REFRIGERANT FORCED- CONVECTION **CONDENSATION** INSIDE HORIZONTAL **TUBES**

Soonhoon Bae John **S.** Maulbetsch Warren M. Rohsenow

Report No. DSR **72591-71**

American Society of Heating, Refrigerating and Air Conditioning Engineers Contract No. ASHRAE RP **⁶³**

Engineering Projects Laboratory Department of Mechanical Engineering Massachusetts Institute of Technology
Cambridge, Massachusetts 02139 Cambridge, Massachusetts **02139**

December **1, 1970**

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Sponsored **by:-**

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Soonhoon Bae, York Division, Borg-Warner Corporation John **S.** Maulbetsch, Dynatech Corporation Warren M. Rohsenow, Massachusetts Institute of Technology

ABSTRACT

High vapor velocity condensation inside a tube was studied theoretically. The heat transfer coefficients were calculated **by** the momentum and heat transfer analogy. The Von Karman universal velocity distribution was applied to the condensate flow. Pressure drop was calculated **by** the Lockhart-Martinelli method and the Zivi void fraction equation.

Experimental data was obtained for the mass velocities from **150,000** to $555,000$ lbm/ft² hr for R-12 and R-22 condensing in a 0.493 " I.D. **18** ft. long test section. The measured heat transfer coefficients agreed with the prediction within **10%** except a few points in the very low quality region.

INTRODUCTION

When condensation takes place inside a horizontal tube with high vapor velocity, condensate flows in an annular shape on the tube wall and vapor flows in the core.

Many investigators have studied this subject both experimentally and analytically. Empirical correlations involving non-dimensional groups were not quite successful because the correlations did not include all the flow variables **[1],** [2], **[5], [7],** [12]. Carpenter and Colburn **[6]** considered only the laminar sublayer of condensate flow and derived an equation with an empirical constant. This method was modified **by** later investigators **[3], [15].** For a small range of the Prandtl Number, this equation gives good agreement with empirical data. But the equation has no general applicability. Rohsenow, Webber and Ling [14] analyzed the liquid film on the vertical plate and the heat transfer coefficient was obtained **by** the heat and momentum transfer analogy. **A** similar approach appeared in later papers [8], **[9].** This method will be developed further for the annular flow regime in this paper.

THEORY

Flow Model

For condensation inside a horizontal tube with high vapor velocities, annular flow is the predominent flow pattern and slug flow may appear at very low vapor qualities. Annular flow with a uniform liquid layer thickness around the circumference of a tube is assumed to exist in the parameter ranges of interest. The condensate accumulation at the bottom of a horizontal tube has a negligible effect except at very low vapor flow rates. At very high vapor flow rates entrainment

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of liquid droplet in the vapor may occur but this will be neglected in the following analysis.

Consider a length element of a tube, shown in Fig. **1.** For the entire cross-section the momentum equation is

$$
-(\frac{dP}{dz})A_{z} - T_{0}S + A_{z} \frac{a}{g_{0}} [\alpha \rho_{v} + (1 - \alpha)\rho_{\ell}] = \frac{1}{g_{0}} \frac{d}{dz} (U_{v}W_{v} + U_{\ell}W_{\ell}) \quad (1)
$$

where a is an acceleration due to the external body force. Rearranging **Eq. (1)** yields

$$
-(\frac{dP}{dz}) = \tau_o \frac{S}{A_z} - \frac{a}{g_o} [\alpha \rho_v + (1 - \alpha) \rho_g] + \frac{1}{g_o A_z} \frac{d}{dz} (U_v W_v + U_g W_g) (2)
$$

The above equation shows that the total static pressure gradient is the sum of pressure gradients due to friction,gravity and momentum change.

$$
\left(\frac{\text{dP}}{\text{d}z}\right) = \left(\frac{\text{dP}}{\text{d}z}\right)_{\text{f}} + \left(\frac{\text{dP}}{\text{d}z}\right)_{\text{g}} + \left(\frac{\text{dP}}{\text{d}z}\right)_{\text{m}}
$$
\n(3)

