

MIT Open Access Articles

Logarithmic Fluctuations for Internal DLA

The MIT Faculty has made this article openly available. **Please share** how this access benefits you. Your story matters.

Citation: Jerison, David, Lionel Levine, and Scott Sheffield. "Logarithmic Fluctuations for Internal DLA." *Journal of the American Mathematical Society* 25.1 (2011).

As Published: <http://dx.doi.org/10.1090/s0894-0347-2011-00716-9>

Publisher: American Mathematical Society (AMS)

Persistent URL: <http://hdl.handle.net/1721.1/71522>

Version: Author's final manuscript: final author's manuscript post peer review, without publisher's formatting or copy editing

Terms of use: Creative Commons Attribution-Noncommercial-Share Alike 3.0



Logarithmic Fluctuations for Internal DLA

David Jerison* Lionel Levine† Scott Sheffield‡

July 6, 2011

Abstract

Let each of n particles starting at the origin in \mathbb{Z}^2 perform simple random walk until reaching a site with no other particles. Lawler, Bramson, and Griffeath proved that the resulting random set $A(n)$ of n occupied sites is (with high probability) close to a disk \mathbf{B}_r of radius $r = \sqrt{n/\pi}$. We show that the discrepancy between $A(n)$ and the disk is at most logarithmic in the radius: i.e., there is an absolute constant C such that with probability 1,

$$\mathbf{B}_{r-C \log r} \subset A(\pi r^2) \subset \mathbf{B}_{r+C \log r} \quad \text{for all sufficiently large } r.$$

*Partially supported by NSF grant DMS-1069225

†Supported by an NSF Postdoctoral Research Fellowship.

‡Partially supported by NSF grant DMS-0645585.

2010 Mathematics Subject Classification: 60G50, 60K35, 82C24.

Contents

1	Introduction	2
2	Logarithmic fluctuations	3
2.1	No thin tentacles	6
2.2	A priori bound on probability of large fluctuations	7
2.3	Brownian exit times	9
2.4	The discrete harmonic function $H_\zeta(z)$	10
2.5	The martingale $M_\zeta(t)$	16
2.6	Early points imply late points	21
2.7	Late points imply early points	24
3	Concluding Remarks	28
A	Appendix: Proof of Lemmas 2 and 3	31
A.1	Bounds on the probability of thin tentacles in all dimensions	31
A.2	Bound on the probability of very late points	36

1 Introduction

In the process known as internal diffusion limited aggregation (IDLA) one constructs for each integer time $n \geq 0$ an **occupied set** $A(n) \subset \mathbb{Z}^2$ as follows: begin with $A(1) = \{0\}$, and for each n let $A(n+1)$ be $A(n)$ plus the first point at which a simple random walk from the origin hits $\mathbb{Z}^2 \setminus A(n)$.

IDLA was first proposed by Meakin and Deutch [MD86] as a model of industrial chemical processes such as electropolishing, corrosion and etching. After discussing these applications, they write “it is also of some fundamental significance to know just how smooth a surface formed by diffusion limited processes may be.” Their numerical findings indicated that such surfaces are astonishingly smooth: in two and higher dimensions, the fluctuations appeared to be at most logarithmic in the number of particles.

In the more than two decades since, this finding has resisted mathematical proof. Diaconis and Fulton [DF91] identified IDLA as a special case of a “smash sum” operation on subsets of \mathbb{Z}^2 . Lawler, Bramson and Griffeath [LBG92] proved that the asymptotic shape of the domain $A(n)$ is a disk, and Lawler [Law95] showed that the fluctuations from circularity are at most of order $r^{1/3}$ up to logarithmic factors, where $r = \sqrt{n/\pi}$ is the radius of the disk of area n . Later work [GQ00, LP10] related internal DLA with multiple sources to variational problems in PDE. However, the basic issue raised by Meakin and Deutch and again by Lawler (“A more interesting question... is whether the errors are $o(n^\alpha)$ for some $\alpha < 1/3$.” [Law95]) stood until Asselah and Gaudillière [AG10a] (with an appendix by Blachère) improved the bound of $r^{1/3}$ in dimensions $d \geq 3$ to $r^{1/(d+1)}$.

Our main result shows that in dimension 2, the fluctuations from circularity are at most order $\log r$. In fact we show somewhat more, namely that the probability that the fluctuations at time $n = \pi r^2$ exceed $a \log r$ decays like $r^{-\gamma}$ (where γ can be taken to be arbitrarily large, but the constant a may depend on γ). Let

$$\mathbf{B}_r = \{z \in \mathbb{Z}^2 \mid z_1^2 + z_2^2 < r^2\}$$

be the set of lattice points lying in the open disk of radius r centered at 0. For convenience, we define the IDLA cluster $A(t)$ for all $t \in \mathbb{R}$ by setting $A(t) = A(\lfloor t \rfloor)$, where $\lfloor t \rfloor$ is the greatest integer $\leq t$ (and $A(t) = \emptyset$ if $t < 1$).

Theorem 1. *For each γ there exists an $a = a(\gamma) < \infty$ such that for all sufficiently large r ,*

$$\mathbb{P}\left(\{\mathbf{B}_{r-a \log r} \subset A(\pi r^2) \subset \mathbf{B}_{r+a \log r}\}^c\right) \leq r^{-\gamma}.$$

The proof will show that $a(\gamma)$ can be taken to be an affine function of γ . Taking $C = a(\gamma)$ for large enough γ , Theorem 1 in particular implies (by Borel-Cantelli) that the fluctuations from circularity are $O(\log r)$ almost surely:

$$\mathbb{P}\{\mathbf{B}_{r-C \log r} \subset A(\pi r^2) \subset \mathbf{B}_{r+C \log r} \text{ for all sufficiently large } r\} = 1. \quad (1)$$

This is the first part of a three part article. In the second part, entitled *Internal DLA in higher dimensions* [JLS10], we show that internal DLA in \mathbb{Z}^d for $d \geq 3$ has even smaller fluctuations than in dimension 2: the higher-dimensional analogue of (1) holds with $\log r$ replaced by $\sqrt{\log r}$. We believe that these orders ($\log r$ in dimension 2 and $\sqrt{\log r}$ in dimensions $d \geq 3$) are the best possible. We discuss a possible approach to proving matching lower bounds, as well as several other open questions, in Section 3. Using methods rather different from ours, Asselah and Gaudillière [AG10b] have recently announced an independently obtained upper bound of order $\log^2 r$ in all dimensions $d \geq 2$.

In the third part, entitled *Internal DLA and the Gaussian free field* [JLS11], we show that the fluctuations from circularity have a weak limit which is a variant of the Gaussian free field. While both this paper and the sequel address fluctuations of IDLA, the results are of a rather different nature: the present paper is concerned with fluctuations at the fine scale of individual lattice points. The sequel addresses average fluctuations, and will show that in a certain sense the fluctuation from circularity *averaged* over a constant fraction of the boundary (here inner and outer fluctuations cancel each other out) is only *constant* order. In fact, the fluctuations from circularity converge in law to a particular Gaussian random distribution on the circle, whose time evolution is related to the Gaussian free field.

2 Logarithmic fluctuations

We begin in Section 2.1 with a lemma that rules out “thin tentacles” in the IDLA cluster (Lemma 2). More precisely, it is unlikely that $z \in A(n)$ and less than a

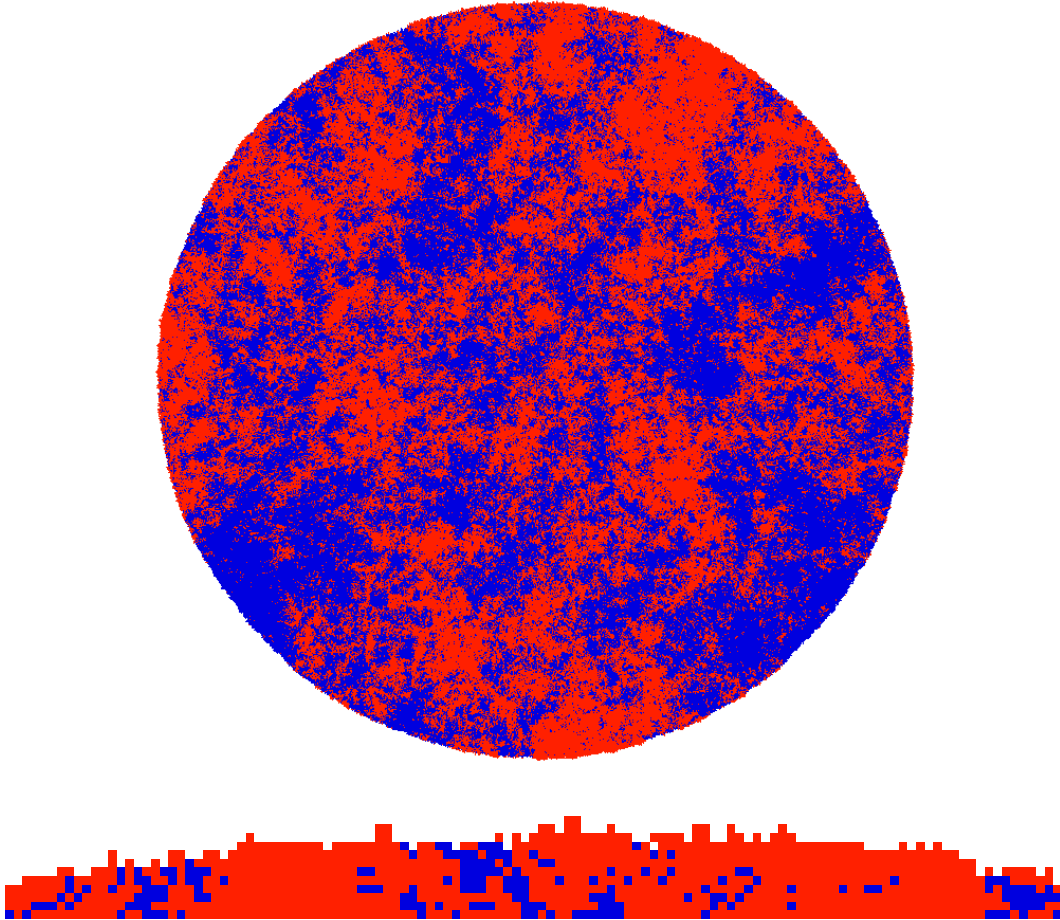


Figure 1: Top: IDLA cluster $A(n)$ for $n = 4 \times 10^5$. Early points (see Section 2.2) are colored red, and late points are colored blue. Bottom: closeup of a portion of the boundary of the cluster.

constant fraction bm^2 of lattice sites in the disk $\mathbf{B}(z, m)$ belong to $A(n)$. Section 2.2 establishes a rather rough a priori bound on the probability of certain types of very large fluctuations. Section 2.3 gives some large deviation results for the exit time of Brownian motion from an interval. These three sections comprise the background results needed for our argument.

In order to control more precisely the deviation from circularity near a point $\zeta \in \mathbb{Z}^2$, we define in Section 2.4 a discrete harmonic function $H_\zeta : \mathbb{Z}^2 \rightarrow \mathbb{R}$ which approximates the harmonic function

$$F_\zeta(z) := \operatorname{Re} \left(\frac{\zeta/|\zeta|}{\zeta - z} \right), \quad (2)$$

where z and ζ are viewed as complex variables. (Discrete harmonicity will fail for H_ζ only at a few points near ζ .) The function $H_\zeta(z) - \frac{1}{2|\zeta|}$ approximates the discrete Poisson kernel for the disk $\mathbf{B}_{|\zeta|}$, that is, the probability that simple random walk started at z first exits $\mathbf{B}_{|\zeta|}$ at ζ . We may extend the function H_ζ linearly to the grid \mathcal{G} of horizontal and vertical line segments joining vertices of \mathbb{Z}^2 . If we begin with a collection of particles on the grid \mathcal{G} , and we allow some of them to move according to Brownian motions on \mathcal{G} (see Section 2.5), according to determined starting and stopping rules, then the sum of H_ζ over these particle locations is a continuous martingale (provided that no particle continues to move after reaching a point where discrete harmonicity fails).

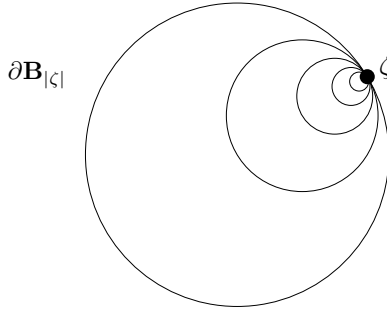


Figure 2: Level sets of F_ζ . Note that $F_\zeta = \frac{1}{2|\zeta|}$ on $\partial\mathbf{B}_{|\zeta|}$ and $F_\zeta(0) = \frac{1}{|\zeta|}$.

By the martingale representation theorem (see [RY05, Theorem V.1.6]), we can represent any continuous martingale as a time change of a standard Brownian motion. Using the representation theorem and large deviations for Brownian exit times, we can control the likelihood that this martingale is larger or smaller than its mean by estimating its quadratic variation.

We will ultimately show that large deviations of $A(n)$ from circularity are unlikely without large deviations of the martingale for some ζ . For example, the fact that thin tentacles are unlikely (Lemma 2) means that with high probability, a point near some ζ cannot belong to $A(n)$ unless many points near ζ belong to $A(n)$. Since $H_\zeta(z)$ is large for z close to ζ , this can be used to show that (when ζ is outside the typical range of $A(n)$) the martingale is likely to be large if a point near ζ belongs to $A(n)$. Similarly, we can show that if ζ fails to be hit (long after we expect it to have been hit) then the martingale has to be small at that time. (For this step we actually use a modified version of the process in which particles are frozen when they first exit a set called Ω_ζ that approximates $\mathbf{B}_{|\zeta|}$, see Figure 5.)

