Inclusive Search for Squark and Gluino Production in \( p(p)\overline{p} \) Collisions at \( \sqrt{s} = \) TeV

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Inclusive Search for Squark and Gluino Production in $p\bar{p}$ Collisions at $\sqrt{s} = 1.96$ TeV


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We report on a search for inclusive production of squarks and gluinos in $p\bar{p}$ collisions at $\sqrt{s} = 1.96$ TeV, in events with large missing transverse energy and multiple jets of hadrons in the final state. The study uses a CDF Run II data sample corresponding to $2fb^{-1}$ of integrated luminosity. The data are in good agreement with the standard model predictions, giving no evidence for any squark or gluino component. In an $R$-parity conserving minimal supergravity scenario with $A_0 = 0$, $\tan\beta = 5$, 95% C.L. upper limits on the production cross sections in the range between 0.1 and 1 pb are obtained, depending on the squark and gluino masses considered. For gluino masses below 280 GeV/c$^2$, arbitrarily
Supersymmetry (SUSY) [1] is regarded as a possible extension of the standard model (SM) that naturally solves the hierarchy problem and provides a possible candidate for dark matter in the Universe. SUSY introduces a new symmetry that relates fermionic and bosonic degrees of freedom, and doubles the SM spectrum of particles by introducing a new supersymmetric partner (sparticle) for each particle in the SM. Results on similar inclusive searches for SUSY using Tevatron data have been previously reported by both the CDF and D0 experiments in Run I [2] and by the D0 experiment in Run II [3]. This Letter presents new results on an inclusive search for squarks and gluinos, supersymmetric partners of quarks and gluons, based on data collected by the CDF experiment in Run II and corresponding to 2.0 fb\(^{-1}\) of integrated luminosity. The analysis is performed within the framework of minimal supergravity (mSUGRA) [4] and assumes R-parity conservation where sparticles are produced in pairs and the lightest supersymmetric particle (LSP) is stable, neutral, and weakly interacting. The expected signal is characterized by the production of multiple jets of hadrons from the cascade decays of squarks and gluinos and large missing transverse energy \(E_T\) [5] from the presence of two LSPs in the final state. In a scenario with squark masses \(M_{\tilde{q}}\) significantly larger than the gluino mass \(M_{\tilde{g}}\), at least four jets in the final state are expected, while for \(M_{\tilde{g}} > M_{\tilde{q}}\) dijet configurations dominate. Separate analyses are carried out for events with at least two, three, and four jets in the final state and with different requirements on the minimum \(E_T\).

The CDF II detector is described in detail elsewhere [6]. The detector has a charged particle tracking system that is immersed in a 1.4 T solenoidal magnetic field coaxial with the beam line, and provides coverage in the pseudorapidity \(|\eta| \leq 2\). Segmented sampling calorimeters, arranged in a projective tower geometry, surround the tracking system and measure the energy of interacting particles for \(|\eta| < 3.6\). Cherenkov counters in the region \(3.7 < |\eta| < 4.7\) measure the number of inelastic \(p\bar{p}\) collisions to determine the luminosity [7].

Samples of simulated QCD-jets, \(t\bar{t}\) production, and diboson (WW, ZW, and ZZ) processes are generated using the PYTHIA 6.216 [8] Monte Carlo generator with Tune A [9]. The normalization of the QCD-jets sample is extracted from data in a low \(E_T\) region, while \(t\bar{t}\) and diboson samples are normalized to next-to-leading order (NLO) predictions [10,11]. Samples of simulated Z/\(\gamma\) + jets and W + jets events are generated using the ALPGEN 2.1 program [12] where exclusive subsamples with different jet multiplicities are combined, and the resulting samples are normalized to the measured Z and W inclusive cross sections [13]. Finally, samples of single top events are produced using the MADEVENT program [14] and normalized using NLO predictions [15]. In mSUGRA, the mass spectrum of sparticles is determined by five parameters: the common scalar and gaugino masses at the GUT scale, \(M_0\) and \(M_1/2\), respectively; the common trilinear coupling at the GUT scale, \(A_0\); the sign of the Higgsino mixing parameter, \(\mu\); and the ratio of the Higgs vacuum expectation values, \(\tan\beta\). The mSUGRA samples are generated using the ISASUGRA implementation in PYTHIA with \(A_0 = 0, \mu < 0\), and \(\tan\beta = 5\), as inspired by previous studies [16]. A total of 132 different squark and gluino masses are generated via variations of \(M_0\) and \(M_1/2\) in the range \(M_0 < 600\text{ GeV}/c^2\) and \(50 < M_1/2 < 220\text{ GeV}/c^2\). At low \(\tan\beta\), the squarks from the first two generations are nearly degenerate, whereas the mixing of the third generation leads to slightly lighter sbottom masses and much lighter top squark masses. In this analysis, top squark pair production processes are not considered. The contribution from hard processes involving sbottom production is almost negligible, and is not included in the calculation of the signal efficiencies to avoid a dependency on the details of the model for squark mixing. The mSUGRA samples are normalized using NLO cross sections as determined by PROSPINO 2.0 [17], with input parameters provided by ISAJET 7.74 [18]. CTEQ61M parton distribution functions (PDFs) [19] are used, and renormalization and factorization scales are set to the average mass [20] of the sparticles produced in the hard interaction. The Monte Carlo events are passed through a full CDF II detector simulation (based on GEANT3 [21] and GFLASH [22]) and reconstructed and analyzed with the same analysis chain as for the data.

