Search for the production of ZW and ZZ boson pairs decaying into charged leptons and jets in p[bar over p] collisions at s = 1.96TeV

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<td>As Published</td>
<td><a href="http://dx.doi.org/10.1103/PhysRevD.88.092002">http://dx.doi.org/10.1103/PhysRevD.88.092002</a></td>
</tr>
<tr>
<td>Publisher</td>
<td>American Physical Society</td>
</tr>
<tr>
<td>Version</td>
<td>Final published version</td>
</tr>
<tr>
<td>Accessed</td>
<td>Mon Mar 11 17:15:28 EDT 2019</td>
</tr>
<tr>
<td>Citable Link</td>
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Search for the production of \(ZW\) and \(ZZ\) boson pairs decaying into charged leptons and jets in \(p\bar{p}\) collisions at \(\sqrt{s} = 1.96\) TeV

1550-7998/2013/88(9)/092002(17) 092002-1 © 2013 American Physical Society

(CDF Collaboration)
We present a measurement of the production cross section for $ZW$ and $ZZ$ boson pairs in final states with a pair of charged leptons, from the decay of a $Z$ boson, and at least two jets, from the decay of a $W$ or $Z$ boson, using the full sample of proton-antiproton collisions recorded with the CDF II detector at the Tevatron, corresponding to $8.9 \text{ fb}^{-1}$ of integrated luminosity. We increase the sensitivity to vector-boson decays into pairs of quarks using a neural-network discriminant that exploits the differences between the

\[ \text{SEARCH FOR THE PRODUCTION OF } ZW \text{ AND ZZ} \ldots \]
I. INTRODUCTION

The standard model (SM) offers precise predictions of production rates associated with self-interactions of the gauge bosons [1]. Differences between these predictions and measured diboson production cross sections may indicate the presence of non-SM physics [2,3], even specifically in hadronic final states [4]. Additionally, since hadronic final states in diboson production are similar to those from associated Higgs boson production ($p\bar{p} \rightarrow VH + X$ where $V = W, Z$), the analysis techniques used to measure diboson production in partially hadronic final states are relevant to searches for associated Higgs boson production.

Measurements of diboson production are typically difficult due to the small production cross sections of the order of 10 pb or less [1]. Furthermore, measurements of decay channels where one $W$ or $Z$ boson decays hadronically are particularly challenging at hadron colliders: although expected event yields are larger than those in purely leptonic decay channels due to the higher hadronic decay ($V \rightarrow q\bar{q}'$) branching ratio, the expected backgrounds from QCD multijet processes and $V +$ jets production are also much greater. Experiments at the Tevatron have previously measured the cross sections for pair production of two leptons to originate from the decay of a boson and $V +$ jets production are also much greater. Experiments at the Tevatron have previously measured the cross sections for pair production of gauge bosons in partially hadronic decay channels [5–8], but all of these measurements have included contributions from $WW$ production, which has a higher cross section than that for combined $ZW$ and $ZZ$ production. Searches using identification of $b$-quark decays in the final states ($b$-tagging) to increase sensitivity to events with $Z \rightarrow b\bar{b}$ decays have been performed [9], but have not yet provided observations of $ZZ$ production in partially hadronic decay channels.

We present a study of $ZV$ production from a final state with two leptons and at least two jets [10]. We require the two leptons to originate from the decay of a $Z$ boson and search for associated $V \rightarrow q\bar{q}'$ decays by performing a fit to the dijet invariant mass ($m_{jj}$) spectrum of the two leading-$E_T$ jets. To maximize sensitivity to diboson production, we separate events into three channels: a heavy-flavor-tagged channel, largely sensitive to $ZZ \rightarrow \ell^+\ell^- b\bar{b}$ decays; a light-flavor-tagged channel, which uses a new artificial-neural-network-based discriminant to preferentially select events with quark-like jets over gluon-like jets; and an untagged channel, which contains the remaining events that pass the event-selection requirements.

The final fit to the $m_{jj}$ spectra is performed simultaneously across all three channels.

This paper is organized as follows: in Sec. II we briefly describe the CDF II detector; in Sec. III we describe the data sets and event-selection requirements that are used in the $ZV$ search; in Sec. IV we show the derivation of new jet energy corrections to Monte Carlo simulations that account for differences in the observed energy scales of quark and gluon jets; in Sec. V we provide details of a new neural-network-based discriminant that identifies jets more likely to originate from quarks than from gluons; and, in Sec. VI, we describe the signal-extraction method, and report the results of the $ZV$ search.

II. THE CDF II DETECTOR

The CDF II detector is described in detail elsewhere [12]. The detector is cylindrically symmetric around the Tevatron beam line. Tracking detectors are installed around the interaction point to reconstruct the trajectories of charged particles (tracks). The tracking systems are located within a superconducting solenoid that produces a 1.4 T magnetic field aligned with the $p\bar{p}$ beams. Around the outside of the solenoid, calorimeter modules arranged in a projective-tower geometry measure the energies of charged and neutral particles. Drift chambers outside the calorimeter are used to detect muons, which typically deposit little energy in the calorimeter.

The central outer tracker (COT) is a 3.1-m-long open-cell drift chamber that has 96 measurement layers in the region between 0.40 and 1.37 m from the beam axis, providing full track coverage in the pseudorapidity region $|\eta| < 1.0$. Sense wires are arranged in eight alternating axial and $\pm 2^\circ$ stereo “superlayers” with 12 wires each. The position resolution of a single drift-time measurement is about 140 $\mu$m. A five-layer double-sided silicon microstrip detector (SVX) covers the region between 2.5 to 11 cm from the beam axis. Three separate SVX barrel modules along the beam line cover a length of 96 cm, approximately 90% of the luminous beam interaction region. Three of the five layers combine an $r$-$\phi$ measurement on one side and a 90° stereo measurement on the other, and the remaining two layers combine an $r$-$\phi$ measurement with a small-angle (±1.2°) stereo measurement. The typical silicon hit resolution is 11 $\mu$m. An intermediate silicon...
layer at a radius of 22 cm from the beam line in the central region links tracks in the COT to hits in the SVX. The fiducial range of the silicon detector extends to pseudorapidity magnitude $|\eta| < 2.0$.

