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We present a measurement of $D^0$-$\bar{D}^0$ mixing and CP violation using the ratio of lifetimes simultaneously extracted from a sample of $D^0$ mesons produced through the flavor-tagged process $D^+ \rightarrow D^0 \pi^+$, where $D^0$ decays to $K^- \pi^+$, $K^- K^+$, or $\pi^- \pi^+$, along with the untagged decays $D^0 \rightarrow K^+ \pi^-$ and $D^0 \rightarrow K^- K^+$. The lifetimes of the CP-even, Cabibbo-suppressed modes $K^- K^+$ and $\pi^- \pi^+$ are compared to that of the CP-mixed mode $K^- \pi^+$ in order to measure $\gamma_{CP}$ and $\Delta Y$. We obtain $\gamma_{CP} = [0.72 \pm 0.18 \text{(stat)} \pm 0.12 \text{(syst)}] \%$ and $\Delta Y = [0.09 \pm 0.26 \text{(stat)} \pm 0.06 \text{(syst)}] \%$, where $\Delta Y$ constrains possible CP violation. The $\gamma_{CP}$ result excludes the null mixing hypothesis at 3.3$\sigma$ significance. This analysis is based on an integrated luminosity of 468 fb$^{-1}$ collected with theBABAR detector at the PEP-II asymmetric-energy $e^+e^-$ collider.

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I. INTRODUCTION

Several measurements [1–6] show evidence for mixing in the $D^0$-$\bar{D}^0$ system consistent with predictions of possible Standard Model (SM) contributions [7–11]. These results also constrain many new physics models [12–16]. An observation of CP violation (CPV) in the $D^0$-$\bar{D}^0$ system at the present experimental sensitivity would provide possible evidence for physics beyond the SM [17–21].
One manifestation of $D^0$-$\bar{D}^0$ mixing is differing $D^0$ decay time distributions for decays to different $CP$ eigenstates Ref. [22]. We present a measurement of charm mixing using the ratio of lifetimes obtained from the decays of neutral $D$ mesons to $CP$-even and $CP$-mixed two-body final states. We also present a search for indirect $CP$ violation arising from a difference in $D^0$ and $\bar{D}^0$ partial decay widths to $CP$-even eigenstates. Recently the LHCb Collaboration has reported evidence for $CPV$ in the difference of the time-integrated $CP$ asymmetries in $D^0 \rightarrow K^- K^+$ and $D^0 \rightarrow \pi^- \pi^+$ decays [23]. This measurement is primarily sensitive to direct $CPV$. As explained in Appendix A, we are not sensitive to effects of direct $CP$ violation at the level of the result reported by LHCb, and we therefore assume no direct $CPV$ in our baseline model.

We measure the effective $D^0$ lifetimes in three different two-body final states: $K^+\pi^\pm$, $K^-K^+$, and $\pi^-\pi^+$. We make no distinction between the Cabibbo-favored $D^0 \rightarrow K^-\pi^+$ and doubly Cabibbo-suppressed $D^0 \rightarrow K^+\pi^-$ modes; in other words, we analyze and describe them together. Given the current experimental evidence indicating a small mixing rate, the lifetime distribution for all two-body final states is exponential to a good approximation. Decays in the $K^+\pi^-$ mode are to a $CP$-mixed final state, and are assumed to be described by the average $D^0$ width $\Gamma$. The singly Cabibbo-suppressed decays $D^0 (\bar{D}^0)$ to the $CP$-even $K^-K^+$ and $\pi^-\pi^+$ final states are described by the partial decay rate $\Gamma^+ (\Gamma^-)$, where $+$ indicates the $CP$ of the final state. We present in Appendix A a discussion of the mixing formalism leading to the expressions that are used to extract the mixing parameter $y_{CP}$ and the $CPV$ parameter $\Delta Y$,

$$y_{CP} = \frac{\Gamma^+ + \Gamma^-}{2\Gamma} - 1,$$

$$\Delta Y = \frac{\Gamma^+ - \Gamma^-}{2\Gamma},$$

from the experimentally measured $CP$-mixed and $CP$-even lifetimes. This definition of $\Delta Y$ is opposite in sign to that in our previous measurement [2] and is now consistent with that used by the Heavy Flavor Averaging Group [24].

Tagged decays refer to $D^0$ mesons coming from $D^{*+} \rightarrow D^0\pi^+$ decays [25], while untagged decays refer to $D^0$ mesons where no $D^{*+}$ parent was found. The charge of the $D^{\pm}$ is used to split the $K^-K^+$ and $\pi^-\pi^+$ samples into those originating from $D^0$ and from $\bar{D}^0$ mesons in order to measure the $CP$-violating parameter $\Delta Y$. The requirement of a $D^{*+}$ parent strongly suppresses backgrounds; hence untagged decays are reconstructed only in $K^+\pi^-$ and $K^-K^+$ because of the relatively poor signal-to-background ratio in the untagged $\pi^-\pi^+$ final state. In summary, we study seven modes: two untagged and five tagged.

