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# Constraints on Lunar Crustal Porosity from the Gravitational Signature of Impact Craters

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#### **Key Points:**

• Crater gravity anomalies are useful for constraining spatial variations in the crustal porosity.

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- Western nearside maria and the southern South Pole-Aitken (SP-A) basin are associated with relatively low crustal porosity.
- Southwestern periphery of the nearside maria and surroundings of most large impact basins show relatively high crustal porosity.

#### Abstract

Based on Gravity Recovery and Interior Laboratory (GRAIL) observations and the Lunar Orbiter Laser Altimeter (LOLA) crater database, we constrain the spatial variations in the Moon's crustal porosity ( $\phi_c$ ) to a depth of several kilometers using the gravity anomalies of 4,864 mid-sized craters (those with a diameter of 20–100 km). For each crater, we estimated the local  $\phi_c$  by quantifying the gravitational effects of impact-induced porosity change and postimpact breccia infill. The crustal porosity model was uniquely determined assuming a global mean  $\phi_c$  of 12%, although a trade-off exists between the porosity of post-impact breccia infill and the (pre-impact) crustal porosity that results in zero net porosity change underlying the crater floor. At this crustal porosity the effects of impact bulking and compaction compensate each other to yield zero crater residual Bouguer anomaly (RBA), when the effect of post-impact infill is excluded. For lower crustal porosities, bulking dominates to produce a negative RBA; for higher crustal porosities, compaction dominates to produce a positive RBA. Spatial kriging and statistical tests suggest lower-than-average  $\phi_c$  in the western nearside maria and southern South Pole-Aitken (SP-A) basin, in contrast to higher-than-average values in the southwestern periphery of the nearside maria. Most of the large impact basins, presumably formed before the majority of the mid-sized craters we analyzed, show reduced porosities relative to their immediate surroundings.

Because the Moon lacks plate tectonics and surface erosion, its crustal structure has been preserved since the Late Heavy Bombardment ~ 3.9 billion years ago. The porosity structure of the lunar crust provides critical information on the early cratering history and the formation of major impact basins. Anomalously low crustal porosity implies the existence of relatively young volcanic deposits. In addition, the early porosity evolution of the Moon sheds lights on the origin of life as similar processing acting on the Earth created heat and pore space (as water conduits) to support life on Earth. This study uses 4,864 mid-sized craters as probes into the internal porosity structure of the lunar crust, which is one of the primary objectives of the Gravity Recovery and Interior Laboratory mission. For each mid-sized crater, we inverted the observed crater gravity signature for local crustal porosity. We then applied spatial statistical technique to map the regional-scale variations in the crustal porosity and related the porosity variations with the major geologic units, including the nearside maria, highlands, South Pole-Aitken basin and other impact basins. This study shows that mid-sized craters are useful for detecting regional-scale porosity structure that has long been established in the lunar evolution history.

#### **1** Introduction

Spatial variations in the crustal porosity of the Moon provide critical information for its thermal evolution and cratering history, which can be extrapolated to the terrestrial planets. Variations in the porosity of the lunar highland crust provide insight into the Moon's cumulative cratering history since the pre-Nectarian period, the thermal conductivity of the crust and global

thermal evolution (Ziethe et al., 2009; Laneuville et al., 2013), and the interpretation of heat flux measurements (Langseth et al., 1976; Rasmussen & Warren, 1985; Siegler & Smrekar, 2014) and seismic profiles (Lognonné et al., 2003; Khan et al., 2013). Crustal porosity also influences the chemical and mechanical reaction rates that control geological and ecological evolution (Navarre-Sitchler & Brantley, 2007). Spatial variations in the crustal porosity of the lunar highlands, however, were poorly constrained until the availability of the high quality and resolution lunar gravity model provided by the GRAIL mission (Goossens et al., 2015). Wieczorek et al. (2013) attributed the observed gravity over spherical harmonic degrees 150–310 to surface relief, and solved for crustal density and porosity variations by minimizing the correlation between the topography and Bouguer gravity over 6°-radius cap windows.

