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*Measurement of inclusive jet cross sections  
in pp and PbPb collisions at  $\sqrt{s_{NN}} = 2.76$  TeV*

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**Measurement of inclusive jet cross sections in  $pp$  and PbPb collisions at  $\sqrt{s_{NN}} = 2.76$  TeV**V. Khachatryan *et al.*\*  
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Inclusive jet spectra from  $pp$  and PbPb collisions at a nucleon-nucleon center-of-mass energy of 2.76 TeV, collected with the CMS detector at the CERN Large Hadron Collider, are presented. Jets are reconstructed with three different distance parameters ( $R = 0.2, 0.3, \text{ and } 0.4$ ) for transverse momentum ( $p_T$ ) greater than 70 GeV/c and pseudorapidity  $|\eta| < 2$ . Next-to-leading-order quantum chromodynamic calculations with nonperturbative corrections are found to overpredict jet production cross sections in  $pp$  for small distance parameters. The jet nuclear modification factors for PbPb compared to  $pp$  collisions, show a steady decrease from peripheral to central events, along with a weak dependence on the jet  $p_T$ . They are found to be independent of the distance parameter in the measured kinematic range.

DOI: [10.1103/PhysRevC.96.015202](https://doi.org/10.1103/PhysRevC.96.015202)**I. INTRODUCTION**

Heavy ion collisions at the CERN Large Hadron Collider (LHC) can generate a hot and dense deconfined state of matter, also known as the quark-gluon plasma (QGP). In these collisions, hard scattered partons are expected to be attenuated due to elastic and inelastic interactions with the produced medium [1–3]. This phenomenon is also known as “jet quenching”, originally proposed in Ref. [4], and is indirectly confirmed by measurements of spectra and correlations of high transverse momenta ( $p_T$ ) hadrons at the BNL Relativistic Heavy Ion Collider (RHIC) [5–8] and LHC [9–11]. In these measurements, jet quenching is observed to have a dependence on event multiplicity and hadron  $p_T$ , and has provided significant insights, including the color opaqueness of the QGP. However, these findings are limited by intrinsic biases. For example, the leading hadron measurements are preferentially from the population of jets that have the least interaction with the medium. These measurements are also not sufficient to discriminate quantitatively between partonic energy loss formalisms or to extract key parameters such as the transport coefficient of the hot medium to precisely measure the stopping-power of the QGP (see Refs. [12,13] for reviews). As jet quenching is intrinsically a partonic process, studies using hadronic observables blur essential physics due to the complexity of the theoretical description of hadronization and the sensitivity to nonperturbative effects. The measurement of jet structure and its modification in terms of energy flow rather than hadronic distributions promises a much closer connection to the underlying theory. Therefore a quantitative picture of jet quenching with respect to theoretical assumptions can be obtained through a full reconstruction of underlying parton kinematics, i.e., jet reconstruction [14,15].

Complementary and robust jet measurements in heavy ion collisions became feasible with the beginning of the LHC heavy ion program. For example, measurements showed that the  $p_T$  of back-to-back dijet pairs becomes increasingly unbalanced as the centrality of the event increases (smaller impact parameters) [16–18]. In these collisions jet pairs are also observed to be undeflected, i.e., their azimuthal angular correlations are independent of the collision centrality. Furthermore, measurements of jet shape, fragmentation functions, jet-track correlations, and missing  $p_T$  find that a significant fraction of the “lost” jet energy is observed to be radiated via low- $p_T$  particles far outside the jet cone [17,19–22]. The comparison of inclusive jets in heavy ion collisions with those in  $pp$  collisions can differentiate between competing models of parton energy loss mechanisms [23–25]. Initial measurements of jet yields in central heavy ion collisions were compared to a  $pp$  baseline, and they are found to have a weak dependence on the jet  $p_T$ , with the low  $p_T$  region suffering slightly larger modification compared to the high  $p_T$  region [26,27]. However the interpretation of the jet modification results in nucleus-nucleus collisions and the understanding of their relation to the properties of the QGP requires detailed knowledge of all nuclear effects that could influence the comparisons with the  $pp$  system. The shape of the jet spectrum in proton-lead collisions is similar to that observed in  $pp$  collisions [28–30]. This suggests the modification of the jet spectra observed in PbPb collisions is indeed an effect of the hot medium produced in these collisions.

For this analysis, the jet measurements are performed as a function of three experimental observables: the jet reconstruction distance parameter [31], the jet  $p_T$ , and the event centrality (related to the impact parameter of the incoming nuclei) of the collisions. The reference  $pp$  jet cross section is also measured and is compared to perturbative quantum chromodynamic (pQCD) calculations. The observable of interest is the jet nuclear modification factor ( $R_{AA}$ ), defined as

$$R_{AA} = \frac{d^2 N_{\text{jets}}^{AA} / dp_T d\eta}{\langle T_{AA} \rangle d^2 \sigma_{\text{jets}}^{pp} / dp_T d\eta}, \quad (1)$$

where  $N_{\text{jets}}^{AA}$  is the jet spectrum measured in PbPb,  $\sigma_{\text{jets}}^{pp}$  is the jet cross section from  $pp$  collisions, and  $\langle T_{AA} \rangle$  is the nuclear

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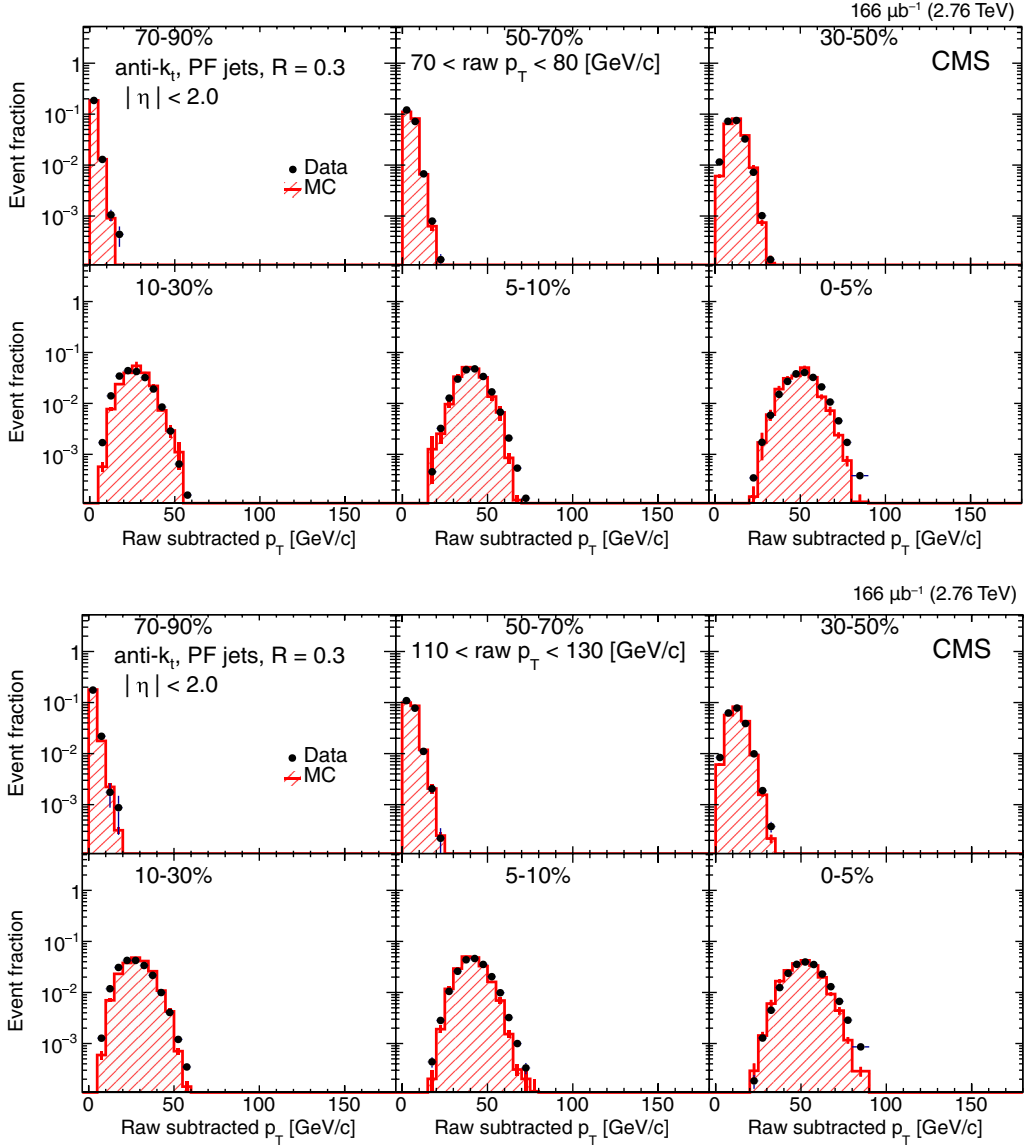


