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# Viability of $\Delta m^2 \sim 1 \text{ eV}^2$ sterile neutrino mixing models in light of MiniBooNE electron neutrino and antineutrino data from the Booster and NuMI beamlines

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This paper examines sterile neutrino oscillation models in light of recently published results from the MiniBooNE Experiment. The new MiniBooNE data include the updated neutrino results, including the low-energy region, and the first antineutrino results, as well as first results from the off-axis NuMI beam observed in the MiniBooNE detector. These new global fits also include data from LSND, KARMEN, NOMAD, Bugey, CHOOZ, CCFR84, and CDHS. Constraints from atmospheric oscillation data have been imposed. We test the validity of the three-active plus one-sterile (3 + 1) and two-sterile (3 + 2) oscillation hypotheses, and we estimate the allowed range of fundamental neutrino oscillation parameters in each case. We assume *CPT*-invariance throughout. However, in the case of (3 + 2) oscillations, *CP* violation is allowed. We find that, with the addition of the new MiniBooNE data sets, a (3 + 2) oscillation hypothesis provides only a marginally better description of all short-baseline data over a (3 + 1) oscillation hypothesis. In the case of (3 + 2) *CP*-violating models, we obtain good  $\chi^2$ -probabilities in general due to the large number of fit parameters. However, we find large incompatibilities among appearance and disappearance experiments, consistent with previous analyses. Aside from LSND, the data sets responsible for this tension are the MiniBooNE neutrino data set, CDHS, and the atmospheric constraints. In addition, new incompatibilities are found between the appearance experiments themselves (MiniBooNE, LSND, KARMEN and NOMAD), independent of *CP*-violation assumptions. On the other hand, fits to antineutrino-only data sets, including appearance and disappearance experiments, are found significantly more compatible, even within a (3 + 1) oscillation scenario.

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## I. INTRODUCTION

Sterile neutrino oscillation models were proposed more than a decade ago as an explanation for the LSND anomaly [1–5], an excess of events consistent with  $\bar{\nu}_\mu \rightarrow \bar{\nu}_e$  oscillations at high  $\Delta m^2$ . These models relate  $\nu_e$  appearance ( $\nu_\mu \rightarrow \nu_e$ ) with  $\nu_\mu$  and  $\nu_e$  disappearance ( $\nu_\mu \rightarrow \nu_\mu$  and  $\nu_e \rightarrow \nu_e$ ), motivating combined fits in all three oscillation channels. Relatively early in the discussion of models, it was demonstrated [5,6] that a three-active plus one-sterile (3 + 1) neutrino oscillation model could not reconcile the LSND result with existing null results from other short-baseline (SBL) experiments, including KARMEN [7], NOMAD [8], Bugey [9], CHOOZ [10], CCFR84 [11], and CDHS [12], which had similar high  $\Delta m^2$  sensitivity. However, it was shown that a three-active plus two-sterile neutrino (3 + 2) oscillation scenario provided a better description of these data sets [5].

In 2001, the MiniBooNE experiment began running with the goal to test the LSND result using both neutrino and antineutrino beams. This is a short-baseline appearance and disappearance experiment located at Fermi National Accelerator Laboratory (Fermilab). MiniBooNE's first results, reported in 2007, described a search for  $\nu_\mu \rightarrow \nu_e$  oscillations [13]. These results were incompatible with a simple two-neutrino oscillation interpretation of the LSND signal and, within this model, MiniBooNE excluded the LSND result at the 98% CL. However, this same analysis reported a  $3.7\sigma$  excess of electron neutrino candidate events at low energies, between 300–475 MeV, which remains unexplained. Reference [14] included the MiniBooNE first result in a global fit to all SBL experiments under the (3 + 1) and (3 + 2) oscillation scenarios. The analysis built on an earlier study, which introduced the possibility of *CP* violation ( $P(\nu_\mu \rightarrow \nu_e) \neq P(\bar{\nu}_\mu \rightarrow \bar{\nu}_e)$ ) within (3 + 2) fits [15]. Including the first MiniBooNE results into the global fit led to two observations in Ref. [14]: 1) MiniBooNE, LSND, and the null appearance experiments (KARMEN and NOMAD) are compatible under a (3 + 2) sterile neutrino oscillation scenario with large *CP* violation. 2) There is severe tension between appearance and disappearance experiments, at a level of

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more than  $3\sigma$ . In this paper we will consider both observations in light of new appearance data. Also, we will show that the incompatibility between appearance and disappearance experiments arises mainly from two  $\nu_\mu$  disappearance data sets: CDHS and atmospheric constraints.

Motivated by three new results from MiniBooNE, this paper reexamines the  $(3 + 1)$  and  $(3 + 2)$  global fits to the SBL data. These new results are: 1) an updated  $\nu_\mu \rightarrow \nu_e$  result [16]; 2) first results for a  $\bar{\nu}_\mu \rightarrow \bar{\nu}_e$  search [17]; and 3) first  $\nu_\mu \rightarrow \nu_e$  results from the NuMI off-axis beam at MiniBooNE [18]. We consider these new results in combination with seven SBL data sets. These provide constraints on:  $\nu_\mu$  disappearance (from the CCFR84 and CDHS experiments),  $\bar{\nu}_e$  disappearance (from the Bugey and CHOOZ experiments),  $\nu_\mu \rightarrow \nu_e$  oscillations (from the NOMAD experiment), and  $\bar{\nu}_\mu \rightarrow \bar{\nu}_e$  oscillations (from the LSND and KARMEN experiments). Furthermore, we have taken into account atmospheric constraints based on the analysis of Ref. [19]. These constraints have been incorporated in our analysis following the method described in Ref. [15], and are included in fits to all SBL experiments, null SBL experiments, or as explicitly stated. Table I summarizes all SBL data sets used in the fits presented in this paper.

In this work, we do not discuss experimental constraints on sterile neutrino models other than SBL and atmospheric neutrino ones. Constraints from the measurement of the electron spectrum near the endpoint in beta-decay experiments are relatively weak as long as the mostly-sterile mass states are heavier than the mostly-active ones, because of the small electron flavor of the former (see Refs. [5,20]). We make this assumption throughout the paper, by requiring that the heavier sterile neutrino mass eigenstates,  $m_5$  and  $m_4$ , obey  $m_5 > m_4 > m_1$ . Constraints on the energy density (and mass) in the Universe carried by sterile neutrinos from cosmic microwave background, matter power spectrum, and supernovae data have been studied in Ref. [21]. While relevant, these constraints are found to

be weaker than SBL ones, since sterile neutrino states do not necessarily feature thermal abundances at decoupling. Constraints on the number of relativistic degrees of freedom from the observations of cosmological abundances of light elements produced at the epoch of big bang nucleosynthesis may also be relevant, although model-dependent. For such a study involving one sterile neutrino species participating in the mixing, see, for example, Ref. [22].

The paper is organized as follows. In Sec. II, we provide a short description of the MiniBooNE experiment and the new data sets. In Sec. III, we specify the formalism used in this analysis to describe  $(3 + n)$  oscillations, where  $n$  is the number of sterile neutrinos. In Sec. IV, we discuss the analysis method followed, and describe in detail the way in which the three MiniBooNE data sets have been incorporated. In Sec. V, we present the results obtained for the  $(3 + 1)$  ( $CP$ -conserving only), and  $(3 + 2)$   $CP$ -conserving and  $CP$ -violating hypotheses. For each hypothesis, we quote the compatibility between various sets of SBL experiments and report the best-fit neutrino mass and mixing parameters derived from the combined analysis of all experimental data sets. In the  $(3 + 2)$   $CP$ -violating case, we discuss the constraints on the  $CP$  violation phase, inferred from a combined analysis of all SBL oscillation results. Finally, in Sec. VI, we discuss constraints to the  $(3 + 2)$   $CP$ -violating models from each of the SBL experiments considered in this analysis. The goal of this particular study was to investigate whether the source of tension between appearance and disappearance experiments [14] is a result of a single experiment, other than LSND.

## II. THE NEW MINIBOONE DATA SETS

The MiniBooNE experiment uses a muon neutrino beam produced by 8 GeV protons from the Fermilab Booster Neutrino Beamline (BNB) impinging on a beryllium target. The target is located within a magnetic focusing horn [23]. The current of the horn can be reversed for running neutrinos or antineutrinos, allowing MiniBooNE to perform both neutrino and antineutrino oscillation searches. The detector [24] is located  $L = 541$  m from the primary target, and the neutrino flux has an average energy of  $\sim 0.75$  GeV. This design maintains the LSND  $L/E$  of  $\sim 1$  m/MeV. The detector consists of a spherical tank with a 610-cm active radius, instrumented with 1520 8-inch photomultipliers. This is filled with 800 tons of pure mineral oil. An outer veto region rejects cosmic rays and neutrino events producing particles which cross the detector boundaries.

The MiniBooNE neutrino data set used in this analysis corresponds to the updated results recently reported by the MiniBooNE collaboration [16]. Compared to the first MiniBooNE result which was released in 2007 [13], the new result involves a reanalysis of the MiniBooNE low-energy excess events and several updates to the Monte Carlo prediction. These updates include a new

TABLE I. Short-baseline oscillation data sets considered in this paper, and oscillation channel that each data set constrains.

Data set	Channel
<i>Appearance experiments:</i>	
LSND	$\bar{\nu}_\mu \rightarrow \bar{\nu}_e$
BNB-MB( $\nu$ )	$\nu_\mu \rightarrow \nu_e$
BNB-MB( $\bar{\nu}$ )	$\bar{\nu}_\mu \rightarrow \bar{\nu}_e$
NUMI-MB	$\nu_\mu \rightarrow \nu_e$
KARMEN	$\bar{\nu}_\mu \rightarrow \bar{\nu}_e$
NOMAD	$\nu_\mu \rightarrow \nu_e$
<i>Disappearance experiments:</i>	
Bugey	$\bar{\nu}_e \rightarrow \bar{\nu}_e$
CHOOZ	$\bar{\nu}_e \rightarrow \bar{\nu}_e$
CCFR84	$\nu_\mu \rightarrow \nu_\mu$
CDHS	$\nu_\mu \rightarrow \nu_\mu$

model of photonuclear effects, incorporation of new data on  $\pi^0$  production and a better treatment of pion reinteraction in the detector following decay, an improved estimate and rejection method of the background from interactions outside the detector, and improvements to the determination of systematic errors. The updated low-energy analysis has resulted in a reduction to the significance of the excess from  $3.7\sigma$  in the original analysis to  $3.4\sigma$ , along with some slight modification to the shape of the energy spectrum; specifically, the peak of the excess has shifted slightly to higher neutrino energies. In addition, the new analysis extends in energy down to 200 MeV, compared to 300 MeV in the original analysis, which offers additional  $L/E$  information. The new result also corresponds to modestly higher statistics, corresponding to the total data collected during the experiment's neutrino running of  $6.46 \times 10^{20}$  protons on target (POT), compared to  $5.58 \times 10^{20}$ , previously.

More recently, the MiniBooNE Collaboration reported its first results from a search for  $\bar{\nu}_\mu \rightarrow \bar{\nu}_e$  oscillations, using a muon antineutrino beam [17]. The antineutrino analysis performed by MiniBooNE mirrors the updated neutrino analysis [16], and includes the Monte Carlo prediction updates of the latter. The total antineutrino data set used in the analysis corresponds to  $3.39 \times 10^{20}$  POT. However, due to meson production and cross-section effects, the antineutrino event sample, unlike the neutrino event sample, is statistically limited. Unlike the neutrino search, the MiniBooNE antineutrino search provides a direct test of the LSND result, similar to the search performed by KARMEN. The MiniBooNE sensitivity to  $\bar{\nu}_\mu \rightarrow \bar{\nu}_e$  extends into the low- $\Delta m^2$  region allowed by a combined analysis of KARMEN and LSND data. Nevertheless, the MiniBooNE antineutrino search has observed no conclusive signal, and a limit has been set, which is considerably weaker than the sensitivity, and comparable to the KARMEN limit. The limit degradation with respect to the sensitivity is due to a  $2.8\sigma$  fluctuation of data above expected background observed in the 475–675 MeV energy region. Thus, at present, the MiniBooNE antineutrino result is inconclusive with respect to oscillations allowed by LSND. However, MiniBooNE is in the process of collecting more antineutrino data. This is expected to improve the experiment's sensitivity to  $\bar{\nu}_\mu \rightarrow \bar{\nu}_e$  oscillations. Updated results are expected after about three years of running.

The third new data set [18] arises from the fact that the MiniBooNE detector is illuminated by the off-axis (110 mrad) neutrino flux from the NuMI beamline at Fermilab. This analysis has reported a  $1.2\sigma$  excess of  $\nu_e$ -like events in the neutrino energy range below 900 MeV. The NuMI data set corresponds to a mean  $L/E$  that is approximately the same as those of the MiniBooNE and LSND data sets, and therefore probes the same  $\Delta m^2$  range, providing complementary information in oscillation fits with MiniBooNE and LSND.

### III. (3 + n) STERILE NEUTRINO OSCILLATION FORMALISM

The formalism used in this paper follows that which was presented in Ref. [15]. We provide a brief summary here.

