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Citation: Su, GY, Moreira, TA, Lee, D, Jena, A, Wang, G et al. 2022. "Wettability and CHF limits of Accident-Tolerant nuclear fuel cladding materials in light water reactor conditions." Applied Thermal Engineering, 216.

As Published: 10.1016/J.APPLTHERMALENG.2022.119018

Publisher: Elsevier BV

Persistent URL: <https://hdl.handle.net/1721.1/147088>

Version: Original manuscript: author's manuscript prior to formal peer review

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1 Wettability and CHF limits of Accident-Tolerant Nuclear Fuel Cladding Materials in 2 Light Water Reactor Conditions

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8 **ABSTRACT**

9 We present the results of experimental investigations aimed at evaluating the thermal-hydraulic performance of chromium-coated
10 zircaloy, i.e., one of the most promising accident tolerant fuel (ATF) cladding material for light water nuclear reactors. Precisely,
11 we investigate the wettability and critical heat flux (CHF) limits of chromium-coated and conventional zircaloy surfaces in
12 prototypical reactor conditions. For both surface types, we measure the contact angle in a vapor-saturated atmosphere from ambient
13 pressure to the operating pressure of pressurized water reactors (PWRs), i.e., ~15 MPa. We measure the ambient-pressure steady-
14 state flow boiling CHF with a spatially uniform heat flux. We measure the high-pressure steady-state flow boiling CHF with a
15 cosine shape heat flux (up to 20 MPa) and with a uniform heat flux (up to 15 MPa), also exploring the effect of CRUD deposits on
16 the chromium-coated surface. Our results reveal that the chromium surface and the bare zircaloy surface have similar wettability
17 and both become super-hydrophilic in PWR conditions, and that there is practically no difference in the steady-state CHF limits,
18 both at low-pressure and high-pressure conditions, also when the chromium-coated surface is covered by a CRUD deposit.
19 However, while the chromium-coating does not improve the CHF compared to the bare zircaloy surface, it improves the post-CHF
20 behavior. The chromium coating prevents the reaction between zircaloy and steam, which results in the formation of a brittle
21 zirconium oxide through the surface of the cladding. We also measure the transient CHF under exponentially escalating heat flux
22 inputs of a nano-smooth and a rough surface mimicking a chromium-coated zircaloy cladding. Interestingly, the results of the
23 transient heat flux tests suggest that the CHF limit for very short periods (i.e., fast transients) is independent of the surface finish,
24 being the same for a rough chromium surface or a nano-smooth surface.

25 **Keywords**

26 CHF, wettability, contact angle, high-pressure, nuclear fuel cladding, zircaloy, chromium, ATF, LWR, RIA

27 **1. Introduction**

28 Since the 2011 Fukushima accident, significant research has been devoted to developing accident-tolerant fuel (ATF) cladding
29 materials [1,2]. These investigations have mainly focused on the ability of potential ATF materials to resist runaway steam
30 oxidation under extreme temperature (> 1200 °C) and retain their mechanical strength and structural integrity under thermal shocks
31 [3-5]. However, knowledge of the thermal-hydraulic performance of ATF materials is relatively lacking, particularly under light
32 water reactor (LWR) operating conditions. Precisely, there is an urgent need to clarify whether ATF materials could lead to a
33 deterioration of the CHF limit compared to traditional zirconium alloy (i.e., zircaloy) claddings, as the CHF limit determines,
34 among other factors, the thermal power of LWRs [6].

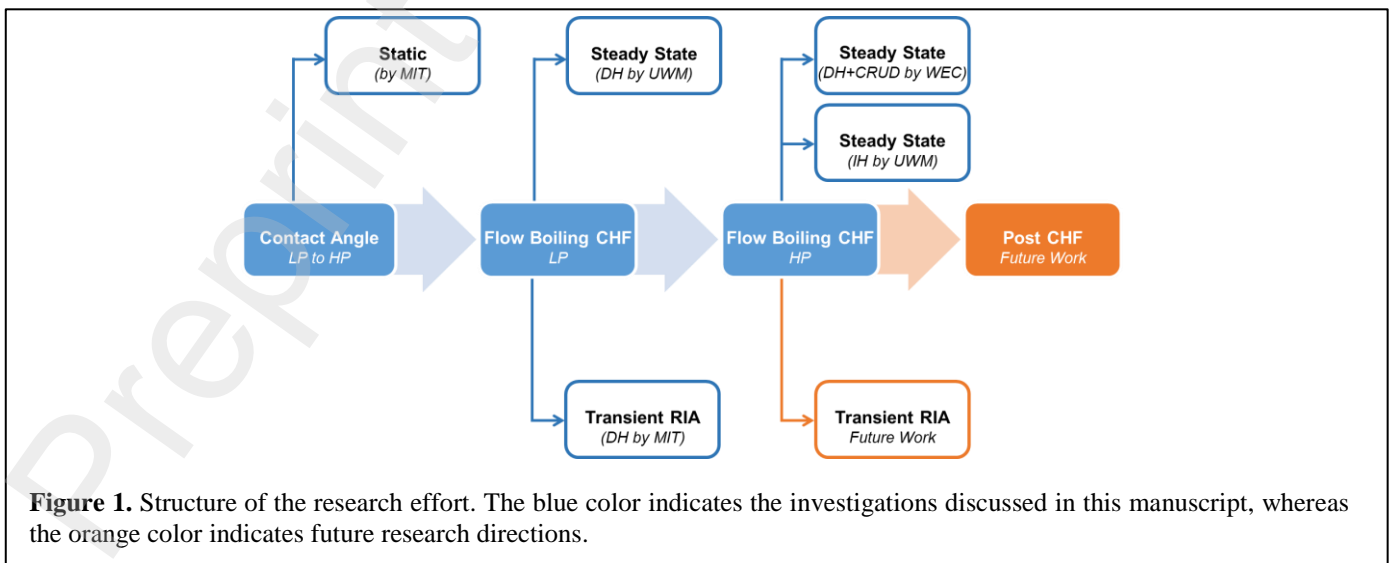
35 In pressurized water reactors (PWRs), for instance, water enters the core highly subcooled, and while its bulk temperature does
36 not reach saturation conditions, subcooled nucleate boiling may occur on the fuel claddings in the upper part of the core. This
37 subcooled nucleate boiling regime is very effective. However, the risk of a transition from a nucleate boiling to a film boiling

38 regime (i.e., the departure from nucleate boiling, DNB) is a major concern in the operation of PWRs, as it may also result in the
 39 melting of the fuel cladding. Such transition occurs if the local heat flux exceeds the CHF limit. However, while there are several
 40 pool and flow boiling CHF studies on zircalloys, the number of studies addressing the CHF performance of ATF materials is still
 41 low and limited to low-pressure conditions [7-12]. To support the successful rollout of ATF cladding materials, it is important to
 42 thoroughly investigate the CHF limit of these materials and understand how they compare to zircalloys, especially in LWRs
 43 conditions.

44 It is widely recognized that the CHF limit depends on operating conditions (e.g., pressure, temperature, and flow rate, if any)
 45 and surface properties. Surface wettability (i.e., surface contact angle with liquid) is known to be a dominant factor, at least in pool
 46 boiling at low pressure [13-17]. Higher surface wettability can increase the CHF by affecting the boiling dynamics (e.g., by
 47 reducing the bubble departure radius and increasing the bubble departure frequency). Since the wettability is directly correlated to
 48 the surface tension and thus a strong function of temperature, it is essential yet very challenging to measure the surface wettability
 49 of fuel cladding materials in LWR conditions. There is hitherto no systematic measurement of contact angles on traditional zircaloy
 50 up to such high pressure and temperature in vapor-saturated atmospheres, let alone ATF materials. Most measurements were
 51 obtained at near atmospheric pressure, and low temperatures [7-12, 18-24]. Only a few studies were performed in pressurized
 52 conditions, but in an inert gas environment, which leads to strong evaporation effect especially under elevated temperature
 53 conditions [25-27]. Practically, contact angles values input in thermal-hydraulics models used for LWR analyses often rely on data
 54 and correlations developed in non-prototypical conditions, or hypothesis on the presumed surface wettability trends [28].

55 Other than normal steady state conditions, CHF under transient conditions is another major concern in nuclear reactors. Such
 56 off-normal scenario might be caused if a rapid insertion of reactivity (e.g., an accidental removal of control rods from the core)
 57 leads the nuclear reactor to a prompt-critical state (i.e., a reactivity-initiated accident, or RIA). Under such conditions, the reactor
 58 power, \dot{Q} , may increase exponentially, i.e., $\dot{Q} \propto e^{t/\tau}$. Depending on the reactivity value, the exponential period τ can vary from a
 59 few millisecond (for large insertion of reactivity) to hundreds of milliseconds (for small insertion of reactivity). The delay between
 60 the power escalation and the void feedback (i.e., the boiling of water) may lead to a large increase in the reactor power, potentially
 61 triggering a DNB event [29-43]. Thus, it is also key to understand the DNB mechanisms and quantify the CHF limit of ATF
 62 materials in transient conditions.

63 Here, we present the results of a research effort involving Massachusetts Institute of Technology (MIT), University of
 64 Wisconsin-Madison (UWM), and Westinghouse Electric Company LLC (WEC). This research effort has focused on investigating
 65 the thermal-hydraulic performance of chromium coated zircaloy as a candidate ATF cladding concept. Figure 1 shows the general
 66 structure of the present research project (in blue) and expected future works (in orange).

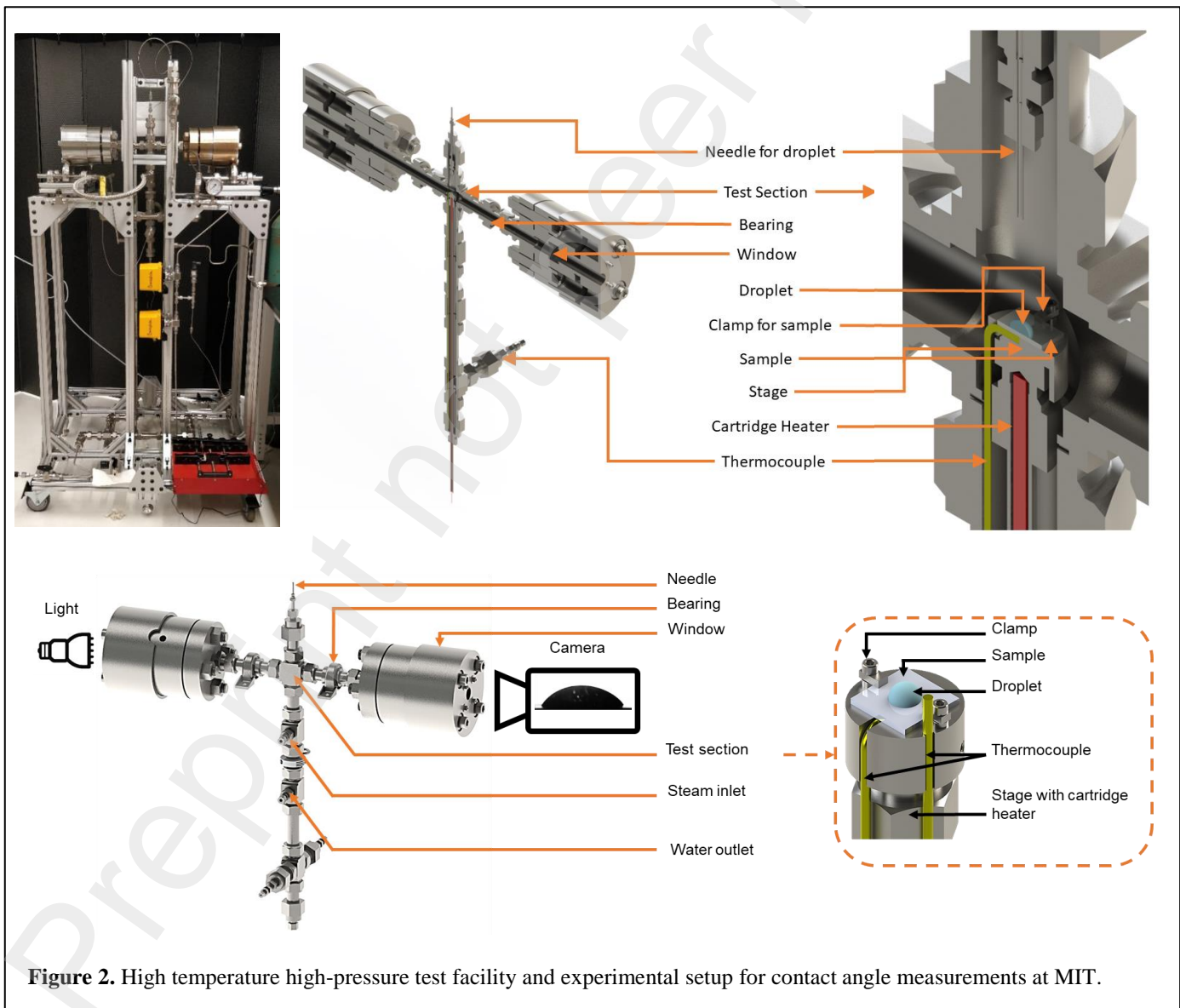


68 Precisely, we discuss methodology and results of experimental investigations aimed at characterizing the wettability and CHF
 69 limits in low pressure conditions (including transient power escalations) and high-pressure conditions (up to 20 MPa). The
 70 wettability of the selected materials is characterized using a special autoclave designed and built by MIT to measure the contact
 71 angles of water in a vapor-saturated environment up to the critical pressure of water. Low-pressure (LP), steady-state CHF flow
 72 boiling tests on actual fuel claddings samples are conducted by UWM. Transient CHF tests with power escalation periods ranging
 73 from 5 to 500 ms are conducted by MIT using a cladding-simulant heater. High-pressure flow boiling tests are conducted by both
 74 UWM and WEC. WEC investigates CHF with uniform heat flux profiles, also considering the effect of CRUD deposits on the
 75 surface of the cladding. UWM investigate CHF with a cosine-shape heat flux mimicking the axial heat flux distribution of PWRs.
 76

77 **2. Surface wettability**

78 *2.1. High temperature high pressure autoclave facility*

79 Figure 2 shows the MIT high-pressure autoclave facility and the optical setup. This facility allows measuring the contact angle
 80 between a liquid droplet and a surface in a vapor-saturated environment (i.e., without noncondensable gases), from ambient to
 81 critical pressure conditions. The autoclave features a cruciform geometry. The test sample is accommodated on a stage installed at
 82 the crossing joint of the autoclave, aligned with the axis of two sapphire windows that enable optical access to it.



