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Flux qubit noise spectroscopy using Rabi oscillations under strong driving conditions

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We infer the high-frequency flux noise spectrum in a superconducting flux qubit by studying the decay of Rabi oscillations under strong driving conditions. The large anharmonicity of the qubit and its strong inductive coupling to a microwave line enabled high-amplitude driving without causing significant additional decoherence. Rabi frequencies up to 1.7 GHz were achieved, approaching the qubit's level splitting of 4.8 GHz, a regime where the rotating-wave approximation breaks down as a model for the driven dynamics. The spectral density of flux noise observed in the wide frequency range decreases with increasing frequency up to 300 MHz, where the spectral density is not very far from the extrapolation of the $1/f$ spectrum obtained from the free-induction-decay measurements. We discuss a possible origin of the flux noise due to surface electron spins.

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Flux noise has been investigated for decades to improve stability and sensitivity in superconducting flux-based devices. Its power spectral density (PSD) has been studied in superconducting quantum interference devices (SQUIDs) [1,2] and in various types of superconducting qubits, such as charge [3], flux [4–10], and phase qubits [11–14]. The spectra typically follow $1/f$ frequency dependence with a spectral density of 1–10 $\mu\Phi_0/\sqrt{\text{Hz}}$ at 1 Hz, where $\Phi_0 = h/2e$ is the superconducting flux quantum. The accessible frequency range of the PSD was limited to approximately 10 MHz in spin-echo measurements [4,5,9,15] and was extended to a few tens of megahertz using Carr-Purcell-Meiboom-Gill pulse sequences [9]. Recently, spin-locking measurements provided the PSD up to approximately 100 MHz [16], and a study of qubit relaxation due to dressed dephasing in a driven resonator revealed the PSD at approximately 1 GHz [17]. The spectrum in a higher-frequency range would give further information for better understanding of the microscopic origin of the flux fluctuations.

Decay of Rabi oscillations has also been used as a tool to characterize the decoherence in superconducting qubits. PSDs of fluctuating parameters, such as charge, flux, or coupling strength to an external two-level system, at the Rabi frequency Ω_R can be detected [3,9,18,19]. The Rabi frequency is proportional to the amplitude of the driving field for weak to moderate driving at the qubit transition frequency, and Rabi frequencies in the gigahertz range have been achieved under a strong driving field [20–22]. However, the decay was not systematically studied because of the presence of extrinsic decoherence mechanisms under the strong driving conditions.

To induce fast Rabi oscillations without significant extra decoherence, we choose a flux qubit having strong inductive coupling to a microwave line and large anharmonicity,

$|(\omega_{12} - \omega_{01})/\omega_{01}|$, to avoid unwanted excitations to the higher energy levels, where ω_{ij} is the transition frequency between the $|i\rangle$ and $|j\rangle$ states. We measured Rabi oscillations in a wide range of $\Omega_R/2\pi$ from 2.7 MHz to 1.7 GHz and evaluated the PSD of flux fluctuations at each Ω_R. The observed PSD decreases up to 300 MHz, where the spectral density is approximately 10^{-20} $(\Phi_0)^2\text{rad}^{-1}$ s. Above 300 MHz, the PSD scatters and slightly increases. We discuss a possible origin of the flux fluctuations due to surface electron spins.

The Hamiltonian of a flux qubit with a flux drive and in the presence of fluctuations can be written in the persistent current basis as

$$\mathcal{H}_{\text{pc}} = -\frac{\hbar}{2}[(\Delta\sigma_x + \varepsilon\sigma_z) + \varepsilon_{\text{mw}}\cos(\omega_{\text{mw}}t)\,\sigma_z \\ + \delta\Delta(t)\sigma_x + \delta\varepsilon(t)\sigma_z], \quad (1)$$

