**Measurement of the top-quark pair-production cross section in events with two leptons and bottom-quark jets using the full CDF data set**

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Measurement of the top-quark pair-production cross section in events with two leptons and bottom-quark jets using the full CDF data set

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RAPID COMMUNICATIONS

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RAPID COMMUNICATIONS


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We present a measurement of the top-quark pair production cross section in proton-antiproton collisions at $\sqrt{s} = 1.96$ TeV. The data were collected at the Fermilab Tevatron by the CDF II detector and correspond to an integrated luminosity of 8.8 fb$^{-1}$, representing the complete CDF Run II data set. We select events consistent with the production of top-quark pairs by requiring the presence of two reconstructed leptons, an imbalance in the total event transverse momentum, and jets. At least one jet is required to be identified as consistent with the fragmentation of a bottom quark using a

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The transverse energy and momentum are defined as \( E_T (\muon) \) with events that contain at least one isolated electron by the tracking system. In the off-line analysis, we select \( p_T \) measured in the calorimeter and are rejected. The same energy requirements, but without isolation requirement, are imposed. The result supersedes a previous analysis [5] if the identified leptons have the same flavor and dilepton mass consistent with the \( Z \) resonance. We also require the dilepton mass to be larger than 5 GeV/\( c^2 \) to remove events from low-mass dimuon resonances. The resulting sample is referred to as events meeting the dilepton selection. Jets are identified in the laboratory frame using a modified cone algorithm [14], and are defined as having \( E_T > 15 \) GeV and pseudorapidity in the lab frame satisfying \( |\eta| < 2.5 \). Events satisfying the dilepton selection that contain exactly zero or one jet are used as control samples for background estimation. The pretag sample contains events passing the dilepton selection with at least two jets, summed transverse energy over all particles (\( H_T \)) satisfying \( H_T > 200 \) GeV, and whose two leptons are of opposite electric charge. The data sample corresponds to an integrated luminosity of 9.1 fb\(^{-1} \), slightly higher than the signal sample, because the detector quality requirements for \( b \) tagging are not imposed. The pretag sample is used to validate the signal and background models. We measure the \( t\bar{t} \) production cross section using the \( \tag \) sample, which is the subset of pretag events in which at least one of the jets in the event is \( b \) tagged by the SECVTX algorithm [15].

The neutrinos from dilepton top-quark pair decays escape detection, so signal events are expected to produce a large imbalance in the event total transverse-energy (\( E_T \)) [13] compared to other SM dilepton processes that do not contain neutrinos. We require \( E_T > 25 \) GeV to reduce contamination from processes that do not involve neutrinos from vector-boson decays. Events in which \( E_T \) originates from instrumental effects typically feature a small angle between the direction of a lepton or jet and the direction of \( E_T \). If this angle is smaller than 20\(^\circ \), we require \( E_T > 50 \) GeV to reject these backgrounds. To specifically reject events from \( Z/\gamma^* \) production, we require high \( E_T \) significance [5] if the identified leptons have the same flavor and dilepton mass consistent with the \( Z \) resonance. Hadronic decays of tau leptons are not considered. The signal yield is measured as the number of selected events in the data after subtraction of the background expectation from other SM sources, and the cross section is measured by correcting the signal yield for acceptance, efficiency, and luminosity. This analysis uses the full CDF Run II data set collected between March 2002 and September 2011, which corresponds to 8.8 fb\(^{-1} \) of integrated luminosity after data-quality requirements are imposed. The result supersedes a previous analysis [5] by exploiting a threefold increase in data set and improved \( t\bar{t} \) signal-to-background ratio. Improved sample purity is obtained by requiring the presence of jets consistent with the fragmentation of \( b \) quarks (\( b \)-tagged jets) from the top quark decay.