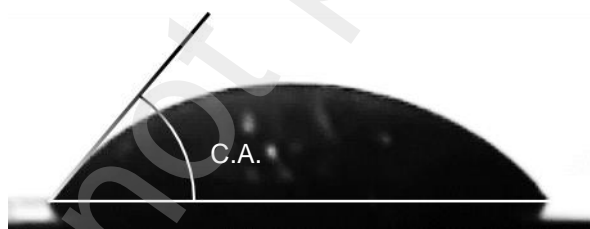
83 **Figure 2.** High temperature high-pressure test facility and experimental setup for contact angle measurements at MIT.

84 Two K-type thermocouples (shown in yellow in Fig. 2) are used to measure the temperatures underneath the test sample and in
85 the vapor atmosphere near it and to verify that they are in equilibrium. The test section can accommodate rectangular samples of
86 different size, from 5 mm × 5 mm to 12.5 mm × 12.5 mm with thickness of 1 to 3 mm. A needle connected to an external syringe
87 pump is used to dispense the droplet on top of the specimen. The needle is in thermal equilibrium with the autoclave atmosphere
88 to ensure that the fluid released from the needle is at the same temperature as the surrounding steam and the sample itself. After a
89 contact angle measurement is made, the surface is flushed with water and a cartridge heater underneath the sample holder is used
90 to evaporate the residual water. Then, the sample is let thermalize before a new droplet is dispensed to make a new measurement.
91 The autoclave pressure and the droplet line pressure are monitored with two pressure transducers separately.

92 2.2. Experimental procedure

93 Before its installation in the autoclave, the test sample is solvent-cleaned in a sonicated bath using soapy water, DI water,
94 acetone, ethanol, and DI water again, for seven minutes each. At the beginning of each test, the test facility is purged with argon
95 gas for at least five minutes to remove all the air from the autoclave. Then, the autoclave is vacuumed to remove the non-
96 condensable gas until the pressure drops below 10 kPa and filled with de-aerated DI water to initialize a two-phase equilibrium
97 atmosphere. The autoclave is externally heated using tape heaters, and the temperature and pressure inside the autoclave are
98 carefully controlled and monitored to ensure that thermal equilibrium conditions are attained at the desired pressure setpoint. Once
99 the autoclave has reached equilibrium, a droplet (also in thermal equilibrium) is released from the needle on top of the sample
100 surface.

101 We measure the static contact angle of deionized (DI) water with selected surfaces using the sessile-drop method, as shown in
102 Fig. 3. We use a high-resolution camera with a resolution of 106 $\mu\text{m}/\text{pixel}$ to take backlit images of sessile drops on the sample
103 surface. Then, we use an image processing algorithm to identify the droplet profile and the tangent to this profile at the solid-liquid-
104 vapor contact line. The contact angle value is then calculated from the tangent as shown in Fig. 3. For more details about the image
105 processing technique, the reader is directed to the supplementary material of Ref. [44].



106 **Fig. 3.** Illustration of sessile-drop method for static contact angle measurement on smooth Zr-4 surface at $T = 213.9\text{ }^\circ\text{C}$,
107 $P = 20.6\text{ bar}$.

108 Different from tests in inert gas environments, our droplets do not evaporate. They maintain their shape and dimensions for an
109 indefinite period of time. This is possible thanks to the operation in a practically vapor-saturated environment. After each
110 measurement, additional water from the needle is poured to clean the sample surface. Then, the cartridge heater is turned on to
111 evaporate the residual water. After the surface is completely dry, the cartridge heaters are shut down and we wait for the sample to
112 cool down until its temperature stabilizes within $\pm 1^\circ\text{C}$ of the autoclave environment temperature. Then, we repeat the same
113 measurement. At least five measurements are made for each saturation pressure and temperature condition. Then, we change the
114 operating conditions, and collect a new set of measurements. Importantly, we take these measurements for increasing and
115 decreasing pressure and temperature, to explore the possibility of contact angle hysteresis.

115 2.3. Contact angle measurements

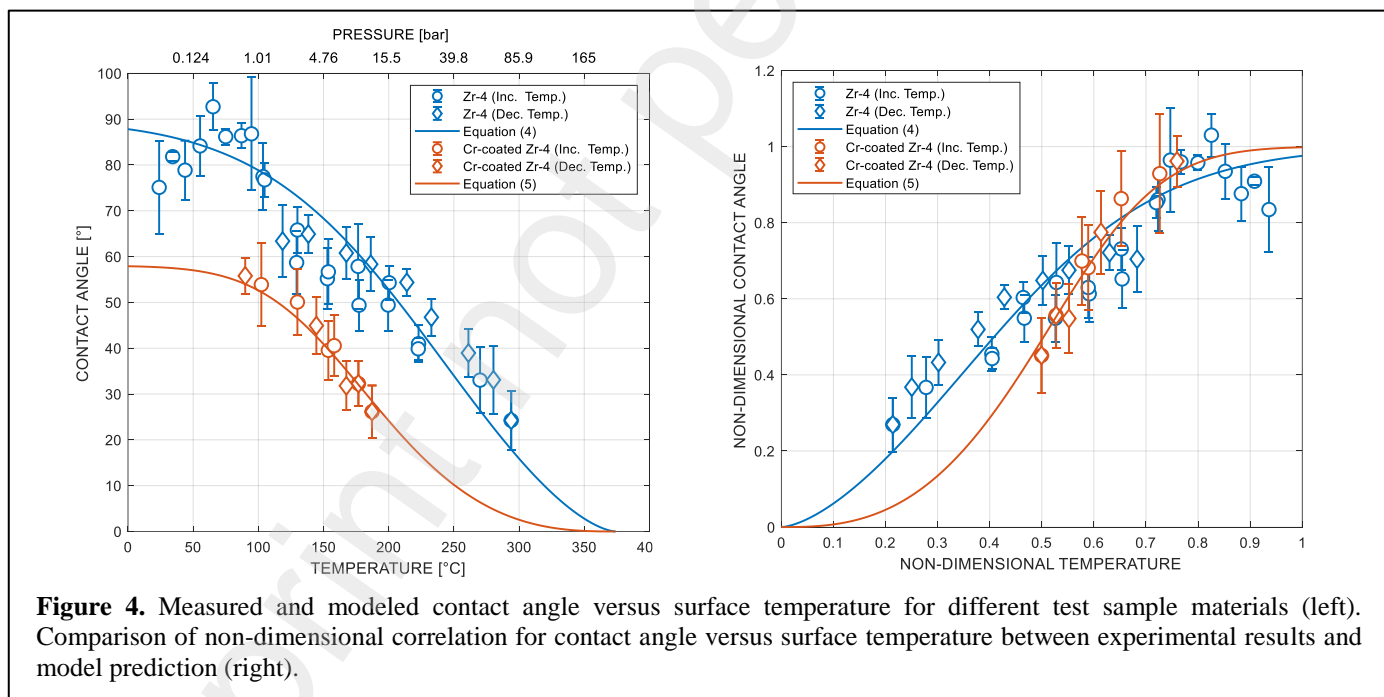
116 The static contact angle measurements are conducted on two different mirror-polished surfaces, i.e., zircaloy-4 (Zr-4) and
117 chromium. The roughness R_a of the test samples is summarized in

118 Table 1. The chromium surface is created by coating a mirror polished Zr-4 sample using physical vapor deposition (PVD). The
 119 PVD-coated chromium layer is about 1 μm , i.e., it is thick enough to eliminate the impact of Zr-4 substrate on the contact angle
 120 [48]. Note that we measure the wettability of mirror polished surfaces with a roughness much lower than commercial surface to
 121 eliminate as much as possible surface effects, e.g., due to random roughness, from the analysis. In other words, what we measure
 122 here is the intrinsic wettability of the surface.

123 **Table 1.** Summary of test samples and the corresponding surface roughness.

Material	Surface finish, Ra [μm]
Zr-4	Mirror-polished, 0.033 ± 0.011
Cr-coated Zr-4	Mirror-polished, 0.050 ± 0.011

124 Figure 4 (left) shows how the contact angle on the two surfaces change with temperature and pressure. The error bar represents
 125 the standard deviation of the contact angle values for the set of five measurements taken at each test condition. Several sub-
 126 atmospheric data points for Zr-4 previously obtained by our group (see Kossolapov et al. [45]) are also shown in the plot. The
 127 results shown in Figure 4 demonstrate that the contact angle decreases as temperature and pressure increase, which is consistent
 128 with the decrease of the liquid-vapor surface tension [46]. Note that, when the contact angle decreases below 20° , measurements
 129 become challenging and inaccurate due to light distortion effects created by the warm vapor in the optical path. Thus, we report
 130 the data only until the point when the contact angle is approximately 20° . However, we have observed that the surfaces become
 131 super-hydrophilic above these temperatures (i.e., the droplet spreads over the entire sample surface as soon as it is dispensed, or,
 132 in other words, the contact angle is practically 0°).



133 **Figure 4.** Measured and modeled contact angle versus surface temperature for different test sample materials (left).
 134 Comparison of non-dimensional correlation for contact angle versus surface temperature between experimental results and
 135 model prediction (right).
 136 The decreasing trend of contact angle versus temperature has implications on the boiling process. It is often argued that the CHF
 137 limit depends on the wettability of the boiling surface and that, improving the surface wettability (i.e., decreasing the contact angle),
 138 increases the CHF. This statement is mostly supported by many pool boiling data, as wettability effects seem to be less pronounced
 139 in flow boiling conditions. However, Zr-4 and chromium have similar contact angles at low temperature conditions, and they both
 tend to become super-hydrophilic in nuclear reactor conditions, suggesting that a different boiling performance, if any, may not be
 caused to a wettability effect.

140 2.4. Empirical correlations for contact angle

141 There have been attempts to model the temperature dependency of contact angle using approaches such as Neumann's
142 theoretical model [46, 47] based on the equation of the state and Adamson's semi-empirical correlation [48]. Unfortunately, there
143 is hitherto no universal model that can accurately capture the temperature dependency of contact angle on different surface materials.
144 In this study, we have chosen a practical approach and developed material-specific empirical correlations for the contact angle of
145 DI water in saturated atmospheres. We define the non-dimensional contact angle θ^* and temperature T^* as follows

$$\theta^* = \frac{\theta}{\theta_0} \quad (2)$$

$$T^* = \frac{T_c - T}{T_c - T_0} \quad (3)$$

146 Here, θ is the contact angle at any given temperature; θ_0 is the contact angle at 0 °C; T is any given temperature; T_c is critical
147 temperature of water, i.e., 374.15 °C, and T_0 is 0 °C. In configuring the correlation, two widely accepted hypotheses are used, i.e.,
148 that a) the contact angle becomes null at the critical point of water and b) the contact angle at 0 °C is approximately the same as at
149 room temperature, i.e., 90° for Zr-4 and 58° for Cr-coated Zr-4. With these two assumptions, the final form of the contact angle-
150 temperature non-dimensional correlation is obtained by fitting the measured data points with hyperbolic tangent functions.
151 Equations (4) and (5) are the correlations for smooth Zr-4 and chromium, respectively.

$$\theta^* = \tanh(2.2 T^{*1.55}) \quad (4)$$

$$\theta^* = \tanh(3.5 T^{*2.70}) \quad (5)$$

152 Dimensional and non-dimensional contact angle values predicted by these correlations are compared to experimental results in
153 Fig.4 (left) and Fig. 4 (right), respectively.

