Synthesis of Lumped Transmission-Line Analogs

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Synthesis of Lumped Transmission-Line Analogs
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Abstract—Transmission lines and their lumped approximating networks have long been incorporated into radio-frequency power amplifiers to improve efficiency and shape circuit waveforms and are beginning to perform a similar role in high-frequency switched-mode power electronics. Though lumped line-simulating networks are often preferred to their distributed exemplars for reasons of design flexibility and manufacturability, the impedance peaks and nulls of such lumped networks must be aligned in a precise, harmonic manner to minimize loss and symmetrize converter waveforms. This paper addresses the issue of harmonic frequency alignment in line-simulating networks, presenting new analytic results for predicting the impedance-minimum and impedance-maximum frequencies of networks in a ladder form. Two means of correcting for the observed harmonic misalignment in practical structures will be presented, corroborated by measurements of laminar structures built into the thickness of printed-circuit boards. These structures comprise inductances and capacitances whose dimensions are largely decoupled, such that the simulated line can be accurately analyzed and designed on a lumped basis. The presented techniques will be placed within a power-electronics setting by a representative application incorporating a lumped, line-simulating network.

Index Terms—Distributed system, lumped model, transmission line.

I. INTRODUCTION

The application of transmission lines to power electronics is a promising technique for miniaturizing DC–DC and DC–AC converters, and is considered here and in a complementary paper (see [1]) at the component and converter scale. This paper will focus on component-level techniques, presenting methods for modelling and constructing networks that mimic transmission lines. Such transmission-line analogs can introduce delay into switching cells in a manner that conveniently reduces the total amount of inductance or capacitance required to realize an energy-processing function. Such a decrease in “bulk” energy storage offers unique implementation advantages, such as compatibility with available laminar construction techniques, and a shift to low-loss, air-core magetics. Moreover, whereas reduction of high-frequency parasitics is a major preoccupation of circuit and component design, the delay networks considered in this paper incorporate parasitics.

Tolerance for these unavoidable component resonances encourages faster switching and may further decrease a converter’s energy-storage requirements.

Section II will present, as motivation for this work, the waveform-shaping and symmetrizing properties of transmission lines that are useful in tuned power circuits. Section III will consider the boundary between lumped and distributed systems for line-simulating networks in ladder form, with particular care given to the frequencies at which such networks’ driving-point impedance is at a local minimum or maximum. Measurements of iterated networks—compact, laminar structures embedded in the thickness of printed-circuit boards—will demonstrate the predicted deviation of impedance-extrema frequencies from harmonic coincidence. Sections IV and V present two techniques which compensate for these deviations and achieve more accurate harmonic alignment of impedance peaks and nulls. The analytic and synthetic techniques of Sections IV and V are again verified by measurements of high-order LC structures constructed into the thickness of printed circuit boards. Section VI previews the application of transmission-line analogs to high-frequency inverters, and orients the results within the context of power-circuit miniaturization.

II. BACKGROUND

Open- and short-circuited quarter-wave transmission lines have reactive driving-point impedances with harmonically related maxima and minima; see Fig. 1(a), and [2, Ch. 2]. These aligned peaks and nulls—whether developed by a distributed transmission line or one of its lumped analogs—have long been utilized to shape waveforms and improve efficiency in radio-frequency power amplifiers (e.g., [3, Ch. 14] and [4]–[8]). Similar applications to the design of high-frequency switched-mode power electronics are gradually emerging [9]–[14], and can be regarded as a high-order extension of single-resonant techniques [11], [12] for reducing passive-component bulk.

The waveform-shaping behavior of a transmission-line resonator terminated in a short circuit is clarified in Fig. 1. With a length $\ell$ and some distributed inductance ($L'$) and capacitance ($C'$) per unit length, the input impedance $Z_{\text{in}}$ of a lossless line is a transcendental function with an infinite number of $j$-axis poles and zeros [2]. The zeros of $Z_{\text{in}}$ lie at $s = j\omega_p$, where

$$\omega_p = \frac{\nu \pi}{2\sqrt{LC'}}$$

for $\nu = 0, 2, 4, \ldots$ and

$$Z_{\text{in}} = j\sqrt{\frac{L'}{C'}} \tan(\omega_p\sqrt{LC'}).$$

(1)

Impedance maxima, likewise, are located at odd multiples ($\nu = 1, 3, 5, \ldots$) of the principal quarter-wave resonance $\omega_1$. Because the line will not sustain applied voltages at even harmonics of...
the voltage waveform at the input port contains only odd multiples of the fundamental frequency and is half-wave symmetric for excitation periodic in \( T = \frac{2\pi}{\omega_1} \), as depicted in Fig. 1(b). At odd multiples of \( \omega_1 \) the line does not draw significant current. The terminal current has only even harmonics and is half-wave repeating [cf. Fig. 1(c), where the line’s input current is periodic with one-half the period of the source].