An alternative method to constrain the lunar crustal porosity is to use the gravitational signature of impact craters (Dvorak & Phillips, 1977, 1979; Soderblom et al., 2015; Bierson et al., 2016). Impact craters into typical low to moderate porosity surfaces will increase porosity through brecciation, fracturing and dilatancy (Pilkington & Grieve, 1992; Collins, 2014), whereas impacts into relatively high porosity targets will reduce porosity via localized heating and compaction (Melosh, 1989; Milbury et al., 2015a). The net porosity change determines the observed crater gravity anomaly, or residual Bouguer anomaly (RBA, defined as area-weighted average Bouguer anomaly interior to the crater rim with respect to the average Bouguer anomaly

within a background annulus, see 2.1). A regional RBA value of ~ 0 mGal implies a crustal porosity  $\phi_r$ , in which the impact bulking and compaction are in balance (i.e., gravitational effects of the co-existing compacted and bulked crustal structures compensate each other, Milbury et al., 2015a). Regions of the crust with such a porosity will experience minimal impactrelated porosity change with time and so can be regarded as in a steady or equilibrium state for a given impactor size-frequency distribution. A large-magnitude non-zero RBA indicates deviation from  $\phi_r$  (ignoring post-impact effects), which implies that the regional crust has not yet reached the steady state. Milbury et al. (2015a) used iSALE impact hydrocode simulations to establish a linear relationship between the crater RBA and the background crustal porosity, yielding a  $\phi_r$  of ~ 7%. Soderblom et al. (2015) compared the observed crater RBA values with local crustal porosities from Wieczorek et al. (2013), supporting this linear relationship but yielding an apparent  $\phi_r$  of ~ 15%. We show that the inconsistency in  $\phi_r$  is largely due to the gravitational effects of highly porous post-impact breccia infill on the observed crater RBA, and therefore reconcile the hydrocode models and observations.

In this study, we used the linear relationship between crater RBA and background crustal porosity, after considering the effects of Bouguer correction and post-impact breccia infill, to estimate local pre-impact ambient crustal porosities ( $\phi_c$ ) for 4,864 craters identified in Lunar Orbiter Laser Altimeter (LOLA) data (Head et al., 2010). We then spatially interpolated  $\phi_c$  to map their spatial variations and applied Student's t-test to confirm anomalous regions at a spatial resolution of 5° for the highlands and 10° for the western nearside maria. We also interpolated

the ratio RBA/ $D_c$  (where  $D_c$  is the crater diameter) to directly show spatial variations in the crater gravity anomaly, where dividing the crater RBA by  $D_c$  removes the crater-size dependence (see Methodology section). It is worth noting that the majority of the observed mid-sized craters (diameters of 20–100 km) must postdate the early formation of regional-scale crustal structure, including the major impact basins (Fassett et al., 2012). In addition, impact-induced porosity changes under the crater floors for these mid-sized craters are distributed almost completely within the crater rims (Collins, 2014; Milbury et al., 2015a). Therefore, the pre-existing regional porosity structure is unlikely to be influenced significantly by these mid-sized craters. The midsized craters can thus be used as probes of the regional porosity structure ( $\phi_c$ ) that was long been established early in the evolutionary history of the Moon.

# 2 Methodology

We first calculated the RBA values for 4,864 craters with a crater diameter ( $D_c$ ) of 20–100 km from the LOLA crater database (Head et al., 2010; Kadish et al., 2011) (Figure 1a). Following Soderblom et al. (2015), the crater RBA is defined as the area-weighted average Bouguer anomaly interior to the crater rim minus the average Bouguer anomaly within a background annulus that extends radially from the outer rim flank (Pike, 1977) (Table 1) to one crater radius beyond the crater rim. The calculated crater RBA was found not to be biased by the existence of ejecta deposits by testing the sensitivity of the crater RBA to the radius of background annulus (Soderblom et al., 2015). We used the GRAIL Bouguer gravity model JGGRAIL\_1200C12A\_BOUGUER (Goossens et al., 2015), which was derived by assuming a

uniform Bouguer correction density ( $\rho_0$ ) of 2,550 kg/m<sup>3</sup> (corresponding to a nominal mean crustal porosity of 12%). We filtered the Bouguer gravity model to include spherical harmonic degrees 33–600 and referenced to local topography for each crater. The topographic model is from the LOLA data (Smith et al., 2010). Our crater RBA estimates are consistent with those estimated by applying high-pass filters (Bierson et al., 2016), especially in the spatial pattern (comparing Figure 5a with Figure S2a).

We then forward modeled the Bouguer gravity for each crater, with a focus on the model dependence on the local ambient crustal porosity,  $\phi_c$ . The corresponding local crustal density ( $\rho_c$ ) is equal to  $\rho_g(1 - \phi_c)$ , where  $\rho_g$  is the grain density. Table 1 lists the model parameters and scaling relationships. With a focus on impact-induced porosity change (Region c in Figure 2e), we sequentially modeled three sources of the crater Bouguer anomaly: (a) Bouguer correction using a uniform correction density  $\rho_0$  (equal to the assumed global mean crustal density  $\overline{\rho_c}$  and the physical domain is indicated by the green area in Figure 2e), which introduces a non-zero gravity anomaly if  $\rho_c \neq \rho_0$  (Figure 2a); (b) post-impact breccia infill, which is typically more porous than the ambient crust (Kiefer et al., 2012, 2015), thus inducing a negative gravity anomaly (Figure 2b); and (c) impact-induced porosity change (the focus of this work). For each crater, we estimated the dimensions of the post-impact infill (blue area in Figure 2e) as the difference between the observed topography and the expected topography for a fresh crater. The fresh-crater topography was determined using the scaling relationships of Pike (1977) and Kalynn et al. (2013) (Table 1, Figure 1c) in a similar manner as Evans et al. (2016). While

