

# MIT Open Access Articles

# Synthesis and Characteristics of Transferrable Single#Crystalline AlN Nanomembranes

The MIT Faculty has made this article openly available. *Please share* how this access benefits you. Your story matters.

**Citation:** Honglae Sohn, Jong-Dae Lee, Intramolecularly Stabilized o-Carboranyl Aluminum Complexes: Synthesis, Characterization, and X-ray Structural Studies, Crystals, 10.3390/ cryst13060877, 13, 6, (877), (2023).

As Published: 10.1002/aelm.202201309

Publisher: Wiley

Persistent URL: https://hdl.handle.net/1721.1/153558

**Version:** Final published version: final published article, as it appeared in a journal, conference proceedings, or other formally published context

Terms of use: Creative Commons Attribution





# Synthesis and Characteristics of Transferrable Single-Crystalline AIN Nanomembranes

Jiarui Gong, Jie Zhou, Ping Wang, Tae-Hyeon Kim, Kuangye Lu, Seunghwan Min, Ranveer Singh, Moheb Sheikhi, Haris Naeem Abbasi, Daniel Vincent, Ding Wang, Neil Campbell, Timothy Grotjohn, Mark Rzchowski, Jeehwan Kim, Edward T. Yu,\* Zetian Mi,\* and Zhenqiang Ma\*

Single-crystalline inorganic semiconductor nanomembranes (NMs) have attracted great attention over the last decade, which poses great advantages to complex device integration. Applications in heterogeneous electronics and flexible electronics have been demonstrated with various semiconductor nanomembranes. Single-crystalline aluminum nitride (AIN), as an ultrawidebandgap semiconductor with great potential in applications such as highpower electronics has not been demonstrated in its NM forms. This very first report demonstrates the creation, transfer-printing, and characteristics of the high-quality single-crystalline AIN NMs. This work successfully transfers the AIN NMs onto various foreign substrates. The crystalline quality of the NMs has been characterized by a broad range of techniques before and after the transfer-printing and no degradation in crystal quality has been observed. Interestingly, a partial relaxation of the tensile stress has been observed when comparing the original as-grown AIN epi and the transferred AIN NMs. In addition, the transferred AIN NMs exhibits the presence of piezoelectricity at the nanoscale, as confirmed by piezoelectric force microscopy. This work also comments on the advantages and the challenges of the approach. Potentially, the novel approach opens a viable path for the development of the AIN-based heterogeneous integration and future novel electronics and optoelectronics.

# 1. Introduction

Single-crystalline inorganic semiconductor nanomembranes (NMs), which have a thickness of less than a few hundred nanometers, have gained increasing interest since 2005 due to their remarkable use for a variety of applications. These free-standing NMs can be transferred to any arbitrary foreign substrates by using elastomeric stamps, which was first demonstrated by Meitl et al.<sup>[1]</sup> Importantly, these NMs are compatible with the traditional semiconductor manufacture approach for fabricating various heterogeneously integrated and flexible electronic devices,<sup>[2-4]</sup> including the seamless integration of semiconductor NMs with the complementary-metal-oxidelow-cost, semiconductor-compatible silicon (Si) substrates. Until now, several applications of transferrable monocrystalline NMs based on Si,<sup>[4-20]</sup> germanium (Ge),<sup>[4,21]</sup> gallium arsenide (GaAs)/aluminum gallium

J. Gong, N. Campbell, M. Rzchowski Department of Physics University of Wisconsin-Madison Madison, WI 53706, USA J. Gong, J. Zhou, S. Min, R. Singh, M. Sheikhi, H. N. Abbasi, D. Vincent, Z. Ma Department of Electrical and Computer Engineering University of Wisconsin-Madison Madison, WI 53706, USA E-mail: mazq@engr.wisc.edu

The ORCID identification number(s) for the author(s) of this article can be found under https://doi.org/10.1002/aelm.202201309.

© 2023 The Authors. Advanced Electronic Materials published by Wiley-VCH GmbH. This is an open access article under the terms of the Creative Commons Attribution License, which permits use, distribution and reproduction in any medium, provided the original work is properly cited.

#### DOI: 10.1002/aelm.202201309

P. Wang, D. Wang, Z. Mi Department of Electrical Engineering and Computer Science University of Michigan Ann Arbor, MI 48109, USA E-mail: ztmi@umich.edu T.-H. Kim, E. T. Yu Microelectronics Research Center Department of Electrical Engineering and Computer Science University of Texas at Austin Austin, TX 78758, USA E-mail: ety@ece.utexas.edu K Lu L Kim Department of Mechanical Engineering Department of Materials Science and Engineering Massachusetts Institute of Technology Cambridge, MA 02139, USA T. Grotiohn Department of Electrical and Computer Engineering Michigan State University East Lansing, MI 48824, USA

(1)

ADVANCED SCIENCE NEWS www.advancedsciencenews.com

arsenide (AlGaAs),<sup>[4,22–26]</sup> gallium nitride (GaN)/aluminum gallium nitride (AlGaN),<sup>[27–30]</sup> silicon carbide (SiC),<sup>[31]</sup> gallium oxide (Ga<sub>2</sub>O<sub>3</sub>),<sup>[32–35]</sup> etc. have been reported during the last few years. These transferrable narrow, wide, and ultrawide bandgap NMs have enabled a variety of device applications over the past nearly two decades.

In comparison to the above semiconductors that have been released in the form of NMs, single-crystalline aluminum nitride (AlN) has the unique properties of ultrawide direct bandgap (6.2 eV),<sup>[36]</sup> high breakdown electric field (15.4 MV cm<sup>-1</sup>),<sup>[37]</sup> good thermal conductivity (3.19 W cm<sup>-1</sup> K<sup>-1</sup>),<sup>[38]</sup> strong spontaneous polarization, and good piezoelectricity. These properties enable AlN for applications in ultraviolet wavelength sensing,<sup>[39,40]</sup> power electronics,<sup>[41,42]</sup> piezoelectric devices,<sup>[43-48]</sup> etc. Furthermore, the polarization properties of single-crystalline AlN can induce free charges at the heterointerface and influence mobile carrier distributions via polarization-induced bound charges, an approach which is expected to be much more effective than doping via impurities.<sup>[49]</sup> Previously, several attempts have been made to grow and release the high-quality free-standing AlN NMs. For instance, Mei et al. used metal-organic vapor phase epitaxy to grow porous crystalline AlN/GaN thin films over the Si (111) substrate and further, released it by selective Si wet etching.<sup>[50]</sup> In another report, Yang et al. sputtered polycrystalline AlN on the silicon dioxide (SiO<sub>2</sub>) substrate by magnetron sputtering and released the NMs through selective SiO<sub>2</sub> wet etching.<sup>[51]</sup> Recently, AlN thin film was grown on Ge and Si substrates by radio frequency magnetron reactive sputtering and molecular beam epitaxy (MBE), respectively, which was released by selective gas-phase dry etching of Ge and Si.<sup>[52,53]</sup> However, these AlN materials were not transferred to a foreign semiconductor substrate, which greatly limits their potential applications. Thus, it is very important not only to fabricate the electronic devices but also from a fundamental point of view to carry out a detailed study on the transfer of free-standing AlN NMs.