FIG. 1. Raw subtracted  $p_T$  for jets reconstructed with the anti- $k_T$  algorithm and a distance parameter of  $R = 0.3$ , in the ranges  $70 < \text{jet } p_T < 80$  [GeV/c] (top panels) and  $110 < \text{jet } p_T < 130$  [GeV/c] (bottom panels). This quantity is found by taking the difference of the sum of PF candidates within the jet cone and raw jet  $p_T$ . Solid symbols show data, and the histogram is from PYTHIA + HYDJET generated events.

overlap function averaged over the event class studied. The quantity  $\langle T_{AA} \rangle$  is related to the mean number of nucleon-nucleon ( $NN$ ) collisions  $\langle N_{\text{coll}} \rangle$ , and  $\sigma_{\text{inel}}^{NN}$ , the nucleon-nucleon inelastic cross section, through  $\langle N_{\text{coll}} \rangle = \langle T_{AA} \rangle \sigma_{\text{inel}}^{NN}$ , and is calculated with a Monte Carlo Glauber model description of the nuclear collision geometry (for a review see Ref. [32]).

## II. THE CMS DETECTOR AND EVENT SELECTION

The central feature of the CMS apparatus is a superconducting solenoid providing a magnetic field of 3.8 T. Charged-particle trajectories are measured with the silicon tracker that allows a transverse impact parameter resolution of  $\sim 15 \mu\text{m}$  and a  $p_T$  resolution of  $\sim 1.5\%$  for particles with  $p_T = 100 \text{ GeV}/c$ . A  $\text{PbWO}_4$  crystal electromagnetic calorimeter (ECAL) and a brass and scintillator hadron

calorimeter (HCAL) surround the tracking volume. The forward regions are instrumented with iron and quartz-fiber hadron forward calorimeters (HF). A set of beam scintillator counters (BSC), used for triggering and beam halo rejection, is mounted on the inner side of the HF calorimeters. The very forward angles are covered at both ends by zero-degree calorimeters (ZDC). A more detailed description of the CMS detector, together with a definition of the coordinate system used and the relevant kinematic variables, can be found in Ref. [33].

The first level of the CMS trigger system, composed of custom hardware processors, uses information from the calorimeters to select the most interesting events in a fixed time interval of less than  $4 \mu\text{s}$ . The high-level trigger (HLT) processor farm further decreases the event rate, from around 100 kHz to less than 1 kHz, before data storage. The PbPb analysis uses minimum bias triggered and single-jet HLT

data sets. The minimum bias events are characterized by the coincidence of signals in the two HF detectors or the forward and backward BSCs. The triggers used in the analysis are constructed from ECAL and HCAL energies requiring a single jet with  $p_T > 55, 65, \text{ and } 80 \text{ GeV}/c$ . For  $pp$  collisions, the triggers require at least one jet with  $p_T > 40, 60, \text{ and } 80 \text{ GeV}/c$ . The objects used in the HLT are jets reconstructed using the iterative-cone algorithm [34] with distance parameter  $R = 0.5$ . The soft background in PbPb collisions is removed with the iterative pileup subtraction technique [35]. In order to extend the reach of the jet spectra, data sets from the high- $p_T$  single-jet triggers are combined together in both  $pp$  and PbPb. To reach lower jet  $p_T$  in the PbPb data set, the minimum bias triggered events are added.

This analysis uses  $166 \mu\text{b}^{-1}$  of PbPb collisions at  $\sqrt{s_{NN}} = 2.76 \text{ TeV}$  recorded by CMS during the 2011 heavy ion run, as well as  $5.43 \text{ pb}^{-1}$  of  $pp$  collisions at the same collision energy recorded in early 2013. The event selection techniques developed for Ref. [20] are employed. These include the identification of a primary vertex and the removal of contamination from beam background, ultraperipheral and HCAL noise events. The primary reconstructed vertex of selected events in the  $z$  direction (beam axis) is constrained to be within  $\pm 15 \text{ cm}$  of the center of the detector. After these selections, events with more than one PbPb collision occurring in the same beam crossing remain and are later referred to as pileup. Utilizing the sensitivity of the ZDC to spectator nucleons and of the HF to particles produced in the collisions, these pileup events (0.2%) are removed by comparing the energy deposited in the ZDC to the HF. This is further substantiated by counting the number of fully reconstructed jets with  $p_T > 50 \text{ GeV}/c$  and comparing this to the number of tracker pixel hit counts, since pileup events tend to have large pixel counts for the same number of jets. The selection for pileup events in data does not remove any events from the simulation. This procedure was checked by individually studying a representative sample of the rejected events.

Simulated dijet events are generated using PYTHIA 6.4.23 Tune Z2 [36] for  $pp$  collisions at 2.76 TeV center-of-mass energy. For comparison to PbPb data, these PYTHIA events are embedded into a simulated PbPb event, generated by HYDJET (version 1.8) [37]. The HYDJET simulations are generated with jet quenching enabled in order to match the distribution of high- $p_T$  jets in a minimum-bias data set. The HYDJET simulations are tuned to represent a minimum bias background measured in CMS collisions of PbPb at  $\sqrt{s_{NN}} = 2.76 \text{ TeV}$ . Collision centrality is classified with the standard CMS heavy ion technique [20] using the total sum of the transverse energy in the HF towers, divided in percentiles according to the minimum bias samples. This distribution is divided into centrality bins, each representing 0.5% of the total nucleus-nucleus interaction cross section. For this analysis, the results are collected in six bins corresponding to the most central (i.e., smallest impact parameter) 5% of the events, denoted 0%–5%, as well as bins of 5%–10%, 10%–30%, 30%–50%, 50%–70%, and 70%–90%. The centrality of an event can be correlated with the impact parameter, as well as with  $\langle N_{part} \rangle$ , the average number of nucleons in the nuclei that participate in the collision, using MC Glauber model calculations [32].

