

Potential Energy: Solid Sphere + Particle

Particle m

Homogeneous solid sphere M, R

Case I: $r \geq R$

Force F on shell depends only on dist. to center

$$\vec{F} = -\frac{GmM}{r^2} \hat{r} \quad (r \geq R)$$

$$U = -GmM/r \quad (r \geq R)$$

Case II: $r < R$

From shell result, force due only to mass between $r=0$ and $r=r$

$$\vec{F} = \frac{GmM_a(r)}{r^2} \hat{r}$$

$M_a(r) \equiv$ mass in $r=0 \rightarrow r$

$$M_a(r) = M \left(\frac{r}{R} \right)^3$$

$$\therefore \vec{F} = -\frac{GmM}{r^2} \hat{r} \quad (r \geq R)$$

$$U(r) = -\int_{\infty}^r \vec{F} \cdot d\vec{r}$$

$$= \int_{\infty}^R \frac{GmM}{r^2} dr + \int_R^r \frac{GmM r' dr'}{R^3}$$

$$= -\frac{GmM}{2R^2} [3R^2 - r^2] \quad (r \leq R)$$



Forces between Two Spherical Objects

Force on A due to (m_i) part of B acts on CM of A.

Sum over all m_i

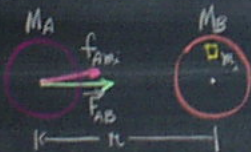
$$\vec{F}_{AB} = \frac{Gm_A m_B}{r^2} \hat{r}$$

Force on A at CM A
Force on B at CM B

If B is not a sphere:

Force on B is NOT at CM B

Force on A is at CM of A



Gravitational Force on an Extended Object.

Earth acting on particle m_i :

$$\vec{F}_i = m_i \vec{g} \quad \vec{g} = \text{accel. due to gravity}$$

$$\vec{F} = \sum \vec{F}_i = \sum m_i \vec{g} = (\sum m_i) \vec{g} = M \vec{g}$$

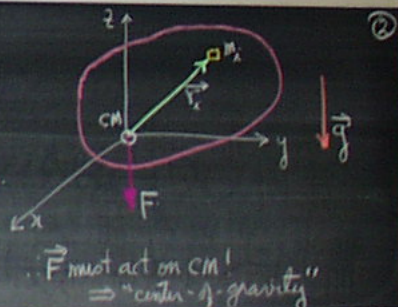
\uparrow total mass.

Calculate torques about CM:

$$\sum \vec{\tau}_{cm} = \sum \vec{r}_i \times \vec{F}_i = \sum \vec{r}_i \times m_i \vec{g} = (\sum m_i \vec{r}_i) \times \vec{g}$$

$$\therefore \sum \vec{\tau}_{cm} \equiv 0!$$

\uparrow $\vec{R}_{cm} \equiv 0$



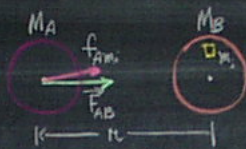
Forces between Two Spherical Objects

Force on A due to (m_i) part of B acts on CM of A.
Sum over all m_i

$$\vec{F}_{AB} = \frac{GM_A M_B}{r^2} \quad \begin{array}{l} \text{Force on A at CM A} \\ \text{Force on B at CM B} \end{array}$$

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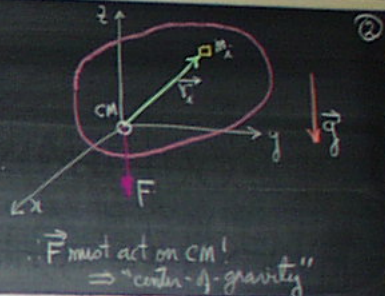
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$$\therefore \sum \vec{\tau}_{cm} = 0!$$

\uparrow total mass.
 $\uparrow \vec{R}_{cm} = 0$

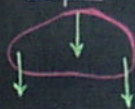


If \vec{g} varies over an extended object, CG cannot be defined except for spherical objects.

Tidal Forces.

- Non-uniform gravity
- Differential Forces.
- Object stretched or compressed.

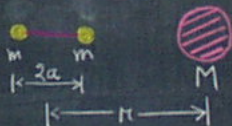
Uniform



Non-Uniform



Earth Tides:



Tidal Force \equiv difference between gravitation force at both ends

$$F_T = \frac{GMm}{(r-a)^2} - \frac{GMm}{(r+a)^2} = GMm \frac{(r+a)^2 - (r-a)^2}{(r-a)^2 (r+a)^2}$$



$$F_T = \frac{4GMm a}{r^3} \quad (\text{tidal force})$$

For $a \ll r$

Falls off as r^{-3}

• Moon, though smaller, more important than sun.

