

2.29 Numerical Fluid MechanicsFall 2009 – Lecture 13

REVIEW Lecture 12:

- Classification of Partial Differential Equations (PDEs) and examples with finite difference discretizations
	- Parabolic PDEs
	- Elliptic PDEs
	- Hyperbolic PDEs

• Error Types and Discretization Properties: $L(\phi) = 0$, $\hat{L}_{\alpha x}(\hat{\phi}) = 0$

- *L* ˆ $-$ Consistency: $\vert L(\phi)-L_{\Lambda x}(\phi)\vert\rightarrow 0\;\;\text{ when }\Delta x\rightarrow 0$
-
-
-
-

– Truncation error: $\tau_{\Delta x} = L(\phi) - \hat{L_{\Delta x}}(\phi) \rightarrow O(\Delta x^p)$ for $\Delta x \rightarrow 0$ – Error equation: $\tau_{\lambda x} = L(\phi) - \hat{L}_{\lambda x}(\hat{\phi} + \varepsilon) = -\hat{L}_{\lambda x}(\varepsilon)$ (for linear systems) ˆ $\|L_{\Delta x}^{\gamma-1}\|$ < Const. (for linear systems) ˆ $\|\mathcal{E}\| \leq \|L_{\Delta x}^{-1}\| \cdot \|\tau_{\Delta x}\| \leq \alpha \ O(\Delta x^p)$

Numerical Fluid Mechanics **PEJL Lecture 13, 1**

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REVIEW Lecture 12, Cont'd:

- Classification of PDEs and examples
- Error Types and Discretization Properties
- Finite Differences based on Taylor Series Expansions
	- Higher Order Accuracy Differences, with Examples
		- Incorporate more higher-order terms of the Taylor series expansion than strictly needed and express them as finite differences themselves (making them function of neighboring function values)
		- If these finite-differences are of sufficient accuracy, this pushes the remainder to higher order terms => increased order of accuracy of the FD method
		- General approximation:

$$
\left(\frac{\partial^m u}{\partial x^m}\right)_j - \sum_{i=-r}^s a_i u_{j+i} = \tau_{\Delta x}
$$

- Taylor Tables or Method of Undetermined Coefficients (Polynomial Fitting)
	- Simply a more systematic way to solve for coefficients a_i

FINITE DIFFERENCES – Outline for Today

- Classification of Partial Differential Equations (PDEs) and examples with finite difference discretizations (Elliptic, Parabolic and Hyperbolic PDEs)
- \bullet Error Types and Discretization Properties
	- Consistency, Truncation error, Error equation, Stability, Convergence
- • Finite Differences based on Taylor Series Expansions
	- Higher Order Accuracy Differences, with Example
	- Taylor Tables or Method of Undetermined Coefficients (Polynomial Fitting)
- \bullet Polynomial approximations
	- Newton's formulas
	- Lagrange polynomial and un-equally spaced differences
	- Hermite Polynomials and Compact/Pade's Difference schemes
	- Boundary conditions
	- Un-Equally spaced differences
	- Error Estimation: order of convergence, discretization error, Richardson's extrapolation, and iterative improvements using Roomberg's algorithm

- Part 8 (PT 8.1-2), Chapter 23 on "Numerical Differentiation" and Chapter 18 on "Interpolation" of "Chapra and Canale, Numerical Methods for Engineers, 2010/2006."
- Chapter 3 on "Finite Difference Methods" of "J. H. Ferziger and M. Peric, Computational Methods for Fluid Dynamics. Springer, NY, 3rd edition, 2002"
- Chapter 3 on "Finite Difference Approximations" of "H. Lomax, T. H. Pulliam, D.W. Zingg, Fundamentals of Computational Fluid Dynamics (Scientific Computation). Springer, 2003"

FINITE DIFFERENCES: Interpolation Formulas for Higher Order Accuracy

2nd approach: Generalize Taylor series using interpolation formulas

- Fit the unknown function solution of the (P)DE to an interpolation curve and differentiate the resulting curve. For example:
	- Fit a parabola to data at points x_{i-1}, x_i, x_{i+1} $(\Delta x_i = x_i x_{i-1})$ differentiate to obtain:

$$
f'(x_i) = \frac{f(x_{i+1})(\Delta x_i)^2 - f(x_{i-1})(\Delta x_{i+1})^2 + f(x_i) \left[(\Delta x_{i+1})^2 - (\Delta x_i)^2 \right]}{\Delta x_{i+1} \Delta x_i (\Delta x_i + \Delta x_{i+1})}
$$