Note that the \( V = I \) symmetry relations are also valid for a half-period of effort by the source. With reference to Fig. 1(b), consider the case of a line excited by a switch which closes from 0 to \( \pi \) radians of the fundamental period \( T \). When the switch imposes a voltage waveform during the first half of the cycle, the transmission line becomes energized so as to impose a half-wave symmetric voltage at the input terminal during the second half cycle. The line stores the voltage waveform in a travelling wave along its length, which returns delayed by one-half fundamental period and inverted, because of the power-reflection condition at the short-circuit termination. The applied current wave also returns, delayed \( \pi/\omega_1 \) seconds but not inverted, so that the line attempts to do the same work on the input network that was done on the line in the first half of the cycle.

Particular care must be given to the harmonic coincidence of the impedance peaks and nulls of line-simulating networks to preserve the waveform-symmetrizing properties of their distributed exemplars. As is discussed in Section VI, such symmetry can be leveraged to reduce switch stresses, improve converter efficiency, and—significantly for miniaturization and control bandwidth—eliminate bulky blocking elements in converter topologies. The next three sections present methods for achieving such harmonic alignment, and clarify the boundary between lumped and distributed systems for a class of \( LC \)-ladder networks of the so-called Cauer form.\(^2\)

### III. Iterated Networks

A compact approximation of a transmission delay can be constructed by cascading \( L \)-sections, as exemplified by the three-, four-, and fivefold iterates of Fig. 2. In the case of such iterated, artificial lines terminated in their characteristic impedance, published bounds are available to guide a designer in selecting the order of an approximating network \([15], [16, Ch. 5]\). A designer can choose how many iterated \( L \)-sections, and of what electrical length, need to be concatenated for some tolerable impedance mismatch with a distributed line. No quantitative guidelines of general applicability, however, appear to have been published for unmatched lines, i.e., for approximating the frequencies of successive \( \lambda/4 \)-wave resonances of artificial lines when they are terminated with an open or short circuit.\(^3\) This unmatched case

\(^1\)The line is a quarter-wave transformer of the short-circuit termination at \( \omega_1 \), with a large impedance at that frequency and all its odd multiples. Recall that, for every additional \( \lambda/4 \) length of line, the driving-point impedance is inverted. In the lossless case, the line transforms the termination from a short --- open --- short --- ... for successive zero --- pole --- zero --- ... at each of which, successively, the line is electrically one quarter-wavelength longer.

\(^2\)Though cascades of resonators, notably, can also simulate line impedance, we focus upon Cauer realizations because of their relative compactness beyond the fourth order \([11], [12]\).

\(^3\)The advice that the electrical length \( \ell \) of component networks be shorter than a some fraction of the smallest signal wavelength of interest is often repeated \([16]\), where “some fraction” is usually set somewhat arbitrarily as the upper bound \( \ell < \lambda/10 \).
is particularly important for the filtering and symmetrizing functions of transmission lines as applied to power-electronics circuits.

To understand how minima and maxima in \( Z_{\text{in}} \) shift for a given degree of discretization in an iterated network with a short-circuit termination, first consider the impedances of the normalized networks of Fig. 2. As shown in the topmost network, all inductors have a value of 1 H, and all capacitors a value of 1 F. The input impedances \( Z_{\text{in}} \) for the normalized networks are a ratio of a numerator \( n_k \) (an odd polynomial in \( s \)) and a denominator \( d_k \) (an even polynomial in \( s \)) with subscripts \( k \) equal to the number of inductors in the network.

