Reduced-order, trajectory piecewise-linear models for nonlinear computational fluid dynamics

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Abstract—A trajectory piecewise-linear (TPWL) approach is developed for a computational fluid dynamics (CFD) model of the two-dimensional Euler equations. The approach uses a weighted combination of linearized models to represent the nonlinear CFD system. The proper orthogonal decomposition (POD) is then used to create a reduced-space basis, onto which the TPWL model is projected. This projection yields an efficient reduced-order model of the nonlinear system, which does not require the evaluation of any full-order system residuals. The method is applied to the case of flow through an actively controlled supersonic diffuser. With an appropriate choice of linearization points and POD basis vectors, the method is found to yield accurate results, including cases with significant shock motion.

I. INTRODUCTION

Computational fluid dynamics (CFD) is now widely used throughout the fluid dynamics community. It produces accurate models for many problems of interest, although the cost of obtaining the solution may be prohibitive for some applications. In particular, this cost becomes critical for multidisciplinary applications such as aeroelasticity or active flow controller design. Model order reduction techniques provide a way to systematically determine low-order models that capture the relevant dynamics of the CFD model while being computationally very efficient. These techniques have been applied successfully for a range of fluid dynamic applications [1], [2], [3], [4], [5].

While model reduction is now a well established approach for large linear systems, addressing the problems that arise for consideration of nonlinearity remains a challenging task. A number of linear reduction techniques have been extended to the nonlinear case with varying success. One approach to generate reduced-order models for nonlinear systems is a polynomial (Taylor) expansion of system's nonlinearity, and subsequent application of Krylov projection methods [6], [7], [8]. However, the main drawbacks of those methods are that they are limited to applications with "small" input disturbances, or more generally called weakly nonlinear systems, and that the quadratic and higher order expansion terms are very expensive to compute.

The proper orthogonal decomposition (POD) is a widely used method of reduction for CFD applications. The approach is to compute a set of empirical eigenfunctions using flow solutions collected from a simulation of the CFD model. A reduced-order model is then obtained by projecting the governing equations onto the reduced space spanned by the POD basis vectors. POD has been applied to nonlinear systems; however, in these applications, the issue of an efficient representation of the nonlinearity in the reduced-order model is inadequately addressed. While the resulting nonlinear models do have a reduced number of states, they still require flux evaluations of the original high-order CFD model.

In Rewienski, a trajectory piecewise-linear (TPWL) scheme, is developed [9], [10]. The basic idea is to represent a nonlinear system as a weighted combination of linear models. The linear models are obtained by linearizing the nonlinear system at various points along a training trajectory. This technique aims to address some fundamental issues presented earlier, i.e. overcoming restrictions of weak nonlinearity and creating a cost-efficient representation of the system's nonlinearity. By using a weighted combination of various linear models, a broader range of the nonlinear space is spanned compared with using a single model. In addition, the TPWL system allows an efficient representation of the reduced-order model. This technique has been successfully applied to nonlinear analogue circuits and micromachined devices [9], [10].

This paper considers the TPWL approach in conjunction with a POD-based reduction for CFD applications. These two techniques can be combined naturally, since both are based upon a training simulation of the system. In the next section, the CFD model is described, considering in particular the case of flow through an actively controlled supersonic diffuser. The general model reduction framework is then established, followed by a description of the TPWL approach and its application to the reduced-order models. Finally, results are presented and conclusions are drawn.

II. COMPUTATIONAL MODEL

The computational model is based on the case of flow through a supersonic diffuser; however, the TPWL methodology is general and could be applied to any CFD model. Figure 1 shows the Mach contours at steady-state conditions inside the fixed geometry of a supersonic diffuser that operates at a freestream Mach number of 2.2. In steady-state operation, a shock forms downstream of the throat; however in practice, the incoming supersonic flow is subject to perturbations, such as atmospheric density disturbances. Such perturbations in the flow may cause the shock to move upstream of the throat, and eventually to be expelled from the diffuser. This phenomenon, known as inlet unstart, causes huge losses in engine performance and thus is highly undesirable. In order to prevent inlet unstart, an active control mechanism of the shock is required.



