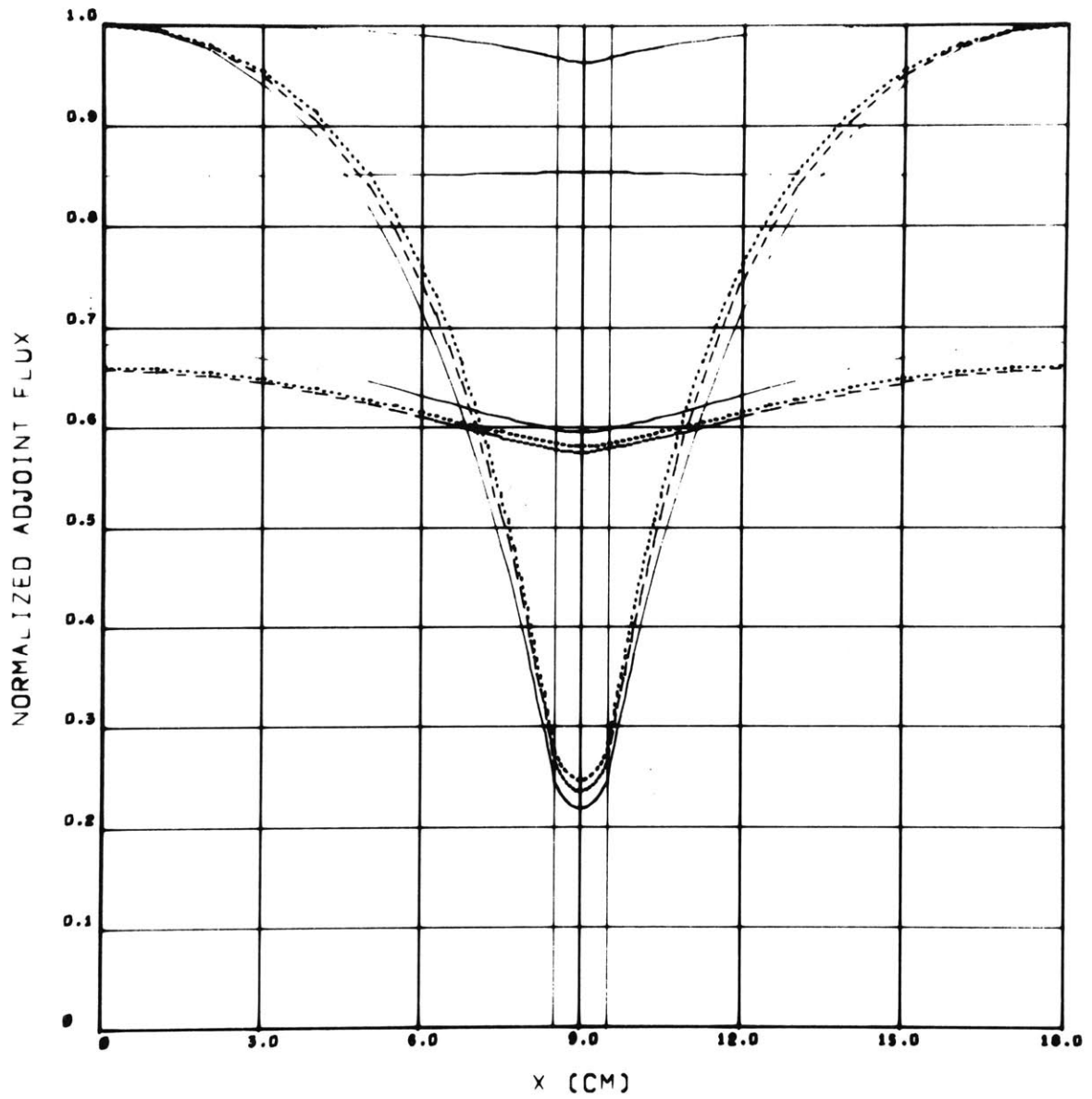


Figure 5.5. Subassembly Detailed Adjoint Flux Solutions for the Two-Group Case

Subassembly Type A ————— (lower curves)  
 Subassembly Type B - - - - -  
 Subassembly Type C ·······  
 Subassembly Type D ————— (upper curves)

Table 5.3. Homogenized Subassembly One-Group  
Nuclear Constants.

Subassembly	$\langle \Sigma_T \rangle$	$\langle \nu \Sigma_f \rangle$	$\langle D \rangle$	$1/\langle 1/D \rangle$
A	.04149392	.01905379	1.134047	1.133253
B	.04341140	.02335046	1.191397	1.189765
C	.04431869	.02602374	1.220054	1.217887
D	.03184731	.01881458	1.100401	1.040479

Table 5.4. Homogenized Subassembly Two-Group  
Nuclear Constants.  
 $\chi_1 = 1.0$ ;  $\chi_2 = 0.0$ .

Sub-assembly	$\langle \Sigma_T \rangle$	$\langle \nu \Sigma_f \rangle$	$\langle D \rangle$	$\langle \Sigma_{21} \rangle$	$1/\langle 1/D \rangle$
A	.02688787	.004601752	1.379526	.01698379	1.371911
	.06812834	.06255182	.3980863		.3919533
B	.02698495	.005246314	1.379480	.01631765	1.37185
	.08647802	.1002221	.3996499		.3931874
C	.02708207	.006251160	1.379433	.01593619	1.371787
	.09893614	.1266822	.3980142		.391310
D	.02657213	.004587110	1.412467	.0189895	1.410790
	.05034835	.05932253	.3804001		.3775656

Fast group constants appear first for each subassembly,  
followed by thermal group constants.

Before applying the proposed approximation methods to complex reactor geometries, test runs were performed in order to evaluate the differences between using either flux or adjoint flux weighting, and using current solutions of either constant or linear form in each subassembly interval, as described in section 4.2. The test problem consisted of three consecutive Type A subassemblies with symmetry boundary conditions imposed on each end so that the converged eigenvalue  $\lambda$  ( $k_{\text{eff}}$ ) should be identical to that of the detailed flux solution of subassembly A. Entire subassemblies were chosen as the mesh regions so that the proposed synthesis methods should converge to flux values of unity, and current values of zero. The numerical results of these tests for the one- and two-group cases are summarized in Table 5.5. Although the choice of weighting function did not influence the results for either approximation, use of current solutions of the linear form enables better eigenvalue accuracy. In addition, the results when using currents of linear form converged to flux values of unity and current values of zero, as expected, while results using the constant current form produced errors of about 0.5% in the converged flux and 0.01% in the converged current at interior points. Although the difference in accuracy between the use of these different current forms is small, the small flux and current errors resulting from the use of the constant current form may lead to larger errors in larger and more complex problems. For the above reasons, the linear current form was used in the following case studies. Adjoint weighting was also used. Although the use of adjoint weighting has not been shown to guarantee the success of Cholesky's method in the numerical solution scheme, no difficulties with its use were ever encountered.

Table 5.5. Test Results Using Three Consecutive Type A Subassemblies.

$$\% \lambda = [ (\lambda_{\text{Sub. A}} - \lambda_{\text{Conv.}}) / \lambda_{\text{Sub. A}} ] \times 100\%.$$

Synthesis Approximation	Weighting Function	Form of Currents	Converged $\lambda$	$\% \lambda$
ONE GROUP: $\lambda_{\text{Sub. A}} = 0.459194$				
Linear	FLUX	Constant	.459363	-.036%
Linear	FLUX	Linear	.459254	-.013%
Cubic	FLUX	Constant	.459363	-.036%
Cubic	FLUX	Linear	.459254	-.013%
TWO GROUPS: $\lambda_{\text{Sub. A}} = 0.751095$				
Linear	FLUX	Linear	.751284	-.025%
Linear	ADJOINT	Constant	.7513818	-.038%
Linear	ADJOINT	Linear	.751284	-.025%
Cubic	FLUX	Linear	.751284	-.025%
Cubic	ADJOINT	Constant	.7513818	-.038%
Cubic	ADJOINT	Linear	.751284	-.025%

### 5.3 Case Studies and Results

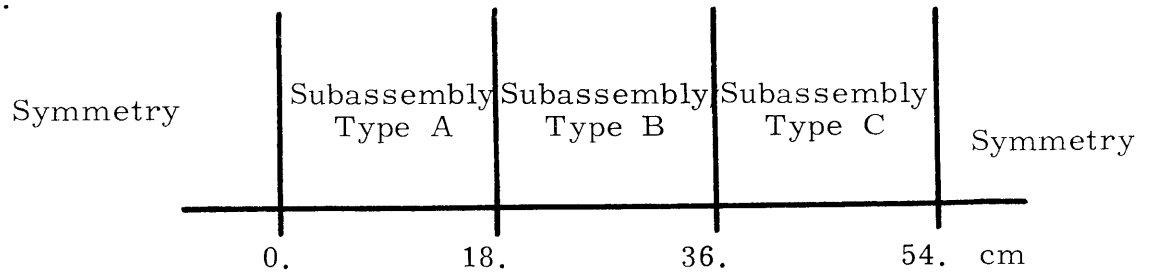
Four one-dimensional reactor configurations, each made up of different combinations of types of subassemblies, are considered in the case studies below. One-group calculations were performed only for the first case, and two-group calculations were performed in all cases. Entire 18-cm subassemblies were used as coarse mesh regions in each case, while the effect of using half-subassembly mesh regions was also included in Case 1. The geometry and subassembly configurations of the case studies are shown in Figure 5.6.

Three separate approximation methods were used to calculate converged detailed flux solutions for comparison in each case. They are:

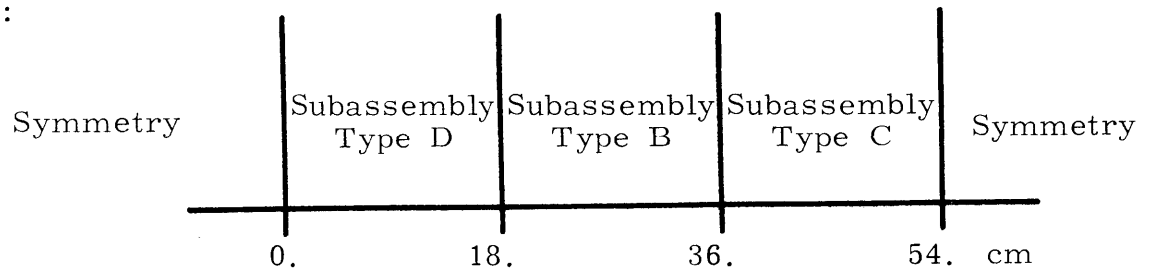
1. The proposed approximation methods using heterogeneous nuclear constants and subassembly detailed flux solutions for coarse mesh solutions.
2. The finite element methods using subassembly homogenized nuclear constants for coarse mesh solutions.
3. The linear finite element method for fine mesh reference solutions.

Calculations of both the proposed approximation and the coarse mesh finite element method using linear basis functions were performed using program LINEAR, while the corresponding cubic Hermite basis function approximations were performed using program CUBIC. The fine mesh reference solutions were calculated using program REF2G, and the results of these approximations were compared and analyzed by program ANALYZE. Descriptions of these programs are given in Appendix D.

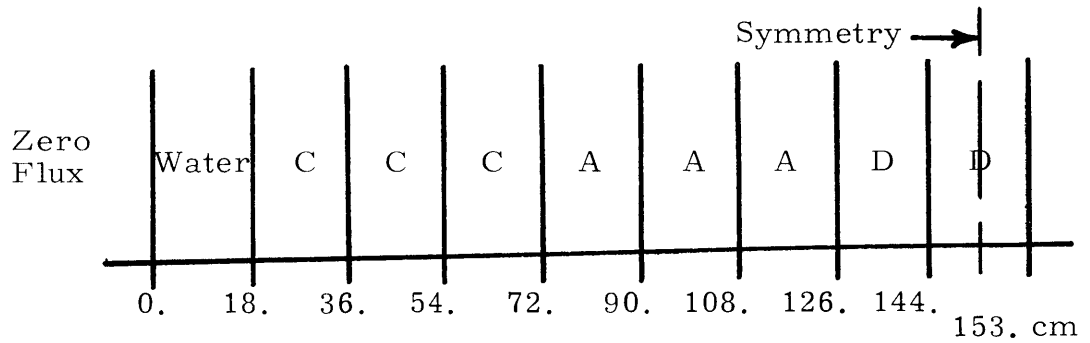
Case 1 :



Case 2 :



Case 3 :



Case 4 :

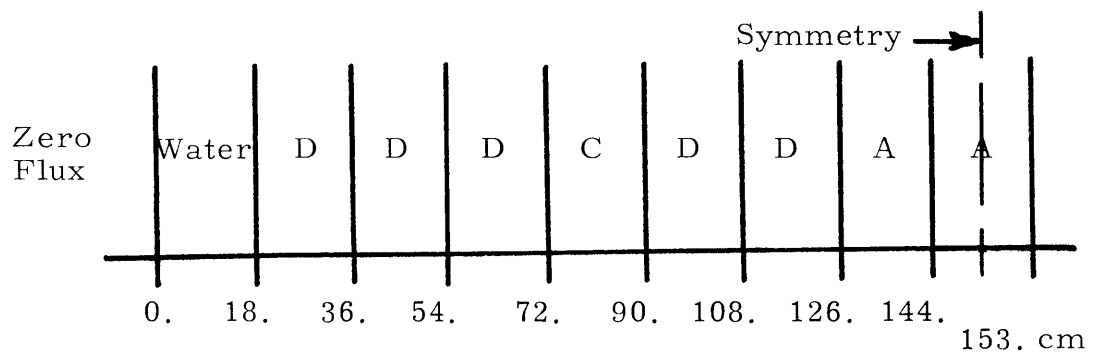


Figure 5.6. Geometry of the Four Case Studies Composed of Types of Subassemblies

The results of each case study are divided into the two approximation method categories as defined below.

1. The linear basis function approximation
  - A. Linear FEM:  
(The linear finite element method using homogenized coarse mesh nuclear constants)
  - B. Linear Synth:  
(The proposed approximation method using heterogeneous coarse mesh nuclear constants and detailed coarse mesh solutions)
2. The cubic Hermite basis function approximations
  - A. Cubic FEM
  - B. Cubic Synth

The results of the approximations in each category are compared to the reference solution by examining:

1. The converged eigenvalues  $\lambda(k_{\text{eff}})$  and their percent normalized eigenvalue error,
 
$$\% \lambda = (\lambda_{\text{Ref}} - \lambda_{\text{Conv}}) / \lambda_{\text{Ref}} \times 100\%$$
2. Composite graphs of the converged detailed group flux solutions  $U_g(z)$  normalized to equivalent power levels
3. The fractional normalized power levels  $P(k)$  calculated for each 18-cm subassembly  $k$  by

$$P(k) = \frac{\int_{\text{Sub. } k} \sum_{g=1}^G \nu \Sigma_{fg}(z) U_g(z) dz}{\int_{\text{All Subs.}} \sum_{g=1}^G \nu \Sigma_{fg}(z) U_g(z) dz} \quad (5.2)$$

and their percent normalized errors

$$\%P(k) = (P(k)_{\text{Ref}} - P(k)_{\text{Conv}}) / P(k)_{\text{Ref}} \times 100\% \quad (5.3)$$

5.3.1. Case 1: Three different subassemblies of Types A, B, and C with symmetry boundary conditions.

The graphical results of the one-group approximation methods for this case are shown in Figures 5.7 and 5.8, while the results of the two-group approximation methods are presented in Figures 5.9 - 5.12. Only the coarse mesh boundaries are labeled in the figures, which indicates that entire 18-cm subassemblies were used as the coarse mesh regions. Two-group results using only half-subassemblies as the coarse mesh regions are shown in Figures 5.13 - 5.16. The reference solutions were calculated using 150-mesh regions, as defined by the symmetric partitioning

$$5(1 \text{ cm}) + 4(.5 \text{ cm}) + 4(.25 \text{ cm}) + 4(.125 \text{ cm}) + 8(.0625 \text{ cm})$$

in each of the three subassemblies. The converged approximation eigenvalues and fractional normalized subassembly power levels for the one- and two-group calculations are summarized in Table 5.6. The fractional powers,  $P(k)$ , for each subassembly are listed in the reference solution column, while the percent errors,  $\%P(k)$  are listed in the approximation columns.



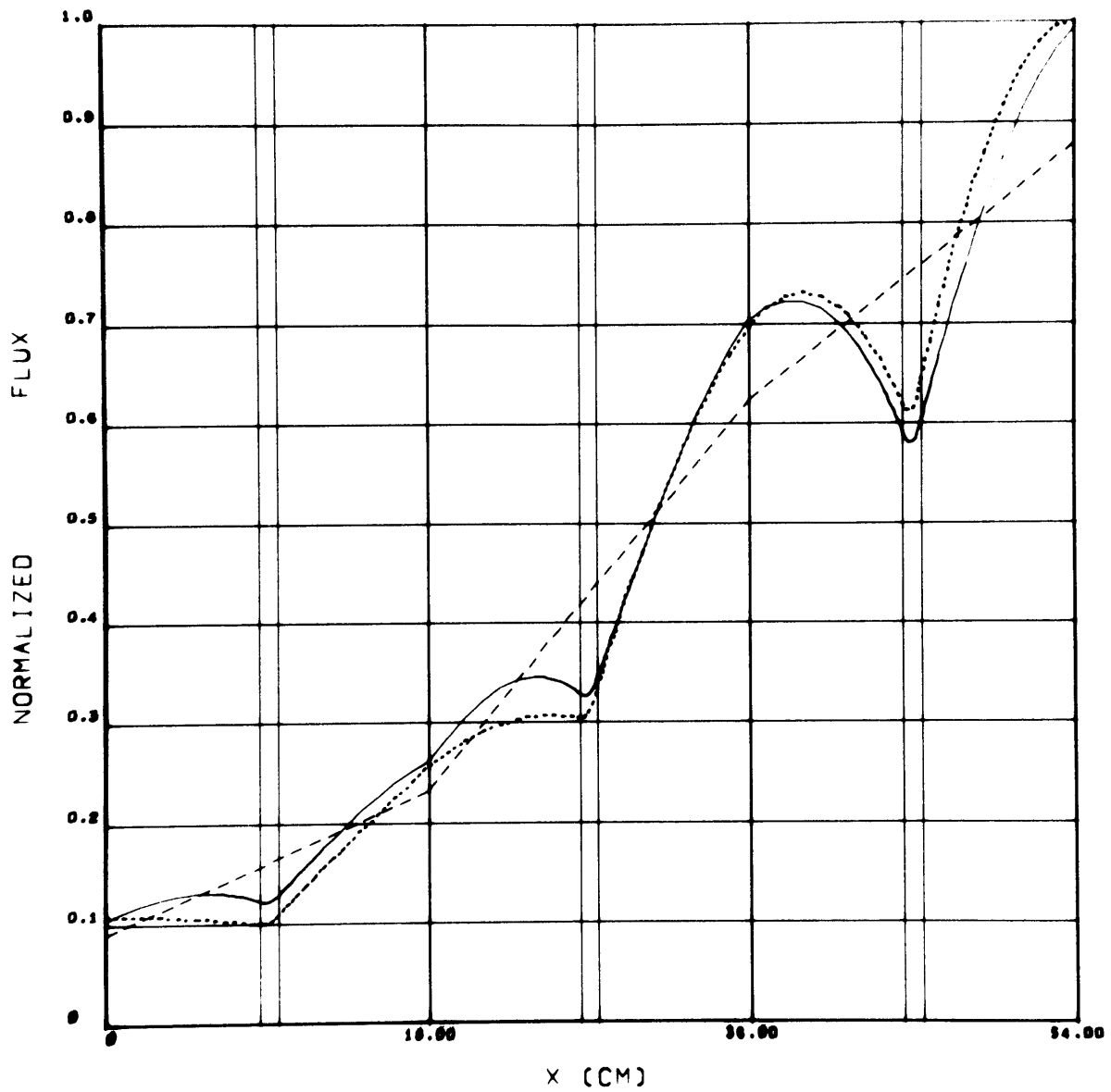


Figure 5.7. Case 1: One-Group Results Using Linear Basis Function Approximations and 18-cm Coarse Mesh Regions

<u>Method</u>		<u><math>\lambda</math></u>
Reference	.....	.559045
Linear FEM	-----	.556943
Linear Synth	—————	.557154

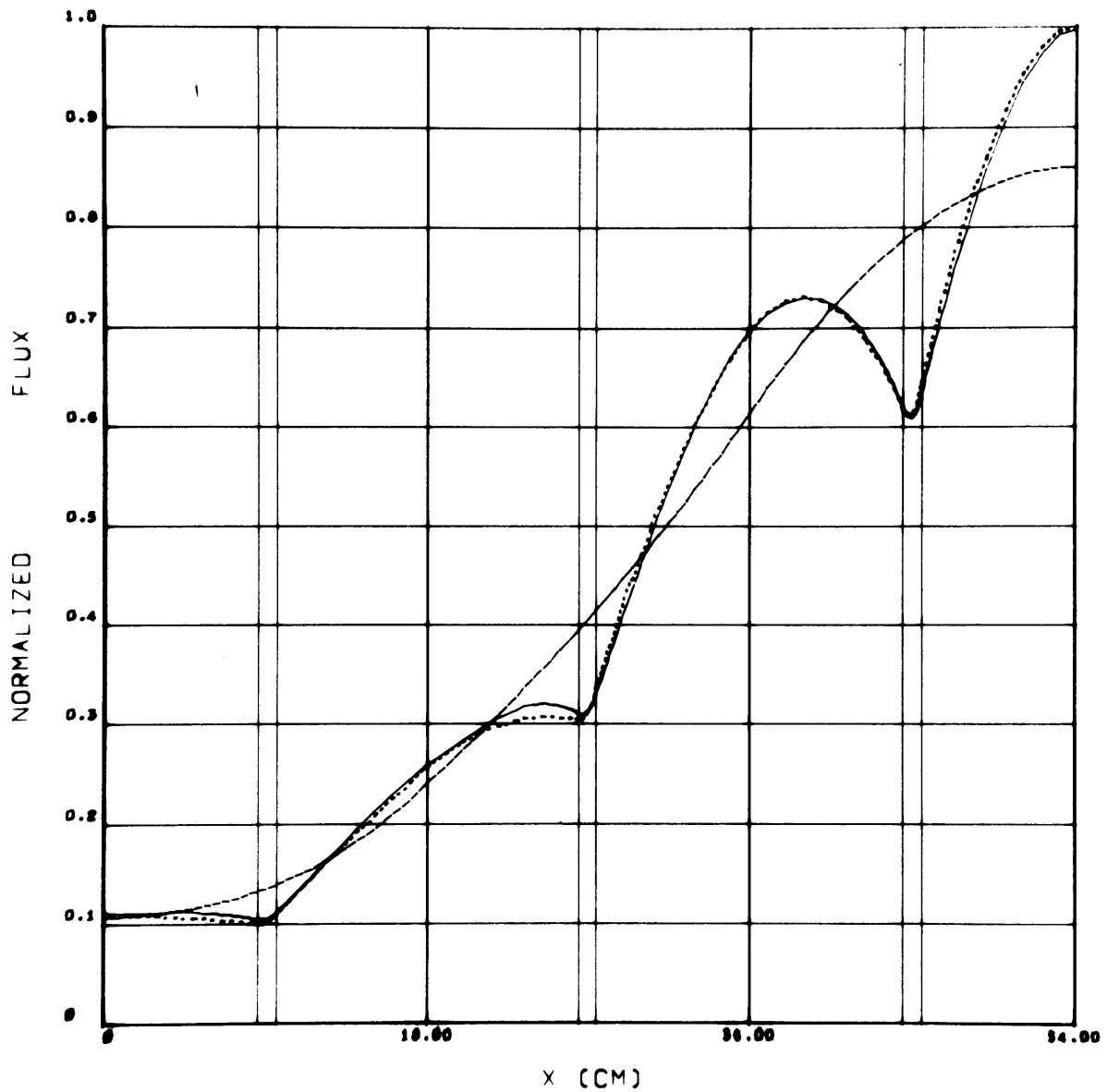


Figure 5.8. Case 1: One-Group Results Using Cubic Hermite Basis Function Approximations and 18-cm Coarse Mesh Regions

<u>Method</u>		<u><math>\lambda</math></u>
Reference	.....	.559045
Cubic FEM	-----	.558647
Cubic Synth	—————	.558761

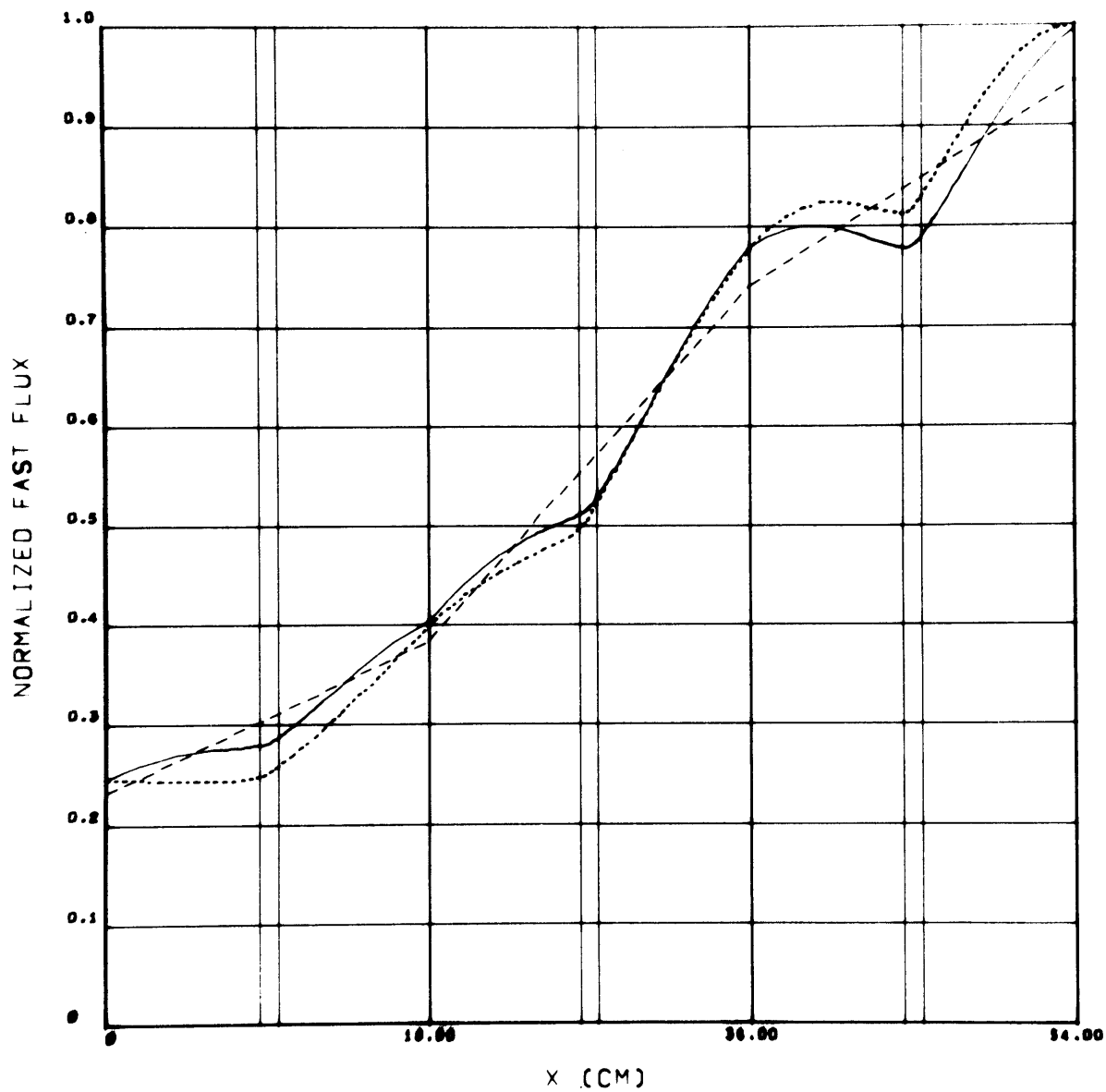


Figure 5.9. Case 1: Two-Group Fast Results Using Linear Basis Function Approximations and 18-cm Coarse Mesh Regions

Method		$\lambda$
Reference	.....	.917267
Linear FEM	-----	.914489
Linear Synth	—————	.915221

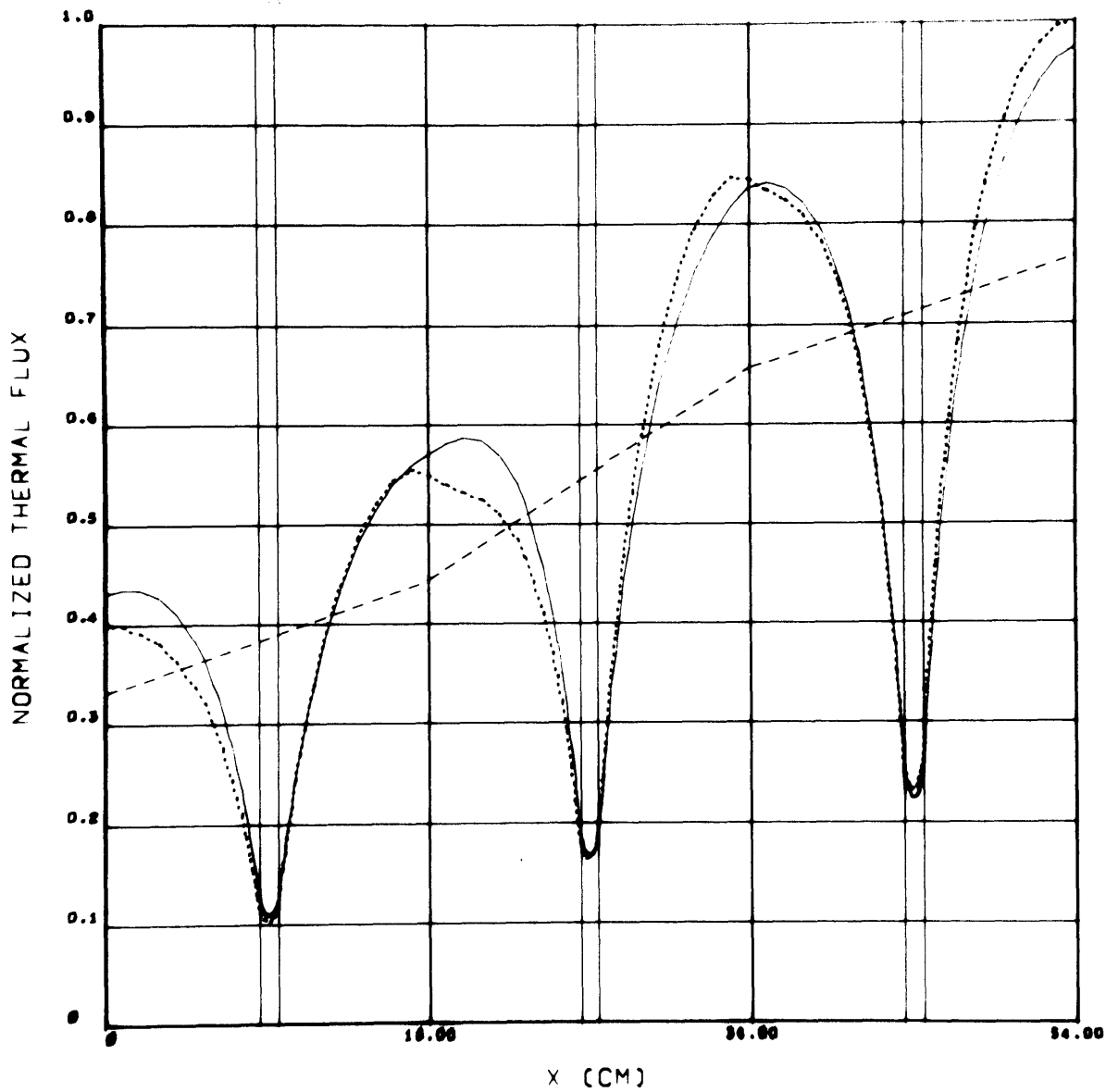


Figure 5.10. Case 1: Two-Group Thermal Results Using Linear Basis Function Approximations and 18-cm Coarse Mesh Regions

Method		$\lambda$
Reference	.....	.917267
Linear FEM	-----	.914489
Linear Synth	————	.915221

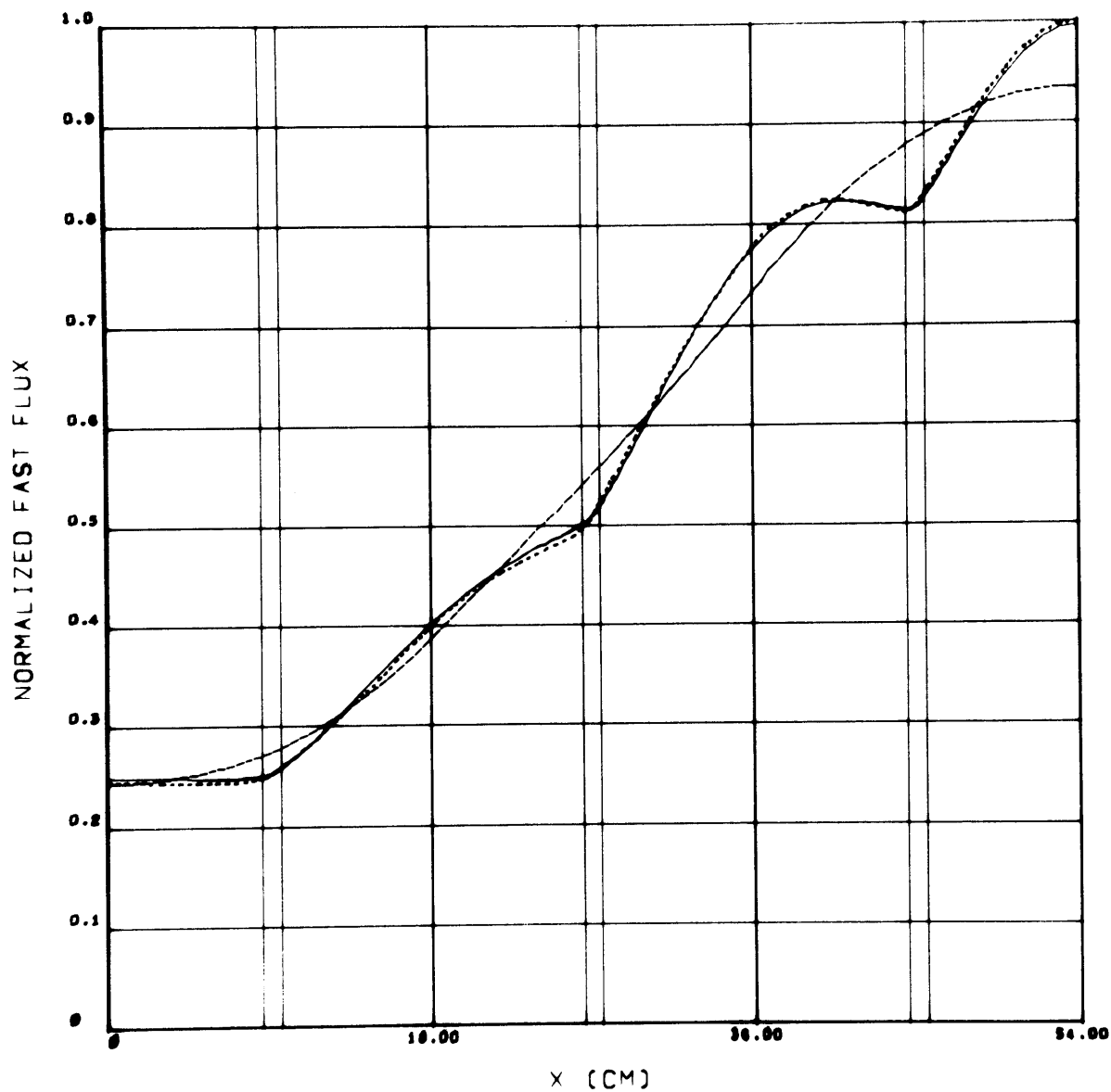


Figure 5.11. Case 1: Two-Group Fast Results Using Cubic Hermite Basis Function Approximations and 18-cm Coarse Mesh Regions

Method		$\lambda$
Reference	.....	.917267
Cubic FEM	-----	.916717
Cubic Synth	—————	.917059

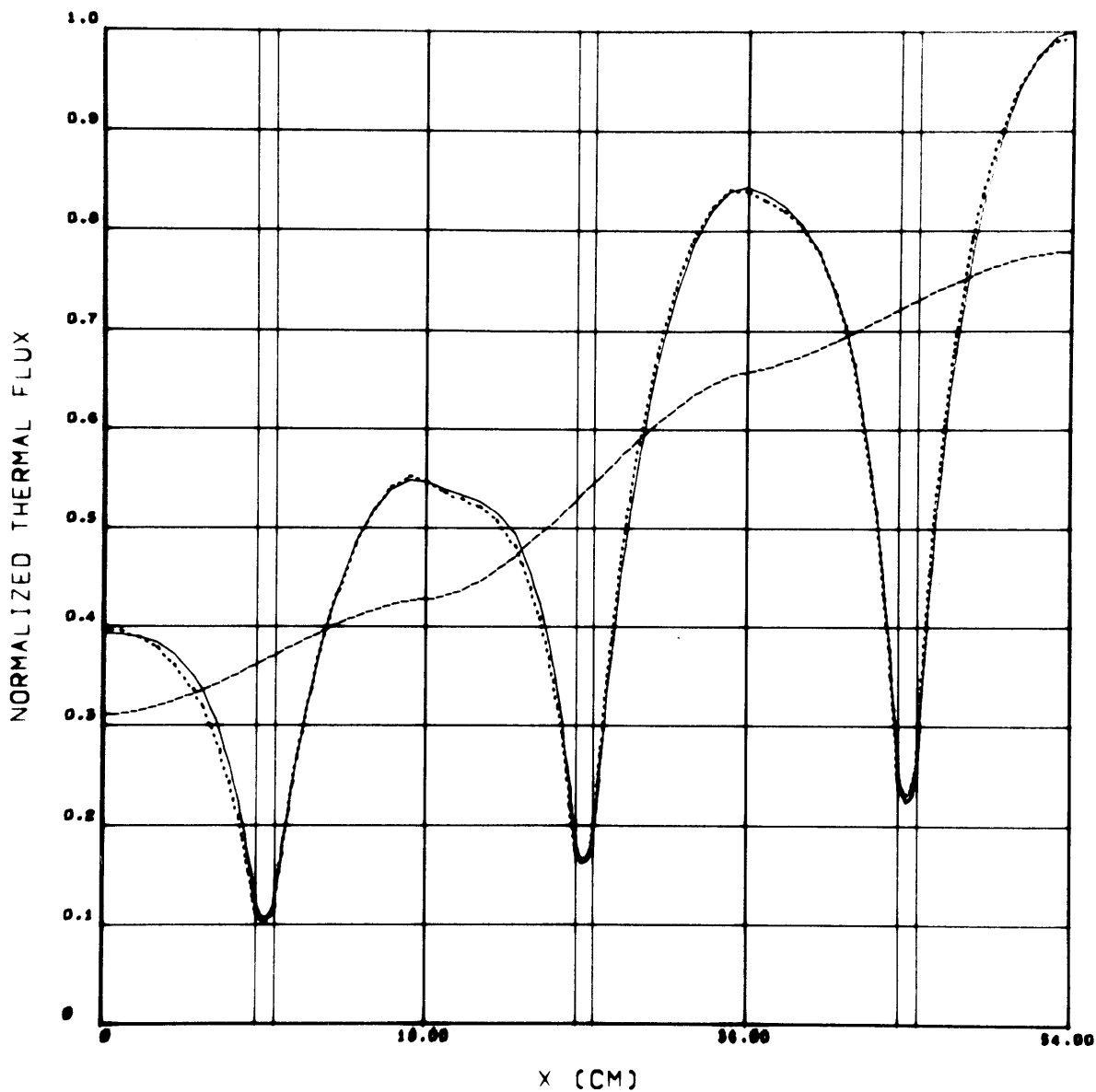


Figure 5.12. Case 1: Two-Group Thermal Results Using Cubic Hermite Basis Function Approximations and 18-cm Coarse Mesh Regions

Method		$\lambda$
Reference	.....	.917267
Cubic FEM	-----	.916717
Cubic Synth	—————	.917059

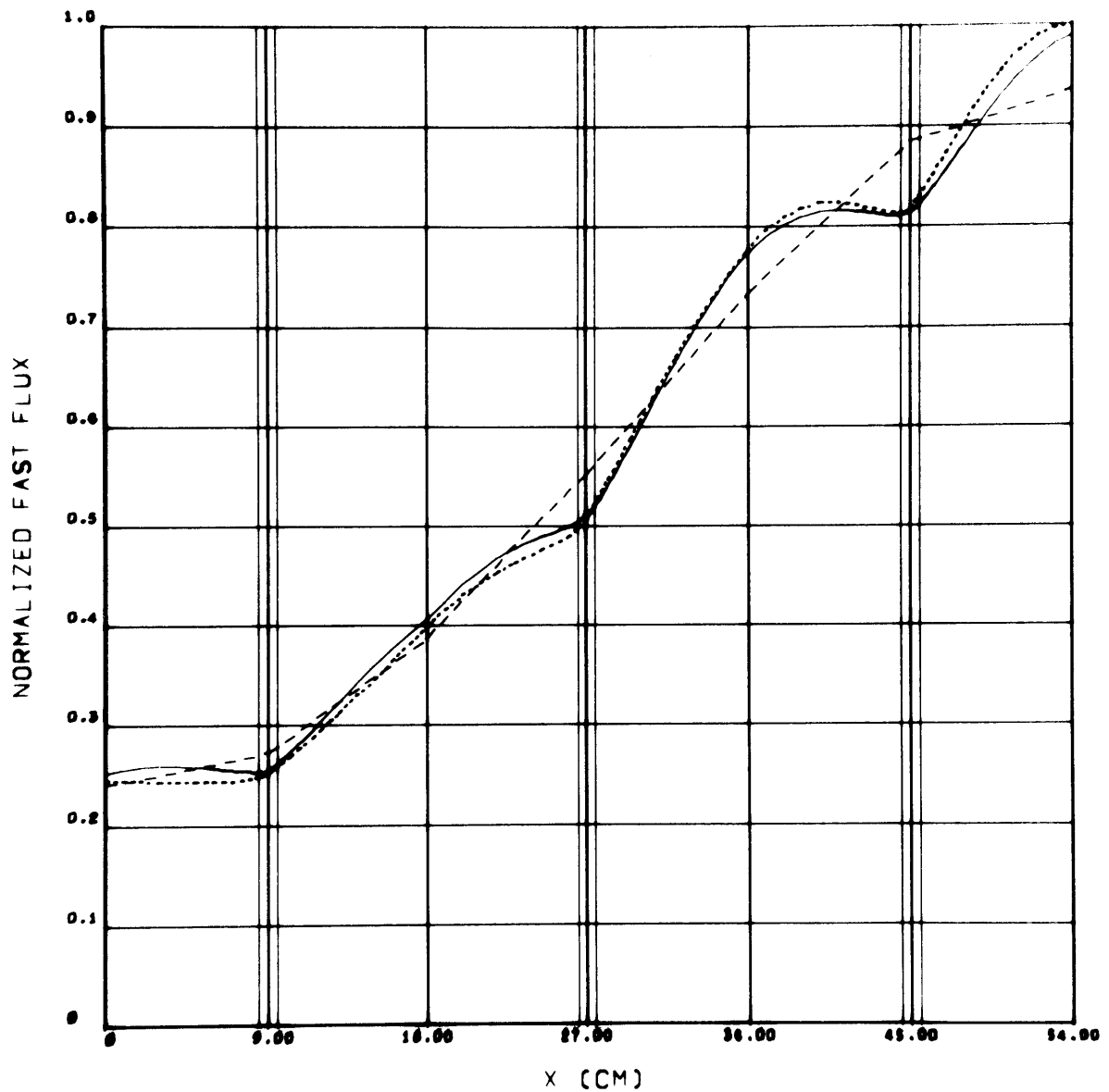


Figure 5.13. Case 1: Two-Group Fast Results Using Linear Basis Function Approximations and 9-cm Coarse Mesh Regions

Method		$\lambda$
Reference	.....	.917267
Linear FEM	-----	.916356
Linear Synth	—————	.916427

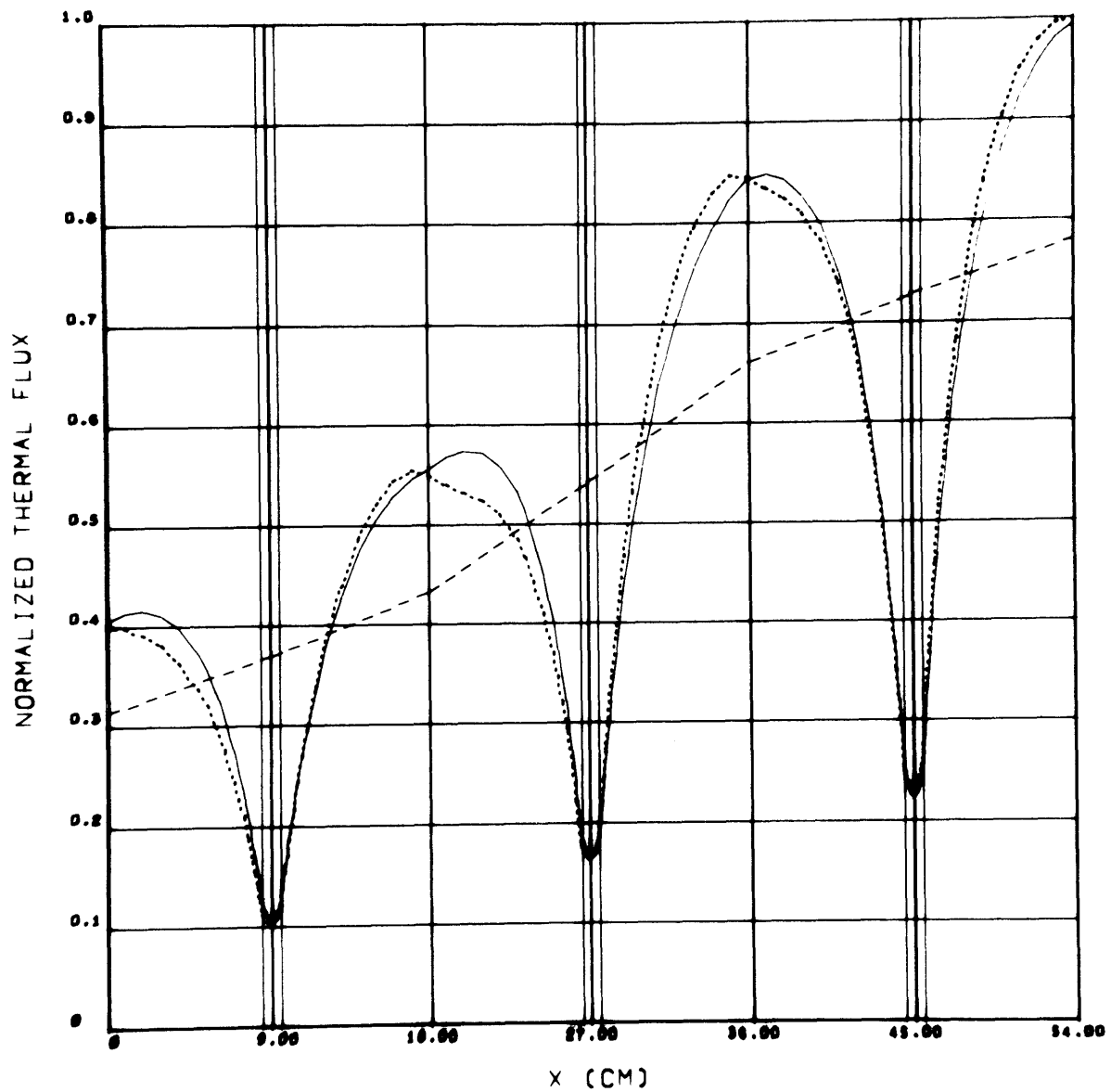


Figure 5.14. Case 1: Two-Group Thermal Results Using Linear Basis Function Approximations and 9-cm Coarse Mesh Regions

Method		$\lambda$
Reference	.....	.917267
Linear FEM	-----	.916356
Linear Synth	—————	.916427



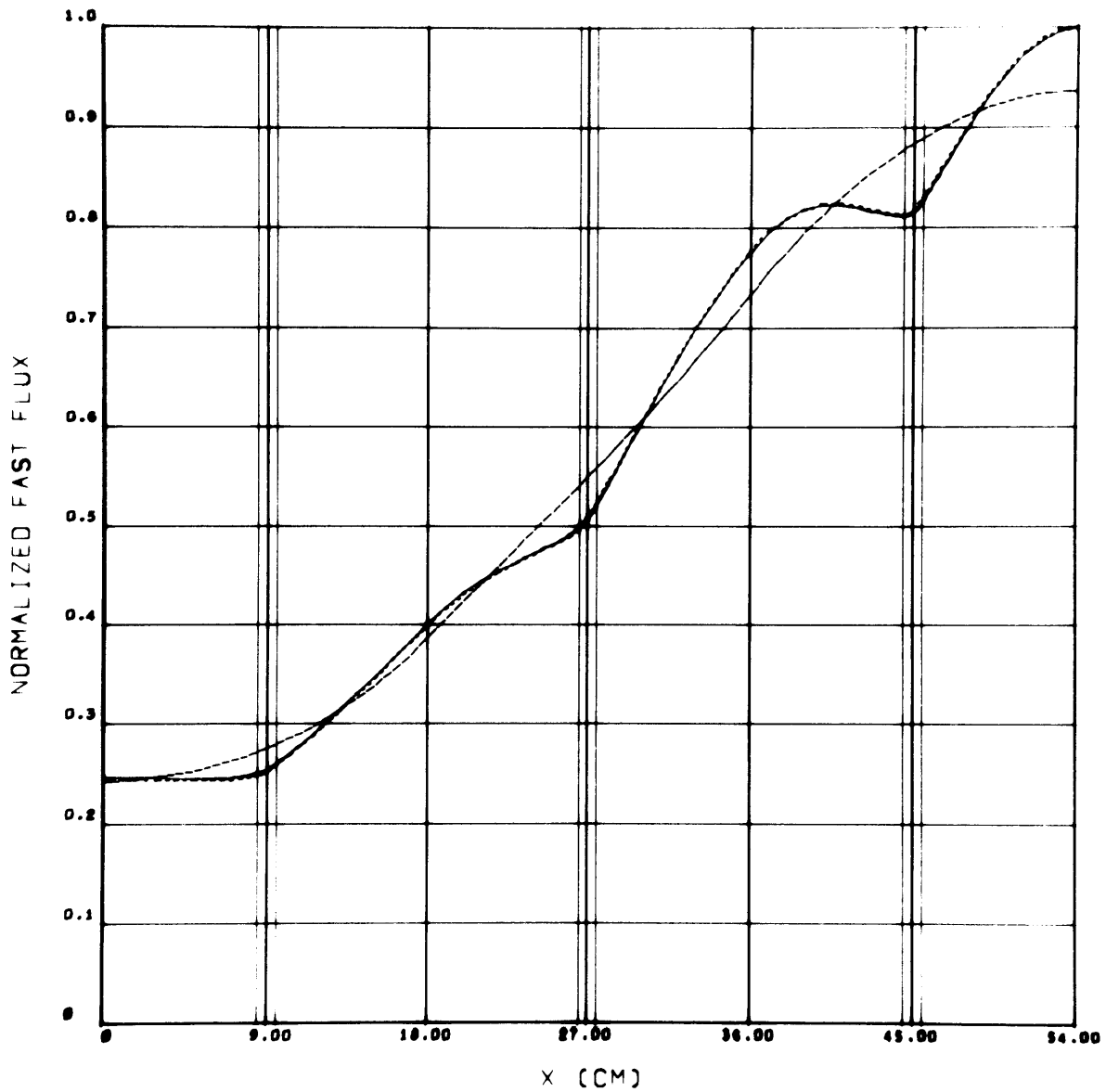


Figure 5.15. Case 1: Two-Group Fast Results Using Cubic Hermite Basis Function Approximations and 9-cm Coarse Mesh Regions

<u>Method</u>		<u><math>\lambda</math></u>
Reference	.....	.917267
Cubic FEM	-----	.916669
Cubic Synth	—————	.917294

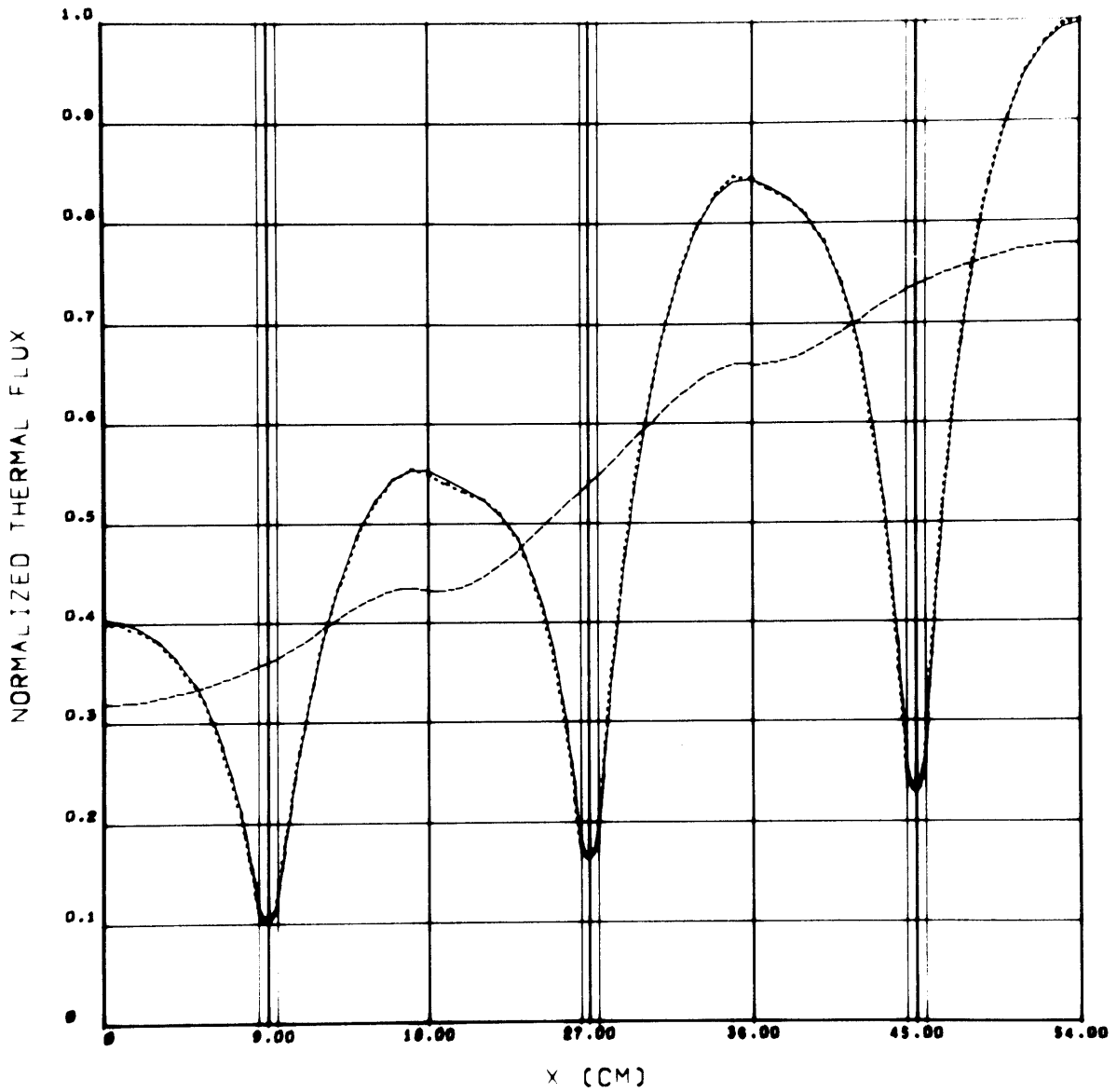


Figure 5.16. Case 1: Two-Group Thermal Results Using Cubic Hermite Basis Function Approximations and 9-cm Coarse Mesh Regions

<u>Method</u>		<u><math>\lambda</math></u>
Reference	.....	.917267
Cubic FEM	-----	.916669
Cubic Synth	—————	.917294

Table 5.6. Results of Case 1.

Method Results	Reference	Linear FEM	Linear Synth	Cubic FEM	Cubic Synth
ONE-GROUP RESULTS:					
$\lambda$	.559045	.556943	.557154	.558647	.558761
$\sigma\lambda$	--	.376%	.338%	.072%	.051%
P(1)	.084	-12.1%	-11.9%	-3.07%	-2.44%
P(2)	.294	-4.29%	-3.93%	-.241%	-.434%
P(3)	.622	+3.67%	+3.47%	+5.28%	+5.34%
TWO-GROUP RESULTS:					
$\lambda$	.917267	.914489	.915221	.916717	.917059
$\sigma\lambda$	--	.302%	.223%	.060%	.023%
P(1)	.134	-6.93%	-5.63%	-1.43%	-1.13%
P(2)	.315	-1.63%	-1.39%	-.072%	-.120%
P(3)	.549	+2.63%	+2.17%	+3.91%	+3.34%
Two-Group Results Using Half-Subassembly Mesh Regions					
$\lambda$	.917267	.916356	.916427	.916669	.917294
$\sigma\lambda$	--	.093%	.092%	.065%	.003%
P(1)	.134	-2.38%	-2.69%	-1.51%	-.461%
P(2)	.315	-.475%	-.602%	-.063%	-.047%
P(3)	.549	+8.51%	-2.63%	-3.22%	+1.40%

It is apparent from these results that the proposed approximation methods, and in particular the method utilizing the cubic Hermite basis functions, approximate to a high degree of accuracy the detailed reference spatial flux. Comparison of the eigenvalue and fractional power results in Table 5.6 indicates that comparable if not superior measurements are obtained using the proposed methods in this case.

It is interesting to note the effects of employing the given subassembly heterogeneous nuclear constants rather than subassembly homogenized nuclear constants for use in the finite element method calculations. Under such conditions, the finite element method becomes identical to the proposed methods in which the heterogeneous nuclear constants and constant or flat subassembly solutions are used. Two-group calculations using the cubic Hermite approximation method were performed for Case 1 and are presented in Figures 5.17 and 5.18. This scheme was found to give very poor detailed flux results, converge to an eigenvalue 21% in error, and yield an average of 20% error in the fractional normalized power levels in each subassembly. This example clearly illustrates the necessity for the use of homogenized constants in the finite element method, or equivalently, the importance of the subassembly detailed solutions in the proposed approximations.

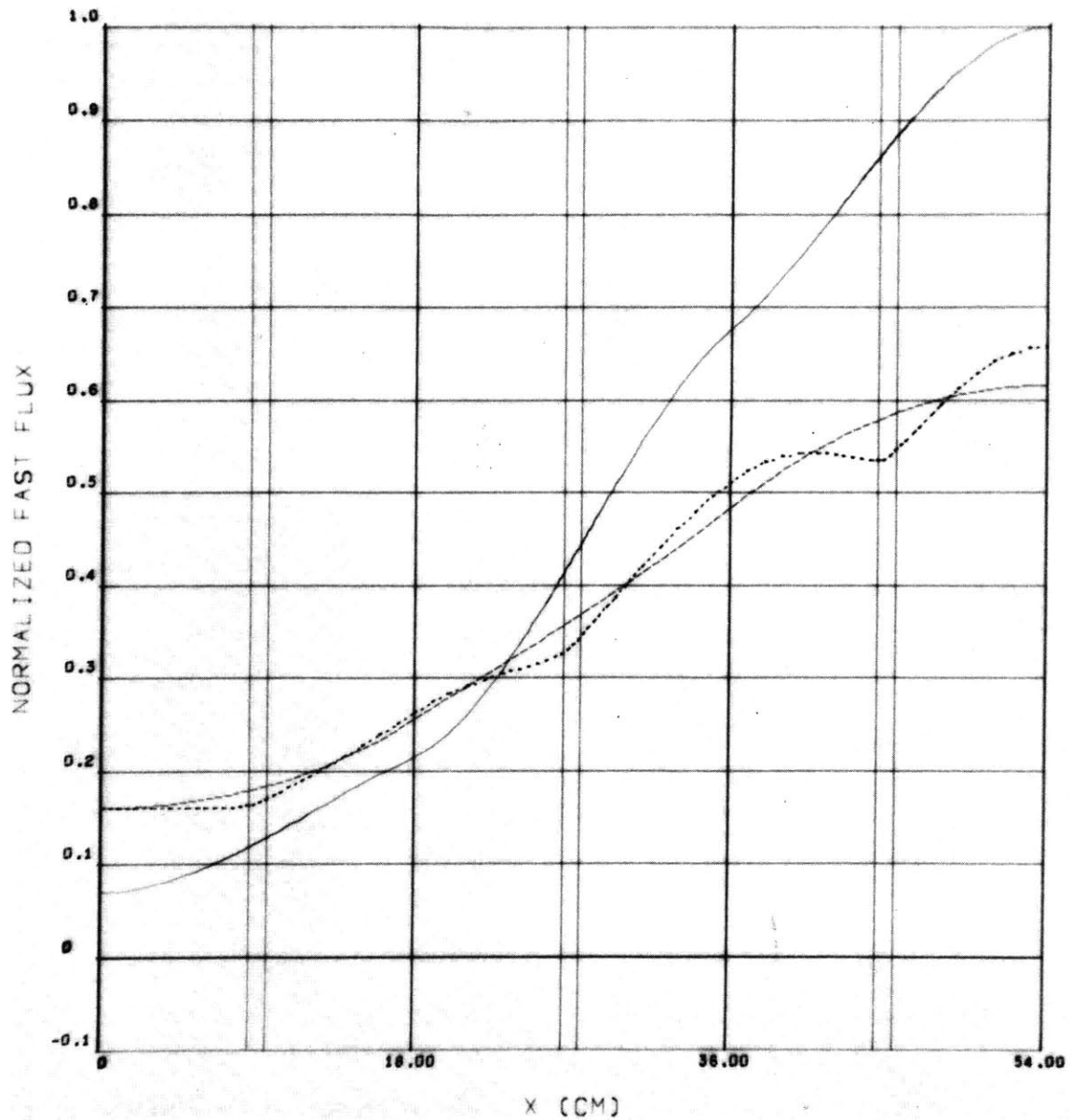


Figure 5.17. Case 1: Two-Group Fast Results Using Cubic Hermite Finite Element Approximations and 18-cm Coarse Mesh Regions

Method		$\lambda$
Reference	.....	.917267
Cubic FEM + Homogenized Consts.	-----	.916717
Cubic FEM + Detailed Consts.	—————	.720422

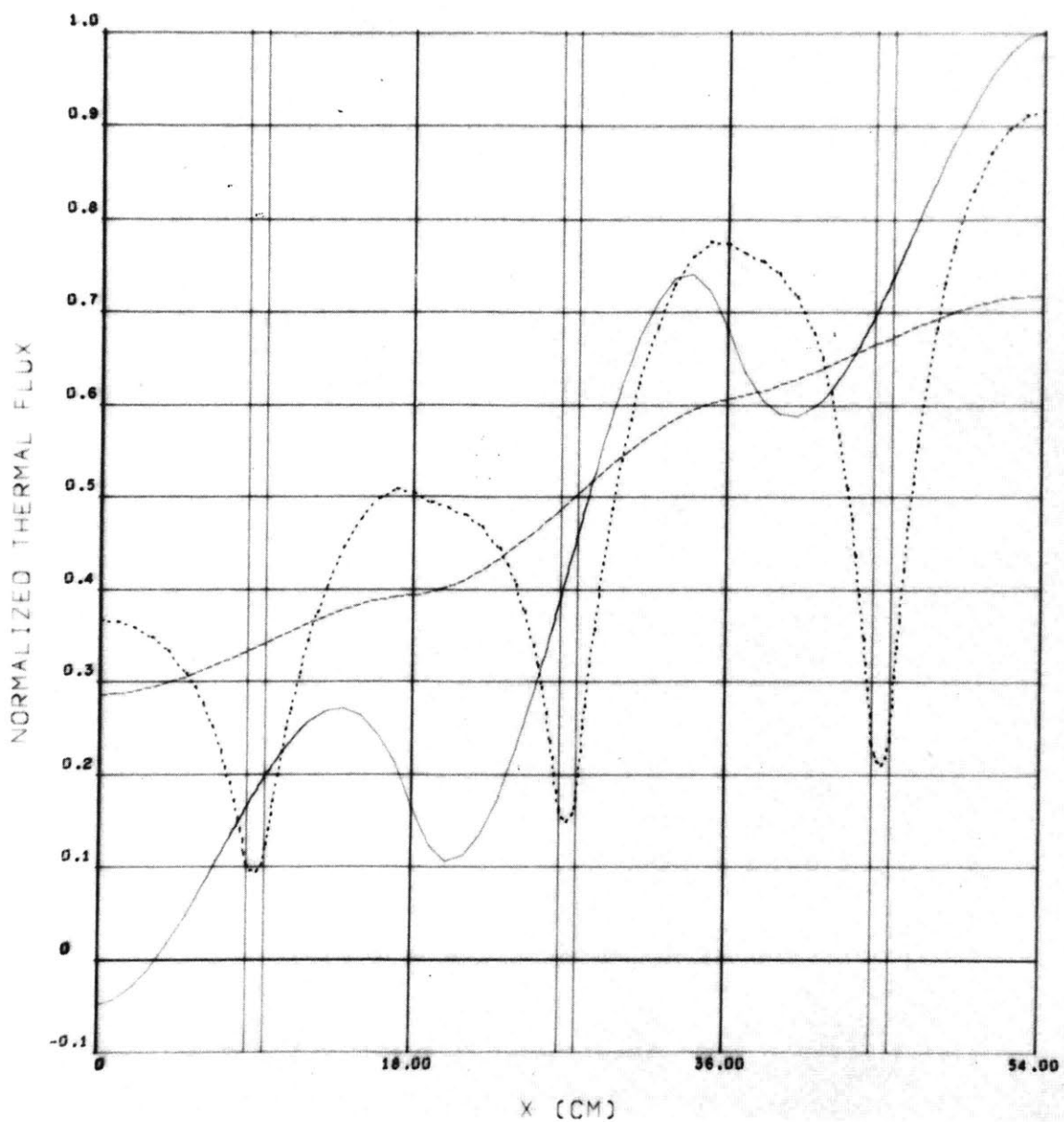


Figure 5.18. Case 1: Two-Group Thermal Results Using Cubic Hermite Finite Element Approximations and 18-cm Coarse Mesh Regions

Method		$\lambda$
Reference	.....	.917267
Cubic FEM + Homogenized Consts.	-----	.916717
Cubic FEM + Detailed Consts.	—————	.720422

5.3.2. Case 2: Three different subassemblies of Types D, B, and C with symmetric boundary conditions.

The results of the two-group approximations for Case 2 are presented in Figures 5.19 - 5.22, where entire subassemblies were taken as the coarse mesh regions. The reference solutions were calculated using the same reference mesh geometry as in Case 1. The converged eigenvalues and fractional normalized power levels in each subassembly are summarized in Table 5.7. These results better illustrate the superiority of the cubic Hermite basis function approximations over the linear basis function approximations, and the superiority of the proposed approximations over the finite element method in all aspects.

Table 5.7. Two-Group Results of Case 2.

Method Results	Reference	Linear FEM	Linear Synth	Cubic FEM	Cubic Synth
$\lambda$	.969986	.965260	.970236	.966816	.969578
$\% \lambda$	--	.487%	-.026%	.326%	.042%
P(1)	.381	+11.26%	+3.46%	+6.07%	+.643%
P(2)	.296	-11.59%	-5.59%	-3.08%	-.157%
P(3)	.322	-2.66%	+1.03%	-4.34%	-.616%

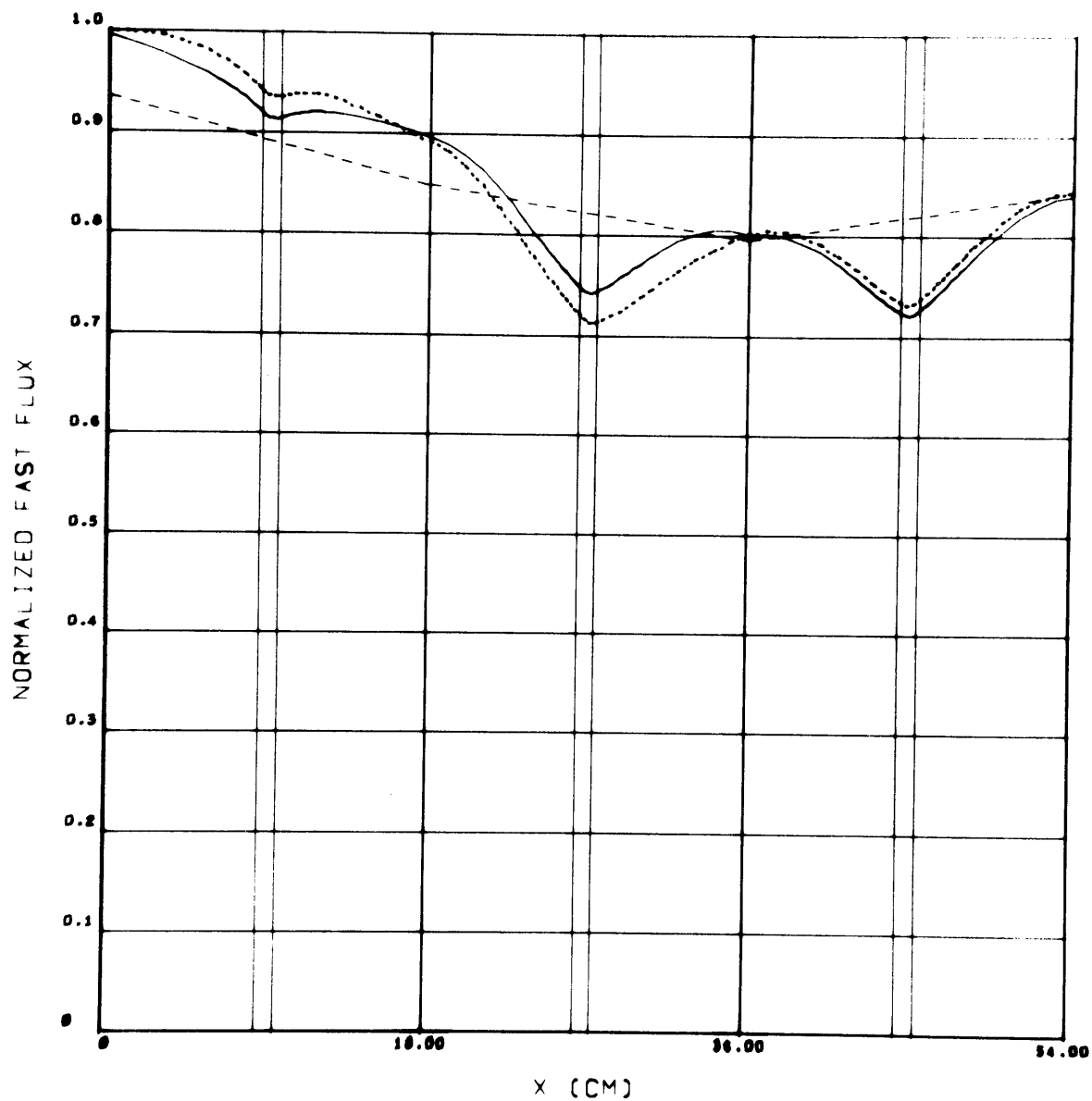


Figure 5.19. Case 2: Two-Group Fast Results Using Linear Basis Function Approximations and 18-cm Coarse Mesh Regions

Method		$\lambda$
Reference	.....	.969986
Linear FEM	-----	.965260
Linear Synth	—————	.970236



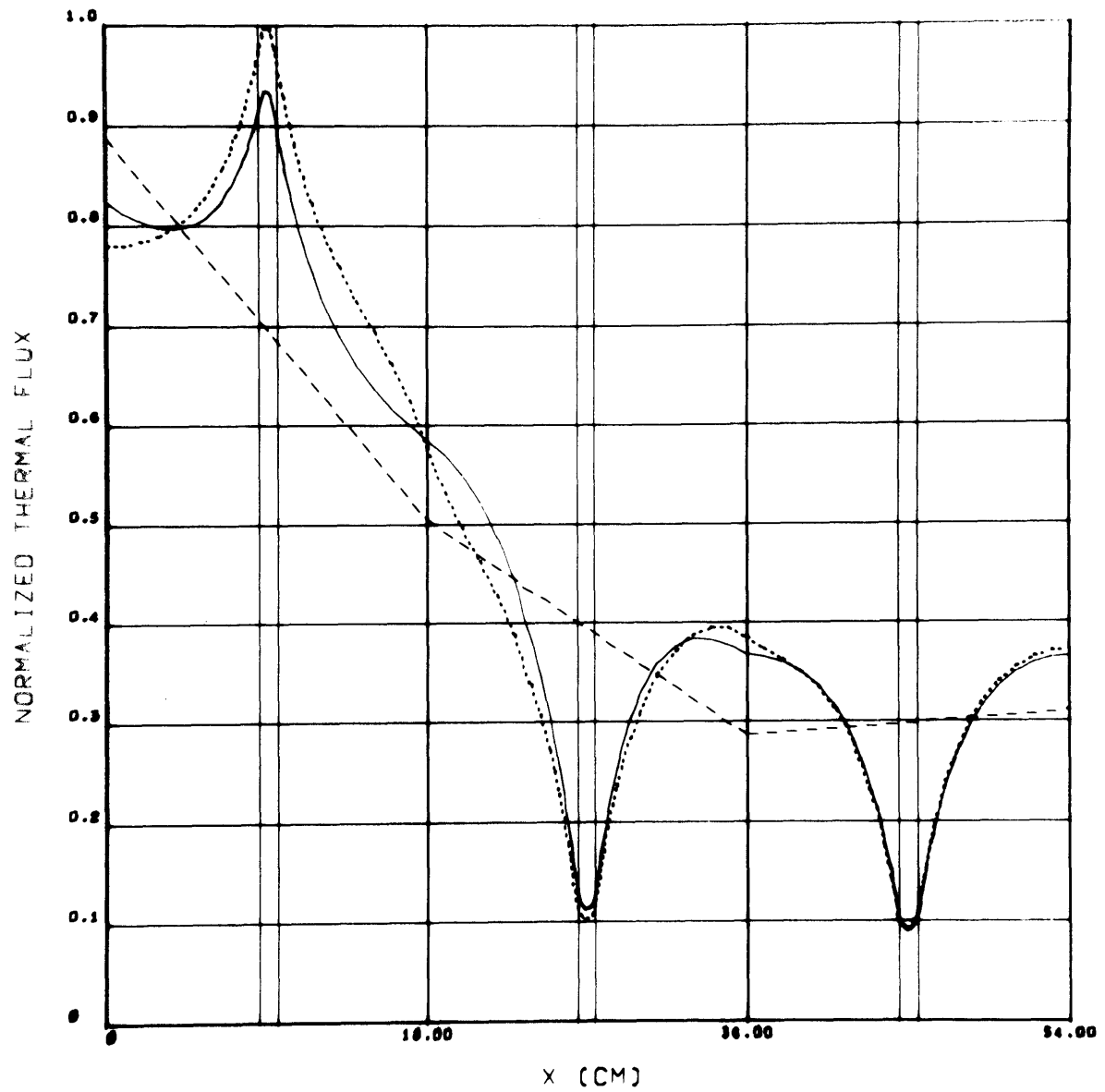


Figure 5.20. Case 2: Two-Group Thermal Results Using Linear Basis Function Approximations and 18-cm Coarse Mesh Regions

Method	$\lambda$
Reference	.969986
Linear FEM	.965260
Linear Synth	.970236

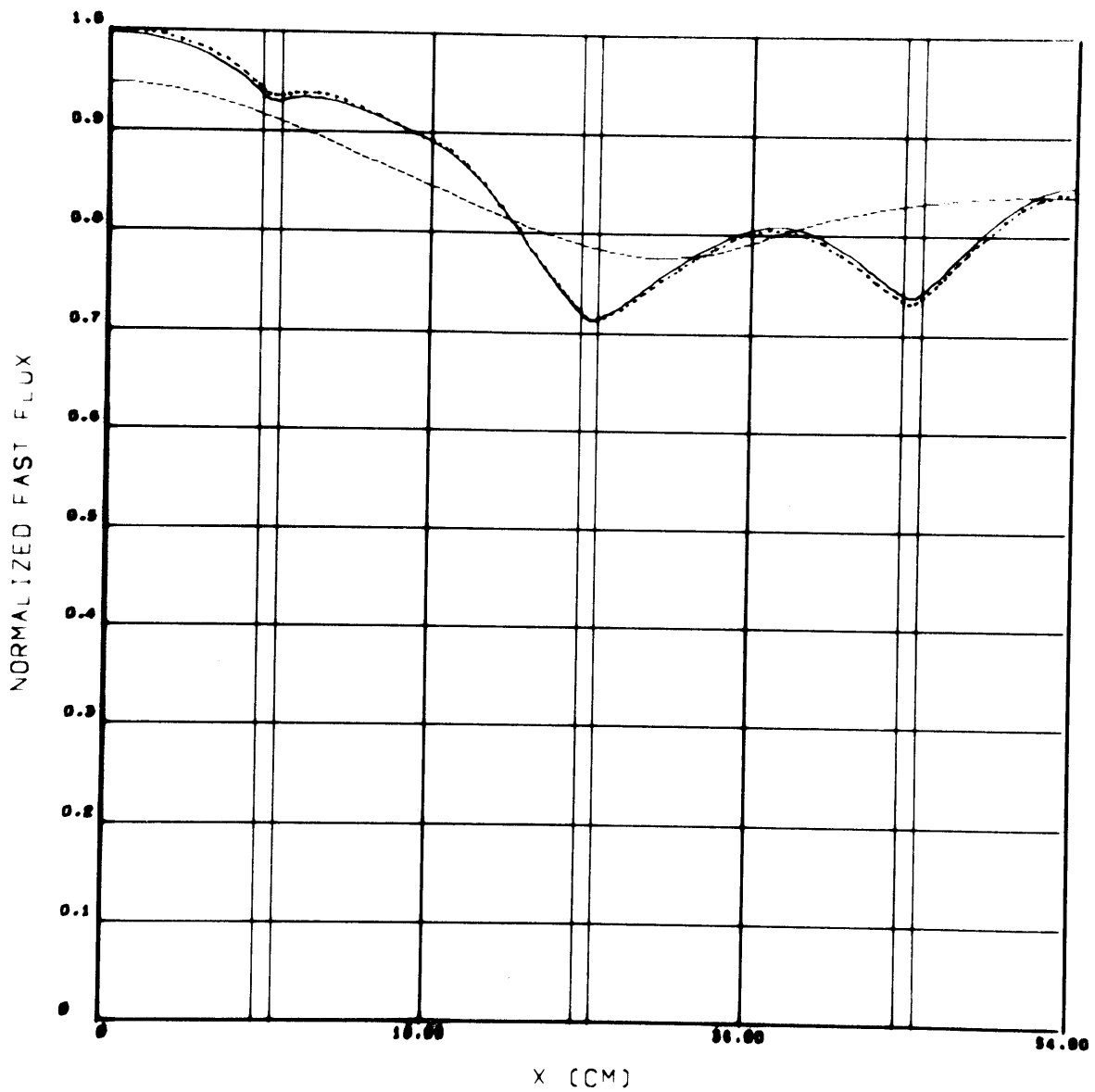


Figure 5.21. Case 2: Two-Group Fast Results Using Cubic Hermite Basis Function Approximations and 18-cm Coarse Mesh Regions

Method		$\lambda$
Reference	.....	.969986
Cubic FEM	-----	.966816
Cubic Synth	—————	.969578

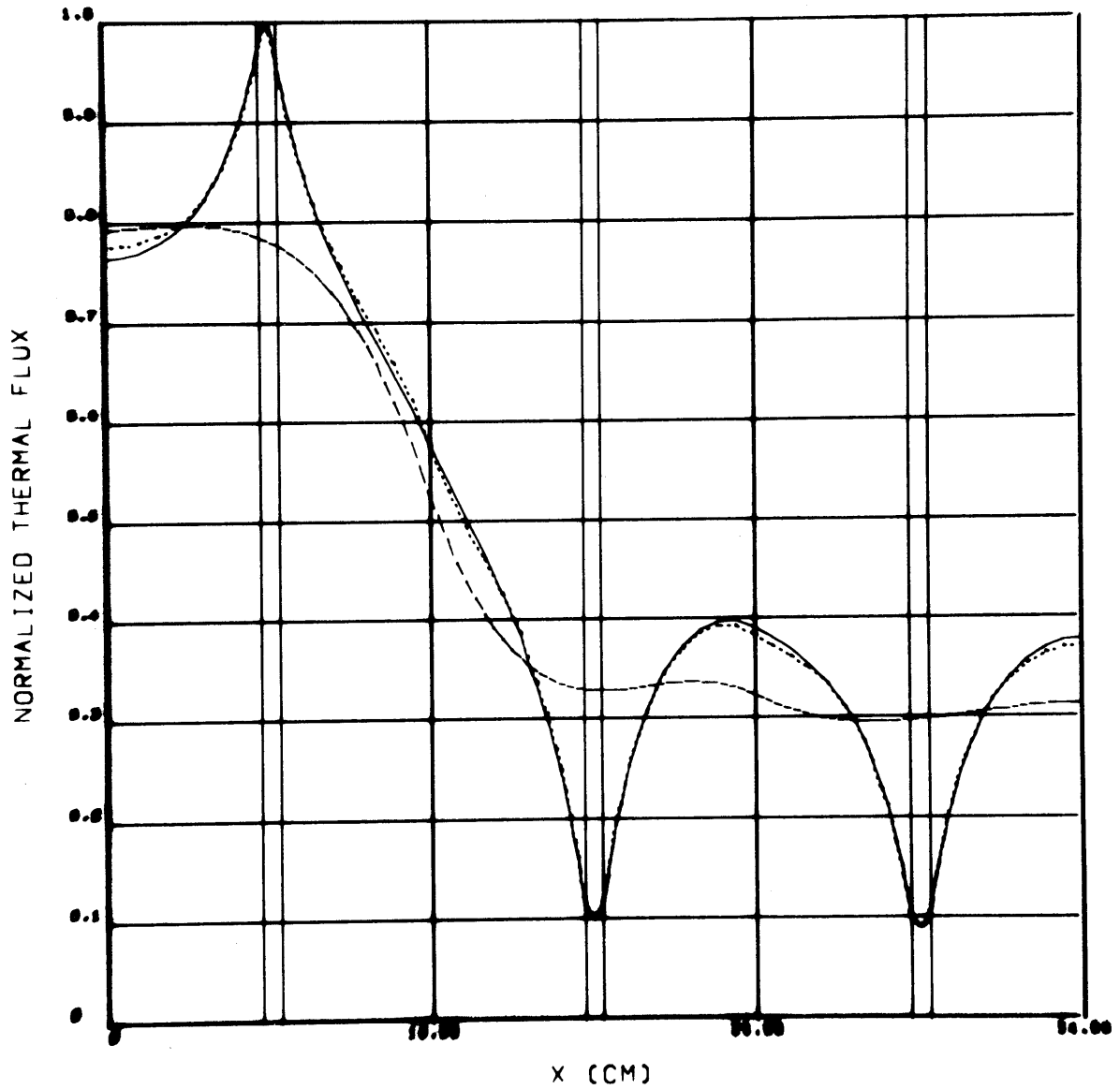


Figure 5.22. Case 2: Two-Group Thermal Results Using Cubic Hermite Basis Function Approximations and 18-cm Coarse Mesh Regions

Method		$\lambda$
Reference	.....	.969986
Cubic FEM	-----	.966816
Cubic Synth	————	.969578

5.3.3. Case 3: Half-core reflected PWR composed of an 18-cm water reflector, the seven subassemblies C,C,C,A,A,A,D, and half of subassembly D. Zero flux boundary conditions are imposed outside the reflector, and symmetry is imposed in the center of the last D-type subassembly.

