

## MIT Open Access Articles

*Thermospheric poleward wind surge at  
midlatitudes during great storm intervals*

The MIT Faculty has made this article openly available. **Please share**  
how this access benefits you. Your story matters.

**Citation:** Zhang, Shun-Rong, Philip J. Erickson, John C. Foster, John M. Holt, Anthea J. Coster, Jonathan J. Makela, John Noto, et al. "Thermospheric Poleward Wind Surge at Midlatitudes during Great Storm Intervals." *Geophysical Research Letters* 42, no. 13 (July 16, 2015): 5132–40.

**As Published:** <http://dx.doi.org/10.1002/2015GL064836>

**Publisher:** American Geophysical Union (AGU)

**Persistent URL:** <http://hdl.handle.net/1721.1/98250>

**Version:** Author's final manuscript: final author's manuscript post peer review, without publisher's formatting or copy editing

**Terms of Use:** Article is made available in accordance with the publisher's policy and may be subject to US copyright law. Please refer to the publisher's site for terms of use.



1 **Thermospheric Poleward Wind Surge at Mid-Latitudes During Great**  
2 **Storm Intervals**

3

4 Shun-Rong Zhang<sup>1,\*</sup>, Philip J. Erickson<sup>1</sup>, John C. Foster<sup>1</sup>, John M. Holt<sup>1</sup>, Anthea J.  
5 Coster<sup>1</sup>, Jonathan J. Makela<sup>2</sup>, John Noto<sup>3</sup>, John W. Meriwether<sup>4</sup>, Brian J. Harding<sup>2</sup>,  
6 Juanita Riccobono<sup>3</sup> and Robert B. Kerr<sup>3</sup>

7

8 *1 MIT Haystack Observatory, Westford, MA, USA*

9 *2 Department of Electrical and Computer Engineering, University of Illinois at*  
10 *Urbana-Champaign, Urbana, Illinois, USA*

11 *3 Scientific Solution, Inc., Chelmsford, MA, USA*

12 *4 Department of Physics and Astronomy, Clemson University, Clemson, SC, USA*

13

14 **\* Correspondence:**

15 Dr. Shunrong Zhang (shunrong@haystack.mit.edu)

16 MIT Haystack Observatory

17 Off Route 40

18 Westford, MA 01886

19 USA

20

21 **Key points**

22 1. Pre-midnight poleward wind surge at 100m/s, preceded by 300m/s westward  
23 winds

24 2. No equatorward wind surge developed throughout the night

25 2. Disturbance wind driven by Subauroral Polarization Streams and Coriolis force

26

27 **Abstract**

28

29 We report a significant poleward surge in thermospheric winds at subauroral and mid  
30 latitudes following the 17-18 March 2015 great geomagnetic storm. This pre-midnight  
31 surge is preceded by strong westward winds. These disturbances were observed over  
32 three sites with geodetic latitudes 35-42°N in the American sector by Fabry-Perot  
33 interferometers at 630-nm wavelength. Prior to the wind disturbances, subauroral  
34 polarization streams (SAPS), were measured by the Millstone Hill incoherent scatter  
35 radar between 20-02 UT. We identify the observed neutral wind variations as driven by  
36 SAPS, through a scenario where strong ion flows cause a westward neutral wind,  
37 subsequently establishing a poleward wind surge due to the poleward Coriolis force on  
38 that westward wind. These regional disturbances appear to have prevented the well-  
39 known storm time equatorward wind surge from propagating into low latitudes, with the  
40 consequence that the classic disturbance dynamo mechanism failed to occur.

41

42

43

## 44 **1. Introduction**

45

46 An equatorward surge of thermospheric meridional neutral wind at mid-latitudes is a  
47 common dynamic feature of Earth's neutral atmosphere during geomagnetic storms.  
48 Numerous studies have shown in detail how storm time momentum and energetic inputs  
49 perturb the high-latitude ionosphere and thermosphere, with subsequent equatorward  
50 propagation of disturbances [see reviews, Matuura, 1972; Prolss, 1995; Buonsanto, 1999;  
51 Mendillo, 2006]. In particular, the storm time equatorward wind generated by impulsive  
52 polar latitude heating processes can be very strong at night (1) when disturbance effects  
53 add to the quiet time equatorward neutral circulation [Rishbeth, 1989], and (2) when  
54 high-latitude convection is enhanced and expanded equatorward causing a significant  
55 anti-sunward ion flow [e.g., Straus and Schulz, 1976]. Due to the Coriolis force, the  
56 equatorward wind surge can subsequently drive a westward neutral wind disturbance as it  
57 reaches mid- and low-latitudes. These wind disturbances are at the heart of the important  
58 ionospheric dynamo effect [Blanc and Richmond, 1980; Fuller-Rowell et al., 2002] and  
59 equatorward wind surges have been the subject of prior studies [e.g., Meriwether, 2008].

60

61 However, equatorward wind surges do not necessarily occur during every storm. In fact,  
62 the disturbance wind undergoes substantial variability with local time, season and solar  
63 activity for a given location. Fejer et al. [2002] showed that at Millstone Hill (MH,  
64 42.6°N, 71.5°W; geodetic), storm time winds sometimes turn poleward following a  
65 midnight or post-midnight equatorward surge (in particular for solar minimum). The pre-  
66 midnight meridional wind disturbance, however, is generally weak and equatorward,  
67 occasionally with a very small and brief poleward turning. The poleward wind surge can  
68 be also seen as traveling atmosphere disturbances (TADs) from the opposite hemisphere  
69 [e.g., Shiokawa et al., 2003].

70

71 This paper reports a fundamentally different neutral wind scenario observed over three  
72 sites at mid and subauroral latitudes during the 17-18 March 2015 great storm. During  
73 this event, the pre-midnight meridional wind turns poleward and remains in this direction  
74 for a few hours. Using a joint analysis of incoherent scatter radar (ISR) and Fabry-Perot  
75 interferometer (FPI) observations along with first-principles model simulations, we show  
76 in this paper that the anomalous poleward wind originated from the Subauroral  
77 Polarization Stream (SAPS) [Foster and Vo, 2002], a characteristic storm time  
78 magnetosphere-ionosphere coupling feature within a relatively narrow region of low  
79 ionospheric conductivity between the auroral precipitation zone and the plasmasphere  
80 boundary layer [Carpenter and Lemaire, 2004]. Strong SAPS flows overlap the  
81 plasmasphere edge and provide a significant convective force moving plasma against  
82 corotation from the dusk sector toward the noontime cusp in storm enhanced density  
83 (SED) features [Foster, 1993; Kelly et al., 2004], also known as the dusk effect  
84 [Mendillo, 2006]. Sunward ion flow in the late afternoon sector (with a poleward  
85 component) also causes SED plasma to drift upwards at subauroral latitudes into regions  
86 where recombination rates are significantly reduced, leading to high electron density  
87 values [Heelis, 2008]. Strong plasma flow has reportedly produced enhanced neutral  
88 westward winds [e.g., Wang H. et al., 2011] due to strong ion drag effects. TIEGCM  
89 simulations [Wang W. et al., 2012] confirm expected ion drag effects in the  
90 thermospheric temperature and zonal winds on a global scale.

