Search for Dark Matter Candidates and Large Extra Dimensions in Events with a Photon and Missing Transverse Momentum in pp Collision Data at $s=7$TeV with the ATLAS Detector.

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Search for Dark Matter Candidates and Large Extra Dimensions in Events with a Photon and Missing Transverse Momentum in \( pp \) Collision Data at \( \sqrt{s} = 7 \) TeV with the ATLAS Detector

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Results of a search for new phenomena in events with an energetic photon and large missing transverse momentum in proton-proton collisions at \( \sqrt{s} = 7 \) TeV are reported. Data collected by the ATLAS experiment at the LHC corresponding to an integrated luminosity of 4.6 fb\(^{-1}\) are used. Good agreement is observed between the data and the standard model predictions. The results are translated into exclusion limits on models with large extra spatial dimensions and on pair production of weakly interacting dark matter candidates.

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Events with an energetic photon and large missing momentum in the final state constitute a clean and distinctive signature in searches for new physics at colliders. In particular, monophoton, and monojet final states have been studied [1–8] in the context of searches for supersymmetry and large extra spatial dimensions (LED), aiming to provide a solution to the mass hierarchy problem, and the search for weakly interacting massive particles (WIMPs) as candidates for dark matter (DM).

The Arkani-Hamed, Dimopoulos, and Dvali (ADD) model for LED [9] explains the large difference between the electroweak unification scale \( O(10^2) \) GeV and the Planck scale \( M_{Pl} \sim O(10^{19}) \) GeV by postulating the presence of \( n \) extra spatial dimensions of size \( R \), and defining a fundamental Planck scale in \( 4 + n \) dimensions, \( M_D \), given by \( M_D^2 \sim M_{Pl}^2 / R^n \). The extra spatial dimensions are compactified, resulting in a Kaluza-Klein tower of massive graviton modes. At hadron colliders, these graviton modes may escape detection and can be produced in association with an energetic photon or a jet, leading to a monophoton or monojet signature.

The presence of a nonbaryonic DM component in the Universe is inferred from the observation of its gravitational interactions [10], although its nature is otherwise unknown. A WIMP \( \chi \) with mass \( m_\chi \) in the range between 1 GeV and a few TeV is a plausible candidate for DM. It could be detected via its scattering with heavy nuclei [11], the detection of cosmic rays (energetic photons, electrons, positrons, protons, antiprotons, or neutrinos) from \( \chi\bar{\chi} \) annihilation in astrophysical sources [10], or via \( \chi\bar{\chi} \) pair production at colliders where the WIMPs do not interact with the detector and the event is identified by the presence of an energetic photon or jet from initial-state radiation. The interaction of WIMPs with standard model (SM) particles is assumed to be driven by a mediator with mass at the TeV scale and described using a nonrenormalizable effective theory [12] with several operators. The vertex coupling is suppressed by an effective cutoff mass scale \( M_* \sim M/\sqrt{8182} \), where \( M \) denotes the mass of the mediator and \( g_1 \) and \( g_2 \) are the couplings of the mediator to the WIMP and SM particles.

This Letter reports results of the search for new phenomena in the monophoton final state, based on \( \sqrt{s} = 7 \) TeV proton-proton collision data corresponding to an integrated luminosity of 4.6 fb\(^{-1}\) collected with the ATLAS detector at the LHC during 2011. The ATLAS detector is described in detail elsewhere [13]. The data are collected using a three-level trigger system that selects events with missing transverse momentum greater than 70 GeV. In the analysis, events are required to have a reconstructed primary vertex and an isolated (\( p_T^\text{miss} > 150 \) GeV), where \( p_T^\text{miss} \) is computed as the magnitude of the vector sum of the transverse momentum of all noise-suppressed calorimeter topological clusters with \( |\eta| < 4.9 \) [14,15]. A photon is also required with transverse momentum \( p_T > 150 \) GeV and \( |\eta| < 2.37, \) excluding the calorimeter barrel or endcap transition regions \( 1.37 < |\eta| < 1.52 \) [13]. With these criteria, the trigger selection is more than 98% efficient, as determined using events selected with a muon trigger. The cluster energies are corrected for the different response of the calorimeters to hadronic jets, \( \tau \) leptons, electrons or photons, as well as dead material and out-of-cluster energy losses. The photon candidate must pass tight identification criteria [16] and is required to be isolated: the energy not associated with the photon cluster in a cone of radius \( \Delta R = \sqrt{(\Delta \eta)^2 + (\Delta \phi)^2} = 0.4 \) around the candidate is required to be less than 5 GeV. Jets are defined using the anti-\( k_t \) jet algorithm [17] with the distance parameter set to \( R = 0.4 \). The measured jet \( p_T \) is corrected for detector

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events and for contributions from multiple proton-proton interactions per beam bunch crossing (pileup) [18].