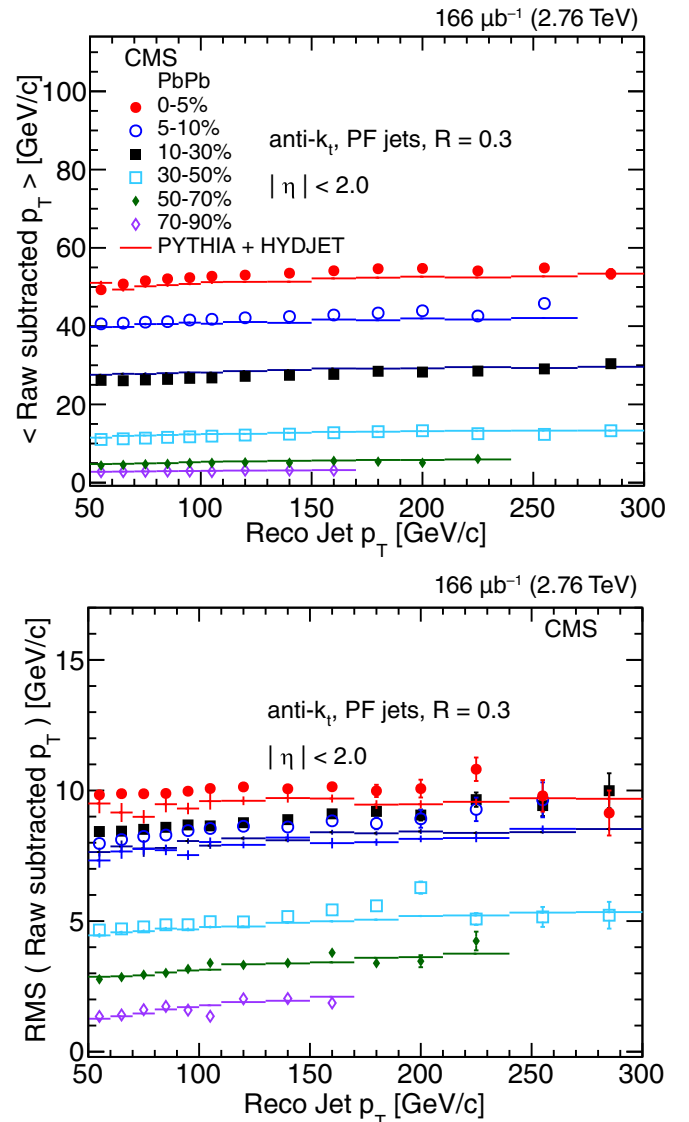


FIG. 2. Average raw subtracted  $p_T$  (top) and its rms (bottom) for PF jets reconstructed with the anti- $k_T$  algorithm, with a distance parameter  $R = 0.3$ . Symbols represent data, and lines show PYTHIA + HYDJET simulated events.

### III. JET RECONSTRUCTION AND SELECTION

Similar to Refs. [17,18,20,38], jet reconstruction in heavy ion collisions in CMS is performed with the sequential anti- $k_T$  clustering algorithm via the FASTJET framework [31]. The jet clustering is performed using particle-flow (PF) [39,40] candidates that combine information from the individual CMS detector systems. Different particle types (charged and neutral hadrons, electrons, muons, and photons) are reconstructed. The anti- $k_T$  distance parameters used are  $R = 0.2, 0.3, \text{ and } 0.4$ .

For PbPb collisions, the soft underlying event (background) is removed from the jets with an iterative subtraction technique described in Ref. [35]. In this procedure, the PF candidates are grouped in towers that correspond to the calorimeter geometry. Jets are selected with  $|\eta| < 2$  to ensure that they are fully contained within the CMS tracker up to a distance parameter of 0.4. Detector-based  $\eta$  and  $p_T$  dependent energy correction

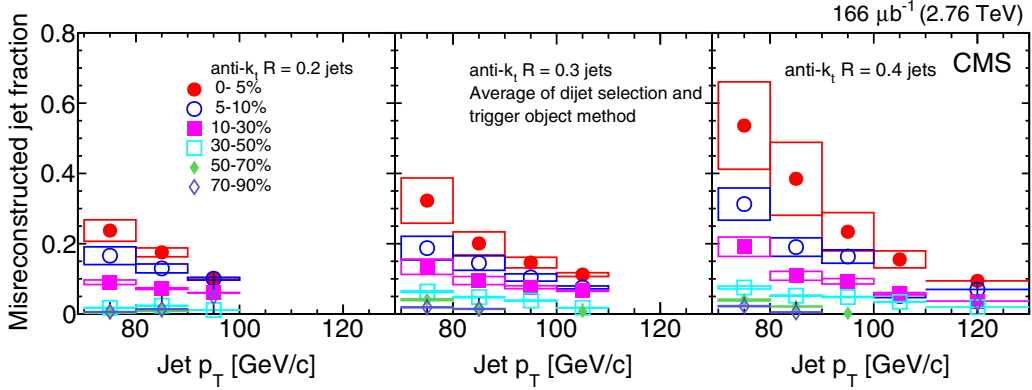


FIG. 3. Misreconstructed jet fraction of the inclusive jet spectra, derived from the minimum bias sample, as a function of reconstructed jet  $p_T$ , for various centralities and three different distance parameters (left:  $R = 0.2$ , center:  $R = 0.3$ , and right:  $R = 0.4$ ). The correction factor is the average of the dijet selection and trigger object methods discussed in the text.

factors [41] are applied to the jets. The raw jet  $p_T$  of a jet is the  $p_T$  before any of the detector-based corrections are applied. To study the background in PbPb events, data and PYTHIA + HYDJET simulations are compared. The correction to the jet  $p_T$  obtained from this iterative subtraction technique (called “raw subtracted  $p_T$ ”), for a jet with distance parameter  $R^{\text{jet}}$  is estimated by taking the difference between the sum of all the PF candidate  $p_T$  in a  $\Delta R < R^{\text{jet}}$  cone and the raw jet  $p_T$ . The  $\Delta R$  is defined as the distance of the PF candidate from the reconstructed jet axis in the  $\eta$ - $\phi$  plane:

$$\Delta R = \sqrt{(\Delta\phi_{\text{candidate, jet}})^2 + (\Delta\eta_{\text{candidate, jet}})^2}. \quad (2)$$

The distributions of raw subtracted  $p_T$  for  $R = 0.3$  jets, from peripheral to central collisions are shown in Fig. 1 for two different reconstructed jet  $p_T$  selections. Data are shown with filled circles and simulations with histograms. There is a good agreement between the two in all centralities and jet  $p_T$  bins. A similar level of agreement is also seen for  $R = 0.2$  and  $R = 0.4$ .

The average raw subtracted  $p_T$  and its root mean square (rms) values are shown in Fig. 2 as a function of the

reconstructed jet  $p_T$ , from central to the most peripheral collisions. Data are shown with markers and are compared with the PYTHIA + HYDJET generated events shown as histograms. The average raw subtracted  $p_T$  decreases, from the most central to peripheral events, as expected, and distributions show reasonable agreement between data and PYTHIA + HYDJET.

#### A. Data driven correction

Although the soft background is primarily removed with the iterative-pileup subtraction, fluctuations in this background can result in misreconstructed jets that do not originate from hard scattering. A method to remove this contamination, used in other experiments [26,27], is to select jets with a requirement on the leading charged-particle track or calorimeter energy deposit among the constituents of the jet. However, this method can bias to preferentially select jets with hard fragmentation, distorting the low- $p_T$  region. In CMS, tracks are reconstructed with a minimum  $p_T$  of 0.15 GeV/c, thus removing any such potential bias.

TABLE I. Summary of the systematic uncertainties in the PbPb jet yield for the central (0–5%), peripheral (70–90%) bins, and the  $pp$  jet cross section. Each column showcases the total systematic uncertainties for the corresponding source for the different  $R$  and two jet  $p_T$  ranges, i.e.,  $70 < \text{jet } p_T < 80$  [GeV/c] and  $250 < \text{jet } p_T < 300$  [GeV/c]. The  $T_{AA}$  uncertainties are not shown in the table. Other sources mentioned in the text that are smaller than 1% are not listed explicitly below.

