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*Inverse-Designed Lithium Niobate Nanophotonics*

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1 **Title**

2 Inverse-designed lithium niobate nanophotonics

3 **Authors**

4 Chengfei Shang,<sup>1†</sup> Jingwei Yang,<sup>1†</sup> Alec M. Hammond,<sup>2</sup> Zhaoxi Chen,<sup>1</sup> Mo Chen,<sup>3</sup> Zin Lin,<sup>4</sup> Steven  
5 G. Johnson,<sup>3\*</sup> and Cheng Wang<sup>1\*</sup>

6 **Affiliations**

7 <sup>1</sup>Department of Electrical Engineering & State Key Laboratory of Terahertz and Millimeter Waves,  
8 City University of Hong Kong, Kowloon, Hong Kong, China

9 <sup>2</sup>School of Electrical and Computer Engineering, Georgia Institute of Technology, Atlanta, GA 30308,  
10 USA

11 <sup>3</sup>Department of Mathematics, Massachusetts Institute of Technology, Cambridge, MA 02139, USA

12 <sup>4</sup>Department of Electrical and Computer Engineering, Virginia Polytechnic Institute and State  
13 University, Arlington, VA 22203, USA

14 <sup>†</sup>*These authors contributed equally to this article*

15 <sup>\*</sup>[stevenj@math.mit.edu](mailto:stevenj@math.mit.edu)

16 <sup>\*</sup>[cwang257@cityu.edu.hk](mailto:cwang257@cityu.edu.hk)

17  
18  
19 **Abstract**

20 Lithium niobate-on-insulator (LNOI) is an emerging photonic platform that exhibits favorable material  
21 properties (such as low optical loss, strong nonlinearities, and stability) and enables large-scale  
22 integration with stronger optical confinement, showing promise for future optical networks, quantum  
23 processors, and nonlinear optical systems. However, while photonics engineering has entered the era  
24 of automated “inverse design” via optimization in recent years, the design of LNOI integrated photonic  
25 devices still mostly relies on intuitive models and inefficient parameter sweeps, limiting the accessible  
26 parameter space, performance, and functionality. Here, we implement a 3D gradient-based inverse-  
27 design model tailored for topology optimization based on the LNOI platform, which not only could  
28 efficiently search a large parameter space but also takes into account practical fabrication constraints,  
29 including minimum feature sizes and etched sidewall angles. We experimentally demonstrate a spatial-  
30 mode multiplexer, a waveguide crossing, and a compact waveguide bend, all with low insertion losses,  
31 tiny footprints, and excellent agreement between simulation and experimental results. The devices,  
32 together with the design methodology, represent a crucial step towards the variety of advanced device  
33 functionalities needed in future LNOI photonics, and could provide compact and cost-effective  
34 solutions for future optical links, quantum technologies and nonlinear optics.

35  
36 **Keywords**

37 Inverse design, topology optimization, lithium niobate, integrate photonics

38

## 39 Introduction

40 The lithium niobate-on-insulator (LNOI) platform has seen rapid development<sup>(1)</sup> in recent years and  
41 shows great potential in future advanced photonics systems, owing to the excellent optical properties  
42 of lithium niobate (LN), including large nonlinear susceptibilities, a wide optical transparency window,  
43 and great stability, as well as strong optical confinement that allows compact and scalable integrated  
44 photonic devices and circuits to be built on wafer scales. A wide range of on-chip nanophotonic devices,  
45 including high-speed electro-optic modulators<sup>(2-4)</sup>, high-Q micro resonators<sup>(5, 6)</sup>, broadband frequency  
46 comb generators<sup>(7-9)</sup>, efficient frequency convertors<sup>(10-13)</sup>, entangled photon pair generators<sup>(14)</sup>,  
47 spectrometer<sup>(15)</sup>, optical isolator<sup>(16)</sup>, ultrafast all-optical switches<sup>(17)</sup> and frequency shifters<sup>(18)</sup> have  
48 been demonstrated, making LN-based photonic integrated circuits a promising solution for future high-  
49 speed optical networks, quantum information processing, and nonlinear optics. However, to date the  
50 design of most LNOI devices still relies heavily on human intuition augmented by simple analytical  
51 models, and can only access limited parameter spaces by scanning a few hand-selected parameters  
52 (such as widths, gaps and radii of curvature), either manually or through inefficient parameter sweeps  
53 (whose cost scales exponentially with the number of parameters).