The martingale is a Brownian motion parameterized by its quadratic variation time, but the amount of quadratic variation time elapsed while $A(n)$ is constructed is random. Note that the conditional expected amount of time elapsed when depositing one particle (given the locations of all previous particles) is the conditional variance

of the value of H_ζ at the position that the particle is deposited. This value is easy to approximate on the event that $A(n)$ is approximately circular. If, as a heuristic, we *assume* that $A(n)$ is approximately circular for all n , then a back of the envelope calculation shows that the total time lapse in constructing $A(\pi|\zeta|^2)$ should be of order $\log|\zeta|$ in dimension two (and of constant order in all higher dimensions). On the event that the total time lapse is about right, the probability of having a martingale fluctuation of $C\sqrt{\log|\zeta|}$ standard deviations should decay like $e^{-\frac{C^2}{2}\log|\zeta|} = |\zeta|^{-C^2/2}$. If C is large enough then we should be able to sum over all ζ and show that we are unlikely to see fluctuations of this magnitude anywhere.

This heuristic actually suggests the basic strategy of our argument: we assume some amount of regularity for the growth of $A(n)$ and use that assumption to show that the quadratic variation time lapse is unlikely to be large, hence large martingale fluctuations are unlikely, and hence the growth of $A(n)$ is *even more regular* than we assumed. This strategy is pursued in Sections 2.6 and 2.7, which contain the heart of proof. Using the ideas above, we will show that the presence of an “early point” (a point z visited at time less than $\pi|z|^2$) implies, with high probability, the presence of a comparably late point at a previous time (Lemma 12); similarly, the presence of a late point implies, with high probability, the presence of a *significantly earlier* point at a previous time (Lemma 13). By iterating these two lemmas, we will see that if there were a substantial probability to have even a slightly early or late point (off by more than $C \log r$), then there would be a substantial probability of having a much larger fluctuation, which would contradict our a priori estimate.

In general, we use C_0, C_1, \dots to denote large constants and c_0, c_1, \dots to denote small constants. Unless explicitly specified otherwise, these are all positive absolute constants.

2.1 No thin tentacles

Let $A(n)$ be the internal DLA cluster formed from n particles started at the origin in \mathbb{Z}^2 . For $z \in \mathbb{Z}^2$ and a positive real number m , let

$$\mathbf{B}(z, m) = \{y \in \mathbb{Z}^2 : |y - z| < m\}$$

where $|z| = (z_1^2 + z_2^2)^{1/2}$ is the Euclidean norm.

The following lemma is a variant of [LP10, Lemma 5.12]. The basic estimates involved are similar to those used by Lawler, Bramson and Griffeath in their proof of the outer bound; see in particular [LBG92, Lemma 5b].

Lemma 2. *There are positive absolute constants b, C_0 , and c_0 such that for all real numbers $m > 0$ and all $z \in \mathbb{Z}^2$ with $|z| \geq m$,*

$$\mathbb{P} \{z \in A(n), \#(A(n) \cap \mathbf{B}(z, m)) \leq bm^2\} \leq C_0 e^{-c_0 m}. \quad (3)$$

In other words, it is unlikely that z belongs to the cluster $A(n)$ but less than a constant fraction of sites in the ball $\mathbf{B}(z, m)$ belong to the cluster. Figure 3 depicts this unlikely scenario.

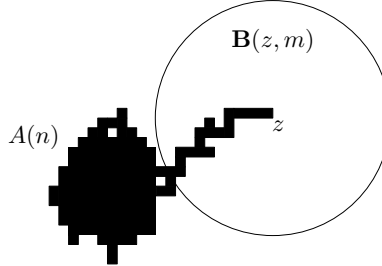


Figure 3: A thin tentacle.

In fact, the exponent on the right side of (3) can be improved to order $m^2/\log m$ (and to order m^2 in dimensions $d \geq 3$). In appendix A.1, we prove these stronger bounds in all dimensions. We anticipate using these bounds for $d \geq 3$ in our sequel paper *Internal DLA in higher dimensions*. For purposes of the present paper, we will use only the bound (3) in dimension $d = 2$.

2.2 A priori bound on probability of large fluctuations

We expect a point $z \in \mathbb{Z}^2$ to first join the IDLA cluster $A(n)$ at time about $n = \pi|z|^2$. This motivates the following two definitions, which are illustrated in Figure 4.

Definition. For a positive real number m , we say that $z \in \mathbb{Z}^2$ is *m -early* if $z \in A(\pi(|z| - m)^2)$. For a time N , let

$$\mathcal{E}_m[N] := \bigcup_{z \in A(N)} \{z \in A(\pi(|z| - m)^2)\}$$

be the event that some point in $A(N)$ is m -early.

Definition. For a positive real number ℓ , we say that $z \in \mathbb{Z}^2$ is *ℓ -late* if $z \notin A(\pi(|z| + \ell)^2)$. For a time N , let

$$\mathcal{L}_\ell[N] = \bigcup_{z \in \mathbf{B}_{\sqrt{N/\pi} - \ell}} \{z \notin A(\pi(|z| + \ell)^2)\}$$

be the event that some point in $\mathbf{B}_{\sqrt{N/\pi} - \ell}$ is ℓ -late.

The next lemma gives the a priori estimate on the probability of very late points that we will use to get the proof started in Sections 2.6 and 2.7. It follows easily from Lemma 6 of [LBG92]; for completeness, we include a proof in appendix A.2.

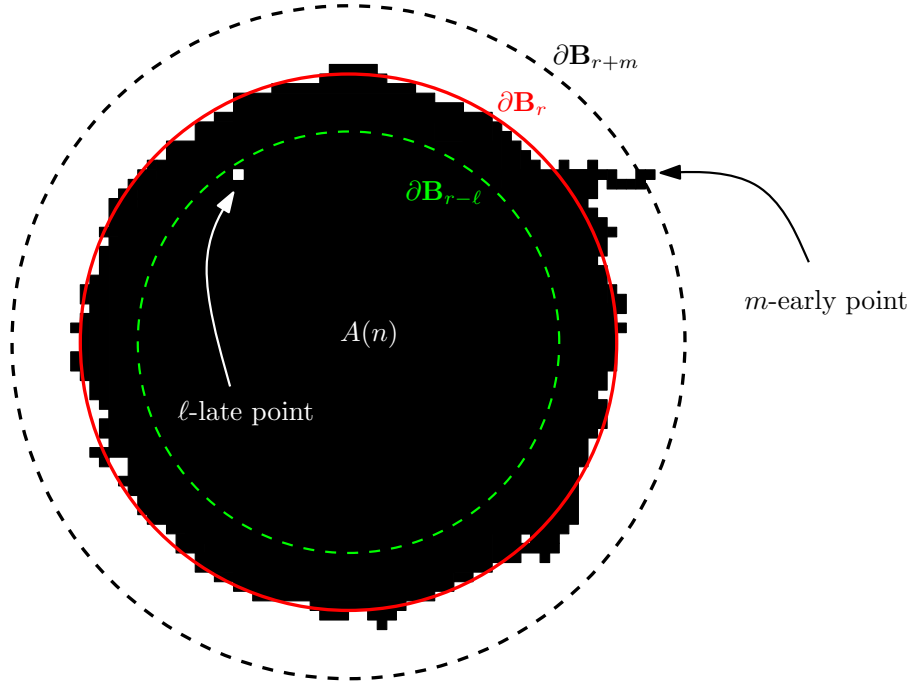


Figure 4: Illustration of an IDLA cluster $A(n)$ (union of the black squares), an m -early point and an ℓ -late point. The early point belongs to $A(n)$ despite lying outside the disk \mathbf{B}_{r+m} , where $r = \sqrt{n/\pi}$; the late point does not belong to $A(n)$ despite lying inside the disk $\mathbf{B}_{r-\ell}$.

Lemma 3. *There are absolute constants $C_0, c_0 > 0$ such that for all real $\ell > 0$,*

$$\mathbb{P}(\mathcal{L}_\ell[100\pi\ell^2]) \leq C_0 e^{-c_0\ell}. \quad (4)$$

To avoid referring to too many unimportant constants, we take C_0 large enough and c_0 small enough so that the same C_0 and c_0 work in Lemmas 2 and 3.

To prove Theorem 1 we will show that for each γ there is a constant $a = a(\gamma)$ such that if r is sufficiently large and $m = \ell = a \log r$ and $n = \pi r^2$, then

$$\mathbb{P}(\mathcal{E}_m(n) \cup \mathcal{L}_\ell(n)) \leq r^{-\gamma}. \quad (5)$$

The next lemma is an easy observation which shows that (5) implies Theorem 1.

Lemma 4. *For all $\ell, m, N > 0$ we have*

$$\begin{aligned} \mathcal{E}_m[N]^c &= \bigcap_{n \leq N} \left\{ A(n) \subset \mathbf{B}_{\sqrt{n/\pi}+m} \right\} \\ \mathcal{L}_\ell[N]^c &= \bigcap_{n \leq N} \left\{ \mathbf{B}_{\sqrt{n/\pi}-\ell} \subset A(n) \right\}. \end{aligned}$$

Proof. Taking complements in the definition, we obtain

$$\begin{aligned}
\mathcal{E}_m[N]^c &= \bigcap_{z \in A(N)} \{z \notin A(\pi(|z| - m)^2)\} \\
&= \bigcap_{z \in A(N)} \bigcap_{n \leq \min(N, \pi(|z| - m)^2)} \{z \notin A(n)\} \\
&= \bigcap_{n \leq N} \bigcap_{z: |z| \geq \sqrt{n/\pi} + m} \{z \notin A(n)\} \\
&= \bigcap_{n \leq N} \left\{ A(n) \subset \mathbf{B}_{\sqrt{n/\pi} + m} \right\}.
\end{aligned}$$

Likewise,

$$\begin{aligned}
\mathcal{L}_\ell[N]^c &= \bigcap_{z \in \mathbf{B}_{\sqrt{N/\pi} - \ell}} \{z \in A(\pi(|z| + \ell)^2)\} \\
&= \bigcap_{z \in \mathbf{B}_{\sqrt{N/\pi} - \ell}} \bigcap_{\pi(|z| + \ell)^2 \leq n \leq N} \{z \in A(n)\} \\
&= \bigcap_{n \leq N} \bigcap_{z: |z| \leq \sqrt{n/\pi} - \ell} \{z \in A(n)\} \\
&= \bigcap_{n \leq N} \left\{ \mathbf{B}_{\sqrt{n/\pi} - \ell} \subset A(n) \right\}. \quad \square
\end{aligned}$$

2.3 Brownian exit times

For $x \in \mathbb{R}$, write \mathbb{E}_x for the expectation on a probability space supporting a one-dimensional standard Brownian motion $\{B(s)\}_{s \geq 0}$ started at $B(0) = x$. For $a, b > 0$, let

$$\tau(-a, b) = \inf\{s > 0 \mid B(s) \notin [-a, b]\}$$

be the first exit time of Brownian motion from the interval $[-a, b]$.

Lemma 5. *Let $0 < a \leq b$ and $\lambda > 0$. If $\sqrt{\lambda}(a + b) \leq 3$, then*

$$\mathbb{E}_0 e^{\lambda \tau(-a, b)} \leq 1 + 10\lambda ab.$$

Proof. Let $c = (a + b)/2$, and for $x \in [-c, c]$ let

$$f(x) = \mathbb{E}_x e^{\lambda \tau(-c, c)}$$

By translating the interval $[-a, b]$ to $[-c, c]$, it suffices to show that if $\sqrt{\lambda}c \leq 3/2$, then

$$f(x) \leq 1 + 10\lambda(c + x)(c - x). \quad (6)$$

The function $e^{\lambda t} f(x)$ obeys the heat equation in the strip $[0, \infty) \times [-c, c]$, hence

$$e^{\lambda t} f''(x) = -\lambda e^{\lambda t} f(x).$$

Since f is even and $f(c) = 1$, we conclude that

$$f(x) = \frac{\cos(\sqrt{\lambda}x)}{\cos(\sqrt{\lambda}c)}.$$

Since $\cos x$ is concave in the interval $[0, \frac{3}{2}]$, we have

$$\cos(\sqrt{\lambda}x) \leq \cos(\sqrt{\lambda}c) + \sqrt{\lambda}(c-x)\sin(\sqrt{\lambda}c)$$

which gives

$$f(x) \leq 1 + \sqrt{\lambda}(c-x)\tan(\sqrt{\lambda}c).$$

Since $\tan \theta < 10\theta$ for $\theta \in [0, \frac{3}{2}]$, we conclude that for $x \in [0, c]$,

$$f(x) \leq 1 + 10\lambda c(c-x) \leq 1 + 10\lambda(c+x)(c-x)$$

which establishes (6). □

The following well-known estimate appears as Proposition II.1.8 in [RY05].

Lemma 6. *Let $\{B(s)\}_{s \geq 0}$ be a standard Brownian motion with $B(0) = 0$. Then for any $k, s > 0$*

$$\mathbb{P} \left\{ \sup_{s' \in [0, s]} B(s') \geq ks \right\} \leq e^{-k^2 s/2}.$$

2.4 The discrete harmonic function $H_\zeta(z)$

Fix a lattice point $\zeta \in \mathbb{Z}^2$. In this section we define the discrete harmonic function H_ζ resembling F_ζ in (2), which we will use to detect early points near ζ .