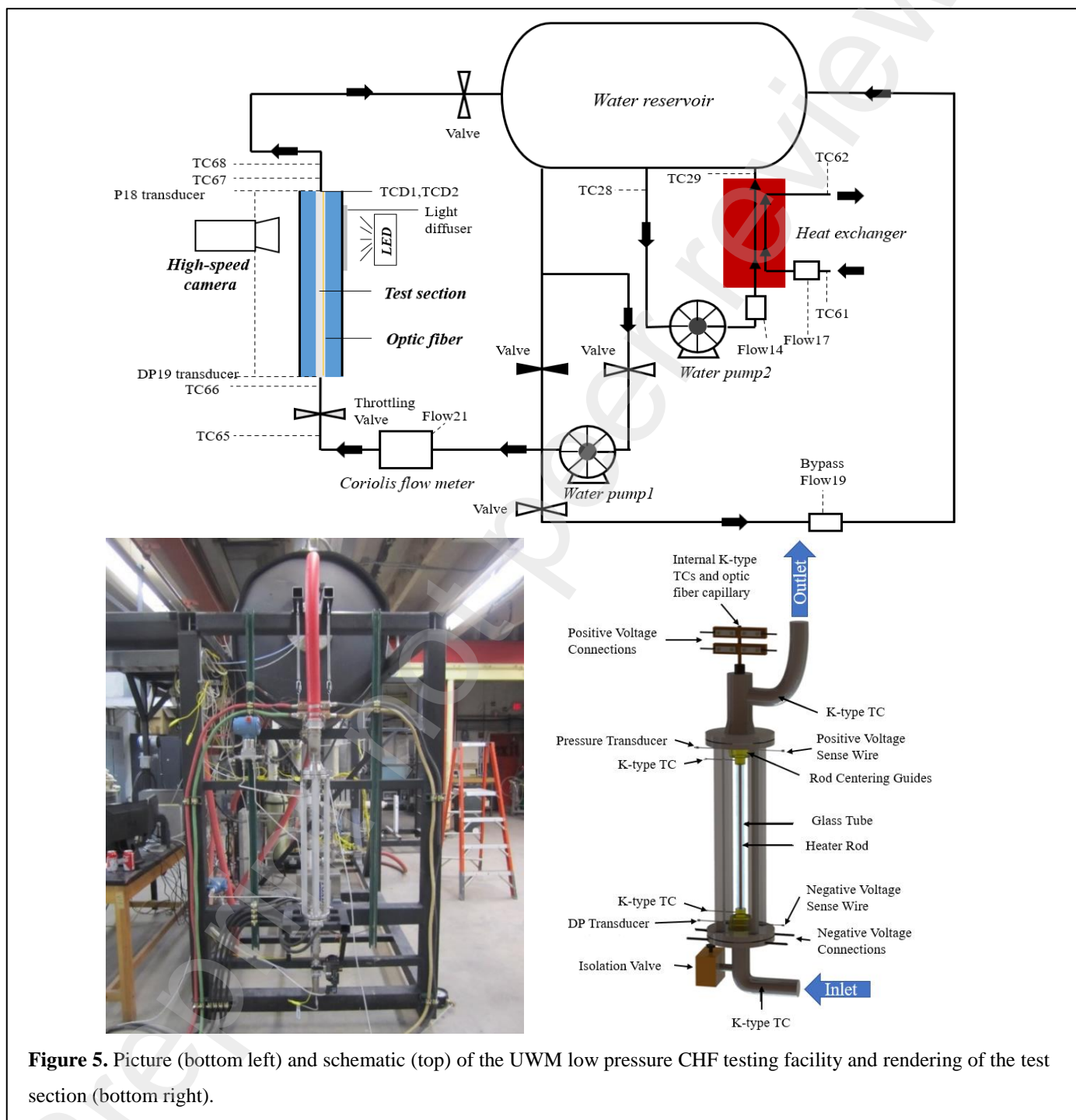
154 3. Low-pressure, steady-state CHF tests

155 3.1. Low-pressure flow boiling test setup

156 Figure 5 shows a schematic of the UWM experimental facility (top) and test section (bottom right) for the near-atmospheric,
157 low pressure experiments (i.e., the low-pressure loop). DI water from the reservoir (i.e., a 1.5 m³ stainless steel tank) is driven
158 through the loop by a centrifugal pump; the frequency of which is controlled by a variable frequency drive (VFD). Downstream
159 the pump the water flows through a Coriolis mass flow meter, a throttling valve (used to mitigate thermal instabilities associated
160 to the bubble nucleation and growth inside the test section during the boiling experiments) and enters the test section. The test
161 section consists of a vertical concentric annular channel (in which the fluid flows upwards), defined by a heater rod (internal
162 boundary) and a glass tube (external boundary), as shown in Figure 1 (bottom right). The rod has a heated length of 456 mm and
163 an outer diameter (OD) of 9.5 mm, while the glass tube has an inner diameter (ID) of 20 mm and is 500 mm long. Centering pins
164 are used to ensure that the heater rod is concentric with the glass tube. During the experiments, the heater rod is directly heated by
165 Joule effect using an 80V-1200A DC power supply. Downstream the test section the water returns to the tank, closing the circuit.
166 To remove the heat added to the water in the loop at the test section and control the temperature in the reservoir, a separate water
167 stream is driven directly from the reservoir through a secondary loop by a VFD-controlled centrifugal pump. In this secondary
168 loop, water flows through a plate-type heat exchanger for heat removal and then returns to the reservoir. On the secondary side of
169 the heat exchanger, a chilled glycol solution provides the ultimate heat sink. Water temperature measurements are performed at
170 the inlet and outlet of the test section through K-type thermocouples. The absolute pressure is measured at the outlet of the test
171 section, and pressure drop across the test channel through a differential pressure transducer. The heater rod is filled with boron
172 nitride to add thermal inertia. It features wall temperature measurements near the end of the heated section (i.e., the location where
173 CHF is expected to occur since the heat flux obtained by Joule heating is uniform) through a K-type thermocouple and axially
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along the test section length through an optical fiber inserted into a capillary. The disposition of the thermocouple and optical fiber along the cross section of the heater rod is illustrated in Figure 6 (left). The optical fiber temperature measurement technique uses pulsed infrared (IR) light to measure the Rayleigh backscattering from grain boundaries and other random imperfections in a section of standard optical fiber, as illustrated in Figure 6 (right). The optical fiber system has a maximum sampling frequency of 250Hz with a spatial resolution of 2.5 mm. The reader interested in the details of optical fiber temperature measurement is directed to Refs. [12, 49].



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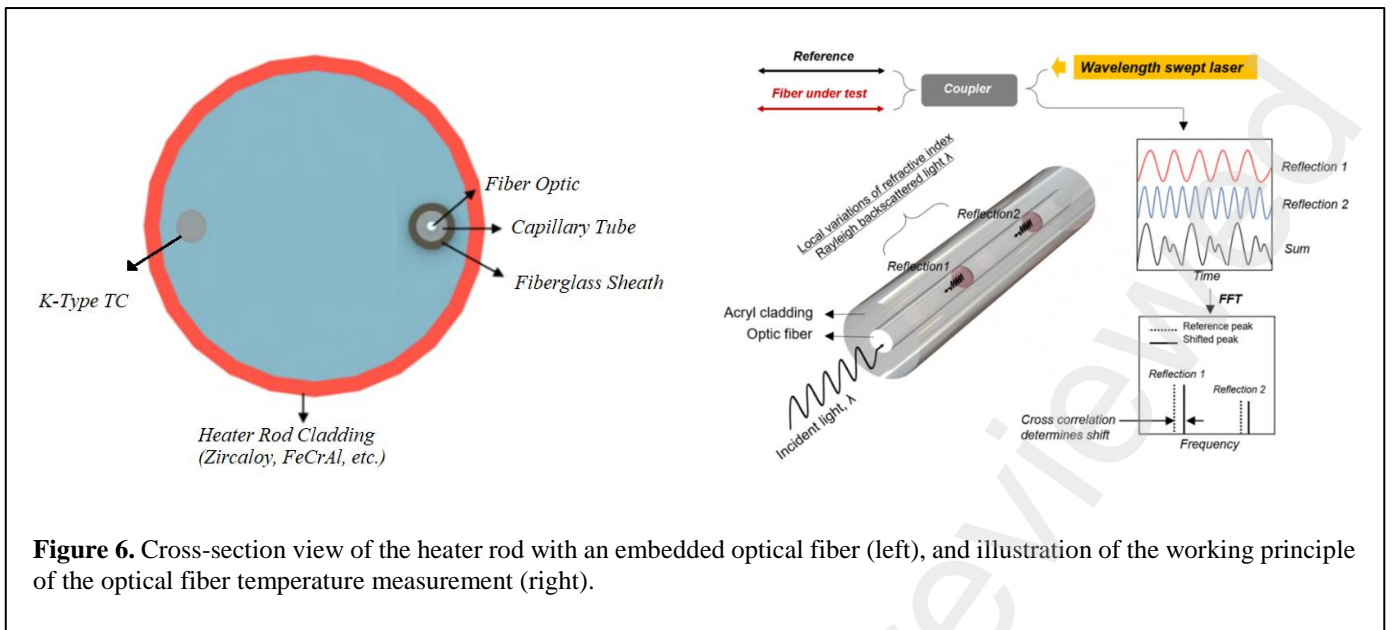


Figure 6. Cross-section view of the heater rod with an embedded optical fiber (left), and illustration of the working principle of the optical fiber temperature measurement (right).

In this study, we test the low-pressure flow boiling CHF limit of heater rods made of four different cladding materials, i.e., bare zircaloy-4, two types (CS-R and CS-S) of cold-sprayed Cr-coated zircaloy, and PVD (Physical Vapor Deposition) Cr-coated zircaloy. CS-R and CS-S have a different surface roughness. CS-S is polished after the cold spray deposition process in order to have a surface roughness similar to one of the bare zircaloy rod. CS-R is instead tested as it is after the cold spray depositions process. The average roughness and the contact angle (of DI water with air at ambient pressure) for each cladding, both measured in a previous study [50], are summarized in Table 2.

Table 2. Summary of the surface properties for all tested rod materials.

Test Material	Ra [μm]	Contact angle [$^\circ$]
Zircaloy	0.43	75
Cr-Coated Zircaloy (CS-R)	1.22	45
Cr-Coated Zircaloy (CS-S)	0.30	50
Cr-Coated Zircaloy (PVD)	0.45	25

Note that the contact angles on the PVD coated chromium surface is fairly different from the value measured at 1 bar in the high-pressure autoclave. This difference may arise from the gaseous atmosphere, surface finish (as the rod tested here is much rougher), PVD deposition parameters, or contamination of the surface. However, the results in Table 2 seem to show that the cold-sprayed surfaces have very similar wettability, even though CS-R is much rougher than CS-S (note that the suffix R stands for “rough”, whereas the suffix S stands for “smooth”).

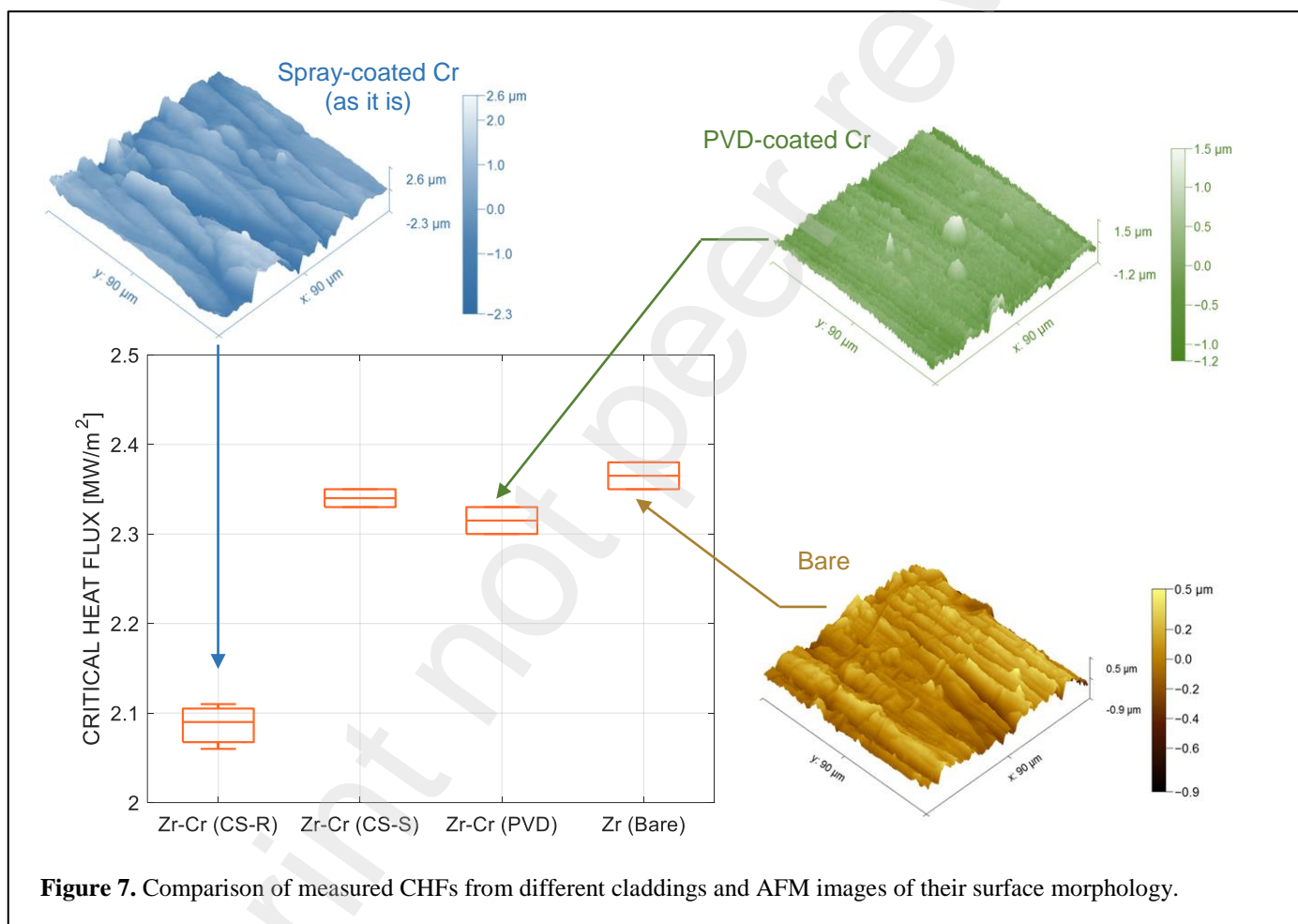
3.2. Test procedure and DNB criterion

The heat flux is gradually increased until departure from nucleate boiling (DNB) occurs. In the early stage of the experiment, the heat flux is increased in steps of $350 \text{ kW/m}^2\text{s}$. Inlet temperature, pressure and mass flow rate are monitored and compared to their measuring uncertainties to make sure that steady-state conditions are attained at each step. Such a procedure is repeated until the heat flux reaches $1750 \text{ kW/m}^2\text{s}$. Beyond this point, the heat flux is increased in steps of 1.25 kW/m^2 every 10 seconds until DNB occurs. The DNB is detected by monitoring either the temperature of the boiling surface (measured by the optical fiber and the thermocouple) or the heater rod electrical resistance. As there is only one optical fiber in these low-pressure tests, the optical fiber alone may not timely detect the DNB temperature rise if it occurs in a circumferential location far from it. The same concern exists for the single thermocouple located 25 mm from the end of the heated length. Therefore, the heater rod resistance is employed as an additional indicator to detect the DNB occurrence in these low-pressure tests. The temperature threshold for the DNB