Fig. 1. Mach contours for steady flow through supersonic diffuser. Steadystate inflow Mach number is 2.2.

Figure 2 presents the schematic of the actuation mechanism. Incoming flow with possible disturbances enter the inlet and is sensed using pressure sensors. The controller then adjusts the bleed upstream of the throat in order to control the position of the shock and to prevent it from moving upstream.



Fig. 2. Supersonic diffuser active flow control problem setup.

A. Nonlinear CFD Model

The full nonlinear solution of the entire flow distribution in the inlet can be obtained using a CFD model. Here, the problem is assumed to be two-dimensional, compressible and inviscid, thus the solution is governed by the Euler equations. The discrete Euler equations are derived from the integral form of the unsteady, two-dimensional equations, which are the usual statements of mass, momentum, and energy:

$$\frac{\partial}{\partial t} \iint \rho dV + \oint dm = 0$$

$$\frac{\partial}{\partial t} \iint \rho \vec{Q} dV + \oint \vec{Q} dm + \oint p d\vec{A} = 0$$

$$\frac{\partial}{\partial t} \iint \rho E dV + \oint H dm = 0$$
(1)

where the flow variables are the density, ρ , the total velocity vector, \vec{Q} , the pressure, p, the energy, E, and the total enthalpy, H. The quantity $dm = \rho \vec{Q} \cdot d\vec{A}$ is the mass flux element across the conservation cell boundary, $d\vec{A} = dA \cdot \hat{n}$, where dA is a surface element and \hat{n} is a unit vector pointing outward from the control volume. The discrete Euler equations approximate the integral form of the continuous Euler equations on small control volumes or control cells. The flow solver is fully described in Drela [11] and Lassaux [4], and uses as state variables q, q_{\perp}, ρ , and H, where q and q_{\perp} are the streamwise and normal components of the velocity, respectively. Using a structured grid for spatial discretization, the discrete Euler equations can be represented as a nonlinear dynamical system of the form:

$$\dot{x}(t) = f(x(t), u(t))$$

$$y(t) = g(x(t))$$

$$(2)$$

where $x(t) \in \mathbb{R}^n$ is a generalized state vector containing the *n* unknown flow quantities, q, q_{\perp}, ρ and *H*, at each point in the computational grid, *f* is a nonlinear vectorvalued function, $u(t) \in \mathbb{R}^l$ is the input to the system, and $y(t) \in \mathbb{R}^k$ contains the system outputs, which are defined by the nonlinear function *g*.

B. Reduced Space Basis

A reduced-order model can be obtained by considering a projection of the state vector x

$$x(t) = V\hat{x}(t) \tag{3}$$

where $\hat{x}(t) \in \mathbb{R}^m$ is the reduced-order state vector, containing the time-dependent amplitudes of m basis vectors, contained in the columns of the matrix V. The basis vectors must be selected appropriately, so that the state x can be accurately represented in the reduced space. In this work, POD is used to determine the basis vectors as follows.

First, N snapshots are obtained from a CFD calculation, where each snapshot corresponds to a flow solution at a particular instant in time. The correlation matrix R is formed by computing the inner product between every pair of snapshots

$$R_{ij} = \frac{1}{N} \left(x^{(i)}, x^{(j)} \right)$$
(4)

where $x^{(i)}$ is the flow solution at a time $t^{(i)}$ and $(x^{(i)}, x^{(j)})$ denotes the inner product between $x^{(i)}$ and $x^{(j)}$. The eigenvalues λ_i and eigenvectors $\psi^{(i)}$ of R are then computed. The magnitude of the j^{th} eigenvalue, λ_i , describes the relative importance of the j^{th} POD basis vector, V_j , which is computed by

$$V_j = \sum_{i=1}^{N} \psi_i^{(j)} x^{(i)}$$
(5)

where $\psi_i^{(j)}$ denotes the i^{th} component of the j^{th} eigenvector.

