LIGHT-SCATTERING; A TECHNIQUE FOR STUDYING SOOT IN FLAMES

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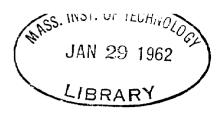
Dear Professor Franklin:

I herewith submit, in partial fulfillment of the requirements for the Degree of Doctor of Science in Chemical Engineering, a thesis entitled, "Light-Scattering; A Technique for Studying Soot in Flames".

Respectfully submitted,

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LIGHT-SCATTERING; A TECHNIQUE FOR STUDYING SOOT IN FLAMES



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NOTATION

$\mathbf{a}_{\mathbf{n}}$	light-scattering coefficient, defined by eq. A.3
$b_{\mathbf{n}}$	light-scattering coefficient, defined by eq. A.4
đ	particle diameter
e(⊖)	measured voltage at angle 0
¹ 1	light-scattering function for perpendicular polarization, defined by eq. A.1
¹ 2	light-scattering function for parallel polarization, defined by eq. A.2
k	wave number, defined as $2\pi/\lambda$
ı	optical path through the cloud of scattering particles
m.	index of refraction of soot particles, m = n - n'i
n	real part of index of refraction
n'	imaginary part of index of refraction
r	distance from the scattering particles to the point of measurement of the scattered-light intensity
s	sensitivity of the multiplier phototube system
x	size parameter, defined as $\pi d/\lambda$
У	product of x and m
α	attenuation factor for the three-mirror system
β	number fraction of agglomerates in soot particle system which contains agglomerates and ultimate particles
7	absorption coefficient
•	angle at which the scattered-light intensity is measured referred to the forward direction of the incident light
λ	wavelength of light used

7

λ

- μ mean particle size for a normal distribution
- σ standarå deviation

Subscripts

- o refers to incident light
- refers to the polarized component whose electric field vector is perpendicular to the plane of observation
- refers to the polarized component whose electric field vector is parallel to the plane of observation

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Submitted to the Department of Chemical Engineering on December 8, 1961 in partial fulfillment of the requirements for the Degree of Doctor of Science.

ABSTRACT

An investigation has been undertaken to develop a light-scattering technique for measuring the size and concentration of soot particles in flames. To this end, the theoretical light-scattering functions for arbitrary size spherical particles have been calculated for a range of parameters which correspond to those expected for soot particles. The numerical results were obtained with the help of an IBM 709 Computer. These light-scattering functions and associated coefficients are presented for the size parameter, $\pi d/\lambda$, ranging from 0.05 up to 2.00 in increments of 0.05 for angles from the incident light, 0, ranging from 0° to 180° in increments of 5°, and 14 values of the index of refraction which cover the expected range and includes the most likely value of m = 1.71 - 0.76 i for $\lambda = 4358 \text{Å}$.

A light-scattering apparatus has been constructed and used to obtain experimental data from a premixed laminar benzene-air flame. Data were also obtained from an amyl acetate diffusion flame and a Cambridge City Gas flame. The experimental light-scattering data indicate that the apparent goot particle size in the premixed benzene-air flame ranges from 1250A to 1800A in diameter, the larger sizes corresponding to larger elevations above the burner tip. Electron micrographs of soot particles collected on a cool metallic probe, however, show that the ultimate soot particles are only about 250A. This difference is believed to be due to the presence of agglomerated units of the ultimate soot particles. Additional experiments coupled with calculations, strongly indicate that the soot particle system within the flame consists of a very large number fraction of ultimate soot particles mixed with a very small fraction of agglomerates.

Further development of this technique is necessary before its full potential as a research tool for studying combustion problems can be realized. Recommendations for additional work are offered.

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I. SUMMARY

The object of this investigation was to develop a technique for determining the size and concentration of soot particles at various positions in laminar flames of hydrocarbons in air. This entailed the design and construction of an apparatus for measuring the polar diagram of the intensity of the light scattered by a flame illuminated by a strong incident light, the calculation of the theoretical light-scattering intensities which were applied to the experimental results to deduce soot particle size and concentration, and an evaluation of the results.

A schematic drawing of the light-scattering apparatus which was used for measuring the scattered-light intensities at various positions in a flame is shown in Figure 8.1. The scattered-light intensity was measured at 10° intervals between 30° and 150° from the direction of the incident light beam and in both planes of polarization at various elevations above the burner tip. The majority of the data were obtained from a fuel-rich premixed benzene-air flame. Additional data were also obtained from an amyl acetate diffusion flame and a Cambridge City Gas diffusion flame.

A machine program for calculating the desired theoretical lightscattering functions was developed using the Fortran coding system.

The calculated light-scattering functions for the perpendicular and
parallel components and the corresponding light-scattering coefficients
are presented in tabular form for practical ranges in the size

parameter and index of refraction which correspond to those expected for soot particles in flames. These calculated results are presented for angles from the incident light direction ranging from 0° to 180° in increments of 5° . The size parameter, $x(=\pi d/\lambda)$, ranges from 0.05 to 2.00 in increments of 0.05, while the range in the index of refraction covered the expected values for soot particles based on a study of the literature. Calculations were made for 14 values of the index of refraction and the above combinations of the range of the size parameters and angles. These calculations are for homogeneous, monodisperse spherical particles.

The procedure for calculating the light-scattering functions for a distribution of sizes of spheres is presented and a Fortran program for obtaining numerical results. A few results for a normal distribution of sizes are also presented but no attempt to use these results in the analysis was made.

The experimental light-scattering data were compared to the theoretical scattering patterns for a uniform particle system for various values of the size parameter and indices of refraction. A rather good fit of the theoretical scattering pattern to the experimental data was always possible for the perpendicular component with an appropriate value of the size parameter for a given value of the index of refraction. The comparison between the theoretical scattering pattern for a uniform size particle system for the parallel component and the corresponding experimental data is quite poor when the value

of the index of refraction is kept reasonably close to the expected value. A good fit of the theory to experiment was found for both the perpendicular and parallel components when the index of refraction was taken to be considerably different from the expected value.

Figure 10.8 shows the results for a typical run where the experimental light-scattering data are compared to theory for various values of the index of refraction m and the size parameter x for a premixed benzene-air flame burning with a fuel-air ratio of 2.39 times stoichiometric. The flame height is approximately 3.15 inches and the elevation above the burner tip at which the data were obtained is 1.25 inches. In this experiment the wavelength was 4358 Å. Figure 10.9 shows these same data compared to theory for a uniform size particle system for a value of the index of refraction of 1.40 - 1.00 i, which is considerably different from the expected value of 1.71 - 0.76 i.

The experimental light-scattering data indicate that the soot particles in the premixed benzene-air flame range from about 1250 Å to 1800 Å in diameter, depending on the elevation above the burner tip at which the data were obtained. In general, larger apparent particle sizes were obtained at progressively higher elevations in the flame.

Electron micrographs of soot particles collected on cool metallic probes, passed through the same flame and at the same elevation for which the light-scattering measurements were obtained, show that the ultimate soot particles from the premixed benzene-air flame are only

about 250 Å. It is expected that this difference between the apparent size deduced from the light-scattering measurements and that observed from the electron micrographs is due to agglomeration of the ultimate soot particles within the flame. The apparent increase in particle size with increasing elevation above the tip of the burner is consistent with the thought of agglomerated particles within the flame, since the agglomerated soot particles could be expected to become larger at progressively higher elevations due to additional agglomeration.

A pair of runs on a flame which had essentially the same test conditions but different incident light wavelengths for each run, show that the ratio of the values of the apparent size parameters for the two runs is quite different from the inverse ratio of the respective wavelengths, as would be expected for a uniform size particle system. Instead of the value of 1.25, which is the inverse ratio of the respective wavelengths used, a ratio of 1.11 was found. In one attempt to explain this result, calculations were made for an assumed soot particle system which was made up of a small number of large particles (representing the agglomerated soot particles) mixed with a very large number of small particles (representing the small ultimate soot particles). The light-scattering patterns were calculated for the same two wavelengths as were used in the experiment and based on this supposed particle system. These were then compared to the theoretical scattering patterns for various uniform size particle systems

to find the best fit for the two cases. The results of these calculations show that the ratio of the apparent size parameters for the two different wavelengths could be as found by experiment for the proper choice of the number ratio of large-to-small particles and the respective absolute sizes.

The experimental results obtained herein from an amyl acetate diffusion flame are consistent with the results of Senftleben and Benedict. Here again, there is a strong indication that the soot particles exist in the flame as agglomerates mixed with a rather large fraction of small ultimate spherical soot particles.

The results of the concentration measurements are in question because the nature of the soot particle distribution is not known. The apparent number concentrations of the soot particles deduced from experiments are about 10^8 particles per cc but are likely orders of magnitude greater.

Additional development of the light-scattering technique is necessary before its full potential as a research tool for studying practical combustion problems can be realized. It should be pointed out that this technique has potential application to several other problems involving study of growth and decay rates of liquid and solid particles in reacting systems in addition to the soot formation process in flames.

II. INTRODUCTION

The formation of soot particles in flames is important in many combustion applications. The commercial production of carbon black is one example. Carbon black produced by incomplete combustion of hydrocarbons is used extensively in rubber compounding, as a pigment in inks and paints, and has many other uses. A problem in the operation of a turbojet combustion chamber is the formation of free soot with the subsequent fouling of the fuel nozzles and blocking of the air ports. Still another example is the formation and subsequent combustion of soot particles in a luminous flame such as is required for the high heat transfer in industrial furneces. The governing physical and chemical processes involved for the various conditions in which soot is formed are not well understood.

A great deal of effort has been directed toward the understanding of the soot formation process over the years as evidenced by the appearance of more than a hundred publications dealing with this problem directly. While working with a miner's safety lamp, Davy, in 1816, concluded that the luminosity of flames is caused by the production and ignition of solid matter in flames (1). At the present time it is generally agreed that the luminosity of hydrocarbon flames is due to the presence of soot particles. It is well known that a sample of soot can be collected by passing a cold probe through a luminous but nonsmoking flame.

Coward and Woodhead (2) measured the luminosities of flames of a rather large variety of chemical compounds to determine the effect of chemical structure of the fuel on the luminosity. The results show that chemical structure of the fuel has a very marked effect on the luminosity of a flame. Hunt (5) has also examined a large number of compounds in order to determine the effect of molecular structure on scot formation. In his work the smoke points of the various fuels were determined. (The smoke point is defined as the height in millimeters of the highest flame produced without smoking when the fuel in question is burned in a standard burner.) In addition to making a systematic study of the smoking tendency of many different hydrocarbon fuel types, Schalla, Clark, and McDonald (4) also studied the effect of pressure, oxygen supply, degree of preheat, and fuel-flow rate in premixed flames.

Fenimore, Jones, and Moore (5) studied scot formation in quenched flat flames of propane and butane. They found that the yellow carbon luminescence of a fuel-rich flat flame fades abruptly as the pressure is lowered to a critical pressure at a fixed mass flow and fuel-air ratio. This critical pressure was higher for lower fuel-air ratios or for greater mass flows per unit area of the burner.

Arthur and Napier (6) examined the physical properties of soot collected from various hydrocarbon flames. They find marked changes in the nature of the collected soot by suitable modification of the combustion conditions.

Stehling, Frazee, and Anderson (7) find from electron microscope observations of collected acetylene soot that the small particles are grouped into aggregates to form chain-like structures. These small particles ranged from 800 to 2000 Å in diameter. Parker and Wolfhard (8) measured, under an electron microscope, the size of soot particles collected from an amyl acetate diffusion flame. They found that the soot particles from a nonsmoky flame were nearly all the same size and approximately 100 Å in diameter. These particles were found to be built together into filaments or chains. They also report that the particles collected near the top of the flame were similar in size to those near the base. Their study of a smoky amyl acetate flame, however, showed some particles with diameters of about 500 Å. Gaydon and Wolfhard (9) indicate that the size of soct particles in flames range from 100 to 2000 Å in diameter.

Several mechanisms have been proposed for the soot formation process, but a firm understanding of this subject is certainly lacking. The question of a soot formation mechanism has recently been reviewed by Gaydon and Wolfhard (9).

One method for determining the size of soot particles in flames involves removing a sample of soot from the flame and examining it under an electron microscope. This type of study has been made by Tesner, Robinovitch, and Rafalkes (10) wherein two methods were used to remove soot particles from different parts of a methane diffusion flame with and without addition of other compounds. One method

entailed the collection of soot on a cold surface introduced into the flame at various heights above the burner. The other method involved collecting the soot particles in a bag filter which was preceded by a quenching lattice made of water-cooled metallic capillaries. The mean diameters of the particles removed from different parts of the flame were in the range of 200 to 400 Å.

The method of determination of soot particle sizes by physically probing a flame with subsequent examination under an electron microscope has the obvious undesirable feature of probe interference. It would be desirable to have a method for determining soot particle sizes in an undisturbed flame.

In 1919, Senftleben and Benedict (11) measured the scattered-light intensity generated by focusing light from a carbon-arc source in the luminous zone of a Hefner lamp burning amyl acetate. The scattered-light intensity was measured at various angles from the incident light beam. From a comparison of the measured scattered-light intensities with the light-scattering theory for arbitrary size spheres due to Mie (12), Senftleben and Benedict deduced a soot particle size of 1750 Å in their experiment. The scattering volume in the flame was comparable with the total flame volume so that their results represent some sort of average size weighted over a large portion of the flame. A method for determinating the soot particle size based on light-scattering results avoids the complication and uncertainty caused by probe disturbances.

It was the object of the present study to develop a lightscattering technique for determining the soot particle size and concentration in flames. This entailed the design and construction of an
apparatus for measuring the scattered-light intensity from a laminar
flame illuminated by a strong incident light, the calculation of the
theoretical light-scattering intensities which, when applied to the
experimentally determined intensities, permit deduction of soot
particle size and concentration, and finally an evaluation of this
technique for studying the soot formation process.

The scope of this work includes a presentation of the lightscattering theory for the expected size range of soot particles. This
is followed by a discussion of the application of light-scattering
theory to experiment for determining the soot particle size and concentration. Also, a discussion of the design considerations and
description of an apparatus for measuring the light-scattering intensity from a luminous flame are given. The experimental results are
presented, followed by a discussion and evaluation of these results.
Finally, the conclusions drawn from this study and some recommendations
for future work on this problem are presented.

While this work has the objective of developing a light-scattering technique for determining the soot particle size and concentration in flames, there are certainly several other practical systems to which this type of study could be applied. For example, the application of the light-scattering technique to a study of the growth

and decay of particles within rocket exhausts looks promising.

Durbin (13) has presented a method for measuring the size and concentration of particles in a supercooled hypersonic stream. Quite recently, Kaskan (14) has reported some preliminary light-scattering measurements on particles of liquid boron exide condensing in a trimethyl borate-air flame. The scattering measurements presented in reference 14 were made entirely by determining the extinction as a function of wavelength.

III. LIGHT-SCATTERING THEORY

Light can be represented as an electromagnetic wave and thus has a transverse oscillating electric field associated with it. When a beam of light passes through a particle, this oscillating electric field, which is perpendicular to the direction of the incident light, causes the electric charges contained within this particle to be set into forced oscillation with a frequency equal to that of the incident light. These oscillating electric charges, in turn, are each sources of electromagnetic radiation, scattered light.

Physical Description of a Simple Scattering Process

Consider the simple case shown in Figure 3.1 where a beam of parallel light is incident on a single scattering particle which is very much smaller than the wavelength of the incident light. Unpolarized parallel light enters from the left along the z-axis and is partially scattered by the very small particle at point p. Since the oscillating electric field of the incident light is always perpendicular to the direction of the incident light, the oscillating charges move only in a plane which is perpendicular to the z-axis and passes through point p. This plane is defined in Figure 3.1.

A reference plane containing the z-axis is also shown and is arbitrarily chosen to correspond to the plane of the page. This second plane is denoted as the plane of observation. (This designation will become apparent later.)

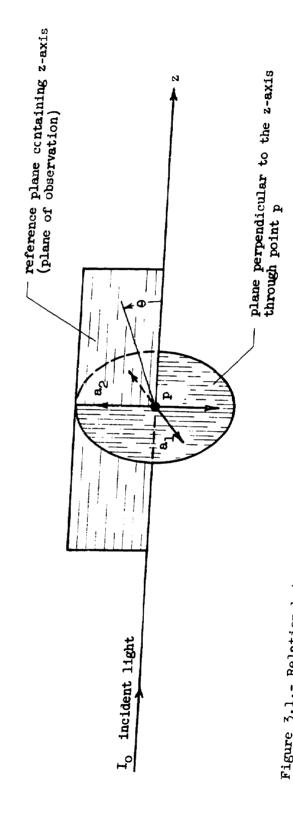


Figure 3.1.- Relation between incident light, plane of observation, and plane in which charges

The components a₁ and a₂ of the vector which defines the direction and amplitude of the periodic motion of the oscillating charges are indicated in Figure 3.1 and lie in the plane perpendicular to the z-axis through point p. The direction of a₁ is perpendicular to the plane of observation and a₂ is parallel to the plane of observation.

If the incident light is unpolarized, the amplitude of al will equal a2. On the other hand, if the incident light is plane polarized with its electric vector perpendicular or parallel to the plane of observation, the amplitude a2 or al will be respectively zero.

The amplitude and direction of the two components of the electric field vector associated with the scattered light, E_1 and E_2 , which are generated by the oscillating charges, are proportional to the projected amplitudes of a_1 and a_2 which lie perpendicular to the direction at which the scattered light is observed. Referring again again to Figure 3.1, it is seen that the amplitude a_1 , which is perpendicular to the plane of observation, does not vary with the angle θ measured from the forward direction of the incident light. Since the intensity of light is proportional to the square of the amplitude of the associated electric field, the plane polarized scattered-light intensity I_1 whose electric field vector is perpendicular to the plane of observation, is independent of the angle of observation θ .

The plane polarized scattered-light intensity, I_2 , whose electric field vector is parallel to the plane of observation, however, does depend on θ because of the variation with θ of the projected amplitude of a_2 , that is a_2^i , which lies perpendicular to the direction of observation as shown in Figure 3.2. It is clear from Figure 3.2 that the projected amplitude of a_2 , that is a_2^i , which lies perpendicular to the direction of observation of the scattered light, is proportional to $\cos \theta$. Since E_2 is proportional to a_2^i , and the intensity I_2 is proportional to the square of E_2 , the scattered-light intensity I_2 , whose electric field vector is parallel to the plane of observation, is proportional to $\cos^2 \theta$. If one simultaneously observes the total scattered-light intensity, $I_1 + I_2$, the angular dependence is $1 + \cos^2 \theta$.

Figure 3.3 shows the polar distribution of the scattered-light intensity for the two plane polarized components, I_1 and I_2 , and for the total scattered light $I_1 + I_2$ as a function of θ when the scattering particle is very much smaller than the wavelength of the incident light. The perpendicular component I_1 is independent of θ , but the parallel component I_2 is a maximum and equal to I_1 at $\theta = 0^\circ$, decreases to zero at $\theta = 90^\circ$, and increases again to a maximum at $\theta = 180^\circ$. The variation in the total light-scattering intensity is the average of the separate component intensities I_1 and I_2 .

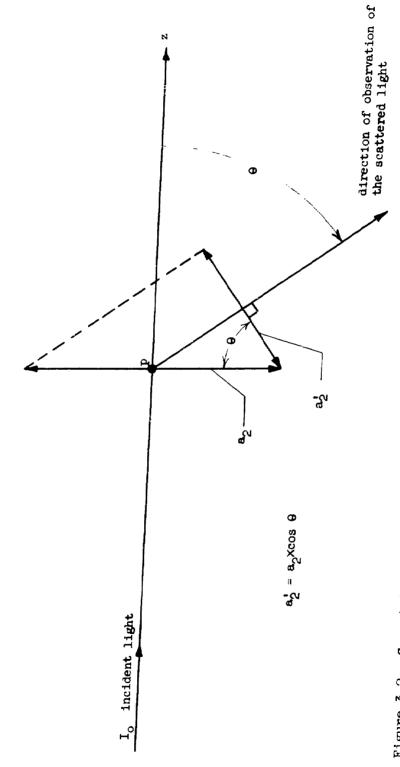


Figure 3.2.- Geometric relation between the component of the oscillating vector in the plane of observation and the amplitude of its component which is perpendicular to the direction of observation of the scattered light.

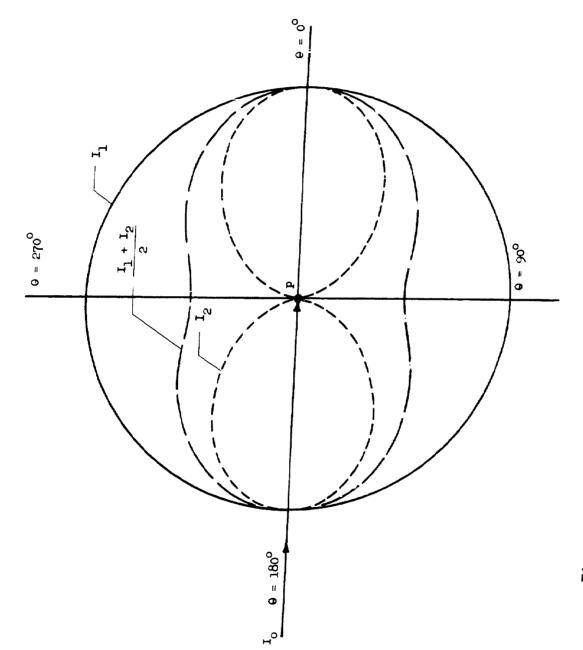


Figure 3.3.- Scattered-light intensity distribution for a << λ .

Rayleigh Scattering Equations

The analytical expressions for the intensity of scattered light from particles much smaller than the wavelength of the incident light were given by Lord Rayleigh in 1899. The equations for the two plane polarized components of the scattered-light intensity for particles much smaller than the wavelength of the incident light are, according to Sinclair and LaMer (15)

$$I_{1} = \frac{9\pi^{2}}{r^{2}} \left| \frac{m^{2} - 1}{m^{2} + 2} \right|^{2} \frac{v^{2}}{\lambda^{4}} I_{0,1}$$
 3.1

for the perpendicular component, and

$$I_2 = \frac{9\pi^2}{r^2} \left| \frac{m^2 - 1}{m^2 + 2} \right|^2 \frac{v^2}{\lambda^4} \cos^2 \theta I_{0,2}$$
 3.2

for the parallel component. In these equations, I_1 and I_2 are the plane polarized scattered-light intensities measured at a distance r from the scattering particle, whose electric field vectors are perpendicular and parallel to the plane of observation, respectively. The corresponding plane polarized incident light intensities are denoted by $I_{0,1}$ and $I_{0,2}$. The distance from the scattering particle to the point of observation of the scattered-light intensity is r. The index of refraction of the scattering particle relative to the surrounding medium is m, the volume of the scattering particle is V, and the wavelength of the incident light is denoted by λ .

The total intensity of the scattered light when the incident light is unpolarized is

$$I = \frac{9\pi^2}{2r^2} \left| \frac{m^2 - 1}{m^2 + 2} \right|^2 \frac{v^2}{\lambda^4} (1 + \cos^2\theta) I_0$$
 3.3

where I is the total scattered-light intensity and $I_{\rm O}$ is the incident-light intensity. It is seen that the angular dependence of the scattered-light intensity as given by the above equations is the same as that already deduced by physical argument.

In the foregoing equations and in the light-scattering equations to follow, the intensity I is understood to be the energy flux per unit area. In c.g.s. units the intensity is in ergs per cm² per sec. It should be pointed out that this definition of intensity, which has been used in light-scattering work in the past and which will be used herein, is called illuminance in optics.

The consequence of the limitation of the scattering particle size to much less than the wavelength of the incident light is the absence of interference due to phase differences between the scattered light generated by oscillating charges located at various positions in the particle. There will be interference between the scattered light generated by the various oscillating charges contained in a scattering particle which can no longer be treated as very small compared to the wavelength. In addition to this effect, the applied electric field from the incident light, which causes the charges within the particle

to oscillate, will be distorted by a scattering particle which is not very much less than the wavelength. As particles of larger size are considered, the effects of interference and distortion become more important and change the angular distribution of the scattered-light intensity.

Mie Scattering Theory

The rigorous scattering theory for spheres of arbitrary size was developed by Mie (12) in 1908. A brief derivation of the Mie theory is presented by Van de Hulst (16). The scattered-light intensity for the two plane polarized components and the total scattered light are given by the following equations. For perpendicular polarization

$$I_1 = \frac{i_1}{k^2 r^2} I_{0,1}$$
 3.4

for parallel polarization

$$I_2 = \frac{i_2}{k^2 r^2} I_{0,2}$$
 3.5

and for the total scattered light when the incident light is unpolarized

$$I = \frac{i_1 + i_2}{2k^2r^2} I_0$$
 3.6

where $k(=2\pi/\lambda)$ is the wave number, i_1 and i_2 are light-scattering functions which are discussed below, and the other variables have the same meaning as before.

The scattering functions i_1 and i_2 which are directly proportional to the scattered-light intensity, are rather involved functions of x, y, and θ . Formally,

$$i_1 = f_1(x, y, \theta)$$
 3.7

$$i_2 = f_2(x, y, \theta)$$
 3.8

where x is the size parameter defined as the ratio of the circumference of a particle divided by the wavelength of light, y is the product of x and the index of refraction of the scattering particle relative to the surrounding medium, and θ is the angle at which the scattered-light intensity is measured relative to the forward direction of the incident light. The rather involved functional relationships for i_1 and i_2 are given in appendix A.

It is seen from equations 3.4 through 3.8 that the angular distribution of the scattered-light intensity, normalized with respect to the incident light intensity, is a function only of x, y, and the state of polarization if λ and r are fixed. This in turn means that the shape of the scattering diagram (I_1 , I_2 , or I plotted as a function of θ) is determined by the values of the particle

diameter d and the particle index of refraction m. If the value of m is known, the particle diameter alone determines the shape of the scattering diagram.

Scattering by a Cloud of Particles

Thus far only the scattering by a single particle has been considered. Since this present work is concerned with measuring the scattered-light intensity from a small volume of scot particles within a flame, it is necessary to know the relationship between the scattering from a many particle system and the single particle case. Provided the concentration of the scattering particles and the distances through which the incident and scattered light must travel are not too large, the intensity of light scattered by a cloud of particles is directly proportional to the number of scattering particles. In other words, the proportionality between the scattered-light intensity and the number of particles holds only if the light incident on each particle is essentially the light of the original incident beam.

As the concentration of the particles increases or the distance through which the incident and scattered light travel increases, more and more of the incident light undergoes extinction by the other particles and the more each particle will be exposed to the scattered light from the other particles. When these effects become important, the proportionality between the scattered-light intensity and the number of particles is no longer a valid assumption. This type of scattering is known as multiple scattering.

If the extinction of the incident light is negligible in passing through a cloud of particles, it may be assumed that multiple scattering effects are unimportant. In order to determine whether or not the effects of multiple scattering are important in this work, the extinction was calculated for a given flame size and for the expected range of soot particle size and concentration. The results indicate that multiple scattering effects may be neglected.

The next section will be concerned with the actual calculation of the light-scattering functions.

IV. CALCULATION OF LIGHT-SCATTERING FUNCTIONS

One of the first steps in a light-scattering investigation of this kind is to calculate the scattering functions over the expected ranges of x and y for various values of θ . As before, $x = \frac{\pi d}{\lambda}$ and y = mx. The expected diameter d of soot particles, according to Gaydon and Wolfhard (9), ranges from 100 to 2000 Å. A combination of a light source and an interference filter was chosen, for reasons given in a later section, to give approximately monochromatic light of 4358 Å. The range of the size parameter x for this combination of d and λ is about 0.07 to 1.44. The values of x over which the scattering functions are calculated range from x = 0.05 to x = 2.00 in increments of 0.05.

The second parameter required, y, depends directly on the index of refraction, m, of the scattering particle. There appear to be no data available on the index of refraction of soot particles in a flame specifically, however, there is a limited amount of data for carbon. In this work it is assumed that the index for carbon closely approximates that of soot. Carbon particles are somewhat opaque and therefore absorb a portion of the incident light as well as scattering it. This absorption must be accounted for in the calculation of the light-scattering functions. This is done by representing the index of refraction m as a complex number.

The real part n has to do only with the velocity of light in the particle and is equivalent to the usual index of refraction for transparent materials, that is, the ratio of the velocity of light in a reference medium, usually vacuum, to that in the material in question. The imaginary part n' has only to do with the absorption of light by the particle and is known as the absorption constant. The absorption constant n' is related to the usual absorption coefficient γ by the relation

$$\gamma = \frac{4\pi n!}{\lambda}$$

as shown in appendix B.

Senftleben and Benedict (17) measured the index of refraction of what they call amorphous carbon at room temperature for various wavelengths. Table 4.1 lists the values of m reported and indicates that m is only a weak function of λ .

TABLE 4.1.- INDEX OF REFRACTION OF AMORPHOUS CARBON FROM SENFTLEBEN AND BENEDICT (17).

λ .	m
4360 A	1.90 - 0.68 i
4920	1.94 - 0.66 i
5460	1.96 - 0.66 i
5780	1.97 - 0.65 i
6230	2.00 - 0.66 i

In order to take into account the effect of temperature on m, Stull and Plass (18) have derived a dispersion equation, based in part on the room temperature data of Senftleben and Benedict, from which the index of refraction at 2250° K can be calculated. Values of m calculated from the dispersion equation given by Stull and Plass for two wavelengths are shown in Table 4.2.

TABLE 4.2.- INDEX OF REFRACTION OF AMORPHOUS CARBON CALCULATED FROM EQUATION BY STULL AND PLASS (18)

λ	m
5348 R	1.71 - 0.76 i
546 1	1.79 - 0.79 i

In addition to the above data, Wartenberg (19) has measured the index of refraction at room temperature of a freshly cut surface of graphite. He reports an index of refraction of m=2.98-1.74 i for a wavelength of $\lambda=5790$ Å.

The values of the index of refraction m for carbon which were determined in both (17) and (19) are based on a reflection method which measures the state of polarization of light reflected from a surface of the sample carbon as a function of the angle of incidence. A general description of the method for determining the values of n and n' from reflection measurements is given in reference 20. Tool (21) points out that the condition of the surface effects the values of n and n'. In addition to the question of surface effects,

there is uncertainty concerning the similarity of the structure of soot particles compared to the actual carbon samples examined.

Warren (22) has made X-ray diffraction studies of carbon black. He indicates that carbon black or soot is not a truly amorphous form of carbon. Carbon black does have a graphitic form but differs from graphite in that the successive layers are not parallel but are oriented with respect to each other in a random manner. The extent of the crystalline or amorphous character of soot depends on the fuel and conditions under which the soot is formed.

without actually measuring the index of refraction of soot particles in a flame, the correct value of m can only be estimated from the above data. The most recent work of Stull and Plass (18) was assumed best to represent the optical properties of soot. The particular values of m, on which the light-scattering calculations are based, were taken to include a broad possible range. The 14 values of m are listed later along with the calculational procedure.

The last independent parameter in the scattering expressions, θ , is varied from 0° to 180° in 5° increments for each combination of x and m (or y).

Gucker and Cohn (23) have developed a general scheme for calculating the desired scattering functions by starting with the expressions already given in appendix A. The scheme is shown in appendix C.

Based on the general scheme of Gucker and Cohn, a detailed system of equations was developed for the numerical evaluation of the

desired scattering functions and coefficients. A machine program, using the Fortran programing system, was then developed which applied this system of equations for the chosen values of x, m, and θ . The detailed system of equations and the Fortran program are also present in appendix C. The calculations were carried out at the MIT Computation Center on an IBM 709 computer.

V. NUMERICAL RESULTS OF LIGHT-SCATTERING CALCULATIONS

The light-scattering functions i_1 and i_2 for spherical particles have been calculated with the help of an IBM 709 computer for values of the size parameter x ranging from x = 0.05 to x = 2.00 in increments of 0.05 for values of θ ranging from $\theta = 0^{\circ}$ to $\theta = 180^{\circ}$ in increments of 5° and for the following values of θ :

1.60 - 0.60 i	1.80 - 1.00 i	
1.60 - 0.80 i	1.90 - 0.68 i	
1.60 - 1.00 i	1.96 - 0.66 i	
1.71 - 0.76 i	2.00 - 0.60 i	
1.79 - 0.79 1	2.00 - 0.80 i	
1.80 - 0.60 i	2.00 - 1.00 i	
1.80 - 1.00 i	2.98 - 1.74 i (graphite)

In addition to the functions i_1 and i_2 , the scattering coefficients a_n and b_n which are intermediate parameters in the program for calculating i_1 and i_2 , were also printed as output results from the machine computations. The numerical results for the above mentioned combinations of x, m, and θ are presented in tabular form in appendix D. A tabulation of a_n and b_n for the corresponding values of x and m for n up to θ are also given in appendix D.

Based on these results, polar scattering diagrams of i_1 and i_2 plotted against θ for m = 1.71 - 0.76 i and various selected values of x are presented in Figures 5.1 through 5.4.

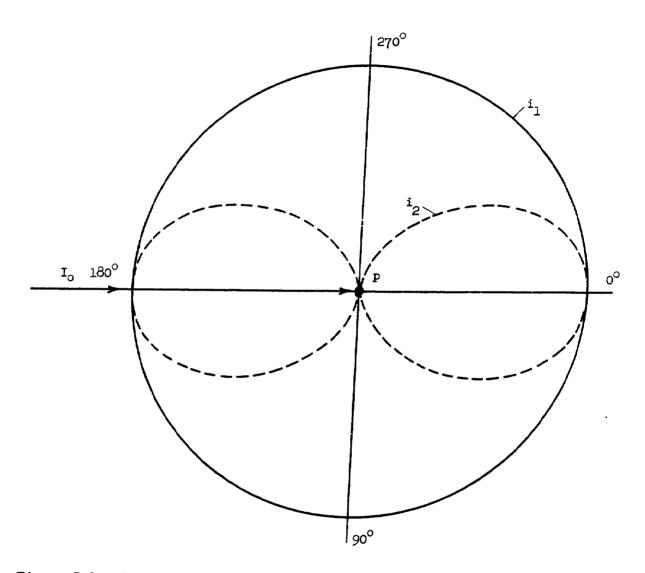


Figure 5.1.- Polar light-scattering diagram for m = 1.71 - 0.76 i and x = 0.05.

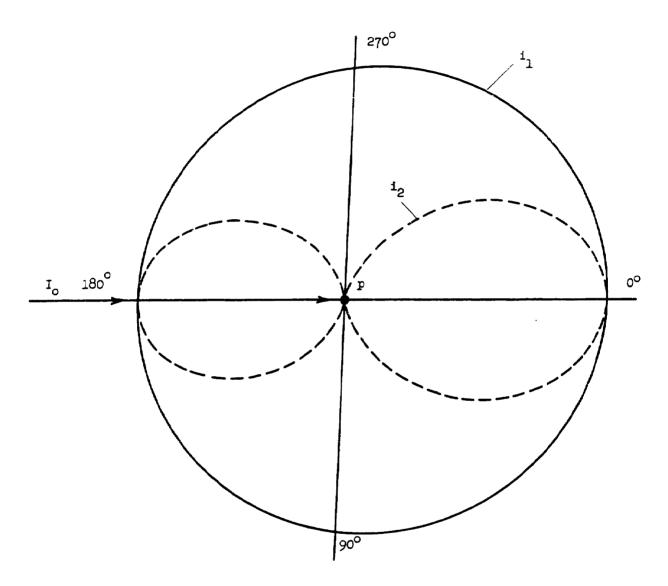


Figure 5.2.- Polar light-scattering diagram for m = 1.71 - 0.76 i and x = 0.50.

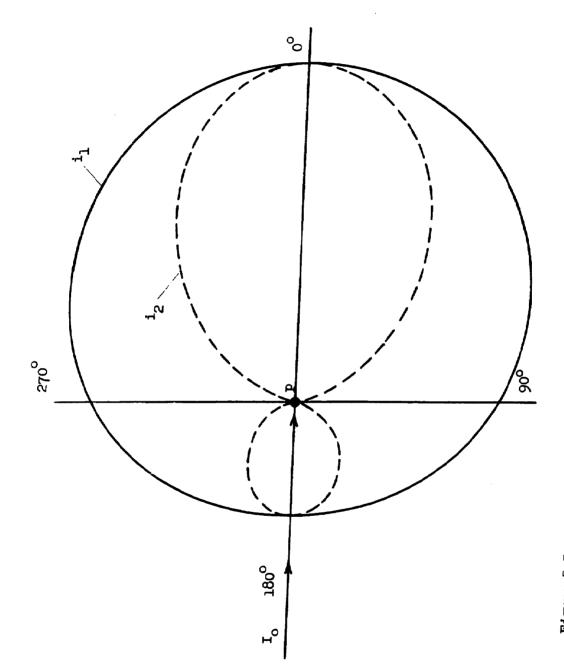


Figure 5.3.- Polar light-scattering diagram for m = 1.71 - 0.76 i and x = 1.00.

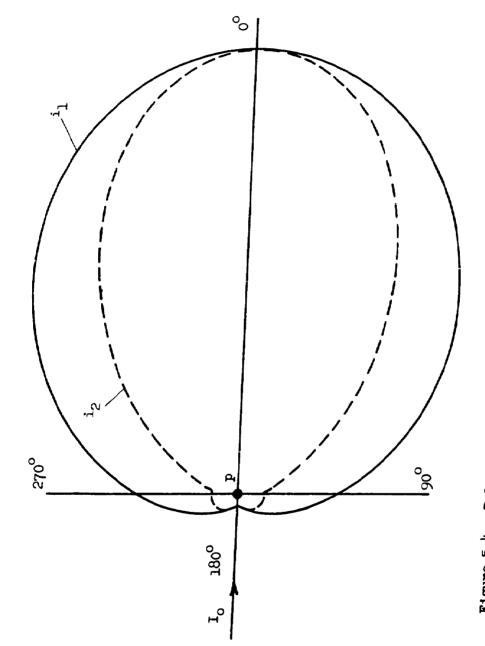


Figure 5.4. Polar light-scattering diagram for m = 1.71 - 0.76 i and x = 1.50.

Figure 5.1 is based on x = 0.05 which represents a case where the particle size is much less than the wavelength of the light used. If the wavelength λ is 4358 Å, the particle diameter is only about 70 Å when x = 0.05. This plot shows that the scattered-light intensity for this value of x is symmetrical about the 90° to 270° axis as it is for the case for all small particles which satisfy the condition $d \ll \lambda$. The shape of the light-scattering diagram of i_1 and i_2 against θ for this case (x = 0.05) is the same as that given in Figure 3.3 for any very small particle for which $d \ll \lambda$. It is also of interest to compare the numerical values of i_1 and i_2 obtained by machine calculation with the values that would be found by using the simple scattering equations based on $d \ll \lambda$.

Equate the right-hand side of equation 3.1 to that of equation 3.4. This gives

$$i_1 = \left(\frac{m^2 - 1}{m^2 + 2}\right)^2 x^6$$
 5.1

where V in equation 3.1 has been replaced by $\frac{\pi d^3}{6}$ and $\frac{\pi d}{\lambda}$ has been replaced by x. This is the equation for i_1 only if $d \ll \lambda$. The corresponding equation for i_2 is

$$i_2 = \left(\frac{m^2 - 1}{m^2 + 2}\right)^2 x^6 \cos^2\theta$$
 5.2

The value of i_1 as calculated from equation 5.1 for x = 0.05 and

m = 1.71 - 0.76 i is approximately 0.5220×10^{-8} and is independent of θ . In the way of comparison, the more exact value of i_1 determined by the machine computation shows a very slight dependence on θ and ranges from $i_1 = 0.5232 \times 10^{-8}$ at $\theta = 0^{\circ}$ down to $i_1 = 0.5219 \times 10^{-8}$ at $\theta = 180^{\circ}$ as shown in appendix D. The results are essentially the same and agree within 0.23 percent.

A similar comparison for i_2 can be made at a given value of θ , say 60° . It follows from equation 5.2 that $i_2 = 0.1305 \times 10^{-8}$. The machine computation result for $\theta = 60^{\circ}$ is $i_2 = 0.1307 \times 10^{-8}$. This is likewise good agreement.

Figure 5.2 is a scattering diagram for the same index of refraction but x = 0.50. It can be seen that the light-scattering intensity in the forward direction ($\theta < 90^{\circ}$) is somewhat greater than that in the backward direction ($\theta > 90^{\circ}$). As the size parameter x is increased further to x = 1.00 for Figure 5.3 and to x = 1.50 for Figure 5.4, the dissymmetry of the scattering diagram increases, that is, the forward scattered-light intensity increases relative to the backward intensity with increasing x.

The effect that the size parameter x has on this dissymmetry is quite clearly shown by plotting the ratio of the forward scattered-light intensity at a given angle θ divided by the backward scattered-light intensity corresponding to the angle $(180^{\circ} - \theta)$. Dissymmetry plots for m = 1.71 - 0.76 i are presented in Figures 5.5 through 5.8 which show the ratios $i_1(\theta)/i_1(180 - \theta)$ and $i_2(\theta)/i_2(180 - \theta)$ plotted against x for four values of θ .

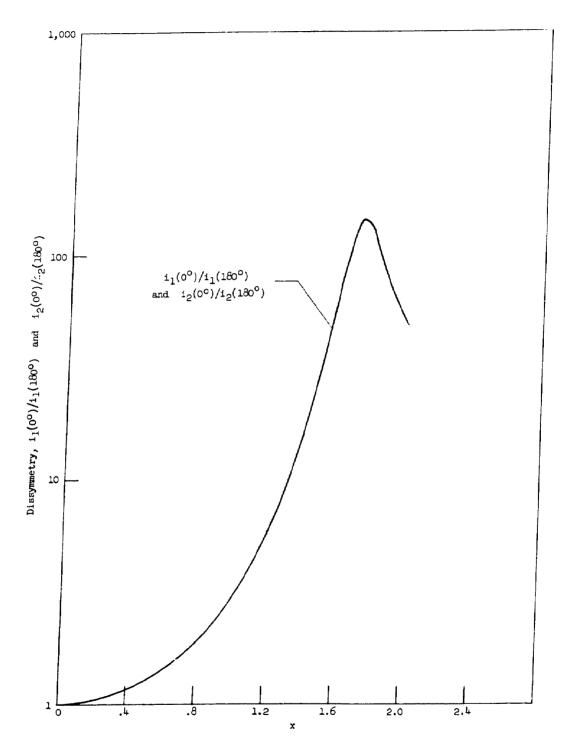


Figure 5.5.- Dissymmetry, $i_1(0^\circ)/i_1(180^\circ)$ and $i_2(0^\circ)/i_2(180^\circ)$ plotted against x and for m = 1.71 - 0.76 i.

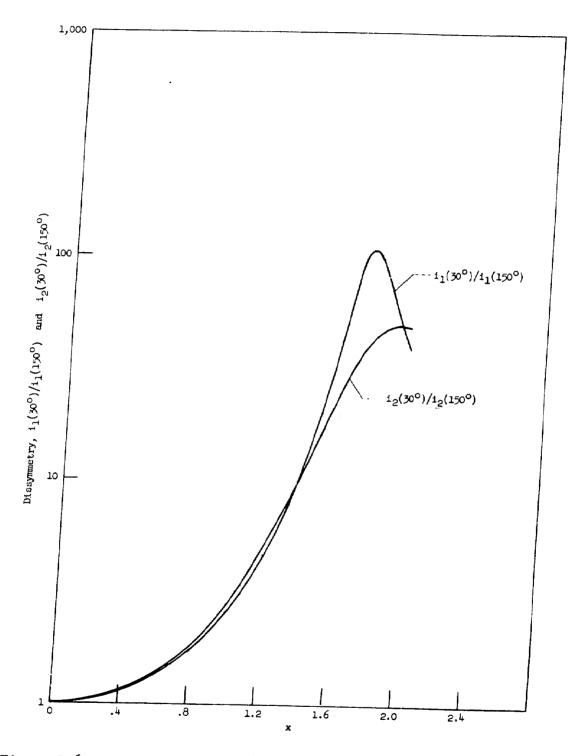


Figure 5.6.- Dissymmetry, $i_1(30^\circ)/i_1(150^\circ)$ and $i_2(30^\circ)/i_2(150^\circ)$ plotted against x for m = 1.71 - 0.76 i.

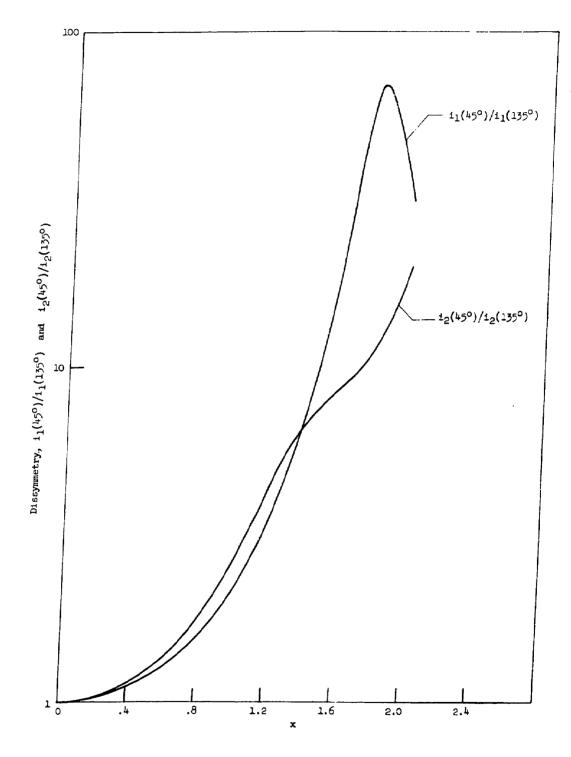


Figure 5.7.- Dissymmetry, $i_1(45^\circ)/i_1(135^\circ)$ and $i_2(45^\circ)/i_2(135^\circ)$ plotted against x for m = 1.71 = 0.76 i.

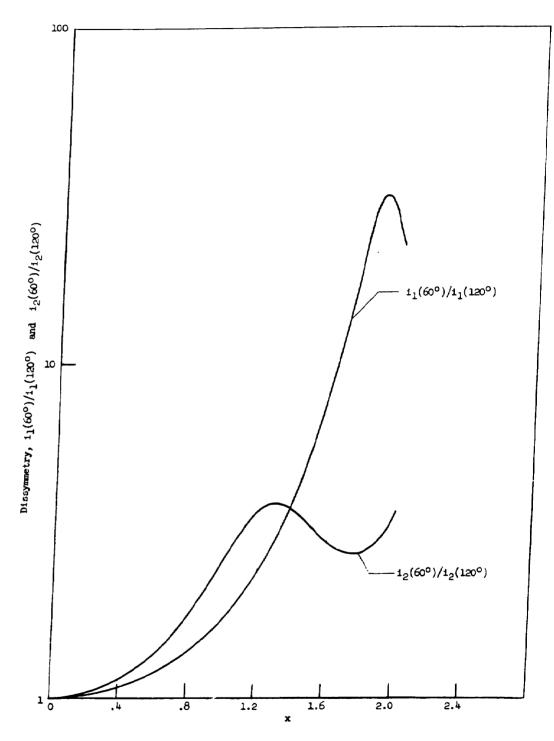


Figure 5.8.- Dissymmetry, $i_1(60^\circ)/i_1(120^\circ)$ and $i_2(60^\circ)/i_2(120^\circ)$ plotted against x for m = 1.71 - 0.76 i.

VI. APPLICATION OF THE CALCULATED LIGHT-SCATTERING RESULTS

In the previous section, the theoretical light-scattering functions were determined and tabulated for the range of soot particle sizes expected in flames. It is the purpose of this section to show how these theoretical results can be applied to light-scattering measurements for determining the size and concentration of soot particles in flames. Before discussing various possible approaches, it should be remembered that, while the primary objective of this work was to develop a technique for measuring soot particle size and concentration in flames, it was done with the hope of eventually being able to follow the growth and decay of soot particles within a flame. Here it is important to keep the size of the observed scattering volume as small as possible in order to reduce size averaging effects due to the presence of different particle sizes at different locations in the flame.

It is assumed that the soot particles are spherical and monodisperse so that the theoretical results of the previous section can be applied directly.

Particle Size Determination

There are several different light-scattering measurements, which, when combined with the theoretical scattering results, can be made to determine particle size, if the particle diameter falls in a given range.

It can be observed from Figures 5.1 through 5.4 that the shape of the light-scattering diagram is strongly dependent on the particle diameter so long as x is not too small. Stacey (24) suggests that the approximate practical lower limit of x for which dissymmetry can be measured corresponds to a particle which has one dimension greater than 1/20 of the wavelength of the incident light, that is, x of about 0.16 for a sphere.

For soot particles which are large enough to produce a measurable unsymmetrical scattering pattern, the size parameter x can be determined by measuring the light-scattering as a function of θ and comparing the resulting pattern with the theoretical scattering diagrams corresponding to different values of x. The value of x which corresponds to the theoretical pattern which best fits the experimental data is then the correct value of x. This method of course assumes a knowledge of the correct values of m.

This is the method used in this work for determining the soot particle size. The two plane polarized scattered-light intensities I_1 and I_2 were measured with the aid of a polarization filter and a relatively narrow band pass optical interference filter. For a cloud of N scattering particles which are monodispersed and not too concentrated, the two plane polarized light-scattering intensities I_1 and I_2 follow from equations 3.4 and 3.5.

$$I_1 = N \frac{I_1}{k^2 r^2} I_{0,1}$$
 6.1

and

$$I_2 = N \frac{i_2}{k^2 r^2} I_{0,2}$$
 6.2

In a given experiment, k is fixed and known from the narrow band pass optical interference filter, the state of polarization is fixed and known from the setting of the polarization filter, N is fixed, r is fixed and known, $I_{0,1}$ and $I_{0,2}$ are fixed, and I_1 and I_2 are measured. Then, if m is known, the shape of the scattering patterns of I_1 against θ and I_2 against θ will be determined by d only, since i_1 and i_2 are functions only of x, m, and θ . A comparison of the measured scattered-light intensity distribution with the various calculated distributions of i_1 and i_2 for a given value of m but various values of x then gives the particle diameter d since λ is known.

An alternative method for determining the value of x is to measure the ratio of the scattered-light intensities at two angles symmetric about $\theta = 90^{\circ}$, that is, at $\theta < 90^{\circ}$ and also at an angle of $(180^{\circ} - \theta)$. It is noted by referring back to Figures 5.5 through 5.8, that the dissymmetry, $i(\theta)/i(180^{\circ} - \theta)$ for either perpendicular or parallel polarization is a smooth function of x. The experimental determination of the dissymmetry for the two plane polarized components followed by a comparison with the theoretical dissymmetry yields the particle size. This method in which the dissymmetry is measured is just a special case of the previous method.

Still another method of size determination involves measuring the transmittance of a beam of light through the flame at several different wavelengths. The transmittance of light passing through a cloud of particles is the ratio of the intensity of the transmitted light intensity I to the incident intensity $I_{\rm O}$ and is given by the equation

$$\frac{I}{I_0} = \exp(-\gamma l) \tag{6.3}$$

where l is the optical path through the scattering cloud and γ is the extinction or absorption coefficient which is defined as

$$\gamma = \left(\frac{N}{V}\right) c_{\text{ext}}$$
 6.4

In this expression, $\left(\frac{N}{V}\right)$ is the number of particles per unit volume in the optical path and $C_{\rm ext}$ is the extinction cross section per particle which in turn is a function of x and λ . It should be remarked that the transmittance is independent of the state of polarization.

The extinction cross section is defined as

$$C_{\text{ext}} = \frac{2\pi}{k^2} \sum_{n=1}^{\infty} (2n + 1) \text{Re}(a_n + b_n)$$
 6.5

where $Re(a_n + b_n)$ denotes the real part of $a_n + b_n$. The scattering coefficients a_n and b_n are given in appendix D for various values of x and m for values of n up to six. The extinction

coefficient $C_{\rm ext}$ can therefore be calculated from equation 6.3 for various values of x at different wavelengths for a known value of m. The quantities I, I_O, and l can be measured at several different wavelengths so that γ can be determined at these wavelengths from equation 6.3. A plot of the logarithm of the measured value of γ against the various values λ used is then compared to the calculated curve of $C_{\rm ext}$ against λ for various values of d to find the best fit without regard to the absolute magnitude of the γ or $C_{\rm ext}$ since $\binom{N}{V}$ in equation 6.4 is constant for a given experiment. The value of d which corresponds to the theoretical curve which best fits the experimental curve is the correct value.

This method, which involves measurement of the extinction coefficient, was not used in this work because the optical path, l, would correspond to the whole width of the flame. This would give some average size, weighted in an unknown manner over a rather large distance in the flame. As already mentioned, it is desired to have as small a scattering volume as possible in order to reduce any size averaging effects due to particle size variation with position in the flame.

None of the mentioned particle size determination methods which are based on light-scattering measurements alone can be applied when the particle size is so small that essentially Rayleigh scattering occurs. This can be seen by rewriting equation 3.1 for N scattering particles which satisfies the condition of $d \ll \lambda$.

$$\frac{I_1}{I_{0,1}} = \frac{9\pi^2}{r^2\lambda^4} \left| \frac{m^2 - 1}{m^2 + 2} \right|^2 (N \times V^2)$$
 6.6

It is seen from this equation which applies when $d \ll \lambda$, that the number of scattering particles N always appears in a product with the square of the volume of the scattering particle V^2 . Since N and V are both unknowns in the experiment and always appear as the product $N \times V^2$ when $d \ll \lambda$, the measurement of an additional quantity, which depends only on N or V or any combination of N and V which is different from $N \times V^2$, is necessary to find either N and/or V. This second measurement could be the refractive index of the composite medium of soot particles and flame gases which depends on the product $N \times V$.

Concentration Determination

The most common way to measure particle concentrations by light scattering after the particle size has been determined is to measure the extinction coefficient at a given wavelength. With a known particle size, wavelength, and value of m, C_{ext} can be calculated from equation 6.5 after which N/V is found directly from equation 6.4. As already indicated, extinction measurements were not used in this work.

The method which was used herein for evaluating the soot particle concentration involved the measurement of the ratio $I_2/I_{0,2}$ at a given angle θ . (The particular choice of using the parallel plane

polarized component is discussed later.) After the size parameter x has been determined by comparing the shape of the experimental and theoretical scattering patterns, the scattering function i_2 is known. Equation 6.2 gives the parallel plane polarized light-scattering intensity for a monodispersed system of N scattering particles. For a given measurement in which θ is fixed, the value of i_2 is known from the previous determination of particle size, the values of k and r are also known, and the ratio $I_2/I_{0,2}$ is measured. Substitution of these values into equation 6.2 gives the number of particles N in the scattering volume. The scattering volume V can be fixed and measured so that the soot particle concentration N/V can be found.

VII. EFFECT OF A DISTRIBUTION OF PARTICLE SIZES ON LIGHT-SCATTERING MEASUREMENTS

Up to this point, the soot particles within a given small volume of a flame have been assumed to be of uniform size so that the scattered-light intensity can be represented by equations 6.1 and 6.2 for N identical scattering particles. It is possible, but difficult, to determine the distribution of sizes from very accurate light-scattering measurements by comparing these measured results with the calculated results made for various assumed size distributions. This was not done in this thesis. This discussion is limited to showing the effect that a distribution of particle sizes may have on light-scattering measurements.

For simplicity, a normal distribution of sizes is used to show this effect. The general procedure is the same for other distributions. The density function as given by Frazer (25) for the normal distribution is

$$f(x) = \frac{1}{\sqrt{2\pi\sigma}} \exp\left[-1/2\left(\frac{x-\mu}{\sigma}\right)^2\right]$$
 7.1

where x is the random variable, μ is the mean about which x is distributed, and σ is the standard deviation. The random variable x in equation 7.1 is taken to be identical to the size parameter x which in turn is equal to $\frac{\pi d}{\lambda}$. It is clear that the density function f(x), defined in this way is dependent on λ as well as on d, so

that a distribution of λ will have an effect which is similar to a distribution of sizes.

The two plane polarized scattered-light intensities for a cloud of particles having a distribution in the size parameter x are given by equations 6.1 and 6.2 if the scattering functions i_1 and i_2 for a single particle are replaced by appropriately averaged scattering functions \bar{i}_1 and \bar{i}_2 , respectively. These weighted scattering functions will depend on μ , σ , m, and θ , whereas the corresponding functions for a single particle or a monodispersed system depend on x, m, and θ . The weighted scattering functions for the two components are defined as

$$\bar{i}_1 = \frac{\int_0^\infty i_1 f(x) dx}{\int_0^\infty f(x) dx}$$
7.2

$$\tilde{i}_2 = \frac{\int_0^\infty i_2 f(x) dx}{\int_0^\infty f(x) dx}$$
7.3

The denominator $\int_0^\infty f(x)dx$ is simply the normalization factor which is constant for a given set of μ and σ and is essentially unity when μ is not very near zero. The exact value of the integral can be found from tables such as given by Pearson and Hartly (26).

The evaluation of the integrals $\int_0^\infty i_1 f(x) dx$ and $\int_0^\infty i_2 f(x) dx$ for various combinations of μ and σ as a function of θ at given values of m was carried out by modifying the Fortram program previously used to calculate i_1 and i_2 . Further discussion of these calculations are presented in appendix E along with a listing of the modified Fortram program and some numerical results for different combinations of μ , σ , and θ .

For all cases considered, the shape of the scattered-light intensity pattern for a given mean size parameter μ and deviation σ corresponds best to a monodispersed scattering pattern for which x is larger than μ . As the standard deviation σ is increased, the difference between x and μ becomes larger. This is so because a given beam of light is scattered more by larger particles as long as the size parameter is not greater than a critical value. Beyond this critical value of x, the intensity oscillates with increasing x. For very small particles such that equations 3.1 and 3.2 apply, it is seen that the scattered-light intensity increases with particle size as d⁶. It follows that the effect of a large number of very small scattering particles is dominated by the effect of a small number of larger particles. Consider a cloud of particles in the size range such that $d \ll \lambda$ in which 10 percent have a diameter d and 90 percent have a diameter of d/4. The light-scattering intensity from the 10 percent which have a diameter d is $\frac{(4)^6}{9}$ or 455 times that scattered from the 90 percent which have a diameter d/4. For

larger particles which are not small compared to λ , the scattered-light intensity does not increase as rapidly with particle size.

Light-scattering apparatus have been construed of determining the size and concentration of suspense a large variety of systems. The basic components scattering apparatus include a light source, a continuident light beam, the particle system to be mation system for the scattered light, a light transcret, and an electronic system for measuring the stransducer. The discussion here will be limited of scattering apparatus which is to be used for determined to the scattering apparatus which is to be used for determined to the scattering apparatus of soot particles in a laminar flame.

Design Considerations

The intensity of the scattered light from a clis many orders of magnitude less than the incident. Therefore, in order to develop scattered light of most it is necessary to use a light source and associated tem which produces and intense beam of incident light be remembered that the light-scattering theory, which is based on the assumption of a parallel beam of incident light into account by making the angle of convergence of the small so that the error due to a nonparallel incident serious. A small collection angle for the scattered

desirable so that the measured scattered-light intensity at a given angular setting is that from only a small angle increment. On the other hand, the angle of convergence of the incident beam and the collection angle of the scattered light must still be large enough to give a measurable signal.

The background light due to the hot soot particles in the flame, whose spectrum is mainly continuous, must also be considered. Since the determination of soot particle size depends on measuring the shape of the scattering diagram or the ratio of the scattered-light intensity at two angles, the signal associated with the background flame light intensity must be essentially eliminated or by some means subtracted from the total signal due to the sum of the scattered light and the background.

As mentioned earlier, a distribution in the wavelength of the incident light will have an effect on the scattered-light intensity which is similar to the effect of a particle size distribution. It is therefore necessary to use a narrow wavelength band in order to be able to interpret the light-scattering measurements.

Since the broad objective of this work is to develop a technique for following the rate of growth and decay of soot particles within a flame, it is desirable to keep the observed scattering volume or the examined flame volume as small as possible. The undesirable size distribution effect due to the variation in particle size with location in the flame will be reduced as the scattering volume is made smaller.

In general, the shorter the wavelength of light used, the greater the scattered-light intensity and the greater the size parameter x for a given particle size. The determination of the size of very small particles which produce nearly a Rayleigh-type scattering pattern at a given wavelength may be made possible by decreasing the wavelength and thereby shifting to a size parameter x to a larger value which has measurable dissymmetry. On the other hand, for ease of alinement of the optical system and checking on equipment operation, it is desirable to keep the light in the visible range. It should be pointed out that, while the use of much shorter wavelengths does increase the value of x, the detection systems required for measuring the scattered-light intensity in this range are limited. The amount of atmospheric absorption for the shorter wavelengths will also be an important factor to consider.

Since the response of the detection system for measuring the scattered-light intensity (i.e., a multiplier phototube for visible illumination) depends on the wavelength, this also must be considered. For illumination with extremely short wavelengths, counting devices would be required.

It follows that the design of the light-scattering apparatus involves a certain degree of optimization, so an effort was made to keep the design flexible enough to permit easy changes after making initial experiments.

Figures 8.1 and 8.2 show schematically the light-scattering apparatus with the various components and associated electronic equipment. Figure 8.1 shows the arrangement for the particle size determination measurements and Figure 8.2 shows the arrangement for measuring the incident light intensity relative to the scattered-light intensity at a given angle which is needed for the concentration measurements. Referring first to Figure 8.1, a portion of the light generated by the light source passes through a small circular aperture which is placed almost on the surface of the source. This light then passes through a series of stops which defines the usable solid angle from the source and eliminates small angle reflections. This system of stops is followed by a short focal length achromatic lens which is located at a distance from the light source equal to its focal length so that it collects the light passing the system of stops and renders it nearly parallel. The light then passes through a stop, the function of which is not to define the light beam but only to reduce reflection from the walls of the tube which contains the lenses and stops. A relatively long focal length lens receives this light which in turn produces a slightly convergent beam which serves as the incident light beam. This lens is located at a distance from the center of the flame which is equal to its focal length. The angle of convergence of the incident light is determined by the light stop located at the long focal length lens.

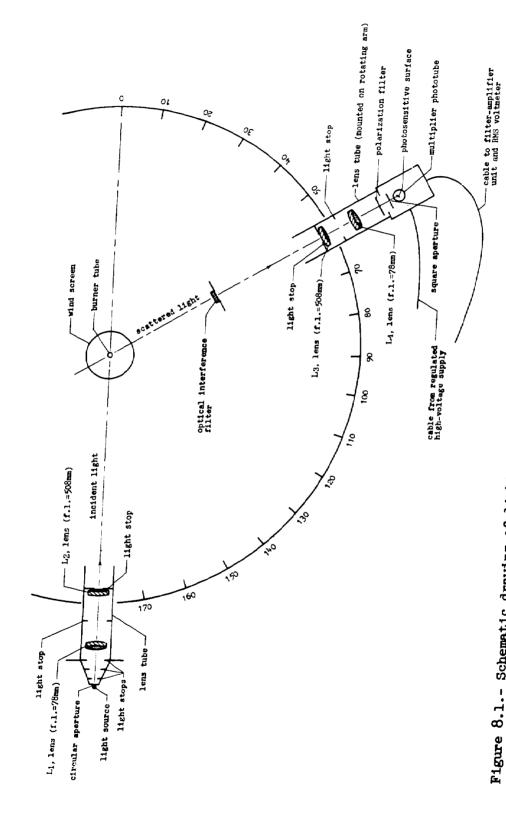


Figure 8.1.- Schematic drawing of light-scattering apparatus for measuring the light-scattering intensity at various angular positions.

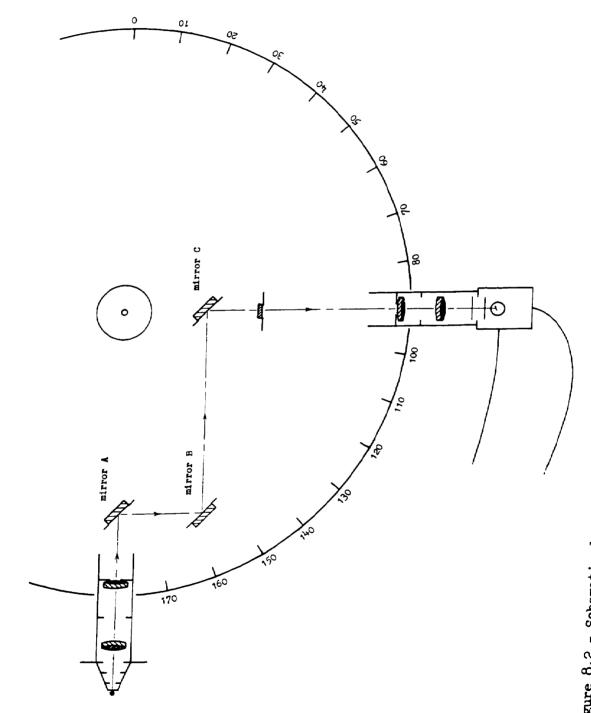


Figure 8.2.- Schematic drawing of light-scattering apparatus for measuring the reflected inci-dent light intensity.

The incident light passes through an opening in a wind screen placed about the flame and generates an approximately cylindrical scattering volume in the flame due to the interaction of this incident light with the soot particles. All of the components, mentioned thus far, of this optical system, are centered on a single axis which intersects and is perpendicular to the axis of the burner tube.

A desired portion of the scattered light is collected by a similar optical system, the components of which are centered on an axis which passes through the intersection of the axes of the burner tube and the optical system for the incident light. Scattered light is generated of course in all directions from the scattering volume, but it is desired to measure the scattered-light intensity at various angles for a small collection angle. The optical system for collecting the scattered light is mounted on an arm which can be rotated about the axis of the burner tube. The angle of observation θ is determined with the help of the graduated circle. With the optical system for collecting the scattered light located at a given angle θ from the direction of the incident light, a portion of the scattered light passes through an optical interference filter which defines the wavelength of the light. The divergence angle or collection angle of this filtered light is determined by the scattering volume and the light stop located at the lens which serves to collect the scattered light. This lens has a relatively long focal length and is located at a distance from the axis of the burner which is equal to its focal length.

This lens produces approximately parallel light which then passes through a light stop to reduce small angle reflections from the walls of the containing tube. This light is focused by a short focal length lens on the plane of a small square aperture. The function of this square aperture is to define the observed scattering volume and eliminate as much of the flame light as possible. An adjustable polarization filter is located between this aperture and short focal length lens so that the intensity of either component of the scattered light can be measured.

A multiplier phototube is located a short distance behind the square aperture and is alined with the optical system in such manner that a spot of light strikes the photosensitive surface of the multiplier phototube. The multiplier phototube in turn generates a signal which is proportional to the intensity of the spot. This signal is passed through a filter-amplifier unit (a General Radio 760-B sound analyzer) which is turned to the frequency of the modulated light source. This electronic arrangement essentially eliminates the steady background signal due to the flame light. The filtered and amplified signal, which is proportional to the intensity of the scattered light, is measured with an RMS voltmeter.

The apparatus for measuring the incident light intensity relative to the scattered-light intensity, which is required for the determination of particle concentration, is shown in Figure 8.2. The apparatus and adjustment of the various components is identical to that used for

the soot particle size determination with a single exception. Instead of the incident light beam passing through the flame, its path is diverted by a series of three black-glass mirrors, after which it is collected by the same optical system, positioned at $\theta = 90^{\circ}$, which was used to measure the scattered-light intensity. The incident light is turned 90° by the first mirror, turned back parallel to its original direction by the second mirror, and turned 90° again by the third mirror so that the attenuated beam is collected and measured by the same optical and electronic systems used to measure the scattered-light intensity. It should be noted that the distance through which the light passes is the same in Figure 8.1 as in Figure 8.2.

A more detailed description of the various apparatus components is presented in appendix F along with the bases used for selecting the various dimensions and locations of the components.

IX. EXPERIMENTAL TECHNIQUE

The method used herein for determining the size of soct particles in various laminar flames entailed measuring the two plane polarized components of the scattered-light intensity for a series of angles in increments of 10° between 30° and 150° from the incident light direction. In order to determine the number concentration of soot particles, the incident light intensity relative to the scatteredlight intensity at a given angle for parallel polarization was measured. The first step toward gathering this experimental data is the alinement of the optical system. The alinement procedure is quite tedious but essentially fixes the axis of the incident light beam and the axis of the collection system for the observed scatteredlight beam in the same plane and further, in such a way that the two intersect at the desired location in the flame and at the center of rotation of the rotating arm on which the collection system is mounted. The procedure for alining the optical system is given in appendix G.

After the apparatus has been optically alined, the mercury arc lamp is started and allowed to operate for a time sufficient to reach steady conditions. At the same time, the test flame is lighted and the desired fuel and airflow rates are adjusted and recorded. Operation of the electronic system, including the high voltage supply to the multiplier phototube, the sound analyzer, and vacuum tube voltmeter, is also begun at this time to insure steady operation during a test.

In all tests, at least 1 hour was allowed for this phase before collecting data.

In order to operate the RCA 931-A multiplier phototube with maximum stability, the anode current should not exceed 10 microamperes (27). It has already been mentioned that a 20 KΩ load resistor is used across the multiplier phototube output so that it follows from the above current limitation that the voltage across this 20 $K\Omega$ load resistor should not exceed 200 millivolts. Since the scattered-light intensity is the largest at angles nearest to $\theta = 0^{\circ}$ and for perpendicular polarization, the rotating arm was positioned at $\theta = 30^{\circ}$ and the polarization filter was set for measuring the perpendicular plane polarized scattered-light intensity while the voltage to the multiplier phototube from the high voltage supply was adjusted in a manner such that the unamplified and unfiltered signal from the multiplier phototube did not exceed 200 millivolts. The voltage setting at the high voltage supply is determined by the combination of the scattered-light intensity and the steady flame light intensity, each of which in turn depends on the size and concentration of soot in the observed volume of the flame. This voltage setting varied from 750 volts to 1250 volts, depending on the flame, but was held constant during any given run.

Following this adjustment, the amplifier of the sound analyzer was set and remained constant during each run. The adjustment and alinement of the entire light-scattering apparatus was unchanged over

the course of a run. Only the angular position of the rotating arm θ and the setting of the polarization filter were changed during a run. The filtered and amplified signal which is directly proportional to the scattered-light intensity was measured with the RMS voltmeter for various values of θ for the two plane-polarized components. This was done by setting the polarization filter to measure the vertical component and measuring the voltage at successive angles beginning at $\theta = 30^{\circ}$ up to $\theta = 150^{\circ}$. Sufficient time at each angular position was allowed for the instruments to reach steady operation. Each reading was checked by repeating this procedure. The parallel plane polarized component of the scattered light was likewise measured at the same angular positions.

The intensity of the incident light relative to the scattered-light intensity at a given angle θ was measured with the same apparatus and alinement directly after completing the scattering measurements at the various angles. A plan view of the apparatus is shown in Figure 8.2. Mirrors A and C are moved to their proper positions as shown in Figure 8.2 which were determined by a previous alinement technique which is discussed in appendix G. The polarization filter is set to observe the parallel plane polarized component after which the rotating arm on which the collection system for the scattered light is mounted is set at $\theta = 90^{\circ}$ so that the incident light is deflected into the collection system by the three parallel mirrors. The same procedure as mentioned already is used to limit

the signal generated by the multiplier phototube to insure its stability. It was often necessary to operate the multiplier phototube at a lower voltage for the incident light measurement than had been used for the corresponding scattering measurements, so it was necessary to measure the scattered-light intensity at one angle for the same supply voltage, amplifier setting, and the same setting of the polarization filter as was used to measure the incident light in order that the two measurements be on the same basis. A scattering angle of $\theta = 60^{\circ}$ was used for this additional measurement.

In order to obtain the actual incident light intensity relative to the scattered light at $\theta = 60^{\circ}$ for parallel polarization $I_{\rm O,2}/I_{\rm 2}(60^{\circ})$, it is necessary to multiply the ratio of the measured reflected incident light to the measured scattered light at $\theta = 60^{\circ}$ by the reciprocal of the total attenuation factor for the three black-glass mirrors and also take into account the geometry of the collection system. The incident light intensity, $I_{\rm O,2}$, is the intensity of light striking the soot particles and has the dimensions of joules/cm². The reflected incident light, however, was measured at the same distance from the scattering soot particles as was the scattered-light intensity. It is therefore necessary to multiply the measured ratio by an additional factor which increases this measured ratio by an amount which would be equivalent to measuring the incident light with this system at the point where the incident light strikes the scattering soot particles. This is more clearly shown with the aid

of Figures 9.1 and 9.2. Figure 9.1 shows the optical arrangement for measuring the intensity of the incident light after it has been attenuated by reflection through the series of three mirrors. The circular cross-sectional area of the focused incident beam striking the scattering particles is denoted by A_0 , the circular cross-section of the light at the plane of the stop S_3 is denoted by A_2^1 , and the circular cross-sectional area of the stop S_3 is denoted by A_3 . It should be noted that A_2^1 is slightly greater than A_3 and was made so in order to be able to visually aline this stop with the incident light.

The distance from the center of rotation of the revolving arm to the lens, L₃, of the collection system is r. The distance r is also equal to $x_1 + x_2$ which is the distance from the focal plane of the incident light where the circular cross section of the beam is A_0 to lens L₃. For the sake of discussion, an optically equivalent system for Figure 9.1 is given in Figure 9.2. Assume that the sensitivity of the collection system is s volts/joule of radiation passing through stop S₃ and let e_0^1 , be the measured voltage when the incident light is attenuated by the three mirrors and $e_2^1(60^\circ)$ be the measured voltage related to the scattered light received at $\theta = 60^\circ$.

It is clear from Figure 9.2 that all of the energy which passes through A_0 also strikes the area A_3^1 so that the ratio of the intensity of the focused incident beam $I_{0,2}$ at p to the light

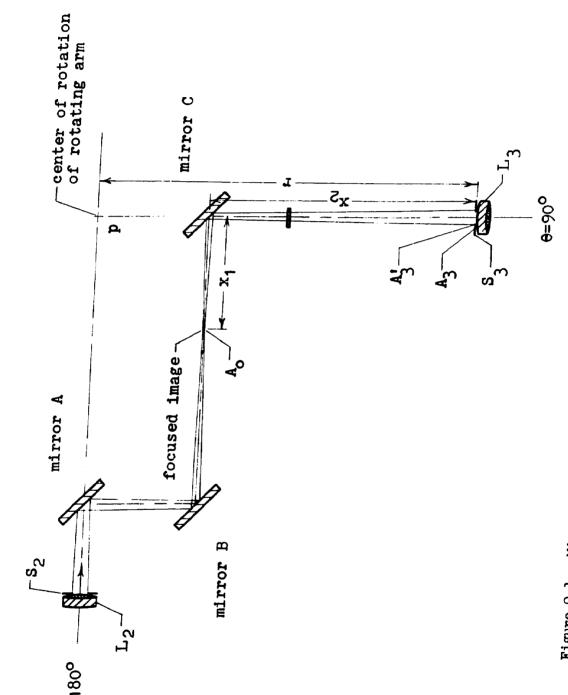


Figure 9.1.- Mirror arrangement for measuring the reflected incident light.

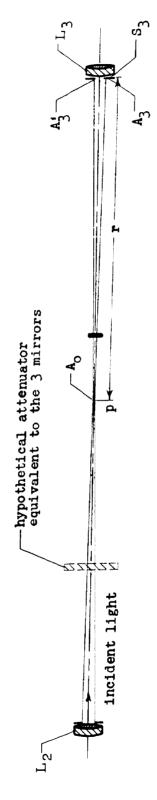


Figure 9.2. Equivalent optical system for measuring the incident light intensity.

intensity I_3^* at the stop S_3 is

$$I_{0,2}/I_3' = A_3'/A_0$$
 9.1

The energy which strikes an imaginary plane placed at S_3 would of course be $I_3^1A_3^1$, but the energy actually received by the collection system is reduced in the ratio, A_3/A_3^1 , since A_3 is smaller than A_3^1 . If we denote the product of the reflectivities of the series of three mirrors as α , the measured voltage for the reflected incident light is

$$e'_{0,2} = \alpha s I_{3}^{i} A_{3}^{i} \left(\frac{A_{3}}{A_{3}^{i}}\right)$$
 9.2

By combining this expression with equation 9.1, the desired incident light intensity is

$$I_{0,2} = \frac{1}{s} \frac{A_3'}{\alpha A_0 A_3} e_{0,2}'$$
 9.3

On the other hand, the measured voltage for the scattered light at $\theta = 60$ is

$$e_2'(60^\circ) = sI_2(60)A_3$$
 9.4

Since the light spot on the photocathode surface of the multiplier phototube is very nearly the same size and is at the same location for the measurement of $e'_{0,2}$ and $e'_{2}(60^{\circ})$, the sensitivity of the

collection system s is the same in equation 9.4 as in equation 9.2. The ratio $I_{0,2}/I_{2}(60^{\circ})$ needed for the concentration measurement is

$$I_{0,2}/I_{2}(60^{\circ}) = \frac{A_{3}^{\prime}}{\alpha A_{0}} \frac{e_{0,2}^{\prime}}{e_{2}^{\prime}(60^{\circ})}$$
 9.5

The diameter of the circular cross section corresponding to A_5^1 was 0.875 in., the diameter of the focused incident light corresponding to A_0 was 0.078 in., and the attenuation factor α was 2.57 \times 10⁻⁶ so that $A_3^1/\alpha A_0 = 4.9 \times 10^7$.

Simple rearrangement of equation 6.2 gives the expression for the number of particles N in the observed scattering volume V as

$$N = \frac{I_2(60^\circ)}{I_{0,2}} \frac{k^2 r^2}{i_2(60^\circ)}$$
 9.6

where the light-scattering function $i_2(60^\circ)$ is known from the particle size determination. (Once the correct value of x has been found from the measured shape of the light-scattering pattern, the value of $i_2(60^\circ)$ for a single particle can be obtained from the tables in appendix D.) The distance from the scattering particles to the point of measurement of the scattered-light intensity r was 508 mm for all runs. Substitution of the value of $k\left(=\frac{2\pi}{\lambda}\right)$ for $\lambda = 4358$ Å, the value for r, and equation 9.5 with the numerical value of A_3/aA_0 into equation 9.6 yields

$$N = 1.10 \times 10^{6} \frac{e'_{2}(60^{\circ})}{e'_{0,2}i_{2}(60^{\circ})}$$
9.7

The observed scattering volume for the system of apertures used and already discussed in the previous section is $\frac{\pi}{4}(0.078)^2(0.094) = 4.49 \times 10^{-4}$ cubic inches or 2.74×10^{-3} cc for $\theta = 90^{\circ}$. For $\theta = 60^{\circ}$, the volume is increased to $2.74 \times 10^{-3}/0.866 = 3.16 \times 10^{-3}$ cc, so that the number concentration of soot particles in particles/cc is

$$(N/V) = 3.48 \times 10^8 \frac{e_2'(60^\circ)}{e_0' \cdot 2^1 2^{(60^\circ)}}$$
 9.8

The essential measurement for determining the number concentration of soot particles N/V after finding the size, is therefore $e_2^*(60^\circ)/e_{0.2}^*$.

Various elevations in the flame above the tip of the burner tube were examined by moving the burner tube up or down while the rest of the apparatus remained fixed.

X. EXPERIMENTAL RESULTS

Preliminary light-scattering tests were made using a rich premixed laminar benzene-air flame by measuring the ratio of the scattered-light intensity at several pairs of symmetric angles about $\theta = 90^{\circ}$ for both the perpendicular and parallel components. An optical interference filter with a peak transmittance at 4358 Å was used. These measured ratios were then compared to the theoretical dissymmetry plots which were previously given in Figures 5.6 through 5.8. The results of this comparison indicated that the size parameter, x, ranged from about 0.9 up to about 1.3, depending on the elevation above the burner tip at which the data were obtained. In general, the greater the elevation, the greater was the value of x. Since λ is approximately 4358 Å, the indicated particle diameter is about 1250 Å to 1800 Å. It should be realized that the correct value of x would be obtained by this method only if the examined soot particle system is essentially monodisperse.