impacts into more porous crust generate slightly smaller craters (Wünnemann et al., 2006; Collins et al., 2011), the effects on the crater gravitational signature will be minor. The physical dimensions of the post-impact infill can be roughly represented by a size-independent parameter, the degradation ratio (*R*) (Forsberg-Taylor et al., 2004), defined as the post-impact infill thickness divided by the expected fresh-crater depth. Figure 1b shows the spatial distribution of *R*. The gravity anomaly of the post-impact infill is controlled by its porosity ( $\Delta \phi_f$ ) relative to that of the local ambient crust, and  $\rho_g$ . We calculated the gravity anomalies due to (a) and (b) using Parker's forward modeling method (Parker, 1973) and show the model sensitivity to  $\rho_g$  in Figure S3.

For plausible values of  $\phi_c$ , R,  $\Delta \phi_f$ , and  $\rho_g$ , the major contributor to the crater Bouguer anomaly is (c) impact-induced porosity change. To estimate the gravitational signature from (c), we re-calculated the crater Bouguer anomalies (Figures 2c) for three impact hydrocode models (Milbury et al., 2015a) with  $D_c$  of 46 or 54 km, close to the mean crater diameter of 40 km for our analyzed craters. This re-calculation separates the gravitational signature of (c) impactinduced porosity change due to impact bulking and compaction from that of (a) Bouguer correction, which is necessary for taking into account (b) post-impact infill. We require the Bouguer correction density to vary and be equal to local  $\rho_c$  to quantify the gravitational effect solely from (c). We calculated the gravity anomaly by integrating the gravitational attraction from point masses with porosity changes (orange area in Figure 2e), from which we derived the modeled crater RBA. Instead, Milbury et al. (2015a) assumed a uniform Bouguer correction density of 2,650 kg/m<sup>3</sup>, and their calculated RBA is to the first order equal to the sum of (a) and (c). Therefore, Milbury et al. (2015a)'s RBA values are notably larger than our re-calculated RBA values solely from (c) (Figure 3a). Further, the dependence of the crater RBA on the integrated crustal column depth reveals that 95% of the crater RBA comes from a depth of 0.15– 0.2  $D_c$  (Figure 3b), yielding a depth scale for the local  $\phi_c$  estimate. By further assuming that the RBA is linearly proportional to  $D_c$  (Milbury et al., 2015a; Soderblom et al., 2015; Bierson et al., 2016), we derived a linear relationship for the gravity signature from (c):  $\frac{RBA}{D_c} = C\rho_g(\phi_c - \phi_r)$ , where the model-based parameters  $\phi_r$  and C are estimated to be ~7% and  $1.1 \times 10^{-6}$ mGal·m<sup>2</sup>/kg. These two parameters depend on the material strength, porous dilatancy and compaction models, temperature profile, and impact velocity used in the hydrocode model. We could not quantify the associated uncertainty, however, as only three hydrocode simulations were used. Figure 2d shows the total crater Bouguer anomalies for varied  $\phi_c$  values after adding (a), (b), and (c), for a typical crater with R = 0.5,  $\Delta \phi_f = 10\%$ ,  $\rho_g = 2,900$  kg/m<sup>3</sup>, and  $D_c = 54$  km. The bold black line in Figure 4a shows the corresponding linear RBA/ $D_c - \phi_c$  relationship. Increasing *R*,  $\Delta \phi_f$ , or  $\rho_g$  reduces the theoretical RBA/ $D_c$  (colored lines in Figure 4a).

