Search for neutral Higgs bosons decaying to tau pairs in pp collisions at s = 7 TeV
Search for neutral Higgs bosons decaying to tau pairs in pp collisions at √s = 7 TeV

CERN Collaboration

CERN, Switzerland

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ABSTRACT

A search for neutral Higgs bosons decaying to tau pairs at a center-of-mass energy of 7 TeV is performed using a dataset corresponding to an integrated luminosity of 4.6 fb⁻¹ recorded by the CMS experiment at the LHC. The search is sensitive to both the standard model Higgs boson and to the neutral Higgs bosons predicted by the minimal supersymmetric extension of the standard model (MSSM). No excess of events is observed in the tau-pair invariant-mass spectrum. For a standard model Higgs boson in the mass range of 110–145 GeV upper limits at 95% confidence level (CL) on the production cross section are determined. We exclude a Higgs boson with m_H = 115 GeV with a production cross section 3.2 times that of the standard model. In the MSSM, upper limits on the neutral Higgs boson production cross section times branching fraction to tau pairs, as a function of the pseudoscalar Higgs boson mass, m_A, sets stringent new bounds in the parameter space, excluding at 95% CL values of tanβ as low as 7.1 at m_A = 160 GeV in the m_{max} benchmark scenario.

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1. Introduction

An important goal of the LHC physics program is to ascertain the mechanism of electroweak symmetry breaking, through which the W and Z bosons attain mass, while the photon remains massless. In the standard model (SM) [1–3], this is achieved via the Higgs mechanism [4–9], which also predicts the existence of a scalar Higgs boson. However, this particle has not yet been observed by experiments. Moreover, the mass of the Higgs boson is quadratically divergent at high energies [10]. Supersymmetry [11] is a well known extension to the SM which allows the cancellation of this divergence.

The minimal supersymmetric standard model (MSSM) contains two Higgs doublets, giving rise to five physical states: a light neutral CP-even state (h), a heavy neutral CP-even state (H), a neutral CP-odd state (A), and a pair of charged states (H⁺, H⁻) [12–15]. The mass relations between these particles depend on the MSSM parameter tanβ, the ratio of the Higgs fields vacuum expectation values. We focus on the m_{max} [16,17] benchmark scenario in which M_{STOP} = 1 TeV; X_t = 2 M_{STOP}; M_{H_2} = 800 GeV; M_{h_2} = 200 GeV; and A_h = A_t. Here, M_{STOP} denotes the common soft-SUSY-breaking squark mass of the third generation; X_t = M_{TOP} − μ/ tanβ is the stop mixing parameter; A_t and A_h are the stop and sbottom trilinear couplings, respectively; μ the Higgsino mass parameter; M_{3} the gluino mass; and M_{2} is the SU(2)-gaugino mass parameter. The value of M_{1} is fixed via the unification relation M_{1} = (5/3) M_{2} sinθ_{W} / cosθ_{W}. In this scenario for values of tanβ ≥ 15, if m_{h} ≤ 130 GeV the masses of the h and A are almost degenerate, while the mass of the H is around 130 GeV. Conversely, if m_{A} ≥ 130 GeV, the masses of the A and H are almost degenerate, while the mass of the H remains near 130 GeV. This will thus always lead to one neutral Higgs boson at 130 GeV and two neutral Higgs bosons with almost degenerate mass of m_{h}.

Direct searches for the SM Higgs boson at the Large Electron–Positron Collider (LEP) set a limit on the mass of m_{H} > 114.4 GeV at 95% confidence level (CL) [18]. The Tevatron collider experiments exclude the SM Higgs boson in the mass range 162–166 GeV [19], and the ATLAS experiment in the mass ranges 112.9–115.5, 131–238, and 251–466 GeV [20]. Precision electroweak data constrain the mass of the SM Higgs boson to be less than 158 GeV [21]. Direct searches for neutral MSSM Higgs bosons have been reported by LEP, the Tevatron, and both LHC experiments, and set limits on the MSSM parameter space in the tanβ–mA plane [22–26].

This Letter reports a search for the SM and the neutral MSSM Higgs bosons using final states with tau pairs in proton–proton colli- sions at √s = 7 TeV at the LHC. We use a data sample collected in 2011 corresponding to an integrated luminosity of 4.6 fb⁻¹ recorded by the Compact Muon Solenoid (CMS) [27] experiment. Three independent tau-pair final states where one or both taus decay leptonically are studied: eτ + X, μτ + X, and eμ + X, where we use the symbol τ_b to indicate a reconstructed hadronic decay of a tau.
In the case of the SM Higgs boson, the gluon-fusion production mechanism has the largest cross section. However, in the mass region of interest, background from Drell–Yan production of tau pairs overwhelms the expected Higgs boson signal. This search therefore relies upon the signature of Higgs bosons produced via vector boson fusion (VBF) or in association with a high-\(p_T\) jet. In the former case, the distinct topology of two jets with a large rapidity separation greatly reduces the background. In the latter, requiring a high-\(p_T\) jet both suppresses background, and improves the measurement of the tau-pair invariant mass.

