

# Magnetic Materials

- The inductor



$$\nabla \times E = -\frac{1}{c} \frac{\partial B}{\partial t} \text{ (CGS)}$$

$$\iint \nabla \times E dS = -\frac{1}{c} \frac{\partial}{\partial t} \left( \iint B dS \right) = -\frac{1}{c} \frac{\partial \Phi_B}{\partial t}$$

$\Phi_B \equiv$  magnetic flux density

$$\iint \nabla \times E dS = \oint E \cdot d\ell \text{ (Green's Theorem)}$$

$$V = \int E \cdot d\ell = -\frac{1}{c} \frac{\partial \Phi_B}{\partial t} \text{ (explicit Faraday's Law)}$$

$$\Phi_B = LI \text{ (Q = CV)}$$

$$\frac{\partial \Phi_B}{\partial t} = L \frac{\partial I}{\partial t}$$

$$V_{EMF} = -\frac{\partial N\Phi_B}{\partial t} = -L \frac{\partial I}{\partial t}$$

$$V = L \frac{\partial I}{\partial t} \text{ (recall } I = C \frac{\partial V}{\partial t} \text{ for the capacitor)}$$

$$\text{Power} = VI = LI \frac{\partial I}{\partial t}$$

$$\text{Energy} = \int \text{Power} \cdot dt = \int LI dI = \frac{1}{2} LI^2 = \frac{1}{2} N\Phi_B I$$

$$\left( \text{capacitor } \frac{1}{2} CV^2 \right)$$



# The Inductor

$$\nabla \times B = \frac{4\pi}{c} J + \frac{1}{c} \frac{\partial E}{\partial t}$$

$$\iint \nabla \times B dS = \oint B \cdot d\ell = \frac{4\pi}{c} \iint J \cdot dS = \frac{4\pi}{c} I$$

$$B = \frac{4\pi}{c} In$$

$$N = n \cdot \text{length} = nl$$

$$L = \frac{N\Phi_B}{I} = \frac{N(BA)}{I} = \frac{4\pi}{c} n^2 lA$$



## Magnetic Permeability and Susceptibility

Insert magnetic material

Magnetic dipoles in material can line-up in magnetic field

$$B = H + 4\pi\chi H = H + 4\pi M$$

$$M = \chi H \quad \frac{\partial M}{\partial H} = \chi \quad \mu = 1 + 4\pi\chi$$

$$B = 4\pi M + H \quad B = \mu H$$

B magnetic induction  
 $\chi$  magnetic susceptibility  
 H magnetic field strength (applied field)  
 M magnetization

MKS:

$$B = \mu_0(H+M) = \mu_0 nI + \mu_0 M$$

$$\mu_r = B / \mu_0 H = 1 + (M/H) = 1 + \chi_m$$



## Maxwell and Magnetic Materials

- Ampere's law  $\oint H \cdot d\ell = I = 0$
- For a permanent magnet, there is no real current flow; if we use B, there is a need for a fictitious current (magnetization current)
- Magnetic material inserted inside inductor increases inductance

$$\Phi_B = BA \sim 4\pi MA = 4\pi\chi HA = 4\pi\chi \left(\frac{4\pi}{c} In\right) A$$

$$L = \frac{N\Phi_B}{I} = \frac{(4\pi)^2}{c} n^2 l A \chi$$

L increased by  $\sim\chi$  due to magnetic material

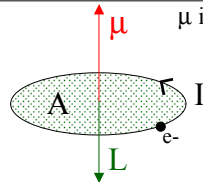
Material Type	$\chi$
Paramagnetic	$+10^{-5} - 10^{-4}$
Diamagnetic	$-10^{-8} - 10^{-5}$
Ferromagnetic	$+10^5$



# Microscopic Source of Magnetization

- No monopoles
- magnetic dipole comes from moving or spinning electrons

## Orbital Angular Momentum



$\mu$  is the magnetic dipole moment

$$Energy = E = -\vec{\mu} \cdot \vec{H} = -|\mu||H|\cos\theta$$

What is  $\mu$ ? For  $\theta=0$ ,  $E = -\mu H \approx -\Phi_B I$  since energy  $\sim LI^2$  and for 1 loop  $L = \frac{\Phi_B}{I}$

$$\Phi_B = \iint H \cdot dS \sim HA$$

$$\therefore \mu H = \Phi_B I = HAI \quad \text{and} \quad \boxed{\mu = IA}$$

$$I = -\frac{e}{c} \frac{\omega}{2\pi} \quad A = \pi r^2$$

$$\mu = -\frac{e}{2c} \omega r^2$$



3.225

© E. Fitzgerald-1999

5

# Microscopic Source of Magnetization

- Classical mechanics gives orbital angular momentum as:

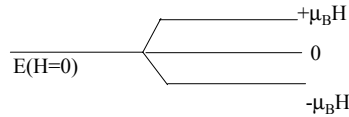
$$\vec{L} = \vec{r} \times \vec{p} = mr^2\omega$$

$$\mu_L = -\frac{e}{2mc} L_{QM} = -\frac{e\hbar}{2mc} L_Z = -\mu_B L_Z$$

Example for  $l=1$ :

$$\left( \mu_B = \frac{e\hbar}{2mc} \right)$$

$$L_Z = m_l \hbar = -\hbar, 0, \hbar$$

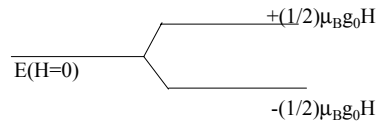


## Spin Moment



$$\mu_s = -\frac{e}{mc} S^{QM} = -g_0 \frac{e\hbar}{2mc} S_z = -g_0 \mu_B S_z$$

$$S_z = m_s \hbar = \pm \frac{1}{2} \hbar \quad g_0 = 2 \quad \text{for electron spin}$$



3.225

© E. Fitzgerald-1999

6

# Exchange

$$E \sim -JS_1S_2$$

J negative,  $E \sim +S_1S_2 \rightarrow$  Energy  $\downarrow$  if  $\downarrow \uparrow$

J positive,  $E \sim -S_1S_2 \rightarrow$  Energy  $\downarrow$  if  $\uparrow \uparrow$

Fe, Ni, Co  $\rightarrow$  J positive!  
Other elements J is negative

Rule of Thumb:

$$\frac{r}{2r_a} \equiv \frac{\text{interatomic distance}}{2(\text{atomic radius})} > 1.5$$

J is a function of distance!

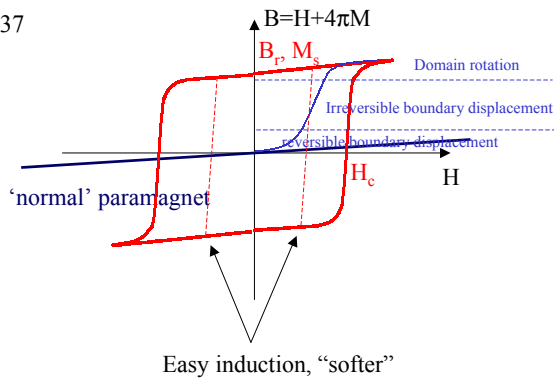
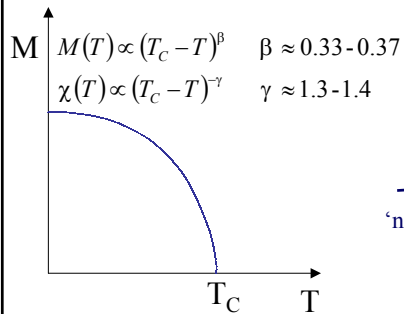


3.225

© E. Fitzgerald-1999

7

# Ferromagnetism



Magnetic anisotropy  
hardness of loop dependent on crystal direction  
comes from spin interacting with bonding

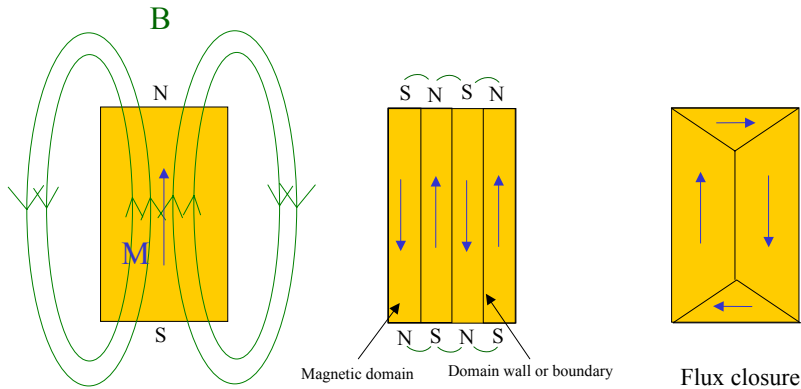


3.225

© E. Fitzgerald-1999

8

# Domains in Ferromagnetic Materials



Magnetic energy

$$= \frac{1}{8} \int B^2 dV$$

Flux closure  
No external field