54 Inverse-design methods, also as known as topology optimization<sup>(19, 20)</sup>, on the other hand, have been  
55 developed and attracted considerable attention in nanophotonics in the last two decades. Such methods  
56 could automatically search for the optimal topological structure of a pre-specified objective in a certain  
57 design region, exploring a large parameter space through meticulously developed optimization  
58 algorithms. Inverse design can reveal highly non-intuitive device designs with extremely compact sizes  
59 (several micrometers in diameter) and unprecedented functionalities. Many inverse-designed  
60 functional devices have been demonstrated in recent years, such as photonic crystal structures<sup>(21-23)</sup>,  
61 mode multiplexers and convertors<sup>(24, 25)</sup>, wavelength multiplexers<sup>(26, 27)</sup>, meta-surfaces<sup>(28, 29)</sup>, nonlinear  
62 wavelength convertors<sup>(29-31)</sup>, dispersion-engineered microresonators<sup>(32, 33)</sup>, together with system  
63 applications such as massively parallel optical transmitters<sup>(34)</sup>, particle accelerators<sup>(35)</sup>, and chip-based  
64 light detection and ranging (LiDAR) systems<sup>(36, 37)</sup>.

65 Thus far, most inverse-designed photonic devices have been demonstrated in silicon (Si) and other  
66 CMOS-compatible platforms such as silicon nitride<sup>(28)</sup> because of the mature fabrication technologies  
67 available for those materials. Compared with Si photonics, the challenges in the design and  
68 implementation of inverse-designed LNOI devices mainly arise from practical fabrication constraints.  
69 Due to the difficulty in dry etching LN and electro-optic overlap considerations, typical waveguides  
70 in LNOI feature a rib structure with an unetched slab underneath, further reducing the already smaller  
71 effective index contrast compared with that of Si photonics. Moreover, dry-etched LNOI waveguides  
72 usually exhibit a substantial sidewall angle, which needs to be taken into consideration during the  
73 optimization process to achieve accurate modeling, and which also limits the minimally achievable  
74 feature sizes. More recently, inverse-design algorithms that take into account specific geometric  
75 constraints<sup>(38)</sup> and sidewall angles<sup>(39)</sup> have led to designs and demonstrations compatible with standard  
76 foundry services<sup>(40)</sup> and in more exotic material platforms such as diamond<sup>(41)</sup> and silicon carbide<sup>(33)</sup>.  
77 However, the realization of inverse-designed LNOI devices is still missing and could substantially  
78 benefit the development of compact LNOI photonic integrated circuits, especially when dealing with  
79 complex design problems with multiple objective figures of merit.