A few studies on transferring AlN thin films to flexible adhesive substrates have been reported recently.<sup>[54–56]</sup> However, there exist three issues: 1) the AlN crystalline quality, e.g., grain size, crystal orientation, in the above-mentioned studies is rather low due to the growing methods, which are limited by the growth approaches; 2) the adhesive handling materials during the transfer can leave residues on the AlN film; 3) only adhesive flexible substrates have been demonstrated as the final host substrate, which eliminate the potential of applications such as AlN-based vertical devices.<sup>[4]</sup> Hence, it is desirable to overcome the above issues and achieve a high-quality transferrable singlecrystalline AlN NM that can be transferred to a broader range of substrates.

In this work, we demonstrate the synthesis and characterizations of the transferrable single-crystalline AlN NMs. We have successfully transferred the AlN NMs onto flat sapphire, flat Si, and pillar-structured Si substrates. The crystal features are characterized by X-ray diffraction (XRD), Raman spectroscopy, and scanning transmission electron microscope (STEM) techniques. Importantly, no degradation has been observed in the crystallinity of the AlN NMs before and after transfer on the different substrates, highlighting the quality of the transferred NMs. Further, atomic force microscopy (AFM) is employed to monitor the surface roughness of as-grown Al-polar and ADVANCED ELECTRONIC MATERIALS

www.advelectronicmat.de

posttransferred N-polar AlN NMs, confirming the ultrasmooth surface of the NMs. Piezoresponse force microscopy (PFM) confirms the presence of the piezoelectricity in the transferred AlN NMs. X-ray photoelectron spectroscopy (XPS) was employed to characterize the valence band maximum (VBM) on both Al-polar and N-polar surfaces, without and with various atomic-layer-deposited (ALD) ultrathin oxide (UO) coatings. Based on the above experimental results, the present study not only opens a new insight for the fundamental device fabrication but also provides a novel platform for a broad range of electronic, optoelectronic, and photonic device applications.

### 2. Results and Discussion

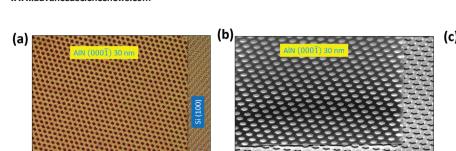
Initially, the AlN layer was grown using MBE and further, it was successfully transferred onto a flat sapphire (Figure S1, Supporting Information), a flat Si (Figures 2a,b and 3a), and a hexagon pillar-structured Si (Figure 1a-c) substrates. Figure 1a exhibits the optical image of the transferred AlN NMs on the pillar-structured Si substrate. The pillar-structured Si substrate was prepared via photolithography and reactive ion etching (RIE) and the height of the pillars is about 4  $\mu$ m. The pillars can be seen through the AlN NM due to the high optical transparency of AlN. Figure 1b,c presents the tilt-view scanning electron microscope (SEM) images of transferred AlN NMs on the pillar-structured Si substrate. These SEM images clearly confirm the uniformity of the NMs over the whole substrate without formation of cracks. The high-resolution SEM image (Figure 1c) confirms the smoothness of the NMs edges, which are highly desired from the device fabrication point of view. It is noted from Figure 1c that the AlN NM has sufficient rigidity even with a thickness of only 30 nm (Figure 2a). The schematics of as-grown AlN and transferred AlN on flat Si substrate have been illustrated in Figure 1d.

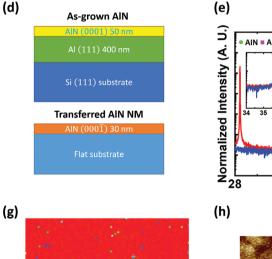
A desirable feature to use NMs for device fabrication is to have great crystalline quality and it is well-known that the crystallinity of a layer strongly depends on the growth process. Therefore, XRD measurements are performed to study the crystalline guality of the AlN NMs before and after the transferprinting process. 2theta-omega spectra of the epi wafer and AlN NM transferred on the sapphire (0001) substrate were characterized. Figure 1e shows the XRD pattern of the AlN epi sample and the transferred AlN on sapphire sample. Of note, the spectrum of the AlN epi sample was aligned by the Si (111) peak while the one of the transferred AlN NM sample was aligned by the Al<sub>2</sub>O<sub>3</sub> (0006) peak during the XRD measurement. From Figure 1e, it is evident that both the samples are high crystalline in nature and the AlN (0002) peak are found to be located at 36.11° and 35.92°. The decrease in the  $2\theta$  (as is seen from the inset of Figure 1e) indicates the increase of the interplane spacing d in the corresponding direction, i.e., AlN (0002), according to the Bragg condition:

$$n\lambda = 2d\sin\theta$$

where *n* is the diffraction order and  $\lambda$  is the wavelength. In addition, the measured full width at half maximums (FWHM) of the AlN (0002) peaks from the AlN epi sample and the transferred







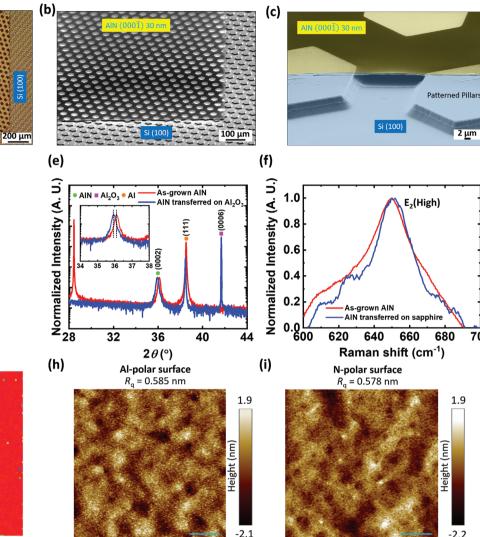


Figure 1. a) Optical image (a top view) and b,c) scanning electron microscope (SEM) images (tilted views) of single-crystalline aluminum nitride (AIN) nanomembrane (NM) transferred on a silicon (Si) (100) substrate that has an array of patterned hexagon shaped pillars. d) Schematics of the AIN (0001)/Al (111)/Si (111) epi wafer (top) and the AlN (0001) NM transferred onto the flat substrate (bottom). e) X-ray diffraction (XRD) 2theta-omega scans and f) Raman spectra of the as-grown AIN epi and the transferred AIN NM on sapphire (0001) substrate. XRD peaks of Si (111), AIN (0002), AI (111), and Al<sub>2</sub>O<sub>3</sub> (0006) planes are observed and labeled in (e). The inset of (e) shows the enlarged spectra around the AlN (0002) peaks. AlN E<sub>2</sub>(high) Raman peak is observed and labeled in (f). g) EBSD map of the transferred AIN NM on the flat Si substrate. h,i) The atomic force microscopy (AFM) images of h) Al-polar surface (as-grown) and i) N-polar surface (posttransferred AlN NM on the Si (100) substrate) of the single-crystalline AlN NM. The surface roughness is about the same for both surfaces.