The CDF II detector is a solenoidal spectrometer surrounded by a sampling calorimeter and muon detectors [11]. CDF uses a cylindrical coordinate system with the origin at the center of the detector and the \( +z \) direction defined by the proton beam. \( \theta \) is the polar angle with respect to the \( z \) axis and \( \phi \) is the azimuthal angle. The events were selected for analysis during data taking with an inclusive selection that required the presence of an electron (a muon) with \( E_T > 18 \) GeV (\( p_T > 18 \) GeV/\( c \)). The transverse energy and momentum are defined as \( E_T = E \sin(\theta) \) and \( p_T = p \sin(\theta) \) where \( E \) is the energy measured in the calorimeter and \( p \) is the momentum measured by the tracking system. In the off-line analysis, we select events that contain at least one isolated [12] electron (muon) with \( E_T > 20 \) GeV (\( p_T > 20 \) GeV/\( c \)). We additionally require the presence of a second lepton with the same energy requirements, but without isolation requirements. Events with more than two reconstructed leptons are rejected.
B-hadron decays, but are used to estimate the false-tag rate due to instrumental sources [15].

Selection efficiency for $t\bar{t}$ events is estimated using the PYTHIA Monte Carlo event generator [16] combined with a detailed simulation of the CDF II detector [17]. The $t\bar{t}$ signal is simulated assuming a top-quark mass of 172.5 GeV/c$^2$ and only contains events in which both $W$ bosons produced by the decay of the top quarks subsequently decay into a charged lepton ($e, \mu, \tau$) and a neutrino. Only simulated events with a primary vertex reconstructed within $\pm 60$ cm of the nominal CDF detector center are retained. This requirement has an efficiency of $(97.47 \pm 0.02$ (stat))% of the full CDF luminous region. The total acceptance for the $b$-tagged (pretag) dilepton signal events is $[0.461 \pm 0.003$ (stat)]% ($([0.756 \pm 0.004$ (stat)]%), including the branching fraction to leptons. This acceptance must be corrected to account for the efficiency of the inclusive lepton triggers, which are measured using data samples selected by an independent set of triggering criteria and are in the range 85%–95%. The simulation is also corrected for imperfect modeling of the lepton identification efficiencies by measuring these efficiencies using $Z/\gamma^{\ast} \rightarrow \ell \ell$ events in the data, where one lepton ($\ell = e$ or $\mu$) is fully identified and the other is used for the efficiency measurement. We use these efficiencies to derive multiplicative correction factors (in the range 0.8–1.0) to apply to the simulated efficiencies. We correct for the difference in efficiency for the $b$-tagging algorithm between data and simulated samples by using a multiplicative correction factor, $S_b = 0.96 \pm 0.05$. This correction accounts for the differences between properties of jets in $t\bar{t}$ events and jets in the $b$-tagging calibration sample [15].

The relevant background processes yielding prompt lepton pairs are diboson (WW, WZ, and ZZ) production and $Z/\gamma^{\ast}$ production. Processes in which a photon or hadronic fractionation are identified as a lepton are also considered, such as $W\gamma$ and $W$ boson production in association with multiple jets. The signal sample contamination is predicted to predominantly comprise $Z/\gamma^{\ast}$ and $W +$ jets production processes, so their normalizations are estimated using data samples enriched in these processes. The contamination for the remaining backgrounds is predicted using the same detector simulation and corrections used for the signal.

Diboson production is simulated with PYTHIA, normalized to the production cross sections from the next-to-leading order calculations using MCFM [18] and MSTW2008 [19] parton distribution functions (PDF). The predicted cross sections are $\sigma_{WW} = 11.34 \pm 0.68$ pb, $\sigma_{WZ} = 3.47 \pm 0.21$ pb, and $\sigma_{ZZ} = 3.62 \pm 0.22$ pb [20]. The $Z/\gamma^{\ast} \rightarrow \ell^{+} \ell^{-}$ production is simulated using the alpgen + pythia event generator [21]. The $Z/\gamma^{\ast} \rightarrow e^{+} e^{-}, \mu^{+} \mu^{-}$ samples, which are only selected due to instrumental mismeasurements, are normalized in a data derived process. The process $Z/\gamma^{\ast} \rightarrow \tau^{+} \tau^{-}$ has significant $E_{T}$ from neutrinos and is treated separately; it is normalized to the alpgen production rate and corrected for next-to-leading-order contributions [5]. The $W\gamma$ decays are simulated with the BAURevent generator [22], assuming a leading-order production cross section of $\sigma_{W\gamma} = 32 \pm 3$ pb and correcting for higher-order effects [23]. This process is observed to be relevant in low jet-multiplicity control samples and negligible in the signal sample.

