Search for heavy lepton partners of neutrinos in proton–proton collisions in the context of the type III seesaw mechanism

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Search for heavy lepton partners of neutrinos in proton–proton collisions in the context of the type III seesaw mechanism

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A search is presented in proton–proton collisions at \( \sqrt{s} = 7 \) TeV for fermionic triplet states expected in type III seesaw models. The search is performed using final states with three isolated charged leptons and an imbalance in transverse momentum. The data, collected with the CMS detector at the LHC, correspond to an integrated luminosity of 4.9 fb\(^{-1}\). No excess of events is observed above the background predicted by the standard model, and the results are interpreted in terms of limits on production cross sections and masses of the heavy partners of the neutrinos in type III seesaw models. Depending on the considered scenarios, lower limits are obtained on the mass of the heavy partner of the neutrino that range from 180 to 210 GeV. These are the first limits on the production of type III seesaw fermionic triplet states reported by an experiment at the LHC.

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1. Introduction

Experiments on neutrino oscillations [1–4] indicate that neutrinos have mass and their masses are much smaller than those of the charged leptons. However, the origin of neutrino mass is still unknown. An interesting possibility is provided by the seesaw mechanism, in which a small Majorana mass can be generated for each of the known neutrinos by introducing massive states with Yukawa couplings to leptons and to the Higgs field. Seesaw models called type I [5,6], type II [7–11], and type III [12,13] introduce heavy states of mass \( M \), that involve, respectively, weak-isospin singlets, scalar triplets, and fermion triplets. The neutrino masses are generically reduced relative to charged fermion masses by a factor \( v/M \), where \( v \) is the vacuum expectation value of the Higgs field. For sufficiently large \( M \) (of the order of \( 10^{14} \) GeV), small neutrino masses are generated even for Yukawa couplings of \( \approx 1 \). On the other hand, either smaller Yukawa couplings or extended seesaw mechanisms, such as those of the inverse seesaw models [14], are required to obtain small neutrino masses while keeping \( M \) close to a few hundreds of GeV. At the Large Hadron Collider (LHC), type II and III states can be produced through gauge interactions, so that the possible smallness of the Yukawa couplings does not affect the production cross section of the heavy states. In particular, the possibility of discovering a type III fermion at a proton–proton centre-of-mass energy of \( \sqrt{s} = 14 \) TeV is discussed in Refs. [15–17]. Recently, a leading-order (LO) computation of the signal expected at \( \sqrt{s} = 7 \) TeV has become available as a computer program for simulating such final states [18].

Given the electric charges of the lepton triplet, hereafter referred to as \( \Sigma^+, \Sigma^0, \) and \( \Sigma^- \), the most promising signature for finding a \( \Sigma \) state with a mass \( M_\Sigma \) of the order of a few hundreds of GeV is in production through quark–antiquark annihilation \( q\bar{q} \rightarrow \Sigma^0\Sigma^+ \), followed by the decays \( \Sigma^0 \rightarrow \ell^+\ell^–\nu \) and \( \Sigma^+ \rightarrow \ell^+\nu\Sigma^- \). The mass differences among the three electric charge states are assumed to be negligible. The mass range relevant for this analysis is bounded by the present lower limits (\( \approx 100 \) GeV) from the L3 experiment [19] and by the CMS loss of sensitivity near \( \approx 200 \) GeV because of the very steep decrease of the expected cross section with mass. Since there are twice as many u as d quark pairs in the proton, the production of \( \Sigma^+\Sigma^0 \) via virtual W\(^\pm\) bosons in the s-channel (Fig. 1) has the highest cross section of all the \( \Sigma \) charge combinations. (The cross section for the charge conjugate intermediary W\(^–\) is expected to be about a factor two smaller.) Selecting \( W^\pm \rightarrow \ell^\pm\nu \) decays (where \( \ell \) is an electron or muon) as the final states for the search, offers a very clean signature of three charged, isolated leptons. The decay \( \Sigma^+ \rightarrow \ell^+Z \), with \( Z \rightarrow \nu\bar{\nu} \) or \( Z \rightarrow q\bar{q} \), can also contribute significantly to the three-lepton final state, especially since its relative yield grows with \( M_\Sigma \). The \( \tau \) lepton also contributes to the three-lepton final states through \( \tau \rightarrow \ell\nu\nu \) decays. Details of the phenomenology and the different contributions to the final state of interest can be found in Ref. [18].