Source		$70 < \text{jet } p_T < 80$ [GeV/c]			$250 < \text{jet } p_T < 300$ [GeV/c]		
		$R = 0.2$	$R = 0.3$	$R = 0.4$	$R = 0.2$	$R = 0.3$	$R = 0.4$
PbPb: (0–5%)	Data driven correction	13%	20%	27%	...	...	...
	JES & unfolding	32%	32%	48%	19%	19%	21%
	JER	3%	3%	3%	3%	3%	3%
	Underlying event	5%	5%	5%	...	...	...
PbPb: (70–90%)	Data driven correction	8%	10%	12%	...	...	...
	JES & unfolding	16%	16%	18%	...	...	...
	JER	3%	3%	3%	...	...	...
	Underlying event	5%	5%	5%	...	...	...
$pp$ :	JES & unfolding	7%	7%	6%	5%	4%	5%
	JER	3%	3%	3%	2%	2%	2%
	Integrated luminosity	3.7%	3.7%	3.7%	3.7%	3.7%	3.7%

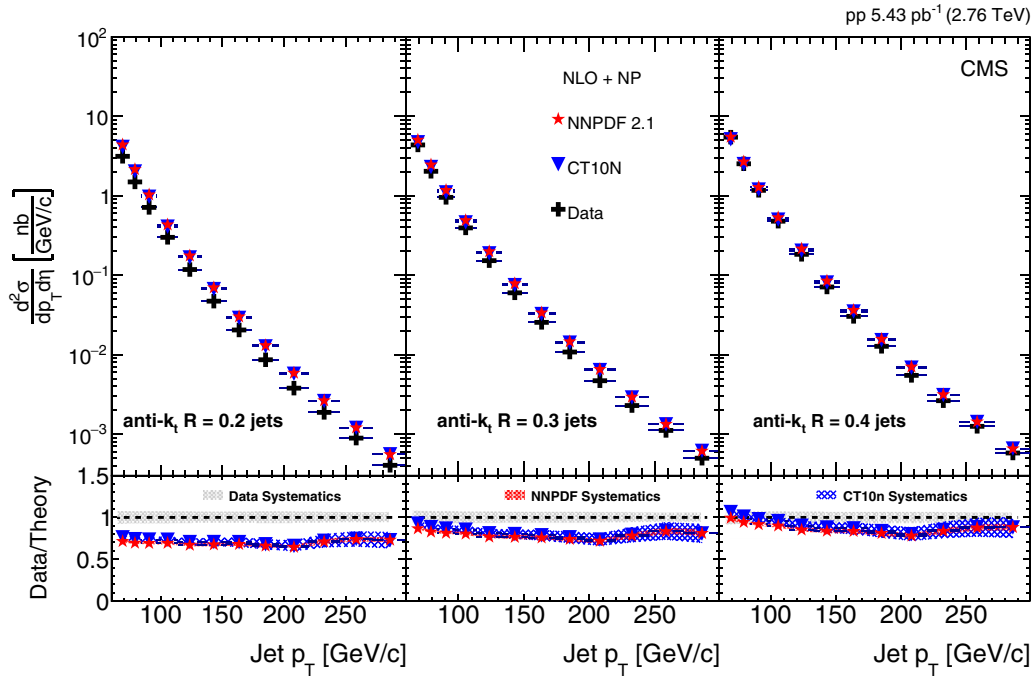


FIG. 4. Comparison of the inclusive jet cross section for anti- $k_T$  jets with distance parameters of  $R = 0.2$  (left),  $0.3$  (middle), and  $0.4$  (right), measured for  $pp$  collisions at 2.76 TeV (black plus markers), and NLO calculations, at the same collision energy, with NNPDF 2.1 (red star) and CT10N (blue triangle), with their respective NP corrections added. The bottom panels show the ratio of measured cross section to theory calculations. The systematic uncertainties for data are shown in the gray shaded band, while the systematic uncertainties in the NLO calculations are shown with the respective color shaded bands.

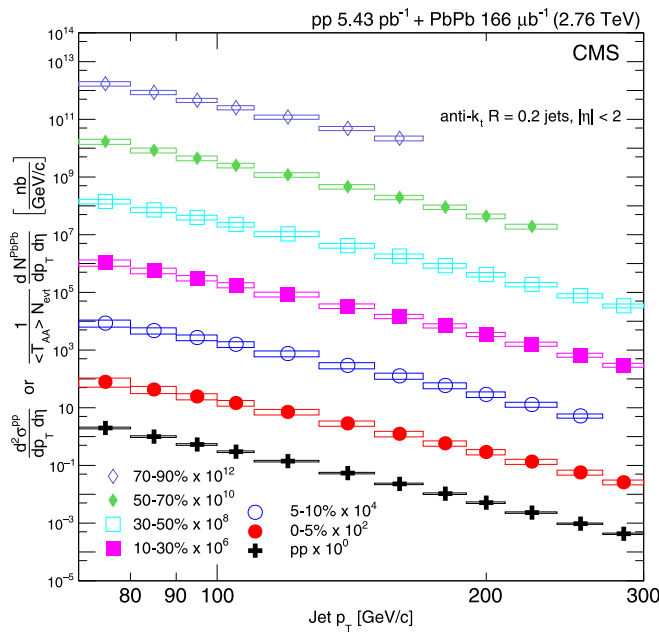


FIG. 5. Inclusive jet spectra for PbPb jets of distance parameter  $R = 0.2$ , in different centrality bins, and  $pp$  reference data. The PbPb jet spectra for different centrality classes are scaled by  $\langle T_{AA} \rangle$  and multiplied by a different factor for better visualization. Vertical bars represent statistical uncertainty (too small to see on this scale) with the systematical uncertainty in the colored boxes around the data points.

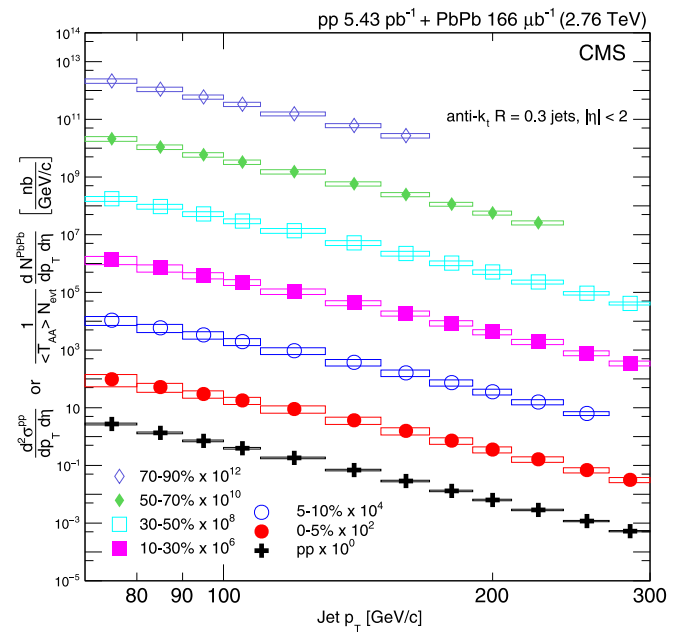


FIG. 6. Inclusive jet spectra for PbPb jets of distance parameter  $R = 0.3$ , in different centrality bins, and  $pp$  reference data. The PbPb jet spectra for different centrality classes are scaled by  $\langle T_{AA} \rangle$  and multiplied by a different factor for better visualization. Vertical bars represent statistical uncertainty (too small to see on this scale) with the systematical uncertainty in the colored boxes around the data points.



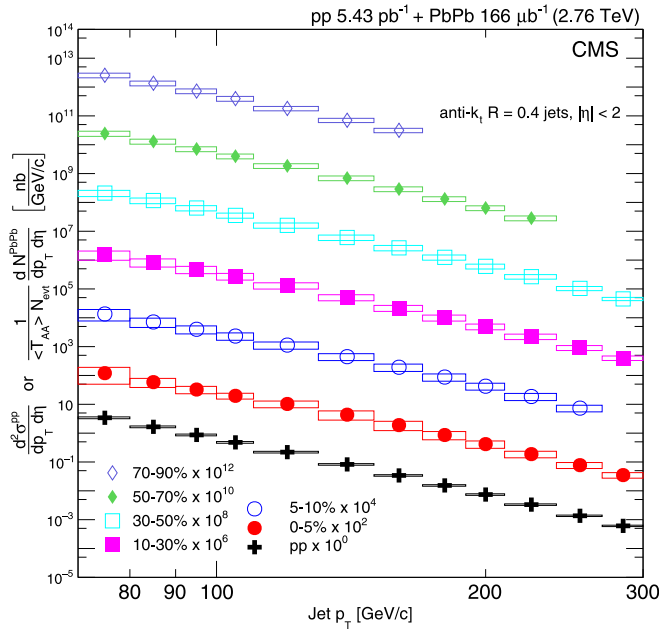


FIG. 7. Inclusive jet spectra for PbPb jets of distance parameter  $R = 0.4$ , in different centrality bins, and  $pp$  reference data. The PbPb jet spectra for different centrality classes are scaled by  $\langle T_{AA} \rangle$  and multiplied by a different factor for better visualization. Vertical bars represent statistical uncertainty (too small to see on this scale) with the systematical uncertainty in the colored boxes around the data points.