In order to check the values of the soot particle diameters deduced from these preliminary light-scattering measurements, a sample of soot was collected on a probe placed in a benzene-air flame burning under nearly identical conditions and at the same elevation at which a light-scattering measurement indicated a particle diameter of about 1400 Å. This collected sample of soot was then examined under an electron microscope. An electron micrograph of this collected soot is shown in Figure 10.1 where the magnification is $50,000 \times 50$ that

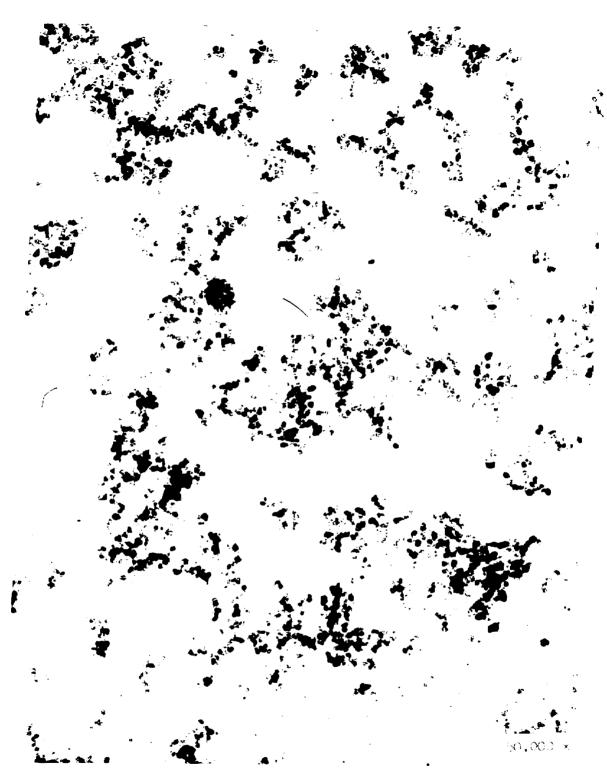


Figure 10.1.- Electron micrograph of soot particles from a premixed laminar benzene-air flame burning with a fuel-air ratio of 2.45 times stoichiometric.

the ultimate soot particle diameter is approximately 250 Å. This figure also shows considerable agglomeration of the ultimate soot particles, but it is not possible to determine how much of this agglomeration took place on the cool probe.

The combination of these preliminary light-scattering measurements and the evidence from the electron micrograph of the collected soot particles indicate that some agglomerates exist in the flame which cause the apparent soot particle size obtained from light-scattering measurements to be greater than the ultimate particle size.

Senftleben and Benedict (11) reported a soot particle diameter of 1750 Å based on light-scattering measurements on an amyl acetate flame burning in a Hefner candle. In order to check their soot particle size determination, a sample of soot was collected on a probe placed in an amyl acetate flame burning in a Hefner candle. An electron micrograph of this soot sample is shown in Figure 10.2 where the magnification is again 50,000 × so that the ultimate particle diameter is approximately 400 Å. As in the previous case, there is considerable agglomeration. Here again it appears that the existence in the flame of agglomerates of the ultimate particles causes the apparent particle size from light-scattering measurements to be much larger than the ultimate size observed in the electron micrograph.

The variation of the intensity of both the perpendicular and parallel plane polarized components of the scattered light with angle from the incident light beam were experimentally determined for



Figure 10.2.- Electron micrograph of soot particles from an amyl acetate diffusion flame burning on a Hefner candle.

several flame conditions. The intensity of the incident light relative to the intensity of the parallel component of the scattered light for $\theta = 60^{\circ}$ was also measured in order to estimate the number concentration of soot particles. The majority of the experimental data were obtained from a laminar premixed benzene-air flame. Some data were also obtained from an amyl acetate diffusion flame and a Cambridge City Gas diffusion flame. These data, along with the test conditions under which they were obtained, are tabulated in appendix H.

Before presenting these light-scattering measurements and their comparison with the theory to determine soot particle size and concentration, it is worthwhile to discuss the effect of the values of x and m on the light-scattering diagram. Figure 10.3 shows the comparison of the scattering diagrams for various values of x for a fixed value of m which is equal to 1.71 - 0.76 i. The general trends shown in this figure for which m = 1.71 - 0.76 i is the same for other values of m. In this figure, the scattered-light intensity for both components have been normalized with respect to the perpendicular component at $\theta = 90^{\circ}$, $i_1(90)$. For very small values of x (which are not shown in this figure) the normalized perpendicular component, $i_1(\theta)/i_1(90)$, would be unity for all angles. For x = 0.80 the normalized values of the perpendicular component are greater than unity for $\theta < 90^{\circ}$ and are less than unity for $\theta > 90^{\circ}$. As the value of x increases, the normalized perpendicular component shows an even greater preponderance of the value of $i_1(\theta)/i_1(90)$ for $\theta < 90^{\circ}$

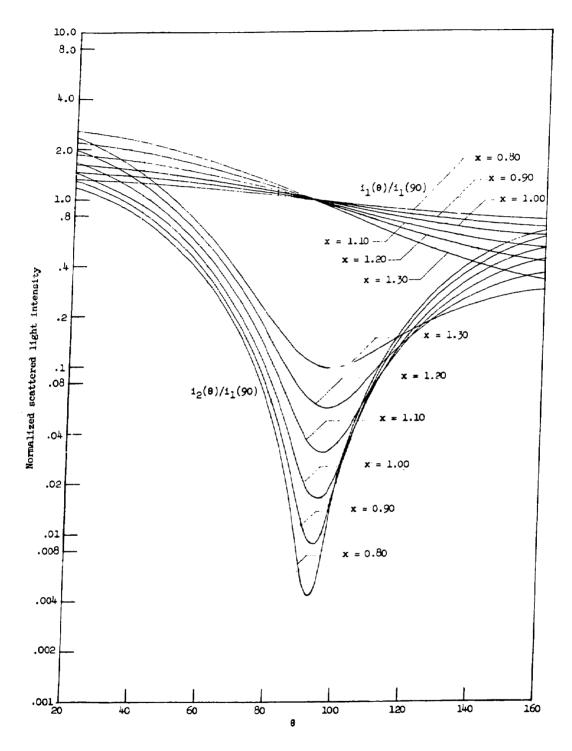


Figure 10.3.- Comparison of the scattering diagrams for various values of x for a given value of m where m = 1.71 - 0.76 i.

compared to values for $\theta > 90^\circ$. The trend of the normalized parallel component, $i_2(\theta)/i_1(90)$, with increasing values of x is also shown. Figure 10.3 shows that the larger the value of x, the shallower is the scattering pattern for the parallel component. In addition, the larger values of x produce a shift in the angular position of the minimum value of $i_2(\theta)/i_1(90)$ to larger values of θ .

The effect which the index of refraction has on the lightscattering pattern is somewhat more difficult to show in that both the real and the imaginary parts of m each have an effect. Figure 10.4 shows the comparison of the scattering diagrams for various values of m where the real part, n, is equal to 1.60 and the imaginary part, n', is taken as 0.60, 0.80, and 1.00 for a fixed value of x = 1.00. The patterns of the perpendicular component for these various values of m differ only slightly and for this work they may be considered essentially the same. The effect on the parallel component is more pronounced, as shown in Figure 10.4, where an increase in the value of the imaginary part shows a shallowing of the scattering pattern with a simultaneous shift of the minimum value of $i_2(\theta)/i_1(90)$ toward $\theta = 90^{\circ}$. It should also be noticed that the parallel component of the forward scattered light is essentially unaffected by this variation of m, whereas there is a noticeable effect on those values where $\theta > 90^{\circ}$. Figure 10.5 is a similar plot in which the real part is equal to 2.00. The trends established in Figure 10.4 are also indicated in this figure.

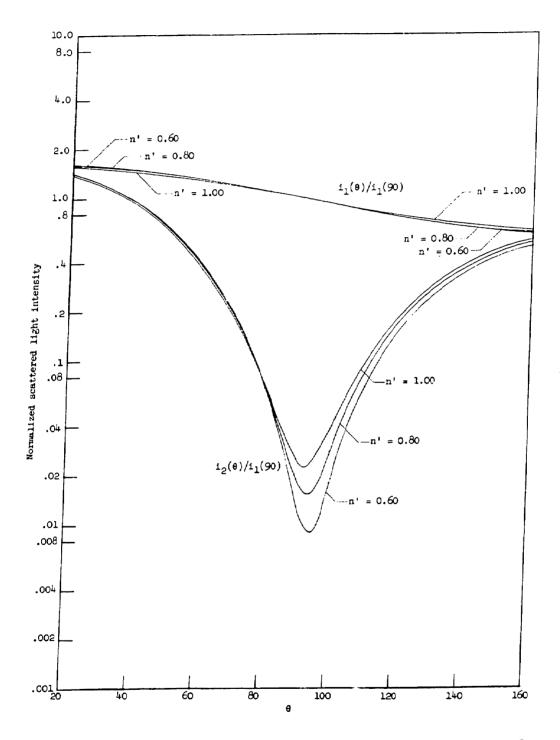


Figure 10.4.- Comparison of the scattering diagrams for various values of m where n=1.60 and n'=0.60, 0.80, and 1.00 for a given value of x where x=1.00.

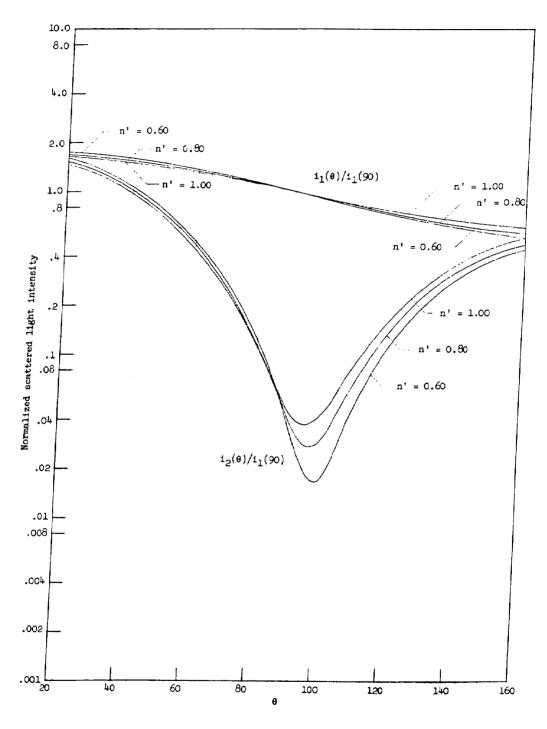


Figure 10.5.- Comparison of the scattering diagrams for various values of m where n=2.00 and n'=0.60, 0.80, and 1.00 for a given value of x where x=1.00.

The effect of varying the real part of m while the imaginary part is held constant at a fixed value of x is shown in Figure 10.6, where the n=1.60, 1.80, and 2.00 for n'=0.80 and x=1.00. Here again the perpendicular component pattern is essentially unaffected. However, the pattern of the parallel component is affected by increasing the value of n in a similar manner in which an increasing value of n affects the parallel component scattering pattern for fixed m.

The values of the real and imaginary parts of m used in the above comparisons cover the range which includes the expected value of m, while the range in x presented in Figure 10.3 corresponds to that in which the experimental data appear to lie.

A series of four runs were carried out experimentally to determine the light-scattering diagram for both the perpendicular and parallel components as a function of elevation above the tip of the burner for a fuel-rich benzene-air flame. The fuel-air ratio, f/a, was held at 2.35 ± 2 percent times stoichiometric for each of the four runs. Run 1 was made at an elevation of 1.25 inches from the tip of the burner on the axis of the flame. The elevation of 1.25 inches corresponds to a position in the flame which is just above the apex of the intercone. (Light-scattering measurements within the intercone and on the axis of the flame indicate the absence of soot in that region of the flame.) Figure 10.7 shows the experimental light-scattering data for run 1. (Refer to table H.1 for the particular test conditions for each run.) The total normalized light-scattering

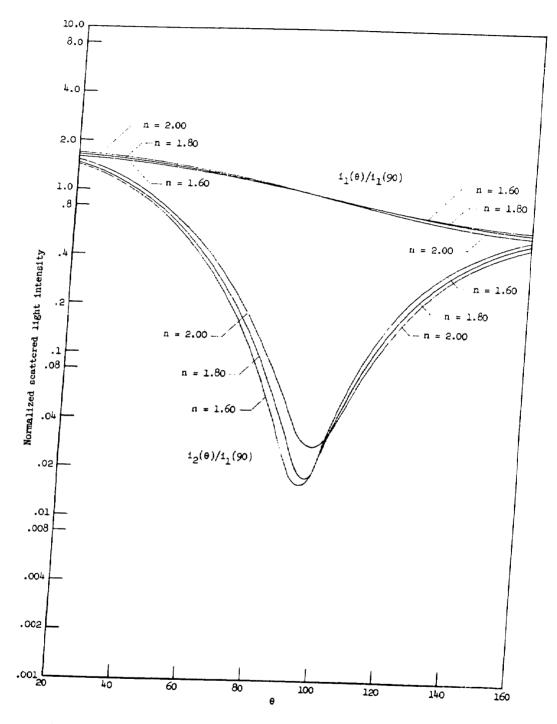


Figure 10.6.- Comparison of the scattering diagrams for various values of m where n' = 0.80 and n = 1.60, 1.80, and 2.00 for a given value of x where x = 1.00.

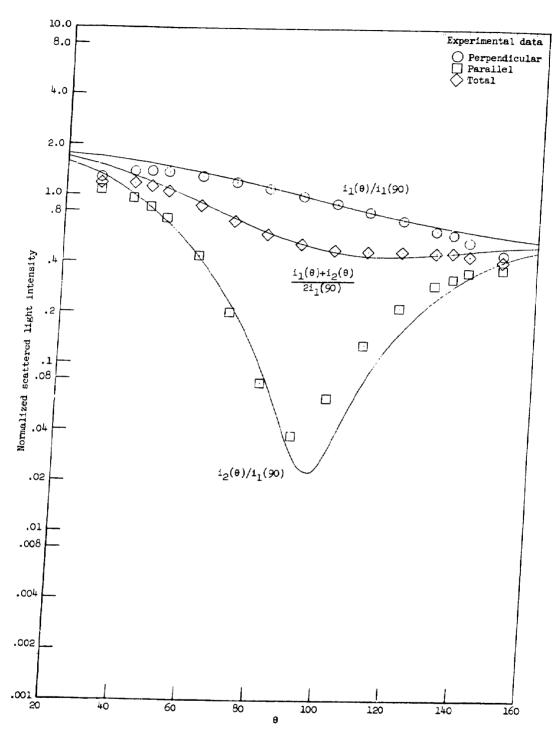


Figure 10.7.- Experimental light-scattering data for run 1 compared to theory for m = 1.71 - 0.76 i and x = 1.05.

pattern, which is simply the arithmatic average of the perpendicular and the parallel components, $\frac{i_1(\theta) + i_2(\theta)}{2i_1(90)}$, is shown along with the other two components.

The data presented in Figure 10.7 for this run and in the remaining figures for all succeeding runs, excluding run 10, have been corrected by multiplying the observed intensities by the sine of the corresponding angle, θ , in order to account for the increase in the observed scattering volume with increasing departure from $\theta = 90^{\circ}$. The greater departure of the data from the theoretical scattering pattern at values of θ near 30° and near 150° is believed to be due to a decrease in soot particle size and/or concentration at the larger distances from the center of the flame. Since the observed scattering volume extends to these larger distances from the center of the flame at the angles in question, there is a decrease in the average intensity for these larger volumes. It should be noted that this explanation requires that the departure at each of a pair of symmetric angles about $\theta = 90^{\circ}$ should be of the same magnitude. This is the case as shown in Figure 10.7. Compare this departure of the data from theory at $\theta = 30^{\circ}$ and at $\theta = 150^{\circ}$.

These data were compared to the normalized theoretical light-scattering pattern for m = 1.71 - 0.76 i and various values of x in order to establish the best fit between experiment and theory for the perpendicular component. For this run, the best fit for the $i_1(\theta)/i_1(90)$ against θ pattern was obtained by setting x = 1.05.

It should be remembered that the theoretical scattering patterns were only obtained for values of x in intervals of 0.05 so that a somewhat better fit may have been found with a slightly different value of x. It is felt that the interval in x of 0.05 is fine enough for this work. After the value of x has been established by a comparison with theory for the perpendicular component, which is x = 1.05 for this run, the corresponding theoretical scattering pattern for the parallel component for x = 1.05 is also superimposed on the plot. The relative position of the theory and the experimental data for the parallel component is fixed once the value of x is established by the perpendicular fit. A rather good agreement between experiment and theory is found for the parallel component where $\theta < 90^{\circ}$, but the comparison is poor for angles of $\theta > 90^{\circ}$. This difference between experiment and theory is much too large to be explained simply as being due to experimental error.

The agreement between experiment and theory for the total light-scattering pattern is about as good as that for the perpendicular component. However, it must be remembered that the total scattering is simply the arithmetic average of $i_1(\theta)$ and $i_2(\theta)$ and that $i_1(\theta)$ is considerably larger than $i_2(\theta)$ in the region of θ where the agreement between experiment and theory is the poorest for the parallel component, so that quite large relative differences for the parallel component would have only slight influence on the total scattering pattern. For this reason, the total scattering pattern is omitted on the remaining plots.

In Figure 10.7 the index of refraction, m = 1.71 - 0.76 i, which corresponds to that given by Stull and Plass for $\lambda = 4358 \text{ Å}$, was used for the comparison. While this value of m for soot is still the most likely value, in view of the lack of the true value of m for soot at the flame conditions for which the scattering measurements were made, several other values of m were also used to see if better agreement between experiment and theory could be obtained. Figure 10.8 shows the same experimental light-scattering data for run 1 compared to theory for various values of m and values of x which correspond to the best fit between experiment and theory for the perpendicular scattering diagram. The four indices used in Figure 10.8, excluding m = 1.71 - 0.76 i, represent a reasonable range in which the true value of m of soot should lie. The limits on the real part of m are 1.6 and 2.0, while the limits on the imaginary part are 0.6 and 1.0. For the values of m considered in Figure 10.8, the best comparison between experiment and theory is obtained for m = 1.60 - 1.00 i with x = 1.05. This shows fair agreement but the difference between experiment and theory is still larger than can be attributed to experimental error.

It is seen by referring back to Figure 10.6, that the parallel light-scattering pattern is shifted toward the left as the real part, n, of m is decreased for a fixed value of the imaginary part, n'. This is the direction in which the theoretical pattern for the parallel component must be shifted in order to produce a better agreement

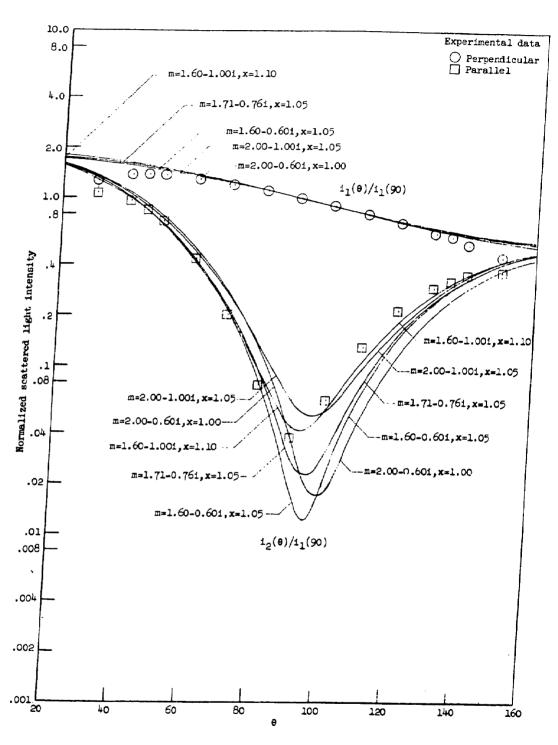


Figure 10.8.- Experimental light-scattering data for run 1 compared to theory for various values of m and x.

between experiment and theory for run 1. In order to see if there is a combination of m and x which gives a theoretical scattering pattern that shows good agreement with this data of run 1, a value of m = 1.40 - 1.00 i was used, based on the trend given in Figure 10.6. Figure 10.9 only shows that it is possible to find good agreement between theory and experiment by allowing m to be a variable and is not intended to show that these are the correct values of m and x.

Three other runs were also made in this series at elevations of 1.50, 1.75, and 2.00 inches above the burner tip. Figures 10.10, 10.11, and 10.12 show the experimental light-scattering data for run 2, run 3, and run 4, respectively. As was the case for run 1, the data for the parallel component do not agree well with theory for the expected range of m. There is, however, an increase in the apparent value of x as the observed elevation is increased, which indicates increasing agglomeration with increasing elevation.

Before the number concentration of soot particles can be determined by the method given earlier, the correct value of x must have been found. For the sake of showing the procedure for determing the number concentration, it is assumed that the value of m is 1.71 - 0.76 i and that only the perpendicular component is to be used to find x. Further, only the experimental and theoretical ratios of $i_1(60)/i_1(120)$ will be used to find x for run 1 through run 4 in order that a number concentration of soot can be calculated. Table 10.1 lists the values of x and corresponding values of an apparent soot

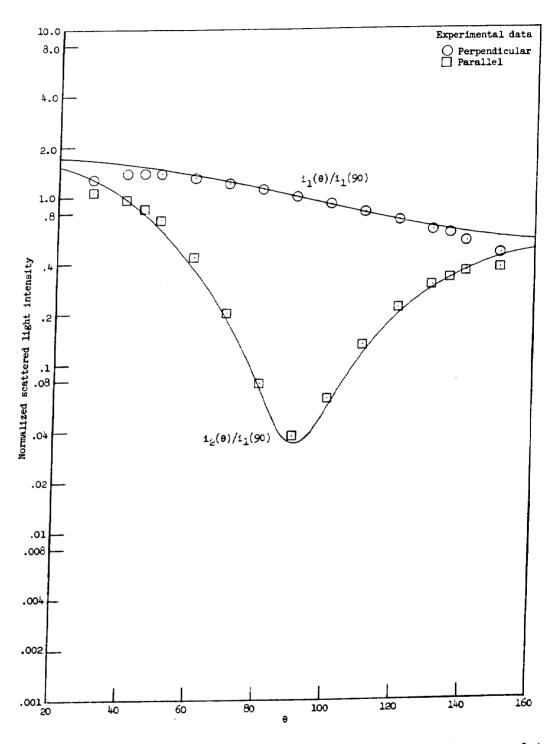


Figure 10.9.- Experimental light-scattering data for run 1 compared to theory for m = 1.40 - 1.00 i and x = 1.10.

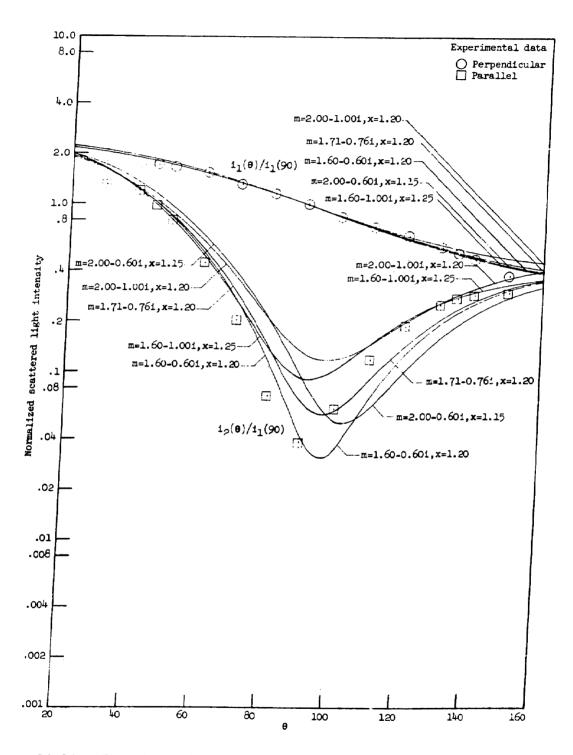


Figure 10.10.- Experimental light-scattering data for run 2 compared to theory for various values of m and x.

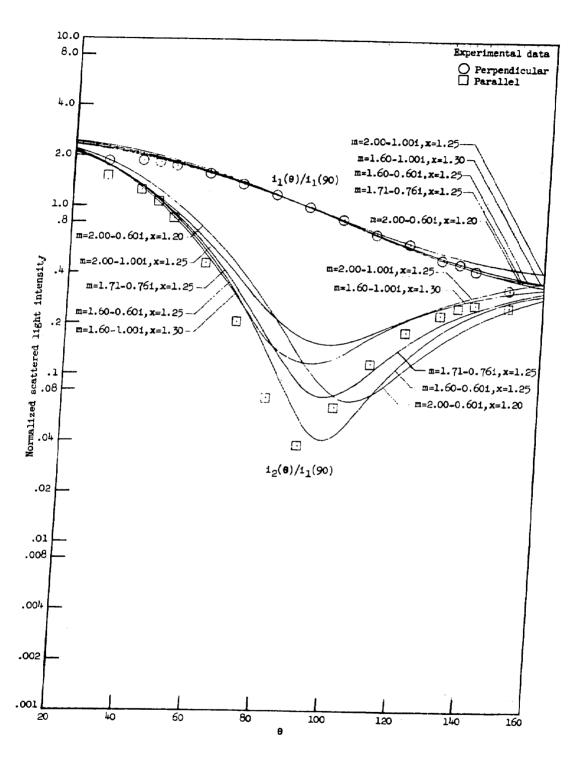


Figure 10.11.- Experimental light-scattering data for run 3 compared to theory for various values of m and x.

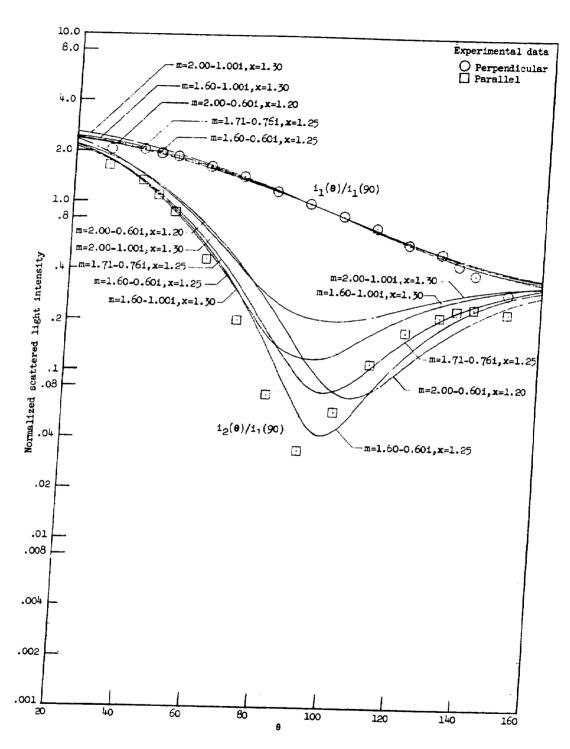


Figure 10.12.- Experimental light-scattering data for run 4 compared to theory for various values of m and x.

particle diameter deduced by comparing the experimental ratio of $e_1(60)/e_1(120)$ to the theoretical dissymmetry plot of $i_1(60)/i_1(120)$ against x, given in Figure 5.8, for run 1 through run 4. Also shown in Table 10.1 is the number concentration for each run based on these apparent values of x and assuming uniform size spherical particles. The number concentration of soot particles was calculated from equation 9.8, in which the value of $i_2(60)$ was found by interpolation from the light-scattering functions given in appendix D.

Table 10.1.- Size and concentration of soot at successive elevations in a fuel-rich benzene-air flame for run 1 through run 4*.

Run	Elevation (in.)	e ₁ (60)/e ₁ (120)	Apparent	x	Apparent diameter, d	Concentration, N/V (particles/cc)
1	1.25	1.80	1.40		1440	2. 7 5 × 10 ⁸
2	1.50	2.36	1.19		1650	2.14 × 10 ⁸
3	1.75	2.54	1.23		1700	1.92 x 10 ⁸
ī	2.00	2.83	1.28		1770	1.42 × 10 ⁸

^{*}These results are based on the assumption that m = 1.71 - 0.76 i, that the soot particles are uniform in size and are homogeneous spheres, and that only dissymmetry between $\theta = 60^{\circ}$ and $\theta = 120^{\circ}$ is needed for particle size determination.

The results presented in Table 10.1 indicate that the apparent soot particle diameter increases and the number concentration decreases with increasing elevation in the flame. The values of x given in

Table 10.1 are in good agreement with the values of x determined by the curve-fitting technique for each run. The absolute values of d and of N/V are of course in question due to reasons previously given.

The size parameter, x, for spherical particles is defined as the ratio of the circumference of the particle to the wavelength of light used in the scattering experiment or $\pi d/\lambda$. It is clear from this definition that the value of x which is determined from light-scattering measurements for a given particle is simply inversely proportional to λ . It follows that two light-scattering experiments in which all test conditions are made identical and only the wavelength λ is changed, would give two different values of x which are in a ratio which is equal to the inverse ratio of the respective values of λ . Two runs with nearly identical test conditions were made with two different wavelengths. A value of $\lambda = 4358$ Å was used for run 5 and a value of $\lambda = 5461$ Å was used for run 6 with both runs having nearly identical test conditions. As before, a premixed benzene-air flame was used with a fuel-air ratio of about 2.45 times stoichiometric at an elevation of 1.25 inches above the tip of the burner.

Figure 10.13 shows the experimental light-scattering data for run 5, in which $\lambda = 4358$ Å, compared to theory for various values of m and x. The value of x obtained by comparing the data with theory for different values of m are given on the figure. For m = 1.71 - 0.76 i, the value of x is approximately 1.00.

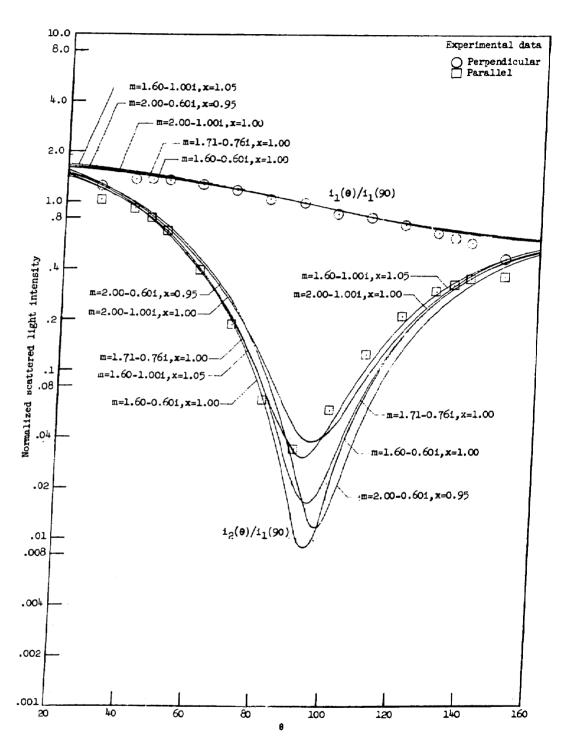


Figure 10.13.- Experimental light-scattering data for run 5 compared to theory for various values of m and x.

Figure 10.14 shows the experimental data for run 6, where $\lambda = 5461$ Å, compared to theory for the various values of m. The values of x obtained from the different values of m are nearly the same. The value of x for this run is 0.90 when m = 1.79 - 0.79 i, the most likely value of the index of refraction for soot when $\lambda = 5461$ Å.

A comparison of these last two runs, which were carried out for the same flame conditions and elevation, shows that the ratio of the values of x for the run where $\lambda = 4358$ Å to that for the run where $\lambda = 5461$ Å is 1.00/0.90 = 1.11 rather than the expected value of 5461/4358 = 1.25. The experimental value of this ratio 1.11 shows the correct trend, in that it is greater than unity, but further explanation is necessary.

It should be remembered at this point that the electron micrograph of the benzene soot shown in Figure 10.1 shows an ultimate particle size of only 250 Å, while the light-scattering results give an apparent size many times larger. It is of interest to assume that the soot within the flame actually is a mixture of small soot particles which have this diameter of 250 Å and much larger particles or elements which are agglomerates of the ultimate particles. It is further assumed for the sake of argument that the light-scattering behavior of the agglomerates can be assigned an equivalent sphere size. The respective values of x for the ultimate soot particles which have a diameter of 250 Å are 0.180 and 0.144 when $\lambda = 4358$ Å and $\lambda = 5461$ Å. The

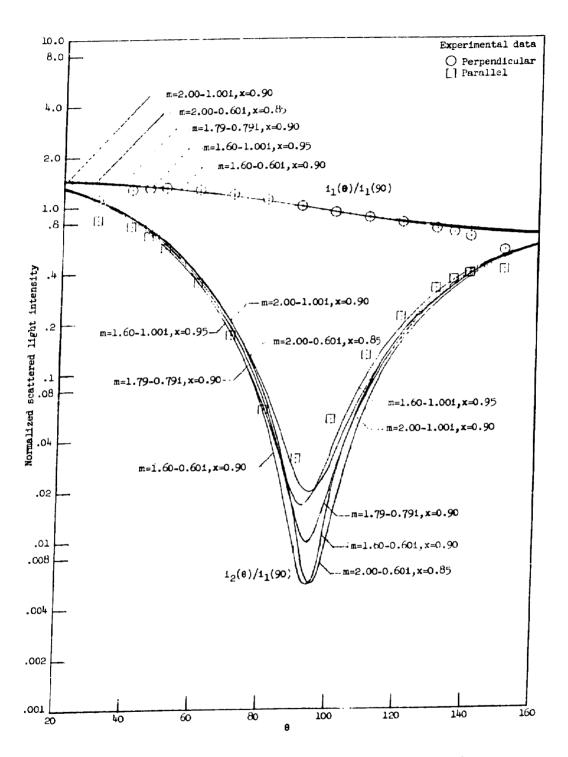


Figure 10.14. - Experimental light-scattering data for run 6 compared to theory for various values of m and x.

closest values of x to these numbers for which the light-scattering functions have been calculated are x = 0.20 and x = 0.15, respectively, and these were used in the following calculation. The effect which these small particles will have on the entire scattering pattern will be such as to reduce the apparent value of x below that which would be obtained from the agglomerates alone. In view of the data obtained from run 5 and run 6, a value of x = 1.25 was assumed for the agglomerates when $\lambda = 4358$ Å. This is equivalent to assuming that the agglomerates are spheres with an effective diameter of 1730 Å, so that the value of x is 1.00 when $\lambda = 5461$ Å. The choice of the effective agglomerate size is arbitrary of course and was selected to simplify calculations, while at the same time supposing a realistic situation based on the data from run 5 and run 6. A value of m = 1.71 - 0.761 was also assumed.

The method of calculation and the results for this hypothetical system which includes a mixture of ultimate particles with a diameter of 250 Å and agglomerates with an effective diameter of 1730 Å are given in appendix I. In these calculations, the number fraction of agglomerates in the total number of particles was taken as 5×10^{-5} . In other words, it is assumed that there are 2×10^{4} individual soot particles with a diameter of 250 Å for each agglomerate with a diameter of 1730 Å. This number fraction was, in part, based on a comparison of the scattered-light intensity from a single particle for the two sizes. Assuming the agglomerates to be of close-packed spheres,

this would correspond to a mass fraction of agglomerates of about 5×10^{-3} . Much smaller values of this number fraction would tend to produce scattering diagrams for each wavelength which would be essentially Rayleigh patterns, but this is not in accord with the data for run 5 and run 6. On the other hand, much larger values of this number fraction of agglomerates would result in a scattering pattern which corresponds to x = 1.25 for $\lambda = 4358$ A and to x = 1.00 for $\lambda = 5461$ A, which is again contrary to the experimental data.

For the assumed conditions, a different scattering diagram was predicted for each wavelength. The size parameter, x, was determined for each of the two diagrams by finding the best fit for the perpendicular component using the theoretical scattering patterns for uniform sizes and m = 1.71 - 0.76 i. In short, the ratio of the resulting values of the apparent values of x for the case where $\lambda = 4358$ Å to the case where $\lambda = 5461$ Å for the hypothetical system of soot particles is approximately 1.1. The above bi-disperse system of soot particles therefore offers a possible explanation for the seemingly peculiar behavior of the experimental light-scattering data exposed by comparing run 5 and run 6. (See appendix I for the details of these calculations and resulting plots.)

The light-scattering data presented up to this point have been obtained from premixed benzene-air flames burning at high fuel-air ratios just short of expelling free soot above the flame. Since the rather large values of x obtained from the foregoing experiments,

relative to the value of x that is expected from the electron micrographs, were tentatively explained as being due to agglomeration, and effort was made to see if a more nearly Rayleigh-type scattering pattern would be obtained from a flame with a lower fuel-air ratio. It is supposed that a smaller amount of soot would result from a lower fuel-air ratio, which in turn may reduce the degree of agglomeration. Run 7 through run 9 were made to this end.

Figure 10.15 through Figure 10.17 present the experimental data compared to theory for various values of m and x for run 7 through run 9, respectively. These runs were made under the same test conditions with a fuel-air ratio of 2.16 times stoichiometric at closer elevation intervals than were used for run 1 through run 4. The values of x for these three runs are slightly less than the values deduced from run 1 through run 4, but it should be pointed out that the flame is shorter in this second series of runs than it was for the previous series. (See Table H.1 for the detail test conditions for these runs.)

As in the previous cases, the experimental data are compared to theory for different values of m with the result that x is about the same for all values of m which were used. The values of m used, and the corresponding values of x resulting from the comparison, are given on each figure. The values of x at successive elevations in the flame for the case of m = 1.71 - 0.76 i are listed in table 10.2. The resulting number concentrations are also presented in table 10.2.

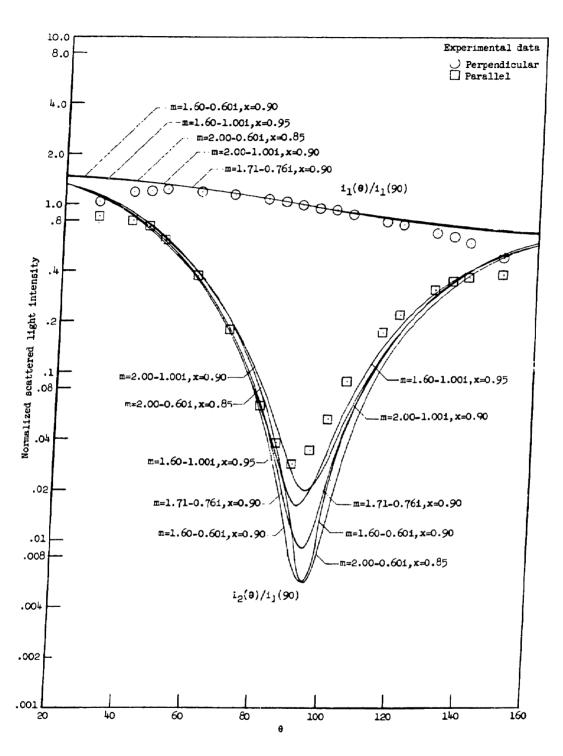


Figure 10.15.- Experimental light-scattering data for run 7 compared to theory for various values of m and x.

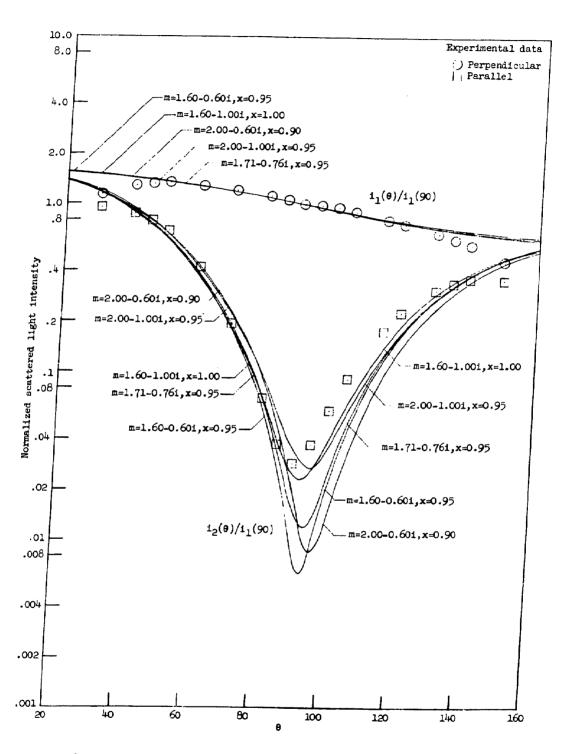


Figure 10.16.- Experimental light-scattering data for run 8 compared to theory for various values of m and x.

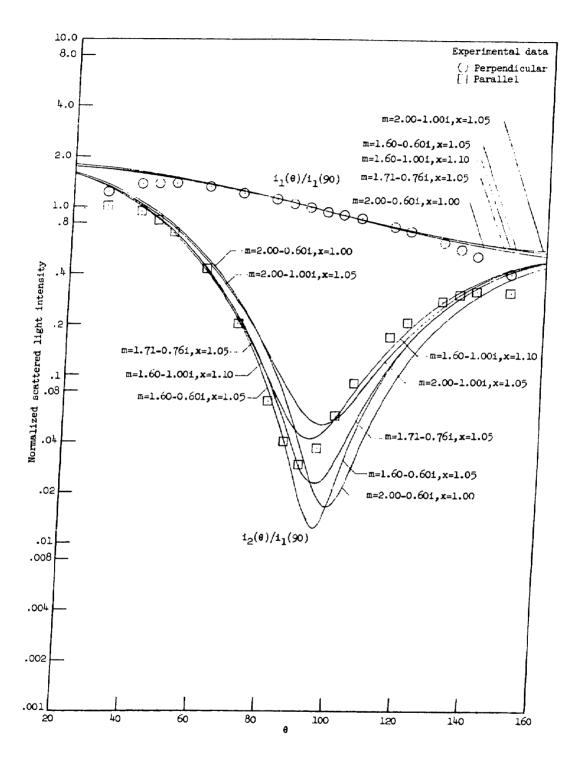


Figure 10.17.- Experimental light-scattering data for run 9 compared to theory for various values of m and x.

Table 10.2.- Apparent value of x and concentration of soot at successive elevations in a fuel-rich flame for run 7 through run 9.

Run	Elevation (in.)	Apparent x	Concentration, N/V particles/cc
7	1.03	0.90	5.04 × 10 ⁸
8	1.10	0.95	4.90 × 10 ⁸
9	1.23	1.05	3.48 × 10 ⁸

The absolute elevations for this series of runs are less than those for run 1 through run 4, but the relative elevation in the flame based on the height of the respective flames are quite comparable. The relative elevation for run 2 is 47.7 percent of the flame height, while the relative elevation for run 8 is 47.8 percent. The apparent values of x for this somewhat leaner, but still very fuel-rich flame, are lower than the values of x found for the more fuel-rich flame at comparable relative elevations. As shown in table 10.2, the number concentration is about twice as large for the leaner flame when compared to the concentrations listed in table 10.1 for the more fuel-rich flame. This may be due to less agglomeration which would allow more separate and smaller soot particles. An attempt to obtain data for an even leaner, but still luminous flame, resulted in a shorter and narrower flame which could not easily be examined with the present apparatus.

Light-scattering measurements were also made on an amyl acetate diffusion flame burning on a Hefner candle with three objectives in mind. The first objective was to gather light-scattering data which could be compared to the data presented by Senftleben and Benedict (11). Secondly, the electron micrograph of amyl acetate soot showed an ultimate particle size of about 400 Å whereas the electron micrograph of benzene soot showed an ultimate particle size of about 250 Å, therefore, it is of interest to see what effect this difference may have on the experimental light-scattering results. Finally, it was desired to see what effect there would be on the experimental data for values of θ near 30° and near 150° (which showed rather definite and systematic departure from theory for the previous runs when corrected by the volume factor, $\sin \theta$) when the observed scattering volume remains constant for all values of θ . The observed scattering volume was made invariant with θ by using a larger aperture in front of the light source which in turn produced an incident light beam with a circular cross section at the flame that was larger than the width of the flame. With this arrangement, the observed scattering volume is defined solely by the width of the flame at the elevation in question and the crosssectional area of the square aperture in front of the multiplier phototube. Both of these are invariant with the angle θ .

The experimental light-scattering data for run 10 compared to theory for various values of m and x are presented in Figure 10.18. The value of x obtained by comparing the data with theory is

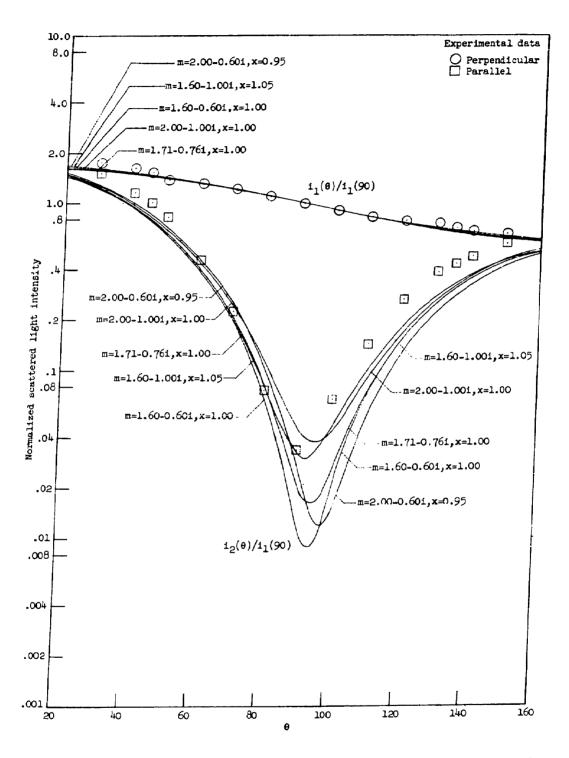


Figure 10.18.- Experimental light-scattering data for run 10 compared to theory for various values of m and x.

approximately 1.00, so that the apparent soot particle diameter is about 1400 Å since $\lambda = 4358$ Å for this run. Senftleben and Benedict reported an apparent soot particle diameter of 1750 Å based on their light-scattering data for which $\lambda = 4910$ Å, that is, their apparent value of x was 1.12. This difference may be due to a possible difference in adjustment of the Hefner candle for the two separate experiments.

The apparent soot particle size deduced from this run is not sufficiently different from those in which benzene was used to be able to indicate an effect of the size of the ultimate particles on the experimental scattering patterns.

As shown in Figure 10.18, the departure of experimental data from theory at angles of θ near 30° and near 150° is absent in this run in which the observed scattering volume was made independent of the angle θ .

One more run was made on a Cambridge City Gas diffusion flame which contains about 94 percent methane. Due to the relatively small amount of soot in this flame, as evidenced by its rather feeble luminosity, it was necessary to operate the multiplier phototube with the maximum allowable supply voltage of 1250v. At this condition, the output voltage dropped quite rapidly with time at a given angular position, due to fatigue of the multiplier phototube, so that a reasonably good experimental scattering pattern was not obtained. The rough data from this run, however, did indicate much smaller dissymmetry than was

produced by the benzene-air and amyl acetate flames. This could be the result of smaller agglomerate sizes in this flame for this run.

XI. DISCUSSION

The experimental light-scattering data indicate that the soot particles in the premixed benzene-air flame range from 1250 Å to 1800 Å in diameter, depending on the elevation above the burner tip at which the data were obtained and the operating conditions of the flame. In general, larger soot particles were found at progressively larger elevations. These particle sizes were deduced by comparing the experimental light-scattering patterns with the theoretical patterns which were calculated for homogeneous carbon spheres of uniform diameter. A rather good fit of the experimental scattering patterns to the theoretical scattering patterns for uniform particle size for the perpendicular component was possible for all runs. The comparison between experiment and theory for the parallel component, however, was rather poor. By proper selection of a value for the index of refraction. a fairly good comparison between experiment and the theory for uniform particle size was found for a given run. The value of the index of refraction which gave this good fit was considerably different from the expected value of the index of refraction found in the literature. There is, of course, some doubt as to the correct value of the index of refraction of the hot soot particles within the flame, but in view of the results shown by the electron micrographs of the collected soot particles, it does not appear correct to assume that the soot particles are monodispersed and allow the index to vary a great deal from the expected value in order to find a fit of experiment to theory.

Electron micrographs of soot particles collected on cool metallic probes, passed through the same flame from which the light-scattering data were obtained, show that the ultimate soot particles from the premixed benzene-air flame are only about 250 A. A possible explanation for this difference may be that the soot within the flame is in an agglomerated state so that the experimental light-scattering data correspond to the relatively large agglomerates of the small ultimate soot particles. The apparent increase in soot particle size with increasing elevation above the burner tip is consistent with the idea of agglomerated soot particles within the flame in that the agglomerated soot particles could be expected to increase in size due to additional agglomeration at progressively higher elevations. Indeed, the electron micrographs show a rather high degree of agglomeration but it is hardly possible to distinguish between the fraction of agglomeration which took place on the cool probe due to its presence in the flame and that fraction which would naturally occur in an undisturbed flame.

The experimental results from a pair of runs on a flame, which had essentially the same test conditions but different wavelengths for each run, show that the ratio of the values of the size parameters for the two runs is not equal to the inverse ratio of the respective wavelengths as would be expected for a monodispersed system. Instead of the value of 1.25, which is the inverse ratio of the two wavelengths used, a ratio of 1.11 was found. In connection with this result,

calculations were made for a supposed soot particle system which was made up of a small number of large particles (corresponding to the agglomerated soot particles) and a very large number of small particles (corresponding to the small ultimate particles). The light-scattering patterns were calculated for the same two wavelengths as used in the experiments based on this supposed particle system. These were then compared to the theoretical scattering patterns for various monodispersed systems to find the best fit for the two cases. The results of these calculations show that the ratio of the apparent size parameters for the two different wavelengths is about 1.06 rather than 1.25 for the particular assumed fraction of large to small particles and their respective absolute sizes. A proper choice of the absolute sizes of the large and small particles and the relative amounts of each would yield a ratio for the apparent size parameters for the two different wavelengths which corresponds to the experimental ratio of 1.11. The point is that the ratio of the values of the size parameters obtained from the same particle system but for two different wavelengths does not necessarily equal the inverse ratio of the respective wavelengths for a particle system which is not monodispersed.

The technique for calculating the light-scattering diagrams for a smooth distribution of particle sizes was presented. A particle system with a normal distribution of sizes was taken and a fixed mean particle size was assumed with several different values for the standard deviation. The results of these calculations show that the apparent particle size increases as the standard deviation increases even though the mean particle size remains fixed. The reason for the increase in the apparent particle size with an increasing value of the standard deviation is that larger particles scatter more light than smaller particles.

For all runs it was seen that the apparent value of the size parameter deduced from a given set of data is only weakly dependent on the particular value of the index of refraction used when the comparison between experiment and theory is made for the perpendicular component of the scattered light. The theoretical scattering diagram for the parallel component is more strongly dependent on the value of the index of refraction.

For all calculations herein, the soot particles have been assumed to be homogeneous spheres. Other particle shapes have not been considered in the analysis and would give different theoretical scattering patterns than presented in this work. This question, however, does not change the rather strong indication that the soot particles, within the flames studied, are agglomerates of smaller spherical units which continue to grow in size by further agglomeration as they travel up through the flame.

The number concentrations of the soot particles which are presented herein are based on the assumption that the observed light-scattering volume is made up of soot particles of uniform size. If the actual soot particle system in the flame contains a very large

fraction of small particles mixed with a rather small fraction of relatively large particles, such as supposed for the example presented in appendix I, the total number concentration of soot particles would be several orders of magnitude greater than indicated and would consist almost completely of the small ultimate particles, even though a sizeable portion of the scattered-light intensity is due to a very small number of larger particles, that is, agglomerates.

If the soot particles in the flame had not been in a. agglomerated state, but instead existed in the flame only as ultimate particles of the size shown in the electron micrographs, the experimental light-scattering diagrams would have been more nearly of the Rayleigh-type showing very little dissymmetry.

In conjunction with this, light-scattering measurements were made on a diffusion flame of Cambridge City Gas (94 percent methane) with the expectation of less agglomeration of the ultimate soot particles, which in turn would give a more nearly Rayleigh-type scattering pattern. The data of this run, using Cambridge City Gas, show much less dissymmetry than given by the benzene-air or amyl acetate flames, but is of a rather poor quality due to apparatus limitations. The quantity and/or condition of the soot within this flame was such that the scattered-light intensity was too low to permit the collection of good data.

An additional quantity which should have been measured, but was not in this work, is the total weight of soot contained in the observed

light-scattering volume of the flame. This would involve collecting a quantity of soot for a known period of time during steady flame conditions. The mass of soot particles contained in the observed light-scattering volume as calculated from the size and concentration of soot particles deduced from light-scattering measurements should agree with the mass of soot determined by directly collecting and weighing it.

Another additional measurement which may help to determine the state of the soot particles in the flame is that of radiant emission from the luminous soot particles in the flame.

The experimental results obtained herein from an amyl acetate diffusion flame burning on a Hefner candle are consistent with the results of Senftleben and Benedict in that good agreement between theory for a uniform particle system and experiment is found for the perpendicular component of the scattered light and the parallel component for $\theta < 90^{\circ}$, but a rather poor comparison exists for the parallel component for $\theta > 90^{\circ}$.

Further development of this technique is of course necessary before its full potential as a research tool can be realized. Recommendations for additional work on this problem are presented in the next section.

It should be pointed out that this light-scattering technique, when developed more fully, has practical application to several other problems in addition to the soot formation problem. For example, the

rate of growth of TiO2 particles, which are used for pigments, might be studied by this technique. As already mentioned, the rate of growth and decay of rocket exhaust particles may also be studied by this light-scattering technique.

XII. CONCLUDING REMARKS AND RECOMMENDATIONS

Concluding Remarks

The main object of this investigation was to develop a light-scattering technique for determining the size and concentration of soot particles in laminar flames. This work is incomplete in that further development of this technique is necessary before its full potential as a research tool can be realized. Based on the results of the present investigation, the following conclusions are drawn:

- 1. The experimental light-scattering results, when analyzed in the light of the electron micrographs of the collected soot particles, indicate that the soot particles within the flames studied are in an agglomerated state. These agglomerated units, which are responsible for the rather large dissymmetry of the experimental light-scattering patterns, are believed to be made up of smaller ultimate spherical particles. The ultimate particle size for benzene soot was about 250 Å, whereas amyl acetate soot was about 400 Å, based on the electron micrographs.
- 2. There is a strong indication that the soot particle system within the examined flames consists mainly of the smaller ultimate soot particles mixed with only a very small fraction of agglomerated units. Based on a comparison between the light-scattering data obtained from the same flame for two different wavelengths and on a comparison of calculated results for a supposed particle system using two different wavelengths, the number fraction of agglomerated units is

several orders of magnitude less than the number fraction of small ultimate soot particles.

- 3. The results indicate that the agglomerated units continue to grow by agglomeration or increase in number as they move up through the flame.
- 4. The value of the apparent size parameters deduced herein are only weakly affected by the value of the index of refraction when the perpendicular component of the scattered-light intensity is used for the analysis. The parallel component, however, is strongly dependent on the value of the index of refraction.
- 5. A light-scattering apparatus has been constructed and used to demonstrate that soot in a flame can be studied by application of the light-scattering technique. Within limits, the supporting analytical work has been developed.
- 6. With additional effort, the light-scattering technique could be applied to several other practical combustion problems.

Recommendations

The following recommendations for future work on this problem are offered:

1. An experiment in which the flame conditions are held constant but in which light-scattering measurements are made at many different wavelengths should be revealing. In conjuction with this, various soot particle systems would be supposed and the resulting calculated light-scattering patterns, based on the same wavelengths for which the

experimental data were taken, would be compared. Electron micrographs might also be taken to establish the size of the ultimate soot particles in this experiment.

- 2. The apparatus should be improved to permit data to be obtained from flames containing smaller amounts of soot than the flames presently investigated. Flames which contain smaller amounts of soot may have less agglomeration of the soot particles or none at all. This would simplify the interpretation of the light-scattering results.
- 3. An attempt should be made to develop a flame that has a more suitable shape for light-scattering measurements than was used herein.
- 4. An investigation in which the weight of soot contained in the observed scattering volume is measured by actually collecting and weighing it, in addition to the light-scattering measurements of size and number concentration, should be undertaken. This would establish a comparison between the soot particle weight in the flame as determined by the apparent size and concentration of soot from the light-scattering measurements and the independent method of collecting and weighing a portion of the soot.
- 5. A theoretical investigation into the effect that particle shape has on the light-scattering patterns is needed. Chains of spherical particles or groups of spherical particles with various packing densities may be realistic cases to consider. The size of these basic spherical particles which make up the bigger units might fruitfully be chosen to be equal to the size of the ultimate particles shown by the electron micrographs.

6. A study with the objective of determining the actual value of the index of refraction of hot soot particles might also be undertaken in support of this work.

APPENDIX A

Light-Scattering Relations for Spheres of Arbitrary Size

The functional relationships for the light-scattering functions,

il and i2, used in equations 3.4 to 3.6 for a sphere of arbitrary

size are

$$i_1 = \left[\sum_{n=1}^{\infty} \frac{2n+1}{n(n+1)} \left(a_n \pi_n(\cos \theta) + b_n \tau_n(\cos \theta)\right)\right]^2$$
 A.1

and

$$i_2 = \left[\sum_{n=1}^{\infty} \frac{2n+1}{n(n+1)} \left\{ b_n \pi_n(\cos \theta) + \mathbf{a}_n \tau_n(\cos \theta) \right\} \right]^2$$
 A.2

In these equations, a_n and b_n are scattering coefficients which are functions of x and y only and are defined as

$$\mathbf{a}_{n} = \frac{\mathbf{S}_{n}'(\mathbf{y})\mathbf{S}_{n}(\mathbf{x}) - \mathbf{m}\mathbf{S}_{n}'(\mathbf{x})\mathbf{S}_{n}(\mathbf{y})}{\mathbf{S}_{n}'(\mathbf{y})\mathbf{\Phi}_{n}(\mathbf{x}) - \mathbf{m}\mathbf{\Phi}_{n}'(\mathbf{x})\mathbf{S}_{n}(\mathbf{y})}$$
 A.3

and

$$b_{n} = \frac{mS'_{n}(y)S_{n}(x) - S'_{n}(x)S_{n}(y)}{mS'_{n}(y)\Phi_{n}(x) - \Phi'_{n}(x)S_{n}(y)}$$
A.4

where

$$S_n(z) = \sqrt{(\pi z/2)} J_{n+1/2}(z)$$
 A.5

$$C_n(z) = (-1)^n \sqrt{(\pi z/2)} J_{-(n+1/2)}(z)$$
 A.6

and

$$\Phi_{n}(z) = S_{n}(z) + iC_{n}(z)$$
 A.7

In these last three equations, the dummy argument z takes on either x or y. The primes denote the partial derivative with respect to the argument, $i \equiv \sqrt{-1}$, and J is the ordinary Bessel function. The angular functions $\pi_n(\cos\theta)$ and $\tau_n(\cos\theta)$ are given as

$$\pi_{n}(\cos \theta) = \frac{1}{\sin \theta} P_{n}^{1}(\cos \theta)$$
 A.8

and

$$\tau_{n}(\cos \theta) = \frac{d}{d\theta} P_{n}^{1}(\cos \theta)$$
 A.9

where $P_n^{l}(\cos \theta)$ is an associated Legendre polynomial.

With this system of equations, it is possible to calculate the necessary light-scattering functions i_1 and i_2 which are required for theoretically determining the scattered-light intensity as a function of x, y, and θ .

APPENDIX B

The Relation Between the Absorption Constant and the Absorption Coefficient

The amplitude of the electric field vector E associated with light traveling through a medium with an index of refraction m in the z direction is

$$E = E_0 \exp \left[i(\omega t - mkz)\right]$$
 B.1

where z is the distance in the direction of propagation away from the point where $E=E_0$ at a given time t. The wave number $k=2\pi/\lambda=\omega c$, where c is the velocity of light in the medium, ω is the angular frequency, and λ is the wavelength in the space surrounding the medium.

Substituting equation 4.1 for m into equation B.1 gives

$$E = E_0 \exp \left[i(\omega t - nkz) - n'kz\right]$$
 B.2

and since the light intensity I is proportional to the square of the amplitude E

$$I = I_0 \exp 2[i(\omega t - nkz) - n'kz]$$
 B.3

where I_O is the light intensity at the point where $E = E_O$.

On the other hand, the usual absorption coefficient γ is defined by the expression

$$I = I_0 \exp(-\gamma z)$$
 B.4

By comparing the real part of the exponent of equation B.3 with that of equation B.4, it is seen that

$$\gamma = 2n'k = \frac{4\pi n'}{\lambda}$$
 B.5

APPENDIX C

Development of a System of Equations for Evaluating
the Light-Scattering Functions and
a Listing of the Fortran Program

The light-scattering functions i₁ and i₂, which are defined by equations A.1 and A.2, are required for the theoretical determination of the scattered-light intensity. Gucker and Cohn (23) have developed a general scheme for evaluating these functions. They start with equations A.3 and A.4 and replace $\Phi_n(x)$ according to equation A.7. This substitution with a little rearrangement yields

$$a_{n} = \left[1 + i \frac{S_{n}'(y)C_{n}(x) - mC_{n}'(x)S_{n}(y)}{S_{n}'(y)S_{n}(x) - mS_{n}'(x)S_{n}(y)}\right]^{-1} = \left[1 + iG_{n}(x,y)\right]^{-1}C.1$$

$$b_{n} = \left[1 + i \frac{mS_{n}'(y)C_{n}(x) - C_{n}'(x)S_{n}(y)}{mS_{n}'(y)S_{n}(x) - S_{n}'(x)S_{n}(y)}\right]^{-1} = \left[1 + iH_{n}(x,y)\right]^{-1} C.2$$

The derivatives of the Bessel functions $S_n'(z)$ and $C_n'(z)$ are eliminated by means of the relation between the derivatives and the functions themselves

$$\frac{S_{n}'(z)}{S_{n}(z)} = \frac{S_{n-1}(z)}{S_{n}(z)} - \frac{n}{z}$$
 C.3

and

$$\frac{C_n'(z)}{C_n(z)} = \frac{C_{n-1}(z)}{C_n(z)} - \frac{n}{z}$$

where z represents either the argument x or y. Substitution of these relations into equation C.2 to eliminate $S_n^i(z)$ and $C_n^i(z)$ yields

$$H_{n}(x,y) = \frac{mR_{n}(y)C_{n}(x) - C_{n-1}(x)}{mR_{n}(y)S_{n}(x) - S_{n-1}(x)}$$
C.5

where

$$R_n(y) \equiv \frac{S_{n-1}(y)}{S_n(y)}$$
 C.6

For the particular case of concern here, where $\,m\,$ and therefore $\,y\,$ is complex, Gucker and Cohn indicate that the simplest way to evaluate $R_n(y)$ is by the recurrence relation for Bessel functions of three successive orders

$$R_n(y) = \left[\frac{2n-1}{y} - R_{n-1}(y)\right]$$
 C.7

For the case n = 0

$$R_0(y) = \cot y$$
 c.8

In a similar manner, the expression for $G_n(x,y)$ in equation C.1 reduces to

$$G_n(x,y) = \frac{W_n(y)C_n(x) - C_{n-1}(x)}{W_n(y)S_n(x) - S_{n-1}(x)}$$
 C.9

where

$$W_n(y) = \frac{1}{m} R_n(y) + \frac{n}{x} \left(1 - \frac{1}{m^2}\right)$$
 C.10

A detailed development of a system of equations for calculating the required scattering functions based on these expressions follows.

The first step, after choosing values for x, y, and θ is to find $R_O(y)$ according to equation C.8. Since y is a complex number, tables of $\cot y$ are generally not available so it is more convenient to express $\cot y$ in exponential form as

$$\cot y = \frac{i(\exp iy + \exp - iy)}{\exp iy - \exp - iy}$$
 C.11

It can be shown by replacing y with $x(N - iN')^*$, that

$$\cot y = \frac{\sin 2xN + i \sinh 2xN'}{\cosh 2xN' - \cos 2xN}$$
 C.12

The real and imaginary parts of $R_{\rm O}(y)$ defined by the following equation

$$R_0(y) = R_0^r + iR_0^i$$
 C.13

are therefore.

$$R_0^{r} = \frac{\sin 2xN}{\cosh 2xN' - \cos 2xN}$$
 C.14

^{*}In this section, the real and imaginary parts of the index of refraction are denoted respectively as N and N' rather than n and n' to avoid confusion with the order n.

and

$$R_0^{i} = \frac{\sinh 2xN'}{\cosh 2xN' - \cos 2xN}$$
 C.15

In these equations and in those to follow, the superscripts r and i denote the real and imaginary parts, respectively.

Before writing equation C.7 in its component parts, it should be noted that

$$1/m = \frac{N + 1N'}{N^2 + (N')^2}$$
 C.16

so that

$$(1/m)^r = \frac{N}{N^2 + (N')^2}$$
 C.17

and

$$(1/m)^{i} = \frac{N'}{N^{2} + (N')^{2}}$$
 C.18

From equation C.7 then

$$1/R_{n}(y) = \frac{2n-1}{x} [(1/m)^{r} + i(1/m)^{i}] (R_{n-1}^{r} + iR_{n-1}^{i})$$
 C.19

so that

$$(1/R_n)^r = \frac{2n-1}{x}(1/m)^r - R_{n-1}^r$$
 C.20

and

$$(1/R_n)^i = \frac{2n-1}{x}(1/m)^i - R_{n-1}^i$$
 C.21

Since

$$R_n(y) = R_n^r + iR_n^i$$
 C.22

it follows that

$$R_n^r = \frac{(1/R_n)^r}{\left[(1/R_n)^{\frac{1}{r}}\right]^2 + \left[(1/R_n)^{\frac{1}{2}}\right]^2}$$
 c.23

and

$$R_{n}^{i} = -\frac{(1/R_{n})^{i}}{\left[(1/R_{n})^{r}\right]^{2} + \left[(1/R_{n})^{i}\right]^{2}}$$
 C.24

Also,

$$mR_n(y) = (N - iN')(R_n^r + iR_n^i)$$
 C.25

so that

$$(mR_n)^r = NR_n^r + N'R_n^i$$
 C.26

and

$$(mR_n)^i = NR_n^i - N'R_n^r$$
 C.27

Now from equation C.5, let

$$mR_n(y)C_n(x) - C_{n-1}(x) = D_n = D_n^r + iD_n^i$$
 C.28

so that

$$(mR_n)^r C_n(x) - C_{n-1}(x) = D_n^r$$
 C.29

and

$$(mR_n)^{1}C_n(x) = D_n^{1} \qquad C.30$$

Likewise, let

$$mR_n(y)S_n(x) - S_{n-1}(x) = E_n = E_n^r + iE_n^i$$
 C.31

so that

$$(mR_n)^r S_n(x) - S_{n-1}(x) = E_n^r$$
 C.32

and

$$(\mathbf{m}\mathbf{R}_{\mathbf{n}})^{\mathbf{1}}\mathbf{S}_{\mathbf{n}}(\mathbf{x}) = \mathbf{E}_{\mathbf{n}}^{\mathbf{1}} \tag{C.33}$$

It follows from equations C.5 and C.28 through C.33 that

$$H_n(x,y) = H_n^r + iH_n^i = \frac{D_n^r + iD_n^i}{E_n^r + iE_n^i}$$
 C.34

so that

$$H_{n}^{r} = \frac{D_{n}^{r}E_{n}^{r} + D_{n}^{i}E_{n}^{i}}{(E_{n}^{r})^{2} + (E_{n}^{i})^{2}}$$
 C.35

and

$$H_{n}^{i} = \frac{D_{n}^{i}E_{n}^{r} - D_{n}^{r}E_{n}^{i}}{(E_{n}^{r})^{2} + (E_{n}^{i})^{2}}$$
 C.36

From equation C.2

$$1/b_n = 1 + iH_n(x,y) = (1 - H_n^i) + H_n^r$$
 C.37

or

$$b_{n} = \frac{(1 - H_{n}^{1}) - iH_{n}^{r}}{(1 - H_{n}^{1})^{2} + (H_{n}^{r})^{2}}$$
 C.38

Since

$$b_n = b_n^r + ib_n^i$$
 C.39

it follows from equation C.38 that

$$b_n^r = \frac{(1 - H_n^i)}{(1 - H_n^i)^2 + (H_n^r)^2}$$
 C.40

and

$$b_{n}^{1} = -\frac{H_{n}^{r}}{(1 - H_{n}^{1})^{2} + (H_{n}^{r})^{2}}$$
 C.41

The scattering coefficient b_n can be evaluated numerically for various values of n by this system of equations for a given set of x, N, and N' if $C_n(x)$ and $S_n(x)$ are known for the appropriate values of n. It should be noted that the terms containing Bessel functions with complex arguments have be eliminated. The numerical values of $C_n(x)$ and $S_n(x)$ as defined by equations A.5 and A.6 can be obtained from tables given in reference 28 with only slight modification for various values of x over a sufficiently large range of n.

A similar expression for an is found by beginning with

$$R_n(y)/m = (R_n/m)^r + i(R_n/m)^i = [(1/m)^r + i(1/m)^i](R_n^r + iR_n^i)$$

C.42

from which it follows that

$$(1/m)^r R_n^r - (1/m)^i R_n^i = (R_n/m)^r$$
 C.43

and

$$(1/m)^{i}R_{n}^{r} + (1/m)^{r}R_{n}^{i} = (R_{n}/m)^{i}$$
 C.44

Also, since

$$1/m^2 = \left[(1/m)^r + i(1/m)^{\frac{1}{2}} \right]^2$$
 c.45

it follows that

$$1/m^{2} = \left[(1/m)^{r} \right]^{2} - \left[(1/m)^{\frac{1}{2}} \right]^{2} + i \left[2(1/m)^{r} (1/m)^{\frac{1}{2}} \right]$$
 c.46

Now let

$$(n/x)(1 - 1/m^2) = F_n = F_n^r + iF_n^i$$
 C.47

Substituting equation C.46 into C.47 yields

$$F_n^r = \frac{n}{x} \left\{ 1 - \left[(1/m)^r \right]^2 + \left[(1/m)^{\frac{1}{2}} \right]^2 \right\}$$
 c.48

and

$$F_n^i = -\frac{2n}{x}(1/m)^r(1/m)^i$$
 C.49

By combining equations C.10, C.42 through C.44, and C.47 through C.49, the relation for $W_{\rm n}(y)$ is

$$W_n(y) = W_n^r + iW_n^i = (R_n/m)^r + F_n^r + i[(R_n/m)^i + F_n^i]$$
 C.50

so that

$$W_n^r = (R_n/m)^r + F_n^r \qquad C.51$$

and

$$W_n^i = (R_n/m)^i + F_n^i$$
 c.52

Now let

$$W_n(y)C_n(x) - C_{n-1}(x) = K_n = K_n^r + iK_n^i$$
 C.53

so that

$$K_n^r = W_n^r C_n(x) - C_{n-1}(x)$$
 C.54

and

$$K_n^i = W_n^i C_n(x)$$
 C.55

Likewise, let

$$W_n(y)S_n(x) - S_{n-1}(x) = I_n = I_n^r + iI_n^i$$
 c.56

so that

$$W_n^r S_n(x) - S_{n-1}(x) = L_n^r$$
 C.57

and

$$W_n^{i}S_n(x) = I_n^{i}$$
 c.58

Now combine equations C.53 through C.58 with equation C.9 to obtain

$$G_n(x,y) = G_n^r + iG_n^i = \frac{K_n^r + iK_n^i}{L_n^r + iL_n^i}$$
 C.59

so that

$$\frac{K_{n}^{r}L_{n}^{r} + K_{n}^{i}L_{n}^{i}}{(L_{n}^{r})^{2} + (L_{n}^{i})^{2}} = G_{n}^{r}$$
 c.60

and

$$\frac{K_n^{i}L_n^{r} - K_n^{r}L_n^{i}}{(L_n^{r})^2 + (L_n^{i})^2} = G_n^{i}$$
 c.61

From equation C.1

$$1/a_n = 1 + i(G_n^r + iG_n^i)$$
 C.62

It follows from

$$a_n = a_n^r + ia_n^i .c.63$$

and from equation C.62, that

$$a_n^r = \frac{(1 - G_n^i)}{(1 - G_n^i)^2 + (G_n^r)^2}$$
 c.64

and

$$a_n^i = -\frac{{G_n}^r}{(1 - {G_n}^i)^2 + ({G_n}^r)^2}$$
 C.65

The scattering coefficient a_n , like b_n , can be evaluated numerically for various values of n for a given set of x, N, and N' if $C_n(x)$ and $S_n(x)$ are known for the appropriate values of n.

The numerical values listed in reference 28 for various values of the real argument x over a large range of n, are the spherical Bessel functions, $\sqrt{\pi/2x} J_{n+1/2}(x)$ and $\sqrt{\pi/2x} J_{-(n+1/2)}(x)$. A comparison of this definition of the tabulated values with equations A.5 and A.6 shows that $S_n(x)$ is obtained by multiplying $\sqrt{\pi/2x} J_{n+1/2}(x)$ by x and that $C_n(x)$ is obtained by multiplying $\sqrt{\pi/2x} J_{-(n+1/2)}(x)$ by x and the term $(-1)^n$. It can be noted in the foregoing system of equations that the term $\sqrt{\pi/2x}$ which appears in the definitions of $S_n(x)$ and $C_n(x)$ appears in equations C.34 and C.59 in both the numerator and denominator in such a way to exactly cancel each other, so that the final result for a_n and b_n is the same whether or not the tabulated spherical Bessel functions are multiplied by x. However, the term $(-1)^n$ in the equation for $C_n(x)$ must be taken into account and has the effect of changing only the sign of the tabulated values for odd values of n. The tabulated values given in reference 28 were therefore used directly as $S_n(x)$ and $C_n(x)$ with the sign of $C_n(x)$ changed for the odd values of n.

The angular scattering functions, $\pi_n(\cos \theta)$ and $\tau_n(\cos \theta)$, which are also needed to evaluate equations A.1 and A.2 for ill and i2, respectively, are tabulated by Gucker and Cohn (23) for values of $\theta = 0^{\circ}(2.5^{\circ})$ 180° for n = 1(1)32. For the present calculations,

 θ takes on all values from 0° to 180° in increments of 5° for each combination of x and m.

It would have been possible to calculate both the Bessel functions $S_n(x)$ and $C_n(x)$ and the angular scattering functions $\pi_n(\cos\theta)$ and $\tau_n(\cos\theta)$ in the course of the following machine program, but it is more efficient from the standpoint of machine time to introduce these quantities as input data to the program.

A machine program to be executed on an IBM 709 computer was developed using the Fortran programing system. This system is discussed in reference 29 and only a listing of the working program will be given here.

For the purpose of writing the Fortran program, the following symbols used in the above equations were redefined by alphabetic groups.