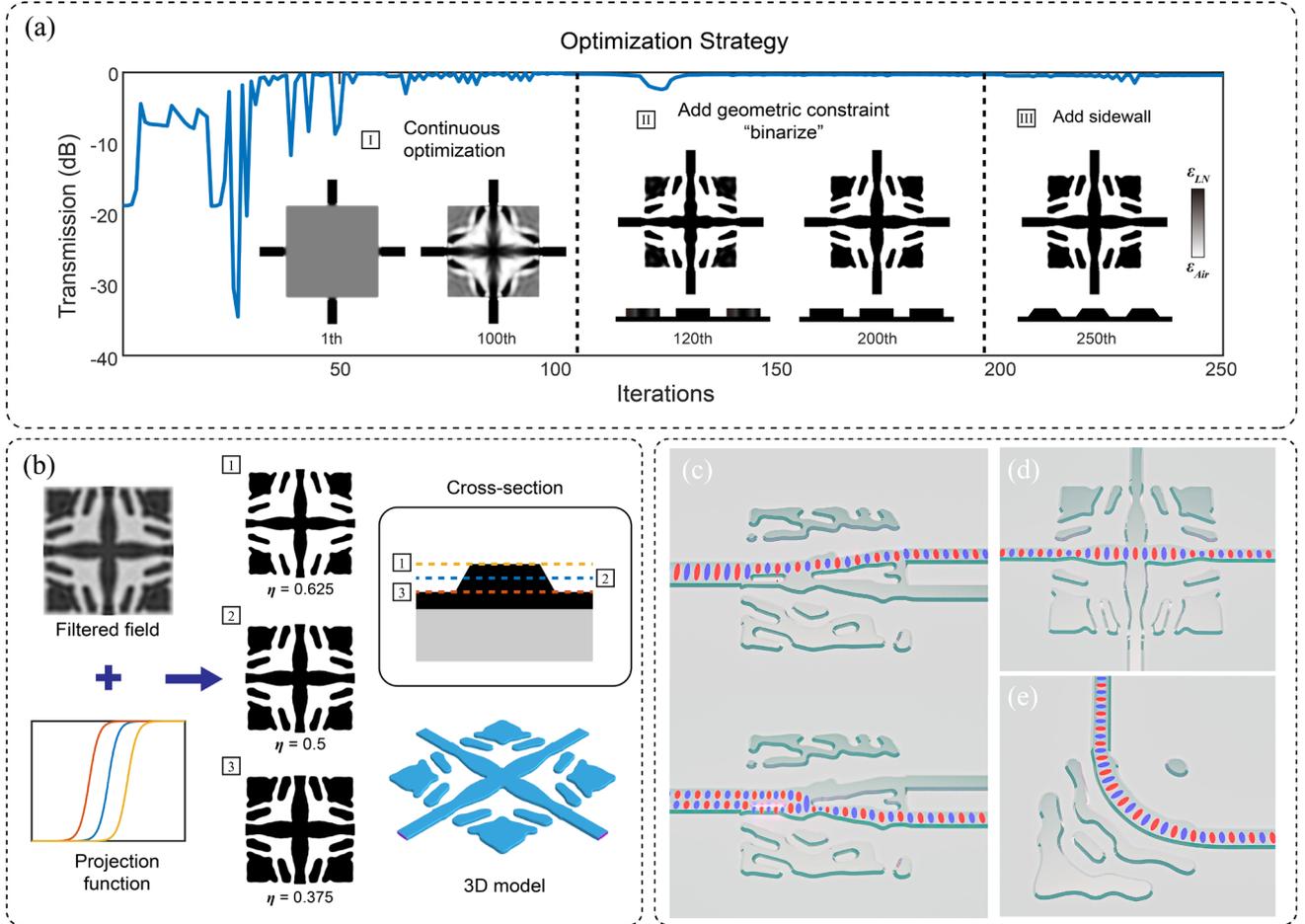
80 Here, we overcome these fabrication and design challenges and demonstrate a series of compact high-  
81 performance inverse-designed LNOI photonic devices based on the open-source Meep package<sup>(42)</sup>.

82 The design algorithm takes into full consideration of the fabrication constraints in the LNOI platform,  
83 including the rib structures, minimum feature sizes and etched sidewalls. We design and fabricate a  
84 spatial-mode multiplexer that separates the fundamental transverse-electric ( $TE_0$ ) mode and second-  
85 order TE ( $TE_1$ ) mode, with an insertion loss  $\sim 1.5$  dB and a crosstalk  $< -15.8$  dB, a waveguide crossing  
86 with a low loss of 0.48 dB and a crosstalk  $< -36$  dB, and a compact waveguide bend that turns the  
87 propagation of light by  $90^\circ$  with a radius of curvature of  $6 \mu\text{m}$  and a loss of 0.41 dB. The devices show  
88 excellent agreement between theoretical and experimental performances, as well as broad operation  
89 bandwidths from 1500 nm to 1600 nm wavelength.

## 90 Results

91 **Design framework.** Figure 1(a) shows an overview of our inverse-design strategy specially  
92 tailored for the LNOI platform and a representative iteration curve. The inverse design relies on a  
93 hybrid time/frequency-domain topology-optimization algorithm<sup>(43)</sup> with adjoint sensitivity analysis  
94 and multiple constraint functions. Starting with a homogeneous design region (typically initialized to  
95 “gray”—halfway between air and LNOI) and a given objective function, e.g., the transmission  
96 coefficients of a particular waveguide mode, the optimization solver efficiently calculates the objective  
97 performance and the gradients in each iteration step by two FDTD (finite-difference time domain)  
98 simulations regardless of the number of pixelated design parameters. Compared with optimization  
99 problems in SOI, the slanted sidewalls in the LNOI platform require special optimization algorithm  
100 design, whereas the existence of an unetched slab substantially increases the difficulty in achieving  
101 good optimization performances due to a lowered effective-index contrast. In this paper we use a 400  
102 nm z-cut LN device layer with a 250 nm etch depth and a sidewall angle of 45 degree for all designs  
103 (see device fabrication in supplementary information (S.I.)), similar to other devices previously  
104 demonstrated in our group<sup>(4, 44)</sup>.

