

## MIT Open Access Articles

The r-process Pattern of a Bright, Highly r -processenhanced Metal-poor Halo Star at [Fe/H ] ~ -2

The MIT Faculty has made this article openly available. *Please share* how this access benefits you. Your story matters.

**Citation:** Sakari, Charli M., Vinicius M. Placco, Terese Hansen, Erika M. Holmbeck, Timothy C. Beers, Anna Frebel, Ian U. Roederer, et al. "The r-Process Pattern of a Bright, Highly r-Process-Enhanced Metal-Poor Halo Star at [Fe/H] # –2." The Astrophysical Journal 854, no. 2 (February 14, 2018): L20.

**As Published:** http://dx.doi.org/10.3847/2041-8213/AAA9B4

**Publisher:** American Astronomical Society

Persistent URL: http://hdl.handle.net/1721.1/118126

**Version:** Final published version: final published article, as it appeared in a journal, conference proceedings, or other formally published context

**Terms of Use:** Article is made available in accordance with the publisher's policy and may be subject to US copyright law. Please refer to the publisher's site for terms of use.





# The r-process Pattern of a Bright, Highly r-process-enhanced Metal-poor Halo Star at $[Fe/H] \sim -2$

Charli M. Sakari , Vinicius M. Placco<sup>2,3</sup>, Terese Hansen , Erika M. Holmbeck<sup>2,3</sup>, Timothy C. Beers<sup>2,3</sup>, Anna Frebel , Ian U. Roederer<sup>3,6</sup>, Kim A. Venn, George Wallerstein, Christopher Evan Davis, Elizabeth M. Farrell, and David Yong Department of Astronomy, University of Washington, Seattle, WA 98195-1580, USA; sakaricm@u.washington.edu

Department of Physics, University of Notre Dame, Notre Dame, IN 46556, USA

Joint Institute for Nuclear Astrophysics Center for the Evolution of the Elements (JINA-CEE), USA

Observatories of the Carnegie Institution of Washington, 813 Santa Barbara Street, Pasadena, CA 91101, USA

Department of Physics and Kavli Institute for Astrophysics and Space Research, Massachusetts Institute of Technology, Cambridge, MA 02139, USA

Department of Astronomy, University of Michigan, 1085 S. University Avenue, Ann Arbor, MI 48109, USA

Department of Physics and Astronomy, University of Victoria, Victoria, BC, Canada

Research School of Astronomy and Astrophysics, Australian National University, Canberra, ACT 2611, Australia Received 2017 November 26; revised 2018 January 19; accepted 2018 January 21; published 2018 February 14

#### Abstract

A high-resolution spectroscopic analysis is presented for a new highly r-process-enhanced ([Eu/Fe] = 1.27, [Ba/Eu] = -0.65), very metal-poor ([Fe/H] = -2.09), retrograde halo star, RAVE J153830.9–180424, discovered as part of the R-Process Alliance survey. At V = 10.86, this is the brightest and most metal-rich r-II star known in the Milky Way halo. Its brightness enables high-S/N detections of a wide variety of chemical species that are mostly created by the r-process, including some infrequently detected lines from elements like Ru, Pd, Ag, Tm, Yb, Lu, Hf, and Th, with upper limits on Pb and U. This is the most complete r-process census in a very metal-poor r-II star. J1538–1804 shows no signs of s-process contamination, based on its low [Ba/Eu] and [Pb/Fe]. As with many other r-process-enhanced stars, J1538–1804's r-process pattern matches that of the Sun for elements between the first, second, and third peaks, and does not exhibit an actinide boost. Cosmo-chronometric age-dating reveals the r-process material to be quite old. This robust main r-process pattern is a necessary constraint for r-process formation scenarios (of particular interest in light of the recent neutron star merger, GW170817), and has important consequences for the origins of r-II stars. Additional r-I and r-II stars will be reported by the R-Process Alliance in the near future.

Key words: Galaxy: formation - stars: abundances - stars: atmospheres - stars: fundamental parameters

#### 1. Introduction

The very metal-poor stars ([Fe/H] < -2) are believed to be some of the oldest objects in the Milky Way (MW). These stars retain the chemical signatures of the few stars that evolved and died before them (e.g., Frebel et al. 2015). The subset of stars that are highly enhanced in the heavy elements that form via the rapid (r-) neutron-capture process are of particular interest, as their abundances trace the yields from early r-process events. The signatures of the r-process are seen throughout the Galaxy (Roederer 2013), but r-processenhanced stars enable measurements of a wide assortment of neutron-capture elements. These stars are classified according to their Eu abundances: r-I stars have  $+0.3 < [Eu/Fe] \le +1$ , while r-II stars have [Eu/Fe] > +1 (Christlieb et al. 2004). There are currently only  $\sim 30$  r-II and  $\sim 100$  r-I stars known (see the JINAbase compilation; Abohalima & Frebel 2017). Studies of these stars have found a nearly identical main r-process pattern (for Ba to Hf) in all types of stars and in all environments, with variations among the lightest and heaviest elements (e.g., Sneden et al. 1994; Roederer et al. 2014b; Siqueira Mello et al. 2014; Ji et al. 2016; Placco et al. 2017).

Neutron star mergers (NSMs) have long been suspected to be a site of the *r*-process (e.g., Lattimer & Schramm 1974; Rosswog et al. 2014; Lippuner et al. 2017). The recent detection of GW170817 (Abbott et al. 2017) and subsequent *r*-

process nucleosynthesis (e.g., Chornock et al. 2017) demonstrates that these conditions can indeed be met in NSMs. Galactic chemical evolution models (e.g., Côté et al. 2017) also suggest that NSMs can produce all the observed Eu in the MW. However, problems such as coalescence time and NSM rates still remain, prompting the question: Is there only a single site for r-process nucleosynthesis, and are the physical conditions always identical? Standard core-collapse supernovae seem to have been ruled out as the site for most of the r-process elements (though they may form the lighter elements; e.g., Arcones & Thielemann 2013), but the magneto-rotational supernovae (e.g., Winteler et al. 2012) remain another option. Observations of the detailed r-process pattern in large samples of stars will be useful for constraining the fundamental physics and site(s) behind the r-process, through determinations of the relative abundances of second- versus third-peak elements, the presence of actinide boosts (e.g., Schatz et al. 2002), and the behavior of the first-peak elements.