The total width of the \( \Sigma \) states and their decay branching fractions to SM leptons depend on the mixing matrix element for the
leptons \( V_\alpha \), where \( \alpha \) labels each of the e, \( \mu \), and \( \tau \) generations of leptons. Constraints on the mixing parameters and their products are available in Refs. [18,20].

The \( \Sigma \Sigma \) production cross section does not depend on the matrix elements \( V_\alpha \), which enter only in the \( \Sigma \) decays. The fraction of \( \Sigma \) decays to the lepton \( \alpha \) is proportional to:

\[
b_\alpha = \frac{|V_{\alpha e}|^2}{|V_{\alpha e}|^2 + |V_{\alpha \mu}|^2 + |V_{\alpha \tau}|^2}.
\]

If all three \( V_{\alpha e} \) values are less than \( \approx 10^{-6} \), the \( \Sigma \) states can have sufficiently long lifetimes to produce leptons at secondary vertices, a possibility not considered in this analysis.

This Letter reports on a search for fermionic triplet states expected in type III seesaw models, in final states with three charged leptons and an imbalance in transverse momentum (\( E_{\text{T}}^{\text{miss}} \)). The data sample corresponds to an integrated luminosity of 4.9 fb\(^{-1}\), collected in proton–proton collisions at \( \sqrt{s} = 7 \text{ TeV} \) with the Compact Muon Solenoid (CMS) detector at the LHC in 2011. The analysis is based on the model described in Ref. [15], using the implementation of Ref. [18]. Three possibilities are considered for the ratios \( b_\alpha \), defined in Eq. (1): first, \( b_\alpha = b_\beta = b_\mu = 1/3 \), hereafter referred to as the flavor-democratic scenario (FDS); second, \( b_\alpha = 1, b_\beta = 0, b_\mu = 1 \), and third, \( b_\alpha = 1 \) and \( b_\beta = 1, b_\mu = 0 \), hereafter referred to as the muon scenario (\( \mu S \)) and the electron scenario (\( eS \)), respectively.

2. The CMS detector

A detailed description of the CMS detector can be found in Ref. [21]. The central feature of the CMS apparatus is a superconducting solenoid that provides an axial magnetic field of 3.8 T. A silicon tracker, a lead-tungstate crystal electromagnetic calorimeter (ECAL), and a brass/scintillator hadron calorimeter (HCAL) reside within the magnetic field volume. Muons are identified using the central tracker and a muon system consisting of gas-ionization detectors embedded in the steel return yoke outside of the solenoid.

The directions of particles in the CMS detector are described using the azimuthal angle \( \phi \) and the pseudorapidity \( \eta \), defined as \( \eta = -\ln(\tan(\theta/2)) \), where \( \theta \) is the polar angle relative to the anticlockwise proton beam. All objects are reconstructed using a particle-flow (PF) algorithm [22–24]. The PF algorithm combines information from all subdetectors to identify and reconstruct particles detected in the collision, namely charged hadrons, photons, neutral hadrons, muons, and electrons. Jets are reconstructed using the anti-\( k_T \) jet clustering algorithm with a distance parameter of 0.5 [25]. Jet energies are corrected for non-uniformity in calorimeter response and for differences found between jets in simulation and in data [26]. An imbalance in transverse momentum (\( E_{\text{T}}^{\text{miss}} \)) is defined by the magnitude of the vectorial sum of the transverse momenta (\( p_T \)) of all particles reconstructed through the PF algorithm.