The WW and $Z/\gamma^{\ast} \rightarrow \tau\tau$ jet multiplicity spectra are corrected to account for discrepancies observed between data and simulation in $Z$ boson decays, using jet-multiplicity-dependent correction factors. The corrections are applied to processes in which jets are produced by initial-state radiation, rather than from final-state partons in the hard scattering. The uncertainties on the acceptances of the simulated background processes come from the convolution of the uncertainties due to finite simulation sample size and uncertainties on the jet-multiplicity correction factors, lepton identification, and jet energy scale [24].

The sample contamination from $Z/\gamma^{\ast}$ to $ee$ and $\mu\mu$ decays with instrumental missing energy is estimated in a data sample in which the dilepton mass for all events is consistent with the $Z$ resonance, but all other selection criteria are that of the pretag sample. We subtract the contributions from other processes, and then extrapolate the observed rate outside of the resonance region by using simulated samples, independently for each lepton type and jet multiplicity. The uncertainty on this background contribution is dominated by the limited number of $Z/\gamma^{\ast}$ data events with high $E_{T}$ used to normalize the overall prediction, from the finite size of the $Z/\gamma^{\ast}$ events that meet the selection, and from the uncertainty on the jet-energy scale. The normalization of the $Z/\gamma^{\ast} + u, d, s, g$ event yield in the tagged sample is determined by applying the scalings determined in the pretag sample, and applying false $b$-tag rates [15] as weights to the alpgen + pythia events. We obtain the $Z/\gamma^{\ast} + b, c$ event yield normalization in the tagged sample by requiring events with dilepton mass consistent with the $Z$ resonance, but all other selection criteria as that of the tag sample. After subtracting the estimated $Z/\gamma^{\ast} + u, d, s, g$ component and other backgrounds, the multiplicative heavy-flavor-specific $Z/\gamma^{\ast}$ normalization corrections are found to be $1.8 \pm 0.1$.

We estimate a small contribution to the sample of events with one electron and one muon from $Z/\gamma^{\ast} \rightarrow \mu\mu$ events, in which bremsstrahlung associated with one muon mimics an electron signature. These events are described using the $Z/\gamma^{\ast} \rightarrow \ell^{+} \ell^{-}$ simulation sample. The background from jets misidentified as leptons is estimated by using data events with exactly one identified lepton and additional leptonlike candidates that satisfy less restrictive identification criteria (called the “$W +$ jet” sample, although other sources of misidentified leptons contribute to this sample as well). The probability that a leptonlike candidate is reconstructed as a lepton is parametrized in terms of the candidate’s transverse energy and isolation, and measured in large
quantum chromodynamics dijet dominated samples triggered by the presence of at least one jet of $E_T > 50$ GeV (called the “jet samples”) [5]. Misidentified leptons are modeled by applying these probabilities as weights to the events in the $W + $jet sample with only one high transverse-energy reconstructed lepton and a second electronlike or muonlike candidate. To remove events with two good leptons from this sample, the leptonlike candidate is required to fail at least one lepton identification requirement. The uncertainty on the misidentified-lepton background model is dominated by the differences observed between identification rates determined in jet samples triggered by jets with $E_T$ greater than 20, 50, 70, or 100 GeV.

