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# Chemical Abundance Signature of J0023+0307: A Second-generation Main-sequence Star with $[\text{Fe}/\text{H}] < -6^*$

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## Abstract

We present a chemical abundance analysis of the faint halo metal-poor main-sequence star J0023+0307, with  $[\text{Fe}/\text{H}] < -6.3$ , based on a high-resolution ( $R \sim 35,000$ ) *Magellan*/MIKE spectrum. The star was originally found to have  $[\text{Fe}/\text{H}] < -6.6$  based on a Ca II K measurement in an  $R \sim 2500$  spectrum. No iron lines could be detected in our MIKE spectrum. Spectral lines of Li, C, Na, Mg, Al, Si, and Ca were detected. The Li abundance is close to the Spite Plateau,  $\log \epsilon(\text{Li}) = 1.7$ , not unlike that of other metal-poor stars, although in stark contrast to the extremely low value found, e.g., in HE 1327–2326 at a similar  $[\text{Fe}/\text{H}]$  value. The carbon G-band is detected and indicates strong C-enhancement, as is typical for stars with low Fe abundances. Elements from Na through Si show a strong odd–even effect, and J0023+0307 displays the second-lowest known  $[\text{Ca}/\text{H}]$  abundance. Overall, the abundance pattern of J0023+0307 suggests that it is a second-generation star that formed from gas enriched by a massive Population III first star exploding as a fallback supernova. The inferred dilution mass of the ejecta is  $10^{5 \pm 0.5} M_{\odot}$  of hydrogen, strongly suggesting J0023+0307 formed in a recollapsed minihalo. J0023+0307 is likely very old because it has a very eccentric orbit with a pericenter in the Galactic bulge.

**Key words:** Galaxy: halo – nuclear reactions, nucleosynthesis, abundances – stars: abundances – stars: individual (J0023+0307) – stars: Population II

## 1. Introduction

The most metal-poor stars are tracers of the physical and chemical conditions of the early universe. In their atmospheres, they carry a record of this early time that was governed by the first stars, the first supernovae, and the formation of the first galaxies (Beers & Christlieb 2005; Frebel & Norris 2015). Much progress has been made in the past two decades to systematically uncover the extremely rare most metal-poor stars through systematic searches, and to chemically characterize them using large telescopes equipped with high-resolution spectrographs. There are now about  $\sim 30$  stars known with  $[\text{Fe}/\text{H}] \lesssim -4.0$ , of which  $\sim 10$  have  $[\text{Fe}/\text{H}] \lesssim -4.5$  (Christlieb et al. 2002; Frebel et al. 2005; Norris et al. 2007; Caffau et al. 2011; Hansen et al. 2014; Keller et al. 2014; Allende Prieto et al. 2015; Bonifacio et al. 2015; Frebel et al. 2015; Aguado et al. 2018b; Starkenburg et al. 2018). Given their small Fe abundances, they are believed to be second-generation objects, i.e., only one progenitor was responsible for the elements now observed in each of these stars. This, in turn, enables studies of the first supernovae and the associated nucleosynthesis by studying the stellar chemical abundance patterns and comparing them with theoretical predictions of the yields of these first explosions (Umeda & Nomoto 2003; Iwamoto et al. 2005; Heger & Woosley 2010; Placco et al. 2015a, 2016). This provides one of the few means to gain clues on the nature of properties of the first stars. One striking feature besides the low Fe abundance is that nearly all these stars are enhanced in

carbon, suggesting that the first stars produced copious amounts of carbon during late stages of stellar evolution and/or their explosions (Umeda & Nomoto 2003; Meynet et al. 2006; Choplin et al. 2018). Recent searches for these most iron-poor stars include those with the SkyMapper telescope (Keller et al. 2014; Jacobson et al. 2015), the Pristine survey (Starkenburg et al. 2017), the Best and Brightest survey (Schlaufman & Casey 2014), and ToPoS (Caffau et al. 2013), and those based on Sloan Digital Sky Survey (SDSS) data (Aguado et al. 2016).

SDSS J002314.00+030758.0 (hereafter J0023+0307) was originally discovered by the SDSS/Baryon Oscillation Spectroscopic Survey (BOSS; Eisenstein et al. 2011; Dawson et al. 2013). It was then recognized as a particularly interesting metal-poor star by Aguado et al. (2018a), who reobserved it with medium-resolution spectroscopy. A very weak Ca K line was detected (together with a significant contribution of interstellar Ca), but no Ca abundance was presented. Instead, just an upper limit on the Fe abundance was listed,  $[\text{Fe}/\text{H}] < -6.6$ . No Fe lines were detected in their spectrum. Then, François et al. (2018) found  $[\text{Ca}/\text{H}] = -5.7$ , inferring  $[\text{Fe}/\text{H}] < -6.1$  under the assumption that  $[\text{Ca}/\text{Fe}] = 0.4$ , because no Fe lines were detected. In this paper, we report new high-resolution spectroscopic observations of this star. Still, no Fe lines could be detected. The local thermal equilibrium (LTE) limit (obtained from the strongest line at 3859 Å) is  $[\text{Fe}/\text{H}] < -5.6$  (the non-local thermodynamic equilibrium, or NLTE, limit is  $[\text{Fe}/\text{H}] < -5.2$ ), and our Ca abundance derived limit is  $[\text{Fe}/\text{H}] < -6.3$ . These results confirm that J0023+0307 is part of the group of stars with the lowest Fe abundances known.

\* This paper includes data gathered with the 6.5 m *Magellan* telescopes located at Las Campanas Observatory, Chile.

<sup>5</sup> Hubble Fellow.

## 2. Observations

J0023+0307 (R.A. = 00:23:14.0, decl. = +03:07:58.1,  $g = 17.9$ ) was observed with the Magellan Inamori Kyocera Echelle (MIKE) spectrograph (Bernstein et al. 2003) on the Magellan-Clay telescope at Las Campanas Observatory on 2018 July 6–8, 15, and 24–25, and August 31–September 2, using a  $0''.7$  slit and  $2 \times 2$  binning. The total exposure time was  $\sim 24$  hr but only  $\sim 20$  hr yielded useful data, i.e., with seeing of  $\lesssim 1''.3$ . Otherwise, the seeing was  $\lesssim 0''.8$  for  $\sim 11$  hr and  $\lesssim 1''.1$  for  $\sim 9$  hr. The spectral resolution is  $\sim 30,000$  in the red and  $\sim 35,000$  in the blue wavelength regimes. The spectrum covers 3050–9000 Å but is usable only above  $\sim 3700$  Å. Data were reduced with the CarPy pipeline (Kelson 2003).<sup>6</sup> The resulting signal-to-noise ratio (S/N) per pixel is  $\sim 40$  at  $\sim 4000$  Å,  $\sim 60$  at  $\sim 4700$  Å,  $\sim 35$  at  $\sim 5200$  Å, and  $\sim 75$  at  $\sim 6700$  Å. We show portions of the spectrum around the Li I doublet at 6706 Å, the Fe I K line at 3859 Å, the Ca II K line at 3933 Å, the Mg b lines at 5170 Å, and the G-band head region around 4313 Å in Figure 1. We also show the spectrum of HE 1327–2326 (Frebel et al. 2005) for comparison purposes, which has a very similar effective temperature ( $T_{\text{eff}} = 6180$ –).

The heliocentric radial velocity obtained from the position of the Mg b lines is  $-194.6 \pm 1.2$  km s<sup>-1</sup>. We see no evidence for velocity variations in our MIKE data over the course of two months. Aguado et al. (2018a) quoted  $v = -110 \pm 9$  km s<sup>-1</sup> from the BOSS discovery spectrum, but this appears to be due to a missed conversion between vacuum and air wavelengths. BOSS reports a velocity of  $-180 \pm 7$  km s<sup>-1</sup> (Dawson et al. 2013), but this velocity relies on broad Balmer lines, uses an inappropriate template (a metal-rich B star), and may be affected by interstellar medium (ISM) absorption lines. Our analysis of several H lines in the BOSS spectrum yields a velocity of  $-205 \pm 10$  km s<sup>-1</sup>, consistent with the MIKE velocity. There are also two X-Shooter (Vernet et al. 2011) observations available on the Very Large Telescope (VLT) archive<sup>7</sup> (François et al. 2018). We independently reduced these data with the default X-Shooter pipeline (Goldoni et al. 2006). The two spectra were observed 2017 June 30 and 2017 July 21, with velocities of  $-198.7$  and  $-192.5$  km s<sup>-1</sup>, showing no significant deviation from our MIKE velocities.

## 3. Stellar Parameters

J0023+0307 has no detected Ti, Fe, or Ca I lines. It is thus not possible to determine stellar parameters through standard spectroscopic methods. We determined an effective temperature from SDSS photometry (Ahn et al. 2012). J0023+0307 has  $g$ ,  $r$ ,  $i$ , =  $17.90 \pm 0.01$ ,  $17.62 \pm 0.01$ ,  $17.52 \pm 0.01$ . The dust maps from Schlafly & Finkbeiner (2011) give reddening values  $A(g) = 0.108$ ,  $A(r) = 0.075$ ,  $A(i) = 0.056$ , with a color uncertainty of 0.02 mag. We converted the dereddened magnitudes to  $B - V$  and  $V - I$  colors with the conversions in Jordi et al. (2006), then applied the color–temperature relations from Casagrande et al. (2010) assuming  $[\text{Fe}/\text{H}] = -5$ . For  $B - V$  and  $V - I$ , we obtained  $T_{\text{eff}} = 6260 \pm 140$  K and  $T_{\text{eff}} = 5997 \pm 130$  K, respectively. The reddening uncertainty dominates the temperature uncertainty. Since the  $B - V$  color–temperature relation is extremely sensitive to metallicity and the color–temperature relations are not well calibrated at

the lowest metallicity (Casagrande et al. 2010), we adopted the temperature from the  $V - I$  color, which has minimal dependence on metallicity ( $\approx 50$  K from  $[\text{Fe}/\text{H}] = -5$  to  $[\text{Fe}/\text{H}] = -3$ ).