In sterile neutrino oscillation models, under the assumptions of  $CPT$  invariance and negligible matter effects, the probability for a neutrino produced with flavor  $\alpha$  and energy  $E$ , to be detected as a neutrino of flavor  $\beta$  after traveling a distance  $L$ , is given by [25,26]:

$$P(\nu_\alpha \rightarrow \nu_\beta) = \delta_{\alpha\beta} - 4 \sum_{i>j} \mathcal{R}(U_{\alpha i}^* U_{\beta i} U_{\alpha j} U_{\beta j}^*) \sin^2 x_{ij} + 2 \sum_{i>j} \mathcal{I}(U_{\alpha i}^* U_{\beta i} U_{\alpha j} U_{\beta j}^*) \sin 2x_{ij}, \quad (1)$$

where  $\mathcal{R}$  and  $\mathcal{I}$  indicate the real and imaginary parts of the product of mixing matrix elements, respectively;  $\alpha, \beta \equiv e, \mu, \tau$ , or  $s$ , ( $s$  being the sterile flavor);  $i, j = 1, \dots, 3 + n$  ( $n$  being the number of sterile neutrino species); and  $x_{ij} \equiv 1.27 \Delta m_{ij}^2 L/E$ . In defining  $x_{ij}$ , we take the neutrino mass splitting  $\Delta m_{ij}^2 \equiv m_i^2 - m_j^2$  in  $\text{eV}^2$ , the neutrino baseline  $L$  in km, and the neutrino energy  $E$  in GeV. For antineutrinos, the oscillation probability is obtained from Eq. (1) by replacing the mixing matrix  $U$  with its complex-conjugate matrix. Therefore, if the mixing matrix is not real, neutrino and antineutrino oscillation probabilities can differ.

For  $3 + n$  neutrino species, there are, in general,  $2 + n$  independent mass splittings,  $(3 + n)(2 + n)/2$  independent moduli of parameters in the unitary mixing matrix, and  $(2 + n)(1 + n)/2$  Dirac  $CP$ -violating phases that may be observed in oscillations. In SBL neutrino experiments that are sensitive only to  $\nu_\mu \rightarrow \nu_\mu$ ,  $\nu_e \rightarrow \nu_e$ , and  $\nu_\mu \rightarrow \nu_e$  transitions, the set of observable parameters is reduced considerably. In this case, the number of observable parameters is restricted to  $n$  independent mass splittings,  $2n$  moduli of mixing matrix parameters, and  $n - 1$   $CP$ -violating phases. Therefore, for  $(3 + 2)$  sterile neutrino models ( $n = 2$  case), for example, there are two independent mass splittings,  $\Delta m_{41}^2$  and  $\Delta m_{51}^2$ , both defined to be greater than zero, four moduli of mixing matrix parameters  $|U_{e4}|$ ,  $|U_{\mu 4}|$ ,  $|U_{e5}|$ ,  $|U_{\mu 5}|$ , and one  $CP$ -violating phase. The convention used for the  $CP$ -phase is:

$$\phi_{45} = \arg(U_{\mu 5}^* U_{e5} U_{\mu 4} U_{e4}^*). \quad (2)$$

In that case, the general oscillation formula in Eq. (1) becomes:

$$P(\nu_\alpha \rightarrow \nu_\alpha) = 1 - 4[(1 - |U_{\alpha 4}|^2 - |U_{\alpha 5}|^2) \cdot (|U_{\alpha 4}|^2 \sin^2 x_{41} + |U_{\alpha 5}|^2 \sin^2 x_{51}) + |U_{\alpha 4}|^2 |U_{\alpha 5}|^2 \sin^2 x_{54}] \quad (3)$$

and

$$\begin{aligned}
P(\nu_\alpha \rightarrow \nu_{\beta \neq \alpha}) &= 4|U_{\alpha 4}|^2|U_{\beta 4}|^2 \sin^2 x_{41} \\
&+ 4|U_{\alpha 5}|^2|U_{\beta 5}|^2 \sin^2 x_{51} \\
&+ 8|U_{\alpha 5}||U_{\beta 5}||U_{\alpha 4}||U_{\beta 4}| \\
&\times \sin x_{41} \sin x_{51} \cos(x_{54} + \phi_{45}). \quad (4)
\end{aligned}$$

The formulas for antineutrino oscillations are obtained by substituting  $\phi_{45} \rightarrow -\phi_{45}$ .

For the case of  $(3 + 1)$  sterile neutrino models ( $n = 1$  case), the corresponding oscillation probabilities are obtained from Eqs. (3) and (4) by setting  $x_{51} = x_{54} = 0$  and  $|U_{\alpha 5}| = 0$ . Note that, under the above assumptions, no  $CP$  violation is allowed for  $(3 + 1)$  models.

#### IV. ANALYSIS METHOD

In this section, we first provide an overview of the fitting technique. We then focus on the method followed for including the MiniBooNE data sets. The physics- and statistical-assumptions for the other null SBL experiments and LSND, which are also included in the fit, are described in detail in Ref. [5]. The constraints from atmospheric experiments, according to Ref. [19], have been incorporated as described in Ref. [15].

##### A. General technique

The Monte Carlo method used to apply the oscillation formalism described in Sec. III closely follows the one described in Ref. [15]. We start by randomly varying sets of oscillation parameters:  $\Delta m_{41}^2$ ,  $|U_{e4}|$ ,  $|U_{\mu 4}|$  for the case of  $(3 + 1)$ ;  $\Delta m_{41}^2$ ,  $|U_{e4}|$ ,  $|U_{\mu 4}|$ ,  $\Delta m_{51}^2$ ,  $|U_{e5}|$ ,  $|U_{\mu 5}|$ ,  $\phi_{45}$  for the case of  $(3 + 2)$ . Without loss of generality, we take  $\Delta m_{51}^2 > \Delta m_{41}^2$ . In  $CP$ -conserving models,  $\phi_{45}$  is set to 0 or  $\pi$  by default, whereas in  $CP$ -violating models  $\phi_{45}$  is allowed to vary within the full  $(0, 2\pi)$  range. For each set of oscillation parameters, a signal prediction is obtained and compared to observed data for each SBL experiment, in the form of a  $\chi^2$  for each experiment. For each set of oscillation parameters that is generated, the various  $\chi^2$ 's are linearly summed together to form  $\chi_{\text{SBL}}^2$ , which is then used to extract the best-fit parameters and allowed regions.

A  $\chi^2$  minimization is carried out using a Markov Chain [27]. This minimization procedure relies on calculating the  $\chi^2$  difference between successive sets of parameters and using that as a measure of whether the new point in parameter space is a ‘‘good’’ point to step to, or whether a new set of parameters needs to be drawn again. This is realized in the form of a probability of accepting a new set of parameters,  $P(x_i \rightarrow x_{i+1})$ , given by

$$P(x_i \rightarrow x_{i+1}) = \min(1, e^{-(\chi_{i+1}^2 - \chi_i^2)/T}), \quad (5)$$

where  $x_i$  and  $x_{i+1}$  are two successive points in parameter space, and  $T$  is a temperature parameter. Larger  $T$  values allow for larger  $\Delta\chi^2$  jumps on the  $\chi^2$  surface, and therefore by varying the  $T$  value, one can avoid local minima, as well

as finely scan the parameter space. This minimization method is particularly preferred in fits with large parameter space dimensionality, as in the case of  $(3 + 2)$  oscillation fits, due to its higher efficiency.

In extracting the various confidence level contours, we marginalize over the parameter space and report results obtained with  $\Delta\chi^2$  levels corresponding to 1 degree of freedom for exclusion limits, and 2 degrees of freedom for allowed regions.

To quantify the statistical compatibility between various data sets, we use the Parameter Goodness-of-fit (PG) test introduced in [28]. In this test one quantifies how well various data sets are in agreement, by comparing the minimum  $\chi^2$  obtained by a fit where all data sets have been included as constraints,  $\chi_{\text{min,all}}^2$ , to the sum of the  $\chi^2$  minima obtained by independent fits for each experiment, i.e.,

$$\chi_{\text{PG}}^2 = \chi_{\text{min,all}}^2 - \sum_i \chi_{\text{min},i}^2, \quad (6)$$

where  $i$  runs over experimental data sets in consideration. The PG is obtained from  $\chi_{\text{PG}}^2$  based on the number of common underlying fit parameters,  $ndf_{\text{PG}}$ :

$$\text{PG} = \text{prob}(\chi_{\text{PG}}^2, ndf_{\text{PG}}). \quad (7)$$

For example, for testing the compatibility between KARMEN and LSND for the  $(3 + 2)$   $CP$ -conserving oscillation hypothesis, we fit for both KARMEN and LSND simultaneously to extract  $\chi_{\text{min,K+L}}^2$ , and for KARMEN and LSND separately to extract  $\chi_{\text{min,K}}^2$ , and  $\chi_{\text{min,L}}^2$ , respectively, and obtain:

$$\chi_{\text{PG}}^2(\text{K, L}) = \chi_{\text{min,K+L}}^2 - (\chi_{\text{min,K}}^2 + \chi_{\text{min,L}}^2). \quad (8)$$

As these are appearance experiments, there are 4 common fit parameters for a  $CP$ -conserving  $(3 + 2)$  model ( $\Delta m_{41}^2$ ,  $\Delta m_{51}^2$ ,  $|U_{e4}||U_{\mu 4}|$ , and  $|U_{e5}||U_{\mu 5}|$ ); therefore,

$$\text{PG}(\text{K, L}) = \text{prob}(\chi_{\text{PG}}^2(\text{K, L}), 4). \quad (9)$$

It should be noted that  $\chi^2$ -probabilities and PG tests can lead to drastically different results [28]. This is often a consequence of a large data set simultaneously fitted with small data sets, where the large data set dominates the  $\chi^2$  of the simultaneous fit.

##### B. Inclusion of the MiniBooNE neutrino and antineutrino data sets

The MiniBooNE neutrino data set (BNB-MB( $\nu$ )), described in Sec. II, is included in the fits in the form of two side-by-side distributions of  $\nu_e$  and  $\nu_\mu$  charged-current quasielastic (CCQE) events. Each distribution is a function of neutrino energy, reconstructed under the hypothesis of CCQE neutrino interaction kinematics,  $E_\nu^{\text{QE}}$ . The full 200–3000 MeV range of  $\nu_e$  CCQE data is used in the fit. The observed event distributions are compared to the corre-

sponding Monte Carlo predicted distributions, and a  $\chi^2$  is calculated using a covariance matrix which includes systematic and statistical uncertainties as well as systematic correlations between the predicted  $\nu_e$  and  $\nu_\mu$  distributions.

During the fit, we vary the  $\nu_e$  distribution according to the sterile neutrino oscillation parameters, but keep the  $\nu_\mu$  distribution unchanged. The  $\nu_\mu$  distribution remains unchanged during the fit, despite the possibility of  $\nu_\mu$  disappearance in the MiniBooNE data. In fact, MiniBooNE has released results from  $\nu_\mu$  and  $\bar{\nu}_\mu$  disappearance searches at  $\Delta m^2 \sim 1 \text{ eV}^2$  [29]. These results are relevant as constraints to sterile neutrino mixing parameters under consideration, but they have been purposefully omitted in this analysis, due to the fact that the  $\nu_\mu$  and  $\bar{\nu}_\mu$  CCQE samples used in the disappearance analysis [29] and the (different)  $\nu_\mu$  and  $\bar{\nu}_\mu$  CCQE samples used as constraints in the appearance analyses [16,17] are highly correlated samples, and these correlations have not yet become available. We assume that including MiniBooNE  $\nu_\mu$  disappearance would have a small effect on sterile neutrino fit results, given the large overlap of the  $\nu_\mu$  disappearance limit from MiniBooNE with the corresponding limits from CDHS and CCFR84 [29]. However, the impact of the MiniBooNE disappearance results on the fits considered in this paper will be discussed. Nevertheless, we employ this side-by-side fitting method as it takes advantage of correlations in the  $\nu_\mu$  and  $\nu_e$  predictions and in order to effectively constrain the  $\nu_e$  prediction and reduce systematic uncertainties in the  $\nu_\mu \rightarrow \nu_e$  search.

The fit method follows the details described in [16], except that it uses a different definition for the covariance matrix used in the  $\chi^2$  calculation. Ref. [16] involves an iterative fit method where the  $\chi^2$  calculation for each point on the parameter space being probed uses the covariance matrix calculated according to the *best-fit* signal prediction. Instead, in the MiniBooNE fits presented here, the  $\chi^2$  surface is estimated using the covariance matrix calculated according to the signal prediction at *each point* of the parameter space under consideration. The two fit methods yield similar results, although, by definition, the iterative method of [16] results in a relative shift of the allowed region to the left, i.e. towards smaller oscillation amplitudes.

The MiniBooNE antineutrino data set (BNB-MB( $\bar{\nu}$ )), described in Sec. II, is included in the fits in the same way as the BNB-MB( $\nu$ ) data set, in the form of two side-by-side  $E_\nu^{\text{QE}}$  distributions of  $\bar{\nu}_e$  and  $\bar{\nu}_\mu$  CCQE events. In this case, the disappearance limit obtained using the MiniBooNE  $\bar{\nu}_\mu$  CCQE sample provides substantial coverage of so-far unexplored sterile neutrino mass and mixing parameter space [29]. Even though we do not explicitly fit the MiniBooNE  $\bar{\nu}_\mu$  CCQE distribution for disappearance, we comment on the effect of the limit from [29] in Sec. V, and justify that excluding the MiniBooNE  $\bar{\nu}_\mu$  disappearance information

from the fits does not substantially affect the parameter space of interest. The full 200–3000 MeV range of  $\bar{\nu}_e$  CCQE data is used in the fit. The BNB-MB( $\bar{\nu}$ ) data fit method also follows the details described in [17], except that it uses the definition for the covariance matrix described above.

In fits where both neutrino and antineutrino data are included, it has been assumed that the two data sets are fully uncorrelated. In reality, the two data sets have large systematic correlations. However, neglecting the effects of these correlations is a reasonable approximation, given that the antineutrino data set is statistics limited.

