

# 56] Angular Momentum Recollection + Application

$$m \frac{d^2 \vec{r}}{dt^2} = \vec{F}$$

$$m \vec{r} \times \frac{d^2 \vec{r}}{dt^2} = \vec{r} \times \vec{F} \equiv \text{torque} \equiv \vec{\tau}$$

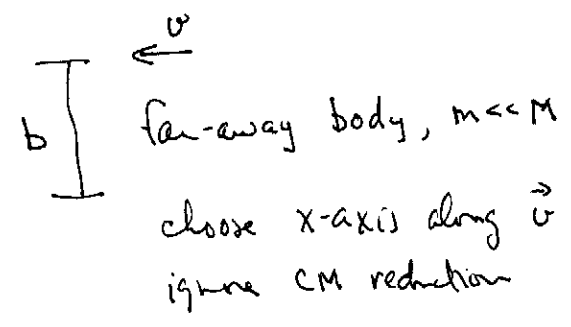
$$\frac{d}{dt} (m \vec{r} \times \frac{d\vec{r}}{dt}) = \frac{d}{dt} \vec{L} \Rightarrow \frac{d\vec{L}}{dt} = \vec{\tau}$$

$$\vec{L} = m \vec{r} \times \vec{v} = \vec{r} \times \vec{p} = \text{angular momentum}$$

Analogue to  $\frac{d\vec{p}}{dt} = \vec{F}$  for linear momentum.

In central forces  $\vec{\tau} = 0 \Rightarrow \vec{L}$  is conserved. We used this, ~~found~~ recently (Kepler's 2<sup>nd</sup>).

example: Gravitational capture



$$|L| = m b v$$
$$E = \frac{1}{2} m \dot{r}^2 + \frac{L^2}{2mr^2} - \frac{GMm}{r} = \frac{1}{2} m \dot{r}^2 + \frac{m b^2 v^2}{2r^2} - \frac{GMm}{r}$$
$$= E_0 = \frac{1}{2} m v^2$$

At close Is  $r=R$  allowed?

Requires ( $\dot{r}^2 > 0$ )  $\frac{1}{2} m v^2 - \frac{m b^2 v^2}{2R^2} + \frac{GMm}{R} \geq 0$

$$\boxed{v^2 \left( \frac{b^2}{R^2} - 1 \right) \leq \frac{2GM}{R}} \quad \text{or} \quad \boxed{\left| \frac{b^2}{R^2} \leq 1 + \frac{2GM}{Rv^2} \right|}$$

This is  $\rightarrow b \approx R$  as  $v \rightarrow \infty$ , but is enhanced over (75)  
 "geometrical" capture for finite  $v$ .

### 57) Angular Momentum for Systems

$$m_i \frac{d^2 \vec{r}_i}{dt^2} = \vec{F}_i + \sum_j \vec{F}_{ij}$$

← external forces

← internal forces

$$\sum_i m_i \vec{r}_i \times \frac{d^2 \vec{r}_i}{dt^2} = \sum_i \vec{r}_i \times \vec{F}_i + \underbrace{\sum_i \sum_j \vec{r}_i \times \vec{F}_{ij}}$$

work on this

$$\sum_i \sum_j \vec{r}_i \times \vec{F}_{ij} = \frac{1}{2} (\sum_i \sum_j + \sum_j \sum_i) (\vec{r}_i \times \vec{F}_{ij})$$

$$= \frac{1}{2} \sum_i \sum_j (\vec{r}_i \times \vec{F}_{ij} + \vec{r}_j \times \vec{F}_{ji})$$

$$\stackrel{3^{rd} \text{ law}}{=} \frac{1}{2} \sum_i \sum_j (\vec{r}_i - \vec{r}_j) \times \vec{F}_{ij}$$

Now if  $\vec{F}_{ij} \propto \vec{r}_i - \vec{r}_j$  (central force), this = 0!

Then

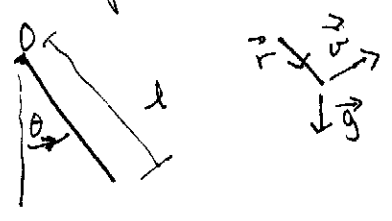
$$\sum_i m_i \vec{r}_i \times \frac{d^2 \vec{r}_i}{dt^2} = \sum_i \vec{r}_i \times \vec{F}_i$$

$$\sum_i \frac{d \vec{L}_i}{dt} = \sum_i \vec{c}_i$$

$$\text{or } \frac{d \vec{L}}{dt} = \vec{c} \quad \text{with } \vec{L} = \sum_i \vec{L}_i, \quad \vec{c} = \sum_i \vec{c}_i$$

The result is more general than this derivation.

