REFRIGERANT FORCED-CONVECTION CONDENSATION INSIDE HORIZONTAL TUBES

Soonhoon Bae John S. Maulbetsch Warren M. Rohsenow

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Engineering Projects Laboratory Department of Mechanical Engineering Massachusetts Institute of Technology Cambridge, Massachusetts 02139

December 1, 1970

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ABSTRACT

High vapor velocity condensation inside a tube was studied theoretically. The heat transfer coefficients were calculated by the momentum and heat transfer analogy. The Von Karman universal velocity distribution was applied to the condensate flow. Pressure drop was calculated by the Lockhart-Martinelli method and the Zivi void fraction equation.

Experimental data was obtained for the mass velocities from 150,000 to 555,000 lbm/ft² hr for R-12 and R-22 condensing in a 0.493" I.D. 18 ft. long test section. The measured heat transfer coefficients agreed with the prediction within 10% except a few points in the very low quality region.

INTRODUCTION

When condensation takes place inside a horizontal tube with high vapor velocity, condensate flows in an annular shape on the tube wall and vapor flows in the core.

Many investigators have studied this subject both experimentally and analytically. Empirical correlations involving non-dimensional groups were not quite successful because the correlations did not include all the flow variables [1], [2], [5], [7], [12]. Carpenter and Colburn [6] considered only the laminar sublayer of condensate flow and derived an equation with an empirical constant. This method was modified by later investigators [3], [15]. For a small range of the Prandtl Number, this equation gives good agreement with empirical data. But the equation has no general applicability. Rohsenow, Webber and Ling [14] analyzed the liquid film on the vertical plate and the heat transfer coefficient was obtained by the heat and momentum transfer analogy. A similar approach appeared in later papers [8], [9]. This method will be developed further for the annular flow regime in this paper.

THEORY

Flow Model

For condensation inside a horizontal tube with high vapor velocities, annular flow is the predominent flow pattern and slug flow may appear at very low vapor qualities. Annular flow with a uniform liquid layer thickness around the circumference of a tube is assumed to exist in the parameter ranges of interest. The condensate accumulation at the bottom of a horizontal tube has a negligible effect except at very low vapor flow rates. At very high vapor flow rates entrainment

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of liquid droplet in the vapor may occur but this will be neglected in the following analysis.

Consider a length element of a tube, shown in Fig. 1. For the entire cross-section the momentum equation is

$$-\left(\frac{\mathrm{dP}}{\mathrm{dz}}\right)A_{z} - \tau_{o}S + A_{z}\frac{a}{g_{o}}\left[\alpha\rho_{v} + (1-\alpha)\rho_{l}\right] = \frac{1}{g_{o}}\frac{\mathrm{d}}{\mathrm{dz}}\left(U_{v}W_{v} + U_{l}W_{l}\right) \quad (1)$$

where a is an acceleration due to the external body force. Rearranging Eq. (1) yields

$$-\left(\frac{\mathrm{d}P}{\mathrm{d}z}\right) = \tau_{o} \frac{\mathrm{S}}{\mathrm{A}_{z}} - \frac{\mathrm{a}}{\mathrm{g}_{o}} \left[\alpha \rho_{v} + (1 - \alpha)\rho_{\ell}\right] + \frac{1}{\mathrm{g}_{o}\mathrm{A}_{z}} \frac{\mathrm{d}}{\mathrm{d}z} \left(\mathrm{U}_{v}\mathrm{W}_{v} + \mathrm{U}_{\ell}\mathrm{W}_{\ell}\right) (2)$$

The above equation shows that the total static pressure gradient is the sum of pressure gradients due to friction, gravity and momentum change.

$$\left(\frac{\mathrm{dP}}{\mathrm{dz}}\right) = \left(\frac{\mathrm{dP}}{\mathrm{dz}}\right)_{\mathrm{f}} + \left(\frac{\mathrm{dP}}{\mathrm{dz}}\right)_{\mathrm{g}} + \left(\frac{\mathrm{dP}}{\mathrm{dz}}\right)_{\mathrm{m}}$$
(3)

Comparing Eq. (2) and (3),

$$\left(\frac{\mathrm{dP}}{\mathrm{dz}}\right)_{\mathrm{f}} = -\tau_{\mathrm{o}} \frac{\mathrm{S}}{\mathrm{A}_{\mathrm{z}}} \tag{4}$$

$$\left(\frac{dP}{dz}\right)_{g} = \frac{a}{g_{o}} \left[\alpha \rho_{v} + (1-\alpha)\rho_{l}\right]$$
(5)

$$\left(\frac{dP}{dz}\right)_{m} = -\frac{1}{g_{o}A_{z}} \frac{d}{dz} \left(U_{v}W_{v} + U_{\ell}W_{\ell}\right)$$
(6)

The friction pressure drop was obtained by an approximation of the Lockhart-Martinelli method [11] as follows:

$$\begin{split} \left(\frac{dP}{dz}\right)_{f} & \frac{g_{o}^{D}}{G^{2}/\rho_{v}} = -\tau_{o} \frac{S}{A_{z}} \frac{g_{o}^{D}}{G^{2}/\rho_{v}} = -0.09 \frac{GD}{\mu_{v}}^{-0.2} \left[x^{1.8} + 5.7 \left(\frac{\mu_{\ell}}{\mu_{v}}\right)^{0.0523} (1-x)^{0.470} x^{1.33} \left(\frac{\rho_{v}}{\rho_{\ell}}\right)^{0.261} \\ &+ 8.11 \left(\frac{\mu_{\ell}}{\mu_{v}}\right)^{0.105} (1-x)^{0.94} x^{0.86} \left(\frac{\rho_{v}}{\rho_{\ell}}\right)^{0.522} \right] \end{split}$$
(7)

This empirical approximation was developed by Soliman et al [15] who used this equation to calculate τ_v . Here it is used to calculate τ_o as Lockhart-Martinelli suggest.

The gravity term, Eq. (5), can be rewritten in the following form:

$$\left(\frac{dP}{dz}\right)_{g} \frac{g_{o}^{D}}{G^{2}/\rho_{v}} = \frac{1}{Fr^{2}} \left[\frac{\rho_{\ell}}{\rho_{v}} - B\alpha\right]$$
(8)

where

$$Fr^{2} = \frac{\left(G/\rho_{v}\right)^{2}}{aD}$$
(9)

is the Froude number based on the total flow and

$$B \equiv \frac{\rho_{\ell} - \rho_{\mathbf{v}}}{\rho_{\mathbf{v}}}$$
(10)

is the buoyancy modulus. In the gravity field

$$\mathbf{a} = \mathbf{g} \, \sin \, \theta \tag{11}$$

The Zivi equation for local void fraction [17] is recommended for use in Eq. (8), as in reference [15].

$$\alpha = \frac{1}{1 + \left(\frac{1-x}{x}\right) \left(\frac{\rho_{v}}{\rho_{\ell}}\right)^{2/3}}$$
(12)

The momentum term, Eq. (6), can be written with Eq. (12) as follows:

$$\left(\frac{\mathrm{d}P}{\mathrm{d}z}\right)_{\mathrm{m}} \frac{\mathrm{g}_{\mathrm{o}}^{\mathrm{D}}}{\mathrm{G}^{2}/\rho_{\mathrm{v}}} = -\mathrm{D}\left(\frac{\mathrm{d}x}{\mathrm{d}z}\right) \left[2\mathrm{x} + (1-2\mathrm{x})\left(\frac{\rho_{\mathrm{v}}}{\rho_{\mathrm{g}}}\right)^{1/3} + (1-2\mathrm{x})\left(\frac{\rho_{\mathrm{v}}}{\rho_{\mathrm{g}}}\right)^{2/3} - 2(1-\mathrm{x})\left(\frac{\rho_{\mathrm{v}}}{\rho_{\mathrm{g}}}\right)\right]$$

$$(13)$$

where in Eq. (6)

$$W_{v} = GA_{z}x = A_{z}\alpha U_{v}\rho_{v}$$

$$W_{l} = GA_{z}(1-x) = A_{z}(1-\alpha)U_{l}\rho_{l}$$
(14)

The momentum equation for the entire liquid layer element, Fig.2, is

$$- \left(\frac{dP}{dz}\right) A_{z_{\ell}} + \tau_{v}S_{v} - \tau_{o}S + \frac{a}{g_{o}} \rho_{\ell}A_{z_{\ell}}$$
$$= \frac{1}{g_{o}} \left[\frac{d(U_{\ell}W_{\ell})}{dz} - U_{i} \frac{dW_{\ell}}{dz}\right]$$
(15a)

or

$$\tau_{o} = F_{o} \frac{A_{z_{l}}}{S} + \tau_{v} \frac{S_{v}}{S}$$
(15b)

where

$$F_{o} \equiv -\frac{dP}{dz} + \frac{a}{g_{o}} \rho_{\ell} - \frac{1}{g_{o}^{A}z_{\ell}} \left[\frac{d(U_{\ell}W_{\ell})}{dz} - U_{i}\frac{dW_{\ell}}{dz}\right]$$
(16)

Since for most of the tube length the liquid film is thin, a simple flat plate analysis will suffice for the heat transfer coefficient derivation.