### C. Inclusion of the NuMI-beam data set

The NuMI-MB data set [18], described in Sec. II, is included in the fits in the form of a distribution of  $\nu_e$  CCQE events as a function of reconstructed neutrino energy,  $E_\nu^{\text{QE}}$ . The predicted  $\nu_e$  distribution is obtained by adding to the expected  $\nu_e$  CCQE background any contribution from  $\nu_\mu$  to  $\nu_e$  oscillations. The contribution is estimated as follows: First, a fully (100%) oscillated NuMI-MB  $\nu_\mu \rightarrow \nu_e$  sample is obtained by reweighting the BNB-MB fully oscillated  $\nu_\mu \rightarrow \nu_e$  Monte Carlo predicted sample according to the ratio of the NuMI-beam flux from [18] to the BNB-MB flux [23], as a function true neutrino energy. As the oscillation parameters vary during the fit, a signal prediction is calculated by rescaling the number of events in this fully oscillated sample by the corresponding oscillation probability, according to the true neutrino energy and distance travelled, from production to detection, of each event. We assume a constant  $L$  of 700 meters. The prediction is compared to the observed  $\nu_e$  CCQE events as a function of 10 bins of  $E_\nu^{\text{QE}}$ . The background and signal prediction are assumed to have the same fractional systematic uncertainties, and a statistical uncertainty is calculated for each point in the parameter space according to the signal prediction of each point under consideration. The data and background central value and systematic uncertainty per  $E_\nu^{\text{QE}}$  bin have been estimated from [18]. Unlike the systematic uncertainties of the BNB-MB  $\nu_e$  and  $\bar{\nu}_e$  CCQE data sets, the NuMI-MB  $\nu_e$  CCQE systematics have not been constrained using information from the  $\nu_\mu$  CCQE spectrum from the NuMI beam line. Furthermore, we have not considered potential systematic correlations among the  $\nu_e$  CCQE bins as a function of  $E_\nu^{\text{QE}}$ .

## V. (3 + 1) AND (3 + 2) MODELS AFTER THE NEW MINIBOONE $\nu$ , $\bar{\nu}$ , AND NUMI RESULTS

### A. (3 + 1) fit results

In this section, the new MiniBooNE results are examined under a (3 + 1) oscillation hypothesis and compared to LSND and other null SBL experiments. The new data sets are studied first within the context of appearance-only experiments, and subsequently in fits involving both ap-

TABLE II. Short-baseline oscillation fits considered in this paper.

Fit	Data sets										
	BNB-MB( $\nu$ )	BNB-MB( $\bar{\nu}$ )	LSND	NUMI-MB	KARMEN	NOMAD	CHOOZ	Bugey	CCFR	CDHS	atm
APP	✓	✓	✓	✓	✓	✓					
DIS							✓	✓	✓	✓	✓
$\nu$	✓			✓		✓			✓	✓	✓
$\nu$ APP	✓			✓		✓			✓	✓	✓
$\bar{\nu}$		✓	✓		✓	✓	✓				
$\bar{\nu}$ APP		✓	✓		✓	✓	✓				
signal	✓	✓	✓		✓	✓					
signal APP	✓	✓	✓		✓	✓					
null				✓	✓	✓	✓	✓	✓	✓	✓
null APP				✓	✓	✓	✓	✓	✓	✓	✓

pearance and disappearance data. Fits to only antineutrino and only neutrino SBL experiments are also explored.

Table II provides a reference for all the different data set combinations explored in fits in this paper.

### 1. Studies with appearance-only experiments

Figure 1 shows the allowed regions obtained by independent fits to each of the following three data sets: BNB-MB( $\nu$ ), BNB-MB( $\bar{\nu}$ ), and LSND. The regions are estimated using a two-dimensional global scan of the  $(3 + 1)$  parameter space  $(\sin^2 2\theta_{\mu e}, \Delta m_{41}^2)$ . Each contour is drawn by applying a flat  $\Delta\chi^2 = \chi^2(\sin^2 2\theta_{\mu e}, \Delta m_{41}^2) - \chi_{\min}^2$  cut over the  $\chi^2$  surface, with respect to the global  $\chi^2$  minimum returned by the fit. The figure shows that, similarly to LSND, both BNB-MB data sets yield contours which exclude the no-oscillations (null) hypothesis at 90% CL. The null  $\chi^2$ 's correspond to 22.2 and 24.5 for BNB-MB( $\nu$ ) and BNB-MB( $\bar{\nu}$ ), respectively. The closed contours reflect a contradiction to the oscillation results published by the MiniBooNE collaboration; this is a con-

sequence of the different  $\chi^2$  definition involved in the fit method used here, as pointed out in Sec. IV B. All three data sets, BNB-MB( $\nu$ ), BNB-MB( $\bar{\nu}$ ), and LSND, yield similar best-fit parameters, indicated by the stars on the three graphs, of  $\Delta m_{41}^2$  of order a few  $\text{eV}^2$  and  $\sin^2 2\theta_{\mu e}$  of order  $10^{-2}$ – $10^{-3}$ . The minimum  $\chi^2$  and best-fit parameters returned by each experiment are summarized in Table III.

In light of the above BNB-MB results and the already established LSND anomaly, we find it interesting to study sterile neutrino oscillations with the LSND, BNB-MB( $\nu$ ), and BNB-MB( $\bar{\nu}$ ) data sets assumed to be (positive) “signal” experiments under both the  $(3 + 1)$  and the  $(3 + 2)$  models. This classification is based on the fact that all three data sets exclude the null result at 90% confidence level under a  $(3 + 1)$  oscillation hypothesis. Figure 2 shows the BNB-MB( $\nu$ ) and BNB-MB( $\bar{\nu}$ ) event distributions for both the null and the best-fit  $(3 + 1)$  oscillation hypothesis for each data set. In the case of the BNB-MB( $\nu$ ) data set, even though the best-fit hypothesis provides a better description of the event spectrum at 90% CL ( $\Delta\chi^2 = \chi_{\text{null}}^2 - \chi_{\text{best-fit}}^2 = 4.7$ , for 2 fit parameters), it fails to fully explain

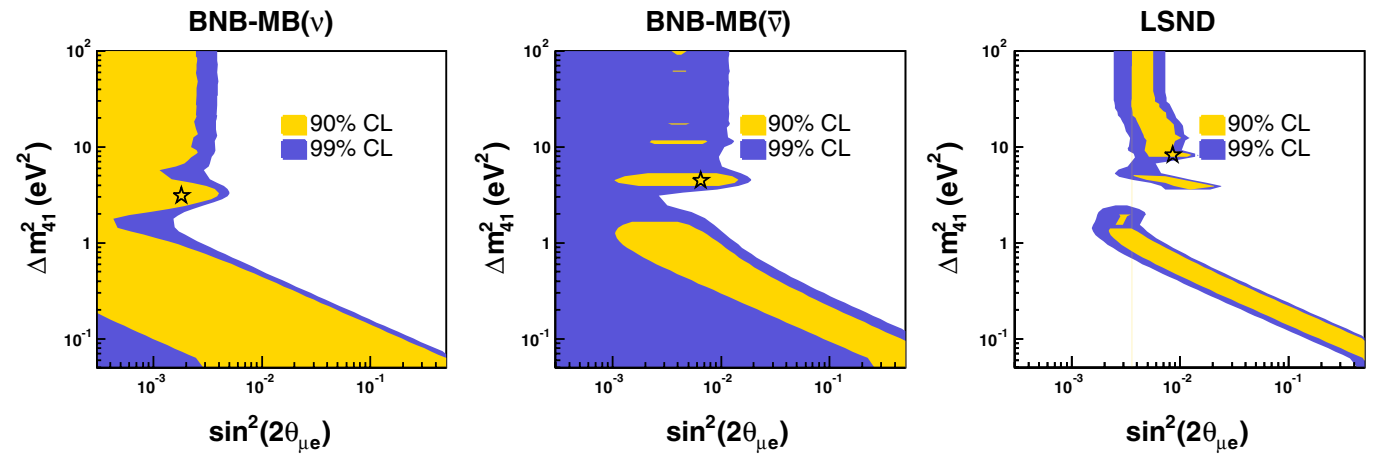


FIG. 1 (color online). Allowed regions (filled areas) at 90% and 99% CL from BNB-MB( $\nu$ )-only, BNB-MB( $\bar{\nu}$ )-only, and LSND-only  $(3 + 1)$  fits. These fits are, by construction,  $CP$ -conserving. The stars indicate the three respective best-fit points. All three data sets show closed contours at 90% CL. See the text for more details.

TABLE III. Comparison of best-fit values for mass splittings and mixing angles obtained from (3 + 1) fits to appearance data sets and appearance + disappearance data sets. Mass splittings are shown in  $\text{eV}^2$ . The minimum  $\chi^2$  from each fit, as well as the  $\chi^2$ -probability are also given. The signal appearance (APP) data sets include BNB-MB( $\nu$ ), BNB-MB( $\bar{\nu}$ ) and LSND. The null APP data sets include KARMEN, NOMAD, and NUMI-MB; the maximal best-fit  $\sin^2 2\theta_{\mu e}$  in this case is inconsequential, as it corresponds to a best-fit  $\Delta m^2$  region of very poor sensitivity. See the text for more details.

Data set	$\chi^2$ (dof)	$\chi^2$ -probability	$\Delta m_{41}^2$	$\sin^2 2\theta_{\mu e}$	$\sin^2 2\theta_{\mu\mu}$	$\sin^2 2\theta_{ee}$
<i>Appearance-only fits:</i>						
LSND	3.4 (3)	34%	8.19	0.0085	-	-
BNB-MB( $\nu$ )	17.5 (16)	35%	3.12	0.0018	-	-
BNB-MB( $\bar{\nu}$ )	17.6 (16)	35%	4.46	0.0065	-	-
NUMI-MB	2.0 (8)	98%	6.97	0.020	-	-
KARMEN	6.0 (7)	54%	6.81	0.00096	-	-
NOMAD	33.3 (28)	31%	53.3	0.00012	-	-
signal APP	50.3 (39)	11%	0.045	0.98	-	-
signal APP <sup>a</sup>	50.4 (39)	10%	0.15	0.090	-	-
null APP	46.6 (47)	49%	0.040	1.00	-	-
APP	97.1 (88)	24%	0.045	1.00	-	-
APP <sup>a</sup>	97.2 (88)	24%	0.15	0.090	-	-
LSND + MB-BNB( $\bar{\nu}$ )	22.3 (21)	38%	4.55	0.0074	-	-
LSND + MB-BNB( $\bar{\nu}$ ) <sup>a</sup>	22.3 (21)	38%	4.55	0.0074	-	-
LSND + MB-BNB( $\bar{\nu}$ ) + KARMEN	33.6 (30)	29%	0.57	0.0097	-	-
BNB-MB( $\nu$ ) + NUMI-MB + NOMAD	57.8 (56)	40%	0.033	1.00	-	-
<i>Appearance and disappearance fits:</i>						
all SBL <sup>a</sup>	197.4 (196)	46%	0.92	0.0025	0.13	0.073
$\nu$	90.5 (90)	47%	0.19	0.031	0.031	0.034
$\bar{\nu}$	87.9 (103)	86%	0.91	0.0043	0.35	0.043

<sup>a</sup>In these fits, the electron and muon content of the sterile neutrino mass eigenstate have been explicitly constrained to  $<0.3$ .

the low-energy excess. Therefore, the (3 + 1) oscillation hypothesis alone seems inadequate as an explanation for the low-energy excess, as also reported by the MiniBooNE collaboration [13,16]. In the case of the BNB-MB( $\bar{\nu}$ ) data set, the best-fit hypothesis provides a better description of

the data in the 500–700 MeV range. However, the statistical uncertainties are too large to allow for a strong conclusion.

The allowed regions obtained by a joint analysis of BNB-MB( $\nu$ ) + BNB-MB( $\bar{\nu}$ ) + LSND, as well as a joint

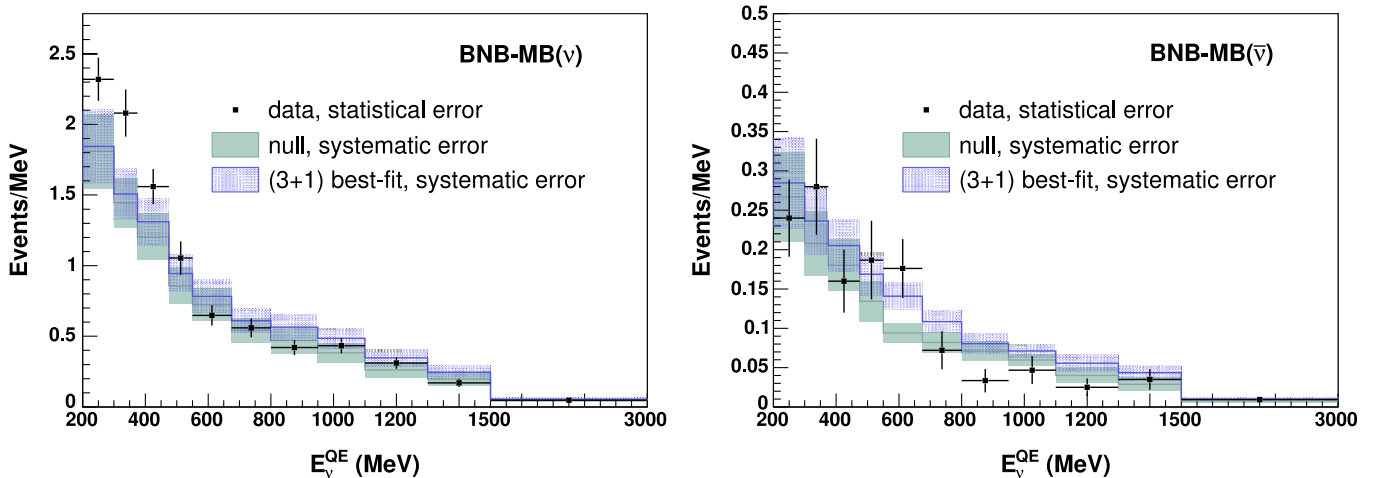


FIG. 2 (color online). Left: Null and (3 + 1) best-fit predicted event distributions  $(\Delta m_{41}^2, \sin^2(2\theta_{\mu e})) = (3.12, 0.0018)$  for BNB-MB( $\nu$ ). Right: Null and (3 + 1) best-fit predicted event distributions  $(\Delta m_{41}^2, \sin^2(2\theta_{\mu e})) = (4.46, 0.0065)$  for BNB-MB( $\bar{\nu}$ ). The event distributions are shown as functions of reconstructed neutrino energy,  $E_{\nu}^{\text{QE}}$ . The data are shown in black points with statistical uncertainty. The null (no-oscillation) prediction is shown by the light gray histogram with (solid) systematic error band. The best-fit prediction (signal and background) is shown by the blue (dark gray) histogram with (shaded) systematic error band.