Calorimeter modules are located outside the central tracking volume and solenoid. The inner electromagnetic layers consist of lead sheets interspersed with scintillators, while the outer hadronic layers consist of scintillators sandwiched between steel sheets. The calorimeter is split between central barrel ($|\eta| < 1.1$) and forward end-plug ($1.1 < |\eta| < 3.6$) sections. Individual towers in the central barrel subtend 0.1 in $|\eta|$ and 15$^\circ$ in $\phi$. The sizes of the towers in the end-plug calorimeter vary with $|\eta|$, subbending 0.1 in $|\eta|$ and 7.5$^\circ$ in $\phi$ at $|\eta| = 1.1$, and 0.5 in $|\eta|$ and 15$^\circ$ in $\phi$ at $|\eta| = 3.6$. The energy resolution in the electromagnetic calorimeters is 14%/\sqrt{E_T} in the central barrel and 16%/\sqrt{E} @ 1% in the forward end-plug section, with the energies in units of GeV. The single-particle energy resolution in the hadronic calorimeters, measured using pions, ranges from 75%/\sqrt{E} in the central barrel to 80%/\sqrt{E} @ 5% in the forward end-plug section, with the energies expressed in units of GeV.

The hadronization of quarks and gluons produced in the interaction leads to collimated groups of high-momentum particles called jets. These jets, along with photons and electrons, leave isolated energy deposits in contiguous groups of calorimeter towers, which can be summed together into an energy cluster. Electrons and photons are identified as isolated, mostly electromagnetic clusters, and quality requirements may be placed on the presence of a high-$p_T$ track geometrically matched to the cluster to more accurately identify electrons. Jets are identified as electromagnetic and hadronic clusters with the combined electromagnetic (EM) fraction $E_{EM}/E_{total} = E_{EM}/(E_{EM} + E_{had}) < 0.9$, clustered using the JETCLU cone algorithm [13] with a fixed cone size of $\Delta R = \sqrt{(\Delta \eta)^2 + (\Delta \phi)^2} = 0.4$.

Outside the calorimeters, drift chambers detect muons. A four-layer stack of planar drift chambers detect muons with $p_T > 1.4$ GeV/c, and another four layers of drift chambers behind 60 cm of steel detect muons with $p_T > 2.0$ GeV/c. Both systems cover a region of $|\eta| < 0.6$, though they have different structure and their geometrical coverages do not overlap exactly. Muons in the region between $0.6 < |\eta| < 1.0$ pass through at least four drift layers arranged within a conic section outside of the central calorimeter. Muons are identified as either COT tracks that extrapolate to hits in the muon detectors, or as isolated tracks unmatched to hits in the muon detectors that satisfy tighter tracking-quality requirements and extrapolate to calorimeter energy depositions consistent with a minimum ionizing particle.

III. DATA SET AND EVENT SELECTION

We analyze the full data set of $p\bar{p}$ collisions collected by the CDF II detector. We require events to be collected from periods when the tracking systems, calorimeters, and muon detectors were all functioning properly, corresponding to an integrated luminosity of 8.9 fb$^{-1}$. Events are selected via a number of high-$E_T$ electron and high-$p_T$ muon online event-selection requirements (triggers). The majority of these triggers require at least one electron (muon) with $E_T > 18$ GeV ($p_T > 18$ GeV/c). We require the events to contain two electrons (muons) with $E_T > 20$ GeV ($p_T > 20$ GeV/c) and determine the trigger selection and event reconstruction efficiencies by comparing the number of data and simulated $Z \rightarrow \ell\ell$ events containing exactly one jet with $E_T > 20$ GeV.

For the final analysis, we select events with at least two leptons, and two or more jets. In the unlikely case that more than two charged leptons are reconstructed, we select the two leptons with the highest $p_T$. In addition to the $p_T$ requirements on the leptons, we require leptons associated with well-reconstructed tracks (central electrons, $|\eta| < 1$, and all muons) to be of opposite charge, and a reconstructed dilepton invariant mass, $m_{\ell\ell}$, consistent with the mass of the Z boson, $76 < m_{\ell\ell} < 106$ GeV/c$^2$. We require both leading-$E_T$ jets to have $E_T > 25$ GeV and $|\eta| < 2.0$, and to be spatially separated from the reconstructed leptons ($\Delta R > 0.4$). Additionally, the two jets must be separated by $\Delta R > 0.7$. Finally, as the final state should contain no particles that are not reconstructed in the detector, we also require that the missing transverse energy, $E_T$ [14], is less than 20 GeV.

After this selection, three major sources of background contribute. The dominant background comes from the production of a Z boson, decaying to an $e^+e^-$ or $\mu^+\mu^-$ pair, in association with two jets. Simulated events generated using ALPGEN [15], interfaced with PYTHIA [16] for showering, are used to estimate this background. The production cross sections for $Z + b\bar{b}$ processes are normalized to experimental measurements [17].

Another significant background results from jets misidentified as leptons. The contributions from these lepton fakes are estimated using data-driven methods. For muons, we use events with same-sign muon pairs (rather than opposite-sign) that otherwise satisfy all event-selection requirements. For electrons, we derive a misidentification rate representing the likelihood for a jet to be misidentified as an electron, as a function of jet $E_T$ and $\eta$, using jet-triggered data with minimal contributions from events with electrons. This rate is then applied to all possible electron-jet pairs in events from the high-$p_T$ electron data set, where the jet is then treated as a second electron, and the event-selection requirements are otherwise applied normally.

While the requirement to have low $E_T$ reduces its total contribution, top-quark-pair production, where each top quark decays into a leptonic final state ($t\bar{t} \rightarrow W^+bW^-\bar{b} \rightarrow \ell^+\nu_\ell b\ell^-\bar{\nu}_\ell \bar{b}$), contributes events to the final event sample, especially in the heavy-flavor-tagged channel. We estimate $t\bar{t}$ contributions using PYTHIA with $\sigma_{tt} = 7.5$ pb and
$m_t = 172.5$ GeV/$c^2$. Finally, ZW and ZZ signal events are also modeled using PYTHIA.