Data are collected using a three-level trigger system that selects events with \(E_T > 35\text{ GeV}\) and at least two calorimeter clusters with \(E_T\) above 10 GeV. The events are then required to have a primary vertex with a \(z\) position within 60 cm of the nominal interaction. Jets are reconstructed from the energy deposits in the calorimeter towers using a cone-based jet algorithm [23] with cone radius \(R = \sqrt{\Delta \eta^2 + \Delta \phi^2} = 0.7\), and the measured \(E_T^{\text{jet}}\) is corrected for detector effects and contributions from multiple \(p\bar{p}\) interactions per crossing at high instantaneous luminosity, as discussed in Ref. [24]. The events are required to have at least two, three, or four jets (depending on the final state considered), each jet with corrected transverse energy \(E_T^{\text{jet}} > 25\text{ GeV}\) and pseudorapidity in the range \(|\eta^{\text{jet}}| < 2.0\), and at least one of the jets is required to have \(|\eta^{\text{jet}}| < 1.1\). Finally, the events are required to have \(E_T > 70\text{ GeV}\). For the kinematic range in \(E_T\) and the \(E_T^{\text{jet}}\) of the jets considered in this analysis, the trigger selection is 100% efficient. Beam-related backgrounds and cosmic rays

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are removed by requiring an average jet electromagnetic fraction $f_{em} = \sum_{jets} E_{T, em}^{jet}/\sum_{jets} E_{T}^{jet} > 0.15$, where $E_{T, em}^{jet}$ denotes the electromagnetic component of the jet transverse energy, and the sums run over all the selected jets in the event. In addition, the events are required to have an average charged particle fraction $f_{ch} = \sum_{jets} p_T^{track}/\sum_{jets} E_{T}^{jet} > 0.15$, where, for each selected jet with $|\eta|^{jet} < 1.1$, $p_T^{track}$ is computed as the scalar sum of the transverse momenta $p_T^{track}$ of tracks with $p_T^{track} > 0.3$ GeV/c and within a cone of radius $R = 0.4$ around the jet axis.

The dominant QCD-jets background with large $E_T$ originates from the misreconstruction of the jet energies in the calorimeters. In such events, the $E_T$ direction tends to be aligned, in the transverse plane, with one of the leading jets in the event. This background contribution is suppressed by requiring an azimuthal separation $\Delta \phi (E_T - jet) > 0.7$ for each of the selected jets in the event. In the case of the four-jets analysis, the requirement for the least energetic jet is limited to $\Delta \phi (E_T - jet) > 0.3$. Finally, in the two-jets analysis case, the events are rejected if they contain a third jet with $E_T^{jet} > 25$ GeV, $|\eta|^{jet} < 2.0$, and $\Delta \phi (E_T - jet) < 0.2$. The SM background contributions with energetic electrons [25] in the final state from $W$ bosons decaying into muon or tau leptons are vetoed to reject backgrounds with energetic electron candidates in the final state, the lower thresholds on $E_T$, $H_T$, and $E_T^{jet}$ is computed as the scalar sum of the transverse momenta $p_T^{track}$ of tracks with $p_T^{track} > 0.3$ GeV/c and within a cone of radius $R = 0.4$ around the jet axis.