The Case 3 results of the two-group approximations using full 18-cm coarse mesh regions in all but the last 9-cm region are presented in Figures 5.23 - 5.26, and summarized in Table 5.8. The reference solutions were obtained using 198 mesh regions given by the symmetric partitioning

$$2(2 \text{ cm}) + 2(1 \text{ cm}) + 4(.5 \text{ cm}) + 2(.25 \text{ cm}) + 2(.25 \text{ cm})$$

in each of the subassemblies, and 18 (1 cm) regions in the reflector.

The use of many subassemblies containing absorption rods throughout the reactor, except in the center subassemblies where water channels are present, results in central peaked fluxes with large gradients and, by comparison, a relatively small thermal neutron peak in the reflector.

Both coarse mesh methods were found to overestimate the flux in the subassemblies near the reflector, and underestimate the flux in the central subassembly regions regardless of the type of basis function approximations used. The larger inaccuracies of the linear basis function methods can be in part attributed to the fact that these methods cannot approximate the peaked thermal flux in the reflector, and result in large flux values in the subassemblies nearest the reflector. The cubic Hermite basis function approximations, however, are better able to approximate both the thermal flux reflector peak and the complex

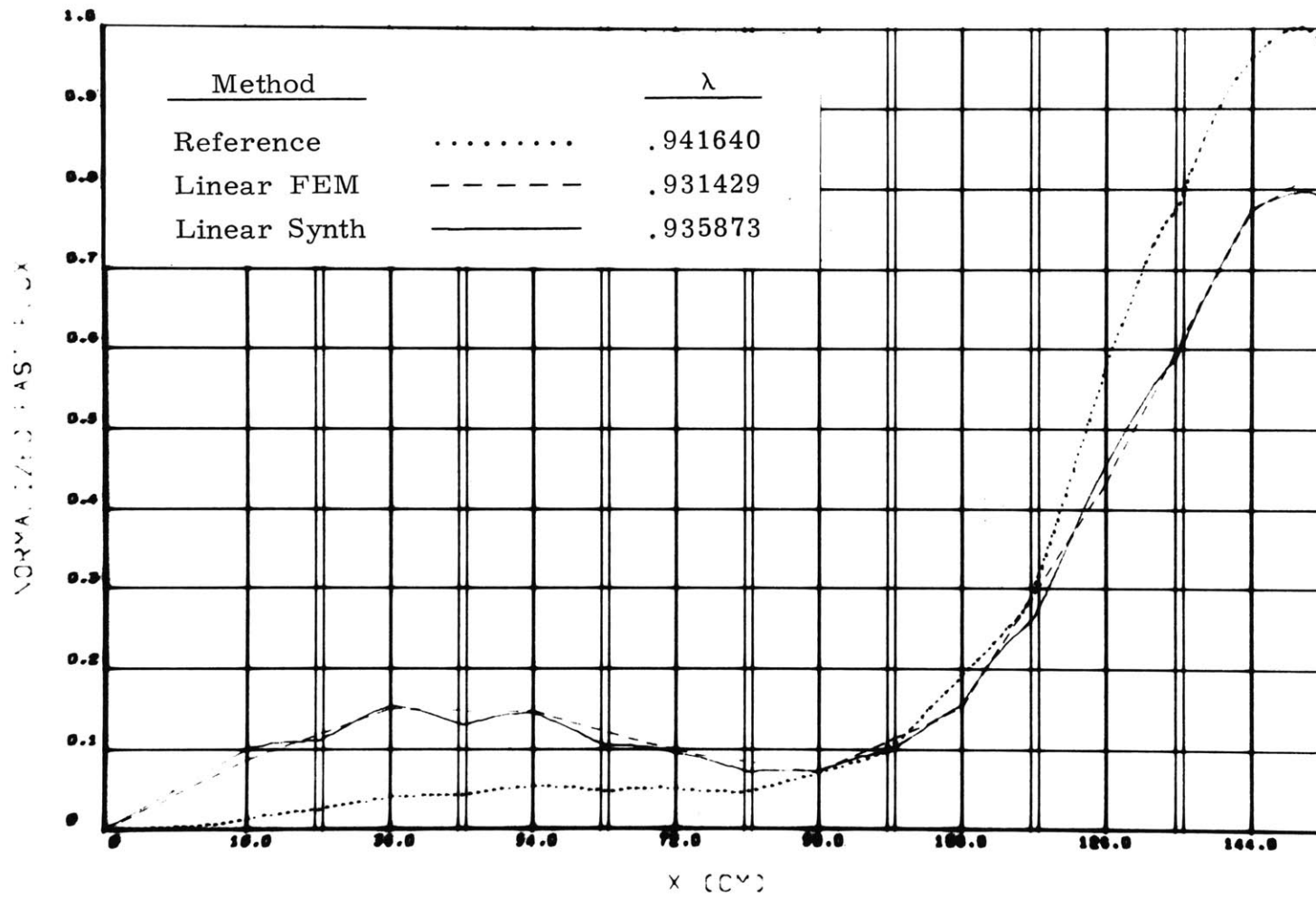


Figure 5.23. Case 3: Two-Group Fast Results Using Linear Basis Function Approximations and 18-cm Coarse Mesh Regions

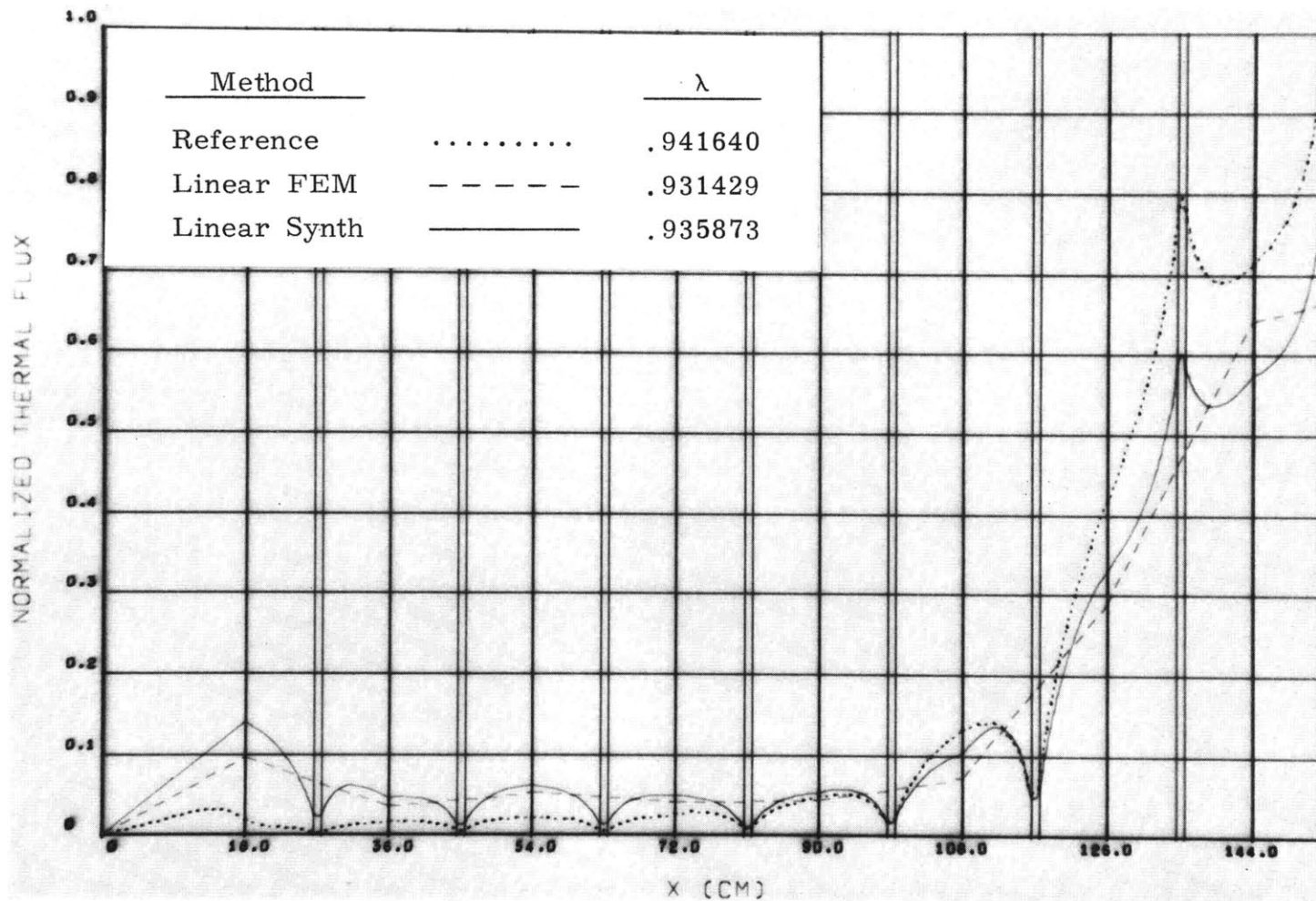


Figure 5.24. Case 3: Two-Group Thermal Results Using Linear Basis Function Approximations and 18-cm Coarse Mesh Regions

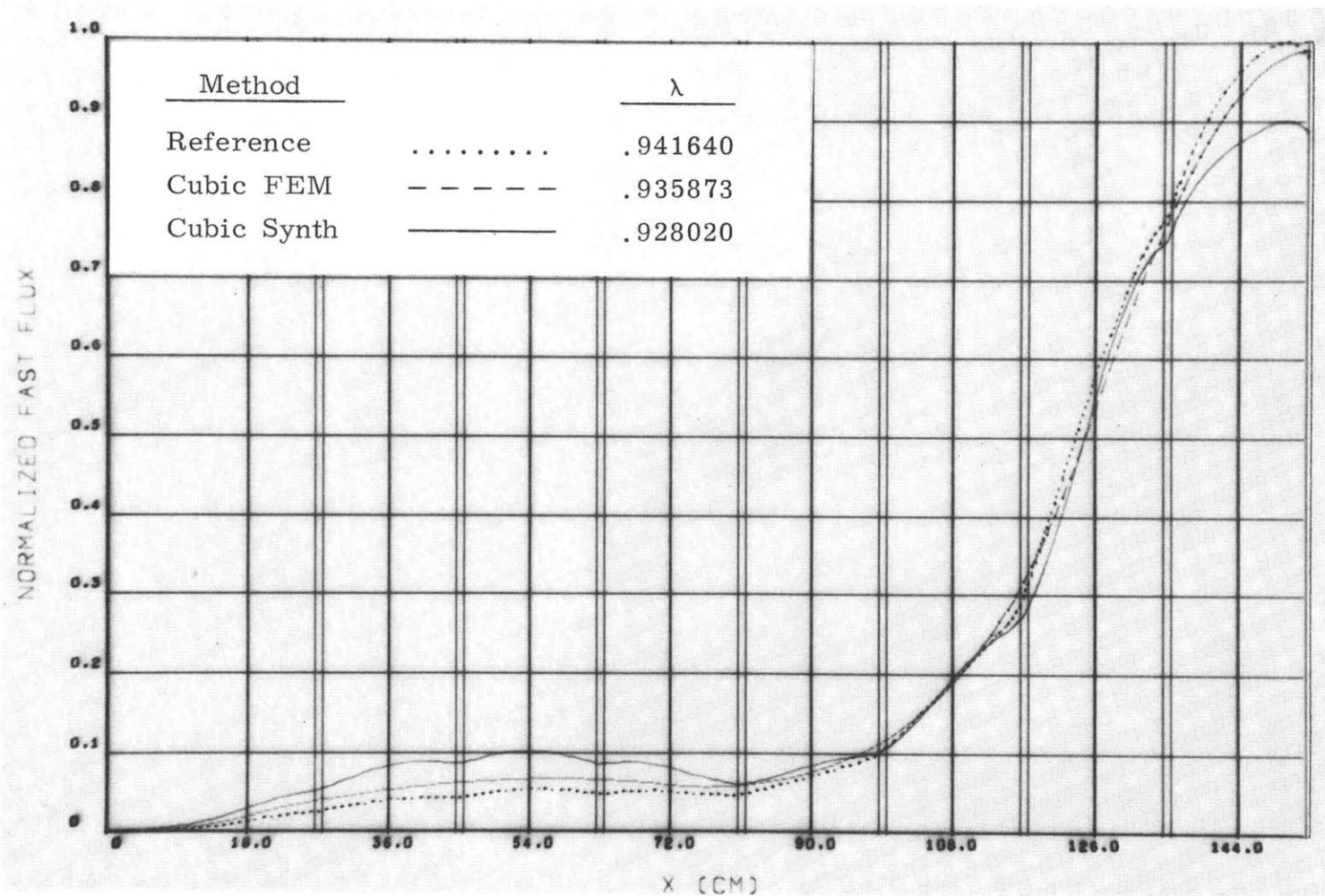


Figure 5.25. Case 3: Two-Group Fast Results Using Cubic Hermite Basis Function Approximations and 18-cm Coarse Mesh Regions

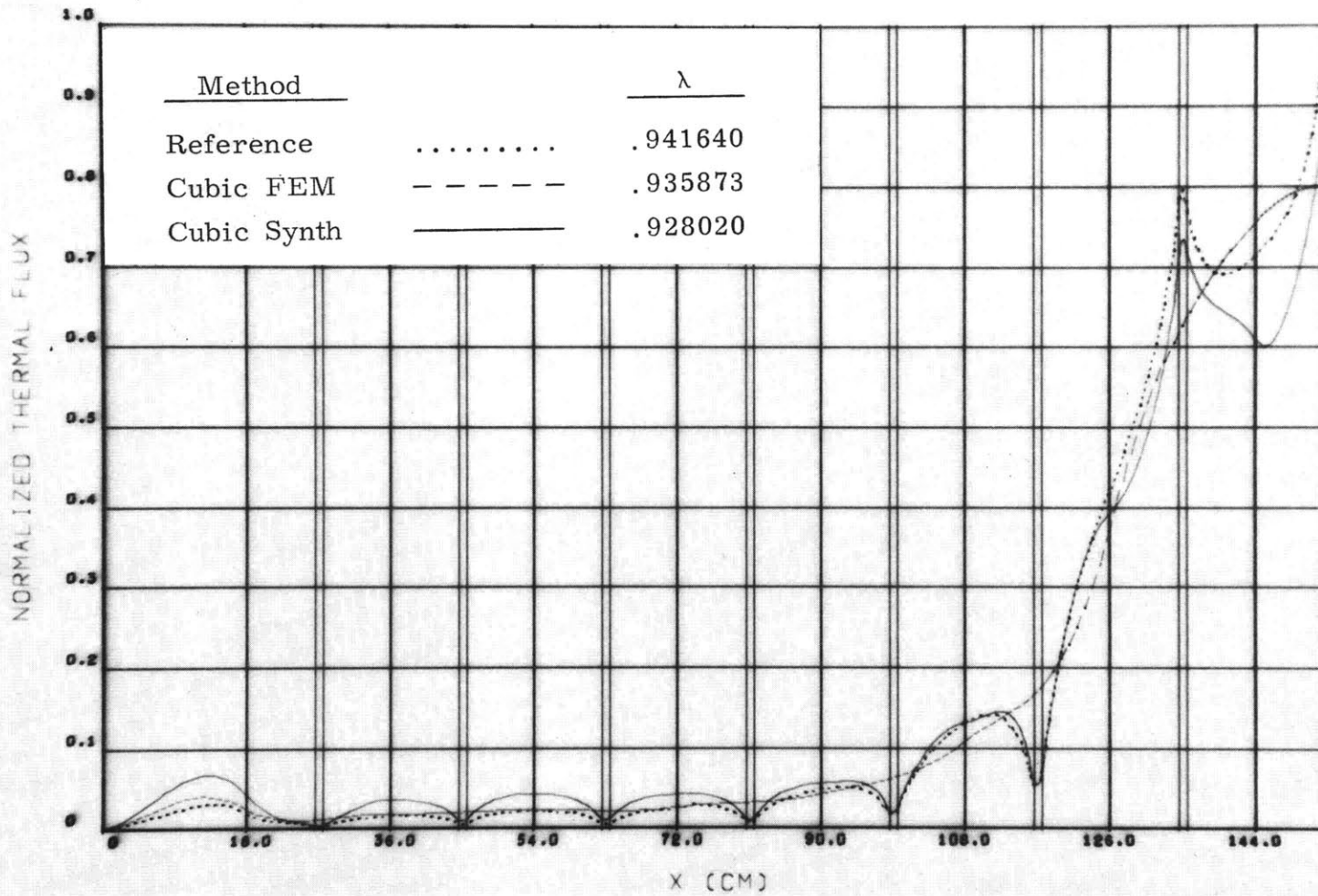


Figure 5.26. Case 3: Two-Group Thermal Results Using Cubic Hermite Basis Function Approximations and 18-cm Coarse Mesh Regions



Table 5.8. Results of Case 3.

Results	Method Reference	Linear FEM	Linear Synth	Cubic FEM	Cubic Synth
$\lambda$	.941640	.931429	.938561	.935873	.928020
$\% \lambda$	--	1.08%	.32%	.61%	1.44%
P(1)	.01699	-461. %	-514. %	-37.4%	-85.1%
P(2)	.02513	-192. %	-179. %	-32.8%	-98.6%
P(3)	.02890	-149. %	-133. %	-26.0%	-79.6%
P(4)	.0224	-50.2%	-36.9%	-13.5%	-31.5%
P(5)	.04969	8.52%	8.24%	-2.78%	-7.63%
P(6)	.1465	7.47%	13.7%	-.494%	.259%
P(7)	.4362	26.4%	22.9%	5.01%	8.40%
P(8)	.2740	18.6%	20.1%	1.99%	13.1%

neutron leakage across the core, and give better results. Table 5.8 indicates that the cubic Hermite basis function approximations better approximate the detailed reference solutions, and that results obtained using the cubic Hermite finite element method were for this case better than those obtained using either of the proposed approximations. The ability of these methods to approximate large thermal flux peaks in the reflector regions is considered in the next case.

5.3.4. Case 4: Half-core reflected PWR composed of an 18-cm water reflector, the seven subassemblies D,D,D,C,D,D,A, and half of subassembly Type A. Zero flux boundary conditions are imposed in the center of the last Type A subassembly.

The Case 4 geometry produces a large but detailed thermal flux in the half-core region and a large thermal peak in the reflector region, as seen from the results in Figures 5.27 - 5.30. The reference solutions were calculated using the reference mesh geometry as given in Case 3. The results of the approximations are summarized in Table 5.9.

The results show that the linear basis function approximations cannot approximate accurately the thermal flux reflector peak and result in large flux and fractional power errors in the subassemblies near the reflector. The cubic Hermite basis function approximations, on the other hand, are better able to approximate this thermal peak and result in much more accurate power levels, especially in the first subassembly region.

The Case 4 results typify the approximation accuracy of both the finite element method and the proposed approximation method. In general, the cubic Hermite basis function approximations are superior to the linear basis function approximations, and the proposed methods give comparable or superior results as compared to those obtained from the finite element method using the same class of basis functions. In this case, the proposed method using cubic Hermite basis functions was able to estimate the reference eigenvalue within 0.04%, closely approximate the detailed reference flux solution to within a few percent at all spatial points, and result in fractional normalized power levels in each subassembly with less than 5% error.

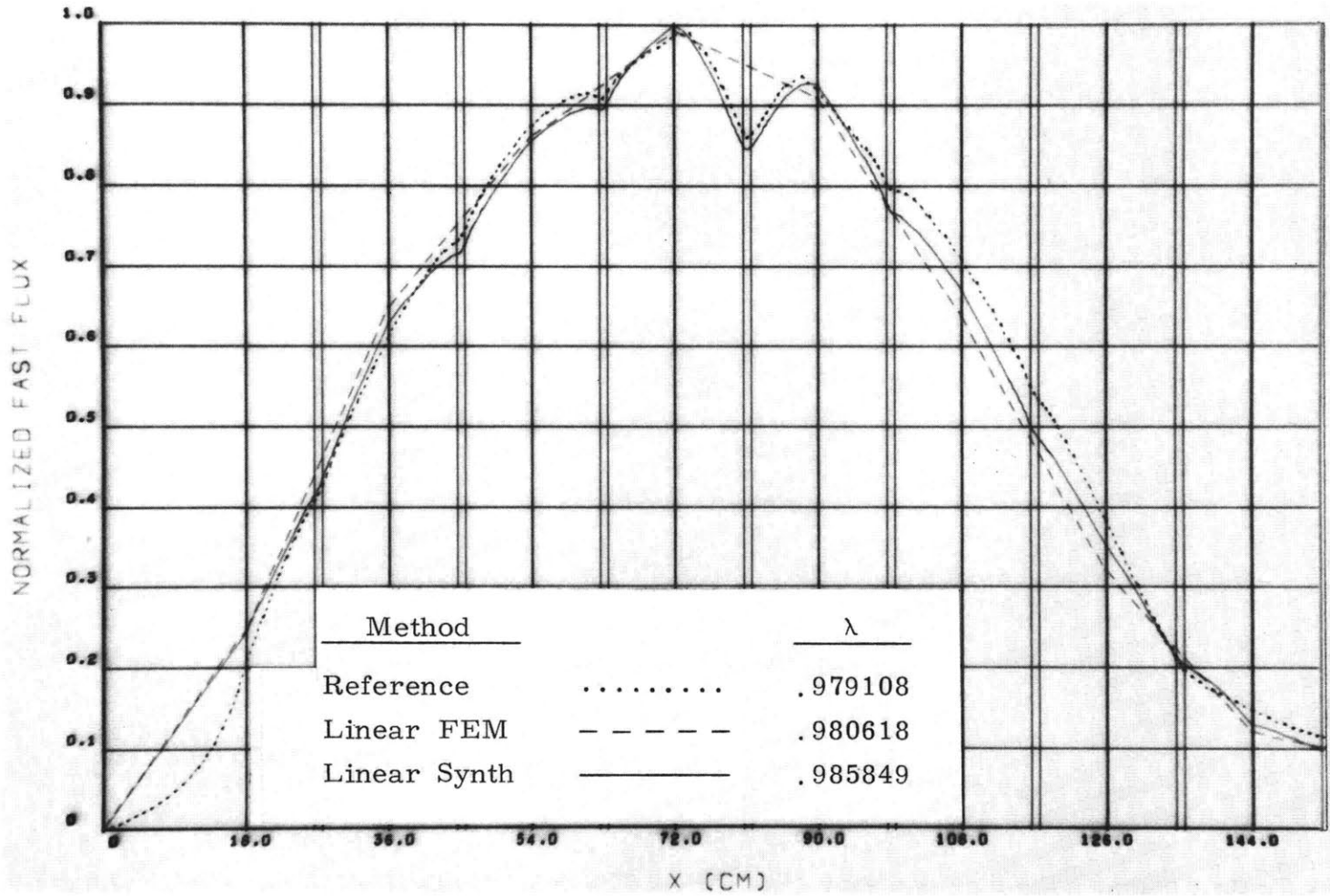


Figure 5.27. Case 4: Two-Group Fast Results Using Linear Basis Function Approximations and 18-cm Coarse Mesh Regions

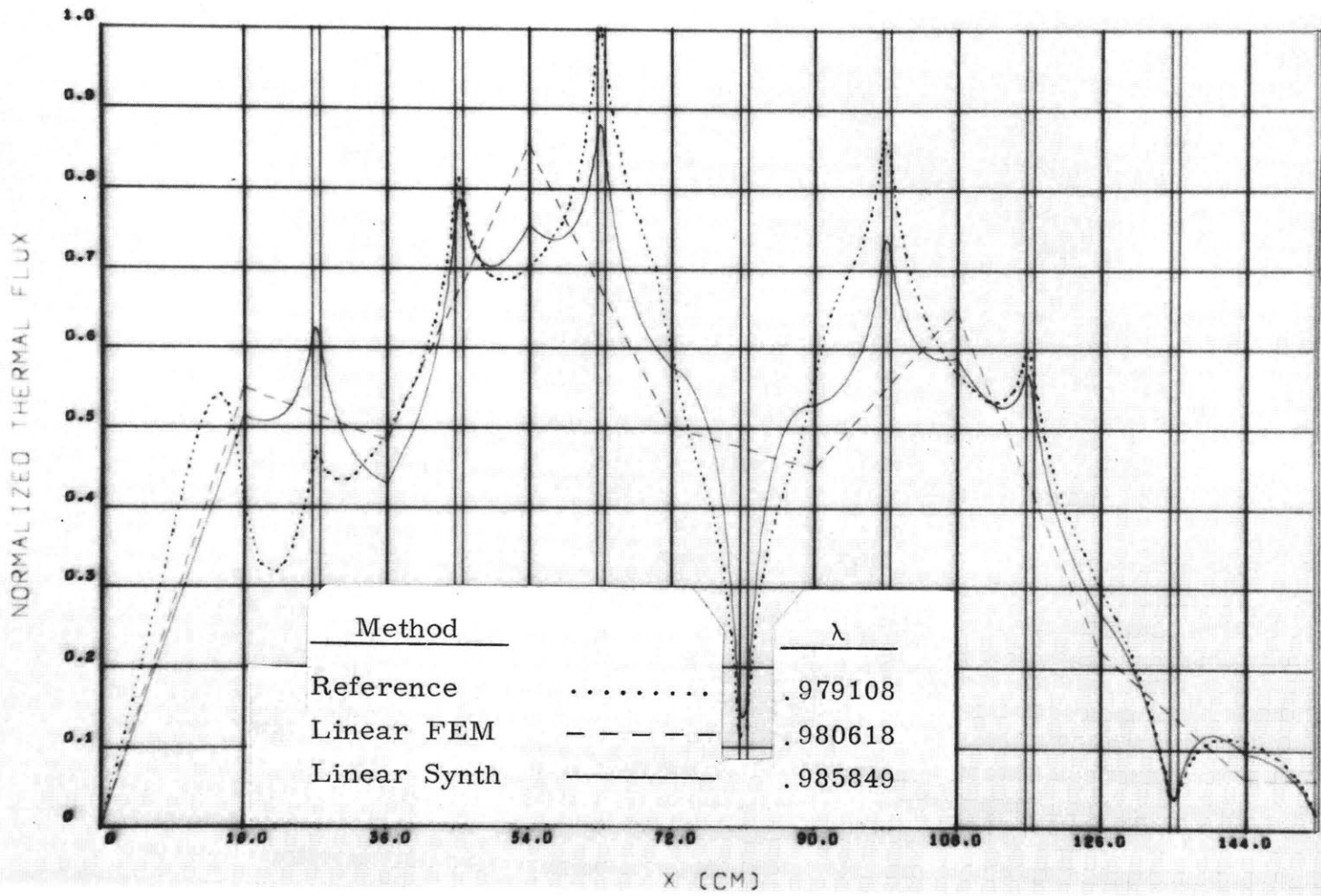


Figure 5.28. Case 4: Two-Group Thermal Results Using Linear Basis Function Approximations and 18-cm Coarse Mesh Regions

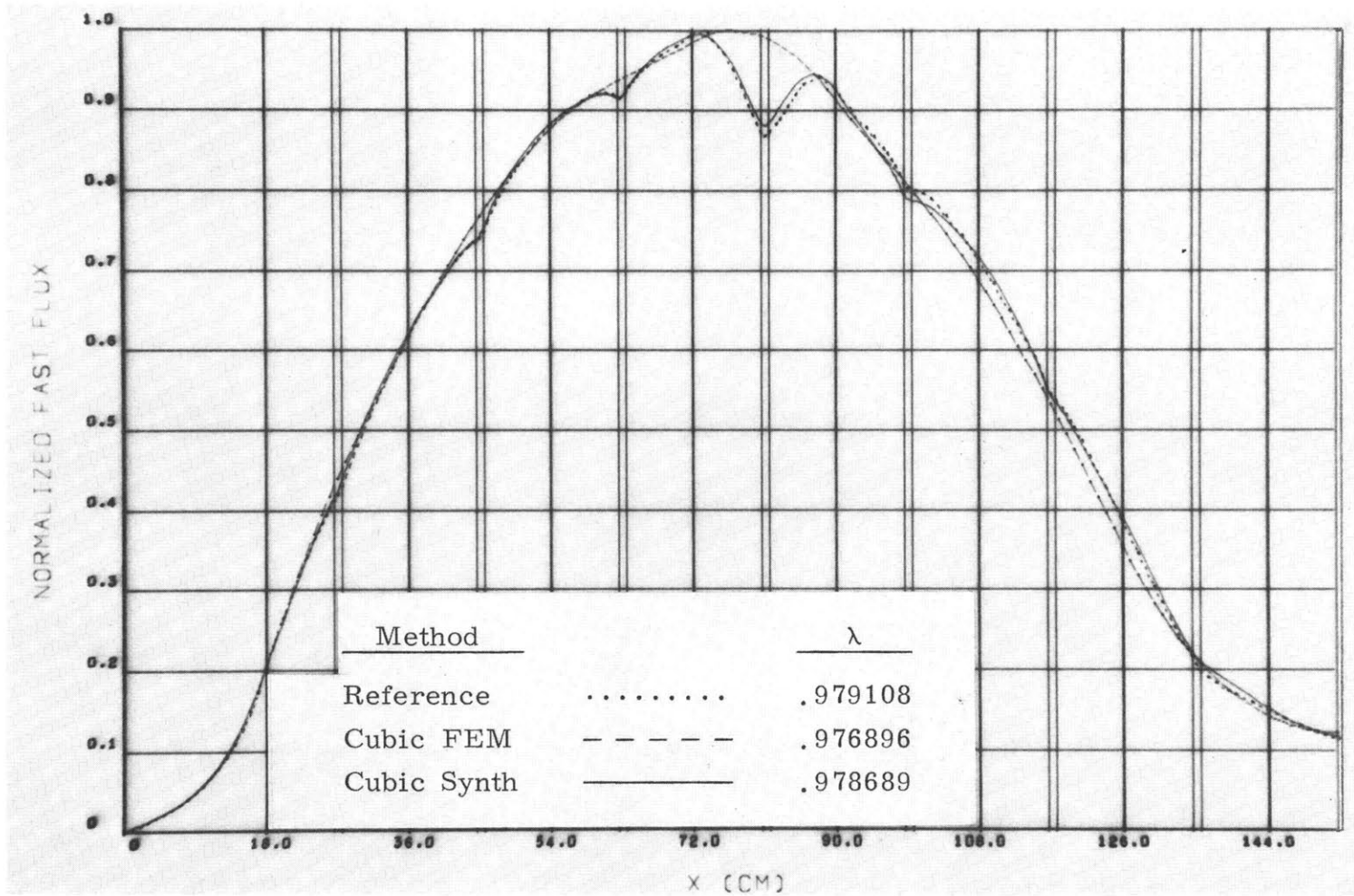


Figure 5.29. Case 4: Two-Group Fast Results Using Cubic Hermite Basis Function Approximations and 18-cm Coarse Mesh Regions

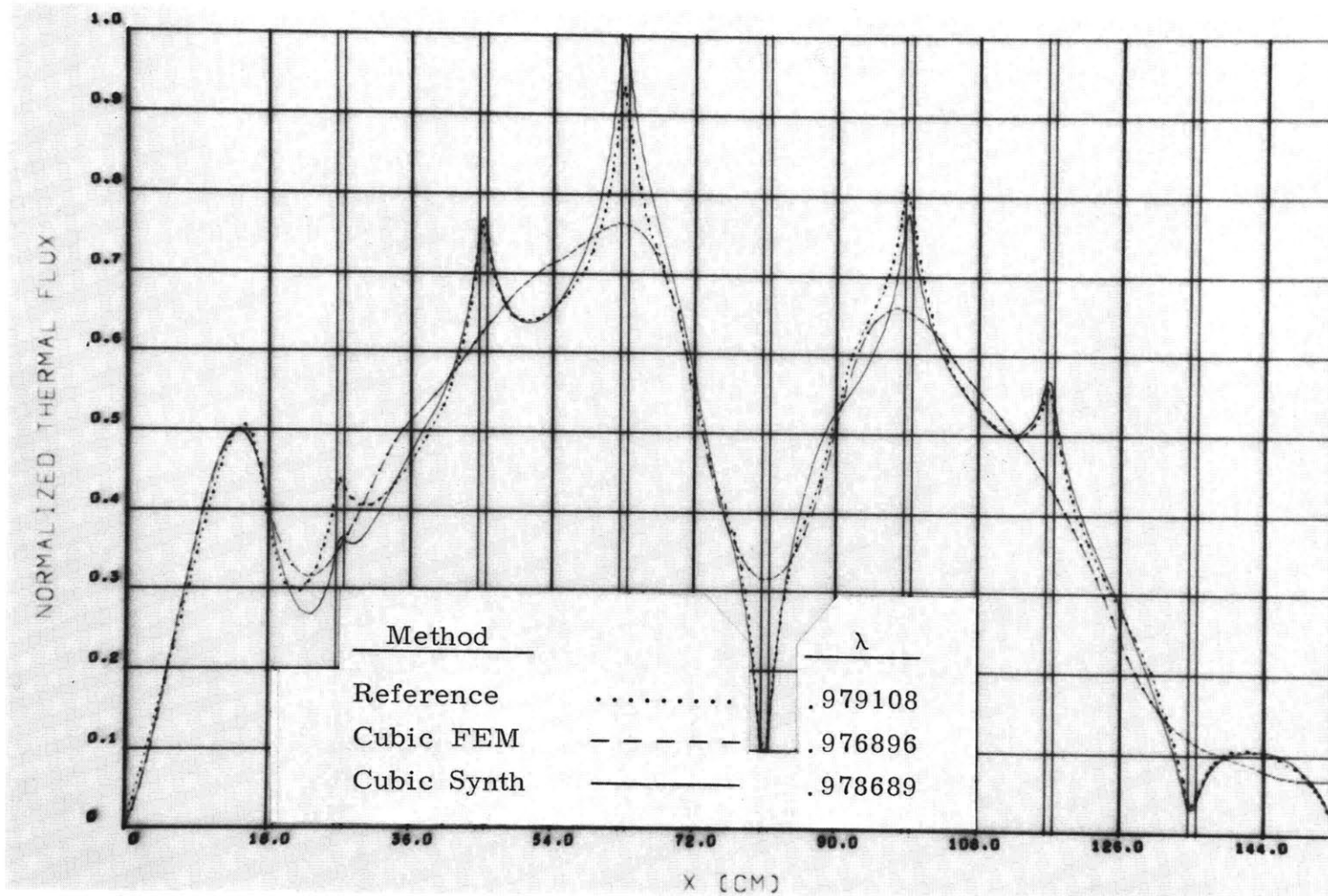


Figure 5.30. Case 4: Two-Group Thermal Results Using Cubic Hermite Basis Function Approximations and 18-cm Coarse Mesh Regions

Table 5.9. Results of Case 4.

Method Results	Reference	Linear FEM	Linear Synth	Cubic FEM	Cubic Synth
$\lambda$	.979108	.980618	.985849	.976896	.978689
$\% \lambda$	--	-.15%	-.69%	.22%	.04%
P(1)	.098	-23.4%	-21.4%	-13.3%	5.54%
P(2)	.160	-1.53%	2.39%	-.893%	-1.08%
P(3)	.193	11.5%	7.11%	1.83%	-3.35%
P(4)	.211	-19.5%	-10.6%	-7.23%	-1.60%
P(5)	.168	17.4%	10.3%	3.86%	3.79%
P(6)	.118	10.6%	4.96%	5.97%	-1.01%
P(7)	.039	1.45%	3.44%	1.65%	3.25%
P(8)	.010	18.1%	11.8%	2.30%	-2.72%

## Chapter 6

## CONCLUSIONS AND RECOMMENDATIONS

6.1. Characteristics of the Proposed Approximation Methods

The use of detailed subassembly flux solutions or other a priori flux shapes directly in the spatial shape or trial function form of flux approximations in reactor physics has resulted in many coarse mesh approximation schemes which are classified in the broad area of overlapping multichannel synthesis. The proposed approximation methods are similar to existing synthesis methods of this kind, but are unique in that they reduce to conventional and well understood approximation methods in regions where little or no spatial flux information is given, or in completely homogeneous regions. In contrast, the overlapping synthesis methods proposed to date do not. This characteristic is especially important in calculations involving homogeneous regions, of which reflector regions are a prime example.

The proposed approximations are very similar to coarse mesh finite element method approximations in which detailed flux behavior has been used to flux-weight the nuclear constants in each region. The methods are conceptually different and become equivalent only when all of the coarse mesh regions are homogeneous.

The matrix equations resulting from the use of the proposed methods are identical in form to those resulting from the finite element method utilizing similar basis functions. In addition, the matrix



elements of the proposed methods are curiously different from those of the finite element methods using detailed flux-weighted nuclear constants. Although the spatial mass and stiffness matrices of the proposed methods for each group have been proven to be positive definite only for the case of Galerkin flux weighting, the use of adjoint weighting in all of the cases considered did not alter these properties. In addition, the proposed methods were found always to converge to a positive eigenvalue and to flux shapes which were everywhere positive.

The numerical results indicate that the proposed methods are able to predict accurate criticality or  $k_{\text{eff}}$  measurements and regional power levels as well as to approximate the reference detailed flux shapes for each group with a high degree of accuracy. The results indicate that in general, use of the proposed methods results in superior criticality estimates over those obtained by the use of the finite element method with flux-weighted constants; this behavior was observed for each type of basis function approximation. Moreover, each of the proposed methods is in general vastly superior to its finite element method counterparts in approximating the actual detailed flux behavior and regional as well as total power levels.

Detailed flux behavior could be reintroduced into the results of the homogenized finite element methods by normalizing the detailed subassembly solutions in each coarse mesh region to match the power levels of the converged results in each region. The detailed solutions resulting from such a procedure would be discontinuous at the region boundaries and may, to some extent, exhibit the fine flux structure present in the results of the proposed methods. However, the results

are not expected to be as good as an approximation as those of the proposed methods, since the current coupling or diffusion approximation is not made until after the coarse mesh homogenization procedure.

## 6.2 Applicability and Limitations

Because the matrix forms of the equations which result from the use of the proposed methods are identical to those which result from the use of the finite element methods, the proposed approximations can be incorporated into existing finite element approximation schemes. Although additional integrations must be performed in the proposed methods, they can be reduced to sums of known products so that little additional computation time is required.

As in any coarse mesh approximation method, inaccurate results can occur when the coarse mesh region sizes chosen are too large. For a given region size, the accuracy of the results for any approximation scheme is unknown. The accuracy of the finite element methods is known to improve geometrically as the mesh size is decreased, resulting in a useful error criterion for the method. A disadvantage of the proposed methods is that no such error criterion has been developed. The inability to predict error estimates has always been a major drawback of synthesis techniques. However, the use of such methods, and use of the proposed methods, has been shown to be justified through proper physical insight and experience.

### 6.3. Recommendations for Future Work

Obviously the next step is the application of these proposed methods to two-dimensional diffusion problems. However, the one-dimensional problem still contains areas which may deserve closer attention. One such area is the examination of the matrix properties of both the finite element method and the proposed approximation methods which are necessary in order to guarantee convergence to a positive eigenvalue and an everywhere positive flux solution. Another area is the development of error criteria for the proposed methods. The close similarity between the proposed methods and the finite element methods may allow an extension or generalization of characteristics which hitherto have belonged only to the finite element methods.

The usefulness of the proposed methods depends on their applicability and accuracy in two- and three-dimensional diffusion problems. Just as the finite element approximations can be derived in two- or three-dimensions using variational modal-nodal techniques, so can the proposed methods for multidimensional problems. The proposed trial functions could be defined as continuous at mesh nodes, but may in general be discontinuous along mesh line interfaces. In order that the flux and current trial functions not be allowed to be discontinuous at identical spatial points, the current trial functions would then have to be defined as continuous across these interfaces. The use of the proposed class of trial function forms in the two-dimensional problem will raise the challenge of extending the spatial overlapping synthesis methods of this type to multidimensional reactor problems.

## REFERENCES

1. T. J. Thompson and J. G. Beckerley, The Technology of Nuclear Reactor Safety, Volume 1, M.I.T. Press (1964).
2. A. M. Weinberg and E. P. Wigner, The Physical Theory of Neutron Chain Reactors, University of Chicago Press (1958).
3. J. R. Lamarsh, Introduction to Nuclear Reactor Theory, Addison-Wesley (1966).
4. F. B. Hildebrand, Finite Difference Equations and Simulations, Prentice-Hall (1968).
5. S. Kaplan, "Synthesis Methods in Reactor Analysis," in Advances in Nuclear Science and Technology, Volume 3, Academic Press, 233 (1966).
6. P. M. Morse and H. Feshbach, Methods of Theoretical Physics, Chap. 6, McGraw-Hill (1953).
7. E. L. Wachspress, Iterative Solution of Elliptic Systems and Applications to the Neutron Diffusion Equations of Reactor Physics, Prentice-Hall (1966).
8. M. Clark, Jr. and K. F. Hansen, Numerical Methods of Reactor Analysis, Academic Press (1964).
9. R. Avery, "Theory of Coupled Reactors," Proc. 2nd U. N. Intern. Conf. Peaceful Uses of Atomic Energy, Geneva, 12, 182 (1958).
10. W. H. Köhler, "Summary of Derivations of Coupled Point Reactor Kinetics Equations," Proc. National Topical Meeting on Coupled Reactor Kinetics, Texas A and M Univ., 192 (1967).
11. M. W. Dyes and K. F. Hansen, "Approximate Methods for Analyzing Coupled Core Kinetics," WARD-35 (1968).
12. M. Komata, "On the Derivation of Avery's Coupled Reactor Kinetics Equations," Nucl. Sci. Eng., 38, 193 (1969).
13. J. B. Yasinsky and A. F. Henry, "Some Numerical Experiments Concerning Space-Time Reactor Kinetics Behavior," Nucl. Sci. Eng., 22, 171 (1965).
14. P. T. Antonopoulos, "Large Mesh Model Development Study," Sc.D. Thesis, M.I.T., April, 1972.

15. M. G. Steveson, "Investigations on the Macroscopic Nodal Approach to Space-Dependent Nuclear Reactor Kinetics," Ph.D. Thesis, University of Texas at Austin (1968).
16. E. L. Wachspress, "Digital Computations of Space-Time Variations of Neutron Fluxes," KAPL-2090 (1960).
17. W. M. Stacey, Jr., Modal Approximations: Theory and an Application to Reactor Physics, Research Monograph No. 41, M.I.T. Press (1967).
18. A. Foderaco and H. L. Garabedian, "Two-Group Reactor Kinetics," Nucl. Sci. Eng., 14, 22 (1962).
19. D. R. Harris, S. Kaplan, and S. G. Marbolis, "Modal Analysis of Flux Tilt: Transients in a Nonuniform Reactor," Trans. Am. Nucl. Soc., 2 (2), 178 (1959).
20. A. F. Henry, "The Application of Inhour Modes to the Description of Nonseparable Reactor Transients," Nucl. Sci. Eng., 13, 22 (1964).
21. E. L. Wachspress, "Variational Methods and Neutron Flux Synthesis," CONF-690401, 271 (1969).
22. E. L. Wachspress, R. D. Burgess, and S. Baron, "Multi-channel Flux Synthesis," Nucl. Sci. Eng., 12, 381 (1962).
23. A. J. Buslik, "Interface Conditions for Few Group Neutron Diffusion Problems with Flux-Adjoint Weighted Constants," Nucl. Sci. Eng., 32, 233 (1968).
24. W. B. Terney, "Interface Conditions for Discontinuous Flux Synthesis Methods," Nucl. Sci. Eng., 41, 303 (1970).
25. J. B. Yasinsky and S. Kaplan, "Synthesis of Three-Dimensional Flux Shapes Using Discontinuous Sets of Trial Functions," Nucl. Sci. Eng., 28, 426 (1967).
26. J. B. Yasinsky and L. R. Foulke, "Improved Spatial Differencing for the Solution of Space-Time Problems Arising from the Movement of Reactor Materials," Nucl. Sci. Eng., 44, 72 (1971).
27. W. M. Stacey, Jr., "Variational Flux Synthesis Methods for Multigroup Neutron Diffusion Theory," Nucl. Sci. Eng., 47, 449 (1972).
28. M. Becker, The Principles and Applications of Variational Methods, Research Monograph No. 27, M.I.T. Press (1964).

29. A. F. Henry, "Variational Techniques as a Means of Deriving Approximation Methods," in M.I.T. Course 22.25 Notes, Advanced Reactor Physics, Chap. 15, p. 104 (Spring, 1969).
30. V. Luco and P. Lambropoulos, "Functionals for Discontinuous Trial Function Flux Synthesis," J. Nucl. Energy A/B, 24, p. 551 (1970).
31. G. C. Pomraning, "A Variational Description of Dissipative Processes," J. of Nuclear Energy A/B, Vol. 20, p. 617 (1966).
32. C. M. Kang and K. F. Hansen, "Finite Element Methods for the Neutron Diffusion Equations," Trans. Am. Nucl. Soc., Vol. 14, No. 1, p. 199 (1971).
33. C. M. Kang and K. F. Hansen, "Finite Element Methods for Space-Time Reactor Analysis," Sc.D. Thesis, Dept. of Nucl. Eng., M.I.T., MITNE-135 (1971).
34. B. W. Roos and W. C. Sangren, "Some Aspects of the Use of Digital Computers in Nuclear Reactor Design," in Advances in Nuclear Science and Technology, Vol. 2, p. 303 (1964).
35. A. F. Henry, "Demonstration That Homogenized Constants for a Cell Depend on Properties of the Surrounding Medium," in M.I.T. Course 22.25 Notes, Advanced Reactor Physics, Chap. 12, p. 63 (Spring, 1969).
36. G. Strang, "The Finite Element Method and Approximation Theory," Numerical Solution of Partial Differential Equations - II, SYNSPADE 1970, Ed. B. Hubbard, p. 547, Academic Press (1971).
37. G. Fix and G. Strang, "Fourier Analysis of the Finite Element Method in Ritz-Galerkin Theory," Studies in Applied Math., 48, 265 (1969).
38. G. Birkhoff, C. deBoor, B. Swartz, and B. Wendroff, "Rayleigh-Ritz Approximation by Piecewise Cubic Polynomials," J. SIAM Numer. Anal., Vol. 3, No. 2, p. 188 (1966).
39. J. J. Goël, "Construction of Basic Functions for Numerical Utilization of Ritz's Method," Numer. Math., 12, p. 435 (1968).
40. A. Radkowsky, Editor, Naval Reactors Physics Handbook, Vol. 1, Selected Basic Techniques, "Homogenization Techniques," Chap. 4.3, USAEC, p. 620 (1964).
41. R. E. Alcouffe and R. W. Albrecht, "A Generalization of the Finite Difference Approximation Method with an Application to Space-Time Nuclear Reactor Kinetics," Nucl. Sci. Eng., 39, 1 (1970).

42. S. Nakamura, "A Variational Rebalancing Method for Linear Iterative Convergence Schemes of Neutron Diffusion and Transport Equations," *Nucl. Sci. Eng.*, 39, 278 (1970).
43. S. Nakamura, "Coarse Mesh Acceleration of Iterative Solution of Neutron Diffusion Equations," *Nucl. Sci. Eng.*, 43, 116 (1971).
44. A. Radkowsky, Editor, Naval Reactors Physics Handbook, Vol. 1, Selected Basic Techniques, "Reactor Physics Computations," Chap. 7.2, USAEC, p. 1414 (1964).
45. G. Bilodeau and L. A. Hageman, "A Survey of Numerical Methods in the Solution of Diffusion Problems," WAPD-TM-64, (July, 1957).
46. G. I. Bell and S. Glasstone, Nuclear Reactor Theory, Van Nostrand Reinhold Company (1960).
47. W. W. Little, Jr. and R. W. Hardie, "2DB User's Manual - Revision 1," BNWL-831 REV 1, UC-32: Mathematics and Computers, Battelle Memorial Institute, Pacific Northwest Laboratories, Richland, Washington (August, 1969).
48. L. A. Hageman, "Numerical Methods and Techniques Used in the Two-Dimensional Neutron-Diffusion Program PDQ-5," WAPD-TM-364, UC-32: Mathematics and Computers, Bettis Atomic Power Company, Pittsburgh, Pennsylvania (February, 1963).
49. E. Isaacson and H. B. Keller, Analysis of Numerical Methods, Wiley (1966).
50. R. S. Varga, Matrix Iterative Analysis, Prentice-Hall (1962).
51. G. Forsythe and C. B. Moler, Computer Solution of Linear Algebraic Systems, Prentice-Hall (1967).
52. L. A. Hageman and J. B. Yasinsky, "Comparison of Alternating-Direction Time-Differencing Methods with Other Implicit Methods for the Solution of the Neutron Group-Diffusion Equations," *Nucl. Sci. Eng.*, 38, p. 8 (1969).
53. A. Radkowsky, Editor, Naval Reactors Physics Handbook, Vol. 1, Selected Basic Techniques, "Few-Group Approximations," Chap. 2.8, USAEC, p. 197 (1964).
54. M.I.T. Information Processing Center, "Stromberg-Carlson 4020," Application Program Series, AP-36 Revision 2, November 30, 1971.

## BIOGRAPHICAL NOTE

The author is a native Californian. Born in San Francisco August 19, 1945, he was raised in Orange, and upon graduation from Orange High School in 1963, he received the Bank of America Award in Mathematics and the Orange County Industrial Education Association Award in Electronics.

He then attended the University of California at Berkeley, and supported by special as well as Regents' Scholarship awards, graduated in 1967 from the School of Engineering with a Bachelor of Science degree in Physics. He was elected to Tau Beta Pi, Phi Beta Kappa, and received a four-year letter in varsity gymnastics.

The author then enrolled at the Massachusetts Institute of Technology. Supported by an Atomic Energy Commission Special Fellowship in Nuclear Engineering, he received the Master of Science degree from the Department of Nuclear Engineering in 1969.

As a research associate in the Space-Time Kinetics Project in the same department, the author completed this work and obtained the degree of Doctor of Philosophy in 1972.

While at M.I.T. the author was active in many organizations, most notable of which was the M.I.T. Rugby Football Club.



APPENDICES

## Appendix A

## TABLE OF SYMBOLS

$g$	Energy group index which runs from the highest to the lowest energy group as $g = 1$ to $G$ .
$\phi_g(r)$	Scalar neutron flux in energy group $g$ (neutrons/cm <sup>2</sup> · sec).
$\underline{J}_g(r)$	Vector neutron current in energy group $g$ (neutrons/cm <sup>2</sup> · sec).
$D_g(r)$	Diffusion coefficient for neutrons in energy group $g$ (cm).
$\Sigma_g(r)$	Macroscopic total removal cross section in energy group $g$ (cm <sup>-1</sup> ).
$\nu\Sigma_{fg}(r)$	Macroscopic fission-production cross section in energy group $g$ (cm <sup>-1</sup> ).
$\Sigma_{gg'}(r)$	Macroscopic transfer cross section from energy group $g'$ to energy group $g$ (cm <sup>-1</sup> ).
$\chi_g$	Fission spectrum yield in energy group $g$ .
$\lambda$	The eigenvalue or criticality of the diffusion problem.
$\Phi(r), \Phi^*(r)$	Scalar group flux column vector of length $G$ and its adjoint.
$\underline{J}(r), \underline{J}^*(r)$	Vector group current column vector of length $G$ and its adjoint.

$\Lambda(r)$	$G \times G$ group material removal, scattering, and production matrix.
$D(r)$	$G \times G$ diagonal group diffusion coefficient matrix.
$U(r), U^*(r)$	Scalar group flux and weighting flux trial function column vectors of length $G$ .
$\underline{V}(r), \underline{V}^*(r)$	Vector group current and weighting current trial function column vectors of length $G$ .
$k$	One-dimensional spatial index which runs from the leftmost first region to the rightmost $K$ -th region, as $k = 1$ to $K$ .
$z$	The one-dimensional axis variable divided into $K$ regions such that each region $k$ is bounded by nodes $z_k$ and $z_{k+1}$ .
$x$	A dimensionless variable defined in each region $k$ as $x = (z - z_k)/(z_{k+1} - z_k)$ , such that $0 \leq x \leq 1$ as $z_k \leq z \leq z_{k+1}$ .
$F_k$	Approximate one-dimensional group flux solution at node $z_k$ .
$G_k$	Approximate one-dimensional group current solution at node $z_k$ .
$\psi_k(z), \psi_k^*(z)$	Detailed one-dimensional subassembly flux and weighting flux solutions in coarse mesh region $k$ whose form is linear within each homogeneous subassembly interval.
$\eta_k(z), \eta_k^*(z)$	Detailed one-dimensional subassembly current and weighting current solutions in coarse mesh region $k$ whose form is <u>constant</u> within each homogeneous subassembly interval.

$\tilde{\eta}_k(z), \tilde{\eta}_k^*(z)$	Detailed one-dimensional subassembly current and weighting current solutions in coarse mesh region k whose form is <u>linear</u> within each homogeneous subassembly interval.
$\mathbb{A}$	Discretized matrix form of the $G \times G$ group diffusion, absorption, and scattering matrices.
$\mathbb{B}$	Discretized matrix form of the $G \times G$ group fission-production matrix.
$\underline{F}$	The unknown approximate group flux solution vector which may contain group current unknowns.
$\% \lambda$	Normalized eigenvalue percent error: $\% \lambda = (\lambda_{\text{Reference}} - \lambda_{\text{Method}}) / \lambda_{\text{Reference}} \times 100\%.$
$P(k)$	Fractional power produced in coarse mesh region k when the total power produced has been normalized to unity.
$\% P(k)$	Normalized fractional power percent error: $\% P(k) = [P(k)_{\text{Reference}} - P(k)_{\text{Method}}] / P(k)_{\text{Reference}} \times 100\%.$

## Appendix B

DIFFERENCE EQUATION COEFFICIENTS RESULTING FROM  
USE OF THE FINITE ELEMENT  
APPROXIMATION METHODS

The  $G \times G$  matrix coefficients resulting from the conventional finite difference approximation, the linear finite element approximation, and the cubic Hermite finite element approximation in one-dimensional multigroup diffusion theory are defined below in sections B.1, B.2, and B.3, respectively. The coefficients are given in terms of assumed homogeneous regional nuclear constants through the use of the  $G \times G$  group matrices  $\mathbb{D}_k$  and  $\mathbf{\Lambda}_k$ , where  $\mathbf{\Lambda}_k = \mathbb{M}_k - \mathbf{T}_k - \frac{1}{\lambda} \mathbb{B}_k$ , which are defined in Chapter 2 and are constant for each region  $k$ , where  $k = 1$  to  $K$ .

More general definitions of these coefficients may be found from the coefficients resulting from the use of the proposed approximations, given in Appendix C, by requiring that  $\psi_k(z)$  be constant and  $\eta_k(z)$  be zero in each region  $k$ .

B.1. Coefficients of the Conventional Finite Difference Equations  
(as defined by Eqs. 2.16)

Interior Coefficients;  $k = 2$  to  $K$ :

$$a_k = -\mathbb{D}_{k-1}/h_{k-1}$$

$$b_k = \frac{1}{2}(\mathbf{\Lambda}_{k-1}h_{k-1} + \mathbf{\Lambda}_k h_k) + \mathbb{D}_{k-1}/h_{k-1} + \mathbb{D}_k/h_k$$

$$c_k = -\mathbb{D}_k/h_k$$

Symmetry Boundary Condition Coefficients:

$$b_1 = \frac{1}{2} \mathbf{\Lambda}_1 h_1 + \mathbb{D}_1/h_1$$

$$c_1 = -\mathbb{D}_1/h_1$$

$$a_{K+1} = -\mathbb{D}_K/h_K$$

$$b_{K+1} = \frac{1}{2} \mathbf{\Lambda}_K h_K + \mathbb{D}_K/h_K$$

B. 2. Coefficients of the Linear Finite Element Method Equations  
(as defined by Eqs. 2.30)

Interior Coefficients;  $k = 2$  to  $K$ :

$$a_k = \frac{1}{6} \mathbf{\Lambda}_{k-1} h_{k-1} - \mathbb{D}_{k-1}/h_{k-1}$$

$$b_k = \frac{1}{3} [\mathbf{\Lambda}_{k-1} h_{k-1} + \mathbf{\Lambda}_k h_k] + \mathbb{D}_{k-1}/h_{k-1} + \mathbb{D}_k/h_k$$

$$c_k = \frac{1}{6} \mathbf{\Lambda}_k h_k - \mathbb{D}_k/h_k$$

Symmetry Boundary Condition Coefficients:

$$b_1 = \frac{1}{3} \mathbf{\Lambda}_1 h_1 + \mathbb{D}_1/h_1$$

$$c_1 = \frac{1}{6} \mathbf{\Lambda}_1 h_1 - \mathbb{D}_1/h_1$$

$$a_{K+1} = \frac{1}{6} \mathbf{\Lambda}_K h_K - \mathbb{D}_K/h_K$$

$$b_{K+1} = \frac{1}{3} \mathbf{\Lambda}_K h_K + \mathbb{D}_K/h_K$$

B.3. Coefficients of the Cubic Hermite Finite Element Method Equations (as defined by Eq. 2.33)

Interior Coefficients;  $k = 2$  to  $K$ :

$$a1_k = \frac{9}{70} \mathbf{\Lambda}_{k-1} h_{k-1} - \frac{6}{5} \mathbb{D}_{k-1} / h_{k-1}$$

$$a2_k = \left( -\frac{13}{420} \mathbf{\Lambda}_{k-1} h_{k-1}^2 \mathbb{D}_{k-1}^{-1} + \frac{1}{10} \right) \theta$$

$$b1_k = \frac{13}{35} \left( \mathbf{\Lambda}_{k-1} h_{k-1} + \mathbf{\Lambda}_k h_k \right) + \frac{6}{5} \left( \mathbb{D}_{k-1} / h_{k-1} + \mathbb{D}_k / h_k \right)$$

$$b2_k = \frac{11}{210} \left( \mathbf{\Lambda}_{k-1} h_{k-1}^2 \mathbb{D}_{k-1}^{-1} - \mathbf{\Lambda}_k h_k^2 \mathbb{D}_k^{-1} \right) \theta$$

$$c1_k = \frac{9}{70} \mathbf{\Lambda}_k h_k - \frac{6}{5} \mathbb{D}_k / h_k$$

$$c2_k = \left( \frac{13}{420} \mathbf{\Lambda}_k h_k^2 \mathbb{D}_k^{-1} - \frac{1}{10} \right) \theta$$

$$a3_k = \left( \frac{13}{420} \mathbb{D}_{k-1}^{-1} h_{k-1}^2 \mathbf{\Lambda}_{k-1} - \frac{1}{10} \right) \theta$$

$$a4_k = \left( -\frac{1}{140} \mathbb{D}_{k-1}^{-1} h_{k-1}^3 \mathbf{\Lambda}_{k-1} \mathbb{D}_{k-1}^{-1} - \frac{1}{30} h_{k-1} \mathbb{D}_{k-1}^{-1} \right) \theta^2$$

$$b3_k = \frac{11}{210} \left( \mathbb{D}_{k-1}^{-1} h_{k-1}^2 \mathbf{\Lambda}_{k-1} - \mathbb{D}_k^{-1} h_k^2 \mathbf{\Lambda}_k \right) \theta$$

$$b4_k = \left[ \frac{1}{105} \left( \mathbb{D}_{k-1}^{-1} h_{k-1}^3 \mathbf{\Lambda}_{k-1} \mathbb{D}_{k-1}^{-1} + \mathbb{D}_k^{-1} h_k^3 \mathbf{\Lambda}_k \mathbb{D}_k^{-1} \right) + \frac{2}{15} \left( h_{k-1} \mathbb{D}_{k-1}^{-1} + h_k \mathbb{D}_k^{-1} \right) \right] \theta^2$$

$$c3_k = \left( -\frac{13}{420} \mathbb{D}_k^{-1} h_k^2 \mathbf{\Lambda}_k + \frac{1}{10} \right) \theta$$

$$c4_k = \left( -\frac{1}{140} \mathbb{D}_k^{-1} h_k^3 \mathbf{\Lambda}_k \mathbb{D}_k^{-1} - \frac{1}{30} h_k \mathbb{D}_k^{-1} \right) \theta^2$$

Zero Flux Boundary Condition Coefficients:

$$b4_1 = \left( \frac{1}{105} \mathbb{D}_1^{-1} h_1^3 \Lambda_1 \mathbb{D}_1^{-1} + \frac{2}{15} h_1 \mathbb{D}_1^{-1} \right) \theta^2$$

$$c3_1 = \left( -\frac{13}{420} \mathbb{D}_1^{-1} h_1^2 \Lambda_1 + \frac{1}{10} \right) \theta$$

$$c4_1 = \left( -\frac{1}{140} \mathbb{D}_1^{-1} h_1^3 \Lambda_1 \mathbb{D}_1^{-1} - \frac{1}{30} h_1 \mathbb{D}_1^{-1} \right) \theta^2$$

$$a3_{K+1} = \left( \frac{13}{420} \mathbb{D}_K^{-1} h_K^2 \Lambda_K - \frac{1}{10} \right) \theta$$

$$a4_{K+1} = \left( -\frac{1}{140} \mathbb{D}_K^{-1} h_K^3 \Lambda_K \mathbb{D}_K^{-1} - \frac{1}{30} h_K \mathbb{D}_K^{-1} \right) \theta^2$$

$$b4_{K+1} = \left( \frac{1}{105} \mathbb{D}_K^{-1} h_K^3 \Lambda_K \mathbb{D}_K^{-1} + \frac{2}{15} h_K \mathbb{D}_K^{-1} \right) \theta^2$$

Symmetry Boundary Condition Coefficients:

$$b1_1 = \frac{13}{35} \Lambda_1 h_1 + \frac{6}{5} \mathbb{D}_1 / h_1$$

$$c1_1 = \frac{9}{70} \Lambda_1 h_1 - \frac{6}{5} \mathbb{D}_1 / h_1$$

$$c2_1 = \left( \frac{13}{420} \Lambda_1 h_1^2 \mathbb{D}_1^{-1} - \frac{1}{10} \right) \theta$$

$$a1_{K+1} = \frac{9}{70} \Lambda_K h_K - \frac{6}{5} \mathbb{D}_K / h_K$$

$$a2_{K+1} = \left( -\frac{13}{420} \Lambda_K h_K^2 \mathbb{D}_K^{-1} + \frac{1}{10} \right) \theta$$

$$b1_{K+1} = \frac{13}{35} \Lambda_K h_K + \frac{6}{5} \mathbb{D}_K / h_K$$



## Appendix C

DIFFERENCE EQUATION COEFFICIENTS RESULTING FROM  
USE OF THE PROPOSED APPROXIMATION METHODS

The  $G \times G$  matrix coefficients resulting from the proposed approximation methods using (1) linear basis functions, and (2) cubic Hermite basis functions, in one-dimensional multigroup diffusion theory are defined below in sections C.1 and C.2, respectively.

The coefficients are given as integrands of functions of  $x$  where the integration of every coefficient-integrand over a region

$$\text{Coefficient} = \int_0^1 [\text{Coefficient-Integrand}(x)] dx$$

is understood.

In order to simplify the forms of the coefficient-integrands, it is convenient to define the following  $G \times G$  matrices:

$$\begin{aligned} \mathbb{K}_k(x) &= \psi_k^*(x) \mathbb{A}_k(x) h_k \psi_k(x) \\ \mathbb{L}_k(x) &= \eta_k^*(x) \mathbb{D}_k^{-1}(x) h_k \eta_k(x) \\ \mathbb{P}_k(x) &= \psi_k^*(x) \eta_k(x) \\ \mathbb{Q}_k(x) &= \eta_k^*(x) \psi_k(x) \\ \mathbb{R}_k(x) &= \frac{1}{h_k} \psi_k^*(x) \mathbb{D}_k(x) \psi_k(x) \end{aligned}$$

for each region  $k$ . In each approximation below, two sets of polynomial functions  $p_1(x) \dots p_{2N}(x)$  and  $q_1(x) \dots q_{2N}(x)$  are given

which represent the basis functions of the approximation and their negative derivatives, where  $N = 1$  for the linear basis function approximations and  $N = 2$  for the cubic Hermite basis function approximations.

The  $G \times G$  coefficient-integrands are then listed in terms of these matrices and polynomials by the  $G \times G$  collapsed matrices  $\mathbb{E}_k^{i,j}(x)$  defined as

$$\begin{aligned} \mathbb{E}_k^{i,j}(x) = & p_i(x)p_j(x)\mathbb{K}_k(x) - p_i(x)p_j(x)\mathbb{L}_k(x) + q_i(x)p_j(x)\mathbb{P}_k(x) \\ & - p_i(x)q_j(x)\mathbb{Q}_k(x) + q_i(x)q_j(x)\mathbb{R}_k(x) \end{aligned}$$

for given values of  $i$  and  $j$  for each region  $k$ , where  $k = 1$  to  $K$ . It should be noted that  $\mathbb{E}_k^{i,j}(x)$  is not symmetric about  $i$  and  $j$ ; i. e.:

$$\mathbb{E}_k^{i,j}(x) \neq \mathbb{E}_k^{j,i}(x), \quad \text{for } i \neq j.$$

### C.1. Coefficient-Integrands of the Proposed Approximation Method Equations Using Linear Basis Functions (as defined by Eq. 3.6)

These coefficients are given in terms of the polynomial functions

$$p_1(x) = (1-x)$$

$$p_2(x) = x$$

$$q_1(x) = 1$$

$$q_2(x) = -1$$

for use in the  $\mathbb{E}_k^{i,j}(x)$  below.

Interior Coefficient-Integrands;  $k = 2$  to  $K$ :

$$a_k(x) = \psi_{k-1}^{*T(1)} \mathbb{E}_{k-1}^{2,1}(x) \psi_{k-1}^{-1}(0)$$

$$b_k(x) = \psi_{k-1}^{*T(1)} \mathbb{E}_{k-1}^{2,2}(x) \psi_{k-1}^{-1}(1) + \psi_k^{*T(0)} \mathbb{E}_k^{1,1}(x) \psi_k^{-1}(0)$$

$$c_k(x) = \psi_k^{*T(0)} \mathbb{E}_k^{1,2}(x) \psi_k^{-1}(1)$$

Symmetry Coefficient-Integrands:

$$b_1(x) = \psi_1^{*T(0)} \mathbb{E}_1^{1,1}(x) \psi_1^{-1}(0)$$

$$c_1(x) = c_k(x); \text{ where } k = 1$$

$$a_{K+1}(x) = a_k(x); \text{ where } k = K + 1$$

$$b_{K+1}(x) = \psi_K^{*T(1)} \mathbb{E}_K^{2,2}(x) \psi_K^{-1}(1)$$

Implied Zero Flux Boundary Condition Coefficient-Integrands

(Corresponding with the modified trial functions of the type in Eqs. 3.9)

$$b_2(x) = \psi_1^{*T(1)} [\mathbb{K}_1(x) - \mathbb{L}_1(x)] \psi_1^{-1}(1) + \psi_2^{*T(0)} \mathbb{E}_2^{1,1}(x) \psi_2^{-1}(0)$$

$$b_K(x) = \psi_{K-1}^{*T(1)} \mathbb{E}_{K-1}^{2,2}(x) \psi_{K-1}^{-1}(1) + \psi_K^{*T(0)} [\mathbb{K}_K - \mathbb{L}_K] \psi_K^{-1}(0)$$

C. 2. Coefficient-Integrands of the Proposed Approximation Method  
Equations Using Cubic Hermite Basis Functions (as defined  
in Eq. 3.16)

These coefficients are given in terms of the polynomials  $p_1(x)$  through  $p_4(x)$  and  $q_1(x)$  through  $q_4(x)$ , previously defined in Eqs. 3.11 and 3.12, for use in the  $\mathbb{E}_k^{i,j}(x)$  below.

