

2.29 Numerical Fluid Mechanics Fall 2011 – Lecture 22

REVIEW Lecture 21:

• End of Time-Marching Methods: higher-order methods

- Runge-Kutta Methods
	- Additional points are between t_n and t_{n+1}
- Multistep/Multipoint Methods: Adams Methods
	- Additional points are at past time steps
- Practical CFD Methods
- Implicit Nonlinear systems
- Deferred-correction Approach
- Complex Geometries
	- Different types of grids
	- Choice of variable arrangements
		- Velocity vectors: Cartesian or Grid-oriented
		- Staggered or Collocated variables
- Grid Generation
	- Basic concepts and structured grids

$$
\phi^{n+1} - \phi^n = \int\limits_{t_n}^{t_{n+1}} f(t,\phi) dt
$$

TODAY (Lecture 22): Grid Generation and Intro to FV Complex Geometries

- Complex Geometries
- Grid Generation
	- Basic concepts and structured grids
		- Stretched grids
		- Algebraic methods (strecthed grids)
		- General coordinate transformation
		- Differential equation methods
		- Conformal mapping methods
	- Unstructured grid generation
		- Delaunay Triangulation
		- Advancing Front method
- Finite Volume on Complex geometries
	- Computation of convective fluxes
	- Computation of diffusive fluxes
	- Comments on 3D
- Solution of the Navier-Stokes Equations

References and Reading Assignments

- Chapter 8 on "Complex Geometries" of "J. H. Ferziger and M. Peric, Computational Methods for Fluid Dynamics. Springer, NY, 3rd edition, 2002"
- Chapter 9 on "Grid Generation" of T. Cebeci, J. P. Shao, F. Kafyeke and E. Laurendeau, *Computational Fluid Dynamics for* Engineers. Springer, 2005.
- Ref on Grid Generation only:
	- Thompson, J.F., Warsi Z.U.A. and C.W. Mastin, "Numerical Grid Generation, Foundations and Applications", North Holland, 1985

Classes of Grid Generation

- An arrangement of discrete set of grid points or cells needs to be generated for the numerical solution of PDEs (fluid conservation equations)
	- Finite volume methods:
		- Can be applied to uniform and non-uniform grids
	- Finite difference methods:
		- Require a coordinate transformation to map the irregular grid in the spatial domain to a regular one in the computational domain
		- Difficult to do this in complex 3D spatial geometries
		- So far, only used with structured grid (could be used with unstructured grids with polynomials ϕ defining the shape of ϕ around a grid point)
- Three major classes of grid generation: i) algebraic methods, ii) differential equation methods and iii) conformal mapping methods
- Grid generation and solving PDE can be independent
	- A numerical (flow) solver can in principle be developed independently of the grid
	- A grid generator then gives the metrics (weights) and the one-to-one correspondence between the spatial-grid and computational-grid

Grid Generation: Basic Concepts for Structured Grids

- Structured Grids (includes curvilinear or non-orthogonal grids)
	- Often utilized with FD schemes
	- Methods based on coordinate transformations
- Consider irregular shaped physical domain (x, y) in Cartesian coordinates and determine its mapping to the computational domain in the (ξ, η) Cartesian coordinates
	- Increase *ξ or η* monotonically in physical domain along "curved lines"
	- Coordinate lines of the same family do not cross
	- Lines of different family don't cross more than once
	- errors are expected as a rectangle in the computational plane. Physical grid refined where large

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– Physical grid refined where large a simply-connected irregular shape in the physical plane is mapped

- Mapped (computational) region has a rectangular shape:
	- Coordinates (*ξ*, *η)* can vary from 1 to (I, J), with mesh sizes taken equal to 1
- Boundaries are mapped to boundaries

Grid Generation: Basic Concepts for Structured Grids, Cont'd

- The example just shown was the mapping of an irregular, simply connected, region into a rectangle.
- Other configurations are of course possible
- For example, a L-shape domain can be mapped into:
	- a regular L-shape

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Grid Generation for Structured Grids: Stretched Grids