In this manner, we establish the RBA/ $D_c - \phi_c$  relationship for each crater based on its  $D_c$  and topography (and thus *R*), assuming a constant  $\Delta \phi_f$  of 10%, and a  $\rho_g$  of 2,900 kg/m<sup>3</sup> for highland craters and 3,300 kg/m<sup>3</sup> for the mare craters (in the nearside maria). We then inverted the observed RBA/ $D_c$  to obtain the local ambient  $\phi_c$  for each crater (black dots in Figure 4b). The scatter in the inverted  $\phi_c$  is primarily due to the variations in *R*. The global mean RBA/ $D_c$  of - 0.07 mGal/km corresponds to a mean  $\phi_c$  of 12%, and the standard deviation of 0.29 mGal/km to 9%. If we assumed a lower  $\Delta \phi_f$  of 5%, the corresponding mean  $\phi_c$  becomes 10% with a standard deviation of 8% (Figure 4c). To be consistent with the previous global mean  $\phi_c$ estimate of 12% based on gravity analysis (Wieczorek et al., 2013; Besserer et al., 2014; Han et al., 2014) and Apollo sample and meteorite measurements (Kiefer et al., 2012), we adopted a  $\Delta \phi_f$  of 10% for the following spatial analysis. However, sensitivity tests showed that our porosity model is only sensitive to the assumed mean crustal porosity, not to the specific choice of  $\Delta \phi_f$  and  $\phi_r$  (Section 3.3).

We applied the moving-window block kriging technique (Cressie, 1993; Hengl, 2007) (Text S1) to map the spatially correlated pattern in the observed RBA/ $D_c$  and derived  $\phi_c$  (Figure 5) with standard errors (Figures 6a–b). We chose a cap window radius of 5° for the highlands and 10° for the western nearside maria (encircled by gray curve), based on the density of craters available for our analysis. We excluded mare regions whose spatial extent is too limited for analysis (gray patches). We centered the moving windows at 72×36 grid nodes with a grid distance of 5°. Each window incorporates an average of 10 craters for the highlands and 5 craters for the western nearside maria (Figure S1). Thus, our results for the western-maria region have larger uncertainties. We applied Student's t-test after decoupling transformation (Borgman, 1988) to test the null hypothesis that the local RBA/ $D_c$  or  $\phi_c$  are equal to the global means (Figures 6c, 7c). We tested the model sensitivity to the window size by doubling the window radii (Figures S4–5). The kriged maps were validated by deterministic interpolation methods

(Ding et al., 2016), including natural neighbor, inverse distance weighted, and Abel-Poisson spline interpolation.

#### **3 Results and Discussion**

#### **3.1 Spatial variations in the highlands**

Figure 5 shows the kriged maps of RBA/ $D_c$  and  $\phi_c$ , and Figure 6 shows the associated standard errors and p-values of the Student's t-tests. The global means of the kriged  $RBA/D_c$  and  $\phi_c$  are the same as those of original crater data, while the standard deviations decrease to 0.14 mGal/km and 4%, as the kriging processes have removed the spatially random signals. In the highlands, the southern part of SP-A is associated with the largest negative RBA/ $D_c$  of - $0.44\pm0.10$  mGal/km and the lowest  $\phi_c$  of  $4\pm3\%$  (all errors are the standard errors of the kriged values), suggesting that the local craters probe the impact melt sheet (Vaughan et al., 2013; Hurwitz & Kring, 2014) underlying a relatively thin porous top layer (Besserer et al., 2014), although we cannot exclude the effects of impact-induced porosity decrease and post-impact magmatism (Whitten & Head, 2015; Shearer et al., 2015). The north-south asymmetry of the SP-A basin might be due to enhanced impact melting in the southern half of the basin caused by an oblique impact from north to south with its first contact near the Ingenii basin (Garrick-Bethell & Zuber, 2009; Schultz & Crawford, 2011; Petro, 2012). Alternatively, the positive RBA/ $D_c$ anomalies near the Ingenii basin could be caused by post-impact mare deposits (Figure 2b), and thus the inferred high (pre-impact) porosity might be erroneous. In contrast to the SP-A basin, a region immediately to the southwest of the western nearside maria is associated with the highest RBA/ $D_c$  of 0.22±0.11 mGal/km and  $\phi_c$  of 21±4 %, probably due to kilometer-scale ejecta deposits emplaced during basin-forming events (Hiesinger, 2006; Petro & Pieters, 2008). The original and kriged RBA/ $D_c$  and  $\phi_c$  data are provided in Table S1 and S2.

Other notable anomalies are associated with impact basins 300–900 km in radius (circles in Figure 5). The Moscoviense basin reveals a more negative RBA/ $D_c$  and lower  $\phi_c$  interior to the crater rim than the surrounding region (from the crater rim to one crater radius beyond the crater rim), which we confirmed by applying a paired-sample *t*-test at a significance level of 0.05. Paired-sample *t*-tests also suggest porosity reduction for the Apollo, Korolev, Hertzsprung, and Coulomb-Sarton basins. This porosity reduction is consistent with impact-induced compaction and heating (Melosh, 1989; Milbury et al., 2015a) and impact-generated melting (Cintala & Grieve, 1998), while the porosity increase immediately exterior to the basin corresponds to impact ejecta. In contrast, positive RBA/ $D_c$  signals in the Humbolditianum, Australe, Orientale, and Ingenii basins are likely due to mare deposits (gray patches, Head, 1976) emplaced after the formation of impact craters, and not a large (pre-impact)  $\phi_c$ .