In addition to the increased integrated luminosity of the new dataset compared to that used in our earlier results [2,3], this analysis benefits from improved charged-particle track reconstruction, and a more inclusive and optimized event selection. The particle identification selection efficiency was sizably increased both for pions and kaons in the high-momentum-spectrum range by improving the algorithms that combine the information coming from the detector. We implement an improved background model, and we simultaneously fit both the tagged and untagged datasets.

II. EVENT RECONSTRUCTION AND SELECTION

We use 468 fb$^{-1}$ of $e^+e^-$ colliding-beam data recorded at, and slightly below, the $Y(4S)$ resonance ($e^+e^-$ center-of-mass [CM] energy $\sqrt{s} \sim 10.6$ GeV) with the BABAR detector [26] at the SLAC National Accelerator Laboratory PEP-II asymmetric-energy $B$ Factory. To avoid potential bias, we finalize our data selection criteria, as well as the procedures for fitting, extracting statistical limits, and determining systematic uncertainties, prior to examining the results.

We reconstruct charged tracks and vertices with a five-layer, double-sided silicon vertex tracker (SVT) and a 40-layer drift chamber. We select $D^0$ candidates by pairing oppositely charged tracks, requiring each track to satisfy particle identification criteria based on specific ionization energy loss ($dE/dx$) from the SVT and drift chamber, and Cherenkov angle measurements from a ring-imaging Cherenkov detector. We then refit the $D^0$ daughter tracks, requiring them to originate from a common vertex. To reduce contributions from $D^0$'s produced via $B$-meson decay to a negligible level, we require each $D^0$ to have momentum in the CM frame $p_{CM} > 2.5$ GeV/c.

For tagged decays, we reconstruct $D^{*+}$ candidates by combining a $D^0$ candidate with a slow pion track $\pi^+_s$, requiring them to originate from a common vertex constrained to the $e^+e^-$ interaction region. We require the $\pi^+_s$ momentum to be greater than 0.1 GeV/c in the laboratory frame and less than 0.45 GeV/c in the CM frame. We reject a positron that fakes a $\pi^+_s$ candidate by using $dE/dx$ information and veto any $\pi^+_s$ candidate that may have originated from a reconstructed photon conversion or $\pi^0$ Dalitz decay. The distribution of the difference $\Delta m$ between the reconstructed $D^{*+}$ and $D^0$ masses peaks near $\Delta m \sim 0.1455$ GeV/$c^2$. Backgrounds are suppressed by retaining only tagged candidates in the range $0.1447 < \Delta m < 0.1463$ GeV/$c^2$.

To determine the proper time $t$ and its error $\sigma_t$, for each $D^0$ candidate, we perform a combined fit to the $D^0$ production and decay vertices. We constrain the production point to be within the $e^+e^-$ interaction region, which we determine using Bhabha and di-muon events from triggers close in time to any given signal candidate event. We retain only candidates with a $\chi^2$-based probability for the fit
modes) or $D^{*+}$ (for the tagged modes) candidates in an event share one or more tracks, we retain only the candidate with the highest $P(\chi^2)$. The fraction of events with multiple $D^0$ candidates with overlapping daughter tracks is $\ll 1\%$ for all final states.

### III. INVARIANT-MASS FITS

We characterize the $D^0$ invariant-mass ($M$) distribution for each of the seven modes with an extended unbinned maximum likelihood fit to $D^0$ and $\bar{D}^0$ samples. We allow the parameters governing the shapes of the probability density functions (PDFs), as well as the expected signal and background candidate yields, to vary in the fits. For the tagged $CP$-even modes we fit the $D^0$ and $\bar{D}^0$ samples simultaneously, sharing all parameters except for the expected signal and background candidate yields.

We fit the tagged $\pi^-\pi^+$ invariant-mass distribution in the fit range $1.82 < M_{\pi\pi} < 1.93$ GeV/c$^2$ using a sum of two Gaussians with independent means and widths for the signal PDF, along with a first-order Chebychev polynomial for the total background.

The fit model for the tagged $K^-K^+$ invariant-mass distribution is similar to $\pi^-\pi^+$, except that the fit range is $1.82 < M_{KK} < 1.91$ GeV/c$^2$, and the signal PDF is the sum of two independent Gaussians and a modified Gaussian with a power-law tail [28], which aids in better modeling of the lower tail of the distribution.