We determined  $\log g$  with the equation (Venn et al. 2017):

$$\log g = 4.44 + \log M_{\star} + 4 \log(T_{\text{eff}}/5780) + 0.4(g_0 - \mu + \text{BC}(g) - 4.75).$$

We used  $g_0 = 17.79$  and  $T_{\text{eff}} = 5997 \pm 130$  K from above. We adopted  $M_{\star} = 0.5\text{--}0.6 M_{\odot}$  as a reasonable range for a 12-Gyr-old main-sequence turnoff star. For the distance modulus, we used the parallax-based distance from *Gaia* Data Release 2 (*Gaia* DR2) inferred by Bailer-Jones et al. (2018), giving  $\mu = 11.68^{+0.54}_{-0.49}$  (*Gaia* Collaboration et al. 2016, 2018). The bolometric correction for the SDSS G-band from Casagrande & Vandenberg (2014) is  $\text{BC}(g) = -0.33 \pm 0.03$ , where the error covers differences from varying the temperature within its error bars. In total, this gives  $\log g = 4.61 - 4.69 \pm 0.25$ , where the range is changing the mass from 0.5 to  $0.6 M_{\odot}$  and the uncertainty is dominated by the parallax uncertainty. Both of these values suggest that J0023+0307 is a main-sequence star. In the following, we thus adopted  $\log g = 4.6$ , taking into account the unevolved nature and hence a lower mass for the star. The surface gravity has little influence on any of the derived chemical abundances for such warm stars, so this choice is not critical. For the microturbulence, we adopted  $v_{\text{mic}} = 1.5$ , a value typical for warm stars (Barklem et al. 2005; Aoki et al. 2013). Given that only weak lines are available and that the star is fairly warm, this choice hardly affects any of the determined abundances.

For comparison, Aguado et al. (2018a) adopted  $T_{\text{eff}} = 6188 \pm 84$  K,  $\log g = 4.9 \pm 0.5$ , and  $v_{\text{mic}} = 2.0$  from their analysis of a  $R \sim 2500$  Intermediate-dispersion Spectrograph and Imaging System (ISIS) spectrum. François et al. (2018) found  $T_{\text{eff}} = 6160 \pm 100$  K from the  $(g-z)$  color, and adopted a canonical main-sequence surface gravity of  $\log g = 4.5$ , in very good agreement with our results.

## 4. The Chemical Abundance Pattern of J0023+0307—Iron-poor but Not as Metal-poor

### 4.1. LTE Abundance Determination

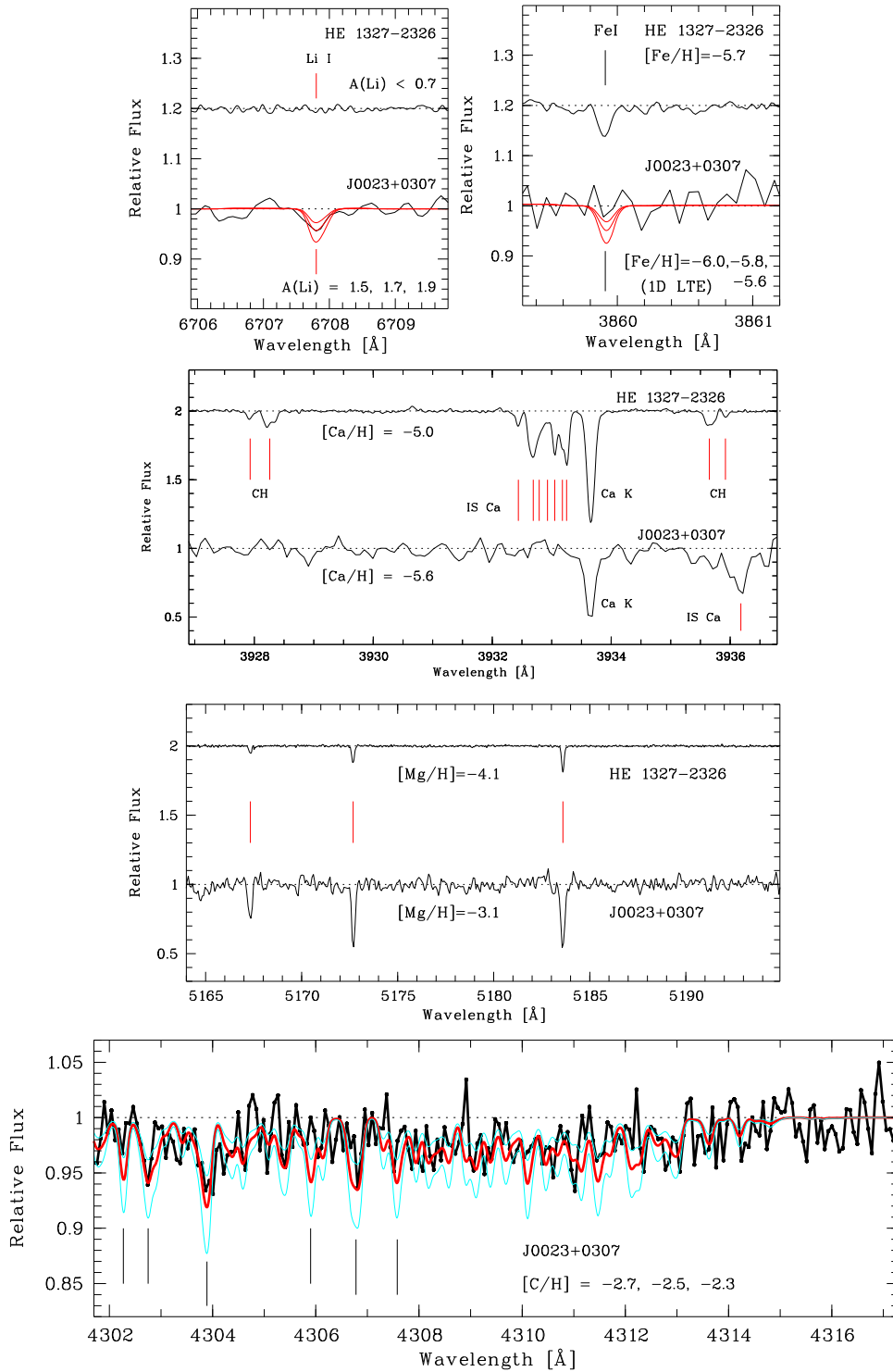
To calculate abundances, we used the 2017 version of the spectrum synthesis computer program MOOG including Rayleigh scattering (Snedden 1973; Sobeck et al. 2011) and using Barklem damping coefficients.<sup>8</sup> Our analysis software (first described in Casey 2014) employs a 1D plane-parallel model atmosphere with  $\alpha$ -enhancement (Castelli & Kurucz 2004) and assumes LTE. Equivalent widths were measured for lines that were detected (Na, Mg, Al, Si, Ca) and abundances were calculated correspondingly. For Li, C, and upper limits of other elements (e.g., Ti, Fe, Sr, Ba), a spectrum synthesis approach was chosen. Line abundances are listed in Table 1. Lower limits on  $[\text{X}/\text{Fe}]$  abundances were calculated using solar abundances from Asplund et al. (2009), and are given in Table 2. In the following, we comment on selected elements and their abundance determinations.

We tentatively detected Li with an abundance  $A(\text{Li}) = 1.7$ . In Figure 1, we show the Li line and synthetic spectra for  $A$

<sup>6</sup> Available at <https://code.obs.carnegiescience.edu/>.

<sup>7</sup> Based on observations collected at the European Southern Observatory (ESO) under ESO programme(s) 099.D-0576(A).

<sup>8</sup> <https://github.com/alexji/moog17scat>



**Figure 1.** Portions of the *Magellan*/MIKE spectrum of J0023+0307 in comparison with HE 1327–2326 (except for the CH region) near the Li I doublet at 6707 Å and the Fe I line at 3859 Å (top), the Ca II K line at 3933 Å, around the Mg b lines at 5170 Å, and the CH G-band at ~4313 Å (bottom). Some absorption lines are indicated. See the text for a discussion.

(Li) =  $1.7 \pm 0.2$ . There is clearly a significant feature at the position of the Li doublet at 6707 Å, but the line profile is somewhat distorted, and the overall depth is only  $\approx 3\%$ . Several other apparent absorption features of similar depth occur near the  $\lambda$  6707 line, but all these are clearly associated with imperfect sky subtraction residuals. Indeed, the sky spectrum shows no features at the Li line position, which suggests that the line is real. Still, we caution that the Li measurement is

rather uncertain and deserves further study based on additional data.

For C, we closely investigated the CH G-band region between 4300 and 4325 Å. The S/N is moderate, but several of the six strongest, most isolated CH features in this region (4302.27, 4302.75, 4303.89, 4305.90, 4306.78, and 4307.58 Å) are well detected. Other regions (4310–4313 Å) are consistent with the abundance from these strong lines. The observed

**Table 1**  
Equivalent Widths Measurements

$\lambda$ (Å)	Species	$\chi$ (eV)	$\log gf$ (dex)	EW (mÅ)	$\log \epsilon$ (X) [dex LTE]	$\log \epsilon$ (X) [dex NLTE]
6707	Li I	...	...	17	1.70 <sup>a</sup>	...
4313	CH	...	...	syn	5.89	...
5889.95	Na I <sup>a</sup>	0.00	0.11	19.5:	2.26:	2.32:
5895.92	Na I <sup>a</sup>	0.00	-0.19	10.0:	2.23:	2.31:
3829.36	Mg I	2.71	-0.21	105.4	4.52	4.64
3832.30	Mg I	2.71	0.27	168.5	4.62	4.81
3838.29	Mg I	2.72	0.49	207.9	4.65	4.78
4702.99	Mg I	4.34	-0.38	16.5	4.74	4.98
5167.32	Mg I	2.71	-1.03	65.1	4.68	4.82
5172.68	Mg I	2.71	-0.45	98.1	4.49	4.60
5183.60	Mg I	2.72	-0.24	128.6	4.60	4.65
5528.41	Mg I	4.34	-0.50	12.9	4.73	4.95
3944.06	Al I	0.00	-0.63	29.3	2.40	2.80
3961.52	Al I	0.01	-0.34	27.4	2.21	2.97
3905.52	Si I	1.91	-1.09	51.1	3.94	4.49
3933.66	Ca II	0.00	0.11	98.0	0.57	0.62
4226.73	Ca I	0.00	0.24	<12	<1.4	...
4246.82	Sc II	0.32	0.24	<8	<-0.4	...
3761.32	Ti II	0.57	0.18	<12	<0.6	...
4254.33	Cr I	0.00	-0.11	<8	<1.4	...
4030.75	Mn I	0.00	-0.48	<10	<1.5	...
3859.11	Fe I	0.00	-0.71	<10	<1.9	<2.3
3858.30	Ni I	0.42	-0.95	<10	<2.2	...
4077.71	Sr II	0.00	0.15	<10	<-1.5	...
4554.03	Ba II	0.00	0.16	<8	<-1.3	...

