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## High-Power, Few-Cycle, Phase-Stabilized 2.2-µm Optical **Parametric Chirped Pulse Amplifier**

Shu-Wei Huang<sup>1</sup>, Jeffrey Moses<sup>1</sup>, Kyung-Han Hong<sup>1</sup>, Edilson L. Falcão-Filho<sup>1</sup>, Andrew Benedick<sup>1</sup>, Jeremy Bolger<sup>2</sup>, Benjamin Eggleton<sup>2</sup>, and Franz X. Kärtner<sup>1</sup>

> <sup>1</sup> Department of Electrical Engineering and Computer Science and Research Laboratory of Electronics, Massachusetts Institute of Technology, Cambridge, MA 02139, USA <sup>2</sup>CUDOS ARC Centre of Excellence School of Physics, University of Sydney, Sydney, Australia klkla@mit.edu

Abstract: We demonstrate a high-peak-power, 1-kHz, 2.2-µm OPCPA for long-wavelengthdriven high harmonic generation that produces 9-GW, 3-optical-cycle, CEP-stabilized pulses, allowing tunneling-ionization-threshold intensity with low Guov phase shift. ©2009 Optical Society of America

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Since the prediction of high-yield soft X-ray photon generation through high harmonic generation (HHG) with longwavelength drive pulses [1, 2], the development of high-power, few-cycle, carrier-envelope phase- (CEP-) stabilized sources in the infrared has attracted great attention. Due to the unavailability of broadband laser gain media in this wavelength range, this technology has widely relied upon some combination of parametric frequency conversion and nonlinear spectral broadening, using a broadband Ti:sapphire amplifier as a front end. Such techniques work well but are ultimately limited in average power by the Ti:sapphire amplifier, and the external compression technique based on either hollow-core fiber or filamentation in gas as a final step is inherently challenging to scale to the multi-mJ energy level [3, 4]. In contrast, ultra-broadband optical parametric chirped pulse amplification (OPCPA) provides a clear route to overcoming this limitation by direct amplification of a few-cycle seed pulse [5-8]. There is a distinct advantage to building an OPCPA system for driving HHG centered around 2 µm, as opposed to other mid-IR wavelengths: high-average-power ultra-broadband amplification is feasible because it is degenerate with the highest average power pump sources available to date [9, 10] and the zero-dispersion wavelength of standard nonlinear crystals such as lithium niobate and lithium tantalate is around 2 µm.

In this paper, we report the development of a 2.2-um OPCPA which generates 9-GW, 3-optical-cycle pulses (200 µJ, 23 fs) at 1 kHz. It is, to date, the highest peak power for few-cycle pulses at 2-µm wavelength obtained by direct parametric amplification. Implementing several methods of suppressing superfluorescence noise introduced in our previous work [8], our results verify the prediction of suitable noise suppression at the multi-100-µJ level, surpassing the 80-µJ level demonstrated previously [7]. Scaling to multi-100-µJ energy is an important and necessary step for efficient HHG; below 100 µJ, even few-cycle pulses require tight focal volumes with high Guoy phase shift (which scales with the inverse of the beam waist squared), impeding good HHG phase matching. With 200  $\mu$ J and a loose focus of ~50  $\mu$ m waist, allowing reasonable phase matching, the intensity of our source exceeds  $10^{14}$  W/cm<sup>2</sup> and we have clearly observed a ~2-mm-long plasma column in air. Our current results suggest that our OPCPA design can be scaled to the multi-mJ energy level while maintaining low noise, which will allow absorption-limited HHG. Such steps are valuable to the pursuit of the next generation of high-flux, HHG-based soft X-ray sources.

The schematic of our OPCPA system is shown in Fig. 1. An octave-spanning Ti:sapphire oscillator is used to generate passively CEP-stabilized broadband seed pulses by intra-pulse DFG. We have developed a Nd:YLF chirped pulse amplifier (CPA) as a pump source, which consists of a chirped fiber Bragg grating (CFBG) stretcher, a regenerative amplifier followed by two multi-pass slab amplifiers, and a diffraction grating compressor. It is seeded



Figure 1. Schematic of the OPCPA system

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by the 1047-nm component of the residual Ti:sapphire oscillator spectrum and provides 12-ps, 4-mJ pump pulses for the OPCPA stages. The OPCPA seed pulse from DFG is first stretched by a bulk silicon block to 5 ps and preamplified in OPA1 to  $1.5 \,\mu$ J. The pre-amplified pulse is then further stretched to  $9.5 \,ps$  by an infrared acousto-optic programmable dispersive filter (AOPDF: DAZZLER, Fastlite), allowing for simultaneous optimization of energy extraction, amplification bandwidth, and signal-to-noise ratio in the power amplifier stage [11]. Losses from the AOPDF (~90%) and spatial filters are compensated by a second preamplifier, OPA2. Finally, the pulse is amplified to 220 µJ in OPA3 and compressed in a broadband anti-reflection coated quartz glass block (Suprasil 300), which introduces ~10% loss. The AOPDF also compensates for the residual dispersion from the bulk stretcher/compressor unit, allowing the pulse duration to reach near transform limit after the bulk compressor. Fig. 2 shows the amplified spectrum (a) and the corresponding interferometric autocorrelation trace (b) of the OPCPA system. The pulse is compressed nearly to its transform limit, i.e., 23 fs, or 3 cycles in FWHM.

Due to the low available seed energy from the DFG process, superfluorescence suppression plays a key role in the pulse energy scalability of the OPCPA. Four design features are implemented in our setup to suppress superfluorescence. First, our stretcher/compressor scheme maximizes the initial signal-to-noise ratio by avoiding placement of the lossy AOPDF before the amplifier chain. Second, we implement both hard and soft apertures after OPA1 to filter out superfluorescence-dominated high-order spatial modes. An aperture also eliminates spatial chirp. a side-effect of the AOPDF. Third, we use a seed with a spectrum spanning 1.6-2.5 µm, which fully covers the amplifier gain bandwidth. Fourth, we have carefully tailored the signal pulse durations in each amplification stage to fully seed the temporal gain profiles of the amplifiers while maintaining high conversion efficiency and wide signal bandwidth. We observe a clean and stable amplified signal spectrum as well as spatial profile. The energy fluctuation of amplified pulses is as low as 1.5% (rms). We predict the system should be scalable to the multi-mJ level without the superfluorescence overtaking the signal.



Figure 2. (a) Amplified spectrum, measuring 500-nm FWHM bandwidth. (b) IAC of the compressed pulse, measuring 23 fs (3 cycles). (c) f-3f spectral interferogram, measuring 150 mrad rms CEP fluctuation.

The CEP stability is characterized using an f-to-3f spectral interferometer with a sapphire plate and a BBO crystal [7]. Fig. 2(c) shows that the rms phase fluctuation is  $\sim$ 150 mrad over 10 s, where the residual phase excursion at time =  $\sim 2$  s is attributed to the amplitude-phase noise coupling in the *f*-to-3*f* interferometer while any significant drift is not observed during 10 s.

In conclusion, we have demonstrated a 9-GW, 3-optical-cycle (200 µJ, 23 fs), phase-stabilized 2.2-µm OPCPA system with suppressed superfluorescence. Scaling up to multi-100-µJ is an important and necessary step for efficient HHG because tight focusing to a few tens of microns is not required. Our current results suggest that scaling should be possible to the multi-mJ level, which will allow absorption-limited HHG.

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