An optimization is carried out to determine, for each final state, the lower thresholds on $E_T$, the $E_T^{jet}$ of the individual jets, and $H_T$, defined as $H_T = \sum_{jets} E_T^{jet}$ [26]. For each mSUGRA sample, the procedure maximizes $S/\sqrt{B}$, where $S$ denotes the number of SUSY events and $B$ is the total SM background. The results from the different mSUGRA samples are then combined to define, for each final state, a single set of lower thresholds that maximizes the search sensitivity in the widest range of squark and gluino masses (see Table I). As an example, for $M_{\tilde{g}} = M_{\tilde{q}}$ and masses between 300 and 400 GeV/$c^2$, values for $S/\sqrt{B}$ in the range between 20 and 6 are obtained, corresponding to SUSY selection efficiencies of 4% to 12%, respectively.

<table>
<thead>
<tr>
<th>Final state</th>
<th>Lower thresholds (GeV) $E_T^1$, $H_T^1$, $E_T^{jet(1)}$, $E_T^{jet(2)}$, $E_T^{jet(3)}$, $E_T^{jet(4)}$</th>
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<tbody>
<tr>
<td>$E_T + \geq 2$ jets</td>
<td>180, 330, 165, 100, 80, 70</td>
</tr>
<tr>
<td>$E_T + \geq 3$ jets</td>
<td>120, 330, 140, 100, 25, 70</td>
</tr>
<tr>
<td>$E_T + \geq 4$ jets</td>
<td>90, 280, 95, 55, 55, 25</td>
</tr>
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A number of control samples in data are considered to test the validity of the SM background predictions for the different processes, as extracted from simulated events [27]. The samples are defined by reversing the logic of the selection criteria described above. Good agreement is observed between the data and the SM predictions in each of the control regions for all the final states considered.

A detailed study of the systematic uncertainties is carried out for each final state [27,28]. A 3% uncertainty on the absolute jet energy scale [24] in the calorimeter introduces an uncertainty in the background prediction that varies between 24% and 34%, and an uncertainty on the mSUGRA signal efficiencies between 15% and 17%. Uncertainties related to the modeling of the initial- and final-state soft gluon radiation in the simulated samples translate into a 3% to 6% uncertainty on the mSUGRA signal efficiency, and uncertainties on the background predictions that vary between 8% and 10%. An additional 10% uncertainty on the diboson and top quark contributions accounts for the uncertainty on the predicted cross sections at NLO. A 2% uncertainty on the measured Drell-Yan cross sections, relevant for $Z/\gamma^* + jets$ and $W + jets$ processes, is also included. The total systematic uncertainty on the SM predictions varies between 31% and 35% as the jet multiplicity increases. Various sources of uncertainty in the mSUGRA cross sections at NLO, as

![FIG. 1](color online). Measured $H_T$ and $E_T$ distributions (black dots) in events with at least two (bottom), three (middle), and four (top) jets in the final state compared to the SM predictions (solid lines) and the SM + mSUGRA predictions (dashed lines). The shaded bands show the total systematic uncertainty on the SM predictions, and the arrows indicate optimized lower thresholds.
determined using PROSPINO, are considered. The uncertainty on the PDFs varies between 10% and 24%, depending on the mSUGRA point and within the mass range considered. Variations of the renormalization and factorization scales by a factor of 2 change the theoretical cross sections by 20% to 25%. For each mSUGRA point considered, observed and expected limits are computed separately for each of the three analyses, and the one with the best expected limit is adopted as the nominal result. Cross sections in the range between 0.1 and 1 pb are excluded by this analysis, depending on the masses considered [27]. The observed numbers of events in data are also translated into 95% C.L. upper limits for squark and gluino masses, for which the uncertainties on the theoretical cross sections are included in the limit calculation, and where the three analyses are combined in a similar way as for the cross section limits. Figure 2 shows the excluded region in the squark-gluino mass plane. For the mSUGRA scenario considered, all squark masses are excluded for $M_{\tilde{g}} < 280 \text{ GeV}/c^2$, while for $M_{\tilde{q}} = M_{\tilde{g}}$ masses up to 392 GeV/$c^2$ are excluded. Finally, for $M_{\tilde{q}} < 400 \text{ GeV}/c^2$ gluinos with $M_{\tilde{q}} < 340 \text{ GeV}/c^2$ are excluded. This analysis extends the previous Run I limits from the Tevatron by 80 to 140 GeV/$c^2$.

