Observation of $\bar{\Lambda}_B[0(s)] J/f_{1}(1285)$ Decays and Measurement of the $f_{1}(1285)$ Mixing Angle

The MIT Faculty has made this article openly available. Please share how this access benefits you. Your story matters.
Observation of $B^0 \to J/\psi f_1(1285)$ Decays and Measurement of the $f_1(1285)$ Mixing Angle

R. Aaij et al. *(LHCb Collaboration)*

(Received 8 October 2013; published 4 March 2014)

Decays of $B^0$ and $B^0$ mesons into $J/\psi \pi^+ \pi^- \pi^+ \pi^-$ final states, produced in $p p$ collisions at the LHC, are investigated using data corresponding to an integrated luminosity of 3 fb$^{-1}$ collected with the LHCb detector. $B^0 \to J/\psi f_1(1285)$ decays are seen for the first time, and the branching fractions are measured. Using these rates, the $f_1(1285)$ mixing angle between strange and nonstrange components of its wave function in the $q\bar{q}$ structure model is determined to be $\pm(24.0^{+3.1+0.6}_{-2.6-0.8})^\circ$. Implications on the possible tetraquark nature of the $f_1(1285)$ are discussed.

DOI: 10.1103/PhysRevLett.112.091802

PACS numbers: 14.40.Be, 13.20.He, 13.25.-k, 14.65.Bt

Light flavorless hadrons, $f$, are not entirely understood as $q\bar{q}$ states. Some states with the same quantum numbers such as the $\eta$ and $\eta'$ exhibit mixing [1]. Others, such as the $f_0(500)$ and the $f_0(980)$, could be mixed $q\bar{q}$ states, or they could be comprised of tetraquarks [2]. In addition, some states, such as the $f_0(1500)$, are discussed as being made solely of gluons [3]. Understanding if the $f$ states are indeed explained by the quark model is crucial to identifying other exotic structures. Previous investigations of $B^0$ and $B^0$ decays (called generically $\bar{B}$) into a $J/\psi$ meson and a $\pi^+ \pi^-\pi^+ \pi^-$ [4,5] or $K^+K^-$ [6,7] pair have revealed the presence of several light flavorless meson resonances including the $f_0(500)$ and the $f_0(980)$. Use of $\bar{B} \to J/\psi f$ decays has been suggested as an excellent way of both measuring mixing angles and discerning if some of the $f$ states are tetraquarks [8,9]. In this Letter the $J/\psi \pi^+ \pi^-\pi^+ \pi^-$ final state is investigated with the aim of seeking additional $f$ states. (Mention of a particular process also implies the use of its charge conjugated decay.)

Data are obtained from 3 fb$^{-1}$ of integrated luminosity collected with the LHCb detector [10] using $p p$ collisions. One third of the data was acquired at a center-of-mass energy of 7 TeV, and the remainder at 8 TeV. The LHCb detector is a single-arm forward spectrometer covering the pseudorapidity range $2 < \eta < 5$, designed for the study of particles containing $b$ or $c$ quarks. The detector includes a high precision tracking system consisting of a silicon-strip vertex detector surrounding the $pp$ interaction region, a large-area silicon-strip detector located upstream of a dipole magnet with a bending power of about 4 Tm, and three stations of silicon-strip detectors and straw drift tubes placed downstream. The combined tracking system provides a momentum measurement with relative uncertainty that varies from 0.4% at 5 GeV to 0.6% at 100 GeV. (We work in units where $c = 1$.) The impact parameter (IP) is defined as the minimum track distance with respect to the primary vertex. For tracks with large transverse momentum, $p_T$, with respect to the proton beam direction, the IP resolution is approximately 20 $\mu$m. Charged hadrons are identified using two ring-imaging Cherenkov (RICH) detectors. Photon, electron, and hadron candidates are identified by a calorimeter system consisting of scintillating-pad and preshower detectors, an electromagnetic calorimeter, and a hadronic calorimeter. Muons are identified by a system composed of alternating layers of iron and multiwire proportional chambers.