**Note.**

<sup>a</sup> See the text for a discussion.

spectrum with synthetic spectra (using the linelist of Masseron et al. 2014) overlaid is shown in Figure 1 (bottom panel). The resulting best-fit abundance is  $[C/H] = -2.5$ . The CH feature at 4323 Å can be fit with the same abundance. We also tested to what extent the C abundance depends on the O abundance. We found no change in carbon abundance up to an assumed  $[O/H] = -1.0$ . For higher O values, small changes (0.05 dex) begin to appear that would reduce the C abundance, but such a high O abundance is not likely to be correct. This is in line with a detailed investigation by Gallagher et al. (2016), who found that the O abundance has almost no effect in 1D, but instead significantly affects 3D C abundance determinations. Hence, J0023+0307 appears to be another carbon-enhanced hyper-metal-poor star. Our measurement is also in agreement with that of Aguado et al. (2018a), whose upper limit is  $[C/H] < -2.3$ . François et al. (2018) reported a higher value of  $[C/H] = -2.0$ . In fact, they found systematically higher abundances for C, Mg, and Si, by 0.3 and 0.5 dex. The exception is Ca, which agrees well. Reasons for that might include the lower resolving power and S/N of the spectrum, model atmosphere techniques used, or both.

Both Na D lines were detected, but their weak lines presented as unusually broad. We thus consider the abundances as uncertain, even though both lines yield nearly identical abundances. We note that at least three, likely four, prominent interstellar Na components were found, but they were well separated from the stellar lines. The interstellar features for Na and Ca II K have essentially identical shapes. While it is possible that an additional interstellar feature could blend with the stellar Na lines, it appears unlikely given that such gas

would need to have the same high velocity as J0023+0307 ( $-195 \text{ km s}^{-1}$ ). Higher S/N data might thus yield a more accurate NA abundance.

Aguado et al. (2018a) indirectly suggested that their  $[Ca/H]$  measurement would be  $[Ca/H] = [Fe/H] + [Ca/Fe] = -6.2$ . We derived  $[Ca/H] = -5.8$ , in agreement when taking into account that their value was derived from a spectrum with much lower resolution, which could have been affected by the subtraction of an interstellar Ca component. The interstellar component is well separated from the stellar Ca feature in our high-resolution spectrum; see Figure 1. François et al. (2018) reported  $[Ca/H] = -5.7$ , in very good agreement with our measurement.

Since no Fe lines were detected, we stacked portions of the spectrum that each contain a strong isolated Fe I line to investigate whether a signal would emerge in the composite spectrum. We followed the procedure outlined in Frebel et al. (2007b). The same was done using synthetic spectra of varying abundances. We verified this procedure by stacking the spectral regions of four detected Mg I lines to reproduce the Mg abundance obtained from individual lines. Stacking 19 Fe lines between 3750 and 4050 Å did not yield a signal. Adding more lines from noisier, bluer parts of the spectrum did not change the outcome. By matching the noise level in the composite spectrum with the corresponding composite synthetic spectrum, we derived an upper limit of  $[Fe/H] < -5.8$ . This is shown in Figure 2. We also stacked portions of the spectrum containing 10 Ti II lines. No signal was produced, and the upper limit is  $[Ti/H] < -4.5$ .

Considering that the other three stars with  $[Fe/H] < -5.0$  (based on detected Fe lines) all have Ca abundances in excess



**Table 2**  
Chemical Abundances of J0023+0307 and HE 1327–2326

Species	J0023+0307					HE 1327–2326				
	[X/H] (LTE)	[X/Fe]	[X/H] (NLTE)	[X/Fe]	$\sigma$	[X/H] (LTE)	[X/Fe] (NLTE)	[X/H]	[X/Fe]	
Li I	...	1.70 <sup>a</sup>	...	...	0.20	...	<0.70 <sup>a</sup>	...	...	
CH	–2.54	>3.76	...	...	0.20	–1.53	4.18	...	...	
Na I	–4.00:	>2.30	–3.92:	>1.98	0.30	–3.25	2.46	–3.30	1.90	
Mg I	–2.97	>3.33	–2.82	>3.08	0.10	–4.06	1.65	–3.80	1.40	
Al I	–4.15	>2.15	–3.56	>2.34	0.20	–4.55	1.16	–4.02	1.18	
Si I	–3.55	>2.75	–3.02	>2.88	0.20	–4.50	0.70	–4.31	0.89	
Ca II	–5.77	>0.53	–5.72	>0.18	0.20	–5.00	0.71	...	...	
Sc II	<–3.5	...	...	...	...	<–1.68	<0.98	...	...	
Ti II	<–4.3 <sup>b</sup>	...	...	...	...	–5.04	0.67	...	...	
Ti II	<–4.5 <sup>c</sup>	...	...	...	...	...	...	...	...	
Cr I	<–4.2	...	...	...	...	<–5.19	<0.52	...	...	
Mn I	<–4.0	...	...	...	...	<–4.59	<1.12	...	...	
Fe I	<–5.6 <sup>b</sup>	...	<–5.2	...	...	–5.71	...	–5.2 <sup>d</sup>	...	
Fe I	<–5.8 <sup>c</sup>	...	<–5.4	...	...	...	...	...	...	
Fe I	<–6.3 <sup>c</sup>	...	<–5.9	...	...	...	...	...	...	
Ni I	<–4.0	...	...	...	...	–5.49	0.22	...	...	
Sr II	<–4.3	...	...	...	...	–4.63	1.08	...	...	
Ba II	<–3.5	...	...	...	...	–4.32	1.39	...	...	

**Notes.** HE 1327–2326 abundances taken from Collet et al. (2006) and Frebel et al. (2008), except for Si, taken from Ezzeddine & Frebel (2018).

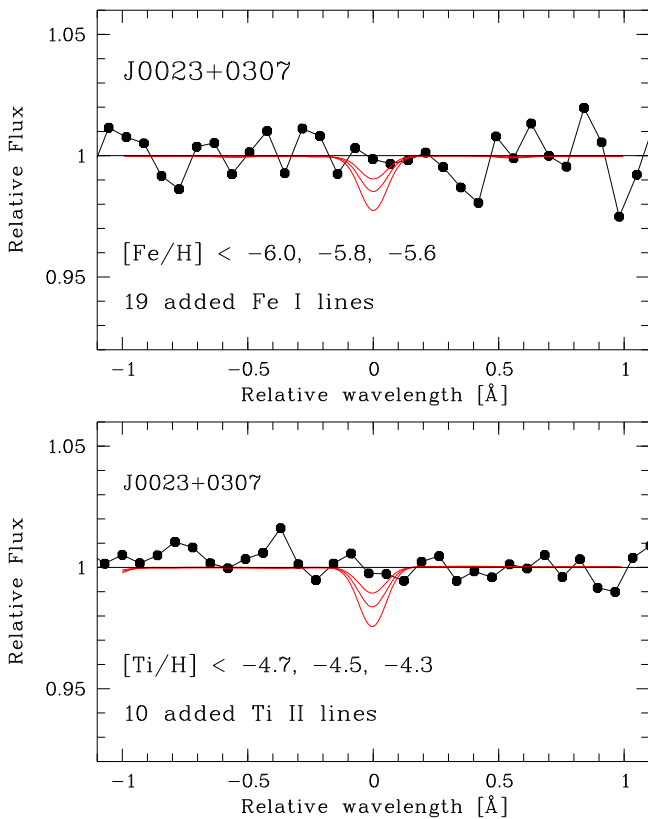
<sup>a</sup>  $A(\text{Li})$  is given instead of  $[\text{X}/\text{Fe}]$ . See the text for a discussion.

<sup>b</sup> Obtained from the Ti II line at 3761 Å, and Fe I line at 3859 Å, respectively.

<sup>c</sup> Obtained from the composite spectrum of eight Ti II lines and 16 Fe I lines, respectively.

<sup>d</sup> Adopted from Ezzeddine & Frebel (2018).

<sup>e</sup> Final adopted Fe abundance. See the text for a discussion. This value is used to calculate  $[\text{X}/\text{Fe}]$ .



