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Measurement of “pretzelosity” asymmetry of charged pion production in semi-inclusive deep inelastic scattering on a polarized $^3$He target


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Studies of nucleon structure have been and still are at the frontier of understanding how quantum chromodynamics (QCD) works in the nonperturbative region. It has been known for decades that the nucleon is composed of quarks and gluons. However, how quarks and gluons contribute to the elementary properties of the nucleon is still an open question. Among these properties, the nucleon spin has been at the center of interest for more than two decades since the European Muon Collaboration’s discovery that quark spins were found to contribute only a small portion to the nucleon spin [1]. In last two decades, polarized deep-inelastic scattering (DIS) experiments [2] have confirmed that the quark spin only contributes to about 25% of the nucleon spin with significantly improved precision. In more recent years, efforts have also been devoted to the determination of the gluon’s intrinsic contribution to the nucleon spin both from fixed-target polarized DIS and from polarized proton-proton collision measurements [3]. New results [4–6] from the Relativistic Heavy Ion Collider spin program suggest that the gluon spin may only contribute to the proton spin at a level comparable to those of quark spins. These findings suggest that the orbital angular momentum (OAM) of quarks and gluons, the most elusive piece, may actually be the largest contributor.

In recent years, major theoretical and experimental efforts have focused on accessing the OAM of quarks. The development of the general parton distribution functions (GPDs) [7] and the transverse-momentum-dependent parton distribution functions (TMDs) [8] provides not only the three-dimensional imaging of the nucleon, but also promising ways to access OAM. By investigating correlations between the quark position and the momentum, GPDs supply a new way to characterize the contribution of the quarks’ orbital motion to the spin of the nucleon. On the other hand, TMDs investigate the parton distributions in three-dimensional momentum space and provide information about the relationship between the quark momenta and the spin of either the nucleon or the quark. Most TMDs are expected to vanish in the absence of quark orbital motion. Thus they supply important and complementary (to GPDs) ways to access the OAM’s contribution to the nucleon spin.

Among the eight leading-twist TMDs, there are only three that remain nonzero after integrating over the parton transverse momentum [8]. They are the unpolarized parton distribution function (PDF) $f_1$, the longitudinally polarized PDF $g_1$ (helicity), and the transversely polarized PDF $h_1$ (transversity). The distribution $f_1$ has been extensively studied for several decades. The distribution $g_1$ is also relatively well understood by continuous efforts since the 1970s [2]. For the $h_1$, although less known than the former two, pioneering studies were made in recent years, both theoretically and experimentally [9]. One of the least known TMDs, $h_{1T}$, referred to as pretzelosity, has drawn significant attention recently [10–14] due to its intuitive relation to the quark OAM. As one of the eight leading-twist TMDs, it has a probabilistic interpretation as in a transversely polarized nucleon the parton number density of which is transversely polarized in a direction perpendicular to the nucleon polarization direction, subtracted by the parton number density with the opposite parton-polarization direction. As with transversity, pretzelosity also has an odd chirality, which leads to an important consequence that there are only quark pretzelosity distributions, with no gluonic counterparts.

In a class of relativistic quark models [13,14], pretzelosity can be expressed as the difference between the helicity and the transversity. In the light cone the difference of quark polarization between the longitudinal and transverse direction is due to the fact that boost and rotation operators do not commute. A nonzero value of the pretzelosity is thus a direct consequence of this relativistic nature of quark motion. Another interesting feature is that pretzelosity emerges from the interference of quark wave-function components differing by two units of orbital angular momentum [15]. Pretzelosity is the only leading-twist TMD carrying this unique feature. In certain models, the quark OAM can be directly accessed via pretzelosity [13,14]. This finding was first obtained in a quark-diquark model [16] and a bag model [12], and then confirmed in a large class of quark models respecting spherical symmetry [14].

In experiments, pretzelosity is suppressed in the inclusive DIS processes due to its chiral-odd nature. However,
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combined with another chiral-odd object such as the Collins fragmentation function [17], it leads to a measurable effect in the semi-inclusive DIS (SIDIS) [18] in which a leading hadron is detected in addition to the scattered lepton. Specifically, with an unpolarized lepton beam scattered from a transversely polarized nucleon target, a nonzero $h_{1T}^\perp$ would produce an azimuthal angular dependent single-spin asymmetry (SSA).