In this analysis, a novel data-driven technique, based on control regions in data, is introduced to derive the spectrum of misreconstructed jets from the minimum bias sample. This spectrum is then subtracted from the jet-triggered sample. Two methods, operating in different kinematic regimes, are combined to get a correction factor. The first method (labeled the trigger object method) selects all events with a leading HLT jet  $p_T$  of less than 60 GeV/c as a control sample potentially containing misreconstructed jets. This  $p_T$  threshold is chosen based on analysis of random cones in minimum bias events, with the leading and subleading jets removed. The second method (labeled the dijet method), performed in parallel with the first method, selects minimum bias events with dijets, which can originate either from a hard scattering or fluctuating background. There are two thresholds defined in this method, one for the leading jet ( $p_T^{\text{min1}}$ ) and another for the subleading jet ( $p_T^{\text{min2}}$ ) in the reconstructed event. If an event fails any of the following selections, it is tagged as a background event. An event is tagged as a signal if it passes all of the criteria: Leading jet  $p_T > p_T^{\text{min1}}$  and  $\Delta\phi_{j1,j2} > 2\pi/3$  and subleading jet  $p_T > p_T^{\text{min2}}$ . To choose the thresholds for the dijet selection, the mean and rms of the subtraction step in the iterative subtraction algorithm are mimicked by applying a cutoff on the transverse energies of the PF towers used in the random cone study. The rms of the background subtracted event energy distribution is used as an estimate of the fluctuation. The thresholds are set as follows:  $p_T^{\text{min1}} = 3 \text{ rms}$  for the leading

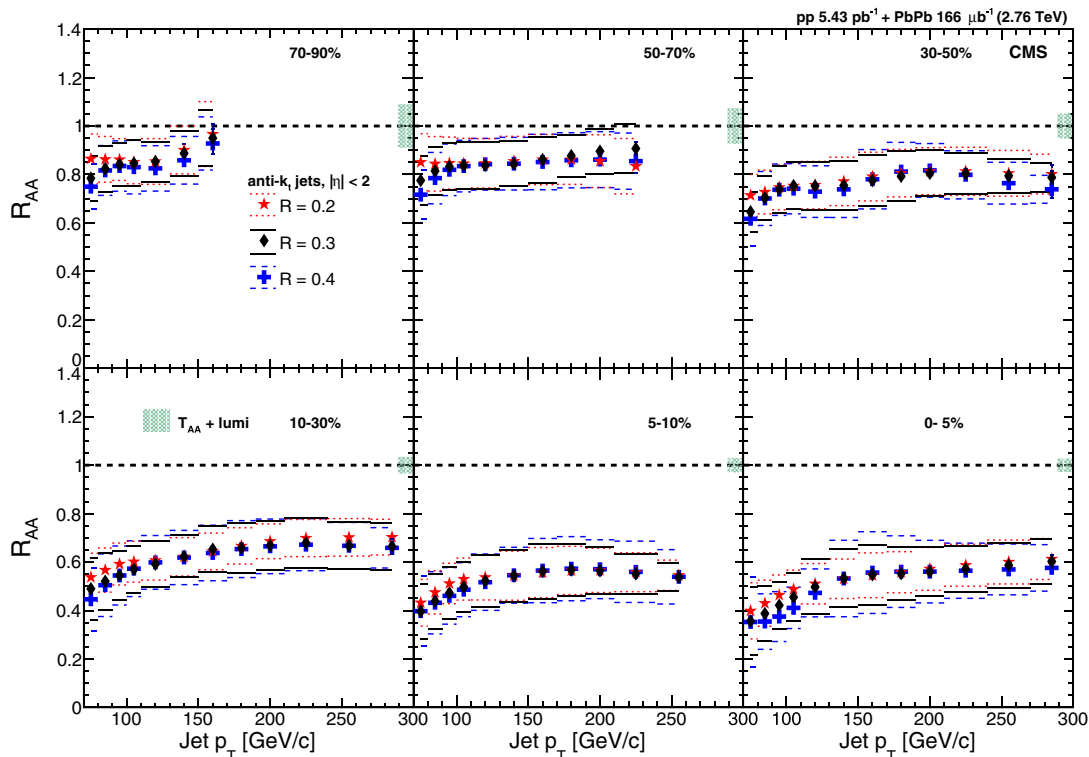


FIG. 8. Inclusive jet  $R_{AA}$  as a function of the jet  $p_T$ , for anti- $k_T$  jets with distance parameters  $R = 0.2$  (red stars),  $0.3$  (black diamonds), and  $0.4$  (blue crosses) for different centrality bins. The vertical bars (smaller than the markers) indicate the statistical uncertainty and the systematic uncertainty is represented by the bounds of the dotted, solid, and dashed horizontal lines. The uncertainty boxes at unity represent the  $T_{AA}$  and luminosity uncertainty.

jet, and  $p_T^{\text{min}2} = 1.8$  rms for the subleading jet, to allow for jet modification in the medium.

Since these two methods operate in different kinematic regimes, the average of the two is used to estimate the data driven correction factor for misreconstructed jet rates as can be seen in Fig. 3, as a function of the jet  $p_T$ . These rates for different distance parameters are shown in the different panels (left:  $R = 0.2$ , center:  $R = 0.3$ , and right:  $R = 0.4$ ). The symbols correspond to the centrality bins in the analysis. The minimum bias background jet spectra are then normalized to a per-event yield and the background is removed from the measured jet spectra, resulting in an inclusive jet spectrum without fragmentation bias. The correction, estimated in a similar way from PYTHIA dijet events, where one does not expect any background, is added as an additional systematic uncertainty, ranging from 6% at 70 GeV to 1% at 100 GeV. The data driven method was also applied to PYTHIA+HYDJET simulations without quenching and, using the same  $p_T$  threshold, this yielded a recovery efficiency of greater than 98% for signal jets, which is well within systematic uncertainties as described in Sec. IV.

### B. Unfolding studies

An unfolding method is required to remove the smearing and bin migration in jet  $p_T$  due to detector resolution, and to extract the jet cross section measurement. Three different techniques are used to determine the final jet  $p_T$  spectra: Single value decomposition (SVD), Bayesian, and a bin-by-bin unfolding technique [42–46]. Results presented here are based on the SVD technique, while the others are used as a crosscheck, giving consistent results within their respective uncertainties. The three aforementioned procedures use a response matrix from PYTHIA + HYDJET of reconstructed jets, matched to generator-level jets in the  $\eta$ - $\phi$  space, that originate from the PYTHIA QCD hard scattering.

The SVD unfolding is performed with a regularization parameter, which is optimized for each centrality class and each jet resolution. The simulation and data used in unfolding have a reconstructed jet  $p_T$  larger than 50 GeV/ $c$  for all distance parameters, with unfolded results reported for jets larger than 70 GeV/ $c$ .