207 detection is ~ 25 °C above the one at the previous steady-state heat flux step, while the resistance threshold is set at 0.03Ω above
208 the normal operating resistance of the heater rod. The CHF limit is the heat flux at the last stable step.

209 3.3. Low-pressure experimental results (near atmospheric conditions)

210 Low-pressure experiments are performed with an inlet subcooling of 80 °C, a mass flux of $750 \text{ kg/m}^2\text{s}$, and an outlet pressure
211 of 116 kPa . For the same type of cladding, experiments are repeated on 2 separate heater rods to reduce random uncertainty and
212 guarantee the repeatability of the results. The measured CHF values from all the tests are plotted in Figure 7, together with AFM
213 (Atomic-Force Microscopy) images of the different surfaces (except CS-S). These results reveal that there is no significant
214 difference in CHF for surfaces with the same roughness, despite the different contact angles (see Table 2). Instead, the CHF on the
215 rougher surface is slightly lower (approximately -10.8%) than the other surfaces. We suspect that the deterioration of the CHF on
216 the rough surface (accompanied by an increase of the boiling heat transfer coefficient) is due to a decrease of the onset of nucleate
217 boiling temperature and an increase of the nucleation site density, consistent with the findings of previous studies from our group
218 [58].



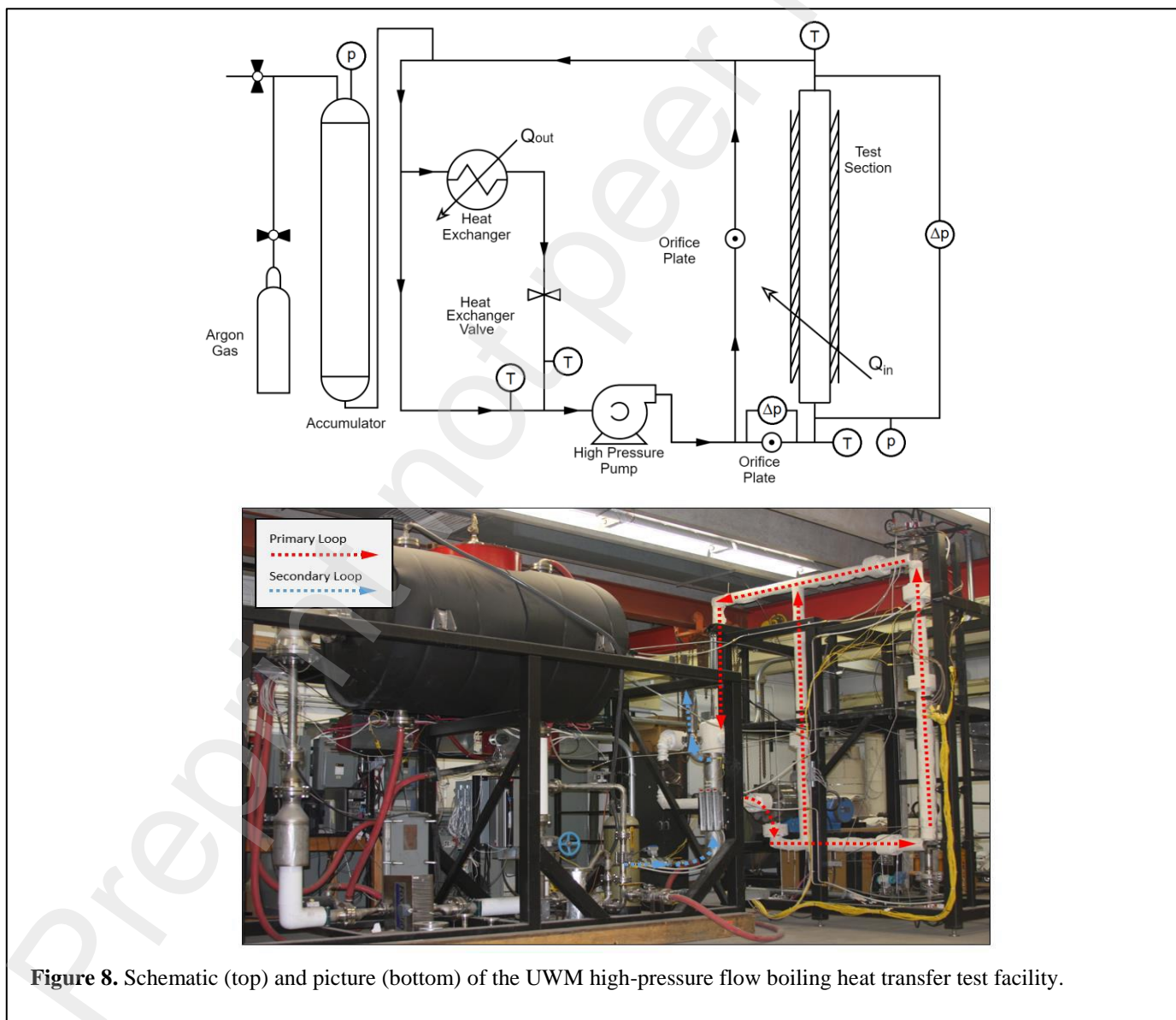
219 **Figure 7.** Comparison of measured CHF values from different claddings and AFM images of their surface morphology.
220 In summary, the results of the low-pressure experiments confirm that the surface wettability has a minor effect in the flow
221 boiling performance. Instead, surface roughness and morphology may play a bigger role, presumably due to their effect on
222 nucleation sites size and density.
223

224 **4. High-pressure, steady-state CHF tests**

225 High pressure flow boiling tests are performed using two different approaches and experimental facilities. Precisely, the UWM
226 high-pressure flow boiling facility, to explore CHF limits with a cosine shape heat flux distribution obtained using an indirect
227 heating technology, and the Westinghouse WALT loop, which also allows testing the effect of CRUD deposit on the boiling
228 surface, to measure the CHF limit with uniform heat fluxes.

229 *4.1. High-pressure flow boiling experiments with a cosine shape heat flux distribution*

230 Figure 8 shows a schematic and a picture of the UWM high-pressure loop. DI water is driven through the loop by a high-pressure
231 high-temperature centrifugal pump (Chempump™, GCT-5k 36I). Downstream the pump the water flow splits between two
232 branches, one that goes upward through the test section, where the working fluid is heated and boiled, and the other one that by-
233 passes it. Orifice-type mass flow meters are installed in each one of these two branches. After the test section, the two streams
234 merge into a single channel, and flows toward a shell-and-tube heat exchanger. A by-pass valve is used to control the flow rate
235 and, consequently, the heat removed from the working fluid at the heat exchanger in order to control the temperature of the flow
236 at the inlet of the test section. After the heat exchanger, the chilled water stream mixes back with the by-passed stream before
237 returning to the pump. An accumulator with an argon gas plenum is connected to the top region of the loop and it is used to set the
238 system pressure.

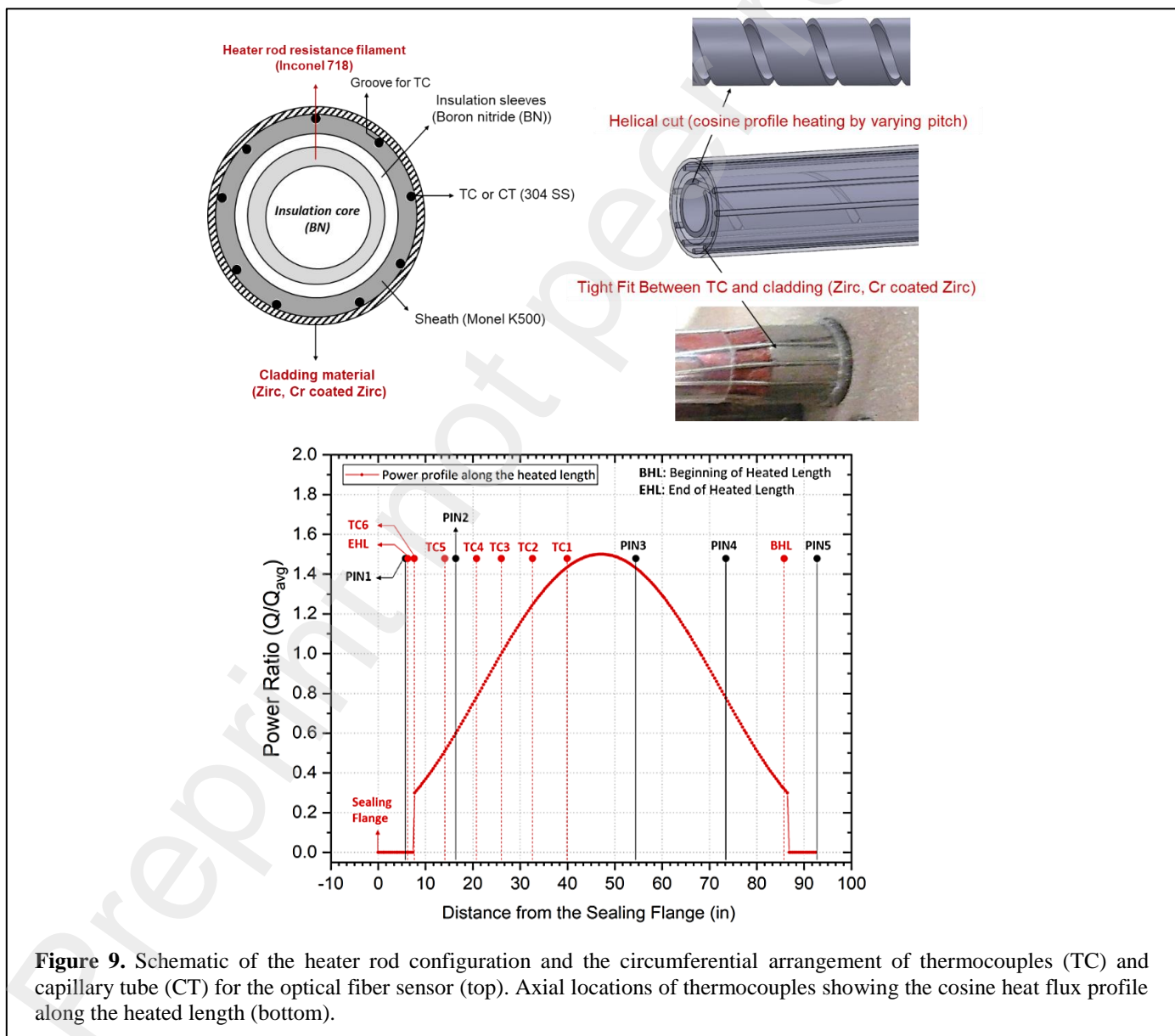


239 **Figure 8.** Schematic (top) and picture (bottom) of the UWM high-pressure flow boiling heat transfer test facility.

