FIRST HARD X-RAY DETECTION OF THE NON-THERMAL EMISSION AROUND THE ARCHES CLUSTER: MORPHOLOGY AND SPECTRAL STUDIES WITH NuSTAR

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Received 2013 September 30; accepted 2013 December 7; published 2014 January 15

ABSTRACT

The Arches cluster is a young, densely packed massive star cluster in our Galaxy that shows a high level of star formation activity. The nature of the extended non-thermal X-ray emission around the cluster remains unclear. The observed bright Fe Kα line emission at 6.4 keV from material that is neutral or in a low ionization state can be produced either by X-ray photoionization or by cosmic-ray particle bombardment or both. In this paper, we report on the first detection of the extended emission around the Arches cluster above 10 keV with the NuSTAR mission, and present results on its morphology and spectrum. The spatial distribution of the hard X-ray emission is found to be consistent with the broad region around the cluster where the 6.4 keV line is observed. The interpretation of the hard X-ray emission within the context of the X-ray reflection model puts a strong constraint on the luminosity of the possible illuminating hard X-ray source. The properties of the observed emission are also in broad agreement with the low-energy cosmic-ray proton excitation scenario.

Key words: cosmic rays – Galaxy: center – ISM: general – X-rays: individual (Arches cluster)

Online-only material: color figures

1. INTRODUCTION

The Arches cluster (G0.12+0.02; Cotera et al. 1996; Serabyn et al. 1998) is a massive star cluster with a core that is about 9" (~0.35 pc at 8 kpc) in radius (Figer et al. 1999), with an average mass density of ~3 x 10^5 M☉ pc^-3, containing more than 160 O-type stars (Figer et al. 2002). The Arches cluster is located in the inner Galactic center (GC) region at the projected angular distance of 11° from the dynamic center of the Galaxy—Sagittarius A* (Sgr A*).

The first serendipitous Chandra observations of the Arches cluster region revealed two bright sources (A1 and A2) located in the cluster’s core which are surrounded by the cluster’s diffuse X-ray emission (Yusef-Zadeh et al. 2002). A1 was later resolved into two distinct bright sources: A1N and A1S by Law & Yusef-Zadeh (2004) and Wang et al. (2006). Subsequent dedicated campaigns revealed a complicated picture of the Arches cluster X-ray emission, using spectral and morphological studies to establish the presence of thermal and non-thermal emission components. The thermal emission is thought to originate from multiple collisions between strong winds of massive stars (Chlebowski & Garmany 1991; Yusef-Zadeh et al. 2002; Wang et al. 2006; Capelli et al. 2011a). Diffuse non-thermal X-ray emission has been detected from a broad region around the cluster (Wang et al. 2006; Tsujimoto et al. 2007; Capelli et al. 2011b; Tatischeff et al. 2012). The non-thermal nature of this extended radiation is revealed by its bright fluorescent Fe Kα 6.4 keV line emission. The Fe Kα line at 6.4 keV, which is due to Fe that is neutral or in a low ionization state, is produced by ejecting a K-shell electron, either by hard X-ray photoionization from an external X-ray source, or by the collisional ionization induced by cosmic-ray (CR)-accelerated electrons or ions (mainly protons and α-particles).

The mechanism producing the fluorescent line emission around the Arches cluster is not completely understood. The fluorescence may be the result of irradiation of the cloud by hard X-ray photons with energies above 7.1 keV (the K-edge of neutral iron). The required X-ray illumination might be associated with a nearby transient X-ray source (e.g., Churazov et al. 1993, for 1E1740.7-2942) or with past activity of Sgr A*, as suggested by Sunyaev et al. (1993) to explain the fluorescent line emission of the giant molecular cloud Sgr B2 in the GC region (Koyama et al. 1996; Sunyaev & Churazov 1998; Murakami et al. 2000; Revnivtsev et al. 2004a; Terrier et al. 2010). Alternatively, the excitation of cold material might be through the collisional ionization by CR particles (e.g., Predehl et al. 2003; Yusef-Zadeh et al. 2007; Dogiel et al. 2009). Tatischeff et al. (2012, hereafter T12) argued that bright 6.4 keV line structures observed around the Arches cluster are very likely produced by bombardment of molecular gas by energetic ions, and they disfavored the possibility of accelerated electrons. T12 suggest that the required large flux of low-energy cosmic ray (LECR) ions could be produced in the ongoing supersonic collision between the star cluster and a nearby molecular cloud.
2. OBSERVATIONS AND DATA REDUCTION

We observed Arches cluster during the GC region campaign with the NuSTAR (Harrison et al. 2013) in 2012 October. NuSTAR operates in wide energy band from 3 to 79 keV, carrying two identical co-aligned X-ray telescopes with an angular resolution of 18′ (FWHM). The focal planes of each telescope, referred to as focal plane module A and B (FPMA and FPMB), provide a spectral resolution of 400 eV (FWHM) at 10 keV.

The first part of the survey was carried out in the inner GC region covering the area between Sgr A∗ and the persistent low-mass X-ray binary 1E 1743.1–2843 located 19′ away (Porquet et al. 2003b; Del Santo et al. 2006). The Arches cluster serendipitously appears in the three NuSTAR observations listed in Table 1.

We processed the data from both modules using the NuSTAR Data Analysis Software (NuSTARDAS) v1.2.0, the 2013 June 28 version of the NuSTAR Calibration Database (CALDB), and HEASOFT v6.13. The data were filtered for periods of Earth occultation, high instrumental background due to South Atlantic Anomaly passages, and known bad/noisy detector pixels.

The NuSTAR detectors are not entirely shielded from X-rays that pass outside the X-ray optics modules and fall directly onto the detectors. Such an unfocused flux or “stray-light” can be significant if there are bright X-ray sources within 2–3 deg of the NuSTAR field of view. Stray-light from a point source has a uniform distribution on the detector and can be easily removed based on a geometrical model of the telescopes. We masked out the stray-light region of the FPMB detector (FPMA is clean) caused by two bright sources: GX 3+1 and the transient Swift J174510.8-262411 (see, e.g., Vovk et al. 2012 and references herein), which was active at the time of the NuSTAR observations. The contaminated detector pixels have been flagged as “bad” and removed from the analysis at the stage of data screening in the NuSTARDAS pipeline run. The observation 40010003001 (FPMB) has been completely disregarded because stray-light completely covered the Arches cluster region.

