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STOCK REMOVAL RATES IN INTERNAL GRAINDING: A MODEL OF THE PROCESS

by

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Submitted in Partial Fulfillment of the Requirements for the Degree of Master of Science at the Massachusetts Institute of Technology

August 1966

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STOCK REMOVAL RATES **IN INTERNAL GRINDING: A** MODEL OF THE **PROCESS**

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Richard P. Lindsay

Submitted to the Department of Mechanical Engineering on August 22, **1966** in partial fulfillment of the requirements for the degree of Master of Science

A model of the grinding process is given in which the cutting and sliding regions of a worn grain are separated from each other. It is assumed that each of these processes can be represented **by** use of constant stresses in the normal and tangential directions resulting in four unknown constants. Using four experimental data points and the relation between the tangential and normal force, the four constants are determined. The theory then is seen to fit a total of fifty-nine experimental results.

The constants thus determined are seen to be compatable with results obtained with single-grain sliding tests.

Thesis Supervisor: Nathan H. Cook Professor of Mechanical Engineering **CONTENTS**

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INTRODUCTION

It has been found that the cutting ability of a grinding wheel diminishes with time from a high initial value to a lower relatively-constant rate, when the constant applied force between the wheel and the work is low enough to prevent gross breakdown of the wheel. Microscopic examination of the wheel surface, at various stages in the process, reveal that the flats worn on the wheel grains continually grow with time. Thus, if the area of these flats in contact with the work is measured and the normal force of grinding is known, the relation between cutting ability of the wheel and actual applied normal stress may be determined empirical **ly.**

When this was done it was found that the removal ability, $\overline{\textbf{w}}$, when divided by applied force, F_n, varied linearly with the applied normal stress, σ

The purpose of this thesis then is to derive theoretically the above relationship or to show how such a linear approximation can represent the true relationship.

Definition of Terms Used

THATE.

 k_1 , k_2 , \equiv Constants

 $\overline{5}$

A. Models of the Process

I. Constant-Stress

Consider Figure **I:** It is assumed that the stock-removal process can be divided into two zones: the cutting region where the chip is produced and the sliding region under the worn grain flat. Then the following assumptions will be made:

a.- Chip formation is continuous and a zone of rigidplasticity exists. Then the forces of cutting may be written.

$$
F_V = k_1 * bt \cot \emptyset
$$

\n
$$
F_H = k_3 * bt
$$
 (1)

that is the forces In the cutting region are proportional to the depth and width of cut, t.

b.- The sliding forces are assumed as:
\n
$$
\frac{\pi}{4} \frac{2}{4}
$$
\n
$$
F_V = k_2 \frac{2}{4}
$$
\n
$$
F_H = k_4 \frac{2}{4}
$$
\n(2)

That Is, constant stresses exist under the grain flat and are independent of the depth of cut.

Thus the normal force may be written as:

$$
F_N = F \text{ cutting} + F \text{sliding}
$$

$$
F_{Nq} = k_1 * \text{bt} \cot \emptyset + k_2 * \frac{\pi a^2}{4}
$$
 (3)

Letting
$$
k_1 = k_1 * \cot \emptyset
$$

and
$$
k_2 = \frac{\pi}{4} k_2^*
$$

(3) becomes:

$$
F_{Ng} = k_1 bt + k_2 \bar{a}^2
$$
 (4a)

or

$$
i = \frac{1}{F_{Ng}} \left[k_1 b + k_2 \bar{a}^2 \right]
$$
 (4b)

The forces in the tangential direction may be written as:

$$
F_{sg} = k_3 * bt + k_4 * \frac{\pi a^2}{4}
$$
 (5)

Letting
$$
k_3 = k_3
$$

and $k_4 = k_4 * (\frac{\pi}{4})$
and defining $F_{sg} = \mu F_{Ng}$

(5) becomes:

$$
u F_{Ng} = k_3 bt + k_4 \frac{a^2}{a}
$$
 (6a)

or
$$
M = \frac{1}{F_{Ng}} [k_3 b + k_4 \bar{a}^2]
$$
 (6b)

Note similarity between (4b) and **(6b)** where **only/4** and constants vary.

2. Elastic Rebound:

Again considering Figure I, assume the following:

a.- At the cutting region:

$$
F_V = k_1^* \text{ bt cot } \emptyset
$$

$$
F_H = k_3 * bt
$$

Same as the first model

b.- Under the grain flat:

$$
F_V = k_2 * \overline{a} \text{ bt}
$$

$$
F_H = k_4 * \overline{a} \text{ bt}
$$

it being now assumed that the forces imposed on the grain flat are proportional to the length, width and depth of cut. That is, these forces are proportional to the elastic rebound of the distorted metal in the shear zone. Then, as before, writing the normal force as:

$$
F_{Ng} = k_1 * bt \cot \emptyset + k_2 * \bar{a} bt
$$

Letting $k_1 = k_1 * \cot \emptyset$ and $k_2 = k_2$ ^{*}

Then $F_{\text{Nq}} = b + (k_1 + k_2 \bar{a})$

or
$$
I = \frac{bt}{F_{Nq}} (k_j + k_2 \bar{a})
$$

This model was found to produce a tensile constant, **k**₁ meaning that the normal stress under the grain is too large under this assumption. Appendix **8** will follow the development of this assumption.

The next sections will be concerned with the determining the quantities necessary for the deriving of working formulas from the relations 4b and **6b.**

B. Analytical Quantities

I. Random Array of Grains

In order to establish the quantities G , grain density of the wheel and **S,** effective grain spacing around the wheel periphery in a plane, the grain array of Figure 2 is used.

The wheel surface is assumed to have the array shown where: "M" is the distance between grains on any ring, **k:** "L" is the distance between rings; a is the diameter of the flat worn onto the grain at any time.

2. Spacing of Effective Grains, 5:

Consider a line drawn from the origin, **"0",** in a radial direction. The average distance travelled along such a line is when the probability of intersecting another grain flat equals the probability of missing anofher flat (i.e. the escaping probability equals **0.5).** It may be shown (Appendix **1)** that the average distance so travelled is given **by** (Appendix formula **1.2.6):**

$$
\overline{S} = \frac{L^2}{a}
$$
 (7)

which states that as a becomes larger (as the real area grows) the effective grain spacing becomes smaller since L is a constant, which will now be determined.

Figures 4 and 5, while they present a picture of the cutting surface of the wheel, also give the distance between effective grains, because if the grain has become worn it was in contact with

the work and so was an "effective grain". Thus an average of the distance between flats from Figure **5** will be a measure of L, when any one flat is considered as the origin. From Figure **5** measurements, for the **60** grit wheel, the average spacing is 0.018 inch. Since the average size of a **60** grit size stone is **0.016** inch, assume the spacing may be given as:

$$
L \approx d \tag{8}
$$

Then **(7)** becomes

$$
\bar{5} = \frac{d^2}{a}
$$
 (9)

Now the grain density, (**,** may be determined.

3. Grain Density,

Referring to Figure 2, in the first ring, the number of grains per area is:

$$
G = \frac{\text{number of grains}}{\text{area}}
$$

$$
G = \frac{1 + (1/2) (6)}{\pi L^2}
$$

$$
C = \frac{1.26}{L^2}
$$
 (10)

or using **(8):**

$$
\mathbf{C} = \frac{1.26}{d^2} \tag{11}
$$

Thus the grain density increases with smaller grain sizes.

i)

4. Length of Contact, L_c:

The assumption that the workpiece is infinitely stiff, relative to the wheel is made to begin. It is also found that due to the assumed array, the number of grains in contact with the work Increases linearly with the length of contact. Thus, as the normal force is increased more and more grains come into contact with the work and if it assumed that the bonding-agent holding the grains in the wheel acts as a spring, then the applied forced, F_{N} , is being resisted **by** more and more spring-loaded grains. Hence the deflection of the wheel in the direction parallel to F_N is not linear with the applied force, F_N, which it would be if only the original number of grains were present for all deflections.

From the assumption that the slope of the force-deflection curve is always a multiple of the original slope (the slope increasing as more and more grains come into contact at higher forces) the following relationshipsis derived (see Appendix 2):

$$
L_c = \left[\frac{d \text{DD}_w F_v}{m \cdot k_g (D_w - D)}\right]^{1/3}
$$

Where:

 M_{60} grit wheel $= 0.2$ dimensionless $M_{90} = 0.167$ k_{60} = 1.025 (10⁵) $\frac{LB}{INCH}$ LB k_{QQ} **= 0.266 (10) INCH**

 $||$

5. Diameter of Worn Grain Flat (a):

The total instantaneous real area in contact, A_{RL} , is merely the sum of the flats worn onto the number of grains in contact. Thus, A_{RL} may be written as:

$$
A_{RL} = (No. grains in contact) \left(\frac{area}{grain}\right)
$$
\n
$$
A_{RL} = C \cdot L_c \cdot W \left(\frac{\pi \bar{a}^2}{4}\right) \qquad \frac{grains}{inch^2} \cdot inch \cdot inch \cdot \frac{inch^2}{grain}
$$

Considering the worn flat to be circular. Using (11) and (12) , the above becomes:

$$
A_{RL} = \left(\frac{1.26}{d^{2}}\right) W \left(\frac{\pi a^{2}}{4}\right) \left[\frac{d_{DD_{w}}F_{x}}{mk_{g}(D_{w}-D)}\right]^{1/3}
$$

$$
\left(\frac{\bar{a}}{d}\right)^{2} = \frac{4 A_{RL}}{1.26\pi W} \left[\frac{m k_{g}(D_{w}-D)}{d_{DD_{w}}F_{x}}\right]^{1/3}
$$

$$
\left(\frac{\bar{a}}{d}\right) = \left(\frac{A_{RL}}{W}\right)^{1/2} \left[\frac{m k_{g}(D_{w}-D)}{d_{DD_{w}}F_{x}}\right]^{1/6}
$$

or

$$
\bar{a} = \left(\frac{A_{\text{RL}}}{W}\right)^{1/2} d^{5/6} \left[\frac{mk_9 (D_w - D)}{DD_w F_N}\right]^{1/6}
$$
 (13b)