Interior Coefficient-Integrands;  $k = 2$  to  $K$ :

$$a1_k(x) = \psi_{k-1}^{*-1}(1) \mathbb{E}_{k-1}^{2,1}(x) \psi_{k-1}^{-1}(0)$$

$$a2_k(x) = \psi_{k-1}^{*-1}(1) \mathbb{E}_{k-1}^{2,3}(x) \psi_{k-1}^{-1}(0) \mathbb{D}_{k-1}^{-1}(0) \theta$$

$$b1_k(x) = \psi_{k-1}^{*-1}(1) \mathbb{E}_{k-1}^{2,2}(x) \psi_{k-1}^{-1}(1) + \psi_k^{*-1}(0) \mathbb{E}_k^{1,1}(x) \psi_k^{-1}(0)$$

$$b2_k(x) = \psi_{k-1}^{*-1}(1) \mathbb{E}_{k-1}^{2,4}(x) \psi_{k-1}^{-1}(1) \mathbb{D}_{k-1}^{-1}(1) \theta$$

$$+ \psi_k^{*-1}(0) \mathbb{E}_k^{1,3}(x) \psi_k^{-1}(0) \mathbb{D}_k^{-1}(0) \theta$$

$$c1_k(x) = \psi_k^{*-1}(0) \mathbb{E}_k^{1,2}(x) \psi_k^{-1}(1)$$

$$c2_k(x) = \psi_k^{*-1}(0) \mathbb{E}_k^{1,4}(x) \psi_k^{-1}(1) \mathbb{D}_k^{-1}(1) \theta$$

$$a3_k(x) = \psi_{k-1}^{*-1}(1) \mathbb{D}_{k-1}^{-1}(1) \mathbb{E}_{k-1}^{4,1}(x) \psi_{k-1}^{-1}(0) \theta$$

$$a4_k(x) = \psi_{k-1}^{*-1}(1) \mathbb{D}_{k-1}^{-1}(1) \mathbb{E}_{k-1}^{4,3}(x) \psi_{k-1}^{-1}(0) \mathbb{D}_{k-1}^{-1}(0) \theta^2$$

$$\begin{aligned}
b3_k(x) &= \psi_{k-1}^{*T(1)} \mathbb{D}_{k-1}^{-1}(1) \mathbb{E}_{k-1}^{4,2}(x) \psi_{k-1}^{-1}(1) \theta \\
&\quad + \psi_k^{*T(0)} \mathbb{D}_k^{-1}(0) \mathbb{E}_k^{3,1}(x) \psi_k^{-1}(0) \theta \\
b4_k(x) &= \psi_{k-1}^{*T(1)} \mathbb{D}_{k-1}^{-1}(1) \mathbb{E}_{k-1}^{4,4}(x) \psi_{k-1}^{-1}(1) \mathbb{D}_{k-1}^{-1}(1) \theta^2 \\
&\quad + \psi_k^{*T(0)} \mathbb{D}_k^{-1}(0) \mathbb{E}_k^{3,3}(x) \psi_k^{-1}(0) \mathbb{D}_k^{-1}(0) \theta^2 \\
c3_k(x) &= \psi_k^{*T(0)} \mathbb{D}_k^{-1}(0) \mathbb{E}_k^{3,2}(x) \psi_k^{-1}(1) \theta \\
c4_k(x) &= \psi_k^{*T(0)} \mathbb{D}_k^{-1}(0) \mathbb{E}_k^{3,4}(x) \psi_k^{-1}(1) \mathbb{D}_k^{-1}(1) \theta^2
\end{aligned}$$

Zero Flux Boundary Condition Coefficient-Integrands:

$$\begin{aligned}
b4_1(x) &= \psi_1^{*T(0)} \mathbb{D}_1^{-1}(0) \mathbb{E}_1^{3,3}(x) \psi_1^{-1}(0) \mathbb{D}_1^{-1}(0) \theta^2 \\
c3_1(x) &= c3_k(x); \text{ where } k = 1 \\
c4_1(x) &= c4_k(x); \text{ where } k = 1 \\
a3_{K+1}(x) &= a3_k(x); \text{ where } k = K+1 \\
a4_{K+1}(x) &= a4_{K+1}(x); \text{ where } k = K+1 \\
b4_{K+1}(x) &= \psi_K^{*T(1)} \mathbb{D}_K^{-1}(1) \mathbb{E}_K^{4,4}(x) \psi_K^{-1}(1) \mathbb{D}_K^{-1}(1) \theta^2
\end{aligned}$$

Symmetry Boundary Condition Coefficient-Integrands:

$$b1_1(x) = \psi_1^{*\text{T}(0)} \mathbb{E}_1^{1,1}(x) \psi_1^{-1}(0)$$

$$c1_1(x) = c1_k(x); \text{ where } k = 1$$

$$c2_1(x) = c2_k(x); \text{ where } k = 1$$

$$a1_{K+1}(x) = a1_k(x); \text{ where } k = K+1$$

$$a2_{K+1}(x) = a2_k(x); \text{ where } k = K+1$$

$$b1_{K+1}(x) = \psi_K^{*\text{T}(1)} \mathbb{E}_K^{2,2}(x) \psi_K^{-1}(1)$$

Implied Zero Flux Boundary Condition Coefficient-Integrands  
(corresponding with the modified trial functions of the type in  
Eq. 3.17):

$$\left. \begin{aligned} c3_1(x) &= c3_k(x); \text{ where } k = 1 \\ a2_2(x) &= a2_k(x); \text{ where } k = 2 \\ b1_2(x) &= b1_k(x); \text{ where } k = 2 \\ b2_2(x) &= b2_k(x); \text{ where } k = 2 \\ b3_2(x) &= b3_k(x); \text{ where } k = 2 \end{aligned} \right\} \text{ and where } \begin{cases} p_1(x) = 0 \\ p_2(x) = 1 \\ q_1(x) = 0 \\ q_2(x) = 0 \end{cases}$$

$$\left. \begin{aligned} b1_K(x) &= b1_k(x); \text{ where } k = K \\ b2_K(x) &= b2_k(x); \text{ where } k = K \\ c2_K(x) &= c2_k(x); \text{ where } k = K \\ b3_K(x) &= b3_k(x); \text{ where } k = K \\ a3_{K+1}(x) &= a3_k(x); \text{ where } k = K+1 \end{aligned} \right\} \text{ and where } \begin{cases} p_1(x) = 1 \\ p_2(x) = 0 \\ q_1(x) = 0 \\ q_2(x) = 0 \end{cases}$$

## Appendix D

## DESCRIPTION OF THE COMPUTER PROGRAMS

The computer programs REF2G, LINEAR, CUBIC, and ANALYZE are described respectively in the following four sections. The programs are written in FORTRAN IV, allow double precision calculations, and were used with the I.B.M. 360/65 and 370/155 FORTRAN G compilers at the M.I.T. Information Processing Center. Sample storage requirements and execution times of the programs are summarized in Table D. 1.

The power method employed in the first three programs allows a maximum of 300 iterations to converge, and program execution continues after this limit. Initial group flux shapes are sinusoidal or flat, depending upon the boundary conditions chosen.

The input and output data of each program are divided into data blocks for ease of representation as described below.

#### D.1. Description of Program REF2G

REF2G finds the reference solutions of the one-dimensional, two-group diffusion equations of each case study, or the detailed subassembly solutions of each subassembly, using the linear finite element approximation method. The program allows up to a total of two hundred homogeneous fine mesh regions and employs combinations of both zero flux and symmetry boundary conditions. Identical material regions can be automatically repeated with no additional input.

Options for plotting graphically the history of the converging spatial flux as well as the converging eigenvalue are also available. In addition, the program allows the calculation of the adjoint flux and current solutions.

The approximate current solutions are linear within each mesh region and are calculated from the converged flux solutions using Eqs. 4.29 and 4.31. The converged flux and current solutions and the converged adjoint solutions can be punched out for future use as described below.

#### A. Reference Solution Input Block

Card Type 1: Format (20A4)

An Appropriate Problem Title

Card Type 2: Format (2I5, 3E10.3, 5I5)

KR    Total number of homogeneous fine mesh regions.  
KR ≤ 200.

IBC    Boundary Condition Option

1. Zero flux on both boundaries
2. Zero flux on the left, symmetry on the right
3. Symmetry on the left, zero flux on the right
4. Symmetry on both boundaries

EPS1    Iteration tolerance to be met by differences between elements of successive iteration solution vectors:

$$|F_j^{(i)} - F_j^{(i-1)}| < \epsilon_1; \text{ for all } j$$

EPS2    Iteration tolerance to be met by the mean square error between successive iteration solution vectors:

$$\left\{ \sum_j [F_j^{(i)} - F_j^{(i-1)}]^2 \right\}^{\frac{1}{2}} < \epsilon_2$$



Table D.1. Sample Storage Requirements and Execution Times of the Programs for Two-Group Results. Obtained using the M.I.T. I.B.M. 360/155.

Storage Requirements in Bytes (without overlays):

REF2G:	260 K
LINEAR:	200 K
CUBIC:	250 K
ANALYZE:	205 K

C.P.U. Execution Times in Minutes:

REF2G:	Detailed Subassembly Solutions (68 regions) <sup>a</sup> :	.120	
	Case 1 Reference Solution (150 regions) <sup>b</sup> :	.238	
	Case 4 Reference Solution (198 regions) <sup>b</sup> :	.644	
LINEAR:	Case 1 Synthesis (Homogenized <sup>c</sup> ) Method (3 regions):	.284	(.157)
	Case 4 Synthesis (Homogenized <sup>c</sup> ) Method (9 regions):	.296	(.209)
CUBIC:	Case 1 Synthesis (Homogenized <sup>c</sup> ) Method (3 regions):	1.227	(.183)
	Case 4 Synthesis (Homogenized <sup>c</sup> ) Method (9 regions):	1.464	(.328)
ANALYZE:	Case 1 Linear (Cubic Hermite) Basis Functions:	.087	(.108)
	Case 4 Linear (Cubic Hermite) Basis Functions:	.122	(.139)

a. Including adjoint flux and current calculations.

b. Not including adjoint calculations.

c. Including .126 minutes for calculation of the two-group homogenized constants.

- EPS3 Iteration tolerance to be met by the difference between successive iteration eigenvalues:
- $$|\lambda^{(i)} - \lambda^{(i-1)}| < \epsilon_3$$
- IPLOT Allows printed graphical display of the converging flux solution:
- 0 No display
  - 1 Plot only the resultant normalized flux
  - 2 Plot a normalized history of the converging flux
- JPLOT Allows printed graphical display of the history of the converging eigenvalue when JPLOT = 1.
- IPUNCH Allows punched output when IPUNCH = 1.
- ISEE Allows printing of storage information:
- 0 No information printed
  - 1 Input regional properties are printed
  - 2 Input regional properties as well as the Common/B5/ storage arrays and the Common/B3/ power method matrices are printed.
- NOADJ Adjoint calculations are performed when NOADJ = 0, and bypassed if NOADJ = 1.

Card Type 3: Format (25I2)

- ITF(k) The consecutive type-number of each region from left to right as k = 1 to KR. Allows for repeating identical regions with no additional input.

Card Type 3 is repeated KR/25 times (rounded off to the next highest integer).

Card Type 4: Format (2F10.5)

- CHI(1), CHI(2) The fission yields  $\chi_1$  and  $\chi_2$  for the fast and thermal groups, respectively.

\* An Input Region Data Block:

Repeated for each different material region;  $\max_k [\text{ITF}(k)]$  times.

Card Type 5: Format (I5)

k The consecutive mesh region number (counting from from left to right) for identification purposes.

Card Types 6, 7: Format (3F10.5, 4E10.3, /, 30X, 3E10.3)

The geometry and nuclear constants for region k:

z(1) Beginning spatial coordinate of region k (cm)  
 z(2) Ending spatial coordinate of region k (cm)  
 H Width of region k (cm)  
 A(1) Fast-group macroscopic total cross section in region k ( $\text{cm}^{-1}$ )  
 F(1) Fast-group macroscopic production cross section,  $\nu\Sigma_f$ , in region k ( $\text{cm}^{-1}$ )  
 D(1) Fast-group diffusion coefficient in region k (cm)  
 S Fast-to-thermal macroscopic scattering cross section in region k ( $\text{cm}^{-1}$ )  
 A(2) Thermal-group macroscopic total cross section in region k ( $\text{cm}^{-1}$ )  
 F(2) Thermal-group macroscopic production cross section,  $\nu\Sigma_f$ , in region k ( $\text{cm}^{-1}$ )  
 D(2) Thermal-group diffusion coefficient in region k (cm)

\* End of an Input Region Data Block.

\*\* Power Method Input Block: Optional

Card Type 8: Format (F10.5)

$\omega$  Outer iteration overrelaxation parameter  $1 \leq \omega \leq 2$ .  
 Default is  $\omega = 1.25$ .

Card Type 9: Format (D25.14)

$\lambda^{(0)}$  Initial eigenvalue guess. Default is  $\lambda^{(0)} = 1.0$ .

Card Type 10: Format (4E20.10)

((F(g, i), i = 1 to N), g = 1 to 2) Initial group flux solution guess without zero flux boundary values. Default is  $\underline{F} = 1.0$ .

\*\* End of Power Method Input Block.

#### B. Reference Solution Output Block

When IPUNCH = 1, REF2G punches out the number of fine mesh regions, KR, under Format (I5) followed by the converged flux solutions  $\psi(g, k)$  and corresponding current solutions  $\tilde{\eta}(g, k)$  for each group g and spatial node k including boundary conditions. When adjoint calculations are included, the results are punched out under Format (4D20.10) as

$$((\psi(g, k), \tilde{\eta}(g, k), \psi^*(g, k), \tilde{\eta}^*(g, k), k = 1 \text{ to } KR + 1), g = 1 \text{ to } 2)$$

where the notation denotes case reference solutions as well as detailed subassembly solutions. When the adjoint calculations have been bypassed, the results are punched out under Format (2D20.10) as

$$((\psi(g, k), \tilde{\eta}(g, k), k = 1 \text{ to } KR + 1), g = 1 \text{ to } 2)$$

A total of  $2 \text{ KR} + 3$  cards are punched out.

#### D.2. Description of Program LINEAR

Program LINEAR forms and solves the difference equations resulting from the proposed approximation method using the linear basis functions. The program allows up to twenty-five coarse mesh regions, each of which is allowed to be broken into not more than one hundred homogeneous intervals. Combinations of both zero flux and

symmetry boundary conditions as well as use of the modified trial function forms in the boundary regions are allowed. Spatial flux and eigenvalue iteration history plots are also available.

The program allows a choice of the type of weighting, Galerkin or adjoint, to be used in the approximation. Also, either form of the detailed subassembly current solutions  $\eta_k(x)$  or  $\tilde{\eta}_k(x)$  is allowed. In addition, identical coarse mesh regions with identical detailed subassembly solutions can be repeated implicitly.

LINEAR also calculates results of the linear finite element method when suitable input is used. Such results can be obtained by using homogenized coarse mesh region nuclear constants and defining the detailed group flux solutions to be constant and the detailed currents to be zero (or by setting  $ITC = 0$ ).

Punched results using detailed subassembly solutions constitute a Synthesis Method Output Block, while punched output resulting from the reduction to the finite element method with homogenized regional constants constitutes a Homogenized Method Output Block.

#### A. Homogenized or Synthesis Method Input Block

Undefined input parameters are identical to those previously defined in the REF2G input.

Card Type 1: Format (20A4)

An Appropriate Problem Title

Card Type 2: Format (2I5, 3E10.3, 6I5)

KR Total number of coarse mesh regions.  $KR \leq 25$ .

IBC 1 - 4 As previously defined  
 5 Modified trial function (no tilting) in the first region, symmetry on the right  
 6 Zero flux on the left, modified trial function in the last region  
 7 Modified trial functions in both boundary regions

EPS1  
 EPS2  
 EPS3  
 IPLOT  
 JPLOT  
 IPUNCH  
 ISEE

ITW Type of approximation weighting desired:  
 0 Flux (Galerkin)  
 1 Adjoint

ITC Form of the detailed current solutions in all sub-assemblies:  
 0  $\eta_k(x), \eta_k^*(x)$  as calculated by Fick's laws  
 1  $\tilde{\eta}_k(x), \tilde{\eta}_k^*(x)$  as given from REF2G output

Card Type 3: Format (25I2)

ITF(k) The consecutive type-number of each coarse mesh region from left to right as  $k=1$  to  $KR$ . Allows for repeating identical subassemblies with no additional input.

Card Type 3 is repeated  $KR/25$  times (rounded off to the next highest integer).

Card Type 4: Format (2F10.5)

CHI(1), CHI(2)

\* An Input Subassembly Region Data Block:

Repeated for each different coarse mesh region;  $\max_k [\text{ITF}(k)]$   
times.

Card Type 5: Format (2I5)

- k The consecutive coarse mesh region number  
(from left to right).
- N The number of homogeneous intervals in sub-  
assembly k.  $N \leq 100$ .

Card Types 6, 7: Format (3F10.5, 4E10.3, /, 30X, 3E10.3)

The subassembly geometry and nuclear constants  
within each interval corresponding to the detailed  
subassembly solutions.

Repeated for each interval as  $i = 1$  to  $N$ :

- z(i) Beginning spatial coordinate of interval  $i$  (cm)
- z(i+1) Ending spatial coordinate of interval  $i$  (cm)
- H(i) Width of interval  $i$  (cm)
- A(1,i) Fast-group macroscopic total cross section in  
interval  $i$  ( $\text{cm}^{-1}$ )
- F(1,i) Fast-group macroscopic production cross section,  
 $\nu\Sigma_f$ , in interval  $i$  ( $\text{cm}^{-1}$ )
- D(1,i) Fast-group diffusion coefficient in interval  $i$  (cm)
- S(i) Fast-to-thermal macroscopic scattering cross  
section in interval  $i$  ( $\text{cm}^{-1}$ )
- A(2,i) Thermal-group macroscopic total cross section  
in interval  $i$  ( $\text{cm}^{-1}$ )
- F(2,i) Thermal-group macroscopic production cross  
section,  $\nu\Sigma_f$ , in interval  $i$  ( $\text{cm}^{-1}$ )
- D(2,i) Thermal-group diffusion coefficient in interval  $i$  (cm)

Card Type 8: Format (4D20.10)

The detailed subassembly solutions.

$((\psi(g,k), \tilde{\eta}(g,k), \psi^*(g,k), \tilde{\eta}^*(g,k), k=1 \text{ to } KR+1), g=1 \text{ to } 2)$

A subassembly's Reference Solution Output Block without the first card.

\* END of an Input Subassembly Region Data Block.

\*\* Expected Solution Input Block: Optional

Card Type 9: Format (D25.14)

$\lambda_{\text{REF}}$  Expected eigenvalue solution. Default is  $\lambda_{\text{REF}} = 1.0$ .

Card Type 10: Format (4E20.10)

$((F(i,g), i=1 \text{ to } N), g=1 \text{ to } 2)$  Expected group flux solution without zero flux boundary values. Default is  $\underline{F} = 1.0$ .

\*\* END of the Expected Solution Input Block.

\*\*\* Power Method Input Block: Optional

As previously defined in the REF2G input.

\*\*\* END of the Power Method Input Block.

## B. Homogenized or Synthesis Method Output Block

When IPUNCH = 1, LINEAR punches out the total number of coarse mesh regions, KR, under Format (I5) followed by the resultant flux solutions including boundary conditions. The flux solutions are punched out under Format (2E20.7) as

$(F(1,k), F(2,k), k=1 \text{ to } KR+1)$

These cards represent either a Homogenized or Synthesis Method Output Block, depending upon the type and form of input data used.



### D.3 Description of Program CUBIC

Program CUBIC forms and solves the difference equations resulting from the proposed approximation method using the cubic Hermite basis functions. The program is very similar in form to program LINEAR and uses similar input.

#### A. Homogenized or Synthesis Method Input Block

The input to CUBIC is identical to that of LINEAR except for the following:

1. The boundary condition options are restricted by  $1 \leq IBC \leq 4$ .
2. The normalization constant  $\theta$  can be included on Card Type 4 after CHI(2) under Format (3F10.5). Default is  $\theta = 1.0$ .
3. Both the expected group solutions and the initial group solutions of the Expected Solution and Power Method Input Blocks, respectively, are of the form  $((F(g, i), i = 1 \text{ to } N), g = 1 \text{ to } 2)$  without either zero flux or zero current (or symmetry) boundary conditions. The solution vector is made up of alternating flux and current values as described in section 3.3 of Chapter 3. Default values are flux values of unity and current values of zero.

#### B. Homogenized or Synthesis Method Output Block

When IPUNCH = 1, CUBIC punches out the total number of coarse mesh regions, KR, under Format (I5) followed by the resultant flux and current solutions including boundary conditions. The solutions are punched out under Format (4E20.7) as

$$(F(1, k), F(2, k), G(1, k), G(2, k), k = 1 \text{ to } KR + 1)$$

where  $F(g,k)$  represents the flux, and  $G(g,k)$  the current solution of group  $g$  at node  $k$ .

As in the case of LINEAR, these KR+2 output cards represent either a Synthesis or Homogenized Method Output Block, depending upon the type and form of input data used.

#### D.4 Description of Program ANALYZE

ANALYZE compares the results of the reference solution, homogenized finite element method, and the proposed synthesis method for each case study where either linear or cubic Hermite basis functions have been used in the latter methods. For each of these three methods, the program first forms the complete detailed flux solution and then normalizes the flux distributions for each method such that their total power levels are unity. The fractional (normalized) power levels produced in each coarse mesh region are then calculated, compared, and listed. Finally, the detailed group fluxes of each method are plotted graphically relative to one another using the Stromberg-Carlson Computer Recorder, SC-4020, facility at M.I.T.<sup>54</sup> The graphic results for each group are normalized by the largest group-flux value such that the equivalent total power levels are preserved.

##### A. ANALYZE Input

The input to ANALYZE is read from five device units: 1, 2, 3, 11, 12, 13, and 5. Input and output data of the reference and approximation programs are read from the former six units while the standard input unit, 5, is reserved for SC-4020 plotting information.

The input is described by "Header Cards" and previously defined Input and Output Blocks. Header cards consist of one or more cards defined as follows:

Header Card 1: Format (4I5)

Method	Indicates the type of basis function approximation:
	1 Linear
	2 Cubic Hermite
NK	Total number of coarse mesh regions involved.
NR	Total number of fine mesh regions involved. NR = NK except for reference solution calculations.
NAP	Number of additional points to be plotted within each coarse mesh region. Used with the homogenized finite element method calculations. NAP < 0 denotes that the additional points are to be used in the first region (reflector) only.

Header Card 2: For use in device unit 3 input when NR ≠ NK.  
Format (16I5).

NRNK(k)	The number of fine mesh regions which make up each coarse mesh region k, as k = 1 to NK.
---------	--

The program is dimensioned to accept up to 200 fine mesh regions (or intervals) per coarse mesh region, up to 25 coarse mesh regions, and up to a grand total of 1000 fine mesh regions in each case study.

The form of the ANALYZE input is given as follows:

Input Data for Unit 1:

Header Cards

[ Homogenized Method Input Block ]

Input Data for Unit 2:

Header Cards

[ Synthesis Method Input Block ]

Input Data for Unit 3:

Header Cards

[ Reference Solution Input Block ]

Input Data for Unit 11:

[ Homogenized Method Output Block ]

Input Data for Unit 12:

[ Synthesis Method Output Block ]

Input Data for Unit 13:

[ Reference Solution Output Block ]

Input Data for Unit 5:

No SC-4020 plots are generated if this data is omitted.

Card 1: Format (20A4)

An appropriate title written above each plotted graph.

Card 2: Format (2F10.5)

XINCH      Total width of the graph in inches including labels  
(limited to 7.45").

YINCH      Total height of the graph in inches including labels  
(limited to 7.45").

Card 3: Format (I10, F10.5)

NCELL      Total number of coarse mesh regions.  $NCELL \leq 25$ .  
( $NCELL < 0$  indicates that the last region is of width  
 $\frac{1}{2}$  WCELL.)

WCELL      Width of each coarse mesh region in cm.

Card 4: Optional. Format (I10, 7F10.5)

NLL      Number of vertical light lines to be added to the plotted  
graphs.  $NLL \leq 100$ .

XL(i)      Spatial location (cm) of the light lines;  $i = 1$  to 7.

Card 5: Format (8F10.5)

XL(i)      As above when  $NLL > 7$ .

## Appendix E

SAMPLE INPUT AND OUTPUT DATA BLOCKS  
FOR PROGRAMS REF2G, LINEAR, CUBIC, AND ANALYZE

(Included in only the first six copies of this report.)

E.1. REF2G SAMPLE INPUT AND OUTPUT DATA BLOCKS

SAMPLE REF2G REFERENCE SOLUTION INPUT BLOCK:

CASE 1 STUDY: THREE DIFFERENT SUBASSEMBLIES. 150 FINE MESH REFERENCE SOLUTION.

```

150 4 1.E-5 1.E-5 1.E-8 1 1 1 0 1
1 1 1 1 2 2 2 2 3 3 3 3 4 4 4 4 5 5 5 5 5 5 5 5
5 5 5 5 5 5 5 4 4 4 4 3 3 3 3 2 2 2 2 1 1 1 1 1
6 6 6 6 6 7 7 7 7 8 8 8 8 9 9 9 9 10 10 10 10 10 10 10
10 10 10 10 10 10 10 9 9 9 9 8 8 8 8 7 7 7 7 6 6 6 6 6
11 11 11 11 11 12 12 12 12 13 13 13 13 14 14 14 14 15 15 15 15 15 15 15
15 15 15 15 15 15 14 14 14 14 13 13 13 13 12 12 12 12 11 11 11 11 11 11
1.0 0.0
1 1
0.0 1.0 1.0 2.59 D-2 4.85 D-3 1.396 D 0 1.79 D-2
5.32 D-2 6.36 D-2 3.88 D-1
6 1
0.0 0.5 0.5 2.59 D-2 4.85 D-3 1.396 D 0 1.79 D-2
5.32 D-2 6.36 D-2 3.88 D-1
10 1
0.0 0.25 0.25 2.59 D-2 4.85 D-3 1.396 D 0 1.79 D-2
5.32 D-2 6.36 D-2 3.88 D-1
14 1
0.0 0.125 0.125 2.59 D-2 4.85 D-3 1.396 D 0 1.79 D-2
5.32 D-2 6.36 D-2 3.88 D-1
18 1
0.0 0.0625 0.0625 4.52 D-2 0.0 D 0 1.0 D 0 0.0 D 0
9.59 D-1 0.0 D 0 1.0 D 0
51 1
0.0 1.0 1.0 2.60 D-2 5.53 D-3 1.397 D 0 1.72 D-2
7.10 D-2 1.02 D-1 3.89 D-1
56 1
0.0 0.5 0.5 2.60 D-2 5.53 D-3 1.397 D 0 1.72 D-2
7.10 D-2 1.02 D-1 3.89 D-1

```

60	1										
0.0	0.25	0.25	2.60	D-2	5.53	D-3	1.397	D 0	1.72	D-2	
			7.10	D-2	1.02	D-1	3.89	D-1			
64	1										
0.0	0.125	0.125	2.60	D-2	5.53	D-3	1.397	D 0	1.72	D-2	
			7.10	D-2	1.02	D-1	3.89	D-1			
68	1										
0.0	0.0625	0.0625	4.52	D-2	0.0	D 0	1.0	D 0	0.0	D 0	
			9.59	D-1	0.0	D 0	1.0	D 0			
101	1										
0.0	1.0	1.0	2.61	D-2	6.59	D-3	1.399	D 0	1.68	D-2	
			8.32	D-2	1.29	D-1	3.87	D-1			
106	1										
0.0	0.5	0.5	2.61	D-2	6.59	D-3	1.399	D 0	1.68	D-2	
			8.32	D-2	1.29	D-1	3.87	D-1			
110	1										
0.0	0.25	0.25	2.61	D-2	6.59	D-3	1.399	D 0	1.68	D-2	
			8.32	D-2	1.29	D-1	3.87	D-1			
114	1										
0.0	0.125	0.125	2.61	D-2	6.59	D-3	1.399	D 0	1.68	D-2	
			8.32	D-2	1.29	D-1	3.87	D-1			
118	1										
0.0	0.0625	0.0625	4.52	D-2	0.0	D 0	1.0	D 0	0.0	D 0	
			9.59	D-1	0.0	D 0	1.0	D 0			



SAMPLE REF2G SUBASSEMBLY SOLUTION INPUT BLOCK:

-----  
 CELL A TWO GROUJP CELL SOLUTION: FUEL A + CRUCIFORM ROD. TWO GROUP CONSTANTS.

68	4	1.E-8	1.E-8	1.E-8	1	1	1	0	0		
1 2 3 3 3 3 4 4 5 5 5 5 5 5 6 6 6 6 7 7 7 7 7 7 7											
7 8 8 8 8 8 8 8 8 8 8 8 8 8 8 8 7 7 7 7 7 7 7 7 7											
6 6 6 6 5 5 5 5 5 5 4 4 3 3 3 3 2 1											
1.0	0.0										
1	1										
0.0		0.0625	0.0625	2.59	D-2	4.85	D-3	1.396	D 0	1.79	D-2
				5.32	D-2	6.36	D-2	3.88	D-1		
2	1										
0.0		0.9375	0.9375	2.59	D-2	4.85	D-3	1.396	D 0	1.79	D-2
				5.32	D-2	6.36	D-2	3.88	D-1		
3	1										
0.0		1.0	1.0	2.59	D-2	4.85	D-3	1.396	D 0	1.79	D-2
				5.32	D-2	6.36	D-2	3.88	D-1		
7	1										
0.0		0.5	0.5	2.59	D-2	4.85	D-3	1.396	D 0	1.79	D-2
				5.32	D-2	6.36	D-2	3.88	D-1		
9	1										
0.0		0.25	0.25	2.59	D-2	4.85	D-3	1.396	D 0	1.79	D-2
				5.32	D-2	6.36	D-2	3.88	D-1		
15	1										
0.0		0.125	0.125	2.59	D-2	4.85	D-3	1.396	D 0	1.79	D-2
				5.32	D-2	6.36	D-2	3.88	D-1		
19	1										
0.0		0.0625	0.0625	2.59	D-2	4.85	D-3	1.396	D 0	1.79	D-2
				5.32	D-2	6.36	D-2	3.88	D-1		
27	1										
0.0		0.0625	0.0625	4.52	D-2	0.0	D 0	1.0	D 0	0.0	D 0
				9.59	D-1	0.0	D 0	1.0	D 0		

SAMPLE REF2G REFERENCE SOLUTION OUTPUT BLOCK (SAMPLE SUBASSEMBLY SOLUTION):

---

68

0.10000000	01	0.0	D 0	0.6840883D	00	0.0	D 0
0.99999060	00		0.2062319D-02	0.6840819D	00		-0.1410808D-02
0.99760400	00		0.6351434D-02	0.6824493D	00		-0.4344942D-02

63 ADDITIONAL FAST GROUP DATA CARDS

0.99760400	00		-0.6351434D-02	0.6824493D	00		0.4344942D-02
0.99999060	00		-0.2062319D-02	0.6840819D	00		0.1410808D-02
0.10000000	01	0.0	D 0	0.6840883D	00	0.0	D 0
0.30901450	00	0.0	D 0	0.10000000	01	0.0	D 0
0.30900720	00		0.4484312D-03	0.9999762D	00		-0.1451165D-02
0.30711030	00		0.1418028D-02	0.9938377D	00		-0.4588871D-02

63 ADDITIONAL THERMAL GROUP DATA CARDS

0.30711030	00		-0.1418028D-02	0.9938377D	00		0.4588871D-02
0.30900720	00		-0.4484312D-03	0.9999762D	00		0.1451165D-02
0.30901450	00	0.0	D 0	0.10000000	01	0.0	D 0

E.2. LINEAR SAMPLE INPUT AND OUTPUT DATA BLOCKS

SAMPLE LINEAR JR CUBIC HOMOGENIZED METHOD INPUT BLOCK:

CASE 1 STUDY: THREE DIFFERENT SUBASSEMBLIES. HOMOGENIZED FINITE ELEMENT METHOD

3	4	1.E-5	1.E-5	1.E-8	1	1	1	0	1	1
1 2 3										
1.0	0.0									
1	1									
0.0	18.0	18.0			.268878D-1.460175D-2.137952D	1.169837D-1				
					.681283D-1.625518D-1.398086D	0				
	1.0	D 0			0.0	D 0	1.0	D 0	0.0	D 0
	1.0	D 0			0.0	D 0	1.0	D 0	0.0	D 0
	1.0	D 0			0.0	D 0	1.0	D 0	0.0	D 0
	1.0	D 0			0.0	D 0	1.0	D 0	0.0	D 0
2	1									
0.0	18.0	18.0			.269849D-1.524631D-2.137948D	1.163176D-1				
					.864780D-1.100222D	0.399649D	0			
	1.0	D 0			0.0	D 0	1.0	D 0	0.0	D 0
	1.0	D 0			0.0	D 0	1.0	D 0	0.0	D 0
	1.0	D 0			0.0	D 0	1.0	D 0	0.0	D 0
	1.0	D 0			0.0	D 0	1.0	D 0	0.0	D 0
3	1									
0.0	18.0	18.0			.270820D-1.625116D-2.137943D	1.159361D-1				
					.989361D-1.126682D	0.398014D	0			
	1.0	D 0			0.0	D 0	1.0	D 0	0.0	D 0
	1.0	D 0			0.0	D 0	1.0	D 0	0.0	D 0
	1.0	D 0			0.0	D 0	1.0	D 0	0.0	D 0
	1.0	D 0			0.0	D 0	1.0	D 0	0.0	D 0

SAMPLE LINEAR METHOD OUTPUT BLOCK:

3

0.2444087D 00	0.6758041D-01
0.4060866D 00	0.9039887D-01
0.7838067D 00	0.1339626D 00
0.1000000D 01	0.1563839D 00

E.3. CUBIC SAMPLE INPUT AND OUTPUT DATA BLOCKS

SAMPLE LINEAR OR CUBIC PROPOSED SYNTHESIS METHOD INPUT BLOCK:

CASE 1 STUDY: THREE DIFFERENT SUBASSEMBLIES. SUBASSEMBLY SYNTHESIS.  
 3 4 1.E-5 1.E-5 1.E-8 1 1 1 0 1 1  
 1 2 3  
 1.0 0.0  
 1 68  
 0.0 0.06250 0.06250 0.259E-01 0.485E-02 0.140E 01 0.179E-01  
 0.532E-01 0.636E-01 0.388E 00  
 0.06250 1.00000 0.93750 0.259E-01 0.485E-02 0.140E 01 0.179E-01  
 0.532E-01 0.636E-01 0.388E 00

64 ADDITIONAL CARD PAIRS OF TYPE A SUBASSEMBLY INTERVAL MATERIAL INPUT

17.00000 17.93750 0.93750 0.259E-01 0.485E-02 0.140E 01 0.179E-01  
 0.532E-01 0.636E-01 0.388E 00  
 17.93750 18.00000 0.06250 0.259E-01 0.485E-02 0.140E 01 0.179E-01  
 0.532E-01 0.636E-01 0.388E 00  
 0.1000000D 01 0.0 D 0 0.6840883D 00 0.0 D 0  
 0.9999906D 00 0.2062319D-02 0.6840819D 00 -0.1410808D-02  
 0.9976040D 00 0.6351434D-02 0.6824493D 00 -0.4344942D-02

63 ADDITIONAL FAST GROUP DATA CARDS

0.9976040D 00 -0.6351434D-02 0.6824493D 00 0.4344942D-02  
 0.9999906D 00 -0.2062319D-02 0.6840819D 00 0.1410808D-02  
 0.1000000D 01 0.0 D 0 0.6840883D 00 0.0 D 0  
 0.3090145D 00 0.0 D 0 0.1000000D 01 0.0 D 0  
 0.3090072D 00 0.4484312D-03 0.9999762D 00 -0.1451165D-02  
 0.3071103D 00 0.1418028D-02 0.9938377D 00 -0.4588871D-02

63 ADDITIONAL THERMAL GROUP DATA CARDS

	0.30711030 00	-0.14180280-02	0.99383770 00	0.45888710-02
1	0.30900720 00	-0.44843120-03	0.99997620 00	0.14511650-02
	0.30901450 00	0.0	0.10000000 01	0.0
2	68			
0.0	0.06250	0.06250	0.260E-01 0.553E-02 0.140E 01	0.172E-01
			0.710E-01 0.102E 00 0.389E 00	
0.06250	1.00000	0.93750	0.260E-01 0.553E-02 0.140E 01	0.172E-01
			0.710E-01 0.102E 00 0.389E 00	

64 ADDITIONAL CARD PAIRS OF TYPE B SUBASSEMBLY INTERVAL MATERIAL INPUT

17.00000	17.93750	0.93750	0.260E-01 0.553E-02 0.140E 01	0.172E-01
			0.710E-01 0.102E 00 0.389E 00	
17.93750	18.00000	0.06250	0.260E-01 0.553E-02 0.140E 01	0.172E-01
			0.710E-01 0.102E 00 0.389E 00	
	0.10000000 01	0.0	0.65791580 00	0.0
	0.99999120 00	0.19386980-02	0.65791000 00	-0.12755000-02
	0.99774740 00	0.59789340-02	0.65643380 00	-0.39336350-02

63 ADDITIONAL FAST GROUP DATA CARDS

0.99774740 00	-0.59789340-02	0.65643380 00	0.39336350-02
0.99999120 00	-0.19386980-02	0.65791000 00	0.12755000-02
0.10000000 01	0.0	0.65791580 00	0.0
0.22944880 00	0.0	0.10000000 01	0.0
0.22944420 00	0.27926490-03	0.99998010 00	-0.12171120-02
0.22826300 00	0.88680160-03	0.99483190 00	-0.38649210-02

63 ADDITIONAL THERMAL GROUP DATA CARDS

0.22826300 00	-0.88680160-03	0.99483190 00	0.38649210-02
0.22944420 00	-0.27926490-03	0.99998010 00	0.12171120-02
0.22944880 00	0.0	0.10000000 01	0.0
3	68		
0.0	0.06250	0.06250	0.261E-01 0.659E-02 0.140E 01
			0.168E-01

0.06250 1.00000 0.93750 0.832E-01 0.129E 00 0.387E 00  
 0.261E-01 0.659E-02 0.140E 01 0.168E-01  
 0.832E-01 0.129E 00 0.387E 00

64 ADDITIONAL CARD PAIRS OF TYPE C SUBASSEMBLY INTERVAL MATERIAL INPUT

17.00000 17.93750 0.93750 0.261E-01 0.659E-02 0.140E 01 0.168E-01  
 0.832E-01 0.129E 00 0.387E 00  
 17.93750 18.00000 0.06250 0.261E-01 0.659E-02 0.140E 01 0.168E-01  
 0.832E-01 0.129E 00 0.387E 00  
 0.1000000D 01 0.0 D 0 0.6615058D 00 0.0 D 0  
 0.9999916D 00 0.1837899D-02 0.6615002D 00 -0.1215781D-02  
 0.9978665D 00 0.5673004D-02 0.6600944D 00 -0.3752725D-02

63 ADDITIONAL FAST GROUP DATA CARDS

0.9978665D 00 -0.5673004D-02 0.6600944D 00 0.3752725D-02  
 0.9999916D 00 -0.1837899D-02 0.6615002D 00 0.1215781D-02  
 0.1000000D 01 0.0 D 0 0.6615058D 00 0.0 D 0  
 0.1937796D 00 0.0 D 0 0.1000000D 01 0.0 D 0  
 0.1937762D 00 0.2081898D-03 0.9999824D 00 -0.1074364D-02  
 0.1928897D 00 0.6627953D-03 0.9954074D 00 -0.3420356D-02

63 ADDITIONAL THERMAL GROUP DATA CARDS

0.1928897D 00 -0.6627953D-03 0.9954074D 00 0.3420356D-02  
 0.1937762D 00 -0.2081898D-03 0.9999824D 00 0.1074364D-02  
 0.1937796D 00 0.0 D 0 0.1000000D 01 0.0 D 0

SAMPLE CUBIC METHOD OUTPUT BLOCK:

3

0.2503191D 00	0.0	0.7648905D-01	0.0
0.4035480D 00	-0.2202874D-01	0.1062932D 00	0.5636188D-03
0.7785649D 00	-0.2731031D-01	0.1642353D 00	0.3579210D-03
0.1000000D 01	0.0	0.1946404D 00	0.0



E.4. ANALYZE SAMPLE INPUT AND OUTPUT

-----  
SAMPLE ANALYZE INPUT (CASE 1 CUBIC METHODS RESULTS):  
-----

```
//G.FT01F001 DD *
  2   3   3   35
CASE 1 HOMOGENIZED LINEAR FINITE ELEMENT METHOD INPUT BLOCK
.
/*

//G.FT02F001 DD *
  2   3   3   0
CASE 1 LINEAR SYNTHESIS METHOD INPUT BLOCK
.
/*

//G.FT03F001 DD *
  1   3 150   0
 50  50  50
CASE 1 REFERENCE SOLUTION INPUT BLOCK
.
/*

//G.FT11F001 DD *
CASE 1 HOMOGENIZED LINEAR FINITE ELEMENT METHOD OUTPUT BLOCK
.
/*
```

//G.FT12F001 DD \*  
CASE 1 LINEAR SYNTHESIS METHOD OUTPUT BLOCK

.  
.\*

//G.FT13F001 DD \*  
CASE 1 REFERENCE SOLUTION OUTPUT BLOCK

.  
.\*

//G.SYSIN DD \*  
TWO GROUP CASE 1 CUBIC RESULTS.

6.0	3	18.0					
	6	8.5	9.5	26.5	27.5	44.5	45.5

.\*

ANALYZE PRINTED OUTPUT: Case 1 with Linear Basis Functions.

RESULTS OF THE INTEGRATED POWER IN EACH OF THE 3 REGIONS:

CALCULATED POWER LEVELS, AND NUMBER OF SUBREGIONS PER REGION:

REGION:	HOMOGENIZED RESULTS:		SYNTHESIZED RESULTS:		REFERENCE RESULTS:	
1	1	0.1158776E 00	68	0.1086056E 00	50	0.1024175E 00
2	1	0.2585564E 00	68	0.2447225E 00	50	0.2404295E 00
3	1	0.4313925E 00	68	0.4112111E 00	50	0.4187305E 00
TOTALS:	3	0.8058265E 00	204	0.7645392E 00	150	0.7615775E 00

FRACTIONAL POWER LEVELS:

REGION:	HOMOGENEOUS RESULTS:		SYNTHESIZED RESULTS:		REFERENCE RESULTS:	
1		0.1437997E 00		0.1420537E 00		0.1344807E 00
2		0.3208586E 00		0.3200914E 00		0.3156993E 00
3		0.5353416E 00		0.5378548E 00		0.5498199E 00
TOTALS:		0.9999999E 00		0.9999999E 00		0.9999999E 00

FRACTIONAL POWER NORMALIZED PERCENT ERRORS:

REGION:	(REF-HOMO)/REF %	(REF-SYNTH)/REF %	(SYNTN-HOMO)/SYNTH %
1	-0.6929624E 01	-0.5631252E 01	-0.1229154E 01
2	-0.1634251E 01	-0.1391243E 01	-0.2396725E 00
3	0.2633273E 01	0.2176184E 01	0.4672581E 00

ANALYZE PRINTED OUTPUT: Case 1 with Linear Basis Functions.

EXECUTING GENERAL ANALYSIS AND FLUX PLOTTING PROGRAM:

TITLE OF PLOTTING RUN IS: | THREE DIFFERENT SUBASSEMBLYS PROBLEM. |

REACTOR GEOMETRY PARAMETERS:

NCELL = 3

WCELL = 18.00000

XMIN = 0.0

XMAX = 54.00000

YMIN = 0.0

YMAX = 1.00000

NLL = 6

(XL(I),I=1,NLL) =

8.50000 9.50000 26.50000 27.50000 44.50000 45.50000

## Appendix F

## SOURCE LISTINGS OF THE PROGRAMS

FORTTRAN source listings of programs REF2G, LINEAR, CUBIC, and ANALYZE are listed in only the first six copies of this report in the following four sections.

A figure of a subroutine overlay structure precedes each listing in order to indicate the construction of each program.

F.1. SOURCE LISTING of Program REF2G

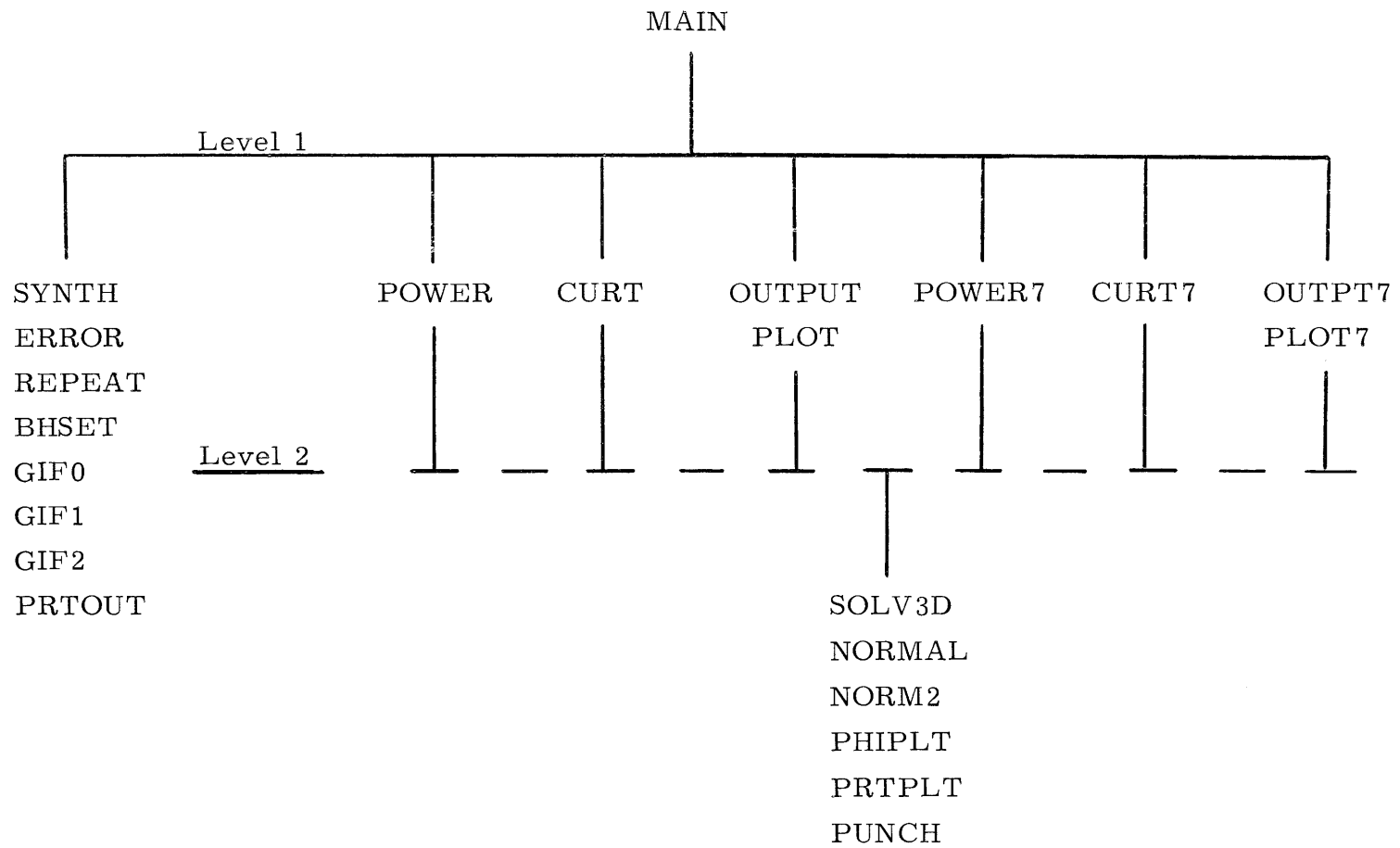


Figure F.1. Structure of Program REF2G.

```

C PROGRAM REF2G:
C TWO GROUP DETAILED REFERENCE AND SUBASSEMBLY SOLUTION PROGRAM.
C
CALL TIMING(I1)
CALL SYNTH
CALL TIMING(I2)
CALL POWER
CALL TIMING(I3)
CALL CURT
CALL TIMING(I4)
CALL OUTPUT
CALL TIMING(I5)
CALL POWER7
CALL TIMING(I6)
CALL CURT7
CALL TIMING(I7)
CALL OUTPUT7
CALL TIMING(I8)
C TIMING EXECUTION
WRITE (6,30)
30 FORMAT (1H1,'TIMING PROGRAM EXECUTION:',/)
J=I2-I1
WRITE(6,701) J
J=I3-I2
WRITE (6,702) J
J=I4-I3
WRITE (6,703) J
J=I5-I4
WRITE (6,704) J
J=I6-I5
WRITE (6,705) J
J=I7-I6
WRITE (6,706) J
J=I8-I7
WRITE(6,707) J
701 FORMAT (1H , ' SYNTH HAS TAKEN',I6,' /100 SECONDS.')
```

```

REF20001
REF20002
REF20003
REF20004
REF20005
REF20006
REF20007
REF20008
REF20009
REF20010
REF20011
REF20012
REF20013
REF20014
REF20015
REF20016
REF20017
REF20018
REF20019
REF20020
REF20021
REF20022
REF20023
REF20024
REF20025
REF20026
REF20027
REF20028
REF20029
REF20030
REF20031
REF20032
REF20033
REF20034
REF20035
REF20036
```



```
702 FORMAT (1H , ' POWER HAS TAKEN',I6,' /100 SECONDS.')
703 FORMAT (1H , ' CURT HAS TAKEN',I6,' /100 SECONDS.')
704 FORMAT (1H , ' OUTPUT HAS TAKEN',I5,' /100 SECONDS.')
705 FORMAT (1H , ' POWER7 HAS TAKEN',I6,' /100 SECONDS.')
706 FORMAT (1H , ' CURT7 HAS TAKEN',I6,' /100 SECONDS.')
707 FORMAT (1H , ' JUTPUT7 HAS TAKEN',I5,' /100 SECONDS.')
      CALL TIMING(I20)
      J=I20-I1
      WRITE(6,720) J
720 FORMAT (1H0,' THIS RUN HAS TAKEN',I6,' /100 SECONDS TO RUN.')
      STOP
      END
```

```
REF20037
REF20038
REF20039
REF20040
REF20041
REF20042
REF20043
REF20044
REF20045
REF20046
REF20047
REF20048
```

	SUBROUTINE SYNTH	SYNT0001
	LINEAR FINITE ELEMENT METHOD:	SYNT0002
C	*****	SYNT0003
C	ADJOINT QUANTITIES OF VARIABLES ARE DENOTED BY 7 RATHER THAN *.	SYNT0004
C	THUS: PHI7 (RATHER THAN PHI*) IS THE ADJOINT OF PHI. ETC.	SYNT0005
	IMPLICIT REAL*8 (A-H,K-Z)	SYNT0006
	COMMON /B1/ IBC,IPLT,JPLT,IPUNCH,ISEE,NOADJ	SYNT0007
	COMMON /B2/ KR,NN	SYNT0008
	COMMON /B3/ L1(201,3), L2(201,3), F1(201,3), F2(201,3), T(201,3)	SYNT0009
	COMMON /B5/ KA0(2,200),KA1(2,200),KA2(2,200),KB0(2,200),	SYNT0010
X	KB1(2,200),KB2(2,200),LA0(2,200),LA1(2,200),LA2(2,200),	SYNT0011
X	SR0(2,200),SR1(2,200),SR2(2,200),P(2,200),P1(2,200),	SYNT0012
X	Q(2,200),Q1(2,200),R(2,200),P0(2,200),P07(2,200),PH(2,200),	SYNT0013
X	PH7(2,200),AL(2,200),BL(2,200),CL(2,200),AF(2,200),BF(2,200),	SYNT0014
X	CF(2,200),AT(200),BT(200),CT(200),	SYNT0015
X	BLO(2),  CLO(2),  BFO(2),  CFO(2),  BTO(2),	SYNT0016
X	CTO(2),	SYNT0017
X	ALK(2),  BLK(2),  AFK(2),  BFK(2),  ATK(2),	SYNT0018
X	BTK(2)	SYNT0019
	COMMON /B7/ HH(200),DD(2,200)	SYNT0020
	COMMON /CHIF/ CHI(2)	SYNT0021
	COMMON /BH/ X(2), H(1)	SYNT0022
	COMMON /ER/ EPS1,EPS2,EPS3	SYNT0023
	DIMENSION PHI(2,2),PHI7(2,2),CUR(2,2),CUR7(2,2),	SYNT0024
X	A(2,1),F(2,1),D(2,1),S(2,1),DI(2,1),XU(2,2)	SYNT0025
	DIMENSION ITF(200), KTF(200)	SYNT0026
	REAL TITLE(20)	SYNT0027
	INTEGER KR,K,KS,KS1,KR0,NN	SYNT0028
	INTEGER NUMITF, KTF,NOADJ	SYNT0029
	READ (5,200) TITLE	SYNT0030
200	FORMAT (20A4)	SYNT0031
	WRITE (6,201) TITLE	SYNT0032
201	FORMAT (1H1,20A4,//)	SYNT0033
C	READ IN THE NUMBER OF REGION TRIAL FUNCTIONS AND TYPE OF B.C.S.	SYNT0034
C	AS WELL AS THE TOLERANCES AND THE OUTPUT TYPES DESIRED.	SYNT0035
	READ (5,1) KR,IBC,EPS1,EPS2,EPS3,IPLT,JPLT,IPUNCH,ISEE,NOADJ	SYNT0036

```

1 FORMAT (2I5,3D10.3,5I5)
C   READ IN THE TYPE-NUMBER OF EACH TF REGION:
   READ (5,100) (ITF(I),I=1,KR)
100 FORMAT (25I2)
C   READ IN THE FISSION YIELD FOR EACH GROUP:
   READ (5,101) CHI(1), CHI(2)
101 FORMAT (2F10.5)
   KRO=KR-1
   WRITE (6,2) KR, IBC
2 FORMAT ('OVARIATIONAL SYNTHESIS PROGRAM #2G(200):',5X,'USING ',I3,
X ' SUBREACTOR REGIONS, OR TRIAL FUNCTIONS.',/,
X ' OBOUNDRY CONDITION NUMBER (IBC) IS ',I1,'.',/,/,
X ' OMATERIAL PROPERTIES AND TRIAL FUNCTIONS FOR EACH SUBREGION FO
X LLOW:',/,
X ' OMATERIAL PROPERTIES ARE HOMOGENEOUS IN THE INDICATED REGIONS.
X',/,
X ' OFLUX TRIAL FUNCTIONS ARE LINEAR IN EACH SEGMENT OF THE SUBREG
X IONS.',/,
X ' OCURRENT TRIAL FUNCTIONS ARE FLAT IN EACH OF THE ',
X ' SUBREGIONS.')
   WRITE (6,20) EPS1, EPS2, EPS3, IPLOT, JPLOT, IPUNCH, ISEE, NOADJ
20 FORMAT (//, 'OTOLERANCES TO POWER ARE : EPS1 = ',1PD10.3,/,
X 28X, 'EPS2 = ',1PD10.3,/,28X, 'EPS3 = ',1PD10.3,/,
X ' OOUTPUT PARAMETERS TO POWER ARE: IPLOT = ',I1,/,
X 34X, 'JPLOT = ',I1,/,34X, 'IPUNCH = ',I1,/,
X 34X, 'ISEE = ',I1,/,
X 34X, 'NOADJ = ',I1,'.')
   WRITE (6,22) CHI(1), CHI(2)
22 FORMAT (/, 'OFISSION YIELDS ARE: CHI(1) =',F10.5,/,
X 22X, 'CHI(2) =',F10.5)
   IF ((KR.LE.2).AND.(IBC.EQ.1)) CALL ERROR(1,KR)
   IF (KR.GT.200) CALL ERROR(2,KR)
   IF (EPS1.LT.1.0E-16) CALL ERROR(6,1)
   IF (EPS2.LT.1.0E-16) CALL ERROR(6,2)
   IF (EPS3.LT.1.0E-16) CALL ERROR(6,3)
   IF ((IBC.LT.1).OR.(IBC.GT.4)) CALL ERROR(7,IBC)

```

```

SYNT0037
SYNT0038
SYNT0039
SYNT0040
SYNT0041
SYNT0042
SYNT0043
SYNT0044
SYNT0045
SYNT0046
SYNT0047
SYNT0048
SYNT0049
SYNT0050
SYNT0051
SYNT0052
SYNT0053
SYNT0054
SYNT0055
SYNT0056
SYNT0057
SYNT0058
SYNT0059
SYNT0060
SYNT0061
SYNT0062
SYNT0063
SYNT0064
SYNT0065
SYNT0066
SYNT0067
SYNT0068
SYNT0069
SYNTC070
SYNT0071
SYNT0072

```

C	DUMMY NORMAL VECTOR: XU = UNITY. (FOR THE INTEGRATION FUNCTIONS.)	SYNT0073
	DO 21 IG=1,2	SYNT0074
	DO 21 II=1,2	SYNT0075
21	XU(IG,II)=1.0	SYNT0076
C	SET FLUXES TO UNITY FOR SYNTH 26:	SYNT0077
	DO 25 IG=1,2	SYNT0078
	DO 25 II=1,2	SYNT0079
	PHI(IG,II)=1.0	SYNT0080
25	PHI7(IG,II)=1.0	SYNT0081
C	COUNTER OF THE NUMBER OF TYPE-NUMBERS OF EACH TF REGION:	SYNT0082
	NUMITF=1	SYNT0083
	WRITE (6,9)	SYNT0084
9	FORMAT ('1')	SYNT0085
C	BEGIN TO READ IN THE TF REGION DATA AND FILL THE ARRAYS,	SYNT0086
C	DEPENDING ON THE TYPE-NUMBER OF EACH TF REGION.	SYNT0087
	DO 50 I=1,KR	SYNT0088
	IF (ITF(I).EQ.NUMITF) GO TO 110	SYNT0089
C	FILL THE ARRAYS FROM OLD TF REGION TYPES:	SYNT0090
	J=ITF(I)	SYNT0091
	CALL REPEAT(I,KTF(J))	SYNT0092
	GO TO 50	SYNT0093
C	READ IN THE TF REGION'S DATA FOR NEW TF REGION TYPE-NUMBERS:	SYNT0094
110	NUMITF=NUMITF+1	SYNT0095
	KTF(NUMITF-1)=I	SYNT0096
C	READ THE SUBREGION NUMBER AND THE NUMBER OF REGIONS IN THE SUBREGION.	SYNT0097
	READ (5,1) K	SYNT0098
	KS=1	SYNT0099
	IF (KS.GT.100) CALL ERROR(3,I)	SYNT0100
	KS1=KS+1	SYNT0101
C	CHECK FOR IMPROPER SEQUENCING OF INPUT DATA:	SYNT0102
	IF (I.NE.K) CALL ERROR(4,I)	SYNT0103
C	READ IN THE GEOMETRY AND THE MATERIAL PROPERTIES OF THIS REGION:	SYNT0104
	READ (5,3) (X(J),X(J+1),H(J),A(1,J),F(1,J),D(1,J),S(1,J),	SYNT0105
	X A(2,J),F(2,J),D(2,J),J=1,KS)	SYNT0106
3	FORMAT (3F10.5,4E10.3,/,30X,3E10.3)	SYNT0107
C	WRITING OUT THE INPUT INFORMATION:	SYNT0108

```

      IF (ISEE.EQ.0) GO TO 14
      WRITE (6,10) K,KR,KS,(J,X(J),X(J+1),H(J),A(1,J),F(1,J),D(1,J),
X     S(1,J),A(2,J),F(2,J),D(2,J),J=1,KS)
10 FORMAT ('INPUT MATERIAL PROPERTIES FOR SUBREGION NUMBER ',I3,
X     ', OF THE ',I3,' USED.',//,
X     5X,'THIS SUBREGION IS DIVIDED INTO ',I3,' HOMOGENEOUS SEGMENTS
XAS FOLLOWS:',//,
X     5X,'FAST GROUP CONSTANTS APPEAR FIRST:',//,
X     ' REGION #',5X,'INTERNAL BOUNDARIES',13X,'WIDTH',3X,
X     'ABSORB. CX (1/CM)',3X,'FISSION CX (1/CM)',6X,'DIFFUSION (CM)',
X     4X,'SCATT. CX (1/CM)',/,
X     5X,'I',11X,'X(I)',9X,'X(I+1)',11X,'H(I)',13X,'A(IG,I)',13X,
X     'F(IG,I)',13X,'D(IG,I)',14X,'S(1,I)',//,
X     (I6,3F15.4,4D20.8,/,51X,3D20.8))
C     END OF THE IN-OUT SECTION.
14 CONTINUE
C     DEFINING MISC. ARRAYS FOR THE INTEGRATION FUNCTIONS:
C     LEGNTH OF THE SUBREGION: HT
      HT=X(KS1)-X(1)
      HH(K)=HT
      DD(1,K)=D(1,1)
      DD(2,K)=D(2,1)
C     INVERSE OF D ARRAYS:
      DO 13 J=1,KS
      DI(1,J)=1./D(1,J)
13 DI(2,J)=1./D(2,J)
C     FORMATION OF THE INTEGRATION FUNCTIONS:
      CALL BHSET(KS)
C     DO FOR ALL ENERGY GROUPS:
      DO 50 IG=1,2
      KA0(IG,K)=GIFO(IG,PHI7,PHI,A,KS)
      KA1(IG,K)=GIF1(IG,PHI7,PHI,A,KS)
      KA2(IG,K)=GIF2(IG,PHI7,PHI,A,KS)
      KB0(IG,K)=GIFO(IG,PHI7,PHI,F,KS)
      KB1(IG,K)=GIF1(IG,PHI7,PHI,F,KS)
      KB2(IG,K)=GIF2(IG,PHI7,PHI,F,KS)

```

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SYNT0109
SYNT0110
SYNT0111
SYNT0112
SYNT0113
SYNT0114
SYNT0115
SYNT0116
SYNT0117
SYNT0118
SYNT0119
SYNT0120
SYNT0121
SYNT0122
SYNT0123
SYNT0124
SYNT0125
SYNT0126
SYNT0127
SYNT0128
SYNT0129
SYNT0130
SYNT0131
SYNT0132
SYNT0133
SYNT0134
SYNT0135
SYNT0136
SYNT0137
SYNT0138
SYNT0139
SYNT0140
SYNT0141
SYNT0142
SYNT0143
SYNT0144

```

	R(IG,K)=GIF0(IG,PHI7,PHI,D,KS)/(HT*HT)	SYNT0145
C	NO SCATTERING IN THE LOWEST GROUP:	SYNT0146
	IF (IG.EQ.2) GO TO 50	SYNT0147
	SR0(IG,K)=GIF0(IG,PHI7,PHI,S,KS)	SYNT0148
	SR1(IG,K)=GIF1(IG,PHI7,PHI,S,KS)	SYNT0149
	SR2(IG,K)=GIF2(IG,PHI7,PHI,S,KS)	SYNT0150
50	CONTINUE	SYNT0151
	NUMITF=NUMITF-1	SYNT0152
	WRITE (6,51) NUMITF	SYNT0153
51	FORMAT ('1THERE ARE ONLY',I3,' DIFFERENT TRIAL FUNCTION REGIONS.')	SYNT0154
	WRITE (6,52) (I,ITF(I),I=1,KR)	SYNT0155
52	FORMAT (/, 'OTABLE OF THE TRIAL FUNCTION REGION TYPES:',//,	SYNT0156
	X 3X, 'TF REGION',4X, 'REGION TYPE-NUMBER',//,	SYNT0157
	X (17,12X,I7))	SYNT0158
C	DETERMINATION OF THE B.C. OPTION PARAMETERS:	SYNT0159
C	NN IS THE MM AND FF MATRIX BLOCK SIZE.	SYNT0160
	IF (IBC.EQ.1) NN=KR-1	SYNT0161
	IF ((IBC.EQ.2).OR.(IBC.EQ.3)) NN=KR	SYNT0162
	IF (IBC.EQ.4) NN=KR+1	SYNT0163
C	FORMATION OF THE COEFFICIENT VECTORS:	SYNT0164
C	THE INTERIOR COEFFS:	SYNT0165
	DO 60 IG=1,2	SYNT0166
	DO 60 K=2,KR	SYNT0167
	J=K-1	SYNT0168
	AL(IG,K)=KA1(IG,J)-KA2(IG,J)-R(IG,J)	SYNT0169
	BL(IG,K)=KA2(IG,J)+R(IG,J)+KA0(IG,K)-2.*KA1(IG,K)+KA2(IG,K)	SYNT0170
	X +R(IG,K)	SYNT0171
	CL(IG,K)=KA1(IG,K)-KA2(IG,K)-R(IG,K)	SYNT0172
	AF(IG,K)=KB1(IG,J)-KB2(IG,J)	SYNT0173
	BF(IG,K)=KB2(IG,J)+KBO(IG,K)-2.*KB1(IG,K)+KB2(IG,K)	SYNT0174
	CF(IG,K)=KB1(IG,K)-KB2(IG,K)	SYNT0175
	AT(K)=SR1(1,J)-SR2(1,J)	SYNT0176
	BT(K)=SR2(1,J)+SR0(1,K)-2.*SR1(1,K)+SR2(1,K)	SYNT0177
	CT(K)=SR1(1,K)-SR2(1,K)	SYNT0178
60	CONTINUE	SYNT0179
C	THE ZERO FLUX COEFFS:	SYNT0180

	DO 61 IG=1,2	SYNT0181
	BLO(IG)=KA0(IG,1)-2.*KA1(IG,1)+KA2(IG,1)+R(IG,1)	SYNT0182
	BFO(IG)=KB0(IG,1)-2.*KB1(IG,1)+KB2(IG,1)	SYNT0183
	CLO(IG)=KA1(IG,1)-KA2(IG,1)-R(IG,1)	SYNT0184
61	CF0(IG)=KB1(IG,1)-KB2(IG,1)	SYNT0185
	BT0(1)=SR0(1,1)-2.*SR1(1,1)+SR2(1,1)	SYNT0186
	CT0(1)=SR1(1,1)-SR2(1,1)	SYNT0187
C	THE ZERO CURRENT COEFFS:	SYNT0188
	K=KR	SYNT0189
	DO 62 IG=1,2	SYNT0190
	ALK(IG)=KA1(IG,K)-KA2(IG,K)-R(IG,K)	SYNT0191
	BLK(IG)=KA2(IG,K)+R(IG,K)	SYNT0192
	AFK(IG)=KB1(IG,K)-KB2(IG,K)	SYNT0193
62	BFK(IG)=KB2(IG,K)	SYNT0194
	ATK(1)=SR1(1,K)-SR2(1,K)	SYNT0195
	BTK(1)=SR2(1,K)	SYNT0196
C	ZERO MATRICES:	SYNT0197
	L1(1,1)=0.	SYNT0198
	L2(1,1)=0.	SYNT0199
	F1(1,1)=0.	SYNT0200
	F2(1,1)=0.	SYNT0201
	T(1,1)=0.	SYNT0202
	L1(NN,3)=0.	SYNT0203
	L2(NN,3)=0.	SYNT0204
	F1(NN,3)=0.	SYNT0205
	F2(NN,3)=0.	SYNT0206
	T(NN,3)=0.	SYNT0207
C	FILL ALL THE MATRICES FOR POWER:	SYNT0208
	J=1	SYNT0209
C	DETERMINE THE LEFT BOUNDARY CONDITIONS:	SYNT0210
	IF (IBC.LT.3) GO TO 67	SYNT0211
	L1(J,2)=BLO(1)	SYNT0212
	L2(J,2)=BLO(2)	SYNT0213
	F1(J,2)=BFO(1)	SYNT0214
	F2(J,2)=BFO(2)	SYNT0215
	T(J,2)=BT0(1)	SYNT0216

L1(J, 3)=CLO(1)	SYNT0217
L2(J, 3)=CLO(2)	SYNT0218
F1(J, 3)=CF0(1)	SYNT0219
F2(J, 3)=CF0(2)	SYNT0220
T(J, 3)=CT0(1)	SYNT0221
J=J+1	SYNT0222
C FOR ALL THE INTERIOR EQUATIONS:	SYNT0223
67 DO 70 K=2,KR	SYNT0224
IF (J.EQ.1) GO TO 69	SYNT0225
L1(J, 1)=AL(1,K)	SYNT0226
L2(J, 1)=AL(2,K)	SYNT0227
F1(J, 1)=AF(1,K)	SYNT0228
F2(J, 1)=AF(2,K)	SYNT0229
T(J, 1)=AT(K)	SYNT0230
69 L1(J,2)=BL(1,K)	SYNT0231
L2(J,2)=BL(2,K)	SYNT0232
F1(J,2)=BF(1,K)	SYNT0233
F2(J,2)=BF(2,K)	SYNT0234
T(J,2)=BT(K)	SYNT0235
L1(J, 3)=CL(1,K)	SYNT0236
L2(J, 3)=CL(2,K)	SYNT0237
F1(J, 3)=CF(1,K)	SYNT0238
F2(J, 3)=CF(2,K)	SYNT0239
T(J, 3)=CT(K)	SYNT0240
J=J+1	SYNT0241
70 CONTINUE	SYNT0242
C DETERMINE THE RIGHT BOUNDARY CONDITIONS:	SYNT0243
IF ((IBC.EQ.1).OR.(IBC.EQ.3)) GO TO 80	SYNT0244
L1(J, 1)=ALK(1)	SYNT0245
L2(J, 1)=ALK(2)	SYNT0246
F1(J, 1)=AFK(1)	SYNT0247
F2(J, 1)=AFK(2)	SYNT0248
T(J, 1)=ATK(1)	SYNT0249
L1(J,2)=BLK(1)	SYNT0250
L2(J,2)=BLK(2)	SYNT0251
F1(J,2)=BFK(1)	SYNT0252



```
F2(J,2)=BFK(2)
T(J,2)=BTK(1)
80 CONTINUE
C PRINTS OUT THE SYNTH K ARRAYS, AND THE MATRICES GIVEN TO POWER
C FOR ISEE = 2.
IF (ISEE.EQ.2) CALL PRTOUT
RETURN
END
```

```
SYNT0253
SYNT0254
SYNT0255
SYNT0256
SYNT0257
SYNT0258
SYNT0259
SYNT0260
```

<pre> SUBROUTINE ERROR(I,J) ANNOUNCES INPUT ERRORS AND TERMINATES PROGRAM EXECUTION: GO TO (1,2,3,4,5,6,7,8,9),I 1 WRITE (6,101) GO TO 10 2 WRITE (6,102) J GO TO 10 3 WRITE (6,103) J GO TO 10 4 WRITE (6,104) J GO TO 10 5 WRITE (6,105) J GO TO 10 6 WRITE (6,106) J GO TO 10 7 CONTINUE 8 CONTINUE 9 CONTINUE 10 WRITE (6,110) 101 FORMAT ('1MUST HAVE &gt; 2 SUBREGIONS FOR ZERO FLUX B.C.S. INVALID.')</pre>	<pre> 102 FORMAT ('1NUMBER OF SUBREGIONS =',I3,' &gt; 25. INVALID.')</pre>	<pre> 103 FORMAT ('1SUBREGION NUMBER',I3,' HAS &gt; 25 SECTIONS. INVALID.')</pre>	<pre> 104 FORMAT ('1INPUT ERROR IN REGION SEQUENCING AT REGION',I5,'.')</pre>	<pre> 105 FORMAT ('1Z(I) = 0. IN REGION I =',I3,'. INVALID.')</pre>	<pre> 106 FORMAT ('1THE TOLERANCE: EPS',I1,' IS &lt; 1.0E-16. INVALID.')</pre>	<pre> 107 FORMAT ('1BOUNDARY CONDITION OPTION =',I2,' &lt; 1 OR &gt; 4. INVALID.')</pre>	<pre> 110 FORMAT (1H0,'PROBLEM TERMINATED.')</pre>	<pre> CALL EXIT RETURN END</pre>	<pre> ERRO0001 ERRO0002 ERRO0003 ERRO0004 ERRO0005 ERRO0006 ERRO0007 ERRO0008 ERRO0009 ERRO0010 ERRO0011 ERRO0012 ERRO0013 ERRO0014 ERRO0015 ERRO0016 ERRO0017 ERRO0018 ERRO0019 ERRO0020 ERRO0021 ERRO0022 ERRO0023 ERRO0024 ERRO0025 ERRO0026 ERRO0027 ERRO0028 ERRO0029 ERRO0030</pre>
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SUBROUTINE REPEAT(K,L)
C   SETS THE /B5/ ARRAYS (K) EQUAL TO PAST STORED ARRAYS (L):
      IMPLICIT REAL*8 (A-Z)
      COMMON /B5/ KA0(2,200),KA1(2,200),KA2(2,200),KB0(2,200),
X     KB1(2,200),KB2(2,200),LA0(2,200),LA1(2,200),LA2(2,200),
X     SR0(2,200),SR1(2,200),SR2(2,200),P(2,200),P1(2,200),
X     Q(2,200),Q1(2,200),R(2,200),P0(2,200),P07(2,200),PH(2,200),
X     PH7(2,200)
      COMMON /B7/ HH(200),DD(2,200)
      INTEGER K,L,G
      DO 10 G=1,2
        KA0(G,K)=KA0(G,L)
        KA1(G,K)=KA1(G,L)
        KA2(G,K)=KA2(G,L)
        KB0(G,K)=KB0(G,L)
        KB1(G,K)=KB1(G,L)
        KB2(G,K)=KB2(G,L)
        LA0(G,K)=LA0(G,L)
        LA1(G,K)=LA1(G,L)
        LA2(G,K)=LA2(G,L)
        IF (G.EQ.2) GO TO 5
        SR0(G,K)=SR0(G,L)
        SR1(G,K)=SR1(G,L)
        SR2(G,K)=SR2(G,L)
5     CONTINUE
        P(G,K)=P(G,L)
        P1(G,K)=P1(G,L)
        Q(G,K)=Q(G,L)
        Q1(G,K)=Q1(G,L)
        R(G,K)=R(G,L)
        P0(G,K)=P0(G,L)
        P07(G,K)=P07(G,L)
        PH(G,K)=PH(G,L)
        PH7(G,K)=PH7(G,L)
10    CONTINUE
      HH(K)=HH(L)

```

```

REPE0001
REPE0002
REPE0003
REPE0004
REPE0005
REPE0006
REPE0007
REPE0008
REPE0009
REPE0010
REPE0011
REPE0012
REPE0013
REPE0014
REPE0015
REPE0016
REPE0017
REPE0018
REPE0019
REPE0020
REPE0021
REPE0022
REPE0023
REPE0024
REPE0025
REPE0026
REPE0027
REPE0028
REPE0029
REPE0030
REPE0031
REPE0032
REPE0033
REPE0034
REPE0035
REPE0036

```

DD(1,K)=DD(1,L)  
DD(2,K)=DD(2,L)  
RETURN  
END

REPE0037  
REPE0038  
REPE0039  
REPE0040

```
C      SUBROUTINE BHSET(K)
        FIRST OF 4 ANALYTICAL INTEGRATION ROUTINES.
        IMPLICIT REAL*8 (A-H,L-Z)
        COMMON /BH/ X(2),H(1),H2(1),H3(1),H4(1),H5(1)
        DO 1 I=1,K
          H2(I)=X(I+1)**2-X(I)**2
          H3(I)=X(I+1)**3-X(I)**3
          H4(I)=X(I+1)**4-X(I)**4
          H5(I)=X(I+1)**5-X(I)**5
1      CONTINUE
        RETURN
        END
```

```
BHSE0001
BHSE0002
BHSE0003
BHSE0004
BHSE0005
BHSE0006
BHSE0007
BHSE0008
BHSE0009
BHSE0010
BHSE0011
BHSE0012
```

```
DOUBLE PRECISION FUNCTION GIFO(IG,Y,Z,C,K)
IMPLICIT REAL*8 (A-H,L-Z)
COMMON /BH/ X(2),H(1),H2(1),H3(1),H4(1),H5(1)
DIMENSION Y(2,2), Z(2,2), C(2,1)
SUM = 0.0
DO 1 I=1,K
SUM=C(IG,I)*(Y(IG,I)*Z(IG,I)*H(I)+H(I)*(Z(IG,I)*
X (Y(IG,I+1)-Y(IG,I))+Y(IG,I)*(Z(IG,I+1)-Z(IG,I)))/2.
X +H(I)*(Y(IG,I+1)-Y(IG,I))*(Z(IG,I+1)-Z(IG,I))/3.) + SUM
1 CONTINUE
GIFO = SUM
RETURN
END
```

```
GIF00001
GIF00002
GIF00003
GIF00004
GIF00005
GIF00006
GIF00007
GIF00008
GIF00009
GIF00010
GIF00011
GIF00012
GIF00013
```