- Consider a viscous flow solution on a given body, where the velocity varies rapidly near the surface of the body (Boundary Layer)
- For efficient computation, a finer grid near the body and coarser grid away from the body is effective (aims to maintain constant accuracy)
- Possible coordinate transformation: a scaling "*η* = log (*y*)" ↔ "*y* = exp(*η*)"

\n The following equation:\n
$$
\text{as a scaling } \eta = \log(y) \to \gamma
$$
\n with the following equation:\n $\text{as a scaling function: } \eta = \log(y) \to \gamma$ \n with the following equation:\n $\text{as a scalar function: } \rho = \log(y) \to \gamma$ \n with the following equation:\n $\text{as a scalar function: } \rho = \log(y) \to \gamma$ \n with the following equation:\n $\text{as a scalar function: } \rho = \log(y) \to \gamma$ \n with the following equation:\n $\text{as a scalar function: } \rho = \log(y) \to \gamma$ \n with the following equation:\n $\text{as a scalar function: } \rho = \log(y) \to \gamma$ \n with the following equation:\n $\text{as a scalar function: } \rho = \frac{\beta + 1}{\beta - 1}$ \n with the following equation:\n $\text{as a scalar function: } \rho = \frac{\beta + 1}{\beta - 1}$ \n with the following equation:\n $\text{as a scalar function: } \rho = \frac{\beta + 1}{\beta - 1}$ \n with the following equation:\n $\text{as a scalar function: } \rho = \frac{\beta + 1}{\beta - 1}$ \n with the following equation:\n $\text{as a scalar function: } \rho = \frac{\beta + 1}{\beta - 1}$ \n with the following equation:\n $\text{as a scalar function: } \rho = \frac{\beta + 1}{\beta - 1}$ \n with the following equation:\n $\text{as a scalar function: } \rho = \frac{\beta + 1}{\beta - 1}$ \n with the following equation:\n $\text{as a scalar function: } \rho = \frac{\beta + 1}{\beta - 1}$ \n with the following equation:\n $\text{as a scalar function: } \rho = \frac{\beta + 1}{\beta - 1}$ \n with the following equation:\n $\text{as a vector function: } \rho = \frac{\beta + 1}{\beta - 1}$ \n with the following equation:\n $\text{as a vector function: } \rho = \frac{\beta + 1}{\beta - 1}$ \n with the following equation:\n $\text{as a vector function: } \rho = \frac{\beta + 1}{\beta - 1}$ \n with the following equation:\n $\text{as a vector function: } \rho = \frac{\beta + 1}{\beta - 1}$ \n with the following equation:\n $\text{as a vector function: } \$

The parameter β (1 < β < ∞) is the stretching parameter. As *β* gets close to 1, more grid points are clustered to the wall in the physical domain.

• Inverse transformation is needed to map solutions back from *ξ*, *η* domain:

$$
x = \xi
$$

$$
\frac{y}{h} = \frac{(\beta + 1) - (\beta - 1)B^{1-\eta}}{1 + B^{1-\eta}}
$$

Fig. 9.4. One-dimensional stretching transformation. (a) Physical plane, (b) computational plane.

Grid Generation for Structured Grids: Stretched Grids, Cont'd

- How do the conservation equations change?
- Consider the continuity equation for steady state flow in physical (*x, y*) space:

$$
\nabla \cdot (\rho \vec{v}) = 0 \quad \Rightarrow \quad \frac{\partial \rho u}{\partial x} + \frac{\partial \rho v}{\partial y} = 0
$$