The LHCb trigger [11] consists of a hardware stage, based on information from the calorimeter and muon systems, followed by a software stage that applies event reconstruction. Events selected for this analysis are triggered by a candidate $J/\psi \to \mu^+ \mu^-$ decay, required to be consistent with coming from the decay of a $b$ hadron by using either IP requirements or detachment from the associated primary vertex. Simulations are performed using PHYTHIA [12] with the specific tuning given in Ref. [13], and the LHCb detector description based on GEANT4 [14] described in Ref. [15]. Decays of $b$ hadrons are based on EVTGEN [16].

Events are preselected and then are further filtered using a multivariate analyzer based on the boosted decision tree (BDT) technique [17]. In the preselection, all charged track candidates are required to have $p_T > 250$ MeV, while for muon candidates the requirement is $p_T > 550$ MeV. Events must have a $\mu^+ \mu^-$ combination that forms a common vertex with $\chi^2 < 20$, an invariant mass between $-48$ and $+43$ MeV of the $J/\psi$ meson mass, and are constrained to the $J/\psi$ mass. The four pions must have a vector summed $p_T > 1.0$ GeV, form a vertex with $\chi^2 < 50$ for 5 degrees of freedom, and a common vertex with the

* Full author list given at the end of the article.
$J/\psi$ candidate with $\chi^2 < 90$ for 9 degrees of freedom. The angle between the $B$ momentum and the vector from the primary vertex to the $B$ decay vertex is required to be smaller than 2.56°. Particle identification [18] requirements are based on the difference in the logarithm of the likelihood, DLL$(h_1 - h_2)$, to distinguish between the hypotheses $h_1$ and $h_2$. We require DLL$(\pi - \mu) > -10$ and DLL$(\pi - K) > -10$. We also explicitly eliminate candidate $\psi(2S)|X(3872)| \rightarrow J/\psi\pi^+\pi^-$ events by rejecting any candidate where one $J/\psi\pi^+\pi^-$ combination is within 23 MeV of the $\psi(2S)$ or 9 MeV of the $X(3872)$ meson masses. Other resonant contributions such as $\bar{B} \rightarrow \psi(4160)\pi^+\pi^-$ are searched for, but not found.

The BDT uses 12 variables that are chosen to separate signal and background: the minimum DLL$(\pi - \mu)$ of the $\mu^+$ and $\mu^-$; the scalar $p_T$ sum of the four pions, and the vector $p_T$ sum of the four pions; relating to the $B$ candidate: the flight distance, the vertex $\chi^2$, the $p_T$, and the $\chi^2_{IP}$, which is defined as the difference in $\chi^2$ of a given primary vertex reconstructed with and without the considered particle. In addition, considering the $\pi^+\pi^+$ and $\pi^-\pi^-$ as pairs of particles, the minimum $p_T$, and the minimum $\chi^2_{IP}$ of each pair are used. The signal sample used for BDT training is based on simulation, while the background sample uses the sideband 200–250 MeV above the $B^0$ mass peak from 1/3 of the available data. The BDT is then tested on independent samples from the same sources. The BDT selection is optimized by taking the signal, $S$, and background, $B$, events within $\pm20$ MeV of the $B^0$ peak from the preselection and maximizing $S/(S + B)$ by using the signal and background efficiencies provided as a function of BDT.