DOUBLE PRECISION FUNCTION GIF1(IG,Y,Z,C,K)	GIF10001
IMPLICIT REAL*8 (A-H,L-Z)	GIF10002
COMMON /BH/ X(2),H(1),H2(1),H3(1),H4(1),H5(1)	GIF10003
DIMENSION Y(2,2), Z(2,2), C(2,1)	GIF10004
SUM = 0.0	GIF10005
DO 1 I=1,K	GIF10006
SUM=C(IG,I)*((H2(I)/2.-X(1)*H(I))*Y(IG,I)*Z(IG,I)	GIF10007
X  +(Z(IG,I)*(Y(IG,I+1)-Y(IG,I))+Y(IG,I)*	GIF10008
X  (Z(IG,I+1)-Z(IG,I))*(1./H(I))*(H3(I)/3.-H2(I)*(X(I)+X(1)))/2.	GIF10009
X  +X(I)*X(1)*H(I)+(Y(IG,I+1)-Y(IG,I))*(Z(IG,I+1)-Z(IG,I))	GIF10010
X  *(H4(I)/4.-H3(I)*(2.*X(I)+X(1))/3.+H2(I)*(X(I)*X(I)+2.*X(I)	GIF10011
X  *X(1))/2.-X(1)*X(I)*X(I)*H(I))/(H(I)*H(I))) + SUM	GIF10012
1 CONTINUE	GIF10013
GIF1 = SUM/(X(K+1)-X(1))	GIF10014
RETURN	GIF10015
END	GIF10016

DOUBLE PRECISION FUNCTION GIF2(IG,Y,Z,C,K)	GIF20001
IMPLICIT REAL*8 (A-H,L-Z)	GIF20002
COMMON /BH/ X(2),H(1),H2(1),H3(1),H4(1),H5(1)	GIF20003
DIMENSION Y(2,2), Z(2,2), C(2,1)	GIF20004
SUM = 0.0	GIF20005
DO 1 I=1,K	GIF20006
SUM=C(IG,I)*(Y(IG,I)*Z(IG,I)*(H3(I)/3.-X(1)*H2(I)+X(1)*X(1)*H(I))	GIF20007
X +(1./H(I))*(Z(IG,I)*(Y(IG,I+1)-Y(IG,I))+Y(IG,I)*(Z(IG,I+1)	GIF20008
X -Z(IG,I))*(H4(I)/4.-H3(I)*(2.*X(1)+X(I))/3.+H2(I)*(X(1)*X(1)	GIF20009
X +2.*X(1)*X(I))/2.-X(1)*X(1)*X(I)*H(I))+(1./(H(I)**2))	GIF20010
X *(Y(IG,I+1)-Y(IG,I))*(Z(IG,I+1)-Z(IG,I))*(H5(I)/5.	GIF20011
X -H4(I)*(X(1)+X(I))/2.+H3(I)*(X(1)*X(1)+4.*X(1)*X(I)+X(I)*X(I))	GIF20012
X /3.-H2(I)*(X(1)*X(I)*X(I)+X(1)*X(1)*X(I))	GIF20013
X +X(1)*X(1)*X(I)*X(I)*H(I)) + SUM	GIF20014
1 CONTINUE	GIF20015
GIF2 = SUM/((X(K+1)-X(1))**2)	GIF20016
RETURN	GIF20017
END	GIF20018



```

SUBROUTINE PRTOUT
PRINTS OUT THE /B5/ ARRAYS AND THE MATRICES GIVEN TO POWER:
IMPLICIT REAL*8 (A-H,K-Z)
COMMON /B2/ KR, N
COMMON /B3/ L1(201,3), L2(201,3), F1(201,3), F2(201,3), T(201,3)
COMMON /B5/ KAO(2,200),KAL(2,200),KA2(2,200),KBO(2,200),
X KB1(2,200),KB2(2,200),LA0(2,200),LA1(2,200),LA2(2,200),
X SR0(2,200),SR1(2,200),SR2(2,200),P(2,200),P1(2,200),
X Q(2,200),Q1(2,200),R(2,200),P0(2,200),P07(2,200),PH(2,200),
X PH7(2,200)
INTEGER KR, G, N
C KA AND KB ARRAYS:
WRITE (6,10)
10 FORMAT ('1 G',4X,'I',12X,'KAO(G,I)',12X,'KAL(G,I)',12X,
X 'KA2(G,I)',12X,'KBO(G,I)',12X,'KB1(G,I)',12X,'KB2(G,I)')
DO 11 G=1,2
WRITE (6,12)
11 WRITE (6,15) (G,I,KAO(G,I),KAL(G,I),KA2(G,I),KBO(G,I),KB1(G,I),
X KB2(G,I),I=1,KR)
12 FORMAT (' ')
15 FORMAT (2I5,6D20.7)
C LA AND SR ARRAYS:
WRITE (6,20)
20 FORMAT ('1 G',4X,'I',12X,'LA0(G,I)',12X,'LA1(G,I)',12X,
X 'LA2(G,I)',12X,'SR0(G,I)',12X,'SR1(G,I)',12X,'SR2(G,I)')
G=1
WRITE (6,12)
WRITE (6,15) (G,I,LA0(G,I),LA1(G,I),LA2(G,I),SR0(G,I),SR1(G,I),
X SR2(G,I),I=1,KR)
G=2
WRITE (6,12)
WRITE (6,25) (G,I,LA0(G,I),LA1(G,I),LA2(G,I),I=1,KR)
25 FORMAT (2I5,3D20.7)
C P, Q, AND R ARRAYS:
WRITE (6,30)
30 FORMAT ('1 G',4X,'I',14X,'P(G,I)',13X,'P1(G,I)',14X,'Q(G,I)',

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PRTO0001
PRTO0002
PRTO0003
PRTO0004
PRTO0005
PRTO0006
PRTO0007
PRTO0008
PRTO0009
PRTO0010
PRTO0011
PRTO0012
PRTO0013
PRTO0014
PRTO0015
PRTO0016
PRTO0017
PRTO0018
PRTO0019
PRTO0020
PRTO0021
PRTO0022
PRTO0023
PRTO0024
PRTO0025
PRTO0026
PRTO0027
PRTO0028
PRTO0029
PRTO0030
PRTO0031
PRTO0032
PRTO0033
PRTO0034
PRTO0035
PRTO0036

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X	13X,'Q1(G,I)',14X,'R(G,I)'	PRT00037
	DO 31 G=1,2	PRT00038
	WRITE (6,12)	PRT00039
31	WRITE (6,35) (G,I,P(G,I),P1(G,I),Q(G,I),Q1(G,I),R(G,I),I=1,KR)	PRT00040
35	FORMAT (2I5,5D20.7)	PRT00041
C	PO AND PH ARRAYS:	PRT00042
	WRITE (6,40)	PRT00043
40	FORMAT ('1 G',4X,'I',13X,'P0(G,I)',12X,'P07(G,I)',13X,'PH(G,I)',	PRT00044
X	12X,'PH7(G,I)')	PRT00045
	DO 41 G=1,2	PRT00046
	WRITE (6,12)	PRT00047
41	WRITE (6,45) (G,I,P0(G,I),P07(G,I),PH(G,I),PH7(G,I),I=1,KR)	PRT00048
45	FORMAT (2I5,4D20.7)	PRT00049
C	PRINT OUT THE /B3/ MATRICES:	PRT00050
	WRITE (6,50)	PRT00051
50	FORMAT ('1MATRIX L1:',/)	PRT00052
	WRITE (6,55) ((L1(I,J),J=1,3),I=1,N)	PRT00053
55	FORMAT (3E12.3,7X,3E12.3,7X,3E12.3)	PRT00054
	WRITE (6,60)	PRT00055
60	FORMAT ('1MATRIX L2:',/)	PRT00056
	WRITE (6,55) ((L2(I,J),J=1,3),I=1,N)	PRT00057
	WRITE (6,70)	PRT00058
70	FORMAT ('1MATRIX F1:',/)	PRT00059
	WRITE (6,55) ((F1(I,J),J=1,3),I=1,N)	PRT00060
	WRITE (6,80)	PRT00061
80	FORMAT ('1MATRIX F2:',/)	PRT00062
	WRITE (6,55) ((F2(I,J),J=1,3),I=1,N)	PRT00063
	WRITE (6,90)	PRT00064
90	FORMAT ('1MATRIX T:',/)	PRT00065
	WRITE (6,55) ((T(I,J),J=1,3),I=1,N)	PRT00066
	RETURN	PRT00067
	END	PRT00068

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SUBROUTINE POWER
C SOLVES THE 2*N MULTIGROUP EQUATIONS: M*PHI = (1/LAMDA)*F*PHI
C BY THE FISSION SOURCE POWER METHOD
C USING SIMULTANEOUS OVERRELAXATION.
C WHERE: M AND F ARE DOUBLE PRECISION 2N BY 2N BLOCK MATRICES;
C AND: PHI IS THE 2N FLUX (FAST AND THERMAL) VECTOR.
C L1*PHI1 = CHI1*(F1*PHI1 + F2*PHI2)
C -T*PHI1 + L2*PHI2 = CHI2*(F1*PHI1 + F2*PHI2)
C METHOD FOLLOWS WACHPRESS, PAGE 83. SOLUTION BY GROUP ITERATION.
IMPLICIT REAL*8 (A-H,L-Z)
COMMON /B1/ IBC,IPLOT,JPLOT,IPUNCH,ISEE
COMMON /B2/ KR,N
COMMON /B3/ L1(201,3), L2(201,3), F1(201,3), F2(201,3), T(201,3)
COMMON /B4/ PHI(2,201), PSI(2,201), LAMDA, ICOUT
COMMON /B5/ S(201), ERROR(2,201), Z(201)
COMMON /B6/ TE1(2,5),TE2(2,5),TE3(5),IN(5)
COMMON /B7/ HH(200)
COMMON /CHIF/ CHI(2)
COMMON /ER/ EPS1,EPS2,EPS3
COMMON /T/ I1,I4
COMMON /FSTR/ PHISTR(2,201,6)
COMMON /ESTR/ LAMSTR(300), EFSTR(2,300), EFMSTR(2,300), ERLAM(300)
COMMON /READ5/ R5
DIMENSION PSI1(201), PSI2(201), SQ(2), DPHI(2), ERRMAX(2)
INTEGER N
R5=1.
C DEFAULT OPTIONS FOR POWER PARAMETERS:
ALPHA=1.25
LAMDA=1.0
HX=0.0
DO 505 I=1,KR
505 HX=HX+HH(I)
DO 555 IG=1,2
IF (IBC.NE.4) GO TO 551
DO 550 I=1,N
550 PHI(IG,I)=1.0

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POW 0001
POW 0002
POW 0003
POW 0004
POW 0005
POW 0006
POW 0007
POW 0008
POW 0009
POW 0010
POW 0011
POW 0012
POW 0013
POW 0014
POW 0015
POW 0016
POW 0017
POW 0018
POW 0019
POW 0020
POW 0021
POW 0022
POW 0023
POW 0024
POW 0025
POW 0026
POW 0027
POW 0028
POW 0029
POW 0030
POW 0031
POW 0032
POW 0033
POW 0034
POW 0035
POW 0036

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	GO TO 555	POW 0037
551	X=3.1415926/HX	POW 0038
	IF (IBC.NE.1) X=X/2.0	POW 0039
	SUM1=0.0	POW 0040
	DO 552 K=1,KR	POW 0041
	SUM1=SUM1+H(K)	POW 0042
552	PHI(IG,K)=DSIN(SUM1*X)	POW 0043
555	CONTINUE	POW 0044
C	READ IN: OVERRELAXATION PARAMETERS ; ALPHA (OUTER ITERATION)	POW 0045
C	INITIAL GUESS AT EIGENVALUE; LAMDA	POW 0046
C	INITIAL NORMALIZED FLUX ; PHI(1-N)	POW 0047
	READ (5,506,END=510) ALPHA	POW 0048
	READ (5,502,END=510) LAMDA	POW 0049
	READ (5,503) (PHI(1,I),I=1,N)	POW 0050
	READ (5,503) (PHI(2,I),I=1,N)	POW 0051
506	FORMAT (F10.5)	POW 0052
502	FORMAT (E25.14)	POW 0053
503	FORMAT ((4E20.10))	POW 0054
	GO TO 511	POW 0055
510	R5=0.	POW 0056
511	CONTINUE	POW 0057
C	STORING FOR PRINTING THE MULTIGRUP FLUX SHAPE.	POW 0058
	DO 11 IG=1,2	POW 0059
	DO 10 I=1,N	POW 0060
10	PHISTR(IG,I,2)=PHI(IG,I)	POW 0061
C	FILL RUNNING COORD IN PHISTR	POW 0062
	KR1=KR+1	POW 0063
	DO 11 I=1,KR1	POW 0064
11	PHISTR(IG,I,1)=DFLOAT(I)	POW 0065
C	IK IS THE FLUX PLOTTING CCOUNTER.	POW 0066
	IK=1	POW 0067
C	STORES THE ITERATION NUMBER FOR FLUX HISTORY PLOTTING:	POW 0068
	IN(1)=0	POW 0069
C	STORES TEMPDRARY ERRORS FOR FLUX HISTORY PLOTTING:	POW 0070
	TE1(1,1)=0.	POW 0071
	TE1(2,1)=0.	POW 0072

	TE2(1,1)=0.	POW 0073
	TE2(2,1)=0.	POW 0074
	TE3(1)=0.0	POW 0075
C	EIGENVALUE OF THE PREVIOUS ITERATION:	POW 0076
	LAMB4=LAMDA	POW 0077
C	THE MAXIMUM NUMBER OF ALLOWED ITERATIONS: ICMAX	POW 0078
	ICMAX=300	POW 0079
C	PRINT OUT THE POWER METHOD PARAMETER INFORMATION:	POW 0080
	WRITE (6,700) ICMAX,ALPHA,LAMDA,(PHI(1,I),I=1,N)	POW 0081
	WRITE (6,701) (PHI(2,I),I=1,N)	POW 0082
700	FORMAT ('EXECUTING MULTIGROUP FISSION SOURCE POWER ITERATION METH	POW 0083
	XOD.',///,	POW 0084
X	5X,'MAXIMUM NUMBER OF ALLOWABLE ITERATIONS:',/,	POW 0085
X	10X,'ICMAX =',I4,///,	POW 0086
X	5X,'OUTER ITERATION RELAXATION PARAMETER:',/,	POW 0087
X	10X,'ALPHA =',F7.3,///,	POW 0088
X	5X,'INITIAL GUESS AT EIGENVALUE:',/,	POW 0089
X	10X,'LAMBDA =',E22.14,///,	POW 0090
X	5X,'INITIAL GUESS AT THE GROUP FLUX SHAPE CONNECTION POINTS:',	POW 0091
X	//,8X,'FAST GROUP:',/,	POW 0092
X	10X,'F(K)''S =',4E25.14,/, (18X,4E25.14))	POW 0093
701	FORMAT ('0',7X,'THERMAL GROUP:',/,	POW 0094
X	10X,'F(K)''S =',4E25.14,/, (18X,4E25.14))	POW 0095
C	BEGIN ITERATION LOOP.	POW 0096
	ICOUT=0	POW 0097
C	ICOUT IS THE OUTER ITERATION COUNTER.	POW 0098
20	ICOUT=ICOUT+1	POW 0099
	IF (ICOUT.GT.ICMAX) GO TO 100	POW 0100
C	FORM THE ITERATION SOURCE VECTOR, S; AND ITS L-2 NORM, SUM1:	POW 0101
	SUM1=0.	POW 0102
	DO 15 I=1,N	POW 0103
	S(I)=0.	POW 0104
	I0=1	POW 0105
	I1=3	POW 0106
	IF (I.EQ.1) I0=2	POW 0107
	IF (I.EQ.N) I1=2	POW 0108

	DO 14 J=10,11	POW 0109
	K=I-2+J	POW 0110
	14 S(I)=S(I)+F1(I,J)*PHI(1,K)+F2(I,J)*PHI(2,K)	POW 0111
	15 SUM1=SUM1+S(I)**2	POW 0112
	SUM1=DSQRT(SUM1)	POW 0113
	SUM1=SUM1*(CHI(1)+CHI(2))	POW 0114
C	SOLVE FOR THE NEW GROUP FLUX VECTORS: PSI:	POW 0115
C	FAST GROUP; SOURCE VECTOR:	POW 0116
	DO 25 I=1,N	POW 0117
	25 Z(I)=CHI(1)*S(I)	POW 0118
C	FAST FLUX:	POW 0119
	CALL SOLV3D(N,L1,PSI1,Z)	POW 0120
C	THERMAL GROUP; SOURCE VECTOR:	POW 0121
	DO 27 I=1,N	POW 0122
	Z(I)=0.	POW 0123
	I0=1	POW 0124
	I1=3	POW 0125
	IF (I.EQ.1) I0=2	POW 0126
	IF (I.EQ.N) I1=2	POW 0127
	DO 26 J=10,11	POW 0128
	K=I-2+J	POW 0129
	26 Z(I)=Z(I)+T(I,J)*PSI1(K)	POW 0130
	27 Z(I)=Z(I)+CHI(2)*S(I)	POW 0131
C	THERMAL FLUX:	POW 0132
	CALL SOLV3D(N,L2,PSI2,Z)	POW 0133
C	FORM NEW SOURCE VECTOR FROM THE NEW UNNORMALIZED FLUXES; PSI:	POW 0134
	SUM2=0.	POW 0135
	DO 29 I=1,N	POW 0136
	S(I)=0.	POW 0137
	I0=1	POW 0138
	I1=3	POW 0139
	IF (I.EQ.1) I0=2	POW 0140
	IF (I.EQ.N) I1=2	POW 0141
	DO 28 J=10,11	POW 0142
	K=I-2+J	POW 0143
	28 S(I)=S(I)+F1(I,J)*PSI1(K)+F2(I,J)*PSI2(K)	POW 0144

29	SUM2=SUM2+S(I)**2	POW 0145
	SUM2=DSQRT(SUM2)	POW 0146
	SUM2=SUM2*(CHI(1)+CHI(2))	POW 0147
C	CALCULATION OF THE EIGENVALUE:	POW 0148
	LAMDA=SUM2/SUM1	POW 0149
	LAMSTR(ICOUT)=LAMDA	POW 0150
	ERRLAM=DABS(LAMDA-LAMB4)	POW 0151
C	PUT PSI1 AND PSI2 INTO BIGGER PSI:	POW 0152
	DO 30 I=1,N	POW 0153
	PSI(1,I)=PSI1(I)	POW 0154
30	PSI(2,I)=PSI2(I)	POW 0155
C	POINT BY POINT SIMULTANEOUS RELAXATION FLUX ITERATION:	POW 0156
	X=ALPHA	POW 0157
C	DO NOT RELAX DURING THE FIRST THREE ITERATIONS:	POW 0158
	IF (ICOUT.LE.3) X=1.0	POW 0159
C	CALCULATE THE NEW GROUP FLUX ITERATES AND GROUP ERRORS:	POW 0160
	DO 40 IG=1,2	POW 0161
	SQ(IG)=0.	POW 0162
	DO 40 I=1,N	POW 0163
	ERROR(IG,I)=PSI(IG,I)/LAMDA-PHI(IG,I)	POW 0164
	SQ(IG)=SQ(IG)+ERROR(IG,I)**2	POW 0165
	PHI(IG,I)=PHI(IG,I) + X*ERROR(IG,I)	POW 0166
C	AND FOR PLOTTING PURPOSES:	POW 0167
	PSI(IG,I)=PHI(IG,I)	POW 0168
40	CONTINUE	POW 0169
	DO 34 IG=1,2	POW 0170
34	SQ(IG)=DSQRT(SQ(IG))	POW 0171
C	NORMALIZE PSI:	POW 0172
C	NORMALIZES BOTH ARRAY GROUPS TO 1.0:	POW 0173
	CALL NORM2(PSI,N)	POW 0174
	DO 36 IG=1,2	POW 0175
C	ERRMAX(IG) = THE MAX ERROR BETWEEN THE GROUP ITERATION FLUXES:	POW 0176
	ERRMAX(IG)=ERROR(IG,1)	POW 0177
	DO 36 I=2,N	POW 0178
	IF (DABS(ERROR(IG,I)).GT.ERRMAX(IG)) ERRMAX(IG)=DABS(ERROR(IG,I))	POW 0179
36	CONTINUE	POW 0180

	IF (IPL0T.NE.2) GO TO 45	POW 0181
C	THE FOLLOWING IS FOR NICELY PLOTTING THE GROUP FLUX HISTORY.	POW 0182
	DO 41 IG=1,2	POW 0183
	DO 41 I=1,N	POW 0184
41	ERROR(IG,I)=PSI(IG,I)	POW 0185
C	ERROR NOW CONTAINS THE NEW NORMALIZED FLUX ITERATE PHI.	POW 0186
	JK=IK	POW 0187
	IF (IK.EQ.0) JK=5	POW 0188
	DO 42 IG=1,2	POW 0189
	DO 42 I=1,N	POW 0190
	IF (DABS(ERROR(IG,I)-PHISTR(IG,I,JK+1)).GE.0.01) GO TO 43	POW 0191
42	CONTINUE	POW 0192
C	FLUX HAS NOT CHANGED ENOUGH FOR PLOTTING.	POW 0193
	GO TO 45	POW 0194
C	SAVE THE NORMALIZED FLUX FOR PLOTTING:	POW 0195
43	IK=IK+1	POW 0196
	IN(IK)=ICOUT	POW 0197
	TE3(IK)=ERRLAM	POW 0198
	DO 44 IG=1,2	POW 0199
	TE1(IG,IK)=ERRMAX(IG)	POW 0200
	TE2(IG,IK)=SQ(IG)	POW 0201
	DO 44 I=1,N	POW 0202
44	PHISTR(IG,I,IK+1)=ERROR(IG,I)	POW 0203
	IF (IK.NE.5) GO TO 45	POW 0204
C	PLOT THE LAST FIVE SAVED FLUXES:	POW 0205
	CALL PHIPLT(5)	POW 0206
	IK=0	POW 0207
45	CONTINUE	POW 0208
C	ERROR CRITERIA FOR ACCEPTANCE OF CONVERGENCE.	POW 0209
	IFLAG1=0	POW 0210
	IFLAG2=0	POW 0211
	IFLAG3=0	POW 0212
C	STORE THE ERRORS FOR COMPARISON:	POW 0213
C	ERROR BETWEEN ITERATION EIGENVALUES:	POW 0214
	ERLAM(ICOUT)=ERRLAM	POW 0215
	DO 46 IG=1,2	POW 0216



C	MAXIMUM ERROR BETWEEN ITERATION FLUXES:	POW 0217
	EFSTR(IG,ICOUT)=ERRMAX(IG)	POW 0218
C	MEAN SQUARE ERROR BETWEEN ITERATION FLUXES:	POW 0219
	EFMSTR(IG,ICOUT)=SQ(IG)	POW 0220
46	CONTINUE	POW 0221
	IF ((ERRMAX(1).LT.EPS1).AND.(ERRMAX(2).LT.EPS1)) IFLAG1=1	POW 0222
	IF ((SQ(1).LT.EPS2).AND.(SQ(2).LT.EPS2)) IFLAG2=1	POW 0223
	IF (ERRLAM.LT.EPS3) IFLAG3=1	POW 0224
	IFLAG4=IFLAG1*IFLAG2*IFLAG3	POW 0225
	IF (IFLAG4.EQ.1) GO TO 50	POW 0226
C	OTHERWISE CONTINUE THE ITERATION.	POW 0227
	LAMB4=LAMDA	POW 0228
	GO TO 20	POW 0229
50	CONTINUE	POW 0230
C	CONVERGENCE ACCOMPLISHED.	POW 0231
C	NORMALIZE THE CONVERGED FLUX VECTOR:	POW 0232
	CALL NORMAL(PHI,N)	POW 0233
C	PLOT ANY LEFT OVER FLUX HISTORY PLOTS:	POW 0234
	IF ((IPLJT.EQ.2).AND.(IK.NE.0)) CALL PHIPLT(IK)	POW 0235
C	BOUNDRY CONDITION INSERTICNS.	POW 0236
	IER=0	POW 0237
C	IER ALLOWS B.C. INSERTIONS FOR YES AND NO CONVERGENCE:	POW 0238
55	IF (IBC.EQ.4) GO TO 90	POW 0239
	IF (IBC.NE.3) GO TO 60	POW 0240
	PHI(1,KR+1)=0.	POW 0241
	PHI(2,KR+1)=0.	POW 0242
	GO TO 90	POW 0243
60	DO 70 I=1,N	POW 0244
	J=N+1-I	POW 0245
	PHI(1,J+1)=PHI(1,J)	POW 0246
70	PHI(2,J+1)=PHI(2,J)	POW 0247
	PHI(1,1)=0.	POW 0248
	PHI(2,1)=0.	POW 0249
	IF (IBC.NE.1) GO TO 90	POW 0250
	PHI(1,KR+1)=0.	POW 0251
	PHI(2,KR+1)=0.	POW 0252

90	IF (IER.EQ.1) GO TO 102	POW 0253
	RETURN	POW 0254
C	NO CONVERGENCE ACCOMPLISHED:	POW 0255
100	CONTINUE	POW 0256
C	NORMALIZE THE UNCONVERGED FLUX:	POW 0257
	CALL NORMAL(PHI,N)	POW 0258
	ICOUT=ICOUT-1	POW 0259
	WRITE (6,101) ICOUT	POW 0260
101	FORMAT (1H1,'POWER METHOD DID NOT CCNVERGE FOR THIS CASE AFTER',	POW 0261
	X I4,' ITERATIONS.',//,1X,'EXECUTION TERMINATED ')	POW 0262
	IER=1	POW 0263
	GO TO 55	POW 0264
102	CONTINUE	POW 0265
C	FOR PRINTING OUT THE EIGENVALUE HISTORY AND THE FINAL FLUX SHAPE:	POW 0266
	IF (IPL0T.EQ.0) IPL0T=1	POW 0267
	IF (JPL0T.EQ.0) JPL0T=1	POW 0268
	RETURN	POW 0269
	END	POW 0270

	SUBROUTINE CURT	CUR 0001
C	SOLVES FOR THE CURRENT FROM THE INPUT H(K)'S AND D(K)'S	CUR 0002
C	USING F(K)'S FROM POWER:	CUR 0003
C	CURRENT IS LINEAR (LEAST SQUARES - VARIATIONAL) AND PUT INTO ARRAY C.	CUR 0004
	IMPLICIT REAL*8 (A-H,O-Z)	CUR 0005
	COMMON /B2/ KR	CUR 0006
	COMMON /B4/ F(2,201), C(2,201)	CUR 0007
	COMMON /B5/ T(201,3), S1(201), S2(201), C1(201), C2(201)	CUR 0008
	COMMON /B7/ H(200), D(2,200)	CUR 0009
C	FROM THE MATRIX PROBLEM FOR LINEAR FIT OF STEP DATA:	CUR 0010
	M=KR	CUR 0011
	N=KR+1	CUR 0012
	T(1,1)=0.	CUR 0013
	T(N,3)=0.	CUR 0014
	T(1,2)=H(1)/3.	CUR 0015
	T(1,3)=H(1)/6.	CUR 0016
	T(N, 1)=H(M)/6.	CUR 0017
	T(N,2)=H(M)/3.	CUR 0018
	S1(1)=D(1,1)*(F(1,1)-F(1,2))/2.	CUR 0019
	S2(1)=D(2,1)*(F(2,1)-F(2,2))/2.	CUR 0020
	S1(N)=D(1,M)*(F(1,M)-F(1,M+1))/2.	CUR 0021
	S2(N)=D(2,M)*(F(2,M)-F(2,M+1))/2.	CUR 0022
	DO 20 I=2,M	CUR 0023
	J=I-1	CUR 0024
	T(I, 1)=H(J)/6.	CUR 0025
	T(I,2)=(H(J)+H(I))/3.	CUR 0026
	T(I, 3)=H(I)/6.	CUR 0027
	S1(I)=(D(1,J)*(F(1,J)-F(1,I))+D(1,I)*(F(1,I)-F(1,I+1)))/2.	CUR 0028
	S2(I)=(D(2,J)*(F(2,J)-F(2,I))+D(2,I)*(F(2,I)-F(2,I+1)))/2.	CUR 0029
20	CONTINUE	CUR 0030
	CALL SOLV3D(N,T,C1,S1)	CUR 0031
	CALL SOLV3D(N,T,C2,S2)	CUR 0032
	DO 30 I=1,N	CUR 0033
	C(1,I)=C1(I)	CUR 0034
30	C(2,I)=C2(I)	CUR 0035
	RETURN	CUR 0036

---

END

CUR 0037

PAGE 214

	SUBROUTINE OUTPUT	OUT 0001
C	PRINTS THE RESULTS OF THE METHOD.	OUT 0002
	IMPLICIT REAL*8 (A-H,L-Z)	OUT 0003
	COMMON /B1/ IBC,IPLT,JPLT,IPUNCH,ISEE,NOADJ	OUT 0004
	COMMON /B2/ KR,N	OUT 0005
	COMMON /B4/ PHI(2,201), CUR(2,201), LAMDA, ICOUT	OUT 0006
	COMMON /ER/ EPS1, EPS2, EPS3	OUT 0007
	COMMON /ESTR/ LAMSTR(300), EFSTR(2,300), EFMSTR(2,300), ERLAM(300)	OUT 0008
	INTEGER N, NOADJ	OUT 0009
	KRO=KR-1	OUT 0010
	KR1=KR+1	OUT 0011
	WRITE (6,1)	OUT 0012
	1 FORMAT ('RESULTS OF THE MULTIGROUP METHOD:')	OUT 0013
	WRITE (6,10) ICOUT	OUT 0014
	10 FORMAT (//, ' PROBLEM TERMINATED AFTER', I5,	OUT 0015
	X ' OUTER (POWER) ITERATIONS TO:')	OUT 0016
	WRITE (6,20) LAMDA	OUT 0017
	20 FORMAT (/, 10X, 'LAMDA = ', 1PE21.14)	OUT 0018
C	PRINT OUT EIGENVALUES.	OUT 0019
	CALL PLOT	OUT 0020
	WRITE (6,30)	OUT 0021
	30 FORMAT ('RESULTS AFTER PROBLEM TERMINATION:', /,	OUT 0022
	X 'NUMBER', 9X, 'THERMAL FLUX', 4X, 'THERMAL CURRENT', 12X,	OUT 0023
	X 'FAST FLUX', 7X, 'FAST CURRENT', /)	OUT 0024
	WRITE (6,50) (K, PHI(2,K), CUR(2,K), PHI(1,K), CUR(1,K), K=1, KR1)	OUT 0025
	50 FORMAT (17, 1PE21.7, 0PE19.7, 1PE21.7, 0PE19.7)	OUT 0026
C	PRINT OUT THE STORED ITERATION ERRORS:	OUT 0027
	WRITE (6,110) EPS1, (EFSTR(2,I), I=1, ICOUT)	OUT 0028
	WRITE (6,111) EPS1, (EFSTR(1,I), I=1, ICOUT)	OUT 0029
	WRITE (6,112) EPS2, (EFMSTR(2,I), I=1, ICOUT)	OUT 0030
	WRITE (6,113) EPS3, (EFMSTR(1,I), I=1, ICOUT)	OUT 0031
	WRITE (6,114) EPS3, (ERLAM(I), I=1, ICOUT)	OUT 0032
	110 FORMAT ('MAXIMUM ERRORS BETWEEN THE THERMAL FLUX ITERATIONS:',	OUT 0033
	X 25X, 'TOLERANCE USED = ', 1PE12.4, //, (1P5E20.5))	OUT 0034
	111 FORMAT ('MAXIMUM ERRORS BETWEEN THE FAST FLUX ITERATIONS:',	OUT 0035
	X 25X, 'TOLERANCE USED = ', 1PE12.4, //, (1P5E20.5))	OUT 0036

112	FORMAT ('1MEAN SQUARE ERROR BETWEEN THE THERMAL FLUX ITERATIONS:',	OUT 0037
X	18X,'TOLERANCE USED = ',1PE12.4, '//, (1P5E20.5))	OUT 0038
113	FORMAT ('1MEAN SQUARE ERROR BETWEEN THE FAST FLUX ITERATIONS:',	OUT 0039
X	18X,'TOLERANCE USED = ',1PE12.4, '//, (1P5E20.5))	OUT 0040
114	FORMAT ('1ERROR BETWEEN THE ITERATION EIGENVALUES:',	OUT 0041
X	28X,'TOLERANCE USED = ',1PE12.4, '//, (1P5E20.5))	OUT 0042
	IF (NDADJ.EQ.0) RETURN	OUT 0043
C	OTHERWISE ADJOINT CALCULATIONS ARE NOT EXECUTED:	OUT 0044
	WRITE (6,120)	OUT 0045
120	FORMAT ('1ADJOINT CALCULATIONS HAVE BEEN BYPASSED.', '//,	OUT 0046
X	' PROGRAM TERMINATED.')	OUT 0047
C	IPUNCH = 1 PUNCHES OUT THE FAST FLUX FOR SYNTH 1G INPUTS:	OUT 0048
	IF (IPUNCH.EQ.1) WRITE (7,124) KR	OUT 0049
	IF (IPUNCH.EQ.1) WRITE (7,125) (PHI(1,I),CUR(1,I),I=1,KR1)	OUT 0050
	IF (IPUNCH.EQ.1) WRITE (7,125) (PHI(2,I),CUR(2,I),I=1,KR1)	OUT 0051
124	FORMAT (15)	OUT 0052
125	FORMAT (2D20.10)	OUT 0053
	CALL EXIT	OUT 0054
	RETURN	OUT 0055
C		OUT 0056
	END	OUT 0057

	SUBROUTINE PLOT	PLT 0001
C	PLOTS OUT THE EIGENVALUE HISTORY AS A TABLE AND A GRAPH,	PLT 0002
C	AS WELL AS PLOTTING OUT THE FINAL MULTIGROUP FLUX SHAPES.	PLT 0003
	IMPLICIT REAL*8 (A-H,L-Z)	PLT 0004
	COMMON /B1/ IBC,IPLOT,JPLOT,IPUNCH	PLT 0005
	COMMON /B2/ KR	PLT 0006
	COMMON /B4/ PHI(2,201), PSI(2,201), LAMDA, ICOUT	PLT 0007
	COMMON /B5/ B(300,2)	PLT 0008
	COMMON /ESTR/ LAMSTR(300)	PLT 0009
	DIMENSION C(201,3)	PLT 0010
C	IN ORDER TO SAVE SOME SPACE:	PLT 0011
	EQUIVALENCE (B(1),C(1))	PLT 0012
	INTEGER ND	PLT 0013
	ND=201	PLT 0014
	WRITE (6,1) (LAMSTR(I),I=1,ICOUT)	PLT 0015
	1  FORMAT ('O TABLE OF EIGENVALUES DURING THE POWER ITERATION:',	PLT 0016
	X      //,(1P5E25.14))	PLT 0017
	IF (JPLOT.EQ.0) GO TO 20	PLT 0018
	DO 10  I=1,ICOUT	PLT 0019
	B(I,1)=I	PLT 0020
	10  B(I,2)=LAMSTR(I)	PLT 0021
	CALL PRTPLT(1,B,ICOUT,2,ICOUT,0,300,2,1)	PLT 0022
	WRITE (6,11)	PLT 0023
	11  FORMAT ('O PLOT OF THE EIGENVALUE HISTORY THROUGH THE ITERATIONS.')	PLT 0024
	20  IF (IPLJT.EQ.0) RETURN	PLT 0025
	KR1=KR+1	PLT 0026
	DO 30  I=1,KR1	PLT 0027
	C(I,1)=I	PLT 0028
	C(I,2)=PHI(1,I)	PLT 0029
	30  C(I,3)=PHI(2,I)	PLT 0030
	CALL PRTPLT(2,C,KR1,3,KR1,0,ND,3,2)	PLT 0031
	WRITE (6,31)	PLT 0032
	31  FORMAT ('O FINAL CONVERGED CONNECTING FLUX POINTS; F(K).',//,	PLT 0033
	X   5X,'FAST FLUX:      .',/,5X,'THERMAL FLUX:  -')	PLT 0034
	RETURN	PLT 0035
	END	PLT 0036

	SUBROUTINE POWER7	POW70001
	*** ADJOINT PROBLEM ***	POW70002
C	SOLVES THE 2*N MULTIGROUP ADJOINT EQUATIONS:	POW70003
C	M*PHI = (1/LAMDA)*F*PHI	POW70004
C	BY THE FISSION SOURCE POWER METHOD	POW70005
C	USING SIMULTANEOUS OVERRELAXATION.	POW70006
C	WHERE: M AND F ARE DOUBLE PRECISION 2N BY 2N BLOCK MATRICES;	POW70007
C	AND:    PHI IS THE 2N ADJOINT (FAST AND THERMAL) VECTOR.	POW70008
C	L1*PHI1 - T*PHI2 = CHI1*F1*PHI1 + CHI2*F1*PHI2	POW70009
C	L2*PHI2 = CHI1*F2*PHI1 + CHI2*F2*PHI2	POW70010
	IMPLICIT REAL*8 (A-H,L-Z)	POW70011
	COMMON /B1/ IBC,IPLOT,JPLOT,IPUNCH,ISEE	POW70012
	COMMON /B2/ KR,N	POW70013
	COMMON /B3/ L1(201,3), L2(201,3), F1(201,3), F2(201,3), T(201,3)	POW70014
	COMMON /B47/ PHI(2,201), PSI(2,201), LAMDA, ICOUT	POW70015
	COMMON /B5/ S(201), ERROR(2,201), Z(201)	POW70016
	COMMON /B6/ TE1(2,5),TE2(2,5),TE3(5),IN(5)	POW70017
	COMMON /B7/ HH(200)	POW70018
	COMMON /CHIF/ CHI(2)	POW70019
	COMMON /ER/ EPS1,EPS2,EPS3	POW70020
	COMMON /T/ I1,I4	POW70021
	COMMON /FSTR/ PHISTR(2,201,6)	POW70022
	COMMON /ESTR/ LAMSTR(300), EFSTR(2,300), EFMSTR(2,300), ERLAM(300)	POW70023
	COMMON /READ5/ R5	POW70024
	DIMENSION PSI1(201), PSI2(201), SQ(2), DPHI(2), ERRMAX(2)	POW70025
	INTEGER N	POW70026
C	DEFAULT OPTIONS FOR POWER PARAMETERS:	POW70027
	ALPHA=1.25	POW70028
	LAMDA=1.0	POW70029
	HX=0.0	POW70030
	DO 505 I=1,KR	POW70031
505	HX=HX+HH(I)	POW70032
	DO 555 IG=1,2	POW70033
	IF (IBC.NE.4) GO TO 551	POW70034
	DO 550 I=1,N	POW70035
550	PHI(IG,I)=1.0	POW70036



	GO TO 555	POW70037
551	X=3.1415926/HX	POW70038
	IF (IBC.NE.1) X=X/2.0	POW70039
	SUM1=0.0	POW70040
	DO 552 K=1,KR	POW70041
	SUM1=SUM1+HH(K)	POW70042
552	PHI(IG,K)=DSIN(SUM1*X)	POW70043
555	CONTINUE	POW70044
	IF (R5.EQ.0.) GO TO 510	POW70045
C	READ IN: OVERRELAXATION PARAMETERS ; ALPHA (OUTER ITERATION)	POW70046
C	INITIAL GUESS AT EIGENVALUE; LAMDA	POW70047
C	INITIAL NORMALIZED FLUX ; PHI(1-N)	POW70048
	READ (5,506,END=510) ALPHA	POW70049
	READ (5,502,END=510) LAMDA	POW70050
	READ (5,503) (PHI(1,I),I=1,N)	POW70051
	READ (5,503) (PHI(2,I),I=1,N)	POW70052
506	FORMAT (F10.5)	POW70053
502	FORMAT (E25.14)	POW70054
503	FORMAT ((4E20.10))	POW70055
510	CONTINUE	POW70056
C	STORING FOR PRINTING THE MULTIGROUP FLUX SHAPE.	POW70057
	DO 11 IG=1,2	POW70058
	DO 10 I=1,N	POW70059
10	PHISTR(IG,I,2)=PHI(IG,I)	POW70060
C	FILL RUNNING COORD IN PHISTR	POW70061
	KR1=KR+1	POW70062
	DO 11 I=1,KR1	POW70063
11	PHISTR(IG,I,1)=DFLOAT(I)	POW70064
C	IK IS THE FLUX PLOTTING COUNTER.	POW70065
	IK=1	POW70066
C	STORES THE ITERATION NUMBER FOR FLUX HISTORY PLOTTING:	POW70067
	IN(1)=0	POW70068
C	STORES TEMPORARY ERRORS FOR FLUX HISTORY PLOTTING:	POW70069
	TE1(1,1)=0.	POW70070
	TE1(2,1)=0.	POW70071
	TE2(1,1)=0.	POW70072

TE2(2,1)=0.	POW70073
TE3(1)=0.0	POW70074
C EIGENVALUE OF THE PREVIOUS ITERATION:	POW70075
LAMB4=LAMDA	POW70076
C THE MAXIMUM NUMBER OF ALLOWED ITERATIONS: ICMAX	POW70077
ICMAX=300	POW70078
C PRINT OUT THE POWER METHOD PARAMETER INFORMATION:	POW70079
WRITE (6,700) ICMAX,ALPHA,LAMDA,(PHI(1,I),I=1,N)	POW70080
WRITE (6,701) (PHI(2,I),I=1,N)	POW70081
700 FORMAT ('1EXECUTING MUTIGROUP ADJOINT FISSION SOURCE POWER ITERATI	POW70082
XON METHOD.',///,	POW70083
X 5X,'MAXIMUM NUMBER OF ALLOWABLE ITERATIONS:',/,	POW70084
X 10X,'ICMAX =',I4,///,	POW70085
X 5X,'OUTER ITERATION RELAXATION PARAMETER:',/,	POW70086
X 10X,'ALPHA =',F7.3,///,	POW70087
X 5X,'INITIAL GUESS AT ADJOINT EIGENVALUE:',/,	POW70088
X 10X,'LAMBDA =',E22.14,///,	POW70089
X 5X,'INITIAL GUESS AT THE GROUP FLUX SHAPE CONNECTION POINTS:',	POW70090
X //,8X,'FAST ADJOINT GROUP:',/,	POW70091
X 10X,'F(K)''S =',4E25.14,/, (18X,4E25.14))	POW70092
701 FORMAT ('0',7X,'THERMAL ADJOINT GROUP:',/,	POW70093
X 10X,'F(K)''S =',4E25.14,/, (18X,4E25.14))	POW70094
C BEGIN ITERATION LOOP.	POW70095
ICOUT=0	POW70096
C ICOUT IS THE OUTER ITERATION COUNTER.	POW70097
20 ICOUT=ICOUT+1	POW70098
IF (ICOUT.GT.ICMAX) GO TO 100	POW70099
C FORM THE GROUP TOTAL SOURCE S, AND ITS L-2 NORM SUM1:	POW70100
C AND THE THERMAL ADJOINT SOURCE VECTOR Z:	POW70101
SUM1=0.	POW70102
DO 15 I=1,N	POW70103
Z(I)=0.	POW70104
S(I)=0.	POW70105
I0=1	POW70106
I1=3	POW70107
IF (I.EQ.1) I0=2	POW70108

IF (I.EQ.N) I1=2	POW70109
DO 14 J=I0,I1	POW70110
K=I-2+J	POW70111
Z(I)=Z(I)+F2(I,J)*(CHI(1)*PHI(1,K)+CHI(2)*PHI(2,K))	POW70112
14 S(I)=S(I)+F1(I,J)*(CHI(1)*PHI(1,K)+CHI(2)*PHI(2,K))	POW70113
S(I)=S(I)+Z(I)	POW70114
15 SUM1=SUM1+S(I)**2	POW70115
SUM1=DSQRT(SUM1)	POW70116
C SOLVE FOR THE NEW GROUP ADJOINT FLUX VECTORS; PSI:	POW70117
C THERMAL ADJOINT FLUX:	POW70118
CALL SOLV3D(N,L2,PSI2,Z)	POW70119
C FAST ADJOINT GROUP; SOURCE VECTOR:	POW70120
DO 27 I=1,N	POW70121
S(I)=0.	POW70122
Z(I)=0.	POW70123
I0=1	POW70124
I1=3	POW70125
IF (I.EQ.1) I0=2	POW70126
IF (I.EQ.N) I1=2	POW70127
DO 26 J=I0,I1	POW70128
K=I-2+J	POW70129
S(I)=S(I)+CHI(1)*F1(I,J)*PHI(1,K)+CHI(2)*F1(I,J)*PSI2(K)	POW70130
26 Z(I)=Z(I)+T(I,J)*PSI2(K)	POW70131
27 Z(I)=Z(I)+S(I)	POW70132
C FAST ADJOINT FLUX:	POW70133
CALL SOLV3D(N,L1,PSI1,Z)	POW70134
C FORM NEW GROUP TOTAL SOURCE S FROM PSI'S, AND ITS L-2 NORM SUM2:	POW70135
SUM2=0.	POW70136
DO 29 I=1,N	POW70137
S(I)=0.	POW70138
I0=1	POW70139
I1=3	POW70140
IF (I.EQ.1) I0=2	POW70141
IF (I.EQ.N) I1=2	POW70142
DO 28 J=I0,I1	POW70143
K=I-2+J	POW70144

28	S(I)=S(I)+(F1(I,J)+F2(I,J))*(CHI(1)*PSI1(K)+CHI(2)*PSI2(K))	POW70145
29	SUM2=SUM2+S(I)**2	POW70146
	SUM2=DSQRT(SUM2)	POW70147
C	CALCULATION OF THE EIGENVALUE:	POW70148
	LAMDA=SUM2/SUM1	POW70149
	LAMSTR(ICDUT)=LAMDA	POW70150
	ERRLAM=DABS(LAMDA-LAMB4)	POW70151
C	PUT PSI1 AND PSI2 INTO BIGGER PSI:	POW70152
	DO 30 I=1,N	POW70153
	PSI(1,I)=PSI1(I)	POW70154
30	PSI(2,I)=PSI2(I)	POW70155
C	POINT BY POINT SIMULTANEOUS RELAXATION FLUX ITERATION:	POW70156
	X=ALPHA	POW70157
C	DO NOT RELAX DURING THE FIRST THREE ITERATIONS:	POW70158
	IF (ICDUT.LE.3) X=1.0	POW70159
C	CALCULATE THE NEW GROUP FLUX ITERATES AND GROUP ERRORS:	POW70160
	DO 40 IG=1,2	POW70161
	SQ(IG)=0.	POW70162
	DO 40 I=1,N	POW70163
	ERROR(IG,I)=PSI(IG,I)/LAMDA-PHI(IG,I)	POW70164
	SQ(IG)=SQ(IG)+ERROR(IG,I)**2	POW70165
	PHI(IG,I)=PHI(IG,I) + X*ERROR(IG,I)	POW70166
C	AND FOR PLOTTING PURPOSES:	POW70167
	PSI(IG,I)=PHI(IG,I)	POW70168
40	CONTINUE	POW70169
	DO 34 IG=1,2	POW70170
34	SQ(IG)=DSQRT(SQ(IG))	POW70171
C	NORMALIZE PSI:	POW70172
C	NORMALIZES BOTH ARRAY GROUPS TO 1.0:	POW70173
	CALL NORM2(PSI,N)	POW70174
	DO 36 IG=1,2	POW70175
C	ERRMAX(IG) = THE MAX ERROR BETWEEN THE GROUP ITERATION FLUXES:	POW70176
	ERRMAX(IG)=ERROR(IG,1)	POW70177
	DO 36 I=2,N	POW70178
	IF (DABS(ERROR(IG,I)).GT.ERRMAX(IG)) ERRMAX(IG)=DABS(ERROR(IG,I))	POW70179
36	CONTINUE	POW70180

	IF (I PLOT.NE.2) GO TO 45	POW70181
C	THE FOLLOWING IS FOR NICELY PLOTTING THE GROUP FLUX HISTORY.	POW70182
	DO 41 IG=1,2	POW70183
	DO 41 I=1,N	POW70184
41	ERROR(IG,I)=PSI(IG,I)	POW70185
C	ERROR NOW CONTAINS THE NEW NORMALIZED FLUX ITERATE PHI.	POW70186
	JK=IK	POW70187
	IF (IK.EQ.0) JK=5	POW70188
	DO 42 IG=1,2	POW70189
	DO 42 I=1,N	POW70190
	IF (DABS(ERROR(IG,I)-PHISTR(IG,I,JK+1)).GE.0.01) GO TO 43	POW70191
42	CONTINUE	POW70192
C	FLUX HAS NOT CHANGED ENOUGH FOR PLOTTING.	POW70193
	GO TO 45	POW70194
C	SAVE THE NORMALIZED FLUX FOR PLOTTING:	POW70195
43	IK=IK+1	POW70196
	IN(IK)=ICOUT	POW70197
	TE3(IK)=ERRLAM	POW70198
	DO 44 IG=1,2	POW70199
	TE1(IG,IK)=ERRMAX(IG)	POW70200
	TE2(IG,IK)=SQ(IG)	POW70201
	DO 44 I=1,N	POW70202
44	PHISTR(IG,I,IK+1)=ERROR(IG,I)	POW70203
	IF (IK.NE.5) GO TO 45	POW70204
C	PLOT THE LAST FIVE SAVED FLUXES:	POW70205
	CALL PHIPLT(5)	POW70206
	IK=0	POW70207
45	CONTINUE	POW70208
C	ERROR CRITERIA FOR ACCEPTANCE OF CONVERGENCE.	POW70209
	IFLAG1=0	POW70210
	IFLAG2=0	POW70211
	IFLAG3=0	POW70212
C	STORE THE ERRORS FOR COMPARISON:	POW70213
C	ERROR BETWEEN ITERATION EIGENVALUES:	POW70214
	ERLAM(ICOUT)=ERRLAM	POW70215
	DO 46 IG=1,2	POW70216

C	MAXIMUM ERROR BETWEEN ITERATION FLUXES:	POW70217
	EFSTR(IG,ICOUT)=ERRMAX(IG)	POW70218
C	MEAN SQUARE ERROR BETWEEN ITERATION FLUXES:	POW70219
	EFMSTR(IG,ICOUT)=SQ(IG)	POW70220
46	CONTINUE	POW70221
	IF ((ERRMAX(1).LT.EPS1).AND.(ERRMAX(2).LT.EPS1)) IFLAG1=1	POW70222
	IF ((SQ(1).LT.EPS2).AND.(SQ(2).LT.EPS2)) IFLAG2=1	POW70223
	IF (ERRLAM.LT.EPS3) IFLAG3=1	POW70224
	IFLAG4=IFLAG1*IFLAG2*IFLAG3	POW70225
	IF (IFLAG4.EQ.1) GO TO 50	POW70226
C	OTHERWISE CONTINUE THE ITERATION.	POW70227
	LAMB4=LAMDA	POW70228
	GO TO 20	POW70229
50	CONTINUE	POW70230
C	CONVERGENCE ACCOMPLISHED.	POW70231
C	NORMALIZE THE CONVERGED FLUX VECTOR:	POW70232
	CALL NORMAL(PHI,N)	POW70233
C	PLOT ANY LEFT OVER FLUX HISTORY PLOTS:	POW70234
	IF ((IPLT.EQ.2).AND.(IK.NE.0)) CALL PHIPLT(IK)	POW70235
C	BOUNDARY CONDITION INSERTIONS.	POW70236
	IER=0	POW70237
C	IER ALLOWS B.C. INSERTIONS FOR YES AND NO CONVERGENCE:	POW70238
55	IF (IBC.EQ.4) GO TO 90	POW70239
	IF (IBC.NE.3) GO TO 60	POW70240
	PHI(1,KR+1)=0.	POW70241
	PHI(2,KR+1)=0.	POW70242
	GO TO 90	POW70243
60	DO 70 I=1,N	POW70244
	J=N+1-I	POW70245
	PHI(1,J+1)=PHI(1,J)	POW70246
70	PHI(2,J+1)=PHI(2,J)	POW70247
	PHI(1,1)=0.	POW70248
	PHI(2,1)=0.	POW70249
	IF (IBC.NE.1) GO TO 90	POW70250
	PHI(1,KR+1)=0.	POW70251
	PHI(2,KR+1)=0.	POW70252

90	IF (IER.EQ.1) GO TO 102	POW70253
	RETURN	POW70254
C	NO CONVERGENCE ACCOMPLISHED:	POW70255
100	CONTINUE	POW70256
C	NORMALIZE THE UNCONVERGED FLUX:	POW70257
	CALL NORMAL(PHI,N)	POW70258
	ICOUT=ICOUT-1	POW70259
	WRITE (6,101) ICOUT	POW70260
101	FORMAT (1H1,'POWER METHOD DID NOT CONVERGE FOR THIS CASE AFTER',	POW70261
	X 14,' ITERATIONS.',//,1X,'EXECUTION TERMINATED ')	POW70262
	IER=1	POW70263
	GO TO 55	POW70264
102	CONTINUE	POW70265
C	FOR PRINTING OUT THE EIGENVALUE HISTORY AND THE FINAL FLUX SHAPE:	POW70266
	IF (IPLST.EQ.0) IPLST=1	POW70267
	IF (JPLOT.EQ.0) JPLOT=1	POW70268
	RETURN	POW70269
	END	POW70270

	SUBROUTINE CURT7	CUR70001
C	SOLVES FOR THE ADJOINT CURRENT FROM THE INPUT H(K)'S AND D(K)'S	CUR70002
C	USING F(K)'S FROM POWER7:	CUR70003
C	CURRENT IS LINEAR (LEAST SQUARES - VARIATIONAL) AND PUT INTO ARRAY C.	CUR70004
	IMPLICIT REAL*8 (A-H,O-Z)	CUR70005
	COMMON /B2/ KR	CUR70006
	COMMON /B47/ F(2,201), C(2,201)	CUR70007
	COMMON /B5/ T(201,3), S1(201), S2(201), C1(201), C2(201)	CUR70008
	COMMON /B7/ H(200), D(2,200)	CUR70009
C	FROM THE MATRIX PROBLEM FOR LINEAR FIT OF STEP DATA:	CUR70010
	M=KR	CUR70011
	N=KR+1	CUR70012
	T(1,1)=0.	CUR70013
	T(N,3)=0.	CUR70014
	T(1,2)=H(1)/3.	CUR70015
	T(1,3)=H(1)/6.	CUR70016
	T(N, 1)=H(M)/6.	CUR70017
	T(N,2)=H(M)/3.	CUR70018
	S1(1)=D(1,1)*(F(1,1)-F(1,2))/2.	CUR70019
	S2(1)=D(2,1)*(F(2,1)-F(2,2))/2.	CUR70020
	S1(N)=D(1,M)*(F(1,M)-F(1,M+1))/2.	CUR70021
	S2(N)=D(2,M)*(F(2,M)-F(2,M+1))/2.	CUR70022
	DO 20 I=2,M	CUR70023
	J=I-1	CUR70024
	T(I, 1)=H(J)/6.	CUR70025
	T(I,2)=(H(J)+H(I))/3.	CUR70026
	T(I, 3)=H(I)/6.	CUR70027
	S1(I)=(D(1,J)*(F(1,J)-F(1,I))+D(1,I)*(F(1,I)-F(1,I+1)))/2.	CUR70028
	S2(I)=(D(2,J)*(F(2,J)-F(2,I))+D(2,I)*(F(2,I)-F(2,I+1)))/2.	CUR70029
20	CONTINUE	CUR70030
	CALL SOLV3D(N,T,C1,S1)	CUR70031
	CALL SOLV3D(N,T,C2,S2)	CUR70032
	DO 30 I=1,N	CUR70033
	C(1,I)=-C1(I)	CUR70034
30	C(2,I)=-C2(I)	CUR70035
	RETURN	CUR70036



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END

CUR70037

	SUBROUTINE OUTPT7	OUT70001
C	PRINTS THE RESULTS OF THE ADJOINT METHOD:	OUT70002
	IMPLICIT REAL*8 (A-H,L-Z)	OUT70003
	COMMON /B1/ IBC,IPLT,JPLT,IPUNCH	OUT70004
	COMMON /B2/ KR,N	OUT70005
	COMMON /B47/ PHI(2,201), CUR(2,201), LAMDA, ICOUT	OUT70006
	COMMON /ER/ EPS1, EPS2, EPS3	OUT70007
	COMMON /ESTR/ LAMSTR(300), EFSTR(2,300), EFMSTR(2,300), ERLAM(300)	OUT70008
	INTEGER N	OUT70009
	KRO=KR-1	OUT70010
	KR1=KR+1	OUT70011
	WRITE (6,1)	OUT70012
	1 FORMAT ('RESULTS OF THE MULTIGROUP ADJOINT METHOD:')	OUT70013
	WRITE (6,10) ICOUT	OUT70014
	10 FORMAT (//, ' PROBLEM TERMINATED AFTER', I5,	OUT70015
	X ' OUTER (POWER) ITERATIONS TO:')	OUT70016
	WRITE (6,20) LAMDA	OUT70017
	20 FORMAT (/ ,10X, 'ADJOINT LAMBDA = ', 1PE21.14)	OUT70018
C	PRINT OUT EIGENVALUES.	OUT70019
	CALL PLOT7	OUT70020
	WRITE (6,30)	OUT70021
	30 FORMAT ('RESULTS AFTER PROBLEM TERMINATION:', /,	OUT70022
	X 5X, 'ADJOINTS:', /,	OUT70023
	X 'ONUMBER', 9X, 'THERMAL FLUX', 4X, 'THERMAL CURRENT', 12X,	OUT70024
	X 'FAST FLUX', 7X, 'FAST CURRENT', /)	OUT70025
	WRITE (6,50) (K, PHI(2,K), CUR(2,K), PHI(1,K), CUR(1,K), K=1, KR1)	OUT70026
	50 FORMAT (I7, 1PE21.7, 0PE19.7, 1PE21.7, 0PE19.7)	OUT70027
C	PRINT OUT THE STORED ITERATION ERRORS:	OUT70028
	WRITE (6,110) EPS1, (EFSTR(2,I), I=1, ICOUT)	OUT70029
	WRITE (6,111) EPS1, (EFSTR(1,I), I=1, ICOUT)	OUT70030
	WRITE (6,112) EPS2, (EFMSTR(2,I), I=1, ICOUT)	OUT70031
	WRITE (6,113) EPS3, (EFMSTR(1,I), I=1, ICOUT)	OUT70032
	WRITE (6,114) EPS3, (ERLAM(I), I=1, ICOUT)	OUT70033
	110 FORMAT ('MAXIMUM ERRORS BETWEEN THE THERMAL FLUX ITERATIONS:',	OUT70034
	X 25X, 'TOLERANCE USED = ', 1PE12.4, //, (1P5E20.5))	OUT70035
	111 FORMAT ('MAXIMUM ERRORS BETWEEN THE FAST FLUX ITERATIONS:',	OUT70036

X	25X,'TOLERANCE USED = ',1PE12.4, '//, (1P5E20.5))	OUT70037
112	FORMAT ('1MEAN SQUARE ERROR BETWEEN THE THERMAL FLUX ITERATIONS:',	OUT70038
X	18X,'TOLERANCE USED = ',1PE12.4, '//, (1P5E20.5))	OUT70039
113	FORMAT ('1MEAN SQUARE ERROR BETWEEN THE FAST FLUX ITERATIONS:',	OUT70040
X	18X,'TOLERANCE USED = ',1PE12.4, '//, (1P5E20.5))	OUT70041
114	FORMAT ('1ERROR BETWEEN THE ITERATION EIGENVALUES:',	OUT70042
X	28X,'TOLERANCE USED = ',1PE12.4, '//, (1P5E20.5))	OUT70043
C	CHECK FOR CALL TO PUNCH:	OUT70044
	IF (IPUNCH.EQ.1) CALL PUNCH	OUT70045
	RETURN	OUT70046
	END	OUT70047

```

SUBROUTINE PLOT7
C   PLOTS OUT THE EIGENVALUE HISTORY AS A TABLE AND A GRAPH,
C   AS WELL AS PLOTTING OUT THE FINAL MULTIGROUP FLUX SHAPES.
C   FOR THE ADJOINTS:
      IMPLICIT REAL*8 (A-H,L-Z)
      COMMON /B1/ IBC,IPLOT,JPLOT,IPUNCH
      COMMON /B2/ KR
      COMMON /B47/ PHI(2,201), CUR(2,201), LAMDA, ICOUT
      COMMON /B5/ B(300,2)
      COMMON /ESTR/ LAMSTR(300)
      DIMENSION C(201,3)
C   IN ORDER TO SAVE SOME SPACE:
      EQUIVALENCE (B(1),C(1))
      INTEGER ND
      ND=201
      WRITE (6,1) (LAMSTR(I),I=1,ICOUT)
1  FORMAT ('O TABLE OF EIGENVALUES DURING THE POWER ITERATION:',
X      //,(1P5E25.14))
C
      IF (JPLOT.EQ.0) GO TO 20
      DO 10 I=1,ICOUT
      B(I,1)=1
10  B(I,2)=LAMSTR(I)
      CALL PRTPLT(1,B,ICOUT,2,ICOUT,0,300,2,1)
      WRITE (6,11)
11  FORMAT ('O PLOT OF THE EIGENVALUE HISTORY THROUGH THE ITERATIONS.')
20  IF (IPLJT.EQ.0) RETURN
      KR1=KR+1
      DO 30 I=1,KR1
      C(I,1)=I
      C(I,2)=PHI(1,I)
30  C(I,3)=PHI(2,I)
      CALL PRTPLT(2,C,KR1,3,KR1,0,ND,3,2)
      WRITE (6,31)
31  FORMAT ('O FINAL CONVERGED CONNECTING FLUX POINTS; F(K).',//,
X      5X,'FAST FLUX:      .',/,5X,'THERMAL FLUX:  -')

```

```

PLT70001
PLT70002
PLT70003
PLT70004
PLT70005
PLT70006
PLT70007
PLT70008
PLT70009
PLT70010
PLT70011
PLT70012
PLT70013
PLT70014
PLT70015
PLT70016
PLT70017
PLT70018
PLT70019
PLT70020
PLT70021
PLT70022
PLT70023
PLT70024
PLT70025
PLT70026
PLT70027
PLT70028
PLT70029
PLT70030
PLT70031
PLT70032
PLT70033
PLT70034
PLT70035
PLT70036

```

---

RETURN  
END

PLT70037  
PLT70038

C	SUBROUTINE SOLV3D(N,A,X,Y)	SOLV0001
C	SOLVES THE N DOUBLE PRECISION MATRIX EQUATIONS: $A * X = Y$ ,	SOLV0002
C	FOR X - GIVEN THE N BY N TRIDIAGONAL MATRIX A	SOLV0003
C	AND THE SOURCE VECTOR Y.	SOLV0004
C	METHOD IS FORWARD ELIMINATION FOLLOWED BY BACKWARD SUBSTITUTION.	SOLV0005
C	CF - WACHPRESS, PAGE 23.	SOLV0006
	REAL*8 A, X, Y, H, P, D	SOLV0007
	DIMENSION A(201,3), X(201), Y(201), H(201), P(201)	SOLV0008
	IF (A(1,2).EQ.0.0) GO TO 10	SOLV0009
	H(1)=-A(1,3)/A(1,2)	SOLV0010
	P(1)=Y(1)/A(1,2)	SOLV0011
	DO 1 M=2,N	SOLV0012
	D=A(M,2)+A(M, 1)*H(M-1)	SOLV0013
	IF (D.EQ.0.0) GO TO 20	SOLV0014
	P(M)=(Y(M)-A(M, 1)*P(M-1))/D	SOLV0015
	IF (M.EQ.N) GO TO 1	SOLV0016
	H(M)=-A(M, 3)/D	SOLV0017
	1 CONTINUE	SOLV0018
	X(N)=P(N)	SOLV0019
	DO 2 I=2,N	SOLV0020
	M=N+1-I	SOLV0021
	2 X(M)=P(M)+H(M)*X(M+1)	SOLV0022
	RETURN	SOLV0023
C	IN CASE OF ANY IMPENDING ZERO DIVISORS:	SOLV0024
	10 WRITE (6,11)	SOLV0025
	11 FORMAT ('FIRST ELEMENT OF A, A(1,1), IS ZERO.',/,	SOLV0026
	X 5X,'BETTER FIX IT BOSS.')	SOLV0027
	GO TO 30	SOLV0028
	20 WRITE (6,21) M	SOLV0029
	21 FORMAT ('ZERO DIVISOR ENCOUNTERED IN EQUATION M =',I3,'.',/,	SOLV0030
	X 5X,'BETTER FIX IT BOSS.')	SOLV0031
	30 WRITE (6,31)	SOLV0032
	31 FORMAT ('EXECUTION TERMINATED.')	SOLV0033
	CALL EXIT	SOLV0034
	RETURN	SOLV0035
	END	SOLV0036

```
C  SUBROUTINE NORMAL(PHI,N)
    NORMALIZES THE GROUP FLUXES TO ONE. NOT BOTH GROUPS.
    REAL*8 PHI(2,201), A
    A=DABS(PHI(1,1))
    DO 1 IG=1,2
    DO 1 I=1,N
    IF (DABS(PHI(IG,I)).GT.A) A=DABS(PHI(IG,I))
1  CONTINUE
    DO 2 IG=1,2
    DO 2 I=1,N
2  PHI(IG,I)=PHI(IG,I)/A
    RETURN
    END
```

```
NORL0001
NORL0002
NORL0003
NORL0004
NORL0005
NORL0006
NORL0007
NORL0008
NORL0009
NORL0010
NORL0011
NORL0012
NORL0013
```

C       SUBROUTINE NORM2(PSI,N)  
          NORMALIZES BOTH ENERGY GROUPS OF PSI TO 1.0.  
          REAL\*8 PSI(2,201), A(2)  
          DO 1 IG=1,2  
          A(IG)=DABS(PSI(IG,1))  
          DO 1 I=1,N  
          IF (DABS(PSI(IG,I)).GT.A(IG)) A(IG)=DABS(PSI(IG,I))  
1       CONTINUE  
          DO 2 IG=1,2  
          DO 2 I=1,N  
2       PSI(IG,I)=PSI(IG,I)/A(IG)  
          RETURN  
          END

NOR20001  
NOR20002  
NOR20003  
NOR20004  
NOR20005  
NOR20006  
NOR20007  
NOR20008  
NOR20009  
NOR20010  
NOR20011  
NOR20012  
NOR20013



	SUBROUTINE PHIPLT(L)	PHIP0001
C	PLOTS THE GROUP FLUX HISTORY, WITH UP TO 5 GROUP FLUXES PER PLOT.	PHIP0002
C	FAST AND THERMAL GROUP FLUXES ARE PLOTTED SEPERATELY.	PHIP0003
C	L IS THE NUMBER OF FLUXES TO BE PLOTTED.	PHIP0004
C	L IS BETWEEN 1 AND 5.	PHIP0005
	IMPLICIT REAL*8 (A-H,O-Z)	PHIP0006
	COMMON /B1/ IBC	PHIP0007
	COMMON /B2/ KR,N	PHIP0008
	COMMON /B5/ S(201), A(201,6), B(201,6)	PHIP0009
	COMMON /B6/ TE1(2,5),TE2(2,5),TE3(5),IN(5)	PHIP0010
	COMMON /ER/ EPS1,EPS2,EPS3	PHIP0011
	COMMON /FSTR/ PHISTR(2,201,6)	PHIP0012
	DIMENSION SYMBOL(5)	PHIP0013
	INTEGER SYMBOL /'.' , '-' , '+' , '#' , '*' /	PHIP0014
	ND=201	PHIP0015
	KR1=KR+1	PHIP0016
C	SET UP B.C. CONDITIONS	PHIP0017
	IF (IBC.EQ.4) GO TO 5	PHIP0018
	IF (IBC.EQ.3) GO TO 3	PHIP0019
	DO 2 IG=1,2	PHIP0020
	DO 2 K=1,L	PHIP0021
	DO 1 I=1,N	PHIP0022
	J=N+1-I	PHIP0023
	1 PHISTR(IG,J+1,K+1)=PHISTR(IG,J,K+1)	PHIP0024
	2 PHISTR(IG,1,K+1)=0.	PHIP0025
	3 IF (IBC.EQ.2) GO TO 5	PHIP0026
	DO 4 IG=1,2	PHIP0027
	DO 4 K=1,L	PHIP0028
	4 PHISTR(IG,KR1,K+1)=0.	PHIP0029
	5 CONTINUE	PHIP0030
C	FLUXES IN PHISTR HAVE BEEN NORMALIZED IN POWER.	PHIP0031
C	PUT THE FAST FLUX IN A, AND THE THERMAL FLUX IN B:	PHIP0032
	L1=L+1	PHIP0033
	DO 10 K=1,L1	PHIP0034
	DO 10 I=1,KR1	PHIP0035
	A(I,K)=PHISTR(1,I,K)	PHIP0036

10	B(I,K)=PHISTR(2,I,K)	PHIP0037
C	PLOT THE L FAST FLUX SHAPES ON ONE GRAPH:	PHIP0038
	CALL PRTPLT(0,A,KR1,L1,KR1,0,ND,6,2)	PHIP0039
	WRITE (6,20)	PHIP0040
20	FORMAT (/, 'OFAST FLUX ITERATION HISTORY PLOT.', /)	PHIP0041
	WRITE (6,30)	PHIP0042
30	FORMAT (	PHIP0043
	X 'OKEY:', 5X, 'SYMBOL', 5X, 'ITERATION NUMBER:', 7X, 'ERROR CRITERIA',	PHIP0044
	X 11X, 'ERROR', 13X, 'TOLERANCE')	PHIP0045
	DO 35 I=1,L	PHIP0046
35	WRITE (6,40) SYMBOL(I), IN(I), TE1(1, I), EPS1, TE2(1, I), EPS2,	PHIP0047
	X                    TE3(I), EPS3	PHIP0048
40	FORMAT (/, 12X, A1, 15X, I3, 16X, 'FLUX', 14X, 1PD15.5, 5X, 1PD15.5, /,	PHIP0049
	X 47X, 'MEAN SQ. FLUX', 5X, 1PD15.5, 5X, 1PD15.5, /,	PHIP0050
	X 47X, 'EIGENVALUE', 8X, 1PD15.5, 5X, 1PD15.5)	PHIP0051
C		PHIP0052
C	PLOT THE L THERMAL FLUX SHAPES ON THE OTHER GRAPH:	PHIP0053
	CALL PRTPLT(0,B,KR1,L1,KR1,0,ND,6,2)	PHIP0054
	WRITE (6,50)	PHIP0055
50	FORMAT (/, 'OTHERMAL FLUX ITERATION PLOT.', /)	PHIP0056
	WRITE (6,30)	PHIP0057
	DO 55 I=1,L	PHIP0058
55	WRITE (6,40) SYMBOL(I), IN(I), TE1(2, I), EPS1, TE2(2, I), EPS2,	PHIP0059
	X                    TE3(I), EPS3	PHIP0060
	RETURN	PHIP0061
	END	PHIP0062

```

SUBROUTINE PRTPLT(NO,B,N,M,NL,NS,KX,JX,ISP)
C*****MODIFIED VERSION FROM THAT OF SSP OR ANY OTHER SOURCE *****
C      CONVERTS DOUBLE PRECISION B ARRAY TO REAL*4.
C      PLOT SEVERAL CROSS-VARIABLES VERSUS A BASE VARIABLE
C      NO - CHART NUMBER (3 DIGITS MAXIMUM)
C      B - MATRIX OF DATA TO BE PLOTTED. FIRST COLUMN REPRESENTS
C      BASE VARIABLE AND SUCCESSIVE COLUMNS ARE THE CROSS-
C      VARIABLES (MAXIMUM IS 9).
C      N - NUMBER OF ROWS IN MATRIX B
C      M - NUMBER OF COLUMNS IN MATRIX B (EQUAL TO THE TOTAL
C      NUMBER OF VARIABLES). MAXIMUM IS 10.
C      NL - NUMBER OF LINES IN THE PLOT. IF 0 IS SPECIFIED, 50
C      LINES ARE USED. THE NUMBER OF LINES MUST BE EQUAL TO
C      OR GREATER THAN N
C      (USUALLY USE NL=N, AND ISP FOR SPACING.)
C      NS - CODE FOR SORTING THE BASE VARIABLE DATA IN ASCENDING
C      ORDER
C           0 SORTING IS NOT NECESSARY (ALREADY IN ASCENDING
C           ORDER).
C           1 SORTING IS NECESSARY.
C      KX- DIMENSION OF B MATRIX FROM DIMENSION STATEMENT.
C      IT MUST BE OF THE FORM B(KX,JX)
C      JX- DIMENSION OF B MATRIX FROM DIMENSION STATEMENT.
C      IT MUST BE OF THE FORM B(KX,JX)
C      ISP- CODE FOR SPACING LINES WHILE PLOTTING:
C           1 SINGLE SPACE
C           2 DOUBLE SPACE
C           3 TRIPLE SPACE
C           ...ETC.
REAL*8 B
DIMENSION OUT(101),YPR(11),IANG(9),A( 1500),B(KX,JX)
INTEGER IDUM/'1'/,IANG/'.' , '-' , '+' , '#' , '*' , 'A' , 'B' , 'C' , 'D' /
INTEGER OUT
I=1
DO 39 J=1,M
DO 39 K=1,N

```

```

P RTP0001
P RTP0002
P RTP0003
P RTP0004
P RTP0005
P RTP0006
P RTP0007
P RTP0008
P RTP0009
P RTP0010
P RTP0011
P RTP0012
P RTP0013
P RTP0014
P RTP0015
P RTP0016
P RTP0017
P RTP0018
P RTP0019
P RTP0020
P RTP0021
P RTP0022
P RTP0023
P RTP0024
P RTP0025
P RTP0026
P RTP0027
P RTP0028
P RTP0029
P RTP0030
P RTP0031
P RTP0032
P RTP0033
P RTP0034
P RTP0035
P RTP0036

```

	A(I)=B(K,J)	P RTP0037
	I=I+1	P RTP0038
39	CONTINUE	P RTP0039
	1 FORMAT(1H1,60X,7H CHART ,I3,//)	P RTP0040
	2 FORMAT(1H ,F11.4,5X,101A1)	P RTP0041
	3 FORMAT(1H )	P RTP0042
	7 FORMAT(1H ,16X,101H+           +           +           +           +           +           +           +)	P RTP0043
	1   +           +           +           +           +           +           +           +)	P RTP0044
	8 FORMAT(1H ,9X,11F10.4)	P RTP0045
	NLL=NL	P RTP0046
	IF(NS) 16, 16, 10	P RTP0047
C	SORT BASE VARIABLE DATA IN ASCENDING ORDER	P RTP0048
10	DO 15 I=1,N	P RTP0049
	DO 14 J=I,N	P RTP0050
	IF(A(I)-A(J)) 14, 14, 11	P RTP0051
11	L=I-N	P RTP0052
	LL=J-N	P RTP0053
	DO 12 K=1,M	P RTP0054
	L=L+N	P RTP0055
	LL=LL+N	P RTP0056
	F=A(L)	P RTP0057
	A(L)=A(LL)	P RTP0058
12	A(LL)=F	P RTP0059
14	CONTINUE	P RTP0060
15	CONTINUE	P RTP0061
C	TEST NLL	P RTP0062
16	IF(NLL) 20, 18, 20	P RTP0063
18	NLL=50	P RTP0064
C	PRINT TITLE	P RTP0065
20	WRITE(6,1)ND	P RTP0066
C	DEVELOP BLANK AND DIGITS FOR PRINTING	P RTP0067
	BLANK=0	P RTP0068
C	FIND SCALE FOR BASE VARIABLE	P RTP0069
	XSCAL=(A(N)-A(1))/(FLOAT(NLL-1))	P RTP0070
C	FIND SCALE FOR CROSS-VARIABLES	P RTP0071
	YMIN=1.0E75	P RTP0072

YMAX=-1.0E75	P RTP0073
M1=N+1	P RTP0074
M2=M*N	P RTP0075
DO 40 J=M1,M2	P RTP0076
IF (A(J).GT.YMAX) YMAX=A(J)	P RTP0077
IF (A(J).LT.YMIN) YMIN=A(J)	P RTP0078
40 CONTINUE	P RTP0079
YSCAL=(YMAX-YMIN)/100.0	P RTP0080
C CHECK TO SEE IF THE SPREAD IN Y IS TOO SMALL FOR PLOTTING:	P RTP0081
IF (YSCAL.EQ.0.0) GO TO 100	P RTP0082
C OTHERWISE, A DIVIDE CHECK WILL OCCUR AFTER STATEMENT 56.	P RTP0083
C FIND BASE VARIABLE PRINT POSITION	P RTP0084
XB=A(1)	P RTP0085
L=1	P RTP0086
MY=M-1	P RTP0087
IT=ISP-1	P RTP0088
DO 80 I=1,NLL	P RTP0089
F=I-1	P RTP0090
XPR=XB+F*XSCAL	P RTP0091
IF(A(L)-XPR-XSCAL*0.5) 50,50,70	P RTP0092
C FIND CROSS-VARIABLES	P RTP0093
50 DO 55 IX=1,101	P RTP0094
55 OUT(IX)=BLANK	P RTP0095
57 CONTINUE	P RTP0096
DO 60 J=1,MY	P RTP0097
56 LL=L+J*N	P RTP0098
JP=((A(LL)-YMIN)/YSCAL)+1.0	P RTP0099
OUT(JP)=IANG(J)	P RTP0100
60 CONTINUE	P RTP0101
C PRINT LINE AND CLEAR, OR SKIP	P RTP0102
IF(L.EQ.N) GO TO 61	P RTP0103
L=L+1	P RTP0104
IF(A(L)-XPR-XSCAL*0.5) 57,57,61	P RTP0105
61 CONTINUE	P RTP0106
WRITE(6,2) XPR, (OUT(IZ), IZ=1,101)	P RTP0107
IF (IT.EQ.0) GO TO 65	P RTP0108