91

## 92 **2. Observations and Analysis Procedure**

93

94 Following the arrival of a Coronal Mass Ejection (CME) and under the influence of high-  
95 speed solar wind streams, severe/great geomagnetic disturbances occurred for an  
96 extended period starting on 17 March 2015 (Figures 1a-b). Interplanetary Magnetic Field  
97 (IMF) northward component  $B_z$  hourly values in Geocentric Solar Magnetospheric  
98 (GSM) coordinates fell to -14 nT between 05 - 08 UT, and then underwent a 12-hour  
99 long sustained negative disturbance of  $\sim -17$  nT between 12 - 24 UT on the 17<sup>th</sup>. During  
100 this time, the hourly Dst index dropped to a minimum of -227 nT at 23 UT. The 3-hourly

101 Kp value jumped to 5- with the initial Bz negative excursion, reaching between 8- and 7+  
102 during the large and long-lasting drop in Bz.

103

104 These geospace disturbances caused a series of significant changes in earth's ionosphere  
105 and thermosphere. Figure 1(c) plots TEC disturbances, as derived from global GPS  
106 ground receiver data by MIT Haystack Observatory MAPGPS software [Rideout and  
107 Coster, 2006]. The TEC disturbances are defined here as the mean TEC value in each 3°  
108 latitude bin over 30-50° N geodetic latitudes over the MH longitude (70-80° W) after  
109 undisturbed background TEC has been subtracted using a monthly-average based  
110 empirical model [Chen et al., 2014]. An initial positive disturbance of more than 50% is  
111 seen poleward of MH following the first Bz equatorward turning. A narrow enhancement  
112 zone at 18 UT is visible, and subsequent enhancements extended and expanded from MH  
113 latitudes at 21 UT to lower latitudes at 23 UT. Subsequent sharp decreases of ~ 50%  
114 expanded from 50° N at 20 UT to 37° N within 4 hours.

115

116 This equatorward moving density reduction zone eventually remained at 38-42°N  
117 between 24-05 UT. It is during these pre-midnight hours, and within the midlatitude  
118 trough, that both SAPS and neutral wind disturbances were observed (Section 3). The  
119 presence of a SAPS flow channel is further evidenced by MH ISR observations (Section  
120 3). The ~50% TEC enhancement seen equatorward of the trough is a typical SED  
121 characteristic, and is a result of plasmaspheric erosion associated with SAPS. Later  
122 sections will present these radar observations and will discuss the connection between  
123 SAPS and neutral wind disturbances.

124

125 During this storm, an international observation campaign along the meridian circle of  
126 60°W/120°E longitude was conducted. In particular, excellent ionospheric plasma  
127 information from ISRs and neutral atmosphere information from FPIs is available in the  
128 western hemisphere, along with other observational facilities supported by the Chinese  
129 Meridian Project [Wang, 2010] and other institutions in the eastern hemisphere.

130

131 At mid-latitudes, the MH ISR operated in an experimental sequence using both zenith  
132 and steerable antennas providing local, regional, and wide coverage. The wide coverage  
133 is provided by low ( $6^\circ$ ) elevation radar scans from north to west of Millstone Hill while  
134 regional coverage uses  $45^\circ$  elevation positions in north and west directions combined  
135 with zenith observations. Combining line-of-sight (LOS) ion velocities in these multiple  
136 directions allows derivation of key components of F-region ion velocity perpendicular to  
137 the magnetic field:  $V_{\text{perE}}$ , eastward perpendicular to magnetic field  $\mathbf{B}$ , and  $V_{\text{perN}}$ ,  
138 poleward perpendicular to  $\mathbf{B}$ .  $V_{\text{perE}}$  is determined to the northwest of MH at  $\sim(47^\circ\text{N},$   
139  $89^\circ\text{W})$  from the west-looking data using azimuths in the  $(-100, -45)$  degree range.  $V_{\text{perN}}$   
140 is determined for MH north at  $\sim(53^\circ\text{N}, 75^\circ\text{W})$  from low elevation north-looking data such  
141 that LOS is nearly perpendicular to the magnetic field. These observations provide  
142 evidence of SAPS.

143

144 FPIs observing the thermospheric 630.0-nm emission arising from dissociative  
145 recombination of  $\text{O}_2^+$  provide LOS measurements of neutral winds at typical emission  
146 heights of  $\sim 250$  km altitude. Data from three sites in North America were used in this  
147 study. At MH, the emission is detected from look directions in the north, east, south,  
148 west (all with  $45^\circ$  elevation) and zenith. The other two FPI redline sites are from the  
149 North American Thermosphere-Ionosphere Observing Network (NATION) [Makela et  
150 al., 2014]: Urbana Atmospheric Observatory (UAO,  $40.13^\circ\text{N}, 88.20^\circ\text{W}$ ) and Pisgah  
151 Astronomical Research Institute (PAR,  $35.2^\circ\text{N}, 82.85^\circ\text{W}$ ) [Makela et al., 2011]. They  
152 have similar look angles to MH, with the addition of a pointing direction upward along  $\mathbf{B}$ .  
153 A typical observational cycle takes 12 min for NATION sites (15 min for MH), and data  
154 are analyzed using the methodology described in Harding et al. [2014]. During the 17-18  
155 March night, the skies over PAR became partly cloudy after 0630 UT on 18 March, and  
156 the skies over UAO became partly cloudy after  $\sim 03$  UT, but these facts do not  
157 significantly affect results for the key period addressed by this work. Additionally, the  
158 apparent vertical wind conditions described in Makela et al. [2014] did not seem to occur  
159 during the observations reported here, making the small vertical wind assumption,  
160 typically used when analyzing FPI observations, valid. To estimate the neutral wind

161 vector based on LOS FPI observations, we follow a procedure where the vector velocity  
162 is determined through a least-squared fit algorithm using the combined LOS data from all  
163 directions within 36 min (45 min for MH), weighted by the LOS errors. Standard  
164 deviations on the resulting vector neutral winds are also given.