3. Simulation of signal and background

To estimate signal efficiency, \( \Sigma^+ \Sigma^0 \) events are generated using the FeynRules and MadGraph computer programs described in Ref. [18], while parton showers and hadronization are implemented using the Pythia generator (v6.420) [27]. The detector simulation is based on the Geant4 program [28]. Given the number of \( M_\Sigma \) mass points to be generated, part of the detector simulation is performed using the CMS Fast Simulation framework [29, 30]. Several background sources are considered in this analysis, the most relevant one being WZ production with both bosons decaying into leptons. A smaller contribution to the background comes from ZZ production, where the Z bosons decay leptonically, and one of the leptons is either outside of the detector acceptance or is mis-reconstructed. These two-boson events, calculated at next-to-LO with MCFM [31], are generated with Pythia. Backgrounds from the production of three EW bosons are generated with MadGraph 5 [32]. Backgrounds from jets and photons that are misidentified as leptons are also taken into account, including events from Drell–Yan \( \ell^+\ell^- + \) jets sources [33], \( W + \) jets, \( Z + \) jets, \( t\bar{t} \), and Drell–Yan \( \ell^+\ell^- + \gamma \) conversions to \( \ell^+\ell^- \). (The Drell–Yan process consists of qq \( \rightarrow \gamma^* / Z \rightarrow \ell^+\ell^- \) production, with \( \gamma^* \) and Z intermediares representing virtual \( \gamma \) or Z bosons.)

The presence of additional simultaneous pp interactions (pileup) is incorporated by simulating and mixing additional interactions with a multiplicity matching that observed in data.

4. Event selection criteria

The online trigger and the offline selection criteria are analogous to those used in other multi-lepton analyses performed by the CMS Collaboration [34,35]. Events are selected through two-lepton triggers in which two muons, two electrons, or one electron and one muon are required to be present. Because of the steady increase in instantaneous luminosity in 2011, some of the lepton \( p_T \) thresholds were increased over time to keep the trigger rates within the capabilities of the data acquisition system. For the two-muon trigger, the \( p_T \) requirements evolved from 7 GeV for each muon to asymmetric requirements of 17 GeV for the highest-\( p_T \) (leading) muon and 8 GeV for the second-highest \( p_T \) muon. For the two-electron trigger, the requirement is asymmetric, with a threshold applied to the energy of an ECAL cluster projected onto the plane transverse to the beam line (\( E_T = E \sin \theta \)). The cluster of the leading electron is required to have \( E_T > 17 \text{ GeV} \), and that of the next-to-leading electron to have \( E_T > 8 \text{ GeV} \). For the electron–muon trigger, the thresholds are either \( E_T > 17 \text{ GeV} \) for the electron and \( p_T > 8 \text{ GeV} \) for the muon, or \( E_T > 8 \text{ GeV} \) for the electron and \( p_T > 17 \text{ GeV} \) for the muon. The selected events must contain at least two lepton candidates with trajectories that have a transverse impact parameter of less than 0.2 mm relative to the principal interaction vertex. The chosen vertex is defined as the one with the largest value for the sum of the \( p_T^2 \) of the emanating tracks.

Muon candidates are reconstructed from a fit performed to hits in both the silicon tracker and the outer muon detectors, thereby defining a “global muon.” The specific selection requirements for a muon are: (i) \( p_T > 10 \text{ GeV} \), (ii) \( |\eta| < 2.4 \), (iii) more than 10 hits in the silicon tracker, and (iv) a global-muon fit with \( \chi^2 / \text{dof} < 10 \), where dof is the number of degrees of freedom.
Electron candidates are reconstructed using clusters of energy depositions in the ECAL that match the extrapolation of a reconstructed track. The electron track is fitted using a Gaussian-sum filter [36], with the algorithm taking into account the emission of bremsstrahlung photons in the silicon tracker. The specific requirements for a reconstructed electron are: (i) $p_T > 10$ GeV, (ii) $|\eta| < 1.44$, within the fully instrumented part of the central barrel, or $1.57 < |\eta| < 2.5$ for the endcap regions, (iii) not being a candidate for photon conversion, and (iv) the tracks reconstructed using three independent algorithms [23] to give the same sign for the electric charge.