where σ_x and σ_z are Pauli matrices, Δ is the tunnel splitting between two states with opposite persistent current direction along the qubit loop I_p, and $\hbar\varepsilon = 2I_p\Phi_0 n_\phi$ is the energy bias between the two states. Here the flux bias through the loop Φ_{ex} is normalized by Φ_0 as $n_\phi = \Phi_{\text{ex}}/\Phi_0 - 0.5$. The first term on the right-hand side of Eq. (1) represents the flux qubit with a static flux bias. The transition frequency can be written as $\omega_{01} = \sqrt{\Delta^2 + \varepsilon^2}$. We find $\omega_{01} = \Delta$ and $\partial\omega_{01}/\partial n_\phi = 0$ at $n_\phi = 0$; this is the optimal flux bias condition where dephasing due to fluctuations of n_ϕ is minimal. The second term is an ac drive at frequency ω_{mw} with the amplitude ε_{mw} to induce Rabi oscillations. The third and fourth terms represent fluctuations of Δ and ε, respectively. In the present sample, ε is tunable via n_ϕ while Δ is fixed.

There exist a few dominant contributors to the decay of Rabi oscillations: the quasistatic noise; the noise at ω_{01}, which causes the qubit energy relaxation; and the noise at Ω_R [3,9,23]. The resulting decay envelope $A_{\text{env}}(t)$ is described as

$$A_{\text{env}}(t) = A_{\text{st}}(t)\exp\left(-\Gamma_R^{\text{exp}}t\right), \quad (2)$$

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where $A_{st}(t)$ is the contribution from the quasistatic noise, which is usually nonexponential, and Γ_R^{exp} is the decay rate of the exponentially decaying term. As we are interested in the flux fluctuations at the Rabi frequency, contributions from other sources are to be separated out.

The quasistatic noise, which results in $A_{st}(t)$ in Eq. (2), is attributed to the fluctuations of the time-averaged values of $\delta\varepsilon(t)$ and $\delta\Delta(t)$ during a single decoherence measurement trial. The variances of the quasistatic flux noise, $\sigma_{\delta\varepsilon}^2$, and the Δ noise, $\sigma_{\delta\Delta}^2$, are determined from the result of free-induction-decay (FID) measurements [24], where we find $\sigma_{\delta\varepsilon}^2 \gg \sigma_{\delta\Delta}^2$. To evaluate the decay envelope $A_{st}(t)$ due to the quasistatic flux noise, we numerically calculate the time evolution of the density matrix of the qubit $\rho_{qubit}(t)$ under $\mathcal{H}_{pc}$.

The exponentially decaying component of the envelope is caused by the fluctuations at ω_{01} and Ω_R, and the rate is written as [3]

$$\Gamma_R^{exp} = \frac{(3-\cos^2\zeta)\Gamma_1}{4} + \Gamma_{\Omega_R}, \qquad (3)$$

where

$$\Gamma_1 = \frac{2\pi}{\hbar^2} \sum_\lambda S_\lambda(\omega_{01}) \left| \left\langle 1 \left| \frac{\partial\mathcal{H}_{pc}}{\partial\lambda} \right| 0 \right\rangle \right|^2 \qquad (4)$$

and

$$\Gamma_{\Omega_R} = \sin^2\zeta \, \frac{\pi}{2\hbar^2} [(2I_p\Phi_0)^2 S_{n_\phi}(\Omega_R)\cos^2\eta + \hbar^2 S_\Delta(\Omega_R)\sin^2\eta]. \qquad (5)$$

Here $\zeta = \arccos(\delta\omega_{mw}/\Omega_R)$, $\eta = \arctan(\Delta/\varepsilon)$, $\delta\omega_{mw} \equiv \omega_{mw} - \omega_{01}$ is the detuning frequency, and $S_\lambda(\omega) = \frac{1}{2\pi}\int_{-\infty}^{\infty} d\tau \langle \delta\lambda(t)\lambda(t+\tau) \rangle \exp(-i\omega\tau)$ denotes the PSD of a fluctuating parameter $\delta\lambda$ such as flux, charge, and critical current of the Josephson junctions. Γ_1 is the rate of the energy relaxation induced by the fluctuations at ω_{01} and can be independently measured as the decay rate of the qubit population after a π-pulse excitation [24]. Strictly speaking, the first term in Eq. (3) is written as $[\Gamma_1' + 2\Gamma_1 + (\Gamma_1' - 2\Gamma_1)\cos^2\zeta]/4$ [3], where Γ_1' is the average of the energy relaxation rates at $\omega_{01} \pm \Omega_R$ and is usually close to Γ_1. Γ_{Ω_R} is the decay rate due to fluctuations at Ω_R. Therefore, by analyzing experimental results using Eqs. (2)–(5), $S_{n_\phi}(\Omega_R)$ and $S_\Delta(\Omega_R)$ can be evaluated from the Rabi oscillation measurements at $\varepsilon \ll \Delta$ and $\varepsilon \approx \Delta$.