Events with more than one jet with $p_T > 30$ GeV and $|\eta| < 4.5$ are rejected. Events with one jet are retained to increase the signal acceptance and reduce systematic uncertainties related to the modeling of initial-state radiation. The reconstructed photon, $E_T^{\text{miss}}$ vector, and jets (if found) are required to be well separated in the transverse plane with $\Delta \phi (\gamma, E_T^{\text{miss}}) > 0.4$, $\Delta R (\gamma, \text{jet}) > 0.4$, and $\Delta \phi (\text{jet}, E_T^{\text{miss}}) > 0.4$. Additional quality criteria [19] are applied to ensure that jets and photons are not produced by noisy calorimeter cells, and to avoid problematic detector regions. Events with unidentified electrons or muons are vetoed to reject mainly $W/Z + \text{jets}$ and $W/Z + \gamma$ background processes with charged leptons in the final state. Electron (muon) candidates are required to have $p_T > 20$ GeV and $|\eta| < 2.47$ ($p_T > 10$ GeV and $|\eta| < 2.4$), and to pass the medium (combined) criteria [20]. The final data sample contains 116 events, where 88 and 28 events have zero and one jet, respectively.

The SM background to the monophoton signal is dominated by the irreducible $Z(\rightarrow \nu \bar{\nu}) + \gamma$ process, and receives contributions from $W/Z + \gamma$ events with unidentified electrons, muons or hadronic $\tau$ decays, and $W/Z + \text{jets}$ events with an electron or jet misreconstructed as a photon. In addition, the monophoton sample receives small events with identified electrons or jet misreconstructed as a photon candidates and/or photon candidates passing a photon veto in the nominal event selection criteria discussed above. According to the simulation, the sample contains a 71% (19%) contribution from $W + \gamma (Z + \gamma)$ processes. This control sample is used to normalize separately the $W + \gamma$ and $Z + \gamma$ MC predictions determined in a data control sample, resulting in a significant reduction of the background uncertainties. A $\gamma + \mu + E_T^{\text{miss}}$ control sample with an identified muon is defined by inverting the muon veto in the nominal event selection criteria discussed above. According to the simulation, the sample contains a 71% (19%) contribution from $W + \gamma (Z + \gamma)$ processes. This control sample is used to normalize separately the $W + \gamma$ and $Z + \gamma$ MC predictions determined in a data control sample, resulting in a significant reduction of the background uncertainties. A $\gamma + \mu + E_T^{\text{miss}}$ control sample with an identified muon is defined by inverting the muon veto in the nominal event selection criteria discussed above. According to the simulation, the sample contains a 71% (19%) contribution from $W + \gamma (Z + \gamma)$ processes. This control sample is used to normalize separately the $W + \gamma$ and $Z + \gamma$ MC predictions determined in a data control sample.

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identifications requirements are used to determine the rate of jets identified as photons in the signal region, after the contribution from $W/Z + \gamma$ processes has been subtracted. This gives an estimate of $4.3 \pm 1.9 \ W/Z + \text{jet}$ background events.

The $\gamma + \text{jet}$ and multijet background contributions to the signature of a photon and large $E_T^{\text{miss}}$ originate from the misreconstruction of the energy of a jet in the calorimeter. The direction of the $E_T^{\text{miss}}$ vector therefore tends to be aligned with the jet. These background contributions are determined from data using a control sample with the nominal selection criteria and at least one jet with $p_T > 30 \ \text{GeV}$ and $\Delta \phi(\text{jet}, E_T^{\text{miss}}) < 0.4$. After the subtraction of electroweak boson and top-quark production processes, a linear extrapolation of the measured $p_T$ spectrum to $p_T < 30 \ \text{GeV}$ leads to an estimate of $1.0 \pm 0.5$ background events in the signal region, where the uncertainty is due to the ambiguity in the functional form used in the extrapolation. Background contributions from top-quark, $\gamma \gamma$, and diboson production processes, determined using MC samples, are small. Finally, noncollision backgrounds are negligible.