Comparing **Eq.** (2) and **(3),**

$$
\left(\frac{\mathrm{dP}}{\mathrm{d}z}\right)_{\mathrm{f}} = -\tau_{\mathrm{o}} \frac{\mathrm{S}}{\mathrm{A}_{z}}
$$
 (4)

$$
\left(\frac{dP}{dz}\right)_g = \frac{a}{g_o} \left[\alpha \rho_v + (1-\alpha)\rho_g\right]
$$
 (5)

$$
\left(\frac{dP}{dz}\right)_m = -\frac{1}{g_o A_z} \frac{d}{dz} (U_v W_v + U_g W_g)
$$
 (6)

The friction pressure drop was obtained **by** an approximation of the Lockhart-Martinelli method **(11]** as follows:

$$
\left(\frac{dP}{dz}\right)_{f} \frac{g_{o}D}{c^{2}/\rho_{v}} = -\tau_{o} \frac{S}{A_{z}} \frac{g_{o}D}{c^{2}/\rho_{v}} = -0.09 \frac{GD}{\mu_{v}} \frac{-0.2}{x} \left[x^{1.8} + \frac{0.0523}{x}\right]
$$

+ 5.7
$$
\left(\frac{\mu_{\ell}}{\mu_{v}}\right)^{0.0523} (1 - x) \frac{0.470}{x} \frac{1.33}{x} \left(\frac{\rho_{v}}{\rho_{\ell}}\right)^{0.261}
$$

+ 8.11
$$
\left(\frac{\mu_{\ell}}{\mu_{v}}\right)^{0.105} (1 - x) \frac{0.94}{x} \left(\frac{\rho_{v}}{\rho_{\ell}}\right)^{0.522} (1 - x) \frac{0.94}{x} \frac{0.86}{x} \left(\frac{\rho_{v}}{\rho_{\ell}}\right)^{0.522} (7)
$$

This empirical approximation was developed **by** Soliman et al **[15]** who used this equation to calculate τ_{α} . Here it is used to calculate **T0** as Lockhart-Martinelli suggest.

The gravity term, **Eq. (5),** can be rewritten in the following form:

$$
\left(\frac{\mathrm{d}P}{\mathrm{d}z}\right)_{g} \frac{\mathrm{g}_{o}D}{\mathrm{g}^{2}/\rho_{v}} = \frac{1}{\mathrm{Fr}^{2}} \left[\frac{\rho_{g}}{\rho_{v}} - \mathrm{B}\alpha\right]
$$
 (8)

where

$$
Fr^{2} = \frac{(G/\rho_{v})^{2}}{aD}
$$
 (9)

is the Froude number based on the total flow and

$$
B = \frac{\rho_{\ell} - \rho_{\mathbf{v}}}{\rho_{\mathbf{v}}} \tag{10}
$$

is the buoyancy modulus. In the gravity field

$$
a = g \sin \theta \tag{11}
$$

The Zivi equation for local void fraction **[17]** is recommended for use in **Eq. (8),** as in reference **[15].**

$$
\alpha = \frac{1}{1 + \left(\frac{1-x}{x}\right)\left(\frac{\rho_v}{\rho_{\chi}}\right)^{2/3}}
$$
(12)

The momentum term, **Eq. (6),** can be written with **Eq.** (12) as follows:

$$
\left(\frac{dP}{dz}\right)_{m} \frac{g_{o}D}{c^{2}/\rho_{v}} = -D\left(\frac{dx}{dz}\right) \left[2x + (1 - 2x)\left(\frac{\rho_{v}}{\rho_{g}}\right)\right]^{1/3}
$$

$$
+ (1 - 2x)\left(\frac{\rho_{v}}{\rho_{g}}\right)^{2/3} - 2(1 - x)\left(\frac{\rho_{v}}{\rho_{g}}\right)1
$$
(13)

where in **Eq. (6)**

$$
W_{\mathbf{v}} = GA_{\mathbf{z}} \mathbf{x} = A_{\mathbf{z}} \alpha U_{\mathbf{v}} \rho_{\mathbf{v}}
$$

(14)

$$
W_{\mathbf{v}} = GA_{\mathbf{z}} (1 - \mathbf{x}) = A_{\mathbf{z}} (1 - \alpha) U_{\mathbf{v}} \rho_{\mathbf{v}}
$$

The momentum equation for the entire liquid layer element, Fig.2, is

$$
-\left(\frac{dP}{dz}\right) A_{z} + \tau_v S_v - \tau_o S + \frac{a}{g} \rho_{\ell} A_{z} \n= \frac{1}{g} \left[\frac{d (U_{\ell} W_{\ell})}{dz} - U_i \frac{d W_{\ell}}{dz} \right]
$$
\n(15a)

or

$$
T_o = F_o \frac{A_z}{S} + T_v \frac{S_v}{S}
$$
 (15b)

where

$$
F_o = -\frac{dP}{dz} + \frac{a}{g_o} \rho_\ell - \frac{1}{g_o A_{z_\ell}} \left[\frac{d(U_\ell W_\ell)}{dz} - U_i \frac{dW_\ell}{dz} \right]
$$
(16)

Since for most of the tube length the liquid film is thin, a simple flat plate analysis will suffice for the heat transfer coefficient derivation.

