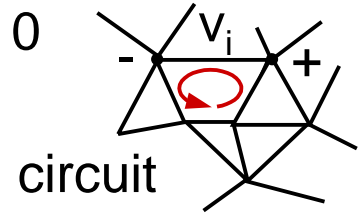


KIRCHOFF'S VOLTAGE LAW

Kirchoff's Voltage Law: (KVL)

Around any loop: $\sum_i V_i = 0$



Faraday's Law:

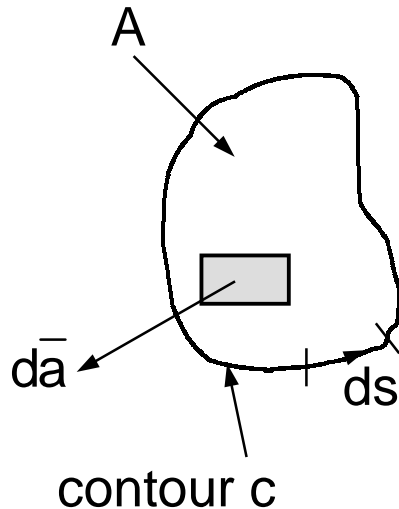
$\nabla \times \bar{E} = -\partial \bar{B} / \partial t$ [differential form]

Integral form is:

$$\oint_C \bar{E} \cdot d\bar{s} = -\frac{d}{dt} \int_A \bar{B} \cdot d\bar{a}$$

[Recall Stoke's Theorem: $\int_A (\nabla \times \bar{G}) \cdot \hat{n} da = \oint_C \bar{G} \cdot d\bar{s}$]

If $d\bar{B}/dt = 0$ in A , then Kirchoff's voltage law must be satisfied

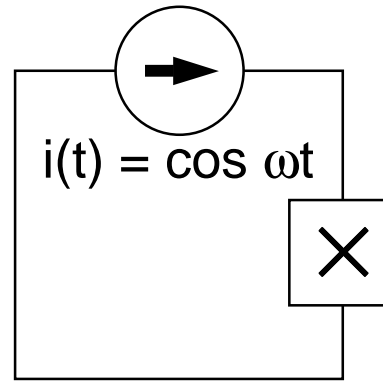


More generally:

KVL assumes all magnetic energy is stored inside circuit elements

KIRCHOFF'S VOLTAGE LAW (2)

Undefined Circuit:

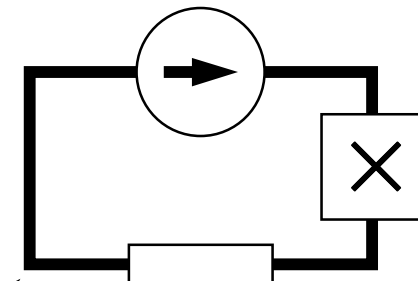


Current in loop varies, so $\frac{d}{dt} \int_A \bar{B} \cdot d\bar{a} \neq 0 \Rightarrow$ voltage!

Defined Circuit:

We assume voltage drops occur only across elements, and ignore H fields generated by currents through them.

We use small loops, thick wires, and high-impedance lumped elements



Lumped inductor L

“parasitic inductance”

KIRCHOFF'S CURRENT LAW (KCL)

**Kirchoff's Current Law:
(KCL)**

$$\sum_i I_i = 0 = \text{Total current into any node}$$

Ampere's Law:

$$\nabla \times \bar{H} = \bar{J} + \partial \bar{D} / \partial t \quad [\text{differential form}]$$

$$\nabla \cdot (\nabla \times \bar{H}) = \nabla \cdot (\bar{J} + \partial \bar{D} / \partial t) = 0$$

$$= \nabla \cdot \bar{J} + \partial (\nabla \cdot \bar{D}) / \partial t$$

$$= \nabla \cdot \bar{J} + \partial \rho / \partial t$$

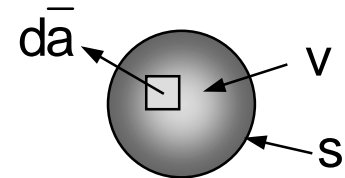
Conservation of Charge:

$$\nabla \cdot \bar{J} = -\partial \rho / \partial t = 0 \text{ at each node}$$

$$[\text{Recall Gauss's Divergence Theorem: } \int_V (\nabla \cdot \bar{G}) dv = \int_S \bar{G} \cdot d\bar{a}]$$

