Search for microscopic black hole signatures at the Large Hadron Collider

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A search for microscopic black hole signatures at the Large Hadron Collider

CMS Collaboration

CERN, Switzerland

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A B S T R A C T

A search for microscopic black hole production and decay in pp collisions at a center-of-mass energy of 7 TeV has been conducted by the CMS Collaboration at the LHC, using a data sample corresponding to an integrated luminosity of 35 pb\(^{-1}\). Events with large total transverse energy are analyzed for the presence of multiple high-energy jets, leptons, and photons, typical of a signal expected from a microscopic black hole. Good agreement with the standard model backgrounds, dominated by QCD multijet production, is observed for various final-state multiplicities and model-independent limits on new physics in these final states are set. Using simple semi-classical approximation, limits on the minimum black hole mass are derived as well, in the range 3.5–4.5 TeV. These are the first direct limits on black hole production at a particle accelerator.

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One of the exciting predictions of theoretical models with extra spatial dimensions and low-scale quantum gravity is the possibility of copious production of microscopic black holes in particle collisions at the CERN Large Hadron Collider (LHC) [1,2]. Models with low-scale gravity are aimed at solving the hierarchy problem, the puzzlingly large difference between the electroweak and the Planck scales.

In this Letter we focus on microscopic black hole production in a model with large, flat, extra spatial dimensions, proposed by Arkani-Hamed, Dimopoulos, and Dvali, and referred to as the ADD model [3,4]. This model alleviates the hierarchy problem by introducing \(n\) extra dimensions in space, compactified on an \(n\)-dimensional torus or sphere with radius \(r\). The multidimensional space–time is only open to the gravitational interaction, while the spatial dimensions and low-scale quantum gravity is the possibility. Such models are derived as well, in the range 3.5–4.5 TeV. These are the first direct limits on black hole production at a particle accelerator.

Extra dimensions

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The parton-level cross section of black hole production is derived from geometrical considerations and is given by \(\sigma \sim \pi r^2\) [1,2]. At LHC energies, this cross section can reach 100 pb for \(M_D\) of 1 TeV. The exact cross section cannot be calculated without knowledge of the underlying theory of quantum gravity and is subject to significant uncertainty. It is commonly accepted [1,2] that the minimum black hole mass \(M_{\text{BH}}\) cannot be smaller than \(M_D\); although the formation threshold can be significantly larger than this. When a black hole is formed, some fraction of the colliding parton energy may not be trapped within the event horizon and will be emitted in the form of gravitational shock waves, which results in energy, momentum, and angular momentum loss.

The Schwarzschild radius of a black hole with mass \(M_{\text{BH}}\) embedded in a \(4+n\) space–time can be found by solving Einstein’s general relativity equations and is given by [6–8]:

\[
 r_S = \frac{1}{\sqrt{\pi} M_D} \left[ \frac{M_{\text{BH}} 8\pi (\frac{n+3}{2})^{\frac{3}{2}}}{M_D (n+2)} \right]^{\frac{1}{n+1}}.
\]

The parton-level cross section of black hole production is derived from geometrical considerations and is given by \(\sigma \sim \pi r^2\) [1,2]. At LHC energies, this cross section can reach 100 pb for \(M_D\) of 1 TeV. The exact cross section cannot be calculated without knowledge of the underlying theory of quantum gravity and is subject to significant uncertainty. It is commonly accepted [1,2] that the minimum black hole mass \(M_{\text{BH}}\) cannot be smaller than \(M_D\); although the formation threshold can be significantly larger than this. When a black hole is formed, some fraction of the colliding parton energy may not be trapped within the event horizon and will be emitted in the form of gravitational shock waves, which results in energy, momentum, and angular momentum loss.

This effect is particularly model-dependent for black hole masses close to \(M_D\). In general, black holes in particle collisions are produced with non-zero angular momentum, which also affects their properties and production cross section.