1 μm

AlN on sapphire sample are 0.210° and 0.216°, respectively, indicating no significant degradation of the AlN crystalline quality throughout the transfer process. Further, the crystallite size can be calculated using the Scherrer equation<sup>[57,58]</sup>

1010

2110

0001

$$L = \frac{K\lambda}{\beta\cos\theta} \tag{2}$$

where L is the average size of the crystallite domains, K is the shape factor (which is taken as 1 for estimation),  $\lambda$  is the wavelength of the X-ray (1.5418 Å),  $\beta$  is the FWHM in radians, and  $\theta$  is the Bragg angle. The estimated crystallite size of the transferred AlN NMs is found to be ≈44 nm. The obtained value of crystallite size is better than the ones of the AlN thin films transferred on the flexible substrates.<sup>[55,56]</sup> It is further noticed that the grain size is mainly limited by the epi film thickness (50 nm). In Figure 1e, the peak observed at  $2\theta = 38.46^{\circ}$  in the XRD spectrum of the transferred AlN on sapphire sample is related to the Al residue after plasma etching (Figure S1, Supporting Information). As mentioned in the experiment section, the exposure of the AlN in the transparent region has been confirmed by XPS survey scan (Figure S2, Supporting Information).

To obtain more details on the crystalline structure, the electron backscatter diffraction (EBSD) map is collected on the

www.advelectronicmat.de

2 <u>µ</u>m

680

1.9

Height (nm

-2.2

1 µm

700

3 µm



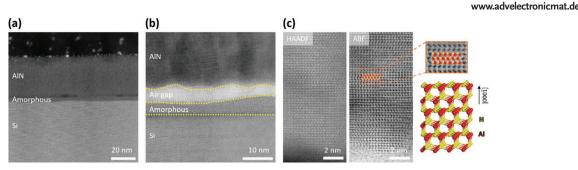


Figure 2. Microstructural analysis of the transferred aluminum nitride (AIN) nanomembrane (NM). a) Cross-sectional high-angle annular dark-field scanning transmission electron microscope (HAADF–STEM) image of the transferred AIN NM on the silicon (Si) (100) substrate. The AIN layer thickness is 30 nm. b) Annular bright-field STEM (ABF–STEM) image of the interface between the AIN NM and the Si (100) substrate. A local air gap was observed. c) High resolution HAADF–STEM image (left) and the corresponding ABF-STEM image (right) of the AIN membrane. Al atoms (the red balls) and N atoms (the yellow balls) are embedded in (c) to visualize the atomic stacking sequence, showing a clear N-polar lattice. Reconstructed atomic structure from the STEM images is shown in the bottom right of (c).

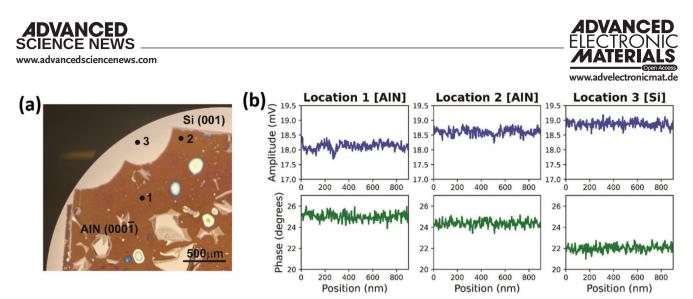
transferred AlN NM as shown in Figure 1g. The map shows (0001) in-plane orientation (red region), which matches the results from XRD spectra (Figure 1e). Some isolated points with other colors are likely due to the low signal-to-noise ratio for this 50 nm AlN NM during the EBSD measurement in SEM and mistakenly generated by the fitting software.

Further, for in-depth study about vibrational properties, the micro-Raman measurements are carried out on the as-grown and transferred AlN NMs to confirm the formation of AlN and measure the stress inside the NMs. It has been confirmed that no major peaks were located between 650 and 660 cm<sup>-1</sup> from the spectra of the Al and the sapphire substrates (Figure S3, Supporting Information). As shown in Figure 1f, the Raman shift appearing at 650 cm<sup>-1</sup> on the epi sample and 651 cm<sup>-1</sup> on the transferred AlN NMs on the sapphire substrate is associated to the  $E_2$  (high) Raman mode of AlN, which corresponds to a biaxial tensile stress of 1.17 and 1.02 GPa, respectively.<sup>[59]</sup> The existence of tensile stress aligns with the curling phenomenon observed during the transfer-printing process, as described in the Experiment Section. In addition, no significant change has been observed in the Raman spectra of AlN NMs before and after the transfer-printing process, confirming the high quality of the transferred AlN NMs. The measured FWHM of the  $E_2$ (high) peak is 34 cm<sup>-1</sup> for the AlN epi sample and 26 cm<sup>-1</sup> for the AlN transferred on sapphire sample. The decrease in the FWHM of the E<sub>2</sub>(high) peak after the Al removal indicates the partial relaxation of the stress gradient in the AlN NMs,<sup>[60]</sup> which also agrees well with the Raman  $E_2$ (high) peak shift and the XRD AlN (0002) peak shift.

It is well-known that a smooth interface is highly desirable to fabricate the high performing electronic devices. Therefore, the surface roughness of the AlN NMs has been measured from both sides by AFM. Figure 1h,i depicts the morphological images of AlN epi sample (Al-polar surface) and transferred AlN NMs (N-polar surface), respectively. It is interesting to note that the AlN NMs from both the polarities show a smooth surface having root mean square (RMS) of 0.585 and 0.578 nm. After transferring the AlN NMs (N-polar surface), the RMS is found to be comparable to that of Al-polar surface even with plasma etching process of Al layer.

To understand the microstructural and crystalline nature of the transferred AlN NMs, cross-sectional high-angle annular dark-field STEM (HAADF-STEM) technique was employed. The sample was an AlN NM transferred onto a Si substrate that was coated with amorphous Al<sub>2</sub>O<sub>3</sub> using atomic layer deposition (ALD). The ALD-Al<sub>2</sub>O<sub>3</sub> coating was to mimic the situation of the grafted heterostructures.<sup>[4]</sup> The sample was annealed at 350 °C for 5 min in N2 ambient via rapid thermal annealing (RTA) after the NM transfer-printing to achieve chemical bonding between the NM and Si substrate. Figure 2a shows the cross-sectional HAADF-STEM image of the transferred AlN membrane on the Si substrate, which not only confirms the uniformity of the AlN NMs but also reveals the thickness of the NMs (≈30 nm). Importantly, no other defects, damages, stacking faults, or threading dislocations were observed in the transferred AlN NMs, confirming the high quality of the NMs. And the top Al layer was removed completely during the excess dry etching process. Originally, the thickness of the AlN layer is 50 nm whereas after transferring it is found to be 30 nm. The thickness of AlN layer decreases because of the aggressive plasma etching process (as shown in Figure 5h).