**Figure 2.** Top: composite spectrum of 19 Fe I lines overlaid with composite synthetic spectra of varying abundances as illustration. Bottom: composite spectrum of 10 Ti II lines overlaid with composite synthetic spectra of varying abundances.

of iron, we used this as an additional constraint on the iron abundance of J0023+0307. HE 1327–2326 has  $[\text{Ca}/\text{Fe}] = 0.71$  (Frebel et al. 2008), HE 0107–5240 has  $[\text{Ca}/\text{Fe}] = 0.46$  (Christlieb et al. 2004), and SMSS0313–6707 has  $[\text{Ca}/\text{Fe}] > 0.34$  (Bessell et al. 2015; considering 1D LTE values). We thus assumed a ratio of  $[\text{Ca}/\text{Fe}] \gtrsim 0.5$  for J0023+0307, given the many chemical similarities in these four stars compared to those with iron abundances of  $[\text{Fe}/\text{H}] \gtrsim -5.0$ , which begin to show more regular halo abundance patterns. We thus adopted  $[\text{Fe}/\text{H}] < -6.3$  as our final upper limit for the iron abundance. This places J0023+0307 firmly among the most iron-poor stars yet discovered. Future additional data on J0023+0307 should confirm this value, or even push it to lower values, as in the case of SMS0313–6707. Following the same argument, Bonifacio et al. (2015) detected no Fe lines in the spectrum of star SDSS J1035+0641, but scaling the very low measured Ca abundance suggests an Fe upper limit of  $[\text{Fe}/\text{H}] < -5.6$  (consistent with their  $1\sigma$  Fe limit).

We also determined additional upper limits for Sc, Cr, Mn, Ni, Sr, and Ba, if values were below  $[\text{X}/\text{H}] < -3.5$ . They are given in Tables 1 and 2. Abundance uncertainties are also listed in Table 2. Mg is very well determined from eight lines, so its uncertainty is significantly smaller than for the other abundances derived from just one line or feature.

#### 4.2. NLTE Abundances

We also calculated non-LTE (NLTE) abundances for J0023+0307 (and the effects of using NLTE over LTE upper limits) for relevant lines of the elements Na, Mg, Al, Si, Ca, and Fe.

**Table 3**  
Model Atoms Used in the NLTE Calculations

Species	$N_{\text{levels}}$	$N_{\text{rad}}^{\text{bb}}$	H Collision Rates Ref.
Na I	139	443	Barklem et al. (2010)
Mg I	229	1629	Belyaev et al. (2012) and Guitou et al. (2015)
Al I	136	223	Belyaev (2013)
Si I	296	9503	Belyaev et al. (2014)
Ca II	69	579	Drawin (1969), $S_{\text{H}} = 0.1$

#### 4.2.1. NLTE Methods

NLTE abundances were computed for Na, Mg, Al, Si, Ca, and Fe using the radiative transfer code `MULTI` in its 2.3 version (Carlsson 1986, 1992) and `MARCS` model atmospheres (Gustafsson et al. 1975, 2008) interpolated to the adopted stellar parameters. Atomic models used in the NLTE calculations for Na, Al, Mg, Si, and Ca were built uniformly using the code `Formato2.0`<sup>9</sup> (T. Merle et al. 2019, in preparation). We provide below a brief description of the atomic data used in the atoms.

#### 4.2.2. Model Atoms

The atomic models for Na, Mg, Al, Si, and Ca used in our NLTE calculations include up-to-date atomic data from various databases. Energy levels and radiative bound-bound (*bb*) transitions for the neutral and first ionized species for each element (except for Ca II, where only the first ionized species levels were included) were extracted from the NIST,<sup>10</sup> VALD3,<sup>11</sup> and Kurucz<sup>12</sup> atomic databases. The number of energy levels and radiative *bb* lines used in each atom are found in Table 3. R-matrix radiative photoionization transitions were also included for all the transitions, when available, from the TOPBASE<sup>13</sup> database.

All levels in our atoms are also coupled via inelastic collisional transitions by electrons and neutral hydrogen atoms. For collisional rates by electrons, we used the empirical impact parameter approximation by Seaton (1962). For inelastic hydrogen collisions, which can have important effects on NLTE abundance calculations and possibly produce large uncertainties if not properly treated (Barklem et al. 2010), we used available rate coefficients computed via *ab initio* quantum calculations for all our atoms, except for Ca II, for which we used the classical Drawin (1969) equation. The rates for Ca II were scaled with a scaling factor  $S_{\text{H}} = 0.1$  following Mashonkina (2013). Regardless, the Ca II K line at 3933 Å used here is only little affected by NLTE (Sitnova et al. 2018). The references for the hydrogen collisional data used for each element species are found in Table 3. For Fe, we used a comprehensive Fe I/Fe II model atom described in detail in Ezzeddine et al. (2016), with hydrogen collision rates incorporated from quantum calculations by Barklem (2018).

#### 4.2.3. NLTE Abundance Results

We determined an Fe NLTE upper limit abundance of  $[\text{Fe}/\text{H}] < -5.2$  for J0023+0307 from an upper limit equivalent width (EW) of  $<10 \text{ m}\text{\AA}$  for the strongest Fe I line at 3859 Å. We also determined a differential abundance correction, NLTE–LTE, a value determined self-consistently in `MULTI2.3`. At this  $[\text{Fe}/\text{H}]$ , the correction is +0.4 dex, which we then also applied to the LTE Fe abundance obtained from the composite spectrum using `MOOG`, leading to  $[\text{Fe}/\text{H}] < -5.4$ . Repeating this for our final upper limit of  $[\text{Fe}/\text{H}] < -6.3$  (for which the correction does not significantly increase), shows that the final NLTE upper limit is  $[\text{Fe}/\text{H}] < -5.9$ . We note that the true iron abundance of J0023+0307 is in all likelihood even lower, which would naturally further increase the NLTE effect and associated correction. By not guessing such a value, we note that we simply produced a more conservative upper limit. However, our calculations do robustly confirm that J0023+0307 has an iron abundance of  $[\text{Fe}/\text{H}] < -5.9$ .

We determined NLTE abundances for the  $\alpha$ -elements:  $[\text{Na}/\text{H}] = -3.92$ ,  $[\text{Mg}/\text{H}] = -2.82$ ,  $[\text{Al}/\text{H}] = -3.56$ ,  $[\text{Si}/\text{H}] = -3.02$ , and  $[\text{Ca}/\text{H}] = -5.72$  from 2 Na I, 8 Mg I, 2 Al I, 1 Si I, and 1 Ca II lines, respectively. NLTE effects, determined self-consistently in `MULTI`, are generally small, to within 0.1 dex for Na I, 0.2 dex for Mg I, and 0.05 dex for Ca II. They are slightly higher for Al I and Si I, up to 0.7 dex and 0.5 dex, respectively. For comparison, we also determined NLTE abundances for Na, Mg, Al, Ca, and Si for HE 1327–2326 (Frebel et al. 2005). NLTE abundances for both J0023+0307 and HE 1327–2326 are found in Table 2.

#### 4.3. The Abundance Pattern

In the following, we focus on discussing  $[\text{X}/\text{H}]$  results because Fe is not detected and its upper limit is very low. However, J0023+0307 does show similarities in its abundances compared to four other stars with  $[\text{Fe}/\text{H}] < -5.0$  (when considering 1D LTE values based on a high-resolution spectrum): SDSS J1035+0641 (Bonifacio et al. 2015), HE 0107–5240 (Christlieb et al. 2002), HE 1327–2327 (Frebel et al. 2005), and SMSS 0313–6707 (Keller et al. 2014), whose abundances are characterized by the extremely low Fe abundances paired with high C, N, and O values. In addition, there are similarities to CS 29498-043 (Aoki et al. 2002), which shares large Mg and Si overabundances, albeit having a much higher iron abundance of  $[\text{Fe}/\text{H}] = -3.75$ . A comparison of the 1D LTE elemental patterns of J0023+0307, CS 29498-043, and HE 1327–2327 is shown in Figure 3.

The abundance of Li ( $A(\text{Li}) = 1.7$ ) is below the value of the Spite Plateau ( $A(\text{Li}) \sim 2.20$ , Sbordone et al. 2010) for J0023+0307 at  $[\text{Fe}/\text{H}] < -6.3$ . The only other warm star with  $[\text{Fe}/\text{H}] < -5.0$  is HE 1327–2326 with  $A(\text{Li}) < 0.60$  (Frebel et al. 2008), but it is worthwhile to also consider SDSS J102915+172927 with  $[\text{Fe}/\text{H}] = -4.7$  and  $A(\text{Li}) < 1.1$  (Cafau et al. 2011) in this context. Both of these have Li abundances far below the Spite Plateau, and thus very different from what is found for J0023+0307. This is shown in Figure 3. The meltdown of the Spite Plateau below  $[\text{Fe}/\text{H}] \lesssim -3.5$  continues to be discussed in the literature (Hansen et al. 2015b; Li et al. 2015; Matsuno et al. 2017; Bonifacio et al. 2018) while explanations for this behavior are sought. J0023+0307 adds to the body of data in an  $[\text{Fe}/\text{H}]$  range that is not well populated

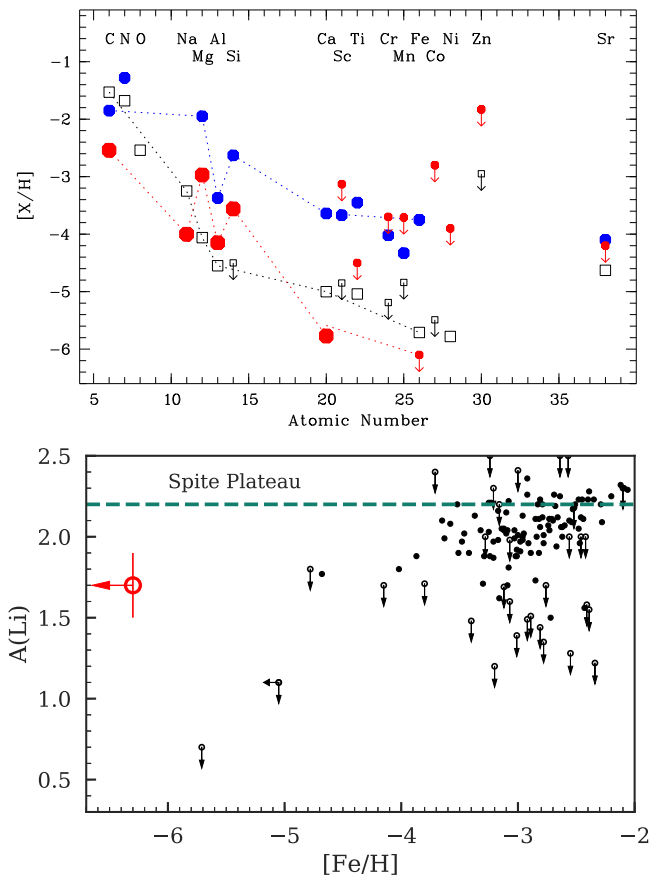
<sup>9</sup> <https://github.com/thibaultmerle/formato>