Once produced, the microscopic black holes would decay thermally via Hawking radiation [12], approximately democratically (with equal probabilities) to all standard model (SM) degrees of freedom.
freedom. Quarks and gluons are the dominant particles produced in the black hole evaporation (≈ 75%) because they have a large number of color degrees of freedom. The remaining fraction is accounted for by leptons, W and Z bosons, photons, and possibly Higgs bosons. Emission of gravitons by a black hole in the bulk space is generally expected to be suppressed [13], although in some models it can be enhanced for rotating black holes for the case of large n [10,11,14]. In some models the evaporation is terminated earlier, when the black hole mass reaches $M_D$, with the formation of a stable non-interacting and non-accepting remnant [15]. The Hawking temperature for a black hole in $4 + n$ space–time is given by [1,2,7,8]: $T_H = \frac{\pi}{m c} (in \ Planck \ units \ h = c = k_B = 1$, where $k_B$ is the Boltzmann constant) and is typically in the range of a few hundred GeV. The lifetime for such a microscopic black hole is $\sim \tau_{\text{MD}}^{−2\eta}$ [1,2,8].

Here we consider semi-classical black holes, whose properties are similar to those for classical black holes described by general relativity and whose mass is close enough to $M_D$ so that quantum effects can not be ignored completely. There are also models [16–18] of quantum black holes that decay before they thermalize, mainly into two-jet final states. We do not consider this signature here, leaving it for dedicated searches in the dijet channel [19,20].

In what follows, we further assume that the semi-classical approximation, which is strictly valid only for $M_{BH} \gg M_D$, still holds even for the BH masses as low as $M_D$. While we expect that unknown quantum corrections to the black hole production and decay may become very important, if not dominant, for $M_{BH} \approx M_D$, we still use semi-classical approximation as a benchmark due to the lack of a better, quantum model of black hole production and decay.

The microscopic black holes produced at the LHC would be distinguished by high multiplicity, democratic, and highly isotropic decays with the final-state particles carrying hundreds of GeV of energy. Most of these particles would be reconstructed as jets of hadrons. Observation of such spectacular signatures would provide direct information on the nature of black holes as well as the structure and dimensionality of space–time [1].

The search for black holes is based on semi-classical black hole signatures described by general relativity and whose mass is close enough to $M_D$, so that quantum effects can not be ignored completely. There are also models [16–18] of quantum black hole models that decay before they thermalize, mainly into two-jet final states. We do not consider this signature here, leaving it for dedicated searches in the dijet channel [19,20].

In what follows, we further assume that the semi-classical approximation, which is strictly valid only for $M_{BH} \gg M_D$, still holds even for the BH masses as low as $M_D$. While we expect that unknown quantum corrections to the black hole production and decay may become very important, if not dominant, for $M_{BH} \approx M_D$, we still use semi-classical approximation as a benchmark due to the lack of a better, quantum model of black hole production and decay.

The search for black holes is based on $\sqrt{s} = 7$ TeV pp collision data recorded by the Compact Muon Solenoid (CMS) detector at the LHC between March and October 2010, which correspond to an integrated luminosity of $34.7 \pm 3.8 \text{ pb}^{-1}$. A detailed description of the CMS experiment can be found elsewhere [21]. The central feature of the CMS detector is the 3.8 T superconducting solenoid enclosing the silicon pixel and strip tracker, the electromagnetic calorimeter (ECAL), and the brass-scintillator hadronic calorimeter (HCAL). For triggering purposes and to facilitate jet reconstruction, the ECAL has an ultimate energy resolution better than 0.5% for unconverted photons or electrons with transverse energies above 100 GeV [26]. In 2010 collision data, for $E_T > 20$ GeV, this resolution is better than 1% in the barrel.

Muons are required to have matched tracks in the central tracker and the muon spectrometer, to be within $|\eta| < 2.1$, be consistent with the interaction vertex to suppress backgrounds from cosmic ray muons, be isolated from other tracks, and have transverse momentum $p_T$ above 20 GeV. The combined fit using tracks measured in the central tracker and the muon spectrometer results in $p_T$ resolution between 1% and 5% for $p_T$ values up to 1 TeV.