This orthonormal set of POD basis vectors can be used to project the solution onto the reduced-space basis using (3). The size of \hat{x} , m, will depend on the number of components taken in the basis V. This number can be chosen using a heuristic criterion based on capturing a sufficiently large amount of the "energy" contained in the snapshot collection. The relative energy e_j captured by each mode j is given by the POD eigenvalues as

$$e_j = \frac{\lambda_j}{\sum_{i=1}^N \lambda_i} \tag{6}$$

Applying the projection (3) to the nonlinear system (2), the resulting reduced-order model is of the form

$$\hat{x}(t) = V^T f(V\hat{x}(t), u(t))$$

$$\hat{y}(t) = g(V\hat{x}(t))$$
(7)

While the system (7) has a reduced number of states, it still requires evaluation of the full order flux term $f(\cdot)$. To obtain a truly reduced model, a more efficient representation of the nonlinearity in the reduced space is required.

C. Linearized Models

Efficient linearized models can be extracted from the system (2) by using a polynomial expansion of the nonlinearity, or more specifically a Taylor expansion about some state (x_i, u_i) , which, following Phillips [12], expands f as:

$$f(x,u) = f(x_i, u_i) + A_i(x - x_i) + B_i(u - u_i) + \frac{1}{2}W_i(x - x_i) \otimes (x - x_i) + \dots$$
(8)

where \otimes is the Kronecker product, and A_i and W_i are, respectively, the Jacobian and the Hessian of $f(\cdot)$ evaluated at the state (x_i, u_i) . The matrix $B_i = \frac{\partial f}{\partial u}$ is also evaluated at (x_i, u_i) . Dropping the quadratic and higher terms of (8), the nonlinear system (2) can be linearized about a given state to yield a state-space model of the form:

$$\dot{x}(t) = A_i x(t) + B_i u(t) + (f(x_i, u_i) - A_i x_i(t) - B_i u_i)$$

$$y(t) = C_i x(t) \tag{9}$$

where $C_i = \frac{\partial g}{\partial x}$ is also evaluated at (x_i, u_i) . The vector of unknowns x(t) can be written as

$$x(t) = x_i + x_i'(t) \tag{10}$$

where x_i , fixed in time, is the value of state vector x at the linearization point i, and $x'_i(t)$ contains the perturbation of the n unknown flow quantities about that linearization point x_i . The linearized equation (9) can then be expressed as

$$\dot{x}_{i}'(t) = A_{i}x_{i}'(t) + B_{1i}u(t) + B_{2i}$$

$$y(t) = C_{i}x_{i}'(t) + C_{0i}$$
(11)

where $B_{2i} = f(x_i, u_i) - B_i u_i$ and $C_{0i} = C_i x_i$.

The linearized system (11) is efficient for time computations, but remains too large for applications such as controller design. A reduced-order linearized model can be obtained by applying the projection (3) to the system (11) yielding

$$\frac{d}{dt}\hat{x}'_{i}(t) = \hat{A}_{i}\hat{x}'_{i}(t) + \hat{B}_{1i}u(t) + \hat{B}_{2i}
\hat{y}_{i}(t) = \hat{C}_{i}\hat{x}'_{i}(t) + C_{0i}$$
(12)

where the reduced-order matrices are given by

$$\hat{A}_{i} = V^{T}A_{i}V$$

$$\hat{B}_{1i} = V^{T}B_{1i}$$

$$\hat{B}_{2i} = V^{T}B_{2i}$$

$$\hat{C}_{i} = C_{i}V$$
(13)

The system (12) is truly reduced since the projections can be carried out to calculate the reduced-order matrices a priori and no CFD-order computations are required for simulation. However, the linearized models do not accurately capture nonlinear behavior. The next section will therefore focus on finding a suitable way to capture nonlinear behavior within the reduction framework.

D. Trajectory Piecewise-Linear Scheme

In Rewienski [9], an efficient, approximate method to represent nonlinear circuit systems is presented and tested. It is proposed that by using a weighted combination of multiple linear models, nonlinear behavior can be modelled. The linear models are obtained via linearization of the nonlinear system at different solutions in time. An approximation to the nonlinear system can then be obtained by using a weighted combination of the closest linear models to the current solution in time.



