XIX. SPONTANEOUS RADIOFREQUENCY EMISSION FROM HOT ELECTRON PLASMAS^{*}

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RESEARCH OBJECTIVES

The production of a hot plasma usually results in an anisotropic velocity distribution function for the constituent particles. Such distribution functions are known to be unstable (velocity-space instabilities), and experimentally are often identified by the radiofrequency emissions that they produce.

The pulsed electron-cyclotron discharge, which we have studied in the past, produces a hot electron plasma. In the "afterglow" (in between the applied microwave pulses) we have observed strong emission near the electron-cyclotron frequency and its harmonics.

We plan to continue our theoretical studies on velocity-space instabilities, and the experimental study of the enhanced cyclotron emission observed in the afterglow of our electron-cyclotron discharge.

A. Bers

A. ELECTROSTATIC PLASMA INSTABILITIES AT CYCLOTRON HARMONICS

We have continued to study instabilities in a neutral unbounded magnetized plasma composed of infinitely massive ions and electrons that are assumed to have the relativistic velocity distribution

 $f_0(v) = \delta(v_{\perp} - v_{\perp 0}) \delta(v_{\parallel})/2\pi v_{\perp 0},$

in which the subscripts \perp and \parallel stand for perpendicular and parallel to the static magnetic field. Under the assumption of $\exp\{j\omega t - j\overline{k} \cdot \overline{r}\}$ dependence for the field quantities, a simultaneous solution of the linearized Vlasov equation and Maxwell's equations has been shown to result in the electrostatic (\overline{k} along \overline{E}) dispersion relation

$$\left(\frac{k_{\parallel}}{k_{\perp}}\right)^{2} = -\frac{1 - \frac{1}{2}\left(\frac{\omega_{p}}{\omega_{b}}\right)^{2} \sum_{n=-\infty}^{\infty} \frac{J_{n-1}^{2} - J_{n+1}^{2}}{\left(\frac{\omega}{\omega_{b}} - n\right)^{2}} + \left(\frac{\beta_{\perp}\omega_{p}}{\omega_{b}}\right)^{2} \sum_{n=-\infty}^{\infty} \frac{n^{2}J_{n}^{2}}{\left(\frac{\omega}{\omega_{b}} - n\right)^{2}},$$

$$\left(\frac{k_{\parallel}}{k_{\perp}}\right)^{2} = -\frac{1 - \left(\frac{\omega_{p}}{\omega_{b}}\right)^{2} \sum_{n=-\infty}^{\infty} \frac{J_{n}^{2}}{\left(\frac{\omega}{\omega_{b}} - n\right)^{2}}}{1 - \left(\frac{\omega_{p}}{\omega_{b}}\right)^{2} \sum_{n=-\infty}^{\infty} \frac{J_{n}^{2}}{\left(\frac{\omega}{\omega_{b}} - n\right)^{2}},$$

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where ω_p and ω_b are the relativistic plasma and cyclotron frequencies, respectively, β_\perp is the ratio of the electron velocity to the speed of light in free space, and J_n is the n^{th} -order Bessel function of argument $p = k_\perp v_{\perp 0}/\omega_b$.¹ The nonrelativistic form of this equation, which is obtained by setting $\beta_\perp = 0$, has been studied by Bers and Gruber² who investigated the small-signal energy implications of the resulting waves, and by Gruber, Klein, and Auer³ who showed that for large enough ω_p/ω_b waves propagating nearly along the static magnetic field were unstable. Dory, Guest, and Harris⁴ have claimed that the nonrelativistic equation predicts that waves propagating perpendicular to the magnetic field are stable unless ω_p/ω_b exceeds 2.7. As we have shown in a previous report,⁵ these instabilities occur near $\frac{\omega}{\omega_b} = n + \frac{1}{2}$, where n is an integer. Our present work, including the relativistic term, predicts unstable waves propagating nearly perpendicular to the magnetic field at values of ω_p/ω_b considerably less than the threshold stated by other authors. In contrast to the previous reported instabilities, these new modes, which are purely relativistic in origin, occur at frequencies close to integer multiples of the

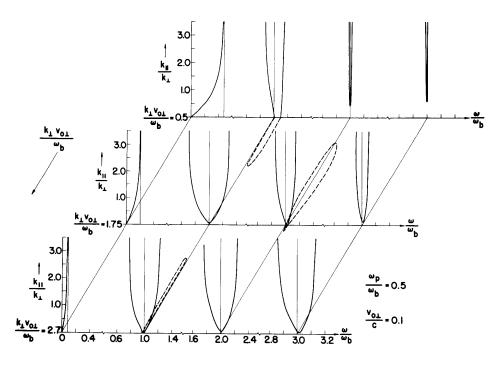


Fig. XIX-1. Dispersion diagram.

cyclotron frequency. Figure XIX-1 illustrates the results of a numerical analysis of the dispersion relation. The curves were obtained for $\omega_p/\omega_b = 0.5$ and $v_{0\perp}/c = 0.1$. Note that near each cyclotron harmonic there are two branches of the dispersion relation. At a given value of $k_{\perp}v_{0\perp}/\omega_b$, the curves are similar to those reported by Bers and Gruber,

except that near some cyclotron harmonics the two branches comprising the particular mode have coupled so as to require complex frequency for all k_{\parallel}/k_{\perp} less than a critical value. The figure illustrates that the particular modes that are unstable depend upon the value of $k_{\perp}v_{0\perp}/\omega_{\rm b}$. For convenience, the root locus of purely real roots has been drawn on the $k_{\parallel}/k_{\perp} = 0$ plane. We have reported a more complete analysis of these roots previously.⁵

Physically, the instabilities have been interpreted in terms of small-signal energy arguments. Evaluation of the small-signal time-average energy density⁶ has shown that the lower frequency wave at each uncoupled harmonic mode carries negative energy, whereas the higher frequency wave carries positive energy. The instabilities occur when the plasma parameters are such as to result in an internal coupling of the two waves.

We are now investigating an exact electromagnetic solution of the waves in this plasma and are considering the physical mechanism that results in the relativistic instability.

C. Speck, A. Bers

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