At the leading twist and following the Trento convention [19], the azimuthal angular dependence of the target SSA can be written as

$$A_{UT}(\phi_h, \phi_s) = \frac{1}{P_{He}} \frac{Y(\phi_h, \phi_s) - Y(\phi_h, \phi_s + \pi)}{Y(\phi_h, \phi_s) + Y(\phi_h, \phi_s + \pi)}$$

$$\approx A^C \sin(\phi_h + \phi_s) + A^S \sin(\phi_h - \phi_s) + A^P \sin(3\phi_h - \phi_s), \quad (1)$$

where the subscripts $U$ and $T$ stand for the unpolarized beam and the transversely polarized target, respectively. $P_{He}$ is the polarization of the target, $Y$ is the normalized yield, $\phi_h$ is the angle between the lepton plane and the hadron plane, which is defined by the hadron momentum direction and the virtual photon momentum direction, and $\phi_s$ is the angle between the target spin direction and the lepton plane. The three leading-twist asymmetries [20] correspond to the Collins asymmetry ($A^C$), the Sivers asymmetry ($A^S$), and the pretzelosity asymmetry ($A^P$). The Collins asymmetry is the transversity distribution function convoluted with the Collins fragmentation function, while the Sivers asymmetry is the Sivers distribution function convoluted with the unpolarized fragmentation function. The last term, referred to as the pretzelosity asymmetry, is the pretzelosity distribution function convoluted with the Collins fragmentation function. As shown in Eq. (1), these three terms have different azimuthal angular dependences, therefore it is possible to simultaneously determine all three terms by studying the angular dependence.

The HERMES collaboration carried out the first measurement of Collins and Sivers asymmetries [21] with electron and positron beams on a transversely polarized proton target. The COMPASS collaboration performed measurements with a muon beam on transversely polarized proton [22] and deuteron targets [23]. In Hall A at Jefferson Laboratory (JLab), an exploratory experiment E06-010 [24,25] was carried out, for the first time using an electron beam on a transversely polarized $^3$He target. The extracted Collins and Sivers asymmetries were published [24]. In extracting these asymmetries, the pretzelosity term was not included. Its uncertainty was estimated and included in the systematic uncertainties.

**II. EXPERIMENT**

In this paper, we present the results of the pretzelosity asymmetry extracted from the JLab E06-010 data. As shown in Fig. 1, a 5.9-GeV electron beam was incident on a transversely polarized gaseous $^3$He target with an average current of 12 $\mu$A. The target [26] was polarized by a spin-exchange optical pumping [27] of a Rb/K mixture, with which an average polarization is 55.4 $\pm$ 2.8%. The scattered electrons were detected using the BigBite spectrometer [26] at beam right with a solid-angle acceptance of $\sim$64 msr. Three sets of drift chambers with 18 wire planes in total were used for tracking. Lead-glass preshower and shower detectors were used to identify electrons. The hadron contamination of the electron sample in the SIDIS process was suppressed to below 2% in the momentum range of 0.6–2.5 GeV. The produced hadrons were detected in the left arm of the high resolution spectrometers [26] (LHRS). A gas Cherenkov detector and two layers of lead-glass detectors provided a clean separation of pions from electrons. An aerogel Cherenkov detector and the coincident time-of-flight technique (about 25 m from the target to the LHRS focal plane) were employed to distinguish pions from kaons and protons.

To extract moments of the SSA, it is important to have the azimuthal angular coverage as complete as possible. In the case of pretzelosity asymmetry, the azimuthal angle is $(3\phi_h - \phi_s)$ and the range is $[0, 2\pi]$. In the experiment, the BigBite and the LHRS spectrometer covered only part of the 2$\pi$ angular range. To increase the angular coverage, four different target spin orientations orthogonal to the beam direction, transverse left, transverse right, vertical up, and vertical down, were used. For each target spin orientation the spectrometers covered only a section of the phase space as shown in the left panel of Fig. 2 (target spin vertical up). However, data from all four orientations, when combined, covered the full angular range as shown in the right panel of Fig. 2, where magenta, green, red, and blue are for horizontal beam left, horizontal beam right, vertical up, and vertical down, respectively. In order to achieve target polarizations in these four orientations, three pairs of mutually orthogonal Helmholtz coils were employed. During the experiment, the target spin direction was flipped every 20 min using the adiabatic fast passage technique, in which the magnetic holding field direction and strength remained unchanged.
is defined as each panel the notation in Ref. [19].

volume) of N$_2$ gas was mixed with 3He gas to reduce depolarization effects. The nitrogen nuclei also contributed to the total measured yield and thus diluted the raw asymmetries. To obtain the asymmetries on 3He, a correction for the nitrogen dilution was applied to the raw asymmetries, as shown in Eq. (2):

$$A_{\text{raw}} = A_{\text{raw}}^{3\text{He}} \left( 1 - \frac{N_{N_2} \sigma_{N_2}}{N_{N_2} \sigma_{N_2} + N_{3\text{He}} \sigma_{3\text{He}}} \right).$$