The
$$A_{z_{\ell}}/S \approx \delta$$
 and $S_{v}/S \approx 1$; so Eq. (15b) becomes
 $\tau_{o} = F_{o}\delta + \tau_{v}$
(17)

The quantity F_0 may be expressed in terms of x and α by substituting Eqs. (12) and (14) into Eq. (16). Further from the universal velocity, Eq. (B-1), distribution $(U_1/U_l) \equiv \beta$ may be obtained as a unique function of δ^+ as shown in Fig. 3. Then Eq. (17) becomes

$$F_{o} = -\left(\frac{dP}{dz}\right) + \frac{a}{g_{o}}\rho_{\ell} - \frac{G^{2}}{g_{o}\rho_{v}}\frac{dx}{dz}\left[\frac{1}{1-\alpha}\left(\frac{\rho_{v}}{\rho_{\ell}}\right)^{1/3} - \left(\frac{(1-x)(2-\beta)}{(1-\alpha)^{2}}\right)\left(\frac{\rho_{v}}{\rho_{\ell}}\right)\right]$$
(18)

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where $\beta \equiv U_i/U_l$ is given by Fig. 3.

To determine the heat transfer coefficient it is assumed that the Karman momentum-heat transfer analogy analysis is applicable in the liquid layer. Then

$$\tau = \frac{\rho_{\ell}}{g_{0}} \left(v_{\ell} + \varepsilon_{m} \right) \frac{dv_{z}}{dy}$$
(19a)

$$q/A = \rho_{\ell} c_{\ell} (\alpha_{\ell} + \varepsilon_{h}) \frac{dT}{dy}$$
(19b)

The universal velocity distribution is used to determine dv_z/dy , ε_m is assumed equal to ε_h , and dT/dy and T(y) is determined by combining the above two equations. The procedure is identical with that presented in detail by Rohsenow, Webber and Ling [14] with two exceptions. In determining this temperature distribution, the momentum equation for the element (δ -y) of Fig. 2 was approximated as

$$\tau = F_{o} (\delta - y) + \tau_{v}$$

and no criterion for transition from laminar to turbulent flow in the film was used. The suggestion of Dukler [8] to integrate the equations and let the universal velocity distribution establish the flow regime was adopted. The details of the analysis are outlined in Appendix B. The results are

$$Nu_{z} \equiv \frac{h_{z}^{D}}{k_{\ell}} = \frac{\rho_{\ell} c_{\ell}^{C} D u_{\tau}}{k_{\ell} F_{2}}$$
(21a)

or

$$St_{z}^{*} \equiv \frac{h_{z}}{\rho_{\ell} c_{\ell} u_{\tau}} = \frac{1}{F_{2}}$$
(21b)

where

$$u_{\tau} = \sqrt{\frac{g_{o}^{\tau} o}{\rho_{\ell}}}$$
(22)

and

for
$$0 < \delta^+ < 5$$
: $F_2 = \delta^+ Pr$ (23a)

for
$$5 < \delta^+ < 30$$
: $F_2 = 5Pr + 5ln[1 + Pr(\frac{\delta^+}{5} + 1)]$ (23b)

for
$$\delta^+ > 30$$
: $F_2 = 5Pr + 5ln(1 + 5Pr) + \frac{2.5}{\sqrt{1 + \frac{10}{Pr} \frac{M}{\delta^+}}} x$

$$x \ln \left[\frac{2M-1+\sqrt{1+\frac{10}{Pr}\frac{M}{\delta^{+}}}}{2M-1-\sqrt{1+\frac{10}{Pr}\frac{M}{\delta^{+}}}}, \frac{\frac{60}{\delta^{+}}M-1-\sqrt{1+\frac{10}{Pr}\frac{M}{\delta^{+}}}}{\frac{60}{\delta^{+}}M-1+\sqrt{1+\frac{10}{Pr}\frac{M}{\delta^{+}}}} \right]$$
(23c)

Here

$$M \equiv \frac{F_{o} \delta^{+} v_{\ell}}{\tau_{o} u_{\tau}}$$
(24)

and

$$\delta^{+} \equiv \delta u_{\tau} / v_{\ell}$$
 (25)

Further Re_{ℓ} defined as

$$\operatorname{Re}_{\mathfrak{L}} = \frac{(1-\mathbf{x})GD}{\mu_{\mathfrak{L}}} = \frac{4\Gamma}{\mu_{\mathfrak{L}}} = \frac{4}{\mu_{\mathfrak{L}}} \int_{0}^{\delta} \rho_{\mathfrak{L}} \mathbf{v}_{\mathbf{z}} d\mathbf{y} = 4 \int_{0}^{\delta^{\mathsf{T}}} \mathbf{v}_{\mathbf{z}}^{\mathsf{T}} d\mathbf{y}^{\mathsf{T}}$$
(26)

is evaluated from the velocity distribution Eq. (B-1) with the following results:

$$\delta^{+} < 5$$
 $\operatorname{Re}_{\ell} = 2(\delta^{+})^{2}$ (27a)

$$5 < \delta^{+} < 30$$
 Re_l = 50 - 32.2 $\delta^{+} + 20\delta^{+} \ln \delta^{+}$ (27b)

$$\delta^+ > 30$$
 $\operatorname{Re}_{\ell} = -256 + 12\delta^+ + 10\delta^+ \ln\delta^+$ (27c)

A plot of Re_{ℓ} vs δ^+ is shown in Fig. 4.

For any assumed magnitude of Pr, δ^+ and M, calculate Re_{ℓ} from Eq. (27), F_2 from Eq. (23) and St^{*} from Eq. (21b). Then curves of St^{*} vs Re_l for various M can be constructed. Fig. 5 for Pr = 1 and 5 was drawn by this procedure.

The calculation procedure starts by dividing the tube length in increments of changes in quality x and for a given flow rate and fluid conditions calculate the increment of length required to accomplish this quality change. The calculation is a step-wise one requiring trial-and-error at each step. The procedure is outlined in a sample calculation in Appendix A.

Average Heat Transfer Coefficient

For the case of uniform wall temperature a mean heat transfer coefficient h_m may be defined by $h_m = (1/L) \int_0^L h_z dz$. Then

$$q = \Gamma_L \pi Dh_{fg} = h_m \Delta T \pi DL$$
 (28)

For an element of length dz

$$dq = \pi D h_{gh} d\Gamma = h_z \Delta T \pi D dz$$
 (29)

Rearrange Eq. (29) and integrate

$$\int_{0}^{\Gamma} \frac{d\Gamma}{h_{z}} = \frac{\Delta T}{h_{fg}} \int_{0}^{L} dz = \frac{\Delta T}{h_{fg}} L = \frac{\Gamma}{h_{m}}$$
(30)

Then from Eq. (30) with Eq. (26)

$$\frac{1}{h_{m}} = \frac{1}{\Gamma_{e}} \int_{0}^{\Gamma_{e}} \frac{d\Gamma}{h_{z}} = \frac{1}{Re_{l,e}} \int_{0}^{Re_{l,e}} \frac{dRe_{l}}{h_{z}}$$
(31)

or since $\operatorname{Re}_{\ell} \sim (1-x)$ from Eq. (26), this becomes

$$\frac{1}{h_{m}} = \frac{1}{x_{e}} \int_{x_{e}}^{1} \frac{dx}{h_{z}}$$
(32)

From the h_z calculated along the length, this length mean heat transfer coefficient may be calculated integrating with respect to quality as an alternative.

EXPERIMENT

The basic apparatus, schematically shown in Figure 6, consists of a closed-loop refrigerant flow circuit driven by a mechanicalsealed rotor pump. Upstream of the test section, an electrically heated boiler produces vapor, which passes through a flow meter and a throttle valve to the test section. Downstream of the test section, an after-condenser was provided to ensure fully condensed refrigerant at the pump inlet. The pump flow was set for any test run and flow rate in the test section was controlled by the by-pass loop. The pressure level was set by adjusting the heat input to the boiler and the throttle value.