Stray-light photons from the extended sources and known celestial X-ray backgrounds cannot be so easily removed because they fill the whole detector area and have a non-uniform pattern depending on the emission morphology. We treated this background differently for the imaging and spectral analysis.

Regarding the former, the non-uniform background of the individual observations is effectively averaged when it is combined into the mosaic, and therefore, can be considered as nearly flat (however, see the discussion in Section 3.1). In contrast, the non-uniformity of the background has a noticeable effect on the spectral analysis since it is done for individual observations. To overcome this issue we applied the approach described in Section 4.

After running the pipeline and screening the data, we end up with event lists for five NuSTAR observations containing the Arches cluster at different off-axis angles (Table 1). Note that the source appears at large off-axis distances, which causes the reduction of the efficiency (vignetting) and Point Spread Function (PSF) distortion. The total nominal (effective) dead-time corrected exposure time of 122.5 (76.7) ks comprises 73.8 (44.7) ks and 48.7 (32.1) ks for FPMA and FPMB, respectively.

3. MORPHOLOGY

The data from both detector modules were used to construct the mosaic images (Figure 1) around the Arches cluster in different energy bands. We used the 3–10 keV NuSTAR band as a proxy for the standard 2–10 keV energy range; 6–7 keV band as containing Fe Kα emission lines at 6.4 and 6.7 keV, which are unresolved by NuSTAR; and 10–20 keV as a hard X-ray band specific to NuSTAR. The mosaics in the energy bands 3–79, 20–40, and 40–79 keV are also shown in Figure 1. For each image an average (local) background level was determined within an annular region centered on the Arches cluster, as defined in Table 2 and illustrated in Figure 2. The background estimate was subsequently subtracted from each individual image, prior to the construction of the image mosaic. This is done because FPMB has a higher background count rate than FPMA, requiring that we adjust the relative detector background normalizations between modules. Note that the visible gradient in Figure 2 is produced mainly by the emission from the GC molecular clouds (see, e.g., Ponti et al. 2010).

The exposure map, utilized for mosaic creation, has not been corrected for vignetting effects. We checked that the vignetting correction does not significantly improve the results of the imaging analysis. Hereafter, we assume no vignetting correction unless otherwise stated.

The primary goal of the current paper is to study the non-thermal emission surrounding the Arches cluster. The emission is revealed by fluorescent Fe Kα 6.4 keV line, as has been observed by XMM-Newton and reported in the recent papers by Capelli et al. (2011a) and T12. We adopted T12 sky regions describing the core of the Arches cluster prominent

### Table 1

Summary of the NuSTAR Observations Available for the Arches Cluster

<table>
<thead>
<tr>
<th>Observation ID</th>
<th>Date</th>
<th>Exposure(^a) (ks)</th>
<th>Off-axis(^b) (arcmin)</th>
</tr>
</thead>
<tbody>
<tr>
<td>40010003001</td>
<td>Oct 14</td>
<td>24.2 [17.0]</td>
<td>3.7</td>
</tr>
<tr>
<td>40010005001</td>
<td>Oct 15</td>
<td>26.0 [15.7]</td>
<td>4.1</td>
</tr>
<tr>
<td>40010006001</td>
<td>Oct 16</td>
<td>23.6 [12.0]</td>
<td>7.5</td>
</tr>
</tbody>
</table>

Notes.

\(^a\) Contains dead-time corrected nominal exposure, whereas the effective exposure, estimated for vignetting at 8 keV, is listed in parentheses.

\(^b\) The off-axis angle is calculated for the Arches cluster position with respect to the optical axis position on the detector.

\(\alpha\) is revealed by fluorescent Fe K\(\alpha\) line, as has been... X-ray emissions surrounding the Arches cluster. The emission is revealed by fluorescent Fe K\(\alpha\) 6.4 keV line, as has been observed by XMM-Newton and reported in the recent papers by Capelli et al. (2011a) and T12. We adopted T12 sky regions describing the core of the Arches cluster prominent...
Figure 1. NuSTAR mosaic images of the Arches cluster in different energy bands. The images are background-subtracted, and adaptively smoothed with task dmimgadapt from CIAO-4.5 using a tophat kernel. The region used for local background estimation is shown in Figure 2. The cyan circle (in black on the two upper images) denotes the Arches cluster emission prominent at 6.7 keV (a He-like Fe Kα line), and this region is denoted as “Cluster” in Table 2. The green ellipse represents the surrounding region characterized by the bright 6.4 keV line emission (“Cloud” region in Table 2). Both regions have been adopted from T12 for reference. The images are shown in the same log-scale color map shown on the right. The color bar is in units of counts s⁻¹ pixel⁻¹, and the angular size of the image pixel is 2′. The white grid indicates Galactic coordinates in degrees. The compass sign shows the image alignment in equatorial coordinates, with north up and east to the left. The image orientation remains unchanged throughout the paper.

(A color version of this figure is available in the online journal.)

Table 2
Definitions of the Sky Regions

<table>
<thead>
<tr>
<th></th>
<th></th>
<th></th>
</tr>
</thead>
<tbody>
<tr>
<td>17°45′50.3″</td>
<td>−28°49′19″</td>
<td>15″ (cluster circle)</td>
</tr>
<tr>
<td>17°45′51″</td>
<td>−28°49′16″</td>
<td>25°, 59°, 155° (Cloud ellipse)</td>
</tr>
<tr>
<td>17°45′50″</td>
<td>−28°49′19″</td>
<td>15″ (excl.)</td>
</tr>
</tbody>
</table>

Notes. Central position and radius for circular regions, and semi-minor/major axes and rotation angle for the elliptical regions. The rotation is defined counter clockwise relative to North (upward). “Cluster” and “Cloud” regions are illustrated in Figure 1. The excluded regions are marked as “excl.”

Figure 2. Definition of the local background region around the Arches cluster. The image shows the NuSTAR mosaic of the Arches cluster in the 3–10 keV energy band (see Figure 1 for the reference). The yellow annulus denotes the local background region used in the mosaic build up and fitting procedure (Section 3.1). The parameters of the annulus are listed in Table 2.

(A color version of this figure is available in the online journal.)