Also important is the rate at which *r*-process events occur, as well as where and when. The detection of *r*-process-enhanced stars in ultra-faint dwarf galaxies (Ji et al. 2016), the age-dating of *r*-process-enhanced stars (Sneden et al. 1996; Cayrel et al. 2001), trends with metallicity or other chemical abundances (Macias & Ramirez-Ruiz 2016), the relative numbers of *r*-I and *r*-II stars (Barklem et al. 2005), and the amount of *r*-process material in a given environment (e.g., Tsujimoto & Nishimura 2015; Beniamini et al. 2016) are all important for deciphering the site of the *r*-process. Obtaining high-precision, detailed abundance patterns and understanding

 $<sup>\</sup>frac{1}{9}$  Both have [Ba/Eu] < 0 to minimize contamination from the slow (s-) process.

Table 1
Target Information

Parameter	Value			
Aliases	TYC 6189-285-1			
	2MASS J15383085-1804242			
R.A. (J2000)	15:38:30.85			
Decl. (J2000)	-18:04:24.2			
V	10.86			
K	8.484			
Observation Date	2017 May 4 (MJD = 57877)			
Exposure Time (s)	1800			
S/N, 3700 Å <sup>a</sup>	287			
$v_{\text{helio}} \text{ (km s}^{-1}\text{)}$	$131.3 \pm 0.5$			
$U (\text{km s}^{-1})^{\text{b}}$	$64^{+6874}_{-104}$			
$V (\text{km s}^{-1})^{\text{b}}$	$-774^{+536}_{-4725}$			
$W (\text{km s}^{-1})^{\text{b}}$	$147^{+85}_{-8776}$			
$T_{\rm eff}$ (K)	$4752 \pm 30$			
$\log g$	$1.63 \pm 0.1$			
$\xi  (\mathrm{km}  \mathrm{s}^{-1})$	$1.51 \pm 0.15$			
[Fe/H]	$-2.09 \pm 0.02$			
[C/Fe] <sup>c</sup>	$+0.26 \pm 0.10$			
$[\alpha/\text{Fe}]^d$	$+0.34 \pm 0.03$			

#### Notes.

the r-process-enhanced stars as a stellar population in a statistical sense requires a much larger sample of r-I and r-II stars.

The R-Process Alliance is a collaboration with the aim of identifying the site(s) of the *r*-process. The first phase of the Alliance is dedicated to discovering larger samples of *r*-I and *r*-II stars in the MW through medium- and high-resolution spectroscopy (V. M. Placco et al. 2018, in preparation; T. Hansen et al. 2018, in preparation; C. M. Sakari et al. 2018, in preparation). This Letter presents the detailed *r*-process abundances of an *r*-II star, RAVE J153830.9–180424 (hereafter J1538–1804) that was discovered in the northern hemisphere sample of C. M. Sakari et al. (2018 in preparation, hereafter Paper I). Future papers will present additional *r*-I and *r*-II stars discovered by the R-Process Alliance.

#### 2. Observations, Data Reduction, and Analysis Techniques

J1538–1804 was identified as a metal-poor star in the reanalyzed data from the RAdial Velocity Experiment (RAVE) survey (Steinmetz et al. 2006) by Matijevič et al. (2017). The star was followed up at medium resolution ( $R \sim 2000$ ) in the blue (3300–5000 Å) to determine atmospheric parameters (V. M. Placco et al. in preparation), and was subsequently identified as an r-II star in the northern hemisphere, high-resolution component of the R-Process Alliance (Paper I), based on echelle spectroscopy with the Astrophysical Research Consortium (ARC) 3.5 m telescope at Apache Point Observatory (APO). The target was then followed up at higher resolution ( $R \sim 83,000$  in the blue and  $R \sim 65,000$  in the red, with the 0″35 slit and  $1 \times 1$  binning) on 2017 May 4 with the Magellan Inamori Kyocera Echelle (MIKE) spectrograph (Bernstein et al. 2003) on the Magellan-Clay

Telescope at Las Campanas Observatory. A wavelength coverage of 3200–5000 Å was obtained in the blue and 4900–10000 Å in the red. The details for the MIKE observations are listed in Table 1; the conditions were photometric, with 0.14 seeing. The spectra were reduced using the Image Reduction and Analysis Facility program (IRAF) and the MTOOLS package.

All abundances were derived with spectrum syntheses, using the 2017 version of MOOG (Sneden 1973) with an appropriate treatment of scattering (Sobeck et al. 2011). 12 The atmospheric parameters of the stars were determined by flattening trends in Fe I lines with wavelength, reduced equivalent width, and excitation potential (EP), and by forcing agreement between Fe I and Fe II abundances. For each Fe I line, a  $\langle 3D \rangle$ , non-local thermodynamic equilibrium (NLTE) correction (Amarsi et al. 2016) was applied to the LTE abundance. Paper I will demonstrate that the NLTE parameters are generally in better agreement with photometric temperatures; however, at [Fe/H]  $\sim -2$ , the differences between NLTE and LTE atmospheric parameters are generally negligible. Using LTE atmospheric parameters also has a slight impact on the derived abundances; see Table 3.

The line lists were generated with the linemake code<sup>13</sup> (C. Sneden 2018, private communication) with additions from Cowan et al. (2005), Xu et al. (2007), and Sneden et al. (2009), and include hyperfine structure, isotopic splitting, and molecular lines from CH,  $C_2$ , and CN. The atmospheric parameters and abundances of the light and iron-peak elements are provided in Paper I, though Table 1 lists the final, adopted atmospheric parameters, the [C/Fe], and the average [ $\alpha$ /Fe]. Note that this star is not C-enhanced. This star was also included in the first *Gaia* data release (Brown et al. 2016). Its proper motion, parallax, and velocity demonstrate that it is a retrograde halo star (see Table 1), even accounting for the large uncertainty in the parallax.