<sup>10</sup> <https://www.nist.gov/pml/atomic-spectra-database>

<sup>11</sup> <http://vald.astro.uu.se>

<sup>12</sup> <http://kurucz.harvard.edu/linelists.html>

<sup>13</sup> <http://cdsweb.u-strasbg.fr/topbase/topbase.html>



**Figure 3.** Top: abundance patterns of J0023+0307 (red circles), HE 1327–2326 (black squares), and CS 29498-043 (blue circles). The same abundances measured in all stars are connected by dotted lines to indicate the similarity of the elemental patterns. Upper limits are denoted with arrows. Bottom: Li abundances vs.  $[\text{Fe}/\text{H}]$ . Comparison halo stars are compiled from JINABase (Abomalima & Frebel 2017); Aoki et al. (2008, 2009, 2012); Behara et al. (2010); Bonifacio et al. (2012, 2015); Caffau et al. (2011); Frebel et al. (2008); Fulbright (2000); Hansen et al. (2015a); Li et al. (2015); Lucatello et al. (2003); Meléndez et al. (2010); Roederer et al. (2010, 2014); Sivarani et al. (2006); Smiljanic et al. (2009). Stars with  $T_{\text{eff}} < 5900$  K and horizontal branch stars have been excluded. Upper limits are denoted with arrows.

to date. It helps to address the important issue of whether there is a meltdown at the lowest  $[\text{Fe}/\text{H}]$ , or perhaps a more varied behavior with no single, well-defined trend.

The high C abundance of  $[\text{C}/\text{H}] = -2.5$  ( $[\text{C}/\text{Fe}] > 3.0$ ) of J0023+0307 is in line with the increased fraction of carbon-enhanced metal-poor stars with decreasing  $[\text{Fe}/\text{H}]$  (Placco et al. 2014). This means that all five stars with  $[\text{Fe}/\text{H}] < -5.0$  are all strongly carbon enhanced, with  $[\text{C}/\text{Fe}] > 3$ .

Spite et al. (2013), and Bonifacio et al. (2018) using the latest data, found that carbon enhancement at the lowest  $[\text{Fe}/\text{H}]$  values lies between  $A(\text{C}) \sim 5.5$  and 7.5, or  $[\text{C}/\text{H}] \sim -2.9$  and  $-0.9$ . The five stars nicely scatter around this value, with J0023+0307, SMSS 0313–6708, HE 0107–5240, SDSS J1035+0641, and HE 1327–2326 having  $A(\text{C}) = 5.9, 6.0, 6.8, 6.9,$  and  $6.9$ , respectively. In a similar way, Yoon et al. (2016) postulated that at the lowest  $[\text{Fe}/\text{H}]$ , stars should have  $A(\text{C}) > 6.3$ . However, both J0023+0307 and SMSS 0313–6708 with  $[\text{Fe}/\text{H}] < -6.0$  do not show quite this level of enhancement.

High carbon abundances have previously been explained with carbon (and oxygen) being present in large enough

amounts in the natal cloud ( $[\text{C}/\text{H}] > -3.5$ ), as provided by the massive first stars. This would have sufficiently cooled the gas to enable the formation of the first low-mass stars in the universe (Bromm & Loeb 2003; Frebel et al. 2007a). For J0023+0307, this threshold is met even if any systematic 3D effects on the C abundance were of order  $-1$  dex (Collet et al. 2006; Gallagher et al. 2017).

J0023+0307 is the second-most calcium-poor star after SM0313+6708. It also has the highest  $[\text{Mg}/\text{H}]$  abundances of all stars with  $[\text{Fe}/\text{H}] < -4.5$ , and the second-highest  $[\text{Mg}/\text{Fe}]$  ratio of all metal-poor stars, after SM0313+6708. These effects remain when considering our NLTE abundances. Overall, J0023+0307 shows a strong odd–even effect in the Na through Si abundances. There are also differences between the hydrostatic and explosive  $\alpha$ -elements: Mg and Si are significantly enhanced in J0023+0307 ( $[\text{Mg}/\text{H}] \sim -3.0$ ,  $[\text{Si}/\text{H}] \sim -3.5$ ), whereas Ca and Ti abundances are much lower ( $[\text{Ca}/\text{H}] \sim -5.8$ ,  $[\text{Ti}/\text{H}] < -4.5$ ).

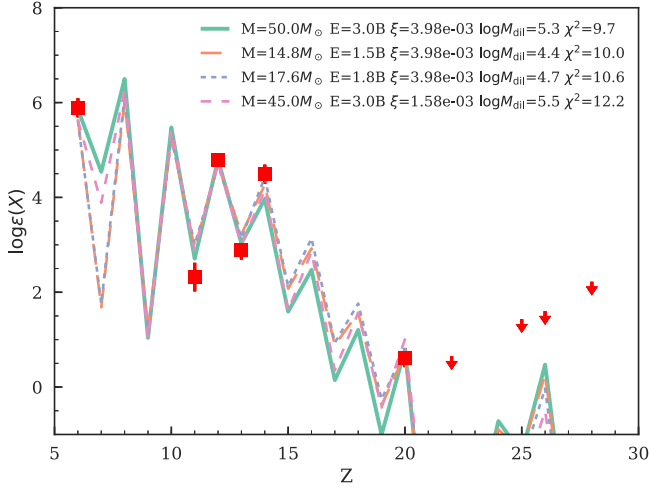
Such strong odd–even abundance ratios are often associated with pair-instability supernovae (PISN, e.g., Heger & Woosley 2002; Aoki et al. 2014). However, the lack of significant amounts of Ca and Fe in J0023+0307 suggests that a PISN could not have been responsible for the observed elements. For completeness, we also note that overall, the large C, Mg, and Si abundances in combination with the nondetection of Fe clearly point to J0023+0307 not being a low-mass ( $< 1 M_{\odot}$ ) Population III star.

High abundances of Mg and Si have also been found for CS 29498-043 (Aoki et al. 2002) and J2217+2104 (Aoki et al. 2018; and also CS 22949-037 but to a lesser extent, Norris et al. 2001; Aoki et al. 2002; Depagne et al. 2002), with a very similar  $\alpha$ -element pattern to J0023+0307. Also, both CS 29498-043 and J0023+0307 are C-enhanced and share a high A abundance (see Figure 3). HE 1327–2326 and SM0313–6708 also display similarly large Mg abundances ( $[\text{Mg}/\text{H}] \sim -4$ ), and much lower Ca ( $[\text{Ca}/\text{H}] \sim -5$  and  $-7$ , respectively), all while being C-enhanced.

Aoki et al. (2018) found that a supernova with a progenitor mass of  $25 M_{\odot}$  best fits the abundance patterns of stars with high Mg and Si abundances, especially J2217+2104. This is, however, not different from any progenitor masses thought to explain the patterns of regular metal-poor stars. Accordingly, they suggested that the unusual overabundances are not related to the mass but to other effects in the progenitors, such as rotation, mixing, and/or fallback. Alternatively, material released through a stellar wind coming off a rotating massive first star might also explain some of these features (Meynet et al. 2006; Choplin et al. 2018). A nitrogen abundance constraint is not possible with our data, but high N could suggest that rotation played a significant role.

Some type of massive progenitor, possibly rotating given the large C abundance, must thus have enriched the natal gas cloud of J0023+0307 through either its supernova yields or a stellar wind, with the possibility of added supernova yields upon explosion. While the total metallicity of J0023+0307 is of course rather high (around  $[\text{M}/\text{H}] \sim -2$  when considering that O and N are likely enhanced at a similar level as C), the low Fe and Ca abundances do clearly point to a single Population III star progenitor, as Population II supernovae will rapidly erase any low-Fe signature (e.g., Ji et al. 2015).





**Figure 4.** The four supernova yield models from Heger & Woosley (2010) (colored lines) that best match the NLTE abundance pattern (red squares). Model parameters are listed.

#### 4.4. Fitting Population III Supernova Yields to the Abundance Pattern

We fit Population III supernova yields from the nonrotating models of Heger & Woosley (2010) to the abundance pattern of J0023+0307. We used NLTE abundances because they are closest to absolute abundances, which are required for a comparison with the yield predictions.

To find the best-fit models, we determined the best scaling factor in  $\log \epsilon(X)$  for all 16,800 models from Heger & Woosley (2010)<sup>14</sup> to minimize  $\chi^2$  for our measured element abundances (C, Na, Mg, Al, Si, Ca), i.e.,

$$\chi^2(M, E, \xi, \delta) = \sum_X \left( \frac{\log \epsilon(X) - \log \mu_{M,E,\xi}(X) + \delta}{\sigma_X} \right)^2, \quad (1)$$

where  $\mu_{M,E,\xi}(X)$  is the model abundance of element  $X$  in units of  $M_\odot/\text{amu}$ , with progenitor mass  $M$ , energy  $E$ , and mixing parameter  $\xi$ ; and  $\sigma_X$  is the abundance uncertainty. Analytically solving for the minimum gives

$$\delta = - \frac{\sum_X \frac{\log \epsilon(X) - \log \mu_{M,E,\xi}(X)}{\sigma_X^2}}{\sum_X \frac{1}{\sigma_X^2}}. \quad (2)$$

Note that  $\delta < 0$ . Since  $\log \epsilon(X) = \log(N_X/N_H) - 12$ , the scaling factor  $\delta$  can also be converted to a dilution mass  $\log M_{\text{dil}}$  of hydrogen implied by the model:

$$\log M_{\text{dil}} = 12 + \delta, \quad (3)$$

where the unit of  $M_{\text{dil}}$  is  $M_\odot$ , and we use the fact that the atomic mass of hydrogen is 1 amu. It is unphysical for this dilution mass to be less than the supernova ejecta mass, so we rejected all models with  $\log M_{\text{dil}} < 2$ , as well as all models conflicting with our upper limits (Ti, Mn, Fe, Ni).