```

DO 64 IV=1,IT
64 WRITE (6,3)
65 GO TO 80
70 WRITE(6,3)
80 CONTINUE
C      PRINT CROSS-VARIABLES NUMBERS
      WRITE(6,7)
      YPR(1)=YMIN
      DO 90 KN=1,9
90 YPR(KN+1)=YPR(KN)+YSCAL*10.0
      YPR(11)=YMAX
      WRITE(6,8)(YPR(IP),IP=1,11)
      RETURN
100 WRITE (6,101)
101 FORMAT ('NO PLOT IS GENERATED BECAUSE THE SPREAD IN THE Y VARIABLE
X IS TOO SMALL.',/,10X,'(I.E. - EQUALS ZERO UNDER REAL*4.)',/,
X 5X, 'EXECUTION CONTINUING.')
      RETURN
      END

```

```

P RTP0109
P RTP0110
P RTP0111
P RTP0112
P RTP0113
P RTP0114
P RTP0115
P RTP0116
P RTP0117
P RTP0118
P RTP0119
P RTP0120
P RTP0121
P RTP0122
P RTP0123
P RTP0124
P RTP0125
P RTP0126
P RTP0127

```

```

SUBROUTINE PUNCH
C   PUNCHES THE FLUX AND CURRENT AND ADJOINTS OUT AFTER CONVERGENCE.
C   CALLED BY IPUNCH=1.
  IMPLICIT REAL*8 (A-H,O-Z)
  COMMON /B2/ KR
  COMMON /B4/ F(2,201), C(2,201)
  COMMON /B47/ F7(2,201), C7(2,201)
  WRITE (7,1) KR
1  FORMAT (I5)
  N=KR+1
C   PUNCH OUT THE FAST FLUX:
  WRITE (7,10) (F(1,J),C(1,J),F7(1,J),C7(1,J),J=1,N)
C   PUNCH OUT THE THERMAL FLUX:
  WRITE (7,10) (F(2,J),C(2,J),F7(2,J),C7(2,J),J=1,N)
10 FORMAT (4E20.7)
  RETURN
  END

```

```

PNCH0001
PNCH0002
PNCH0003
PNCH0004
PNCH0005
PNCH0006
PNCH0007
PNCH0008
PNCH0009
PNCH0010
PNCH0011
PNCH0012
PNCH0013
PNCH0014
PNCH0015
PNCH0016
PNCH0017

```

F. 2. SOURCE LISTING of Program LINEAR



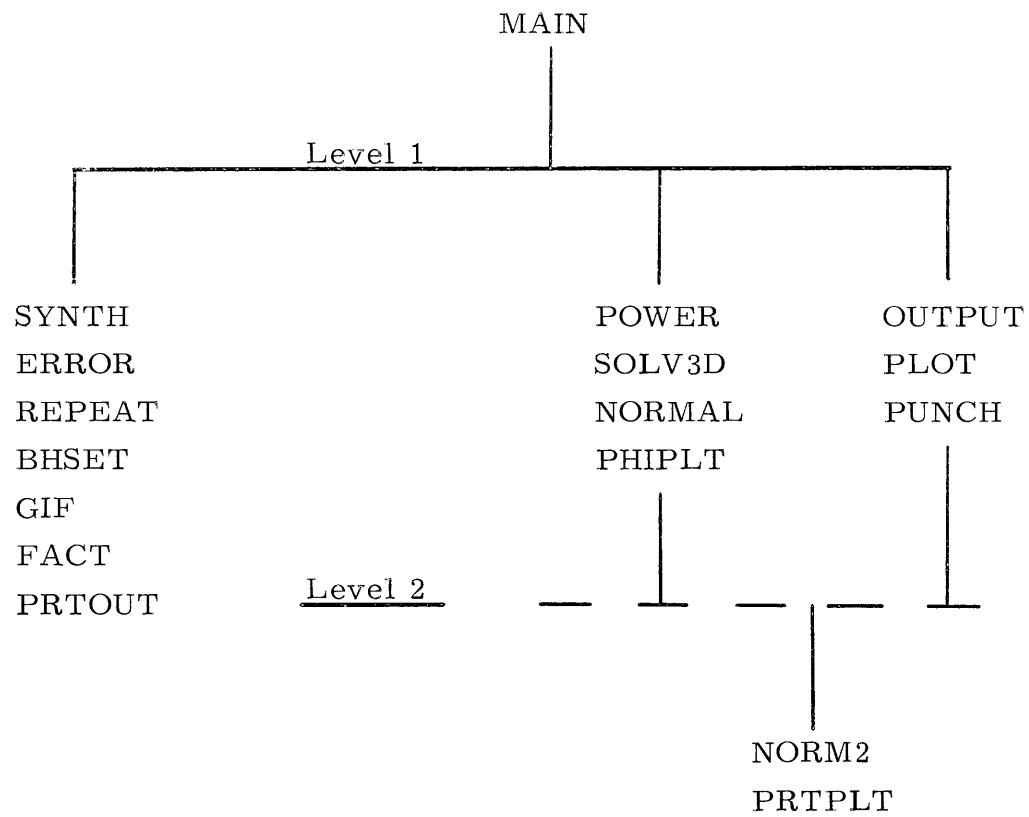


Figure F.2. Structure of Program LINEAR.

C	PROGRAM LINEAR:	LINE0001
C	TWO GROUP PROPOSED METHOD USING LINEAR BASIS FUNCTIONS.	LINE0002
	CALL TIMING(I1)	LINE0003
	CALL SYNTH	LINE0004
	CALL TIMING(I4)	LINE0005
	CALL POWER	LINE0006
	CALL TIMING(I6)	LINE0007
	CALL TIMING(I7)	LINE0008
	CALL OUTPUT	LINE0009
	CALL TIMING(I8)	LINE0010
C	TIMING EXECUTION	LINE0011
	WRITE (6,30)	LINE0012
	30 FORMAT (1H1,'TIMING PROGRAM EXECUTION:',/)	LINE0013
	J=I4-I1	LINE0014
	WRITE(6,701) J	LINE0015
	J=I6-I4	LINE0016
	WRITE(6,704) J	LINE0017
	J=I7-I6	LINE0018
	WRITE(6,706) J	LINE0019
	J=I8-I7	LINE0020
	WRITE(6,707) J	LINE0021
	701 FORMAT (1H , ' SYNTH HAS TAKEN',I6,' /100 SECONDS.')	LINE0022
	704 FORMAT (1H , ' POWER HAS TAKEN',I6,' /100 SECONDS.')	LINE0023
	706 FORMAT (1H , ' CURENT HAS TAKEN',I5,' /100 SECONDS.')	LINE0024
	707 FORMAT (1H , ' OUTPUT HAS TAKEN',I5,' /100 SECONDS.')	LINE0025
	CALL TIMING(I20)	LINE0026
	J=I20-I1	LINE0027
	WRITE(6,720) J	LINE0028
	720 FORMAT (1H0, ' THIS RUN HAS TAKEN',I6,' /100 SECONDS TO RUN.')	LINE0029
	STOP	LINE0030
	END	LINE0031

```

SUBROUTINE SYNTH
C      PROPOSED LINEAR SYNTHESIS METHOD:
C * * * * *
C      ADJOINT QUANTITIES OF VARIABLES ARE DENOTED BY 7 RATHER THAN *.
C      THUS: PHI7 (RATHER THAN PHI*) IS THE ADJOINT OF PHI. ETC.
      IMPLICIT REAL*8 (A-H,K-Z)
      COMMON /B1/ IBC, IPLOT, JPLOT, IPUNCH, ISEE
      COMMON /B2/ KR, NN
      COMMON /B3/ L1(26,26), L2(26,26), F1(26,26), F2(26,26),
X          F3(26,26), F4(26,26), T(26,26)
      COMMON /B5/ KA0(2,25), KA1(2,25), KA2(2,25), KB0(2,25), KB1(2,25),
X          KB2(2,25), LA0(2,25), LA1(2,25), LA2(2,25), SR0(1,25),
X          SR1(1,25), SR2(1,25), KC0(1,25), KC1(1,25), KC2(1,25),
X          KD0(1,25), KD1(1,25), KD2(1,25),
X          P(2,25), P1(2,25), Q(2,25),
X          Q1(2,25), R(2,25), P0(2,25), P07(2,25), PH(2,25),
X          PH7(2,25), AL(2,25), BL(2,25), CL(2,25), AF(4,25),
X          BF(4,25), CF(4,25), AT(25), BT(25), CT(25),
X          ALK(2), BLK(2), AFK(4), BFK(4), ATK(2),
X          BTK(2), BLO(2), CLO(2), BFO(4), CFO(4), BTO(2), CTO(2),
X          CO(2), CH(2)
      COMMON /CHIF/ CHI(2)
      COMMON /XAXIS/ HX, HR(25)
      COMMON /BH/ X(101), H(101)
      COMMON /ER/ EPS1, EPS2, EPS3
      DIMENSION PHI(2,101), PHI7(2,101), CUR(2,101), CUR7(2,101),
X          A(2,100), F(2,100), D(2,100), S(2,100), DI(2,100),
X          XU(2,100)
      DIMENSION V(2), V1(2), V2(2), V3(2)
      DIMENSION ITF(25), KTF(25)
C      IN ORDER TO SAVE SPACE:
      EQUIVALENCE (PHI(1), L1(1)), (PHI7(1), L1(301)),
X          (CUR(1), L2(1)), (CUR7(1), L2(301)),
X          (XU(1), F1(1)), (A(1), F1(301)),
X          (F(1), F2(1)), (D(1), F2(301)),
X          (S(1), T(1)), (DI(1), T(301))

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SYNT0001
SYNT0002
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SYNT0036

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REAL TITLE(20)
INTEGER KR,K,KS,KS1,KRO,NN,NUMITF,KTF
READ (5,200) TITLE
200 FORMAT (20A4)
WRITE (6,201) TITLE
201 FORMAT (1H1,20A4,/)
C      READ IN THE NUMBER OF REGION TRIAL FUNCTIONS AND TYPE OF B.C.S.
C      AS WELL AS THE TOLERANCES AND THE OUTPUT TYPES DESIRED:
READ (5,1) KR,IBC,EPS1,EPS2,EPS3,IPLCT,JPLOT,IPUNCH,ISEE,ITW,ITC
1 FORMAT (2I5,3D10.3,6I5)
IF (IBC.EQ.3) IBC=2
C      READ IN THE TYPE-NUMBER OF EACH TF REGION:
READ (5,100) (ITF(I),I=1,KR)
100 FORMAT (25I2)
C      READ IN THE FISSION YIELDS FOR EACH GROUP:
READ (5,101) CHI(1), CHI(2)
101 FORMAT (2F10.5)
KRO=KR-1
WRITE (6,2) KR, IBC, ISEE, ITW, ITC
2 FORMAT ('0DNE DIMENSIONAL TWO GROUP LINEAR SYNTHESIS PROGRAM:',//,
X 5X,'NUMBER OF COARSE MESH REGIONS: KR = ',I2,/,
X 5X,'BOUNDARY CONDITION NUMBER: IBC = ',I2,/,
X 5X,'AMOUNT OF OUTPUT REQUESTED: ISEE = ',I2,//,
X 5X,'TYPE OF WEIGHTING FUNCTIONS: ITW = ',I2,/,
X 5X,'TYPE OF CURRENT FUNCTIONS: ITC = ',I2,//,
X 5X,'REGIONAL INPUT MATERIAL PROPERTIES AND FLUX SHAPES FOLLOW',
X /,5X,'IF ISEE > 0:',//,
X 5X,'FLUX SHAPES ARE LINEAR IN EACH INDICATED SUBREGION.')
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IF (ITC.EQ.0) WRITE (6,16)
IF (ITC.EQ.1) WRITE (6,17)
16 FORMAT (5X,'CURRENTS ARE CONSTANT IN EACH INDICATED SUBREGION.')
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17 FORMAT (5X,'CURRENTS ARE LINEAR IN EACH INDICATED SUBREGION.')
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IF (ITW.EQ.0) WRITE (6,116)
IF (ITW.EQ.1) WRITE (6,117)
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116 FORMAT (/,5X,'WEIGHTING FLUX = FLUX;',/,5X,'WEIGHTING CURRENT = -
XCURRENT.')
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SYNT0037
SYNT0038
SYNT0039
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SYNT0065
SYNT0066
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SYNT0068
SYNT0069
SYNT0070
SYNT0071
SYNT0072
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117	FORMAT (/ ,5X, 'WEIGHTING FLUX = ADJOINT FLUX; ', / ,5X, 'WEIGHTING CURR XENT = ADJOINT CURRENT.')	SYNT0073
	WRITE (6,20) EPS1, EPS2, EPS3, IPLOT, JPLOT, IPUNCH	SYNT0074
20	FORMAT (/ / , 'TOLERANCES TO POWER ARE : EPS1 = ', 1PD10.3, / , X 28X, 'EPS2 = ', 1PD10.3, / , 28X, 'EPS3 = ', 1PD10.3, / , X 'OUTPUT PARAMETERS TO POWER ARE: IPLOT = ', I1, / , X 34X, 'JPLOT = ', I1, / , 34X, 'IPUNCH = ', I1)	SYNT0075
	WRITE (6,22) CHI(1), CHI(2)	SYNT0076
22	FORMAT (/ , 'FISSION YIELDS ARE: CHI(1) = ', F10.5, / , X 22X, 'CHI(2) = ', F10.5)	SYNT0077
	IF ((KR.LE.2).AND.(IBC.EQ.1)) CALL ERROR(1,KR)	SYNT0078
	IF (KR.GT.25) CALL ERROR(2,KR)	SYNT0079
	IF (EPS1.LT.1.0E-16) CALL ERROR(6,1)	SYNT0080
	IF (EPS2.LT.1.0E-16) CALL ERROR(6,2)	SYNT0081
	IF (EPS3.LT.1.0E-16) CALL ERROR(6,3)	SYNT0082
	IF ((IBC.LT.1).OR.(IBC.GT.7)) CALL ERROR(7,IBC)	SYNT0083
C	DUMMY NORMAL VECTOR XU = UNITY. (FOR THE INTEGRATION FUNCTIONS)	SYNT0084
	DO 21 IG=1,2	SYNT0085
	DO 21 II=1,100	SYNT0086
21	XU(IG,II)=1.0	SYNT0087
	ITC0=2	SYNT0088
	ITC1=2	SYNT0089
	IF (ITC.EQ.1) GO TO 23	SYNT0090
	ITC0=0	SYNT0091
	ITC1=1	SYNT0092
C	COUNTER OF THE NUMBER OF TYPE-NUMBERS OF EACH TF REGION:	SYNT0093
23	NUMITF=1	SYNT0094
	HX=0.0	SYNT0095
C	BEGIN TO READ IN THE TF REGION DATA AND FILL THE ARRAYS,	SYNT0096
C	DEPENDING ON THE TYPE-NUMBER OF EACH TF REGION.	SYNT0097
	DO 50 I=1,KR	SYNT0098
	IF (ITF(I).EQ.NUMITF) GO TO 110	SYNT0099
C	FILL THE ARRAYS FROM OLD TF REGION TYPES:	SYNT0100
	J=ITF(I)	SYNT0101
	CALL REPEAT(I,KTF(J))	SYNT0102
	GO TO 50	SYNT0103
		SYNT0104
		SYNT0105
		SYNT0106
		SYNT0107
		SYNT0108

C	READ IN THE TF REGION'S DATA FOR NEW TF REGION TYPE-NUMBERS:	SYNT0109
110	NUMITF=NUMITF+1	SYNT0110
	KTF(NUMITF-1)=I	SYNT0111
C	READ THE SUBREGION NUMBER AND THE NUMBER OF REGIONS IN THE SUBREGION:	SYNT0112
	READ (5,1) K, KS	SYNT0113
	IF (KS.GT.100) CALL ERROR(3,I)	SYNT0114
C	CHECK FOR IMPROPER SEQUENCING OF INPUT DATA:	SYNT0115
	IF (I.NE.K) CALL ERROR(4,I)	SYNT0116
C	READ IN THE GEOMETRY AND THE MATERIAL PROPERTIES:	SYNT0117
	READ (5,3) (X(J),X(J+1),H(J),A(1,J),F(1,J),D(1,J),S(1,J),	SYNT0118
	X                                  A(2,J),F(2,J),D(2,J),J=1,KS)	SYNT0119
3	FORMAT (3F10.5,4D10.3,/,30X,3D10.3)	SYNT0120
C	READ IN THE REGIONAL GROUP TRIAL FUNCTIONS:	SYNT0121
	KS1=KS+1	SYNT0122
	READ (5,4) (PHI(1,J),CUR(1,J),PHI7(1,J),CUR7(1,J),J=1,KS1)	SYNT0123
	READ (5,4) (PHI(2,J),CUR(2,J),PHI7(2,J),CUR7(2,J),J=1,KS1)	SYNT0124
4	FORMAT (4D20.7)	SYNT0125
	IF (ITW.EQ.1) GO TO 120	SYNT0126
C	FORM WEIGHTING FUNCTIONS FROM THE GIVEN FUNCTIONS:	SYNT0127
	DO 119 IG=1,2	SYNT0128
	DO 119 J=1,KS1	SYNT0129
	PHI7(IG,J)=PHI(IG,J)	SYNT0130
119	CUR7(IG,J)=-CUR(IG,J)	SYNT0131
120	IF (ITC.EQ.1) GO TO 5	SYNT0132
C	FORM THE REGION CONSTANT CURRENTS:	SYNT0133
	DO 7 IG=1,2	SYNT0134
	DO 6 J=1,KS	SYNT0135
	CUR(IG,J)=-D(IG,J)*(-PHI(IG,J)+PHI(IG,J+1))/H(J)	SYNT0136
6	CUR7(IG,J)=+D(IG,J)*(-PHI7(IG,J)+PHI7(IG,J+1))/H(J)	SYNT0137
	CUR(IG,KS1)=0.0	SYNT0138
7	CUR7(IG,KS1)=0.0	SYNT0139
C	WRITING OUT THE INPUT INFORMATION:	SYNT0140
5	IF (ISEE.EQ.0) GO TO 14	SYNT0141
	WRITE (6,10) K,KR,KS,(J,X(J),X(J+1),H(J),A(1,J),F(1,J),D(1,J),	SYNT0142
	X    S(1,J),A(2,J),F(2,J),D(2,J),J=1,KS)	SYNT0143
10	FORMAT ('INPUT MATERIAL PROPERTIES FOR REGION NUMBER ',I3,	SYNT0144

X	' , OF THE ',I3,' USED.',//,	SYNT0145
X	5X,' THIS REGION IS DIVIDED INTO ',I3,' HOMOGENEOUS SUBREGIONS A	SYNT0146
XS	FOLLOWS:',//,	SYNT0147
X	5X,' FAST GROUP CONSTANTS APPEAR FIRST:',//,	SYNT0148
X	' SUBREGION #',5X,' INTERNAL BOUNDARIES',10X,' WIDTH',3X,	SYNT0149
X	' TOTAL CX (1/CM)',3X,' FISSION CX (1/CM)',6X,' DIFFUSION (CM)',	SYNT0150
X	4X,' SCATT.CX (1/CM)',/,	SYNT0151
X	5X,' I',11X,' X(I)',9X,' X(I+1)',11X,' H(I)',13X,' A(IG,I)',13X,	SYNT0152
X	' F(IG,I)',13X,' D(IG,I)',14X,' S(1,I)',//,	SYNT0153
X	(I6,3F15.4,4D20.8,/,51X,3D20.8))	SYNT0154
C		SYNT0155
	DO 15 IG=1,2	SYNT0156
15	WRITE (6,11) IG,K,KR,{J,X(J),PHI(IG,J),CUR(IG,J),PHI7(IG,J),	SYNT0157
X	CUR7(IG,J),J=1,KS1)	SYNT0158
11	FORMAT ('INPUT TRIAL FUNCTIONS FOR GROUP',I2,' FOR REGION',I3,	SYNT0159
X	' OUT OF THE ',I3,' USED:',//,	SYNT0160
X	' INDEX',5X,' COORD',16X,' FLUX',13X,' CURRENT',8X,' WEIGHT FLUX',	SYNT0161
X	5X,' WEIGHT CURRENT',//,(I6,F10.5,4D20.7))	SYNT0162
14	CONTINUE	SYNT0163
C	END OF THE IN-OUT SECTION:	SYNT0164
C	DEFINING MISC. ARRAYS FOR THE INTEGRATION FUNCTIONS:	SYNT0165
C	LENGTH OF THE SUBREGION: HT	SYNT0166
	HT=X(KS1)-X(1)	SYNT0167
	HR(K)=HT	SYNT0168
	HX=HX+HR(K)	SYNT0169
C	INVERSE OF THE D ARRAYS:	SYNT0170
	DO 13 J=1,KS	SYNT0171
	DI(1,J)=1./D(1,J)	SYNT0172
13	DI(2,J)=1./D(2,J)	SYNT0173
C	FORMATION OF THE INTEGRATION FUNCTIONS:	SYNT0174
	CALL BHSET(KS)	SYNT0175
C	DO FOR ALL ENERGY GROUPS:	SYNT0176
	DO 50 IG=1,2	SYNT0177
	KA0(IG,K)=GIF(0,IG,PHI7,IG,A,PHI,KS,2)	SYNT0178
	KA1(IG,K)=GIF(1,IG,PHI7,IG,A,PHI,KS,2)	SYNT0179
	KA2(IG,K)=GIF(2,IG,PHI7,IG,A,PHI,KS,2)	SYNT0180

KB0(IG,K)=GIF(0,IG,PHI7,IG,F,PHI,KS,2)	SYNT0181
KB1(IG,K)=GIF(1,IG,PHI7,IG,F,PHI,KS,2)	SYNT0182
KB2(IG,K)=GIF(2,IG,PHI7,IG,F,PHI,KS,2)	SYNT0183
LA0(IG,K)=GIF(0,IG,CUR7,IG,DI,CUR,KS,ITC0)	SYNT0184
LA1(IG,K)=GIF(1,IG,CUR7,IG,DI,CUR,KS,ITC0)	SYNT0185
LA2(IG,K)=GIF(2,IG,CUR7,IG,DI,CUR,KS,ITC0)	SYNT0186
P(IG,K) =GIF(0,IG,PHI7,IG,XU,CUR,KS,ITC1)/HT	SYNT0187
P1(IG,K) =GIF(1,IG,PHI7,IG,XU,CUR,KS,ITC1)/HT	SYNT0188
Q(IG,K) =GIF(0,IG,PHI,IG,XU,CUR7,KS,ITC1)/HT	SYNT0189
Q1(IG,K) =GIF(1,IG,PHI,IG,XU,CUR7,KS,ITC1)/HT	SYNT0190
R(IG,K) =GIF(0,IG,PHI7,IG,D,PHI,KS,2)/HT**2	SYNT0191
C STORE THE TERMINAL POINTS FOR LATER USE:	SYNT0192
PO(IG,K)=PHI(IG,1)	SYNT0193
PO7(IG,K)=PHI7(IG,1)	SYNT0194
PH(IG,K)=PHI(IG,KS1)	SYNT0195
PH7(IG,K)=PHI7(IG,KS1)	SYNT0196
IF (K.EQ.1) CO(IG)=CUR(IG,1)	SYNT0197
IF (NUMITF-1.EQ.ITF(KR).AND.ITC.EQ.C) CH(IG)=CUR(IG,KS)	SYNT0198
IF (NUMITF-1.EQ.ITF(KR).AND.ITC.EQ.1) CH(IG)=CUR(IG,KS1)	SYNT0199
C FOR THE OFF DIAGONAL MATRIX ELEMENTS:	SYNT0200
IF (IG.EQ.2) GO TO 50	SYNT0201
SR0(IG,K)=GIF(0,2,PHI7,1,S,PHI,KS,2)	SYNT0202
SR1(IG,K)=GIF(1,2,PHI7,1,S,PHI,KS,2)	SYNT0203
SR2(IG,K)=GIF(2,2,PHI7,1,S,PHI,KS,2)	SYNT0204
KC0(IG,K)=GIF(0,1,PHI7,2,F,PHI,KS,2)	SYNT0205
KC1(IG,K)=GIF(1,1,PHI7,2,F,PHI,KS,2)	SYNT0206
KC2(IG,K)=GIF(2,1,PHI7,2,F,PHI,KS,2)	SYNT0207
KD0(IG,K)=GIF(0,2,PHI7,1,F,PHI,KS,2)	SYNT0208
KD1(IG,K)=GIF(1,2,PHI7,1,F,PHI,KS,2)	SYNT0209
KD2(IG,K)=GIF(2,2,PHI7,1,F,PHI,KS,2)	SYNT0210
50 CONTINUE	SYNT0211
NUMITF=NUMITF-1	SYNT0212
WRITE (6,51) NUMITF	SYNT0213
51 FORMAT ('1THERE ARE ONLY',I3,' DIFFERENT TRIAL FUNCTION REGIONS.')	SYNT0214
WRITE (6,52) (I,ITF(I),I=1,KR)	SYNT0215
52 FORMAT (/, 'OTABLE OF THE TRIAL FUNCTION NUMBER TYPES:',//,	SYNT0216



X	3X,'TF REGION',4X,'REGION TYPE-NUMBER',//,	SYNT0217
X	(I7,12X,I7))	SYNT0218
C	PRINTS OUT THE /B5/ ARRAYS:	SYNT0219
	IF (ISEE.GE.2) CALL PRTOUT(1)	SYNT0220
C	NN IS THE MATRIX BLOCK SIZE:	SYNT0221
	NN=KR	SYNT0222
	IF (IBC.EQ.1.OR.IBC.GE.6) NN=KR-1	SYNT0223
	IF (IBC.EQ.4) NN=KR+1	SYNT0224
C	FORMATION OF THE COEFFICIENT VECTORS:	SYNT0225
C	THE INTERIOR COEFFICIENTS:	SYNT0226
	IX=2	SYNT0227
	IY=KR	SYNT0228
	IF (IBC.EQ.5.OR.IBC.EQ.7) IX=3	SYNT0229
	IF (IBC.EQ.1.OR.IBC.GE.6) IY=KR-1	SYNT0230
	DO 60 IG=1,2	SYNT0231
	DO 60 K=IX,IY	SYNT0232
	J=K-1	SYNT0233
	V(IG)=1./(PH7(IG,J)*P0(IG,J))	SYNT0234
	V1(IG)=1./(PH7(IG,J)*PH(IG,J))	SYNT0235
	V2(IG)=1./(P07(IG,K)*P0(IG,K))	SYNT0236
	V3(IG)=1./(P07(IG,K)*PH(IG,K))	SYNT0237
	AL(IG,K)=(KA1(IG,J)-KA2(IG,J)-R(IG,J)+LA2(IG,J)-LA1(IG,J)	SYNT0238
X	-Q1(IG,J)+P1(IG,J)-P(IG,J))*V(IG)	SYNT0239
	BL(IG,K)=(KA2(IG,J)-LA2(IG,J)+Q1(IG,J)-P1(IG,J)+R(IG,J))*V1(IG)	SYNT0240
X	+(KA0(IG,K)-2.*KA1(IG,K)+KA2(IG,K)-LA0(IG,K)+2.*LA1(IG,K)	SYNT0241
X	-LA2(IG,K)+Q1(IG,K)-Q(IG,K)+P(IG,K)-P1(IG,K)+R(IG,K))*V2(IG)	SYNT0242
	CL(IG,K)=(KA1(IG,K)-KA2(IG,K)-R(IG,K)+LA2(IG,K)-LA1(IG,K)+Q(IG,K)	SYNT0243
X	-Q1(IG,K)+P1(IG,K))*V3(IG)	SYNT0244
	AF(IG,K)=(KB1(IG,J)-KB2(IG,J))*V(IG)	SYNT0245
	BF(IG,K)=KB2(IG,J)*V1(IG)+(KB0(IG,K)-2.*KB1(IG,K)+KB2(IG,K))	SYNT0246
X	*V2(IG)	SYNT0247
	CF(IG,K)=(KB1(IG,K)-KB2(IG,K))*V3(IG)	SYNT0248
	IF (IG.EQ.2) GO TO 60	SYNT0249
	AT(K)=(SR1(1,J)-SR2(1,J))/(PH7(2,J)*P0(1,J))	SYNT0250
	AF(3,K)=(KC1(IG,J)-KC2(IG,J))/(PH7(1,J)*P0(2,J))	SYNT0251
	AF(4,K)=(KD1(IG,J)-KD2(IG,J))/(PH7(2,J)*P0(1,J))	SYNT0252

	BT(K)=SR2(1,J)/(PH7(2,J)*PH(1,J))	SYNT0253
X	+(SR0(1,K)-2.*SR1(1,K)+SR2(1,K))/(P07(2,K)*P0(1,K))	SYNT0254
	BF(3,K)=KC2(IG,J)/(PH7(1,J)*PH(2,J))	SYNT0255
X	+(KC0(IG,K)-2.*KC1(IG,K)+KC2(IG,K))/(P07(1,K)*P0(2,K))	SYNT0256
	BF(4,K)=KD2(IG,J)/(PH7(2,J)*PH(1,J))	SYNT0257
X	+(KD0(IG,K)-2.*KD1(IG,K)+KD2(IG,K))/(P07(2,K)*P0(1,K))	SYNT0258
	CT(K)=(SR1(1,K)-SR2(1,K))/(P07(2,K)*PH(1,K))	SYNT0259
	CF(3,K)=(KC1(IG,K)-KC2(IG,K))/(P07(1,K)*PH(2,K))	SYNT0260
	CF(4,K)=(KD1(IG,K)-KD2(IG,K))/(P07(2,K)*PH(1,K))	SYNT0261
60	CONTINUE	SYNT0262
C	THE ZERO FLUX COEFFS:	SYNT0263
C	NONE NEEDED AS F(1,1) AND F(2,1) BOTH = 0.0.	SYNT0264
	IF (IBC.EQ.1) GO TO 64	SYNT0265
C		SYNT0266
	IF (.NOT.(IBC.EQ.5.OR.IBC.EQ.7)) GO TO 66	SYNT0267
C	ZERO FLUX COEFFICIENTS FOR NO TILTING ON THE LEFT:	SYNT0268
	DO 61 IG=1,2	SYNT0269
	V(IG)=1./(PH7(IG,1)*PH(IG,1))	SYNT0270
	V1(IG)=1./(P07(IG,2)*P0(IG,2))	SYNT0271
	V2(IG)=1./(P07(IG,2)*PH(IG,2))	SYNT0272
	BLO(IG)=(KA0(IG,1)-LA0(IG,1))*V(IG)+(KA0(IG,2)-2.*KA1(IG,2)	SYNT0273
X	+KA2(IG,2)-LA0(IG,2)+2.*LA1(IG,2)-LA2(IG,2)+P(IG,2)-P1(IG,2)	SYNT0274
X	-Q(IG,2)+Q1(IG,2)+R(IG,2))*V1(IG)	SYNT0275
	CLO(IG)=(KA1(IG,2)-KA2(IG,2)+LA2(IG,2)-LA1(IG,2)+P1(IG,2)	SYNT0276
X	+Q(IG,2)-Q1(IG,2)-R(IG,2))*V2(IG)	SYNT0277
	BF0(IG)=KB0(IG,1)*V(IG)+(KB0(IG,2)-2.*KB1(IG,2)+KB2(IG,2))*V1(IG)	SYNT0278
61	CF0(IG)=(KB1(IG,2)-KB2(IG,2))*V2(IG)	SYNT0279
	BT0(1)=SR0(1,1)/(PH7(2,1)*PH(1,1))	SYNT0280
X	+(SR0(1,2)-2.*SR1(1,2)+SR2(1,2))/(P07(2,2)*P0(1,2))	SYNT0281
	BF0(3)=KC0(1,1)/(PH7(1,1)*PH(2,1))	SYNT0282
X	+(KC0(1,2)-2.*KC1(1,2)+KC2(1,2))/(P07(1,2)*P0(2,2))	SYNT0283
	BF0(4)=KD0(1,1)/(PH7(2,1)*PH(1,1))	SYNT0284
X	+(KD0(1,2)-2.*KD1(1,2)+KD2(1,2))/(P07(2,2)*P0(1,2))	SYNT0285
	CT0(1)=(SR1(1,2)-SR2(1,2))/(P07(2,2)*PH(1,2))	SYNT0286
	CF0(3)=(KC1(1,2)-KC2(1,2))/(P07(1,2)*PH(2,2))	SYNT0287
	CF0(4)=(KD1(1,2)-KD2(1,2))/(P07(2,2)*PH(1,2))	SYNT0288

66	IF (IBC.LE.5) GO TO 69	SYNT0289
C	FOR THE LAST REGION TO BE NOT TILTED:	SYNT0290
	K=KR-1	SYNT0291
	DO 67 IG=1,2	SYNT0292
	V(IG)=1./(PH7(IG,K)*PO(IG,K))	SYNT0293
	V1(IG)=1./(PO7(IG,KR)*PO(IG,KR))	SYNT0294
	V2(IG)=1./(PH7(IG,K)*PH(IG,K))	SYNT0295
	AL(IG,KR)=(KA1(IG,K)-KA2(IG,K)-LA1(IG,K)+LA2(IG,K)-P(IG,K)	SYNT0296
X	+P1(IG,K)-Q1(IG,K)-R(IG,K))*V(IG)	SYNT0297
	BL(IG,KR)=(KA0(IG,KR)-LA0(IG,KR))*V1(IG)+(KA2(IG,K)-LA2(IG,K)	SYNT0298
X	-P1(IG,K)+Q1(IG,K)+R(IG,K))*V2(IG)	SYNT0299
	AF(IG,KR)=(KB1(IG,K)-KB2(IG,K))*V(IG)	SYNT0300
	BF(IG,KR)=KB0(IG,KR)*V1(IG)+KB2(IG,K)*V2(IG)	SYNT0301
	IF (IG.EQ.2) GO TO 67	SYNT0302
	AT( KR)=(SR1(IG,K)-SR2(IG,K))/(PH7(2,K)*PO(1,K))	SYNT0303
	AF(3,KR)=(KC1(IG,K)-KC2(IG,K))/(PH7(1,K)*PO(2,K))	SYNT0304
	AF(4,KR)=(KD1(IG,K)-KD2(IG,K))/(PH7(2,K)*PO(1,K))	SYNT0305
	BT(KR)=SRO(IG,KR)/(PO7(2,KR)*PO(1,KR))	SYNT0306
X	+SR2(IG,K)/(PH7(2,K)*PH(1,K))	SYNT0307
	BF(3,KR)=KC0(IG,KR)/(PO7(1,KR)*PO(2,KR))	SYNT0308
X	+KC2(IG,K)/(PH7(1,K)*PH(2,K))	SYNT0309
	BF(4,KR)=KDO(IG,KR)/(PO7(2,KR)*PO(1,KR))	SYNT0310
X	+KD2(IG,K)/(PH7(2,K)*PH(1,K))	SYNT0311
67	CONTINUE	SYNT0312
	GO TO 64	SYNT0313
C	THE ZERO CURRENT COEFFS:	SYNT0314
69	K=KR	SYNT0315
	DO 62 IG=1,2	SYNT0316
	V(IG)=1./(PH7(IG,K)*PO(IG,K))	SYNT0317
	V1(IG)=1./(PH7(IG,K)*PH(IG,K))	SYNT0318
	ALK(IG)=(KA1(IG,K)-KA2(IG,K)-R(IG,K)+P1(IG,K)-P(IG,K)+LA2(IG,K)	SYNT0319
X	-LA1(IG,K)-Q1(IG,K))*V(IG)	SYNT0320
	BLK(IG)=(KA2(IG,K)+R(IG,K)-P1(IG,K)-LA2(IG,K)+Q1(IG,K))*V1(IG)	SYNT0321
X	+CH(IG)/PH(IG,KR)	SYNT0322
	AFK(IG)=(KB1(IG,K)-KB2(IG,K))*V(IG)	SYNT0323
62	BFK(IG)=KB2(IG,K)*V1(IG)	SYNT0324

	ATK(1)=(SR1(1,K)-SR2(1,K))/(PH7(2,K)*PO(1,K))	SYNT0325
	AFK(3)=(KC1(1,K)-KC2(1,K))/(PH7(1,K)*PO(2,K))	SYNT0326
	AFK(4)=(KD1(1,K)-KD2(1,K))/(PH7(2,K)*PO(1,K))	SYNT0327
	BTK(1)=SR2(1,K)/(PH7(2,K)*PH(1,K))	SYNT0328
	BFK(3)=KC2(1,K)/(PH7(1,K)*PH(2,K))	SYNT0329
	BFK(4)=KD2(1,K)/(PH7(2,K)*PH(1,K))	SYNT0330
	IF (IBC.NE.4) GO TO 64	SYNT0331
C	ZERU CURRENT ON THE LEFT COEFFS:	SYNT0332
	K=1	SYNT0333
	DO 63 IG=1,2	SYNT0334
	V2(IG)=1./(PO7(IG,K)*PO(IG,K))	SYNT0335
	V3(IG)=1./(PO7(IG,K)*PH(IG,K))	SYNT0336
	BLO(IG)=	SYNT0337
	X +(KA0(IG,K)-2.*KA1(IG,K)+KA2(IG,K)-LA0(IG,K)+2.*LA1(IG,K)	SYNT0338
	X -LA2(IG,K)+Q1(IG,K)-Q(IG,K)+P(IG,K)-P1(IG,K)+R(IG,K))*V2(IG)	SYNT0339
	X -CO(IG)/PO(IG,1)	SYNT0340
	CLO(IG) =(KA1(IG,K)-KA2(IG,K)-R(IG,K)+LA2(IG,K)-LA1(IG,K)+Q(IG,K)	SYNT0341
	X -Q1(IG,K)+P1(IG,K))*V3(IG)	SYNT0342
	BF0(IG)= V2(IG)*(KB0(IG,K)-2.*KB1(IG,K)+KB2(IG,K))	SYNT0343
63	CF0(IG) =(KB1(IG,K)-KB2(IG,K))*V3(IG)	SYNT0344
	BT0(1)=(SR0(1,K)-2.*SR1(1,K)+SR2(1,K))/(PO7(2,K)*PO(1,K))	SYNT0345
	BF0(3)=(KC0(1,K)-2.*KC1(1,K)+KC2(1,K))/(PO7(1,K)*PO(2,K))	SYNT0346
	BF0(4)=(KD0(1,K)-2.*KD1(1,K)+KD2(1,K))/(PO7(2,K)*PO(1,K))	SYNT0347
	CT0(1)=(SR1(1,K)-SR2(1,K))/(PO7(2,K)*PH(1,K))	SYNT0348
	CF0(3)=(KC1(1,K)-KC2(1,K))/(PO7(1,K)*PH(2,K))	SYNT0349
	CF0(4)=(KD1(1,K)-KD2(1,K))/(PO7(2,K)*PH(1,K))	SYNT0350
C	ZERU MATRICES:	SYNT0351
64	DO 65 I=1,NN	SYNT0352
	DO 65 J=1,NN	SYNT0353
	L1(I,J)=0.0	SYNT0354
	L2(I,J)=0.0	SYNT0355
	F1(I,J)=0.0	SYNT0356
	F2(I,J)=0.0	SYNT0357
	F3(I,J)=0.0	SYNT0358
	F4(I,J)=0.0	SYNT0359
65	T(I,J)=0.0	SYNT0360

C FILLING THE MATRICES FOR POWER:  
C DETERMINE THE LEFT BC:

J=1  
IF (IBC.LE.2.OR.IBC.EQ.6) GO TO 75  
L1(J,1)=BL0(1)  
L2(J,1)=BL0(2)  
F1(J,1)=BF0(1)  
F2(J,1)=BF0(2)  
F3(J,1)=BF0(3)  
F4(J,1)=BF0(4)  
T(J,1)=BT0(1)  
L1(J,2)=CL0(1)  
L2(J,2)=CL0(2)  
F1(J,2)=CF0(1)  
F2(J,2)=CF0(2)  
F3(J,2)=CF0(3)  
F4(J,2)=CF0(4)  
T(J,2)=CT0(1)  
J=J+1

C FOR ALL INTERNAL EQUATIONS:

75 DO 70 K=IX,IY  
IF (J.EQ.1) GO TO 76  
L1(J,J-1)=AL(1,K)  
L2(J,J-1)=AL(2,K)  
F1(J,J-1)=AF(1,K)  
F2(J,J-1)=AF(2,K)  
F3(J,J-1)=AF(3,K)  
F4(J,J-1)=AF(4,K)  
T(J,J-1)=AT(K)  
76 L1(J,J)=BL(1,K)  
L1(J,J+1)=CL(1,K)  
L2(J,J)=BL(2,K)  
L2(J,J+1)=CL(2,K)  
F1(J,J)=BF(1,K)  
F1(J,J+1)=CF(1,K)  
F2(J,J)=BF(2,K)

SYNT0361  
SYNT0362  
SYNT0363  
SYNT0364  
SYNT0365  
SYNT0366  
SYNT0367  
SYNT0368  
SYNT0369  
SYNT0370  
SYNT0371  
SYNT0372  
SYNT0373  
SYNT0374  
SYNT0375  
SYNT0376  
SYNT0377  
SYNT0378  
SYNT0379  
SYNT0380  
SYNT0381  
SYNT0382  
SYNT0383  
SYNT0384  
SYNT0385  
SYNT0386  
SYNT0387  
SYNT0388  
SYNT0389  
SYNT0390  
SYNT0391  
SYNT0392  
SYNT0393  
SYNT0394  
SYNT0395  
SYNT0396

```

F2(J,J+1)=CF(2,K)
F3(J,J)=BF(3,K)
F3(J,J+1)=CF(3,K)
F4(J,J)=BF(4,K)
F4(J,J+1)=CF(4,K)
T(J,J)=BT(K)
T(J,J+1)=CT(K)
J=J+1
70 CONTINUE
IF (IBC.EQ.1.OR.IBC.EQ.5) GO TO 500
C   FOR THE LAST EQUATION: (K=NN):
K=NN
L1(K ,K -1)=ALK(1)
L1(K ,K )=BLK(1)
L2(K ,K -1)=ALK(2)
L2(K ,K )=BLK(2)
F1(K ,K -1)=AFK(1)
F1(K ,K )=BFK(1)
F2(K ,K -1)=AFK(2)
F2(K ,K )=BFK(2)
F3(K,K-1)=AFK(3)
F3(K,K)=BFK(3)
F4(K,K-1)=AFK(4)
F4(K,K)=BFK(4)
T(K ,K -1)=ATK(1)
T(K ,K )=BTK(1)
500 CONTINUE
C   PRINTS OUT THE /B3/ MATRICES:
IF (ISEE.GE.2) CALL PRTOUT(2)
RETURN
END

```

```

SYNT0397
SYNT0398
SYNT0399
SYNT0400
SYNT0401
SYNT0402
SYNT0403
SYNT0404
SYNT0405
SYNT0406
SYNT0407
SYNT0408
SYNT0409
SYNT0410
SYNT0411
SYNT0412
SYNT0413
SYNT0414
SYNT0415
SYNT0416
SYNT0417
SYNT0418
SYNT0419
SYNT0420
SYNT0421
SYNT0422
SYNT0423
SYNT0424
SYNT0425
SYNT0426
SYNT0427

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<pre> SUBROUTINE ERROR(I,J) C   ANNOUNCES INPUT ERRORS AND TERMINATES PROGRAM EXECUTION: GO TO (1,2,3,4,5,6,7,8,9),I 1 WRITE (6,101) GO TO 10 2 WRITE (6,102) J GO TO 10 3 WRITE (6,103) J GO TO 10 4 WRITE (6,104) J GO TO 10 5 WRITE (6,105) J GO TO 10 6 WRITE (6,106) J GO TO 10 7 CONTINUE 8 CONTINUE 9 CONTINUE 10 WRITE (6,110) 101 FORMAT ('1MUST HAVE &gt; 2 SUBREGIONS FOR ZERO FLUX B.C.S. INVALID.')</pre>	<pre> 102 FORMAT ('1NUMBER OF SUBREGIONS =',I3,' &gt; 25. INVALID.')</pre>	<pre> 103 FORMAT ('1SUBREGION NUMBER',I3,' HAS &gt; 25 SECTIONS. INVALID.')</pre>	<pre> 104 FORMAT ('1INPUT ERROR IN REGION SEQUENCING AT REGION',I5,'.')</pre>	<pre> 105 FORMAT ('1Z(I) = 0. IN REGION I =',I3,'. INVALID.')</pre>	<pre> 106 FORMAT ('1THE TOLERANCE: EPS',I1,' IS &lt; 1.0E-16. INVALID.')</pre>	<pre> 107 FORMAT ('1BOUNDARY CONDITION OPTICN =',I2,' &lt; 1 OR &gt; 7. INVALID.')</pre>	<pre> 110 FORMAT (1H0,'PROBLEM TERMINATED.')</pre>	<pre> CALL EXIT RETURN END</pre>	<pre> ERRO0001 ERRO0002 ERRO0003 ERRO0004 ERRO0005 ERRO0006 ERRO0007 ERRO0008 ERRO0009 ERRO0010 ERRO0011 ERRO0012 ERRO0013 ERRO0014 ERRO0015 ERRO0016 ERRO0017 ERRO0018 ERRO0019 ERRO0020 ERRO0021 ERRO0022 ERRO0023 ERRO0024 ERRO0025 ERRO0026 ERRO0027 ERRO0028 ERRO0029 ERRO0030</pre>
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SUBROUTINE REPEAT(K,L)
C   SETS THE /B5/ ARRAYS (K) EQUAL TO PAST STORED ARRAYS (L):
    IMPLICIT REAL*8 (A-Z)
    COMMON /B5/ KA0(2,25),KA1(2,25),KA2(2,25),KBO(2,25),KB1(2,25),
X      KB2(2,25),LA0(2,25),LA1(2,25),LA2(2,25),SR0(1,25),
X      SR1(1,25),SR2(1,25),KCO(1,25),KC1(1,25),KC2(1,25),
X      KDO(1,25),KD1(1,25),KD2(1,25),
X      P(2,25),P1(2,25),Q(2,25),
X      Q1(2,25), R(2,25), P0(2,25), P07(2,25),PH(2,25),
X      PH7(2,25)
    COMMON /XAXIS/ HX,HR(25)
    INTEGER K,L,G
    HR(K)=HR(L)
    HX=HX+HR(K)
    DO 10 G=1,2
    KA0(G,K)=KA0(G,L)
    KA1(G,K)=KA1(G,L)
    KA2(G,K)=KA2(G,L)
    KBO(G,K)=KBO(G,L)
    KB1(G,K)=KB1(G,L)
    KB2(G,K)=KB2(G,L)
    LA0(G,K)=LA0(G,L)
    LA1(G,K)=LA1(G,L)
    LA2(G,K)=LA2(G,L)
    IF (G.EQ.2) GO TO 5
    SR0(G,K)=SR0(G,L)
    SR1(G,K)=SR1(G,L)
    SR2(G,K)=SR2(G,L)
    KCO(G,K)=KCO(G,L)
    KC1(G,K)=KC1(G,L)
    KC2(G,K)=KC2(G,L)
    KDO(G,K)=KDO(G,L)
    KD1(G,K)=KD1(G,L)
    KD2(G,K)=KD2(G,L)
5  CONTINUE
    P(G,K)=P(G,L)

```

```

REPE0001
REPE0002
REPE0003
REPE0004
REPE0005
REPE0006
REPE0007
REPE0008
REPE0009
REPE0010
REPE0011
REPE0012
REPE0013
REPE0014
REPE0015
REPE0016
REPE0017
REPE0018
REPE0019
REPE0020
REPE0021
REPE0022
REPE0023
REPE0024
REPE0025
REPE0026
REPE0027
REPE0028
REPE0029
REPE0030
REPE0031
REPE0032
REPE0033
REPE0034
REPE0035
REPE0036

```



P1(G,K)=P1(G,L)  
Q(G,K)=Q(G,L)  
Q1(G,K)=Q1(G,L)  
R(G,K)=R(G,L)  
P0(G,K)=P0(G,L)  
P07(G,K)=P07(G,L)  
PH(G,K)=PH(G,L)  
PH7(G,K)=PH7(G,L)  
10 CONTINUE  
RETURN  
END

REPE0037  
REPE0038  
REPE0039  
REPE0040  
REPE0041  
REPE0042  
REPE0043  
REPE0043  
REPE0044  
REPE0045  
REPE0046  
REPE0047

```
C      SUBROUTINE BHSET(K)
        SETS UP THE /BH/ ARRAYS FOR GIF:
        IMPLICIT REAL*8 (A-H,L-Z)
        COMMON /BH/ X(101), H(101), Z(101)
        DO 1 I=1,K
1      Z(I)=X(I)-X(1)
        RETURN
        END
```

```
BHSE0001
BHSE0002
BHSE0003
BHSE0004
BHSE0005
BHSE0006
BHSE0007
BHSE0008
```

```

C      DOUBLE PRECISION FUNCTION GIF(N,G1,F,G2,C,G,K,ITC)          GIF 0001
C      INTEGRATES: F(G1,J)*C(G2,J)*G(G2,J) * (Z/H)**N          GIF 0002
C      OVER ALL K SUBREGIONS J                                   GIF 0003
C      WHERE Z RUNS FROM 0 TO X(K+1)-X(1) IN THIS REGION.       GIF 0004
C      WHERE THE FORM OF F AND G IN REGION J IS GIVEN BY ITC:  GIF 0005
C      ITC = 0: F AND G ARE BOTH CONSTANT.                       GIF 0006
C      ITC = 1: F IS LINEAR AND G IS CCNSTANT.                  GIF 0007
C      ITC = 2: F AND G ARE BOTH LINEAR.                         GIF 0008
C      IMPLICIT REAL*8 (A-H,O-Z)                                  GIF 0009
C      COMMON /BH/ X(101),H(101),Z(101)                          GIF 0010
C      DIMENSION F(2,101),G(2,101),C(2,100)                     GIF 0011
C      INTEGER G1,G2                                             GIF 0012
C      N1=N+1                                                    GIF 0013
C      SUMJ=0.000                                                GIF 0014
C      IF (ITC.EQ.0) GO TO 40                                     GIF 0015
C      IF (ITC.EQ.1) GO TO 20                                     GIF 0016
C      LINEAR F AND G IN REGIONS J:                               GIF 0017
C      DO 10 J=1,K                                              GIF 0018
C      A=F(G1,J)*G(G2,J)                                         GIF 0019
C      B=F(G1,J)*G(G2,J+1)+F(G1,J+1)*G(G2,J)                   GIF 0020
C      D=F(G1,J+1)*G(G2,J+1)                                     GIF 0021
C      SUML=0.000                                                GIF 0022
C      DO 5 I=1,N1                                              GIF 0023
C      L=I-1                                                      GIF 0024
C      M=N-L                                                       GIF 0025
C      IF (H(J).EQ.0.0.AND.L.EQ.0) GO TO 1                       GIF 0026
C      EH=H(J)**L                                                 GIF 0027
C      GO TO 2                                                    GIF 0028
1  EH=1.0                                                         GIF 0029
2  IF (Z(J).EQ.0.0.AND.M.EQ.0) GO TO 3                           GIF 0030
C      EZ=Z(J)**M                                                 GIF 0031
C      GO TO 4                                                    GIF 0032
3  EZ=1.0                                                         GIF 0033
4  SUML=SUML+FACT(N)/(FACT(M)*FACT(L))*EH*EZ*                   GIF 0034
C      X (2.DO*A/DFLJAT((L+3)*(L+2)*(L+1))+B/DFLOAT((L+3)*(L+2)) GIF 0035
C      X +D/DFLOAT(L+3))                                          GIF 0036

```

5	CONTINUE	GIF 0037
	SUMJ=SUMJ+H(J)*C(G2,J)*SUML	GIF 0038
10	CONTINUE	GIF 0039
	GO TO 100	GIF 0040
C	LINEAR F AND CONSTANT G IN REGIONS J:	GIF 0041
20	DO 30 J=1,K	GIF 0042
	SUML=0.000	GIF 0043
	DO 25 I=1,N1	GIF 0044
	L=I-1	GIF 0045
	M=N-L	GIF 0046
	IF (H(J).EQ.0.0.AND.L.EQ.0) GO TO 21	GIF 0047
	EH=H(J)**L	GIF 0048
	GO TO 22	GIF 0049
21	EH=1.0	GIF 0050
22	IF (Z(J).EQ.0.0.AND.M.EQ.0) GO TO 23	GIF 0051
	EZ=Z(J)**M	GIF 0052
	GO TO 24	GIF 0053
23	EZ=1.0	GIF 0054
24	SUML=SUML+FACT(N)/(FACT(M)*FACT(L))*EH*EZ*	GIF 0055
	X (F(G1,J)/DFLOAT((L+1)*(L+2))+F(G1,J+1)/DFLOAT(L+2))	GIF 0056
25	CONTINUE	GIF 0057
	SUMJ=SUMJ+H(J)*C(G2,J)*G(G2,J)*SUML	GIF 0058
30	CONTINUE	GIF 0059
	GO TO 100	GIF 0060
C	CONSTANT F AND G IN REGIONS J:	GIF 0061
40	DO 50 J=1,K	GIF 0062
	SUML=0.000	GIF 0063
	DO 55 I=1,N1	GIF 0064
	L=I-1	GIF 0065
	M=N-L	GIF 0066
	IF (H(J).EQ.0.0.AND.L.EQ.0) GO TO 51	GIF 0067
	EH=H(J)**L	GIF 0068
	GO TO 52	GIF 0069
51	EH=1.0	GIF 0070
52	IF (Z(J).EQ.0.0.AND.M.EQ.0) GO TO 53	GIF 0071
	EZ=Z(J)**M	GIF 0072

```
GO TO 54
53 EZ=1.0
54 SUML=SUML+FACT(N)/(FACT(M)*FACT(L))*EH*EZ*
  X (1./DFLOAT(L+1))
55 CONTINUE
  SUMJ=SUMJ+H(J)*F(G1,J)*C(G2,J)*G(G2,J)*SUML
50 CONTINUE
100 GIF=SUMJ/(X(K+1)-X(1))*N
  RETURN
  END
```

```
GIF 0073
GIF 0074
GIF 0075
GIF 0076
GIF 0077
GIF 0078
GIF 0079
GIF 0080
GIF 0081
GIF 0082
```

```
C  DOUBLE PRECISION FUNCTION FACT(N)
    COMPUTES N FACTORIAL:
    FACT=1.0D0
    IF (N.LE.1) RETURN
    DO 1 I=2,N
1  FACT=FACT*DFLOAT(I)
    RETURN
    END
```

```
FACT0001
FACT0002
FACT0003
FACT0004
FACT0005
FACT0006
FACT0007
FACT0008
```

	SUBROUTINE PRTOUT(IP)	PRT00001
C	IP = 1: PRINTS OUT THE /B5/ ARRAYS.	PRT00002
C	IP = 2: PRINTS OUT THE /B3/ MATRICES GIVEN TO POWER.	PRT00003
	IMPLICIT REAL*8 (A-H,K-Z)	PRT00004
	COMMON /B2/ KR, N	PRT00005
	COMMON /B3/ L1(26,26),L2(26,26),F1(26,26),F4(26,26),	PRT00006
X	F2(26,26),F3(26,26),T(26,26)	PRT00007
	COMMON /B5/ KAO(2,25),KA1(2,25),KA2(2,25),KBO(2,25),KB1(2,25),	PRT00008
X	KB2(2,25),LA0(2,25),LA1(2,25),LA2(2,25),SR0(1,25),	PRT00009
X	SR1(1,25),SR2(1,25),KCO(1,25),KC1(1,25),KC2(1,25),	PRT00010
X	KDO(1,25),KD1(1,25),KD2(1,25),	PRT00011
X	P(2,25),P1(2,25),Q(2,25),	PRT00012
X	Q1(2,25), R(2,25), PO(2,25), P07(2,25),PH(2,25),	PRT00013
X	PH7(2,25)	PRT00014
	COMMON /XAXIS/ HX,HR(25)	PRT00015
	INTEGER KR, G, N	PRT00016
	GO TO (1001,1002),IP	PRT00017
C	KA AND KB ARRAYS:	PRT00018
	1001 WRITE (6,10)	PRT00019
	10 FORMAT ('1 G',4X,'I',12X,'KAO(G,I)',12X,'KA1(G,I)',12X,	PRT00020
X	'KA2(G,I)',12X,'KBO(G,I)',12X,'KB1(G,I)',12X,'KB2(G,I)')	PRT00021
	DO 11 G=1,2	PRT00022
	WRITE (6,12)	PRT00023
	11 WRITE (6,15) (G,I,KAO(G,I),KA1(G,I),KA2(G,I),KBO(G,I),KB1(G,I),	PRT00024
X	KB2(G,I),I=1,KR)	PRT00025
	12 FORMAT (' ')	PRT00026
	15 FORMAT (2I5,6D20.7)	PRT00027
C	KC AND KD ARRAYS:	PRT00028
	WRITE (6,16)	PRT00029
	16 FORMAT ('1 G',4X,'I',12X,'KCO(G,I)',12X,'KC1(G,I)',12X,	PRT00030
X	'KC2(G,I)',12X,'KDO(G,I)',12X,'KD1(G,I)',12X,'KD2(G,I)')	PRT00031
	G=1	PRT00032
	WRITE (6,12)	PRT00033
	17 WRITE (6,15) (G,I,KCO(G,I),KC1(G,I),KC2(G,I),KDO(G,I),KD1(G,I),	PRT00034
X	KD2(G,I),I=1,KR)	PRT00035
C	LA AND SR ARRAYS:	PRT00036

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WRITE (6,20)
20 FORMAT ('1  G',4X,'I',12X,'LA0(G,I)',12X,'LA1(G,I)',12X,
X  'LA2(G,I)',12X,'SRO(G,I)',12X,'SR1(G,I)',12X,'SR2(G,I)')
G=1
WRITE (6,12)
WRITE (6,15) (G,I,LA0(G,I),LA1(G,I),LA2(G,I),SRO(G,I),SR1(G,I),
X  SR2(G,I),I=1,KR)
G=2
WRITE (6,12)
WRITE (6,25) (G,I,LA0(G,I),LA1(G,I),LA2(G,I),I=1,KR)
25 FORMAT (2I5,3D20.7)
C   P, Q, AND R ARRAYS:
WRITE (6,30)
30 FORMAT ('1  G',4X,'I',14X,'P(G,I)',13X,'P1(G,I)',14X,'Q(G,I)',
X  13X,'Q1(G,I)',14X,'R(G,I)')
DO 31 G=1,2
WRITE (6,12)
31 WRITE (6,35) (G,I,P(G,I),P1(G,I),Q(G,I),Q1(G,I),R(G,I),I=1,KR)
35 FORMAT (2I5,5D20.7)
C   PO, PH, AND HR ARRAYS:
WRITE (6,40)
40 FORMAT ('1  G',4X,'I',13X,'P0(G,I)',12X,'P07(G,I)',13X,'PH(G,I)',
X  12X,'PH7(G,I)',15X,'HR(I)')
DO 41 G=1,2
WRITE (6,12)
41 WRITE (6,45) (G,I,P0(G,I),P07(G,I),PH(G,I),PH7(G,I),HR(I),I=1,KR)
45 FORMAT (2I5,5D20.7)
GO TO 100
C/  PRINT OUT THE /B3/ MATRICES:
1002 WRITE (6,50)
50 FORMAT ('1MATRIX L1:',/)
DO 51 I=1,N
51 WRITE (6,55) (L1(I,J),J=1,N)
55 FORMAT (10D12.3,/, (2X,10D12.3))
WRITE (6,60)
60 FORMAT ('1MATRIX L2:',/)

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PRT00037
PRT00038
PRT00039
PRT00040
PRT00041
PRT00042
PRT00043
PRT00044
PRT00045
PRT00046
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PRT00048
PRT00049
PRT00050
PRT00051
PRT00052
PRT00053
PRT00054
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PRT00061
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PRT00064
PRT00065
PRT00066
PRT00067
PRT00068
PRT00069
PRT00070
PRT00071
PRT00072

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```
DO 61 I=1,N
61 WRITE (6,55) (L2(I,J),J=1,N)
   WRITE (6,70)
70 FORMAT ('MATRIX F1:',/)
   DO 71 I=1,N
71 WRITE (6,55) (F1(I,J),J=1,N)
   WRITE (6,80)
80 FORMAT ('MATRIX F2:',/)
   DO 81 I=1,N
81 WRITE (6,55) (F2(I,J),J=1,N)
   WRITE (6,82)
82 FORMAT ('MATRIX F3:',/)
   DO 83 I=1,N
83 WRITE (6,55) (F3(I,J),J=1,N)
   WRITE (6,84)
84 FORMAT ('MARTIX F4:',/)
   DO 85 I=1,N
85 WRITE (6,55) (F4(I,J),J=1,N)
   WRITE (6,90)
90 FORMAT ('MATRIX T:',/)
   DO 91 I=1,N
91 WRITE (6,55) (T(I,J),J=1,N)
100 RETURN
END
```

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PRT00073
PRT00074
PRT00075
PRT00076
PRT00077
PRT00078
PRT00079
PRT00080
PRT00081
PRT00082
PRT00083
PRT00084
PRT00085
PRT00086
PRT00087
PRT00088
PRT00089
PRT00090
PRT00091
PRT00092
PRT00093
PRT00094
PRT00095
PRT00096
```

	SUBROUTINE POWER	POWE0001
C	SOLVES THE 2*N MULTIGROUP EQUATIONS: $M*PHI = (1/LAMDA)*F*PHI$	POWE0002
C	BY THE FISSION SOURCE POWER METHOD	POWE0003
C	USING SIMULTANEOUS OVERRELAXATION.	POWE0004
C	WHERE: M AND F ARE DOUBLE PRECISION 2N BY 2N BLOCK MATRICES;	POWE0005
C	AND: PHI IS THE 2N FLUX (FAST AND THERMAL) VECTOR.	POWE0006
C	$L1*PHI1 = CHI1*(F1*PHI1 + F2*PHI2)$	POWE0007
C	$-T*PHI1 + L2*PHI2 = CHI2*(F3*PHI1 + F4*PHI2)$	POWE0008
C	METHOD FOLLOWS WACHPRESS, PAGE 83. SOLUTION BY GROUP ITERATION.	POWE0009
	IMPLICIT REAL*8 (A-H,L-Z)	POWE0010
	COMMON /B1/ IBC, IPLOT, JPLOT, IPUNCH, ISEE	POWE0011
	COMMON /B2/ KR, N	POWE0012
	COMMON /B3/ L1(26,26), L2(26,26), F1(26,26), F4(26,26),	POWE0013
X	F2(26,26), F3(26,26), T(26,26)	POWE0014
	COMMON /B4/ PHI(2,26), PSI(2,26), LAMDA, ICOUT	POWE0015
	COMMON /B5/ S(26), ERROR(2,26), Z(26)	POWE0016
	COMMON /B6/ TE1(2,5), TE2(2,5), TE3(5), IN(5)	POWE0017
	COMMON /CHIF/ CHI(2)	POWE0018
	COMMON /XAXIS/ HX, HR(25)	POWE0019
	COMMON /ER/ EPS1, EPS2, EPS3	POWE0020
	COMMON /FSTR/ PHISTR(2,26,6)	POWE0021
	COMMON /ESTR/ LAMSTR(300), EFSTR(2,300), EFMSTR(2,300), ERLAM(300)	POWE0022
	COMMON /TRUE/ TRULAM, TRUPHI(2,26), PHICON(2,300), LAMCON(300), IFT	POWE0023
	DIMENSION PSI1(26), PSI2(26), SQ(2), DPHI(2), ERRMAX(2)	POWE0024
	INTEGER N	POWE0025
C	DEFAULT OPTIONS FOR THE TRUE EIGENVALUE AND FLUXES:	POWE0026
	TRULAM=1.0	POWE0027
	DO 5 IG=1,2	POWE0028
	DO 5 I=1,N	POWE0029
5	TRUPHI(IG,I)=1.0	POWE0030
C	DEFAULT OPTIONS FOR POWER PARAMETERS:	POWE0031
	ALPHA=1.25	POWE0032
	LAMDA=1.0	POWE0033
	DO 555 IG=1,2	POWE0034
	IF (IBC.NE.4) GO TO 551	POWE0035
	DO 550 I=1,N	POWE0036

550	PHI(IG,I)=1.0	POWE0037
	GO TO 555	POWE0038
551	X=3.1415926/HX	POWE0039
	IF (IBC.NE.1) X=X/2.0	POWE0040
	SUM1=0.0	POWE0041
	DO 552 K=1,KR	POWE0042
	SUM1=SUM1+HR(K)	POWE0043
552	PHI(IG,K)=DSIN(SUM1*X)	POWE0044
555	CONTINUE	POWE0045
C	READ IN THE TRUE (EXPECTED) EIGENVALUE AND FLUX VECTOR (MINUS 0 BC'S):	POWE0046
	IFT=0	POWE0047
	READ (5,500,END=501) TRULAM, (TRUPHI(1,I),I=1,N)	POWE0048
	READ (5,503,END=501) (TRUPHI(2,I),I=1,N)	POWE0049
	IFT=1	POWE0050
500	FORMAT (E25.14,/, (4E20.10))	POWE0051
C	READ IN: OVERRELAXATION PARAMETERS ; ALPHA (OUTER ITERATION)	POWE0052
C	INITIAL GUESS AT EIGENVALUE; LAMDA	POWE0053
C	INITIAL NORMALIZED FLUX ; PHI(1-N)	POWE0054
501	READ (5,506,END=510) ALPHA	POWE0055
	READ (5,502,END=510) LAMDA	POWE0056
	READ (5,503) (PHI(1,I),I=1,N)	POWE0057
	READ (5,503) (PHI(2,I),I=1,N)	POWE0058
506	FORMAT (F10.5)	POWE0059
502	FORMAT (E25.14)	POWE0060
503	FORMAT ((4E20.10))	POWE0061
510	CONTINUE	POWE0062
C	STORING FOR PRINTING THE MULTIGRUP FLUX SHAPE.	POWE0063
	DO 11 IG=1,2	POWE0064
	DO 10 I=1,N	POWE0065
10	PHISTR(IG,I,2)=PHI(IG,I)	POWE0066
C	FILL RUNNING COORD IN PHISTR	POWE0067
	KR1=KR+1	POWE0068
	DO 11 I=1,KR1	POWE0069
11	PHISTR(IG,I,1)=DFLOAT(I)	POWE0070
C	IK IS THE FLUX PLOTTING COUNTER.	POWE0071
	IK=1	POWE0072

C	STORES THE ITERATION NUMBER FOR FLUX HISTORY PLOTTING:	POWE0073
	IN(1)=0	POWE0074
C	STORES TEMPORARY ERRORS FOR FLUX HISTORY PLOTTING:	POWE0075
	TE1(1,1)=0.	POWE0076
	TE1(2,1)=0.	POWE0077
	TE2(1,1)=0.	POWE0078
	TE2(2,1)=0.	POWE0079
	TE3(1)=0.0	POWE0080
C	EIGENVALUE OF THE PREVIOUS ITERATION:	POWE0081
	LAMB4=LAMDA	POWE0082
C	THE MAXIMUM NUMBER OF ALLOWED ITERATIONS: ICMAX	POWE0083
	ICMAX=300	POWE0084
C	PRINT OUT THE POWER METHOD PARAMETER INFORMATION:	POWE0085
	WRITE (6,700) ICMAX,ALPHA,LAMDA,(PHI(1,I),I=1,N)	POWE0086
	WRITE (6,701) (PHI(2,I),I=1,N)	POWE0087
	700 FORMAT ('1EXECUTING MULTIGROUP FISSION SOURCE POWER ITERATION METH	POWE0088
	XOD.',///,	POWE0089
	X 5X,'MAXIMUM NUMBER OF ALLOWABLE ITERATIONS:',/,	POWE0090
	X 10X,'ICMAX =',I4,///,	POWE0091
	X 5X,'OUTER ITERATION RELAXATION PARAMETER:',/,	POWE0092
	X 10X,'ALPHA =',F7.3,///,	POWE0093
	X 5X,'INITIAL GUESS AT EIGENVALUE:',/,	POWE0094
	X 10X,'LAMBDA =',E22.14,///,	POWE0095
	X 5X,'INITIAL GUESS AT THE GROUP FLUX SHAPE CONNECTION POINTS:',	POWE0096
	X //,8X,'FAST GROUP:',/,	POWE0097
	X 10X,'F(K)''S =',4E25.14,/, (18X,4E25.14))	POWE0098
	701 FORMAT ('0',7X,'THERMAL GROUP:',/,	POWE0099
	X 10X,'F(K)''S =',4E25.14,/, (18X,4E25.14))	POWE0100
C	BEGIN ITERATION LOOP.	POWE0101
	ICOUT=0	POWE0102
C	ICOUT IS THE OUTER ITERATION COUNTER.	POWE0103
	20 ICOUT=ICOUT+1	POWE0104
	IF (ICOUT.GT.ICMAX) GO TO 100	POWE0105
C	SOLVE FOR THE NEW GROUP FLUX VECTORS: PSI:	POWE0106
C	FAST GROUP; SOURCE VECTOR:	POWE0107
	DO 25 I=1,N	POWE0108

	S(I)=0.0	POWE0109
	DO 24 J=1,N	POWE0110
	24 S(I)=S(I)+F1(I,J)*PHI(1,J)+F2(I,J)*PHI(2,J)	POWE0111
	25 S(I)=CHI(1)*S(I)	POWE0112
C	FAST FLUX:	POWE0113
	CALL SOLV3D(N,L1,PSI1,S)	POWE0114
C	THERMAL GROUP; SOURCE VECTOR:	POWE0115
	DO 27 I=1,N	POWE0116
	S(I)=0.0	POWE0117
	Z(I)=0.0	POWE0118
	DO 26 J=1,N	POWE0119
	S(I)=S(I)+F3(I,J)*PSI1(J)+F4(I,J)*PHI(2,J)	POWE0120
	26 Z(I)=Z(I)+T(I,J)*PSI1(J)	POWE0121
	27 Z(I)=Z(I)+CHI(2)*S(I)	POWE0122
C	THERMAL FLUX:	POWE0123
	CALL SOLV3D(N,L2,PSI2,Z)	POWE0124
C	CALCULATION OF THE EIGENVALUE:	POWE0125
	SUM1=0.000	POWE0126
	SUM2=0.000	POWE0127
	DO 28 I=1,N	POWE0128
	SUM2=SUM2+PSI1(I)*PSI1(I)+PSI2(I)*PSI2(I)	POWE0129
	28 SUM1=SUM1+PSI1(I)*PHI(1,I)+PSI2(I)*PHI(2,I)	POWE0130
	LAMDA=SUM2/SUM1	POWE0131
	LAMSTR(ICOUT)=LAMDA	POWE0132
	ERRLAM=DABS(LAMDA-LAMB4)	POWE0133
C	PUT PSI1 AND PSI2 INTO BIGGER PSI:	POWE0134
	DO 30 I=1,N	POWE0135
	PSI(1,I)=PSI1(I)	POWE0136
	30 PSI(2,I)=PSI2(I)	POWE0137
C	POINT BY POINT SIMULTANEOUS RELAXATION FLUX ITERATION:	POWE0138
	X=ALPHA	POWE0139
C	DO NOT RELAX DURING THE FIRST THREE ITERATIONS:	POWE0140
	IF (ICOUT.LE.3) X=1.0	POWE0141
C	CALCULATE THE NEW GROUP FLUX ITERATES AND GROUP ERRORS:	POWE0142
	DO 40 IG=1,2	POWE0143
	DO 40 I=1,N	POWE0144

40	PSI(IG,I)=PHI(IG,I)+X*(PSI(IG,I)/LAMDA-PHI(IG,I))	POWE0145
	CALL NORMAL(PSI,N)	POWE0146
	DO 39 IG=1,2	POWE0147
	ERRMAX(IG)=0.0	POWE0148
	SQ(IG)=0.0	POWE0149
	DO 38 I=1,N	POWE0150
	ERROR(IG,I)=DABS((PSI(IG,I)-PHI(IG,I))/PSI(IG,I))	POWE0151
	IF (ERROR(IG,I).GT.ERRMAX(IG)) ERRMAX(IG)=ERROR(IG,I)	POWE0152
	SQ(IG)=SQ(IG)+ERROR(IG,I)**2	POWE0153
C	UPDATE THE FLUX ITERATE:	POWE0154
38	PHI(IG,I)=PSI(IG,I)	POWE0155
39	SQ(IG)=DSQRT(SQ(IG))	POWE0156
C	NORMALIZE PSI GROUPS TO UNITY:	POWE0157
	CALL NORM2(PSI,TRUPHI,N)	POWE0158
	IF (IFT.EQ.0) GO TO 37	POWE0159
	DLAM=LAMDA-TRULAM	POWE0160
	DO 36 IG=1,2	POWE0161
	DPHI(IG)=0.0	POWE0162
	DO 35 I=1,N	POWE0163
35	DPHI(IG)=DPHI(IG)+(PSI(IG,I)-TRUPHI(IG,I))*2	POWE0164
36	DPHI(IG)=DSQRT(DPHI(IG))	POWE0165
37	IF (IPLT.NE.2) GO TO 45	POWE0166
C	THE FOLLOWING IS FOR NICELY PLOTTING THE GROUP FLUX HISTORY.	POWE0167
	DO 41 IG=1,2	POWE0168
	DO 41 I=1,N	POWE0169
41	ERROR(IG,I)=PSI(IG,I)	POWE0170
C	ERROR NOW CONTAINS THE NEW NORMALIZED FLUX ITERATE PHI.	POWE0171
	JK=IK	POWE0172
	IF (IK.EQ.0) JK=5	POWE0173
	DO 42 IG=1,2	POWE0174
	DO 42 I=1,N	POWE0175
	IF (DABS(ERROR(IG,I)-PHISTR(IG,I,JK+1)).GE.0.01) GO TO 43	POWE0176
42	CONTINUE	POWE0177
C	FLUX HAS NOT CHANGED ENOUGH FOR PLOTTING.	POWE0178
	GO TO 45	POWE0179
C	SAVE THE NORMALIZED FLUX FOR PLOTTING:	POWE0180