165

### 166 **3. Superstorm Thermospheric and Ionospheric Response**

167

168 FPI observations during the evening of 17-18 March began just after the main phase of  
169 the storm at the time of maximum ring current intensity / minimum Dst value, as Bz  
170 started to return to positive values. At MH (Figure 2), a strong westward zonal wind of  
171 300 m/s was already established, at a 350 m/s offset from its monthly mean in the  
172 eastward direction. The strong westward wind then weakened after 02 UT, and stayed at  
173 200 m/s for 2 hours. During the next 3 hours, the westward zonal wind eventually  
174 returned to its monthly mean and the disturbance vanished after local midnight. The  
175 meridional wind component started slightly equatorward then subsequently became  
176 significantly poleward, eventually reaching 100 m/s at 0230-0300 UT (~22 LT). The  
177 poleward wind surge was significant between 2030 and 2230 LT, primarily before  
178 midnight. Following this period, the meridional wind returned to its monthly mean value  
179 and remained in line with monthly averages without any further equatorward or poleward  
180 surges.

181

182 FPI observations at UAO at latitudes comparable to MH show very similar features. In  
183 particular, a strong westward zonal wind lasting ~2 hours was followed by a 2-hour long  
184 poleward meridional wind surge, reaching 100 m/s at 0230 UT. The PAR FPI station is  
185 5° to the south of UAO and MH, and the poleward wind abatement, reaching 100 m/s at  
186 0300 UT, is preceded by a strong westward enhancement for ~2 hours. The westward  
187 wind enhancement is not as strong at PAR as at the other two higher latitude sites.

188

189 Prior to the occurrence of these neutral wind disturbances, strong ionospheric  
190 disturbances were observed such as the midlatitude trough, SEDs, and associated SAPS



191 flows. Figure 3 shows  $V_{perE}$  (perpendicular eastward, a) in the west and  $V_{perN}$   
192 (perpendicular northward, b) in the north of MH on 17-18 March as well as later on 19  
193 March, where observations for the latter half of the 18th and first half of the 19th can be  
194 considered representative of normal ionospheric conditions due to relatively low  
195 magnetic activity.

196

197 The magnetically zonal ion drift,  $V_{perE}$ , turned westward at 20 UT. In the next 5 hours,  
198 this westward drift remained very strong at 500-750 m/s. We identified the high-speed  
199 ion flow as SAPS since it is situated near the low density (midlatitude trough) region as  
200 evidenced in TEC (Figure 1c) and also in the ISR measured electron content IEC (Figure  
201 3d). SAPS features are also associated with SED passage over MH, prior to the SAPS  
202 local onset, as shown in the TEC plume equatorward of the trough (Figure 1c) and in the  
203 IEC peak at 20UT (Figure 3c). Since the ion speed is faster than the speed of the neutrals,  
204 SAPS will accordingly drive the neutrals in the same direction. MH and UAO zonal  
205 winds in Figure 2 are shown in Figure 3(e) to highlight this neutral and ion velocity  
206 connection. Even though the exact onset time of the westward neutral wind  
207 enhancements are unknown due to lack of data prior to 00 UT, the enhancements  
208 weakened around 02-03UT when the large westward ion drift disappeared. The largest  
209 westward wind surge ( $\sim 400$  m/s) is at UAO where the TEC drop appears the most  
210 significant, and therefore SAPS peak velocities are expected to be the strongest.

211

212 The neutral wind disturbance, and in particular the poleward surge, can be seen in the  
213 upward ion drift  $V_z$  at 250 km (Figure 3c). Between 01-05 UT on 18 March when  
214 electric fields are quiet,  $V_z$  is more negative (downward) compared to the reference,  
215 suggesting that the poleward wind surge makes a significant contribution.

216

217 The radar data shows also a strong poleward ion drift  $V_{perN}$  at  $\sim 55^\circ N$ , in particular,  
218 between 20-23 UT. Combining these northward and westward components, it is likely  
219 that the sunward drift is a significant factor in producing the large observed SED plumes  
220 in TEC and IEC data.

221

## 222 **4. Simulation of Wind Effects**

223

### 224 **4.1 March 17th Event Summary**

225 The observations described in Section 3 are consistent with the following timeline of  
226 ionosphere and thermosphere processes during the 17-18 March superstorm: (1) At ~21  
227 UT on 17 March, SED plumes are present, observed as TEC and IEC enhancements,  
228 along with onset of strong SAPS westward ion drift of >500 m/s peaking in the  
229 midlatitude trough of the American sector; (2) Starting at about 22 UT, a strong westward  
230 neutral wind appears (with greater amplitude in the trough region of ~40 - 42°N latitude  
231 as compared to 35°N) and lasts for 4 hours until 02 UT on 18 March. After 02 UT, the  
232 zonal wind amplitude decreased dramatically (at higher latitudes) or gradually (at lower  
233 latitudes); and (3) At 03 UT, a poleward wind surge of ~100 m/s occurs. We posit that  
234 the chain of events from (1) to (3) are a result of ion-neutral coupling and thermospheric  
235 dynamical processes. In particular, the westward neutral wind characteristics described in  
236 (2) are attributable to the well-known ion drag effect following the onset of SAPS  
237 described in (1). Our proposition is further indicated by the observation that westward  
238 neutral wind anomalies disappeared when SAPS forcing disappeared, with the effects  
239 strongest in the area of the deepest midlatitude trough. Subsequently, the poleward wind  
240 surge described in (3) is produced following strong westward winds as described in (2)  
241 because of the poleward Coriolis force arising from significant westward wind  
242 amplitudes.

243

### 244 **4.2 Simulated F-region westward ion drift effects**

245 To explore the validity of this scenario, an ion-neutral coupling simulation was conducted  
246 with parameters characteristic of mid-latitudes near Millstone Hill. This relatively simple  
247 simulation was designed to study the basic proposed mechanism and its general features,  
248 rather than reproducing all observational details. The numerical experiment utilizes a  
249 local ionosphere-thermosphere wind coupling model. The ionospheric model [Zhang and  
250 Huang, 1995; Zhang et al., 2003] solves the equation of mass continuity for multiple ions

251 and the equation of momentum for  $O^+$ . Ionospheric temperatures are set to the empirical  
252 ISR Ionospheric Model (ISRIM) for MH [Zhang et al., 2005; 2007b]. The non-  
253 disturbance electric field is specified by the local electric fields derived from the  
254 statistical Millstone Hill – Sondrestrom convection model [Zhang et al., 2007a]. Neutral  
255 densities and temperatures are specified by the NRL-MSIS model [Picone et al., 2002].  
256 The neutral wind calculation is based on solving the momentum equation for the neutral  
257 gas where pressure gradients are derived from the NRL-MSIS model. A resolution of 2-  
258 min in time and 2-km in altitude is used. Other details of the simulation are described