In the MSSM case, two main production processes contribute to \(pp \rightarrow h + X\), where \(h = h, H, or A\): gluon fusion through a b-quark loop and direct \(bb\) annihilation from the b-quark content of the beam protons. In the latter case, there is a significant probability that a b-quark jet is produced centrally in association with the beam protons. In the former, requiring a b-quark jet increases the sensitivity of the search by reducing the \(Z + jets\) background.

2. CMS detector

The CMS detector is described in detail elsewhere [27]. The central feature of the CMS apparatus is a superconducting solenoid of 6 m internal diameter, providing a magnetic field of 3.8 T. Within the solenoid are the silicon pixel and strip tracker, which cover a pseudorapidity region of \(|\eta| < 2.5\). Here, the pseudorapidity is defined as \(\eta = -\ln(\tan(\theta/2))\), where \(\theta\) is the polar angle of the trajectory of the particle with respect to the direction of the counterclockwise beam. The lead tungstate crystal electromagnetic calorimeter and the brass-scintillator hadron calorimeter surround the central feature of the CMS apparatus is a superconducting solenoid cover a pseudorapidity region of \(|\eta| < 5\). Muons are measured in gas-ionization detectors embedded in the steel return yoke, with a coverage of \(|\eta| < 2.4\).

3. Trigger and event selection

The analysis makes use of the three independent tau-pair final states, \(e\tau + X, \mu\tau + X,\) and \(e\mu + X\). In all three channels, there is substantial background, both from processes with similar experimental signatures, and from unrelated hadronic activity in the detector.

The trigger selection required a combination of electron, muon and tau trigger objects [28–30]. The identification criteria and \(p_T\) thresholds of these objects were progressively tightened as the LHC instantaneous luminosity increased over the data-taking period.

A particle-flow algorithm [31–33] is used to combine information from all CMS subdetectors to identify and reconstruct individual particles in the event, namely muons, electrons, photons, and charged and neutral hadrons. From the resulting particle list, jets, hadronically-decaying taus, and missing transverse energy (\(E_T^{\text{miss}}\)), defined as the negative of the vector sum of the transverse momenta, are reconstructed. The jets are identified using the anti-\(k_T\) jet algorithm [34,35] with a distance parameter of \(R = 0.5\). Hadronically-decaying taus are reconstructed using the hadron plus strips (HPS) algorithm, which considers candidates with one or three charged pions and up to two neutral pions [36].

For the \(e\tau + X\) and \(\mu\tau + X\) final states, in the region \(|\eta| < 2.1\), we select events with an electron of \(p_T > 20\) GeV or a muon of \(p_T > 17\) GeV, together with an oppositely charged \(\tau\) of \(p_T > 20\) GeV within the range \(|\eta| < 2.3\). For the \(e\mu + X\) final state, we select events with an electron of \(|\eta| < 2.3\) and an oppositely charged muon of \(|\eta| < 2.1\), requiring \(p_T > 20\) GeV for the highest-\(p_T\) lepton and \(p_T > 10\) GeV for the next-to-highest-\(p_T\) lepton.

In the case of the SM Higgs boson, the gluon-fusion production mechanism has the largest cross section. However, in the mass region of interest, background from Drell–Yan production of tau pairs overwhelms the expected Higgs boson signal. This search therefore relies upon the signature of Higgs bosons produced via vector boson fusion (VBF) or in association with a high-\(p_T\) jet. In the former case, the distinct topology of two jets with a large rapidity separation greatly reduces the background. In the latter, requiring a high-\(p_T\) jet both suppresses background, and improves the measurement of the tau-pair invariant mass.

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Table 1
Numbers of expected and observed events in the event categories as described in the text for the $e\tau_\ell + X$ channel. Also given are the expected signal yields and efficiencies for an MSSM Higgs boson with $m_A = 120$ GeV and $\tan\beta = 10$, and for an SM Higgs boson with $m_h = 120$ GeV. Combined statistical and systematic uncertainties on each estimate are reported. For the yield estimates for the Higgs signal the production cross sections for $h$ and $A$, which have almost degenerate masses, are taken into account. The quoted efficiencies do not include the branching fraction into $\tau\tau$.