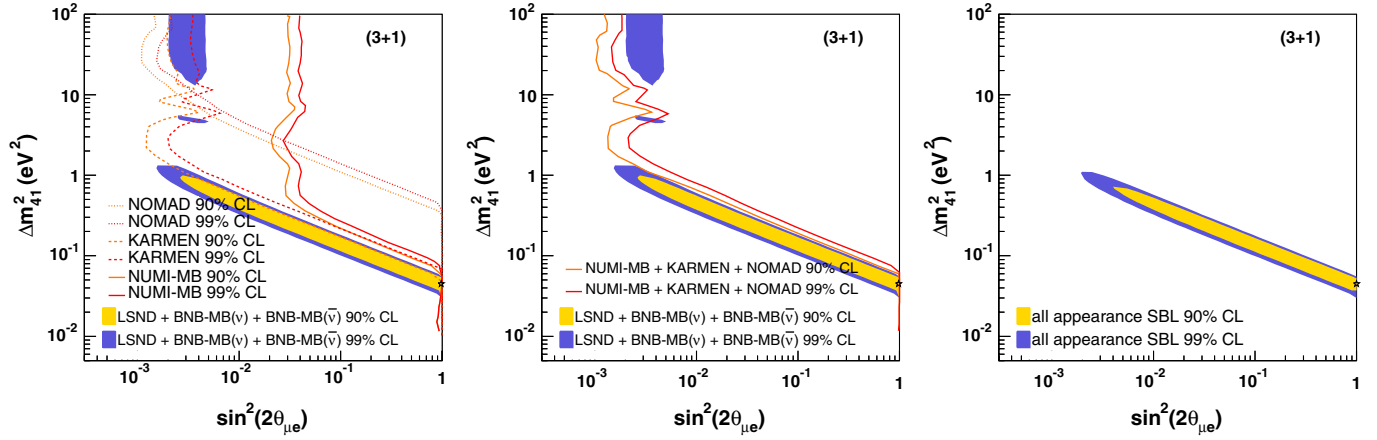


FIG. 3 (color online). Left: Allowed 90% and 99% CL regions (light and dark filled areas, respectively) from a combined analysis of BNB-MB( $\nu$ ), BNB-MB( $\bar{\nu}$ ) and LSND data sets, and 90% and 99% exclusion limits (light and dark curves, respectively) from each of the null appearance experiments, NUMI-MB (solid curves), KARMEN (dashed curves) and NOMAD (dotted curves). Middle: The same allowed region with overlaid 90% and 99% exclusion limits from a combined analysis of all null appearance experiments. Right: Allowed region obtained by a combined analysis of all appearance data sets, signal and null. See the text for more details.

analysis of BNB-MB( $\bar{\nu}$ ) + LSND are shown on the left panels of Figs. 3 and 4, respectively. In the case of the combined fit of all three data sets (Fig. 3), due to the difference in preferred mixing amplitudes mentioned in the previous paragraph, the best-fit point ends up shifting from an intermediate  $\Delta m^2$  and small mixing amplitude to a smaller  $\Delta m^2$  and maximal mixing amplitude of 0.98. Obviously a maximal mixing amplitude is unphysical in the case of sterile neutrino oscillations. If the fits are repeated with the electron and muon content of the sterile mass eigenstate limited to values less than 0.3 [30], the returned  $\chi^2$ -probabilities are 10% and 38%, for BNB-MB( $\nu$ ) + BNB-MB( $\bar{\nu}$ ) + LSND and BNB-MB( $\bar{\nu}$ ) + LSND, respectively; the reduction in  $\sin^2 2\theta_{\mu e}$  space has

essentially no effect on these results. The best-fit parameters from these fits are also given in Table III.

Perhaps a more interesting observation regarding Fig. 1 is the striking similarity of BNB-MB( $\bar{\nu}$ ) and LSND 90% CL allowed regions and best-fit oscillation parameters, keeping in mind that both data sets describe antineutrino oscillations. It should be noted that in a (3 + 1) oscillation scenario, under the assumption of  $CPT$  invariance, there can be no difference between neutrino and antineutrino oscillation probabilities. However, a PG test, as described in Sec. IVA, suggests a significantly higher compatibility (49%) between BNB-MB( $\bar{\nu}$ ) and LSND, rather than for all three signal experiments (BNB-MB( $\nu$ ), BNB-MB( $\bar{\nu}$ ) and LSND) combined (0.26%). This is also supported by the

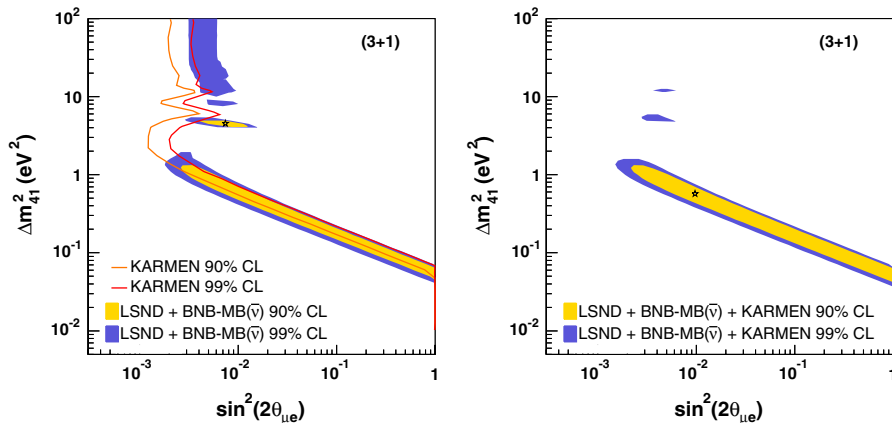


FIG. 4 (color online). Left: The allowed 90% and 99% CL regions (light and dark filled areas, respectively) from a combined analysis of BNB-MB( $\bar{\nu}$ ) and LSND data sets, and 90% and 99% exclusion limits (light and dark curves, respectively) from KARMEN. A comparison of only these three experiments is interesting, as these three experiments have searched for antineutrino oscillations at short baselines. Right: The allowed regions obtained from a combined analysis of all three experiments (BNB-MB( $\bar{\nu}$ ), LSND, and KARMEN). See the text for more details.

TABLE IV. Summary of  $\chi^2$ -probabilities for (3 + 1) fits with different combinations of SBL data sets, and PG results testing compatibility among different data sets. See the text for more details.

Data set	$\chi^2$ -probability (%)	PG (%)	
APP	24	PG(BNB-MB( $\nu$ ), BNB-MB( $\bar{\nu}$ ), LSND, NUMI-MB, KARMEN, NOMAD)	= Prob(17.3, 2) = $1.7 \times 10^{-2}$
signal APP	11	PG(BNB-MB( $\nu$ ), BNB-MB( $\bar{\nu}$ ), LSND)	= Prob(11.9, 2) = 0.26
LSND + MB-BNB( $\bar{\nu}$ )	38	PG(BNB-MB( $\bar{\nu}$ ), LSND)	= Prob(1.4, 2) = 49
$\bar{\nu}$ APP	29	PG(BNB-MB( $\bar{\nu}$ ), LSND, KARMEN)	= Prob(6.7, 2) = 3.4
$\nu$ APP	40	PG(BNB-MB( $\nu$ ), NUMI-MB, NOMAD)	= Prob(4.9, 2) = 8.8
all SBL <sup>a</sup>	46	PG(BNB-MB( $\nu$ ), BNB-MB( $\bar{\nu}$ ), NUMI-MB, LSND, KARMEN, NOMAD, Bugey, CHOOZ, CCFR84, CDHS, ATM)	= Prob(42.0, 2) = $7.6 \times 10^{-8}$
		PG(APP, DIS)	= Prob(14.8, 2) = $6.2 \times 10^{-2}$
		PG( $\nu$ , $\bar{\nu}$ )	= Prob(18.8, 2) = $8.1 \times 10^{-3}$
$\nu$	47	PG(BNB-MB( $\nu$ ), NUMI-MB, NOMAD, CCFR84, CDHS, ATM)	= Prob(14.7, 2) = $6.3 \times 10^{-2}$
$\bar{\nu}$	86	PG(BNB - MB( $\bar{\nu}$ ), LSND, KARMEN, Bugey, CHOOZ)	= Prob(8.43, 2) = 1.5

<sup>a</sup>In these fits, the electron and muon content of the sterile neutrino mass eigenstate have been explicitly constrained to  $< 0.3$ .

$\chi^2$ -probabilities returned by the fits: 11% in the case of the BNB-MB( $\nu$ ) + BNB-MB( $\bar{\nu}$ ) + LSND fit, and 38% in the case of the BNB-MB( $\bar{\nu}$ ) + LSND fit. This incompatibility is due to the fact that the BNB-MB( $\nu$ ) data set prefers a mixing amplitude  $\sim 3$  times smaller than the amplitude preferred by LSND or BNB-MB( $\bar{\nu}$ ), and excludes the LSND and BNB-MB( $\bar{\nu}$ ) best fits at 99% CL. Table IV provides a summary of the above  $\chi^2$ -probabilities and PG test results.

Figs. 3 and 4 also illustrate the limits from various combinations of the remaining three (null) SBL appearance experiments, KARMEN, NOMAD, and NUMI-MB, under a (3 + 1) oscillation scenario, overlaid on the allowed regions described above.

The 90% and 99% CL limits obtained by each of the null appearance experiments are shown on the left panel of Fig. 3. These limits correspond to the upper  $\sin^2 2\theta_{\mu e}$  values allowed at each  $\Delta m_{41}^2$ , estimated using a one-sided raster scan of the parameter space. It is interesting to point out that, despite the indication of a slight excess ( $1.2\sigma$ ) of observed  $\nu_e$ -like events for neutrino energies below 900 MeV found in the NuMI analysis [18], the currently assumed NUMI-MB systematic and statistical uncertainties are quite large, resulting in a limit that is much weaker relative to the limits of KARMEN and NOMAD. In fact, due to this excess and the large systematic uncertainties, the NUMI-MB data set provides a very good fit to (3 + 1) models, with a  $\chi^2$ -probability of 98%. The event distributions for the null and best-fit (3 + 1) oscillation hypothesis for the NUMI-MB data set are shown in Fig. 5. The observed distribution fits nicely to an oscillation signal at  $(\Delta m_{41}^2, \sin^2(2\theta_{\mu e})) = (7.36, 0.019)$ . Such large signal, however, would be in disagreement with the BNB-MB( $\nu$ ) results. Additional data and reduced systematic uncertain-

ties in the NUMI-MB analysis are necessary for higher sensitivity and more conclusive results. This is currently an ongoing effort and new results are expected soon. The limits from a combined NUMI-MB + KARMEN + NOMAD analysis are shown on the middle panel of Fig. 3. Both panels illustrate that the null appearance experiments provide essentially no constraints to the parameter space allowed by the BNB-MB and LSND data sets, except at higher  $\Delta m^2$ .

The best-fit parameters obtained independently from the NUMI-MB and KARMEN data sets, shown in Table III,

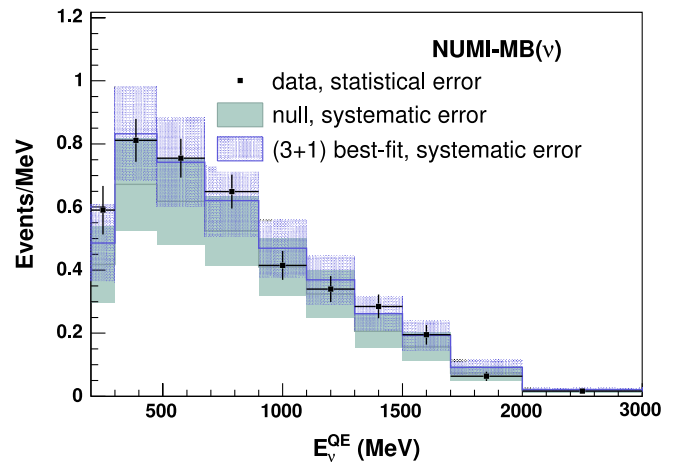


FIG. 5 (color online). Null and (3 + 1) best-fit predicted event distributions  $(\Delta m_{41}^2, \sin^2(2\theta_{\mu e})) = (7.36, 0.019)$  for NUMI-MB. The data are shown in black points with statistical uncertainty. The null (no-oscillation) prediction is shown by the light gray histogram with (solid) systematic error band. The best-fit prediction (signal and background) is shown by the blue (dark gray) histogram with (shaded) systematic error band.

are similar to those of LSND, BNB-MB( $\nu$ ), and BNB-MB( $\bar{\nu}$ ). The NOMAD data set, on the other hand, prefers a much larger  $\Delta m_{41}^2 \sim 50 \text{ eV}^2$ , and a much smaller  $\sin^2 2\theta_{\mu e} \sim 10^{-4}$ .