The A_z/S
$$
\approx
$$
 δ and S_v/S \approx 1; so Eq. (15b) becomes
\n
$$
\tau_o = F_o \delta + \tau_v
$$
\n(17)

The quantity F_{o} may be expressed in terms of x and α by substituting Eqs. (12) and (14) into **Eq. (16).** Further from the universal velocity, Eq. (B-1), distribution $(U_i/U_{\ell}) = \beta$ may be obtained as a unique function of **6+** as shown in Fig. **3.** Then **Eq. (17)** becomes

$$
F_o = -\left(\frac{dP}{dz}\right) + \frac{a}{g_o} \rho_g - \frac{G^2}{g_o \rho_v} \frac{dx}{dz} \left[\frac{1}{1 - \alpha} \left(\frac{\rho_v}{\rho_g}\right)^{1/3} - \left(\frac{(1 - x)(2 - \beta)}{(1 - \alpha)^2}\right) \left(\frac{\rho_v}{\rho_g}\right)\right]
$$
(18)

where $\beta \equiv U_i/U_{\ell}$ is given by Fig. 3.

To determine the heat transfer coefficient it is assumed that the Karman momentum-heat transfer analogy analysis is applicable in the liquid layer. Then

$$
\tau = \frac{\rho_{\ell}}{g_o} \left(v_{\ell} + \varepsilon_m \right) \frac{dv_z}{dy}
$$
 (19a)

$$
q/A = \rho_{\ell} c_{\ell} \left(q_{\ell} + \epsilon_{h} \right) \frac{dT}{dy}
$$
 (19b)

The universal velocity distribution is used to determine dv_z/dy , ϵ_m is assumed equal to ϵ_h , and dT/dy and T(y) is determined by combining the above two equations. The procedure is identical with that presented in detail **by** Rohsenow, Webber and Ling [14] with two exceptions. In determining this temperature distribution, the momentum equation for the element **(6-y)** of Fig. 2 was approximated as

$$
\tau = F_0 (\delta - y) + \tau_v
$$

and no criterion for transition from laminar to turbulent flow in the film was used. The suggestion of Dukler **[8]** to integrate the equations and let the universal velocity distribution establish the flow regime was adopted. The details of the analysis are outlined in Appendix B. The results are

$$
Nu_{z} \equiv \frac{h_{z}D}{k_{\ell}} = \frac{\rho_{\ell}c_{\ell}D u_{\tau}}{k_{\ell}F_{2}}
$$
 (21a)

$$
\text{St}_{\mathbf{z}}^* \equiv \frac{\mathbf{h}_{\mathbf{z}}}{\rho_{\ell} \mathbf{c}_{\ell} \mathbf{u}_{\tau}} = \frac{1}{\mathbf{F}_{\mathbf{z}}} \tag{21b}
$$

where

$$
u_{\tau} = \sqrt{\frac{g_o \tau_o}{\rho_{\ell}}} \tag{22}
$$

and

$$
\text{for } 0 < \delta^+ < 5 \colon \mathbf{F}_2 = \delta^+ \text{ Pr} \tag{23a}
$$

for
$$
5 < \delta^+
$$
 < 30: F₂ = 5Pr + 5ln[1 + Pr $(\frac{\delta^+}{5} + 1)$] (23b)

for
$$
\delta^+
$$
 > 30: $F_2 = 5Pr + 5ln(1 + 5Pr) + \frac{2.5}{\sqrt{1 + \frac{10}{Pr} \frac{M}{\delta^+}}}$ x

$$
x \ln \left[\frac{2M-1 + \sqrt{1 + \frac{10}{Pr} \frac{M}{\delta^+}}}{2M-1 - \sqrt{1 + \frac{10}{Pr} \frac{M}{\delta^+}} + \frac{60}{\delta^+} M - 1 + \sqrt{1 + \frac{10}{Pr} \frac{M}{\delta^+}}} \right]
$$
(23c)