- This is a 2^{nd} order approximation
- For uniform spacing, reduces to centered difference seen before
- In general, approximation of first derivative has truncation error of the order of the polynomial
- All types of polynomials or numerical differentiation methods can be used to derive such interpolations formulas
	- Polynomial fitting, Method of undetermined coefficients, Newton's interpolating polynomials, Lagrangian and Hermite Polynomials, etc

Taylor Tables: Convenient way of forming linear combinations of Taylor Series on a term-by-term basis

The Taylor table for a centered three point Lagrangian approximation to a second derivative.

What we are
\nlooking for
\n
$$
\left(\frac{\partial^2 u}{\partial x^2}\right)_j - \frac{1}{\Delta x^2} (a u_{j-1} + b u_j + c u_{j+1}) = ?
$$
\nlooking for
\n
$$
u_j \quad \Delta x \left(\frac{\partial u}{\partial x}\right)_j \quad \Delta x^2 \left(\frac{\partial^2 u}{\partial x^2}\right)_j \quad \Delta x^3 \left(\frac{\partial^3 u}{\partial x^3}\right)_j \quad \Delta x^4 \left(\frac{\partial^4 u}{\partial x^4}\right)_j
$$
\n
\nTaylor
\n
$$
\Delta x^2 \left(\frac{\partial^2 u}{\partial x^2}\right)_j
$$
\n
$$
j = 1 \quad -a \cdot u_{j-1} \quad -a \quad -a \cdot (-1) \cdot \frac{1}{1} \quad -a \cdot (-1)^2 \cdot \frac{1}{2} \quad -a \cdot (-1)^3 \cdot \frac{1}{6} \quad -a \cdot (-1)^4 \cdot \frac{1}{24}
$$
\n
$$
j = -b \cdot u_j \quad -b
$$
\n
$$
j = -b \cdot u_j \quad -c \quad -c \cdot (1) \cdot \frac{1}{1} \quad -c \cdot (1)^2 \cdot \frac{1}{2} \quad -c \cdot (1)^3 \cdot \frac{1}{6} \quad -c \cdot (1)^4 \cdot \frac{1}{24}
$$

Sum each column starting from left, force the sums to zero and so choose a, b, c, etc

FINITE DIFFERENCESHigher Order Accuracy: Taylor Tables Cont'd

The Taylor table for a centered three point Lagrangian approximation to a second derivative.

$$
\left(\frac{\partial^2 u}{\partial x^2}\right)_j - \frac{1}{\Delta x^2} (a u_{j-1} + b u_j + c u_{j+1}) = ?
$$

Sum each column starting from left and force the sums to be zero by proper choice of a, b, c, etc:

$$
\begin{bmatrix} -1 & -1 & -1 \ 1 & 0 & -1 \ -1 & 0 & -1 \ \end{bmatrix} \begin{bmatrix} a \ b \ c \end{bmatrix} = \begin{bmatrix} 0 \ 0 \ -2 \end{bmatrix} \implies [a \quad b \quad c] = \begin{bmatrix} 1 & -2 & 1 \end{bmatrix} \quad \text{= Familiar 3-point central difference} \\ \text{central difference}
$$

Truncation error is first column in the table that does not vanish, here fifth column of table:

$$
\tau_{\Delta x} = \frac{1}{\Delta x^2} \left[\frac{-a}{24} + \frac{-c}{24} \right] \Delta x^4 \left(\frac{\partial^4 u}{\partial x^4} \right)_j = -\frac{\Delta x^2}{12} \left(\frac{\partial^4 u}{\partial x^4} \right)_j
$$

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FINITE DIFFERENCESHigher Order Accuracy: Taylor Tables Cont'd

The Taylor table for a backward three point Lagrangian approximation to a second derivative.