### 3. r-process Patterns

Abundances or upper limits are derived for 27 neutroncapture elements; sample syntheses are shown in Figure 1. Table 2 shows the line-by-line r-process abundances, while Table 3 shows the final mean abundances. When a spectral line was sufficiently weak and unblended in the Kurucz solar spectrum, <sup>14</sup> the solar abundance for that line was derived using the same atomic data (see Table 2; otherwise, the Asplund et al. 2009 solar values are used). All [X/Fe] ratios are relative to the [Fe I/H] ratios in Table 1 (note that the adoption of NLTE corrections during the atmospheric parameter determination has ensured that the Fe I and Fe II ratios are equal). The random errors are based on the quality of the data and the fitted synthetic spectra. A minimum random error of 0.05 dex was adopted, but weak or blended lines in low-S/N regions could have random errors as large as 0.1–0.2 dex. The systematic errors are due to uncertainties in the atmospheric parameters. Table 3 also provides the abundance offsets that occur when LTE parameters are adopted. Figure 2 shows the *r*-process

<sup>&</sup>lt;sup>a</sup> S/N is per resolution element.

b The large uncertainty is due to the uncertain Gaia DR1 parallax.

<sup>&</sup>lt;sup>c</sup> This is the "natal" [C/Fe], calculated with the evolutionary corrections of Placco et al. (2017). The measured [C/Fe] is -0.10.

 $<sup>^{\</sup>rm d}$  This is an average of [Mg/Fe], [Si/Fe], and [Ca/Fe] from Paper I.

<sup>10</sup> IRAF is distributed by the National Optical Astronomy Observatory, which is operated by the Association of Universities for Research in Astronomy, Inc., under cooperative agreement with the National Science Foundation.

<sup>&</sup>lt;sup>11</sup> http://www.lco.cl/telescopes-information/magellan/instruments/mike/iraf-tools/iraf-mtools-package

<sup>12</sup> https://github.com/alexji/moog17scat

https://github.com/vmplacco/linemake

<sup>14</sup> http://kurucz.harvard.edu/Sun.html

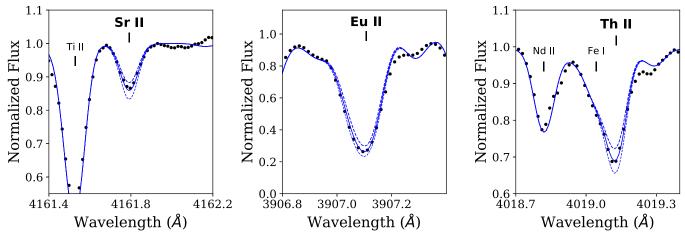


Figure 1. Best-fit syntheses to Sr, Eu, and Th lines (solid lines), along with  $\pm 0.1$  dex abundances (dashed lines).

pattern in J1538–1804, along with reference patterns in the Sun, a very metal-poor, actinide boost r-II star (Hill et al. 2002), and an extremely metal-poor r-I star (Roederer et al. 2014a). The pattern in J1538–1804 is well fit by the solar r-process pattern and does not agree with the solar s-process pattern. The various groups of elements are discussed further below.

#### 3.1. Barium and Europium

Barium and europium are the elements used for classification of r-I and r-II stars. The Ba lines are quite strong in this fairly metalrich star; only the 5853 and 6141 Å lines were used. Multiple Eu II lines were detected, including the weak line at 6645 Å . The derived subsolar value of [Ba/Eu] ( $-0.65 \pm 0.08$ ) suggests that the star has minimal contamination from the s-process despite its moderate [Fe/H], while its high [Eu/Fe] ( $+1.27 \pm 0.05$ ) makes it an r-II star.

J1538–1804 is on the metal-rich end of known r-II stars. There are only four r-II stars in JINAbase that have  $[Fe/H] \gtrsim -2.1$ , and only one has a detailed r-process pattern that covers from Sr to U and provides an age determination. Two of these four stars are associated with dwarf galaxies (Reticulum II and Ursa Minor; Ji et al. 2016; Aoki et al. 2007), while the other two are associated with the MW bulge (Howes et al. 2016). J1538–1804 is therefore the most metal-rich r-II star known in the MW halo.

#### 3.2. Lighter r-process Elements

The lighter *r*-process elements Sr, Y, and Zr were derived with 1, 4, and 6 lines, respectively (see Figure 1 for the fit to the Sr line). These elements are commonly derived in many abundance analyses. Abundances of Ru and Rh were determined from 5 and 2 lines, while Pd and Ag were derived from single lines in the noisier blue end of the spectrum. Sr, Y, and Zr are in agreement with the solar pattern and the other *r*-I and *r*-II stars. Ru and Ag are slightly low in J1538–1804 compared to the Sun; CS 31082–001 also has similarly low Ag relative to the Sun.

These lighter *r*-process elements are thought to form in the main *r*-process and in a "weak" or "limited" *r*-process (A. Frebel 2018, in preparation) that could occur in traditional core-collapse supernovae (e.g., Arcones & Thielemann 2013). The lighter *r*-process pattern may therefore vary between stars

depending on age, environment, etc. Both Siqueira Mello et al. (2014) and Placco et al. (2017) have noted first-peak variations in *r*-I and *r*-II stars. J1538–1804's abundance pattern is robust with respect to the Sun, suggesting that the relative contribution from the "limited" and main *r*-processes are similar.

#### 3.3. The Lanthanides and Hafnium

There are many detectable, relatively unblended lines from rare earth elements (La through Lu) and Hf. There is only a single, strong line available for Yb, which requires hyperfine and isotopic structure. The pattern of these elements agrees well with the Sun and with the other *r*-I and *r*-II stars. This robust pattern among the rare earth elements is a feature that is seen in all *r*-process-enhanced stars (e.g., Sneden et al. 2008), and is therefore an essential requirement for models of *r*-process nucleosynthesis.

#### 3.4. Third-peak Elements

Os and Ir are the only third-peak elements available in this spectrum. Os has three lines, though Ir only has one. Despite the paucity of lines, both elements agree very well with the other patterns in Figure 2 (and with other *r*-II stars; e.g., Placco et al. 2017).