We need to pay attention to the drive-induced frequency shift of the qubit in the Rabi oscillation measurements under strong driving. We resort to numerical calculations to study the shift of the resonant frequency $\delta\omega$ as a function of ε_{mw}. At each ε_{mw}, Ω_R is calculated as a function of ω_{mw} and fitted with an analytic form,

$$\Omega_R = \sqrt{\left(\frac{\Delta}{2}\frac{\varepsilon_{mw}}{\omega_{01}}\right)^2 + [\omega_{mw} - (\omega_{01} + \delta\omega)]^2}. \qquad (6)$$

The first term, $(\frac{\Delta}{2}\frac{\varepsilon_{mw}}{\omega_{01}})^2 \equiv \Omega_{R0}^2$, is the square of the Rabi frequency at the new resonance condition ($\omega_{mw} = \omega_{01} + \delta\omega$), and the second term is the square of the detuning from the resonance. For simplicity, we use the linear approximation, $\Omega_{R0} \propto \varepsilon_{mw}/\omega_{01}$. This approximation is numerically validated
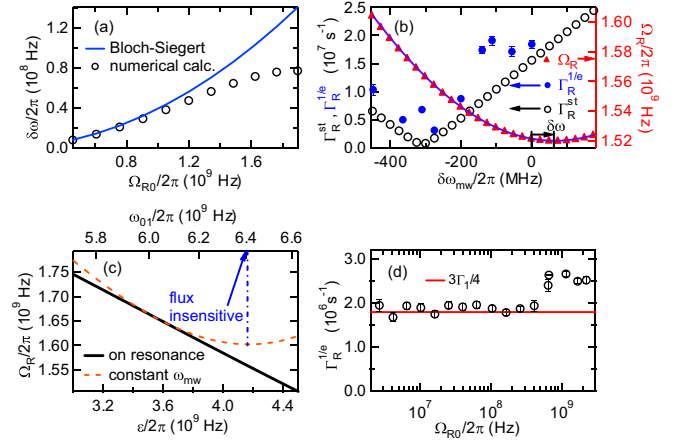

FIG. 1. (Color online) (a) Numerically calculated shift of the resonant frequency $\delta\omega$ (black open circles) and the Bloch–Siegert shift $\delta\omega_{BS}$ (blue line). (b) Numerically calculated decay rate Γ_R^{st} (black open circles) and Rabi frequency Ω_R (red solid triangles) as functions of the detuning $\delta\omega_{mw}$ from ω_{01}. The purple solid line is a fit based on Eq. (6). The measured $1/e$ decay rates $\Gamma_R^{1/e}$ at $\varepsilon/2\pi = 4.16$ GHz for the range of Rabi frequencies $\Omega_R/2\pi$ between 1.5 and 1.6 GHz (blue solid circles) are also plotted. (c) Calculated Rabi frequency Ω_R based on Eq. (6), as a function of ε for the cases (i) $\omega_{mw} = \omega_{01} + \delta\omega$ (black solid line) and (ii) $\omega_{mw}/2\pi = 6.1$ GHz (red dashed line). The upper axis indicates ω_{01}, corresponding to ε in the bottom axis. (d) The measured $1/e$ decay rate of the Rabi oscillations $\Gamma_R^{1/e}$ at $\varepsilon = 0$ and as a function of Ω_{R0}. The red solid line indicates $\frac{3}{4}\Gamma_1$ obtained independently.

within the range of parameters Ω_{R0} and ε in most cases in the present study.