Figure 2b shows an amorphous layer between the AlN membrane and Si substrate, as marked by yellow dotted lines in Figure 2b. Even though we had removed the native oxide of the Si substrate by HF solution before the ALD-Al<sub>2</sub>O<sub>3</sub> deposition, the sample was unavoidably exposed to air before it was sent into the ALD chamber, resulting in the formation of Si native oxide over the Si substrate.<sup>[20,61]</sup> Furthermore, the sample underwent the RTA process after the NM transfer, which may induce oxygen diffusion from the interface into the substrate and thus the formation of Si native oxide.<sup>[20]</sup> Therefore, the amorphous layer is expected to consist of both ALD-Al2O3 and Si native oxide (SiO<sub>x</sub>). In addition, we also observed an air gap locally in some regions between the substrate and the transferred AlN NM. It is speculated that the local air gaps were generated due to the rigidity/stiffness of the AlN NM,<sup>[62,63]</sup> which caused its nonintimate contacts with the substrate at various locations. The RTA process can exacerbate the air gaps due to the large mismatch of thermal expansion coefficients between Si and AlN.<sup>[64]</sup> Figure 2c presents the high-resolution HAADF-STEM



**Figure 3.** a) Photograph of aluminum nitride (AIN)/silicon (Si) sample and location for which piezoresponse force microscopy (PFM) measurements are shown. Darker regions correspond to the AIN (0001) thin film, and light regions to the Si (001) substrate. b) Amplitude (blue lines, top) and phase (green lines, bottom) of surface modulation measured during PFM measurement. Locations 1–3 are indicated in (a).

image and the corresponding annular bright-field STEM (ABF– STEM) image acquired from the transferred AlN membrane. From Figure 2c, it is found that the transferred AlN NMs is highly crystalline and has a highly ordered wurtzite lattice without in-plane rotation or cubic phase, agreeing well with the XRD results. It is well known that in the ABF–STEM image, the darkest spots correspond to the heaviest atoms.<sup>[65,66]</sup> Therefore, we are able to configure the atomic stacking sequence of Al and N atoms, which is embedded in the ABF–STEM image (Figure 2c). Comparing with the atomic structure of wurtzite AlN, we confirmed that the lattice-polarity of the transferred AlN membrane is N-polar. In other words, the as-grown AlN on Al/Si template has an Al-polar lattice. These results unambiguously confirm that the transferred AlN membrane has a singlecrystalline wurtzite phase with a N-polar lattice.

The intrinsic polarization in III-nitride semiconductors plays an important role to fabricate the high electron mobility or high hole mobility devices. Therefore, the piezoelectric response of the transferred AlN ( $000\overline{1}$ ) NMs on the flat Si (001) substrate, according to the conductivity requirement for the substrate, was confirmed using PFM.

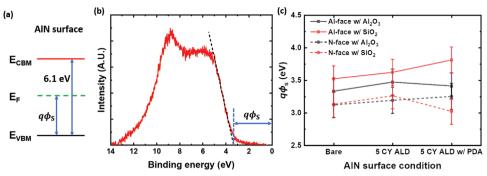
Figure 3a shows a digital photograph of an AlN  $(000\overline{1})/Si$ (001) sample which was used to characterize the piezoelectric response. As shown in Figure 3a, the dark brown region and bright region correspond to the AlN NM and Si substrate, respectively. The colorful features shown in the optical image of Figure 3a are explained in details in Figure S4, Supporting Information. For PFM measurement, the sample was annealed at 350 °C for 5 min in N2 ambient via RTA to ensure good contact between the AlN NM and the Si substrate. In Figure 3a, three different locations on the sample were marked, from which the PFM amplitude and phase data were obtained. Locations 1 and 2 were situated on the AlN layer, while location 3 was on the Si substrate, for which no piezoelectric response should be present. PFM measurements were also performed at several other locations on both the AlN NM and Si substrate to confirm their uniformity, and the results obtained are very similar to those shown in Figure 3b. Figure 3b shows the amplitude and phase of surface modulation induced by applied bias voltages  $V_{dc} = 0$  V and  $V_{ac} = 4$  V at locations 1–3

indicated in Figure 3a. The electromechanical response associated specifically with the AlN NM layer was then determined using a background subtraction approach, described in detail elsewhere,<sup>[67,68]</sup> with the PFM measurements performed using values for  $V_{dc}$  and  $V_{ac}$  of 0-4 and 3-5 V, respectively, at a frequency of 60 kHz. Negligible variation in measurement results was observed over these ranges of voltages and across multiple samples. The resulting value of  $|d_{33,eff}|$  for the AlN NM was found to be 0.65  $\pm$  0.12 pm V<sup>-1</sup>. We assume based on STEM results reported here (Figure 2) that the AlN NM after transfer is N-polar; because of parasitic tip-sample capacitance effects,<sup>[69]</sup> the polarity of the nanomembrane cannot be determined definitively from these PFM measurements. We can then correct for the deformation of the Si substrate during the PFM measurement using the analysis in the reference,<sup>[70]</sup> values for the Young's modulus of AlN and Si of 340 and 170 GPa, respectively, and Poisson's ratio of AlN and Si of 0.25 and 0.27, respectively.<sup>[71–73]</sup> The resulting value for  $|d_{33}|$  of the AlN NM ( $|d_{33(AIN)}|$ ) was =  $5.2 \pm 1.1 \text{ pm V}^{-1}$ , in excellent agreement with values reported in the literature of 5.0-5.6 pm V<sup>-1.[74-76]</sup>

The bandgap of AlN has been reported to be ≈6.1 eV by measuring the optical absorption spectra (Tauc plot) and photoluminescence (PL) spectra.<sup>[77,78]</sup> Further, to construct the energy band diagram and explore the surface charge states, the VBM using the XPS measurements was calculated on both Al-polar and N-polar, as shown in Figure 4a. The charge correction was conducted using the C 1s peak at 284.6 eV<sup>[79]</sup> and a Shirley background<sup>[80]</sup> was subtracted. The energy difference between the VBM and Fermi level,  $q\phi_s$ , was extracted as the X-intercept of the linear fitting curve on the XPS spectrum around the VBM (Figure 4b). The AlN NMs coated with an ultrathin ALD deposited SiO<sub>2</sub> or Al<sub>2</sub>O<sub>3</sub> were also measured to investigate the effect of UO coating on the surface charge. The uncertainty of the measurement was estimated to be 0.2 eV.[81,82] As shown in Figure 4c, the bare Al-polar surface of the AlN NM sample has a larger  $q\phi_S$  than the bare N-polar surface, indicating that there is more positive charge on the bare Al-polar surface than the bare N-polar surface. If only considering the polarization charge, under the tensile strain, the spontaneous polarization and piezoelectric polarization would be parallel to each other.<sup>[83]</sup>







**Figure 4.** a) Schematic band diagram at the surface of aluminum nitride (AIN). b) An example of AIN valence band X-ray photoelectron spectroscopy (XPS) spectrum. The black dash line is the linear fitting curve to extract the value of  $q\phi_S$  as shown in the graph. c) The summary of  $q\phi_S$  measured on both Al-polar (solid) and N-polar (dash) surfaces, with ultrathin ALD-Al<sub>2</sub>O<sub>3</sub> (black) or ALD-SiO<sub>2</sub> (red) coating.