Here
$$
M = \frac{F_o \delta^+ \nu_e}{\tau_o u_\tau}
$$
 (24)

and

$$
\delta^+ \equiv \delta u_\tau / v_\ell \tag{25}
$$

Further Re_{ℓ} defined as

$$
\text{Re}_{\ell} = \frac{(1-x)\text{GD}}{\mu_{\ell}} = \frac{4\Gamma}{\mu_{\ell}} = \frac{4}{\mu_{\ell}} \int_{0}^{\delta} \rho_{\ell} v_{z} dy = 4 \int_{0}^{\delta^{+}} v_{z}^{+} dy^{+} \tag{26}
$$

is evaluated from the velocity distribution **Eq.** (B-1) with the following results:

$$
\delta^+ < 5 \qquad \qquad \text{Re}_\ell = 2(\delta^+)^2 \qquad (27a)
$$

$$
5 < \delta^+
$$
 < 30 $Re_{\ell} = 50 - 32.2 \delta^+ + 20 \delta^+ 1 n \delta^+$ (27b)

$$
\delta^+ > 30 \qquad \text{Re}_\ell = -256 + 12\delta^+ + 10\delta^+ \text{ln}\delta^+ \qquad (27c)
$$

A plot of Re_ℓ vs δ^+ is shown in Fig. 4.

For any assumed magnitude of Pr, δ^+ and M, calculate Re_{ℓ} from Eq. (27), F_2 from Eq. (23) and St^{*} from Eq. (21b). Then curves of St^{*} vs Re_c for various M can be constructed. Fig. 5 for Pr = 1 and **5** was drawn **by** this procedure.

The calculation procedure starts **by** dividing the tube length in increments of changes in quality x and for a given flow rate and fluid conditions calculate the increment of length required to accomplish this quality change. The calculation is a step-wise one requiring trial-and-error at each step. The procedure is outlined in a sample calculation in Appendix **A.**

Average Heat Transfer Coefficient

For the case of uniform wall temperature a mean heat transfer L coefficient **h_m** may be defined by $h_m = (1/L)$ | $h_z dz$. Then **.0**

$$
q = \Gamma_L \pi Dh_{fg} = h_m \Delta T \pi DL \qquad (28)
$$

For an element of length dz

$$
dq = \pi D h_{gh} d\Gamma = h_{z} \Delta T \pi D dz
$$
 (29)

Rearrange **Eq. (29)** and integrate

$$
\int_{0}^{\Gamma_{\mathbf{L}}} \frac{d\Gamma}{h_{\mathbf{z}}} = \frac{\Delta T}{h_{\mathbf{f}}g} \int_{0}^{\Gamma_{\mathbf{L}}} dz = \frac{\Delta T}{h_{\mathbf{f}}g} \mathbf{L} = \frac{\Gamma_{\mathbf{L}}}{h_{\mathbf{m}}} \tag{30}
$$

Then from **Eq. (30)** with **Eq. (26)**

$$
\frac{1}{h_m} = \frac{1}{\Gamma_e} \int_0^{\Gamma_e} \frac{d\Gamma}{h_z} = \frac{1}{Re_{\ell,e}} \int_0^{Re_{\ell,e}} \frac{dRe_{\ell}}{h_z}
$$
 (31)

or since $Re_{\ell} \sim (1-x)$ from Eq. (26), this becomes

$$
\frac{1}{h_m} = \frac{1}{x_e} \int_{x_e}^{1} \frac{dx}{h_z}
$$
 (32)

From the h_z calculated along the length, this length mean heat transfer coefficient may be calculated integrating with respect to quality as an alternative.

EXPERIMENT

The basic apparatus, schematically shown in Figure **6,** consists of a closed-loop refrigerant flow circuit driven **by** a mechanicalsealed rotor pump. Upstream of the test section, an electrically heated boiler produces vapor, which passes through a flow meter and a throttle valve to the test section. Downstream of the test section, an after-condenser was provided to ensure fully condensed refrigerant at the pump inlet. The pump flow was set for any test run and flow

rate in the test section was controlled **by** the by-pass loop. The pressure level was set **by** adjusting the heat input to the boiler and the throttle valve.