As is evident in Table I, the numerator and denominator coefficients of \( n_k \) and \( d_k \) are diagonal sequences from Pascal’s triangle (cf. the left-justified triangle on the right of Table I, in which the diagonals for the polynomials in \( k = 3, 4, \) and 5 are shown in boxes). These diagonal sequences are indexed by increasing \( m \), beginning with \( m = 1 \) for the zeroth-order polynomial, 1, topmost in the triangle. When expressed in the staggered-order form of \( n_k \) and \( d_k \), these polynomials are known as the Fibonacci polynomials \( F_m \). The roots of the Fibonacci polynomials are derived from hyperbolic functions in [17], though Möbius expressed the roots of collapsed-order versions of the polynomials while deriving a periodicity condition for Möbius transforms [18]. The coefficients from the \( m = 11 \) diagonal in Table I, for instance, correspond, in Möbius’ formulation, to the polynomial

\[
\sigma^5 + 9\sigma^4 + 28\sigma^3 + 35\sigma^2 + 15\sigma + 1 = 0
\]

which has five roots given by

\[
\sigma = -4\cos\left(\frac{k'\pi}{m}\right) = 4\cos\left(\frac{k'\pi}{11}\right) \quad \text{for } k' = 1, 2, \ldots, 5. 
\]

These roots can also be expressed in terms of the roots of unity; i.e., if \( \tau \) denotes any \( m \)th root of 1, then the corresponding root \( \sigma \) is \( \sigma = (1 + \tau^2)/r \), e.g., for the eleventh root of unity with angle \( \angle(180/11) \)

\[
4\cos\left(\frac{\pi}{11}\right)^2 = \left[1 + 4\cos\left(180/11\right)^2\right] \approx 3.6825.
\]

Pole and zero frequencies of \( Z_{\text{in}}(s) \) can be recovered from Möbius’ analytic roots in \( \sigma \) by the substitution \( \sigma = s^2 = -\omega^2 \), from which the peak frequencies \( \omega_p \) are

\[
\omega_p = \sqrt{4\cos^2\left(\frac{k'\pi}{m_p}\right)} = 2\cos\left(\frac{k'\pi}{m_p}\right) \quad \text{for } m_p \text{ odd and } k' = 1, \ldots, \left\lfloor \frac{m_p}{2} \right\rfloor
\]

where \([\cdot]\) is the floor function. Zero frequencies, likewise, are given by

\[
\omega_z = \sqrt{4\cos^2\left(\frac{k'\pi}{m_z}\right)} = 2\cos\left(\frac{k'\pi}{m_z}\right)
\]

for \( m_z \) even and

\[
k' = 1, \ldots, \left\lfloor \frac{m_z}{2} \right\rfloor = \left\lfloor \frac{m_p}{2} \right\rfloor.
\]

In (2) and (3), \( m_p \) is the Pascal-diagonal \( m \)-index for the denominator (pole) polynomial, and \( m_z \) is the \( m \)-index for the numerator (zero) polynomial. Note that \( m_z \) is always one greater than \( m_p \). These results provide analytic expressions for the pole and zero locations of uniform lumped-element transmission lines with short-circuit terminations. The impedance extrema for open-circuit terminations are given by the expressions above, after exchanging pole and zero locations appropriately.

Fig. 3(a) shows peak and null frequencies for increasing numbers of cascaded L-sections, demonstrating the alignment of conjugate impedance poles with the range of \( 2 \cdot \cos(k'\pi/m) \) over \( k' = 1, \ldots, [m/2] \). Critical frequencies crowd beneath the cut-off frequency, departing further from harmonic coincidence with increasing \( m \). Impedance nulls, moreover, are never equidistant between adjacent peak frequencies as in an ideal, lossless transmission line, but are always closer to the adjacent, lower frequency pole.

A short-circuit impedance measurement of a network that approximates 30 iterated sections is shown in Fig. 4(a). The frequencies of local maxima and minima in driving-point impedance are depicted by the \( \times \) and \( \circ \) markers in Fig. 4(b), matching well with the predicted pole/zero locations from (2) and (3) [the dashed line shows the locus of analytical roots from (2)]. Further details of the construction and dimensions of this iterated network are provided in Section IV.

Whereas a transmission line terminated in an open or short is capacitive and inductive over equally broad ranges of frequency, the iterated-network \( Z_{\text{in}} \) is capacitive over an increasingly narrow band following each conjugate pole, so that the phase envelope of Fig. 4(a) has an inductive bias at higher frequencies. This tendency is quantified for the measured impedance of Fig. 4(a) by the modal coupling depicted in Fig. 4(c). As developed in the Appendix, the coupling coefficient \( k \) is a measure of the separation between a modal frequency and the zero introduced by exciting it. The \( k \) for lumped-line modes [dashed line, Fig. 4(c)] can be calculated from the theoretical pole and zero frequencies of (2) and (3), and is lower than the coupling of distributed-line modes (solid line). Measured coupling coefficients (triangle markers) for a network with 30 iterated sections closely follow the expected trend, and quantifies the uneven spacing of pole and zero frequencies observed in Fig. 4(a). Methods of improving both harmonic coincidence [Fig. 4(b)] and coupling [Fig. 4(c)] to approximate the impedance of a distributed line will be considered in Sections IV and V.

