Precision Measurement of the $[0 \over b]$ Baryon Lifetime

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Precision Measurement of the $\Lambda_b^0$ Baryon Lifetime

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The ratio of the $\Lambda_b^0$ baryon lifetime to that of the $B^0$ meson is measured using 1.0 fb$^{-1}$ of integrated luminosity in 7 TeV center-of-mass energy $pp$ collisions at the LHC. The $\Lambda_b^0$ baryon is observed for the first time in the decay mode $\Lambda_b^0 \rightarrow J/\psi pK^-$, while the $B^0$ meson decay used is the well known $B^0 \rightarrow J/\psi \pi^+K^-$ mode, where the $\pi^+K^-$ mass is consistent with that of the $K^{\pm}(892)$ meson. The ratio of lifetimes is measured to be $0.976 \pm 0.012 \pm 0.006$, in agreement with theoretical expectations based on the heavy quark expansion. Using previous determinations of the $B^0$ meson lifetime, the $\Lambda_b^0$ lifetime is found to be $1.482 \pm 0.018 \pm 0.012$ ps. In both cases, the first uncertainty is statistical and the second systematic.

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Evaluations from experimental data of fundamental parameters, such as Cabibbo-Kobayashi-Maskawa (CKM) matrix elements [1], and limits on physics beyond that described by the standard model, often rely on theoretical input [2]. One of the most useful models, the heavy quark expansion (HQE) [3–5], is based on the operator product expansion [6]; it is used, for example, to extract values for $|V_{ub}|$ and $|V_{cb}|$ from measurements of inclusive semileptonic $B$ meson decays [7]. In the free quark model, the lifetimes of all $b$-flavored hadrons are equal, because the decay width is determined by the $b$ quark lifetime. This model is too naive, since effects of other quarks in the hadron are not taken into account [8]. Early predictions using the HQE, however, supported the idea that $b$-hadron lifetimes were quite similar, due to the absence of correction terms $O(1/m_b)$.

The ratio of the $\Lambda_b^0$ baryon lifetime to that of the $B^0$ meson is measured using 1.0 fb$^{-1}$ of integrated luminosity accumulated by the LHCb experiment in 7 TeV center-of-mass energy $pp$ collisions. The $\Lambda_b^0$ baryon is detected in the $J/\psi pK^-$ decay mode, while the $B^0$ meson is found in $J/\psi \pi^+K^-$ decays. Mention of a particular decay channel implies the additional use of the charge-conjugate mode. This $\Lambda_b^0$ decay mode has not been observed before. (Measurement of the branching fraction is under study, and will be reported in a subsequent publication.) On the other hand, the $B^0$ decay is well known, and we impose the further requirement that the invariant mass of the $\pi^+K^-$ combination be within $\pm 100$ MeV of the $K^{\pm}(892)$ mass (we work in units where $c = 1$) in order to simplify the simulation and reduce systematic uncertainties. These decays have the same decay topology into four charged tracks, thus, facilitating the cancellation of uncertainties.

The LHCb detector is a single-arm forward spectrometer covering the pseudorapidity range $2 < \eta < 5$, described in detail in Ref. [17]. Events selected for this analysis are triggered [18] by a $J/\psi \rightarrow \mu^+\mu^-$ decay, where the $J/\psi$ is required at the software level to be consistent with coming from the decay of a $b$ hadron by use either of impact parameter (IP) requirements or detachment of the $J/\psi$ from the associated primary vertex. The simulated events used in this analysis are produced using the software described in Ref. [19].

Events are preselected and then are further filtered using a multivariate analyzer based on the boosted decision tree (BDT) technique [20]. In the preselection, all hadron track candidates are required to have $p_T$ larger than 250 MeV, while for muon candidates, the requirement is more than 550 MeV. Events must have a $\mu^+\mu^-$ combination that forms a common vertex with $\chi^2 < 16$, and an invariant mass between $-48$ and $+43$ MeV of the $J/\psi$ mass. Candidate $\mu^+\mu^-$ combinations are then constrained to the $J/\psi$ mass for subsequent use in event selection. The two charged final state hadrons must have a vector summed