3.1. Cluster Emission

The Arches cluster has been relatively well studied in the standard 2–10 keV energy band with the Chandra and XMM-Newton observatories. Close attention has been paid to the central part of the cluster, where thermal emission, thought to be from colliding winds of giant stars, is observed. With data from NuSTAR, we can obtain new insight into the hard X-ray emission of the Arches cluster. However, even with angular resolution that
Given the spatial extent of the emission around the Arches Cluster as found by Chandra and XMM-Newton, we want to confirm whether we can resolve this same emission with NuSTAR. This will give us confidence that we may also be able to spatially resolve the Arches cluster above 10 keV.

We represented the Arches cluster core emission by a 2D Gaussian. Due to the low count rate per pixel, we used Poisson statistics in the fitting procedure. The background estimation, depending on the adopted energy used, or inside of the inner radius of the annular region used for the background estimation, depending on the adopted energy used. Using this fact, we estimated the background count rate in the background term has been constrained in the annulus region shown in Figure 2. After running the fitting procedure with the 2D Gaussian model, we found that the 3–10 keV emission from the Arches cluster was not acceptably fitted with a single point source, and an extended “halo” around the cluster had to be considered. This finding is in line with the detection of extended emission surrounding the Arches cluster as observed by Chandra and XMM-Newton (Yusef-Zadeh et al. 2002; Law & Yusef-Zadeh 2004; Wang et al. 2006).

We constrained the fitting area to a circle centered on the core and extending to a radius of 60′′. The model includes (1) a narrow 2D Gaussian of fixed width (4′′ FWHM, chosen to emulate the PSF smearing effect caused by incomplete aspect reconstruction), representing the cluster’s core, (2) a wide 2D Gaussian describing the halo, and (3) a constant background term. Note that the visible in Figure 2 background gradient has positional uncertainty is less than 2′′ and 6′′–8′′ for the core and halo components, respectively; see 2D confidence regions in Figure 6 for details. Model normalization factor is marked as “Norm.”

not acceptably fitted with a single point source, and an extended “halo” around the cluster had to be considered. This finding is in line with the detection of extended emission surrounding the Arches cluster as observed by Chandra and XMM-Newton (Yusef-Zadeh et al. 2002; Law & Yusef-Zadeh 2004; Wang et al. 2006).

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To demonstrate a surface brightness distribution of the Arches cluster halo, we show in Figure 5 the residuals after subtracting the radius of 60′′ around the Arches cluster (see below). Note that 90% of the averaged PSF encircled energy is situated at the centroid position of the core, we fitted a spatial model convolved with an averaged PSF on the combined data set using the Sherpa package (Freeman et al. 2001), which is part of the CIAO-4.5 software (Fruscione et al. 2006).

To make a PSF function representing the average over the different source offsets in the ObsIDs, we used the inflight-calibrated PSF shape stored in the CALDB. The PSF images were weighted according to the exposure time of each data set. The resulting averaged PSF is shown in Figure 3. Its shape is strongly distorted because of the large off-axis angles (on-axis PSF would be symmetric). Importantly, the PSF elongation is not co-aligned with the 6.4 keV line cloud region. For the fitting procedure, we assume no significant change of PSF shape inside the radius of 60′′ around the Arches cluster (see below). Note that 90% of the averaged PSF encircled energy is situated at or inside of the inner radius of the annular region used for the background estimation, depending on the adopted energy used. Using this fact, we estimated the background count rate in the separate fitting procedure without PSF convolution.

Given the spatial extent of the emission around the Arches Cluster as found by Chandra and XMM-Newton, we want to confirm whether we can resolve this same emission with NuSTAR. This will give us confidence that we may also be able to spatially resolve the Arches cluster above 10 keV.

We represented the Arches cluster core emission by a 2D Gaussian. Due to the low count rate per pixel, we used Poisson statistics in the fitting procedure. The background term has been constrained in the annulus region shown in Figure 2. After running the fitting procedure with the 2D Gaussian model, we found that the 3–10 keV emission from the Arches cluster was not acceptably fitted with a single point source, and an extended “halo” around the cluster had to be considered. This finding is in line with the detection of extended emission surrounding the Arches cluster as observed by Chandra and XMM-Newton (Yusef-Zadeh et al. 2002; Law & Yusef-Zadeh 2004; Wang et al. 2006).

We constrained the fitting area to a circle centered on the core and extending to a radius of 60′′. The model includes (1) a narrow 2D Gaussian of fixed width (4′′ FWHM, chosen to emulate the PSF smearing effect caused by incomplete aspect reconstruction), representing the cluster’s core, (2) a wide 2D Gaussian describing the halo, and (3) a constant background term. Note that the visible in Figure 2 background gradient has no significant effect on the fitting procedure. The best-fit model parameters are listed in Table 3. Figure 4 shows the radial profile of the Arches cluster intensity distribution centered at the best-fit position of the core.

To demonstrate a surface brightness distribution of the Arches cluster halo, we show in Figure 5 the residuals after subtracting off a best-fit point source corresponding to the core emission.

<Figure 3. Average NuSTAR PSF in the 3–10 keV energy range. The blue contours denote PSF encircled energy levels for 10%, 30%, 50%, 80%, and 90%, respectively, going from the center outwards. The PSF was averaged over five observations listed in Table 1. The green ellipse and black circle represent Arches cluster emission regions; see Figure 1 for details. (A color version of this figure is available in the online journal.)>

<Table 3: Best-fit Model Parameters for the PSF Fitting Procedure of the 3–10 keV Arches Cluster Core and Halo Emissions>

<table>
<thead>
<tr>
<th>Parameter</th>
<th>Value</th>
</tr>
</thead>
<tbody>
<tr>
<td>Background (constant)</td>
<td>1.18±0.01 counts s⁻¹ pixel⁻¹</td>
</tr>
<tr>
<td>Core (2D Gaussian)</td>
<td></td>
</tr>
<tr>
<td>Center R.A.</td>
<td>17°50′50″.43</td>
</tr>
<tr>
<td>Center Decl.</td>
<td>−28°49′23″.07</td>
</tr>
<tr>
<td>FWHM</td>
<td>4′′ (fixed)</td>
</tr>
<tr>
<td>Norm. (×10⁻⁵)</td>
<td>2.2±0.3 counts s⁻¹ pixel⁻¹</td>
</tr>
<tr>
<td>Halo (2D Gaussian)</td>
<td></td>
</tr>
<tr>
<td>Center R.A.</td>
<td>17°45′50″.62</td>
</tr>
<tr>
<td>Center Decl.</td>
<td>−28°49′47″.17</td>
</tr>
<tr>
<td>FWHM</td>
<td>72″ ± 46″</td>
</tr>
<tr>
<td>Norm. (×10⁻⁵)</td>
<td>1.77±0.45 counts s⁻¹ pixel⁻¹</td>
</tr>
<tr>
<td>Cash statistic</td>
<td>4661.91</td>
</tr>
<tr>
<td>Degrees of freedom</td>
<td>2820</td>
</tr>
</tbody>
</table>