\footnote{The impedance levels and critical frequencies of the cascaded sections can be denormalized without affecting the relative frequency relationships among poles and zeros.}

\footnote{\( F_m(x) \) evaluated at \( x = 1 \) yields the \( m \)th Fibonacci number, as defined by the recurrence relation \( F_m \equiv F_{m-2} + F_{m-1} \), where \( m = 3, 4, \ldots \) and \( F_1 = F_2 = 1 \)}

\footnote{Fig. 4(b) shows all frequencies as a deviation \( \Delta \omega \) from an integral multiple of the fundamental resonance \( \omega_1 \), e.g., \( \Delta \omega / 2\omega_1 \) for the first zero, and \( \Delta \omega / 3\omega_1 \) for the second impedance peak.}
IV. ITERATED NETWORK WITH MUTUAL INDUCTANCES

The first method of compensating for the critical-frequency shortfall observed in Section III is to introduce adjacent-section mutual inductance into the iterated networks of Fig. 2. Because the calculation of mutual inductances depends on a precise specification of geometry, this section will begin by detailing the structure through which we now propose to transfer energy by both conduction current and mutual flux. This specification will provide an opportunity to discuss how the magnetically uncoupled measurements of Fig. 4(a) were made. The inductance matrices introduced in this section were calculated using FastHenry [19], a freely available program which extracts inductances and resistances of 3-D conductor geometries on a quasistatic basis. The section will conclude with the impedance measurement of a magnetically coupled network, with pole and zero locations summarized in the manner of Fig. 4(b).

The effect of mutual inductances will be introduced by the structure at the top of Fig. 5 built in a printed circuit board, from which the the impedance data of Fig. 4(a) were collected. The toroidal structure comprises 30 section inductances of one turn that are connected in series and tapped at interior nodes with 29 equal capacitances. The total tap capacitance is 915 pF and the total low-frequency self-inductance of the toroid is 355 nH. The structure has the dimensions shown in the top- and bottom-copper masks (reproduced in Fig. 5), and is built on two-sided 62-mil printed circuit board. Fig. 6 depicts a cross-section through the central axis of the toroid. A mylar film is applied over the tap-capacitor plates,

<table>
<thead>
<tr>
<th>$d_1$</th>
<th>$d_2$</th>
<th>$d_3$</th>
<th>$d_4$</th>
<th>$d_5$</th>
<th>$n_1$</th>
<th>$n_2$</th>
<th>$n_3$</th>
<th>$n_4$</th>
<th>$n_5$</th>
</tr>
</thead>
<tbody>
<tr>
<td>1</td>
<td>1</td>
<td>1</td>
<td>1</td>
<td>1</td>
<td>2</td>
<td>3</td>
<td>4</td>
<td>5</td>
<td>6</td>
</tr>
<tr>
<td>$m = 1$</td>
<td>$m = 2$</td>
<td>$m = 3$</td>
<td>$m = 4$</td>
<td>$m = 5$</td>
<td>$m = 6$</td>
<td>$m = 7$</td>
<td>$m = 8$</td>
<td>$m = 9$</td>
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<td>$1$</td>
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<td>$7$</td>
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<td>$9$</td>
<td>$10$</td>
</tr>
</tbody>
</table>

Conjugate pole and zero frequencies for networks with an increasing number of iterated sections

Fig. 3. Analytical prediction of the pole and zero frequencies for iterated L-section networks.

<table>
<thead>
<tr>
<th>Frequency ($\text{rad/s}$)</th>
<th>$2 \cdot \cos \left( \frac{\pi \cdot n}{m} \right)$</th>
</tr>
</thead>
<tbody>
<tr>
<td>0</td>
<td>0</td>
</tr>
<tr>
<td>0.5</td>
<td>0.5</td>
</tr>
<tr>
<td>1</td>
<td>1</td>
</tr>
<tr>
<td>1.5</td>
<td>1.5</td>
</tr>
<tr>
<td>2</td>
<td>2</td>
</tr>
</tbody>
</table>

order of $Z_{in}$ denominator polynomial

Fig. 4. Tabular arrangement of impedance-function coefficients.
with 1-mil adhesive copper foil forming the common node for the taps (a capacitance of 157.6 pF/in² was measured for this method of construction). The structures were designed for a characteristic impedance of 20 Ω and a λ/4 mode at 13.56 MHz. $Z_0 = 19.7\,\Omega$ was measured, with the first quarter-wave mode at 13.41 MHz.