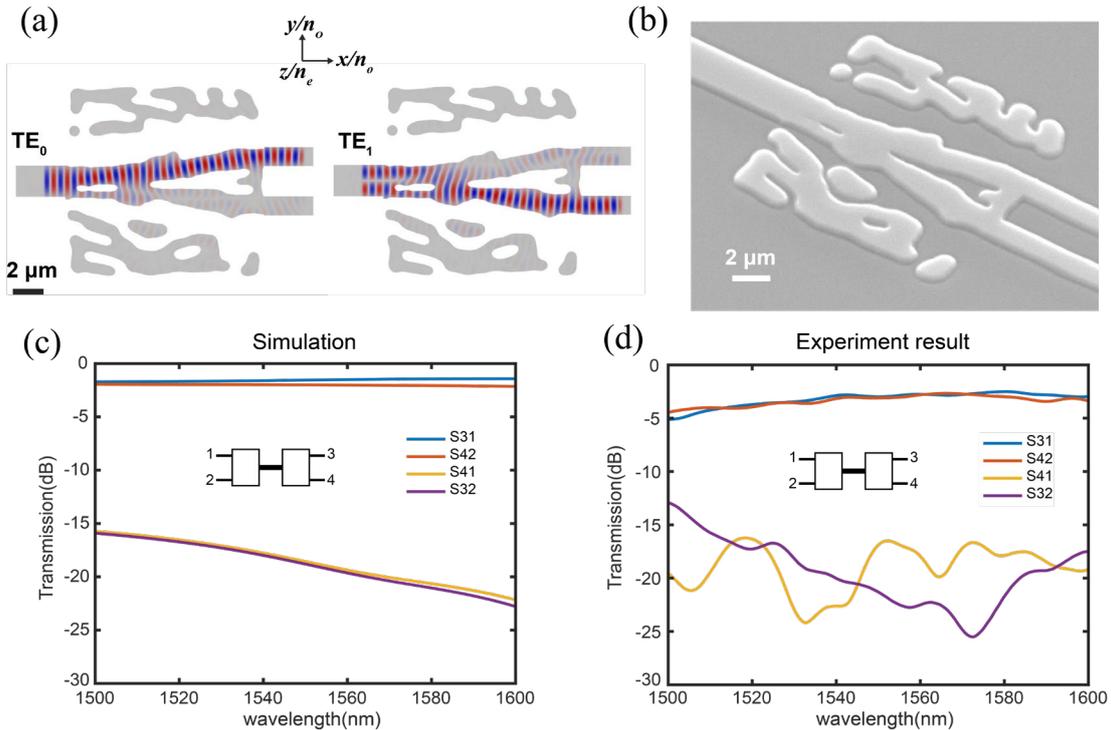
105 The optimization strategy shown in Fig. 1(a) contains three main steps as outlined in Ref. <sup>(43)</sup>. In the  
106 first step, permittivity in the design region is allowed to vary continuously, while a convolution and  
107 projection function is introduced to smoothen the design region and improve the dynamic range. The  
108 projection strength  $\beta$  is increased gradually during the iterations to avoid a dramatic change in the  
109 design parameters, such that the optimizer focuses on improving the figure of merit while gradually  
110 pushing the design towards a more binarized permittivity distribution. In the second step, we  
111 incorporate a geometric constraint<sup>(38)</sup> to eliminate features smaller than the indicated minimum length  
112 scales, during which the design parameters are further binarized. Finally, we introduce sidewall  
113 features by using linearly shifted threshold values for the projection function at different vertical  
114 slices<sup>(39)</sup>, so that the pattern is “eroded” as a function of the height above the slab surface (Fig. 1(b)).  
115 In our simulations, we use a threshold value interval from 0.375 (bottom) to 0.625 (top), within which  
116 the length scale roughly changes linearly with the threshold according to previous studies<sup>(45, 46)</sup> and our  
117 experience. In each step, the optimizer runs for several tens of iterations before convergence, the  
118 precise number depending on the specific problem. Importantly, all simulations are performed in 3D,  
119 taking fully into account the practical rib/slab thicknesses, crystal anisotropy, and sidewall angles,  
120 which are crucial to realize experimentally achievable designs. More mathematical details regarding  
121 the inverse design algorithms can be found in the device design section and Fig. S1 in S.I., as well as  
122 in previous references<sup>(38, 39, 47)</sup>. Figure 1(c-e) shows schematic views of the designed structures and  
123 their corresponding functions of our mode multiplexer, waveguide crossing and compact waveguide  
124 bend, respectively.