Fig. 3. Collection of linearization points x_0 , x_1 , x_2 and x_3 in a 2D state space. Circles denotes suitable region for use of each linearization point. Trajectory A is called the training trajectory. Figure from Rewienski [9].

Figure 3 presents a two-dimensional conceptual view of a series of linearized models. Plotted are four linearization points, x_0, x_1, x_2 and x_3 , along a "training trajectory", which is obtained using a simulation of the nonlinear system. The range of validity of each of the corresponding linearized models is denoted by the circles. In order to capture the most relevant dynamics of the system, the range of inputs simulated for the training trajectory should reflect dynamics of interest for the application at hand. For instance, in Figure 3, trajectories such as B and C will be well represented by the set of linear models, while trajectories D and E may demonstrate poor results, since they lie beyond the range of validity.

The linearization points can be chosen using the following approach. Consider N snapshots, taken from the training trajectory. The algorithm compares each pair of snapshots by computing the two-norm of the distance between them. When this difference is larger than a specified criterion, δ_{min} , a new linearization point is selected. The value of δ_{min} sets the distance between subsequent linearization points; therefore, lowering its value implies increasing the number of models in the system. This approach is described by the pseudoalgorithm below, which takes as inputs δ_{min} and the matrix U containing N CFD snapshots

$$U = \{x^{(0)}, x^{(1)}, \dots, x^{(N-1)}\}$$
(14)

and returns the vector linPt, which contains the column index in U of the selected linearization points.

Algorithm 1

(Choice of linearization points on the fly) Function linPt = linearizationPoint(U, δ_{min})

$$N = size(U, 2)$$

$$linPt = [0]$$
for $i = 1 : N$

$$k = size(linPt)$$

$$\delta = \infty$$
for $j = 1 : k$

$$\delta' = \frac{\|U^{(i)} - U^{(linPt(j))}\|}{\|U^{(linPt(j))}\|}$$

$$\delta = min[\delta, \delta']$$
end
if $(\delta > \delta_{min})$

$$linPt = [linPt \ i]$$
end
end

With the set of linear models created, a TPWL scheme can be assembled in order to model nonlinearity. Consider a weighted combination of s linearized models of system (11)

$$\sum_{i=0}^{s-1} \tilde{\omega}_i(x) \Big\{ \dot{x}'_i(t) = A_i x'_i(t) + B_{1i} u(t) + B_{2i} \Big\}$$

$$\sum_{i=0}^{s-1} \tilde{\omega}_i(x) \Big\{ y_i(t) = C_i x'_i(t) + C_{0i} \Big\}$$
(15)

where $\tilde{\omega}_i(x)$ are weights depending on the value of the perturbation about the linearization point x_i . It is assumed that for all x, $\sum_{i=0}^{s-1} \tilde{\omega}_i(x) = 1$. The weights $\tilde{\omega}_i(x)$ are then obtained using the distance $||x(t) - x_i||$ between the linearization point x_i and the current solution x(t). The procedure below, following Rewienski [9], ensures that the "dominant" model *i* is that corresponding to the linearization point x_i that is the closest to the current state of the system:

Algorithm 2 (Weights computation)

1) For
$$i = 0, ..., (s - 1)$$
 compute:
 $d_i = ||x(t) - x_i||_2.$
2) $[m, k] = min\{d_i : i = 0, ..., (s - 1)\}.$
3) a) For $i = 0, ..., (s - 1)$ compute:
 $\tilde{\omega}_i = (exp(d_i)/m)^{-25}.$
or
b) For $i = 0, ..., (s - 1), \tilde{\omega}_i = 0$
 $\tilde{\omega}_k = 1.$

4) Normalize
$$\tilde{\omega}_i$$

First, Algorithm 2 obtains the difference d_i between the current state x(t) and the linearization point x_i . The minimum distance is given by m and corresponds to the model with index k. Then, the weights can be computed in two different ways. The first method shows a weighted sum strongly concentrated on the closest model, while the second uses only the closest model at the time. As will be presented later, each formulation yields slightly different results. The last step in the algorithm ensures that the summation of the s weights is unity.