The separation between any two objects (jet, lepton, or photon) is required to be

$$\Delta R = \sqrt{\Delta \phi^2 + \Delta \eta^2} > 0.3.$$
Table 1
Monte Carlo signal points for some of the model parameters probed, corresponding leading order cross sections (σ), and the minimum required values for the event multiplicity (N ≥ Nmin) and ST (ST > Smin), as well as the signal acceptance (A), the expected number of signal events (nexp), the number of observed events (ndata), and the observed (terr) and expected (σterr) limits on the signal cross section at 95% confidence level.

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<th>M0 (TeV)</th>
<th>M0min (TeV)</th>
<th>n</th>
<th>σ (pb)</th>
<th>Nmin</th>
<th>Smin (TeV)</th>
<th>A (%)</th>
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<td>0.002</td>
<td>0</td>
<td>0.03±0.07</td>
<td>0.26</td>
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The two generators yield different values of total cross section, as BlackMax introduces an additional n-dependent factor applied on top of the geometrical cross section. The CHARYBDIS cross sections are a factor of 1.36, 1.59, and 1.78 smaller than those from BlackMax for n = 2, 4, and 6, respectively. In addition, CHARYBDIS has been used to simulate black hole evaporation resulting in a stable non-interacting remnant with mass M0 (this model is not implemented in BlackMax). In the generation, we use the Particle Data Group [5] definition of the Planck scale Mp. (Using another popular choice for M0 from Dimopoulos and Landsberg [1] would result in a suppression of the production cross section by a factor of 1.35, 5.21, or 9.29 for n = 2, 4, or 6, respectively.)

We employ a selection based on total transverse energy to separate black hole candidate events from the backgrounds. The variable ST is defined as a scalar sum of the ET of the N individual objects (jets, electrons, photons, and muons) passing the above selections. Only objects with ET > 50 GeV are included in the calculation of ST, in order to suppress the SM backgrounds and to be insensitive to jets from pile-up, while being fully efficient for black hole decays. Further, the missing transverse energy in the event is added to ST, if the missing transverse energy value exceeds 50 GeV. Note that while ET is counted toward ST, it is not considered in the determination of N.

The main background to black hole signals arises from QCD multijet events. Other backgrounds from direct photon, W/Z+jets, and tt production were estimated from MC simulations, using the MadGRAPH [34] leading-order parton-level event generator with CTEQ6L PDF set [35]; followed by PYTHIA [28] parton showering and full CMS detector simulation via GEANT4 [30]. These additional backgrounds are negligible at large values of ST and contribute less than 1% to the total background after the final selection.

The dominant multijet background can only be estimated reliably from data. For QCD events, ST is almost completely determined by the hard 2 → 2 parton scattering process. Further splitting of the jets due to final-state radiation, as well as additional jets due to initial-state radiation – most often nearly collinear with either incoming or outgoing partons – does not change the ST value considerably. Consequently, the shape of the ST distribution is expected to be independent of the event multiplicity N, as long as ST is sufficiently above the turn-on region (i.e., much higher than N × 50 GeV). This shape invariance offers a direct way of extracting the expected number of background events in the search for black hole production.

We confirmed the assumption of the ST shape invariance of N up to high multiplicities using MC generators capable of simulating multijet final states from either matrix elements [36] or parton showers [28]. The conjecture that the ST shape is independent of the multiplicity has been also checked with data using the exclusive multiplicities of N = 2 and N = 3. Even in the presence of a black hole signal with a mass of a few TeV, the decays of these black holes result in events with half-a-dozen objects in the final state. Hence, the signal contribution to the N = 2 and N = 3 data is expected to be small and only seen at large values of ST, so these samples still can be used for the background prediction at higher multiplicities. Moreover, since dedicated analyses of the dijet invariant mass spectrum have been conducted [19,20], we know that there are no appreciable contributions from new physics to the dijet final state up to invariant masses of about 1.5 TeV, which, for central jets, translates to a similar range of ST.