6. Shape of Cut in Internal Grinding

a. Condition for Scalloped-Shaped Cut:

Consider Figure **6** wherein Is shown a plane section of the wheel, with grits spaced on the plane periphery, S distance apart. The normal force will cause a penetration into the work as shown but a grain may or may not be present at the interference region, **Lc'**

The shape of the chip produced, as motions **M** and V are introduced, will be scalloped as shown if the following condition exists:

$$
\frac{N}{\sqrt[n]{\sigma}} < \frac{S}{\sqrt[n]{\sqrt[3]{\sigma}}}
$$
\n
$$
\frac{N}{\sqrt[3]{\sigma}} > \frac{N}{\sqrt[3]{\sigma}}
$$
\n
$$
\frac{N}{\sqrt[3]{\sigma}} < \frac{S}{\sqrt[3]{\sigma}}
$$
\n
$$
\frac{N}{\sqrt[3]{\sigma}} <
$$

Experimental data will prove this condition to be valid. Also general ly, from the geometry of Figure **6:**

$$
N = 2\sqrt{10}
$$
 (15)

also

or

$$
A_{\text{cut}} = \frac{3}{2} t^{3/2} D^{1/2} \tag{16}
$$

b. Determination of Depth of Cut t:

Considering continuity of the stock-removal process from Figure **6,** the area removal per unit time may be written two ways:

I. from the fact that \overline{v} is the rate of radius growth, then the area removed per unit time is:

$$
\overline{\sigma}:\pi\circ_{_{\!W}}\!\!)
$$

2. Also the amount of material removed must be equal to the number of cuts per unit time multiplied by the area removed per cut, or:

$$
\left(\frac{\nu}{5}\right)\cdot A\text{ cut}
$$

Thus, for continiuty these two rates must be equal:

$$
\mathbf{T} \, \mathbb{D}_{\mathsf{W}} \, \overline{\mathbf{A}} \qquad = \left(\frac{\mathsf{V}}{\overline{\mathsf{S}}} \right) \cdot \mathsf{A} \, \text{cut}
$$

Using (9) and (16):

$$
\pi D_w \overline{w} = \frac{\nabla \overline{a}}{d^2} \left(\frac{3}{2} \, \mathbf{t}^{3/2} \, D^{1/2} \right)
$$

$$
t = \left[\frac{2\pi}{3} \frac{D_w}{D'^2} \frac{\overline{d}^2}{\overline{a}}\right]^{2/3}
$$
 (17)

Using (13a) and dividing both sides by "d":

$$
\left(\frac{t}{d}\right) = \left\{\frac{2\pi D_w \overline{\sigma} d}{3 D^{v_2} V d^{3/z}} \left(\frac{w}{A_{RL}}\right)^{v_2} \left[\frac{dDD_w F_N}{mR_g (D_w D)}\right]^{v_2}\right\}^{2/3}
$$

yielding:

$$
\left(\frac{t}{d}\right) = 1.635 \left(\frac{\overline{v}}{V}\right)^{2/3} \left(\frac{w}{A_{R1}}\right)^{1/3} \left[\frac{F_{N} D_{w}^{7}}{m f_{q} (D_{w} - D)(D_{d})^{2}}\right]^{1/9}
$$

 $\ddot{\kappa}$

C. Determination of the Width of Cut **(b):**

Consider Figure **7** where **b** is the width of cut of a ploughed groove made by a grain with flat, a. From continuity the volume removed per unit time may be written two ways:

1. Considering $\overline{\omega}$, the rate of hole radius growth, the volume removed is:

$$
\overline{\sigma}:\pi\circ_{\scriptscriptstyle W})\ \ (\text{w})
$$

2. Also the total volume removed must be equal to the volume removed per cut times the number of cuts per unit time:

$$
WV \bullet (A_{\text{cut}} b) \quad (\frac{\text{CUTTERS}}{\text{SECOND}}) \cdot (\frac{\text{VOLUME}}{\text{CUT}})
$$

These must be equal. Using **(11)** and **(16):**

$$
b = \frac{\sqrt{v} \pi D_w w}{w V \left(\frac{1.26}{d^2}\right) \left[\frac{3}{2} t^{3/z} D^{v_2}\right]}
$$

Using **(17):**

b =
$$
\frac{\pi}{1.26} \left(\frac{2}{3} \right) \frac{\pi D_w d^2}{V D^{\frac{1}{2}} \left\{ \left[\frac{2}{3} \frac{\pi D_w \sqrt{v} d^2}{V \cdot v} \right\}^{2/3} \right\}^{3/2}}
$$

b = \bar{a} /1.26
b = 0.8 \bar{a} (19)

7. Force per Grain, F_{Nq}

Define

$$
F_{Ng} \equiv \frac{F_N}{No. of grains in contact}
$$

$$
F_{Ng} = \frac{F_N}{(G \cup (L_c) (W)}}
$$

Using (11) and (12) :

$$
\left(\frac{F_{Ng}}{F_{N}}\right) = \frac{d^{2}}{1.26 \text{W}} \left[\frac{m \cancel{R}_{q} (D_{W}-D)}{d D D_{W} F_{N}}\right]^{1/3}
$$

or

$$
\left(\frac{F_{Ng}}{F_N}\right) = \frac{0.294}{W} \left[\frac{m\cancel{k}_g (D_w - D)d^6}{DD_w \cancel{F_N} d}\right]^{93} \tag{20}
$$

C. Mathematical Development of Constant Stress Model

I. General Solution in Normal Direction:

Since all the necessary parameters are now described mathematically, the development of the theory may be carried out. Copying formula 4b below:

$$
1 = \frac{1}{F_{\text{ng}}} \left[k_1 bt + k_2 \bar{\alpha}^2 \right]
$$
 (4b)

Using **(19),** (20), (13a) and **(18)** (see Appendix **3)** it is found that

$$
1=\frac{4.45 k_1 w^{5/6}}{F_n^{5/9} \sigma^{1/6}}\left[\frac{D_w^{17}}{m^5 k_g^5 (D_w \cdot D)^5 (D_d)}\right] \left(\frac{\pi}{V}\right)^{2/3}+3.4 \frac{k_2}{\sigma}
$$

2. General Sotution in Tangential Direction:

Comparing the form of 4b and **6b,** noting the only changes are and the constants, the solution in the tangential direction may be written from (21) as:

$$
\mu = \frac{4.45 k_3 w^{5/6}}{F_0^{5/9} \sigma^{1/6}} \left[\frac{D_w^{17}}{m^5 k_y^5 (D_w - D)^5 (dD)} \right] \cdot \left(\frac{\pi}{V} \right)^{2/3} + 3.4 \frac{k_4}{\sigma} \tag{22}
$$

3. General Solution for Rate of Radius Growth,
$$
\overline{w}
$$
 :
\nReturning now to (21) and solving for $\left(\frac{\overline{w}}{\sqrt{w}}\right)^{2/3}$ yields:

$$
\left(\frac{\overline{v}}{v}\right)^{2/3} = \left\{1 - 3.4 \frac{k_2}{\sigma}\right\} \left\{\frac{\sigma^{1/6} F_{\mu}^{5/9}}{4.45 \ k, w^{5/6}} \left[\frac{d D (m k_4 [D_w - D])}{D_w^{17}}\right]^{1/8}\right\}
$$

Revising this as per Appendix 4, yields:

$$
\left(\frac{\overline{v}}{V}\right) = \frac{0.106}{k^{3/2}w^{5/4} D_{w}^{1/12} F_{w}^{5/12} A_{RL}^{1/4}} \left(F_{w} - 3.4 k_{z} A_{RL}\right)^{3/2}
$$

where the last parenthesis is the potential force available for cutting, F_N , minus the force being absorbed under the flat. Thus, as the grain flats grow the potential cutting force is diminished by the force under the flat, causing the cutting rate \overline{iv} to diminish.

From (24) it may be noted that

$$
\left(\frac{\overline{w}}{v}\right) = 0 \qquad \text{when } \left\{F_{N} - 3.4 F_{z} A_{R} u\right\} = 0
$$

or
$$
\sigma_{0} = 3.4 F_{z}
$$
 (25)

Thus σ , the stress at which the wheel will cease to cut, is **3.4** times the constant normal stress assumed to exist under the grain flat.

Summarizing then: a model has been assumed for which the cutting forces are proportional to the depth and width of cut and the sliding forces are independent of the depth of cut but proportional to the worn flat area on the grains.

Using this model, formulas have been derived which contain various constants. Using four experimental results, these constants will be evaluated and the resultant formulas used to predict other experimental results. Specifically, formula (21) will be solved using two data points for each wheel, **60** and **90** grit. Since **If** and **G** appear in (21) it is necessary to have grinding data for which \overline{v} and σ are known. Using the evaluated constants, formula (24) may be used to predict other results.

It should be remembered that the scallop-shaped cut assumption remains to be proven. This will be done when the experimental procedure and data have been given.

D. Experimental Quantities

1. Normal and Tangential Forces:

The test grinding was done under constant normal force, plunge-cut conditions. Thus the normal force was an adjustable and independent variable. Tangential forces were approximated from torque measurements. Generally, for a sharp wheel, $\overline{\mathbf{f}}$ **2** 0.5; and for a wheel which has become dull, $\overline{\mathbf{f}}$ **2** 0.3. **FN FN** Figure **8** illustrates these forces.