Note. a J2000 epoch, the 1σ positional uncertainty is less than 2′′ and 6′′–8′′ for the core and halo components, respectively; see 2D confidence regions in Figure 6 for details. Model normalization factor is marked as “Norm.”
The position of the cluster’s core is spatially consistent with the brightest source, A1, detected by *Chandra* in the Arches cluster (Yusef-Zadeh et al. 2002) and later resolved into the two sources A1S/N (Law & Yusef-Zadeh 2004; Wang et al. 2006). Figure 6 shows confidence regions for the position uncertainty of the core and halo components.

Note that due to the *NuSTAR* optics design, the celestial coordinates of each incoming photon are determined during the post-processing. This task is complicated by distortions due to thermal bending and external forces acting on the mast during orbit. The systematic offset of the source position determination can be as high as $10''$. However, the fact that the position of the Arches cluster core emission detected with *NuSTAR* is in agreement (within the angular separation of $\sim 1''.5$) with the catalogued coordinates of the two bright *Chandra* sources A1S/N, tells us that the described systematic offsets are negligible for these observations.

Our ability to centroid a weak source with *NuSTAR* is also subject to uncertainties based on how positions are probabilistically assigned to photons in the *NuSTARDAS* pipeline and the finite spatial resolution of the *NuSTAR* pixels. To estimate any scatter introduced by this process, we ran the *NuSTARDAS* pipeline 50 times, repeated out analyses, and found that the sky position of the Gaussian core has an rms scatter of $0''.9$. We adopt this as the systematic error on the cluster core position.

The width of the 2D Gaussian halo model component does not have strong physical meaning since this component is used specifically to describe local background of the core. However, we should note that the center position of the halo component does not coincide with the core centroid, and the trend suggests that the halo centroid falls outside the Arches cluster (Figure 6).

### 3.2. Wavelet Decomposition

We studied in Section 3.1 the morphology of the Arches cluster in the $3–10$ keV band fitting data with a parametric model. We constrained the central part of the cluster, determined its centroid position, flux, and revealed strong excess outside the cluster core. Since we observe extended emission with a complicated projected morphology, it becomes quite difficult to represent this emission with an analytical model, i.e., to parameterize it. In this case, non-parametric methods can be used for morphological studies. One such method is a wavelet image transform.

For the decomposition of *NuSTAR* images in different energy bands, we used the so called *à trous* digital wavelet transform algorithm (Starck & Murtaugh 1994; Slezak et al. 1994; Vikhlinin et al. 1997; Krivonos et al. 2010). The main advantage of this algorithm over other wavelet methods is the fact that it decomposes the image into a linear combination of layers or “scales,” each of them emphasizes the structures with a characteristic size of $\approx 2^J-1$ pixels, which is $2^J$ arcsec in the case of our *NuSTAR* mosaic image resolution (1 pixel $= 2''$). Thus, low (high) $J$ values correspond to small (large) spatial variations. We used the *WVDDECOMP* wavelet decomposition method for the morphological analysis of the Arches cluster.
Figure 7. Wavelet decomposition of NuSTAR mosaic images of the Arches cluster in 3–10, 6–7, 10–20, and 20–40 keV energy bands. All structures with detection significance below 5σ have been removed from the images. The units of the color bar on the right are counts s⁻¹ pixel⁻¹, where the pixel size is 2″.

(A color version of this figure is available in the online journal.)

Figure 8. Stacked wavelet layers (see Figure 7) of the Arches cluster NuSTAR mosaic images in the (from left to right) 3–10, 6–7, and 10–20 keV energy bands. All structures with detection significance below 5σ have been removed from the images. The units of the images are counts s⁻¹ pixel⁻¹ and the pixel size is 2″.

(A color version of this figure is available in the online journal.)

code developed by Vikhlinin et al. (1995, 1998) as a part of ZHTOOLS software package.¹³

We ran WVDECOMP on mosaic images for scales 1–7 (2″–128″), detecting all significant structures above the 5σ level. To keep the photon statistics on the images, we provided counts and exposure maps to WVDECOMP. As mentioned in Section 3, FPMB has a higher background count rate than FPMA. In fact, this can mimic extended structures on the exposure-corrected image. To diminish this effect, we built the background map for a given observation by multiplying the exposure map and background count rate measured in the annulus region shown in Figure 2. The individual background maps were combined into the mosaic and used as a part of the WVDECOMP algorithm.

Figure 7 shows wavelet scales 3–6 of the Arches cluster mosaic images in 3–10, 6–7, and 10–20 keV. The lowest scales (1, 2) do not contain significantly detected (5σ) structures. The highest scale (7) contains large-scale diffuse background attributed to the Galactic center molecular complex zone (Ponti et al. 2010) and Galactic Ridge X-ray Emission (GRXE; see, e.g., Revnivtsev et al. 2006). We combined scales 3–6, which only contain structures related to the Arches cluster into the mosaics shown in Figure 8. The main results from the WT decomposition can be summarized as follows (refer to Figures 7 and 8).

3–10 keV. The core of the Arches cluster is clearly visible on the scales J = 3–4 and the halo emission is revealed on scale J = 5. The apparent position of the halo is consistent with fitting results in Section 3.1 showing a shifted position with respect to the cluster core.

6–7 keV. This energy band contains two blended Fe Kα lines at 6.4 and 6.7 keV. We speculate that these lines appear at scales J = 4 (6.7 keV) and J = 5 (6.4 keV). The 6.7 keV line is

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spatially consistent with core emission, while the 6.4 keV line emission follows the cloud region as expected. These results are in full agreement with XMM-Newton observations (T12).

In contrast to the 3–10 keV band, the 6.7 keV line emission distribution does not demonstrate high-frequency detection at \( J = 3 \), which may indicate its broader spatial distribution in comparison with the 3–10 keV continuum emission of the Arches cluster. However, we should note that the 6–7 keV mosaic image is not continuum subtracted.