Current Law follows from:

$$\int_S \bar{J} \cdot d\bar{a} = \sum_i I_i = 0$$



More generally:

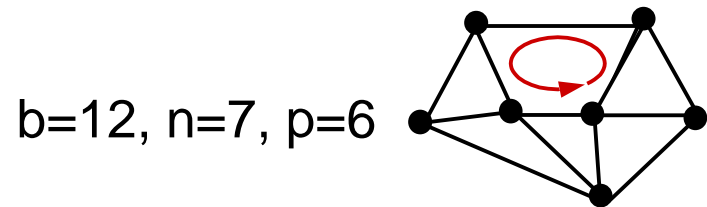
KCL assumes nodes and wires store no charge ($\partial \rho / \partial t = 0$), or that all electric energy is stored inside circuit elements

“parasitic capacitance”

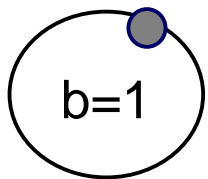
SOLVING CIRCUITS

Generic Circuit Topology:

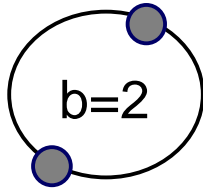
Assume b branches, n nodes, and p unique loops



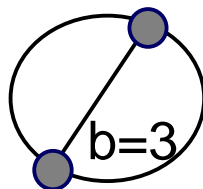
We can show $b = n + p - 1$.



$$n = p = 1$$



$$n = 2$$



$$n = p = 2$$

Every time we add a branch, the number of nodes or loops (meshes) increases by one.

Voltage sources are branches characterized by $v = V_o$
Current sources are branches characterized by $i = I_o$

SOLVING CIRCUITS (2)

Number of Unknowns: Number of unknowns is $2b$ (v, i for each branch)

Number of Equations: We have one device equation for each branch relating or specifying v and i

We also have $n - 1$ independent node equations (KCL) and p loop equations (KVL), for a total of:

$$b + (n - 1) + p = 2b \text{ equations} = \text{number of unknowns}$$

(Recall $b = n + p - 1$)

Given initial conditions and all sources we can solve the equations analytically for simple circuits, or by simulation for any circuit.

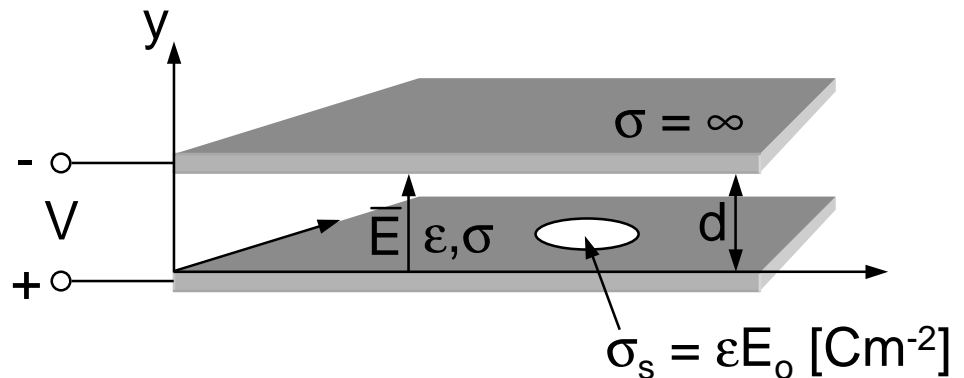
(Non-unique exceptions: indeterminate flip-flop or chaotic circuits)

SIMPLE CIRCUIT ELEMENTS, R AND C

Electrostatics:

In general:	Everywhere	$\nabla \times \bar{E} = 0, \quad \nabla \cdot \bar{E} = \rho/\epsilon$
	At conductor	$\hat{n} \cdot \bar{E} = \sigma_S/\epsilon, \quad \bar{E}_{//} = 0$
	Between conductors,	$\nabla \cdot \bar{E} = \rho/\epsilon = 0$

Parallel-Plate Devices:

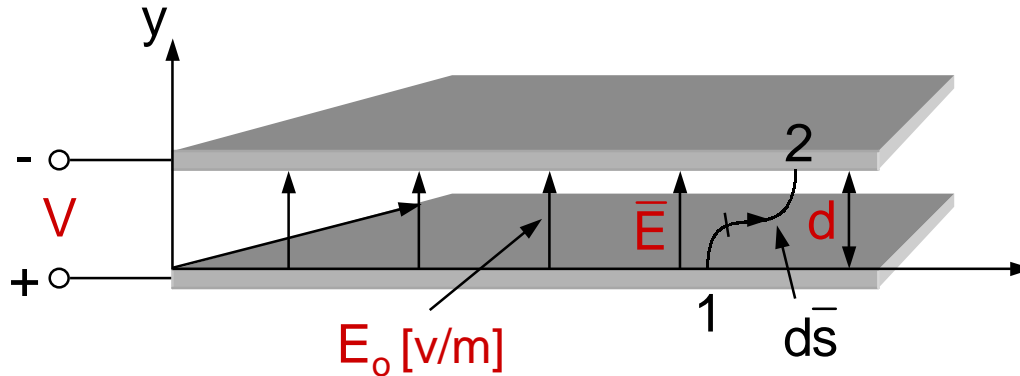


Field Solution:

$\bar{E} = \hat{y}E_0$ for rectangular geometry here

$$\sigma_S = \epsilon E_0$$

SIMPLE CAPACITOR



Relating Fields to Potentials:

Since: $\nabla \times \bar{E} = 0$

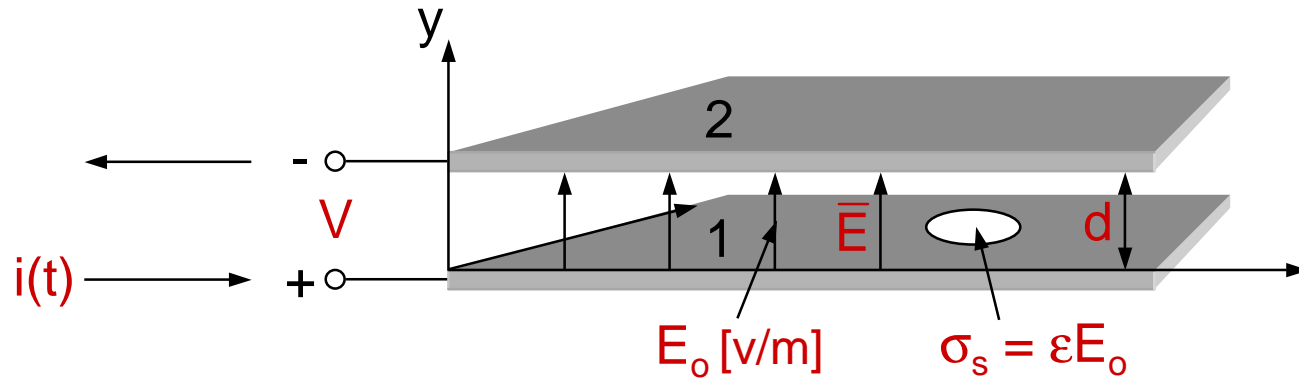
We define $\bar{E} = -\nabla\Phi$ $\Phi = \int_1^2 \bar{E} \cdot d\bar{s}$

Where Φ is the electrostatic potential of [1] relative to [2]

Therefore here: $E_o d = V$, $E_o = V/d$

In general, absolute potential $\Phi = 0$ at infinity, by definition

SIMPLE CAPACITOR (2)