The relative strength of the second and third r-process peaks is also important for constraining the physical conditions of the r-process (see Section 4.1). In J1538–1804, in the other r-I and r-II stars, and in the Sun, the pattern between the second and third peaks appears to be consistent across  $\sim$ 4 dex in metallicity.

#### 3.5. Lead

Only an upper limit of  $\log \epsilon$  (Pb) = 0.36 can be derived from the line at 4057.8 Å. Pb is a significant product of the *s*-process. This upper limit therefore supports the assertion from the [Ba/Eu] that any contamination from the *s*-process in this star is minimal. This agrees with Roederer et al. (2010), who find minimal *s*-process contributions to MW stars with metallicities as high as [Fe/H]  $\sim -1.4$ .

#### 3.6. The Actinides

Th and U are highly desirable elements for r-process studies for two reasons: (1) a handful of r-II stars exhibit an

Table 2 Abundances per Line

Table 2 (Continued)

	Wavelength	EP	log gf	Solar	J1538-1804	
	(Å)	(eV)		$\log\epsilon$	$\log\epsilon$	
r II	4161.792	2.938	-0.50	A09	$1.20 \pm 0.05$	
/ II	3747.556	0.104	-0.91	A09	$0.22 \pm 0.10$	
7 П 7 П	4398.013	0.129	-1.00	A09	$0.22 \pm 0.10$	
II II	4682.324 4883.680	0.408 1.083	-1.51 $0.07$	A09 A09	$0.24 \pm 0.05$ $0.42 \pm 0.10$	
r II	3998.954	0.558	-0.39	A09	$0.42 \pm 0.10$ $0.65 \pm 0.05$	
r II	4050.316	0.713	-1.00	A09	$1.02 \pm 0.05$	
r II	4317.299	0.713	-1.38	A09	$1.01 \pm 0.05$	
u I	3436.736	0.148	-0.02	1.67	$0.51 \pm 0.10$	
lu I	3498.942	0.000	0.33	1.60	$0.46 \pm 0.10$	
lu I	3798.898	0.148	-0.09	1.75	$0.56 \pm 0.10$	
tu I	3799.349	0.000	-0.07	1.75	$0.51 \pm 0.10$	
lh I	3396.819	0.000	0.05	0.91	$0.07 \pm 0.10$	
h I	3692.358	0.000	0.17	0.91	$-0.08 \pm 0.10$	
d I	3404.579	0.813	0.32	1.37	$0.08 \pm 0.10$	
g I	3382.889	0.000	-0.38	0.64	$0.55 \pm 0.10$	
a II	5853.675 <sup>a</sup>	0.604	-1.00	A09	$0.69 \pm 0.10$	
a II a II	6141.713 <sup>a</sup> 3988.515 <sup>a</sup>	0.704 0.403	-0.08 $0.21$	A09 A09	$0.69 \pm 0.10$ $-0.34 \pm 0.05$	
a II	4086.709 <sup>a</sup>	0.000	-0.07	A09	$-0.31 \pm 0.05$	
a II	4123.230 <sup>a</sup>	0.321	0.13	A09	$-0.31 \pm 0.05$ $-0.31 \pm 0.05$	
a II	4333.750 <sup>a</sup>	0.173	-0.06	A09	$-0.34 \pm 0.05$	
a II	5303.520 <sup>a</sup>	0.321	-1.35	A09	$-0.14 \pm 0.05$	
a II	6390.460 <sup>a</sup>	0.321	-1.41	A09	$-0.04 \pm 0.05$	
e II	3940.330	0.318	-0.27	A09	$-0.14 \pm 0.05$	
е п	3999.237	0.295	0.06	A09	$-0.03 \pm 0.05$	
e II	4014.897	0.529	-0.20	A09	$0.09 \pm 0.05$	
е п	4072.918	0.327	-0.64	A09	$0.19 \pm 0.05$	
e II	4073.474	0.477	0.21	A09	$0.07 \pm 0.05$	
e II	4083.230	0.700	0.27	A09	$0.16 \pm 0.05$	
e II e II	4120.840 4137.645	0.320 0.516	-0.37 $0.40$	A09 A09	$0.19 \pm 0.05$ $-0.06 \pm 0.05$	
e II	4138.096	0.924	-0.08	A09	$-0.00 \pm 0.05$ $0.17 \pm 0.05$	
e II	4165.599	0.909	0.52	A09	$0.09 \pm 0.05$	
e II	4222.597	0.122	-0.15	A09	$0.09 \pm 0.05$	
ЕΠ	4418.790	0.863	0.27	A09	$0.06 \pm 0.05$	
еп	4449.330	0.608	0.04	A09	$0.04 \pm 0.05$	
e II	4486.910	0.295	-0.18	A09	$0.01 \pm 0.05$	
е п	4562.370	0.477	0.21	A09	$0.09 \pm 0.05$	
e II	4628.160	0.516	0.14	A09	$0.19 \pm 0.10$	
e II	5274.230	1.044	0.13	A09	$0.19 \pm 0.05$	
II	3964.812ª	0.055	0.07	A09	$-0.52 \pm 0.05$	
r II	4179.393ª	0.204	0.46	A09	$-0.40 \pm 0.05$	
r II r II	4222.934 <sup>a</sup> 4408.819 <sup>a</sup>	0.055	0.23 0.05	A09 A09	$-0.48 \pm 0.05$ $-0.48 \pm 0.05$	
. II	5259.740 <sup>a</sup>	0.000 0.633	0.03	A09 A09	$-0.48 \pm 0.03$ $-0.47 \pm 0.10$	
d II	3862.566	0.033	-0.76	A09	$0.18 \pm 0.10$	
d II	3863.408	0.000	-0.01	A09	$0.03 \pm 0.10$	
d II	3900.215	0.471	0.10	A09	$0.05 \pm 0.10$	
d II	4021.728	0.182	-0.31	A09	$0.18 \pm 0.05$	
d II	4051.139	0.380	-0.30	A09	$0.18 \pm 0.05$	
d II	4061.080	0.471	0.55	A09	$0.23 \pm 0.05$	
d II	4069.270	0.064	-0.57	A09	$0.20 \pm 0.05$	
d II	4177.320	0.064	-0.10	A09	$0.12 \pm 0.05$	
d II	4178.635	0.182	-1.03	A09	$0.17 \pm 0.05$	
1 II	4179.580	0.182	-0.64	A09	$0.13 \pm 0.10$	
d II	4232.380	0.064	-0.47	A09	$0.18 \pm 0.05$	
d II	4446.390	0.204	-0.35	A09	$0.15 \pm 0.05$	
l П 1 п	4462.990	0.559	0.04	A09	$0.23 \pm 0.05$	
d II d II	4959.120 4989.950	0.064 0.630	-0.80 $-0.50$	A09 A09	$0.25 \pm 0.05$ $0.28 \pm 0.05$	
lп	5092.790	0.380	-0.50 -0.61	A09 A09	$0.28 \pm 0.03$ $0.19 \pm 0.05$	
і п 1 П	5130.590	1.303	0.45	A09	$0.19 \pm 0.05$ $0.09 \pm 0.05$	
dП	5212.350	0.204	-0.96	A09	$0.03 \pm 0.03$ $0.23 \pm 0.10$	
d II	5249.590	0.975	0.20	A09	$0.13 \pm 0.10$	
d II	5319.820	0.550	-0.14	A09	$0.23 \pm 0.05$	
m II	3896.970	0.040	-0.67	A09	$-0.25 \pm 0.05$	
	4188.128	0.543	-0.44	A09	$-0.05\pm0.05$	
m II						
m II m II	4318.926	0.277	-0.25	A09	$-0.13 \pm 0.05$	
	4318.926 4421.126 4424.337	0.277 0.378 0.484	-0.25 $-0.49$ $0.14$	A09 A09 A09	$-0.13 \pm 0.05$ $-0.03 \pm 0.05$ $-0.13 \pm 0.05$	