Weight and Gravitational Force:

$$W = mg \text{ near earth.}$$

$$mg = \frac{GM_E m}{R_E^2} \quad \vec{F} = m\vec{a}$$

$$g = \frac{GM_E}{R_E^2} \text{ definition of } g!$$

Know $G, R_E \Rightarrow$ Weigh Earth

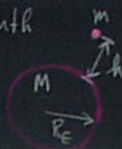
Suppose object not near earth

$$r = h + R_E \quad h = \text{height}$$

$$mg' = \frac{GM_E m}{r^2} = \frac{GM_E m}{(h + R_E)^2}$$

$$\therefore g' = \frac{GM_E}{r^2} = \frac{GM_E}{(h + R_E)^2}$$

g' decreases with altitude!



$$g' \rightarrow 0 \quad \text{as } h \rightarrow \infty$$

PE of m at radius r :

$$U(r) = -\frac{GM_E m}{r}$$

$$U(R_E) = -\frac{GM_E m}{R_E}$$

$$\Delta U = U(R_E) - U(r)$$

$$= -GM_E m \left[\frac{1}{r} - \frac{1}{R_E} \right]$$

$$\Delta U = -GM_E m \left[\frac{r - R_E}{r R_E} \right]$$

$$\text{let } r - R_E = z \quad \left. \begin{array}{l} \\ r R_E \sim R_E^2 \end{array} \right\} \text{ for } z \ll R_E \quad (4)$$

$$\Delta U = \frac{GM_E m}{R_E^2} z = (mg) z$$

Change in PE for particle m due to change in elevation near earth's surface.

Planetary Motion

Case I: Circular Orbits

$$F = \frac{GMm}{r_0^2} \text{ Force on planet.}$$

$$m \ll M$$

uniform speed: circles.

$$a_c = \frac{v^2}{r_0}$$

$$\frac{GMm}{r_0^2} = \frac{m v^2}{r_0}$$

$$v^2 = \frac{GM}{r_0} \quad \text{Radius of orbit fixes speed.}$$

$$T = \frac{2\pi r_0}{v} = \frac{\text{Circumference}}{\text{Speed}} = \text{Period - one revolution}$$



$$T^2 = \frac{4\pi^2 r_0^3}{v^2} = \frac{4\pi^2}{GM} r_0^3$$

(Kepler's Law)

$$K = \frac{1}{2} m v^2 = \frac{1}{2} \frac{GMm}{r_0} \text{ (KE)}$$

$$U = -\frac{GMm}{r_0} \text{ (PE)}$$

$$E = K + U \text{ Total Energy.} \quad (5)$$

$$= \frac{1}{2} \frac{GMm}{r_0} - \frac{GMm}{r_0}$$

$$= -\frac{1}{2} \frac{GMm}{r_0} \quad [E < 0]$$

$$L^2 = (m v r_0)^2 = GMm^2 r_0$$

$$r_0 = \frac{L^2}{GMm^2}$$

$$r_0 = \frac{4\pi^2 a^3}{GM} \text{ where } \left(a = \frac{h}{\lambda} \right)$$

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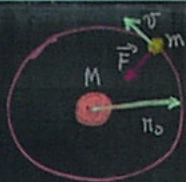
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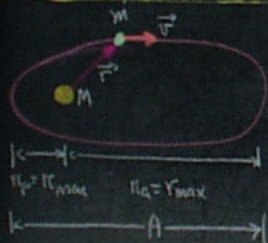
$$= \frac{1}{2} \frac{GMm}{r_0} - \frac{GMm}{r_0}$$

$$= -\frac{1}{2} \frac{GMm}{r_0} [E < 0]$$

$$L^2 = (m v r_0)^2 = GM m^2 r_0$$

$$r_0 = \frac{L^2}{GM m^2}$$

Case II: Elliptical Orbits



$m \ll M$

Easy to understand at extremes of orbit. Then $\vec{v} \perp \vec{r}$!

$$E = \frac{1}{2} m v^2 - \frac{GMm}{r} \text{ (} E = \text{constant)}$$

$$L = m v r \text{ (Total } L = \text{constant)}$$

$$K = \frac{L^2}{2I} = \frac{L^2}{2m r^2}$$

$$\therefore \frac{L^2}{2m r^2} - \frac{GMm}{r} = E$$

$$r^2 + \frac{GMm}{E} r - \frac{L^2}{2mE} = 0$$

$$r = \frac{-\frac{GMm}{E} \pm \sqrt{\left(\frac{GMm}{E}\right)^2 + \frac{4L^2}{2mE}}}{2}$$

$$r_a = r_{\max} = \left[\left(\right) + \sqrt{\quad} \right] / 2 \textcircled{6}$$

$$r_p = r_{\min} = \left[\left(\right) - \sqrt{\quad} \right] / 2$$

$$A = r_a + r_p = -\frac{GMm}{E} \text{ length of maj. axis indep. } \theta$$

same $A \Rightarrow$ same total E .

$$\text{Can show that: } E = \sqrt{1 + \frac{2EL^2}{m(GMm)^2}} \text{ eccentricity}$$

$$\text{Also: } T^2 = \frac{4\pi^2 a^3}{GM} \text{ where } (a = \frac{A}{2})$$

General Planetary Motion

Inertial Frame:

$$\vec{F}_{12} = m_1 \frac{d^2 \vec{r}_1}{dt^2} \quad (1)$$

$$\vec{F}_{21} = m_2 \frac{d^2 \vec{r}_2}{dt^2} \quad (2)$$

\vec{R} = Pos. Vector of CM.