N by U	b _n r by BREAL (N)
N' by V	B _n i by BIMAG (N)
sinh 2xN' by SINH	$(R_{n}/m)^{r}$ by TREAL
cosh 2xN' by COSH	$(R_{\rm D}/m)^{1}$ by TIMAG
Ror by GREAL (1)	Fnr by UREAL
R _O ⁱ by GIMAG (1)	Fn ⁱ by UIMAG
(1/m)r by DREAL	Wnr by VREAL
(1/m) by DIMAG	W _n i by VIMAG
$(1/R_n)^n$ by EREAL	$K_{\mathbf{n}}^{\mathbf{r}}$ by WREAL
$(1/R_n)^1$ by EIMAG	K _n i by WIMAG

```
Rnr by GREAL (N)
                               Lnr by XREAL
 Rn<sup>1</sup> by GIMAG (N)
                               Lni by XIMAG
 (mRn)r by HREAL
                                Gnr by YREAL
 (mR<sub>n</sub>)<sup>i</sup> by HIMAG
                                Gni by YIMAG
 Dnr by QREAL
                               an' by AREAL (N)
Dn by QIMAG
                                an i by AIMAG (N)
Enr by RREAL
                                \pi_n(\cos \theta) by PI (N,M)
Eni by RIMAG
                               T_n(\cos \theta) by TAU (N,M)
Hnr by SREAL
                               x by ALPHA
Hn i by SIMAG
                               θ by ANGLE
```

A listing of the Fortran program which is based on the above system of equations follows:

```
DIMENSION PI(10,75), TAU(10,75), C(10), S(10), GREAL(10), GIMAG(10),
      L(10), P(10), AREAL(10), AIMAG(10), BREAL(10), BIMAG(10),
      ANGLE(75), PRI(75), PLI(75), TLI(75), Z(10)
 10 FORMAT (216, F.2)
 20 FORMAT (6E12.5)
 30 FORMAT (5E14.8)
 40 FORMAT (2F10.3)
50 FORMAT (II)
   READIO, NFINAL, MFINAL, ANGSIZ
   READ20, ((PI(N,M), N=2, NFINAL), M=1, MFINAL),((TAU(N,M), N=2,
     NFINAL), M=1, MFINAL)
   GOTO 900
60 READ30, ALPHA, (C(N), N=1, NFINAL), (S(N), N=1, NFINAL)
           (70, 80, 70, 1000), IND
   GOTO
70 READ40, U, V
80 DCONS=U**2+V**2
  DREAL=U/DCONS
  DIMAG=V/DCONS
  SINH=(EXPF(2.*ALPHA*V)-EXPF(-2.*ALPHA*V))/2.
  COSH=(EXPF(2.*ALPHA*V)+EXPF(-2.*ALPHA*V))/2.
  GCONS=COSH-COSF(2.*ALPHA*U)
  GREAL(1)=SINF(2.*ALPHA*U)/GCONS
```

```
GIMAG(1)=SINH/GCONS
        DO90N=2.NFINAL
        P(N)=N-1
        L(N)=N-1
        EREAL=(2.*P(N)-1.)*DREAL/ALPHA-GREAL(N-1)
        EIMAG=(2.*P(N)-1.)*DIMAG/ALPHA-GIMAG(N-1)
       GCONS=EREAL**2+EIMAG**2
       GREAL(N)=EREAL/GCONS
       GIMAG(N)=-EIMAG/GCONS
       HREAL=U*GREAL(N)+V*GIMAG(N)
       HIMAG=U*GIMAG(N)-V*GREAL(N)
       QREAL=HREAL*C(N)-C(N-1)
       QIMAG=HIMAG*C(N)
       RREAL=HREAL*S(N)-S(N-1)
       RIMAG=HIMAG*S(N)
       SCONS=RREAL**2+RIMAG**2
      SREAL=(QREAL*RREAL+QIMAG*RIMAG)/SCONS
      SIMAG=(QIMAG*RREAL-QREAL*RIMAG)/SCONS
      BCONS=(1.-SIMAG)**2+SREAL**2
      BREAL(N)=(1.-SIMAG)/BCONS
      BIMAG(N)=-SREAL/BCONS
      TREAL=DREAL*GREAL(N)-DIMAG*GIMAG(N)
      TIMAG=DIMAG*GREAL(N)+DREAL*GIMAG(N)
      UREAL=P(N)*(1.-DREAL**2+DIMAG**2)/ALPHA
     UIMAG=-2.*P(N)*DREAL*DIMAG/ALPHA
     VREAL-TREAL+UREAL
     VIMAG=TIMAG+UIMAG
     WREAL=VREAL*C(N)-C(N+1)
     WIMAG=VIMAG*C(N)
     XREAL=VREAL*S(N)-S(N-1)
     XIMAG=VIMAG*S(N)
     YCONS=XREAL**2+XIMAG**2
    YREAL=(WREAL*XREAL+WIMAG*XIMAG)/YCONS
    YIMAG=(WIMAG*XREAL-WREAL*XIMAG)/YCONS
    ACONS=(1.-YIMAG)**2+YREAL**2
    AREAL(N)=(1.-YIMAG)/ACONS
    AIMAG(N)=-YREAL/ACONS
 90 CONTINUE
100 FORMAT(10H1
                     ALPHA=F5.3,6H
                                         U=F5.3,6H
                                                        V=F5.3)
110 FORMAT(72HO
                              A(REAL)
                                                  A(IMAG)
                                                               B(REAL)
            B(IMAG))
120 FORMAT(1H013,4E17.5)
   PRINT100, ALPHA, U, V
   PRINT110
   PRINT 120,(L(N), AREAL(N), AIMAG(N), BREAL(N), BIMAG(N), N=2,
     NFINAL)
   DO14OM=1, MFINAL
   G=M-l
```

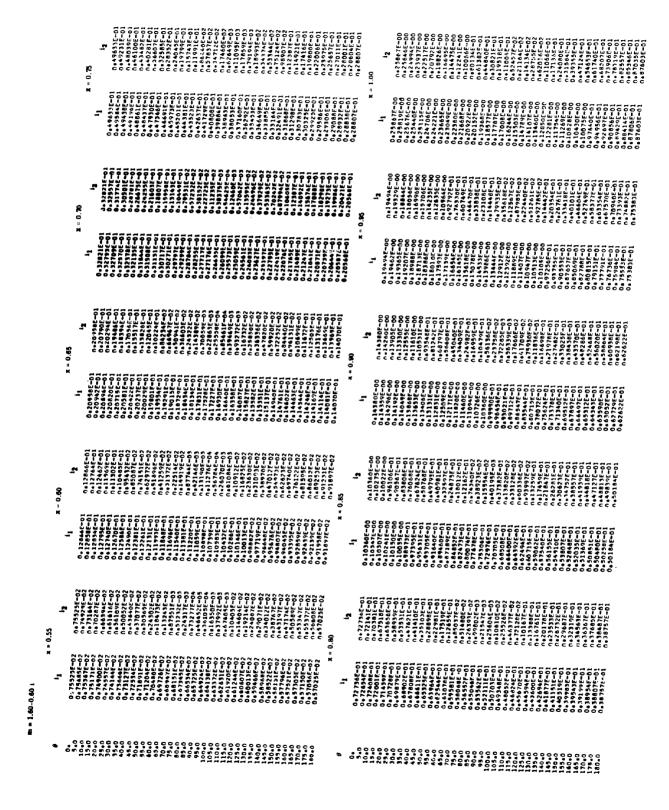
```
ANGLE(M)=ANGSIZ*G
         SARP=0.0
         SART=0.0
         SBRP=0.0
        SBRT=0.0
        SAIP=0.0
        SAIT=0.0
        SBIP=0.0
        SBIT=0.0
        DO130N=2, NFINAL
        P(N)=N-1
        Z(N)=(2.*P(N)+1.)/(P(N)*(P(N)+1.))
        ARP=AREAL(N)*Z(N)*PI(N,M)
        SARP=SARP+ARP
       ART=AREAL(N)*Z(N)*TAU(N,M)
       SART=SART+ART
       BRP=BREAL(N)*Z(N)*PI(N,M)
       SBRP=SBRP+BRP
       BRT=BREAL(N)*Z(N)*TAU(N,M)
       SBRT=SBRT+BRT
       AIP=AIMAG(N)*Z(N)*PI(N,M)
       SAIP=SAIP+AIP
      AIT=AIMAG(N)*Z(N)*TAU(N,M)
      SAIT=SAIT+AIT
      BIP=BIMAG(N)*Z(N)*PI(N,M)
      SBIP=SBIP+BIP
      BIT=BIMAG(N)*Z(N)*TAU(N,M)
      SBIT=SBI+BIT
  130 CONTINUE
      PRI(M)=(SARP+SBRT)**2.+(SAIP+SBIT)**2.
      PLI(M)=(SART+SBRP)**2.+(SAIT+SBIP)**2.
      TLI(M)=(PRI(M)+PLI(M))/2.
 140 CONTINUE
 150 FORMAT(72HO
                      ANGLE
                                      PERPENDICULAR
                                                               PARALLEL
              TOTAL)
 160 FORMAT(1HOF8.1,3E21.5)
     PRINT150
     PRINT160, (ANGLE(M), PRI(M), PLI(M), TLI(M), M=1, MFINAL)
 900 READ50, IND
     GOTO(60,60,70,1000), IND
1000 CALL EXIT
    END
```

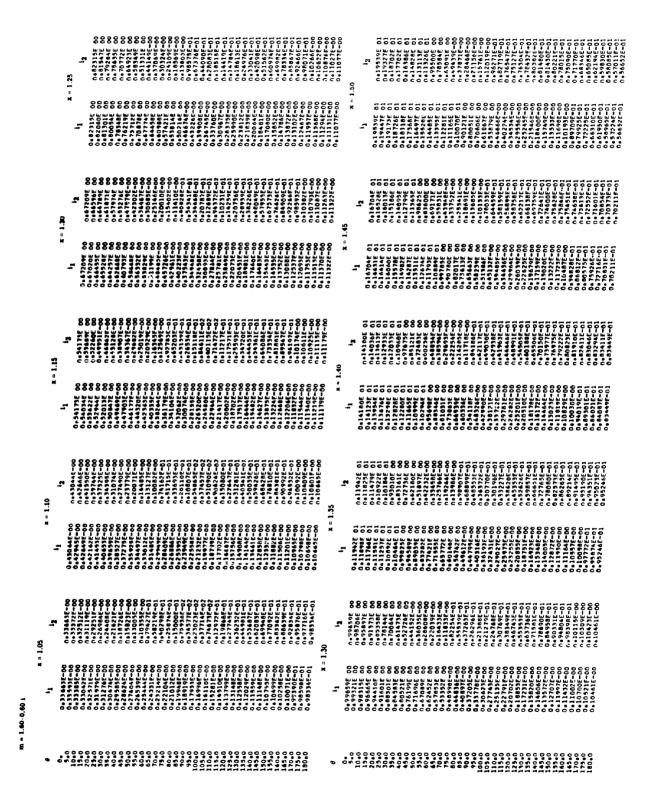
APPENDIX D

Tabulation of the Calculated Light-Scattering Functions and Coefficients for Spheres of Arbitrary Size

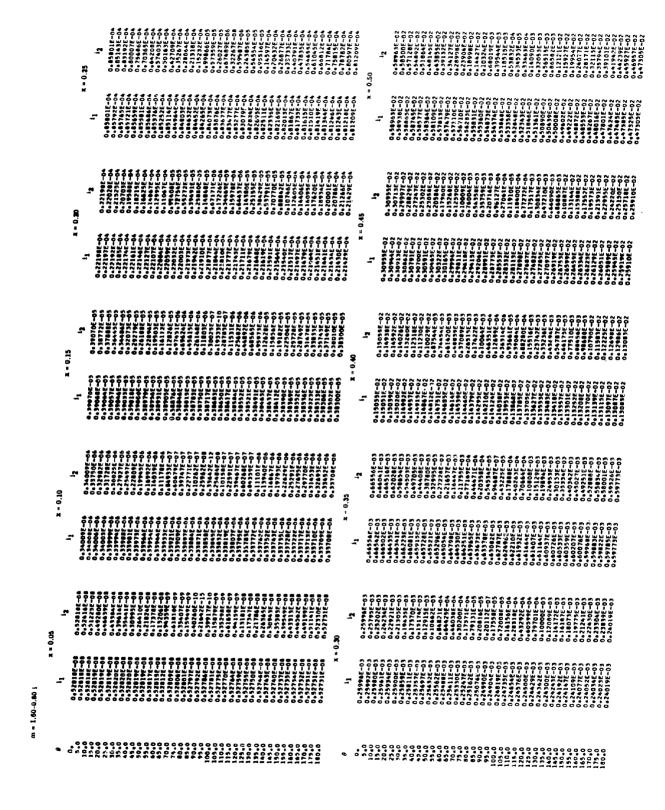
Two sets of tables are presented. The first set lists the calculated light-scattering functions i_1 and i_2 for the size parameter x ranging from 0.05 in increments of 0.05 up to 2.00 for θ ranging from 0° to 180° in increments of 5° for 14 values of the index of refraction m. Following this set of tables, the light-scattering coefficients a_n and b_n are listed for values of n up to 6 for the same range of x and m. All values of i_1 , i_2 , a_n , and b_n are presented to five significant figures. The exponent associated with these tabulated values is denoted by E, the value of which is determined by the sign and two digits following it. For example, the number followed by E-08 is to be multiplied by 10^{-8} .

			= 0.25	ă"	70-314669*0	0.591085-04	0.548466-04	70-196667°U	0.456258-04	0.35649E=04	0.303545-04	0-19952E-04	0.151595-04	0-710745-05	0-3250000	0.493625-06	0.153035-02	0.171585-06	0-388405-05	00-103206-05	0 1 4 4 6 4 E - 0 4	0.238915-04	0.288935-04	0.387215-04	0.43250E-04	0.508625-04	0.537205-04	A-57110E-04	0-5/544E-04	0.50	Ţ,		0.417505-02	0.40772E-02	0.37019E-02	0.343665-02	0-27943E-02	0.207185-02	0.17085E-02	0.103256-02	0.407456-03	0.284306-03	0.38904E-04	0.2176345-05	0-965676-04	0.395245-03	0-60739E-03	0.112095-02	0.14055E-02	0.19864E-02	0.252546-02	0427592E-02	0.31226E-02	0.32420F=02 0.33148E+02	0.33394E~02
			×	:	0.509515-04	0.609345-04	40-368809*0	0+607975=04	0.606456-04	0.605525-04	0-903385-04	0.602185-04	0-559575-04	0.598186-04	0+59528E=04	0-3980E-04	0.590825+04	0.58935E-04	0.58790E=04	0.585135-04	0.583825-04	0.561425-04	0.579355-04	0.578465-04	0 57700F-04	0.5764E=04	0.575695-04	0.575505-04	*0-3***	×	ı,	0+42080F-02	0-420615-02	0.419175-02	0-417926-02	0.414435-02	0.41222E-02	0-406795-02	0-400421-02	0.397465-02	0-35032E-02	0.36660E=02	0.379035-02	0.371475-02	0.36777E-02 0.36415F-01	0-36065E-02	0.357285-02	0.351055-02	0+3+561E-02	0.343245-02	0.33924E-02	0.336326-02	0.3352BE-02	0.33409E-02	0.33394E=02
		97.0	.	04157705-04	0-15659E-04	0-147176-04	0.134255-04	0-11818E-04	0-105695-04	0.786855-05	0.649965-05	0+39312E-05	0-260936-05	0.105725-05	047882E=06	00173956-06	0-112515-06	0-101865-05	0-17829E-05	0.381505-05	0-301998-05	0e76239F=05	A-89436E-05	0-114215-04	0-12503E-04	0.134365-04	0-147505-04	0-15206E-04	à	}	<u>.</u>	0.22037E-02	0=218655=02	0-20525E-02	0.194001-02	0-164215-02	7012794E-02	0.597645-03	0.71376F-03	0.388235-03	0-155735-03	10-360169°U	0.60626E=04	0-12219E-04	0-122195-03	0-216025-03	0046478E-03	0.01147E-03	0.92633F-03	0-12377E-02	0.150905-02	0.16195E-02	nel7746E-02	0.181485-02	1
		Ä	ī,	0-157795-04	0-15778F-04	0-157696-04	0-157526-04	0.157405-04	04197115-04	0-156945-04	0.156565-04	0-156345-04	0-155895-04	0-195656-04	0.195166-04	0.154916-34	0.134665-04	0.15417E-04	0.153935-04	0-153485-04	0.153276-04	0-152896-04	0415273E-04	0-152445-04	0.152315-04	0.152166-04	0.152115-04	0-15206E-04	**	-	. ::	04220376-02	0.220075-02	0.219605-02	00218486-02	0.216735-02	0.21567F-02	0.21322E-02	0.210425-02	0.20892E-02	0-205768-02	0.204135-02	0-20086E-02	0.19762E-02	00196055-02	0-19306F-02	0e19166E-02 0e19034F-02	0-189105-02	0 1069ZF-02	0-185985-02	0-18446F-02	0.183425-02	0.18309E-02	0.182636-02	
	;	6.15	. .	0.277896-05	0.26948E-05	0.245295-05	0-226134-05	3+14627E-05	0.130725-05	0-114626-05	0.912485-06	0-499426-06	0+32465E-06	0-840855-07	0.21431E-07	0.20365-11	A-81791E-07	0-182316-06	0-486996-06	0.661755-06	0-112656-05	C-136295-05	0.18281E-05	0.223676-05	0.240-1E-05	0.25398E-05	0-27009E-05		9	4	0.10702E-02	9-10-13E-02	0.997175-02	0.94274E-03	0.178505-03	0.71322E-03	De52968E-03	0.347635-03	0.26417E-03	0.12432E-03	0.33236-04	0-919676-05	0.633255-05	0-271585-04	0-10893E-03	0.16699E-03	0-307865-03	0.46669E-03	0.524775-03	0.69642E-03	0.0161546-03	0.86296E-03	0.916715-03	0.92358E-03	
	- 1	*	-	0.277885-05	0-277795-05	0.277716-35	0-27750E-05	0.277366-05	0.277046-05	0.27686E-05	3-276455-05	0.27622E-05	0-27575E-05	G+27551E-05	0+27501E=05	0-27476E-05	0.274.275-05	0+27403E=05	0=27380E=05	0-27337E-05	0.272995-05	0-27282E-05	0.27267E=05 0.27254F=05	0.27242E-05	0.272335-05	0+27220E-05	0-27217E-05 0-27216F-05		H H	.	0-10702E-02	0-106915-02	0-10676E-02	0-106555-02	0+105986-02	0-105215-02	0.104765-02	0-10374E-02	0+10318E-02	0+101995-02	0-100745-02	0.100105-02	0.988235-03	0.975795-03	0.96981 E-03	0.95853E-03	0.953325-03	0.94392E-03	0.936105-03	0+932846-03	0.92774E-03	0.925936-03	0+92385E-03	0.763381-03	
	ð.10	5	0.242015-06	0-240175-06	7-22578E-06	0-193745-06	0-101456-06	0.141946-06	70120935-05	0-75553E-07	0.50450E-07	A-2829-E-07	7-142135-07	0-3-0cs1-0	21-197/201-0	74 72 19 35 -03	-16052c-07	70-1025E-07	10-30266500	0e76954E-37	9-11997E-96	0-14375E-06	A-179926-05	0-21177E-06	n.2237-E-06	7.23796F-06	n.23978E-06	0.35	2	•	0447280E-03	0045831E-03	00410065-03	0-34248-03	0.353176-03	0.27561E-03	0-16261-03	0-13402F-03	0.83688E-04	0-317356-04	0.145678-04	0-479756-07	0.24642E-05	0428271E-04	76159E-04	0-106695-03	0.176145-03	0.21291E-03	0.28495E-03	0.347956-03	0e37367E-03	7.40982E-03	04419206-03		
	*	=	0-242615-06	0-241396-06	0.241975-06	3-241916-06	0.241815-06	3-24175E-06	3.24162E-C6	0-24154E+06 0-24144F+04	3-241376-06	0.241285-06	0-24103=-06	0.241005-05	C+24030E=34	0.242726-05	3-543526-06	3.243435-36	3+2+03+1-05	3-243135-36	0.24.011.c-06 0.24.0047.e06	20-36641300	0.233895-06	3-237456-36	0-239505-06	0.27775E-06	90	=		0.477405.62	0-472705-03	0-472405-03	0.47121E-03	0.469375-03	0-46804E-03	0-465216-03	0-463436-03	0-459735-03	0-455665-03	0-453548-03	0449176-03	0.446975-03	0-442615-03	0.436426-03	0.436425-03	0.432715-03	0.431016-03	0-428016-03	04425725-03	0.424616-03	0-42317E-03	0.422726-03	0.42236E-03		
	8 0		0.37616±-03	7-364825-03	7+33215E-08	7-30677E-08 0-24212E-08	9-252495-06	7-18808E-08	0+10042E-04	7.940485-09	9-671975-09	7-25215E-09	0-113545-09	7-179575-10	10234135-10	7.251295-09	7-63467-03	3+670355-09	0.12350E-08	7*155105*04	7- 42023E-04	0-251505-08	0-36635-0	7.331485-08	7-36407E-08	7-372545-09		•	<u>.</u>	1.18465E-03	00183235-03	1-17214E-03	0-15282E-03	7-136075-03	0.123416-03	0-917795-04	0-200 09 eu	0.327545-04	0+21901E-04	0.565205-05	0-14949E-05	0-12175E-05	0.13805-05	\$0-362661°0	0-47817E-04	0.56335E=04	0.65487E-04	0.100236-03	0-12785E-03	0-150285-03	0.158695-03	0.16867E-03	0+16995E-03		
1 09:0	å 	0.374.45	0.37616E-08	0-376165-09	0.376155-08	0-376155-08	0+376135-08	3.376125-06	0.376085-08	0-376065-08	0-376005-0	04375965-08	0-375-06-08	0-375852-08	0.17576F=0a	0-37572E-06	0.375635-08	0+37560E-08	0+375575-08	0.37551E-06	0.375485-06	0-375445-08	0-375426-08	0-375405-09	0.375405-08	0.375395-08		-		0+184656-03	0-184535-03	0.184376-03	0-18393E-03	0.183435-03	0-162875-03	0.181945-03	0-181425-09	0-10029E-03	0.179695-03	0=178456-03	0-177175-03	0e176536-03	0-175275-03	0-174675-03	0-17352E-03	0417250E-03	0.172036-03	0-171236-03	0.170516-03	0.17037E-03	0-17005E-03	0.169976-03	50-30-00-0		
3.1.e	•	•	0.0	0.00	25.0	35.0	0	30.0	0.60	95.0	70.0	0.00	85.0	95.0	0.001	105.0	115.0	120.0	130.0	135.0	145.0	150.0	150.0	165.0	175.0	180,0		•	4	0 n	000	2040	25.0	39.0	0.04	000	60.00	6.0° 0.0°	75.0	000	0	100.0	105.0	115.0	129.0	130.0	140.0	145.0	199.0	165.0	170.0	180.0			

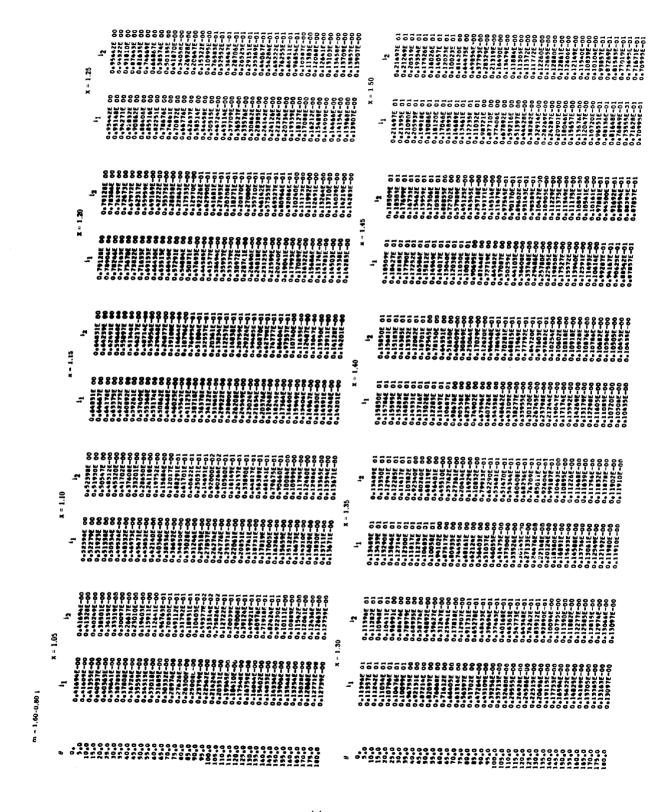


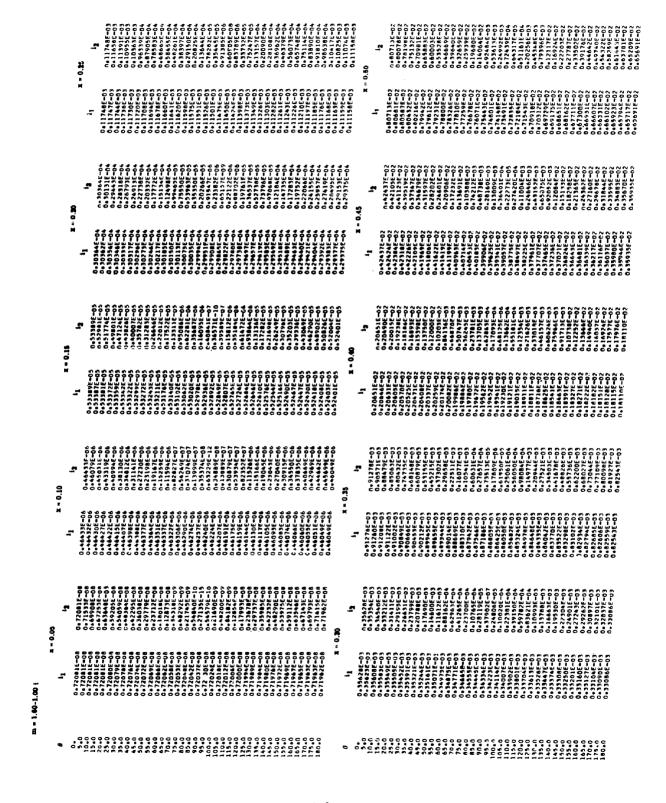


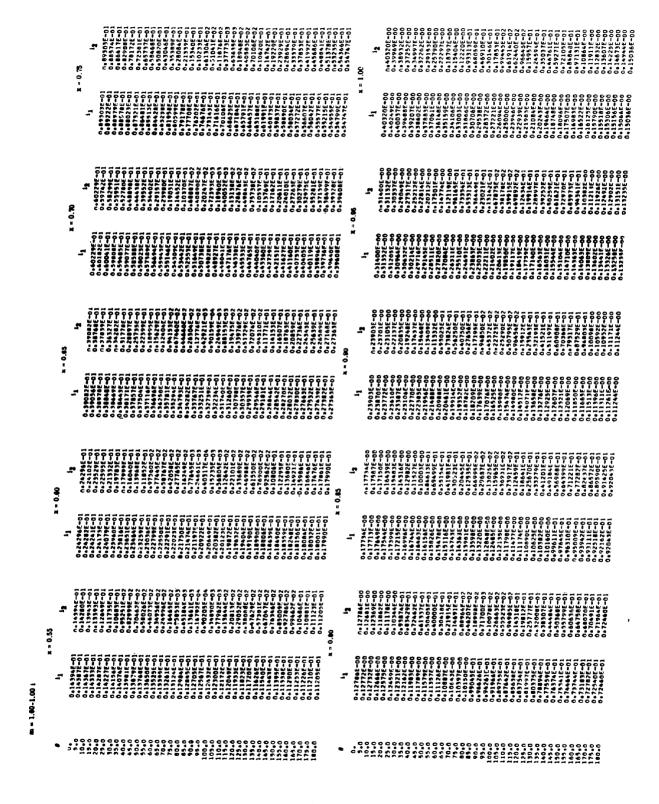
	1.78	_		0.39036E 01	0.37910E 01	0.31955£ 01	0.28308E 01	0.20429E 01	0+16628E 01	0413166E 01	0. 76TOTE 00	00 3511950	0.326115-00	0.26512E-00	0.22965E-00	00412835-00	0-20326E-00	0.203495-00	00198675-00	0.17587E-00	0.15862E-00	00139095-00	0+98800E-01	0.80638E C1	0.549465-01	0421456-01	0.34869E-01	0-298845-01	0.250375-01	0.24501E-01		7 .00	77	Treat.	0.70885E 01	0.67603E 01	0.555465	0.48315E 01	0.40408E 01	0.0256785 01	0.19221E 01	0.14067E 01	0.718955 00	0.52692E 00	0.353025-00	0.329505-00	0.32334F=00	0.31862E-00	04305365-00	0.25185E-00	0.216215-00	04144946-00	0+11976E-00	0.797395-01	0.734826-01	0.799705-01	0.08002E-01	0.973667-01	0.11286+00	20201110
	×		0.194475 01	0.39290E 01	0.36501E 01	0.3550ZE 01	0.33405E 01	0.20401E 01	0.228715 01	0.20110E 01	0.17443E 01	0.12596F 01	0.1040ZE 01	00 380000	0.444315 00	0.43804E-00	0.340395-00	00-10076-00	0.148615-00	0+111166-00	0.630416-01	0-46062E-01	0-36180E-01	0.317326-01	0=236475-01	0-2447F-01	0.2411E-51	0.24.1255-01	0.244405-01	0.24501E-01	3		-	C. 72031F 01	0.71358E 01	0.694338 01	0.622185 01	0.57279E 01	0458615 01	0.39854E 01	0.33944E 01	04231135 01	O.18443E 01	0.108115 01	0. 80630E 00	0457865E 00	0+271505-00	04176115-00	0.762155-01	0.59016E-01	0.466496-01	0,52028E-01	0.60250E-01	0.78826E-01	0.873246-01	0.10070E-01	0.10546E-00	0.111426-00	0-11282E-00	
	1.70	<u>.</u>	0+34698E 01	0.34218E 01	0.30806.01	062824E 01	0.21769E 01	001835E 01	0.120346 01	0.93732E 00	00936345 00	00402235-00	0-304958-00	00-375-00	0+16762E-00	0+17941E-00	00173316-00	0.17229E-00	0-166136-00	04198365-00	0.12796E-00	0-111335-00	0.784976-01	0.64436E-01	0.52261E-01	0-42369E-01	0-34736E-01	0-294846-01	0-23370E-01	0.420BbE-01	8:1	د	<u>.</u>	0464425E 01	10 38488 01	0456124F 01	0+50371E 01	0.437858 01	0.299790	0.23595E 01	0417975E 01	0.96277E 00	0.69327E 00	00-39779E-00	04337736-00	0.30562E-00	00-369696-00	0.303235-00	0.27621E-00	0.250496-00	04185575-00	0+152695-00	0.99316F-01	D-82263E-01	0.671752E-01	0.677616-01	0.714295-01	0-81628E-01	0-89225E-01 0-86523E-01	
		-	0.3465BF 01	0.337926 01	0.92741E 01	0.29594E 01	0-27609E 01	04231435 01	0.20796F 01	0.161226 01	0+14004E 01	0-11979E 01	0 84480F 00	0-696856 00	0.368115 00	00.363086.00	0.28812E-00	0.22932E-00	0-17-005-00	G+10412E-00	(1+80516E=01	0-48956-01	0.405746+91	0.34659E-01	0.296316-01	0-26/125-01	0-23739E-01	0.23162E-01	0-227825-01	10-3001	=	4	-	0.64429E 01	0.62240E 01	0.59611E 01	0.51896F 01	0+47148E 01	0.42079E 01	0431707F 01	0a26751F 01	0.22126E 01	0.14217E 01	0. 11022E 01	0461964E 00	0.44230E-00	0-304116-00	0.14016E-00	0.92885E-01	0-471346-01	0+403456-01	0-34046-0	0.492705-01	0.55445E-01	C. 68900E-01	0.74309E-01	0.82192E-01	0.046146-01	0.86523E-01	
8	9		0.30241E 01	0.28790E 01	0.27069E 01	0.22191E 01	0019335 01	0.13572E 02	0.10935E 01	00.0000000	0.499176-00	0028416-00	0.222695-00	00184326-00	00-10070100	0.148765-00	0.14741E-00	00142775-00	0.13596E-00	0.12696E-00	0-119775-00	0.69696E-01	0.76909E-01	0.648398-61	0-45176-0	0-379786-01	0.325576-01	0.268116-01	0.25698E-01	8	Ŗ	12	-	0.565485 01	0.540938 01	0.50223E 01	0.39510E 01	0.334396 01	0.21758F 01	0.16739E 01	0412517E 01	0465856	0+492245-00	04382475-00	0-293425-00	00-284825-00	0-28332E-00	0.277185-00	04242645-00	0-21594E-00	04185995-00	0-12715E-00	0+10272E-00	0.69839[-0]	0-615395-01	0.57279E-01	0.587705-01	0.60930E-01 0.62695E-01	0-633645-01	
- I			0.300636 01	0.29596 01	0.27919£ 01	0.24444	0.22470E 01	0.20764E U1	0.16035E 01	0.149046 01	0.113015 01	00.000000	0-82080E 00	0.472245 00	0-470995-00	0.384175-00	00356000	0-149916-00	0.15877E-00	0.126185-00	0.803936-01	0.450056-01	0.333385-01	04 38302 F=01	0.33723E-01	0.30502E-01	0.263146-01	0+26149E-01	0.25898E-01	1			10.477.00	0.96924E C.	0.555500 01	0450375E 01	0.46794E 01	0+42796E 01	0.33431E 01	0.2945BE 01	0023126E 01	0.17969E 01	0.13957E 01	0.654105 00	0.64687E 00	0.346165-00	0.24445-00	0.10898E-00	0.784685-01	0.55183F=01	0.354845-01	0.343726-01	0.361705-01	0.441255-01	0-486546-01	0.965586-01	0.595055-01	0.62732E-01	0.693646-01	
8 .	<u>.</u>	0.24271F A1	0.25964E 01	0.236175 01	0-21729E 01	0.1709FE 01	0.14998E 01	0.98710E 00	0.78193E 00	00000000	0-34845-00	0.240765-00	00-100985-00	0.13896-00	0-12715E-00	00122275-00	0-12037E-00	0.11902E-C0	00115876-00	0-103451-00	0.94613E-01	0+85315E-01	0.662368-01	0-976746-01	0.502316-01	10-3640400	0.359896-01	0.33990E-01	0.333286-01	.95	4		0.50905E 01	0.50184E 01	0.48078E 01	0.4049E 01	0.35502E 01	0.2495E 01	0.19979E 01	00155228 01	0.868695	0.63657E 00	00-36276-00	0.30375E-00	0.273215-00	0.240125-00	0.25967E-00	04245695-00	00-229106-00	0.205926-00	0-153786-00	0.12730E-00	0.834146-01	0.67902E-01	0.000978-01	0.462126-01	0.44736-01	0.447646-01	0.44906E+01	
x = 1	1	0.26273E 01	0.26129E 01	0.29006E 01	0.224085 01	0.219796 01	0.18329F 01	0.14902E 01	0.192400	0-12067F 01	0.10979E 01	04 784125 00	0.47243E 00	0.56853E 00	00-11-10-00	0.32636F-00	04270165-00	00-34612290	0+14770E-00	0-150016-00	0-345640	0.667435-01	0.58350E-01	0-90412F-01	0+400386-01	0.369475-01	0.34885E-01	0633709E-01	10-3076566	T = K	4		0-50905E 01	0.50516E 01	0.47509E 01	0.45021E 01	0.36592E 01	0.34902E 01	0.310705 01	0.23464F 01	0.19895E 01	0.135945 01	0.10945E 01	0.84524E 00	0451076E 00	0.301005-00	0.199955.00	0-14122E-00	0.48859E-01	0.503525-01	0.387416-01	0.30319F=01	0-30506F-01	04321646-01	0-37274E-01	0.39796E-01	0-43566E-01	0.445675-01	10-1004-4-0	
	<u>.</u>	0+22718E 61	0.21707E 01	6.20504E 01	0-17061E 01	0.15010E 01	0.108166 01	0.88445E 00	0.549936 00	0-419796-00	01234347-00	0-179116-00	n-14052E-00	00-110011-00	0.90018E-01	D-309696-01	0-31/40/00	0.971046-01	0.952366-01	0.918095-01	0.81266F-01	0.749716-01	0.68600E-01	0-57064F-01	0.524245-01	0.48747E-01	0.46096E-01	0-13460-01		8	<u>.</u>		0-44466 01	0.425946 01	0.397116 01	0430298 01	0.27208E 01	0.226268 01	0.14329E 01	0.10945E 01	0.618345 00	0.451116-06	0.347755-00	0022185F-00	0-237885-06	04233285-00	0+23071E-00	0+22360E-00	0.193176-00	0-171598-00	0-14798-00	00-36610100	0.823485-01	0.53678E-01	0.44672E-01	0.35031E-01	0.330465-01	0.31216-01		
.t. x :1.		0.227185 01	0-22296	0.20919F 01	0.20004E 01	0e17721F 01	0.164365 01	0.13740F 01	0.123898 01	0.110705 01	0.86138E 00	0.75057E 00	0004903E 00	0=4752BF=00	0.40300E-00	00-33666-00	0.23929E-00	0+20009E=00	04187355-00	0-117985-00	0-940065-01	0.83598-01	0.667105-01	0-376946-01	0.52398E-01	0-48553F-01	0.444516-01	0.439405-01			, ,	0.44444	O. 44643E 01	0.43687E 01	0.42140E 01	0.37943E 01	0.346735 01	0.28307E 01	0.25018E 01	0.21786F 01	0.19789E 01	0.13136E 01	0.867375 00	00.000000	0.536885 00	0.91112E-00	04231255-00	0-122896-00	0.88762E-01	0+645776-01	0-375446-01	0-31274E-01	0.268756-01	0.26995E-01	0.269655-01	0.301035-01	0.310416-01	0.31860E-01		
4	ь,	0.0	10.0	20.0	29.0	35.0	0.04	0.00	0000	000	70.0	0.6	0000	90.0	0.50	10501	110.0	115.0	129.0	130.0	135.0	0001	150.0	199.0	160.0	170.0	175.0	180.0			0	•	9.0	000	2002	25.0	0000	0.0	000	0000	0.09	0.0	75.0	0.0	9000	0.66	10500	110.0	119.0	125.0	139.0	1950	145.0	150.0	160.0	169.0	175.0	100.0		



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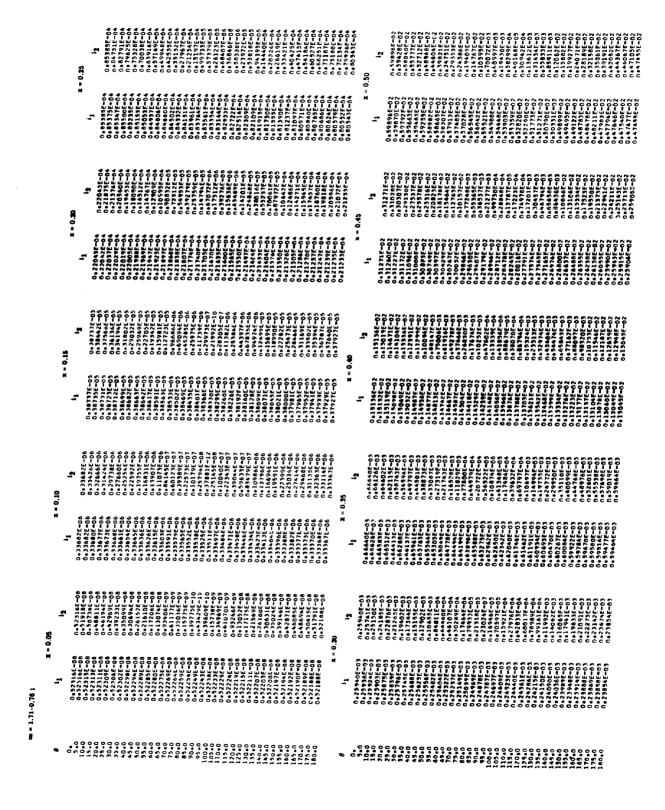


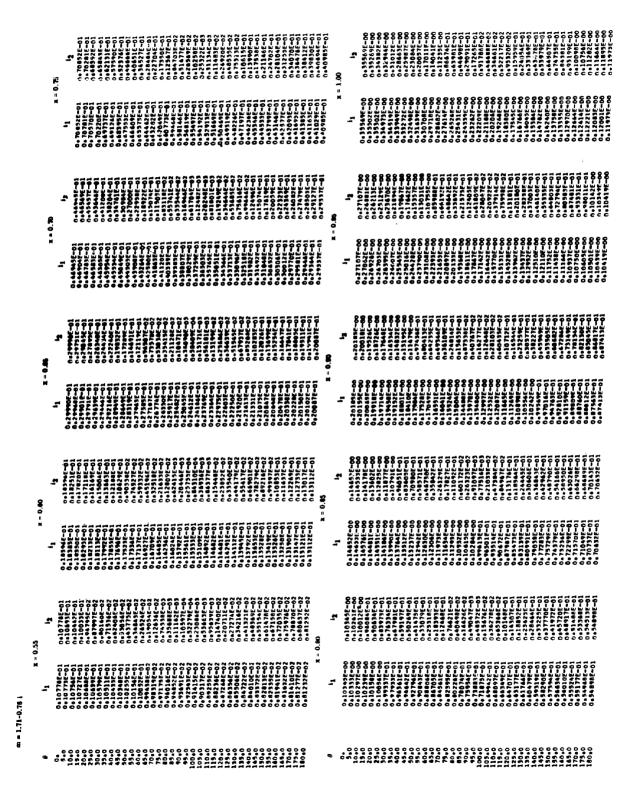


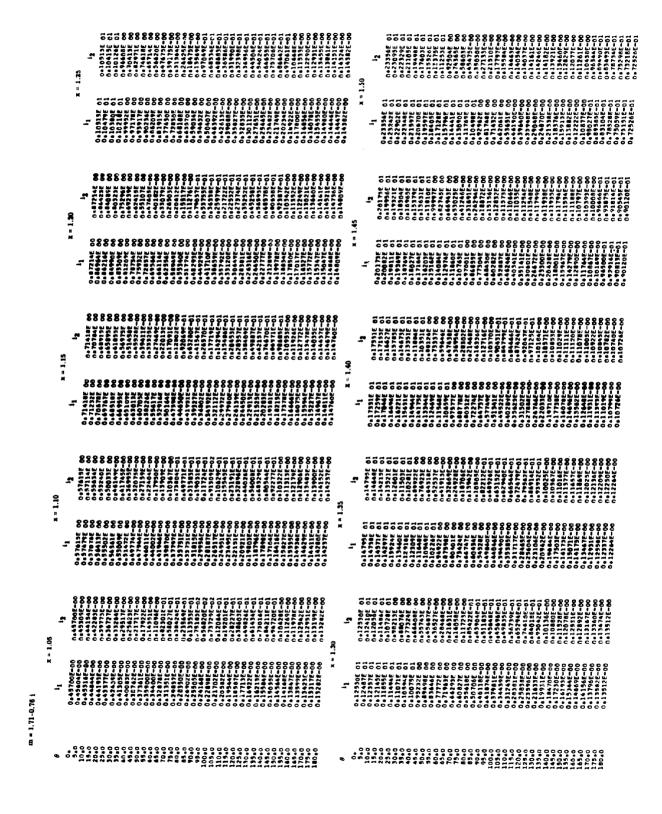


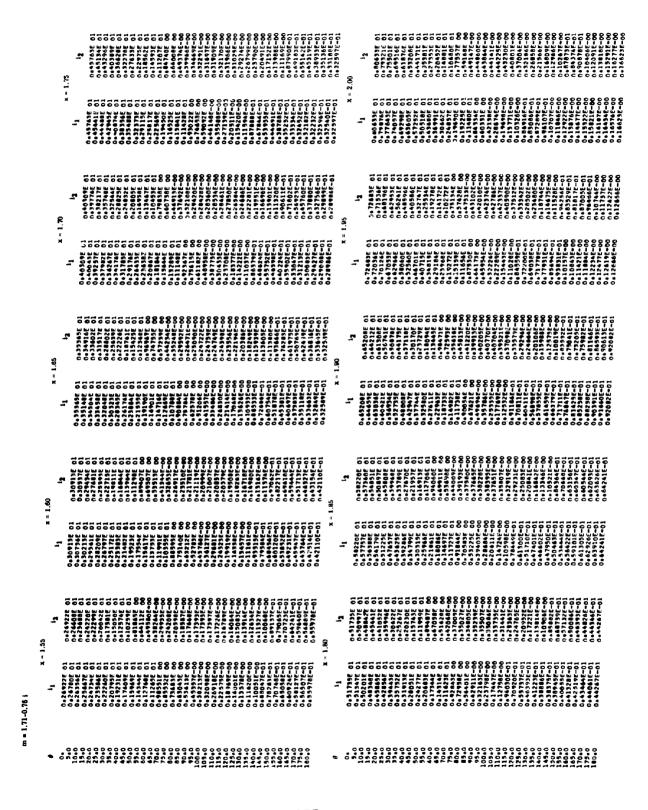
			- 25	•			01 0410949E 01																							1.5																						0-11906F-00			
				τ,			0.104838 01																							*	4																					0+10248E-00			
-		•	X = 1.30				0.78838E 00																						-			•	-0	•	90	•		00	•	•	00	0 (90	00	0	00	Ö	Ö	ŏ	ŏ	ŏĠ	0-116356-00	ŏŏ	ŏć	
				0421365	0-910846 00	0.0911288	0.881895	0.8354E DO	0.176000	0.74264E 00	0+67153E 00	0.43488	0.561916 00	045244E 00 0449212E-00	045923E-00	0.396595-00	0+37112E-00	0.345675-00	00.300875-00	0.28149E-00	0.248315-00	0.23478E-00	0.21239F-00	0.203596-00	0.190385-00	0.185861-00	0-180765-00	0+160125-00				0.205386 01	0.Z0Z68E 01	0+19228F 01	0.18486E 01	0.16663E 01	0-15626E 01	0-134276 01	0.11205E 01	0+10134E 01	0-81431E 00	0.64136E 00	0.56585E 00	00+3648E+00	0.383185-00	0.29486E-00	0+25936E=00 0+22693E=00	0.20302E-00	0+16285E-00	0.14768E-00	0e12534E-00	0-11756E-00 0-11174F-00	0+10770E-00	0e10454E-00	
		= 1,15	r ,	00 30892																								1	1.40	<u>.</u> .	0.177645 01	0.17565E 01	0.16048E 01	0+14621E 01	0.117735 01	0.647395 00	0.55045E 00	00-12766-00	0+32526E-00 0+24418F-00	0.18370E-00	0.114305-00	0-103595-00	00-104216-00	0.120765-00	0412971E-00	0-143/06-00	0-146416-00	00147946-00	0e14565E-00 0e14281E-00	0+13945E-00	0.13598E-00 0.13274F-00	0 + 1 3 0 0 E - 00	00) 2662E-00	0.1261BE-00	
				00 765305 00	0e79775E 00	0-735716 00	00 71979E 00	00 36462 900	0-63057E 00	0+6038BE 00	0.54805E 00	0+51969E 00	04-6376E-D0	00-10765-00	0-30590E-00	0.340185-00	00-31951E-00	00-282776-00	0.26671E-00	042321BE-00	0.22754E-00	0-217346-00	0.20094E-00	0019464E-00	0.185635-00	0a18286E-00	0.180655-00	1	"	-	0+17764E 01	0017479E 01	0.17130E 01	0.16065E 01	0015378E 51	0013777E 01	0.119976 01	0411085E 01	0.92934E 00	00 762815 00	0.686516 00	00.55036E 00	00490946-00	0.389305-00	0.309125-00	0.27630E-00	0822324E-00	0e 20226E=00	0.16960E=00	0.15731E-00	0-139496-00	0.13356E-00 0.12943E-00	0.126995-00	00-1819710	
	. 1.10	£		0.622115 00																								1.35	ŗ,	0.152245 01	04150600 01	04138115 01	0+12793E D1	00102575 01	0.88641E 00	0.614486 00	0.384496-00	0+29322E-00	0.163055-00	0412277E-00 Ca96736E-01	0-826226-01	0+80030E-01	0.86722E-01	0-10619E-00	0-116165-00	0-13254E-00	0.142505-00	0-1452E-00	0-147435-00	0.147485-00	04146755-00	0-146345-00	0-14593E-00		
	Ħ	- -	0.62782£ 00	00 422276 00	00 61543E 00 00 60 60 3E 00	0.59425E 00	0+564426 00	0.527975 00	04507978 00	0 + 6719E-00 0 + 6590F-00	0-444386-00	0.401645-00	0-34084E-00	0+341316+00	0.32283E-00	0+26895E-00	0-27367E-00	0.246565-00	0-234755-00	08214576-00	0.20615E-00	0.192425.00	0+18726E-00	0.18299E-00 0.17949E-00	0-177365-00	0 17596E-00	1	H	<u>.</u> -	0.15224E 01	0.15167E 01	0.14726F 01	0el3889F 01	0.13346E 01	0.12/37E 01	0=11375£ 01	0.99125E 00	0.845245 00	0.774945 00	0.64372E 00	00528316 00	00-1723E-00	0.36856E=00	0.350786-00	0.287445-00	0+26137E-00	0.21906F-00	0.202305-00	0-17631E-00	0-166675-00	0-193205-00	0.14914E-00 0.14673E-00	0+14593E-00		
	1.05	*	0.502996 00	0.48968E-00	0-43930E-00	0-354816-00	0.322146-00	0.2339&E-00	0-15181E-00	0-11604E-00	0-59075E-01	0.2386875-01	0-14-76-01	0.962755-02	0-13441E-01	0+301775-01	0-417695-01	0.547275-01	0-82559E-01	0.10031E-01	0-121465-00	0-13251E-00	0+15028F-00	0-15671E-00	0.164185-00	0.16512E-00	1.30		•	00129775 01	00124456	00109//2 01	00 337 30E 00	00 17710000	00 7792400	00+324/1-00	00-24T04E-00	0019475E-00	00104895-00	Consider to	10-20198C+0	10-34746000	10-1x0xcx00	10-1007/200	00-3+440400	20-241-2	001222200	00-1400+140	00122212	U. 10.53.48.0	00-10750-00	00-10-00-00	00-317		
- 1.60-1.00 i	Ä	0.507505 00	003040600	0-498555-00	00483216-00	0.473095-00	00-44 B 74 E - 00	0-434916-00	0.4049ZE-00	0-373196-00	0-357156-00	0-325956-00	0-310296-00	0+281396-00	0-255255-00	0.24336E-00	0.2221F=00	0+212775-00	0-104305-00	0-169935-00	0-184046-00	0-174705-00	0-171235-00	0.166645-00	0-165508-00	00-17F-00		<i>-</i> 7	0.129/SE 01	0012929E 01	0 14 2 d / 0 1	0.122706 01	0.11510E 01	0011032E 01	0.99971E 00	0001726E 00	0.82016E 00	00 10407 00	U-54924E CO	00 366736 00	U-45689E-U	30-36-44-0	201775	00.31047e=00. 00.2921.50=00	00-1149976	00-346-00	0041501E-00	0617075	Uelelane-un	0010414500	00-10-10-00	0016221E-00		•	
E.	•		10.0	15.0 20.0	25.0	35.0	0.04	0	60.00	65.0	75.0	0.0	90.0	100.0	105.0	115.0	120.0	130.0	135.0	C 0 1	150.0	155.0	165.0	170.0	180.0			90	•0	10.01	15.0	250	30.0	0	0 0 0 0	950	6.5.0	7000	0	000	950	1090	110.0	12000	130.0	13960	149.0	150.0	160,0	170-0	175.0	16000			

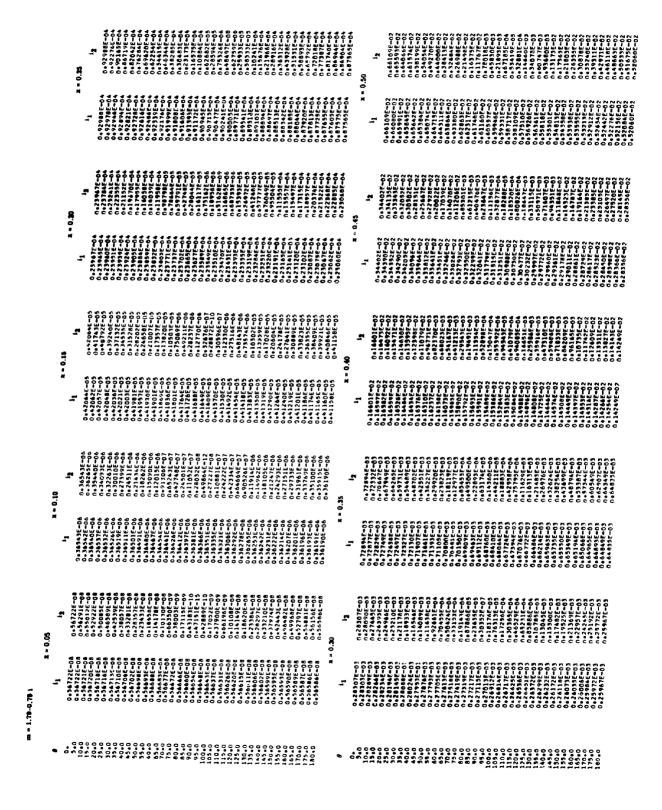
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	1	5	<u>.</u> r	0.4044ZE 01 0.40354E 01	0.38643E 01	0.32440E 01	0.20404E 01	0-198126 01	0-121576 01	0.91130E 00	000 30100	0+372696-00	0+30273E-00	0.266696-00	00-305.2	0.309165-00	04317426-00	04305155-00	0.28496F-00	00-23415-00	00193235-00	0-140936-00	00103865-00	0.04946-01	0=520505-01	0.4.002E-01	0.434416-01	0-403426-01	ă	g	<u>-</u> -	0.724985 01	0.71272E 01	0.427076 01	0.55017E 61	0.304116 01	0+10923E 01	0.16954E 01	0.119996 01	0059261E 00	0.451938-00	0-3794SE-00	0-34036-00	0.43018E-00	0.416895-00	0.38914E-00	0424035-00	04238426-00	0+191175-00	0.12499E-00	0-104526-30	0-11120E-00	0+12384E-00 0+14029E-00	0.15719E-00	0-16142F-00	00-14481-00
			1.	0.40401E 01	0.300000	0.34997	0.32527E 01	0.272215 01	0.2441E 01	0.21676E 01	0.16439E 01	0.14063E 01	00 34160	0.62217E 00	0+54540	0.43821E-00	0-349476-00	0.21 900E-00	04173046-00	0-109935-00	0.89298E-01	0.62946	0.55080E-01	0.49604E=01	0=43446-01	0.41901E-01	0-404916-01	04403425-01	•	•		0.72498E 01	0.71876E 01	0.670401	0.58431F 01	0. 53122E 01	0.41640F	0.35870€ 01	0.252426 01	0.20588E 01	0-129456 01	00 72634F 00	00.56479E 00	00-307426-00	0.22952E-00	0.143825-00	0+12577E-00	0+11852E-00	0-123456-00	0.19137E-00	0-149641-0	0+15840E-00 0+16622E-00	00172835-00	04181826-00	0.1840BE-00	
	1.65	7.	0.36002E 01	0.3506E 01	0.31758E 01	0.25330£ 01	0.21637E 01	0-144005 01	0.0112245 01	0.63299E 00	00-306-00	0-284476-00	0.249416-00	0*245416-00	0+2980E-00	0+282135-00	0.28434E-00	00-264146-00	0.24340E-00	0.218015-00	00-140185-00	0-135156-0	0.111046-00	0.736546-01	0.606956-01	0.44993E-01	0-413476-01	10-20-10-10	8.	.51		0.65277E 01	0.61092E 01	0056208E 01	0042873E 01	0+26252E 01	0.21660E 01	0.11448E 01	0.80916E 00	00434176-00	0.37071E-00	0.348115-00	0.39137E-00 0.41128E-00	00-102610-0	0.36728F-00	0435029E-00	0.304635-00	0.20918E-00	0-14815E-00	0+11364E-00	0.10162E-00	0.10262E-00	0.11105E-00 0.12119E-00	0-130525-00	0-139326-00	
	il H	-	0.36002E 01	0.39149E 01	0e32714E 01	C-31004E 01	C+24685E 01	0.224602E 01	0.199105 01	00176146 01	0.13347E 01	0.11430E 01	0.01974E 00	0.4792E 00	00 460396-00	0-375186-00	0.245887.00	0-1983E-00	0+14055E-00	0.10669F-00	0-3+0+G+0	0.7440E=01	0.962206-01	0-304366-01	0.434426-01	0-419736-01	0.403806-01		- K	1,	0.663776 0.	0.66767E	0.63182E 01	0.57291E 01	Co+8647F 01	0.437316 01	06336115 01	0.207316 01	0.19942E 01	0.14189£ 01	0.10130E 01	0.78152E 00	0044575E-00	0+33211E-00 0+24772E-00	0.1874E-00	0.146465-00	0 10 5 0 1 E - 00	0.993046-01	0+10128E-00	0-106265-00	0+11850E-00	0.12438E-00	0013374E-00	0 1 3 6 6 2 E - 0 0	0. 13932E-00	
Ş	. <u>.</u> .		0.31517E 01 0.31100E 01	0.27934E 01	0.25470E 01	0.1935E 01	0.150466 01	0-10312E 01	00 34035 00	00+4304E-00	00.33505E-00	0-22729E-00	00-21276E-00	0+223486-00	0-235445-00	00-350036-00	04248265-00	00.27775-00	0.206305-00	0+144396-00	0-161276-00	0 - 1 1 7 2 /k = 00	0.9446E-01	0.694145-01	0=60322E=01	0.49761E-01	004#502E-01	.85	4	•	0.50512E 01	0.548767 01	0-50624E 01	0.3493F 01	0.32428E 01	0.201415 01	0+15035E 01	0.77335£ 00	0.55505E 00	0.33345-00	0.34795-00	0.366268-00	0.3967565-00	0.39754E-00	0.350805-00	00-311346-00	0.22195-00	0.180388-00	0.14521E-00	0.998016-01	0.09948-01	0.87308E-01	0.937635-01	0.992435-01	00-37200100	
# #		0.11417	0631342E 01	0029974E 01	0027629E 01	0029805E 01	0622123E 01	0.20163E 01 0.18189F 01	0.16243E 01	0.12645	0.10925E 01	00 94010E 00	0147984E 00	0057200E 00	04398025-00	0.329895-00	0.272766-00	0-100416-00	0.19473E-00	00129175-00	0-94646-01	0.600016-01	0-627765-01	0-572186-01	00505065-01	0-440046-01	0.48502E-01	x = 1.	-5	- !!	0.566126 01	0.56736	0.54594E 01	0.48250E 01	0.408978 01	0.35704E 01	0.270155	0.22949E 01	0.197936 01	0.12792E 01	0.001095	0.42011E 00	0.35666-00	0.2700BE-00	0-196316-00	0.123375-00	0.889800	0.82392E-01	0-80826-01	0.633236-01	0.902296-01	0.936195-01	0.967065-01	0.10006F-00	*	
1.55	ŗ.	0.27479£ 01	0.271265 01	00244778 61	0-198016 01	0-17203E 01 0-14474F 01	0-118436 01	04 730405 00	0.55429€ 00	0+314118-00	0.24546E-00	00106336-00	0.10319E-00	0-189046-00	0.206476-00	0+21504E-00	0621351E-00	0.205016-00	00172305-00	00-3/16/100	00-141306-00	0.108345-00	0.94642E-01	0.83395-01	0-663696-01	0.646996-01		8	٠.	0.52199F 01	0.513986 01	0.45374F 01	0.40646	0.35228 01	0.238745 01	0.140455 01	0.10366 01	0.539258 00	00-346-00	0+31402E-00	0.319378-00	0-398965-00	0+379908-00	0.365996-00	0.34283E-00	0-270+96-00	0.22891E-00	0-152646-00	0=122775-00	0.84232E-01	0-747176-01	0-667936-01	0.692146-01	0.704328-01		
1 = x	. 7	0.27475E 01	0.26911E 01	0.29294 01	0.241505 01	04213626 01	0.10175	0.16527E 01	0.13306E 01	0e11788E 01	0-90360E 00	00 399286	0.5765E 00	0-491065-00	00-14436-00	0.297715-00	0+29194E-00	00180616-00	0-194046-00	0-1142416-00	0.10075-00	10-364 568-0	0-7-2265-01	0+693521-01	000000000000000000000000000000000000000	0-634726-01	-	•	-	0.52199E 01	04507036 01	0.488985 01	0.454748 01	0.40192E 01	0.327978 01	0.28988E 01	0.21479F 01	0.1834£ 01	0.12576	0.101925 01	00 26220	0+30064E	0.295416-00	0.22936E-00	0-133845-00	0-106446-00	0.758255-01	0+687696-01	0.43982E-01	0.64299E-01	0.468005-01	0-682326-01	0.701685-01	0.70492E-01		
4	ъ,		19.0	20.0	0.00	0.04	9	9.66	0.0	70.07	75.0	0.00	0	0.5	105.0	0.01	120.0	125.0		140.0	145.0	155.0	140.0	169.0	179.0	180.0		q	s	.00	10.01	15.0	25,00	30.0	0.0	0 0	0.56	0.00	2000	90.08	0.60	0.56	100.0	110,0	1150	129.0	130.0	140.0	145.0	199.00	0.00	170.0	175.0	2		

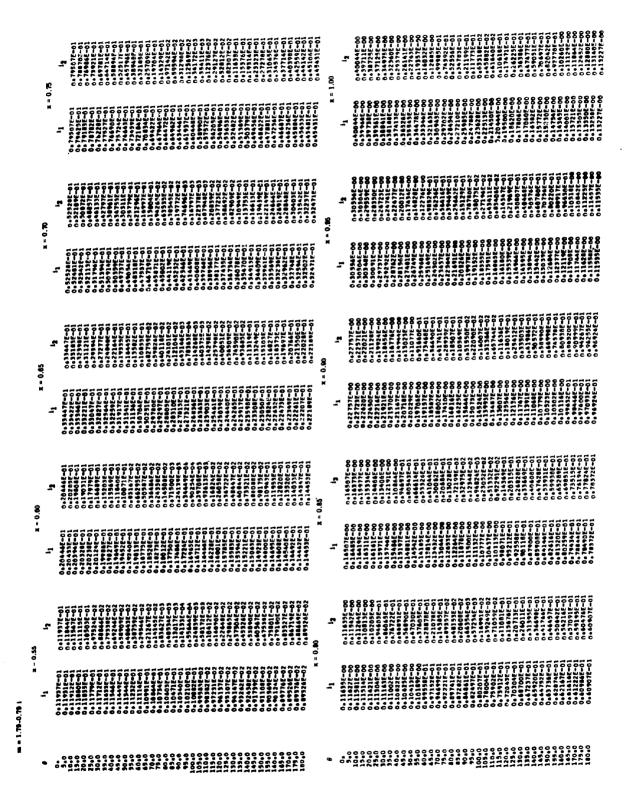


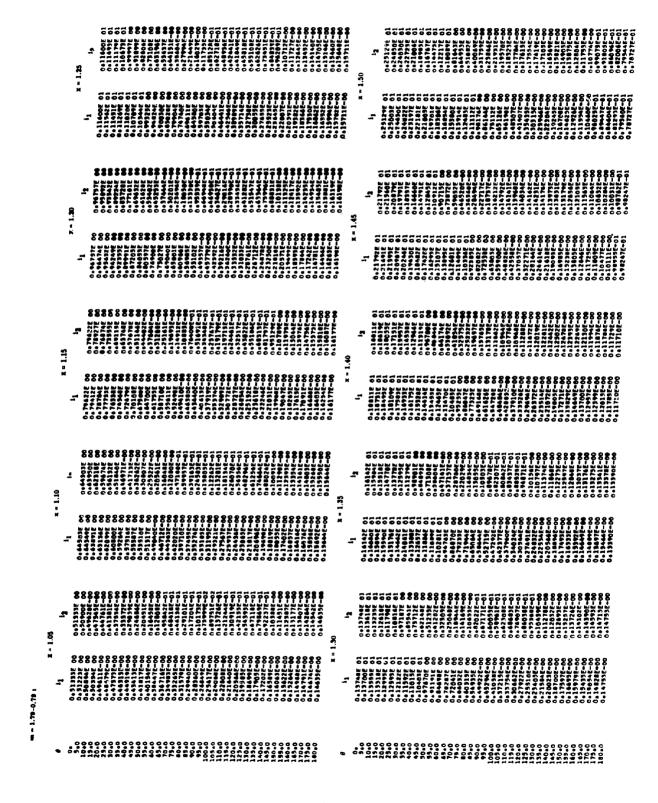


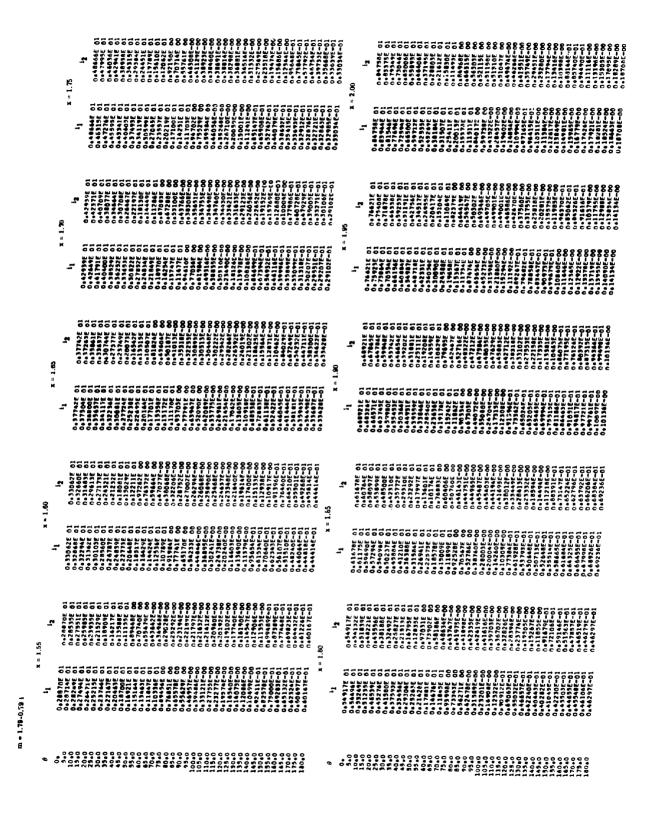


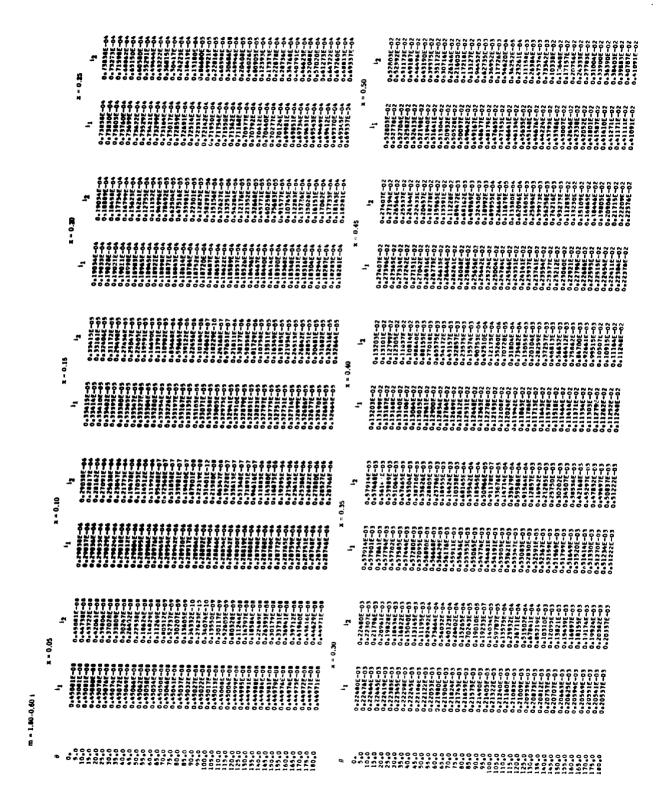


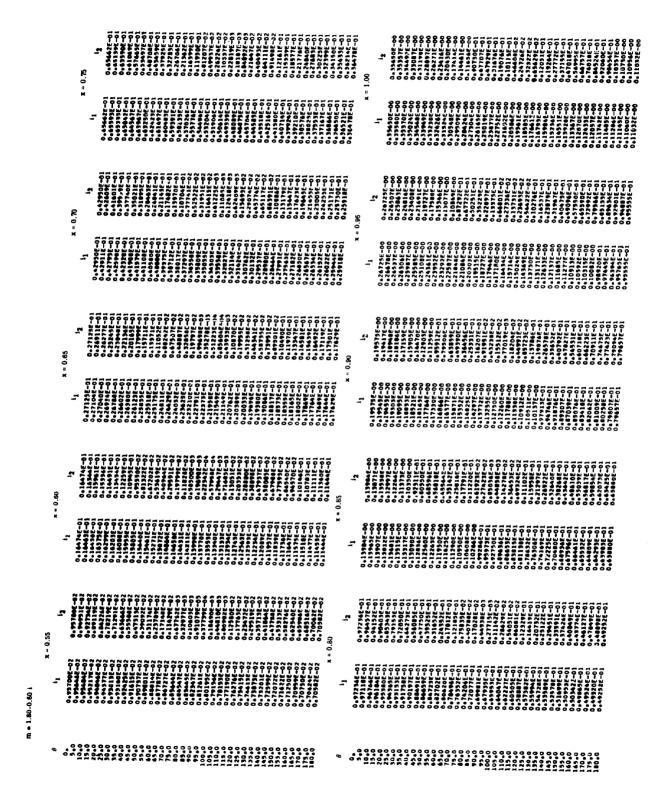


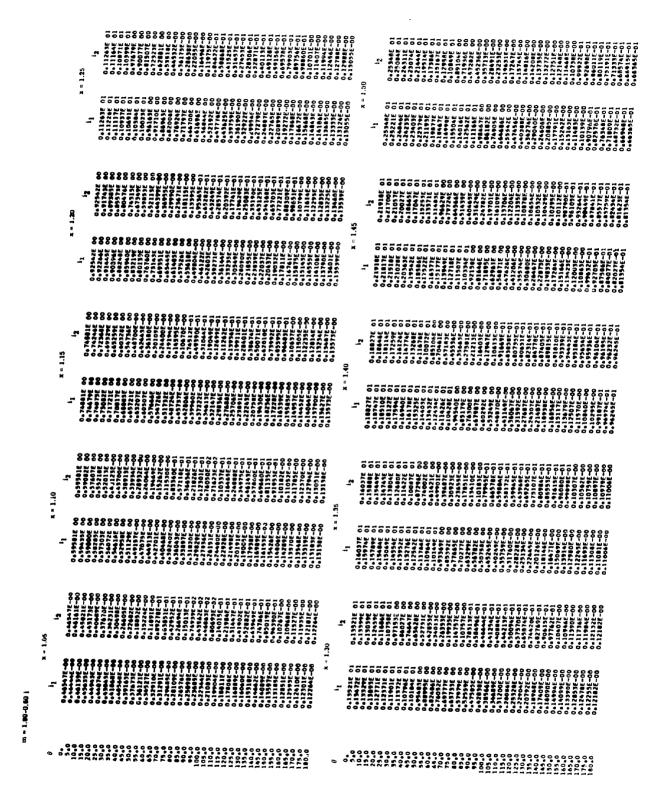


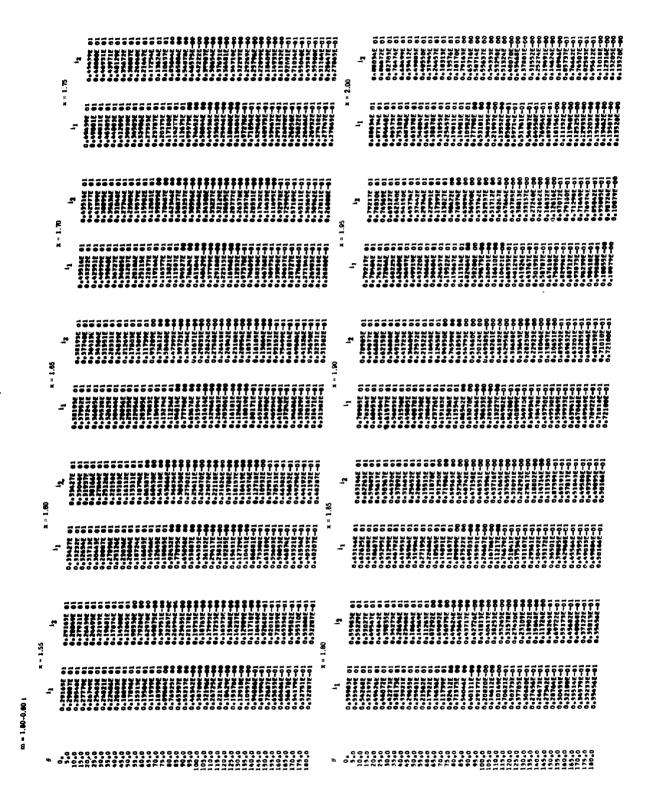


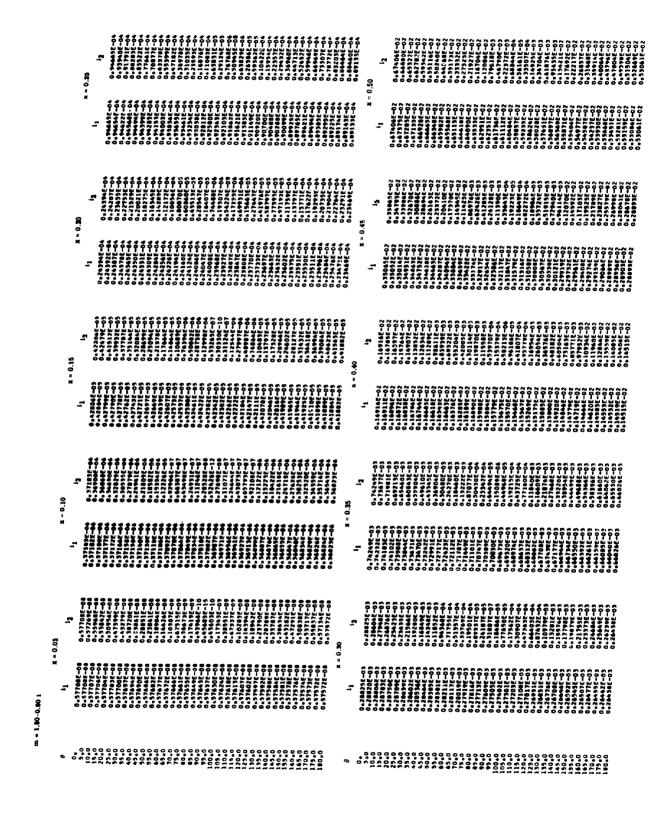


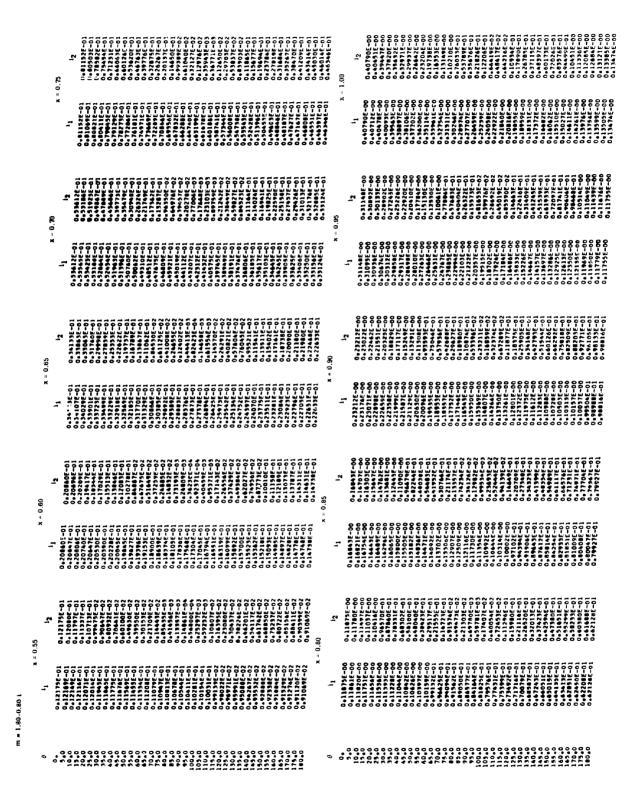


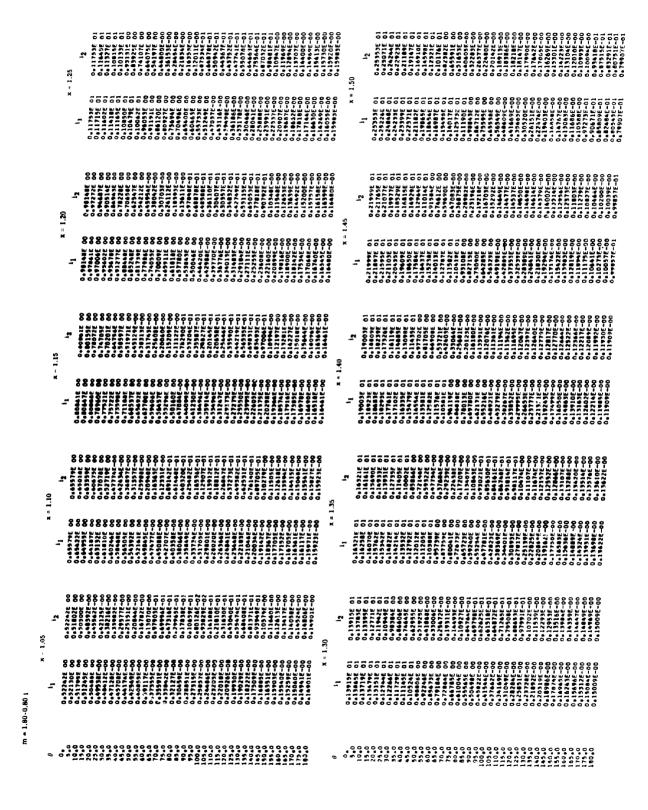


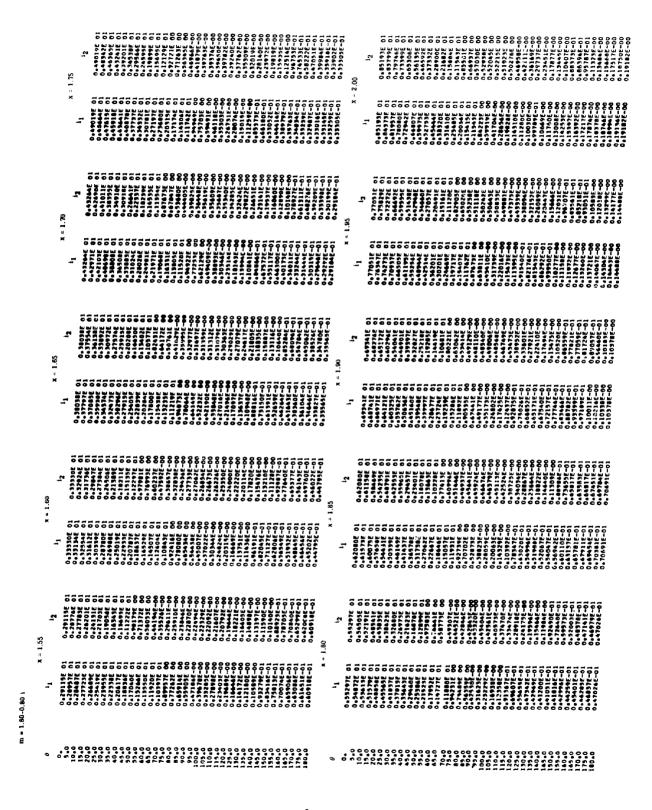


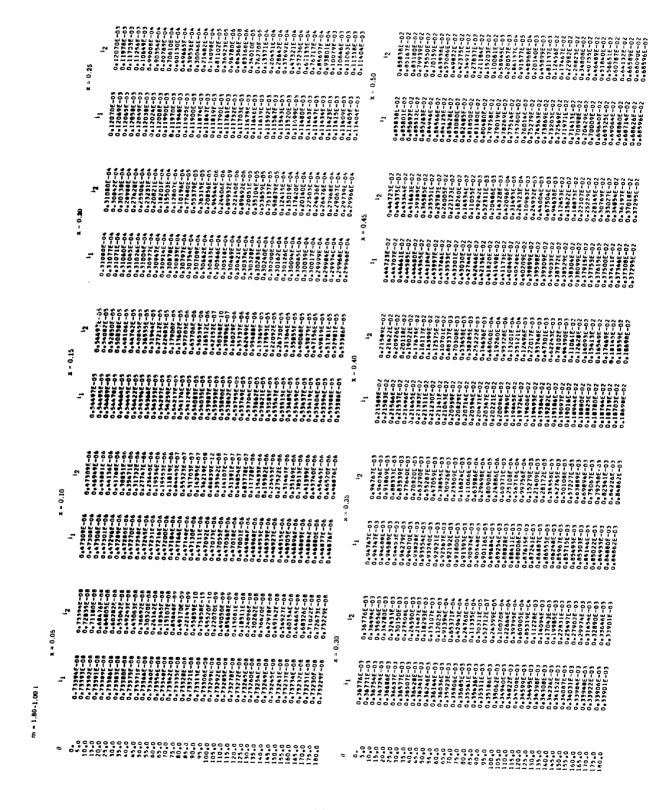


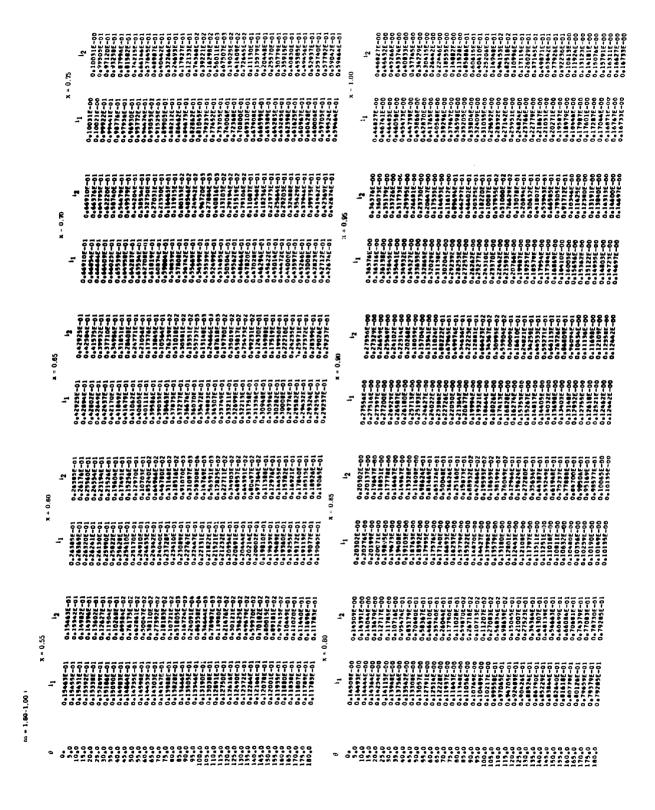




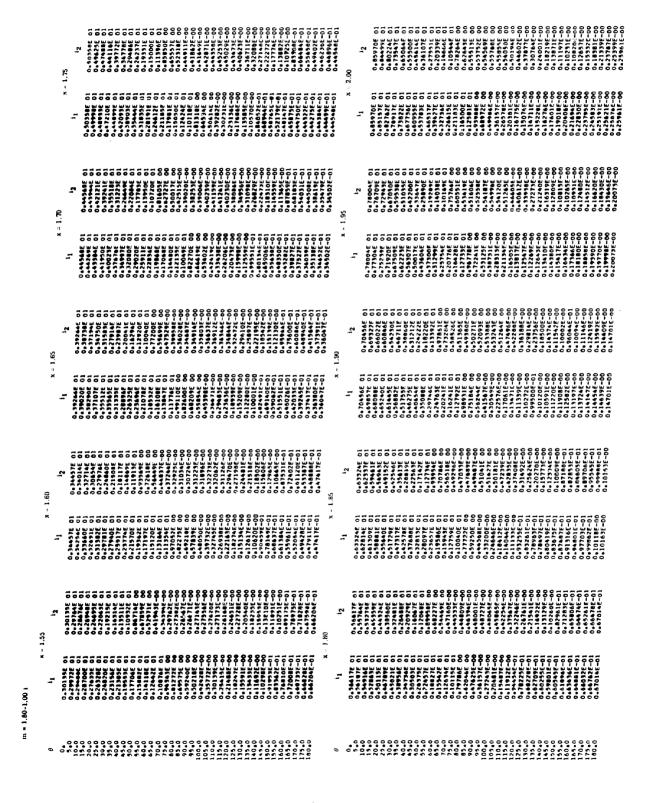






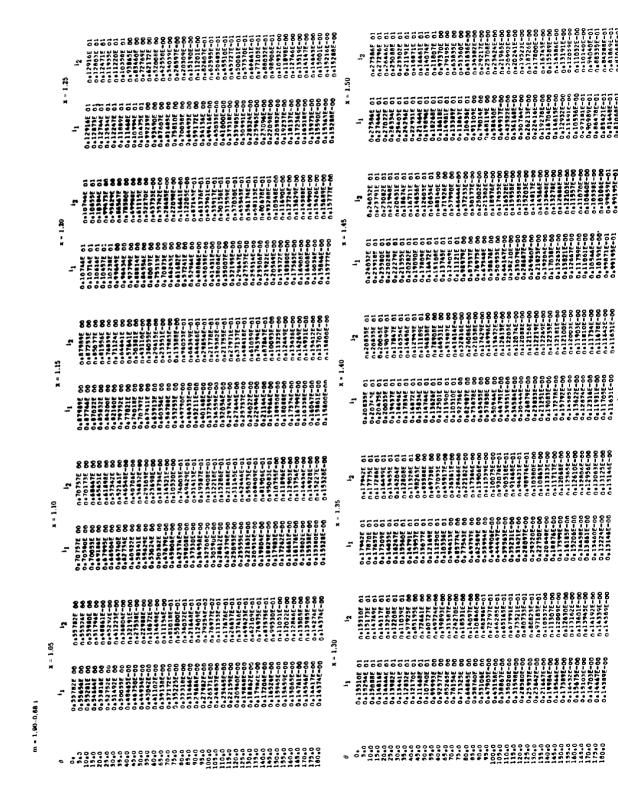


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			- -	0+12513E 01	0.12473E 01	0.12359E 01	O. 1167BE O1	0-115692 01	0.107486 01	0.10272E 01	0.97656E 00	0-049778 00	0.815325 00	0. 76130E 00	0.70837E 00	0.65711E 00	0.541485	00517775 00	0.4770BE-00	0-439925-00	0-405125-00	0-3686760	0432036-00	0-297895-90	0.27807E-00	0.260756-00	0.24575E-00	00-32126-00	0-21328-00	0.20610F-00	0.2006&E-00	0.196835-00	0.193805-00		X = 1		-	0.26276E 01	0.26144E 01	0.25754E 01	0.24366 01	0e23199F 01	0.219735 01	0.20615E 01	0-19696 01	0-16120E 01	0-149976 01	0-191126 01	0.10347F 01	0.91004E 00	0+79570E 00	0.5892196 00	0.51765E 00	0-445905-00	0-363656-00	00330146-00	0.24592F-00	0-319615-00	0.18646E-00	00-14-09E-00	0-13074E-00	0+11670E-00	0-105185-00	0.96238E-01	0.92605E-01	0.90404E-01	0+89682E-01
		1.20	ŭ	0.10574E 01	0-10-77	0.97273€ 00	0.911106 00	0.73370F 00	0.064798 00	0.57384E 00	0-344056-00	0-350656-00	04250976-00	0-141276-00	0-103405	0.796715-01	0-370015-01	0-46737E-01	0-36165-01	0-54646-01	0.64269E-01	0.771655-01	0.91383E-01	0-106146-00	0-14696-00	0-140116-00	0-160696-00	0.17069E-00	0+17978E-00	00-382/81-0	0-147485-00	0-19991E-00	0.2007&E-00		1.45	5	7	0.22844E 01	0.21645F 01	0.20634E 01	0-19057E 01	0-17199E 01	0.15165E 01	0.109956 01	0.90501E 00	0.72978E 00	00378835 00	0.356465-00	0-264376-00	0-20-316-00	0-188225-00	0.182495-00	0.183165-00	0-186916-00	0.194176-00	0-19480E-00	0-1926BE-00	0-167695-00	0-140076-00	0-16276E-00	0-19309E-00	0-149825-00	0-120295-00	0+12259E-00	0.11847E-00	0-115975-00	00-3112100
		= #	.	0.105746 01	0.10494E 01	0-103046 01	0-96512E 00	00 3686600	0.884916 00	0-84359£ 00	0.80432E 00	00 70100	0-676086 00	0+5382E 00	0.99262E 00	0.55290E DO	0.514948 00	00-71-00	0-14105-00	0.38318E-00	0-35869E-00	00-30000-00	00-36831-00	0.27609E-00	0.260B1E-00	0.247936-00	00-34096-00	00-31897-00	0-211946-01	0-20701E-00	0.20352E-00	0-201445-00	00-38/00780		d H	.	0.238465 61	0.22739£ 01	0.224286 01	0-219196 01	0.21220	0.199876.01	0-102075 01	0-171056 01	0-198708 01	0 1 3 4 7 6 0 1	0+121005 01	04109128 01	0-977325 00	0.771626 00	0.68112E 00	0.599216 00	0.329856 00	0.403586-00	0-3572E-00	0-310596-00	04273576-00	0.214105-00	0-192#7E-00	0-174176-00	0.158745-00	00-134116-00	0.12826E-00	0-122385-00	0-119326-00	0+115146-00	
	1.15	_	2.	0.87545€ 00	0.03214E 00	0.74100 00	0.78349E 00	0-341347 00	0-486016-00	0-411095-00	04272476-00	0.212926-00	0.16124E-00	0-1162BE-00	0.50.7E-01	0.42074-01	0.328345-01	0.300536-01	0+35824E-01	0-401005-01	0-900000	0.76835F-01	0.944765-01	0-110305-00	0.125765-00	0-140436-00	0-144100	0-176625-00	0-182405-00	0-19232E-00	00-19732E-00	0.201336-00		1.40		<u>.</u> 2	0.19797E 01	0.19585 01	0.189415 01	Coldador of	0.15094E 01	0.13302E 01	0.11599€ 01	0.983045 00	00.461395.00	0.52690F 00	0-11916-00	0424236-00	0.20495E-00	0-172655-00	0-146995	0-145025-00	0-14841E-00	0.1538ZE-00	0.159465-00	0-167245-00	0.1684E-00	0.16785E-00	0-10-14-00	0-158475-00	0.15456E-00	0.1505BE-00	0.14.01E-00	0-141905-00	0.14056E-00	0-140116-00	
	×		0.68114	0-081085	0.043175 00	0.047905	0-628816 00	0.78072E 00	0-79261E 00	0460746 00	04698015 00	0.62468E 00	00 341246 00	0.42446 00	00-20-00	0.463695-00	00-3004600	0.40792E-00	00-396186	0-146776-00	0+316186-00	0.297926-00	0.28145E-00	00-211-00	0.282216-00	0-232306-00	0.229885-00	0+21686E-00	0+21120E-00	00-368602-0	00.201948.000	0-201936-00		- x	-	-	197976		ò	361501	17832E	70406	42016	10199	13169	121336	101126	91961E	2512E	40346	3908 E	32607E	44797E		330706	20507	265995	70205	200002	184626	7194E	704	0-147635-00	14342E	366001	3110	
	1.10	12	0.72748E 00	0.721335 00	0.67220E 00	0.582385 00	0.52470E 00	0-464735-00	0.342916-00	0.28328E-00	0.22759E-00	041111111111111111111111111111111111111	0.96609E-01	0067269E-01	0.45351E-01	0 30361E-01	0.200946-01	0-329216-01	0.299895-01	0-40454E-01	0.539776-01	0-3460346	0.997785-01	0+11943E-00	0-130416-00	0-14496-00	00-137605-00	0-178105-00	0.185846-00	0.19123E-00	0-19449E-00	0.19558E-00	.35	!	7	0.170895 01	0.16916	0.16398E 01	0.15970E 01	0.10.70.01	0.117416 01	0.102336 01	0.87240E 00	0.72750E 00	0.0393626 00	0.372895-00	0.28957E-00	0-224506-00	0.143785-00	0-12982E-00	00-114076-00	0-111925-00	0.12120F-00	O. 1ZBBE-00	0e 13679E-00	0.14409E-00	0.15514F-00	0+15648E-00	0.16104E-00	0.162455-00	0.163325-00	0.16322E-00	0.16300E-00	0-162565-00	0.162535-00	•	
	×	Ι,	0.72768E 00	0.72098E 00	00.701415 00	0+60722E 00	0-470436 00	00.43030£ 00	0.60764E 00	0-948816 00	0.53354E 00	0.50794E 00	0.42476-00	00-3667690	00-310-00-00	04384375-00	0.364776-00	0.34440E-00	0-353346-00	0.201926-00	0.276396-00	04262846-00	0.250635-00	00-36266200	0-231886-00	0-214776-00	0.208825-00	0.204016-00	0-200305-00	00-3/9/6/6	00-104-00		X = X		-	0.17089E 01	0.17029E 01	0.16820F 01	0-140745	0.19936F 01	0.149065 01	0.14201E 01	0.134376 01	0.117005	0.109556 01	0.10116F 01	00 3046760	0 77410	0-70259E 00	00.39886 00	00 31367	0.46464	0420436-00	0.3799&E-00	0-343645-00	0.203095-00	0.259756-00	0.23901E-00	0.221356-00	0.194126-00	0.1840&F-00	0-17610F-00	0.165855-00	0.163346-00	0-16253E-00		
. VE	_	2.	0.58443E 00	0.54951E 00	0.51299£ 00	0.473655-00	0+360765-00	0-330475-00	0.241825-00	0.104125-00	0.14437E-00	0-10011E-00	0+524735-01	0.3414BE-01	0.21626E-01	0.14762E-01	0-130916-01	0-22-00	0-124705-01	0-96956-01	0.585365-01	0.735435-01	0-890216-01	0-119325-00	0+133295-00	0-14-016-00	0-197186-00	00-100001-00	0-17-00-00	0+162656-00	0.163956-00		2	2	9	6.7AE	35.5	100	3.		226E	7	1035	492E	236E	7,5	5646	3600	ž	9805	996E	935E	705		371E	3006	3005		\$15E	1996	707	900	0-179416-00	0.18051E-00	04180825-00		
,		0.989687	00.588265	0.570635 00	0.57540E 00	0.547828 00	0.539848 00	0.501735 00	0-404105-00	04469716-00	0-0-31-00	0-408986-00	04389995-00	0-310985-00	00-32565-00	0-11070	0.302985-00	0.28810E-00	0.27420E-00	0-26132E-00	00-34046700	0.228865-00	0+220105-00	0-21234E-00	0+20994E-00	0-16-61-0	00-3/804100	0-18785-00	0.18968E-00	0-166396-00	0-163956-00	•	- v	<u>-</u> -		0.146745 01	0 14470F 01	0-142215 01	0.13081E 01	0-1245/E 01	0-12-0-1	0.11798E 01	0.11196E 01	0-104895 01	0.91302F 00	0.84595E 00	0.78066E 00	0.446135 00	0.00208F 00	0.549785 00	0.301500	0-427316-00	0-361036-0	0.34870E-00	0.32000E-00	04272975-00	0.25338E-00	0.23689E-00	0.222885-00	0-201575-00	0-19393E-00	0-180125-00	0-184046-00	0-100825-00			
	6	•	0.0	15.0	23.0	0	0.04	0.64	9000	0.04	0.00	70.0	0.0			95.0	10000	000	2	120.0	129.0	130.0	135.0	0.00	0.051	155.0	160.0	165.0	0.071	000			,	θ	•0	0.0	1000	0.00	0.5	30.0	35.0	0		99.0	0.00	0.0	2 2	0.08	0.5	0.06	100	10,01	11000	0.00	125.0	130.0	135.0	0	120.0	155.0	160,0	0.021	175.0	180.0			