43	IK=IK+1	POWE0181
	IN(IK)=ICOUT	POWE0182
	TE3(IK)=ERRLAM	POWE0183
	DO 44 IG=1,2	POWE0184
	TE1(IG,IK)=ERRMAX(IG)	POWE0185
	TE2(IG,IK)=SQ(IG)	POWE0186
	DO 44 I=1,N	POWE0187
44	PHISTR(IG,I,IK+1)=ERROR(IG,I)	POWE0188
	IF (IK.NE.5) GO TO 45	POWE0189
C	PLOT THE LAST FIVE SAVED FLUXES:	POWE0190
	CALL PHIPLT(5)	POWE0191
	IK=0	POWE0192
45	CONTINUE	POWE0193
C	ERROR CRITERIA FOR ACCEPTANCE OF CONVERGENCE.	POWE0194
	IFLAG1=0	POWE0195
	IFLAG2=0	POWE0196
	IFLAG3=0	POWE0197
C	STORE THE ERRORS FOR COMPARISON:	POWE0198
C	ERROR BETWEEN ITERATION EIGENVALUES:	POWE0199
	ERLAM(ICOUT)=ERRLAM	POWE0200
	DO 46 IG=1,2	POWE0201
C	MAXIMUM ERROR BETWEEN ITERATION FLUXES:	POWE0202
	EFSTR(IG,ICOUT)=ERRMAX(IG)	POWE0203
C	MEAN SQUARE ERROR BETWEEN ITERATION FLUXES:	POWE0204
	EFMSTR(IG,ICOUT)=SQ(IG)	POWE0205
C	MEAN SQUARE ERROR BETWEEN THE ITERATION FLUX AND GIVEN TRUE FLUX:	POWE0206
	PHICON(IG,ICOUT)=DPHI(IG)	POWE0207
46	CONTINUE	POWE0208
C	ERROR BETWEEN THE ITERATION EIGENVALUE AND GIVEN TRUE EIGENVALUE:	POWE0209
	LAMCON(ICOUT)=DLAM	POWE0210
	IF ((ERRMAX(1).LT.EPS1).AND.(ERRMAX(2).LT.EPS1))	IFLAG1=1
	IF ((SQ(1).LT.EPS2).AND.(SQ(2).LT.EPS2))	IFLAG2=1
	IF (ERRLAM.LT.EPS3) IFLAG3=1	POWE0213
	IFLAG4=IFLAG1*IFLAG2*IFLAG3	POWE0214
	IF (IFLAG4.EQ.1) GO TO 50	POWE0215
C	OTHERWISE CONTINUE THE ITERATION.	POWE0216

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LAMB4=LAMDA
GO TO 20
50 CONTINUE
C     CONVERGENCE ACCOMPLISHED.
C     NORMALIZE THE CONVERGED FLUX VECTOR:
CALL NORMAL(PHI,N)
C     PLOT ANY LEFT OVER FLUX HISTORY PLOTS:
IF ((IPLJT.EQ.2).AND.(IK.NE.0)) CALL PHIPLT(IK)
C     BOUNDRY CONDITION INSERTIONS.
IER=0
C     IER ALLOWS B.C. INSERTIONS FOR YES AND NO CONVERGENCE:
55 IF (IBC.EQ.4) GO TO 90
   IF (IBC.NE.3) GO TO 60
   PHI(1,KR+1)=0.
   PHI(2,KR+1)=0.
   GO TO 90
60 DO 70 I=1,N
   J=N+1-I
   PHI(1,J+1)=PHI(1,J)
70 PHI(2,J+1)=PHI(2,J)
   IF (IBC.EQ.5.OR.IBC.EQ.7) GO TO 71
   PHI(1,1)=0.0
   PHI(2,1)=0.0
   GO TO 72
71 PHI(1,1)=PHI(1,2)
   PHI(2,1)=PHI(2,2)
72 IF (IBC.NE.1) GO TO 73
   PHI(1,KR+1)=0.0
   PHI(2,KR+1)=0.0
   GO TO 90
73 IF (IBC.LT.6) GO TO 90
   PHI(1,KR+1)=PHI(1,KR)
   PHI(2,KR+1)=PHI(2,KR)
90 IF (IER.EQ.1) GO TO 102
   RETURN
C     NO CONVERGENCE ACCOMPLISHED:

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POWE0217
POWE0218
POWE0219
POWE0220
POWE0221
POWE0222
POWE0223
POWE0224
POWE0225
POWE0226
POWE0227
POWE0228
POWE0229
POWE0230
POWE0231
POWE0232
POWE0233
POWE0234
POWE0235
POWE0236
POWE0237
POWE0238
POWE0239
POWE0240
POWE0241
POWE0242
POWE0243
POWE0244
POWE0245
POWE0246
POWE0247
POWE0248
POWE0249
POWE0250
POWE0251
POWE0252

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100 CONTINUE	POWE0253
C     NORMALIZE THE UNCONVERGED FLUX:	POWE0254
CALL NORMAL(PHI,N)	POWE0255
ICOUT=ICOUT-1	POWE0256
WRITE (6,101) ICOUT	POWE0257
101 FORMAT (1H1,'POWER METHOD DID NOT CONVERGE FOR THIS CASE AFTER',	POWE0258
X 14,' ITERATIONS.',//,1X,'EXECUTION TERMINATED ')	POWE0259
IER=1	POWE0260
GO TO 55	POWE0261
102 CONTINUE	POWE0262
C     FOR PRINTING OUT THE EIGENVALUE HISTORY AND THE FINAL FLUX SHAPE:	POWE0263
IPLOT=1	POWE0264
JPLOT=1	POWE0265
RETURN	POWE0266
END	POWE0267

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SUBROUTINE SOLV3D(N,A,X,Y)
C SOLVES THE N DOUBLE PRECISION MATRIX EQUATIONS:  A*X = Y,
C FOR X - GIVEN THE N BY N TRIDIAGONAL MATRIX A
C AND THE SOURCE VECTOR Y.
C METHOD IS FORWARD ELIMINATION FOLLOWED BY BACKWARD SUBSTITUTION.
C CF - WACHPRESS, PAGE 23.
REAL*8 A, X, Y, H, P, D
DIMENSION A(26,26), X(26), Y(26), H(26), P(26)
IF (A(1,1).EQ.0.0) GO TO 10
H(1)=-A(1,2)/A(1,1)
P(1)=Y(1)/A(1,1)
DO 1 M=2,N
D=A(M,M)+A(M,M-1)*H(M-1)
IF (D.EQ.0.0) GO TO 20
P(M)=(Y(M)-A(M,M-1)*P(M-1))/D
IF (M.EQ.N) GO TO 1
H(M)=-A(M,M+1)/D
1 CONTINUE
X(N)=P(N)
DO 2 I=2,N
M=N+1-I
2 X(M)=P(M)+H(M)*X(M+1)
RETURN
C IN CASE OF ANY IMPENDING ZERO DIVISORS:
10 WRITE (6,11)
11 FORMAT ('FIRST ELEMENT OF A, A(1,1), IS ZERO.',/,
X 5X,'BETTER FIX IT BOSS.')
GO TO 30
20 WRITE (6,21) M
21 FORMAT ('ZERO DIVISOR ENCOUNTERED IN EQUATION M =',I3,'.',/,
X 5X,'BETTER FIX IT BOSS.')
30 WRITE (6,31)
31 FORMAT ('EXECUTION TERMINATED.')
CALL EXIT
RETURN
END

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SOLV0001
SOLV0002
SOLV0003
SOLV0004
SOLV0005
SOLV0006
SOLV0007
SOLV0008
SOLV0009
SOLV0010
SOLV0011
SOLV0012
SOLV0013
SOLV0014
SOLV0015
SOLV0016
SOLV0017
SOLV0018
SOLV0019
SOLV0020
SOLV0021
SOLV0022
SOLV0023
SOLV0024
SOLV0025
SOLV0026
SOLV0027
SOLV0028
SOLV0029
SOLV0030
SOLV0031
SOLV0032
SOLV0033
SOLV0034
SOLV0035
SOLV0036

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C  SUBROUTINE NORMAL(PHI,N)
   NORMALIZES THE GROUP FLUXES TO ONE. NOT BOTH GROUPS.
   REAL*8 PHI(2,26), A
   A=DABS(PHI(1,1))
   DO 1 IG=1,2
   DO 1 I=1,N
   IF (DABS(PHI(IG,I)).GT.A) A=DABS(PHI(IG,I))
1  CONTINUE
   DO 2 IG=1,2
   DO 2 I=1,N
2  PHI(IG,I)=PHI(IG,I)/A
   RETURN
   END
```

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NORL0001
NORL0002
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NORL0010
NORL0011
NORL0012
NORL0013
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SUBROUTINE PHIPLT(L)
C   PLOTS THE GROUP FLUX HISTORY, WITH UP TO 5 GROUP FLUXES PER PLOT.
C   FAST AND THERMAL GROUP FLUXES ARE PLOTTED SEPERATELY.
C   L IS THE NUMBER OF FLUXES TO BE PLOTTED.
C   L IS BETWEEN 1 AND 5.
IMPLICIT REAL*8 (A-H,O-Z)
COMMON /B1/ IBC
COMMON /B2/ KR,N
COMMON /B5/ S(26), A(26,6), B(26,6)
COMMON /B6/ TE1(2,5),TE2(2,5),TE3(5),IN(5)
COMMON /ER/ EPS1, EPS2, EPS3
COMMON /FSTR/ PHISTR(2,26,6)
DIMENSION SYMBOL(5)
INTEGER SYMBOL /'.' , '-' , '+' , '#' , '*' /
KR1=KR+1
C   SET UP B.C. CONDITIONS
IF (IBC.EQ.4) GO TO 5
IF (IBC.EQ.3) GO TO 3
DO 2 IG=1,2
DO 2 K=1,L
DO 1 I=1,N
J=N+1-I
1 PHISTR(IG,J+1,K+1)=PHISTR(IG,J,K+1)
2 PHISTR(IG,1,K+1)=0.
3 IF (IBC.EQ.2) GO TO 5
DO 4 IG=1,2
DO 4 K=1,L
4 PHISTR(IG,KR1,K+1)=0.
5 CONTINUE
C   FLUXES IN PHISTR HAVE BEEN NORMALIZED IN POWER.
C   PUT THE FAST FLUX IN A, AND THE THERMAL FLUX IN B:
L1=L+1
DO 10 K=1,L1
DO 10 I=1,KR1
A(I,K)=PHISTR(1,I,K)
10 B(I,K)=PHISTR(2,I,K)

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PHIP0001
PHIP0002
PHIP0003
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PHIP0014
PHIP0015
PHIP0016
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PHIP0019
PHIP0020
PHIP0021
PHIP0022
PHIP0023
PHIP0024
PHIP0025
PHIP0026
PHIP0027
PHIP0028
PHIP0029
PHIP0030
PHIP0031
PHIP0032
PHIP0033
PHIP0034
PHIP0035
PHIP0036

```

C	PLOT THE L FAST FLUX SHAPES ON ONE GRAPH:	PHIP0037
	CALL PRTPLT(0,A,KR1,L1,KR1,0,26,6,2)	PHIP0038
	WRITE (6,20)	PHIP0039
20	FORMAT (/,'OFAST FLUX ITERATION HISTORY PLOT.',/)	PHIP0040
	WRITE (6,30)	PHIP0041
30	FORMAT (	PHIP0042
	X 'OKEY:',5X,'SYMBOL',5X,'ITERATION NUMBER:',7X,'ERRCR CRITERIA',	PHIP0043
	X 11X,'ERROR',13X,'TOLERANCE')	PHIP0044
	DO 35 I=1,L	PHIP0045
35	WRITE (6,40) SYMBOL(I),IN(I),TE1(1,I),EPS1,TE2(1,I),EPS2,	PHIP0046
	X TE3(I),EPS3	PHIP0047
40	FORMAT (/ ,12X,A1,15X,I3,16X,'FLUX',14X,1PD15.5,5X,1PD15.5,/,	PHIP0048
	X 47X,'MEAN SQ. FLUX',5X,1PD15.5,5X,1PD15.5,/,	PHIP0049
	X 47X,'EIGENVALUE',8X,1PD15.5,5X,1PD15.5)	PHIP0050
C	PLOT THE L THERMAL FLUX SHAPES ON THE OTHER GRAPH:	PHIP0051
	CALL PRTPLT(0,B,KR1,L1,KR1,0,26,6,2)	PHIP0052
	WRITE (6,50)	PHIP0053
50	FORMAT (/,'OTHERMAL FLUX ITERATION PLOT.',/)	PHIP0054
	WRITE (6,30)	PHIP0055
	DO 55 I=1,L	PHIP0056
55	WRITE (6,40) SYMBOL(I),IN(I),TE1(2,I),EPS1,TE2(2,I),EPS2,	PHIP0057
	X TE3(I),EPS3	PHIP0058
	RETURN	PHIP0059
	END	PHIP0060

	SUBROUTINE OUTPUT	OUTP0001
C	PRINTS THE RESULTS OF THE METHOD.	OUTP0002
	IMPLICIT REAL*8 (A-H,L-Z)	OUTP0003
	COMMON /B1/ IBC,IPL0T,JPL0T,IPUNCH	OUTP0004
	COMMON /B2/ KR,N	OUTP0005
	COMMON /B4/ PHI(2,26),PSI(2,26),LAMDA,IC0UT	OUTP0006
	COMMON /ER/ EPS1,EPS2,EPS3	OUTP0007
	COMMON /ESTR/ LAMSTR(300),EFSTR(2,300),EFMSTR(2,300),ERLAM(300)	OUTP0008
	COMMON /TRUE/ TRULAM,TRUPHI(2,26),PHIC0N(2,300),LAMCON(300),IFT	OUTP0009
	INTEGER N	OUTP0010
	KR0=KR-1	OUTP0011
	KR1=KR+1	OUTP0012
	WRITE (6,1)	OUTP0013
	1 FORMAT ('RESULTS OF THE MULTIGROUP METHOD:')	OUTP0014
	WRITE (6,10) IC0UT	OUTP0015
	10 FORMAT (//,' PROBLEM TERMINATED AFTER',I5,	OUTP0016
	X ' OUTER (POWER) ITERATIONS TO:')	OUTP0017
	WRITE (6,20) LAMDA	OUTP0018
	20 FORMAT (/,10X,'LAMDA = ',1PE21.14)	OUTP0019
C	PRINT OUT EIGENVALUES.	OUTP0020
	CALL PLOT	OUTP0021
	WRITE (6,30)	OUTP0022
	30 FORMAT ('RESULTS AFTER PROBLEM TERMINATION:',/,	OUTP0023
	X 'ONUMBER',5X,'THERMAL FLUX POINTS',5X,'FAST FLUX POINTS')	OUTP0024
	WRITE (6,50) (K,PHI(2,K),PHI(1,K),K=1,KR1)	OUTP0025
	50 FORMAT (I5,1PE26.7,1PE21.7)	OUTP0026
	IF (IPUNCH.EQ.1) CALL PUNCH	OUTP0027
C	CALCULATE THE FINAL TO EXPECTED FLUX RATIOS:	OUTP0028
C	NORMALIZE BOTH PHI GROUP FLUXES FOR TRUPHI COMPARISON:	OUTP0029
	CALL NORM2(PHI,TRUPHI,KR1)	OUTP0030
	K1=1	OUTP0031
	K2=KR1	OUTP0032
	IF (IBC.LE.2) K1=2	OUTP0033
	IF ((IBC.EQ.1).OR.(IBC.EQ.3)) K2=KR	OUTP0034
	DO 60 IG=1,2	OUTP0035
	IF (IBC.LE.2) PSI(IG,1)=1.0	OUTP0036

IF ((IBC.EQ.1).OR.(IBC.EQ.3)) PSI(IG,KR1) = 1.0	OUTP0037
I=0	OUTP0038
DO 60 K=K1,K2	OUTP0039
I=I+1	OUTP0040
60 PSI(IG,K)=PHI(IG,K)/TRUPHI(IG,I)	OUTP0041
WRITE (6,70) (I,PSI(2,I),PSI(1,I),I=1,KR1)	OUTP0042
70 FORMAT ('RATIOS OF THE TERMINATED GROUP FLUX TO THE EXPECTED GROU	OUTP0043
XP FLUX:',//,	OUTP0044
X 10X,'- AN INDICATION OF THE ACCURACY OF THE CONVERGENCE -',///,	OUTP0045
X ' K',12X,'THERMAL RATIO',15X,'FAST RATIO',//,(I5,2E25.10))	OUTP0046
C PRINT OUT THE STORED ITERATION ERRORS:	OUTP0047
WRITE (6,110) EPS1,(EFSTR(2,I),I=1,ICOUT)	OUTP0048
WRITE (6,111) EPS1,(EFSTR(1,I),I=1,ICOUT)	OUTP0049
WRITE (6,112) EPS2,(EFMSTR(2,I),I=1,ICOUT)	OUTP0050
WRITE (6,113) EPS3,(EFMSTR(1,I),I=1,ICOUT)	OUTP0051
WRITE (6,114) EPS3,(ERLAM(I),I=1,ICOUT)	OUTP0052
110 FORMAT ('MAXIMUM NORMALIZED ERRORS BETWEEN THE THERMAL FLUX ITERA	OUTP0053
XTIONS:',	OUTP0054
X 25X,'TOLERANCE USED = ',1PE12.4,//, (1P5E20.5))	OUTP0055
111 FORMAT ('MAXIMUM NORMALIZED ERRORS BETWEEN THE FAST FLUX ITERATIO	OUTP0056
XNS:',	OUTP0057
X 25X,'TOLERANCE USED = ',1PE12.4,//, (1P5E20.5))	OUTP0058
112 FORMAT ('MEAN SQUARE NORMALIZED ERROR BETWEEN THE THERMAL FLUX IT	OUTP0059
ERATIONS:',	OUTP0060
X 18X,'TOLERANCE USED = ',1PE12.4,//, (1P5E20.5))	OUTP0061
113 FORMAT ('MEAN SQUARE NORMALIZED ERROR BETWEEN THE FAST FLUX ITERA	OUTP0062
XTIONS:',	OUTP0063
X 18X,'TOLERANCE USED = ',1PE12.4,//, (1P5E20.5))	OUTP0064
114 FORMAT ('ERROR BETWEEN THE ITERATION EIGENVALUES:',	OUTP0065
X 28X,'TOLERANCE USED = ',1PE12.4,//, (1P5E20.5))	OUTP0066
C PRINT OUT THE GIVEN TRUE EIGENVALUE AND FLUX:	OUTP0067
IF (IFT.EQ.0) RETURN	OUTP0068
WRITE (6,115) TRULAM,((TRUPHI(3-J,I),J=1,2),I=1,N)	OUTP0069
115 FORMAT ('THE GIVEN TRUE EIGENVALUE:',//,15X,	OUTP0070
X 'TRULAM =',E22.14,///,	OUTP0071
X 'OTHE GIVEN MULTIGROUP FLUXES:',//,	OUTP0072

<pre> X 13X, 'THERMAL', 16X, 'FAST', //, (2D20.10) C PRINT OUT THE STORED CONVERGENCE ERRORS:   WRITE (6,120) (PHICON(2,I), I=1, ICOUT)   WRITE (6,121) (PHICON(1,I), I=1, ICOUT)   WRITE (6,122) (LAMCON(I), I=1, ICOUT) 120 FORMAT ('1MEAN SQUARE ERROR BETWEEN THE THERMAL ITERATION FLUX AND X THE GIVEN TRUE THERMAL FLUX:', //, (1P5E20.5)) 121 FORMAT ('1MEAN SQUARE ERROR BETWEEN THE FAST ITERATION FLUX AND TH XE GIVEN TRUE FAST FLUX:', //, (1P5E20.5)) 122 FORMAT ('1ERROR BETWEEN THE ITERATION EIGENVALUES AND THE GIVEN TR XUE EIGENVALUE:', //, (1P5E20.5))   RETURN   END </pre>	<pre> OUTPUT0073 OUTPUT0074 OUTPUT0075 OUTPUT0076 OUTPUT0077 OUTPUT0078 OUTPUT0079 OUTPUT0080 OUTPUT0081 OUTPUT0082 OUTPUT0083 OUTPUT0084 OUTPUT0085 </pre>
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<pre> SUBROUTINE PLOT C   PLOTS OUT THE EIGENVALUE HISTORY AS A TABLE AND A GRAPH, C   AS WELL AS PLOTTING OUT THE FINAL MULTIGROUP FLUX SHAPES. IMPLICIT REAL*8 (A-H,L-Z) COMMON /B1/ IBC,IPLT,JPLT,IPUN COMMON /B2/ KR COMMON /B4/ PHI(2,26), PSI(2,26), LAMDA, ICUT COMMON /B5/ B(300,2) COMMON /ESTR/ LAMSTR(300) DIMENSION C(26,3) C   IN ORDER TO SAVE SOME SPACE: EQUIVALENCE (B(1),C(1)) WRITE (6,1) (LAMSTR(I),I=1,ICUT) 1 FORMAT ('TABLE OF EIGENVALUES DURING THE POWER ITERATION:', X      '//,(1P5E25.14)) IF (JPLT.EQ.0) GO TO 20 DO 10 I=1,ICUT B(I,1)=I 10 B(I,2)=LAMSTR(I) CALL PRTPLT(1,B,ICUT,2,ICUT,0,300,2,1) WRITE (6,11) 11 FORMAT ('PLOT OF THE EIGENVALUE HISTORY THROUGH THE ITERATIONS.') 20 IF (IPLT.EQ.0) RETURN KR1=KR+1 DO 30 I=1,KR1 C(I,1)=I C(I,2)=PHI(1,I) 30 C(I,3)=PHI(2,I) CALL PRTPLT(2,C,KR1,3,KR1,0,26,3,2) WRITE (6,31) 31 FORMAT ('FINAL CONVERGED CONNECTING FLUX POINTS; F(K).',//, X 5X,'FAST FLUX:      .',/,5X,'THERMAL FLUX:  -') RETURN END </pre>	<pre> PLOT0001 PLOT0002 PLOT0003 PLOT0004 PLOT0005 PLOT0006 PLOT0007 PLOT0008 PLOT0009 PLOT0010 PLOT0011 PLOT0012 PLOT0013 PLOT0014 PLOT0015 PLOT0016 PLOT0017 PLOT0018 PLOT0019 PLOT0020 PLOT0021 PLOT0022 PLOT0023 PLOT0024 PLOT0025 PLOT0026 PLOT0027 PLOT0028 PLOT0029 PLOT0030 PLOT0031 PLOT0032 PLOT0033 PLOT0034 </pre>
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C  SUBROUTINE PUNCH
    PUNCHES OUT INPUT AND OUTPUT DATA.
    COMMON /B2/ KR
    COMMON /B4/ F(2,26)
    REAL*8 F
    KR1=KR+1
    WRITE (7,1) KR, (F(1,I),F(2,I),I=1,KR1)
  1  FORMAT (I5,/, (2E20.7))
    WRITE (6,100)
100 FORMAT (///, ' THE OUTPUT HAS BEEN PUNCHED OUT ONTO CARDS ')
    RETURN
    END
```

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PNCH0001
PNCH0002
PNCH0003
PNCH0004
PNCH0005
PNCH0006
PNCH0007
PNCH0008
PNCH0009
PNCH0010
PNCH0011
PNCH0012
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C SUBROUTINE NORM2(PHI,TRUPHI,N)  
C NORMALIZES BOTH ENERGY GROUPS OF PHI TO 1.0.  
C DITTO FOR TRUPHI ON THE FIRST CALL.  
REAL\*8 PHI(2,26), TRUPHI(2,26), A(2)  
DATA K /0/  
K=K+1

C DO 1 IG=1,2  
A(IG)=DABS(PHI(IG,1))  
DO 1 I=1,N  
IF (DABS(PHI(IG,I)).GT.A(IG)) A(IG)=DABS(PHI(IG,I))  
1 CONTINUE  
DO 2 IG=1,2  
DO 2 I=1,N  
2 PHI(IG,I)=PHI(IG,I)/A(IG)  
IF (K.NE.1) RETURN  
DO 5 IG=1,2  
A(IG)=0.  
DO 5 I=1,N  
IF (TRUPHI(IG,I).GT.A(IG)) A(IG)=TRUPHI(IG,I)  
5 CONTINUE  
DO 6 IG=1,2  
DO 6 I=1,N  
6 TRUPHI(IG,I)=TRUPHI(IG,I)/A(IG)  
RETURN  
END

NOR20001  
NOR20002  
NOR20003  
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NOR20022  
NOR20023  
NOR20024  
NOR20025  
NOR20026

C  
C  
C

SUBROUTINE PRTPLT(NO,B,N,M,NL,NS,KX,JX,ISP)

\* IDENTICAL TO SUBROUTINE PRTPLT PREVIOUSLY LISTED IN PROGRAM REF2G.

RETURN  
END

P RTP0001  
P RTP0002  
P RTP0003  
P RTP0004  
P RTP0005  
P RTP0006

F.3. SOURCE LISTING of Program CUBIC

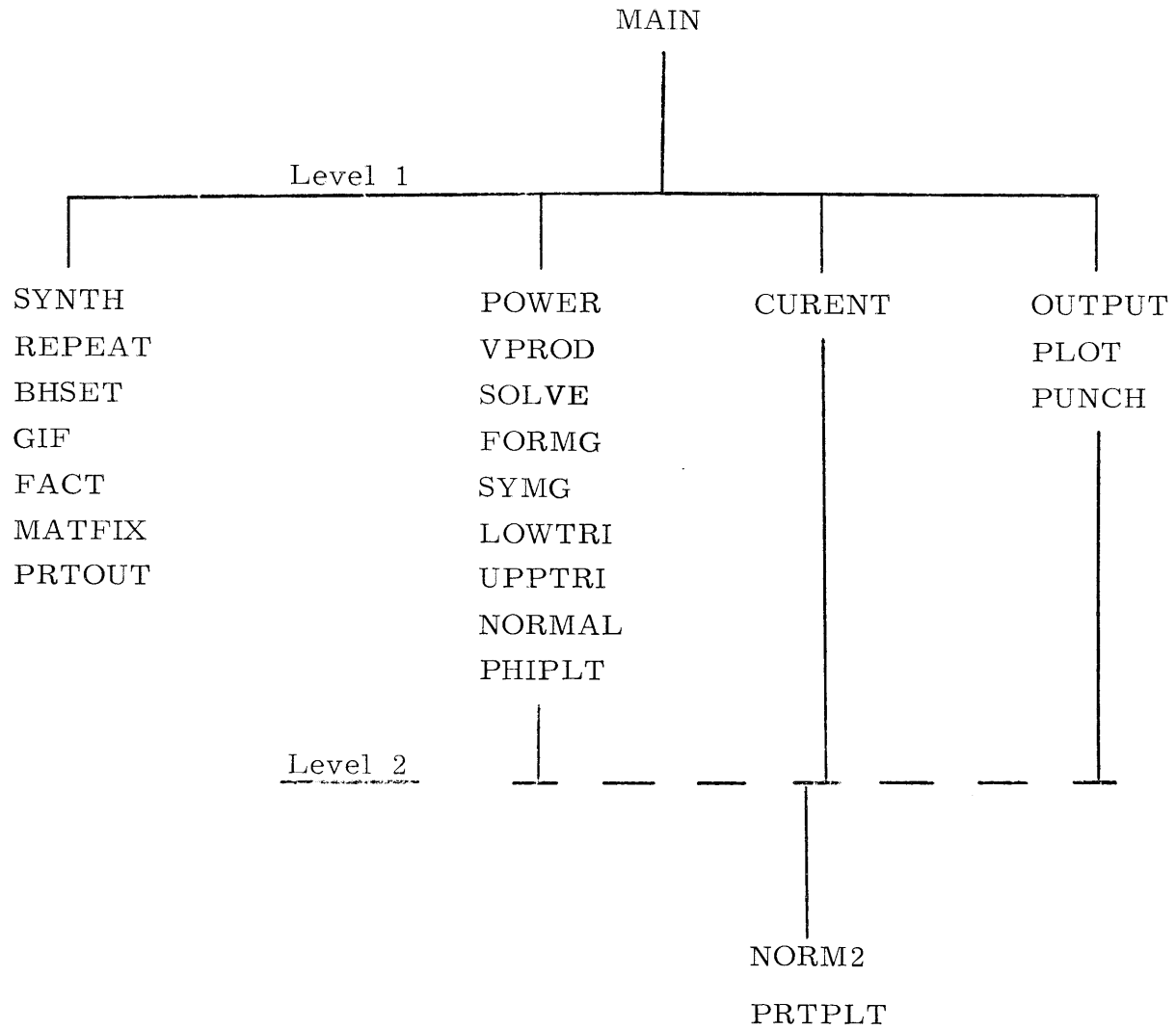


Figure F.3. Structure of Program CUBIC.

C	PROGRAM CUBIC:	CUBI0001
C	TWO GROUP PROPOSED METHOD USING CUBIC HERMITE BASIS FUNCTIONS.	CUBI0002
	CALL TIMING (11)	CUBI0003
	CALL SYNTH	CUBI0004
	CALL TIMING (14)	CUBI0005
	CALL POWER	CUBI0006
	CALL TIMING (16)	CUBI0007
	CALL CURENT	CUBI0008
	CALL TIMING (17)	CUBI0009
	CALL OUTPUT	CUBI0010
	CALL TIMING (18)	CUBI0011
C	TIMING EXECUTION	CUBI0012
	WRITE (6,30)	CUBI0013
	30 FORMAT (1H1,'TIMING PROGRAM EXECUTION:',/)	CUBI0014
	J=14-11	CUBI0015
	WRITE(6,701) J	CUBI0016
	J=16-14	CUBI0017
	WRITE(6,704) J	CUBI0018
	J=17-16	CUBI0019
	WRITE(6,706) J	CUBI0020
	J=18-17	CUBI0021
	WRITE(6,707) J	CUBI0022
	701 FORMAT (1H , ' SYNTH HAS TAKEN',I6,' /100 SECONDS.')	CUBI0023
	704 FORMAT (1H , ' POWER HAS TAKEN',I6,' /100 SECONDS.')	CUBI0024
	706 FORMAT (1H , ' CURENT HAS TAKEN',I5,' /100 SECONDS.')	CUBI0025
	707 FORMAT (1H , ' OUTPUT HAS TAKEN',I5,' /100 SECONDS.')	CUBI0026
	CALL TIMING (120)	CUBI0027
	J=120-11	CUBI0028
	WRITE(6,720) J	CUBI0029
	720 FORMAT (1H0,' THIS RUN HAS TAKEN',I6,' /100 SECONDS TO RUN.')	CUBI0030
	STOP	CUBI0031
	END	CUBI0032

```

SUBROUTINE SYNTH
C      PROPOSED CUBIC HERMITE SYNTHESIS METHOD:
C *****
C      ADJOINT QUANTITIES OF VARIABLES ARE DENOTED BY 7 RATHER THAN *.
C      THUS: PHI7 (RATHER THAN PHI*) IS THE ADJOINT OF PHI. ETC.
      IMPLICIT REAL*8 (A-H,K-Z)
      COMMON /B1/ IBC, IPLOT, JPLOT, IPUNCH, ISEE
      COMMON /B2/ KR, NN
      COMMON /B3/ LT(50,6,2), FT(50,6,4), T(50,6)
      COMMON /B5/
X     KA0(2,25), KA1(2,25), KA2(2,25), KA3(2,25), KA4(2,25), KA5(2,25),
X     KA6(2,25), KB0(2,25), KB1(2,25), KB2(2,25), KB3(2,25), KB4(2,25),
X     KB5(2,25), KB6(2,25), LA0(2,25), LA1(2,25), LA2(2,25), LA3(2,25),
X     LA4(2,25), LA5(2,25), LA6(2,25), P0(2,25) , P1(2,25) , P2(2,25) ,
X     P3(2,25) , P4(2,25) , P5(2,25) , P6(2,25) , Q0(2,25) , Q1(2,25) ,
X     Q2(2,25) , Q3(2,25) , Q4(2,25) , Q5(2,25) , Q6(2,25) , R0(2,25) ,
X     R1(2,25) , R2(2,25) , R3(2,25) , R4(2,25) , SR0(1,25), SR1(1,25),
X     SR2(1,25), SR3(1,25), SR4(1,25), SR5(1,25), SR6(1,25), KC0(1,25),
X     KC1(1,25), KC2(1,25), KC3(1,25), KC4(1,25), KC5(1,25), KC6(1,25),
X     KD0(1,25), KD1(1,25), KD2(1,25), KD3(1,25), KD4(1,25), KD5(1,25),
X     KD6(1,25),
X     P0(2,25) , PH(2,25) , P07(2,25), PH7(2,25), D0(2,25) , DH(2,25) ,
X     C0(2), CH(2), TITLE(20), ITF(25), KTF(25)
      COMMON /CHIF/ CHI(2)
      COMMON /XAXIS/ HX, HR(25)
      COMMON /BH/ X(101), H(101)
      COMMON /ER/ EPS1, EPS2, EPS3
      DIMENSION PHI(2,101), PHI7(2,101), CUR(2,101), CUR7(2,101)
      DIMENSION A(2,100), F(2,100), D(2,100), S(2,100), DI(2,100), XU(2,100)
C      IN ORDER TO SAVE SPACE:
      EQUIVALENCE (XU(1),LT(1)), (A(1),LT(201)), (F(1),LT(401)),
X     (DI(1),FT(1)), (D(1),FT(201)), (S(1),FT(401))
      REAL TITLE
      INTEGER KR, K, KS, KS1, KRO, NN, NUMITF, KTF, N
      READ (5,200) TITLE
200 FORMAT (20A4)

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SYNT0001
SYNT0002
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SYNT0035
SYNT0036

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	WRITE (6,201) TITLE	SYNT0037
201	FORMAT (1H1,20A4,//)	SYNT0038
C	READ IN THE NUMBER OF REGION TRIAL FUNCTIONS AND TYPE OF B.C.S.	SYNT0039
C	AS WELL AS THE TOLERANCES AND THE OUTPUT TYPES DESIRED:	SYNT0040
	READ (5,1) KR,IBC,EPS1,EPS2,EPS3,IPLCT,JPLOT,IPUNCH,ISEE,ITW,ITC	SYNT0041
1	FORMAT (2I5,3D10.3,6I5)	SYNT0042
	IF (IBC.EQ.3) IBC=2	SYNT0043
C	READ IN THE TYPE-NUMBER OF EACH TF REGION:	SYNT0044
	READ (5,100) (ITF(I),I=1,KR)	SYNT0045
100	FORMAT (25I2)	SYNT0046
C	READ IN THE FISSION YIELDS FOR EACH GROUP:	SYNT0047
C	AND THE MATRIX NORMALIZATION PARAMETER: THETA (DEFAULT = 1.0):	SYNT0048
	READ (5,101) CHI(1), CHI(2), THETA	SYNT0049
101	FORMAT (3F10.5)	SYNT0050
	IF (THETA.EQ.0.0) THETA=1.0	SYNT0051
	KRO=KR-1	SYNT0052
	WRITE (6,2) KR, IBC, ISEE, ITW, ITC	SYNT0053
2	FORMAT ('0ONE DIMENSIONAL TWO GROUP CUBIC SYNTHESIS PROGRAM:',//,	SYNT0054
X	5X,'NUMBER OF COARSE MESH REGIONS: KR = ',I2,/,	SYNT0055
X	5X,'BOUNDARY CONDITION NUMBER: IBC = ',I2,/,	SYNT0056
X	5X,'AMOUNT OF OUTPUT REQUESTED: ISEE = ',I2,//,	SYNT0057
X	5X,'TYPE OF WEIGHTING FUNCTIONS: ITW = ',I2,/,	SYNT0058
X	5X,'TYPE OF CURRENT FUNCTIONS: ITC = ',I2,//,	SYNT0059
X	5X,'REGIONAL INPUT MATERIAL PROPERTIES AND FLUX SHAPES FOLLOW',	SYNT0060
X	/,5X,'IF ISEE > 0:',//,	SYNT0061
X	5X,'FLUX SHAPES ARE LINEAR IN EACH INDICATED SUBREGION.')	SYNT0062
	IF (ITC.EQ.0) WRITE (6,16)	SYNT0063
	IF (ITC.EQ.1) WRITE (6,17)	SYNT0064
16	FORMAT (5X,'CURRENTS ARE CONSTANT IN EACH INDICATED SUBREGION.')	SYNT0065
17	FORMAT (5X,'CURRENTS ARE LINEAR IN EACH INDICATED SUBREGION.')	SYNT0066
	IF (ITW.EQ.0) WRITE (6,116)	SYNT0067
	IF (ITW.EQ.1) WRITE (6,117)	SYNT0068
116	FORMAT (/,5X,'WEIGHTING FLUX = FLUX;',/,5X,'WEIGHTING CURRENT = -	SYNT0069
	XCURRENT.')	SYNT0070
117	FORMAT (/,5X,'WEIGHTING FLUX = ADJOINT FLUX;',/,5X,'WEIGHTING CURR	SYNT0071
	XENT = ADJOINT CURRENT.')	SYNT0072

	WRITE (6,20) EPS1, EPS2, EPS3, IPLOT, JPLOT, IPUNCH	SYNT0073
20	FORMAT (//, '0TOLERANCES TO POWER ARE : EPS1 = ', 1PD10.3, /,	SYNT0074
X	28X, 'EPS2 = ', 1PD10.3, /, 28X, 'EPS3 = ', 1PD10.3, /,	SYNT0075
X	'0OUTPUT PARAMETERS TO POWER ARE: IPLOT = ', I1, /,	SYNT0076
X	34X, 'JPLOT = ', I1, /, 34X, 'IPUNCH = ', I1)	SYNT0077
	WRITE (6,22) CHI(1), CHI(2), THETA	SYNT0078
22	FORMAT (/, '0FISSION YIELDS ARE: CHI(1) =', F10.5, /,	SYNT0079
X	22X, 'CHI(2) =', F10.5, /,	SYNT0080
X	'0INPUT THETA PARAMETER (FOR MATRICES) =', D15.7)	SYNT0081
	IF ((KR.LE.2).AND.(IBC.EQ.1)) CALL ERROR(1,KR)	SYNT0082
	IF (KR.GT.25) CALL ERROR(2,KR)	SYNT0083
	IF (EPS1.LT.1.0E-16) CALL ERROR(6,1)	SYNT0084
	IF (EPS2.LT.1.0E-16) CALL ERROR(6,2)	SYNT0085
	IF (EPS3.LT.1.0E-16) CALL ERROR(6,3)	SYNT0086
	IF ((IBC.LT.1).OR.(IBC.GT.4)) CALL ERROR(7,IBC)	SYNT0087
C	DUMMY NORMAL VECTOR XU = UNITY. (FOR THE INTEGRATION FUNCTIONS)	SYNT0088
	DO 21 IG=1,2	SYNT0089
	DO 21 II=1,100	SYNT0090
21	XU(IG,II)=1.0	SYNT0091
	ITC0=2	SYNT0092
	ITC1=2	SYNT0093
	IF (ITC.EQ.1) GO TO 23	SYNT0094
	ITC0=0	SYNT0095
	ITC1=1	SYNT0096
C	COUNTER OF THE NUMBER OF TYPE-NUMBERS OF EACH TF REGION:	SYNT0097
23	NUMITF=1	SYNT0098
	HX=0.0	SYNT0099
C	BEGIN TO READ IN THE TF REGION DATA AND FILL THE ARRAYS,	SYNT0100
C	DEPENDING ON THE TYPE-NUMBER OF EACH TF REGION.	SYNT0101
	DO 50 I=1,KR	SYNT0102
	IF (ITF(I).EQ.NUMITF) GO TO 110	SYNT0103
C	FILL THE ARRAYS FROM OLD TF REGION TYPES:	SYNT0104
	J=ITF(I)	SYNT0105
	CALL REPEAT(I,KTF(J))	SYNT0106
	GO TO 50	SYNT0107
C	READ IN THE TF REGION'S DATA FOR NEW TF REGION TYPE-NUMBERS:	SYNT0108

110	NUMITF=NUMITF+1	SYNT0109
	KTF(NUMITF-1)=1	SYNT0110
C	READ THE SUBREGION NUMBER AND THE NUMBER OF REGIONS IN THE SUBREGION:	SYNT0111
	READ (5,1) K, KS	SYNT0112
	IF (KS.GT.100) CALL ERROR(3,I)	SYNT0113
	KS1=KS+1	SYNT0114
C	CHECK FOR IMPROPER SEQUENCING OF INPUT DATA:	SYNT0115
	IF (I.NE.K) CALL ERROR(4,I)	SYNT0116
C	READ IN THE GEOMETRY AND THE MATERIAL PROPERTIES:	SYNT0117
	READ (5,3) (X(J),X(J+1),H(J),A(1,J),F(1,J),D(1,J),S(1,J),	SYNT0118
	X A(2,J),F(2,J),D(2,J),J=1,KS)	SYNT0119
3	FORMAT (3F10.5,4D10.3,/,30X,3D10.3)	SYNT0120
C	READ IN THE REGIONAL GROUP TRIAL FUNCTIONS:	SYNT0121
	READ (5,4) (PHI(1,J),CUR(1,J),PHI7(1,J),CUR7(1,J),J=1,KS1)	SYNT0122
	READ (5,4) (PHI(2,J),CUR(2,J),PHI7(2,J),CUR7(2,J),J=1,KS1)	SYNT0123
4	FORMAT (4D20.7)	SYNT0124
	IF (ITW.EQ.1) GO TO 120	SYNT0125
C	FORM WEIGHTING FUNCTIONS FROM THE GIVEN FUNCTIONS:	SYNT0126
	DO 119 IG=1,2	SYNT0127
	DO 119 J=1,KS1	SYNT0128
	PHI7(IG,J)=PHI(IG,J)	SYNT0129
119	CUR7(IG,J)=-CUR(IG,J)	SYNT0130
120	IF (ITC.EQ.1) GO TO 5	SYNT0131
C	FORM THE REGION CONSTANT CURRENTS FROM THE FLUXES:	SYNT0132
	DO 7 IG=1,2	SYNT0133
	DO 6 J=1,KS	SYNT0134
	CUR(IG,J)=-D(IG,J)*(-PHI(IG,J)+PHI(IG,J+1))/H(J)	SYNT0135
6	CUR7(IG,J)=+D(IG,J)*(-PHI7(IG,J)+PHI7(IG,J+1))/H(J)	SYNT0136
	CUR(IG,KS1)=0.0	SYNT0137
7	CUR7(IG,KS1)=0.0	SYNT0138
C	WRITE OUT THE INPUT INFORMATION IF ISEE .GE.2:	SYNT0139
5	IF (ISEE.LE.1) GO TO 14	SYNT0140
	WRITE (6,10) K,KS,(J,X(J),X(J+1),H(J),A(1,J),F(1,J),D(1,J),	SYNT0141
	X S(1,J),A(2,J),F(2,J),D(2,J),J=1,KS)	SYNT0142
10	FORMAT ('INPUT MATERIAL PROPERTIES FOR REGION NUMBER ',I3,	SYNT0143
	X ', OF THE ',I3,' USED.',//,	SYNT0144

X	5X,'THIS REGION IS DIVIDED INTO ',I3,' HOMOGENEOUS SUBREGIONS A	SYNT0145
XS	FOLLOWS:',//,	SYNT0146
X	5X,'FAST GROUP CONSTANTS APPEAR FIRST:',//,	SYNT0147
X	' SUBREGION #',5X,'INTERNAL BOUNDARIES',10X,'WIDTH',3X,	SYNT0148
X	' TOTAL CX (1/CM)',3X,'FISSION CX (1/CM)',6X,'DIFFUSION (CM)',	SYNT0149
X	4X,'SCATT.CX (1/CM)',/,	SYNT0150
X	5X,'I',11X,'X(I)',9X,'X(I+1)',11X,'H(I)',13X,'A(IG,I)',13X,	SYNT0151
X	'F(IG,I)',13X,'D(IG,I)',14X,'S(1,I)',//,	SYNT0152
X	(I6,3F15.4,4D20.8,/,5I1X,3D20.8))	SYNT0153
	DO 15 IG=1,2	SYNT0154
15	WRITE (6,11) IG,K,KR,(J,X(J),PHI(IG,J),CUR(IG,J),PHI7(IG,J),	SYNT0155
X	CUR7(IG,J),J=1,KS1)	SYNT0156
11	FORMAT ('INPUT TRIAL FUNCTIONS FOR GROUP',I2,' FOR REGION',I3,	SYNT0157
X	' OUT OF THE',I3,' USED:',//,	SYNT0158
X	' INDEX',5X,'COORD',16X,'FLUX',13X,'CURRENT',8X,' WEIGHT FLUX',	SYNT0159
X	5X,' WEIGHT CURRENT',//,(I6,F10.5,4D20.7))	SYNT0160
14	CONTINUE	SYNT0161
C	END OF THE IN-OUT SECTION:	SYNT0162
C	DEFINING MISC. ARRAYS FOR THE INTEGRATION FUNCTIONS:	SYNT0163
C	LEGNTH OF THE SUBREGION: HT	SYNT0164
	HT=X(KS1)-X(1)	SYNT0165
	HR(K)=HT	SYNT0166
	HX=HX+HR(K)	SYNT0167
C	INVERSE OF THE D ARRAYS:	SYNT0168
	DO 13 J=1,KS	SYNT0169
	DI(1,J)=1./D(1,J)	SYNT0170
13	DI(2,J)=1./D(2,J)	SYNT0171
C	FORMATION OF THE INTEGRATION FUNCTIONS:	SYNT0172
	CALL BHSET(KS)	SYNT0173
C	DO FOR ALL ENERGY GROUPS:	SYNT0174
	DO 50 IG=1,2	SYNT0175
	KA0(IG,K)=GIF(0,IG,PHI7,IG,A,PHI,KS,2)	SYNT0176
	KA1(IG,K)=GIF(1,IG,PHI7,IG,A,PHI,KS,2)	SYNT0177
	KA2(IG,K)=GIF(2,IG,PHI7,IG,A,PHI,KS,2)	SYNT0178
	KA3(IG,K)=GIF(3,IG,PHI7,IG,A,PHI,KS,2)	SYNT0179
	KA4(IG,K)=GIF(4,IG,PHI7,IG,A,PHI,KS,2)	SYNT0180

KA5(IG,K)=GIF(5,IG,PHI7,IG,A,PHI,KS,2)  
KA6(IG,K)=GIF(6,IG,PHI7,IG,A,PHI,KS,2)  
KB0(IG,K)=GIF(0,IG,PHI7,IG,F,PHI,KS,2)  
KB1(IG,K)=GIF(1,IG,PHI7,IG,F,PHI,KS,2)  
KB2(IG,K)=GIF(2,IG,PHI7,IG,F,PHI,KS,2)  
KB3(IG,K)=GIF(3,IG,PHI7,IG,F,PHI,KS,2)  
KB4(IG,K)=GIF(4,IG,PHI7,IG,F,PHI,KS,2)  
KB5(IG,K)=GIF(5,IG,PHI7,IG,F,PHI,KS,2)  
KB6(IG,K)=GIF(6,IG,PHI7,IG,F,PHI,KS,2)  
LA0(IG,K)=GIF(0,IG,CUR7,IG,DI,CUR,KS,ITC0)  
LA1(IG,K)=GIF(1,IG,CUR7,IG,DI,CUR,KS,ITC0)  
LA2(IG,K)=GIF(2,IG,CUR7,IG,DI,CUR,KS,ITC0)  
LA3(IG,K)=GIF(3,IG,CUR7,IG,DI,CUR,KS,ITC0)  
LA4(IG,K)=GIF(4,IG,CUR7,IG,DI,CUR,KS,ITC0)  
LA5(IG,K)=GIF(5,IG,CUR7,IG,DI,CUR,KS,ITC0)  
LA6(IG,K)=GIF(6,IG,CUR7,IG,DI,CUR,KS,ITC0)  
P0(IG,K)=GIF(0,IG,PHI7,IG,XU,CUR,KS,ITC1)/HT  
P1(IG,K)=GIF(1,IG,PHI7,IG,XU,CUR,KS,ITC1)/HT  
P2(IG,K)=GIF(2,IG,PHI7,IG,XU,CUR,KS,ITC1)/HT  
P3(IG,K)=GIF(3,IG,PHI7,IG,XU,CUR,KS,ITC1)/HT  
P4(IG,K)=GIF(4,IG,PHI7,IG,XU,CUR,KS,ITC1)/HT  
P5(IG,K)=GIF(5,IG,PHI7,IG,XU,CUR,KS,ITC1)/HT  
P6(IG,K)=GIF(6,IG,PHI7,IG,XU,CUR,KS,ITC1)/HT  
Q0(IG,K)=GIF(0,IG,PHI,IG,XU,CUR7,KS,ITC1)/HT  
Q1(IG,K)=GIF(1,IG,PHI,IG,XU,CUR7,KS,ITC1)/HT  
Q2(IG,K)=GIF(2,IG,PHI,IG,XU,CUR7,KS,ITC1)/HT  
Q3(IG,K)=GIF(3,IG,PHI,IG,XU,CUR7,KS,ITC1)/HT  
Q4(IG,K)=GIF(4,IG,PHI,IG,XU,CUR7,KS,ITC1)/HT  
Q5(IG,K)=GIF(5,IG,PHI,IG,XU,CUR7,KS,ITC1)/HT  
Q6(IG,K)=GIF(6,IG,PHI,IG,XU,CUR7,KS,ITC1)/HT  
R0(IG,K)=GIF(0,IG,PHI7,IG,D,PHI,KS,2)/HT\*\*2  
R1(IG,K)=GIF(1,IG,PHI7,IG,D,PHI,KS,2)/HT\*\*2  
R2(IG,K)=GIF(2,IG,PHI7,IG,D,PHI,KS,2)/HT\*\*2  
R3(IG,K)=GIF(3,IG,PHI7,IG,D,PHI,KS,2)/HT\*\*2  
R4(IG,K)=GIF(4,IG,PHI7,IG,D,PHI,KS,2)/HT\*\*2  
STORE THE TERMINAL POINTS FOR LATER USE:

SYNT0181  
SYNT0182  
SYNT0183  
SYNT0184  
SYNT0185  
SYNT0186  
SYNT0187  
SYNT0188  
SYNT0189  
SYNT0190  
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SYNT0198  
SYNT0199  
SYNT0200  
SYNT0201  
SYNT0202  
SYNT0203  
SYNT0204  
SYNT0205  
SYNT0206  
SYNT0207  
SYNT0208  
SYNT0209  
SYNT0210  
SYNT0211  
SYNT0212  
SYNT0213  
SYNT0214  
SYNT0215  
SYNT0216

C

	PO(IG,K)=PHI(IG,1)	SYNT0217
	PO7(IG,K)=PHI7(IG,1)	SYNT0218
	PH(IG,K)=PHI(IG,KS1)	SYNT0219
	PH7(IG,K)=PHI7(IG,KS1)	SYNT0220
	DO(IG,K)=D(IG,1)	SYNT0221
	DH(IG,K)=D(IG,KS)	SYNT0222
	IF (K.EQ.1) CO(IG)=CUR(IG,1)	SYNT0223
	IF (NUMITF-1.EQ.ITF(KR).AND.ITC.EQ.0) CH(IG)=CUR(IG,KS)	SYNT0224
	IF (NUMITF-1.EQ.ITF(KR).AND.ITC.EQ.1) CH(IG)=CUR(IG,KS1)	SYNT0225
C	FOR THE OFF DIAGONAL MATRIX ELEMENTS:	SYNT0226
	IF (IG.EQ.2) GO TO 50	SYNT0227
	SR0(1,K)=GIF(0,2,PHI7,1,S,PHI,KS,2)	SYNT0228
	SR1(1,K)=GIF(1,2,PHI7,1,S,PHI,KS,2)	SYNT0229
	SR2(1,K)=GIF(2,2,PHI7,1,S,PHI,KS,2)	SYNT0230
	SR3(1,K)=GIF(3,2,PHI7,1,S,PHI,KS,2)	SYNT0231
	SR4(1,K)=GIF(4,2,PHI7,1,S,PHI,KS,2)	SYNT0232
	SR5(1,K)=GIF(5,2,PHI7,1,S,PHI,KS,2)	SYNT0233
	SR6(1,K)=GIF(6,2,PHI7,1,S,PHI,KS,2)	SYNT0234
	KC0(1,K)=GIF(0,1,PHI7,2,F,PHI,KS,2)	SYNT0235
	KC1(1,K)=GIF(1,1,PHI7,2,F,PHI,KS,2)	SYNT0236
	KC2(1,K)=GIF(2,1,PHI7,2,F,PHI,KS,2)	SYNT0237
	KC3(1,K)=GIF(3,1,PHI7,2,F,PHI,KS,2)	SYNT0238
	KC4(1,K)=GIF(4,1,PHI7,2,F,PHI,KS,2)	SYNT0239
	KC5(1,K)=GIF(5,1,PHI7,2,F,PHI,KS,2)	SYNT0240
	KC6(1,K)=GIF(6,1,PHI7,2,F,PHI,KS,2)	SYNT0241
	KD0(1,K)=GIF(0,2,PHI7,1,F,PHI,KS,2)	SYNT0242
	KD1(1,K)=GIF(1,2,PHI7,1,F,PHI,KS,2)	SYNT0243
	KD2(1,K)=GIF(2,2,PHI7,1,F,PHI,KS,2)	SYNT0244
	KD3(1,K)=GIF(3,2,PHI7,1,F,PHI,KS,2)	SYNT0245
	KD4(1,K)=GIF(4,2,PHI7,1,F,PHI,KS,2)	SYNT0246
	KD5(1,K)=GIF(5,2,PHI7,1,F,PHI,KS,2)	SYNT0247
	KD6(1,K)=GIF(6,2,PHI7,1,F,PHI,KS,2)	SYNT0248
50	CONTINUE	SYNT0249
	NUMITF=NUMITF-1	SYNT0250
	WRITE (6,51) NUMITF	SYNT0251
51	FORMAT ('1THERE ARE ONLY',I3,' DIFFERENT TRIAL FUNCTION REGIONS.')	SYNT0252

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WRITE (6,52) (1,ITF(I),I=1,KR)
52 FORMAT (/, 'OTABLE OF THE TRIAL FUNCTION NUMBER TYPES:', //,
X 3X, 'TF REGION', 4X, 'REGION TYPE-NUMBER', //,
X (17,12X,I7))
C TO PRINT OUT THE /B5/ ARRAYS:
C IF (ISEE.GE.2) CALL PRTOUT(1)
C DETERMINATIONS OF THE B.C. OPTION PARAMETERS:
C NN IS THE BLOCK SIZE OF THE POWER MATRICES.
NN=2*KR
C FORMATION OF THE COEFFICIENT VECTORS:
C FILLING THE MATRICES FOR POWER:
C FOR BOTH ENERGY GROUPS:
DO 60 IG=1,2
C THE MATRIX ROW INDEX:
I=1
C FOR ALL THE INTERIOR COEFFICIENTS:
DO 60 K=2,KR
I=I+1
56 J=K-1
V=1./(PH7(IG,J)*PO(IG,J))
V1=1./(PH7(IG,J)*PH(IG,J))
V2=1./(PO7(IG,K)*PO(IG,K))
V3=1./(PO7(IG,K)*PH(IG,K))
LT(I,1,IG)=(3.*KA2(IG,J)-2.*KA3(IG,J)-9.*KA4(IG,J)+12.*KA5(IG,J)
X -4.*KA6(IG,J)-(3.*LA2(IG,J)-2.*LA3(IG,J)-9.*LA4(IG,J)+12.*
X LA5(IG,J)-4.*LA6(IG,J))-(6.*P1(IG,J)-6.*P2(IG,J)-18.*P3(IG,J)
X +30.*P4(IG,J)-12.*P5(IG,J))-(18.*Q3(IG,J)-30.*Q4(IG,J)+12.*
X Q5(IG,J))-(36.*R2(IG,J)-72.*R3(IG,J)+36.*R4(IG,J))*V
LT(I,2,IG)=(-3.*KA3(IG,J)+8.*KA4(IG,J)-7.*KA5(IG,J)+2.*KA6(IG,J)
X -(-3.*LA3(IG,J)+8.*LA4(IG,J)-7.*LA5(IG,J)+2.*LA6(IG,J))
X -(-6.*P2(IG,J)+18.*P3(IG,J)-18.*P4(IG,J)+6.*P5(IG,J))
X -(3.*Q2(IG,J)-14.*Q3(IG,J)+17.*Q4(IG,J)-6.*Q5(IG,J))
X -(6.*R1(IG,J)-30.*
X R2(IG,J)+42.*R3(IG,J)-18.*R4(IG,J))*V*HR(J)/DO(IG,J)
LT(I,3,IG)=(9.*KA4(IG,J)-12.*KA5(IG,J)+4.*KA6(IG,J)-(9.*LA4(IG,J)
X -12.*LA5(IG,J)+4.*LA6(IG,J))-(18.*P3(IG,J)-30.*P4(IG,J)+12.*

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SYNT0253
SYNT0254
SYNT0255
SYNT0256
SYNT0257
SYNT0258
SYNT0259
SYNT0260
SYNT0261
SYNT0262
SYNT0263
SYNT0264
SYNT0265
SYNT0266
SYNT0267
SYNT0268
SYNT0269
SYNT0270
SYNT0271
SYNT0272
SYNT0273
SYNT0274
SYNT0275
SYNT0276
SYNT0277
SYNT0278
SYNT0279
SYNT0280
SYNT0281
SYNT0282
SYNT0283
SYNT0284
SYNT0285
SYNT0286
SYNT0287
SYNT0288

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X	P5(IG,J))+(18.*Q3(IG,J)-30.*Q4(IG,J)+12.*Q5(IG,J))	SYNT0289
X	+36.*R2(IG,J)-72.*R3(IG,J)+36.*R4(IG,J))*V1 +	SYNT0290
X	(KA0(IG,K)-6.*KA2(IG,K)+4.*KA3(IG,K)+9.*KA4(IG,K)-12.*KA5(IG,K)	SYNT0291
X	+4.*KA6(IG,K)-(LA0(IG,K)-6.*LA2(IG,K)+4.*LA3(IG,K)+9.*LA4(IG,K)	SYNT0292
X	-12.*LA5(IG,K)+4.*LA6(IG,K))+6.*P1(IG,K)-6.*P2(IG,K)-18.*	SYNT0293
X	P3(IG,K)+30.*P4(IG,K)-12.*P5(IG,K)-(6.*Q1(IG,K)-6.*Q2(IG,K)	SYNT0294
X	-18.*Q3(IG,K)+30.*Q4(IG,K)-12.*Q5(IG,K))+36.*R2(IG,K)-72.*	SYNT0295
X	R3(IG,K)+36.*R4(IG,K))*V2	SYNT0296
	LT(I,4,IG)=(3.*KA4(IG,J)-5.*KA5(IG,J)+2.*KA6(IG,J)-(3.*LA4(IG,J)	SYNT0297
X	-5.*LA5(IG,J)+2.*LA6(IG,J))-6.*P3(IG,J)-12.*P4(IG,J)+6.*	SYNT0298
X	P5(IG,J))-(-6.*Q3(IG,J)+13.*Q4(IG,J)-6.*Q5(IG,J))-	SYNT0299
X	-12.*R2(IG,J)	SYNT0300
X	+30.*R3(IG,J)-18.*R4(IG,J))*V1*HR(J)/DH(IG,J)	SYNT0301
X	+(-KA1(IG,K)+2.*KA2(IG,K)+2.*KA3(IG,K)-8.*KA4(IG,K)+7.*	SYNT0302
X	KA5(IG,K)-2.*KA6(IG,K)-(-LA1(IG,K)+2.*LA2(IG,K)+2.*LA3(IG,K)	SYNT0303
X	-8.*LA4(IG,K)+7.*LA5(IG,K)-2.*LA6(IG,K))-6.*P2(IG,K)+18.*	SYNT0304
X	P3(IG,K)-18.*P4(IG,K)+6.*P5(IG,K)-(Q0(IG,K)-4.*Q1(IG,K)+14.*	SYNT0305
X	Q3(IG,K)-17.*Q4(IG,K)+6.*Q5(IG,K))+6.*R1(IG,K)-30.*R2(IG,K)	SYNT0306
X	+42.*R3(IG,K)-18.*R4(IG,K))*V2*HR(K)/DO(IG,K)	SYNT0307
	LT(I,5,IG)=(3.*KA2(IG,K)-2.*KA3(IG,K)-9.*KA4(IG,K)+12.*KA5(IG,K)	SYNT0308
X	-4.*KA6(IG,K)-(3.*LA2(IG,K)-2.*LA3(IG,K)-9.*LA4(IG,K)+12.*	SYNT0309
X	LA5(IG,K)-4.*LA6(IG,K))+18.*P3(IG,K)-30.*P4(IG,K)+12.*P5(IG,K)	SYNT0310
X	+(6.*Q1(IG,K)-6.*Q2(IG,K)-18.*Q3(IG,K)+30.*Q4(IG,K)-12.*	SYNT0311
X	Q5(IG,K))-36.*R2(IG,K)-72.*R3(IG,K)+36.*R4(IG,K))*V3	SYNT0312
	LT(I,6,IG)=(KA2(IG,K)-KA3(IG,K)-3.*KA4(IG,K)+5.*KA5(IG,K)-2.*	SYNT0313
X	KA6(IG,K)-(LA2(IG,K)-LA3(IG,K)-3.*LA4(IG,K)+5.*LA5(IG,K)-2.*	SYNT0314
X	LA6(IG,K))+6.*P3(IG,K)-12.*P4(IG,K)+6.*P5(IG,K)-(-2.*Q1(IG,K)+	SYNT0315
X	3.*Q2(IG,K)+6.*Q3(IG,K)-13.*Q4(IG,K)+6.*Q5(IG,K))-12.*R2(IG,K)	SYNT0316
X	+30.*R3(IG,K)-18.*R4(IG,K))*V3*HR(K)/DH(IG,K)	SYNT0317
	FT(I,1,IG)=(3.*KB2(IG,J)-2.*KB3(IG,J)-9.*KB4(IG,J)+12.*KB5(IG,J)	SYNT0318
X	-4.*KB6(IG,J))*V	SYNT0319
	FT(I,2,IG)=(-3.*KB3(IG,J)+8.*KB4(IG,J)-7.*KB5(IG,J)+2.*KB6(IG,J))	SYNT0320
X	*V*HR(J)/DO(IG,J)	SYNT0321
	FT(I,3,IG)=(9.*KB4(IG,J)-12.*KB5(IG,J)+4.*KB6(IG,J))*V1	SYNT0322
X	+(KB0(IG,K)-6.*KB2(IG,K)+4.*KB3(IG,K)+9.*KB4(IG,K)	SYNT0323
X	-12.*KB5(IG,K)+4.*KB6(IG,K))*V2	SYNT0324



FT(I,4,IG)=(3.*KB4(IG,J)-5.*KB5(IG,J)+2.*KB6(IG,J))	SYNT0325
X *V1*HR(J)/DH(IG,J)+(-KB1(IG,K)+2.*KB2(IG,K)+2.*KB3(IG,K)	SYNT0326
X -8.*KB4(IG,K)+7.*KB5(IG,K)-2.*KB6(IG,K))*V2*HR(K)/DO(IG,K)	SYNT0327
FT(I,5,IG)=(3.*KB2(IG,K)-2.*KB3(IG,K)-9.*KB4(IG,K)+12.*KB5(IG,K)	SYNT0328
X -4.*KB6(IG,K))*V3	SYNT0329
FT(I,6,IG)=(KB2(IG,K)-KB3(IG,K)-3.*KB4(IG,K)+5.*KB5(IG,K)	SYNT0330
X -2.*KB6(IG,K))*V3*HR(K)/DH(IG,K)	SYNT0331
IF (IG.EQ.2) GO TO 57	SYNT0332
T(I,1) =(3.*SR2(IG,J)-2.*SR3(IG,J)-9.*SR4(IG,J)+12.*SR5(IG,J)	SYNT0333
X -4.*SR6(IG,J))/(PH7(2,J)*PO(1,J))	SYNT0334
T(I,2) =(-3.*SR3(IG,J)+8.*SR4(IG,J)-7.*SR5(IG,J)+2.*SR6(IG,J))	SYNT0335
X *HR(J)/(PH7(2,J)*PO(1,J)*DO(1,J))	SYNT0336
T(I,3) =(9.*SR4(IG,J)-12.*SR5(IG,J)+4.*SR6(IG,J))	SYNT0337
X /(PH7(2,J)*PH(1,J))	SYNT0338
X +(SR0(IG,K)-6.*SR2(IG,K)+4.*SR3(IG,K)+9.*SR4(IG,K)	SYNT0339
X -12.*SR5(IG,K)+4.*SR6(IG,K))/(PO7(2,K)*PO(1,K))	SYNT0340
T(I,4) =(3.*SR4(IG,J)-5.*SR5(IG,J)+2.*SR6(IG,J))	SYNT0341
X *HR(J)/(PH7(2,J)*PH(1,J)*DH(1,J))	SYNT0342
X +(-SR1(IG,K)+2.*SR2(IG,K)+2.*SR3(IG,K)	SYNT0343
X -8.*SR4(IG,K)+7.*SR5(IG,K)-2.*SR6(IG,K))*HR(K)	SYNT0344
X /(PO7(2,K)*PO(1,K)*DO(1,K))	SYNT0345
T(I,5) =(3.*SR2(IG,K)-2.*SR3(IG,K)-9.*SR4(IG,K)+12.*SR5(IG,K)	SYNT0346
X -4.*SR6(IG,K))/(PO7(2,K)*PH(1,K))	SYNT0347
T(I,6) =(SR2(IG,K)-SR3(IG,K)-3.*SR4(IG,K)+5.*SR5(IG,K)	SYNT0348
X -2.*SR6(IG,K))*HR(K)/(PO7(2,K)*PH(1,K)*DH(1,K))	SYNT0349
FT(I,1,3)=(3.*KD2(IG,J)-2.*KD3(IG,J)-9.*KD4(IG,J)+12.*KD5(IG,J)	SYNT0350
X -4.*KD6(IG,J))/(PH7(2,J)*PO(1,J))	SYNT0351
FT(I,2,3)=(-3.*KD3(IG,J)+8.*KD4(IG,J)-7.*KD5(IG,J)+2.*KD6(IG,J))	SYNT0352
X *HR(J)/(PH7(2,J)*PO(1,J)*DO(1,J))	SYNT0353
FT(I,3,3)=(9.*KD4(IG,J)-12.*KD5(IG,J)+4.*KD6(IG,J))	SYNT0354
X /(PH7(2,J)*PH(1,J))	SYNT0355
X +(KD0(IG,K)-6.*KD2(IG,K)+4.*KD3(IG,K)+9.*KD4(IG,K)	SYNT0356
X -12.*KD5(IG,K)+4.*KD6(IG,K))/(PO7(2,K)*PO(1,K))	SYNT0357
FT(I,4,3)=(3.*KD4(IG,J)-5.*KD5(IG,J)+2.*KD6(IG,J))	SYNT0358
X *HR(J)/(PH7(2,J)*PH(1,J)*DH(1,J))	SYNT0359
X +(-KD1(IG,K)+2.*KD2(IG,K)+2.*KD3(IG,K)	SYNT0360

X	-8.*KD4(IG,K)+7.*KD5(IG,K)-2.*KD6(IG,K))*HR(K)	SYNT0361
X	/(PO7(2,K)*PO(1,K)*DO(1,K))	SYNT0362
	FT(1,5,3)=(3.*KD2(IG,K)-2.*KD3(IG,K)-9.*KD4(IG,K)+12.*KD5(IG,K)	SYNT0363
X	-4.*KD6(IG,K))/(PO7(2,K)*PH(1,K))	SYNT0364
	FT(1,6,3)=(KD2(IG,K)-KD3(IG,K)-3.*KD4(IG,K)+5.*KD5(IG,K)	SYNT0365
X	-2.*KD6(IG,K))*HR(K)/(PO7(2,K)*PH(1,K)*DH(1,K))	SYNT0366
	FT(1,1,4)=(3.*KC2(IG,J)-2.*KC3(IG,J)-9.*KC4(IG,J)+12.*KC5(IG,J)	SYNT0367
X	-4.*KC6(IG,J))/(PH7(1,J)*PC(2,J))	SYNT0368
	FT(1,2,4)=(-3.*KC3(IG,J)+8.*KC4(IG,J)-7.*KC5(IG,J)+2.*KC6(IG,J))	SYNT0369
X	*HR(J)/(PH7(1,J)*PO(2,J)*DO(2,J))	SYNT0370
	FT(1,3,4)=(9.*KC4(IG,J)-12.*KC5(IG,J)+4.*KC6(IG,J))	SYNT0371
X	/(PH7(1,J)*PH(2,J))	SYNT0372
X	+(KC0(IG,K)-6.*KC2(IG,K)+4.*KC3(IG,K)+9.*KC4(IG,K)	SYNT0373
X	-12.*KC5(IG,K)+4.*KC6(IG,K))/(PO7(1,K)*PO(2,K))	SYNT0374
	FT(1,4,4)=(3.*KC4(IG,J)-5.*KC5(IG,J)+2.*KC6(IG,J))	SYNT0375
X	*HR(J)/(PH7(1,J)*PH(2,J)*DH(2,J))	SYNT0376
X	+(-KC1(IG,K)+2.*KC2(IG,K)+2.*KC3(IG,K)	SYNT0377
X	-8.*KC4(IG,K)+7.*KC5(IG,K)-2.*KC6(IG,K))*HR(K)	SYNT0378
X	/(PO7(1,K)*PO(2,K)*DO(2,K))	SYNT0379
	FT(1,5,4)=(3.*KC2(IG,K)-2.*KC3(IG,K)-9.*KC4(IG,K)+12.*KC5(IG,K)	SYNT0380
X	-4.*KC6(IG,K))/(PO7(1,K)*PH(2,K))	SYNT0381
	FT(1,6,4)=(KC2(IG,K)-KC3(IG,K)-3.*KC4(IG,K)+5.*KC5(IG,K)	SYNT0382
X	-2.*KC6(IG,K))*HR(K)/(PO7(1,K)*PH(2,K)*DH(2,K))	SYNT0383
57	I=I+1	SYNT0384
	LT(1,1,IG)=(KA2(IG,J)-KA3(IG,J)-3.*KA4(IG,J)+5.*KA5(IG,J)-2.*	SYNT0385
X	KA6(IG,J)-(LA2(IG,J)-LA3(IG,J)-3.*LA4(IG,J)+5.*LA5(IG,J)-2.*	SYNT0386
X	LA6(IG,J))-2.*P1(IG,J)+3.*P2(IG,J)+6.*P3(IG,J)-13.*P4(IG,J)+6.*	SYNT0387
X	P5(IG,J)-(6.*Q3(IG,J)-12.*Q4(IG,J)+6.*Q5(IG,J))-12.*R2(IG,J)	SYNT0388
X	+30.*R3(IG,J)-18.*R4(IG,J))*V*HR(J)/DH(IG,J)	SYNT0389
	LT(1,2,IG)=(-KA3(IG,J)+3.*KA4(IG,J)-3.*KA5(IG,J)+KA6(IG,J)	SYNT0390
X	-(-LA3(IG,J)+3.*LA4(IG,J)-3.*LA5(IG,J)+LA6(IG,J))+2.*P2(IG,J)	SYNT0391
X	-7.*P3(IG,J)+8.*P4(IG,J)-3.*P5(IG,J)-(Q2(IG,J)-5.*Q3(IG,J)	SYNT0392
X	+7.*Q4(IG,J)-3.*Q5(IG,J))-2.*R1(IG,J)+11.*R2(IG,J)-18.*	SYNT0393
X	R3(IG,J)+9.*R4(IG,J))*V*HR(J)*HR(J)/(DO(IG,J)*DH(IG,J))	SYNT0394
	LT(1,3,IG)=(3.*KA4(IG,J)-5.*KA5(IG,J)+2.*KA6(IG,J)-(3.*LA4(IG,J)	SYNT0395
X	-5.*LA5(IG,J)+2.*LA6(IG,J))-6.*P3(IG,J)+13.*P4(IG,J)-6.*	SYNT0396

X	P5(IG,J)+(6.*Q3(IG,J)-12.*Q4(IG,J)+6.*Q5(IG,J))-(-12.*R2(IG,J)	SYNT0397
X	+30.*R3(IG,J)-18.*R4(IG,J))*V1*HR(J)/DH(IG,J)	SYNT0398
X	+(-KA1(IG,K)+2.*KA2(IG,K)+2.*KA3(IG,K)-8.*KA4(IG,K)+7.*	SYNT0399
X	KA5(IG,K)-2.*KA6(IG,K)-(-LA1(IG,K)+2.*LA2(IG,K)+2.*LA3(IG,K)-8.	SYNT0400
X	*LA4(IG,K)+7.*LA5(IG,K)-2.*LA6(IG,K))+P0(IG,K)-4.*P1(IG,K)+14.*	SYNT0401
X	P3(IG,K)-17.*P4(IG,K)+6.*P5(IG,K)-(-6.*Q2(IG,K)+18.*Q3(IG,K)	SYNT0402
X	-18.*Q4(IG,K)+6.*Q5(IG,K))+6.*R1(IG,K)-30.*R2(IG,K)+42.*	SYNT0403
X	R3(IG,K)-18.*R4(IG,K))*V2*HR(K)/DO(IG,K)	SYNT0404
	LT(I,4,IG)=(KA4(IG,J)-2.*KA5(IG,J)+KA6(IG,J)-(LA4(IG,J)-2.*	SYNT0405
X	LA5(IG,J)+LA6(IG,J))-2.*P3(IG,J)+5.*P4(IG,J)-3.*P5(IG,J)-(-2.*	SYNT0406
X	Q3(IG,J)+5.*Q4(IG,J)-3.*Q5(IG,J))+4.*R2(IG,J)-12.*R3(IG,J)+9.*	SYNT0407
X	R4(IG,J))*V1*(HR(J)/DH(IG,J))*2	SYNT0408
X	+(KA2(IG,K)-4.*KA3(IG,K)+6.*KA4(IG,K)-4.*KA5(IG,K)+KA6(IG,K)	SYNT0409
X	-(-LA2(IG,K)-4.*LA3(IG,K)+6.*LA4(IG,K)-4.*LA5(IG,K)+LA6(IG,K))	SYNT0410
X	-P1(IG,K)+6.*P2(IG,K)-12.*P3(IG,K)+10.*P4(IG,K)-3.*P5(IG,K)	SYNT0411
X	-(-Q1(IG,K)+6.*Q2(IG,K)-12.*Q3(IG,K)+10.*Q4(IG,K)-3.*Q5(IG,K))	SYNT0412
X	+R0(IG,K)-8.*R1(IG,K)+22.*R2(IG,K)-24.*R3(IG,K)+9.*R4(IG,K))*	SYNT0413
X	V2*(HR(K)/DO(IG,K))*2	SYNT0414
	LT(I,5,IG)=(-3.*KA3(IG,K)+8.*KA4(IG,K)-7.*KA5(IG,K)+2.*KA6(IG,K)	SYNT0415
X	-(-3.*LA3(IG,K)+8.*LA4(IG,K)-7.*LA5(IG,K)+2.*LA6(IG,K))	SYNT0416
X	+3.*P2(IG,K)-14.*P3(IG,K)+17.*P4(IG,K)-6.*P5(IG,K)+(-6.*	SYNT0417
X	Q2(IG,K)+18.*Q3(IG,K)-18.*Q4(IG,K)+6.*Q5(IG,K))-6.*R1(IG,K)	SYNT0418
X	-30.*R2(IG,K)+42.*R3(IG,K)-18.*R4(IG,K))*V3*HR(K)/DO(IG,K)	SYNT0419
	LT(I,6,IG)=(-KA3(IG,K)+3.*KA4(IG,K)-3.*KA5(IG,K)+KA6(IG,K)	SYNT0420
X	-(-LA3(IG,K)+3.*LA4(IG,K)-3.*LA5(IG,K)+LA6(IG,K))	SYNT0421
X	+P2(IG,K)-5.*P3(IG,K)+7.*P4(IG,K)-3.*P5(IG,K)-(2.*Q2(IG,K)	SYNT0422
X	-7.*Q3(IG,K)+8.*Q4(IG,K)-3.*Q5(IG,K))-2.*R1(IG,K)+11.*R2(IG,K)	SYNT0423
X	-18.*R3(IG,K)+9.*R4(IG,K))*V3*HR(K)*HR(K)/(DO(IG,K)*DH(IG,K))	SYNT0424
	FT(I,1,IG)=(KB2(IG,J)-KB3(IG,J)-3.*KB4(IG,J)+5.*KB5(IG,J)	SYNT0425
X	-2.*KB6(IG,J))*V*HR(J)/DH(IG,J)	SYNT0426
	FT(I,2,IG)=(-KB3(IG,J)+3.*KB4(IG,J)-3.*KB5(IG,J)+KB6(IG,J))	SYNT0427
X	*V*HR(J)*HR(J)/(DO(IG,J)*DH(IG,J))	SYNT0428
	FT(I,3,IG)=(3.*KB4(IG,J)-5.*KB5(IG,J)+2.*KB6(IG,J))	SYNT0429
X	*V1*HR(J)/DH(IG,J)+(-KB1(IG,K)+2.*KB2(IG,K)+2.*KB3(IG,K)	SYNT0430
X	-8.*KB4(IG,K)+7.*KB5(IG,K)-2.*KB6(IG,K))*V2*HR(K)/DO(IG,K)	SYNT0431
	FT(I,4,IG)=(KB4(IG,J)-2.*KB5(IG,J)+KB6(IG,J))	SYNT0432