Though a fully populated inductance matrix would be most accurate for modelling turn-to-turn coupling in a toroid, a suitable approximation can be made by only considering adjacent mutual terms (cf. the couplings in the upper schematic of Fig. 5). Because of the large flux leakage in a nonpermeable printed circuit, the mutual inductances between any two ports (where a port is two adjacent tap terminals) falls off rapidly with distance around the toroid. For the structure at the top of Fig. 5, the single-turn self-inductance is 9.97 nH. The mutual inductances decrease rapidly, from 1.47 nH between adjacent turns, to 0.35 nH, 0.12 nH, moving clockwise or counterclockwise along the toroid. Mutual values eventually become small and negative when the turn-sections have antiparallel axes at opposite ends of the toroid.

The lower network in Fig. 5 demonstrates, by means of T-network transformations of the largest mutual terms, how inductance in tap leads offsets the mutual induction between adjacent meshes and diagonalizes the overall $L$-matrix. In the structure measured in Fig. 4(a), 20×80-mil long inductor-traces in series with the capacitor branches cancel 1.47 nH of off-diagonal inductance to approximate the uncoupled, iterated meshes of Fig. 2. Magnetic coupling between input and output was also reduced in the measured structure by introducing a gap of two turns in the full toroid. The lower tableau in Fig. 5 schematically depicts the inductance matrix for this case, in which the inductances in the upper right and lower left of the $L$ matrix have been eliminated by introducing a gap.

By shortening the length of the tap leads, a designer can offset the subcoincident⁸ alignment of critical frequencies observed in an iterated $LC$ network (cf. the dashed locus in Fig. 4(b)). The effect of this change in tap inductance can be explained by a frequency-dependent cancellation of the capacitance loading the toroid. Fig. 7(a) presents the measured critical frequencies of a toroid without tap extensions—but with the same dimensions and tap capacitances—as the toroid in Section III. Compared to the diagonalized critical frequencies, which are shown in grey for comparison, the impedance peaks of the “undiag-
Fig. 5. For a network with adjacent-section mutual coupling (upper schematic), Δ-Y circuit transformations reflect the mutual terms as negative inductances in the tap paths.

Fig. 6. Cross section of the toroidal structures from Fig. 5, shown with exaggerated thicknesses. The toroidal flux path is indicated between pairs of vias on either side of central axis.

Modal coupling falls below the desired transmission-line locus, and has shifted by less than \( k = 0.03 \) from the values in Fig. 4(c).

Though capacitance offset improves harmonic alignment of higher frequencies, the alignment of the lowest critical frequencies is of the greatest significance for enforcing waveform symmetries. These low-frequency poles and zeros are hardly shifted by adjusting the tap inductances \( L_{C_i} \). As with the diagonalized structure of Section III, moreover, low modal coupling prevents pole and zero frequencies from aligning simultaneously at harmonics of the first quarter-wave resonance. Without harmonic zero alignment, the waveform-symmetrizing properties of transmission lines as applied to power electronics (see Section VI) are compromised.

V. CAUER-SYNTHESIZED NETWORK

From the measurements presented in Figs. 4(b) and Fig. 7, a salient problem of iterated networks the subcoincident align-
ment of zero frequencies when impedance poles are near the desired harmonic frequencies. Considered separately, the poles and zeros of iterated networks have good harmonic alignment for little design effort, less than $\pm 1\%$ over the first four critical frequencies. Incidence is poorer (about $\pm 3\%$ of frequency) for poles and zeros considered jointly, and a designer may need to consider more iterated $LC$ divisions—possibly with a larger overall size—to preserve the filtering and symmetrizing functions of the lumped line.

A means of correcting the subcoincident alignment of both poles and zeros without resorting to higher-order networks is to abandon uniform sections in favor of $LC$-ladder elements designed using Cauer synthesis. Cauer-derived networks preserve the basic form shown in Fig. 2, but each series-path inductor and shunt-path capacitance is specified independently to match the impedance of a lossless line over some desired bandwidth. The mathematics of synthesis is described elsewhere [20], [21], and synthesized $L$ and $C$ values are assumed to be given for purposes of this discussion.