<table>
<thead>
<tr>
<th>Process</th>
<th>SM</th>
<th>MSSM</th>
</tr>
</thead>
<tbody>
<tr>
<td>$Z \rightarrow \tau\tau$</td>
<td>0/1-Jet</td>
<td>Boosted</td>
</tr>
<tr>
<td>$Z \rightarrow \tau\tau$</td>
<td>$13 438 \pm 977$</td>
<td>$190 \pm 14$</td>
</tr>
<tr>
<td>Multijets</td>
<td>$63 55 \pm 299$</td>
<td>$27 \pm 3$</td>
</tr>
<tr>
<td>W+jets</td>
<td>$2983 \pm 216$</td>
<td>$62 \pm 4$</td>
</tr>
<tr>
<td>Z $\rightarrow$ $\ell$$\ell$</td>
<td>$5170 \pm 464$</td>
<td>$28 \pm 4$</td>
</tr>
<tr>
<td>t$\bar{t}$</td>
<td>$63 \pm 7$</td>
<td>$42 \pm 6$</td>
</tr>
<tr>
<td>Dibosons</td>
<td>$68 \pm 21$</td>
<td>$5 \pm 2$</td>
</tr>
<tr>
<td>Total background</td>
<td>$28 087 \pm 1142$</td>
<td>$354 \pm 17$</td>
</tr>
<tr>
<td>H $\rightarrow$ $\tau\tau$</td>
<td>$53 \pm 9$</td>
<td>$2.7 \pm 0.6$</td>
</tr>
<tr>
<td>Data</td>
<td>$27727$</td>
<td>$318$</td>
</tr>
</tbody>
</table>

Table 2
Numbers of expected and observed events in the event categories as described in the text for the $\mu\tau_\ell + X$ channel. Also given are the expected signal yields and efficiencies for an MSSM Higgs boson with $m_A = 120$ GeV and $\tan\beta = 10$, and for an SM Higgs boson with $m_h = 120$ GeV. Combined statistical and systematic uncertainties on each estimate are reported. For the yield estimates for the Higgs signal the production cross sections for $h$ and $A$, which have almost degenerate masses, are taken into account. The quoted efficiencies do not include the branching fraction into $\tau\tau$.

<table>
<thead>
<tr>
<th>Process</th>
<th>SM</th>
<th>MSSM</th>
</tr>
</thead>
<tbody>
<tr>
<td>$Z \rightarrow \tau\tau$</td>
<td>0/1-Jet</td>
<td>Boosted</td>
</tr>
<tr>
<td>$Z \rightarrow \tau\tau$</td>
<td>$28 955 \pm 2054$</td>
<td>$295 \pm 22$</td>
</tr>
<tr>
<td>Multijets</td>
<td>$7841 \pm 141$</td>
<td>$36 \pm 2$</td>
</tr>
<tr>
<td>W+jets</td>
<td>$5827 \pm 392$</td>
<td>$65 \pm 4$</td>
</tr>
<tr>
<td>$Z \rightarrow$ $\ell$$\ell$</td>
<td>$777 \pm 70$</td>
<td>$5 \pm 1$</td>
</tr>
<tr>
<td>t$\bar{t}$</td>
<td>$147 \pm 15$</td>
<td>$94 \pm 12$</td>
</tr>
<tr>
<td>Dibosons</td>
<td>$178 \pm 55$</td>
<td>$9 \pm 4$</td>
</tr>
<tr>
<td>Total background</td>
<td>$43 725 \pm 2097$</td>
<td>$504 \pm 26$</td>
</tr>
<tr>
<td>H $\rightarrow$ $\tau\tau$</td>
<td>$96 \pm 17$</td>
<td>$3.9 \pm 0.8$</td>
</tr>
<tr>
<td>Data</td>
<td>$43 612$</td>
<td>$500$</td>
</tr>
</tbody>
</table>

Signal efficiency

<table>
<thead>
<tr>
<th>Process</th>
<th>SM</th>
<th>MSSM</th>
</tr>
</thead>
<tbody>
<tr>
<td>$gg \rightarrow \phi$</td>
<td>–</td>
<td>–</td>
</tr>
<tr>
<td>$gg \rightarrow b\bar{b}$</td>
<td>–</td>
<td>–</td>
</tr>
<tr>
<td>$gg \rightarrow H$</td>
<td>$1.7 \cdot 10^{-2}$</td>
<td>$3.9 \cdot 10^{-4}$</td>
</tr>
<tr>
<td>$qq \rightarrow q\bar{q}H$</td>
<td>$8.6 \cdot 10^{-3}$</td>
<td>$2.6 \cdot 10^{-3}$</td>
</tr>
<tr>
<td>$qq \rightarrow H\ell\ell$ or VH</td>
<td>$1.5 \cdot 10^{-2}$</td>
<td>$3.3 \cdot 10^{-3}$</td>
</tr>
</tbody>
</table>

selected events into several mutually exclusive categories based on the jet multiplicity and b-content.

In the MSSM case, there is a large probability for having a b-tagged jet in the central region. We use an algorithm based on the impact parameter of the tracks associated to the event vertex to identify b-tagged jets [40]. The MSSM search has three categories:

**b-Tag category**: We require at least one b-tagged jet with $p_T > 30$ GeV and at least one b-tagged jet with $p_T > 20$ GeV.