	W141-	ED		C - 1	I1520 1004
	Wavelength (Å)	EP (eV)	log gf	Solar	J1538-1804
		(CV)		$\log \epsilon$	$\log \epsilon$
Eu II	3724.931 <sup>a</sup>	0.000	-0.09	0.42	$-0.47 \pm 0.05$
Eu II	3907.107 <sup>a</sup>	0.207	0.17	0.42	$-0.52 \pm 0.05$
Eu II	4129.725 <sup>a</sup>	0.000	0.22	0.40	$-0.42 \pm 0.05$
Eu II	4435.578 <sup>a</sup>	0.207	-0.11	0.40	$-0.37 \pm 0.05$
Eu II	6645.064 <sup>a</sup>	1.379	0.12	0.52	$-0.27 \pm 0.05$
Gd II	3549.359	0.240	0.29	0.97	$-0.32 \pm 0.10$
Gd II	3697.733	0.032	-0.34	0.87	$-0.02 \pm 0.10$
Gd II	3768.396	0.078	0.21	A09	$-0.03 \pm 0.10$
Gd II	3796.384	0.032	0.02	A09	$0.03 \pm 0.10$
Gd II	3844.578	0.144	-0.46	A09	$0.23 \pm 0.10$
Gd II	4191.075	0.427	-0.48	A09	$0.13 \pm 0.05$
Gd II	4215.022	0.427	-0.44	A09	$0.13 \pm 0.05$
Gd II	4251.731	0.382	-0.22	A09	$0.05 \pm 0.05$
ТЬ ІІ	3702.850 <sup>a</sup>	0.126	0.44	A09	$-0.74 \pm 0.05$
ТЬ ІІ	3747.380 <sup>a</sup>	0.401	0.04	A09	$-0.64 \pm 0.10$
ТЬ ІІ	3848.730 <sup>a</sup>	0.000	0.28	A09	$-0.67 \pm 0.05$
ТЬ ІІ	3874.168 <sup>a</sup>	0.000	0.27	A09	$-0.77 \pm 0.05$
ТЬ ІІ	4002.566 <sup>a</sup>	0.641	0.10	A09	$-0.74 \pm 0.05$
Dy II	3757.368	0.100	-0.17	A09	$0.01 \pm 0.10$
Dy II	3944.680	0.000	0.11	A09	$0.21 \pm 0.10$
Dy II	3996.689	0.590	-0.26	A09	$0.26 \pm 0.05$
Dy II	4050.565	0.590	-0.47	A09	$0.36 \pm 0.05$
Dy II	4073.120	0.540	-0.32	A09	$0.38 \pm 0.05$
Dy II	4077.966	0.100	-0.04	A09	$0.51 \pm 0.05$
Dy II	4103.306	0.100	-0.38	A09	$0.46 \pm 0.05$
Ho II	3796.730 <sup>a</sup>	0.000	0.16	A09	$-0.44 \pm 0.05$
Ho II	3810.738ª	0.000	0.19	A09	$-0.56 \pm 0.10$
Ho II	3890.738ª	0.079	0.46	A09	$-0.61 \pm 0.10$
Er II	3692.649	0.055	0.28	A09	$-0.12 \pm 0.10$
Er II	3729.524	0.000	-0.59	A09	$-0.02 \pm 0.10$
Er II	3786.836	0.000	-0.52	A09	$-0.07 \pm 0.10$
Er II	3830.481	0.000	-0.22	A09	$-0.17 \pm 0.10$
Er II	3896.233	0.055	-0.12	A09	$-0.07 \pm 0.10$
Er II	3906.311	0.000	0.12	A09	$-0.17 \pm 0.10$
Er II	3938.626	0.000	-0.52	A09	$-0.27 \pm 0.10$
Tm II	3700.255	0.029	-0.38	A09	$-0.74 \pm 0.10$
Tm II	3701.362	0.000	-0.54	A09	$-0.89 \pm 0.10$
Tm II	3795.759	0.029	-0.23	A09	$-1.04 \pm 0.10$
Tm II	3848.019	0.000	-0.14	A09	$-0.67 \pm 0.10$
Tm II	3996.510	0.000	-1.20	A09	$-0.79 \pm 0.10$
Yb II	3694.190 <sup>a</sup>	0.000	-0.30	0.54	$0.08 \pm 0.10$
Lu II Hf II	3507.380 <sup>a</sup> 3719.276	0.000	-1.16	0.10	$-0.99 \pm 0.10$ $-0.34 \pm 0.10$
Hf II		0.608	-0.81	A09	$-0.34 \pm 0.10$ $-0.34 \pm 0.10$
Hf II	3918.090	0.452	-1.14	A09	
	4093.150	0.452	-1.15	A09	$-0.39 \pm 0.05$
Os I Os I	4135.775	0.515	-1.26	A09	$0.56 \pm 0.05$ $0.31 \pm 0.05$
Os I	4260.848 4420.520	0.000	-1.44 -1.53	A09 A09	$0.31 \pm 0.05$ $0.38 \pm 0.05$
Os 1 Ir I	4420.520 3800.120	0.000	-1.53 $-1.43$	A09 A09	$0.38 \pm 0.05$ $0.69 \pm 0.10$
Ir i Pb i	4057.807 <sup>a</sup>	1.319	-1.43 -0.17	A09 A09	$0.69 \pm 0.10$ < 0.36
Pb I Th II					
	4019.129	0.000	-0.23	A09	$-1.07 \pm 0.05$
Th II U II	4086.521 3859.571	0.000	-0.93 $-0.10$	A09 A09	$-0.87 \pm 0.05$
— II	3037.371	0.030	-0.10	A09	<-1.23