E. Reduced-Order TPWL Model

Using the TPWL representation of the nonlinear system, an efficient reduced-order model can now be obtained using the projection (3) applied to (15), yielding a reduced-order TPWL model as follows.

$$\sum_{i=0}^{s-1} \tilde{\omega}_i(\hat{x}) \left\{ \frac{d}{dt} \hat{x}'_i(t) = \hat{A}_i \hat{x}'_i(t) + \hat{B}_{1i} u(t) + \hat{B}_{2i} \right\}$$
$$\sum_{i=0}^{s-1} \tilde{\omega}_i(\hat{x}) \left\{ \hat{y}_i(t) = \hat{C}_i \hat{x}'_{0i}(t) + C_{0i} \right\}$$
(16)

where the reduced-order matrices are defined as before in (13). As in the linear case, this representation is efficient, since all reduced-order matrices in (16) can be precomputed. Note also that the weights $\tilde{\omega}_i$ are computed as a function of the reduced-order state \hat{x} . The TPWL approach fits well within the context of POD-based model reduction, since a simulation of the nonlinear system can provide both the snapshots for computation of the POD basis vectors and also a set of instantaneous flow states from which to select the linearization points.

The final TPWL reduction approach can be summarized as follows. First, simulate the nonlinear CFD model for a range of forcing functions and conditions that are representative of the application at hand. Second, from the resulting snapshot collection, calculate a set of POD basis vectors. Third, from the same snapshot collection, select a set of linearization points using Algorithm 1. Fourth, using the dominant POD basis vectors, project each linearized model to obtain a set of reducedorder linear state-space systems. Finally, combine these loworder state-space systems using the TPWL representation and a set of weights from Algorithm 2.

This approach will now be demonstrated for the case of flow through the supersonic diffuser shown in Figure 2. Both fullorder and reduced-order TPWL models will be constructed, and the results compared with full nonlinear CFD outputs.

III. RESULTS

A number of test cases will be presented to demonstrate the TPWL methodology. In all cases, the input considered is an incoming density disturbance and the output of interest is the average Mach number at the throat of the diffuser. The six different temporal distributions considered for the input are presented in Figure 4, and vary temporally either with a Gaussian pulse or a sinusoidal distribution of various frequencies and amplitudes. The Gaussian distribution is described by

$$\rho'(t) = -\Lambda \rho_0 e^{-\alpha (t - t_{peak}/f_0)^2}$$
(17)

while the sinusoidal distribution is described by

$$\rho'(t) = -\Lambda \rho_0 \sin \omega_0 t \tag{18}$$

where the nominal frequency f_0 equals a_0/h , the inlet speed of sound divided by the height of the inlet, $\omega_0 = 2\pi f/f_0$, and the non-dimensional time, t_{peak} , sets the time at which the perturbation peaks. The parameter α sets the sharpness of the



Fig. 4. Incoming density disturbances. Top: Gaussian distributions. Bottom: sinusoidal distributions.

Case	Λ	t_{peak}	α
1	1%	20	$0.03f_0^2$
2	2%	20	$0.03f_0^2$
3	3%	20	$0.03f_0^2$

TABLE I Data used for the Gaussian distribution.

perturbation (i.e. its frequency content), Λ corresponds to the amplitude of the perturbation, and ρ_0 is the nominal value of freestream density. The parameter values corresponding to the six different input functions are presented in Tables I and II.

Nonlinear CFD results are obtained by simulation of the full system, and snapshots at each timestep are saved. Using Algorithm 1 for different values of δ_{min} and the snapshots just obtained, various sets of models are found. Table III shows the number of models as a function of the choice of δ_{min} for four of the cases, where Algorithm 1 was applied to each case separately. For each case, it can be seen by how much the number of models grows as the distance between linearization points is decreased. By comparing the number of models for a given δ_{min} , one gains some insight to the importance of nonlinearity in each case. For example, a Gaussian distribution of 3% can be seen to introduce more nonlinearity into the system than one of 1%, requiring substantially more models for a given δ_{min} .