Figure 7 shows the original and kriged RBA/ $D_c$  and  $\phi_c$  data, as well as kriging errors and pvalues, along a great-circle cross-section A–F (Figure 5). The results using smaller windows (red curves) are consistent with those using larger windows (blue curves) to the first order. Smaller windows resolve shorter-wavelength variations, e.g., in the Moscoviense basin, while a larger window yields more robust regional estimates with lower kriging errors and p-values. Nevertheless, the p-values at both window sizes (Figures 6c, 7c) suggest that at least the southern SP-A basin is associated with anomalous RBA/ $D_c$  at a significance level of 0.05 (i.e., p < 0.05). This distinction of the SP-A basin from the rest of the highlands was previously found by gravity admittance analysis (Besserer et al., 2014). We also confirmed a statistical deviation from the global mean in the southwestern periphery of the nearside maria (Figure S5c). The statistical test for  $\phi_c$  showed consistent results.

#### 3.2 Western nearside maria

We estimated an average RBA/ $D_c$  of ~ -0.39 mGal/km and  $\phi_c$  of ~ 6% (Figure 5) in the western mare region (including the Mare Imbrium, Oceanus Procellarum, and Mare Nubium), while the craters in the eastern maria are too sparse (Figure 1a) for a reliable estimate. The associated kriging errors (Figures 6a–b, 7a–b) and p-values (Figures 6c, 7c) for the western nearside maria are larger than the highlands due to a sparser crater distribution. We conclude, however, that the western mare region as a whole (with a total of 74 craters) show statistically lower RBA/ $D_c$  and  $\phi_c$  values than the global mean values, with a p-value on the order of 10<sup>-4</sup>.

The largest negative RBA/ $D_c$  of  $-0.70\pm0.34$  mGal/km and the lowest  $\phi_c$  of  $0\pm9$  % after kriging exists in Mare Imbrium, probably due to the existence of a relatively thick low-porosity top mare layer and/or low-porosity underlying feldspathic crust. A preliminary hydrocode simulation (Milbury et al., 2015b) suggests that a non-porous mare layer of 5 km could reduce the RBA/ $D_c$  by ~ 0.3 mGal/km. Therefore, a mare layer of several kilometers should be sufficient to explain the mean RBA/ $D_c$  in the Mare Imbrium. However, if the regional mare thickness is only ~ 2 km (De Hon, 1979; Thomson et al., 2009), low-porosity underlying crustal materials would be expected as suggested by Gong et al. (2016). In contrast, the Mare Nubium is associated with slightly positive RBA/ $D_c$  and thus relatively high  $\phi_c$ , indicating a thinner mare layer with a higher porosity and/or more porous underlying crust. Post-impact mare deposits may also introduce positive RBA (Evans et al., 2016). We did not find RBA/ $D_c$  extremes in Marius Hills, Oceanus Procellarum, where the thickest mare has been suggested (Gong et al., 2016). This is probably due to the limited spatial coverage of mare craters (Figure 1a). The correlation between thick mare basalts (and thus low porosities) and young surface ages (Hiesinger et al., 2011) in the western Mare Imbrium and eastern Oceanus Procellarum implies a scenario that the volcanic eruption starts synchronously (relative to the rate of impacts) and then maintains a constant eruption rate until the observed regional surface age. This correlation also implies that the impact flux during and after mare emplacement was inadequate to notably increase the porosity of the mare layer.

#### **3.3 Model sensitivity and implications**

Our porosity ( $\phi_c$ ) model depends on the crater diameter range considered (20–100 km for our study), because the depth to which a crater probe is proportional to its size. Our  $\phi_c$  map also relies on the assumed  $\phi_r$  and  $\Delta \phi_f$  values. We adopted a  $\phi_r$  of 7% from impact hydrocode simulations (Milbury et al., 2015a) and a  $\Delta \phi_f$  of 10% to match the observed global mean crustal porosity  $\overline{\phi_c}$  of 12% (Kiefer et al., 2012; Wieczorek et al., 2013).  $\overline{\phi_c} > \phi_r$  implies that most of the crust has a high porosity to induce a positive RBA through impact compaction. This is consistent with the existence of a ~10 km thick megaregolith mainly due to basin-scale ejecta