We fit the ST distributions between 600 and 1100 GeV, where no black hole signal is expected, for data events with N = 2 and N = 3 using an ansatz function \( \frac{p_0\exp\left(-\frac{ST}{p_0}\right)}{(ST)^{1/2} + (ST)^{1/2}} \), which is shown with the solid line in Fig. 1. To check the systematic uncertainty of the fit, we use two additional ansatz functions, \( \frac{p_0\exp\left(-\frac{ST}{p_0}\right)}{(ST)^{1/2} + (ST)^{1/2}} \) and \( \frac{p_0\exp\left(-\frac{ST}{p_0}\right)}{(ST)^{1/2} + (ST)^{1/2}} \) [19], which are shown as the upper and lower boundaries of the shaded band in Fig. 1. The default choice of the ansatz function was made based on the best-fit to the ST distribution for N = 2. Additional systematic uncertainty arises from a slight difference between the best-fit shapes for N = 2 and N = 3. Nevertheless, the fits for these two exclusive multiplicities agree with each other within the uncertainties, demonstrating that the shape of the ST distribution is independent of the final-state multiplicity.

The ST distributions for data events with multiplicities N ≥ 3, 4, and 5 are shown in Figs. 2a, b, and c, respectively. The solid curves in the figures are the predicted background shapes, found by normalizing the fits of the N = 2 ST distribution to the range...
Fig. 1. Total transverse energy $S_T$, for events with the multiplicities of a) $N = 2$, and b) $N = 3$ objects in the final state. Data are depicted as solid circles with error bars; the shaded band is the background prediction obtained from data (solid line) with its uncertainty. Non-multijet backgrounds are shown as colored histograms. Also shown is the predicted black hole signal for three different parameter sets.

$$S_T = 1000–1100 \text{ GeV},$$ where no black hole signal contribution is expected.

Since no excess is observed above the predicted background, we set limits on the black hole production. We assign a systematic uncertainty on the background estimate of 6% to 125% for the $S_T$ range used in this search. This uncertainty comes from the normalization uncertainty (4–12%, dominated by the statistics in the normalization region) added in quadrature to the uncertainties arising from using various ansatz fit functions and the difference between the shapes obtained from the $N = 2$ and $N = 3$ samples. The integrated luminosity is measured with an uncertainty of 11% [22]. The uncertainty on the signal yield is dominated by the jet-energy-scale uncertainty of $\approx 5\%$ [24] which translates into a 5% uncertainty on the signal. An additional 2% uncertainty on
the signal acceptance comes from the variation of PDFs within the CTEQ6 error set [35]. The particle identification efficiency does not affect the signal distribution, since an electron failing the identification requirements would be classified either as a photon or a jet; a photon failing the selection would become a jet; a rejected muon would contribute to the $E_T$. In any case the total value of $S_T$ is not affected.

We set limits on black hole production with the optimized $S_T$ and $N$ selections by counting events with $S_T > S_T^{\text{min}}$ and $N > N^{\text{min}}$. We optimized the signal ($S$) significance in the presence of background ($B$) using the ratio $S/\sqrt{S+B}$ for each set. The optimum choice of parameters is listed in Table 1, as well as the predicted number of background events, the expected number of signal events, and the observed number of events in data. Note that the background uncertainty, dominated by the choice of the fitting function, is highly correlated for various working points listed in Table 1 and also bin-to-bin for the $S_T$ distributions shown in Figs. 1 and 2.

We set upper limits on the black hole production cross section using the Bayesian method with flat signal prior and log-normal prior for integration over the nuisance parameters (background, signal acceptance, luminosity) [5,37]. These upper limits at the 95% confidence level (CL) are shown in Fig. 3, as a function of $M_{\text{BH}}^{\text{min}}$. For the three model parameter sets shown in the figure, the observed (expected) lower limits on the black hole mass are 3.5, 4.2 and 4.5 TeV (3.2, 4.0, and 4.5 TeV), respectively.

Translating these upper limits into lower limits on the parameters of the ADD model, we can exclude the production of black holes with minimum mass of 3.5–4.5 TeV for values of the multidimensional Planck scale up to 3.5 TeV at 95% CL. These limits, shown in Fig. 4, do not exhibit significant dependence on the details of the production and evaporation within the set of models we studied. These are the first limits of a dedicated search for black hole production at a particle accelerator.