The test section itself is an annular shaped heat exchanger with refrigerant flowing through the inner tube and cooling water flowing in the outer annulus counter-currently. The 0.493 in. ID. smooth nickel tube test section was divided into six **3** ft-long sections. Each section has a separate cooling water circuit and the sections are connected smoothly with specially made stainless steel fittings in order not to disturb the condensate flow.

Each of the six sections except the third section from the inlet was separately and identically instrumented to give basic data on the condensing refrigerant. Two thermocouples are placed in the middle of the **3** ft section at the side; one at the outside of the condenser tube and the other one in the vapor at the center of the tube. Two differential thermocouples between the inlet and the outlet of the cooling water circuit are located in two different radial positions in order to detect any possible non-uniformity in temperature. On the third section, in addition to the above thermocouples, two more thermocouples are placed at the top and bottom of the tube wall to measure any circumferential variation of the wall temperature.

All the thermocouples were made of **0.005** in. **O.D.** nylon-sheathed copper and constantan wire.

Seven pressure taps were installed at every connection between the **3** ft sections for measurement of local pressure gradients.

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All the loop except the part from the pump to the boiler was insulated with fiberglass. The heat transfer between the test section and the atmosphere was not measurable in a blanked off run with no vapor flow.

Data were taken after steady state had been attained for one hour in the system. The heat flux to the coolant was obtained from the coolant flow rate and the temperature change. The condensing wall temperature was determined from the outside tube wall temperature and the heat flux. **All** the measurements were done on one **3** ft section at a time from up-stream to down-stream. The coolant flow was regulated such that the wall temperatures were kept almost constant through the test section and the temperature change of the coolant was in the range of **1** to **3*F.**

Heat balance was checked with total enthalpy change from the inlet of the test section to the outlet of the after-condenser. In most runs, except Run **1,** the heat balance error was less than **+ 6%.**

The data for both R-12 and R-22 are tabulated in the Appendix. Pressure drop data was taken only for R-22. Figures **7, 8, 9** are samples of the plot of the data but are representative of all of the data. Additional plots of the data are presented in references **[18]** and **[19].**

DISCUSSION OF **RESULTS**

Since the theoretical analysis was based on the annular flow model, the results are applicable only to the case where annular flow is developed. To date no successful investigation has been made of condensation flow regimes. For gas and oil mixtures,

a flow regime map was drawn **by** Baker [4], but it may not be applicable to two-phase flow with condensation. However, it surely gives an approximate view of the flow regime boundaries of condensation. Quandt **[13]** analyzed qualitatively the force field of gas-liquid flow. Still a quantitative figure of the flow regime boundaries cannot be obtained from an analysis. Therefore, until more reliable information about flow regimes of vapor-liquid flow with condensation is available, it is recommended that the Baker plot be used for determining probable flow regimes.

In most cases of practical forced-convection flow the regime appears to be annular except at the very low quality region. This analysis is not applicable to the very low quality region because the flow regime may be different and because the condensate film is so thick that the flat plate analysis is no longer valid for a tube. The present method is therefore not suggested for use when the vapor quality is less than 20%. **A** fared curve between the present correlation at $x = 0.20$ and McAdams equation for single phase flow (x **= 0)** will give useful information for the low quality region.

Entrainment of liquid in the vapor core was neglected in the analysis. Since thermal resistance is mainly offered **by** the laminar sublayer and the buffer layer, the entrainment effect is not significant when the condensate film thickness is larger than that of the high thermal resistance layers $(\delta^+ > 30)$. However, as expected, the effect appears to be significant at the very high vapor quality region where a very thin film exists $(\delta^+ < 5)$, as shown in some of the test runs. In a few runs at very high vapor flow rate when a considerable amount of entrainment was produced, the theory predicted

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lower values of h than those measured.

As the total flow rate decreases to low values, the thickness of the liquid film on the wall of a horizontal tube may be changed significantly. Even though the flow shape becomes an eccentric annulus, the analysis may give a good prediction because the heat transfer coefficient increases at the top and decreases at the bottom of the tube when this happens. However when truly stratefied flow exists another theory should be used.

The agreement with the present data is within **10%** except for a few low quality points. In general predictions are slightly lower than the experimental data within the range of measurement accuracy, Figs. **7, 8, 10,** and **11.** The pressure drop measurements, Figs. **9** and 12, also show good agreement except for Run **8.** It is interesting to note that at the upstream end of the condensing tube the predicted pressure gradient has a negative slope. However, the measurement shows the opposite trend. Except for Runs **5** and **8,** the pressure drop of the first section is always higher than that of the other sections.