2. Total Instantaneous Real Area in Contact:

The used wheel was mounted on the bed of a 175-power microscope and the wheel surface was brought into focus. **A** Polaroid picture was taken at this location. The bed was moved axially the width of the photograph, and another picture was taken at this new location. In this manner a group of photos were obtained, which, when mounted together consecutively, gave a picture of the wheel surface at **175** magnification. Figure **9** illustrates this method; Figure 4 is a reproduction of one such photo and Figure **5** Is a reproduction of three such strips, (at a smaller scale of reproduction).

However, the length of the resulting photo graphic strip, " ℓ " in Figure **5,** is not necessarily the length of the zone of wheelwork contact under the normal force. Hence this length, L_c, must be determined. **A** workpiece was finished to a smooth surface **by** polishing and a wheel was placed in the hole and the normal forces used in grinding, were placed on the wheel-work system (all on the

test machine). Now the wheel was moved **0.010** Inch axially causing the wheel to scratch the work surface. The average width of the "scratch band" was then taken as the width of contact, L_{c} . See Figure **10,** a, **b.**

Using two of these data points and the deflection hypothesis **6f** this report the formula for the length of contact has been found and is repeated here:

$$
L_c = \left[\frac{d \text{D}D_w F_w}{m + k_g (D_w - D)}\right]^{1/3}
$$
 (12, repeated)

or

$$
L_c = \left[\frac{d F_{\rm N}}{m k_{\rm T} \left(\frac{1}{D} - \frac{1}{D_{\rm W}}\right)}\right]^{1/3}
$$

The theoretical and experimental results are given in Figure **18D.**

Therefore, from the photographic strips (as Figure **5)** of the wheel surface, the area in the strip may be measured using a planimeter. Then the real area in the length of contact, assuming the picture strip is representative of any strip on the wheel, is:

$$
A_{\rm RL} = \frac{L_{\rm c}}{L} A_{\rm RP}
$$

When the troichoidal equations of motion of the wheel are considered, the length of contact L_c is changed less than 1%, therefore the above manner of measuring A_{DI} was used.

3. Instantaneous Rate of Radius Growth $\overline{\textbf{U}}$:

The rate of growth of the work diameter was measured as shown in Figure II. A work-riding finger probe, spring-loaded against the work surface was used and its motion was sensed with a Microtrol Electronic gage **(I** Division **= 10** x **10** maximum sensitivity). The motion of the Microtrol gage was fed into a recorder to obtain a tape of the position of the probe versus time. Then, as Figure II Illustrates, the slope of the position versus time plot, is $\overline{\omega}$. The system was sufficiently sensitive to Microinches allow **IT~** measurements to within **+ 5** Second

Tables **I,** 2 and **3** gives the experimental data obtained. Thus the experimental data was collected, and values of σ , for various \overline{N} rates, were obtained.

The theory hinges on the assumption of the scallop-shaped cut and now that the experimental data has been introduced, this assumption will be proven.

4. Proof of Scallop-Shaped Cut Assumption

The condition for scallop-shaped cuts is copied for convenience:

$$
\frac{N}{\sqrt{N}} > \frac{N}{\overline{S}}
$$
 (14, repeated)

Using already derived relations (14) becomes (see Appendix **5):**

(1) HAHN , ROBERT **S.** " **ON** THE **NATURE** OF THE **GRINDING PROCESS"** PERGAMON PRESS **,, LONDON** , **1963.**

$$
\frac{dV}{V} > 2.56 \left[\frac{(DD_w)^2 m k_q (D_w - D)}{F_w d^2} \right] \cdot \left(\frac{A_{RL} \overline{dV}}{W V} \right)^{1/3}
$$

From the experiments:

Case I: When the wheel is cutting slowly:

$$
\overline{\text{10}} = 75 \times 10^{-6} \frac{\text{inch}}{\text{sec}^2} \text{ A}_{\text{RL}} = 728 \times 10^{-6} \text{ inch}^2
$$

\n
$$
d_{60} = 0.016 \text{ inch}; \quad D = 1.87 \text{ inch}; \quad D_W = 2.37 \text{ inch}
$$

\n
$$
M_{60} = 0.2; \quad k_{60} = 1.025 \text{ (10}^5) \frac{\text{LB}}{\text{NCH}}; \quad F_N = 15 \text{ LB}
$$

\n
$$
W = 0.250 \text{ inch}; \quad V = 1430 \frac{\text{inch}}{\text{sec}}; \quad \text{10} = 224 \frac{\text{inch}}{\text{sec}}
$$

 U_{α}

Into **(25):**

$$
\frac{224}{1430} > 2.56 \left[\frac{(1.87[2.37])^{2} 0.2 (1.025 \times 10^{5})(2.37 - 1.87)}{15 (0.016)^{7}} \right]^{19}
$$

$$
\left[\frac{(728 \times 10^{-6})(75 \times 10^{-6})}{(\frac{1}{4})} \right]^{1/3}
$$

Yields:

 0.156 > 0.128

Case 2: When wheel is cutting at a rapid rate
\n
$$
\overline{or} = 440 \times 10^{-6} \frac{\text{inch}}{\text{sec; A}_{\text{RL}}} = 126 \times 10^{-6} \text{ inch}^2
$$

All other values the same

$$
0.156 > 2.56 \left[\frac{(1.87[2.37])^{2} 0.2 (1.025 \times 10^{5})(2.37 - 1.87)}{15 (0.016)^{7}} \right]^{1/9} \cdot \left[\frac{(440 \times 10^{-6}) (126 \times 10^{6})}{(\frac{1}{4}) (1430)} \right]^{1/3}
$$

0.156 > 0.129

Thus for both fast and slow stock removal the scallop-shaped assumption holds.

 \hat{V}

E. SOLUTION OF **CONSTANT-STRESS** MODEL

1. Tests Using **60** Grit Wheel

a. Normal Direction:
$$
k_1^{60}
$$
 and k_2^{60} :

Since (21) contains the two unknowns, k_1 and k_2 , two experimental results must be used to evaluate these constants.

Case I: When
\n
$$
\overline{w} = 440 \times 10^{-6} \frac{10 \text{CH}}{5 \text{ECOND}}; \sigma = 118000 \text{ psi};
$$

\n $F_N = 15 \text{ LB}; W = .250 \text{ IN}; D_N = 2.37 \text{ inch};$
\n $M = 0.2; k_g = 1.025(110^5) \frac{LB}{1111}; D = 1.87 \text{ inch};$
\n $d = 0.016 \text{ in}; V = 1430 \frac{IN}{5 \text{EC}}; (21) \text{ becomes:}$
\n $1 = \frac{4.45 \text{ ft}}{(15)^{5/9}} \left(\frac{1}{4}\right)^{5/6} \left(\frac{2.37}{(118,000)^{1/6}}\right) \left(\frac{2.37^{17}}{(0.2)^{5}(1.025 \times 10^{5})^{5}(2.37 \cdot 1.87)^{5}(1.87)(.016)}\right)$
\n $\cdot \left(\frac{440 \times 10^{-6}}{1430}\right)^{2/3} + 3.4 \frac{R_2}{118,000}$

yielding

$$
f_{2} + 14.8(\sqrt{10}^{3}) f_{1} = 3.47(\sqrt{10}^{4})
$$
 (26)

Case 2: Wheel is dulled and cutting slowly, σ is reduced **by** the growth of flat area: $\sqrt{ }$ = 75×10^{-6} SEC; σ = 20630 psi other values as above, (21) becomes:

$$
1 = \frac{4.45 \cdot k}{(15)^{5/9}} \left(\frac{1}{4}\right)^{5/6} \frac{2.37^{17/8}}{(20630)^{1/6}} \left[\frac{1}{(0.2)(1.025 \times 10^5)(2.37 - 1.87)}\right] \frac{(75 \times 10^{6})^{2/3}}{[1.87(\omega 16)]^{1/8}}
$$

$$
\cdot \left(\frac{1}{1430}\right)^{2/3} + 3.4 \cdot \frac{k}{20630}
$$

yielding

$$
\hat{\mathcal{R}}_2 + 1.06 \left(10^{-3} \right) \hat{\mathcal{R}}_1 = 0.608 \left(10^4 \right) \tag{27}
$$

solving **(26)** and **(27)** simultaneously gives the constants:

 $k_1 = 2.08$ (10) psi $k_0 = 3900 \text{ psi}$

b. Tangential Direction: k_3^{60} and k_4^{60} :

Equation (22) contains the unknown constants k_3 and k_4 . Using the same cases as just used and substituting numerical values into (22) yields:

 $\frac{\text{Case 1: }}{\text{4}}$ = 0.5 for fast cutting, other values as previously given; (22) becomes:

$$
k_4 + 14.8 \ (10^{-3}) \ k_3 = 1.735 \ (10^4)
$$
 (28)

Case 2: $M = 0.3$ for slow-cutting or a dull wheel, other values as before; (22) becomes:

$$
k_4 + 1.06 (10^{-3}) k_3 = 0.182 (10^{4})
$$
 (29)

Solving **(28)** and **(29)** simultaneously yields the desired constants:

$$
k_3^{60} = 1.13 (10^{-6}) \text{ psi}
$$

\n
$$
k_4^{60} = 630 \text{ psi}
$$

Summarizing then, it is found that the constant stresses which exist are:

At the cutting zone:

In the normal direction: $k_1^{60} = 2.08$ (10⁶) psi

In the tangential direction: $k_3^{60} = 1.13$ (10⁶) psi

Under the flat on the worn grain:

- In the normal direction: k_2^{60} = 3900 psi
- In the tangential direction: $k_4^{60} = 630$ psi

2. Tests Using the **90** Grit Wheel

a. Normal Direction:
$$
k_1^{90}
$$
 and k_2^{90} :

In a similar manner (21) will be solved for k_1^{90} and k_2^{90} by using two experimental points from the **90** grit wheel tests.