The signal PDF for the untagged $K^-K^+$ mode and for both tagged and untagged $K^\pm\pi^\mp$ modes is a sum of three independent Gaussians: the background is modeled using a second-order Chebychev polynomial. The mass fit range is $1.82 < M_{KK} < 1.91$ GeV/c$^2$ for the untagged $K^-K^+$ mode, $1.81 < M_{K\pi} < 1.92$ GeV/c$^2$ for the untagged $K^\pm\pi^\mp$ mode, and $1.80 < M_{K\pi} < 1.93$ GeV/c$^2$ for the tagged $K^+\pi^-$ mode. In these modes, we do not distinguish $D^0$ from $\bar{D}^0$ candidates, and therefore determine only the total signal and total background yields, in addition to the signal and background shape parameters.

The reconstructed $D^0$ invariant-mass distributions and the fit results are shown in Fig. 1, together with a plot of the corresponding normalized Poisson pulls [29].

### IV. SIGNAL AND SIDEBAND REGIONS

For the lifetime fit, we determine the regions in two-body invariant mass that maximize signal significance, minimize systematic effects due to backgrounds, and minimize the effect of the correlation between the $D^0$ invariant mass and proper time. We refer to these regions as the lifetime-fit mass regions. Based on these studies, the optimal lifetime-fit mass region is $34$ MeV/c$^2$ wide for all tagged modes and untagged $K^\pm\pi^\mp$ events, $1.847 < M < 1.881$ GeV/c$^2$. Because of the smaller signal-to-background ratio for the untagged $K^-K^+$ events, the lifetime-fit mass region for this mode is only $24$ MeV/c$^2$ in width.
1.852 < M < 1.876 GeV/c². For the tagged modes, a mass difference sideband 0.151 < Δm < 0.159 GeV/c² is used, along with a low- (high-) invariant-mass sideband, 1.819 (1.890) < M < 1.839 (1.910) GeV/c². The low- (high-) mass sideband used for the untagged modes, 1.810 (1.899) < M < 1.830 (1.919) GeV/c², is displaced from the tagged sideband in order to reduce the signal component there. The contribution of the signal events in the sideband regions is in general very small compared to the background; however, it has been considered when extracting the combinatorial-background PDF. The signal purities in the lifetime-fit mass regions range from ~75% for the untagged K⁻K⁺ sample to ~99.8% for the tagged K⁺π⁻ events.

We classify D⁰ candidate decays in the lifetime-fit mass region as follows: D⁰ signal decays; misreconstructed-charm decays, i.e., those in which the candidate-D⁰ daughter tracks are decay products of a nonsignal weak charm decay; and random combinatorial background. Table I gives the composition of the misreconstructed-charm backgrounds expected from simulated events [30] in each final state.

V. LIFETIME FIT

The lifetimes are determined from an extended unbinned maximum likelihood fit to t and σₚ for candidates in the lifetime-fit mass region. All modes are fit simultaneously using shared signal-resolution-function parameters. The signal, misreconstructed-charm, and combinatorial components are described by their own set of PDFs, which in the tagged modes can also depend on the charm flavor.

The lifetime PDF for the signal is an exponential function convolved with a resolution function, which is the sum of three Gaussian functions whose widths are proportional to σₚ. The explicit form of the signal lifetime PDF is

\[ R_{f,L}^P(t, \sigma_p) = f_{i1} D(t, \sigma_p; S'_p F_{p}s_1, t_0, \tau_L) + (1 - f_{i1}) [ f_{i2} D(t, \sigma_p; S'_p F_{p}s_2, t_0, \tau_L) ] + (1 - f_{i2}) D(t, \sigma_p; S'_p F_{p}s_3, t_0, \tau_L) \]

where \( f_{ij} \) (with \( i = 1, 2 \)) parametrizes the contribution of each individual Gaussian, \( \sigma_i \) (with \( i = 1, 2, 3 \)) is a scaling factor associated with each Gaussian, and \( t_0 \) is an offset of the mean of the resolution function. The function \( D(t, \sigma_p; s, t_0, \tau) \) is given by

\[ D(t, \sigma_p; s, t_0, \tau) = C_{\sigma_p} \int \exp(-t/t_{\text{true}}/\tau) \times \exp\left(-\frac{(t-t_{\text{true}}+t_0)^2}{2(s^2+\sigma_p^2)}\right) dt_{\text{true}}, \]

where the normalization coefficient \( C_{\sigma_p} \) is chosen such that

\[ \int D(t, \sigma_p; s, t_0, \tau) dt = 1 \quad \text{for each } \sigma_p. \]
The $f_{D^0}$ are varied as part of the systematic error estimate for $\gamma_{CP}$ and $\Delta Y$. All five tagged and two untagged signal-lifetime PDFs are explicitly given in Appendix B.