This means that more negative charge is expected on the Alpolar surface than N-polar according to the polarization of AlN, which is contrary to what we observed through XPS. Therefore, one reasonable speculation is that the strong Fermi level pinning effect from the surface states, which has been observed in nitride materials through XPS measurements,<sup>[79,84]</sup> has dominated the polarization effect, and more positive charge is observed on the Al-polar surface. In addition, the deposition of five cycles of ALD-SiO<sub>2</sub> and postdeposition annealing in N<sub>2</sub> ambient at 350 °C for 5 min by RTA generates more positive charge to the Al-polar surface. On the other hand, ALD-Al<sub>2</sub>O<sub>3</sub> tends not to have much impact on the surface charge of AlN NMs. It is expected that an optimized ultrathin interfacial dielectric layer can potentially reduce the interface state density,<sup>[4]</sup> thus promising for AlN-based heterostructure vertical devices.

#### 3. Conclusion

We have successfully created, transferred, and characterized single-crystalline AlN NMs. The AlN NM crystalline quality has been retained after the creation and transfer-printing processes. The crystalline quality has been characterized and confirmed by a broad range of techniques. The single-crystalline pure wurtzite phase *c*-axis orientation has been confirmed by STEM and XRD. The piezoelectric response has been confirmed by PFM measurements. The energy band-bending properties of both Al-polar and N-polar surfaces have been characterized by detailed XPS measurements. The excellent material, electronic, and piezoelectric properties obtained from the AlN NMs collectively are expected to enable a new range of AlN-based heterogeneous electronics and optoelectronics on both rigid and flexible foreign substrates in the future.

#### 4. Experimental Section

Nanomembrane Fabrication and Transfer: The 50 nm/400 nm unintentionally doped (UID) AlN (0001)/Al (111) epi layer was grown on the Si (111) substrate via a Veeco GENxplor MBE system. The thickness of the AlN epi layer has been confirmed by STEM (Figure S5, Supporting Information). Detailed study of the growth technique will be reported separately. In brief, the single crystalline Al (111) of 400 nm was firstly grown on a Si (111) substrate. On top of the epitaxial Al (111) layer, a

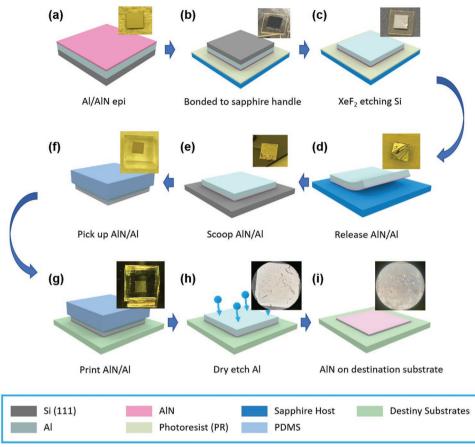
single-crystalline wurtzite AlN of 50 nm was grown. The finished AlN epi layer is Al-polar.

After fabricating the epitaxial wafer of AIN, it was first diced into 4×4 mm<sup>2</sup> pieces (Figure 1a). To release the AIN NMs from the epitaxial wafer, a standard semiconductor cleaning process was first performed on the AlN/Al/Si epi using acetone, isopropyl alcohol, and deionized (DI) water sequentially in the ultrasonic bath for 10 min each and then dried using N2. The epi wafer was (with AIN facing) attached to a temporary sapphire substrate, which had been spin-coated by the photoresist (PR) AZ5214 (Figure 5b). The whole stack of Si/Al/AlN/PR/sapphire (from top to bottom) was then baked together at 120 °C for 15 min on a hotplate. The Si was selectively removed by the SPTS Xactix E1 XeF<sub>2</sub> etcher, with the NM form of AlN/Al left on the temporary sapphire substrate (Figure 5c). Further, the remaining stack of AIN/AI/PR/sapphire was soaked into acetone to undercut the PR layer and release the AIN/AI NM from the sapphire substrate (Figure 5d). It is noteworthy that the freestanding AIN/AI NMs could curl down, as AIN layer is inside and AI layer is outside, in the acetone due to strain originated from the lattice mismatch between AIN and AI when released from the host substrate. To flatten the AIN/AI NM, acetone was gradually drained and meanwhile DI water was gradually added. The DI water, having stronger surface tension than acetone,<sup>[85-87]</sup> helped to unfold the curled AIN/AI NMs and make it float on the liquid surface. As a next step, a Si substrate was used to "scoop" the floating AIN/AI NMs from the DI water (Figure 5e).<sup>[10]</sup> Further, the DI water residues were dried out using N2. Moving forward, the AIN/AI NMs, sitting on the Si substrate, was picked up by a polydimethylsiloxane (PDMS) stamp (shown in Figure 5f) and was transfer-printed onto the destination substrate (e.g., Si, and sapphire) as shown in Figure 5g. It is noted that the existence of the Al (111) interlayer helps "soften" the rigidity of AIN NM, which makes the PDMS-assisted transfer process feasible as the Al/AlN NM becomes more compliant to the substrate. The PDMS stamp was also surface-textured, rather than flat, to reduce its surface adhesion to the NM to make the transfer possible. Finally, the free-standing single-crystalline AIN NM was obtained by removing the top 400 nm-thick Al layer using Plasma-Therm RIE (Figures 5h and 5i). For the etching process, a mixture of BCl<sub>3</sub>, Cl<sub>2</sub>, and Ar gases was used which has high selectivity between Al and AlN layers. XPS survey spectra were measured before and after the RIE etching to confirm removal of the Al layer (Figure S2, Supporting Information). The remaining AlN NM thickness after the Al layer etching was 30 nm.

Compared with deposition of AIN thin films, the transfer-printing method allows to inherit the high single-crystalline quality of AIN grown using state-of-the-art MBE technique regardless of the final host substrates, which is beneficial to applications such as AIN-based flexible electronics<sup>[83]</sup> and III-V/Si integration.<sup>[88]</sup> In addition, via substrate flipping the top surface of the as-grown epi wafer was utilized as the contact surface, which had better crystal quality than the bottom surface of AIN due to the nucleation at the bottom surface. In comparison to the exfoliation method of the remote epitaxy growth, the method does not require the deposition of a stressor layer or handling layer on top of







**Figure 5.** Fabrication and transfer-printing process of the aluminum nitride (AIN) nanomembrane (NM). a) Structure of the AlN/Al/silicon (Si) epi with the photo of the diced epi sample. b) AlN/Al/Si epi is flipped and bonded to the sapphire handle with the photo of the bonded sample. c) Si substrate is removed by XeF<sub>2</sub> dry etching with the photo of the bonded sample after etching. d) AlN/Al NM is released by undercutting photoresist with the photo of the AlN/Al NM floating in acetone solution and e) scooped by a Si handle wafer with the photo of the AlN/Al NM on the Si handle wafer. f) The AlN/Al NM is picked up by a polydimethylsiloxane (PDMS) stamp with the photo of the AlN/Al NM on the PDMS stamp, and g) then transfer-printed to a destiny substrate with the photo of the AlN/Al NM being printed by the PDMS stamp. h,i) The free-standing AlN NM is finalized by dry etching of the Al layer with the photos of the AlN/Al NM h) before and i) after etching. The final thickness of the AlN NM is only 30 nm resulting from overetching in the Al etching step.

the as-grown  $\mathsf{NM}^{[89]}$  Therefore, this clean and smooth contact surface provides the possibility to create a firmly bonded interface with the substrate.