### IV. JET-ENERGY CALIBRATION

The energies of jets, measured in the calorimeter, are corrected to account for a number of effects that distort the true jet energy. These effects include changes in calorimeter performance as a function of $|\eta|$ and time, contributions from multiple $p\bar{p}$ interactions per beam crossing (pileup), contributions from the other partons in the interacting proton and antiproton (underlying event), the nonlinear response of the calorimeter, and energy radiated outside of the jet cone. These jet-energy-scale (JES) corrections are described in detail in Ref. [18].

These energy corrections, however, do not attempt to account for potential differences in the modeled calorimeter response to jets originating from quarks and gluons. For example, the largest corrections modify the energy scale of the jets to more accurately match that of the initial parton energies and their resulting particle jets, and are derived using PYTHIA [16] dijet Monte Carlo (MC) simulations, but independently of the initiating type of parton. Differences in the response of gluon and quark jets between simulation and data lead to differences in the measured energies for these jets that are not covered by the previously assigned systematic uncertainties on the JES [19].

We derive a data-driven correction for the response to quark and gluon jets in simulated events using two independent samples of jets with different compositions of quark- and gluon-initiated jets. In these samples, we derive a correction to the jet energy by matching the transverse energy of the jets against particles of known transverse momentum. We use events where a jet recoils against a high-$E_T$ photon—a sample rich in quark jets (based on simulations modeled using PYTHIA)—and utilize the significant number of $Z \rightarrow \ell^+ \ell^-$ jet events available in the full CDF data set, which have a larger fraction of gluon jets. The quark and gluon content of these two samples differ due to the difference in mass between the Z boson and the photon: because the Z boson mass is higher, the initial partons of the production process typically carry a higher fraction of the momentum of the proton than those involved in the production of high-energy photons. This leads to a difference in the quark and gluon content of these samples.

To derive a correction, we construct the balance of the jet with these more accurately measured reference particles,

$$K_{Z,\gamma} = (E_T^{\text{jet}} / p_T^{Z,\gamma}) - 1.$$  

(1)

For unbiased measurements of the jet energy, the balance will equal zero. Samples of jets with nonzero balance could be corrected with a jet-energy correction factor of $1/(K_{Z,\gamma} + 1)$. However, rather than derive independent JES corrections for quark and gluon jets in data and simulation, we compare the balance in data and simulation and derive an additional correction to be applied to simulated jets, dependent upon whether these jets are matched to quarks or gluons. The correction to simulated quark-jet energies is $+1.4 \pm 2.7\%$, while the correction to gluon-jet energies is much larger: $-7.9 \pm 4.4\%$.

#### A. Data set and event selection

The data set and event selection for the Z-jet balancing sample largely follow those described in Sec. III. We require two leptons consistent with resulting from the decay of a Z boson and exactly one jet with $E_T > 15$ GeV and no other jets with (uncorrected) $E_T > 3$ GeV within $|\eta| < 2.4$. Additionally, we ensure that the Z boson and jet are azimuthally opposite (back-to-back) by requiring their azimuthal separation to exceed 2.8 radians, and we require that $p_T^Z > 10$ GeV/$c$.

For the $\gamma$-jet balancing sample, we closely mirror the selection requirements described in Ref. [18]. We use events collected with an isolated-central-photon trigger over the same period of time as that of the high-$p_T$ lepton samples. We compare these data to PYTHIA simulations of both $\gamma +$ jet production as well as dijet production, which also contributes to the $\gamma$-jet balancing sample.

To match the selection requirements of the isolated-central-photon trigger, we require $E_T^\gamma > 27$ GeV and $0.2 < |\eta_\gamma| < 0.6$ in both data and MC simulation. To decrease the contribution from dijet production, where a jet mimics the photon selection, we require the energy in the calorimeter and momentum in the tracking system contained within a cone of $R = 0.4$ around the photon to be less than 1 GeV and 2 GeV/$c$, respectively. As in the Z-jet balancing sample, we require events to have exactly one measured jet with $E_T > 15$ GeV and no other jets with (uncorrected) $E_T > 3$ GeV within $|\eta| < 2.4$. We also demand the $\Delta \phi$ between the jet and photon to be larger than 3.0 radians. We further reduce contamination by vetoing events with more than one reconstructed interaction point, and by removing events with $E_T^\gamma / E_T^\gamma > 0.8$, which likely contain activity from cosmic rays.

#### B. Determination of corrections

We derive separate corrections for the quark- and gluon-jet energy scales in data and simulation using the Z-jet and $\gamma$-jet balancing samples in the following way. The balances of the Z-jet and $\gamma$-jet systems ($K_Z$ and $K_\gamma$, respectively) can be expressed as linear combinations of independent quark- and gluon-balance variables ($K_q$ and $K_g$, respectively), weighted by the sample-specific quark and gluon fractions ($F^{q,g}_{Z,\gamma}$),

$$K_Z = F^q_Z K_q + F^g_Z K_g = F^q_Z K_q + (1 - F^q_Z) K_g.$$  

(2)

$$K_\gamma = F^q_\gamma K_q + F^g_\gamma K_g = F^q_\gamma K_q + (1 - F^q_\gamma) K_g.$$  

(3)

Rewriting these expressions by solving for $K_q$ and $K_g$, we find

...
\[ K_q = \frac{1}{F_q^2 - F_Z^2} \left[ (1 - F_{Zq}^2) K_{Zq} - (1 - F_{Zg}^2) K_{Zg} \right]. \]  
\[ K_g = \frac{1}{F_g^2 - F_Z^2} \left[ F_{Zq}^2 K_{Zq} - F_{Zg}^2 K_{Zg} \right]. \]

These expressions apply separately to experimental data and simulated data, yielding a different balance in data and Monte Carlo simulation (\( K_{q\text{data}} \) and \( K_{q\text{MC}} \)), respectively, and may include a dependence on the energy of the jet, with \( F_{Zq}^\gamma = F_{Zq}^{\gamma}(E_T^{\text{jet}}) \) and consequently \( K = K(E_T^{\text{jet}}) \).