The $\sigma_t$ PDF for signal candidates is obtained directly from data by subtracting the sum of the background $\sigma_t$ distributions from that of all candidates in the lifetime-fit mass region. These one-dimensional $\sigma_t$ distributions are used to model the $H_{t\sigma}^{D^0} (\sigma_t)$ PDF discussed previously.

We determine the $t$ versus $\sigma_t$ misreconstructed-charm signal-like PDF shape parameters and yields by fitting simulated events in the lifetime-fit mass region and then fix these parameters in the lifetime fit to data. We vary the lifetimes and yields as part of the study of systematic effects.

The largest background in the lifetime-fit mass region is due to random combinations of tracks. The PDF describing the two-dimensional combinatorial background in $t$ and $\sigma_t$ in the lifetime-fit mass region is characterized as a weighted average of the two-dimensional PDFs extracted from the mass sideband regions. The weights for the low and high sidebands are obtained from simulated events. The $(t, \sigma_t)$ combinatorial PDF in each sideband and for each mode, except for the untagged $K^- K^+$ mode, is extracted as a two-dimensional histogram from the sideband samples. From these histograms we subtract the contribution of signal and misreconstructed-charm backgrounds, each of which is estimated from simulated events, to obtain the final combinatorial PDF in each sideband. For the untagged $K^- K^+$ mode, a similar procedure is used but, instead of histograms, analytic signal-like PDFs are used. For the background PDFs the offsets and the lifetimes are allowed to be different for each Gaussian. The signal and misreconstructed-charm PDF parameters are extracted by fitting simulated events, and then fixed, along with the expected candidate yields, in the fit that extracts the combinatorial PDFs in each sideband.

For the untagged $K^- K^+$ mode both the expected signal and combinatorial yields are free parameters in the lifetime fit. The expected combinatorial background yields in the other modes are determined by integrating the total background PDF extracted from the mass fit in the lifetime-fit mass region, and then subtracting the expected misreconstructed-charm background yields, which are determined from samples of simulated events. A small bias on these fit yields is observed in fits to simulated events. To correct for this, we scale the data yields based on the simulated-event fits and vary the mode-dependent scale factors as a systematic uncertainty. Table II gives

<table>
<thead>
<tr>
<th>Tagged Modes</th>
<th>Untagged Modes</th>
</tr>
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<tbody>
<tr>
<td>$\pi^+ \pi^-$</td>
<td>$K^- K^+$</td>
</tr>
<tr>
<td>$K^- \pi^+$</td>
<td>$K^- \pi^-$</td>
</tr>
</tbody>
</table>

| Signal | 65 430 | 136 870 | 1 487 000 | 496 200 | 5 825 300 |
|        | $\pm 260$ | $\pm 370$ | $\pm 1200$ | $\pm 1200$ | $\pm 2600$ |
| Comb. Bkgd. | 3760 | 653 | 2849 | 165 000 | 1 044 552 | $\pm 1000$ |
| Charm Bkgd. | 97 | 309 | 642 | 5477 | 4645 |

FIG. 2 (color online). Proper-time $t$ distribution for each decay mode with the fit results overlaid. The combinatorial distribution (indicated as “Comb.” in light gray) is stacked on top of the misreconstructed-charm distribution (indicated as “Charm” in dark gray). The normalized Poisson pulls for each fit are shown under each plot: “unt” refers to the untagged datasets. The bottom right plot shows the individual lifetimes (with statistical uncertainties only): the gray band indicates the PDG $D^0$ lifetime $\pm 1 \sigma_{[27]}$. 

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the event-class yields plus uncertainties obtained from the lifetime fit and indicates the yields that are fixed.

The simultaneous fit to all events in the lifetime-fit mass region has 20 floating parameters: the seven signal yields and three signal lifetimes; the yield of untagged $K^-K^+$ combinatorial candidates; the offset $t_0$; the parameters $f_{SD}$ and $f_{12}$ characterizing the weight of each Gaussian in the signal resolution mode; and the proper-time error scaling parameters $s_1, s_2, s_3, S_{KK}, S_{\pi\pi},$ and $S'_{bg}$. After extracting the three signal lifetimes, using their reciprocals in the computation of $\gamma_{CP}$ and $\Delta Y$ as defined in Eqs. (1) and (2), respectively, we find

$$\gamma_{CP} = [0.72 \pm 0.18\text{(stat)}]\%,$$
$$\Delta Y = [0.09 \pm 0.26\text{(stat)}]\%.$$

The statistical errors are computed using the covariance matrix returned by the fit. The lifetime-fit mass region proper-time distributions and projections of the lifetime fit for the seven different decay modes are shown in Fig. 2.