enhancement in the actinides (e.g., Schatz et al. 2002) and (2) U and Th are radioactive, and relative abundance ratios with respect to stable elements like Eu are therefore useful for agedating (see Section 4.2).

There are two clean Th II lines in this star, enabling a highprecision Th measurement (see Figure 1). The strongest Th II line at 4019 Å is often blended, but this problem is reduced at low metallicities. The U II line at 3859 Å is severely blended with an Fe I in J1538-1804, providing only an upper limit.

Line has hyperfine structure and/or isotopic splitting.

Table 3
Mean Neutron-capture Abundances and Errors

Element <sup>a</sup>	N	$\log \epsilon^{b}$	$\sigma_{ m rand}$	$\sigma_{\mathrm{Total}}^{}}$	$\sigma_{ m LTE}^{ m d}$	[X/Fe I]	$\sigma_{\mathrm{Total}}^{\mathrm{c}}$	$\sigma_{ m LTE}^{ m d}$
Sr II	1	1.20	0.10	0.12	-0.08	0.44	0.20	0.07
Y II	4	0.27	0.05	0.14	-0.11	0.17	0.13	0.04
Zr II	3	0.89	0.07	0.14	-0.09	0.47	0.13	0.06
Ru I	5	0.57	0.05	0.28	0.10	0.89	0.07	0.09
Rh I	2	-0.01	0.06	0.28	0.10	1.15	0.08	0.09
Pd I	1	0.28	0.10	0.36	0.07	0.78	0.15	0.06
Ag I	1	-0.25	0.10	0.30	0.09	0.88	0.11	0.08
Ba II	2	0.69	0.07	0.20	-0.16	0.62	0.08	-0.01
La II	6	-0.24	0.05	0.14	-0.11	0.77	0.12	0.04
Ce II	17	0.08	0.05	0.13	-0.07	0.61	0.13	0.08
Pr II	5	-0.47	0.05	0.13	-0.07	0.92	0.13	0.08
Nd II	20	0.17	0.05	0.13	-0.08	0.86	0.12	0.07
Sm II	5	-0.12	0.05	0.13	-0.08	1.03	0.12	0.07
Eu II	5	-0.32	0.05	0.15	-0.16	1.27	0.11	-0.01
Gd II	8	0.06	0.05	0.14	-0.10	1.10	0.12	0.05
Tb II	5	-0.71	0.05	0.13	-0.09	1.10	0.12	0.06
Dy II	7	0.31	0.05	0.20	-0.15	1.32	0.07	0.0
Ho II	3	-0.54	0.05	0.21	-0.15	1.09	0.07	0.0
Er II	7	-0.13	0.05	0.20	-0.14	1.06	0.07	0.01
Tm II	5	-0.83	0.05	0.18	-0.08	1.18	0.08	0.07
Yb II	1	0.25	0.10	0.24	-0.24	1.52	0.11	-0.09
Lu II	1	-0.99	0.10	0.20	-0.06	0.92	0.12	0.09
Hf II	3	-0.36	0.05	0.13	-0.06	0.90	0.13	0.09
Os I	3	0.42	0.05	0.28	0.09	1.13	0.07	0.08
Ir I	1	0.44	0.10	0.25	0.02	1.17	0.10	0.01
Pb I	1	< 0.36	•••	•••	•••	< 0.72	•••	•••
Th II	2	-0.97	0.07	0.14	-0.05	1.12	0.14	0.10
U II	1	<-1.23				< 1.42		

#### Notes.

Based on its Th and U abundances, J153830.9–180424 does not appear to be an actinide boost star.

#### 4. Discussion

#### 4.1. The r-process Pattern

Despite its high metallicity relative to other *r*-II stars, J1538–1804 is a typical *r*-II star. Its *r*-process abundance pattern matches that of the Sun and the *r*-II star CS 31082–001 (except that J1538–1804 does not exhibit an actinide boost like CS 31082–001). Its Pb abundance and [Ba/Eu] ratio suggest that the *s*-process has not contributed to its abundance pattern. It does not exhibit the light *r*-process variations that have been observed in other *r*-I and *r*-II stars (e.g., Honda et al. 2006), and therefore does not require excessive contributions from the "limited" (or "weak") *r*-process (see A. Frebel et al. 2018, in preparation).

Patterns such as these are essential for identifying the site(s) of the r-process. The abundance patterns from models are highly sensitive to uncertainties in nuclear masses,  $\beta$ -decay rates, fission cycling, neutrino properties, etc. (e.g., Surman et al. 2017), as well as the physical conditions of the

environment, such as temperature, the electron fraction, and the density (Hoffman et al. 1997). The pattern in J1538–1804 sets strong requirements for r-process models.