10–20 keV. The Arches cluster does not show any significant (>5σ) point-like structures in the core region, which is expected due to the rapid fall of the ~1.7 keV thermal emission above 10 keV. On the other hand, the hard X-ray emission is significantly detected at scales \( J = 5–6 \). Note that the observed extended structure cannot be caused by the wide wings of the PSF—as is demonstrated in Figure 3, the PSF shape is not aligned with the cloud region morphology. The observed 10–20 keV continuum emission apparently coincides with Fe K\( ^{+} \) fluorescent 6 keV line emission observed with XMM-Newton (T12). According to T12, the 6.4 keV line emission demonstrates brightening in the south-east part of the ellipse, which is also noticeable at scale \( J = 5 \).

20–40 keV. The mosaic map of the Arches cluster shown in Figure 1 does not allow to make a definitive conclusion about the source detection. However, wavelet decomposition shows weak excess at scale \( J = 6 \). This fact is implicitly confirmed by detection of the emission up to 40 keV in the spectral analysis of Section 4. The apparent hard X-ray emission is located outside the cluster and consistent with general extent of 10–20 keV.

4. SPECTROSCOPY

We extracted spectral information for the Arches cluster from the circular region with radius of 50" centered at the cluster’s core (Table 2), including the core and extended emission. The spectra have been extracted from all observations using the nuproducts task applying PSF and ghost ray corrections to the effective area Ancillary Response File.

The extraction of the corresponding background spectra is complicated by the strong CXB and GRXE stray-light illumination causing a non-uniform pattern on the detectors with a gradient of a factor of two. For that reason, the background spectrum extracted from the source-free region over- or underpredicts the background count rate at the position of the source. This difficulty (not relevant for spectra of bright sources) causes strong systematic noise for the background-subtracted spectrum of weak sources such as Arches. To overcome this, we use the fact that the stray-light background does not significantly change for sky pointings separated by only 10–20', and therefore we utilize nearby observations. Note that this method differs from that used in the imaging analysis. In the paired observation, we extract the background spectrum from the same detector region where the Arches cluster was observed. Obviously, this region should not contain any other diffuse or point-like source or stray-light from nearby bright sources. Table 4 lists the observations used for source and background spectrum extraction. For ObsID 40010006001, we had three observations which were suitable for background spectrum extractions and thus averaged together for better statistics.

Using the above-mentioned method, we have constructed one FPMB and three FPMA spectra. The latter have been further grouped together using the FTOOL task MATHPHA. The corresponding background spectra and responses have been averaged accordingly. Finally, we have grouped spectrum channels to have at least 30 counts bin\(^{-1} \) using GRPPHA tool. Figure 9 shows the FPMA and FPMB spectra of the Arches cluster.

According to the morphological studies presented in Section 3, we attribute the hard spectral component above 10 keV to the fluorescing cloud around the Arches cluster.

We first approximated the observed spectra with the collisionally ionized plasma emission model APEC representing the cluster’s core thermal emission, and a non-thermal power-law continuum with a Gaussian line at 6.4 keV, modeling the cloud emission outside the cluster (Figure 9). Due to the fact that NuSTAR has a limited spectral resolution, the centroid energy of the Gaussian line and its width have been fixed at 6.4 keV and 0.1 keV, respectively. All the emission components were subject to a line-of-sight photoelectric absorption model WABS in XSPEC. The fitting procedure does not allow us to constrain the absorption, and we fixed it at \( N_H = 9.5 \times 10^{22} \) cm\(^{-2} \) measured by T12 (also using WABS) in the core of the Arches cluster. The metallicity of the X-ray emitting plasma was not constrained either, and we fixed it to \( Z = 1.7 \) Z_⊙ as per T12. The best-fit results obtained with this model (referred to as model 1 in the discussion) are shown in Table 5. All fitted parameters in this table and in the following discussion contain uncertainty estimation at the 90% confidence level.

Model 1 provides a reliable fit to the data by constraining the temperature of the core emission, the power-law component and the intensity of the 6.4 keV line. The 1.76±0.36 keV temperature of the thermal component is in agreement with that recently measured by T12 in the cluster’s core (1.79±0.05 keV from “model 1” in T12). The power-law slope \( \Gamma = 1.62 \pm 0.31 \) is consistent with \( \Gamma = 1.6^{+0.06}_{-0.05} \) from T12, measured in the cloud region. T12 finds the slope and 6.4 keV line equivalent width (EW_{6.4 keV}) of the cloud region compatible with the LECR ion model, where collisional ionization of the cloud by CR protons produces the 6.4 keV fluorescence line. According to the XMM-Newton study of the Arches cluster by T12, the LECR electron model (LECR_e) requires an ambient Fe abundance \( \geq 3 \) times the solar value to account for the measured EW of 1.2 ± 0.2 keV, which makes the measured properties of the non-thermal cloud region emission hardly compatible with the LECR electron model. We consider only the LECR proton (LECR_p) model in the following analysis.

We fitted the spectra with the XSPEC model WABS \times (APEC + LECR_p), henceforth “model 2.” The model depends

<table>
<thead>
<tr>
<th>Source Spectrum (ObsID(^a))</th>
<th>Background Spectrum (ObsID(^a))</th>
</tr>
</thead>
<tbody>
<tr>
<td>40010003001A(^b)</td>
<td>40010001001A</td>
</tr>
<tr>
<td>40010005001A</td>
<td>40010003001A(^b)</td>
</tr>
<tr>
<td>40010006001A(^b)</td>
<td>40010001002A/B</td>
</tr>
<tr>
<td>40010003001A/B</td>
<td>40010005001A/B</td>
</tr>
</tbody>
</table>

Notes.

\(^a\) The name of the focal plane module (A or B) is attached to the end of the ObsID.

\(^b\) The corresponding FPMB data has strong stray-light contamination, see Section 2.

\(^c\) The background spectrum has been averaged over three observations.
on the LECR path length in the ambient medium, Λ, the minimum energy of the CR particles entering the cloud, \( E_{\text{min}} \), the power-law index of the CR source energy spectrum, \( s \), and the metallicity of the X-ray emission region, Z. We fixed Λ and \( E_{\text{min}} \) parameters of the LECRp model according to T12, making free slope \( s \) and normalization \( N_{\text{LECR}} \). The best-fit results obtained with model 2 are given in Table 5 and the spectral components are shown in Figure 9.