Capacitor Charge Q:

We define:

$$Q = \int_A \sigma_s da = A\epsilon E_0 = A\epsilon V/d$$

We also define:

$$Q = CV$$

Capacitance C:

Therefore:

$$C = A\epsilon/d$$

We know:

$$q(t) = \int_{-\infty}^t i(t) dt = Cv(t)$$

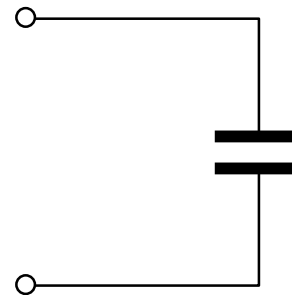
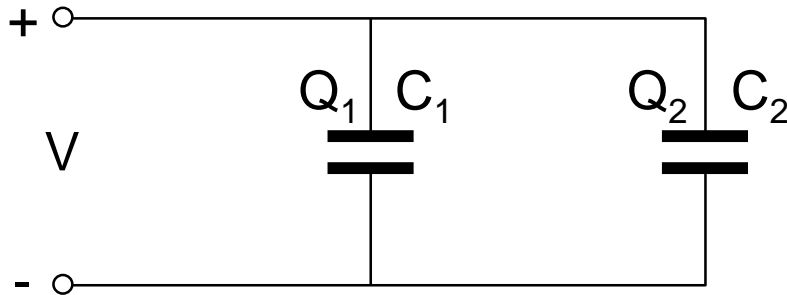
Therefore:

$$v(t) = (1/C) \int_{-\infty}^t i(t) dt$$

Also,

$$i = C dv/dt$$

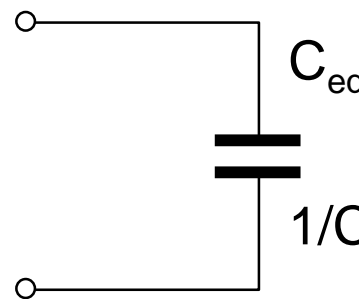
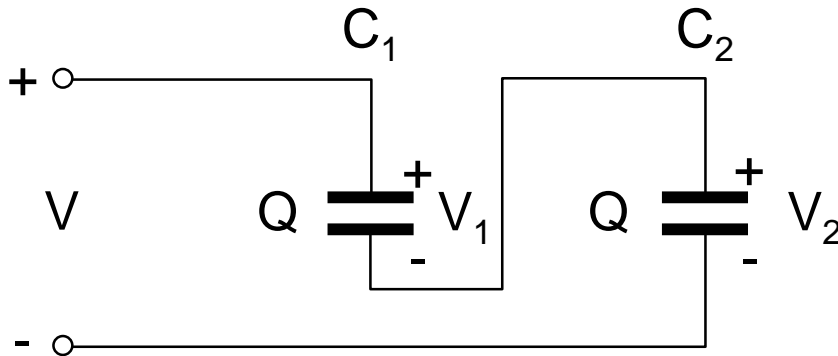
CAPACITORS IN SERIES AND PARALLEL



$$C_{eq} = Q_{eq}/V = (Q_1 + Q_2)/V = C_1 + C_2$$

Capacitors in Parallel:

$$C_{eq} = C_1 + C_2$$



$$1/C_{eq} = V/Q = (V_1 + V_2)/Q$$

Capacitor in Series:

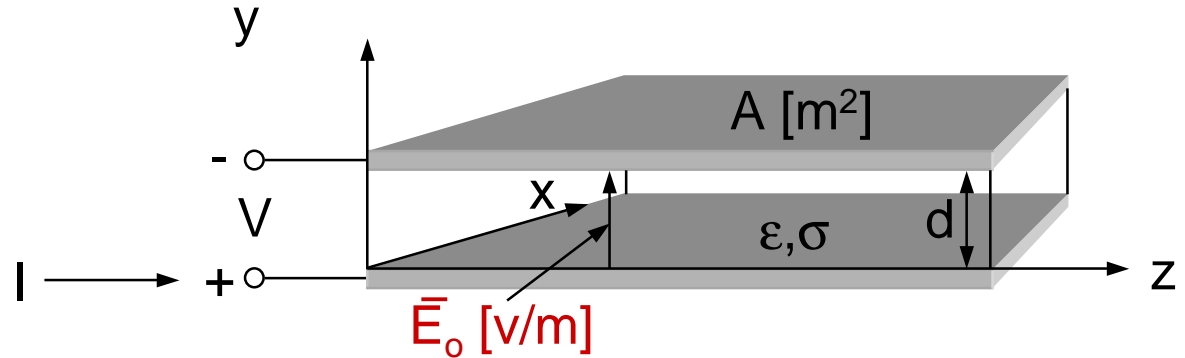
$$1/C_{eq} = 1/C_1 + 1/C_2$$

SIMPLE RESISTORS

Conductance σ :

$$\bar{J} = \sigma \bar{E} \quad [\text{am}^{-2}]$$

$$I = AJ = A\sigma E_0 = A\sigma V/d$$



Resistance R:

$$R = V/I = d/A\sigma$$

Resistors in Series:

$$R_{\text{eq}} = R_1 + R_2$$

$$R_{\text{eq}} = (V_1 + V_2)/I$$

Resistors in Parallel:

$$1/R_{\text{eq}} = 1/R_1 + 1/R_2$$

$$1/R_{\text{eq}} = (I_1 + I_2)/V$$

CHARGE RELAXATION

RC Circuits:

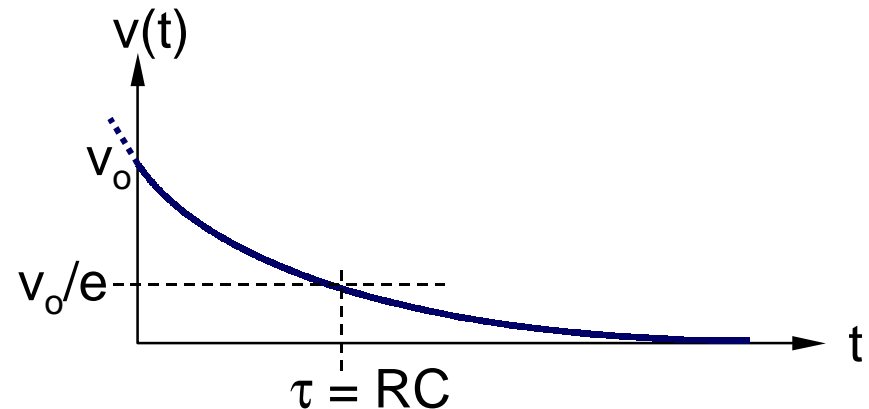
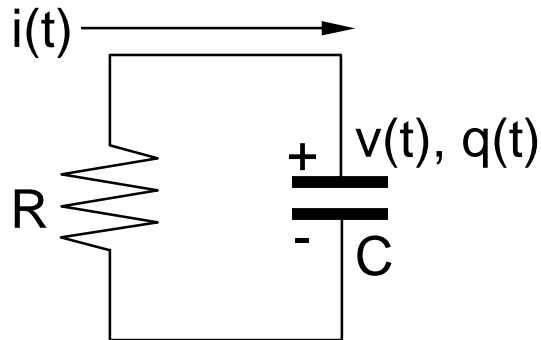
$$q(t) = Cv(t) = \int_0^t i(t)dt + q_0$$

$$v(t) = -i(t)R = -RC dv(t)/dt$$

Try:

$$v(t) = v_0 e^{-t/\tau} \quad [v_0 \text{ is initial condition}]$$

It works if $\tau = RC$: $v_0 e^{-t/\tau} = (-RC)(-1/\tau)v_0 e^{-t/\tau}$



Dielectric Relaxation:

$$R = d/A\sigma, \quad C = \epsilon A/d$$

$$\tau = RC = \epsilon/\sigma \text{ seconds "Relaxation time constant"}$$

independent
of geometry