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			*		0.87764E-04	0.87793E-04	0-876675-04	0.875935-04	0 873855-04	0.67253E-04	0+869385-04	0.867985-04	0-863585-04	0.861425-04	0.856875+04	0+85451F-04	0.849725-04	0.847325-04	0.842625-04	0.840345-04	0.838156-04	0.934055-04	0.83217E-04	0.828355-04	0.82741E-04	0.825075-04	0-82417E-04	0.8229AF=04	0.822665-04	0.82256F-04	**			0.636535-02	0.63561E-02	0-632005-02	0.629345-02	0.62243E-02	0.61365F-07	0.60865E-02	0.597685-02	0-59181E-02	0.579536-02	0.57324E-02	0.560586-02	0.554315-02	0.542145-02	0.536338-02	0.52543F=02	0.52042E=02	0.515746-02	0.50750E-02	0.50090E-02	0.49826F-02	0-49437E-02	0.49315E-02	0.49216E=02
		i c	R .	<u>.</u>	0.225796-04	0-210935-0	0.210616-04	0+18594E-04	0.169196-0	0-132326-04	0-112736-04	0.741916-05	0.564135-05	0-26496E-05	0-152486-05	0-18129E-06	0.522936-09	0.640745-06	0.143935-05	0-304406-09	0.541826-05	0.71946E-05	0-10840E-04	0-127276-04	0.162595-04	0-178025-04	0.202135-04	0-210085-04	0.216595-04	0.45		12	0.32945E-02	0.32690E-02	0.30702E-02	0429032E-02	0+246095-02	0.19956-02	0.16370E-02	0.107925-02	0-82306E-03	0.393135-03	0.231145-03	0.33475E-04	0+18917E-05	0.729335-04	0-171416-03	0.47503F-03	0.66952E-03	0.088415E-03	0.13461E-02	0-18040E-02	0.20143E-02	0.236645-02	0.25955E-02	0-26544E-02	0476744E-02
		•	.=	0.278786.0	0-225776-04	0.225436-04	0.225516-64	0.225145-04	0.224946-04	0.224496-04	0-24126-04	0.223605-04	0.223106-04	0.222728-04	0.22195E-04	0.221155-04	0+22079E-04	0.220355-04	0-219586-04	0.21921E-04	0-218525-04	0-21821E-04	0.217416-04	0.21741E-04	0-21719E-04	0-216665-04	0-216745-04	0-216416-04	0.216595-04	= *			0.329435-02	0.32693E-02	0.32830E-02	0+32629F-02	0.323385-02	0.32162E-02	0.317565-0.3	0.315305-02	0.310415-02	0.307835-02	0.302506-02	0.29979E-02	0.294395-02	0.29174E=02	0.286645-02	0.284235-02	0.27975F-02	0+27772E-02	0.274145-02	0.272605-02	0.27010F-02	0.26915F-02	0-26767E-02	0.267555-02	*
		- 0.15	, ,	0.39576E-05	0.383805-05	0-36920E-05	0.3249376-05	0-29667E-05	0.232065-05	0-197715-05	0.163385-05	0-98903E-06	0.463765-06	0.26631E-06	0.310115-06	0.29192E-10	0.115335-06	0-25780E-06	0.491446-06	0.96682E-06	0.127286-05	0-193465-05	0-227036-05	0.290095-05	0-31766F-05	0-14146-05	0-37495E-05	0.3865E-05		0.00	12	0.158205-03	0.15698F-02	0.153376-02	0-139475-02	0.129635-02	0-10572F-02	0.78692F-03	0.65041F-03	0.39509E-03	0.283825-03	0.109795-03	0.517455-04	0.569595-06	0.821455-05	0+86864E-04	0.15475E-03	0.335916-03	0.443236-03	0.674465-03	0.79124E-03	0.10097E-03	0.11048F-02	0-12530E-02	0-13017E-02 0-13316E-02	0.134165-02	
		H	! ,	0.395765-05	0-395695-05	0.395485-05	0-395325-05	0-394926-05	0-3946BE-05	3-39440E-05	3+39379E-05	3-39345E-05 3-39309F-04	3-19272F-05	0.391946-05	0-19155E-05	0.390755-0	7.39035E-05	0.389585-05	0-38921E-05	0.368865-05	0.388215-05	0.38791F-05	0.387405-05	0.18719F-05	0.387015-05	0.36674E-05	0-38665E-05 0-38660F-05	0.386585-05	•		-	0.15820E-02	0-158156-02	0-157765-02	0.157635-02	0-19649E-02	0+15589E-02	0.15447E-32	0-15279E-02	0.15187E-02	0 1 4 9 9 2 E - 0 2	0.14890E-02	0.146816-02	0.145766-02	0.143695-02	0.14268E-02	0-14076E-02	0.13986E-02	0.138215-02	0+13748E-02	0 - 13620E-02	0-13567E-02	0.134846-02	0-13454E-02	0-134205-02	0.13416E-02	
	010		0.343526-06	0.340916-06	0.32047E-06	0.303306-06	0-25759E-06	0.230446-06	0-171705-06	0-141885-06	0.658675-06	0-61366E-07	0.230585-07	0-10409E-07	0.50305F-12	0.25469E-08	0.22727E-0;	0-39736E-01	0.850708E-01	0-111686-06	0-140516-06	0-199566-06	0.228185-06	0-279285-06	n. 30022E-06	0.329715-06	0-37576-0	0.339956-06	0.35	¹ 2		0.68655F-03	0.67078E-03	0.61015F-03	0.56721E-03	0-462735-03	0440458-03	1-28-71E-03	0.172775-03	0-12394E-03	0-818145-04	0-22179E-04	0.624896-05	0.39217E-05	0.173235-04	0.70592E-04	0.15265E-03	0.201285-03	0-752998-03	0.35921E-03	0-410505-03	0.501806-03	0.539105-03	0.59154E-03	0.609725-03		
	Ä		0-343526-06	0.343505-06	0.343465-06	0-34336E-00	0.343286-06	0.343106-06	0.343005-06	0.342765-06	0-34263E-06	0.342495-06	0-34220E-06	0.141895-06	0-341736-06	0.341425-06	0-34127E-06	0.340946-06	0.34084E-06	0 340716-06	0-340475-06	0.340376-06	0-340196-06	0-34012E-06	0.340015-06	0.33998E-06	0.339955-06		*	1,1	0.691875-01	0.69170E-03	0.690395-03	0.68926F-03	0-68609F-03	0.684075-03	0.67928F-03	0.676546-03	0.67050E-03	0.66724E-03	0-86039E-03	0.658865-03	0-64969E-03	0.642608-03	0+63914E-03	0.635785-03	0.62944E-03	0.62650F-03	0-621205-03	0.61888E-03	0.61499E-03	0.613376-03	0.611056-03	0.610316-03	0.60972E-03		
	0.05	2	0.53294E-08 0.52888E-08	0.51686F-08	0-470386-08	0-43773E-08	0+35757E-08	0.312705-08	0-22017E-08	0-17531E-06 0-13322F-06	0.95178E-09	0.623435-09	0-16083E-09	0.45036F-10	0-40276E-10	0 - 3560175 - 09	0.62184F-0"	0.449358-00	0 17492E - 08	0.21968E-04	0.31199F-08	0+15674L-0B	0 - 3/8 / 3E - 08	0-469426	0.496005-08	0.52755E-08	0.531605-08	0.30	5	7	0-26775-03	0-25969E-03	0.249755-03	0.21968E-03	0.20048E-03	0-156716-03	0.133496-03	0.87892E-04	0-479275-04	0-315696-04	0.84419F-05	0.230875-05	0.163955-07	0.70.58E-05	0.283375-04	0-43510E-04	0.804946-04	0.10115E-03	0.122415-03	0.164155-03	0.183385-03	0.215696-03	0.22781E-03	0.24218E-03	0-370****		
90-0-06	#	1.	0-532936-08	0.537916-08	0.532895-08	0-932846-08	0.53281E-08 0.53278F-08	0.532745-08	0.532695-08	0-53260E-08	0.532555-08	0-532446-08	0.53238E~08	0.532276-08	0.532216-08	0.53210t-0P	0.231996-06	0-531946-08	0.531898-08	0.531801-08	0.53176E-08	0.531695-08	0.531666-00	0.531646-08	0.531616-08	0.531605-08	80-3001cc-0	16	1,	0.267835-03	0+26778E-03	0.267605-03	0.26708E-03	0.26675-03	0+26960E-03	0.264946-03	0-26345-03	0.261705-03	0.26076F-03	0.258785-03	0-25776E-03	0.255695-03	0.25465E-03	0.25262F-03	0.25165E-03	0.24980F-03	0.248945-03	0-247395-03	C+24671E-01	0.24556-03	0.24510F-03	0.24416-03	0.244205-03	0.244025-03			
m • 1,	90	• 0	5.0 10.0	15.0	25.0	35.0	0.04	44.00 0.00	55.0	0 0 0	70.0	0.0	0.45	95.0	100.0	110.0	11:	125.0	0.061	0.65	245.0	0.061	160.0	165.0	175.0	180.0			0	0.	0.01	15.0	20.0	30.0	0.04	0.5	0.00	0.09	70.0	75.0	0.00	0.00	100.00	105.0	115.0	120.0	130.0	135.0	145.0	150.0	160.0	165.0	175.0	180.0			

		Į.	;	3+80221E-01	0.795875-01	0-744535-01	0.69467E-01	0.59655E-01	04532805-01	0.39693E-01	0.32904E-01	0.26389E+01	0-703316-01	0.101935-01	0.63430E-02	0+34064E-02	0.141665-02	0-31-316-03	0.100495-02	0.254355-02	0.47765E-02	0.75937E-02	0-144025-01	0.18317F-01	0.22235-01	0.260855-01	0.29789E-01	0.33229E-01	0-389816-01	0.410645-01	0.42619E-01	0-43567F-01	0.4438875-01	1.00		61	2 0000	00-37606-00	0.416586-00	0-399855-00	04377346-00	0.318435-00	0.284535-00	0.248865-00	0.177415	0.143755-00	0.112685-00	0.84906E-01	0.411685-01	0.257065-01	0.145635-01	0.451575-02	0.496895-02	0.85109E-02	0.228965-01	0+32725E-01	0.436425-01	0.668936-01	0.78390E-01	0.893275-01	0.993956-01	0.11592E-00	00-122015-00	0.12644E-00	0+13004E-00
		•		0.80221E-01	0.79873E-01	0.788485-01	0.786965-01	0 761625-01	0.75002E-01	0.73733E-01	0.7287E-01	0.69410F=01	0.67850E-01	0.66255E-01	0.646416-01	0.61412F-01	0.59824F-01	0.58270E-01	0.56761E-01	0.533068-01	0.525945-01	0.513515-01	0.50192E-01	0.49121E-01	0.481435-01	0-472605-01	0-457975-01	0.49213F-01	0.447375-01	0-44365E-01	0.441005-01	04438875-01	10-1100	*		ι,	0.430375-00	0-429505-00	0-42691E-00	00-1675-00	0040936F-00	0.40059E-00	0.39056F-00	00-367636-00	00-35946640	0.34130E-00	0 32 7555-00	0.299525-00	0.28555E-00	0.271795-00	0.245385-00	0.23291E-00	0.221045-00	0+19932F=00	0.18953E-00	0.180516-00	0.164745-00	0.15801E-00	0-152035-00	0 1 1 2 3 0 F = 0 0	0-138526-00	0.13545E-00	0 131395-00	0+13038E-00	0.13004E-00
		0.70	7	0.519846-01	0.507625-01	0.46081E-01	0.389916-01	0.348265-01	0.304275-01	0.21480F-01	0-17203E-01	0.13221E-01	0.454245-02	0.402925-02	0.211085-02	0.82683E-03	0.182235-03	0-1919190	0-184546-05	0+34369E-02	0.543391-02	0.77547E-02	0.136186	0 - 1 5 0 2 0 2 0 1	0-184246-01	0.21157F-01	0.236025-01	0.25794E-01	7.27671E-01	10-3036201	0.30973E-01	0.31201F-01	•	5	3	2,	0.324426-00	0.371815-00	04301555-00	0.28466F-00	0.26405E-00	0-3840486-00	0.187825-00	0.160506-00	0.133645-00	0.843365-00	n+63124F-01	0.448176-01	04179495-01	0.956925-02	0.44820F-02	0.330125-02	0.672925-02	0.122376-01	0.194865-01	0.375775-01	0.476246-01	0.678415-01	0.774635-01	0.86297E-01	0.941545-01	04106215-00	04110125-00	0-11251E-06	00-118811-0
		×	1,000,000	0+52947E-01	0.519565-01	0.516195-01	0.506815-01	0.500926-01	0.487065-01	0.47925E-01	0.47096E-01	0.453285-01	0+4409F=01	0.43474E-01	0.42535E-01	0.41598E-01	0-147676	0.388765-01	0.380195-01	0.37197E-01	0.36415F-01	0.350/85-03	0+34348F~01	0.337635-01	0.332346-01	0.32763E-01	0.32352E-01	0.320026-01	0.314905-01	0.313305-01	0.312335-01	0.312016-01				0.134437	0.37383500	0.32209E-00	0.31921F-00	0.315246-00	0.304325-00	0.297535-00	0.28998E-00	0.27045-00	0.26391E-00	0.254455-00	0.244805-00	0.223355-00	0.21573E-00	0.206305-00	0.188305-00	0.179A5F-00	0.1718ZE-00	0-157195-00	0.15063E-00	0.14460F-00	0 1 3 4 1 4 5 - 00	0 12972E-00	0+12584E-00	0.122495-00	0+117376-00	0.115595-00	0.114325-00	0.11331F-00	1
	0.65	1	7. 0.33018E-01	0.32759E-01 0.31991E-01	0.307416-01	0.269765-01	0.219635-01	0-191885-01	0.16349E-01	0.108265-01	0.03029E-02	0.603416-02	0.407715-02	20-308/-02	0.469265-03	n.83978E-04	0.105575-03	0.31440E-03	20-31001-0	0+371615-02	0.526645-02	0.70175E-02	0.050016-02	0-10/225-01	04143715-01	04160345-01	0.175265-01	0.18806E-01	0.198395-01	042105975-01	0421215F-01		0.50		<u>.</u> 9	0.23854E-00	0.236635-00	0-386067*0	0.221805-00	0.194316-00	1.17699E-00	0+15608E-00	0 1 1 80 3 F - 00	0.981525-01	0.79162E-01	0+457925-01	0+32176F-01	0.20955E-01	0.6122648-01	0.25336E-02	0+13665E-02	0.540325-02	0+10110E-01	0.16213E-01	0.313676-01	0.39737E=01	0.48283E-01	0.56699E-01	0.72164F-01	0.78782E-01	0.844255-01	0.922705-01	0.94285E-01	0.949645-01	
	n x		0.330185-01	0.92911E-01	0.32796-01	0.32364E-01	0.31766F-01	0.314066-01	0.303835-01	0.30129E-01	0.296925-01	04291985-01	04281345-01	0.276135-01	0.270935-01	0.265766-01	0.246775-01	0.230915-01	0.246295-01	0+241885-01	0.23770E-01	0-34126-0	0.226825-01	0.223316-01	0.22111E-01	0.21876E-01	0.216756-01	0.215105-01	0.212895-01	0.212315-01	0421215E-01		×	1,	•	0.238515-00	0.237035-00	0423513F-00	0.23254E-00	0.22927E-00	0.330036-00	0.215946-00	0+21052E-00	04204745-00	0.192376-00	0+18592E-00	0 - 17939E-00	041/2846-00	0.159936-00	0.153695-00	0.141875-00	0.136285-00	0+131036-00	0.1215105-00	0.11727E-00	04113395-00	04109885+00	0.10397F-00	0.10157E-00	0.99546E-01	0.966105-01	0.95695E-01	0.951476-01	0+74764E=01	
	09.0	7.	0.19997E-01 0.19840E-01	0.193776-01	0-17601E-01	0.149028-01	0+139136-01	0.990816-02	0.619775-02	0.50105-02	0436446-02	0.24420E-02	0.14695E-02	0.737246-03	0.346305-03	0-69590	0.34469E-03	0.84537E-03	0 154 /2E-02	70-36174740	044519540	0.5735E-02	0.69477E-02	0.81506E-02	0.93087E-02	0-10-2005-01	0.121895-01	0.128616-01	0.13354E-01	0-136556-01	0.137565-01	0.85		27	00-10705-00	00-36446-00	0+16340E-00	0.158855-00	0.15001E-00	04126025-00	0-113276-00	n.99009E=01	0 - 84 500E - 01	0.564895-01	0.43779E-01	0.32376E-0)	0.14435F-01	0.819495-02	0.38572E-02	0.752965-03	0-17794E-02	0.43135E-02	0.130775-01	0.18846-01	0.25215F-01	0.319476-01	0.45572E-01	0.520435-01	0.56033E-01	0.67943E=01	0.71614E-01	0.742995-01	0.764845-01	•	
	H	-	0-19997E-01 0-19983E-01	0-198736-01	0.19779E-01 0.19659F-01	0.195166-01	0-101635-01	0-189986-01	0.18736[-0]	0+162505-01	0.179925-01	0.17726E-01	0-17-01	0 169065-01	0.166335-01	0-163645-01	0.161015-01	0-15-00-01	0+15362F-01	0-151386-01	0.14928E-01	0.147326-01	0-142526-01	0 1 4 2 4 3 5 5 1	0.141167-01	0+1+007F+01	0-139178-01	0-130476-01	0.137976-01	0-1375/610	10-30	- x		-	0.170795-00	0-170546-00	0+16987F-00	0.166965.00	0+16+915-00	0.162445-00	0-159595-00	0.142045-00	0-149261-00	0.14536E-00	0-141305-00	0.132895-00	0.12863E-00	0.12438F-00	0+116085-00	0+11208E-00	0 10 8225-00	0.101025-00	0.977136-01	0-946216-01	0.891226-01	0.86730E-01	0.845655-01	0-855885-01	0-796505-01	0.78510F-01	0.776235-01	0.76611E-01	0.764845-01		
0.55			0.114905-01	0-107875-01	0.94723E-02	0.771675-02	0.674105-02	0.074265-02	0-37923E-02	0.28990E-02	0.209475-02	0.635785-03	0.411616-03	0.136846-03	0.146415-04	0.21895-04	0.52900E-03	0.96053E~03	0-14966E-02	0.211785-02	0.352895-03	0.427335-02	0.50129E-02	0.57257E-02	0.639058-02	0-918969*U	20-3010670	0.82214F-02	0.84076E-02	-	0.80		77	• !	0.11885F-00	0+11790F-00	0.110576-00	0.1044E-00	0.969345-01	0.0003175-01	0.68934F-01	0.58800E-01	0.487885-01	04302795-01	0.222775-01	0-153705-01	0.5349815-62	0.236585-02	0.741185-03	0.134265-03	0.336785-02	0.636265-02	0-10167F-01	0.195116-01	C+24692E-01	0.299765-01	0.40194F-01	0+440295-01	0.48972E-01	0.553645-01	0.57451E-01	0.58722F-01	0-30576-01		
= #		0.11580E-01	0-11973E-01	0.119205-01	0.114166-01	0.11264E-01	0.110725-01	0-10964E-01	0.10847E-01	0-10-256-01	0+104465-01	0.103326-01	0-101965-01	0.100395-01	00978895-02	0.965705-02	0.95286F-02	20-36-04-00	0.917275-02	0.90664F=02	0.896735-02	0.88761E-02	0.6315-02	0.841400	20=3485900 0-84884500	0.035265-03	0.851675-02	0.849105-02	0.847555-02	0.04/035-02	- ×		-	0.118855-00	0-118705-00	0.11626F-00	0-117535-00	00-110325-00	0-113735-00	0-111986-00	0.110035-00	0.1078455-00	0.103185-00	0.10066F-00	0.98036E-01	0.92732E-01	0.90058E-01	0.474075-01	0.822585-01	0. 79796E-01	0.774316-01	0.730425-01	0.71040E-01	0.69178E-01	0-01-01-01	0.64494E-01	0.63Z48E-01	0.62164E-01	0-604926-01	0-35066-01	0.594865-01	0-391506-0	•		
	Q	•	00.5	200	30.0	0.0	0	0	0.00	69.0	70.0	0.0		0.06	95.0	10000	110.0	113.0	120.0	129.0	130.0	140.0	145.0	150.0	155.0	160.0	165.0	176.0	180.0	:		9	5	• 0	0	0.01	20.0	25.0	30.0	000	0.54	2000	55.0	0.00	70.07	75.0	0.00	0.06	95.0	100.0	110.0	115.0	120.0	130.0	135.0	1,000	130.0	155.0	160.0	170.0	175.0	180.0			



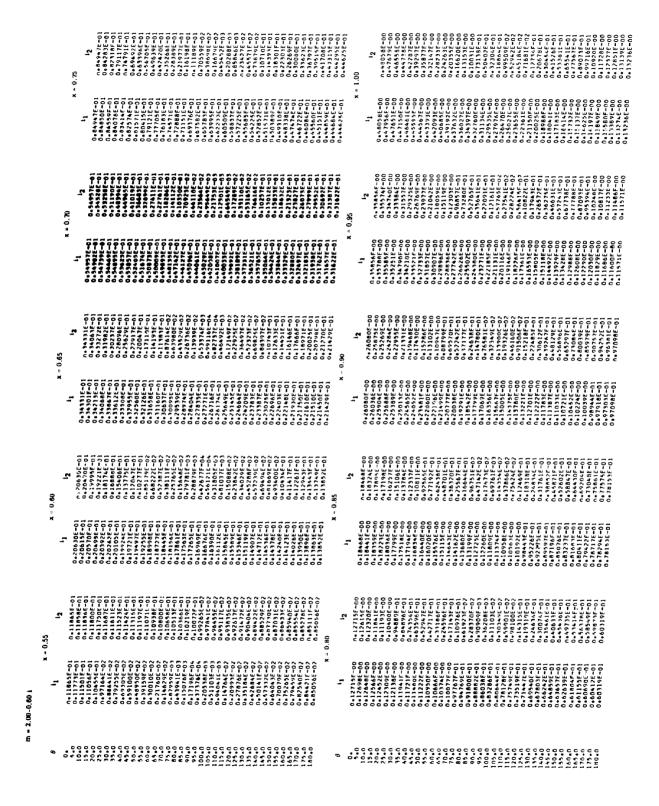
			,	G.																										0.840655-04		nc.n	75	0-65686E-02	04653785-02	0.61400E-02	0.53946.02	0-49216E-02	0.384545	0.327696-02	0.21652E-02	0.165495-02	0.797895-03	0447406E-03	0 - 74 5 70E = 04	0.634815-05	0-130905-03	0.570855-03	0.88864E-03	0.166405-02	0.2539595-02	04298045-02	0.340515-02	0.416135-02	0.471785-02	0.50124F-02	0.505015.03
			•		-	0.696705-04	0-89836E-04	0.897025-04	0-89602E-04	0+893425-04	0+89184E-04	0.888195-04	0.886146-04	0.481695-04	0.879925-04	0+674395-04	0.87187E-04	0.866805-04	0.864295-04	0.85361036-04	0-857116-04	0.854895-04	0 85080E-04	0.84897E=04	0.845785-04	0.84445	0-842365-04	0.84162E-04	0.840765-04	0-840655-04				0.658885-02	0469758E-02	0.65378-02	0.65089E-02	0.647485-02	0.63908E-02	0.62885-02	0.62317F-02	0.610925-02	0.604465-02	0.59116E-02	0.58441E-02	0.571025-02	0.5564485-02	0.551915-02	0.940335-02	0+53500E-02	0.52546E-02	0.521295-02	0.51427E-52	0.50916F-02	0.507355-02	0.505276-02	0.505015-02
			0.30	<u>.</u> r						0.154846-04																					0.45	<u>,,</u>	0.339925-02	0.337296-02	0.316625-02	0.29963E-02	04254065-02	0-227136-02	0.169175-02	0413991E-02	0.05235E-03	0.61409E-03	0+24087E-03	0.356235-03	0.21914E-05	0.131885-04	00173625-03	0.464165-03	0.683405-03	0-119716-02	0.15772E-02	0 = 18472E-02	0.22575E-02	0.242475-02	0426596E-02	0.272046-02	4
			×		0.230995-04	0-210916-04	0.230825-04	0.23033E-04	0.230335-04	0.22984E=04	00229986-00	0.22889E-0	0.228335-04	0.22776E-04	0.227355-04	0.226916-04	0.226095-04	0.22525E-04	0-22484E-04	0.22405F=04	0-223685-04	04223325-04	0-222685-04	0.222405-04	0-221925-04	0.22173E-04	0-221445-04	0.22135E-04 0.22130F-04	0.221285-04		# #	-	0.33992E-02	0.33978E-02	0-33670E-02	0.336565-02	0.335135-02	0.331605-02	0.329536-02	0.3248E-02	0.372346-02	0+31695E-02	0-311285-02	0-308416-02	0.30267E=02	0.29986E-02	0.294465-02	0.29190E-02	0.287155-02	0.28500F-02 0.28301F-03	0.28120E-02	0427814E-02	0.27692E-02	0-27912E-02	0427455E-02 0427421F-02	0.2 7410E-07	
		= 0.15	ζ,							0.202146-05																				0.40	Į,	0.14141.0	0.161595-02	0-15787F-02	0-143595-02	0-13348E-02	0.100895-02	0.951876-03	0.67059E-03	0.53487E-03	0+29329E-03	0 0 1 1 3 9 1 E - D 3	0.539435-04	0.646035-06	0.811936-05	0.881416-04	0.157506-03	0-342861-03	0.56958E-03	0.689685-03	0.92505E-03	0410333E-02	0.121485-02	0.12828E-02	0-13634E-02	0-137376-62	
			1,1	0-404558-05	0.404486-05	0.40439E=05	0.40404040	0.40390E-05	0-403416-05	0.40312E-05	0-402476-05	0-402116-05	0.401355-05	0-400946-05	0.40031.05	0-399696-03	0.399276-05	0.39844E-05	0.398046-05	0.397285-05	0.396925-05	0-39628E-05	0.396005-05	0.3955746-05	0-19532E-05	0.395165-05	0-394955-05	0.39489F-05		Ħ	-	0+16284E-02	0-162795-02	0+152385-02	0+16202E-02	0 16102E-02	0.160395-02	0-15886E-02	0 1 5 7 1 0 E - 0 2	0-156136-02	0+154065-02	0a15298E=02	0-190775-02	0.149695-02	0.14746F-02	0.14639E-02	0-144356-02	0 1 4 2 5 0 F - 0 2	0-141655-02	0 1 4 0 1 6 E - 0 2	0.139525-02	0-1304R-02	0.136065-02	0-13755E-02	0 13737E-02	1	
	# D 10	•		0.350965-06																									0.35	-	,	0.704776-03	0.68910E-03	0.662696-03	0.58281E-03	0.531865-03	0-41576E-03	0.292845-03	0.23347E-03	0-12770E-03	0.843945-04	0-53006-04	0.655005-05	0-39092E-05	0-175246-04	0.71697E-04	0.110775-03	0.205575-03	0431294E-03	0.367245-03	0-468936-03	0.55146F-03	0.582416-03	0.61910E-03	7+62380E-03		
	H		0.350965-04	0.35095E-06	0.350908-0	0-350855-06	0-350715-06	0.35062E=06	0+350416-06	0435029E-06	0.35002E-06	0.349875-06	0.34956E-06	0.349405-06	0.349086-0	0.348916-06	0.3485.E-0¢	0.348435-06	0.346146-06	G-34800E-06	0.347875-06	0+3+76+E-06	0.347546-06	0 - 347376 - 06	0.34731E-06	0.347235-06	0.347216-06	90-3071-00	*	ι,	0.710715-01	0.71056F=03	0.710045-03	0.70797E-03	0.704615-03	0.70248F-03	0.697405-03	0.69450E-03	0.686105-03	0.68465E-03	0.677405-03	0.67366E-03	0.566085-03	0.65305-03	0+654916-03	0.647936-03	0.64465E-03	0.638645-03	0.639946-03	0+63127E-03	0.627665-03	0. 62628E-03	0.624436-03	0.62396E-03	0.023805-03		
	= 0.05		0.54437E-08																									0.30		-	2	99		2.2	8	6 9	÷:	3	66		28	9		3	38	Š,	:	6		5 6	6	0.232915-03	Š				
1,96-0.66 +	.	1.	0.5:4376-08	0.544365-0	0-54433E-08	0.344305-08	0.544246-08	0.544205-08	0.54411E-08	0.544065-08	0.543956-08	0.543895-08	0.543776-08	0.543445-00	0-543576-08	0.543516-08	0.543385-08	0.54325-08	0.543216-08	0.543165-08	0.54307F±0H	0.54303E-0B	0-54296-08	0-54294E-08	0.542926-08	0-242901-08	0.54290£-08	K		0.37449	0.274628-03	0-274476-03	0.27388E-03	0.27344E-03	0-272316-03	0.27162E-03	0.27002E-03	0.269135-03	0.267195-03	0-265165-03	0.26402E-03	0.261035-03	0.2456074E=03	0.25696=03	0.29756E-03	0.255616-03	0.254705-03	0-29306E-03	0.252346-03	0.251136-03	0.250235-03	0.249925-03	0+24955E-03	0.24950E-03			
E .	93	• •	0.0	15.0	20.0	30.0	0.04	0	0.00	0.09	70.0	75.0	0 0	0.06	100.0	105.0	110.0	120.0	125.0	135.0	140.0	0.000	155.0	160.0	170.0	175.0	?		θ	•0	0.0	13.0	20.0	10.0	0.66	0.6	0	0.00	0.07	75.0	000	000	100.0	105.0	119.0	120.0	130.0	140.0	149.0	199.0	160.0	170.0	18000				

			= 0.75		0.884569E-01	0.839046-01	0 - 78734E-01	0.69105E-01	0.563276-01	0.49261E-01	0.420535-01	0.280595-01	0-159305-01	0.109615-01	0.37390E-02	0.159285-02	0.249376-03	0.96877E-03	04479595-02	0.76895E-02	0.148125-01	0.18775E-01	0.266386-01	0.306895-01	0.374715-01	0.40214E-01	0.44046E-01	0.45034E-01	10-320/E-01	1.00		2	0-45897F-00	0.44804E-00	0.40639E-00	C+37727E-00	0-30771E-00	0+23122E-00	0.193416-00	0.12398E-00	0.681196-01	0.46564F-01	0.171956-01	0.91255E-02 0.52606F-02	0.516095-02	0.14385F-01	0.22670	0.43919E-01	0.558565-01	n.80034E-01	0.91479E-01	0-111436-00	0.12583E-00	0.130505-00	0-134305-00
			×		0-849698-01	0.84192E-01	0-83726E-01	0.82267E-01	0-80173E-01	0.775436-01	0.760665-01	0.728675-01	0.71180E-01	0-57713E-01	0.65965E-01	0.625145-01	0.60839E-01	0-57646-01	0.361455-01	0.53386F-01	0.521395-01	0.499365-01	0.46988F-01	0.47438-01	0.467905-01	0.456805-01	0 4 5 5 9 5 E = 0 1	0.45367E-01		*	Ι,	0.462675-00	0.46171F-CO	0-454156-00	0.439445-00	0-13862-0	0. * 0. # 0. # 0. 0. 0. 0. 0. 0. 0. 0. 0. 0. 0. 0. 0.	0.393416-00	0 - 364 70E - 00	0.334275-00	0+31887E-00	0.288505-00	0.27381E-00	0.24599E-00	0.22082F-00	0-209376-00	0.18892E-00	0.17994E-00	0.16451E-00	0.158046-00	0.14752E-00	0-14345E-00 0-14013E-00	0-137576-00	00134665-00	0.134305-00
			- 0.70		0.54.570E-01	0.532936-01	0.48404E-01	0-409956-01	0.32039E-01	0.273365-01	0.181906-01	0.140125-01	0. 70001E-02	04433186402	0.92067E-03	0.211976-03	0.71267E-03	0.183795-02	0.55153E-02	0.1966E-02	0-13948E-01	04162126-01	0.217785-01	0424316E-01	0.285405-01	0.301136-01	0.319706-01	0.32208E-01	0.95	-	ŗ,	0-347696-00	0.336755-00	0-32346E-00	0.283705-00	00-388695-00	0.20272F-00	0-144996-00	0.92246F-01	0.694505-01	0+33324E-01	04204645-01	0-53605E-02	0.284566-02	0-653725-02	0.193675-01	0.281655-01	0+48404E-01	0.590416-01	0.7951E-01	0.868125-01	0.104095-00	0.113865-00	0.11637E-00 0.11721E-00	
			× -1	0.55001 6-01	0.549475-01	0-345255-01	0.537606-01	0.5251496-01	0.91800E-01	0.30177E-01	0.492846-01	C+47380E-01	0.463895-01	0-443725-01	0.423645-01	0+41390E-01	0.395176-01	0.386395-01	0.377966-01	0+36264E-01	3.160416.01	0-3696960	0.338805-01	0.339655-01	0.32798E-01	0.323456-01	0.32262E-01	10-10-1	*	-	0.143404	0.347055-00	0+34513E-00	0.337615-00	0.332135-00	0.31815F-00	0.300685-00	0.29131E-00	0.2 7092E-00	0.249705-00	0.23906E-00	0+21825E-00	0.20825E-00 0.19861E-00	0 18940F-00	0.172435-00	0-16474E-00	0-15167E-00	0-145116-00	0-134996-00	0.127135-00	0.124065-00	0.121605-00	0.110316-00	0.11721E-00	
		- 0.65		0.345176-01	0-33448	0.303076-01	0.2576-01 0.25740E-01	0-23005E-01	9-17152E-01	0-142146-01	0.07502E-02	0.432485-02	0.26443E-02	0.51753E-03	0.96998E-04	0.50459E-03	0.128136-62	0.377785-02	0.53940E-02	0.907025-02	0-11007E-01	04147785-01	0.164995-01	0.180445-01	0.204395-01	0.212246-01	0.21865E-01	0.0	ì	<u>.</u>	0.25471E-00	00252695-00	0+23701E-00	0.223925-00	0+189595-00	0-148486-00	0+12705E-00	C. 85691E-01	0.500355-01	0.23303F-01	0413652E-01	0.303625-02	0+19048E-02	0.524345-02	0.99757E-02	0.235945-01	0.318135-01	0.49429E-01	0.582135-01	0.74394E-01	0.872526-01	0.920126-01	0.976105-01	0.98323E-01	
		*	. 7	0.34.317E-01 0.34469E-01	0.344035-01	0.34046-01	0.33516E-01	0.327846-01	0.323595-01	0-31412E-01	0.309016-01	0+29825E-01	0+207126-01	0.28154E-01	0.270575-01	0.26525E-01	0.255156-01	0.250436-01	04241785-01	0.237895-01	0.23433E-01	0-228235-01	0.22571E-01	0422181E-01	0.220436-01	0.21885E-01	0.21865E-01	**		-• :	0.254715-00	0+25304E-00	0.24815-00	00-24455-00	0.235405-00	0.22996E-00	0.21772E-00	0.20419F-00	0+197145-00	0-18287E-00	04179785-00	0.16200E-00	0.155416-00	0-14306E-00	0+13201E-00	0.12703E-00	0-118235-00	0-11443E-00	0-108045-00	00-103455-00	0-10148E-00	0.991126-00	0.98323E-01		
		8.0	12	0.206635-01	0-19399E-01	0-17038E-01	0.15538E-01 0.13887E-01	0-121396-01	0-85707E-02	0.525865-02	0.381855-02	0-159745-02	0-78873E-03	0-419876-04	0.0403525-04	0-3956490	0.24405	0.35057E-02	0-465575-02	0.712505-02	0.03661E-02	0.106765-02	0.116765-01	0.12536E-01	0-13740E-01	0.140516-01	0.85	3	2 ,	0.181615-00	04160185-00	0-175935-00	0415971F-00	0-135235-00	0-120916-00	0.904885-01	0.75315E-01	0.47272F-01	0.246056-01	0.15912E-01 0.491677E-02	0.44218E-02	0.165735-02	0+17038E-02	0.809735-02	0-131395-01	0.25679E-01	0.376685-01	0.46860E-01	0.59875E-01	0.65469E-01	0.740935-61	0.76906E-01	0.79197E-01		
	i	# -	0.208296-01	0-20116-01	0.205926-01	0.204646-01	0-20132E-01	0.19712E-01	0e19474E-61 0e19221F-01	0-189555-01	0.18678E-01	0.18104F-01	0.17517E-01	0.17226E-01	0.16656E-01	0.163835-01	0-150676-01	0-196285-01	0-151965-01	0+19004E-01	0.146795-01	0.145395-01	0.144235-01	0.142526-01	0.141996-01	0+1+156F-01	*		-	0-18161F-00	0-181395-00	0.17926E-00	0+17520F-00	0-17250E-00	0.165945-00	0.162175-00	0.153885-00	0.144926-00	0-140916-00	0-13104F-00	0.122035-00	0+11769E-00	0.109516-00	0-10570E-00	0.987776-01	0.956776-01	0.90248E-01	0.850856-01	0.04110E-01	0.813806-01	0+80424E-01	0+7934E-01	04 /9197E-01		
	0.55	ζ,	0.12018E-01	0-116485-01	0-10507E-01	0-89718E-02	0.70095E-02	0.59743E-02	0-395255-02	0+21903E=02	0.146925-02	04437095-03	0.148145-03	0-422446-04	0.217825-03	0-97456E-03	0.152415-02	0-28658E-02	0+36127E-02	0.51403E-02	0-38746-02	0.717525-02	D+77044E-02	0.844655-02	0.86384E-02	70-31ED/9*	0.80	-	2.1380.0	0.123845-00	0.12191E-00	0.110696-00	0.10279E-00	0.837895-01	0.733116-01	0.520586-01	0.419185-01	0.23997E-01	0.105985-01	0.26861F-02	0-870475-03	0.12857E-03	0.6360936-52	0.102615-01	0.148435-01	0-25290E-01	0.307825-01	0-414226-01	0.50579E-01	0.54276E-01	چة	0.607655-01	•		
1,96-0,66 1	# #	. -	0.120105-01	0.11990F-01	0.119055-01	0411768E-01	0.115835-01	0.11360E-01	0.11236F-01	0.109696-01	0.106856-01	0-105406-01	0 10250F-01	0.10107E-01	0.982936-02	0.947045-02	0-945036-02	0.933715-02	20-394616-0	0.904645-02	0.889835-02	0.883938-02	0 - 87574F=02	0.87251E-02	0.870315-02		×	<u>-</u> -	0+12584E-00	0412568E-00	0+124405-00	0.12331F-00	0.12027E-00	0+11637E-00	0.113936-00	0.11445-00	0-106076-00	0-100365-00	0.97479E-01	0.917116-01	0.66138F-01	0.83476E-01	0.78482E-01	0.76179E-01	0-72010E-01	0.701615-01	0.66962E-01	0-644945-01	0.634665-01	0+62024E-01	0-615736-01	0-612136-01			
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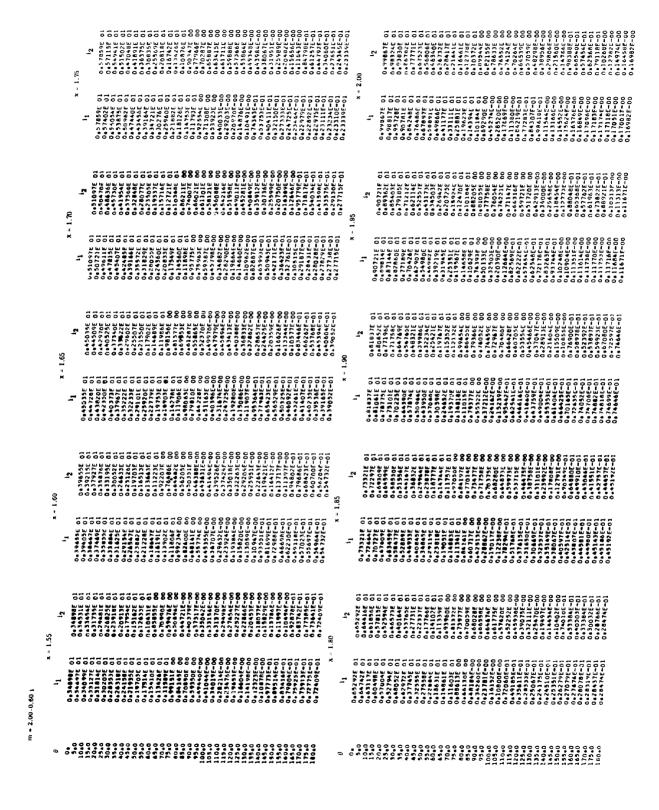
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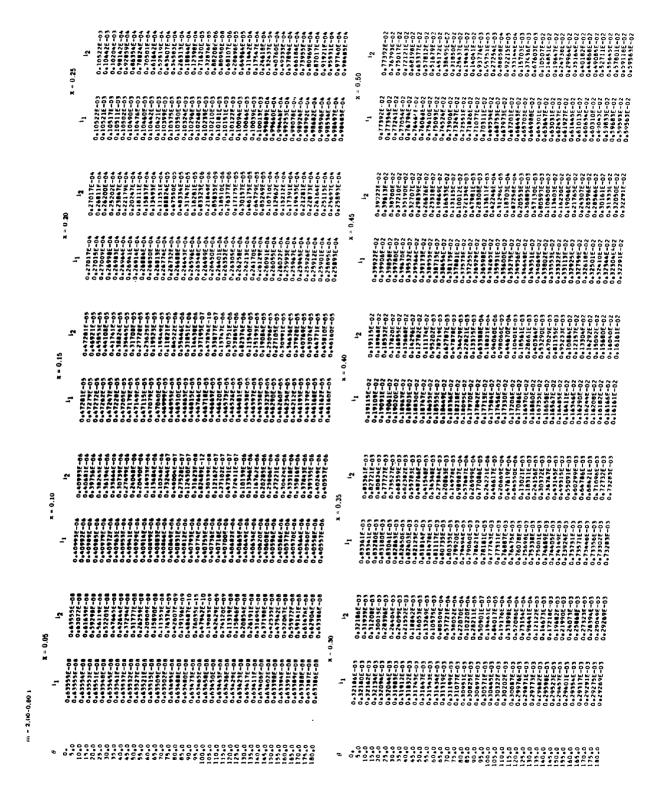
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	x - 1.70	<u>.</u> 7	0.48777E 01	004616*E 01	0.43534F 01	0.3558nF 01	0.30986E 01	0.219145 01	0=17881E 01	0411516E 01	0.92748E 00	0.649846 00	043/408E 00	0+51304F 00	0.49893E-00	0.469245-00	0+449335-00	0.374405-00	A 33010E-00	0.23616E-00	0.19165F-00	0.117085-00	0.668725-01	0.50603F-01	0.31992718-01	0.27976E-01	10-110/070	x - 1.95	19	* *************************************	0.85200E 01	0.81365E 01	0.67628E 01	0.58832E 01	0.405:9E 01	0.24967E 01	0-19025E 01	0.11200E 01	0.90679F 00 0.78231E 00	0.71994E 00	G.68330E 00	0.6461EE OD	0.50460F 00	0.47429E-00	0.39483E-00	0-239316-00	04124675-00	0.903276-01	0.66812E-01	0.72653E-01 0.85287E-01	0.1005PE-00	0-124526-00	0-128035-00	
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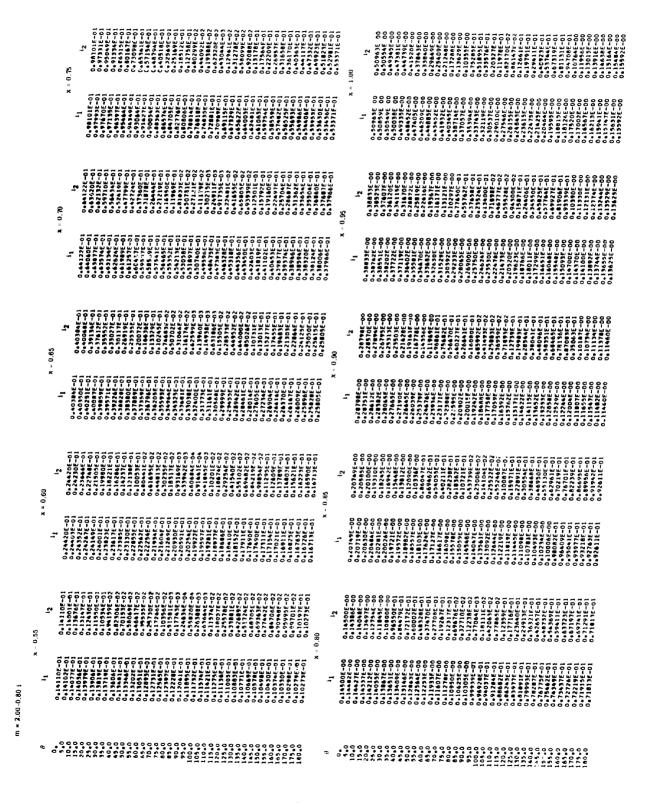
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E .	0	•0	5.0	15.0	75.0	30.0	0.04	30.0	0.44	0.00	70.0	80.0	9.0	0.00	100.0	105.0	115.0	120.0	125.0	135.0	140.0	0.041	155.0	0.041	170.0	175.0	0.001		•		• • •	10.0	0.0	23.0	30.0	0.04	45.0	95.0	0.09	70.0	75.0	9.00	90.0	100.0	105.0	0.01	120.0	130-0	135.0	140.0	150.0	160.0	165.0	175.0	180.0			



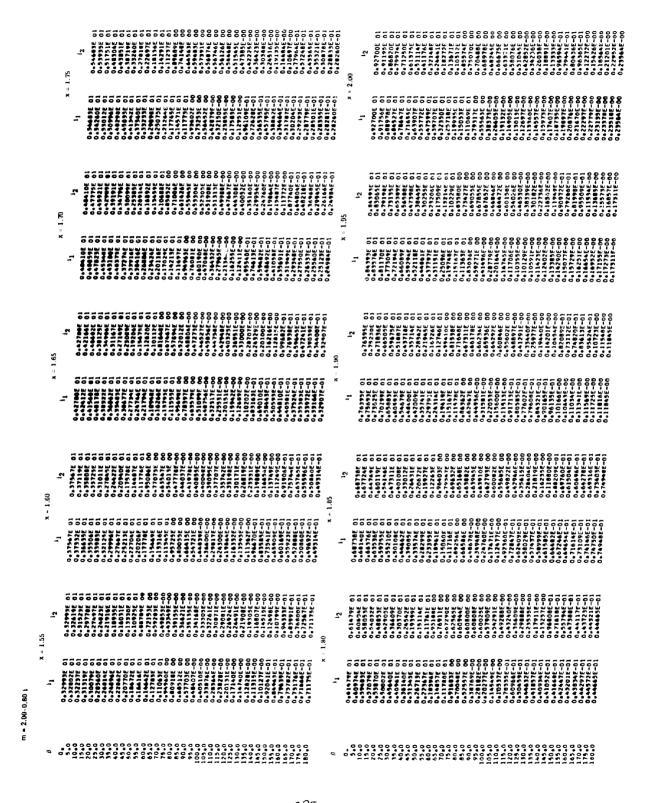
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		0. 79505E 00	0.77810E 00	0.10015E 01	0.99628E 00	0.12298F 01	0.122985 01	0+14778E 01	2,
		0.78476E 00	0.74743E 00	0.97869E 00	0.97242E 00	0 12144E 01	0+12198E 01	0.14727E 01	0414655 01
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		0.66201E 00	0.472105-00	0.87395E 00	00.00000	0.10958F 01	0.91233E 00	04130455 01	0.11980f 01
		0003382E 00	04406495-00	0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0	0.50301E 00	0.10084E 01	0.71001E 00	0.12904E 01	0.976405 00
		0.59329E 00	0.281705-00	0.75992E 00	0-428486-00	0.90515F 00	0+61981E 00	0.112425 01	0.657928 00
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		0.498695-00	04130925-00	0.63422E 00	0+22261E-00	0479608E 00	0435455-00	0.91740E 00	0.51900E O
		0.46771E-00	0.94084E-01	0.551665 00	0.16895E-00	0.68767E 00	0.276635-00	0.84837E 00	0.33567E-0
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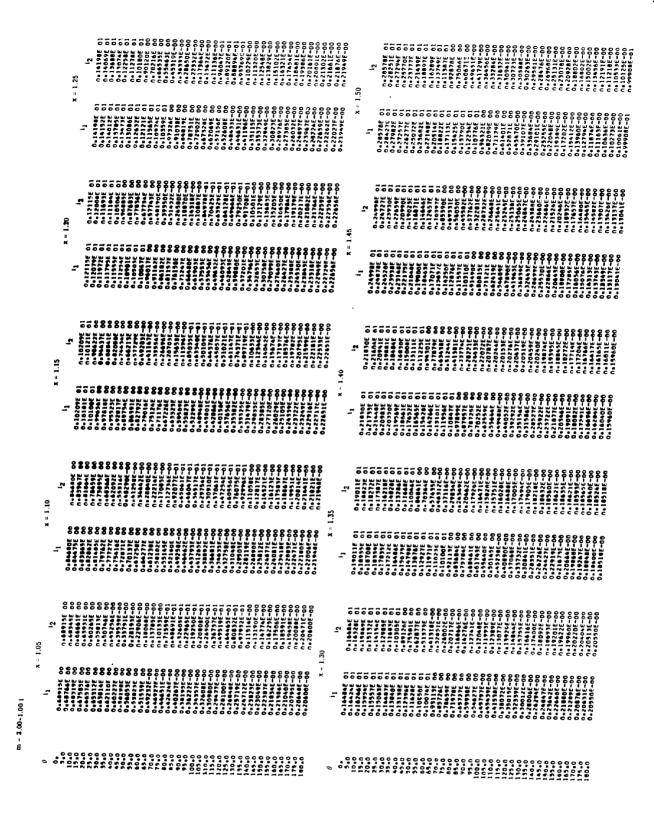


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		×		0.142075 01	0.14017E 01	0.13464E 01	0-13066E UI	0-120795 01	0-11503F 01	0-102426 01	0.96198E (1)	0.832556 00	0.76989E 00	0.651855 00	0.59749E 00	00460715-00	0456625-00	0.41741E-00	00-320286-0	0.32219E-00	0.29737E-00	0.27686F-0n	0.24072E-00	0.227035-00	04206245	041981E-00	0-193165-00	0.186855-00	0+18608E-00		*	1,	0.2891AF 01	0.28783£ 01	0.289276 01	0-205000	0.25346E 01	0.22950E 01	0.20675E 01	0-171916 01	0+1546ZE 01	0-121945 01	0.10704E 01	0.933205 00	0.697085 00	0-59845E 00	0-37855-0	0.374225-00	04275185-00	0-23762E-00	0.181225-00	0-160596-00	0.13063F-00	0+12011E-00	04105675-00	0-101096-00	0.979705-01	0.93541E-01	
		1.20	i2 0=12002f ci	0-11900E 01	0.111126 01	0.10461E 01	0.87#716 00	0.58493F 00	0.56709E 00	0440405-00	0.52492E-00	0.25493E-00	0.147116-00	0-109365-00	0.631895-01	0-527505-01	0.491446-01	0-574836-01	0.670926-01	0.78945E-01	0-106035-00	0-11990E-00	0-13330E-00	0.157186-00	0.16712E-00	0-17546E-00	0 1 8 6 8 8 ± 00	0.18977E-00	0+19074E-00	1.45		2	0.25336E 01	0.25076E 01	0.230936 01	0.108.75 01	0-17495E 01	0.15317E 01	0.11095E 01	0.9209BE 00	0.614255 00	0.49965E-00	0.343376-00	0.29594E-00	0+24402E-00	0-231306-00	0.21633500	0+21343E-00	0.20786E-00	0.192945-00	00-183765-00	0.16398E-00	0.15439E-00	0+13784E-00	0.131405-00	0.122816-00	0+12066E-00	0-119956-00	
		# # ·	11 0.12002E 01	0-11945E 01	0-11676E 01	0.11129E 01	0.10761E 01	0.99071E 00	0-943086 00	0.04221E 00	0.790425 00	0.68879E DO	0-63977E 00	04542686 00	0-5057&E 00	0.46642E-00	0+346416-00	0-366185-00	0-33870E-00	0.292146-00	0-372B0E-00	0.255905-00	04278776-00	0-21826F-00	04209605-00	0.197395-00	0-1936BE-00	0-191406-00		X = 1			0.25336E 01	0.24846E 01	0.24250E 01	0.22443E 01	0.2128BE 01	0.18635F 01	0.17207E 01	0.14310F 01	0-12899E 01	0-11545E 01	0-907456 00	0.598055 00	0.6099TE 00	0.53122E 00	00-40229E-00	0+35060E-00	0.269535-00	0.23832E-00	0.190835-00	0-173216-00	0 1 9 7 9 F - 00	0+13837E-00	0-131406-00	0-12270E-00	0.120436-00	*	
	1.15	1,2	7	0.963456 00	0.92571E 00 0.87247F 00	0.80801E 00	0.456075 00	0-574355 00	0+413336-00	0+398825-00	0.210816-00	0-19950E-00	0.043416-01	0-001196-01	0446386-01	0.328436-01	0.35731E-01	0-511111-01	0-66117E-01	0-804095-01	0-110446-00	0.12942E-00	0.139336-00	0+163196-00	0-172616-00	0-180115-00	0-16555-00	0+18995E-00	į	1.40	12	0.221246 01	0.21908E 01	0.21271E 01	0.18904E 01	0-17302E 01	0-13528E 01	0.11603E 01	0.83448F 00	0.68568E On	0.557276 00	0.36502E-00	3.29942E-00	0.217965-00	0-196205-00	0-18309E-00	0-17250E-00	04170146-00	0.16915E-00	0-167605-00	0.16296E-00	0.159935-00	0-192686-00	0.149675-00	0.14504E-00	0.14356E-00	0.14237E-00		
	Ħ	1,	0.99833E 00	0.987275 00	0-95036 00	0.931745 00	0.87916E 00	0.80239E 00	0.76397E 00	0.689195 00	0.64389E 00	0.3604095	0.52757E 00	0.491955-00	00-35256-00	0.395296-00	0.342145-00	0.318976-00	0-298036-00	04279255-00	0.24784E-00	04235035-00	0+2400E-00	0.206935-00	0.20073E-00	0-193616-00	0.19042E-00	0.189956-00		×	 -	0.22124E 01	0.220268 01	0.21294F 01	0.20607E 01	0.19805E 01	0.17836€ 01	0.16719E 01	0.14354F 01	0,131556 01	0.108336 01	0.97438E 00	0.176905 00	0.48971E 00	0.440716 00	0447666-00	0421216-00	0+3004E-00	0+29500E-00	0.238515-00	0.21681500	04183946-00	0-17189E-00	0.15480E-00	0.149225-00	0.143116-00	0.14237E-00		
	1.10	7	0-81536E 00 0-80867E 00	0.75693E 00	0.714005 00	0.60263E 00	0.53849E 00	00-00-20	0.237436-00	0.2207EE-00	0-170446-00	0.91545E-01	0.63957E-01	0-296206-01	0.22625E-01	0.248775-01	0.3242E-01	0-430936-01	0.306060E-01	0-858656-01	0-101376-00	0-116535-00	0-144036-00	0.19564E-00	0.15347E-00	0-179036-00	0-182505-00	0+18366t-00	1.35	5	,7	0.19234E 01	0-180345 01	0+17671E 01	0.16542E 01	0-136900 01	0.121035 01	0.89402E 00	0.74775E 00	0.498961-00	0-400995-00	0.257245-00	0.20996E-00	0 15 4 1 F - 00	0.14088E-00	0.134356-00	0.13377E-00	0.13674E-00	0-144336-00	0-14792E-00	0.153635-00	0-155746-00	0.156725-00	0-19572E-00	0.16101E-00	0-16131E-00	00-32-191-0		
	*	1.	0-81930E 00 0-81930E 00	0-79711E 00	0.76600€ 00	0.74559E 00	0.49675E 00	0.66929E 00 0.6402BF 00	0.41027E 00	0.54865 00	0.51823E 00	0-46813E-00	0-430625-00	0-403485-00	0.35420F-00	0.331685-00	0-311246-00	0-274076-00	0.259516-00	0.245575-00	04222386-00	0.212995-00	0.205015-00	0-193015-00	0-168891-0	0-185976-00	0.18424E-00 0.18366F-00		x . 1	-		0.191566 01	0.18924£ 01	0.18544F 01	0-17367£ 01	0.166406 01	0.149035 01	0.19953F 01	0.11988F 01	0.110100 01	0.91384F 01	0.82676E 00	0.44513E 00	0.60012E 00	0.5371BE 00	0.430106-00	0.38964E-00	0-312675-00	0.28339E-00	0.2370316-00	0+21913E-00	0.191995-00	0.182196-00	0-17441E-00	0-164956-00	0-142205-00	000		
1 05	Ž,	4	0-6466E 00 0-63129E 00	55	2	4 8	۳.	žΞ	2	ž		ξΞ.	¥ :	2	7		3	9	2	3 5	3	6 5	=	3	ö	2 2	8	OF.	}	ç	•	٠.	•	•	•	-0	21	0 -0	٠.	• •	•		•	~ 9	2 9	2 4	으 으	=:			•	•	•	0-17204E-00	0.17434E-00	.0.17618E-00			
108		0045214E 00	0.45046E 00 0.44624E 00	0.628986 00	0461645E 00	0.58471E 00	0.545918 00	0.524615 00	0479795-00	0-456B&F-00	0-411305-00	0-3616E-00	00-387068-00	0-32741E-00	0-308845-00	0-275236-00	0.26027E-00	0-346576-00	0-222896-00	0.21288E-00	0-194196-00	0+18981F-00	0-18432E-00	0-179885-00	0+17403F-00	0-17238E-00	0-17210E-00	x - 1		-1	0.166/05 01	0-165486 01	0-16066E 01	0.15697E 01	0.14556F 01	0-138916 01	0.19167E 01	0.11611E 01	0.108085 01	0.97179E 00	0.845395 00	0.703145 00	0.63846£ 00	0.52367E 00	0.47378E-00	0.428846-00	0.353106-00	0.32180E-00	0.270835-00	0.250541-00	0.21876E-00	0=20675E-00	0-169275-00	0+1834E-00	0+17697E-00	0.17618E-00			
D-00'7 • III	0	•	000	20.0	10.0	0,00	0.0	0.00	0	0.69	0,67	0.08	90.0	0 60	105.0	110.0	0.021	125.0	130.0	135.0	145.0	150.0	199.0	165.0	170.0	0.641	•		*		•	10.0	13.0	0000	30.0	0.0	0	2000	0.09	0.00	0.00	90.0	0.00	95.0	100.0	110.0	115.0	125.0	130.0	140.0	145.0	155.0	160.0	170.0	179.0	163.0			



m = 2.00-1.00

2.00-1.00



			= 0.25		7	0.206926-03	0.202225-03	0-18413E-03	0.15645E-03	0.14003E-03	0-10453E-03	0.65215-04	0.52594E-04	0.378865-04	0.14641E-04	0.687565-05	ne12029E-06	0.543835-05	0-12420E-04	0.33868E-04	0.47646E-04	0.029035-04	0.059248-04	0-12890E-03	0.14412E-03	0.16972E-03	0.179335-03	0.19074E=03		0.50	r,	0.16761E-01	0.16633E-01	0415636E-01	0.147995-01	0.12590E-01	0.987576-02	0.701925-02	0-564146-02	0.31863E-02	0.237776-02	0.729895-03	0.141255-03	0.180916-03	A.80260E-03	0.232995-02	0.32582E-02	0.537845-02	0.650265-02	0.971445-02	0.10625E-01	0.11409E-01	0-125098-01	0.12891E-01
			*	ı,	0.20850E-03	0+208475-03	0-200215-03	0.207985-03	0.207356-03	0.20650E=03	0420600E-03	0.20488E-03	0.204265-03	0.20294E=03	0.202255-03	0.20084E-03	0-1994ZE-03	0-17872E-03	0.19737F-03	0-19672E-03	0.19552E=03	0.194988-03	0+19401E-03	0.193605-03	0 - 19292E-03	0.192668-03	0.197315-03	0+19220E-03		×	. 	0+16761E-01	0-167258-01	0.16679E-01	0.16538E=01	0.16334F-01	0.162125-01	0-15935-01	0.156.0E-01	04154546+01	0-15113E-01	0.14940F-01 0.1476RE-01	0-145995-01	0.142716-01	0.13966F=01	0.138248-01	0+175676-01	0.13452E-01 0.13447E-01	0-13252E-01	0.13095E-01	0.129825-01	0-129426-01	0.1259136-01	0412891E-01
		į	8.	2	0.52839E-04	0.51247E-04	0.455615-04	0-634095-04	0.354775-04	0.310395-04	0-21894E-04	0.134625-04	0.954316-05	0.629265-05	0.168645-05	0.126245-06	04343425-06	0.327045-05	0.57700E-05	0-12448E-04	0-16418E-04	0.250135-04	0.293735-04	0.37565E-04	04411456-04	0.467416-04	0-49715E-04	0.50096E-04	0.45		7	0.642695-02	0.62352F-02	0.74984E-02	0.69769E-02	0.5708RE-02	0.42750E-02	0.28487F-02	0+21915E-02	0-106301-02	0.344915-03	0-13908E-03	0.722216-04	0.40798E-03	0.78492E-03	0.169215-02	0.22324E-02	0.33996E-02	1.39899E-02	0.50949E-02	0.59907E-02	0.63265E-02 0.65733E-02	0.672435-02	
			4	.	0.528345-04	0.928185-04	0-327946-04	0.52707E-04 0.52650F-04	0-52584E-04	0.524265-04	0.523365-04	0.521365-04	0.520285-04	0.51801E-04	0.5156835-04	0-31446-04	0.51255	0.51091E-04	0+308695-04	0-507645-04	0.505725-04	0.50466-04	0.503365-04	0.502745-04	0.501765-04	0.501416-04	0.5010!E-04	0.270765-04) = #	-1-	0.84916	0.84877E-02	0-84568F-02	0-843021-02	0.83561E-02	0.825696-02	0.819935-02	0.80710E-02	0 79295E-02	0.778015-02	0.77044E-02	0.755345-02	0.747946-02	0.73372E-02	0.720585-02	0.714525-02	0-703575-02	0.698746-02	0-690475-02	0.684175-02	0.68379E-02	0.67859E-02	0.67779E-02 0.67752E-02	
		0.15	ţ	0.91517F-05	0.908216-05	0.65368E-05	0.408146-05	0.68651E-05	0.537365-05	0.45805E-05	0.30186E-05	0.154495-05	0.10814E-05	0.623565-06	0.75237E-07	0.630905-09	0-260116-06	0.102865-05	0-15761E-05	0.291346-05	0-36622E-05	0.52061E-05	0.595456-05	0.72907E-05	0.783795-05	0-86089E-05	0.88766E-05		3	.	0.398905-02	0.39666-02	0.372185-02	0-32776E-02	0.299436-02	0-23486E-02	0.1664E-02	0.102356-02	0.74300E-03	0.30138E-03	0.15209E-03	00144126-04	0.256035-04	0.21368E-03	0.582435-03	0.108395-02	0.136475-02	0-19423E-02	0.24629E-02	0.27187E-02	0+30864E-02	0.32874E-02	0+33064E-02	
			1 1	0-91517E-05	0.914956-05	0.914695-05	0.913856-05	0.913296-05	0.911896-05	0.910176-05	0.909216-05	0.90711E-05	0.906005-05	0.903676-05	0.902485-05	0.900086-05	0.69989E-05	0-896598-05	0.89443E-05	0.89343E-05	0.891625-05	0.890825-05	0.88947E-05	0.888926-05	0.888125-05	0.88786E-05	0.88766E-05	0 = H			0.398755-02	0+39830E-02	0.396546-02	0.395295-02	0.39190E-02	0.387655-02	0.385236-02	0.37996E-02	0.374245-02	0.37128E-02	0.365296-02	0.36230E-02 0.35935F-02	0.356465-02	0.350936-02	0.348346-02	0-34356E-02	0-339426-02	0.336025-02	0.33461E-02	0.332426-02	0.33165E=02	0-33076E-02	0.330645-02	
	0.10		<u>.</u>	0.78754E-06	0-763005-06	0-695438-06	0.646915-06	0.52853E+06	0.393955-06	0-325635-06	0.19724E-06	0.141065-06	0.53156E-07	0.240835-07	0-119476-10	0.231116-08	0.515576-07	0-13836F-06	0+193865-06	0.320736-06	0.300235-06	0.521166-06	0.58256F-06	0.68591E-06	0.724765-06	0-770695-06	0.77550t ±06	0.35	<u>.</u>	0.11.0725.402	0+16942E-02	0.15558E-02	0.15077E-02	0.128145-02	0-114735-02	0-857685-03	0.568885-03	0.435615-03	0-210705-03	0.615365-04	0.206965-04	0-107615-04	0.941355-04	0.167195-03	0-36370E-03	0.605415-03	0.733685-03	0.986346-03	0+11027E-02 0+12077F-02	0.129825-02	0.13/136-02	0.145855-02	70-184-101	
	H			0.767528-06	0.787361-06	0-787216-06	0.78681E-06	0.78656E=06 0.78627E=06	0.785965-06	0.785235-06	0.78484E-06	0.783995-06	0.783545-06	0.782616-06	0.78156-06	0-781216-06	0.780316-06	0-179885-06	0.77907F=06	0-778705-06	0.778046-06	0-777765-06	0-17730E-06	0.777126-06	0-7768BE-06	0.77682E-06		Ħ	.	0+17072E-02	0417067E-02 0417052F-02	0-170275-02	0.169495-02	0.16897E-02	0.16767E-02	0.16691E=02 0.16609E=02	0.16522E-02	0.163336-02	0.161315-02	0 1 6027E-02	0.15819E-02	0.15716E-02 0.19613E-02	0.155176-02	0.153296-02	0.15242E-02	0-190835-02	0.149485-02	0+14890E-02	0 14 796E-02	0.147336-02	0+147126-02	0.14696E-02		
	0.05	ζ,	0.171556-07	0.11789F-07	0-11341E-07	0-9846-08	0.91169F-08	0-713366-08	0.50235-08	0-30-036-08	0.217261-08	0.142365-08	0-36765-0	0.13676-10	0-912421-10	0.364005-09	0.141546-08	0.302705-08	0.398415-08	0.605608-08	0-71078E-08	0.908495-08	0.79498E-08	0-11302E-07	0-117485-07	0-121135-07	0.30	2		647735	02820E	60436E	93219E	43915E	0.36091E-03	26924E	16444E	11855E	46436	66971E	76493E	16186E	37066E	10155	188996	23796E 28843F	336696	433465	47477E	53929E	57356E	37795E		
98-1.74 i	Ħ	.	0-121555-07	0-12155E-07	0-121546-07	0.121548-07	0-121526-07	0-12151E-07 0-12149E-07	0-121486-07	0-121456-07	0.121435-07	0-121-06-07	0.121305-07	0-121346-07	0-121336-07	2-121296-07	0-123365-07	0-121248-07	0-121216-07	0-121205-07	0-121175-07	0-12116E-07	0-121156-07	0-121146-07	0-121136-07	0-121135-07) = #	1,	0.647735-03	0.647585-03	0.64645F=03	0.64946-03	0-642705-03	0=64096E=03	0+636818-03	0.631916-03	0.6248-03	0.623545-03	0.617555-03	0.61450E-03	0.608425-03	0.60544F-03 0.60252F-03	0.599695-03	04594365-03	0.541905-03	0.587485-03	0.583816-03	0.562285-03	0.57989E-03	0.578435-03	0.57807E-03	50-37,4,50		
B = 2.9	æ	, ه	5.0	10.0	20.0	20.0	0.04	0	0.00	0.00	0.07	75.0	85.0	0.06	00001	105.0	115.0	125.0	130.0	140.0	145.0	155.0	160.0	170.0	175.0	•		θ	•	000	0.61	2000	0.0	40.0	0.00	90	0.69	0.07	0	0.00	95.0	105.0	115.0	120.0	130.0	140.0	145.0	155.0	160.0	170.0	180.0			