125  
126  
127 **Figure 1.** Overview of the optimization strategy and device schematic. (a) General optimization strategy and a typical  
128 iteration curve of a waveguide crossing optimization problem consisting of three main steps. In the first step, the  
129 permittivity in each pixel is allowed to vary continuously; second, we add geometric constraints to binarize the structure  
130 and achieve a minimum length scale; finally, slanted sidewall is introduced while the solver reoptimizes the design to  
131 “compensate” for the sidewall-induced performance degradation. The cross-section views in step 2 and step 3 are drawn  
132 along the diagonal of the design region. (b) Method of adding sidewall features. We introduce linearly shifted threshold  
133 values  $\eta$  for the hyperbolic tangent projection function at different heights (left), which leads to eroded ( $\eta = 0.625$ ),  
134 normal ( $\eta = 0.5$ ) and dilated ( $\eta = 0.375$ ) structures at the top, center and bottom slices of the rib structure, respectively  
135 (right). We design and fabricate (c) a  $TE_0/TE_1$  mode multiplexer, (d) a waveguide crossing, and (e) a compact waveguide  
136 bend to verify the design strategy, all of which show good performance and excellent fabrication compatibility.

137  
138 **Spatial mode multiplexer.** First, we consider an LNOI spatial mode multiplexer, an important  
139 component for mode-division multiplexing (MDM) technology in future high-volume data  
140 transmission systems<sup>(48)</sup>. In our design, the fundamental  $TE_0$  mode in a  $2\ \mu\text{m}$  wide input waveguide  
141 will be coupled into the upper output arm, whereas the  $TE_1$  mode will be converted into  $TE_0$  mode,  
142 and output from the bottom arm (Fig. 2(a)). The two output waveguides are both  $1\ \mu\text{m}$  wide and are  
143 separated by a  $2\ \mu\text{m}$  gap. Figure 2(a) shows the final inverse-designed pattern together with simulated  
144 field evolution ( $E_y$ ) for  $TE_0$  and  $TE_1$  input, whereas Figure 2(b) shows the scanning electron  
145 microscope (SEM) image of the fabricated mode multiplexer. The footprint of the final design is  $12 \times$   
146  $12\ \mu\text{m}^2$ , orders of magnitude smaller compared with mode multiplexers based on traditional  
147 asymmetrical directional couplers (ADCs)<sup>(49)</sup> or cascaded Mach–Zehnder interferometers (MZI)<sup>(50)</sup>.  
148 The minimum feature size is set to be  $0.2\ \mu\text{m}$  (at the middle slice of the rib) to balance between the  
149 potential degrees of freedom in design and the fabrication constraints. Our simulation results show that,