For $z \in \mathbb{Z}^2$, let $P_n(z)$ be the probability that an n -step simple random walk started at the origin in \mathbb{Z}^2 ends at z . The *recurrent potential kernel* for simple random walk in \mathbb{Z}^2 is defined by

$$g(z) = \sum_{n=0}^{\infty} (P_n(0) - P_n(z)). \tag{7}$$

This sum is finite for all $z \in \mathbb{Z}^2$, and its discrete Laplacian

$$\Delta g(z) = \frac{g(z+1) + g(z-1) + g(z+i) + g(z-i)}{4} - g(z) \tag{8}$$

vanishes for $z \neq 0$ and equals 1 if $z = 0$. Here we have identified \mathbb{Z}^2 with $\mathbb{Z} + i\mathbb{Z} \subset \mathbb{C}$. Fukai and Uchiyama [FU96] and Kozma and Schreiber [KS04] develop an asymptotic

expansion for $g(z)$. We will use only the following estimate: there are absolute constants λ and C_0 (rendered explicit in [KS04]; see also [LL10, §4.4]) such that

$$\left| g(z) - \frac{2}{\pi} \log |z| - \lambda \right| \leq C_0 |z|^{-2}. \quad (9)$$

We will use this sharp asymptotic estimate of the potential kernel to show that the discrete harmonic function H_ζ defined below is close to the continuum harmonic function F_ζ (Lemma 7c), which in turn gives a discrete mean value property: the average value of H_ζ on the discrete disk \mathbf{B}_r is very close to $H_\zeta(0)$ (Lemma 8c).

Evidently, $g(0) = 0$ and g is symmetric under the dihedral symmetries generated by the reflections across the axes $x+iy \mapsto \pm x \pm iy$, and the diagonal, $x+iy \mapsto y+ix$. Moreover, $g(z) > 0$ for all $z \neq 0$. Exact values of g near the origin can be computed using the McCrea-Whipple algorithm, described in [KS04]. In particular,

$$g(1) = 1, \quad g(1+i) = \frac{4}{\pi}.$$

(We will use these exact values for convenience in specifying constants. But what we really need is not the actual values, but the property that $g(1+i) > g(1)$, which is not surprising since $1+i$ is farther from the origin than 1.)

We construct the function H_ζ for $\zeta = x+iy$ in the sector $S = \{\zeta \mid 0 \leq y \leq x\}$. (For $\zeta \notin S$, we choose a dihedral symmetry ϕ such that $\phi\zeta \in S$, and define $H_\zeta(z) := H_{\phi\zeta}(\phi z)$.) Let

$$\rho = |\zeta|$$

and write

$$\zeta/\rho = \alpha_1 + \alpha_2(1+i)$$

with $\alpha_2 = y/\rho$ and $\alpha_1 = (x-y)/\rho$. The property we will use about the coefficients α_j for $j = 1, 2$ is that they are positive and bounded. Indeed $0 \leq \alpha_j \leq 1$.

Define

$$H_\zeta(z) = \frac{\pi}{2} [\alpha_1 g(z - \zeta - 1) + \alpha_2 g(z - \zeta - 1 - i) - (\alpha_1 + \alpha_2) g(z - \zeta)].$$

Note that (using the symmetries)

$$\begin{aligned} H_\zeta(\zeta) &= \frac{\pi}{2} [\alpha_1 g(1) + \alpha_2 g(1+i)] > 0 \\ H_\zeta(\zeta+1) &= \frac{\pi}{2} [\alpha_2 g(1) - (\alpha_1 + \alpha_2) g(1)] < 0 \\ H_\zeta(\zeta+1+i) &= \frac{\pi}{2} [\alpha_1 g(1) - (\alpha_1 + \alpha_2) g(1+i)] < 0. \end{aligned} \quad (10)$$

The last inequality follows from the fact that $g(1+i) > g(1)$.

Extend H_ζ linearly along each edge of the grid

$$\mathcal{G} = \{x+iy \mid x \in \mathbb{Z} \text{ or } y \in \mathbb{Z}\}.$$

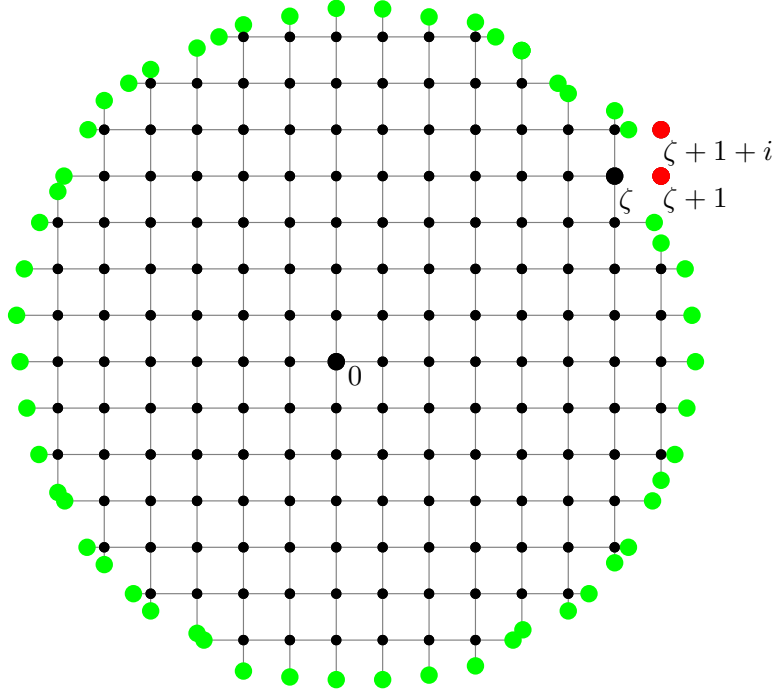


Figure 5: Illustration of the set Ω_ζ for $\zeta = 6 + 4i$. The function H_ζ is discrete harmonic except at the three points ζ , $\zeta + 1$, and $\zeta + 1 + i$. Its value is positive at ζ and negative at $\zeta + 1$ and $\zeta + 1 + i$. For general $z \in \mathcal{G}$, we have $H_\zeta(z) = \mathbb{E}H_\zeta(w)$ where w is the first of these three points hit by a Brownian motion in \mathcal{G} started at z . Large green dots on the edges of \mathcal{G} represent points $z \in \mathcal{G}$ for which $H_\zeta(z) = 1/2\rho$. Together with ζ , these points form the boundary of Ω_ζ . Small black dots indicate points in $\mathbb{Z}^2 \cap \Omega_\zeta$. For $z \in \Omega_\zeta$, the value $H_\zeta(z) - 1/2\rho$ is proportional to the probability that a Brownian motion in \mathcal{G} started at z first exits Ω_ζ at ζ .

(By an edge of \mathcal{G} we mean a line segment of length 1 whose endpoints lie in \mathbb{Z}^2 .) Now define Ω_ζ as the connected component of the origin in the set

$$\{z \in \mathcal{G} - \{\zeta\} : H_\zeta(z) > 1/2\rho\}.$$

Recall the discrete Laplacian Δ on \mathbb{Z}^2 , defined in (8). We say that a function f defined on a subset $\Lambda \subset \mathcal{G}$ is *grid-harmonic* if it is continuous and

- For each edge e of \mathcal{G} , if $I \subset e \cap \Lambda$ is an interval, then f is linear on I .
- For each vertex of \mathbb{Z}^2 in the interior of Λ , the sum of the slopes of the four line segments starting at z (directed away from z) is zero.

For example, this holds if $\Lambda = \mathcal{G}$ and f is discrete harmonic on vertices and linear on edges.

Lemma 7.

- (a) H_ζ is grid-harmonic on Ω_ζ , and $H_\zeta \geq 1/2\rho$ on Ω_ζ .
- (b) $\zeta \in \partial\Omega_\zeta$, and for all $z \in \partial\Omega_\zeta \setminus \{\zeta\}$ we have $H_\zeta(z) = 1/2\rho$.
- (c) There is an absolute constant $C_1 < \infty$ such that

$$|H_\zeta(z) - F_\zeta(z)| \leq C_1|z - \zeta|^{-2}.$$

In particular,

$$|H_\zeta(0) - 1/\rho| \leq C_1\rho^{-2}.$$

- (d) $1 \leq H_\zeta(\zeta) \leq 2$.

Proof.

- (a) By (10) we have $H_\zeta(\zeta + 1) < 1/2\rho$ and $H_\zeta(\zeta + 1 + i) < 1/2\rho$, so the points $\zeta + 1$ and $\zeta + 1 + i$ are not in Ω_ζ . Furthermore by definition $\zeta \notin \Omega_\zeta$. Since H_ζ is discrete harmonic at all other lattice points, its linear extension is grid-harmonic in Ω_ζ . Consequently, the inequality $H_\zeta \geq 1/2\rho$ follows from the maximum principle.
- (b) This is immediate from the definition of Ω_ζ .
- (c) The approximation for $|z - \zeta| \leq 10$ is trivial (choosing C_1 sufficiently large). For $|z - \zeta| > 10$, the bound (9) implies

$$H_\zeta(z) - \alpha_1(\log|z - \zeta - 1| - \log|z - \zeta|) - \alpha_2(\log|z - \zeta - 1 - i| - \log|z - \zeta|) = O\left(\frac{1}{|z - \zeta|^2}\right).$$

By the fundamental theorem of calculus,

$$\begin{aligned} & \alpha_1(\log|z - \zeta - 1| - \log|z - \zeta|) + \alpha_2(\log|z - \zeta - 1 - i| - \log|z - \zeta|) \\ &= -\alpha_1 \int_0^1 \operatorname{Re} \frac{1}{z - \zeta - t} dt - \alpha_2 \int_0^1 \operatorname{Re} \frac{1 + i}{z - \zeta - t(1 + i)} dt \\ &= -\alpha_1 \operatorname{Re} \frac{1}{z - \zeta} - \alpha_2 \operatorname{Re} \frac{1 + i}{z - \zeta} + O\left(\frac{1}{|z - \zeta|^2}\right) \\ &= F_\zeta(z) + O\left(\frac{1}{|z - \zeta|^2}\right). \end{aligned}$$

- (d) Using the exact values $g(1) = 1$ and $g(1 + i) = 4/\pi$, we have

$$H_\zeta(\zeta) = \frac{\pi}{2} \left(\alpha_1 + \frac{4}{\pi} \alpha_2 \right) = \frac{\pi}{2} \alpha_1 + 2\alpha_2$$

with the constraints $\alpha_i \geq 0$ and

$$(\alpha_1 + \alpha_2)^2 + \alpha_2^2 = 1.$$

Hence $\alpha_1 + \alpha_2 \leq 1$ and

$$(\pi/2)\alpha_1 + 2\alpha_2 \leq 2(\alpha_1 + \alpha_2) \leq 2.$$

Thus $H_\zeta(\zeta) \leq 2$. For the lower bound, one checks that α_1 and hence $(\pi/2)\alpha_1 + 2\alpha_2$ is a concave function of α_2 . At the extreme points $(\alpha_1, \alpha_2) = (1, 0)$ and $(\alpha_1, \alpha_2) = (0, 1/\sqrt{2})$, the values of $H_\zeta(\zeta)$ are $\pi/2$ and $\sqrt{2}$. Therefore, $H_\zeta(\zeta) \geq \sqrt{2} \geq 1$. \square

A corollary of Lemma 7 is the following.

Lemma 8. *There is an absolute constant $C_2 < \infty$ such that*

- (a) $\mathbf{B}_{\rho-C_2} \subset \Omega_\zeta \subset B_{\rho+C_2}$, where B_r is the disk of radius r in \mathbb{R}^2 .
- (b) $1/2\rho \leq H_\zeta(z) \leq 1/(\rho - r - C_2)$ for $0 \leq r \leq \rho - C_2$ and $z \in \mathbf{B}_r$.
- (c) For all $0 < r \leq \rho$,

$$\left| \sum_{z \in \mathbf{B}_r} (H_\zeta(z) - H_\zeta(0)) \right| + \left| \sum_{z \in \Omega_\zeta \cap \mathbb{Z}^2} (H_\zeta(z) - H_\zeta(0)) \right| \leq C_2 \log \rho.$$

Proof.

- (a) Recall that Lemma 7 (c) says

$$|H_\zeta(z) - F_\zeta(z)| \leq C_1/|z - \zeta|^2.$$

We will show that

$$F_\zeta(z) \leq \frac{1}{2\rho} - \frac{C}{|z - \zeta|^2}, \quad z \in \partial B_{\rho+2C} \quad (11)$$

$$F_\zeta(z) \geq \frac{1}{2\rho} + \frac{C}{|z - \zeta|^2}, \quad z \in \partial B_{\rho-2C}. \quad (12)$$

Since the level curve $\{z : F_\zeta(z) = 1/2\rho\}$ is the circle $\partial B_\rho \setminus \{\zeta\}$, this confirms part (a) with $C_2 = 2C_1$.

To prove (11) and (12), note first that

$$|\nabla F_\zeta(z)| = \frac{1}{|z - \zeta|^2}$$

Indeed, by the Cauchy-Riemann equations, this is the modulus of the (complex) derivative of $1/(\zeta - z)$. Any curve from ∂B_ρ to $\partial B_{\rho+C}$ has length at least C . In particular following the integral curves vector field ∇F_ζ we see that F_ζ changes by at least

$$\frac{C}{\max_z |z - \zeta|^2}$$

for z along the curve. These integral curves are the circles through ζ centered on the tangent line to ∂B_ρ at ζ , and $|z - \zeta|$ is nearly constant along the portion of the integral curve between ∂B_ρ and nearby concentric circles. Therefore the change in F_ζ is bounded as in (11), and similarly for (12).