### IV. SYSTEMATIC UNCERTAINTIES

The systematic uncertainty is calculated from a number of sources and is shown in Table I. For  $R = 0.3$  jets, in the low  $p_T < 80$  GeV/ $c$  region, a large contribution to the jet yield uncertainty in PbPb collisions is from the data driven corrections (20%). The data driven systematic uncertainty is estimated from the overlap of the two different methods (trigger object and dijet methods as described in Sec. III A) along with an additional uncertainty of 1–6% across all jet  $p_T$ , centrality ranges, and jet distance parameters determined from its application on a PYTHIA sample. The jet energy scale (JES) uncertainty ranges from 6–32% (from peripheral to central events), varying due to the uncertainty in the heavy ion tracking and the quark/gluon fragmentation. The fragmentation difference is included in the JES uncertainty for  $pp$ , but is extended for PbPb jets due to expected asymmetric jet quenching

effects for quark and gluon jets. The jet response matrix is smeared by 1%, at both the generator and reconstructed levels to account for variations in the simulations. Separately the regularization parameter used for the unfolding is varied between 4 and 8 resulting in at most 8% systematic uncertainty for the PbPb jet yield and at most 2% for the  $pp$  jet cross section.

A residual jet energy correction, using the dijet balance method [41], is derived and applied to the jets from  $pp$  collisions. It corresponds to less than 1% correction to the jet  $p_T$ . The jet energy resolution (JER) uncertainty is estimated for each  $p_T$  bin in the analysis and is found to be at most 3%, for both  $pp$  and PbPb. Studies of the underlying event fluctuations

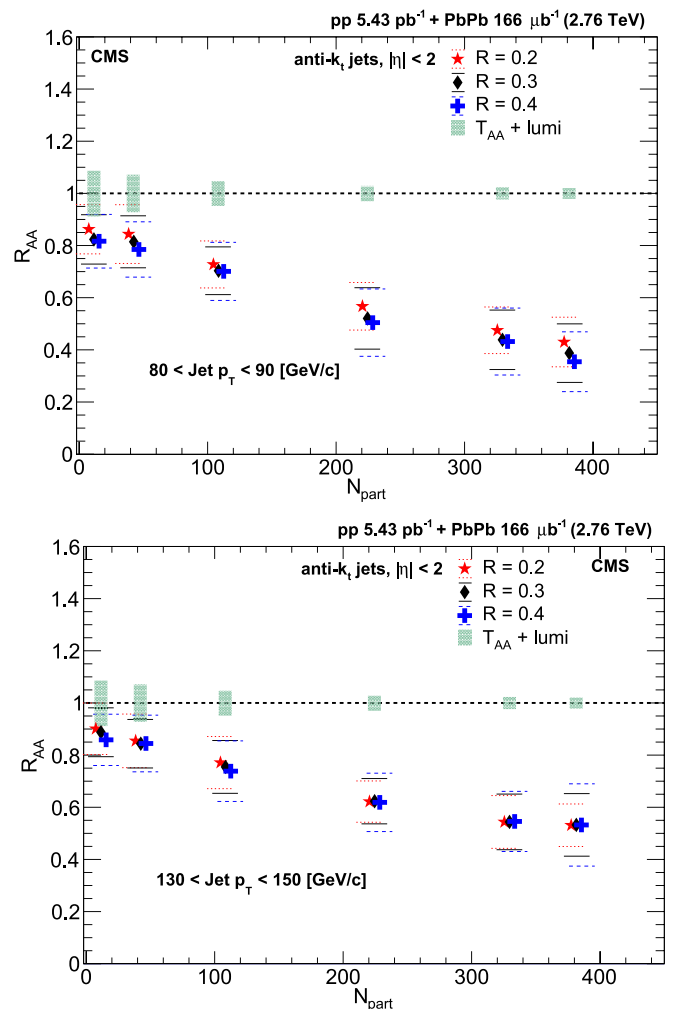


FIG. 9. Inclusive jet  $R_{AA}$  for anti- $k_T$  jets with distance parameters  $R = 0.2$  (red stars),  $0.3$  (black diamonds), and  $0.4$  (blue crosses), as a function of the average  $N_{\text{part}}$  for each collision centrality, for jets of  $80 < p_T < 90$  and  $130 < p_T < 150$  [GeV/ $c$ ], in the top and bottom panels, respectively. Points are shifted to the left ( $R = 0.2$ ) and right ( $R = 0.4$ ) for clarity. The statistical uncertainty is indicated by colored vertical lines (smaller than the markers). The systematic uncertainty is represented by the bounds of the dotted, solid, and dashed horizontal lines for the corresponding distance parameters. The uncertainty boxes at unity represent the  $T_{AA}$  and luminosity uncertainty.



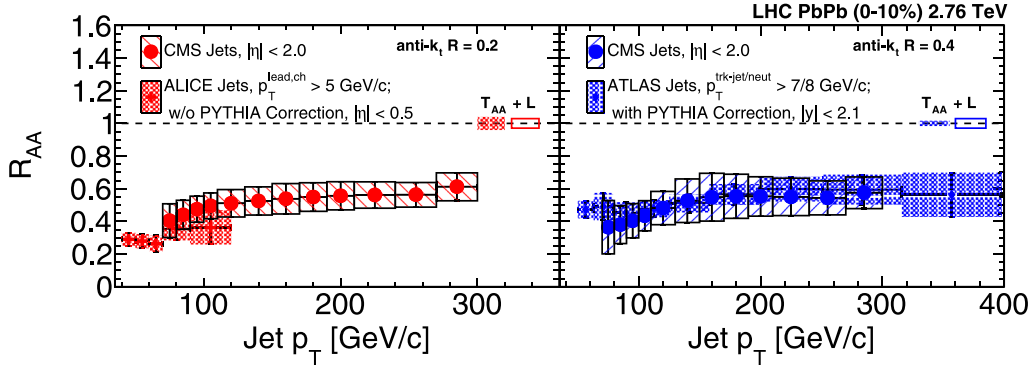


FIG. 10. Left panel: Inclusive jet  $R_{AA}$  as a function of the jet  $p_T$ , for anti- $k_T$  jets with distance parameter  $R = 0.2$  in the 0%–10% centrality bin for CMS (closed circles) and ALICE (pluses) [27]. Right panel: Inclusive jet  $R_{AA}$  as a function of the jet  $p_T$ , for anti- $k_T$  jets with distance parameter  $R = 0.4$  in the 0%–10% centrality bin for CMS (closed circles) and ATLAS (diamonds) [26]. The vertical bars indicate the statistical uncertainty. The systematic uncertainty is represented by the bounds of the boxes. The uncertainty boxes at unity represent the  $T_{AA}$  and luminosity uncertainty, open for CMS and shaded for ALICE and ATLAS. See text for a further discussion of differences in the analyses used by the three collaborations.

in jet-triggered and minimum bias events show a contribution of up to 5% to the uncertainty of reconstructed jet yields based on differences between data and PYTHIA + HYDJET quantified in the right side of Fig. 2. The contributions due to jet reconstruction efficiency, detector noise, and unfolding response matrix smearing are about 1% each.

Since in PbPb, the per-event jet yield is being measured, there is a 3% uncertainty on the number of minimum bias events and there is no uncertainty quoted for the luminosity. For the  $pp$  cross section, there is a 3.7% uncertainty in the integrated luminosity [47]. Systematic uncertainties, from different contributions to the jet  $R_{AA}$ , are summed in quadrature with an overall uncertainty of 19–40%, from peripheral to central collisions for  $R = 0.3$  jets. Detailed systematic uncertainties for different  $R$  and two representative jet  $p_T$  ranges are shown in Table I.