**Non-b-Tag category**: We require at least one jet with $p_T > 30$ GeV and no b-tagged jet with $p_T > 20$ GeV.

The SM search has three categories:

**VBF category**: We require at least two jets with $p_T > 30$ GeV, $|\Delta\eta_{jj}| > 4.0$, $\eta_1 \cdot \eta_2 < 0$, and a dijet invariant mass $m_{jj} > 400$ GeV, with no other jet with $p_T > 30$ GeV in the rapidity region between the two jets.

**Boosted category**: We require one jet with $p_T > 150$ GeV, and, in the $\mu\mu$ channel, no b-tagged jet with $p_T > 20$ GeV.

**0/1-Jet category**: We require no more than one jet with $p_T > 30$ GeV, and if such a jet is present, it must have $p_T < 150$ GeV.
The normalization for this process is determined from the measurement of the Z → ee and Z → μμ cross section [41]. Another significant source of background is multijet events in which there is one jet misidentified as an isolated electron or muon, and a second jet misidentified as a τ, W + jets in which there is a jet misidentified as a τ, and a second jet misidentified as a τ. Other background processes include tt production and Z → ee/μμ events, particularly in the ττ + X channel due to the 2–3% probability for electrons to be misidentified as τ. The small background from W + jets and multijet events for the eμ channel where jets are misidentified as isolated leptons is derived by measuring the number of events with one good lepton and a second which passes relaxed selection criteria, but fails the nominal lepton selection. This sample is extrapolated to the signal region using the efficiencies for such loose lepton candidates to pass the nominal lepton selection. These efficiencies are measured in data using multijet events. Backgrounds from tt and di-boson production are estimated from simulation using the MADGRAPH [42] event generator to simulate the shapes for tt events and PYTHIA 6.424 [43] to simulate the shapes for di-boson events. The event yields are normalized to the inclusive cross sections: $$σ_{tt} = 164.4 \pm 14.3 \text{ pb}$$ and $$σ_{WZ} = 55.3 \pm 8.3 \text{ pb}$$ as measured with an analysis similar to that described in [44,45] using a larger data sample.

To model the MSSM and SM Higgs boson signals the event generators PYTHIA and POWHEG [46] are used, respectively. The TAUOLA [47] package is used for tau decays in all cases. Additional next-to-next-to-leading order (NNLO) K-factors from FeHiPro [48,49] are applied to the Higgs boson pT spectrum from Higgs boson events produced via gluon fusion.

The presence of pile-up is incorporated by simulating additional interactions and then reweighting the simulated events to match the distribution of additional interactions observed in data. The events in the embedded Z → ττ sample and in other background samples obtained from data contain the correct distribution of pile-up interactions. The missing transverse energy response from simulation is corrected using a prescription, based on data, developed for inclusive W and Z cross section measurements [41], where Z bosons are reconstructed in the dimuon channel, and the missing transverse energy scale and resolution calibrated as a function of the Z boson transverse momentum.

4. Tau-pair invariant mass reconstruction

To distinguish the Higgs boson signal from the background, we reconstruct the tau-pair mass using a maximum likelihood technique [26]. The algorithm estimates the original momentum components of the two taus by maximizing a likelihood with respect to free parameters corresponding to the missing neutrino momenta, subject to kinematic constraints. Other terms in the likelihood take into account the tau-decay phase space and the probability density in the tau transverse momentum, parametrized as a function of the tau-pair mass. This algorithm yields a tau-pair mass with a mean consistent with the true value, and a distribution with a nearly Gaussian shape. The standard deviation of the mass resolution is estimated to be 21% at a Higgs boson mass of 130 GeV, compared with 24% for the (non-Gaussian) distribution of the invariant mass spectrum reconstructed from the visible tau-decay products in the inclusive selection. The resolution improves to 15% in the b-Tag category in the MSSM analysis and in the Boosted and VBF categories in the SM analysis where the Higgs boson is produced with significant transverse momentum.

5. Systematic uncertainties

Various imperfectly known or simulated effects can alter the shape and normalization of the invariant mass spectrum. The main contributions to the normalization uncertainty include the uncertainty in the total integrated luminosity (4.5%) [50], jet energy scale (2–5% depending on η and pT), background normalization (Tables 1–3), Z boson production cross section (2.5%) [41], lepton identification and isolation efficiency (1.0%), and trigger efficiency (1.0%). The tau-identification efficiency uncertainty is estimated to be 6% from an independent study using a tag-and-probe technique [41]. The lepton identification and isolation efficiencies are stable as a function of the number of additional interactions in the bunch crossing in data and in Monte Carlo simulation. The b-tagging efficiency carries an uncertainty of 10%, and the b-mis-tag rate is accurate to 30% [51]. Uncertainties that contribute to mass spectrum shape variations include the tau (3%), muon (1%), and electron (1% in the barrel region, 2.5% in the endcap region) energy scales. The effect of the uncertainty on the $$E_{T}^{miss}$$ scale, mainly due to pile-up effects, is incorporated by varying the mass spectrum shape as described in the next section.