The test section itself is an annular shaped heat exchanger with refrigerant flowing through the inner tube and cooling water flowing in the outer annulus counter-currently. The 0.493 in. ID. smooth nickel tube test section was divided into six 3 ft-long sections. Each section has a separate cooling water circuit and the sections are connected smoothly with specially made stainless steel fittings in order not to disturb the condensate flow.

Each of the six sections except the third section from the inlet was separately and identically instrumented to give basic data on the condensing refrigerant. Two thermocouples are placed in the middle of the 3 ft section at the side; one at the outside of the condenser tube and the other one in the vapor at the center of the tube. Two differential thermocouples between the inlet and the outlet of the cooling water circuit are located in two different radial positions in order to detect any possible non-uniformity in temperature. On the third section, in addition to the above thermocouples, two more thermocouples are placed at the top and bottom of the tube wall to measure any circumferential variation of the wall temperature.

All the thermocouples were made of 0.005 in. 0.D. nylon-sheathed copper and constantan wire.

Seven pressure taps were installed at every connection between the 3 ft sections for measurement of local pressure gradients.

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All the loop except the part from the pump to the boiler was insulated with fiberglass. The heat transfer between the test section and the atmosphere was not measurable in a blanked off run with no vapor flow.

Data were taken after steady state had been attained for one hour in the system. The heat flux to the coolant was obtained from the coolant flow rate and the temperature change. The condensing wall temperature was determined from the outside tube wall temperature and the heat flux. All the measurements were done on one 3 ft section at a time from up-stream to down-stream. The coolant flow was regulated such that the wall temperatures were kept almost constant through the test section and the temperature change of the coolant was in the range of 1 to 3°F.

Heat balance was checked with total enthalpy change from the inlet of the test section to the outlet of the after-condenser. In most runs, except Run 1, the heat balance error was less than $\pm 6\%$.

The data for both R-12 and R-22 are tabulated in the Appendix. Pressure drop data was taken only for R-22. Figures 7, 8, 9 are samples of the plot of the data but are representative of all of the data. Additional plots of the data are presented in references [18] and [19].

DISCUSSION OF RESULTS

Since the theoretical analysis was based on the annular flow model, the results are applicable only to the case where annular flow is developed. To date no successful investigation has been made of condensation flow regimes. For gas and oil mixtures, a flow regime map was drawn by Baker [4], but it may not be applicable to two-phase flow with condensation. However, it surely gives an approximate view of the flow regime boundaries of condensation. Quandt [13] analyzed qualitatively the force field of gas-liquid flow. Still a quantitative figure of the flow regime boundaries cannot be obtained from an analysis. Therefore, until more reliable information about flow regimes of vapor-liquid flow with condensation is available, it is recommended that the Baker plot be used for determining probable flow regimes.

In most cases of practical forced-convection flow the regime appears to be annular except at the very low quality region. This analysis is not applicable to the very low quality region because the flow regime may be different and because the condensate film is so thick that the flat plate analysis is no longer valid for a tube. The present method is therefore not suggested for use when the vapor quality is less than 20%. A fared curve between the present correlation at x = 0.20 and McAdams equation for single phase flow (x = 0) will give useful information for the low quality region.

Entrainment of liquid in the vapor core was neglected in the analysis. Since thermal resistance is mainly offered by the laminar sublayer and the buffer layer, the entrainment effect is not significant when the condensate film thickness is larger than that of the high thermal resistance layers ($\delta^+ > 30$). However, as expected, the effect appears to be significant at the very high vapor quality region where a very thin film exists ($\delta^+ < 5$), as shown in some of the test runs. In a few runs at very high vapor flow rate when a considerable amount of entrainment was produced, the theory predicted

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lower values of h than those measured.

As the total flow rate decreases to low values, the thickness of the liquid film on the wall of a horizontal tube may be changed significantly. Even though the flow shape becomes an eccentric annulus, the analysis may give a good prediction because the heat transfer coefficient increases at the top and decreases at the bottom of the tube when this happens. However when truly stratefied flow exists another theory should be used.

The agreement with the present data is within 10% except for a few low quality points. In general predictions are slightly lower than the experimental data within the range of measurement accuracy, Figs. 7, 8, 10, and 11. The pressure drop measurements, Figs. 9 and 12, also show good agreement except for Run 8. It is interesting to note that at the upstream end of the condensing tube the predicted pressure gradient has a negative slope. However, the measurement shows the opposite trend. Except for Runs 5 and 8, the pressure drop of the first section is always higher than that of the other sections.

Other Comparisons

Figure 13 shows the present data plotted on coordinates suggested by Akers and Rosson [2]. The solid lines represent their recommended correlation equations. Practically all of the data fall well above this recommendation. A plot of this same data [18] on coordinates suggested by Brauser [5] shows an equally large scatter. It is not surprising that such scatter should exist. In Fig. 13 the h for a given ΔT and pressure is essentially a function of G_v independent of quality. For the same G_v the liquid layer thickness, which

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offers the primary heat transfer resistance, is greatly different at qualities of, say, 10% and 90%; hence h should be quite different.

The present data along with the data of Altman et al [3] was compared [18] with a prediction equation suggested by Boyko and Kruzhilin [20] and was found to scatter badly. In general, the data fell as much as 250% above and 100% below the suggested prediction.

Figure 5 shows a comparison of the present predicted results with the predictions of Carpenter and Colburn [6] and Kunz and Yerazunis [9]. The Carpenter-Colburn equation was derived considering only a laminar sub-layer and shows essentially no effect of the liquid Reynolds number. The coefficients were determined empirically for a limited range of data. The Kunz and Yerazunis study omitted the effect of D, gravitational effect and the momentum pressure gradient. Their result shows a discrepency from the present analysis at liquid Reynolds numbers above around 1000.

CONCLUSION

The proposed prediction method for forced convection condensation heat transfer involves a combination of and modification of several previous analyses, [14][15][8], and agrees with the present and other data to within \pm 10% for refrigerants in the range of conditions commonly found in commercial refrigeration equipment.

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Nomenclature

A _z	cross section area ft ²
a	actual gravitational acceleration in the axial direction, g sin θ ft/hr^2
В	buoyancy modulus, Eq. (10)
С _р	specific heat Btu/lbm °F
D	tube inner diameter ft
Fo	defined in Eq. (16) lbf/ft ² /ft
F ₂	defined in Eq.(23a, b, c)
Fr	Froude number, Eq. (9)
g	gravitational acceleration ft/hr ²
go	constant, 4.17 x 10^8 lbm ft/lbf hr ²
G	total mass velocity lbm/ft ²
h _z	local heat transfer coefficient Btu/hr ft ² °F
h m	mean heat transfer coefficient Btu/hr ft ² °F
K	conductivity of the liquid Btu/ft hr °F
М	defined in Eq. (24)
(dP/dz)	Pressure Gradient lbf/ft ² /ft
Pr	Prandtl number µ _l C _p /K
(q/A)	heat flux Btu/ft ² hr
Rel	local liquid Reynolds number $\frac{G(1-x)D}{\mu_{g}}$
S	perimeter ft
St [*]	Stanton Number $\frac{n_z}{\rho_{\ell}c_{p}u_{\tau}}$
Т	temperature °F
U	mean velocity ft/hr
u _t	friction velocity $\sqrt{\frac{\kappa_0'o}{\rho_l}}$ ft/hr

vz	local velocity in the axial direction ft/hr
Wl	liquid flow rate lbm/hr
W	vapor flow rate lbm/hr
x	quality
у	radial distance from the wall ft
Z	axial distance from the condensation starting point ft
α	void fraction
al	thermal diffusivity $K/\rho_{l}c_{p}$ ft ² hr
β	U _i /U _L
δ	thickness of the condensate film ft
ε	eddy diffusivity ft ² hr
θ	angle of inclination
μ	viscosity lbm/ft hr
ν	kinematic viscosity ft ² /hr
ρ	density lbm/ft ³
τ	shear stress lbf/ft ²
τ _v	vapor shear stress on the liquid film lbf/ft ²
Г	liquid flow rate per unit circumference lbm/ft hr
Subscript	
δ	liquid vapor interface
e	exit
f	friction
g	gravity
h	thermal
i	interface
l	liquid