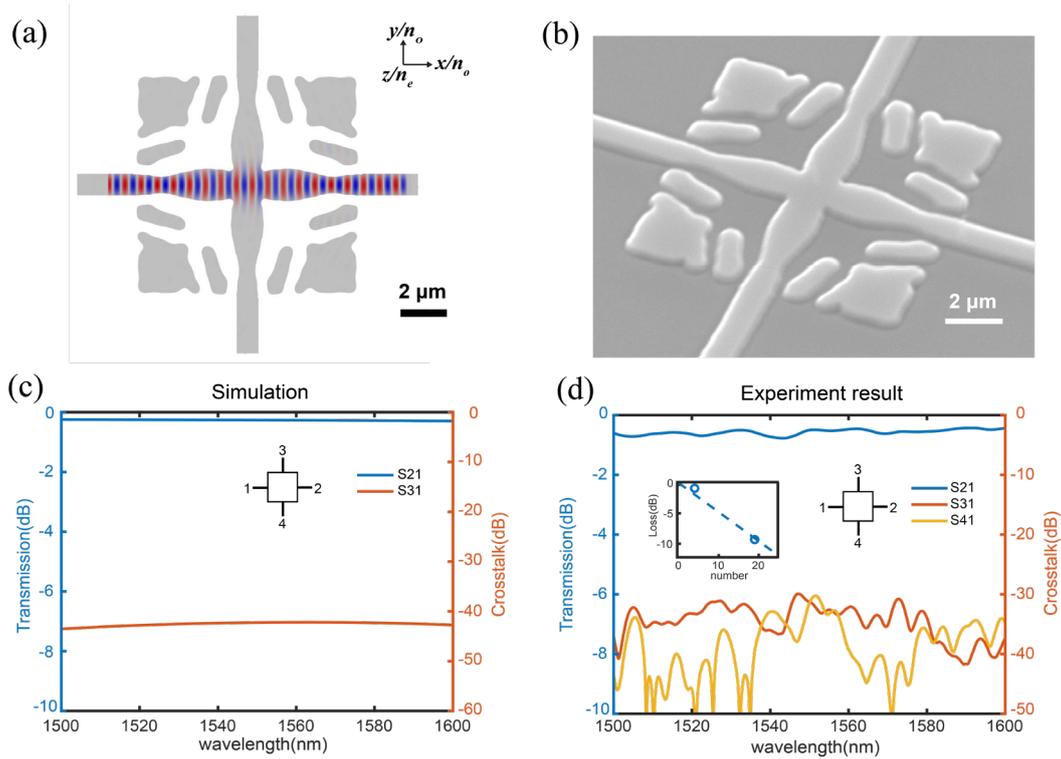
150 for both  $TE_0$  and  $TE_1$  input scenarios, a single mode multiplexer features an average insertion loss less  
 151 than 0.9 dB and crosstalk less than -17 dB over a 100-nm wavelength window (1500 nm-1600 nm). In  
 152 our actual experiments, we fabricate two multiplexers back to back connected by a 50- $\mu\text{m}$ -long, 2- $\mu\text{m}$ -  
 153 wide multimode waveguide, such that all input/output signals are in  $TE_0$  mode and are easier to be  
 154 precisely calibrated. The simulation result (Fig. 2(c)) for such a cascaded mode multiplexer pair shows  
 155 average insertion losses at the desired outputs ( $S_{31}$  and  $S_{42}$ ) of less than 2 dB (twice the loss of a single  
 156 device), and crosstalk ( $S_{41}$  and  $S_{32}$ ) of less than -16.5 dB. In the experimental test (Fig. 2(d)), we  
 157 measure average insertion losses of less than 3 dB and crosstalk less than -15.8 dB for the mode-  
 158 multiplexer pair within a wavelength range between 1520 nm and 1600 nm, consistent with the  
 159 simulation results. The measured insertion loss values are further corroborated by comparing the losses  
 160 of two and four cascaded mode-multiplexers (see performance estimation in S.I.). From the results we  
 161 estimate that a single mode multiplexer features an insertion loss  $\sim 1.5$  dB and an upper-bound  
 162 crosstalk value of -15.8 dB within the 80-nm wavelength window (see performance estimation for a  
 163 detailed analysis). The remaining differences between simulation and experimental values could result  
 164 from excessive scattering losses due to deviations in the actually fabricated device parameters and  
 165 fabrication imperfections especially on the small features (see origin of excessive insertion loss and  
 166 Fig. S2 in S.I. for more details). We also note that the peripheral structures in our design (e.g. the top  
 167 and bottom bars in Fig.2(a)) do not substantially contribute to the overall field evolution (Fig. S3 in  
 168 S.I.), potentially allowing the achievement of similar device performances within even smaller  
 169 footprints by adopting customized shapes of design region.



170  
 171 **Figure 2.** Inverse-designed spatial mode multiplexer. (a) The optimized design pattern and simulated field ( $E_y$ ) evolution  
 172 of the mode multiplexer with a  $12 \times 12 \mu\text{m}^2$  footprint. The gray and white areas correspond to LN rib and slab regions,  
 173 respectively. (b) SEM image of the fabricated mode multiplexer. (c) Simulated transmission coefficients of a back-to-  
 174 back mode multiplexer pair between the four input/output ports as shown in the inset. (d) Experimentally measured  
 175 transmission coefficients of a fabricated mode multiplexer pair, showing broadband low-loss and low-crosstalk operation  
 176 consistent with simulation results.