#### 4.2. Age

The Th abundance and the upper limit on U indicate that J1538-1804 is an old star. The relative abundances between Th and all the second- and third- peak elements in Table 3 give a mean age of  $11.2 \pm 3.9$  Gyr when the Schatz et al. (2002) initial production ratios (from waiting-point calculations) are adopted, and 17.2  $\pm$  7.2 Gyr when the Hill et al. (2017) values (from a high-entropy wind model) are adopted (see Table 7 in Placco et al. 2017). The quoted uncertainties represent the standard deviations from different chronometer pairs. The upper limit in the U abundance yields a lower limit from U/Eu of 5.4-5.7 Gyr, depending on the production ratio. The dominant sources of uncertainty in these ages are the choices of production ratios and chronometer pairs; while U/Th would be a better ratio, the upper limit on U is not very constraining in this case. These ratios strongly suggest that the r-process material in J1538–1804 is indeed ancient.

<sup>&</sup>lt;sup>a</sup> In order of atomic number.

<sup>&</sup>lt;sup>b</sup> A mean [X/H] abundance was calculated with a straight mean of the differential [X/H] ratios (using the solar values in Table 2); the mean [X/H] was then converted to  $\log \epsilon$  with the Asplund et al. (2009) solar abundance.

<sup>&</sup>lt;sup>c</sup> The total error refers to the combination of random and systematic errors (where the latter are due to uncertainties in the atmospheric parameters), calculated according to Equations (A1), (A4), and (A5) in McWilliam et al. (2013). Errors in  $\log \epsilon$  and [X/Fe] are listed separately.

<sup>&</sup>lt;sup>d</sup> The LTE error refers to the offset that results when LTE atmospheric parameters are used (see Paper I). Errors in log  $\epsilon$  and [X/Fe] are listed separately. These offsets are not included in  $\sigma_{\text{Total}}$ .

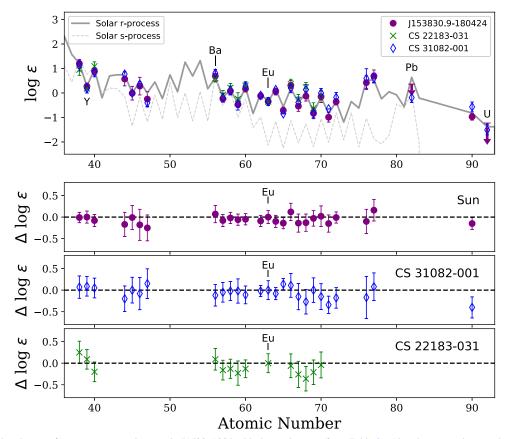


Figure 2. Top panel: abundances of neutron-capture elements in J1538–1804 with the total errors (from Table 3). Also shown are the r- and s-process patterns in the Sun (gray line; from Arlandini et al. 1999), and the abundances of an extremely metal-poor r-II (CS 31082–001; from Hill et al. 2002; Sneden et al. 2009; Siqueira Mello et al. 2013) and r-I (CS 22183-031; from Roederer et al. 2014b) stars. The solar r-process pattern and the metal-poor star abundances are shifted to the Eu abundance in J1538–1804; the solar s-process pattern is shifted to match the Ba abundance. Upper limits are shown for Pb and U. Bottom panels: abundance offsets between J1538–1804 and the other stars, where  $\Delta \log \epsilon(X) = \log \epsilon_{J1538}(X) - \log \epsilon_{\rm star}(X)$ . The offsets from the Sun (second panel) are relative to the solar r-process residuals.

#### 4.3. J1538-1804 and the Site of the r-process

Given that J1538–1804 is a typical r-II star, what is gained from these observations? First, this adds another r-II star to the known sample of  $\sim$ 30, an important step for statistically analyzing the r-process patterns as a function of stellar properties (metallicity, age, location in the Galaxy, etc.). Second, this Letter has demonstrated that the main r-process pattern at  $[Fe/H] \sim -2$  is very similar to the pattern at  $[Fe/H] \sim -4$  and  $[Fe/H] \sim 0$ , a powerful result when combined with the old age of the r-process material. This either implicates a single site for the r-process, or requires that r-process nucleosynthesis in different sites leads to the same final abundance pattern.

The very metal-poor *r*-I and *r*-II stars without *s*-process contamination also provide constraints on the birth environments of these stars, such as star formation rates, timescales relative to the onset of type Ia supernovae, etc. The recent discovery of *r*-process-enhanced stars in ultra-faint dwarfs (e.g., Ji et al. 2016) has led to speculation that all *r*-process-enhanced stars may originate in dwarf galaxies. Indeed, J1538–1804's retrograde orbit in the MW halo does hint at a possible extragalactic origin. The chemical enrichment of lower-mass systems proceeds more slowly than in more massive galaxies; in particular, dwarf galaxies cannot form as many metal-rich stars, and the onset of the *s*-process occurs at a lower [Fe/H] than in the MW (e.g., Tolstoy et al. 2009). As more *r*-I and *r*-II stars are discovered, particularly at higher

metallicities, the general properties of the r-process-enhanced stellar population will place limits on the nature of the birth environments of these stars.

In the near future, the R-Process Alliance will provide data for many more *r*-process-enhanced stars. This will enable *r*-I and *r*-II stars to be studied as stellar populations, and will provide *r*-process patterns that can be used to tease out any variations as a function of stellar properties such as metallicity, location, and more.

The authors wish to the thank the anonymous referee for helpful comments. C.M.S., G.W., C.E.D., and E.M.F. acknowledge funding from the Kenilworth Fund of the New York Community Trust. V.M.P., E.M.H., T.C.B., and I.U.R. acknowledge partial support from grant PHY 14-30152 (Physics Frontier Center/JINA-CEE), awarded by the US National Science Foundation. Australian access to the *Magellan* Telescopes was supported through the National Collaborative Research Infrastructure Strategy of the Australian Federal Government.