240 The high-pressure loop can operate in both subcooled and saturated conditions at maximum pressure and temperature of 25 MPa
 241 and 400°C separately. The test section of the high-pressure loop consists of an annular flow channel defined by a heater rod with
 242 OD of 9.5 mm and an outer tube of 20 mm ID, similar to the low-pressure one. In this case, however, the heater rod is ~3 m long
 243 with a heated length of 2 m, and the outer tube is made of Inconel. The rod is held centered in place though alignment pins
 244 distributed along the test section length. The location of these pins is carefully picked to not interfere with the DNB occurrence.
 245 The heater rods are designed ad-hoc by Stern lab to provide a cosine heat flux profile that simulates the axial heat flux distribution
 246 along an actual LWRs fuel rod. The heat flux distribution follows the equation

$$\frac{q''(x)}{q''_{avg}} = \theta_0 + \theta_1 \cos \left[2\theta_2 \left(\frac{x}{L_{HL}} - 0.5 \right) \right] \quad (6)$$

247 where $\theta_0 = 0.819$, $\theta_1 = 0.681$, $\theta_2 = 2.436$, q''_{avg} is the average heat flux along the rod, defined by the ratio between the power
 248 applied and the area of the heated surface, and L_{HL} represents the total heated length, i.e., 2 m. The cosine profile heat flux is
 249 obtained by a specially tailored helical Joule-heating filament made of Inconel 718 (illustrated in Figure 9 (top)) to heat up the
 250 heater rod cladding via thermal conduction from the internal surface.



251
 252

253 Differently from the low-pressure loop, the test section for the high-pressure loop experiments is indirectly heated. Two different
 254 claddings are used for the experiments in the high-pressure loop, one bare zircaloy and one Cr-coated zircaloy with the same
 255 surface properties as the CS-S used for the low pressure experiments. The surface properties of the two tested claddings are
 256 consistent with those for the low-pressure tests summarized in Table 2.

257 The heater rods are equipped with 6 K-type thermocouples and 3 optical fibers at different axial and circumferential positions,
 258 as illustrated in Figure 9 (top), to characterize and identify the correct location where the DNB starts. Precisely, the temperature
 259 threshold for the DNB detection is ~ 25 °C above the one at the previous steady-state heat flux step. Note that one cannot use the
 260 cladding resistivity to detect DNB, as there is no electric current circulating through the cladding. Also, the CHF may not occur at
 261 the very end of the test section (as expected for directedly-heated low-pressure experiments where the heat flux is uniform), because
 262 of the cosine-shape heat flux profile. The thermocouples are axially distributed in the downstream half of the heater rod, as shown
 263 in Figure 9 (bottom), where the DNB is expected to occur, and the optical fibers equally spaced of 120° along the rod circumference,
 264 see Figure 9 (top). The sampling rate of the optical fiber system is 100 Hz at a spatial resolution of 2.5 mm. The test procedure of
 265 these high-pressure tests is similar to the one of the low-pressure tests. The high-pressure test matrix is summarized in
 266 Table 3. These conditions are selected similar to those investigated by Wheatherhead [59].

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Table 3. Nominal operating conditions for high-pressure CHF tests.

Pressure [MPa]	Mass Flux [$\text{kg}/\text{m}^2/\text{s}$]	Subcooling [$^\circ\text{C}$]
10.3	1965	10, 30, 40
10.3	2712	10, 20, 30
15.2	1965	10, 30, 50
15.2	2712	10, 20, 30
20.0	1965	20, 30, 50
20.0	2712	20, 30, 50

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270 Figure 10 shows a comparison between the total power delivered through the heated rod when the DNB is detected. Note that the
 271 total power is a proxy for the average surface heat flux. These results show quite convincingly that there are no practical differences
 272 between the bare and chromium coated surfaces. In some cases, the results are perfectly overlapped (i.e., it is impossible to
 273 distinguish the marker for the bare and the coated surface). Such conclusion is corroborated by the comparison between coated and
 274 uncoated surfaces proposed in Figure 11 (left). Figure 11 (right), instead, shows a comparison of the local CHF values (i.e., the
 275 heat flux at the location where DNB starts) obtained for the bare and Cr-coated heater rods. The results shown in Figure 11 are
 276 obtained from the optical fibers temperature data. The local CHF is estimated from the cosine heat flux profile (Eq. 6), considering
 277 the location of the CHF as the one associated to a temperature excursion nearest to the inlet of the test section. At a pressure of
 278 10.3 MPa, the local CHF percentage difference, defined as the absolute difference divided by the bare zircaloy value, varies from
 279 9.8% to 82.3%. Most tests, i.e., 6 out of 7, show slightly higher CHF for the Cr-coated heater rod. At 15.2 MPa, a reduction in the
 280 local CHF difference range occurs, ranging from 4.5% to 44.2%. The tests show either close or slightly higher CHF for the Cr-
 281 coated heater rod, see Figure 11. At the highest test pressure of 20 MPa, the local CHF difference range reduces even further, to
 282 values from 5.6% to 29.7%. Most tests, i.e., 4 out of 6, show close CHF values between the bare and Cr-coated heater rod. In
 283 summary, it can be indicated that CHF values from Cr-coated heater rods are either close to or slightly higher than the values for the bare
 284 zircaloy heater rods with an average difference of 18.8%. It is worth to highlight that, under certain conditions, different CHF
 285 location are noticed based on the optical fibers temperature data. This indicates that the dryout starts in one spot and spreads
 286 downstream. Such a phenomenon leads to an uncertainty of the local CHF estimative, possibly being the culprit for the differences
 287 shown in Figure 11 (right). More details about these high-pressure experiments can be found in a parallel communication [60].

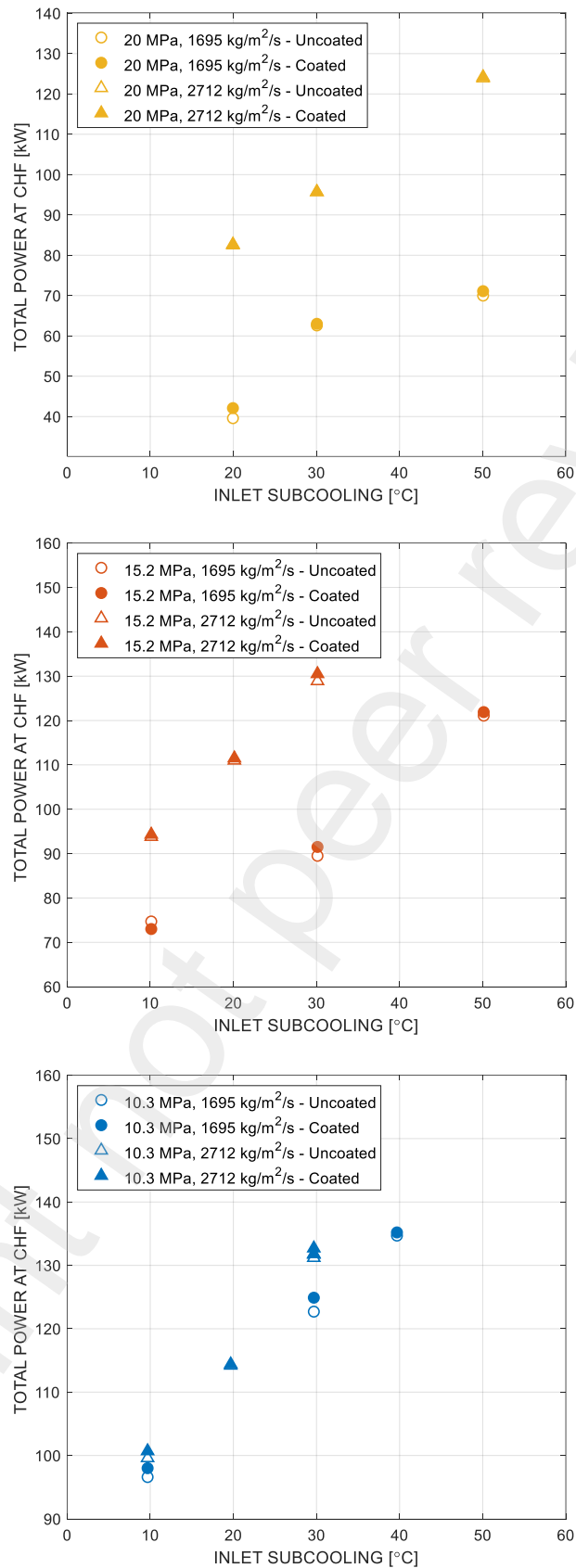


Figure 10. Comparison of total power at CHF between bare and coated heater rod: (bottom) 10.3 MPa (middle) 15.2 MPa and (total) 20 MPa nominal pressure condition.

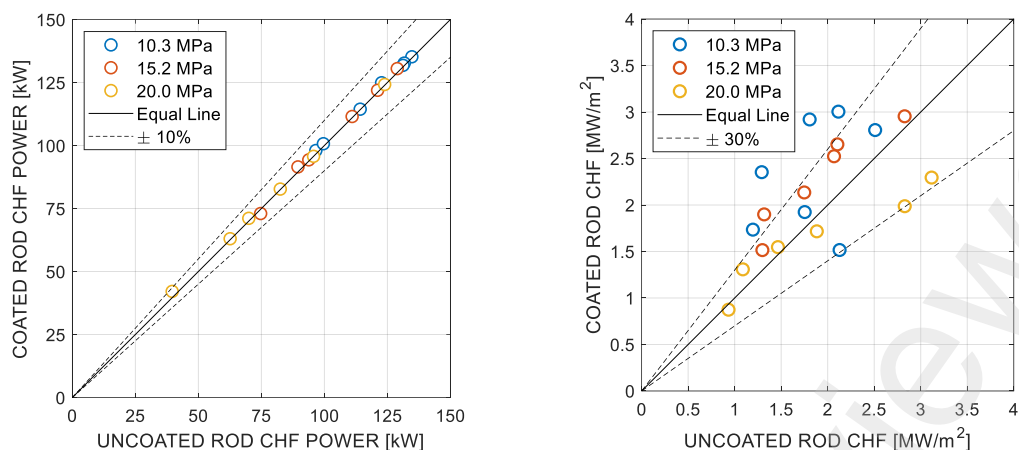


Figure 11. Comparison of local CHF (right) and total CHF Power (left) between bare and coated heater rod: (a) 10.3 MPa (b)15.2 MPa and (c) 20 MPa nominal pressure condition. Dotted lines represent the $\pm 30\%$ deviation.

4.2. High-pressure flow boiling experiments with a uniform heat flux distribution

The Westinghouse Advanced Loop Tester (WALT) loop allows investigations of flow boiling DNB at prototypical PWR pressure, temperature, and water chemistry. The schematic of WALT loop is shown in Figure 12. A centrifugal pump drives the specially treated water into the autoclave from top. The water first flows downwards through the downcomer and then changes direction to flow vertically upwards through the inner chimney, which comprises the WALT loop test section. The test section features an annular flow path, which is formed by the concentric inner chimney with an ID of 20.96 mm and the heater rod with an OD of 9.5 mm. The WALT loop is capable of operating up to a 17.22 MPa system pressure, a 6 m/s linear velocity, and close to local saturation temperature at the inlet of the inner chimney. One unique capability of WALT loop is to simulate the PWR water chemistry in addition to the prototypical pressure and temperature. The coolant used during the DNB tests consists of 1000 ppm of Boric Acid, and 4 ppm of LiOH dissolved in high purity DI water, which is consistent with the PWR reactor coolant system (RCS). Hydrogen was added to the coolant at concentrations between 25 and 50 cc/kg. The WALT heater rod is shown in Figure 13. It has a total length of 445 mm, with a central heated length of 133 mm. The heater rod is connected to a DC power supply and generates by Joule effect a uniform heat flux. Four thermocouples attached to the inner surface of the heater rod (see Figure 13) are used to monitor the cladding temperature and detect the occurrence of DNB events. Two sets of heater rods, i.e., uncoated Optimized ZIRLO rods and cold spray Cr-coated Optimized ZIRLO rods, are used in these experiments. In addition, the heater rods have been coated with a simulated crud deposit to evaluate the crud effect on CHF. The crud deposits are created by adding FeEDTA, NiEDTA, and colloidal crud precursors to the WALT loop during medium heat flux boiling [49].

After the system reaches the desired system pressure, chimney inlet temperature, and mass flux, the DNB test starts from an arbitrary single-phase convective heat transfer point, with the power being gradually increased until DNB is achieved. The heat flux step increment is gradually reduced when DNB conditions are approached. The heater rod power is cut off based on the cladding temperature increasing rate. Precisely, the power shutdown occurs if the cladding temperature increases at a rate of 20 °C per second or faster. The DNB occurrence is also confirmed by a rapid change in the rod electrical resistance, due to a rapid change of the cladding temperature. The CHF is the average heat flux at the operating condition prior to the heater rod temperature spike.