Other Comparisons

Figure **13** shows the present data plotted on coordinates suggested **by** Akers and Rosson [2]. The solid lines represent their recommended correlation equations. Practically all of the data fall well above this recommendation. **A** plot of this same data **[18]** on coordinates suggested **by** Brauser **[5]** shows an equally large scatter. It is not surprising that such scatter should exist. In Fig. **13** the h for a given ΔT and pressure is essentially a function of $G_{\mathbf{v}}$ independent of quality. For the same G_v the liquid layer thickness, which

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offers the primary heat transfer resistance, is greatly different at qualities of, say, **10%** and **90%;** hence h should be quite different.

The present data along with the data of Altman et al **[3]** was compared **[18]** with a prediction equation suggested **by** Boyko and Kruzhilin [20] and was found to scatter badly. In general, the data fell as much as **250%** above and **100%** below the suggested prediction.

Figure **5** shows a comparison of the present predicted results with the predictions of Carpenter and Colburn **[6]** and Kunz and Yerazunis **[9].** The Carpenter-Colburn equation was derived considering only a laminar sub-layer and shows essentially no effect of the liquid Reynolds number. The coefficients were determined empirically for a limited range of data. The Kunz and Yerazunis study omitted the effect of **D,** gravitational effect and the momentum pressure gradient. Their result shows a discrepency from the present analysis at liquid Reynolds numbers above around **1000.**

CONCLUSION

The proposed prediction method for forced convection condensation heat transfer involves a combination of and modification of several previous analyses, [14][15][8], and agrees with the present and other data to within **+ 10%** for refrigerants in the range of conditions commonly found in commercial refrigeration equipment.

ACKNOWLEDGEMENTS

This work was supported **by** Technical Committee **1.3** of ASHRAE. Valuable suggestions and advice were received during the course of the work from Professors **A. E.** Bergles and **A. A.** Sonin and to Mr. Donald Traviss.

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Nomenclature

- total condensing length $\mathbf L$
- momentum $\mathbf m$
- $wall$ \mathbf{o}
- vapor $\mathbf v$
- $local$ \mathbf{z}

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APPENDIX A

Sample Calculation

Given Conditions

 $G = 250,000$ 1bm/ft^2 hr T_{sat} = 86°F $T_o = 76^\circ F$ physical properties (from Du Pont Table of F-22) Viscosity $\mu_0 = 0.557$ lbm/hr ft $\mu_{v} = 0.0322$ 1bm/hr ft Conductivity $K_0 = 0.0495$ Btu/hr ft^oF Specific heat $C_n = 0.305 \text{ Btu/lbm}$ ^oF Latent heat **h** $h_{fg} = 76.470 \text{ Btu/lbm}$ Density $\rho_g = 73.278$ lbm/ft³ $p_{y} = 3.1622$ 1bm/ft^3

Pr **=** 3.43

D = 0.493 in

Assuming that complete condensation occurs in the tube, the quality change is divided into 20 steps. **A** sample calculation will be done for the quality change from **72.5%** to **67.5%.** The local heat transfer coefficient at x **= 0.7** will be considered as the average value in this quality change.

From **Eq. (7)**

$$
\left(\frac{\mathrm{dP}}{\mathrm{d}z}\right)_{\mathrm{f}} = -16.96 \mathrm{lbf/ft}^2/\mathrm{ft}
$$

From Eq. (4) with $(S/A_z) = D/4 = 0.0103$ ft, $\tau_o = 0.174$ $1bf/ft^2$ From Eq. (22), $u_{\tau} = 992 \text{ ft/hr}$ Take for a first trial $D(dx/dz) = -0.001$.