VI. CROSS CHECKS AND SYSTEMATICS

We have performed numerous cross checks to search for potential problems, in addition to quantitative studies that yield the systematic uncertainties given in Table III, discussed below. Initially, we tested the fit model by generating large ensembles of data sets randomly drawn from the underlying total PDF, and observed no biases in the $\gamma_{CP}$ and $\Delta Y$ results obtained. In addition, we have fit an ensemble of four simulated data sets, each equivalent in luminosity to the data, and found no evidence of bias in $\gamma_{CP}$ or $\Delta Y$.

In fitting the data, we find that the tagged and untagged extracted lifetimes for $K^-K^+$, and separately for $K^{\pm}\pi^{\mp}$, are compatible within the statistical uncertainties. We performed a simultaneous fit to the tagged channels, and a separate simultaneous fit to the untagged channels, and find the lifetimes to be compatible within the statistical uncertainties. We repeated the fit allowing the $K^-K^+$ and $\pi^-\pi^+$ final states to have separate $\tau^+$ and $\bar{\tau}^+$ lifetimes, and observed no statistically significant difference between the $K^-K^+$ and $\pi^-\pi^+$ results. We estimated the effects of the SVT misalignment to be negligible.

We varied the lifetime-fit mass region width by $\pm 4$ and $\pm 2$ MeV/$c^2$. We adopt as the systematic uncertainty half the rms of the differences $|\Delta[\gamma_{CP}]|$ and $|\Delta[\Delta Y]|$ from the nominal-fit central values. We also shifted the position of each mass region by centering each of them at the most probable value for the signal PDF obtained in the invariant-mass fits. These systematic uncertainties are given in the first two lines of Table III.

For the untagged $K^-K^+$ mode, the combinatorial yield is a parameter determined in the lifetime fit. However, it is also needed to determine the signal $\sigma$, PDF. We first use the total background yield determined from the mass fit to extract a signal $\sigma$, PDF, which is employed in an initial simultaneous lifetime fit. The combinatorial yield from this fit is used to construct an improved $\sigma$, signal PDF and a second fit is performed (the nominal fit). We estimate the systematic error on $\gamma_{CP}$ and $\Delta Y$ associated with the determination of the signal $\sigma$, PDF for the untagged $K^-K^+$ mode to be the difference in the values obtained from an additional iteration of the fit and the nominal fit.

We vary the nominal mistag rate of 0.2% by $\pm 0.04\%$, a 20% relative variation, and find no significant change in the nominal fit values. Instead of assuming equal fractions of $D^0$ and $\bar{D}^0$ in the untagged $K^-K^+$ mode, we adopt the latest CDF result for direct CPV [32], and find negligible change in $\gamma_{CP}$ and $\Delta Y$.

We rely on simulated events to determine both the PDF shapes and yields for the misreconstructed-charm backgrounds. To account for the model dependence, we vary the effective lifetime of these events by $\pm 5\%$, except for the tagged $\pi^-\pi^+$ mode where the variation is $\pm 15\%$ due to the small number of simulated events that pass the selection criteria for this mode. We also vary the expected misreconstructed-charm yields by $\pm 10\%$ in the tagged channels, and $\pm 5\%$ in the untagged channels. Each variation is simultaneously applied to all modes. These are $\geq 2\sigma$ variations relative to the statistical uncertainties of the simulated data sets.

We vary the yields, weighting parameters, and fitting strategy used to obtain the two-dimensional lifetime PDF for combinatorial-background events in the lifetime-fit mass region from the mass sidebands. The yields for the tagged combinatorial-background events are varied by $\sim 5\%$ in the $\pi^-\pi^+$ mode, 15% in tagged $K^-K^+$, and 20% in $K^{\pm}\pi^{\mp}$. The untagged $K^{\pm}\pi^{\mp}$ combinatorial-background yield is varied using the value extracted from an alternative lifetime-fit model in which the yield is allowed to vary.
We exclude the null mixing hypothesis at \(3\) parameter consistent with the world average value of the mixing significance, and find no evidence for \(y_{CP}\) used in both cases. The new weighted average of all previous measurements \([24]\), when the previous \(\text{BABAR}\) results are excluded. We obtain

\[
y_{CP} = [0.72 \pm 0.18(\text{stat}) \pm 0.12(\text{syst})]\%.
\]

\[
\Delta Y = [0.09 \pm 0.26(\text{stat}) \pm 0.06(\text{syst})]\%.
\]

Finally, we vary the \(\sigma_i\) criteria by \(\pm 0.1\) ps from the nominal \(\sigma_i < 0.5\) ps, and take as the systematic uncertainty the rms of the deviations from the nominal-fit central value divided by \(\sqrt{2}\). We also consider two variations in how multiple candidates sharing one or more daughter tracks are treated. In the first variation, we retain all multiple candidates if each candidate passes all the other selection criteria. In the second variation, we reject all multiple candidates sharing one or more daughter tracks. We fit these data sets using the nominal-fit model, and assign the largest observed deviation from the nominal \(y_{CP}\) and \(\Delta Y\) central values as the systematic uncertainty in Table III. The total \(y_{CP}\) and \(\Delta Y\) systematic uncertainties are calculated by summing the contributions from all sources in quadrature, and are reported in the last row of Table III.