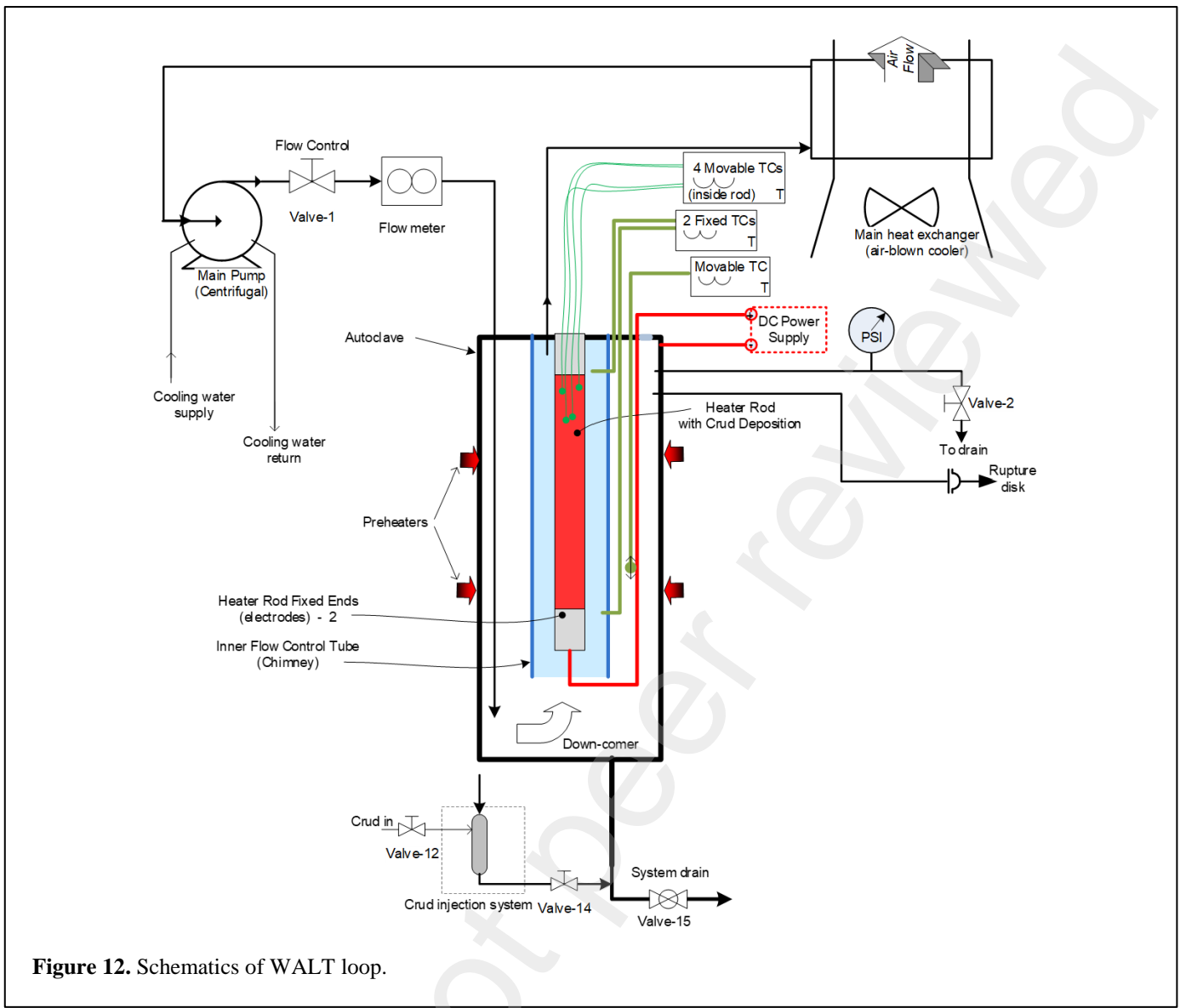


Figure 12. Schematics of WALT loop.

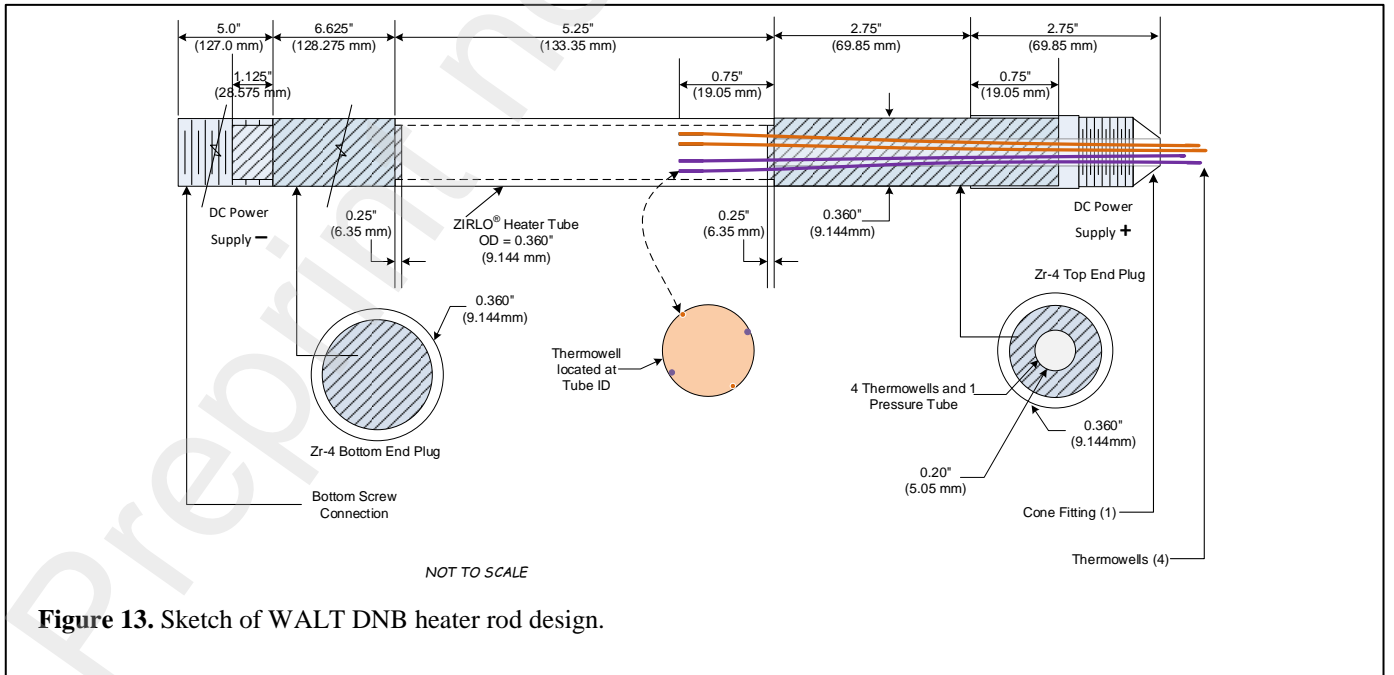


Figure 13. Sketch of WALT DNB heater rod design.

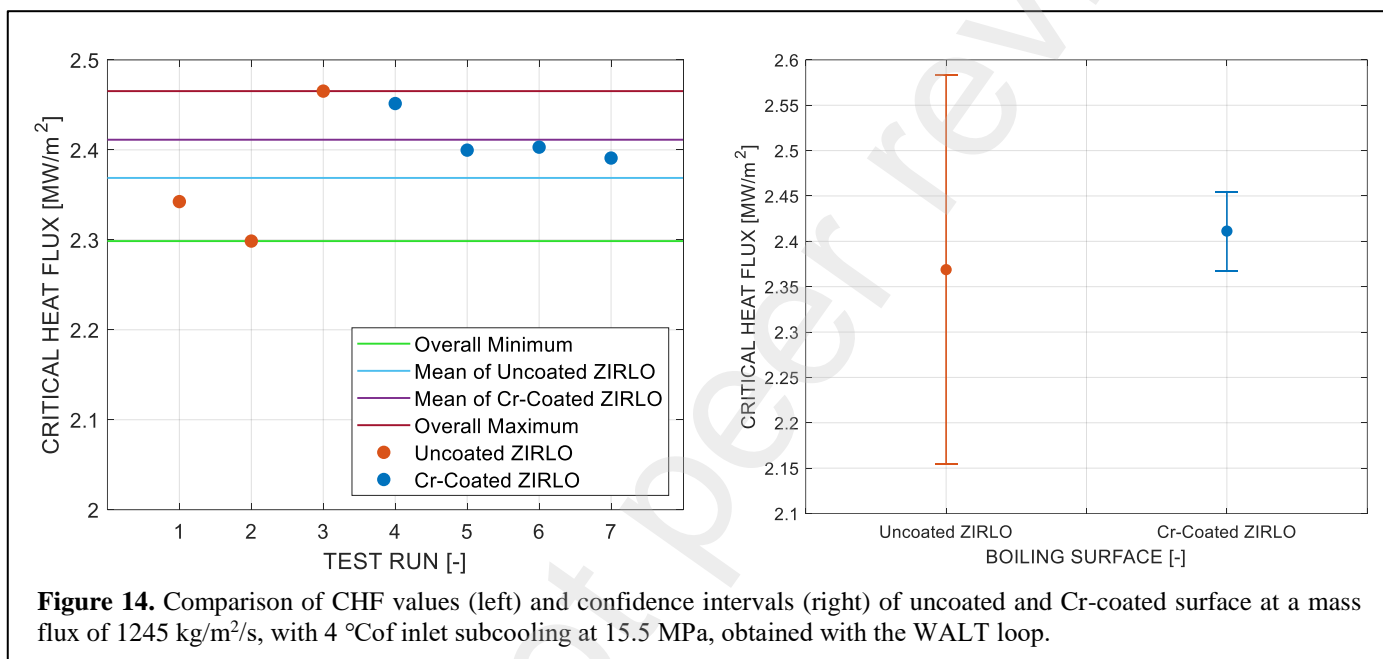
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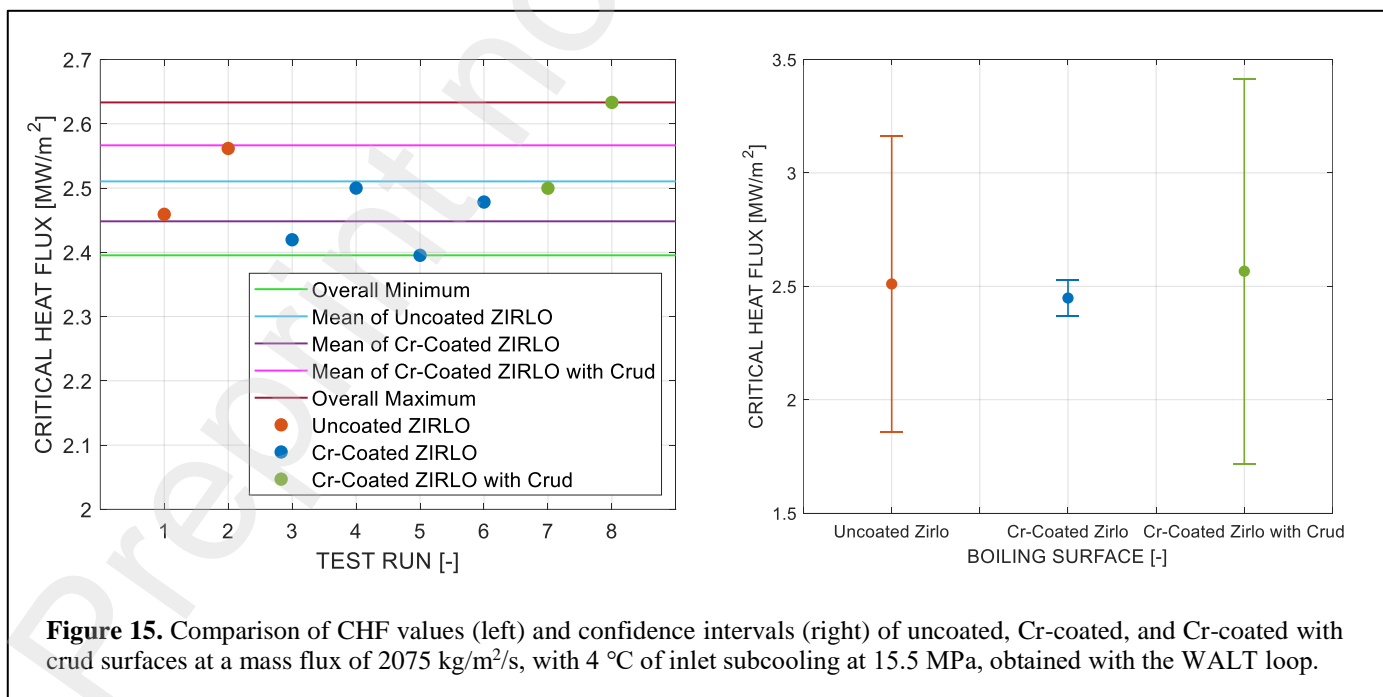
Experiments are performed for uncoated Optimized ZIRLO and Cr-coated Optimized ZIRLO rods, at a nominal system pressure of 15.5 MPa, a chimney inlet subcooling 4 °C, and a mass fluxes of 1245 kg/m²/s. To understand better the effect of crud, test results for Cr-coated Optimized ZIRLO with crud [49] are also discussed in this analysis. The tests to evaluate the effect of CRUD were performed with a higher mass flux (2075 kg/m²/s). The thickness of the crud deposits is about 40 μm. The repeatability of the CHF experiments is confirmed by testing different heater rods for each cladding surface condition.

CHF experimental results of the run at mass flux of 1245 kg/m²/s are shown in Figure 14 (left). This figure shows that the CHF values for the Cr-coated Optimized ZIRLO rods are similar to those for uncoated Optimized ZIRLO rods. Assuming that the scatter of CHF values follows the t-distribution, it is possible to calculate the 95% confidence interval of the mean for uncoated and Cr-coated test results as shown in Figure 14 (right). The confidence interval of the uncoated results completely covers the confidence interval of the Cr-coated results. Practically, there is no statistical difference in the mean CHF values for uncoated and Cr-coated tests. Similar conclusion can be drawn from the result at mass flux of 2075 kg/m²/s shown in Figure 15 (left).