The $J/\psi\pi^+\pi^-\pi^+\pi^-$ invariant mass distribution is shown in Fig. 1. Multiple combinations are at the 6% level and a single candidate is chosen based on vertex $\chi^2$ and $J/\psi$ mass. We fit the mass distribution using the same signal function shape for both $B^0$ and $\bar{B}^0$ peaks. This shape is a double Crystal Ball function [19] with common means and radiative tail parameters obtained from simulation. The combinatorial background is parameterized with an exponential function. There are 1193 $\pm$ 46 $B^0$ and 839 $\pm$ 39 $\bar{B}^0$ decays. Possible backgrounds caused by particle misidentification, for example $\bar{B}^0 \rightarrow J/\psi\pi^+K^-\pi^+\pi^-$ decays, would appear as signal if the particle identification incorrectly assigns the $K^-$ as a $\pi^-$. In this case the invariant mass is always below the $B^0$ signal region. Evaluating all such backgrounds shows negligible contributions in the signal regions. These and other low-mass backgrounds are described by a Gaussian distribution.

In order to improve the four-pion mass resolution we kinematically fit each candidate with the constraints that the $\mu^+\mu^-$ be at the $J/\psi$ mass and that the $J/\psi\pi^+\pi^-\pi^+\pi^-$ be at the $B$ mass. The four-pion invariant mass distributions for $B^0$ and $\bar{B}^0$ decays within $\pm20$ MeV of the $B$ mass peaks are shown in Fig. 2. The backgrounds, determined from fits to the number of events in the region 40–80 MeV above the $B^0$ mass, are subtracted.

There are clear signals around 1285 MeV in both $B^0$ and $\bar{B}^0$ decays with structures at higher masses. The $J/\psi$ decay angular distribution is used to probe the spin of the recoiling four-pion system. We examine the distribution of the helicity angle $\theta$ of the $\mu^-$ with respect to the $B$ direction in the $J/\psi$ rest frame, after correcting for the angular acceptance using simulation. The resulting distribution is then fit by the sum of shapes $f_1(1285)$.

There is also a large and wider peak near 1450 MeV in the $B^0$ channel. Previously we observed a structure at a mass near 1475 MeV using $B^0 \rightarrow J/\psi\pi^+\pi^-\pi^0$ decays that we attributed to $f_0(1370)$ decay. However, it could equally well be the $f_0(1500)$ meson, an interpretation favored by Ochs [3]. While the $f_0(1500)$ is known to decay into four pions, the structure observed in our data cannot be pure spin 0 because of the significant helicity $\pm1$ component in this mass region. We do not pursue further the composition of the higher mass regions in either $B^0$ or $\bar{B}^0$ decays in this Letter.

We use the measured branching fractions of $B^0 \rightarrow J/\psi\pi^+\pi^-\pi^0$ [4] and $B^0 \rightarrow J/\psi\pi^+\pi^-\pi^\sigma$ [5] for normalizations. The data selection is updated from that used in previous publications to more closely follow the procedure in this analysis. We find signal yields of 22476 $\pm$ 177 $B^0$ events and 16016 $\pm$ 187 $\bar{B}^0$ events within $\pm20$ MeV of the signal.
peaks. The overall efficiencies determined by simulation are $(1.411 \pm 0.015)\%$ and $(1.317 \pm 0.015)\%$, respectively, for $B_s^0$ and $B^0$ decays, where the uncertainty is statistical only. The relative efficiencies for the $J/\psi \pi^+ \pi^- \pi^+ \pi^-$ final states with respect to $J/\psi \pi^+ \pi^-$ are $14.3\%$ and $14.5\%$ for $B_s^0$ and $B^0$ decays, with small statistical uncertainties. We compute the overall branching fraction ratios

$$B(B_s^0 \to J/\psi \pi^+ \pi^- \pi^+ \pi^-)/B(B_s^0 \to J/\psi \pi^+ \pi^-) = 0.371 \pm 0.015 \pm 0.022,$$

$$B(B^0 \to J/\psi \pi^+ \pi^- \pi^+ \pi^-)/B(B^0 \to J/\psi \pi^+ \pi^-) = 0.361 \pm 0.017 \pm 0.021.$$
TABLE I. Fit results for $\bar{B}_s^0 \to J/\psi f_1(1285)$ and $\bar{B}_s^0 \to J/\psi f_1(1285)$ decays.