In order to solve for \( K_{q} \) and \( K_{g} \), we require knowledge of \( K_{Zq} \) and \( F_{Zq}^\gamma \). We extract a value of \( K_{Zq} \) as a function of \( E_T^{\text{jet}} \) by constructing the balancing distribution, as defined in Eq. (1), in ranges of \( E_T^{\text{jet}} \), and fit the core of the distribution with a Gaussian shape. We perform these fits separately in data and simulation and use the mean and uncertainty on the mean of the fitted Gaussian shape as the corrections that need to be applied to simulated jets in order to best match the energy scale of the data. These corrections that need to be applied to simulated jets in order to best match the energy scale of the data. These MC simulation corrections are defined as \( (K_{q\text{data}} + 1)/(K_{q\text{MC}} + 1) \) for quark jets and \( (K_{g\text{data}} + 1)/(K_{g\text{MC}} + 1) \) for gluon jets, as shown in Fig. 3. Due to the photon trigger used to select the \( \gamma \)-jet balancing sample, reliable balancing information is not available for jets with transverse energies smaller than 27.5 GeV in that sample, limiting the range of applicability of the corrections. Since we are interested in jets of energies extending down to 20 GeV, we extrapolate the quark-jet energy correction derived for jets with \( E_T > 27.5 \) GeV to lower jet energies and use the \( Z \)-jet balancing sample to extract a gluon correction assuming this extrapolated quark correction.

The quark and gluon corrections’ dependence on jet energy are accurately modeled by a constant for jets with \( E_T > 15 \) GeV. Quark jet energies in simulation should be increased by approximately 1.4% to more accurately match the data, while gluon jet energies should be decreased by approximately 7.9%.

![Z-jet balance](image1.png)
![\( \gamma \)-Jet balance](image2.png)

**FIG. 1** (color online). Balancing distributions, (a) \( K_Z \) and (b) \( K_g \), in data and MC simulation as a function of \( E_T^{\text{jet}} \). The uncertainties include solely the contribution from the fluctuations in the mean of a Gaussian fit to the balancing distributions in bins of \( E_T^{\text{jet}} \).
FIG. 2 (color online). Derived balancing variable for (a) quark jets, $K_q$, and (b) gluon jets, $K_g$, in data and MC simulation as a function of $E_T^{\text{jet}}$. The uncertainties on each point are from the uncertainties from the mean of the Gaussian fit and the uncertainties on the quark fractions, added in quadrature.

FIG. 3 (color online). Derived correction for simulated quark jets and gluon jets as a function of $E_T^{\text{jet}}$. The uncertainties shown are statistical only. The short dashed lines represent corrections derived using both quark and gluon templates from simulation. We take the average deviation of the value determined for this quark-gluon discriminant from the nominal MC simulation value as a systematic uncertainty on $F_q$, constant across jet $E_T$. This uncertainty is approximately 10%. Here we vary the calculated quark fraction in the data $Z$-jet balancing sample by $\pm 10\%$ and recalculate the corrections for quark and gluon jets. This translates to an uncertainty of $\pm 0.6\%$ for quark-jet energies and $\pm 2.1\%$ for gluon-jet energies.

(3) $F_q^2$: We follow a similar procedure of fitting the quark-gluon discriminant parameter in the $\gamma$-jet sample and obtain a similar uncertainty of $\pm 10\%$ on the quark fraction. This translates to an uncertainty of $\pm 1.8\%$ for quark-jet energies, and $\pm 2.7\%$ for gluon-jet energies.

(4) Low $E_T$ extrapolation: We check the dependence of the gluon-jet energy corrections on the assumed quark-jet corrections for low-$E_T$ jets by varying the quark-jet $E_T$ for these jets by $\pm 2\%$. We see a small change in the gluon-energy corrections, which translates to $\pm 0.4\%$ of the gluon-jet energy.

(5) Number of interaction vertices dependence: The $\gamma$-jet balancing sample incorporates a requirement on the number of reconstructed interaction vertices to reduce contamination from pileup. The $Z$-jet balancing sample does not have such a requirement, due to a much smaller background contribution and in order to retain as many events as possible. We check for any bias in the corrections resulting from the effect of this requirement by looking for any shift in the corrections when the requirement is placed on the $Z$-jet balancing sample. We see a change in the quark-jet energies of $\pm 0.2\%$, and the gluon-jet energies of $\pm 1.2\%$.

The uncertainties are summarized in Table I. Because the corrections shift the energy response in the simulation to more accurately match the data, the quark-jet and gluon-jet...
energy correction uncertainties are anticorrelated. The uncertainties are similar in magnitude to the default CDF jet energy scale uncertainties [18].

V. ARTIFICIAL NEURAL NETWORK QUARK-TO-GLUON DISCRIMINANT

In this analysis, we search for two high-\(p_T\) leptons from the decay of a \(Z\) boson and two jets from a \(W \rightarrow q\bar{q}\) or \(Z \rightarrow q\bar{q}\) decay. Thus, the two signal jets are quark jets. Conversely, the dominant background, two jets produced in association with a \(Z \rightarrow \ell^+\ell^-\) decay, contains a significant fraction of gluon jets (of the order of 50%). The ability to separate quark jets from gluon jets is therefore useful for increasing sensitivity to \(ZW\) and \(ZZ\) production.

For a given energy, gluon jets, due to their higher color charge, tend to feature a higher particle multiplicity and be spatially broader in the detector than light-quark (\(u, d,\) and \(s\)) jets. We attempt to quantify the spatial spread of jets using a collection of artificial neural networks (NNs) trained to separate gluon jets from light-flavor quark jets. We refer to the output of the final NN as the jet quark-to-gluon discriminant value (or jet QG value). We calibrate the response of the final NN in MC simulation to match the response in data based on a \(W \rightarrow \ell\nu + 1\) jet event sample. The tagging efficiency and mistag rate associated to a jet are obtained from two NN values are then used as inputs to a third NN.