• In the computational plane, this equation becomes (chain rule)

$$
\frac{\partial \rho u}{\partial x} = \frac{\partial \rho u}{\partial \xi} \frac{\partial \xi}{\partial x} + \frac{\partial \rho u}{\partial \eta} \frac{\partial \eta}{\partial x} \Bigg|_{\text{in } \mathcal{F}} = \frac{\partial \rho u}{\partial \xi} \xi_x + \frac{\partial \rho u}{\partial \eta} \eta_x + \frac{\partial \rho v}{\partial \xi} \xi_y + \frac{\partial \rho v}{\partial \eta} \eta_y = 0
$$

$$
\frac{\partial \rho u}{\partial y} = \frac{\partial \rho v}{\partial \xi} \frac{\partial \xi}{\partial y} + \frac{\partial \rho v}{\partial \eta} \frac{\partial \eta}{\partial y} \Bigg|_{\text{in } \mathcal{F}} = \frac{\partial \rho u}{\partial \xi} \xi_x + \frac{\partial \rho u}{\partial \eta} \eta_x + \frac{\partial \rho v}{\partial \xi} \xi_y + \frac{\partial \rho v}{\partial \eta} \eta_y = 0
$$

• For our stretching transformation, one obtains:

$$
\xi_x = 1
$$
, $\eta_x = 0$, $\xi_y = 0$, $\eta_y = \frac{2\beta}{h \ln(B)} \frac{1}{\beta^2 - (1 - y/h)^2}$

• Therefore, the continuity equation becomes:

$$
\frac{\partial \rho u}{\partial \xi} + \frac{\partial \rho v}{\partial \eta} \eta_{y} = 0
$$

2.29 Numerical Fluid Mechanics PFJL Lecture 22, 8 – This equation can be solved on a uniform grid (slightly more complicated eqn. system), and the solution mapped back to the physical domain using the inverse transform

 $\xi = x$

 $\eta = \frac{y}{2}$ *x* 2

Grid Generation for Structured Grids: Algebraic Methods

• Algebraic Method = Generalization of stretching method (2 & 3D)

D

E F G

- Consider fitting a diverging nozzle:
	- Let's assume a nozzle defined by: $y = x^2$ $1 \le x \le 2$
- Now, choose a curvilinear system:
	- Define the transformation:

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2.29 Numerical Fluid Mechanics PFJL Lecture 22, 9 2 2 1 0, where 1, 0, , *^x ^x ^y ^y ^x ^y ^x ^y u u v v*

Grid Generation for Structured Grids: Algebraic Methods: Transfinite Interpolation

- Multi-directional interpolation (Transfinite Interpolation)
	- To generate algebraic grids within more complex domains or around more complex configurations, multi-directional interpolations can be used
		- They consist of a suite of unidirectional interpolations
- Unidirectional Interpolations (1D curve)
	- The Cartesian coordinate vector of each point on a curve **r**(x,y) is obtained as an interpolation between points that lie on the boundary curves
	- How to interpolate? the regulars:
		- Lagrange Polynomials: match function values

$$
\vec{r}(i) = \sum_{k=0}^{n} L_k(i) \vec{r}_k
$$
 with $L_k(i) = \prod_{j=0, j \neq k}^{n} \frac{i - i_j}{i_k - i_j}$,

• Hermite Polynomials: match both function and 1st derivative values

$$
\vec{r}(i) = \sum_{k=1}^{n} a_k(i) \, \vec{r}_k + \sum_{k=1}^{m} b_k(i) \, \vec{r} \, \, \vec{r}
$$

Grid Generation for Structured Grids: Algebraic Methods: Transfinite Interpolation, Cont'd

- Unidirectional Interpolations (1D curve), Cont'd
	- Lagrange and Hermite Polynomials fit a single polynomial from one boundary to the next => for long boundaries, oscillations may occur
	- Alternative, use set of lower order polynomials to form a piece-wise continuous interpolation:
		- Spline interpolation (match as many derivatives as possible at interior point junctions), Tension-spline (more localized curvature) and B-splines (allows local modification of the interpolation)
	- Use interpolation functions that are not polynomials, usually "stretching functions": exp, tanh, sinh, etc \mathbf{r}_2
- Multi-directional or Transfinite Interpolation
	- Extends 1D results to 2D or 3D by

successive applications of 1D interpolations

– For example, *i* then *j*.