Case **I:** Fast Cutting, high stress condition:

$$
\overline{M} = 460 \times 10^{-6} \frac{10}{\text{sec}}; \quad \overline{J} = 68000 \text{ psi}; \quad F_N = 15 \text{ LB};
$$

W = .250 in; D = 1.75 in; D_w = 2.5 in; M = 0.167;

$$
k_g = 2.66 (10^4) \frac{LB}{1N}
$$
; $d = 0.0085$ in; $V = 1340 \frac{in}{sec}$;
(21) becomes:

$$
k_2 + 15.1 (10^{-3}) k_{\overline{1}} = 2.02 (10^4)
$$
 (30)

Case 2: Slow cutting, lowsstress condition:

 \overline{W} = 75 x 10⁻⁶ $\frac{10}{\sec}$; σ = 19300 psi other values as above; (21) becomes:

$$
k_2 + 1.56 (10^{-3}) k_1 = 0.568 (10^{4})
$$
 (31)

Solving **(30)** and **(31)** yields the normal-direction constants:

$$
k_1^{90} = 1.071 (10^6) \text{ psi}
$$

$$
k_2^{90} = 4000 \text{ psi}
$$

b. Tangential Direction

Solving (22) for the two cases above will yield the unknown constants **k**₃ and **k**₄.

Case \vert : As before, fast cutting, so $\mathcal{M} = 0.5$; (22) becomes: k_4 + 15.1 (10^{-3}) k_3 = 1.01 (10^4) **(32)**

 $\frac{\text{Case 2:}}{\text{Out}}$ Dull wheel, so $\mathcal{M} = 0.3$; (22) becomes: k_4 + **1.56 (10⁻³)** k_3 **= 0.1704 (10⁴) (33)**

Solving **(32)** and **(33)** simultaneously yields the tangentialdirection constants:

$$
k_3^{90} = 0.62 (10^6) \text{ psi}
$$

$$
k_4^{90} = 750 \text{ psi}
$$

Summarizing then for the **90** grit wheel, the constant stresses are found to be:

At the cutting zone:

In the **normal direction:** $k_1^{90} = 1.071 (10^6)$ psi In the tangential direction: $k_3^{90} = 0.62$ (10⁶) psi

Under the worn grain flat:

In the normal direction: $k_2^{\overline{90}}$ = 4000 psi In the tangential direction: k_4^{90} = 750 psi

3. Solution of $(\overline{\overline{W}})$ Using Constants:

Copying (24) for convenience:

 $\left(\frac{\pi}{V}\right)$ = 0.106 $\left[AD(mk_3\left[D_w-D\right]\right)^5\right]^{1/2}$ $\left(F_N-3.4k_2A_R\right)^{3/2}$

(24 tepeated)

Where, as previously noted, the last parenthesis represents the available force for cutting, F_N , minus the force absorbed under the grain flat which will increase as A_{RI} increases. Thus the rate of cutting, \overline{W} , will be diminished as the wheel wears larger flats due to the potential force available being reduced **by** the force absorbed under the continually growing wear flat area.

The only variable in the above equation, as grinding is begun and continued, for any wheel, is A_{RL} . (F_N is varied independently, for test runs only.) Thus, for the **60** grit wheel tests, the only terms which could vary or be changed were F_N and A_{RI} . The same was true for the 90 grit wheel tests. Then In Appendix **6.1** is given the numerical derivation of the final forms of (24) for the **60** grit wheel. These results are:

$$
\overline{v}_{60}^{7.7518} = \frac{4.0}{6^{-5/4}} \left(\frac{\sigma}{\hbar \epsilon^6} - 3.4 \right)^{3/2}
$$
 (34)

$$
\overline{N}_{60}^{15LB} = \frac{6.91}{\sigma^{5/4}} \left(\frac{\sigma}{R_2^{60}} - 3.4 \right)^{3/2}
$$
 (35)

$$
\overline{W}_{60}^{3018} = \frac{12.3}{\sigma^{5/4}} \left(\frac{\sigma}{k_2^{60}} - 3.4 \right)^{3/2}
$$
 (36)

 $\overline{\textbf{W}}$ tabulations for various assumed σ values are given in TAble 4 and the theoretical formulas (34), **(35), (36)** are shown in Figure 12 with the other experimental data for the **60** grit wheel.

In Appendix **6.2** the same is done for the **90** grit wheels. The results are:

$$
\overline{v}_{30}^{5.08} = \frac{4.0}{0.5/4} \left(\frac{\sigma}{k_2^{90}} - 3.4 \right)^{3/2}
$$
 (37)

$$
\overline{v}_{90}^{\text{15LB}} = \frac{9.98}{6.5/4} \left(\frac{6}{k_2^{\text{90}}} - 3.4\right)^{3/2} \tag{38}
$$

$$
\bar{v}_{90}^{2519} = \frac{15.2}{6.5/4} \left(\frac{6}{k_2^{90}} - 3.4\right)^{3/2}
$$
 (39)

tabulations for various assumed **C"** values are given in Table **5** and the theoretical formulas **(37), (38), (39)** are shown in Figure **13** with the other experimental data for the **⁹⁰** grit wheel.

4. $\left(\begin{array}{c|c}\n\overline{\mathbf{w}} \\
\hline\n\overline{\mathbf{w}}\n\end{array}\right)$ <u>vs. σ </u> As A Linear Function

As stated in the introduction, originally the data obtained was plotted as \leftarrow vs. σ and a linear relation seemed to $r_{\scriptscriptstyle\mathsf{N}}$ exist. Taking equations (34) thru **(39)** and revising them, using **k2** values, yields:

$$
\frac{\sqrt{15}}{\sqrt{60}} = \frac{2.1 \times 10^{-6}}{5^{5/4}} \left(\sigma - 13350 \right)^{3/2}
$$
\n
$$
\frac{\sqrt{15}}{\sqrt{60}} = \frac{1.87 \times 10^{-6}}{0^{-5/4}} \left(\sigma - 13350 \right)^{3/2}
$$
\n(35a)
\n
$$
\frac{\sqrt{15}}{\sqrt{60}} = \frac{1.87 \times 10^{-6}}{0^{-5/4}} \left(\sigma - 13350 \right)^{3/2}
$$

$$
\frac{\overline{w}_{60}}{\overline{F}_{N}} = \frac{1.67 \times 10^{-6}}{\sigma^{5/4}} \left(\sigma - 13350 \right)^{3/2}
$$
\n(36a)

$$
\frac{\overline{05}_{20}}{\overline{F}_{N}} = \frac{3.18 \times 10^{-6}}{5.5/4} \left(\sigma - 13400 \right)^{3/2}
$$
 (37a)

$$
\frac{\overline{U}_{90}}{\overline{F}_{N}} = \frac{2.62 \times 10^{-6}}{0.5/4} \left(\sigma - 13400 \right)^{3/2}
$$
 (38a)

$$
\frac{\overline{v}_{90}}{\overline{F}_{N}} = \frac{2.4 \times 10^{-6}}{0.5/4} \left(\overline{0} - 13400 \right)^{3/2}
$$
 (39a)

Values of these functions are computed for various σ and V listed in Table 6. These results are plotted in Figures 14 and 15 with the experimental data.

5 with the experimental data.
The approximate linearity of $\left(\frac{\overline{w}}{\overline{h}}\right)$ vs. σ , initially assumed, is easily seen.

 $\bar{\epsilon}$

F. COMPARISON OF THE **CONSTANTS** WITH **SINGLE GRAIN SLIDING TESTS**

Using the apparatus shown in Figure **16** single grains of aluminum oxide were slid on hardened steel discs **(AISI** 4150, Rockwell **C - 60)** at speeds of **3000 - 5000** surface feet-pet minute. The results of these tests are shown on Figure **17** along with the results of the conventional grinding tests listed in Tables I, 2 and **3.** (The wheel results have been normalized **by** considering the fact that a grain in a wheel only contacts the work a fraction of the gross grinding time listed in Tables I, 2 and **3.)** For the condition of stress after **39000** inches of sliding, if it is assumed that after this much sliding, the cutting action has ceased then all the forces (and resulting stresses) measured are from the sliding of the grain flat over the work, then the stresses are measured as:

From the analytical work in this paper, the constant stresses determined to exist under the grain flat are:

Comparing these results, the order of magnitudes are correct, and indeed for the longest sliding test, the normal stress measured is almost exactly that derived, 4000 psi.

It is noted In Figure **17** that the continuation of the curve for the wheel tests into the single-grit data is smooth leading to the belief that the wheel merely acts as a group of individual grains and the action of a wheel may be predicted from the behavior of single grains.

This continuation-curve result was the basis for isolating a single grain from the wheel, and assuming a model of stock removal based on the single grain as done in this paper.

35 **G. CONCLUSIONS**

Based on single-grain sliding results it was assumed that the action of a grinding wheel could be analyzed **by** considering a single grain. The grain, with a worn flat, was assumed to have constant stresses acting on it in the normal and tangential directions at the cutting zone where the chip is produced and under the worn flat. The model was analyzed mathematically and the constants found **by** use of four experimental data points. The sliding constants seemed to agree closely with those found **by** the single-grain tests. The resulting equations then seemed to predict the results of **57** grinding tests.

Based on these facts the following conclusions may be stated:

1. The cutting potential of a wheel is the normal force and actual stress which it can apply to the work, in the absence of wheel breakage.

2. This normal force available for cutting is gradually reduced by the growth of worn-flat area at the grain-work interface. Since it appears that this junction supports a constant normal stress, the flat area growth gradually absorbs more and more of the available normal force until the area becomes so large it absorbs allitthe normal force at which time no stock removal would be possible.

3. In the grinding tests conducted the wheel continually removed metal and seemed to reach a situation where N was

-6 inch constant at a low value (about **75** x **10** second). The fact that the wheel cutting ability reached a plateau means that the real area of contact, A_{RI} , ceased to grow. Since it was also found that the wheel continually wore down (reaching a low-wear rate plateau at the same time as the low-rate \overline{w} plateau) means that the'grains were continually wearing down (at nottime did the total wheel wear exceed one grain diameter). Then the only way for the grain to wear down but the area A_{RI} , not to increase must be for some of the area flats to crack out through some mechanism, perhaps thermal. In any one wheel rotation the grain is in contact with the work and absorbs heat, then is plunged into a voilently churning coolant bath, possibly causing tensile stresses to develop, causing breakout of the flat area.