- L total condensing length
- m momentum
- o wall
- v vapor
- z local

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APPENDIX A

Sample Calculation

Given Conditions

 $G = 250,000 \ 1 \text{bm/ft}^2 \text{hr}$ $T_{sat} = 86^{\circ}F$ $T_{0} = 76^{\circ}F$ physical properties (from Du Pont Table of F-22) $\mu_0 = 0.557 \, 1 \, \text{bm/hr}$ ft Viscosity $\mu_{ij} = 0.0322 \text{ lbm/hr ft}$ $K_0 = 0.0495 \text{ Btu/hr ft}^{\circ}\text{F}$ Conductivity $C_p = 0.305 \text{ Btu/lbm }^\circ \text{F}$ Specific heat $h_{fg} = 76.470 \text{ Btu/lbm}$ Latent heat $\rho_0 = 73.278 \ \text{lbm/ft}^3$ Density $\rho_{\rm M} = 3.1622 \ \rm 1 bm/ft^3$

Pr = 3.43

D = 0.493 in

Assuming that complete condensation occurs in the tube, the quality change is divided into 20 steps. A sample calculation will be done for the quality change from 72.5% to 67.5%. The local heat transfer coefficient at x = 0.7 will be considered as the average value in this quality change.

From Eq. (7)

$$\left(\frac{dP}{dz}\right)_{f}$$
 = -16.96 lbf/ft²/ft

From Eq. (4) with $(S/A_z) = D/4 = 0.0103 \text{ ft}, \tau_o = 0.174 \text{ lbf/ft}^2$ From Eq. (22), $u_\tau = 992 \text{ ft/hr}$ Take for a first trial D(dx/dz) = -0.001. From Eq. (13), $\left(\frac{dP}{dz}\right)_{m} = 1.36 \text{ lbf/ft}^2/\text{ft}$ For a horizontal tube $\left(\frac{dP}{dz}\right)_{z} = 0$ From Eq. (3), $\frac{dP}{dz} = -16.96 + 1.36 = -15.60 \text{ lbf/ft}^2/\text{ft}$ From Eq. (12), $\alpha = 0.95$ From Eq. (26), $\operatorname{Re}_{\ell} = \frac{(1 - 0.7)(250,000)(0.493)}{(0.557)(12)} = 5532$ From Eq. (27c), $5532 = -256 + 12\delta^{+} + 10\delta^{+} \ln \delta^{+}$ By trial and error calculate $\delta^+ = 99.7$ Then from Fig. 3 at δ^+ = 99.7, β = 1.25 From Eq. (18), $F_0 = 19.20 \, \text{lbf/ft}^2 / \text{ft}$ From Eq. (24), M = 0.084From Eq. (23c), $F_2 = 34.82$ From Eq. (21b), $h_z = \frac{(73.278)(0.305)(992)}{(34.82)} = 637 \text{ Btu/hr ft}^2 \text{F}$ Since $\frac{q}{A} = h_z \Delta T = \frac{\pi}{4} \frac{D^2 G h_{fg}}{\pi D} \frac{\Delta x}{\Lambda z}$ $D_{\Delta z}^{\Delta x} = \frac{4 h_z^{\Delta T}}{G h_{f_z}} = \frac{4(637)(10)}{(250,000)(76.470)} = 0.00133$

Recalculate using this magnitude instead of 0.001. The final results $h_z = 637$, convergence is very rapid. Then

$$\Delta z = \frac{(\Delta x)(D)}{0.00133} = \frac{(0.05)(0.493)}{(0.00133)(12)} = 1.53 \text{ ft}$$

the increment of length required to change the quality from 72.5% to 67.5%.

A similar calculation should be made for each Δx of 5% to determine the corresponding h_z and Δz . A plot of h_z and x vs z may be constructed. Also P vs x or z may be plotted.

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APPENDIX B

Heat Transfer Analysis

The universal velocity was assumed in the liquid layer

$$0 < \delta^{+} < 5 \qquad v_{z}^{+} = y^{+}$$

$$5 < \delta^{+} < 30 \qquad v_{z}^{+} = -3.05 + 5 \ln y^{+} \qquad (B-1)$$

$$30 < y^{+} \qquad v_{z}^{+} = 5.5 + 2.5 \ln y^{+}$$

where

$$v_z^+ = v_z^{\prime} \sqrt{g_o^{\tau} o^{\prime \rho}} = v_z^{\prime} u_{\tau}; \delta^+ = \frac{\delta}{v} \sqrt{\frac{g_o^{\tau} o}{\rho}}$$

Rewrite Eq. (19a) as follows

$$\tau = \frac{\rho_{\ell}}{g_{0}} \left(1 + \frac{\varepsilon_{m}}{\nu_{\ell}}\right) u_{\tau}^{2} \frac{dv_{z}^{+}}{dy^{+}}$$
(B-2)

Solve this for $\boldsymbol{\epsilon}_m$ with Eq. (B-1)

$$0 < \delta^{+} < 5, \tau \tilde{\tau}_{0} \qquad \varepsilon_{m} = 0$$

$$5 < \delta^{+} < 30, \tau \tilde{\tau}_{0} \qquad \varepsilon_{m} = v_{\ell} (\frac{y^{+}}{5} - 1)$$

$$30 < \delta^{+}, \tau = F_{0} (\delta - y) + \tau_{v} \text{ and } v << \varepsilon_{m}$$

$$\varepsilon_{m} = \frac{v_{\ell}}{2.5} \left[y^{+} - \frac{M}{\delta^{+}} (y^{+})^{2} \right]$$
(B-3)

where

$$M \equiv \frac{F_o}{\tau_o} \frac{\delta^+ v}{u_{\tau}}$$

Rewrite Eq. (19b) in the following form assuming q/A \sim (q/A)_o:

$$\frac{1}{h_z} = \frac{T_{\delta} - T_o}{(q/A)_o} = \int_0^{\delta^+} \frac{v_{\ell}}{\rho_{\ell} C_{\ell} (q + \varepsilon_h) u_{\tau}} dy^+$$
(B-4)

Taking $\varepsilon_h = \varepsilon_m$, substitute Eq. (B-3) into (B-4) and obtain Eq. (21) where F_2 is given by Eq. (23) in the three zones.

The results of this analysis can be put in an alternative form. Eq. (21) can be rewritten as follows:

$$h_{z}^{*} = \frac{Pr}{F_{2}} \left(\frac{\delta^{+}}{M}\right)^{1/3}$$
(B-5)

$$h_{z}^{*} = \frac{h_{z}}{k} \left(\frac{\nu_{\ell} \mu_{\ell}}{g_{o} F_{o}} \right)^{1/3}$$
(B-6)

The results can be plotted as shown in Fig. 14 and involve

$$\tau_{\mathbf{v}}^{*} \equiv \frac{\tau_{\mathbf{v}}}{F_{o}} \left(\frac{\nu_{\ell} \mu_{\ell}}{g_{o} F_{o}} \right)^{-1/3}$$

$$\delta^{*} \equiv \delta \left(\frac{\nu_{\ell} \mu_{\ell}}{g_{o} F_{o}} \right)^{-1/3}$$
(B-7)

where

$$\delta^{+} = \delta^{*} (\delta^{*} + \tau_{v}^{*})^{1/2}$$
 (B-8)

Then

$$M = \frac{1}{1 + \tau_v^* / \delta^*}$$
(B-9)

The momentum equation for the vapor core, Fig. 1, is

$$- \frac{dP}{dz} A_{v} - \tau_{v}S_{v} + \frac{a}{g_{o}} \rho_{v}A_{v} = \frac{1}{g_{o}} \frac{d}{dz} (U_{v}W_{v}) - U_{i} \frac{dW_{v}}{dz}$$
(B-10)

Again substituting Eq. (12) and (14) in Eq. (B-10)

$$\tau_{\mathbf{v}} \frac{4}{\alpha D} = -\frac{dP}{dz} + \frac{a}{g_{o}} \rho_{\mathbf{v}} - \frac{G^{2}/\rho_{\mathbf{v}}}{g_{o}D} D \frac{dx}{dz} \left[2 \frac{x}{\alpha} + \frac{1-2x}{\alpha} \left(\frac{\rho_{\mathbf{v}}}{\rho_{\boldsymbol{\ell}}} \right)^{2/3} - \frac{\beta(1-x)}{\alpha(1-\alpha)} \left(\frac{\rho_{\mathbf{v}}}{\rho_{\boldsymbol{\ell}}} \right) \right]$$
(B-11)

For assumed magnitudes of δ^+ , Pr and τ_v^* , calculate Re_{ℓ} from Eq. (28), M from Eq. (B-9), δ^* from Eq. (B-8), F₂ from Eq. (26) and h_z^* from Eq. (B-5). With these calculations, Fig. 14 can be drawn and is an alternative presentation of results.