In Fig. 1(a), $\delta\omega$ as a function of Ω_{R0} is plotted together with the well-known Bloch-Siegert shift [25,26], $\delta\omega_{BS} = \frac{1}{4}\frac{\Omega_{R0}^2}{\omega_{01}}$, obtained from the second-order perturbation theory. Fixed parameters for the calculation are $\Delta/2\pi = 4.869$ and $\varepsilon/2\pi = 4.154$ GHz ($\omega_{01}/2\pi = 6.400$ GHz). We find that $\delta\omega_{BS}$ overestimates $\delta\omega$ when $\Omega_{R0}/2\pi \gtrsim 800$ MHz. The deviation from the Bloch-Siegert shift is due to the component of the ac flux drive that is parallel to the qubit's energy eigenbasis; this component is not averaged out when Ω_R is comparable to ω_{mw}.

We next calculate the decay of Rabi oscillations due to quasistatic flux noise [24] and examine its dependence on $\delta\omega_{mw}$. In Fig. 1(b), Ω_R and the decay rate Γ_R^{st}, defined as the inverse of the $1/e$ decay time, are plotted as functions of $\delta\omega_{mw}$. Fixed parameters, $\varepsilon_{mw}/2\pi = 4.100$ GHz and $\sigma_{\delta\varepsilon}/2\pi = 27.8$ MHz, are chosen. Interestingly, neither the minimum of Ω_R nor that of Γ_R^{st} is located at $\delta\omega_{mw} = 0$, but at $\delta\omega_{mw}/2\pi = 66.5$ MHz for Ω_R and at $\delta\omega_{mw}/2\pi = -311$ MHz for Γ_R^{st}. For the ac flux drive, the frequency offset that minimizes the Rabi frequency is a consequence of the amplitude-dependent frequency shift $\delta\omega$, as can be observed in Eq. (6). Since the fluctuations of Ω_R causes the decay of Rabi oscillations [27], the minimum of Γ_R^{st} is understood by considering the flux sensitivity of Ω_R, which is expressed as

$$\frac{\partial\Omega_R}{\partial\varepsilon} = \frac{-\varepsilon}{\omega_{01}\Omega_R}\left[\omega_{mw} - \left(\omega_{01} + \delta\omega - \frac{\Omega_{R0}^2}{\omega_{01}}\right)\right]. \qquad (7)$$

The condition, $\partial\Omega_R/\partial\varepsilon = 0$, is satisfied when $\varepsilon = 0$ or $\delta\omega_{mw} = \delta\omega - \Omega_{R0}^2/\omega_{01}$. For $\Omega_{R0}/2\pi = 1.52$ GHz and $\omega_{01}/2\pi = 6.400$ GHz, the latter condition is calculated to be $\delta\omega_{mw}/2\pi = -295$ MHz, slightly different from the minimum of Γ_R^{st} seen in Fig. 1(b). The difference is due to the deviation from the linear approximation in Eq. (6), $\Omega_{R0} \propto \varepsilon_{mw}/\omega_{01}$. Figure 1(c) shows the calculation of Ω_R as a function of ε, based on Eq. (6). The Rabi frequency Ω_{R0} at the shifted resonance decreases as ε increases, while Ω_R, for a fixed microwave frequency of $\omega_{mw}/2\pi = 6.1$ GHz, has a minimum of approximately $\omega_{01}/2\pi = 6.4$ GHz. Here in the first order, Ω_R is insensitive to the fluctuation of ε.

The experiments were performed with a sample fabricated by electron-beam lithography and shadow evaporation of Al films, with a thickness of 13 nm for the first layer and 30 nm for the second, on an undoped Si substrate covered with a 300-nm-thick SiO_2 layer [24]. The qubit is a superconducting loop intersected by four Josephson junctions, among which one is smaller than the others by a factor of 0.5, nominally. The loop area is larger than that of flux qubits that we previously used [4], yielding a large mutual inductance between the qubit and the microwave line (1.2 pH) and facilitating strong driving.

We first measured ω_{01} as a function of ε and determined the qubit parameters. A 1-μs microwave pulse is applied to the qubit, followed by a bias current pulse of the readout SQUID (readout pulse). When the microwave frequency hits a transition of the qubit, the excitation is detected as a change in the SQUID switching probability P_{sw}. The flux qubit under study was cooled twice in between, up to room temperature with a thermal cycling. We noticed that Δ decreased by 1% after the thermal cycling: $\Delta/2\pi = 4.87$ GHz during the first cooldown and $\Delta/2\pi = 4.82$ GHz during the second. $I_p = 235$ nA was the same for both cooldowns. Unless explicitly mentioned below, we present the data from the first cooldown.