Each of the NNs is trained using simulated samples of jets matched to either a light-flavor quark or gluon with \(p_T > 20\) GeV/c within \(\Delta R = 0.4\) of the center of the jet and further requiring that no additional partons with transverse momenta exceeding 8 GeV/c are present within \(\Delta R = 0.7\). These jets are selected from a \(Z \rightarrow \mu^+\mu^- + 2\) parton ALPGEN sample, interfaced with PYTHIA showering. Each NN is a feed-forward multilayer perceptron with a hyperbolic-tangent-like response function [20]. The networks are trained on 100 000 quark and gluon jets and tested for biases in overtraining on samples containing 500 000 quark and gluon jets. Gluon-jet distributions are reweighted to match the \(E_T\) and \(\eta\) distributions of the quark jets to remove any discrimination power coming solely from these variables.

For each jet we obtain a list of the calorimeter towers within a cone of \(\Delta R = 0.7\). Each tower has a location coordinate, \((\eta, \phi)\), and energy deposition \(E\) associated with it. We construct a distribution of the distance, \(\Delta R\), between all pairs of towers within the jet and weight each tower pair by its relevance in terms of energy to obtain a distribution that characterizes the spatial spread of the energy within each jet. The weight is given by

\[
\frac{E_i E_j}{0.5[(\Sigma E)^2 - \Sigma E^2]}
\]

where \(E_i\) and \(E_j\) are the energies detected in the two towers of the pair, \(\Sigma E\) is the sum of the energy in all towers within a cone of \(\Delta R = 0.7\) around the jet, and \(\Sigma E^2\) is the sum of the square of the energies of each tower in that same cone. The denominator is chosen to normalize the sum of the weights of all tower pairs to unity. We sample this distribution in 56 intervals (bins) of size \(\Delta R_{\text{bin}} = 0.025\) for \(0.0 < \Delta R < 1.4\), where the contents of the first three bins are empty due to the segmentation of the calorimeter. Typical distributions of the weighted \(\Delta R\) between tower pairs for quark and gluon jets are shown in Fig. 4, using a larger bin size. The outputs of the tower NN for quark and gluon jets using the training and testing samples are shown in Fig. 5.

We follow a similar prescription using tracks within a cone of \(\Delta R = 0.7\) around each jet, using the tracks’ locations in \((\eta, \phi)\) (with respect to the primary vertex) and momenta \(p\) to obtain a distribution of the distance between pairs of tracks (in \(\Delta R\)), with each pair weighted by the momentum carried by that pair, or

\[
\frac{p_i p_j}{0.5[(\Sigma p)^2 - \Sigma p^2]}\]

where \(p_i\) and \(p_j\) are the magnitude of the momenta of the charged particles in the pair, \(\Sigma p\) is the scalar sum of the momenta carried by all charged particles within a cone of \(\Delta R = 0.7\) around the jet, and \(\Sigma p^2\) is the sum of the square of the momenta of each charged particle within that same cone. We require all contributing charged particles to come from the primary vertex and have \(p_T > 0.4\) GeV/c. We split the \(\Delta R\) between track pairs distribution into the same
56 intervals (bins) as used in the tower NN, and the content of each bin is used as an input into the track NN.

Typical distributions of $\Delta R$ between track pairs for quark and gluon jets are shown in Fig. 4. Light-flavor quark jets tend to peak at low $\Delta R$, indicating that they are rather collimated, while gluon jets tend to have a higher mean-valued $\Delta R$ distribution. The bin contents of these $\Delta R$ distributions are used as inputs into NNs that discriminate between quark and gluon jets.

The outputs of the track NN for quark and gluon jets using the training and testing samples are shown in Fig. 5. Higher NN scores indicate jets that are more quark-like. We see good performance in both the tower and track NNs. The cusps in the track NN distribution are associated to jets containing only two charged-particle tracks located inside a cone of $\Delta R = 0.7$, and thus have only one nonzero bin in their distributions of $\Delta R$ between track pairs.

The final NN uses both the tower and track NN values as inputs, along with other jet variables that provide discrimination power between quark jets and gluons: the ratio of $\Sigma E$ associated to towers within a cone of $\Delta R = 0.4$ to $\Sigma E$ associated to towers within a cone of $\Delta R = 0.7$; the ratio of $\Sigma p$ associated to charged particles within a cone of $\Delta R = 0.4$ to $\Sigma p$ associated to charged particles within a cone of $\Delta R = 0.7$; the number of contributing towers with nonzero energy in cones of $\Delta R = 0.4$ and 0.7; the number of contributing charged-particle tracks in cones of $\Delta R = 0.4$ and 0.7; and the jet EM fraction. Additionally, other variables that affect the shape of the $\Delta R$ distributions, independent of whether the jet originates from a quark or gluon, are included: the jet $E_T$, the jet $\eta$, and the number of reconstructed interaction vertices in the event. The output of this final NN is shown in Fig. 6 for light-flavor quark and gluon jets using the training and testing samples. In simulated jets, we see significant separation between quark and gluon jets using this discriminant.

B. Jet QG calibration and performance

The response of the NN quark-to-gluon discriminant may differ between data and MC simulation, especially since uncorrected tower energies are used in the construction of

![FIG. 4. Typical distributions of energy (momentum) content of jets as a function of the $\Delta R$ (a) between pairs of towers and (b) between pairs of tracks in light-flavor quark and gluon jets in simulation.](image)

![FIG. 5. Distributions of the outputs of the NNs processing (a) tower information and (b) track information in light-flavor quark and gluon jets in simulation. Higher NN scores indicate jets that are more quark-like.](image)
the tower NN. Since the signal and most backgrounds are modeled with simulated data, we calibrate the simulation response to match the response in data. We use a control region of independent events with a jet composition similar to that of the final state, $W \rightarrow \ell \nu + 1$ jet events. We then validate the calibration and establish uncertainties on the modeling using control samples of data with features similar to the signal and samples enriched in the dominant backgrounds: $t\bar{t}$ decays in lepton + jets final states and $W \rightarrow \ell \nu + 2$ jet events, respectively.