<table>
<thead>
<tr>
<th>Source</th>
<th>$\bar{B}_s^0$</th>
<th>$\bar{B}_s^0$</th>
<th>$\bar{B}_s^0$</th>
<th>$\bar{B}_s^0$</th>
<th>$\bar{B}_s^0$</th>
<th>$\bar{B}_s^0$</th>
</tr>
</thead>
<tbody>
<tr>
<td></td>
<td>Yield</td>
<td>$-2\Delta \ln L$</td>
<td>Significance ($\sigma$)</td>
<td></td>
<td></td>
<td></td>
</tr>
<tr>
<td>$\bar{B}_s^0$</td>
<td>110.2 ± 15.0</td>
<td>58.1</td>
<td>7.2</td>
<td></td>
<td></td>
<td></td>
</tr>
<tr>
<td>$\bar{B}_s^0$</td>
<td>49.2 ± 11.4</td>
<td>29.5</td>
<td>5.2</td>
<td></td>
<td></td>
<td></td>
</tr>
</tbody>
</table>

$f_1(1285)$ in $b$-hadron decays. As a consistency check, we also perform a simultaneous fit to both $\bar{B}_s^0$ and $B^0$ samples letting the mass and width vary in the fit. We find the mass and width of the $f_1(1285)$ to be $1284.2 \pm 2.2$ MeV and $32.4 \pm 5.8$ MeV, respectively, where the uncertainties are statistical only, consistent with the known values.

To determine the systematic uncertainty in the yields we redo the fits allowing $\pm 1\sigma$ variations of the mass and width values independently. We assign $\pm 2.7\%$ and $\pm 2.0\%$ for the systematic uncertainties on the $\bar{B}_s^0$ and $B^0$ yields, respectively, from this source.

We obtain the branching fraction ratios, using an efficiency of $0.1820 \pm 0.0036\%$, determined by simulation, for the $J/\psi f_1(1285)$ final state as

$$\frac{B[\bar{B}_s^0 \to J/\psi f_1(1285), f_1(1285) \to \pi^+\pi^-\pi^-]}{B[\bar{B}_s^0 \to J/\psi \pi^+\pi^-]} = (3.82 \pm 0.52^{+0.29}_{-0.32})\%,$$

$$\frac{B[\bar{B}_s^0 \to J/\psi f_1(1285), f_1(1285) \to \pi^+\pi^-\pi^-]}{B[B^0 \to J/\psi f_1(1285)]} = (0.32 \pm 0.54 \pm 0.11)\%,$$

$$\frac{B[B^0 \to J/\psi f_1(1285)]}{B[B^0 \to J/\psi f_1(1285)]} = (11.6 \pm 3.1^{+0.7}_{-0.8})\%.$$  

For the latter ratio we use a $\bar{B}_s^0/B^0$ production ratio of $0.259 \pm 0.015$ [20]; this uncertainty is taken as systematic. The other systematic uncertainties are listed in Table II. The shape of the high-mass tail is changed in the case of $\bar{B}_s^0$ decays from a single Gaussian to two relativistic Breit-Wigner shapes corresponding to the mass and width values of the $f_1(1420)$ and the $f_0(1500)$ mesons. For the $B^0$ high mass shape we change from a Gaussian shape to a second order polynomial. The decay model reflects the allowed variation in the fraction of $\rho^0\rho^0$ and $\rho^0\pi^+\pi^-$ decays. The total uncertainties are ascertained by adding the individual components in quadrature separately for the positive and negative values.