Additionally, such MBE-grown AIN thin films have good scaling and doping controllability. In the current experiments, the thickness of the as-grown AIN film could be controlled from even monolayer up to 200 nm that was crack-free and with a smooth surface. Furthermore, to develop high-performance transistors that were based on such an AIN membrane, the doping profile of the AIN film could be precisely modified by introducing efficient dopants, such as beryllium (Be) and silicon (Si).<sup>[90,91]</sup>

Topography and Structural Measurements: SEM images (Figure 1a–c) of the transferred AIN NM were taken by Zeiss LEO 1530 field emission SEM system. The AIN NM was transferred to a patterned Si substrate that has an array of hexagon shaped pillars. XRD 2theta-omega spectra (Figure 1e) were measured on both as-grown and transferred AIN samples by a Malvern Panalytical Empyrean X-ray diffractometer with a Cu K- $\alpha$  X-ray source. Raman spectra (Figure 1f) were obtained from both as-grown and transferred AIN samples by a Horiba LabRAM HR Evolution Raman spectrometer using a 532-nm laser. Considering the fact that the AIN NM is only 54 nm thick, the confocal hole size of the spectrometer and the acquisition time were set to be 25  $\mu$ m and 20 min,

respectively, to increase the axial resolution and the layer discrimination while maintaining a moderate signal-to-noise ratio. The EBSD map (Figure 1g) was collected on the transferred AIN NM at 20 kV and 10 nA with a Zeiss Merlin High-resolution SEM system equipped with an EBSD detector. AFM images was taken by a Bruker Dimension Icon AFM system.

*Microstructural Measurements*: STEM specimen was prepared by a Thermo Fisher Scientific Helios G4 UXe focus ion beam. HAADF– STEM and ABF–STEM images were collected using a Cs corrected JEOL 3100R05 microscope operated at 300 kV.

Piezoresponse Force Microscopy Measurements: PFM measurements were performed using a Bruker Dimension Icon AFM system. In this measurement, an AFM with a conductive probe was used to create a vertical electric field through the sample material by applying a driving voltage between the tip and substrate while the AFM operated in contact mode, as shown schematically in Figure S6 (Supporting Information). In the presence of piezoelectricity in the sample, the AFM tip would be deflected by the expansion and contraction of the material induced by the converse piezoelectric effect. A lock-in amplifier was used to detect the amplitude and phase of the response. The effective piezoelectric response,  $d_{33,eff}$ , was then obtained via an analysis procedure described in detail elsewhere.<sup>[67,68]</sup> Because of the substantial difference in elastic

DVANCED ENCE NEWS www.advancedsciencenews.com

modulus between the AIN NM and Si (100) substrate, it was also necessary to account for deformation of the Si substrate beneath the AIN NM, which would alter the surface deflection measured in PFM.<sup>[70]</sup> The true piezoelectric coefficient  $d_{33}$  of the AIN layer could then be determined based on the analysis given in the reference,<sup>[70]</sup> and was found to be substantially larger than  $d_{33,eff}$ .

Valance Band Maximum Measurements: The XPS characterizations were performed on as-grown and transferred AIN samples with a Thermo Scientific K Alpha X-ray Photoelectron Spectrometer with Al  $K_{\alpha}$ X-ray source (hv = 1486.6 eV). The detailed setup of the XPS instrumental parameters has been reported in the previous works.<sup>[81,82]</sup>

## **Supporting Information**

Supporting Information is available from the Wiley Online Library or from the author.

#### Acknowledgements

J.G., J.Z., and P.W. contributed equally to this work. The work was supported by the Defense Advanced Research Projects Agency H2 program under the grant HR00112190107 and partially by Air Force Office of Scientific Research under grant FA9550-21-1-0081. Part of this work was supported by the National Science Foundation under award No. DMR-1905287.

## **Conflict of Interest**

The authors declare no conflict of interest.

# **Data Availability Statement**

The data that support the findings of this study are available from the corresponding author upon reasonable request.

### **Keywords**

aluminum nitride (AIN), piezoresponse force microscopy, scanning transmission electron microscopy, single-crystalline nanomembranes, transfer-printing, ultra-wide bandgap (UWBG), X-ray photoelectron spectroscopy

> Received: December 10, 2022 Revised: February 4, 2023 Published online: March 4, 2023

- [1] M. A. Meitl, Z.-T. Zhu, V. Kumar, K. J. Lee, X. Feng, Y. Y. Huang, I. Adesida, R. G. Nuzzo, J. A. Rogers, Nat. Mater. 2006, 5, 33.
- [2] J. Rogers, M. Lagally, R. Nuzzo, Nature 2011, 477, 45.
- [3] A. Carlson, A. M. Bowen, Y. Huang, R. G. Nuzzo, J. A. Rogers, Adv. Mater. 2012, 24, 5284.
- [4] D. Liu, S. J. Cho, J.-H. Seo, K. Kim, M. Kim, J. Shi, X. Yin, W. Choi, C. Zhang, J. Kim, M. A. Baboli, J. Park, J. Bong, I.-K. Lee, J. Gong, S. Mikael, J. H. Ryu, P. K. Mohseni, X. Li, S. Gong, X. Wang, Z. Ma, arXiv preprint arXiv:1812.10225, 2018.
- [5] D.-Y. Khang, H. Jiang, Y. Huang, J. A. Rogers, Science 2006, 311, 208.
- [6] Y. Sun, W. M. Choi, H. Jiang, Y. Y. Huang, J. A. Rogers, Nat. Nanotechnol. 2006, 1, 201.