X	*V1*(HR(J)/DH(IG,J))**2+(KB2(IG,K)-4.*KB3(IG,K)+6.*KB4(IG,K)	SYNT0433
X	-4.*KB5(IG,K)+KB6(IG,K))*V2*(HR(K)/DO(IG,K))**2	SYNT0434
	FT(1,5,IG)=(-3.*KB3(IG,K)+8.*KB4(IG,K)-7.*KB5(IG,K)+2.*KB6(IG,K))	SYNT0435
X	*V3*HR(K)/DO(IG,K)	SYNT0436
	FT(1,6,IG)=(-KB3(IG,K)+3.*KB4(IG,K)-3.*KB5(IG,K)+KB6(IG,K))	SYNT0437
X	*V3*HR(K)*HR(K)/(DO(IG,K)*DH(IG,K))	SYNT0438
	IF (IG.EQ.2) GO TO 60	SYNT0439
	T(1,1)=(SR2(IG,J)-SR3(IG,J)-3.*SR4(IG,J)+5.*SR5(IG,J)	SYNT0440
X	-2.*SR6(IG,J))*HR(J)/(PH7(2,J)*PC(1,J)*DH(2,J))	SYNT0441
	T(1,2)=(-SR3(IG,J)+3.*SR4(IG,J)-3.*SR5(IG,J)+SR6(IG,J))	SYNT0442
X	*HR(J)**2/(PH7(2,J)*PO(1,J)*DO(1,J)*DH(2,J))	SYNT0443
	T(1,3)=(3.*SR4(IG,J)-5.*SR5(IG,J)+2.*SR6(IG,J))	SYNT0444
X	*HR(J)/(PH7(2,J)*PH(1,J)*DH(2,J))	SYNT0445
X	+(-SR1(IG,K)+2.*SR2(IG,K)+2.*SR3(IG,K)	SYNT0446
X	-8.*SR4(IG,K)+7.*SR5(IG,K)-2.*SR6(IG,K))*HR(K)	SYNT0447
X	/(PO7(2,K)*PO(1,K)*DO(2,K))	SYNT0448
	T(1,4)=(SR4(IG,J)-2.*SR5(IG,J)+SR6(IG,J))	SYNT0449
X	*HR(J)**2/(PH7(2,J)*PH(1,J)*DH(2,J)*DH(1,J))	SYNT0450
X	+(SR2(IG,K)-4.*SR3(IG,K)	SYNT0451
X	+6.*SR4(IG,K)-4.*SR5(IG,K)+SR6(IG,K))	SYNT0452
X	*HR(K)**2/(PO7(2,K)*PO(1,K)*DO(2,K)*DO(1,K))	SYNT0453
	T(1,5)=(-3.*SR3(IG,K)+8.*SR4(IG,K)-7.*SR5(IG,K)+2.*SR6(IG,K))	SYNT0454
X	*HR(K)/(PO7(2,K)*PH(1,K)*DO(2,K))	SYNT0455
	T(1,6)=(-SR3(IG,K)+3.*SR4(IG,K)-3.*SR5(IG,K)+SR6(IG,K))	SYNT0456
X	*HR(K)**2/(PO7(2,K)*PH(1,K)*DO(2,K)*DH(1,K))	SYNT0457
	FT(1,1,3)=(KD2(IG,J)-KD3(IG,J)-3.*KD4(IG,J)+5.*KD5(IG,J)	SYNT0458
X	-2.*KD6(IG,J))*HR(J)/(PH7(2,J)*PC(1,J)*DH(2,J))	SYNT0459
	FT(1,2,3)=(-KD3(IG,J)+3.*KD4(IG,J)-3.*KD5(IG,J)+KD6(IG,J))	SYNT0460
X	*HR(J)**2/(PH7(2,J)*PO(1,J)*DO(1,J)*DH(2,J))	SYNT0461
	FT(1,3,3)=(3.*KD4(IG,J)-5.*KD5(IG,J)+2.*KD6(IG,J))	SYNT0462
X	*HR(J)/(PH7(2,J)*PH(1,J)*DH(2,J))	SYNT0463
X	+(-KD1(IG,K)+2.*KD2(IG,K)+2.*KD3(IG,K)	SYNT0464
X	-8.*KD4(IG,K)+7.*KD5(IG,K)-2.*KD6(IG,K))*HR(K)	SYNT0465
X	/(PO7(2,K)*PO(1,K)*DO(2,K))	SYNT0466
	FT(1,4,3)=(KD4(IG,J)-2.*KD5(IG,J)+KD6(IG,J))	SYNT0467
X	*HR(J)**2/(PH7(2,J)*PH(1,J)*DH(2,J)*DH(1,J))	SYNT0468

X	+(KD2(IG,K)-4.*KD3(IG,K)+6.*KD4(IG,K)	SYNT0469
X	-4.*KD5(IG,K)+KD6(IG,K))*HR(K)**2	SYNT0470
X	/(PO7(2,K)*PO(1,K)*DO(2,K)*DO(1,K))	SYNT0471
	FT(1,5,3)=(-3.*KD3(IG,K)+8.*KD4(IG,K)-7.*KD5(IG,K)+2.*KD6(IG,K))	SYNT0472
X	*HR(K)/(PO7(2,K)*PH(1,K)*DO(2,K))	SYNT0473
	FT(1,6,3)=(-KD3(IG,K)+3.*KD4(IG,K)-3.*KD5(IG,K)+KD6(IG,K))	SYNT0474
X	*HR(K)**2/(PO7(2,K)*PH(1,K)*DO(2,K)*DH(1,K))	SYNT0475
	FT(1,1,4)=(KC2(IG,J)-KC3(IG,J)-3.*KC4(IG,J)+5.*KC5(IG,J)	SYNT0476
X	-2.*KC6(IG,J))*HR(J)/(PH7(1,J)*PC(2,J)*DH(1,J))	SYNT0477
	FT(1,2,4)=(-KC3(IG,J)+3.*KC4(IG,J)-3.*KC5(IG,J)+KC6(IG,J))	SYNT0478
X	*HR(J)**2/(PH7(1,J)*PO(2,J)*DO(2,J)*DH(1,J))	SYNT0479
	FT(1,3,4)=(3.*KC4(IG,J)-5.*KC5(IG,J)+2.*KC6(IG,J))	SYNT0480
X	*HR(J)/(PH7(1,J)*PH(2,J)*DH(1,J))	SYNT0481
X	+(-KC1(IG,K)+2.*KC2(IG,K)+2.*KC3(IG,K)	SYNT0482
X	-8.*KC4(IG,K)+7.*KC5(IG,K)-2.*KC6(IG,K))*HR(K)	SYNT0483
X	/(PO7(1,K)*PO(2,K)*DO(1,K))	SYNT0484
	FT(1,4,4)=(KC4(IG,J)-2.*KC5(IG,J)+KC6(IG,J))	SYNT0485
X	*HR(J)**2/(PH7(1,J)*PH(2,J)*DH(1,J)*DH(2,J))	SYNT0486
X	+(KC2(IG,K)-4.*KC3(IG,K)+6.*KC4(IG,K)	SYNT0487
X	-4.*KC5(IG,K)+KC6(IG,K))*HR(K)**2	SYNT0488
X	/(PO7(1,K)*PO(2,K)*DO(1,K)*DO(2,K))	SYNT0489
	FT(1,5,4)=(-3.*KC3(IG,K)+8.*KC4(IG,K)-7.*KC5(IG,K)+2.*KC6(IG,K))	SYNT0490
X	*HR(K)/(PO7(1,K)*PH(2,K)*DO(1,K))	SYNT0491
	FT(1,6,4)=(-KC3(IG,K)+3.*KC4(IG,K)-3.*KC5(IG,K)+KC6(IG,K))	SYNT0492
X	*HR(K)**2/(PO7(1,K)*PH(2,K)*DO(1,K)*DH(2,K))	SYNT0493
60	CONTINUE	SYNT0494
	IF (IBC.EQ.4) GO TO 63	SYNT0495
C	ZERO FLUX COEFFICIENTS ON THE LEFT:	SYNT0496
	DO 61 IG=1,2	SYNT0497
	V2=1./(PO7(IG,1)*PO(IG,1))	SYNT0498
	V3=1./(PO7(IG,1)*PH(IG,1))	SYNT0499
	LT(1,4,IG)=(KA2(IG,1)-4.*KA3(IG,1)+6.*KA4(IG,1)-4.*KA5(IG,1)	SYNT0500
X	+KA6(IG,1)-(LA2(IG,1)-4.*LA3(IG,1)+6.*LA4(IG,1)-4.*LA5(IG,1)	SYNT0501
X	+LA6(IG,1))-P1(IG,1)+6.*P2(IG,1)-12.*P3(IG,1)+10.*P4(IG,1)	SYNT0502
X	-3.*P5(IG,1)-(-Q1(IG,1)+6.*Q2(IG,1)-12.*Q3(IG,1)+10.*Q4(IG,1)	SYNT0503
X	-3.*Q5(IG,1))+R0(IG,1)-8.*R1(IG,1)+22.*R2(IG,1)-24.*R3(IG,1)	SYNT0504

X	+9.*R4(IG,1))*V2*(HR(1)/DO(IG,1))*2	SYNT0505
LT(1,5,IG)	=(-3.*KA3(IG,1)+8.*KA4(IG,1)-7.*KA5(IG,1)+2.*KA6(IG,1)	SYNT0506
X	-(-3.*LA3(IG,1)+8.*LA4(IG,1)-7.*LA5(IG,1)+2.*LA6(IG,1))	SYNT0507
X	+3.*P2(IG,1)-14.*P3(IG,1)+17.*P4(IG,1)-6.*P5(IG,1)+(-6.*	SYNT0508
X	Q2(IG,1)+18.*Q3(IG,1)-18.*Q4(IG,1)+6.*Q5(IG,1))-6.*R1(IG,1)	SYNT0509
X	-30.*R2(IG,1)+42.*R3(IG,1)-18.*R4(IG,1))*V3*HR(1)/DO(IG,1)	SYNT0510
LT(1,6,IG)	=(-KA3(IG,1)+3.*KA4(IG,1)-3.*KA5(IG,1)+KA6(IG,1)	SYNT0511
X	-(-LA3(IG,1)+3.*LA4(IG,1)-3.*LA5(IG,1)+LA6(IG,1))+P2(IG,1)	SYNT0512
X	-5.*P3(IG,1)+7.*P4(IG,1)-3.*P5(IG,1)-(2.*Q2(IG,1)-7.*Q3(IG,1)	SYNT0513
X	+8.*Q4(IG,1)-3.*Q5(IG,1))-2.*R1(IG,1)+11.*R2(IG,1)-18.*R3(IG,1)	SYNT0514
X	+9.*R4(IG,1))*V3*HR(1)*HR(1)/(DO(IG,1)*DH(IG,1))	SYNT0515
FT(1,4,IG)	=(KB2(IG,1)-4.*KB3(IG,1)+6.*KB4(IG,1)-4.*KB5(IG,1)	SYNT0516
X	+KB6(IG,1))*V2*(HR(1)/DO(IG,1))*2	SYNT0517
FT(1,5,IG)	=(-3.*KB3(IG,1)+8.*KB4(IG,1)-7.*KB5(IG,1)+2.*KB6(IG,1))	SYNT0518
X	*V3*HR(1)/DO(IG,1)	SYNT0519
FT(1,6,IG)	=(-KB3(IG,1)+3.*KB4(IG,1)-3.*KB5(IG,1)+KB6(IG,1))	SYNT0520
X	*V3*HR(1)*HR(1)/(DO(IG,1)*DH(IG,1))	SYNT0521
IF (IG.EQ.2)	GO TO 61	SYNT0522
T(1,4)	=(SR2(IG,1)-4.*SR3(IG,1)+6.*SR4(IG,1)-4.*SR5(IG,1)	SYNT0523
X	+SR6(IG,1))*HR(1)**2/(PO7(2,1)*PO(1,1)*DO(2,1)*DO(1,1))	SYNT0524
T(1,5)	=(-3.*SR3(IG,1)+8.*SR4(IG,1)-7.*SR5(IG,1)+2.*SR6(IG,1))	SYNT0525
X	*HR(1)/(PO7(2,1)*PH(1,1)*DO(2,1))	SYNT0526
T(1,6)	=(-SR3(IG,1)+3.*SR4(IG,1)-3.*SR5(IG,1)+SR6(IG,1))	SYNT0527
X	*HR(1)**2/(PO7(2,1)*PH(1,1)*DO(2,1)*DH(1,1))	SYNT0528
FT(1,4,3)	=(KD2(IG,1)-4.*KD3(IG,1)+6.*KD4(IG,1)-4.*KD5(IG,1)	SYNT0529
X	+KD6(IG,1))*HR(1)**2/(PO7(2,1)*PO(1,1)*DO(2,1)*DO(1,1))	SYNT0530
FT(1,5,3)	=(-3.*KD3(IG,1)+8.*KD4(IG,1)-7.*KD5(IG,1)+2.*KD6(IG,1))	SYNT0531
X	*HR(1)/(PO7(2,1)*PH(1,1)*DO(2,1))	SYNT0532
FT(1,6,3)	=(-KD3(IG,1)+3.*KD4(IG,1)-3.*KD5(IG,1)+KD6(IG,1))	SYNT0533
X	*HR(1)**2/(PO7(2,1)*PH(1,1)*DO(2,1)*DH(1,1))	SYNT0534
FT(1,4,4)	=(KC2(IG,1)-4.*KC3(IG,1)+6.*KC4(IG,1)-4.*KC5(IG,1)	SYNT0535
X	+KC6(IG,1))*HR(1)**2/(PO7(1,1)*PO(2,1)*DO(1,1)*DO(2,1))	SYNT0536
FT(1,5,4)	=(-3.*KC3(IG,1)+8.*KC4(IG,1)-7.*KC5(IG,1)+2.*KC6(IG,1))	SYNT0537
X	*HR(1)/(PO7(1,1)*PH(2,1)*DO(1,1))	SYNT0538
FT(1,6,4)	=(-KC3(IG,1)+3.*KC4(IG,1)-3.*KC5(IG,1)+KC6(IG,1))	SYNT0539
X	*HR(1)**2/(PO7(1,1)*PH(2,1)*DO(1,1)*DH(2,1))	SYNT0540

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61 CONTINUE
GO TO 65
C ZERO CURRENT COEFFICIENTS ON THE LEFT:
63 K=1
I=1
DO 64 IG=1,2
V2=1./(PO7(IG,K)*PO(IG,K))
V3=1./(PO7(IG,K)*PH(IG,K))
LT(I,4,IG)=
X (KA0(IG,K)-6.*KA2(IG,K)+4.*KA3(IG,K)+9.*KA4(IG,K)-12.*KA5(IG,K)
X +4.*KA6(IG,K)-(LA0(IG,K)-6.*LA2(IG,K)+4.*LA3(IG,K)+9.*LA4(IG,K)
X -12.*LA5(IG,K)+4.*LA6(IG,K))+6.*P1(IG,K)-6.*P2(IG,K)-18.*
X P3(IG,K)+30.*P4(IG,K)-12.*P5(IG,K)-(6.*Q1(IG,K)-6.*Q2(IG,K)
X -18.*Q3(IG,K)+30.*Q4(IG,K)-12.*Q5(IG,K))+36.*R2(IG,K)-72.*
X R3(IG,K)+36.*R4(IG,K))*V2 -CO(IG)/PO(IG,1)
LT(I,5,IG)=(3.*KA2(IG,K)-2.*KA3(IG,K)-9.*KA4(IG,K)+12.*KA5(IG,K)
X -4.*KA6(IG,K)-(3.*LA2(IG,K)-2.*LA3(IG,K)-9.*LA4(IG,K)+12.*
X LA5(IG,K)-4.*LA6(IG,K))+18.*P3(IG,K)-30.*P4(IG,K)+12.*P5(IG,K)
X +(6.*Q1(IG,K)-6.*Q2(IG,K)-18.*Q3(IG,K)+30.*Q4(IG,K)-12.*
X Q5(IG,K))-(36.*R2(IG,K)-72.*R3(IG,K)+36.*R4(IG,K))*V3
LT(I,6,IG)=(KA2(IG,K)-KA3(IG,K)-3.*KA4(IG,K)+5.*KA5(IG,K)-2.*
X KA6(IG,K)-(LA2(IG,K)-LA3(IG,K)-3.*LA4(IG,K)+5.*LA5(IG,K)-2.*
X LA6(IG,K))+6.*P3(IG,K)-12.*P4(IG,K)+6.*P5(IG,K)-(-2.*Q1(IG,K)+
X 3.*Q2(IG,K)+6.*Q3(IG,K)-13.*Q4(IG,K)+6.*Q5(IG,K))-12.*R2(IG,K)
X +30.*R3(IG,K)-18.*R4(IG,K))*V3*HR(K)/DH(IG,K)
FT(I,4,IG)=
X +(KB0(IG,K)-6.*KB2(IG,K)+4.*KB3(IG,K)+9.*KB4(IG,K)
X -12.*KB5(IG,K)+4.*KB6(IG,K))*V2
FT(I,5,IG)=(3.*KB2(IG,K)-2.*KB3(IG,K)-9.*KB4(IG,K)+12.*KB5(IG,K)
X -4.*KB6(IG,K))*V3
FT(I,6,IG)=(KB2(IG,K)-KB3(IG,K)-3.*KB4(IG,K)+5.*KB5(IG,K)
X -2.*KB6(IG,K))*V3*HR(K)/DH(IG,K)
IF (IG.EQ.2) GO TO 64
T(I,4)=
X +(SR0(IG,K)-6.*SR2(IG,K)+4.*SR3(IG,K)+9.*SR4(IG,K)
X -12.*SR5(IG,K)+4.*SR6(IG,K))/(PO7(2,K)*PO(1,K))

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SYNT0541
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SYNT0576

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T(I,5)	=(3.*SR2(IG,K)-2.*SR3(IG,K)-9.*SR4(IG,K)+12.*SR5(IG,K)	SYNT0577
X	-4.*SR6(IG,K))/(PO7(2,K)*PH(1,K))	SYNT0578
T(I,6)	=(SR2(IG,K)-SR3(IG,K)-3.*SR4(IG,K)+5.*SR5(IG,K)	SYNT0579
X	-2.*SR6(IG,K))*HR(K)/(PO7(2,K)*PH(1,K)*DH(1,K))	SYNT0580
FT(I,4,3)=		SYNT0581
X	+(KDO(IG,K)-6.*KD2(IG,K)+4.*KD3(IG,K)+9.*KD4(IG,K)	SYNT0582
X	-12.*KD5(IG,K)+4.*KD6(IG,K))/(PO7(2,K)*PO(1,K))	SYNT0583
FT(I,5,3)=	(3.*KD2(IG,K)-2.*KD3(IG,K)-9.*KD4(IG,K)+12.*KD5(IG,K)	SYNT0584
X	-4.*KD6(IG,K))/(PO7(2,K)*PH(1,K))	SYNT0585
FT(I,6,3)=	(KD2(IG,K)-KD3(IG,K)-3.*KD4(IG,K)+5.*KD5(IG,K)	SYNT0586
X	-2.*KD6(IG,K))*HR(K)/(PO7(2,K)*PH(1,K)*DH(1,K))	SYNT0587
FT(I,4,4)=		SYNT0588
X	+(KCO(IG,K)-6.*KC2(IG,K)+4.*KC3(IG,K)+9.*KC4(IG,K)	SYNT0589
X	-12.*KC5(IG,K)+4.*KC6(IG,K))/(PO7(1,K)*PO(2,K))	SYNT0590
FT(I,5,4)=	(3.*KC2(IG,K)-2.*KC3(IG,K)-9.*KC4(IG,K)+12.*KC5(IG,K)	SYNT0591
X	-4.*KC6(IG,K))/(PO7(1,K)*PH(2,K))	SYNT0592
FT(I,6,4)=	(KC2(IG,K)-KC3(IG,K)-3.*KC4(IG,K)+5.*KC5(IG,K)	SYNT0593
X	-2.*KC6(IG,K))*HR(K)/(PO7(1,K)*PH(2,K)*DH(2,K))	SYNT0594
64 CONTINUE		SYNT0595
C	FIX UP THE TWO FIRST COLUMN ENTRIES TO MATCH F(1) (NOT G(1)):	SYNT0596
	DO 70 IG=1,2	SYNT0597
	DO 70 I=2,3	SYNT0598
	LT(I,2,IG)=LT(I,1,IG)	SYNT0599
	FT(I,2,IG)=FT(I,1,IG)	SYNT0600
	IF (IG.EQ.2) GO TO 70	SYNT0601
	T(I,2)=T(I,1)	SYNT0602
	FT(I,2,3)=FT(I,1,3)	SYNT0603
	FT(I,2,4)=FT(I,1,4)	SYNT0604
70 CONTINUE		SYNT0605
65	I=2*KR	SYNT0606
	J=KR	SYNT0607
	K=KR	SYNT0608
	IF (IBC.EQ.2.OR.IBC.EQ.4) GO TO 74	SYNT0609
C	ZERO FLUX COEFFICIENTS ON THE RIGHT:	SYNT0610
	DO 72 IG=1,2	SYNT0611
	V=1./(PH7(IG,J)*PO(IG,J))	SYNT0612



VI=1./(PH7(IG,J)*PH(IG,J))	SYNT0613
LT(I,1,IG)=(KA2(IG,J)-KA3(IG,J)-3.*KA4(IG,J)+5.*KA5(IG,J)-2.*	SYNT0614
X KA6(IG,J)-(LA2(IG,J)-LA3(IG,J)-3.*LA4(IG,J)+5.*LA5(IG,J)-2.*	SYNT0615
X LA6(IG,J))-2.*P1(IG,J)+3.*P2(IG,J)+6.*P3(IG,J)-13.*P4(IG,J)+6.*	SYNT0616
X P5(IG,J)-(6.*Q3(IG,J)-12.*Q4(IG,J)+6.*Q5(IG,J))-12.*R2(IG,J)	SYNT0617
X +30.*R3(IG,J)-18.*R4(IG,J))*V*HR(J)/DH(IG,J)	SYNT0618
LT(I,2,IG)=(-KA3(IG,J)+3.*KA4(IG,J)-3.*KA5(IG,J)+KA6(IG,J)	SYNT0619
X -(-LA3(IG,J)+3.*LA4(IG,J)-3.*LA5(IG,J)+LA6(IG,J))+2.*P2(IG,J)	SYNT0620
X -7.*P3(IG,J)+8.*P4(IG,J)-3.*P5(IG,J)-(Q2(IG,J)-5.*Q3(IG,J)	SYNT0621
X +7.*Q4(IG,J)-3.*Q5(IG,J))-2.*R1(IG,J)+11.*R2(IG,J)-18.*	SYNT0622
X R3(IG,J)+9.*R4(IG,J))*V*HR(J)*HR(J)/(DO(IG,J)*DH(IG,J))	SYNT0623
LT(I,3,IG)=(KA4(IG,J)-2.*KA5(IG,J)+KA6(IG,J)-(LA4(IG,J)-2.*	SYNT0624
X LA5(IG,J)+LA6(IG,J))-2.*P3(IG,J)+5.*P4(IG,J)-3.*P5(IG,J)-(-2.*	SYNT0625
X Q3(IG,J)+5.*Q4(IG,J)-3.*Q5(IG,J))+4.*R2(IG,J)-12.*R3(IG,J)+9.*	SYNT0626
X R4(IG,J))*V1*(HR(J)/DH(IG,J))**2	SYNT0627
FT(I,1,IG)=(KB2(IG,J)-KB3(IG,J)-3.*KB4(IG,J)+5.*KB5(IG,J)	SYNT0628
X -2.*KB6(IG,J))*V*HR(J)/DH(IG,J)	SYNT0629
FT(I,2,IG)=(-KB3(IG,J)+3.*KB4(IG,J)-3.*KB5(IG,J)+KB6(IG,J))	SYNT0630
X *V*HR(J)*HR(J)/(DO(IG,J)*DH(IG,J))	SYNT0631
FT(I,3,IG)=(KB4(IG,J)-2.*KB5(IG,J)+KB6(IG,J))	SYNT0632
X *V1*(HR(J)/DH(IG,J))**2	SYNT0633
IF (IG.EQ.2) GO TO 72	SYNT0634
T(I,1)=(SR2(IG,J)-SR3(IG,J)-3.*SR4(IG,J)+5.*SR5(IG,J)	SYNT0635
X -2.*SR6(IG,J))*HR(J)/(PH7(2,J)*PC(1,J)*DH(2,J))	SYNT0636
T(I,2)=(-SR3(IG,J)+3.*SR4(IG,J)-3.*SR5(IG,J)+SR6(IG,J))	SYNT0637
X *HR(J)**2/(PH7(2,J)*PO(1,J)*DH(2,J)*DO(1,J))	SYNT0638
T(I,3)=(SR4(IG,J)-2.*SR5(IG,J)+SR6(IG,J))	SYNT0639
X *HR(J)**2/(PH7(2,J)*PH(1,J)*DH(2,J)*DH(1,J))	SYNT0640
FT(I,1,3)=(KD2(IG,J)-KD3(IG,J)-3.*KD4(IG,J)+5.*KD5(IG,J)	SYNT0641
X -2.*KD6(IG,J))*HR(J)/(PH7(2,J)*PC(1,J)*DH(2,J))	SYNT0642
FT(I,2,3)=(-KD3(IG,J)+3.*KD4(IG,J)-3.*KD5(IG,J)+KD6(IG,J))	SYNT0643
X *HR(J)**2/(PH7(2,J)*PO(1,J)*DH(2,J)*DO(1,J))	SYNT0644
FT(I,3,3)=(KD4(IG,J)-2.*KD5(IG,J)+KD6(IG,J))	SYNT0645
X *HR(J)**2/(PH7(2,J)*PH(1,J)*DH(2,J)*DH(1,J))	SYNT0646
FT(I,1,4)=(KC2(IG,J)-KC3(IG,J)-3.*KC4(IG,J)+5.*KC5(IG,J)	SYNT0647
X -2.*KC6(IG,J))*HR(J)/(PH7(1,J)*PC(2,J)*DH(1,J))	SYNT0648

	FT(I,2,4)={-KC3(IG,J)+3.*KC4(IG,J)-3.*KC5(IG,J)+KC6(IG,J)}	SYNT0649
X	*HR(J)**2/(PH7(1,J)*PD(2,J)*DH(1,J)*DO(2,J))	SYNT0650
	FT(I,3,4)={KC4(IG,J)-2.*KC5(IG,J)+KC6(IG,J)}	SYNT0651
X	*HR(J)**2/(PH7(1,J)*PH(2,J)*DH(1,J)*DH(2,J))	SYNT0652
	72 CONTINUE	SYNT0653
C	FIX UP THE LAST TWO COLUMNS TO MATCH G(K+1) (NOT F(K+1)):	SYNT0654
	I1=2*KR-2	SYNT0655
	I2=I1+1	SYNT0656
	DO 73 IG=1,2	SYNT0657
	DO 73 I=I1,I2	SYNT0658
	LT(I,5,IG)=LT(I,6,IG)	SYNT0659
	FT(I,5,IG)=FT(I,6,IG)	SYNT0660
	IF (IG.EQ.2) GO TO 73	SYNT0661
	T(I,5)=T(I,6)	SYNT0662
	FT(I,5,3)=FT(I,6,3)	SYNT0663
	FT(I,5,4)=FT(I,6,4)	SYNT0664
	73 CONTINUE	SYNT0665
	GO TO 80	SYNT0666
C	ZERO CURRENT COEFFICIENTS ON THE RIGHT:	SYNT0667
	74 DO 62 IG=1,2	SYNT0668
	V=1./(PH7(IG,K)*PO(IG,K))	SYNT0669
	V1=1./(PH7(IG,K)*PH(IG,K))	SYNT0670
	LT(I,1,IG)={3.*KA2(IG,K)-2.*KA3(IG,K)-9.*KA4(IG,K)+12.*KA5(IG,K)	SYNT0671
X	-4.*KA6(IG,K)-(3.*LA2(IG,K)-2.*LA3(IG,K)-9.*LA4(IG,K)+12.*	SYNT0672
X	LA5(IG,K)-4.*LA6(IG,K))-(6.*P1(IG,K)-6.*P2(IG,K)-18.*P3(IG,K)	SYNT0673
X	+30.*P4(IG,K)-12.*P5(IG,K))-(18.*Q3(IG,K)-30.*Q4(IG,K)+12.*	SYNT0674
X	Q5(IG,K))-(36.*R2(IG,K)-72.*R3(IG,K)+36.*R4(IG,K))*V	SYNT0675
	LT(I,2,IG)={-3.*KA3(IG,K)+8.*KA4(IG,K)-7.*KA5(IG,K)+2.*KA6(IG,K)	SYNT0676
X	-(-3.*LA3(IG,K)+8.*LA4(IG,K)-7.*LA5(IG,K)+2.*LA6(IG,K))-(-6.*	SYNT0677
X	P2(IG,K)+18.*P3(IG,K)-18.*P4(IG,K)+6.*P5(IG,K))-(3.*Q2(IG,K)	SYNT0678
X	-14.*Q3(IG,K)+17.*Q4(IG,K)-6.*Q5(IG,K))-(6.*R1(IG,K)-30.*	SYNT0679
X	R2(IG,K)+42.*R3(IG,K)-18.*R4(IG,K))*V*HR(K)/DO(IG,K)	SYNT0680
	LT(I,3,IG)={9.*KA4(IG,K)-12.*KA5(IG,K)+4.*KA6(IG,K)-(9.*LA4(IG,K)	SYNT0681
X	-12.*LA5(IG,K)+4.*LA6(IG,K))-(18.*P3(IG,K)-30.*P4(IG,K)+12.*	SYNT0682
X	P5(IG,K))+(18.*Q3(IG,K)-30.*Q4(IG,K)+12.*Q5(IG,K))	SYNT0683
X	+36.*R2(IG,K)-72.*R3(IG,K)+36.*R4(IG,K))*V1 +CH(IG)/PH(IG,KR)	SYNT0684

	FT(1,1,IG)=(3.*KB2(IG,K)-2.*KB3(IG,K)-9.*KB4(IG,K)+12.*KB5(IG,K)	SYNT0685
X	-4.*KB6(IG,K))*V	SYNT0686
	FT(1,2,IG)=(-3.*KB3(IG,K)+8.*KB4(IG,K)-7.*KB5(IG,K)+2.*KB6(IG,K))	SYNT0687
X	*V*HR(K)/DO(IG,K)	SYNT0688
	FT(1,3,IG)=(9.*KB4(IG,K)-12.*KB5(IG,K)+4.*KB6(IG,K))*V1	SYNT0689
	IF (IG.EQ.2) GO TO 62	SYNT0690
	T(1,1)=(3.*SR2(IG,K)-2.*SR3(IG,K)-9.*SR4(IG,K)+12.*SR5(IG,K)	SYNT0691
X	-4.*SR6(IG,K))/(PH7(2,K)*PO(1,K))	SYNT0692
	T(1,2)=(-3.*SR3(IG,K)+8.*SR4(IG,K)-7.*SR5(IG,K)+2.*SR6(IG,K))	SYNT0693
X	*HR(K)/(PH7(2,K)*PO(1,K)*DO(1,K))	SYNT0694
	T(1,3)=(9.*SR4(IG,K)-12.*SR5(IG,K)+4.*SR6(IG,K))	SYNT0695
X	/(PH7(2,K)*PH(1,K))	SYNT0696
	FT(1,1,3)=(3.*KD2(IG,K)-2.*KD3(IG,K)-9.*KD4(IG,K)+12.*KD5(IG,K)	SYNT0697
X	-4.*KD6(IG,K))/(PH7(2,K)*PO(1,K))	SYNT0698
	FT(1,2,3)=(-3.*KD3(IG,K)+8.*KD4(IG,K)-7.*KD5(IG,K)+2.*KD6(IG,K))	SYNT0699
X	*HR(K)/(PH7(2,K)*PO(1,K)*DO(1,K))	SYNT0700
	FT(1,3,3)=(9.*KD4(IG,K)-12.*KD5(IG,K)+4.*KD6(IG,K))	SYNT0701
X	/(PH7(2,K)*PH(1,K))	SYNT0702
	FT(1,1,4)=(3.*KC2(IG,K)-2.*KC3(IG,K)-9.*KC4(IG,K)+12.*KC5(IG,K)	SYNT0703
X	-4.*KC6(IG,K))/(PH7(1,K)*PO(2,K))	SYNT0704
	FT(1,2,4)=(-3.*KC3(IG,K)+8.*KC4(IG,K)-7.*KC5(IG,K)+2.*KC6(IG,K))	SYNT0705
X	*HR(K)/(PH7(1,K)*PO(2,K)*DO(2,K))	SYNT0706
	FT(1,3,4)=(9.*KC4(IG,K)-12.*KC5(IG,K)+4.*KC6(IG,K))	SYNT0707
X	/(PH7(1,K)*PH(2,K))	SYNT0708
62	CONTINUE	SYNT0709
80	CONTINUE	SYNT0710
C	INCLUDE THETA AND PHI (PHI = -1) IN THE MATRIX FORMATIONS:	SYNT0711
C	THE ABOVE EQUATIONS ARE DERIVED USING PHI = -1.	SYNT0712
	PHIPHI=+1.000	SYNT0713
	IF (THETA.NE.1.0) CALL MATFIX(THETA,PHIPHI)	SYNT0714
C	TO PRINT OUT THE /B3/ MATRICES:	SYNT0715
	IF (ISEE.GE.2) CALL PRTOUT(2)	SYNT0716
	RETURN	SYNT0717
	END	SYNT0718

<pre> SUBROUTINE ERROR(I,J) ANNOUNCES INPUT ERRORS AND TERMINATES PROGRAM EXECUTION: GO TO (1,2,3,4,5,6,7,8,9),I 1 WRITE (6,101) GO TO 10 2 WRITE (6,102) J GO TO 10 3 WRITE (6,103) J GO TO 10 4 WRITE (6,104) J GO TO 10 5 WRITE (6,105) J GO TO 10 6 WRITE (6,106) J GO TO 10 7 CONTINUE 8 CONTINUE 9 CONTINUE 10 WRITE (6,110) 101 FORMAT ('MUST HAVE &gt; 2 SUBREGIONS FOR ZERO FLUX B.C.S. INVALID.') 102 FORMAT ('NUMBER OF SUBREGIONS =',I3,' &gt; 25. INVALID.') 103 FORMAT ('SUBREGION NUMBER',I3,' HAS &gt; 100 SECTIONS. INVALID.') 104 FORMAT ('INPUT ERROR IN REGION SEQUENCING AT REGION',I5,'.') 105 FORMAT ('IZ(I) = 0. IN REGION I =',I3,'. INVALID.') 106 FORMAT ('THE TOLERANCE: EPS',I1,' IS &lt; 1.0E-16. INVALID.') 107 FORMAT ('BOUNDARY CONDITION OPTION =',I2,' &lt; 1 OR &gt; 4. INVALID.') 110 FORMAT (1H0,'PROBLEM TERMINATED.') CALL EXIT RETURN END </pre>	<pre> ERR00001 ERR00002 ERR00003 ERR00004 ERR00005 ERR00006 ERR00007 ERR00008 ERR00009 ERR00010 ERR00011 ERR00012 ERR00013 ERR00014 ERR00015 ERR00016 ERR00017 ERR00018 ERR00019 ERR00020 ERR00021 ERR00022 ERR00023 ERR00024 ERR00025 ERR00026 ERR00027 ERR00028 ERR00029 ERR00030 </pre>
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<pre> SUBROUTINE REPEAT(K,L) C   SETS THE /B5/ ARRAYS (K) EQUAL TO PAST STORED ARRAYS (L):       IMPLICIT REAL*8 (A-Z)       COMMON /B5/ X     KA0(2,25),KA1(2,25),KA2(2,25),KA3(2,25),KA4(2,25),KA5(2,25), X     KA6(2,25),KB0(2,25),KB1(2,25),KB2(2,25),KB3(2,25),KB4(2,25), X     KB5(2,25),KB6(2,25),LA0(2,25),LA1(2,25),LA2(2,25),LA3(2,25), X     LA4(2,25),LA5(2,25),LA6(2,25),P0(2,25) ,P1(2,25) ,P2(2,25) , X     P3(2,25) ,P4(2,25) ,P5(2,25) ,P6(2,25) ,Q0(2,25) ,Q1(2,25) , X     Q2(2,25) ,Q3(2,25) ,Q4(2,25) ,Q5(2,25) ,Q6(2,25) ,R0(2,25) , X     R1(2,25) ,R2(2,25) ,R3(2,25) ,R4(2,25) ,SR0(1,25),SR1(1,25), X     SR2(1,25),SR3(1,25),SR4(1,25),SR5(1,25),SR6(1,25),KC0(1,25), X     KC1(1,25),KC2(1,25),KC3(1,25),KC4(1,25),KC5(1,25),KC6(1,25), X     KD0(1,25),KD1(1,25),KD2(1,25),KD3(1,25),KD4(1,25),KD5(1,25), X     KD6(1,25), X     PO(2,25) ,PH(2,25) ,PO7(2,25),PH7(2,25),D0(2,25) ,DH(2,25)       COMMON /XAXIS/ HX,HR(25)       INTEGER K,L,G       HR(K)=HR(L)       HX=HX+HR(K)       DO 10 G=1,2       KA0(G,K)=KA0(G,L)       KA1(G,K)=KA1(G,L)       KA2(G,K)=KA2(G,L)       KA3(G,K)=KA3(G,L)       KA4(G,K)=KA4(G,L)       KA5(G,K)=KA5(G,L)       KA6(G,K)=KA6(G,L)       KB0(G,K)=KB0(G,L)       KB1(G,K)=KB1(G,L)       KB2(G,K)=KB2(G,L)       KB3(G,K)=KB3(G,L)       KB4(G,K)=KB4(G,L)       KB5(G,K)=KB5(G,L)       KB6(G,K)=KB6(G,L)       LA0(G,K)=LA0(G,L) </pre>	<pre> REPE0001 REPE0002 REPE0003 REPE0004 REPE0005 REPE0006 REPE0007 REPE0008 REPE0009 REPE0010 REPE0011 REPE0012 REPE0013 REPE0014 REPE0015 REPE0016 REPE0017 REPE0018 REPE0019 REPE0020 REPE0021 REPE0022 REPE0023 REPE0024 REPE0025 REPE0026 REPE0027 REPE0028 REPE0029 REPE0030 REPE0031 REPE0032 REPE0033 REPE0034 REPE0035 REPE0036 </pre>
--	--

LA1(G,K)=LA1(G,L)  
LA2(G,K)=LA2(G,L)  
LA3(G,K)=LA3(G,L)  
LA4(G,K)=LA4(G,L)  
LA5(G,K)=LA5(G,L)  
LA6(G,K)=LA6(G,L)  
P0(G,K)=P0(G,L)  
P1(G,K)=P1(G,L)  
P2(G,K)=P2(G,L)  
P3(G,K)=P3(G,L)  
P4(G,K)=P4(G,L)  
P5(G,K)=P5(G,L)  
P6(G,K)=P6(G,L)  
Q0(G,K)=Q0(G,L)  
Q1(G,K)=Q1(G,L)  
Q2(G,K)=Q2(G,L)  
Q3(G,K)=Q3(G,L)  
Q4(G,K)=Q4(G,L)  
Q5(G,K)=Q5(G,L)  
Q6(G,K)=Q6(G,L)  
R0(G,K)=R0(G,L)  
R1(G,K)=R1(G,L)  
R2(G,K)=R2(G,L)  
R3(G,K)=R3(G,L)  
R4(G,K)=R4(G,L)  
IF (G.EQ.2) GO TO 5  
SR0(G,K)=SR0(G,L)  
SR1(G,K)=SR1(G,L)  
SR2(G,K)=SR2(G,L)  
SR3(G,K)=SR3(G,L)  
SR4(G,K)=SR4(G,L)  
SR5(G,K)=SR5(G,L)  
SR6(G,K)=SR6(G,L)  
KC0(G,K)=KC0(G,L)  
KC1(G,K)=KC1(G,L)  
KC2(G,K)=KC2(G,L)

REPE0037  
REPE0038  
REPE0039  
REPE0040  
REPE0041  
REPE0042  
REPE0043  
REPE0044  
REPE0045  
REPE0046  
REPE0047  
REPE0048  
REPE0049  
REPE0050  
REPE0051  
REPE0052  
REPE0053  
REPE0054  
REPE0055  
REPE0056  
REPE0057  
REPE0058  
REPE0059  
REPE0060  
REPE0061  
REPE0062  
REPE0063  
REPE0064  
REPE0065  
REPE0066  
REPE0067  
REPE0068  
REPE0069  
REPE0070  
REPE0071  
REPE0072

KC3(G,K)=KC3(G,L)  
KC4(G,K)=KC4(G,L)  
KC5(G,K)=KC5(G,L)  
KC6(G,K)=KC6(G,L)  
KD0(G,K)=KD0(G,L)  
KD1(G,K)=KD1(G,L)  
KD2(G,K)=KD2(G,L)  
KD3(G,K)=KD3(G,L)  
KD4(G,K)=KD4(G,L)  
KD5(G,K)=KD5(G,L)  
KD6(G,K)=KD6(G,L)  
5 CONTINUE  
PO(G,K)=PO(G,L)  
PO7(G,K)=PO7(G,L)  
PH(G,K)=PH(G,L)  
PH7(G,K)=PH7(G,L)  
DO(G,K)=DO(G,L)  
DH(G,K)=DH(G,L)  
10 CONTINUE  
RETURN  
END

REPE0073  
REPE0074  
REPE0075  
REPE0076  
REPE0077  
REPE0078  
REPE0079  
REPE0080  
REPE0081  
REPE0082  
REPE0083  
REPE0084  
REPE0085  
REPE0086  
REPE0087  
REPE0088  
REPE0088  
REPE0089  
REPE0090  
REPE0091  
REPE0092  
REPE0093

```
C      SUBROUTINE BHSET(K)
        SETS UP THE /BH/ ARRAYS FOR GIF:
        IMPLICIT REAL*8 (A-H,L-Z)
        COMMON /BH/ X(101), H(101), Z(101)
        DO 1 I=1,K
1      Z(I)=X(I)-X(1)
        RETURN
        END
```

```
BHSE0001
BHSE0002
BHSE0003
BHSE0004
BHSE0005
BHSE0006
BHSE0007
BHSE0008
```



C  
C  
C

DOUBLE PRECISION FUNCTION GIF(N,G1,F,G2,C,G,K,ITC)

\* IDENTICAL TO SUBROUTINE GIF PREVIOUSLY LISTED IN PROGRAM LINEAR.

RETURN  
END

GIF 0001  
GIF 0002  
GIF 0003  
GIF 0004  
GIF 0005  
GIF 0006

```
C    DOUBLE PRECISION FUNCTION FACT(N)
      COMPUTES N FACTORIAL:
      FACT=1.0D0
      IF (N.LE.1) RETURN
      DO 1 I=2,N
1     FACT=FACT*DFLOAT(I)
      RETURN
      END
```

```
FACT0001
FACT0002
FACT0003
FACT0004
FACT0005
FACT0006
FACT0007
FACT0008
```

```

SUBROUTINE MATFIX(THEATA,PHI)
C   MODIFYS THE MATRIX ELEMENTS OF THE /B3/ MATRICES BY THEATA AND PHI.
C   PROPER CHOICE OF THEATA PROVIDES EASIER INVERSION OF THE MATRICES.
C   MATRIX SOLUTION SHOULD BE INDEPENDENT OF PHI. HOWEVER:
C   USE OF PHI > 0 RESULTS IN POSITIVE DEFINITE MATRICES.
COMMON /B1/ IBC
COMMON /B2/ KR,NN
COMMON /B3/ A(50,6,7)
REAL*8 A,THEATA,PHI,X,Y,Z
NO=NN-1
DO 10 L=1,7
DO 5 I=2,NO
X=THEATA
Y=1.000
IF (MOD(I,2).EQ.0) GO TO 1
X=THEATA*PHI
Y=PHI
1 DO 2 M=2,6,2
2 A(I,M,L)=X*A(I,M,L)
DO 3 M=1,5,2
3 A(I,M,L)=Y*A(I,M,L)
5 CONTINUE
C   BOUNDARY CONDITION EQUATIONS:
X=1.000
Y=1.000
Z=THEATA
IF (IBC.GT.2) GO TO 6
X=THEATA*PHI
Y=PHI
Z=X
6 A(1,4,L)=X*A(1,4,L)
A(1,5,L)=Y*A(1,5,L)
A(1,6,L)=Z*A(1,6,L)
X=1.000
Y=THEATA
Z=X

```

```

MATF0001
MATF0002
MATF0003
MATF0004
MATF0005
MATF0006
MATF0007
MATF0008
MATF0009
MATF0010
MATF0011
MATF0012
MATF0013
MATF0014
MATF0015
MATF0016
MATF0017
MATF0018
MATF0019
MATF0020
MATF0021
MATF0022
MATF0023
MATF0024
MATF0025
MATF0026
MATF0027
MATF0028
MATF0029
MATF0030
MATF0031
MATF0032
MATF0033
MATF0034
MATF0035
MATF0036

```

```
IF (IBC.EQ.2.OR.IBC.EQ.4) GO TO 7
X=PHI
Y=THEATA*PHI
Z=Y
7 A(NN,1,L)=X*A(NN,1,L)
  A(NN,2,L)=Y*A(NN,2,L)
  A(NN,3,L)=Z*A(NN,3,L)
10 CONTINUE
RETURN
END
```

```
MATF0037
MATF0038
MATF0039
MATF0040
MATF0041
MATF0042
MATF0043
MATF0044
MATF0045
MATF0046
```

```

SUBROUTINE PRTOUT(IP)
C   IP = 1 PRINTS OUT THE /B5/ ARRAYS USED IN MATRIX FORMATIONS.
C   IP = 2 PRINTS OUT THE /B3/ MATRICES GIVEN TO POWER.
IMPLICIT REAL*8 (A-H,K-Z)
COMMON /B2/ KR, N
COMMON /B3/ L1(50,6), L2(50,6), F1(50,6), F4(50,6), F3(50,6),
X   F2(50,6), T(50,6)
COMMON /B5/
X   KA0(2,25),KA1(2,25),KA2(2,25),KA3(2,25),KA4(2,25),KA5(2,25),
X   KA6(2,25),KB0(2,25),KB1(2,25),KB2(2,25),KB3(2,25),KB4(2,25),
X   KB5(2,25),KB6(2,25),LA0(2,25),LA1(2,25),LA2(2,25),LA3(2,25),
X   LA4(2,25),LA5(2,25),LA6(2,25),P0(2,25),P1(2,25),P2(2,25),
X   P3(2,25),P4(2,25),P5(2,25),P6(2,25),Q0(2,25),Q1(2,25),
X   Q2(2,25),Q3(2,25),Q4(2,25),Q5(2,25),Q6(2,25),R0(2,25),
X   R1(2,25),R2(2,25),R3(2,25),R4(2,25),SR0(1,25),SR1(1,25),
X   SR2(1,25),SR3(1,25),SR4(1,25),SR5(1,25),SR6(1,25),KC0(1,25),
X   KC1(1,25),KC2(1,25),KC3(1,25),KC4(1,25),KC5(1,25),KC6(1,25),
X   KD0(1,25),KD1(1,25),KD2(1,25),KD3(1,25),KD4(1,25),KD5(1,25),
X   KD6(1,25),
X   P0(2,25),PH(2,25),P07(2,25),PH7(2,25),D0(2,25),DH(2,25)
COMMON /XAXIS/ HX,HR(25)
INTEGER KR, G, N, K
GO TO (1001,1002), IP
C   KA'S:
1001 WRITE (6,10)
DO 11 G=1,2
11 WRITE (6,100) (G,K,KA0(G,K),KA1(G,K),KA2(G,K),KA3(G,K),KA4(G,K),
X   KA5(G,K),KA6(G,K),K=1,KR)
C   KB'S:
WRITE (6,20)
DO 21 G=1,2
21 WRITE (6,100) (G,K,KB0(G,K),KB1(G,K),KB2(G,K),KB3(G,K),KB4(G,K),
X   KB5(G,K),KB6(G,K),K=1,KR)
C   KC'S:
WRITE (6,22)
G=1

```

```

PRTO0001
PRTO0002
PRTO0003
PRTO0004
PRTO0005
PRTO0006
PRTO0007
PRTO0008
PRTO0009
PRTO0010
PRTO0011
PRTO0012
PRTO0013
PRTO0014
PRTO0015
PRTO0016
PRTO0017
PRTO0018
PRTO0019
PRTO0020
PRTO0021
PRTO0022
PRTO0023
PRTO0024
PRTO0025
PRTO0026
PRTO0027
PRTO0028
PRTO0029
PRTO0030
PRTO0031
PRTO0032
PRTO0033
PRTO0034
PRTO0035
PRTO0036

```

```

23 WRITE (6,100) (G,K,KC0(G,K),KC1(G,K),KC2(G,K),KC3(G,K),KC4(G,K),
X   KC5(G,K),KC6(G,K),K=1,KR)
C   KD'S:
    WRITE (6,24)
    G=1
25 WRITE (6,100) (G,K,KD0(G,K),KD1(G,K),KD2(G,K),KD3(G,K),KD4(G,K),
X   KD5(G,K),KD6(G,K),K=1,KR)
C   LA'S:
    WRITE (6,30)
    DO 31 G=1,2
31 WRITE (6,100) (G,K,LA0(G,K),LA1(G,K),LA2(G,K),LA3(G,K),LA4(G,K),
X   LA5(G,K),LA6(G,K),K=1,KR)
C   SR'S:
    WRITE (6,40)
    G=1
41 WRITE (6,100) (G,K,SRO(G,K),SR1(G,K),SR2(G,K),SR3(G,K),SR4(G,K),
X   SR5(G,K),SR6(G,K),K=1,KR)
C   P'S:
    WRITE (6,50)
    DO 51 G=1,2
51 WRITE (6,100) (G,K,P0(G,K),P1(G,K),P2(G,K),P3(G,K),P4(G,K),
X   P5(G,K),P6(G,K),K=1,KR)
C   Q'S:
    WRITE (6,60)
    DO 61 G=1,2
61 WRITE (6,100) (G,K,Q0(G,K),Q1(G,K),Q2(G,K),Q3(G,K),Q4(G,K),
X   Q5(G,K),Q6(G,K),K=1,KR)
C   R'S:
    WRITE (6,70)
    DO 71 G=1,2
71 WRITE (6,101) (G,K,RO(G,K),R1(G,K),R2(G,K),R3(G,K),R4(G,K),K=1,KR)
C   BOUNDARY VALUES:
    WRITE (6,80)
    DO 81 G=1,2
81 WRITE (6,100) (G,K,PO(G,K),PH(G,K),PD7(G,K),PH7(G,K),DO(G,K),
X   DH(G,K),HR(K),K=1,KR)

```

```

PRT00037
PRT00038
PRT00039
PRT00040
PRT00041
PRT00042
PRT00043
PRT00044
PRT00045
PRT00046
PRT00047
PRT00048
PRT00049
PRT00050
PRT00051
PRT00052
PRT00053
PRT00054
PRT00055
PRT00056
PRT00057
PRT00058
PRT00059
PRT00060
PRT00061
PRT00062
PRT00063
PRT00064
PRT00065
PRT00066
PRT00067
PRT00068
PRT00069
PRT00070
PRT00071
PRT00072

```

```

RETURN
1002 WRITE (6,90)
      WRITE (6,110) ((L1(I,J),J=1,6),I=1,N)
      WRITE (6,91)
      WRITE (6,110) ((L2(I,J),J=1,6),I=1,N)
      WRITE (6,92)
      WRITE (6,110) ((F1(I,J),J=1,6),I=1,N)
      WRITE (6,93)
      WRITE (6,110) ((F2(I,J),J=1,6),I=1,N)
      WRITE (6,94)
      WRITE (6,110) ((F3(I,J),J=1,6),I=1,N)
      WRITE (6,95)
      WRITE (6,110) ((F4(I,J),J=1,6),I=1,N)
      WRITE (6,96)
      WRITE (6,110) ((T(I,J),J=1,6),I=1,N)
10  FORMAT ('1  G',4X,'K',7X,'KA0(G,K)',7X,'KA1(G,K)',7X,'KA2(G,K)',
X       7X,'KA3(G,K)',7X,'KA4(G,K)',7X,'KA5(G,K)',7X,'KA6(G,K)',/)
20  FORMAT ('1  G',4X,'K',7X,'KB0(G,K)',7X,'KB1(G,K)',7X,'KB2(G,K)',
X       7X,'KB3(G,K)',7X,'KB4(G,K)',7X,'KB5(G,K)',7X,'KB6(G,K)',/)
22  FORMAT ('1  G',4X,'K',7X,'KC0(G,K)',7X,'KC1(G,K)',7X,'KC2(G,K)',
X       7X,'KC3(G,K)',7X,'KC4(G,K)',7X,'KC5(G,K)',7X,'KC6(G,K)',/)
24  FORMAT ('1  G',4X,'K',7X,'KD0(G,K)',7X,'KD1(G,K)',7X,'KD2(G,K)',
X       7X,'KD3(G,K)',7X,'KD4(G,K)',7X,'KD5(G,K)',7X,'KD6(G,K)',/)
30  FORMAT ('1  G',4X,'K',7X,'LA0(G,K)',7X,'LA1(G,K)',7X,'LA2(G,K)',
X       7X,'LA3(G,K)',7X,'LA4(G,K)',7X,'LA5(G,K)',7X,'LA6(G,K)',/)
40  FORMAT ('1  G',4X,'K',7X,'SR0(G,K)',7X,'SR1(G,K)',7X,'SR2(G,K)',
X       7X,'SR3(G,K)',7X,'SR4(G,K)',7X,'SR5(G,K)',7X,'SR6(G,K)',/)
50  FORMAT ('1  G',4X,'K',8X,'P0(G,K)',8X,'P1(G,K)',8X,'P2(G,K)',8X,
X       'P3(G,K)',8X,'P4(G,K)',8X,'P5(G,K)',8X,'P6(G,K)',/)
60  FORMAT ('1  G',4X,'K',8X,'Q0(G,K)',8X,'Q1(G,K)',8X,'Q2(G,K)',8X,
X       'Q3(G,K)',8X,'Q4(G,K)',8X,'Q5(G,K)',8X,'Q6(G,K)',/)
70  FORMAT ('1  G',4X,'K',8X,'R0(G,K)',8X,'R1(G,K)',8X,'R2(G,K)',8X,
X       'R3(G,K)',8X,'R4(G,K)',/)
80  FORMAT ('1  G',4X,'K',8X,'PO(G,K)',8X,'PH(G,K)',7X,'PO7(G,K)',7X,
X       'PH7(G,K)',8X,'DO(G,K)',8X,'DH(G,K)',10X,'HR(K)',/)
100  FORMAT (2I5,7D15.7)

```

```

PRTO0073
PRTO0074
PRTO0075
PRTO0076
PRTO0077
PRTO0078
PRTO0079
PRTO0080
PRTO0081
PRTO0082
PRTO0083
PRTO0084
PRTO0085
PRTO0086
PRTO0087
PRTO0088
PRTO0089
PRTO0090
PRTO0091
PRTO0092
PRTO0093
PRTO0094
PRTO0095
PRTO0096
PRTO0097
PRTO0098
PRTO0099
PRTO0100
PRTO0101
PRTO0102
PRTO0103
PRTO0104
PRTO0105
PRTO0106
PRTO0107
PRTO0108

```

```
101 FORMAT (2I5,5D15.7)
 90 FORMAT ('1MATRIX L1:',/)
 91 FORMAT ('1MATRIX L2:',/)
 92 FORMAT ('1MATRIX F1:',/)
 93 FORMAT ('1MATRIX F2:',/)
 94 FORMAT ('1MATRIX F3:',/)
 95 FORMAT ('1MATRIX F4:',/)
 96 FORMAT ('1MATRIX T:',/)
110 FORMAT (6D20.10)
    RETURN
    END
```

```
PRT00109
PRT00110
PRT00111
PRT00112
PRT00113
PRT00114
PRT00115
PRT00116
PRT00117
PRT00118
PRT00119
```



C	SUBROUTINE POWER	POWE0001
C	SOLVES THE 2*N MULTIGROUP EQUATIONS: $M*PHI = (1/LAMDA)*F*PHI$	POWE0002
C	BY THE FISSION SOURCE POWER METHOD	POWE0003
C	USING SIMULTANEOUS OVERRELAXATION.	POWE0004
C	WHERE: M AND F ARE DOUBLE PRECISION 2N BY 2N BLOCK MATRICES;	POWE0005
C	AND: PHI IS THE 2N FLUX (FAST AND THERMAL) VECTOR.	POWE0006
C	$L1*PHI1 = CHI1*(F1*PHI1 + F2*PHI2)$	POWE0007
C	$-T*PHI1 + L2*PHI2 = CHI2*(F3*PHI1 + F4*PHI2)$	POWE0008
C	METHOD FOLLOWS WACHPRESS, PAGE 83. SOLUTION BY GROUP ITERATION.	POWE0009
	IMPLICIT REAL*8 (A-H,L-Z)	POWE0010
	COMMON /B1/ IBC,IPLT,JPLT,IPUNCH,ISEE	POWE0011
	COMMON /B2/ KR,N	POWE0012
	COMMON /B3/ L1(50,6),L2(50,6),F1(50,6),F4(50,6),F3(50,6),F2(50,6),	POWE0013
X	T(50,6)	POWE0014
	COMMON /B4/ PHI(2,52), PSI(2,52), LAMDA, ICOUT	POWE0015
	COMMON /B5/ S(52), ERROR(2,52), Z(52), G(50,6), GT(50)	POWE0016
	COMMON /B6/ TE1(2,5),TE2(2,5),TE3(5),IN(5)	POWE0017
	COMMON /CHIF/ CHI(2)	POWE0018
	COMMON /XAXIS/ HX,HR(25)	POWE0019
	COMMON /ER/ EPS1,EPS2,EPS3	POWE0020
	COMMON /FSTR/ PHISTR(2,26,6)	POWE0021
	COMMON /ESTR/ LAMSTR(300), EFSTR(2,300), EFMSTR(2,300), ERLAM(300)	POWE0022
	COMMON /TRUE/ TRULAM, TRUPHI(2,52), PHICON(2,300), LAMCON(300),IFT	POWE0023
	DIMENSION PSI1(52), PSI2(52), SQ(2), DPHI(2), ERRMAX(2)	POWE0024
	INTEGER N, NN	POWE0025
C	DEFAULT OPTIONS FOR THE TRUE EIGENVALUE AND FLUX-CURRENT VECTOR:	POWE0026
	TRULAM=1.0	POWE0027
	DO 5 IG=1,2	POWE0028
	DO 4 I=1,N	POWE0029
	TRUPHI(IG,I)=1.0	POWE0030
	IF (MOD(I,2).EQ.1) TRUPHI(IG,I)=0.0	POWE0031
4	CONTINUE	POWE0032
	IF (IBC.EQ.1) TRUPHI(IG,N)=0.0	POWE0033
	IF (IBC.EQ.4) TRUPHI(IG,1)=1.0	POWE0034
5	CONTINUE	POWE0035
C	DEFAULT OPTIONS FOR POWER PARAMETERS:	POWE0036

ALPHA=1.25	POWE0037
LAMDA=1.0	POWE0038
DO 555 IG=1,2	POWE0039
IF (IBC.NE.4) GO TO 551	POWE0040
DO 550 I=1,N	POWE0041
PHI(IG,I)=1.0	POWE0042
IF (MOD(I,2).EQ.1) PHI(IG,I)=0.0	POWE0043
550 CONTINUE	POWE0044
PHI(IG,1)=1.0	POWE0045
GO TO 555	POWE0046
551 X=3.1415926/HX	POWE0047
IF (IBC.NE.1) X=X/2.0	POWE0048
PHI(IG,1)=-X	POWE0049
SUM1=0.0	POWE0050
DO 552 K=2,KR	POWE0051
I=2*K-2	POWE0052
J=I+1	POWE0053
SUM1=SUM1+HR(K-1)	POWE0054
PHI(IG,I)=DSIN(SUM1*X)	POWE0055
552 PHI(IG,J)=-X*DCOS(SUM1*X)	POWE0056
PHI(IG,N)=1.0	POWE0057
IF (IBC.EQ.1) PHI(IG,N)=X	POWE0058
555 CONTINUE	POWE0059
C READ IN THE TRUE (EXPECTED) EIGENVALUE AND FLUX VECTOR (MINUS 0 BC'S):	POWE0060
IFT=0	POWE0061
READ (5,500,END=501) TRULAM, (TRUPHI(1,I),I=1,N)	POWE0062
READ (5,503,END=501) (TRUPHI(2,I),I=1,N)	POWE0063
IFT=1	POWE0064
500 FORMAT (E25.14,/, (4E20.10))	POWE0065
C READ IN: OVERRELAXATION PARAMETERS ; ALPHA (OUTER ITERATION)	POWE0066
C INITIAL GUESS AT EIGENVALUE; LAMDA	POWE0067
C INITIAL NORMALIZED FLUX ; PHI(1-N)	POWE0068
501 READ (5,506,END=510) ALPHA	POWE0069
READ (5,502,END=510) LAMDA	POWE0070
READ (5,503) (PHI(1,I),I=1,N)	POWE0071
READ (5,503) (PHI(2,I),I=1,N)	POWE0072

506	FORMAT (F10.5)	POWE0073
502	FORMAT (E25.14)	POWE0074
503	FORMAT ((4E20.10))	POWE0075
510	CONTINUE	POWE0076
C	STORING FOR PRINTING THE MULTIGROUP FLUX SHAPE.	POWE0077
	K=0	POWE0078
	IF (IBC.EQ.4) K=1	POWE0079
	DO 11 IG=1,2	POWE0080
	DO 10 I=1,KR	POWE0081
	II=I+K	POWE0082
10	PHISTR(IG,II,2)=PHI(IG,2*I)	POWE0083
	IF (IBC.EQ.4) PHISTR(IG,1,2)=PHI(IG,1)	POWE0084
C	FILL RUNNING COORD IN PHISTR	POWE0085
	KR1=KR+1	POWE0086
	DO 11 I=1,KR1	POWE0087
11	PHISTR(IG,I,1)=DFLOAT(I)	POWE0088
C	IK IS THE FLUX PLOTTING COUNTER.	POWE0089
	IK=1	POWE0090
C	STORES THE ITERATION NUMBER FOR FLUX HISTORY PLOTTING:	POWE0091
	IN(1)=0	POWE0092
C	STORES TEMPORARY ERRORS FOR FLUX HISTORY PLOTTING:	POWE0093
	TE1(1,1)=0.	POWE0094
	TE1(2,1)=0.	POWE0095
	TE2(1,1)=0.	POWE0096
	TE2(2,1)=0.	POWE0097
	TE3(1)=0.0	POWE0098
C	EIGENVALUE OF THE PREVIOUS ITERATION:	POWE0099
	LAMB4=LAMDA	POWE0100
C	THE MAXIMUM NUMBER OF ALLOWED ITERATIONS: ICMAX	POWE0101
	ICMAX=300	POWE0102
C	PRINT OUT THE POWER METHOD PARAMETER INFORMATION:	POWE0103
	WRITE (6,700) ICMAX,ALPHA,LAMDA,(PHI(1,I),I=1,N)	POWE0104
	WRITE (6,701) (PHI(2,I),I=1,N)	POWE0105
700	FORMAT ('EXECUTING MULTIGROUP FISSION SOURCE POWER ITERATION METH	POWE0106
	XOD.',///,	POWE0107
	X 5X,'MAXIMUM NUMBER OF ALLOWABLE ITERATIONS:',/,	POWE0108

X	10X,'ICMAX =',I4,///,	POWE0109
X	5X,'OUTER ITERATION RELAXATION PARAMETER:',/,	POWE0110
X	10X,'ALPHA =',F7.3,///,	POWE0111
X	5X,'INITIAL GUESS AT EIGENVALUE:',/,	POWE0112
X	10X,'LAMBDA =',E22.14,///,	POWE0113
X	5X,'INITIAL GUESS AT THE GROUP FLUX SHAPE CONNECTION POINTS:',	POWE0114
X	//,8X,'FAST GROUP:',/,	POWE0115
X	10X,'F(K)''S =',4E25.14,/, (18X,4E25.14))	POWE0116
701	FORMAT ('0',7X,'THERMAL GROUP:',/,	POWE0117
X	10X,'F(K)''S =',4E25.14,/, (18X,4E25.14))	POWE0118
C	BEGIN ITERATION LOOP.	POWE0119
	ICOUT=0	POWE0120
C	ICOUT IS THE OUTER ITERATION COUNTER.	POWE0121
20	ICOUT=ICOUT+1	POWE0122
	IF (ICOUT.GT.ICMAX) GO TO 100	POWE0123
C	SOLVE FOR THE NEW GROUP FLUX VECTORS: PSI:	POWE0124
C	THE GROUP FLUX ITERATES:	POWE0125
	DO 21 I=1,N	POWE0126
	PSI1(I)=PHI(1,I)	POWE0127
21	PSI2(I)=PHI(2,I)	POWE0128
C	THE FAST ITERATION SOURCE S:	POWE0129
	CALL VPRJD(F1,PSI1,S,N)	POWE0130
	CALL VPROD(F2,PSI2,Z,N)	POWE0131
	DO 22 I=1,N	POWE0132
22	S(I)=CHI(1)*(S(I)+Z(I))	POWE0133
C	FAST FLUX:	POWE0134
	CALL SOLVE(L1,PSI1,S,N,G,GT)	POWE0135
C	THE THERMAL ITERATION SOURCE S:	POWE0136
	CALL VPROD(F3,PSI1,S,N)	POWE0137
	CALL VPROD(F4,PSI2,Z,N)	POWE0138
	DO 25 I=1,N	POWE0139
25	S(I)=CHI(2)*(S(I)+Z(I))	POWE0140
	CALL VPROD(T,PSI1,Z,N)	POWE0141
	DO 27 I=1,N	POWE0142
27	S(I)=S(I)+Z(I)	POWE0143
C	THERMAL FLUX:	POWE0144

	CALL SOLVE(L2,PSI2,S,N,G,GT)	POWE0145
C	CALCULATION OF THE EIGENVALUE:	POWE0146
	SUM1=0.000	POWE0147
	SUM2=0.000	POWE0148
	DO 28 I=1,N	POWE0149
	SUM1=SUM1+PSI1(I)*PHI(1,I)	POWE0150
28	SUM2=SUM2+PSI1(I)*PSI1(I)	POWE0151
	DO 29 I=1,N	POWE0152
	SUM1=SUM1+PSI2(I)*PHI(2,I)	POWE0153
29	SUM2=SUM2+PSI2(I)*PSI2(I)	POWE0154
	LAMDA=SUM2/SUM1	POWE0155
	LAMSTR(ICOUT)=LAMDA	POWE0156
	ERRLAM=DABS(LAMDA-LAMB4)	POWE0157
C	PUT PSI1 AND PSI2 INTO BIGGER PSI:	POWE0158
	DO 30 I=1,N	POWE0159
	PSI(1,I)=PSI1(I)	POWE0160
30	PSI(2,I)=PSI2(I)	POWE0161
C	POINT BY POINT SIMULTANEOUS RELAXATION FLUX ITERATION:	POWE0162
	X=ALPHA	POWE0163
C	DO NOT RELAX DURING THE FIRST THREE ITERATIONS:	POWE0164
	IF (ICOUT.LE.3) X=1.0	POWE0165
C	CALCULATE THE NEW GROUP FLUX ITERATES AND GROUP ERRORS:	POWE0166
	DO 40 IG=1,2	POWE0167
	DO 40 I=1,N	POWE0168
	PSI(IG,I)=PHI(IG,I) + X*(PSI(IG,I)/LAMDA-PHI(IG,I))	POWE0169
40	CONTINUE	POWE0170
C	NORMALIZE THE FLUX ITERATE (ONE GROUP):	POWE0171
	CALL NORMAL(PSI,N)	POWE0172
	NN=N	POWE0173
	IF (IBC.EQ.1) NN=N-2	POWE0174
C	NORMALIZED ERRORS OF THE FLUX ONLY:	POWE0175
	DO 39 IG=1,2	POWE0176
	ERRMAX(IG)=0.0	POWE0177
	SQ(IG)=0.0	POWE0178
	IF (IBC.NE.4) GO TO 37	POWE0179
	ERROR(IG,1)=DABS((PSI(IG,1)-PHI(IG,1))/PSI(IG,1) )	POWE0180

ERRMAX(IG)=ERROR(IG,1)	POWE0181
SQ(IG)=ERROR(IG,1)**2	POWE0182
37 DO 38 I=2,NN,2	POWE0183
ERROR(IG,I)=DABS((PSI(IG,I)-PHI(IG,I))/PSI(IG,I) )	POWE0184
IF (ERROR(IG,I).GT.ERRMAX(IG)) ERRMAX(IG)=ERROR(IG,I)	POWE0185
38 SQ(IG)=SQ(IG)+ERROR(IG,I)**2	POWE0186
SQ(IG)=DSQRT(SQ(IG))	POWE0187
C        UPDATE THE FLUX ITERATE:	POWE0188
DO 39 I=1,N	POWE0189
39 PHI(IG,I)=PSI(IG,I)	POWE0190
IF (IFT.EQ.0) GO TO 31	POWE0191
DLAM=LAMDA-TRULAM	POWE0192
DO 36 IG=1,2	POWE0193
DPHI(IG)=0.0	POWE0194
DO 35 I=1,N	POWE0195
35 DPHI(IG)=DPHI(IG)+(PSI(IG,I)-TRUPHI(IG,I))**2	POWE0196
36 DPHI(IG)=DSQRT(DPHI(IG))	POWE0197
31 IF (IPLT.NE.2) GO TO 45	POWE0198
C        THE FOLLOWING IS FOR NICELY PLOTTING THE GROUP FLUX HISTORY.	POWE0199
CALL NORM2(PSI,TRUPHI,N)	POWE0200
K=0	POWE0201
IF (IBC.EQ.4) K=1	POWE0202
KBC=KR	POWE0203
IF (IBC.EQ.1) KBC=KR-1	POWE0204
IF (IBC.EQ.4) KBC=KR+1	POWE0205
DO 41 IG=1,2	POWE0206
DO 41 I=1,KR	POWE0207
II=I+K	POWE0208
41 ERROR(IG,II)=PSI(IG,2*I)	POWE0209
IF (IBC.EQ.4) ERROR( 1,1)=PSI( 1,1)	POWE0210
IF (IBC.EQ.4) ERROR( 2,1)=PSI( 2,1)	POWE0211
C        ERROR NOW CONTAINS THE NEW NORMALIZED FLUX ITERATE PHI.	POWE0212
JK=IK	POWE0213
IF (IK.EQ.0) JK=5	POWE0214
DO 42 IG=1,2	POWE0215
DO 42 I=1,KBC	POWE0216

	IF (DABS(ERROR(IG,I)-PHISTR(IG,I,JK+1)).GE.0.01) GO TO 43	POWE0217
42	CONTINUE	POWE0218
C	FLUX HAS NOT CHANGED ENOUGH FOR PLOTTING.	POWE0219
	GO TO 45	POWE0220
C	SAVE THE NORMALIZED FLUX FOR PLOTTING:	POWE0221
43	IK=IK+1	POWE0222
	IN(IK)=ICOUT	POWE0223
	TE3(IK)=ERRLAM	POWE0224
	DO 44 IG=1,2	POWE0225
	TE1(IG,IK)=ERRMAX(IG)	POWE0226
	TE2(IG,IK)=SQ(IG)	POWE0227
	DO 44 I=1,KBC	POWE0228
44	PHISTR(IG,I,IK+1)=ERROR(IG,I)	POWE0229
	IF (IK.NE.5) GO TO 45	POWE0230
C	PLOT THE LAST FIVE SAVED FLUXES:	POWE0231
	CALL PHIPLT(5)	POWE0232
	IK=0	POWE0233
45	CONTINUE	POWE0234
C	ERROR CRITERIA FOR ACCEPTANCE OF CONVERGENCE.	POWE0235
	IFLAG1=0	POWE0236
	IFLAG2=0	POWE0237
	IFLAG3=0	POWE0238
C	STORE THE ERRORS FOR COMPARISON:	POWE0239
C	ERROR BETWEEN ITERATION EIGENVALUES:	POWE0240
	ERLAM(ICOUT)=ERRLAM	POWE0241
	DO 46 IG=1,2	POWE0242
C	MAXIMUM ERROR BETWEEN ITERATION FLUXES:	POWE0243
	EFSTR(IG,ICOUT)=ERRMAX(IG)	POWE0244
C	MEAN SQUARE ERROR BETWEEN ITERATION FLUXES:	POWE0245
	EFMSTR(IG,ICOUT)=SQ(IG)	POWE0246
C	MEAN SQUARE ERROR BETWEEN THE ITERATION FLUX AND GIVEN TRUE FLUX:	POWE0247
	PHICON(IG,ICOUT)=DPHI(IG)	POWE0248
46	CONTINUE	POWE0249
C	ERROR BETWEEN THE ITERATION EIGENVALUE AND GIVEN TRUE EIGENVALUE:	POWE0250
	LAMCON(ICOUT)=DLAM	POWE0251
	IF ((ERRMAX(1).LT.EPS1).AND.(ERRMAX(2).LT.EPS1)) IFLAG1=1	POWE0252