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APPENDIX C

Tables of Data

Run 1	G = 303,000	lbm/ft ² hr	T = \$6°1 sat	F R-22			
		M	leasured				
Sec No	T vapor	T _o out	Wwater	∆t _w	dP/dz		
1 2 3 4 5 6	83.17 82.97 82.40 82.31 82.04 81.88	71.25 71.45 69.47 70. 8 9 70.06 67.92	2400 2134 2718 1807 1953 2254	1.482 1.450 1.250 1.531 1.473 1.186	15.3 24.7 21.2 14.1 15.3 9.4		
	Calculated						
Sec No	Q/A	T _o in	Δτ	h	×m		
1 2 3 4 5 6	9200 7990 8450 7150 7450 6820	72.73 72.74 70.83 71.05 71.26 69.03	10.44 10.23 11.57 10.26 10.78 12.85	880 780 730 696 692 530	93.8 82.5 72.2 62.4 53.3 44.3		
Heat Ba	lance Error -	-8.5%					
Run 2	G = 485,000	1bm/ft ² hr	T = 8 1° sat Measured	°F R-22			
Sec No	T vapor	T _o out	Wwater	∆T _w	dP/dz		
1 2 3 4 5 6	81.64 81.19 81.01 80.19 79.62 79.44	72.81 70.79 70.12 68.84 68.13 67.89	1690 2800 1920 1510 1730 1910	2.61 1.58 2.29 2.85 2.50 2.20	73.2 62.5 57.8 47.2 53.0 43.6		
		Ca	alculated				
Sec No	Q/A	T _o in	Δτ	h	x m		
1 2 3 4 5 6	9,440 11,400 11,350 11,100 11,200 10,850	74.33 72.64 71.95 70.63 69.93 69.64	7.31 7.55 9.06 9.56 9.69 9.80	1290 1330 1250 1160 1150 1100	96.4 88.3 79.4 70.6 61.8 53.3		
Heat Bal	lance Error	+1.6%					

Run 3	G = 250,000	lbm/ft ² hr	$T_{sat} = 86$	°F R-22	
			Measured		
Sec No	T vapor	out	Wwater	$\Delta \mathbf{T}_{\mathbf{w}}$	dP/d z
1 2 3 4 5 6	86.57 86.39 85.96 85.77 85.69 85.58	70.37 71.02 70.12 69.38 69.67 71.04	3310 2560 2490 2910 2790 1750	1.27 1.47 1.37 1.07 1.02 1.36	24.2 20.5 17.0 13.0 13.0 7.1
		C	alculated		
Sec No	Q/A	T _o in	ΔΤ	h	x m
1 2 3 4 5 6	10,850 9,730 8,810 8,050 7,350 6,150	72.12 73.09 71.54 70.68 70.85 72.03	14.45 13.30 14.42 15.09 14.84 13.55	750 730 610 535 495 455	91.1 74.3 59.6 46.7 34.1 23.0
Heat B	alance Error	+4.65%			
Run 4	G = 470,000	1bm/ft ² hr	T = 85 sat = 85	•F R-22	
			Measured		
Sec No	T vapor	To out	W water	∆T _w	dP/d z
1 2 3 4 5 6	85.53 85.01 85.00 84.54 84.19 83.64	74.23 75.10 71.96 73.92 72.22 71.28	2660 1670 3500 1970 2470 3090	1.77 2.27 1.42 1.88 1.54 1.24	62.3 60.9 57.0 47.4 47.2 38.5
		C	Calculated		
Sec No	Q/A	^T o _{in}	$\Delta \mathbf{T}$	h	× m
1 2 3 4 5 6	12,200 9,700 12,800 9,610 9,840 9,900	76.19 76.68 74.02 75.47 73.81 72.87	9.34 8.33 10.98 9.07 10.38 10.77	1,300 1,160 1,165 1,060 950 920	95.1 85.6 76.0 66.9 59.0 51.0

Heat Balance Error +6.05%

Run 5	G = 270,000	lbm/ft ² hr	T = 92° sat	°F R-22	
		M	leasured		
Sec No	T vapor	out	W water	∆t _w	dP/dz
1 2 3 4 5 6	92.22 92.07 91.81 91.77 90.90 90.94	74.98 75.30 71.90 71.71 73.73 72.03	2300 20 8 0 2770 2690 2060 2340	1.87 2.02 1.69 1.54 1.67 1.35	18.9 24.8 20.1 14.2 14.2 9.0
		Ca	alculated		
Sec No	Q/A	T _o in	Δτ	h	x m
1 2 3 4 5 6	11,100 10,850 12,100 10,700 8,900 8,160	76.77 77.05 73.85 73.44 75.16 73.35	15.45 15.02 17.94 18.33 15.74 17.59	718 721 674 584 566 465	92.0 76.2 59.7 43.3 29.2 17.0
Heat Ba	lance Error	+4.45%			
Run 6	G = 240,000	lbm/ft ² hr	$T_{sat} = 92^{\circ}$	F R-22	
		1	Measured		
S ec No	T vapor -	out	W water	∆t _w	dP/dz
1 2 3 4 5 6	93.04 92.46 91.98 91.88 91.66 91.80	80.34 80.97 79.72 79.28 78.29 77.91	2440 2180 2300 2310 2600 2580	1.16 1.26 1.19 1.15 1.04 1.01	20.1 21.2 16.5 11.8 11.8 7.1
-		Ca	alculated		
Sec No	Q/A	T _o in	Δ T	h	×m
1 2 3 4 5 6	7550 7100 7090 6870 7000 6740	81.52 82.10 80.86 80.39 79.42 79.00	11.52 10.36 11.12 11.49 12.24 12.80	655 685 636 599 571 526	93.9 82.1 70.6 59.4 48.2 37.1
11 b D-	1	1 159			

Run 7	G = 308,000	lbm/ft ² hr	T = 92 sat	°F R-22	
		1	Measured		
Sec No	T vapor	T _o out	Wwater	$\Delta \mathbf{T}_{\mathbf{w}}$	dP/dz
1 2	98.07 97.80	80.96 82.12	2510 1860	1.83 2.13	25.4 23.6
3 4 5	97.37 97.02 96.98	80.06 79.26 78.58	2410 2800 2840	1.72 1.69 1.39	21.2 16.5 16.5
6	96.93	77.26	2080	1.58	11.8
		Ca	alculated		
Sec No	Q/A	To in	ΔT	h	× m
1 2 3 4	11,850 10,200 10,700 12,200	82.87 83.77 81.79 81.23	15.20 14.03 15.58 15.79	780 726 687 770	92.4 78.1 64.6 49.8
5 6	10,200 8,500	80.23 78.63	16.75 18.80	610 452	35.3 23.2
Heat Bal	lance Error	+0.9%			
		_			
Run 8	G = 316,000	1bm/ft ² hr	$T_{sat} = 100$	3°F R-22	
Run 8	G = 316,000	lbm/ft ² hr	T _{sat} = 103 Measured	3°F R-22	
Run 8 Sec No	$G = 316,000$ T vapor	lbm/ft ² hr M	T _{sat} = 103 Measured ^W water	δ°F R-22 ΔT _w	dP/dz
Run 8 Sec No 1 2 3 4 5 6	G = 316,000 Tvapor 103.12 102.91 101.79 101.33 100.98 100.98	lbm/ft ² hr M Toout 84.63 86.73 85.33 83.72 80.99 79.66	T _{sat} = 103 feasured ^W water 2530 1390 2480 3000 3070 3330	ΔT w 2.10 2.89 1.58 1.50 1.40 1.31	dP/dz 15.3 28.3 27.1 20.0 20.0 15.3
Run 8 Sec No 1 2 3 4 5 6	G = 316,000 Tvapor 103.12 102.91 101.79 101.33 100.98 100.98	lbm/ft ² hr M Toout 84.63 86.73 85.33 83.72 80.99 79.66 Ca	T _{sat} = 103 feasured ^W water 2530 1390 2480 3000 3070 3330 alculated	ΔT w 2.10 2.89 1.58 1.50 1.40 1.31	dP/dz 15.3 28.3 27.1 20.0 20.0 15.3
Run 8 Sec No 1 2 3 4 5 6 Sec No	G = 316,000 Tvapor 103.12 102.91 101.79 101.33 100.98 100.98	lbm/ft ² hr N Toout 84.63 86.73 85.33 83.72 80.99 79.66 Ca Tooin	$T_{sat} = 103$ Measured Wwater 2530 1390 2480 3000 3070 3330 Alculated ΔT	ΔT _w 2.10 2.89 1.58 1.50 1.40 1.31	dP/dz 15.3 28.3 27.1 20.0 20.0 15.3
Run 8 Sec No 1 2 3 4 5 6 Sec No 1 2 3 4 5 6	G = 316,000 Tvapor 103.12 102.91 101.79 101.33 100.98 100.98 Q/A 13,700 10,400 10,100 11,600 11,100 11,300	lbm/ft ² hr N Toout 84.63 86.73 85.33 83.72 80.99 79.66 Ca Toin 86.85 88.41 86.96 87.59 82.78 81.48	T _{sat} = 103 feasured Wwater 2530 1390 2480 3000 3070 3330 alculated ΔT 16.27 14.50 14.83 15.74 18.20 19.50	ΔT w 2.10 2.89 1.58 1.50 1.40 1.31 h 842 717 680 737 610 610	dP/dz 15.3 28.3 27.1 20.0 20.0 15.3 xm 91.2 75.8 62.6 48.7 34.1 19.8