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of the $C_{22}$ DLL and shown in Fig. 1. The shapes are smooth and do not peak near the $\pi$ reflections, we take each of the candidates in the five degrees of freedom. This $b$-hadron candidate must have a momentum vector that, when parity inverted, points to the primary vertex within an angle smaller than 2.56$^\circ$. Particle identification requirements differ in the two modes. We use the difference in the logarithm of the likelihood, $\text{DLL}(h_1 - h_2)$, to distinguish between the two hypotheses: $h_1$ and $h_2$ as described in [21]. In the $\Lambda^0_b$ decay, the kaon candidate must have $\text{DLL}(K - \pi) > 4$ and $\text{DLL}(K - p) > -3$, while the proton must have $\text{DLL}(p - \pi) > 10$ and $\text{DLL}(p - K) > -3$. For the $B^0$ decay, the requirements on the pion candidate are $\text{DLL}(\pi - \mu) > -10$ and $\text{DLL}(\pi - K) > -10$, while $\text{DLL}(K - p) > 0$ is required for the kaon.

The BDT selection is based on the minimum DLL($\mu - \pi$) of the $\mu^+$ and $\mu^-$ candidates, the $p_T$ of each of the two charged hadrons, and their sum, the $\Lambda^0_b$'s $p_T$, the $\Lambda^0_b$ vertex $\chi^2$, and the impact parameter $\chi^2$ of the $\Lambda^0_b$ candidate, where the latter results from calculating the difference in $\chi^2$ by using the hypothesis that the IP is zero. These variables are chosen with the aim of having the selection efficiency be independent of decay time. The BDT is trained on a simulated sample of either $\Lambda^0_b \rightarrow J/\psi pK^-$ signal events and a background data sample from the mass sidebands of the $\Lambda^0_b$ signal peak. It is then tested on independent samples from the same sources. The BDT selection is implemented to maximize $S^2/(S + B)$, where $S$ indicates the signal and $B$ the background event yields. This optimization includes the requirement that the $\Lambda^0_b$ baryon decay time be greater than 0.5 ps. The same BDT selection is used for the $B^0 \rightarrow J/\psi \pi^- K^+$ mode.

The $J/\psi pK^-$ mass distribution after the BDT selection is shown in Fig. 1. There is a large and significant signal. Backgrounds can be combinatorial in nature, but can also be formed by reflections from $B$ meson decays where the particle identification fails. As long as these backgrounds do not peak near the $\Lambda^0_b$ mass they cannot cause incorrect determinations of the $\Lambda^0_b$ signal yield. The shapes of the main $B$ meson reflections are determined from simulation and shown in Fig. 1. The shapes are smooth and do not peak in the signal region. To estimate the contributions of the reflections, we take each of the candidates in the $J/\psi pK^-$ sideband regions 60–200 MeV on either side of the $\Lambda^0_b$ mass peak, reassign proton to kaon and pion mass hypotheses, respectively, and fit the resulting signal peaks determining signal yields of 5576 ± 95 $B^0$ and 1769 ± 192 $B^0$ decays. To translate these yields to those within ±20 MeV of the $\Lambda^0_b$ peak, we use simulations of $\bar{B}^0 \rightarrow J/\psi K^+ K^-$ with the $K^+ K^-$ mass distribution matched to that obtained in our previous analysis of this final state [22], and a simulation of $B^0 \rightarrow J/\psi \pi^+ K^-$ decays, leading to $1186 \pm 35$ $J/\psi K^+ K^-$ and $308 \pm 33$ $J/\psi \pi^+ K^-$ reflected decays, respectively.

To determine the $\Lambda^0_b$ signal yield, we perform an unbinned maximum likelihood fit to the $J/\psi pK^-$ invariant mass spectrum shown in Fig. 1 in the region between 5500 and 5750 MeV. The fit function is the sum of the $\Lambda^0_b$ signal component, combinatorial background, and the contribution from the $B^0 \rightarrow J/\psi K^+ K^-$ and $B^0 \rightarrow J/\psi \pi^+ K^-$ reflections. The signal is modeled by a triple-Gaussian function with common means; the effective rms width is 5.5 MeV. The combinatorial background is described by an exponential function. The event yields of the reflections are included in the fit as Gaussian constraints. The fit gives 15 581 ± 178 signal and 5535 ± 50 combinatorial background candidates together with 1235 ± 35 $\bar{B}^0 \rightarrow J/\psi K^+ K^-$ and 313 ± 26 $B^0 \rightarrow J/\psi \pi^+ K^-$ reflection candidates within ±20 MeV of the $\Lambda^0_b$ mass peak.