To form the $W + 1$ jet calibration sample, we choose data collected with the standard high-$E_T$ ($p_T$) central electron (muon) triggers and select events with exactly one central (|$\eta$| < 1.0) electron (muon) with $E_T (p_T) > 20$ GeV/c. To select events consistent with a $W \rightarrow \ell \nu$ decay, we also require a significant missing transverse energy, $E_T > 25$ GeV, and a reconstructed transverse mass [21] consistent with leptonic $W$-boson decays, $m_T > 25$ GeV/c$^2$. To further suppress any contributions from multijet events where a jet mimics the lepton + $E_T$ signature, we require that the $E_T$ is not aligned with any reconstructed jet [$\Delta \phi(E_T, \text{jet}) > 0.2$ radians] and that the $E_T$ significance—a dimensionless quantity comparing the observed $E_T$ against the energy resolution of jets, soft unclustered particles, and the event topology (see Ref. [5])—be larger than four (one) for events with electrons (muons). We also require that the events in this calibration sample have exactly one jet with $E_T > 20$ GeV and |$\eta$| < 2.0.

We consider various processes that contribute to this sample, listed in Table II, and model them using a combination of the PYTHIA, ALPGEN, and MADGRAPH [22] event generators interfaced with PYTHIA for showering. The dominant contribution is $W \rightarrow \ell \nu$ production in association with one jet, which is modeled using ALPGEN. As we are largely concerned with the agreement in shapes between data and simulation, we scale the simulation distributions to match the data. Additionally, we reweight the simulation to match the jet $E_T$ and $\eta$ distributions in data to remove these variables as a possible cause for mismodeling of the jet QG value.

We observe poor modeling of the tower NN values, where the jets in data appear more gluon-like than those in simulated events. The fact that jets in data appear more spatially spread than jets in simulation is consistent with the observed differences in jet energy scales for data and simulation, described in Sec. IV: the fraction of the jet energy contained within a cone of $\Delta R = 0.4$ is higher in simulated gluon jets than in gluon jets in data. We correct for these discrepancies by applying a linear shift to the tower NN values observed in simulation in order to match with data using the $W + 1$ jet sample. We apply different linear shifts for jets in the central and plug calorimeters, and for jets in events with different levels of pileup. We apply further corrections to the response of the final NN to more accurately match the correlations of these calibrated tower-NN values with other jet quantities: the number of towers in the jets and the ratio of $\Sigma E$ in a cone of $\Delta R = 0.4$ to $\Sigma E$ in a cone of $\Delta R = 0.7$. The modeling is more accurate in the track NN than in the tower NN, though we still introduce a similar linear shift in simulated track NN values to more accurately match data. The calibrated variables are input directly into the final NN, without retraining the network.

We further validate the response of the jet QG value by comparing data and MC simulation in a $W \rightarrow \ell \nu + 2$ jets event sample and in an event sample dominated by $t\bar{t}$ production where two quark jets originate from the hadronic decay of a $W$ boson. Table III summarizes the requirements used to select these two samples: the $W + 2$ jet sample is similar to the previously described $W + 1$ jet sample, except for modified jet selections to
match those used in the signal region of the $ZV \rightarrow \ell\ell jj$ search. The $t\bar{t}$ selection eschews the $E_T$ significance and $m_T(W)$ requirements, used to reduce multijet backgrounds, in favor of requirements on the minimum scalar sum of transverse quantities ($jets'$ $E_T$, $6ET$, and the charged lepton's $p_T$) in the event, which is effective in removing both multijet and $W + jets$ backgrounds. Because we are interested in selecting the two jets in the $t\bar{t}$ candidate events that come from the decay of a $W$ boson, as opposed to the $b$ jets produced in the $t \rightarrow Wb$ decays, we make use of the jet-bness tagger [23]. This multivariate $b$-jet identification algorithm exploits properties of individual charged-particle tracks within a jet, looking at properties characteristic of charged particles originating from $B$-hadron decays. The final score, the output of a NN discriminant that ranges between $-1$ and $1$, is called the jet bness, where higher scores identify jets that are more likely to originate from $B$-hadron decays. We classify the two jets with the highest bness scores in the event as the two $b$ jets, and the remaining two jets as those resulting from a $W \rightarrow q\bar{q}'$ decay.

Because we are looking for jet QG shape differences between data and simulation that induce acceptance uncertainties when a requirement on the jet QG value is applied, we scale the number of $W + jet$ events in simulation to match the yield observed in the $W + 2$ jets data. The number of events in each sample is shown in Table II. The distributions of the maximum and minimum QG values of the two jets are shown in Figs. 7 and 8. We see fairly good modeling in the $t\bar{t}$ sample, but poorer modeling in the $W + 2$ jet sample where, even after calibrations, the jets in simulation appear more gluon-like than the jets in data. We account for this remaining discrepancy between data and simulation below.

We enhance the sensitivity to the signal when forming a light-flavor-tagged channel where the minimum jet QG value is greater than 0.0. We determine a probability for a quark jet to meet this requirement (efficiency), and for a gluon jet to be misidentified as a quark jet (mistag rate), with the $t\bar{t}$ and $W + 2$ jet samples. The efficiency measured in data, $e_D$, is a function of the QG requirement, $q$, and may be expressed as
We can write a similar expression for the mistag rate in data; and where $\sigma_\text{MF}_q$ is the uncertainty on the mistag rate, which may be expressed in an analogous fashion. The uncertainty includes a statistical uncertainty on the data, uncertainties on the mistag rate and efficiency, and uncertainties on the relative difference in the contributions from the simulation. We take the uncertainties on the normalizations of the $t\bar{t}$, single top quark, diboson, $V +$ jets, and $V + b$ jets to be 6%, 10%, 6%, 10%, and 40%, respectively, based on the uncertainties in their production cross sections.

We measure the efficiency in the $t\bar{t}$ sample, where the fraction of gluon jets is small, and measure the mistag rate in the $W + 2$ jets sample, where the gluon fraction is much larger and similar to the fraction in the $Z + 2$ jets signal region. The efficiency, mistag rate, and their uncertainties are determined using an iterative procedure. We first calculate the mistag rate in data assuming that the efficiency in data equals the efficiency in simulation. We then calculate the efficiency in data assuming that value for the mistag rate and proceed to update the mistag rate assuming the new value for the efficiency from data. We observe rapid convergence on robust values for the efficiency and mistag rate. Table IV shows the efficiency and mistag rate for the given requirement of minimum QG $> 0.0$, measured in both data and MC simulation. The simulation underestimates the rate.