From Eq. (13), $\left(\frac{dP}{dz}\right)_m = 1.36$ lbf/ft²/ft For a horizontal tube $\left|\frac{du}{dx}\right| = 0$ \int_{0}^{1} $\frac{1}{2}$ From Eq. (3), $\frac{dP}{dz}$ = -16.96 + 1.36 = -15.60 lbf/ft²/ft From Eq. (12) , $\alpha = 0.95$ From Eq. (26), $Re_{\ell} = \frac{(1 - 0.7)(250,000)(0.493)}{(0.557)(12)} = 5532$ From Eq. (27c), $5532 = -256 + 126^{+} + 106^{+}1n6^{+}$ By **trial and error calculate** δ^+ = 99.7 Then from Fig. 3 at $\delta^+ = 99.7$, $\beta = 1.25$ From Eq. (18), $F_0 = 19.20 \text{ lbf/ft}^2/\text{ft}$ From **Eq.** (24), M **=** 0.084 From Eq. $(23c)$, F₂ = 34.82 From Eq. (21b), $h_z = \frac{(73.278)(0.305)(992)}{(34.82)} = 637$ Btu/hr ft²F D^2 G h Since $\frac{A}{A} = h_2 \Delta T = \frac{h}{4} \frac{A T}{\pi D} \frac{B}{\Delta z}$ $\Delta x = \frac{4(637)(10)}{2} = 0$ \overline{G} **h**_{fg} \overline{G} (250,000) (76.470) \overline{G} 0.00133

Recalculate using this magnitude instead of **0.001.** The final results $h_z = 637$, convergence is very rapid. Then

$$
\Delta z = \frac{(\Delta x)(D)}{0.00133} = \frac{(0.05)(0.493)}{(0.00133)(12)} = 1.53 \text{ ft}
$$

the increment of length required to change the quality from **72.5%** to **67.5%.**

A similar calculation should be made for each Ax of **5%** to determine the corresponding $h_{\mathbf{z}}$ and $\Delta \mathbf{z}$. A plot of $h_{\mathbf{z}}$ and \mathbf{x} vs z may be constructed. Also P vs x or z may be plotted.

 $\bar{\mathbf{r}}$

APPENDIX B

Heat Transfer Analysis

The universal velocity was assumed in the liquid layer

$$
0 < \delta^{+} < 5
$$

\n
$$
v_{z}^{+} = y^{+}
$$

\n
$$
5 < \delta^{+} < 30
$$

\n
$$
v_{z}^{+} = -3.05 + 5 \ln y^{+}
$$

\n
$$
30 < y^{+}
$$

\n
$$
v_{z}^{+} = 5.5 + 2.5 \ln y^{+}
$$

\n
$$
(B-1)
$$

where

$$
v_z^+ = v_z/v_{6\sigma_0}^T/\rho = v_z/u_{\tau}; \delta^+ = \frac{\delta}{\nu}\sqrt{\frac{g_0^T \sigma_0}{\rho}}
$$

Rewrite **Eq.** (19a) as follows

$$
\tau = \frac{\rho_{\ell}}{g_0} \left(1 + \frac{\epsilon_{m}}{\nu_{\ell}} \right) u_{\tau}^{2} \frac{dv_{z}^{+}}{dy^{+}}
$$
 (B-2)

Solve this for ϵ_m with Eq. $(B-1)$

$$
0 < \delta^{+} < 5, \tau \tilde{z}_{0} \qquad \epsilon_{m} = 0
$$

\n
$$
5 < \delta^{+} < 30, \tau \tilde{z}_{0} \qquad \epsilon_{m} = \nu_{\ell} (\frac{y^{+}}{5} - 1)
$$

\n
$$
30 < \delta^{+}, \tau = F_{0} (\delta - y) + \tau_{v} \text{ and } \nu < \epsilon_{m}
$$

\n
$$
\epsilon_{m} = \frac{\nu_{\ell}}{2.5} \left[y^{+} - \frac{M}{\delta^{+}} (y^{+})^{2} \right]
$$
 (B-3)

where

$$
M \equiv \frac{F_o}{\tau_o} \frac{\delta^+ \nu}{\mu_{\tau}}
$$

Rewrite Eq. (19b) in the following form assuming $q/A \approx (q/A)_{\text{o}}$:

$$
\frac{1}{h_z} = \frac{T_\delta - T_o}{(q/A)_o} = \int_0^{\delta^+} \frac{v_\ell}{\rho_\ell C_\ell (q + \epsilon_h) u_\tau} dy^+
$$
 (B-4)

Taking $\epsilon_h = \epsilon_m$, substitute Eq. (B-3) into (B-4) and obtain Eq. (21) where F_2 is given by Eq. (23) in the three zones.