A. Full-Order TPWL Models

Once the linearization points have been determined, the validity of the TPWL representation must be tested. This was

Case	Λ	ω_0
4	1.5%	0.65
5	2%	0.35
6	3%	0.65
TABLE II		

DATA USED FOR THE SINUSOIDAL DISTRIBUTION.

	Case number			
δ_{min}	1	2	3	6
∞	1	1	1	1
0.030	1	3	4	7
0.020	2	4	8	16
0.015	3	6	16	31
0.012	4	11	23	41
0.010	5	16	28	50
0.008	6	20	34	70
0.006	12	29	48	100
0.005	15	36	56	118
0.004	20	42	69	147
TABLE III				

Number of models given by different values of δ_{min} for four of the input cases.



Fig. 5. Nonlinear response plotted against various TPWL model combinations for a Gaussian incoming disturbance of 3% amplitude. Training trajectory obtained from the same simulation.

done by comparing nonlinear CFD results with those obtained using a full-order TPWL approximation as in Equation (15). The results using different sets of models from Table III are shown on Figure 5, where the average Mach number at the throat is plotted against time. Here, both the training trajectory and the disturbance were a Gaussian distribution of 3% amplitude. Figure 5 shows the number of models needed to accurately represent the nonlinear behavior. It can be seen that only one linearized model cannot capture the nonlinear behavior of a shock. As the value of δ_{min} is decreased, the match improves with increasing number of models. It can be seen in Figure 5 that with 28 models ($\delta_{min} = 0.01$), the nonlinear CFD results are matched by the combination of fullorder linear models.

Figures 6 and 7 show TPWL results for all of the Gaussian amplitudes, using values of δ_{min} equal to 0.01 and 0.005, respectively. For each case, the training trajectory corresponds to the desired incoming disturbance. Comparing these figures, one gains insight to the value of δ_{min} required in order to obtain a good match between the piecewise-linear combination of models and the nonlinear CFD. As Figure 5 shows, a



Fig. 6. Nonlinear response plotted against TPWL models for a Gaussian incoming disturbance of various amplitudes. From top, amplitude of 1%, 2% and 3%. The training trajectory for each case was the same as the simulation. Linearized models were selected using $\delta_{min} = 0.01$.



Fig. 7. Nonlinear response plotted against TPWL models for a Gaussian incoming disturbance of various amplitudes. From top, amplitude of 1%, 2% and 3%. The training trajectory for each cases was the same as the simulation. Linearized models were selected using $\delta_{min} = 0.005$.

minimum number of models is needed to capture a sufficiently high degree of nonlinearity. However, as Figure 7 demonstrates, taking too many models may cause undesirable results. In particular, oscillations may be observed or behavior may be inaccurately captured in sensitive regions. This is observed in the lower plot of Figure 7 at a time $t/T_0 \approx 50$ corresponding to the point at which the shock returns to a position downstream of the throat. These problems are further demonstrated in Figure 8, where even a small increase in the number of models leads to oscillations and inaccuracies in sensitive regions. Systematic strategies to avoid this behavior are the subject of ongoing research.

In the context of finding a reduced-order model that is valid over a range of flow conditions, the different input cases would not be considered separately. Rather, the snapshots from each would be combined to find a TPWL system that captures all



Fig. 8. Nonlinear response plotted against various TPWL model combinations for a Gaussian incoming disturbance of 3% amplitude. Training trajectory obtained from the same simulation.



Fig. 9. Nonlinear CFD response plotted against TPWL model with 74 linearization points chosen from three different training trajectories. From top: Gaussian amplitude of 1.5% and 2.5%. Training trajectories obtained from incoming density disturbances of Gaussian amplitudes of 1%, 2% and 3%.

training trajectories. To achieve this, all snapshots obtained from the three different training trajectories of 1%, 2% and 3% Gaussian disturbances were combined to form one large data set. Linearization points were then selected from the complete set using Algorithm 1 with $\delta_{min} = 0.005$, which resulted in the selection of 74 points. Results for simulations of this TPWL model for Gaussian amplitudes of 1.5% and 2.5% are shown on Figure 9. Note that these cases were not considered as part of the training trajectory set; however, they would be expected to fall within the range of validity of the existing ensemble. Very good agreement between the full nonlinear CFD and the set of combined models can be seen for the 1.5% case. For the larger amplitude 2.5% case, some discrepancy with the CFD is observed, but the TPWL approach shows a dramatic improvement over using a single linearized system.