### Table 3

Numbers of expected and observed events in the event categories as described in the text for the $$E_{T}^{miss} + X$$ channel. Also given are the expected signal yields and efficiencies for an MSSM Higgs boson with $$m_{h} = 120$$ GeV and tanβ = 10, and for an SM Higgs boson with $$m_{h} = 120$$ GeV. Combined statistical and systematic uncertainties on each estimate are reported. For the yield estimates for the Higgs signal the production cross sections for h and A, which have almost degenerate masses, are taken into account. The quoted efficiencies do not include the branching fraction into $$ττ$$.

<table>
<thead>
<tr>
<th>Process</th>
<th>SM</th>
<th>MSSM</th>
</tr>
</thead>
<tbody>
<tr>
<td></td>
<td>0/1-Jet</td>
<td>Boosted</td>
</tr>
<tr>
<td>Z → ττ</td>
<td>11 787 ± 790</td>
<td>98 ± 11</td>
</tr>
<tr>
<td>Multijet and W + jets</td>
<td>483 ± 145</td>
<td>9 ± 3</td>
</tr>
<tr>
<td>tt</td>
<td>427 ± 41</td>
<td>70 ± 8</td>
</tr>
<tr>
<td>Dibosons</td>
<td>570 ± 91</td>
<td>21 ± 4</td>
</tr>
<tr>
<td>Total background</td>
<td>13 267 ± 809</td>
<td>197 ± 14</td>
</tr>
<tr>
<td>H → ττ</td>
<td>36 ± 6</td>
<td>1.0 ± 0.3</td>
</tr>
<tr>
<td>Data</td>
<td>13 152</td>
<td>189</td>
</tr>
</tbody>
</table>

### Signal efficiency

<table>
<thead>
<tr>
<th>Process</th>
<th>SM</th>
<th>MSSM</th>
</tr>
</thead>
<tbody>
<tr>
<td>gg → φ</td>
<td>–</td>
<td>–</td>
</tr>
<tr>
<td>bb → bbb</td>
<td>–</td>
<td>–</td>
</tr>
<tr>
<td>gg → H</td>
<td>6.3 · 10⁻³</td>
<td>1.8 · 10⁻⁴</td>
</tr>
<tr>
<td>qq → qqq</td>
<td>3.0 · 10⁻³</td>
<td>8.1 · 10⁻⁴</td>
</tr>
<tr>
<td>qq → Ht or VH</td>
<td>3.8 · 10⁻³</td>
<td>6.8 · 10⁻⁴</td>
</tr>
</tbody>
</table>
The various production cross sections and branching fractions for SM and MSSM Higgs bosons and corresponding uncertainties are taken from [52–77]. Theoretical uncertainties on the Higgs production cross section are included in the SM and the MSSM search. For the SM signal, these uncertainties range from 12 to 30% for gluon fusion, depending on the event category, and 10% for VBF production. The uncertainty for the MSSM signal depends on \( \tan \beta \) and \( m_A \) and ranges from 20 to 25%.

6. Maximum likelihood fit

To search for the presence of a Higgs boson signal in the selected events, we perform a binned maximum likelihood fit to the tau-pair invariant-mass spectrum, \( m_{\tau\tau} \). The fit is performed jointly across the three SM and two MSSM event categories, but independently in the two cases.

Systematic uncertainties are represented by nuisance parameters in the fitting process. We assume log-normal priors for normalization parameters, and Gaussian priors for mass-spectrum shape uncertainties. The uncertainties that affect the shape of the mass spectrum, mainly those corresponding to the energy scales, are represented by nuisance parameters whose variation results in a continuous perturbation of the spectrum shape [78].

7. Results

Figs. 1 and 2 show for the SM and MSSM, respectively, the distributions of the tau-pair mass \( m_{\tau\tau} \) summed over the three search channels, for each category, compared with the background prediction. The background mass distributions show the results of the fit using the background-only hypothesis.

The invariant mass spectra for both the MSSM and SM categories show no evidence for the presence of a Higgs boson signal, and we therefore set 95% CL upper bounds on the Higgs boson cross section times the branching fraction into a tau pair. For calculations of exclusion limits, we use the modified frequentist construction CL [79–81]. Theoretical uncertainties on the Higgs boson production cross sections are taken into account as systematic uncertainties in the limit calculations.

7.1. Limits on MSSM Higgs boson production

For the \( m_{h_b} \) benchmark scenario as described above we set a 95% CL upper limit on \( \tan \beta \) as a function of the pseudoscalar Higgs boson mass \( m_{h_b} \) from the observed di-tau mass distributions in the b-Tag and non-b-Tag event categories (see Table 4). Signal contributions from \( h, H \) and \( A \) production are considered. The mass values of \( h \) and \( H \), as well as the ratio between the gluon fusion process and the associated production with b quarks, depend on the value of \( \tan \beta \). To account for this, we perform a scan of \( \tan \beta \) for each mass hypothesis, using the Higgs boson cross sections as a function of \( m_{h_b} \) as reported by the LHC Cross Section Working Group [52].