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Figure 14. Comparison of CHF values (left) and confidence intervals (right) of uncoated and Cr-coated surface at a mass flux of 1245 kg/m²/s, with 4 °C of inlet subcooling at 15.5 MPa, obtained with the WALT loop.



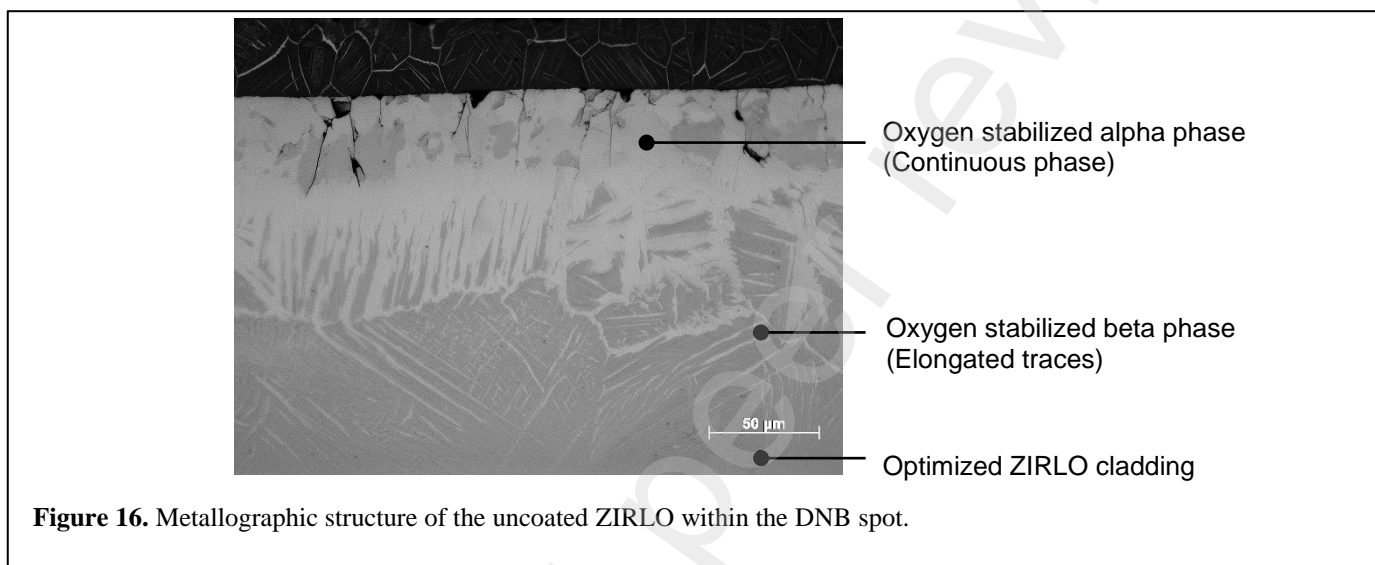
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Figure 15. Comparison of CHF values (left) and confidence intervals (right) of uncoated, Cr-coated, and Cr-coated with crud surfaces at a mass flux of 2075 kg/m²/s, with 4 °C of inlet subcooling at 15.5 MPa, obtained with the WALT loop.

330 Neither the Cr-coating nor the crud deposits deteriorate CHF compared to uncoated test results. The confidence interval of the
331 uncoated results largely overlaps with the confidence interval of the Cr-coated and Cr-coated with crud results as shown in Figure
332 15 (right). Note that the large confidence intervals for uncoated and Cr-coated with crud results in Figure 15 (right) come from the
333 small number of the data points rather than the scattering of the data themselves.

334 4.3. Post DNB metallography analyses

335 The microscopic structure of the uncoated and Cr-coated heater rods used in the WALT loop testing is examined at the DNB
336 location through a Scanning Electron Microscope (SEM). The metallographic structure of the uncoated ZIRLO rods is shown in
337 Figure 16. Oxygen rapidly penetrates the uncoated ZIRLO surface, forming a continuous brittle oxygen-stabilized alpha phase near
338 the surface that appears brighter in the SEM image. Deeper into the cladding, there is an elongated oxygen-stabilized beta phase
339 penetrating along the grain boundaries of metallic phase. This microstructure is typical of a zircaloy cladding heated to high
340 temperatures in steam or water.



342 Different from the uncoated ZIRLO rods, the cold-sprayed Cr coatings remain well bonded to the underlying ZIRLO and protect
343 it from excessive oxidation and cracking, as shown in Figure 17. The brittle zirconium oxide phase is inhibited as long as an intact
344 metallic chromium layer exists. The extent of cracking in the Cr-coated cladding after DNB events is much less than in uncoated
345 cladding. However, once all metallic chromium is converted into a Cr-Zr intermetallic phases, oxidation proceeds and cracking in
346 the brittle oxygen-containing phases and intermetallic phases is similar to that found in uncoated ZIRLO. It is worth of mention
347 that damages to the external surface of the rod are visually observed on the bare zircaloy rods tested in the UWM loop, whereas
348 such damages are not visible on the Cr-coated surfaces [12].

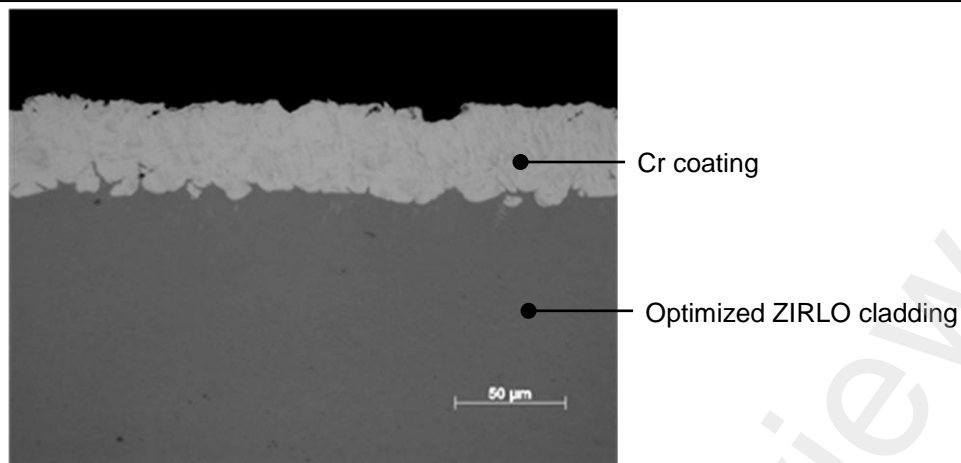


Figure 17. Metallographic structure of the Cr-coated ZIRLO within the DNB spot.

5. Transient CHF test

5.1. Flow boiling facility and infrared heaters

Transient CHF experiments are performed in subcooled flow boiling conditions with an exponentially increasing heat flux (i.e., $q'' \propto e^{t/\tau}$, where τ is the period of the power escalation), simulating the reactor power increase during a hypothetical reactivity accident. These tests are run using the MIT Pether platform. The test section features a rectangular cross-section channel ($1 \times 3 \text{ cm}^2$). It stands on top of a developing channel with the same cross-section, and more than 60 hydraulic diameters long to guarantee fully-developed turbulent flow at the actual test section. Each side of the test section channel has an opening (see an exploded view in Figure 18). Three sides accommodate quartz windows that provide optical access to the flow. The remaining side accommodates a Shapal™ cartridge, which, in turn, houses the heating element. This heating elements enable infrared (IR) measurements of the time-dependent temperature and heat flux distribution on the boiling surface with a high temporal (0.4 ms) and spatial resolution (123 μm/pixel). Heater, cartridge, quartz windows, and test section, are all accurately flush with each other to form a rectangular channel with perfectly flat walls. The reader interested in the details of the apparatus is directed to Ref. [52].

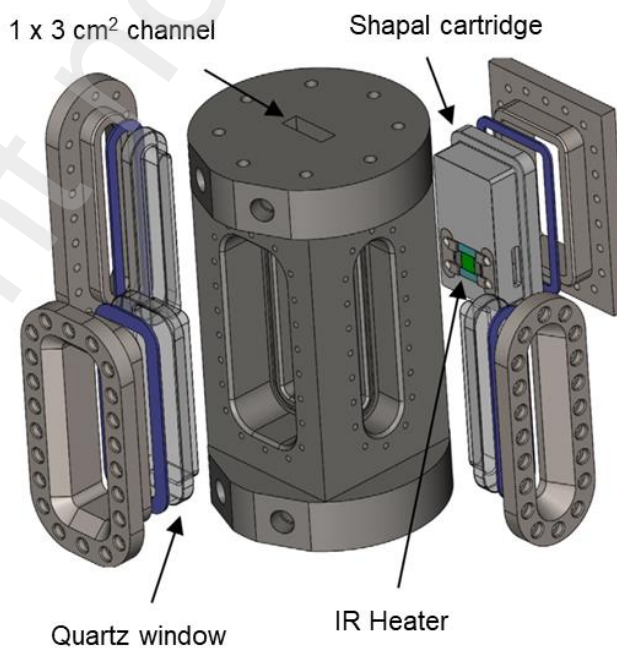
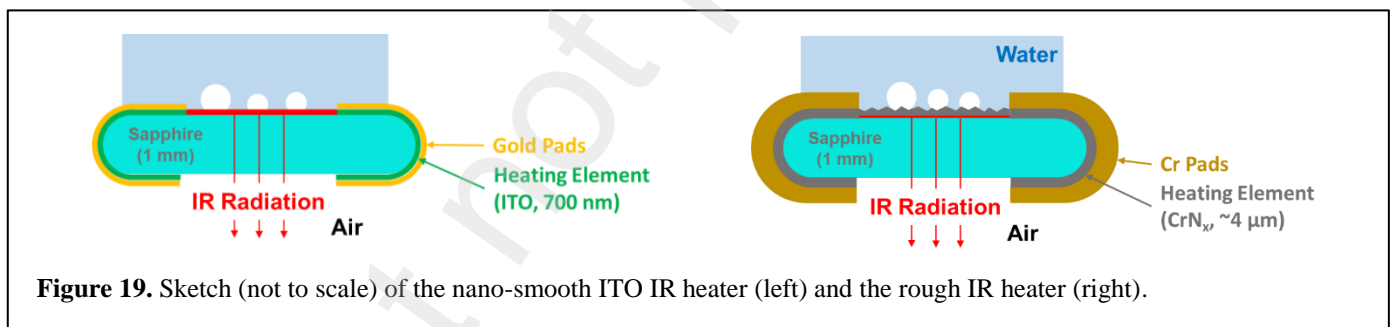


Figure 18. Exploded view of the 3D design for the flow boiling test section at MIT.

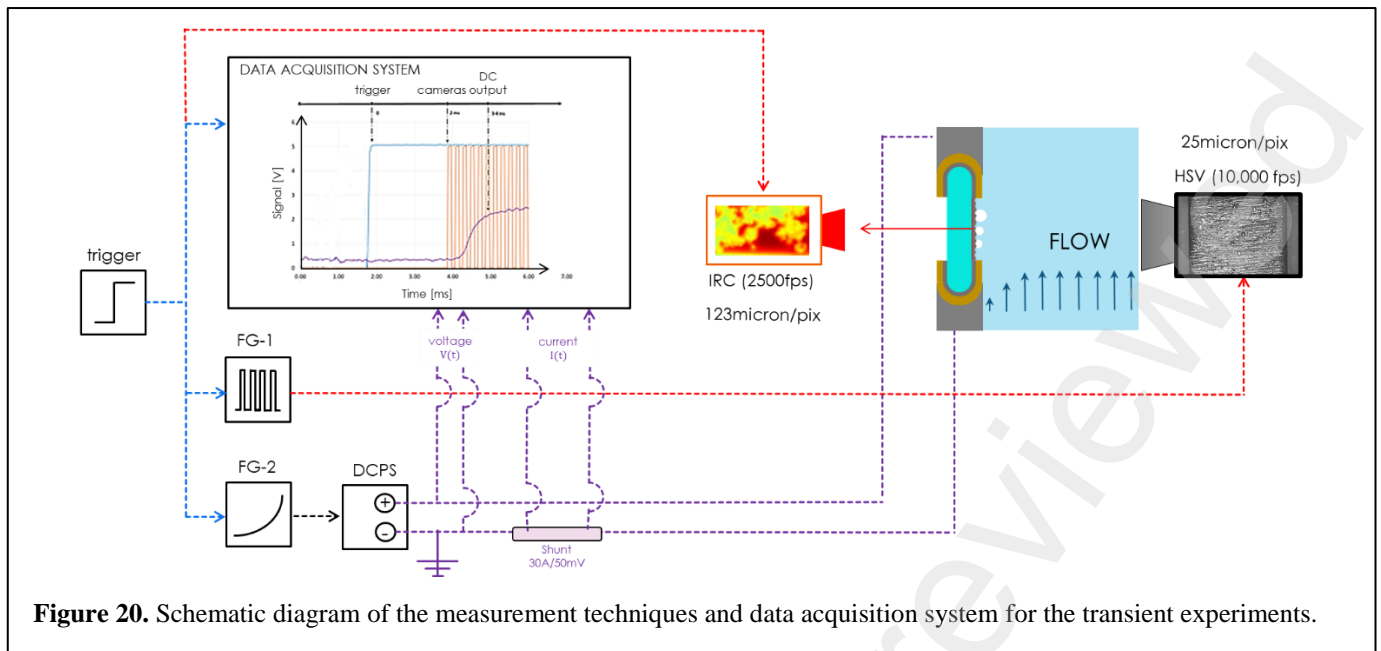
364 We conduct transient flow boiling experiments on a surface mimicking the chromium coated nuclear reactor cladding and
365 compare the results to those previously obtained using a nano-smooth surface, which is supposed to have a radically different
366 boiling behavior. The nano-smooth IR heater, and consists of a 1 mm thick sapphire substrate, coated with a 700 nm thick indium
367 tin oxide (ITO) layer, as shown in Figure 19 left. The ITO, in contact with water, has an active heating area $\sim 1 \times 1 \text{ cm}^2$ (red treat in
368 Figure 19 left). The ITO is very thin and has a negligible thermal resistance and heat capacity. Practically, the temperature at the
369 sapphire-ITO interface coincides with the temperature at the ITO-water interface, i.e., the boiling surface. Importantly, the ITO
370 layer is completely IR opaque. Thus, the radiation that it emits can be used to measure the time-dependent temperature and heat
371 flux distributions on the boiling surface. Practically, this task requires the solution of an inverse problem coupling optical radiation
372 and conduction heat transfer in the sapphire substrate. The reader interested in the details of this technique is directed to Refs.
373 [40,53].