4. The above explains why a finely-dressed wheel (diamond dress lead of 50 x 10⁻⁶ wheel rev. say) cuts very slowly. A large contact area, A_{RI} is dressed onto the wheel and hence cutting, if possible, is at a slow rate, similar to when a large area has been created **by** attritious wear.

5. The constant normal and tangential stresses under the grain flats were found to be nearly identical for each wheel **(3900** and **630** psi for the **60** grit; 4000 and **750** psi for the **90** grit wheel). Since for the single grit studies values of 4070 and **1270** psi were found for the longest test and since these tests were run dry (no coolant) the use of coolant to reduce sliding friction under the wear flat is questionable in grinding.

6. From wel. known metal cutting theory see Appendix **7** and Figure **19':**

$$
\mathcal{C}_{\mathbf{p}} = k_3 \cos \emptyset \sin \emptyset - k_1 \sin^2 \emptyset
$$

Solving this for the **60** and **90** grit constants obtained in this paper, gives:

$$
\sum_{\text{plane}} = 140,000 \text{ psi} \quad \emptyset = 14^{\circ} \text{ for } 60 \text{ grit}
$$
\n
$$
\sum_{\text{p}} = 81,000 \text{ psi} \quad \emptyset \quad \emptyset = 14^{\circ} \text{ for } 90 \text{ grit}
$$

The work material ground in the tests was **AISI 52100** with Rockwell C \sim 60 having a tensile yield strength of about 300,000 psi. If $\tau_{\text{max}} = \frac{\sigma_{\text{y}}}{2}$ = 150,000 spi, then the computed results for the **60** grit wheel tests appear to be correct. **Why** the **90** grit test!compdtation is low is not understood.

7. The constant-stress model thus seems to be valid for the following reasons:

- **a.** The $\overline{\boldsymbol{w}}$ vs. $\sum_{k=1}^{\infty}$ curves shown theoretically and experimentally on Figures 12 and **13** have the proper "shape" and magnitude change with changing F_N .
- **b.** The values of sliding constants under the grain flat determined from the grinding data and theory give results remarkably similar to single-grit test results after long sliding distances.

c. For the **60** grit wheel, the constants at the cutting region can be related to the shear strength of the material along an average shear plane. For a shear angle of about 14° it is found that the shear strength is nearly predicted using these constants. **If** the average grain is assumed to have a zero or negative rake angle, a value of $\emptyset = 14^{\circ}$ is reasonable.

APPENDIX

I. Spacing of Effective Grits, S:

I.I Probability of Intersecting

Using Figure 2. Consider a line drawn from the origin "O" in a radial line. The probability of the line intersecting a first-ring grain flat is "P"; the probability of the line getting through or escaping the first ring without intersecting a flat is then (I-P). So that, the following table may be constructed.

1.2 Spacing of Effective Grits, 5:

The average distance travelled along this line is obtained when the probability of intersecting equals the probability of escaping. Thus the escaping probability must be equal **0.5** or:

$$
(1 - P_1)^k = 0.5 \t(1.2.1)
$$

ln 0.5 =
$$
k \ln (1 - P_1)
$$

\n $k = \ln 0.5$
\n $k = \ln (1 - P_1)$ (1.2.2)
\nwhere P₁ = $\frac{3}{\pi L}$

Figure 3 is a plot of k vs. P₁ and the curve may be approximated as

$$
P_1k = 1
$$
 (1.2.3)

so that
$$
k = \frac{1}{P_1} = \frac{\pi L}{3 \bar{a}}
$$
 (1.2.4)

Then the average distance travelled **by** a line from **"0"** is:

 $\overline{S} = kL$ (1.2.5)

using (1.2.4):

$$
\overline{S} = \frac{L}{P_1}
$$
\n
$$
\overline{S} = \frac{\pi L^2}{3a}
$$
\n
$$
\overline{S} \approx L^2 / \frac{1}{a}
$$
\n(1.2.6)

2. Length of Contact

24J Linearity of Grains in Contact with Length of Contact:

If the array of grinding grits given in Figure 2 is drawn onto a width equal to the workpiece width (i.e. o.250 inch) and the **"S"** direction is determined, the resultant picture is given as in Figure **18A.** Counting, perpendicular to the S direction, the number of grain centerlines intercepted in length $2X = L_c = d$ is found that there are fifty grains intercepted. In a length $2X1 \div 5D$, eighty three centerlines are encountered. Hence the following data may be tabulated:

Then the approximate relationship may be given as:

number of grains in contact for any
$$
L_c = nd
$$
 $L_c = mc$ (2.1.1)

2.2 Actual Deflection Parallel to
$$
F_N : N_{TRUE}
$$

Refer to Figure 18B. For the origin at the wheel centerline (axes X'y'), the X, y relation is:

$$
x^2 + y^2 = \left(\frac{D}{2}\right)^2
$$

moving the coordinates to the rigid-work and wheel interface is accomplished as:

$$
x^{2}
$$
 + $\left(y - \frac{D}{2}\right)^{2} = \left(\frac{D}{2}\right)^{2}$
 x^{2} = $yD - y^{2}$

For small y's (i.e. $y < .0006$ say), the y^2 term may be neglected, yielding

$$
X = \sqrt{\gamma D} \tag{2.2.1}
$$

- or

$$
y = \frac{x^2}{D}
$$
 (2.2.2)

Also from Figure 18B it may be seen that the actual deflected amount of the wheel under some F_N is

$$
y_T = y - y_{WORK}
$$

=
$$
\frac{x^2}{D} - \frac{x_W^2}{D_W}
$$

but for small X values, $X = X_W$ so

$$
y_T = \left(\frac{1}{D} - \frac{1}{D_W}\right)x^2
$$
\n
$$
y_T = \left(\frac{1}{D} - \frac{1}{D_W}\right) \cdot \left(\frac{L_c}{2}\right)^2
$$

$$
(F_N - \mathbf{L}_c)
$$
 Relation:

Now, taking as an assumption that any row of grains may be represented as Figure 18B (i.e. a single grain under a spring, k_{18}

Then

$$
dF_N = k_{18} dy_T \text{ for } 0 \lt L_c \lt d
$$

and

$$
dF_N = 2k_{18} dy_T \text{ for } d \leq L_c < 2d
$$

or generally

$$
dF_N = nk_{18} d(y_T) \text{ for } (n-1)d \quad \langle L_c \rangle
$$

Thus the force-displacement diagram of Figure **18C** may be taken as representing the wheel-work system, for an infinitelyrigid workpiece.

Assume an $F - y_T$ relation as:

$$
F_N = \mathbf{B} \mathsf{y}_T^n \tag{2.3.1}
$$

then

$$
\frac{dF_N}{dy_T} = Bn (v_T)^{n-1}
$$
 (2.3.2)

The boundary conditions for finding B and n above will vary depending upon "d" (i.e. depending upon the grain size of the wheel being considered).

Then, for the measured data:

Thus the ranges being considered for 60 and 90 grit wheels are shown on Figure 18C.

Then the boundary conditions for (2.3.2) above:

60 grit wheel:

$$
\mathcal{C}\frac{\mathcal{U}\tau}{\left(\frac{1}{D}-\frac{1}{D_{w}}\right)}=\frac{6.25}{4}d^{2}\left(ie.\ L_{c}=2.5d\right)
$$

$$
\mathcal{C}\frac{\mathcal{U}\tau}{\left(\frac{1}{D}-\frac{1}{D_{w}}\right)}=\frac{12.25}{4}d^{2}\left(ie.\ L_{c}=3.5d\right)
$$

90 grit wheel:

$$
\begin{aligned}\n\text{(iii)} \left(\frac{\mathrm{d}F}{\mathrm{d} \eta \tau} \right) &= 5 \, \hat{R}_{18} & \omega \frac{\eta \tau}{\left(\frac{1}{D} - \frac{1}{D_{\mathrm{w}}} \right)} = \frac{25}{4} \, d^2 \quad \left(i e \, L_e = 5d \right) \\
\text{(iv)} \left(\frac{\mathrm{d}F}{\mathrm{d} \eta \tau} \right) &= 7 \, \hat{R}_{18} & \omega \frac{\eta \tau}{\left(\frac{1}{D} - \frac{1}{D_{\mathrm{w}}} \right)} = \frac{49}{4} \, d^2 \quad \left(i e \, L_e = 7d \right)\n\end{aligned}
$$

Then 60 grit wheel:

(i)
$$
3k_{18} = Bn[(\frac{1}{D} - \frac{1}{D_w}) \frac{6.25}{4}d^2]^{n-1}
$$

(ii) $4k_{18} = Bn[(\frac{1}{D} - \frac{1}{D_w}) \frac{12.25}{4}d^2]^{n-1}$

Solving (i) for k_1 and inserting it into (ii) yields

$$
4\left\{\frac{8n}{3}\left(\left[\frac{1}{D}-\frac{1}{D_{w}}\right]\frac{6.25}{4}d^{2}\right)^{n-1}\right\} = 8n\left\{\left(\frac{1}{D}-\frac{1}{D_{w}}\right)\frac{12.25}{4}d^{2}\right\}
$$

$$
\frac{4}{3}\left(6.25\right)^{n-1} = (12.25)^{n-1}
$$

$$
\ln\frac{4}{3} + (n-1)\ln 6.25 = (n-1)\ln 12.25
$$

0.288 + (n-1) (1.83) = (n-1) (2.5)

$$
n = 1.4
$$

Say:

 $n = 1.5$

Solving for B from (i) using $n = 1.5$ and substituting into $(2.3.1)$ yields

$$
F_{\text{logrit}} = 1.6 \left(\frac{k_{60}}{d}\right) \left\{ \left(\frac{DD_w}{D_w - D}\right) \gamma_f^3 \right\}^{1/2}
$$
 (2.3.3)

or since

$$
\gamma_{\text{T}} = \left(\frac{D_{\text{w}} - D}{D D_{\text{w}}}\right) \left(\frac{L_{\text{c}}}{2}\right)^{2}
$$
\n
$$
F_{N_{60}} = 1.6 \frac{R_{60}}{d} \left\{\left(\frac{DD_{\text{w}}}{D_{\text{w}} - D}\right) \left(\frac{D_{\text{w}} - D}{D D_{\text{w}}}\right) \left(\frac{L_{\text{c}}}{2}\right)^{2}\right\}^{3}
$$
\n
$$
F_{N_{60}} = 1.6 \frac{R_{60}}{d} \left\{\left(\frac{D_{\text{w}} - D}{D D_{\text{w}}}\right)^{2} \left(\frac{L_{\text{c}}}{2}\right)^{6}\right\}^{1/2}
$$

$$
F_{N_{60}} = 0.2 \frac{R_{60}}{4} \left[\frac{D_w - D}{DD_w} \right] L_c^3
$$
 (2.3.4)

In a similar manner for the **90** grit wheel:

(111)
$$
5\hat{\mathcal{H}}_{18} = Bn \left[\left(\frac{1}{D} - \frac{1}{D_w} \right) \frac{25}{4} d^2 \right]^{n-1}
$$

\n(11) $7\hat{\mathcal{H}}_{18} = Bn \left[\left(\frac{1}{D} - \frac{1}{D_w} \right) \frac{49}{4} d^2 \right]^{n-1}$

yielding:

$$
\frac{7}{5} (25)^{n-1} = (49)^{n-1}
$$

n= 1.5

and as before:

$$
F_{N_{90}} = 1.33 \frac{k_{90}}{d} \left[\left(\frac{DD_{W}}{D_{W}-D} \right)^{2/2} N_{T}^{3} \right]^{1/2}
$$
 (2.3.5)

$$
F_{N_{90}} = 0.167 \frac{k_{90}}{d} \left(\frac{D_w - D}{D D_w} \right) L_c^3
$$
 (2.3.6)

(2.3.4) and **(2.3.6)** may be written in the form:

$$
F_{N_{\alpha}} = \frac{m_{\alpha} k_{\alpha}}{2} \left(\frac{D_{w} - D}{DD_{w}} \right) L_{c}^{3}
$$
 (2.3.7)

$$
L_c = \left[\frac{d \text{ DD}_w F_n}{m_q k_q (D_w - D)}\right]^{1/3}
$$
\n(2.3.8)

2.4 Determination of Spring Constants k_{18}^{60} and k_{18}^{90} :

Obviously k_{18}^{60} and k_{18}^{90} must be determined. For the 60 grit wheel, a measured point was:

$$
L_c = 0.047
$$
 inch $\ell F_N = 15$ LB

Then this into $(2.3.4)$ using $D_W = 2.37$ inch and $D = 1.87$ inch:

$$
15 = \frac{0.2 \times \frac{66}{18}}{0.016} \left[\frac{2.37 - 1.87}{(2.37)(1.87)} \right] (0.047)^{3}
$$

\n
$$
+ \frac{66}{18} = \frac{15 (0.016) (4.44)}{0.2 (0.5) (104 \times 10^{-6})}
$$

\n
$$
+ \frac{66}{18} = 1.025 (10^{5}) \frac{1.8}{18}
$$
 (2.4.1)

For the 90 grit wheel, a measured point was:

 $L_c = 0.005$ inch e F_N 15 LB
This into (2.3.6) Using $D_w = 2.5$ inch and $D = 1.75$ inch.

$$
15 = \frac{0.167 \hat{R}_{18}^{90}}{0.0085} \left[\frac{2.5 - 1.75}{2.5 (1.75)} \right] (0.055)^{3}
$$

$$
\hat{R}_{18}^{90} = \frac{15 (0.0085) (4.37)}{0.167 (0.75) (167 \times 10^{-6})}
$$

$$
\hat{R}_{18}^{90} = 2.66 (10^{4}) \frac{18}{18}
$$
 (2.4.2)

or

Inserting these values of k into (2.3.8) using D, D_W, d values as given above, yields

$$
L_c^{60} = 0.0192 \quad \sqrt[3]{F_N}
$$
 (2.4.3)

$$
L_{c}^{90} = 0.02 \qquad \sqrt[3]{F_{N}}
$$
 (2.4.4)

Taking an average of these as

$$
L_c = 0.0196 F_N^{1/3}
$$
 (2.4.5)

This curve and the remaining experimental data **are shown** plotted in Figure **180.**

1, General Stress in Normal Direction

Copying (4b):
\n
$$
I = \frac{1}{F_N} \left[k_1 bt + k_2 \frac{2}{a} \right]
$$

Using **(19)** and (20) yields

$$
1 = \frac{W}{0.294 \text{ F}_{\text{N}}}\left(\frac{DD_{\text{W}}F_{\text{N}}d}{m.k_{\text{P}}d^{6}[D_{\text{W}}-D]}\right)^{1/3}\left(0.8\bar{a}k_{\text{I}}t + k_{\text{Z}}\bar{a}^{2}\right)
$$

rewritting:

$$
1 = \frac{3.4 \text{ w}}{F_{\text{N}}^{2/3}} \left[\frac{DD_{\text{w}}d}{mk_{\text{g}}(D_{\text{w}}-D)} \right]^{1/3} \left[0.8 \text{ k}_{1} \left(\frac{\bar{\alpha}}{d} \right) \left(\frac{t}{d} \right) + k_{2} \left(\frac{\bar{\alpha}}{d} \right)^{2} \right]
$$

Using (13a) and (18):

$$
1 = \frac{3.4 \text{ W}}{F_{\text{N}}^{2/3}} \left[\frac{DD_{\text{W}}d}{mF_{\text{G}}(D_{\text{W}}-D)} \right]^{1/3} \left[0.8 \text{ ft} \left(\frac{A_{\text{R}}}{\text{W}} \right)^{1/2} \left\{ \frac{mF_{\text{G}}(D_{\text{W}}-D)}{dDD_{\text{W}}F_{\text{N}}} \right\} \frac{(635 \text{ ft})}{(Dd)^{2/9}} \cdot \left(\frac{\overline{\text{W}}}{\text{W}} \right)^{2/3} \left\{ \frac{N}{mF_{\text{Q}}(D_{\text{W}}-D)} \right\}^{1/9} + \frac{1}{\text{W}} \left\{ \frac{A_{\text{R}}}{\text{W}} \left(\frac{mF_{\text{G}}(D_{\text{W}}-D)}{dDD_{\text{W}}F_{\text{N}}} \right)^{1/3} \right\}
$$

Regrouping:

$$
1 = \frac{4.45 \cdot k_{1}w^{5/6}A_{R1}^{1/6} \left[\frac{1}{mk_{3}(D_{w}-D)} \right]^{5/18} \left(DD_{w}d \right)^{3/18} \left(\frac{D_{w}^{7}}{D^{2}d^{2}} \right)^{7/3} + \frac{3.4 k_{2}A_{R1}}{F_{1}}.
$$

Using definition of σ , yields:

 \mathcal{A}

$$
1 = \frac{4.45 k_{1} w^{56}}{F_{N}^{5/9} \sigma^{1/6}} \left[m^{5} k_{3}^{5} (D_{w} - D)^{5} (D_{w}) \right]^{1/8} + 3.4 \frac{k_{2}}{\sigma}
$$

4. General Equation for $(\frac{\overline{w}}{\overline{V}})$

$$
\left(\frac{\overline{v}}{\overline{v}}\right)^{2/3}\left[\left[1-3+\frac{k}{\sigma}\right]\left(\frac{\sigma^{1/6}F_{N}^{5/9}}{4.45 k_{N}W^{5/6}}\left[\frac{dD(mk_{\theta}[D_{w}-D])}{D_{w}^{17}}\right]^{5}\right]^{1/8}\right]
$$
\n
$$
\left(\frac{\overline{v}}{\overline{v}}\right)=\frac{\sigma^{1/4}F_{N}^{15/10}}{9.4 k_{1}^{3/2}W^{5/4}}\left[\frac{dD(mk_{\theta}[D_{w}-D])}{D_{w}^{17}}\right]^{1/2}\left(1-3+\frac{k_{2}}{\sigma}\right)^{3/2}
$$

writing
$$
\sigma = \frac{F_N}{A_{\text{RL}}}
$$
 and rearranging the last parentheses is:

$$
\left(\frac{\overline{N}}{V}\right) = \frac{0.106}{R_{1}^{3/2}} \left(\frac{F_{N}}{A_{RL}}\right)^{1/4} \frac{F_{N}^{15/18}}{W^{5/4}} \left[\frac{d D \left(mR_{3} [D_{W} - D]\right)^{5}}{D_{W}^{17}}\right]^{1/2} \left(\frac{F_{N}}{F_{N}} - \frac{3.4 R_{2} A_{RL}}{F_{N}}\right)^{3/2}
$$

$$
\left(\frac{\overline{w}}{V}\right) = \frac{0.106}{\hat{R}_{1}^{3/2} \text{ W}^{5/4} \text{ A} \hat{\kappa}_{1}^{1/4}} \left[\frac{dD \left(m\hat{R}_{3} \left[\overline{D_{w}} \cdot \overline{D}\right]\right)^{5}}{D_{w}^{17}}\right]^{1/2} \cdot \left(F_{N} - 3.4 \hat{R}_{2} \text{ A} \hat{\kappa}_{2}\right)^{3/2}
$$

$$
\left(\frac{\overline{w}}{v}\right) = \frac{0.106 \left[dD(mk_9[D_w-D])\right]^5 J^{1/2}}{k_1^{3/2} W^{5/4} A_{RL}^{1/4} F_w^{15/36} D_w^{17/12}} \left(F_w - 3.4 k_2 A_{RL}\right)^{3/2}
$$

$$
\left(\frac{\overline{w}}{v}\right) = \frac{0.106 \left[d D(m k_{\text{g}} \left[D_{w} - D \right]) \right]^{5}}{k_{1}^{3/2} w^{5/4} D_{w}^{1/7/2} E_{n}^{5/12} A_{\text{RL}}^{1/4}} \left(F_{N} - 3.4 k_{\text{g}} A_{\text{RL}}\right)^{3/2} \tag{4.1.2}
$$

$$
\frac{1}{18^{3/2}} \left[\frac{1}{18^{3/2}} \left(\frac{1}{18} \right)^5 \right]^{1/2} = \frac{18^{3/2}}{18^{3/2}} = \frac{18^{3/2}}{18^{3/2}} = \frac{36+2}{12} = \frac{38}{12}
$$

dimensions check.