A common systematic uncertainty for signal and simulated background estimates comes from the uncertainty on the lepton identification correction factors, which is measured to be 2.2%. The 3.3% uncertainty due to the jet-energy scale affects all simulated samples and is estimated by varying the jet-energy corrections by $\pm 1\sigma$ of their systematic uncertainty and measuring the shift in signal and background acceptance. We consider several other sources of systematic uncertainties predominantly affecting the signal efficiency: difference in $t\bar{t}$ modeling by various simulation generators, simulation of initial- and final-state radiation, color reconnection, and PDF [19] uncertainty. These are determined by comparing the uncorrected simulation acceptance of the default $t\bar{t}$ PYTHIA sample to specialized simulation samples. The uncertainty due to each of these sources is estimated to be less than 2.0%. The systematic uncertainty due to the $b$-tagging efficiency correction is 5.0%, dominated by the light-flavor modeling. All simulated backgrounds have uncertainty due to the jet-multiplicity correction factor. Uncorrelated sources of systematic uncertainties affecting individual backgrounds include the 30% systematic uncertainty on the misidentified lepton contamination and individual theoretical uncertainties, ranging from 2% to 10%, on the production cross sections of diboson and $Z/\gamma^{\ast} \rightarrow \tau\tau$ processes. Each of these effects contributes to only a small fraction of the resulting 2.1% (7.1%) background systematic uncertainty for the $b$-tagged (pretag) sample. Table I summarizes the systematic uncertainties that affect the signal acceptance and background model [5].

The expected and observed background events that are $b$-tagged in the 1-jet sample are used as a control sample. The final sample of events with two or more jets passing all candidate selection criteria is given in Table II. The signal purity in the tag sample is 91%, which can be compared to the 73% achieved in the pretag sample. In Fig. 1, we present the jet $E_T$ spectrum for the leading two jets in events with at least two jets, and at least one $b$-tag. The signal yield in the figure is normalized to the measured cross section, and the shape of the distribution is well described by the prediction.

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<th>Source</th>
<th>Uncertainty (%)</th>
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<td>Lepton identification efficiency</td>
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<tr>
<td>Jet energy scale</td>
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<tr>
<td>Simulated event generator</td>
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</tr>
<tr>
<td>Initial- and final-state radiation</td>
<td>1.3</td>
</tr>
<tr>
<td>Color reconnection</td>
<td>1.2</td>
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<tr>
<td>PDF</td>
<td>0.6</td>
</tr>
<tr>
<td>$b$-tagging</td>
<td>5.0</td>
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<td>Background model</td>
<td>2.1</td>
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The expected and observed background events that are $b$-tagged. The total systematic uncertainty is the sum in quadrature of each independent contribution.

<table>
<thead>
<tr>
<th>Source</th>
<th>1 jet</th>
<th>$\geq$ 2 jets ($H_T + OS$)</th>
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<tr>
<td>WW</td>
<td>0.8 ± 0.2</td>
<td>0.6 ± 0.2</td>
</tr>
<tr>
<td>WZ</td>
<td>0.2 ± 0.0</td>
<td>0.1 ± 0.0</td>
</tr>
<tr>
<td>ZZ</td>
<td>0.1 ± 0.0</td>
<td>0.3 ± 0.1</td>
</tr>
<tr>
<td>$Z/\gamma^{\ast} + u, d, s, g$</td>
<td>2.1 ± 0.2</td>
<td>2.8 ± 0.3</td>
</tr>
<tr>
<td>$Z/\gamma^{\ast} + b, c$</td>
<td>1.8 ± 0.2</td>
<td>2.5 ± 0.2</td>
</tr>
<tr>
<td>Other</td>
<td>1.9 ± 0.7</td>
<td>16 ± 5</td>
</tr>
<tr>
<td>Total background</td>
<td>6.9 ± 0.9</td>
<td>22 ± 5</td>
</tr>
<tr>
<td>$t\bar{t}$ ($\sigma = 7.09$ pb)</td>
<td>20.2 ± 1.4</td>
<td>224 ± 15</td>
</tr>
<tr>
<td>Total SM expectation</td>
<td>27.1 ± 2.2</td>
<td>246 ± 20</td>
</tr>
<tr>
<td>Observed</td>
<td>29</td>
<td>246</td>
</tr>
</tbody>
</table>
Data corresponds to an integrated luminosity 3 times greater than that of the previous publication [5], producing a result with a statistical uncertainty that is smaller than the systematic uncertainty. The use of b-jet identification further improves the signal purity from 73% to 91%, and the total uncertainty of the measurement has been improved from 1.04 pb in the previous CDF publication to the current value of 0.84 pb.