m = 2.98-1.74

m = 2.98-1,74

						x = 0.05								x = 0.55			
		n	a _n r		4 ⁱ n	i	P _L	ρį			ar n		ai n		b <u>r</u>	LÍ	
		1 0.225, 2 0.8800 3 ~0.434; 6 5	72E-09	2.462	00E-04	0.4688 -0.7694 -0.7694	2F -0d	0 • 1 2 2 4 6 6 • 0 • 7 6 2 5 5 £ • 0 • 3 2 4 4 3 £	-10	0.3445 0.4446 0.3515 0.1584 0.4566 =0.2168	52E=01 54E=03 58E=05 5E=07	0.76 0.614 0.284 0.104	75E-0 75E-0 75E-0 75E-0	1 0.2209 3 0.1854 5 0.8855 7 0.3285	1E-02 4E-04 1E-07 7E-09 4E-11	bin 0+11907E- 0+10901E- 0+53287E- 0+19637E- 0+21823E- 0+63448E-	07
					1	= 0.10								x = 0.60			
	3 4 5	0.91229 -0.91190	E-07	0:2726 0:1>>2 0:9403	15-10 05-09	0.42447 -0.17648 -0.10372	E-04	0 • 26686E = -0 • 22854E = -0 • 280356 =	09	0.45633 0.71913 0.64489 0.34504 0.10626 ~0.65632	E-03 E-05 E-07 E-09	0.126 0.116 0.116 0.6204 0.62529 0.1363	95E-02 2E-04 7E-07 9E-09	0.34314 0.34097 0.19340 0.77467 -0.80912 -0.92426	E=04 E=06 E=09 E=12	0.178665=0 0.197925=0 0.197925=0 0.115795=0 0.520145=0 0.417935=1 0.8852625=1	اد اد اد
	1	3.614365				- 0.15							,	t = 0.65			
	3 4 9 6	0.61435E 0.69930E 0.53798E	-06	0.91915 0.11808 0.66713	E-05	0+32453E 0+21828E 0+11375E	-08	0.20092E=0 0.90911E=0 =0.76454E=1;	7	0.59102F 0.107558 0.11261E 0.70599E 0.27273E -0.15548E	-02 -04 -07 -09	0 = 66690 0 = 17338 0 = 19577 0 = 12676 0 = 53488 0 = 32512	E-02 E-04 E-06 E-09	0.51482E 0.59709E 0.39635E 0.17323E -0.31941E -0.92690E	-04 -06 -08	0.25747F-02 0.34175E-04 0.23574E-06 0.11085E-08 0.36081E-10 0.63279E-11	•
					x =	0.20							x	- 0.70			
	3 4 5 6	0:14690E- 0:29469E- 0:32404E- -0:37778E- -0:17100E- -0:47152E-	05 08 11	0.21741E 0.49684E 0.55675E 0.91327E 0.35580E 0.97231E	-05 -08 -11 -14 -15	0:137036- 0:167106- 0:350746- 0:408426- -0:181346- -0:777336-1	07 09 10 11	0.84192E=05 0.94254E=08 0.19203E=09 0.12517E=09 0.12523E=11 0.48441E=15		0.73004E- 0.15625E- 0.14881E- 0.13701E- 0.59434E- -0.44104E-	-02 -04 -06 -09	0.775798 0.249358 0.326888 0.245758 0.125508 0.844488	-02 -04 -06	0-74977E= 0-10031E= 0-77100E= 0-38943E= -0-23623E= -0-24317E=	09 06 08	C+35782E-02 O+56518E-04 O+5685E-06 O+25532E-08 O+76950E-10 O+13768E-10	
	1 2	C+29019E=0	,	0.423156-	# • 0								x =	0.75			
	2 3 4 5 6	0.89880E-0 0.14093E-0 -0.484701E-1(-0.19890E-1) -0.89994E-1	5 7 0	0.15130E- 0.24930E- 0.35780E- 0.41342E-1 0.17647E-1	04 07 10 13	0.41896E-0 0.73472E-0 -0.16691E-0 -0.17421E-10 -0.26448E-10 -0.90334E-14	7	0.255075-0. 0.459675-07 -0.15285E-09 -0.10527E-09 0.59910E-11 0.56228E-14		0493406E=0 0422144E=0 0430475E=0 0425394E=0 0413710E=0 =0422833E=1	6 8	0.70763E- 0.34730E- 0.52636E- 0.45425E- 0.25663E- 0.14705C-	02 04 06 08	0-10640E=0 0-16260E=0 0-14308E=0 0-81778E=0 0-30210E=1: -0-10464E=1(3 5 9	0.48086E=02 0.90028E=04 0.883633E=06 0.49664E=08 0.64309E=10 0.10935E=10	
					E = 0.3	10							x = 0	.80			
	1 2 3 4 9 5	C=50844E=02 0=22364E=04 0=51270E=07 C=45390E=10 =0=14559E=12 =0=91351E=16		0.72732E=0 0.37546E=0 0.90284E=0 0.21039E=0 0.30448E=12	•	0-10452E-03 0-26759E-06 0-67095E-09 0-27693E-10 0-17779E-10 0-66552E-13		0.628485-04 0.143985-06 0.442225-0+ 0.118845-05 0.228415-10	-	0-11427E-00 0-30717E-02 0-47725E-04 0-45189E-06 0-27984E-08	!	0.10446E-0 0.47613E-0 0.82121E-0 0.80701E-0 0.51537E-0 0.25966E-1	2 4 6 8	0.14756E-01 0.25549E-03 0.25520E-05 0.16453E-07 0.38466E-10	į	04625995-02 04138745-02 04147745-05 04996376-08 04102745-09 04689415-11	
					x = 0.35	i							x = 0.1	15			
3 4 5 6		0.67049E-02 0.48351E-04 0.15029E-06 0.27967E-09 0.76198E-12 0.67965E-15	0	-11462E-01 -80892E-04 -26439E-06 -63866E-09 -16358E+11 -13923F-14	-0	-22660E-03 -7853-E-06 -19548E-08 -64775E-10 -20406E-10 -35542E-12	9 0:	e13423E-03 e47804E-06 e11469E-06 e13364E-09 e45408E-10 e21976E-12	0	0+19745E-00 0+1820E-02 1+72716E-04 1+77634E-06 1+54383E-08 119822E-10	0	0.11760E-00 0.64118E-02 0.12460E-03 0.13834E-05 0.99277E-03 0.48928E-10		0-20043E=01 0-39060E=03 0-43937E=05 0-31896E=07 0-13458E=07 0-14304E=10	0 0 0	• 78980F=02 • 20760E=03 • 25163E=03 • 25163E=03 • 13342E=09 • 35864E=12	
					= 0.40								= 0.90)			
1 2 3 4 5 6	0	0412471E=01 094291E=04 0938112E=06 092092E=09 031976E=11 037915E=16	0.	169325-01 157135-03 669955-06 163285-08 713595-11 788995-14	0 e 0 e	44338E-03 19961E-05 50609E-08 21068E-10 52879E-11 19680E-12	0.1 0.1 -0.1	29795F-03 12089E-05 50551E-08 64622E-11 72408E-12 3999E-11	0.	16265E-00 956013E-02 10816E-03 12926E-05 10100E-07 29425E-10	0.	13003E-00 94417E-02 18445E-03 22981E-05 18552E-07 11387E-09	0	0-26709E-01 0-58269E-03 0-7328E-05 0-59774E-07 0-31060E-09 0-31356E-10	0 a d 0 a d 0 a d	96543f=G1 30249f=Of 41503f=O5 95281f=O7 91887f=Of 3490f=10	
_				х -	0.45							x	0.95				
2	0 0 -0	18108E=01 17001E=03 86752E=06 26422E=08 12469E=11 17513E=13	0 • 2 • 0 • 1 ! 0 • 1 ! 0 • 20	3777E-01 8197E-03 5223E-05 7726E-08 061E-10 i127E-13	0.6 0.8 0.9 -0.2	0226E=03 5527E=05 4777E=07 1411E=10 1328E=12 2396E=11	0 = 21 0 = 81 0 = 71 0 = 25	5682E-01 7929E-09 1936E-08 1936E-10 300E-10 669E-11	0 • 1 0 • 2 0 • 1	18950E-00 19940E-02 5737E-05 0921E-05 8275E-37 5825E-10	0 • 1 0 • 2 0 • 3	14136E-00 10930E-01 16708E-05 7109E-05 3321E-07 1686E-09	0.	34961E-01 85119E-03 11899E-04 10783E-06 66064E-09 23024E-10	0 a 6 a 0 a 6 a 0 a 6 a	1417E-01 3019E-03 550ZE-05 2949E-07 7921E-05 554E-10	
		!9954E+01	0.13	X = (039E=0]		649E=02						x •	1.00				
	0.2 0.1 0.6	8619E-03 8092E-05 7263E-08 5266E-10 1375E-13	0.47 0.31 0.123 0.334	526E-03 669E-03 120E-07 526E-10	0.45 0.437 0.10 0.40	649E-02 167E-05 563E-07 326E-09 752E-12	0e565 0e221 0e656 0e426	7775-05 7775-05 7776-07 5885-10 5535-11 426-11	0.22	750E-00 368E-02 469E-03 027E-09 023E-07 104E-09	0.13 0.37 0.58	6133F-00 1938E-01 1911E-03 425E-05 066E-07 274E-09	0 • 1 0 • 1 0 • 1	4989E=01 2191E=02 8831E=04 8668E=06 3081E=08 5737E=12	0.590 0.103 0.109 0.802	025E-01 825E-03 982E-06 03E-06 06E-09 72E-11	

						x = 1.05						x = 1.55					
		п	ar n		a i n		ρ <mark>u</mark>	þ	1		ar n		an n		br n		.i
		1 0.2461 2 0.12463 3 0.15151 4 0.5095 5 0.5466 6 0.3717	5E=01 0E=03 5E=05 0E=07	0.159 0.175; 0.5285 0.8990 0.9870 0.7745	23E+01 35E+03 17E+05 13E+07 0E+09	0.171 0.291 0.3218 0.2480 -0.1633	6E=04 12E=06 12E=08	6E=02 0.81476E=0; 6E=04 0.15827E=04 2E=06 0.18436E=04 EE=08 0.15948E=06		0.4742 0.9643 0.4745 0.1575 0.3696 0.60626	225-00 185-01 15-02 55-03 75-05		01E-co 36E-0: 36E-0: 36E-0:	0 • 2621 0 • 2541 0 • 9376 0 • 2225 0 • 3725 0 • 4513	64E-00 51E-01 57E-05 55E-04	-0:4054 0:61151 0:4034: 0:11748 0:20221 0:26340	05-07 65-03 85-04 L5-01
	1	0.27484	c-aa			= 1.10							,	= 1.60			
	2 3 6 5	0 415826 0 449509 0 477036 0 490351 0 470610	E-01 E-03 E-05 E-07	0 - 1666; 0 - 21740 0 - 72493 0 - 13550 0 - 15336 0 - 13935	E-01 E-03 E-04 E-06	0.7092: 0.23759 0.44162 0.53467 0.45195 -0.51621	7E=02 ?E=04 ?E=06 E=08	0.148917 0.103795 0.236095 0.303355 0.279105 0.566605	-02 -04 -06 -08	0.48661 0.11329 0.59361 0.21110 0.52014 0.90971	E-00 E-02 E-03	0 • 1781; 0 • 9867; 0 • 87810; 0 • 31501; 0 • 91639; 0 • 16435	7E=01)E=02 .E=03 !E=05	0.28584 0.31431 0.12423 0.31444 0.55967 0.73065	E=01 E=02 E=04 E=06	-0.53165 0.65488 0.51575 0.155889 0.300208 0.412348	-02 -03 -04
	1					1.15							x	- 1.65			
	3 4 5 6	0.30313E 0.19987E 0.59227E 0.11428E 0.14645E 0.12679E	-01 -03 -04	0.171898 0.266378 0.979516 0.200366 0.264166 0.242646	-01 -03 -04 -06	0.869468 0.3?4248 0.636958 0.868028 0.796195 0.197448	-02 -04 -06 -08	0.146185= 0.142535= 0.142535= 0.45235= 0.475285= 0.475285= 0.478505=	02 04 06	0.49697E Ca19194E Ca19194E Ca73789E Ca72389E Ca12369E	-00 -02 -03	0:177178 0:107698 0:107068 0:462328 0:127118 0:242778	-00 -01 -03 -04	0.308528 0.384368 0.163118 0.438648 0.630018 0.115158	-01 -02 -04 -05	-0.66167E- 0.67251E- 9.65108E- 0.21344E- 0.44026E-	02
	1				X = :	1.20							X =	1.70			
	3 4 5	0.390525= 0.250165= 0.775826= 0.156666= 0.232465= 0.222865=(01 03 06	0.17574E= 0.32247E= 0.13054E= 0.29117E= 0.41848E= 0.41992E=	01 02 04 06	0:10492E- 0:43654E- 0:9606EE- 0:13807E- 0:13819E- 0:48737E-	02 04 05 07	0.131655-0 0.18331E-0; 0.45953E-0; 0.7673:T-0; 0.81814E-08 0.973405-10	; ; ;	0.50536E 0.15231E= 0.91200E= 0.35975E= 0.99719E= 0.19690E=	00	0.17618E- 0.11611E- 0.12935E- 0.5966E- 0.17448E- 0.35419E-	00 01 03	0.330625- 0.465375- 0.212316- 0.605616- 0.121606- 0.179536-	01 02 04	~0.79213E- 0.65249F-(0.81170E-(0.28890E-0 0.63714E-0	02 03 04
	1	0.156625+0	0.17633E		× = 1.								x = 1	.75			
	2 3 4 5 6	0:31051E=0 0:10567E=0 0:23919E=0 0:36192E=0 0:38138E=0	i d 2 0 4 0	0 = 38589E=0 0 = 17179E=0 0 = 41639E=0 0 = 64997E=0 0 = 70533E=0	1 2 4	0.12468E-0 0.58036E-0 0.13829E-0 0.21536E-0 0.23368E-0	2 3 5 7	0:102835-01 0:231475-02 0:699115-04 0:118345-05 0:135565-07 0:122135-09		0.51187E 0. 0.17421E=0. 0.11212E=0. 0.46412E=0. 0.13604E=0.4 0.28466E=0.6	0 1 3	0.17515E-0 0.12366E-0 0.15561E-0 0.76424E-0 0.23716E-0 0.51085E-0	0 1 3	0.35213E-0 0.55783E-0 0.27409E-0 0.82608E-0 0.17614E-0: 0.27596E-0	: 4	-0.919945-01 0.381275-02 0.999175-03 0.386755-04 0.911215-06 0.151585-07	•
	1	0-381135-00			X = 1.3	י							x = 1.8	0			
	2 3 6	0-38237E-01 0-13880E-02 0-39838E-04 0-55365E-06 0-63461E-08	0.	17989E-00 45655E-01 22142E-02 58673E-04 99182E-06 11700E-07	į	0:14595E=00 0:76251E=02 0:19518E=03 0:13014E=05 0:18922E=07 0:10:609E=09	0	*578425-02 *286935-02 *97049E-04 *17915E-09 *22255E-07 *25256E-09		0.51658E 00 0.19740E-00 0.13714E-01 0.57448E-03 0.18392E-04 0.40709E+06		0.17404E-00 0.13010E-00 0.13557E-01 0.97098E-03 0.31934E-04 0.72874E-06		0.37305E=00 0.66190F=01 0.35107E=02 0.11219E=03 0.25236E=05 0.41853E=07		0+104245-00 0+444215-02 0+121405-02 0+512285-04 0+128835-05 0+227365-07	
1		0:403835-00			= 1.35							×	- 1,85				
3 6 5 6		0.46724E-01 0.18091E-02 0.47230E-04 0.89277E-06 0.10304E-07	0.9 0.2 0.8 0.1	18062E-00 13413E-01 18740E-02 1548E-04 4884E-05 8911E-07	0000	•15839E=00 •99069E=02 •27460E=03 •49771E=05 •53144E=07 •53183E=09	0 a 0 a 0 a	42576E=03 34896E=02 13269E=03 25666E=05 35788E=07 35923E=06		0.51956E 00 0.22158E-00 0.13694E-01 0.75632E-03 0.24651E-04 0.57632E-06	0	•17283E-00 •13524E-00 •21973E-01 •12245E-02 •42623E-04 •10289E-05	0000	.39342E-00 .77735E-01 .44627E-02 .15068E-03 .35795E-05	0	115718-00 226576-07 145528-01 671638-04 160158-05 337098-07	
1	0	•42455E=00			1.40							x ·	1.90				
3 4 5 6	0	•55460E-01 •23263E-02 •45109E-04 •12337E-05 •16425E-07	0.61 0.36 0.11 0.21	073E-00 790E-01 594E-02 191E-03 994E-05 179E-07	0.1	19163E=00 12735E=01 17759E=03 73901E=05 0679E=06 2335E=09	0.4 0.1 0.3 0.5	33305+02 15945+02 78806+03 90576+05 54056+07 42506+09	0	0-52093E 00 0-24639E-00 0-20230E-01 0-95612E-03 0-32772E-04 0-60791E-06	0. 0.	17145E-00 13896E-00 25835E-01 15335E-02 56414E-04 14386E-05	0 • 0 •	41329E-00 9035EE-01 56315E-02 20071E-03 50279E-05 93055E-07	-0.1 0.1 0.6	126235-00 954935-01 171945-01 171815-04 49345-07	
1	٥.	449215=00		x = :								x =	1.95				
2 3 4 5 5 5	0. 0.1 0.1	68185E-01 29728E-02 88725E-04 18020E-05 25837E-07	0.705 0.4614 0.1917 0.3203 0.4704	75E-03 14E-05	0.14 0.51 0.15 0.15	15305=00 2025=01 18645=03 18185=04 8345=06 6595=09	0.48 0.23 0.56 0.87	814F+01 912E-02 769F+0T 941E-05 795E-07 166E-08	0.	52076E 00 27143E=00 24405E=01 12015E=02 43234E=04 11223E=05	0 • 1 0 • 1 0 • 1	6986E-00 6121E-00 0163E-01 9074E-02 4075E-04 9927E-05	0 • 1 0 • 2 0 • 6	3269E-00 0398E-00 0558E-02 6525F-03 9977E-05 3649E-06	0 - 20 0 - 20 0 - 11 0 - 34	35645-00 03565-02 00035-02 2075-03 01665-05 8215-07	
	04459775-00 04179805-00					204E-02						x = 2.	.00				
	0 • 1 : 0 • 2 !	1416E-01 7697E-02 1961E-03 5970E-05 9889E-07	0.7992 0.57672 0.20352 0.45038 0.72404	5E-01 2E-02 2E-03	0.204 0.700 0.156 0.244	706E-00 107E-01 193E-03 27E-04 64E-06 20E-08	0.311	285-02 645-03 15-05 125-06	0 • 2 0 • 2 0 • 1 0 • 5	11917E 00 9632E-00 9314E-01 5014E-02 6620E-04 5454E-05	0.14 0.34 0.23 0.96	800E-00 200E-00 965E-01 374E-02 329E-04 157E-05	0.11 0.87 0.34 0.76	165E=00 839E=00 789E=02 795E=03 839E=05	-0:10: 0:228 0:142 0:463	9825-00 9585-01 1745-02 705-03 665-05 985-06	

				x = 0.05						0 55		
		n ;	ar n	a _n i	b r	ьi		ar n	ai	x = 0.55	r	
	3 4 5	-0.4008	4E-08 0.42	3895-08 -0.26	 199E-07 571E-08 188E-09	0 = 11335E=(-0 = 06696E=(-0 = 68430E=0	07 0.4503	4E-01 9E-03 4E-05 4E-07 2E-10	21 0.456125-01 0.864626-03 0.698495-05 0.4323586-07 0.101426-09 0.385316-12	0.24516	E-02 E-04 E-06 E-09 E-10	bi 0.76944E-0 0.78250E-0 0.39503E-0 0.17469E-0 0.11808E-1 0.48444E-1
				x = 0.10					:	c = 0.60		
	3 4 5 6	0,236 86 ; 0,11765; -0,215928	E-06 0.177	94E-09 0.951 34E-09 0.101	09E-06 11E-09 78E-09	0-20126E-06 0-23826E-06 0-23133E-09	0.91951	E-03 0 E-05 0 E-07 0 E-09 0	*96968E-01 *13275E-02 *12777E-04 *70422E-07 *24975E-09 *10190E-11	0.444651 0.449928 0.256058 0.918128 -0.138608 -0.471148	-04 (-06 (-09 (-10 (0.11020E-02 0.14003E-04 0.85048E-07 0.37325E-09 0.79442E-11
	,	0.798345	.49	x = 0.15					×	= 0.65		
	2 3 4 5 6	0488819E- 0-50484E-	-06 0-1162	5E-05 0.3102	4E-08	0-15259E-05 0-10451E-08 0-92649E-10	0.76826E 0.15770E 0.14217E 0.88522E 0.3498E -0.14765E	-02 0. -04 0.: -07 0.:	59112E-01 19673E-02 22246E-04 14399E-06 10845E-05 1965E-11	0.66235E- 0.76621E- 0.52472E- 0.22530E- -0.65953E- -0.62264E-	04 0. 06 0. 08 0.	14987E-02 23760E-04 17159E-06 82045E-09 17950E-10 39594E-12
				x = 0.20					x =	0.70		
	2 3 4 5 6	0-14112E-0 0-37433E-0 0-37421E-0 -0-34435E-1 -0-19747E-1 -0-43553E-1	05 0.564738 06 0.631778 1 0.758188 4 0.305678	F-05 0.21741 -08 0.15851 -11 -0.4028 -14 -0.167378 -18 -0.718038	E-07 E-09 E-10 [-11	C.69447E-05 0.74926E-06 0.24994E-09 0.889071E-10 0.56175E-12 0.25656E-15	0.970126- 0.200346- 0.298266- 0.171816- 0.773686- -0.498226-1	0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0	16375-01 12765-02 71465-04 79025-06 14525-08 4235-11	0.95639E-0 0.13176E-0 0.1019IE-0 0.49958E-0 -0.23761E-1 +0.10700E-1	3 0.3 5 0.3 6 0.1 7 0.1	.9359E-02 .8586E-04 .2794E-06 .8186E-08 .8913E-10 .6986E-11
	1	0,37806E-0;		x = 0.25					x = (.75		
	3 6	0411425E-04 0418015E-01 0418015E-01 0410465E-10 -0418213E-13 -0479061E-17	0-171946- 0-284716- 0-435243- 0-353876-	-04 0.99219E- -07 0.16846E- -10 0.70845E- -13 -0.36550E-	-07 -09 -11 -11	0.19029E-04 0.34650E-07 0.58799E-10 0.22102E-10 0.33712E-11 0.29494E-14	0.12002E-0 0.28495E-0 0.38524E-0 0.31857E-0 0.17028E-0 0.27304E-1	2 0.39 0.59 5 0.51 0.20	072E-01 578E-02 813E-04 605E-05 103E-09	0.13435E-01 0.21299E-03 0.18895E-05 0.16640E-07 0.51761E-11 -0.51497E-11	0.50 0.59 0.36 0.83	0680E-02 0135E-04 0897E-05 368E-09 368E-11 327E-10
				x = 0.30					x = 0.	80		
		0.66331E-02 0.2844E-04 0.66237E-07 0.15956E-10 -0.13592E-12 -0.84381E-16	0.61903E-(0.42668E-0 0.10267E-0 0.10161E-0 0.26096E-1 0.16153E-1	0.35565E-0 0.01896E-0 0.06941E-1 2 -0.24077E-1	06 d 09 d 10 d	0-46269E-04 0-12364E-06 0-53074E-09 0-93510E-10 0-29655E-11 0-21539E-13	0.14563E-00 0.39505E-02 0.50374E-04 0.56735E-06 0.38659E-08 0.14905E-10	0.105 0.541 0.933 0.916 0.5786 0.1509	6E-06 9E-08	0.418412E-01 0.33362E-03 0.33655E-05 0.21653E-07 0.666225E-10 0.30647E-11	0 = 903 0 = 104 0 = 717	112E-02 176E-04 20E-05 32E-08 31E-10 49E-11
				x = 0.35					x = 0.9	5		
3 4 5	-	0a10718E-01 0a61523E-04 0a18876E-06 0a25059E-09 0a74109E-12 0a62422E-15	0412864E-91 0491921E-04 0430642E-06 0460605E-09 0413728E-11 0411944E-14	0+10430E-05 0+21127E-08 -0+70487E-10	0.	97205E-04 35828E-06 10652E-08 89191E-10 31989E-12 10726E-12	0.17347E-00 0.53874E-02 0.92060E-04 0.97519E-06 0.67801E-08 0.28135E-10	0-1167 0-7248 0-1415 0-15718 0-11251 0-90098	7E-02 7E-03 1E-05 E-07	0.24668E-01 0.50825E-03 0.57896E-05 0.42168E-07 0.18897E-09 0.69350E-11	0.2888 0.1313 0.1752 0.1370 0.8449 -0.3461	7E-03 DE-05 IE-07 4E-10
	_			K = 0.40					x = 0.90			
2 9 4 5	0.	016309E-01 012007E-03 048011E-06 011040E-08 018588E-12 036921E-14	0.18926E-01 0.17894E-03 0.76147E-06 0.18663E-08 0.50490E-11 0.67289E-14	0.56519E-03 0.26503E-05 0.67419E-08 0.21042E-10 0.12983E-11 -0.46880E-12	0.8 0.2 0.6 0.2	8303E-03 9691E-06 3103E-08 1570E-11 9987E-11	0 = 20301E = 00 0 = 72277E = 02 0 = 13702E = 03 0 = 16242E = 05 0 = 12658E = 07 0 = 62773E = 10	0.126178 0.952648 0.209548 0.26108E 0.20958E 0.10362E	-02 0. -03 0. -05 0.	32360E-01 75539E-03 96376E-05 78713E-07 39584E-09 10612E-10	0:27257 0:18509 0:28494 0:252218 0:161856 -0:967465	E-03 E-05 E-07
			×	= 0.45					x = 0.95			
1 2 9 9 9 9 9 9 9 9 9 9 9 9 9 9 9 9 9 9	0 . 0 . -0 .	29695E-01 21667E-03 10990E-05 32202E-08 00186E-11 6376E-13	0+26+29E-01 0+3203+E-03 0+17902E-05 0+3843E-08 0+11425E-10 0+30989E-13	0.10959E-02 0.60371E-05 0.1942E-07 0.63222E-11 -0.19381E-10 -0.34365E-11	0 • 20 0 • 68 0 • 67 0 • 15	812E-03 090E-05 445E-08 613E-10 698E-11 320E-12	0.23365E-00 0.95572E-02 0.19958E-03 0.26304E-05 0.22879E-07 0.13852E-09	0.13393E- 0.12309E- 0.30336E- 0.42161E- 0.37736E- 0.21984E-	01 0.1 03 0.1 05 0.1 07 0.6	01614E-01 10981E-02 5609E-04 4204E-06 4443E-07 9022E-11	0.20968E 0.25312E 0.44950E 0.4467IE 0.29183E -0.12601E-	-03 -05 -07 -09
	0.3	91 74E -01	X = 0-35361E-01	0.50					x = 1.00			
	0.3 0.2 0.8	67402-03 28082-03 37402-08 23682-11 11082-13	0.53940E-03 0.34017E-05 0.13805E-07 0.93903E-10 0.12006E-12	Gel7681F-02 0-12602E-04 0-49962E-07 0-99666E-10 -0-20549E-10 -0-49073E-11	0.411 0.171 0.107	92E-07 48E-09 52E-11	0-26473E-00 0-12473E-01 0-28915E-03 0-41549E-05 0-40053E-07 0-4053E-07	0.13994E-0 0.15654E-0 0.43053E-0 0.66380E-0 0.65950E-0 0.44743E-0	1 0.15 3 0.24 5 0.24 7 0.17		0.67778E-0.33815E-0.68960E-0.0.76733E-0.0.99160E-0.0.93813E-1	03 05 07

						x = 1.05		x = 1.55								
	n ar an ain					br n	b <u>i</u>			a,r		-1	X = 1.05			
	2		40E-01 40E-01	0.10 0.60 0.10 0.11	422E-00 416E-01 909E-09 214E-04 196E-08 220E-09	01 0-21889E-0 05 0-38047E-0 04 0-42324E-0 06 0-31801F-0		2 0.43343E-0 4 0.10313E-0 6 0.12836E-0 9 0.10606F-0		0.12 0.61 0.20 0.44	71E-07	0.905		0.284 0.117 0.287 0.486	br 751E-00 752E-01 759E-02 752E-04 18E-06 82E-08	bi -0.9415dE-01 -0.19077E-02 0.1765dE-03 0.65763E-03 0.1912eE-04 0.17307E-08
	1	0.1914	95 -00										×	- 1.60		
	2 3 4 9	0.2056 0.5556 0.6700 0.11314 0.91769	1E-01 0E-01 DE-05 DE-04	0.242 0.822 0.158		0.7921 0.3011 0.3751 0.7025 0.3883 0.1450	5E-02 8E-04 3E-08 3E-08	-0.50479E- 0.50189E- 0.1505EE- 0.20681E- 0.19189E- -0.12605E-	04 04 04	G.519: O.192: G.769: G.2676 O.4894 G.1139	6E-02	0.1450 0.9801 0.9861 0.40261 0.16404 0.18656	7E-01 1E-02 1E-01	0.2672 0.9510 0.1947 0.4042 0.7291 0.9628	3E-01 0E-02 3E-04 0E-04	-0+10781E-00 -0+27078E-02 0+2078E-05 0+8947E-05 0+19207E-06 0+27257E-08
	1	C . 35448(T-00			x = 1.15							x ·	= 1.65		
	,	0.25959; 0.79959; 0.79950E 0.14900E 0.14959E G.14424E	2-01 2-09 3-04	0.1483 0.2999 0.1113 0.2276 0.2999 0.27320	4E-01 1E-02 1E-04	0.9488 0.40792 0.85319 0.1393 0.134716 0.538486	E-02 E-34 E-05	-0.98201E-0 0.69489E-0 0.21497E-0 0.33148E-0 0.33114E-0 0.486559E-1) •	0.452701 0.16464 0.95249 0.35143 0.91141 0.16877	E-00 E-02 E-09	0+14582 0+10466 0+11993 0+124086 0+144308 0+279648	-00 -01 -09	0.20096 0.42104 0.20169 0.56207 0.10794 0.151888	E-01 E-02 E-04	-0+12190E-90 -0+0424E-02 0+29603E-03 0+11816E-04 0+22723E-04 0+42697E-08
	1	0.361565			x =	1.20							x = 1	1.70		
	3	0.92902E- 0.10168E- 0.21010E- 0.29162E- 0.28412E-	01 02 04 04	0.14872 0.39546 0.14804 0.33073 0.47516 0.472585	-04 -06 -06	0:111845 0:54440E 0:124375 0:18090E 0:18174E 0:10961E	-02 -03 -05 -07	-C.16007E-01 0.76144E-03 0.30042E-04 0.51445E-04 0.56472E-08 0.26989E-13		0.53248E 0.18772E 0.11789E 0.45739E 0.12542E 0.24668E	-00 -01 -03	0.14549E- 0.11019E- 0.14449E- 0.67628E- 0.19809E- 0.40214E-	00 01 03	0.30453E 0.4995E 0.26043E 0.7733E 0.19782E 0.23594E	-01 -02 -04 -05	-0.19438E-00 -0.72144E-02 C.29849E-03 U.19469E-04 O.39470E-06 C.69457E-08
	1	0-404716-0	00	0.148395	* = 1. -00	45 0.12985E-							X = 1.	75		
	2 5 6	0=40333E=0 0=13492E=0 0=30178E=0 0=45386E=0: 0=48009E=0:	2	0.42257E- 0.19470E- 0.47293E- 0.73819E- 0.79770E-	-01 -02 -04 -05 09	0.71721E-0.17641E-0.0.26173E-0.0.30737E-0.0.20543E-0	72 33 15	-0.23678E-01 0.84419E-03 0.41128E-04 0.76245E-04 0.76245E-04 0.44090E-08		0.99600E 0.2120E= 0.1490IE= 0.99064E= 0.17149E=(0.99699E=0	00 01 39	0:14829E-0 0:11442E-0 0:17918E-0 0:46548E-0 0:26915E-0	0 1 3	0.32602E-0 0.58645E-0 0.39334E-0 0.10534E-0 0.22811E-0 0.34154E-0)1)2 3	-0.14688E-00 -0.10521E-01 0.26899E-03 0.19941E-04 0.99451E-05 0.99451E-05
1		0-42968E-00	,	0.147636-0	x = 1.3								× = 1.80	0		
3 4 5		0.49619E-01 0.17745E-02 0.42721E-04 0.69454E-08 0.79712E-08	0	0+49370E-0 0+29304E-0 0+66490E-0 011264E-0 013229E-0	2 4 5	0.14868E-00 0.9323EF-02 0.25215E-03 0.43096E-05 0.51046E-07 0.42119E-09	0	-32801E-01 -57798E-03 -55187E-04 -11657E-05 -1253E-07 -12911E-09		0.53773E 0 0.23715E-0 0.17749E-0 0.75739E-0 0.23199E-04		0-149775-00 0-117295-00 0-209465-01 0-109895-02 0-362365-04 0-827325-06	0	0.34547E-00 3.48139E-01 9.42300E-02 9.14213E-03 9.92609E-09 9.54748E-07		Col3854E-00 Col4649E-01 Col3758E-03 Col3758E-04 Col465E-06 Col4658E-07
ļ		0.4503 8 E-00	0.	14670E-00	x = 1.35	. 148678						x	- 1.85			
3 4 9	9	0+60509E-01 0+23109E-02 0+39674E-04 0+10453E-05 0+12950E-07	0. 0. 0.	.97410E-01 32923E-02 92995E-04 16903E-05 21426E-07	0	-14807E-00 -11972E-01 -39148E-03 -64837E-05 -82867E-07 -75221E-09	0 • 1 0 • 2 0 • 2	43262E-01 82816E-03 72996E-04 17054E-05 24297E-07 1932E-09	0	•99777E 00 •26292E-00 •21619E-01 •98459E-03 •31115E-04 •72327E-06	0	015141E-00 011859E-00 024234E-01 013849E-02 08955E-04 11681E-05	0.	• 34489E-00 • 78343E-01 • 53226E-02 • 19006E-03 • 6138E-05 • 81894E-07	-0. 0. 0.	-16928E-00 -16928E-01 -21400E-03 -31652E-04 -10525E-05 21580E-07
1	3.	44874E-00	0.1	4361E-00		187615-00	-4.4					х -	1.90			
9	0.	73194E-01 29823E-02 82924E-04 15492E-05 20638E-07	0.4 0.1 0.2	5694E-01 1369E-02 2704E-03 4977E-05 4093E-07	0 e 1 0 e 1 0 e 1	15188E-01 18366E-03 16052E-05 3212E-06 2937E-08	0.64 0.91 0.24 0.38	0878E-01 0901E-05 0603E-06 0512E-05 1049E-07 156E-09	0. 0. 0.	59622E 00 28780E-00 28187E-01 12207E-02 1393E-04 0144E-05	0. 0.1 0.1	15305E-00 11654E-00 18330E-01 7529E-02 3986E-04 6329E-05	0 • 6 0 • 6	38429E-00 19253E-01 16433E-02 15197E-03 4637E-05 2117E-06	-0.2 0.1 0.3 0.1	17893E-00 19617E-01 2332E-03 8910E-04 4436E-05 1178E-07
1	040	84 8 4€-20		x = 9126-00	1.45							X = .	1.95			
	0.3 0.1 0.2	7463E-01 0165E-02 1226E-03 1629E-05 1386E-07	0.740 0.521 0.172 0.363	772E-00 053E-01 104E-02 125E-03 179E-05 177E-07	0.19 0.65 0.14 0.20 0.22	771E-00 047E-01 794E-03 024E-04 724E-06 989E-08	0.202 0.118 0.346	19E-01 199E-03 29E-03 39E-05 67E-07 89E-09	0.31 0.31 0.15 0.54	3317E 00 1254E-00 1571E-01 1354E-02 444E-04 094E-05	0 • 11 0 • 32 0 • 21 0 • 63	9461E-00 721E-00 856E-01 534E-02 995E-04 617E-05	0.10 0.82 0.33 0.49	0949E-00 0072E-00 0251E-02 132E-03 713E-05 727E-05	-0.32 -0.33 0.46	1739E-00 514E-01 716E-04 979E-04 029E-05 722E-07
	0.49	865E-00 409E-00	0+144	74E-00	0.221	'64E-00	-0.8059	l ar -ar				# = 2.	00			
	0.19	409E-00 469E-02 150E-03 144E-05 174E-07	0.624: 0.6502 0.2309 0.5227 0.6220	10E-02 14E-03 19E-05	0.234 0.883 0.202 0.319	15E-0] 82E-03 09E-04 85E-06 79E-08	-0.3420 0.1461 0.4817 0.8820 0.10939	14E-03 1E-03 3E-03 4E-07	8:375	75E 00 54E-00 178E-81 17E-04 21E-05	8:273	97E-00 81E-00 3 5E-02 42E-03 9E-03	0.112	06E-00 03E-03 41E-04 93E-06	-8:399	12E-02 19E-05

						x = 0.05								x = 0.55		
		n ;	ar n		ai n		b r	b_n^i		2			,r	⊾i.		
	1 2 3 4 5 6		7E-08	0.458	485E-G: 394E-0: 255E-1:	-0.126	11E-07 67E-08	0.60698E-	-oa	0.55050 0.69575 0.51205 0.22669 0.59597	8E-01 8E-03 8E-05 8E-07		4E-03 4E-05 3E-07 0E-07		-5-02 15-04 15-04 15-09 15-11	
						x = 0.10							,	ı = 0.60		
	3 4 5 6	0.28118 0.13613: 0.39931!	E-06	0.3556 0.2036 0.6723	05E-06 78E-10	0-710A 0-2413 0-1271	6L-09	u•12199F⇒ 0•36369E= 0•58117E=1	11	0.728205 0.107995 0.239995 0.495012 0.17845E	-07 -05 -07 -39	0.63345 0.15251 0.14629 0.80365 0.26721 0.13754	E-02 E-04 E-07 E-09	0.536739 0.555008 0.17448 0.106938 -0.137735 0.439865	-04 -04 -04 -11	0.26297E=03 0.466690E=05 0.465789E=07 0.16447E=09 =0.18429E=10 =0.35917E=11
		4 0174				= 0.15						×		- 0.65		
	3 6 5 6	0.95704E- 0.10338E- 0.53199E-	-05	0:11975 0:15404 0:96453	£=05	0.53962 0.36448 -3.15061	-08	-08 C-74562E-0		04939218-01 04162018-02 04164348-04 04101488-05 04417138-09 00403788-11		0.759385 0.22601E 0.254785 0.164435 0.665815 0.625245	-02 -04 -06	0.79290E=0: 0.95723E=04 0.64987E=06 0.27075E=08 0.57207E=1:2 0.39612E=11		0:19823E-03 0:10657E+04 0:89902E-07 0:41592E-09 -0:17685E-10 -0:42619E-11
						- 0.20							x =	0.70		
	2 3 4 5 6	0.22955E=0 0.43593E=0 0.41561E=0 =0.33427E=1 =0.13848E=1 =0.38173E=1	05 08 11 .4	0.282926 0.648075 0.706565 0.651395 0.273145 0.744425	-05 -08 -11 -14 -15	0.22723E 0.26676E -0.73827E -0.77138E -0.15238E -0.52946E	-07 -10 -13	0.36852E-05 0.50955E-08 0.27431E-09 0.10059E-10 0.15448E-12 0.10867E-15		0+11834E-0 C+23618E-0 D+27562E-0 O+19710E-0 C+92763E-0 -0+30149E-1) 2 6 5	0.88421E= 0.32483E= 0.42545E= 0.31971E= 0.15030E= 0.13493E=	02 04 05	0.11339E-0 0.16161E-0 0.12605E-0 0.6119E-0 -0.24562E-1 0.51464E-1	5 8 1	-C.15435F-04 0.16078F-04 0.15673F-06 0.95137E-09 -0.28074F-10 -0.99464E-11
	,	0 444135	_			0.25							x = 0).75		
5	2	0.45513E-0; 0.13314E-04 0.20715E-07 0.23936E-10 -0.15716E-13 -0.69293E-17	,	0.54932E- 0.19732E- 0.32735F- 0.38808E- 0.31828E- 0.13510E-	04 07 10	0.692656- 0.123486- 0.142376- 0.861536- 0.418396-1 -0.723366-1	10	0:107125-04 0:203785-07 -0:343965-11 0:209505-11 0:243235-11 0:127765-14		0 - 14584E - 00 0 - 33595E - 02 0 - 44600E - 04 0 - 36575E + 06 C - 19938E - 08 0 - 16025E - 10	!	0.10019E-0 0.45459E-0 0.46118E-0 0.55969E-0 0.17200E-1	2	0.15753E+0: 0.26034E-0: 0.26034E-0: 0.23336E-0: 0.13220E-0: 0.52685E-1: 0.69552E-1:		-0.47450E-03 D.22730E-04 0.29263E-06 0.18592E-08 -0.18520E-10 -0.30352C-11
					x = 0	.30							x = 0.	80		
1 2 3 4 5 6	:	0.80070E=02 0.33168E=04 0.73956E=07 0.68997E=10 0.12310E=12 0.73963E=16		0.94074E-0 0.48972E-0 0.11714E-0 0.15120E-0 0.23089E-1 0.14415E-1	6 9 2	0.17212E-0 0.44216E-0 0.60152E-0 -0.32189E-1 -0.50335E-11	,	0.251075-04 0.71709E-07 0.27092F-09 0.54775E-11 -0.67388E-11 0.83900E-14		0.17598E-00 0.46756E-02 0.69959E-04 0.65158E-04 0.40150E-08 0.24051E-10		0-11069F-00 0-62142E-03 0-1069[F-03 0-10479F-04 0-66082E-08 0-28460E-10		0.21317E-01 0.40622E-03 0.41491E-05 0.24855E-07 0.11330E-09 0.47805E-11		-0.130375-02 0.301945-04 0.490655-06 0.371575-08 -0.263195-11 -0.428855-11
					x = 0.3	15						,	= 0.8	5		
1 2 3 4 5 6	- C	0-12976E-01 7-71870E-04 0-21742E-06 0-33819E-09 0-71553C-12	000	•14747E-01 •10552E-03 •34350E-06 •61451E-07 •12150E-11 •10660E-14	-	0.37132E-03 0.12977E-05 0.24458E-08 0.58677E-10 0.95468E-11 0.30856E-12	-(0.500875-04 0.203975-05 0.622545-09 0.652545-09 0.652545-0 0.6372635-10 0.6372635-13		0.20814E-00 0.61965F-02 0.10678E-03 0.11204E-05 0.77555E-08 0.45570F-10	0	0-11945E-00 1-83179E-07 1-16222E-03 117967E-05 12816E-07 16255EC-10	0	0.29143E-01 0.61617E-03 0.71179E-04 0.52109E-07 0.25420E+07		0.26559E-02 C.2713IF-04 O.78733E-06 0.49551E-08 O.49551E-08 0.43594F-11
					c = 0.40	•						x	= 0.90			
1 2 3 4 5	0. 0.	-19803E-01 -14026E+03 -55337E-06 -12826E-08 -35010E-11 -31090E-14	0 • 0 • 0 • :	21626E-01 20498E-03 87152E-06 21108E-08 25133C-12 50364F-14	-0 -0	1.722136-03 1.32950E-05 1.83304E-03 1.95201E-11 1.24578E-11 1.15480F-11	0 0 0 -0	.08993E=04 .49826E=06 .14215E=08 .69390E=12 .59706E=11 .11971E=11	000	0 = 24155E = 00 0 = 86046E = 02 0 = 15909E = 03 = 18671E = 05 + 14536E = 07 - 87732E = 10	0. 0.	126206-00 10724E-01 24015E-03 29849E-05 23909E-07 13252E-09	0.	36367E-01 91133E-03 911830E-04 97256E-07 34316E-03 20204E-11	0 0 0	+71015-02 +41045-04 +121165-05 +124625-07 514735-10 553015-11
				x	• 0.45							x =	0.95			
2 3 4 9 6	0.1 0.1 0.2	78853E-01 75178E-03 12504E-C5 17165E-08 15495E-11 4448E-13	0+3 0+1 0+60	0072E-01 6784E-03 9802E-05 0895F-08 2997E-14 7880E-13	0. 0. -0.	12967E=02 74932E=05 24029E=07 10478F=10 41784E=11 17242E=12	0 • : 0 • : 0 • :	13863E+03 10845E-05 39171E-08 54479E-11 0013E-10 7716E-11	0.	27539E=00 2114105=0; 23199E=03 30256E=03 26239E=07 16921E=09	0.1 0.4	3085E-00 4098F-01 4772E-03 8209E-05 3086E-07 7063E-09	0 • 1 0 • 1 0 • 1	45968E-01 13176E-02 19117E-04 17533E-06 1125E-08 3671E-10	0 ± 1 0 ± 2 0 ± 1	76506E-02 17704E-04 17704E-05 21444E-07 2022E-09 1872E-11
	0-4	x = 0.50 «40492E=01 0.40003E=01 0.21#446=03 0.10447								x = 1.00						
	0 • 2 6 0 • 2 6 0 • 2 7	7492E-01 9041E-03 9146E-05 336E-08 778E+10 341E-13	0.61 0.41 0.15 0.29	003E-01 999E-03 227E-05 592E-07 569E-10 739E-12	1 • 0 6 • 0 1 • 0 2 • 0=	1866E-92 5614E-04 1912E-07 2608E-09 1786E-12 5292E-11	0.2 0.95 0.12	96695-03 14535-05 36535-08 17495-10 1666-11 5775-11	0 • 1 0 • 3 0 • 4 0 • 4	0687E-00 4935E-01 3165E-03 7613F-05 5891E-07 1444E-09	0+17 0+49 0+75 0+75	13515-00 19015-01 1354E-03 9105-05 2785-07 0055-09	0 · 16 0 · 30 0 · 30 0 · 21	3659E-02 0102E-04 0649E-06 526E-09	0.25 0.25 0.27	1361E-01 1590E-04 4402E-05 1718E-07 1621E-09 560E-12

					x = 1.05										
		n	ar n		a ⁱ n	br n	$\mathbf{b}_{\mathbf{n}}^{i}$		ar n			x = 1.55			
		2 0.193 3 0.4669		0.1344 0.2238 0.6879 0.1168 0.1278 0.9719	0.4F-00 0.6 13E-01 0.2 18E-03 0.49 2E-04 0.53 4E-05 0.40	9366E-01 5927E-02 5327E-04 7072E-06 1256E-08	-0.16900E -0.20445E 0.34513E 0.57515E- 0.50035E- -0.40371E-	-04 0 -05 0 -07 0	"n 542755 c 146905-0 725255-0 735155-0 734955-0 73975-0	0 0 9479 2 0 9190 3 0 3503 5 0 8476	73E-02 73E-03	0.2333 0.3056 0.1363 0.3448 0.5918 0.7426	9E-01 7E-02 0E-04	51 70+137556- 70+837486- 70+843836- 0+946046- 0+431266- 0+707878	0.0
					x = 1.10						x	= 1.60			٠.
	1 2 3 4 5	0.37210 0.24730 0.64589 0.11179 0.12960 0.10725	E-01 E-03 E-04	0.13404 0.27586 0.94337 0.17607 0.21108 0.17549	E=01 0.334 E=03 0.698 E=04 0.882 E=06 0.728 E=08 0.426	932E-01 400E-02 311E-04 157E-06 118E-08 24E-10	-0:23482E-c -0:99630E-c 0:44682E-c 0:89786E-c 0:490179E-c -0:15342E-1	0.17 0.91 7 0.31 9 0.75	1995E 90 1009E-00 031E-02 158E-03 359E-05 052E-06		E-00 E-01 E-03 E-04	0.25200 0.36496 0.17821 0.48303 0.68553 0.118488	E-01 E-02 E-04 E-06	-0.14674E-0 -0.11231C-0 -0.14229E-0 0.81403E-0 0.58279E-0 0.10923E-0	3 6 7
	1	0.400866			x = 1.15						x =	1.65			
	2 0,31332 3 0,00155 4 0,16606 5 0,21032 6 0,18940		2E=01 0.33503 5E=03 0.12742 5E=04 0.26038 5E=04 0.34256		-01 0.4759 -02 0.1031 -04 0.1394	0E-02 9E-03 2E-05 4E-07	-0.31464E-0: -0.2367IE-0: 0.54595E-0: 0.13595E-0: 0.13595E-0: 0.13400E-0: -0.34290E-12	0:194 0:113 0:409 0:105	74E-00 74E-00 59E-01 31E-03 01E-04	0.152218 0.105458 0.136688 0.59982E 0.14493E 0.31469E	-00 -01 -03	0.27044E-00 C.45092E-01 0.23050E-02 0.66900E-04 0.13077E-05 0.18588E-07		-0:16C53E-00 -0:14754E-01 -0:22754E-03 0:36872E-06 0:76576E-07	
	1	•			x = 1.20						x = 1	1.70			
	2 3 4 5	0.42729E= 0.439313E= 0.11879E= 0.424248E= 0.39409E= 0.32704E=	01 02 04 06	0.13086E- 0.40132E- 0.16974E- 0.37839E- 0.354269E- 0.54090E-	01 0.62896 02 0.14985 04 0.22124 06 0.22453	E-02 E-03 E-05 E-07	-0.40834E-01 -0.45811E-01 0461569E-05 0419832E-06 0x25456E-08 0x51809E-11	0.55d3 0.4204 0.1408 0.5333 0.1448 0.2828	5E-00 0E-01 •E-03 5E-04	0 • 13550E- C • 10 866E- C • 16457E- O • 77395E- O • 22638E- O • 45914E-	00 01 03	0.28914E- 0.5033E- 0.29320E- 0.91001E- 0.19073E- 0.28810E-	01 02 04 05	-0e17373E-00 -0e19012E-01 -0e34987E-03 -0e32308E-06 0e97070E-07 0e24001E-08	
	1	0.451246-0			x = 1.25						x = 1.	75			
	2 3 4 5	0.4865E-0 0.15620E-0 0.15620E-0 0.34857E-0 0.52060E-0 0.55297E-0	1 0 2 0 4 0	•12889E-0 •47404E-0 •22321E-0 •54113E-0 •64915E-0 •91359E-0	1 0.81992E 2 0.21407E 6 0.34377E 6 0.37910E 7 0.31699E	-02 -03 -05 -07	-0.51512E-01 -0.79805E-03 0.61269E-05 0.28166E-06 0.41210E-08 0.17458E-10	0.55963 0.24678 0.17360 0.68969 0.19789 0.40934	E-00 E-01 E-03 F-04	0.13939E=0 0.11039E=0 0.19650E=0 0.99032E=0 0.30765E=0 0.66221E=0	0 1 3	0.30813E=0 0.58203E=0 0.37448E=0; 0.12430E=0; 0.27495E=0; 0.44133E=07	1	-0.18619E-00 -0.224062E-01 -0.52155E-03 -0.52022E-05 0.11788E-06 0.34147E-08	
1	ı.	0:47264E-00	_		× = 1.30						x = 1.80)			
3		0440171E-01 2020845E-02 0049388E-04 0079690E-05 0091317E-08	0 • 0 • 0 • 1	12711E-00 55209E-01 25005E-02 76247E-04 12867E-05 15109E-07	0 = 14592E= 0 = 10542E= 0 = 30114E= 0 = 52460E= 0 = 64764E= 0 = 25604E=(01 - 03 6	0.63356E-01 0.12996E-02 1.46544E-05 0.1719E-06 0.65302E-08 0.58685E-10	0.55919E 0.27327E 0.21288E 0.88955E 0.26790E 0.58583E	-00 -01 -03	0 - 14949E-00 0 - 11038E-00 0 - 23267E-01 0 - 12572E-02 0 - 41421E-04 0 - 94468E-06	' 6	0 • 32741E-00 0 • 6664E-01 0 • 47068E-02 0 • 16694E-03 0 • 39199E-05 1 • 66602E-07	-	0:19774E~00 0:29941E~01 0:75783E~03 0:50466E~05 0:13596E~06 0:48275E~08	
1					= 1.35					x	- 1.85				
3 4 5	į	0.49150E-00 0.73394E-01 0.73394E-02 0.49048E-04 0.21999E-05 0.14836E-07	0 • 6: 0 • 3: 0 • 10: 0 • 10:	2576E-00 3385E-01 7269E-02 3597E-03 3310E-05 4476E-07	0.16301E-0 0.13373E-0 0.41777E-0 0.41777E-0 0.78719E-0 0.10169E-0 0.97763E-09	1 -0	76175E-01 20151E-02 65207E-06 551482E-06 99641E-08 90487E-10	0.59707E 0.29947E= 0.29965E= 0.11294E= 0.39961E= 0.83903E=	00 01 02	0.14775E-00 0.10903E-00 0.27317E-01 0.15839E-02 0.55275E-06 0.13337E-05	0.	.34698E-00 .75611E-01 .58631E-02 .22211E-03 .55297E-05 99423E-07	-0 -0	•20827E-00 •3665E-01 •10772E-02 •96038E-05 •14446E-06	
1	o	•50786E 00			1.40					x	1.90				
2 3 4 5 6	0.	88656E=01 35170E=02 95350E=04 17794E=05 23635E=07	0 • 71 0 • 4 71 0 • 1 4 5 0 • 2 6 5	501E-00 711E-01 567E-02 540E-03 537E-05	0+18039E-00 0+16748E-01 0+57183E-03 0+11630E-04 0+16187E-06 0+16580E-08	-0.1 -0.1 0.6	89746E-01 8066E-02 14864E-05 15750E-06 4955E-07 7029E-09	0.55937E Q 0.32499E-0 0.51500E-0 0.14315E-0: 0.47881E-0: 0.11448E-0:	0 0	0.15205E-00 0.10645E-00 1.31797E-01 1.1981ZE-02 0.73141E-06 0.18646E-05	0 • 8 0 • 7 0 • 2 0 • 7	36686E-00 35052E-01 72397E-02 19291E-03 17233E-05 4682E-06	-0. -0.	21765E-00 44230E-01 15020E-02 16646E-04 13287E-06	
1	0.	52179E 00		X = 1						x =	1.95				
2 3 4 5 5 6 6 6 6 6 6 6 6 6 6 6 6 6 6 6 6 6	0 e 4 0 e 2 0 e 2	10602E-00 95103E-02 13017E-03 4014E-05 7041E-07	0.7992 0.5945 0.1971 0.4154 0.6092	20E-01 32E-02 15E-03 19E-05 9E-07	0.19799E-00 0.20715E-01 0.77294E-03 0.16932E-04 0.25328E-06 0.27685E-08	-0.42 -0.22 0.80	0383E-00 1446E-02 1052E-04 034E-06 873E-07 720E-09	0.54818E 00 0.34950E-00 0.38011E-01 0.18039E-02 0.43266E-04 0.16197E-05	0.0	15623E-00 10290E-00 36466E-01 26610E-02 96010E-06 25627E-05	0.94 0.88 0.38	1700E-00 1916E-01 1916E-02 298E-03 685E-06	-0.57 -0.27 -0.27	2578E-00 2508E-01 0581E-02 7173E-04 111E-07 383E-07	
	0.5	1339E 00	0.12976	X = 1.9 F=00	0.21580E-00					x • 2	.00				
	0.12 0.51 0.17 0.37	1547E-00 1397E-02 1583E-03 1537E-05 164E-07	0.0123706 0.7930 0.26434 0.59742 0.93832	E-01 E-02 E-03 E-05	0.21980E=00 0.25913E=01 0.10325E=02 0.24918E=04 0.39009E=08 0.49620E=08	-0.610 -0.462 0.913 0.311	118E-00 67E-02 37E-04 59E-04 66E-07 73E-09	00-34163E 00 00-37272E-00 00-3525E-01 00-35370E-02 00-3270T-04 00-4340E-03	0 0 1 0 0 0 0 0 0 0 0 4	06135-00 d7605-01 17465-01 03065-02 43071-03 34305-05	0e107 0e107 0e1e6	002-01 002-01	-0004 -0027 -00422	2075-00 7015-01 7055-02 306-04 356-07 706-07	

x = 0.05														x = 0.55			
		n a	n		a ⁱ n		br n	b_n^i		á	ar n	4	ı n	b	5	1	b <mark>i</mark>
		1 0.2537 2 0.3995 3 -0.3490 4 5	4E-08	0.4098 0.6179 0.8437	99E-08	0•2022 0•9027 0•1491	0E-09	0.7041JE- 0.204945- 0.21609E-	01	0:4060 0:5116 0:3805 0:1696 0:50866 0:81508	7E-03 7E-05 7E-07 SE-10	0.5021	1E-01 3E-03 1E-05 8E-07	0.30063 0.25191 0.11986	E-02 E-04 E-06 E-09 E-11	0:123 0:118 0:5892 0:1778 -0:2505 0:1703	94E-02 95E-04 92E-07 12E-09 1E-12
					,	= 0.10							,	x = 0.60			
	1 2 3 4 5 6	0+20394 0+1004[0+250158	-06	0.32773 0.18427 -0.40118	7E-06	0.57756 0.42671 -0.47647	E-11	1 -0.370625-10		0.54160 0.79352 0.698318 0.369898 0.113028 -0.748308	-03 -05 -07	0.63221 0.13914 0.13259 0.72786 0.255318	E=02 E=04 E=07	0.46590E=0 0.46344E=0 0.26208E=0 0.91959E=0 =0.15527E=1 =0.43122E=1		0.18204 0.21461 0.12769 0.53039 0.56557	E=04 E=06 E=09 E=11
					x	0.15	0.15						x	= 0.65			
	1 2 3 4 5 6	0.69381E 0.75349E 0.41390E	-06	0-110558 0-139928 0-847028	-05	0.439778 0.284825 -0.109786	-0 a	d 0.15332F=0A		0-70596E-01 0-11699E-02 0-1203E-04 0-75814E-07 0-24843E-09 -0-12256E-11		0.77572E 0.20646E 0.23096E 0.14884E 0.62026E 0.15958E	-02 -04 -06 -09	0.70001E-02 0.81199E-04 0.83783E-06 0.22797E-08 -0.42929E-11 -0.32890E-11		04255431 0e367875 0e258756 0e117175 0e929276 -0e120306	-04 -06 -06 -11
					x =	0.20							x =	0.70			
	1 2 3 4 5 6	0.16635E-0 0.32191E-0 0.32531E-0 -0.28351E+1 -0.18850E-1 -0.38188E-1	1 4	0,26169E- 0.56872E- 0.64300E- 0.78035E- 0.32063E- 0.87549E-	05 08 11	0.18570E- 0.21870E- 0.11916E- -0.20228E- -0.14197E- -0462957E-	C7 09 10 11	0.93609E-05 0.10984E-07 0.19445E-09 0.56923E-10 0.71931E-12 0.32463E-15		0.90093E=0 0.17328E=0 0.20457E=0 0.14727E=0 0.648777E=0 =0.32240E=1	02 04 06	0.92243E= 0.29717E= 0.38579E= 0.2884EE=(0.14034E=(0.51754E=)	02 34 36	0:10180E=0 0:13648E=0 0:13646E=0 0:51817E=0 0:13293E=1:-0;37941E=1	9	0.343245- 0.603295- 0.496792- 0.256925- 0.231105- 0.174765-	06 08 10
					x • 0	.25							x = ().75			
	1 2 3 4 5 6	9.32978E=0: 0.98273E=0: 0.15425E=0: 0.57371E=11 =0.16010E=13 =0.69322E=17	3	0:50976E-0 0:17930E-0 0:29722E-0 0:40433E-1 0:37070E-1 0:15889E-1	0 3	0.56815E-0 0.10106E-0 0.15490E-0 -0.17653E-1 -0.35977E-1 -0.72957E-1	6 9 1 I	0.28290E-04 0.51269E-07 0.10885E-09 0.22661E-10 0.24229E-11 0.37286E-14		0.11269E=0. 0.24631E=0; 0.33065E=0. 0.27302E=0. 0.14806E=0.6 0.66861E=11	2	0.10731E-0 0.41667E-0 0.62149E-0 0.53370E-0 0.29553E-0 0.96162E-1	2 6 5	0.14410E-01 0.22133E-03 0.19435E-05 0.10952E-07 0.37859E-10 0.51313E-12		0.44195E=(0.95171E=(0.951010E=0 0.539329E=0 0.14947E=1	6
					x = 0.	30							x = 0.	āO			
9		0.58037E=02 0.24471E=04 0.59115E=07 0.65912E=10 -0.11613E=12 -0.74001E=16		0.67692E-02 0.44511E-04 0.10604E-06 0.15198E-09 0.25810E-12 0.16952E-19		0.10182E-03 0.36134E-06 0.53240E-09 0.40023E-11 0.424894E-11 0.53527E-13		0.69325E-04 0.18235E-05 0.28963E-09 0.11014E-10 0.10574E-11 0.24067E-13	0	0+19825E=00 0+34261E=02 0+51868E=04 0+48423E=06 0+30007E=08 0+14601E=10		0.12200E-00 0.57089E-02 0.97006E-04 0.97841E-06 0.59843E-08 0.24126E-10		0:19907E=01 0:34790E=03 0:34683E=05 0:22243E=07 0:92594E=10 0:81768E=12		0.54439E-0: 0.14503E-0: 0.15993E-0: 0.10719E-0: 0.40699E-10)
				:	x = 0.35	i				x = 0,65							
3 4 5	:	0:94152E=02 0:52944E=04 0:16194E=06 0:28918E=09 =0:39946E=12 =0:54771E=15	000	0 13829E-01 0 95938E-04 0 31100E-06 0 56607E-09 0 11595E-11 0 12522E-14	-0	0.30770E+03 0.10630E+05 0.20874E+08 0.26265E+11 0.22187E+11 0.23469E+12	0	0+14701E=03 1+53129E=06 +10696E=08 +11912E=10 +19993E=11 +76513E=13	0 0 0	.14648E-00 .46854E-02 .79121E-04 .83581E-06 .98114E-08 .25064E-10	0	0.13572E-00 0.76632E-02 0.14726E-03 0.14265E-05 0.11610E-07 0.54582E-10		0-26893E-01 0-53202E-03 0-59747E-05 0-69747E-07 0-20251E-09 0-33609E-11	000	#63876E-02 #21421E-03 #27092E-05 #20879E:07 #11264E-05	
				×	= 0.40							x	- 0.90				
1 2 3 4 5 6	0	0.14397E-01 0.10336E-03 0.41170E-06 0.497239E-09 0.58730E-11	0.	20436E=01 18645E=03 78432E=06 19108E=08 15864E=11 71339E=14	0	160249E=03 127075E=05 16036E=06 59325E=11 13194E=11 27564E=11	-0. -0.	280001-07 199816-05 34994E-03 507916-11 31124E-12 14489E-12	0 a l	19487E-00 89021E-02 11781E-03 13923E-05 10868E-07 11034E-10	000	.14796E=00 .19099E=01 .21810E=03 .27024E=05 .21631E=07 .11619E=09	0	+35574E=01 +79398E=03 +99741E=05 +60900E=07 +44027E=09 +84770E=12	0 0 0	70795E=02 30743E=03 44607E=05 36156E=07 21213E=09 45274E=11	
				x =	0.45							х =	0.95				
1 2 3 4 5 5 5	0	021042E=01 01860E=03 993794E=06 023054E=08 034519E=11 014011E=13	0 • 3 0 • 1 0 • 9	8700E-01 3476E-03 7938E-05 5426E-08 3720E-10 2000E-13	0 • 0 0 • 1 0 • 5	10909E=02 11777E=05 9794E=07 0568E=10 7182E=12 9184E=12	0+3 0+9 0+3 0+4	90415-03 01275-05 87995-05 29035-10 96445-11 49515-12	0 • 8 0 • 1 0 • 2 0 • 1	2878E-00 3585E-02 7149E-03 1594E-05 1594E-07 219E-09	0 a 1 0 a 4 0 a 3	15629E-00 13087E-01 31598E-03 3658E-05 19002E-07 3889E-09	0.	46121E-01 11593E-02 14193E-04 14635E-06 89623E-09 52667E-11	0.7	2934E-02 2953E-03 0466E-05 8230E-07 7125E-09 8734E-11	
	*											x = 1	.00				
	0.3 0.1 0.7 0.3	29670E=01 \$1676E=03 19565E=05 1956E=08 19957E=10 4699E=13	0.56 0.37 0.14 0.30	646E-01 460E-03 349E-05 178E-07 274E-10 129E-12	0.12 0.50 0.12 0.92	1569E=02 921E=04 756E=07 302E=09 966E=11 142E=11	0 62 0 25 0 17 -0 98	2976-08 0716-05 1696-07 9016-10 9366-11 2326-12	0.245 0.245 0.356 0.343	148E-00 149E-01 149E-03 129E-05 129E-07 44E-09	0.14	6650E-00 6701E-01 4678E-03 8759E-05 3097E-07	0 e 1 0 e 2 0 e 1	8642E-01 6597E-02 5615E-04 5646E-05 7789E-08	0.10 0.11 0.65	996E=02 909E=03 947E=04 787E=06 816E=09 397E=11	

			x = 1.05		x = 1.55							
	n ar	\mathbf{a}_{n}^{i}	b_n^r	b <u>i</u>	a_n^r	\mathbf{a}_{n}^{i}	$\mathfrak{b}_{\mathbf{n}}^{\mathbf{r}}$	b <mark>i</mark>				
	1	01 0.21009E 03 0.62607E 05 0.10584E 07 0.11567E	-01 0.233596- -03 0.396746- -04 0.437235- -05 0.333526-	02 0.77737E-03 04 0.16546E-04 06 0.19827E-06 08 0.16120E-00		0 0.10367E- 2 0.85244E- 3 0.11931E- 5 0.76857E-	00 0.32813E= 02 0.12712E= 03 0.30368E= 05 0.50737E=	01 0.16670E-02 02 0.35326E-03 04 0.11211E-04 06 0.21077E-06				
			x = 1.10				x = 1.60					
	1 0.32625E-0 2 0.18169E-0 3 0.47719E-0 4 0.83239E-0 5 0.96940E-0 6 0.79071E-0	0 • 26075E== 0 • 85924E== 5 0 • 15956E== 7 0 • 19136E==	01 0+32283E=0 03 0+60148E=0 04 C+72683E=0 06 0+60974E=0	2 0.100795-07 6 0.204465-07 6 0.324725-06 8 0.288635-08	0.52833E 0 0.13702E-0 0.67147E-0 0.23103E-0 0.436196E-0 0.97667E-0	0 + 11323E + (2 0 + 10481E + (3 0 + 41988E + (5 0 + 10783E + (00 0:39991E=0 0:16811E=0 0:42820E=0 4 0:76222E=0	01 0.540828-03 02 0.435998-03 04 0.153468-04 06 0.311158-06				
			x = 1.15				x = 1.65					
	2 0.29110E=0 3 0.65082E=0 4 0.12559E=0 5 0.15726E=0	1 0.319478=6 3 0.11618E=6 4 0.23603E=6 6 0.31000E=0	0:43979E-0: 0:43979E-0: 0:49490E-0: 4 0:11609E-0: 6 0:10891E-0:	04127475-02 04353625-04 04519245-04 04506215-05	0.93904E 00 0.19957E-00 0.45851E-02 0.45851E-03 0.78269E-03 0.78269E-05 0.14463E-06	0.16377E-0: 0.12192E-0: 0.12785E-0: 0.54708E-0: 0.14962E-04	0.48175E=0 0.22020E=0 0.59714E=0 0.11305E=0	1 -0.12103E-02 2 0.52603E-03 4 0.20727E-04 5 0.45304E-06				
			x = 1.20				x = 1.70					
1 2 3 4 5	G-38589E-00 D-29111E-01 O-87646E-03 O-18040E-04 G-24972E-06 D-24391E-08	0.17899E-00 0.38656E-01 0.15499E-02 0.34313E-04 0.49115E-06 0.48899E-08	0.99068E-02 0.13087E-03 0.18788E-05 0.18795E-07	-0.87117E-02 0.15704E-02 0.50150E-04 0.81220E-06 0.86414E-08 0.57008E-10	0.53940E CC 0.16367E-00 0.10414E-01 0.3952E-03 0.10795E-04 0.21162E-06	0.16288E-00 0.12938E-00 0.12938E-00 0.12479E-01 0.70669E-03 0.20369E-04 0.41571E-06	0.57357E-01 0.28579E-02 0.82411E-04 0.16564E-09	-0.37256€-02 0.62649₹-03 0.27835€-04				
		,	c = 1.25			,	r = 1,75					
1 2 3 4 5	0.41261E=00 0.36370E=01 0.11667E=02 0.25921E=04 0.38900E=06 0.41225E=08	0 •17892E-00 0 • 6204E-01 0 •20405E-02 0 •49089E-04 0 •76321E-06 0 •82583E-08	0-14798E-00 0-78277F-02 0-18837E-03 0-29320E-05 0-31841E-07 0-24106E-09	-0:16947E=01 0:18806E=02 0:69798E=04 0:12447E=05 0:14467E=07 0:11160E=09	0.54212E 00 0.20956E-00 0.12868E-01 0.51095E-03 0.14763E-04 0.30611E-06	0.16225E-00 0.13534E-00 0.18539E-01 0.90530E-03 0.27934E-04 0.59971E-06	0.35568E-00 0.67498E-01 0.36766E-02 0.11262E-03 0.23992E-05 0.37664E-07	-0.152605-00 -0.715715-02 0.726325-01 0.363795-04 0.923285-06 0.158365-07				
		x	= 1.30			х	- 1.80					
1 2 3 4 5	0.43689E-00 0.45077E-01 0.15366E-02 0.36711E-04 0.59532E-06 0.66311E-06	0.17682E-00 0.54559E-01 0.26557E-02 0.69199E-04 0.11649E-05 0.13655E-07	0.16935E-00 0.10241E-01 0.26718E-03 0.44968E-05 0.52820E-07 0.446730E-09	-0.270025-01 0.218095-02 0.954205-04 0.167135-05 04296635-07 0429635-07	0.50273E 00 0.23621E-00 0.15623E-01 0.65594E-03 0.19954E-04 0.43808E-06	0:16182E00 0:1395TE00 0:22163E01 0:11508E02 0:37627E04 0:85576E05	0.37485E=00 0.478520E=01 0.46894E=02 0.15247E=03 0.34370E=05 0.57138E=07	-0.164745-00 -0.115595-01 0.819645-01 0.472995-04 0.12936-05 0.236585-07				
		x e	1.35			x = 1.85						
1 2 3 4 5 6	0.49899E-00 0.49942EE-01 0.20042E-02 0.51902E-04 0.89613E-06 0.11086E-07	C.174835-00 0.636406-01 0.34188E-02 0.96220E-04 0.17486E-05 0.22131E-07	0.19105E-00 0.13236E-01 0.3738ZE-03 0.67779E-05 0.85905E-07 0.78510E-09	-0.38703E-01 0.74345E-01 0.12821E-01 0.27638E-05 0.37986E-07 0.36445E-09	0.54156E 00 0.24333E=00 0.19366E=01 0.493652E=03 0.24777E=04 0.62054E=06	0.16151E-00 0.14195E-00 0.2622E-01 0.14921E-02 0.50239E-04 0.12085E-05	0.39388E-00 0.90316E-01 0.57314E-02 0.20458E-03 0.48750E-05 0.85616E-07	-0.17570E-00 -0.17074E-01 0.89489E-03 0.60719E-04 0.17921E-05 0.34927E-07				
		x =	1.40			х =	1.90					
3 4 5 6	0.47766E-00 0.47601E-01 0.29910E-02 0.70814E-04 0.13265E-05 0.17679E-07	0.17261E-00 0.73115E-01 0.43364E-02 0.13210E-01 0.25846E-05 0.35213E-07	0.21279E-00 0.15902E-01 0.51643E-03 0.10065E-04 0.13716E-06 0.13555E-08	+0.6517775-01 0.258925-02 0.169395-03 0.401385-05 0.597895-07 0.614695-09	0.53874E 00 0.29039E-00 0.23597E-01 0.10602E-02 0.35645E-04 0.87059E-06	0 +16123E-00 0 +14251E-00 0 +30806E-01 0 +18194E-02 0 +66519E-04 0 +16901E-05	0.41280E-00 0.10275E-00 0.74410E-07 0.27220E-03 0.68462E-05 0.12687E-06	-0:18537E-00 +0:23728E-01 0:93588E-03 0:76902E-04 0:24545E-05 0:36926E+07				
		x = 1	1.45			x = 1	95					
2	0:49410E=00 0:81764E=01 0:39226E=02 0:96634E=04 0:19618E=05 0:27730E=07	0 + 17038E-00 0 + 83390E-01 0 + 54979E-02 0 + 17922E-03 0 + 37659E-05 0 + 55120E-07	0.23433E=00 0.21332E=01 0.70500E=03 0.14735E=04 0.21546E=06 0.23153E=08	-C.49806E-01 0.25792E-07 0.22009E-03 0.57370E-05 0.92462E-07 0.10463E-01	0-59439E 00 0-31691E-00 0-28625E-01 0-1339E-02 0-67067E-04 D-12102E-05	0.16089E-00 0.14137E-00 0.35910E-01 0.22642E-02 0.87375E-04 0.23417E-05	0.43163E-00 0.11570E-00 0.92595E-02 0.35924E-03 0.95259E-05 0.18601E-06	-0.193645-00 -0.315195-01 0.921345-03 0.963615-04 0.332555-05 0.735035-07				
		x = 1,	50		x = 2.00							
	0.50796E 00 0.98035E-01 0.42289E-02 0.13047E-03 0.28008E-05 0.42792E-07	0.16830E-00 0.93612E-01 0.68736E-02 0.24047E-03 0.54138E-05 0.84874E-07	C+25553E-00 C+2661E-01 C+95171E-03 C+21266E-C4 C+33299E-C6 C+36173E-08	-0.80658E-01 0.23037E-02 0.2812AE-03 0.80781E-05 0.14071E-06 0.16986E-00	0.34245E-00 0.34569E-01 0.16742E-02 0.41755E-04	0-16040E-00 0-13873E-00 0-41533E-01 0-27997E-02 0-11590E-03 0-32159E-05	0:45037E-00 0:12980E-00 0:12980E-01 0:47045E-03 0:13137E-04 0:27802E-06	-0.20045E-00 -0.40403E-01 0.62299E-07 0.11994E-03 0.44576E-05 0.10477E-06				

					x	= 0.05							,	c = 0.55		
		n a	r n		\mathbf{a}_{n}^{i}	ь	r n	b _n i	ь <mark>і</mark> n		ar n			P.		b <u>i</u>
	2 3 4 5	0.3132 -0.30563	1E-08	0 = 4 3 8 0 = 50 7 0 = 8 4 6	/4E-04 09E-0d 07E-17	0 • 1 7 7 1 0 • 1 2 5 7 0 0 • 1 8 7 0 1	0E − 0∀	0+10391r =0+302+25 0+394415	= -07 = -10 = -11	0.40348 0.42865 0.36097 0.16055 0.42448 ~0.15002	5E-03 F-03 F-05 E-07	0.96067	E-01 E-03 E-05 E-07	0.33166 0.275968 0.131078 0.424666 -0.163878 -0.199115	-07 -04 -06 -07 -12	0.145745-02 0.130865-04 0.692445-07 0.243135-09 0.104145-10
					x =	0.10							x =	0.60		
	1 2 3 4 5 6	C+196551 C+759865 C+144425	-07	C + 3507 C + 1953 C + 1032	0E-05 4E-09	0.63023 0.5256 0.70921	-09	0-47370:-10		0.541579 0.758545 0.662558 0.349848 0.109338 -0.615858	-03 -05 -07	0.58837E 0.14777E 0.13997E 0.76642E 0.27211E 0.10250E	-02 -04 -07 -09	0.51638E 0.50811F 0.28647F 0.1024E- -0.10424E- +0.36152E-	-04 -05 -09 -10	0.213878-02 0.252405-04 0.459985-00 0.611655-09 0.893185-11 -0.112705-11
	1	4 ((6335				0.15						X =	0.65			
	1 0.6667325-00 2 0.726605-06 3 0.428445-09 5 6			C+11338-92 O+14809E-05 O+d5252E-09		0+478925-05 0+299305-08 0+909345-11		0+263/46-0 0+15+276-0 -0+337206-1) a	0.71058E=01 0.11384E=02 0.11582E=04 0.71732E=07 0.28365E=09 =0.77144E=12		0.84440E=01 0.21941E=02 0.24387E=04 0.15673E=06 0.65305E=09 0.16953E=11		0.776198-0 0.891448-0 0.588268-0 0.249208-0 -0.103448-1 -0.260848-1	14 16 12 1	0.29958E-02 0.43274E-04 0.30411F-06 0.13663E-07 0.10177E-16 0.53564F-12
					x = 0	x = 0.20							x = 0	.70		
	1 2 3 4 5 5	0+10071E-0 0+30636E-0 0+30561E-0 -0+79905E-1 -0+12315E-1 -0+333926E-1	5 8 2 4	0.28047E- 0.52325E- 0.55923E- 0.55260E- 0.31726E- 0.86715E-	05 09 11	0.202/15-0 0.230955-0 0.182/15-1 -0.43209E-1 -0.82221E-1 -0.55953E-1	7 0 1 2	0.12017F-04 0.12759E-07 0.33865-10 0.52029E-11 0.24408F-12 0.31019F-15	,	0.91277E=0 0.16604E=0 0.19427E=0 0.13931E=0 0.04491E=0 =0.49126E=1	02 04 06 19	0.10087E-0 0.31608E-0 0.40746E-0 0.10382E-0 0.14637E-0 0.31608E-1	2 4 6	0.11316E-0 0.15005E-0 0.11453E-0 0.56054E-0 0.41542E-11	3 5 1	0.401395-01 0.709785-07 0.584265-05 0.302335-08 0.15215E-10 0.204235-11
	1	0.110015-01			x = 0.2								x = 0.7	75		
	2 3 4 5 6	0.31925E-02 0.93560E-05 0.14621E-07 0.52801E-10 -0.14247E-13 -0.61584E-17		0.54597E= 0.18986E= 0.31184E= 0.24795E=1 0.37023E=1 0.15738E=1	04 07 10	0:61992E=0: 0:10949E=3: 0:99214E=10 0:17981E=10 0:10421E=10 0:64305E=14	5))	0.332325=04 0.579715=07 =0.409845=10 =0.19465=10 0.625415=11 0.363225=14		J-11490E-0 7-23637E-0 0-21474E-0 0-25839E-0 0-13925E-0 0-44665E-11	2 4 6	0.11754E-00 0.44359E-02 0.45659E-04 0.5622E-06 0.31120E-08 0.99748E-11		0.16057E=01 0.24371E=01 0.21295E=05 0.11979E=07 0.36787F=16 0.16998E=11	0	-51436E-02 -11197E-03 -10703E-05 -62997E-04 -25546E-10 -18453E-11
					x = 0.30)						,	= 0.80)		
1 2 3 4 5 6		Q=56334E=02 O=23306E=04 O=52242E=07 O=47885E=10 O=10305E=12 O=65733E=16	0	0.94229E=0: 0.47145E=0: 0.47145E=0: 0.17432E=0: 0.26795E=12: 0.16792E=15	-0	-15495E-03 -39435E-06 -63780E-09 -12362E-10 -81976E-11 -47947E-13		0.914635-0/ 9.214205-06 0.444015-09 0.397845-10 0.371815-11 0.255325-11		0.14181E-00 0.32962E-02 0.49285E-04 0.46030E-05 0.28327E-08 0.411389E-10		0.13175E-00 0.500041E-02 0.10252E-03 0.99922F-06 0.4304BE-08 0.27082E-10	0	0.22230E=01 0.36370E=03 0.36036E=05 0.24369E=07 0.10438E=09 0.10419E=11	0. 0. 0.	528865-02 170625-03 168105-05 128215-07 609885-10 647405-13
					x = 0.35							*	= 0.85			
3 6	-	0.91697E-02 0.50440E-06 0.15349E-06 0.27766E-09 0.91185E-11 0.48548E-15	0	+14886E=01 +10164E=03 +12801E=06 +61438E=09 +11456E=11 +12386E=14	0. 0. -0.	33668E-03 11596E-05 22647E-08 57625E-11 35203E-12 14472E-12	000	0-172325-03 0-623985-04 0-122595-04 0-62515-11 0-135585-11 0-779925-13		0:171675-00 0:451436-02 0:752186-04 0:791405-05 0:549246-03 0:235216-10		0.14881E-00 0.81763E-07 0.15568E-03 0.17138E-05 0.12221E-07 0.59693E-10	0. 0.	30060E-01 58778E-03 65591E-05 47405E-07 22909E-09 35788E-11	042 043 041	29315-02 51925-01 18675-05 43025-07 40155-09 12705-11
				×	= 0.40							x =	0.90			
123456	0 0 -0	•14078E-01 •98520E-04 •39032E-06 •90520E-09 •12670E-11 •27583E-14	0 • 2 0 • 2 0 • 4	22043E-01 19760E-03 89279E-06 80440E-08 4680E-11 80036E-14	0:2 0:7 2:7 -0:6	66037E=03 9565E=05 95373E=09 6474E=11 5278E=11 2986E=12	0 e 0 e	329275=05 157205=05 414205=06 367755=16 434205=11 182145=12		0-20393E-00 0-50865E-02 0-1326E-03 0-13187E-05 0-10265E-07 0-52500E-10	0000	0.16209E=00 0.10789E=01 0.23066E=03 0.26482E=05 0.22779F=07 0.12568F=09	0 • 5 0 • 1 0 • 8	398785-01 378775-01 109625-04 187005-07 88950-09 08415-11	0 • 36 0 • 52 0 • 44 0 • 27	7342E=02 134E=03 241E=05 911E=07 767E=09 441E=11
				x a	0.45							x = (0.95			
2 2 3 4 5 5 5	0. 0.1	20679E=01 17799E=03 88894E=06 6380E=08 11013E=12 .2702E=13	0.15	1022E-01 5493F-03 3926F-05 1626E-08 246E-10 095E-13	0.67 0.21 0.54 -0.80	780E=02 512E=05 703E=07 079E=10 466E=11 787E=11	0.3 0.1 0.7	76845-03 54035-05 17035-07 49565-10 07325-10 31375-12	0 0 0	#237845-00 #309505-07 #163425-03 #213565-05 #185335-07 #110945-09	0.	17312E-00 14001E-01 33401E-03 46017E-05 41020E-07 25174E-09	0 • 12 0 • 17 0 • 16 0 • 93	639E=01 855E=02 815E=04 035E=06 826E=09 160E=12	0.831	
	x = 0.50 0-293085=01											x = 1.	00			
	0.0 0.6 0.1	93085-01 02275-03 95546-05 92376-08 94315-10 71485-13	0.596 0.396 0.149	873E-01 891E-03 814E-05 870E-07 826E-10 87E-12	0.141 0.555 0.142	37F=02 36F=04 56F=07 25F=09 32E=17 49F=11	0+72 0+29 0+75 0+73	380E-03 962E-05 668E-07 556E-10 535E-12 428E-12	0 • 0 •	272365-00 106375-01 233785-03 337635-05 324615-07 219775-09	0 • 1 0 • 4 0 • 7	8159E=00 7895E=01 7503E=03 2490E=05 1657E=67 9459E=09	0.184 0.282 0.281 0.195	590E-01 439E-02 238E-04 33E-06 77E-08 66E-11	0.6920 0.6559 0.1287 0.1386 0.1004 0.5769	7E-01 9E-04 0E-06 7E-08