 $i_2=I$

 \mathbf{r}_1

j

Grid Generation for Structured Grids:

Algebraic Methods: Transfinite Interpolation, Cont'd

- Multi-directional or Transfinite Interpolation, Cont'd
	- In 2D, the transfinite interpolation can be implemented as follows
		- Interpolate position vectors **r** in *i*-direction => leads to points $f_1 = \Box_i(r)$ and *i*-lines
		- Evaluate the difference between **r** and this result on the j-lines that will be used in the *j*-interpolation (e.g. difference with curved $i=0$ and $i=1$): $\mathbf{r} - \mathbf{f}_1$
		- Interpolation of the discrepancy in the j-direction: $f_2 = \Box_j (r f_1)$
		- Addition of the results of this *j*-interpolation to the results of the *i*-interpolation: \mathbf{r} (*i*, *j*)= $\mathbf{f}_1 + \mathbf{f}_2$
- •Of course, Lagrange, Hermite Polynomials, Spline and non-polynomial (stretching) functions can be used for transfinite interpolations
- •In 2D, inputs to program are 4 boundaries
- •Issues: Propagates discontinuities in the interior and grid lines can overlap in some situations
- => needs to be refined by grid generator solving a PDE

Examples:

Fig. 9.12. (a) C-grid around ellipse: Unidirectional Lagrange Interpolation, (b) C-grid around ellipse: Unidirectional Hermite Interpolation, (c) C-grid around ellipse: Unidirectional Lagrange Interpolation with Hyperbolic Tangent Spacing, (d) C-grid around ellipse: Unidirectional Hermite Interpolation with Hyperbolic Tangent Spacing.

Grid Generation, structured Grids: General Coordinate transformation

- In general, coordinates are defined by a transformation: $x_i = x_i (\xi_i)$ (*i* and $j = 1, 2, 3$)
- All transformations are characterized by their Jacobian determinant *J.*

$$
J = \det\left(\frac{\partial x_i}{\partial \xi_j}\right) = \begin{vmatrix} \frac{\partial x_1}{\partial \xi_1} & \frac{\partial x_1}{\partial \xi_2} & \frac{\partial x_1}{\partial \xi_3} \\ \frac{\partial x_2}{\partial \xi_1} & \frac{\partial x_2}{\partial \xi_2} & \frac{\partial x_2}{\partial \xi_3} \\ \frac{\partial x_3}{\partial \xi_1} & \frac{\partial x_3}{\partial \xi_3} & \frac{\partial x_3}{\partial \xi_3} \end{vmatrix}
$$

– For Cartesian vector components, one only needs to transform derivatives. One has: For Cartesian vector comp
derivatives. One has:
 $\frac{\partial \phi}{\partial x} = \frac{\partial \phi}{\partial \xi} \frac{\partial \xi_j}{\partial x} = \frac{\partial \phi}{\partial \xi} \frac{\beta^{ij}}{I}$, where β^{ij} r - For Cartesian vector components, one only nee
derivatives. One has:
 $\frac{\partial \phi}{\partial x} = \frac{\partial \phi}{\partial \xi} \frac{\partial \xi_j}{\partial x} = \frac{\partial \phi}{\partial \xi} \frac{\partial^y}{\partial x}$, where β^y represents the cofactor of $\frac{\partial \phi}{\partial \xi}$

ector components, one only needs to transform
 $\frac{\partial x_i}{\partial \xi_j}$ (element *i*, *j* of Jacobian matrix) *ij j ij i* $\frac{\partial}{\partial \xi_i} = \frac{\partial \phi}{\partial \xi_j} \frac{\partial \xi_j}{\partial x_i} = \frac{\partial \phi}{\partial \xi_j} \frac{\beta^{ij}}{J}$, where β^{ij} represents the cofactor of $\frac{\partial x_i}{\partial \xi_j}$ *x* **sform**
i, *j* of Jacobian mate *For Cartesian vector comporterior derivatives. One has:*
 $\frac{\partial \phi}{\partial x_i} = \frac{\partial \phi}{\partial \xi_j} \frac{\partial \xi_j}{\partial x_i} = \frac{\partial \phi}{\partial \xi_j} \frac{\beta^{ij}}{J}$, where β^{ij} represents r Cartesian vector components, one only need
rivatives. One has:
 $\frac{\partial \phi}{\partial \xi_j} \frac{\partial \xi_j}{\partial x_i} = \frac{\partial \phi}{\partial \xi_j} \frac{\beta^{ij}}{J}$, where β^{ij} represents the cofactor of $\frac{\partial x_i}{\partial \xi_j}$ For Cartesian vector derivatives. One has:
= $\frac{\partial \phi}{\partial \xi_j} \frac{\partial \xi_j}{\partial x_i} = \frac{\partial \phi}{\partial \xi_j} \frac{\beta^{ij}}{J}$, when - For Cartesian vector components, one only nee
derivatives. One has:
 $\frac{\partial \phi}{\partial x_i} = \frac{\partial \phi}{\partial \xi_j} \frac{\partial \xi_j}{\partial x_i} = \frac{\partial \phi}{\partial \xi_j} \frac{\beta^{ij}}{J}$, where β^{ij} represents the cofactor of $\frac{\partial}{\partial \theta}$