#### www.advelectronicmat.de

- [7] J.-H. Ahn, H.-S. Kim, K. J. Lee, S. Jeon, S. J. Kang, Y. Sun, R. G. Nuzzo, J. A. Rogers, Science 2006, 314, 1754.
- [8] D.-H. Kim, J.-H. Ahn, W. M. Choi, H.-S. Kim, T.-H. Kim, J. Song, Y. Y. Huang, Z. Liu, C. Lu, J. A. Rogers, Science 2008, 320, 507.
- [9] J. Yoon, A. J. Baca, S.-I. Park, P. Elvikis, J. B. Geddes, L. Li, R. H. Kim, J. Xiao, S. Wang, T.-H. Kim, Nat. Mater. 2008, 7, 907.
- [10] A. M. Kiefer, D. M. Paskiewicz, A. M. Clausen, W. R. Buchwald, R. A. Soref, M. G. Lagally, ACS Nano 2011, 5, 1179.
- [11] J. H. Seo, T. Y. Oh, J. Park, W. Zhou, B. K. Ju, Z. Ma, Adv. Funct. Mater. 2013, 23, 3398.
- [12] S.-K. Kang, R. K. Murphy, S.-W. Hwang, S. M. Lee, D. V. Harburg, N. A. Krueger, J. Shin, P. Gamble, H. Cheng, S. Yu, Nature 2016, 530.71.
- [13] J.-H. Seo, K. Zhang, M. Kim, W. Zhou, Z. Ma, npj Flexible Electron. 2017, 1, 1.
- [14] D. Liu, S. J. Cho, J. Park, J.-H. Seo, R. Dalmau, D. Zhao, K. Kim, J. Gong, M. Kim, I.-K. Lee, Appl. Phys. Lett. 2018, 112, 081101.
- [15] D. Liu, S. J. Cho, J. Park, J. Gong, J.-H. Seo, R. Dalmau, D. Zhao, K. Kim, M. Kim, A. R. Kalapala, Appl. Phys. Lett. 2018, 113, 011111.
- [16] Y. Liu, L. Wang, L. Wang, X. Wu, Z. Hao, J. Yu, Y. Luo, C. Sun, Y. Han, B. Xiong, Semicond. Sci. Technol. 2019, 34, 105023.
- [17] Y. Zheng, M. N. Hasan, J. H. Seo, Adv. Mater. Technol. 2021, 6, 2100254.
- [18] K. Kim, J. Jang, H. Kim, Results Phys. 2021, 25, 104279.
- [19] C. Chen, C. Li, S. Min, Q. Guo, Z. Xia, D. Liu, Z. Ma, F. Xia, Nano Lett. 2021, 21, 8385.
- [20] M. N. Hasan, Y. Zheng, J. Lai, E. Swinnich, O. G. Licata, M. A. Baboli, B. Mazumder, P. K. Mohseni, J. H. Seo, Adv. Mater. Interfaces 2022, 9, 2101531.
- [21] Z. Xia, H. Song, M. Kim, M. Zhou, T.-H. Chang, D. Liu, X. Yin, K. Xiong, H. Mi, X. Wang, Sci. Adv. 2017, 3, e1602783.
- [22] S.-I. Park, Y. Xiong, R.-H. Kim, P. Elvikis, M. Meitl, D.-H. Kim, J. Wu, J. Yoon, C.-J. Yu, Z. Liu, Science 2009, 325, 977.
- [23] J. Yoon, S. Jo, I. S. Chun, I. Jung, H.-S. Kim, M. Meitl, E. Menard, X. Li, J. J. Coleman, U. Paik, Nature 2010, 465, 329.
- [24] J. Lee, J. Wu, M. Shi, J. Yoon, S. I. Park, M. Li, Z. Liu, Y. Huang, J. A. Rogers, Adv. Mater. 2011, 23, 986.
- [25] K. Xiong, H. Mi, T. H. Chang, D. Liu, Z. Xia, M. Y. Wu, X. Yin, S. Gong, W. Zhou, J. C. Shin, Energy Sci. Eng. 2018, 6, 47.
- [26] S. J. Cho, D. Liu, A. Hardy, J. Kim, J. Gong, C. J. Herrera-Rodriguez, E. Swinnich, X. Konstantinou, G.-Y. Oh, D. G. Kim, AIP Adv. 2020, 10, 125226.
- [27] S. H. Park, G. Yuan, D. Chen, K. Xiong, J. Song, B. Leung, J. Han, Nano Lett. 2014, 14, 4293.
- [28] T.-H. Chang, K. Xiong, S. H. Park, G. Yuan, Z. Ma, J. Han, Sci. Rep. 2017, 7, 6360.
- [29] Y.-Y. Zhang, Y.-X. Zheng, J.-Y. Lai, J.-H. Seo, K. H. Lee, C. S. Tan, S. An, S.-H. Shin, B. Son, M. Kim, ACS Nano 2021, 15, 8386.
- [30] Y. Y. Zhang, S. An, Y. Zheng, J. Lai, J. H. Seo, K. H. Lee, M. Kim, Adv. Electron. Mater. 2022, 8, 2100652.
- [31] M. Kim, J.-H. Seo, D. Zhao, S.-C. Liu, K. Kim, K. Lim, W. Zhou, E. Waks, Z. Ma, J. Mater. Chem. C 2017, 5, 264.
- [32] E. Swinnich, M. N. Hasan, K. Zeng, Y. Dove, U. Singisetti, B. Mazumder, J. H. Seo, Adv. Electron. Mater. 2019, 5, 1800714.
- [33] Y. Zheng, E. Swinnich, J.-H. Seo, ECS J. Solid State Sci. Technol. 2020, 9, 055007.
- [34] J. Lai, M. N. Hasan, E. Swinnich, Z. Tang, S.-H. Shin, M. Kim, P. Zhang, J.-H. Seo, J. Mater. Chem. C 2020, 8, 14732.
- [35] Y. Zheng, Z. Feng, A. A. U. Bhuiyan, L. Meng, S. Dhole, Q. Jia, H. Zhao, J.-H. Seo, J. Mater. Chem. C 2021, 9, 6180.
- [36] M. Feneberg, R. A. Leute, B. Neuschl, K. Thonke, M. Bickermann, Phys. Rev. B 2010, 82, 075208.
- [37] J. Tsao, S. Chowdhury, M. Hollis, D. Jena, N. Johnson, K. Jones, R. Kaplar, S. Rajan, C. Van de Walle, E. Bellotti, Adv. Electron. Mater. 2018, 4, 1600501.