All accepted lepton candidates are required to be isolated from other particles. In particular, selected muons must have $(\sum p_T)^{\mu}/p_T^{\mu} < 0.15$, where the sum over scalar $p_T$ includes all other PF objects within a cone of radius $\Delta R = \sqrt{(\Delta \eta)^2 + (\Delta \phi)^2} = 0.3$ of the muon track, where $\Delta \eta$ and $\Delta \phi$ are the differences in pseudorapidity and azimuthal angle between the lepton axis and the positions of other particles. Similarly, an electron candidate is accepted if $(\sum p_T)^{e}/p_T^{e} < 0.20$ within a cone of $\Delta R = 0.3$.

The candidate events used for the search are required to have: (i) three isolated charged leptons originating from the same primary vertex, as defined above, (ii) sum of the lepton charges equal to $-1$, (iii) $E_{\text{miss}} \geq 30$ GeV, (iv) $p_T > 18, 15, 10$ GeV for the lepton of highest, next-to-highest, and lowest $p_T$, and (v) $H_T < 100$ GeV, where $H_T$ is the scalar sum of the transverse momenta of jets with $p_T > 30$ GeV and $|\eta| < 2.4$, which reduces the background from $t\bar{t}$ events.

The selected events are classified into six categories that depend on lepton flavour and electric charge: $\mu^+\mu^-e^+$, $\mu^+e^-\mu^+$, $\mu^-\mu^+\mu^+$, $e^-\mu^+\mu^+$, $e^+\mu^-\mu^+$, and $e^+e^-e^-$. Except for the first and fourth categories, such configurations can also result from $W^+Z$ events. Fig. 2 shows the distributions of the $\mu^+\mu^-$ invariant mass for $\mu^+e^-\mu^+$ and $\mu^-\mu^+\mu^+$ events in data, before applying any requirement on the $\mu^+\mu^-$ mass, compared to the sum of SM background contributions. A peak in the $\mu^+\mu^-$ effective mass close to that of the Z boson is evident in both simulated events and in data. To reduce the background from $W^+Z$ events, a $Z$ veto is added to the selection requirements for the corresponding categories as follows. Events with at least one $\ell^+\ell^-$ mass combination in the range $82 < m_{\ell^+\ell^-} < 102$ GeV are rejected. To reject lepton pairs from decays of heavy-flavour quarks, events with $m_{\ell^+\ell^-} < 12$ GeV are also discarded.

Other sources of background in final states with three leptons arise from conversions of photons into additional $\ell^+\ell^-$ pairs through the process $Z \rightarrow \ell^+\ell^-\gamma \rightarrow \ell^+\ell^-\ell^+\ell^-\gamma^-$. If one of these additional leptons carries most of the momentum of the photon, the final state can appear as a three-lepton event. In such cases, the invariant mass of the $\ell^+\ell^-\ell^-$ state peaks close to the mass of the Z boson [34]. Since the probability of a photon conversion to electrons is higher than to muons, an additional Z veto of $82 < m_{\ell^+\ell^-\gamma} < 102$ GeV is applied to the $\mu^+e^-\mu^+$ and $e^+e^-e^-$ categories to reject such events. This is discussed further in the next section.

5. Background estimation

Three types of SM processes can produce a three-lepton final state: (i) events containing three or more prompt leptons from production and leptonic decays of two or three EW bosons. This is referred to as irreducible background, since it corresponds to the same final states as the signal from $\Sigma$ production, (ii) $V + \gamma$ and $V + \gamma^*$ events, where $V$ represents any EW boson, with the accompanying photons converting to $\ell^+\ell^-$, and (iii) events with one or two prompt leptons and additional non-prompt leptons that arise from lepton decays of hadrons within jets, called “misidentified jets”.