To view the background subtracted $pK^-$ mass spectrum, we perform fits, as described above, to the $m(J/\psi pK^-)$ distributions in bins of $m(pK^-)$ and extract the signal yields within ±20 MeV of the $\Lambda^0_b$ mass peak. The resulting $pK^-$ mass spectrum is shown in Fig. 2. A distinct peak is observed in the $pK^-$ invariant mass distribution near 1520 MeV, together with the other resonant and nonresonant structures over the entire kinematical region. The peak corresponds to the $\Lambda(1520)$ resonance [23]. Simulations of the $\Lambda^0_b$ decay are weighted to reproduce this mass distribution.

The $J/\psi \pi^+ K^-$ mass spectrum, after the BDT selection, is shown in Fig. 3. There is a large signal peak at the $B^0$ mass and a much smaller one at the $\bar{B}^0$ mass. Triple-Gaussian functions, each with common means, are used to fit the signal peaks; the effective rms width is 6.7 MeV. An exponential function is used to fit the combinatorial background.
FIG. 2. Background subtracted $m(pK^-)$ distribution obtained by fitting the $m(J/\psi pK^-)$ distribution in bins of $m(pK^-)$.

The mass fit gives $97,506\pm 447$ signal and $3660\pm 74$ background candidates within $\pm 20$ MeV of the $B^0$ mass peak. Reflections are possible from both $B^0 \rightarrow J/\psi K^+ K^-$ and $\Lambda^0_b \rightarrow J/\psi pK^-$ decays. Following the same procedure as outlined above using the sidebands of the $B^0$ signal, we find no evidence of a reflection from the $B^0$ state and a small, nonpeaking, contribution of $506\pm 19$ events from the $\Lambda^0_b$ state, in the $B^0$ signal region, that is ignored.

The decay time for each candidate is given by $t = m d \cdot \vec{p}/|\vec{p}|^2$, where $m$ is the mass, $d$ the distance vector from the primary vertex to the decay point, and $\vec{p}$ is the measured $b$ hadron momentum. Here, we do not constrain the two muons to the $J/\psi$ mass to avoid systematic biases. The decay time resolutions are 40 fs for the $\Lambda^0_b$ decay and 37 fs for the $B^0$ decay. In addition, the decay time acceptances are also almost equal. For equal acceptances, the ratio of events, $R(t)$, as a function of decay time is given by

$$R(t) = \frac{N_{B^0}(0) e^{-t/\tau_{B^0}}}{N_{B^0}(0) e^{-t/\tau_{B^0}}} = R(0) e^{-t/\Delta_{AB}}$$

where $\Delta_{AB} = (1/\tau_{B^0} - 1/\tau_{B^0})$. Effects of the different decay time resolutions in the two modes are negligible above 0.5 ps. First order corrections for a decay time dependent acceptance ratio can be taken into account by modifying Eq. (1) with a linear function

$$R(t) = R(0)[1 + at] e^{-t/\Delta_{AB}}$$

where $a$ represents the slope of the acceptance ratio as a function of decay time.

The decay time acceptances for both modes are determined by simulations that are weighted to match either the $pK^-$ or $\pi^+ K^-$ invariant mass distributions seen in data, as well as to match the measured $p$ and $p_T$ distributions of the $b$ hadrons. In addition, we further weight the samples so that the simulation matches the hadron identification efficiencies obtained from $D^{*+} \rightarrow \pi^+(D^0 \rightarrow \pi^+ K^-)$ events for pions and kaons, and $\Lambda^0 \rightarrow p\pi^-\pi^0$ for protons.

The ratio of the decay time acceptances is shown in Fig. 4. Here, we have removed the minimum requirement on decay time so we can view the distributions in the region close to zero time. The individual acceptances in both cases can be described with a linear function above 0.5 ps. In order to minimize possible systematic effects, we use candidates with decay times larger than 0.6 ps. We also choose an upper time cut of 7.0 ps, because the acceptance is poorly determined beyond this value. The acceptance ratio is fitted with a linear function between 0.6 and 7.0 ps. The slope is $a = 0.0033 \pm 0.0024$ ps$^{-1}$, and the $\chi^2$/number of degrees of freedom (NDF) of the fit is 81/62.