deposits (Hörz et al., 1991). However, we find that reducing either  $\Delta \phi_f$  or  $\phi_r$  lowers the inferred  $\overline{\phi_c}$  (Figure 8). If both  $\Delta \phi_f$  and  $\phi_r$  are allowed to vary and  $\overline{\phi_c}$  is required to be 12%, the porosity model becomes uniquely determined as it now solely relies on the assumed  $\overline{\phi_c}$ . This does, however, create a trade-off between  $\Delta \phi_f$  and  $\phi_r$  (Figure 8b). It is therefore possible that most of the highland crust has achieved a steady-state in porosity (i.e.,  $\phi_r = \overline{\phi_c} = 12\%$ ) if there is a systematic bias in the hydrocode simulations of Milbury et al. (2015a). In this case, the expected porosity of the post-impact breccia infill relative to the ambient crust ( $\Delta \phi_f$ ) is 2%, corresponding to an absolute breccia infill porosity of 14%. This porosity value is entirely possible as it lies within the porosity range of 11–21% from the Apollo impact breccia measurements (Kiefer et al., 2015). We therefore cannot conclude if the highland crust has achieved a steady state in porosity solely from the observed crater gravity anomalies, if a steady state is expected after a prolonged bombardment history.

Our crustal porosity model is more robust in a relative sense given the uncertainty in the input parameters of the iSALE hydrocode models (Milbury et al., 2015a). Our assumption of a depth-independent porosity, instead of a porosity that decreases with depth (Besserer et al., 2013; Wieczorek et al., 2013) may introduce additional errors. Preliminary hydrocode simulations that include a porosity gradient suggest that the crater RBA is most sensitive to the porosity structure of the top few kilometers (Milbury et al., 2015b), implying that the weighted crater probe depth might be less than  $0.15-0.2 D_c$ .

In addition to crustal porosity, we have implicitly considered the effects of crater degradation state and post-impact mare deposits on the crater RBA. Other properties are probably less significant: (1) Hydrocode impact models (Milbury et al., 2015a) showed that the crustal thickness is not particularly influential. (2) Regression analysis by Bierson et al. (2016) found no correlation between the crater RBA and the crustal thickness, as well as the grain density (Huang & Wieczorek, 2012; Wieczorek et al., 2013) and crater age. (3) Impact melt sheets have limited influence for  $D_c$  less than 100 km (Cintala & Grieve, 1998). (4) Hotter mantle temperatures in the Procellarum KREEP Terrane (Wieczorek & Phillips, 2000; Laneuville et al., 2013; Miljković et al., 2013) may reduce the regional  $\phi_r$  because the hotter temperature reduces the material strength and suppress the creation of pore space (Collins et al., 2004). This implies less porous nearside maria than in this study. However, the spatial pattern should stay the same. (5) Impact characteristics, including impactor composition (material strength and density), size, velocity and angle may contribute significantly to the standard deviations of the RBA/ $D_c$  and  $\phi_c$ before kriging. The impact characteristics (e.g., Silber et al., 2017) also influence the fresh crater topography and thus our estimated post-impact infill. However, the theoretical spatial randomness (i.e., lack of correlation) of these impactor parameters means that their effects should be removed by the kriging technique.

A comparison of our porosity model with that of Wieczorek et al. (2013) (Figure S2b) yields a correlation coefficient of 0.59, statistically distinct from zero based on Student's t-test and Fisher transformation. Wieczorek et al. (2013) solved for Bouguer correction densities and corresponding porosities that minimize the correlation between the regional Bouguer anomaly and topography, equivalent to considering only the gravitational source (a) Bouguer correction in our crater gravity method. Thus, our method and the Bouguer correlation method are not entirely independent. The Bouguer correlation method may overestimate the porosity variations for densely cratered regions, as it ignores the gravitational effects of (c) impact bulking and compaction. This overestimation exists in the regions immediately exterior to the Moscoviense and other basins (comparing Figure 5b and S2b). However, the model by Wieczorek et al. (2013) is more accurate for less densely cratered regions such as the Orientale basin (Figures 1, S1a), because our crater gravity-based model depend strongly on the local crater densities (reflected in the standard errors in Figure 6).

The porosity gradient model of Besserer et al. (2014) may overestimate the porosity variations of the pre-impact crust, as they do not consider the gravitational source (c) neither. Although our crater gravity method cannot resolve the spatial variations in the porosity gradient due to the large data scattering, a comparison between Besserer et al. (2014) and our model helps to confirm the depth scale of our porosity model. We calculated the averaged crustal porosity of Besserer's linear porosity model to various depths and found a best-fit depth of  $\sim$  7 km that minimizes the difference between Besserer et al. (2014) and our porosity models (Figure S6). This depth scale is consistent with a mean crater probe depth of 6–8 km, estimated as 0.15–0.2 of the mean crater diameter of 40 km.