For the gluon-fusion process these cross sections have been obtained from the GGH@NNLO [56–83] and HIGLU [84] programs. For the \( b \to \tau \bar{\tau} \) process, the four-flavor calculation [85,86] and the five-flavor calculation as implemented in the BBH@NNLO [87] program have been combined using the Santander scheme [88]. Rescaling of the corresponding Yukawa couplings by the MSSM factors calculated with FeynHiggs [89–91] has been applied.

Fig. 3 also shows the region excluded by the LEP experiments [22]. The results reported in this Letter considerably extend the exclusion region of the MSSM parameter space and supersede limits reported by CMS using a smaller data sample collected in 2010 [26].

Fig. 1. Distribution of the tau-pair invariant mass, \( m_{\tau\tau} \), in the MSSM Higgs boson search categories: Non-b-Tag category (top), b-Tag category (bottom). The background labeled ‘electroweak’ combines the contribution from W + jets, Z \( \rightarrow \) \( \ell \ell \), and diboson processes.

7.2. Limits on SM Higgs boson production

The 0/1-Jet, VBF and Boosted categories are used to set a 95% CL upper limit on the product of the Higgs boson production cross section and the \( H \to \tau \tau \) branching fraction, \( \sigma_H \times \text{BR}(H \to \tau \tau) \), with respect to the SM Higgs expectation, \( \sigma/\sigma_{\text{SM}} \). Fig. 4 shows the observed and the mean expected 95% CL upper limits for Higgs boson mass hypotheses ranging from 110 to 145 GeV. The bands represent the one- and two-standard-deviation probability intervals around the expected limit. Table 5 shows the results for selected mass values. We set a 95% upper limit on \( \sigma/\sigma_{\text{SM}} \) in the range of 3–7.

8. Summary

We have reported a search for SM and neutral MSSM Higgs bosons, using a sample of CMS data from proton–proton collisions at a center-of-mass energy of 7 TeV at the LHC, corresponding to an integrated luminosity of 4.6 fb\(^{-1}\). The tau-pair decay mode in final states with one e or \( \mu \) plus a hadronic decay of a tau, and...
Table 4

Expected range and observed 95% CL upper limits for $\tan\beta$ as a function of $m_A$, for the MSSM search.

<table>
<thead>
<tr>
<th>MSSM Higgs $m_A$ [GeV]</th>
<th>Expected tan$\beta$ limit</th>
<th>Obs. tan$\beta$ limit</th>
</tr>
</thead>
<tbody>
<tr>
<td></td>
<td>$-2\sigma$</td>
<td>$-1\sigma$</td>
</tr>
<tr>
<td>90</td>
<td>5.19</td>
<td>7.01</td>
</tr>
<tr>
<td>100</td>
<td>6.49</td>
<td>7.45</td>
</tr>
<tr>
<td>120</td>
<td>4.50</td>
<td>6.47</td>
</tr>
<tr>
<td>130</td>
<td>5.37</td>
<td>6.71</td>
</tr>
<tr>
<td>140</td>
<td>5.62</td>
<td>6.63</td>
</tr>
<tr>
<td>160</td>
<td>5.57</td>
<td>6.99</td>
</tr>
<tr>
<td>180</td>
<td>6.75</td>
<td>8.14</td>
</tr>
<tr>
<td>200</td>
<td>7.84</td>
<td>9.12</td>
</tr>
<tr>
<td>250</td>
<td>10.3</td>
<td>12.3</td>
</tr>
<tr>
<td>300</td>
<td>13.5</td>
<td>15.7</td>
</tr>
<tr>
<td>350</td>
<td>17.7</td>
<td>20.1</td>
</tr>
<tr>
<td>400</td>
<td>21.9</td>
<td>24.3</td>
</tr>
<tr>
<td>450</td>
<td>25.0</td>
<td>29.2</td>
</tr>
<tr>
<td>500</td>
<td>30.3</td>
<td>35.7</td>
</tr>
</tbody>
</table>

The $e\mu$ final state are used. The observed tau-pair mass spectra reveal no evidence for neutral Higgs boson production. In the SM case we determine a 95% CL upper limit in the mass range of 110–145 GeV on the Higgs boson production cross section. We exclude a Higgs boson with $m_A = 115$ GeV with a production cross section 3.2 times of that predicted by the standard model. In the MSSM

Fig. 2. Distribution of the tau-pair invariant mass, $m_{\tau\tau}$, in the SM Higgs boson search categories: 0/1-Jet (top row, linear and log vertical scale), VBF (lower left), and Boosted (lower right). The background labeled ‘electroweak’ combines the contribution from $W+\text{jets}$, $Z\rightarrow\ell\ell$, and diboson processes.
Table 5

Expected range and observed 95% CL upper limits on the cross section, divided by the expected SM Higgs cross section as a function of $m_A$, for the SM search.