	IF ((SQ(1).LT.EPS2).AND.(SQ(2).LT.EPS2))	IFLAG2=1	POWE0253
	IF (ERRLAM.LT.EPS3) IFLAG3=1		POWE0254
	IFLAG4=IFLAG1*IFLAG2*IFLAG3		POWE0255
	IF (IFLAG4.EQ.1) GO TO 50		POWE0256
C	OTHERWISE CONTINUE THE ITERATION.		POWE0257
	LAMB4=LAMDA		POWE0258
	GO TO 20		POWE0259
50	CONTINUE		POWE0260
C	CONVERGENCE ACCOMPLISHED.		POWE0261
C			POWE0262
C	NORMALIZE THE CONVERGED FLUX VECTOR:		POWE0263
	CALL NORMAL(PHI,N)		POWE0264
C	PLOT ANY LEFT OVER FLUX HISTORY PLOTS:		POWE0265
	IF ((IPLJT.EQ.2).AND.(IK.NE.0)) CALL PHIPLT(IK)		POWE0266
C	BOUNDARY CONDITION INSERTIONS.		POWE0267
	IER=0		POWE0268
C	IER ALLOWS B.C. INSERTIONS FOR YES AND NO CONVERGENCE:		POWE0269
55	DO 70 IG=1,2		POWE0270
	PHI(IG,N+2)=0.0		POWE0271
	DO 60 I=1,N		POWE0272
	J=N+1-I		POWE0273
60	PHI(IG,J+1)=PHI(IG,J)		POWE0274
	IF (IBC.NE.4) GO TO 65		POWE0275
	PHI(IG,2)=0.0		POWE0276
	GO TO 70		POWE0277
65	PHI(IG,1)=0.0		POWE0278
	IF (IBC.NE.1) GO TO 70		POWE0279
	PHI(IG,N+2)=PHI(IG,N+1)		POWE0280
	PHI(IG,N+1)=0.0		POWE0281
	70 CONTINUE		POWE0282
	90 IF (IER.EQ.1) GO TO 102		POWE0283
	RETURN		POWE0284
C	NO CONVERGENCE ACCOMPLISHED:		POWE0285
100	CONTINUE		POWE0286
C	NORMALIZE THE UNCONVERGED FLUX:		POWE0287
	CALL NORMAL(PHI,N)		POWE0288



```
      ICOUT=ICOUT-1
      WRITE (6,101) ICOUT
101  FORMAT (1H1,'POWER METHOD DID NOT CONVERGE FOR THIS CASE AFTER',
      X   I4,' ITERATIONS.',//,1X,'EXECUTION TERMINATED ')
      IER=1
      GO TO 55
102  CONTINUE
C      FOR PRINTING OUT THE EIGENVALUE HISTORY AND THE FINAL FLUX SHAPE:
      IPLOT=1
      JPLOT=1
      RETURN
      END
```

```
POWE0289
POWE0290
POWE0291
POWE0292
POWE0293
POWE0294
POWE0295
POWE0296
POWE0297
POWE0298
POWE0299
POWE0300
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SUBROUTINE VPROD(A,X,S,N)
C   FORMS THE VECTOR S = PRODUCT OF NXN MATRIX A AND VECTOR X:
C   WHERE A IS THE CUBIC HERMITE (I,6) STORAGE MATRIX.
REAL*8 A(50,6),X(50),S(50)
DO 1 I=1,N
S(I)=0.0
DO 1 M=1,6
K=2*(I/2)+M-3
IF ((K.LT.1).OR.(K.GT.N)) GO TO 1
S(I)=S(I)+A(I,M)*X(K)
1 CONTINUE
C
RETURN
END

```

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VPRO0001
VPRO0002
VPRO0003
VPRO0004
VPRO0005
VPRO0006
VPRO0007
VPRO0008
VPRO0009
VPRO0010
VPRO0011
VPRO0012
VPRO0013
VPRO0014

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	SUBROUTINE SOLVE(A,X,Y,N,G,Z)	SOLV0001
C	SOLVES $A*X = Y$ USING CHOLESKY'S METHOD OF FACTORAZATION	SOLV0002
C	FOR POSITIVE DEFINITE REAL AND SYMMETRIC MATRICES A.	SOLV0003
C	REFERENCE: FORSYTHE & MOLER.	SOLV0004
C	G AND Z ARE TEMPORARY WORK AREAS.	SOLV0005
C	MODIFIED FOR THE CRAZY CUBIC HERMITE (I,6) MATRICES:	SOLV0006
	IMPLICIT REAL*8 (A-H,O-Z)	SOLV0007
	DIMENSION A(50,6),G(50,6),X(50),Y(50),Z(50)	SOLV0008
C	FORM THE MATRIX FACTORAZATION TO G:	SOLV0009
	CALL FORMG(A,G,N)	SOLV0010
C	SOLVE: $G*Z = Y$ :	SOLV0011
	CALL LOWTRI(G,Z,Y,N)	SOLV0012
C	FORM G AS SYMMETRIC MATRIX:	SOLV0013
	CALL SYMG(G,N)	SOLV0014
C	SOLVE: $G-TRANSPOSE*X = Z$ :	SOLV0015
	CALL UPPTRI(G,X,Z,N)	SOLV0016
	RETURN	SOLV0017
	END	SOLV0018

```

C      SUBROUTINE FURMG(A,G,N)
        FORMS MATRIX G FROM A:
        IMPLICIT REAL*8 (A-H,O-Z)
        DIMENSION A(50,6),G(50,6)
        DO 20 J=1,N
            K=3+MOD(J,2)
            KO=K-1
            LO=1
            IF (J.LE.3) LO=2
            IF (J.EQ.1) LO=4
            SUM=0.0
            IF (LO.GT.KO) GO TO 2
            DO 1 L=LO,KO
1          SUM=SUM+G(J,L)**2
2          SUM=A(J,K)-SUM
            IF (SUM.LT.0.0) GO TO 100
            G(J,K)=DSQRT(SUM)
            IF (J.EQ.N) GO TO 20
            I1=J+1
            I2=J+6-K
            IF (I2.GT.N) I2=N
            M=2
            IF (K.EQ.3) M=3
            DO 10 I=I1,I2
                SUM=0.0
                LO=1
                IF (I.LE.3) LO=2
                MO=M-1
                IF (LO.GT.MO) GO TO 7
                DO 5 L=LO,MO
                    JL=L
                    IF (M.EQ.2) JL=3
5                SUM=SUM+G(I,L)*G(J,JL)
7                G(I,M)=(A(I,M)-SUM)/G(J,K)
                IF (M.EQ.3) M=1
10           CONTINUE

```

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FORM0001
FORM0002
FORM0003
FORM0004
FORM0005
FORM0006
FORM0007
FORM0008
FORM0009
FORM0010
FORM0011
FORM0012
FORM0013
FORM0014
FORM0015
FORM0016
FORM0017
FORM0018
FORM0019
FORM0020
FORM0021
FORM0022
FORM0023
FORM0024
FORM0025
FORM0026
FORM0027
FORM0028
FORM0029
FORM0030
FORM0031
FORM0032
FORM0033
FORM0034
FORM0035
FORM0036

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```
20 CONTINUE
   RETURN
100 WRITE (6,101) J,K
101 FORMAT ('OERROR IN FORMG:',//,
X 5X,'A(',I2,',',I2,') < 0.0',//,
X 5X,'CHOLESKY METHOD HAS FAILED.',//,
X 5X,'MATRIX A MAY NOT BE POSITIVE DEFINITE OR SYMMETRIC ',//,
X 'OEXECUTION TERMINATED ')
   CALL EXIT
   RETURN
   END
```

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FORM0037
FORM0038
FORM0039
FORM0040
FORM0041
FORM0042
FORM0043
FORM0044
FORM0045
FORM0046
FORM0047
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	SUBROUTINE SYMG(G,N)		SYMG0001
C	FORMS SYMMETRIC G FROM G LOWER TRIANGULAR:		SYMG0002
	REAL*8 G(50,6)		SYMG0003
	N2=N-2		SYMG0004
C	FILL THE UPPER PORTION OF SYMMETRIC G:		SYMG0005
10	DO 20 I=1,N2		SYMG0006
	IF (MOD(I,2).EQ.1) GO TO 15		SYMG0007
C	I IS EVEN:		SYMG0008
	G(I,4)=G(I+1,3)		SYMG0009
	G(I,5)=G(I+2,1)		SYMG0010
	G(I,6)=G(I+3,1)		SYMG0011
	GO TO 20		SYMG0012
C	I IS ODD:		SYMG0013
15	G(I,5)=G(I+1,2)		SYMG0014
	G(I,6)=G(I+2,2)		SYMG0015
20	CONTINUE		SYMG0016
	G(N-1,5)=G(N,2)		SYMG0017
	RETURN		SYMG0018
	END		SYMG0019

```
C      SUBROUTINE LQWTRI(G,Z,Y,N)
C      SOLVES:  G*Z = Y;  FOR Z
C      WHERE G IS NXN LOWER TRIANGULAR.
      REAL*8 G(50,6),Z(50),Y(50),SUM
      DO 10 I=1,N
      K=3+MOD(I,2)
      K0=K-1
      SUM=0.0
      LT=2*(I/2)
      DO 5 M=1,K0
      L=LT+M-3
      IF ((L.LT.1).OR.(L.GT.N)) GO TO 5
      SUM=SUM+G(I,M)*Z(L)
5  CONTINUE
10  Z(I)=(Y(I)-SUM)/G(I,K)
      RETURN
      END
```

```
LOWT0001
LOWT0002
LOWT0003
LOWT0004
LOWT0005
LOWT0006
LOWT0007
LOWT0008
LOWT0009
LOWT0010
LOWT0011
LOWT0012
LOWT0013
LOWT0014
LOWT0015
LOWT0016
LOWT0017
```

```

SUBROUTINE UPPTRI(G,X,Z,N)
C   SOLVES: G*X = Z; FOR X
C   WHERE G IS NXN UPPER TRIANULAR.
REAL*8 G(50,6),X(50),Z(50),SUM
DO 10 J=1,N
I=N+1-J
K=3+MOD(I,2)
K1=K+1
SUM=0.0
LT=2*(I/2)
DO 5 M=K1,6
L=LT+M-3
IF ((L.LT.1).OR.(L.GT.N)) GO TO 5
SUM=SUM+G(I,M)*X(L)
5 CONTINUE
10 X(I)=(Z(I)-SUM)/G(I,K)
RETURN
END

```

```

UPPT0001
UPPT0002
UPPT0003
UPPT0004
UPPT0005
UPPT0006
UPPT0007
UPPT0008
UPPT0009
UPPT0010
UPPT0011
UPPT0012
UPPT0013
UPPT0014
UPPT0015
UPPT0016
UPPT0017
UPPT0018

```



C SUBROUTINE NORMAL(PHI,N)  
NORMALIZES THE GROUP FLUXES AND CURRENTS BY THE LARGEST FLUX VALUE:  
COMMON /B1/ IBC  
REAL\*8 PHI(2,52), A  
NN=N  
IF (IBC.EQ.1) NN=N-2  
A=0.0  
IF (IBC.NE.4) GO TO 5  
A=DABS(PHI(1,1))  
IF (DABS(PHI(2,1)).GT.A) A=DABS(PHI(2,1))  
5 DO 1 IG=1,2  
DO 1 I=2,NN,2  
IF (DABS(PHI(IG,I)).GT.A) A=DABS(PHI(IG,I))  
1 CONTINUE  
DO 2 IG=1,2  
DO 2 I=1,N  
2 PHI(IG,I)=PHI(IG,I)/A  
RETURN  
END

NORL0001  
NORL0002  
NORL0003  
NORL0004  
NORL0005  
NORL0006  
NORL0007  
NORL0008  
NORL0009  
NORL0010  
NORL0011  
NORL0012  
NORL0013  
NORL0014  
NORL0015  
NORL0016  
NORL0017  
NORL0018  
NORL0019

	SUBROUTINE PHIPLT(L)	PHIP0001
C	PLOTS THE GROUP FLUX HISTORY, WITH UP TO 5 GROUP FLUXES PER PLOT.	PHIP0002
C	FAST AND THERMAL GROUP FLUXES ARE PLOTTED SEPERATELY.	PHIP0003
C	L IS THE NUMBER OF FLUXES TO BE PLOTTED.	PHIP0004
C	L IS BETWEEN 1 AND 5.	PHIP0005
	IMPLICIT REAL*8 (A-H,O-Z)	PHIP0006
	COMMON /B1/ IBC	PHIP0007
	COMMON /B2/ KR,N	PHIP0008
	COMMON /B5/ A(26,6), B(26,6)	PHIP0009
	COMMON /B6/ TE1(2,5),TE2(2,5),TE3(5),IN(5)	PHIP0010
	COMMON /ER/ EPS1,EPS2,EPS3	PHIP0011
	COMMON /FSTR/ PHISTR(2,26,6)	PHIP0012
	DIMENSION SYMBOL(5)	PHIP0013
	INTEGER SYMBOL /'.', '-.', '+.', '#', '*'/	PHIP0014
	KR1=KR+1	PHIP0015
C	SET UP B.C. CONDITIONS	PHIP0016
	IF (IBC.EQ.4) GO TO 5	PHIP0017
	IF (IBC.EQ.3) GO TO 3	PHIP0018
	DO 2 IG=1,2	PHIP0019
	DO 2 K=1,L	PHIP0020
	DO 1 I=1,KR	PHIP0021
	J=KR+1-I	PHIP0022
	1 PHISTR(IG,J+1,K+1)=PHISTR(IG,J,K+1)	PHIP0023
	2 PHISTR(IG,1,K+1)=0.	PHIP0024
	3 IF (IBC.EQ.2) GO TO 5	PHIP0025
	DO 4 IG=1,2	PHIP0026
	DO 4 K=1,L	PHIP0027
	4 PHISTR(IG,KR1,K+1)=0.	PHIP0028
	5 CONTINUE	PHIP0029
C	FLUXES IN PHISTR HAVE BEEN NORMALIZED IN POWER.	PHIP0030
C	PUT THE FAST FLUX IN A, AND THE THERMAL FLUX IN B:	PHIP0031
	L1=L+1	PHIP0032
	DO 10 K=1,L1	PHIP0033
	DO 10 I=1,KR1	PHIP0034
	A(I,K)=PHISTR(1,I,K)	PHIP0035
10	B(I,K)=PHISTR(2,I,K)	PHIP0036

C	<pre> PLOT THE L FAST FLUX SHAPES ON ONE GRAPH: CALL PRTPLOT(0,A,KR1,L1,KR1,0,26,6,2) WRITE (6,20) 20 FORMAT (/, 'FAST FLUX ITERATION HISTORY PLOT.', /) WRITE (6,30) 30 FORMAT ( X   'OKEY:', 5X, 'SYMBOL', 5X, 'ITERATION NUMBER:', 7X, 'ERRCR CRITERIA', X   11X, 'ERROR', 13X, 'TOLERANCE') DO 35 I=1,L 35 WRITE (6,40) SYMBOL(I),IN(I),TE1(1,I),EPS1,TE2(1,I),EPS2, X           TE3(I),EPS3 40 FORMAT (/, 12X, A1, 15X, I3, 16X, 'FLUX', 14X, 1PD15.5, 5X, 1PD15.5, /, X   47X, 'MEAN SQ. FLUX', 5X, 1PD15.5, 5X, 1PD15.5, /, X   47X, 'EIGENVALUE', 8X, 1PD15.5, 5X, 1PD15.5) </pre>	<pre> PHIP0037 PHIP0038 PHIP0039 PHIP0040 PHIP0041 PHIP0042 PHIP0043 PHIP0044 PHIP0045 PHIP0046 PHIP0047 PHIP0048 PHIP0049 PHIP0050 PHIP0051 PHIP0052 PHIP0053 PHIP0054 PHIP0055 PHIP0056 PHIP0057 PHIP0058 PHIP0059 PHIP0060 </pre>
C	<pre> PLOT THE L THERMAL FLUX SHAPES ON THE OTHER GRAPH: CALL PRTPLOT(0,B,KR1,L1,KR1,0,26,6,2) WRITE (6,50) 50 FORMAT (/, 'OTHERMAL FLUX ITERATION PLOT.', /) WRITE (6,30) DO 55 I=1,L 55 WRITE (6,40) SYMBOL(I),IN(I),TE1(2,I),EPS1,TE2(2,I),EPS2, X           TE3(I),EPS3 RETURN END </pre>	

```
C      SUBROUTINE CURENT
C      FORMS THE SEPERATE FLUX AND CURRENT VECTORS
C      FROM THE COMBINED ELEMENTS OF PHI.
C      THEN: FLUX=PHI; CURRENT=CUR.
COMMON /B2/ KR,N
COMMON /B4/ PHI(2,52),CUR(2,52)
REAL*8 PHI,CUR
KR1=KR+1
DO 2 IG=1,2
DO 1 K=1,KR1
1 CUR(IG,K)=PHI(IG,2*K)
DO 2 K=1,KR1
2 PHI(IG,K)=PHI(IG,2*K-1)
RETURN
END
```

```
CURE0001
CURE0002
CURE0003
CURE0004
CURE0005
CURE0006
CURE0007
CURE0008
CURE0009
CURE0010
CURE0011
CURE0012
CURE0013
CURE0014
CURE0015
```

C	SUBROUTINE OUTPUT PRINTS THE RESULTS OF THE METHOD. IMPLICIT REAL*8 (A-H,L-Z) COMMON /B1/ IBC,IPLT,JPLT,IPUNCH COMMON /B2/ KR,N COMMON /B4/ PHI(2,52), CUR(2,52), LAMDA, ICOUT COMMON /B5/ PSI(2,26) COMMON /ER/ EPS1, EPS2, EPS3 COMMON /ESTR/ LAMSTR(300), EFSTR(2,300), EFMSTR(2,300), ERLAM(300) COMMON /TRUE/ TRULAM, TRUPHI(2,52), PHICCN(2,300), LAMCCN(300), IFT INTEGER N KRO=KR-1 KR1=KR+1 WRITE (6,1) 1 FORMAT ('RESULTS OF THE MULTIGROUP METHOD:') WRITE (6,10) ICOUT 10 FORMAT (//, ' PROBLEM TERMINATED AFTER', I5, X ' OUTER (POWER) ITERATIONS TO:') WRITE (6,20) LAMDA 20 FORMAT (/, 10X, 'LAMDA = ', 1PE21.14) C PRINT OUT EIGENVALUES. CALL PLOT WRITE (6,30) 30 FORMAT ('RESULTS AFTER PROBLEM TERMINATION:', /, X ' OINDEX', 8X, ' THERMAL FLUX', 11X, ' FAST FLUX', 5X, X ' THERMAL CURRENT', 8X, ' FAST CURRENT', /) WRITE (6,50) (K, PHI(2,K), PHI(1,K), CUR(2,K), CUR(1,K), K=1, KR1) 50 FORMAT (I6, 4D20.7) IF (IPUNCH.EQ.1) CALL PUNCH C PRINT OUT GROUP NORMALIZED RESULTS: DO 52 IG=1,2 A=0.0 DO 51 I=1, KR1 IF (PHI(IG,I).GT.A) A=PHI(IG,I) 51 CONTINUE DO 52 I=1, KR1	OUTP0001 OUTP0002 OUTP0003 OUTP0004 OUTP0005 OUTP0006 OUTP0007 OUTP0008 OUTP0009 OUTP0010 OUTP0011 OUTP0012 OUTP0013 OUTP0014 OUTP0015 OUTP0016 OUTP0017 OUTP0018 OUTP0019 OUTP0020 OUTP0021 OUTP0022 OUTP0023 OUTP0024 OUTP0025 OUTP0026 OUTP0027 OUTP0028 OUTP0029 OUTP0030 OUTP0031 OUTP0032 OUTP0033 OUTP0034 OUTP0035 OUTP0036
---	--	--

PHI(IG,I)=PHI(IG,I)/A	OUTP0037
52 CUR(IG,I)=CUR(IG,I)/A	OUTP0038
WRITE (6,55) (K,PHI(2,K),PHI(1,K),CUR(2,K),CUR(1,K),K=1,KR1)	OUTP0039
55 FORMAT (//,'OGROUP NORMALIZED RESULTS:',//,(I6,4D20.7))	OUTP0040
C CALCULATE THE FINAL TO EXPECTED FLUX RATIOS:	OUTP0041
K1=1	OUTP0042
K2=KR1	OUTP0043
IF (IBC.LE.2) K1=2	OUTP0044
IF ((IBC.EQ.1).OR.(IBC.EQ.3)) K2=KR	OUTP0045
DO 60 IG=1,2	OUTP0046
IF (IBC.LE.2) PSI(IG,1)=1.0	OUTP0047
IF ((IBC.EQ.1).OR.(IBC.EQ.3)) PSI(IG,KR1) = 1.0	OUTP0048
I=0	OUTP0049
DO 60 K=K1,K2	OUTP0050
I=I+2	OUTP0051
60 PSI(IG,K)=PHI(IG,K)/TRUPHI(IG,I)	OUTP0052
WRITE (6,70) (I,PSI(2,I),PSI(1,I),I=1,KR1)	OUTP0053
70 FORMAT ('RATIOS OF THE TERMINATED GROUP FLUX TO THE EXPECTED GROU	OUTP0054
XP FLUX:',//,	OUTP0055
X 10X,'- AN INDICATION OF THE ACCURACY OF THE CONVERGENCE -',///,	OUTP0056
X ' K',12X,'THERMAL RATIO',15X,'FAST RATIO',//,(I5,2E25.10))	OUTP0057
C PRINT OUT THE STORED ITERATION ERRORS:	OUTP0058
WRITE (6,110) EPS1,(EFSTR(2,I),I=1,ICOUT)	OUTP0059
WRITE (6,111) EPS1,(EFSTR(1,I),I=1,ICOUT)	OUTP0060
WRITE (6,112) EPS2,(EFMSTR(2,I),I=1,ICOUT)	OUTP0061
WRITE (6,113) EPS3,(EFMSTR(1,I),I=1,ICOUT)	OUTP0062
WRITE (6,114) EPS3,(ERLAM(I),I=1,ICOUT)	OUTP0063
110 FORMAT ('MAXIMUM NORMALIZED ERRORS BETWEEN THE THERMAL FLUX ITERA	OUTP0064
XTIONS:',	OUTP0065
X 25X,'TOLERANCE USED = ',1PE12.4,//, (1P5E20.5))	OUTP0066
111 FORMAT ('MAXIMUM NORMALIZED ERRORS BETWEEN THE FAST FLUX ITERATIO	OUTP0067
XNS:',	OUTP0068
X 25X,'TOLERANCE USED = ',1PE12.4,//, (1P5E20.5))	OUTP0069
112 FORMAT ('MEAN SQUARE NORMALIZED ERRCR BETWEEN THE THERMAL FLUX IT	OUTP0070
ERATIONS:',	OUTP0071
X 18X,'TOLERANCE USED = ',1PE12.4,//, (1P5E20.5))	OUTP0072

113	FORMAT ('MEAN SQUARE NORMALIZED ERROR BETWEEN THE FAST FLUX ITERATIONS:',	OUTP0073
	X 18X, 'TOLERANCE USED = ', 1PE12.4, '//, (1P5E20.5))	OUTP0074
114	FORMAT ('ERROR BETWEEN THE ITERATION EIGENVALUES:',	OUTP0075
	X 28X, 'TOLERANCE USED = ', 1PE12.4, '//, (1P5E20.5))	OUTP0076
	IF (IFT.EQ.0) RETURN	OUTP0077
C	PRINT OUT THE GIVEN TRUE EIGENVALUE AND FLUX:	OUTP0078
	WRITE (6,115) TRULAM, ((TRUPHI(3-J,I), J=1,2), I=1,N)	OUTP0079
115	FORMAT ('THE GIVEN TRUE EIGENVALUE:', '//, 15X,	OUTP0080
	X 'TRULAM = ', E22.14, '///,	OUTP0081
	X 'OTHE GIVEN MULTIGROUP FLUXES:', '//,	OUTP0082
	X 13X, 'THERMAL ', 16X, 'FAST ', '//, (2D20.10))	OUTP0083
C	PRINT OUT THE STORED CONVERGENCE ERRORS:	OUTP0084
	WRITE (6,120) (PHICON(2,I), I=1, ICCUT)	OUTP0085
	WRITE (6,121) (PHICON(1,I), I=1, ICCUT)	OUTP0086
	WRITE (6,122) (LAMCCN(I), I=1, ICCUT)	OUTP0087
120	FORMAT ('MEAN SQUARE ERROR BETWEEN THE THERMAL ITERATION FLUX AND	OUTP0088
	X THE GIVEN TRUE THERMAL FLUX:', '//, (1P5E20.5))	OUTP0089
121	FORMAT ('MEAN SQUARE ERROR BETWEEN THE FAST ITERATION FLUX AND TH	OUTP0090
	X E GIVEN TRUE FAST FLUX:', '//, (1P5E20.5))	OUTP0091
122	FORMAT ('ERROR BETWEEN THE ITERATION EIGENVALUES AND THE GIVEN TR	OUTP0092
	X UE EIGENVALUE:', '//, (1P5E20.5))	OUTP0093
C	RETURN	OUTP0094
	END	OUTP0095
		OUTP0096
		OUTP0097

<pre> SUBROUTINE PLOT C   PLOTS OUT THE EIGENVALUE HISTORY AS A TABLE AND A GRAPH, C   AS WELL AS PLOTTING OUT THE FINAL MULTIGROUP FLUX SHAPES. IMPLICIT REAL*8 (A-H,L-Z) COMMON /B1/ IBC, IPLOT, JPLOT, IPUNCH COMMON /B2/ KR COMMON /B4/ PHI(2,52), PSI(2,52), LAMDA, ICOUT COMMON /B5/ B(300,2) COMMON /ESTR/ LAMSTR(300) DIMENSION C(26,3) C   IN ORDER TO SAVE SOME SPACE: EQUIVALENCE (B(1),C(1)) WRITE (6,1) (LAMSTR(I),I=1,ICOUT) 1  FORMAT ('TABLE OF EIGENVALUES DURING THE POWER ITERATION:', X      '//,(1P5E25.14)) IF (JPLOT.EQ.0) GO TO 20 DO 10 I=1,ICOUT B(I,1)=I 10 B(I,2)=LAMSTR(I) CALL PRTPLT(1,B,ICOUT,2,ICOUT,0,300,2,1) WRITE (6,11) 11 FORMAT ('PLOT OF THE EIGENVALUE HISTORY THROUGH THE ITERATIONS.')</pre>	<pre> PLOT0001 PLOT0002 PLOT0003 PLOT0004 PLOT0005 PLOT0006 PLOT0007 PLOT0008 PLOT0009 PLOT0010 PLOT0011 PLOT0012 PLOT0013 PLOT0014 PLOT0015 PLOT0016 PLOT0017 PLOT0018 PLOT0019 PLOT0020 PLOT0021 PLOT0022 PLOT0023 PLOT0024 PLOT0025 PLOT0026 PLOT0027 PLOT0028 PLOT0029 PLOT0030 PLOT0031 PLOT0032 PLOT0033 PLOT0034</pre>
<pre> 20 IF (IPLOT.EQ.0) RETURN KR1=KR+1 DO 30 I=1,KR1 C(I,1)=I C(I,2)=PHI(1,I) 30 C(I,3)=PHI(2,I) CALL PRTPLT(2,C,KR1,3,KR1,0,26,3,2) WRITE (6,31) 31 FORMAT ('FINAL CONVERGED CONNECTING FLUX POINTS; F(K).',//, X 5X,'FAST FLUX: .',/,5X,'THERMAL FLUX: -') RETURN END</pre>	



```
C SUBROUTINE PUNCH
   PUNCHES OUT THE CUBIC RESULTS FOR PLOT2G ROUTINE INPUT:
   REAL*8 F,C
   COMMON /B2/ KR
   COMMON /B4/ F(2,52), C(2,52)
   KR1=KR+1
   WRITE (7,1) KR,(F(1,I),C(1,I),F(2,I),C(2,I),I=1,KR1)
  1 FORMAT (15,/, (4D20.7))
   WRITE (6,100)
100 FORMAT (///, ' THE OUTPUT HAS BEEN PUNCHED OUT ONTO CARDS ')
   RETURN
   END
```

```
PNCH0001
PNCH0002
PNCH0003
PNCH0004
PNCH0005
PNCH0006
PNCH0007
PNCH0008
PNCH0009
PNCH0010
PNCH0011
PNCH0012
```

```

SUBROUTINE NORM2(PSI,TRUPHI,N)
NORMALIZES BOTH ENERGY GROUP FLUXES IN PSI TO 1.0:
DITTO FOR TRUPHI ON THE FIRST CALL.
COMMON /B1/ IBC
REAL*8 PSI(2,52), TRUPHI(2,52), A(2)
DATA K /0/
K=K+1
NN=N
IF (IBC.EQ.1) NN=NN-2
DO 1 IG=1,2
A(IG)=0.0
IF (IBC.NE.4) GO TO 7
A(IG)=DABS(PSI(IG,1))
7 DO 1 I=2,NN,2
IF (DABS(PSI(IG,I)).GT.A(IG)) A(IG)=DABS(PSI(IG,I))
1 CONTINUE
DO 2 IG=1,2
DO 2 I=1,N
IF (A(IG).EQ.0.0) GO TO 2
PSI(IG,I)=PSI(IG,I)/A(IG)
2 CONTINUE
IF (K.NE.1) RETURN
DO 5 IG=1,2
A(IG)=0.
IF (IBC.NE.4) GO TO 8
A(IG)=DABS(TRUPHI(IG,1))
8 DO 5 I=2,NN,2
IF (TRUPHI(IG,I).GT.A(IG)) A(IG)=TRUPHI(IG,I)
5 CONTINUE
DO 6 IG=1,2
DO 6 I=1,N
IF (A(IG).EQ.0.0) GO TO 6
TRUPHI(IG,I)=TRUPHI(IG,I)/A(IG)
6 CONTINUE
RETURN
END

```

```

NOR20001
NOR20002
NOR20003
NOR20004
NOR20005
NOR20006
NOR20007
NOR20008
NOR20009
NOR20010
NOR20011
NOR20012
NOR20013
NOR20014
NOR20015
NOR20016
NOR20017
NOR20018
NOR20019
NOR20020
NOR20021
NOR20022
NOR20023
NOR20024
NOR20025
NOR20026
NOR20027
NOR20028
NOR20029
NOR20030
NOR20031
NOR20032
NOR20033
NOR20034
NOR20035
NOR20036

```

C  
C  
C

SUBROUTINE PRTPLT(NO,B,N,M,NL,NS,KX,JX,ISP)

\* IDENTICAL TO SUBROUTINE PRTPLT PREVIOUSLY LISTED IN PROGRAM REF2G.

RETURN  
END

P RTP0001  
P RTP0002  
P RTP0003  
P RTP0004  
P RTP0005  
P RTP0006

F.4. SOURCE LISTING of Program ANALYZE

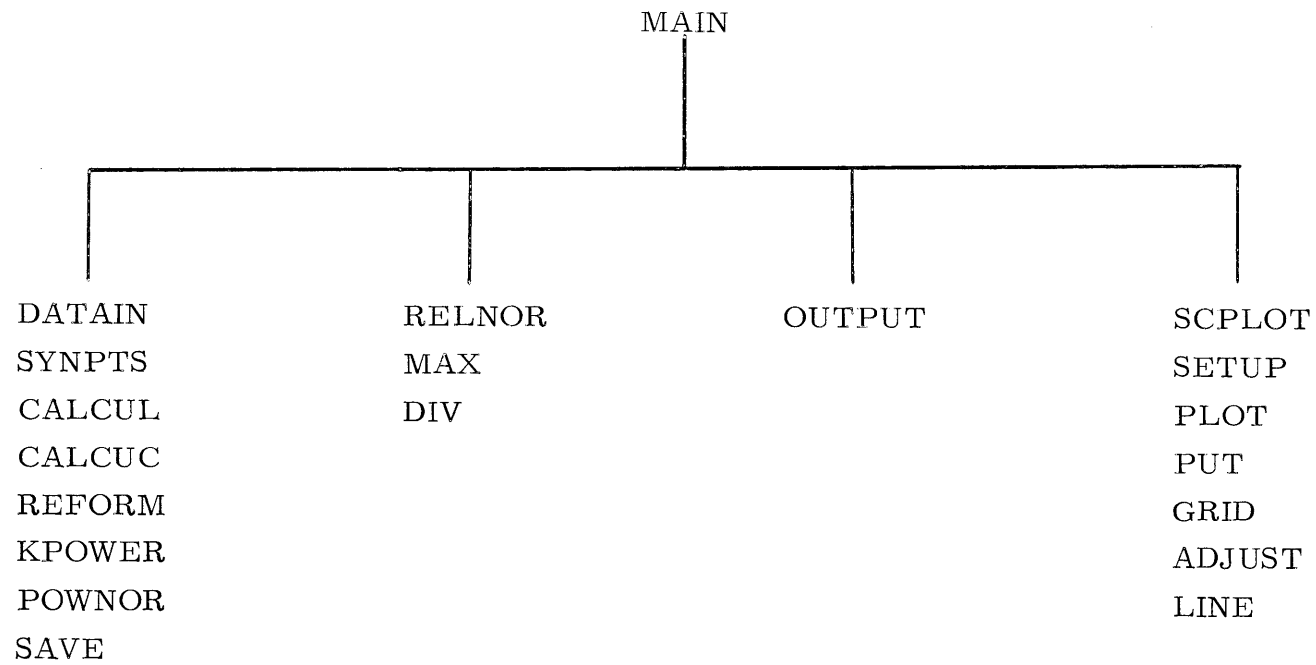


Figure F.4. Structure of Program ANALYZE.  
Not including the M.I.T. SC-4020 Subroutine Package.<sup>54</sup>

C	PROGRAM ANALYZE:	ANAL0001
C	DATA ANALYSIS AND COMPARISON PROGRAM:	ANAL0002
	DIMENSION NP(200),CHI(2),R(9,26),F(2,201),G(2,201),H(1000),	ANAL0003
X	X(1001),PHI(2,1001),SF(2,1000),D(2,1000),	ANAL0004
X	U(2,1001),XB(1001),UB(2,1001),XC(1001),UC(2,1001),	ANAL0005
X	XS(1001),NSR(3,26),NRNK(25)	ANAL0006
C	THETA IS DEFINED SUCH THAT THE CURRENT (0) = THETA * G, FOR EACH K.	ANAL0007
C	DO FOR THE 3 FLUX DATA BLOCKS:	ANAL0008
	DO 10 IT=1,3	ANAL0009
C	READ IN THE MATERIALS INPUT DATA BLOCK:	ANAL0010
	CALL DATAIN(IT,METHOD,NK,NR,NRNK,NP,N,XS,PHI,SF,D,CHI,THETA,NAP,H)	ANAL0011
C	READ IN THE CONVERGED FLUX POINTS DATA BLOCK:	ANAL0012
	CALL SYNPTS(IT,METHOD,NK,NR,F,G)	ANAL0013
C	CALCULATE THE DETAILED FLUX SHAPES:	ANAL0014
	IF (METHOD.EQ.2) GO TO 4	ANAL0015
	CALL CALCUL(IT,METHOD,NK,NR,NRNK,NP,N,X,U,XS,PHI,SF,D,CHI,	ANAL0016
X	THETA,F,G)	ANAL0017
	GO TO 5	ANAL0018
4	CALL CALCUC(IT,METHOD,NK,NR,NRNK,NP,N,X,U,XS,PHI,SF,D,CHI,	ANAL0019
X	THETA,F,G,NAP)	ANAL0020
C	TRANSFORM THE FLUX POINTS INTO NK DISTINCT REGIONS:	ANAL0021
5	CALL REFORM(IT,METHOD,NK,NR,NRNK,NP,N,X,U,SF,D)	ANAL0022
C	CALCULATE THE POWER IN EACH OF THE NK REGIONS FOR EACH FLUX:	ANAL0023
	CALL KPOWER(IT,METHOD,NK,NP,N,X,U,SF,D,CHI,THETA,NAP,F,G,H,	ANAL0024
X	NSR,R,TPOWER)	ANAL0025
C	NORMALIZE THE REGION POWERS AND THE FLUX POINTS BY TPPOWER)	ANAL0026
	CALL POWNDR(IT,METHOD,NK,NP,N,X,U,R,TPOWER)	ANAL0027
C	SAVE THESE RESULTS FOR PLOTTING:	ANAL0028
	CALL SAVE(IT,NK,N,X,U,NB,XB,UB,NC,XC,UC)	ANAL0029
10	CONTINUE	ANAL0030
C	NORMALIZE THE FLUX RESULTS TOGETHER FOR EACH GROUP TO 1.0:	ANAL0031
	CALL RELNDR(NK,N,X,U,NB,XB,UB,NC,XC,UC)	ANAL0032
C	PRINT OUT THE POWER AND ARRAY RESULTS:	ANAL0033
	CALL OUTPUT(NK,N,X,U,NB,XB,UB,NC,XC,UC,NSR,R)	ANAL0034
C	READ IN THE PLOTTING INFORMATION AND PLOT THE RESULTS ON THE SC 4020:	ANAL0035
	CALL SCPLLOT(NK,N,X,U,NB,XB,UB,NC,XC,UC)	ANAL0036

```
WRITE (6,20)  
20 FORMAT (/, 'ONORMAL PROGRAM TERMINATION.')
```

```
STOP  
END
```

```
ANAL0037  
ANAL0038  
ANAL0039  
ANAL0040
```

```

SUBROUTINE DATAIN(IT,METHOD,NK,NR,NRNK,NP,N,X,PHI,SF,C,CHI,THETA,
X      NAP,H)
C      READS IN THE MATERIAL DATA INPUT BLOCKS:
DIMENSION NP(200),K(200),L(200),ITF(200),KTF(200),CHI(2),
X      X(1001),PHI(2,1001),SF(2,1000),D(2,1000),H(1000),
X      NRNK(25)
C      DATA BLOCK ID CARDS:
READ (IT,20) METHODD, NK, NR, NAP
C      NRNK = NUMBER OF NR REGIONS PER EACH NK REGION:
IF (NR.NE.NRK) READ (IT,20) (NRNK(I),I=1,NK)
20 FORMAT (16I5)
XSTART=0.0
K(1)=1
READ (IT,11) NR
READ (IT,12) (ITF(I),I=1,NR)
READ (IT,17) CHI(1),CHI(2), THETA
IF (THETA.EQ.0.) THETA=1.
NUMITF=1
DO 10 I=1,NR
IF (ITF(I).LT.NUMITF) GO TO 2
C      NEW REGION DATA:
NUMITF=NUMITF+1
KTF(NUMITF-1)=I
READ (IT,13) NS
IF (NS.EQ.0) NS=1
NP(I)=NS+1
IF (I.NE.1) K(I)=K(I-1)+NP(I-1)
L(I)=NP(I)
IF (I.NE.1) L(I)=L(I-1)+NP(I)
KO=K(I)
LO=L(I)
L1=LO-1
L2=LO-2
IF (NS.GT.1) READ (IT,14) (X(J),H(J),SF(1,J),D(1,J),
X      SF(2,J),D(2,J),J=KO,L2)
READ (IT,15) X(L1),X(L0),H(L1),SF(1,L1),D(1,L1),SF(2,L1),D(2,L1)

```

```

DATA0001
DATA0002
DATA0003
DATA0004
DATA0005
DATA0006
DATA0007
DATA0008
DATA0009
DATA0010
DATA0011
DATA0012
DATA0013
DATA0014
DATA0015
DATA0016
DATA0017
DATA0018
DATA0019
DATA0020
DATA0021
DATA0022
DATA0023
DATA0024
DATA0025
DATA0026
DATA0027
DATA0028
DATA0029
DATA0030
DATA0031
DATA0032
DATA0033
DATA0034
DATA0035
DATA0036

```



```

XSTART=XSTART-X(KO)
DO 1 J=KO,LO
1 X(J)=X(J)+XSTART
XSTART=X(LO)
IF (IT.EQ.3) GO TO 10
READ (IT,16) (PHI(1,J),J=KO,LO)
READ (IT,16) (PHI(2,J),J=KO,LO)
GO TO 10
C OLD REPEATED REGION DATA:
2 M=ITF(I)
M=KTF(M)
NP(I)=NP(M)
K(I)=K(I-1)+NP(I-1)
L(I)=L(I-1)+NP(I)
KO=K(I)
LO=L(I)
LI=LO-1
KM=K(M)
IB=KO-KM
XSTART=XSTART-X(KM)
DO 5 J=KO,LO
NJ=J-IB
X(J)=X(NJ)+XSTART
H(J)=H(NJ)
DO 5 IG=1,2
D(IG,J)=D(IG,NJ)
SF(IG,J)=SF(IG,NJ)
PHI(IG,J)=PHI(IG,NJ)
5 CONTINUE
XSTART=X(LO)
10 CONTINUE
11 FORMAT (/ ,15)
12 FORMAT (25I2)
13 FORMAT (5X,15)
14 FORMAT (F10.5,10X,F10.5,10X,2E10.3,/,40X,2E10.3)
15 FORMAT (3F10.5,10X,2E10.3,/,40X,2E10.3)

```

```

DATA0037
DATA0038
DATA0039
DATA0040
DATA0041
DATA0042
DATA0043
DATA0044
DATA0045
DATA0046
DATA0047
DATA0048
DATA0049
DATA0050
DATA0051
DATA0052
DATA0053
DATA0054
DATA0055
DATA0056
DATA0057
DATA0058
DATA0059
DATA0060
DATA0061
DATA0062
DATA0063
DATA0064
DATA0065
DATA0066
DATA0067
DATA0068
DATA0069
DATA0070
DATA0071
DATA0072

```

```
16 FORMAT (E20.7)
17 FORMAT (3F10.5)
  N=L(NR)
  IF (IT.NE.3) RETURN
  DO 50 I=1,N
  DO 50 IG=1,2
50 PHI(IG,I)=1.0
  RETURN
  END
```

```
DATA0073
DATA0074
DATA0075
DATA0076
DATA0077
DATA0078
DATA0079
DATA0080
DATA0081
```

	SUBROUTINE SYNPTS(IT,METHOD,NK,NR,F,G)	SYNP0001
C	READS IN CONVERGED FLUX AND CURRENT POINTS:	SYNP0002
	DIMENSION F(2,201),G(2,201)	SYNP0003
	JT=IT+10	SYNP0004
	READ (JT,1) NR	SYNP0005
1	FORMAT (I5)	SYNP0006
	NR1=NR+1	SYNP0007
	IF (METHOD.NE.1) GO TO 10	SYNP0008
C	FOR LINEAR SYNTHESIS:	SYNP0009
	IF (IT.NE.3) READ (JT,2) (F(1,I),F(2,I),I=1,NR1)	SYNP0010
	IF (IT.EQ.3) READ (JT,3) (F(1,I),I=1,NR1),(F(2,I),I=1,NR1)	SYNP0011
2	FORMAT (2E20.7)	SYNP0012
3	FORMAT (E20.7)	SYNP0013
	GO TO 20	SYNP0014
10	READ (JT,11) (F(1,I),G(1,I),F(2,I),G(2,I),I=1,NR1)	SYNP0015
11	FORMAT (4E20.7)	SYNP0016
20	RETURN	SYNP0017
	END	SYNP0018

	SUBROUTINE CALCUL(IT,METHOD,NK,NR,NRNK,NU,XU,U,XS,PHI,SF,D,CHI,	CALL0001
X	THETA,F,G)	CALL0002
C	CALCULATES THE FLUX TRIAL FUNCTION U FOR LINEAR SYNTHESIS:	CALL0003
	DIMENSION NP(200),XU(1001),U(2,1001),PHI(2,1001),SF(2,1000),	CALL0004
X	D(2,1000),CHI(2),F(2,201),G(2,201),XS(1001),NRNK(25)	CALL0005
	ISYNTH=20	CALL0006
C	FOR EACH K*TH REGION:	CALL0007
	N=0	CALL0008
	LL=0	CALL0009
	DO 20 K=1,NR	CALL0010
	L=LL+1	CALL0011
	LL=LL+NP(K)	CALL0012
	DO 7 IG=1,2	CALL0013
	IF (PHI(IG,L).EQ.0.0) PHI(IG,L)=1.0E-6	CALL0014
	IF (PHI(IG,LL).EQ.0.0) PHI(IG,LL)=1.0E-6	CALL0015
	7 CONTINUE	CALL0016
C	FOR ALL POINTS IN THIS REGION:	CALL0017
	DO 10 I=L,LL	CALL0018
	N=N+1	CALL0019
	XU(N)=XS(I)	CALL0020
	X=(XU(N)-XS(L))/(XS(LL)-XS(L))	CALL0021
	IF ((K.EQ.1).AND.(ISYNTH.EQ.21.OR.ISYNTH.EQ.24)) X=1.0	CALL0022
	IF ((K.EQ.NR).AND.(ISYNTH.EQ.22.OR.ISYNTH.EQ.24)) X=0.0	CALL0023
	DO 1 IG=1,2	CALL0024
	1 U(IG,N)=PHI(IG,I)*(F(IG,K)*(1.-X)/PHI(IG,L)	CALL0025
X	+F(IG,K+1)*X/PHI(IG,LL))	CALL0026
	10 CONTINUE	CALL0027
	20 CONTINUE	CALL0028
	NU=N	CALL0029
	RETURN	CALL0030
	END	CALL0031

	SUBROUTINE CALCUC(IT,METHOD,NK,NR,NRNK,NP,NU,XU,U,XS,PHI,SF,D,CHI,	CALC0001
X	THETA,F,G,NAPT)	CALC0002
C	CALCULATES THE FLUX TRIAL FUNCTION U FOR CUBIC HERMITE SYNTHESIS:	CALC0003
C	NAP = # OF ADDITIONAL POINTS TO BE CALCULATED IN EACH REGION	CALC0004
C	AND IS NEGATIVE IF NAP APPLIES ONLY TO THE FIRST REGION.	CALC0005
	DIMENSION NP(200),XU(1001),U(2,1001),PHI(2,1001),SF(2,1000),	CALC0006
X	D(2,1000),CHI(2),F(2,201),G(2,201),XS(1001),NRNK(25)	CALC0007
	NAP=NAPT	CALC0008
	NL=0	CALC0009
	DO 50 K=1,NR	CALC0010
50	NL=NL+NP(K)	CALC0011
C	FOR EACH K' TH REGION:	CALC0012
	N=0	CALC0013
	LL=0	CALC0014
	DO 20 K=1,NR	CALC0015
	L=LL+1	CALC0016
	LL=LL+NP(K)	CALC0017
	H=XS(LL)-XS(L)	CALC0018
	DO 7 IG=1,2	CALC0019
	IF (PHI(IG,L).EQ.0.0) PHI(IG,L)=1.0E-6	CALC0020
	IF (PHI(IG,LL).EQ.0.0) PHI(IG,LL)=1.0E-6	CALC0021
7	CONTINUE	CALC0022
C	FOR ALL POINTS IN THIS REGION:	CALC0023
	DO 10 I=L,LL	CALC0024
	N=N+1	CALC0025
	XU(N)=XS(I)	CALC0026
	X=(XU(N)-XS(L))/(XS(LL)-XS(L))	CALC0027
	DO 1 IG=1,2	CALC0028
1	U(IG,N)=PHI(IG,I)*(F(IG,K)*(1.-3.*X**2+2.*X**3)/PHI(IG,L)	CALC0029
X	+F(IG,K+1)*(3.*X**2-2.*X**3)/PHI(IG,LL)	CALC0030
X	+H*(G(IG,K)/(D(IG,L)*PHI(IG,L))*(-X+2.*X**2-X**3)	CALC0031
X	+G(IG,K+1)/(D(IG,LL-1)*PHI(IG,LL))*(X**2-X**3))*THETA)	CALC0032
	NAPP=IABS(NAP)	CALC0033
	IF (NAP.EQ.0) GO TO 10	CALC0034
	IF (I.EQ.LL) GO TO 10	CALC0035
	DX=(XS(I+1)-XS(I))/FLOAT(NAPP+1)	CALC0036

```

DO 5 INP=1,NAPP
N=N+1
XU(N)=XU(N-1)+DX
X=(XU(N)-XS(L))/(XS(LL)-XS(L))
DO 5 IG=1,2
5 U(IG,N)=PHI(IG,I)*(F(IG,K)*(1.-3.*X**2+2.*X**3)/PHI(IG,L)
X +F(IG,K+1)*(3.*X**2-2.*X**3)/PHI(IG,LL)
X +H*(G(IG,K)/(D(IG,L)*PHI(IG,L))*(-X+2.*X**2-X**3)
X +G(IG,K+1)/(D(IG,LL-1)*PHI(IG,LL))*(X**2-X**3))*THETA)
10 CONTINUE
IF (NAPP.EQ.0) GO TO 15
IF (IT.EQ.1) GO TO 15
C SHIFTING SF AND D ARRAYS TO COMPENSATE FOR NAP .NE. 0:
NS=NP(K)-1
NNP=NS*NAPP
NPS=0
K1=K+1
IF (K1.GT.NR) GO TO 60
DO 51 I=K1,NR
51 NPS=NPS+NP(I)
DO 55 J=1,NPS
I=NL+1-J
II=I+NNP
DO 55 IG=1,2
SF(IG,II)=SF(IG,I)
55 D(IG,II)=D(IG,I)
60 NL=NL+NNP
L=NL-NPS
L1=L-NNP-NP(K)
DO 70 I=1,NS
II=NS+1-I+L1
NAPP1=NAPP+1
DO 65 J=1,NAPP1
L=L-1
DO 65 IG=1,2
SF(IG,L)=SF(IG,II)

```

```

CALC0037
CALC0038
CALC0039
CALC0040
CALC0041
CALC0042
CALC0043
CALC0044
CALC0045
CALC0046
CALC0047
CALC0048
CALC0049
CALC0050
CALC0051
CALC0052
CALC0053
CALC0054
CALC0055
CALC0056
CALC0057
CALC0058
CALC0059
CALC0060
CALC0061
CALC0062
CALC0063
CALC0064
CALC0065
CALC0066
CALC0067
CALC0068
CALC0069
CALC0070
CALC0071
CALC0072

```

```
65 D(IG,L)=D(IG,II)
70 CONTINUE
15 NP(K)=NP(K)+NAPP
   IF (NAP.LT.0) NAP=0
20 CONTINUE
   NU=N
   RETURN
   END
```

```
CALC0073
CALC0074
CALC0075
CALC0076
CALC0077
CALC0078
CALC0079
CALC0080
```

	SUBROUTINE REFORM(IT,METHOD,NK,NR,NRNK,NP,N,X,U,SF,D)	REF00001
C	REFORMS THE NR GIVEN REGIONS INTO NK DESIRED REGIONS:	REF00002
	DIMENSION NP(200),X(1001),U(2,1001),SF(2,1000),D(2,1000),NRNK(25)	REF00003
	IF (NK.EQ.NR) RETURN	REF00004
C	IT SHOULD BE 3 ONLY.	REF00005
C	DELETE ALL DOUBLE ENTRIES WITHIN EACH NK DATA BLOCK:	REF00006
	M=-1	REF00007
	I=0	REF00008
	DO 5 K=1,NK	REF00009
	L=NRNK(K)	REF00010
	DO 4 J=1,L	REF00011
	I=I+1	REF00012
	M=M+2	REF00013
	X(I)=X(M)	REF00014
	DO 4 IG=1,2	REF00015
	U(IG,I)=U(IG,M)	REF00016
	SF(IG,I)=SF(IG,M)	REF00017
	4 D(IG,I)=D(IG,M)	REF00018
	I=I+1	REF00019
	X(I)=X(M+1)	REF00020
	DO 5 IG=1,2	REF00021
	5 U(IG,I)=U(IG,M+1)	REF00022
	N=I	REF00023
C	NP(K) = # OF FLUX POINTS IN REGION K = # NS SUBREGIONS + 1:	REF00024
	DO 10 K=1,NK	REF00025
10	NP(K)=NP(K)+NRNK(K)-1	REF00026
	NR=NK	REF00027
	RETURN	REF00028
	END	REF00029



	SUBROUTINE KPOWER(IT,METHOD,NK,NP,N,X,U,SF,D,CHI,THETA,NAP,F,G,	KPCW0001
X	H,NSR,R,TPOWER)	KPCW0002
C	CALCULATES THE POWER IN EACH REGION BY INTEGRATIONS:	KPCW0003
	DIMENSION NP(200),X(1001),U(2,1001),SF(2,1000),D(2,1000),CHI(2),	KPCW0004
X	F(2,201),G(2,201),H(1000),R(9,26),NSR(3,26)	KPCW0005
	TPOWER=0.0	KPCW0006
	IF (IT.NE.1) GO TO 20	KPCW0007
C	FOR THE HOMOGENEOUS CASE:	KPCW0008
C	ASSUMING THAT NP IS ALWAYS 2 AND THUS D'S ARE CONSTANT.	KPCW0009
	DO 10 K=1,NK	KPCW0010
	M=2*K-1	KPCW0011
	NSR(IT,K)=1	KPCW0012
	R(IT,K)=0.0	KPCW0013
	DO 8 IG=1,2	KPCW0014
	IF (METHOD.NE.1) GO TO 5	KPCW0015
C	LINEAR FLUX:	KPCW0016
	R(IT,K)=(F(IG,K)+F(IG,K+1))*H(M)*SF(IG,M)/2.0+R(IT,K)	KPCW0017
	GO TO 8	KPCW0018
C	CUBIC FLUX:	KPCW0019
	5 R(IT,K)=((F(IG,K)+F(IG,K+1))*H(M)/2.0+THETA*(-G(IG,K)	KPCW0020
X	/D(IG,M)+G(IG,K+1)/D(IG,M))*H(M)**2/12.0)*SF(IG,M)+R(IT,K)	KPCW0021
	8 CONTINUE	KPCW0022
	10 TPOWER=TPOWER+R(IT,K)	KPCW0023
	RETURN	KPCW0024
C	SYNTHESIS OR REFERENCE DETAILED FLUX CASE:	KPCW0025
20	M=-1	KPCW0026
	DO 50 K=1,NK	KPCW0027
	M=M+1	KPCW0028
	NS=NP(K)-1	KPCW0029
	NSR(IT,K)=NS	KPCW0030
	R(IT,K)=0.0	KPCW0031
	DO 40 J=1,NS	KPCW0032
	M=M+1	KPCW0033
	DO 40 IG=1,2	KPCW0034
40	R(IT,K)=(U(IG,M)+U(IG,M+1))*(X(M+1)-X(M))*SF(IG,M)/2.0 + R(IT,K)	KPCW0035
50	TPOWER=TPOWER+R(IT,K)	KPCW0036

RETURN  
END

KPOW0037  
KPOW0038

C       SUBROUTINE POWNOR(IT,METHOD,NK,NP,N,X,U,R,TPOWER)  
          NORMALIZES THE FLUXES AND REGION POWERS BY TPOWER:  
          DIMENSION NP(200),X(1001),U(2,1001),R(9,26)  
          DO 1 K=1,NK  
1       R(IT+3,K)=R(IT,K)/TPOWER  
          DO 2 IG=1,2  
          DO 2 K=1,N  
2       U(IG,K)=U(IG,K)/TPOWER  
          RETURN  
          END

POWN0001  
POWN0002  
POWN0003  
POWN0004  
POWN0005  
POWN0006  
POWN0007  
POWN0008  
POWN0009  
POWN0010

```

SUBROUTINE SAVE(IT,NK,N,X,U,NB,XB,UB,NC,XC,UC)
C SAVES THE FLUX DISTRIBUTIONS IN THE IT LOOP BY PLACING:
C   IT = 1:  HOMOGENEOUS RESULTS IN UC;
C   IT = 2:  SYNTHESIS RESULTS  IN UB;
C   IT = 3:  REFERENCE RESULTS  IN U.
DIMENSION X(1001),U(2,1001),XB(1001),UB(2,1001),
X          XC(1001),UC(2,1001)
IF (IT.NE.1) GO TO 2
C   HOMOGENEOUS RESULTS:
DO 1 I=1,N
XC(I)=X(I)
DO 1 IG=1,2
1 UC(IG,I)=U(IG,I)
NC=N
RETURN
2 IF (IT.EQ.3) RETURN
C   SYNTHESIS RESULTS:
DO 3 I=1,N
XB(I)=X(I)
DO 3 IG=1,2
3 UB(IG,I)=U(IG,I)
NB=N
RETURN
END

```

```

SAVE0001
SAVE0002
SAVE0003
SAVE0004
SAVE0005
SAVE0006
SAVE0007
SAVE0008
SAVE0009
SAVE0010
SAVE0011
SAVE0012
SAVE0013
SAVE0014
SAVE0015
SAVE0016
SAVE0017
SAVE0018
SAVE0019
SAVE0020
SAVE0021
SAVE0022
SAVE0023
SAVE0024

```

```
C  SUBROUTINE RELNOR(NK,N,X,U,NB,XB,UB,NC,XC,UC)
    NORMALIZES THE SET OF U, UB, UC; TO UNITY FOR EACH GROUP:
    DIMENSION X(1001),U(2,1001),XB(1001),UB(2,1001),
X     XC(1001),UC(2,1001)
    DIMENSION A(2),B(2),C(2),Z(2,3)
    CALL MAX(N,U,A)
    CALL MAX(NB,UB,B)
    CALL MAX(NC,UC,C)
    DO 1 IG=1,2
    Z(IG,1)=A(IG)
    Z(IG,2)=B(IG)
1  Z(IG,3)=C(IG)
    CALL MAX(3,Z,A)
    CALL DIV(N,U,A)
    CALL DIV(NB,UB,A)
    CALL DIV(NC,UC,A)
    RETURN
    END
```

```
RELN0001
RELN0002
RELN0003
RELN0004
RELN0005
RELN0006
RELN0007
RELN0008
RELN0009
RELN0010
RELN0011
RELN0012
RELN0013
RELN0014
RELN0015
RELN0016
RELN0017
RELN0018
```

C      SUBROUTINE MAX(N,U,A)  
         FINDS THE MAXIMUM POSITIVE ELEMENT OF U FOR EACH GROUP:  
         DIMENSION U(2,1001),A(2)  
         DO 1 IG=1,2  
         A(IG)=0.  
         DO 1 I=1,N  
         IF (U(IG,I).GT.A(IG)) A(IG)=U(IG,I)  
1 CONTINUE  
         RETURN  
         END

MAX 0001  
MAX 0002  
MAX 0003  
MAX 0004  
MAX 0005  
MAX 0006  
MAX 0007  
MAX 0008  
MAX 0009  
MAX 0010

```
C      SUBROUTINE DIV(N,U,A)
        DIVIDES A INTO U FOR EACH GROUP:
        DIMENSION U(2,1001),A(2)
        DO 1 IG=1,2
        DO 1 I=1,N
1      U(IG,I)=U(IG,I)/A(IG)
        RETURN
        END
```

```
DIV 0001
DIV 0002
DIV 0003
DIV 0004
DIV 0005
DIV 0006
DIV 0007
DIV 0008
```

	SUBROUTINE OUTPUT(NK,N,X,U,NB,XB,UB,NC,XC,UC,NSR,R)	OUTP0001
C	PRINTS OUT THE ANALYSIS RESULTS:	OUTP0002
	DIMENSION X(1001),U(2,1001),XB(1001),UB(2,1001),	OUTP0003
X	XC(1001),UC(2,1001),R(9,26),NSR(3,26)	OUTP0004
C	SUM UP THE REGION RAW POWERS AND REGION FRACTIONAL POWERS:	OUTP0005
	DO 5 IT=1,6	OUTP0006
	IF (IT.LE.3) NSR(IT,26)=0	OUTP0007
	R(IT,26)=0.0	OUTP0008
	DO 5 K=1,NK	OUTP0009
	IF (IT.LE.3) NSR(IT,26)=NSR(IT,26)+NSR(IT,K)	OUTP0010
5	R(IT,26)=R(IT,26)+R(IT,K)	OUTP0011
C	THE PERCENT NORMALIZED DIFFERENCES OF THE REGION FRACTIONAL POWERS:	OUTP0012
	DO 60 K=1,NK	OUTP0013
	DO 50 IT=7,9	OUTP0014
50	R(IT,K)=0.0	OUTP0015
	IF (R(6,K).EQ.0.0) GO TO 60	OUTP0016
	R(7,K)=(R(6,K)-R(4,K))*100./R(6,K)	OUTP0017
	R(8,K)=(R(6,K)-R(5,K))*100./R(6,K)	OUTP0018
	IF (R(5,K).EQ.0.0) GO TO 60	OUTP0019
	R(9,K)=(R(5,K)-R(4,K))*100./R(5,K)	OUTP0020
60	CONTINUE	OUTP0021
C	THE PERCENT NORMALIZED DIFFERENCES OF THE TOTAL RAW POWER PRODUCED:	OUTP0022
	R(7,26)=(R(3,26)-R(1,26))*100./R(3,26)	OUTP0023
	R(8,26)=(R(3,26)-R(2,26))*100./R(3,26)	OUTP0024
	R(9,26)=(R(2,26)-R(1,26))*100./R(2,26)	OUTP0025
	WRITE (6,10) NK,(K,(NSR(IT,K),R(IT,K),IT=1,3),K=1,NK)	OUTP0026
10	FORMAT ('RESULTS OF THE INTEGRATED POWER IN EACH OF THE',I3,	OUTP0027
X	' REGIONS:',///,	OUTP0028
X	' CALCULATED POWER LEVELS, AND NUMBER OF SUBREGIONS PER REGION:	OUTP0029
X	' ,//,3X,'REGION:',10X,'HOMOGENIZED RESULTS:',5X,	OUTP0030
X	' SYNTHESIZED RESULTS:',7X,'REFERENCE RESULTS:',//,	OUTP0031
X	(I10,10X,I3,E17.7,I8,E17.7,I10,E15.7))	OUTP0032
	WRITE (6,12) (NSR(IT,26),R(IT,26),IT=1,3)	OUTP0033
12	FORMAT (/ ,3X,'TOTALS:',10X,I3,E17.7,I8,E17.7,I10,E15.7)	OUTP0034
	WRITE (6,20) (K,(R(IT,K),IT=4,6),K=1,NK)	OUTP0035
20	FORMAT (///,'FRACTIONAL POWER LEVELS:',//,3X,'REGION:',10X,	OUTP0036



X	'HOMOGENEOUS RESULTS:',5X,	OUTP0037
X	'SYNTHESIZED RESULTS:',7X,'REFERENCE RESULTS:',//,	OUTP0038
X	(I10,5X,3E25.7))	OUTP0039
	WRITE (6,11) (R(IT,26),IT=4,6)	OUTP0040
11	FORMAT (/,3X,'TOTALS:',5X,3E25.7)	OUTP0041
	WRITE (6,30) (K,(R(IT,K),IT=7,9),K=1,NK)	OUTP0042
30	FORMAT (//,'OFRACTIONAL POWER NORMALIZED PERCENT ERRORS:',//,	OUTP0043
X	3X,'REGION:',14X,'(REF-HOMO)/REF %',	OUTP0044
X	8X,'(REF-SYNTH)/REF %',5X,'(SYNTN-HOMO)/SYNTH %',//,	OUTP0045
X	(I10,5X,3E25.7))	OUTP0046
	WRITE (6,40) (R(IT,26),IT=7,9)	OUTP0047
40	FORMAT (//,'ORAW PRODUCTION POWER NORMALIZED PERCENT ERRORS:',//,	OUTP0048
X	24X,'(REF-HOMO)/REF %',	OUTP0049
X	8X,'(REF-SYNTH)/REF %',5X,'(SYNTN-HOMO)/SYNTH %',//,	OUTP0050
X	(15X,3E25.7))	OUTP0051
	RETURN	OUTP0052
	END	OUTP0053

	SUBROUTINE SCPLOT(NK,N,X,U,NB,XB,UB,NC,XC,UC)	SCPL0001
C	READS IN PLOTTING INFORMATION AND PLOTS THE FLUX COMPARISONS:	SCPL0002
	COMMON /SC/ TITLE(20),XINCH,YINCH,NCELL,WCELL,NLL,XL(100)	SCPL0003
	COMMON /MAX/ XMIN,XMAX,YMIN,YMAX	SCPL0004
	COMMON /RASTER/ IXS,IYS,IXE,IYE,IXT,IYT,IXLX,IYLY,	SCPL0005
X	IXLY,IYLY(3),LH,LW,IS,IR	SCPL0006
	COMMON /MARGIN/ ML,MR,MB,MT	SCPL0007
	DIMENSION X(1001),U(2,1001),XB(1001),UB(2,1001),XC(1001),	SCPL0008
X	UC(2,1001)	SCPL0009
	READ (5,1,END=20) TITLE	SCPL0010
	READ (5,4,END=20) XINCH, YINCH	SCPL0011
C	NEGATIVE NCELL SPECIFIES THAT THE LAST CELL IS A HALF CELL:	SCPL0012
	READ (5,3,END=20) NCELL, WCELL	SCPL0013
	READ (5,3,END=5) NLL, (XL(I),I=1,7)	SCPL0014
	IF (NLL.GT.7) READ (5,4) (XL(I),I=8,NLL)	SCPL0015
1	FORMAT (20A4)	SCPL0016
3	FORMAT (I10,7F10.5)	SCPL0017
4	FORMAT (8F10.5)	SCPL0018
	GO TO 10	SCPL0019
5	NLL=0	SCPL0020
10	WRITE (6,11) TITLE	SCPL0021
11	FORMAT ('EXECUTING GENERAL ANALYSIS AND FLUX PLOTTING PROGRAM:',	SCPL0022
X	//,5X,'TITLE OF PLOTTING RUN IS:   ',20A4,'  ')	SCPL0023
	XMIN=0.0	SCPL0024
	XMAX=WCELL*FLOAT(NCELL)	SCPL0025
	IF (NCELL.LT.0) XMAX=WCELL*(FLOAT(-NCELL)-0.5)	SCPL0026
	YMIN=0.	SCPL0027
	YMAX=1.0	SCPL0028
	WRITE (6,13) NCELL,WCELL,XMIN,XMAX,YMIN,YMAX,NLL	SCPL0029
13	FORMAT ('REACTOR GEOMETRY PARAMETERS:',/,5X,'NCELL =',I5,/,	SCPL0030
X	5X,'WCELL =',F10.5,/,5X,'XMIN =',F10.5,/,5X,'XMAX =',F10.5,/,	SCPL0031
X	5X,'YMIN =',F10.5,/,5X,'YMAX =',F10.5,/,	SCPL0032
X	5X,'NLL =',I5)	SCPL0033
	IF (NLL.GT.0) WRITE (6,14) (XL(I),I=1,NLL)	SCPL0034
14	FORMAT (5X,'(XL(I),I=1,NLL) =',/, (10X,10F10.5))	SCPL0035
	CALL SETUP	SCPL0036

```
CALL PLOT(N,X,U,NB,XB,UB,NC,XC,UC)  
20 RETURN  
END
```

```
SCPL0037  
SCPL0038  
SCPL0039
```

	SUBROUTINE SETUP	SETU0001
C	FINDS THE RASTER LOCATIONS FOR THE GRID CORNERS, LABELS, AND TITLE	SETU0002
C	ENTIRE PLOTTING AREA (INCLUDING LABELS) IS XINCH X YINCH.	SETU0003
	COMMON /SC/ TITLE(20),XINCH,YINCH,NCELL,WCELL,NLL,XL(100)	SETU0004
	COMMON /RASTER/ IXS,IYS,IXE,IYE,IXT,IYT,IXLX,IYLY,	SETU0005
X	IXLY,IYLY(3),LH,LW,IS,IR	SETU0006
	COMMON /MARGIN/ ML,MR,MB,MT	SETU0007
C	SET THE LETTER SIZE FOR THE LABELS AND TITLES:	SETU0008
	LH=2	SETU0009
	LW=2	SETU0010
	IS=5*LW+3	SETU0011
	IR=7*LH+5	SETU0012
C	PLOTTING AREA:	SETU0013
C	ONE INCH = 137 RASTERS. PLOTTING AREA IS 1023 RASTERS SQUARE.	SETU0014
	IXE=XINCH*137.+0.5	SETU0015
	IYE=YINCH*137.+0.5	SETU0016
	IF (IXE.GT.1020) IXE=1020	SETU0017
	IF (IYE.GT.1020) IYE=1020	SETU0018
	IXS=(1023-IXE)/2.+1	SETU0019
	IYS=(1023-IYE)/2+1	SETU0020
	IXE=IXS+IXE-1	SETU0021
	IYE=IYS+IYE-1	SETU0022
C	COORDS FOR THE LETTERS ARE FOR THEIR CENTERS:	SETU0023
C	TITLE (UPPER LEFT CORNER):	SETU0024
	IXT=8	SETU0025
	IYT=1014	SETU0026
C	X AXIS LABEL	SETU0027
	IXLX=(IXS+60+IXE-6*IS)/2+1+8	SETU0028
	IYLY=IYS + 10	SETU0029
C	Y AXIS LABELS (FOR EACH FLUX PLOT TYPE):	SETU0030
	IXLY=IXS + 10	SETU0031
	IYLY(1)=(IYS+43+IYE-20*IS)/2+1 +4	SETU0032
	IYLY(2)=(IYS+43+IYE-23*IS)/2+1 +4	SETU0033
	IYLY(3)=(IYS+43+IYE-26*IS)/2+1 +4	SETU0034
C	SET MARGIN SPACING FOR GRID:	SETU0035
	ML=IXS + 2*18 -1	SETU0036

MR=1023-IXE  
MB=IYS + 2\*18 -1  
MT=1023-IYE  
RETURN  
END

SETU0037  
SETU0038  
SETU0039  
SETU0040  
SETU0041

C	SUBROUTINE PLOT(NA,XA,UA,NB,XB,UB,NC,XC,UC)	PLOT0001
C	USES MIT-IPC'S SC-4020 SUBROUTINE PLOTTING PACKAGE.	PLOT0002
C	PLOTS THE COMPARISON FLUX DISTRIBUTIONS:	PLOT0003
C	KEY: HOMOGENEOUS - DASHED LINE.	PLOT0004
C	SYNTHESIS - SOLID LINE.	PLOT0005
C	REFERENCE - DOTTED LINE.	PLOT0006
	COMMON /RASTER/ IXS,IYS,IXE,IYE,IXT,IYT,IXLX,IYLY,	PLOT0007
X	IXLY,IYLY(3)	PLOT0008
	DIMENSION XA(1001),UA(2,1001),XB(1001),UB(2,1001),	PLOT0009
X	XC(1001),UC(2,1001), X(1001),T(1001)	PLOT0010
C	HARDCOPY INITIALIZATION:	PLOT0011
	CALL STOIDV('M7788-6571',9,2)	PLOT0012
C	DO FOR EACH ENERGY GROUP:	PLOT0013
	DO 10 IG=1,2	PLOT0014
	NTH=1	PLCT0015
	CALL GRID	PLOT0016
	IF (IG.EQ.1) CALL RITE2V(IXLY,IYLY(1),1024,90,1,20,NTH,	PLOT0017
X	'NORMALIZED FAST FLUX',N)	PLOT0018
	IF (IG.EQ.2) CALL RITE2V(IXLY,IYLY(2),1024,90,1,23,NTH,	PLOT0019
X	'NORMALIZED THERMAL FLUX',N)	PLOT0020
C	HOMOGENEOUS RESULTS:	PLOT0021
	CALL PUT(IG,NC,XC,UC,N,X,T)	PLOT0022
	CALL ADJUST(N,X,T,1)	PLCT0023
	CALL LINE(N,X,T,1)	PLOT0024
C	SYNTHESIS RESULTS:	PLOT0025
	CALL PUT(IG,NB,XB,UB,N,X,T)	PLOT0026
	CALL ADJUST(N,X,T,0)	PLOT0027
	CALL LINE(N,X,T,0)	PLOT0028
C	REFERENCE RESULTS:	PLOT0029
	CALL PUT(IG,NA,XA,UA,N,X,T)	PLOT0030
	CALL ADJUST(N,X,T,2)	PLOT0031
	CALL LINE(N,X,T,2)	PLOT0032
10	CONTINUE	PLOT0033
	CALL PLTND(NTH)	PLOT0034
	WRITE (6,20) NTH	PLOT0035
20	FORMAT ('1',I5,' SC 4020 PLOTS HAVE BEEN PLOTTED.')	PLOT0036

RETURN  
END

PLOT0037  
PLOT0038

C      SUBROUTINE PUT(IG,NA,XA,UA,N,X,T)  
          TRANSFERS NA,XA,UA OF IG INTO N,X,T:  
          DIMENSION XA(1001),UA(2,1001),X(1001),T(1001)  
          N=NA  
          DO 1 I=1,N  
          X(I)=XA(I)  
1        T(I)=UA(IG,I)  
          RETURN  
          END

PUT 0001  
PUT 0002  
PUT 0003  
PUT 0004  
PUT 0005  
PUT 0006  
PUT 0007  
PUT 0008  
PUT 0009



	SUBROUTINE GRID	GRID0001
C	SETS UP A NEW FRAME AND PLOTS THE GRID	GRID0002
C	AS WELL AS RUN TITLE, LABELS, AND LIGHT LINES.	GRID0003
	COMMON /SC/ TITLE(20),XINCH,YINCH,NCELL,WCELL,NLL,XL(100)	GRID0004
	COMMON /MAX/ XMIN, XMAX, YMIN, YMAX	GRID0005
	COMMON /RASTER/ IXS,IYS,IXE,IYE,IXT,IYT,IXLX,IYLY,	GRID0006
X	IXLY,IYLY(3),LH,LW,IS,IR	GRID0007
	COMMON /MARGIN/ ML,MR,MB,MT	GRID0008
	DATA IT /1/	GRID0009
	IT=IT+1	GRID0010
	IF (IT.GT.3) IT=3	GRID0011
	NTH=1	GRID0012
	IF (IT.GT.2) GO TO 10	GRID0013
	CALL CHSIZV(LW,LH)	GRID0014
	CALL RITSTV(IS,IR)	GRID0015
	CALL SETMIV(ML,MR,MB,MT)	GRID0016
10	CALL GRID1V(IT,XMIN,XMAX,YMIN,YMAX,WCELL,0.1,1,0,-1,-1,5,3)	GRID0017
	CALL PRINTV(80,TITLE,IXT,IYT,1)	GRID0018
	CALL RITE2V(IXLX,IYLY,1024,0,1,6,NTH,'X (CM)',N)	GRID0019
20	IF (NLL.EQ.0) GO TO 5	GRID0020
	IY1=IYV(YMIN)	GRID0021
	IY2=IYV(YMAX)	GRID0022
	DO 1 I=1,NLL	GRID0023
	IX=IXV(XL(I))	GRID0024
	1 CALL LINEV(IX,IY1,IX,IY2)	GRID0025
5	RETURN	GRID0026
	END	GRID0027

```

SUBROUTINE ADJUST(N,X,Y,L)
C   ADJUSTS THE X AND Y ARRAYS BY DELETING ANY (X,Y) POINTS
C   WITHIN OR EQUAL TO NR RASTERS DISTANCE OF EACH OTHER.
C   HOPEFULLY ELIMINATES DARK SPOTS ON THE PLOTS.
C   L = 0: SOLID LINE PLOTTING: NR = 2.
C   L = 1: DASHED LINE PLOTTING: NR = 2; INCRV(10,5).
C   L = 2: DOTTED LINE PLOTTING: NR = 4; INCRV( 2,2).
DIMENSION X(1), Y(1)
NR=2
IF (L.EQ.2) NR=4
IF (N.LE.1) RETURN
K=1
IX1=IXV(X(1))
IY1=IYV(Y(1))
DO 1 I=2,N
IX2=IXV(X(I))
IY2=IYV(Y(I))
ID=SQRT(FLJAT((IX2-IX1)**2+(IY2-IY1)**2))+0.5
IF (ID.LE.NR) GO TO 1
K=K+1
X(K)=X(I)
Y(K)=Y(I)
IX1=IX2
IY1=IY2
1 CONTINUE
N=K
RETURN
END

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ADJU0001
ADJU0002
ADJU0003
ADJU0004
ADJU0005
ADJU0006
ADJU0007
ADJU0008
ADJU0009
ADJU0010
ADJU0011
ADJU0012
ADJU0013
ADJU0014
ADJU0015
ADJU0016
ADJU0017
ADJU0018
ADJU0019
ADJU0020
ADJU0021
ADJU0022
ADJU0023
ADJU0024
ADJU0025
ADJU0026
ADJU0027
ADJU0028

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	SUBROUTINE LINE(N,X,Y,K)	LINE0001
C	PLOTS: A LINE (K=0), OR A DASHED LINE (K=1), OR A DOTTED LINE (K=2)	LINE0002
C	THROUGH Y(X) DATA POINTS.	LINE0003
	DIMENSION X(1), Y(1)	LINE0004
	IX1=IXV(X(1))	LINE0005
	IY1=IYV(Y(1))	LINE0006
	DO 10 I=2,N	LINE0007
	IX2=IXV(X(I))	LINE0008
	IY2=IYV(Y(I))	LINE0009
	IF (K.NE.0) GO TO 1	LINE0010
C	SOLID LINE PLOT:	LINE0011
	CALL LINEV(IX1,IY1,IX2,IY2)	LINE0012
	GO TO 5	LINE0013
C	DOTTED OR DASHED LINE PLOT:	LINE0014
	1 IF (K.EQ.1) CALL INCRV(10,5)	LINE0015
	IF (K.EQ.2) CALL INCRV(2,2)	LINE0016
	CALL DOTLNV(IX1,IY1,IX2,IY2)	LINE0017
	5 IX1=IX2	LINE0018
10	IY1=IY2	LINE0019
	RETURN	LINE0020
	END	LINE0021