5. Scallop-Shaped Assumption

$$
\frac{N}{V} \qquad \frac{N}{S} \tag{14}
$$

Using (20) and (14) yields:

$$
\frac{a\sigma}{\sqrt{2}} > \frac{2\sqrt{10}}{d^2/\bar{a}}
$$

Using (22)

$$
\frac{d\sigma}{d\theta} > \frac{2 D^{1/2} \overline{a}}{d^{2}} \left[\left(\frac{2 \pi D_{w} \overline{c} d^{2}}{3 D^{1/2} \overline{v} \overline{a}} \right)^{2/3} \right]^{1/2}
$$
\n
$$
\frac{d\sigma}{d\theta} > \frac{2 D^{1/3} \overline{a}^{2/3}}{d^{4/3}} \left(\frac{2 \pi}{3} \cdot \frac{D_{w} \overline{c} d}{\overline{v}} \right)^{1/3}
$$
\n
$$
\frac{d\sigma}{d\theta} > 2.56 \left(\frac{DD_{w} \overline{c} d}{\overline{v} d^{2}} \right)^{1/3} \left(\frac{\overline{a}}{d} \right)^{2/3}
$$

Using (18a)

$$
\frac{dV}{dV} > 2.56 \left(\frac{DD_{w} \overline{v}}{V} \frac{d^{2}}{d^{2}} \right)^{1/2} \left\{ \left(\frac{A_{RL}}{W} \right)^{1/2} \left[\frac{mk_{g}(D_{w}-D)}{d^{2}D_{w}F_{w}} \right]^{1/6} \right\}^{4/3}
$$
\n
$$
\frac{dV}{dV} > 2.56 \left(\frac{DD_{w} \overline{v}}{V} \frac{d^{2}}{d^{2}} \right)^{1/3} \left(\frac{A_{RL}}{W} \right)^{1/3} \left[\frac{mk_{g}(D_{w}-D)}{d^{2}D_{w}F_{w}} \right]^{1/9}
$$
\n
$$
\frac{dV}{dV} > 2.56 \left(\frac{D^{2}D_{w}^{2}}{d^{7}} \right)^{1/9} \left(\frac{A_{RL} \overline{v}}{W} \right)^{1/3} \left(\frac{mk_{g}(D_{w}-D)}{F_{w}} \right)^{1/9}
$$
\n
$$
\frac{dV}{dV} > 2.56 \left[\frac{(DD_{w})^{2} m k_{g}(D_{w}-D)}{F_{w} d^{7}} \right]^{1/9} \left[\frac{A_{RL} \overline{v}}{W} \right]^{1/9}
$$

Copying (24):

$$
\left(\frac{\overline{w}}{v}\right) = \frac{0.106}{k^{3/2}} \frac{[dD(mk_3[D_w - D])^5]}{w^{5/4} D_w^{17/12} E_w^{5/12} A_w^{1/4}} \left(F_w - 3.4 k_2 A_w\right)^{3/2}
$$
 (24)

6.1 60 Grit Wheel Tests:

For the 60 grit wheel the constant terms were as follows: $V = 1430 \frac{\text{in}}{\text{sec}}$; d = 0.016 in; D = 1.87 in; M = 0.2; k_g = 1.025 (10⁵) $\frac{16}{10}$ $D_W = 2.37$ in; kpo = 2.08 (10⁶)psi; W = .250 in; k₂⁶⁰ = 3900 psi. Then insertion of these into (24):

$$
\overline{U}_{60} = \frac{0.106 (1430) \{0.016 (1.87)\}^{1/2} [0.2(1.025 \times 10^5) (2.37 - 1.87)]^{3/2}}{(2.08 \times 10^6)^{3/2} (0.250)^{5/4} (2.37)^{17/2} F_n^{5/12} A_n^{1/4}} \cdot \left(\frac{1}{1}\right)^{3/2}
$$
\n
$$
\overline{U}_{60} = \frac{0.106 (1430) \{0.746\} [47.0]}{(3 \times 10^9) (\frac{1}{5.65}) (3.4) F_n^{5/12} A_n^{1/4}} \left(\frac{1}{1}\right)^{3/2}
$$
\n
$$
\overline{U}_{60} = \frac{2.94 \times 10^{-6}}{F_n^{5/12} A_{RL}^{1/4}} [F_n - 3.4 f_{R2} A_{RL}]^{3/2}
$$

writing $A_{RL} = F_{N/\sigma}$ yields:

$$
\overline{N}_{60} = \frac{2.94 \times 10^{-6} \text{ C}^{1/4}}{F_{N}^{5/12} F_{N}^{1/4}} \left[F_{N} - 3.4 \overline{R}_{2} F_{N} \right]^{3/2}
$$

$$
\overline{N}_{60} = \frac{2.94 \times 10^{-6} \sigma^{1/4} F_N^{3/2}}{F_N^{2/3}} \left[\sigma - 3.4 \overline{R}_z \right]^{3/2} (6.1.1)
$$

where the 3.4k₂ stress is recognized as σ _o or that stress where \overline{N} = 0 as defined previously.

Rearranging:

$$
\overline{N}_{60} = \frac{2.94 \times 10^{-6} \text{ F}_{\text{N}}^{-5/6} \text{ A}_{\text{2}}^{3/2}}{\sigma^{-5/4}} \left[\frac{\sigma}{\overline{R}_{2}} - 3.4 \right]^{3/2} \tag{6.1.2}
$$

Then for tests at various F_N 's:

$$
\overline{v}_{60}^{7.7518} = \frac{4.0}{\sigma^{5/4}} \left[\frac{\sigma}{\hbar_{2}^{60}} - 3.4 \right]^{3/2}
$$

$$
\overline{v}_{60}^{1518} = \frac{6.91}{\sigma^{5/4}} \left[\frac{\sigma}{\hbar_{2}^{60}} - 3.4 \right]^{3/2}
$$

$$
\overline{v}_{60}^{308} = \frac{12.3}{\sigma^{5/4}} \left[\frac{\sigma}{\hbar_{2}^{60}} - 3.4 \right]^{3/2}
$$

 $(6.1.3)$

6.2. 90 Grit Wheel Tests

For the 90 grit wheel the constant terms were as follows: $V = 1340 \frac{\text{in}}{\text{sec}}$; d = 0.0085 in; D = 1.75 in; M = 0.167; k_g = 2.66(10⁴); $D_W = 2.5$ in; $k_1^{90} = 1.071$ (10⁶) psi; $W = 1/4$ in; $k_2^{90} = 4000$ psi. Then insertion of these into (24):

$$
\overline{U}_{90} = \frac{0.106 (1340) \{0.0085 (1.75)\}^{1/2} [0.167 (0.266 \times 10^5) (2.5 - 1.75)]^{5/2}}{(1.071 \times 10^6)^{3/2} (\frac{1}{4})^{5/4} (2.5)^{1/2} F_N^{5/12} A_N^{1/4}} (\cdot \cdot)^{3/2}
$$
\n
$$
\overline{U}_{90} = \frac{0.106 (1340) (0.705) (29.5)}{(1.11 \times 10^9) (\frac{1}{5.65}) (3.65) F_N^{5/12} A_{R1}^{1/4}} (\cdot \cdot \cdot)^{3/2}
$$
\n
$$
\overline{U}_{90} = \frac{4.13 \times 10^{-6}}{F_N^{5/12} A_{R2}^{1/4}} [F_N - 3.4 F_Z A_{R2}]^{3/2}
$$
\n
$$
\overline{U}_{90} = \frac{4.13 \times 10^{-6}}{F_N^{5/12} A_{R2}^{1/4}} [F_N - 3.4 F_Z A_{R2}]^{3/2}
$$
\n
$$
\overline{U}_{90} = \frac{4.13 \times 10^{-6}}{F_N^{5/12} A_{R2}^{1/4}} [F_N - 3.4 F_Z A_{R1}]^{3/2}
$$
\n
$$
\overline{U}_{90} = \frac{4.13 \times 10^{-6}}{F_N^{5/12} A_{R2}^{1/4}}
$$

$$
\overline{10}_{90} = \frac{4.13 \times 10^{-6} \text{ F}_{\text{N}}^{5/6} \text{ F}_{\text{N}}^{3/2}}{\sigma^{5/4}} \left[\frac{\sigma}{\text{F}_{\text{Z}}} - 3.4 \right]^{3/2} \tag{6.2.2}
$$

Then for tests at various F_N 'S

$$
\bar{N}_{90}^{548} = \frac{4.0}{\sigma^{5/4}} \left[\frac{\sigma}{k_2^{90}} - 3.4 \right]^{3/2}
$$