$$
\left(\frac{\partial u}{\partial x}\right)_j - \frac{1}{\Delta x} (a_2 u_{j-2} + a_1 u_{j-1} + b u_j) = ?
$$
\n
$$
u_j \qquad \Delta x \left(\frac{\partial u}{\partial x}\right)_j \qquad \Delta x^2 \left(\frac{\partial^2 u}{\partial x^2}\right)_j \qquad \Delta x^3 \left(\frac{\partial^3 u}{\partial x^3}\right)_j \qquad \Delta x^4 \left(\frac{\partial^4 u}{\partial x^4}\right)_j
$$
\n
$$
\Delta x \left(\frac{\partial u}{\partial x}\right)_j \qquad 1
$$
\n
$$
-a_2 \cdot u_{j-2} \qquad -a_2 \qquad -a_2 \cdot (-2) \cdot \frac{1}{1} \qquad -a_2 \cdot (-2)^2 \cdot \frac{1}{2} \qquad -a_2 \cdot (-2)^3 \cdot \frac{1}{6} \qquad -a_2 \cdot (-2)^4 \cdot \frac{1}{24}
$$
\n
$$
-a_1 \cdot u_{j-1} \qquad -a_1 \qquad -a_1 \cdot (-1) \cdot \frac{1}{1} \qquad -a_1 \cdot (-1)^2 \cdot \frac{1}{2} \qquad -a_1 \cdot (-1)^3 \cdot \frac{1}{6} \qquad -a_1 \cdot (-1)^4 \cdot \frac{1}{24}
$$
\n
$$
\frac{-b \cdot u_j \qquad -b}{\Rightarrow \quad [a_2 \quad a_1 \quad b] = [1 \quad -4 \quad 3]/2 \qquad \text{and} \qquad \tau_{xx} = \frac{1}{\Delta x} \left[\frac{8a_2}{6} + \frac{a_1}{6}\right] \Delta x^3 \left(\frac{\partial^3 u}{\partial x^3}\right)_j = \frac{\Delta x^2}{3} \left(\frac{\partial^3 u}{\partial x^3}\right)_j
$$
\n(as in lecture 12) Numerical Fluid Mechanics

Finite Differences using Polynomial approximations Numerical Interpolation: Newton's Iteration Formula

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Finite Differences using Polynomial approximations Equidistant Newton's Interpolation

Equidistant Sampling **Divided Differences**

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Numerical Differentiation using Newton's algorithm for equidistant sampling

x

n=1

Numerical Differentiation using Newton's algorithm for equidistant sampling, Cont'd

Second order

 $n=2$

$$
f(x) = f_0 + \frac{\Delta f_0}{h}(x-x_0) + \frac{\Delta^2 f_0}{2!h^2}(x-x_0)(x-x_1) + \frac{f'''(\xi)}{3!}(x-x_0)(x-x_1)(x-x_2) + \cdots
$$

$$
f'(x) = \frac{\Delta f_0}{h} + \frac{\Delta^2 f_0}{2h^2}(x - x_0) + \frac{\Delta^2 f_0}{2h^2}(x - x_1) + O(h^2)
$$

$$
f'(x_0) = \frac{f_1 - f_0}{h} - \frac{1}{2h}(f_2 - 2f_1 + f_0) + O(h^2)
$$

\n
$$
= \frac{2f_1 - 2f_0 - f_2 + 2f_1 - f_0}{2h} + O(h^2)
$$

\n
$$
= \left[\frac{1}{h} \left(-\frac{3}{2}f_0 + 2f_1 - \frac{1}{2}f_2 \right) + O(h^2) \right]
$$
 Forward Difference
\n
$$
f'(x_1) = \frac{f_1 - f_0}{h} + \frac{1}{2h}(f_2 - 2f_1 + f_0) + O(h^2)
$$