In the Rabi oscillation measurements, a microwave pulse is applied to the qubit followed by a readout pulse, and P_{sw} as a function of the microwave pulse length is measured. First, we measure the Rabi oscillation decay at $\varepsilon = 0$, where the quasistatic noise contribution is negligible. Figure 1(d) shows the measured $1/e$ decay rate of the Rabi oscillations $\Gamma_R^{1/e}$ as a function of Ω_{R0}. For $\Omega_{R0}/2\pi$ up to 400 MHz, $\Gamma_R^{1/e}$ is approximately $3\Gamma_1/4$, limited by the energy relaxation, and $S_\Delta(\Omega_{R0})$ is negligible. For $\Omega_{R0}/2\pi$ from 600 MHz to 2.2 GHz, $\Gamma_R^{1/e} > 3\Gamma_1/4$. A possible origin of this additional decoherence is fluctuations of ε_{mw}, $\delta\varepsilon_{mw}$: Ω_{R0} is first order sensitive to $\delta\varepsilon_{mw}$, which is reported to be proportional to ε_{mw} itself [28]. Next, the decay for the case $\varepsilon \approx \Delta$ is studied. To observe the contribution from quasistatic flux noise, the Rabi oscillation decay as a function of ω_{mw} is measured, where the contribution from the other sources is expected to be almost constant. Figure 1(b) shows $\Gamma_R^{1/e}$ at $\varepsilon/2\pi = 4.16$ GHz as a function of $\delta\omega_{mw}$ while keeping $\Omega_R/2\pi$ between 1.5 and 1.6 GHz. Besides the offset and scatter, the trend of $\Gamma_R^{1/e}$ agrees with that of the simulated Γ_R^{st}. This result indicates that numerical calculation properly evaluates $\delta\omega_{mw}$ minimizing Γ_R^{st}. Finally, the decay for the case $\varepsilon \approx \Delta$ as a function of ε_{mw}, covering a wide range of Ω_R, is measured (Fig. 2). At each Ω_R, $\delta\omega_{mw}$ is chosen to minimize dephasing due to quasistatic flux noise, which is numerically calculated as $A_{st}(t)$ in Eq. (2). After dividing
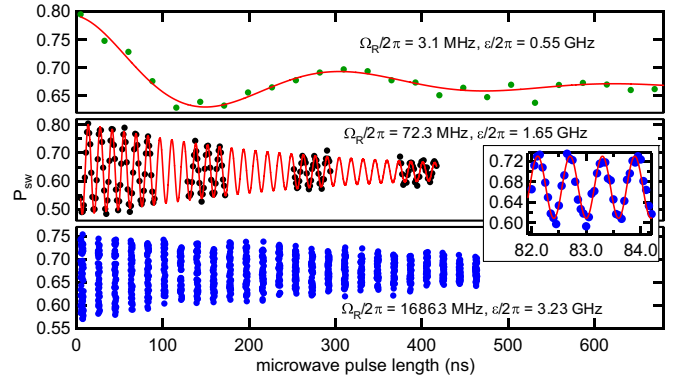

FIG. 2. (Color online) Rabi oscillation curves with different Rabi frequencies Ω_R measured at different static flux bias ε. At each Ω_R, $\delta\omega_{mw}$ is chosen to minimize dephasing due to quasistatic flux noise. The red lines are the fitting curves [24]. In the measurements shown in the middle and bottom panels, only parts of the oscillations are monitored, so that we can save measurement time while the envelopes of Rabi oscillations are captured. The inset is a magnification of the data in the bottom panel together with the fitting curve.

$A_{env}(t)$ by $A_{st}(t)$ in Eq. (2) and subtracting the decay rates obtained by Γ_1 and $S_\Delta(\Omega_R)$ from Γ_R^{exp} using Eqs. (3)–(5), $S_{n_\phi}(\Omega_R)$ is extracted [24]. Parameters in calculations and measurements are summarized in Table I.