 $(6.2.3)$

$$
\bar{w}_{90}^{'545} = \frac{9.98}{\sigma^{5/4}} \left[\frac{\sigma}{k_2^{90}} - 3.4 \right]^{3/2}
$$

$$
\overline{v}_{90}^{2548} = \frac{15.2}{\sigma^{5/4}} \left[\frac{\sigma}{k_2^{90}} - 3.4 \right]^{3/2}
$$

(6.2.3 cond't)

7. Metal-Cutting Theory

7.1 Stresses in the Shear Zone:

See Figure **19.** From sum of forces along the shear plane direction:

$$
k_{3}^{*} \text{ bt } \cos \emptyset - k_{1}^{*} \text{ bt } \cot \emptyset \sin \emptyset = \mathcal{E}_{p} \frac{\text{bt}}{\sin \emptyset}
$$
\n
$$
\mathcal{C}_{p} = k_{3}^{*} \cos \emptyset \sin \emptyset - k_{1}^{*} \cos \emptyset \sin \emptyset \qquad (7.1.1)
$$

but the definition of **k*** and **k*** used in this paper is **3**

$$
k_3 = k_3^*
$$

$$
k_1 = k^* \cot \emptyset
$$

then **(7.1.1)** becomes

$$
c_p = k_3 \cos \emptyset \sin \emptyset - k_1 \sin^2 \emptyset \qquad (7.1.2)
$$

Solving **(7.1.2)** for **0** plane gives maximum values of:

$$
\begin{array}{rcl}\n\mathbf{C}_p^{60} & = & 140,000 \text{ psi at } \emptyset \neq 14^{\circ} \\
\hline\n\end{array}
$$
\n
$$
\mathbf{C}_p^{60} = 8
$$
\n
$$
\mathbf{C}_p^{60} = 81,000 \text{ psi at } \emptyset = 14^{\circ}
$$

8. Elastic-Rebound Model

8.1 General Equation in Normal Direction

The force equation in the normal direction is:

$$
1 = \frac{bt}{F_{Ng}}
$$
 $(k_1 + k_2 \bar{a})$

Using (19) and (20), the above becomes:

$$
1 = \frac{0.8\overline{a} \, \text{tw}}{0.294 \, \text{F}_N} \left[\frac{D \, D_w \, F_n \, d}{m \, k_g \, (D_w - D) \, d^6} \right]^{\frac{1}{3}} \cdot \left\{ k_i + k_z \, \overline{a} \right\}
$$
\n
$$
1 = \frac{2.72 \, \text{w}}{\text{F}_N} \left(\frac{\overline{a}}{d} \right) \left(\frac{t}{m} \right) \left[\frac{D D_w \, F_N \, d}{m k_g \, [D_w - D]} \right]^{\frac{1}{3}} \cdot \left(k_i + k_z \overline{a} \right)
$$

using (13a), (13b) and (18):

$$
1 = \frac{2.72 \text{w}}{F_{N}} \left[\left(\frac{A_{R1}}{\text{w}} \right)^{12} \left(\frac{mk_{3} \left[D_{w} - D \right]}{\text{d} D D_{w} F_{N}} \right)^{16} \right] \left\{ 1.635 \left(\frac{\pi}{V} \right)^{2} \left(\frac{W}{A_{R1}} \right)^{13} \left[\frac{D_{w} F_{N}}{mk_{3} \left[D_{w} - D \right] \left(D_{d} \right)^{2}} \right] \right\}
$$
\n
$$
\left(\frac{DD_{w} F_{N}}{mk_{3} \left[D_{w} - D \right]} \right)^{1/3} \left\{ k_{1} + k_{2} \left(\frac{A_{R1}}{\text{w}} \right)^{1/2} \left(\frac{mk_{3} \left[D_{w} - D \right]}{\text{d} D D_{w} F_{N}} \right) \right\}
$$

combining and using the definition of σ , yields

$$
1 = \frac{4.45 \text{w}}{F_{\text{N}}^{5/9} \sigma^{1/6}} \left[\left(\frac{\bar{w}}{V} \right)^{2/3} \left\{ \frac{D_{\text{w}}}{m \hat{k}_{3} (D_{\text{w}} \cdot \overline{D})} \right\} \left(\frac{D_{\text{w}}^{12}}{d \overline{D}} \right)^{1/8} \left\{ \hat{k}_{1} + \frac{\hat{k}_{2} F_{\text{w}}^{12}}{(\sigma \text{w})^{1/2}} \left[\frac{m \hat{k}_{3} [D_{\text{w}} \cdot \overline{D}]}{d \overline{D} D_{\text{w}}} \right]^{1/8} \right\}
$$

8.2 Solution in Normal Direction Using Data Points

The solution of this model using the cases already described yields the following equations:

$$
k_1^{60} + 1.129 (10^{-3}) k_2^{60} = 232 (10^4)
$$

Slow cutting:

$$
k_1^{60} + 2.69 (10^{-3}) k_2^{60} = 566 (10^{4})
$$

Solving these simultaneously yields:

These results dictate that the normal stress applied **by** the grain on the work in the cutting region is tensile so as to compensate for the stress under the grain flat which is very large. This model is thus discarded.

TABLE 1

(90 GRIT WHEEL)

TABLE 3

FORMULAS FROM TEXT:

 $\overline{10}_{7.7518} = \frac{4.0}{\sigma^{5/4}} \left\{ \frac{\sigma}{\frac{6}{5}} - 3.4 \right\}^{3/2} \dots (34)$ $\overline{v}_{1518} = \frac{6.91}{\sigma^{5/4}} \left\{ \frac{\sigma}{\overline{h}_{2}^{6}} - 3.4 \right\}^{3/2} \dots (35)$ $\overline{v}_{3018} = \frac{12.3}{\sigma^{5/4}} \left\{ \frac{\sigma}{\frac{46}{5}} - 3.3 \right\}^{3/2} \dots (36)$

> (60 GRIT WHEEL) TABLE 4

WHERE: \hat{k}_2^6 = 3900 psi

FORMULAS FROM THE TEXT:

 $\overline{N}_{518} = \frac{4.0}{\sigma^{5/4}} \left\{ \frac{\sigma}{\phi_2} - 3.4 \right\}^{3/2} \cdot \cdot \cdot$ (37) $\overline{w}_{1515} = \frac{9.98}{\sigma^{5/4}} \left\{ \frac{\sigma}{4^{3/2}} - 3.4 \right\}^{3/2} \cdots \cdots (38)$
 $\overline{w}_{2515} = \frac{15.2}{\sigma^{5/4}} \left\{ \frac{\sigma}{4^{3/2}} - 3.4 \right\}^{3/2} \cdots \cdots (39)$

WHERE: $\frac{1}{2}$ ⁹⁰ = 4000 psi

(90 GRIT WHEEL) TABLE 5

TABLE 6

 $F = k_2 * a b t$ $F_A = \frac{1}{2} \hbar \Delta b t$

FOR CONSTANT- STRESS MODEL: AT THE SHEAR ZONE: $F_v =$ A_i^* bt cor ϕ
 $F_H =$ A_3^* bt AT THE SLIDING ZONE: $F_v = \frac{24}{3} (\pi a^2/4)$ $F_H = k_A^* (\pi \bar{a}^2 / 4)$

FIG.1 GRAIN-CHIP GEOMETRY

FIG.3

WHEEL ROTATION AND GRINDING DIRECTION

(ABOVE IS A DRAWING TRACED FROM AN ACTUAL PHOTOGRAPH)

FIG. 4 PHOTO OF WORN WHEEL

SERIES Nº 19: \overline{N} = 360 x 10⁻⁶ $\frac{INCH}{SEC}$; F_N= 30 UB; ARL= 684 x 10⁻⁶ IN²; σ = 43,861 psi

SERIES Nº 17: \bar{w} = 110 × 10⁻⁶ INCH; F_N = 30LB; ARL=1472 × 10⁻⁶ IN²; σ = 20, 377 psi ABOVE DRAWINGS WERE TRACED FROM ACTUAL PHOTOGRAPHIC "STRIPS. ORIGINAL STRIPS WERE REDUCED IN SIZE. PHOTO "STRIPS" OF THREE WORN

FIG.6 CUTTING GEOMETRY

WORK WIDTH WAS 0.250 INCH; WHEELS WERE 3/8 INCH WIDE, AND SINCE ONLY PLUNGE GRINDING WAS PERFORMED, THE WORN PORTION OF THE WHEEL WAS EASILY DISCERNABLE

FIG. 9 OBTAINING A "STRIP"

FIG. 11 MEASURING NT

FIG. 12 EXP. AND THEOR. N VS.

FIG. 13 EXP. AND THEOR. N vs.

FIG. IG NSF LATHE ARRANGED
FIG. IG NSF LATHE ARRANGED

FIG. 18A ARRAY OF GRAINS

THE ABOVE FIGURE IS FOR 60 GRIT (D= 0.016 INCH). SHOWN 25X SIZE

ABOVE PICTURE REPRESENTS A PLANE THRU THE WHEEL AND WORK. EACH GRAIN SHOWN REPRESENTS A ROW OF 18 GRAINS. THEN EACH ROW CAN BE CONSIDERED TO GRAIN MOUNTED ON A **BE** ONE R_{18} SPRING OF STIFFNESS THEN FOR THREE ROWS IN CONTACT $(L_c=2x=3D)$: $\frac{1}{2}$ dry dF= $3k_{18}$ dry SPRING MODEL OF WHEEL $FIG.18$

FIG. 18D EXP. AND THEOR. L. VS.F.

 R_3^* bt cosd - R_1^* bt cord sind - C_p bt = 0
 $C_p = R_3^*$ cosd sind - R_1^* cosd sind BUT $k_3 = k_3^*$ AND $k_1 = k_1^*$ cord \therefore $x_p = k_3 \cos\phi \sin\phi - k_1 \cos\phi \sin\phi$
 $x_p = k_3 \cos\phi \sin\phi - k_1 \sin^2\phi$ FIG. 19 SHEAR-PLANE RELATIONS