The NuSTAR-measured spectral index of the injected low-energy protons \( s = 1.65^{+0.59}_{-0.55} \) is in general agreement with the T12 value \((1.9^{+0.5}_{-0.6})\). The best-fit LECR model normalization for the fixed \( E_{\text{min}} = 10 \text{ MeV nucleon}^{-1} \) \( N_{\text{LECR}} = (8.08^{+1.82}_{-1.62}) \times 10^{-8} \text{ erg cm}^{-2} \text{ s}^{-1} \) is somewhat higher than that measured by T12 \((N_{\text{LECR}} = (5.6^{+0.7}_{-0.5}) \times 10^{-8} \text{ erg cm}^{-2} \text{ s}^{-1})\). The corresponding power injected by LECR protons lies in the range \((5–8) \times 10^{38} \text{ erg s}^{-1}\).

To test the X-ray reflection scenario of the Arches cluster non-thermal emission, we applied the self-consistent reflection model REFLIONX in XSPEC describing the reflected

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### Table 5
Spectral Analysis of the Arches Cluster X-Ray Emission.

<table>
<thead>
<tr>
<th>Parameter</th>
<th>(Unit)</th>
<th>Model 1</th>
<th>Model 2</th>
<th>Model 3</th>
</tr>
</thead>
<tbody>
<tr>
<td>( C )</td>
<td>(10^{22} \text{ cm}^{-2})</td>
<td>9.5 (fixed)</td>
<td>9.5 (fixed)</td>
<td>9.5 (fixed)</td>
</tr>
<tr>
<td>( N_{\text{H}} )</td>
<td>(10^{22} \text{ cm}^{-2})</td>
<td>1.7 (fixed)</td>
<td>1.7 (fixed)</td>
<td>1.7 (fixed)</td>
</tr>
<tr>
<td>( kT )</td>
<td>(keV)</td>
<td>1.76^{+0.36}_{-0.29}</td>
<td>1.69^{+0.31}_{-0.26}</td>
<td>1.64^{+0.37}_{-0.23}</td>
</tr>
<tr>
<td>( I_{\text{TR}} )</td>
<td>(see notes)</td>
<td>36.0^{+17.8}_{-12.9}</td>
<td>39.6^{+16.7}_{-11.9}</td>
<td>43.2^{+16.9}_{-15.0}</td>
</tr>
<tr>
<td>( E_{6.4\text{keV}} )</td>
<td>(keV)</td>
<td>6.4 (fixed)</td>
<td>...</td>
<td>...</td>
</tr>
<tr>
<td>( F_{6.4\text{keV}} )</td>
<td>(10^{-5} \text{ ph cm}^{-2} \text{ s}^{-1})</td>
<td>1.24 ± 0.5</td>
<td>...</td>
<td>...</td>
</tr>
<tr>
<td>( \Gamma )</td>
<td></td>
<td>1.62 ± 0.31</td>
<td>...</td>
<td>...</td>
</tr>
<tr>
<td>( I_{6\text{keV}} )</td>
<td>(10^{-12} \text{ erg cm}^{-2} \text{ s}^{-1})</td>
<td>1.49^{+0.26}_{-0.24}</td>
<td>...</td>
<td>...</td>
</tr>
<tr>
<td>( E_{6.4\text{keV}} )</td>
<td>(keV)</td>
<td>1.1^{+0.5}_{-0.3}</td>
<td>...</td>
<td>...</td>
</tr>
<tr>
<td>( \Lambda )</td>
<td>(H-atoms cm^{-2})</td>
<td>...</td>
<td>5 × 10^{22} (fixed)</td>
<td>...</td>
</tr>
<tr>
<td>( s )</td>
<td></td>
<td>...</td>
<td>1.65^{+0.59}_{-0.55}</td>
<td>...</td>
</tr>
<tr>
<td>( E_{\text{min}} )</td>
<td>(keV)</td>
<td>...</td>
<td>10^{4} (fixed)</td>
<td>...</td>
</tr>
<tr>
<td>( N_{\text{LECR}} )</td>
<td>(10^{-6} \text{ erg cm}^{-2} \text{ s}^{-1})</td>
<td>...</td>
<td>8.08^{+1.82}_{-1.62}</td>
<td>...</td>
</tr>
<tr>
<td>( \Gamma_{\text{RX}} )</td>
<td></td>
<td>...</td>
<td>...</td>
<td>2.93 ± 0.31</td>
</tr>
<tr>
<td>( \xi )</td>
<td>(erg cm s^{-1})</td>
<td>...</td>
<td>...</td>
<td>10 (fixed)</td>
</tr>
<tr>
<td>( I_{\text{RX}} )</td>
<td>(10^{-5})</td>
<td>...</td>
<td>...</td>
<td>6.09 ± 5.00</td>
</tr>
<tr>
<td>Reduced ( \chi^2 ) (d.o.f.)</td>
<td></td>
<td>0.91 (134)</td>
<td>0.89 (136)</td>
<td>0.91 (140)</td>
</tr>
</tbody>
</table>

**Notes.** Model 1: \( C \times \text{WABS} \times (\text{APEC} + \text{Gaussian} + \text{power law}) \); Model 2: \( C \times \text{WABS} \times (\text{APEC} + \text{LECRp}) \); Model 3: \( C \times \text{WABS} \times (\text{APEC} + \text{REFLIONX}) \). The cross-calibration constant term C is fixed at unity for FPMA data and fitted for FPMB. The absorption column density is noted as \( N_{\text{H}} \). APEC thermal plasma model is characterized by \( kT \), \( Z/Z_{\odot} \), and \( I_{\text{TR}} \) parameters, respectively, describing temperature, metallicity relative to solar, and normalization in units of \( 10^{-18} \int n_e n_H dV/(4\pi D^2) \), where \( n_e \) and \( n_H \) are the electron and proton number densities in units of \( \text{cm}^{-3} \), and \( D \) is the distance in cm. \( E_{6.4\text{keV}} \) and \( F_{6.4\text{keV}} \) : centroid energy and flux of the neutral or low-ionization Fe Kα line. \( \Gamma \) and \( I_{6\text{keV}} \) : index and 3–20 keV pegged normalization of the power-law component (model PEGPWRLW). \( E_{6.4\text{keV}} \) : EW of the 6.4 keV line with respect to the power-law continuum. \( \Lambda \), \( s \), \( E_{\text{min}} \), and \( N_{\text{LECR}} \) : LCR ion path length, source spectrum index, minimum energy, and model normalization. By definition \( dW/dt = 4\pi D^2N_{\text{LECR}} \) is the power injected in the interaction region by primary CR protons with energies between \( E_{\text{min}} \) and \( E_{\text{max}} = 1 \text{ GeV} \), where \( D \) is the distance to the source. \( \Gamma_{\text{RX}} \), \( \xi \), and \( I_{\text{RX}} \): the power-law slope of the incident radiation, ionization parameter (defined as \( \xi = 4\pi F_{\text{ion}}/n_{\text{H}} \) where \( F_{\text{ion}} \) is the total illuminating flux, and \( n_{\text{H}} \) is the density of the reflector), and normalization of REFLIONX model (Ross & Fabian 2005). The metallicity parameter \( Z \) for LCR and REFLIONX models is fixed at 1.7 \( Z_{\odot} \). Reduced \( \chi^2 \) (d.o.f.): reduced \( \chi^2 \) and degrees of freedom.