VII. CONCLUSIONS

In summary, we measured \(y_{CP}\) and \(\Delta Y\) to a precision significantly better than our previous measurements \([2,3]\). Both results are more precise than, and consistent with, the weighted average of all previous measurements \([24]\), when the previous \(\text{BABAR}\) results are excluded. We obtain

\[
y_{CP} = [0.72 \pm 0.18(\text{stat}) \pm 0.12(\text{syst})]\%.
\]

\[
\Delta Y = [0.09 \pm 0.26(\text{stat}) \pm 0.06(\text{syst})]\%.
\]

We exclude the null mixing hypothesis at \(3.3\sigma\) significance, and find no evidence for \(CPV\). Our results are consistent with the world average value of the mixing parameter \(y\) obtained from \(D^0 \rightarrow K^0_s h^- h^+\) (where \(h = K, \pi\)) \([24]\), as expected in the absence of \(CPV\). The \(y_{CP}\) measurement is the most precise single measurement to date, with significant improvements on the statistical and systematic error with respect to the previous most precise measurement \([3]\) \(y_{CP} = (1.16 \pm 0.22 \pm 0.18)\%\).

The value of \(\Delta Y\) obtained here is consistent with our previously published result \([2]\) when the same definition is used in both cases. The new \(y_{CP}\) value is consistent with our previous result \([3]\) with a probability of \(\approx 2\%\), assuming that the systematics for both the old and new measurements are fully correlated, and taking into account the fact that \(~40\%) of the events in the current sample are also present in the samples used in the previous measurements \([2,3]\). The results here supersede the previous \(\text{BABAR}\) results for these modes \([2,3]\).

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APPENDIX A: MIXING FORMALISM

AND CONSIDERATIONS ON THE ROLE OF DIRECT CP VIOLATION

In the following we briefly review the mixing formalism \([27]\) considering the possible effects of direct \(CPV\) on the \(y_{CP}\) and \(\Delta Y\) observables.

The time evolution of the flavor eigenstates \(D^0\) and \(\bar{D}^0\) is governed by the Schrödinger equation:

\[
i \frac{\mathrm{d}}{\mathrm{d}t} \begin{pmatrix} D^0(t) \\ \bar{D}^0(t) \end{pmatrix} = \begin{pmatrix} M & -i/2 \Gamma \\ -i/2 \Gamma & M \end{pmatrix} \begin{pmatrix} D^0(t) \\ \bar{D}^0(t) \end{pmatrix} \tag{A1} \]

The mass eigenstates \(D_1\) and \(D_2\) are obtained from the diagonalization of the effective Hamiltonian \(\mathcal{H}_{\text{eff}} = M - i \Gamma/2\). Under the hypothesis of \(CPT\) conservation the two mass eigenstates can be written in terms of the flavor eigenstates as

\[
|D_1\rangle = p|D^0\rangle + q|\bar{D}^0\rangle, \quad |D_2\rangle = p|D^0\rangle - q|\bar{D}^0\rangle, \tag{A2}\]

where
\[ \left( \frac{q^2}{p} \right) = \frac{M_{12}^* - i \Gamma_{12}^*}{M_{12} - i \Gamma_{12}} \quad \text{and} \quad |p|^2 + |q|^2 = 1. \]  
(A3)

We choose the positive root for \( q/p \): choosing the negative one just means exchanging \( D_1 \) with \( D_2 \). If \( CP[D^0] = +|D^0| \), in the case of no \( CPV \), \( D_1 \) is the \( CP \)-even state and \( D_2 \) the \( CP \)-odd state.

It is traditional to quantify the size of \( D^0 \)-\( D^0 \) mixing in terms of the parameters \( x = \Delta m / \Gamma \) and \( y = \Delta \Gamma / 2 \Gamma \), where \( \Delta m = m_1 - m_2 \) (\( \Delta \Gamma = \Gamma_1 - \Gamma_2 \)) is the difference in mass (width) of the states defined in Eq. (A2) and \( \Gamma = (\Gamma_1 + \Gamma_2) / 2 \) is the average width. If either \( x \) or \( y \) is nonzero, mixing will occur. While most SM expectations for the size of both are \( \approx 10^{-3} \) [10,33], values as high as \( 10^{-2} \) or even higher are predicted by certain models [13,15].