177  
 178

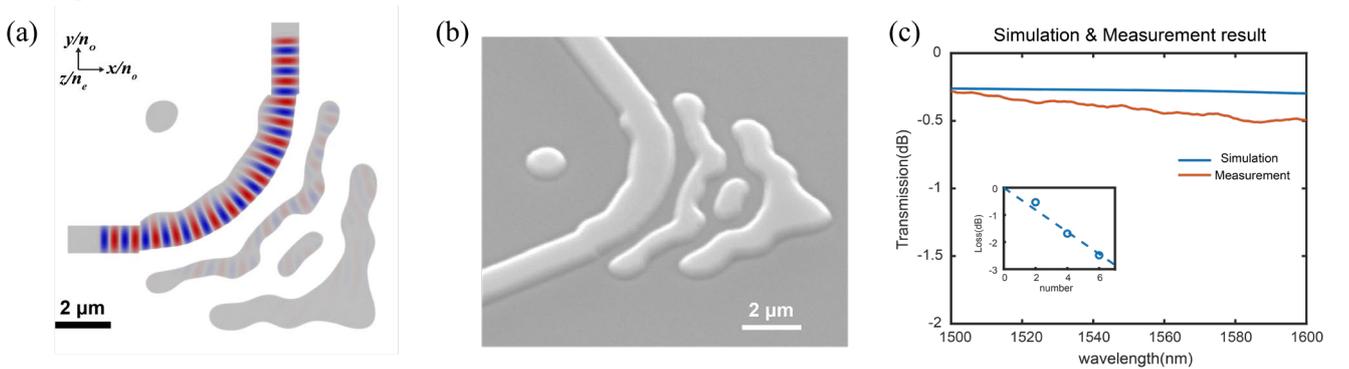
179 **Waveguide crossing.** A waveguide crossing<sup>(51-56)</sup> is an essential component for signal routing in  
 180 large-scale and high-density photonic integrated circuits. Traditional waveguide crossing designs  
 181 typically rely on heuristic shaped taper or multimode interferometer (MMI) structures combined with  
 182 exhaustive parameter sweeps<sup>(54, 55)</sup>. Here we design a compact waveguide crossing using our inverse-  
 183 design algorithm without the need for an initial guess. The design region (Fig. 3(a)) again has a  
 184 footprint of  $12 \times 12 \mu\text{m}^2$ , and is set to have mirror symmetry along both horizontal and vertical  
 185 directions, since the  $x$ - and  $y$ -crystal orientations are isotropic in our  $z$ -cut LNOI wafer. The objective  
 186 function is designed to maximize the output power in fundamental  $\text{TE}_0$  mode. Here the minimum  
 187 feature size is set to be  $0.5 \mu\text{m}$  as this optimization problem is easier to converge. Figure 3(b) shows  
 188 the SEM image of the fabricated waveguide crossing. The simulated average insertion loss over a 100-  
 189 nm wavelength band (Fig. 3(c)) is 0.22 dB, with low crosstalk of less than -40 dB. In the experiment,  
 190 we measure the insertion loss by cascading five and twenty crossing structures and comparing with a  
 191 single waveguide crossing (as indicated in the inset of Fig. 3(d)), showing a low fitted average insertion  
 192 loss of 0.48 dB per crossing over the tested wavelength range (1500-1600 nm). The crosstalk values  
 193 measured from the two vertical output waveguides (port 3 and port 4, Fig. 3(d)) are -36 dB and -39 dB,  
 194 respectively, consistent with the simulation results. Similar to the top and bottom features in the mode  
 195 multiplexer, we find that the corners of the design region also do not influence much on the device  
 196 function (see Fig. S3(c-d) in S.I).



197  
 198 **Figure 3.** Inverse-designed LNOI waveguide crossing. (a) The optimized design pattern and simulated field ( $E_y$ )  
 199 evolution of the waveguide crossing with a  $12 \times 12 \mu\text{m}^2$  footprint. The gray and white areas correspond to LN rib and  
 200 slab regions, respectively. (b) SEM image of the fabricated crossing device. (c) Simulated optical transmission ( $S_{21}$ ) and  
 201 crosstalk ( $S_{31}$ ) of the designed waveguide crossing. (d) Experimentally measured device performance, showing a low  
 202 insertion loss and crosstalk from 1500 nm to 1600 nm. Inset shows cut-back loss measurement results of 5 and 20  
 203 cascaded crossing structures.

204  
 205 **Compact waveguide bend.** Finally, we demonstrate a compact bending waveguide that rotates  
 206 the propagation direction of the fundamental  $\text{TE}_0$  mode by  $90^\circ$  within a tight bending radius of  $6 \mu\text{m}$ .