 $\ddot{\bullet}$

$$
= \frac{1}{2h}(f_2 - f_0) + O(h^2)
$$
Central Difference
Second Derivatives

$$
n=2 \t f''(x_0) = \frac{\Delta^2 f_0}{h^2} + O(h) = \frac{1}{h^2}(f_0 - 2f_1 + f_2) + O(h)
$$
Forward Difference

$$
n=3 \t f''(x_1) = \frac{1}{h^2}(f_0 - 2f_1 + f_2) + O(h^2)
$$
Central Difference

x

Finite Differences using Polynomial approximations Numerical Interpolation: Lagrange Polynomials (Reformulation of Newton's polynomial)

$$
p(x) = \sum_{k=0}^{n} L_k(x) f(x_k) = \sum_{k=0}^{n} L_k(x) f_k
$$

\n
$$
L_k(x) = \sum_{i=0}^{n} \ell_{ik} x^{i}
$$

\n
$$
L_k(x_i) = \delta_{ki} = \begin{cases} 0 & k \neq i \\ 1 & k = i \end{cases}
$$

\nDifferentiate the program
\n
$$
L_k(x) = \prod_{j=0, j \neq k}^{n} \frac{x - x_j}{x_k - x_j}
$$

\n
$$
L_k(x) = \prod_{j=0, j \neq k}^{n} \frac{x - x_j}{x_k - x_j}
$$

\n
$$
L_k(x) = \prod_{j=0, j \neq k}^{n} \frac{x - x_j}{x_k - x_j}
$$

Important for numerical integration

Hermite Interpolation Polynomials and Compact / Pade' Difference Schemes

- Use the values of the function and its derivative(s) at given points *k*
	- For example, for values of the function and of its first derivatives at pts *k*

$$
u(x) = \sum_{k=1}^{n} a_k(x) u_k + \sum_{k=1}^{m} b_k(x) \left(\frac{\partial u}{\partial x}\right)_k
$$

• General form for implicit/explicit schemes (here focusing on space)

$$
\sum_{i=-r}^{s} b_i \left(\frac{\partial^m u}{\partial x^m} \right)_{j+i} - \sum_{i=-p}^{q} a_i u_{j+i} = \tau_{\Delta x}
$$

– Generalizes the Lagrangian approach by using Hermitian interpolation

- Leads to the "Compact difference schemes" or " Pade' schemes "
- Are implemented by the use of efficient banded solvers

FINITE DIFFERENCES: Higher Order Accuracy Taylor Tables for Pade' schemes

Taylor table for a central three point Hermitian approximation to a first derivative.

$$
d\Bigg(\frac{\partial u}{\partial x}\Bigg)_{j-1}+\Bigg(\frac{\partial u}{\partial x}\Bigg)_j+e\Bigg(\frac{\partial u}{\partial x}\Bigg)_{j+1}-\frac{1}{\Delta x}\Big(au_{j-1}+bu_j+cu_{j+1}\Big)=?
$$

Image by MIT OpenCourseWare.

FINITE DIFFERENCES: Higher Order Accuracy Taylor Tables for Pade' schemes, Cont'd

Taylor table for a central three point Hermitian approximation to a first derivative.

$$
d\left(\frac{\partial u}{\partial x}\right)_{j-1}+\left(\frac{\partial u}{\partial x}\right)_j+e\left(\frac{\partial u}{\partial x}\right)_{j+1}-\frac{1}{\Delta x}\left(au_{j-1}+bu_j+cu_{j+1}\right)=?
$$

Image by MIT OpenCourseWare.

Sum each column starting from left and force the sums to be zero by proper choice of a, b, c, etc:

$$
\begin{bmatrix} -1 & -1 & -1 & 0 & 0 \\ 1 & 0 & -1 & 1 & 1 \\ -1 & 0 & -1 & 2 & 2 \\ 1 & 0 & -1 & 3 & 2 \\ -1 & 0 & -1 & -4 & 4 \end{bmatrix} \begin{bmatrix} a \\ b \\ c \\ d \\ e \end{bmatrix} = \begin{bmatrix} 0 \\ -1 \\ 0 \\ 0 \\ 0 \end{bmatrix} \Rightarrow [a \quad b \quad c \quad d \quad e] = \frac{1}{4} [-3 \quad 0 \quad 3 \quad 1 \quad 1]
$$

Truncation error is first column in the table that does not vanish, here sixth column:

$$
\tau_{\Delta x} = \frac{\Delta x^4}{120} \left(\frac{\partial^5 u}{\partial x^5} \right)_j
$$