A combined analysis of all appearance data yields a  $\chi^2$ -probability of 24% for the best-fit hypothesis, both in the case where maximal mixing is allowed in the fit, and in the case where the electron and muon content of the sterile mass eigenstate has been limited to small values ( $< 0.3$ ). The allowed region obtained by a joint analysis of all appearance experiments under a (3 + 1) oscillation scenario is shown in the right panel of Fig. 3.

Similarly, Fig. 4 (left) corresponds to the allowed region obtained by a joint analysis of BNB-MB( $\bar{\nu}$ ) + LSND. The limit shown is that of the KARMEN experiment, which is the only other SBL experiment to perform an appearance search with antineutrinos. The KARMEN limit provides substantial coverage of the joint LSND and BNB-MB( $\bar{\nu}$ ) allowed region, excluding the best-fit point of the LSND + BNB-MB( $\bar{\nu}$ ) fit at  $>99\%$  CL. However, KARMEN imposes little constraint to the lower- $\Delta m^2$  allowed solutions. A joint analysis of all three data sets yields a  $\chi^2$ -probability of 29% for the best-fit hypothesis, and an allowed region shown on the right panel of Fig. 4. The  $\chi^2$ -probability remains the same for fits where the electron and muon content of the sterile mass eigenstates have been limited to values less than 0.3. Nevertheless, as shown in Table IV, the three data sets are compatible at only 3.4%. New results from MiniBooNE with increased antineutrino statistics should be able to provide more information to these fits [17].

## 2. Studies with appearance and disappearance experiments

Much stronger constraints than those of the null appearance experiments are provided by the null disappearance experiments (CCFR84, CDHS, CHOOZ, and Bugey) and atmospheric constraints, under the assumptions of *CPT* conservation and unitarity of the neutrino mixing matrix. The 90% and 99% CL exclusion limits from a combined analysis of all null data sets (NUMI-MB, KARMEN, NOMAD, Bugey, CHOOZ, CCFR84, CDHS, and atmospheric constraints) are shown in Fig. 6. The figure shows that the parameter space jointly-allowed by BNB-MB( $\nu$ ) + BNB-MB( $\bar{\nu}$ ) + and LSND at 99% CL is excluded by a combined analysis of all null SBL experiments, appearance and disappearance, at 99% CL. The severe tension between LSND and the null SBL experiments [14] continues to exist and in fact increases further with the addition of BNB-MB( $\nu$ ) and BNB-MB( $\bar{\nu}$ ) data. The signal results show low (0.15%) compatibility with null results. The LSND result remains to be mostly responsible for the low compatibility, as the BNB-MB( $\nu$ ) and BNB-MB( $\bar{\nu}$ ) experiments show 14% and 3.7% compatibility with the null experiments, respectively.

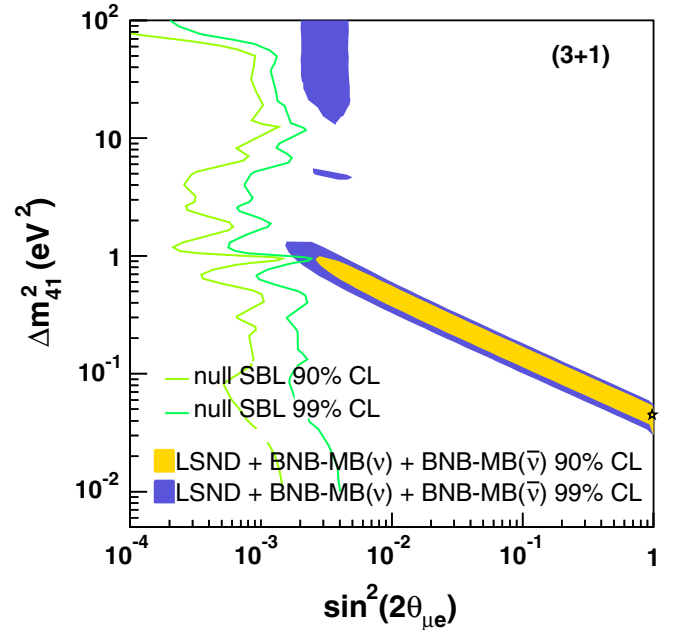


FIG. 6 (color online). The allowed 90% and 99% CL regions (light and dark filled areas, respectively) from a combined analysis of BNB-MB( $\nu$ ) + BNB-MB( $\bar{\nu}$ ) + LSND data sets, and 90% and 99% exclusion limits (light and dark curves, respectively) from a combined analysis of all remaining (null, appearance and disappearance) SBL experiments. The null fit includes atmospheric constraints. The null SBL experiments exclude the joint 99% CL allowed region at 99% CL.

Figure 7 shows the allowed region obtained by the joint BNB-MB( $\bar{\nu}$ ) + LSND + KARMEN + Bugey + CHOOZ analysis. Here, the  $\bar{\nu}_e$  disappearance constraints from Bugey and CHOOZ are interesting to consider from the perspective of a joint analysis of only antineutrino SBL experiments. In a joint fit, all of the above (antineutrino) experiments yield a high  $\chi^2$ -probability of 86%, and 1.5% compatibility. In these fits, Bugey and CHOOZ constrain  $|U_{e4}|$ , but provide no direct constraints on  $|U_{\mu 4}|$ . However, a joint analysis with the LSND, BNB-MB( $\bar{\nu}$ ), and KARMEN appearance experiments, which are sensitive to the product of  $|U_{e4}||U_{\mu 4}|$ , provides indirect constraints to the  $|U_{\mu 4}|$  mixing element. Figure 7 (left) also shows that a fit to all antineutrino experiments without LSND yields similar closed contours at 90% CL, which include the best-fit point. Current constraints from MiniBooNE on  $\bar{\nu}_\mu$  disappearance [29] provide relatively small constraints to the  $\sin^2 2\theta_{\mu\mu}$  allowed space, as illustrated in the right panel of Fig. 7. However, new results from a joint MiniBooNE and SciBooNE [31]  $\bar{\nu}_\mu$  disappearance search, which are expected soon [32], may be able to probe this region with higher sensitivity, and will be interesting within the context of *CPT*-violating models. According to the best-fit oscillation parameters from a fit to only antineutrino SBL data, MiniBooNE should observe muon antineutrino disappearance with an amplitude of  $\sin^2 2\theta_{\mu\mu} \sim 0.35$ , at

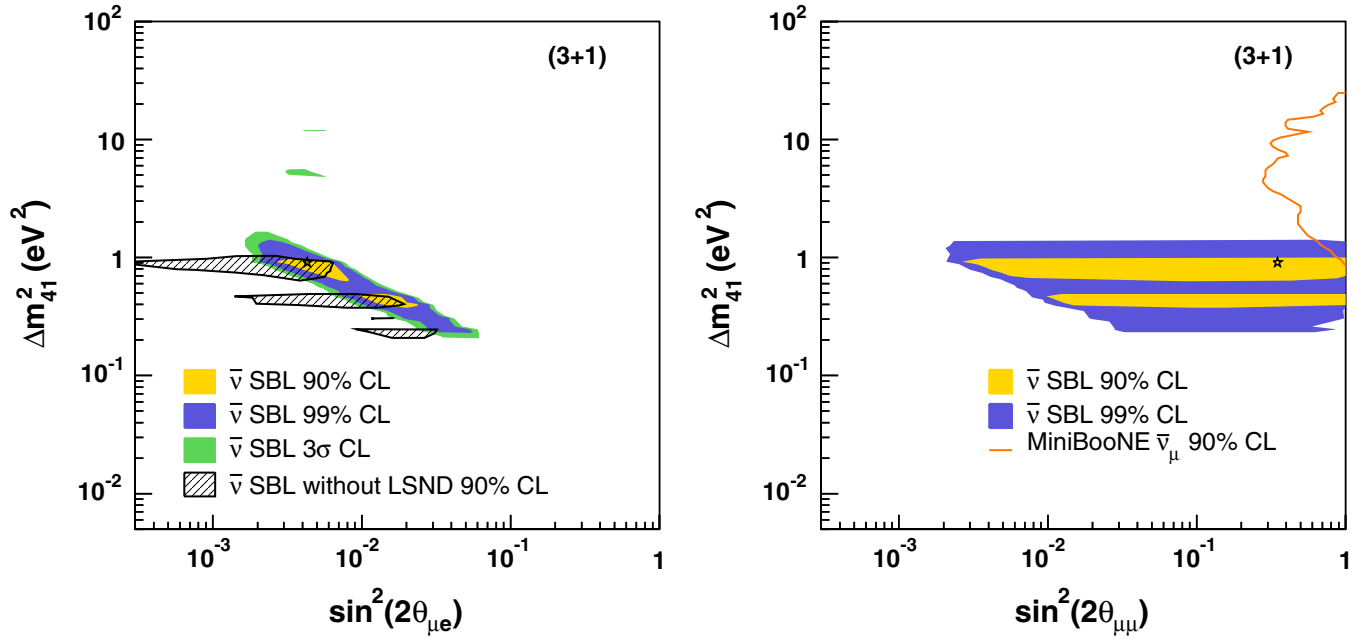


FIG. 7 (color online). The allowed 90%, 99%, and  $3\sigma$  CL regions from a combined analysis of all antineutrino SBL data sets. The left plot also shows the 90% CL allowed region obtained from a combined analysis of all antineutrino experiments except LSND (KARMEN, BNB-MB( $\bar{\nu}$ ), Bugey, and CHOOZ). The right plot also shows the 90% CL exclusion limit from [29]. The MiniBooNE  $\bar{\nu}_\mu$  disappearance search excludes the parameter space to the right of the line at 90% CL. See the text for more details.

$\Delta m^2 \sim 0.91 \text{ eV}^2$ . The MINOS experiment [33] should also have sensitivity to these oscillation parameters in antineutrino running mode; muon antineutrino disappearance search results from MINOS are expected soon [34]. Incorporation of the upcoming MiniBooNE and MINOS

disappearance results in these fits is currently being investigated.

Neutrino-only fits also yield a reasonably high  $\chi^2$ -probability of 47%; the corresponding allowed regions are shown in Fig. 8. Current constraints from MiniBooNE

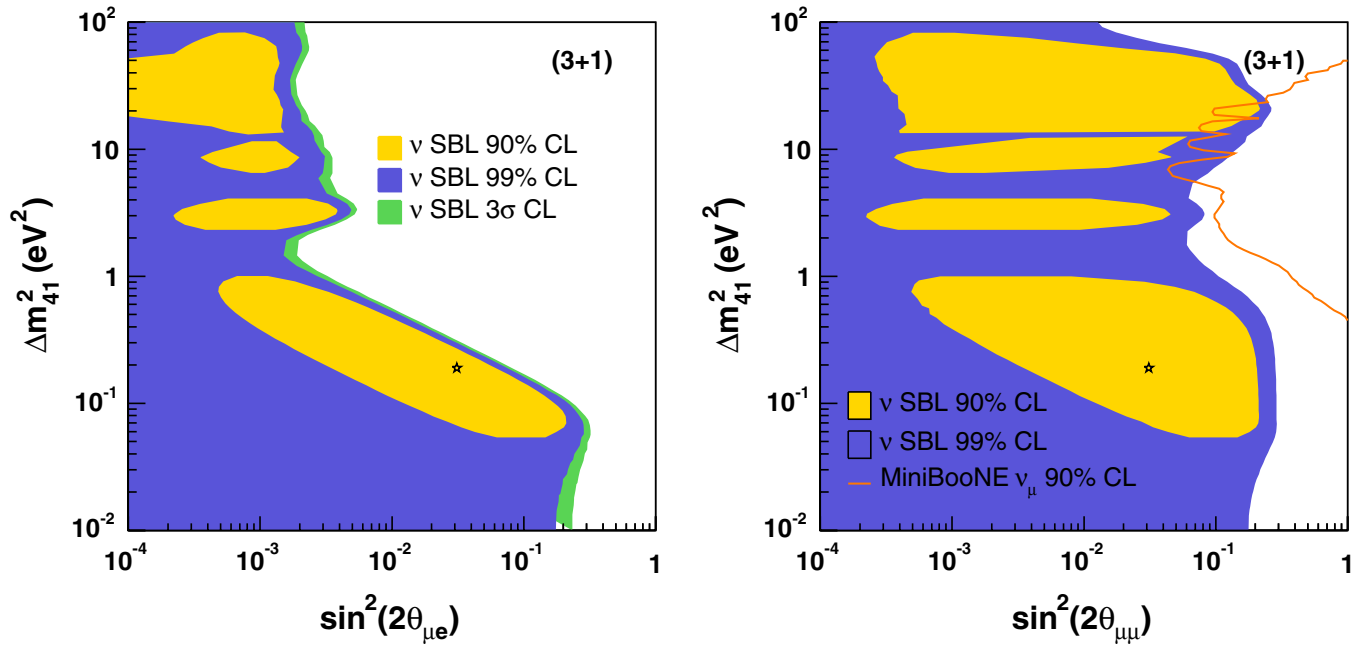


FIG. 8 (color online). The allowed 90%, 99%, and  $3\sigma$  CL regions from a combined analysis of all neutrino SBL data sets. The right plot also shows the 90% CL exclusion limit from [29]. The MiniBooNE  $\nu_\mu$  disappearance search excludes the parameter space to the right of the line at 90% CL. See the text for more details.