					,	= 1.05										
		n	\mathbf{a}_{n}^{r}		a i		br n	ı	i O		.r			x = 1.55		
		2 G-19 3 0-32 4 0-92 9 0-55	720E-00 828E-01 876E-03 154E-05 212E-07	0.22 0.66 0.11 0.12	746E=00 549E=01 301E=03 161E=04 175E=06 222E=09	0.254 0.437 0.480 0.364 0.712	0446-01 788-02 978-04 978-04 958-08 858-08 778-11	0 - 4562	8E-02 7E-05 4E-04 IE-06	0:130 0:513 0:166 0:376	174E 00 195E+00 195E+00 195E-03 197E-05 62E-07	0.11 0.91 0.33 0.81	ai 086E-0 314E-0 001E-0 782E-0 30E-0 58E-0	0 0.367 2 0.142 3 0.336 5 0.559	hr n 42E-00 47E-01 18E-02 80E-04 68E-06 18E-08	-C-11291E-0 G-1201E-0 O-40920E-0 C-19161E-0 C-24796E-0(C-32562E-0)
		1 0.160	92 E-00			1.10								x = 1.60		
		2 0.178: 3 0.499;	03E=01 21E=03 84E=05 72E=07	0.2801 0.9104 0.1683 0.2016 0.1645	3E-03 0E-04	049962 043600 046642 047987 046700 042746	1E-02 4E-04 9E-06 4E-08	0451303 0411758 0428751 0438203 0438203 043828	-02 -04 -06	0e5413 0e1413 0e6518 0e2204 0e5331 Ce9238	3E-00 3E-02 2E-03 1E-05	0.168 0.123 0.1120 0.444 0.1137 0.2029	2E-00 2E-01 3E-03 1E-04	0.3090 0.6466 0.1882 0.4754 0.84153 0.10993	9E-01 7E-02 8E-04	-0.128785-00 -0.389805-03 0.502415-03 0.180005-04 0.385975-05 0.513365-05
	1	0.37303	E-00	0,19208									×	- 1.65		
	3 4 5	0.62155	E-03 E-04 E-06	0.34415 0.12317 0.24903 0.326348 0.298048	E-01 E-02 E-04	0:11924; 0:491338; 0:989608 0:12988; 0:119058 0:700586	-02 -04 -05 -07	-0.549235- 0.148065- 0.415765- 0.610865- 0.610885- 0.423885-	02 04 06	0:54557 0:14529 0:461694 0:29973 0:743016 0:134856	E-00 E-02 E-03	0 = 16661 0 = 13260 0 = 13682 0 = 37939 0 = 15782 0 = 299860	E-00 E-01 E-03	0+32836: 0+53617: 0+24689: 0+663818: 0+124938: 0+173618	-01 -02 -04 -05	-0.14400E-00 -0.27698E-02 0.60439E-02 0.24285E-04 0.53268E-08 0.80152E-08
	1 2	0.402958 0.287538	-00	0 - 1 92 3 9 €	-00	0+14009E-	.00						x =	1.70		
	5 6	0.89799E 0.89799E 0.17121E 0.23626E 0.23003E	-03 -04	0.41725E 0.16438E 0.36212E 0.51712E 0.51455E	-01 -02 -04 -06	0.66101E- 0.14492E- 0.20682E- C.20632E- 0.14008E-	02 03 05	-0,19995-0 04181325-0 0,989305-0 04955565-0 04101895-0 04742175-10	2	0.34452E 0.19119E 0.10171E 0.37781E 0.10251E 0.20030E	-00 -01 -03	0.16511E 0.14079E 0.16582E 0.74672E 0.21676E	-00 -01 -03	0.34734E- 0.63550E- 0.32075E- 0.91716E- 0.18321E- 0.27030E-	01 02 04	-0:158345-00 -0:608405-01 0:710775-02 0:32355-04 0:765035-06 0:123115-07
	1	0.43030E-	20	0,18966E-0	x = 1.25								x = 1	.75		
	3 4 5 6	0.34105E-0 0411149E-0 0.24615E-0 0.36815E-0 0.38941E-0) 1) 2 4	0.49975E-0 0.21662E-0 0.51822E-0 0.60369E-0 0.66901E-0	01 0 2 0 4 0 6 0 8 0	-16172E-0 -87725E-0 -20888E-0 -32308E-0 -34982E-01 -26815E-09	9	-0.23785E-01 0.21532E-02 0.81957E-04 0.14642E-05 0.17021E-07 0.13672E-09		0.55036E C.21841E-0 0.12611E-0 0.48901E-0 0.14007E-0 0.28983E-0	11	0.16398E- 0.14686E- 0.19947E- 0.9997ZE- 0.29480E- 0.63151E-	00 21 33	0.36613E=0 0.74385E=0 0.41297E=0 0.12548E=0 0.26562E=0 0.41529E=0	i . 2 3	-0.17160E-06 -0.10466E-01 0.61432E-03 0.42496E-04 0.10648E-05 0.18643E-07
	1	0:454825-0:	, ,	•18689E-00	x = 1.30	183775-00							x = 1.8	0		
	3 6	0.44490E=0: 0.14730E=0: 0.34875E=0: 0.54354E=0:6 0.64487E=0:8	0000	.59131E-01 .28215E-02 .73074E-04 .12269E-05 .14365E-07	0. 0.	18372E-00 11491E-01 29670E-03 49376E-05 56095E-07 48767E-09	0	0035901E-01 0024671E-02 0011193E-03 0022012E-05 002704E-07 027864E-07		0.54925E 00 0.24460E-00 0.15564E-01 0.62855E-03 0.18969E-04	1 4	0.14314E-00 0.15046E-00 0.23817E-01 0.12207E-02 0.39723E-04 0.90129E-06	,	0.38481E-00 0.86001E-01 0.92704E-02 0.17008E-03 0.36093E-05 0.63040E-07	-000	0.18362E-00 -16006E-01 -90410E-03 -55192E-04 -15194E-05 -27667E-07
	,	0-47639E-00			- 1.35								- 1.85		Í	
4 9	!	0-55430E-01 0-19244E-02 0-48745E-04 0-84854E-06 0-10478E-07	0 • · 0 • · 0 • ·	18361E-00 59108E-01 34352E-02 10164E-03 8420E-05 3284E-07	0.1 0.4 0.7 0.9	0570E=00 4863E=01 1574E=03 4838E=05 482E=07 482E=09	0. 0. 0.	49649E-01 27057E-02 15018E-03 32509E-05 44680E-07 43516E-01	į	0.54634E 00 0.27512E-00 0.19126E-01 0.60267E-09 0.25471E-04 0.58794E-06	0.	16250E-00 15275E-00 28227E-01 15412E-02 53053E-04 12731E-05	0	•403425-00 •982505-01 •866895-07 •228485-09 •340795-05 •945285-07	-0. 0. 0.	19429E-00 22765E-01 9642E-03 70605E-04 21033E-05 1116E-07
1		0+49500E-00	0.18	 1016E-00		716E-00						x :	1.90			
5 6	0	0:462416-01 0:249226-02 0:673546-04 0:129846-05 0:167046-07	0.79 0.46 0.13 0.27	7995-01 3625-02 9595-03 2325-09 0515-07	0.181 0.575 0.111 0.150	966E-01 18E-03 29E-04 96E-06 84E-08	0 a 2 i 0 a 1 i 0 a 4 i 0 a 7 i	4647E-01 8002E-02 9806E-02 7193E-02 1370E-07 893E-00	0.	94178E 00 90998E=00 929404E=01 10188E=02 39928E=04 82919E=06	0 • 1 0 • 1 0 • 7	4197E-00 9292E-00 9209E-01 9329E-02 9267E-94 7407E-05	0 ± 5 0 ± 6 0 ± 3	42200E=00 11098E=00 31650E=02 10432E=03 6014E=05 4018E=06	0.9°	0991E-00 1742E-01 7384E-00 7384E-05 788E-05 967E-07
1	0.	51069E 00		Z = 1	1.45							X.	1 95			
2 3 4 5	0.0	.09015E=01 .32021E=02 91974E=04 18400E=05 26212E=07	0.938 0.585 0.189		0.2486 0.2395 0.7863 0.1630 0.2373 0.2546	6E=01 4E=03 5E=04 2E=0A	0.256 0.256 0.674 0.108	469E-01 509E-02 579E-09 127E-05 79E-06 04E-08	0.2 0.1 0.4	19973E 00 19076E+00 18917E+01 2896E+02 4892E+04 1475E+05	0.19 0.38 0.24	1436-00 0386-00 7436-01 0446-02 9296-04 9776-05	0.12 0.10 0.40 0.10	0985-00 14025-00 14045-01 2065-01 9875-04 9665-05	0.904 0.111 0.389	1216-00 1865-01
:		2357E 00	0 - 1 734	# = 1.5 4E-00	-	•						x = 2.	00			
•	0.1 0.4 0.1	0009E-00 0849E-02 2427E-03 8950E-03 0470E-07	0.1021 0.7331 0.25430 0.57064 0.69330	0E=00 BE=02 DE=01 E=09	0.26927 0.29855 0.10630 0.23581 0.367148 0.421988	E=01 E=02 E=04 E=05	0.9564 0.2173 0.3270 0.9488 0.1459 1.2026	3E-02 5E-03 9E-05 1E-06	0.39	894E 00 484E=00 998E=01 141E=02 867E=04	0.146	0E-03	0.137 0.128 0.527	24E-00 31E-01 03E-03 15E-04	0.2179 0.5009 0.7192 0.1968 0.5217	7E-01 4E-09 9E-09 9E-0f

					x = 0.05						x = 0.55	
		n a	r n	í	a ⁱ n	h <u>r</u>	$\mathfrak{b}_{\mathbf{n}}^{\mathbf{i}}$		ar n	a		pr pi
	3	0.2562	IE-08	0.40501 0.54714 0.93394	6E-08 0.15	900E-07 069E-09 063E-10		-09 0.364 -10 0.286 9.127 0.355	77E-01 00E-03 20E-05 95E-07 61E-10 44E-12		9E-01 0.2598 0E-03 0.2128 0E-05 0.1036 0E-07 0.3405 0E-09 0.3062	12E-0? 0:19409E-0; 16E-04 0:17983E-04 17E-06 0:84426E-07 13E-09 0:25089E-09 19E-11 0:19279E-10
					x = 0.10						x = 0.60	
	3 4 5 6	0,19244 0,74904 -0,704298	-07	0-32420 0-17 99 5 0-7 96 626	E-04 -0.907 E-10 -0.194	09E-04 06E-09 02E-09	0.41622E- 0.76848E- -C.11647E-	10 045957	9E-03 4E-05 4E-07 4E-10	0.650248 0.136466 0.129146 0.70741E 0.26720E 0.12990E	-02 0.39291 -04 0.22034 -07 0.84950	1E-04
	1				x = 0.15						x = 0.65	
	3 4 5 6	0.91992E- 0.9732PE- 0.99404E-	-04	0.10947E- 0.19669E- 0.7722EE-	-05 0.2299	80-38	0.31576E=0 0.16426E=0: -0.50169E=1:	0.89335	E-03 E-05 E-07 E-09	0.81579E- 0.20268E- 0.22508E- 0.14479E- 0.81596E- 0.29962E-	02 0.69028E 04 0.45277E 06 0.19337E 09 -0.47914E	-04
					x = 0.20						x = 0.70	
	2 3 4 5 6	0a12467E-0 0a24167E-0 0a24763E-0 -0a93600E-1 -0a12012E-1 -0a33169E-1	5 6 2	0.23995E-0 0.57923E-0 0.60991E-0 0.85661E-1 0.93352E-1 0.96525E-1	0-17606 18 0-70440 1 0-99281 4 -0-83211	E-07 E-10 E-11 E-12	0.13259E-0.0 0.15110E-0.7 0.20273E-10 0.10072E-10 0.11191E-11 0.47423E-15	0.72050E 0.15019E 0.15381E 0.11089E 0.52050E -0.16012E	-02 -04 -06	0.98637E-0 0.29212E-0 0.57408E-0 0.26032E-0 0.13734E-0 0.65710E-1	2 0.11651E- 4 0.48246E- 4 0.49459E- 8 0.49014E-	03 0.91160E-04 06 0.72416E-06 08 0.37267E-08 11 0.42053E-10
					x = 0.25						x = 0.75	
3		0.24753E-02 0.73778E-05 0.11596E-07 0.14759E-10 -0.13659E-13 -0.40196E-17	0	•50674E=02 •17921E=04 •28842E=07 •25930E=10 •41634E=13 •17537E=16	0.81968E 0.89057E 0.53662E	-07 -10 -11 -11	0.40247E-04 0.72141E-07 0.47880E-10 -0.42352E-11 0.14544E-12 0.56196E-14	0+92020E- 0418320E- 0+24878E- 0+10541E- 0411158E- 0+19602E-	02 04 06	0-11656E-00 0-41020E-02 0-41020E-02 0-51903E-06 0-26917E-08 0-13651E-10	0-189825-0	0-145916-05 5 0-199968-05 8 0-772888-08
					x = 0.30					1	x = 0.80	
3 4 5		0.49663E-02 0.18371E-04 0.41356E-07 0.54347E-10 -0.99647E-13 -0.66248E-16	0 • 0 • 0 •	87441E-02 49507E-04 10328E-06 17504E-09 30036E-12 18712E-15	0.11902E- 0.30207E- 0.56904E- 0.17328E- -0.37016E- -0.46234E-]	06 99 10	0.99977E-04 0.29771E-06 0.49239E-09 0.45236E-10 0.10132E-10 0.39899E-13	0.11478E-0 0.29802E-0 0.39899E-0 0.34811E-0 0.22558E-0 0.55399E-1	2	0.13477E-00 0.56304E-02 0.94429E-04 0.92247E-06 0.58407E-08 0.28550E-10	0=18376E-01 0=29995E-03 0=29995E-03 0=1881E-07 0=83542E-10 -0=41575E-11	0-22599E-03 0-2363E-05 0-13596E-07 0-94444E-10
				x	0.35					х	= 0.85	
2 9 4 5 6	-	0.71070E-02 0.39743E-04 0.12189E-06 0.23984E-09 0.35513E-12 0.47480E-15	0.9 0.3 0.5 0.1	3845E-01 3803E-04 0275E-06 8268E-09 5792E-11 9832E-14	0.25924E-0 0.88894E-0 0.18195E-0 0.24928E-1 0.1864E-1 -0.21116E-1		0.21357E-03 0.73460E-08 0.14737E-08 0.14737E-10 0.74676E-11 0.20690E-12	0.14045E-00 0435309E-02 0.59478E-04 0442926E-06 049879E-04 0-18623E-10	0	0-15263E-00 0-75762E-02 0-16371E-03 0-15821E-05 0-11292E-07 0-58865E-10	0-29304E-01 0-46139E-03 0-50786E-05 0-36950E-07 9-18692E-09 -0-16860E-11	0.13965E-01 0.33997E-03 0.40613E-03 0.3004E-07 0.17201E-09 0.52210E-11
	_				0.40					x =	0.90	
2 3 4 5 6	-0.	•10916E-01 •7796E-04 •9097EE-06 •73978E-09 •13778E-12 •26790E-14	0.18: 0.766 0.186	\$46E-01 297E-03 056E-06 099E-08 194E-11 195E-14	0.51003E=03 0.22684E=03 0.58131E=08 0.26184E=10 -0.14948E=12 -0.61802E=12	0 . 0 . 0 .	41243E-05 19119E-05 49103E-08 30879E-10 78466E-11 62324E-12	0.16939E-00 0.47574E-02 0.48562E-04 0.10462E-05 0.81912E-08 0.38921E-10	0. 0. 0.	16946E-00 10007E-01 21295E-03 24293E-05 21057E-07 12303E-09	0.34188E-01 0.69309E-03 0.85051E-05 0.68494E-07 0.39194E-09 -0.34691E-11	0.16706E=01 0.69023E=03 0.68688E=03 0.53896E=07 0.53999E=09 0.11003E=10
				x = (0.45					x = (0.95	
2	0.	16049E-01 16006E-03 70923E-06 21240E-08 74113E-11 2170E-13	0.327	96-19 96-19	0.92878E-03 0.51846E-05 0.16623E-07 0.59178E-10 0.17547E-11 -0.10017E-11	0.4 0.1 0.5	79498E-09 93323E-09 4039E-07 1872E-10 93400E-11 4768E-11	0.20039E-00 0.69236E-02 0.12907E-03 0.16978E-03 0.16978E-07 0.03910E-10	0.1 0.3 0.6 0.3	18469E-00 3008E-01 1086E-03 2482E-05 7905E-07 4024E-09	0-49953E-01 0-10194E-02 0-13854E-04 0-12993E-06 0-77653E-09 -0-29668E-11	0.1984E-01 0.71292E-03 0.10751E-04 0.10038E-04 0.4079FE-09 0.11972E-10
				x • 0.						x = 1.	00	
	0.2: 0.1: 0.54	2795E-01 3775E-03 6725E-05 6451E-08 6716E-10 6289E-13	0.3944 0.5546 0.36376 0.19816 0.36087	5E-03 6E-05 DE-07 PE-10	0.15918E-02 0.10876E-04 0.42629E-07 0.12101E-09 0.41011E-11 0.58884E-12	0.81 0.35 0.66 0.13	2264E-02 1932E-05 1966E-07 1243E-10 749E-11 479E-12	0.23314E-00 0.8306E-02 0.18453E-03 0.25618E-05 0.25674E-07 0.17485E-09	0.16 0.43 0.46 0.66	749E-00 459E-01 648E-03 923E-09 147E-07 832E-09	0.59084E-01 0.14713E-02 0.22013E-04 0.21755E-04 0.13119E-08 0.73415E-11	0.22715E-01 0.77791E-03 0.16845E-06 0.17462E-08 0.1249E-08

			x = 1.05				x = 1.55	
	n an		i br	bl	ar n	a _n	$\mathfrak{b}_{\mathbf{r}}^{\mathbf{r}}$	$\mathbf{b_n^i}$
	1 0.246691 2 0.107981 3 0.259918 4 0.414195 5 0.440278 6 0.329658	E-01 0.21046 E-03 0.61240 E-09 0.10304 E-07 C.11244	E-01 G.208778 E-09 0.942088 E-04 0.371638 E-06 0.289248	E-02	2 0.973778 0.405398 0.131378 0.300418	-01 0.11475E -02 0.64461E -03 0.91239E -09 0.74890E	-00 0.342456 -02 0.116446 -09 0.267296 -09 0.497796	-01 0:10894E-01 -02 0:69063E-03 -04 0:18649E-04 -06 0:32919E-06
			x = 1.10				x = 1.60	
	1 0.30066E 2 0.13906E 3 0.3380E 4 0.5380E 5 0.73052E 6 0.5922EE	-01 C.28290E -09 0.84115E -09 0.1999e -07 0.18420E	-01 0.29169E- -03 0.92089E- -04 0.61962E- -06 0.51946E-	-02 0.18402E-02 -04 0.38409E-04 -04 0.48814E-04 -08 0.42494E-08	0.51099E 0.11647E- 0.51040E- 0.17406E- 0.42357E- 0.73604E-	-00 0.12747E- -02 0.10449E- -03 0.41102E- -09 0.10505E-	00 0.42655E- 01 0.15601E- 03 0.37863E- 64 0.69959E-	01 0.11994g=01 02 0.88801g=03 04 0.29946g=04 06 0.49036g=04
			x = 1.15				x = 1.65	
,	1 0.99985E=0 2 0.1779F=1 3 0.48955E=0 6 0.99921E=0 5 0.11847E=0 6 0.10963E=0	01 0.12349E- 09 0.11363E- 09 0.22994E- 06 0.30139E-	01 0.40183E-0 02 0.77858E-0 04 0.10093E-0 06 0.92207E-0	02 0.24284E-02 04 0.58494E-04 05 0.78498E-08 06 0.74491E-08	G.59993E	00 0.19972E-0 02 0.12702E-0 09 0.59400E-0 09 0.14982E-0	00 0:52009E-0 01 0:20647E-0 0:3079E-0 0:90147E-0	01 0.11491E-01 2 0.11276E-02 4 0.39471E-04 4 0.72099E-04
			x = 1.20				x = 1.70	
1 2 3 4 5	0.36571E=00 0.42508E=0 0.659951E=01 0.13577E=04 0.18814C=00 0.18924E=06	1 0.37409E-0 3 0.15198E-0 6 0.33439E-0 6 0.47753E-0	1 0.54697E-0: 2 0.11444E-0: 4 0.16103E-0: 6 0.15993E-0:	2	0.54517E 0 0.36197E-0 0.79646E-0 0.29797E-0 0.01341E-0 0.15991E-0	0 0.19102E-00 2 0.19932E-03 9 0.69299E-03 0 0.20090E-04	0.63811E-0 0.27099E-0 0.73653E-0 0.14428E-05	0-11064E-01 0-14198E-02 0-68486E-04 0-10467E-05
			x = 1.25				x = 1.75	
1 2 3 4 5 6	0.395676-00 0.283246-01 0.876316-03 0.195096-06 0.275066-06 0.310136-08		0+73392E-02 0+16563E-03 0+25206E-05	0.19946E-01 0.40030E-02 0.11587E-03 0.19219E-05 0.21424E-07 0.18686E-09	G-54867E GO G-18799E-GO G-788G1E-G2 G-38547E-G3 G-11111E-G4 G-23079E-G4	0:18554E-00 0:16092E-00 0:18741E-01 0:88874E-03 0:27244E-04 0:58339E-06	0.40736E-00 0.74600E-01 0.35282E-02 0.10124E-03 0.20971E-05 0.32437E-07	-0.13824E-00 0.99997E-02 0.17900E-02 0.69190E-04 0.13015E-05 0.24607E-07
		,	T = 1.30			x	1.80	
1 2 9 4 9	0.42331E-00 0.35401E-01 0.11579E-02 0.27651E-04 0.44647E-06 0.51402E-08	0.22428E-00 0.56569E-01 0.26114E-02 0.67492E-04 0.11330E-05 0.13262E-07	0.19457E-00 0.97461E-02 0.23634E-03 0.38749E-05 0.45067E-07 0.37725E-09	0.48149E-02 0.49979E-02 0.18149E-03 0.29142E-05 0.39607E-07 0.33239E-09	0.234996E 00 0.21573E-00 0.12205E-01 0.49517E-03 0.15040E-04 0.33021E-06	0.18071E-00 0.18877E-00 0.22458E-01 0.11911E-02 0.34719E-04 0.63266E-06	0.42809E-00 0.90793E-01 0.45583E-02 0.13792E-03 0.30158E-05 0.49308E-07	-0.19026E-00 0.64260E-02 0.21976E-02 0.064732E-04 0.21296E-05 0.37037E-07
		x	= 1.35			x ·	1.85	
1 2 3 4 5	0.44832E-00 0.43959E-01 0.15110E-02 0.38816E-04 0.67506E-06 0.83496E-08	0.22194E-00 0.66670E-01 0.39669E-02 0.93890E-04 0.17011E-05 0.21516E-07	0+22217E-00 0+12801E-01 0+33284E-03 0+38646E-05 0+73411E-07 0+64492E-09	-0.46740E-02 0.61178E-02 0.22214E-03 0.43975E-05 0.57479E-07 0.54567E-09	0.59919E 00 0.26474E-00 0.15020E-01 0.69199E-00 0.20186E-04 0.46779E-06	0-17591E-00 0-17483E-00 0-26731E-01 0-14290E-02 0-49032E-04 0-11762E-09	0.44623E-00 0.10609E-00 0.58451E-02 0.18630E-09 0.42935E-05 0.74086E-07	-0:16276E-00 0:19404E-02 0:25739E-02 0:1362E-03 0:29879E-05 0:55097E-07
		х -	1.40			х •	1.90	
1 2 3 4 5	0.47054E-00 0.54210E-01 0.19555E-02 0.53304E-04 0.10004E-05 0.13314E-07	0.21840E-G0 0.77707E-01 0.42978E-02 0.12897E-03 0.23190E-03 0.34239E-07	0.24948E-00 0.16635E-01 0.46306E-03 0.87446E-05 0.11790E-06 0.11564E-06	-0.18969E-01 0.73939E-02 0.30094E-03 0.66034E-05 0.91127E-07 0.95870E-09	0.34651E 00 0.27438E-00 0.14016E-01 0.40174E-03 0.26877E-03 0.65633E-06	0+17113E-00 0+17827E-00 0+31409E-01 0+17930E-02 0+6483E-04 0+16454E-05	0.46390E=00 0.12234E=00 0.74403E=02 0.24963E=03 0.60344E=05 0.11009E=06	-0.173625-00 -0.421925-02 0.305275-02 0.168955-03 0.418905-03 0.810405-07
		x = :	1.45			x = 1	95	
1 2 3 4 5	0.48989E-00 0.46899E-01 0.25108E-02 0.72794E-04 0.14629E-05 0.20892E-07	0.21454E-00 0.09549E-01 0.5434BE-02 0.17507E-03 0.34653E-05 0.53571E-07	0.27599E-00 0.21392E-01 0.63697E-03 0.12859E-04 0.12859E-06 0.19649E-08	-0.34350E-01 0.49806E-02 0.40218E-03 0.92719E-05 0.14200E-06 0.15894E-08	0.54209E 00 0.30400E-00 0.22902E-01 0.10114E-02 0.395914E-04 0.91248E-08	0-16632E-00 0-17923E-00 0-37110E-01 0-22330E-02 0-85400E-04 0-22805E-05	0.46119E-00 0.13922E-00 0.94020E-02 0.33194E-03 0.94603E-05 0.16186E-05	-0.10278E-00 -0.12120E-01 0.39806E-02 0.19236E-03 0.57049E-05 0.11787E-06
	A 40404F	л = 1.				x = 2.	oc	
	0.90697E 00 0.80725E-01 0.92008E-02 0.92248E-04 0.21099E-05 0.92247E-07	0.21002E-00 0.10199E-00 0.68129E-02 0.23905E-03 0.52708E-05 0.82504E-07	0.30133E-00 0.27220E-01 0.86698E-03 0.18668E-04 0.28659E-06 0.32574E-08	-0.91491E-01 0.97970E-02 0.93040E-03 0.13237E-04 0.21781E-06 0.2998E-08	0-33408E 00 0-33294E-00 0-27399E-01 0-12691E-02 0-46591E-04 0-12577E-05	0.161445-00 0.177815-00 0.432776-01 0.432776-01 0.276855-02 0.111415-03 0.319285-05	0.49814E-00 0.15639E-00 0.11795E-01 0.43817E-03 0.11721E-04 0.23962E-06	-0.19021E-00 -0.21726E-01 0.40758E-02 0.24607F-03 0.77688E-05 0.17688E-05

						x = 0.05										
		n	$\mathbf{a}_{\mathbf{r}}^{\mathbf{n}}$		a_n^i		ρľ	ь	Ĺ		a.r			x = 0.55	_	
		2 0.256	83E-04 74E-08 36E-12	0.58	299E-04 725E-08 184E-12	0:162 -0:514 -0:701	66E-07	0=12229 -0=13884 -0=11844	E-07	0.488 0.360 0.160	27 86E = 01 80E = 03 65E = 05 19E = 07 36E = 10	0.970 0.771 0.354	41 90E-01 14E-03 80E-05 08E-07 76E-09	0+3380 0+2810 0+1334 0+4122 -0+27411	6E-04 6E-06 DE-09	bi 0.14687E-0 0.14139E-0 0.70107E-0 0.24088E-0 0.54576E-1 2.17075E-1
					х	= 0.10							x	≃ 0.60	•	30170735-1
	9	0,9595	5E-07	0.354 0.197; 0.9934		0.6420 0.5579 0.6094	6E-09	0.352798 0.12005E 0.75137E	-09	0.5451 0.7588 0.6620 0.3494 0.1084 -0.54926	8E-01 2E-05 5E-07 IE-09	0.6968 0.1492 0.1412 0.7793 0.27774 0.13584	2E-01 4E-02 6E-04 3E-07	0.52636 0.51772 0.29189 0.105896 -0.945446	E-04 E-06 E-08	0.21520E-02 0.25506E-04 0.15189E-06 0.62748E-09 0.15164E-10
	1	0.670651	-03	1.11841		0.15							x ·	0.85		
	2 3 4 5 6	0.726956 0.421246	-06	0+11961 0+14952 0+91251	E-05	0.487819 0.325299 0.268419	-08	0.26596E- 0.17041E- 0.21229E-	08	0.71582 0.11591 0.11573 0.71650 0.282145 -0.77330E	E-02 E-04 E-07	0.85449 0.22160 0.24612: 0.15814: 0.4653568 0.223228	-02 -04 -06 -09	0.79123E 0.90438E 0.59937E 0.25605E 0.65415E -0.34332E	-04 -06 -08	0+30083E-02 0+437:0E-04 0+30752E-06 0+13933E-08 0+1623ZE-10 0+74643E-12
	1	0.16107E	- 0.2		x = (0.20							X = 1	0.70		
	2 3 4 5 6	0.30631E- 0.30611E- -0.16796E- -0.12111E- -0.39404F-	-05 -08 -11	0.28356E 0.62921E 0.66945E 0.73433E 0.31610E 0.86319E	-05 -08 -11 -14 -18	0.20615E- 0.23735E- 0.67949E- -0.17455E- -0.10256E- -0.55061E-	07 10 11 11	0.11151E-0 0.12853E-0 0.64264E-10 0.19697E-10 0.61212E-12 0.30429E-15	7	0.92013E- 0.18618E- 0.19912E- 0.3864E- 0.1908E-	-02 -04 -06	0.10206E 0.3192EE 0.41123E- 0.1054E- 0.1477EE- 0.35930E-	-02 -04 -06	0+11535E-0 0+15292E-0 0+11669E-0 0+57109E-0 -0+49697E-1 -0-50971E-1	9 9 8 2	0.40198E-02 0.71654E-06 0.59080E-06 0.30553E-08 0.19809E-10 -0.22265E-11
	1	0.320056-0	12	0.55305E-	X = 0.								x = 0.	75		
	2 3 4 5 6	0.93544E=0 0.14576E=0 0.71453E=1 -0.14050E=1 -0.60639E=1	5 7 1 3	0-19149E- 0-31594E- 0-35354E- 0-36388E- 0-15466E-	04 07 10 -	0-63145E-0 0-11200E-0 0-13034E-0 0-45642E-1 0-22614E-1 0-63912E-1	6 9 1	0.39629E-04 0.60998E-07 0.1099E-10 0.06298E-12 0.36596E-14		0:1389E-0 0:23644E-0 0:31403E-0 0:23806E-0 0:13895E-0 0:49020E-1	12 14 16	0.11889E-0 0.44812E-0 0.64269E-0 0.56723E-0 0.31391E-0 0.99900E-1	2 4 6 8	0.16366E-01 0.24839E-03 0.21696E-05 0.12203E-07 0.36794E-10 -0.19171E-11		0.919276-02 0.112976-03 0.10618E-05 0.637656-06 0.26621E-10 0.18990E-11
	1				x = 0.30	0							x = 0.8	0		
	2 9 4	0.56496E-02 0.29302E-04 0.52180E-07 0.49238E-10 -0.10382E-12 -0.64721E-16	0	0.95290E-0 0.4759BE-0 0.11296E-0 0.17527E-0 0.26703E-12 114716E-15	-0	0-157846-03 0-401686-06 0-645426-09 -159386-10 -877946-11 -472906-13	0	0.82912E-09 0.21886E-08 0.46398E-09 0.40407E-10 0.33629E-11 0.25036E-13		0.14310E-00 0.33005E-02 0.49257E-04 0.45949E-05 0.28190E-08 0.11133E-10	!	0:13523E-00 0:61466E-02 0:10347E-03 0:10361E-05 0:63514E-08 0:24133E-10		0+72493E-01 0+39110E-03 0+38758E-05 0+24777E-07 0+91863E-10	0	0+62440E-02 0+17200E-03 0+19009E-05 0+12799E-07 0+57462E-10
1		0.00000			× = 0.35							,	= 0.85			
3 6 9	-	0,92003E-02 0.50036E-00 0.15336E-00 0.25529E-09 0.52681E-12 0.67896E-15	0. 0. 0.	15096E-01 10262E-03 33112E-06 63973E-09 13376E-11 12357E-14	-0.	34299E-03 11816E-05 23715E-08 72360E-11 86707E-11 24886E-12	0. 0.	17475E-03 63169E-06 13703E-08 39749E-10 64216E-11 11140E-12		0-17330E-00 0-45216E-02 0-75181E-04 0-79164E-06 0-54740E-08	0	-15036E-00 -62615E-02 -15714E-03 -17292E-05 -12329E-07 -58412E-10	0.	30440E-01 59916E-03 66837E-05 48249E-07 22466E-09 46748E-11	0. 0. 0.	719018-02 259728-09 321908-09 245948-07 140128-09 895898-12
1		0-14133E-01			= 0.40							x	= 0.90			
3 • 9	-0	-985 ME-04 -98995E-04 -91039E-09 -83942E-12 -27145E-14	0 - 1 0 - 8 0 - 2: 0 - 4:	2299E-01 9952E-03 4034E-06 0590E-08 4389E-11 9653E-14	0+3 0+7 0+1 -0+4	7265E-03 0114E-05 6729E-08 3508E-10 3444E-11 3922E-12	0.1 0.4 0.3 0.4	3277E-03 5906E-05 1596E-08 0623E-10 2776E-11 7993E-12	000	-20591E-00 -60987E-02 -11202E-03 -13171E-05 -10240E-07	0. 0. 0.	16368E-00 10902E-01 23263E-03 26738E-09 12977E-07 12461E-09	0 • 6 0 • 1 0 • 9	0547E-01 19585E-03 1171E-04 0356E-07 8725E-09 5543E-11	0 • 9 0 • 4! 0 • 2!	7346E-02 6351E-03 2749E-05 5445E-07 8465E-09
					0.45							x =	0.95			
2 2 3 6 5 6 6	0. 0.	20768E-01 17797E-03 88803E-06 26247E-08 84176E-12 12639E-13	0.35 0.19 0.50 0.13	986E-01 839E-03 899E-05 979E-08 838E-10 921E-13	0.68 0.22 0.365 -0.874 -0.180	2085-02 7715-05 9225-07 9145-10 9295-11	0.35: 0.114 0.630 0.500	256E-03 914E-05 966E-07 972E-10 162E-11 41E-11	0. 0. 0.	24016E-00 61143E-02 16337E-01 21343E-05 18495E-07 10800E-09	0.1 0.3 0.4 0.4	7066E-00 4130E-01 9747E-03 9432E-05 1986E-07	0.53 0.13 0.16 0.16 0.10	2531E-01 2106E-02 157E-06 341E-06 003E-08	0.83 0.83 0.81	9146-02 6725-03 9076-05 0496-07 7606-09 2126-12
	0.2	9460E-01	0.421	x = (49E-01		34F_A-						x = 1	.00			
	0.10 0.40 0.10	0292E-03 8998E-05 8229E-08 890E-10 8910E-13	0.604 0.397 0.1511 0.4117 0.1197	78E-03 74E-05 17E-07 77E-10	0.1439 0.15666 0.1609 0.1093 -0.8803	2E-07 5E-09 17E-11	0.992: 0.7376 0.3000 0.8863 0.5706 0.9864	02E-07 04E-10 06E-11	0.1 0.2 0.3	7520E-00 0867E-01 9974E-03 9728E-05 1396E-07 366E-09	0.479 0.479 0.731 0.723	903E-00 388E-01 935E-03 43E-05 101E-07	0.187 0.287 0.286 0.198	37E-08	0.688 0.129 0.140 0.104	11E-02 04E-03 91E-04 13E-06 13E-11

						x = 1.05								x = 1.55			
		n	a <mark>r</mark>		$\mathbf{a}_{\mathbf{n}}^{\mathbf{i}}$		ь г	è	n n		ar n				ъ.г	. 4	
	3 4 5	0.1381 0.3281 0.5210	14E-01 14E-03 15E-05 3E-07	0.227 0.669 0.112	73E-00 95E-01 39E-09 62E-04 01E-05	0.826 0.264 0.446 0.489 0.3723 0.1156	79E-01 85E-02 90E-04 13E-06	0.3802; 0.9104; 0.1962; 0.2553 0.19163 0.10636	2E-02 8E-03 1E-04 7E-06 IE-08	0.120 0.518	92E 0 62E-0 73E-0 25E-0 03E-0	0 0.1704 0 0.1143 2 0.9194 9 0.3410	0E-00 0E-03 9E-05	0 - 2899 0 - 3730 0 - 1449 0 - 3433 0 - 5704 0 - 69804	E-01 E-02 E-04	bi -0.116346 0.821998 0.40649E 0.18220E 0.249996 0.327698	-00 -01 -04
	_					1.10							x	= 1.60			
	1 2 3 4 5	G.9440; 0.1787; 0.45525 G.73860 G.91548 G.74735	E-01 E-05 E-07	0.1919 0.2834 0.9192 0.1698 0.28391 0.16931	46-01 06-03 36-04 96-04	G-1010 D-3670 O-6770 G-8138 D-68076 O-34370	2E-02 DE-04 DE-06 DE-08	-0.540531 0.117401 0.289645 0.385705 0.542305 0.173335	E-02 E-04 E-06	0.5432 0:1426 0.6535 0.2204: 0.53264 0.9226	5E-00 4E-02 3E-03 4E-05	0=1681c 0=12475 0=11919 0=44864 0=11476 0=20471	E-00 E-01 E-03 E-04	0-309420 0-452980 0-191898 0-484725 0-857778	-01 -02 -04 -06	-0.13225E- -0.90806E- 0.49738E- 0.18069E- 0.36870E- 0.51748E-	-03 -03 -04 -04
						1.15							x ·	1.65			
	1 2 3 4 5 6	0.376268 0.220168 0.62171E 0.13715E 0.148528 0.13246E	-01 -05 -04 -06	0+19291 0+34900 0+12437 0+29130 0+329216 0+300488	E-01 E-02 E-04 [-06 !-08	0.12076: 0.50085: 0.10088: 0.12248: 0.121248: 0.71243E	-02 -03 -05 -07	-0.68797E- 0.14797E- 0.41850E- 0.61661E- 0.60221E- 0.41191E-	-02 -04 -06	0.54819 0.16486 0.61882 0.28979 0.762016 0.156678	E-00 E-02 E-03 E-05	0.166218 0.13610E 0.13826E 0.15826E 0.15929E	-00 -01 -03	0.32056E- 0.54305E- 0.25159E- 0.67676E- 0.12735E- 0.17692E-	-01 -02 -04	-0.14748E-0 -0.34432E-0 0.39581E-0 0.24344E-0 0.53641E-0 0.80912E-0	12
	1	U-43624E-	.00	0.192126	x = 1								x =	1.70			
	3 4 5	0.28897E- 0.83834E- 0.17104E- 0.23593E- 0.22961E-	01 03 04 96	0.42174E- 0.14578E- 0.36749E- 0.32170E- 0.52170E-	-01 -02 -04 -04	0.14198E- 0.67371E- 0.14774E- 0.21079E- 0.21025E- 0.14264E-	02 03 05	-0.153402-0 0.17979E-0 0.59245E-0 0.96411E-0 0.10292E-0 0.76030E-1) 2 4 6 7	0.59046E 0.19294E 0.10206E 0.37795E 0.10243E 0.20905E	-00 -01 -05	0:16479E- 0:14193E- 0:1679BE- 0:7558BE- 0:21873E- 0:44294E-	00 01 03	0.34743E-0 0.64279E-0 0.32679E-0 0.53508E-0 0.18676E-0 0.27549E-0)1)2)4	-0.18179E-00 -0.49378E-02 0.696982E-03 0.32373E-04 0.74991E-04 0.12409E-07	!
			_		x = 1.2	25							x = 1.	75			
	1 2 3 4 4 5 5 6 6 6 6 6 6 6 6 6 6 6 6 6 6 6 6	0.49357E-0 0.36309E-0 0.1117E-0 0.26595E-0 0.36768E-0 0.38899E-0	1 2 6	0.19002E-0 0.50542E-0 0.21875E-0 0.52297E-0 0.81080E-0	11 12 14	0.16919E-0 0.89991E-0 0.21296E-0 0.32928E-0 0.33640E-0 0.27398E-0	2 3 5 7	-0.29890E-01 0.21232E-02 0.82337E-04 0.10767E-05 0.17184E-07 0.13748E-09	!	0.99142E 0.22046E=(0.1266E=(0.48927E=(0.13998E=0 0.28949E=0	00 01 03	0-16968E-0 0-14790E-0 0-20160E-0 0-96899E-0 0-29749E-04	1	0.36614E-00 0.75130E-01 0.42063E-02 0.12793E-03 0.2707EE-05 0.42335E-07		-0.17099E-00 -0.11516E-01 0.79285E-03 0.42485E-04 0.10910E-05 0.18793E-07	
1		0 48888			x = 1.30								x = 1.8	0			
3 6		0.45802E-00 0.4527ZE-01 0.14743E-02 0.14853E-04 0.56283E-06 0.44370E-08	0. 0. 0.	18704E-00 59804E-01 28494E-02 73747E-04 12378E-05 14488E-07		0.18509E-00 0.11706E-01 0.30249E-03 0.50534E-05 0.59151E-07 0.49428E-09	0	0+38347E-01 0+24152E-02 0+11231E-03 0+22189E-05 0+28144E-07 0+25198E-09		0.55003E 00 0.24685E-00 0.15633E-01 0.62900E-03 0.12959E-04 0.41643E-06	,	0.16292E-00 0:15177E-00 0:24073E-01 0:12325E-02 0:40086E-04 0:90927E-06		0.38476E-00 0.46733E-01 0.653662E-02 0.17339E-03 0.38835E-05 0.64249E-07	0	0-18694E-00 0-17268E-01 0-87188E-03 0-85029E-04 0-15271E-05 0-28075E-07	
					= 1.35							x	= 1.85				
1 2 3 4 5 6	9	0+47946E-00 0+56014E-01 0+19264E-02 0+48739E-04 0+84734E-05 0+10459E-07	0.1 0.1 0.1	10950E-00 19092E-01 14715E-02 10250E-03 8504E-05 3482E-07	0. 0. 0.	20693E-00 15134E-01 42346E-03 76284E-05 96263E-07 47987E-09	0. 0.	92398E=01 26214E=02 15049E=03 32741E=05 45111E=07 43609E=09		0.54684E 00 0.27752E-00 0.19220E-01 0.80340E-03 0.29460E-04 0.58730E-06		0-16237E-00 0-15346E-00 0-28532E-01 0-15542E-02 0-53540E-04 0-12844E-05	0.	•40334E=00 •98939E=01 •47846E=02 •23291E=03 •55128E=05 •96362E=07	-0. 0. 0.	19793E-00 24240E-01 91740E-03 70342E-04 21124E-05 61427E-07	
1	0.	497 89 E-00	0.1	7998E-00	= 1.40	220446-00						х -	1.90				
9 4 5 6	0	048754E-01 124955E-02 147324E-04 12570E-05 14685E-07	0 = 36 0 = 14 0 = 27	0652E-01 1829E-02 1829E-03 1875E-05 1379E-07	0 - 1 0 - 1 0 - 1	19322E-01 00441E-03 11341E-04 5380E-06 5243E-08	0 e i 0 e i 0 e i	67650E-01 26698E-02 19013E-03 07514E-05 0996E-07 4873E-09		0.394202E 00 0.30586E-00 0.23530E-01 0.10199E-02 0.33916E-04 0.62439E-06	0	•16194E-00 •15303E-00 •33564E-01 •19512E-02 •70914E-04 •17966E-05	0.0 0.1 0.1	42193E-00 11199E-00 85092E-02 81021E-03 77495E-05 4287E-06	0.9	0645E-00 2429E-01 0684E-03 8716E-04 888E-05 0971E-07	
1				X =	1.45							x =	1.95				
2	0.9 0.9 0.1	11998E 00 19467E-01 2073E-02 11943E-04 8360E-05 4174E-07	0.914 0.591 0.191	144E-00 155E-01 162E-02 22E-03 41E-05 95E-07	0.24 0.80 0.16 0.24 0.25	947E-00 363E-01 166E-03 622E-04 184E-06 887E-08	0 • 2 • 0 • 6 7 0 • 10	8670E-01 867E-02 634E-03 642E-05 974E-06 634E-08	0	0-59571E 00 0-39329E-00 0-28686E-01 0-12879E-02 0-44881E-08 0-11864E-09	0. 0.	16151E-00 15068E-00 15068E-01 15068E-01 15089F-02 1508E-04 14897E-05	0 = 1 : 0 = 1 : 0 = 1 :	40526-00 24536-00 55786-01 19906-02 17936-04	0.61 0.11 0.39	424E-00 759E-01 071E-03 035E-03 067E-05	
	0.59	401E 00	4.1700	X = 1.		***						x = 2.	.00				
	0.10 0.40 0.12 0.26	0946-00 0946-00 9236-02 4256-03 5226-05 1046-07	0.1732 0.1032 0.7407 0.2566 0.5757 0.9010	0E-00 0E-02 8E-03 7E-05	0+303 0+108 0+240	96E-00 3BE-01 57E-02 39E-04 17E-06 15E-08	0.189 0.329 0.953 0.166	03E-00 61E-02 78E-03 94E-05 92E-06 97E-08	0. 0.	52407E 80 35754E-00 34420E-01 14147E-02 58860E-04 15404E-05	0.1 0.4 0.3 0.1	4098E-00 4472E-00 5302E-01 1072E-02 1192E-03 199E-05	0.13° 0.13° 0.531 0.149	763E-00 036E-01 712E-03 00E-04	0.521 0.591 0.135 0.522	29E-00 41E-01 67E-09 28E-09 57E-05 POE-06	

					x = 0.05						- A 6 F		
		n	a <mark>r</mark>	$\mathbf{a}_{\mathbf{n}}^{i}$	b	r	b <mark>i</mark>	;	ar n	a <u>i</u>	x = 0.55 h	n r	
		1 0.2914 2 0.3448 3 -0.2933 6	94E-08	0.49080E- 0.65627E- 0.76324E-	08 -0.1395	IE-09 -0.	.04617E-08 19536E-09 54303E-10		76E-01 0:6 2E-03 Cal 9E-05 5.8 2E-07 0.3 5E-10 0:1	7n 0020E-01 0718E-02 5039E-05 9019E-07 1416E-09 663E-12	0.41004 0.34739 0.16367 0.49645 -0.98090 -0.23621	#E=02	E-03 E-04 E-07 E-09
					x = 0.10					x =	- 0.60		
	1 2 3 4 5		E-04	0+39279E-0 0+71743E-0 0+39649E-1	6 0-149130	-09 0.9	?7280E-06 99579E-11 0061E-10	0.69904 0.69014 0.76653 0.40223 0.136008 -0.14203	E=03 0.16 E=05 0.15 E=07 0.85 E=09 0.29	190E-01 488E-02 566E-04 024E-07 470E-09 799E-12	0.633312 0.638318 0.361906 0.128276 -0.390396 -0.125926	E-04 0-17703E -06 0-1208E -08 0-43747E -11 -0-2749E	-04 -06 -09
					x = 0.15					x =	0.65		
	1 2 3 4 5 6	0.1997[6 0.846276 0.4461736	-06 0	0•13250E-02 1•10503E-05 1•98053E-09	0.40879E- 0.39275E- -0.73607E-	-08 0.13	1439E-05 1537E-08 1334E-11	0.86396E 0.1338DE 0.13407E- 0.82482E- 0.33417E- 0.34242E-	-02 0.2444 -04 0.2712 -07 0.1739 -09 0.7157	95E-02 94E-04 94E-06 91E-09	0.94275E- 0.11167E- 0.74227E- 0.31252E- 0.42369E- 0.64957E-1	03 0.29607E-0 05 0.22219E-0 08 0.10003E-0 11 -0.31642E-1	04 06 08
				,	= 0.20					x = 0.	.70		•
	2 3 4 5 6	0+19235E- 0+35694E- 0+34853E- -0+21351E- -0+11346E- +0+31262E-	05 0. 08 0. 11 0.0	31407E-02 69456E-05 73663E-08 61619E-11 28769E-14 76530E-18	0.25695E-0 0.29572E-0 0.12712E-1 -0.16756E-1 -0.1153E-1 -0.51555E-1	0 0 100 0 0 971 0 0 991 1 0 177	99E-09 #9E-07 52E-10 89E-11 78E-12 91E-15	0.11071E- 0.19530E- 0.22497E- 0.16026E- 0.74998E-0 0.35880E-1	02 0.3527 04 0.4532 06 0.3972	9E-02 5E-04 NE-06 LE-08	0.13566E-0 0.18737E-0 0.18432E-0 0.70816E-0 0.15053E-1 -0.33196E-1	03	4 6 8 1
	1	0.382905-0			= 0.25					x = 0.7	5		
	2 3 4 5 6	04109056-04 0416457E-03 0416457E-03 0416457E-13 -0412530E-13 -0456765E-17	7 0-34 0-36 0-35	1226E-02 1361E-06 1745E-07 1937E-10 604E-13 252E-16	0.78565E-04 0.19966E-06 0.19267E-09 0.52025E-11 0.31556E-12 -0.58597E-14	0.2515 0.46306 0.30078 0.89997 0.11796 0.20720	8E-07 8E-10 7E-12 9E-11	0413876E-00 0427887E-02 0436425E-04 0429737E-06 0416094E-08 0472234E-11	0.495058 0.730486	-02 -04 -06 -08	0.19016E-01 0.30320E-03 0.26786E-05 0.15150E-07 0.61913E-10 0.16774E-11	0:17293E-02 0:71872E-04 0:75978E-06 0:46029E-08 0:12764E-10 -0:19953E-12	
					0.30					X = 0.80			
		0.67730E-02 0.27179E-04 0.60240E-07 0.60319E-10 -0.10786E-12 -0.60587E-16	0.525 0.124 0.165 0.238	540E-01 550E-04 528E-06 513E-09 187E-12 08E-15	0-19594E-03 0-49995E-06 0-69059E-09 -0-13633E-10 -0-23747E-11 -0-46510E-13	0.605050 0.164238 0.296528 0.290568 -0.232778 0.19217E	E-06 E-09 E-11	0+17020E-00 0+38968E-02 0+37174E-04 0+32989E-06 0+32504E+08 0+14815E-10	0.13871E- 0.67889E- 0.11407E- 0.11095E- 0.69857E- 0.29833E-	02 0 03 0 05 0 06 0	.25911E-01 .47534E-03 .47757E-05 .30735E-07 .13267E-09 .52862E-12	0.13529E-02 0.10510E-03 0.13130E-05 0.92182E-06 0.39687E-10 -0.12703E-21	
				x = 0	.35					x = 0.85			
	2 3 4 -	0+11054E-01 0+58865E-04 0+17715E-06 0+30865E-09 0+78645E-12 0+44925E-15	0:1662 0:1133 0:3645 0:6721 0:9068 0:1125	9E-06 9E-06 3E-09 0E-12	0.42457E-03 0.14703E-05 0.28294E-08 -0.87194E-11 -0.18611E-11 -0.30683E-12	0.12540E- 0.47390E- 0.99336E- 0.29036E- -0.22103E- -0.12007E-	-06 -09 -11 -11	0.20441E-00 0.33489E-02 0.87332E-04 0.91149E-06 0.43111E-08 0.31244E-10	0.15191E-0 0.91206E-0 0.17325E-0 0.19031E-0 0.13542E-0 0.66734E-1	2 0.0 3 0.6 5 0.5 7 0.2	344315-01 724636-03 821706-05 597186-07 298276-09	0.27051E-03 0.1476EE-03 0.21822E-05 0.17529E-07 0.91276E-10 -0.5131BE-12	
	•	•17017E-01		x = 0.4	Ю					x = 0.90			
2 5 4 5 4	-0.	011507E-03 045075E-04 010677E-08 10677E-13 120422E-13 125769E-16	0.20973 0.22033 0.22963 0.22486 0.312276 0.633322	E-09 E-06 E-08 E-11 -	0-82988E-03 0-37452E-05 0-974930E-08 0-97694E-11 0-25024E-11	0.23196E-0 0.11802E-0 0.31432E-0 0.14836E-1 0.84998E-1 -0.25460E-1	5 6 0 3	0-24049E-00 0-72294E-02 0-13029E-03 0-15194E-05 0-11787E-07 0-65230E-10	0-16176E-00 0-12028E-01 0-25472E-03 0-31630E-05 0-25269E-07 0-13956E-09	0.10 0.11 0.11	4675E-01 0774E-02 3700E-04 1160E-06 1391E-09	-0.10809E-02 0.19901E-03 0.35010E-05 0.32056E-07 0.18957E-09 -0.10641E-11	
				x = 0.45					,	· = 0.95			
1 2 3 4 5	0 . 0 . 0 .	29055E-01 20806E-03 10289E-05 90390E-08 93292E-11 14228E-13	0.34480E 0.39583E 0.21039E 0.64586E 0.12097E 0.28808E	-03 0 -05 0 -08 0 -10 -0 -13 -0	-14990E-02 -8534E-05 -27303E-07 -51982E-10 -10017E-11 -72990E-13	0.39142E-03 0.26233E-05 0.88389E-06 0.21526E-10 -0.35651E-12 -0.26098E-12		0+27745E-00 0+96391E-02 0+19010E-03 0+24631E-05 0+21282E-07 0+19161E-09	G:16914E-00 G:15597E-01 O:37213E-03 O:51109E-07 O:45489E-07 O:28160E-09	0.156 0.222 0.201 0.126	660E-01 661E-02 207E-04 193E-06 197E-08	-0.44824E-02 0.29823E-03 0.94971E-05 0.94848E-07 0.97124E-09 -0.74810E-12	
	0.3	5594E-01	0,463446-	x = 0.50	24444				ж	= 1.00			
	0.39 0.29 0.76 0.17	3377E-03 1445E-03 1445E-03 1458E-08 1614E-10 1011E-13	0:46805E- 0:45820E- 0:16621E- 0:43104E- 0:10596E-	03 0. 05 0. 07 0.	19037E-09 10442E-11	0.61092E-09 0.93228E-09 C.22992E-07 0.62110E-10 0.32096E-11 0.48602E-12	0. 0. 0.	91429E-00 12698E-01 27225E-09 38942E-05 37245E-07 25178E-09	0-17365E-00 0-19912E-01 0-52886E-03 0-80520E-05 0-79476E-07 0-54519E-09	0.7030 0.2230 0.3309 0.3330 0.2453 0.1064	03E-02 99E-04 92E-06 14E-08	-0.89059E-02 0.31924E-03 0.81928E-05 0.96019E-07 0.72953E-09 0.14173E-11	

						x = 1.05										
		n	ar n		$\mathbf{a}_{\mathbf{n}}^{i}$		b _n	1	b i		ar.			x = 1,55		
		2 0.1 3 0.3 4 0.6 5 0.6	4987E-00 4549E-01 8331E-03 0191E-05 8426E-07 7454E-09	0.25 0.73 0.12 0.13	7552E-00 050E-01 018E-03 398E-04 502E-06 240E-08	0.31 0.54; 0.60; 0.46; 0.254	4422-01 1665-02 2065-04 1215-06 1755-08	0,3733 0,1199 0,1994	3E-01 0E-03 4E-04 6E-06	0.14 0.61 0.19 0.43	2F 5441E 00 4199E-00 1336E-02 1343E-03 1647E-05 683E-07	0.11 0.10 0.375 0.899	ai 731E-00 570E-00 104E-01 48E-03 42E-05 34E-06	0.37 0.14 0.416 0.67	bf 499E-06 975E-01 629E-02 101E-04 106E-08	-0.79405E-
		1 0.30	9545-00			= 1.10							x.	- 1.50		20110536-
		2 0.211 9 0.591 4 0.911 5 0.109	9985-01 496-03 496-03 906-04 446-09	0+175 0+310 0+1015 0+1861 0+2236 0+1862	8E-04 4E-04 8E-08	0-1018 0-4279 0-8199 0-7984 0-8419 0-4461	2E-04 7E-06 2E-08	-0.22821 0.40544 0.17083 0.25696 0.23984 0.10748	-04 -04 -06	0.164 0.773 0.256 0.615	745E 00 156E-00 193E-02 171E-03 12E-05 11E-06		3E-00 9E-00 5E-01 4E-03 2E-04	0.283 0.452 0.2181 0.5763 0.1039 0.1372	7E-02 1E-04 3E-05	-0.16459E-0 -0.11346E-0 0.07313E-0 0.86528E-0 0.21886E-0: 0.39483E-0:
		0,4141	1E-00	0.1792									x •	1.65		
		2 0.2791 3 0.7265 4 0.1354 5 0.1708 6 0.1928	8E-03 8E-04 9E-04	0.38031 0.13711 0.27669 0.36202 0.33005	7E-01 1E-02 1E-04 1F-04	0:11912 0:97798 0:12167 0:16703 0:14949 0:4978958	E-02 E-03 E-07	-0:32481E- 0:39292E- 0:29484E- 0:40314E- 0:41592E- 0:22898E-	-03 -04 -06 -08	0.5422 0.1930 0.97103 0.33782 0.89812 0.15724	3E-00 3E-02 3E-03	0.15544 0.15249 0.15153 0.64582 0.17519 0.352548	E-00 E-01 E-03	0.30193 0.53208 0.28324 0.80091 0.153679	E-01 [-02 [-04	-0.17896E-00 -0.19406E-01 0.21628E-04 0.10938E-04 0.9107E-06
	1	0.44292	F-00		x = 1	.20							x = 1			************
	9	04 146 11	-01 -03 -04	0.17013E 0.49943E 0.18303E 0.40238E 0.57373E 0.57091E	-01 -02 -04	0.137056 0.768546 0.177406 0.257446 0.258766 0.190906	-02 -03 -05 -07	-0.49867E-0 0.30318E-0 0.31997E-0 0.41732E-0 0.70233E-0 0.47033E-10	, ,	0,94294 0,22991 0,22199 0,440998 0,118478 0,230238	-00 -0 -03	0.15901E 0.13771E 0.1833E 0.8319E- 0.24862E- 0.48340E-	-00 -00 -01 -03	0.32006E 0.61837E 0.36401E 0.11011E 0.22903E 0.33598E	-01 -02 -03	-0.19261E-00 -0.20803E-01 -0.61226E-04 -0.1386E-04 -0.43214E-04 -0.78267E-08
	1 2	0.46799E		0-16651E-		0							x = 1.7	75		
	3 4 5 6	0+43540E- 0413095E- 0+28460E- 0+42337E- 0+44727E-	02 04 06	0.54762E-(0.24118E-0 0.57586E-0 0.49176E-0 0.96371E-0	01 02 04 06	0.10068E-0 0.25447E-0 0.40116E-0 0.43849E-0 0.43849E-0) i 9 5 7	-04567876-01 0-91782E-04 0-41935E-04 0-92592E-06 0-11407E-07 0-95769E-10		G.56174E G.20958E= G.19049E= G.57146E= G.16200E= G.99330E=	00 01 03	0.15706E-0 0.14084E-0 0.22004E-0 0.10662E-0 0.32725E-0 0.70037E-0	1 2	0+33884E-0 0+71043E-0 0+46322E-0 0+14984E-0 0+32526E-0 0+51913E-0) [2 3	-0.20479E-00 -0.26982E-01 -0.1963E-03 0.16103E-06 0.59265E-06 0.11626E-07
	1 2	0-48986E-0		-16200E-00	0	•17975E-00		-79970E-01					x = 1.80			
	5 6	0.54281E-0 0.17299E-0; 0.40388E-0; 0.64835E-0; 0.74034E-0;	2 0	-64405E-01 -31410E-02 -81210E-04 -13614E-05 -15928E-07	0.	12999E-01 35954E-03 61396E-05 72527E-07 63064E-09	-0 0 0	-27939E-09 -53415E-06 -13531E-05 -13531E-05 -18744E-07 -17325E-09		0.55871E 0 0.27844E-0 0.18597E-0 0.73543E-0 0.21954E-0 0.47734E-0	9 6	0.15946E-00 0.14184E-00 0.76203E-01 0.13558E-02 0.44094E-04 0.79958E-06	9	0.39791E-00 0.40740E-01 0.58377E-02 0.20191E-03 0.46498E-05 0.77997E-07	-0	0-21998E-00 0-34159E-01 0-4040E-03 0-18580E-06 0-79980E-06 0-17040E-07
	1	0.508566 00	0.	x 15936E-00	= 1.35							x	= 1.85			
	2 3 4 5	0.67089E-01 0-22698E-02 0.56519E-06 0-97671E-06 0.12026E-07	0 - 1 0 - 1 0 - 2	74713E-01 40459E-02 11296E-03 70441E-05 75814E-07	0.9 0.9 0.1 0.1	19216E-00 16550E-01 10085E-03 12408E-05 1795E-06 1088E-08	-0.0 0.0 0.0	86098E-01 9428ZE-03 558Z9E-04 19404E-03 19595E-07 19594E-09		0.55403E	0. 0. 0.	14204E-00 14084E-00 30954E-01 17112E-02 58890E-04 4120E-09	0.1 0.7 0.2 0.4	37726E-00 P0639E-01 72875E-02 P6956E-09 -5778E-05 1673E-06	-0. -0. 0.	22507E-00 42319E-01 70914E-03 70945E-08 10617E-05
1		9+92422E 00	0.14	x = 9647E-00	1.40							х -	1.90			
3 4 5		0-82198E-01 0-29370E-02 0-78191E-06 0-14492E-05 0-1477E-07	0.89 0.51 0.15 0.30	3447E-01 1582E-02 1515E-03 123E-05 086E-07	0.20 0.68 0.13 0.18 0.18	048E-00 783E-01 649E-03 694E-04 819E-06 964E-08	-0.19 0.74 0.27	01495-00 P2985-02 P4485-04 P3455-09 P395-09	0.	954782E 00 93437E-00 928003E-01 11950E-02 39324E-04 95026E-04	0.1: 0.3: 0.21 0.77	6472E-00 3810E-00 8269E-01 1446E-02 1993E-C4 1750E-05	0 - 10 0 - 90 0 - 35 0 - 92	P690E-00 9126E-00 9124E-02 969E-03 969E-05 269E-06	-0.5 -0.1 0.21 0.13	3437E-00 1998E-01 1408E-02 1100E-04 851E-03 039E-07
1 2		0.53696E QQ 0.99667E-01	0.154	285-00		69E-00	=0.117	#3E-00				x = 1	.95			
3 5 6	0	-37405E-02 -10579E-03 -21200E-05 -30080E-07	0.651		0.994 0.200 0.295 0.3199	2E-04	0.881	80E-02 46E-04 32E-03 80E-07	0.1 0.1 0.5	14023E 00 18045E-00 4119E-01 5101E-02 2028E-04 9222E-05	0.199 0.421 0.266 0.102	720E-00 197E-00 141E-01 196E-02 47E-03 70E-05	0.111 0.110 0.467	79E-00 95E-00 63E-01 83E-03 81E-06	-0.613 -0.173 0.195 0.177	140E-00 139E-01 139E-02 17E-04 72E-05 77E-07
1		54697E 00	0-1528		10 0 - 2 4 4 8 4	1E-00 -						x = 2.0	ю			
į	0. C.	11961E-00 4831E-02 14449E-03 30607E-05 46449E-07	0:1068; 0:8147; 0:2826; 0:63934 0:99046	9E-00 9E-02 9E-03 9E-05	0.32475 0.12516 0.26621 7.45552 0.52630	E-01 E-02 E-06	-0.1994 -0.5318 -0.93096 -0.5063 -0.19596	BE-02 DE-04 DE-05 DE-06	0.98	134E 00 081E-00 365E-01 762E-02 162E-03	0.1699 0.1208 0.4659 0.3301 0.1336 0.37594	5E-00 9E-01 6E-02 2E-03	0.4349; 0.12284 0.13401 0.46857 0.17572 0.36629	7E-00 7E-01 7E-03	-0.2469 -0.7202 -0.2929 0.1434 0.2236 0.68308	1E-01 4E-02 4E-04 1E-05

					x = 0.05						x = 0.55	
		n a	r n	a	i n	t•r n	bi	a i	i	ai a	b,r	b <u>í</u>
			£-34	0=4~554 0=60541 0=68372	E-08 0.449.	24c+07 255=09 14E=09	0:127672+1 +0:269142+0 0:151785+1	9 0.38961	E-03 0.96 E-05 0.76 E-07 0.35 E-10 0.10	7366E-01 856E-03 668E-05 180E-07 725E-09 912E-12	0.315566 0.256248 0.120818	E=0?
					x = 0.10					ı	c = 0.60	
	1 2 3 4 5 6	3.756705	-07	0.356746 0.176126 0.648326	-06 -0.4023	55-09	0+47779E-00 0+19954E-03 -0+76015E-15	0.603136	-05 0.140 -07 0.766 -10 0.277	34E-01 10E-02 34E-04 65E-07 68E-09 55E-11	0.49606E 0.47361E- 0.26473E- 0.97366E- -0.35519E- -0.21342E-	-05 0.20889E-05 -09 0.80758E-09 -12 0.12707E-10
					x = 0.15					х	= 0.65	
	1 2 3 4 5 6	0+334698+ 0+778968+ 0+380746+	-0.9 -0.9	0:12052E= 0:1488E= 0:93202E=	05 0+316108	-08	0.36033E=05 0.21946E=08 0.31914E=10	0.60540E= 0.70301E= 0.92227E= 0.57122E= 0.22950E=(03 0 2215 05 0 2445 07 0 1567 09 0 6544	9E-02 9E-04 9E-06 5E-09	0.75388E-0 0.83414E-0 0.54484E-0 0.23089E-0 0.49345E-1 -0.12065E-1	04
					x = 0.20					х -	0.70	
	1 2 3 4 5	0.12856E-0 0.24399E-0 0.24872E-0 -0.23075E-1 -0.10434E-1 -0.28816E-1	95 96 2	0.28594E-0 0.62660E-0 0.64009E-0 0.85508E-1 0.33573E-1 0.91711E-1	5 0.21274E 6 0.77393E 1 0.11741E 6 -0.69154E	-07 -10 -10	0.15129E-04 0.17184E-07 0.12877E-10 0.1042E-10 0.1040E-11 0.39329E-15	0.78652E-0 0.13269E-0 0.15467E-0 0.11094E-0 0.51940E-0 0.70911E-1	2 0.31960 4 0.40877 6 0.30398 9 0.14667	E-02 E-04 E-06 E-08	0.11130E-0 0.14101E-0 0.1628E-0 0.51953F-0 0.13777E-10 -0.14!73E-11	0 + 10295E=01 0 + 82234E=06 0 + 41489E=08 0 + 18979E=10
	_				x = 0.25			,		x = 1	0.75	
	2 2 3 4 5 5 6	0:25591E-02 0:74512E-05 0:1650E-07 0:16534E-10 -0:1897E-13 -0:52309E-17	0	•55874E-02 •19093E-04 •1352e-07 •36704E-10 •38793E-13 •16645E-16	0.10085E-0	06 09 .1 .2	0.45887E-06 0.82342E-07 0.10095E-09 0.11581E-10 0.25046E-11 0.45198E-14	0:10075E-00 0:18914E-02 0:25028E-04 0:20573E-06 0:11099E-08 0:32160E-11	0.44913E 0.65890E 0.562625	-02 -04 -06 -08	0.16019E=01 0.23014E=03 0.19810E=05 0.11087E=07 0.37812E+10 -0.14345E=11	0.164315-03
				3	x = 0.30					X ≠ 0.	.80	
3 4 5		0.45294E=02 0.18560E=04 0.41609E=07 0.57421E=10 -0.72401E=13 -0.55815E=16	0 • 0 •	96509E-02 47420E-04 11201E-06 16022E-09 27850E-12 17758E-15	0.14305E=0 0.36159E=0 0.55450E=0 0.69841E=1 -0.65306E=1: -0.37156E=1	6 9 !	0 = 113335-05 0 = 293935-04 0 = 428945-05 0 = 816885-11 0 = 241635-11 0 = 316175-12	0.12630E-00 0.26416E-07 0.29261E-04 0.36651E-06 0.22543E-08 0.83787E-11	0.14846E- 0.61699E- 0.1029E- 0.1000E- 0.63043E- 0.27834E-	-02 -03 -05	0.22528E-01 0.36430E-03 0.35474E-05 0.22573E-07 0.97098E-10	0-11249E-01 0-25370E-03 0-26805E-05 0-17676E-07 0-87269E-10 0-11276E-11
				x	= 0.35					x = 0.6	15	
1 2 3 4 5 6	-	0.74032E-02 0.40175E-04 0.12218E-06 0.21940E-09 -0.43150E-13 -0.41767E-15	0 . 1 0 . 6 0 . 6	5298E-01 0227E+03 2862E+06 1120E-09 5193E-12 3063E-14	0.31209E~03 0.10662E-05 0.20616E-09 0.18824E-11 -0.46559E-12 -0.98853E-13	0	0:24267E-03 0:86042E-06 0:16902E-08 0:43000E-11 0:31427E-12 0:14622E-13	0e15536E-00 0e36251E-02 0e59931E-04 0e69019E-06 0e43793E-08 0e19762E-10	0.16731E=0 0.83072E=0 0.15635E=0 0.17156E=0 0.12214E=0 0.61382E=1	13 15 17	0431010E-01 0556140E-03 0.61367E-05 0.43977E-07 0.21952E-09 0.94314E-12	0.13939E-01 0.38025E-03 0.45729E-05 0.4117E-07 0.18478E-05 0.24340E-11
				x =	0.40					x × 0.90)	
1 2 3 4 5 6	-0	9411428E=01 Ce78460E=04 3e31077E=06 0e73062E=19 0e50492E=12	0.19 0.83 0.20 0.49	749E-01 889E-03 440E-06 422E-08 092E-11 993E-14	0.61518E-03 0.27241E-05 0.69371E-09 0.185855-10 -0.22972F-11 -0.56236E-12	0.	46760E-03 71774E-05 54724E-08 32825E-10 55968E-11 38205E-12	0.18791F-00 0.49000E-02 0.89309E-04 0.10501E-05 0.81733E-08 0.41286E-10	0.18459F-00 0.10986F-01 0.23175E-03 0.28517F-05 0.22773E-07 0.12718E-09	000	0.41814E-01 0.84490E-03 0.10299F-04 0.82641E-07 0.45468E-09 0.24957E-11	0.165685-01 7.55477E-05 0.75545E-01 0.65417E-07 0.36314E-09 0.37817E+11
				x = (0.45					x = 0.95		
1	0. 0.	:16d96E=01 :14179E=03 :70769E=06 :21063E=08 :12308E=11 :10775E=13	0.357 0.189 0.584 0.143	84E-01 42E-03 66E-05 17E-08 40E-10 08E-11	0:11227E-02 0:62348E-05 0:19886E-07 0:47467E-10 -0:27786E-11 -0:10587E-11	0 ± 4 0 ± 1 n ± 5 0 ± 5	33046E-03 9304E-05 6053E-07 9549E-10 -1260E+11 8632E-12	0.22214E-00 0.65369E-02 0.13028E-03 0.17017E-05 0.14762E-07 0.87459E-10	0 - 19956E-00 0 - 14294E-01 0 - 13606E-03 0 - 46084E-05 0 - 41019E-07 0 - 25545E-09	0. 0.	55246E=01 12450E=02 16781E=04 14940E=06 92838E=09 97909E=12	0.18767E-01 0.78955E-01 0.12124E-04 0.11375E-06 0.74551E-05 0.765533F-11
				x = 0.					x	= 1.00		
	7 • 1 2 • 1 2 • 1	241485-01 24995-03 147725-05 542755-08 138745-19	0.4371 0.6034 0.1950 0.1496 0.34471 0.1525	3E-03 4E-05 1E-07 8E-10	0.19286E-02 0.13088E-04 0.51090E-07 0.12093F-09 0.98655E-13 0.45935E+12	0.10 0.41 0.95	88135-02 2245-04 3245-07 1346-10 0245-11 2935-12	0.25833E-00 0.86215E-02 0.18645E-03 0.26692E-05 0.25845E-07 0.17364E-09	0.21167E-00 0.18326E-01 0.47779E-03 0.72612E-05 0.71648E-07 0.49394E-09	0 0 2	71507E-01 17507E-01 26700E-04 26261E-06 8117E-09 23616E-11	0=20049E-01 0=109805-02 0=189565-04 0=191855-06 0=141795-08 0=847515-11

						x = 1.05											
		n	a _r		a ⁱ n		br n		b <u>í</u>		_			x = 1.55	;		
		3 0,	29513E-00 11257E-01 26231E-03	0.0	22063E+00 23175E+01 56730E+03	0.	90629E=01		n 9837E-0 9365-0		2 <u>r</u> •543645 :	oo o.	ai 194145-0		br n	$b_{\mathbf{n}}^{i}$	
		. č	41547E-05 41994E-07 12623E-00	0.1	1183F-04 2173F-06 2725E-09	0.4	4920E-06 4920E-06 4201E-08 6873E-10	0.28 0.33 0.26	9395-04 4585-06 5185-06 8325-10	0	107495-0 419145-0 132905-0 301746-0 470555-0	0.00	12564E-0 92860E-0 33998E-0 81204E-05	0 0	• 13171E-0 • 41326E-0 • 14381E-0 • 32647E-0 • 53182E-0 • 64565E-0	0.820985-0 0.712355-0 0.206395-0	3
		1 0.3 2 0.1	31535-00	0.22	617F=00									x = 1.60		01417416=0	B
		3 0.30 0.62 5 0.73	4570E=01 5342E=03 1887E=05 029E=07 828E=09	0 • 2 8 0 • 91 0 • 16 • 0 • 20 1	935E-01 700E-03 868E-04 164E-06 947E-08	0 ± 35 0 ± 63 0 ± 74 0 ± 624	242E=00 810E=02 345E=04 926E=06 882E=08 644E=10	0.198	4E-08	0.1 0.5 0.1 0.4;	5142E 00 2903E-00 2962E-02 7631E-03 2523E-05 1618E-07	C.1 C.4 C.4	8831E-00 3881E-00 1461E-01 758E-03 399E-04 302E-06	0.3 0.5 0.1 0.4 0.8	5165E-00 0963E-01 9215E-02 6310E-04 0213E-06 0362E-07	-0.12924E-00 0.75298E-01 0.03021E-05 0.28594E-04 0.54911E-06 0.75031E-08	
		1 0.36s	662E-00 07E-01	0.229	09E=00		44E=00						x	= 1.65			
		3 0,496 4 0,994 5 0,118	57E-03 30E-05 47E-06 98E-08	0.356 0.124 0.2496 0.3264 0.2983	71E-01 6E-02 6E-04 3E-05 7E-08	0:493 0:948; 0:122; 0:121; 0:7179	96E-02 7E-04 .9E-05 4E-07	0.12607 0.25938 0.63276 0.88792 0.84559 0.61624	E-01 E-04 E-06 E-08	0.15 0.665 0.231 0.592	539E 00 527E=00 682E=02 94E=03 63E=05	0 •15; 0 •140 0 •563 0 •158	08E-00 08E-00 38E-01 91E-03 27E-04 03E-06	0 • 61 ° 0 • 25 ¢ 0 • 64 9 0 • 11 9	051E-00 986E-01 441E-02 989E-04 945E-05 93E-07	-0.14647E-00 0.59945E-02 0.11641E-02 0.39153E-08 0.40567E-05 0.11746E-07	
	1 2	0.3796	3 E-0 o	0.22915	x = 1.								x =			O#11/#6E=07	
	5 5	0:2384 0:6700 0:1364 0:18817 0:18295	E-03 E-04 E-06	0.43514 0.16586 0.36316 0.517388 0.514788	E-01 E-02 E-04	0:16204 0:67227 0:13959 0:19517 0:192979 0:132368	E-02 E-03 E-05 E-07	0.47219E 0.33111E 0.490905E 0.13990E 0.14554E 0.11571E	-02 -04 -07	0.1987 0.1797 0.6332 0.3027 0.81816 0.15961	8E-00 0E-02 BE-03 5E-05	0+1777 0+1618 0+1707 0+7552 0+21746 0+43802	4E-00 8E-00 0E-01 7E-03	0.3887 0.7432 0.3339 0.9028 0.1757 0.2557	4E=01 1E=02 6E=04 7E=05	-0.162385-00 0.104717-02 2.143455-02 0.529565-04 0.116715-05 0.116715-05	
	1 2	0.429966 0.301916	-01	• 22703E-	-00	0•18846E.	-00						x = 1.	7 5			
	3 4 5 6	0.89395E 0.19621E 0.29326E 0.31004E	-03 -04 -06	.52451E- .21880E- .51989E- .80421E- .86489E-	02 04 06	0.903415- 0.202325- 0.305825- 0.327535- 0.251525-	02 0 03 0 05 0	0.62006E-0 0.41364E-0 0.12832E-0 0.21605E-0 0.24628E-0 0.26191E-0) r 3 5 7	0.55868 0.208758 0.103835 0.392256 0.111845 0.231076	-00 -01 -01	0 :173036 0 :170656 0 :206106 0 :969056 0 :295686 0 :632146	-00 -01 -03	0:40646 0:876198 0:876198 0:434558 0:25573E 0:25573E	-01 -02 -03 -05	-0.17679E-00 -0.12431E-02 0.17369E-02 0.7080AE-04 0.16701F-05 0.27632E-07	
	2	0.45723E- 0.37972E-	21	22323E-0	0 4	21492E-0	ıo						x = 1.80				
	3 4 5	0.11805E= 0.27811E=0 0.44687E=0 0.51380E=0	2 0. 16 0.	62512E=0 28534E=0 73340E=0 12279E=0 4362E=07	0.00	11993E-0 28911E-0 47089E-0 54474E-0 6036E-09	1 0. 3 0. 7 0.	19946E-01 50421E-02 17816E-03 32751E-05 60182E-07		0.55640E 0.23902E= 0.12889E= 0.50480E= 0.15153E= 0.33087E=	00	0 • 16870E- 0 • 17695E- 0 • 24712E- 0 • 12338E- 0 • 39885E- 0 • 90239E-0	00 d 01 0 02 0 4 0	0.42388E= 0.10224E= 0.56084E= 0.16941E= 0.36812E= 0.59920E=	00 -0 02 0 03 0	0+18957E-00 •71495E-0; •20704E-0; •93615E-04 •23622E-05	
1		0-48119E-00		18255-00	1.35								x = 1.85			41525E-07	
5 6		0.47647E=01 0.15646E=02 0.36901E=04 0.57614E=06 0.83646E=08	0.36 0.10 0.18	8657E-01 816E-02 206E-03 440E-05 280E-07	0.1 0.4 0.7 0.8	4076E+00 5733E+01 5773E+03 537E+05 1854E+07 440E+09	0 + 59 0 + 24 0 + 48	6049E=01 9909E=02 1348E=01 1796E=05 1820E=07 1821E=09	0	0.55210E 0: 0.27007E-0: 0.15942E-0: 0.64553E-0:3 0.20357E-0:4 0.46896E-0:6	0.	15468E+00 18049E+00 27426E-01 15596E-02 53295E-04	0 • 0 • 0 • 0 • 0 • 5 • 6 • 5 • 6 • 5 • 6 • 5	44110E-00 11731E-00 71792E-03 22905E-03 2459E-05	-0. 0.2 0.3	200545-00 147605-01 241635-02 22416-07 30425-05	
1	o	•50177E 00	0.212		1.40							_	= 1.90		U # D	16775-07	
	0	•58885E+01 •20036E-02 •3751E-04 •10029E-05 15311E-07	0 + 85 77 0 + 4702 0 + 1402 0 + 2726 0 + 3705	78E-01 74E-02 78E-01 8E-01	0.203 0.568 0.106 0.1423 0.1406	95-04 36-04	-0.538 0.6909 0.3277 0.7153 0.1026 0.1067	945-02 736-03 145-05 165-04	0 e 2 0 e 2	54598E 00 30111E-00 19650E-01 32065E-03 :7131E-04 :5833E-06	0.1 0.3 0.1 0.7	6093E=00 8118E+00 4796E=01 9579E+02 1625E=04 7839E=05	0,45 0,13 0,91 0,30 0,740	819F=00 271E=00 151E=02 715F=03 046E=05	70 + 24 1 0 + 27 1 0 + 45 1	9945-06 0825-01 5605-02 9345-03 7325-05 7755-07	
	0.5	51897E 00 72954E=01	0 - 20 642	E-00	0.2889	15						x =	1.95			722 -07	
	0.2 0.7 0.1	5797E-02 3432E-04 4665E-05 0678E-07	0 498688 0 - 59508 0 - 19043 0 - 39748 0 - 579845	E-02	0+26132 0+78232 0+15665 0+22423 0+23784	E-01 E-03 E-04 E-06	-0.72732 0.76908 0.43469 0.10328 0.13979 0.178008	E-02 E-03 E-04 E-06	0.24 0.10 0.35	1827E 00 1136E-00 141E-01 976E-02 887E-04	0.408	737E-00 135E-00 137E-01 17E-02 11E-04	0.475; 0.148; 0.1147 0.4086 0.1035	75-00 75-01 95-03 75-04	-0.2174 -0.3693 0.305 0.2026 0.6264 0.1316	185-01 175-07 05-01 75-01	
	0.86	288E 00	0 • 20022E	-00	Calloss	- 00 -	-0.91948E					z = 2	.00				
	0.32	985E-02 271E-04 166E-05 243E-07	0.11210E- 0.74647E- 0.25576E- 0.57170E- 0.89327E-	02 03 05	0-33071E 0-10660E 0-22753E 0-34785E 0-39593E	-01 -02 -04 -06	0.91948E 0.81875E 0.456828E 0.14700E 0.24472E 0.29207E	-02 -03 -04	0.529; 0.360; 0.295; 0.1305; 0.4713 0.1263	6E-01 2E-02 4E-04	0.1539 0.1747 0.47626 0.30261 0.12116 0.33977	0E=00 5E=01 :E=02 :E=03	0.49210 0.16340 0.163312 0.539758 0.143628 0.287208	E-00 E-01 -01	-0.22316 -0.471871 -0.328141 0.256433 0.849765 0.138896	E=01 E=02 E=03 =05	