Tables of Data

Run 1	G = 316,000	lb/hrft ²	^T water in	= 64.6 R-	12
			Measured		
Sec No	T vapor	T out	W water	∆t _w	
1	94	78.8	1510	2.72	
2	93.2	77.5	1450	2.49	
3	92.4	76.6	1460	2.32	
4	91.9	74.1	1930	1.69	
5	90.2	72.7	1835	1.46	
6	89.0	71.3	2170	1.17	
		(Calculated		
Sec	Т		- I.		х

Sec No	^T o _{in}	$\Delta \mathbf{T}$	Q/A	h	x m
1	80.5	13.5	10,600	785	91.5
2	79.0	14.2	9,350	658	75.2
3	78.0	14.4	8,780	610	60.4
4	75.5	16.4	8,440	514	46.4
5	73.8	16.4	6,930	423	33.9
6	72.4	16.6	6,560	396	22.1

Heat Balance error = 2.9%

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Run 2	G = 354,000	lb/hrft ²	T water in	= 64.4	R-12
]	Measured		
Sec No	T vapor	T Gout	Wwater	ΔT_{w}	
1 2 3 4 5 6	96.0 95.6 95.1 94.6 94.0 93.6	81.9 80.9 78.4 76.0 73.6 73.6 Ca	2070 1120 1580 1600 1590 1375 alculated	3.09 3.42 2.62 2.22 1.72 1.80	
Sec No	^T o _{in}	ΔT	Q/A	h	×m
1 2 3 4 5 6	84.6 82.5 80.1 77.5 74.7 74.6	11.4 13.1 15.0 17.1 19.3 19.0	16,500 9,900 10,700 9,200 7,060 6,400	1,450 755 713 538 366 337	88.0 68.6 53.6 39.1 27.2 17.3

Heat balance error = 0.8%

Run 3	G = 468,000	lb/hrft ²	T water i	.n = 64.4	R-12
			Measured		
Sec No	T vapor	T out	W water	∆t _w	
1	97.0	84.4	2,070	3.47	
2	95.8	84.0	1,120	4.06	
3	95.0	79.7	1,580	2.89	
4	93.9	77.0	1,600	2.42	
5	93.0	75.9	1,590	2.17	
6	91.0	75.4	1,475	2.02	
		C	Calculated		
Sec	T				
No	in	$\Delta \mathbf{T}$	Q/A	h	m
1	87.4	9.6	18,600	1,940	81.6
2	85.9	9.9	11,800	1,190	73.0
3	81.6	13.4	11,800	880	59.9
4	78.6	15.3	9,960	652	46.3
5	77.4	15.6	8,910	572	37.5
6	76.4	14.3	7,700	548	27.8
H eat ba	alance error	= 2.7%			
Run 4	G = 360,000	lb/hrft ²	T _{water} i	n = 67.9	R-12
			Massured		
C • •		T	measureu		
No	^T vapor	out	W water	$\Delta \mathbf{T}_{\mathbf{w}}$	
1	95.0	82.3	2,090	2.54	
2	94.7	81.0	1,490	2.58	
3	94.3	79.6	1,870	2.14	
4	93.7	79.3	1,500	2.22	
5	93.3	77.6	1,900	1.76	
6	93.0	76.5	2,770	1.38	
		C	alculated		
Sec	T _o .				x
No	īn	$\Delta \mathbf{T}$	Q/A	h	m
1	84.5	10.5	13,700	1,300	90.0
2	82.6	12.1	9,930	820	73.1
3	81.3	13.0	10,300	793	58.5
4	80.7	13.0	8,600	662	45.0
5	79.0	13.3	8,630	648	32.6
6	78.1	14.9	9,900	664	13.2

Heat balance error = 13.7%

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Run 5	G = 254,000	lb/hrft ²	^T water i	in = 67.9	R -12
			Measured		
Sec No	T vapor	out ^T	W water	$\Delta \mathtt{T}_{\mathbf{w}}$	
1 2 3 4 5 6	93.0 92.1 91.7 91.4 91.0	79.4 78.6 76.6 75.4 75.0	2,090 1,490 1,870 1,500 1,900 - Calculated	2.03 2.10 1.59 1.48 1.30	
Sec No	To in	Δτ	Q/A	h	×m
1 2 3 4 5 6	81.2 79.9 77.9 76.3 76.1	11.8 13.2 13.8 15.1 14.9	11,000 8,100 7,700 5,740 6,400 -	932 613 558 380 428 -	89.0 69.8 53.7 40.2 27.9 -

Run 6	G = 265,000	lb/hrft ²	T _{water} in = 67.9	R-12
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Measure	d
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Sec No	T vapor	Toout	Wwater	$\Delta \mathbf{T}_{\mathbf{w}}$
1	99.0	81.8	2,090	2.44
2	97.8	80.1	1,490	2.38
3	97.3	77.8	1,870	1.80
4	96.9	76.8	1,500	1.74
5	96.5	75.5	1,900	1.39
6	96	74.0	2,770	1.00

Calculated	l
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Sec No	To in	$\Delta \mathbf{T}$	Q/A	h	×m
1	84.0	15.0	13,200	879	87.0
2	81.6	16.2	9,150	565	65.2
3	79.2	18.1	8,700	478	47.4
4	77.9	19.0	6,750	355	32.2
5	77.6	19.9	6,830	344	19.0
6	75.2	20.8	7,160	344	3.5

Heat balance error = 1.5%

Run 7	G = 155,000	$1b/hrft^2$	T water in ⁼	= 51.4°F	R-12
			Measured		
Sec No	T vapor	^T out	Wwater	$\Delta \mathbf{T}_{\mathbf{w}}$	
1 2 3 4 5 6	84.6 84.2 83.8 83.5 82.0 79.4	67.2 62.7 62.4 60.5 58.2 55.2	1,760 1,500 1,610 1,900 2,260 2,360 alculated	2.32 1.92 1.78 1.40 0.84 0.56	
Sec No	To in	ΔT	Q/A	h	×m
1 2 3 4 5 6	68.9 63.7 63.6 61.6 59.2 55.8	15.7 20.3 20.2 21.9 22.8 23.6	10,600 7,450 7,400 6,870 5,840 3,440	675 367 366 314 256 143	82.8 53.8 29.7 8.8 -
Heat bal	lance error =	= 4.9%			
Run 8	G = 445,000	lb/hrft ²	T water in '	= 51.6	R-12
			Measured		
Sec No	T vapor	^T oout	Wwater	ΔT_w	
1 2 3 4 5 6	87.8 86.7 85.7 85.1 84.5 84.0	72.6 68.9 66.7 66.2 63.4 60.4	2,670 2,470 2,300 1,730 2,050 2,260	2.82 2.30 2.18 2.31 1.78 1.28	
		C	Calculated		
Sec No	To in	$\Delta \mathbf{T}$	Q/A	h	× m
1 2 3 4 5 6	74.9 71.3 68.8 67.9 65.0 61.6	12.9 15.4 16.9 17.2 19.5 22.4	20,150 14,700 13,000 10,350 9,430 7,500	1,560 955 768 602 484 335	88.6 68.7 53.0 39.8 28.8 19.3