![Distribution of (a) the maximum and (b) minimum jet QG values of the two jets in the $t\bar{t}$ sample. The distinction between quark and b jets refers to the lower two helicity jets: events where both jets are matched to non-$b$ quark jets are labeled quark jets, while if one of the jets is matched to a $b$ jet, the event is labeled $b$ jets.](image)

$e_D(q) = \frac{e_{\text{raw}}(q) - s_m(q)m_{\text{MC}}(q)f_g}{1 - f_g}$, \hspace{1cm} (6)

where $e_{\text{raw}}$ is the fraction of data events passing the QG requirement; $m_{\text{MC}}$ is the mistag rate for gluons, as measured in simulation; $s_m$ is a scale factor on the mistag rate in simulated jets to match the mistag rate measured in data; and $f_g$ is the fraction of gluon jets in the sample.

We write a similar expression for the mistag rate from $m_D(q) = \frac{m_{\text{raw}}(q) - s_e(q)e_{\text{MC}}(q)f_q}{1 - f_q}$, \hspace{1cm} (7)

where $m_{\text{raw}}$ is the fraction of data events meeting the QG requirements; $e_{\text{MC}}$ is the efficiency for quarks to pass the requirement, as measured in simulation; $s_e$ is a scale factor on the efficiency in simulated jets to match the mistag rate measured in data; and $f_q$ is the fraction of quark jets in the sample. Squared uncertainties on these quantities may be expressed as

$\sigma^2_q = \frac{1}{(1-f_g)^2} \left[ \frac{e_{\text{raw}}(1-e_{\text{raw}})}{N_D} + (\sigma_m f_g)^2 \right]$

$+ \sum \frac{\sigma^2_X}{X(1-f_g)^2}$

$\times [(e + s_m m_{\text{MC}}) (f_g - f_q^X) + f_q^X (e_{\text{MC}} - e_X)]^2$. \hspace{1cm} (8)

<table>
<thead>
<tr>
<th>MC</th>
<th>Data</th>
<th>MC revised jet QG requirement ($-1\sigma$, nom., $+1\sigma$)</th>
</tr>
</thead>
<tbody>
<tr>
<td>Efficiency</td>
<td>0.241</td>
<td>0.295 ± 0.034</td>
</tr>
<tr>
<td>Mistag rate</td>
<td>0.088</td>
<td>0.087 ± 0.027</td>
</tr>
</tbody>
</table>
for quark jets to meet the jet QG requirement, while correctly predicting the observed mistag rate.

We implement a correction to the MC simulation by varying the requirement on the minimum QG value in order to reproduce the efficiency and mistag rate observed in data. The uncertainties on these quantities are also obtained by varying the jet QG requirement. The alternate thresholds used for simulated quark and gluon jets are listed in Table IV.

VI. SIGNAL EXTRACTION AND RESULTS

We extract the number of signal events using a binned $\chi^2$-minimization fit to data, using the techniques described in Ref. [24]. We create histogram templates for both signal and background samples. The templates, along with the uncertainties we assign to their normalization in the fit procedure, are listed below.

1. ZV signal: We allow the normalization of the signal template to float unconstrained in the fit. We assume that each signal process contributes proportionally to its predicted SM cross section: 3.6 pb for ZW and 1.5 pb for ZZ [1].

2. Z + jets: This is the largest background. We allow its normalization to float in the fit, unconstrained.

3. Z + b jets: We constrain the normalization of this significant background to be within ±40% of its nominal value.

4. $t\bar{t}$: We use a production cross section of $\sigma_{t\bar{t}} = 7.5$ pb, and assign an uncertainty of 6.5% to the normalization of this template, based on the theoretical cross-section uncertainty [25].

5. Misidentified leptons: We use the method described in Sec. III to construct templates for the contribution from jets mimicking one or two leptons. We assign an uncertainty of 50% to the misidentification rate, based on studies using different trigger thresholds for the jet data used to obtain these rates.

We perform a simultaneous fit to data using independent templates for each of three channels. For events passing the basic signal selection requirements described in Sec. III, we first construct a heavy-flavor tag (HF-tag) channel composed of events passing a minimum jet $b$ness requirement (jet $b$ness > 0), using the jet-$b$ness tagger [23]. For events failing this requirement, we then select events passing the minimum jet-QG value requirement described in Sec. V to form a light-flavor tag (LF-tag) channel. Events failing this requirement are then placed in the third, untagged channel, which has a lower signal fraction than the two tagged channels, but still includes a significant amount of signal due to the tight tagging requirements. The predicted and observed numbers of events are shown in Table V.

Additional systematic uncertainties on both the normalization and shapes of the templates used in the fit are also considered. We estimate shape and normalization uncertainties due to mismodeling between data and MC simulation in the jet energy scale (as described in Sec. IV C) and the jet energy resolution, the modeling of the tagging variables, and the lepton energy scale and resolution. Additional shape uncertainties on the Z + jets backgrounds are considered by increasing and decreasing the renormalization and factorization scale, $Q^2$, from the default value in the simulation of $m_Z^2 + p_T^{Z,Z}$. We also consider the effect on the shape of the dijet invariant mass when increasing or decreasing initial- and final-state QCD radiation (ISR/FSR) in the ZV signal model. These systematic uncertainties, along with the normalization constraints described above, are treated as nuisance parameters in the fit, and are included in the $\chi^2$-minimization procedure [24]. They are summarized in Table VI.