### **4** Conclusions

Based on high-resolution GRAIL gravity observations and the extensive LOLA crater database, we developed a method to use the RBA/ $D_c$  ratios of the mid-sized craters to map regional-scale spatial variations in the Moon's crustal porosity ( $\phi_c$ ), giving insights into the effects of older and larger impact events and nearside volcanism:

1. The analyzed 4,864 mid-sized craters are associated with a mean RBA/ $D_c$  of -0.07 mGal/km. The crustal porosity model for the top several kilometers can be uniquely determined assuming a global mean  $\phi_c$  of 12%. Meanwhile, a trade-off exists between the porosity that results in zero net impact-induced porosity change ( $\phi_r$ ) and the post-impact infill porosity: Larger  $\phi_r$  implies less porous post-impact breccia infill.

2. In the highlands, the southern SP-A basin shows the largest negative RBA/ $D_c$  of -0.44±0.10 mGal/km and lowest  $\phi_c$  of 4±3 %, consistent with the existence of an impact melt sheet due to an oblique impact. In contrast, a region immediately to the southwest of the nearside maria is associated with the highest RBA/ $D_c$  of 0.22±0.11 mGal/km and  $\phi_c$  of 21±4 %, probably due to ejecta deposits emplaced during basin-forming events.

3. Most of the impact basins we investigated, including the Moscoviense, Apollo, Korolev, Hertzsprung, and Coulomb-Sarton basins, show lower porosities interior to the crater rims compared to their surroundings. This is consistent with impact-generated melt sheets and outward emplacement of ejecta deposits. Recovery of the impact basin structures validates our premise that mid-sized craters can be used as probes for regional porosity structure.

4. The western nearside maria display negative RBA/ $D_c$ , consistent with the existence of a top mare layer. The largest negative RBA/ $D_c$  of -0.70±0.34 mGal/km exists in the Mare Imbrium, probably due to the combined effects of a thick low-porosity mare layer and low-porosity underlying crust.

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Table 1. Parameter Ranges and Scaling Relationships for Crater Bouguer Gravity Models

Parameter	Description	<b>Range/Equation</b>
$D_c$	Crater diameter (km)	20–100
$D_{f}$	Crater floor diameter (km)	$0.19D_c^{1.25*}$
Н	Fresh crater floor depth in the highlands (km)	$1.558 D_c^{0.254\dagger}$
	Fresh crater floor depth in the maria (km)	$0.87 D_c^{0.352\dagger}$
h	Observed crater floor depth (km)	0–5
R	Crater degradation ratio	(H-h)/H
$W_R$	Width of rim flank (km)	$0.467 D_c^{0.836*}$
$D_R$	Diameter of outer rim flank (km)	$D_c + 2W_R$
$h_R$	Rim height (km)	$0.236D_c^{0.399*}$
$\rho_{g}$	Highland grain density (kg/m <sup>3</sup> )	2,900 <sup>‡</sup>
U	Mare grain density $(kg/m^3)$	3,300 <sup>‡</sup>
$\overline{\rho_c}$	Global mean crustal density (kg/m <sup>3</sup> )	$2,550^{\$}$
$\overline{\boldsymbol{\phi}_{c}}$	Global mean crustal porosity (%)	12 <sup>§</sup>
$\rho_0$	Bouguer correction density (kg/m <sup>3</sup> )	2,550
$\Delta \phi_f$	Porosity of post-impact breccia infill wrt. ambient crust (%)	$0 - 10^{**}$
$\phi_r$	Crustal porosity that results in zero net impact-induced porosity change and crater RBA (%)	$7^{\dagger\dagger}$
*Pike (1977)		

Kalynn et al. (2013)

<sup>‡</sup>Huang & Wieczorek (2012)

<sup>§</sup>Kiefer et al. (2012); Wieczorek et al. (2013); Besserer et al. (2014); Han et al. (2014)

<sup>\*\*</sup>Kiefer et al. (2015)

<sup>††</sup>Milbury et al. (2015a)

**Figure 1.** Crater locations and degradation states based on lunar topography. (a) Topography map and locations of 4,864 mid-sized craters from the LOLA crater database used in this study. (b) Map of the kriged crater degradation ratio (*R*). The grid spacing is 10° for the western nearside maria (encircled by the gray curve) and 5° for the rest of the area (same as in Figures 5 and 6). The global mean *R* value of 0.47 corresponds to a partially degraded state, while lower *R* implies younger local craters with a limited amount of post-impact infill. (c) Observed crater floor depth for highland craters (blue dots) and mare craters (red triangles). The blue and red curves are the expected crater floor depth for fresh craters (i.e., *R* = 0) in the highlands and maria, respectively (Kalynn et al., 2013).