- (b) The lower bound $1/2\rho$ holds for all $z \in \Omega_\zeta$. As we saw in (a), the level curves of F_ζ differ from level sets of H_ζ by at most a fixed distance $2C_1$. The supremum of F_ζ in B_r is $F_\zeta((r/\rho)\zeta) = 1/(\rho - r)$.
- (c) First suppose that $r \leq \rho - 2C_1$. By Lemma 7(c),

$$\left| \sum_{z \in \mathbf{B}_r} (H_\zeta(z) - H_\zeta(0)) - \sum_{z \in \mathbf{B}_r} (F_\zeta(z) - F_\zeta(0)) \right| \leq \sum_{z \in \mathbf{B}_r} \left(\frac{C_1}{|z - \zeta|^2} + \frac{C_1}{\rho^2} \right) \leq 8\pi C_1 \log \rho$$

Let B_r^\square be the union of unit squares around lattice points of \mathbf{B}_r . Since $|\nabla F_\zeta| \leq 1/|z - \zeta|^2$, we have (writing dz for area measure)

$$\left| \int_{B_r^\square} F_\zeta(z) dz - \sum_{z \in \mathbf{B}_r} F_\zeta(z) \right| \leq 4 \int_{B_r} \frac{dz}{|z - \zeta|^2} \leq 8\pi \log \rho.$$

Next,

$$\left| \int_{B_r^\square} F_\zeta(z) dz - \int_{B_r} F_\zeta(z) dz \right| \leq 2 \int_0^{2\pi} \left| \frac{1}{re^{i\theta} - \zeta} \right| r d\theta \leq 8\pi \log \rho. \quad (13)$$

Moreover,

$$\left| \int_{B_r} F_\zeta(0) dz - \sum_{z \in \mathbf{B}_r} F_\zeta(0) \right| = |\pi r^2 - \#\mathbf{B}_r|/\rho \leq 10r/\rho \leq 10.$$

Combining all these inequalities, we have shown that

$$\sum_{z \in \mathbf{B}_r} (H_\zeta(z) - H_\zeta(0))$$

is within $C_2 \log \rho$ of

$$\int_{B_r} (F_\zeta(z) - F_\zeta(0)) dz = 0. \quad (14)$$

The last equation (14) holds because F_ζ is harmonic in B_r .

We now turn to the cases of Ω_ζ and \mathbf{B}_r with $\rho - 2C_1 < r \leq \rho$. This follows from the case $\mathbf{B}_{\rho-2C_1}$, which we have already proved. In fact, these other sums differ only on lattice sites in $\rho - 2C_1 \leq |z| \leq \rho + 2C_1$, and the sum over these sites is majorized the same way as in (13),

$$\sum_{z \in \mathbf{B}_{\rho+2C_1} \setminus \mathbf{B}_{\rho-2C_1}} |H_\zeta(z)| \leq \sum_{z \in \mathbf{B}_{\rho+2C_1} \setminus \mathbf{B}_{\rho-2C_1}} \min\left(\frac{2}{|z - \zeta|}, 4\right) \leq 32\pi C_1 \log \rho. \quad \square$$

Remark. One can replace H_ζ with harmonic functions for which Lemma 8(a) and (b) hold with B_r replaced by shapes other than disks. However, the approximate mean value property Lemma 8(c) relies on B_r being a disk. The main ingredient is the continuum mean value property (14), which can be rephrased as saying that harmonic measure on ∂B_r from the origin is the same as from a uniform point in B_r . Accordingly, the place where the approximate mean value property Lemma 8(c) is used marks a crucial step in the proof: The bound for \mathbf{B}_r is used in the proof of Lemma 12, and the bound for Ω_ζ in the proof of Lemma 13.

2.5 The martingale $M_\zeta(t)$

For each integer n , recall that the set $A(n)$ consists of n lattice points with $A(n) \subset A(n+1)$ growing according to the IDLA scheme. For $\zeta \in \mathbb{Z}^2$, we will define a modified cluster $A_\zeta(n)$ in which particles are stopped on exiting Ω_ζ ; moreover, we will define $A_\zeta(t)$ for all $t \in \mathbb{R}$, not just for integer values.

To define the modified cluster, recall first the notion of Brownian motion $\beta(s)$ on the grid

$$\mathcal{G} = \{x + iy \in \mathbb{C} \mid x \in \mathbb{Z} \text{ or } y \in \mathbb{Z}\}.$$

To construct $\beta(s)$, take a standard Brownian motion $B(s)$ starting at 0 at time $s_0 = 0$ until the first time $s = s_1$ that $B(s) = \pm 1$. For each open interval of the set $S_1 = \{s : 0 < s < s_1, B(s) \neq 0\}$ choose a direction (up, down, left, right) with equal probability 1/4, independently. Define a corresponding motion of a particle $\beta(s) = B(s)$ in an interval in which right was chosen (and $\beta(s) = B(s)i$ in an interval in which up was chosen, and similarly for left and down). On each interval except the final one, the particle returns to the origin. On the final interval it reaches a lattice point adjacent to the origin for the first time. We continue the process (centered at this new lattice point) until it hits a point adjacent to that one, etc.

Let $\tilde{\beta}_0(s), \tilde{\beta}_1(s), \dots$ be independent Brownian motions on the grid \mathcal{G} , and for each integer $n \geq 0$, let

$$s_n^* = \inf \left\{ s \geq 0 \mid \tilde{\beta}_n(s) \in (\mathbb{Z}^2 - A_\zeta(n)) \cup (\mathcal{G} - \Omega_\zeta) \right\}$$

be the first time that $\tilde{\beta}_n$ visits either a lattice point outside the cluster $A_\zeta(n)$, or any point outside the set Ω_ζ . Then for $0 \leq s \leq 1$, define

$$A_\zeta(n+s) = A_\zeta(n) \cup \{\beta_{n+1}(s)\}$$

where

$$\beta_n(s) = \tilde{\beta}_n \left(\min \left(\frac{s}{1-s}, s_n^* \right) \right), \quad 0 \leq s < 1$$

(and by definition, $\beta_n(1) = \tilde{\beta}_n(s_n^*)$). The choice of $\frac{s}{1-s}$ is not important here; for our purposes, any increasing function $f(s)$ satisfying $f(0) = 0$ and $f(s) \uparrow \infty$ as $s \uparrow 1$ would suffice in its place. Effectively, the particle moves at an ever-increasing rate so that it *a.s.* visits the set $(\mathbb{Z}^2 - A_\zeta(n)) \cup (\mathcal{G} - \Omega_\zeta)$ in less than unit time.

The points of $\partial\Omega_\zeta$ where (eventually) many particles accumulate include ζ , but are typically non-lattice points on the grid \mathcal{G} . We regard $A_\zeta(t)$ as a multiset: it may have points of multiplicity greater than one on the boundary $\partial\Omega_\zeta$, or at interior vertices at times when the moving particle is on top of an occupied site. When summing a function over $A_\zeta(t)$, we take these multiplicities into account; thus for a function f defined on $\Omega_\zeta \cup \partial\Omega_\zeta$ and $t \in [n, n+1)$, the notation $\sum_{z \in A_\zeta(t)} f(z)$ means,

$$\sum_{z \in A_\zeta(t)} f(z) := \sum_{j=0}^{n-1} f(\beta_j(1)) + f(\beta_n(t-n)).$$

Recall the function H_ζ defined in Section 2.3. Because H_ζ is grid-harmonic in Ω_ζ (that is, discrete harmonic at lattice points, and linear on edges), the process

$$M_\zeta(t) := \sum_{z \in A_\zeta(t)} (H_\zeta(z) - H_\zeta(0))$$

(summing with multiplicity) is a continuous-time martingale adapted to the filtration

$$\mathcal{F}_t = \sigma\{A_\zeta(s) \mid 0 \leq s \leq t\}.$$

We will refer to M_ζ as the martingale with pole at ζ .

By the martingale representation theorem of Dambis and Dubbins-Schwarz (see, e.g. [RY05, Theorem V.1.6]),

$$M_\zeta(t) = B_\zeta(S_\zeta(t))$$

where $B_\zeta(s)$ is a standard Brownian motion and $S_\zeta(t) = \langle M_\zeta, M_\zeta \rangle_t$ is the quadratic variation of M_ζ .

Note that $S_\zeta(n)$ is a stopping time with respect to the filtration $\{\mathcal{F}_{T_\zeta(s)}\}_{s \geq 0}$, where $T_\zeta(s) = \inf\{t \mid S_\zeta(t) > s\}$; and $B_\zeta(s)$ is adapted to this filtration. By the strong Markov property ([MP10, Theorem 2.16]), for any n , the process

$$\tilde{B}_\zeta^n(u) := B_\zeta(S_\zeta(n) + u) - B_\zeta(S_\zeta(n))$$

is a standard Brownian motion that is independent of \mathcal{F}_n . It follows that we can write

$$B_\zeta(s) = B_\zeta(S_\zeta(n)) + \tilde{B}_\zeta^n(s - S_\zeta(n)), \quad s \in [S_\zeta(n), S_\zeta(n+1)].$$

where $\tilde{B}_\zeta^0(\cdot), \tilde{B}_\zeta^1(\cdot), \dots$ are independent standard Brownian motions started at 0.

For $-a < 0 < b$, let

$$\tau_n(-a, b) = \inf\{u > 0 \mid \tilde{B}_\zeta^n(u) \notin [-a, b]\}$$

be the first exit time of \tilde{B}_ζ^n from the interval $[-a, b]$.

Lemma 9. Fix $\zeta \in \mathbb{Z}^2$. Let

$$-a_n = \min_{z \in \partial A_\zeta(n)} (H_\zeta(z) - H_\zeta(0))$$

$$b_n = \max_{z \in \partial A_\zeta(n)} (H_\zeta(z) - H_\zeta(0)).$$

Then

$$S_\zeta(n+1) - S_\zeta(n) \leq \tau_n(-a_n, b_n).$$

Proof. By the maximum principle, since H_ζ is harmonic on $A_\zeta(n) \subset \Omega_\zeta$, we have

$$-a_n \leq H_\zeta(z) - H_\zeta(0) \leq b_n$$

for all $z \in A_\zeta(n)$. Hence for all $t \in [n, n+1]$, we have

$$-a_n \leq M_\zeta(t) - M_\zeta(n) \leq b_n.$$

That is, for all $s \in [S_\zeta(n), S_\zeta(n+1)]$, we have

$$-a_n \leq B_\zeta(s) - B_\zeta(S_\zeta(n)) \leq b_n.$$

That is, for $u \in [0, S_\zeta(n+1) - S_\zeta(n)]$, we have

$$-a_n \leq \tilde{B}_\zeta^n(u) \leq b_n.$$

Hence the first exit of \tilde{B}^n from the interval $[-a_n, b_n]$ occurs after time $S_\zeta(n+1) - S_\zeta(n)$, which gives $S_\zeta(n+1) - S_\zeta(n) \leq \tau_n(-a_n, b_n)$. \square

Fix $m \geq 1$. To detect the presence of m -early points in the aggregate $A(t)$, we will use the martingales M_ζ for ζ lying just outside the aggregate. The next lemma helps control the elapsed time $S_\zeta(t)$ on the event that there are no $(m+1)$ -early points. Define

$$r_0(t) = \sqrt{t/\pi} + 4m + 2C_2.$$

Lemma 10. For all $t > 0$ and all $\zeta \in \mathbb{Z}^2$ such that $|\zeta| \geq r_0(t)$,

$$\mathbb{E} \left[e^{S_\zeta(t)} \mathbf{1}_{\mathcal{E}_{m+1}[t]^c} \right] \leq t^{40}. \quad (15)$$

Proof. Fix t , and abbreviate $r_0 = r_0(t)$. On the event $\mathcal{E}_{m+1}[t]^c$, we have for all $n \leq t$

$$A(n) \subset \mathbf{B}_{\sqrt{n/\pi}+m+1}.$$

Hence by Lemma 8(b), on the event $\mathcal{E}_{m+1}[t]^c$ we have

$$-\frac{1}{2r_0} \leq H_\zeta(z) - H_\zeta(0) \leq \frac{1}{r_0 - \sqrt{n/\pi} - m - 2 - C_2}$$

for all positive integers $n \leq t$ and all $z \in A_\zeta(n)$.

We obtain from Lemma 9,

$$(S_\zeta(n) - S_\zeta(n-1)) \mathbf{1}_{\mathcal{E}_{m+1}[t]^c} \leq \tau_n(-a, b_n).$$

where

$$a = \frac{1}{2r_0}, \quad b_n = \frac{1}{r_0 - \sqrt{n/\pi} - m - 2 - C_2}.$$

Summing over $n = 1, \dots, t$ yields

$$e^{S_\zeta(t)} \mathbf{1}_{\mathcal{E}_{m+1}[t]^c} \leq e^{S_\zeta(t)} \mathbf{1}_{\mathcal{E}_{m+1}[t]^c} \leq e^{\sum_{n=1}^t \tau_n(-a, b_n)}.$$

The Brownian exit times $\tau_n(-a, b_n)$ for $n = 1, \dots, t$ are independent. Using Lemma 5 with $\lambda = 1$, we obtain

$$\log \mathbb{E} e^{\sum_{n=1}^t \tau_n(-a, b_n)} = \sum_{n=1}^t \log \mathbb{E} e^{\tau_n(-a, b_n)} \leq \sum_{n=1}^t \log(1 + 10ab_n) \leq \sum_{n=1}^t 10ab_n.$$

Setting $r_1 = r_0 - m - 3 - C_2$ we have

$$\begin{aligned} \sum_{n=1}^t b_n &\leq \int_0^t \frac{dn}{r_1 - \sqrt{n/\pi}} \\ &= \int_0^{\sqrt{t/\pi}} \frac{2\pi x dx}{r_1 - x} \\ &= 2\pi \int_{r_1 - \sqrt{t/\pi}}^{r_1} \frac{r_1 - y}{y} dy \\ &\leq 2\pi r_1 \log \frac{r_1}{r_1 - \sqrt{t/\pi}} \leq 2\pi r_0 \log \frac{r_0}{C_2}. \end{aligned}$$

Hence $\sum_{n=1}^t 10ab_n \leq 10\pi \log(r_0/C_2)$. We conclude that

$$\mathbb{E} \left[e^{S_\zeta(t)} \mathbf{1}_{\mathcal{E}_{m+1}[t]^c} \right] \leq (r_0/C_2)^{10\pi} \leq t^{40}. \quad \square$$

In order to study late points in the cluster $A(t)$, we will need a variant of Lemma 10 which helps control the elapsed time $S_\zeta(t)$ when $|\zeta|$ is slightly smaller than $r_0(t)$. In this case the quadratic variation of M_ζ is larger due to particles that accumulate on the boundary $\partial\Omega_\zeta$.