\vec{r} = relative separation.

$$\vec{R} = \frac{m_1 \vec{r}_1 + m_2 \vec{r}_2}{m_1 + m_2}$$

$$M \vec{R} = m_1 \vec{r}_1 + m_2 \vec{r}_2 \quad (3)$$

$$M = m_1 + m_2$$

$$\vec{r} = \vec{r}_2 - \vec{r}_1 \quad (4)$$

Solve: $\vec{r}_1 = \vec{R} - \frac{m_2}{M} \vec{r} \quad (5)$

$$\vec{r}_2 = \vec{R} + \frac{m_1}{M} \vec{r} \quad (6)$$

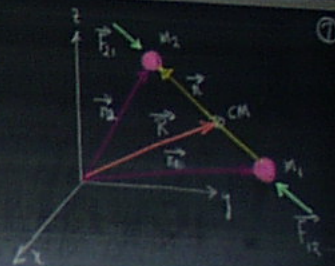
$$\vec{F}_{12} = -\vec{F}_{21} \quad (\text{2nd Law})$$

$$(1) + (2) \quad M \frac{d^2 \vec{R}}{dt^2} = m_1 \frac{d^2 \vec{r}_1}{dt^2} + m_2 \frac{d^2 \vec{r}_2}{dt^2} \equiv 0 \quad \frac{d\vec{R}}{dt} = \vec{V}_{CM} = \text{const.}$$

$$m_2 \times \text{Eq (1)} - m_1 \times \text{Eq (2)} \text{ and use (5) and (6)}$$

$$m_1 m_2 \frac{d^2 \vec{r}}{dt^2} = M \vec{F}_{21}$$

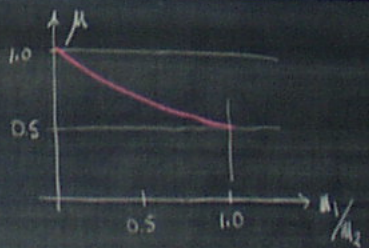
$$\text{Let } \mu = \frac{m_1 m_2}{M} \Rightarrow \mu \frac{d^2 \vec{r}}{dt^2} = \vec{F}_{21}$$



Summary:

Relative coordinate $\vec{r} = \vec{r}_2 - \vec{r}_1$ has the same behavior as the coordinate of a single particle of mass μ moving in a force field $\vec{F}_{21}(\vec{r})$. $\vec{F}_{21}(\vec{r})$ is the actual mutual force between m_1 and m_2 .

$\mu = \frac{m_1 m_2}{m_1 + m_2}$
 If $m_1 \ll m_2$
 $\mu = m_1$
 If $m_1 = m_2$
 $\mu = \frac{1}{2} m_1$

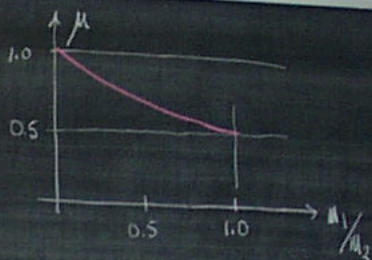


$$(V_{esc})_{Moss} = 61.8 \text{ km/s}$$

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 of a single particle of mass $= \mu$
 moving in a force field $\vec{F}_{21}(\vec{r})$.
 $\vec{F}_{21}(\vec{r})$ is the actual mutual force
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$\mu = \frac{m_1 m_2}{m_1 + m_2}$
 If $m_1 \ll m_2$
 $\mu = m_1$
 If $m_1 = m_2$
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Example:

Initial speed $= v_i$
 Max height, $h = ?$

Energy conserved.

$K_i + U_i = K_f + U_f$

$$\frac{1}{2} m v_i^2 - \frac{GM_e m}{r_i} = \frac{1}{2} m v_f^2 - \frac{GM_e m}{r_f}$$

At $r_f = r_{max}$, $v_f = 0$

$r_i = R_E$
 $\therefore v_i^2 = 2GM_E \left(\frac{1}{R_E} - \frac{1}{r_{max}} \right)$
 $h = r_{max} - R_E$

Escape Velocity

let $r_{max} = \infty$

$$v_{esc} = \sqrt{\frac{2GM_e}{R_E}}$$

$$(v_{esc})_E = \sqrt{\frac{2 \times 6.67 \times 10^{-11} \times 5.98 \times 10^{24}}{6.37 \times 10^6}}$$

$$= 1 \times 10^4 \text{ m/s} \quad 25,000 \text{ mi/h}$$

$(v_{esc})_{Moon} = 618 \text{ km/s}$