\( CP \) violation can manifest in \( D^0 \) decays in three ways:

(i) in decay, when \( |A_f/\tilde{A}_f| \neq 1 \),
(ii) in mixing, when \( r_m = |q/p| \neq 1 \), and
(iii) in the interference between decays with and without mixing, when the weak phase \( \phi_f \) of \( \lambda_f = \frac{q}{p} \tilde{A}_f \) is different from zero, where \( A_f (\tilde{A}_f) \) is the amplitude for \( D^0 (\tilde{D}^0) \) decaying into a final state \( f \), \( A_f = \langle f | \mathcal{H}_D | D^0 \rangle (\tilde{A}_f = \langle f | \mathcal{H}_D | \tilde{D}^0 \rangle) \).

The presence of mixing alters the exponential distribution for the \( D^0 \) decay into a final state \( f \). In particular we have

\[ \Gamma(D^0(t) \to f) = \frac{1}{2} |A_f|^2 e^{-\Gamma_f t} [(1 + |A_f|^2) \cosh \Gamma_f t + (1 - |A_f|^2) \cos \Gamma_f t - 2 \sinh \Gamma_f t \sin \Gamma_f t], \]  
(A4)

\[ \Gamma(\tilde{D}^0(t) \to f) = \frac{1}{2} |\tilde{A}_f|^2 e^{-\tilde{\Gamma}_f t} [(1 + |\tilde{A}_f|^2) \cosh \tilde{\Gamma}_f t + (1 - |\tilde{A}_f|^2) \cos \tilde{\Gamma}_f t - 2 \sinh \Gamma_f t \sin \Gamma_f t], \]  
(A5)

In this analysis we are interested in \( CP \)-even final states \( f = h^+ h^-, h = K, \pi \). If we neglect second-order terms in \( x \Gamma_f t \) and \( y \Gamma_f t \), the decay time distributions can be treated as exponentials with effective widths [34]:

\[ \Gamma(D^0(t) \to f) \approx e^{-\Gamma^+_h t} \quad \text{with} \]
\[ \Gamma^+_h = \Gamma[1 + \alpha(y) \beta(y_h) - \alpha \beta(y_h)], \]  
(A6)

\[ \Gamma(\tilde{D}^0(t) \to f) \approx e^{-\tilde{\Gamma}^+_h t} \quad \text{with} \]
\[ \tilde{\Gamma}^+_h = \Gamma[1 + \alpha(y) \beta(y_{\tilde{h}}) - \alpha \beta(y_{\tilde{h}})], \]  
(A7)

To better understand the effects of \( CP \) violation we introduce two more parameters, one describing \( CPV \) in decay \( (A_f^D) \) and one in mixing \( (A_M) \):

\[ A_f^D = |A_f/\tilde{A}_f|^2 - |\tilde{A}_f/A_f|^2, \]  
(A8)

\[ A_M = \frac{r_m^2 - r_m^{-2}}{r_m^2 + r_m^{-2}}. \]  
(A9)

Since \( f = h^+ h^- \), then \( f = \tilde{f} \). Noting that there is no strong phase in \( \lambda_f \) since the final state is its own \( CP \) conjugate, we can express \( \lambda_{hh} \) in terms of \( A_{hh}^D \), \( A_M \), and the \( CP \)-violating phase \( \phi_{hh} \):

\[ \lambda_{hh} = \left[ \frac{1 - A_{hh}^D}{1 + A_{hh}^D} \right]^{1/4} e^{i \phi_{hh}}. \]  
(A10)

Expanding Eqs. (A6) and (A7), and retaining only terms up to first order in \( A_{hh}^D \) and \( A_M \), we obtain

\[ \Gamma_{hh}^+ = \Gamma \left[ 1 + (y \cos \phi_{hh} - x \sin \phi_{hh}) + \frac{1}{2} (A_M - A_{hh}^D)(y \cos \phi_{hh} - x \sin \phi_{hh}) - \frac{1}{4} A_M A_{hh}^D (y \cos \phi_{hh} - x \sin \phi_{hh}) \right], \]  
(A11)

\[ \Gamma_{hh}^\tilde{+} = \Gamma \left[ 1 + (y \cos \phi_{hh} + x \sin \phi_{hh}) - \frac{1}{2} (A_M - A_{hh}^D)(y \cos \phi_{hh} + x \sin \phi_{hh}) - \frac{1}{4} A_M A_{hh}^D (y \cos \phi_{hh} + x \sin \phi_{hh}) \right]. \]  
(A12)