Lemma 11. *Fix $m \geq 2C_2 + 2$ and $\ell \leq m$. Fix $\zeta \in \mathbb{Z}^2$, and let $t = \pi(|\zeta| + \ell)^2$. If $|\zeta| \geq \ell$, then*

$$\mathbb{E} \left[e^{S_\zeta(t)} \mathbf{1}_{\mathcal{E}_m[t]^c} \right] \leq t^{40} e^{700m}. \quad (16)$$

Proof. Let $\rho = |\zeta|$, and suppose first that $\rho \geq 5m$. Let $t_0 = \lfloor \pi(\rho - 4m - 2C_2 - 1)^2 \rfloor$. Then $r_0(t_0) \leq \rho$, so by Lemma 10,

$$\mathbb{E} \left[e^{S_\zeta(t_0)} \mathbf{1}_{\mathcal{E}_m[t_0]^c} \right] \leq \mathbb{E} \left[e^{S_\zeta(t_0)} \mathbf{1}_{\mathcal{E}_{m+1}[t_0]^c} \right] \leq t_0^{40} \leq t^{40}.$$

(Here we have used the trivial inclusion $\mathcal{E}_m[t_0]^c \subset \mathcal{E}_{m+1}[t_0]^c$.)

Furthermore, because $H_\zeta \leq 2$ on all of $\partial\Omega_\zeta$, we have by Lemma 9 for $n = t_0 + 1, \dots, t$

$$S_\zeta(n) - S_\zeta(n-1) \leq \tau_n \left(-\frac{1}{2\rho}, 2 \right).$$

Hence, using the inclusion $\mathcal{E}_m[t]^c \subset \mathcal{E}_m[t_0]^c$,

$$e^{S_\zeta(t)} \mathbf{1}_{\mathcal{E}_m[t]^c} \leq e^{S_\zeta(t_0)} \mathbf{1}_{\mathcal{E}_m[t_0]^c} \leq e^{S_\zeta(t_0)} \mathbf{1}_{\mathcal{E}_m[t_0]^c} \prod_{n=t_0+1}^t e^{\tau_n \left(-\frac{1}{2\rho}, 2 \right)}.$$

By Lemma 5, we have $\log \mathbb{E} e^{\tau_n \left(-\frac{1}{2\rho}, 2 \right)} \leq \frac{10}{\rho}$ for $t_0 < n \leq t$, hence

$$\sum_{n=t_0+1}^t \log \mathbb{E} e^{\tau_n} \leq \frac{10}{\rho} (t - t_0) \leq 400m.$$

In the last inequality we have used that $t \leq \pi(\rho + m)^2$ and $m \geq 2C_2 + 2$.

Since the exit times τ_n for $n = t_0 + 1, \dots, t$ are independent of one another and of \mathcal{F}_{t_0} , we obtain

$$\begin{aligned} \log \mathbb{E} \left[e^{S_\zeta(t)} \mathbf{1}_{\mathcal{E}_m[t]^c} \right] &\leq \log \mathbb{E} \left[e^{S_\zeta(t_0)} \mathbf{1}_{\mathcal{E}_m[t_0]^c} \right] + \sum_{n=t_0+1}^t \log \mathbb{E} e^{\tau_n} \\ &\leq 40 \log t + 400m. \end{aligned}$$

It remains to consider the case $\ell \leq \rho \leq 5m$. In this case we take $t_0 = 0$, obtaining

$$\log \mathbb{E} e^{S_\zeta(t)} \leq \frac{10}{\rho} t = \frac{10\pi(\rho + \ell)^2}{\rho} \leq 40\pi\rho \leq 700m. \quad \square$$

2.6 Early points imply late points

Now we begin the main part of the proof of Theorem 1. Fix $\gamma \geq 1$. With constants b , C_0 and c_0 given by Lemma 2, and C_2 given by Lemma 7, we set

$$C_3 = \max \{(\gamma + 5 + \log C_0)/c_0, (1000/b)C_2\}.$$

(In particular, if $T \geq 3$ and $m \geq C_3 \log T$, then $C_0 e^{-c_0 m} < T^{-(\gamma+5)}$.)

Recall the events

$$\mathcal{E}_m[T] = \left\{ \text{some } z \in A(T) \text{ is } m\text{-early} \right\}$$

and

$$\mathcal{L}_\ell[T] = \left\{ \text{some } z \in \mathbf{B}_{\sqrt{T/\pi}-\ell} \text{ is } \ell\text{-late} \right\}$$

defined in Section 2.2.

Lemma 12. (*Early points imply late points*) Fix $\gamma \geq 1$, $T \geq 20$ and $m \geq C_3 \log T$. If $\ell \leq (b/1000)m$, then

$$\mathbb{P}(\mathcal{E}_m[T] \cap \mathcal{L}_\ell[T]^c) \leq T^{-(\gamma+1)}.$$

In other words, with high probability either there is no m -early point by time T , or there is a $(b/1000)m$ -late point by time T .

Proof. The proof consists of two main estimates. First, if $z \in \mathbb{Z}^2$ is the first m -early point and joins the cluster at time t , then for a suitably chosen point ζ near z , the quadratic variation $S_\zeta(t)$ is small: equation (18) below. Second, if in addition there are no ℓ -late points by time t , then the martingale $M_\zeta(t)$ is large: equation (19) below. To finish the proof we use the large deviation bound for Brownian exit times to show that these two events (small $S_\zeta(t)$ and large $M_\zeta(t)$) are unlikely both to occur.

Without loss of generality we assume that $\ell = (b/1000)m$. Note that $\mathcal{E}_m[T]$ is empty if $m > T^{1/2}$, so we may also assume $m \leq T^{1/2}$. For each integer $t = 1, 2, \dots, T$ and each lattice point z , let

$$Q_{z,t} = \{z \in A(t) - A(t-1)\} \cap \{z \in A(\pi(|z| - m)^2)\} \cap \mathcal{E}_m[t-1]^c$$

be the event that z first joins the cluster at time t and z is m -early, but no previous point is m -early. Since $A(T) \subset \mathbf{B}_T$,

$$\bigcup_{t \leq T} \bigcup_{z \in \mathbf{B}_T} Q_{z,t} = \mathcal{E}_m[T]. \tag{17}$$

As in Lemma 10, we write

$$r_0 = r_0(t) = \sqrt{t/\pi} + 4m + 2C_2.$$

Fix a lattice point $z \in \mathbf{B}_t$, and consider the unit square with lattice point corners containing $r_0 z/|z|$. Define $\zeta := \zeta_0(z, t)$ to be the corner of that square farthest from the origin. Then

$$r_0 \leq |\zeta_0(z, t)| < r_0 + 2.$$

By Lemma 10, we have for $\zeta = \zeta_0(z, t)$

$$\mathbb{E} e^{S_\zeta(t)} \mathbf{1}_{\mathcal{E}_{m+1}[t]^c} \leq T^{40}.$$

Hence by Markov's inequality, letting $s = (2\gamma + 100) \log T$, we have

$$\mathbb{P}(\mathcal{E}_{m+1}[t]^c \cap \{S_\zeta(t) > s\}) \leq T^{-(2\gamma+100)} \mathbb{E} e^{S_\zeta(t)} \mathbf{1}_{\mathcal{E}_{m+1}[t]^c} \leq T^{-(2\gamma+60)}.$$

On the event $Q_{z,t}$, there are no m -early points before time t ; moreover, z is at unit distance from some point of $A(t-1)$, so z is not $(m+1)$ -early. Therefore,

$$Q_{z,t} \subset \mathcal{E}_{m+1}[t]^c.$$

and hence

$$\mathbb{P}(Q_{z,t} \cap \{S_\zeta(t) > s\}) \leq T^{-(2\gamma+60)}. \quad (18)$$

Next, we show that for $\zeta = \zeta_0(z, t)$,

$$\mathbb{P}(\{M_\zeta(t) < (b/10)m\} \cap \mathcal{L}_\ell[t]^c \cap Q_{z,t}) < T^{-(\gamma+5)}. \quad (19)$$

In other words, on the event $E := \mathcal{L}_\ell[t]^c \cap Q_{z,t}$ it is likely that $M_\zeta(t) \geq (b/10)m$.

On the event $Q_{z,t}$, no $z' \in A(t)$ is $(m+1)$ -early, so we have $|z'| \leq \sqrt{t/\pi} + m + 1 < r_0 - C_2$ for all $z' \in A(t)$. In other words,

$$A(t) \subset \mathbf{B}_{r_0 - C_2}.$$

Hence Lemma 8(a) implies that $A_\zeta(t)$ does not meet $\partial\Omega_\zeta$. In other words, on the event $Q_{z,t}$ we have

$$A_\zeta(t) = A(t).$$

Partition $A_\zeta(t)$ into three disjoint sets as follows.

$$A_1 = A_\zeta(t) \cap \mathbf{B}_{\sqrt{t/\pi} - \ell}; \quad A_2 = A_\zeta(t) \cap \mathbf{B}(z, m); \quad A_3 = A_\zeta(t) \setminus (A_1 \cup A_2).$$

On the event E , the ball of radius $\sqrt{t/\pi} - \ell$ is filled by $A_\zeta(t)$, so

$$A_1 = \mathbf{B}_{\sqrt{t/\pi} - \ell};$$

and since $A_\zeta(t)$ has t points, $A_2 \cup A_3$ has at most $10\sqrt{t}\ell$ points. In particular, $(\#A_3)/r_0 \leq 10\sqrt{\pi}\ell$ and

$$\sum_{z' \in A_3} (H_\zeta(z') - H_\zeta(0)) \geq -(\#A_3)H_\zeta(0) \geq -20\ell = -(b/50)m,$$

since H is non-negative on Ω_ζ and, by Lemma 7(c), $H_\zeta(0)$ is close to $1/r_0$. Furthermore, Lemma 8(c) implies

$$\sum_{z' \in A_1} (H_\zeta(z') - H_\zeta(0)) \geq -C_2 \log r_0 \geq -C_2 \log T > -(b/1000)m,$$

where in the last step we have used that $C_2 < bC_3/1000$ and $C_3 \log T \leq m$. On the other hand, since $|z| \geq m$, it follows from Lemma 2 that

$$\mathbb{P}(\{\#A_2 \leq bm^2\} \cap E) < C_0 e^{-c_0 m} \leq T^{-(\gamma+5)}.$$

Next, recall that on the event $Q_{z,t}$, the point z is m -early but not $(m+1)$ -early, so that $\sqrt{t/\pi} + m \leq |z| \leq \sqrt{t/\pi} + m + 1$. Moreover, ζ is at distance roughly $3m$ from z along the ray 0 to z . $F_\zeta(z) \approx 1/3m$, and using Lemma 7(c) we find that $H_\zeta(z') \geq 1/4m + O(1/m^2)$ for all $z' \in \mathbf{B}(z, m)$. Furthermore, $H_\zeta(0) \approx 1/r_0$ and $r_0 \geq \sqrt{t/\pi} \geq 10m$, so $H_\zeta(z') - H_\zeta(0) \geq 1/8m$. Therefore, on the event $\{\#A_2 > bm^2\} \cap E$, we have

$$\sum_{z' \in A_2} (H_\zeta(z') - H_\zeta(0)) \geq (\#A_2) \frac{1}{8m} > bm^2 \frac{1}{8m} = (b/8)m.$$

Summing the contributions to $M_\zeta(t)$ from A_1 , A_2 and A_3 , we find that on the event $\{\#A_2 > bm^2\} \cap E$ we have $M_\zeta(t) > (\frac{b}{8} - \frac{b}{50} - \frac{b}{1000})m > (b/10)m$. Hence

$$\{M_\zeta(t) \leq (b/10)m\} \subset \{\#A_2 \leq bm^2\} \cup E^c.$$

Intersecting with E , we obtain

$$\mathbb{P}(\{M_\zeta(t) \leq (b/10)m\} \cap E) \leq \mathbb{P}(\{\#A_2 \leq bm^2\} \cap E) < T^{-(\gamma+5)}$$

which proves (19).

Recall that $s = (2\gamma + 100) \log T$. Since $m \geq C_3 \log T$ we have $(b/10)m \geq s$. Using the large deviation bound for the maximum of Brownian motion, Lemma 6, we obtain

$$\begin{aligned} \mathbb{P}(\{S_\zeta(t) \leq s\} \cap \{M_\zeta(t) \geq (b/10)m\}) &\leq \mathbb{P}(\{S_\zeta(t) \leq s\} \cap \{B_\zeta(S_\zeta(t)) \geq s\}) \\ &\leq \mathbb{P}\left\{\sup_{s' \in [0, s]} B_\zeta(s') \geq s\right\} \\ &\leq e^{-s/2} = T^{-(\gamma+50)}. \end{aligned} \tag{20}$$

Finally, we assemble these ingredients to bound

$$\mathbb{P}(\mathcal{E}_m[T] \cap \mathcal{L}_\ell[T]^c) = \sum_{t=1}^T \sum_{z \in \mathbf{B}_T} \mathbb{P}(Q_{z,t} \cap \mathcal{L}_\ell[T]^c).$$

Split each term in the sum on the right into three pieces, with $\zeta = \zeta_0(z, t)$, as indicated below and majorize using (18), (19), and (20), respectively. We obtain with $s = (2\gamma + 100)\log T$,

$$\begin{aligned} \mathbb{P}(Q_{z,t} \cap \mathcal{L}_\ell[T]^c) &\leq \mathbb{P}(Q_{z,t} \cap \{S_\zeta(t) > s\}) \\ &\quad + \mathbb{P}(Q_{z,t} \cap \{M_\zeta(t) < (b/10)m\} \cap \mathcal{L}_\ell[t]^c) \\ &\quad + \mathbb{P}(\{S_\zeta(t) \leq s\} \cap \{M_\zeta(t) \geq (b/10)m\}) < 3T^{-(\gamma+5)} \end{aligned}$$

Summing over $z \in \mathbf{B}_T$ and $t = 1, \dots, T$, since $T \geq 20$ we conclude that

$$\mathbb{P}(\mathcal{E}_m[T] \cap \mathcal{L}_\ell[T]^c) \leq (2\pi T^2) \cdot T \cdot 3T^{-(\gamma+5)} < T^{-(\gamma+1)}.$$

This completes the proof of Lemma 12. \square

2.7 Late points imply early points

Fix $\gamma \geq 1$, and let

$$C_4 = (\gamma + 3)(1500 + \frac{1}{c_0}) + \frac{2\log C_0}{c_0} + 2C_2 + 2.$$

where C_2 is the constant appearing in Lemma 8, and C_0 and c_0 appear in Lemma 3.