**Figure 2.** Modeled crater Bouguer gravity due to (a) Bouguer correction assuming a uniform Bouguer correction density, (b) post-impact breccia infill with variable  $\Delta \phi_f$ , or mare infill with a density of 3,150 kg/m<sup>3</sup>, and (c) impact-induced porosity change. (d) Total crater Bouguer anomaly by adding (a), (b), and (c). All the modeled curves are vertically shifted to ensure that the average gravity within the background annulus (i.e., from outer rim flank to one crater radius) is zero; vertical shifts do not influence the crater RBA values. (e) Physical domains corresponding to (a), (b) and (c), indicated by different colors. Bold black curve is the observed topography. Fixed parameters include R = 0.5,  $D_c = 54$  km,  $\Delta \phi_f = 10\%$ , and  $\rho_g = 2,900$  kg/m<sup>3</sup>. The scaling relationships for *H*,  $h_R$ ,  $D_f$ , and  $D_R$  are listed in Table 1, yielding 4.3 km, 1.2 km, 27.8 km, and 81.8 km, respectively. *h* is the observed crater floor depth, assuming to be 0.5*H*. **Figure 3.** Crater gravity and probe depth based on iSALE hydrocode simulations (Milbury et al., 2015a). (a) Crater RBA/ $D_c$  ratios solely due to impact-induced porosity changes (corresponding to Figure 2c). Black dots and black line are re-calculated ratios using (varied) Bouguer correction densities equal to crustal densities. Blue dots and blue line are original values from Milbury et al. (2015a) where a constant Bouguer correction density was assumed. Gray dots and dashed line are the results for larger crater diameter of 90 or 96 km. The slope for the larger-diameter models is steeper, but this secondary effect is negligible in this study. (b) Cumulative crater RBA, normalized by the total RBA, over increasing depth column (*z*) for four impact hydrocode models in Milbury et al. (2015a). Dots indicate the effective depths when 95% of the total RBA is reached.

**Figure 4.** (a) Modeled RBA/ $D_c$ - $\phi_c$  relationship for variable R,  $\Delta \phi_f$ , and  $\rho_g$ . (b) Observed RBA/ $D_c$  ratios and converted  $\phi_c$  assuming  $\Delta \phi_f = 10\%$ . (c) Histogram of converted  $\phi_c$  assuming  $\Delta \phi_f = 10\%$  (gray) or 5% (orange).

**Figure 5.** Maps of the kriged (a) RBA/ $D_c$  and (b)  $\phi_c$  in a Lambert azimuthal equal-area projection centered over the nearside (left) and farside (right). The grid spacing is 10° for the western nearside maria (encircled by the gray curve) and 5° for the rest of the area. Gray patches indicate mare basalt deposits (Head, 1976) other than the western nearside maria. Circles indicate major highland basins. Initials SP-A, CS, FS and MR represent South Pole-Aitken, Coulomb-Sarton, Freundlich-Scharonov and Mendel-Rydberg basins, respectively.

**Figure 6.** Standard errors of the kriged (a) RBA/ $D_c$  and (b)  $\phi_c$  values in Figure 5. (c) P-values under the null hypothesis that the regional RBA/ $D_c$  are equal to the global mean value of -0.07 mGal/km.

**Figure 7.** (a) Observed RBA/ $D_c$  and (b) converted  $\phi_c$  values before and after kriging along the cross-section A–F (Figure 5). Red solid and dashed lines are the kriged values (corresponding to Figure 5) and errors (Figure 6), respectively, using a window radius of 5° for the highlands and 10° for the western nearside maria. Gray dots are the original data within the local windows. Blue lines correspond to results with doubled window radii (Figures S4–5). (c) P-values under the null hypothesis that the regional RBA/ $D_c$  is equal to the global mean. (d) Crustal structure and inferred crustal porosities. Vertical lines indicate the maximum crater probe depth of  $0.2D_c$  for the craters in (a) and (b). The Moho relief is from Wieczorek et al. (2013).

**Figure 8.** Trade-off between  $\Delta \phi_f$  and  $\phi_r$ . (a) Expected linear relationship between the total RBA/ $D_c$  and crustal porosity for a typical crater with varied  $\Delta \phi_f$  and  $\phi_r$ , assuming R = 0.5,  $D_c = 54$  km, and  $\rho_g = 2,900$  kg/m<sup>3</sup>. (b) Mean crustal porosity for varied  $\Delta \phi_f$  and  $\phi_r$ . The bold black line indicates the solutions with an assumed global mean crustal porosity of 12%. Our porosity model assumed a  $\phi_r$  of 7% (black dot).





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