The irreducible background from more than two leptons is dominated by SM WZ production, but also includes ZZ and three-boson events. The two-boson contribution, which is reduced substantially by the $Z$ mass veto, and the three-boson contribution, which is dominated by the WWW channel, are both evaluated using MC simulation. The contribution from three-boson production is small relative to the other sources, as shown in Table 1.

As mentioned in Section 4, photon conversions in the presence of $W$ or $Z$ bosons can produce isolated leptons that constitute another source of background. External conversions of photons, namely of produced photons that interact with the material in the detector to yield primarily $e^+e^-$ pairs, are evaluated from simulation ($V\gamma$ in Table 1). Internal conversions, involving the direct materialisation of virtual photons into $\mu^-\mu^+$ or $e^-e^+$ pairs, can also provide a similar source of background. Both external and internal conversions can become problematic when one of the two final-state leptons carries off most of the photon energy, and the second lepton is not detected. The contribution of conversions to electrons is reduced by the additional three-lepton-mass rejection applied to the $\mu^-e^-\mu^+$ and $e^-e^-e^+$ categories as discussed above. The
The contribution from internal photon conversions to muons $\gamma^* \to \mu^+\mu^-$ is evaluated according to the method described in Ref. [34], where the ratio of $\ell^+\ell^-\mu^\pm$ to $\ell^+\ell^-\gamma$ events, in which the mass is close to that of a Z boson, defines a conversion factor $C_\gamma$ for muons. The background is estimated from $C_\mu$ and from the number of $\ell^+\ell^-\gamma$ events in data that pass all selections, except the three-lepton requirements. An alternative evaluation is obtained from events in an independent Z-enriched control region, by reversing the $E_\text{miss}$ requirement to $E_\text{miss} < 20$ GeV. As mentioned before, events from $Z$ decays into two muons or two electrons that contain an additional muon from internal photon conversion, produce a peak in the three-lepton invariant mass distribution close to the $Z$ mass. The number of events expected in the final sample is estimated from the ratio of simulated events for $Z$ production with $E_\text{miss} > 30$ GeV to that with $E_\text{miss} < 20$ GeV. This estimate agrees with that of the previous method. The $\gamma^* \to \mu^+\mu^-$ background contribution is small, as can be seen in Table 1. An overall uncertainty of ±50% is assumed for this source of background, which is limited by the statistical precision of both estimates (30%), and has an additional contribution from the choice of normalization criteria (40%).

The largest background, aside from the irreducible backgrounds, arises from the $Z + j$ jets process (including the Drell–Yan contribution), in which the $Z$ boson decays leptonically, and a jet in the event is misidentified as a third lepton. Processes with non-prompt leptons from heavy-flavour decays are not simulated with that of the previous method. The $\gamma^* \to \mu^+\mu^-$ background contribution is small, as can be seen in Table 1. An overall uncertainty of ±50% is assumed for this source of background, which is limited by the statistical precision of both estimates (30%), and has an additional contribution from the choice of normalization criteria (40%).

The overall uncertainty on integrated luminosity is 2.2% [41]. For backgrounds determined from simulation, the systematic uncertainties on efficiency and luminosity are common to all signals.

### Table 1

Summary of the mean number of SM background events expected in each event category, after final selections. $V$ represents a $Z$ or a $W$ bosons and $V_f$ is the contribution from external photon conversions. The column labelled “Misidentified jets” includes backgrounds with non-prompt leptons, the column $\gamma^* \to \mu^+\mu^-$ shows the background expectation from internal photon conversions, where a virtual photon converts to a muon pair, and one muon is lost. The contribution of $\gamma^* \to e^+e^-$ is removed by the rejection criteria on three-lepton masses. Statistical uncertainties are included for the six categories, and systematic uncertainties on normalizations are listed in the last row.