Considering the $f_1(1285)$ as a mixed $q\bar{q}$ state, we characterize the mixing with a $2 \times 2$ rotation matrix containing a single parameter, the angle $\phi$, so that the wave functions of the $f_1(1285)$ and its partner, indicated by $f_1^*$, are given by

$$|f_1(1285)\rangle = \cos \phi |n\bar{n}\rangle - \sin \phi |s\bar{s}\rangle,$$

$$|f_1^*(1285)\rangle = \sin \phi |n\bar{n}\rangle + \cos \phi |s\bar{s}\rangle,$$

where $|n\bar{n}\rangle = \frac{1}{\sqrt{2}}(|u\bar{u}\rangle + |d\bar{d}\rangle).$ (1)

The decay widths can be written as [8]

$$\Gamma(\bar{B}^0 \to J/\psi f_1(1285)) = 0.5|A_1|^2|V_{cd}|^2\Phi_0 \cos^2 \phi,$$

$$\Gamma(B^0 \to J/\psi f_1(1285)) = |A_1|^2|V_{cs}|^2\Phi_1 \sin^2 \phi,$$ (2)

where $A_1$ is the tree level amplitude, $V_{cd}$ and $V_{cs}$ are quark mixing matrix elements, and $\Phi_1$ is phase space factors. The amplitude ratio $|A_1|/|A_0|$ is taken as unity [8]. The width ratio is given by

$$\frac{B[\bar{B}^0 \to J/\psi f_1(1285)]}{B[B^0 \to J/\psi f_1(1285)]} = \frac{\tau_0}{\tau_s} \frac{|V_{cd}|^2 \Phi_0 \cos^2 \phi}{|V_{cs}|^2 \Phi_1 \sin^2 \phi},$$ (3)

where $\tau_s$ is the $B_s^0$ lifetime and $\tau_0$ is the $B^0$ lifetime. The angle $\phi$ is then given by

$$\tan^2 \phi = \frac{1}{2} \frac{B[B^0 \to J/\psi f_1(1285)]}{B[\bar{B}^0 \to J/\psi f_1(1285)]} \frac{|V_{cd}|^2 \Phi_0}{|V_{cs}|^2 \Phi_1} \frac{\tau_0}{\tau_s},$$ (4)

where $\tau_s = 1.508$ ps [21], $\tau_0 = 1.519$ ps, $|V_{cd}| = 0.2245$, and $|V_{cs}| = 0.97345$ [1]. We use the lifetime measured in $\bar{B}_s^0 \to J/\psi \phi$ decays as the helicity components are in approximately the same ratio as in $J/\psi f_1(1285)$. No uncertainties are assigned on these

<table>
<thead>
<tr>
<th>Source</th>
<th>$B^0$</th>
<th>$\bar{B}_s^0$</th>
<th>$\bar{B}_s^0$</th>
<th>$\bar{B}_s^0$</th>
</tr>
</thead>
<tbody>
<tr>
<td>Mass and width of $f_1$</td>
<td>2.0</td>
<td>2.0</td>
<td>2.7</td>
<td>2.7</td>
</tr>
<tr>
<td>Shape of high mass</td>
<td>0.6</td>
<td>0</td>
<td>3.7</td>
<td>0</td>
</tr>
<tr>
<td>Efficiency</td>
<td>2.0</td>
<td>2.0</td>
<td>2.0</td>
<td>2.0</td>
</tr>
<tr>
<td>Tracking</td>
<td>2.0</td>
<td>2.0</td>
<td>2.0</td>
<td>2.0</td>
</tr>
<tr>
<td>Simulation statistics</td>
<td>2.0</td>
<td>2.0</td>
<td>2.0</td>
<td>2.0</td>
</tr>
<tr>
<td>Total</td>
<td>4.0</td>
<td>4.0</td>
<td>4.4</td>
<td>5.7</td>
</tr>
</tbody>
</table>
quantities as they are much smaller than the other errors. The resulting mixing angle is 

$$\phi = \pm (24.0^{+3.1+0.6}_{-2.6-0.8})^\circ.$$  

The systematic uncertainty is computed from the systematic errors assigned to the branching fractions.