– In 2D, $x = x(\xi, η)$ and $\phi = \phi(\xi, η)$, this leads to:

$$
(\xi, \eta) \text{ and } \phi = \phi(\xi, \eta), \text{ this leads to:}
$$
\n
$$
\frac{\partial \phi}{\partial x} = \frac{\partial \phi}{\partial \xi} \frac{\partial \xi}{\partial x} + \frac{\partial \phi}{\partial \eta} \frac{\partial \eta}{\partial x} = \frac{\partial \phi}{\partial \xi} \frac{\beta^{11}}{J} + \frac{\partial \phi}{\partial \eta} \frac{\beta^{12}}{J} = \frac{1}{J} \left(\frac{\partial \phi}{\partial \xi} \frac{\partial y}{\partial \eta} - \frac{\partial \phi}{\partial \eta} \frac{\partial y}{\partial \xi} \right)
$$

Recall: the minor element m_{ij} corresponding to a_{ij} is the determinant of the submatrix that remains after the i^{th} row and the j^{th} column are deleted from ${\bf A}.$ The cofactor c_{ij} of a_{ij} is: $\;\; c_{ij}=(-1)^{i+j}\;m_{ij}$

Grid Generation Structured Grids: Coordinate transformation, Cont'd

• How do the conservation equations transform? The generic conservation equation in Cartesian coordinates:

$$
\frac{\partial \rho \phi}{\partial t} + \nabla \cdot (\rho \phi \vec{v}) = \nabla \cdot (k \nabla \phi) + s_{\phi} \quad \Leftrightarrow \quad \frac{\partial \rho \phi}{\partial t} + \frac{\partial}{\partial x_j} \left(\frac{\rho \phi v_j - k \frac{\partial \phi}{\partial x_j}}{\phi} \right) = s_{\phi}
$$

• becomes:

$$
J\frac{\partial \rho \phi}{\partial t} + \frac{\partial}{\partial \xi_j} \left(\rho \phi U_j - \frac{k}{J} \left(\frac{\partial \phi}{\partial \xi_m} B^{mj} \right) \right) = J s_{\phi}
$$

where:

 $1j$ 1 32 R^2 1 33 R^3 $1^{1}P$ 1^{1} $1^{2}P$ 1^{1} 1^{2} 1^{3} is proportional to the velocity component aligned with (normal to ξ_i = const.) k_j \cdots R^{1j} \cdots R^{2j} \cdots R^{3j} $j = v_k p - v_1 p - v_2 p - v_3 p$ is proportional to the velocity component anglied with ς_j *j* $U_i = v_k \beta^{kj} = v_1 \beta^{lj} + v_2 \beta^{2j} + v_3$ $\beta^{\prime\prime} = v_1 \beta^{1} + v_2 \beta^{2} + v_3 \beta^{3}$ is proportional to the velocity component aligned with ξ ξ $= v_k \beta^{kj} = v_1 \beta^{1j} + v_2 \beta^{2j} +$ $=$