$\nu_\mu$  disappearance are shown on the right panel of Fig. 8. Interestingly, fits to only neutrino SBL data also yield a closed contour at 90% CL. The parameter space, however,

points to smaller mixing amplitudes relative to those preferred by the antineutrino-only fit. Neutrino-only fits and antineutrino-only fits are incompatible, with a PG of

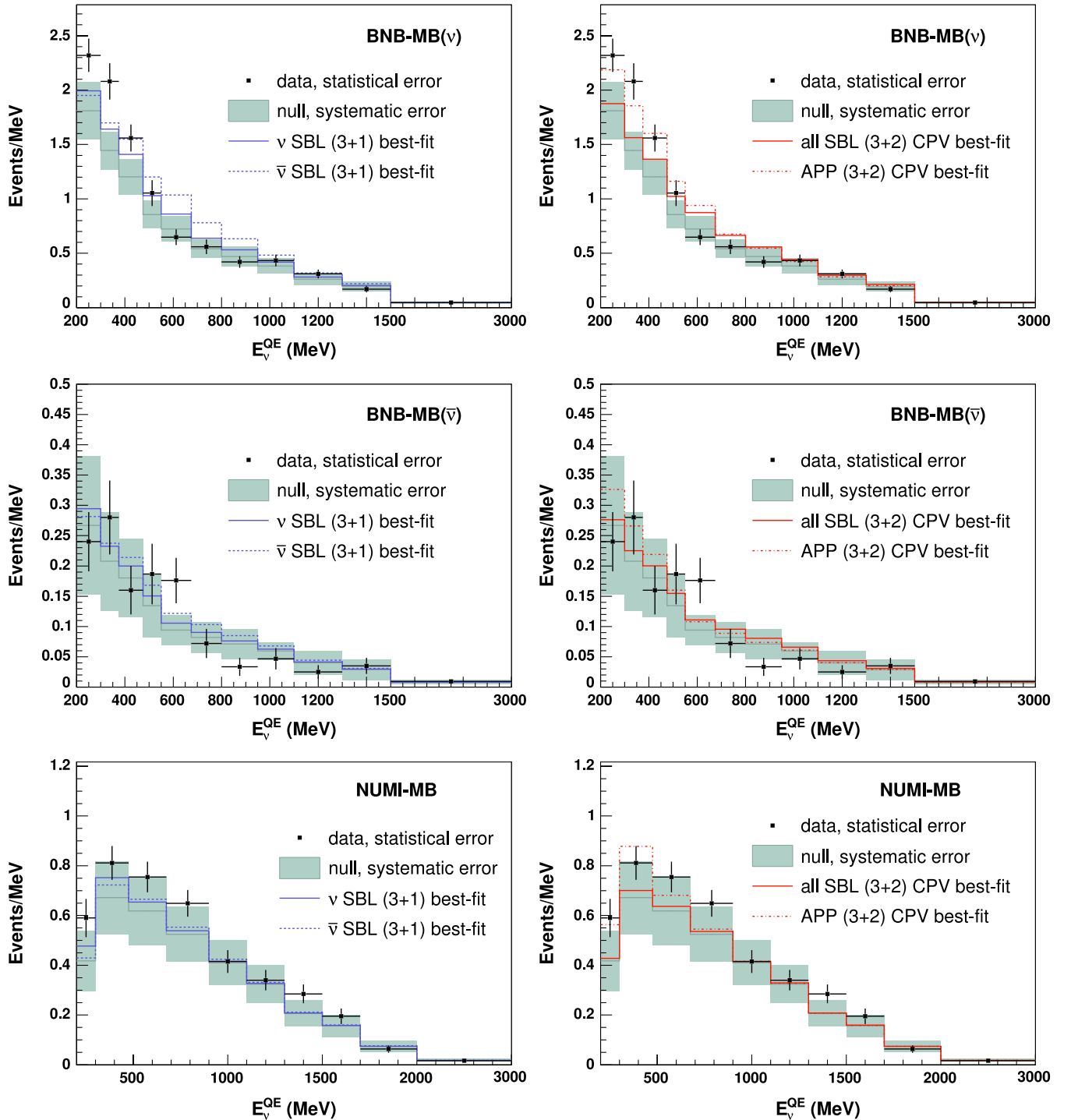


FIG. 9 (color online). Left: MiniBooNE predicted event distributions using the neutrino-only (3 + 1) best-fit parameters  $(\Delta m^2, \sin^2(2\theta)) = (0.19, 0.031)$  in blue (dark gray) solid line and antineutrino-only (3 + 1) best-fit parameters  $(\Delta m^2, \sin^2(2\theta)) = (0.91, 0.0043)$  in blue (dark gray) dashed line. The null predictions are shown in light gray with systematic error bands. The observed data are shown in black points with statistical error bars. Right: MiniBooNE predicted event distributions using the best-fit parameters obtained from a (3 + 2) CP-violating fit to all SBL experiments and appearance-only SBL experiments, in red (dark gray) solid line and red (dark gray) dashed line, respectively.

$8.1 \times 10^{-5}$ , as shown in Table IV. The large incompatibility between antineutrino and neutrino SBL results suggests that the neutrino and antineutrino data sets cannot be accommodated within a  $(3 + 1)$   $CPT$ -conserving sterile neutrino oscillation scenario. However, the constraining power of antineutrino SBL experiments alone on  $\Delta m_{41}^2$  and  $\sin^2 2\theta_{\mu e}$  is remarkable, and invites exploration of models that provide the possibility of different oscillation patterns for neutrinos versus antineutrinos.

Figure 9 (left) shows a comparison of the BNB-MB( $\nu$ ), BNB-MB( $\bar{\nu}$ ), and NUMI-MB event distributions for the neutrino-only best-fit parameters and antineutrino-only best-fit parameters. The neutrino best-fit parameters provide a better description to BNB-MB( $\nu$ ) and NUMI-MB distributions than the antineutrino best-fit parameters, with  $\chi_{\text{BNB-MB}(\nu)}^2 = 18.4$  and  $\chi_{\text{NUMI-MB}}^2 = 4.4$ , compared to  $\chi_{\text{BNB-MB}(\bar{\nu})}^2 = 32.4$  and  $\chi_{\text{NUMI-MB}}^2 = 4.8$ . On the other hand, the antineutrino best-fit parameters provide a better description to BNB-MB( $\bar{\nu}$ ) than the neutrino best-fit parameters ( $\chi_{\text{BNB-MB}(\bar{\nu})}^2 = 19.7$ , compared to  $\chi_{\text{BNB-MB}(\nu)}^2 = 21.7$ ).

The best-fit results from the  $(3 + 1)$  oscillation fit involving all SBL data sets are summarized in Table III. The best-fit parameters from neutrino-only and antineutrino-only fits are also shown.

### B. $(3 + 2)$ fit results

Neutrino oscillation models with more than one sterile neutrino have been of particular interest because they open up the possibility of observable  $CP$  violation effects in short-baseline neutrino oscillations. If  $(3 + n)$  sterile neutrino oscillations are realized in nature, with  $n > 1$ ,  $CP$  violation becomes a natural possibility, which is very appealing from the perspective of theories attempting to explain the matter-antimatter asymmetry in our universe [35].

In this section, the new MiniBooNE results are examined under both a  $CP$ -conserving (CPC) and a  $CP$ -violating (CPV)  $(3 + 2)$  oscillation hypothesis. The new results are studied first within the context of appearance-only experiments, and subsequently in fits involving both appearance and disappearance data.

From the point of view of the data at hand from LSND, BNB-MB( $\nu$ ), and BNB-MB( $\bar{\nu}$ ) (see Fig. 1),  $CP$  violation offers the potential of reconciling two experimental signatures—an excess in LSND data at  $3.8\sigma$  significance and one suggested at 90% CL in BNB-MB antineutrino data, both pointing to relatively large mixing, reconciled with a possible excess found in BNB-MB neutrino data suggesting relatively small mixing, both at a similar  $L/E$ —as manifestations of the same underlying oscillation hypothesis.

It should be noted that in the studies presented in this section, due to the larger dimensionality of the fits, the electron and muon content of the sterile mass eigenstates

have been limited to values less than 0.3. This is a realistic assumption for sterile neutrino oscillation models.

#### 1. Studies with appearance-only experiments

Allowing for  $CP$  violation in  $(3 + 2)$  fits to LSND and BNB-MB( $\nu$  and  $\bar{\nu}$ ) data leads to a reduction in absolute  $\chi^2$  of 4.4, for 1 degree of freedom (dof), corresponding to a best-fit CPV phase  $\phi_{45} = 1.7\pi$ . The  $\chi^2$ -probability of the fit increases from 13% in the CPC case to 21% in the CPV case. The same test can be performed using all appearance data. In this case, allowing for  $CP$  violation leads to a reduction in  $\chi^2$  of 3.2 for 1 dof, with a best-fit CPV phase  $\phi_{45} = 1.1\pi$ . The  $\chi^2$ -probability from the CPV fit is comparable to that of a signal-only fit, at 27%.

The 90% and 99% CL allowed  $(\Delta m_{41}^2, \Delta m_{51}^2)$  parameter space obtained by a combined fit to BNB-MB( $\nu$ ) + BNB-MB( $\bar{\nu}$ ) + LSND is shown in Fig. 10. The figure illustrates that both scenarios, CPC (left) and CPV (right), prefer similar  $\Delta m^2$  parameters at 99% CL; however, the CPV hypothesis is more restrictive in  $\Delta m_{51}^2$  at 90% CL, shown by the shrunk light yellow (light gray) area in the right plot.

A similar 99% CL allowed parameter space is obtained when data from all appearance experiments are included in the fits, as shown in Fig. 11. Here, the effect of  $CP$  violation on constraining  $\Delta m_{51}^2$  is not as strong.

The best-fit parameters for the signal-only and appearance-only fits are summarized in Table V.

#### 2. Studies with appearance and disappearance experiments

A dramatic reduction in the allowed  $(\Delta m_{41}^2, \Delta m_{51}^2)$  parameter space occurs once all SBL data sets are considered in the fit, as shown in Fig. 12. Compared to the CPC hypothesis, the CPV hypothesis fails to provide a substantially better description of the data, reflected by the reduction in  $\chi^2$  of 1.3 for 1 dof. The returned  $\chi^2$ -probability for both CPC and CPV fits is 52%; however, a PG test among all experimental data sets yields  $7.6 \times 10^{-8}\%$  compatibility between all data sets. This large incompatibility in part reflects the large number of data sets being compared, a consequence of the nature of the PG test; however, as discussed in Sec. VI, an underlying source of tension exists due LSND and three other data sets: BNB-MB( $\nu$ ), CDHS, and atmospheric constraints. The best-fit parameters extracted from a fit to all SBL data are also summarized in Table V.

A comparison of Tables III and V suggests that, with the addition of the new data sets from MiniBooNE, the  $(3 + 2)$  oscillation hypothesis now provides only a marginally better description of all data, compared to the  $(3 + 1)$  hypothesis. Compared to  $(3 + 1)$  models,  $(3 + 2)$   $CP$ -conserving models give a reduction of 5.9  $\chi^2$  units for 3 additional fit parameters, while  $(3 + 2)$   $CP$ -violating

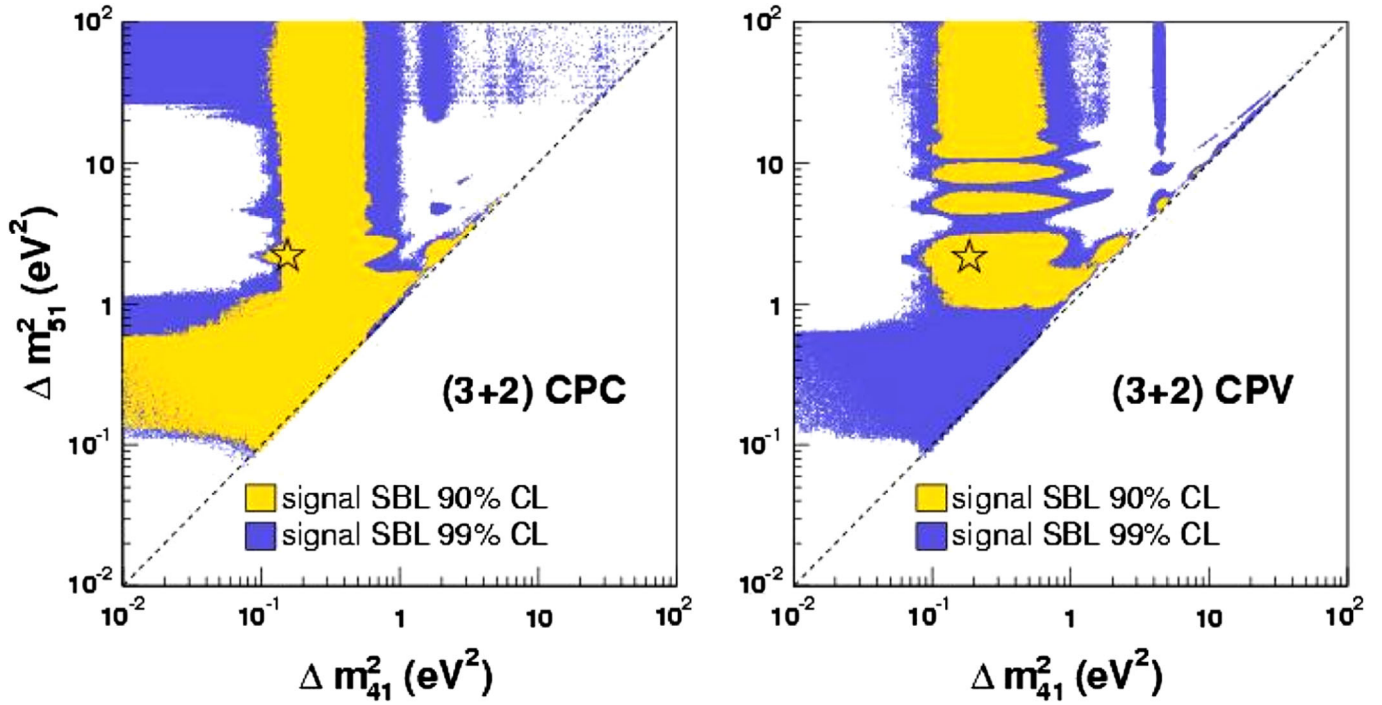


FIG. 10 (color online). Allowed regions in  $(\Delta m_{41}^2, \Delta m_{51}^2)$  space for fits to  $CP$ -conserving (CPC, left) and  $CP$ -violating (CPV, right)  $(3 + 2)$  oscillation models. Only the BNB-MB( $\nu$ ), BNB-MB( $\bar{\nu}$ ) and LSND data sets have been included in the fit.