The $LC$ ladder values computed from Cauer synthesis can be realized with lumped inductances and capacitances to form a transmission-line approximating network. Of more interest than this straightforward approach is a technique to realize a desired Cauer network with the same family of laminar toroids used in the measurements of Sections III and IV. To implement an “integrated” Cauer network with this planar construction technique, tap points can be placed around the base toroid so that the turn-inductances between taps approximate Cauer-synthesized values. Synthesized capacitances can now be added around the periphery of the toroid to complete the Cauer-derived network. The inductance matrix describing the complete turn-by-turn magnetic coupling of the toroid is related to a new, condensed inductance matrix by a block-sum procedure depicted in Fig. 8. The new inductance matrix is described as “condensed” because the gathering of turns yields an equivalent network of lower order than the densely tapped structure depicted at the top Fig. 5.

Referring to Fig. 8, the self-inductance of a group of $n_l$ turns in a series equals the appropriate $n_l \times n_l$ block sum of the full, turn-by-turn inductance matrix. The process of condensing the full $L$-matrix by progressive block sums is illustrated by the partitioned matrices at the top right of the figure. Primed values are self and mutual terms already condensed by summing, and the three next turns along the toroid—for this example of ladder development—are to be combined into one section inductance. The condensed self-inductance and mutual terms are

$$L_{ij}^{\prime} = \sum_{i=3}^{5} \sum_{j=3}^{5} L_{ij}$$
The corresponding self-inductances \( L_k \) of the condensed network likewise increase as the block-sums of \( n_k \times n_k \) inductance submatrices

\[
L_k = 10.0, 22.1, 22.1, 22.1, 22.1, 22.1, 22.1, 22.1, 35.8, 48.9, 103.2 \text{ nH},
\]

The discrepancy between the sum of these self inductance (353.5 nH) and the total toroidal self inductance (355 nH measured, 366 nH calculated) is due to adjacent-section mutual entries.

As in the case of adjacent-inductor coupled discussed in Section IV, mutual inductances \( M \) in an immediate off-diagonal add to the self-inductances of the sections they couple, and appear as an impedance \( -sM \) in the tap between coupled sections. These mutual terms are of the order of 10% of the tap-to-tap self-inductances, and must be taken into account for accurate pole-zero placement in the Cauer-derived toroid. Assuming that the condensed \( L \) matrix is diagonalized with tap extensions (as detailed in Section IV) we have only to consider the contribution of neighboring mutuals to any given condensed section \( L'_k = L_{k,k-1} + L_{k,k+1} \) (i.e., the inductances as designed in Fig. 9). The progressive grouping of turns along the toroid represented by (4) and (5), however, cannot explicitly account for the mutual inductance of condensed sections yet to be designed. The iterated method summarized in Fig. 9 circumvents this difficulty by regrouping turns for each section \( k \) once the section \( k+1 \) is designed, refining the estimate of condensed section-inductance \( L'_k \) with a better approximation for the new mutual terms. The initial guess for \( L'_k \) is based upon the forward mutual inductance to a block with an identical number of turns \( n_k \). The Cauer-derived network measured in Fig. 10 was designed using this method.

Fig. 10 shows the poles and zeros of measured impedance for a Cauer-synthesized toroid based upon the 20 \( \Omega \) layout of Fig. 5, with diagonalizing tap inductances and 11 condensed sections. Though poles and zeros are both in the vicinity of harmonic coincidence, their location seems much more uncertain (\( \pm 3\% \)) than in the iterated networks. The alignment of the lowest frequency poles and zeros is notably worse than the in the iterated networks [cf. Figs. 4(b) and 7], where better alignment was achieved with less design effort. Note that overall harmonic alignment (i.e., considering poles and zeros together) is not appreciably worse than in the iterated cases.

Critical-frequency alignment improves dramatically when pole and zero frequencies are compared, not to harmonically aligned values, but to the values expected after the process of turns gathering outlined previously in this Section. Because of the discrete choices for section-inductance available at any given point in the network realization, the continuously increasing \( L \) and \( C \) values (typified by Fig. 9) can at best be approximated. Normalizing measured poles and zeros to the frequencies computed from the approximated Cauer network yields the alignment depicted in Fig. 10(b): \( \pm 1\% \) alignment of poles and zeros over a broad frequency range, with tighter coincidence when the first five or six critical frequencies are considered alone. This result is a strong endorsement for the accuracy of the inductance matrix computed with FastHenry.
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Fig. 9. Cauer-synthesized $L$ and $C$ values for a 11-section ladder designed to match 21 transmission-line nonzero critical frequencies. Flowchart of an iterative method for determining tap locations along a toroid. The $n_k$ are the number of turns in each Cauer-synthesized section with self-inductance $L_{k}$. 