Heat balance error = 1.9%

Run 9	G = 440,000	lb/hrft ²	T water in ⁼	= 51.4	R-12
]	Measured		
Sec No	T vapor	T _o out	Wwater	$\Delta \mathbf{T}_{\mathbf{w}}$	
1 2 3 4 5 6	88.3 87.7 86.6 86.1 85.8 85.1	72.0 69.1 68.2 67.8 64.7 60.9	2,580 2,350 1,900 1,470 1,460 1,190	2.88 2.65 2.76 2.60 2.21 1.74	
		С	alculated		
Sec No	^T o in	$\Delta \mathbf{T}$	Q/A	h	× m
1 2 3 4 5 6	75.1 71.6 70.6 69.3 66.1 61.8	13.2 16.1 16.0 16.8 19.7 23.3	19,200 15,500 13,500 9,250 8,360 5,350	1,450 962 856 551 425 230	89.2 69.2 52.6 39.5 29.4 21.6
Heat ba	lance error :	= 1.9%			
Run 10	G = 220,000	0 lb/hrft ²	^T water in	= 52.2	R-12
			Measured		
Sec No	T vapor	out	Wwater	Δt _w	
1 2 3 4 5 6	102.0 101.7 101.0 98.8 -	73.6 69.2 66.7 61.6 -	1,970 2,230 2,230 1,860 - - -	3.01 2.50 2.12 1.47 -	
Sec	Т				v
No	oin	$\Delta \mathbf{T}$	Q/A	h	^ m
1 2 3 4 5	76.1 71.6 68.7 62.3	25.9 30.1 32.3 36.5	15,300 14,400 12,200 7,060	590 478 378 194	81.6 42.8 14.5 - -
6	-	-	-	-	-

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Run 11	G = 272,	,000 1b/hr:	ft ² T _{water i}	n = 75.0	R-12
			Measured		
Sec No	^T vapor	Toout	Wwater	$\Delta \mathbf{T}_{\mathbf{w}}$	
1 2 3 4 5 6	94.5 94.2 93.9 93.6 93.0 91.0	84.7 84.1 83.2 80.4 79.3 78.3	1,750 1,420 1,510 1,270 1,915 1,565 Calculated	1.94 1.82 1.66 1.80 0.81 0.68	
Sec No	To in	$\Delta \mathbf{T}$	Q/A	h	× m
1 2 3 4 5 6 Heat ba	86.1 85.2 84.3 81.3 80.0 78.8	$8.4 \\ 9.0 \\ 9.6 \\ 12.3 \\ 13.0 \\ 12.2 \\ or = 0.4\%$	8,780 6.680 6,490 5,900 4,030 2,740	1,045 743 676 482 310 275	91.8 75.1 64.6 54.6 47.1 40.7
Run 12	G = 477,	000 1b/hr	ft ² T _{water i}	n = 64.0	R-12
			Measured		
Sec No	T vapor	Toout	Wwater	$\Delta \mathbf{T}_{\mathbf{w}}$	
1 2 3 4 5 6	93.0 92.2 91.0 90.2 87.0 85.0	82.0 80.5 77.3 72.9 69.4 68.7	1,820 1,030 1,915 2,280 2,070 2,280 Calculated	3.26 3.50 2.36 1.50 0.91 0.80	
Sec No	^T o _{in}	$\Delta \mathbf{T}$	Q/A	h	x m
1 2 3 4 5 6	84.5 82.0 79.2 74.4 70.3	8.5 10.2 11.8 15.8 16.7	15,300 9,350 11,600 8,840 5,350 4,720	1,800 916 984 558 320	91.8 78.6 67.2 56.2 48.6
U	07.0	L).)	4,720	204	43.0

Heat balance error = 2.3%

Run 13	G = 154,000) lb/hrft ²	T water in	= 63.8	R-12	
		Ν	leasured			
Sec No	T vapor	To out	W water	∆t _w		
1 2 3 4 5 6	97.0 96.2 95.7 95.2 95.0 90.0	76.9 73.5 72.6 71.1 70.0 68.0	2,800 2,030 1,590 1,850 2,390 1,730	2.04 1.62 1.67 1.30 1.03 0.78		
		Ca	alculated			
Sec No	^T o _{in}	ΔT	Q/A	h	x m	
1 2 3 4 5 6	79.3 75.0 73.7 72.1 71.1 68.6	17.7 21.2 22.0 23.1 23.9 21.4	14,700 8,850 6,840 6,220 6,360 3,480	831 417 311 269 266 163	75.3 35.6 9.6 - -	
Heat bal	lance error =	= 17.2%				
Run 14	G = 326,000) lb/hrft ²	T water in	= 63.9	R-12	
]	Measured			
Sec No	T vapor	Toout	Wwater	ΔT_{w}		
1 2 3 4 5 6	98.5 98.1 97.9 97.5 97.0 96.5	81.8 80.7 78.9 76.1 73.3 80.6	1,870 1,310 1,660 1,950 1,760 2,280	3.20 3.33 2.80 2.16 1.53 1.12		
Calculated						
Sec No	^T o _{in}	$\Delta \mathbf{T}$	Q/A	h	× m	
1 2 3 4 5 6	84.3 82.6 81.0 77.9 73.5 81.7	14.2 15.5 16.9 19.4 23.5 24.8	15,500 11,300 12,000 10,800 6,990 6,600	1,090 728 710 556 298 266	87.7 66.4 47.8 29.5 15.8 4.4	

Heat balance error = 2.7%

Run 15	G = 425,00	00 1b/hrft ²	T water in	= 64.9	R-12
			Measured		
Sec No	T vapor	To out	Wwater	∆t _w	
1	99.0	83.9	2,240	3.20	
2	98.5	81.4	2,720	2.45	
3	97.8	80.0	2,120	2.60	
4	97.4	77.8	2,870	2.00	
5	96.5	74.7	2,140	1.67	
6	96.4	72.8	2,380	1.30	
		C	alculated		
Sec	Т				v
No	in	$\Delta \mathbf{T}$	Q/A	h	'n
1	86.9	12.1	18,600	1,540	88.7
2	84.2	14.3	17,200	1,200	66.8
3	82.3	15.5	14,200	916	47.5
4	80.2	17.2	14,800	860	29.8
5	76.2	20.3	9,210	454	15.1
6	74.1	22.3	8,000	359	4.5
Heat Ba	lance error	= 3.9%			
Run 16	G = 372,00	00 lb/hrft ²	T water in	= 63.5	R-12
			Measured		
Sec	Ŧ	Т	1.1	٨٣	
No	vapor	out	"water	∆ 1 w	
1	96.0	79.2	2,240	2.65	
2	95.4	76.5	2,720	2.07	
3	94.9	76.6	2,120	2.25	
4	94.4	74.3	2,870	1.67	
5	93.4	70.8	2,140	1.30	
6	92.5	69.7	2,380	1.02	
Sec	т	C	alculated		
No	oin	$\Delta \mathbf{T}$	Q/A	h	× m
1	81.7	14.3	15,350	1,070	89.1
2	78.9	16.5	14,550	882	68.5
3	78.6	16.3	12,300	755	50.0
4	76.3	18.1	12,400	685	32.7
5	81.9	21.5	6,900	322	19.3
6	70.7	21.8	6,270	288	14.9

Heat balance error = 1.06%

Run 17	G = 358,	000 1b/hrft	2 Twater :	= 65 In	R-12
			Measured		
Sec No	T vapor	^T out	Wwater	∆t _w	
1 2 3 4 5 6	101.0 100.8 100.1 99.5 98.8 98.4	83.3 81.6 79.2 75.8 72.5 71.2	2,240 2,720 2,120 2,870 2,140 2,380	3.10 2.62 2.43 1.68 1.31 1.02	
			Calculated		
Sec No	To in	$\Delta \mathbf{T}$	Q/A	h	×m
1 2 3 4 5 6	86.2 84.6 81.3 77.8 73.7 72.2	14.8 16.2 18.7 21.7 25.1 26.2	17,900 18,400 13,300 12,500 7,250 6,270	1,210 1,140 713 576 289 240	87.7 61.3 38.0 19.2 4.8
H eat ba	alance erro	pr = 4%			
Run 18	G = 506	,000 lb/hrft	² T _{water}	in ^{63.9}	R-12
			Measured		
Sec No	T vapor	Toout	W water	$\Delta \mathbf{T}_{\mathbf{w}}$	
1 2 3 4 5 6	91.0 90.3 89.0 88.1 87.0 86.0	79.4 76.7 74.7 75.0 73.0 71.4	2,200 1,970 2,210 1,880 1,490 1,340	2.64 2.24 1.90 1.98 1.74 1.50	
			Calculated		
Sec No	To in	Δτ	Q/A	h	x m
1 2 3 4 5	81.8 78.5 76.4 76.6 74.1	9.2 11.8 12.6 11.5 12.9	15,000 11,400 10,800 9,650 6,700	1,630 965 856 839 519	92.4 79.3 67.7 57.7 49.4
6 Heat ha	72.2 alance erro	13.8 or = 0.9%	5,200	377	43.5