Combining the widths defined above we obtain the two observables \( y_{CP}^{hh} \) and \( \Delta Y^{hh} \) which, in general, depend on the final state because of the \( CPV \) parameters \( A_{hh}^D \) and \( \phi_{hh} \):

\[ y_{CP}^{hh} = \frac{\Gamma_{hh}^+ + \Gamma_{hh}^\tilde{+}}{2 \Gamma} - 1, \]  
(A13)

\[ \Delta Y^{hh} = \frac{\Gamma_{hh}^+ - \Gamma_{hh}^\tilde{+}}{2 \Gamma}. \]  
(A14)

Other experiments characterize the \( CP \)-violating observable as \( A_{hh}^D \),

\[ A_{hh}^D = \frac{\Gamma_{hh}^+ - \Gamma_{hh}^\tilde{+}}{\Gamma_{hh}^+ + \Gamma_{hh}^\tilde{+}}. \]  
(A15)

The relationship between \( A_{hh}^D \), \( \Delta Y^{hh} \), and \( y_{CP}^{hh} \) is

\[ \Delta Y^{hh} = (1 + y_{CP}^{hh}) A_{hh}^D. \]  
(A16)

These quantities are directly related to the fundamental parameters that govern mixing and \( CPV \) in the charm sector.
From the experimental point of view, we measure three lifetimes instead of the partial widths: 

(i) $\tau^+$ for the $D^0 \to K^- K^+$, $\pi^- \pi^+$ decays, 
(ii) $\bar{\tau}^+$ for the $\bar{D}^0 \to K^- K^+$, $\pi^- \pi^+$ decays, 
(iii) $\tau_{K\pi}$ for the $D^0$ (and $\bar{D}^0$) $\to K^\pm \pi^\mp$ decays (the Cabibbo-favored $K^- \pi^+$ and the doubly Cabibbo-suppressed $K^+ \pi^-$ decays are collected in the same sample), 

and use their inverse to compute $\gamma_{CP}$ and $\Delta Y$.

The measured observables constrain the parameters that govern mixing and indirect CPV in the charm sector.

APPENDIX B: SIGNAL-LIFETIME PDFS

The explicit form of the signal-lifetime PDFs based on the prototype PDFs presented in the main text are given below:

$$P_{I\pi\pi}^d(t, \sigma_i) = (1 - f_{tag}) R_{I\pi\pi}^{tag}(t, \sigma_i, S_{\pi\pi} S_{\pi\pi}^s, t_0, \tau^+)$$

$$+ f_{tag} R_{I\pi\pi}^{tag}(t, \sigma_i, S_{\pi\pi} S_{\pi\pi}^s, t_0, \bar{\tau}^+),$$

$$P_{K\pi}^d(t, \sigma_i) = (1 - f_{tag}) R_{K\pi}^{tag}(t, \sigma_i, S_{KK} S_{\pi\pi}^s, t_0, \tau^+)$$

$$+ f_{tag} R_{K\pi}^{tag}(t, \sigma_i, S_{KK} S_{\pi\pi}^s, t_0, \bar{\tau}^+),$$

$$P_{K\pi}^{d-}(t, \sigma_i) = \frac{1}{R_{K\pi}^{tag}(t, \sigma_i, S_{KK} S_{\pi\pi}^s, t_0, \tau^+)}$$

$$+ \frac{1}{R_{K\pi}^{tag}(t, \sigma_i, S_{KK} S_{\pi\pi}^s, t_0, \bar{\tau}^+)}$$

$$P_{KK}^{d-}(t, \sigma_i) = \frac{1}{R_{KK}^{tag}(t, \sigma_i, S_{KK} S_{KK}^s, t_0, \tau_{KK})}$$

$$+ \frac{1}{R_{KK}^{tag}(t, \sigma_i, S_{KK} S_{KK}^s, t_0, \bar{\tau}_{KK})}$$

$$P_{KK}^{d-}(t, \sigma_i) = \frac{1}{R_{KK}^{tag}(t, \sigma_i, S_{KK} S_{KK}^s, t_0, \tau_{KK})}$$

$$+ \frac{1}{R_{KK}^{tag}(t, \sigma_i, S_{KK} S_{KK}^s, t_0, \bar{\tau}_{KK})}$$

$$P_{KK}^{d-}(t, \sigma_i) = \frac{1}{R_{KK}^{tag}(t, \sigma_i, S_{KK} S_{KK}^s, t_0, \tau_{KK})}$$

$$+ \frac{1}{R_{KK}^{tag}(t, \sigma_i, S_{KK} S_{KK}^s, t_0, \bar{\tau}_{KK})}$$

where $f_{tag} = 0.2\%$, $f_{D^0} = 0.5$, and $S_{K\pi} = S_{\pi\pi}^s = 1$ are fixed in the nominal fit.

[25] Charge conjugation is implied throughout.
[30] The detector simulation is based on the GEANT 4 [31] toolkit. The simulated events are reconstructed using the same procedure as for real data.