models give a reduction of 7.2  $\chi^2$  units with 4 additional parameters. This represents a relative improvement that is significantly smaller than that found in Ref. [5] from fits using data sets prior to the new MiniBooNE results. The MiniBooNE event distributions for the  $(3 + 2)$  CPV best-

fit parameters are shown on the right panels of Fig. 9. The best-fit parameters from a  $(3 + 2)$  CPV fit to appearance-only SBL data provide a better description of the MiniBooNE data than those from a  $(3 + 2)$  CPV fit to all SBL data, except for the case of BNB-MB( $\bar{\nu}$ ). In the case

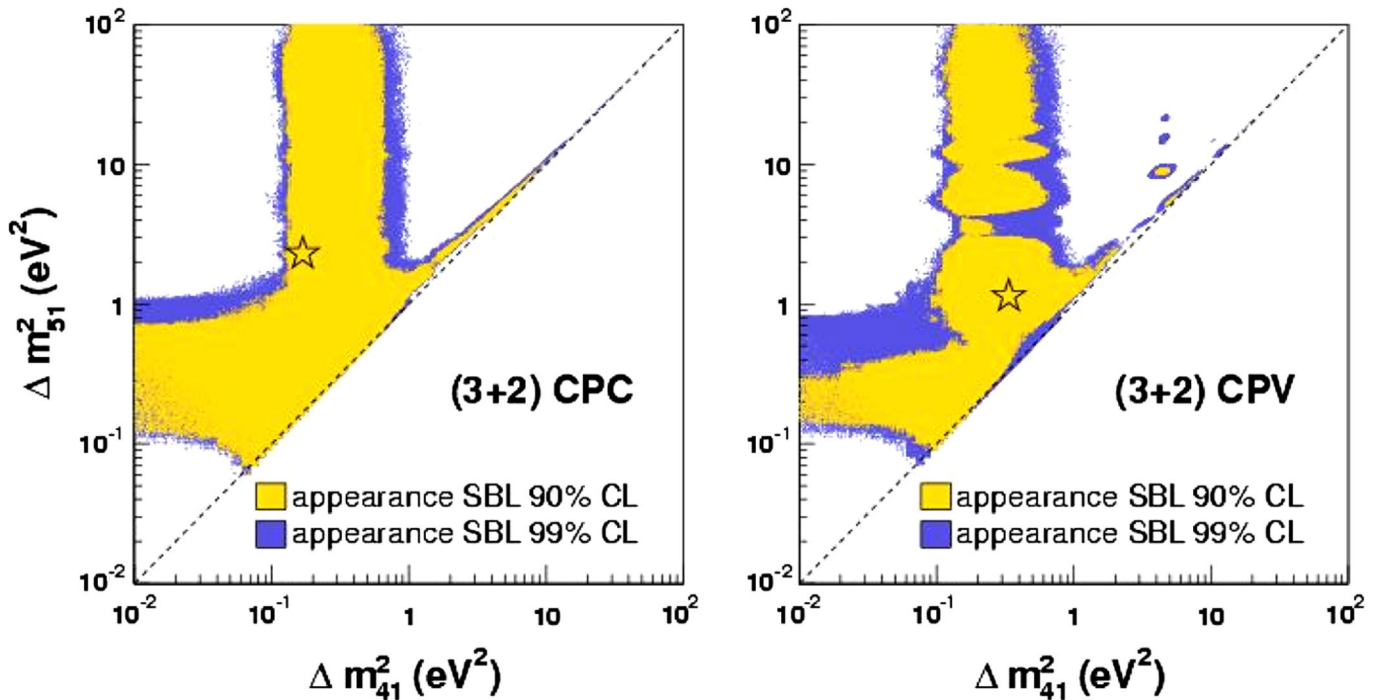


FIG. 11 (color online). Allowed regions in  $(\Delta m_{41}^2, \Delta m_{51}^2)$  space for fits to  $CP$ -conserving (CPC, left) and  $CP$ -violating (CPV, right)  $(3 + 2)$  oscillation models. Only appearance data sets have been included in the fit.

TABLE V. Comparison of best-fit values for mass splittings and mixing parameters for (3 + 2)  $CP$ -conserving (CPC) and  $CP$ -violating (CPV) models. Mass splittings are shown in  $\text{eV}^2$ . The appearance experiments include BNB-MB( $\nu$  and  $\bar{\nu}$ ), LSND, NUMI-MB, KARMEN, and NOMAD. The signal experiments include LSND, BNB-MB( $\nu$ ), and BNB-MB( $\bar{\nu}$ ). See the text for more details.

Data set	Fit	$\chi^2(\text{dof})$	$\chi^2$ -probability	$\Delta m_{41}^2$	$\Delta m_{51}^2$	$ U_{e4} $	$ U_{\mu 4} $	$ U_{e5} $	$ U_{\mu 5} $	$\phi_{45}$
signal APP	CPV	42.5(36)	21%	0.14	2.06	0.35	0.40	0.056	0.18	$1.7\pi$
signal APP	CPC	46.9(37)	13%	2.01	2.22	0.42	0.24	0.33	0.33	0
APP	CPV	92.6(85)	27%	0.28	1.12	0.44	0.20	0.21	0.094	$1.1\pi$
APP	CPC	95.8(86)	22%	0.18	2.31	0.32	0.38	0.086	0.071	0
all SBL	CPV	190.2(192)	52%	0.90	24.5	0.12	0.14	0.065	0.18	$1.8\pi$
all SBL	CPC	191.5(193)	52%	0.92	24.0	0.12	0.14	0.070	0.14	0

of BNB-MB( $\nu$ ) and NUMI-MB, the (3 + 2) CPV appearance-only best-fit distributions, shown in dashed red (dark gray) on the right, have comparable agreement with data as the (3 + 1) neutrino-only best-fit distributions, shown in solid blue (dark gray) on the left. However, in the case of BNB-MB( $\bar{\nu}$ ), the (3 + 2) CPV appearance-only best-fit parameters yield a worse  $\chi_{\text{BNB-MB}(\bar{\nu})}^2 = 21.4$  than the (3 + 2) CPV best-fit parameters from a fit to all SBL data sets,  $\chi_{\text{BNB-MB}(\bar{\nu})}^2 = 21.0$ . Both (3 + 2) CPV fits yield a worse  $\chi_{\text{BNB-MB}(\bar{\nu})}^2$  than the (3 + 1) antineutrino-only best-fit prediction ( $\chi_{\text{BNB-MB}(\bar{\nu})}^2 = 19.7$ ).

Figure 13 shows the projection of  $\Delta\chi^2 = \chi^2 - \chi_{\text{min}}^2$  as a function of the  $CP$ -violating phase  $\phi_{45}$  for the three fits discussed in this section: the appearance-only projection

is shown in the solid orange (light gray) line, the BNB-MB( $\nu$ ) + BNB-MB( $\bar{\nu}$ ) + LSND projection in dashed orange (light gray), and the fit with all SBL experiments is shown in blue (dark gray). Despite the fact that the signal experiments yield a best-fit  $\phi_{45}$  of  $1.1\pi$ , all three (3 + 2) CPV fits favor a similar phase range around  $\phi_{45} \sim 1.7\pi$ , as illustrated by the three overlapping dips in the  $\Delta\chi^2$  distribution.

## VI. CONSTRAINTS TO (3 + 2) $CP$ -VIOLATING FITS FROM EACH SBL EXPERIMENT

In this section we study the constraints to experimentally allowed (3 + 2)  $CP$ -violating oscillations by each of the SBL experiments. This is accomplished through a study

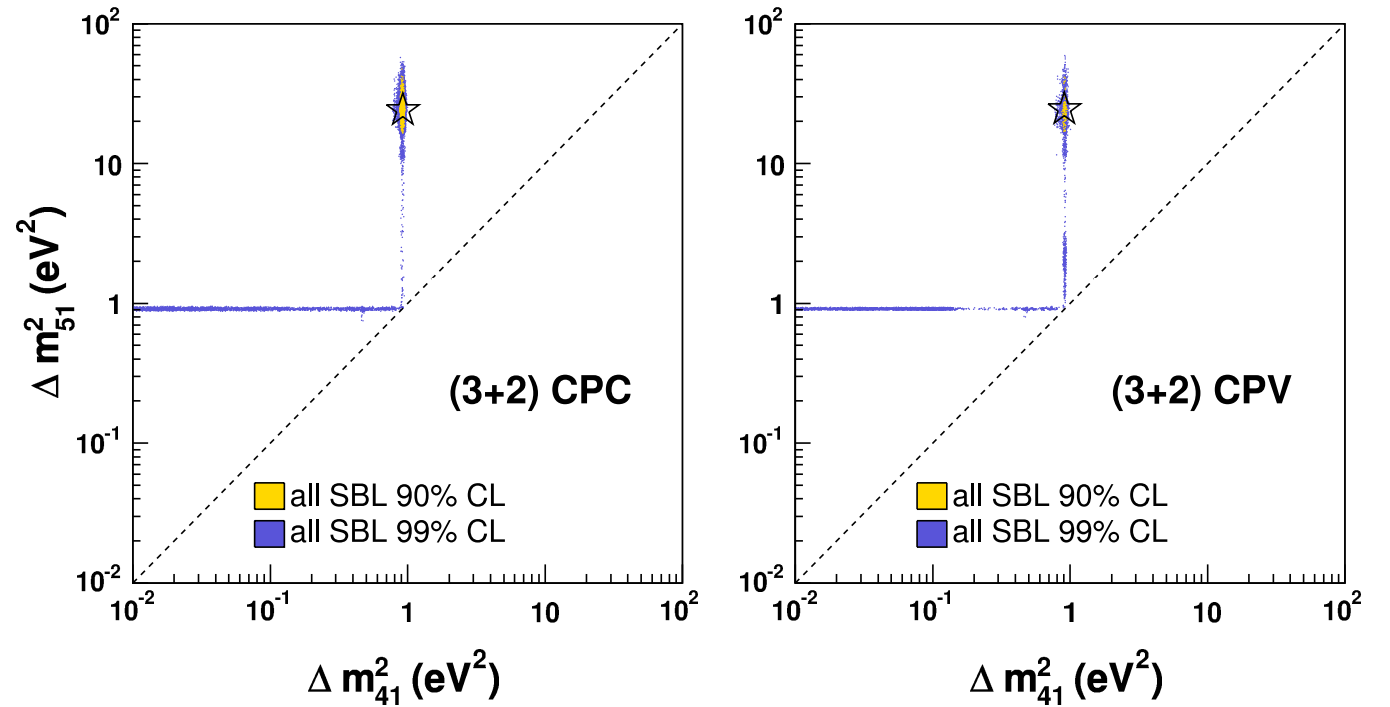


FIG. 12 (color online). Allowed regions in  $(\Delta m_{41}^2, \Delta m_{51}^2)$  space for fits to  $CP$ -conserving (CPC, left) and  $CP$ -violating (CPV, right) (3 + 2) oscillation models. All SBL data sets (appearance and disappearance) and atmospheric constraints have been included in the fit.



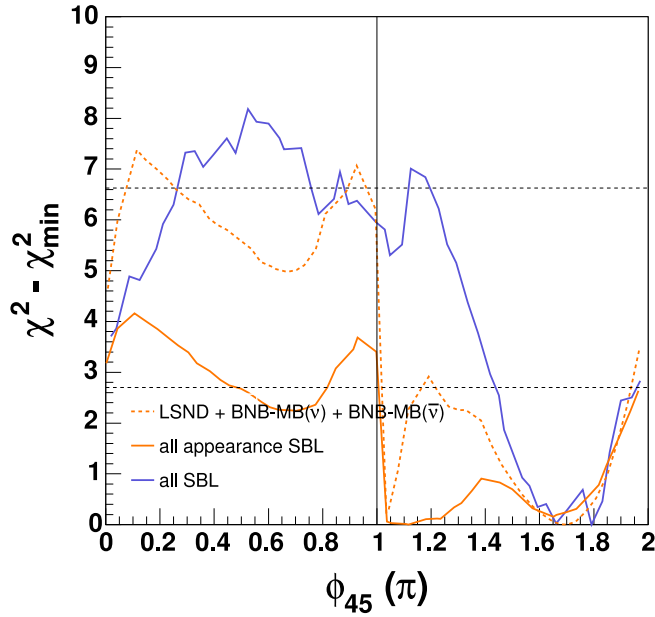


FIG. 13 (color online). Projection of  $\Delta\chi^2 = \chi^2 - \chi^2_{\min}$  as a function of the  $CP$ -violating phase  $\phi_{45}$ . The dashed horizontal lines indicate the 90% and 99% CL  $\Delta\chi^2$ .

where fits are performed using all-but-one experiment at a time. Within this study, we are also interested in examining the source of incompatibility between appearance and disappearance data, as well as testing compatibility between neutrino and antineutrino appearance search results within a  $CP$ -violating scenario. The latter is motivated by the larger incompatibility found in neutrino versus antineutrino fits, as opposed to appearance versus disappearance fits.