<table>
<thead>
<tr>
<th>$\ell^-\ell^-\ell^+$</th>
<th>$\ell^-\ell^-\ell^+$</th>
<th>$\ell^-\ell^-\ell^+$</th>
<th>$\ell^-\ell^-\ell^+$</th>
<th>Misidentified jets</th>
<th>$\gamma^* \to \mu^+\mu^-$</th>
</tr>
</thead>
<tbody>
<tr>
<td>$\mu^- e^+ e^+$</td>
<td>$0.3 \pm 0.1$</td>
<td>$0.09 \pm 0.01$</td>
<td>–</td>
<td>$0.4 \pm 0.4$</td>
<td>–</td>
</tr>
<tr>
<td>$\mu^- e^+ e^+$</td>
<td>$4.0 \pm 0.3$</td>
<td>$0.19 \pm 0.01$</td>
<td>–</td>
<td>$3.1 \pm 1.2$</td>
<td>–</td>
</tr>
<tr>
<td>$\mu^- e^+ e^+$</td>
<td>$4.9 \pm 0.3$</td>
<td>$0.11 \pm 0.01$</td>
<td>–</td>
<td>$5.7 \pm 1.9$</td>
<td>$0.7 \pm 0.2$</td>
</tr>
<tr>
<td>$e^- \mu^+ e^+$</td>
<td>$0.3 \pm 0.1$</td>
<td>$0.09 \pm 0.01$</td>
<td>–</td>
<td>$0.8 \pm 0.5$</td>
<td>–</td>
</tr>
<tr>
<td>$e^- e^+ \mu^+$</td>
<td>$4.9 \pm 0.3$</td>
<td>$0.21 \pm 0.02$</td>
<td>–</td>
<td>$3.0 \pm 1.2$</td>
<td>$0.4 \pm 0.1$</td>
</tr>
<tr>
<td>$e^- e^+ e^+$</td>
<td>$2.5 \pm 0.2$</td>
<td>$0.06 \pm 0.01$</td>
<td>$1.4 \pm 1.0$</td>
<td>$1.1 \pm 0.6$</td>
<td>–</td>
</tr>
</tbody>
</table>

**Normalization uncertainties**: 17% (WZ) 7.5% (ZZ) 50% 13% 50% 50%
Table 2
Uncertainties on signal efficiency for each event category for \( M_\Sigma = 180 \) GeV. Total systematic and total systematic + statistical (fourth and sixth columns) are calculated in quadrature.

<table>
<thead>
<tr>
<th>Source of uncertainty</th>
<th>Trigger</th>
<th>Signal efficiency (Full simulation)</th>
<th>(Fullsim/Fastsim) systematic</th>
<th>Total systematic</th>
<th>(Fullsim/Fastsim) statistical</th>
<th>Total syst. + stat.</th>
</tr>
</thead>
<tbody>
<tr>
<td>( \mu^- e^+ e^+ )</td>
<td>1.0%</td>
<td>6.3%</td>
<td>2.9%</td>
<td>7.0%</td>
<td>3.0%</td>
<td>7.6%</td>
</tr>
<tr>
<td>( \mu^- e^+ \mu^+ )</td>
<td>1.0%</td>
<td>4.5%</td>
<td>6.8%</td>
<td>8.2%</td>
<td>2.3%</td>
<td>8.5%</td>
</tr>
<tr>
<td>( \mu^- \mu^+ \mu^- )</td>
<td>1.0%</td>
<td>3.9%</td>
<td>11.1%</td>
<td>11.8%</td>
<td>3.3%</td>
<td>12.2%</td>
</tr>
<tr>
<td>( e^- e^+ \mu^+ )</td>
<td>1.0%</td>
<td>4.5%</td>
<td>8.5%</td>
<td>9.7%</td>
<td>2.9%</td>
<td>10.1%</td>
</tr>
<tr>
<td>( e^- e^+ e^- )</td>
<td>1.0%</td>
<td>6.3%</td>
<td>4.1%</td>
<td>7.6%</td>
<td>2.4%</td>
<td>7.9%</td>
</tr>
<tr>
<td>( e^- e^+ e^- )</td>
<td>1.0%</td>
<td>7.6%</td>
<td>2.8%</td>
<td>8.0%</td>
<td>4.2%</td>
<td>9.1%</td>
</tr>
</tbody>
</table>

Table 3
Summary of the expected mean number of events for signal as a function of \( M_\Sigma \), for the expected SM background, and the observed number of events in data, after implementing all analysis selections. Each of the three possibilities for mixing (FDS, \( \mu S, eS \)) described in Section 1 are considered separately in the analysis.