A detailed study of systematic uncertainties on the background predictions has been performed. An uncertainty of $0.3\%$ to $1.5\%$ on the absolute photon energy scale [16], depending on the photon $p_T$ and $\eta$, translates into a $0.9\%$ uncertainty on the total background prediction. Uncertainties on the simulated photon energy resolution, photon isolation, and photon identification efficiency introduce a combined $1.1\%$ uncertainty on the background yield. Uncertainties on the simulated lepton identification efficiencies introduce a $0.3\%$ uncertainty on the background predictions. The uncertainty on the absolute jet energy scale [18] and jet energy resolution introduce $0.9\%$ and $1.2\%$ uncertainties on the background estimation, respectively. A $10\%$ uncertainty on the absolute energy scale for low $p_T$ jets and unclustered energy in the calorimeter, and a $6.6\%$ uncertainty on the subtraction of pileup contributions, are taken into account. They affect the $E_T^{\text{miss}}$ determination and translate into $0.8\%$ and $0.3\%$ uncertainties on the background yield, respectively. The dependence of the predicted $W/Z + \gamma$ backgrounds on the parton shower and hadronization model used in the MC simulations is studied by comparing the predictions from SHERPA and ALPGEN. This results in a conservative $6.9\%$ uncertainty on the total background yield. Uncertainties due to the choice of PDFs and the variation of the renormalization and factorization scales in the $W/Z + \gamma$ MC samples introduce an additional $1.0\%$ uncertainty on the total background yields. Other sources of systematic uncertainty related to the trigger selection, the lepton $p_T$ scale and resolution, the pileup description, background normalization of the top quark, $\gamma \gamma$ and diboson contributions, and a $1.8\%$ uncertainty on the total luminosity [37] introduce a combined uncertainty of less than $0.5\%$ on the total predicted yields. The different sources of uncertainty are added in quadrature, resulting in a total $15\%$ uncertainty on the background prediction.

In Table I, the observed number of events and the SM predictions are presented. The data are in agreement with the SM background-only hypothesis with a $p$ value of 0.2. Figure 1 shows the measured $E_T^{\text{miss}}$ distribution compared to the background predictions. The results are expressed in terms of model-independent $90\%$ and $95\%$ confidence level (C.L.) upper limits on the visible cross section, defined as the production cross section times acceptance times efficiency ($\sigma \times A \times \epsilon$), using the CL$_{s}$ modified frequentist approach [38] and considering the systematic uncertainties on the SM backgrounds and on the integrated luminosity. Values of $\sigma \times A \times \epsilon$ above $5.6 \ \text{fb}$ and $6.8 \ \text{fb}$ are excluded at $90\%$ C.L. and $95\%$ C.L., respectively. Typical event selection efficiencies of $\epsilon \sim 75\%$ are found in simulated ADD and WIMP signal samples.

The results are translated into $95\%$ C.L. limits on the parameters of the ADD model. The typical $A \times \epsilon$ of the selection criteria is $20.0 \pm 0.4(\text{stat}) \pm 1.6(\text{syst}) \%$, approximately independent of $n$ and $M_D$. Experimental

<table>
<thead>
<tr>
<th>Background source</th>
<th>Prediction</th>
<th>$\pm(\text{stat})$</th>
<th>$\pm(\text{syst})$</th>
</tr>
</thead>
<tbody>
<tr>
<td>$Z(\rightarrow \nu \bar{\nu}) + \gamma$</td>
<td>93</td>
<td>$\pm 16$</td>
<td>$\pm 8$</td>
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<tr>
<td>$Z/\gamma^*(\rightarrow \ell^+ \ell^-) + \gamma$</td>
<td>0.4</td>
<td>$\pm 0.2$</td>
<td>$\pm 0.1$</td>
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<tr>
<td>$W(\rightarrow \ell \nu) + \gamma$</td>
<td>24</td>
<td>$\pm 5$</td>
<td>$\pm 2$</td>
</tr>
<tr>
<td>$W/Z + \text{jets}$</td>
<td>18</td>
<td>$\cdots$</td>
<td>$\pm 6$</td>
</tr>
<tr>
<td>Top</td>
<td>0.07</td>
<td>$\pm 0.07$</td>
<td>$\pm 0.01$</td>
</tr>
<tr>
<td>$WW, WZ, ZZ, \gamma \gamma$</td>
<td>0.3</td>
<td>$\pm 0.1$</td>
<td>$\pm 0.1$</td>
</tr>
<tr>
<td>$\gamma + \text{jets and multijet}$</td>
<td>1.0</td>
<td>$\cdots$</td>
<td>$\pm 0.5$</td>
</tr>
<tr>
<td>Total background</td>
<td>137</td>
<td>$\pm 18$</td>
<td>$\pm 9$</td>
</tr>
<tr>
<td>Events in data ($4.6 \ \text{fb}^{-1}$)</td>
<td>116</td>
<td></td>
<td></td>
</tr>
</tbody>
</table>