We determine the event yields in both decay modes by fitting the invariant mass distributions in 16 bins of decay time, each bin 0.4 ps wide, using the same signal and
we determine $/C_28$ where we use the world average value of the decay time within a bin determined by the fitted signal yields in both modes are placed at the average $0$.

The decay time ratio difference is $p_B$ shown in Fig.5(b). The distribution fitted with the function given in Eq. (2) is shown in Fig.5(a). Here, the background shapes obtained in the aforementioned mass fits. Since the bin size is approximately ten times the resolution, there is no effect due to the small difference of time resolution ($<7\%$) between the two modes. The resulting distributions are shown in Fig. 5(a). Here, the fitted signal yields in both modes are placed at the average of the decay time within a bin determined by the $B^0$ data in order to correct for the exponential decrease of the decay time distributions across the bin. The decay time ratio distribution fitted with the function given in Eq. (2) is shown in Fig. 5(b). The $\chi^2$/NDF of the fit is 18/14, with a $p$ value of 21%. The fitted value of the reciprocal lifetime difference is

$$\Delta_{AB} = 16.4 \pm 8.2 \pm 4.4 \text{ ns}^{-1}.$$  

Whenever two uncertainties are quoted, the first is the statistical and the second systematic; the latter will be discussed below. Numerically, the ratio of lifetimes is

$$\frac{\tau_{\Lambda_b^0}}{\tau_{B^0}} = \frac{1}{1 + 3.0/\Delta_{AB}} = 0.976 \pm 0.012 \pm 0.006,$$

where we use the world average value $\tau_{B^0} = 1.519 \pm 0.007$ ps [23]. Multiplying the lifetime ratio by this value, we determine

$$\tau_{\Lambda_b^0} = 1.482 \pm 0.018 \pm 0.012 \text{ ps}.$$  

Our result is consistent with, but higher and more accurate than, the current world average of $1.429 \pm 0.024$ ps [23].

The absolute systematic uncertainties are listed in Table I. There is an uncertainty due to the decay time range used because of the possible change of the acceptance ratio at short decay times. This uncertainty is ascertained by changing the fit range to be 1–7 ps and using the difference with the baseline fit. To determine the acceptance slope uncertainty we vary the value of $a$ by its error determined from the fit to the simulation samples and propagate this change to the results. For the signal shape uncertainty, we repeat the measurement of $\Delta_{AB}$ using a double-Gaussian signal shape in the mass fits. The uncertainty in the background parametrization is assigned by letting the background parameters vary in the fits to the time dependent yields and comparing the difference in final results. Effects of changes in the acceptance for the $\Lambda_b^0$ mode due to the angular decay distributions are evaluated by weighting the simulation by the observed $pK^-$ helicity angle in addition to the $pK^-$ invariant mass, and redoing the analysis. The acceptance function uncertainty is evaluated by using a parabola instead of a linear function. The total systematic uncertainty is obtained by adding all of the elements in quadrature.

In conclusion, our value for $\tau_{\Lambda_b^0}/\tau_{B^0} = 0.976 \pm 0.012 \pm 0.006$ shows that the $\Lambda_b^0$ and $B^0$ lifetimes are indeed equal to within a few percent, as the original advocates of the HQE claimed [3,4,9], without any need to find additional corrections. Adding both uncertainties in quadrature, the lifetimes are consistent with being equal at the level of 1.9 standard deviations; thus, we do not exclude that the $\Lambda_b^0$ baryon has a longer lifetime than the $B^0$ meson. Using the world average measured value for the $B^0$ lifetime, we determine $\tau_{\Lambda_b^0} = 1.482 \pm 0.018 \pm 0.012$ ps.

We are thankful for many useful and interesting conversations with Prof. Nikolai Uraltsev (now deceased) who contributed greatly to theories describing heavy hadron lifetimes. We express our gratitude to our colleagues in the CERN accelerator departments for the excellent performance of the LHC. We thank the technical and administrative staff at the LHCb institutes. We acknowledge

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