259

260 A baseline / reference run for medium solar activity levels at equinox, corresponding to  
261 17 March 2015 conditions, was performed without imposing the observed westward ion  
262 drift. A second run included a 500 m/s westward ion drift imposed during 18-21 LT (23-  
263 02 UT). This 500 m/s value represents a modest enhancement in zonal ion velocity  
264 compared to the ISR measured SAPS speed on 17 March (cf. Section 3). The top panel of  
265 Figure 4 plots the calculated model profiles of electron density and neutral winds at 1930  
266 LT, 1.5 hours into the strong westward ion drift injection event. The middle panel shows  
267 differential changes in calculated poleward and eastward winds relative to the reference  
268 run. The differential results clearly show a build-up of westward winds within one hour  
269 following SAPS turn-on, saturating at 300 m/s at ~1930 LT. This ~200 m/s difference in  
270 flow speeds between ions and neutrals created frictional heating leading to neutral  
271 temperature increases (not shown). The poleward wind showed a different behavior,  
272 gradually increasing in speed but at slower rates compared to westward wind increases.  
273 In particular, northern wind differential values peaked at 130 m/s at ~2130 LT, with a  
274 delay of ~2 hours following westward neutral flow saturation and ~3.5 hours following  
275 SAPS initiation. These time constants are governed by background thermospheric density  
276 and temperature, which undergo substantial storm time changes (and therefore the real  
277 response times can be quite different from simulated results.) In the simulation, the  
278 Coriolis force is the only mechanism that connects the zonal and meridional winds, so the  
279 key finding is that poleward wind buildup is caused by Coriolis force effects in the north

280 direction because of westward neutral motion, with the characteristic time delay in storm  
281 time variations between the two components as another sign of Coriolis forcing.

282

283 The third panel of Figure 4 plots ionospheric electron density differential response to  
284 strong westward ion drift. The peak height, hmF2, dropped by ~55 km during the period  
285 of significant poleward wind, and NmF2 dropped by ~80% after the westward ion drift  
286 ceased. In this numerical experiment, the electron density reduction is achieved by the  
287 induced poleward winds only, but an increase in the ion recombination rate caused by  
288 frictional heating can further reduce the F-region electron density [Schunk et al., 1975].

289

## 290 **5. Discussion**

291

292 Reproducing all observational details with the numerical experiment in Section 4 is not  
293 possible because of the simulation's simplified treatment of ion and neutral temperatures  
294 and neutral densities, which in reality are likely subject to large offsets from the empirical  
295 model specifications used in this storm study. In particular, we have ignored the potential  
296 presence of meridional pressure gradients to the north of MH, produced by the same  
297 auroral heating processes that likely would generate an equatorward wind surge. The  
298 simulated poleward surge produced by Coriolis force action is smooth and gradual in  
299 time. This is quite similar to FPI observations at the lower latitude site, PAR. By  
300 contrast, at the higher latitude of MH the observed poleward surge grows faster than in  
301 the simulation. A more precise neutral response simulation would require a better  
302 specification of storm time neutral density and temperature. In general, more  
303 sophisticated modeling is needed to put the regional observations into correct global  
304 context and to better explain observational features. However, the initial simplified  
305 simulation presented here is sufficient to demonstrate that ion-neutral coupling and  
306 Coriolis force effects are likely to play fundamental roles in observed wind dynamics.

307

308 Coriolis force effects on westward neutral wind amplitude were noted in an earlier  
309 numerical experiment by Forbes and Roble [1990]. Hagan and Sipler [1991] also report a

310 similar set of observed ionospheric and thermospheric storm time effects with smaller  
311 westward wind and minor meridional wind abatement during the 7-10 March 1989 storm  
312 at MH. They pointed out that westward winds caused diminishment of equatorward wind  
313 surges due to Coriolis forces. Hernandez et al. [1982] reported a converging wind system  
314 measured by the FPI over Fritz Park (39.9°N, 105.5°W) during two major storms where  
315 similar ion-neutral interaction was speculated for the south of the site. It appears that  
316 large westward winds, the poleward turning of the wind, and low electron densities are all  
317 potential consequences of ion-neutral coupling in some large storm events such as  
318 presented here.

319

320 During large geospace storm intervals, TADs are often observed, and in fact poleward  
321 winds observed at lower mid-latitudes in Asia were explained as resulting from TAD  
322 passage with a disturbance source region in the south [e.g., Shiokawa et al., 2003]. This  
323 study, however, does not show obvious signs of propagating waves as in the TAD  
324 scenario. In particular, the poleward wind surge and westward wind enhancement appear  
325 stronger at subauroral latitudes than at lower latitudes, counter to expectations of a TAD  
326 mechanism to the south generating atmospheric waves. TADs from sources to the north  
327 would cause equatorward wind perturbations, not poleward as observed.

328

## 329 **6. Summary and conclusion**

330

331 Strong geospace disturbances during the 17-18 March 2015 superstorm produced intense  
332 electric fields at subauroral and mid-latitudes. SED plumes shown in TEC and ISR  
333 electron content enhancements over the northeast US are present prior to 21 UT before  
334 passage of the midlatitude trough. A strong westward ion drift, identified as SAPS, then  
335 developed during 21-02 UT, as observed by the Millstone Hill ISR. This drift, drove  
336 neutral particles westward, causing a strong westward neutral wind (~300 m/s) observed  
337 by multiple FPIs between 35-42°N latitudes in the American sector. Later in the event, a  
338 poleward neutral wind response occurred due to Coriolis force effects on the westward  
339 neutral wind. The poleward wind, directly observed by the FPIs, eventually reached 100

340 m/s amplitude in a few hours following the onset of SAPS. A simplified numerical model  
341 with coupling of ionospheric density and thermospheric neutral winds is able to  
342 demonstrate response characteristics of this mechanism that reasonably match general  
343 observational features. We conclude that unusual neutral wind disturbances in both zonal  
344 and meridional directions have their ultimate source in the SAPS electric field which  
345 generates substantial ion flows and leads to strong ion-neutral interaction.

346

347 We point out that the observed pre-midnight storm time surge in the poleward wind  
348 implies a regional circulation background that may prevent propagation of the auroral  
349 heating produced equatorward wind surge to lower latitudes, and would therefore prevent  
350 development of the classic disturbance dynamo. If correct, this mechanism may  
351 substantially influence the storm time low latitude and equatorial electrodynamic  
352 response in some cases [Fejer and Scherliess, 1995; Maruyama et al., 2005; Lu et al.,  
353 2012]. Thus, further investigation is needed to characterize westward ion drifts at  
354 subauroral latitudes, including SAPS, as a significant driving factor, and to understand  
355 more fully how the SAPS driven neutral wind disturbances are affected by storm-time  
356 ionosphere and thermosphere conditions.