						x = 0.05								056			
		n	ar n		a i		br n	bi n			ar n	,	ıi.	x = 0.55	r		
		1 0.180; 2 0.195; 3 -0.2460	1E-08	0 • 6 2 2	725E-04 29E-08 65E+12	, 0:171 -0:2739 -0:5676	385-07 745-09		E-07	0.3184 0.3592 0.2642 0.1172 0.31351	15-01 05-03 15-05 95-07	0.5957 9.9900 0.7815	6E-03 2E-05 3E-07 0E-09	0.3216 0.75848 0.12169 0.37534 -0.12698 -0.11089	7E-02 3E-04 9E-06 E-09 E-11	0 • 2486; 0 • 22298 0 • 10826 0 • 35678 0 • 61940 0 • 69688	05-0 6-0 6-0 6-0
					x	= 0.10							x	= 0.60			- ••
	1 2 3 4 5	0.14515 0.76511 0.24411	E-07	0.3661 0.2002 C.1146	4E-05 0E-07	0.5776 0.58384 0.10863	£-07	0.531895- -0.599715- 0.200715-	-1.0	0.43477 0.55820 0.48504 0.25604 0.880288 ~0.539088	E-03 E-05 E-07	0.76055 0.15246 0.14307 0.780341 0.2780341 0.278081	E-02 E-04 E-07 E-09	0.507701 0.478478 0.266298 0.967155 0.261898 -0.412508	-04 -06 -09 -11	0.375229 0.406358 0.235298 0.872169 0.540928 0.4618079	-04 -06 -09 -11
	1	1.494036	-03			0.15							x ·	0.65			
	3 4 5	0.49492E 0.53243E 0.28503E	-06 -09	0+12375 0+15197 0+3593a	E-05	0.439656 0.260968 -0.300838	-09	0.40439E=0 0.25559E=0 =0.36514E=1	l d	0.58140E 0.83888E 0.84806E 0.52471E 0.21351E 0.28258E	-03 -05 -07	0.94491E 0.22669E 0.24938E 0.15965E 0.66534E 0.23417E	-02 -04 -05 -09	0.775185- 0.84415E- 0.54857E- 0.23253E- 0.83393E- -0.22566E-	06 08 11	0.54345E- 0.470442E- 0.848021E- 0.820754E- 0.898020E- 0.89451E-	04 06 08
					x = (0.20							X = 1	0.70			_
	1 2 3 4 5 6	0.11911E- 0.22456E- 0.22362E- 0.98657E- -0.95636E- -0'226094E-1	05 08 12 15	0.29377E 0.63968E 0.66995E 0.22235E 0.33813F- 0.92028F-	-05 -08 -11 -14 -18	0.18636E- 0.21077E- 0.61829E- -0.67776E- 0.52536E- -0.43225E-	07 11 12 13	0a16983E-04 0a194565-03 0a77825E-11 -0a186095-11 -0a62398E-17 0a40111E-15	•	0.76243E=0 0.12758E=0 0.1275E=0 0.10191E=0 0.47973E=0 0.61138E=1	02 04 06	0.11448E- 0.32710E- 0.41683E- 0.30955E- 0.14956E- 0.57229E-1	02 04 06 08	0:11504E-0 0:14296E-0 0:10717E-0 0:52498E-0 0:19254E-1 -0:72850E-1	3 5 8	0.75786E-0 0.11699E-0 0.92852E-0 0.46640E-0 0.23318E-10 0.15636E-11	3 6 8
	1	0.237455+0	,		x = 0								x = 0.	7 5			
	2 3 4 5 6	0.685915-0 0.10685E-0 0.133495-1(-0.103195-13 -0.473615-13	5 7)	0.57447E-1 0.19493E-0 0.31905E-0 0.30048E-1 0.39777E-1 0.16703E-1	04 07 10	0.57296E-0 0.10090E-0 0.10793E-0 0.39755E-1 0.10349E-1 0.485A7E-1	6 9 !	0.51595E+04 0.92406E+07 0.75062E+10 -0.18830E+11 0.28974E+12 0.447584E+14		0.98099E-0: 0.17491E-0: 0.23022E-0: 0.18899E-0: 0.10228E-0:8	2 4 5	0 • 13941E-01 0 • 45993E-03 0 • 67200E-04 0 • 57297E-06 0 • 31586E-08 0 • 12232E-10	?	0.16652E-01 0.23385E-03 0.17996E-05 0.11167E-07 0.43110E-10 0.58854E-13		0.10191E-01 0.18718E-03 0.17133E-05 0.98572E-08 0.39771E-10 0.17587E-12	
					x = 0.3	0							x = 0.8	0		-	
5		0.42112E-02 0.17088E-04 0.38221E-07 0.55073E-10 -0.57323E-13 -0.50499E-16	0	0.49420E-0: 0.48420E-0: 0.11410E-0: 0.15366E-0: 0.32070E-12	• (5 5 7	0.14390E=03 0.36260E=06 0.51911E=09 0.49401E=11 0.14483E=11 0.37381E=13	-	0 = 12764E-03 0 = 33008E-06 0 = 44695E-09 0 = 37366E-11 0 = 11544E-12 7 = 40507E-13		0.12385E-00 0.24458E-02 0.36127E-04 0.33671E+06 0.20758E+08 0.93053E-11		0.15651E-00 0.632Z4E-02 0.10498E-03 0.10186E-05 0.64121E-08 0.26639E-10		0+23563E+01 0+37102E+03 0+35859E+05 0+22735F+07 0+10402E+09 0+96885E+12		0+13218E-01 0+28977E-05 0+30347E-05 0+19890E-07 0+89454E-10 0+68812E-12	
					x = 0.35							x	= 0.85				
1 2 3 4 5 6	_	0.69013E-02 0.36994E-04 0.11232E-06 0.20337E-09 0.22180E-12 0.37373E-15	0 0	.15764E-01 .10444E-03 .33490E-06 .63192E-09 .13446E-11 .13165E-14	0.	31452E-03 10709E-05 21214E-09 90972E-11 70006E-12 15086E-12	0 0 0	-273875-03 -967035-06 -191545-08 -156535-10 -313385-11 -116975-12		0.15341E-00 0.33614E-02 0.55153E-04 0.57898E-06 0.40209E-08 0.18697E-10	0	0.17690E-00 0.65190E-02 0.15951E-03 0.17475E-05 0.12428E-07 0.62026E-10	0	• 32646E=01 • 57324E=03 • 62100E=05 • 64317E=07 • 22019E=09 • 31574E=12	0	*16505E=01 *43557E=00 *51852E=00 *38503E=07 *20271E=00 *16663E=11	
				x	= 0.40							х :	0.90				
1 2 3 6 5 6	0	-10688E-01 -72282E-04 -28560E-06 -67384E-09 -42776E-12 -20962E-14	0 0 0 0	23478E-01 20315E-03 25029E-06 20663E-08 3098E-11 23861E-14	0 • 2 0 • 6 0 • 1 • 0 • 1	52133E+03 7376E+05 9138E+08 6476E+10 5252E+12 5079E+12	0. 0. 0.	529015-03 244995-05 625635-00 182745-10 226825-11 264975-12	000	•18638E-00 •45515E-02 •52214E-04 •96493E-06 •75123E-08 •39906E-10	0 • 0 •	19566E-00 11276E-01 23649E-03 29049E-05 23174E-07 13011E-09	0 • : 0 • : 0 • 4	44311E-01 96517E-03 10427E-04 93183F-07 16701E-09 14186E-12	0 •	197725-01 637525-03 857965-05 716035-07 120435-09 175815-11	
r		10000			0.45							X = -	0.95				
	0. 0. 0.	15866E-01 13062E-03 65047E-06 19431E-08 32670E-11 92850E-14	0.34 0.15 0.55	3275E-01 5515E-03 3329E-05 7362E-08 228E-10 641E-13	0.62 0.19 0.53 0.10 -0.34	369E=02 728E=05 945E=07 764E=10 348E=11 665E=12	0.5 0.1 0.4 0.3	42245-03 55405-05 79685-07 71585-10 82455-12 73505-12	0. 0.	22211E-00 50847E-02 11997E-03 15538E-05 13557E-07 30582E-10	0.1	1194E-00 4686E-01 4301E-03 6950E-05 1754F-07 6093E-09	0 • 1 : 0 • 1 : 0 • 9 •	9921E-01 2789E-02 7032E-04 6090E-06 136F-09 547E-11	0.9 0.1 0.1 0.6	2583E-01 1052E-03 3793E-04 2866E-06 4013E-09 1125E-11	
	3.2	2783E-01	0.45	x = (2925-01		*05-03						x = 1.	.00				
	0.1 0.1 0.4 0.1	22016-03 35798-05 98758-08 82176-10 88668-13	0.616 0.402 0.153 0.360	04E-03 04E-05 34E-07 17E-10 01E-17	0 • 131 0 • 513 0 • 125 0 • 101	89E-02 84E-04 44E-67 54E-09 51E-11 44E-12	0.11 0.45 0.10 -0.20	7255-02 5315-04 1215-07 7365-09 7805-11 695-11	0.8 0.1 0.2 0.2	5969E=00 J446E=02 7176E=03 4717E=05 3733E=07 5050E=09	0 • 1 6 0 • 4 8 0 • 7 2	25095-00 18505-01 7815-03 9845-05 9285-07 2505-09	0 • 18 0 • 27 0 • 26 0 • 182	709E=01 555E=02 144E=04 647E=06 90E=08 127E=11	0.12 0.21 0.22 0.15	343E-01 711E-02 606E-04 402E-06 930E-08	

			x = 1.05				x = 1.55	
	n a	[1	an bi	Бi h	ar n	ai n	b.r n	$\mathbf{b_{n}^{i}}$
	1 0.29804 2 0.10534 3 0.24176 4 0.38194 5 0.40398 6 0.297218	E=01 0.2387 E=03 0.6814 E=05 0.1139 E=07 0.1239	0E-01 0.26472 9E-03 0.42311 5E-04 0.45458 2E-06 0.34469	0.2434[E-0 1E-02	1 0.55346E 2 0.10595E 4 0.39066E 6 0.12258E 6 0.27753E	00 0.198125 -00 0.198125 -02 0.95371E -03 0.347145 -05 0.82751E	-00 0.344098 -00 0.446128 -02 0.150398 -03 0.335478 -05 0.541828	-00 -0.12177E-00 -01 0.10212E-01 -02 0.86214E-03 -04 0.23808E-04 -06 0.42181E-06
			x = 1.10				x = 1.60	
	1 0.33608E 2 0.13678E 3 0.33513E 4 0.57822E 5 0.67069E 6 0.54806E	-01 0.29847 -03 0.93660 -05 0.17191 -07 0.20529	E-01 0.37194E E-03 0.64640E E-04 0.75918E E-06 0.63089E	-02 0:23230E-02 -04 0:49477E-04 -06 0:62616E-04 -08 0:54160E-08	0.12802E= 0.49672E= 0.16272E= 0.39122E=	00 0.14663E- 02 0.11782E- 03 0.45716E- 05 0.11618E-	00 0.55444E- 01 0.20167E- 03 0.47685E- 04 0.81832E-	01 0.95067E-02 02 0.10997E-02 04 0.33073E-04 06 0.62827E-06
			x = 1.15				x = 1.65	
	1 0.37275E- 2 0.17627E- 3 0.45820E- 4 0.85923E- 5 0.10882E- 6 0.97295E-	01 0.36879E 03 0.12689E 05 0.25448E 06 0.33235E	-01 0,515265- -02 0.969656- -04 0.123956- -06 0.112346-	-02 0.30443E-02 -04 0.72597E-04 -05 0.10095E-05 -07 0.95573E-0F	0.56443E 0 0.15305E-0 0.62948E-0 0.21420E-0 0.54558E-0 0.10023E-0	0 0.15991E-0 2 0.14446E-0 3 0.59661E-0 5 0.16133E-0	00 0.67587E-0 0 0.26802E-0 3 0.67064E-0 4 0.12203E-0	01 0.76512E-02 02 0.13825E-02 04 0.45382E-04 05 0.92324E-06
			x = 1.20				x = 1.70	
3 4 5	0.22563E=0 0.61672E=0 0.12552E=0	0.45047E= 3 0.16957E= 4 0.37023E= 6 0.52679E=	-01 0.70450E=0 -02 0.14305E=0 -04 0.19823E=0 06 0.19528E=0	02 0.39071E-02 03 0.10456E-03 05 0.15928E-06 07 0.16442E-07	0.56561E 00 0.18081E-00 0.78240E-02 0.27977E-03 0.75319E-05 0.14874E-06	0 +17157E-00 0 +17509E-01 0 +77197E-03 0 +22171E-04	0 0.81115E-01 0.35315E-01 0.93380E-04 0.17982E-05	1
			x = 1.25				x = 1.75	
1 2 3 4 5 6	0.43875E-00 0.28699E-01 0.82616E-03 0.18054E-04 0.25945E-06 0.28499E-08	0.54414E-0 0.22380E-0 0.53011E-0 0.81895E-0	0.95133E-0 0.20784E-0 0.31107E-0 0.33166E-0	2 0.49070E-02 3 0.14799E-03 5 0.24637E-05 7 0.27679E-07	0.55423E 00 0.21086E-00 0.97802E-02 0.36283E-03 0.10300E-04 0.21241E-06	0:17139E-00 0:18097E-50 0:21269E-01 0:99087E-03 0:30171E-04 0:54380E-05	0.41887E-00 0.99797E-01 0.46145E-02 0.12880E-03 0.26203E-05 0.40095E-07	-0.18959E-00 -0.6909TE-03 042085E-02 0.82590TE-04 0.192019-05 0.31532E-07
			x = 1.30			x	= 1.80	
1 2 3 4 5 6	0.46695E-00 0.36281E-01 0.10926-02 0.25597E-00 0.41256E-06 0.47202E-08	0.23485E-00 0.65001E-01 0.29200E-02 0.74794E-04 0.12505E-05 0.14617E-07	0.12693E=01 0.29774E=03 0.47967E=05 0.55198E=07	-0.21450E-01 0.60219E-02 0.20606E-03 0.437410E-05 0.45596E-07 0.40518E-09	0.56093E 00 0.24255E-00 0.12183E-01 0.46737E-03 0.13960E-04 0.30420E-06	0-16595E-00 0-18797E-00 0-25534E-01 0-12622E-02 0-40680E-04 0-91914E-06	0.43545E-00 0.11131E-00 0.59801E-02 0.17608E-03 0.37778E-05 0.61060E-07	-0.20225E-0G -0.76136E-02 0.24966E-02 0.10339E-03 0.27207E-0E 0.47460E-07
		,	c = 1.35			x :	- 1.85	
1 2 3 4 5 5	0.49155E-00 0.45586E-01 0.14503E-02 0.35816E-04 0.62138E-05 0.76646E-08	0.22851E-00 0.76781E-d1 0.37695E-02 C.10410E-03 0.18781E-05 0.23691E-07	0.25825E-00 0.16739E-01 0.42104E-03 0.472780E-05 0.90101E-07 0.82569E-09	-0 = 39 782E-01 0 = 72023E-02 0 = 26248E-03 0 = 55835E-05 0 = 73653E-07 0 = 70582E-09	0.59473E 00 0.27908E-00 0.15126E-01 0.59829E-03 0.18762E-04 0.49126E-06	0.16097E-00 0.19302E-00 0.30452E-01 0.15952E-02 0.54370E-04 0.12988E-05	0.45186E-00 0.12729E-00 0.76867E-02 0.23868E-03 0.53925E-05 0.91910E-07	-0.21302E-00 -0.15481E-01 0.29301E-02 0.14348E-02 0.38128E-02 0.70570E-07
		x	= 1.40			x •	1.90	
2 3 4 5 6	0.51243E 00 0.5691&E=01 0.18580E=02 0.49507E=04 0.92130E=06 0.12224E=07	0:22128E-00 0:89655E-01 0:48181E-02 0:143010-03 0:27776E-05 0:37708E-07	U+28327E-00 0+21814E-01 0+58825E-03 0+10879E-04 0+14448E-06 0+1424E-08	-0.597865-01 0.83659E-02 0.38148E-03 0.2007E-05 0.11680E-06 0.12036E-06	0.54709E 00 0.50791E-00 0.18724E-01 0.78148E-03 0.25017E-04 0.60956E-06	0.15641E-00 0.1912ZE-00 0.36071E-01 0.20048E-02 0.72064E-04 0.18174E-05	0.46821E-00 0.14334E-00 0.97992E-02 0.32095E-03 0.76247E-05 0.13683E-06	-0.22190E-00 -0.27213E-01 0.33624E-02 0.18617E-03 0.52872E-05 0.10375E-06
		x =	1.45			x = 1	1,95	
1 2 3 4 4 5 6 6 6 6 6 6 6 6 6 6 6 6 6 6 6 6 6	0:52964E 00 0:70569E-01 0:23957E-02 0:67663E-04 0:13485E-05 0:19175E-07	0.21360E-00 0.10343E-00 0.61015E-02 0.19433E-03 0.4049E-05 0.59020E-07	0.30652E-00 0.28088E-01 0.81269E-03 0.16036E-04 0.22790E-06 0.24102E-08	-0.906085-01 0.93846-07 0.907775-05 0.18646-04 0.182055-06 0.201755-08	0+33785E 00 0+33896E-00 0+23107E-01 0+96409E-03 0+33108E+04 0+84262E-06	0.15216E-00 0.18831E-00 0.42432E-01 0.25017E-02 0.94767E-04 0.25197E-05	0.48454E-00 0.15914E-00 0.12388E-01 0.42827E-03 0.10684E-06 0.20160E-06	+0.22888E-00 -0.39602E-01 0.37565E-02 0.23900E-03 0.72577E-05 0.15084E-06
		X o	1.50			и • 2.	00	
	0.54327E 00 0.886836E=01 0.80684E=02 0.91516E=04 0.91516E=05 0.29610E=07	0.20579E-00 C.11782E-00 0.76598E-02 0.26107E-03 0.58251E-05 0.90930E-07	0.32506E-00 0.35713E-01 0.1110E-02 0.23335E-04 0.35394E-06 0.40044E-08	-0.10148E-00 0.10078E-01 0.66633E-01 0.16920E-04 0.2792ZE-06 0.33166E-08	0.52724E 00 0.36860E-00 0.28431E-01 0.12144E-02 0.43508E-04 0.11624E-05	0.14820E-00 0.18268E-00 0.49558E-01 0.31022E-02 0.12369E-03 0.34626E-05	0.50083E 00 0.17445E-00 0.15525E-01 0.55729E-03 0.14844E-04 0.29408E-06	-0.23395E-00 -0.53350E-01 0.40602E-02 0.3036E-03 0.98655E-05 0.21701E-06

			x = 0.05				x = 0.55	
1	n ar	a i	ρ <mark>ι</mark>	$\mathbf{b_n^i}$	ar n	a_n^i	br	$\mathfrak{b}_{\mathbf{n}}^{\mathbf{i}}$
1 2 3 4 5	0.15999E=04 0.19096E=08 -0.23117E=12	0.47823£-0.	-0.00041E-1	o -0.10162E-09	0.28869E-01 0.31868E-03 0.23433E-05 0.10411E-07 0.37408E-10 -0.48047E-13	0.6C422E-01 0.98994E-03 0.78016E-05 0.35728E-07 0.10835E-09 0.38137E-12	0.30384E-02 0.24146E-04 0.11312E-05 0.35411E-09 0.19807E-11 +0.15400E-12	0.24772E-04 0.11955E-06
			x = 0.10				x = 0.60	
1 2 3 4 5	0.12675E-03 0.62002E-07 -0.25137F-10	0.36701E-03 0.1999E-06 0.24763E-19	0.53430E-06 -0.72522E-10 -0.75119E-10	C.10499E-09	0.439649E-01 0.49536E-03 0.43020E-05 0.22725E-07 0.79416E-10 -0.99583E-13	0.77386E-01 0.15248E-02 0.14284E-04 0.77863E-07 0.28068F-09 0.11432E-11	0.48181E-02 0.44764E-04 0.24829E-05 0.72243E-07 0.58939E-11 -0.22941E-12	0.42816E=02 0.45265E=04 0.26032E=06 0.96755E=09 0.95291F=11 0.17763F=11
			x = 0.15			,	c = 0.65	
1 2 3 4 5	0.43916E-03 0.47233E-06 0.28702E-04	0.12407E-02 0.15181E-05 0.87799E-09	0.40839E-05 0.26125E-08 0.20160E-10	0:444435-05 0:270845-08 -0:272855-10	0.53368E-01 0.74471E-03 0.75217E-05 0.46538E-07 1.147445-09 0.32751E-13	0.96514E-01 0.22678E-02 0.24898E-04 0.15929E-06 0.66421E-09 0.23234E-11	0.73975E-02 0.79110E-04 0.51175E-06 0.21844E-08 0.6156ZE-11 -0.68961E-12	0.62603E-02 0.76709E-04 0.53223E-06 0.23026E-08 0.12101E-10 0.98479E-12
		,	= 0.20			x	= 0.70	
! 2 3 + 5	0.10576F=02 0.19913E=05 0.19905E=03 -0.82216E=11 -0.86322E=15 -0.23736E=18	7.29466F-02 0.51908E-05 0.66685E-03 0.21378E-11 0.34392E-14 0.93871E-18	J.17319F-04 J.19443E-07 -0.16970E-10 -0.28924E-11 -0.28551E-11 -0.39180E-15	0.18496F=04 0.21370E=07 -0.80023E=11 -0.78223E=11 0.38483E=13 0.42945E=15	0.70487E-01 0.10887E-02 0.12617E-04 0.90399E-07 0.42633E-09 0.88861E-12	0.11744E-00 0.32734E-02 0.41620E-04 0.30586E-06 0.14910E-08 0.57214E-11	0.11051E-01 0.13426E-05 6.10007E-05 0.48902E-06 0.18590E-10 -0.35671E-12	0.38291E-02 0.13117E-03 0.10308E-09 0.91500E-08 0.23768E-10 0.14396E-11
•		x	= 0.25			x :	= 0.75	
1 2 3 4 5	0.21107F-02 0.60827E-05 0.94757E-08 0.99977E-11 -0.84733E-14 -0.443078E-17	2.576575-07 0.194765-04 0.315545-07 0.315175-10 0.395965-13 0.170365-15	0.5330 AE - 04 0.93775E - 07 0.10380E - 09 0.2378 2E - 11 0.2643 BE - 12 - 0.4277 2E - 14	0.56911E-04 0.10169E-05 0.99556E-10 0.12160E-11 0.51802E-12 0.49148E-14	0.91392E-01 0.15545E-07 0.20421E-04 0.16765E-06 0.90901E-09 0.32685E-11	0:13959F-00 0:460508-02 0:671048-04 0:571718-06 0:216338-08	0+16119E-01 0+22006E-03 0+18693E-05 0+10430E+07 0+43836E-10 0+74516E-12	0.12035E-01 0.21068E-03 0.19056E-05 0.10911E-07 0.46869E-10 0.12367E-11
		x:	0.30			x =	0.80	
1 2 1 4 5 6	0.37495E-02 0.15155E-04 0.33945E-07 0.50724E-10 -0.63409E-13 -0.45975E-16	0,997755-02 0.48380E-04 0.11392E-06 0.16698E-09 0.29023E-12 0.18177E-15	C+13407E=03 O+32738E=06 O+55819E=09 O+14117E=10 O+75573E=17 -0+21433E=13	0.14114E-03 0.36338E-06 0.53148E-09 0.12801E-10 0.10568E-11 0.25694E-13	0.11632E-00 0.21754E-02 0.32041E-04 0.29867E-06 0.18397E-08 0.718215-11	0.16220E-00 0.63335E-02 0.10484E-03 0.10164E-J5 0.64021E-08 0.29399E-10	0.23014E-01 0.35001E-03 0.33562E-05 0.21237E-07 0.97789E-10 -0.12867E-12	0.15861E-01 0.32759E-03 0.33827E-05 0.22077E-07 0.10840E-07 0.25075E-11
		x =	0.35			x =	0.85	
1 2 3 4 5	0.61584E-02 0.32807E-04 0.99609E-07 0.18719E-09 0.20697E+12 -0.33821E-15	0.15857E-01 0.10436E-03 0.37470E-06 0.61720F-09 0.17886E-11 0.13449E-14	0.79355E-03 0.99573E-06 0.19345E-08 0.11027F-10 0.24140E-11 -0.65484E-13	0.303775-03 0.106605-05 0.203355-08 -0.347125-11 -0.824185-14 0.251805-12	0:14529F-00 0:29929E-02 0:48926E-04 0:51358E-06 0:35695E-08 0:17040E-10	0.18436F-90 0.85395E-02 0.15932E-03 0.17438E-05 0.12396E-07 0.62288E-10	0.327115-01 0.542295-03 0.581955-05 0.413835-07 0.208725-09 0.547555-12	0.201975-01 0.494885-03 0.579315-05 0.427125-07 0.222235-09 0.203055-11
		x =	0.40			x =	0.90	
1 2 3 4 5	0.0553E-02 0.64104E-04 0.2533E-06 0.60703E-07 0.21391E-12 -0.19080E-14	0.23646E-01 0.20303E-03 0.64870E-06 0.20704E-08 0.54482E-11 0.75788E-14	0.58119E-03 0.25487E-05 0.55466E-08 0.31420E-10 0.19366E-11 -0.32288E-12	0.58899E-03 0.27049E-05 0.58916E-08 0.27091E-10 0.56491E-11 0.56437E-i2	0+17602E=00 0+0577E=02 0+7293BE=04 0+85595E=06 0+66675E=08 0+35673E=10	0.20507E-00 0.11312E-01 0.23624E-03 0.286990E-05 0.23119E-07 0.13074F-09	0.44219E-01 0.82105E-03 0.97854E-05 0.77751E-07 0.44130E-09 0.15769E-11	0.24773E+01 0.72843E-03 0.96100E+05 0.79572E+07 0.46632E+09 0.50752E+11
		x = (). 4 5			x = 0	.95	
1 7 3 4 5 6	0.14250E-01 0.11555E-03 0.57591E-05 0.17267E-08 7.37227F-11 -0.81651E-14	1.33576E-01 1.36497E-01 1.17293E-05 0.59197E-03 0.14064F-10 0.34450E-13	0.10662E-02 0.58452E-05 0.18550E-07 0.52422E-10 0.21236E-11 -0.68565E-13	0.10539E-02 0.61434E-05 0.19790E-07 0.66305E-10 0.30192E-11 0.71379E-12	0.21193E-00 0.54331E-02 0.10644E-03 0.13872E-05 0.12031E-07 0.7279GE-10	0.22339F-00 0.14747E-01 0.34271E-03 0.46856E-05 0.41647E-07 0.26039E-09	0.59525E-01 0.12180E-02 0.16007E-04 0.14111E-06 0.88305E-07 0.38744E-11	0.2911E-01 0.10470E-02 0.15493E-04 0.16419E-06 0.92232E-09 0.70624E-11
		x = 0	.50			x = 1	.00	
1	0.20552E-01 0.10593E-03 0.17943E-05 0.44236E-08 0.1993E-10 0.84195E-13	0.45808E-01 0.61644F-03 0.40192E-05 0.15203E-07 0.36897E-10 0.20793E-12	0:18432E-02 0:12300F-04 0:47806E-07 0:11884E-09 0:18694E-12 0:11380E-12	0:17685E-02 0:12787E-04 0:12787E-04 0:50916E-07 0:12624E-09 -0:39375E-12 -0:94849E-13	0.25214E-00 0.71968E-02 0.15242E-03 0.21925E-05 0.21925E-07 0.16245F-09	0.23852E-C0 0.18948E-01 0.48747E-03 0.73843E-05 0.72752E-07 0.501465-09	0.78500E-01 0.17743E-02 0.25558E-04 0.24853E-06 0.17093E-08 0.47093E-11	0.12590E-01 0.14723E-02 0.2434E-04 0.2434E-05 0.17679E-08 0.99981F-1!

			x = 1.05				x = 1.55			
	n a <mark>r</mark>		an br	$\mathbf{b_{n}^{i}}$	$\mathbf{a}_{n}^{\mathbf{r}}$	a i	b_n^r	$b_{\mathbf{n}}^{\mathbf{i}}$		
	1 0.291586 2 0.944558 3 0.214598 4 0.338838 5 0.358438 6 0.266758	F=02 0.2402 F=03 0.6811 F=05 0.11374 F=07 0.12362	5E-01	-02	0 = 10010E= 0 = 34925E= 0 = 10889E= 0 = 24632E=	00 0.13765E 02 0.95604E 03 0.34703E	-00 0.46763E -02 0.14651Z -03 0.31964E -05 0.51136E	-01 0:14952E-01 -02 0:10346E-02 -04 0:27399E-04 -06 0:47717E-06		
			x = 1.10				x = 1.60			
	1 0.33108E 2 0.12298E 3 3.29753E 4 0.51299E 5 0.59503E 6 0.48729E	-01 0.30086 -03 0.33657 -05 0.17161 -07 0.20481	E-01 0.35912E- E-03 0.61105E- E-04 0.71232E- E-06 0.59027E-	-02 0.27391E-02 -04 0.56143E-04 -06 0.70104E-06 -08 0.60179E-05	0.56929E 0 0.12199E-0 0.44302E-0 0.14460E-0 0.34728E-0 0.60021E-0	0 0.15330E- 2 0.11848E- 3 0.45713E- 5 0.1160E-	00 0.58488E= 01 0.19747E= 03 0.45544E= 04 0.77340E=	01 0-15003E-01 02 0-13331E-02 04 0-38227E-04 06 0-71249E-06		
			x = 1.15				x = 1.65			
		0: 0:37244E 03 0:12690E 05 0:25406E 07 0:331585	-01 0.50037E- -02 0.91870E- -04 0.11644E- -05 0.10514C-	02 0.36289E-02 04 0.82713E-04 05 0.11328E-05 07 0.10633E-07	0.57330E 00 0.14712E-00 0.55943E-02 0.19041E-03 0.48434E-05 0.86946E-07	0 .16813E-0 2 0 .14545E-0 0 .59675E-0 0 .16110E-0	00 0.72033E=0 01 0.26390E=0 03 0.64220E=0 0411551E=0	71 0.13775E-01 72 0.16952E-02 74 0.52703E-04 75 0.10498E-05		
			x = 1.20				x = 1.70			
1 2 3 4 5 6	0.40586E-0 0.20426E-0 0.54969E-0 0.1138E-0 0.15339E-0 0.14947E-0	1 0.45592E- 3 0.16963E- 4 0.36966E- 6 0.52562E-	-01 0.68857E+0 -02 0.13587E+0 -04 0.18649E+0 -05 0.18289E+0	2 0447162E-02 3 0411966E-03 5 0417917E-05 7 0418355E-07	0.57422E 00 0.17536E-00 0.70363E-02 0.24882E-03 0.66674E-05 0.13023E-06	0 :17926-00 0:18138E-00 0:17731E-01 0:77241E-01 0:22143E-04 0:44512E-06	0 0.87252E-0 1 0.34983E-0 3 0.89669E-0 4 0.17049E-0	1 0.10849E-01 2 0.21261E+02 0.71835E-04 5 0.15276E-05		
			x = 1.25				x = 1.75			
1 2 3 4 5 6	0.43935E-00 0.26091E-01 0.73433E-03 0.16021E-04 0.23905E-06 0.25281E-08	0.55210E-0 0.22397E-0 0.52936E-0	0.93656E-02 0.19794E-03 4 0.29305E-05 6 0.31987E-07	0.60100E=02 0.17020E=03 0.27786E=05	0.572325 00 0.206325-00 0.681705-02 0.581705-02 0.322835-03 0.714655-05 0.188535-06	0.17166E-00 0.1923IE-00 0.21469E-01 0.99177E-03 0.30137E-04 0.6425IE-06	0.43930E-00 0.10384E-00 0.46011E-02 0.12405E-03 0.24884E-05 0.37808E-07	-0.19549E-00 0.58550E-02 0.26285E-02 0.96844E-04 0.21963E-05 0.35605E-07		
			x = 1.30			x	: = 1.80			
1 2 3 + 5	3.46943E-00 0.33144E-01 0.9712IE-03 0.22717E-04 0.36604E-06 0.41865E-08	0.25237E-0 0.66146E-0 0.29235E-0 0.74700E-0 0.12479E-0 0.14585E-0	0.12507E-01 0.28442E-03 0.45260E-05 0.51798E-07	-0.11197E-01 0.75014E-01 0.23827E-03 0.2382E-05 0.61056E-07 0.6596E-07	0.56784E 00 0.23934E-90 0.11019E-01 0.41609E-03 0.12399E-04 0.27001E-06	0.16399E-00 0.20020E-00 0.25824E-01 0.12639E-02 0.40640E-04 0.91738E-06	0.45474E=00 0.12136E=00 0.60051E=02 0.17014E=03 0.35943E=05 0.57639E=07	-0,20858E-00 -0,14617E-02 0,1397E-02 0,12918E-03 0,3127E-05 0,53705E-07		
		:	x = 1.35			x :	= 1.85			
2 3 4 5 6	0.49573E-00 0.41872E-01 0.12730E-02 0.31791E-04 0.55135E-06 0.66000E+08	0.24505E-00 0.78307E-01 0.37758E-02 0.10399E-03 0.18743E-03 0.23634E-07	0.27933E-00 0.16758E-01 0.46350E-03 0.68788E-05 0.88627E-07 0.77383E-09	-0.31618E+01 0.91541E+02 0.32859E+03 0.63332E+05 0.92603E+07 0.73898E+09	0.56106E 00 0.27355E-00 0.13730E-01 0.53301E-03 0.16668E-04 0.38283E-06	0.15691E-00 0.20457E-00 0.30860E-01 0.15991E-02 0.54326E-04 0.12964E-05	0.47002E-00 0.13928E-00 0.17775E-02 0.23143E-03 0.51404E-05 0.86468E-07	-0.21958E-00 -0.11180E-01 0.38296E-02 0.17054E-03 0.83908E-05 0.80038E-07		
		x	= 1.40			x =	1.90			
1 2 3 4 5	3.51812E 00 0.52597E-01 0.16550E-02 0.43950E-04 0.81745E-06 0.10848E-07	0.23647E-00 0.91894E-01 0.48289E-02 0.14294E-03 0.27722E-05 0.37619E-07	0.30607E-00 0.22049E-01 0.56576E-03 0.10301E-04 0.13566F-06 0.13365E-06	-0.54043E-01 0.10992E-01 0.44670E-03 0.493316E-05 0.13123E-06 0.13447E-08	0.55227E 00 0.30791E-00 0.17063E-01 0.67892E-03 0.22229E-04 0.53761E-05	0:15041E-00 0:20517E-00 0:36637E-01 0:20094E-02 0:72021E-04 0:18143E-05	0.48523E-00 0.15708E-30 0.49951E-02 0.31235E-03 0.72832E-05 0.12949E-05	-0.22850E-00 -0.23182E-01 0.44968E-01 0.22288E-03 0.51112E-05 0.11793E-05		
		x :	- 1.45			x =	1.95			
	0.53655E 00 0.55659E-01 10.21351E-02 0.60079E+04 0.11955E+05 0.17016E-07	0 • 22 712 E - 00 0 • 10 54 7E - 00 0 • 51 192 E - 02 0 • 194 18 E - 03 0 • 494 19 E - 05 2 • 588 79 E - 97	0+23029E-00 0+28662E-01 0+78467E-03 0+15214E-04 0+21453E-06 0+22635E-08	-0:77352E-01 0:12583E-01 0:59895E-03 0:13547E-04 0:20497E-06 0:22526E-08	0.5%17%E 00 0.3413%E-00 0.21153E-01 0.66029E-03 0.29%28E-04 0.7%817E-06	0.1%445E-00 0.20209E-00 0.43209E-01 0.43209E-02 0.94730E-04 0.25156E-05	0.50040E 00 0.17428E-00 0.12743E-01 0.41846E-03 0.10228E-04 0.19105E-06	-0.23537E-UC -0.17163C-U1 0.51652E-U2 0.251652E-U2 0.2838E-U3 0.84219E-U5 0.17169E-U6		
			1.50			x = 2.00				
	0.55112E 00 0.81394E-01 0.27397E-02 2.81276E-04 0.17273E-05 0.26275E-07	0.21739E-00 0.12186E-00 0.76878E-02 0.26093E-03 C.58146E-05 D.90718E-07	0.35215E-00 0.36852E-01 0.10773E-02 0.22184E-04 0.33360E-06 0.37604E-09	-0.10056E-00 0.14023E-01 0.79239E-03 0.19393E-04 0.11509E-06 0.37716E-08	0.52777E 00 0.37287E-00 0.28159E-01 0.10849E-02 0.3688E-04 0.10323E-05	G.13895E-00 G.19571E-00 G.50611E-01 G.31131E-02 G.12367E-03 G.36574E-05	0.51554E 00 0.19057E-00 0.16111E-01 0.55666E-03 0.14243E-04 0.27911E-06	-0.24022E-00 -0.52894E-01 0.57790E-02 0.369942E-03 0.11895E-06 0.24791E-06		

						x = 0.05										
		n	$\mathbf{a}_n^{\mathbf{r}}$		a i		ρ <mark>ι</mark>	b	ı		٩F			x = 0.55		
		2 0.217	567E-04 24 IE-08 311E-12	0.65	9969E-04 015E-08 271E-12	-0.444	" 643E-07 643E-09 79E-10	0-17626	E-07 E-10	0.409 0.295 0.130	2F R 743E-0 27E-0 73E-0 72E-0 96E-10 67E-12	0.10 0.83 0.38	21 398E-01 568E-02 218E-07 091E-07 057E-09 747E-12		12E-04 10E-06 1E-09 5E-11	0.22743E-02 0.21330E-04 0.10477E-06 0.34578E-09 0.35828E-11
	,				x	: ≥ 0.10							,	c = 0.60		
	1 2 3 4 5 6	0.1692 0.1692	SE-07	0.213	18E-03 67E-06 40E-10	0.7143 0.5477 0.6016	35-09	0.57144E 0.10926E 0.44440E	-0 9	0.5019 0.6303 0.53630 0.2855 0.9479 ~0.2695	5E-03 9E-05 0E-07 4E-10	0.162	51E-01 77E-02 36E-04 96E-07	0.62021 0.59016 0.32899 0.11906 -0.50509	E-04 E-06 E-08 E-12	0.33719E-02 0.38649E-04 0.22708E-06 0.86655E-09 0.76874E-11 0.33337E-12
	1	0.56498	E-01	0-1994		0.15							x	= 0.65		
	2 3 4 5 6	0.59877 0.34981	E-05	0+1929 0+1620 0+9759	7E-05	0.54381 0.36344 0.31461	E-08	0.39514E= 0.25072E= 0.17591E=	5 é	0 - 67092 0 - 94827 0 - 94990 0 - 585278 P - 233778 - 0 - 103248	F-03 E-05 E-07	0.9973 0.24208 0.26555 0.16982 0.70708 0.21486	E-02 E-04 E-06 F-09	0.94237E 0.10404E 0.67747E 0.28676E 0.56044E -0.13305E	-03 -06 -08	G-47645E-02 G-6656DE-04 G-46184E-06 G-20327E-08 G-12246E-10 G-54257E-12
	1	0.13611E	-0.3		x =	0.20							х :	0.70		
	2 3 4 5 6	0.25246E 0.25150E 0.85462E -0.92196E -0.25482E	-05 -06 -11	0.31470 0.68219 0.711778 0.993388 0.317518 0.865558	E-05 E-08 E-11 E-14 E-18	0.23040E 0.26028E 0.18812E 0.44626E 0.59100E -0.41964E	-07 -10 -11 -12	0.16552E-0. 0.18926E-0. -0.15379E-10 -0.39288E-11 0.18481E-11 0.30928E-15	3	0.87857E 0.13872E- 0.15940E- 0.11372E- 0.53252E- 0.92658E-	-02 -04 -06	04120118 04349338 0443988 04329278 04154788 04554398	-02 -04 -06 -08	0.13896E- 0.17604E- 0.13229E- 0.64745E- 0.21138E-1	03 05 08 .0	0.64385E-02 0.10970E-03 0.88964E-06 0.49124E-08 0.22101E-10 0.89576E-12
	1	0.27167E-0		0.615386-	x = 0. .o.>	25 0.70816E-0							x = 0	.75		
	2 3 4 5 6	0.77137E-0 0.11949E-0 0.74370E-1 -0.10615E-1 -0.46259E-1	17 1 3	0.20790E- 0.34013E- 0.38860E- 0.38534E- 0.15709E-	04 07 10	0.12505E-0 0.16561E-0 0.65319E-1 0.14567E-1	9 2 1	0.50091E-04 0.90297E-07 0.12985E-09 0.12518E-10 0.15148E-11 0.35143E-14		0.11274E-0 0.19619E-0 0.25610E-0 0.21094E-0 0.11374E-0 0.16023E-1	2 4 6	0.14101E- 0.49125E- 0.71585E- 0.60950E- 0.33691E- 0.12963E-	02 04 06 08	0.19946E-0: 0.20755E-0: 0.24673E-05 0.13800E-07 0.52417E-10		0.83161E-02 0.17397E-03 0.16349E-05 0.95170E-08 0.41238E-10 0.40047E-12
	1				x = 0.3	0							x = 0.8	o		
	2 3 4	0.48251E-02 0.19223E-04 0.42759E-07 0.51540E-10 -0.74926E-13 -0.49374E-16	0	0:10638E-0 0:51648E-0 0:12154E-0 0:17975E-0 0:26477E-12	6 0 9 -0	0.17777E-03 0.44483E-06 0.71955E-09 0.42223E-11 0.25512E-11 0.35176E-13	0	0=12330E-09 0=32163E-06 0=52309E-09 0=22144E-10 0=35013E-11 0=24744E-19		0.14177E-00 0.2175 E-02 0.40515E-04 0.37589E-06 0.23060E-08 0.94194E-11	•	0.16152E-0 0.67533E-0 0.11104E-0 0.10836E-0: 0.68156E-0: 0.29630E-10	2	0-27927E-01 0-45553E-03 0-4423E-05 0-4423E-07 0-28063E-07 0-12123E-09	0 0 0	- 10233E-01 - 26656E-03 - 28827E-05 - 19159E-07 - 90010E-10
1		0.79198E-02			x = 0.35								x = 0.85			
3 4 5 5	-0	0.41631E-04 0.12559E-06 0.22642E-09 0.15245E-12 0.36490E-15	0. 0.	.16877E-01 .11142E-03 .35649E-G6 .66785E-09 .12409E-11 .12375E-14	0. 0. -0. -0.	38#28E-03 13235E-05 25998E-08 87273E-11 48638E-12 13203E-12	0. G.	026286E-03 09974E-06 018533E-08 011252E-10 21911E-11 93564E-13		0:17465E=00 0:38197E=02 0:61897E=04 0:64656E=06 0:44760E=08 0:20519E=10	0	0=18062E-00 1=90996E-02 =16996E-03 =18591E-05 =13210E-07 =65374E-10	0	• 38180E-01 • 70247E-03 • 76536E-05 • 54711E-07 • 26961E-09 • 88151E-12	0. 0. 0.	11912E-01 39591E-03 49004E-05 38997E-07 19834E-09
1	0.	122866-01			= 0.40							х	= 0.90			
3 4 5 6	0.	81386E-04 31941E-06 74601E-09 54741E-13 20869E-14	0 • 2 0 • 9 0 • 2 0 • 4	75119E-01 1679E-03 0522E-06 1970E-08 0061E-11 9511E-14	0.3 0.8 0.1 -0.1	6632E-03 3830E-05 5355E-08 6041E-10 3713E-11	0 • 2 0 • 6 0 • 2 0 • 1	0370E-03 3737E-05 1266E-08 0949E-10 7531E-11 3973E-13	0	-21075E-00 -91804E-02 -92323E-04 -10779E-05 -83675E-08	0 e 0 e 0 e	19735E-00 12045E-01 25201E-03 30908E-05 24633E-07 13746E-09	0. 0. 0.	50989E-01 10577E-02 12041E-04 10265E-06 17545E-09 17774E-12	0 = 5 0 = 61 0 = 40	2954E-01 7139E-03 0624E-05 1842E-07 1991E-09
,	0				0.45							х -	0.95			
1	0.1 0.2 0.3	18266E-01 4716E-03 2764E-06 1579E-08 1455E-11 6231E-14	0.38 0.20 0.630 0.133	560E-01 965E-03 579E-05 077E-08 280E-10 512E-13	0+77- 0+24- 0+526 -0+847 -0+360	002E-02 091E-05 002E-07 049E-10 11E-12 48E-12	0.53 0.17 0.44 0.18	#19E-03 618E-05 489E-07 ##2E-10 093E-11 176E-12	0. 0.	24911E-00 69371E-02 13481E-03 1747E-05 15103E-07 90602E-10	0 a 1 0 a 3 0 a 4	1093E-00 568ZE-01 6557E-03 9956E-05 9579E-07 7521E-09	0 - 15 0 - 26 0 - 15	5513E-01 5590E-02 7558E-04 610E-06 551E-08 727E-11	0.80: 0.128 0.122 0.801	940E-01 272E-03 80E-04 959E-06 36E-09 06E-11
	0.26	265E-01	0.481	x = 4 18E-01		115						x =	00.1			
	0.15 0.55 0.12	029E-03 194E-05 610E-08 897E-10 057E-13	0.658	12E-03 71E-05 12E-07 77E-10	0.1620 0.1620 0.6352 0.1611 -0.1921 -0.1719	1E+07 0E+09 8E-11	0.110	15-11	0.91 0.19 0.27 0.26	8864E-00 1879E-02 1316E-03 1429E-05 446E-07 773E-09	0.20 0.51 0.78 0.775	085E-00 120E-01 193E-03 129E-05 34E-07 65E-09	0.225 0.335 0.327 0.225	7E-08	0.1098 0.1098 0.2003 0.2125 0.1533	96-02 36-04 86-06 76-08

			x = 1.05				x = 1.55					
	n ar	, a	i br	$b_{\mathbf{n}}^{i}$	a _n r	a_n^i	b <mark>r</mark>	$b_{\mathbf{n}}^{\mathbf{i}}$				
	1 0.32804 2 0.12053 3 0.27210 4 0.42711 5 0.45029 6 0.333638	E-01 0.25461 E-03 0.72647 E-05 0.12127 E-07 0.13175	E-01 0.32020E E-03 0.51949E E-04 0.56019E E-06 0.42498E	-02 0.14667E-0. -04 0.30404E-0. -06 0.35849E-0. -08 0.26614E-0.	0.120686 0.445186 0.137916 0.310416	-00 0.13522E -02 0.10159E -03 0.36966E -05 0.88025E	-00 0.479758 -01 0.179218 -03 0.403808 -05 0.664988	-01 0-15944E-02 -02 0-66998E-03 -04 0-20946E-04 -06 0-38741F-06				
			x = 1.10			x = 1.60						
	1 0.36615E 2 0.15677E 3 0.37752E 4 0.64688E 5 0.74776E 6 0.61063E	-01 0.318078 -03 0.998748 -05 0.182968 -07 0.218268	F-01 0.44753E- F-03 0.79254E- F-04 0.93499E- F-06 0.77753E-	-02	0.56296E 0.1m4912- 0.56462E- 0.18322E- 0.43776E- 0.75433E-	00 0.14811E- 02 0.12543E- 03 0.44868E- 05 0.12359E-	03 0.58079E- 01 0.23885E- 03 0.57988E- 04 0.10033E-	01 -0.12592E-02 02 0.82554E-03 04 0.28759E-04 05 0.57354E-06				
			x = 1.15				x = 1.65					
	1 0.40195E- 2 0.20237E- 3 0.51666E- 4 0.96170E- 5 0.12135E- 6 0.10828E-6	01 0.39267E- 03 0.13529E- 05 0.27086E- 06 0.35338E-	-01 0.61607E-6 -02 0.11870E-0 -04 0.15256E-0 -06 0.13849E-0	0.24171E-02 0.65533E-04 0.65533E-04 0.694481E-06 0.90918E-08	0.56422E 0.17100E= 0.71266E= 0.20141E= 0.61001E= 0.11102E=	00 0:15949E-0 02 0:15369E-0 03 0:63530E-0 05 0:17169E-0	0 0.69218E- 1 0.31522E- 3 0.81363E- 4 0.14947E-0	01 -0.53485E-02 02 0.99581E-03 04 0.38955E-04 05 0.83730E-06				
			x = 1.20				x = 1.70					
3 6	0.25940E-0 0.69841E-0 0.14056E-0	1 0.47852E-0 3 0.18081E-0 4 0.39409E-0 6 0.56016E-0	01 0.83594E-0; 02 0.17478E-0; 04 0.24380E-0; 06 0.24060E-0;	2	0.56300E 00 0.20116E-00 0.49559E-02 0.31563E-03 0.88349E-05 0.16376E-06	0 0.16875E-00 2 0.18692E-01 3 0.82230E-03 5 0.23587E-04	0.61209E-0 0.41204E-0 0.11299E-0 0.21999E-0	1 -0.10894E-01 2 0.11728E-02 3 0.52069E-04 5 0.12062E-05				
			x = 1.25				x = 1.75					
3 4 5 6	0.46363E-00 0.33030E-01 0.93362E-03 0.2022EE-04 0.30067E-06 0.31730E-08	0.57659E-0 0.23864E-0 0.56432E-0 0.87087E-0	0.11186E-01 0.25338E-03 0.38223E-05 0.40865E-07	-0.38502E-01 0.35467E-02 0.13045E-03 0.22746E-05 0.26147E-07 0.21608E-09	0.59956E 00 0.23211E-00 0.11210E-01 0.40972E-03 0.11544E-04 0.23713E-06	0.16461E-00 0.17535E-00 0.22569E-01 0.10550E-02 0.32099E-04 0.68442E-06	0-39142E-00 0-93623E-01 0-53351E-02 0-15539E-03 0-32013E-05 0-49244E-07	-0.21159E-00 -0.17966E-01 0.13433E-02 0.68693E-04 0.17157E-05 0.29040E-07				
			x = 1.30			х	= 1.80					
1 2 3 4 5 6	0.48913E-00 0.41782E-01 0.12356E-02 0.28696E-04 0.46052E-06 0.52558E-08	0.21359E-00 0.68659E-01 0.31137E-02 0.79627E-04 0.13299E-05 0.15532E-07	0-14765E-Q1	-0.56012E-01 0.40710E-02 0.17901E-03 0.34587E-05 0.42911E-07 0.38551E-09	0.55413E 00 0.26393E-00 0.13979E-01 0.52832E-03 0.15655E-04 0.33972E-06	0.16249E-00 0.17899E-00 0.27054E-01 0.13437E-02 0.43279E-04 0.97718E-06	0.40871E-00 0.10681E-00 0.668452E-02 0.21172E-03 0.46089E-05 0.74927E-07	-0.22259E-00 -0.24543E-01 0.14870E-02 0.89450E-04 0.24109E-05 0.43559E-07				
		x	= 1.35			х •	1.85					
1 2 3	0.51048E 00 0.52495E-01 0.14204E-02 0.440177E-04 0.469387E-06 0.85384E-08	0.20670E-00 0.80760E-01 0.40193E-02 0.11084E-03 0.19974E-05 0.25176E-07	0.24477E-00 0.19226E-01 0.51048E-03 0.89241E-05 0.11091E-06 0.10196E-08	-0.74679E-01 0.44561E-02 0.24139E-03 0.50957E-05 0.609015E-07 0.66616E-09	0-54695E 00 0-29571E-00 0-17371E-01 0-67706E-03 0-21054E-04 0-48181E-06	0.16088E-00 0.17956E-00 0.32199E-01 0.16990E-02 0.57843E-04 0.13808E-05	0.42615E-00 0.11993E-00 0.86949E-02 0.29592E-03 0.65681E-05 0.11269E-06	-0.29199E-00 -0.36606E-01 0.1979E-02 0.11496E-03 0.394@1E-05 0.66434E-07				
			= 1.40			x =	1.90					
1 2 3 4 5 6	0.52796E 00 0.65478E-01 0.21079E-02 0.55571E-04 0.10297E-05 0.13621E-07	0.19927E-00 0.93805E-01 0.51368E-02 0.15234E-03 0.29542E-05 0.40077E-07	0.26562E-00 0.26692E-01 0.71077E-03 0.13322E-00 0.17775E-06 0.17571E-06	-0.94411E-01 0.45826E-02 0.31991E-03 0.74254E-05 0.10896E-06 0.11398E-08	0.53825E 00 0.32659E-00 0.21512E-01 0.86270E-03 0.28091E-04 0.67682E-06	0.15962E-00 0.17724E-00 0.38045E-01 0.21595E-02 0.76668E-04 0.19323E-05	0.44376E-00 0.13299E-00 0.10942E-01 0.38286E-03 0.92702E-05 0.16762E-06	-0.23955E-00 -0.47933E-01 0.15489E-02 0.48977E-03 0.45969E-05 0.98178E-07				
		х =	1.45			x = 1	95					
1 2 3 4 5 5 5 6 6 6 6 6 6 6 6 6 6 6 6 6 6 6 6	0.50171E 00 0.81025E-01 0.27218E-02 0.76005E-04 0.15067E-05 0.21371E-07	0.19228E-00 0.10751E-00 0.65039E-02 0.20692E-03 0.43071E-05 0.62728E-07	0 = 28536E-00 0 = 31268E-01 0 = 97805E-03 0 = 19610E-04 0 = 28016E-06 0 = 29729E-08	-0.11401E-00 0.42819E-02 0.41869E-03 0.10649E-04 0.16690IE-06 0.19017E-08	0.52824E 00 0.34579E-00 0.26546E-01 0.10934E-02 0.37203E-04 0.94219E-06	0:19857E-00 0:17242E-00 0:44409E-01 0:26615E-02 0:10082E-03 0:26789E-05	0.46134E-00 0.14584E-00 0.13637E-01 0.50849E-03 0.12964E-04 0.24669E-06	-0.24537E-00 -0.40331E-01 0.14174E-02 0.10230E-03 0.62423E-05 0.13611E-06				
		x = 1	1.50			r = 2.	00					
	0.55197E 00 0.99377E-01 0.34913E-02 0.10288E-03 0.217605-05 0.33008E-07	0.18581E-00 0.12149E-00 0.81625E-02 0.27890E-03 0.61961E-05 0.96650E-07	0.30415E-00 0.39023E-01 0.13309E-02 0.28489E-04 0.43479E-06 0.49377E-08	-0.13316E-00 0.33596E-02 0.53324E-03 0.15044E-04 0.25787E-06 0.31205E-08	0.51713E 00 0.38272E-00 0.32640E-01 0.13791E-02 0.48924E-04 0.13004E-05	0.15757E-00 0.16567E-00 0.51900E-01 0.32992E-02 0.13158E-03 0.36815E-05	0.47944E-00 0.15840E-00 0.16825E-01 0.67002E-03 0.17973E-04 0.35949E-05	-0.24940E-00 -0.7955E-01 0.1057E-02 0.2246E-03 0.63855E-05 0.19454E-05				

						x = 0.05											
		a	a_n^r		\mathbf{a}_{n}^{i}		br n	ь	<u>.</u>		a.r			x = 0.55			
		2 0.27	*41 <i>E</i> - 04 957 E - 08 786E - 12	0.59	?869E~0 !869E~0 !859E~!	d -0.19	5915-07 6762-09 359F-11	0•14208 -0•15060 -0•50688	€-07 E-09	0.475 0.343 0.151; 0.4173	27 19E-01 41E-05 61E-05 24E-07	0.115 0.893 0.408 0.120	an 1888-01 1778-02 1888-05 1878-07 1888-09	0.477 0.394; 0.1865 0.5602	br n 47712E-02 0.1 99416E-04 0.1 8637E-06 0.8 6021E-09 0.2 8984E-11 0.1		-01
					x =					-0.1745	116-12	0 • 25 3	24E-12	-0.4988	2E-12	0.14982E- -0.38239E-	12
		1 0.196/ 2 n.9205 3 0.2333 4 6	3E-07	0.229	47E-03 77E-06 63E-10	0 • 314 0 • 839	16E-06 51E-09 58E-11	0.44099E 0.11948E 0.88644F	-09	0-6003 0-7401 0-6912 0-39021 0-11185 -0-12255	7E-03 7E-05 FE-07 SE-09	0+4533 0+1752 0+1636 0+4903 0+3125(0+2761)	0E-01 9E-02 7E-04 4E-07	= 0.60 0 - 74321 0 - 72719 0 - 40798 0 - 14612 -0 - 31794 -0 - 41619	E-04 E-06 E-08	0.22751E=0 0.30289E=0 0.18454E=0 0.70549E=0 0.16489E=1 -0.60018E=1	9
	:	0.67227	'E-03	0.1431		= 0.15							х -	0.65			
	3 4 5 6	0.49920 0.38340	E-06	0.1743 0.1035	4E-05	0.6786 0.4416 0.4478	E-08	0:33197E= 0:21853E= 0:16156E=	o e	0.600248 0.111478 0.11046E 0.67722E 0.27207E 0.72679E	-02 -04 -07 -09	0+10470 0+26062 0+26533 0+18218 0+752698 0+150926	E-02 E-04 E-06 [-09	0:11164E 0:12779E 0:83895E 0:33949E 0:63394E -0:35860E	-03 -06 -08 -11	0.25986E-02 0.91279E-04 0.97230E-04 0.1697E-08 0.49739E-11 -0.71671E-12	
	,	0.15/125	- 22			= 0.20						x = 0.70					
	3 4 5 6	7.294.95E 7.29110F 0.17902E -0.91177E- -0.25179E-	-07 -04 -1.	0.33973; 0.73390; 0.76818; 0.518533 0.794118 0.903518	E-05 -08 -11 -14 -18	0.237176 0.327208 0.402976 =0.318598 =0.457036 =0.414815	-07 -10 -12 -12	0.13802E-0. 0.15976E-0. 0.267515-10. 0.21102E-11. 0.2010E-12.	7	0:104326= 0:16326E= 0:18564E= 0:13161E= 0:61346E=(0:16374E=)	02 04 06 09	0+12485E- 0+37611E- 0+47702E- 0+35329E- 0+16970E- 0+49582E-	-00 -02 -04 -06 -08	0.16292E- 0.21944E- 0.16396E- 0.00071E- 0.21747E-1	13 19 8 0	0.36891E-02 0.82849E-04 0.71091E-06 0.38949E-08 0.14979E-10	
	:	C+32400E=(12	0.6.7.123	x =								x = 0.1	75			
	3 4 5 5	0.901416+0 0.13840E+0 0.910746+1 -0.10700E+1 -0.45712E+1) 5) 7 1 3	0.504126- 0.223685- 0.30500F- 0.361696- 0.336256- 0.145538-	07 19 13	0.001045- 0.100805- 0.160205- 0.160205- 0.299825- -0.104125- -0.464046-1	06 99 11 . t	0.41367E-04 0.756H6E-07 0.10339E-09 0.51424E-11 0.16351E-12 0.22651E-14		0.13304E-0 0.23354E-0 0.30049E-0 0.24424E-0 0.13134E-0 0.51147E-1	12 14 16	0.14489E-0.52886E-0.76917E-0.769400E-0.76981E-0.0.13344E-1	02 04 06 18	0.22914E-0 0.35044E-0 0.3644E-2 0.17090E-0 0.62833E-1(•	0.41735E-02 0.412836E-03 0.412820E-05 0.47230E-08 0.30256E-10	
		1 474315. 1			x = 0	30							x = 0.80	,			
		3+57621E-0: 3+27474E-0+ 1+43649E-03 7+39334-10 *3+81514F-E3 ************************************	:	7-14-74E-0 7-59971E-0 7-13946E-0 7-17931E-0 1-73051E-1 -1998E-1	6 7 .	0.420mpt-0 0.55e-re-0 0.80700t-0 -0.389515-1 -0.15350t-11 -0.16174E-1	9 0	0+10054c+03 1+2571+c+05 +475865+09 +67437+11 +505606+13 +127775+13		0-16596E-00 0-3276E-02 0-47188E-04 0-43596F-06 0-26610E-08 0-10921E-10	0	0-16361E-00 0-72689E-02 0-12018E-03 0-11628E-09 0-73066E-08 0-31261E-10		0-31460E-01 0-55244E-03 0-55244E-05 0-54457E-05 0-34740E-07 0-14857E-09 0-85841E-12	0	0+41959E=02 7+19191E=03 8+22516E=05 8+15446E=07 8+75504E=10 8+26679E=12	
1		0.94705E-02			x = 0.3	-						,	c = 0.85				
2 3 4 5 6		0.48697E-04 0.14572E-06 0.29519E-09 0.29679E-12 -0.36283E-19	0. 0.	18183E-0: 11989E-03 30283E-06 71108E-09 96100E-12 11494E-14		0.48047E-03 0.16474E-05 0.32017E-08 0.30314E-12 0.17374E-11 0.16288E-12	0. 0. 0.	21089E-03 77\$17E-06 15\$06E-08 97773E-11 25167E-12 52650E-15		0 - 20244E - 00 0 - 45134E - 02 0 - 72135E - 04 0 - 74910E - 08 0 - 51712E - 08 0 - 24969E - 10	0 • 0 •	18009E-00 97908E-02 18265E-03 19991E-05 14158E-07 69865E-10	0. 0. 0.	42059E-01 84710E-05 94016E-05 67642E-07 35319E-09 19863E-12	0 • 0 • 0 •	34170E-02 27503E-03 37775E-05 27536E-07 15674E-03 21516E-12	
1		0.10708E-01	9.2	7019E-01	= 0.40							x	• 0.90				
2 3 4 5 6	-	0.95250E-04 0.37078E-06 0.85736E-09 0.85736E-12 0.20785E-14	0.2 0.9 0.2 0.3	3327E-03 7215E-06 3542E-08 5863E-11 5794E-14	0. 0. -0.	994460E=03 92091E=05 10625E=07 11008E=10 27927E=11 40945E=12	0 - 1 · 0 - 5 i 0 - 2 i	99926-03 95016-05 13986-08 14866-10 14866-12 3526-12	0	>24153E-00 •61900E-02 •10767E-03 •12492E-09 •96619E-08 •52239E-10	0+3	9341E-00 2951E-81 7084E-03 3178E-05 6405E-07	0 • 1 0 • 1 0 • 1 0 • 7	4790E-01 2672E-02 3791E-04 2672E-06 287E-09 1678E-11	0 - 5 : 0 - 5 : 0 - 5 :	4340E-82 B173E-03 1239E-05 1350E-07 1547E-09	
1		•21886E-01			0.45							X =	0.95				
3	0.	-17294E-03 -84491E+06 -24887E-08 -30597E-11 10147E-13	0+419 0+221 0+674 0+130	158E-01 939E-03 102E-05 158E-08 176E-10 70E-15	0.9 0.3 0.5 -0.2	7175E-02 6158E-05 0622E-07 8107E-10 5576E-11	0.496 0.1494 0.494	130E-03 146E-05 32E-07 01E-10 57E-12 74E-12	0.2	28190E-00 82203E-02 15733E-03 0240E-05 7450E-07 0583E-09	0 - 164 0 - 392 0 - 534 0 - 475	300E-00 951E-01 190E-03 14E-05 80E-07 58E-09	0 • 145 0 • 255 0 • 229 0 • 1 • 2	807E-01 837E-02 897E-04 42E-04 49E-08 89E-11	0.962 0.964 0.642	748E-02 191E-03 19E-05 19E-07 05E-09 99E-11	
		31481E-01		715-01		358E-02						x = 1.	.00				
	0.6	29335E-03 17649E-05 14278E-08 14369E-10 16064E-13	0.7084 0.4604 0.1739 0.4222 0.8757	3E-03 7E-05 1E-07 6E-10	0.20 0.78 0.19 -0.92	338E-02 163E-04 190E-07 368E-09 358E-12 139E-12	0+1089 0+8928 0+3491 0+1008 0+4228 =0+1685	1E-07 8E-09 7E-12	0 • 10 0 • 21 0 • 30	2242E-00 1902E-01 1960E-03 1047E-05 565E-07 404E-09	0.2086 0.2155 0.5588 0.8450 0.8313	14E-03 10E-05 18E-07	0.862 0.2691 0.4061 0.4028 0.2789 0.1341	70E-02 6E-04 4E-06 1E-08	0.779! 0.6601 0.1456 0.1695 0.1224	72E-03 17E-04 17E-06 8E-08	

				x = 1.05				x = 1.55	
	п	a _n	a	i br	b _n	a ^r	a _n i	b,	, bi
	1 2 3 4 5	0.36156E 0.14317E 0.31807E 0.49560E 0.52053E- 0.38666E-	-01 0.27286 -03 0.78881 -05 0.130170 -07 0.141271	E-01 0.373705 E-03 0.629975 E-04 0.687996 E-06 0.524345	-02 0.819812-0 -04 0.21823E-0 -06 0.27624E-0 -08 0.22719E-0	0-199146 0-525566 0-160966 0-360106	-00 0.13538E -02 0.10876E -09 0.39667E -05 0.94413E	-00 0.466721 -01 0.202596 -03 0.485538 -05 0.805686	E-00 -0.17375E-00 E-01 -0.89412E-02 E-02 0.27479E-03 E-04 0.13279E-04 E-06 0.27702E-06
				× = 1.10				x = 1.60	
	2 3 4 5	0.39829E- 0.18639E- 0.44170E- 0.75097E- 0.86462E- 0.70536E-	01 0.34017E 05 0.10734E 05 0.19639E 07 0.23407E	-01 0.51641E- -02 0.95706E- -04 0.11459E- -06 0.95876E-	02 0.97066E-03 04 0.31608E-04 05 0.44847E-06 08 0.40765E-08	0.16534E- 0.46705E- 0.21400E- 0.50607E-	-00 0:14442E- -02 0:13411E- -03 0:52250E- -05 0:13256E-	00 0.55251E- 01 0.26692E- 03 0.69514E- 04 0.12121E-	-01 -0.13405E-01 -02 0.28772E-03 -04 0.17615E-04 -05 0.40348E-06
				x = 1.15				x = 1.65	
	2 0 3 0 4 0 5 0	0.43182E-0 0.24073E-0 0.60506E-0 0.1170E-0 0.14036E-0 0.12516E-0	1 0.41869E- 2 0.14540E- 4 0.29075E- 6 0.37897E-	01	02 0.10834E-02 02 0.44671E-04 05 0.71009E-05 07 0.71192E-08	0.567026 0.19306- 0.84235E- 0.262175- 0.70922E- 0.12934E-	00 0.15532E-0 02 0.16407E-0 03 0.48147E-0 05 0.18407E-0	00 0.64475E= 01 0.34765E= 03 0.95589E= 0.18000E=	01 -0.18985E-01 02 0.27331E-05 04 0.22941E-04 05 0.57859E-06
				x = 1.20				x = 1.70	
	2 0. 3 0. 4 0. 5 0.	46161E-00 30856E-01 81871E-03 16334E-04 22315E-06 21667E-08	0+50867E-0 0+19429E+0 0+42304E-0	11	0.11127E-02 0.61499E-04 0.10984E-05 0.12111E-07	0.56479E 0 0.22391E-0 0.10588E-0 0.36919E-0 0.98009E-0 0.18949E-0	0 0:14181E-00 1 0:19917E-01 9 0:88190E-03 5 0:25295E-04	0.74215E-0 0.4485ZE-0 0.19193E-0 0.24405E-0	1 -0.25727E-01 2 0.21498E-03 9 0.29508E-04 5 0.81739E-06
				x = 1.25				x = 1.75	
3	0 ± 3 0 ± 2 0 ± 3	18743E-00 19260E-01 10956E-02 13519E-04 4800E-05 6654E-08	0.19444E-00 0.61025E-01 0.25439E-02 0.60578E-04 0.93402E-06 0.10071E-07	0+12355E-01 0+30123E-03 0+4474E-05 0+50188E-07	-0.69098E-01 0.99476E-03 0.83949E-04 0.16621E-05 0.20122E-07 0.16947E-09	0.56063E 00 0.25492E-00 0.13231E-01 0.47959E-03 0.13417E-04 0.27448E-06		0-36073E=00 0-84339E=01 0-57226E=02 0-18024E=03 0-38207E=05 0-59778E=07	-0.39627E-01 0.90732E-04 0.36681E-04
			,	c = 1.30			x	= 1.80	
1 2 3 4 5	0.49 0.14 0.39 0.59	1926E 00 1560E-01 1514E-02 1384E-04 1917E-06 717E-08	0:18830E-00 0:72267E-01 0:33443E-02 0:85476E-64 0:14264E-05 0:16647E-07	0+20089E-00 0+16017E-01 0+42768E-03 0+71972E-05 0+83317E-07 0+71636E-09	-0.84904E-01 0.64431E-03 0.10921E-03 0.24439E-09 0.32698E-07 0.29877E-09	0.55476E D0 0.28600E-00 0.18514E-01 0.61884E-03 0.18204E-04 0.39336E-06	0-14712E-00 0-16465E-00 0-28662E-01 0-14400E-02 0-46406E-04 0-10475E-05	0-37936E-00 0-94723E-01 0-72255E-02 0-24381E-03 0-54909E-05 0-90735E-07	-0.23881E-00 -0.42631E-01 -0.12774E-03 0.44887E-04 0.13615E-05 0.31091E-07
			x	= 1.3 5			x :	1.85	
1 2 3 4 5	0+621 0+190 0+467 0+803	719E 00 118E-01 156E-02 16BE-04 65E-05 39E-08	0 - 18223E-00 0 - 80 438E-01 0 - 43155E-02 0 - 11898E-05 0 - 21026E-05 0 - 26985E-07	0.21926E-00 0.20490E-01 0.59872E-03 0.10780E-04 0.13567E-06 0.12589E-08	-0.10413E-00 -0.44069E-04 0.14002E-03 0.39810E-03 0.52039E-07 0.51794E-09	0.54733E 00 0.31633E-00 0.20498E-01 0.7939E-03 0.24495E-04 0.55808E-06	0+14408E-00 0+16912E-00 0+34024E-01 0+18199E-02 0+2019E-04 0+14601E-09	0.39437E-00 0.10526E-00 0.90267E-02 0.9267E-03 0.77927E-05 0.15611E-06	-0.24757E-00 -0.5264E-01 -0.67547E-03 0.53585E-04 0.21124E-05 0.65581E-07
			х =	: 1.40			х •	1.90	
1 2 3 4 5	0.771	44E 00 63E-01 14E-02 18E-04 51E-05 57E-07	0-17441E-00 0-97280E-01 0-59127E-02 0-16952E-03 0-31467E-05 0-42960E-07	0.23721E-00 0.23719E-01 0.82708E-03 0.16034E-04 0.21698E-06 0.21456E-08	-0.12215E-00 -0.12045E-03 0.1746BE-03 0.51045E-05 0.61207E-07 0.88214E-09	0.5385AE 00 0.34527E+00 0.25337E-01 0.10118E-02 0.32699E-04 0.78429E-06	0-17074E-00 0-16139E-00 0-40041E-01 0-22840E-02 0-821845-04 0-20712E-05	0.41774E-00 0.11589E-00 0.11165E-01 0.433B4E-03 0.10950E-04 0.20193E-06	-0.25474E-00 -0.63544E-01 -0.99468E-03 0.41919E-04 0.28171E-05 0.65272E-07
			X =	1.45			x = 1	1.95	
1 2 3 4 5	0.5529 0.9494 0.92074 0.8858 0.17464 0.24696	9E-01 AE-02 7E-04 AE-05	0.17228E-00 0.11042E-00 0.69754E-02 0.22210E-03 0.46198E-05 0.67241E-07	0.25481E-00 0.51847E-01 0.11282E-02 0.23906E-04 0.34119E-06 0.36552E-08	-0.139996-00 -0.297446-02 0.211916-09 0.714986-05 0.124976-06 0.149916-06	0-52854E 00 0-37223E-00 0-31187E-01 0-12831E-02 0-43327E-04 0-10922E-05	0-17258E-00 0-15597E-00 0-15597E-01 0-2847E-02 0-10805E-03 0-28713E-05	0.03740E-00 0.12557E-00 0.13565E-01 0.5710ZE-03 0.15261E-04 0.29636E-06	-0.26027E-00 -0.75194E-01 -0.17945E-02 0.69061E-04 0.37029E-05 0.92773E-07
			x = 1.				x = 2.	00	
	0.56003 0.11561 0.41182 0.11999 0.25231 0.36148	E-00 E-02 E-03 E-05	0-16877E-00 0-1233E-00 0-87473E-02 0-29835E-03 0-66458E-05 0-10361E-06	0+27221E-00 0-38845E-01 0+15206E-02 0+33998E-04 0+52818E-06 0+60527E-08	-0.15724E-00 -0.55055E-02 0.24452E-03 0.24652E-03 0.16719E-06 0.25786E-08	0-51749E 00 0-39688E-00 0-38213E-01 0-16191E-02 0-57008E-04 0-19079E-05	0-17414E-00 0-14944E-00 0-54085E-01 0-35248E-02 0-14099E-03 0-39456E-05	0-45741E-00 0-13728E-00 0-16591E-01 0-74518E-03 0-2102E-06 0-43054E-06	-0.26412E-00 -0.67435E-01 -0.27494E-02 0.73159E-04 0.47959E-05 0.13026E-06