The $f_1(1285)$ mixing angle has been estimated assuming that it is mixed with the $f_1(1420)$ state. Yang finds $\phi = (15.8^{+4.5}_{-4.6})^\circ$ using radiative decays [22], consistent with an earlier determination of $(15.4^{+3.9}_{-8})^\circ$ [23]. A lattice QCD analysis gives $(31 \pm 2)^\circ$, while another phenomenological calculation gives a range between $(20-30)^\circ$ [24]; see also Ref. [25] for other theoretical predictions. In this analysis we do not specify the other mixed partner.

If the $f_1(1285)$ is a tetraquark state its wave function would be $|f_1\rangle = (1/\sqrt{2})([su] |\bar{s}\bar{u}\rangle + [sd] |\bar{s}\bar{d}\rangle)$ in order for it to be produced significantly in both $B^0$ and $\bar{B}^0$ decays into $J/\psi f_1(1285)$. Using this wave function, the tetraquark model described in Ref. [8] predicts

$$\frac{B(B^0 \to J/\psi f_1(1285))}{B(\bar{B}^0 \to J/\psi f_1(1285))} = \frac{1}{4} \frac{|V_{ct}|^2}{|V_{ct}|^2} = 1.14\%,$$

with small uncertainties. Our measurement of this ratio of $(11.6 \pm 3.1^{+0.7}_{-0.8})\%$ differs by 3.3 standard deviations from the tetraquark interpretation including the systematic uncertainty.

Branching fraction ratios are converted into branching fractions using the previously measured rates listed in Table III. We correct the $B^0$ rates to reflect the updated value of the $B^0$ to $B^0$ production fraction of $0.259 \pm 0.015$ [20]. We determine

$$B(B^0 \to J/\psi f_1(1285)) = 0.36 \pm 0.64 \pm 0.42 \times 10^{-5},$$

where the first uncertainty is statistical, the second and third are systematic, being due to the relative branching fraction measurements and the errors in the absolute branching fraction normalization, respectively. For the $B^0$ decay this normalization error is due to the uncertainty on the production ratio of $B^0$ versus $\bar{B}^0$ and is 5.8% [5]. For the $B^0$ mode the uncertainty is due to the error of 4.1% on $\mathcal{B}(B^- \to J/\psi K^-)$ [6].

In conclusion, we report the first observations of $B^0$ and $\bar{B}^0 \to J/\psi f_1(1285)$ decays. These are also the first observations of the $f_1(1285)$ meson in heavy quark decays.

We express our gratitude to our colleagues in the CERN accelerator departments for the excellent performance of the LHC. We thank the technical and administrative staff at the LHCb institutes. We acknowledge support from CERN and from the national agencies CAPES, CNPq, FAPERJ, and FINEP (Brazil); NSFC (China); CNRS/IN2P3 and Region Auvergne (France); BMBF, DFG, HGF, and MPG (Germany); SFI (Ireland); INFN (Italy); FOM and NWO (Netherlands); SCSR (Poland); MEN/IFA (Romania);.