The PSD of flux fluctuations $S_{n_\phi}(\omega)$, evaluated from the Rabi oscillation measurements in the first and second cooldowns, and PSDs from the spin-echo and energy relaxation measurements [24] in the second cooldown are plotted in Fig. 3. The $1/f$ spectrum extrapolated from the FID measurements in the second cooldown, $S_{n_\phi}(\omega) = (3.2 \times 10^{-6})^2/\omega$ [24], is also plotted. Several points are worth mentioning: (i) $S_{n_\phi}(\omega)$ from the Rabi oscillation measurements in the first and second cooldowns is consistent. (ii) $S_{n_\phi}(\omega)$ from the spin-echo measurements is consistent with that from the Rabi-oscillation measurements. (iii) $S_{n_\phi}(\omega)$ from the energy relaxation measurements is 2.5 times larger than expected for the decay into a 50 Ω microwave line coupled to the qubit by a mutual inductance of 1.2 pH and nominally cooled to 35 mK. (iv) There can be an additional decoherence induced by strong driving as observed in Fig. 1(d), so it is not surprising to see the increased and scattered $S_{n_\phi}(\omega)$ from the Rabi oscillation measurements above 300 MHz. These data points should be

TABLE I. Parameters in calculations and measurements in units of GHz. In the first column, cal: $\delta\omega(\Omega_{R0})$ stands for the calculation to study the shift of the resonant frequency, and cal: $\Gamma_R^{st}(\delta\omega_{mw})$ stands for the calculation to study the decay of Rabi oscillations due to quasistatic flux noise. "Optimal" in the last column means that at each ε_{mw}, ω_{mw} is chosen to minimize dephasing due to quasistatic flux noise.

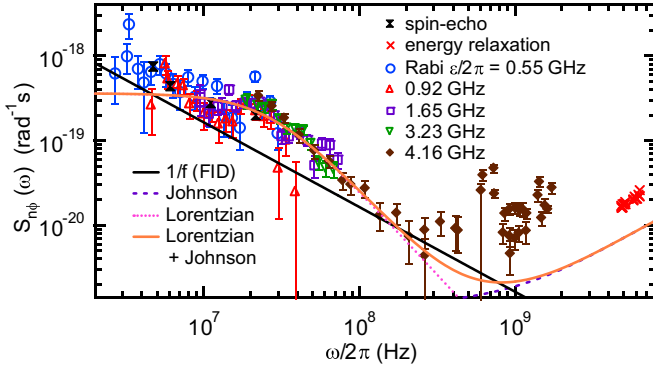|	$\Delta/2\pi$	$\varepsilon/2\pi$	$\varepsilon_{mw}/2\pi$	$\delta\omega_{mw}/2\pi$
cal: $\delta\omega(\Omega_{R0})$	4.869	4.154	1.2–5.0	-0.02–0.12
cal: $\Gamma_R^{st}(\delta\omega_{mw})$	4.869	4.154	4.100	-0.45–0.175
Cooldown 1	4.87	0, 4.16	0.005–4.5	optimal
Cooldown 2	4.82	0.55–3.23	0.02–0.16	optimal

FIG. 3. (Color online) Power spectrum density of flux fluctuations $S_{n_\phi}(\omega)$ extracted from the Rabi oscillation measurements in the first ($\varepsilon/2\pi = 4.16$ GHz) and second cooldowns. The PSDs obtained from the spin-echo and energy relaxation measurements in the second cooldown are also plotted. The black solid line is the $1/f$ spectrum extrapolated from the FID measurements in the second cooldown. The purple dashed line is the estimated Johnson noise from a 50 Ω microwave line coupled to the qubit by a mutual inductance of 1.2 pH and nominally cooled to 35 mK. The pink dotted line is a Lorentzian, $S_{n_\phi}^{\text{model}}(\omega) = S_h \omega_w^2/(\omega^2 + \omega_w^2)$, and the orange solid line is the sum of the Lorentzian and the Johnson noise. Here the parameters are $S_h = 3.6 \times 10^{-19}$ rad^{-1} s and $\omega_w/2\pi = 2.7 \times 10^7$ Hz.

considered as the upper limit of the noise. (v) $S_{n_\phi}(\omega)$ from the Rabi oscillation measurements is roughly parallel to the $1/f$ spectrum extrapolated from the FID measurements but is larger in general and has more structures: the deviation is largest at 25 MHz, and the slope at approximately 100 MHz is steeper than $1/f$. (vi) $S_{n_\phi}(\omega)$ around 300 MHz is approximately 10^{-20} rad^{-1} s, which is (number) orders of magnitude smaller than those reported [17], demonstrating that the noise level is not very far from the extrapolation of the $1/f$ spectrum, even at such high frequencies.