Example: rod-pendulum



$$\lambda = \text{mass/length}$$

$$\vec{L} = \int_0^l dm \vec{r} \times \vec{v} = \hat{z} \int_0^l \lambda dr \cdot r \cdot r \dot{\theta} = \hat{z} \frac{l^3}{3} \lambda \dot{\theta} = \hat{z} \frac{Ml^2}{3} \dot{\theta}$$

↑  
into plane

general theorem for near-Earth gravity

$$\vec{\tau} = \int \vec{r} \times d\vec{F}, \quad d\vec{F} = dm\vec{g}$$

⇓

$$\vec{\tau} = \int dm \vec{r} \times \vec{g} = M \vec{r}_{cm} \times \vec{g} = \vec{r}_{cm} \times M \vec{g}$$

$$\left\{ \begin{aligned} &= (\int dm \vec{r}) \times \vec{g} \\ &= M \vec{r}_{cm} \times \vec{g} \end{aligned} \right.$$

⇒ The torque may be calculated as if the force is acting at the CM

N.B. Near-Earth gravity only!!

In present example:

$$\vec{\tau} = Mg \frac{l}{2} \sin \theta \hat{z} \quad (\text{evaluating } \times\text{-product})$$

so  $\frac{d\vec{L}}{dt} = \vec{\tau}$  reads  $-\frac{Ml^2}{3} \ddot{\theta} = \frac{Mgl}{2} \sin \theta$ , or  $\ddot{\theta} = -\frac{3g}{2l} \sin \theta$

small  $\theta$ :  $\ddot{\theta} = -\frac{3g}{2l} \theta$ ;  $\omega^2 = \frac{3g}{2l}$  oscillation.

# 58 | Statics

If a rigid body is not moving  
(no internal motion!)

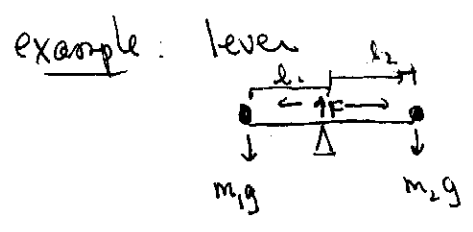
$$i) \sum \vec{F}^{ext.} = 0$$

$$ii) \sum \vec{\tau}^{ext.} = 0$$

since  $\frac{d\vec{p}}{dt} = 0$ ,  $\frac{d\vec{L}}{dt} = 0$ . And if  $\sum \vec{F}^{ext.} \neq 0$  or  $\sum \vec{\tau}^{ext.} \neq 0$ ,

for a rigid body, you will have visible motion. So  
(buckle)

i) and ii) are the fundamental equations of statics. This is a vast subject in engineering (you want buildings, bridges, etc. to be static!)

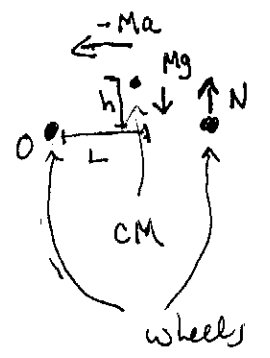


$$F = m_1 g + m_2 g$$

$$\text{around fulcrum } l_2 m_2 g - l_1 m_1 g = 0 \quad \frac{l_1}{l_2} = \frac{m_2}{m_1}$$

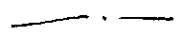
example: motorcycle lift-off

$\xrightarrow{a}$   
accelerate



torques around O:  $hMa - LMg + 2LN = 0$   
with  $N \geq 0$   
possible only for  $a < \frac{L}{h} g$   
with bigger a, "lift off"

Note high CM and limited wheelbase are conducive to liftoff.



Converse? IF  $\Sigma F = 0, \Sigma \tau = 0$  we have  $\vec{P} = \vec{P}_0, \vec{L} = \vec{L}_0$ . The first is removed "trivially" by going to CM frame. The 2nd ( $\vec{L}_0 \neq 0$ ) is much more interesting. It does not correspond to a constant rotation, ~~well~~ in general. We'll treat it down the road.