### Table III. Branching fractions used for normalization.

<table>
<thead>
<tr>
<th>Rate</th>
<th>Value</th>
<th>Ref.</th>
</tr>
</thead>
<tbody>
<tr>
<td>$B(B^0 \to J/\psi \pi^+ \pi^-)/B(\bar{B}^0 \to J/\psi \phi)$</td>
<td>$(19.79 \pm 0.47 \pm 0.52)%$</td>
<td>[4]</td>
</tr>
<tr>
<td>$B(B^0 \to J/\psi \pi^+ \pi^-)$</td>
<td>$(3.97 \pm 0.09 \pm 0.11 \pm 0.16) \times 10^{-5}$</td>
<td>[5]</td>
</tr>
<tr>
<td>$B(B^0 \to J/\psi \phi)$</td>
<td>$(10.50 \pm 0.13 \pm 0.64 \pm 0.82) \times 10^{-4}$</td>
<td>[6]</td>
</tr>
<tr>
<td>$B(B^- \to J/\psi K^-)$</td>
<td>$(10.18 \pm 0.42) \times 10^{-4}$</td>
<td>[6]</td>
</tr>
</tbody>
</table>
MinES, Rosatom, RFBR, and NRC “Kurchatov Institute” (Russia); MinECo, XuntaGal and GENCAT (Spain); SNSF and SER (Switzerland); NAS Ukraine (Ukraine); STFC (United Kingdom); NSF (USA). We also acknowledge the support received from the ERC under FP7. The Tier1 computing centres are supported by IN2P3 (France), KIT, and BMBF (Germany), INFN (Italy), NWO, and SURF (Netherlands), PIC (Spain), GridPP (United Kingdom). We are thankful for the computing resources put at our disposal by Yandex LLC (Russia), as well as to the communities and BMBF (Germany), INFN (Italy), NWO, and SURF computing centres are supported by IN2P3 (France), KIT, and BMBF (Germany), INFN (Italy), NWO, and SURF (Netherlands), PIC (Spain), GridPP (United Kingdom). We are thankful for the computing resources put at our disposal by Yandex LLC (Russia), as well as to the communities


(LHCb Collaboration)

1Centro Brasileiro de Pesquisas Físicas (CBPF), Rio de Janeiro, Brazil
2Universidade Federal do Rio de Janeiro (UFRJ), Rio de Janeiro, Brazil
3Center for High Energy Physics, Tsinghua University, Beijing, China
4LAPP, Université de Savoie, CNRS/IN2P3, Annecy-Le-Vieux, France
5Clermont Université, Université Blaise Pascal, CNRS/IN2P3, LPC, Clermont-Ferrand, France
6CPPM, Aix-Marseille Université, CNRS/IN2P3, Marseille, France
7LAL, Université Paris-Sud, CNRS/IN2P3, Orsay, France
8LPNHE, Université Pierre et Marie Curie, Université Paris Diderot, CNRS/IN2P3, Paris, France
9Fakultät Physik, Technische Universität Dortmund, Dortmund, Germany
10Max-Planck-Institut für Kernphysik (MPIK), Heidelberg, Germany
11Physikalisches Institut, Ruprecht-Karls-Universität Heidelberg, Heidelberg, Germany
12School of Physics, University College Dublin, Dublin, Ireland
13Sezione INFN di Bari, Bari, Italy
14Sezione INFN di Bologna, Bologna, Italy
15Sezione INFN di Cagliari, Cagliari, Italy
16Sezione INFN di Ferrara, Ferrara, Italy
17Sezione INFN di Firenze, Firenze, Italy
18Laboratori Nazionali dell’INFN di Frascati, Frascati, Italy
19Sezione INFN di Genova, Genova, Italy
20Sezione INFN di Milano Bicocca, Milano, Italy
21Sezione INFN di Padova, Padova, Italy
22Sezione INFN di Pisa, Pisa, Italy
23Sezione INFN di Roma Tor Vergata, Roma, Italy
24Sezione INFN di Roma La Sapienza, Roma, Italy
25Henryk Niewodniczanski Institute of Nuclear Physics Polish Academy of Sciences, Kraków, Poland
26Faculty of Physics and Applied Computer Science, AGH—University of Science and Technology, Kraków, Poland
27National Center for Nuclear Research (NCBJ), Warsaw, Poland
28Horia Hulubei National Institute of Physics and Nuclear Engineering, Bucharest-Magurele, Romania
29Petersburg Nuclear Physics Institute (PNPI), Gatchina, Russia
30Institute of Theoretical and Experimental Physics (ITEP), Moscow, Russia
31Institute of Nuclear Physics, Moscow State University (SINR MSU), Moscow, Russia
32Institute for Nuclear Research of the Russian Academy of Sciences (INR RAN), Moscow, Russia
33Budker Institute of Nuclear Physics (SB RAS) and Novosibirsk State University, Novosibirsk, Russia

091802-8