374 It is worth mentioning that the thermal properties of sapphire are similar to zircaloy [52]. Therefore, the thermal response of
375 sapphire and zircaloys to a bubble life cycle is expected to be similar. In summary, the key difference between nano-smooth
376 sapphire-ITO IR heaters and commercial-grade heating surfaces is the surface finish and wettability. In this work, we develop a
377 new infrared heater, whose surface mimics the finish and wettability of a chromium coated fuel cladding. The new heater builds
378 upon the same type of substrate as the nano-smooth IR heater, as shown in Figure 19 right. It features a $\sim 4 \mu\text{m}$ thick chromium
379 nitrate heating layer with an active heating area of $\sim 1 \times 1 \text{ cm}^2$ (red treat in Figure 19 right). Chromium nitride (CrN_x , with $x = 1$ to
380 2) is used to simulate the wettability of metallic chromium. The active CrN_x area is grinded using P360-grit sandpaper until the
381 desired roughness and scratch pattern as the commercial surface roughness is obtained. A sufficiently thick heating element layer,
382 i.e., $\sim 4 \mu\text{m}$, is required to survive the grinding, while still maintaining the uniform heating power distribution and negligible
383 thermal resistance and heat capacity. Note that metallic chromium is not suitable for this purpose due to very low electric resistance
384 at such thicknesses. The chromium nitrate layer is IR opaque. Thus, the radiation that it emits can be used to measure the time-
385 dependent temperature and heat flux distribution on the boiling surface using the same technique used for the nano-smooth ITO
386 heater.



387
388 The nano-smooth heater preserves the crystalline smoothness of the sapphire substrate (with an $R_a < 0.01 \mu\text{m}$) except for small
389 imperfections caused by the sapphire machining process, which may serve as nucleation sites. The rough IR heater has many
390 parallel scratches and a much higher surface roughness ($R_a \sim 0.3 \mu\text{m}$). The contact angle measured on nano-smooth ITO surface is
391 $86^\circ \pm 3^\circ$ in air at ambient pressure. A larger variation is noticed in the contact angle measured for the rough CrN_x surface, showing
392 $72^\circ \pm 3^\circ$ and $61^\circ \pm 5^\circ$ respectively for 2 randomly selected samples.

393 We use the same diagnostics and data acquisition technique as the one used in Refs. [40,43], as illustrated in Figure 20.



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A IRC806HS high-speed infrared (HSIR) camera captures the infrared radiation emitted by the boiling surface with a temporal resolution of 0.4 ms and pixel resolution of 123 μm . A Phantom v.12.1 high-speed video (HSV) camera driven by a function generator (FG-1) is used to capture images of the bubbles on the boiling surface. The HSV operates at an acquisition frequency of 10,000 fps and pixel resolution of 25 μm . A second function generator (FG-2) is used to drive the exponential output of the DC power supply (DCPS). The voltage and current are measured with an Agilent U2542A high speed data acquisition system (HSDAS). At the beginning of each test, the HSIR camera, HSDAS, FG-1 and FG-2 are simultaneously initialized with a single trigger signal. Further, a controlled delay between the trigger and outputs of FG-1 and FG-2 to capture the beginning of the transient is introduced.

The experiments consist in exponentially increasing the heat flux released by the heating element according to the desired power escalation period τ . Experiments are conducted in ambient pressure, with 50 $^{\circ}\text{C}$ of subcooling, and mass flux of 1275 $\text{kg}/\text{m}^2\text{s}$, which corresponds to a Reynolds number of 35000. The power escalation periods covered in the campaign vary from 5 milliseconds to 500 milliseconds. CHF is defined as the value of heat flux that creates the first irreversible dry spot on the boiling surface, which is the same criterion used in Ref. [43]. The identification of such conditions is possible thanks to the infrared measurements. An irreversible dry spot is such that it never gets rewetted once it appears on the boiling surface (unless there is a reduction of the applied heat flux). The reported CHF in this study is the average heat flux transferred to water at the moment when the first irreversible dry spot starts to grow on the surface. Details about the procedure to estimate the uncertainty on the measured CHF values are discussed in Ref. [43]. Briefly, the overall uncertainty is determined by three main contributions: the measurement uncertainty of the applied heat flux that is propagated from voltage and current measurement; the absolute temporal uncertainty that depends on the frame rate of the IR camera; and the standard deviation of CHF values measured in different runs.

The measured transient CHF values for the rough CrN_x heater are shown in Figure 21, compared to the values measured on the nano-smooth heater under the same exact operating conditions [43]. Two clear asymptotic regions emerge from the plot. The first asymptotic region, Asymptotic Region I, is observed for slow transients (i.e., large exponential power escalation periods with $\tau > \sim 100$ ms). The CHF values plateau for both heaters. Under such conditions, the boiling process is practically the same as observed with steady-state heat flux inputs. The second asymptotic region, Asymptotic Region II, corresponds to fast transients (i.e., small power escalation periods with $\tau < \sim 10$ ms). In this case, the CHF increases with decreasing periods (precisely, it is proportional to $1/\sqrt{\tau}$). However, while the CHF for the two different surfaces are, as expected, different under steady-state conditions (i.e., Asymptotic Region I), they converge to similar values as the period is shortened (i.e., Asymptotic Region II).

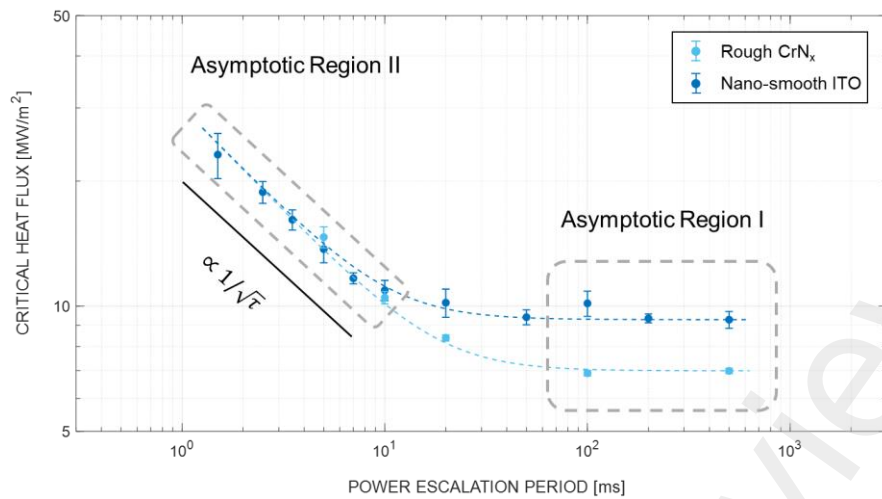


Figure 21. Transient CHF versus power escalation period for rough CrN_x and nano-smooth ITO surfaces. The grey dashed-line boxes indicate two asymptotic regions. The black solid line shows a reference trend of $\propto 1/\sqrt{\tau}$.

At long periods (with $\tau > \sim 100$ ms), the CHF is independent of the exponential escalation period. Under high subcooling conditions, such as the ones in the present tests, bubble growth is limited by the relatively thin superheated liquid layer adjacent to the wall. Bubbles condense whenever they grow beyond the superheated layer. As already reported in several studies [54-58], the CHF during steady state conditions depends on surface conditions, such as roughness and morphology. The difference in surface condition changes bubble parameters, such as nucleation site density, bubble radius, and, therefore, the way they interact with each other [58]. As a result, the CHF may be different. As observed in Figure 21, for Asymptotic Region I, the measured CHF value for the rough heater is about 75% of the value for nano-smooth heater at 500 ms transient power input. Such difference is consistent with the findings of our previous study with steady-state heat inputs [58].

Several studies attempt to model and quantify the transient CHF based on a presumed triggering mechanism of the steady state CHF [34,35,59]. In these studies, the boiling crisis triggering mechanism for transient or steady state boiling is assumed to be the same, no matter how fast is the transient, and the transient CHF value is predicted by multiplying the steady state CHF value with a scaling factor dependent on the power escalation period. The implication of this assumption is that, given different steady state CHF values for different surface conditions, the CHF values at short periods should also be different, since the proposed scaling factor is the same regardless of the surface conditions. In other words, one should expect the light blue dots in Asymptotic Region II of Figure 21 to be in a parallel asymptote, lower than the one for the dark blue dots. Clearly, this prediction is in contradiction with the present data: the CHF measurements for the two different boiling surfaces overlap with each other in Asymptotic Region II, as shown in Figure 21. These results suggest that there is a decorrelation between the steady state and the transient CHF mechanisms, questioning the theoretical basis and the validity of mechanistic models and correlations that aim to predict the transient CHF from steady state value, as discussed before. In addition, it is plausible that the boiling crisis mechanism for fast transient is different from the one under steady state conditions. The relative importance of different physical phenomena may change with power escalation periods. Surface effects, which are key parameters in the prediction of the CHF under steady-state conditions, seems to diminish at fast transient power escalations. CHF trends for different surfaces show the same dependency on the period for fast transient, i.e., inversely proportional to the square root of period ($\propto 1/\sqrt{\tau}$), as shown in Figure 21 and also reported in Ref. [43]. Such dependency is the same as the heat transfer coefficient for single-phase transient conduction. This observation suggests that transient conduction in the liquid may be the limiting heat transfer mechanism in fast transient conditions. In conclusion, while further investigations are necessary to consolidate these findings, it is plausible that chromium coatings and other surface effects do not negatively affect the transient CHF under rapid power escalation conditions.

451 6. Conclusion

452 Chromium-coated nuclear reactor fuel claddings can resist runaway steam oxidation under extreme temperature ($> 1200\text{ }^{\circ}\text{C}$)
453 and retain their mechanical strength and structural integrity under thermal shocks. They have the potential to replace conventional
454 zircaloy claddings. However, it is necessary to demonstrate that the chromium coating does not deteriorate the CHF limit compared
455 to conventional zircaloy claddings, in particular in high-pressure operating conditions.

456 In this paper, we present the results of a collaborative research effort aimed at evaluating the thermal-hydraulics performance
457 of chromium-coated zircaloy claddings. We have developed experimental facilities and conducted experiments to measure the
458 wettability of chromium coated and convention zircaloy cladding materials in vapor-saturated atmosphere from ambient pressure
459 up to the operating pressure of PWRs, as well as the flow boiling CHF limit at near-atmospheric pressure and at pressures close to
460 PWR conditions. The high-pressure flow boiling CHF tests were conducted with a uniform heat flux and a cosine-shape heat flux
461 profile, and exploring (for the uniform heat flux tests) the effect of CRUD deposits on the cladding surface. Our results demonstrate
462 that the chromium coating does not deteriorate the CHF limit of conventional zircaloy claddings, not even when covered by a
463 CRUD deposit. Instead, the chromium coating improves significantly the post-CHF behavior, as the it prevents the reaction
464 between zircaloy and steam, which results in the formation of a brittle zirconium oxide through the surface of the cladding. We
465 also conducted transient CHF under exponentially escalating heat flux inputs. The results of these transient heat flux tests suggest
466 that the CHF limit for very short periods (i.e., fast transients) is independent of the surface finish, being the same for a rough
467 chromium surface or a nano-smooth surface.

468 Future research will focus on consolidating the results of the transient CHF tests under exponential power escalations, as well
469 as anticipated operational occurrences (e.g., a power pulse during a BWR turbine trip), post-LOCA reflooding scenarios, and post-
470 CHF heat transfer at ambient pressure and PWR conditions.

471 Acknowledgments

472 This material is based upon work supported by U.S. Department of Energy Office of Nuclear Energy's Nuclear Energy University
473 Program (NEUP).

474 CRediT authorship contribution statement

475 **G. Su:** Conceptualization, Methodology, Investigation, Visualization, Formal analysis, Writing - Original Draft. **T. A. Moreira:**
476 Conceptualization, Methodology, Investigation, Visualization, Formal analysis, Writing - Review & Editing. **D. Lee:**
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