We point out that the semi-classical approximation used in this search is valid only for the lowest values of the $M_D$, for which the limits on the minimum black hole mass exceed $M_D$ by a factor of a few. For higher values of $M_D$ the limits become comparable with $M_D$, which implies that the approximation is no longer valid and that the BH production cross section may be modified significantly. Nevertheless, due to the exponentially falling nature of production cross section with the black hole mass, even large changes in the cross section translate only in moderate changes in the minimum black hole mass limit, as evident from Fig. 3.

Finally, we produce model-independent upper limits on the cross section times the acceptance for new physics production in high-$S_T$ inclusive final states for $N \geq 3$, 4, and 5. Fig. 5 shows 95% CL upper limits from a counting experiment for $S_T > S_T^{\text{min}}$ as a function of $M_{\text{BH}}^{\text{min}}$, which can be used to test models of new physics that result in these final states. A few examples of such models are production of high-mass $t\bar{t}$ resonances [38] in the six-jet and $t+\bar{t}$ jet final states, $R$-parity violating gluino decay into three jets, resulting in the six-jet final state [39,40], and a class of models with strong dynamics, with a strongly produced resonance decaying into a pair of resonances further decaying into two jets each, resulting in the four-jet final state [41]. In addition, these limits can be used to constrain black hole production for additional regions of the parameter space of the model, as well as set limits on the existence of string balls [42], which are quantum precursors of black holes predicted in certain string models. We have checked that for the black hole model parameters we probed with the dedicated optimized analysis, the sensitivity of the search in terms of the excluded black hole mass range exceeds that from the model-independent cross section limits by as little as 5–8%. Thus, model-independent limits can be used efficiently to constrain the allowed parameter space in an even broader variety of black hole models than we covered in this Letter.

To conclude, we have performed the first dedicated search for microscopic black holes at a particle accelerator and set limits on their production in the model with large extra dimensions in space using simple semi-classical approximation of the black hole production and decay [12]. The lower limits on the black hole mass at 95% CL range from 3.5 to 4.5 TeV for values of the Planck scale up to 3 TeV. Additionally, we have produced model-independent limits on the production of energetic, high-multiplicity final states, which can be used to constrain a variety of models of new physics.

**Acknowledgements**

We wish to congratulate our colleagues in the CERN accelerator departments for the excellent performance of the LHC machine. We thank the technical and administrative staff at CERN and other CMS institutes, and acknowledge support from: FMSR (Austria); FNRS and FWO (Belgium); CNPq, CAPES, FAPERJ, and FAPESP (Brazil); MES (Bulgaria); CERN; CAS, MoST, and NSFC (China); COLCIENCIAS (Colombia); MSES (Croatia); RPF (Cyprus); Academy of Sciences and NICPB (Estonia); Academy of Finland, ME, and HIP (Finland); CEA and CNRS/IN2P3 (France); BMBF, DFG, and HGF (Germany); GSRT (Greece); OTKA (Hungary); DWORD (Ireland); INFN (Italy); NRF and WCU (Korea); LOA, MNiSW, NWO (Netherlands); BRF, SNF (Switzerland); NAS (Serbia) and INFN (Spain); and in the United States by NSF.

**Fig. 3.** The 95% confidence level upper limits on the black hole production cross section (solid lines) and three theoretical predictions for the cross section (dashed lines), as a function of the black hole mass.

**Fig. 4.** The 95% confidence level limits on the black hole mass as a function of the multidimensional Planck scale $M_D$ for several benchmark scenarios. The area below each curve is excluded by this search.
Fig. 5. Model-independent 95% confidence level upper limits on a signal cross section times acceptance for counting experiments with $S_T > S_T^{\text{min}}$ as a function of $S_T^{\text{min}}$ for (a) $N \geq 3$, (b) $N \geq 4$, and (c) $N \geq 5$. The solid (dashed) lines correspond to an observed (expected) limit for nominal signal acceptance uncertainty of 5%.

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