FIG. 1 (color online). The measured $E_T^{\text{miss}}$ distribution (black dots) compared to the SM (solid lines), SM + ADD (dashed lines), and SM + WIMP (dotted lines) predictions, for two particular ADD and WIMP scenarios.
uncertainties related to the photon, jet, and $E_T^{miss}$ scales and resolutions, the photon reconstruction, the trigger efficiency, the pileup description, and the luminosity introduce a 6.8% uncertainty on the signal yield. Uncertainties related to the modeling of the initial- and final-state gluon radiation translate into a 3.5% uncertainty on the ADD signal yield. Systematic uncertainties due to PDFs result in a 0.8% to 1.4% uncertainty on the signal $A \times \epsilon$ and a 4% to 11% uncertainty on the signal cross section, increasing as $n$ increases. Variations of the renormalization and factorization scales by factors of 2 and $\frac{1}{2}$ introduce a 0.6% uncertainty on the signal $A \times \epsilon$ and an uncertainty on the signal cross section that decreases from 9% to 5% as $n$ increases.

Figure 2 shows the expected and observed 95% C.L. lower limits on $M_D$ as a function of $n$, as determined using the CL_s method and considering uncertainties on both signal and SM background predictions. Values of $M_D$ below 1.93 TeV ($n = 2$), 1.83 TeV ($n = 3$ or 4), 1.86 TeV ($n = 5$), and 1.89 TeV ($n = 6$) are excluded at 95% C.L. The observed limits decrease by 3% to 2% after considering the $-1\sigma$ uncertainty from PDFs, scale variations, and parton shower modeling in the ADD theoretical predictions (dashed lines in Fig. 2). These results improve upon previous limits on $M_D$ from LEP and Tevatron experiments [1–3]. In this analysis, no weights are applied for signal events in the phase space region with $\hat{s} > M_D^2$, which is sensitive to the unknown ultraviolet behavior of the theory. For $M_D$ values close to the observed limits, the visible signal cross sections decrease by 15% to 75% as $n$ increases when truncated samples with $\hat{s} < M_D^2$ are considered. This analysis probes a kinematic range for which the model predictions are defined but ambiguous.

Similarly, 90% C.L. upper limits on the pair-production cross section of dark matter WIMP candidates are determined. The $A \times \epsilon$ of the selection criteria are typically 11.0 ± 0.2(stat) ± 1.6(syst)% for the D1 operator, 18.0 ± 0.3(stat) ± 1.4(syst)% for the D5 and D8 operators, and 23.0 ± 0.3(stat) ± 2.1(syst)% for the D9 operator, with a moderate dependence on $m_x$. Experimental uncertainties, as discussed above, translate into a 6.6% uncertainty on the signal yields. Theoretical uncertainties on initial- and final-state gluon radiation introduce a 3.5% to 10% uncertainty on the signal yields. The uncertainties related to PDFs result in 1.0% to 8.0% and 5.0% to 30% uncertainties on the signal $A \times \epsilon$ and cross section, respectively. Variations of the renormalization and factorization scales lead to a change of 1.0% to 2.0% and 8.0% in the signal $A \times \epsilon$ and cross section, respectively. In the case of the D1 (D5) spin-independent operator, values of $M_x$ below 31 and 5 GeV (585 and 156 GeV) are excluded at 90% C.L. for $m_x$ equal to 1 GeV and 1.3 TeV, respectively. Values of $M_x$ below 585 and 100 GeV (794 and 188 GeV) are excluded for the D5 (D9) spin-dependent operator for $m_x$ equal to 1 GeV and 1.3 TeV, respectively. These results can be translated into upper limits on the nucleon-WIMP interaction cross section using the prescription in Refs. [12,39]. Figure 3 shows 90% C.L. upper limits on the nucleon-WIMP cross section as a function of $m_x$. In the case of the D1 (D5) spin-independent interaction, nucleon-WIMP cross sections above $2.7 \times 10^{-39}$ cm$^2$ and $5.8 \times 10^{-34}$ cm$^2$ (2.2 × $10^{-39}$ cm$^2$ and 1.7 × $10^{-36}$ cm$^2$) are excluded at 90% C.L. for $m_x = 1$ GeV and $m_x = 1.3$ TeV, respectively. Spin-dependent interactions cross sections in the range $7.6 \times 10^{-41}$ cm$^2$ to $3.4 \times 10^{-37}$ cm$^2$ (2.2 × $10^{-41}$ cm$^2$ to $2.7 \times 10^{-38}$ cm$^2$) are excluded at 90% C.L. for the D8 (D9) operator and $m_x$ varying between 1 GeV and 1.3 TeV. The quoted observed limits on $M_x$ typically decrease by 2% to 10% if the $-1\sigma$ theoretical uncertainty is considered. This translates into a 10% to 50% increase of the quoted nucleon-WIMP cross section limits. The exclusion in the region 1 GeV < $m_x$ < 3.5 GeV (1 GeV < $m_x$ < 1 TeV) for spin-independent (spin-dependent)