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spectrum for an optically thick atmosphere of constant density, illuminated by radiation with a power-law spectrum (Ross & Fabian 2005), henceforth “model 3.” The REFLIONX spectrum contains modeled fluorescence lines and continuum emission (Figure 9). Developed for the surface of hot accretion disks, REFLIONX can still be applied for cold material around the Arches cluster, fixing the ionization parameter $\xi$ at the lowest allowed value of 10 erg cm s$^{-1}$. The best-fit parameters of model 3 are shown in Table 5.

5. DISCUSSION

In this section, we discuss the results of this work in the context of two scenarios proposed to explain the excitation mechanism of the fluorescent emission around the Arches cluster.

5.1. CR Bombardment Hypothesis

A particle bombardment excitation scenario for the Fe fluorescence line emission involves the interaction of CR ions with cold molecular gas. T12 developed a steady-state slab model in which accelerated particles penetrate a cloud of neutral gas at a constant rate until they either stop or escape from the cloud. The model describes the properties of the neutral Fe Kα line emission excited by impacts of LECR ions. These authors argue that fluorescent 6.4 keV line emission of the Arches cluster is excited by LECR protons with a relativistic soft incident spectrum $s = 1.9^{+0.5}_{-0.5}$. Our NuSTAR spectral analysis (model 2) implies a somewhat harder slope $s = 1.65^{+0.59}_{-0.55}$, however, both measurements are consistent within the error intervals. T12 notes that production of 6.4 keV line photons by LECR collisional ionization is relatively inefficient and that high power in LECRs is needed to produce the observed fluorescent Fe Kα line flux. As the authors suggest, a large flux of LECR ions could be produced in the ongoing supersonic collision between the star cluster and nearby molecular cloud.

Note that LECRp model provides an equally good fit as the reflection model discussed below.

5.2. X-Ray Reflection Hypothesis

By definition, the excitation of the neutral material of the cloud by hard X-ray irradiation implies a transient or persistent X-ray source outside or within the cloud.

Assuming that the irradiation source is located in the Arches cluster, the source luminosity required to produce the observed 6.4 keV line flux can be estimated from the fluorescent model described in Sunyaev & Churazov (1998). Applying this approach on the XMM-Newton data, T12 showed that the luminosity of the source in the cluster must be at least $10^{36}$ erg s$^{-1}$. However, the persistent X-ray luminosity is far below this level. Capelli et al. (2011a) recently reported X-ray flaring activity in the core of the Arches cluster, probably originating from stellar wind interactions in one or more massive binary systems. Such transient activity, in principle, can trigger fluorescent emission of the neutral material outside the Arches cluster, but the peak luminosity of the observed flare is far below the threshold mentioned above.

The illuminating source could also be a transient X-ray binary system. Their typical X-ray luminosity can reach levels on the order of $10^{36}–10^{38}$ erg s$^{-1}$ for high- and low-mass systems. If the bright state of an X-ray binary lasted for more than 8 yr at or above $10^{37}$ erg s$^{-1}$, it would be sufficient for the observed 6.4 keV flux (Capelli et al. 2011b). However, such a long outburst is much longer than typical flaring timescales of X-ray transients (Degenaar & Wijnands 2010), though, at least one exceptional case is known, the neutron star X-ray transient KS 1731-260 (Sunyaev & Kwant Team 1989; Sunyaev et al. 1990) which became quiescent after more than a decade in the bright state (Wijnands et al. 2001; Revnivtsev & Sunyaev 2002).

The NuSTAR-measured EW of the 6.4 keV line with respect to the power-law continuum, $EW_{6.4keV} = 1.1^{+0.7}_{-0.5}$ keV, is in agreement with previous studies performed on spatially resolved fluorescent clouds around the Arches cluster (Capelli et al. 2011b, T12). In the context of an X-ray reflection scenario, the relatively high EW of the Fe Kα line indicates a scattered emission component (assuming normal Fe abundance) or strong photoabsorption (Sunyaev & Churazov 1998), which can explain the fact that the direct source of photoionization is not seen by the observer. However, according to NuSTAR maps of the Arches cluster above 10 keV (Figure 1), for which the efficiency of photoabsorption rapidly falls, no significant source has been detected.

The Arches cluster spectrum is well-fitted with a model composed of an optically thin thermal plasma and an X-ray reflection component (Figure 9). The spectrum of the illuminating X-ray source is estimated to have relatively soft slope $\Gamma_{RX} = 2.93^{+0.31}_{-0.31}$. We may assess the luminosity of the primary source $L_{RX}$ required by the reflection component following the method used by Reis et al. (2013) to estimate the reflection fraction $R$. This parameter is defined as a measure of the ratio of the reflected flux to the continuum power-law flux. REFLIONX calculates $R$ by dividing the extrapolated (1 eV–1000 keV) REFLIONX flux by the 100 eV–1000 keV power-law flux. On the other hand, the reflection scaling factor $R$ is in fact the solid angle $\Omega = \Omega/2\pi$ subtended by the reflecting cloud ($R = 0$ means no reflection). We estimated the unabsorbed 0.3–8 keV luminosity $L_{RX}$ (at 8 kpc) for the tabulated $R$ values—0.5, 1.0 (source above a plane), 1.5, and 2.0 (source inside the cloud) as 2.5, 1.3, 0.8 and $6 \times 10^{36}$ erg s$^{-1}$, respectively. Wang et al. (2006) reported the 0.3–8 keV luminosity of the bright Arches cluster sources A1N/S and A2, respectively, as 7.2 (A1N), 11 (A1S), and $4.6 \times 10^{35}$ erg s$^{-1}$, which is 1–2 orders of magnitude below the luminosity of the illuminating continuum required by the reflected component. This constraint is consistent with what can be drawn from the observed 6.4 keV line flux, using the general idea of Sunyaev & Churazov (1998), as demonstrated by T12.