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Run 19	G = 556,000	0 lb/hrft ²	^T water in	= 62.9	R-12
		I	Measured		
Sec No	T vapor	o _{out}	Wwater	$\Delta \mathbf{T}_{\mathbf{w}}$	
1 2 3	88.0 86.9 85.2	77.0 76.3 74.9	2,200 1,970 2,210	2.39 2.34 2.03	
4 5 6	83.9 81.6 81.2	71.3 68.0 68.1	1,880 1,490 1,340	1.51 1.04 1.00	
		Ca	alculated		
Sec No	^T o _{in}	$\Delta \mathbf{T}$	Q/A	h	× m
1 2 3 4 5	79.2 78.2 76.8 72.5 68.6	8.8 8.7 8.4 11.4 13.0	13,600 11,900 11,600 7,350 4,010	1,550 1,370 1,390 645 319	94.0 82.1 71.6 62.6 57.5
6	68.7	12.7	3,460	272	54.2
Heat ba	lance error =	= 22.8%			
Run 20	G = 257,000) lb/hrft ²	T water in	= 68.8	R-12
		Ν	leasured		
Sec No	^T vapor	To out	Wwater	∆T _w	
1 2 3 4 5 6	94.6 94.2 93.8 93.4 93.0 92.6	82.0 82.3 79.1 79.2 76.3 77.5	2,020 1,790 1,720 1,620 1,540 1,330	2.38 2.50 1.93 2.02 1.44 1.78	
		Ca	alculated		
Sec No	^T o _{in}	$\Delta \mathbf{T}$	Q/A	h	× m
1 2 3 4 5	84.0 84.2 80.5 80.6 77.2	10.6 10.0 13.3 12.8 15.8	12,400 11,600 8,600 8,480 5,720	1,170 1,160 647 660 362	84.6 79.0 58.5 47.7 38.9
6	78.5	14.1	6,130	435	4.2

Heat balance error = 20.8%

Run 21	G = 308,000) $1b/hrft^2$	T water in	= 68.8	R-12
		1	Measured		
Sec No	T vapor	T _{oout}	Wwater	ΔT_{w}	
1 2	104.2 103.9	88.7 90.4	2,020 1,790	3.38 3.26	
3 4 5	103.6 103.2 102.8	83.8 83.8 78.0	1,720 1,620 1,540	2.81 2.47 1.78	
6	102.4	78.9	1,330	2.05	
		(Calculated		
Sec No	T _o in	ΔT	Q/A	h	×m
1 2 3 4 5 6	91.7 92.8 85.8 85.5 79.1 80.0	12.5 11.1 17.8 19.7 23.7 22.4	18,500 15,100 12,500 10,700 7,090 7,050	1,480 1,360 703 543 299 315	84.0 55.4 25.7 14.6 3.0
Heat bal	lance error =	= 22.8%			
Run 22	G = 307.000	$1b/brft^2$	т	= 79	R-12
	,	,	water in		
Sec	_	т	Measured		
No	^T vapor	out	Water	∆T ₩	
1 2 3	109.0	91.8	1 070		
4 5 6	108.1 107.9 107.7 107.5 107.2	91.8 90.9 90.6 89.6 89.2	1,870 1,710 1,450 1,420 1,400 1,270	2.33 2.39 2.32 2.26 2.09 2.07	
4 5 6	108.1 107.9 107.7 107.5 107.2	91.8 90.9 90.6 89.6 89.2 Ca	1,870 1,710 1,450 1,420 1,400 1,270 alculated	2.33 2.39 2.32 2.26 2.09 2.07	
4 5 6 Sec No	108.1 107.9 107.7 107.5 107.2 T _o in	91.8 90.9 90.6 89.6 89.2 ℃a	1,870 1,710 1,450 1,420 1,400 1,270 alculated	2.33 2.39 2.32 2.26 2.09 2.07	×m
4 5 6 No 1 2 3 4 5 6	108.1 107.9 107.7 107.5 107.2 Tonin 93.6 93.5 92.3 91.9 90.8 90.3	91.8 90.9 90.6 89.6 89.2 Ca ΔT 15.2 14.6 15.6 15.8 16.7 16.9	1,870 1,710 1,450 1,420 1,400 1,270 alculated Q/A 11,250 10,600 8,700 8,300 7,560 6,800	2.33 2.39 2.32 2.26 2.09 2.07 h h 740 723 553 525 453 402	xm 90.2 71.2 54.3 39.4 25.5 13.0

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Run 23	G = 314,000	0 lb/hrft ²	T water in	= 79.5	R-
			Measured		
Sec	T	Т	IJ	٨٣	
No	vapor	out	water	Δ Ι w	
1	110.5	95.8	1,870	2.77	
2	110.0	94.5	1,710	2,79	
3	109.6	93.7	1,450	2.78	
4	109.2	92.6	1.420	2.56	
5	108.8	91.4	1,400	2.36	
6	108.4	91.0	1,270	2.36	
		С	alculated		
Sec	T				
No	o _{in}	$\Delta \mathbf{T}$	Q/A	h	хп
1	97.0	13.5	13,400	992	88
2	96.5	13.5	12,300	914	66
3	95.4	14.2	10,400	734	44
4	94.1	15.1	9,400	623	31
5	92.8	16.0	8,550	534	14
6	92.3	16.1	7,750	482	
Run 24	G = 327,000) lb/hrft ²	T water in	= 79.1	R-
]	Measured		
Sec	т	T	IJ	٨٣	
No	'vapor	čout	"water	w w	
1	118.7	97.5	1,870	3.32	
2	118.4	98.2	1,710	3.57	
3	118.1	96.4	1,450	3.38	
4	117.8	92.9	1,420	2.69	
5	117.5	88.9	1,400	2.12	
6	117.1	88.9	1,270	2.00	
		C	alculated		
Sec	T _{o.}	A m	0.14	1	x
NO	in	ΔT	Q/A	n	n
1	100.1	18.6	16,100	865	86
2	100.8	17.6	15,800	896	59
3	98.4	19.7	12,700	643	35
4	94.5	23.3	9,870	424	16
5	90.1	27.4	7,670	280	2
6	90.0	27.1	6,560	242	

Heat balance error = 4.7

Figure Captions

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Figure No.

1	Control Volume of a Tube Element
2	Elemental Volume In the Condensate
3	U _i /U _l vs δ ⁺
4	Dimensionless Film Thickness δ^+
5	Stanton Number St*
6	Schematic Diagram of Apparatus
7	Heat Transfer Data for R-12
8	Heat Transfer Data for R-22
9	Pressure Drop Data for R-22
10	Predicted vs Measured Heat Transfer Data, R-12
11	Predicted vs Measured Heat Transfer Data, R-22
12	Predicted vs Measured Pressure Drop Data, R-22
13	Comparison of Data with Akers-Rosson Recommended Correlation
14	Dimensionless Local Heat Transfer Coefficient







FIG. 2



FIGURE 3 δ^+ vs. B





FIG 5



FIGURE 6 SCHEMATIC DIAGRAM OF APPARATUS



FIG.7 LOCAL HEAT TRANSFER COEFFICIENT FOR R-12 COMPARED WITH ANALYSIS



FIGURE 8 LOCAL HEAT TRANSFER COEFFICIENTS



FIGURE 9 TOTAL STATIC PRESSURE GRADIENTS





FIGURE II DATA COMPARED WITH PRESENT ANALYSI



R-22 PRESSURE DROP DATA COMPARED WITH ANALYSIS

FIG. 12



Fig.13. Data on Akers-Rosson Plot



FIGURE 14 DIMENSIONLESS LOCAL HEAT TRANSFER COEFFICIENTS (Pr=5)