The best four models are shown in Figure 4. Qualitatively, the Heger & Woosley (2010) models can reasonably reproduce the abundances of J0023+0307, fitting the high carbon, the strong odd–even ratios from Na to Si, the very low Ca, and the

low upper limits of heavier elements like Fe. Several progenitor masses appear to fit almost equally well, but the energies are moderate (1–3 B) and mixing parameters  $\xi$  are low. J0023+0307 thus appears to be well described by a fallback-with-little-mixing supernova. Note that we could not measure N, but it is clear that the supernova yields vary greatly in their predicted N abundances, so even a loose constraint on N in the future would be very helpful (also see Placco et al. 2015b). The N abundance would also be useful for determining whether the progenitor star was rotating, because the rotation should greatly enhance N (Meynet et al. 2006; Choplin et al. 2018).

The four best-fitting models do not fully represent the complete set of well-fitting models, nor all the parameter degeneracies (e.g., Placco et al. 2015b; Nordlander et al. 2017). Thus, in Figure 5, we show the best-fit parameters of all models within  $\Delta\chi^2 < 9.7$  of the lowest  $\chi^2$ . This corresponds to a confidence level of 95% (or  $2\sigma$ ), assuming four parameters are fit ( $M$ ,  $E$ ,  $\xi$ , and  $M_{\text{dil}}$ ). To draw the eye, the size of the point is larger for models that fit better (i.e., lower  $\chi^2$ ). It is immediately clear that the progenitor mass is not well constrained, because essentially all masses from 10 to  $100 M_\odot$  have a model consistent within  $2\sigma$  of the measured abundances. However, the other parameters are better constrained: energies are mostly in the range  $(1-3) \times 10^{51}$  erg, the mixing parameters are very low, and the dilution masses are in the range  $\sim 10^{4.5-5.5} M_\odot$ . There are also correlations between mass, explosion energy, and dilution mass. Since the explosion energy and mixing are not true free parameters but expressions of uncertain supernova explosion physics, more realistic 3D simulations (e.g., Chan et al. 2018) may eventually be able to break these degeneracies and constrain the progenitor masses.

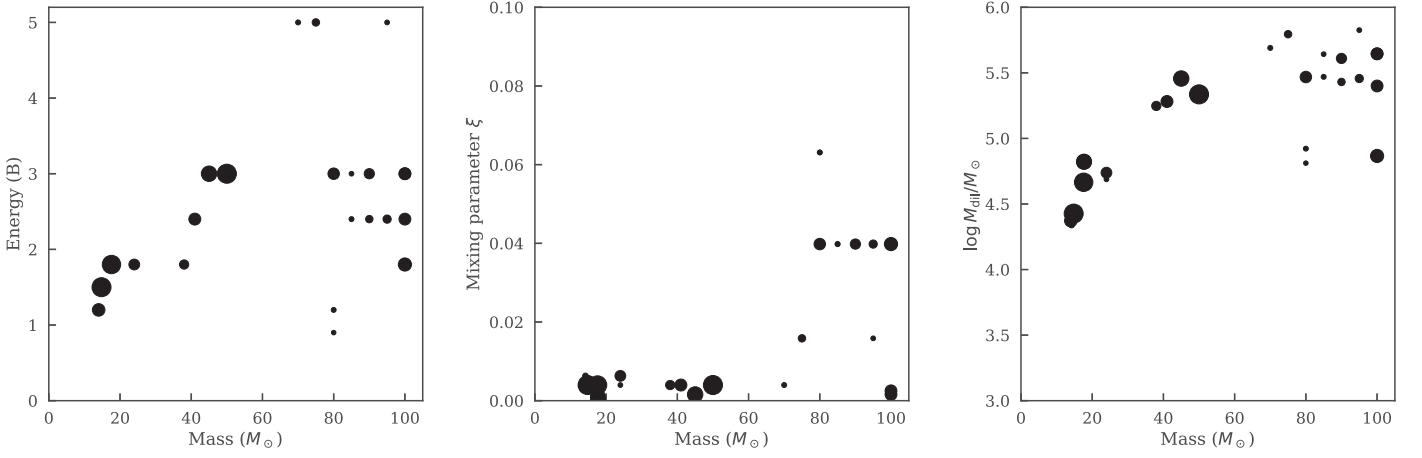
The dilution masses inferred here strongly suggest that J0023+0307 is a second-generation star formed by recollapse in a Population III minihalo (Ritter et al. 2012; Cooke & Madau 2014; Ji et al. 2015). In this model, moderate energy Population III supernovae occurring in  $10^6 M_\odot$  dark matter minihalos do not evacuate the minihalo of gas; instead, some fraction of the ejected metals recollapses into the same minihalo, with typical effective dilution masses  $\sim 10^6 M_\odot$ . The most likely alternative scenarios for second-generation star formation are either external pollution of another minihalo (e.g., Smith et al. 2015; Griffen et al. 2018) or delayed second-generation star formation in an atomic cooling halo (e.g., Wise & Abel 2008; Greif et al. 2010). In either case, we should expect the dilution masses to be much larger,  $> 10^7 M_\odot$ , although for the atomic cooling halo multiple SNe will have contributed metals to the nascent galaxy. We also note that our somewhat uncertain Na detection has a significant effect on the preferred progenitor masses. Removing the Na constraint causes the fit to prefer lower-mass (10–20  $M_\odot$ ) progenitors, while including Na pushes the fit toward more-massive progenitors (also see Ishigaki et al. 2018).

## 5. Kinematic Signature

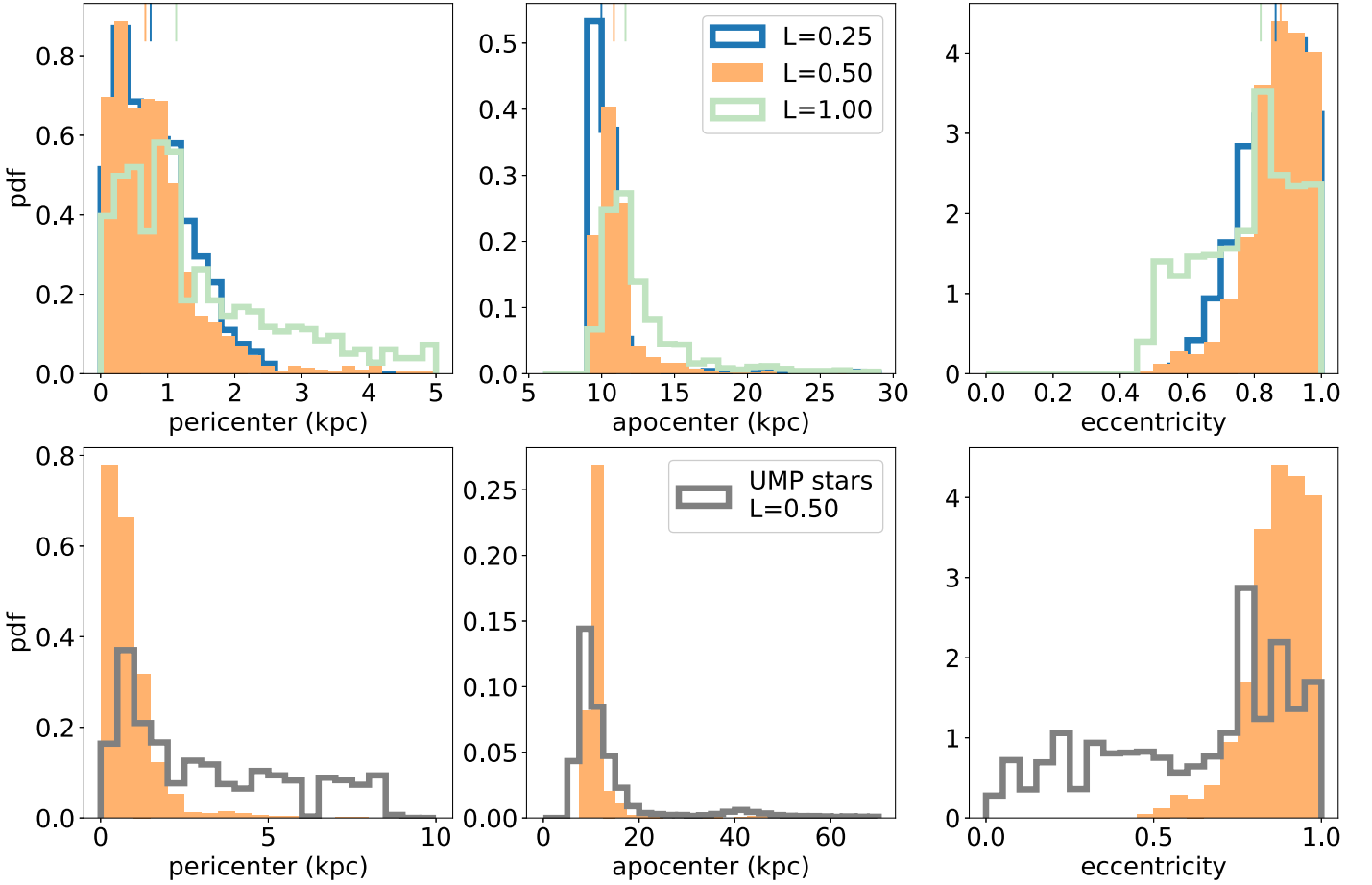
We investigated the Galactic orbit of J0023+0307 using the parallax and proper motion from *Gaia* DR2 (Gaia Collaboration et al. 2018). Following Bailer-Jones et al. (2018) and its extension to proper motions,<sup>15</sup> we simultaneously sampled the distance and tangential velocity posterior with *emcee*

<sup>14</sup> Table `znuc.S4.star.el.y` from [www.2sn.org/starfit](http://www.2sn.org/starfit); note that the Starfit algorithm differs from ours in its treatment of upper limits.