#### **ADVANCED** SCIENCE NEWS

www.advancedsciencenews.com

- [38] G. A. Slack, R. A. Tanzilli, R. Pohl, J. Vandersande, J. Phys. Chem. Solids 1987, 48, 641.
- [39] J. Li, Z. Fan, R. Dahal, M. Nakarmi, J. Lin, H. Jiang, Appl. Phys. Lett. 2006, 89, 213510.
- [40] T. Li, L. Long, Z. Hu, R. Wan, X. Gong, L. Zhang, Y. Yuan, J. Yan, W. Zhu, L. Wang, Opt. Lett. 2020, 45, 3325.
- [41] Y. Irokawa, E. A. G. Villora, K. Shimamura, Jpn. J. Appl. Phys. 2012, 51, 040206.
- [42] H. Fu, I. Baranowski, X. Huang, H. Chen, Z. Lu, J. Montes, X. Zhang, Y. Zhao, IEEE Electron Device Lett. 2017, 38, 1286.
- [43] K. Tonisch, V. Cimalla, C. Foerster, H. Romanus, O. Ambacher, D. Dontsov, Sens. Actuators, A 2006, 132, 658.
- [44] M. D. Williams, B. A. Griffin, T. N. Reagan, J. R. Underbrink, M. Sheplak, J. Microelectromech. Syst. 2012, 21, 270.
- [45] N. Jackson, L. Keeney, A. Mathewson, Smart Mater. Struct. 2013, 22, 115033.
- [46] V. M. Mastronardi, F. Guido, M. Amato, M. De Vittorio, S. Petroni, *Microelectron. Eng.* 2014, 121, 59.
- [47] X. Yang, J. Liang, Y. Jiang, X. Chen, H. Zhang, M. Zhang, W. Pang, IEEE Electron Device Lett. 2017, 38, 1125.
- [48] A. Gao, K. Liu, J. Liang, T. Wu, Microsyst. Nanoeng. 2020, 6, 74.
- [49] J. Simon, V. Protasenko, C. Lian, H. Xing, D. Jena, Science 2010, 327, 60.
- [50] Y. Mei, D. J. Thurmer, C. Deneke, S. Kiravittaya, Y.-F. Chen, A. Dadgar, F. Bertram, B. Bastek, A. Krost, J. r. Christen, ACS Nano 2009, 3, 1663.
- [51] J. Yang, Y. Wang, L. Wang, Z. Tian, Z. Di, Y. Mei, J. Semicond. 2020, 41, 042601.
- [52] A. Khandelwal, N. Athreya, M. Q. Tu, L. L. Janavicius, Z. Yang, O. Milenkovic, J.-P. Leburton, C. M. Schroeder, X. Li, *Microsyst. Nanoeng.* 2022, 8, 27.
- [53] A. Khandelwal, Z. Ren, S. Namiki, Z. Yang, N. Choudhary, C. Li, P. Wang, Z. Mi, X. Li, ACS Appl. Mater. Interfaces 2022, 14, 29014.
- [54] E. Azrak, L. G. Michaud, A. Reinhardt, S. Tardif, M. Bousquet, N. Vaxelaire, J. Eymery, F. Fournel, P. Montmeat, ECS J. Solid State Sci. Technol. 2021, 10, 064001.
- [55] Y. Jia, H. Wu, J. Zhao, H. Guo, Y. Zeng, B. Wang, C. Zhang, Y. Zhang, J. Ning, J. Zhang, *CrystEngComm* **2021**, *23*, 7406.
- [56] Z. Shi, X. Lu, X. Tang, D. Wang, Z. Cong, X. Ma, Y. Hao, Appl. Surf. Sci. 2022, 585, 152378.
- [57] P. Scherrer, Kolloidchemie Ein Lehrbuch, Springer, Berlin 1912, p. 387.
- [58] A. Patterson, Phys. Rev. 1939, 56, 978.
- [59] T. Prokofyeva, M. Seon, J. Vanbuskirk, M. Holtz, S. Nikishin, N. Faleev, H. Temkin, S. Zollner, *Phys. Rev. B* 2001, *63*, 125313.
- [60] V. Lughi, D. R. Clarke, Appl. Phys. Lett. 2006, 89, 241911.
- [61] M. Morita, T. Ohmi, E. Hasegawa, M. Kawakami, M. Ohwada, J. Appl. Phys. 1990, 68, 1272.
- [62] M. Kazan, E. Moussaed, R. Nader, P. Masri, Phys. Status Solidi C 2007, 4, 204.
- [63] D. Zagorac, J. Zagorac, M. Djukic, D. Jordanov, B. Matović, Theor. Appl. Fract. Mech. 2019, 103, 102289.

- [64] W. Yim, R. Paff, J. Appl. Phys. 1974, 45, 1456.
- [65] P. Wang, Y. Yuan, C. Zhao, X. Wang, X. Zheng, X. Rong, T. Wang, B. Sheng, Q. Wang, Y. Zhang, *Nano Lett.* **2016**, *16*, 1328.
- [66] P. Wang, D. Wang, S. Mondal, Y. Wu, T. Ma, Z. Mi, ACS Appl. Mater. Interfaces 2022, 14, 15747.
- [67] C. J. Brennan, R. Ghosh, K. Koul, S. K. Banerjee, N. Lu, E. T. Yu, Nano Lett. 2017, 17, 5464.
- [68] C. J. Brennan, K. Koul, N. Lu, E. T. Yu, Appl. Phys. Lett. 2020, 116, 053101.
- [69] B. J. Rodriguez, A. Gruverman, A. Kingon, R. Nemanich, J. Cryst. Growth 2002, 246, 252.
- [70] J.-H. Li, L. Chen, V. Nagarajan, R. Ramesh, A. Roytburd, Appl. Phys. Lett. 2004, 84, 2626.
- [71] J. Wortman, R. Evans, J. Appl. Phys. 1965, 36, 153.
- [72] A. Wright, J. Appl. Phys. 1997, 82, 2833.
- [73] E. Österlund, J. Kinnunen, V. Rontu, A. Torkkeli, M. Paulasto-Kröckel, J. Alloys Compd. 2019, 772, 306.
- [74] I. Guy, S. Muensit, E. Goldys, Appl. Phys. Lett. 1999, 75, 4133.
- [75] C. Lueng, H. L. Chan, C. Surya, C. Choy, J. Appl. Phys. 2000, 88, 5360.
- [76] F. Bernardini, V. Fiorentini, Appl. Phys. Lett. 2002, 80, 4145.
- [77] P. Wang, D. A. Laleyan, A. Pandey, Y. Sun, Z. Mi, Appl. Phys. Lett. 2020, 116, 151903.
- [78] P. Wang, B. Wang, D. A. Laleyan, A. Pandey, Y. Wu, Y. Sun, X. Liu, Z. Deng, E. Kioupakis, Z. Mi, *Appl. Phys. Lett.* **2021**, *118*, 032102.
- [79] P. Reddy, I. Bryan, Z. Bryan, W. Guo, L. Hussey, R. Collazo, Z. Sitar, J. Appl. Phys. 2014, 116, 123701.
- [80] D. A. Shirley, Phys. Rev. B 1972, 5, 4709.
- [81] J. Gong, K. Lu, J. Kim, T. Ng, D. Kim, J. Zhou, D. Liu, J. Kim, B. S. Ooi, Z. Ma, Jpn. J. Appl. Phys. 2021, 61, 011003.
- [82] J. Gong, Z. Zheng, D. Vincent, J. Zhou, J. Kim, D. Kim, T. K. Ng, B. S. Ooi, K. J. Chen, Z. Ma, J. Appl. Phys. 2022, 132, 135302.
- [83] C. Fei, X. Liu, B. Zhu, D. Li, X. Yang, Y. Yang, Q. Zhou, Nano Energy 2018, 51, 146.
- [84] Y. Jia, J. S. Wallace, E. Echeverria, J. A. GardellaJr, U. Singisetti, Phys. Status Solidi B 2017, 254, 1600681.
- [85] K. S. Howard, R. McAllister, AIChE J. 1957, 3, 325.
- [86] S. Enders, H. Kahl, J. Winkelmann, J. Chem. Eng. Data 2007, 52, 1072.
- [87] N. Wu, X. Li, S. Liu, M. Zhang, S. Ouyang, Appl. Sci. 2019, 9, 1235.
- [88] G. Roelkens, J. Van Campenhout, J. Brouckaert, D. Van Thourhout, R. Baets, P. R. Romeo, P. Regreny, A. Kazmierczak, C. Seassal, X. Letartre, *Mater. Today* 2007, 10, 36.
- [89] H. Kim, C. S. Chang, S. Lee, J. Jiang, J. Jeong, M. Park, Y. Meng, J. Ji, Y. Kwon, X. Sun, Nat. Rev. Methods Primers 2022, 2, 1.
- [90] H. Ahmad, J. Lindemuth, Z. Engel, C. M. Matthews, T. M. McCrone, W. A. Doolittle, *Adv. Mater.* 2021, 33, 2104497.
- [91] H. Ahmad, Z. Engel, C. M. Matthews, S. Lee, W. A. Doolittle, J. Appl. Phys. 2022, 131, 175701.

www.advelectronicmat.de