357

### 358 **Acknowledgements**

359

360 Ground-based observational data from FPIs and the Millstone Hill ISR are all obtained  
361 through the Madrigal database system ([www.openmadrigal.org](http://www.openmadrigal.org)) developed and  
362 maintained by MIT Haystack Observatory. Dst and Kp data are obtained from the World  
363 Data Center at Kyoto University (<http://wdc.kugi.kyoto-u.ac.jp/wdc/Sec3.html>). The  
364 solar wind data are obtained from the Advanced Composition Explorer (ACE) satellite,  
365 and the real-time 1-hour average IMF Bz data sets are obtained from the archive at the  
366 NOAA Space Weather Prediction Center (<ftp://ftp.swpc.noaa.gov/pub/lists/ace2/>). ISR  
367 observations and analysis at MIT Haystack Observatory are supported by cooperative  
368 agreement AGS-1242204 between the National Science Foundation and the  
369 Massachusetts Institute of Technology. SRZ and PJE also acknowledge NASA LWS

370 funding support (NNX15AB83G). FPI measurements at Millstone Hill are supported by  
371 NSF grant AGS-0640745 to the University Corporation for Atmospheric Research, and at  
372 UAO and PAR by National Science Foundation CEDAR grants AGS-1138998 and AGS-  
373 1138931 to the University of Illinois and Clemson University, respectively. SRZ is  
374 indebted to Wenbin Wang, Tim Fuller-Rowell, and members of the NASA LWS topside  
375 team for valuable comments on initial results of this research. GPS TEC data products  
376 and access through the Madrigal distributed data system are provided to the community  
377 by the Massachusetts Institute of Technology under support from US National Science  
378 Foundation grant AGS-1242204. Data for the TEC processing is provided from the  
379 following organizations: UNAVCO, Scripps Orbit and Permanent Array Center, Institut  
380 Geographique National, France, International GNSS Service, The Crustal Dynamics Data  
381 Information System (CDDIS), National Geodetic Survey, Instituto Brasileiro de  
382 Geografia e Estatística, RAMSAC CORS of Instituto Geográfico Nacional del la  
383 República Argentina, Arecibo Observatory, Low-Latitude Ionospheric Sensor Network  
384 (LISN), Topcon Positioning Systems, Inc., Canadian High Arctic Ionospheric Network,  
385 Institute of Geology and Geophysics (Chinese Academy of Sciences), China  
386 Meteorology Administration, Centro di Ricerche Sismologiche, Système d'Observation du  
387 Niveau des Eaux Littorales (SONEL), RENAG : REseau NAtional GPS permanent, and  
388 GeoNet - the official source of geological hazard information for New Zealand.

389

390

391

392

393

394 **References**

395

396 Blanc, M., and A. D. Richmond (1980), The ionospheric disturbance dynamo, *J.*

397 *Geophys. Res.*, 85, 1669–1686.

398

399 Buonsanto, M. J. (1999), Ionospheric storms—A review, *Space Sci. Rev.*, 88, 563–601.

400

401 Buonsanto, M. J., Foster, J. C., Galasso, A. D., Sipler, D. P., Holt, J. M. (1990) Neutral

402 Winds and Thermosphere/Ionosphere Coupling and Energetics During the Geomagnetic

403 Disturbances of March 6–10, 1989. *J. Geophys. Res.* 95: pp. 21033-21050

404

405 Carpenter, D. and J. Lemaire (2004), The plasmasphere boundary layer, *Ann. Geophys.*

406 22, 4291–4298.

407

408 Chen, Z., S.-R. Zhang, A. J. Coster, and G. Fang (2015), EOF analysis and modeling of

409 GPS TEC climatology over North America, *J. Geophys. Res. Space Physics*, 120,

410 doi:10.1002/2014JA020837.

411

412 Fejer, B. G. (2002), Climatology and storm time dependence of nighttime thermospheric

413 neutral winds over Millstone Hill, *J. Geophys. Res.*, 107(A5), 1052,

414 doi:10.1029/2001JA000300.

415

416 Fejer, B. G., and L. Scherliess (1995), Time dependent response of equatorial ionospheric

417 electric fields to magnetospheric disturbances, *Geophys. Res. Lett.*, 22(7), 851–854,

418 doi:10.1029/95GL00390.

419

420 Foster, J. C. (1993), Storm time plasma transport at middle and high latitudes, *J.*

421 *Geophys. Res.*, 98, 1675–1690.

422

423 Foster, J. C., and H. B. Vo (2002), Average Characteristics and Activity Dependence of

424 the Subauroral Polarization Stream, *J. Geophys. Res.*, 107(A12), 1475, doi:

425 10.1029/2002JA009409.

426

427 Forbes, J. M., and R. G. Roble (1990), Thermosphere-ionosphere coupling: An

428 experiment in interactive modeling, *J. Geophys. Res.*, 95(A1), 201,

429 doi:10.1029/JA095iA01p00201.

430

431 Fuller-Rowell, T. J., G. H. Millward, A. D. Richmond, and M. V. Codrescu (2002),

432 Storm-time changes in the upper atmosphere at low latitudes, *J. Atmos. Sol. Terr. Phys.*,

433 64, 1383.

434



435 Hagan, M. E., and D. P. Sipler (1991), Combined incoherent scatter radar and Fabry-  
436 Perot interferometer measurements of frictional heating effects over Millstone Hill during  
437 March 7–10, 1989, *J. Geophys. Res.*, 96(A1), 289, doi:10.1029/90JA02250.  
438

439 Harding, B. J., T. W. Gehrels, and J. J. Makela (2014), Nonlinear regression  
440 method for estimating neutral wind and temperature from Fabry–Perot  
441 interferometer data, *Appl. Opt.*, 53(4), 666.