<sup>15</sup> [https://github.com/agabrown/astrometry-inference-tutorials/blob/master/3d-distance/resources/3D\\_astrometry\\_inference.pdf](https://github.com/agabrown/astrometry-inference-tutorials/blob/master/3d-distance/resources/3D_astrometry_inference.pdf)



**Figure 5.** Supernova model parameters of all fits within  $2\sigma$  of the data (see the text for details). Larger points indicate better-fitting models, to draw the eye. The progenitor masses are not constrained, but the energies are moderate (1–3 B), the mixing is low, and the dilution masses range from  $10^{4.5-6.0} M_{\odot}$ . Degeneracies exist between the progenitor mass and the other parameters.



**Figure 6.** Top row: posterior distributions for pericenter, apocenter, and eccentricity in J0023+0307. Different colors indicate different assumptions for the distance prior scale  $L$ . Lines at top indicate median value. Bottom row: showing comparison with orbits of 25 ultra-metal-poor (UMP) stars (gray) assuming  $L = 0.50$  (orange, same data as top panel with different binning). J0023+0307 clearly passes through the bulge at pericenter, and is on a relatively eccentric orbit compared with other UMP stars.

(Foreman-Mackey et al. 2013b). For the distance, we used an exponentially decreasing volume density prior (Bailer-Jones et al. 2018). The scale parameter is varied from  $L = 250$  pc to  $L = 1000$  pc, where  $L = 372$  pc is the value used by Bailer-Jones et al. (2018) based on a magnitude-limited mock survey of a simulated Milky Way. The total tangential velocity prior is a beta distribution from 0 to  $1800 \text{ km s}^{-1}$  with  $\alpha = 2$  and

$\beta = 8$ . The parameters were chosen to peak at  $\sim 180 \text{ km s}^{-1}$ , but changing this prior makes little difference to final results. Radial velocities were assumed to be normally distributed around  $-195 \text{ km s}^{-1}$  with a standard deviation of  $5 \text{ km s}^{-1}$ . We initialized 400 walkers using the *Gaia* covariance matrix and assuming  $d = 1/\varpi$ ; burned in 500 steps; then ran a chain of length 2500. Visual examination showed that the chains

**Table 4**  
UMP Star Kinematic Data

Star	RA h:m:s	Decl. d:m:s	$v_{\text{hel}}$ ( $\text{km s}^{-1}$ )	Distance (kpc)	$v_{\alpha}$ ( $\text{km s}^{-1}$ )	$v_{\delta}$ ( $\text{km s}^{-1}$ )	Pericenter (kpc)	Apocenter (kpc)	Eccentricity
SDSSJ0023+0307	00:23:14.0	+03:07:58.1	-195	$2.80_{-0.65}^{+0.86}$	$49.7_{-11.7}^{+15.4}$	$-184.8_{-57.4}^{+42.4}$	$0.67_{-0.45}^{+0.61}$	$10.84_{-0.91}^{+1.24}$	$0.88_{-0.09}^{+0.08}$
HE 0233-0343	02:36:29.8	-03:30:06.0	+64	$1.27_{-0.09}^{+0.10}$	$301.9_{-20.2}^{+23.5}$	$-64.0_{-4.8}^{+4.2}$	$0.30_{-0.20}^{+0.40}$	$13.40_{-0.76}^{+4.31}$	$0.95_{-0.05}^{+0.03}$
CS 22963-004	02:56:46.6	-04:51:17.5	+294	$4.01_{-0.40}^{+0.57}$	$412.8_{-41.2}^{+58.1}$	$-50.4_{-7.5}^{+5.4}$	$1.80_{-0.81}^{+1.27}$	$72.82_{-19.11}^{+59.37}$	$0.96_{-0.01}^{+0.01}$
HE 0557-4840	05:58:39.3	-48:39:56.8	+212	$9.90_{-1.28}^{+1.67}$	$33.5_{-4.9}^{+6.3}$	$34.0_{-4.9}^{+6.5}$	$2.94_{-0.29}^{+0.44}$	$14.62_{-1.09}^{+1.43}$	$0.66_{-0.02}^{+0.02}$
HE 1012-1540	10:14:53.5	-15:55:53.2	+226	$0.39_{-0.00}^{+0.00}$	$-191.0_{-2.2}^{+2.2}$	$52.5_{-0.6}^{+0.6}$	$0.88_{-0.07}^{+0.08}$	$13.58_{-0.15}^{+0.14}$	$0.88_{-0.01}^{+0.01}$
HE 1310-0536 <sup>a</sup>	13:13:31.2	-05:52:12.5	+113	$8.61_{-1.44}^{+1.85}$	$-198.1_{-63.5}^{+38.7}$	$-66.4_{-23.8}^{+13.2}$	...	...	...
SDSS J1808-5104	18:08:20.0	-51:04:37.9	+16	$0.60_{-0.01}^{+0.01}$	$-15.9_{-0.4}^{+0.4}$	$-35.8_{-1.0}^{+0.9}$	$5.19_{-0.11}^{+0.11}$	$7.82_{-0.03}^{+0.04}$	$0.20_{-0.01}^{+0.01}$
CS 22891-200	19:35:19.1	-61:42:24.4	+138	$6.05_{-0.83}^{+1.05}$	$-146.3_{-23.6}^{+19.7}$	$22.0_{-3.0}^{+3.7}$	$0.77_{-0.54}^{+0.32}$	$10.22_{-0.49}^{+1.63}$	$0.86_{-0.05}^{+0.10}$
CS 22885-096	20:20:51.2	-39:53:30.2	-249	$5.31_{-0.55}^{+0.59}$	$-111.6_{-12.8}^{+10.9}$	$-173.6_{-20.1}^{+17.9}$	$4.60_{-0.21}^{+0.05}$	$7.53_{-0.36}^{+0.48}$	$0.25_{-0.03}^{+0.04}$
CS 22950-046	20:21:28.4	-13:16:33.6	+107	$8.04_{-1.07}^{+1.35}$	$57.9_{-8.5}^{+10.1}$	$-70.4_{-11.3}^{+9.1}$	$2.82_{-0.77}^{+0.74}$	$10.96_{-0.97}^{+1.45}$	$0.59_{-0.11}^{+0.13}$
CS 30336-049	20:45:23.5	-28:42:35.9	-237	$9.05_{-1.37}^{+1.07}$	$-73.8_{-10.7}^{+9.1}$	$-348.6_{-53.2}^{+43.8}$	$2.91_{-0.15}^{+0.38}$	$9.62_{-2.19}^{+2.19}$	$0.55_{-0.15}^{+0.10}$
HE 2239-5019	22:42:26.8	-50:04:00.9	+370	$3.70_{-0.57}^{+0.69}$	$135.6_{-20.8}^{+25.3}$	$-415.8_{-77.0}^{+63.6}$	$7.03_{-0.07}^{+0.02}$	$34.98_{-9.89}^{+30.02}$	$0.66_{-0.10}^{+0.14}$
CS 22949-037 <sup>a</sup>	23:26:29.8	-02:39:57.9	-125	$8.96_{-1.75}^{+1.91}$	$69.6_{-14.6}^{+53.66}$	$-74.2_{-49.57}^{+15.0}$	...	...	...
BD+44° 493	02:26:49.7	+44:57:46.5	-151	$0.21_{-0.00}^{+0.00}$	$118.0_{-1.5}^{+1.6}$	$-32.1_{-0.5}^{+0.4}$	$1.01_{-0.07}^{+0.08}$	$8.60_{-0.03}^{+0.04}$	$0.79_{-0.02}^{+0.01}$
CD-38° 245	00:46:36.2	-37:39:33.5	+46	$4.16_{-0.44}^{+0.64}$	$300.1_{-31.4}^{+46.4}$	$-148.7_{-22.6}^{+15.5}$	$1.19_{-0.76}^{+1.26}$	$11.83_{-0.98}^{+2.47}$	$0.82_{-0.12}^{+0.11}$
HE 0057-5959	00:59:54.1	-59:43:30.0	+375	$4.70_{-0.46}^{+0.57}$	$53.2_{-5.2}^{+6.6}$	$-234.8_{-28.1}^{+23.6}$	$8.07_{-0.24}^{+0.27}$	$16.38_{-1.23}^{+1.76}$	$0.34_{-0.02}^{+0.03}$
HE 0107-5240	01:09:29.2	-52:24:34.2	+44	$7.55_{-1.00}^{+1.24}$	$86.6_{-11.8}^{+14.0}$	$-134.3_{-21.7}^{+17.1}$	$1.38_{-0.64}^{+0.62}$	$10.48_{-0.60}^{+0.78}$	$0.77_{-0.10}^{+0.11}$
HE 1327-2326	13:30:06.0	-23:41:49.7	+112	$1.13_{-0.03}^{+0.03}$	$-280.7_{-8.0}^{+6.7}$	$243.2_{-6.0}^{+6.0}$	$5.60_{-0.06}^{+0.06}$	$41.96_{-3.38}^{+4.52}$	$0.76_{-0.02}^{+0.02}$
HE 2139-5432	21:42:42.5	-54:18:43.0	+116	$9.16_{-1.23}^{+1.54}$	$107.9_{-14.9}^{+18.5}$	$-195.4_{-33.5}^{+26.5}$	$0.63_{-0.41}^{+0.59}$	$8.11_{-0.85}^{+1.95}$	$0.86_{-0.14}^{+0.09}$
LAMOSTJ 1253+0753	12:53:46.1	+07:53:43.1	+78	$0.71_{-0.02}^{+0.02}$	$71.2_{-1.9}^{+2.2}$	$-198.6_{-6.0}^{+5.3}$	$1.66_{-0.09}^{+0.09}$	$12.13_{-0.17}^{+0.20}$	$0.76_{-0.01}^{+0.01}$
SMSSJ 0313-6708	03:13:00.4	-67:08:39.0	+300	$8.18_{-0.85}^{+0.86}$	$272.2_{-28.4}^{+28.8}$	$41.5_{-4.7}^{+5.2}$	$0.92_{-0.11}^{+0.19}$	$16.39_{-1.98}^{+2.72}$	$0.89_{-0.03}^{+0.02}$
SDSSJ 1204+1201	12:04:41.4	+12:01:11.5	+51	$3.64_{-0.64}^{+0.90}$	$6.9_{-2.2}^{+2.8}$	$-85.1_{-21.8}^{+14.3}$	$4.39_{-0.56}^{+0.42}$	$9.84_{-0.37}^{+0.58}$	$0.38_{-0.06}^{+0.08}$
SDSSJ 1313-0019	13:13:26.9	-00:19:41.5	+267	$3.08_{-0.63}^{+0.83}$	$-55.1_{-15.4}^{+11.1}$	$-95.7_{-25.5}^{+19.1}$	$6.87_{-0.18}^{+0.23}$	$9.08_{-0.62}^{+0.83}$	$0.14_{-0.04}^{+0.04}$
SDSSJ 1442-0015	14:42:56.4	-00:15:42.8	+225	$3.50_{-1.04}^{+1.46}$	$-4.4_{-5.9}^{+4.7}$	$110.4_{-33.6}^{+47.9}$	$6.68_{-0.34}^{+0.19}$	$31.93_{-0.07}^{+17.38}$	$0.65_{-0.07}^{+0.12}$
SDSSJ 1029+1729	10:29:15.2	+17:29:27.9	-34	$1.41_{-0.13}^{+0.18}$	$-72.7_{-9.3}^{+7.3}$	$-27.0_{-3.4}^{+2.5}$	$8.03_{-0.24}^{+0.23}$	$9.03_{-0.09}^{+0.14}$	$0.06_{-0.02}^{+0.02}$
HE 1424-0241	14:26:40.3	-02:54:27.5	+60	$5.42_{-0.93}^{+1.16}$	$-97.9_{-21.2}^{+16.4}$	$-72.8_{-16.2}^{+12.8}$	$2.51_{-0.63}^{+0.51}$	$6.79_{-0.04}^{+0.19}$	$0.46_{-0.08}^{+0.12}$
SDSSJ 1035+0641	10:35:56.1	+06:41:44.0	-78	$3.49_{-1.09}^{+1.42}$	$59.9_{-19.0}^{+26.6}$	$-47.7_{-18.4}^{+14.0}$	$6.50_{-0.65}^{+0.52}$	$19.08_{-2.35}^{+4.08}$	$0.49_{-0.08}^{+0.10}$
SDSSJ 1742+2531	17:42:59.7	+25:31:35.9	-208	$3.65_{-0.81}^{+1.26}$	$-109.5_{-36.6}^{+24.3}$	$-198.4_{-72.2}^{+44.3}$	$0.90_{-0.36}^{+0.18}$	$7.42_{-0.17}^{+1.75}$	$0.79_{-0.04}^{+0.10}$
Pristine_221	14:47:30.7	+09:47:03.7	-149	$4.04_{-0.73}^{+1.03}$	$-149.0_{-38.6}^{+27.8}$	$-1.0_{-2.2}^{+2.1}$	$3.46_{-0.27}^{+0.35}$	$9.54_{-0.19}^{+0.40}$	$0.47_{-0.05}^{+0.04}$