Compact / Pade' Difference Schemes: Examples

We can derive family of compact centered approximations for ϕ up to 6th order using:

$$
\alpha \left(\frac{\partial \phi}{\partial x}\right)_{i+1} + \left(\frac{\partial \phi}{\partial x}\right)_{i} + \alpha \left(\frac{\partial \phi}{\partial x}\right)_{i-1} = \beta \frac{\phi_{i+1} - \phi_{i-1}}{2\Delta x} + \gamma \frac{\phi_{i+2} - \phi_{i-2}}{4\Delta x}
$$

Comments:

- Pade' schemes use fewer computational nodes and thus are more compact than CDS
- Can be advantageous (more banded systems!)

Image by MIT OpenCourseWare.

Higher-Order Finite Difference Schemes **Considerations**

- Retaining more terms in Taylor Series or in polynomial approximations allows to obtain FD schemes of increased order of accuracy
- However, higher-order approximations involve more nodes, hence more complex system of equations to solve and more complex treatment of boundary condition schemes
- Results shown for one variable valid for mixed derivatives
- To approximate other terms that are not differentiated: reaction terms, etc
	- Values at the center node is normally all that is needed
	- However, for strongly nonlinear terms, care is needed (see later)
- Boundary conditions must be discretized

Finite Difference Schemes: Implementation of Boundary conditions

- For unique solutions, information is needed at boundaries
- Generally, one is given either:

i) the variable: $u(x = x_{\text{bnd}}, t) =$ (Dirichlet BCs) ii) a gradient in a specific direction, e.g.: $\left| \frac{\partial u}{\partial x} \right|$ = $(Neumann BCs)$ $\partial x\left|_{(x_{\textrm{bnd}}, t)}\right.$ iii) a linear combination of the two quantities (Robin BCs)

- Straightforward cases:
	- If value is known, nothing special needed (one doesn't solve for the BC)
	- If derivatives are specified, for first-order schemes, this is also straightforward to treat

Finite Difference Schemes: Implementation of Boundary conditions, Cont'd

- Harder cases: when higher-order approximations are used
	- At and near the boundary: nodes outside of domain would be needed
- Remedy: use different approximations at and near the boundary
	- Either, approximations of lower order are used
	- Or, approximations go deeper in the interior and are one-sided. For example,
		- $\frac{\partial u}{\partial x}\Big|_{(x_{\text{bnd}},t)} = 0 \Rightarrow \frac{u_2 u_1}{x_2 x_1} \approx 0 \Rightarrow u_2 = u_1$ • 1st order forward-difference: $\frac{u}{u}$ = 0 \rightarrow $\frac{u_2 - u}{u_1}$
		- Parabolic fit to the bnd point and two inner points:

$$
\frac{\partial u}{\partial x}\Big|_{(x_{\text{bnd}},t)} \approx \frac{-u_3(x_2 - x_1)^2 + u_2(x_3 - x_1)^2 - u_1 \left[(x_3 - x_1)^2 - (x_2 - x_1)^2 \right]}{(x_2 - x_1)(x_3 - x_1)(x_3 - x_2)} \qquad \left(\approx \frac{-u_3 + 4u_2 - 3u_1}{2\Delta x} \text{ for equidistant nodes}\right)
$$

- Cubic fit to 4 nodes (3rd order difference): $\frac{\partial u}{\partial x}$ $\approx \frac{2u_4 9u_3 + 18u_2 11u_1}{5!} + O(\Delta x^3)$ for equidistant nodes $\partial x \big|_{(x_{\text{bnd}}, t)}$ 6 Δx
- $18u_2 9u_3 + 2u_4$ 6 $\Delta x \left(\partial u \right)$ • Compact schemes, cubic fit to 4 pts: $u_{(x_{bnd},t)} = u_1 \approx \frac{16u_2 + 6u_3 + 2u_4}{11} - \frac{6u_4}{11} \left(\frac{6u_5}{\partial x}\right)_1$ for equidistant nodes
- In Open-boundary systems, boundary problem is not well posed =>
	- Separate treatment for inflow/outflow points, multi-scale approach and/or generalized inverse problem (using data in the interior)

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