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**FIG. 2 (color online).** Observed (solid lines) and expected (dashed-dotted lines) 95% C.L. limits on $M_D$ as a function of the number of extra spatial dimensions $n$ in the ADD model. The results are compared with previous results [1,3,6] (other lines). In [6], weights are applied that suppress the region with $\hat{s} > M_D^2$.

**FIG. 3 (color online).** 90% C.L. upper limits on the nucleon-WIMP cross section as a function of $m$ for spin-dependent (left) and spin-independent (right) interactions [12,39]. The results are compared with previous monojet and monophoton results at colliders [4,6,8] and results from direct detection experiments [11].
nucleon-WIMP interactions is driven by the results from collider experiments, with the assumption of the validity of the effective theory, and is still dominated by the monojet results. The cross section upper limits improve upon CDF results [4] and are similar to those obtained by the CMS experiment [5,6].

In summary, we report results on the search for new phenomena in events with an energetic photon and large missing transverse momentum in proton-proton collisions at $\sqrt{s} = 7$ TeV at the LHC, based on ATLAS data corresponding to an integrated luminosity of 4.6 fb$^{-1}$. The measurements are in agreement with the SM predictions for the background. The results are translated into model-independent 90% and 95% confidence level upper limits on $\sigma \times A \times \epsilon$ of 5.6 and 6.8 fb, respectively. The results are presented in terms of improved limits on $M_D$ versus the number of extra spatial dimensions in the ADD model and upper limits on the spin-independent and spin-dependent contributions to the nucleon-WIMP elastic cross section as a function of the WIMP mass.

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[15] ATLAS uses a cylindrical coordinate system about the beam axis with polar angle $\theta$ and azimuthal angle $\phi$. Anticlockwise beam direction defines the positive $z$ axis, while the positive $x$ axis is defined as pointing from the collision point to the center of the LHC ring and the positive $y$ axis points upwards. We define transverse energy $E_T = E \sin \theta$, transverse momentum $p_T = p \sin \theta$, and pseudorapidity $\eta = - \ln (\tan(\theta/2))$.
The strange and charm quark masses (relevant for the $D1$ operator) are set to 0.1 and 1.42 GeV, respectively.
J. Wetter, 161 C. Weydert, 55 K. Whalen, 29 S. J. Wheeler-Ellis, 163 A. White, 8 M. J. White, 86 S. White, 122a, 122b
S. K. Whalen, 29 S. J. Wheeler-Ellis, 163 A. White, 8 M. J. White, 86 S. White, 122a, 122b
S. R. Whitehead, 118 D. Whiteson, 163 D. Whittington, 60 F. Wicek, 115 D. Wicke, 175 F. J. Wickens, 129
A. Wildauer, 99 M. A. Wildt, 42a I. Wilhelm, 126 H. G. Wilkens, 30 J. Z. Will, 98 E. Williams, 35 H. H. Williams, 120
W. Willis, 35 S. Willocq, 84 J. A. Wilson, 15 M. G. Wilson, 143 A. Wilson, 87 I. Wingertner-Seez, 5 F. Winkelnkemper, 48
B. K. Wosiek, 39 J. Wotschack, 30 M. J. Woudstra, 82 K. W. Wozniak, 39 K. Wraight, 53 M. Wright, 53 B. Wrona, 73
S. L. Wu, 173 X. Wu, 49 Y. Wu, 33b, ll E. Wulf, 35 B. M. Wynne, 46 S. Xella, 36 M. Xiao, 136 S. Xie, 48 C. Xu, 33b, 5 D. Xu, 139
B. Yabsley, 150 S. Yacoob, 145a, mm M. Yamada, 65 H. Yamaguchi, 155 A. Yamamoto, 65 K. Yamamoto, 63 S. Yamamoto, 155

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