In conclusion, we have measured the production cross section of top-quark pairs at the Tevatron, using the full CDF Run II data set. This measurement offers a robust addition to global combined measurements of the top-quark production cross section, which can then be used as constraints to theoretical calculations and limits on non-SM contributions in the top-quark sector.

We thank the Fermilab staff and the technical staffs of the participating institutions for their vital contributions. This work was supported by the U.S. Department of Energy and National Science Foundation; the Italian Istituto Nazionale di Fisica Nucleare; the Ministry of Education, Culture, Sports, Science and Technology of Japan; the Natural Sciences and Engineering Research Council of Canada; the National Science Council of the Republic of China; the Swiss National Science Foundation; the A.P. Sloan Foundation; the Bundesministerium für Bildung und Forschung, Germany; the Korean World Class University Program, the National Research Foundation of Korea; the Science and Technology Facilities Council and the Royal Society, UK; the Russian Foundation for Basic Research; the Ministerio de Ciencia e Innovación, and Programa Consolider-Ingenio 2010, Spain; the Slovak R&D Agency; the Academy of Finland; the Australian Research Council (ARC); and the EU community Marie Curie Fellowship Contract No. 302103.

FIG. 1 (color online). Distribution of jet $E_T$ values for the two jets with largest $E_T$ in each event (black points) for the signal sample. The histogram represents the sum of the signal and background estimates, where the signal is normalized to the observed cross section. The hatched area is the total uncertainty on the sum of the signal and background predictions. The lower panel shows the observed data yield divided by the predicted yield.

For the $b$-tagged $t\bar{t}$ dilepton sample, we measure a cross section of $\sigma_{t\bar{t}} = 7.09 \pm 0.49 \text{(stat)} \pm 0.52 \text{(syst)} \pm 0.43 \text{(lumi)}$ pb $= 7.09 \pm 0.84$ pb with the 246 signal candidate events. The systematic uncertainty is the convolution of the acceptance and the background uncertainties shown in Table I. The 6% luminosity uncertainty is kept separate [25]. The results presented here are consistent with the best recent predictions from next-to-leading order theoretical calculations [10], and with previous D0 and CDF publications [4,5]. The current data sample corresponds to an integrated luminosity 3 times greater than that of the previous publication [5], producing a result with a statistical uncertainty that is smaller than the systematic uncertainty. The use of b-jet identification further improves the signal purity from 73% to 91%, and the total uncertainty of the measurement has been improved from 1.04 pb in the previous CDF publication to the current value of 0.84 pb.

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[13] The missing transverse energy $E_T^m$ is defined to be $-\sum_i E_{T_i} \hat{n}_i$ where $i$ is the calorimeter tower number.
with $|\eta| < 3.6$, $\hat{n}_i$ is a unit vector perpendicular to the beam axis and pointing at the $i$th calorimeter tower and $E_i$ is the transverse energy of the $i$th tower.


[20]  The diboson cross sections when $75 < m_{\ell^+\ell^-} < 105 \text{ GeV}/c^2$ are $\sigma_{WZ} = 3.22 \pm 0.22 \text{ pb}$ and $\sigma_{ZZ} = 1.20 \pm 0.06 \text{ pb}$. The values are calculated using MCFM [18]. The WZ and ZZ cross sections given in the main text are larger, because in the corresponding samples the dilepton mass phase space is restricted only to $m_{\ell^+\ell^-} > 2 \text{ GeV}/c^2$. This allows for the appearance of off-shell $Z$ bosons in the samples, which contribute negligibly to the top dilepton background but can be an important component of other CDF analyses, with which these samples are shared.