<table>
<thead>
<tr>
<th>Category</th>
<th>Expected signal for ( M_\Sigma ) (GeV)</th>
<th>( \mu S )</th>
<th>( eS )</th>
<th>Expected background</th>
<th>Observed in data</th>
</tr>
</thead>
<tbody>
<tr>
<td></td>
<td>FDS</td>
<td>120</td>
<td>130</td>
<td>140</td>
<td>180</td>
</tr>
<tr>
<td>( \mu^- e^+ e^+ )</td>
<td>7.9</td>
<td>6.0</td>
<td>4.5</td>
<td>1.7</td>
<td>1.1</td>
</tr>
<tr>
<td>( \mu^- e^+ \mu^+ )</td>
<td>12.3</td>
<td>9.0</td>
<td>7.0</td>
<td>3.0</td>
<td>2.0</td>
</tr>
<tr>
<td>( \mu^- \mu^+ \mu^- )</td>
<td>7.8</td>
<td>5.2</td>
<td>3.6</td>
<td>1.4</td>
<td>0.93</td>
</tr>
<tr>
<td>( e^- e^+ \mu^+ )</td>
<td>8.3</td>
<td>6.2</td>
<td>4.8</td>
<td>1.8</td>
<td>1.2</td>
</tr>
<tr>
<td>( e^- e^+ e^- )</td>
<td>13.2</td>
<td>9.5</td>
<td>6.9</td>
<td>2.7</td>
<td>1.8</td>
</tr>
<tr>
<td>( e^- e^+ e^- )</td>
<td>3.9</td>
<td>2.8</td>
<td>2.0</td>
<td>1.0</td>
<td>0.63</td>
</tr>
</tbody>
</table>

7. Results

Table 3 presents the results of our search for the fermionic \( \Sigma \) triplet states in terms of the expected number of signal events, the expected number of events from SM background, and the number of observed events in each of the analyzed event categories. Each of the three possibilities for mixing (FDS, \( \mu S, eS \)) described in Section 1 is considered in the analysis.

No significant excess of events is observed relative to the SM expectations in any of the six analysis channels. Combining all channels, we set upper limits at the 95% confidence level (CL) on \( \sigma \times B \), the product of the production cross section of \( \Sigma^- \Sigma^0 \) and its branching fraction \( B \) to the three-lepton final states, where the lepton can be an electron, muon or \( \tau \) (contribution through \( \tau \rightarrow \ell \nu \ell \nu \)). The branching fraction to three-lepton final states depends on \( M_\Sigma \) [18], and is predicted to be about 9% for \( M_\Sigma \approx 180 \) GeV, where we extrapolate signal yields to \( M_\Sigma > 180 \) GeV using the results of Ref. [18].