<table>
<thead>
<tr>
<th>Z + jets</th>
<th>8700 ± 1100</th>
<th>93 ± 14</th>
<th>1520 ± 310</th>
<th>7100 ± 970</th>
</tr>
</thead>
<tbody>
<tr>
<td>Z + b jets</td>
<td>710 ± 300</td>
<td>111 ± 48</td>
<td>55 ± 26</td>
<td>550 ± 230</td>
</tr>
<tr>
<td>$t\bar{t}$</td>
<td>9.2 ± 0.9</td>
<td>3.3 ± 0.4</td>
<td>0.7 ± 0.1</td>
<td>5.1 ± 0.6</td>
</tr>
<tr>
<td>Misidentified leptons</td>
<td>330 ± 170</td>
<td>4.8 ± 2.4</td>
<td>41 ± 20</td>
<td>280 ± 140</td>
</tr>
<tr>
<td>Predicted background</td>
<td>9700 ± 1200</td>
<td>212 ± 55</td>
<td>1620 ± 330</td>
<td>7900 ± 1100</td>
</tr>
<tr>
<td>ZW + ZZ</td>
<td>313 ± 29</td>
<td>12.8 ± 1.6</td>
<td>89 ± 12</td>
<td>212 ± 22</td>
</tr>
<tr>
<td>$N_{\text{signal}}/N_{\text{background}}$</td>
<td>3.2%</td>
<td>6.0%</td>
<td>5.2%</td>
<td>2.6%</td>
</tr>
<tr>
<td>Total predicted events</td>
<td>10000 ± 1300</td>
<td>225 ± 55</td>
<td>1710 ± 330</td>
<td>8100 ± 1100</td>
</tr>
<tr>
<td>Data events</td>
<td>9 846</td>
<td>172</td>
<td>1 724</td>
<td>7 950</td>
</tr>
</tbody>
</table>
Figure 9 shows the dijet mass distributions in data with the fit results overlaid. Table VII shows the number of events of each class determined by the fit. We fit for approximately 50% of the expected signal normalization and observe good agreement between data and simulation in the final fit for each of the three channels, with a total $\chi^2$/d.o.f. = 59.8/55.
We do not see significant evidence for $ZW + ZZ$ production in this decay channel. Hence, we set upper limits on the production cross section using likelihood-ratio ordering [26], where we analyze the distribution of observed cross sections in pseudoexperiments generated with a variety of scale factors on the input-signal cross section. When generating pseudoexperiments, we consider additional systematic uncertainties that affect the acceptance, assigning a 2% uncertainty on the signal template from limited knowledge of the parton distribution functions, and 2.5% and 6% uncertainties due to the uncertainties on the lepton-scale factor and integrated luminosity, respectively. The set of input cross sections range from 0.0 to 2.9 times the expected cross section, with a step size of 0.1.

Figure 10 shows the resulting confidence band. Using the $1\sigma$ bands, we determine $\sigma(p\bar{p} \rightarrow ZW + ZZ) = 2.5^{+2.0}_{-1.0}$ pb, compared to the standard model prediction of $\sigma_{SM} = 5.1$ pb. We do not exclude the no-signal hypothesis, and establish a limit of $\sigma_{ZW + ZZ} < 6.1$ pb ($1.25 \times \sigma_{SM}$) at the 95% C.L.

![Image](103x176 to 190x189)

**FIG. 10** (color online). Confidence bands showing the expected range of observed cross sections as a function of the true cross section, with 68% C.L. (black dashed region) and 95% C.L. (solid gray region).

### VII. SUMMARY

In this paper we described a search for $ZW$- and $ZZ$-boson pair production from a final state with two charged, high-transverse-momentum electrons or muons and two hadronic jets. We increased the sensitivity by tagging events with jets likely originating from heavy- and light-flavor quarks and classifying them in separate analysis channels using neural-network-based taggers. These taggers benefited from the large sample of events containing top quarks collected by CDF, allowing a data-driven estimate of the efficiency and mistag rates for jets passing tagging requirements. We also improved the modeling of the Monte Carlo simulations, especially those that describe the $Z+$ jets background, by deriving and incorporating improved energy corrections for simulated jets to more accurately reproduce the phenomenology of jets originating from quarks and gluons in the data.

Using the full CDF Run II proton-antiproton collisions data set, which corresponds to an integrated luminosity of 8.9 fb$^{-1}$, we fit for the normalization of $ZW$, $ZZ \rightarrow \ell^{+}\ell^{-} + q\bar{q}$ events using the dijet invariant mass distribution. We incorporated many of the systematic uncertainties associated with the modeling of signal and background processes as nuisance parameters in the dijet mass fit. We measured a cross section of $\sigma_{ZW + ZZ} = 2.5^{+2.0}_{-1.0}$ pb, which is nonzero at the $1.75\sigma$ level of significance. We also obtained a limit on the cross section of $\sigma_{ZW + ZZ} < 6.1$ pb at the 95% C.L.

### ACKNOWLEDGMENTS

We thank the Fermilab staff and the technical staffs of the participating institutions for their vital contributions. This work was supported by the U.S. Department of Energy and National Science Foundation; the Italian Istituto Nazionale di Fisica Nucleare; the Ministry of Education, Culture, Sports, Science and Technology of Japan; the National Sciences and Engineering Research Council of Canada; the National Science Council of the Republic of China; the Swiss National Science Foundation; the A.P. Sloan Foundation; the Bundesministerium für Bildung und...
Forschung, Germany; the Korean World Class University Program, the National Research Foundation of Korea; the Science and Technology Facilities Council and the Royal Society, United Kingdom; the Russian Foundation for Basic Research; the Ministerio de Ciencia e Innovación, and Programa Consolider-Ingenio 2010, Spain; the Slovak R&D Agency; the Academy of Finland; the Australian Research Council (ARC); and the EU community Marie Curie Fellowship Contract No. 302103.

[11] The coordinate system used to describe the detector has its origin located at the detector’s center. The proton beam direction is defined as the $+z$ direction. The polar angle, $\theta$, is measured from the origin with respect to the $z$ axis, and $\phi$ is the azimuthal angle. Pseudorapidity, transverse energy, and transverse momentum are defined as $\eta = -\ln \tan(\theta/2)$, $E_T = E \sin \theta$, and $p_T = p \sin \theta$, respectively. The rectangular coordinates $x$ and $y$ point radially outward and vertically upward from the Tevatron ring, respectively.
[14] We define the missing transverse momentum $\vec{E}_T = -\sum E'_i n_i$, where $n_i$ is the unit vector in the azimuthal plane that points from the beamline to the $i$th calorimeter tower. We call the magnitude of this vector the missing transverse energy, $E_T$.
[21] We define the transverse mass as $m_T = \sqrt{\sum_{ij} p_{T,ij}^2 (1 - \cos \Delta \phi_{ij})}$, where $\Delta \phi_{ij}$ is the difference in $\phi$ in the direction of the lepton $\vec{p}_T$ and $\vec{E}_T$ vectors.