Fig. 10. Poles and zeros of the measured impedance of a Cauer-derived toroid, based upon a 30-turn gapped toroid. In the left plot, the measured critical frequencies are normalized to the desired transmission-line locus. In the right plot, the frequencies are normalized to expectations from the synthesis model. Note that a redesign of the PCB—accommodating more precisely the Cauer-derived inductances and capacitances—would be expected to achieve ±1% frequency precision and accuracy.

VI. APPLICATIONS

Power-electronics applications of the transmission-line analogs explored in this paper are treated in an adjunct paper [1]. To place this work within a useful context, however, consider the schematic of Fig. 11, in which the dashed box around the lumped transmission line replaces the input choke of a Class E inverter. In the Introduction, it was mentioned that transmission-line techniques can reduce the total amount of inductance or capacitance required to realize an energy-processing function, which in turn can have important manufacturing benefits. For the case at hand, the Class E input choke (replaced in Fig. 11) comprises an air-core solenoid, 21 turns of 18-gauge wire, wound on a plastic former with a 26-mm diameter. Significant for manufacturability, the transmission-analog replacement is constructed directly into the thickness of a four-layer, 2 oz. copper PCB. A 59-mil core was selected for the magnetic thickness dimension, with capacitors constructed across outer prepreg layers and an overall outer diameter of 4.4 in. The transmission-line input network is of the iterated type analyzed in Section III, and is constructed with only 207 nH of planar inductance and about 500 pF of interlayer capacitance. Compare these passive values with 8.02-μH inductance for the Class E design. The inverter with transmission line circulates waveforms internally to accomplish its power conversion function, exchanging large-valued blocking components for high-resonant elements. This tradeoff will be explored in the adjunct paper, as well as the reduced stresses attending the natural square-wave switching introduced by the transmission-line impedance (shown for an previously published case [21] in Fig. 11).

Design decisions regarding the model order, characteristic impedance, and losses of a lumped line analog have so far only been considered with respect to the overall efficiency of particular power-electronics designs. In the inverter example considered in [1], for instance, a 28-turn, 20-Ω toroid based upon the design of Fig. 4 was selected because its first three zeros

Note that stray capacitance (5 pF per tap in addition to the parallel-plate value, as layed out in copper) and lead inductance (10-nH lead inductance in series with the first section) were considered when predicting the critical frequencies for Fig. 10(b).
formed an even-harmonic set within 1% of frequency. Waveform shapes in the inverter—with symmetries providing zero-voltage-switching opportunities—were realized by aligning the switching frequency with this set of even-harmonic zeros. Beyond this constraint on zero frequencies, measured converter efficiency (rather than pole location or characteristic impedance per se) was used to discriminate among designs. Systematic selection of line parameters remains an important topic for further study, as does a general treatment of loss for methods of construction.

**VII. CONCLUSION**

This paper has considered critical-frequency alignment (see Table II) of iterated and Cauer-derived transmission-line analogs with planar, air-core magnets. For the iterated network of Sections III and IV, a relation between the critical frequencies of cascaded \( LC \) sections and the roots of Fibonacci polynomials has been presented. This result quantifies the transition between a distributed line and its lumped ladder approximations, and appears to be new in the literature. The Cauer-derived line-simulating network, with nonuniform placement of taps in approximation of the \( L \) and \( C \) values calculated by Cauer synthesis, exhibited critical frequencies matching the lumped-model prediction within \( \pm 1\% \). Though the Cauer-derived network had poorer harmonic coincidence than the iterated networks, the precision of the MQS model allows this shortfall to be overcome by design. Moreover, since the Cauer network realizes a specified driving-point reactance, it alone—at the three options presented—permits the simultaneous harmonic alignment of poles and zeros necessary for the symmetrizing function of transmission networks.