<table>
<thead>
<tr>
<th>$m_A$ [GeV]</th>
<th>Expected limit</th>
<th>Obs. limit</th>
</tr>
</thead>
<tbody>
<tr>
<td></td>
<td>$-2\sigma$</td>
<td>$-1\sigma$</td>
</tr>
<tr>
<td>110</td>
<td>1.83</td>
<td>2.36</td>
</tr>
<tr>
<td>115</td>
<td>1.61</td>
<td>2.13</td>
</tr>
<tr>
<td>120</td>
<td>1.65</td>
<td>2.17</td>
</tr>
<tr>
<td>125</td>
<td>1.75</td>
<td>2.19</td>
</tr>
<tr>
<td>130</td>
<td>1.82</td>
<td>2.37</td>
</tr>
<tr>
<td>135</td>
<td>2.25</td>
<td>2.96</td>
</tr>
<tr>
<td>140</td>
<td>2.39</td>
<td>2.99</td>
</tr>
<tr>
<td>145</td>
<td>3.06</td>
<td>3.97</td>
</tr>
</tbody>
</table>

Fig. 3. Region in the parameter space of $\tan\beta$ versus $m_A$ excluded at 95% CL in the context of the MSSM $m_{\text{max}}$ scenario. The expected one- and two-standard-deviation ranges and the observed 95% CL upper limits are shown together with the observed excluded region.

Fig. 4. The expected one- and two-standard-deviation ranges are shown together with the observed 95% CL upper limits on the cross section, normalized to the SM expectation for Higgs boson production, as a function of $m_A$.

case, we determine a 95% CL upper bound on the value of $\tan\beta$ as a function of $m_A$, for the $m_{\text{max}}$ scenario. This search excludes a previously unexplored region reaching as low as $\tan\beta = 7.1$ at $m_A = 160$ GeV.

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References

N. Beliy, T. Caebergs, E. Daubie

Université de Mons, Mons, Belgium


Centro Brasileiro de Pesquisas Fisicas, Rio de Janeiro, Brazil


Universidade do Estado do Rio de Janeiro, Rio de Janeiro, Brazil


Instituto de Física Teórica, Universidade Estadual Paulista, São Paulo, Brazil

V. Genchev, P. Iaydjiev, S. Piperov, M. Rodozov, S. Stoykova, G. Sultanov, V. Tcholakov, R. Trayanov, M. Vutova

Institute for Nuclear Research and Nuclear Energy, Sofia, Bulgaria

A. Dimitrov, R. Hadjiiska, A. Karadzhinova, V. Kozhuharov, L. Litov, B. Pavlov, P. Petkov

University of Sofia, Sofia, Bulgaria


Institute of High Energy Physics, Beijing, China


State Key Lab. of Nucl. Phys. and Tech., Peking University, Beijing, China

A. Cabrera, B. Gomez Moreno, A.F. Osorio Oliveros, J.C. Sanabria

Universidad de Los Andes, Bogota, Colombia

N. Godinovic, D. Lelas, R. Plestina, D. Polic, I. Puljak

Technical University of Split, Split, Croatia

Z. Antunovic, M. Dzelalija, M. Kovac

University of Split, Split, Croatia

V. Brigljevic, S. Duric, K. Kadija, J. Luetic, S. Morovic

Institute Rudjer Boskovic, Zagreb, Croatia

A. Attikis, M. Galanti, J. Mousa, C. Nicolaou, F. Ptochos, P.A. Razis

University of Cyprus, Nicosia, Cyprus

M. Finger, M. Finger Jr.

Charles University, Prague, Czech Republic

Y. Assran, A. Ellithi Kamel, S. Khalil, M.A. Mahmoud, A. Radi

Academy of Scientific Research and Technology of the Arab Republic of Egypt, Egyptian Network of High Energy Physics, Cairo, Egypt
S. Belforte $^a$, F. Cossutti $^a$, G. Della Ricca $^{a,b}$, B. Gobbo $^a$, M. Marone $^{a,b}$, D. Montanino $^{a,b,1}$, A. Penzo $^a$

$^a$ INFN Sezione di Trieste, Trieste, Italy
$^b$ Università di Trieste, Trieste, Italy

S.G. Heo, S.K. Nam

Kangwon National University, Chunchon, Republic of Korea


Kyungpook National University, Daegu, Republic of Korea

J.Y. Kim, Zero J. Kim, S. Song

Chonnam National University, Institute for Universe and Elementary Particles, Kwangju, Republic of Korea