The results of this analysis can be put in an alternative form. **Eq.** (21) can be rewritten as follows:

$$
h_{z}^* = \frac{Pr}{F_2} \left(\frac{\delta^+}{M}\right)^{1/3}
$$
 (B-5)

$$
h_{z}^{\star} = \frac{h_{z}}{k} \left(\frac{v_{\ell} \mu_{\ell}}{g_{o} F_{o}} \right)^{1/3}
$$
 (B-6)

The results can be plotted as shown in Fig. 14 and involve

$$
\tau_{\mathbf{v}}^* \equiv \frac{\tau_{\mathbf{v}}}{F_o} \left(\frac{v_{\ell} u_{\ell}}{g_o F_o} \right)^{-1/3}
$$

$$
\delta^* \equiv \delta \left(\frac{v_{\ell} u_{\ell}}{g_o F_o} \right)^{-1/3}
$$
 (B-7)

where

$$
\delta^{+} = \delta^{*} (\delta^{*} + \tau_{v}^{*})^{1/2}
$$
 (B-8)

Then

$$
M = \frac{1}{1 + \tau_v^{\star}/\delta^{\star}}
$$
 (B-9)

The momentum equation for the vapor core, Fig. **1,** is

$$
- \frac{dP}{dz} A_{\mathbf{v}} - \tau_{\mathbf{v}} S_{\mathbf{v}} + \frac{a}{g_{\mathbf{0}}} \rho_{\mathbf{v}} A_{\mathbf{v}} = \frac{1}{g_{\mathbf{0}}} \frac{d}{dz} (U_{\mathbf{v}} W_{\mathbf{v}}) - U_{\mathbf{i}} \frac{dW_{\mathbf{v}}}{dz}
$$
 (B-10)

Again substituting **Eq.** (12)and (14) in **Eq.** (B-10)

$$
\tau_{\mathbf{v}} \frac{4}{\alpha D} = -\frac{dP}{dz} + \frac{a}{g_o} \rho_{\mathbf{v}} - \frac{G^2/\rho_{\mathbf{v}}}{g_o D} D \frac{dx}{dz} \left[2 \frac{x}{\alpha} \right]
$$

$$
+ \frac{1-2x}{\alpha} \left(\frac{\rho_{\mathbf{v}}}{\rho_{\mathbf{g}}} \right)^{2/3} - \frac{\beta(1-x)}{\alpha(1-\alpha)} \left(\frac{\rho_{\mathbf{v}}}{\rho_{\mathbf{g}}} \right) \right]
$$
(B-11)

For assumed magnitudes of δ^+ , Pr and τ_v^* , calculate Re_{ℓ} from Eq. (28), M from Eq. (B-9), δ^* from Eq. (B-8), F₂ from Eq. (26) and $h_{\mathbf{z}}^{\uparrow}$ from Eq. (B-5). With these calculations, Fig. 14 can be drawn and is an alternative presentation of results.

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APPENDIX **C**

Tables of Data

Heat Balance Error **+6.05%**

 $\mathcal{L}^{\text{max}}_{\text{max}}$

Tables of Data

Heat Balance error = 2.9%

 \bar{s}

Heat balance error **= 0.8%**

Heat balance error = 13.7%

 $\mathcal{L}(\mathcal{A})$ and $\mathcal{L}(\mathcal{A})$

Calculated

balance error = **1.5%** Heat

balance error **= 1.9%** Heat

32.3 36.5

68.7 62.3 12,200 **7,060** **378**

14.5

 $\frac{194}{-}$

Heat balance error = 2.3%

Heat balance error = 2.7%

Heat balance error = **1.06%**

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Heat balance error = 20.8%

Heat balance $error = 4$.

Figure Captions

 $\sim 10^{-10}$

Figure No.

 $FIG.2$

FIGURE 5^+ vs. β

 $F \cup 5$

DIAGRAM OF **APPARATUS FIGURE 6 SCHEMATIC**

FIG.7 LOCAL HEAT TRANSFER **COEFFICIENT** FOR R-12 COMPARED WITH **ANALYSIS**

FIGURE LOCAL HEAT TRANSFER 8 **COEFFICIENTS**

FIGURE TOTAL STATIC PRESSURE 9 **GRADIENTS**

FIGURE II DATA COMPARED WITH PRESENT ANALYSI

R-22 PRESSURE DROP **DATA** COMPARED WITH **ANALYSIS**

FIG. 12

 $F1g.13.$ Data on Akers-Rosson Plot

FIGURE 14 DIMENSIONLESS LOCAL HEAT TRANSFER COEFFICIENTS (Pr=5)