**APPENDIX**

Whenever driving effort excites some mode, whether native to the energy-domain of excitation or not, the coupling coefficient \( k \) represents the extent of energy conversion, and is defined over a cycle as

\[
k^2 = \frac{\text{energy delivered to load}}{\text{total energy delivered}} + \frac{\text{energy stored and recovered from load}}{\text{energy delivered}}.
\]

Consider a simple case of energy storage shown in Fig. 12(a): a mass slides on a frictionless plane, and is tied to a mechanical ground through a spring \( k_3 \). A force with infinite authority—i.e., no source impedance—stretches \( k_3 \) and stores energy in it. All of the energy delivered from the source stretches the spring, and the drive is perfectly coupled to \( k_3 \). For this case of zero source impedance, the coupling coefficient \( k \equiv 1 \).

Mode excitation is rarely as simple as the case pictured in Fig. 12(a). Usually, energy stored in \( k_3 \) requires that some en-
nergy be delivered to the exciting structure, as shown in Fig. 12(b). At frequencies far below resonance, the oscillator is compliance-dominated (i.e., the energy stored in \( m \) is negligible) and \( F \) stores energy in \( k_1 \) by stretching \( k_2 \). If \( k_2 \) is stiff, \( F \) is strongly coupled to \( k_1 \); if \( k_2 \) is compliant, it is stretched considerably, \( k_1 \) is relatively undisturbed, and the coupling is weak. \( k^2 \) in this case can be expressed as the fraction of total energy delivered to \( k_1 \)

\[
k^2 = \frac{\text{energy stored in } k_1}{\text{energy stored in } k_1 + \text{energy stored in } k_2}.
\]

In laboratory practice, the frequency separation between resonance and antiresonance (i.e., between a modal resonance and the zero introduced by exciting it) is a measure of the degree of energy coupling. To understand this useful relationship, consider the impedance of Fig. 13(a) for the network in Fig. 12(c)

\[
Z = j \omega L_0 - \frac{j}{\omega C - \frac{1}{\omega L}} = j \omega \left[ \frac{L}{\omega^2 L C - 1} \right] = j \omega^2 \frac{L}{\omega^2 L C - 1} = \omega_p^2 \frac{L}{\omega^2 L C - 1} - \omega_z^2 \frac{L}{\omega^2 L C - 1}
\]

\( \omega_p \) and \( \omega_z \) from Fig. 13(a) are found by setting numerator and denominator to zero and solving for frequency, from which

\[
\omega_p = \sqrt{\frac{1}{LC}} \quad \text{and} \quad \omega_z = \sqrt{\frac{1}{L_0} \cdot \frac{1}{L \cdot C}} = \sqrt{\frac{1}{LC} + \frac{1}{L_0 C}}.
\]

The ratio \( \omega_z^2/\omega_p^2 \) bears a useful relationship to \( k^2 \), and can be expressed as

\[
\frac{\omega_z^2}{\omega_p^2} = \frac{L + L_0}{L_0} \quad \text{since} \quad \frac{\omega_z^2}{\omega_p^2} - 1 = \frac{L}{L_0}.
\]

By the energy definition of coupling coefficient, and for some applied current \( i \) in the low-frequency, inductance-dominated regime

\[
k^2 = \frac{\frac{1}{2} i L_x^2}{\frac{1}{2} i L_x^2 + \frac{1}{2} i L_0 x^2} = \frac{L}{L_0 + L}
\]

which can in turn be related to the ratio \( \omega_z^2/\omega_p^2 \)

\[
1 - k^2 = \frac{L_0}{L + L_0} \quad \Rightarrow \quad k^2 = 1 - \frac{\omega_z^2}{\omega_p^2}.
\]

Using the gross pole-zero coincidence approximation \( \omega_z + \omega_1 \approx 2\omega_z \)

\[
k^2 = \frac{(\omega_z + \omega_p)(\omega_z - \omega_p)}{\omega_z^2} \approx 2 \cdot \frac{\omega_z - \omega_p}{\omega_z^2}. \tag{7}
\]

While (6) precisely determines \( k \) from measured \( \omega_z \) and \( \omega_p \), (7) affords the insight that pole-zero spacing increases linearly with increasing \( k \), to a better approximation as \( \omega_z - \omega_p \) is small with respect to either critical frequency. We can apply this single-resonant treatment of coupling to cascaded oscillators, as long as we can ignore the impedances of neighboring intertia- or compliance-dominated resonators, excited away from their tuned frequencies. This approximation, and the others of this section, are common practice within the transducer field [22, Section 4.8], and are a convenient means of comparing the alignment of adjacent poles and zeros.

REFERENCES


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