H.Y. Jo

Konkuk University, Seoul, Republic of Korea


Korea University, Seoul, Republic of Korea

M. Choi, S. Kang, H. Kim, J.H. Kim, C. Park, I.C. Park, S. Park, G. Ryu

University of Seoul, Seoul, Republic of Korea


Sungkyunkwan University, Suwon, Republic of Korea

M.J. Bilinskas, I. Grigelionis, M. Janulis

Vilnius University, Vilnius, Lithuania


Centro de Investigacion y de Estudios Avanzados del IPN, Mexico City, Mexico

S. Carrillo Moreno, F. Vazquez Valencia

Universidad Iberoamericana, Mexico City, Mexico

H.A. Salazar Ibarguen

Benemérita Universidad Autónoma de Puebla, Puebla, Mexico

E. Casimiro Linares, A. Morelos Pineda, M.A. Reyes-Santos

Universidad Autónoma de San Luis Potosí, San Luis Potosí, Mexico

D. Krofcheck

University of Auckland, Auckland, New Zealand

A.J. Bell, P.H. Butler, R. Doesburg, S. Reucroft, H. Silverwood

University of Canterbury, Christchurch, New Zealand


National Centre for Physics, Quaid-I-Azam University, Islamabad, Pakistan

Universidad de Oviedo, Oviedo, Spain


Instituto de Física de Cantabria (IFCA), CSIC-Universidad de Cantabria, Santander, Spain


CERN, European Organization for Nuclear Research, Geneva, Switzerland


Paul Scherrer Institut, Villigen, Switzerland


Institute for Particle Physics, ETH Zurich, Zurich, Switzerland

E. Aguilo, C. Amsler, V. Chiochia, S. De Visscher, C. Favaro, M. Ivova Rikova, B. Millan Mejias, P. Otiougoua, P. Robmann, H. Snoek, M. Verzetti

Universität Zürich, Zurich, Switzerland


National Central University, Chung-Li, Taiwan


National Taiwan University (NTU), Taipei, Taiwan


Cukurova University, Adana, Turkey

University of California, Los Angeles, Los Angeles, USA


University of California, Riverside, Riverside, USA


University of California, San Diego, La Jolla, USA


University of California, Santa Barbara, Santa Barbara, USA


California Institute of Technology, Pasadena, USA


Carnegie Mellon University, Pittsburgh, USA


University of Colorado at Boulder, Boulder, USA


Cornell University, Ithaca, USA

A. Biselli, D. Winn

Fairfield University, Fairfield, USA


Fermi National Accelerator Laboratory, Batavia, USA


University of Minnesota, Minneapolis, USA

L.M. Cremaldi, R. Godang, R. Kroeger, L. Perera, R. Rahmat, D.A. Sanders, D. Summers

University of Mississippi, University, USA


University of Nebraska-Lincoln, Lincoln, USA


State University of New York at Buffalo, Buffalo, USA


Northeastern University, Boston, USA


Northwestern University, Evanston, USA


University of Notre Dame, Notre Dame, USA


The Ohio State University, Columbus, USA


Princeton University, Princeton, USA


University of Puerto Rico, Mayaguez, USA


Purdue University, West Lafayette, USA

S. Guragain, N. Parashar

Purdue University Calumet, Hammond, USA


Rice University, Houston, USA
B. Betchart, A. Bodek, Y.S. Chung, R. Covarelli, P. de Barbaro, R. Demina, Y. Eshaq, A. Garcia-Bellido, P. Goldenzweig, Y. Gotra, J. Han, A. Harel, D.C. Miner, G. Petrillo, W. Sakimoto, D. Vishnevskiy, M. Zielinski

University of Rochester, Rochester, USA

A. Bhatti, R. Ciesielski, L. Demortier, K. Goulianos, G. Lungu, S. Malik, C. Mesropian

The Rockefeller University, New York, USA


Rutgers, the State University of New Jersey, Piscataway, USA

G. Cerizza, M. Hollingsworth, S. Spanier, Z.C. Yang, A. York

University of Tennessee, Knoxville, USA


Texas A&M University, College Station, USA

N. Akchurin, J. Damgov, P.R. Dudero, C. Jeong, K. Kovitanggoon, S.W. Lee, T. Libeiro, Y. Roh, A. Sill, I. Volobouev, R. Wigmans

Texas Tech University, Lubbock, USA


Vanderbilt University, Nashville, USA


University of Virginia, Charlottesville, USA

S. Gollapinni, R. Harr, P.E. Karchin, C. Kottachchi Kankanamge Don, P. Lamichhane, M. Mattson, C. Milstène, A. Sakharov

Wayne State University, Detroit, USA


University of Wisconsin, Madison, USA

* Corresponding author.
† Deceased.
1 Also at CERN, European Organization for Nuclear Research, Geneva, Switzerland.
2 Also at National Institute of Chemical Physics and Biophysics, Tallinn, Estonia.
3 Also at Universidade Federal do ABC, Santo Andre, Brazil.
4 Also at Caltech, University of California, Pasadena, USA.
5 Also at Laboratoire Leprince-Ringuet, Ecole Polytechnique, IN2P3-CNRS, Palaiseau, France.
6 Also at Suez Canal University, Suez, Egypt.
7 Also at Cairo University, Cairo, Egypt.
8 Also at British University, Cairo, Egypt.
9 Also at Fayoum University, El-Fayoum, Egypt.
10 Now at Ain Shams University, Cairo, Egypt.