Lemma 13. (*Late points imply early points*) *If $\ell \geq C_4 \log T$ and $m \leq \ell^2/C_4 \log T$, then*

$$\mathbb{P}(\mathcal{L}_\ell[T] \cap \mathcal{E}_m[T]^c) \leq T^{-(\gamma+1)}.$$

In other words, with high probability, either there is no ℓ -late point by time T , or there is an $\ell^2/C_4 \log T$ -early point by time T .

Proof. Without loss of generality we take $m = \ell^2/C_4 \log T$. Note that since $\ell \geq C_4 \log T$ this ensures $m \geq \ell$. Let $\zeta \in \mathbb{Z}^2$ be such that $|\zeta| \leq \sqrt{T/\pi} - \ell$, and set

$$\rho = |\zeta|, \quad T_1 = \pi(\rho + \ell)^2.$$

Note that $T_1 \leq T$. Denote by $L[\zeta] = \{\zeta \notin A(T_1)\}$ the event that ζ is ℓ -late. Then

$$\mathcal{L}_\ell[T] = \bigcup_{|\zeta| \leq \sqrt{T/\pi} - \ell} L[\zeta].$$

We will use the martingale $M_\zeta(t)$ with pole at ζ . We first show that on the event $L[\zeta]$,

$$M_\zeta(T_1) \leq -\ell. \tag{21}$$

To prove this, note first that by the maximum principle, the smallest value of H_ζ on Ω_ζ is $1/2\rho$ attained on the boundary $\partial\Omega_\zeta$. This boundary value is achieved

everywhere on $\partial\Omega_\zeta$ except at the point ζ , at which H_ζ has a much larger value. On the event $L[\zeta]$, no particle reaches ζ by time T_1 , so the values of H_ζ at all boundary points reached at time T_1 are $1/2\rho$, which is smaller than the interior values. Therefore, $M_\zeta(T_1)$ is maximized if all the sites of Ω_ζ are occupied, i.e., $\Omega_\zeta \subset A_\zeta(T_1)$. By Lemma 8(a), we have $\Omega_\zeta \subset \mathbf{B}_{\rho+C_2}$, so

$$T_1 - \#\Omega_\zeta \geq \pi(\rho + \ell)^2 - \pi(\rho + C_2 + 1)^2.$$

Since $\ell \geq C_4 \log T$, the right side is at least $\pi\rho\ell$. Hence on the event $L[\zeta]$, we have

$$M_\zeta(T_1) \leq (\pi\rho\ell) \left(\frac{1}{2\rho} - H_\zeta(0) \right) + \sum_{z \in \Omega_\zeta} (H_\zeta(z) - H_\zeta(0)). \quad (22)$$

By Lemma 7(c), we have $\frac{1}{2\rho} - H_\zeta(0) = -\frac{1}{2\rho} + O(\frac{1}{\rho^2})$, which shows that the first term on the right side of (22) is less than $-3\ell/2$. By Lemma 8(c), the second term is at most $C_2 \log \rho \leq C_2 \log T \leq \ell/2$. This concludes the proof of (21).

Now suppose that $|\zeta| \geq \ell$. Recall that $m \geq \ell \geq C_4 \log T$ and $C_4 \geq 1$, so $e^m \geq T$. By Lemma 11,

$$\mathbb{E} e^{S_\zeta(T_1)} \mathbf{1}_{\mathcal{E}_m[T_1]^c} \leq T^{40} e^{700m} \leq e^{740m}. \quad (23)$$

Letting $s = 750m$, we have by Markov's inequality

$$\mathbb{P}(\{S_\zeta(T_1) > s\} \cap \mathcal{E}_m[T_1]^c) \leq e^{-s} \mathbb{E} e^{S_\zeta(T_1)} \mathbf{1}_{\mathcal{E}_m[T_1]^c} \leq e^{-10m}.$$

Now by (21),

$$\begin{aligned} \mathbb{P}(L[\zeta] \cap \mathcal{E}_m[T]^c) &\leq \mathbb{P}(\mathcal{E}_m[T]^c \cap \{S_\zeta(T_1) > s\}) + \mathbb{P}(\{S_\zeta(T_1) \leq s\} \cap L[\zeta]) \\ &\leq e^{-10m} + \mathbb{P}\{S_\zeta(T_1) \leq s, M_\zeta(T_1) \leq -\ell\}. \end{aligned}$$

We use Lemma 6 to bound the second term. Since $M_\zeta(t) = B_\zeta(S_\zeta(t))$, where B_ζ is a standard Brownian motion with $B_\zeta(0) = 0$, we have

$$\mathbb{P}\{S_\zeta(T_1) \leq s, M_\zeta(T_1) \leq -\ell\} \leq \mathbb{P}\left\{ \inf_{s' \in [0, s]} B_\zeta(s') \leq -\ell \right\} \leq e^{-\ell^2/2s}.$$

We conclude that for all $\zeta \in \mathbb{Z}^2$ such that $|\zeta| \leq \sqrt{T/\pi} - \ell$,

$$\mathbb{P}(\mathcal{E}_m[T]^c \cap L[\zeta]) \leq e^{-10m} + e^{-\ell^2/1500m}.$$

Since $m \geq \ell \geq (\gamma + 3) \log T$ and $\ell^2/1500m \geq (\gamma + 3) \log T$ by hypothesis, each term on the right side is at most $T^{-(\gamma+3)}$. Hence for $\ell \leq |\zeta| \leq \sqrt{T/\pi} - \ell$, we have

$$\mathbb{P}(\mathcal{E}_m[T]^c \cap L[\zeta]) \leq 2T^{-(\gamma+3)}. \quad (24)$$

To take care of the case $|\zeta| < \ell$, we use Lemma 3, which gives

$$\mathbb{P}\left(\bigcup_{|\zeta| < \ell} L[\zeta]\right) = \mathbb{P}(\mathcal{L}_\ell[4\pi\ell^2]) \leq C_0 e^{-c_0\ell} < T^{-(\gamma+3)}.$$

Summing (24) over $\zeta \in \mathbf{B}_{\sqrt{T/\pi}-\ell} - \mathbf{B}_\ell$ yields

$$\begin{aligned} \mathbb{P}(\mathcal{E}_m[T]^c \cap \mathcal{L}_\ell[T]) &\leq \mathbb{P}\left(\bigcup_{|\zeta| < \ell} L[\zeta]\right) + \sum_{\ell \leq |\zeta| \leq \sqrt{T/\pi}-\ell} \mathbb{P}(\mathcal{E}_m[T]^c \cap L[\zeta]) \\ &< T^{-(\gamma+3)} + T^{-(\gamma+2)}. \end{aligned}$$

This completes the proof of Lemma 13. \square

The proof of the main result from Lemmas 12 and 13 is a routine iteration. We use the lemmas in their contrapositive direction (no early point implies no late point and no late point implies no early point) for progressively smaller values of m and ℓ , starting from (4).

Fix $\gamma \geq 1$ and $T \geq 20$. Let $\ell_0 = \sqrt{T/100\pi}$. Then (4) implies

$$\mathbb{P}(\mathcal{L}_{\ell_0}[T]) \leq C_0 e^{-c_0\ell_0}.$$

If T is sufficiently large, say $T > T_0(\gamma)$, then the right side is less than $T^{-(\gamma+1)}$.

Now define

$$m_0 = (1000/b)\ell_0.$$

Lemma 12 implies that if $m_0 \geq C_3 \log T$, then

$$\mathbb{P}(\mathcal{E}_{m_0}[T] \cap \mathcal{L}_{\ell_0}[T]^c) \leq T^{-(\gamma+1)}.$$

so that

$$\mathbb{P}(\mathcal{E}_{m_0}[T]) < 2T^{-(\gamma+1)}.$$

For $k \geq 1$, define ℓ_k and m_k recursively by

$$\begin{aligned} \ell_k &= \sqrt{(C_4 \log T) m_{k-1}} \\ m_k &= (1000/b)\ell_k. \end{aligned}$$

By Lemma 13, if $\ell_1 \geq C_4 \log T$, then

$$\mathbb{P}(\mathcal{L}_{\ell_1}[T] \cap \mathcal{E}_{m_0}[T]^c) \leq T^{-(\gamma+1)}$$

hence

$$\mathbb{P}(\mathcal{L}_{\ell_1}[T]) < 3T^{-(\gamma+1)}.$$

By induction on k , we find that

$$\mathbb{P}(\mathcal{L}_{\ell_k}[T]) < (2k+1)T^{-(\gamma+1)}$$

and

$$\mathbb{P}(\mathcal{E}_{m_k}[T]) < (2k+2)T^{-(\gamma+1)}$$

provided $m_k \geq C_3 \log T$ and $\ell_k \geq C_4 \log T$.

The recursion gives

$$\ell_k = \alpha^{1-2^{-k}} \ell_0^{1/2^k}$$

and

$$m_k = \alpha^{1-2^{-k}} m_0^{1/2^k}.$$

where $\alpha = (1000/b)C_4 \log T$.

Let $C_5 = (1000/b) \max(C_3, C_4)$. Then

$$\ell_k, m_k \leq (C_5 \log T) T^{1/2^k}.$$

For $k > (\log \log T)/\log 2$, we have

$$T^{1/2^k} < 3.$$

Therefore, if we choose k largest so that $\ell_k \geq 3C_5 \log T$, then

$$k \leq (\log \log T)/\log 2.$$

Thus, in our final step we may use

$$\begin{aligned} \ell &:= \max(\ell_{k+1}, C_4 \log T) \leq 3C_5 \log T \\ m &:= \max(m_{k+1}, C_4 \log T) \leq 3C_5 \log T \end{aligned}$$

to obtain

$$\begin{aligned} \mathbb{P}(\mathcal{L}_\ell[T]) &< (2k+3)T^{-(\gamma+1)} < \frac{1}{2}T^{-\gamma} \\ \mathbb{P}(\mathcal{E}_m[T]) &< (2k+4)T^{-(\gamma+1)} < \frac{1}{2}T^{-\gamma} \end{aligned}$$

for all sufficiently large T .

Let $r = \sqrt{T/\pi}$ and $a = 10C_5 \log r$. Then $a > 3C_5 \log T \geq \ell, m$. By Lemma 4, we conclude that

$$\begin{aligned} \mathbb{P} \{ \mathbf{B}_{r-a \log r} \subset A(\pi r^2) \subset \mathbf{B}_{r+a \log r} \}^c &\leq \mathbb{P} \{ \mathbf{B}_{r-\ell} \not\subset A(T) \} + \mathbb{P} \{ A(T) \not\subset \mathbf{B}_{r+m} \} \\ &\leq \mathbb{P}(\mathcal{L}_\ell[T]) + \mathbb{P}(\mathcal{E}_m[T]) \\ &< T^{-\gamma} < r^{-\gamma}. \end{aligned}$$

This completes the proof of Theorem 1.

	$d = 1$	$d = 2$	$d \geq 3$
Order of $h(x)$ at typical $x \in \Lambda$	\sqrt{n}	$\sqrt{\log n}$	1
Order of $\max_{x \in \Lambda} h(x) $	\sqrt{n}	$\log n$	$\sqrt{\log n}$
Order of $ \Lambda ^{-1} \sum_{x \in \Lambda} \phi(x/n)h(x)$	\sqrt{n}	1	$n^{1-d/2}$

Table 1: Order of various kinds of fluctuations for the discrete Gaussian free field in \mathbb{Z}^d . In the bottom row, $\phi : [-1, 1]^d \rightarrow \mathbb{R}$ is a smooth test function.

3 Concluding Remarks

We mention here a few open questions and possible extensions of our results, along with a brief overview of our forthcoming sequels.

Gaussian free field

Let $\Lambda = [-n, n]^d \subset \mathbb{Z}^d$ and let h be an instance of the discrete Gaussian free field on Λ with zero boundary conditions. That is

1. write

$$E(f) = \sum_{x, y \in \Lambda, |x-y|=1} |f(x) - f(y)|^2,$$

2. let H be the $(2n - 1)^d$ -dimensional vector space of real-valued functions on Λ that vanish on the boundary $\partial\Lambda$,
3. and choose h from the measure $e^{-E(h)/2} dh$ (normalized to be a probability measure), where dh is Lebesgue measure on H .

Table 1 gives the orders of various observables associated to h . (See [BZ10] for recent results on the maximum of h in two dimensions, and [Sh07] for a general survey on the GFF.) Perhaps surprisingly, the orders in the last two rows do not change if we replace Λ with a fixed co-dimension one subset of Λ . This shows that weighted averages of h are remarkably well concentrated (and become more concentrated in higher dimensions). Appropriately normalized, they scale to Gaussian random variables with an asymptotic variance that depends ϕ .

Next, if we let

$$L(z) = \sqrt{\frac{\min\{n | z \in A(n)\}}{\pi}} - |z|$$

represent the “lateness” of the vertex z , then we expect L and its higher dimensional analogues to behave similarly to h in all dimensions. (This is trivial for dimension $d = 1$.) Our sequel paper, *Internal DLA and the Gaussian free field* [JLS11] addresses the bottom row (for certain test functions ϕ) and shows that L has a variant of the continuum Gaussian free field as a scaling limit.