442 Heelis, R. A. (2008) Low- and Middle-Latitude Ionospheric Dynamics Associated with  
443 Magnetic Storms, in *Midlatitude Ionospheric Dynamics and Disturbances* (eds P. M.  
444 Kintner, A. J. Coster, T. Fuller-Rowell, A. J. Mannucci, M. Mendillo and R. Heelis),  
445 American Geophysical Union, Washington, D. C.. doi: 10.1029/181GM06

446 Hernandez, G., R. G. Roble, E. C. Ridley, and J. H. Allen (1982), Thermospheric  
447 response observed over Fritz Peak, Colorado, during two large geomagnetic storms near  
448 Solar Cycle Maximum, *J. Geophys. Res.*, 87(A11), 9181-9192, doi:10.1029/JA087.  
449

450

451 Lu, G., L. Goncharenko, M. J. Nicolls, A. Maute, A. Coster, and L. J. Paxton (2012),  
452 Ionospheric and thermospheric variations associated with prompt penetration electric  
453 fields, *J. Geophys. Res.*, 117, A08312, doi:10.1029/2012JA017769.  
454

455 Makela, J. J., J. W. Meriwether, Y. Huang, and P. J. Sherwood (2011), Simulation and  
456 analysis of a multi-order imaging Fabry–Perot interferometer for the study of  
457 thermospheric winds and temperatures, *Appl. Opt.*, 50, 4403-4416.  
458

459 Makela, J. J., et al. (2014), Storm time response of the midlatitude thermosphere:  
460 Observations from a network of Fabry-Perot interferometers, *J. Geophys. Res. Space*  
461 *Physics*, 119, 6758–6773, doi:10.1002/2014JA019832.  
462

463 Matuura, N. (1972), Theoretical models of ionospheric storms, *Space Sci. Rev.*, 13, 124–  
464 189.  
465

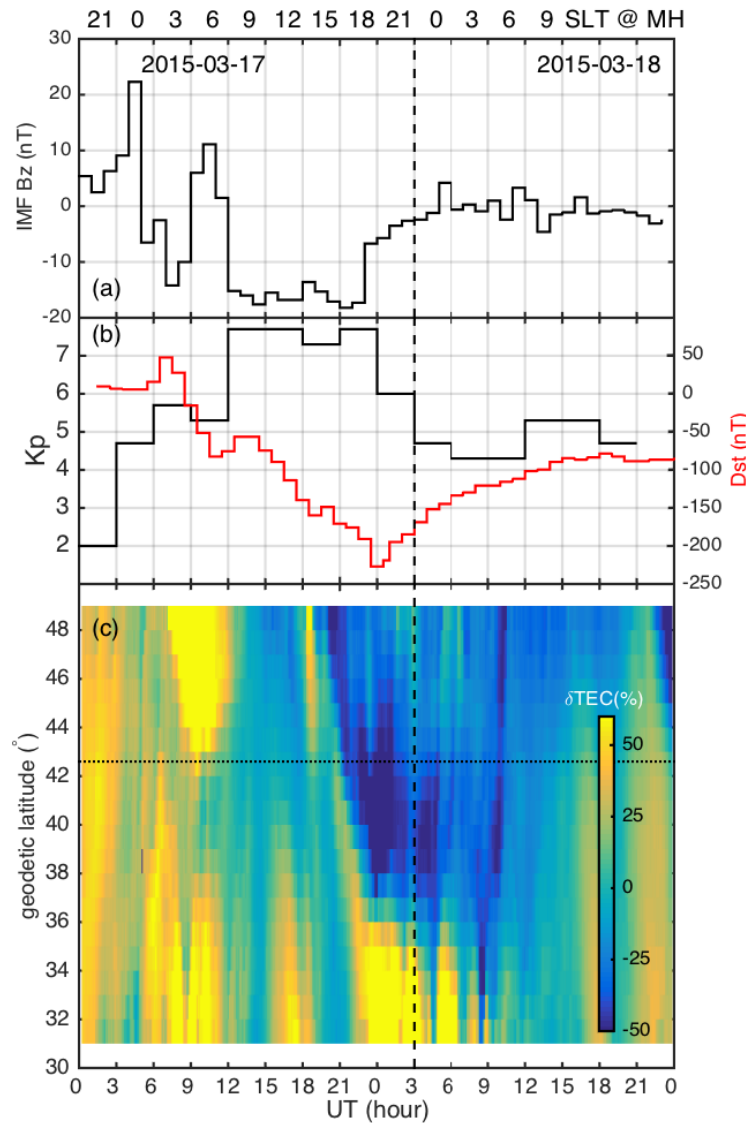
466 Mendillo, M. (2006), Storms in the ionosphere: Patterns and processes for total electron  
467 content, *Rev. Geophys.*, 44, RG4001, doi:10.1029/2005RG000193.  
468

469 Maruyama, N., A. D. Richmond, T. J. Fuller-Rowell, M. V. Codrescu, S. Sazykin, F. R.  
470 Toffoletto, R. W. Spiro, and G. H. Millward (2005), Interaction between direct  
471 penetration and disturbance dynamo electric fields in the storm-time equatorial  
472 ionosphere, *Geophys. Res. Lett.*, 32, L17105, doi:10.1029/2005GL023763.  
473

474 Picone, J. M., A. E. Hedin, D. P. Drob, and A. C. Aikin, NRLMSISE-00 empirical model  
475 of the atmosphere: Statistical comparisons and scientific issues, *J. Geophys. Res.*,  
476 107(A12), 1468, doi:10.1029/2002JA009430, 2002.