The upper limits on \( \sigma B \) as a function of fermion mass \( M_\Sigma \), combining for all channels by multiplying the corresponding likelihood functions, are shown in Fig. 3, 4, and 5, for FDS, \( \mu S \), and \( eS \) possibilities, respectively. The dashed lines correspond to the expected limits obtained from MC pseudo-experiments, and are based on the CLs criterion [42,43]. The observed limits on data are computed following both a Bayesian approach [33, Ch. 33], and a frequentist method also based on the CLs criterion. In the former, the assumed prior is a constant. In both calculations, the uncertainties on efficiencies for detecting the signal, the uncertainty on integrated luminosity and on the expected SM background, are treated as uninteresting “nuisance” parameters with Gaussian or log-normal densities. Upper limits are computed at 95% CL using the RooStats software [44], and the package developed to combine results from searches for the Higgs boson [45]. The two results are similar, as shown in Figs. 3, 4, and 5. The results are stable relative to variations of ±20% on the systematic uncertainties. Finally, we extract lower limits on \( M_\Sigma \) using the theoretical dependence of the cross section on \( M_\Sigma \), as represented by the solid blue lines of Fig. 3, 4, and 5, for the three possibilities for the type III seesaw model for signal. The expected and observed 95% CL limits obtained with the Bayesian method are given in Table 4.

8. Summary

A search has been presented for fermionic triplet states expected in type III seesaw models. The search was performed in events with three isolated leptons (muons or electrons), whose charges sum to +1, and contain jets and an imbalance in transverse
Fig. 4. The expected (dashed line) and observed (asterisks and black points) exclusion limits at 95% confidence level on $\sigma B$ as a function of the fermion mass $M_{\Sigma}$, assuming $b_\nu = 0$, $b_\mu = 1$, $b_\tau = 0$ ($\mu S$) for the signal. The solid (blue) curve represents the predictions of the LO type III seesaw models. The light (yellow) and dark (green) shaded areas represent, respectively, the 1 standard deviation (68% CL) and 2 standard deviations (95% CL) limits on the expected results obtained from MC pseudo-experiments, which reflect the combined statistical and systematic uncertainties of the SM contributions. The asterisks and the black points show, respectively, the observed limits computed following a frequentist method based on the CLs criterion and a Bayesian approach.

Fig. 5. The expected (dashed line) and observed (black points) exclusion limits at 95% confidence level on $\sigma B$ as a function of the fermion mass $M_{\Sigma}$, assuming $b_\nu = 1$, $b_\mu = 0$, $b_\tau = 0$ (eS) for the signal. The solid (blue) curve represents the predictions of the LO type III seesaw models. The light (yellow) and dark (green) shaded areas represent, respectively, the 1 standard deviation (68% CL) and 2 standard deviations (95% CL) limits on the expected results obtained from MC pseudo-experiments, which reflect the combined statistical and systematic uncertainties of the SM contributions. The asterisks and the black points show, respectively, the observed limits computed following a frequentist method based on the CLs criterion and a Bayesian approach.

Table 4

<table>
<thead>
<tr>
<th>Scenario</th>
<th>$95% \text{ CL: } \sigma B$ (fb)</th>
<th>$95% \text{ CL: } M_{\Sigma}$ (GeV)</th>
</tr>
</thead>
<tbody>
<tr>
<td>FDS</td>
<td>22</td>
<td>20</td>
</tr>
<tr>
<td>$\mu S$</td>
<td>13</td>
<td>11</td>
</tr>
<tr>
<td>eS</td>
<td>13</td>
<td>13</td>
</tr>
</tbody>
</table>

The data are from proton–proton collisions at $\sqrt{s} = 7$ TeV, recorded during 2011 by the CMS experiment at the CERN LHC, and correspond to an integrated luminosity of 4.9 fb$^{-1}$.

No evidence for pair production of $\Sigma^+ \Sigma^0$ states has been found, and 95% confidence upper limits are set on the product of the production cross section of $\Sigma^+ \Sigma^0$ and its branching fraction to the examined three-lepton final states. Comparing the results with predictions from type III seesaw models, lower bounds are established at 95% confidence on the mass of the $\Sigma$ states. Limits are reported for three choices of mixing possibilities between the $\Sigma$ states and the three lepton generations. Depending on the considered scenarios, lower limits are obtained on the mass of the heavy partner of the neutrino that range from 180 to 210 GeV. The results are valid only if at least one of the mixing matrix elements is larger than $\sim 10^{-6}$. These are the first limits on the production of type III seesaw fermionic triplet states reported by an experiment at the LHC.

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