Lower bounds on fluctuations

The present paper established the upper bound of $\log n$ for the fluctuations of L corresponding to the middle square of Table 1. Simulations [MM00, FL10] indicate that fluctuations of this order are really present, and proving this is a very natural open problem. One possible approach would be to use the methods of this paper to show that for small enough c , the *expected* number of aberrations of size $c \log n$ (or $c\sqrt{\log n}$ in dimension $d \geq 3$) in $A(n)$ grows like a *positive* power of n — and then using a second moment estimate to show that infinitely many such aberrations occur almost surely. The main new ingredient required would be to control the correlations between fluctuations in two different directions.

Higher dimensions

The $d \geq 3$ sequel to this paper [JLS10] uses variants of the arguments presented here. The function H_ζ is replaced by a discrete Green's function for a ball in \mathbb{Z}^d (with zero boundary conditions). The obvious analogue of H_ζ (constructed from differences of \mathbb{Z}^d Green's functions) does not have spherical level surfaces. This causes some extra complication in the estimates involved in Lemmas 7 and 8. It becomes necessary to estimate the discrete Green's function on the ball directly, instead of building on known estimates for the \mathbb{Z}^d Green's function. The argument requires a few other modifications as well (for example, taking ζ to be a point inside the ball, instead of on the boundary).

Fluctuations in a fixed direction

The bound of order $\log r$ in Theorem 1 measures the “fluctuations in the worst direction.” Another interesting quantity is the fluctuation in a specific direction, say along positive x -axis, as measured by the differences

$$\max\{x > 0 \mid (x, 0) \in A(\pi r^2)\} - r$$

and

$$r - \min\{x > 0 \mid (x, 0) \notin A(\pi r^2)\}.$$

Based on the analogy with the discrete Gaussian free field, we believe that these random variables have standard deviation of order $\sqrt{\log r}$ in two dimensions, and that their higher dimensional analogues have $O(1)$ standard deviation. These orders match those in the top line of Table 1.

In numerical experiments on a related quantity, Meakin and Deutch [MD86] found order $\sqrt{\log r}$ fluctuations in dimension 2 and $O(1)$ fluctuations in dimension 3.

We have not been able to adapt our bootstrapping argument to prove bounds of this type on fluctuations in a fixed direction. An m -early point where $m = \sqrt{\log r}$ contributes only $\sqrt{\log r}$ to the martingale M_ζ , which is swamped by the logarithmic error term in the discrete mean value property, Lemma 8(c).

Derandomized aggregation

James Propp around 2001 proposed a growth process called *rotor-router aggregation*, in which particles started at the origin in \mathbb{Z}^2 perform deterministic analogues of random walks until reaching an unoccupied site. Large-scale simulations show that the aggregates produced in this way are substantially smoother than IDLA, with even smaller fluctuations from circularity [FL10]. Indeed, part of Propp’s motivation was to study how much of the random fluctuation can be removed by passing to the deterministic analogue. The data indicate that the deviation from circularity in rotor-router aggregation may even be bounded independent of the size of the cluster; that is, $\mathbf{B}_{r-C} \subset A(\pi r^2) \subset \mathbf{B}_{r+C}$ for an absolute constant C . Yet the current best upper bounds (order $\log r$ on the inside and $r^{1/2} \log r$ on the outside, [LP09]) remain distant from this goal. In trying to adapt our methods to rotor-router aggregation, we can define a deterministic analogue of the martingale $M_\zeta(t)$ by summing the function H_ζ over the rotor-router aggregate $A(t)$. At this point an interesting question arises: what should take the place of the martingale representation theorem for this deterministic process?

External DLA

External DLA, defined by Witten and Sander [WS81], produces intricate branching structures instead of a smooth spherical surface. Typically it is defined with particles starting from infinity, but there is also a version in which particles start at the origin. In this version, the initial occupied set $A(0)$ is the set of sites outside some extremely large origin-centered ball. At time n we start a random walk at the origin and stop when it first reaches a site z_n adjacent to $A(n)$. Then we set $A(n+1) = A(n) \cup \{z_n\}$.

Both this paper and the sequel (about the Gaussian free field) use a similar technique: namely, they both use geometric arguments to control the quadratic variation of continuous martingales obtained by summing the values of discrete harmonic functions over sets of particle locations. One intriguing aspect of this work is that all of the martingales we utilize are still martingales for external DLA (grown from the origin). Indeed, Kesten’s upper bound on the growth rate of external DLA [Kes90] uses a closely related martingale.

Because the geometry of external DLA is much more intricate, however, it becomes more difficult to control the quadratic variation of the martingales. In our argument, the interlinked inner and outer estimates — indeed, even the notions of early and late points — rely on the fact that the cluster fluctuates around a predictable asymptotic shape (the disk). In external DLA, on the other hand, even the large-scale geometry is random.

A Appendix: Proof of Lemmas 2 and 3

A.1 Bounds on the probability of thin tentacles in all dimensions

This appendix is devoted to proving the following lemma, which strengthens Lemma 2. Since we anticipate using this result to bound the fluctuations of internal DLA in higher dimensions as well, we treat all dimensions $d \geq 2$.

Lemma A. *In dimension $d = 2$, there are positive absolute constants b , C_0 , and c such that for all real numbers $m > 0$ and all $z \in \mathbb{Z}^2$ with $0 \notin \mathbf{B}(z, m)$,*

$$\mathbb{P} \{z \in A(n), \#(A(n) \cap \mathbf{B}(z, m)) \leq bm^2\} \leq C_0 e^{-cm^2 / \log m}. \quad (25)$$

For each dimension $d \geq 3$, there are positive absolute constants b , C_0 , and c (depending only on d) such that for all real numbers $m > 0$ and all $z \in \mathbb{Z}^2$ with $0 \notin \mathbf{B}(z, m)$,

$$\mathbb{P} \{z \in A(n), \#(A(n) \cap \mathbf{B}(z, m)) \leq bm^d\} \leq C_0 e^{-cm^2}. \quad (26)$$

Note that even the thinnest possible tentacle, a path of length m , can be formed using only $m(m+1)/2$ random walk steps: for $j = 1, \dots, m$, the j -th particle to enter $\mathbf{B}(z, m)$ takes j steps along the path and occupies the next site on the path. This shows that the exponent m^2 on the right side of (26) is best possible.

Instead of working with the ball $\mathbf{B}(z, m)$ it will be slightly more convenient to use the box $\tilde{\mathbf{B}}(z, m) := z + [-m, m]^d \subset \mathbb{Z}^d$. Since the volume of a ball is (up to a d dependent constant factor) the same as that of a box of the same radius, Lemma A holds as stated if and only if it holds with $\mathbf{B}(z, m)$ replaced by $\tilde{\mathbf{B}}(z, m)$.

Fix $z \in \mathbb{Z}^d$, and for $0 \leq j \leq m$ let S_j be the set of points $x = (x_1, \dots, x_d) \in \mathbb{Z}^d$ such that

$$\max_{i=1}^d |x_i - z_i| = m - j.$$

The sets S_j partition $\tilde{\mathbf{B}}(z, m)$ into $m+1$ concentric shells, indexed from the outside inward. Since $\mathbb{Z}^2 \setminus S_j$ is disconnected, at least one site in shell S_j must join the cluster $A(n)$ before any site in S_{j-1} can join.

It is enough to show that (25) and (26) hold if we replace $A(n)$ with the set A , where $A = A(n_z) \cap \tilde{\mathbf{B}}(z, m)$, and n_z is defined to be the smallest n for which $z \in A(n)$. Now write $a_j := \#(A \cap S_j)$ for $j \in \{0, 1, \dots, m\}$. The cluster $A(n_z)$ intersects each shell S_j , so $a_j \geq 1$ for all $j = 0, 1, \dots, m$.

Given m , the number of possibilities for the sequence a_j is at most $((2m+1)^d)^m$, which is much smaller than the reciprocal of the right side of (25) and (26). It thus suffices enough to show that, for each fixed sequence $\mathbf{a} = (a_j)_{j=0}^m$ with $a_j \geq 1$ and $\sum_{j=0}^m a_j \leq bm^d$, (25) and (26) hold when the left side is replaced by the probability of the event

$$\mathcal{T}_{\mathbf{a}} := \bigcap_{j=0}^m \{\#(A \cap S_j) = a_j\}.$$

We emphasize that C_0 and c are still absolute constants; in particular, they do not depend on the sequence \mathbf{a} or on m . We will henceforth assume that such a sequence \mathbf{a} has been fixed, and bound the probability $\mathbb{P}(\mathcal{T}_{\mathbf{a}})$.

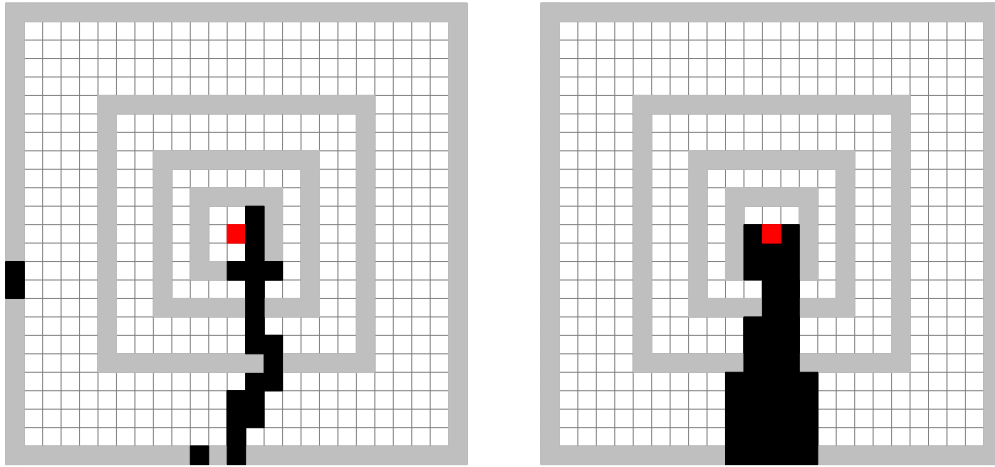


Figure 6: On the left: a possible value for A (black, with z in red at the center). Here $m = 12$ and the sequence a_0, a_1, \dots, a_{12} (giving number of black squares in the square shell S_j , indexed from outside inward) is given by 4, 1, 2, 2, 2, 1, 2, 1, 1, 1, 3, 3, 1. Taking $c' = 1/2$, the corresponding sequence b_0, \dots, b_{12} is given by 5, 5, 5, 5, 5, 3, 3, 3, 2, 2, 3, 3, 3, and a corresponding tower of squares is illustrated on the right. The sizes of the squares are $\beta_1 = 5$, $\beta_2 = 3$, $\beta_3 = 2$, $\beta_4 = 3$. The grey shells indicate the S_{α_i} , where $\alpha_0 = 0$, $\alpha_1 = 5$, $\alpha_2 = 8$, and $\alpha_3 = 10$. Since $c' = 1/2$, the number of black squares in the “annulus” formed by the union of the S_j crossing a given square on the right (other than the innermost square) is always at least twice the number black of squares in the corresponding annulus on the left.

Now we will construct a related sequence b_j indexed by $[0, m]$ with two nice properties:

1. b_j is roughly the kind of sequence a_j that one would see if A were a tower of cubes. More precisely, let β_1, \dots, β_k be positive integers summing to m , representing the side lengths of the cubes. For $i \in \{0, \dots, k\}$, let $\alpha_i = \beta_1 + \dots + \beta_i$. For $j \in \{0, \dots, m\}$, let γ_j be the unique integer i such that $\alpha_{i-1} \leq j < \alpha_i$. Then set

$$b_j = (\beta_{\gamma_j})^{d-1}.$$

Note that $\sum_{j=0}^m b_j = \sum_{i=1}^k \beta_i^d$.

2. There is a small constant c' such that for each $i = 1, \dots, k-1$ (but possibly not for $i = k$)

$$c'(\beta_i/2)^d \leq \sum_{j=\alpha_{i-1}}^{\alpha_i-1} a_j \leq c'(\beta_i)^d.$$

Once we have fixed c' , we can construct the sequence b_j from a_j in a deterministic way by induction: Given the cube side lengths $\beta_1, \dots, \beta_{i-1}$, take β_i to be the smallest possible such that the second condition holds.

Next we make explicit our choice of c' . For $x \in \mathbb{Z}^d$, write \mathbb{P}_x for the probability measure associated to simple random walk $(X(t))_{t=0}^\infty$ in \mathbb{Z}^d with $X(0) = x$. By the local central limit theorem [LL10, §2], if t is even, then for each site y of even parity in the ball $\mathbf{B}(x, \sqrt{t})$, we have

$$\mathbb{P}_x(X(t) = y) \geq c_0 t^{-d/2}$$

for a constant $c_0 > 0$ depending only on d . Let $c_1 = \omega_d 2^{-d-5} c_0$, where ω_d is the volume of the unit ball in \mathbb{R}^d .

By [LL10, Prop. 2.1.2], there is a constant $c_2 > 0$ depending only on d , so that

$$\mathbb{P}_0 \left\{ \max_{0 \leq s \leq t} |X(s)| \geq \sqrt{t}/c_2 \right\} < c_1. \quad (27)$$

Now take $c' = \omega_d 2^{-d-4} c_2^d$ in the above construction of the sequence b_j from a_j . This defines a sequence of shells S_{α_i} as in Figure 6. According to the next lemma, if A' is any subset of \mathbb{Z}^d satisfying $A' \cap S_j \leq a_j$ for $j = 0, \dots, m$, then simple random walk started from any point in the i -th shell S_{α_i} has at least a constant $c_1 > 0$ probability of exiting A' before reaching $S_{\alpha_{i-1}} \cup S_{\alpha_{i+1}}$.