In Eq. (2) the $\sigma$'s are the unpolarized cross sections and the $N$'s are the number densities. In the experiment, the cross section ratio $\sigma_{3\text{He}}/\sigma_{N_2}$ was measured through dedicated data taking with a reference target cell filled with a known amount of 3He and N$_2$ gases. The number densities of 3He and N$_2$ in the polarized target were verified by taking the data of electron elastic scattering on both the reference target and the production 3He target [28]. Another important correction was made due to the pair-produced background electrons (and positions) in the SIDIS electron samples. This is especially significant in the lowest $x$ bin corresponding to the lowest momentum.

Several kinematic cuts were used to select SIDIS events: the negative square of the four-momentum transfer $Q^2 > 1$ GeV$^2$, the invariant mass of the virtual photon-nucleon system $W > 2.3$ GeV, and the invariant mass of the undetected final state particles $W' > 1.6$ GeV. Data were divided into four Björken-$x$ bins with roughly equal statistics. The central kinematics are presented in Table I. To minimize the systematic uncertainties, the data taken between each of two flips of the target spin were divided into two sections. Two adjacent data sets with opposite spin directions formed a local pair, from which a local raw asymmetry was extracted. Throughout the experiment, approximately 2850 of such local raw asymmetries were combined to form the total raw asymmetry. Pretzelosity moments were extracted by a least-$\chi^2$ fit of the total raw asymmetry to Eq. (1), in a two-dimensional ($\phi_h, \phi_s$) histogram which contained 100 bins in the $2\pi$ range for each quantity.

In the polarized 3He target, a small amount (~1% in volume) of N$_2$ gas was mixed with 3He gas to reduce depolarization effects. The nitrogen nuclei also contributed to the total measured yield and thus diluted the raw asymmetries. To obtain the asymmetries on 3He, a correction for the nitrogen dilution was applied to the raw asymmetries, as shown in Eq. (2):

$$A_{\text{raw}}^{3\text{He}} = A_{\text{raw}}^{3\text{He}} \left( 1 - \frac{N_{N_2} \sigma_{N_2}}{N_{N_2} \sigma_{N_2} + N_{3\text{He}} \sigma_{3\text{He}}} \right).$$

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The extracted moments of the pretzelosity asymmetry on the 3He target are shown in the top two panels of Fig. 3 and in Table II. Only statistical uncertainties are included in the error bars. The experimental systematic uncertainties are combined in quadrature and shown as the band labeled as “sys.” All the extracted $\pi^+$ and $\pi^-$ pretzelosity terms, which were cross checked with an unbinned maximum-likelihood fit, are small and consistent with zero within the uncertainties. This observation further supports the assumption in previous.
the two upper panels, the error bars shown only represent the statistical uncertainties, while the bands labeled “sys.” represent the systematic uncertainties. Since both amplitudes of model predictions and differences between the two predictions are hardly visible compared to the statistical uncertainties, the curves in the two panels are multiplied by a factor of 10. The extracted neutron asymmetries of both \(e,e'\pi^+\) and \(e,e'\pi^-\) are again consistent with zero. Compared to the \(\sin(\phi_h + \phi_s)\) terms, the \(\sin(3\phi_h - \phi_s)\) terms are suppressed due to the different azimuthal dependent terms besides the TMDs and the Collins fragmentation functions in the convolution [20].

As suggested in [16], a large \(P_{h,T}\) coverage such as that planned for future experiments [34] with a higher statistical precision, is necessary to observe nonzero pretzelosity asymmetry. It is worth mentioning that the small value for the asymmetry predicted by the quark-diquark model (of the order of \(10^{-3}\)) is mainly due to kinematic suppression and hence does not necessarily imply that \(h_{1T}^T\) is small. In this calculation, \(h_{1T}^T\) is proportional to the OAM of the quarks, originating from a Melosh rotation of the quark spin distribution between the instant and the light-cone frame.

V. CONCLUSION

In summary, we present the first measurement of pretzelosity asymmetries on a transversely polarized \(^3\)He target, utilizing charged pion production in the semi-inclusive deep-inelastic scattering. The asymmetries are consistent with zero within experimental uncertainties, and are also consistent with model expectations. This work demonstrated an experimental approach for studying the \(h_{1T}^T\) TMD and laid a foundation for future high-precision measurements [34].

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