δ_{min}	# of Models	
∞	1	
0.010	27	
0.008	38	
0.006	58	
0.005	74	
0.004	93	

TABLE IV

Number of models given by different values of δ_{min} applied to the three Gaussian distributions combined together.

B. Reduced-Order TPWL Models

The POD basis calculation and the selection of linearization points can be performed efficiently using the same ensemble of snapshots. It is important that this snapshot selection span all operating conditions of interest. For the results presented here, three training trajectories were used, which corresponded to the three Gaussian input pulses given in Table I. For each trajectory, 480 snapshots were collected corresponding to the solution at every timestep, yielding a total of 1440 snapshots. POD basis vectors were then calculated, resulting in the POD eigenvalue spectrum plotted in Figure 10. To capture 99%, 99.9%, 99.99% and 99.999% of the snapshot energy defined by (6), 9, 18, 33 and 61 basis vectors are required, respectively.



Fig. 10. First 50 POD eigenvalues of a total of 1440.

The second step in creating the reduced-order TPWL model is to determine appropriate linearization points using Algorithm 1. This algorithm is applied to the entire set of available training trajectories, i.e. the three Gaussian disturbances from Table I. Table IV presents the resulting number of models as a function of the criterion δ_{min} .

Using $\delta_{min} = 0.005$ and the corresponding 74 linearization points, a set of reduced-order models was created by projection of each linearized model onto the reduced space spanned by the first 50 POD basis vectors. If a sufficient number of POD basis vectors is used to define the reduced space, accurate reduced-order models can be obtained; however, it should be noted that there are no accuracy or stability guarantees associ-



Fig. 11. Transfer function from incoming density disturbance to average throat Mach number for model number 37 out of 74 ($\delta_{min} = 0.005$), with 50 states.

ated to these reduced models. The accuracy can be checked a posteriori by comparing the transfer functions of the full-order and reduced-order models at each linearization point. Figure 11 shows this comparison for the transfer function between an incoming density disturbance and the throat Mach number for one particular model. The reduced model uses 50 states, which corresponds to 99.998% POD energy. The frequency content of the input disturbances used for the training trajectories lies in the range $f/f_0 < 1.3$. The reduced-order model is expected to produce accurate results for behavior contained in the snapshot samples. It can be seen in Figure 11 that a good match is obtained for the frequencies included in the sampling process; however, the figure demonstrates the danger of applying the reduced models outside their range of validity and emphasizes the importance of selecting the training trajectories appropriately.

The 74 models were then combined to form a TPWL system as defined by (16). Simulation results are presented for the second weighting procedure given in Algorithm 2, which uses only the closest model. Figures 12 and 13 present the simulation results of this final model for a range of different incoming density disturbances. For the Gaussian pulses in Figure 12, it can be seen that good agreement is achieved for disturbances of smaller amplitude. For the lower two cases, with 2.5% and 3% amplitudes, some discrepancy is noted. In particular, the point at which the shock returns downstream of the throat continues to show extreme sensitivity. In Figure 13, a similar trend is observed. Results for the smaller amplitude sinusoids are excellent, but discrepancy is again observed for the 3% amplitude in the lower plot.

IV. CONCLUSION

The TPWL methodology has been demonstrated as a viable approach for obtaining accurate, efficient, reduced-order models for nonlinear CFD applications. For many of the cases considered, the performance of the method is excellent. For



Fig. 12. Reduced-order TPWL simulation for Gaussian inputs. From top: Gaussian distribution of 1%, 1.5%, 2%, 2.5%, and 3% amplitude.



Fig. 13. Reduced-order TPWL simulation for sinusoidal inputs. From top: Case 4, Case 5, Case 6.

cases where the nonlinearity is significant and leads to very large shock motion, the results are found to be sensitive to the choices of linearization points and model size. Systematic strategies to reduce this sensitivity are the subject of ongoing research.

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