We consider localized electron spins on the surface of the superconducting loop [29–32] as a possible cause of the PSD of flux fluctuations. The total number of electron spins is estimated to be 9×10^6, adopting the reported surface spin density of 5×10^{17} m^{-2} [29] and the total surface area of $\sim 19 \, \mu$m^2 considering both the top and bottom surfaces of the superconducting loop; the loop of the qubit has a $4.8 \times 6.8 \, \mu$m^2 rectangular shape, and the line width is 400 nm.

The magnetic field perpendicularly applied to the qubit loop was approximately 2 G, and screening due to the superconducting film leads to a variation of the field; the magnetic field at the top and bottom surfaces of the loop is shielded, while the field at the edge of the loop is doubled. Considering that the corresponding Zeeman splitting, at most $h \times 11$ MHz, is much smaller than the thermal energy at 35 mK,

the electron spins are expected to be oriented randomly. Because of the broad spectrum of the Zeeman splitting, a clear signal from the electron spin resonance is not expected in $S_{n_\phi}(\omega)$.

We next consider the case where each electron spin generates a random telegraph signal (RTS). The PSD of flux RTSs generated by N independent electron spins is written as a sum of Lorentzians [33]:

$$S_{n_\phi}^{\text{RTS}}(\omega) = \frac{N}{3} n_{\phi e}^2 \sum_{i=1}^{N} \frac{1}{\pi} \frac{2\gamma_i}{\omega^2 + 4\gamma_i^2}, \tag{8}$$

where γ_i is the mean rate of transition per second between two states of the ith electron spin and $n_{\phi e}$ is a normalized flux through the qubit loop in units of Φ_0. Here $n_{\phi e}$ is induced by an electron spin parallel to the magnetic field generated by the persistent current in the qubit loop. For simplicity, we use a constant normalized flux $n_{\phi e} = 1.3 \times 10^{-8}$ [24].

In the case of a $1/f$ spectrum, the distribution function of γ is expressed as $g(\gamma) \propto 1/\gamma$. On the other hand, we speculate that the steep slope at approximately 100 MHz in $S_{n_\phi}(\omega)$ is a part of a Lorentzian, $S_{n_\phi}^{\text{model}}(\omega) = S_h \omega_w^2/(\omega^2 + \omega_w^2)$, where S_h and ω_w are the height and the width of the Lorentzian peak, respectively. In Fig. 3, an example of $S_{n_\phi}^{\text{model}}(\omega)$ is also plotted. Here we chose $S_h = 3.6 \times 10^{-19}$ rad^{-1} s and $\omega_w/2\pi = 2.7 \times 10^7$ Hz, and $S_{n_\phi}^{\text{model}}(\omega)$ amounts to the PSD generated by 3.6×10^6 independent electron spins with the same transition rate of $\gamma = 8.5 \times 10^7$ s^{-1}. The number of electron spins corresponds to approximately 40% of the total surface spins. The number would be smaller in the case where electron spins form ferromagnetic clusters and the spins in each cluster flip simultaneously [34]. The rest of the surface spins may form a $1/f$ spectrum up to a few megahertz, where $S_{n_\phi}^{\text{model}}(\omega)$ deviates from $S_{n_\phi}(\Omega_R)$. To further investigate the origin of the flux noise, a systematic study of the PSD in the high-frequency domain is required.

In conclusion, we have evaluated the PSD of flux fluctuations in a superconducting flux qubit by measuring the decay of Rabi oscillations. The measured Rabi frequency ranges from 2.7 MHz to 1.7 GHz, close to the qubit's level splitting of 4.8 GHz. The observed PSD decreases up to 300 MHz, where the PSD is approximately 10^{-20} rad^{-1} s, not very far from the $1/f$ spectrum extrapolated from the FID measurements.

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