The interpretation of the observed spectra with the X-ray reflection model confirms that the Arches cluster itself is unlikely to be the source of the cloud fluorescence emission.

The non-detection of Fe Kα line flux variability over more than 8 yr of XMM-Newton observations (Capelli et al. 2011b, T12) puts strong limits on the photoionization excitation of molecular clouds around the Arches cluster by a hard X-ray source. Nevertheless, the plausibility of the scenario with an earlier outburst of Sgr A* is under debate. The large-scale spatial distribution and time variation of the Fe-K fluorescence seen in the GC region reveal the apparent superluminal propagation of a light front illuminating a molecular nebula (Ponti et al. 2010). In this context, it is quite reasonable to anticipate X-ray photoionization of the molecular clouds near the Arches cluster by Sgr A*.” A possible decrease in the 6.4 keV flux with time would be expected in this scenario. For instance, strong variability of the 6.4 keV emission has been detected in Sgr B2 (Terrier et al. 2010) after 10 yr of a persistent state as shown by Revnivtsev et al. (2004a); however, see Inui et al. (2009).
of the Arches cluster clouds consistent with a long outburst from Sgr A∗ at a constant luminosity of \(1.4 \times 10^{39}\) erg s\(^{-1}\) (Ponti et al. 2010). On the other hand, Clavel et al. (2013) possibly revealed a highly variable active phase of Sgr A∗ during the past few hundred years.

The measured power-law slope of the illuminating source in the model 3 (Section 4) is softer (\(\Gamma \approx 3\)) than observed for Sgr A∗ flares (Porquet et al. 2003a, 2008; Nowak et al. 2012). Additionally, the slope is softer than that estimated by Revnivtsev et al. (2004b) for Sgr B2 in the context of Sgr A∗ reflected emission. However, we should note that REFLIONX calculates the reflected spectrum for an optically thick atmosphere while the molecular clouds close to the Arches cluster exhibit transparency at \(\tau < 1\) (Capelli et al. 2011b). A more realistic reflection model at lower values and different source-cloud geometry settings can produce a harder spectrum of the primary emission.

6. SUMMARY

In this paper, we present results on the first focused hard X-ray observation of the Arches cluster performed by NuSTAR. We report on significant detection of the hard X-ray emission around the cluster with a spectrum that extends for the first time up to \(\sim 40\) keV. The main results can be summarized as follows.

1. The NuSTAR mosaic image of the Arches cluster shows a core component and extended halo emission in the 3–10 keV energy band. The centroid position of the core emission coincides with the bright sources A1N/S resolved with Chandra in the Arches cluster. The halo emission around the cluster also follows the morphology observed with Chandra. The centroid position of the halo emission does not coincide with the core and tends to be outside the Arches cluster.

2. Wavelet decomposition of sky mosaics does not show any significant (5\(\sigma\)) detection of point-like sources in the core of the Arches cluster above 10 keV. The continuum 10–20 keV emission is significantly detected well beyond the Arches cluster with a spatial morphology consistent with the Fe Kα fluorescent 6.4 keV line emission observed with XMM-Newton (T12).

3. The spectral analysis of the Arches cluster region, including the extended emission around it, reveals a \(\sim 1.7\) keV temperature plasma, which is consistent with the optically thin thermal plasma measured by T12 in the cluster’s core and associated with thermalization of massive star winds and several colliding stellar wind binaries within the cluster (T12).

4. The power-law slope \(\Gamma = 1.62 \pm 0.31\) and \(EW_{6.4\,keV} = 1.1^{+0.7}_{-0.5}\) keV are in agreement with that measured by T12 around the Arches cluster and expected for the LECR ion model.

5. The X-ray emission of the Arches cluster is well fitted with a model composed of an optically thin thermal plasma and a non-thermal component produced by LECR ions (T12). The best-fit CR spectral index is \(s = 1.65^{+0.59}_{-0.55}\). The LECR model normalization gives the corresponding power of the injected LECR protons in the cloud region at the level of \((5–8) \times 10^{28}\) erg s\(^{-1}\) (with \(D = 8\) kpc and CR minimum energy \(E_{\text{min}} = 10\) MeV nucleon\(^{-1}\)).

6. The luminosity of the primary source required by the X-ray reflection model is 1–2 orders of magnitude higher than that observed in the Arches cluster, which is in agreement with similar constraints drawn from the 6.4 keV line flux.

The X-ray photoionization and CR-induced emission models can reproduce the data equally well. Additional NuSTAR observations with the Arches cluster on-axis would be useful to detect the difference between the emission models above 20 keV. At this point, we can conclude that explaining the non-thermal emission around the Arches cluster with an X-ray reflection scenario in which the illuminating source is located in or close to the cluster is unlikely. However, irradiation by a long outburst of Sgr A∗ can still be considered, and would imply long-term variability of the 6.4 keV line flux should be detected. The excitation of the neutral material around the cluster by CR interactions also looks plausible, though a large flux of LECR protons is needed to produce the observed fluorescent emission.

This work was supported under NASA Contract No. NNG08FD60C, and made use of data from the NuSTAR mission, a project led by the California Institute of Technology, managed by the Jet Propulsion Laboratory, and funded by the National Aeronautics and Space Administration. We thank the NuSTAR Operations, Software, and Calibration teams for support with the execution and analysis of these observations. This research has made use of the NuSTAR Data Analysis Software (NuSTARDAS) jointly developed by the ASI Science Data Center (ASDC, Italy) and the California Institute of Technology (USA). F.E.B. acknowledges support from Basal-CATA (PB-06/2007) and CONICYT-Chile (FONDECYT-Chile 11101024 and Anillo ACT1101). R.K. thanks Eugene Churazov for fruitful discussions and valuable suggestions to the paper.

Facility: NuSTAR

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