## V. RESULTS

The inclusive jet cross sections in  $pp$  collisions at 2.76 TeV are shown in Fig. 4 for three different distance parameters. A comparison is made to next-to-leading-order (NLO) [48] calculations of quantum chromodynamics. These calculations are shown for two parton distribution functions (PDF) sets: NNPDF 2.1 [49] (red stars), and CT10N [50] (purple triangles) including nonperturbative (NP) contributions such as multiparton interactions and hadronization. Contributions to the jet cross section from NP effects are not inherently included in pQCD calculations due to a lower scale cutoff of a few GeV/ $c$ . Thus, the NP correction factors need to be added and are computed as the ratio of cross sections calculated with leading order (LO) + parton shower (PS) + multiparton interactions + hadronization to LO+PS [48]. The bottom panel of Fig. 4 shows the ratio of the data for jet cross sections in  $pp$  collisions to theoretical calculations, with the measured jet cross section from  $pp$  collisions for different distance parameters. The agreement with data gets better at larger distance parameters. In Ref. [51] the ratio tends closer to unity for jets with  $R = 0.7$ . The theoretical uncertainties shown are due to variations of the

strong coupling constant and the parton shower, factorization scales involved in the NLO calculations for the different PDF sets.

The unfolded jet cross sections for PbPb and  $pp$  events are shown in Figs. 5–7 for different distance parameters. The PbPb spectra are normalized by the number of minimum bias events, and are scaled by  $\langle T_{AA} \rangle$ , with each centrality multiplied by a different factor, to separate the spectra for better visualization. The  $pp$  reference data are normalized to the integrated luminosity of the analyzed data set. The high  $p_T$  cutoffs for the spectra (hence also the  $R_{AA}$ ) are dictated by statistical limitations.

The jet  $R_{AA}$ , found from the PbPb and  $pp$  spectra after all corrections including SVD unfolding, are shown for different distance parameters in Fig. 8. The jet  $R_{AA}$  decreases with increasing collision centrality in the range of the measured jet  $p_T$ . Within the systematic uncertainty, the jet  $R_{AA}$  shows the same level of suppression for the three distance parameters. Uncorrelated uncertainties remain too large to further elucidate the hierarchy of the jet distance parameter dependence of this  $R_{AA}$  measurement.

To focus on the centrality dependence of the jet  $R_{AA}$ , two ranges of jet  $p_T$  are selected and the corresponding jet  $R_{AA}$  values are plotted as a function of the average number of participants ( $N_{part}$ ) in Fig. 9, for jets of  $80 < p_T < 90$  and  $130 < p_T < 150$  GeV/ $c$ . The systematic uncertainty is shown in the three bounds of lines for  $R = 0.2$  (dotted), 0.3 (solid), and 0.4 (dashed) jets. The jet  $R_{AA}$  shows a clear trend of increasing suppression as the number of participants in the PbPb collision increases. Overall, in the kinematic range explored, the  $R_{AA}$  show the same level of suppression across the three distance parameters.

An experimental comparison of inclusive anti- $k_T$  jet  $R_{AA}$  for 0–10% centrality is shown in Fig. 10 (left panel for anti- $k_T$  jets with distance parameter  $R = 0.2$  for ALICE [27] and the right panel with  $R = 0.4$  for ATLAS [26]). Uncertainties are represented by the vertical bars for the statistical and boxes for the systematic uncertainties. The  $T_{AA}$  and luminosity uncertainty are shown by the boxes at unity. The collection of

jets for the jet  $R_{AA}$  calculation in these experiments differ, especially for lower jet  $p_T$ , due to the techniques employed to remove or correct the jets that did not originate in a hard scattering but that are purely due to the fluctuations in the heavy-ion underlying event. Some, but not all of the key differences are described here, for more, see ALICE [27], ATLAS [26], and [52] for a review. ALICE requires the leading track constituent of the jet to have  $p_T > 5 \text{ GeV}/c$  and constrains  $R = 0.2$  jets to be within  $|\eta| < 0.9$ . ATLAS requires its  $R = 0.4$  jets in  $|y| < 2.1$  to have a track jet with  $p_T > 7 \text{ GeV}/c$  or a calorimeter cluster with  $p_T > 8 \text{ GeV}/c$  within  $\Delta R = 0.2$ . While ALICE does not apply any correction on this constituent selection, ATLAS corrects for the missing jets due to this selection with correction factors estimated by PYTHIA. In this analysis, as described in Sec. III A, a data-driven background subtraction is introduced and all jets which are using tracks down to a  $p_T$  of  $0.15 \text{ GeV}/c$  and calorimeter deposits down to a  $E_T$  of  $1 \text{ GeV}$  are included in the jet  $R_{AA}$  calculation. Within the current precision of jet  $R_{AA}$  measurements, there is a good agreement in the overlapping  $p_T$  ranges despite the fact that the measured jet collections differ between experiments.

## VI. SUMMARY

The cross section of anti- $k_T$  particle-flow jets has been measured in  $pp$  and PbPb collisions at  $\sqrt{s_{NN}} = 2.76 \text{ TeV}$  for distance parameters  $R = 0.2, 0.3, \text{ and } 0.4$  in  $|\eta| < 2$  and for jet  $p_T$  above  $70 \text{ GeV}/c$ . It is found that next-to-leading order calculations with nonperturbative corrections over predict the  $pp$  cross sections, with a smaller discrepancy for larger distance parameters. The PbPb inclusive jet nuclear modification factors show a steady decrease from peripheral to central events, with a slight rise with jet  $p_T$ . No significant dependence of the jet nuclear modification factor on the distance parameter is found for the jets in the kinematic range measured in this analysis.

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L. Scodellaro,<sup>110</sup> N. Trevisani,<sup>110</sup> I. Vila,<sup>110</sup> R. Vilar Cortabitarte,<sup>110</sup> D. Abbaneo,<sup>111</sup> E. Auffray,<sup>111</sup> G. Auzinger,<sup>111</sup>  
M. Bachtis,<sup>111</sup> P. Baillon,<sup>111</sup> A. H. Ball,<sup>111</sup> D. Barney,<sup>111</sup> P. Bloch,<sup>111</sup> A. Bocci,<sup>111</sup> A. Bonato,<sup>111</sup> C. Botta,<sup>111</sup> T. Camporesi,<sup>111</sup>