In summary, we report results on an inclusive search for squarks and gluinos in $p\bar{p}$ collisions at $\sqrt{s} = 1.96 \text{ TeV}$ in events with large $E_T$ and multiple jets in the final states, based on 2 fb$^{-1}$ of CDF Run II data. The measurements are in good agreement with SM predictions for backgrounds. The results are translated into 95% C.L. upper limits on production cross sections and squark and gluino masses in a given mSUGRA scenario, which significantly extend Run I results.

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Figure 2 (color online). Exclusion plane at 95% C.L. as a function of squark and gluino masses in an mSUGRA scenario with $A_0 = 0$, $\mu < 0$, and $\tan\beta = 5$. The observed (solid line) and expected (dashed line) upper limits are compared to previous results from SPS [30] and LEP [31] experiments at CERN (shaded bands and dotted lines), and from the Run I at the Tevatron [2] (dashed-dotted line). The hatched area indicates the region in the plane with no mSUGRA solution.
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\textsuperscript{a}Deceased
\textsuperscript{b}With visitors from University of Massachusetts Amherst, Amherst, MA 01003, USA.
\textsuperscript{c}With visitors from Universiteit Antwerpen, B-2610 Antwerp, Belgium.
\textsuperscript{d}With visitors from University of Bristol, Bristol BS8 1TL, United Kingdom.
\textsuperscript{e}With visitors from Chinese Academy of Sciences, Beijing 100864, China.
\textsuperscript{f}With visitors from Istituto Nazionale di Fisica Nucleare, Sezione di Cagliari, 09042 Monserrato (Cagliari), Italy.
\textsuperscript{g}With visitors from University of California Irvine, Irvine, CA 92697, USA.
\textsuperscript{h}With visitors from University of California Santa Cruz, Santa Cruz, CA 95064, USA.
\textsuperscript{i}With visitors from Cornell University, Ithaca, NY 14853, USA.
\textsuperscript{j}With visitors from University of Cyprus, Nicosia CY-1678, Cyprus.
\textsuperscript{k}With visitors from University College Dublin, Dublin 4, Ireland.
\textsuperscript{l}With visitors from Royal Society of Edinburgh/Scottish Executive Support Research Fellow.
\textsuperscript{m}With visitors from University of Edinburgh, Edinburgh EH9 3JZ, United Kingdom.
\textsuperscript{n}With visitors from Universidad Iberoamericana, Mexico D.F., Mexico.
\textsuperscript{o}With visitors from Queen Mary, University of London, London, E1 4NS, England.
\textsuperscript{p}With visitors from University of Manchester, Manchester M13 9PL, England.
\textsuperscript{q}With visitors from Nagasaki Institute of Applied Science, Nagasaki, Japan.
\textsuperscript{r}With visitors from University of Notre Dame, Notre Dame, IN 46556, USA.
\textsuperscript{s}With visitors from University de Oviedo, E-33007 Oviedo, Spain.
\textsuperscript{t}With visitors from Simon Fraser University, Vancouver, British Columbia, Canada V6B 5K3.
\textsuperscript{u}With visitors from Texas Tech University, Lubbock, TX 79409, USA.
\textsuperscript{v}With visitors from IFIC(CSIC-Universitat de Valencia), 46071 Valencia, Spain.
\textsuperscript{w}With visitors from University of Virginia, Charlottesville, VA 22904, USA.
\textsuperscript{x}With visitors from On leave from J. Stefan Institute, Ljubljana, Slovenia.
\textsuperscript{y}With visitors from Now at Freiburg University, Germany.

[5] CDF uses a cylindrical coordinate system about the beam axis with polar angle $\theta$ and azimuthal angle $\phi$. We define transverse energy $E_T = E \sin \theta$, transverse momentum $p_T = p \sin \theta$, pseudorapidity $\eta = -\ln(\tan(\theta/2))$, and rapidity $y = \frac{1}{2} \ln(\frac{1 + \eta}{1 - \eta})$. The missing transverse energy $E_F$ is defined as the norm of $-\sum E_T \cdot \hat{n}_i$, where $\hat{n}_i$ is the component in the azimuthal plane of the unit vector pointing from the interaction point to the $i$-th calorimeter tower.
[20] Pole masses are considered. The squark mass is averaged over the first two squark generations.
[25] Charge conjugation is implied throughout the Letter.
[26] The sum runs over the selected jets. In the four-jets case, the first three leading jets are considered.