 $1j \rho 1m$, $\rho 2j \rho 2m$, $\rho 3j \rho 3$ are coefficients, sum of products of cofactors $m_j = \alpha^{kj} \alpha^{km} = \alpha^{1j} \alpha^{1m} + \alpha^{2j} \alpha^{2m} + \alpha^{3j} \alpha^{3m}$ are exercisely sum of products of expectators α^{ij} *B* $\beta^{\kappa} \beta^{\kappa m} = \beta^{\kappa} \beta^{\kappa m} + \beta^{\kappa} \beta^{\kappa m} + \beta^{\kappa} \beta^{\kappa m}$ are coefficients, sum of products of cofactors β $=\beta^{kj}\beta^{km}=\beta^{1j}\beta^{1m}+\beta^{2j}\beta^{2m}+$

- As a result, each 1st derivative term is replaced by a sum of three terms which contains derivatives of the coordinates as coefficients
- Unusual features of conservation equations in non-orthogonal grids:
- 2.29 Numerical Fluid Mechanics PFJL Lecture 22, 15 – Mixed derivatives appear in the diffusive terms and metrics coefficients appear in the continuity eqn.

Structured Grids: Coordinate transformation, Cont'd Some Comments

- Coordinate transformation often presented only as a means of converting a complicated non-orthogonal grid into a simple, uniform Cartesian grid (the computational domain, whose grid-spacing is arbitrary)
- However, simplification is only apparent:
	- Yes, the computational grid is simpler than the original physical one
	- But, the information about the complexity in the computational domain is now in the metric coefficients of the transformed equations
		- i.e. discretization of computational domain is now simple, but the calculation of the Jacobian and other geometric information is not trivial (the difficulty is hidden in the metric coefficients)
- As mentioned earlier, FD method can in principle be applied to unstructured grids: specify a local shape function, differentiate and write FD equations. Has not yet been done.

Grid Generation for Structured Grids: Differential Equation Methods

- Grid transformation relations determined by a finite-difference solution of PDEs
	- For 2D problems, two elliptic (Poisson) PDEs are solved
	- Can be done for any coordinate systems, but here we will use Cartesian coordinates. The 2D transformation is then:
		- From the physical domain (*x*, *y*) to the computational domain (*ξ*, *η*)
		- At physical boundaries, one of *ξ*, *η is* constant, the other is monotonically varying
		- At interior points:

$$
\frac{\partial^2 \xi}{\partial x^2} + \frac{\partial^2 \xi}{\partial y^2} = P(\xi, \eta)
$$

$$
\frac{\partial^2 \eta}{\partial x^2} + \frac{\partial^2 \eta}{\partial y^2} = Q(\xi, \eta)
$$

where $P(\xi, \eta)$ and $Q(\xi, \eta)$ are called the "control functions"

- Their selection allows to concentrate the *ξ*, *η* lines in specific regions
- If they are null, coordinates will tend to be equally spaced away from boundaries
- Boundary conditions: *ξ*, *η* specified on boundaries of physical domain

Grid Generation for Structured Grids: Differential Equation Methods, Cont'd

- Computations to generate the grid mapping are actually carried out in the computational domain (*ξ*, *η*) itself !
	- don't want to solve the elliptic problem in the complex physical domain!
- Using the general rule, the elliptic problem is transformed into:

e general rule, the elliptic problem is transformed into:
\n
$$
\alpha \frac{\partial^2 x}{\partial \xi^2} - 2\beta \frac{\partial^2 x}{\partial \xi \partial \eta} + \gamma \frac{\partial^2 x}{\partial \eta^2} + J^2 \left(P \frac{\partial x}{\partial \xi} + Q \frac{\partial x}{\partial \eta} \right) = 0
$$
\n
$$
\alpha \frac{\partial^2 y}{\partial \xi^2} - 2\beta \frac{\partial^2 y}{\partial \xi \partial \eta} + \gamma \frac{\partial^2 y}{\partial \eta^2} + J^2 \left(P \frac{\partial y}{\partial \xi} + Q \frac{\partial y}{\partial \eta} \right) = 0
$$
\n
$$
\alpha = x_n^2 + y_n^2; \quad \beta = x_\xi x_n + y_\xi y_n; \quad \gamma = x_\xi^2 + y_\xi^2; \quad J = x_\xi y_n - x_n y_\xi \quad \text{(with } x_\xi = \frac{\partial x}{\partial \xi}, \text{ etc.})
$$
\nwhere α is a solution of the PQQs in (1)