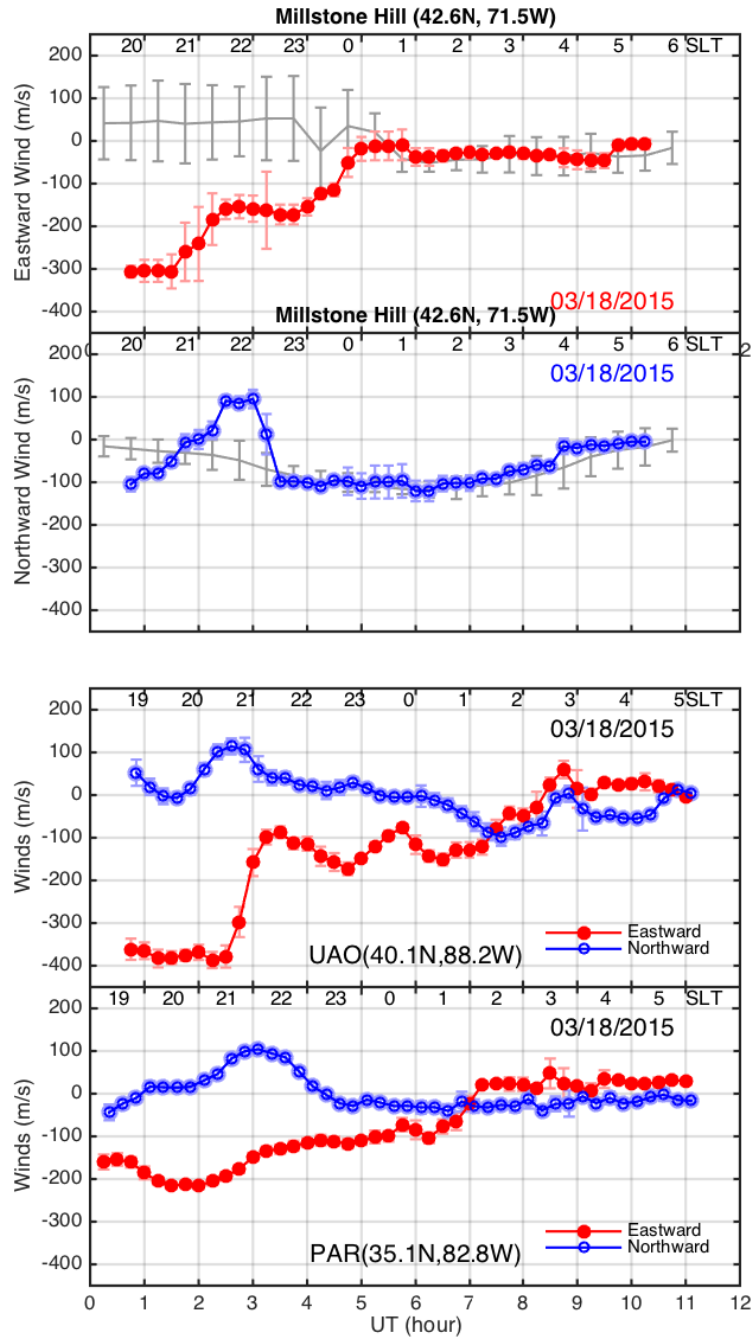
477  
478 Prölss, G. W. (1995), Ionospheric F-region storms, in Handbook of Atmospheric  
479 Electrodynamics, vol. 2, edited by H. Volland, chap. 8, pp. 195–248, CRC Press, Boca  
480 Raton, Fla.  
481  
482 Rideout, W., and A. Coster (2006), Automated GPS processing for global total electron  
483 content data, *GPS Solut.*, 10, 219–228, doi:10.1007/s10291-006-0029-5.  
484  
485 Rishbeth, H. (1998), How the thermospheric circulation affects the ionospheric F2-layer,  
486 *J. Atmo. Sol.-Terr. Phys.*, 60(14), 1385–1402.  
487  
488 Schunk, R. W., W. J. Raitt, and P. M. Banks (1975), Effect of electric fields on the  
489 daytime high-latitude E and F regions, *J. Geophys. Res.*, 80, 3121.  
490  
491 Shiokawa, K., et al. (2003), Thermospheric wind during a storm-time large-scale  
492 traveling ionospheric disturbance, *J. Geophys. Res.*, 108, 1423,  
493 doi:10.1029/2003JA010001, A12.  
494  
495 Straus, J.M., Schulz, M. (1976), Magnetospheric Convection and Upper Atmospheric  
496 Dynamics'. *J. Geophys. Res.* 81, 5822-5832.  
497  
498 Wang, C. (2010), New Chains of Space Weather Monitoring Stations in China. *Space*  
499 *Weather*, 8: n/a. doi: 10.1029/2010SW000603  
500  
501 Wang, W., E. R. Talaat, A. G. Burns, B. Emery, S. Hsieh, J. Lei, and J. Xu (2012),  
502 Thermosphere and ionosphere response to subauroral polarization streams (SAPS):  
503 Model simulations, *J. Geophys. Res.*, 117, A07301, doi:10.1029/2012JA017656.  
504  
505 Wang, H., H. Lühr, K. Häusler, and P. Ritter (2011), Effect of subauroral polarization  
506 streams on the thermosphere: A statistical study, *J. Geophys. Res.*, 116, A03312,  
507 doi:10.1029/2010JA016236.  
508  
509 Zhang S.-R. and X.-Y. Huang (1995), A Numerical Study of Ionospheric Profiles for  
510 Mid-Latitudes, *Ann Geophysicae*, 13, 551-557.  
511  
512 Zhang, S.-R. W. L. Oliver, J. M. Holt, and S. Fukao (2003), Ionospheric data  
513 assimilation: Comparison of extracted parameters using full density profiles and key  
514 parameters, *J. Geophys. Res.* 108, 10.1029/2002JA009521.  
515  
516 Zhang, S.-R., J. M. Holt, A. P. van Eyken, M. McCready, C. Amory-Mazaudier, S.  
517 Fukao, and M. Sulzer (2005), Ionospheric local model and climatology from long-term  
518 databases of multiple incoherent scatter radars, *Geophys. Res. Lett.*, 32, L20102,  
519 doi:10.1029/2005GL023603  
520

521 Zhang, S.-R. and J. M. Holt (2007a), Ionospheric Climatology and Variability from  
522 Long-term and Multiple Incoherent Scatter Radar Observations: Climatology in Eastern  
523 American Sector, *J. Geophys. Res.*, 12, A06328, doi:10.1029/2006JA012206  
524  
525 Zhang, S.-R., J. M. Holt, and M. McReady (2007b), High latitude convection model  
526 based on long-term incoherent scatter radar observations in North America, *J. Atmos.*  
527 *Sol.-Terr. Phys.*, 69, 1273-1291, doi:10.1016/j.jastp.2006.08.017  
528



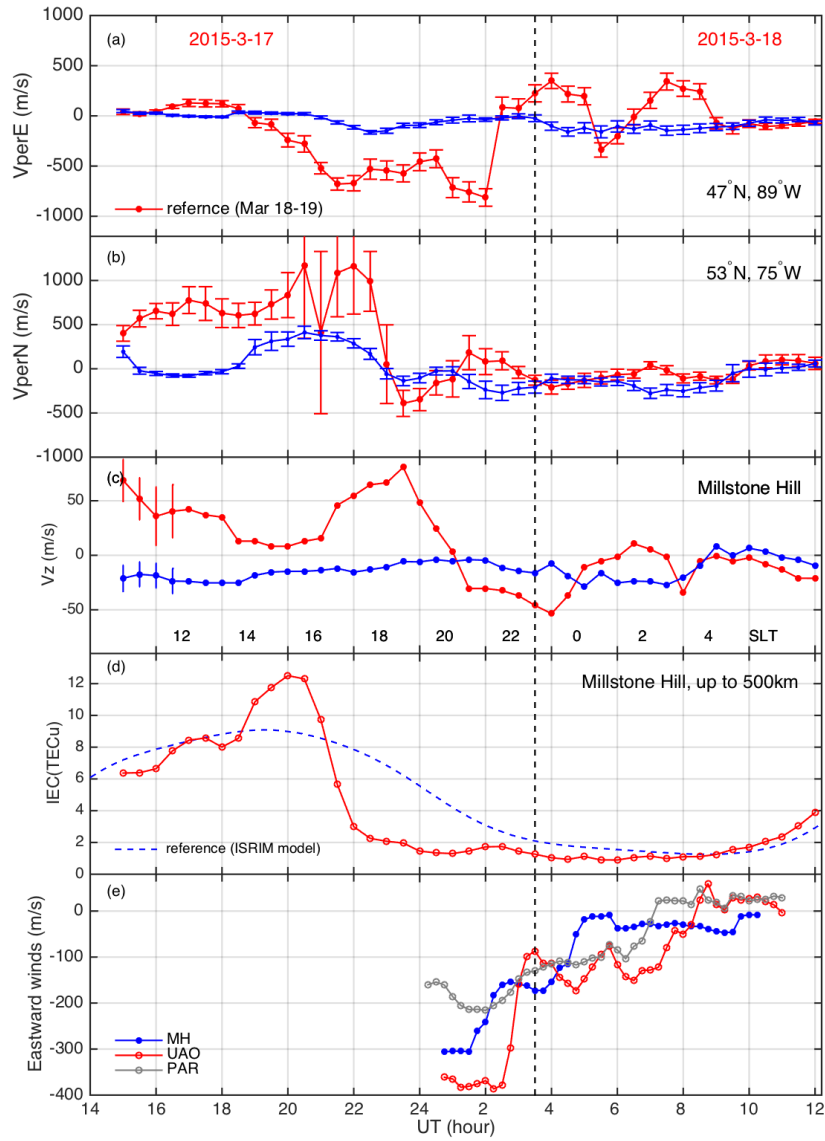
530  
 531  
 532  
 533  
 534  
 535  
 536  
 537