R. Castello,<sup>111</sup> M. Cepeda,<sup>111</sup> G. Cerminara,<sup>111</sup> M. D'Alfonso,<sup>111</sup> D. d'Enterria,<sup>111</sup> A. Dabrowski,<sup>111</sup> V. Daponte,<sup>111</sup> A. David,<sup>111</sup> M. De Gruttola,<sup>111</sup> A. De Roeck,<sup>111</sup> E. Di Marco,<sup>111,ar</sup> M. Dobson,<sup>111</sup> B. Dorney,<sup>111</sup> T. du Pree,<sup>111</sup> D. Duggan,<sup>111</sup> M. Dünser,<sup>111</sup> N. Dupont,<sup>111</sup> A. Elliott-Peisert,<sup>111</sup> S. Fartoukh,<sup>111</sup> G. Franzoni,<sup>111</sup> J. Fulcher,<sup>111</sup> W. Funk,<sup>111</sup> D. Gigi,<sup>111</sup> K. Gill,<sup>111</sup> M. Girone,<sup>111</sup> F. Glege,<sup>111</sup> D. Gulhan,<sup>111</sup> S. Gundacker,<sup>111</sup> M. Guthoff,<sup>111</sup> J. Hammer,<sup>111</sup> P. Harris,<sup>111</sup> J. Hegeman,<sup>111</sup> V. Innocente,<sup>111</sup> P. Janot,<sup>111</sup> J. Kieseler,<sup>111</sup> H. Kirschenmann,<sup>111</sup> V. Knünz,<sup>111</sup> A. 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Seidel,<sup>111</sup> A. Sharma,<sup>111</sup> P. Silva,<sup>111</sup> P. Sphicas,<sup>111,at</sup> J. Steggemann,<sup>111</sup> M. Stoye,<sup>111</sup> Y. Takahashi,<sup>111</sup> M. Tosi,<sup>111</sup> D. Treille,<sup>111</sup> A. Triossi,<sup>111</sup> A. Tsirou,<sup>111</sup> V. Veckalns,<sup>111,au</sup> G. I. Veres,<sup>111,s</sup> N. Wardle,<sup>111</sup> H. K. Wöhri,<sup>111</sup> A. Zagodzinska,<sup>111,ah</sup> W. D. Zeuner,<sup>111</sup> W. Bertl,<sup>112</sup> K. Deiters,<sup>112</sup> W. Erdmann,<sup>112</sup> R. Horisberger,<sup>112</sup> Q. Ingram,<sup>112</sup> H. C. Kaestli,<sup>112</sup> D. Kotlinski,<sup>112</sup> U. Langenegger,<sup>112</sup> T. Rohe,<sup>112</sup> F. Bachmair,<sup>113</sup> L. Bäni,<sup>113</sup> L. Bianchini,<sup>113</sup> B. Casal,<sup>113</sup> G. Dissertori,<sup>113</sup> M. Dittmar,<sup>113</sup> M. Donegà,<sup>113</sup> C. Grab,<sup>113</sup> C. Heidegger,<sup>113</sup> D. Hits,<sup>113</sup> J. Hoss,<sup>113</sup> G. Kasieczka,<sup>113</sup> P. Lecomte,<sup>113,av</sup> W. Lustermaan,<sup>113</sup> B. Mangano,<sup>113</sup> M. Marionneau,<sup>113</sup> P. Martinez Ruiz del Arbol,<sup>113</sup> M. Masciovecchio,<sup>113</sup> M. T. Meinhard,<sup>113</sup> D. Meister,<sup>113</sup> F. Micheli,<sup>113</sup> P. Musella,<sup>113</sup> F. Nessi-Tedaldi,<sup>113</sup> F. Pandolfi,<sup>113</sup> J. Pata,<sup>113</sup> F. Pauss,<sup>113</sup> G. Perrin,<sup>113</sup> L. Perrozzi,<sup>113</sup> M. Quittnat,<sup>113</sup> M. Rossini,<sup>113</sup> M. Schönenberger,<sup>113</sup> A. Starodumov,<sup>113,aw</sup> V. R. Tavolaro,<sup>113</sup> K. Theofilatos,<sup>113</sup> R. Wallny,<sup>113</sup> T. K. Aarrestad,<sup>114</sup> C. Amsler,<sup>114,ax</sup> L. Caminada,<sup>114</sup> M. F. Canelli,<sup>114</sup> A. De Cosa,<sup>114</sup> C. Galloni,<sup>114</sup> A. Hinzmann,<sup>114</sup> T. Hreus,<sup>114</sup> B. Kilminster,<sup>114</sup> J. Ngadiuba,<sup>114</sup> D. Pinna,<sup>114</sup> G. 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Seez,<sup>126</sup> S. Summers,<sup>126</sup> A. Tapper,<sup>126</sup> K. Uchida,<sup>126</sup> M. Vazquez Acosta,<sup>126,bm</sup> T. Virdee,<sup>126,n</sup> J. Wright,<sup>126</sup> S. C. Zenz,<sup>126</sup> J. E. Cole,<sup>127</sup> P. R. Hobson,<sup>127</sup> A. Khan,<sup>127</sup> P. Kyberd,<sup>127</sup> D. Leslie,<sup>127</sup> I. D. Reid,<sup>127</sup> P. Symonds,<sup>127</sup> L. Teodorescu,<sup>127</sup> M. Turner,<sup>127</sup> A. Borzou,<sup>128</sup> K. Call,<sup>128</sup> J. Dittmann,<sup>128</sup> K. Hatakeyama,<sup>128</sup> H. Liu,<sup>128</sup> N. Pastika,<sup>128</sup> O. Charaf,<sup>129</sup> S. I. Cooper,<sup>129</sup> C. Henderson,<sup>129</sup> P. Rumerio,<sup>129</sup> C. West,<sup>129</sup> D. Arcaro,<sup>130</sup> A. Avetisyan,<sup>130</sup> T. Bose,<sup>130</sup> D. Gastler,<sup>130</sup> D. Rankin,<sup>130</sup> C. Richardson,<sup>130</sup> J. Rohlf,<sup>130</sup> L. Sulak,<sup>130</sup> D. Zou,<sup>130</sup> G. Benelli,<sup>131</sup> E. Berry,<sup>131</sup> D. Cutts,<sup>131</sup> A. Garabedian,<sup>131</sup> J. Hakala,<sup>131</sup> U. Heintz,<sup>131</sup> J. M. Hogan,<sup>131</sup> O. Jesus,<sup>131</sup> K. H. M. Kwok,<sup>131</sup> E. Laird,<sup>131</sup> G. Landsberg,<sup>131</sup> Z. Mao,<sup>131</sup> M. Narain,<sup>131</sup> S. Piperov,<sup>131</sup> S. Sagir,<sup>131</sup> E. Spencer,<sup>131</sup> R. Syarif,<sup>131</sup> R. Breedon,<sup>132</sup> G. Breto,<sup>132</sup> D. Burns,<sup>132</sup> M. Calderon De La Barca Sanchez,<sup>132</sup> S. Chauhan,<sup>132</sup> M. Chertok,<sup>132</sup> J. Conway,<sup>132</sup> R. Conway,<sup>132</sup> P. T. Cox,<sup>132</sup> R. Erbacher,<sup>132</sup> C. Flores,<sup>132</sup> G. Funk,<sup>132</sup> M. Gardner,<sup>132</sup> W. Ko,<sup>132</sup> R. Lander,<sup>132</sup> C. Mclean,<sup>132</sup> M. Mulhearn,<sup>132</sup> D. Pellett,<sup>132</sup> J. Pilot,<sup>132</sup> S. Shalhout,<sup>132</sup> J. Smith,<sup>132</sup> M. Squires,<sup>132</sup> D. Stolp,<sup>132</sup> M. Tripathi,<sup>132</sup> S. Wilbur,<sup>132</sup> R. Yohay,<sup>132</sup> R. Cousins,<sup>133</sup> P. Everaerts,<sup>133</sup> A. Florent,<sup>133</sup> J. Hauser,<sup>133</sup> M. Ignatenko,<sup>133</sup> D. Saltzberg,<sup>133</sup> E. Takasugi,<sup>133</sup> V. Valuev,<sup>133</sup> M. Weber,<sup>133</sup> K. Burt,<sup>134</sup> R. Clare,<sup>134</sup> J. Ellison,<sup>134</sup> J. W. Gary,<sup>134</sup> S. M. A. Ghiasi Shirazi,<sup>134</sup> G. Hanson,<sup>134</sup> J. Heilman,<sup>134</sup> P. Jandir,<sup>134</sup> E. Kennedy,<sup>134</sup> F. Lacroix,<sup>134</sup> O. R. Long,<sup>134</sup> M. Olmedo Negrete,<sup>134</sup> M. I. Paneva,<sup>134</sup> A. Shrinivas,<sup>134</sup> W. Si,<sup>134</sup> H. Wei,<sup>134</sup> S. Wimpenny,<sup>134</sup> B. R. Yates,<sup>134</sup> J. G. Branson,<sup>135</sup> G. B. Cerati,<sup>135</sup> S. Cittolin,<sup>135</sup> M. Derdzinski,<sup>135</sup> R. Gerosa,<sup>135</sup> A. Holzner,<sup>135</sup> D. 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Yoo,<sup>136</sup> D. Anderson,<sup>137</sup> A. Apresyan,<sup>137</sup> J. Bendavid,<sup>137</sup> A. Bornheim,<sup>137</sup> J. Bunn,<sup>137</sup> Y. Chen,<sup>137</sup> J. Duarte,<sup>137</sup> J. M. Lawhorn,<sup>137</sup> A. Mott,<sup>137</sup> H. B. Newman,<sup>137</sup> C. Pena,<sup>137</sup> M. Spiropulu,<sup>137</sup> J. R. Vlimant,<sup>137</sup> S. Xie,<sup>137</sup> R. Y. Zhu,<sup>137</sup>

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