Table VI summarizes the  $\chi^2$ -probability and PG results from (3 + 2)  $CP$ -violating fits. The upper rows summarize  $\chi^2$ -probabilities and PG's from fits to all SBL experiments, as well as fits to appearance-only, disappearance-only, neutrino-only, antineutrino-only, neutrino appearance-only, and antineutrino appearance-only data sets. Appearance and disappearance data sets, as well as neutrino and antineutrino data sets, are incompatible with a PG of 0.26% or less. Grouping SBL appearance-only data sets according to whether they are neutrino or antineutrino experiments yields higher compatibilities—56% among  $\bar{\nu}$  appearance experiments, and 18% among  $\nu$  appearance experiments. However, the compatibility between  $\nu$  and  $\bar{\nu}$  appearance-only results is still low, at 2.2%. In the case where disappearance experiments are included in the comparison between neutrino and antineutrino fits, the compatibility of all  $\bar{\nu}$  SBL data sets remains reasonable, at 5.8%, but the compatibility among all  $\nu$  SBL data sets, and the compatibility between  $\bar{\nu}$  and  $\nu$  results are both  $<1\%$ .

The remaining rows of Table VI provide the  $\chi^2$ -probabilities of global fits under the same oscillation scenario where one experiment is excluded from the fit at a

time (as indicated by the “-” sign in the table). The  $\chi^2$  probabilities of all fits are acceptable, ranging between 25.7% for a fit excluding the Bugey data set, and 70.6% for a fit excluding the LSND data set. Aside from LSND, three experiments stand out as having poorest compatibility when compared to a global fit with all other SBL data sets: 1) BNB-MB( $\nu$ ), 2) CDHS and 3) atmospheric constraints (ATM). These three experiments have been identified as the possible source of tension among appearance and disappearance experiments. The remaining combinations yield reasonably high compatibility of at least 45%, with the exception of LSND and BNB-MB( $\bar{\nu}$ ) which are compatible with the remaining data sets at 2.9% and 20%, respectively.

To further test the hypothesis that the tension between appearance and disappearance experiments is a result of the BNB-MB( $\nu$ ) and CDHS data sets and atmospheric constraints, the compatibility between appearance and disappearance experiments is reevaluated several times. Each time, a different combination of these three experiments is excluded from the fits. The results are summarized in Table VII. The compatibility among appearance and disappearance experiments with BNB-MB( $\nu$ ), CDHS, and atmospheric constraints excluded from the fits is high, at 36.4%. The BNB-MB( $\nu$ ) data set alone is not responsible for the disagreement between appearance and disappearance experiments, as suggested by the sixth row of the Table VII. Even with BNB-MB( $\nu$ ) included in the fit, a compatibility of 15.7% can be obtained if CDHS and atmospheric constraints are excluded from the fit. Instead, the tension seems to exist mainly between BNB-MB( $\nu$ ) and CDHS, and BNB-MB( $\nu$ ) and atmospheric constraints.

The same test can be performed between neutrino and antineutrino experiments. The results are summarized in Table VIII. Again, the compatibility between neutrino and antineutrino experiments is reevaluated several times; each time, a different combination of the BNB-MB( $\nu$ ), CDHS, and atmospheric constraint data sets is excluded from the fits. Here, the compatibility among neutrino and antineutrino experiments with BNB-MB( $\nu$ ), CDHS, and atmospheric constraints excluded from the fits is even higher, at 44.0%. In this case, however, the BNB-MB( $\nu$ ) data set is just as responsible for the disagreement between neutrino and antineutrino experiments as the CDHS data set and atmospheric constraints alone. The tension seems to be caused by all three experiments.

It is possible that higher compatibility between BNB-MB( $\nu$ ) and all remaining SBL data sets may be achieved if the fits are to be repeated with the low-energy region ( $200 < E_\nu^{\text{QE}} < 475$  MeV) excluded from the BNB-MB( $\nu$ ) data set.

A global analysis with BNB-MB( $\nu$ ), CDHS, and atmospheric constraints excluded from the fit yields a  $\chi^2$ -probability of 83% but still relatively low compatibility, of 0.4%.

TABLE VI. Comparison of  $\chi^2$ -probabilities for  $(3 + 2)$   $CP$ -violating fits with different combinations of SBL data sets. Also shown are PG results testing compatibility among different data sets. The last 11 rows of the table provide the compatibility (PG) between the experiment being removed from each fit and all remaining experiments. See the text for more details.

Data set	$\chi^2(\text{dof})$	$\chi^2$ -probability (%)	PG (%)	
all SBL	190.2 (192)	52.3	PG(BNB-MB( $\nu$ ), BNB-MB( $\bar{\nu}$ ), NUMI-MB, LSND, KARMEN, NOMAD, Bugey, CHOOZ, CCFR84, CDHS, ATM)	= Prob(49.1, 5) = $2.2 \times 10^{-7}$
			PG(APP, DIS)	= Prob(16.3, 4) = 0.26
APP	92.6 (85)	26.9	PG( $\nu$ , $\bar{\nu}$ )	= Prob(21.0, 5) = $8.2 \times 10^{-2}$
			PG(BNB-MB( $\nu$ ), BNB-MB( $\bar{\nu}$ ), NUMI-MB, LSND, KARMEN, NOMAD)	= Prob(24.6, 5) = $1.7 \times 10^{-2}$
DIS	81.3 (103)	94.4	PG(Bugey, CHOOZ, CCFR84, CDHS, ATM)	= Prob(8.14, 4) = 8.6
$\nu$	86.1 (86)	47.8	PG(BNB-MB( $\nu$ ), NUMI-MB, NOMAD, CCFR84, CDHS, ATM)	= Prob(17.4, 5) = 0.37
$\bar{\nu}$	83.2 (99)	87.3	PG(BNB-MB( $\bar{\nu}$ ), KARMEN, LSND, Bugey, CHOOZ)	= Prob(10.7, 5) = 5.8
$\nu$ APP	50.6 (53)	56.8	PG(BNB-MB( $\nu$ ), NUMI-MB, NOMAD)	= Prob(3.91, 5) = 56
$\bar{\nu}$ APP	28.9 (27)	36.7	PG(BNB-MB( $\bar{\nu}$ ), KARMEN, LSND)	= Prob(7.55, 5) = 18
			PG( $\nu$ APP, $\bar{\nu}$ APP)	= Prob(13.1, 5) = 2.2
all-BNB-MB( $\nu$ )	167.2 (174)	63.2	PG(all SBL-BNB-MB( $\nu$ ), BNB-MB( $\nu$ ))	= Prob(12.1, 5) = 3.4
all-BNB-MB( $\bar{\nu}$ )	168.5 (174)	60.3	PG(all SBL-BNB-MB( $\bar{\nu}$ ), BNB-MB( $\bar{\nu}$ ))	= Prob(7.34, 5) = 20
all-NUMI-MB	184.2 (182)	43.9	PG(all SBL-NUMI-MB, NUMI-MB)	= Prob(4.34, 5) = 50
all-LSND	176.1 (187)	70.6	PG(all SBL-LSND, LSND)	= Prob(12.5, 5) = 2.9
all-KARMEN	180.2 (183)	54.6	PG(all SBL-KARMEN, KARMEN)	= Prob(4.72, 5) = 45
all-NOMAD	154.2 (162)	65.6	PG(all SBL-NOMAD, NOMAD)	= Prob(1.89, 5) = 86
all-Bugey	142.2 (132)	25.7	PG(all SBL-Bugey, Bugey)	= Prob(3.07, 4) = 55
all-CHOOZ	180.9 (178)	42.6	PG(all SBL-CHOOZ, CHOOZ)	= Prob(3.06, 4) = 55
all-CCFR84	174.7 (174)	47.0	PG(all SBL-CCFR84, CCFR84)	= Prob(0.82, 4) = 94
all-CDHS	175.4 (177)	51.9	PG(all SBL-CDHS, CDHS)	= Prob(7.48, 4) = 11
all-ATM	184.4 (190)	60.0	PG(all SBL-ATM, ATM)	= Prob(5.78, 2) = 5.6

## VII. CONCLUSIONS

We have reexamined global fits to sterile neutrino oscillation models, using new data from MiniBooNE. Those include the final MiniBooNE neutrino mode results and the first, low statistics MiniBooNE antineutrino results, as well

 TABLE VII. Comparison of compatibility between appearance (APP) and disappearance (DIS) experiments, within a  $(3 + 2)$   $CP$ -violating scenario. The BNB-MB( $\nu$ ) data set, CDHS data set, and atmospheric constraints (ATM) are removed from the fits as specified in order to establish the source of tension between appearance and disappearance experiments. See the text for more details.

Data sets	PG (%)
APP vs DIS	0.26
APP(no BNB-MB( $\nu$ )) vs DIS(no CDHS + ATM)	36.4
APP(no BNB-MB( $\nu$ )) vs DIS(no CDHS)	0.58
APP(no BNB-MB( $\nu$ )) vs DIS(no ATM)	0.82
APP(no BNB-MB( $\nu$ )) vs DIS	0.10
APP vs DIS(no CDHS + ATM)	15.7
APP vs DIS(no CDHS)	0.76
APP vs DIS(no ATM)	1.4

as first results from the off-axis NuMI beam observed in the MiniBooNE detector.

Within a  $(3 + 1)$   $CP$ - and  $CPT$ -conserving scenario, we have found that the data set collected by MiniBooNE using the NuMI off-axis beam (NUMI-MB) currently provides very weak constraints to sterile neutrino fits, due to large

 TABLE VIII. Comparison of compatibility between neutrino ( $\nu$ ) and antineutrino ( $\bar{\nu}$ ) experiments, within a  $(3 + 2)$   $CP$ -violating scenario. The BNB-MB( $\nu$ ) data set, CDHS data set, and atmospheric constraints (ATM) are removed from the fits as specified in order to establish the source of tension between neutrino and antineutrino experiments. See the text for more details.

Data sets	PG (%)
$\nu$ vs $\bar{\nu}$	$8.2 \times 10^{-2}$
$\nu$ (no BNB-MB( $\nu$ ) + CDHS + ATM) vs $\bar{\nu}$	44.0
$\nu$ (no BNB-MB( $\nu$ ) + CDHS) vs $\bar{\nu}$	1.6
$\nu$ (no BNB-MB( $\nu$ ) + ATM) vs $\bar{\nu}$	1.8
$\nu$ (no BNB-MB( $\nu$ )) vs $\bar{\nu}$	0.36
$\nu$ (no CDHS + ATM) vs $\bar{\nu}$	0.70
$\nu$ (no CDHS) vs $\bar{\nu}$	0.68
$\nu$ (no ATM) vs $\bar{\nu}$	0.36

systematic uncertainties. Updated NuMI results, expected soon, should have a greater impact on these fits. Within the same oscillation framework, the MiniBooNE antineutrino data set (BNB-MB( $\bar{\nu}$ )) is found in agreement with LSND, yielding, in a combined analysis with LSND and KARMEN under a  $(3 + 1)$  oscillation hypothesis, a  $\chi^2$ -probability of 29%, and best-fit parameters similar to those of LSND. Updated MiniBooNE antineutrino appearance results, with almost twice the current statistics are expected in the near future. The MiniBooNE neutrino data set (BNB-MB( $\nu$ )), although suggestive of an excess that could be described by a  $(3 + 1)$  oscillation hypothesis with a  $\chi^2$  probability of 35%, is found incompatible with the signals from the MiniBooNE antineutrino and LSND results. The remaining null appearance and disappearance experiments (BNB-MB( $\nu$ ), KARMEN, NOMAD, Bugey, CHOOZ, CDHS, CCFR84) and atmospheric oscillation data impose strong constraints to the parameter space allowed by a combined analysis of MiniBooNE neutrino and antineutrino and LSND data, excluding the 99% CL allowed region at 99% CL. However, the constraints from antineutrino disappearance experiments on the parameter space allowed by antineutrino appearance experiments (BNB-MB( $\nu$ ), LSND, and KARMEN) are weaker. All antineutrino experiments yield a best-fit  $\chi^2$ -probability of 86%, and exclude the no-oscillations hypothesis at  $>5.0\sigma$ . The best-fit parameters are similar to those of LSND, and correspond to a muon antineutrino disappearance amplitude of 0.35, which may be addressed by upcoming results from MiniBooNE and MINOS on muon antineutrino disappearance. Additionally, fits to all neutrino experiments yield a best-fit  $\chi^2$ -probability of 47% and exclude the null hypothesis at  $>90\%$  CL.

Furthermore, we find that with the addition of the new MiniBooNE data sets, the  $(3 + 2)$  oscillation models provide only a marginally better description of all SBL data

sets compared to  $(3 + 1)$  models. In the case of  $(3 + 2)$  fits,  $CP$  violation only allows for a small improvement in  $\chi^2$ -probability for fits to only BNB-MB( $\nu$ ) + BNB-MB( $\bar{\nu}$ ) + LSND, and fits to only appearance experiments. In the case of global fits, the  $\chi^2$ -probability for the best-fit hypothesis is 52%. The best-fit corresponds to large but not maximal  $CP$  violation ( $\phi_{45} = 1.8\pi$ ). The high incompatibility among appearance and disappearance data remains after the addition of the new MiniBooNE results. This incompatibility is a result of the BNB-MB( $\nu$ ) and CDHS data sets and atmospheric constraints. The compatibility between appearance and disappearance data with these three experiments excluded from the fits is significantly higher, at 36%.

Neutrino and antineutrino data sets are also incompatible within a  $(3 + 2)$   $CP$ -violating scenario, with a PG of less than 0.1%. The compatibility improves slightly to 2.2% in the case of comparing appearance-only neutrino and antineutrino data sets.

Overall, allowing for mixing with multiple sterile neutrino states or  $CP$  violation does not seem sufficient to allow incorporating all SBL experiments within a  $CPT$ -conserving, sterile neutrino oscillation framework. It may be that there is an issue with one or more of the following data sets: LSND, BNB-MB( $\nu$ ), including the low-energy excess, CDHS, or atmospheric constraints; alternatively, theories with  $CPT$ -violating oscillations or effective  $CPT$  violation [36–38] may succeed in reconciling all short-baseline oscillation signatures, and should be explored.

## ACKNOWLEDGMENTS

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