**Figure 1.** Solar geophysical conditions during March 17-18, 2015: hourly IMF Bz (a), 3-hourly Kp and hourly Dst (b), and GPS TEC disturbances (%) at mid- and subauroral latitudes (c). The TEC is a 3° latitudinal bin average over Millstone Hill longitudes (70-80°W), with non-disturbance background subtracted using the monthly average based North America TEC model [NATEC, Chen et al., 2015]. The dashed line is the approximate time of the poleward wind surge.

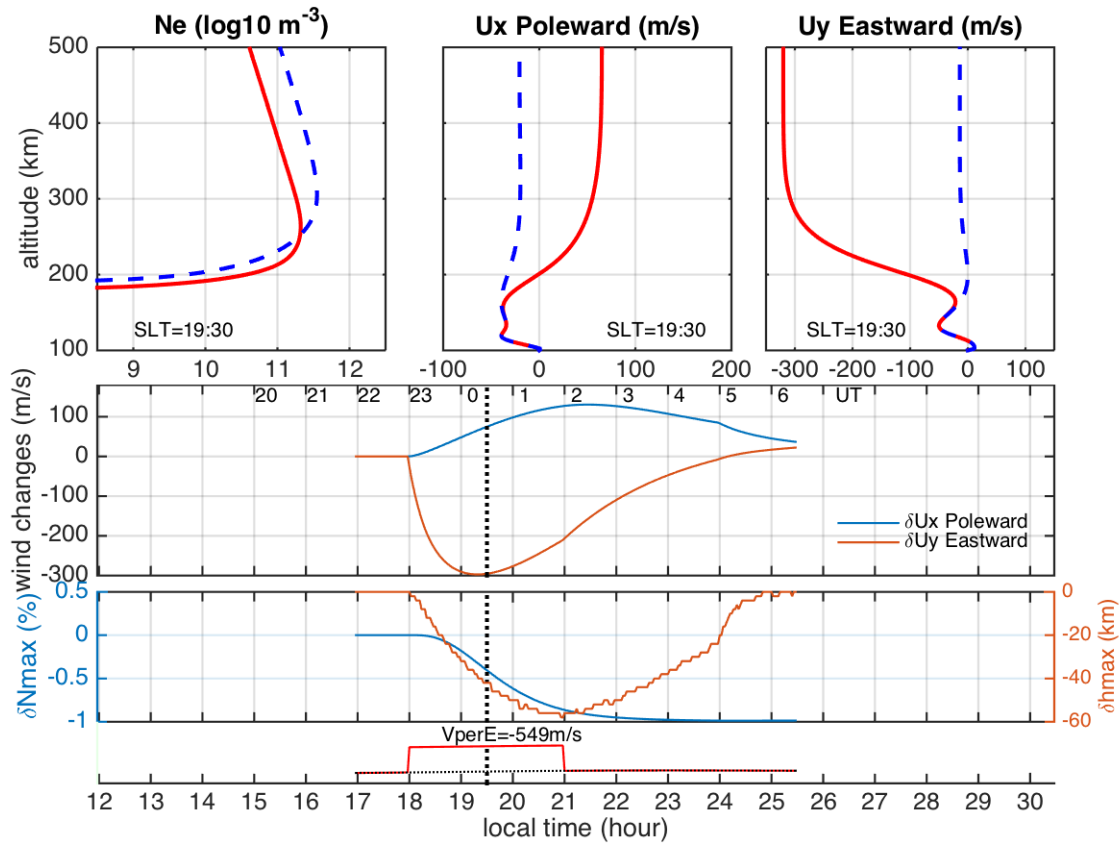


538  
 539  
 540  
 541  
 542

**Figure 2.** FPI redline neutral winds measured over Millstone Hill (top two panels, gray lines being monthly average and corresponding standard deviation; red and blue curves are winds for 18 March); over UAO and PAR (bottom two panels).



543  
544 **Figure 3.** Millstone Hill ISR measurements of plasma drifts  $V_{perE}$  (perpendicular east, a) for  
545  $\sim(89^{\circ}W, 47^{\circ}N)$ ,  $V_{perN}$  (perpendicular north, b) for  $\sim(75^{\circ}W, 53^{\circ}N)$ , vertical upward ion drift  $V_z$  at  
546 250 km above MH (c), and integrated ionospheric content IEC (up to 500 km, d). Eastward winds from Figure 2 are also given in (d). The red curves in (a)-(d) are for 17-18 March, and the blue  
547 ones in (a) - (c) show reference drift patterns from 18-19 March observations. The dashed blue  
548 curve in (d) is a quiet-time reference IEC variation calculated using the ISR empirical ionospheric  
549 model ISIRIM [Zhang et al., 2007a]. The vertical dashed line indicates the approximate peak time  
550 of poleward wind surge.  
551



552

553 **Figure 4.** A numerical experiment demonstrates the mechanism of westward ion drift inducing  
 554 poleward neutral winds over Millstone Hill. A westward ion drift of 500 m/s is applied over 18-  
 555 20 LT (23-01UT). Profiles on the top panels show results with (red) and without (blue dashed)  
 556 the strong westward ion drift at 1930 UT for electron density, poleward winds and eastward  
 557 winds. The middle panel shows corresponding changes in poleward and eastward winds due to  
 558 inclusion of the strong westward ion drift. The third panel plots simulated changes in the peak  
 559 electron density (blue) and the peak density height (red). The bottom shows the timing of the  
 560 injected 500 m/s westward ion drift.