#### ORCID iDs

Charli M. Sakari https://orcid.org/0000-0002-5095-4000
Vinicius M. Placco https://orcid.org/0000-0003-4479-1265
Terese Hansen https://orcid.org/0000-0001-6154-8983
Erika M. Holmbeck https://orcid.org/0000-0002-5463-6800

Timothy C. Beers https://orcid.org/0000-0003-4573-6233 Anna Frebel https://orcid.org/0000-0002-2139-7145 Ian U. Roederer https://orcid.org/0000-0001-5107-8930

#### References

```
Abbott, B. P., Abbott, R., Abbott, T. D., et al. 2017, PhRvL, 119, 161101
Abohalima, A., & Frebel, F. 2017, ApJS, submitted (arXiv:1711.04410)
Amarsi, A. M., Lind, K., Asplund, M., Barklem, P. S., & Collet, R. 2016,
           , 463, 1518
Aoki, W., Honda, S., Sadakane, K., & Arimoto, N. 2007, PASJ, 59, 15
Arcones, A., & Thielemann, F.-K. 2013, JPhG, 40, 3201
Arlandini, C., Käppeler, F., Wisshak, K., et al. 1999, ApJ, 525, 886
Asplund, M., Grevesse, N., Sauval, J. A., & Scott, P. 2009, ARA&A, 47, 481
Barklem, P. S., Christlieb, N., Beers, T. C., et al. 2005, A&A, 439, 129
Beniamini, P., Hotokezaka, K., & Piran, T. 2016, ApJ, 832, 149
Bernstein, R., Shectman, S. A., Gunnels, S. M., Mochnacki, S., & Athey, A. E.
   2003, Proc. SPIE, 4841, 1694
Brown, A. G. A., Vallenari, A., Prusti, T., et al. 2016, A&A, 595, A2
Cayrel, R., Hill, V., Beers, T. C., et al. 2001, Natur, 409, 691
Chornock, R., Berger, E., Kasen, D., et al. 2017, ApJL, 848, L19
Christlieb, N., Beers, T. C., Barklem, P. S., et al. 2004, A&A, 428, 1027
Côté, B., Fryer, C. L., Belczynski, K., et al. 2017, ApJ, submitted (arXiv:1710.
Cowan, J. J., Sneden, C., Beers, T. C., et al. 2005, ApJ, 627, 238
Frebel, A., Chiti, A., Ji, A. P., Jacobson, H. R., & Placco, V. M. 2015, ApJL,
   810, L27
Hill, V., Christlieb, N., Beers, T. C., et al. 2017, A&A, 607, 91
Hill, V., Plez, B., Cayrel, R., et al. 2002, A&A, 387, 560
Hoffman, R. D., Woosley, S. E., & Qian, Y.-Z. 1997, ApJ, 482, 957
Honda, S., Aoki, W., Ishimaru, Y., Wanajo, S., & Ryan, S. G. 2006, ApJ,
   643, 1180
Howes, L. M., Asplund, M., Keller, S. C., et al. 2016, MNRAS, 460, 884
```

```
Ji, A. P., Frebel, A., Simon, J. D., & Chiti, A. 2016, ApJ, 830, 93
Lattimer, J., & Schramm, D. 1974, ApJL, 192, L145
Lippuner, J., Fernández, R., Roberts, L. F., et al. 2017, MNRAS, 472, 904
Macias, P., & Ramirez-Ruiz, E. 2016, ApJL, submitted (arXiv:1609.04826)
Matijevič, G., Chiappini, C., Grebel, E. K., et al. 2017, A&A, 603, 19
McWilliam, A., Wallerstein, G., & Mottini, M. 2013, ApJ, 778, 149
Placco, V. M., Frebel, A., Beers, T. C., & Stancliffe, R. J. 2017, ApJ, 797,
Placco, V. M., Holmbeck, E. M., Frebel, A., et al. 2017, ApJ, 844, 18
Roederer, I. U. 2013, AJ, 145, 26
Roederer, I. U., Cowan, J. J., Karakas, A. I., et al. 2010, ApJ, 724, 975
Roederer, I. U., Cowan, J. J., Preston, G. W., et al. 2014a, MNRAS, 445, 2970
Roederer, I. U., Preston, G. W., Thompson, I. B., et al. 2014b, AJ, 147,
Rosswog, S., Korobkin, O., Arcones, A., Thielemann, F.-K., & Piran, T. 2014,
   MNRAS, 439, 744
Schatz, H., Toenjes, R., Pfeiffer, B., et al. 2002, ApJ, 579, 626
Siqueira Mello, C., Hill, V., Barbuy, B., et al. 2014, A&A, 565, 93
Siqueira Mello, C., Spite, M., & Barbuy, B. 2013, A&A, 550, 122
Sneden, C. 1973, ApJ, 184, 839
Sneden, C., Cowan, J. J., & Gallino, R. 2008, ARA&A, 46, 241
Sneden, C., Lawler, J. E., Cowan, J. J., Ivans, I. I., & Den Hartog, E. A. 2009,
   ApJS, 182, 80
Sneden, C., McWilliam, A., Preston, G. W., et al. 1996, ApJ, 467, 819
Sneden, C., Preston, G. W., McWilliam, A., & Searle, L. 1994, ApJ, 431, 27
Sobeck, J. S., Kraft, R. P., Sneden, C., et al. 2011, AJ, 141, 175
Steinmetz, M., Zwitter, T., Siebert, A., et al. 2006, AJ, 132, 1645
Surman, R., Mumpower, M., & McLaughlin, G. 2017, in Proc. 14th Int. Symp.
  on Nuclei in the Cosmos, ed. S. Kubono et al., 010612
Tolstoy, E., Hill, V., & Tosi, M. 2009, ARA&A, 47, 371
Tsujimoto, T., & Nishimura, N. 2015, ApJL, 811, L10
Winteler, C., Käppeli, R., Perego, A., et al. 2012, ApJL, 750, L22
Xu, H. L., Svanberg, S., Quinet, P., Palmeri, P., & Biémont, É. 2007, JQSRT,
   104, 52
```