207 Due to the existence of unetched slab, waveguide bends in LNOI platforms usually require radii of at  
 208 least  $30\ \mu\text{m}$  to limit the radiation loss. A simple circular  $90^\circ$  bend with a radius of  $6\ \mu\text{m}$  in our current  
 209 platform will lead to a high loss of  $3.7\ \text{dB}$  according to our numerical simulation. Here the design  
 210 region (Fig. 4(a)) is  $10 \times 10\ \mu\text{m}^2$  which is nearly an order of magnitude smaller than a traditional  $30\text{-}\mu\text{m}$   
 211  $\mu\text{m}$  bend, both widths of input and output waveguides are  $1\ \mu\text{m}$ , and the minimum length scale is  $0.4\ \mu\text{m}$ .  
 212 Figure 4(b) shows the SEM image of the fabricated waveguide bend. Physically, this optimized  
 213 bend design could be roughly interpreted as a sharp bend augmented with a Bragg-mirror-like structure  
 214 to suppress radiation loss, qualitatively similar to bends designed by topology optimization in other  
 215 material platforms<sup>(57)</sup>, but determining the precise details requires the power of inverse design. The  
 216 simulated average transmission loss (Fig. 4(c), blue) in the desired band ( $1500\ \text{nm}\sim 1600\ \text{nm}$ ) is  $0.29\ \text{dB}$ ,  
 217 which is equivalent to a simple bend with a radius of  $30\ \mu\text{m}$  according to our simulation. In the  
 218 experimental test, we cascaded two, four, and six bends and compare the measured total insertion  
 219 losses with a reference waveguide (Fig. 4(c), red and inset). The average measured insertion loss for a  
 220 single bend is  $0.41\ \text{dB}$  over the  $100\text{-nm}$  wavelength range. The measured loss is slightly higher than  
 221 expected, possibly due to an under-etched rib height ( $\sim 230\ \text{nm}$ ) resulting in more power leakage  
 222 through the thicker slab. Nevertheless, the measured loss is still more than 8 times lower than that of  
 223 a simple circular bend with the same radius.



224  
 225 **Fig. 4.** Inverse-designed LNOI waveguide bend. (a) The optimized design pattern and simulated field ( $H_z$ ) evolution of  
 226 the bending waveguide with a  $6\text{-}\mu\text{m}$  radius. The gray and white areas correspond to LN rib and slab regions, respectively.  
 227 (b) SEM image of the fabricated waveguide bend. (c) Simulated and experimentally measured optical transmission of the  
 228 waveguide bend, showing a low average insertion loss from  $1500\ \text{nm}$  to  $1600\ \text{nm}$ . Inset shows cut-back loss  
 229 measurement results of two, four and six cascaded bends.

## 230 Conclusions and Discussions

231 In summary, we have successfully implemented an inverse-design algorithm which is compatible with  
 232 the practical fabrication constraints in the LNOI platform, including non-vertical sidewalls, relatively  
 233 large minimum length scales and an unetched-slab structure. Based on this algorithm, we demonstrate,  
 234 for the first time, a series of inverse-designed LNOI devices with small footprints, low losses, low  
 235 crosstalk, broad bandwidths, and good consistency with theoretical prediction. Further relaxing the  
 236 minimum length constraints could lead to designs compatible with stepper photolithography  
 237 processes<sup>(58)</sup> with much better scalability and cost-effectiveness. The same design methodology could  
 238 be readily applied to achieve LNOI devices with more advanced functions, especially those make use  
 239 of the electro-optic and/or nonlinear properties of LN, such as fast-tunable switches<sup>(17)</sup>, dispersion-  
 240 engineered comb generators<sup>(8)</sup>, and efficient nonlinear wavelength converters<sup>(10)</sup>, and be extended to  
 241 other material platforms that share similar fabrication constraints, such as yttrium orthovanadate<sup>(59)</sup>  
 242 and titanium dioxide<sup>(60)</sup>. The devices, together with the design methods, could become important

243 building blocks for future LNOI functional photonic circuits with applications in on-chip optical links,  
244 quantum technologies and nonlinear optics.

245

## 246 **Supporting Information**

247 Device fabrication & characterization; Device design; Origin of excessive insertion loss in fabricated  
248 devices; Performance estimation of a single mode multiplexer; Simulation after removing edge  
249 features

250

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260

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