Lemma 14. *Let A' be any set for which $\#(A' \cap S_j) \leq a_j$ for each $j = 0, \dots, m$. Fix $1 < i < k$ and $x \in S_{\alpha_i}$. For simple random walk in \mathbb{Z}^d started at x , let T be the first hitting time $S_{\alpha_{i-1}} \cup S_{\alpha_{i+1}}$, and let T' be the first hitting time of $(A')^c$. Then*

$$\mathbb{P}_x \{T > T'\} \geq c_1.$$

Proof. Let $h = \min(\beta_i, \beta_{i+1})$, and let t be the greatest even integer $\leq (c_2 h)^2$. Let $r = c_2 h$. By the local central limit theorem, for each site y of even parity in the ball $\mathbf{B}(x, r)$, we have

$$\mathbb{P}_x(X(t) = y) \geq c_0 r^{-d}.$$

Suppose first that $\beta_{i+1} \leq \beta_i$. Let B be the set of points of even parity in the intersection

$$\mathbf{B}(x, r) \cap \bigcup_{j=\alpha_i}^{\alpha_{i+1}-1} S_j.$$

Note that $\#B \geq \#\mathbf{B}(x, r)/2^{d+2} \geq (\omega_d 2^{-d-3})r^d$. Hence by our choice of c' ,

$$\#(A' \cap B) \leq \sum_{j=\alpha_i}^{\alpha_{i+1}-1} a_j \leq c' \beta_{i+1}^d = c' h^d < \#B/2.$$

Hence

$$\mathbb{P}_x\{X(t) \notin A'\} \geq \sum_{y \in B \cap (A')^c} \mathbb{P}_x(X(t) = y) \geq (\#B/2) \frac{c_0}{r^d} > 2c_1.$$

By (27), we have

$$\mathbb{P}_x(T \leq t) < c_1.$$

We conclude that

$$\begin{aligned} \mathbb{P}\{T > T'\} &\geq \mathbb{P}_x\{T > t, X(t) \notin A'\} \\ &\geq 1 - \mathbb{P}_x\{T \leq t\} - \mathbb{P}\{X(t) \in A'\} \\ &> 1 - c_1 - (1 - 2c_1). \end{aligned}$$

The proof in the case $\beta_{i+1} > \beta_i$ is identical, using the shells $\bigcup_{j=\alpha_{i-1}+1}^{\alpha_i} S_j$ in the definition of the set B . \square

During the formation of the IDLA cluster A , each time a particle hits a shell S_{α_i} ($1 < i < k$) for the first time, we begin a “trial” which ends either when the particle stops walking (due to exiting the cluster and occupying a previously unoccupied site) or when it next hits $S_{\alpha_{i-1}} \cup S_{\alpha_{i+1}}$. In the former case, we call the trial a “failure.” Otherwise, we call it a “success.”

Let F_s be the event that the s -th trial fails, and let \mathcal{G}_s be the σ -field generated by all random walks generating the IDLA cluster up until the start of the s -th trial. Let A_s be the occupied cluster at the time when the s -th trial starts, and let $E_s = \bigcap_{j=0}^m \{\#(A_s \cap S_j) \leq a_j\}$. By Lemma 14, we have

$$\mathbb{P}(F_s | \mathcal{G}_s) \mathbf{1}_{E_s} \geq c_1 \mathbf{1}_{E_s}.$$

In other words, on the event E_s , once the s -th trial begins (and we condition on the past of the process — in particular on all previous trials) the trial fails with probability at least probability c_1 .

Since each failure adds a new site to A , there are at most $\#A$ failures. Next we estimate the total number of trials. Each particle that makes it to S_{α_i} has to first pass through all the $S_{\alpha_{i'}}$ for $i' < i$. Thus the process of successful growing a set A , such that $\#(A \cap S_j) = a_j$ for all j , necessarily includes at least $T = \sum a_j \gamma_j$ trials. Hence on the event \mathcal{T}_a , there are at least T trials, at most $\sum a_j$ of which fail; and conditional on all past trials, each trial fails with probability at least c_1 .

Before we put an explicit bound on T , we'd like to find a way to reduce the problem to the case that

$$\#A = \sum a_j < c_1 T/2 \tag{28}$$

(i.e., failure is required to occur at less than half the expected rate). In other words, we'd like to reduce the case that $\sum a_j \gamma_j$ is larger (by a fixed constant factor $2/c_1$) than $\sum a_j$. This will certainly be the case, for example, if more than a fraction $1/4$ of the sum $\sum a_j$ comes from values of j for which $\gamma_j > 8/c_1$.

Let R be a fixed integer greater than $8/c_1$. Let I be the smallest positive integer such that

$$\sum_{j=\alpha_{IR}}^m a_j < 4 \sum_{j=\alpha_{IR+1}}^m a_j.$$

Then consider the smaller box we get by replacing m by $m - \alpha_{IR}$. We can take b small enough so that β_i^d is at most an arbitrarily small constant times m^d , and hence β_i is at most some arbitrarily small constant times m . Indeed, since R is just a constant, we can take b small enough so that $\sum_{i=1}^R \beta_i$ is at most an arbitrarily small constant times m . Since the sequence

$$\sum_{j=\alpha_{iR}}^{\alpha_{(i+1)R}-1} a_j$$

decreases by a factor of at least 4 each time i increases by 1 (up until $i = I$), we also obtain that the sequence

$$\beta_{iR} + \cdots + \beta_{(i+1)R-1}$$

is also bounded above by a sequence that decreases exponentially in i , and hence (by making b small enough) we can arrange so that $\sum_{i=1}^{IR} \beta_i$ is less than $m/2$. Then if we replace m with $m - \alpha_{IR}$ (and replace A by its intersection with the corresponding smaller box), we reduce the problem to one for which (28) holds and m has decreased by at most a factor of 2.

Now assuming (28), the usual large deviations bound (Cramer's theorem) implies that the probability of seeing T trials with only $c_1 T/2$ or fewer failures is exponentially unlikely in T . Note that $T = \sum_{j=0}^m a_j \gamma_j$ is up to a multiplicative constant (between $c'/2^d$ and c') the same as $\sum_{j=0}^m b_j \gamma_j = \sum_{i=1}^k i \beta_i^d$. Hence the proof of Lemma A.1 is completed by the following lower bound on $\sum i \beta_i^d$.

Lemma 15. *Fix $d \geq 2$ and $m \geq 1$, and consider an integer k and sequence of real numbers $\beta_1, \beta_2, \dots, \beta_k$ for which the following hold:*

1. $\beta_i \geq 1$ for all $i \in \{1, \dots, k\}$.
2. $\sum_{i=1}^k \beta_k = m + 1$.
3. $E(\beta) := \sum_{i=1}^k i \beta_i^d$ is minimal among sequences satisfying the two conditions above.

There is a constant $c(d) > 0$, depending only on d , such that

$$\begin{cases} E(\beta) \geq c(d)m^2/\log m, & d = 2; \\ E(\beta) \geq c(d)m^2, & d \geq 3. \end{cases} \quad (29)$$

Proof. The first two conditions imply that $k \leq m + 1$. By Holder's inequality,

$$m + 1 = \sum_{i=1}^k \beta_i \leq \left(\sum_{i=1}^k i \beta_i^d \right)^{1/d} \left(\sum_{i=1}^k i^{-d'/d} \right)^{1/d'}$$

where d' is such that $1/d + 1/d' = 1$. Hence

$$E(\beta) \geq (m + 1)^d \left(\sum_{i=1}^k i^{-d'/d} \right)^{-d/d'}.$$

If $d > 2$, then $d'/d < 1$ and the right side is at least a constant times $m^d k^{(1-d'/d)(-d/d')} = m^d k^{2-d} \geq m^2$. In the case $d = 2$, we have $d'/d = 1$ and the right side is at least a constant times $m^2(\log k)^{-1} \geq m^2(\log m)^{-1}$. \square

A.2 Bound on the probability of very late points

In this section we prove Lemma 3. Lawler, Bramson and Griffeath give an exponential bound on the probability of ℓ -late points in the case that ℓ is comparable to the radius of the cluster. Specifically, they show (see Lemma 6 of [LBG92] and the displayed equation immediately following it) that for any $\epsilon > 0$ there is a constant c_2 depending on ϵ such that for all $z \in \mathbf{B}_{(1-\epsilon)n}$

$$\mathbb{P}(z \notin A((1 + \epsilon)\pi n^2)) < 2e^{-c_2 n}. \quad (30)$$

Fix $r > 0$. Taking $\epsilon' = \epsilon/2$ and $n = (1 - \epsilon')r$ in (30), we have $(1 - \epsilon)r < (1 - \epsilon')n$ and $\pi r^2 > (1 + \epsilon')\pi n^2$, so for all $z \in \mathbf{B}_{(1-\epsilon)r}$,

$$\mathbb{P}(z \notin A(\pi r^2)) < 2e^{-c_2 n} \leq 2e^{-cr} \quad (31)$$

where c is a constant depending on ϵ .

By definition,

$$\mathcal{L}_\ell[100\pi\ell^2] = \bigcup_{z \in \mathbf{B}_{9\ell}} \{z \notin A(\pi(|z| + \ell)^2)\}.$$

Taking $r = |z| + \ell$ in (31) with $\epsilon = \frac{1}{10}$, we have for all $z \in \mathbf{B}_{9\ell}$

$$\mathbb{P}\{z \notin A(\pi r^2)\} < 2e^{-cr} \leq 2e^{-c\ell}.$$

Summing over $z \in \mathbf{B}_{9\ell}$ we obtain

$$\mathbb{P}(\mathcal{L}_\ell[100\pi\ell^2]) \leq 200\pi\ell^2 e^{-c\ell}.$$

The right side is at most $C_0 e^{-c_0 \ell}$ for suitable constants $C_0, c_0 > 0$, which proves (4).

References

- [AG10a] A. Asselah and A. Gaudillièrè, A note on the fluctuations for internal diffusion limited aggregation. [arXiv:1004.4665](#)
- [AG10b] A. Asselah and A. Gaudillièrè, From logarithmic to subdiffusive polynomial fluctuations for internal DLA and related growth models. [arXiv:1009.2838](#)
- [BZ10] M. Bramson and O. Zeitouni, Tightness of the recentered maximum of the two-dimensional discrete Gaussian free field. [arXiv:1009.3443](#)
- [DF91] P. Diaconis and W. Fulton, A growth model, a game, an algebra, Lagrange inversion, and characteristic classes, *Rend. Sem. Mat. Univ. Pol. Torino* **49**(1): 95–119, 1991.
- [FL10] T. Friedrich and L. Levine, Fast simulation of large-scale growth models. [arXiv:1006.1003](#)
- [FU96] Y. Fukai and K. Uchiyama, Potential kernel for two-dimensional random walk. *Ann. Probab.* **24**(4):1979–1992, 1996.
- [GQ00] J. Gravner and J. Quastel, Internal DLA and the Stefan problem, *Ann. Probab.* **28**(4):1528–1562, 2000.
- [JLS10] D. Jerison, L. Levine and S. Sheffield, Internal DLA in higher dimensions. [arXiv:1012.3453](#)
- [JLS11] D. Jerison, L. Levine and S. Sheffield, Internal DLA and the Gaussian free field. [arXiv:1101.0596](#)
- [KL10] W. Kager and L. Levine, Diamond aggregation, *Math. Proc. Cambridge Phil. Soc.* **149**: 351–372, 2010. [arXiv:0905.1361](#)
- [Kes90] H. Kesten, Upper bounds for the growth rate of DLA, *Physica A* **168**, 529–535, 1990.
- [KS04] G. Kozma and E. Schreiber, An asymptotic expansion for the discrete harmonic potential, *Electron. J. Probab.* **9**(1):1–17, 2004. [arXiv:math/0212156](#)
- [LBG92] G. F. Lawler, M. Bramson and D. Griffeath, Internal diffusion limited aggregation, *Ann. Probab.* **20**(4):, 2117–2140, 1992.
- [Law95] G. F. Lawler, Subdiffusive fluctuations for internal diffusion limited aggregation, *Ann. Probab.* **23**(1):71–86, 1995.
- [LL10] G. F. Lawler and V. Limic, *Random Walk: A Modern Introduction*, Cambridge Studies in Advanced Mathematics **123**, 2010.

- [LP09] L. Levine and Y. Peres, Strong spherical asymptotics for rotor-router aggregation and the divisible sandpile, *Potential Anal.* **30**:1–27, 2009. [arXiv:0704.0688](#)
- [LP10] L. Levine and Y. Peres, Scaling limits for internal aggregation models with multiple sources, *J. d'Analyse Math.* **111**: 151–219, 2010. [arXiv:0712.3378](#)
- [MD86] P. Meakin and J. M. Deutch, The formation of surfaces by diffusion-limited annihilation, *J. Chem. Phys.* **85**:2320, 1986.
- [MM00] C. Moore and J. Machta, Internal diffusion-limited aggregation: parallel algorithms and complexity, *J. Stat. Phys.* **99**(3-4): 661–690, 2000. [arXiv:cond-mat/9909233](#)
- [MP10] P. Mörters and Y. Peres, *Brownian Motion*, Cambridge University Press, 2010.
- [RY05] D. Revuz and M. Yor, *Continuous Martingales and Brownian Motion*, Springer, 2005.
- [Sh07] S. Sheffield., Gaussian free fields for mathematicians. *Probab. Theory Related Fields*, 139(3-4):521–541, 2007. [arXiv:math/0312099](#)
- [Wil96] D. B. Wilson, Generating random spanning trees more quickly than the cover time, In *28th Annual ACM Symposium on the Theory of Computing (STOC '96)*, 296–303, 1996.
- [WS81] T. A. Witten and L. M. Sander, Diffusion-limited aggregation, a kinetic critical phenomenon, *Phys. Rev. Lett.* **47**(19):1400–1403, 1981.