where $\alpha = x_{\eta}^2 + y_{\eta}^2$; $\beta = x_{\xi}x_{\eta} + y_{\xi}y_{\eta}$; $\gamma = x_{\xi}^2 + y_{\xi}^2$; $J = x_{\xi}y_{\eta} - x_{\eta}y_{\xi}$ (with $x_{\xi} = \frac{\partial x}{\partial \xi}$, etc) γ گخ ∂

- Boundary conditions are now the transformed values of the BCs in (x, y) domain: they are the values of the positions (*x, y*) of the grid points on the physical domain mapped to their locations in the computational domain
- Equations can be solved by FD method to determine values of every grid point (*x, y*) in the interior of the physical domain
- Method developed by Thomson et al, 1985 (see ref)

Grid Generation for Structured Grids: Differential Equation Methods, Example

 (c)

Fig. 9.13. (a) Starting algebraic C-grid around an airfoil section; 70×30 grid points; inner spacing $\Delta S_1 = 0.015c$, outer spacing $\Delta S_2 = 0.3c$, (b) Elliptic C-grid obtained after smoothing the algebraic grid of (a) by the solution of Poisson equations (50 iterations), (c) Close-up of the C-grid showing the application of orthogonality conditions near the leading edge region.

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Grid Generation for Structured Grids: Conformal Mapping Methods

- Conformal mapping schemes are analytical or partially analytical (as opposed to differential equation methods)
- Restricted to two dimensional flows (based on complex variables): useful for airfoils
- Examples:

Fig. 9.14. Three common grids for airfoils. (a) C-grid, (b) O-grid, and (c) H-grid.

- C-mesh: high density near leading edge of airfoil and good wake
- O-mesh: high density near leading and trailing edge of airfoil
- H-mesh: two sets of mesh lines similar to a Cartesian mesh, which is easiest to generate. Its mesh lines are often well aligned with streamlines

Grid Generation for Structured Grids: Conformal Mapping Methods: Example

- C-mesh example is generated by a parabolic mapping function
- It is essentially a set of confocal, orthogonal parabolas wrapping around the airfoil
- The mapping is defined by:

$$
2(x+iy) = (\xi + i\eta)^2
$$

or

$$
2 x = \xi^2 - \eta^2; \quad y = \xi \eta
$$

• Inverse transformation:

$$
\xi^2 = \sqrt{x^2 + y^2} + x; \quad \eta^2 = \sqrt{x^2 + y^2} - x
$$

- Polar coordinates can be used for easier physical plane to computational plane transformation. The mapping is defined by: $2(x+iy) = (\xi + i\eta)^2$

or
 $2x = \xi^2 - \eta^2$; $y = \xi \eta$

Inverse transformation:
 $\xi^2 = \sqrt{x^2 + y^2} + x$; $\eta^2 = \sqrt{x^2 + y^2} - x$
 ξ sharples *Ratified Secretion, see http://*

Polar coordinates can be used
- In conformal mapping, singular point is point where mapping fails (here, it is

Grid Generation: Unstructured Grids

- Generating unstructured grid is complicated but now relatively automated in "classic" cases
- Involves succession of smoothing techniques that attempt to align elements with boundaries of physical domain
- Decompose domain into blocks to decouple the problems
- Need to define point positions and connections
- Most popular algorithms:
	- Delaunay Triangulation Method
	- Advancing Front Method
- Two schools of thought: structured vs. unstructured, what is best for CFD?

Fig. 9.16. 2D Unstructured grid for Navier-Stokes computations of a multi-element airfoil generated with the hybrid advancing front Delaunay method of Mavriplis [6].

- Structured grids: simpler grid and straightforward treatment of algebraic system, but mesh generation constraints on complex geometries
- Unstructured grids: generated faster on complex domains, easier mesh refinements, but data storage and solution of algebraic system more complex

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