						x = 0.05											
	r	1	a _n r		a i		br n	ьi			ar n		a i	x = 0.55	ьг		
	1 2 3 4 5	0.144	715-04 785-09 345-12	0.21.	118-04 835-08 515-13	0 1 3	925-07	n.3392) -n.39124 -n.405348	E-07 E-10	0.285	76E-01 98E-03 95E-05 92E-08	0 • 148 0 • 113 0 • 511	08E-06 32E-02 32E-04 82E-07 69E-09	0.108 0.502 0.150 0.313	57 22E=01 68E=03 81E=06 76E=08 21E=11 21E=14	0+316+25 0+18420F	-04 -06 -09 -11
	,					- 0.10								K = 0.60			
	2 3 4 5 6	0.1243 0.0353 0.12010	bE-07	0.5763 3.2793 3.7640	±€=05	0+2322 0+6904 0+12302	15-09	0.10585E- 0.32177E- 0.29029E-	-99	0.5815; 0.4501; 0.36252 0.18587 0.63442 0.15348	E-03 E-05 E-07 E-10	0.1297 0.2292 0.2077 0.1116 0.3911 0.95851	2E-02 1E-04 1E-06	0.2014 0.2019 0.1107 0.3942 0.9729 0.14364	4E-03 9E-05 3E-08 3E-11	-0.32406E- 0.51699E- 0.38392E- 0.15336E- 0.40513E- 0.24485E-	06 06 08
					X =	0.15							x	- 0.65		001400324	
	1 2 3 4 5	0.428503 0.407138 0.220298	-Co	0+129216 0+129216	-0:	0+1780cc 0+1131vc 0+42455E	-07	0:/8456g=0 0:714925=0 0:533945=1	5	0+832518 04687198 0+687198 0+381678 0+19276E 0+49366E	-03 -05 -07 -09	0 • 1 6 1 2 6 1 0 • 3 6 2 5 2 6 0 • 3 6 2 5 2 6 0 • 9 3 9 9 7 5 0 • 2 7 0 8 6 6	E-00 E-02 E-04 E-06	0 4 2 9 0 9 9 9 0 4 3 5 6 8 6 0 4 2 2 9 2 8 0 4 9 5 5 0 2 1 0 4 5 7 6 7 5 0 4 5 5 7 6 0 7 5	E=03 E=05 E=0B E=10	-0.30596E-0 0.77543E-0 0.74473E-0 0.35789E-0 0.11195E-1 0.20555E-1	4 6 8 C
	1				x = 0	.20							x =	0.70			
	2 3 4 5 6	0.10491E- 0.17194E- 0.16475E- 0.95391E- -0.22059E-	05 06 12 15	3.46652E- 0.94464E- 0.96718E- 0.58270E- 3.25883E- 0.72404E-	05 08 11	0.76018E- 0.85349E- 0.55099E- 0.11770E- -0.11083E- -0.13466E-	10 12 14	0:31653E-04 0:37973E-07 0:23862E-10 0:7602IE-13 0:95784E-14 0:61265E-16		0:11608E- 0:10234E- 0:10725E- 0:10725E- 0:74309E- 0:34452E- 0:11328E-	02 04 07 09	0.19371E 0.49587E 0.60682E 0.44338E 0.21157E 0.70750E	02 04 06	0.39891E 0.60356E 0.44971E 0.21674E 0.73054E 0.17998E	-03 -05 -07	-0.79966E-02 0.10556E-02 0.1354E-05 0.77907E-08 0.28621E-10 0.69418E-13	
	•		_		x = 0.2								x = 0	.75			
	2 3 4 5	0.21442E=0 0.52635E=0 0.78542E=0 0.71287E=1 0.24643E=1 0.17783E=1	5 8 1 4	0.91846E=0 0.28818E=0 0.46056E=0 0.43069E=1 0.29225E=1 0.13115E=1	7 0 (0.23575E=0 0.40977E=0 0.40331E=0 0.41392E=1 0.82918E=1 0.48458E=1	6 9 2	0.92134E=04 0.17723E=04 0.17627E=09 0.90775E=13 0.10901E=13 0.61490E=15		0.15702E-0 0.14938E-0 0.17462E-0 0.13820E-0 0.73445E-0 0.27682E-1	2 4 6 9	0.22462E- 0.70059E-0 0.97985E-0 0.82138E-0 0.45001E-0 0.17284E-1) 2) 4) 6 8	0.52036E= 0.98156E= 0.84198E= 0.46496E= 0.17978E= 0.51315E=1	03 05 07	-0.19640E-01 0.12639E-01 0.23201E-05 0.15945E-07 0.68303E-10 0.20610E-12	
	, ,				x = 0.30								x = 0.4	90			
5	0 0	0.39323E-02 0.13152E-04 0.28143E-07 0.36452E-10 0.25024E-13 0.21840E-16	0	0.16011E-01 0.71682E-04 0.16477E-06 0.22169E-09 0.19200E-12 0.13790E-15		•59752E-03 •14864E-05 •20889E-08 •18902E-11 •15444E-13 •13411E-14	0	0+21254E-03 0+61698E-06 0+91547E-00 0+10752E-13 0+10905E-13 0+68435E-15		0.20530E=00 0.21456E=02 0.21456E=04 0.24597E=06 0.24597E=08 0.14908E=08 0.63887E=11		0.25101E-00 0.96805E-02 0.15333E-03 0.14614E-05 0.91114E-08	? .	0.64850E=0 0.154C1E=0; 0.15133E=0; 0.94964E=0; 0.41715E=09 0.13523E=11	2	0.26169E=01 0.12113E=03 0.37520E=05 0.30881E=07 0.15333E=09 0.53222E=12	
1	0.	67240E-02			= 0.35								x = 0.85	i			
3 4 5 6	0. 0. 0.1	28584E-04 82812E-07 14599E-09 16379E-12 28649F-15	0 • 0 • 0 •	25654E=01 15488E=03 46381E=05 86571E=09 10494E=11 91282E=15	0 + 1 0 + 1 -0 + 1 -0 + 1	131715-02 439625-05 339775-08 .01985-10 55615-13 22585-14	0 . 0 . 0 .	412695-03 174965-05 359035-05 482315-11 713805-14 161775-14	0	-25869E-00 -30421E-02 -42473E-04 -42624E-06 -26986E-08 -14015E-10	0000	.27025E-00 .13116E-01 .23341E-03 .25095E-05 .17566E-07 .87210E-10	000	0.77670E-01 0.23371E-02 0.25230E-04 0.18573E-06 0.71978E-09 0.33648E-11	0	**39379E-01 *57667E-04 *57245E-05 *55868E-07 *32594E-09 *12847E-11	
1	0.1	09625-01	•		= 0.40							ж	= 0.90				
2 3 4 5 6	0.5 0.2 0.48 0.71	6131E-04 1097E-06 8544E-09 9904E-12 0710F-14	0 . 3 0 . 1 0 . 2 0 . 4	8614E-01 0187E-03 2294E-05 9390E-08 5724E-11 9097E-16	0+11 3+28 0+44 0+44	5175E-02 1311E-04 1076E-07 1805C-10 866E-13 953F-15	0.4 0.1 0.1 5.1	89455-02 25435-05 16465-07 94665-10 89855-12 85547-15	0 • 0 • 0 •	31381E-00 42683E-02 63915E-04 71333E-06 54254E-08 79387E-10	0. 0.	28081E-00 17459E-01 34674E-03 41758E-05 32961E-07 18249E-09	0. 0. 0.	90035E=01 34355E=02 44000E=04 34955E=06 19386E=08 79486E=!1	-0. 0. 0.	54772E-01 11415F-03 82063E-05 99929E-07 55926E-09 29498E-11	
1	0.17	266E - 01	0.10	x = 3385-01	0.45							x =	0.95				
2 3 4 5 6	0+102 0+481 0+139 0+268	209E-03 148E-06 198E-08 192E-11 68E-14	0.54 0.27 0.84 0.16	376E-03 970E-05 612E-08 567E-10 193E-13	0.260 0.815 0.164 0.210 -0.743	049E-02 61E-04 34E-07 47E-09 67E-12 63E-15	0.91 0.32 0.69	575E-03 571E-05 649E-07 962E-16 354E-12 112E-14	0.9 0.9 0.11	8693E-00 9379E-02 4291E-04 1616E-05 9166E-08	0.5 0.5	8279E-00 2870E-01 3397E-03 7559E-05 P426E-07	0.7 0.6 0.3	01 77E-06 8973E-02 1644E-04 3551E-06 9247E-08 7907E-10	0.10 0.10 0.11	1729E-01 6330E-03 1936E-04 5795E-06 757E-08	
į		775-01	0.761	x = 0 405-01	.50 9+8205	55-02	A 10:					X = 1	.00				
2 3 4 5 5	0:1749 0:1007 0:3609 0:8573 0:1389	77E-05 76E-09 45-11	0.920	94E-03 22E-05 76E-07 18E-16	0.5515 0.2118 0.5265 0.4304 0.1146	0E=04 1E=06 3E=09 2E=12	0.178	8E-12	0:820 0:136 0:184 0:172	994E=00 119E=02 170E=03 59E=05 30E=07 95E=09	0.299 0.716 0.106 0.103	7815-00 1156-01 1265-03 585-04 895-06 575-09	0.67 0.111 0.111 0.766	296E-00 756E-02 950C-03 199E-05 13E-08 77E-10	-0.108 0.132 0.270 0.236	555E+01 379E+02 40E+04 30E+06 855+08 945+10	

							x = 1.0	5									
				a,r		\mathbf{a}_{n}^{i}		br n		b _n					x = 1.55		
			3 0.	45595E= 11257E= 19519E= 28700E= 29428E= 21616E=	01 0 03 0 05 0	.26834E. 37562E. 10058E. 16434E. 17665E. 13328E.	01 0 02 0 04 5 06 0	0.12382E-0 0.91055E-0 0.17528E-0 0.19182E-0 0.14472E-0 0.80519E-10	2 -0.20 3 0.13 5 0.41	919E-00 919E-07 761E-04 644E-06 301E-08	0 s 0 s 0 s 1	2F 56924E 0 15748E-0: 40263E-0: 10311E-0: 1477E-0: 13438E-07	2 0	21 n 22843E-00 14402E-00 14332E-01 50671E-03 11894E-04 19656E-06	0.	br n 26804E=0 51202E=0 41860E=0 12689E=0 22305E=0 27316E=0	1
		1 2	0.4	8937E-0		5695E-0									- 1.60		0.485045-08
		3 4 5 6	0.41 0.41	5356E=0 7498E=0: 1750E=0: 1032E=07	0.1 0.2 0.2	7155E-0 3858E-0: 4822E-0: 9286E-06 261E-08	0 • 2 0 • 2 0 • 3 0 • 2	13467E-00 11897E-01 16426E-03 2008E-05 6533E-07 6200E-09	-0:126; -0:3593 0:1009 0:5129 0:7283 0:5260;	13E-02 18E-04 7E-0€ 15-06	0.18 0.52 0.13 0.30	459E 00 604E-00 611E-02 921E-01 564E-05	0.19 0.17 0.66 0.16	3293E-00 9878E-00 7699E-01 767E-03 713E-04 450E-06	0 - 2 9 - 5 0 - 5; 0 - 1; 0 - 3;	8858E=00 6641E=01 2007E=02 7526E=03 1363E=05 106E=07	-0-28761E-00 -0-59681E-01 -0-24557E-01 -0-24171E-04 0-15070E-06 0-68922E-08
		1 2 1 4 5 6	0.208 0.382 0.655 0.798	59E 00 11E-01 14E-03 19E-05 79E-07	0.58 0.18 0.367 0.474	569E-00 104E-01 116E-02 190E-04 48E-06 79E-03	Cal4 Cal5 Cal6 Cal6 Cal6 Cal6 Cal6 Cal6 Cal6 Cal6	581E-00 131E-01 742E-03 140E-05 328E-07 576E-09	-0.14508 -0.57078 -0.16484 0.85794 0.121078 0.982155	E-02 E-05 E-06	9.220	2F=04	0 • 237 0 • 208 0 • 216 0 • 872; 0 • 232; 0 • 4355	81E-00 93E-00 75E-01 20E-03 25E-04	0.623 0.636 0.238	025E-00 195E-01 48E-02 90E-03 71E-05 96E-07	-0s29914E-00 -0s69315E-01 -0s39822E-02 -0s42373E-04 0s8280E-07
		1	0.9359	15 00	0,2359		= 1.20							X = 1		,62-0,	0 • 94362E=0u
		2 3 4 5 6	0.2799: 0.92740 0.92740 0.96949 0.12749 0.12167	E-03 E-05	0.7135 0.2520 0.2529 0.25270 0.75270 0.74271	1E-01 5E-02 1E-04	0.1876 0.5614 0.8304 0.8231 0.5981	9E-03 9E-05 8E-07	-0.163265- -0.653205- -0.274005- 0.113145- 0.17425- 0.177425-	04 05 07	0.55009 0.29312 0.88327 0.24825 0.601021 0.110928	E=00 E=02 E=03 E=05	0 = 24274 0 = 21601 0 = 26330 0 = 11291 0 = 31939 0 = 63633	9E-00 IE-00 IE-01 E-02 E-04	0.3329 0.6839 0.7689 0.3199 0.71450	9E-01 1E-02 1E-03 1E-05	-0.30944E-0G -0.7951E-02 -0.45436E-02 -0.45436E-02 -0.45436E-07 0.12356E-07
	2 3 4 5 6		0.55018: 0.37332E 0.72016E 0.14021E 0.19971E 0.20634E	-01 -03 -04	0.22849 0.859398 0.333498 0.76832E 0.11711E C.12545E	-01 -02 -04 -05	0.16998 0.22740 0.79259 0.12953 0.13988 0.110388	E-01 E-03 E-04 E-06	-0.18105E-0 -0.12137E-0 -0.75656E-0 -0.13802E-0 0.30157E-0 0.31075E-0	1 4 5	0.54043E 0.26586E- 0.21352E- 0.32820E- 0.83151E- 0.16170E-	-00 -01 -03	0.24759E 0.22040E 0.31727E 0.1449E 0.43492E 0.91899E	-00 -01 -02 -04	75 0+356421 0+744078 0+917098 0+42328E 0+102418: 0+16582E	-01 -02 -03	-0.31837E-00 -0.90359E-01 -0.59700E-02 -0.11013E-03 -0.4120AE-06 -0.15279E-07
	1 2 3 4 5 6	0	•56055E •49262E •97522E •20092E •30751E •34276E	01 03 04	0.22351E- 0.10197E- 0.43601E- 0.1085E- 0.17898E- 0.20746E-	00 00 02 03 05	0.18341E 0.26995E 0.10963E 0.19807E 0.23257E 0.19870E	-01 -0 -02 -0 -04 0.	0.19838E-or e16562E-o1 e15691E-o5 e14959E-o5 e45137E-o7 e52798E-or		0.52927E 0 0.31767E=0 0.14507E=0 0.43120E=0 0.11409E=0	9 0	0 - 25198E - 0 0 - 22227E - 0 0 37924E - 0 0 18464E - 0 0 56672E - 0 13128E - 0	00 0: 1 0: 2 0:	-38111E- -31650E- -19815E- -55259E- -14487E- -25069E-0	00 -0 01 -0 01 -0 03 -0	3+325#2E-00 0+101#9E-00 1-7690#E-02 1-16735E-03 1-7402E-07
	1 2 3 4 5	0 • 1 0 • 2 0 • 4	36741E 06 64201E-01 3110E-02 3443E-04 6606E-06 5861E-08	0000	•22093E-00 •11909E-00 •56400E-02 •15127E-03 •26903E-05 •33647E-07	0	0-19793E-0 0-31475E-0 0-14868E-0 0-29722E-0 0-37894E-0 0-34972E-0	1 -0.2 -0.2 -0.1	215146-00 216176-01 269556-05 26266-05 50296-07 71666-09	0	-51673E 00 -34792E-00 -18429E-01 -56318E-03 -15933E-04 -33352E-06	0.2 0.4 0.2 0.7	25576E-00 2208E-00 4969E-01 3342E-02 8450E-04 8561E-05	0.8 0.1; 0.7; 0.20	0633E-00 8973E-01 2625E-01 1225E-03 1234E-04	-0. -0. -0.	33169E-00 11352E-00 77944E-02 24458E-03 0917E-03
1		0.51	1355 00			= 1.40								0.37	410E-06	0.1	72218-07
9 6		0.82 0.17 0.39 0.69	#66E-01 #66E-01 #609E-02 #20E-04 #12E-06	0 . 7 0 . 2 0 . 3	22050E-00 13677E-00 2216E-02 0815E-03 9822E-05 3591E-07	0 a 1 0 a 4 0 a 6 0 a 6	21362E-00 36143E-01 19762E-02 3812E-04 6578E-06 0265E-08	-0.27 -0.48 0.41 0.89	1266-00 8896-01 9006-63 9286-06 7705-07	0 a 7 0 a 7 0 a 7	90292E 00 97617E-00 19266E-01 19140E-03 0992E-04 7238E-06	0 • 22 0 • 52 0 • 29; 0 • 104	#72E-00 03ZE-00 03ZE-00 090E-01 296E-02 401E-03 186E-05	0.944 0.146 0.906 0.279	21+E-00 12+E-01 101E-01 14E-03 15E-04 29E-06	-0.12 -0.12 -0.35 -0.37	1987E-00 980E-00 128E-01 197E-01 1804E-05 132E-07
1		0 • 572	9E 00	0.22	ж = 1856-00	I.45								1.95		0412	1326-07
3 4 5 6	Ċ	0.2324 0.5518	2E-04	0.15 0.91 0.28; 0.58;	432E-00 580E-02 306F-03 105E-05 129E-07	0:40 0:25 0:63 0:95 0:10]	053E-00 981E-01 841E-02 475E-04 119E-05 182E-07	-0.2466 -0.43474 -0.47786 -0.1621 0.1177 0.2180	95-01 86-03 05-05 85-06	0.40 0.29 0.94 0.28	797E-00 211E-00 182E-01 470E-03 173E-04 139E-06	0.2603 0.2179 0.6169 0.3651 0.1368 0.3604	70E-00 31E-00 35E-01 3E-02 0E-01	0.4583 0.1052 0.1675; 0.11416 0.38056 0.80264	6E=00 3E=01 5E=02 5E=04	-0.338 -0.138 -0.148 -0.4965 -0.5399	44E-0G 97E-02 13E-03 14E-08
	0.	57202 12933 30682 75156 169298 216608	E=00 E=02 E=04	0.2246 0.1705 0.1150 0.3807 0.8365 0.1293	52E-00 P9E-00 PE-01 OE-03	-	16E=02 3E=04 4E=05	-0.26126 -0.42360 -0.41841 -0.4359000 0.144665 0.330195	E-01 E-02 E-05	044725 04425 04333 041213 043756 049240	8E-02	0-26155 0-21411 0-71354 0-492111 0-178558 0-495538	E-00 E-01 E-02	00 0 484936 0 114316 0 19086 0 142126 0 512866 0 115506	-00 -01 -02 -04	-0.3367; -0.1515; -0.16063 -0.66235 -0.410335 -0.32214	55-00 5-00 5-01 5-03 5-04

APPENDIX E

Calculation of Light-Scattering Functions for a Normal Distribution of Sizes

The density function for a normal distribution is given by equation 7.1. Figure E.1 is a plot of f(x) against x for given values of the mean μ and the standard deviation σ . It is seen from this figure that μ is the number mean about which x is distributed while the standard deviation σ determines the extent of the distribution. Small values of σ give steep or narrow distributions, while large values of σ give broad distributions.

The calculation of i_1 and i_2 defined by equations 7.2 and 7.3 involves first, the calculation of i_1 and i_2 for a given value of m for a range of x at different values of θ . At a given angle θ , each of these values of i_1 and i_2 which correspond to a particular value of x is multiplied by the appropriate value of f(x) to obtain the integrands, $i_1f(x)$ and $i_2f(x)$ for various values of x. The numerical value of the integrals can then be evaluated by a numerical method. Simpson's rule was used to evaluate these integrals.

Simpson's rule can be written as

$$\int_{x_0}^{x_n} y dx = h/3[y_0 + y_n + 4(y_1 + y_3 + \dots + y_{n-1}) + 2(y_2 + y_4 + \dots + y_{n-2})]$$
E.1

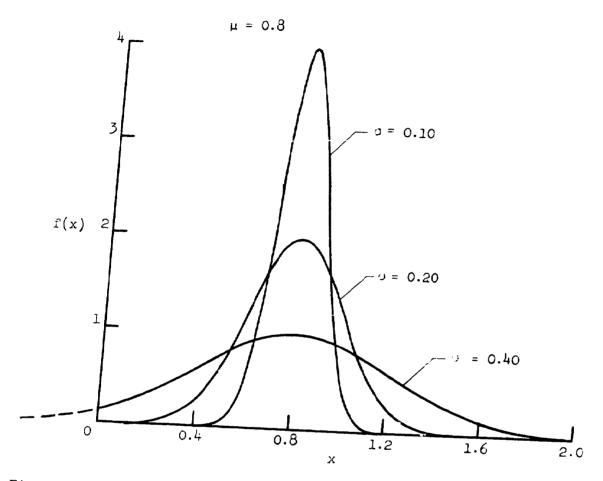


Figure E.1.- Density function, f(x), for a normal distribution plotted against x for μ = 0.8 at various values of the standard deviation.

where h is the difference between successive values of x, y_0 is the first value of the integrand which corresponds to x_0 , and y_n is the last value of the integrand which corresponds to x_n . The value of h was taken as 0.05 to correspond to the available data cards used in the program already presented in appendix C. In those calculations, x was taken in increments of 0.05.

The modified Fortran program which was used to calculate \bar{i}_1 and \bar{i}_2 for various combinations of μ , σ , and θ follows.

For the purpose of calculation, m was set equal to 1.71 - 0.76 i and the functions i_1 and i_2 were calculated for $\theta = 0^{\circ}(5^{\circ})$ 180° for $\mu = 0.80$ and $\sigma = 0.10$, 0.20, and 0.40. The results of these calculations are plotted in Figure E.2 along with the monodispersed functions i_1 and i_2 for x = 0.80 (i.e., $\mu = 0.80$ and $\sigma = 0$). It is seen from Figure E.2 that the effect of increasing the value of σ at a given value of μ is to increase the apparent diameter as determined by light scattering.

A listing of the modified Fortran program for a normal distribution of sizes follows:

```
GOT0900
  60 READ30, ALPHA, (C(N), N=1,7), (S(N), N=1,7)
     GOTO(70,80,70,1000), IND
  70 READ40,U,V
     J=O
  80 DCONS=U**2+V**2
    J=J+1
    DREAL=U/DCONS
    DIMAG=V/DCONS
    SINH=(EXPF(2.*ALPHA*V)-EXPF(-2.*ALPHA*V))/2.
    COSH=(EXPF(2.*ALPHA*V)+EXPF(-2.*ALPHA*V))/2.
    GCONS=COSH-COSF(2.*ALPHA*U)
    GREAL(1)=SINF(2.*ALPHA*U)/GCONS
    GIMAG(1)=SINH/GCONS
   DO90N=2, NFINAL
   P(N)=N-1
   L(N)=N-1
   EREAL=(2.*P(N)-1.)*DREAL/ALPHA-GREAL(N-1)
   EIMAG=(2.*P(N)-1.)*DIMAG/ALPHA-GIMAG(N-1)
   GCONS=EREAL**2+EIMAG**2
   GREAL(N)=EREAL/GCONS
   GIMAG(N)=-EIMAG/GCONS
   HREAL=U*GREAL(N)+V*GIMAG(N)
  HIMAG=U*GIMAG(N)-V*GREAL(N)
  QREAL=HREAL*C(N)-C(N-1)
  QIMAG=HIMAG*C(N)
  RREAL=HREAL*S(N)-S(N-1)
  RIMAG=HIMAG*S(N)
  SCONS=RREAL**2+RIMAG**2
  SREAL=(QREAL*RREAL+QIMAG*RIMAG)/SCONS
  SIMAG=(QIMAG*RREAL-QREAL*RIMAG)/SCONS
 BCONS=(1.-SIMAG)**2+SREAL**2
 BREAL(N)=(1.-SIMAG)/BCONS
 BIMAG(N)=-SPEAL/BCONS
 TREAL=DREAL*GREAL(N)-DIMAG*GIMAG(N)
 TIMAG=DIMAG*GREAL(N)+DREAL*GIMAG(N)
 UREAL=P(N)*(1.-DREAL**2+DIMAG**2)/ALPHA
 UIMAG=-2.*P(N)*DREAL*DIMAG/ALPHA
 VREAL=TREAL+UREAL
 VINAG=TIMAG+UIMAG
WREAL=VREAL*C(N)-C(N-1)
WIMAG=VIMAG*C(N)
XREAL=VREAL*S(N)-S(N-1)
XIMAG=VIMAG*S(N)
YCONS=XREAL**2+XIMAG**2
YREAL=(WREAL*XREAL+WIMAG*XIMAG)/YCONS
YIMAG=(WIMAG*XREAL-WREAL*XIMAG)/YCONS
ACONS=(1.-YIMAG)**2+YREAL**2
```

```
AREAL(N)=(1.-YIMAG)/ACONS
          AIMAG(N) = -YREAL/ACONS
       90 CONTINUE
         DO14CM=1, MFINAL
         G=M-1
         ANGLE (M) = ANGSIZ*G
         SARP=0.0
         SART=0.0
         SBRP=0.0
         SBRT=0.0
         SAIP=0.0
         SAIT=0.0
         SBIP=0.0
        SBTT=0.0
        DO130N=2, NFINAL
        P(N)=N-1
        Z(N)=(2.*P(N)+1.)/(P(N)*(P(N)+1.))
        ARP=AREAL(N)*Z(N)*PI(N,M)
        SARP=SARP+ARP
       ART=AREAL(N)*Z(N)*TAU(N,M)
        SART=SART+ART
       BRP=BREAL(N)*Z(N)*PI(N,M)
       SBRP=SBRP+BRP
       BRT=BREAL(N)*Z(N)*TAU(N,M)
       SBRT=SBRT+BRT
       AIP=AIMAG(N)*Z(N)*PI(N,M)
       SAIP=SAIP+AIP
      AIT=AIMAG(N)*Z(N)*TAU(N,M)
      SAIT=SAIT+AIT
      BIP=BIMAG(N)*Z(N)*PI(N,M)
      SBIP=SBIP+BIP
      BIT=BIMAG(N)*Z(N)*TAU(N,M)
      SBIT=SBIT+BIT
  130 CONTINUE
      PRI(J,M)=(SARP+SBRT)**2.+(SAIP+SBIT)**2.
      PLI(J,M)=(SART+SBRP)**2.+(SAIT+SBIP)**2.
 140 CONTINUE
 900 READ50, IND
     GOTO(60,60,70,1000), IND
1000 JFINAL=J
1030 READ1010, SIGMA,UM
     IF(SIGMA)1001,9999,1001
1001 D01002J=2, JFINAL
     E=J-1
1002 DENS(J)=(.399/SIGMA)*EXPF(-.5((.05*E-UM)/SIGMA)**2)
    DO1005M=1, MFINAL
    SEPRI=0.0
    SOPRI=0.0
```

```
SEPLI=0.0
      SOPLI=0.0
      DO1003J=2, JFINAL
      YPRI(J,M)=PRI(J,M)*DENS(J)
 1003 YPLI(J,M)=PLI(J,M)*DENS(J)
      D01004J=2, JFINAL, 2
      SEPRI=SEPRI+YPRI(J,M)
      SOPRI=SOPRI+YPRI(J+1,M)
      SEPLI=SEPLI+YPLI(J,M)
 1004 SOPLI=SOPLI+YPLI(J+1,M)
      PRIA(M)=.05/3.*(4.*SEPRI+2.*SOPRI-YPRI(JFINAL,M))
1005 PLIA(M)=.05/3.*(4.*SEPLI+2.*SOPLI-YPLI(JFINAL,M))
1010 FORMAT(2F10.3)
1015 FORMAT(7H1SIGMA=F5.3,6H MU=F5.3,6H V=F5.3)
1020 FORMAT(50H
                   ANGLE
                              PERPENDICULAR
                                                 PARALLEL)
1025 FORMAT(F8.1,2E21.5)
     PRINT1015, SIGMA, UM, U, V
     PRINT1020
     PRINT1025, (ANGLE(M), PRIA(M), PLIA(M), M=1, MFINAL)
     GOTO1030
9999 CALL EXIT
    END
```

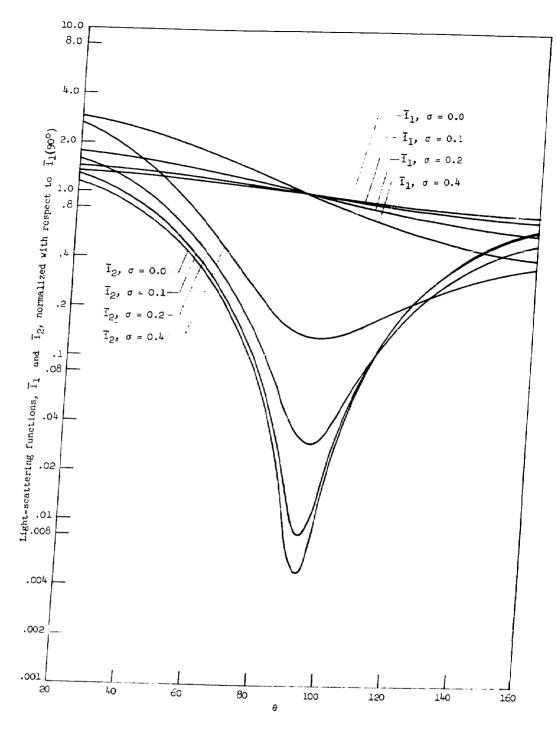


Figure E.2.- Light-scattering functions, \tilde{i}_1 and \tilde{i}_2 , normalized with respect to $\tilde{i}_1(90^\circ)$ for $\mu=0.8$ and various values of the standard deviation for m=1.71-0.76 i.

APPENDIX F

Further Description of the Light-Scattering Apparatus
A brief description of the light-scattering apparatus has already
been given and a schematic drawing is shown in Figures 8.1 and 8.2.
A more detailed description of the various apparatus components than
was given before is presented here. Since the selection of each of the
various components must be based on consideration of other components,
it is necessary to discuss the interdependent components together.

Let us begin with the light source and its interdependent components. A GE B-H6 high brightness mercury arc lamp was chosen for the light source for several reasons. First, the extremely high brightness of this lamp makes it possible to develop a scattered-light intensity which is quite easily measured with an inexpensive multiplier phototube. In addition, the radiation from this lamp is concentrated in relatively narrow spectral bands, characteristic of mercury. Thus, by using an appropriate optical interference filter, a suitably measurable scattered-light intensity, having a fairly narrow wavelength band width, is possible. Another reason for selecting this lamp was that it operates from a 60-cps source and therefore generates modulated light at 120 cps. This, with an appropriate electronic filter, permits easy separation of the scattered-light signal from the steady-flame signal.

The average spectral distribution of the radiation from several B-H6 lamps is given in reference 30. The radiation in watts per

steradian perpendicular to the axis of the lamp is tabulated for a large number of wavelength intervals. Figure F.1 is a histogram based on these data. This figure indicates that the various spectral bands are broadened and there is a slight continuum superimposed in the low pressure mercury spectrum. This broadening and superpositioning of a continuum is due to the high pressure at which this lamp operates. The interal operating pressure of this lamp is about 100 atm. A part of the apparent broadening shown in Figure F.1 is due to the relatively large interval size of the wavelength over which each measurement was made, so that the actual spectral distribution may show steeper and higher peaks. The highest peaks correspond approximately to the principle mercury lines of 3650 Å, 4047 Å, 4358 Å, and 5461 Å.

Since it is desirable to produce a strong high quality signal as a measure of the scattered-light intensity, the selection of the particular spectral band must include a consideration of the spectral sensitivity of the multiplier phototube. The RCA 931-A multiplier phototube which was chosen typically has a spectral sensitivity as shown in Figure F.2 which is based on data from reference 27. It is seen from Figure F.2 that the maximum sensitivity of this type of tube occurs at about 4000 Å. Even though the sensitivity of the tube would be greater for 4047 Å than for 4358 Å, the band about 4358 Å is stronger and more than compensates for the small difference in sensitivity. The band about 4358 Å was selected. In order to define a narrow wavelength band about 4358 Å, a relatively narrow band optical

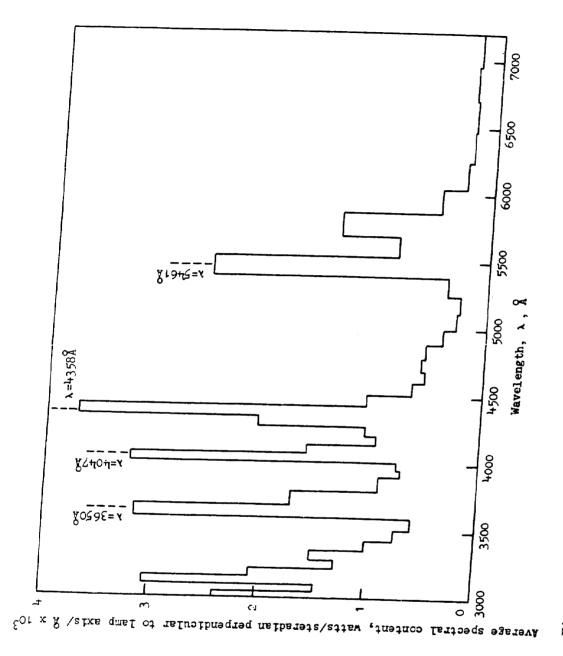


Figure F.1.- Average spectral distribution from a GE B-H6 mercury arc lamp.

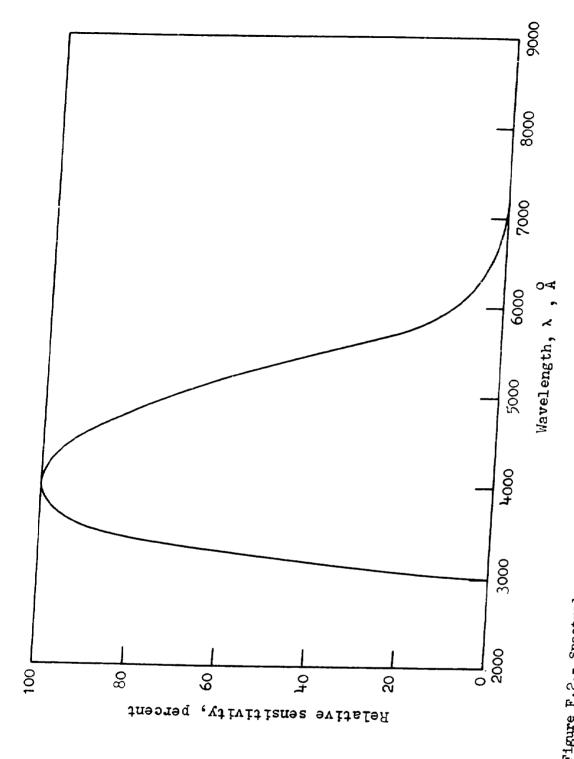


Figure F.2. Spectral sensitivity characteristic of an RCA 931-A multiplier phototube for equal values of radiant flux at all wavelengths.

interference filter was used. The particular filter used was a type B-l optical interference filter with a peak wavelength of 4358 Å manufactured by Baird-Atomic, Inc. This filter has a bandwidth at half of the peak transmission of approximately 65 Å and has a peak transmission of about 50 percent. This bandwidth corresponds to a range in λ about 4358 Å of only ± 0.75 percent which is felt to be quite acceptable for this investigation. In addition to defining a relatively narrow wavelength band, this filter also reduces the steady flame light intensity, particularly wavelengths outside the band pass, when it is placed between the flame and the collection system for the scattered light.

Achromatic lenses of good quality were used to assure that the alinement and focal adjustment of the optical system (made by using easy-to-see unfiltered light from the light source) remained the same for the light-scattering measurements which were made by using roughly monochromatic light. The two short focal length lenses used, each had a length of 78 mm and the two long focal length lenses each had a focal length of 508 mm. It would be possible, of course, to replace both combinations of long and short focal length lenses in each collimation system by a single lens in each system.

The size of the circular aperture at the light source, along with the lenses of the collimation system for the incident light, determine the size of the circular cross section of the cylindrical scattering volume in the flame. As already mentioned, it is desirable to keep this probe volume as small as possible but still large enought to produce a measurable quantity of scattered light. The square aperture between the multiplier phototube and polarization filter (see Fig. 8.1) defines a square-shaped cross-sectional view of the flame which is received by the scattered-light collection system. The intersection of this projected square with the cylindrical scattering volume determines the observation volume or probe volume in the flame.

Based on preliminary measurements with various aperture sizes and the above-mentioned lens system, it was decided to use a circular aperture at the light source which consists of a 0.0135-inch-diameter hole through 0.004-inch copper sheet. Provision is made for alining the aperture with the axis of the lamp. With the aperture alined with the lamp and the lenses of the incident light system properly focused, an image of the circular aperture is formed at the center line of the burner tube which was measured to be 0.078-inch diameter. This produces an approximately cylindrical scattering volume in the flame with a diameter of approximately 0.078 inch.

The square aperture located in front of the multiplier phototube is formed by the intersection of two narrow rectangular slits as shown in Figure F.3. Each slit is 0.015 inch wide and about 0.375 inch long, also made from 0.004-inch copper sheet. The vertical position of the horizontal slit is controlled by moving the stock containing this slit, denoted as H in Figure F.3, vertically up or down. The horizontal position of the vertical slit is likewise controlled by

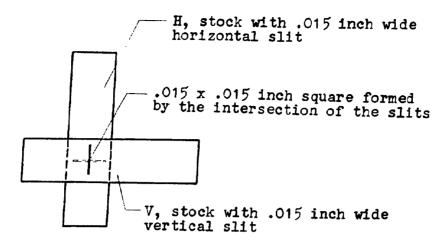


Figure F.3.- Schematic drawing of square aperture system.

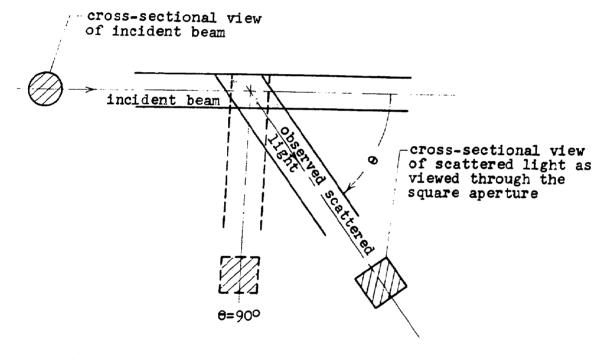


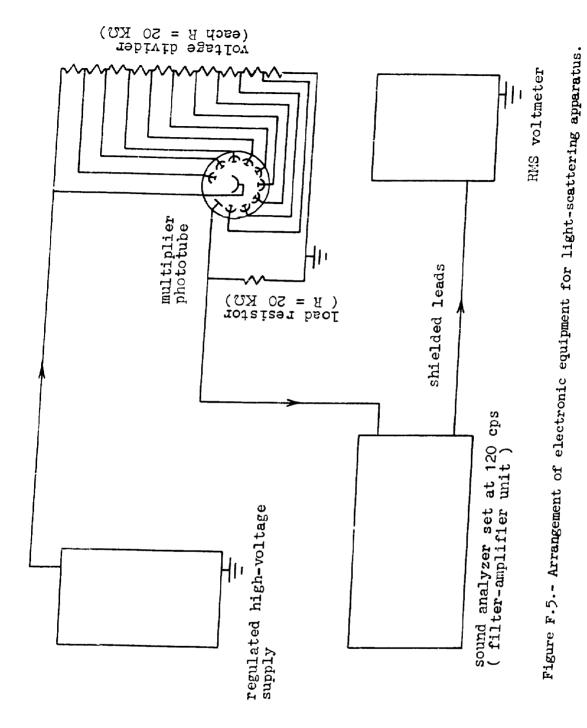
Figure F.4.- Schematic drawing of observed scattering volume at an arbitrary angle θ .

horizontal movement of V. The simultaneous movement of the two slits permits the square aperture to be properly positioned. The viewing area of the flame seen through this square aperture was measured by placing a lamp behind the aperture and observing the image formed at the center line of the burner tube. This image was found to be a 0.094-inch square so that when the scattered light is observed at right angles to the incident light, that is, $\theta = 90^{\circ}$, the observed scattering volume is a right-circular cylinder with a diameter of 0.078 inch and a length of 0.094 inch. It is noted that the square is larger than the diameter of the incident beam through the flame by 0.016 inch. This permits some freedom in the vertical alinement of the scattered-light collection system with respect to the incident beam. The observed scattering volume formed by the intersection of the cylinder and projected square is of course larger at any angle θ different than 90° . It is easily seen from Figure F.4 that the observed scattering volume at any angle θ is equal to the observed scattering volume at $\theta = 90^{\circ}$ multiplied by $1/\sin \theta$.

It is of course very important to have the optical system properly alined. The method of determining the soot particle size by comparing the shape of experimental and theoretical light-scattering patterns requires that the relation between the observed scattering volume and the angle θ be known. This volume is proportional to $1/\sin\theta$ only if the system is properly alined. See appendix G.

Figure F.5 shows the arrangement of the electronic equipment. The high voltage required to operate the RCA 931-A multiplier phototube is supplied by a regulated high-voltage supply, model RE-1602, Northeast Scientific Corp. The ten resistors which make up the voltage divider and the load resistor are each 20 KM. The signal from the multiplier phototube which is produced by that portion of the scattered light which strikes the photocathode is then passed through a sound analyzer, type 760-B, General Radio Co. This unit is turned to 120 cps to match the modulated light from the mercury lamp and has the measured filtration characteristics shown in Figure F.6. Included as an integral part of the analyzer is a variable amplifier which has a maximum amplification factor of 80. The filtered and amplified signal is then measured with a root-mean-square voltmeter, model 400c, Hewlet Packard. Shielded cables were used for all leads to reduce the pick-up of extraneous signals.

This arrangement was very effective in practically eliminating the steady signal due to the flame. A test was made in which the filtered signal was measured on an RMS voltmeter and the unfiltered steady signal was measured on a d-c voltmeter for both the perpendicular scattered light at $\theta = 30^{\circ}$, including of course the flame light, and for the flame light only, without scattering. The ratio of the signal from the flame light only to the signal from the scattered light and flame light was less than 0.1 percent for the filtered RMS signal, but was 8.4 percent for the unfiltered steady signal.



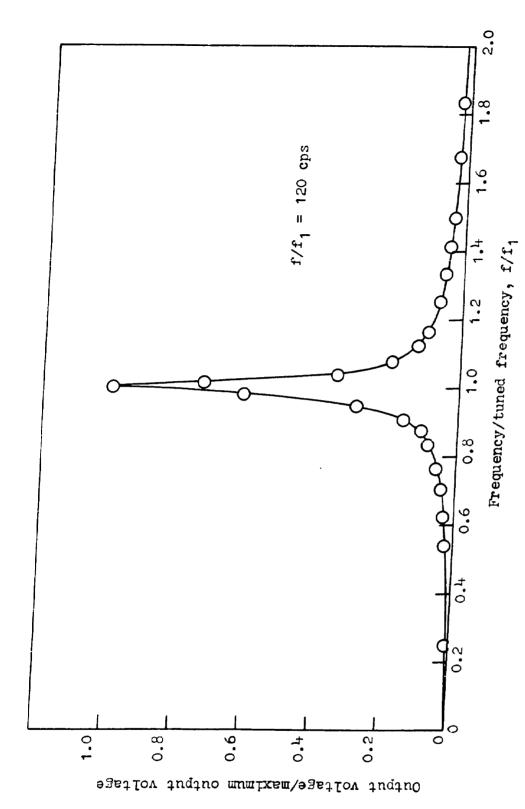


Figure F.6.- Measured filtration characteristics of the sound analyzer tuned to 120 cps.

A 30-inch length of 3/8-inch copper tubing was used for the burner tube. The inside diameter of this tube was 0.317 inch. The corresponding length-to-diameter ratio was approximately 95.

A wind screen was placed concentric about the burner tube and consisted of a 4-inch-diameter by 18-inch-long cylinder constructed from perforated sheet metal around which one layer of 50-mesh copper screen was wrapped. The holes in the perforated sheet were 1/8-inch diameter on 3/16-inch centers. The top of the cylindrical wind screen was open and the bottom was covered with the perforated sheet and copper screen with an opening just large enough for the burner tube. A 5/8-inch slit was cut in the wind screen at 10-1/2 inches above its base to permit the incident and scattered light to enter and leave the flame. It was necessary to mask this slit down with paper to 3/8 inch in order to stabilize the flame.

The benzene-air mixing system is shown in Figure F.7. The supply of benzene is contained in a cylindrical brass reservoir having a diameter of 2.5 inches and a depth of 2.5 inches. This gives a storage volume of approximately 200 cc. A needle valve is located directly below the benzene reservoir for crude adjustment of the flow and for shutting off the flow. Directly following this valve, the benzene flows through a length of capillary with its lower end drawn to a taper to give small drops of uniform size. The tip of this capillary on which the benzene drops form and drop from is visible through the glass tee which encloses the capillary. The benzene flow

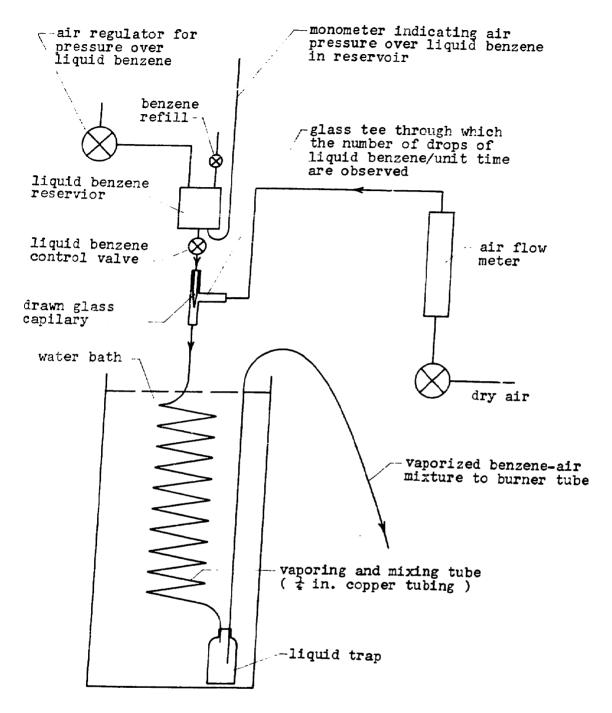


Figure F.7.- Benzene-air mixing system.

rate is determined by counting the drops per unit time and comparing with a previous calibration.

Fine adjustment of the benzene flow is facilitated by regulated air pressure over the benzene in the air-tight reservoir. The change in benzene flow rate with time due to the decrease in the level of benzene in the reservoir was made negligible by holding this applied pressure at about 25 cm of benzene. By taking into account the additional head due to the distance of 16 cm from the base of the reservoir to the tip of the capillary where the benzene exits, the rate of decrease of the total effective benzene head for the largest benzene flow used (0.3 cc/min.) is only about 1 percent per hour. Metered air enters the system at the glass mixing tee and flows in parallel with the liquid benzene into the mixing coil. This mixing coil is made from 14 feet of 1/4-inch copper tube wound with 14 turns and a diameter of 3.75 inches. The coil serves to vaporize all of the benzene and produce a well-mixed benzene-air stream of known composition. A liquid trap is located at the end of the mixing coil to remove any nonvolatile material from the benzene-air stream before it passes through the flexible tubing to the burner tube. A small wad of steel wool is placed between the burner tube and the flexible tubing to act as a flame trap. The benzene-air stream flows up the burner tube to be burned.

The incident light intensity is measured with the apparatus shown in Figure 8.2. The rotating arm is positioned at $\theta = 90^{\circ}$ with the

series of three black-glass mirrors alined as shown in Figure 8.2. Each of these mirrors is a 2-1/2-inch square of "Carrara" black-glass made by the Pittsburgh Plate Glass Co. The reflecting power of a single black-glass mirror was measured by replacing the mercury lamp with a 40-watt frosted bulb. For the vertical setting of the polarization filter, the reflecting power was 0.101 and for the horizontal setting, 0.0137. By using three of these mirrors in series, the incident light is attenuated to a fraction of $(0.101)^3$ or 1.03×10^{-3} for the vertical setting of the polarization filter and to $(0.0137)^3$ or 2.57×10^{-6} for the parallel setting of the polarization filter. As already mentioned, the scattered-light intensity is usually only about 10⁻⁶ times that of the incident light. Therefore, in order to measure the incident light intensity at a level comparable to the scattered-light intensity, either the parallel setting of the polarization filter must be used or three more additional black-glass mirrors in series must be used with the vertical setting. It was decided to use the horizontal setting to avoid using additional mirrors.

APPENDIX G

Procedure for Alining Optical System

The first step in the alinement procedure is to position the circular aperture in front of the mercury vapor lamp so that the center of the aperture coincides with the axis of the lamp and the aperture is filled with light as seen by the first lens L_l as shown in Figure G.1. The axis of the lens tube which holds the lenses and stops for the collimating system for the incident light is then alined with the center of the aperture by using the two alinement jigs shown in Figure G.2. The first jig, J_1 , was accurately machined so as to fit snugly, but able to move freely, inside the lens tube, the inside of which was accurately bored. There is a 0.020-inchdiameter hole through the exact center of this jig. A second jig, J2, was likewise accurately machined to fit in the lens tube but has fitted to it a piece of translucent plastic material on which narrow reference lines at right angles are scribed through its exact center. With the two jigs in place as shown in Figure G.2, the combination of the small hole of J_1 and the intersection of the reference lines on J_2 determines the axis of the lens tube. The axis of this lens tube is alined with the center of the circular aperture by moving the lens tube system on the front surface of the mercury lamp enclosure until the small light spot on jig J2, formed by the narrow beam of light which passes through the centered hole in jig J_1 , coincides with the intersection of the two reference lines on J_2 . When this has been accomplished, the adjustment between the aperture and lens tube is secured.

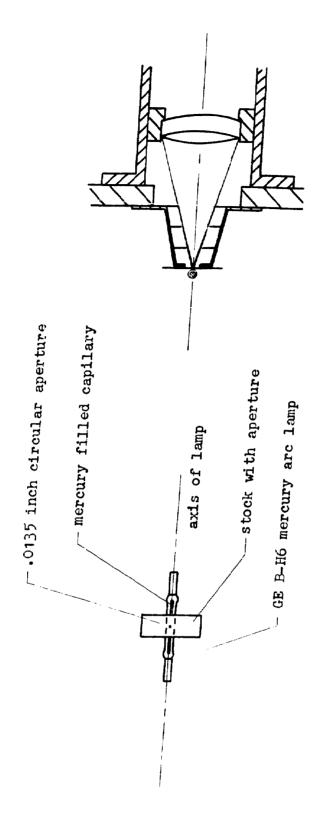


Figure G.1. Alinement of circular aperture with axis of mercury arc lamp. SIDE VIEW

FRONT VIEW

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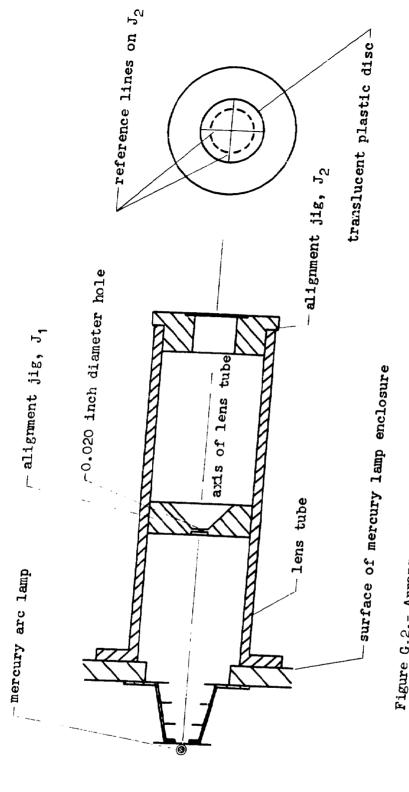


Figure G.2.- Arrangement of alinement jigs for alining lens tube for incident beam.

The two lenses, L_1 and L_2 , which form the incident light are then inserted into the lens tube. (See Fig. 8.1.) These lenses are centered in brass lens holders which have been machined to mate with the inside diameter of the lens tube. The axis of the aperture-lenstube combination for the incident light is then made parallel to a plane which is generated by a point on the rotating arm system. (The plane in which the incident beam lies and which is also parallel to a plane generated by a point on the rotating arm system is called the plane of observation.) This is done by placing the jig J_2 shown in Figure G.2 into the lens tube for the collection system and alternately focusing the incident light on the translucent surface of J_2 when the arm is at $\theta = 0^{\circ}$ and then at $\theta = 180^{\circ}$. When the center of the image of the incident light falls at the same vertical position on J_2 , the axis of the incident beam is in the plane of observation.

The rotating arm is then positioned at $\theta = 0^{\circ}$. With the arm in this position, the lens tube for collecting the scattered light is mounted on the rotating arm with its axis made to coincide with the axis of the incident light beam. This was done by placing the alinement jig, J_2 , shown in Figure G.2, first in one end of the lens tube and simultaneously adjusting the vertical and horizontal position of the tube until the center of the focused image of the incident light was located at the intersection of the right-angle reference lines on J_2 and then repeated with the alinement jig in the other end of the tube. In general, the movement of the lens tube for the alinement of

one end of the tube will alter the alinement of the other so it is necessary to repeat this alinement procedure until both ends are simultaneously alined.

After this has been accomplished, a fixture, which serves to hold and aline the burner tube at the center of rotation of the rotating arm system, is mounted on the spindle about which the arm revolves as shown in Figure G.3. A drawing of the fixture alone is shown in Figure G.4. The outside diameter of the long cylindrical section and the 3/8-inch inside bore were machined in such a way to insure that the center line of the outside diameter coincides with that of the inside bore. The center line of this cylinder which is to hold the burner tube is made to coincide with the axis of rotation of the rotating arm system in the following way. A dial indicator gage is mounted on the rotating arm with its contact point touching the outside diameter of the cylinder at a point which is about an inch above the larger diameter base. The four horizontal set screws are simultaneously adjusted until the indicator gives a constant reading as the arm system is rotated through 360°. The contact point of the dial indicator gage is then made to touch the outside diameter of the cylinder at a point which is about an inch below its top. The indicator reading will be constant through a 360° rotation of the arm system for this location of the contact point only if the center line of this cylinder is perpendicular to the plane in which the arm rotates. In general, however, the center line of the cylinder will be tilted

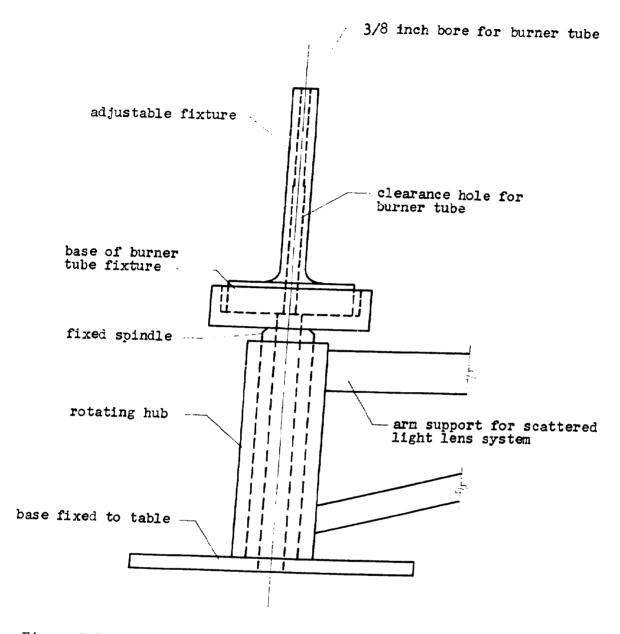
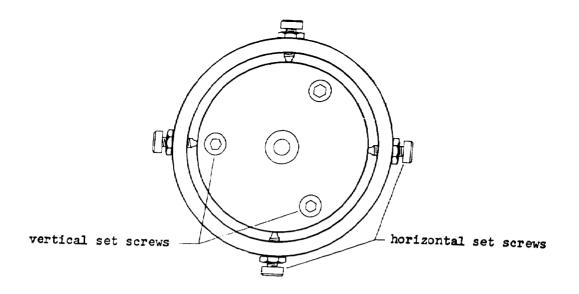


Figure G.3.- Burner tube fixture mounted on spindle of rotating arm system.



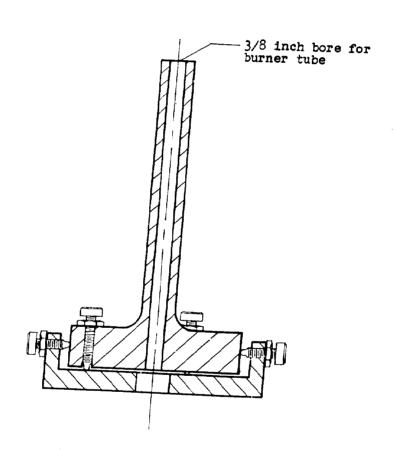


Figure G.4. - Detail view of burner tube fixture.

with respect to this plane so that the indicator readings will vary with angle. Then, depending on the values of these readings, the center line of the cylinder is made more nearly perpendicular to the plane of rotation of the arm by adjustment of the vertical set screws in the base indicated in Figure G.4. This movement will, of course, alter the initial adjustment so that it is necessary to repeat the alinement procedure for the contact point of the indicator in the lower position. The center line of the cylinder coincides with the axis of rotation of the rotating arm when the dial indicator readings remain constant with angular position at both the upper and lower locations of the contact point of the indicator. It was found necessary to repeat this process about three or four times in order to obtain indicator readings at both stations which remained constant to within 0.001 inch.

In order to fix the axis of the incident light beam through the center of rotation of the rotating arm system, the alinement jig shown in Figure G.5 was used. This jig was machined in such a manner that the two diameters have the same center line and the smaller diameter fits snugly into the 3/8-inch-diameter bored hole of the fixture shown in Figure G.4. Referring again to Figure G.5, the 3-inch slot in this jig was milled to a depth of half its diameter so that the center line of this jig lies in the plane of the surface of this slot, midway between the vertical edges. A narrow vertical reference line was carefully scribed on the surface of the slot to physically indicate

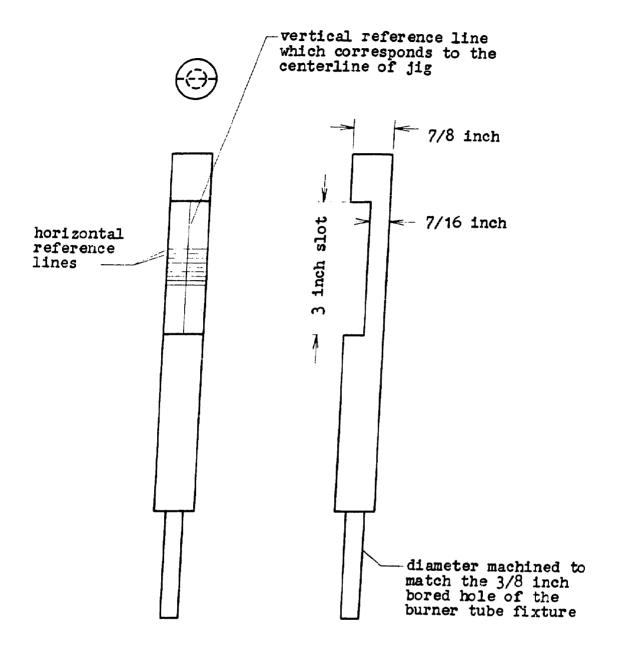


Figure G.5.- Slotted alinement jig for alining the lens systems.

surface for reference. The small diameter section of this jig was fitted into the bored hole of the burner tube fixture and the plane of the surface of the slot was made perpendicular to the incident beam. The incident light system was then moved so as to position the center of the focused image of the incident light on the center line of this jig. During this movement, care was also taken to keep the incident beam in the plane of observation, defined above. Up to this point, then, the incident beam has been alined in the plane of observation and the center of the focused image of the incident light fixed on the axis about which the rotating arm turns.

At this point, the three black-glass mirrors which are used to measure the relative incident light intensity are alined as already shown in Figure 8.2. A more detailed drawing of this arrangement is shown in Figure G.6. Each mirror is mounted on a support system shown in Figure G.7 which permits the plane of the mirror to be rotated about three mutually perpendicular axes. Mirror A and its associated support system is fixed to the stationary optical bench with the approximate center of the mirror made to coincide with the center of the focused incident image and the plane of the mirror adjusted to make an angle of 45° with respect to the incident beam. This is done by setting the plane of the mirror at 45° with the end surface of the lens tube for the incident light by using a 45° square. Then with the other two mirrors not yet mounted, the alinement jig J2 is placed

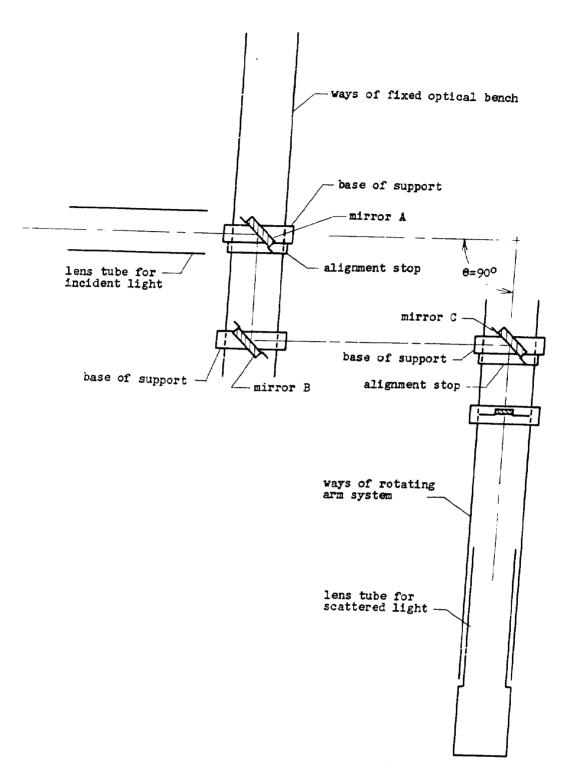


Figure G.6.- Alinement of the three-mirror system.

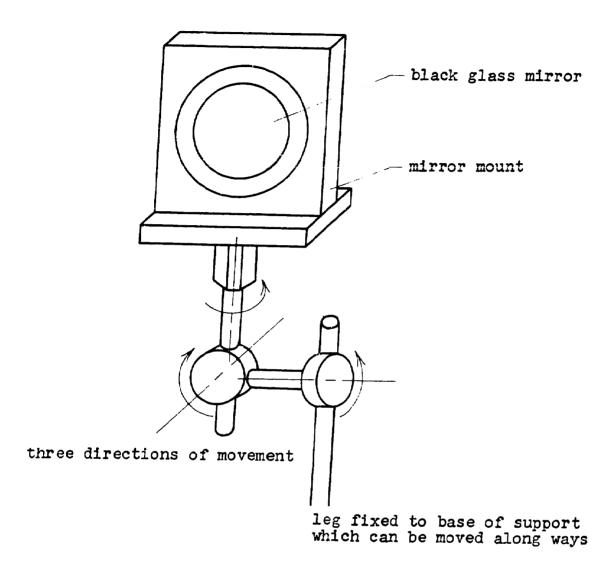


Figure G.7.- Mirror support system.

rotation and the rotating arm system is swung to an angular position such that the focused image of the once reflected incident light strikes the translucent surface of J2. The support system holding mirror A is then adjusted to make the focused incident image coincide with the crossing reference lines on J2, while at the same time, the plane of the mirror is kept at an angle of 45° with respect to the end surface of the incident lens tube. As a result of this alinement of this mirror, the incident light beam is turned 90° from its initial direction while the direction of the reflected beam is kept in the plane of observation.

Mirror B and its associated support system is then also fixed to the stationary optical bench with the plane of the mirror made approximately parallel to the plane of mirror A. The horizontal axis of mirror B is accurately adjusted by swinging the rotating arm system, with the alinement jig J_2 still in place, to an angular position such that the focused image of the twice reflected incident light strikes the vertical reference line on the translucent surface of J_2 . The support system holding mirror B is then adjusted to make the center of the image coincide with the crossing reference lines on J_2 . Since the angular position of the rotating arm will in general not place the intersection of the reference lines on J_2 at the appropriate point for alining the vertical axis of mirror B, that is, turning the reflected light through 90° at mirror B, only partial alinement of

mirror B is possible at this point. The plane in which the twice-reflected beam travels was made to coincide with the plane of observation by simultaneously adjusting the support system for mirror B so that the center of the image of the twice-reflected beam coincides with the intersection of the reference lines on J_2 and moving the rotating arm system through a small range of angles which would bracket to position of the arm which would turn the light 90° at mirror B.

The alinement to this point causes the incident light to be reflected 90° by mirror A while keeping the direction of the reflected beam in the plane of observation. The partial alinement of mirror B insures that the direction of the twice-reflected beam remains in the plane of observation for a range of reflection angles.

Mirror C is placed on the rotating arm system and adjusted so that the plane of the mirror is 45° with respect to the end surface of the collection lens tube by using a 45° square. The rotating arm system is then positioned at approximately $\theta = 90^{\circ}$ and mirror C is moved along the ways of the rotating arm system so as to intercept and reflect the twice-reflected light beam. The direction of the resulting reflected light is made to lie in the plane of observation by moving the rotating arm system and adjusting the horizontal axis of mirror C until the center of the image of the reflected light coincides with the intersection of the reference lines on J_2 . This last adjustment may alter the angle between the surface of mirror C and

again after which the previous alinement also checked. Mirror C and its support system is then secured to the rotating arm system as shown in Figure G.6.

The rotating arm system is positioned at $\theta = 90^{\circ}$ and the final alinement of mirror B is made by the simultaneous movement of the support system for mirror B along the stationary optical bench and slight rotation of the surface of mirror B only about its vertical axis. The center of the image of the thrice-reflected light is compared to the intersection of the reference lines of J_2 when J_2 is first placed at one end of the collecting lens tube and then at the other end. When the center of this image coincides with the intersection of the reference lines of J_2 for both positions of J_2 , the final alinement of mirror B has been made and the entire mirror system is alined.

Since mirror A and mirror C must be moved for the measurement of the scattered-light intensity at various angular positions, stops are fixed to the stationary optical bench and the ways of the rotating arm system at the base of the support systems for these two mirrors which allow them to be replaced to their proper alined positions.

The two lenses, $L_{\bar{5}}$ and L_{4} , and the polarization filter were then placed inside the lens tube. (See Fig. 8.1.) The long focal length lens, $L_{\bar{5}}$, which collects the scattered light, was then adjusted to a distance of 508 mm from the axis of the rotating arm. The square

tube system. Following this, the converging lens L4 was adjusted so as to refocus the incident light at the plane of the square aperture. The square aperture was then adjusted to allow all of the incident light to pass through it.

The alinement of the collection system for the scattered light was then checked by a somewhat more sensitive technique in the following way. The alinement jig shown in Figure G.5 was placed in the burner tube fixture with the surface of its slot normal to the axis of the collection system. A lamp was placed behind the square aperture which in turn formed an image on the slot surface of the alinement jig. The horizontal alinement was checked by comparing the center of the image with the vertically scribed reference line on the slot surface which coincides the axis of rotation. In general, the center of the image was only very slightly misalined with the vertical reference line of the jig and was more accurately alined by very slight adjustment of the vertical slit of the square aperture. This slight misalinement of the square aperture by the first technique is due to the refocused image of the incident light being slightly smaller than the aperture which allows some freedom in adjustment. The vertical alinement was checked by comparing the position of the focused incident image on the jig when the slot surface was normal to the axis of the incident beam with the position of the image of the square aperture when the slot surface was normal to the axis of the

alined and was more accurately alined by very slight adjustment of the horizontal slit of the square aperture. This check was then repeated at several different angles and it was found that once the system was accurately alined at $\theta = 0^{\circ}$ by the second technique, no further adjustment was necessary.

With the alinement jig removed from the burner tube fixture and the rotating arm set at $\theta = 0^{\circ}$, the alinement was again checked to see that no part of the refocused image of the incident was cut off by the square aperture.

The optical interference filter and the multiplier phototube system were then mounted on the rotating arm. The multiplier phototube tube was mounted so that the photocathode surface was approximately 1/2 inch from the plane of the square aperture and the light spot strikes the center of the photocathode surface.

APPENDIX H

Tabulation of Test Conditions for Light-Scattering
Measurements and Experimental Data

The original data and calibrations are on file at the MIT Fuels Research Laboratory. Table H.1 lists the test conditions for which light-scattering measurements were made. The runs numbered 1 through 4 were made for nearly identical flame conditions but at successively greater elevations above the burner tip. The runs numbered 5 and 6 were made for nearly identical flame conditions at the same elevation but for two different wavelengths. Runs 7 through 9 were also made for nearly identical flame conditions at successively greater elevations above the burner tip but with a lower fuel-air ratio and smaller elevation intervals. Run number 10 was made with a Hefner candle burning amyl acetate.

Table H.2 gives the experimental light-scattering data for each run. The measured voltage, which is proportional to the scattered-light intensity for the perpendicular component, $e_1(\theta)$, and for the parallel component, $e_2(\theta)$, is given for the various angular positions, θ . The measured voltage due to the thrice reflected incident light, $e'_{0,2}$, and the voltage due to the parallel component of the scattered-light intensity for $\theta = 60^{\circ}$, $e'_{2}(60)$, for the same apparatus adjustment, are also listed for each run.

	Amplification of signal for incident light measurements	. 225 255	11 11 11	16 16 16	16
	High voltage to multi- plier phototube for incident light meas- urements, volts	750 750 750	750	750 750 750	750
MENTS.	Amplification of signal for scattering measurements	25.88	8 27,	76 16 16	or 90
MEASURE	High voltage to multiplier phototube for scattering meas- urements, volts	1000 950 950	95 95	0001	1000
ATTER INC	Approximate flame height, inches	3.15 3.15 3.15	3.8	2.30	
LIGHT-SCATTERING MEASUREMENTS.	Elevation of above burner, above burner, inches	1.25		1.03	1.00
CONDITIONS FOR	f/a times stoichiometric	2.39 2.38 32.36	2.45 2.44	2.16 2.16 2.16	Diffu- sion flame
TEST COND	Airflow rate, 201 × mim/slom-m3	4.70 4.70 4.70 4.70	4.70 4.70	9.4 9.4 9.60	
H.l	Benzene flow rate, $ m gm-mole/min imes 10^5$	3.16 3.15 3.12 3.07	3.23	8.88 8.89 8.89	
TABLE	Å , A	4358 4358 4358 4358	4358 5461	4358 4358 4358	4358
	Тэич	Benzene Benzene Benzene Benzene	Benzene Benzene	Benzene Benzene Benzene	Amyl acetate
	uny	1 2 5 4	62	- 86	07

TABLE H.2.- EXPERIMENTAL LIGHT-SCAPTERING DATA

Run 4	e ₁ (θ) e ₂ (θ) 580. mv 468. mv 450. 294. 345. 219. 345. 161. 269. 75. 216. 30.2 170. 10.3 141. 4.8 122. 8.3 109. 16.9 95. 28.5 88. 40.0 87. 47.5 88. 65. mv e' ₂ (60) = 25.0 mv
Run 3	e ₁ (θ) e ₂ (θ) 605. mv 500. mv 475. 220. 249. 375. 288. 500. 88. 530. 88. 530. 36. 199. 12.3 166. 6.3 10.9 121. 20.8 118. 34.8 107. 50. 109. 59.5 109. 69. 109. 69. 109. 69. 109. 69.
Run 2	$e_{1}(\theta)$ $e_{2}(\theta)$ 560 mv $t_{2}0$ mv $t_{5}0$ mv $t_{5}0$. 310
Run 1 $e_1(\theta)$	681. mv 570. mv 575. 400. 524. 320. 400. 131. 320. 400. 131. 340. 57.5 298. 20.8 266. 10. 244. 17. 230. 222. 66.8 219. 102. 223. 223. 122. 220. 147. 238. 196. 196. e', 2 = 223. mv e', (60) = 18.7 mv
θ	563 57 57 60 60 80 90 90 90 90 90 90 90 90 90 90 90 90 90

TABLE H.2.- EXPERIMENTAL LIGHT-SCATTERING DATA - Continued.

	1	Run 5	Rı	ın 6
9 40 45 50 60 70 80 90 100 110 120 130 135 140	e ₁ (θ) 621. my 525. 478. 435. 362. 308. 276. 246. 223. 212. 208. 210. 218. 224.	e ₂ (θ) 505. mv 350. 282. 218. 112. 49. 17.5 8.4 15.1 32.5 59.5 94. 111. 132. 175.	Pice 1 (0) 143. mv 127. 116. 109. 93.5 81. 71.8 64.3 60. 58. 59. 61. 62.5 64.5	e ₂ (θ) 108. mv 76.5 61. 47. 25.8 11.5 4. 2. 3.5 8.7 16. 26.7 32.5
	e' _{0,2} = 249 e' ₂ (60) = 2		e' _{0,2} = 41. e' ₂ (60) = 10	

= 24.2 mv

e₂(60)

TABLE H.2- EXPERIMENTAL LIGHT-SCATTERING DATA - Concluded.

θ	e _l (в)	e ₂ (θ)
30 40 45 50 60 70 80 90 100 110 120 130 135 140	575. mv 530. 500. 450. 428. 400. 360. 325. 295. 270. 245. 230. 220. 210.	500. mv 380. 330. 270. 150. 74. 25. 11. 22. 47. 86. 125. 140. 155. 185.
-	- ·	±0)•

A Sample Calculation Showing the Effect That a Large
Number of Small Particles Mixed With a Relatively
Small Number of Large Particles has on the
Light-Scattering Patterns for Two
Different Wavelengths

A comparison of the light-scattering data for run 5, which is for $\lambda = 4358$ Å, to run 6 which is for $\lambda = 5461$ Å, shows that the ratio of the apparent values of x for these respective runs is about 1.11 when the values of x are deduced by comparing the data with theory for uniform size particles. If the soot particles are nearly uniform in size, one would expect to have found the ratio of the two values of x equal to the inverse ratio of the respective wavelengths, that is, 5461/4358 = 1.25, since the test conditions for these two runs are very nearly identical.

It is the object here to see if this apparent discrepancy between the data of run 5 and run 6 could be due to a soot particle system which consists of a relatively large number of small particles, say 250 Å in diameter, and a relatively small number of large particles or agglomerates. The choice of a diameter of 250Å corresponds to the ultimate soot particle size shown in the electron micrograph, Figure 10.1.

The apparent value of x, deduced from the data of run 5 when m = 1.71 - 0.76 i, is 1.00. For the supposed particle system just mentioned, this value of x = 1.00 is expected to lie between the

corresponding value of x for the 250 Å particles and the equivalent value of x for the agglomerates. The value of x for the agglomerates was taken as 1.25 for this example when $\lambda = 4358$ Å.

The apparent value of x deduced from the data of run 6 when m=1.79-0.79 i is 0.90. Since the size parameter x was already taken as 1.25 for the case when $\lambda=4358$ Å, the value of x for the agglomerates when $\lambda=5461$ Å is necessarily 1.00 as determined by the inverse ratio of the respective wavelengths. The choice of these values of x is due in part to the fact that the scattering functions for these values are tabulated in appendix D, so no interpolation is necessary.

The corresponding values of x for the small particles which have a diameter of 250 Å are x=0.180 for $\lambda=4358$ Å and x=0.144 for $\lambda=5461$ Å. Since it is the purpose here to show a rather gross effect, the values of the size parameter for the small particles were taken as x=0.20 for $\lambda=4358$ Å and x=0.15 for $\lambda=5461$ Å, to avoid interpolation between the tabulated light-scattering functions.

Let β be defined as the number fraction of soot particles which are agglomerates and $(1-\beta)$ the number fraction of ultimate soot particles. The effective light-scattering functions for this system of soot particles when $\lambda = 4358$ Å are

$$i_1(\lambda = 4358, \theta) = i_1(x = 1.25, \theta)\beta + i_1(x = 0.20, \theta)(1 - \beta)$$
 I.1

and

$$i_2(\lambda = 4358, \theta) = i_2(x = 1.25, \theta)\beta + i_2(x = 0.20, \theta)(1 - \beta)$$
 I.2

where $i_1(x=1.25,\theta)$, for example, is the perpendicular light-scattering function for x=1.25 at angle θ . The corresponding effective light-scattering functions for this system when $\lambda=5461$ ${\stackrel{\circ}{A}}$ are

$$i_1(\lambda = 5461, \theta) = i_1(x = 1.00, \theta)\beta + i_1(x = 0.15, \theta)(1 - \beta)$$
 I.3

and

$$i_2(\lambda = 5461, \theta) = i_2(x = 1.00, \theta)\beta + i_2(x = 0.15, \theta)(1 - \beta)$$
 I.4

The values of $i_1(x=1.25,\theta)$ to be used in equation I.1, as tabulated in appendix D, range between about 1.0 and 0.15 and the values of $i_1(x=0.20,\theta)$ are approximately 2.0 \times 10⁻⁵. Therefore, in order for the two different particles sizes to each significantly affect the light-scattering data, the value of β should be near the order of 10⁻⁴ or 10⁻⁵. Larger values of β will cause the agglomerated particles to dominate the light-scattering results, which in turn would lead to different scattering patterns for each wavelength that correspond to two values of x, the ratio of which is equal to the reciprocal of the wavelengths. This is contrary to experiment. On the other hand, much smaller values of β would produce a nearly Rayleigh-type scattering pattern, which is also contrary to the data from run 5 and run 6. With this in mind, a value of $\beta = 5 \times 10^{-5}$ was chosen. The light-scattering functions used for the case where $\lambda = 4358$ were taken from appendix D for m = 1.71 - 0.76 i. The

corresponding functions for the case where $\lambda = 5461 \, \text{Å}$ are based on m = 1.79 - 0.79 i.

Figure I.1 shows the results of these calculations based on equations I.1 and I.2 for the case of $\lambda = 4358$ Å. The symbols in this figure represent the calculated scattering pattern for the mixed particle system. This pattern was compared to the theoretical scattering curves for uniform size particles for various values of x at m = 1.71 - 0.76 i to find the best fit for the perpendicular component. The best curve fit was found for x of about 0.95 and is drawn on the plot along with the symbols.

Figure I.2 presents the results based on equations I.3 and I.4 for the case of $\lambda = 5461$ Å, but the same supposed particle system as before. The resulting scattering pattern, again indicated with symbols, is compared to curves for various values of x for m = 1.79 - 0.79 i. The apparent value of x is about 0.90 or slightly less for this case.

For this supposed system of soot particles, the ratio of the size parameters for the case of $\lambda = 4358$ Å to the case of $\lambda = 5461$ Å is only slightly larger than 0.95/0.90 = 1.06, rather than the inverse ratio of the two wavelengths, 5461/4358 or 1.25. This shows the same effect as was indicated by the experimental data from run 5 and run 6 and therefore offers a possible reason for this seeming discrepancy. A different value of β or assumed value of the agglomerate size would certainly change the ratio of the values of x for the two wavelengths, so that the value of 1.11 for this ratio, as found in comparing run 5 and run 6, is not surprising.

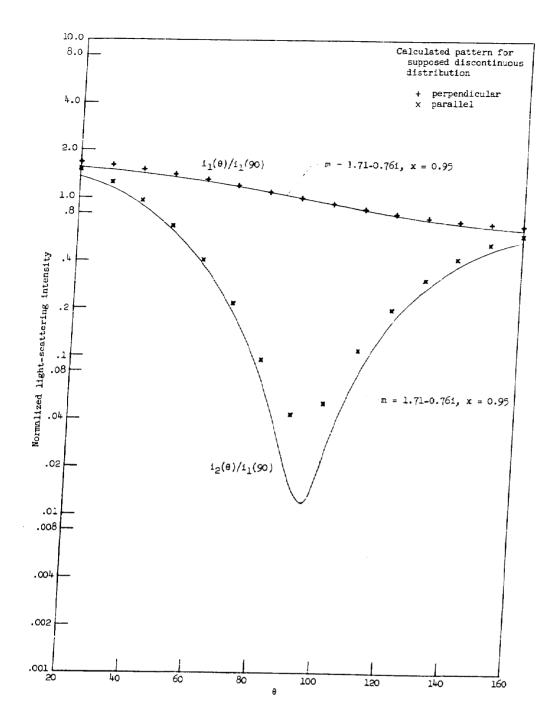


Figure I.1.- Comparison of the calculated light-scattering pattern for a supposed discontinuous particle distribution system, which contains a large fraction of small particles mixed with a small fraction of large particles, to the theoretical scattering pattern for a uniform particle size system where $\lambda = 4358$ Å and m = 1.71 - 0.76 1.

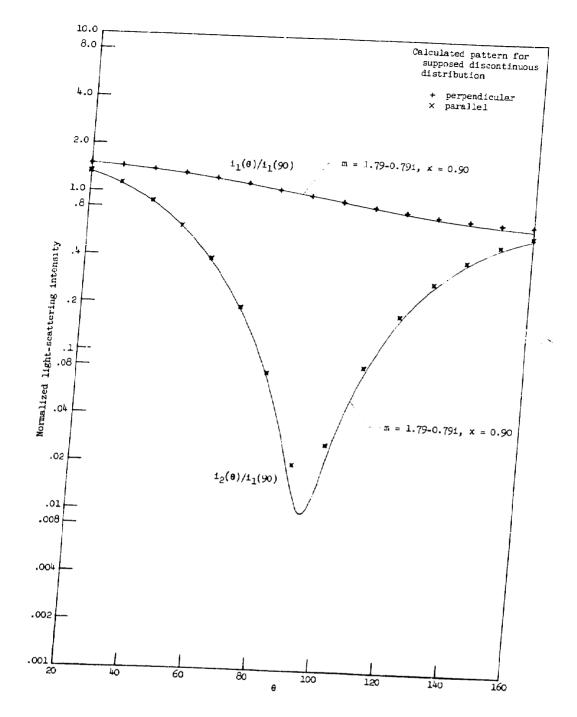


Figure I.2.- Comparison of the calculated light-scattering pattern for the same supposed discontinuous particle distribution as used for Figure I.1 to the theoretical scattering pattern for a uniform particle size system where $\lambda = 5461~\text{Å}$ and m = 1.79 - 0.79 i.

It should be noted that this supposed particle system is an over simplified distribution which consists of a very large number of small particles, each of the same size, mixed with a relatively few larger particles, each of the same size. In the actual situation, one might expect to find a smooth distribution about a small mean particle size which corresponds to the ultimate soot particles and a second smooth distribution about a much larger mean particle size which corresponds to the agglomerates. It is felt that this consideration would not change the general effect already shown.

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BIOGRAPHICAL NOTE

The author was born February 19, 1932 in Lansing, Michigan. He received his elementary and high school education in the public school system of Lansing and entered Michigan State College in March of 1950. He received the degree Bachelor of Science in Chemical Engineering in June of 1954 and the degree Master of Science in Chemical Engineering in the summer of 1955. In June of 1955, the author entered active duty with the United States Air Force and was assigned to work at the former National Advisory Committee for Aeronautics at Langley Field, Virginia. After completing the tour of duty in June of 1957, he entered the Graduate School of the Massachusetts Institute of Technology and received the degree Master of Science in Chemical Engineering Practice in June of 1958. The following month he returned to the National Advisory Committee for Aeronautics/National Aeronautics and Space Administration at the Langley Research Center, Langley Field, Virginia and has been employed there since. For intervals during the past few years, the author was granted graduate study leave to attend the Massachusetts Institute of Technology. He is presently engaged in hypersonic aerodynamic research at NASA, Langley Field, Virginia.

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Erickson, W. D.: Study of Pressure Distributions on Simple Sharp-Nosed Models at Mach Numbers From 16 to 18 in Helium Flow. NACA TN 4113. October 1957.

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