**Note.** Estimate is median of 1000 samples; uncertainty is 16–84th percentiles. All values assume distance prior with  $L = 500$  pc.

<sup>a</sup> Not used due to large uncertainty.

appeared converged. We then integrated  $10^3$  random samples with gala backward for 3 Gyr in the default Milky Way potential (Bovy 2015; Price-Whelan 2017) to derive pericenters, apocenters, and eccentricities. The posterior distributions for the three priors are shown in Figure 6.

For comparison, we performed the same analysis for other ultra-metal-poor (UMP,  $[\text{Fe}/\text{H}] < -4.0$ ) stars. We obtained our list of UMP stars from Abohalima & Frebel (2017), supplementing with stars referenced in Ezzeddine et al. (2017) and Starkenburg et al. (2018). After removing two stars with poor parallaxes, our final list contained 26 stars. Some of these stars are probably binaries, but we did not anticipate radial velocity variations to significantly affect the results for the distribution because the total velocities of most stars are usually relatively high ( $>100 \text{ km s}^{-1}$ ). To somewhat account for this, we included a  $5 \text{ km s}^{-1}$  scatter on the radial velocity for all stars. We sampled the kinematic posterior, integrated 1000 orbits, and show the summed posterior of all stars as a gray histogram in the bottom row of Figure 6. Per-star results are in Table 4.

J0023+0307 is clearly on an eccentric orbit ( $e > 0.8$ ), with a clear pericenter in the bulge ( $\lesssim 1$  kpc). Its eccentricity is higher and its pericenter smaller than that of a typical  $[\text{Fe}/\text{H}] < -4.0$  star. Metal-poor stars from the bulge are thought to be

older in absolute age (Tumlinson 2010; Howes et al. 2015; Starkenburg et al. 2017; El-Badry et al. 2018). J0023+0307 is thus likely one of the oldest stars known in the Milky Way. Note that of the other two stars with no detected iron, SMSS 0313-6707 also has a pericenter  $< 1$  kpc and is probably also old; however, SDSS J1035+0641 does not, with a pericenter at 6.5 kpc. The highly eccentric radial orbit suggests a possible association with the recently discovered *Gaia* sausage/blob structure (Belokurov et al. 2018; Koppelman et al. 2018), although such orbits could also just reflect typical virialized halo star orbits. We caution that our simple orbit integrations do not account for effects like the Galactic bar, which can significantly influence halo star orbits (e.g., Price-Whelan et al. 2016).

## 6. Summary and Conclusions

We have presented a high-resolution spectrum of the iron-poor main-sequence star J0023+0307. Iron lines were not detected, but the S/N of the spectrum is only moderate. Stacking 19 Fe lines into a composite spectrum did not yield a detection either, but led to an upper limit of  $[\text{Fe}/\text{H}] < -5.8$ . Invoking also that the Ca abundances in the most iron-poor stars always exceed the Fe abundance, we deduced an upper

limit of  $[\text{Fe}/\text{H}] < -6.3$ . Lines of other elements were detected, i.e., Li, C, Na, Mg, Al, Si, and Ca. Lithium was challenging to determine, given the somewhat distorted line. We thus cannot draw any firm conclusions on the Li abundance, beyond suggesting that the two warm stars with  $[\text{Fe}/\text{H}] < -5.0$  both have values lower than the Spite Plateau. However, the range covered by those low Li abundances is large ( $\sim 1$  dex), adding to the complex behavior of Li found in the most iron-poor stars. Higher-quality data would be able to refine the Li abundance to gain more insight into the evolution of Li in the early universe.

There is a strong odd-even effect in the abundances from Na to Si. J0023+0307 has the second-lowest Ca abundance of any known star after SMSS0313-6708, by far the highest  $[\text{Mg}/\text{Fe}]$  ratio of any metal-poor stars (together with SMSS0313-6708), and among the highest  $[\text{Mg}/\text{H}]$  of stars with  $[\text{Fe}/\text{H}] < -4.0$ . Carbon was detected and strongly C-enhanced, in line with values found for the other most iron-poor stars. This likely points to fragmentation from early gas cooling by C and O provided by Population III first stars.

The overall metallicity of J0023+0307 is  $[\text{M}/\text{H}] \sim -2.0$ , which is largely driven by the large C (and presumably large N and O) abundance. Nevertheless, the low Fe and Ca abundances clearly point to J0023+0307 being a second-generation star that formed from gas enriched by just one Population III supernova. Fitting the abundance pattern with nucleosynthesis yields of the first stars shows that the mass of the progenitor star is essentially unconstrained given current abundance uncertainties, but intermediate explosion energies of  $(1-3) \times 10^{51}$  erg and a low mixing parameter are likely. The inferred dilution masses from these fits are  $10^{4.5-5.5} M_{\odot}$  of hydrogen, as is expected if J0023+0307 formed in a recollapsed minihalo. J0023+0307 has a rather eccentric orbit ( $e > 0.8$ ), which is among the more eccentric values of 25 stars with  $[\text{Fe}/\text{H}] < -4.0$ . The pericenter takes the star right through the Galactic bulge, suggesting that J0023+0307 is one of the oldest stars.

Additional data should be sought to attempt a detection of Fe lines and to produce even tighter upper limits on various undetected elemental abundances. This will be important for further constraining the nature of the progenitor of J0023+0307. More data would also assist in confirming the Li abundance to firmly establish the range of observed (i.e., physically possible) Li abundances at  $[\text{Fe}/\text{H}] < -5.0$ . Thus far, no radial velocity variations have been detected; future measurements would also help to further constrain the nature and formation scenario of J0023+0307.

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*Facilities:* Magellan-Clay (MIKE, Bernstein et al. 2003).

*Software:* MOOG (Snedden 1973; Sobeck et al. 2011), MIKE Carnegie Python Pipeline (Kelson 2003), MULTI (Carlsson 1986, 1992), MARCS (Gustafsson et al. 1975, 2008), IRAF (Tody 1986, 1993), NumPy (van der Walt et al. 2011), SciPy (Jones et al. 2001), Matplotlib (Hunter 2007), Astropy (Astropy Collaboration et al. 2013; The Astropy Collaboration et al. 2018), emcee (Foreman-Mackey et al. 2013a, 2013b), gala (Price-Whelan 2017; Price-Whelan et al. 2017).

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